

Polarization in a thermal medium

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Based on: E. Speranza, A. Jaiswal, B.F. (to be published)

Space-like and time-like electromagnetic baryonic transitions

ECT* May 8-12, 2017

Polarization without polarization

$$A + B \rightarrow X + \gamma^* \rightarrow X + \ell^+ \ell^-$$

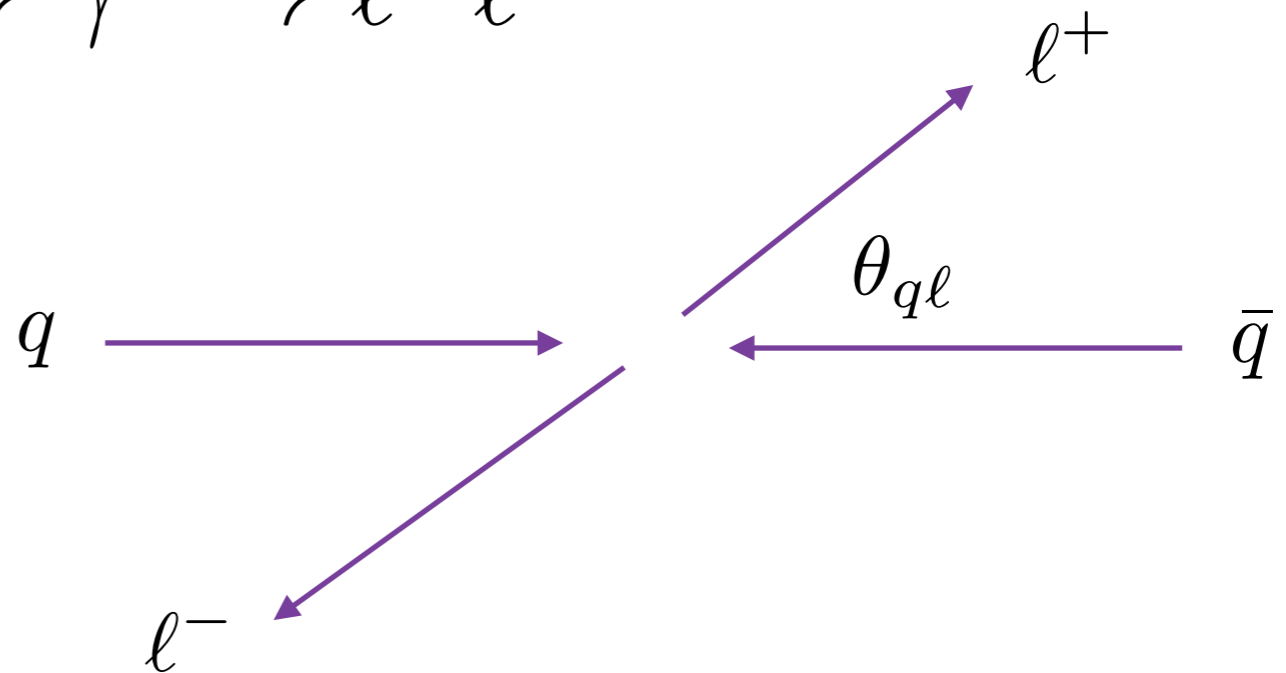
- Polarization of virtual photon depends on the production mechanism
- Polarization of virtual photon is reflected in angular distribution of lepton pair
- Used to constrain production mechanism in elementary reactions (→ Miklos' talk)
- What about nucleus-nucleus collisions?

Drell-Yan process

$$q + \bar{q} \rightarrow \gamma^* \rightarrow \ell^+ \ell^-$$

- In CM frame:

$$\frac{d\sigma}{d\Omega_\ell} \sim 1 + \cos^2 \theta_{q\ell}$$



- Photon polarized along beam axis (quantiz. axis)

$$\pi^+ + \pi^- \rightarrow \gamma^* \rightarrow \ell^+ \ell^-$$

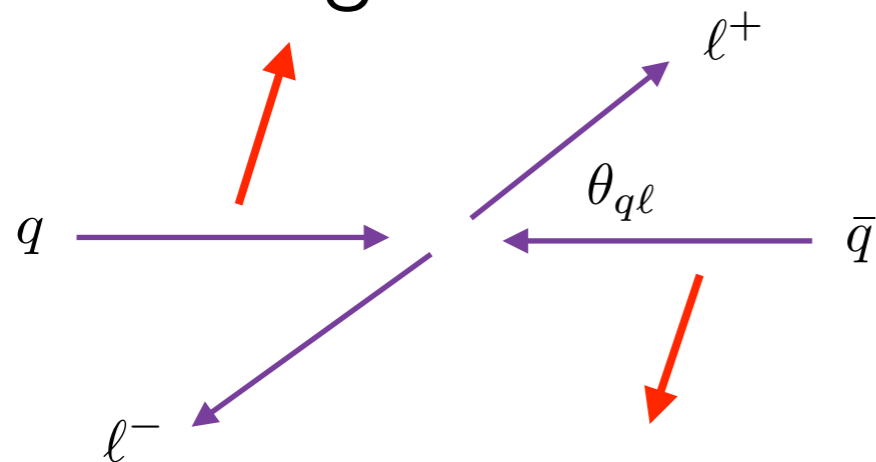
$$\frac{d\sigma}{d\Omega_\ell} \sim 1 - \cos^2 \theta_{q\ell}$$

Emission from a thermal system

- Thermal distribution of initial particles

$$\langle A \rangle = \int \frac{d^3 p_1}{(2\pi)^3 2E_1} \frac{d^3 p_2}{(2\pi)^3 2E_2} (2\pi)^4 \delta^{(4)}(q - p_1 - p_2) \times \frac{1}{e^{(p_1)/T} \pm 1} \frac{1}{e^{(p_2)/T} \pm 1} A$$

- Average over beam axis: polarization vanishes?



P. Hoyer, PLB 187, 162

$$\frac{1}{4\pi} \int d\Omega_q \cos^2 \theta_{ql} = \frac{1}{3}$$

- More detailed study:

E. Bratkovskaya, PLB 348, 283

Polarization of spin-1 particle

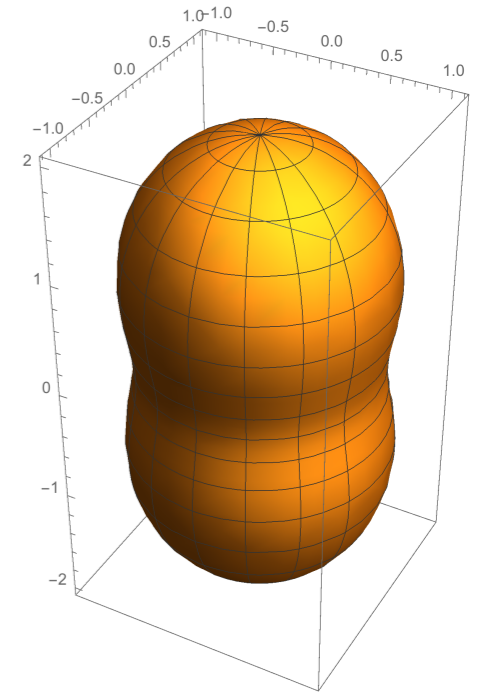
- Three spin states $(\vec{p} = |\vec{p}|\vec{e}_z)$

$$\epsilon(\vec{p}, \lambda = +1) = -\frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ i \\ 0 \end{pmatrix}$$

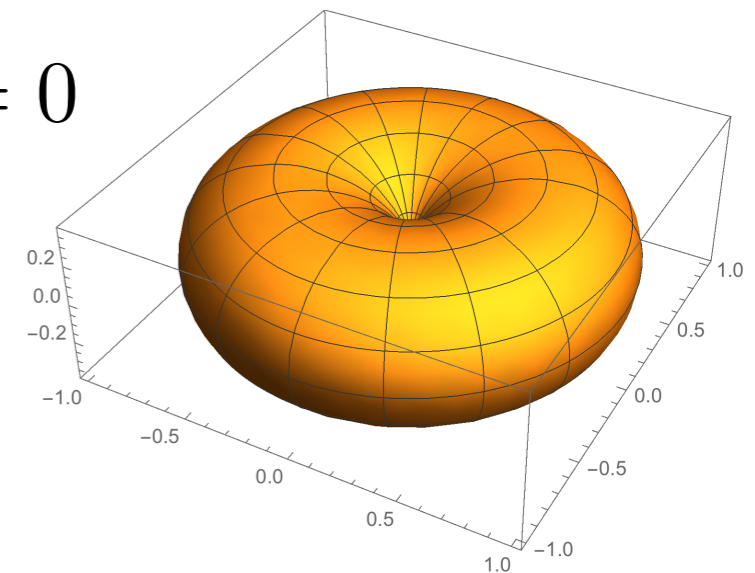
$$\epsilon(\vec{p}, \lambda = 0) = \frac{1}{m} \begin{pmatrix} |\vec{p}| \\ 0 \\ 0 \\ p_0 \end{pmatrix}$$

$$\epsilon(\vec{p}, \lambda = -1) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ -i \\ 0 \end{pmatrix}$$

$$\lambda = \pm 1$$



$$\lambda = 0$$



- Decay of a polarized spin-1 particle

$$\Gamma(\lambda = \pm 1) \sim 1 + \cos^2 \theta_\ell$$

$$\Gamma(\lambda = 0) \sim 2(1 - \cos^2 \theta_\ell)$$

Spin-density matrix

- State of ensemble specified by density matrix $\rho_{\lambda\lambda'}$

- Pure state $|\Psi\rangle = \sum_{\lambda} c_{\lambda} |\lambda\rangle$

- Expectation value of operator \hat{O}

$$\langle \hat{O} \rangle_{\Psi} \equiv \langle \Psi | \hat{O} | \Psi \rangle = \sum_{\lambda\lambda'} c_{\lambda'}^* O_{\lambda'\lambda} c_{\lambda} \quad O_{\lambda'\lambda} = \langle \lambda' | \hat{O} | \lambda \rangle$$

- Mixed state: incoherent mixture of $|\Psi^{(i)}\rangle$

$$\langle \hat{O} \rangle = \sum_i p_i \langle \hat{O} \rangle_{\Psi^{(i)}} = \sum_{\lambda'\lambda} O_{\lambda'\lambda} \rho_{\lambda\lambda'} \quad \rho_{\lambda\lambda'} = \sum_i p_i c_{\lambda}^{(i)} c_{\lambda'}^{*(i)}$$

$$= \text{Tr}(O \rho)$$

Spin-1/2 particle

- 2 x 2 hermitian matrix $\rho = \frac{1}{2}(\mathbf{1} + \vec{\mathcal{P}} \cdot \vec{\sigma})$
- Spin polarization vector $\vec{\mathcal{P}} = \langle \vec{\sigma} \rangle = \text{Tr}(\rho \vec{\sigma})$

- Pure state: $\vec{\mathcal{P}} \cdot \vec{\mathcal{P}} = 1$

- Mixed state: $\vec{\mathcal{P}} \cdot \vec{\mathcal{P}} \leq 1$

$$\rho = \begin{pmatrix} \rho_{++} & \rho_{+-} \\ \rho_{-+} & \rho_{--} \end{pmatrix}$$

- Diagonalize:

$$\rho = \frac{1}{2}(\mathbf{1} + (p_+ - p_-)\sigma_z)$$

$$N_+ = p_+ N$$

$$N_- = p_- N$$

Spin-1 particle

- Spin matrix

$$S_x = \frac{1}{2} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \quad S_y = \frac{1}{2} \begin{pmatrix} 0 & -i & 0 \\ i & 0 & -i \\ 0 & i & 0 \end{pmatrix} \quad S_z = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}$$

- ρ hermitian, trace=1, 3 x 3 matrix: 8 parameters

$$\rho = \frac{1}{3} \left(1 + \frac{3}{2} \vec{\mathcal{P}} \cdot \vec{S} + \sqrt{\frac{3}{2}} T_{ij} (S_i S_j + S_j S_i) \right)$$

- Vector polarization: $\vec{\mathcal{P}}$ (3 parameters)

- Tensor polarization: T_{ij} (5 parameters) $\sum_i T_{ii} = 0$

$$\vec{\mathcal{P}} = \langle \vec{S} \rangle \quad T_{ij} = \frac{1}{2} \sqrt{\frac{3}{2}} \left(\langle S_i S_j + S_j S_i \rangle - \frac{4}{3} \delta_{ij} \right)$$

In Ensemble

- Spin-density matrix can be diagonalized

$$\rho = \frac{1}{3} \begin{pmatrix} 1 + \frac{3}{2}\mathcal{P}_z + \sqrt{\frac{3}{2}}T_{zz} & 0 & 0 \\ 0 & 1 - \sqrt{6}T_{zz} & 0 \\ 0 & 0 & 1 - \frac{3}{2}\mathcal{P}_z + \sqrt{\frac{3}{2}}T_{zz} \end{pmatrix}$$

- In unpolarized system $\mathcal{P}_z = 0$, but often $T_{zz} \neq 0$, i.e., no vector but tensor polarization!
- In general vector and tensor polarization axes can be different.

Lepton angular distribution

- Polarization state of virtual photon reflected in lepton angular distribution
- General form (parity conserved)

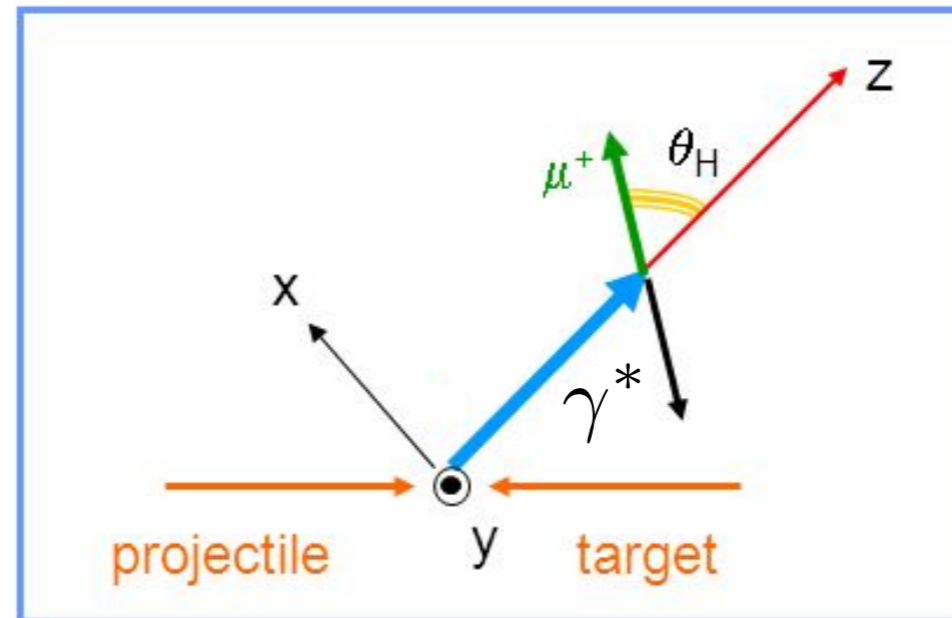
$$\frac{d\Gamma}{d^4q d\Omega_e} \propto \mathcal{N} \left(1 + \lambda_\theta \cos^2 \theta_e + \lambda_\phi \sin^2 \theta_e \cos 2\phi_e + \lambda_{\theta\phi} \sin 2\theta_e \cos \phi_e + \lambda_\phi^\perp \sin^2 \theta_e \sin 2\phi_e + \lambda_{\theta\phi}^\perp \sin 2\theta_e \sin \phi_e \right)$$

- Angles defined in photon rest frame
- Different frames related through rotation: Helicity, Collins-Soper,

Reference frames

- Helicity frame

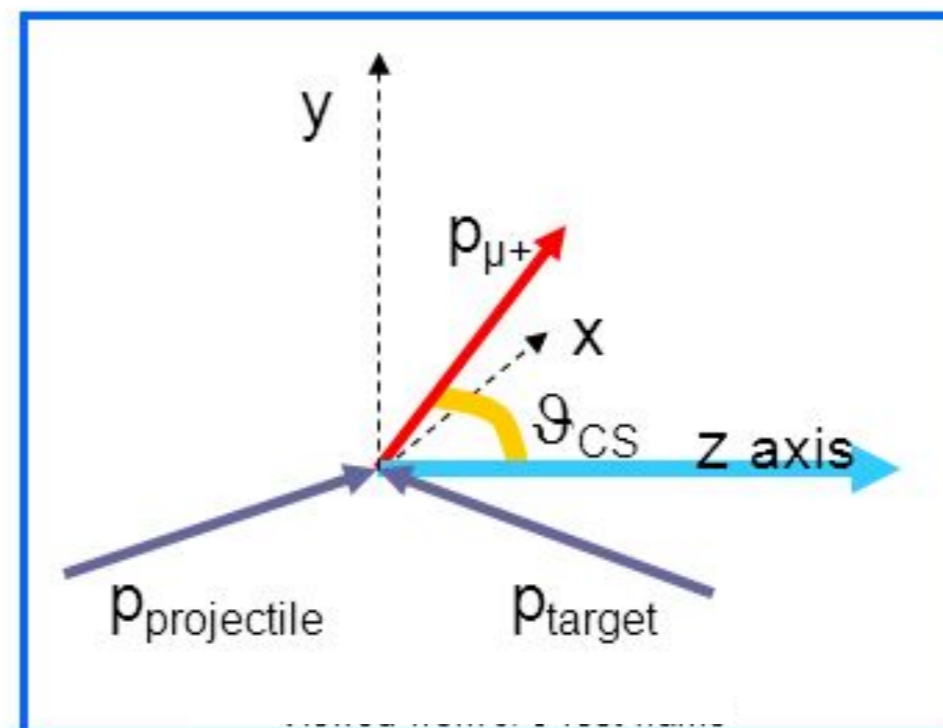
Photon direction
in target-projectile
CM-frame



Rotation

- Collins-Soper frame

Bisector of angle
between “beam”
and “target”



In γ^* rest-frame

Emission from thermal system

- Consider Drell-Yan photon with $q^\mu = (q_0, 0, 0, q_z)$

$$q + \bar{q} \rightarrow \gamma^* \rightarrow \ell^+ \ell^- \quad M^2 = q_0^2 - q_z^2$$

- Quark distribution function $f(p) = \frac{1}{e^{(u \cdot p)/T} + 1}$

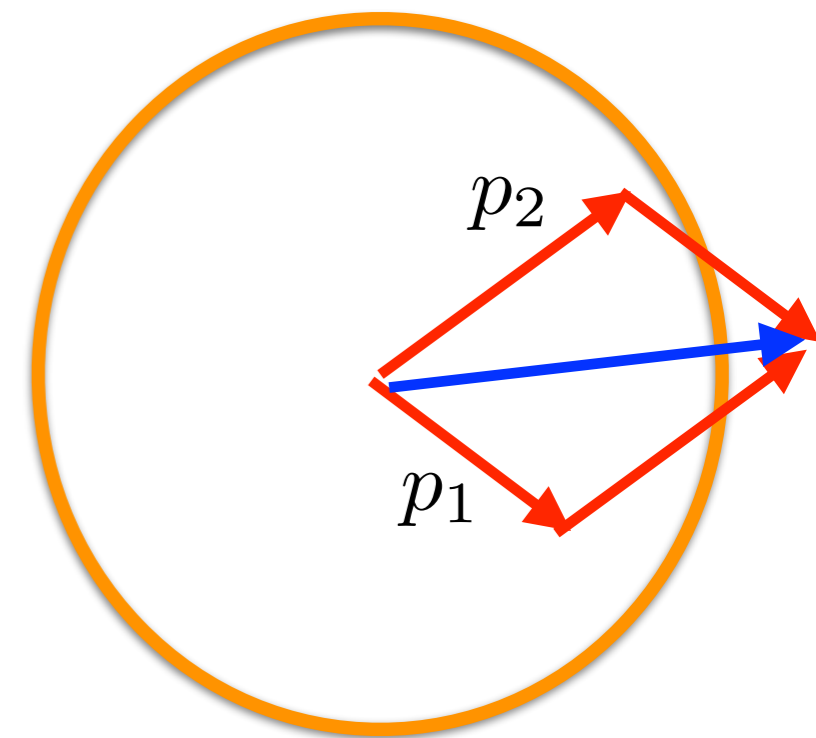
- In fluid rest frame $u^\mu = (1, 0, 0, 0)$

$$q^\mu = p_1^\mu + p_2^\mu$$

- spherically symmetric distribution

- Photon momentum “breaks” spherical symmetry
azimuthal symmetry persists

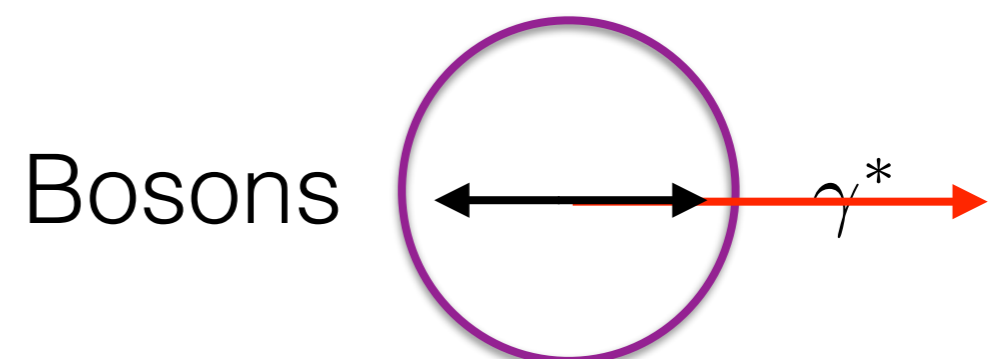
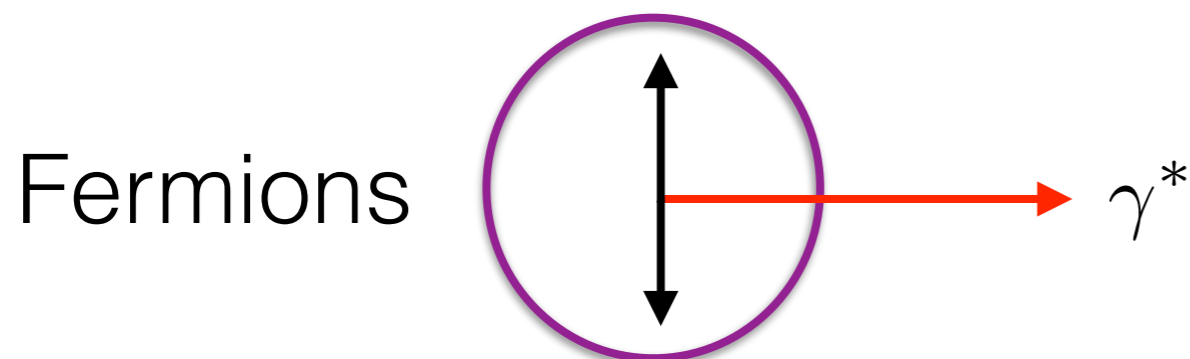
→ tensor polarized photon



Boltzmann limit

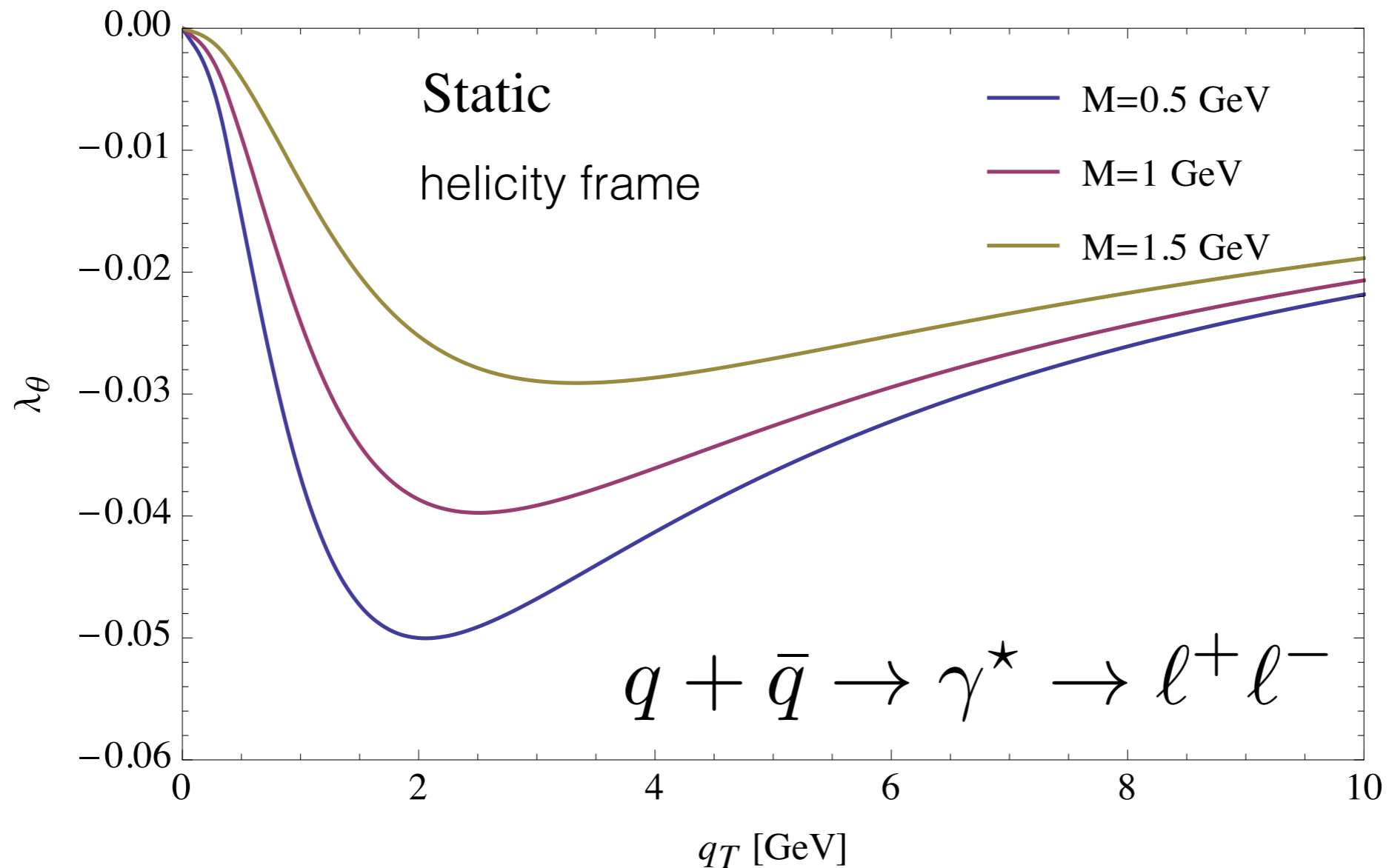
$$\frac{1}{e^{(u \cdot p_1)/T} \pm 1} \frac{1}{e^{(u \cdot p_2)/T} \pm 1} \rightarrow e^{-u \cdot (p_1 + p_2)/T} = e^{-u \cdot q/T}$$

- Independent of direction of quark momenta
 - Yields unpolarized photons
- Non-zero polarization for $0 < |\vec{q}| < \infty$
- Thermal polarization is a quantum statistical effect!



Anisotropy coefficients

- Static homogeneous thermal system



- λ_θ only non-zero coefficient

Effect of hydrodynamic flow

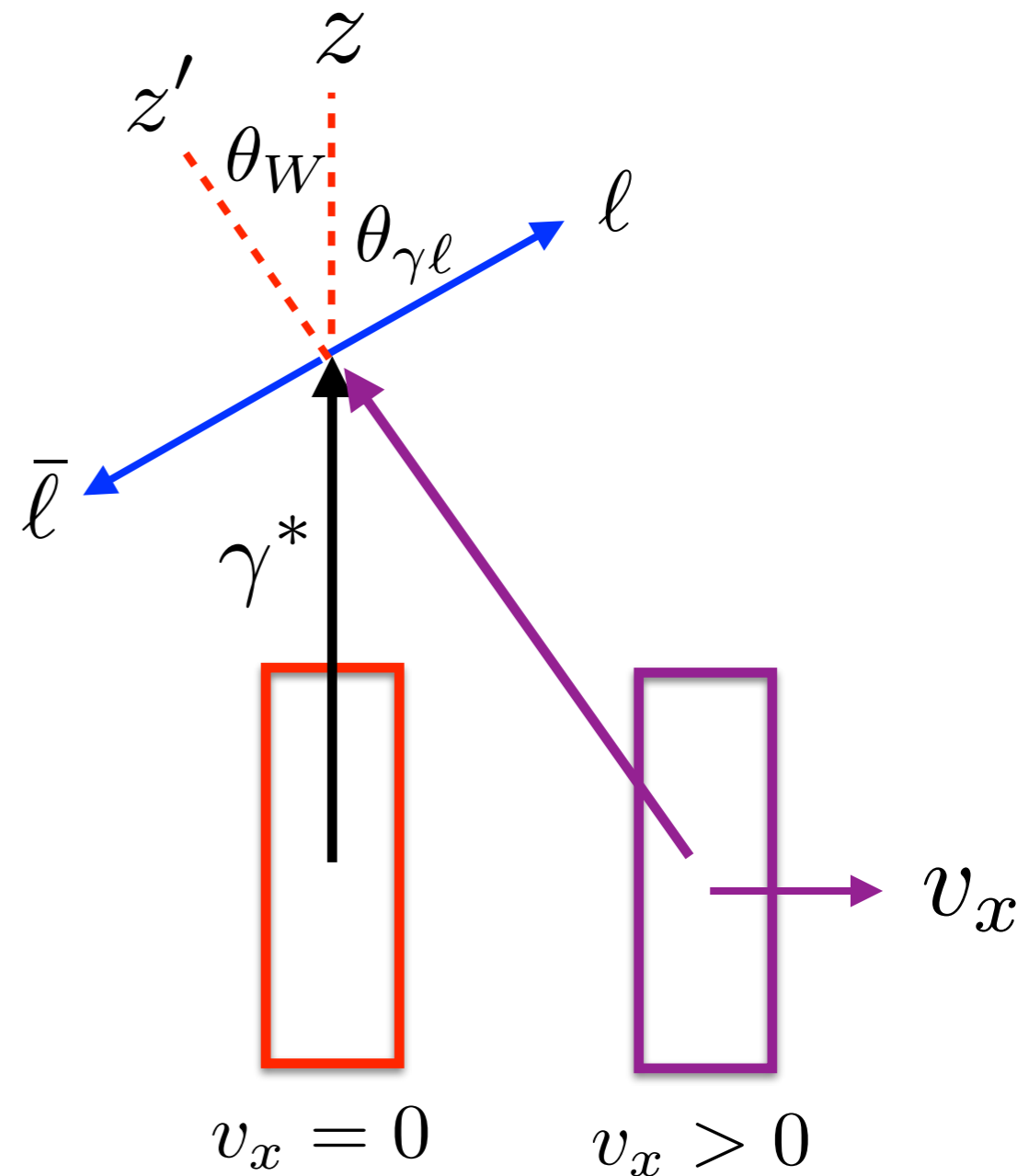
- Bjorken flow along x-axis ($c_s^2 = 1/3$)

$$v_x = x/t$$

$$T(\tau) \sim \tau^{-1/3}$$

$$\tau = \sqrt{t^2 - x^2}$$

- Photon emitted from moving cell: tensor polarized along z' -axis
Wick helicity rotation
- Transform polarization to helicity frame (z).



Anisotropy parameters

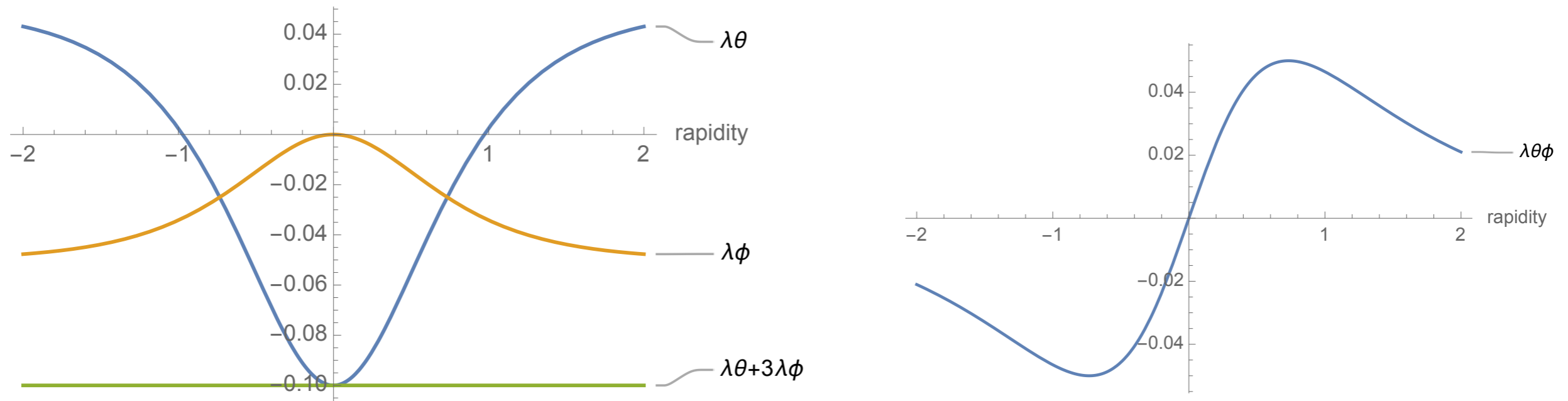
- Tensor polarization along z' :
axial symmetry about z lost!

$$\begin{aligned} \frac{d\Gamma}{d^4q d\Omega_e} \propto \mathcal{N} & \left(1 + \lambda_\theta \cos^2 \theta_e \right. \\ & + \lambda_\phi \sin^2 \theta_e \cos 2\phi_e + \lambda_{\theta\phi} \sin 2\theta_e \cos \phi_e \\ & \left. + \lambda_\phi^\perp \sin^2 \theta_e \sin 2\phi_e + \lambda_{\theta\phi}^\perp \sin 2\theta_e \sin \phi_e \right) \end{aligned}$$

- $\lambda_\theta, \lambda_\phi, \lambda_{\theta\phi} \neq 0, \quad \lambda_\phi^\perp, \lambda_{\theta\phi}^\perp = 0$ (odd under $\phi \rightarrow -\phi$)
- $\lambda_{\theta\phi}$ odd under $v_x \rightarrow -v_x$

Polarization with Bjorken flow

- Anisotropy parameters from moving fluid cell



- Integrate over space-time evolution

$$\frac{dN}{d^2q_T dM^2 dy} = \pi R_A^2 \int_{\tau_i}^{\tau_f} d\tau \tau \int_{-\infty}^{\infty} d\eta \left(\frac{1}{2} \frac{dN}{d^4x d^4q} \right)$$

Drell-Yan

$$T_i = 500 \text{ MeV}$$

$$T_f = 160 \text{ MeV}$$

Pion annihilation

$$T_i = 160 \text{ MeV}$$

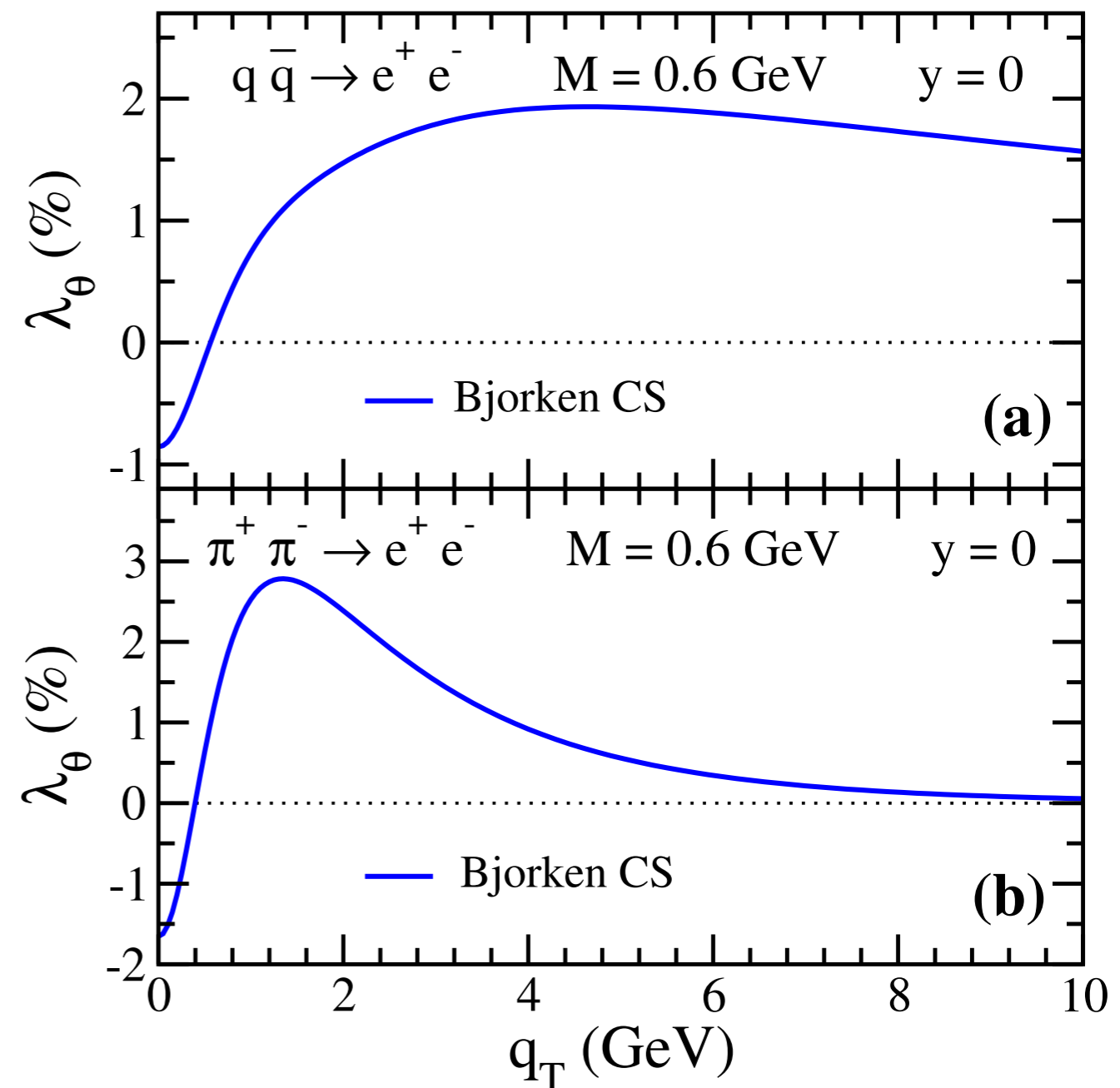
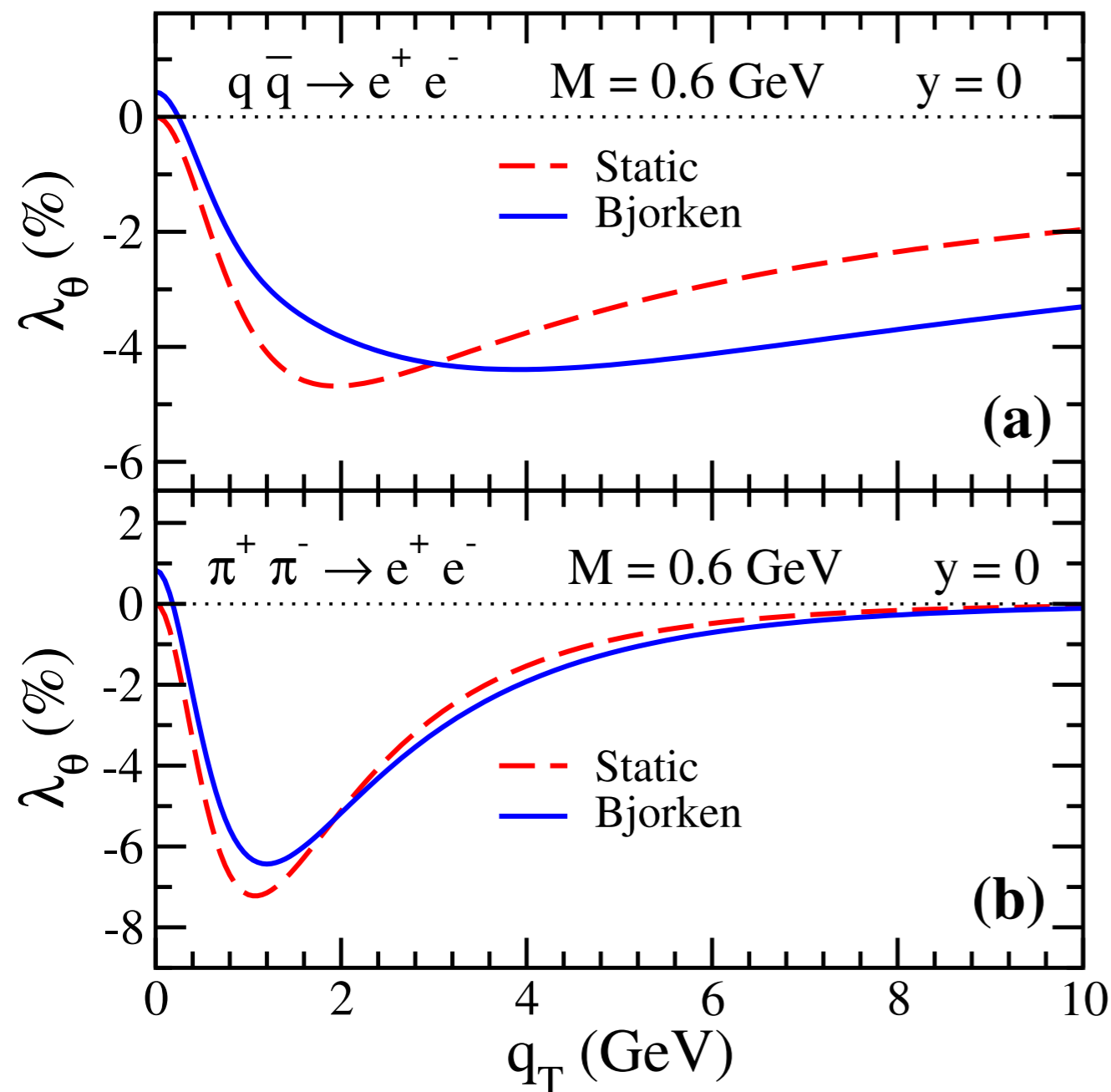
$$T_f = 120 \text{ MeV}$$

Results (Bjorken)

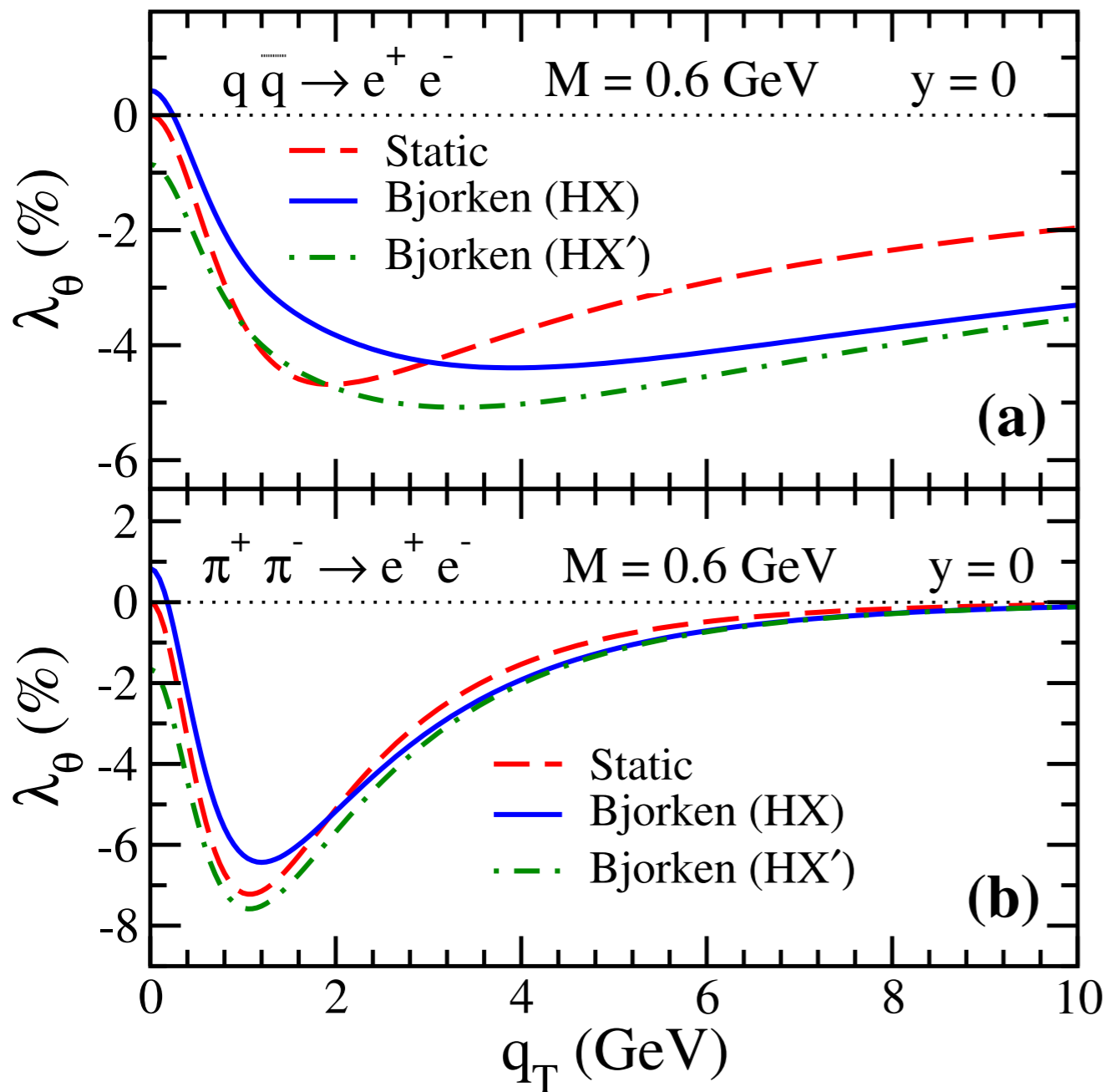
- Photon rapidity $y = 0$ (polarization not boost invariant)

Helicity frame

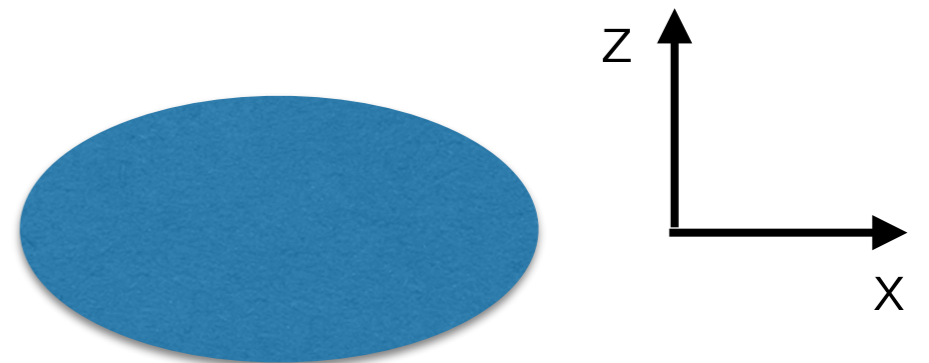
Collins-Soper frame



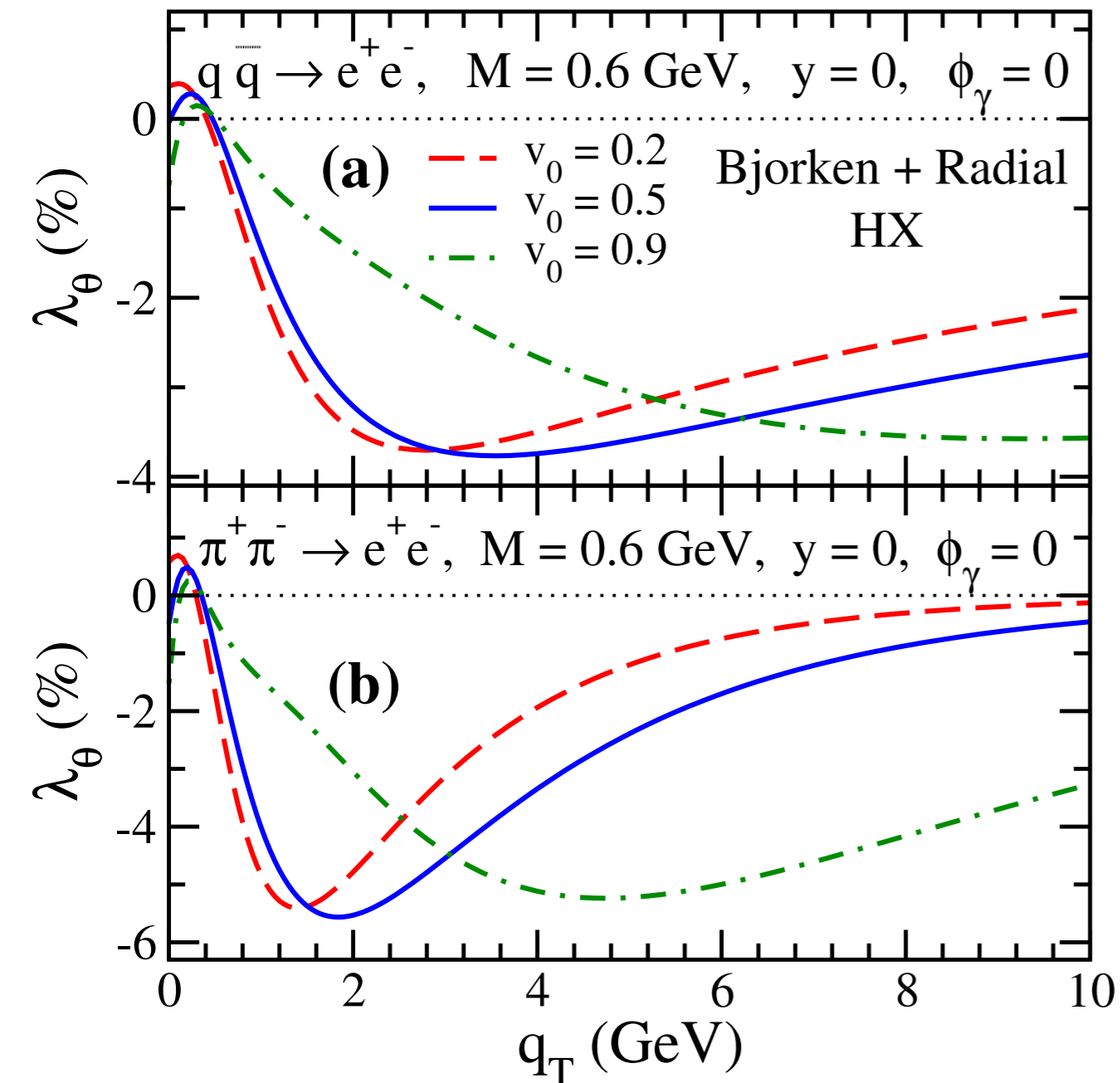
Results ($\lambda_\theta + 3\lambda_\phi$)



- Local thermal equilibrium + longitudinal flow \simeq anisotropic momentum distribution (Baym + Hatsuda + Strickland)
- “Effective” momentum distribution (schematic)

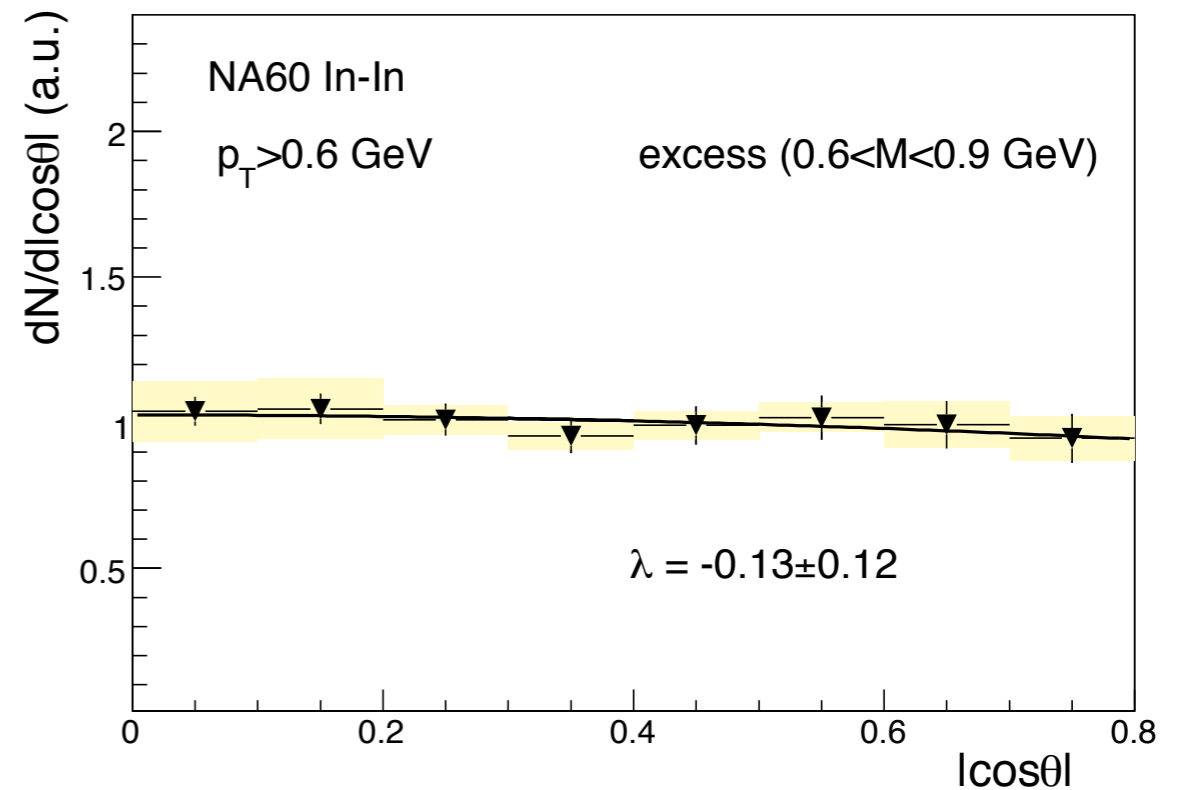
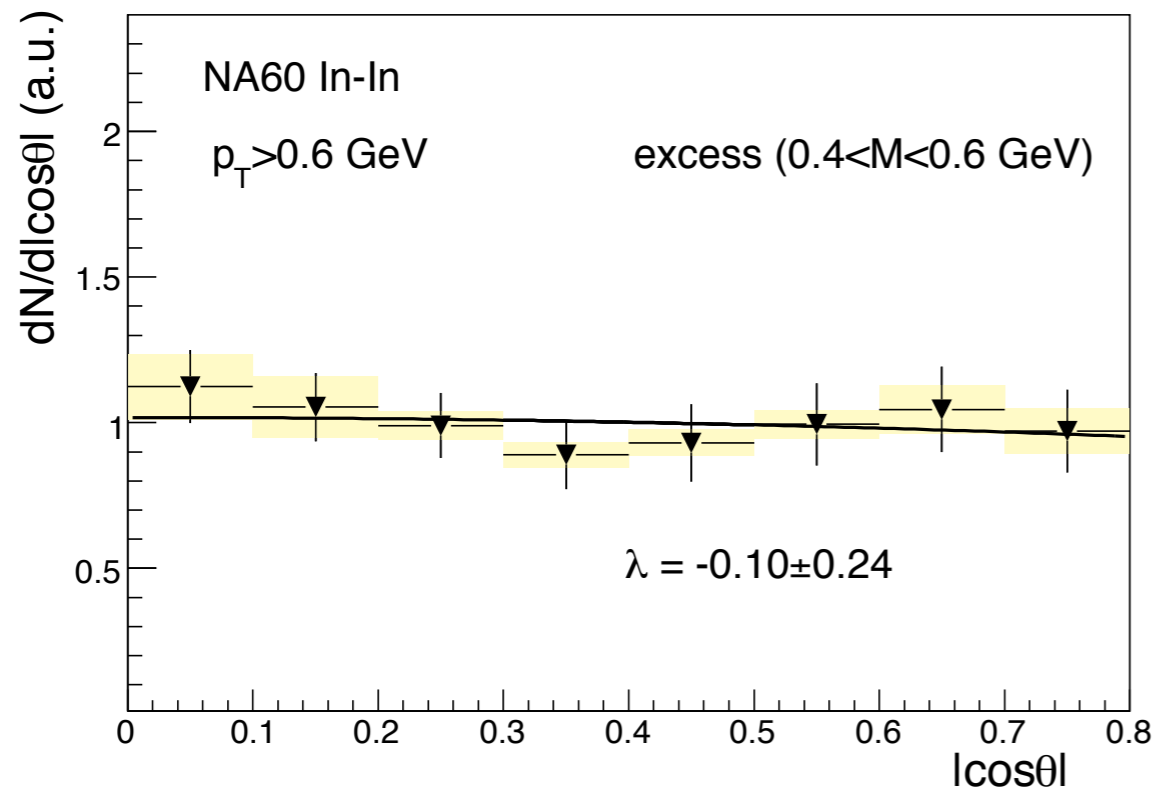


Results (Bjorken + radial flow)



- In general
 $\lambda_\theta, \lambda_\phi, \lambda_{\theta\phi} \neq 0$
- Shape of λ_θ vs. q_T depends on flow velocity

NA-60 data



R. Arnaldi *et al.* PRL 102, 222301

- Analyzed in Collins-Soper frame
- Consistent with $\lambda = 0$
- Need more statistics, better in helicity frame (?)

Summary + Outlook

- Virtual photons from thermal source **are polarized**
- Collective flow modifies shape of λ vs. q_{\perp}
- Small effect, consistent with NA60 (CS vs. HX?)
- Non-equilibrium (deformed momentum distribution)
 - Stronger effect?** Baym + Hatsuda + Strickland
- Polarization due to magnetic fields & rotation?

Two directions:
photon momentum + polarization axis
expect richer structure