

BSM Theory Overview

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Dissecting the LHC results
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The ideal world:

I need to write a paper tonight!

Alas! My model is ruled out!

Theorists

The wise experimentalist



The real world:

Struggling to put together a model in the midst of "buzzers"!

Flavour!

Higgs couplings!

No DM!

EWPTs!

Top couplings!

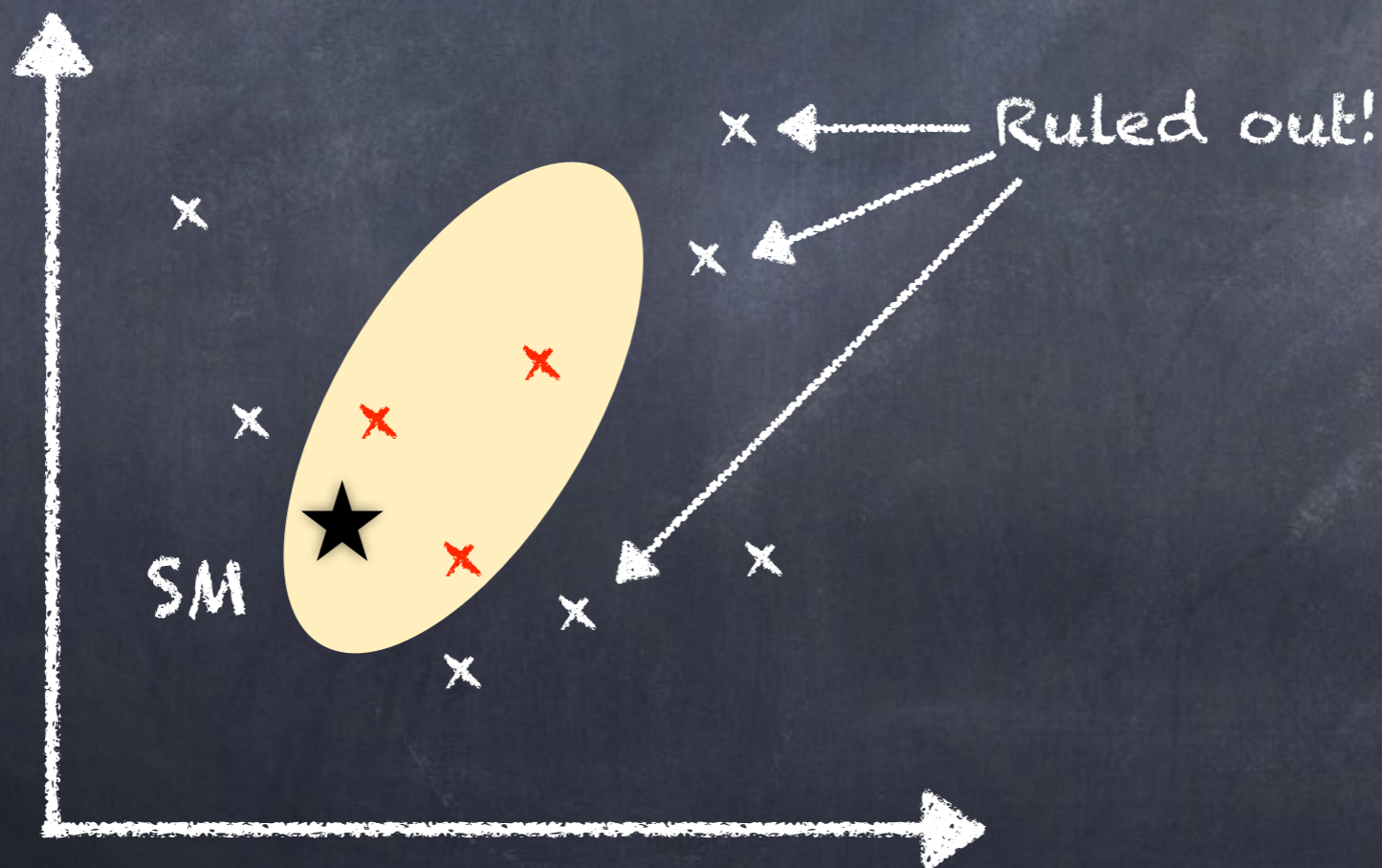
No bumps!

Higgs mass!



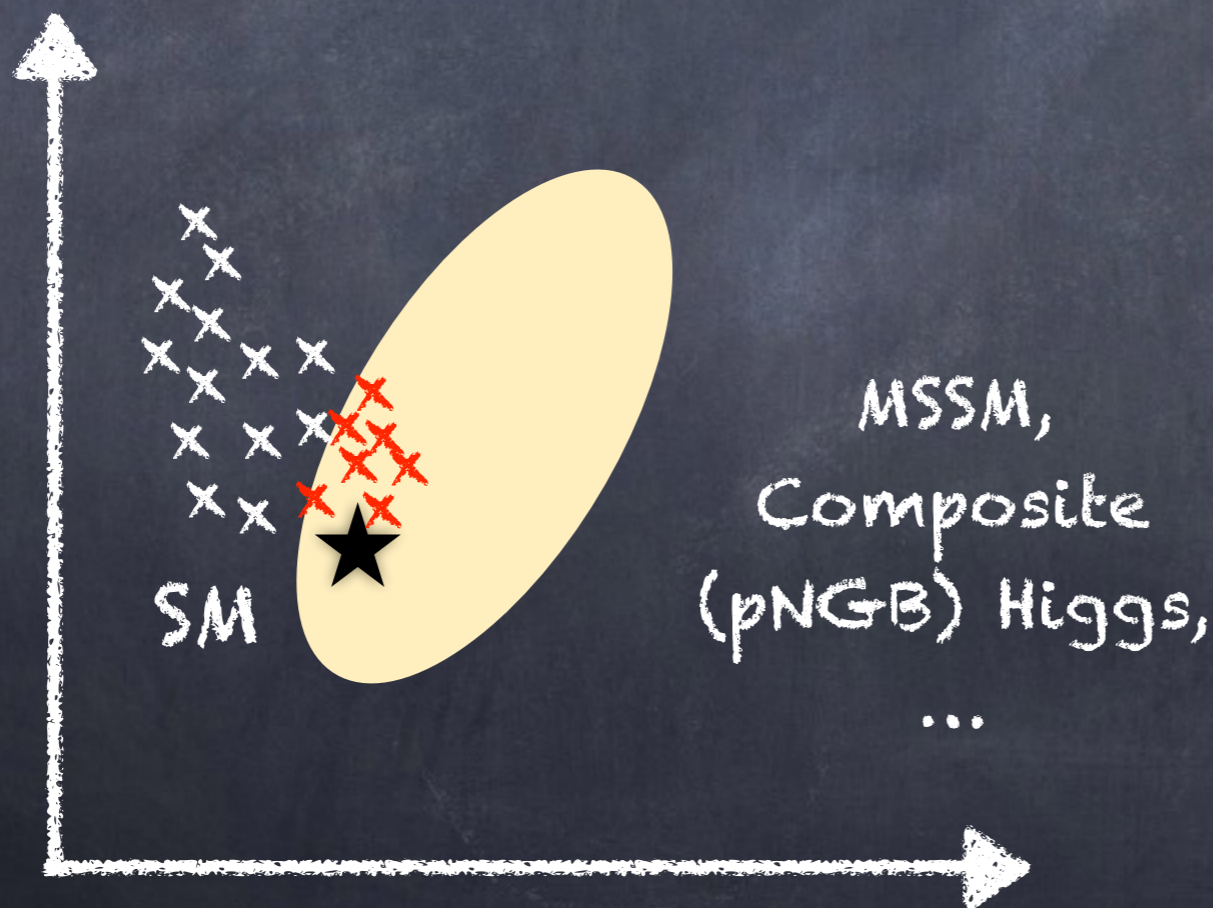
What is our job?

- Models can be ruled out, but cannot be proven right!



What is our job?

- Models can be ruled out, but cannot be proven right!



Class A:

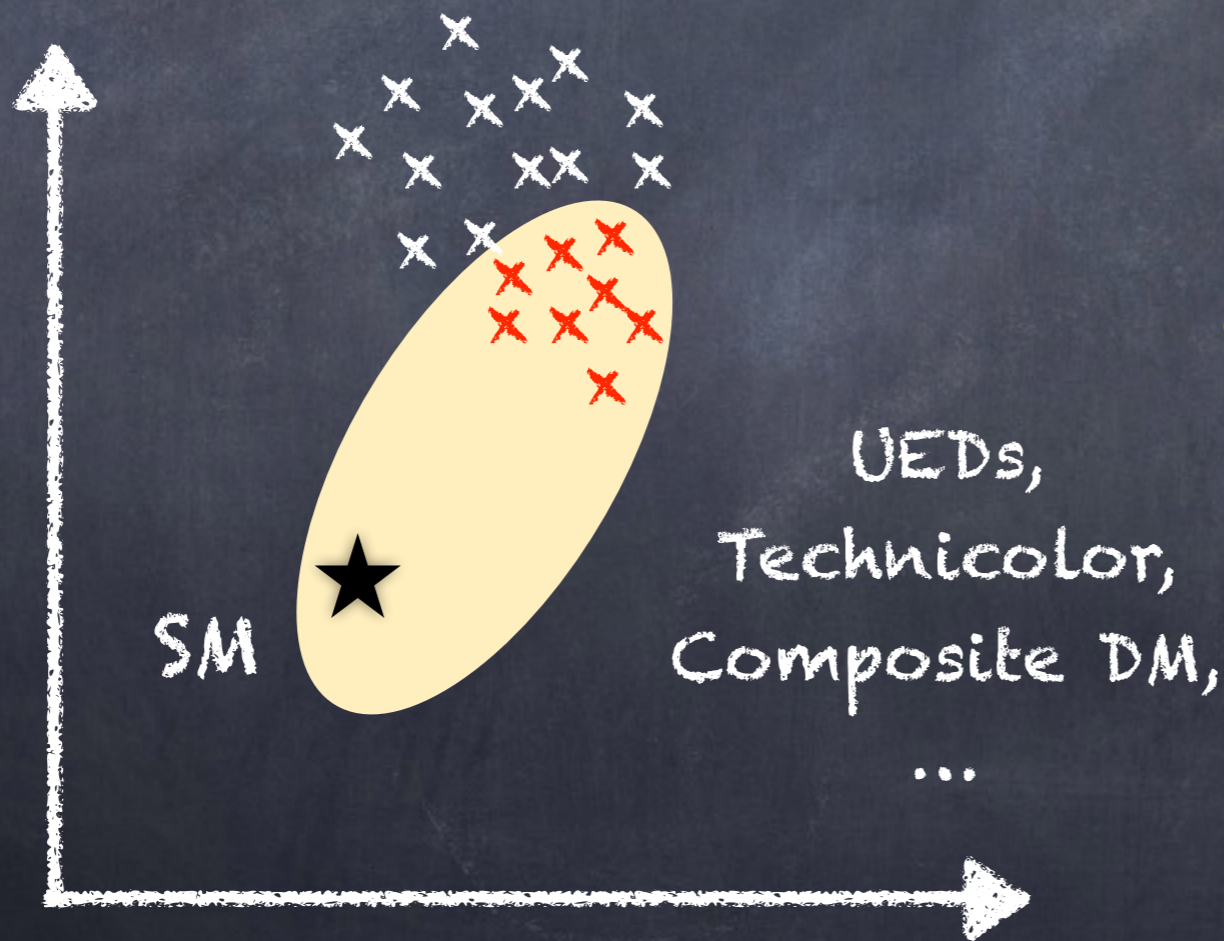
Parameter space
connected to the
SM prediction



Cannot be ruled
out!

What is our job?

- Models can be ruled out, but cannot be proven right!



Class B:

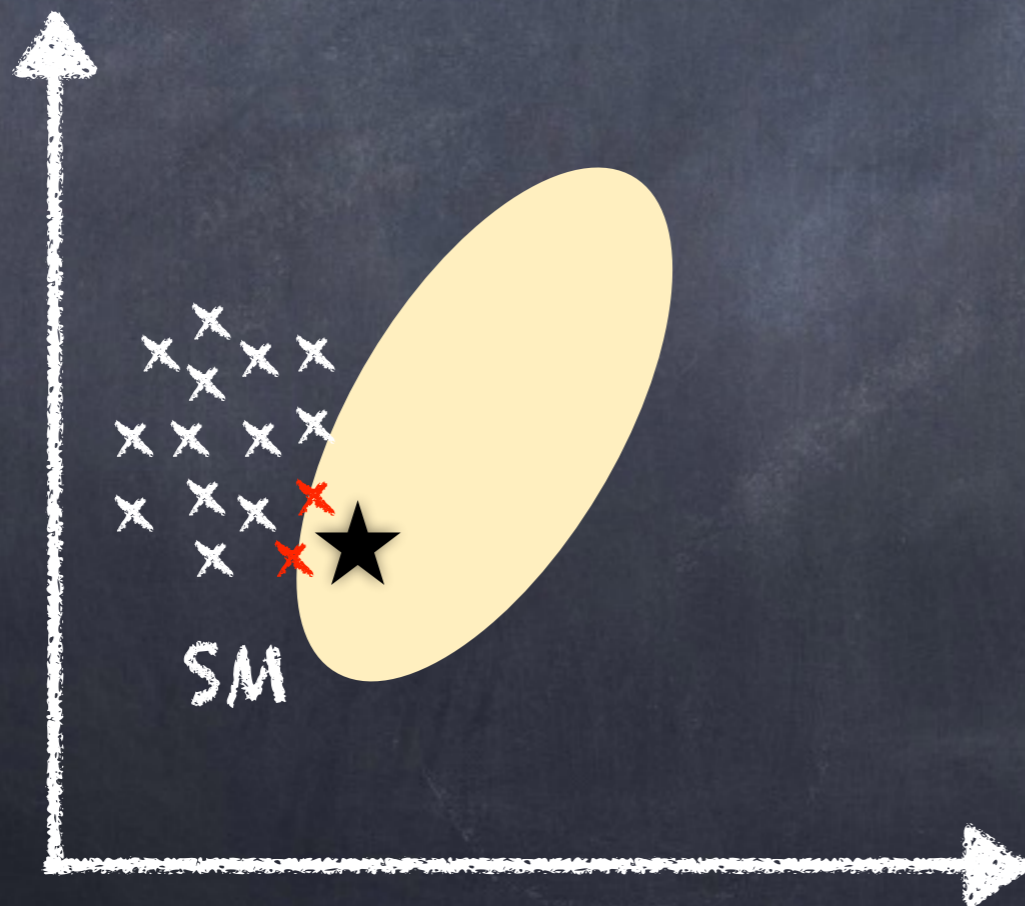
Parameter space disconnected from SM prediction



Can be ruled out!

What is our job?

- Models can be ruled out, but cannot be proven right!



Grey zone:

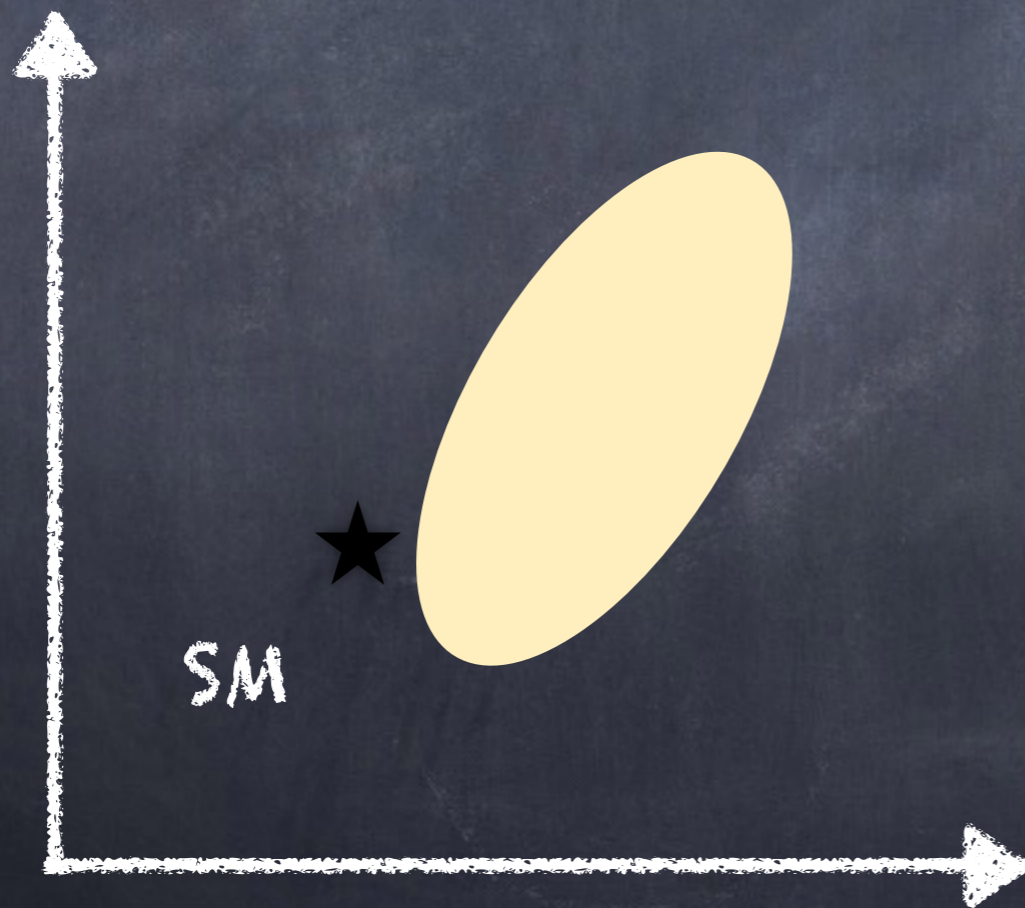
Fine tuning?

Personal taste?

How close to
the decoupling limit?

What is our job?

- Models can be ruled out, but cannot be proven right!



BSM dream:

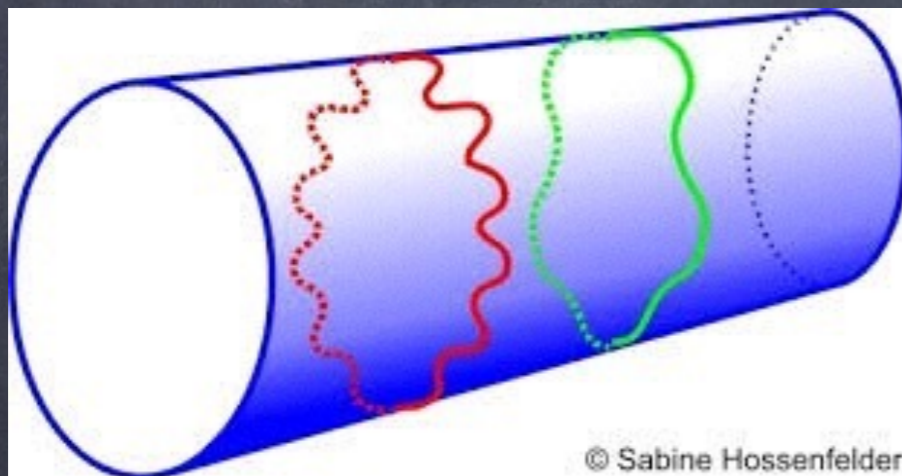
The SM itself can be excluded!

Are we there yet?

No...

Class B: UEDs

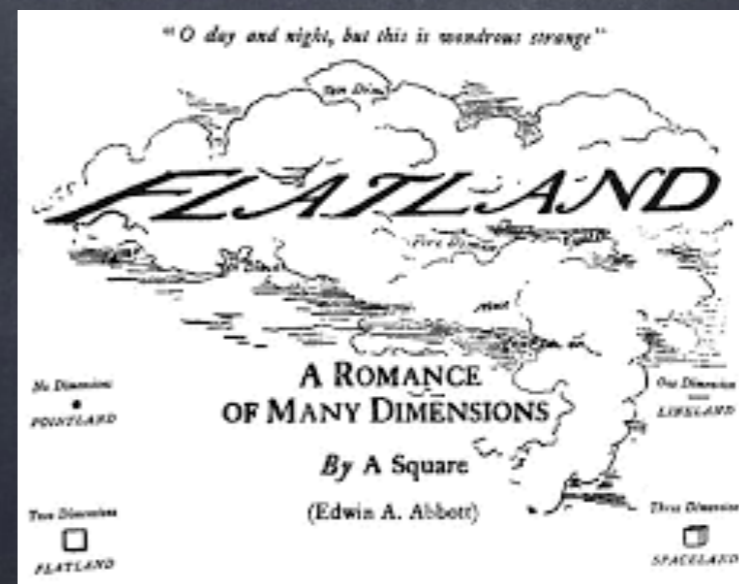
- Model of Dark Matter based on extra dimensions!



XD fields \rightarrow tower of KK states

frequencies \rightarrow KK masses

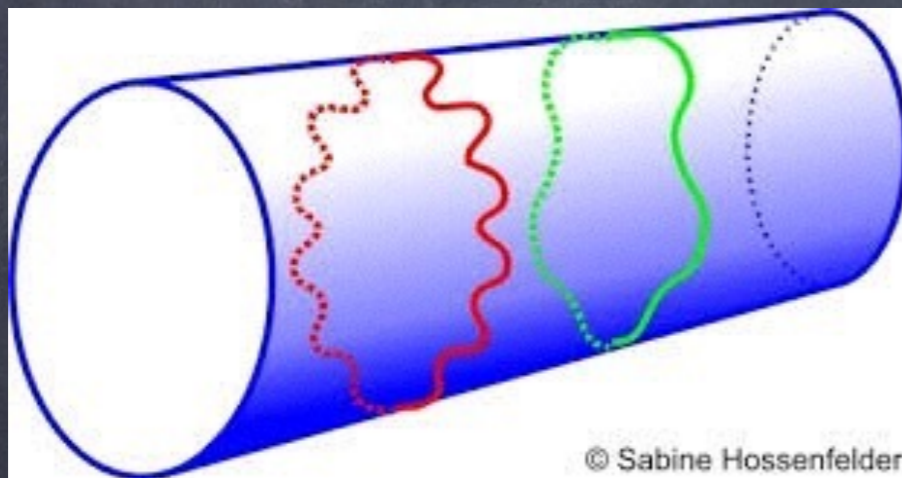
geometry \rightarrow KK parities



Class B: UEDs

- Model of Dark Matter based on extra dimensions!

1702.00410



Mass splitting crucially depends on the "cut-off" of the theory.

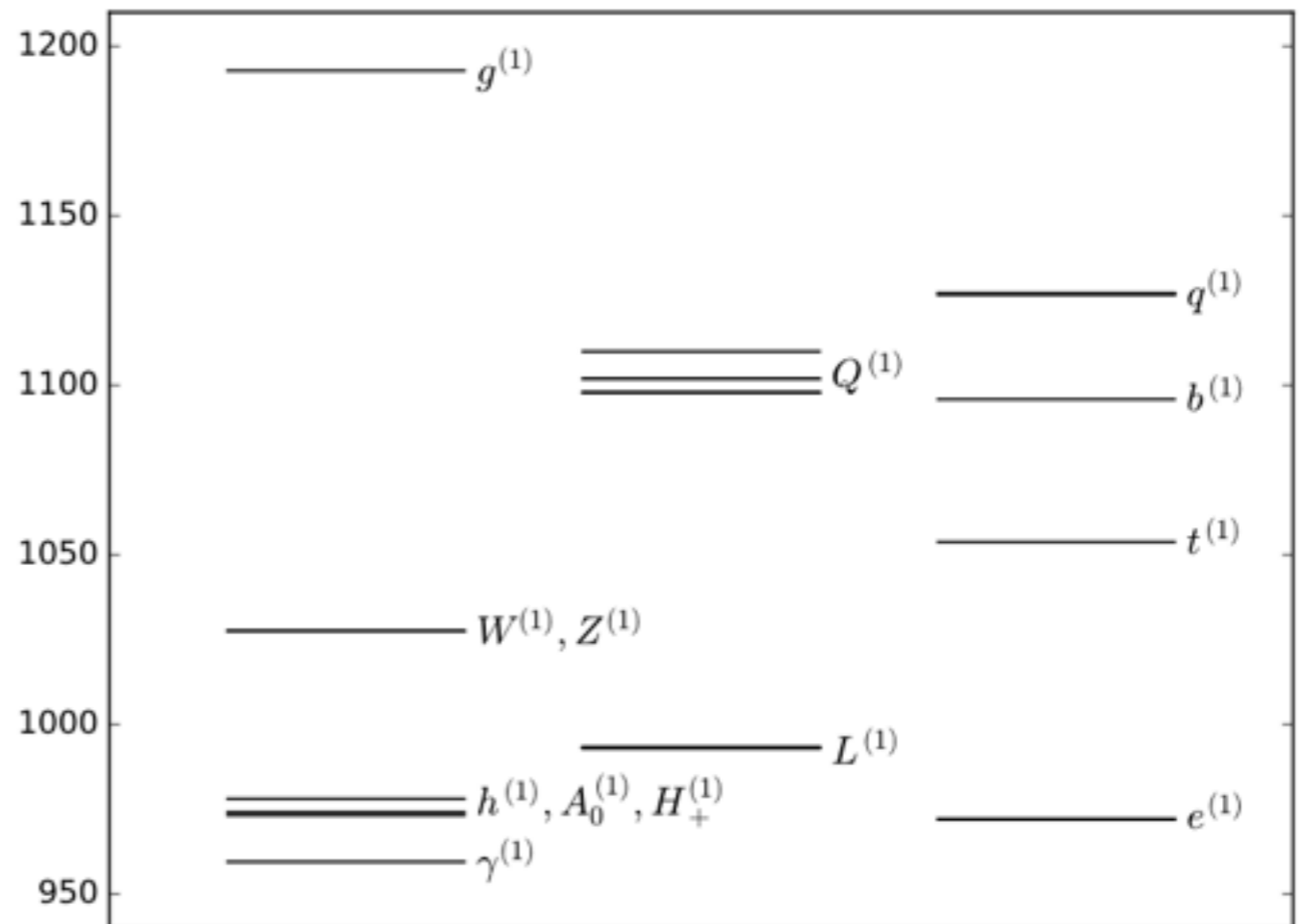
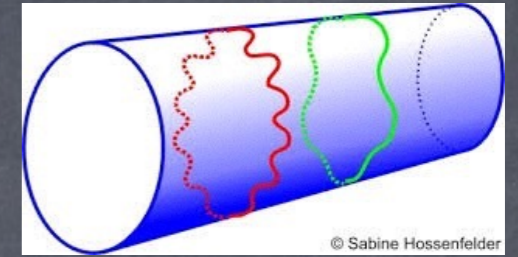


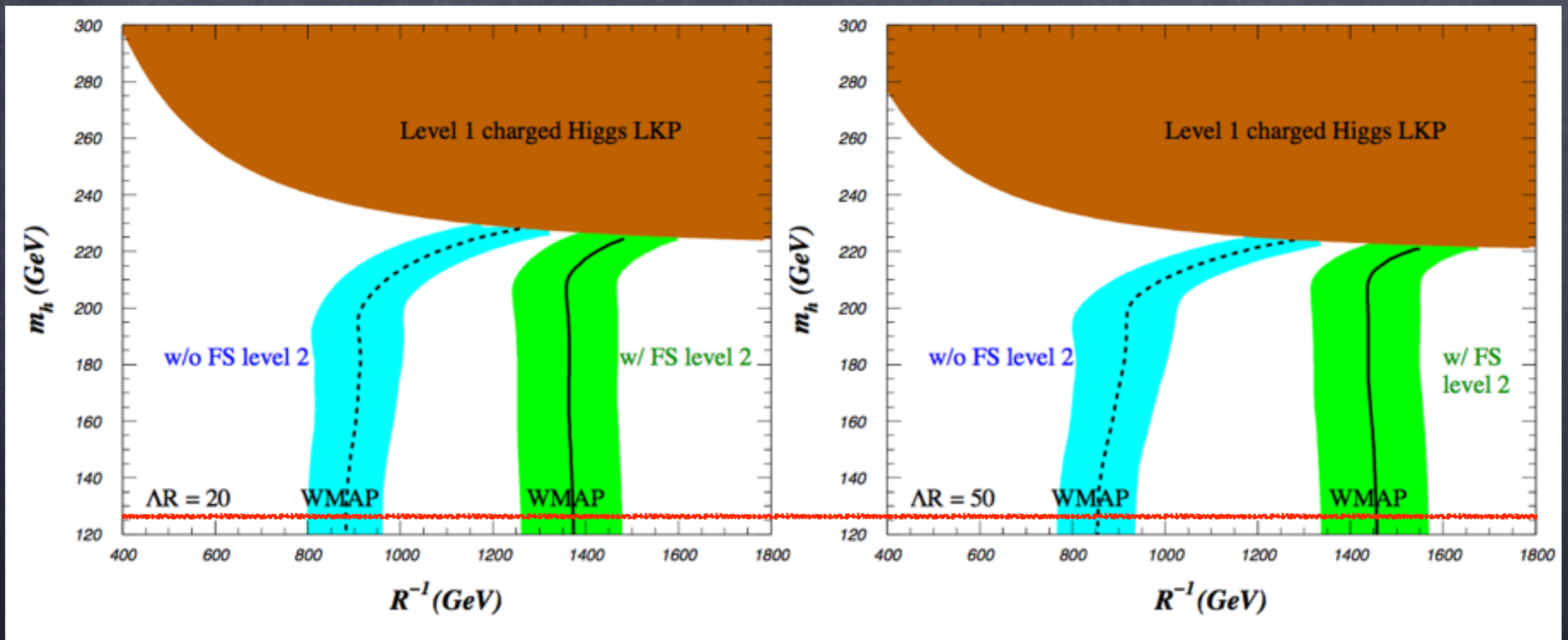
FIG. 1. Loop corrected KK mass spectrum in MUED for $R^{-1}=960$ GeV, $\Lambda R=30$ and $m_h=125$ GeV.

Class B: UEDs

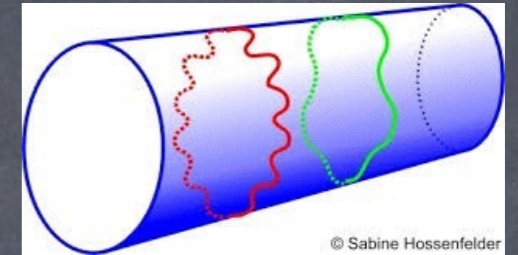


- Model of Dark Matter based on extra dimensions!

1012.2577



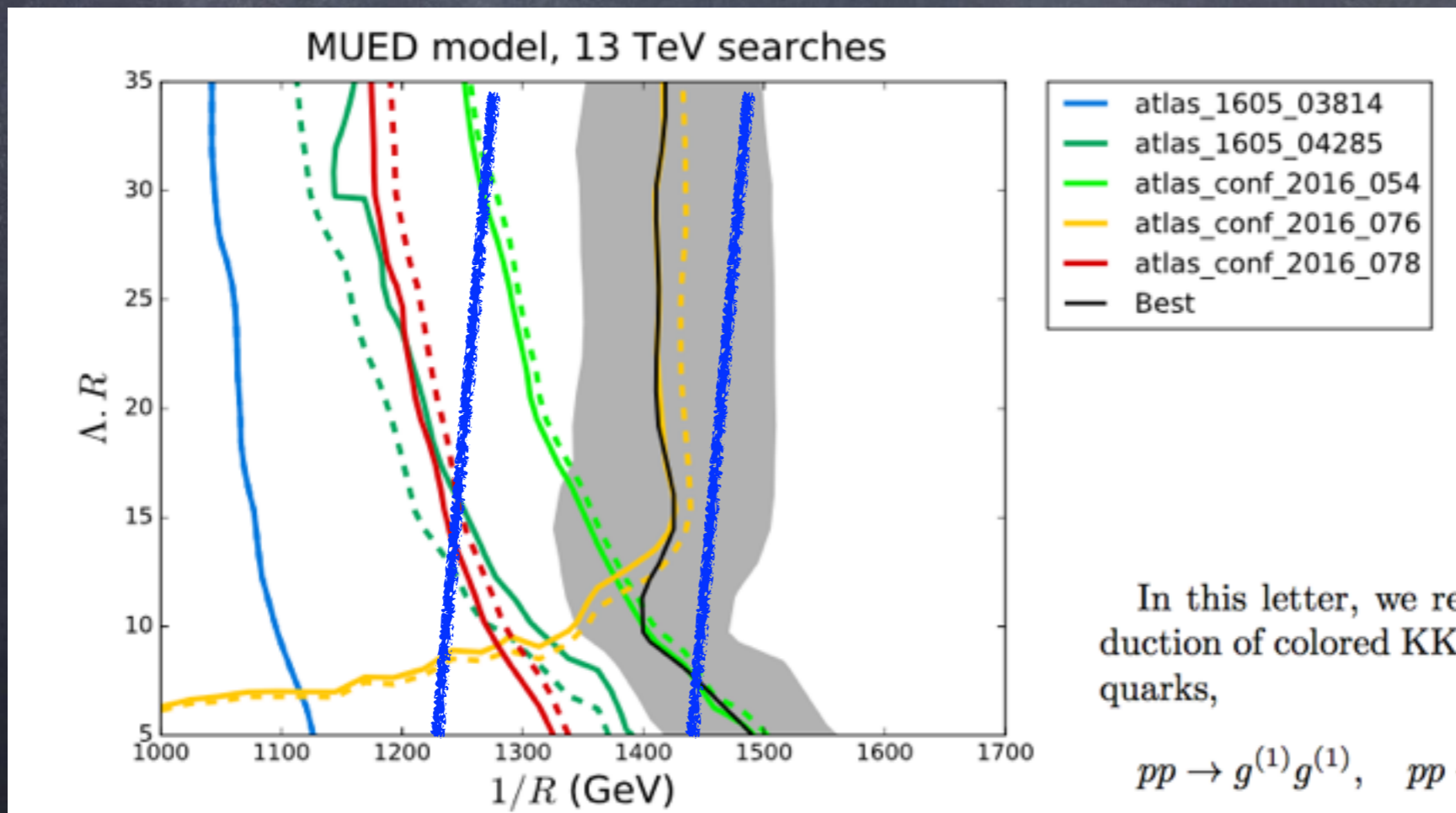
Class B: UEDs



- Model of Dark Matter based on extra dimensions!

1702.00410

On the verge of exclusion!

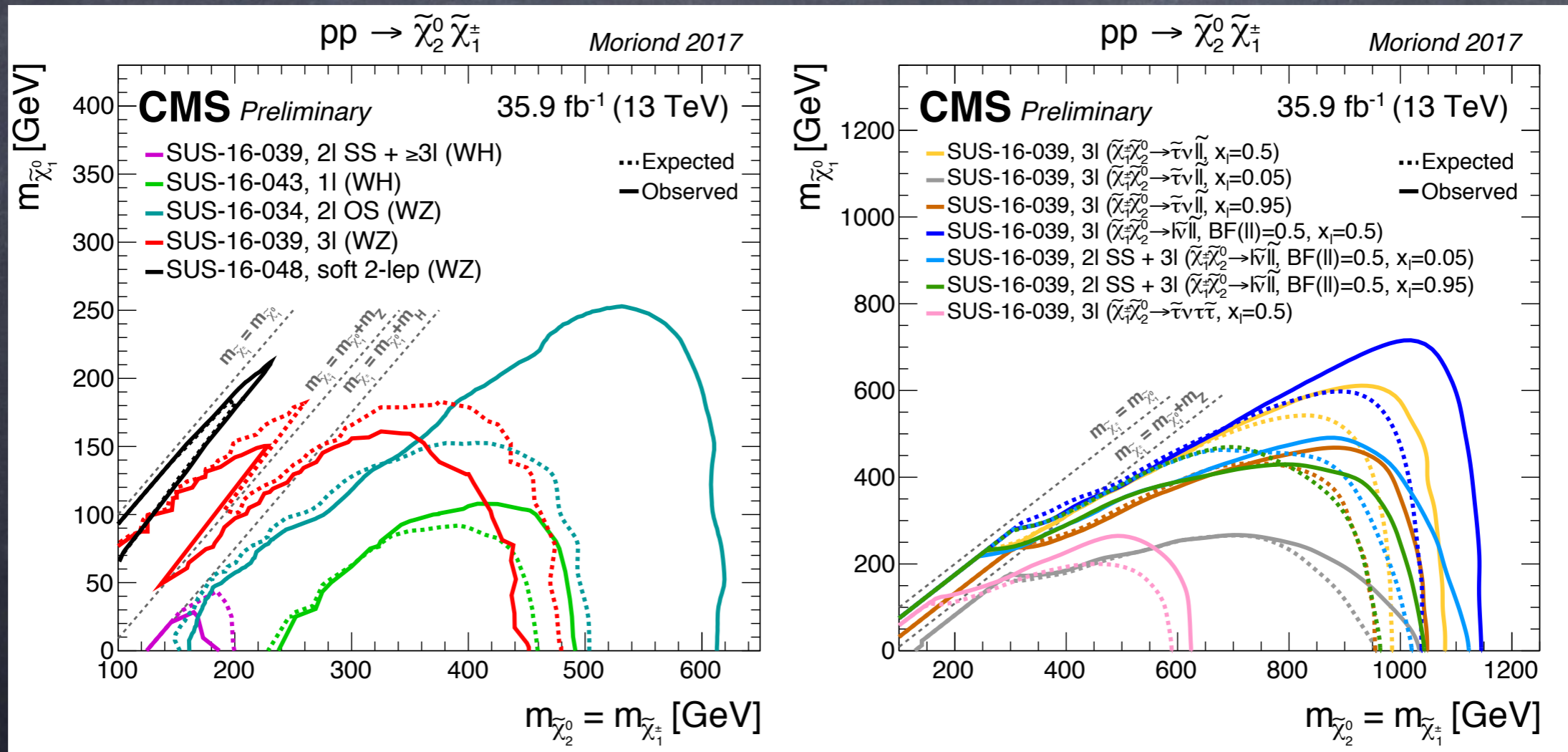


In this letter, we restrict ourselves to the strong production of colored KK modes such as KK gluons and KK quarks,

$$pp \rightarrow g^{(1)}g^{(1)}, \quad pp \rightarrow Q_i^{(1)}Q_j^{(1)}, \quad pp \rightarrow g^{(1)}Q_j^{(1)}, \quad (1)$$

Class A: MSSM

Moriond bounds on EW-inos

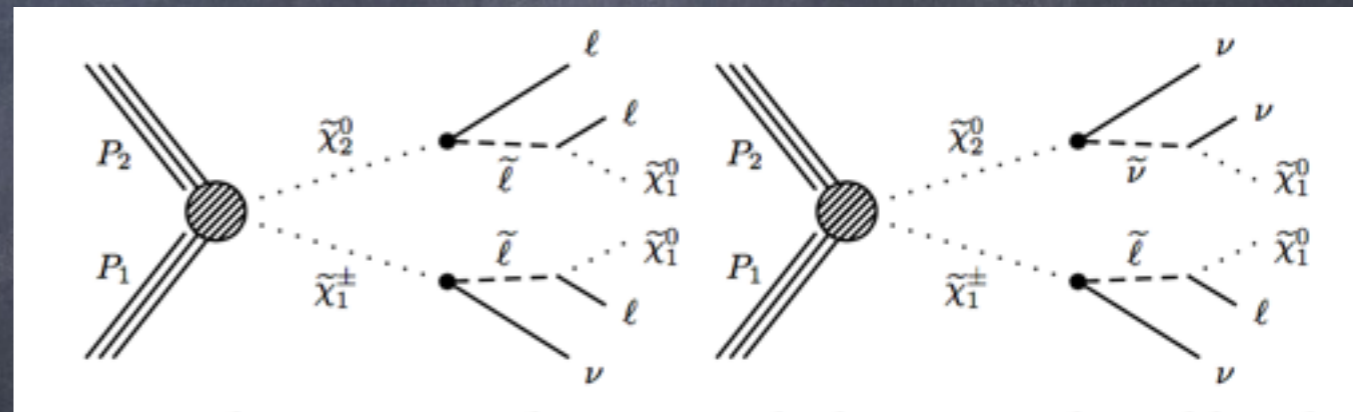
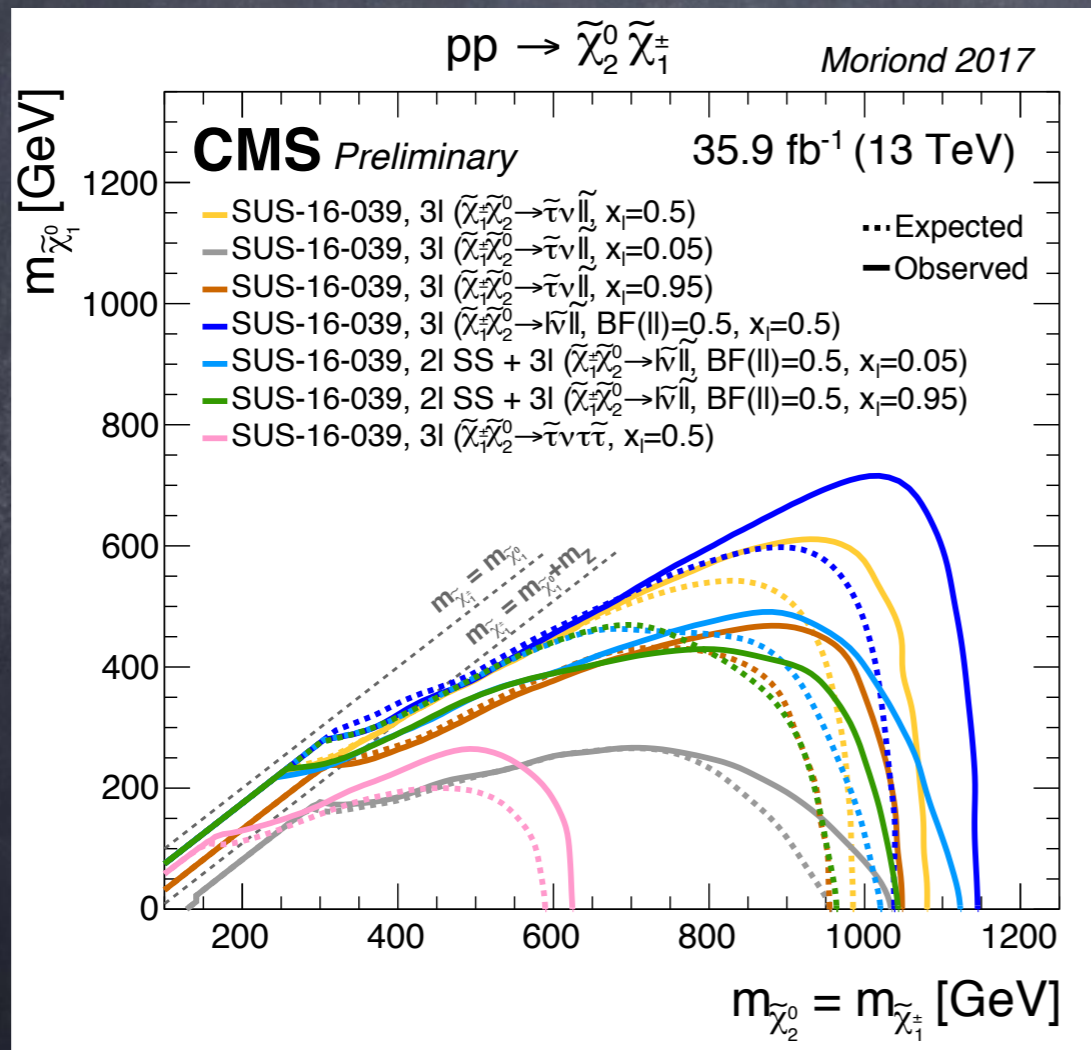


Class A: MSSM

• Moriond bounds on EW-inos

The bounds are impressive,
however...

those are for Simplified Models!



- Production sensitive to the details of the model;
- Decays assumed @ 100%!

Beware of Simplified Models!

- Simplified models are designed for simplifying exp. analyses!
- Bound on S.M.s cannot be used to give general conclusions on models!

The case of VLQs, aka Top partners

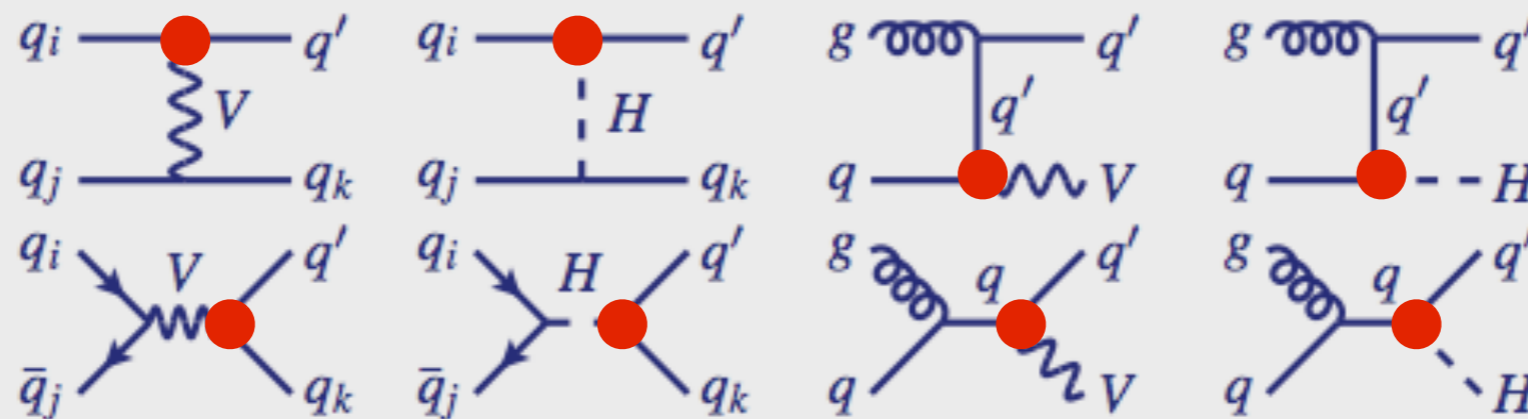
- VLQs are non-chiral quarks that mix to the SM ones via Higgs couplings!

	uW	uZ	uH	dW	dZ	dH
$T(2/3)$		γ	γ	γ		
$B(-1/3)$	γ				γ	γ
$X(5/3)$	γ					
$\gamma(-4/3)$				γ		

The case of VLQs, aka Top partners

- BRs depend on the model.
- QCD pair-production is model-independent!
- Single production is not:

Single production: $pp \rightarrow q' + \{q, V, H\}$



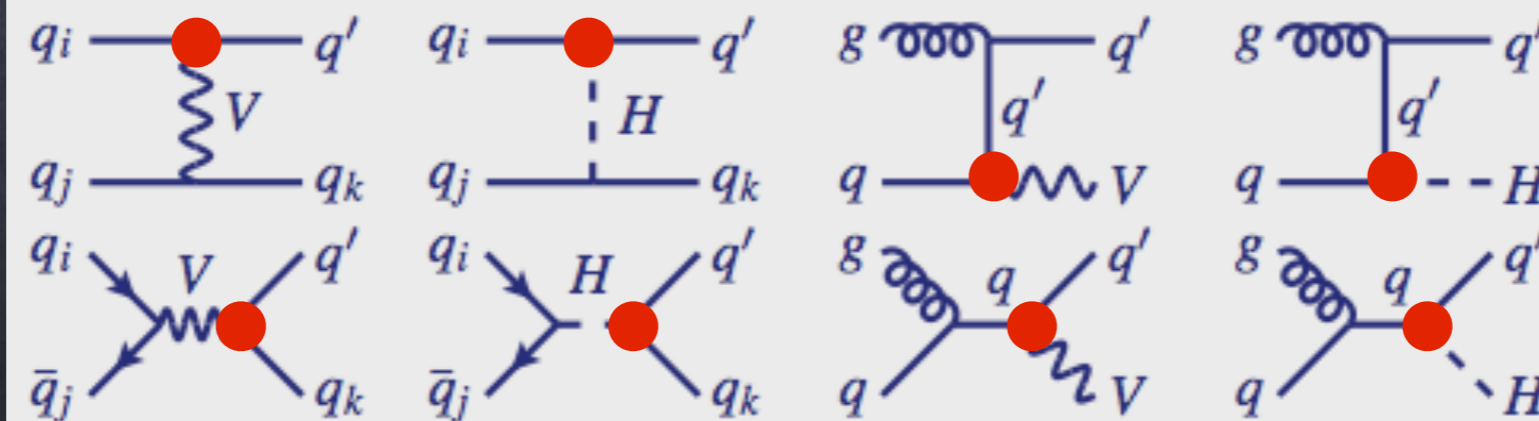
The case of VLQs, aka Top partners

$$-\hat{\kappa}_T H \bar{T}_R u_L + \frac{g}{2 \cos \theta_W} \tilde{\kappa}_T Z_\mu \bar{T}_L \gamma^\mu u_L + \frac{g}{\sqrt{2}} \kappa_T W_\mu^+ \bar{T}_L \gamma^\mu d_L$$

Couplings can be expressed in terms of the BRs,
plus an overall coupling strength!

BRs and Single Productions are correlated!

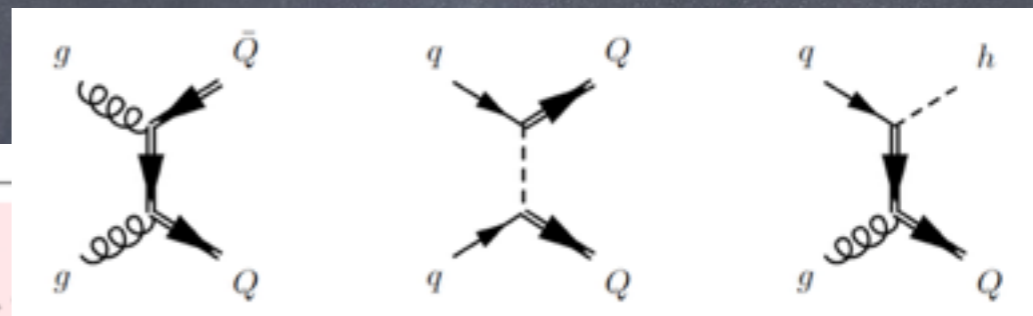
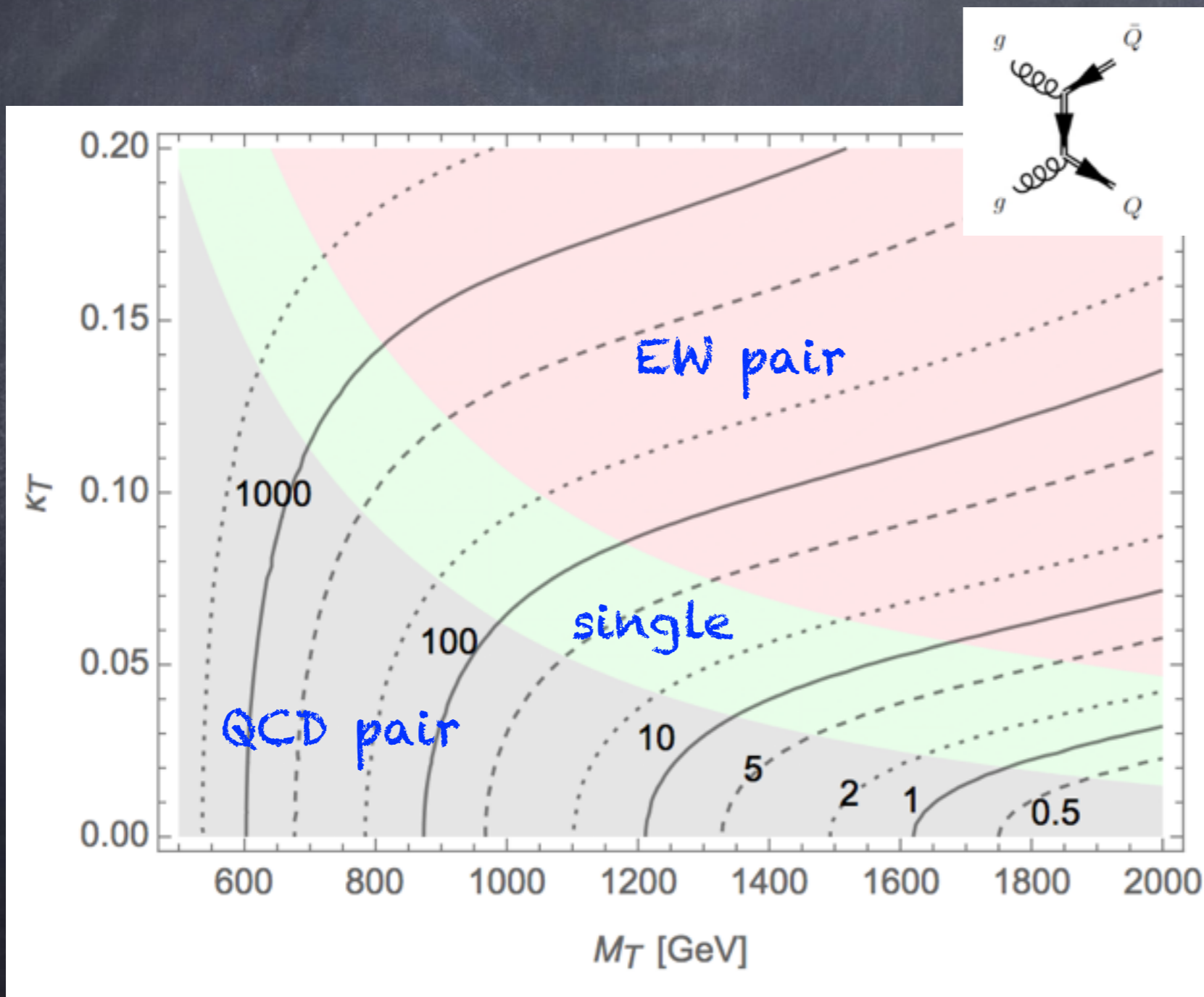
Single production: $pp \rightarrow q' + \{q, V, H\}$



The case of VLQs, aka Top partners

- BRs depend on the model.
- QCD pair-production is model-independent!
- Single production is not:
- EW pair production may be important!

The case of VLQs, aka Top partners



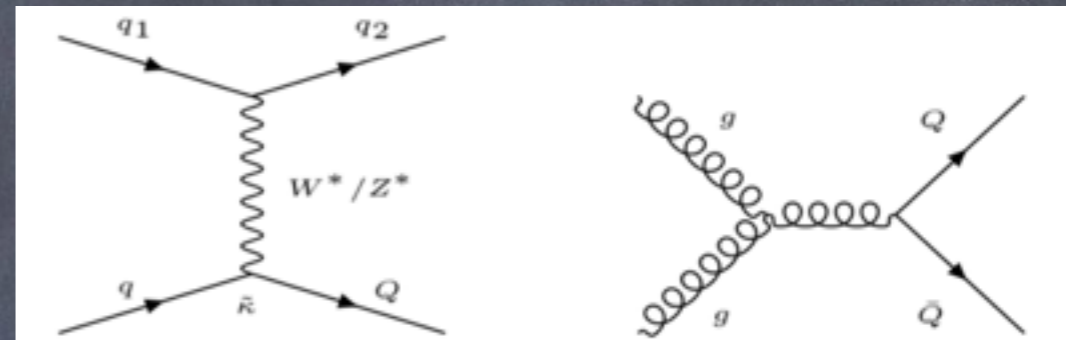
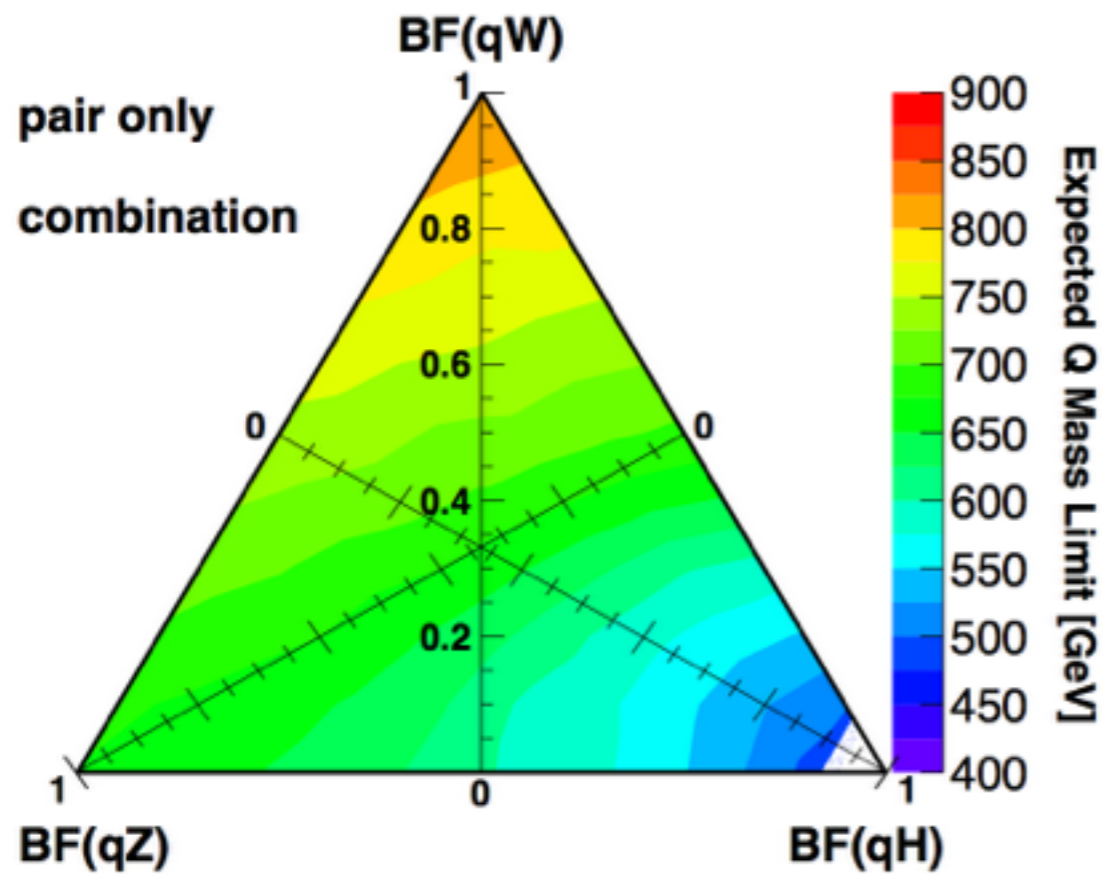
$$\kappa_T \frac{m_T}{v_{\text{SM}}} \bar{Q} h q + h.c.$$

The case of VLQs, aka Top partners

CMS-PAS-B2G-12-016

CMS Preliminary

19.7 fb⁻¹ (8 TeV)

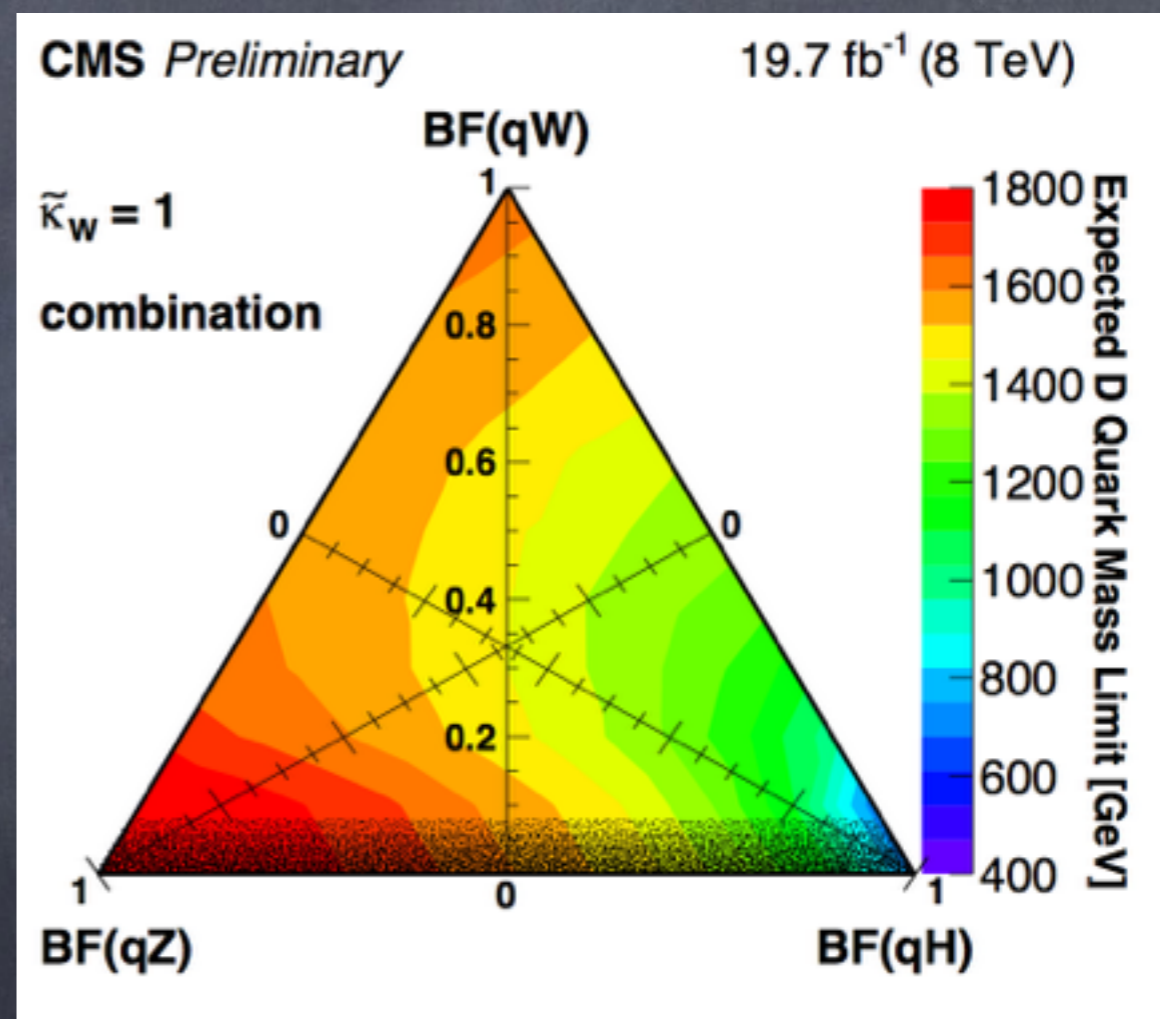
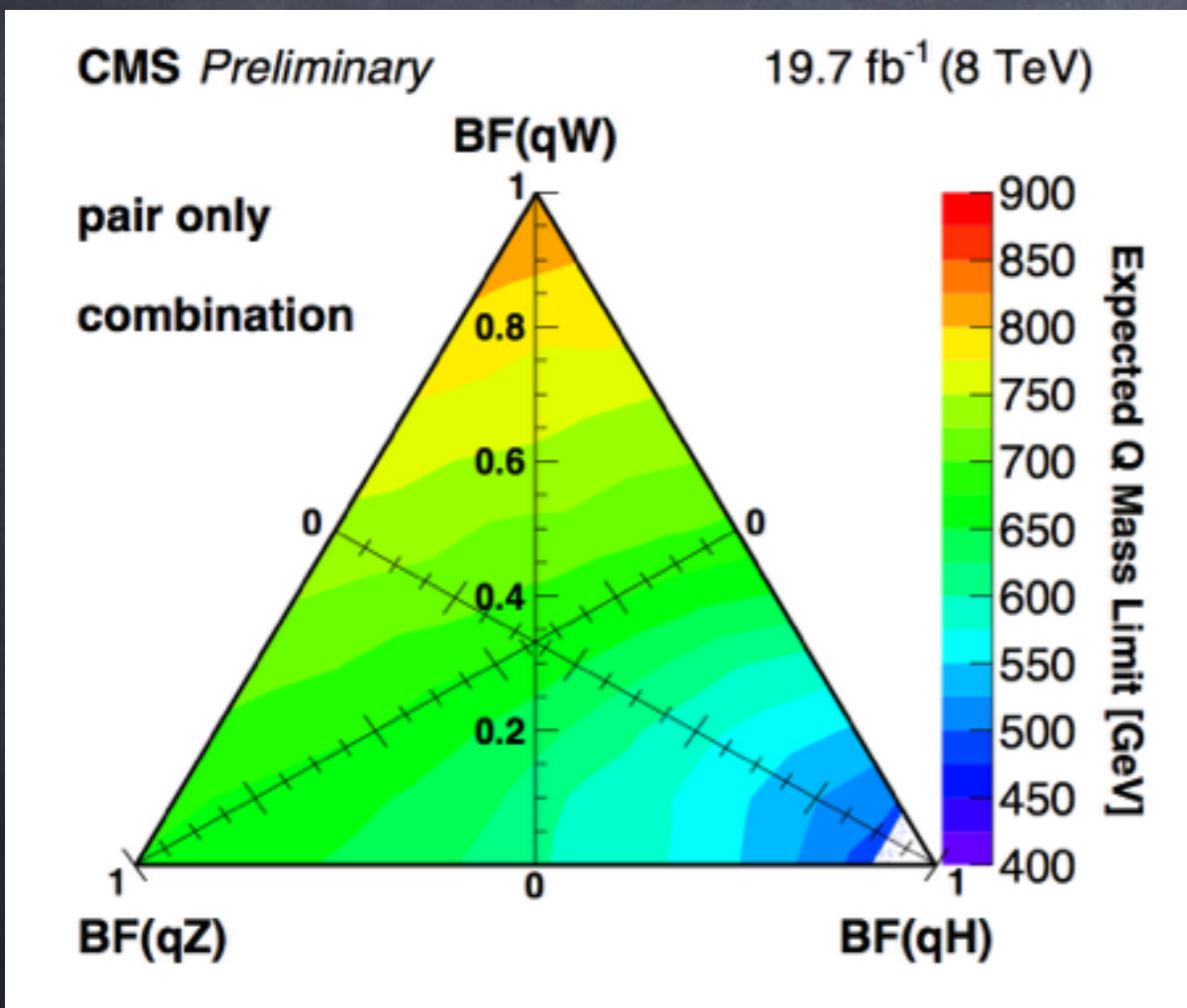


Full coverage of BR combinations.

Bounds are very sensitive to the BRs, ranging from less than 400 GeV to 800 GeV.

The case of VLQs, aka Top partners

CMS-PAS-B2G-12-016



pp \rightarrow Q jet added in.

The case of VLQs, aka Top partners

CMS-PAS-B2G-12-016

- Strategy allows full coverage.
- Small improvements needed (fix overall coupling instead of κ_W)
- QCD@NLO, other single channels, EW pair can be easily included.
- Are we really complete and fool-proof?

The case of VLQs, aka Top partners

- Exotic decay channels may be present:

E.g., non-minimal composite Higgs models.



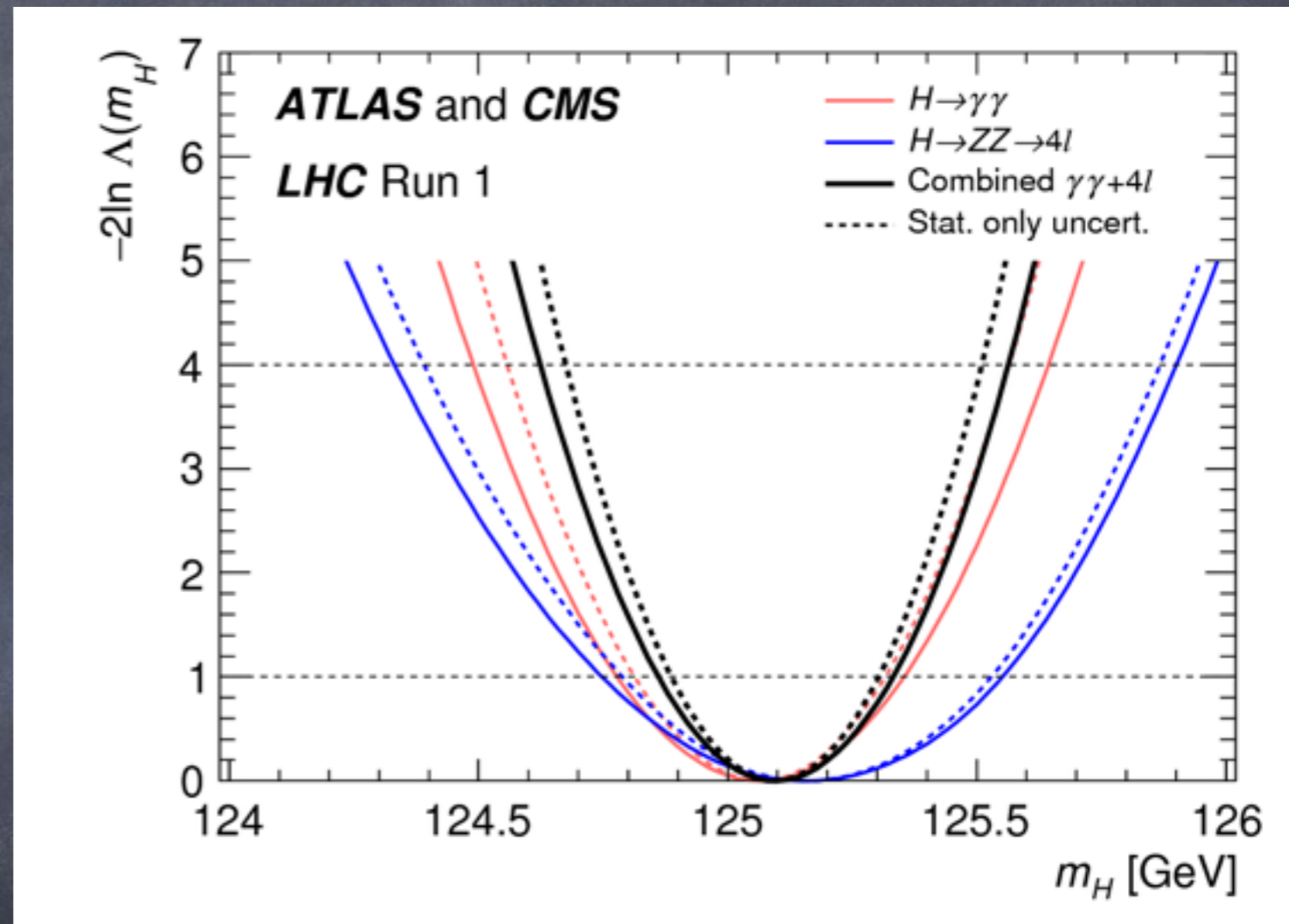
- Long-lived (i.e. MET)
- Decay to pair of gauge bosons (WZW anomaly)
- Decay into di-top

Theory: what have we learnt so far?

- There is a Higgs: what do we really know about it?
- Do we still need BSM physics?
- How can we rule out our favourite model(s)?

What do we know about H?

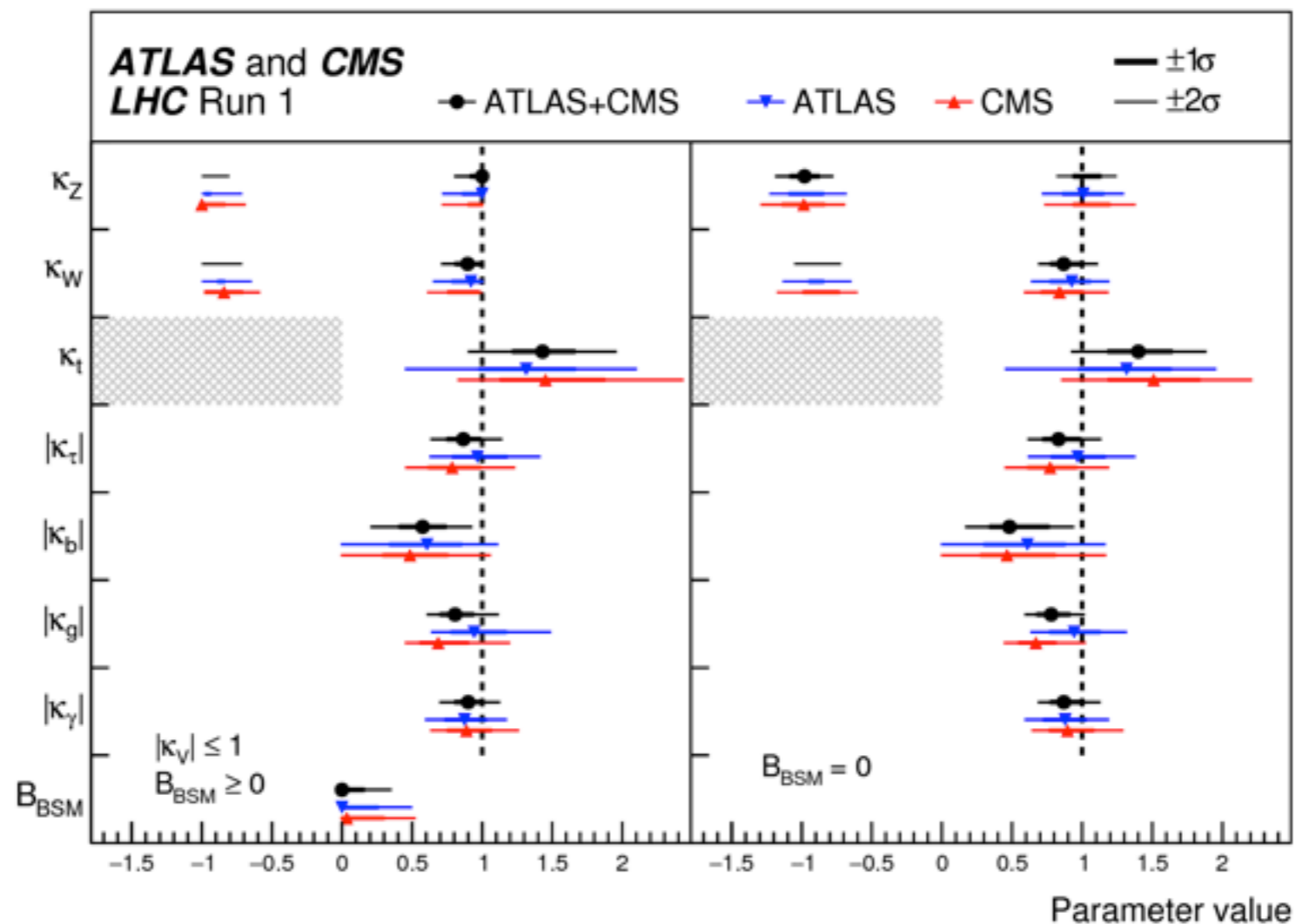
- The mass has been precisely measured!



$$125.09 \pm 0.21(\text{stat}) \pm 0.11(\text{syst})$$

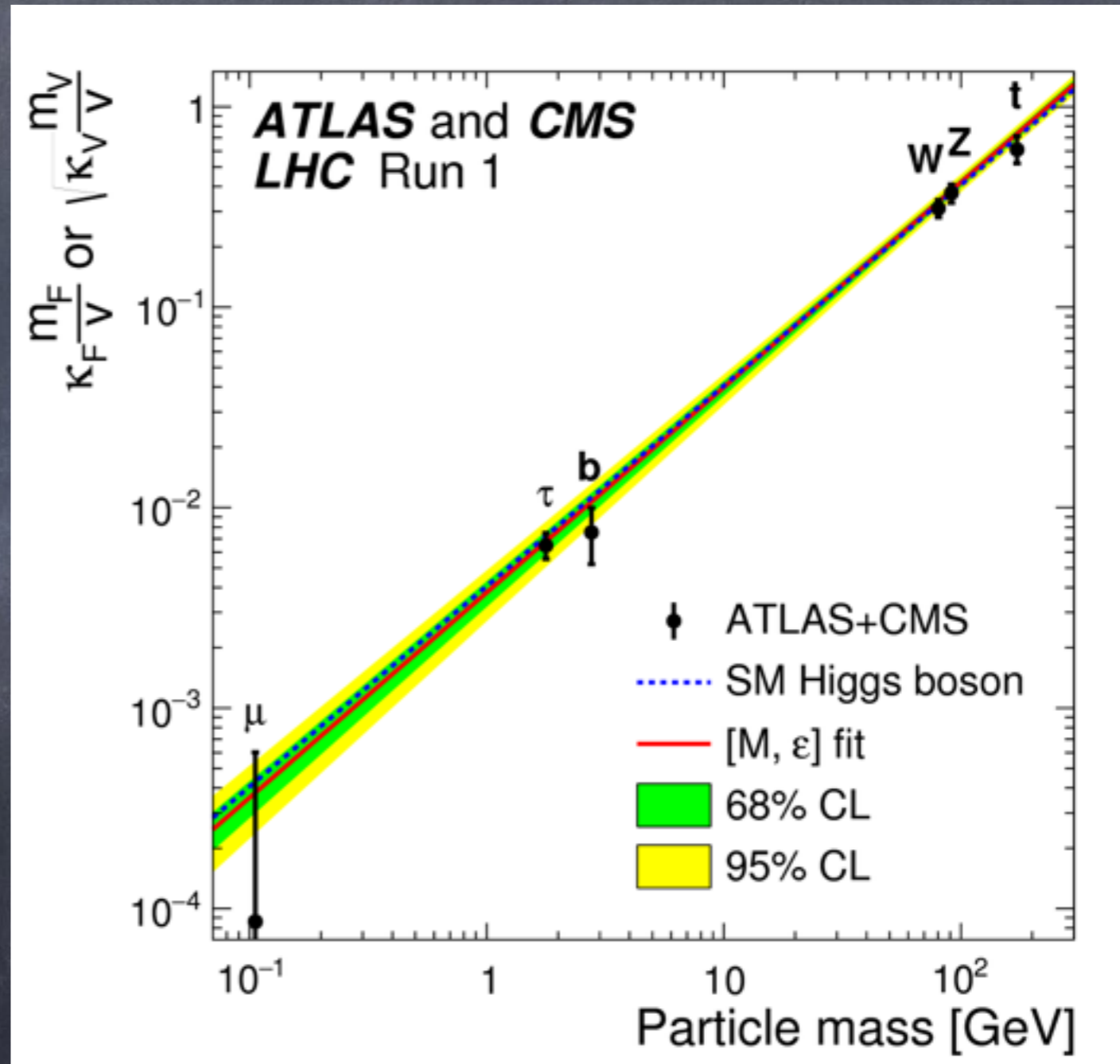
What do we know about H?

- The mass has been precisely measured!
- The couplings follow the SM expectations:



What do we know about H?

- The mass has been precisely measured!
- The couplings follow the SM expectations: being proportional to mass.



What do we know about H?

- The mass has been precisely measured!
- The couplings follow the SM expectations: being proportional to mass.
- The uncertainties are still large!
- Coupling measurements are always subject to model assumptions!!!

Run I

Decay channel	ATLAS+CMS	ATLAS	CMS
$\mu^{\gamma\gamma}$	1.14 $^{+0.19}_{-0.18}$ $\left(\begin{smallmatrix} +0.18 \\ -0.17 \end{smallmatrix} \right)$	1.14 $^{+0.27}_{-0.25}$ $\left(\begin{smallmatrix} +0.26 \\ -0.24 \end{smallmatrix} \right)$	1.11 $^{+0.25}_{-0.23}$ $\left(\begin{smallmatrix} +0.23 \\ -0.21 \end{smallmatrix} \right)$
μ^{ZZ}	1.29 $^{+0.26}_{-0.23}$ $\left(\begin{smallmatrix} +0.23 \\ -0.20 \end{smallmatrix} \right)$	1.52 $^{+0.40}_{-0.34}$ $\left(\begin{smallmatrix} +0.32 \\ -0.27 \end{smallmatrix} \right)$	1.04 $^{+0.32}_{-0.26}$ $\left(\begin{smallmatrix} +0.30 \\ -0.25 \end{smallmatrix} \right)$
μ^{WW}	1.09 $^{+0.18}_{-0.16}$ $\left(\begin{smallmatrix} +0.16 \\ -0.15 \end{smallmatrix} \right)$	1.22 $^{+0.23}_{-0.21}$ $\left(\begin{smallmatrix} +0.21 \\ -0.20 \end{smallmatrix} \right)$	0.90 $^{+0.23}_{-0.21}$ $\left(\begin{smallmatrix} +0.23 \\ -0.20 \end{smallmatrix} \right)$
$\mu^{\tau\tau}$	1.11 $^{+0.24}_{-0.22}$ $\left(\begin{smallmatrix} +0.24 \\ -0.22 \end{smallmatrix} \right)$	1.41 $^{+0.40}_{-0.36}$ $\left(\begin{smallmatrix} +0.37 \\ -0.33 \end{smallmatrix} \right)$	0.88 $^{+0.30}_{-0.28}$ $\left(\begin{smallmatrix} +0.31 \\ -0.29 \end{smallmatrix} \right)$
μ^{bb}	0.70 $^{+0.29}_{-0.27}$ $\left(\begin{smallmatrix} +0.29 \\ -0.28 \end{smallmatrix} \right)$	0.62 $^{+0.37}_{-0.37}$ $\left(\begin{smallmatrix} +0.39 \\ -0.37 \end{smallmatrix} \right)$	0.81 $^{+0.45}_{-0.43}$ $\left(\begin{smallmatrix} +0.45 \\ -0.43 \end{smallmatrix} \right)$
$\mu^{\mu\mu}$	0.1 $^{+2.5}_{-2.5}$ $\left(\begin{smallmatrix} +2.4 \\ -2.3 \end{smallmatrix} \right)$	-0.6 $^{+3.6}_{-3.6}$ $\left(\begin{smallmatrix} +3.6 \\ -3.6 \end{smallmatrix} \right)$	0.9 $^{+3.6}_{-3.5}$ $\left(\begin{smallmatrix} +3.3 \\ -3.2 \end{smallmatrix} \right)$

What do we know about H?

- Theoretical Modelling, i.e. the Standard Model Higgs

$$\mathcal{L}_{\text{Higgs}} = (D_\mu \phi)^\dagger (D^\mu \phi) + \mu^2 \phi^\dagger \phi - \lambda (\phi^\dagger \phi)^2$$

"wrong sign"

It well describes
the symmetry breaking,
but no dynamical
insight!

$$\tau^i = \frac{\sigma^i}{2} \quad \text{Pauli matrices}$$

$$\phi = e^{i\pi^i \tau^i} \cdot \begin{pmatrix} 0 \\ v + \frac{h}{\sqrt{2}} \end{pmatrix}$$

$$v = \frac{\mu}{\sqrt{2\lambda}} \sim 246 \text{ GeV}$$

What do we know about H?

$$\mathcal{L}_{\text{Higgs}} = (D_\mu \phi)^\dagger (D^\mu \phi) + \mu^2 \phi^\dagger \phi - \lambda (\phi^\dagger \phi)^2$$

- Custodial symmetry as a lucky accident:

$$\phi = \begin{pmatrix} \varphi_u \\ \varphi_d \end{pmatrix} \quad \tilde{\phi} = (i\sigma^2) \cdot \phi^* = \begin{pmatrix} \varphi_d^* \\ -\varphi_u^* \end{pmatrix}$$

Both transform as doublets of $SU(2)$
[pseudo-real irrep]

- We can rewrite the Lagrangian as:

$$\Phi = \begin{pmatrix} \tilde{\phi} & \phi \end{pmatrix} = \begin{pmatrix} \varphi_d^* & \varphi_u \\ -\varphi_u^* & \varphi_d \end{pmatrix} \quad \mathcal{L}_{\text{Higgs}} = \frac{1}{2} \text{Tr} [(D_\mu \Phi)^\dagger (D^\mu \Phi)] + \frac{\mu^2}{2} \text{Tr} [\Phi^\dagger \Phi] + \dots$$

$$\Phi \rightarrow U_L \cdot \Phi \cdot U_R^\dagger$$

uncovers a "hidden" invariance
under a global $SU(2)_L \times SU(2)_R$
broken to $SU(2)_D$ by the VEV

What do we know about H?

- Non-linear description:

$$\Sigma = e^{i\pi^i \tau^i} \cdot \begin{pmatrix} 0 \\ v \end{pmatrix} \quad \mathcal{L}_{NL} = f(h) (D_\mu \Sigma)^\dagger (D^\mu \Sigma) - V(h)$$

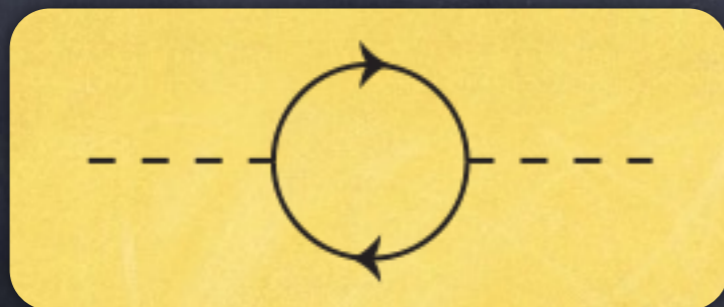
- Goldstones transform as a triplet of SU(2)_D.
- The coupling of h to gauge bosons ARE proportional to the mass (but not determined).
- However: trilinear h coupling is not determined!

Do we still need BSM?



We have a pretty good idea of the mechanism

But, we don't know how to protect it:



$$\delta m_h^2 \sim \frac{g^2}{16\pi^2} M_{\text{NPh}}^2$$

Do we still need BSM?

Fact:

we have been working on the same ideas for the last 30-40 years!

Supersymmetry

(Strong)
dynamics

Do we still need BSM?

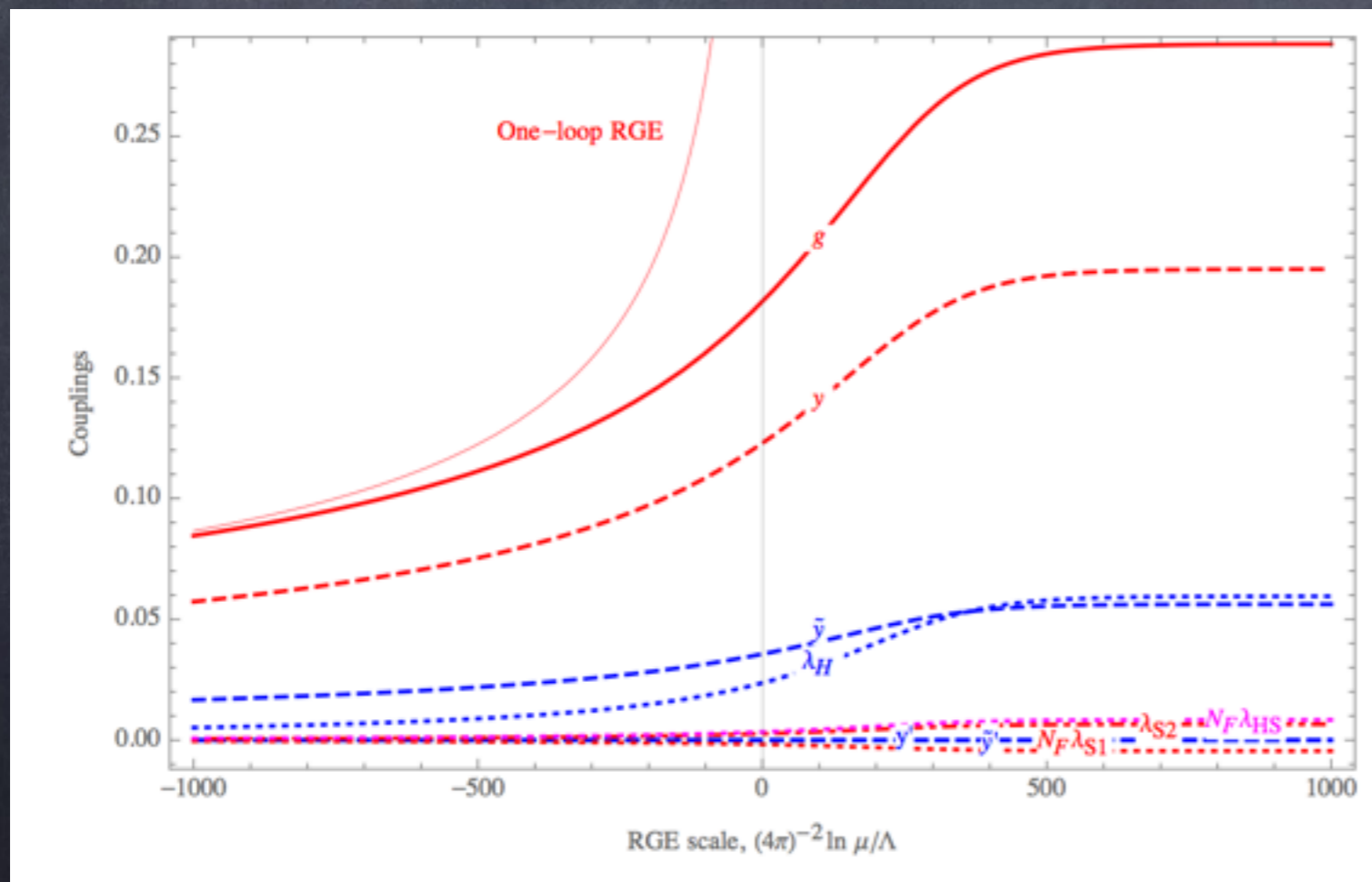
No judgement: I'm playing devil's advocate!

Recent New Ideas:

- Scale invariance: only true at classical level.
- Relaxion: classical field evolution imposed (tautologic fault!)
- Warped extra dimensions: exponentiation of large scale hierarchies.
- Asymptotic Safe theories: still under study (see 1701.01453) - no realistic example!

Do we still need BSM?

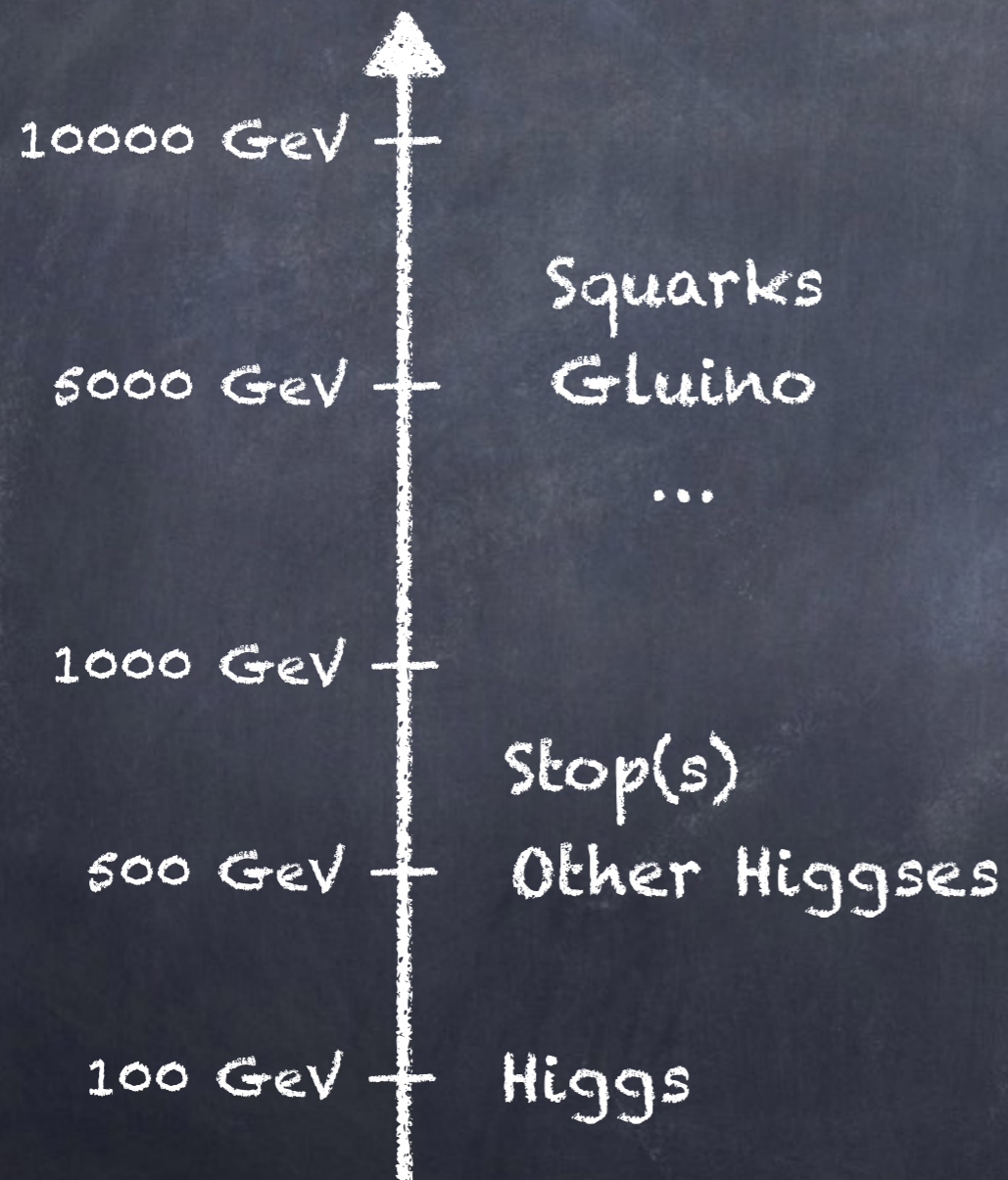
- Asymptotic Safe theories: still under study (see 1701.01453)



Interacting UV fixed point!

Claim that scalar masses can be natural, i.e. quantum corrections proportional to the mass itself!

How can we rule out our favourite model(s)?



"Natural" SUSY spectrum

- Higgs decoupling limit: SM-like light Higgs
- Light-ish stops
- Other sparticles above TeV scale

How can we rule out our favourite model(s)?



Compositeness



- Higgs close to SM-like

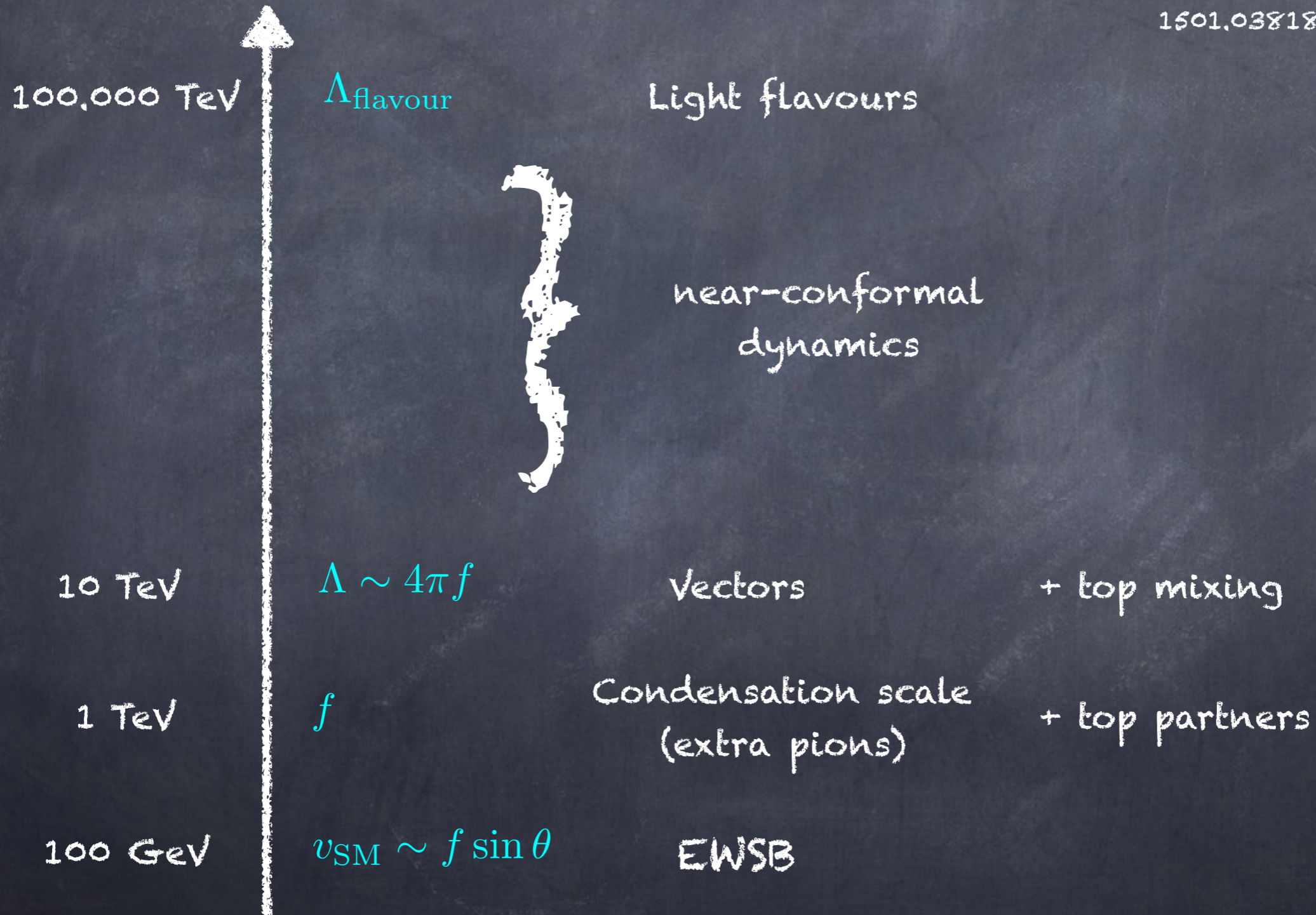
$$\frac{g_{hVV}}{g_{hVV}^{SM}} \sim \cos \theta \sim 1 - \frac{1}{2} \sin^2 \theta$$

$$\sin \theta \sim \frac{v}{f} < 0.2$$

(EW precision)

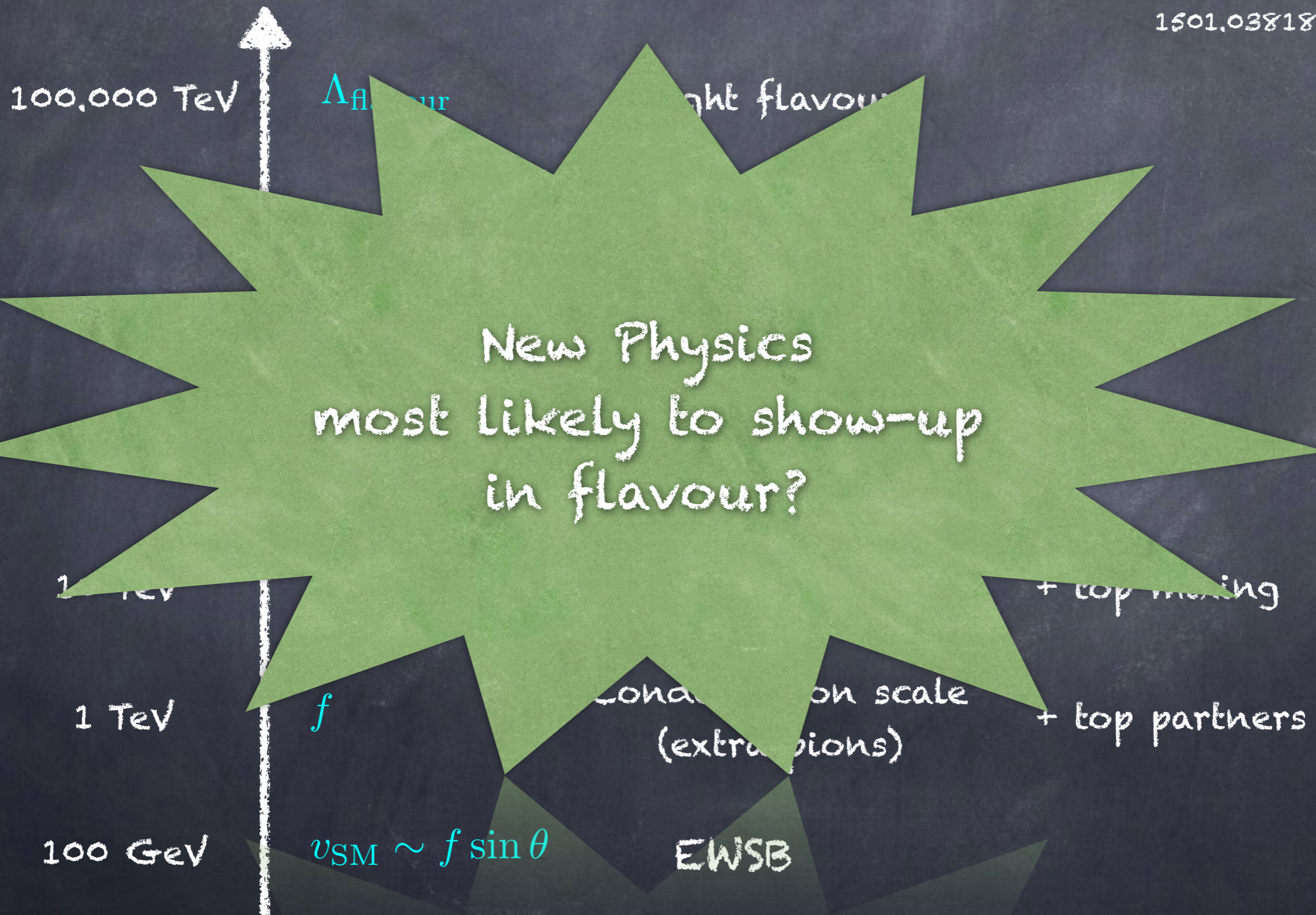
The hot potato: flavour!

1501.03818, ...



The hot potato: flavour!

1501.03818, ...



A fermionic theory of top partners

\mathcal{G}_{TC} : rep R rep R' 1312, 5330, 1604, 06467
 Q χ $T' = QQ\chi$ or $Q\chi\chi$

SM : EW colour + hypercharge

global : $\langle QQ \rangle \neq 0$



pNGB Higgs
 DM?

a) $\langle \chi\chi \rangle \neq 0$

coloured pNGBs
 di-boson

b) $\langle \chi\chi \rangle = 0$

Light top partners
 from \ddagger Hooft anomaly
 conditions?

Global symmetries

More precisely, the global symmetries are:

$$SU(N_Q) \times SU(N_X) \times U(1)_Q \times U(1)_X$$

WZW term:

$$\mathcal{L} \supset \frac{g_i^2}{32\pi^2} \frac{\kappa_i}{f_a} a \epsilon^{\mu\nu\alpha\beta} G_{\mu\nu}^i G_{\alpha\beta}^i,$$

Coefficients depend
on the underlying dynamics!

$G = A, W, Z, g$!!!

1512.04508

Anomalous $U(1) \rightarrow$ heavy η'

Orthogonal $U(1) \rightarrow$ pNGB a

Decays and production
only via WZW anomaly.

Model zoology

G_{HC}	ψ	χ	Restrictions	$-q_\chi/q_\psi$	Y_χ	Non Conformal	Model Name
Real			SU(5)/SO(5) \times SU(6)/SO(6)				
$SO(N_{\text{HC}})$	$5 \times \mathbf{S}_2$	$6 \times \mathbf{F}$	$N_{\text{HC}} \geq 55$	$\frac{5(N_{\text{HC}}+2)}{6}$	1/3	/	
$SO(N_{\text{HC}})$	$5 \times \mathbf{Ad}$	$6 \times \mathbf{F}$	$N_{\text{HC}} \geq 15$	$\frac{5(N_{\text{HC}}-2)}{6}$	1/3	/	
$SO(N_{\text{HC}})$	$5 \times \mathbf{F}$	$6 \times \mathbf{Spin}$	$N_{\text{HC}} = 7, 9$	$\frac{5}{6}, \frac{5}{12}$	1/3	$N_{\text{HC}} = 7, 9$	M1, M2
$SO(N_{\text{HC}})$	$5 \times \mathbf{Spin}$	$6 \times \mathbf{F}$	$N_{\text{HC}} = 7, 9$	$\frac{5}{6}, \frac{5}{3}$	2/3	$N_{\text{HC}} = 7, 9$	M3, M4
Real			Pseudo-Real		SU(5)/SO(5) \times SU(6)/Sp(6)		
$Sp(2N_{\text{HC}})$	$5 \times \mathbf{Ad}$	$6 \times \mathbf{F}$	$2N_{\text{HC}} \geq 12$	$\frac{5(N_{\text{HC}}+1)}{3}$	1/3	/	
$Sp(2N_{\text{HC}})$	$5 \times \mathbf{A}_2$	$6 \times \mathbf{F}$	$2N_{\text{HC}} \geq 4$	$\frac{5(N_{\text{HC}}-1)}{3}$	1/3	$2N_{\text{HC}} = 4$	M5
$SO(N_{\text{HC}})$	$5 \times \mathbf{F}$	$6 \times \mathbf{Spin}$	$N_{\text{HC}} = 11, 13$	$\frac{5}{24}, \frac{5}{48}$	1/3	/	
Real			Complex		SU(5)/SO(5) \times SU(3) ² /SU(3)		
$SU(N_{\text{HC}})$	$5 \times \mathbf{A}_2$	$3 \times (\mathbf{F}, \bar{\mathbf{F}})$	$N_{\text{HC}} = 4$	$\frac{5}{3}$	1/3	$N_{\text{HC}} = 4$	M6
$SO(N_{\text{HC}})$	$5 \times \mathbf{F}$	$3 \times (\mathbf{Spin}, \bar{\mathbf{Spin}})$	$N_{\text{HC}} = 10, 14$	$\frac{5}{12}, \frac{5}{48}$	1/3	$N_{\text{HC}} = 10$	M7
Pseudo-Real			Real		SU(4)/Sp(4) \times SU(6)/SO(6)		
$Sp(2N_{\text{HC}})$	$4 \times \mathbf{F}$	$6 \times \mathbf{A}_2$	$2N_{\text{HC}} \leq 36$	$\frac{1}{3(N_{\text{HC}}-1)}$	2/3	$2N_{\text{HC}} = 4$	M8
$SO(N_{\text{HC}})$	$4 \times \mathbf{Spin}$	$6 \times \mathbf{F}$	$N_{\text{HC}} = 11, 13$	$\frac{8}{3}, \frac{16}{3}$	2/3	$N_{\text{HC}} = 11$	M9
Complex			Real		SU(4) ² /SU(4) \times SU(6)/SO(6)		
$SO(N_{\text{HC}})$	$4 \times (\mathbf{Spin}, \bar{\mathbf{Spin}})$	$6 \times \mathbf{F}$	$N_{\text{HC}} = 10$	$\frac{8}{3}$	2/3	$N_{\text{HC}} = 10$	M10
$SU(N_{\text{HC}})$	$4 \times (\mathbf{F}, \bar{\mathbf{F}})$	$6 \times \mathbf{A}_2$	$N_{\text{HC}} = 4$	$\frac{2}{3}$	2/3	$N_{\text{HC}} = 4$	M11
Complex			Complex		SU(4) ² /SU(4) \times SU(3) ² /SU(3)		
$SU(N_{\text{HC}})$	$4 \times (\mathbf{F}, \bar{\mathbf{F}})$	$3 \times (\mathbf{A}_2, \bar{\mathbf{A}}_2)$	$N_{\text{HC}} \geq 5$	$\frac{4}{3(N_{\text{HC}}-2)}$	2/3	$N_{\text{HC}} = 5$	M12
$SU(N_{\text{HC}})$	$4 \times (\mathbf{F}, \bar{\mathbf{F}})$	$3 \times (\mathbf{S}_2, \bar{\mathbf{S}}_2)$	$N_{\text{HC}} \geq 5$	$\frac{4}{3(N_{\text{HC}}+2)}$	2/3	/	
$SU(N_{\text{HC}})$	$4 \times (\mathbf{A}_2, \bar{\mathbf{A}}_2)$	$3 \times (\mathbf{F}, \bar{\mathbf{F}})$	$N_{\text{HC}} = 5$	4	2/3	/	

1604.06467

Model zoology

G_{HC}	ψ	χ	Restrictions	$-q_\chi/q_\psi$	Y_χ	Non Conformal	Model Name
	Pseudo-Real	Real	SU(4)/Sp(4) \times SU(6)/SO(6)				
$Sp(2N_{\text{HC}})$	$4 \times \mathbf{F}$	$6 \times \mathbf{A}_2$	$2N_{\text{HC}} \leq 36$	$\frac{1}{3(N_{\text{HC}}-1)}$	$2/3$	$2N_{\text{HC}} = 4$	M8
$SO(N_{\text{HC}})$	$4 \times \mathbf{Spin}$	$6 \times \mathbf{F}$	$N_{\text{HC}} = 11, 13$	$\frac{8}{3}, \frac{16}{3}$	$2/3$	$N_{\text{HC}} = 11$	M9

Defines $\tan \zeta$

Theory confines!

$$T' = \psi\psi\chi$$

Note: there is enough baryons to give mass to the top (and bottom) only!

Example of predictions: di-boson resonances

1610.06591

	Pseudo-Real	Real	SU(4)/Sp(4) × SU(6)/SO(6)				
$Sp(2N_{\text{HC}})$	$4 \times \mathbf{F}$	$6 \times \mathbf{A}_2$	$2N_{\text{HC}} \leq 36$	$\frac{1}{3(N_{\text{HC}}-1)}$	2/3	$2N_{\text{HC}} = 4$	M8
$SO(N_{\text{HC}})$	$4 \times \mathbf{Spin}$	$6 \times \mathbf{F}$	$N_{\text{HC}} = 11, 13$	$\frac{8}{3}, \frac{16}{3}$	2/3	$N_{\text{HC}} = 11$	M9

The EFT is the same!

Numerical value of couplings:

Model		κ_g	$\frac{\kappa_W}{\kappa_g}$	$\frac{\kappa_B}{\kappa_g}$	$\frac{C_t}{\kappa_g} (2, 0)$	$\frac{C_t}{\kappa_g} (0, 2)$	$\tan \zeta$
M8	a	-0.77(-0.39)	-1.2(-2.5)	1.5(0.17)	-1.2(-2.5)	0.40(0.40)	-0.41
	η'	1.9(2.0)	0.20(0.096)	2.9(2.8)	0.20(0.096)	0.40(0.40)	
	π_8	7.1	0	1.3	0	0.40	
M9	a	-4.3(-2.7)	-0.55(-2.4)	2.1(0.26)	-0.068(-0.30)	0.18(0.18)	-3.26
	η'	1.3(3.6)	5.8(1.3)	8.5(4.0)	0.73(0.16)	0.18(0.18)	
	π_8	16.	0	1.3	0	0.18	

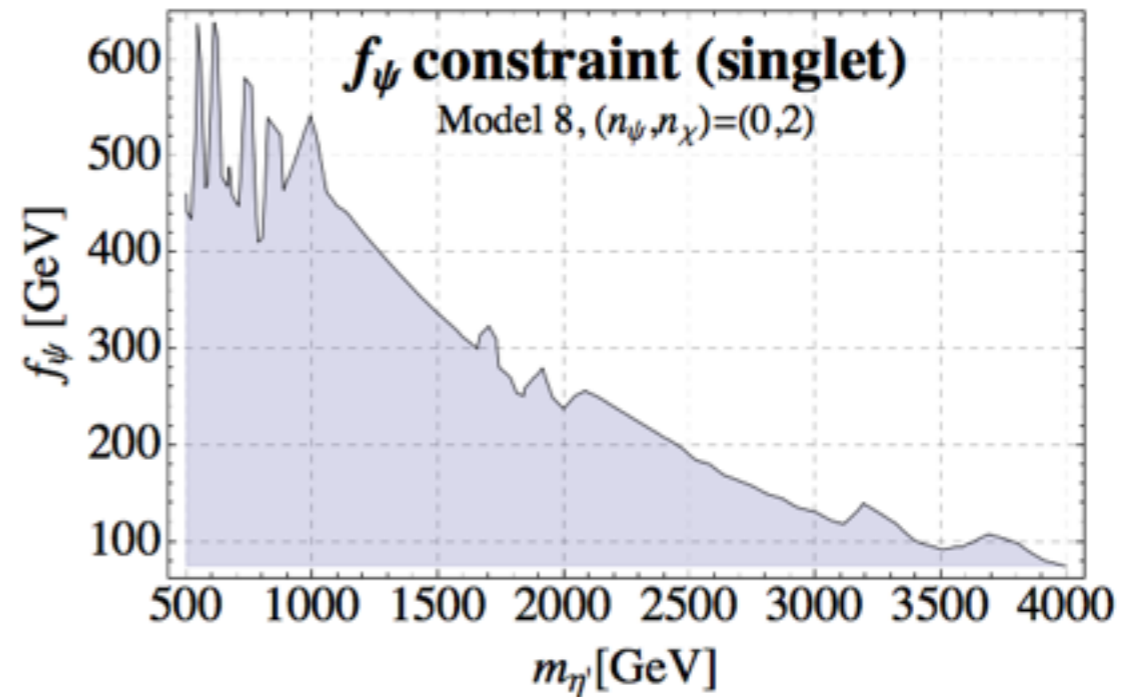
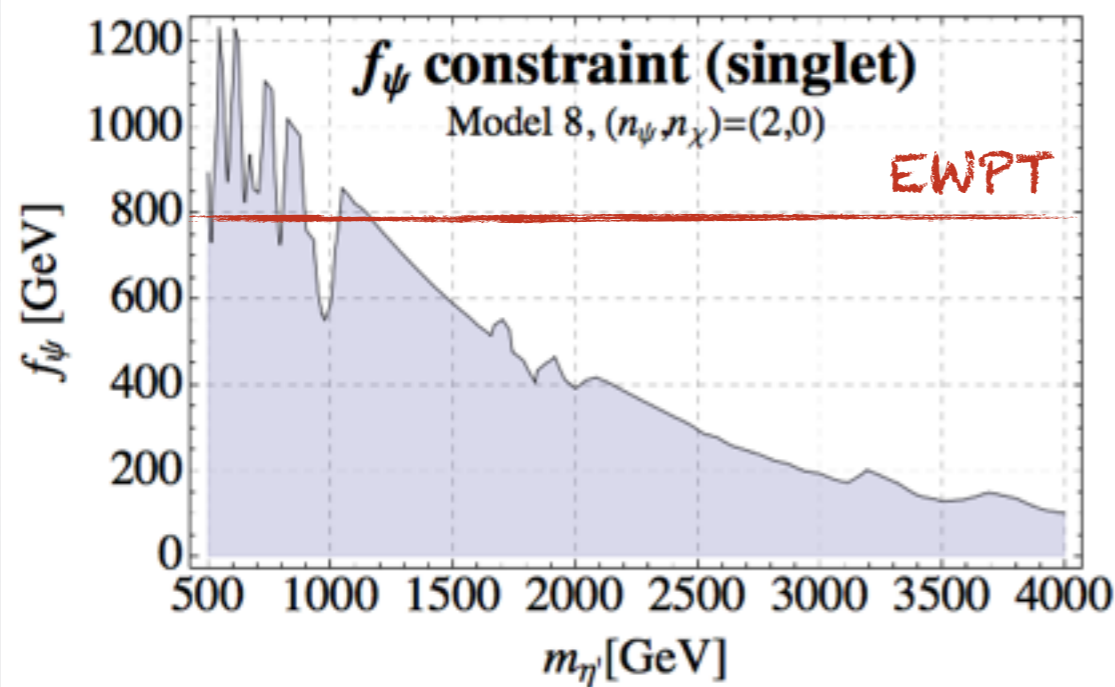
Assuming $f_a = f_\psi = f_\chi$

Model M8

1610.06591

"a" too light for the LHC!

$$\left. \frac{m_a}{m_{\eta'}} \right|_{\max} = 0.20$$



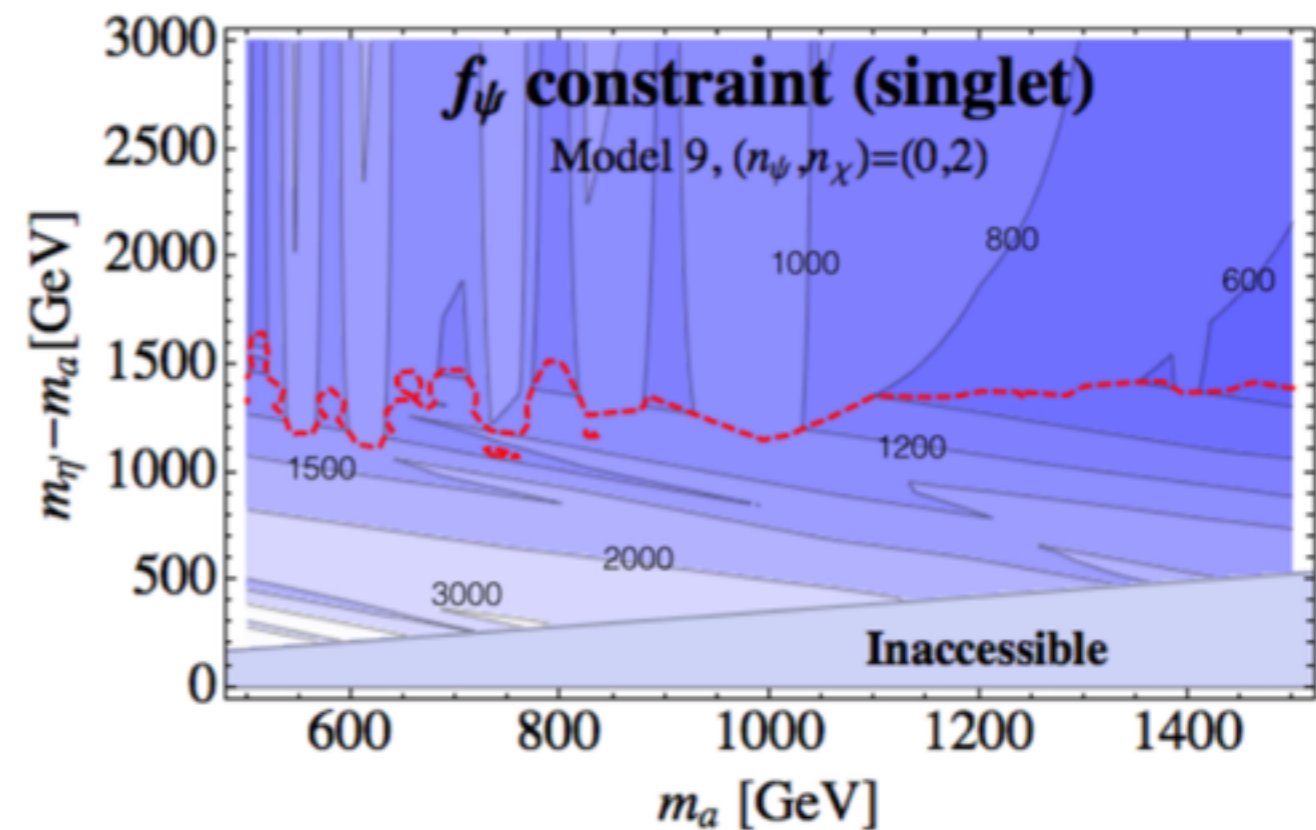
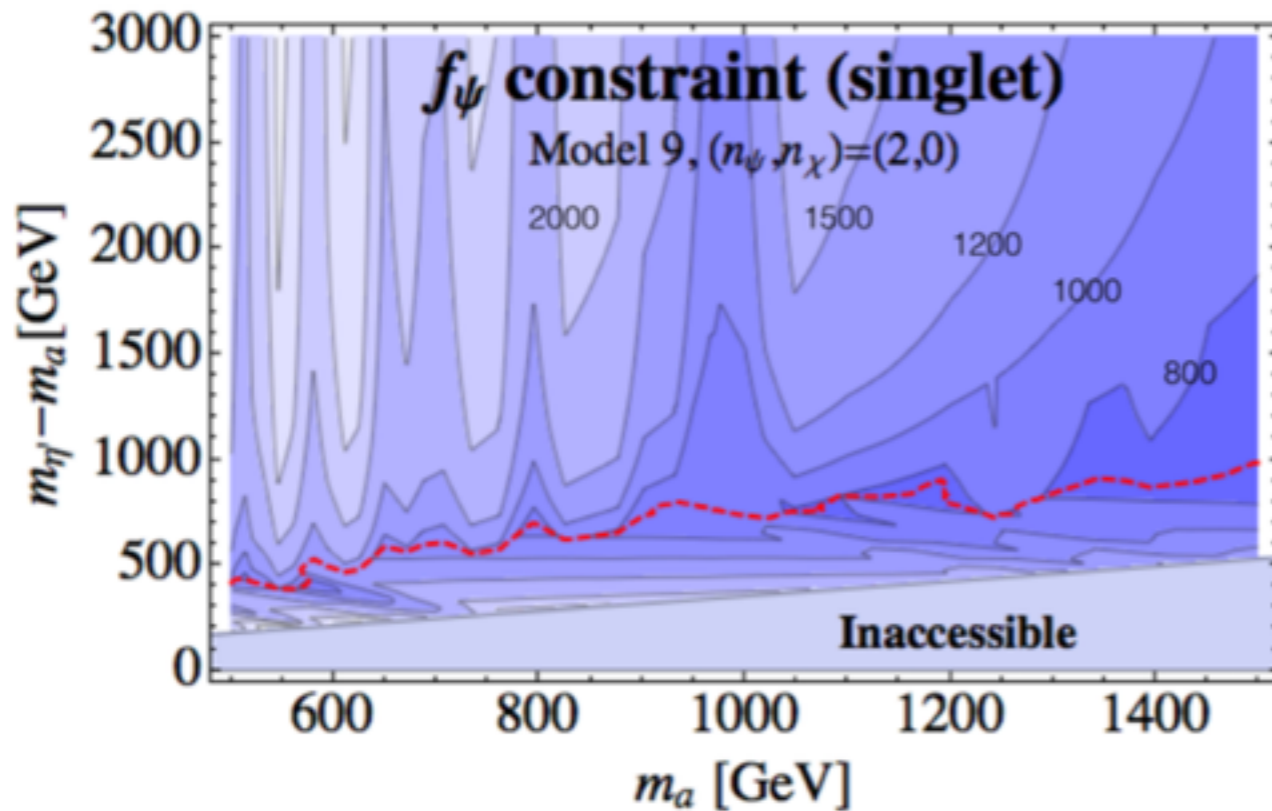
For light masses:
bounds competitive
with EW precision!

Larger top couplings:
reduced diboson rates
due to $t\bar{t}$ BR.

Model M9

$$\left. \frac{m_a}{m_{\eta'}} \right|_{\max} = 0.74$$

1610.06591



Above red line, bound driven by "a"!

Bounds stronger than EW precision
in most of the parameter space!

Let's bet:
which model will be ruled
out first?

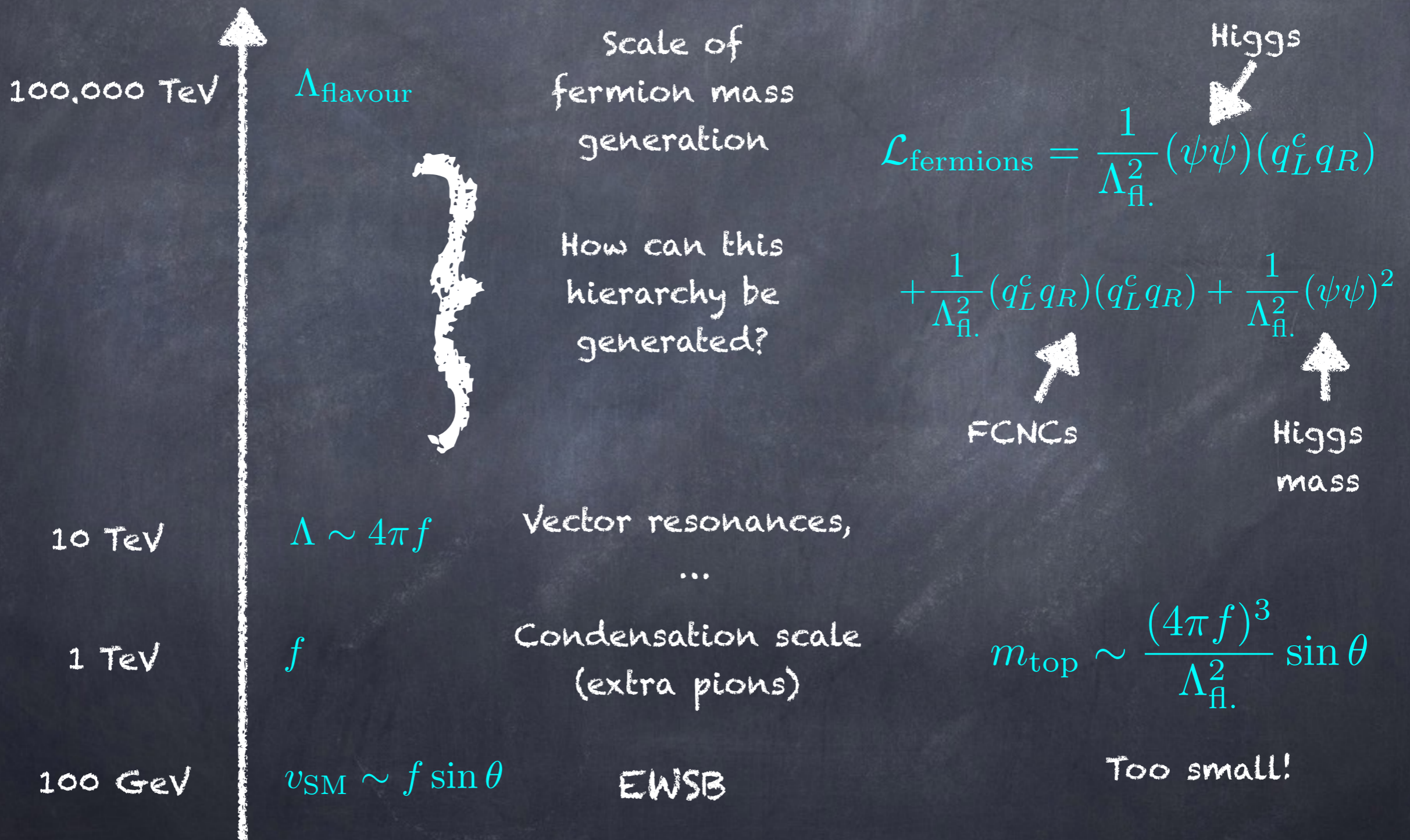
Supersymmetry?

Composite
(pNGB) Higgs?

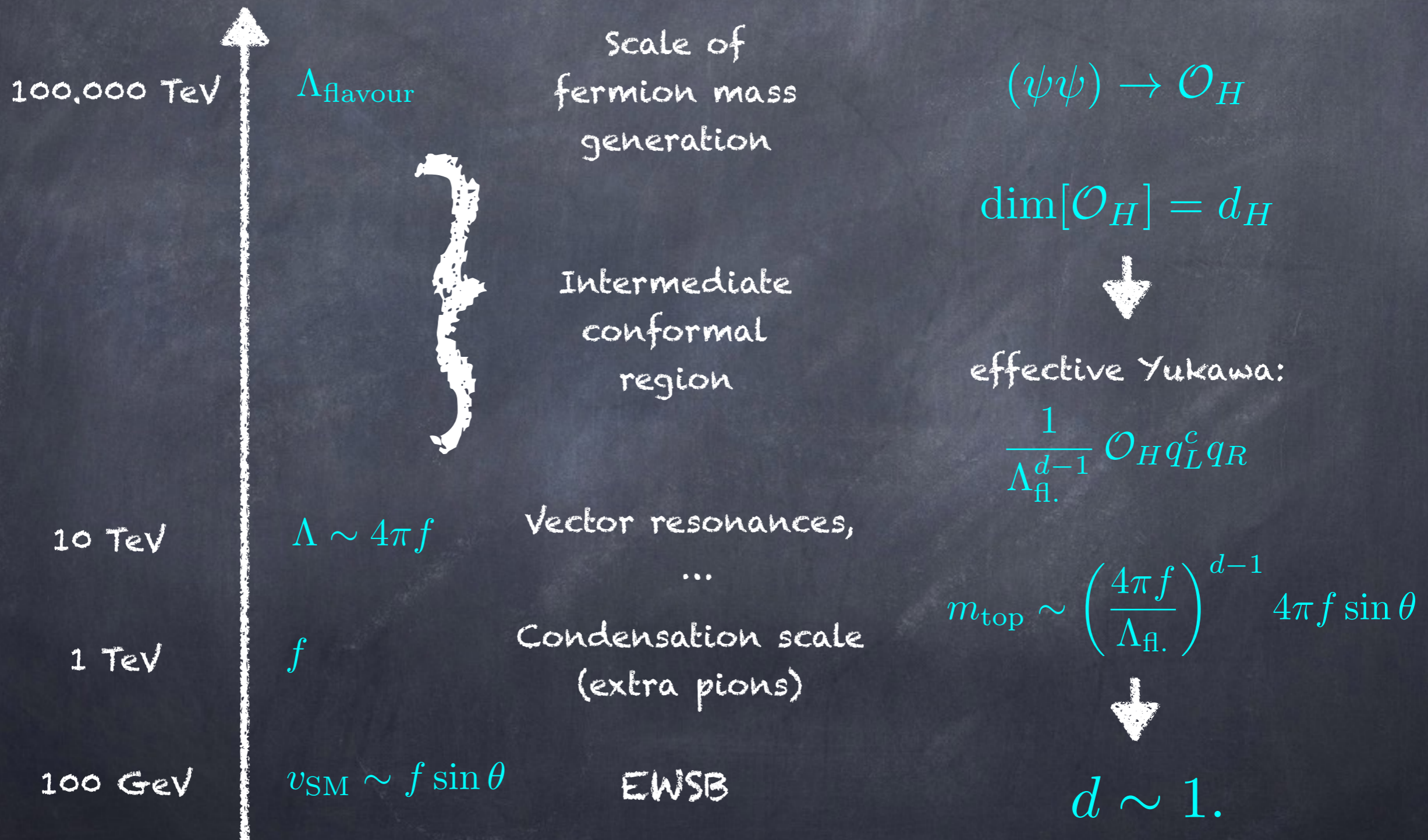
(please, add
your model)

Bonus

The hot potato: flavour!



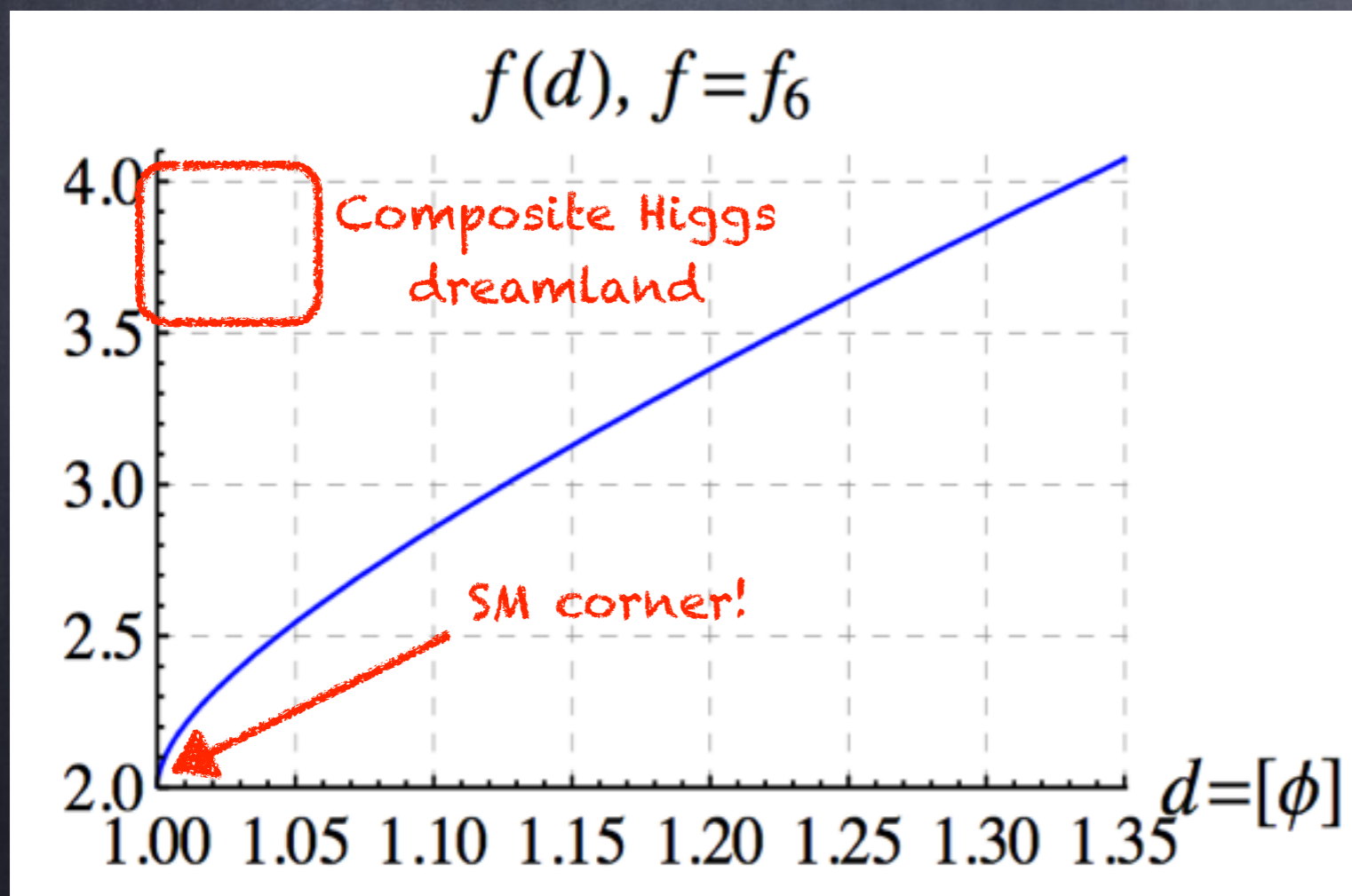
The hot potato: flavour!



A no-go theorem?

Bounds on the dimensions of scalar operators can be extracted using bootstrap techniques!

Rattazzi, Rychkov, Tonni, Vichi 0807.0004



$$\phi \equiv \mathcal{O}_H$$

$$d[\phi^2]_{\min} < f(d)$$

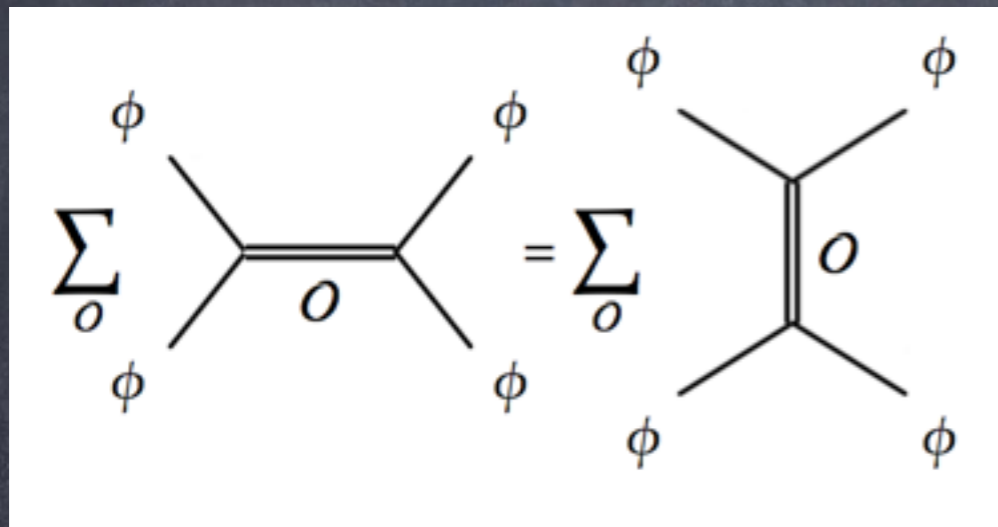


Higgs mass operator!

$$\Delta m_H^2 \sim \left(\frac{4\pi f}{\Lambda_{\text{fl.}}} \right)^{d-4} f^2$$

A no-go theorem?

Q: does the bound apply to the Higgs?



$$(\psi^i \psi^j) = \phi^{ij}$$

The scalar operator has
flavour indices:
many bi-linear ops appear!

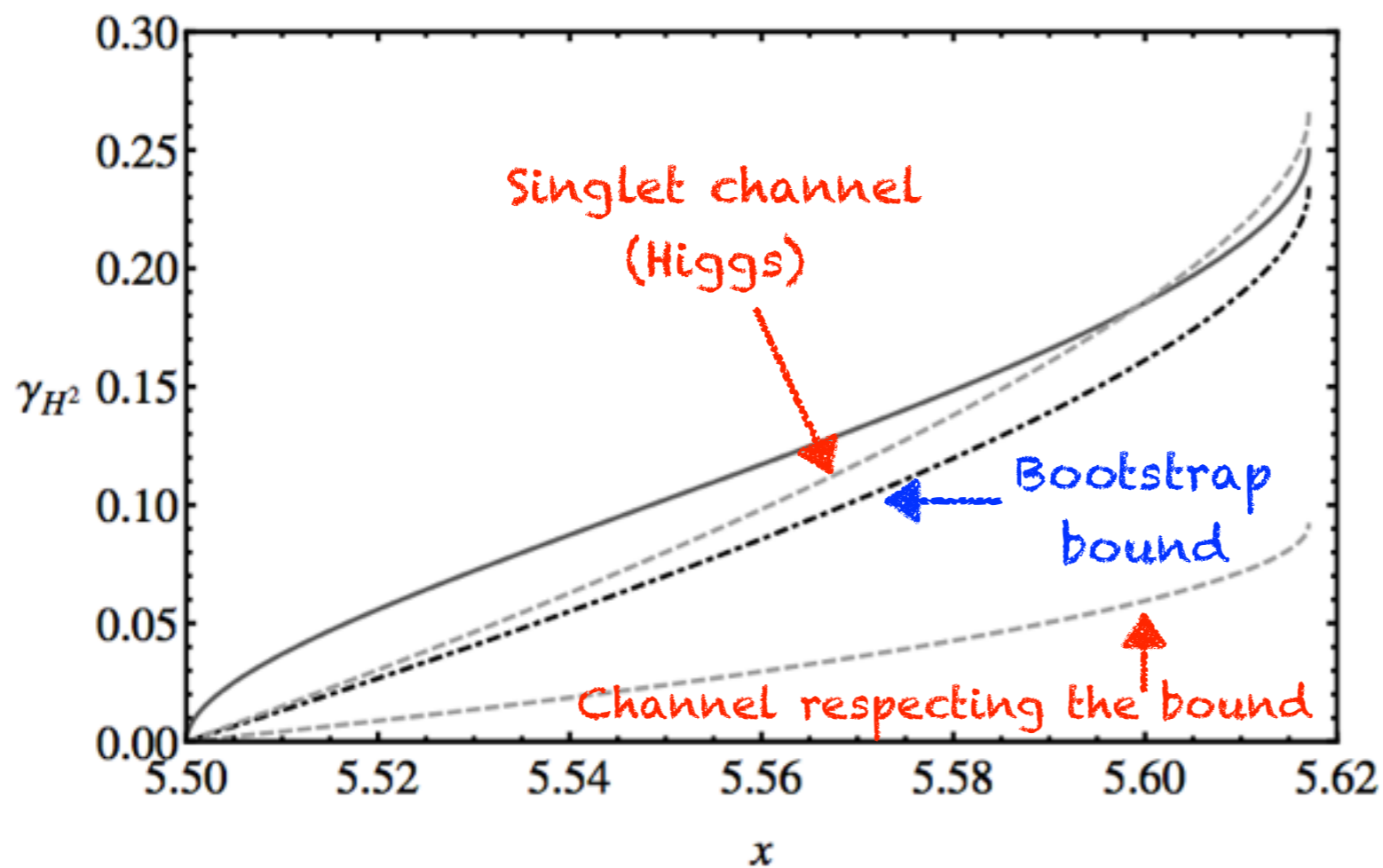


The bound applies to the one
with lowest dimension!

A no-go theorem? No...

Q: does the bound apply to the Higgs?

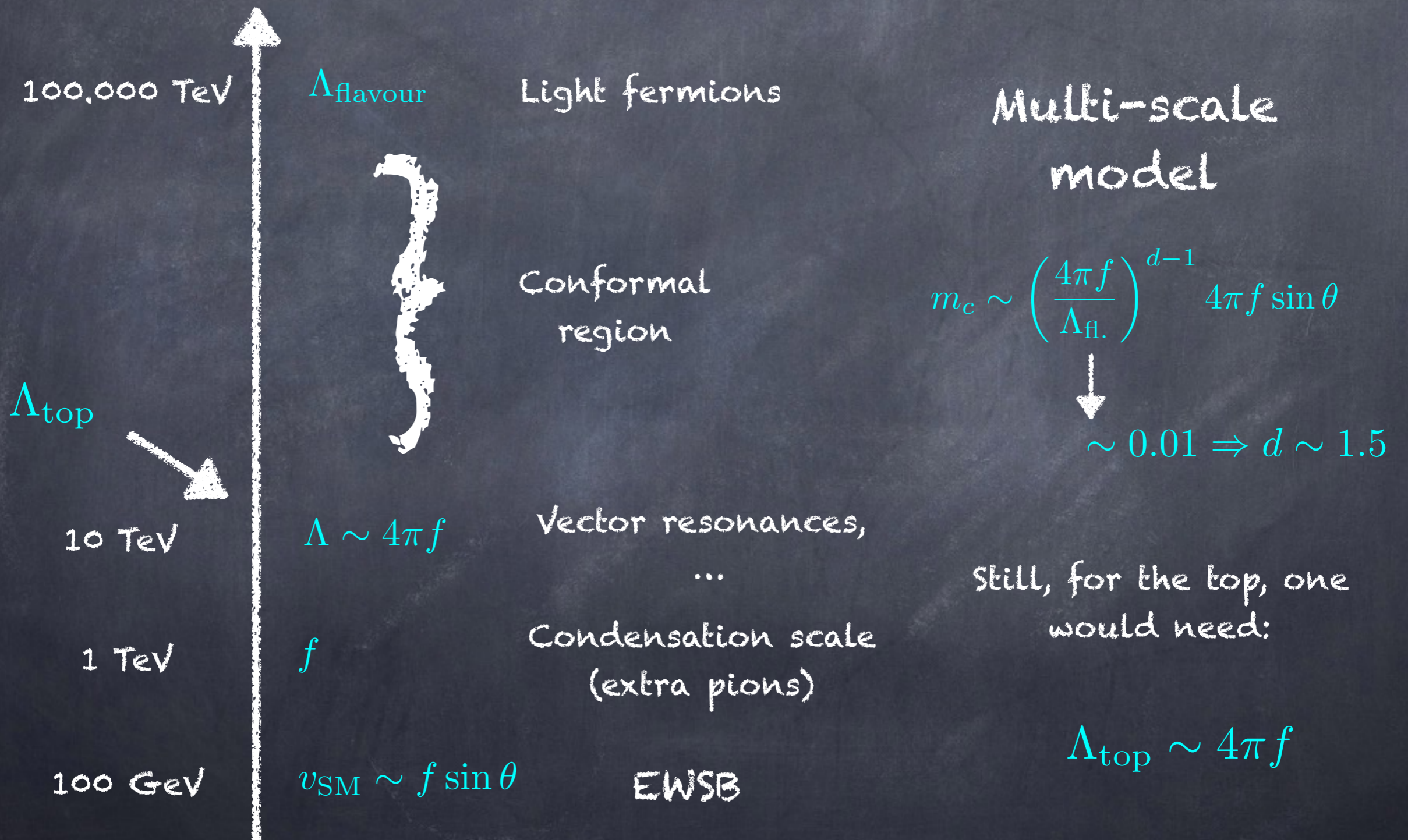
Antipin, Mølgaard, Sannino 1406.6166



Gauge-Yukawa theory
with weakly-coupled
fixed point.

Dimensions are calculable
(but small...)

The hot potato: flavour!



The partial compositeness paradigm

Kaplan Nucl.Phys. B365 (1991) 259

$$\frac{1}{\Lambda_{\text{fl.}}^{d-1}} \mathcal{O}_H q_L^c q_R \quad \Delta m_H^2 \sim \left(\frac{4\pi f}{\Lambda_{\text{fl.}}} \right)^{d-4} f^2 \quad \text{Both irrelevant if}$$

we assume: $d_H > 1$ $d_{H^2} > 4$

Let's postulate the existence of fermionic operators:

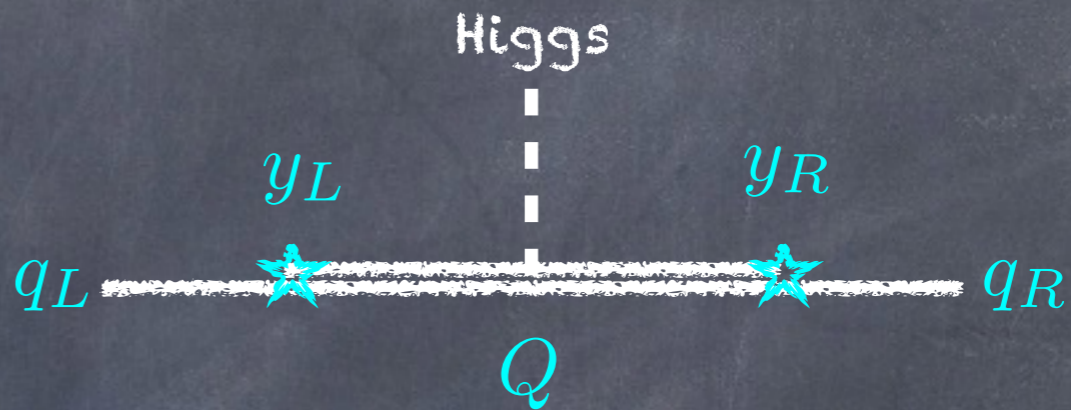
$$\frac{1}{\Lambda_{\text{fl.}}^{d_F-5/2}} (\tilde{y}_L q_L \mathcal{F}_L + \tilde{y}_R q_R \mathcal{F}_R)$$

This dimension is not related to the Higgs!

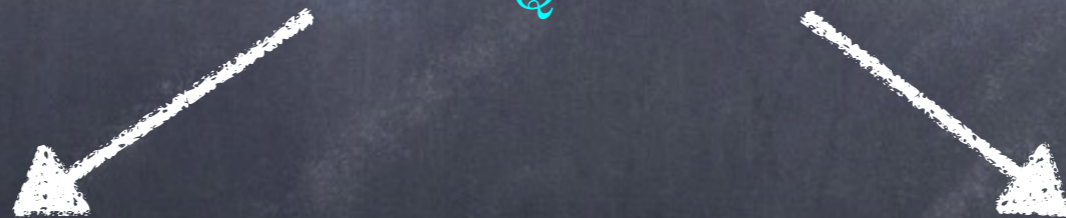
$$f(y_L q_L Q_L + y_R q_R Q_R) \quad \text{with} \quad y_{L/R} f \sim \left(\frac{4\pi f}{\Lambda_{\text{fl.}}} \right)^{d_F-5/2} 4\pi f$$

The partial compositeness paradigm

$$f(y_L q_L Q_L + y_R q_R Q_R)$$



$$m_q \sim \frac{y_L y_R f^2}{M_Q^2} f \sin \theta$$



$$M_Q \sim f \Rightarrow y_L, y_R \sim 1$$

Top can cancel top loop,
PUV

$$M_Q \sim 4\pi f \Rightarrow y_L, y_R \sim 4\pi$$