Locally non-geometric fluxes in M-theory and missing momenta

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Based on Blair, EM arXiv:1412.0635; Günaydin, Lüst, EM arXiv:1607.06474; Lüst, EM, Syväri arXiv:1710.05919

- Geometric spaces in general do not allow all possible brane wrapping modes.
- Duality: wrapping \longleftrightarrow momentum.
- Occurs in locally non-geometric backgrounds (c.f. Dieter Lüst's talk).
- Local non-geometry ⇒ non-associativity [Blumenhagen, Plauschinn], [Lüst], [Blumenhagen, Deser, Lüst, Plauschinn, Rennecke], [Mylonas, Schupp, Szabo],

• Toy model T^3 with *H*-flux [Kachru, Schulz, Tripathy, Trivedi]

$$H_{123} \xrightarrow{T_1} T^1_{23} \xrightarrow{T_2} Q^{12}_3 \xrightarrow{T_3} R^{123}$$

- Generate globally non-geometric (*Q*-flux) and locally non-geometric (*R*-flux) spaces.
- Uplift to M-theory? 4-d model with U-dualities (3 directions).

$$T^{1}_{23} \xrightarrow{U_{234}}$$
 locally non-geometric.

• Dependence on dual $\tilde{x}_{\alpha} \longrightarrow \tilde{x}_{ij}$ coordinates.

M-theory duality chain

- Start with "geometric flux", e.g. $\mathcal{N}_3(\textit{N}) imes S^1$
- Nilmanifold $\mathcal{N}_3(N) = \mathbb{R}^3 / \sim$ with

$$egin{aligned} & \left(x^1,\,x^2,\,x^3
ight) \sim \left(x^1+\,1,\,x^2,\,x^3
ight) \sim \left(x^1,\,x^2,\,x^3+\,1
ight) \ & \sim \left(x^1-\mathit{N}\!x^3,\,x^2+\,1,\,x^3
ight)\,. \end{aligned}$$

Can choose metric

$$ds_{\mathcal{N}_{3}}^{2} = \left(dx^{1} + Nx^{2} dx^{3}\right)^{2} + \left(dx^{2}\right)^{2} + \left(dx^{3}\right)^{2}$$

• Parallelisable \Rightarrow well-defined 1-forms

$$e^{\bar{1}} = dx^1 + Nx^2 dx^3$$
, $e^{\bar{2}} = dx^2$, $e^{\bar{3}} = dx^3$.
 $de^{\bar{i}} = T^{\bar{i}}_{\bar{j}\bar{k}}e^{\bar{j}} \wedge e^{\bar{k}}$.

• Geometric flux: $T^{1}_{23} = N$.

Bariloche

Local non-geometry in M-theory

• U-duality along x_2 , x_3 , x_4 gives

$$ds^{2} = (1 + N\tilde{x}_{34})^{-2/3} \left[(dx^{2})^{2} + (dx^{3})^{2} + (dx^{4})^{2} \right] \\ + (1 + N\tilde{x}_{34})^{1/3} (dx^{1})^{2} ,$$

$$C_{(3)} = \frac{N\tilde{x}_{34}}{1 + N^{2}\tilde{x}_{34}} dx^{2} \wedge dx^{3} \wedge dx^{4} .$$

• (g, C_3) patched by U-duality as $\tilde{x}_{34} \longrightarrow \tilde{x}_{34} + 1$.

$$\mathcal{M}_{MN} = \begin{pmatrix} g_{ij} + C_{ikl}C_j^{kl} & C_i^{jk} \\ C^{ij}_{k} & g^{i[k}g^{l]j} \end{pmatrix} = \begin{pmatrix} \hat{g}_{ij} & \Omega_i^{jk} \\ \Omega^{ij}_{k} & \hat{g}^{i[k}\hat{g}^{l]j} + \Omega^{ijm}\Omega_m^{kl} \end{pmatrix}$$

• Non-geometric parameterisation:

$$d\hat{s}^2 = \left(dx^1
ight)^2 + \left(dx^2
ight)^2 + \left(dx^3
ight)^2 + \left(dx^4
ight)^2\,, \qquad \Omega^{124} = -N ilde{x}_{34}\,.$$

• Well-defined up to Ω gauge transformation: $\Omega^{ijk} \longrightarrow \Omega^{ijk} + 3\partial^{[ij}\chi^{k]}$.

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Local non-geometry in M-theory

- Characterised by M-theory R-flux
- Generalised metric + generalised diffeos:

$$\delta_{\xi}\Omega^{ijk} = L_{\xi}\Omega^{ijk} - 3\partial^{[ij}\xi^{k]}.$$

- $R^{i,jklm} = 4\hat{\partial}^{i[j}\Omega^{klm]}$ is spacetime tensor. (c.f. [Andriot, Hohm, Larfors, Lüst, Patalong])
- Also appears in embedding tensor of 7-D gSUGRA.
- $R^{4,1234} = N$.

$$T^1_{23} \xrightarrow{U_{234}} R^{4,1234}$$

• Beyond 4-d compactifications?

Embedding tensor and non-geometric fluxes

- Fluxes generate gaugings of lower-dim gSUGRAs.
- 4-D embedding tensor \in **912** \oplus **56** of $\mathrm{E}_{7(7)}$.
- Under GL(7),

$$\begin{array}{l} 912 \oplus \mathbf{56} \longrightarrow \mathbf{1}_{-14} \oplus \mathbf{35}_{-10} \oplus \overline{\mathbf{140}}_{-6} \oplus 2 \cdot \overline{\mathbf{7}}_{-6} \oplus \mathbf{224}_{-2} \oplus 2 \cdot \mathbf{21}_{-2} \\ \oplus \mathbf{28}_{-2} \oplus \overline{\mathbf{28}}_{2} \oplus \overline{\mathbf{21}}_{2} \oplus \overline{\mathbf{21}}_{2} \oplus \overline{\mathbf{224}}_{2} \oplus 2 \cdot \mathbf{7}_{6} \oplus \mathbf{140}_{6} \\ \oplus \overline{\mathbf{35}}_{10} \oplus \mathbf{1}_{14} \, . \end{array}$$

Construct from

$$\begin{split} \mathbf{133} &\longrightarrow e^{i}{}_{\overline{i}} \in \mathbf{49}_{0} \,, \quad \Omega^{ijk} \in \mathbf{35}_{4} \,, \quad \Omega^{ijklmn} \in \overline{\mathbf{7}}_{8} \,, \quad \dots \\ \partial_{M} &\longrightarrow \partial_{i} \in \overline{\mathbf{7}}_{-6} \,, \quad \partial^{ij} \in \mathbf{21}_{-2} \,, \quad \partial_{ij} \in \overline{\mathbf{21}}_{2} \,, \quad \partial^{i} \in \mathbf{7}_{6} \,. \end{split}$$

Locally non-geometric fluxes in M-theory

 $\bullet \ \mathrm{E}_{7(7)}$ generalised metric + generalised diffeos:

$$\begin{split} \delta_{\xi} \Omega^{ijk} &= L_{\xi} \Omega^{ijk} - 3\partial^{[ij}\xi^{k]} \,, \\ \delta_{\xi} \Omega^{ijklmn} &= L_{\xi} \Omega^{ijklmn} - \frac{1}{10} \epsilon^{ijklmnp} \partial_{pq}\xi^{q} - 6\Omega^{[ijk}\partial^{lm}\xi^{n]} \,. \end{split}$$

• Spacetime tensors

$$\begin{split} R^{i,jklm} &= 4\hat{\partial}^{i[j}\Omega^{klm]} - e_{\overline{i}}^{[j}\epsilon^{klm]inpq}\hat{\partial}_{pq}e^{\overline{i}}_{n}, \\ R^{ij}{}_{k} &= \hat{\partial}_{kl}\Omega^{ijl} - \frac{1}{72}\epsilon_{klmnpqr}\hat{\partial}^{ij}\Omega^{lmnpqr} + \dots, \\ R^{i} &= \hat{\partial}_{jk}\Omega^{ijk} - 4e_{\overline{i}}^{j}\hat{\partial}^{i}e^{\overline{i}}_{j} - 8e_{\overline{i}}^{i}\hat{\partial}^{j}e^{\overline{i}}_{j}, \\ R^{ijklmnp} &= \hat{\partial}^{[i}\Omega^{jklmnp]} - 2\Omega^{[ijk}\hat{\partial}^{m}\Omega^{lnp]}, \\ R^{ijkl} &= \frac{5}{8}\hat{\partial}^{[i}\Omega^{jkl]} + \frac{1}{2}\hat{\partial}_{pq}\Omega^{pqijkl} + \frac{1}{4}\Omega^{[ijk}\hat{\partial}_{pq}\Omega^{l]pq}. \end{split}$$

Non-unimodular geometric flux

Generate all *R*-fluxes from N₃ × T⁴ and S₂ × T⁵.
S₂ − ℝ²/ ~ with

$$\mathcal{S}_2 = \mathbb{R}^2 / \sim \text{with}$$

$$(x_1, x_2) \sim (x_1 + e^{-Nx_2}, x_2)$$
,

• Parallelisable \Rightarrow well-defined 1-forms

$$\begin{split} e^{\bar{1}} &= dx_1 + Nx_1 dx_2 \,, \\ e^{\bar{2}} &= dx_2 \,. \\ de^{\bar{1}} &= e^{\bar{1}} \wedge e^{\bar{2}} \Rightarrow \mathcal{H}^2_{dR}(\mathcal{S}_2) = 0 \,. \end{split}$$

- Non-compact x_2 direction: ∞ -long cylinder.
- Geometric flux: $T^2_{12} = N$.

- Two kinds of U-duality acting in 7-d:
 - U-duality taken along three directions (U_3)
 - U-duality taken along six directions (U_6)
- Read off "Buscher rules" by acting on generalised metric.

912
$$T_{23} \xrightarrow{U_{234}} R^{4,1234} \xrightarrow{U_{156}} R^{14}{}_7$$
,
912 $T_{23} \xrightarrow{U_{234567}} R^{4567} \xrightarrow{U_{123}} R$,
56 $T_{12}^1 \xrightarrow{U_{234}} Q_2^{234} \xrightarrow{U_{125}} R^{5,1234} = -R^{2,1345} \xrightarrow{U_{267}} R^{25}{}_5$
 $\xrightarrow{U_{134}} R^2 = R^{26}{}_6 = R^{27}{}_7$.

Missing wrapping modes

 $\mathcal{N}_3 \times$ T^4 has missing wrapping modes

- e.g. $e^{\bar{2}} \wedge e^{\bar{3}} = \frac{1}{N} de^{\bar{1}}$ $(de^{\bar{i}} = T^{\bar{i}}_{\bar{j}\bar{k}}e^{\bar{j}} \wedge e^{\bar{k}}),$
- \mathcal{N}_3 closed, compact, orientable $\Rightarrow w^{23} = 0$.
- Similarly $w^{23456} = w^{23457} = w^{23567} = w^{23467} = w^{14567} = 0.$
- $w_{KK}^2 = w_{KK}^3 = 0$ (e.g. from self-duality of $\mathcal{N}_3 \times T^4$ under U_{124}).
- Dual to Freed-Witten anomaly in IIB (T^6 with H-flux).

 $\mathcal{S}_2 \times \mathit{T}^5$ missing wrapping states

• x^2 non-compact.

•
$$w^{12} = w^{12345} = w^{12346} = \ldots = w^2_{KK} = 0.$$

U-duality \Rightarrow missing momentum modes!

Missing momenta

$$(w^{23} = w^{23456} = w^{23457} = w^{23467} = w^{23567} = w^{14567} = w^2_{KK} = w^3_{KK} = 0)$$

$$T_{23}^{1} \xrightarrow{U_{234}} R^{4,1234} \xrightarrow{U_{156}} R^{14}_{7}$$

• $R^{4,1234} = N$ has $p_4 = 0$.
• $R^{14}_{7} = N$ has $p_4 = p_1 = 0$.

Missing momenta

$$(w^{23} = w^{23456} = w^{23457} = w^{23467} = w^{23567} = w^{14567} = w^2_{KK} = w^3_{KK} = 0)$$

$$T_{23}^1 \xrightarrow{U_{234}} R^{4,1234} \xrightarrow{U_{156}} R^{14}_7$$

$$R^{4,1234} = N \text{ has } p_4 = 0.$$

$$R^{14}_7 = N \text{ has } p_4 = p_1 = 0.$$

$$T_{23}^1 \xrightarrow{U_{234567}} R^{4567} \xrightarrow{U_{123}} R^{1234567}$$

$$R^{4567} = N \text{ has } p_4 = p_5 = p_6 = p_7 = 0.$$

$$R^{1234567} = N \text{ has all } p_i = 0.$$

Missing momenta

 $(w^{23} = w^{23456} = w^{23457} = w^{23467} = w^{23567} = w^{14567} = w^2_{KK} = w^3_{KK} = 0)$ $T_{22}^1 \xrightarrow{U_{234}} R^{4,1234} \xrightarrow{U_{156}} R^{14}$ • $R^{4,1234} = N$ has $p_4 = 0$. • $R^{14}_{7} = N$ has $p_4 = p_1 = 0$. $T_{23}^1 \xrightarrow{U_{234567}} R^{4567} \xrightarrow{U_{123}} R^{1234567}$ • $R^{4567} = N$ has $p_4 = p_5 = p_6 = p_7 = 0$. • $R^{1234567} = N$ has all $p_i = 0$. $T_{12}^{1} \xrightarrow{U_{234}} O_{2}^{234} \xrightarrow{U_{125}} R^{5,1234} R^{2,1345} \xrightarrow{U_{267}} R^{25} F_{5} \xrightarrow{U_{134}} R^{2} R^{26} R^{27} T_{5}$ • $R^{5,1234} = -R^{2,1345} = N$ has $p_2 = p_5 = 0$. • $R^{25}{}_{5} = N$ has $p_2 = p_5 = 0$. • $R^2 = R^{26}{}_6 = R^{27}{}_7 = N$ has $p_1 = p_2 = p_3 = p_4 = p_5 = 0$.

• These can be summarised as:

$$\begin{aligned} R^{i,jklm} p_i &= 0, \\ \left(R^{ij}_{\ k} - 2R^{[i}\delta^{j]}_k \right) p_i &= 0, \\ R^{ijkl} p_i &= 0, \\ R^{ijklmnp} p_i &= 0. \end{aligned}$$

- In 4-d $R^{4,1234}p_4 = 0$ reduces to IIA strings as no D0-branes.
- Matches expectations from non-associativity ⇒ minimal volume element ⇒ no point particles.
- In 4-d, this leads to 7-dim phase space ⇒ non-associativity based on imaginary octonions.
- More general consequences for non-associativity?

- Solutions rather than toy model.
- Beyond \mathcal{N}_3 , e.g. not parallelisable. Geometric flux \Rightarrow 1st Chern class of U(1)-fibration.
- New R-branes [Bakhmatov, Berman, Kleinschmidt, Musaev, Otsuki].
- Topological result. More general derivation from EFT? Exceptional cohomology controlling allowed momentum / winding states?
- Implications for non-associative theories? Higher-bracket structures, L_{∞} ?

• Two kinds of U-duality acting in 7-d:

• U-duality taken along three directions (U_3)

★ M-theory
$$\xrightarrow{s^1 \longrightarrow 0}$$
 IIA $\xrightarrow{s^1 \longrightarrow 0}$ IIB
+ M theory $\xrightarrow{T^2 \longrightarrow 0}$ IIB

* M-theory
$$\xrightarrow{}$$
 IIB

★ M-theory
$$\xrightarrow{T^{\circ} \longrightarrow 0}$$
 M-theory

• U-duality taken along six directions (U_6)

- Dual IIB geometry: T^6 with *H*-flux, $H_{123} = N$.
- Freed-Witten: no wrapping states $w^{123} = w^{12345} = \ldots = 0$.
- M-theory \longleftrightarrow IIB duality maps these to missing M-theory states.
- $\mathcal{N}_3 \times \mathcal{T}^n$ is an F-theory description of the Freed-Witten anomaly.