

# New deeply $\gamma$ and $\pi^0$ electroproduction results from Hall A at JLab

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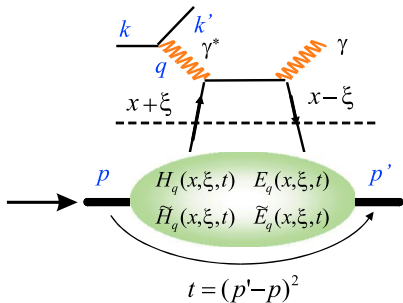
IPN-Orsay, CNRS/IN2P3 (France)

Nucleon and Resonance Structure with Hard Exclusive Processes  
May 29–31, 2017

# Outline

- 1 Very brief experimental introduction to DVCS
- 2 Recent DVCS results from Hall A at JLab
  - DVCS: beam energy dependence of the cross section (arxiv:1703.09442)
  - $\pi^0$ : off the neutron  $\rightarrow$  flavor separation (arXiv:1702.00835)
- 3 Outlook
  - Jefferson Lab at 12 GeV: Hall A  $\rightarrow$  C programs

# Deeply Virtual Compton Scattering (DVCS): $\gamma^* p \rightarrow \gamma p$



High  $Q^2$   
Perturbative QCD

Non-perturbative  
GPDs

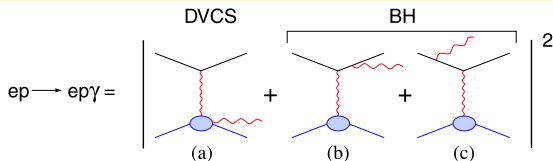
**Handbag diagram**

**Bjorken limit:**

$$Q^2 = \left. \begin{array}{l} -q^2 \rightarrow \infty \\ \nu \rightarrow \infty \end{array} \right\} x_B = \frac{Q^2}{2M\nu} \text{ fixed}$$

- GPDs accesible through DVCS *only* at  $Q^2 \rightarrow \infty$
- Actual value of  $Q^2$  *must* be tested and established **by experiment**

# DVCS experimentally: interference with Bethe-Heitler



At leading twist:

$$d^5 \vec{\sigma} - d^5 \overleftarrow{\sigma} = 2 \Im (T^{BH} \cdot T^{DVCS})$$

$$d^5 \vec{\sigma} + d^5 \overleftarrow{\sigma} = |BH|^2 + 2 \Re (T^{BH} \cdot T^{DVCS}) + |DVCS|^2$$

$$T^{DVCS} = \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi + i\epsilon} + \dots =$$

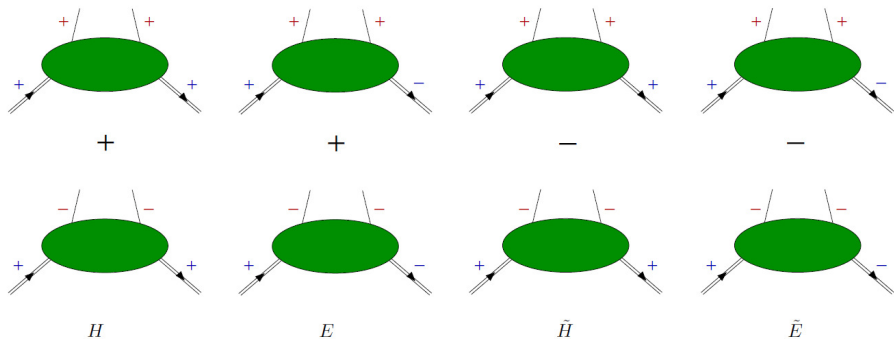
$$\underbrace{\mathcal{P} \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi}}_{\text{Access in helicity-independent cross section}} - \underbrace{i\pi H(x = \xi, \xi, t)}_{\text{Access in helicity-dependent cross-section}} + \dots$$

Access in **helicity-independent cross section**

Access in **helicity-dependent cross-section**

# Leading twist GPDs

8 GPDs related to the different combination of quark/nucleon helicities

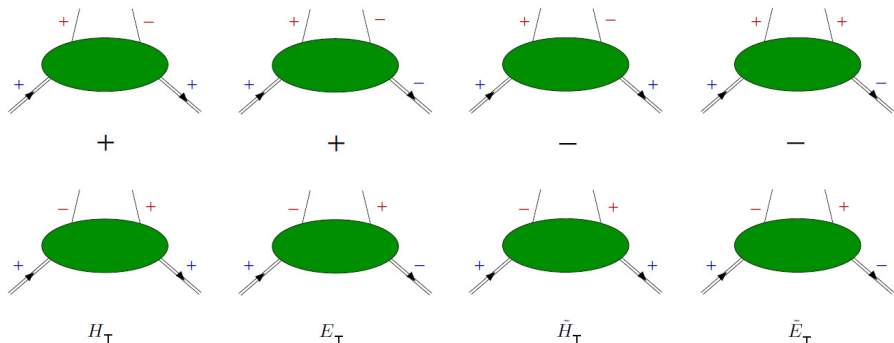


4 chiral-even GPDs: conserve the helicity of the quark

Access through DVCS (and DVMP)

# Leading twist GPDs

8 GPDs related to the different combination of quark/nucleon helicities



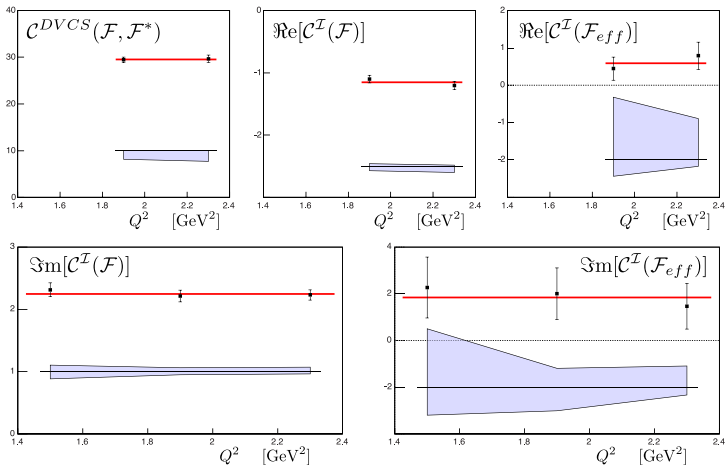
4 chiral-odd GPDs: flip helicity of the quark

“transversity GPDs”

Experimental access more complicated ( $\pi^0$  electroproduction?)

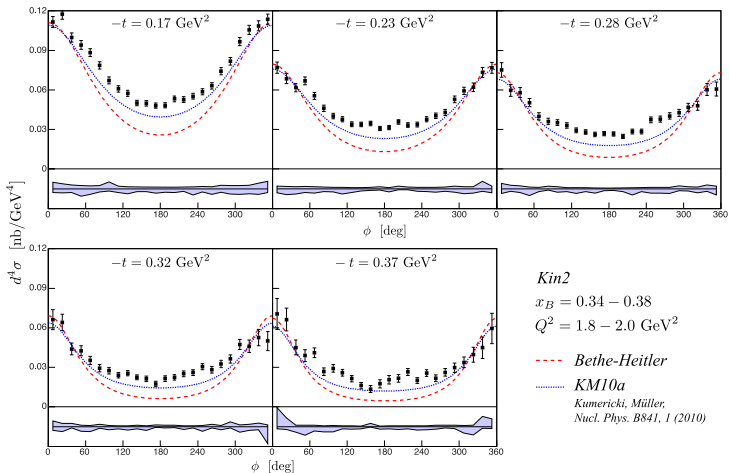
Ahmad, Goldstein, Liuti (2009)

Goloskokov, Kroll (2011)

DVCS cross sections:  $Q^2$ -dependence

No  $Q^2$ -dependence within limited range  $\Rightarrow$  leading twist dominance

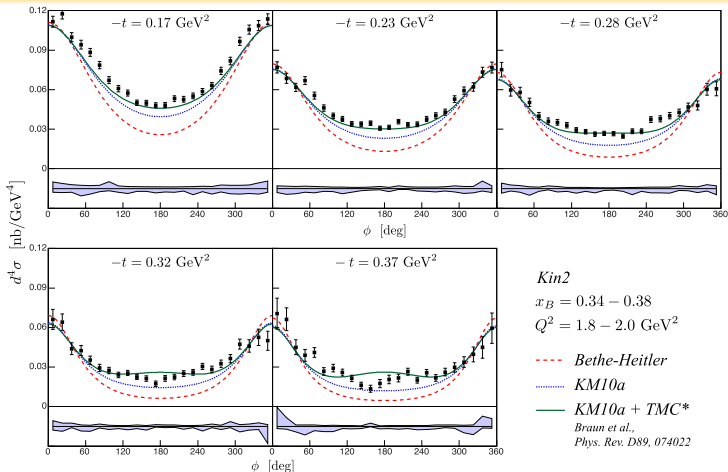
# DVCS cross sections: higher twist corrections



- KM10a: global fit to HERA x-sec & HERMES + CLAS spin asymmetries  
 Kumericki and Mueller (2010)



# DVCS cross sections: higher twist corrections



- KM10a: global fit to HERA x-sec & HERMES + CLAS spin asymmetries  
Kumericki and Mueller (2010)
- Target-mass corrections (TMC):  $\sim \mathcal{O}(M^2/Q^2)$  and  $\sim \mathcal{O}(t/Q^2)$

Braun, Manashov, Mueller and Pirnay (2014)

# DVCS process: leading twist ambiguity

- DVCS defines a preferred axis: light-cone axis
- At finite  $Q^2$  and non-zero  $t$ , there is an ambiguity:
  - 1 Belitsky et al. (“BMK”, 2002–2010): light-cone axis in plane  $(q, P)$
  - 2 Braun et al. (“BMP”, 2014): light-cone axis in plane  $(q, q')$   
easier to account for kin. corrections  $\sim \mathcal{O}(M^2/Q^2)$ ,  $\sim \mathcal{O}(t/Q^2)$

$$\left. \begin{aligned}
 \mathcal{F}_{++} &= \mathbb{F}_{++} + \frac{\chi}{2} [\mathbb{F}_{++} + \mathbb{F}_{-+}] - \chi_0 \mathbb{F}_{0+} \\
 \mathcal{F}_{-+} &= \mathbb{F}_{-+} + \frac{\chi}{2} [\mathbb{F}_{++} + \mathbb{F}_{-+}] - \chi_0 \mathbb{F}_{0+} \\
 \mathcal{F}_{0+} &= -(1 + \chi) \mathbb{F}_{0+} + \chi_0 [\mathbb{F}_{++} + \mathbb{F}_{-+}]
 \end{aligned} \right\} \begin{array}{l} \mathbb{F}_{-+} = 0 \\ \mathbb{F}_{0+} = 0 \end{array} \rightarrow \left\{ \begin{array}{l} \mathcal{F}_{++} = (1 + \frac{\chi}{2}) \mathbb{F}_{++} \\ \mathcal{F}_{-+} = \frac{\chi}{2} \mathbb{F}_{++} \\ \mathcal{F}_{0+} = \chi_0 \mathbb{F}_{++} \end{array} \right.$$

(eg.  $\chi_0 = 0.25$ ,  $\chi = 0.06$  for  $Q^2 = 2 \text{ GeV}^2$ ,  $x_B = 0.36$ ,  $t = -0.24 \text{ GeV}^2$ )

# Rosenbluth-like separation of the DVCS cross section

$$\sigma(ep \rightarrow ep\gamma) = \underbrace{|BH|^2}_{\text{Known to } \sim 1\%} + \underbrace{\mathcal{I}(BH \cdot DVCS)}_{\text{Linear combination of GPDs}} + \underbrace{|DVCS|^2}_{\text{Bilinear combination of GPDs}}$$

$$\mathcal{I} \propto 1/y^3 = (k/\nu)^3,$$

$$|\mathcal{T}^{DVCS}|^2 \propto 1/y^2 = (k/\nu)^2$$

BMK-2010 – at leading twist  $\rightarrow$  7 independent GPD terms:

$$\{\Re, \Im [c^{\mathcal{I}}, c^{\mathcal{I},V}, c^{\mathcal{I},A}] (\mathcal{F})\}, \quad \text{and} \quad c^{DVCS}(\mathcal{F}, \mathcal{F}^*).$$

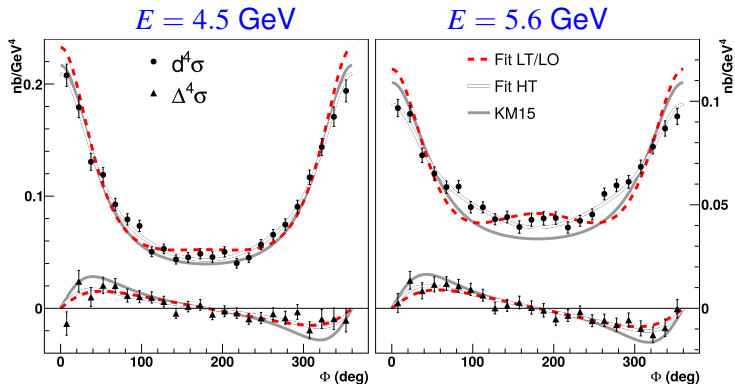
$\varphi$ -dependence provides 5 independent observables:

$$\sim 1, \sim \cos \varphi, \sim \sin \varphi, \sim \cos(2\varphi), \sim \sin(2\varphi)$$

The measurement of the cross section at **two or more beam energies** for exactly the **same  $Q^2, x_B, t$  kinematics**, provides the additional information in order to extract all leading twist observables independently.

## E07-007: DVCS beam-energy dependence

- Cross section measured at 2 beam energies and constant  $Q^2$ ,  $x_B$ ,  $t$



- Leading-twist and LO simultaneous fit of both beam energies (dashed line) does not reproduce the data

**Light-cone axis in the  $(q, q')$  plane (Braun *et al.*):**  $\mathbb{H}_{++}$ ,  $\tilde{\mathbb{H}}_{++}$ ,  $\mathbb{E}_{++}$ ,  $\tilde{\mathbb{E}}_{++}$

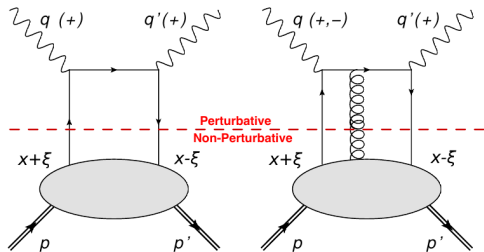
# Beyond Leading Order (LO) and Leading Twist (LT)

## Two fit-scenarios:

Light-cone axis in  
the  $(q, q')$  plane (Braun *et al.*)

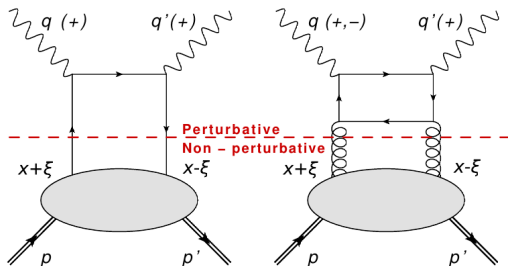
LO/LT + HT

$$\mathbb{H}_{++}, \widetilde{\mathbb{H}}_{++}, \mathbb{H}_{0+}, \widetilde{\mathbb{H}}_{0+}$$



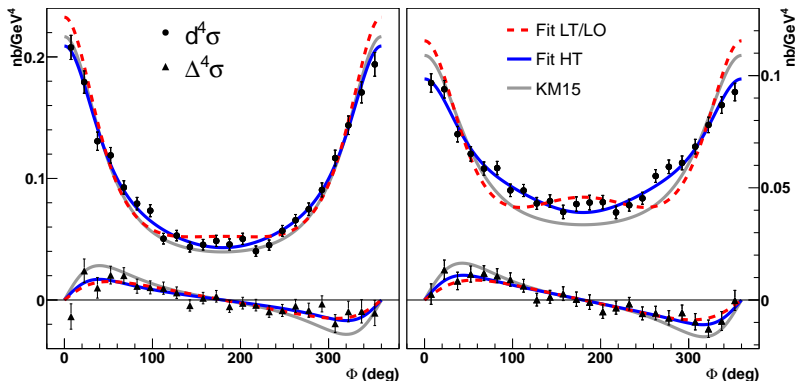
LO/LT + NLO

$$\mathbb{H}_{++}, \widetilde{\mathbb{H}}_{++}, \mathbb{H}_{-+}, \widetilde{\mathbb{H}}_{-+}$$



# E07-007: DVCS beam-energy dependence

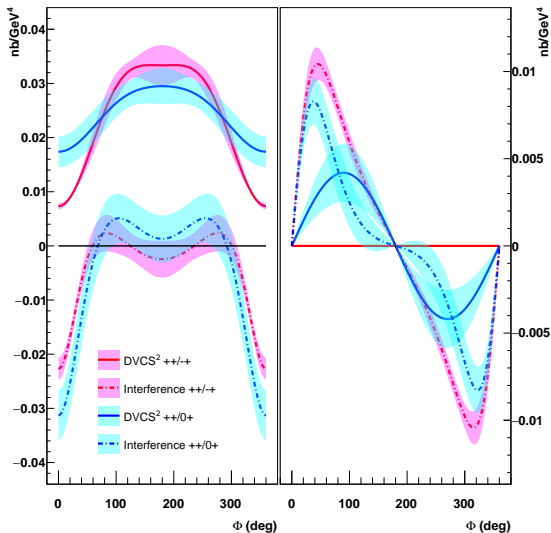
- Cross section measured at 2 beam energies and constant  $Q^2$ ,  $x_B$ ,  $t$



- Leading-twist and LO simultaneous fit of both beam energies (dashed line) does not reproduce the data
- Including either NLO or higher-twist effects (dark solid line) satisfactorily reproduce the angular dependence

# DVCS<sup>2</sup> and $\mathcal{I}$ (DVCS·BH) separation

DVCS<sup>2</sup> and  $\mathcal{I}$  (DVCS·BH) separated in NLO and higher-twist scenarios



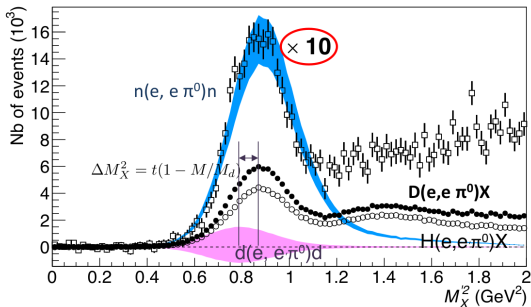
- DVCS<sup>2</sup> &  $\mathcal{I}$  significantly different in each scenario
- Sizeable DVCS<sup>2</sup> contribution in the higher-twist scenario in the helicity-dependent cross section

arXiv:1703.09442

# E08-025: DVCS and $\pi^0$ off quasi-free neutrons

- LD<sub>2</sub> as a target
- Quasi-free  $p$  evts subtracted using the (normalized) data from E07-007
- Concurrent running: switching LH2/LD2 → minimize uncertainties

$$D(e, e \pi^0)X - p(e, e \pi^0)p = n(e, e \pi^0)n + d(e, e \pi^0)d$$

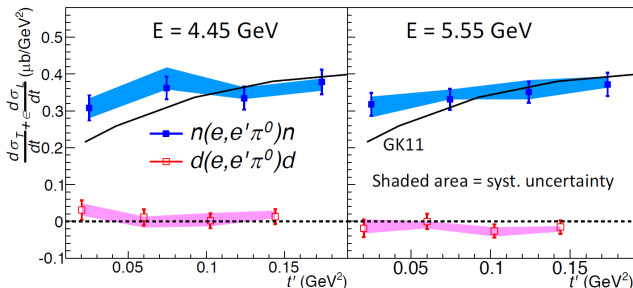
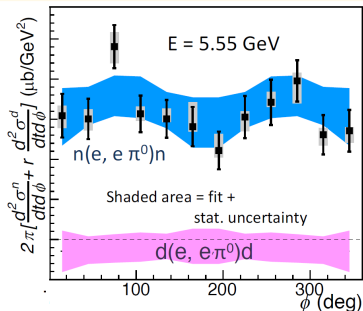


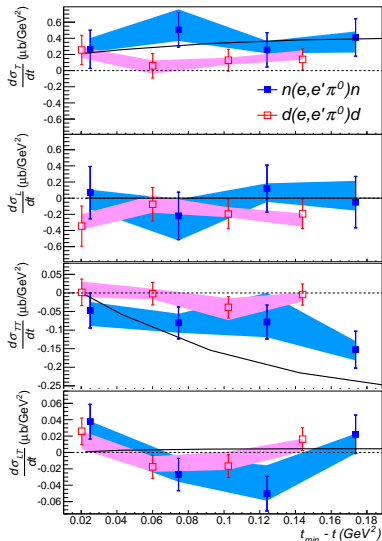
The average momentum transfer to the target is much larger than the  $np$  relative momentum, justifying this **impulse approximation**



# $\pi^0$ electroproduction cross section off the neutron

- Cross section off coherent  $d$  found negligible within uncertainties
- Very low  $E_{beam}$  dependence of the  $n$  cross section  $\rightarrow$  dominance of  $\sigma_T$

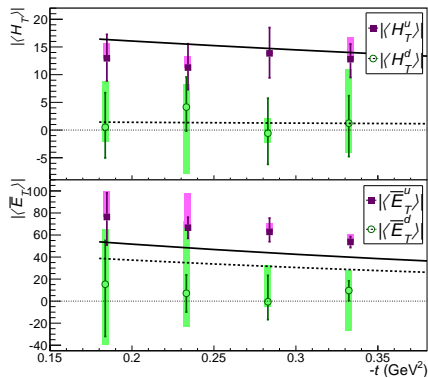


Separated  $\pi^0$  cross section off the neutron

M. Mazouz et al, arXiv:1702.00835  
Phys. Rev. Lett., in press

In the modified factorization approach (KG):

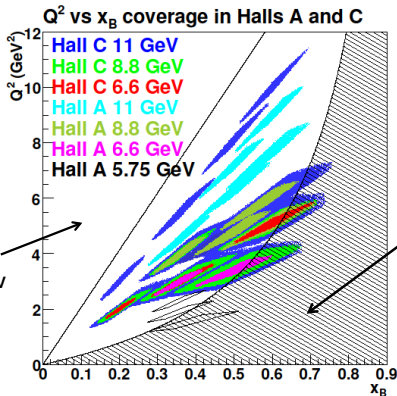
- $d\sigma_T \propto \left[ (1 - \xi^2) |\langle H_T \rangle|^2 - \frac{t'}{8M^2} |\langle \bar{E}_T \rangle|^2 \right]$
- $d\sigma_{TT} \propto \frac{t'}{8M^2} |\langle \bar{E}_T \rangle|^2$



$$|\langle H_T^{p,n} \rangle|^2 = \frac{1}{2} \left| \frac{2}{3} \langle H_T^{u,d} \rangle + \frac{1}{3} \langle H_T^{d,u} \rangle \right|^2$$

# 11 (and 8.8 and 6.6) GeV program in Halls A & C

- Approved high precision DVCS &  $\pi^0$  programs in Halls A & C



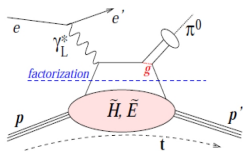
- **Hall A** 50% of data already collected (2014–2016):  
analysis underway...

- **Hall C** program could start ~2019–2020:  
Neutral Particle Spectrometer (NPS) under construction

# Summary

- Recent high precision DVCS cross sections from Hall A at JLab
- Need of higher twist and/or NLO contributions to fully describe the data (eg. in global GPD fits)
- First separation of DVCS<sup>2</sup> and BH-DVCS interference in the  $ep \rightarrow e\gamma p$  cross section
- L/T separation of  $\pi^0$  electroproduction cross section off neutron: dominance of  $\sigma_T$  measured
- Flavor separation of transversity GPD convolutions within the modified factorization approach
- Approved program of experiments in Hall A and C to continue these high precision DVCS measurements at 12 GeV

# Back-up

$\pi^0$  electroproduction ( $ep \rightarrow ep\pi^0$ )

At leading twist:

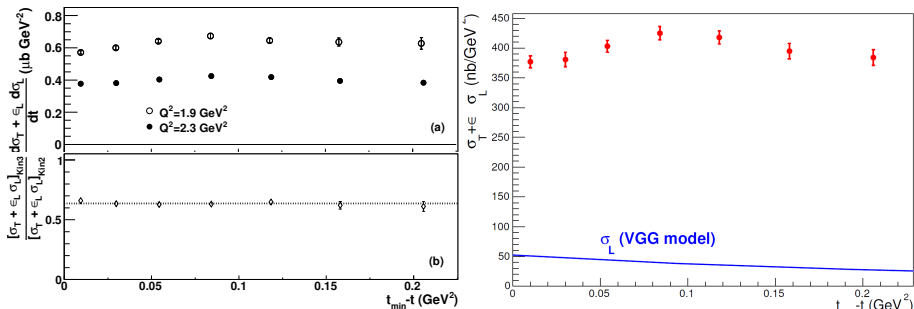
$$\frac{d\sigma_L}{dt} = \frac{1}{2}\Gamma \sum_{h_N, h_{N'}} |\mathcal{M}^L(\lambda_M = 0, h'_N, h_N)|^2 \propto \frac{1}{Q^6} \quad \sigma_T \propto \frac{1}{Q^8}$$

$$\mathcal{M}^L \propto \left[ \int_0^1 dz \frac{\phi_\pi(z)}{z} \right] \int_{-1}^1 dx \left[ \frac{1}{x - \xi} + \frac{1}{x + \xi} \right] \times \left\{ \Gamma_1 \tilde{H}_{\pi^0} + \Gamma_2 \tilde{E}_{\pi^0} \right\}$$

Different quark weights: flavor separation of GPDs

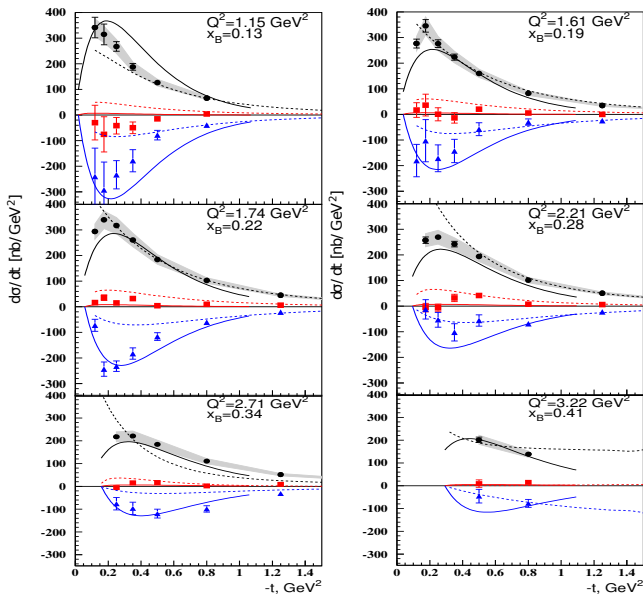
$$|\pi^0\rangle = \frac{1}{\sqrt{2}} \{ |u\bar{u}\rangle - |d\bar{d}\rangle \} \quad \tilde{H}_{\pi^0} = \frac{1}{\sqrt{2}} \left\{ \frac{2}{3} \tilde{H}^u + \frac{1}{3} \tilde{H}^d \right\}$$

$$|p\rangle = |uud\rangle \quad H_{DVCS} = \frac{4}{9} H^u + \frac{1}{9} H^d$$

Exclusive  $\pi^0$  electroproduction cross-sections – Hall A

- $\sigma_T + \epsilon_L \sigma_L \sim Q^{-5}$   
(similar to  $\sigma_T(ep \rightarrow ep\pi^+)$  measured in Hall C)
- GPDs predict  $\sigma_L \sim Q^{-6}$
- $\sigma_T$  likely to dominate at these  $Q^2$ ,  
but L/T separation necessary ( $\rightarrow$  new experiment. . .)

E. Fuchey et al., Phys. Rev. C83 (2011), 025125

Exclusive  $\pi^0$  electroproduction cross-sections – Hall B

- $\sigma_T + \epsilon\sigma_L$
- $\sigma_{TL}$
- $\sigma_{TT}$

- Very large cross section compared to leading chiral even GPD models
- Fair agreement w/ models of trans. photon (modified) factorization (& transversity GPDs)

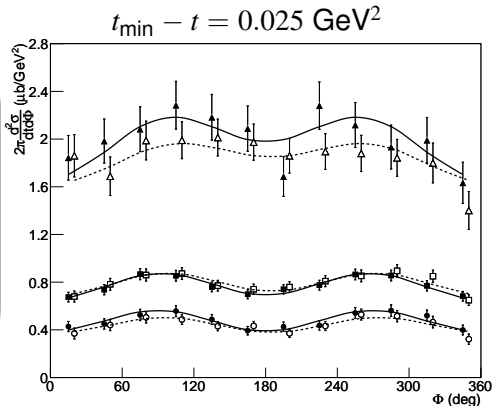


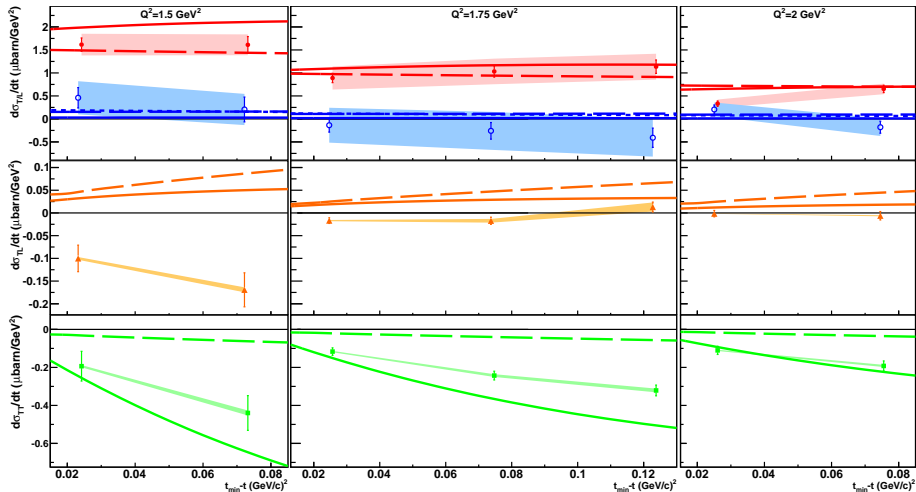
# Rosenbluth separation

$$\frac{d^4\sigma}{dQ^2 dx_B dt d\phi} = \frac{1}{2\pi} \Gamma(Q^2, x_B, E) \left[ \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \sqrt{2\epsilon(1+\epsilon)} \frac{d\sigma_{TL}}{dt} \cos\phi + \epsilon \frac{d\sigma_{TT}}{dt} \cos 2\phi \right]$$

## Kinematics

Setting	$Q^2$ (GeV <sup>2</sup> )	$x_B$	$E^{beam}$ (GeV)	$\epsilon$
Kin1	1.50	0.36	3.355	0.52
			5.55	0.84
Kin2	1.75	0.36	4.455	0.65
			5.55	0.79
Kin3	2.00	0.36	4.455	0.53
			5.55	0.72



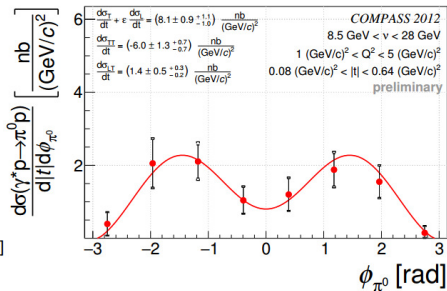
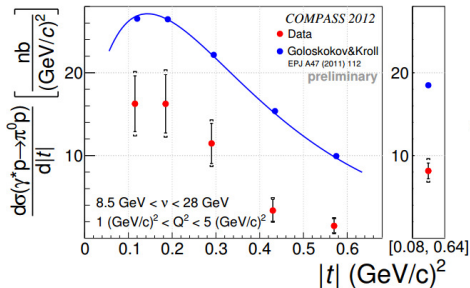
$\pi^0$  separated response functions

- Goloskokov, Kroll (2011)
- - - Goldstein, Hernandez, Liuti (2011)
- · - · Vanderhaeghen, Guichon, Guidal (1999)

# $\pi^0$ L/T separated cross section

- Cross section largely dominated by transverse component  
→ far from asymptotic prediction of QCD
- Fair agreement with models using modified factorization approach  
→ potential access to transversity GPDs
- Indications of small longitudinal response through non-zero  $\sigma_{LT}$

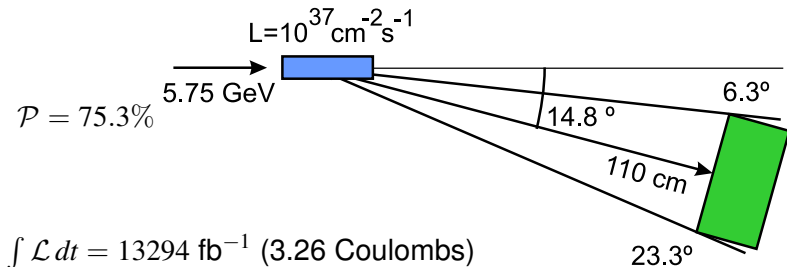
# Exclusive $\pi^0$ cross section from COMPASS 2012 run



Results also in fair agreement with transversity GPD models

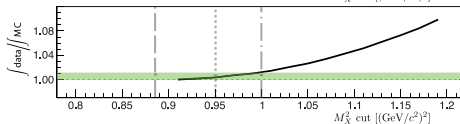
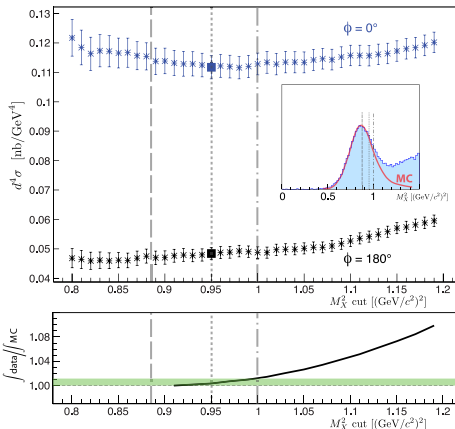
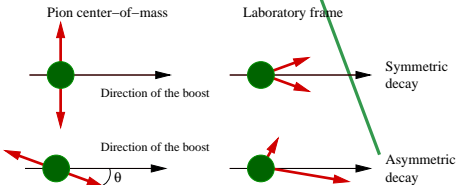
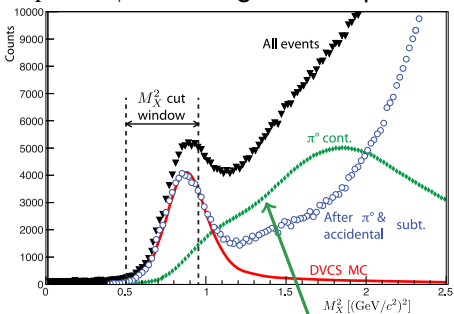
Kinematic settings: testing  $Q^2$ -dependance

Kin	$Q^2$ (GeV <sup>2</sup> )	$x_B$	$\theta_e$ (deg.)	$\theta_{\gamma^*}$ (deg.)	$P_e$ (GeV)
1	<b>1.5</b>	0.36	15.6	<b>22.3</b>	3.6
2	<b>1.9</b>	0.36	19.3	<b>18.3</b>	2.9
3	<b>2.3</b>	0.36	23.9	<b>14.8</b>	2.3



# Data analysis: exclusivity and background subtraction

$ep \rightarrow e\gamma X$  missing mass squared



- Only  $e'$  &  $\gamma$  detected +  $M_X^2$ -cut
- 2-3% uncertainty on exclusivity