


# Observational Constraints on the Supernova Engine

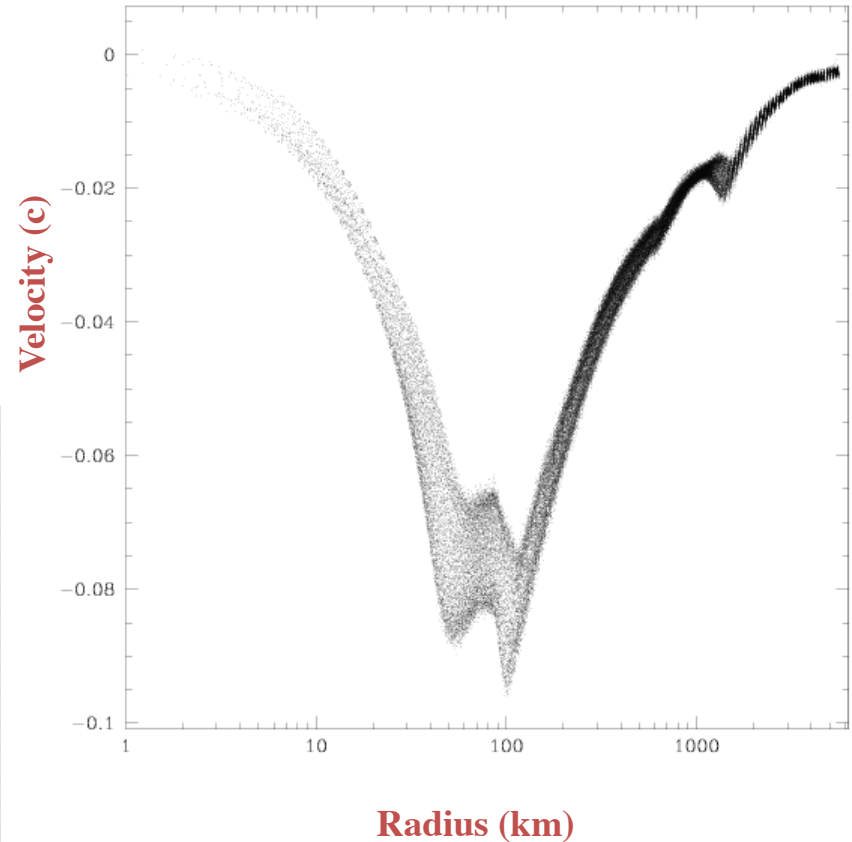
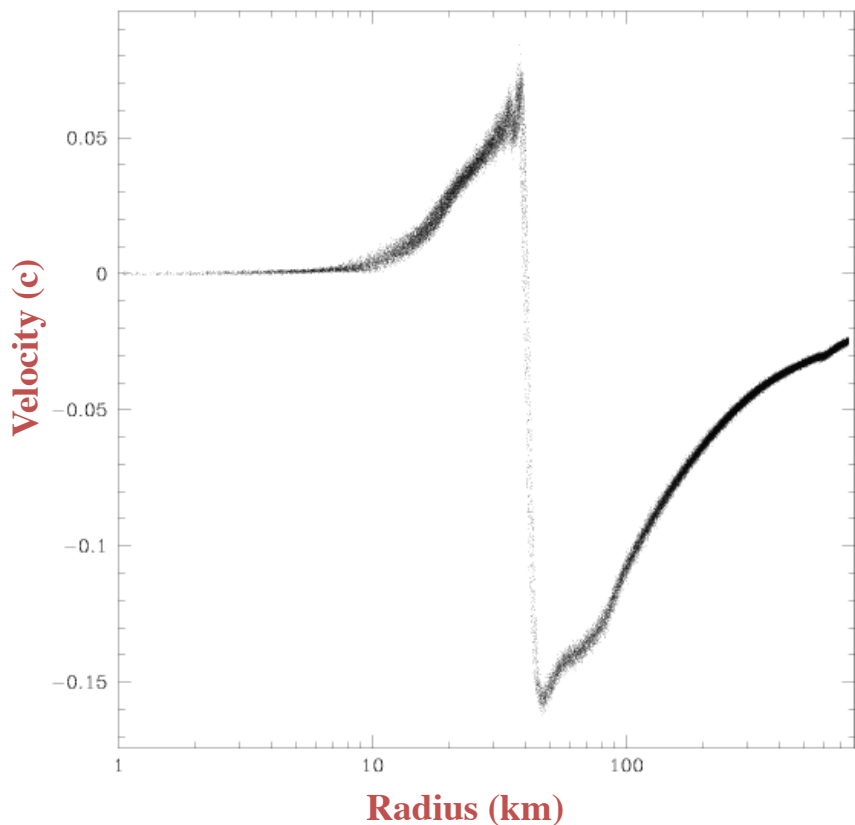
The background of the slide is a composite image of a galaxy, likely the Whirlpool Galaxy (M51), showing its characteristic spiral structure. The galaxy is oriented diagonally across the frame. A bright, glowing region is visible in the central part of the galaxy, which is likely the site of a supernova remnant. The colors are a mix of blue and red, suggesting different wavelengths of light. In the lower-left foreground, there is a bright, multi-pointed star, possibly a supernova remnant or a distant star, with a blue-white hue.

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# Neutrino-Driven Supernova Mechanism

**Temperature and Density of the Core  
Becomes so High that:**

**Iron dissociates into alpha particles**  
**Electrons capture onto protons**  
**Core collapses nearly at freefall!**

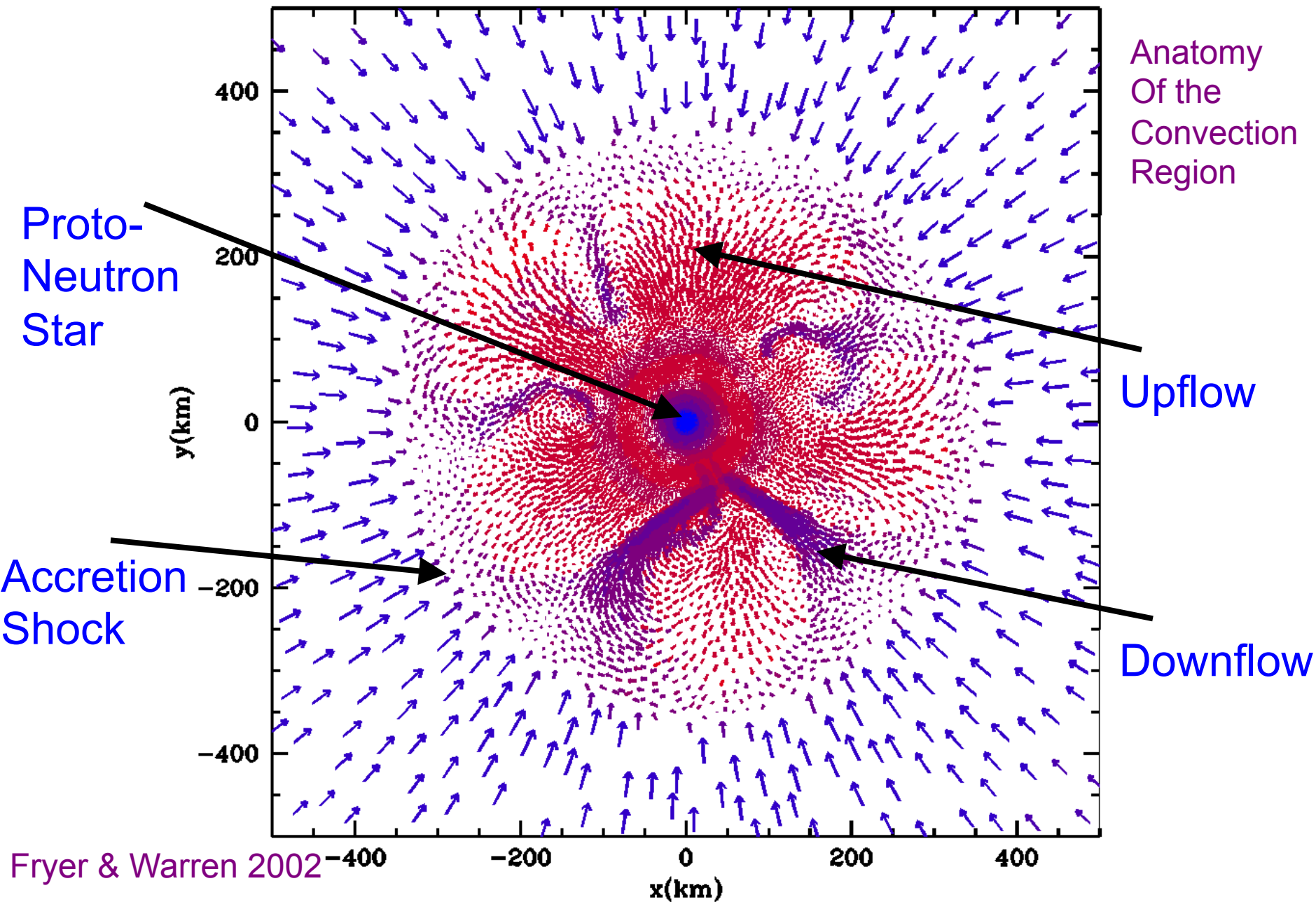


**Core reaches nuclear densities**

**Nuclear forces and neutron  
degeneracy increase pressure**

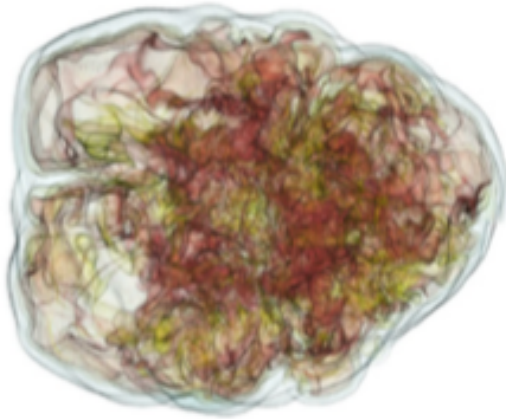
**Bounce!**

# The Herant et al. (1994) Convective Supernova Engine





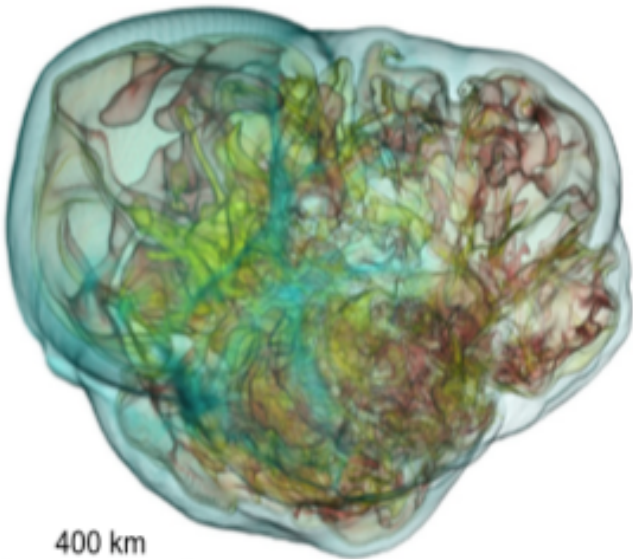
C15-3D 300 ms



400 km

Lentz et al. 2015

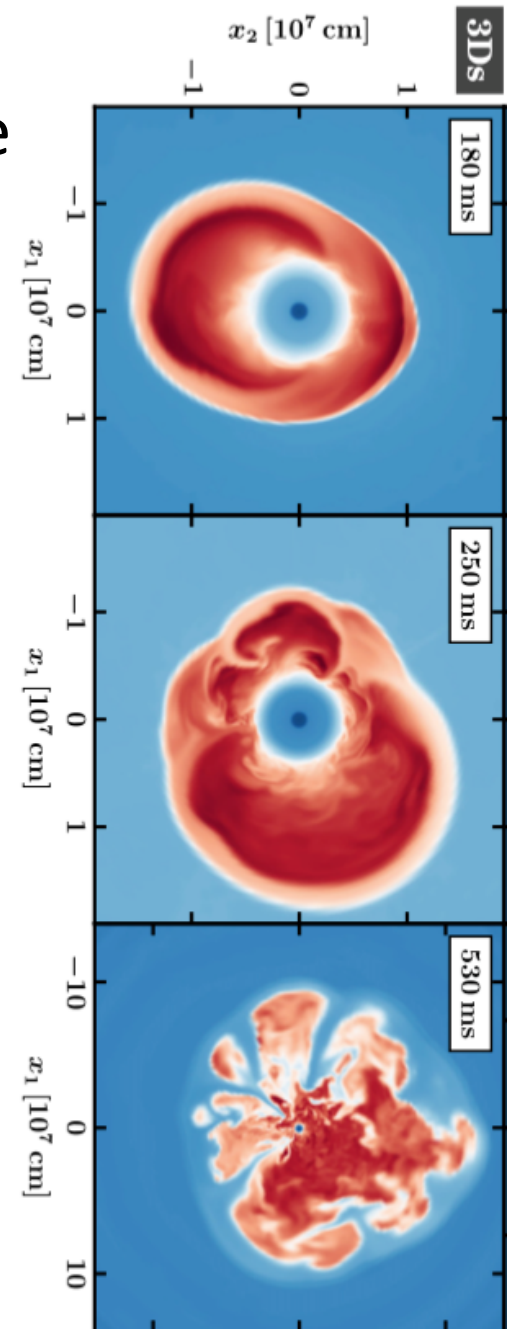
C15-3D 400 ms



400 km

Depending on the physics, most groups now produce explosions with this convective engine. Current arguments focus on the:

- Most important physics
- Source of instabilities



Melson et al. 2015

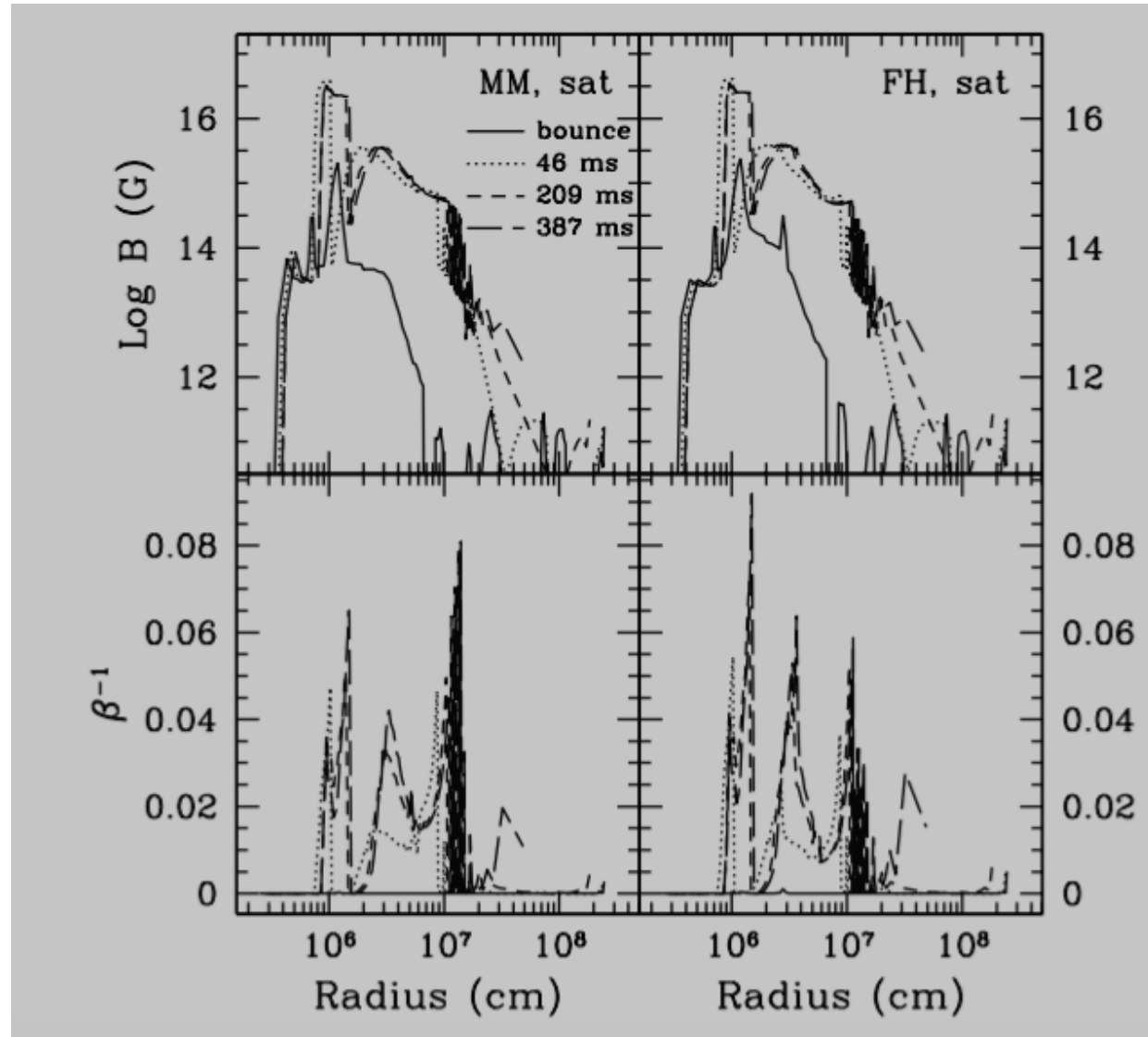
# Alternate engines

For long-duration gamma-ray bursts, a class of black-hole accretion disk engines have been proposed. Here, magnetic fields are wound up in the disk formed with the high-angular momentum layers fall onto a failed supernova.

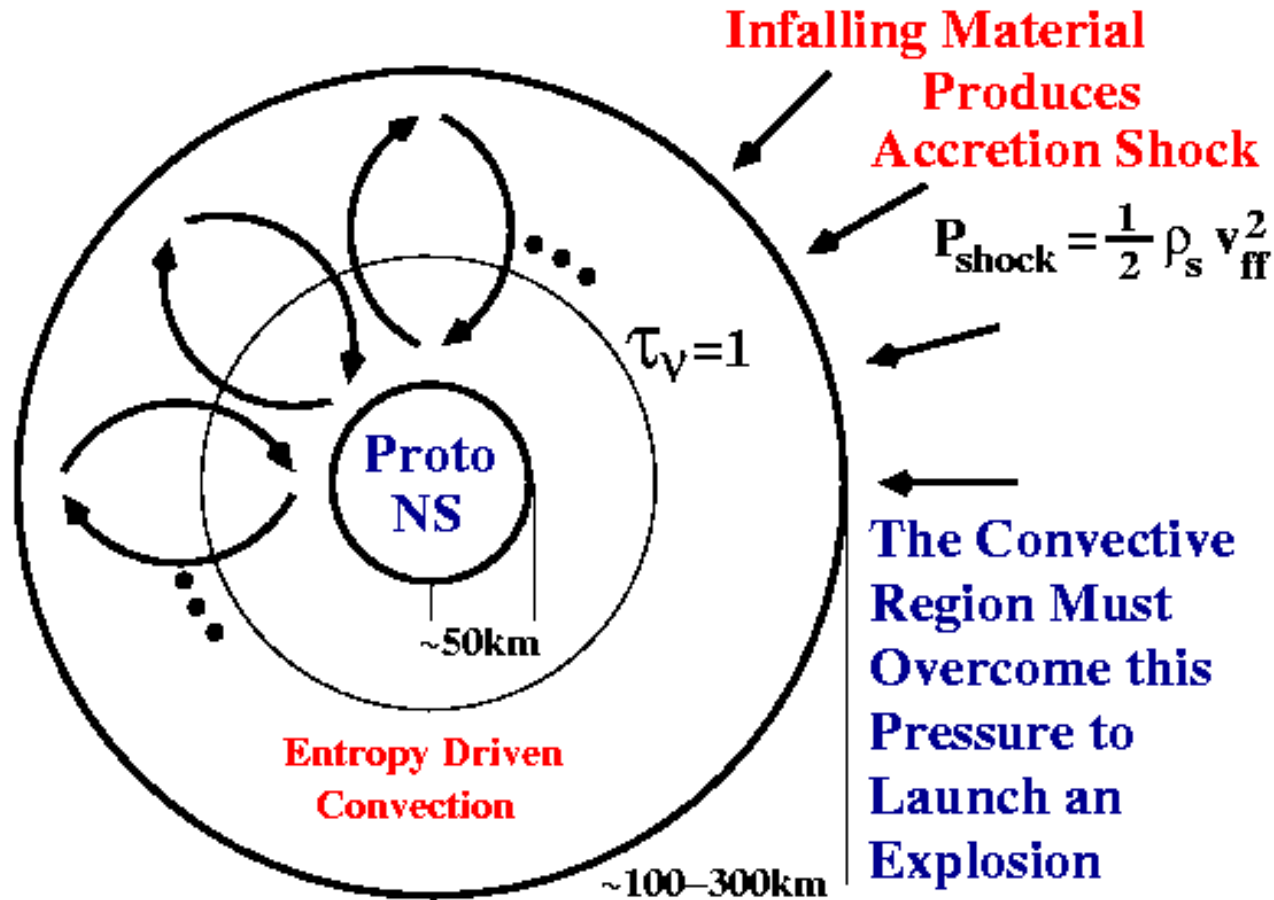
If the core is rapidly-rotating, a disk can form around the proto-neutron star.

If strong magnetic fields form (and not buried) in a rapidly rotating neutron star, magnetar engines may also work.

Both magnetic field engines predict jet-like explosions.



# Neutrino-Driven Supernova Mechanism: Convection

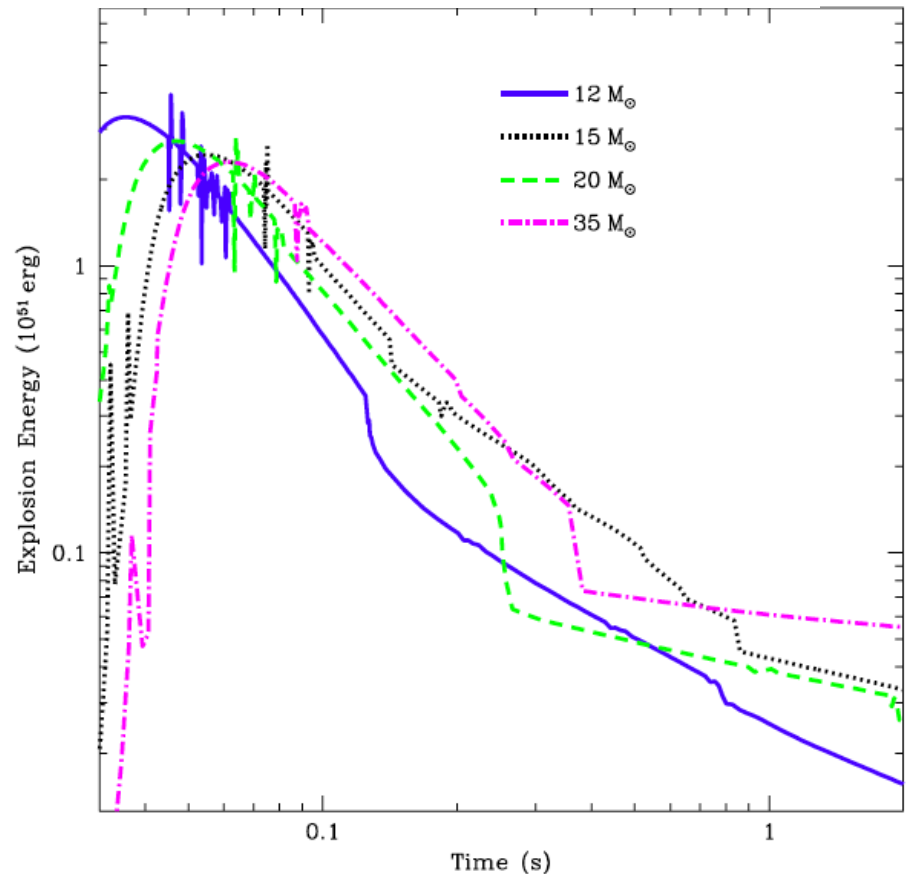


Fryer 1999

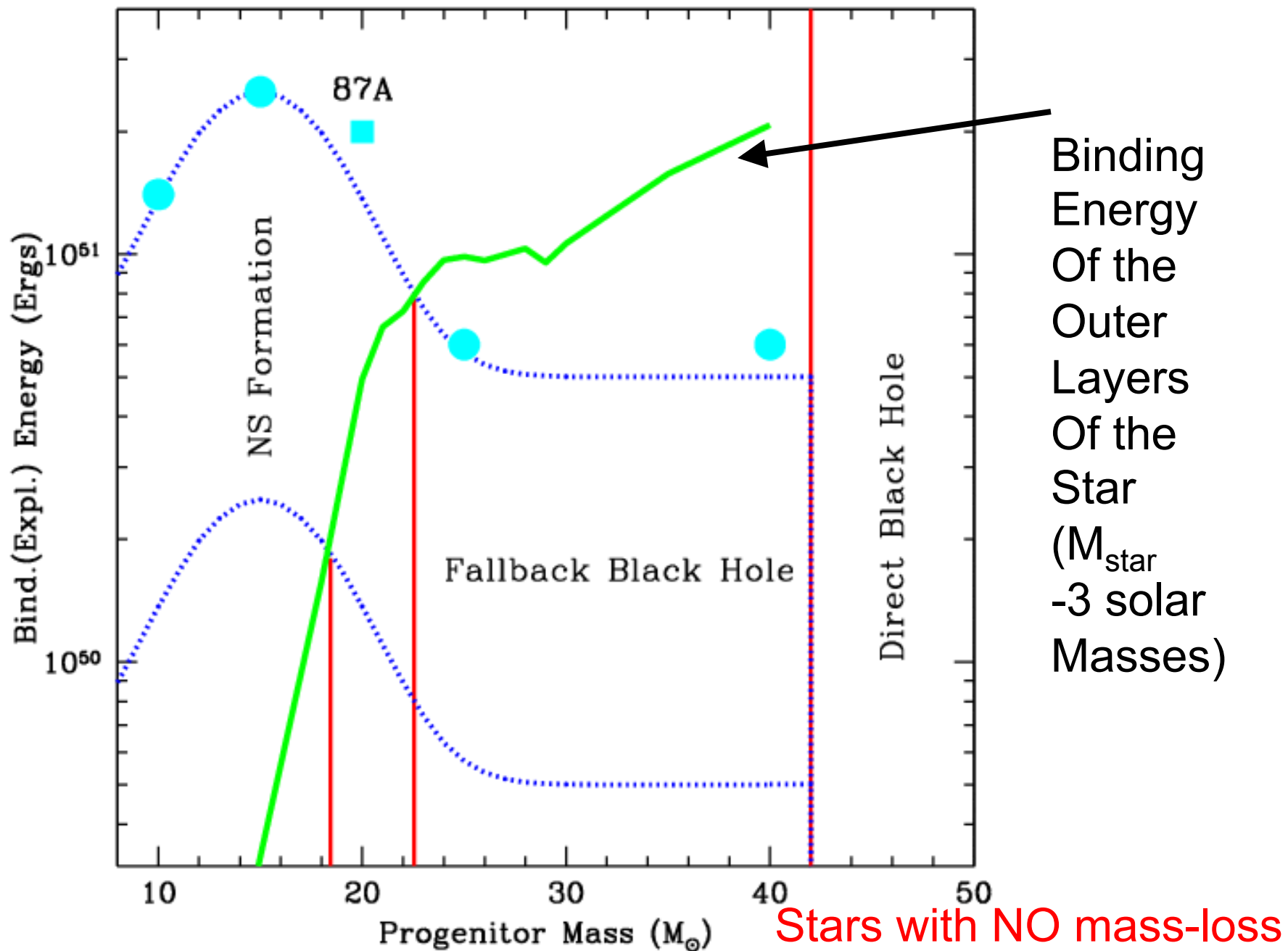
$$P_{\text{shock}}(r) = 1/2\rho_{\text{shock}}v_{\text{free-fall}} = (2GM_{\text{NS}})^{0.5}\dot{M}_{\text{acc}}/(8\pi r_{\text{shock}}^{2.5})$$

$$u_{\text{convection}}(r) = 3 \left[ 4.7 \times 10^8 \frac{M_{\text{NS}}}{M_{\odot}} \frac{10k_{\text{B}}\text{nucleon}^{-1}}{S_{\text{rad}}} \left( \frac{10^6\text{cm}}{r} - \frac{10^6\text{cm}}{r_{\text{shock}}} \right) + 1.2 \times 10^6 \left( \frac{M_{\text{NS}}}{M_{\odot}} \frac{\dot{M}_{\text{acc}}}{M_{\odot}\text{s}^{-1}} \right)^{1/4} \left( \frac{2 \times 10^7\text{cm}}{r_{\text{shock}}} \right)^{5/8} \right]^4 \text{erg cm}^{-3}.$$

- For most stars, the maximum energy is a few times  $10^{51}$  erg. Fallback can increase this value, but not by too much.
- This is a natural explanation for the energy, but it means that this engine can not explain hypernovae.



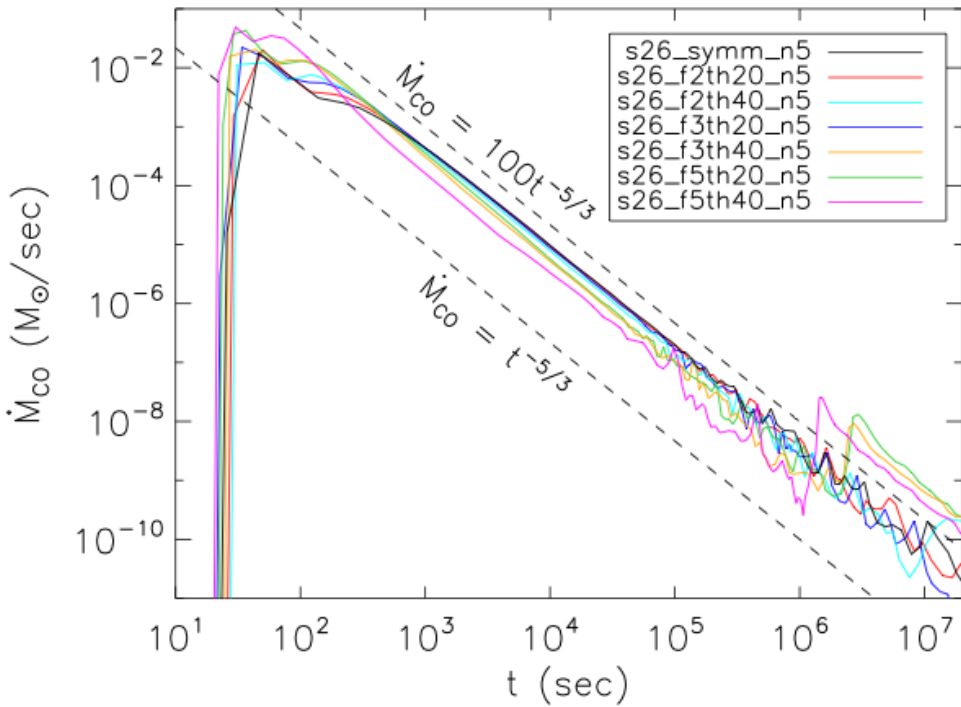
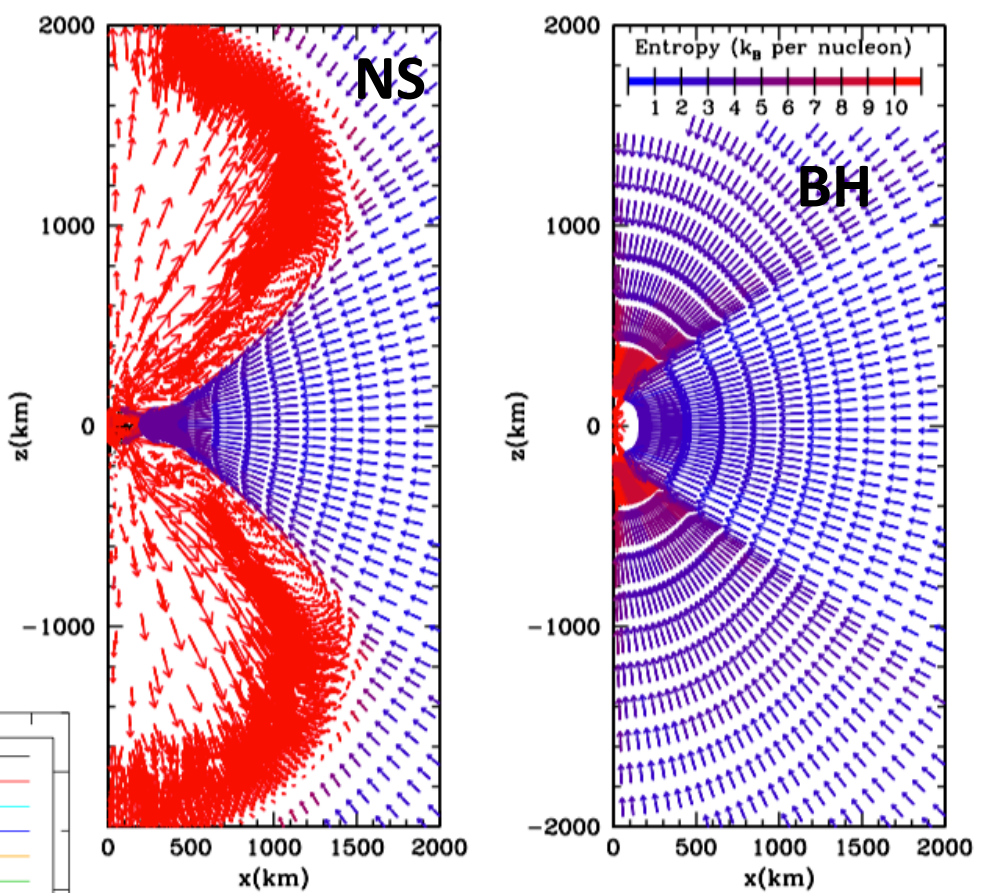
Fryer 1999





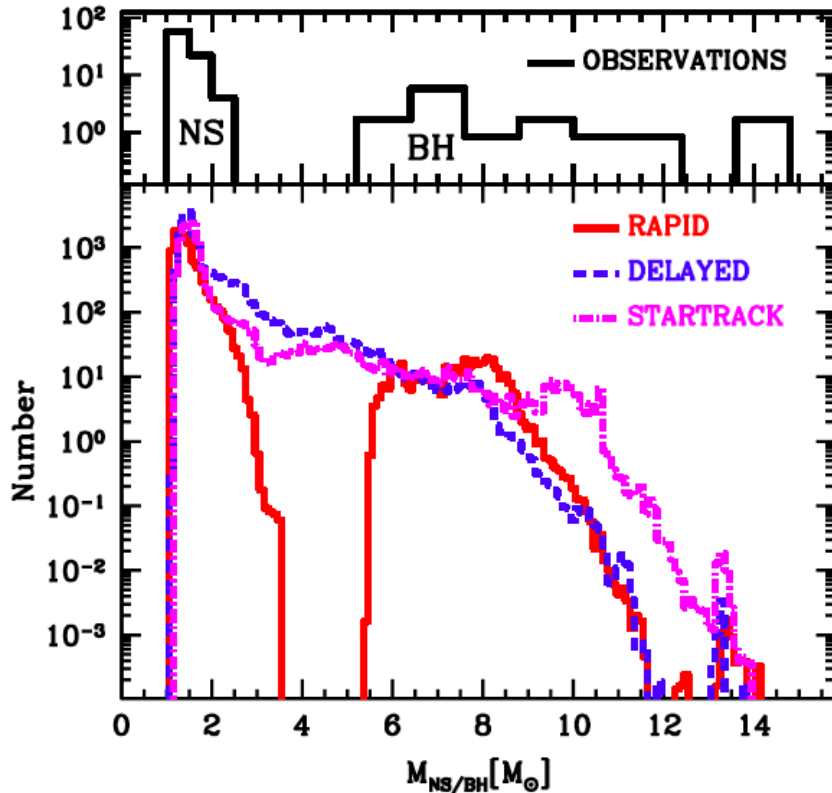
# Fallback rates

It is difficult to avoid fallback. Strong explosions have  $\sim 0.1$  Msun of fallback. Weak explosions allow a range of fallback from  $0.1+$  Msun. Most happens at early times, but at the level of  $10^{-4}$  Msun, this can happen even a year after the explosion.

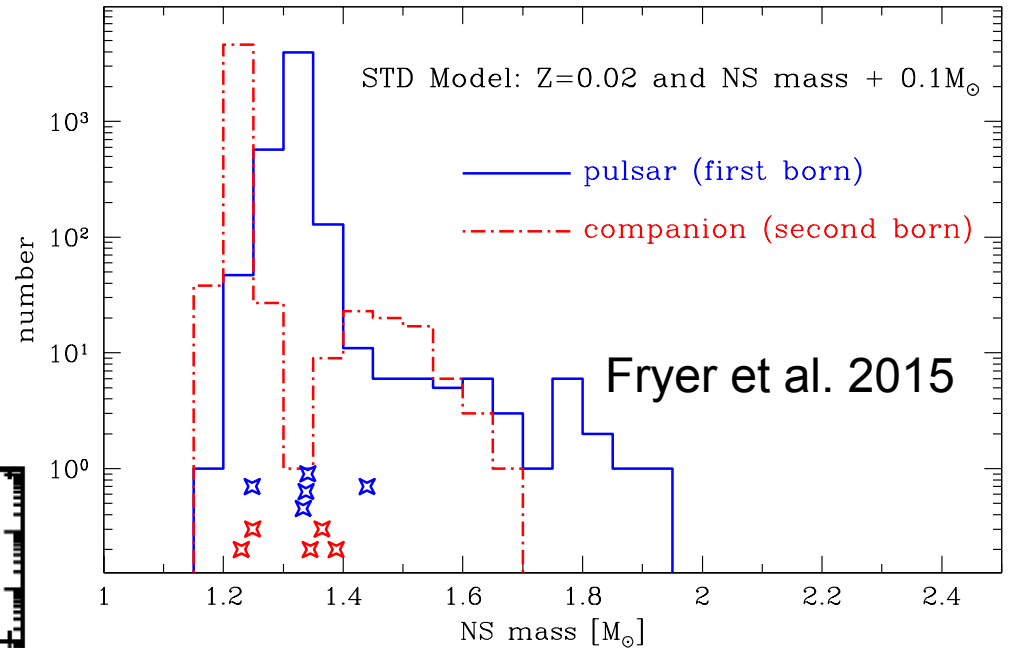


With rotation (even if not enough to form a disk), a significant fraction of the fallback can be re-ejected if the compact remnant is a neutron star.

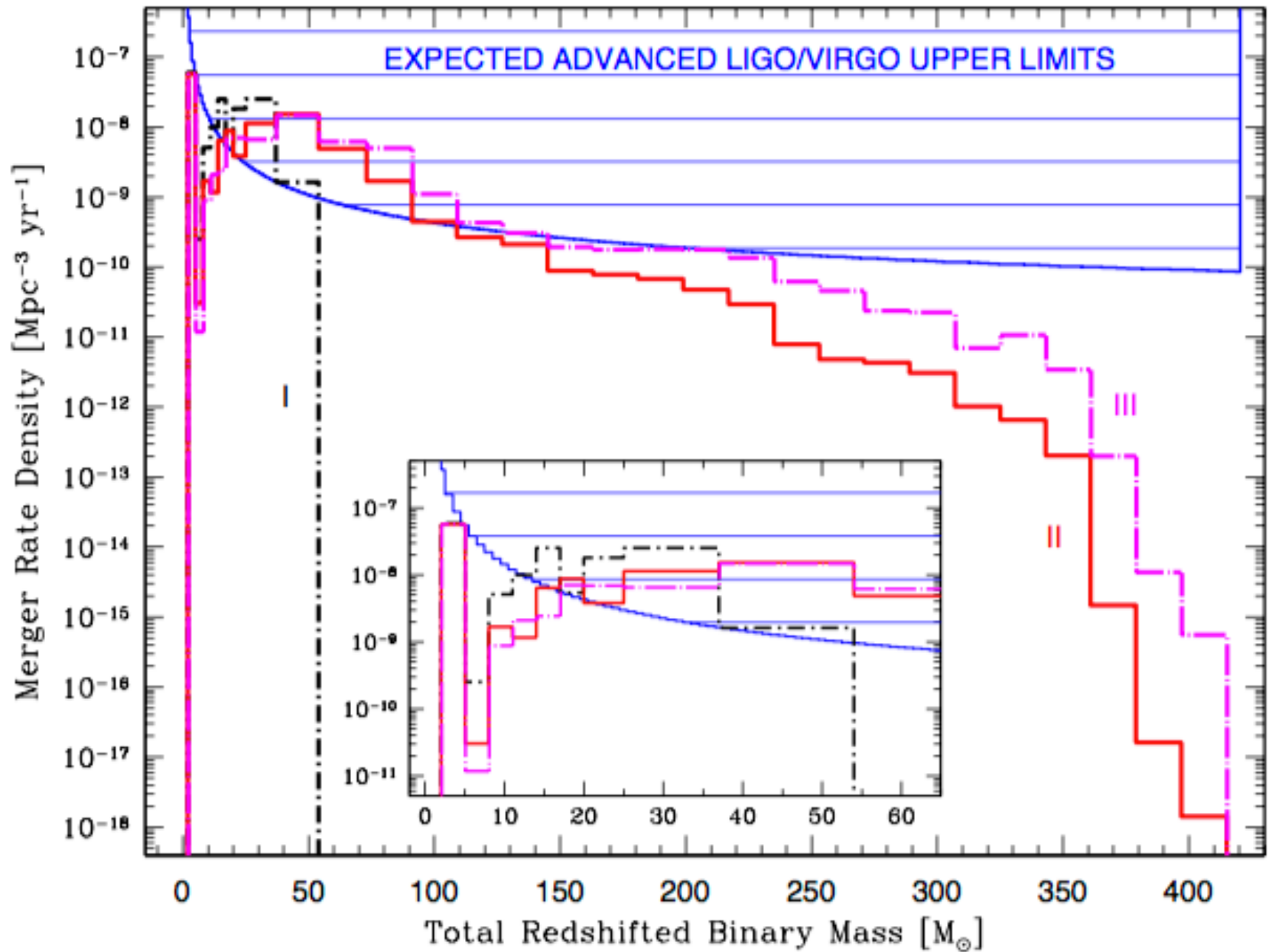
# Distribution of Neutron and Black Hole Masses



Belczynski 2012



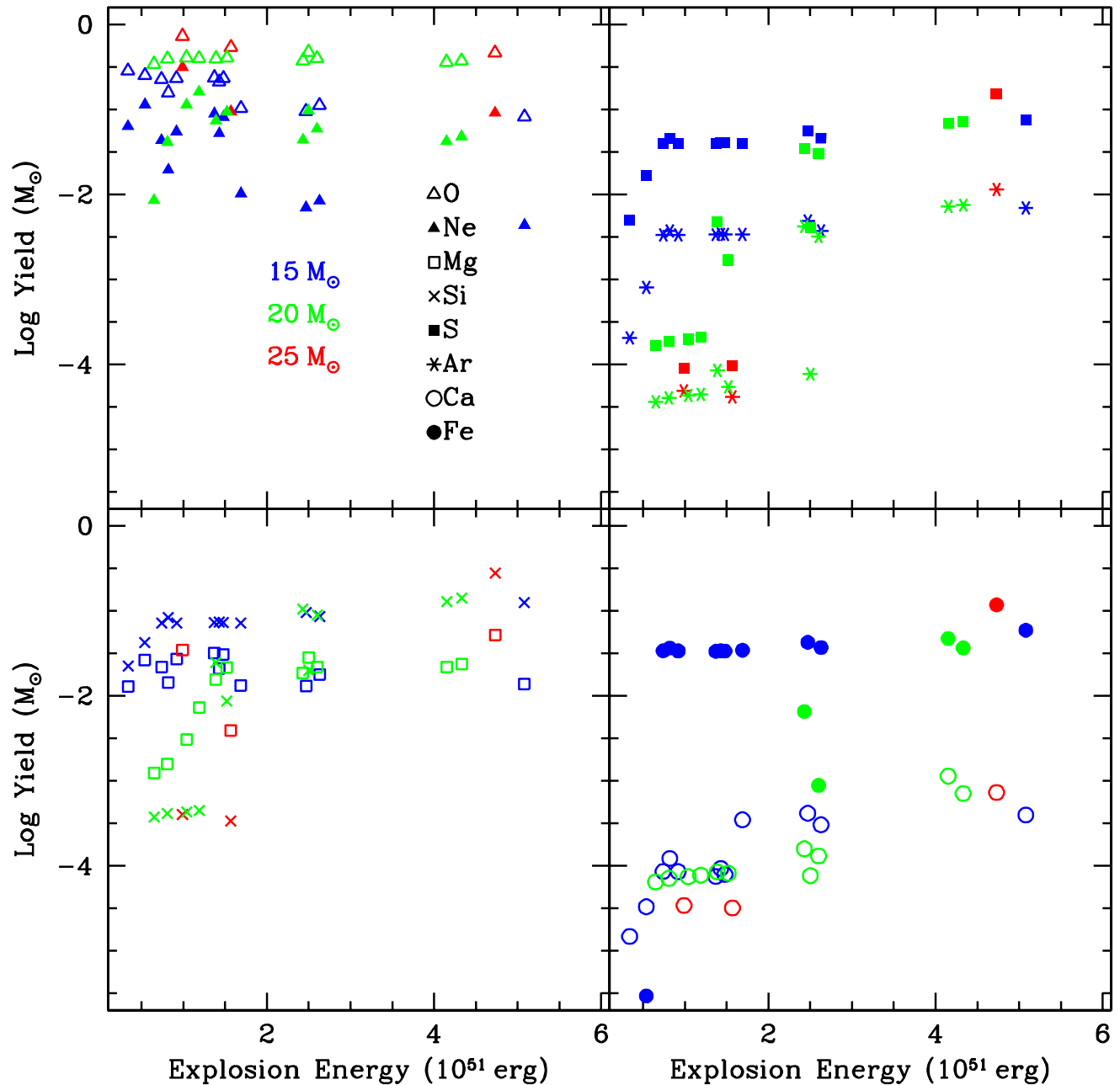
Both the NS-BH mass gap and the distribution of NSs must be fit by any explosion model. Since these masses have all evolved from close binaries, we must also include binary effects.



# Yields

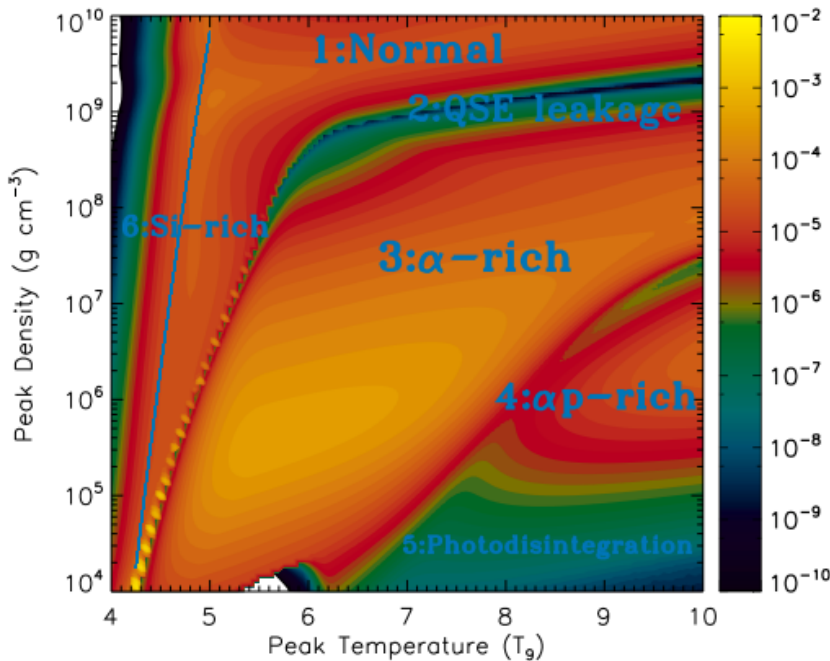
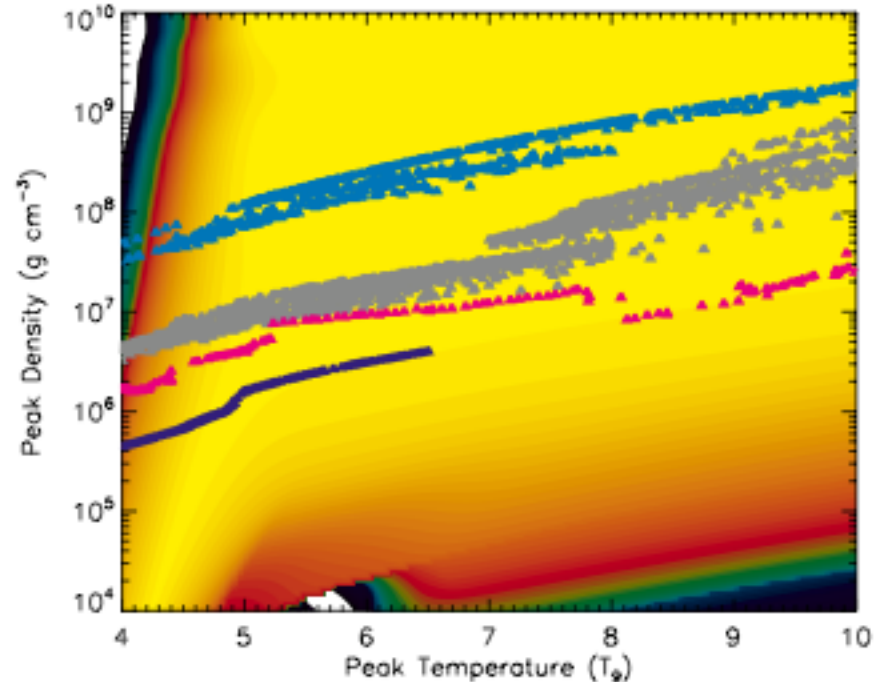
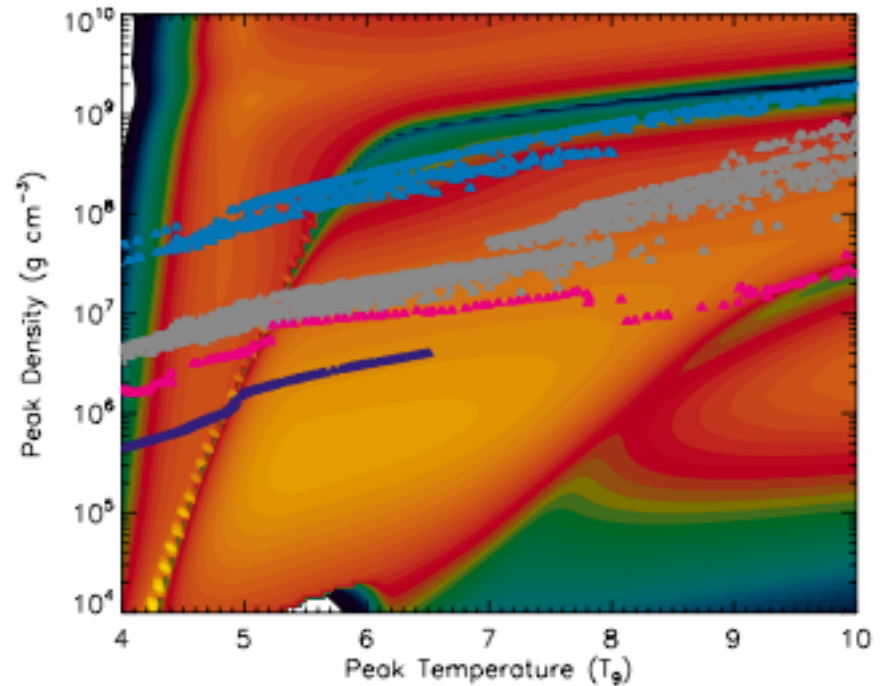
Supernova expel the yields made during stellar evolution. The supernova shock alters these yields, destroying some material and making others.

The elements made in the innermost ejecta are most sensitive to the amount of fallback.



To understand mixing, we must compare elements produced at the engine to remnants, making this a turbulent burn problem.

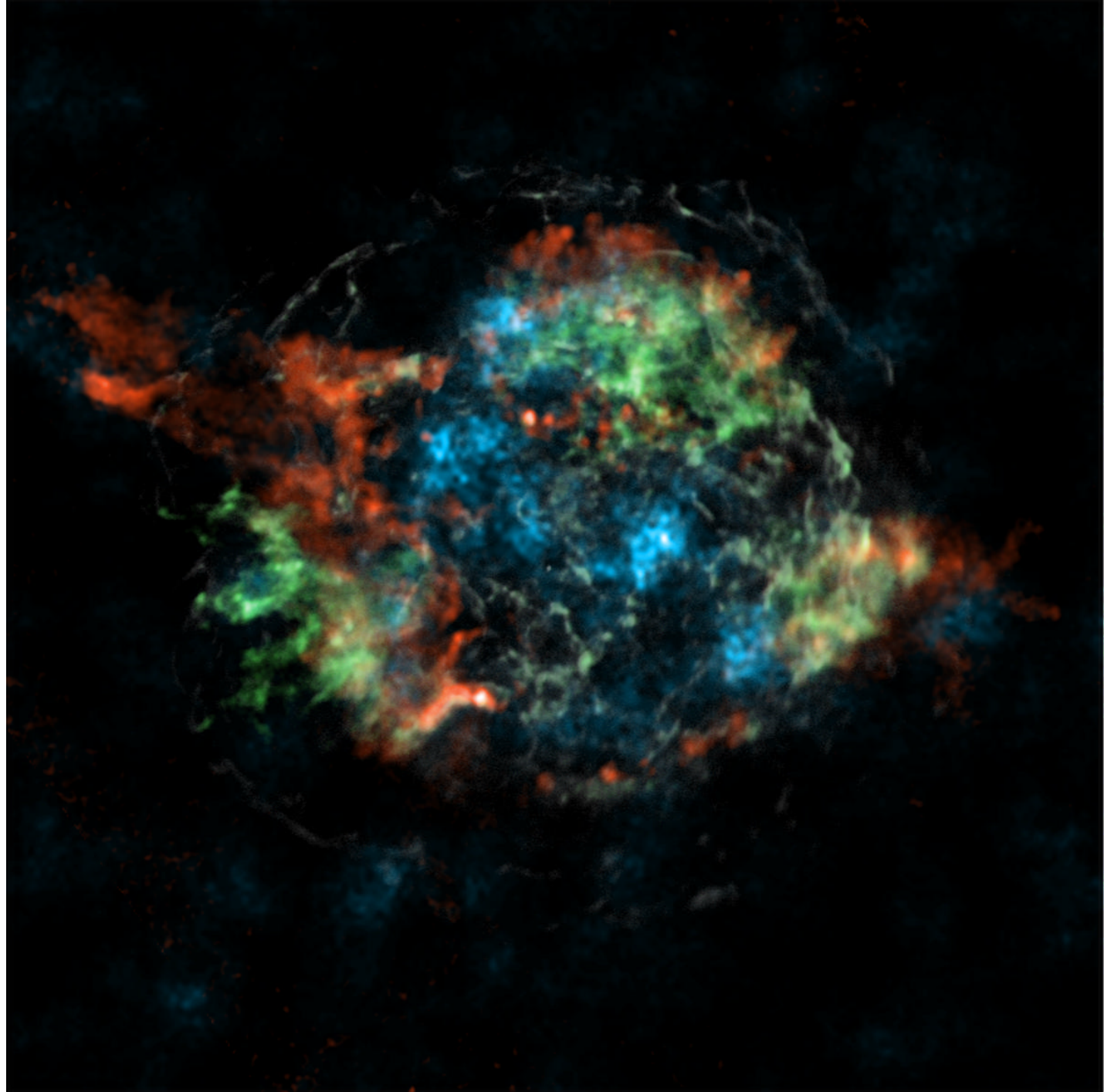
Magkotsios et al. 2010 studied the  $^{56}\text{Ni}$  and  $^{44}\text{Ti}$  yields.





NuSTAR  
demonstrates  
low-mode  
convection in the  
engine  
(Grefenstette et  
al. 2014)

- $^{44}\text{Ti}$  is produced in the core (near the convective engine).
- Because NuSTAR detects the decay emission from  $^{44}\text{Ti}$ , it provides a direct probe of the engine.
- The structure shows a low mode explosion.
- More on observations from Brian Grefenstette and simulations by Janka.



# Conclusions

Convection-enhanced supernova engine explains many aspects of normal supernovae:

- near foe energies
- remnant mass range (yes, but model can predict a range of results)
- yields (but range predicted)
- generic asymmetries
- $^{44}\text{Ti}$  distribution in Cas A

Magnetic field models may explain:

- more energetic explosions (would require tuning to explain peak at  $\sim$ foe energies)
- remnant masses (predictions?)
- yields (with turning)
- generic asymmetries
- **Not Cas A**