

Generalised T-dualities and Integrable deformations

Daniel Thompson

String Dualities and Geometry

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Based roughly on :

1711.0084 (Demulder; Driezen; Sevrin DCT) and

1706.05322+forthcoming (Appadu, Hollowood, Price, DCT)

Overview and Motivation



We will look at classical and quantum features of some classes of integrable sigma models

- Old problem: Which spacetimes give integrable theory?
- Rich interplay of string dualities and holography
- Examples of generalised notions of T-duality
- Triggered advances in DFT (I-modified supergravity)

Plan



λ and η Models



λ , η and Poisson Lie Duality



Quantum aspects of multi-pr η Model



Quantum aspects of multi-pr λ Model



Setting: The Principal Chiral Model



$$S = \int \langle g^{-1} \partial_+ g | g^{-1} \partial_- g \rangle \quad g : \Sigma \rightarrow G$$

- $G_L \times G_R$ global symmetry and classically integrable

$$\mathcal{L}(z) = \frac{1}{1-z^2} J + \frac{z}{1-z^2} \star J, \quad d\mathcal{L} - \mathcal{L} \wedge \mathcal{L} = 0, \quad T(z) = P \exp \int d\sigma \mathcal{L}_\sigma$$

- Factorised scattering S-matrix [\[Zamolodchikov, Zamolodchikov\]](#) and exact results for the mass gap [\[Neidermayer, Hasenfratz\]](#)

$$\mathbb{S}[\theta] = S_{SU(2)} \otimes S_{SU(2)}$$

Non-Abelian Dualisation



- Non-Abelian T-dual [\[de la Ossa, Quevedo\]](#) theory obtained via Buscher procedure
 - Gauge the G_L isometry and add lag. multiplier for flat connection
 - Integrate out dual gauge fields for sigma model
- Useful tool for generating new SUGRA solutions in context of AdS/CFT (see talk of Lozano)
- Some puzzles:— apparent non-compactness of dual geometry, loss of control in global issues, quantum aspects

λ model as a regularisation



“Sfetsos procedure” constructs a new integrable theory

1. Double degrees of freedom: $\kappa^2 S_{PCM}[\tilde{g}] + k S_{WZW}[g]$

2. Gauge G_L in PCM and G_{Diag} in WZW

3. Fix PCM group element to identity

4. Integrate out gauge fields

$$\lambda = \frac{k}{\kappa^2 + k}$$

$$S_\lambda = k S_{WZW} + \frac{k}{2\pi} \int \text{Tr} \left(g^{-1} \partial_+ g \frac{1}{\lambda^{-1} + \text{Ad}_g} \partial_- g g^{-1} \right)$$

λ model limits

- Small λ marginally relevant deformation of WZW

$$S_\lambda|_{\lambda \rightarrow 0} \approx kS_{WZW} + \frac{k}{\pi} \int \lambda J_+^a J_-^a + \mathcal{O}(\lambda^2)$$

- Large deformations in scaling limit recover Non-Abelian T-dual; Sfetsos procedure reduces to Buscher procedure

$$S_\lambda|_{\lambda \rightarrow 1} \approx \frac{1}{\pi} \int \partial_+ X^a (\delta_{ab} + f_{ab}{}^c X_c)^{-1} \partial_- X^b + \mathcal{O}(k^{-1})$$

- An integrable model for all values of λ



Klimcik introduced integrable models based on “modified classical Yang-Baxter equation”

$$S_\eta = \frac{1}{2\pi t} \int_\Sigma d^2\sigma \text{Tr} \left(g^{-1} \partial_+ g, \frac{1}{1 - \eta \mathcal{R}} g^{-1} \partial_- g \right)$$

$$[\mathcal{R}A, \mathcal{R}B] - \mathcal{R}([\mathcal{R}A, B] + [A, \mathcal{R}B]) - [A, B] = 0, \quad \forall A, B \in \mathfrak{g}$$

η model

- Simplest example is the sigma model on squashed S^3

$$S = \frac{1}{2\pi} \int_{\Sigma} d^2\sigma \text{Tr} (g^{-1} \partial_+ g g^{-1} \partial_- g) + C J_+^3 J_-^3$$

- Integrable [\[Cherednik 81\]](#) but with G_R broken to $U(1)$
- Non-local charges recover classical version of $U_q(\text{SU}(2))$
[\[Kawaguchi et al 12\]](#)

$$\{Q_R^+, Q_R^-\}_{P.B.} = \frac{q^{Q_R^3} - q^{-Q_R^3}}{q - q^{-1}}, \quad q = \exp\left(\frac{\sqrt{C}}{1+C}\right)$$

- Same structure for arbitrary group [\[Delduc Magro Vicedo et al\]](#)

Deforming $\text{AdS}_5 \times \text{S}^5$



- Significant interest following discovery of η deformations for cosets and $\text{AdS}_5 \times \text{S}^5$ superstring [\[Delduc Magro Vicedo\]](#)
 - Conjectured q -deformation of holography?
 - TsT can also be thought of as Yang-Baxter deformation
 - Not local Weyl invariant, solves “modified” SUGRA [\[Arutyunov et al\]](#)
 - Relation to DFT established [\[Sakamoto et al; Baguet et al\]](#)
- λ deformations also for cosets $\text{PSU}(2,2|4)$ superstring [\[Hollowood\]](#)
 - Conjectured to be q root-of-unity deformation [\[Hollowood\]](#)
 - Do give rise to genuine (albeit ugly) supergravity solutions [\[Driezen Sfetsos Thompson; Borsato and Wulff\]](#)
 - Squashing of “conformal-coset” (gWZW) backgrounds

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λ and η Models



λ , η and Poisson Lie Duality



Quantum aspects of multi-pr η Model



Quantum aspects of multi-pr λ Model

η λ connection

η and λ models are related via Poisson-Lie T-duality + analytic continuation [Tseytlin Hoare; Siampos Sfetsos DCT; Klimcik]

- η model broke G_R in a very special way:

$$d \star J_a = \tilde{f}^{bc}_a \star J_b \wedge \star J_c$$

- Here we find structure constants for dual algebra:

$$\mathfrak{g}_{\mathcal{R}} \quad [A, B]_{\mathcal{R}} = [\mathcal{R}A, B] + [A, \mathcal{R}B]$$

- So structure constants pure gauge in dual algebra. Suggestive of a duality transformation?

$$\mathfrak{g} \oplus \mathfrak{g}_{\mathcal{R}} \simeq \mathfrak{g}^{\mathbb{C}}$$

- In fact we have the structure of a Drinfeld double

η λ connection

- Not isometric but still dualisable (not via Buscher) in Poisson-Lie sense [\[Klimcik and Severa\]](#)
- Classical canonical transformation between two non-isometric sigma models

$$S[g] = \frac{1}{2\pi t} \int d^2\sigma L_+^T (E - \Pi)^{-1} L_-, \quad g \in \mathcal{G},$$

$$\tilde{S}[\tilde{g}] = \frac{1}{2\pi t} \int d^2\sigma \tilde{L}_+^T (E^{-1} - \tilde{\Pi})^{-1} \tilde{L}_-, \quad \tilde{g} \in \tilde{\mathcal{G}}.$$

$$a_a{}^b = \langle g^{-1} T_a g, \tilde{T}^b \rangle, \quad b^{ab} = \langle g^{-1} \tilde{T}^a g, \tilde{T}^b \rangle, \quad \Pi = b^T a$$

- η model is of this form $E = \eta^{-1} - \mathcal{R}$
- Most examples are pretty ugly.... something nice coming up

Doubled Approach

- PL pairs follow from a parent Doubled Worldsheet sigma model

[[Klimcik Severa, Sfetsos, Hull & Reid-Edwards]

$$S = \int_{\Sigma} d^2\sigma - \mathcal{H}_{AB} \mathbb{L}_{\sigma}^A \mathbb{L}_{\sigma}^B + \eta_{AB} \mathbb{L}_{\sigma}^A \mathbb{L}_{\tau}^B + \int_{\mathcal{M}_3} f_{AB}{}^D \eta_{DC} \mathbb{L}^A \wedge \mathbb{L}^B \wedge \mathbb{L}^C$$

$$\langle T_a | T_b \rangle = \langle \tilde{T}^a | \tilde{T}^b \rangle = 0 \Rightarrow \langle T_A | T_B \rangle = \eta_{AB}$$

- PST covariantisation exists
- Duality manifest one-loop conformal invariance conditions [Avramis, Derendinger, Prezas, Sfetsos-Siampos-DT]
- DFT interpretation later today (Hassler)
- H depends on X and dual directions — beyond section!

η λ connection summary

- After PL T-dualisation of the η theory one gets something...
- A further analytic continuation of some angles and parameters gives the λ model

$$\eta \rightarrow i \frac{1 - \lambda}{1 + \lambda}, \quad t \rightarrow \frac{\pi(1 + \lambda)}{k(1 - \lambda)}$$

$$q = e^{\eta t} \Leftrightarrow q = e^{\frac{i\pi}{k}}$$

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λ , η and Poisson Lie Duality



Quantum aspects of multi-pr η Model



Quantum aspects of multi-pr λ Model

Yang-Baxter + WZ model

- Natural to search for more general models of this type e.g. by inclusion of a WZ term to the η model

$$\mathcal{S} = -\frac{1}{2\pi} \int \langle g^{-1} \partial_+ g, (\alpha \mathbb{I} + \beta \mathcal{R} + \gamma \mathcal{R}^2) g^{-1} \partial_- g \rangle \\ + \frac{k}{24\pi} \int_{M_3} \langle \bar{g}^{-1} d\bar{g}, [\bar{g}^{-1} d\bar{g}, \bar{g}^{-1} d\bar{g}] \rangle$$

- locus of classical integrability [\[Delduc, Magro, Vicedo\]](#)

$$\beta^2 = \frac{\gamma}{\alpha} (\alpha^2 - \alpha\gamma - k^2)$$

- SU(2) case exhibits q -deformed G_R symmetries [\[Kawaguchi Orlando Yoshida\]](#)

$$q = \exp \left[\frac{8\pi\Theta}{\Theta^2 + k^2} \right] \quad \Theta^2 = \frac{\beta^2 \alpha^2}{\gamma^2}$$

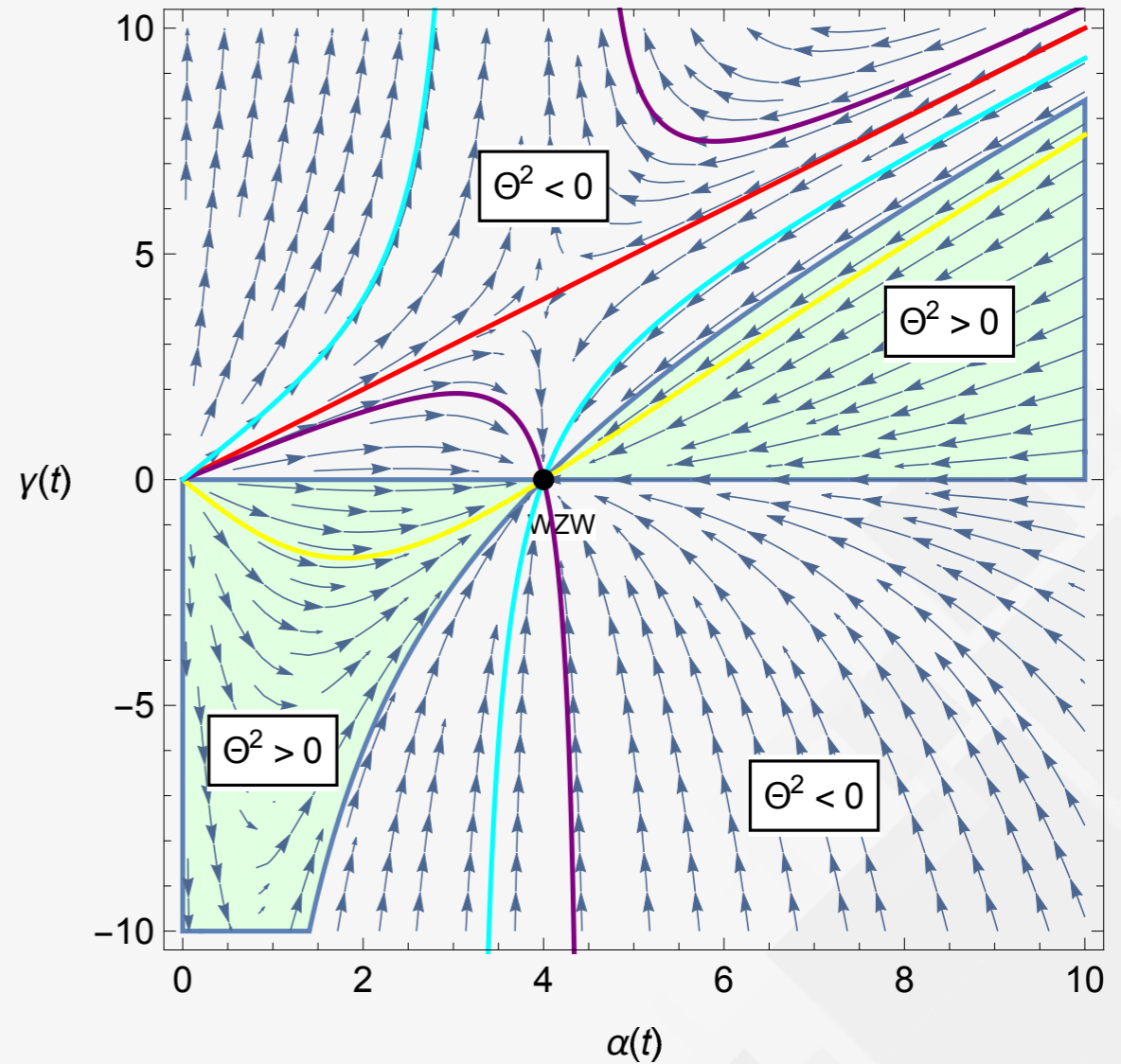
YB+WZ: RG

- Study RG flow equations “geometrically”

$$\mu \frac{d}{d\mu} G_{\mu\nu} = \hat{\beta}_{\mu\nu}^G = \alpha' \left(R_{\mu\nu} - \frac{1}{4} H_{\mu\nu}^2 + \nabla_{(\mu} W_{\nu)} \right) + O(\alpha')^2,$$

$$\mu \frac{d}{d\mu} B_{\mu\nu} = \hat{\beta}_{\mu\nu}^B = \alpha' \left(-\frac{1}{2} \nabla^\lambda H_{\lambda\mu\nu} + (\iota_W H)_{\mu\nu} \right) + O(\alpha')^2$$

- General groups integrable locus is both preserved and required to avoid generating new couplings
- Quantum group parameter q is RG invariant



unlike λ model here WZW is IR fixed point:
solvable marginally irrelevant deformation

YB+WZ: Poisson - Lie

- YB+WZ can, on integrable locus, be obtained as a PL sigma model from its parent 1st order formalism [Klimcik] Same algebra $\mathfrak{g}^{\mathbb{C}}$ but a deformed inner product for the Drinfeld double

$$(z_1, z_2)_{\mathfrak{d}} = \frac{k^2 + \Theta^2}{8\pi\Theta} \text{Im} \left\langle \frac{k + i\Theta}{k - i\Theta} z_1, z_2 \right\rangle$$

- T-dual pairs found by picking out maximal isotropic subspaces
- Messy PL T-duality becomes a really simple radial inversion

$$\alpha \rightarrow \tilde{\alpha} = \frac{k^2}{\alpha},$$

$$\beta \rightarrow \tilde{\beta} = -\beta,$$

$$\gamma \rightarrow \tilde{\gamma} = \frac{k^2 + \alpha\gamma - \alpha^2}{\alpha} = -\frac{\beta^2}{\gamma}.$$

Self dual point == IR conformal invariant fixed point

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λ and η Models



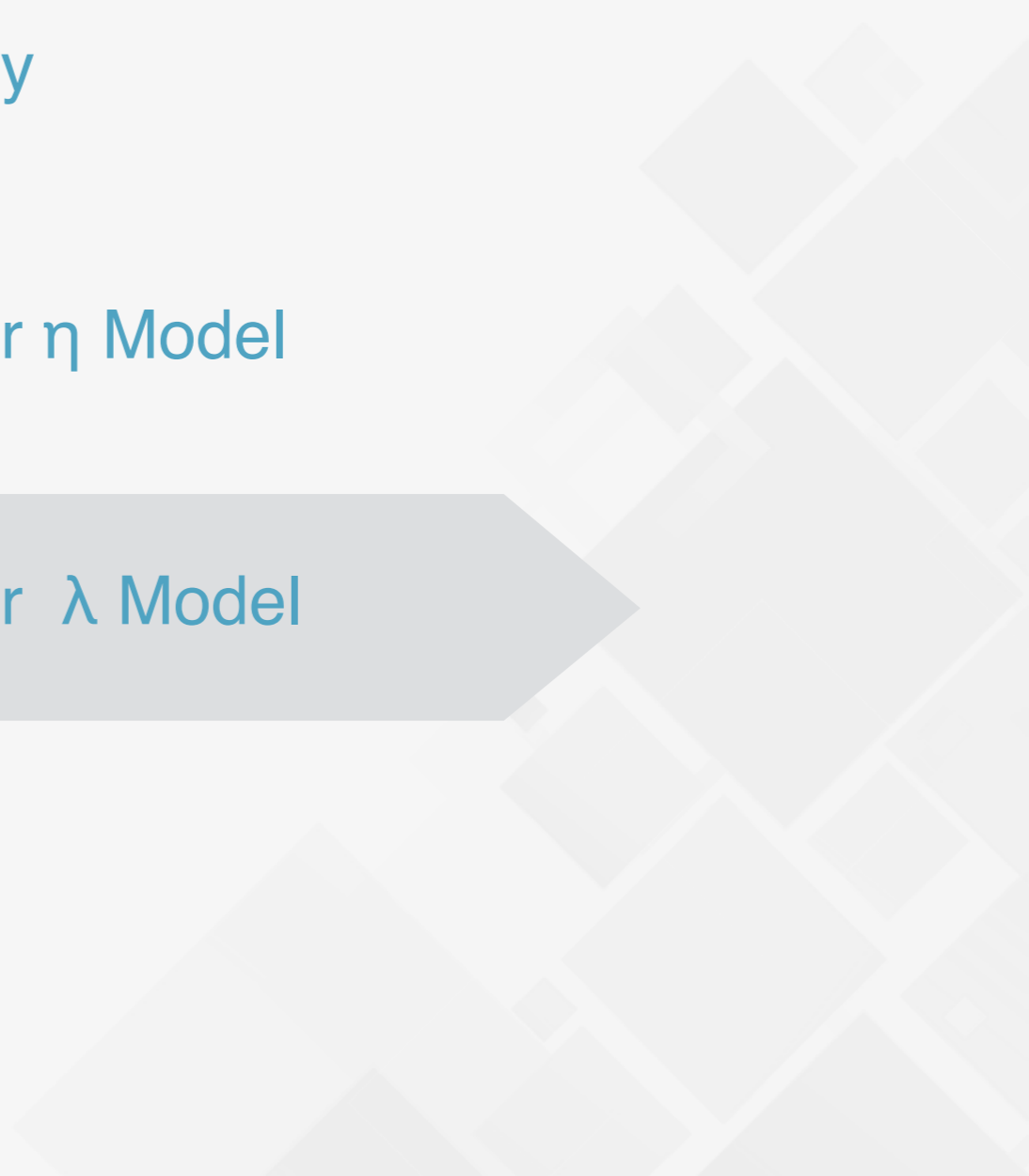
λ , η and Poisson Lie Duality



Quantum aspects of multi-pr η Model



Quantum aspects of multi-pr λ Model



λ XXZ model

- Sfetsos procedure starting not with P.C.M. but any other integrable model gives λ -type with many deformation parameters [\[Siampos Sfetsos Thompson\]](#)

$$S = kS_{WZW}[g] + \frac{k}{2\pi} \int Tr \left(g^{-1} \partial_+ g \frac{1}{\Lambda^{-1} + ad_g} g^{-1} \partial_- g g^{-1} \right)$$

- E.g. start with squashed S^3 we get integrable λ -XXZ model

$$\Lambda = \begin{pmatrix} \xi & & \\ & \xi & \\ & & \lambda \end{pmatrix}$$

- Has a classical $Uq(SU(2))$ (affine — principal gradation)

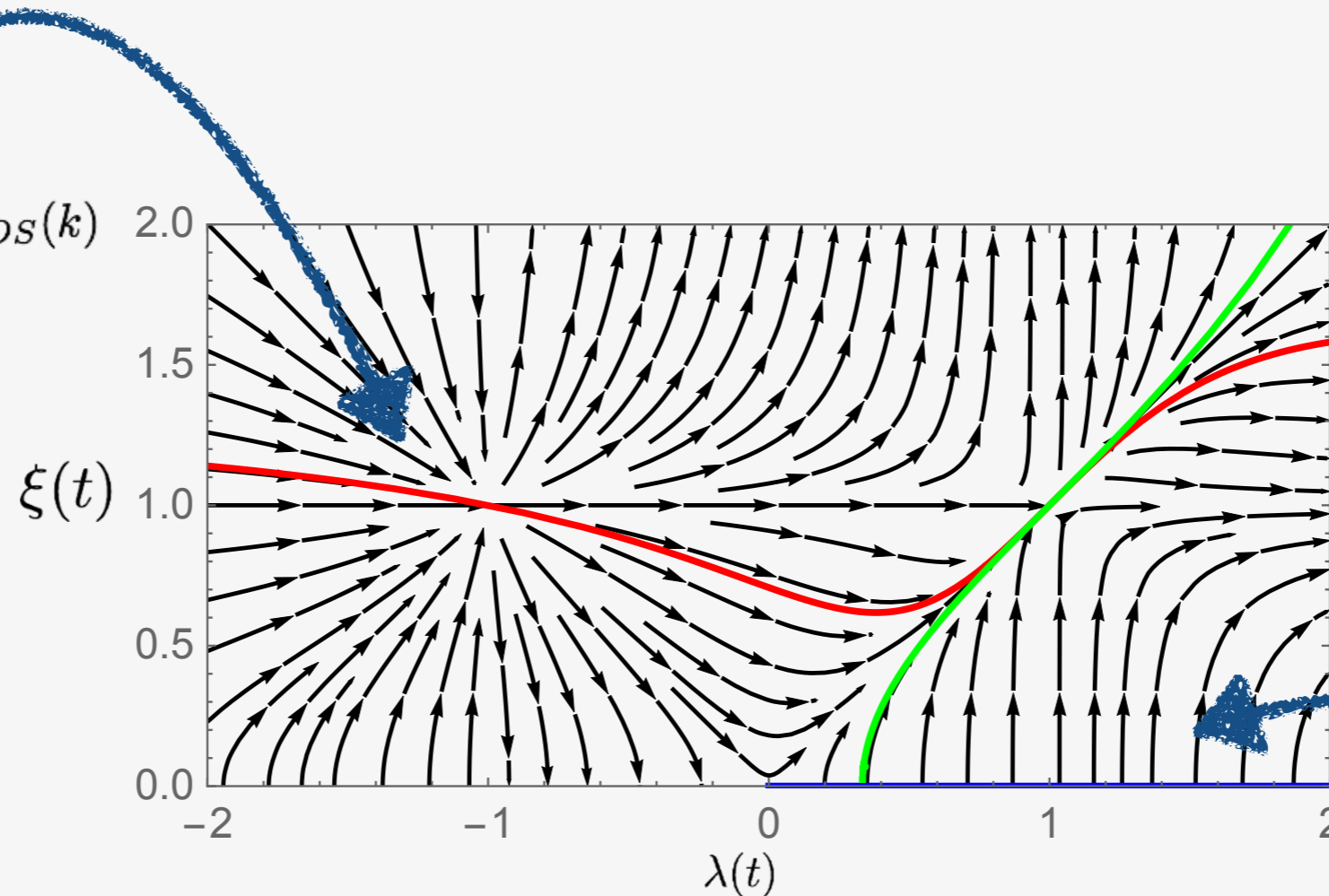
λ XXZ model: RG and S-matrix



Cyclic RG?!?

$$\mathbb{S}[\theta] = S_{cyclic}(\sigma) \times S_{RSOS}(k)$$

$$q = \exp(\pi/\sigma)$$



UV safe regime

$$q = e^{-i\pi/\gamma'}$$

$$\mathbb{S}[\theta] = S_{SG}(\gamma') \times S_{RSOS}(k)$$

λ XXZ as a spin chain



Can we prove conjectured S-matrix?

- Would like to find a lattice regularisation that can be quantised
- Hard to maintain integrability due to “non-ultra-locality” i.e. Schwinger/central term in classical current algebra structure

$$\{J_+^a(\sigma), J_+^b(\sigma')\} = f^{abc} J_+^c(\sigma) \delta(\sigma - \sigma') + \frac{k}{2\pi} \delta'(\sigma - \sigma')$$

- Heuristic idea of Faddeev-Reshetikhin to take a limit

$$k \rightarrow 0, \quad k/\xi, k/\lambda, \text{ fixed}$$

- Central term drops

λ XXZ as a spin chain

- We can now place on a light-cone lattice Destri Devega
- Currents lie on null links at $x_n = n L$ and can be quantised

$$[J_n^a J_m^b] = \frac{i}{L} f^{abc} J_m^c \delta_{mn}$$

- Pick the spin $s = k/2$ representation at each node
- Now able to apply Integrability ideas:
 - Classical Lax Wilson Line uniquely defines quantum R-matrix obeying the Yang-Baxter
 - Spin s XXZ spin chain + impurities

$$H_{XXZ} = \sum_n S_n^1 S_{n+1}^1 + S_n^2 S_{n+1}^2 + \Delta S_n^3 S_{n+1}^3 + \dots$$

$$\Delta = \cos(\gamma)$$

λ XXZ as a spin chain

- Bethe Ansatz techniques can be used to establish the ground state — a sea of 2s-Bethe strings with density ρ
- Excitation above the ground state are holes in the sea with density ρ_H

$$\rho(z) + \rho_h(z) + K_{2s,2s} * \rho(z) = \frac{d}{dz} p_{2s}(z)$$

- Implies scattering Kernel for excitations

$$K_h = -\frac{K_{2s,2s}}{1 + K_{2s,2s}}$$

λ XXZ as a spin chain

- Integrating the scattering kernel gives the S-matrix element

$$K_h = d_z \log S_h(z) \quad S_h(z) = \exp 2i \int \frac{d\omega}{\omega} \sin(\omega z) \left(1 - \frac{\sinh(\pi\omega/2) \tanh(\gamma\omega/2)}{2 \sinh(s\gamma\omega) \sinh((\pi - 2s\gamma)\omega/2)} \right)$$

- Precisely matches conjectured S-matrix $\gamma' = \frac{\pi}{\gamma}$

- Continuum limit:

$$L \rightarrow 0 \quad \nu \rightarrow \infty \quad M = \frac{1}{L} \exp\left(-\frac{\pi\nu}{\gamma}\right), \quad \text{fixed}$$

- Relativistic dispersion relation emerges for particle of mass M
- Note mass generated from cut off and dimensionless parameter


- Beta functions recovered $\mu \frac{d}{d\mu} = \frac{\gamma}{\pi} \quad \mu = L^{-1}$

- *no continuum limit in cyclic regime*



Conclusion



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- η and λ provide paradigms for the interplay of generalised duality and integrability
 - Interesting potential application in holography
 - Simple example where Poisson-Lie is really Buscher rule like
 - Can harness integrability to probe quantum aspects...
 - ...Can we harness this to learn more about duality?