



QUARKONIUM PRODUCTION AT THE LHC

a tool to probe **DPS**

HUA-SHENG SHAO

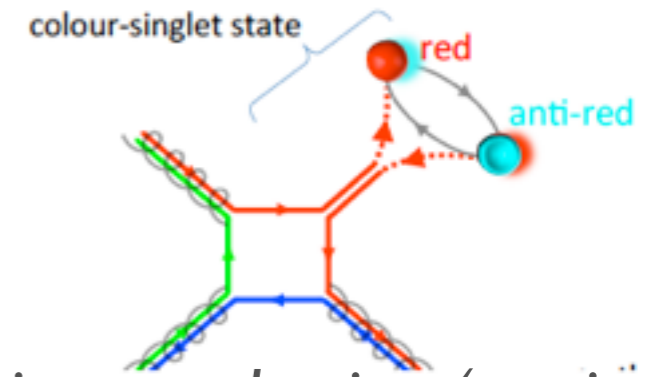


GDR QCD @ IPN ORSAY
01 JUNE 2017

QUARKONIUM PRODUCTION: INTRODUCTION



- **Experiment:** easy to measure and many precise data are available
- **Theory:** various production models
 - Color-Singlet Model (CSM) back in the game
 - *Pro:* good performance; In the game for the total yields
 - *Con:* large QCD corrections; Insufficient to explain inclusive onium production (=onium+jet)
 - Color-Octet Mechanism (COM) predicted by (NR)QCD
 - *Pro:* helps to describe the P_T spectrum of inclusive onium
 - *Con:* debates on its magnitude; only partially works

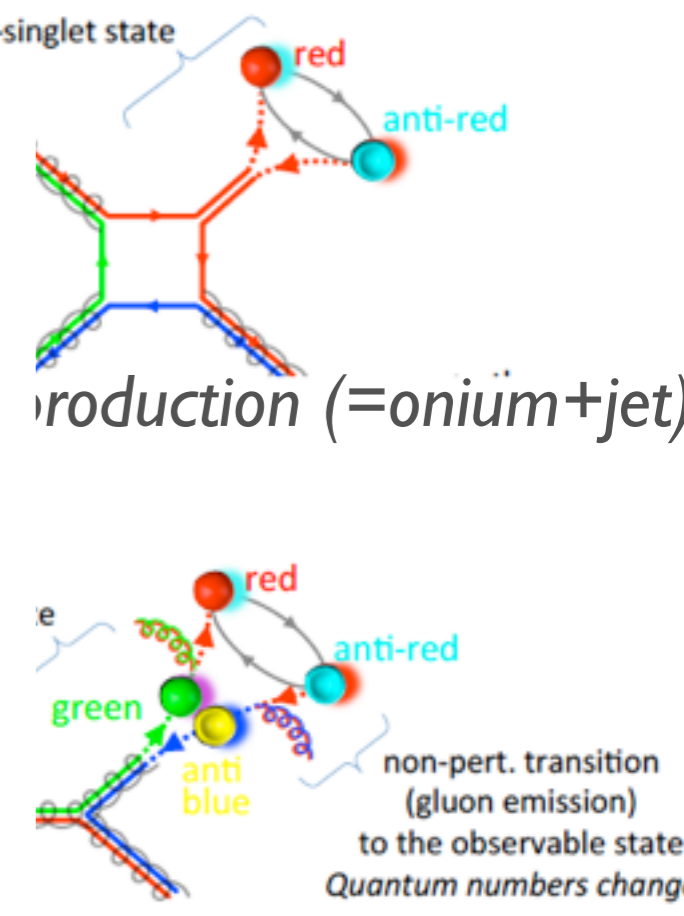
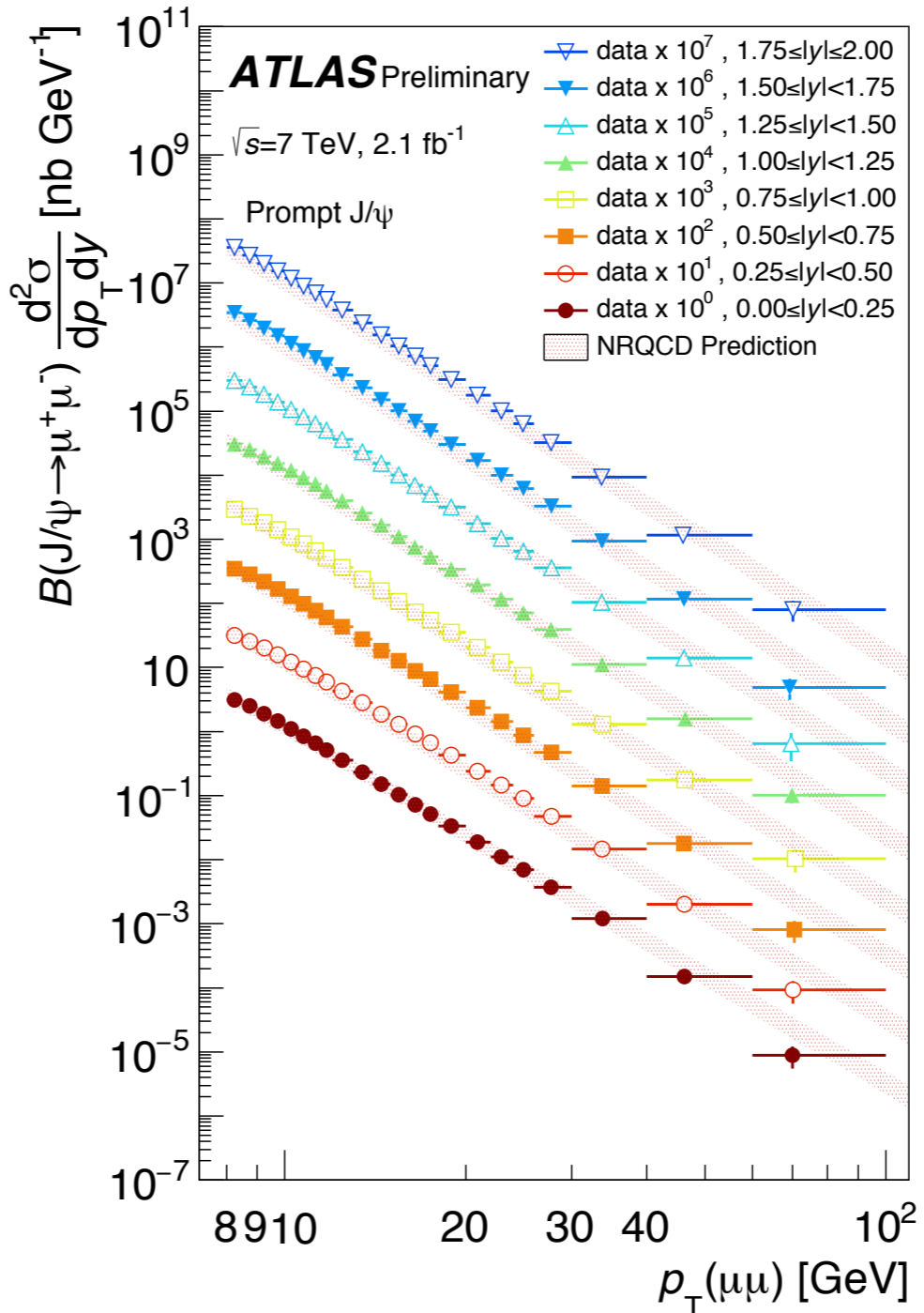
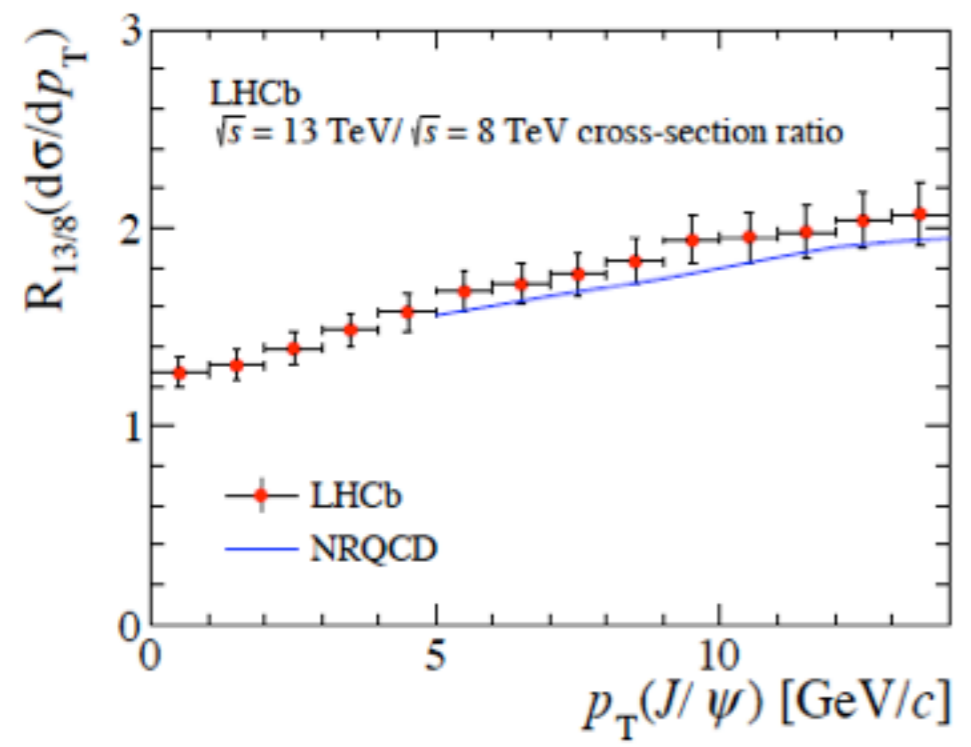


QUARKONIUM PRODUCTION: INTRODUCTION



- **Experiment:** easy to measure and many precise data are available
- **Theory:** various production models

- Color-Singlet Model (CSM)
 - *Pro:* good performance; In the ξ
 - *Con:* large QCD corrections; Ins
- Color-Octet Mechanism (C)
 - *Pro:* helps to describe the P_T sp
 - *Con:* debates on its magnitude;



QUARKONIUM PRODUCTION: INTRODUCTION



- **Experiment:** easy to measure and many precise data are available

- **Theory:** various production models

- Color-Singlet Model (CSM) back in the game

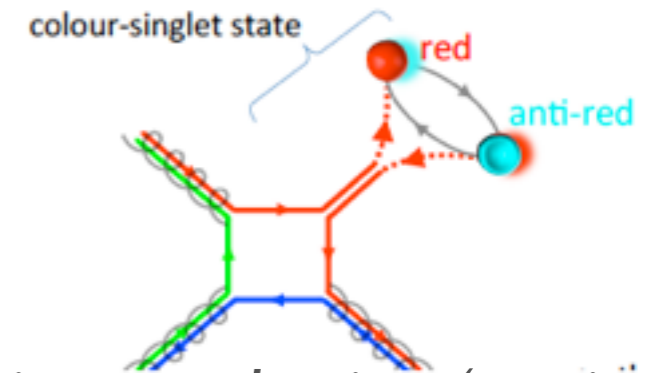
- *Pro:* good performance; In the game for the total yields

- *Con:* large QCD corrections; Insufficient to explain inclusive onium production (=onium+jet)

- Color-Octet Mechanism (COM) predicted by (NR)QCD

- *Pro:* helps to describe the P_T spectrum of inclusive onium

- *Con:* debates on its magnitude; only partially works

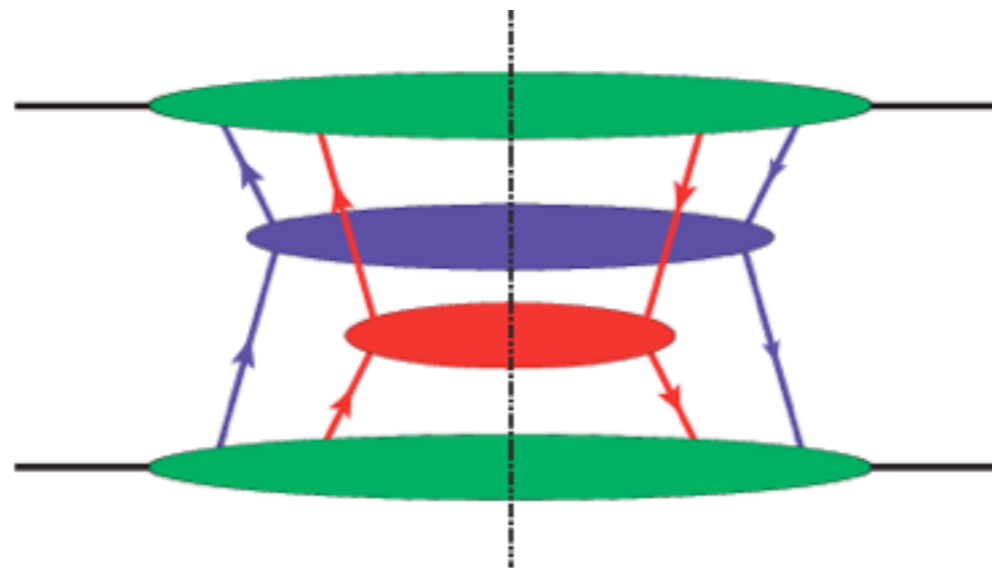


- All approaches have troubles in describing the onium inclusive production data

- This motivates the study of new observables which can be more discriminant for specific effects

- Quarkonium production at the LHC remains a very sensitive probe of the gluon density in the proton.

DOUBLE-PARTON SCATTERINGS



ASSOCIATED (HADRO)PRODUCTION MEASUREMENT

- In associated quarkonium production, many measurements exist

$$J/\psi + J/\psi$$

LHCb, *Phys. Lett. B* 707 (2012) 52; arXiv: 1612.07451

D0, *Phys. Rev. D* 90 (2014) 11, 111101

CMS, *JHEP* 1409 (2014) 094

ATLAS, *Eur. Phys. J. C* (2017) 77:76

$$J/\psi + \Upsilon, \Upsilon + b\bar{b} (= \Upsilon + \text{nonprompt } J/\psi)$$

D0, *Phys. Rev. Lett.* 116 (2016) 082002

$$J/\psi + \text{charm}$$

LHCb, *JHEP* 1206 (2012) 141

$$\Upsilon + \text{charm}$$

LHCb, *JHEP* 1607 (2016) 052

$$J/\psi + W^\pm$$

ATLAS, *JHEP* 1404 (2014) 172

$$J/\psi + Z$$

ATLAS, *Eur. Phys. J. C* 75 (2015) 229

$$\Upsilon + \Upsilon$$

CMS, *JHEP* 1705 (2017) 013

ASSOCIATED (HADRO)PRODUCTION MEASUREMENT

- In associated quarkonium production, many measurements exist

$$J/\psi + J/\psi$$

LHCb, *Phys. Lett. B* 707 (2012) 52; arXiv: 1612.07451

D0, *Phys. Rev. D* 90 (2014) 11, 111101

CMS, *JHEP* 1409 (2014) 094

ATLAS, *Eur. Phys. J. C* (2017) 77:76

$$J/\psi + \Upsilon, \Upsilon + b\bar{b} (= \Upsilon + \text{nonprompt } J/\psi)$$

D0, *Phys. Rev. Lett.* 116 (2016) 082002

$$J/\psi + \text{charm}$$

LHCb, *JHEP* 1206 (2012) 141

$$\Upsilon + \text{charm}$$

LHCb, *JHEP* 1607 (2016) 052

$$J/\psi + W^\pm$$

ATLAS, *JHEP* 1404 (2014) 172

$$J/\psi + Z$$

ATLAS, *Eur. Phys. J. C* 75 (2015) 229

$$\Upsilon + \Upsilon$$

CMS, *JHEP* 1705 (2017) 013

- Challenges in theory:

- taking into account all important **SPS**
- quantifying **DPS**

CALCULATION FRAMEWORK: SPS

- All calculations are performed by the general-purposed matrix-element/event generator *HELAC-Onia* [HSS, CPC '12,'15] with the correct spin-entangled decay.
- **SPS** is calculated in the framework of **NRQCD**:

$$\mathcal{A}_{ab \rightarrow Q_1^{\lambda_1}(P_1) + Q_2^{\lambda_2}(P_2) + X} = \text{WF@Orig or LDME} \quad (10)$$

$$\sum_{s_1, s_2, c_1, c_2} \sum_{s_3, s_4, c_3, c_4} \frac{N(\lambda_1 | s_1, s_2) N(\lambda_2 | s_3, s_4) \delta_{c_1 c_2} \delta_{c_3 c_4} R_1(0) R_2(0)}{\sqrt{M_{Q_1} M_{Q_2}} N_c 4\pi} \mathcal{A}_{ab \rightarrow Q_{c_1}^{s_1} \bar{Q}_{c_2}^{s_2}(p_1=0) + Q_{c_3}^{s_3} \bar{Q}_{c_4}^{s_4}(p_2=0) + X},$$

↑

spin projector

↑

color proj.

↑

four-quark amp.

CALCULATION FRAMEWORK: SPS

- All calculations are performed by the general-purposed matrix-element/event generator *HELAC-Onia* [HSS, CPC '12,'15] with the correct spin-entangled decay.
- SPS is calculated in the framework of NRQCD:

$$\mathcal{A}_{ab \rightarrow Q_1^{\lambda_1}(P_1) + Q_2^{\lambda_2}(P_2) + X} = \sum_{s_1, s_2, c_1, c_2} \sum_{s_3, s_4, c_3, c_4} \frac{N(\lambda_1 | s_1, s_2) N(\lambda_2 | s_3, s_4)}{\sqrt{M_{Q_1} M_{Q_2}}} \frac{\delta_{c_1 c_2} \delta_{c_3 c_4}}{N_c} \frac{R_1(0) R_2(0)}{4\pi} \mathcal{A}_{ab \rightarrow Q_{c_1}^{s_1} \bar{Q}_{c_2}^{s_2}(p_1=0) + Q_{c_3}^{s_3} \bar{Q}_{c_4}^{s_4}(p_2=0) + X} \quad (10)$$

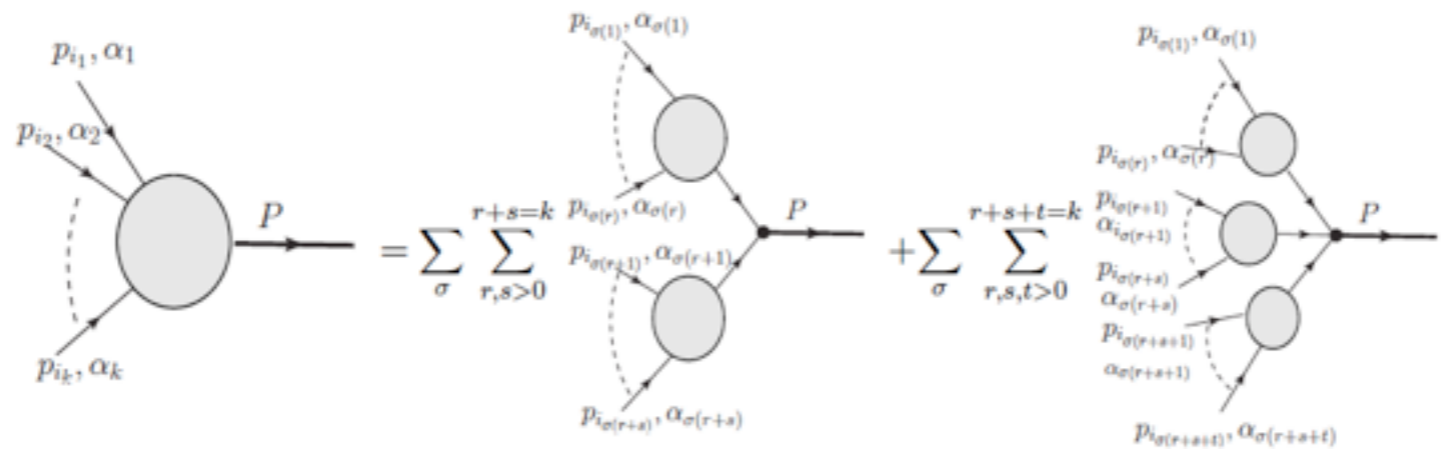
WF@Orig or LDME

spin projector

color proj.

four-quark amp.

- via recursion relations

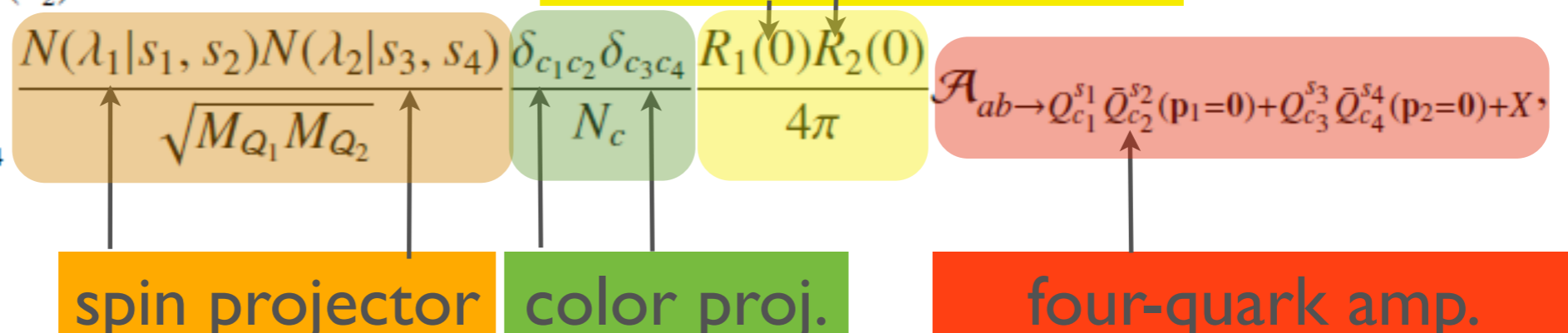


CALCULATION FRAMEWORK: SPS

- All calculations are performed by the general-purposed matrix-element/event generator *HELAC-Onia* [HSS, CPC '12,'15] with the correct spin-entangled decay.
- SPS is calculated in the framework of NRQCD:

$$\mathcal{A}_{ab \rightarrow Q_1^{\lambda_1}(P_1) + Q_2^{\lambda_2}(P_2) + X} = \sum_{s_1, s_2, c_1, c_2} \sum_{s_3, s_4, c_3, c_4} \frac{N(\lambda_1 | s_1, s_2) N(\lambda_2 | s_3, s_4)}{\sqrt{M_{Q_1} M_{Q_2}}} \frac{\delta_{c_1 c_2} \delta_{c_3 c_4}}{N_c} \frac{R_1(0) R_2(0)}{4\pi} \mathcal{A}_{ab \rightarrow Q_{c_1}^{s_1} \bar{Q}_{c_2}^{s_2}(p_1=0) + Q_{c_3}^{s_3} \bar{Q}_{c_4}^{s_4}(p_2=0) + X}, \quad (10)$$

WF@Orig or LDME



The diagram illustrates the decomposition of the NRQCD amplitude into three main components:

- spin projector** (orange box): $\frac{N(\lambda_1 | s_1, s_2) N(\lambda_2 | s_3, s_4)}{\sqrt{M_{Q_1} M_{Q_2}}}$
- color proj.** (green box): $\frac{\delta_{c_1 c_2} \delta_{c_3 c_4}}{N_c}$
- four-quark amp.** (red box): $\mathcal{A}_{ab \rightarrow Q_{c_1}^{s_1} \bar{Q}_{c_2}^{s_2}(p_1=0) + Q_{c_3}^{s_3} \bar{Q}_{c_4}^{s_4}(p_2=0) + X}$

Additional factors in the equation include $R_1(0)R_2(0)/4\pi$ and the label **WF@Orig or LDME** (yellow box) pointing to the overall structure.

- via recursion relations
- potential model or from data (but should be universal)

CALCULATION FRAMEWORK: DPS

- All calculations are performed by the general-purposed matrix-element/event generator *HELAC-Onia* [HSS, CPC '12,'15] with the correct spin-entangled decay.
- **DPS** has the general formula via

$$\sigma_{Q_1 Q_2} = \frac{1}{1 + \delta_{Q_1 Q_2}} \sum_{i,j,k,l} \int dx_1 dx_2 dx'_1 dx'_2 d^2 \mathbf{b}_1 d^2 \mathbf{b}_2 d^2 \mathbf{b} \\ \times \Gamma_{ij}(x_1, x_2, \mathbf{b}_1, \mathbf{b}_2) \hat{\sigma}_{ik}^{Q_1}(x_1, x'_1) \hat{\sigma}_{jl}^{Q_2}(x_2, x'_2) \Gamma_{kl}(x'_1, x'_2, \mathbf{b}_1 - \mathbf{b}, \mathbf{b}_2 - \mathbf{b}),$$

CALCULATION FRAMEWORK: DPS

- All calculations are performed by the general-purposed matrix-element/event generator *HELAC-Onia* [HSS, CPC '12,'15] with the correct spin-entangled decay.
- **DPS** has the general formula via

Generalised double distribution

Single-quarkonium parton-level XS

$$\sigma_{Q_1 Q_2} = \frac{1}{1 + \delta_{Q_1 Q_2}} \sum_{i,j,k,l} \int dx_1 dx_2 dx'_1 dx'_2 d^2 \mathbf{b}_1 d^2 \mathbf{b}_2 d^2 \mathbf{b}$$

$$\times \Gamma_{ij}(x_1, x_2, \mathbf{b}_1, \mathbf{b}_2) \hat{\sigma}_{ik}^{Q_1}(x_1, x'_1) \hat{\sigma}_{jl}^{Q_2}(x_2, x'_2) \Gamma_{kl}(x'_1, x'_2, \mathbf{b}_1 - \mathbf{b}, \mathbf{b}_2 - \mathbf{b}),$$

CALCULATION FRAMEWORK: DPS

- All calculations are performed by the general-purposed matrix-element/event generator *HELAC-Onia* [HSS, CPC '12,'15] with the correct spin-entangled decay.
- **DPS** has the general formula via

Generalised double distribution

Single-quarkonium parton-level XS

$$\sigma_{Q_1 Q_2} = \frac{1}{1 + \delta_{Q_1 Q_2}} \sum_{i,j,k,l} \int dx_1 dx_2 dx'_1 dx'_2 d^2 \mathbf{b}_1 d^2 \mathbf{b}_2 d^2 \mathbf{b}$$

$$\times \Gamma_{ij}(x_1, x_2, \mathbf{b}_1, \mathbf{b}_2) \hat{\sigma}_{ik}^{Q_1}(x_1, x'_1) \hat{\sigma}_{jl}^{Q_2}(x_2, x'_2) \Gamma_{kl}(x'_1, x'_2, \mathbf{b}_1 - \mathbf{b}, \mathbf{b}_2 - \mathbf{b}),$$

- Factorization I $\Gamma_{ij}(x_1, x_2, \mathbf{b}_1, \mathbf{b}_2) = D_{ij}(x_1, x_2) T_{ij}(\mathbf{b}_1, \mathbf{b}_2),$

CALCULATION FRAMEWORK: DPS

- All calculations are performed by the general-purposed matrix-element/event generator *HELAC-Onia* [HSS, CPC '12,'15] with the correct spin-entangled decay.
- **DPS** has the general formula via

Generalised double distribution

Single-quarkonium parton-level XS

$$\sigma_{Q_1 Q_2} = \frac{1}{1 + \delta_{Q_1 Q_2}} \sum_{i,j,k,l} \int dx_1 dx_2 dx'_1 dx'_2 d^2 \mathbf{b}_1 d^2 \mathbf{b}_2 d^2 \mathbf{b}$$

$$\times \Gamma_{ij}(x_1, x_2, \mathbf{b}_1, \mathbf{b}_2) \hat{\sigma}_{ik}^{Q_1}(x_1, x'_1) \hat{\sigma}_{jl}^{Q_2}(x_2, x'_2) \Gamma_{kl}(x'_1, x'_2, \mathbf{b}_1 - \mathbf{b}, \mathbf{b}_2 - \mathbf{b}),$$

- Factorization I

$$\Gamma_{ij}(x_1, x_2, \mathbf{b}_1, \mathbf{b}_2) = D_{ij}(x_1, x_2) T_{ij}(\mathbf{b}_1, \mathbf{b}_2),$$

dPDF

CALCULATION FRAMEWORK: DPS

- All calculations are performed by the general-purposed matrix-element/event generator *HELAC-Onia* [HSS, CPC '12,'15] with the correct spin-entangled decay.
- **DPS** has the general formula via

Generalised double distribution

Single-quarkonium parton-level XS

$$\sigma_{Q_1 Q_2} = \frac{1}{1 + \delta_{Q_1 Q_2}} \sum_{i,j,k,l} \int dx_1 dx_2 dx'_1 dx'_2 d^2 \mathbf{b}_1 d^2 \mathbf{b}_2 d^2 \mathbf{b}$$

$$\times \Gamma_{ij}(x_1, x_2, \mathbf{b}_1, \mathbf{b}_2) \hat{\sigma}_{ik}^{Q_1}(x_1, x'_1) \hat{\sigma}_{jl}^{Q_2}(x_2, x'_2) \Gamma_{kl}(x'_1, x'_2, \mathbf{b}_1 - \mathbf{b}, \mathbf{b}_2 - \mathbf{b}),$$

- Factorization I $\Gamma_{ij}(x_1, x_2, \mathbf{b}_1, \mathbf{b}_2) = D_{ij}(x_1, x_2) T_{ij}(\mathbf{b}_1, \mathbf{b}_2),$

dPDF

- Factorization II $D_{ij}(x_1, x_2) = f_i(x_1) f_j(x_2),$
 $T_{ij}(\mathbf{b}_1, \mathbf{b}_2) = T_i(\mathbf{b}_1) T_j(\mathbf{b}_2),$

CALCULATION FRAMEWORK: DPS

- All calculations are performed by the general-purposed matrix-element/event generator *HELAC-Onia* [HSS, CPC '12,'15] with the correct spin-entangled decay.
- **DPS** has the general formula via

Generalised double distribution

Single-quarkonium parton-level XS

$$\sigma_{Q_1 Q_2} = \frac{1}{1 + \delta_{Q_1 Q_2}} \sum_{i,j,k,l} \int dx_1 dx_2 dx'_1 dx'_2 d^2 \mathbf{b}_1 d^2 \mathbf{b}_2 d^2 \mathbf{b}$$

$$\times \Gamma_{ij}(x_1, x_2, \mathbf{b}_1, \mathbf{b}_2) \hat{\sigma}_{ik}^{Q_1}(x_1, x'_1) \hat{\sigma}_{jl}^{Q_2}(x_2, x'_2) \Gamma_{kl}(x'_1, x'_2, \mathbf{b}_1 - \mathbf{b}, \mathbf{b}_2 - \mathbf{b}),$$

- Factorization I $\Gamma_{ij}(x_1, x_2, \mathbf{b}_1, \mathbf{b}_2) = D_{ij}(x_1, x_2) T_{ij}(\mathbf{b}_1, \mathbf{b}_2),$

dPDF

- Factorization II $D_{ij}(x_1, x_2) = f_i(x_1) f_j(x_2),$
 $T_{ij}(\mathbf{b}_1, \mathbf{b}_2) = T_i(\mathbf{b}_1) T_j(\mathbf{b}_2),$

PDF

CALCULATION FRAMEWORK: DPS

- All calculations are performed by the general-purposed matrix-element/event generator *HELAC-Onia* [HSS, CPC '12,'15] with the correct spin-entangled decay.
- **DPS** has the general formula via

Generalised double distribution

Single-quarkonium parton-level XS

$$\sigma_{Q_1 Q_2} = \frac{1}{1 + \delta_{Q_1 Q_2}} \sum_{i,j,k,l} \int dx_1 dx_2 dx'_1 dx'_2 d^2 \mathbf{b}_1 d^2 \mathbf{b}_2 d^2 \mathbf{b} \\ \times \Gamma_{ij}(x_1, x_2, \mathbf{b}_1, \mathbf{b}_2) \hat{\sigma}_{ik}^{Q_1}(x_1, x'_1) \hat{\sigma}_{jl}^{Q_2}(x_2, x'_2) \Gamma_{kl}(x'_1, x'_2, \mathbf{b}_1 - \mathbf{b}, \mathbf{b}_2 - \mathbf{b}),$$

- Factorization I $\Gamma_{ij}(x_1, x_2, \mathbf{b}_1, \mathbf{b}_2) = D_{ij}(x_1, x_2) T_{ij}(\mathbf{b}_1, \mathbf{b}_2),$

dPDF

- Factorization II $D_{ij}(x_1, x_2) = f_i(x_1) f_j(x_2),$
 $T_{ij}(\mathbf{b}_1, \mathbf{b}_2) = T_i(\mathbf{b}_1) T_j(\mathbf{b}_2),$

PDF

- Assume flavor universality in \bar{T}

$$\sigma_{Q_1 Q_2} = \frac{1}{1 + \delta_{Q_1 Q_2}} \frac{\sigma_{Q_1} \sigma_{Q_2}}{\sigma_{\text{eff}}},$$

$$\sigma_{\text{eff}} = \left[\int d^2 \mathbf{b} F(\mathbf{b})^2 \right]^{-1}.$$

$$F(\mathbf{b}) = \int T(\mathbf{b}_i) T(\mathbf{b}_i - \mathbf{b}) d^2 \mathbf{b}_i,$$

CALCULATION FRAMEWORK: DPS

- All calculations are performed by the general-purposed matrix-element/event generator *HELAC-Onia* [HSS, CPC '12,'15] with the correct spin-entangled decay.
- **DPS** has the general formula via

Generalised double distribution

Single-quarkonium parton-level XS

$$\sigma_{Q_1 Q_2} = \frac{1}{1 + \delta_{Q_1 Q_2}} \sum_{i,j,k,l} \int dx_1 dx_2 dx'_1 dx'_2 d^2 \mathbf{b}_1 d^2 \mathbf{b}_2 d^2 \mathbf{b} \\ \times \Gamma_{ij}(x_1, x_2, \mathbf{b}_1, \mathbf{b}_2) \hat{\sigma}_{ik}^{Q_1}(x_1, x'_1) \hat{\sigma}_{jl}^{Q_2}(x_2, x'_2) \Gamma_{kl}(x'_1, x'_2, \mathbf{b}_1 - \mathbf{b}, \mathbf{b}_2 - \mathbf{b}),$$

- Factorization I $\Gamma_{ij}(x_1, x_2, \mathbf{b}_1, \mathbf{b}_2) = D_{ij}(x_1, x_2) T_{ij}(\mathbf{b}_1, \mathbf{b}_2),$

dPDF

- Factorization II $D_{ij}(x_1, x_2) = f_i(x_1) f_j(x_2),$
 $T_{ij}(\mathbf{b}_1, \mathbf{b}_2) = T_i(\mathbf{b}_1) T_j(\mathbf{b}_2),$

PDF

- Assume flavor universality in T

$$\sigma_{\text{eff}} = \left[\int d^2 \mathbf{b} F(\mathbf{b})^2 \right]^{-1}.$$

$$F(\mathbf{b}) = \int T(\mathbf{b}_i) T(\mathbf{b}_i - \mathbf{b}) d^2 \mathbf{b}_i,$$



$$\sigma_{Q_1 Q_2} = \frac{1}{1 + \delta_{Q_1 Q_2}} \frac{\sigma_{Q_1} \sigma_{Q_2}}{\sigma_{\text{eff}}},$$

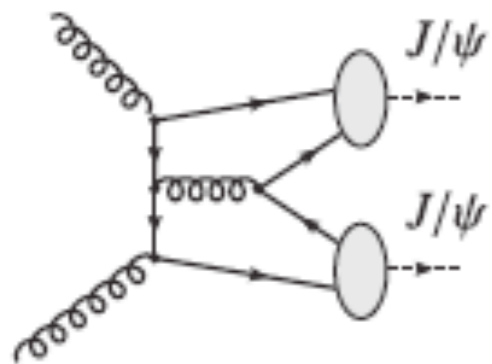
Pocket Formula



LPTHE
LABORATOIRE DE PHYSIQUE
THEORIQUE ET HAUTES ENERGIES

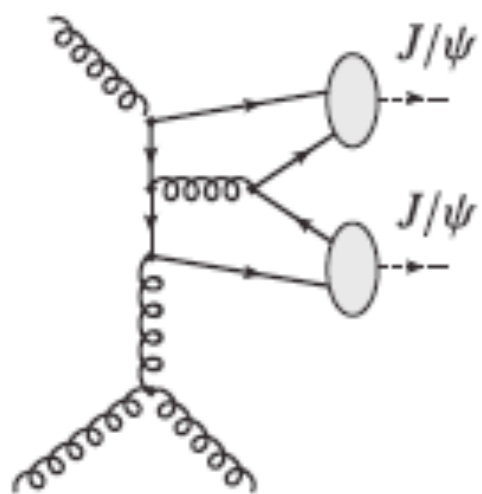
$$J/\psi + J/\psi$$

QCD CORRECTIONS TO SPS



\sim

$$\alpha_s^4 \left(\frac{m_\psi}{P_T^\psi} \right)^8$$

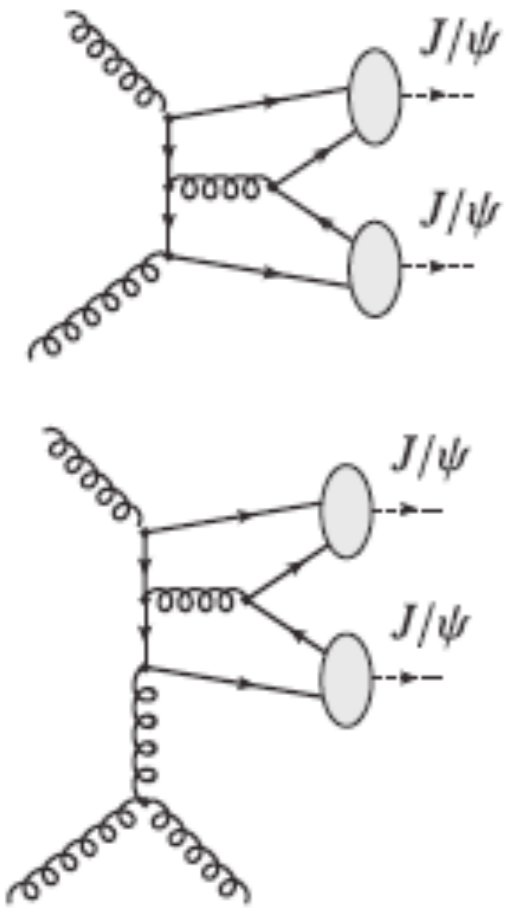


\sim

$$\alpha_s^5 \left(\frac{m_\psi}{P_T^\psi} \right)^6$$

QCD CORRECTIONS TO SPS

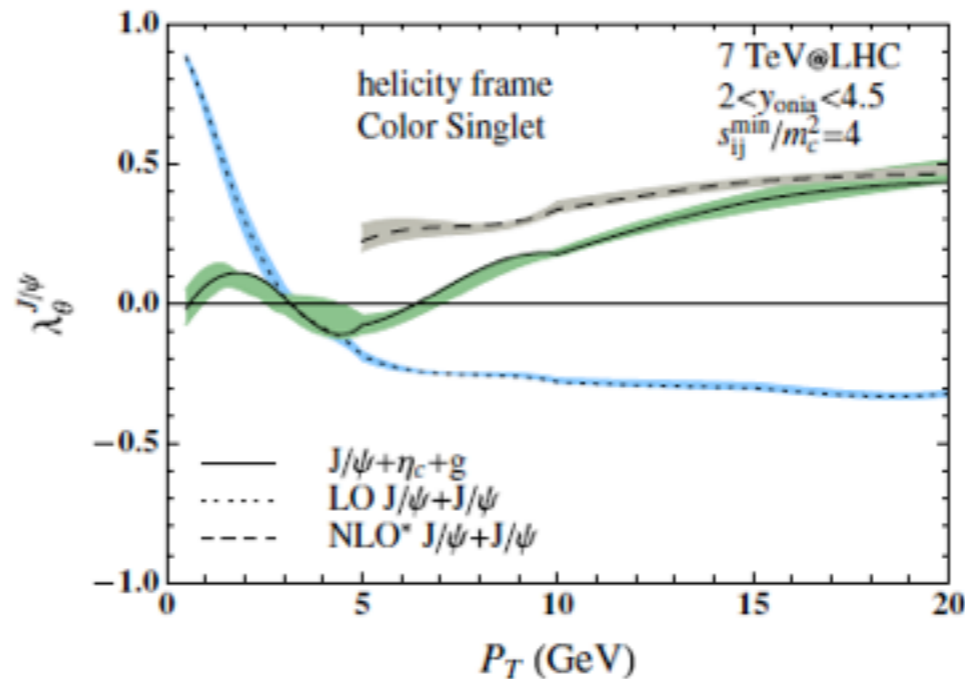
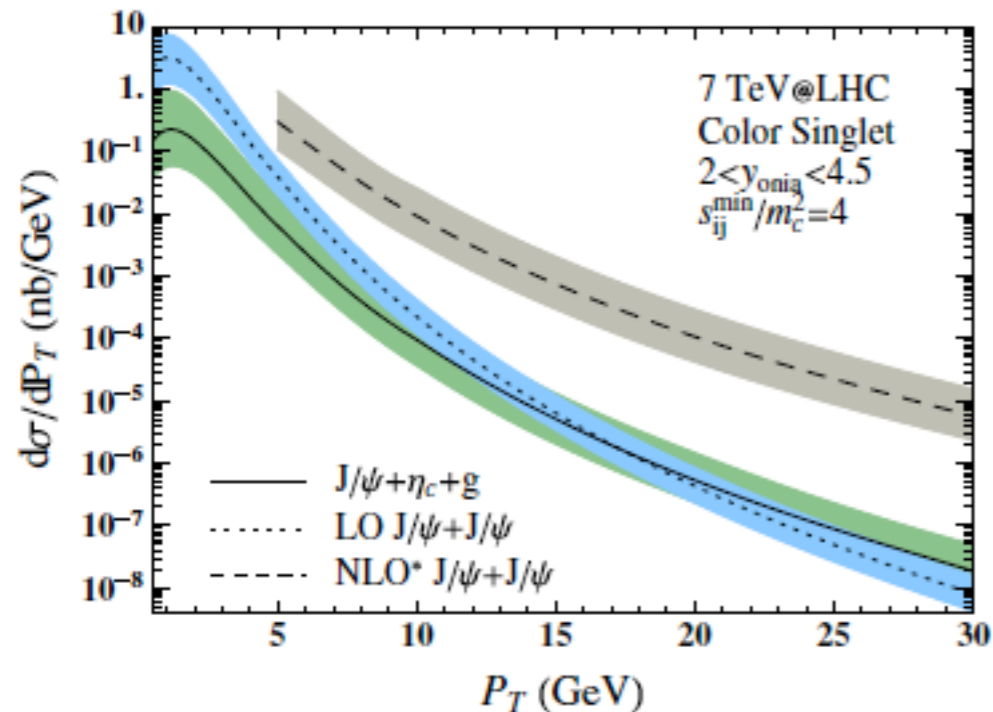
Lansberg, HSS PRL '13



$$\sim \alpha_s^4 \left(\frac{m_\psi}{P_T^\psi} \right)^8$$

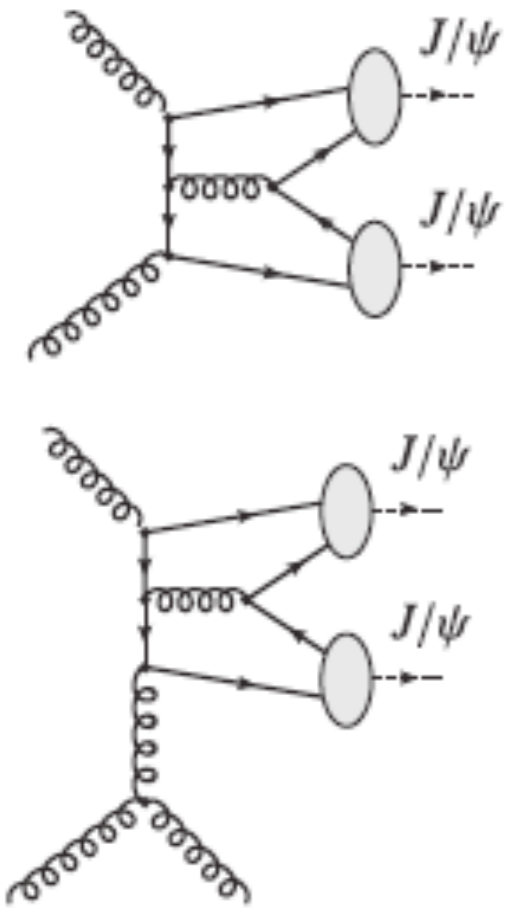
$$\sim \alpha_s^5 \left(\frac{m_\psi}{P_T^\psi} \right)^6$$

- For the first time, we calculated the leading- P_T contribution at α_s^5 with **HELAC-Onia** [HSS, CPC '13,'15].
- It was nicely confirmed by a complete NLO calculation [Sun, Han, Chao, '14].



QCD CORRECTIONS TO SPS

Lansberg, HSS PRL '13

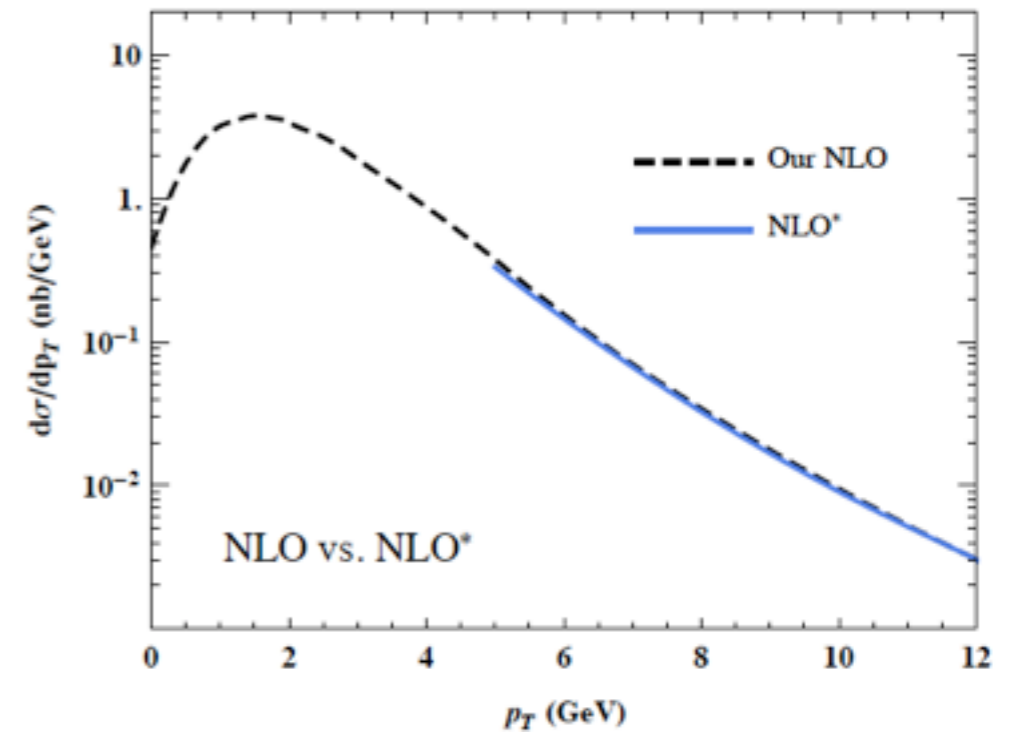
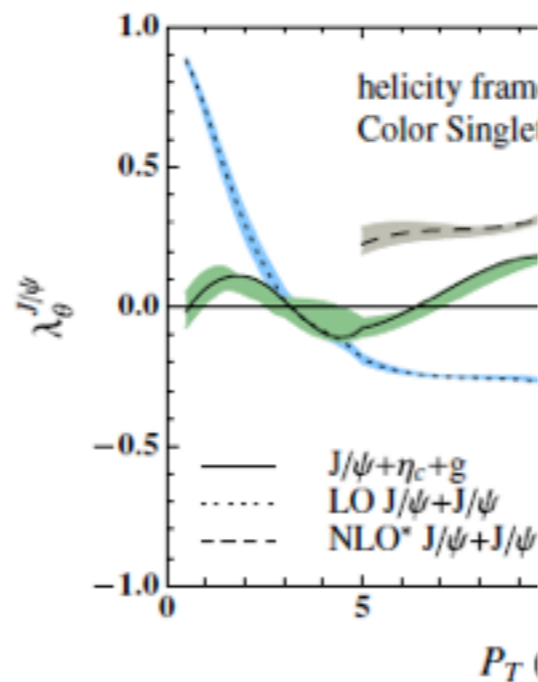
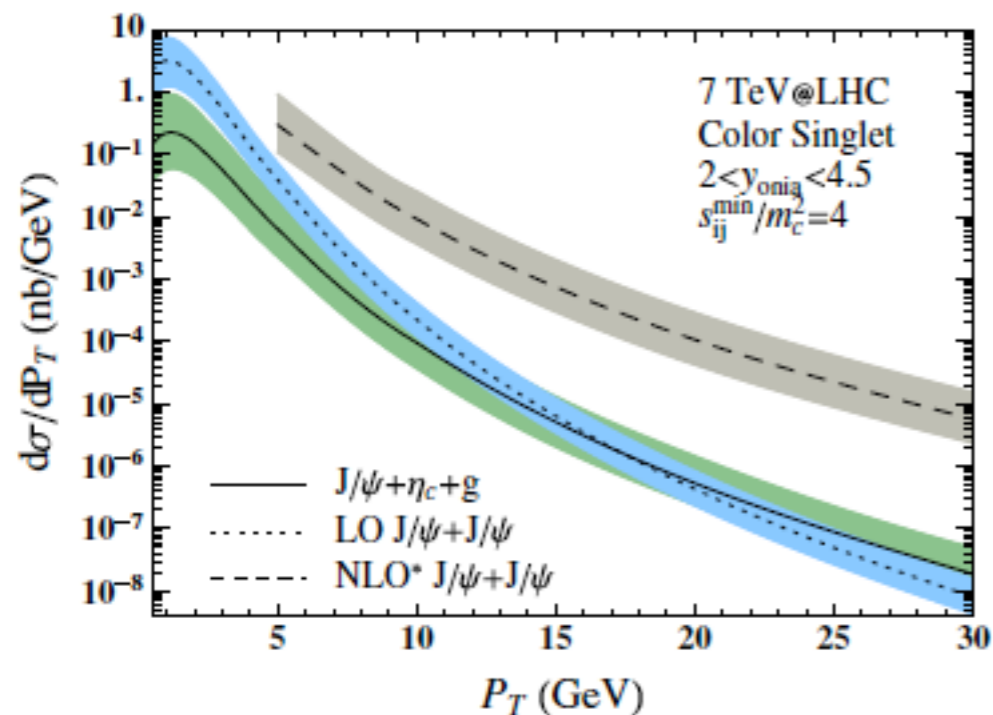


$$\sim \alpha_s^4 \left(\frac{m_\psi}{P_T^\psi} \right)^8$$

$$\sim \alpha_s^5 \left(\frac{m_\psi}{P_T^\psi} \right)^6$$

- For the first time, we calculated the leading- P_T contribution at α_s^5 with **HELAC-Onia** [HSS, CPC '13,'15].

- It was nicely confirmed by a complete NLO calculation [Sun, Han, Chao, '14].



EVIDENCE OF DPS ?



PUBLISHED FOR SISSA BY SPRINGER

RECEIVED: June 2, 2014

ACCEPTED: August 4, 2014

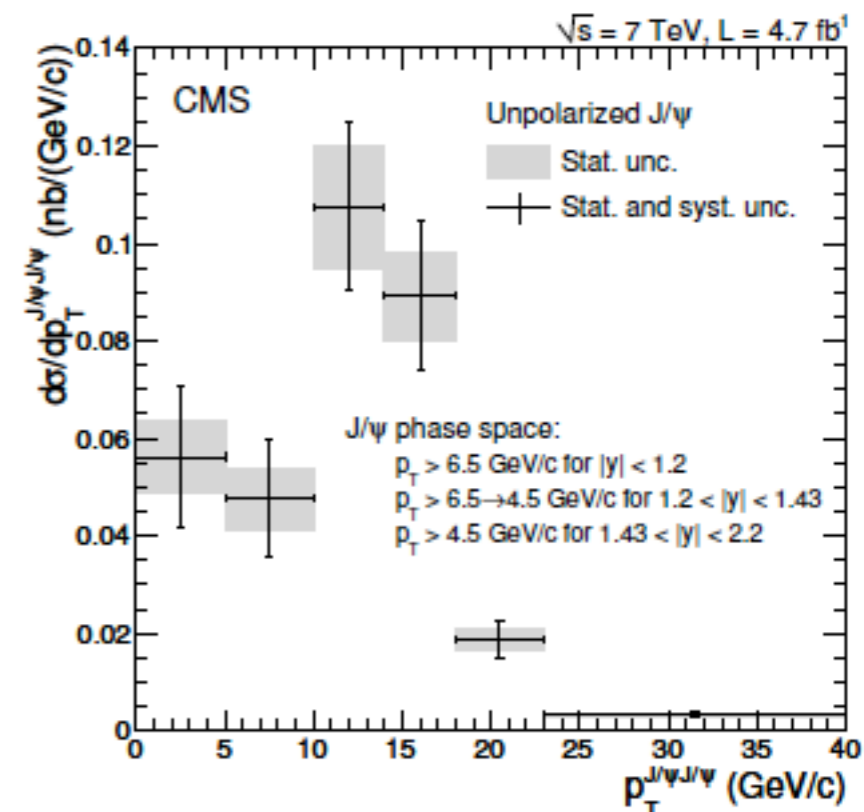
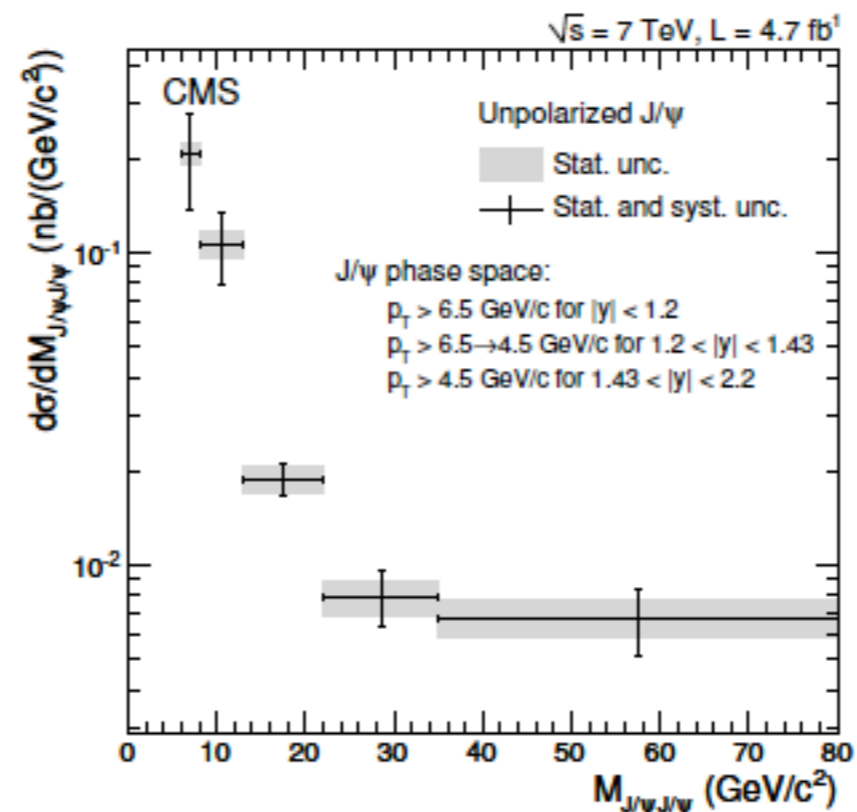
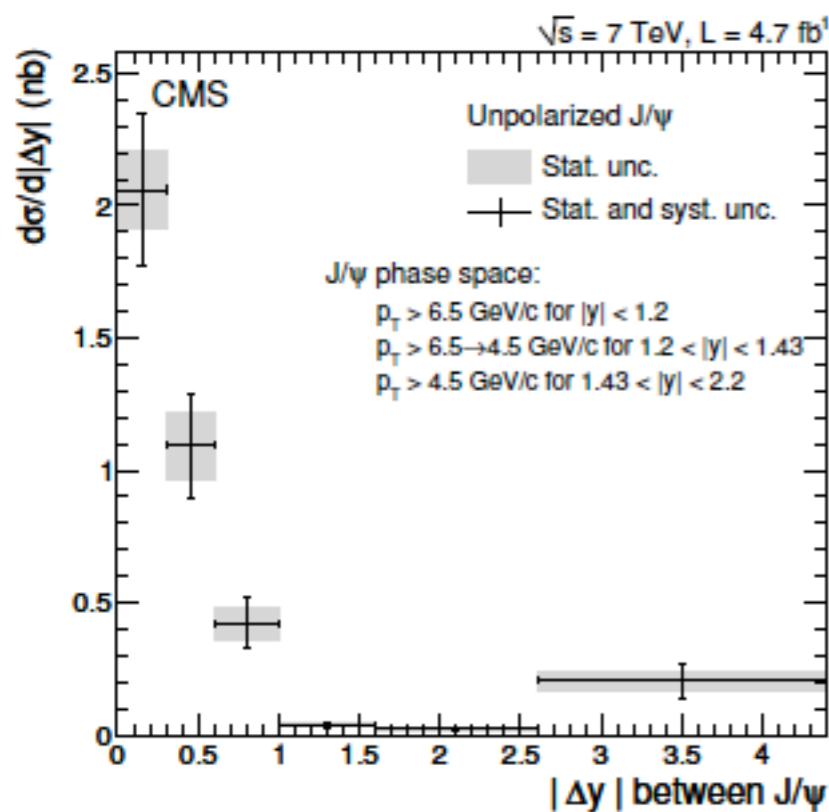
PUBLISHED: September 17, 2014

Measurement of prompt J/ψ pair production in pp collisions at $\sqrt{s} = 7$ TeV



The CMS collaboration

E-mail: cms-publication-committee-chair@cern.ch



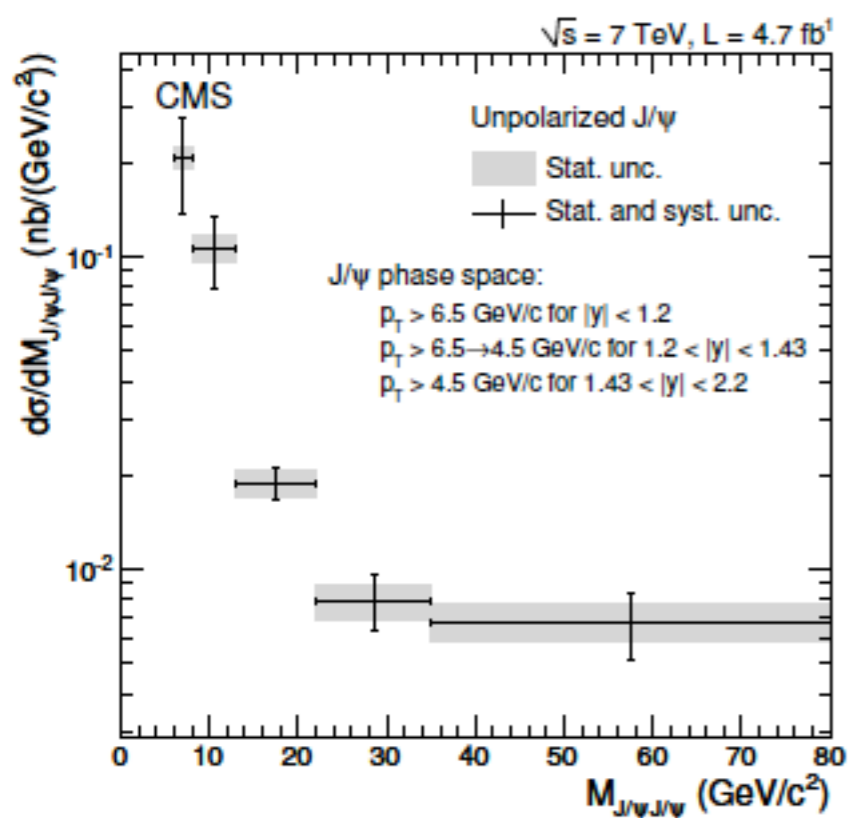
EVIDENCE OF DPS ?

Measurement of prompt J/ψ pair production in pp collisions at $\sqrt{s} = 7$ TeV

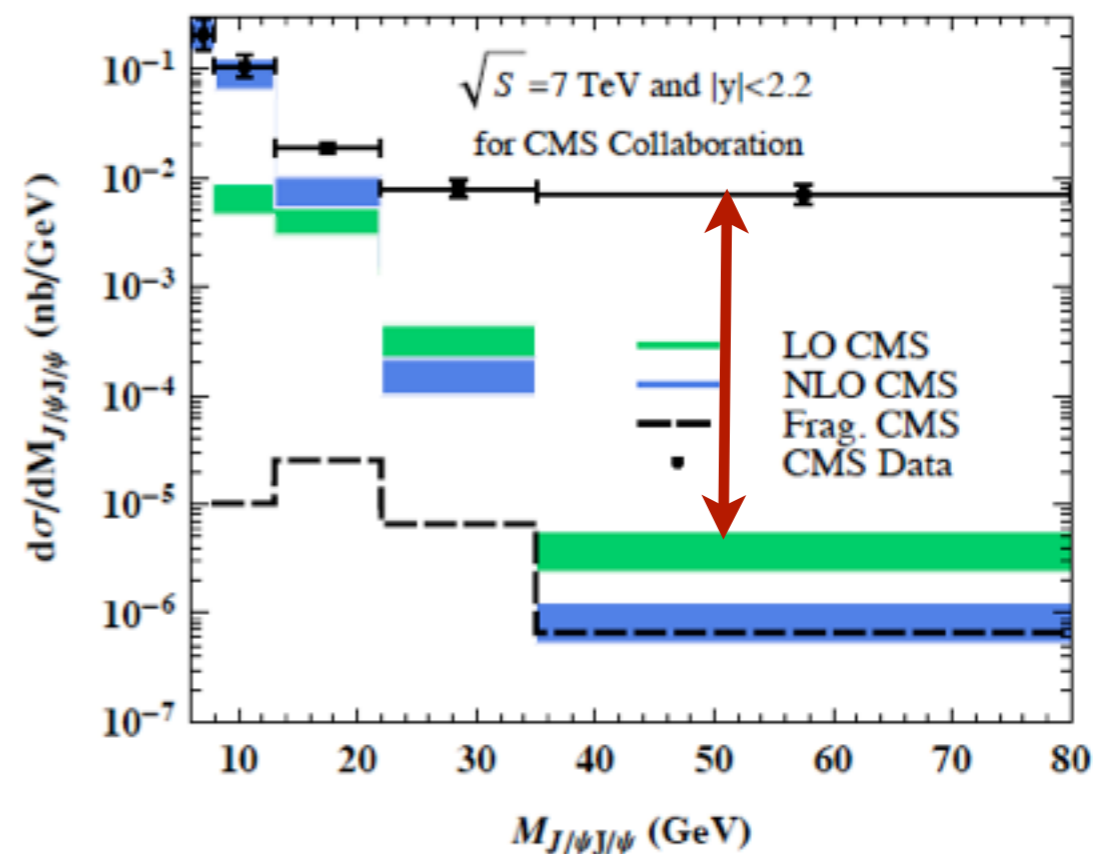


The CMS collaboration

E-mail: cms-publication-committee-chair@cern.ch



- Large discrepancy found with NLO-level Single-Parton Scatterings [Sun, Han, Chao, '14].



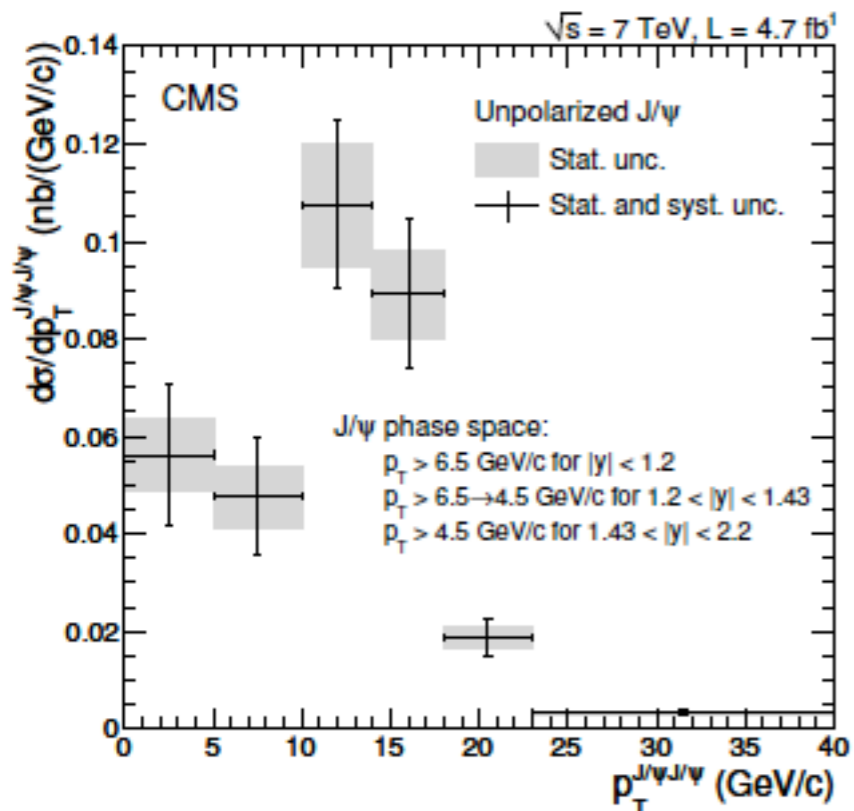
EVIDENCE OF DPS ?

Measurement of prompt J/ψ pair production in pp collisions at $\sqrt{s} = 7$ TeV

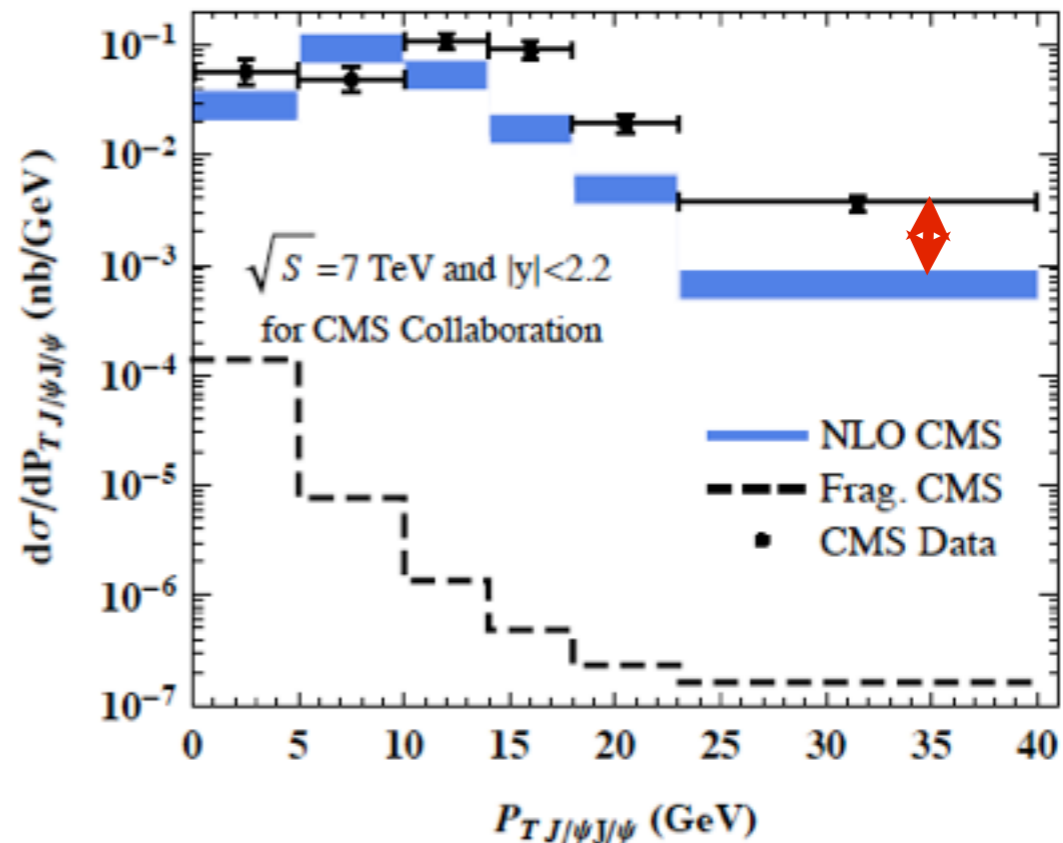


The CMS collaboration

E-mail: cms-publication-committee-chair@cern.ch

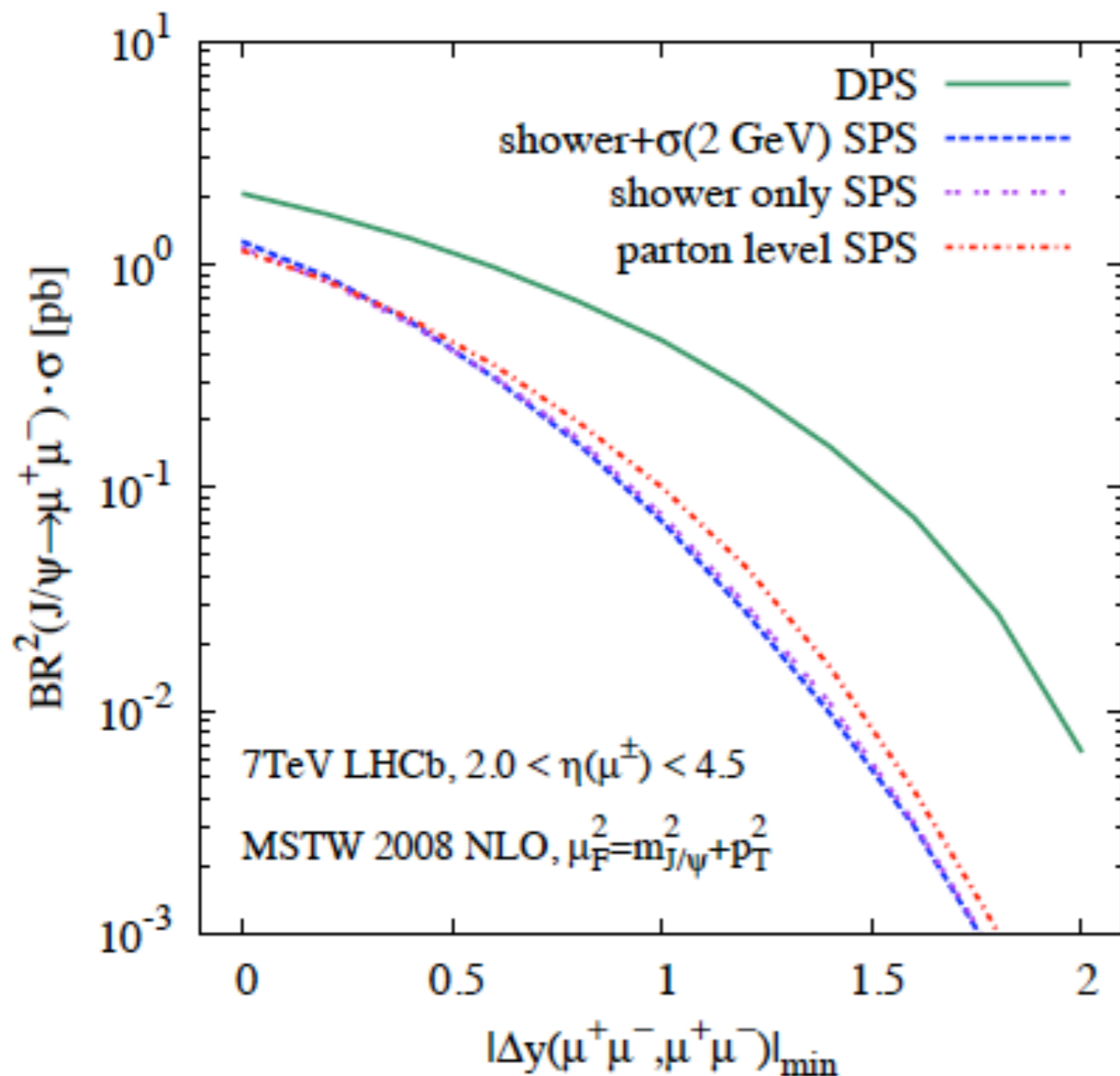


- Large discrepancy found with NLO-level Single-Parton Scatterings [Sun, Han, Chao, '14].



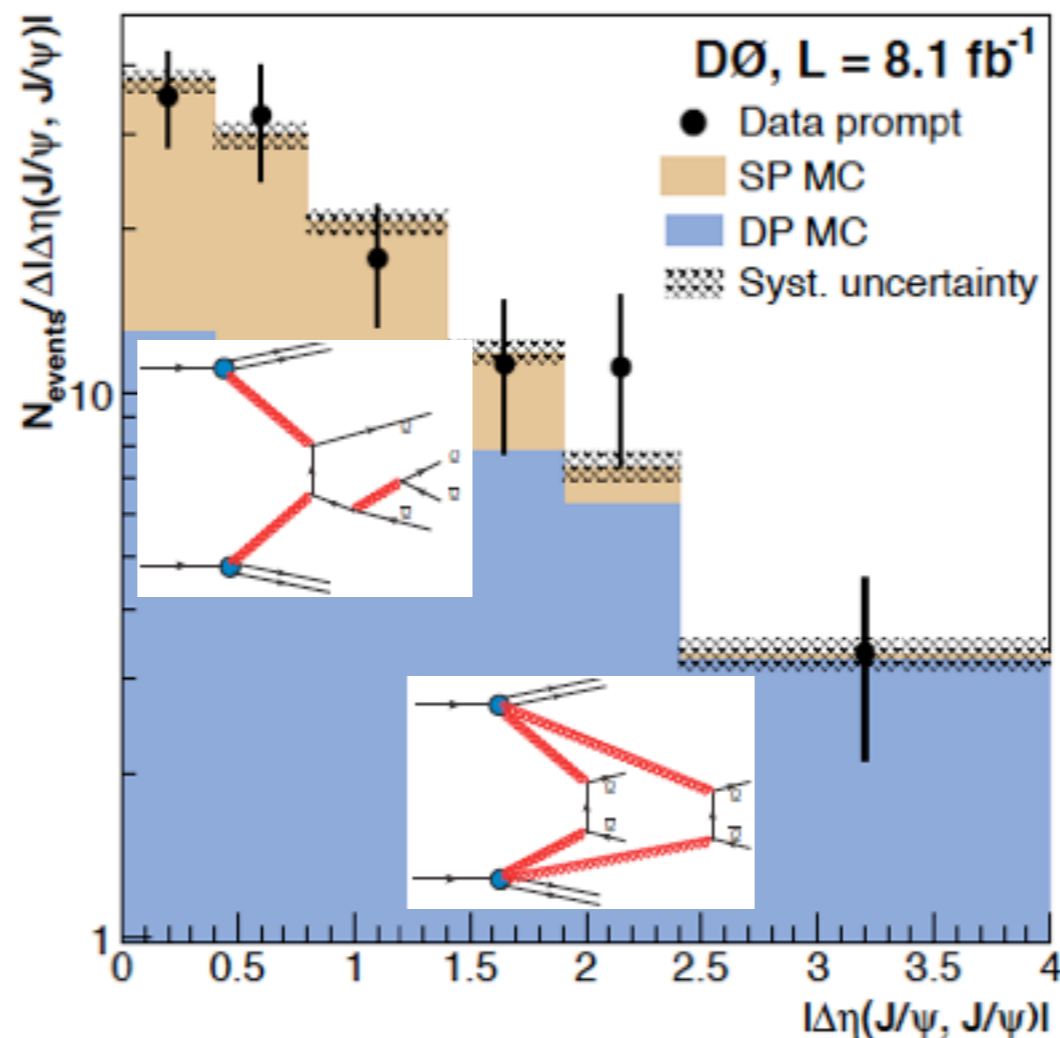
DPS IN DZERO MEASUREMENT

- It was proposed by Kom et al. (2011) rapidity difference can be a good observable to measure **DPS**, which is little dependent on shower and primordial kT smearing.



DPS IN DZERO MEASUREMENT

- It was proposed by Kom et al. (2011) rapidity difference can be a good observable to measure **DPS**, which is little dependent on shower and primordial kT smearing.
- D0 observed double J/ψ at Tevatron and separated **SPS** and **DPS** for the first time.



J/ψ -Pair Production at Large Momenta: Indications for Double-Parton Scatterings and Large α_s^5 Contributions

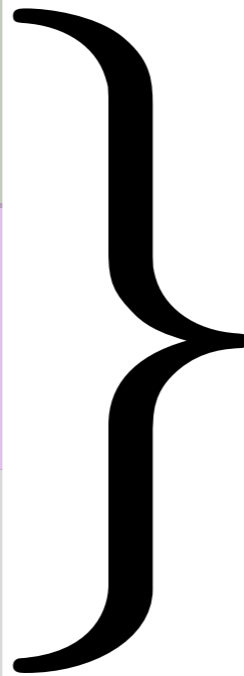
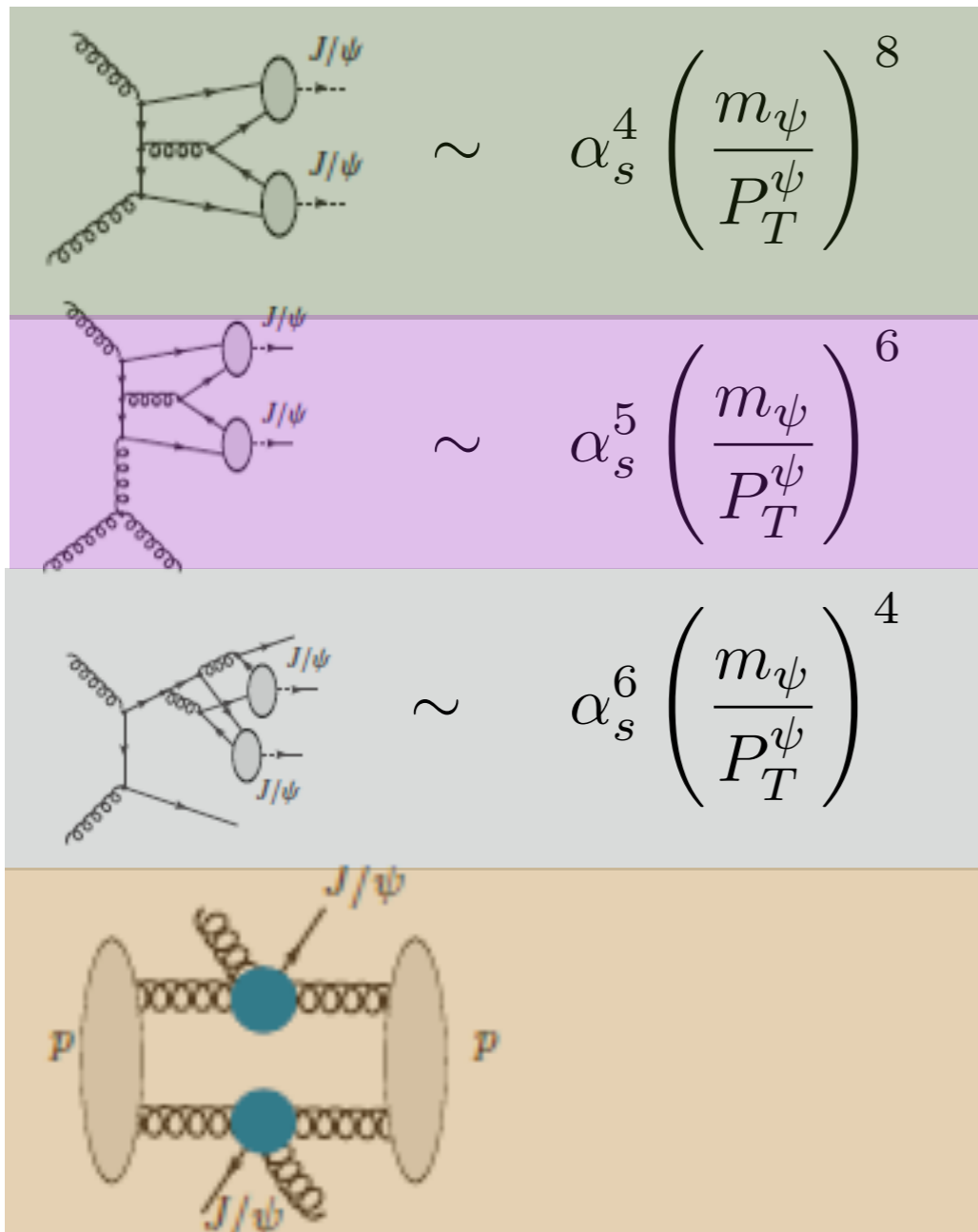
Lansberg, HSS PLB '14

Jean-Philippe Lansberg^a, Hua-Sheng Shao^{b,c}

^a IPNO, Université Paris-Sud, CNRS/IN2P3,

Department of Physics and State Key Laboratory of Nuclear Physics and

^c PH Department, TH Unit, CERN, CH-1211,



S
P
S

D
P
S

	$\sigma_{\text{exp.}}^{\text{prompt}}$	$\sigma_{\text{SPS}}^{\text{LO,direct}}$	$\sigma_{\text{SPS}}^{\text{NLO,direct}}$	$\sigma_{\text{SPS}}^{\text{NLO,prompt}}$	$\sigma_{\text{DPS}}^{\text{prompt}}$
LHCb	18 ± 5.3	$22^{+27.7}_{-13.1}$	$24.3^{+30.6}_{-14.4}$	$46.0^{+58.0}_{-27.3}$	$36.0^{+44.0}_{-12.8}$
D0	SPS: 70 ± 23	$28.9^{+30.7}_{-14.5}$	91^{+177}_{-55}	173^{+335}_{-105}	87^{+106}_{-31}
	DPS: 59 ± 23				
CMS	5.25 ± 0.52	$0.19^{+0.14}_{-0.09}$	$0.82^{+1.18}_{-0.46}$	$1.54^{+2.24}_{-0.87}$	$1.46^{+1.78}_{-0.52}$
ATLAS	N/A	$3.45^{+2.35}_{-1.40}$	$35.5^{+48.9}_{-19.8}$	$67.1^{+92.4}_{-37.6}$	$39.1^{+47.7}_{-13.9}$

TABLE I: $\sigma(pp(\bar{p}) \rightarrow J/\psi + J/\psi + X) \times \mathcal{B}_{\mu\mu}^2$ [Values in units of pb for LHCb and CMS and fb for D0 and ATLAS. The kinematical cuts are given as supplemental material.]

- Using the D0 data to fix the DPS parameter
- If one used the D0 data to fix the DPS yield, DPS and SPS are comparable in the CMS acceptance.
- Large p_T : CMS and D0 measurements imply a large DPS yield.
- Small p_T : LHCb data do NOT imply a large DPS yield

CALCULATION FRAMEWORK: DPS



- All calculations are performed by the general-purposed matrix-element/event generator *HELAC-Onia* [HSS, CPC '12,'15] with the correct spin-entangled decay.

- DPS has

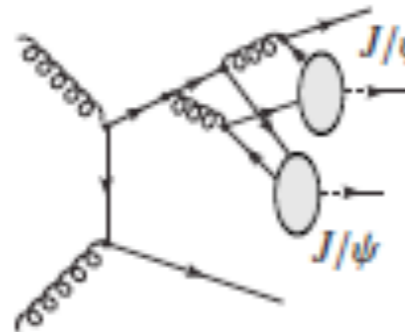
$$\sigma_{Q_1 Q_2} = \frac{1}{1 + \delta_{Q_1 Q_2}} \frac{\sigma_{Q_1} \sigma_{Q_2}}{\sigma_{\text{eff}}},$$

- Normally, σ_{eff} is thought to be universal, i.e. process&energy independent. **However, it is important to be tested ?**
- Since no satisfying solution to describe single-quarkonium production cross sections σ_{ψ} , we decide to use a data-driven way because a lot of single quarkonium data are available.
- By doing so, we assume the amplitude of single quarkonium production in the Crystal-ball function form [Kom et al. (2011)]

$$\overline{|\mathcal{A}_{gg \rightarrow Q+X}|^2} = \begin{cases} K \exp(-\kappa \frac{P_T^2}{M_Q^2}) & \text{when } P_T \leq \langle P_T \rangle \\ K \exp(-\kappa \frac{\langle P_T \rangle^2}{M_Q^2}) \left(1 + \frac{\kappa}{n} \frac{P_T^2 - \langle P_T \rangle^2}{M_Q^2}\right)^{-n} & \text{when } P_T > \langle P_T \rangle \end{cases} \quad K = \lambda^2 \kappa \hat{s} / M_Q^2.$$

OTHER CONTRIBUTIONS: SPS

- Beyond NLO contributions (new fragmentation topology):

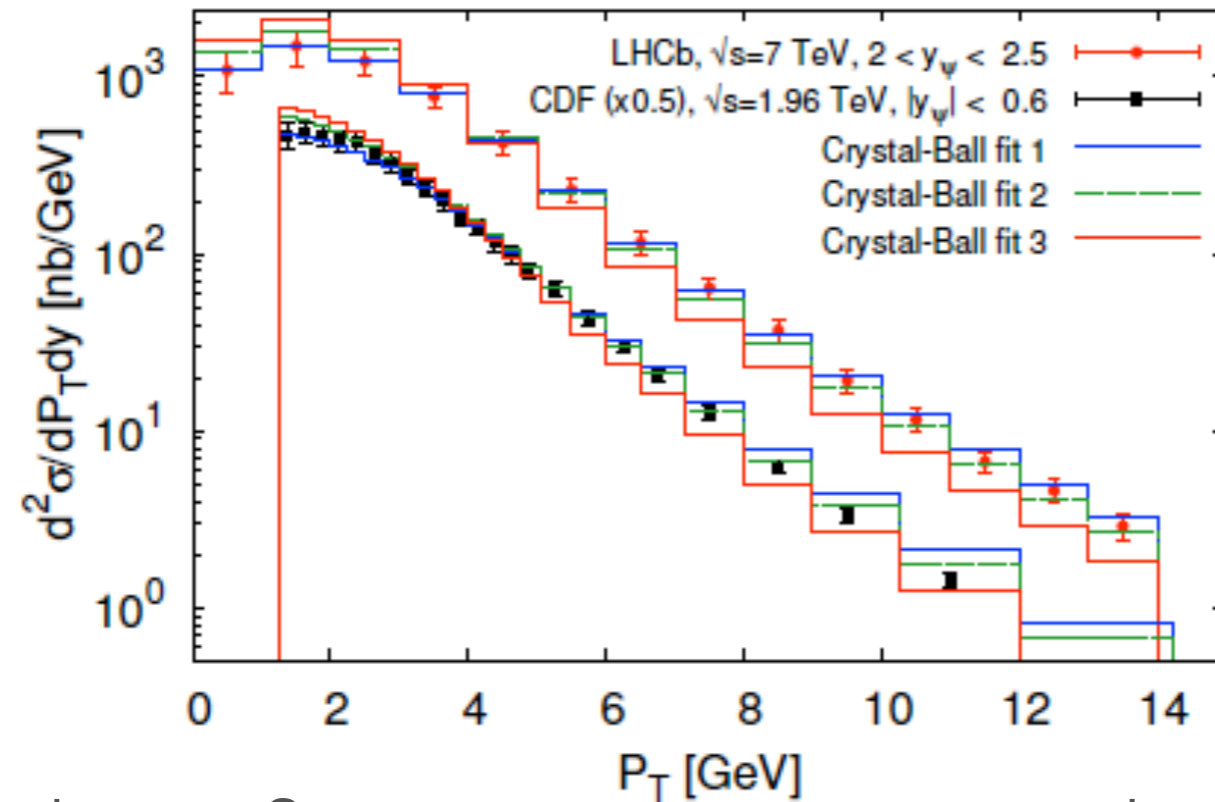

$$\sim \alpha_s^5 \left(\frac{m_\psi}{P_T^\psi} \right)^4$$

- Feeddown from $J/\psi + \psi(2S)$ contributes 46% (i.e. 85% of direct), while others like $J/\psi + \chi_c$ are suppressed.
- CO contributions are also suppressed because of either smallness of CO LDMEs or no p_T -enhanced diagrams.
- In the accessible region, CO to SPS never dominates compared to CS SPS + DPS.

CALCULATION FRAMEWORK: DPS



- Single- J/ψ cross sections input from fits of existing data



- We used three fits to assess systematical uncertainties.
- Together with σ_{eff} , they allow to predict σ_{DPS} .
- Our strategy is therefore to fit σ_{eff} from CMS data via $\sigma_{\text{SPS}} + \sigma_{\text{DPS}}$.

FITTING SIGMA_EFF FROM CMS J/PSI-PAIR DATA

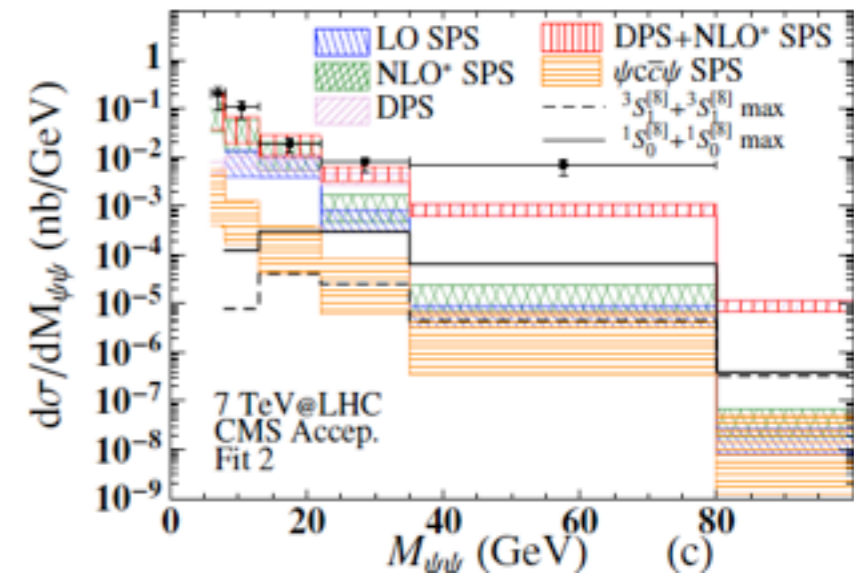
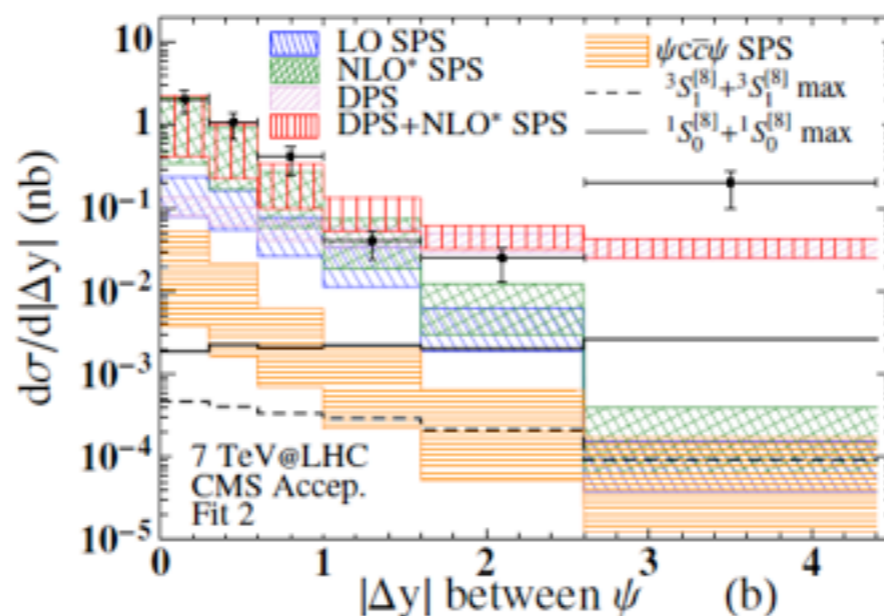
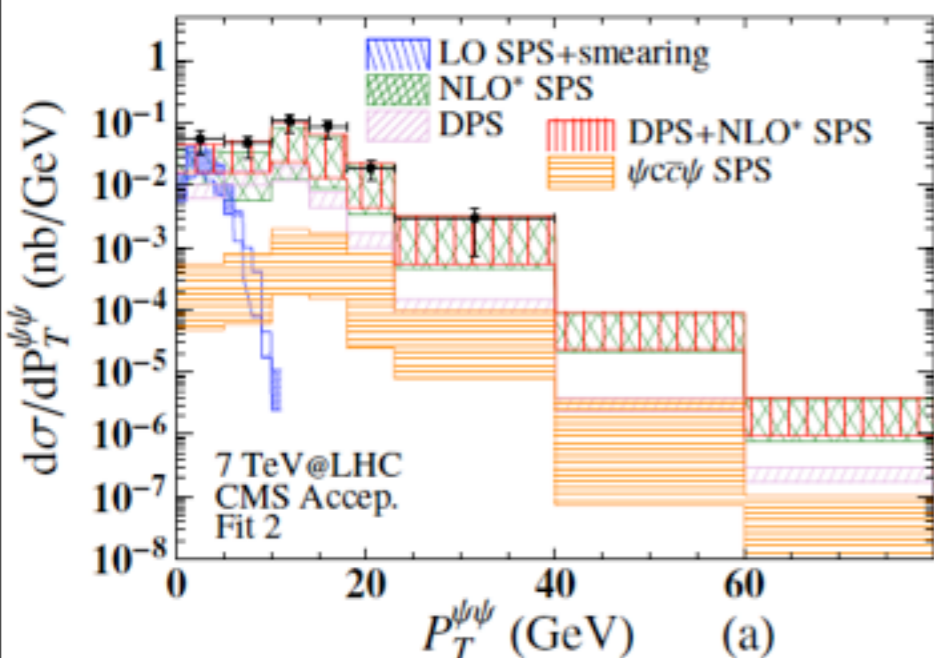


- $p_T^{\psi\psi}$, $|\Delta y_{\psi\psi}|$ & $M_{\psi\psi}$ distributions are fitted

	σ_{eff} [mb]	$\chi^2_{\text{d.o.f.}}$	d.o.f.
σ_ψ Fit 1 [25]	11 ± 2.9	1.9	16
σ_ψ Fit 2	8.2 ± 2.2	1.8	16
σ_ψ Fit 3	5.3 ± 1.4	1.9	16
Only LO SPS	N/A	7.6	17
Only NLO* SPS	N/A	2.6	17

Table 2: Result of the fit of the DPS yield via σ_{eff} on the 18 CMS values.

- Clear need for DPS (LO and NLO* SPS are not sufficient)



OUR EXTRACTION OF SIGMA_EFF



- Combining our three fits, we obtain

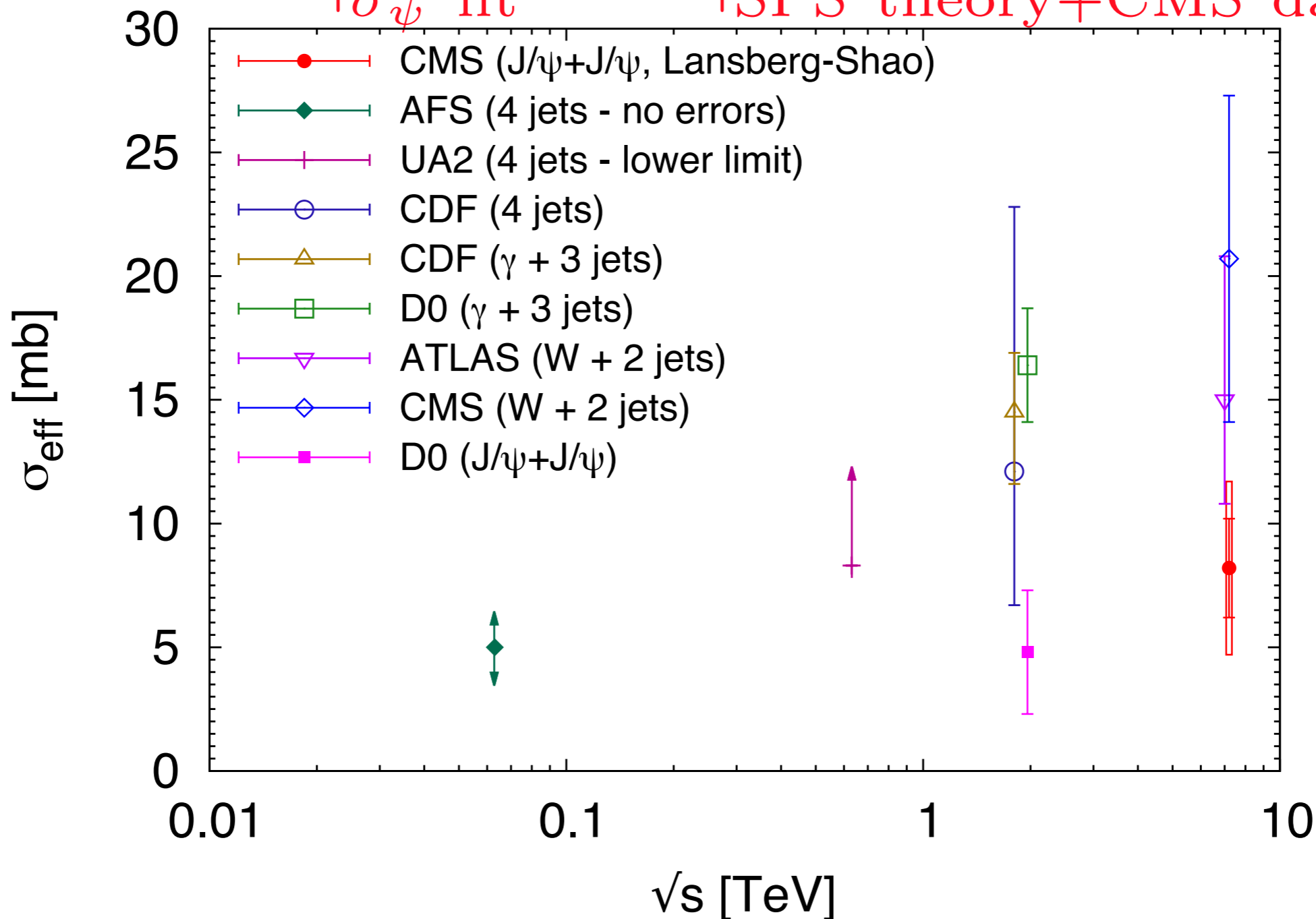
$$\sigma_{\text{eff}} = 8.2 \pm 2.9 |_{\sigma_{\psi} \text{ fit}} \pm 2.0 |_{\text{SPS theory+CMS data}} \text{ mb}$$

- The SPS theory uncertainty can in principle be removed by measuring a DPS cross section (as done by D0).
- The CMS data uncertainty can be reduced with more double quarkonium data.
- The last uncertainty is of course more tricky to deal with.

OUR EXTRACTION OF SIGMA_EFF

- Combining our three fits, we obtain

$$\sigma_{\text{eff}} = 8.2 \pm 2.9 |_{\sigma_{\psi} \text{ fit}} \pm 2.0 |_{\text{SPS theory} + \text{CMS data}} \text{ mb}$$



OUR EXTRACTION OF SIGMA_EFF



- Combining our three fits, we obtain

$$\sigma_{\text{eff}} = 8.2 \pm 2.9 |_{\sigma_{\psi} \text{ fit}} \pm 2.0 |_{\text{SPS theory+CMS data}} \text{ mb}$$

- The SPS theory uncertainty can in principle be removed by measuring a DPS cross section (as done by D0).
- The CMS data uncertainty can be reduced with more double quarkonium data.
- The last uncertainty is of course more tricky to deal with.
- Our extraction is compatible with that of D0.
- Both point at a small σ_{eff} compared to jet-related extraction.
- Does a smaller scale mean a smaller σ_{eff} ?
- Does gluon-induced process mean a smaller σ_{eff} ?

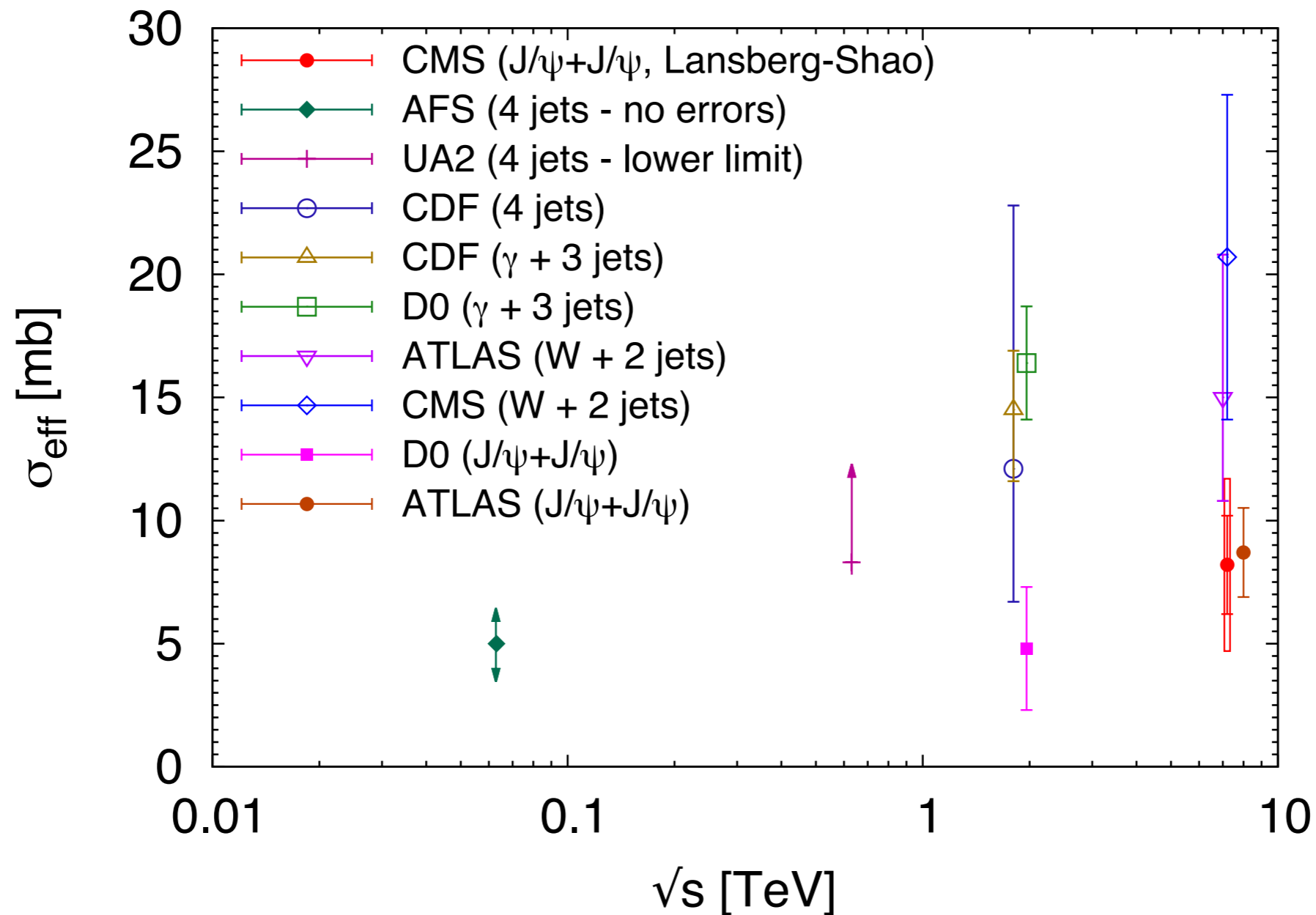
RECENT ATLAS MEASUREMENT



ATLAS-CONF-2016-047

- With a data-driven way, ATLAS extracted the DPS with 8 TeV data

$$\sigma_{\text{eff}} = 8.7 \pm 1.1(\text{stat}) \pm 1.4(\text{syst}) \pm 0.1(\text{BF}) \pm 0.3(\text{lumi}) \text{ mb}$$

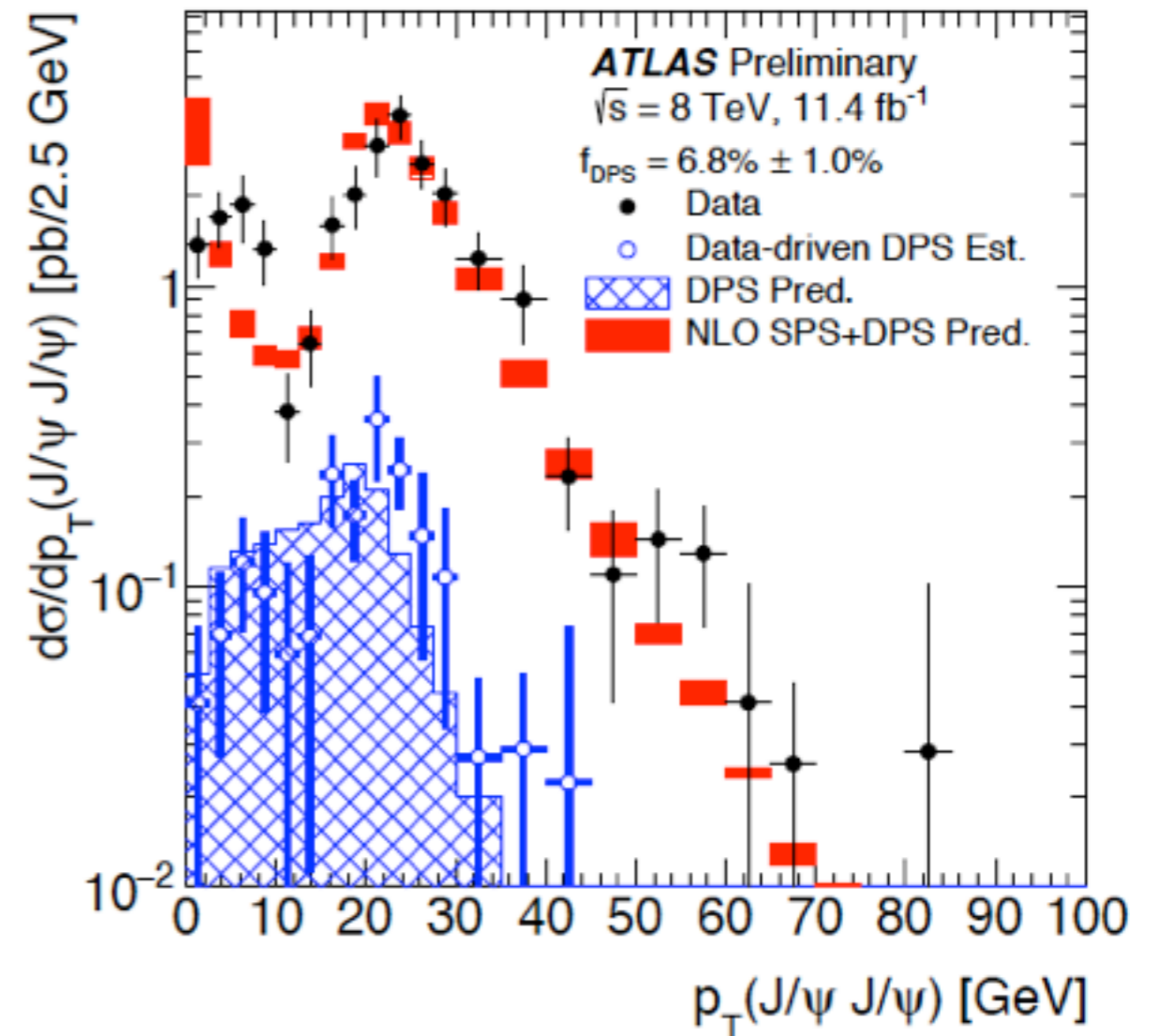
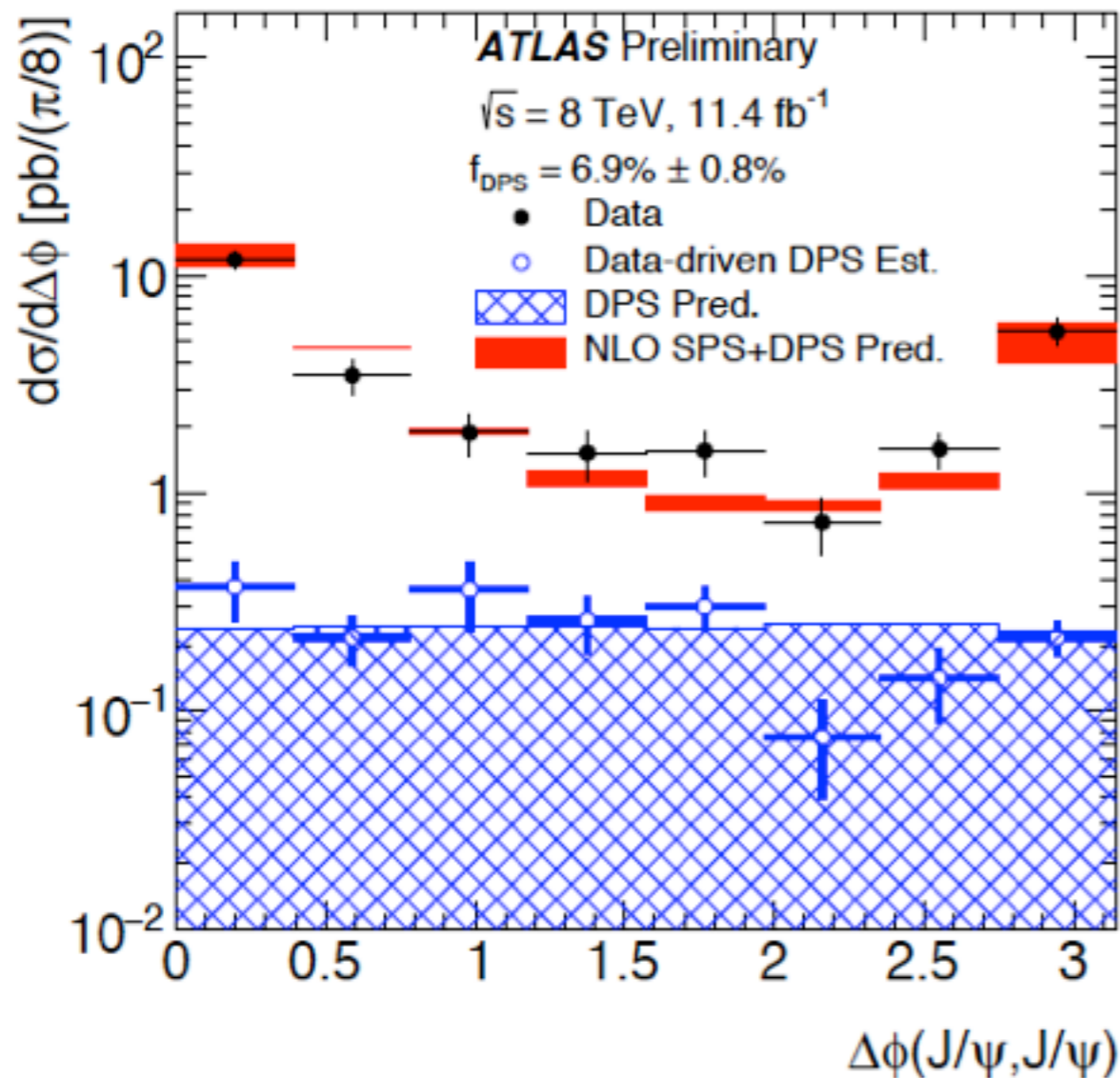


RECENT ATLAS MEASUREMENT



ATLAS-CONF-2016-047

- With a data-driven way, ATLAS extracted the DPS with 8 TeV data
 $\sigma_{\text{eff}} = 8.7 \pm 1.1(\text{stat}) \pm 1.4(\text{syst}) \pm 0.1(\text{BF}) \pm 0.3(\text{lumi}) \text{ mb}$
- It also confirms the CS dominant in the **SPS** of this process

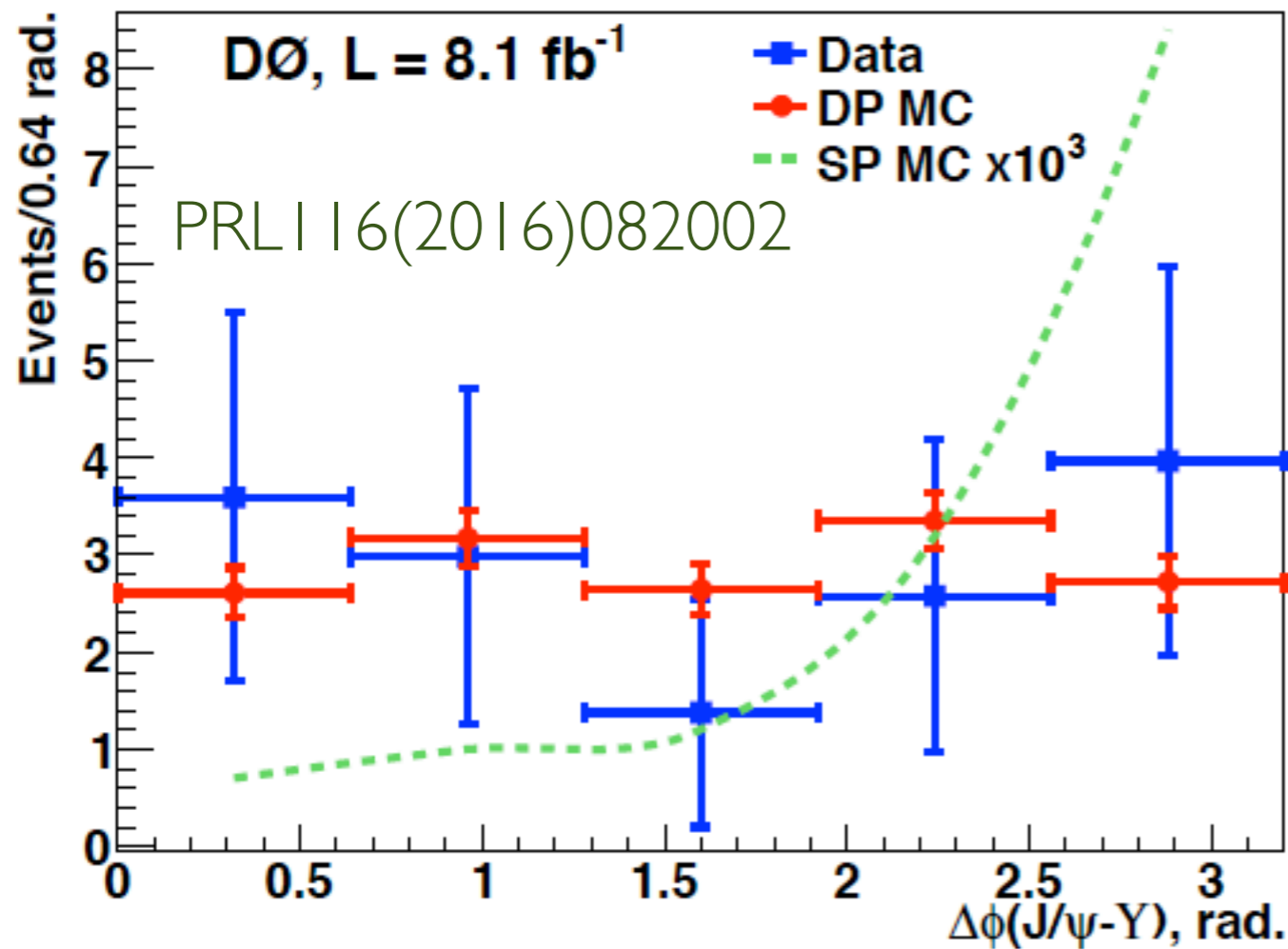




LPTHE
LABORATOIRE DE PHYSIQUE
THEORIQUE ET HAUTES ENERGIES

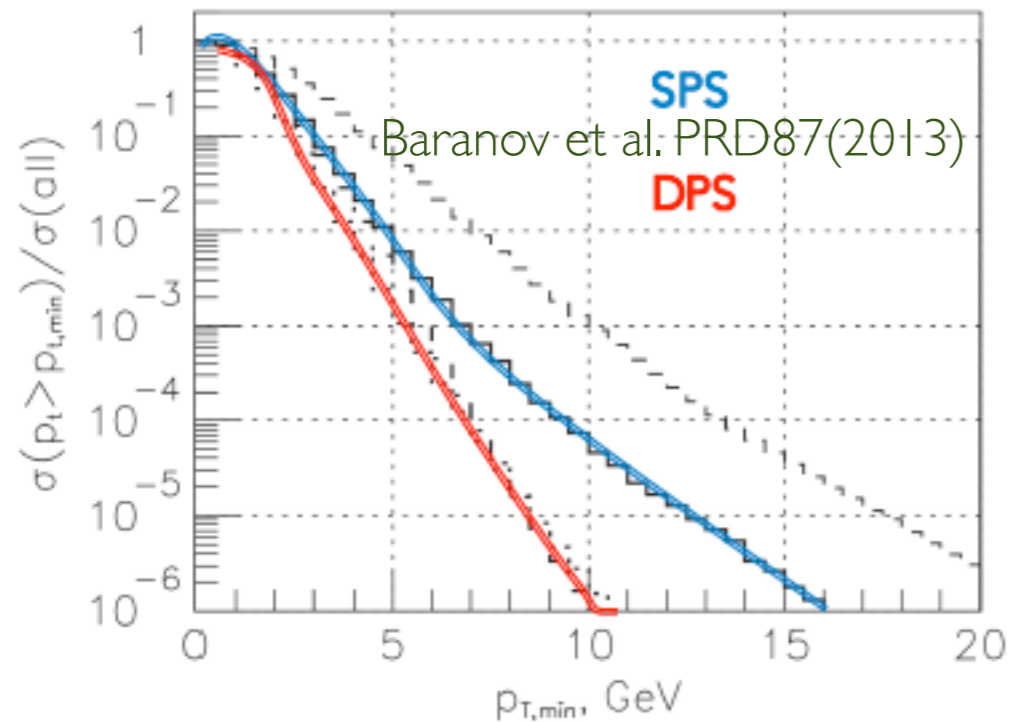
$$J/\psi + \gamma$$

RECENT DZERO MEASUREMENT: J/ ψ +Y



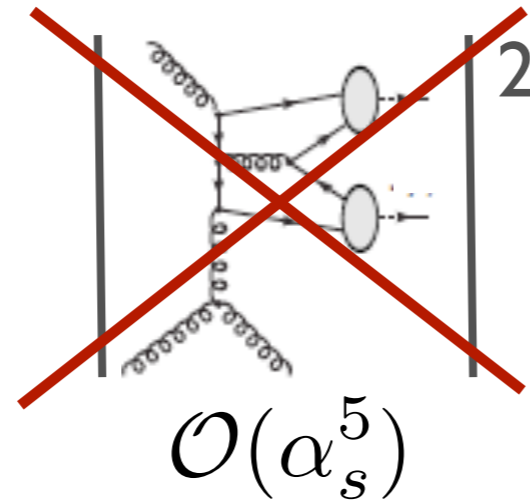
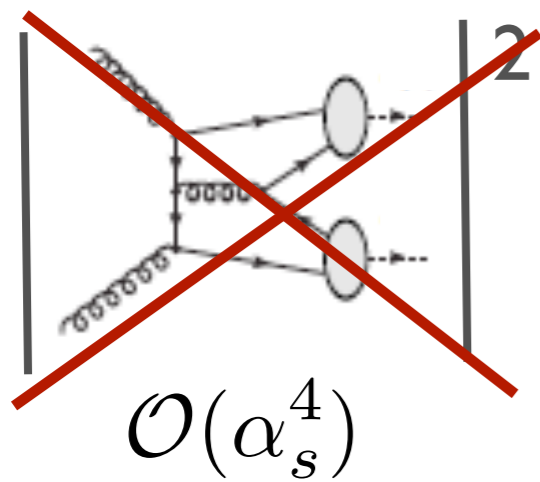
$$\sigma_{\text{eff}} = 2.2 \pm 0.7 \text{ (stat)} \pm 0.9 \text{ (syst)} \text{ mb.}$$

- Does it indicate a significant **SPS** contribution ?
- The value of σ_{eff} is a little bit too small if one assumes all contributions are **DPS** ?
- Would LHC measurements (especially LHCb) give a clarification in the future ?
- Is **SPS** can be completely negligible (even at low p_T) ?



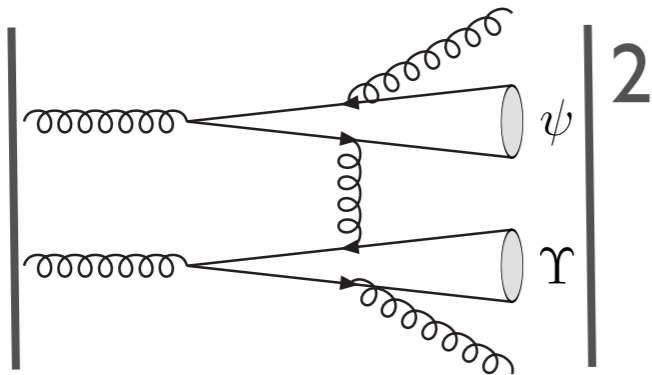
PROMPT J/PSI+Y: SPS

- In pQCD, there are **NO** $\mathcal{O}(\alpha_s^4)$ and $\mathcal{O}(\alpha_s^5)$ **CS** contributions



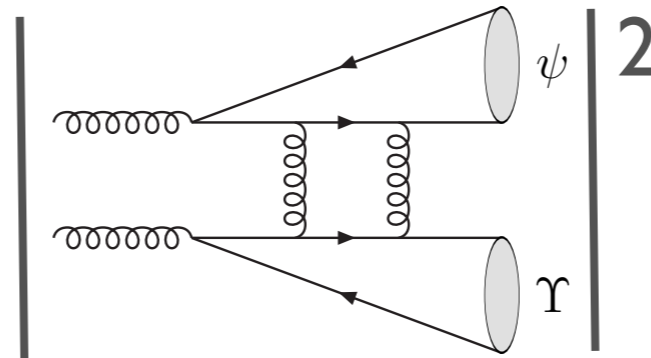
PROMPT J/PSI+Y: SPS

- In pQCD, there are **NO** $\mathcal{O}(\alpha_s^4)$ and $\mathcal{O}(\alpha_s^5)$ **CS** contributions
- The LO in pQCD is $\mathcal{O}(\alpha_s^6)$



$$\mathcal{O}(\alpha_s^6)$$

Double **R**real

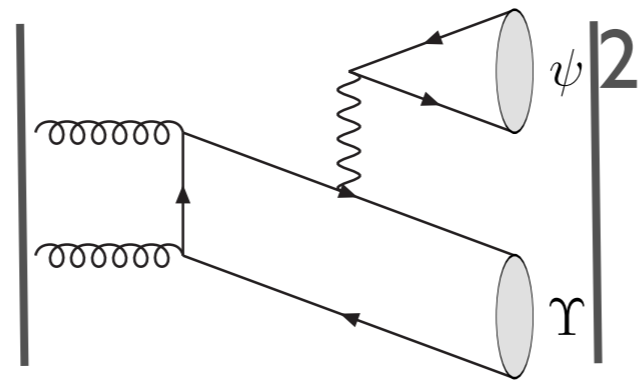


$$\mathcal{O}(\alpha_s^6)$$

Loop **I**nduced

PROMPT J/PSI+Y: SPS

- In pQCD, there are **NO** $\mathcal{O}(\alpha_s^4)$ and $\mathcal{O}(\alpha_s^5)$ **CS** contributions
- The LO in pQCD is $\mathcal{O}(\alpha_s^6)$
- There is also **EW** contribution $\mathcal{O}(\alpha_s^2 \alpha^2)$

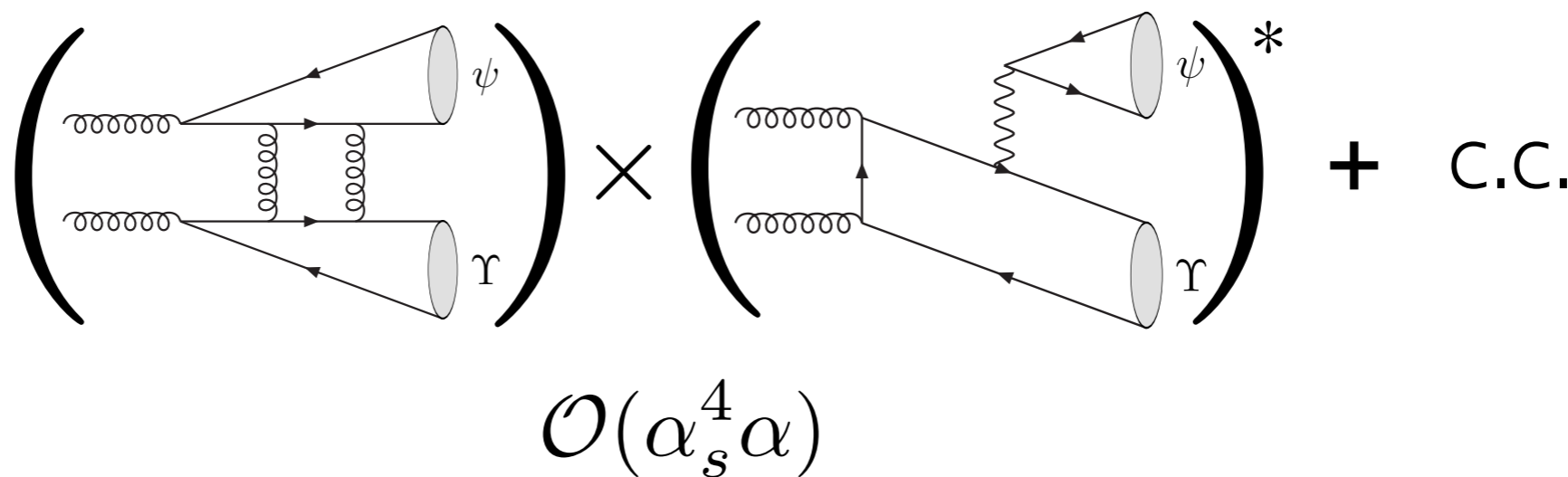


$$\mathcal{O}(\alpha_s^2 \alpha^2)$$

ElectroWeak

PROMPT J/PSI+Y: SPS

- In pQCD, there are **NO** $\mathcal{O}(\alpha_s^4)$ and $\mathcal{O}(\alpha_s^5)$ **CS** contributions
- The LO in pQCD is $\mathcal{O}(\alpha_s^6)$
- There is also **EW** contribution $\mathcal{O}(\alpha_s^2\alpha^2)$
- **LI** and **EW** can have interference $\mathcal{O}(\alpha_s^4\alpha)$



The diagram shows the interference between a LO pQCD contribution and an EW contribution. The LO pQCD part is represented by a large bracketed term containing a diagram with two incoming gluons and two outgoing particles, ψ and γ . The EW part is represented by a large bracketed term containing a diagram with two incoming gluons, a loop of a fermion, and two outgoing particles, ψ and γ . The two terms are multiplied together, and the result is added to its complex conjugate (C.C.).

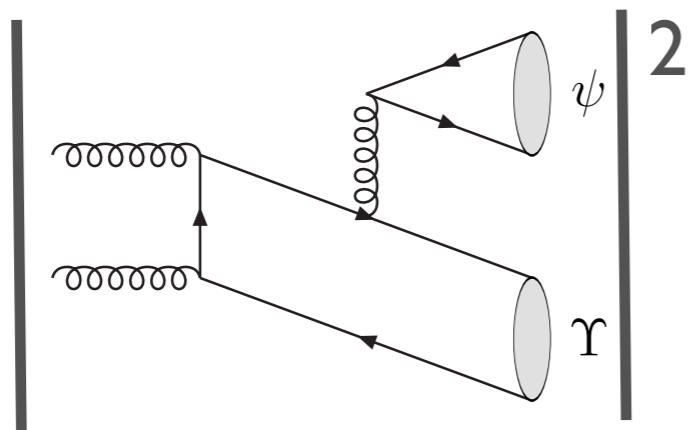
$$\left(\text{LO pQCD} \right) \times \left(\text{EW} \right)^* + \text{C.C.}$$

$\mathcal{O}(\alpha_s^4\alpha)$

INTERference

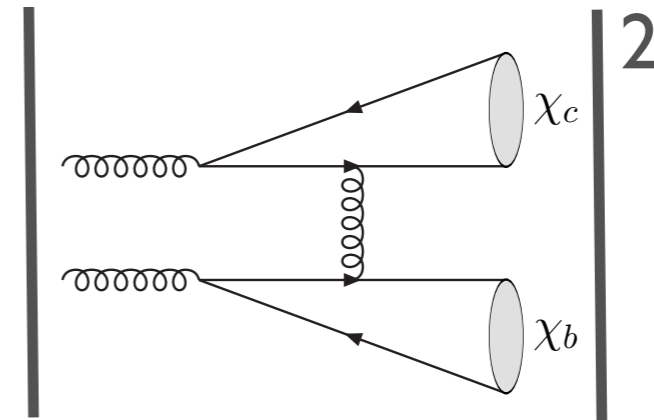
PROMPT J/PSI+Y: SPS

- In pQCD, there are **NO** $\mathcal{O}(\alpha_s^4)$ and $\mathcal{O}(\alpha_s^5)$ **CS** contributions
- The LO in pQCD is $\mathcal{O}(\alpha_s^6)$
- There is also **EW** contribution $\mathcal{O}(\alpha_s^2\alpha^2)$
- **LI** and **EW** can have interference $\mathcal{O}(\alpha_s^4\alpha)$
- **CO** contributes at $\mathcal{O}(\alpha_s^4 v_c^i v_b^j)$, $i + j \geq 4$



$$\mathcal{O}(\alpha_s^4 v_c^i v_b^j), i + j \geq 4$$

Color-**O**ctet **M**echanism



Feeddown

PROMPT J/PSI+Y: SPS



HSS, Zhang PRL'16

- In pQCD, there are **NO** $\mathcal{O}(\alpha_s^4)$ and $\mathcal{O}(\alpha_s^5)$ **CS** contributions
- The LO in pQCD is $\mathcal{O}(\alpha_s^6)$
- There is also **EW** contribution $\mathcal{O}(\alpha_s^2\alpha^2)$
- **LI** and **EW** can have interference $\mathcal{O}(\alpha_s^4\alpha)$
- **CO** contributes at $\mathcal{O}(\alpha_s^4 v_c^i v_b^j), i + j \geq 4$

Label	HELAC-ONIA 2.0 syntax	First order	Description
DR	<code>g g > cc~(3S11) bb~(3S11) g g</code>	$\mathcal{O}(\alpha_s^6)$	Double Real (DR) CS contribution
LI	<code>addon 8</code>	$\mathcal{O}(\alpha_s^6)$	Loop-Induced (LI) CS contribution
EW	<code>p p > cc~(3S11) bb~(3S11)</code>	$\mathcal{O}(\alpha_s^2\alpha^2)$	ElectroWeak (EW) CS contribution
INTER	<code>addon 8</code>	$\mathcal{O}(\alpha_s^4\alpha)$	INTERference (INTER) between LI and EW
COM	<code>g g > jpsi y(1s)</code>	$\mathcal{O}(\alpha_s^4 v_c^i v_b^j), i + j \geq 4$	CO $\mathcal{O}(\alpha_s^4)$ contribution

Similar leading contributions

$$\mathcal{O}(\alpha_s^6) \approx \mathcal{O}(\alpha_s^2\alpha^2) \approx \mathcal{O}(\alpha_s^4\alpha) \approx \mathcal{O}(\alpha_s^4 v_c^i v_b^j)$$

$$\text{Giving } \alpha_s \approx \sqrt{\alpha} \approx v_c^2 \approx v_b^2$$

PROMPT J/PSI+Y: SPS VS DATA



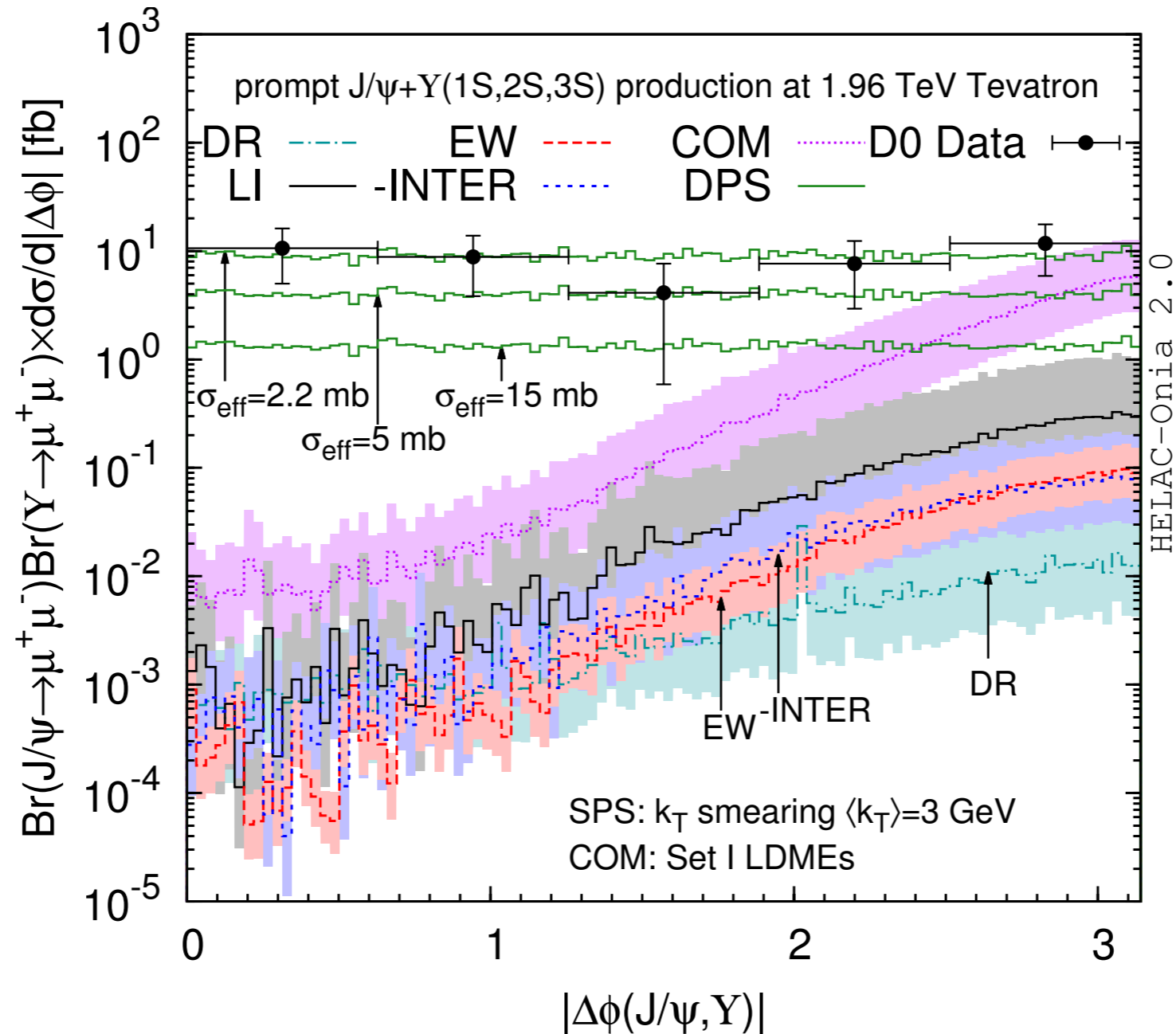
HSS, Zhang PRL '16

exp.	CSM				COM			
	DR	LI	EW	INTER	Set I	Set II	Set III	Set IV
D0: $27 \pm 42.2\%$	$0.0146^{+233\%}_{-66.6\%}$	$0.229^{+264\%}_{-70.4\%}$	$0.065^{+75.5\%}_{-46.6\%}$	$-0.068^{+162\%}_{-62.2\%}$	$2.96^{+135\%}_{-56.2\%}$	$1.41^{+160\%}_{-77.6\%}$	$1.80^{+143\%}_{-58.0\%}$	$0.418^{+144\%}_{-58.3\%}$
LHCb	$0.255^{+391\%}_{-79.7\%}$	$6.05^{+436\%}_{-82.2\%}$	$1.71^{+135\%}_{-65.2\%}$	$-3.23^{+262\%}_{-75.9\%}$	$38.8^{+238\%}_{-73.0\%}$	$21.2^{+243\%}_{-73.6\%}$	$28.1^{+243\%}_{-73.8\%}$	$6.57^{+243\%}_{-73.9\%}$

- SPS is smaller than the central value of D0 data
- SPS is NOT completely negligible
- COM SPS is strongly dependent on CO LDMEs
 - Set I: [Kramer, Prog.Part.Nucl.Phys. 47 \(2001\) 141](#)
 - Set II: [Sharma, Vitev, Phys. Rev. C 87 \(2013\) 044905](#)
 - Set III: [HSS et al, JHEP 1505 \(2015\) 103; Han et al, arXiv:1410.8537 \[hep-ph\]](#)
 - Set IV: [Gong et al, Phys. Rev. Lett. 110 \(2013\) 042002; Feng et al, Chin. Phys. C 39 \(2015\) 123102](#)

PROMPT J/ ψ +Y: THEORY VS DATA

- With SPS, we can derive $\sigma_{\text{eff}} \leq 8.2 \text{ mb}$ at 68% C.L.



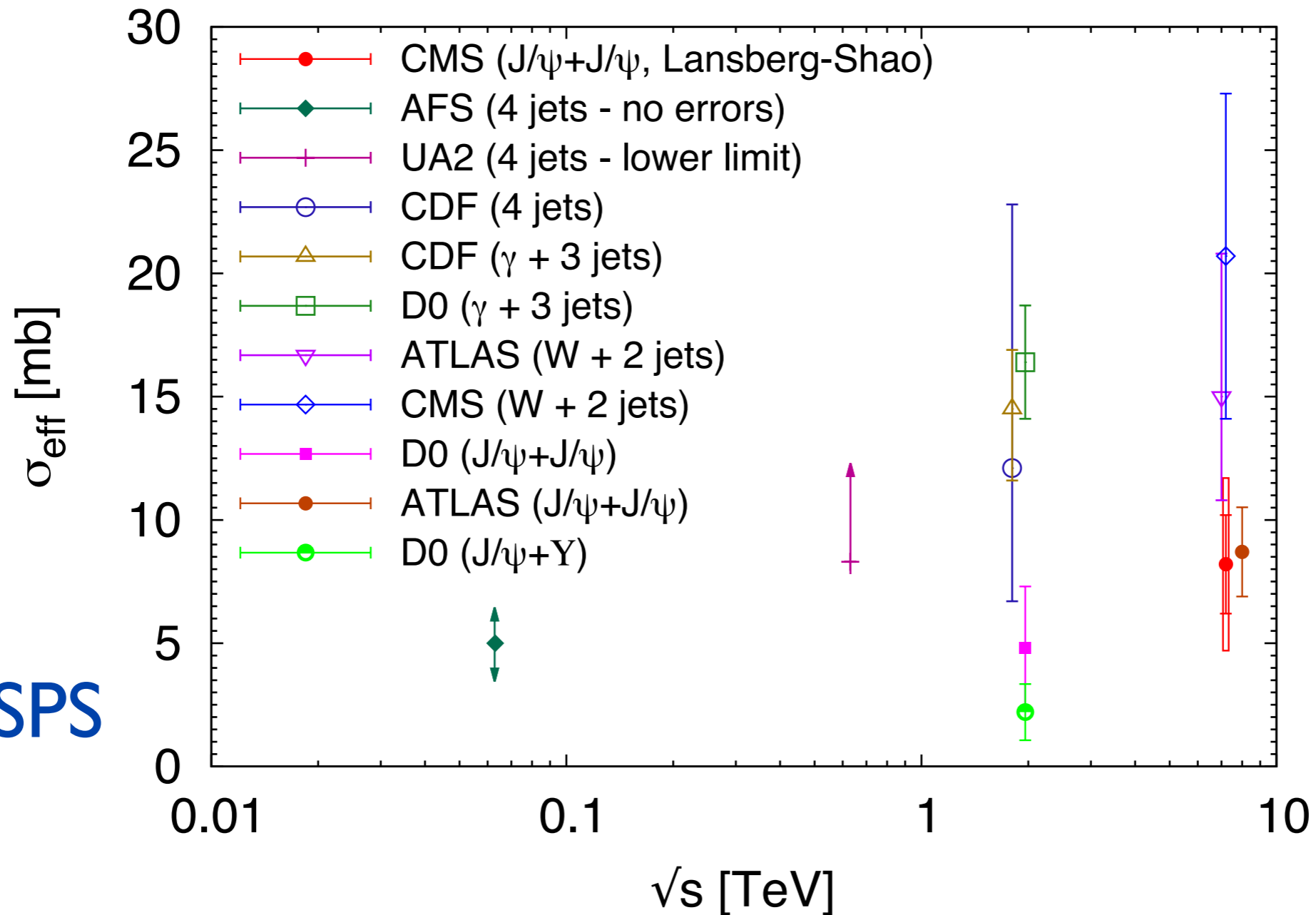
PROMPT J/ ψ +Y: THEORY VS DATA



HSS, Zhang PRL'16

- With SPS, we can derive $\sigma_{\text{eff}} \leq 8.2 \text{ mb}$ at 68% C.L.

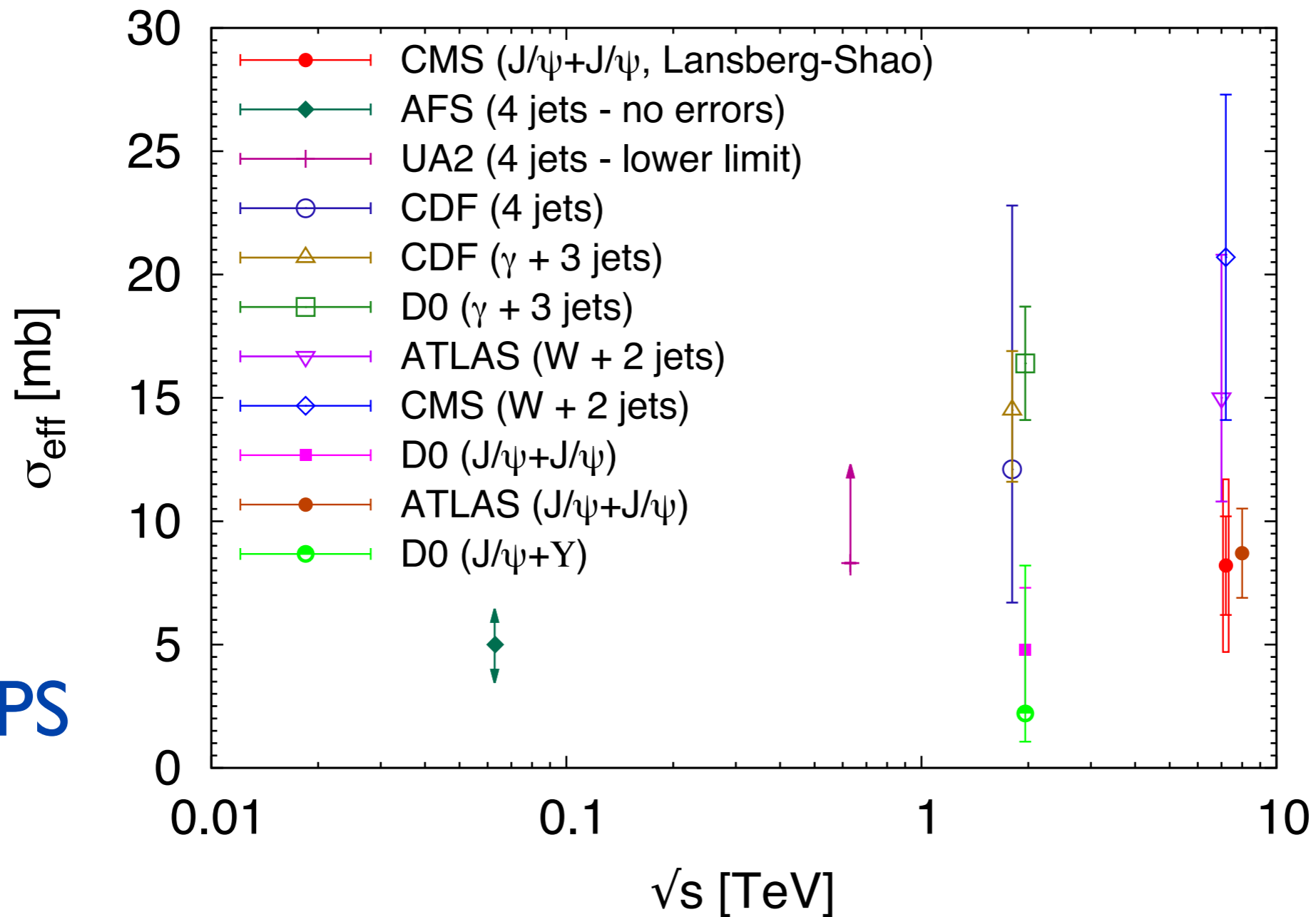
w/o SPS



PROMPT J/ ψ + Υ : THEORY VS DATA

- With SPS, we can derive $\sigma_{\text{eff}} \leq 8.2 \text{ mb}$ at 68% C.L.

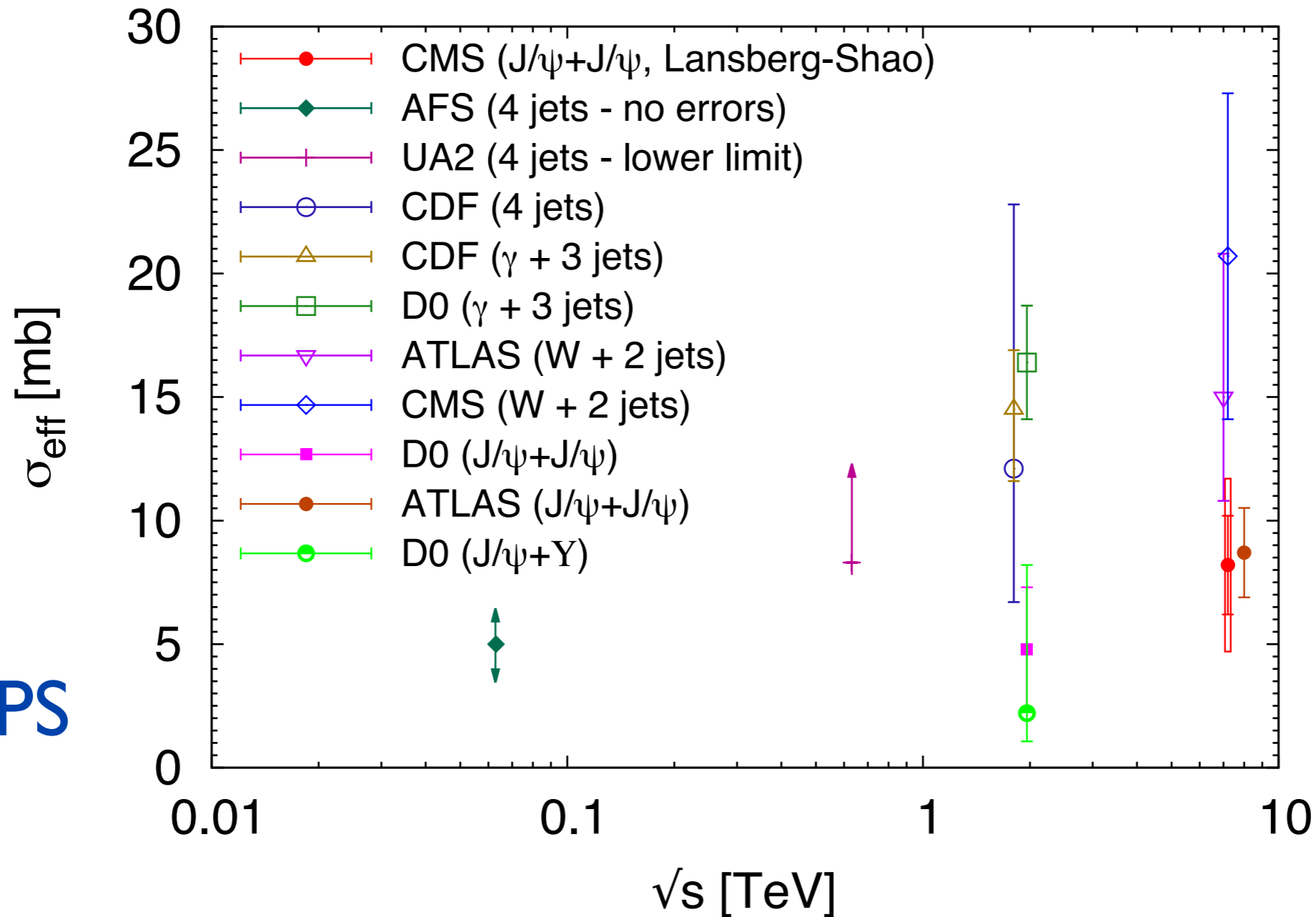
w SPS



PROMPT J/ ψ + Υ : THEORY VS DATA

- With SPS, we can derive $\sigma_{\text{eff}} \leq 8.2 \text{ mb}$ at 68% C.L.

w SPS



- Future measurement at the LHC ?



LPTHE
LABORATOIRE DE PHYSIQUE
THEORIQUE ET HAUTES ENERGIES

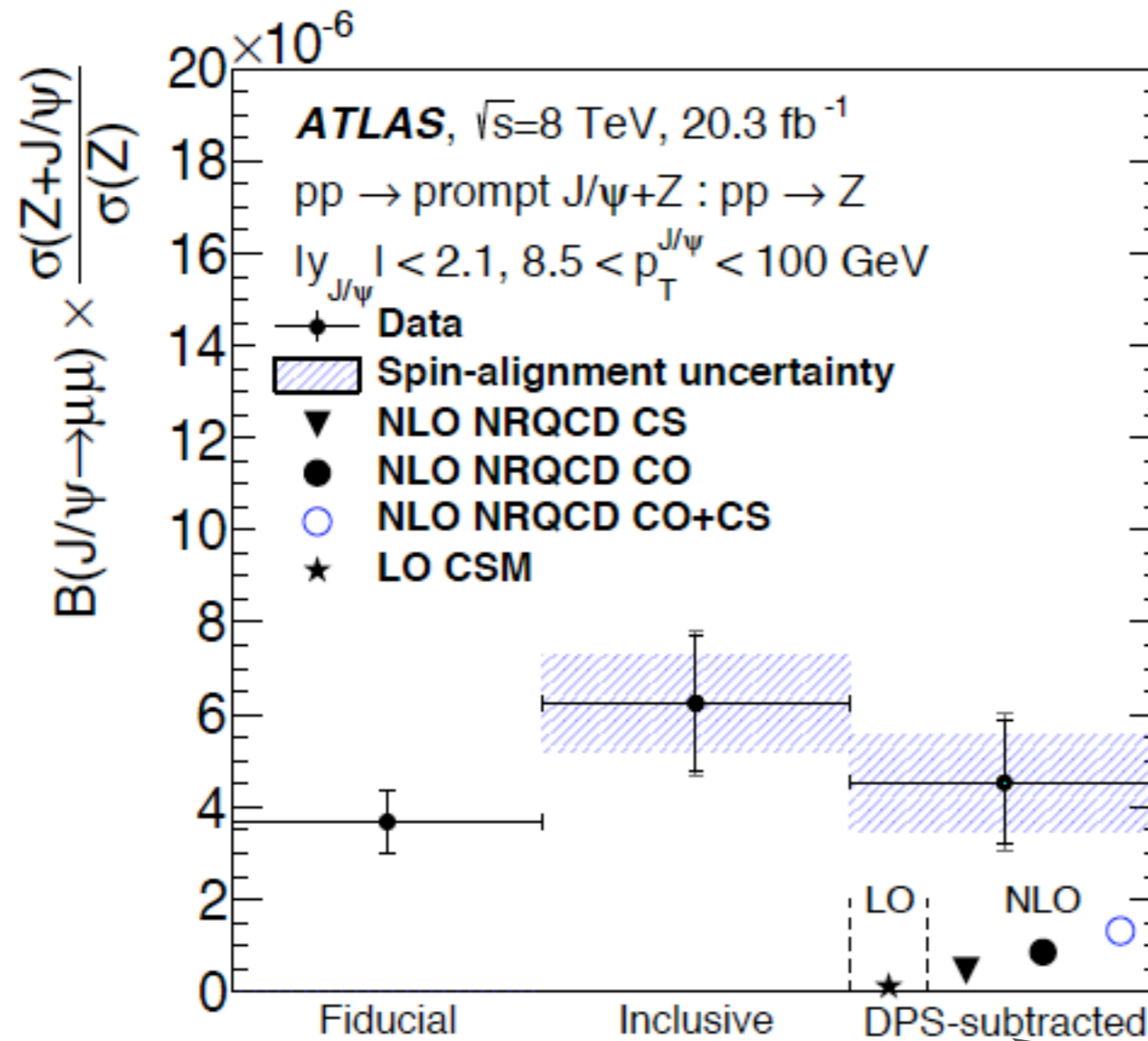
$$J/\psi + Z$$

RECENT ATLAS MEASUREMENT: PROMPT J/PSI+Z



ATLAS EPJC'15

- First measurement by ATLAS
- Compare to theoretical calculations: DPS is dominant



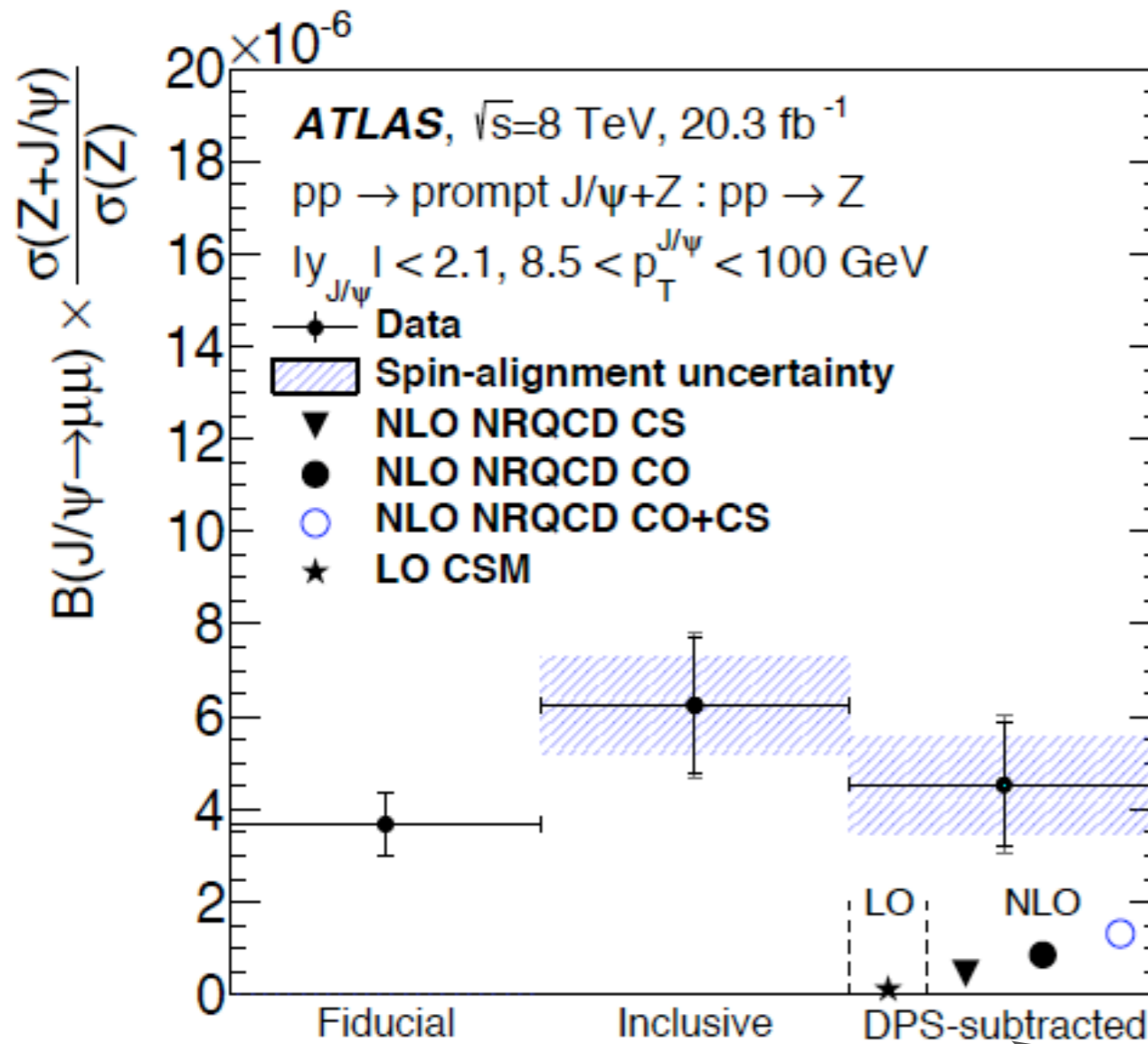
$$\sigma_{\text{eff}} = 15 \text{ mb}$$

RECENT ATLAS MEASUREMENT: PROMPT J/PSI+Z

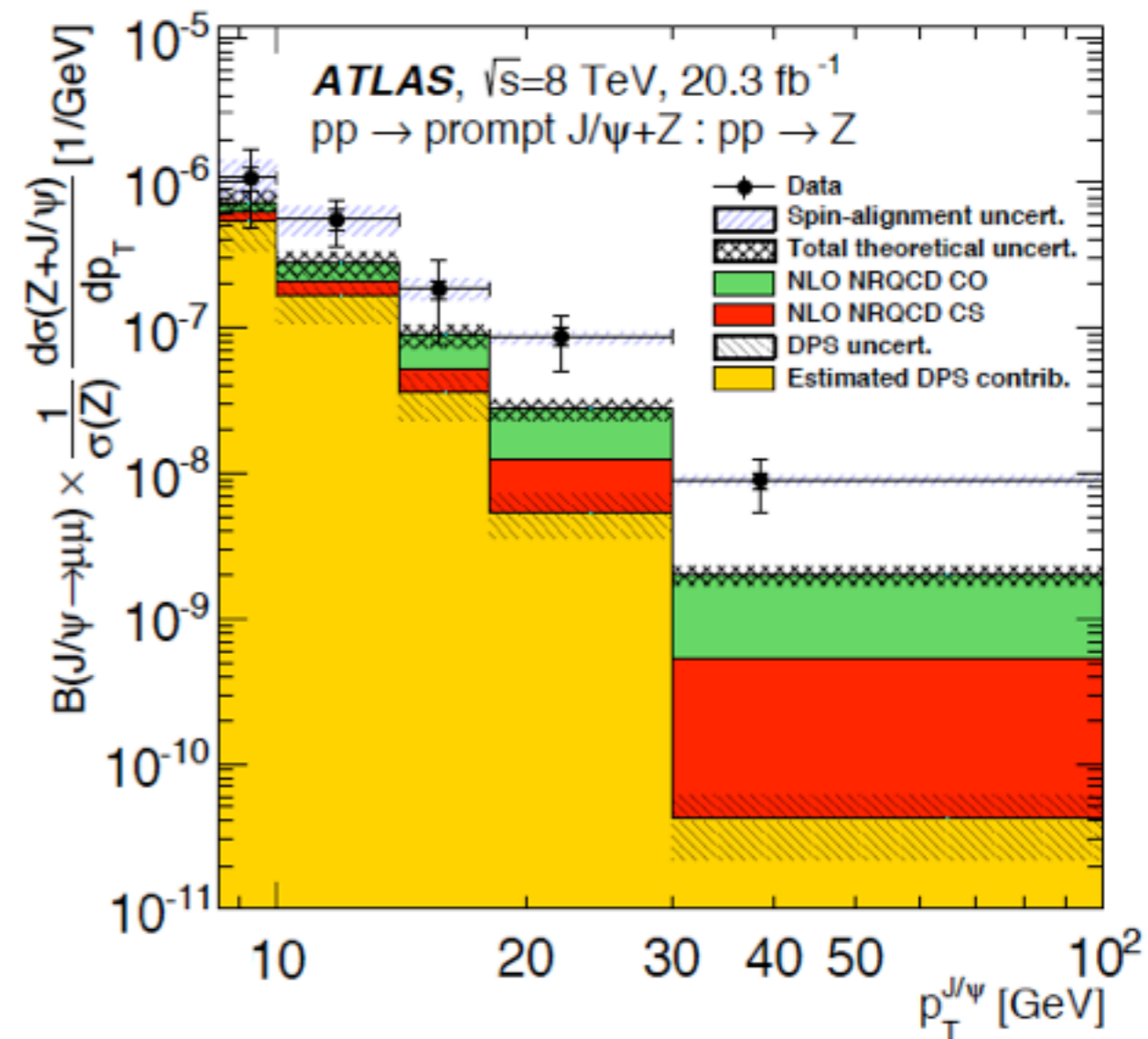


ATLAS EPJC'15

- First measurement by ATLAS
- Compare to theoretical calculations: DPS is dominant
- Compare to theoretical calculations: large discrepancy



$$\sigma_{\text{eff}} = 15 \text{ mb}$$

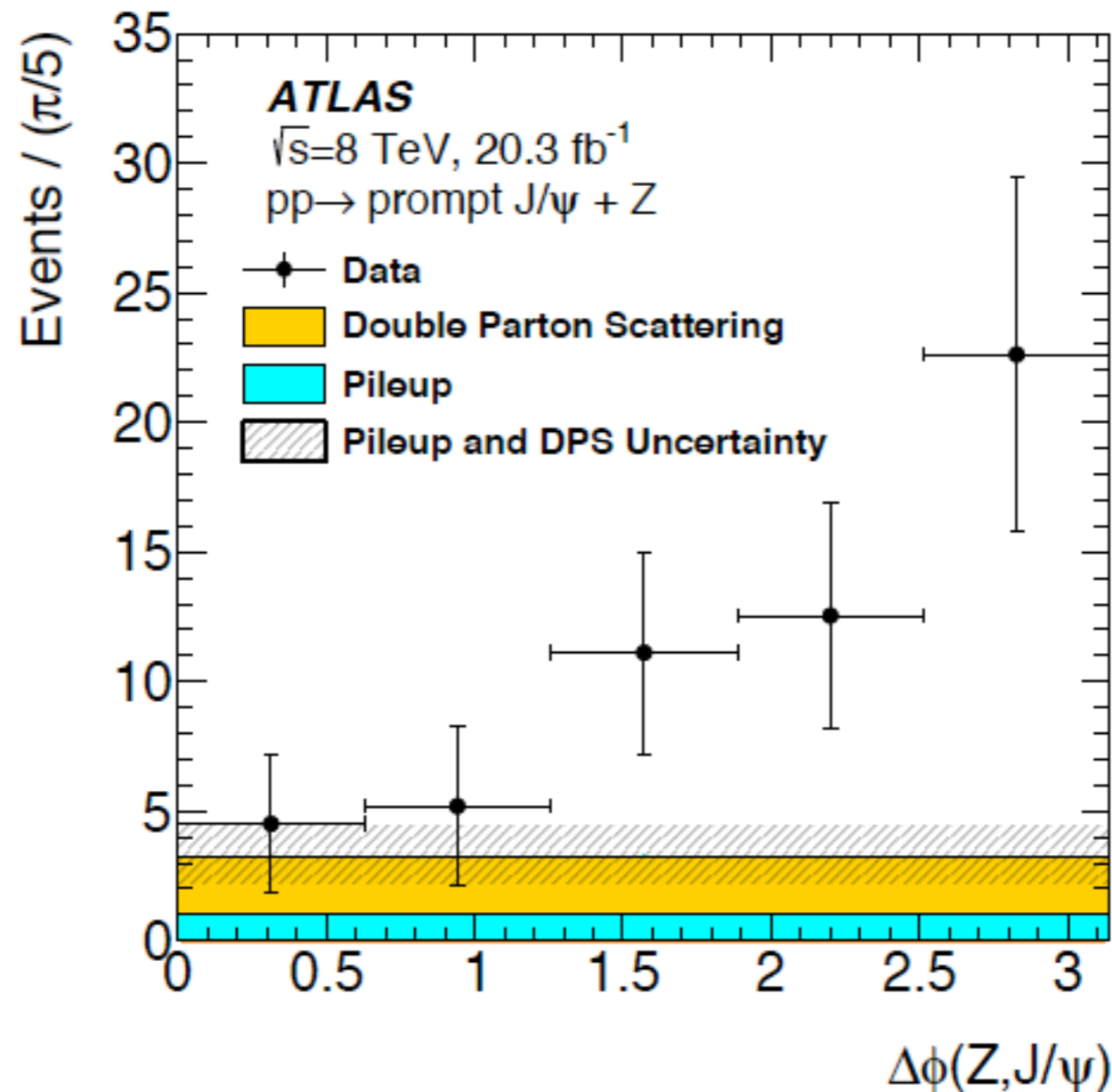


RECENT ATLAS MEASUREMENT: PROMPT J/ ψ +Z



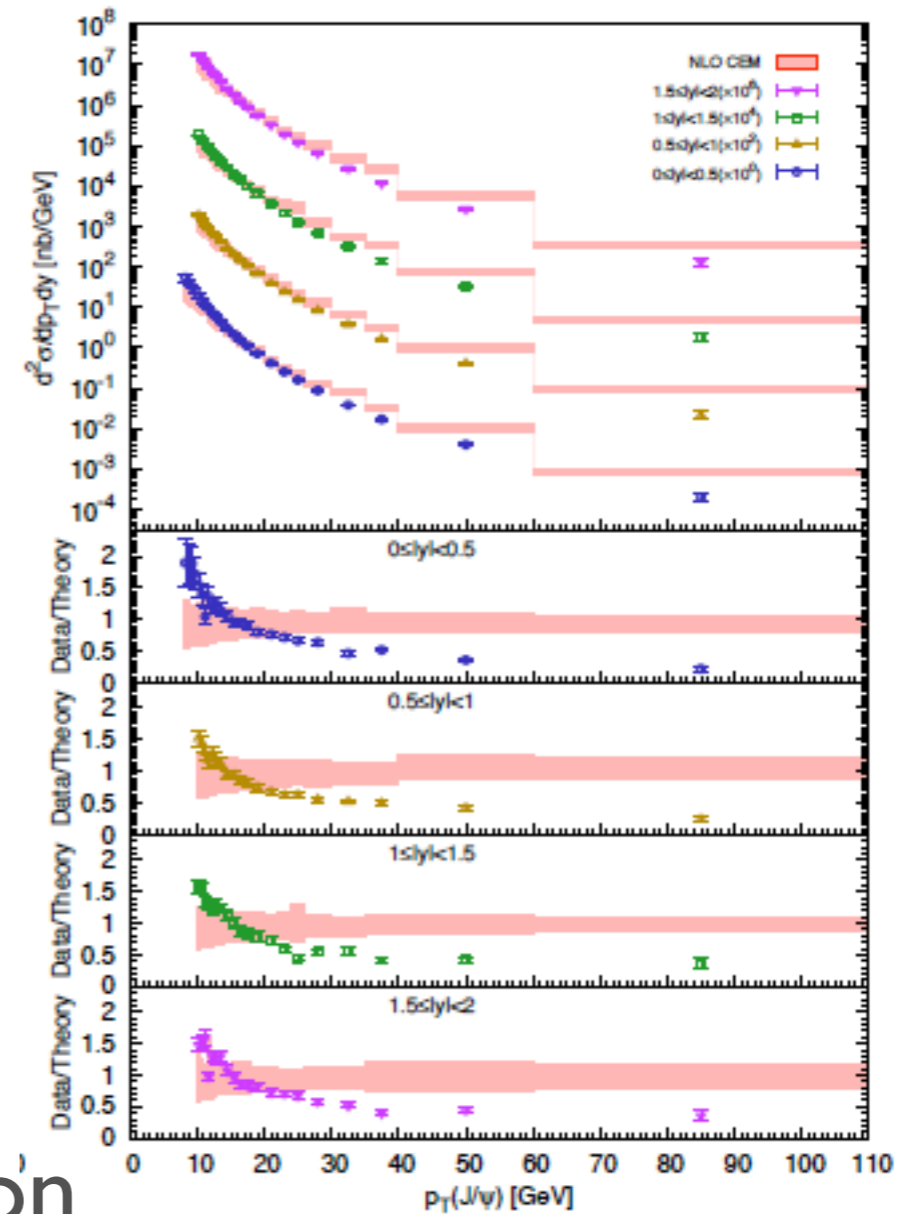
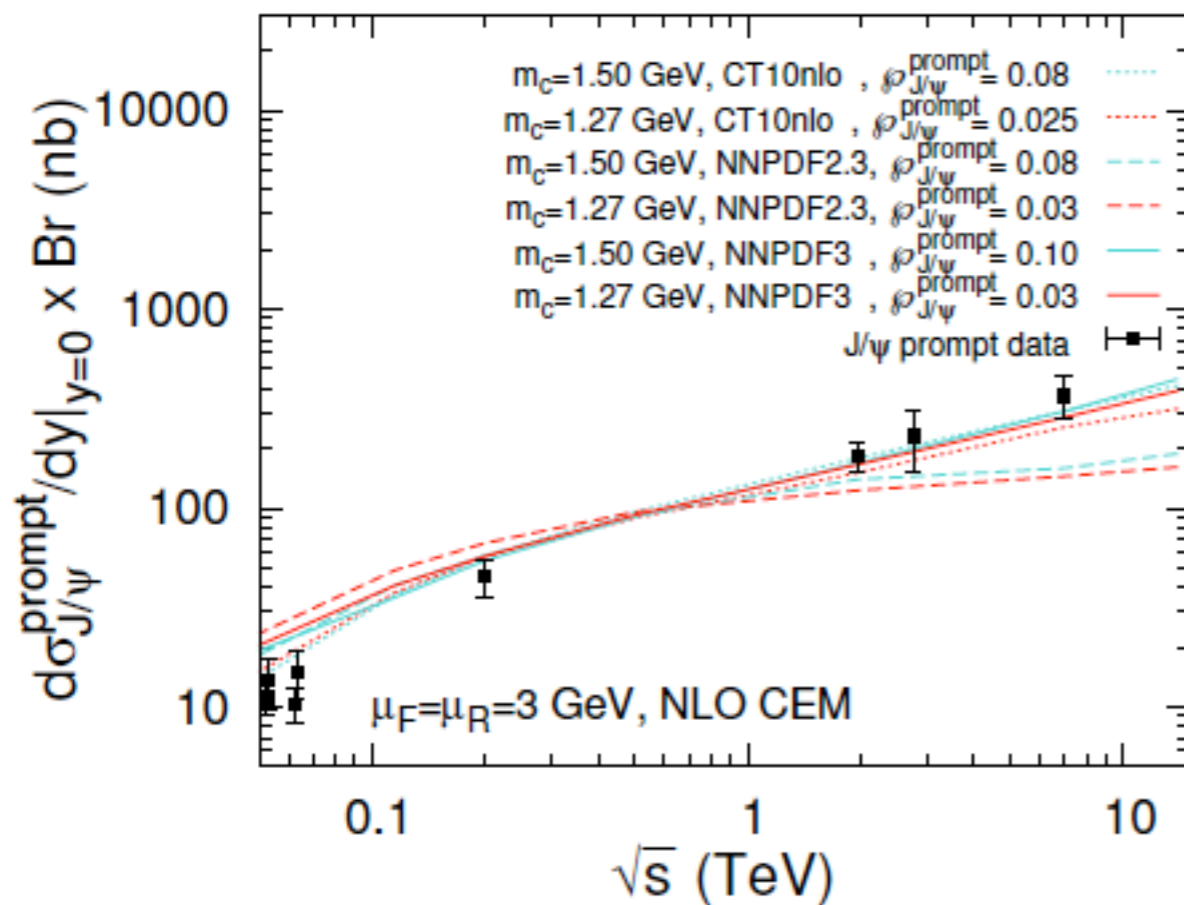
ATLAS EPJC'15

- First measurement by ATLAS
- Compare to theoretical calculations: DPS is dominant
- Compare to theoretical calculations: large discrepancy
- However, it seems one needs SPS



REANALYSIS J/PSI PRODUCTION IN CEM

- Prompt J/psi production in quark-hadron duality (CEM)
- **Pro**: only one parameter and easy to include higher order
- **Con**: not good to describe the pt spectrum Lansberg, HSS '16

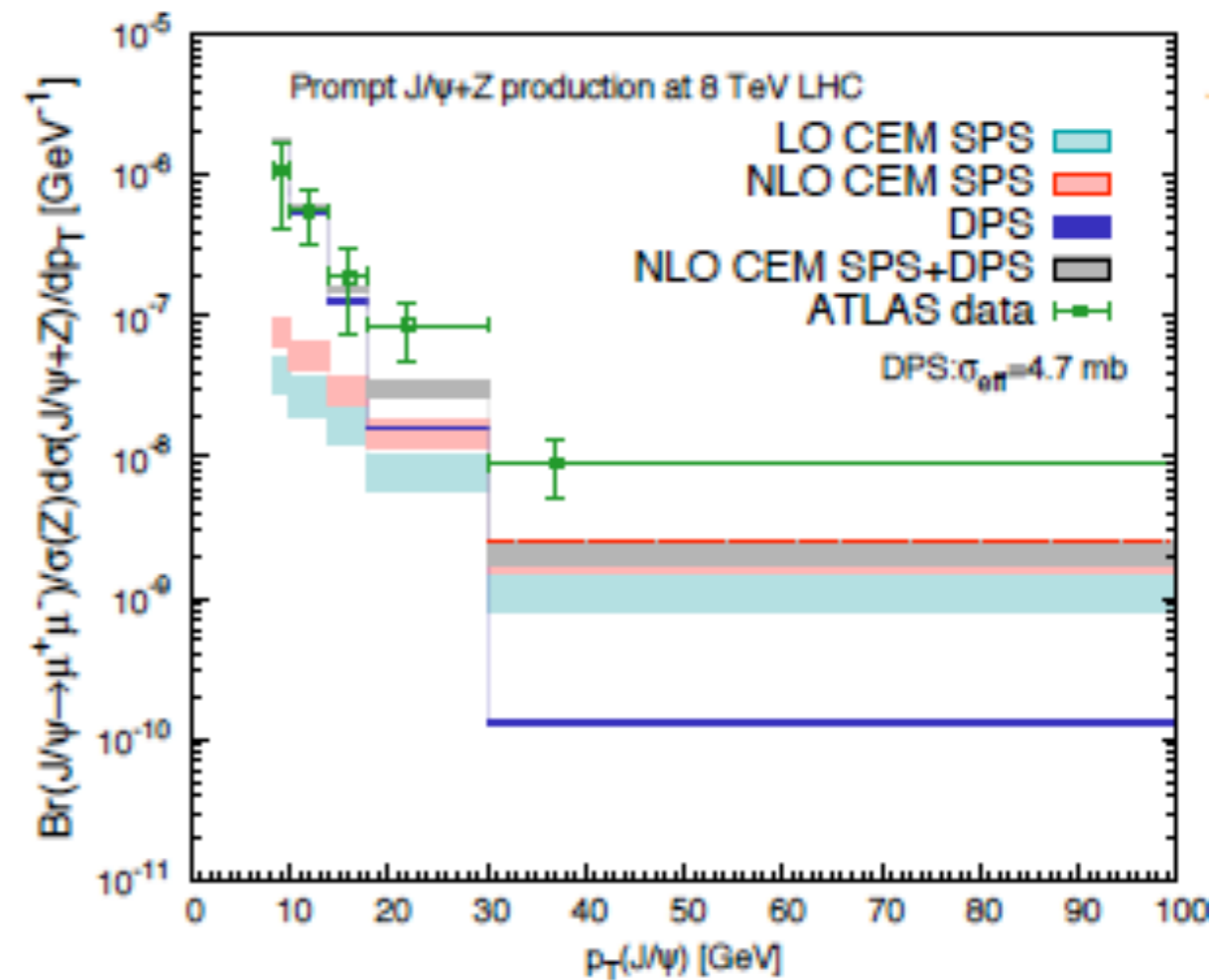
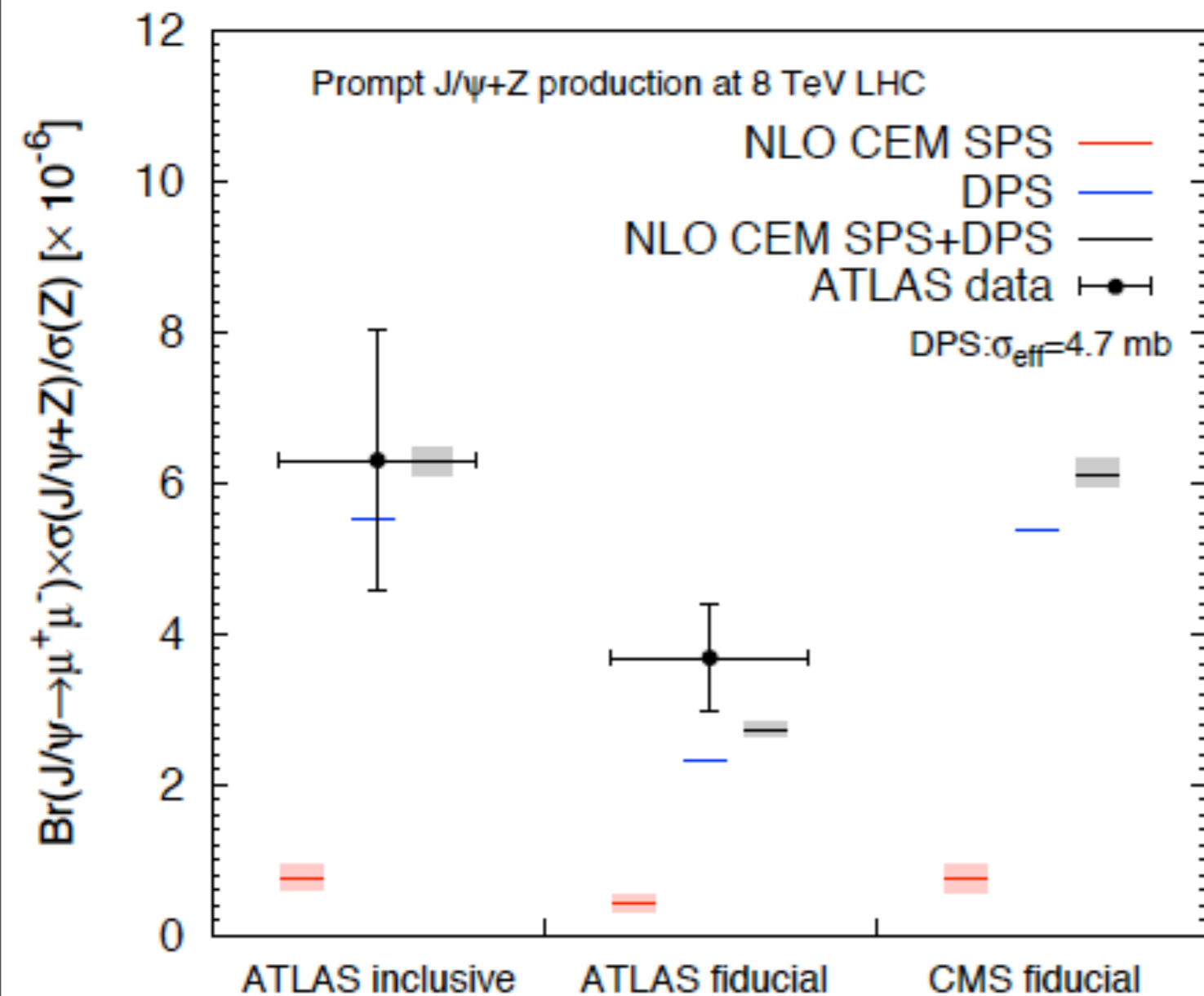


- Upper limit yields for J/psi production

PROMPT J/PSI+Z: NLO CEM VS DATA

Lansberg, HSS '16

- Applying CEM to J/psi+Z: NLO QCD vs ATLAS data
- Confirm DPS is dominant and small effective sigma

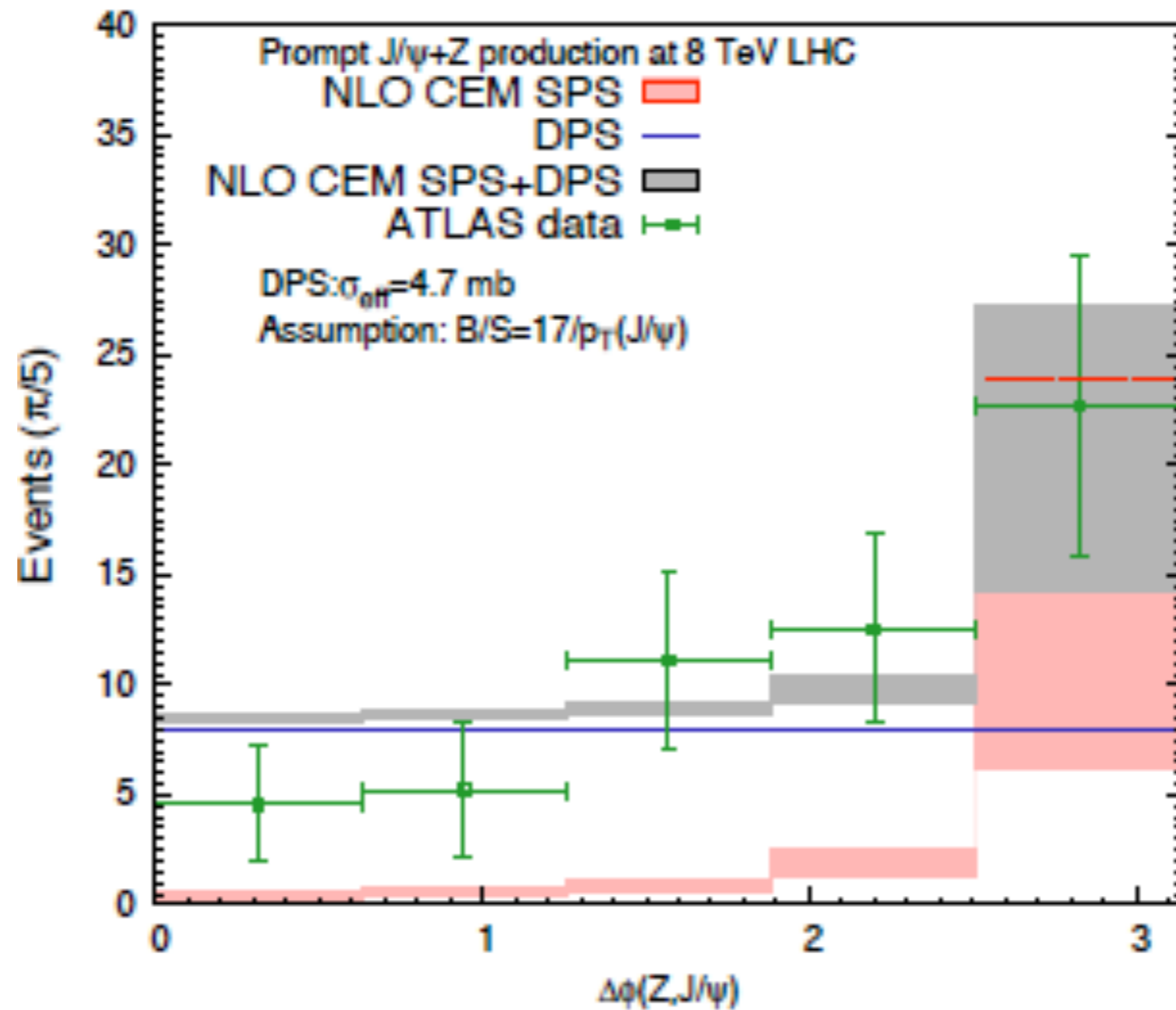


PROMPT J/ ψ +Z: NLO CEM VS DATA

Lansberg, HSS '16



- Azimuthal angular can be described well with SPS+DPS

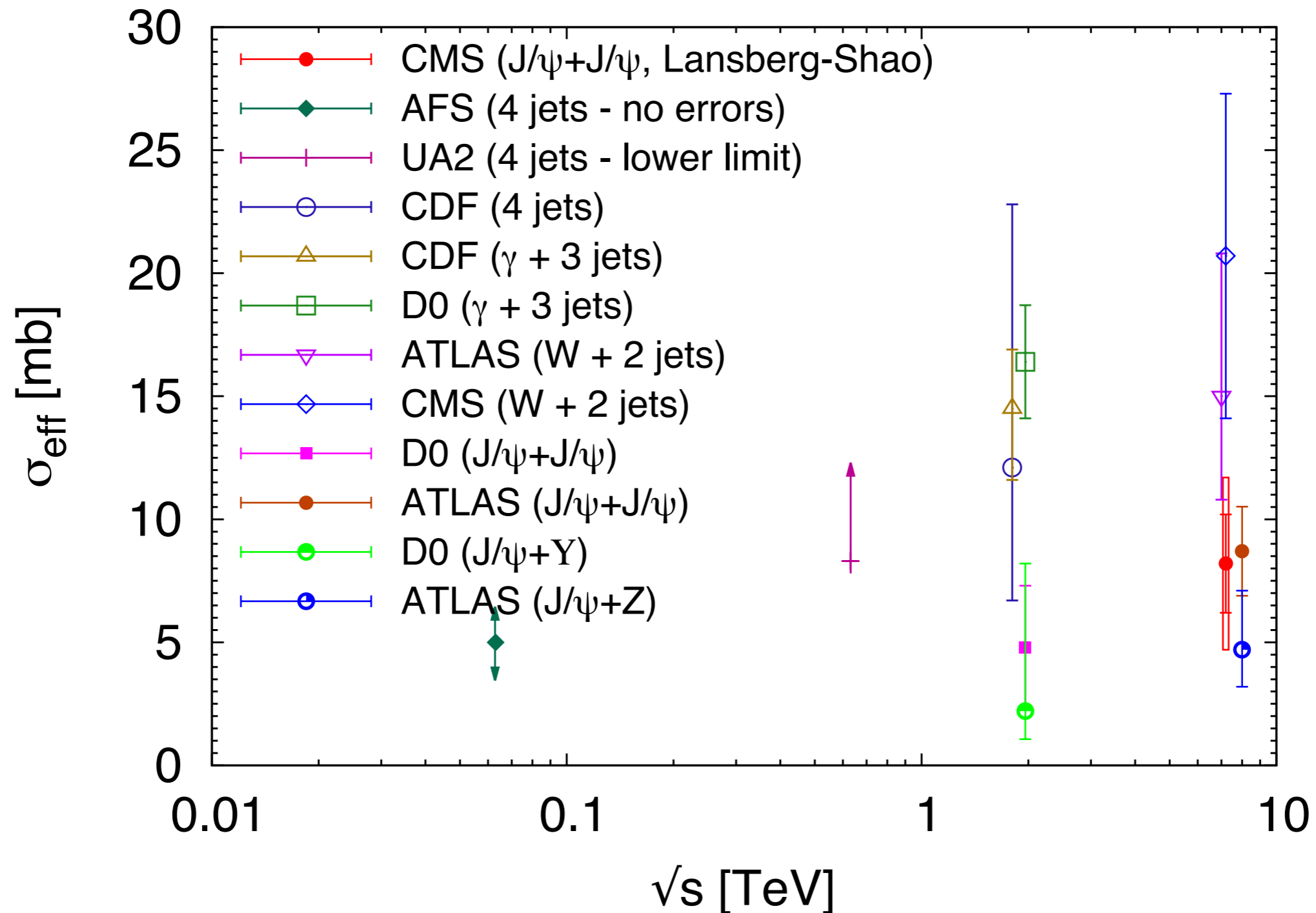


PROMPT J/ ψ +Z: NLO CEM VS DATA

Lansberg, HSS '16



- Azimuthal angular can be described well with SPS+DPS
- Refine the sigma effective

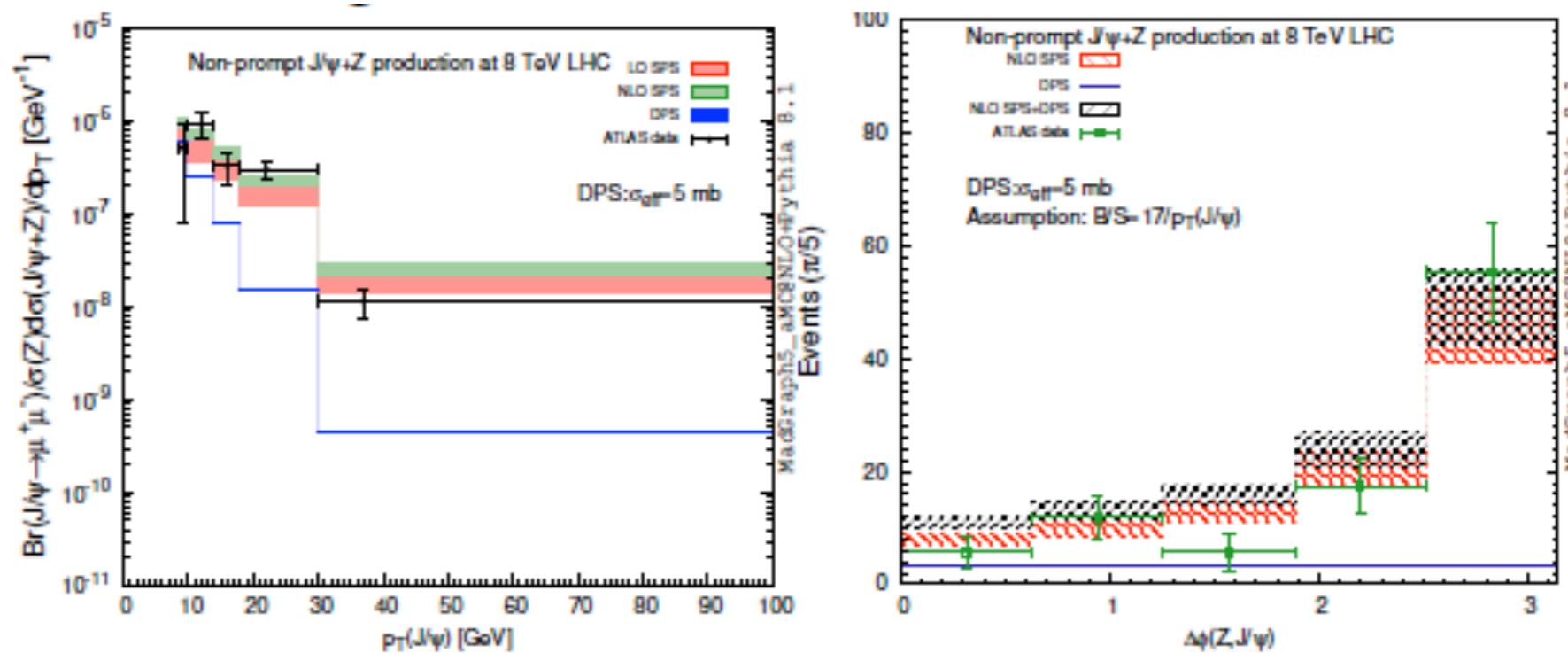


NONPROMPT J/ Ψ +Z: NLO+PYTHIA8 VS DATA



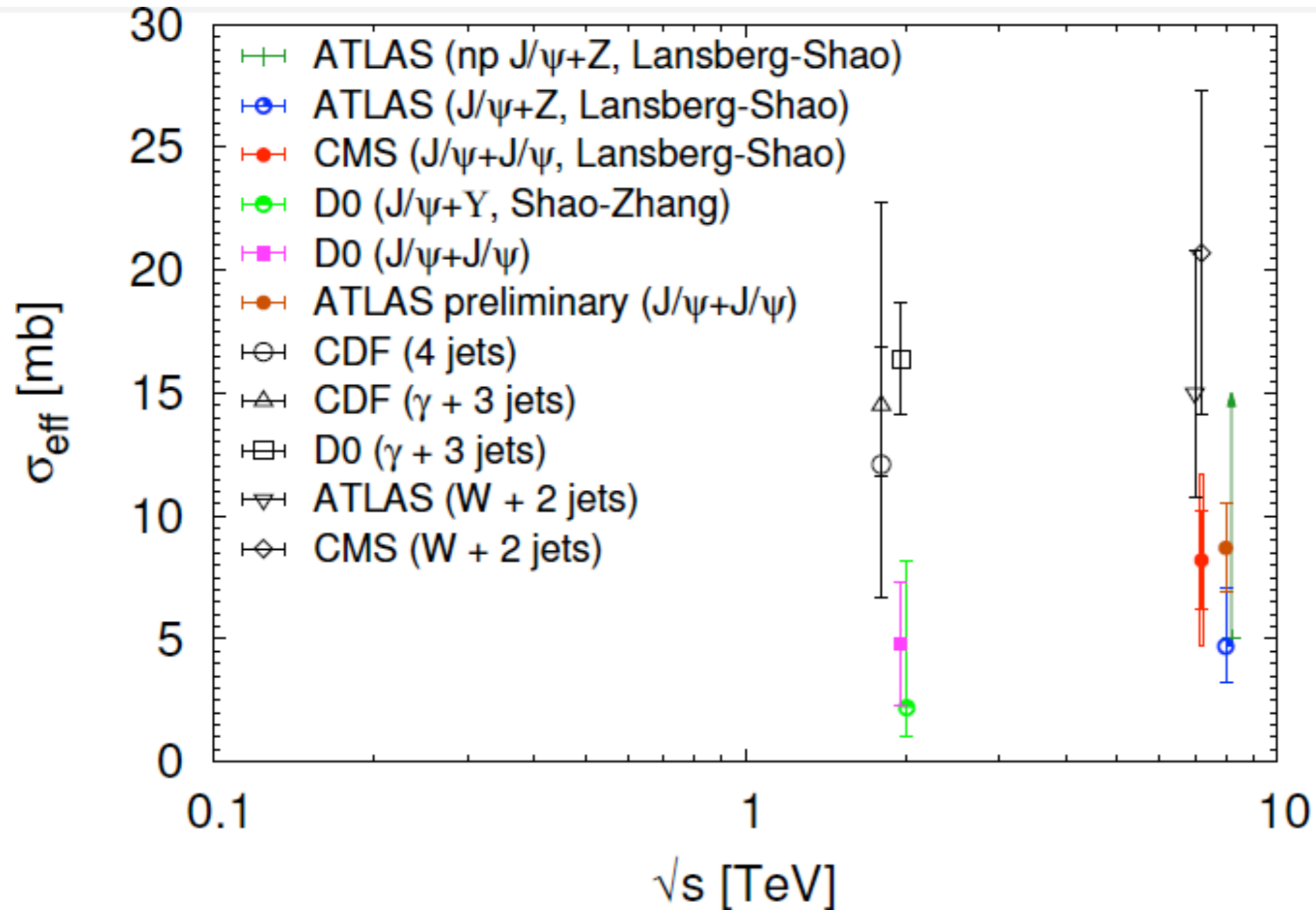
Lansberg, HSS '17

- In the same analysis, ATLAS reported on Z+nonprompt J/ Ψ
- This gives an original handle on Z+b at lower P_T than b-jets
- Interesting check that nothing went wrong with the prompt analysis
- SPS predictions were absent at the time. We filled the gap using MadGraph5_aMC@NLO and Pythia 8.1.



- In general, SPS is dominant here, which helps to refine the lower limit of the effective sigma.

DPS IN QUARKONIUM PRODUCTION: CONCLUSION



4 quarkonium extractions using theory ingredients