

Flavor Anomalies on the Eve of the Run-2 Verdict

Diego Guadagnoli
LAPTh Annecy (France)

Flavor anomalies

① ***A first qualitative observation***

A whole range of $b \rightarrow s$ measurements involving a $\mu\mu$ pair display a consistent pattern:

$\text{Exp} < \text{SM}$

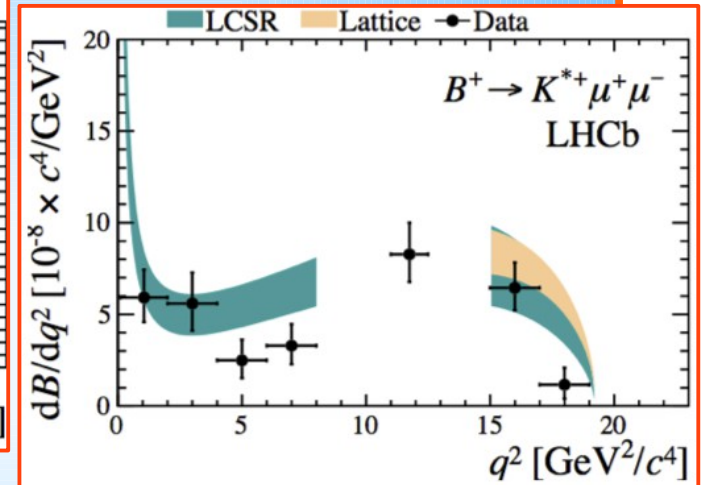
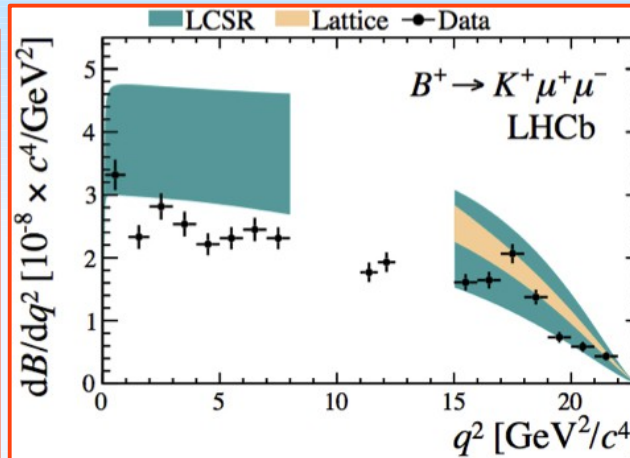
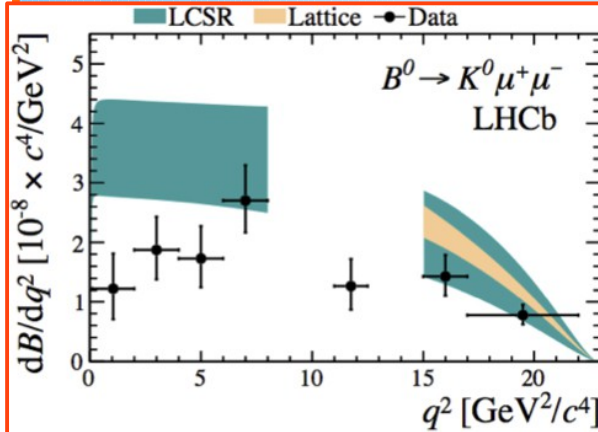
Flavor anomalies

① A first qualitative observation

A whole range of $b \rightarrow s$ measurements involving a $\mu\mu$ pair display a consistent pattern:

$\text{Exp} < \text{SM}$

LHCb 2014



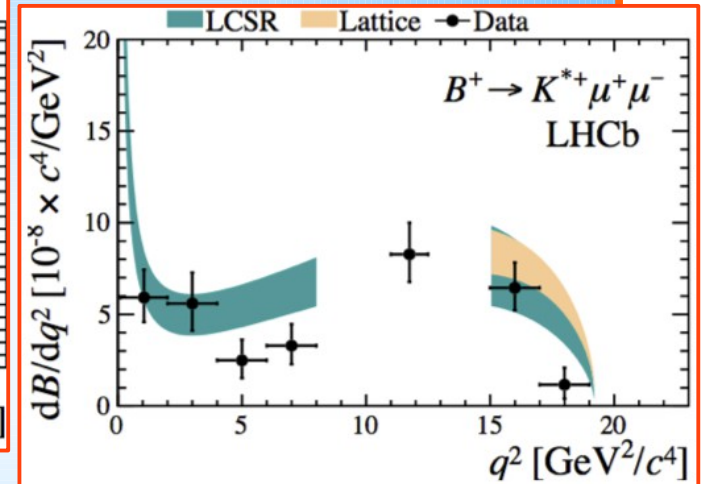
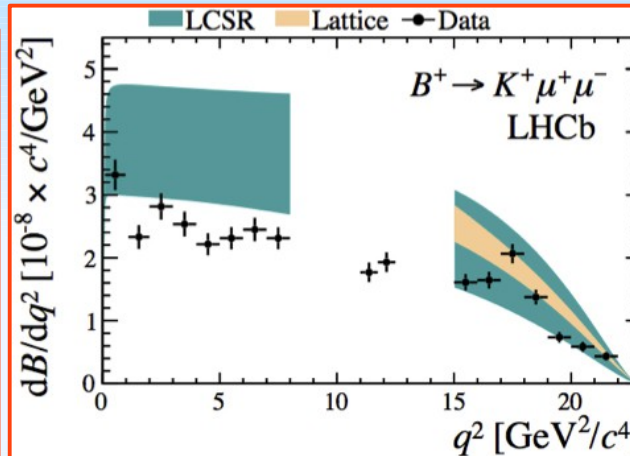
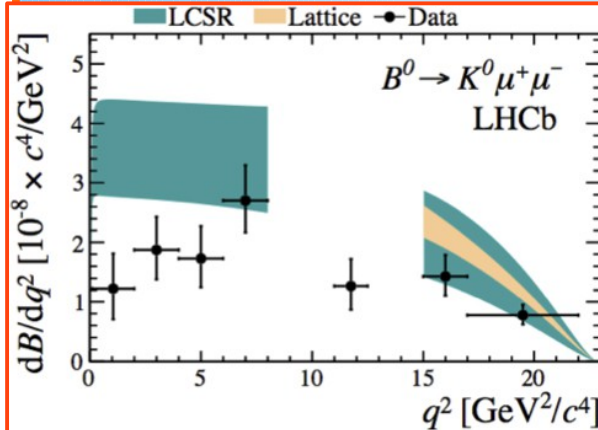
Flavor anomalies

① A first qualitative observation

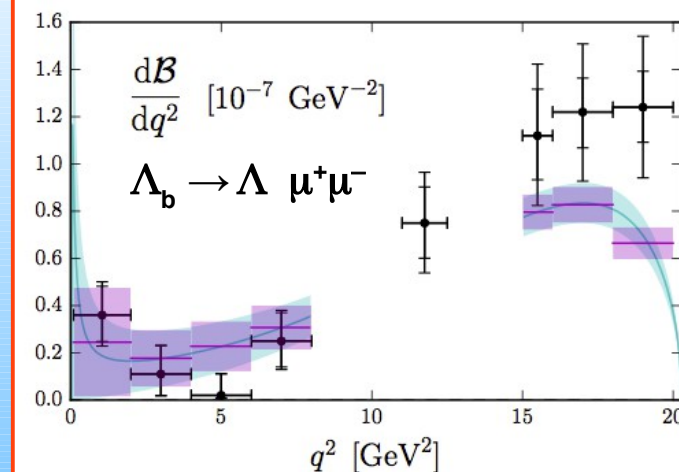
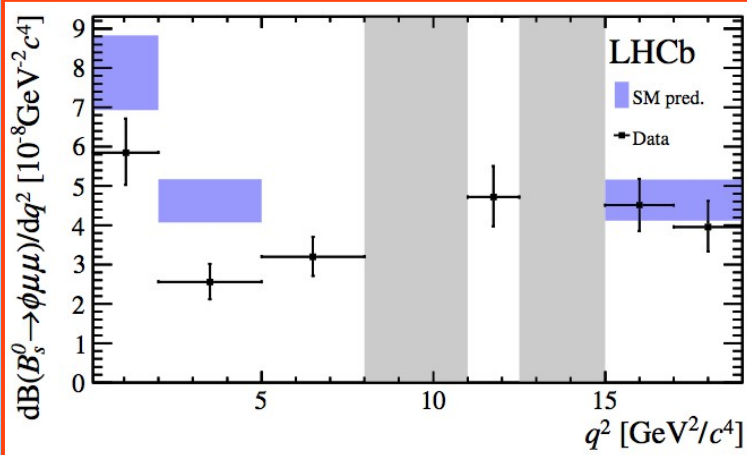
A whole range of $b \rightarrow s$ measurements involving a $\mu\mu$ pair display a consistent pattern:

$\text{Exp} < \text{SM}$

LHCb 2014



LHCb 2015



LHCb 2015
& Detmold+Meinel '16

b → s data

*We know that BR measurements suffer from large f.f. uncertainties.
However, here's a clean quantity:*

$$\textcircled{1} \quad R_K(q_{\min}^2, q_{\max}^2) \equiv \frac{BR(B^+ \rightarrow K^+ \mu \mu)}{BR(B^+ \rightarrow K^+ e e)} \Big|_{[q_{\min}^2, q_{\max}^2]}$$

b → s data

*We know that BR measurements suffer from large f.f. uncertainties.
However, here's a clean quantity:*

$$\textcircled{1} \quad R_K(q_{\min}^2, q_{\max}^2) \equiv \frac{BR(B^+ \rightarrow K^+ \mu \mu)}{BR(B^+ \rightarrow K^+ e e)} \Big|_{[q_{\min}^2, q_{\max}^2]}$$

$$R_K(1\text{GeV}^2, 6\text{GeV}^2) = 0.745 \cdot (1 \pm 13\%)$$

(2.6σ effect)

b → s data

*We know that BR measurements suffer from large f.f. uncertainties.
However, here's a clean quantity:*

$$\textcircled{1} \quad R_K(q_{\min}^2, q_{\max}^2) \equiv \frac{BR(B^+ \rightarrow K^+ \mu \mu)}{BR(B^+ \rightarrow K^+ e e)} \Big|_{[q_{\min}^2, q_{\max}^2]}$$

$$R_K(1\text{GeV}^2, 6\text{GeV}^2) = 0.745 \cdot (1 \pm 13\%)$$

(2.6σ effect)

And here's another (freshly measured) one:

$$\textcircled{2} \quad R_{K^{*0}}(1.1\text{GeV}^2, 6.0\text{GeV}^2) = 0.685_{-0.069}^{+0.113} \pm 0.047$$

(~2.4σ effect)

b → s data

*We know that BR measurements suffer from large f.f. uncertainties.
However, here's a clean quantity:*

$$1 \quad R_K(q_{\min}^2, q_{\max}^2) \equiv \frac{BR(B^+ \rightarrow K^+ \mu \mu)}{BR(B^+ \rightarrow K^+ e e)} \Big|_{[q_{\min}^2, q_{\max}^2]}$$

$$R_K(1\text{GeV}^2, 6\text{GeV}^2) = 0.745 \cdot (1 \pm 13\%)$$

(2.6σ effect)

And here's another (freshly measured) one:

$$2 \quad R_{K^{*0}}(1.1\text{GeV}^2, 6.0\text{GeV}^2) = 0.685_{-0.069}^{+0.113} \pm 0.047$$

(~2.4σ effect)

$$R_{K^{*0}}(0.045\text{GeV}^2, 1.1\text{GeV}^2) = 0.660_{-0.070}^{+0.110} \pm 0.024$$

(~2.2σ effect)

b → s data

*We know that BR measurements suffer from large f.f. uncertainties.
However, here's a clean quantity:*

$$1 \quad R_K(q_{\min}^2, q_{\max}^2) \equiv \frac{BR(B^+ \rightarrow K^+ \mu \mu)}{BR(B^+ \rightarrow K^+ e e)} \Big|_{[q_{\min}^2, q_{\max}^2]}$$

$$R_K(1\text{GeV}^2, 6\text{GeV}^2) = 0.745 \cdot (1 \pm 13\%)$$

(2.6σ effect)

And here's another (freshly measured) one:

$$2 \quad R_{K^{*0}}(1.1\text{GeV}^2, 6.0\text{GeV}^2) = 0.685_{-0.069}^{+0.113} \pm 0.047$$

(~2.4σ effect)

$$R_{K^{*0}}(0.045\text{GeV}^2, 1.1\text{GeV}^2) = 0.660_{-0.070}^{+0.110} \pm 0.024$$

(~2.2σ effect)

- *the electron channel would be an obvious culprit (brems + low stats).
But disagreement is rather in muons*
- *muons are among the most reliable objects within LHCb*

$b \rightarrow s$ data

The other mentioned $b \rightarrow s \mu\mu$ modes fit a coherent picture with R_K and R_{K^}*

$b \rightarrow s$ data

The other mentioned $b \rightarrow s \mu\mu$ modes fit a coherent picture with R_K and R_{K^}*

- ③ **$BR(B_s \rightarrow \phi \mu\mu)$** : $>3\sigma$ below SM prediction. Same kinematical region $m_{\mu\mu}^2 \in [1, 6] \text{ GeV}^2$
Initially found in 1/fb of LHCb data, then confirmed by a full Run-I analysis (3/fb)

$b \rightarrow s$ data

The other mentioned $b \rightarrow s \mu\mu$ modes fit a coherent picture with R_K and R_{K^}*

- ③ **$BR(B_s \rightarrow \phi \mu\mu)$** : $>3\sigma$ below SM prediction. Same kinematical region $m_{\mu\mu}^2 \in [1, 6] \text{ GeV}^2$
Initially found in 1/fb of LHCb data, then confirmed by a full Run-I analysis (3/fb)
- ④ **$B \rightarrow K^* \mu\mu$ angular analysis**: discrepancy in one combination of the angular expansion coefficients, known as P'_5

4

 $B \rightarrow K^* \mu\mu$ angular analysis:*The P'_5 anomaly*

- *From LHCb's full angular analysis of the decay products in $B \rightarrow K^* \mu\mu$, one can construct observables with limited sensitivity to form factors.*

4

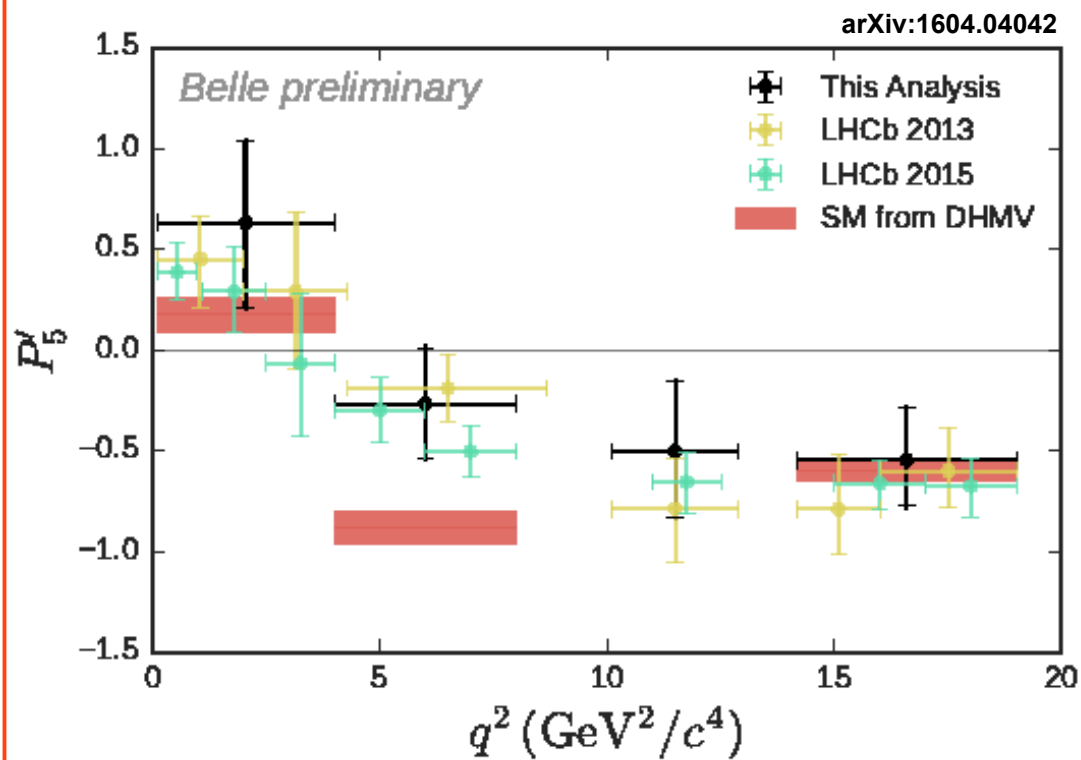
 $B \rightarrow K^* \mu\mu$ angular analysis:*The P'_5 anomaly*

- From LHCb's full angular analysis of the decay products in $B \rightarrow K^* \mu\mu$, one can construct observables with limited sensitivity to form factors.
- One of such “clean” observables is called P'_5

4

 $B \rightarrow K^* \mu\mu$ angular analysis:**The P'_5 anomaly**

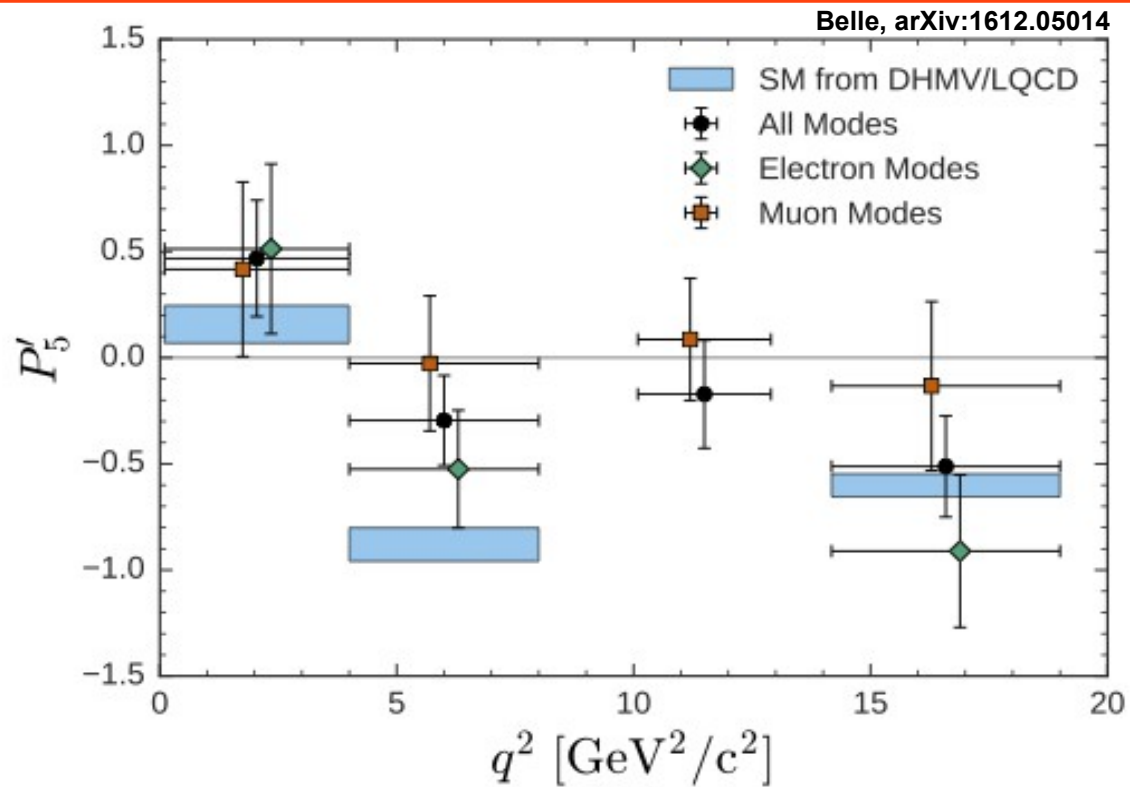
- From LHCb's full angular analysis of the decay products in $B \rightarrow K^* \mu\mu$, one can construct observables with limited sensitivity to form factors.
- One of such “clean” observables is called P'_5



4

 $B \rightarrow K^* \mu\mu$ angular analysis:**The P'_5 anomaly**

- From LHCb's full angular analysis of the decay products in $B \rightarrow K^* \mu\mu$, one can construct observables with limited sensitivity to form factors.
- One of such “clean” observables is called P'_5

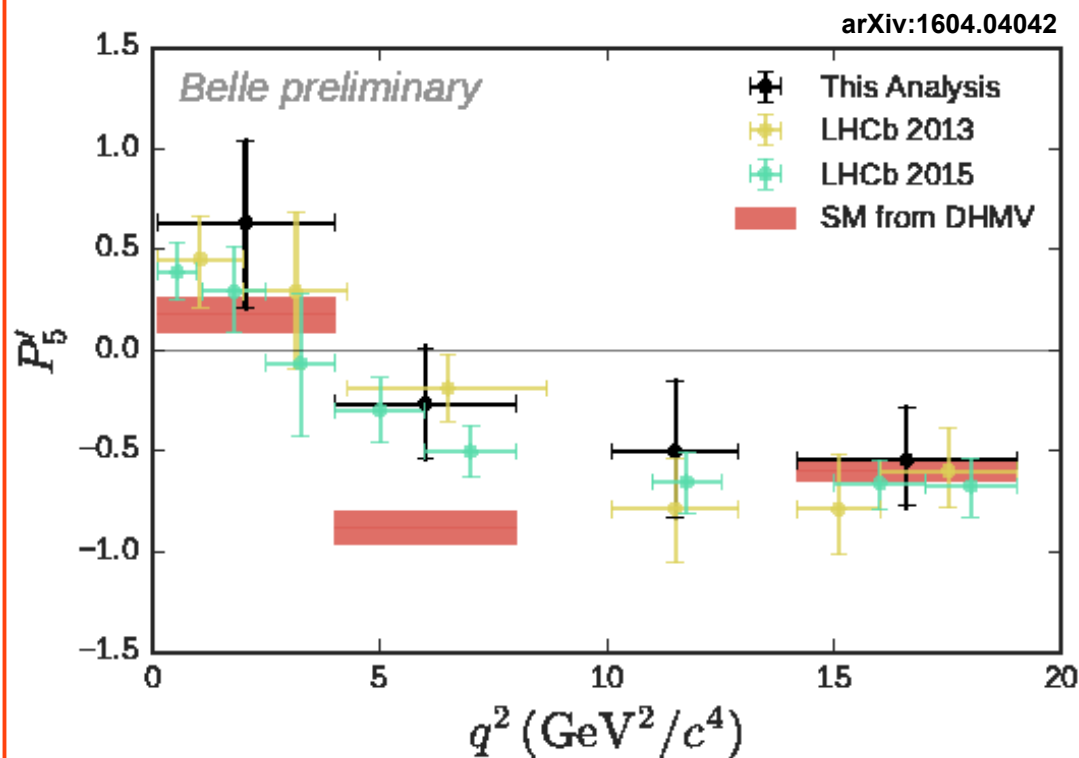


4

$B \rightarrow K^* \mu\mu$ angular analysis:

The P'_5 anomaly

- From LHCb's full angular analysis of the decay products in $B \rightarrow K^* \mu\mu$, one can construct observables with limited sensitivity to form factors.
- One of such "clean" observables is called P'_5



Caveat:

this obs needs be taken cum grano salis

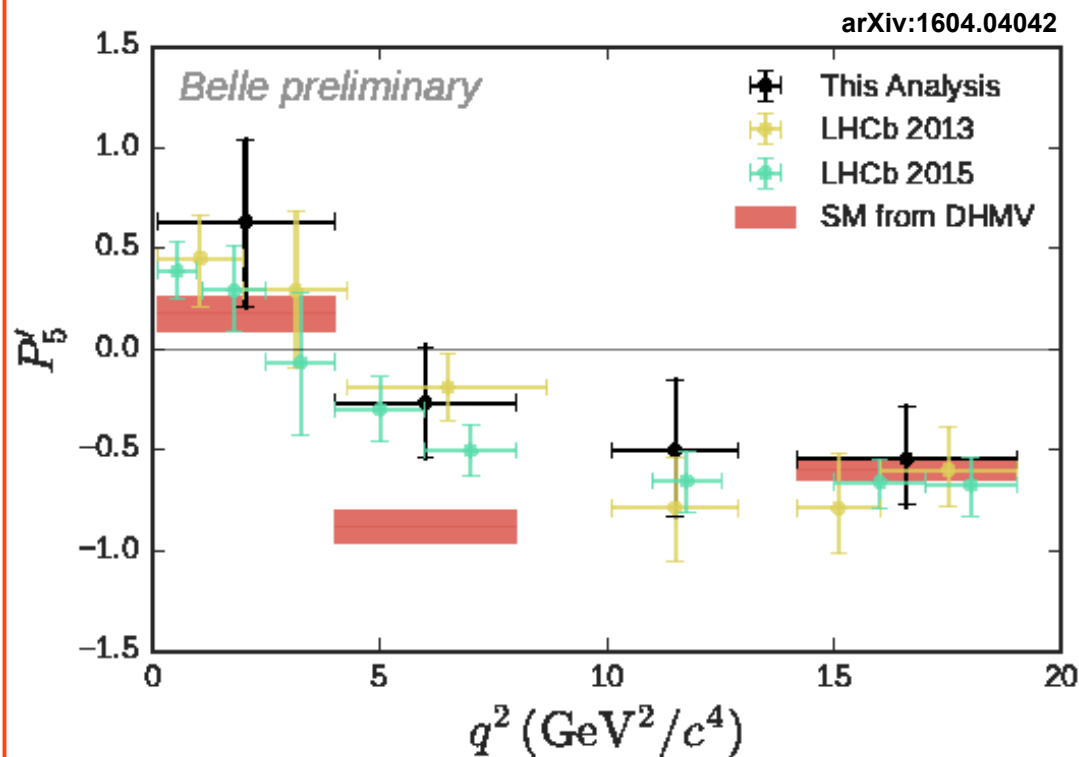
- What cancels is the dependence on the large- m_b form factors.

4

$B \rightarrow K^* \mu\mu$ angular analysis:

The P'_5 anomaly

- From LHCb's full angular analysis of the decay products in $B \rightarrow K^* \mu\mu$, one can construct observables with limited sensitivity to form factors.
- One of such "clean" observables is called P'_5



Caveat:

this obs needs be taken cum grano salis

- What cancels is the dependence on the large- m_b form factors.
- Crucial issue:

How important departures from the infinite- m_b limit are, for q^2 approaching $4 m_c^2$.

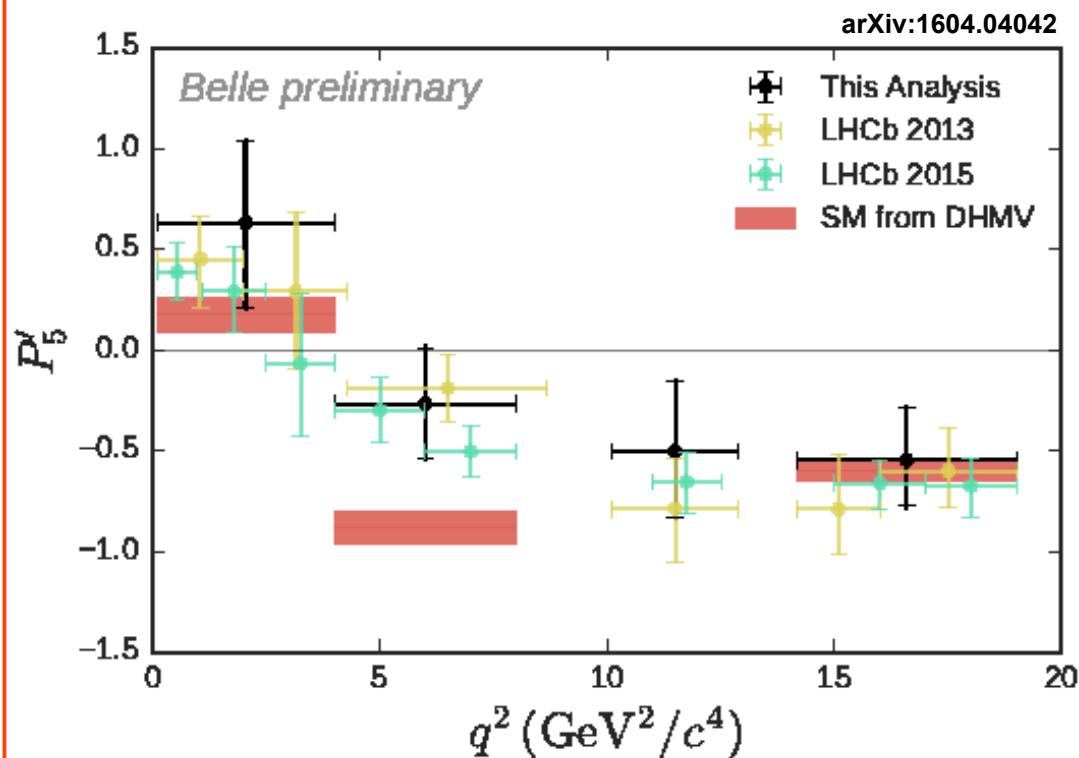
In fact, $c\bar{c}$ contributions are suppressed by $q^2 - 4 m_c^2$.

4

$B \rightarrow K^* \mu\mu$ angular analysis:

The P'_5 anomaly

- From LHCb's full angular analysis of the decay products in $B \rightarrow K^* \mu\mu$, one can construct observables with limited sensitivity to form factors.
- One of such "clean" observables is called P'_5



Caveat:

this obs needs be taken cum grano salis

- What cancels is the dependence on the large- m_b form factors.

- Crucial issue:

How important departures from the infinite- m_b limit are, for q^2 approaching $4 m_c^2$.

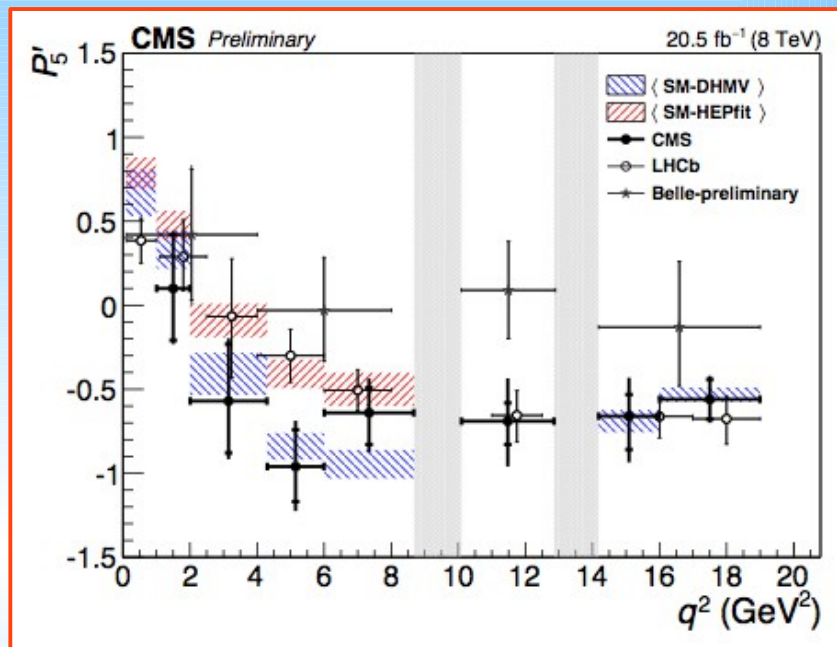
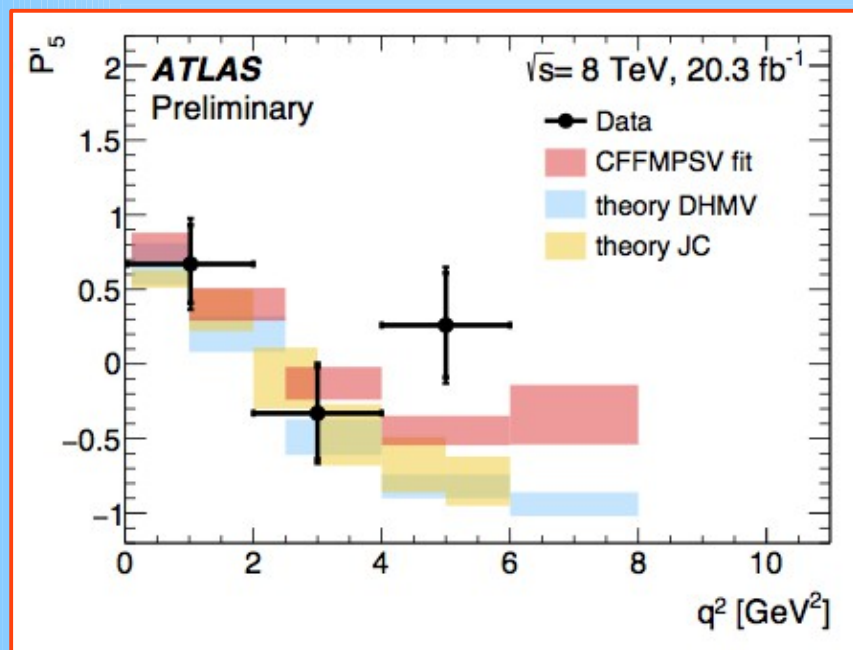
In fact, $c\bar{c}$ contributions are suppressed by $q^2 - 4 m_c^2$.

But interesting nonetheless, because:

- Effect is again in the same region:
 $m_{\mu\mu}^2 \in [1, 6] \text{ GeV}^2$
- Compatibility between 1/fb and 3/fb LHCb analyses and a recent Belle analysis

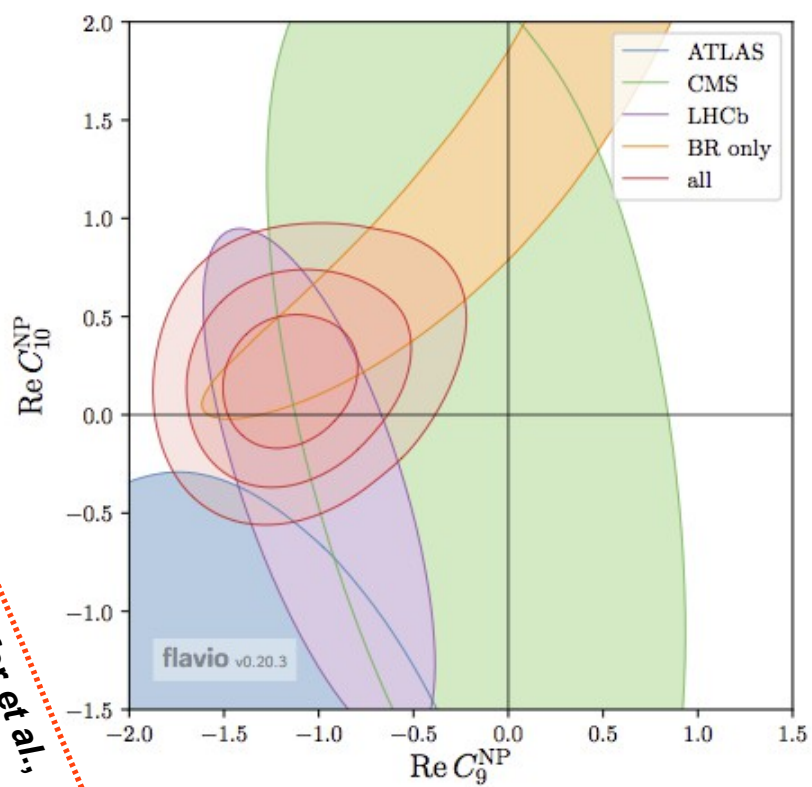
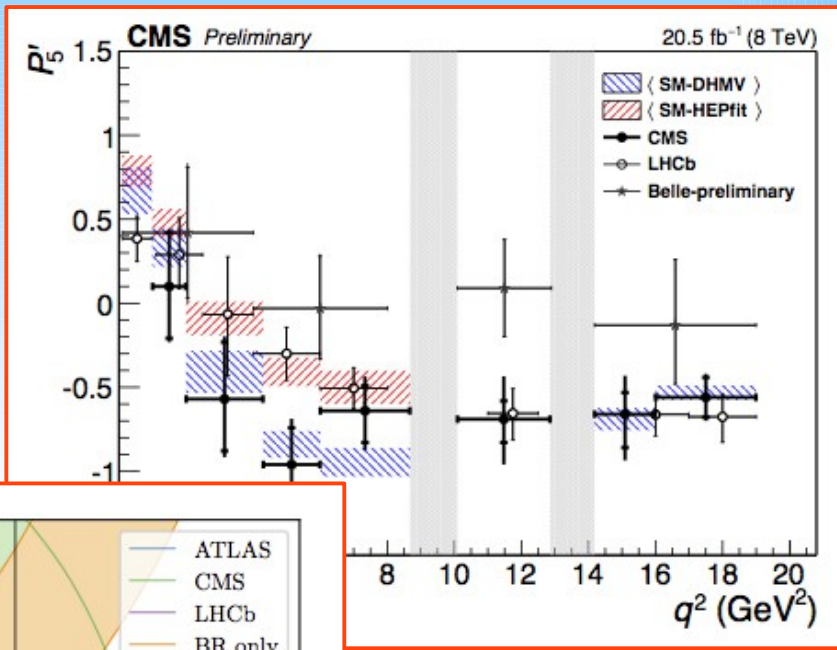
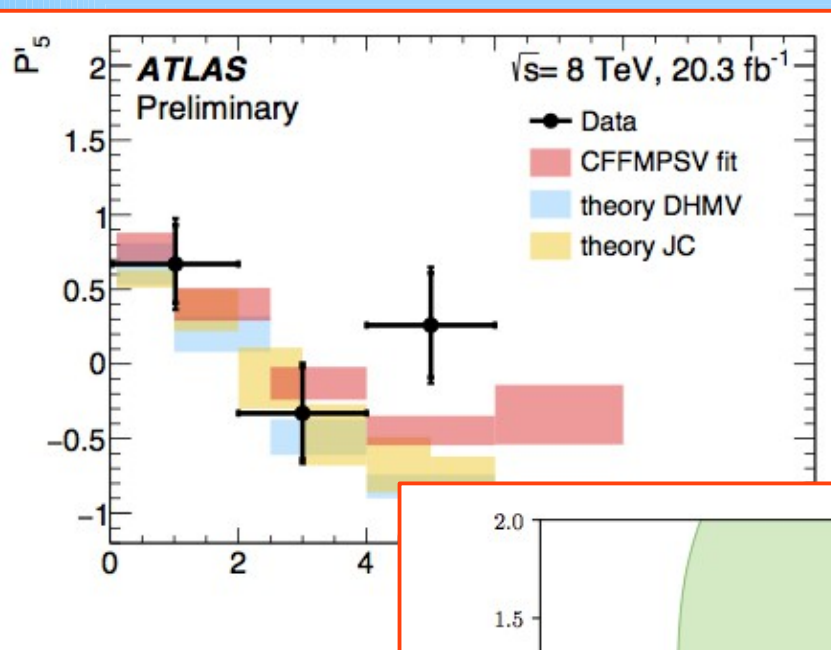
4

$B \rightarrow K^* \mu\mu$ angular analysis:
post-Moriond update



4

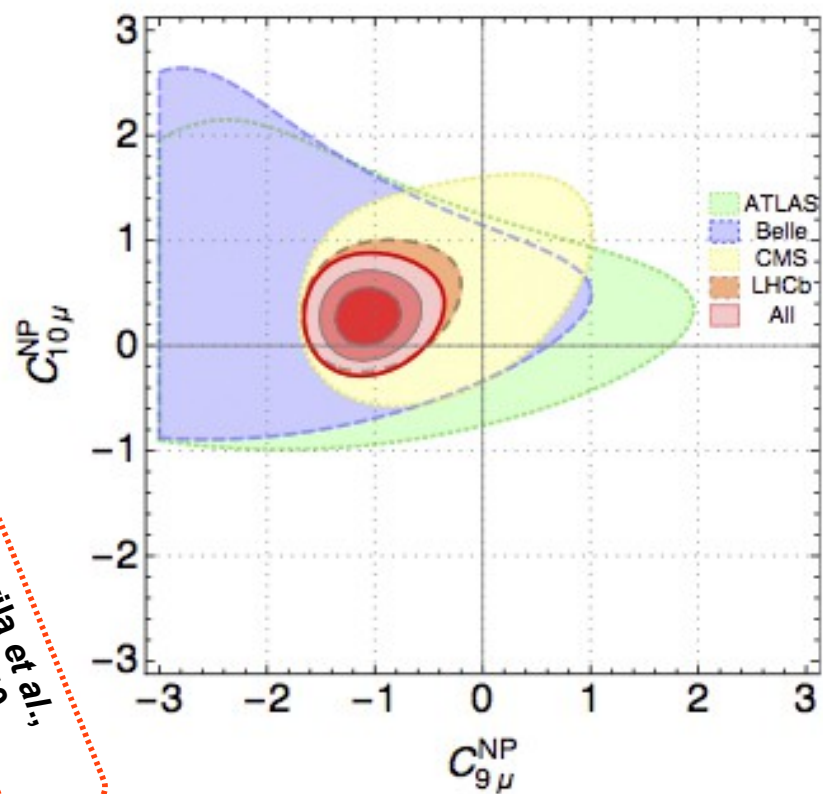
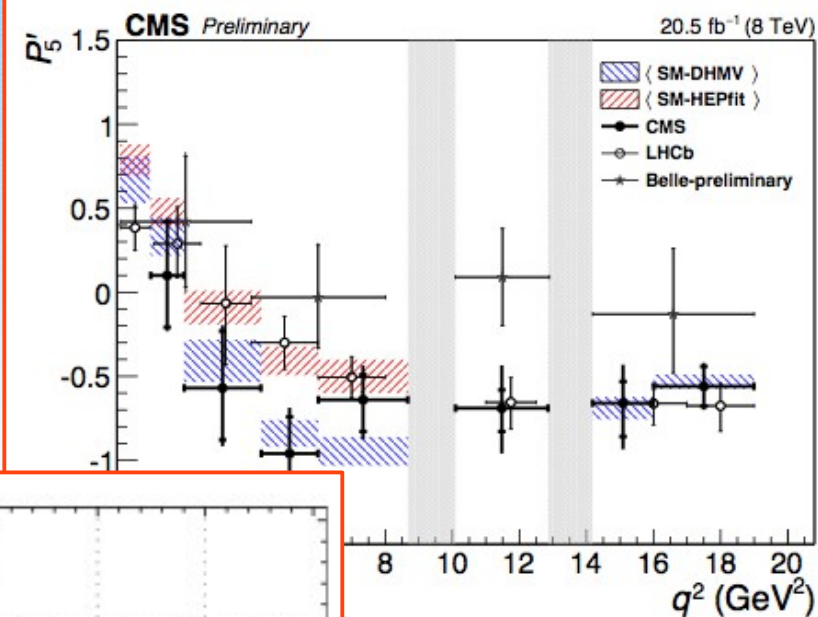
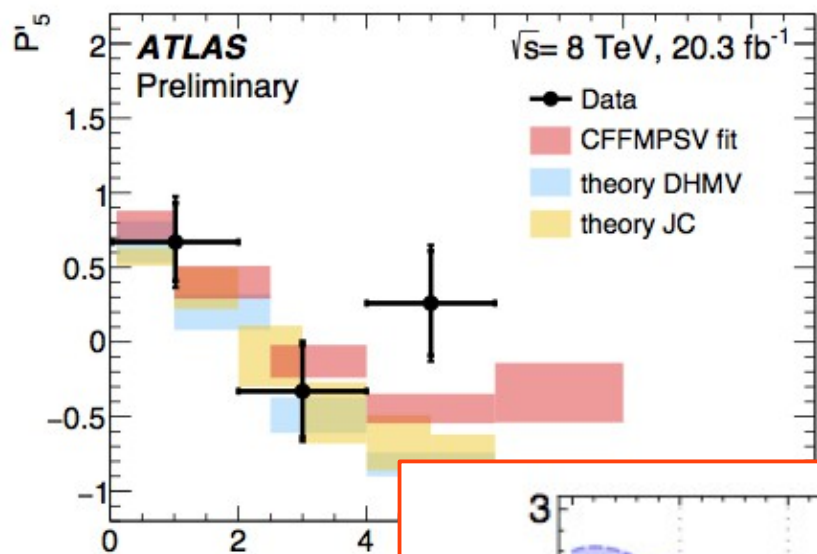
$B \rightarrow K^* \mu\mu$ angular analysis:
post-Moriond update



Altmannshofer et al.,
1703.09189

4

$B \rightarrow K^* \mu\mu$ angular analysis:
post-Moriond update



Capdevila et al.,
1704.05340

b → c data

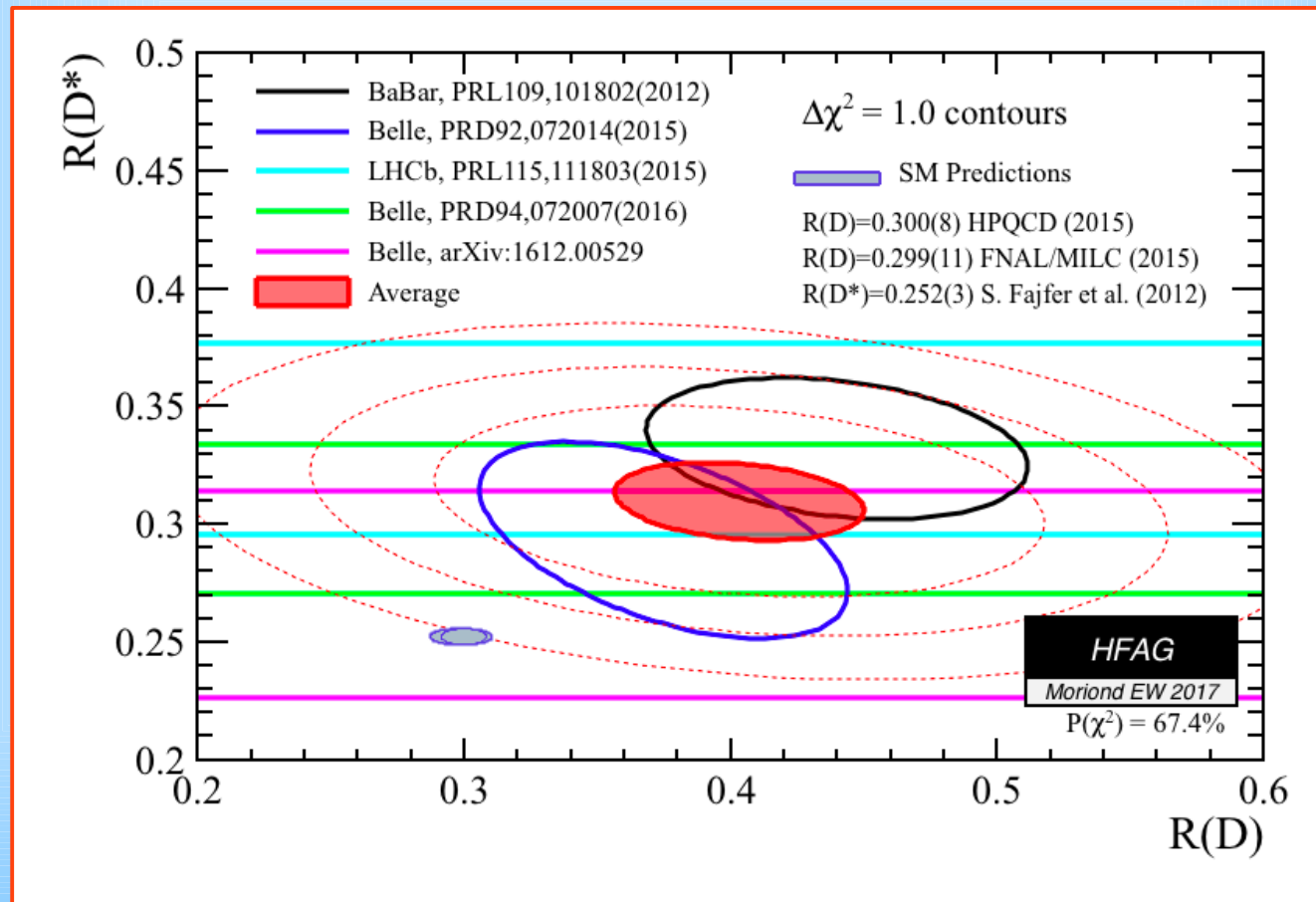
There are long-standing discrepancies in $b \rightarrow c$ transitions as well.

$$R(D^{(*)}) = \frac{BR(B \rightarrow D^{(*)} \tau \nu)}{BR(B \rightarrow D^{(*)} \ell \nu)} \quad (\text{with } \ell = e, \mu)$$

$b \rightarrow c$ data

There are long-standing discrepancies in $b \rightarrow c$ transitions as well.

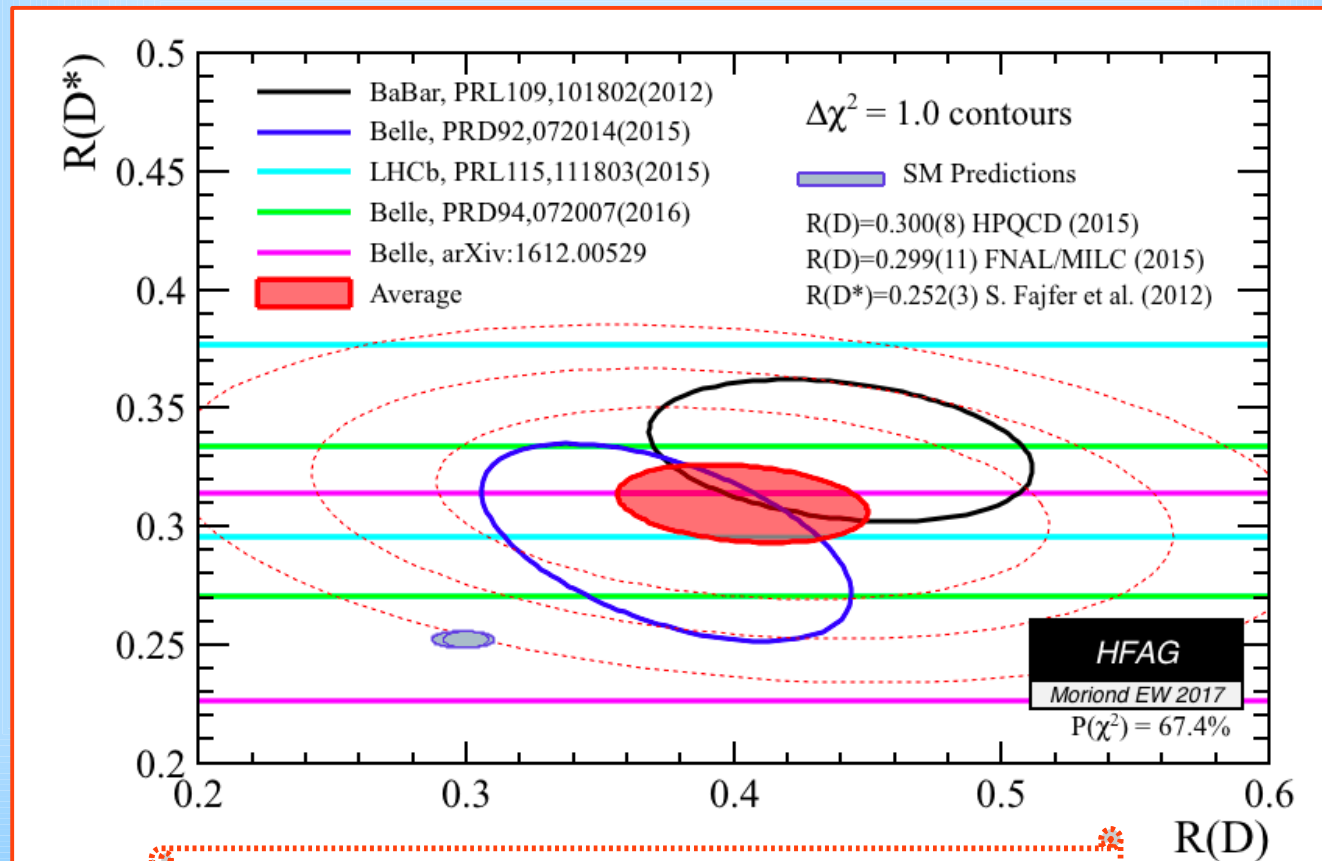
$$R(D^{(*)}) = \frac{BR(B \rightarrow D^{(*)} \tau \nu)}{BR(B \rightarrow D^{(*)} \ell \nu)} \quad (\text{with } \ell = e, \mu)$$



$b \rightarrow c$ data

There are long-standing discrepancies in $b \rightarrow c$ transitions as well.

$$R(D^{(*)}) = \frac{BR(B \rightarrow D^{(*)} \tau \nu)}{BR(B \rightarrow D^{(*)} \ell \nu)} \quad (\text{with } \ell = e, \mu)$$



Simultaneous fit to $R(D)$ & $R(D^*)$ about 4σ away from SM

Wrap-up

- R_K and R_{K^*} hint at Lepton Universality Violation (LUV), the effect being in muons, rather than electrons

Wrap-up

- R_K and R_{K^*} hint at Lepton Universality Violation (LUV), the effect being in muons, rather than electrons
- Also $R(D^{(*)})$ points to LUV. But can we really trust final-state taus?

Wrap-up

- R_K and R_{K^*} hint at Lepton Universality Violation (LUV), the effect being in muons, rather than electrons
- Also $R(D^{(*)})$ points to LUV. But can we really trust final-state taus?
- R_K and R_{K^*} significance fairly low.
Yet interesting that all $b \rightarrow s \mu\mu$ modes go in a consistent direction

Wrap-up

- R_K and R_{K^*} hint at Lepton Universality Violation (LUV), the effect being in muons, rather than electrons
- Also $R(D^{(*)})$ points to LUV. But can we really trust final-state taus?
- R_K and R_{K^*} significance fairly low.
Yet interesting that all $b \rightarrow s \mu\mu$ modes go in a consistent direction
- Focusing for the moment on the $b \rightarrow s$ discrepancies
 - **Q1:** Can we (easily) make theoretical sense of data?

Wrap-up

- R_K and R_{K^*} hint at Lepton Universality Violation (LUV), the effect being in muons, rather than electrons
- Also $R(D^{(*)})$ points to LUV. But can we really trust final-state taus?
- R_K and R_{K^*} significance fairly low.
Yet interesting that all $b \rightarrow s \mu\mu$ modes go in a consistent direction
- Focusing for the moment on the $b \rightarrow s$ discrepancies
 - **Q1:** Can we (easily) make theoretical sense of data?
 - **Q2:** What are the most immediate signatures to expect ?

Concerning Q1: can we easily make theoretical sense of these data?

- *Yes we can. Consider the following Hamiltonian*

$$H_{\text{SM+NP}}(\bar{b} \rightarrow \bar{s} \mu \mu) = -\frac{4G_F}{\sqrt{2}} V_{tb}^* V_{ts} \frac{\alpha_{\text{em}}}{4\pi} \left[\bar{b}_L \gamma^\lambda s_L \cdot \left(C_9^{(\mu)} \bar{\mu} \gamma_\lambda \mu + C_{10}^{(\mu)} \bar{\mu} \gamma_\lambda \gamma_5 \mu \right) \right]$$

Concerning Q1: can we easily make theoretical sense of these data?

- Yes we can. Consider the following Hamiltonian

About equal size & opposite sign
in the SM (at the m_b scale)

$$H_{\text{SM+NP}}(\bar{b} \rightarrow \bar{s} \mu \mu) = -\frac{4G_F}{\sqrt{2}} V_{tb}^* V_{ts} \frac{\alpha_{\text{em}}}{4\pi} \left[\bar{b}_L \gamma^\lambda s_L \cdot \left(C_9^{(\mu)} \bar{\mu} \gamma_\lambda \mu + C_{10}^{(\mu)} \bar{\mu} \gamma_\lambda \gamma_5 \mu \right) \right]$$

Concerning Q1: can we easily make theoretical sense of these data?

- Yes we can. Consider the following Hamiltonian

$$H_{\text{SM+NP}}(\bar{b} \rightarrow \bar{s} \mu \mu) = -\frac{4G_F}{\sqrt{2}} V_{tb}^* V_{ts} \frac{\alpha_{\text{em}}}{4\pi} \left[\bar{b}_L \gamma^\lambda s_L \cdot \left(C_9^{(\mu)} \bar{\mu} \gamma_\lambda \mu + C_{10}^{(\mu)} \bar{\mu} \gamma_\lambda \gamma_5 \mu \right) \right]$$

About equal size & opposite sign
in the SM (at the m_b scale)

- Advocating the same $(V - A) \times (V - A)$ structure also for the corrections to $C_{9,10}^{\text{SM}}$ (in the $\mu\mu$ -channel only!) would account for:
 - R_K and R_{K^*} lower than 1
 - $b \rightarrow s \mu\mu$ BR data below predictions
 - the P_5' anomaly in $B \rightarrow K^* \mu\mu$

Concerning Q1: can we easily make theoretical sense of these data?

- Yes we can. Consider the following Hamiltonian

$$H_{\text{SM+NP}}(\bar{b} \rightarrow \bar{s} \mu \mu) = -\frac{4G_F}{\sqrt{2}} V_{tb}^* V_{ts} \frac{\alpha_{\text{em}}}{4\pi} \left[\bar{b}_L \gamma^\lambda s_L \cdot \left(C_9^{(\mu)} \bar{\mu} \gamma_\lambda \mu + C_{10}^{(\mu)} \bar{\mu} \gamma_\lambda \gamma_5 \mu \right) \right]$$

About equal size & opposite sign
in the SM (at the m_b scale)

- Advocating the same $(V - A) \times (V - A)$ structure also for the corrections to $C_{9,10}^{\text{SM}}$ (in the $\mu\mu$ -channel only!) would account for:
 - R_K and R_{K^*} lower than 1
 - $b \rightarrow s \mu\mu$ BR data below predictions
 - the P_5' anomaly in $B \rightarrow K^* \mu\mu$
- A fully quantitative test requires a global fit.

new physics contributions to the Wilson coefficients. We find that the by far largest decrease in the χ^2 can be obtained either by a negative new physics contribution to C_9 (with $C_9^{\text{NP}} \sim -30\% \times C_9^{\text{SM}}$), or by new physics in the $SU(2)_L$ invariant direction $C_9^{\text{NP}} = -C_{10}^{\text{NP}}$, (with $C_9^{\text{NP}} \sim -12\% \times C_9^{\text{SM}}$). A positive NP contribution to C_{10} alone would also improve the fit, although to a lesser extent. [Altmannshofer, Straub, EPJC '15]

For analogous conclusions, see also [Ghosh, Nardecchia, Renner, JHEP '14]

Model example

- *As we saw before, all $b \rightarrow s$ data are explained at one stroke if:*

- $C_9^{(\ell)} \approx -C_{10}^{(\ell)} \quad (V-A \text{ structure})$
- $|C_{9,\text{NP}}^{(\mu)}| \gg |C_{9,\text{NP}}^{(e)}| \quad (LUV)$

Model example

- As we saw before, all $b \rightarrow s$ data are explained at one stroke if:

- $C_9^{(\ell)} \approx -C_{10}^{(\ell)} \quad (V-A \text{ structure})$
- $|C_{9,\text{NP}}^{(\mu)}| \gg |C_{9,\text{NP}}^{(e)}| \quad (\text{LUV})$

- This pattern can be generated from a purely 3rd-generation interaction of the kind

$$H_{\text{NP}} = G \bar{b}'_L \gamma^\lambda b'_L \bar{\tau}'_L \gamma_\lambda \tau'_L$$

$$\text{with } G = 1/\Lambda_{\text{NP}}^2 \ll G_F$$

Glashow et al., 2015

Model example

- As we saw before, all $b \rightarrow s$ data are explained at one stroke if:

- $C_9^{(\ell)} \approx -C_{10}^{(\ell)} \quad (V-A \text{ structure})$
- $|C_{9,\text{NP}}^{(\mu)}| \gg |C_{9,\text{NP}}^{(e)}| \quad (\text{LUV})$

- This pattern can be generated from a purely 3rd-generation interaction of the kind

$$H_{\text{NP}} = G \bar{b}'_L \gamma^\lambda b'_L \bar{\tau}'_L \gamma_\lambda \tau'_L$$

with $G = 1/\Lambda_{\text{NP}}^2 \ll G_F$

Glashow et al., 2015

expected e.g. in
partial-compositeness
frameworks

Model example

- As we saw before, all $b \rightarrow s$ data are explained at one stroke if:

- $C_9^{(\ell)} \approx -C_{10}^{(\ell)} \quad (V-A \text{ structure})$
- $|C_{9,\text{NP}}^{(\mu)}| \gg |C_{9,\text{NP}}^{(e)}| \quad (\text{LUV})$

- This pattern can be generated from a purely 3rd-generation interaction of the kind

$$H_{\text{NP}} = G \bar{b}'_L \gamma^\lambda b'_L \bar{\tau}'_L \gamma_\lambda \tau'_L$$

$$\text{with } G = 1/\Lambda_{\text{NP}}^2 \ll G_F$$

Glashow et al., 2015

expected e.g. in
partial-compositeness
frameworks

- Note:** primed fields
 - Fields are in the “gauge” basis (= primed)

Model example

- As we saw before, all $b \rightarrow s$ data are explained at one stroke if:

- $C_9^{(\ell)} \approx -C_{10}^{(\ell)} \quad (V-A \text{ structure})$
- $|C_{9,\text{NP}}^{(\mu)}| \gg |C_{9,\text{NP}}^{(e)}| \quad (\text{LUV})$

- This pattern can be generated from a purely 3rd-generation interaction of the kind

$$H_{\text{NP}} = G \bar{b}'_L \gamma^\lambda b'_L \bar{\tau}'_L \gamma_\lambda \tau'_L$$

$$\text{with } G = 1/\Lambda_{\text{NP}}^2 \ll G_F$$

Glashow et al., 2015

expected e.g. in
partial-compositeness
frameworks

- Note:** primed fields

- Fields are in the “gauge” basis (= primed)
- They need to be rotated to the mass eigenbasis



mass
basis

$$b'_L \equiv (d'_L)_3 = (U_L^d)_{3i} (d_L)_i$$

$$\tau'_L \equiv (\ell'_L)_3 = (U_L^\ell)_{3i} (\ell_L)_i$$

Model example

- As we saw before, all $b \rightarrow s$ data are explained at one stroke if:

- $C_9^{(\ell)} \approx -C_{10}^{(\ell)} \quad (V-A \text{ structure})$
- $|C_{9,\text{NP}}^{(u)}| \gg |C_{9,\text{NP}}^{(e)}| \quad (\text{LUV})$

- This pattern can be generated from a purely 3rd-generation interaction of the kind

$$H_{\text{NP}} = G \bar{b}'_L \gamma^\lambda b'_L \bar{\tau}'_L \gamma_\lambda \tau'_L$$

$$\text{with } G = 1/\Lambda_{\text{NP}}^2 \ll G_F$$

Glashow et al., 2015

expected e.g. in
partial-compositeness
frameworks

- Note:** primed fields

- Fields are in the “gauge” basis (= primed)
- They need to be rotated to the mass eigenbasis
- This rotation induces LUV and LFV effects

mass
basis

$$b'_L \equiv (d'_L)_3 = (U_L^d)_{3i} (d_L)_i$$

$$\tau'_L \equiv (\ell'_L)_3 = (U_L^\ell)_{3i} (\ell_L)_i$$

LFV in B decays

As mentioned: if R_K is signaling BSM LUV, then, in general, expect BSM LFV as well

LFV in B decays

As mentioned: if R_K is signaling BSM LUV, then, in general, expect BSM LFV as well

Actually, the expected ballpark of LFV effects can be predicted from $BR(B \rightarrow K \mu\mu)$ and the R_K deviation alone [Glashow et al., 2015]

LFV in B decays

As mentioned: if R_K is signaling BSM LUV, then, in general, expect BSM LFV as well

Actually, the expected ballpark of LFV effects can be predicted from $BR(B \rightarrow K \mu \mu)$ and the R_K deviation alone [Glashow et al., 2015]

$$\boxed{\checkmark} \quad \frac{BR(B^+ \rightarrow K^+ \mu e)}{BR(B^+ \rightarrow K^+ \mu \mu)} = \frac{|\delta C_{10}|^2}{|C_{10}^{\text{SM}} + \delta C_{10}|^2} \cdot \frac{|(U_L^\ell)_{31}|^2}{|(U_L^\ell)_{32}|^2} \cdot 2$$

LFV in B decays

As mentioned: if R_K is signaling BSM LUV, then, in general, expect BSM LFV as well

Actually, the expected ballpark of LFV effects can be predicted from $BR(B \rightarrow K \mu \mu)$ and the R_K deviation alone [Glashow et al., 2015]

$$\checkmark \quad \frac{BR(B^+ \rightarrow K^+ \mu e)}{BR(B^+ \rightarrow K^+ \mu \mu)} = \frac{|\delta C_{10}|^2}{|C_{10}^{\text{SM}} + \delta C_{10}|^2} \cdot \frac{|(U_L^\ell)_{31}|^2}{|(U_L^\ell)_{32}|^2} \cdot 2$$

= 0.159²
according to R_K

LFV in B decays

As mentioned: if R_K is signaling BSM LUV, then, in general, expect BSM LFV as well

Actually, the expected ballpark of LFV effects can be predicted from $BR(B \rightarrow K \mu \mu)$ and the R_K deviation alone [Glashow et al., 2015]

$$\checkmark \quad \frac{BR(B^+ \rightarrow K^+ \mu e)}{BR(B^+ \rightarrow K^+ \mu \mu)} = \frac{|\delta C_{10}|^2}{|C_{10}^{\text{SM}} + \delta C_{10}|^2} \cdot \frac{|(U_L^\ell)_{31}|^2}{|(U_L^\ell)_{32}|^2} \cdot 2$$

$= 0.159^2$
 according to R_K

$\mu^+ e^- \text{ \& } \mu^- e^+$
 modes

LFV in B decays


As mentioned: if R_K is signaling BSM LUV, then, in general, expect BSM LFV as well

Actually, the expected ballpark of LFV effects can be predicted from $BR(B \rightarrow K \mu \mu)$ and the R_K deviation alone [Glashow et al., 2015]

$$\checkmark \quad \frac{BR(B^+ \rightarrow K^+ \mu e)}{BR(B^+ \rightarrow K^+ \mu \mu)} = \frac{|\delta C_{10}|^2}{|C_{10}^{\text{SM}} + \delta C_{10}|^2} \cdot \frac{|(U_L^\ell)_{31}|^2}{|(U_L^\ell)_{32}|^2} \cdot 2$$

= 0.159²
according to R_K

**$\mu^+ e^-$ & $\mu^- e^+$
modes**



$$BR(B^+ \rightarrow K^+ \mu e) < 2.2 \times 10^{-8} \cdot \frac{|(U_L^\ell)_{31}|^2}{|(U_L^\ell)_{32}|^2}$$

The current $BR(B^+ \rightarrow K^+ \mu e)$ limit yields the weak bound

$$|(U_L^\ell)_{31}| / |(U_L^\ell)_{32}| < 3.7$$

LFV in B decays


As mentioned: if R_K is signaling BSM LUV, then, in general, expect BSM LFV as well

Actually, the expected ballpark of LFV effects can be predicted from $BR(B \rightarrow K \mu \mu)$ and the R_K deviation alone [Glashow et al., 2015]

$$\checkmark \quad \frac{BR(B^+ \rightarrow K^+ \mu e)}{BR(B^+ \rightarrow K^+ \mu \mu)} = \frac{|\delta C_{10}|^2}{|C_{10}^{\text{SM}} + \delta C_{10}|^2} \cdot \frac{|(U_L^\ell)_{31}|^2}{|(U_L^\ell)_{32}|^2} \cdot 2$$

**= 0.159²
according to R_K**

**$\mu^+ e^-$ & $\mu^- e^+$
modes**

 $BR(B^+ \rightarrow K^+ \mu e) < 2.2 \times 10^{-8} \cdot \frac{|(U_L^\ell)_{31}|^2}{|(U_L^\ell)_{32}|^2}$

The current $BR(B^+ \rightarrow K^+ \mu e)$ limit yields the weak bound

$$|(U_L^\ell)_{31}|/|(U_L^\ell)_{32}| < 3.7$$

$$\checkmark \quad BR(B^+ \rightarrow K^+ \mu \tau) \quad \text{would be even more promising, as it scales with } |(U_L^\ell)_{33}|/|(U_L^\ell)_{32}|^2$$

LFV in B decays


As mentioned: if R_K is signaling BSM LUV, then, in general, expect BSM LFV as well

Actually, the expected ballpark of LFV effects can be predicted from $BR(B \rightarrow K \mu \mu)$ and the R_K deviation alone [Glashow et al., 2015]

$$\checkmark \quad \frac{BR(B^+ \rightarrow K^+ \mu e)}{BR(B^+ \rightarrow K^+ \mu \mu)} = \frac{|\delta C_{10}|^2}{|C_{10}^{\text{SM}} + \delta C_{10}|^2} \cdot \frac{|(U_L^\ell)_{31}|^2}{|(U_L^\ell)_{32}|^2} \cdot 2$$

= 0.159²
according to R_K

**$\mu^+ e^-$ & $\mu^- e^+$
modes**

 $BR(B^+ \rightarrow K^+ \mu e) < 2.2 \times 10^{-8} \cdot \frac{|(U_L^\ell)_{31}|^2}{|(U_L^\ell)_{32}|^2}$

The current $BR(B^+ \rightarrow K^+ \mu e)$ limit yields the weak bound

$$|(U_L^\ell)_{31}|/|(U_L^\ell)_{32}| < 3.7$$

\checkmark $BR(B^+ \rightarrow K^+ \mu \tau)$ would be even more promising, as it scales with $|(U_L^\ell)_{33}|/|(U_L^\ell)_{32}|^2$

\checkmark An analogous argument holds for purely leptonic modes

Making the interaction G_{SM} - invariant

- *Being defined above the EWSB scale, our assumed operator*

$$\bar{b}'_L \gamma^\lambda b'_L \bar{\tau}'_L \gamma_\lambda \tau'_L$$

*must actually be made invariant
under $SU(3)_c \times SU(2)_L \times U(1)_Y$*


See:
Bhattacharya, Datta, London,
Shivashankara, PLB 15

Making the interaction G_{SM} - invariant

- Being defined above the EWSB scale, our assumed operator

$$\bar{b}'_L \gamma^\lambda b'_L \bar{\tau}'_L \gamma_\lambda \tau'_L$$

must actually be made invariant
under $SU(3)_c \times SU(2)_L \times U(1)_Y$

$SU(2)_L$

inv.

- $\bar{Q}'_L \gamma^\lambda Q'_L \bar{L}'_L \gamma_\lambda L'_L$ [neutral-current int's only]
- $\bar{Q}'^i_L \gamma^\lambda Q'^j_L \bar{L}'^j_L \gamma_\lambda L'^i_L$ [also charged-current int's]


See:
Bhattacharya, Datta, London,
Shivashankara, PLB 15

Making the interaction G_{SM} - invariant

- Being defined above the EWSB scale,
our assumed operator

$$\bar{b}'_L \gamma^\lambda b'_L \bar{\tau}'_L \gamma_\lambda \tau'_L$$

must actually be made invariant
under $SU(3)_c \times SU(2)_L \times U(1)_Y$

$SU(2)_L$

inv.

$$\left\{ \begin{array}{ll} \bullet \bar{Q}'_L \gamma^\lambda Q'_L \bar{L}'_L \gamma_\lambda L'_L & \text{[neutral-current int's only]} \\ \bullet \bar{Q}'^i_L \gamma^\lambda Q'^j_L \bar{L}'^j_L \gamma_\lambda L'^i_L & \text{[also charged-current int's]} \end{array} \right.$$

- Thus, the generated structures are all of:

$$\begin{array}{ll} t't'v'_\tau v'_\tau, & b'b'v'_\tau v'_\tau, \\ t't'\tau'\tau', & b'b'\tau'\tau', \end{array}$$

Making the interaction G_{SM} - invariant

- Being defined above the EWSB scale,
our assumed operator

$$\bar{b}'_L \gamma^\lambda b'_L \bar{\tau}'_L \gamma_\lambda \tau'_L$$

must actually be made invariant
under $SU(3)_c \times SU(2)_L \times U(1)_Y$

$SU(2)_L$
inv.

$$\left\{ \begin{array}{l} \bullet \bar{Q}'_L \gamma^\lambda Q'_L \bar{L}'_L \gamma_\lambda L'_L \quad [\text{neutral-current int's only}] \\ \bullet \bar{Q}'^i_L \gamma^\lambda Q'^j_L \bar{L}'^j_L \gamma_\lambda L'^i_L \quad [\text{also charged-current int's}] \end{array} \right.$$

- Thus, the generated structures are all of:

$$\begin{array}{ll} t't'v'_\tau v'_\tau, & b'b'v'_\tau v'_\tau, \\ t't'\tau'\tau', & b'b'\tau'\tau', \end{array} \quad \text{and}$$

$$t'b'\tau'v'_\tau$$

Making the interaction G_{SM} - invariant

- Being defined above the EWSB scale,
our assumed operator

$$\bar{b}'_L \gamma^\lambda b'_L \bar{\tau}'_L \gamma_\lambda \tau'_L$$

must actually be made invariant
under $SU(3)_c \times SU(2)_L \times U(1)_Y$

$SU(2)_L$
inv.

$$\left\{ \begin{array}{ll} \bullet \bar{Q}'_L \gamma^\lambda Q'_L \bar{L}'_L \gamma_\lambda L'_L & \text{[neutral-current int's only]} \\ \bullet \bar{Q}'^i_L \gamma^\lambda Q'^j_L \bar{L}'^j_L \gamma_\lambda L'^i_L & \text{[also charged-current int's]} \end{array} \right.$$

- Thus, the generated structures are all of:

$$\begin{array}{ll} t't'v'_\tau v'_\tau, & b'b'v'_\tau v'_\tau, \\ t't'\tau'\tau', & b'b'\tau'\tau', \end{array} \quad \text{and}$$

$$t'b'\tau'v'_\tau$$



After rotation to the mass basis (unprimed), the last structure contributes to $\Gamma(b \rightarrow c \tau \bar{\nu})$



i.e. it can explain deviations on $R(D^{(*)})$

Making the interaction G_{SM} - invariant

- Being defined above the EWSB scale, our assumed operator

$$\bar{b}'_L \gamma^\lambda b'_L \bar{\tau}'_L \gamma_\lambda \tau'_L$$

must actually be made invariant under $SU(3)_c \times SU(2)_L \times U(1)_Y$

$SU(2)_L$
inv.

$$\left\{ \begin{array}{ll} \bullet \bar{Q}'_L \gamma^\lambda Q'_L \bar{L}'_L \gamma_\lambda L'_L & \text{[neutral-current int's only]} \\ \bullet \bar{Q}'^i_L \gamma^\lambda Q'^j_L \bar{L}'^j_L \gamma_\lambda L'^i_L & \text{[also charged-current int's]} \end{array} \right.$$

- Thus, the generated structures are all of:

$$\begin{array}{ll} t't'v'_\tau v'_\tau, & b'b'v'_\tau v'_\tau, \\ t't'\tau'\tau', & b'b'\tau'\tau', \end{array} \quad \text{and}$$

$$t'b'\tau'v'_\tau$$



After rotation to the mass basis (unprimed), the last structure contributes to $\Gamma(b \rightarrow c \tau \bar{\nu})$



i.e. it can explain deviations on $R(D^{(*)})$

- But this coin has a flip side.

Through RGE running, one gets also LFU-breaking effects in $\tau \rightarrow \ell \nu \bar{\nu}$ (tested at per mil accuracy)

See:
Bhattacharya, Datta, London,
Shivashankara, PLB 15

Feruglio, Paradisi, Pattori, 2016

Making the interaction G_{SM} - invariant

- Being defined above the EWSB scale, our assumed operator

$$\bar{b}'_L \gamma^\lambda b'_L \bar{\tau}'_L \gamma_\lambda \tau'_L$$

must actually be made invariant under $SU(3)_c \times SU(2)_L \times U(1)_Y$

$SU(2)_L$
inv.

$$\left\{ \begin{array}{l} \bullet \bar{Q}'_L \gamma^\lambda Q'_L \bar{L}'_L \gamma_\lambda L'_L \quad [\text{neutral-current int's only}] \\ \bullet \bar{Q}'^i_L \gamma^\lambda Q'^j_L \bar{L}'^j_L \gamma_\lambda L'^i_L \quad [\text{also charged-current int's}] \end{array} \right.$$

- Thus, the generated structures are all of:

$$\begin{array}{ll} t't'v'_\tau v'_\tau, & b'b'v'_\tau v'_\tau, \\ t't'\tau'\tau', & b'b'\tau'\tau', \end{array} \quad \text{and}$$

$$t'b'\tau'v'_\tau$$



After rotation to the mass basis (unprimed), the last structure contributes to $\Gamma(b \rightarrow c \tau \bar{\nu})$



i.e. it can explain deviations on $R(D^{(*)})$

- But this coin has a flip side.

Through RGE running, one gets also LFU-breaking effects in $\tau \rightarrow \ell \nu \nu$
(tested at per mil accuracy)

Such effects “strongly disfavour an explanation of the $R(D^{(*)})$ anomaly model-independently”

See:
Bhattacharya, Datta, London,
Shivashankara, PLB 15

Feruglio, Paradisi, Pattori, 2016

Further tests

- **Measure more LUV ratios:** $R_{K^*}, R_\phi, R_{Xs}, R_{K_0(1430)}, R_{f_0}$

Further tests

- **Measure more LUV ratios:** $R_{K^*}, R_\phi, R_{X_s}, R_{K_0(1430)}, R_{f_0}$

Interesting test: define $X_H \equiv \frac{R_H}{R_K}$, with $H = K^*, \phi, X_s, K_0(1430), f_0$

Deviations from unity in the double ratios X_H can only come from RH currents

Hiller, Schmaltz, JHEP 2015

Further tests

- **Measure more LUV ratios:** $R_{K^*}, R_\phi, R_{X_s}, R_{K_0(1430)}, R_{f_0}$

Hiller, Schmaltz, JHEP 2015

Interesting test: define $X_H \equiv \frac{R_H}{R_K}$, with $H = K^*, \phi, X_s, K_0(1430), f_0$

Deviations from unity in the double ratios X_H can only come from RH currents

- We now know the first of these double ratios, X_{K^*} (for $q^2 \in [1, 6] \text{ GeV}^2$)

$$X_{K^*} = 0.92 \cdot (1 \pm O(20\%))$$

Further tests

- **Measure more LUV ratios:** $R_{K^*}, R_\phi, R_{X_s}, R_{K_0(1430)}, R_{f_0}$

Hiller, Schmaltz, JHEP 2015

Interesting test: define $X_H \equiv \frac{R_H}{R_K}$, with $H = K^*, \phi, X_s, K_0(1430), f_0$

Deviations from unity in the double ratios X_H can only come from RH currents

- We now know the first of these double ratios, X_{K^*} (for $q^2 \in [1, 6] \text{ GeV}^2$)

$$X_{K^*} = 0.92 \cdot (1 \pm O(20\%)) \quad \Rightarrow \quad \text{within errors, it is compatible with unity}$$

Further tests

- **Measure more LUV ratios:** $R_{K^*}, R_\phi, R_{X_s}, R_{K_0(1430)}, R_{f_0}$

Hiller, Schmaltz, JHEP 2015

Interesting test: define $X_H \equiv \frac{R_H}{R_K}$, with $H = K^*, \phi, X_s, K_0(1430), f_0$

Deviations from unity in the double ratios X_H can only come from RH currents

- We now know the first of these double ratios, X_{K^*} (for $q^2 \in [1, 6] \text{ GeV}^2$)

$$X_{K^*} = 0.92 \cdot (1 \pm O(20\%)) \quad \Rightarrow \quad \text{within errors, it is compatible with unity}$$

In general we have [Hiller, Schmaltz, JHEP 2015]

$$X_{K^*} \simeq 1 - 0.41 \operatorname{Re} \left(C_9^{\prime \mu} - C_{10}^{\prime \mu} - \{\mu \rightarrow e\} \right)$$

Remember

$$O_9^{\prime \ell} = (\bar{s} \gamma^\mu P_R b) (\bar{\ell} \gamma_\mu \ell)$$

$$O_{10}^{\prime \ell} = (\bar{s} \gamma^\mu P_R b) (\bar{\ell} \gamma_\mu \gamma_5 \ell)$$

Further tests

- **Extract LD effects from data**

Recently, LHCb measured $BR(B^+ \rightarrow K^+ \mu\mu)$ including an accurate parameterization of the LD component in the $c\bar{c}$ region

LHCb, 1612.06764

Similar approach as
Lyon, Zwicky, '14

Further tests

- **Extract LD effects from data**

Recently, LHCb measured $BR(B^+ \rightarrow K^+ \mu\mu)$ including an accurate parameterization of the LD component in the $c\bar{c}$ region

Idea: *Sizable LD contributions far from the resonance region could explain away tensions*

LHCb, 1612.06764

Similar approach as
Lyon, Zwicky, '14

Further tests

- **Extract LD effects from data**

Recently, LHCb measured $BR(B^+ \rightarrow K^+ \mu\mu)$ including an accurate parameterization of the LD component in the $c\bar{c}$ region

Idea: *Sizable LD contributions far from the resonance region could explain away tensions*

Method: *Measure $m_{\mu\mu}$ spectrum, including the $c\bar{c}$ resonances as a sum of BW, and fit 'em all*

LHCb, 1612.06764

Similar approach as
Lyon, Zwicky, '14

Further tests

- **Extract LD effects from data**

Recently, LHCb measured $BR(B^+ \rightarrow K^+ \mu\mu)$ including an accurate parameterization of the LD component in the $c\bar{c}$ region

Idea: Sizable LD contributions far from the resonance region could explain away tensions

Method: Measure $m_{\mu\mu}$ spectrum, including the $c\bar{c}$ resonances as a sum of BW, and fit 'em all

Result: BR compatible with previous measurements, and (again) smaller than SM

LHCb, 1612.06764

Similar approach as
Lyon, Zwicky, '14

Further tests

- *Measure new observables sensitive to C_9 and C_{10}*

Further tests

- **Measure new observables sensitive to C_9 and C_{10}**
 - The $B_s \rightarrow \mu\mu \gamma$ decay offers sensitivity to C_7 , C_9 , C_{10} (and its total BR is 10^{-8})
Its direct measurement (= with photon detection) is veeery challenging at hadron colliders

Further tests

- **Measure new observables sensitive to C_9 and C_{10}**
 - The $B_s \rightarrow \mu\mu \gamma$ decay offers sensitivity to C_7 , C_9 , C_{10} (and its total BR is 10^{-8})
Its direct measurement (= with photon detection) is veeery challenging at hadron colliders
- Extract $B_s \rightarrow \mu\mu \gamma$ from $B_s \rightarrow \mu\mu$ event sample, by enlarging $m_{\mu\mu}$ window downwards

Dettori, DG, Reboud, '16

Further tests

- **Measure new observables sensitive to C_9 and C_{10}**
 - The $B_s \rightarrow \mu\mu \gamma$ decay offers sensitivity to C_7 , C_9 , C_{10} (and its total BR is 10^{-8})
Its direct measurement (= with photon detection) is veeery challenging at hadron colliders
- Extract $B_s \rightarrow \mu\mu \gamma$ from $B_s \rightarrow \mu\mu$ event sample, by enlarging $m_{\mu\mu}$ window downwards

Dettori, DG, Reboud, '16

Note in fact:

- ISR and FSR components in $B_s \rightarrow \mu\mu \gamma$ can be treated as independent
(relevant in different regions & interference is negligible)

Further tests

- **Measure new observables sensitive to C_9 and C_{10}**
 - The $B_s \rightarrow \mu\mu \gamma$ decay offers sensitivity to C_7 , C_9 , C_{10} (and its total BR is 10^{-8})
Its direct measurement (= with photon detection) is veeery challenging at hadron colliders
- Extract $B_s \rightarrow \mu\mu \gamma$ from $B_s \rightarrow \mu\mu$ event sample, by enlarging $m_{\mu\mu}$ window downwards

Dettori, DG, Reboud, '16

Note in fact:

- ISR and FSR components in $B_s \rightarrow \mu\mu \gamma$ can be treated as independent
(relevant in different regions & interference is negligible)
- The FSR component can be systematically subtracted from data
(the same way it is in $B_s \rightarrow \mu\mu$)

Further tests

- **Measure new observables sensitive to C_9 and C_{10}**
 - The $B_s \rightarrow \mu\mu \gamma$ decay offers sensitivity to C_7 , C_9 , C_{10} (and its total BR is 10^{-8})
Its direct measurement (= with photon detection) is veeery challenging at hadron colliders
- Extract $B_s \rightarrow \mu\mu \gamma$ from $B_s \rightarrow \mu\mu$ event sample, by enlarging $m_{\mu\mu}$ window downwards

Dettori, DG, Reboud, '16

Note in fact:

- ISR and FSR components in $B_s \rightarrow \mu\mu \gamma$ can be treated as independent
(relevant in different regions & interference is negligible)
- The FSR component can be systematically subtracted from data
(the same way it is in $B_s \rightarrow \mu\mu$)
- So this measurement gives access to the ISR spectrum, to be compared with theory
[Melikhov-Nikitin, '04]

But LQCD calculation of $B \rightarrow \gamma$ f.f.'s required


EFT is ok. But models?

*Up to now: Fermi-like interactions involving SM fields only.
Is there any plausible dynamics generating these interactions?*

EFT is ok. But models?

*Up to now: Fermi-like interactions involving SM fields only.
Is there any plausible dynamics generating these interactions?*


First obstacle towards a common explanation of $b \rightarrow s \ell \ell$ and $b \rightarrow c \tau \nu$:

- $B \rightarrow D(*) \tau \nu$: arises at tree level in the SM, and the effect is $O(25\%)$
  tree-level charged mediators?

EFT is ok. But models?

*Up to now: Fermi-like interactions involving SM fields only.
Is there any plausible dynamics generating these interactions?*


First obstacle towards a common explanation of $b \rightarrow s \ell \ell$ and $b \rightarrow c \tau \nu$:

- $B \rightarrow D(*) \tau \nu$: arises at tree level in the SM, and the effect is $O(25\%)$
 tree-level charged mediators?
- $B \rightarrow K(*) \ell \ell$: again 25% effect, but this is a loop effect in the SM

EFT is ok. But models?

*Up to now: Fermi-like interactions involving SM fields only.
Is there any plausible dynamics generating these interactions?*

First obstacle towards a common explanation of $b \rightarrow s \ell \ell$ and $b \rightarrow c \tau \nu$:

- $B \rightarrow D(*) \tau \nu$: arises at tree level in the SM, and the effect is $O(25\%)$
 tree-level charged mediators?
- $B \rightarrow K(*) \ell \ell$: again 25% effect, but this is a loop effect in the SM


Second obstacle

- The needed NP is of the kind $J_{quark} \times J_{lepton}$
Hard to believe that it leaves no traces in $J_{quark} \times J_{quark}$ and $J_{lepton} \times J_{lepton}$ as well


EFT is ok. But models?

*Up to now: Fermi-like interactions involving SM fields only.
Is there any plausible dynamics generating these interactions?*

First obstacle towards a common explanation of $b \rightarrow s \ell \ell$ and $b \rightarrow c \tau \nu$:

- $B \rightarrow D(*) \tau \nu$: arises at tree level in the SM, and the effect is $O(25\%)$
 tree-level charged mediators?
- $B \rightarrow K(*) \ell \ell$: again 25% effect, but this is a loop effect in the SM


Second obstacle

- The needed NP is of the kind $J_{quark} \times J_{lepton}$
Hard to believe that it leaves no traces in $J_{quark} \times J_{quark}$ and $J_{lepton} \times J_{lepton}$ as well
 Strong constraints from B_s -mixing & purely leptonic LFV or LUV decays


EFT is ok. But models?

*Up to now: Fermi-like interactions involving SM fields only.
Is there any plausible dynamics generating these interactions?*

First obstacle towards a common explanation of $b \rightarrow s \ell \ell$ and $b \rightarrow c \tau \nu$:

- $B \rightarrow D(*) \tau \nu$: arises at tree level in the SM, and the effect is $O(25\%)$
 tree-level charged mediators?
- $B \rightarrow K(*) \ell \ell$: again 25% effect, but this is a loop effect in the SM

Second obstacle

- The needed NP is of the kind $J_{quark} \times J_{lepton}$
Hard to believe that it leaves no traces in $J_{quark} \times J_{quark}$ and $J_{lepton} \times J_{lepton}$ as well
 Strong constraints from B_s -mixing & purely leptonic LFV or LUV decays

Third obstacle


Most (all?) model-building possibilities involve:

- new charged (and possibly colored) states


EFT is ok. But models?

*Up to now: Fermi-like interactions involving SM fields only.
Is there any plausible dynamics generating these interactions?*

First obstacle towards a common explanation of $b \rightarrow s \ell \ell$ and $b \rightarrow c \tau \nu$:

- $B \rightarrow D(*) \tau \nu$: arises at tree level in the SM, and the effect is $O(25\%)$
 tree-level charged mediators?
- $B \rightarrow K(*) \ell \ell$: again 25% effect, but this is a loop effect in the SM

Second obstacle

- The needed NP is of the kind $J_{quark} \times J_{lepton}$
Hard to believe that it leaves no traces in $J_{quark} \times J_{quark}$ and $J_{lepton} \times J_{lepton}$ as well
 Strong constraints from B_s -mixing & purely leptonic LFV or LUV decays

Third obstacle


Most (all?) model-building possibilities involve:

- new charged (and possibly colored) states
- with masses in the TeV region and


EFT is ok. But models?

*Up to now: Fermi-like interactions involving SM fields only.
Is there any plausible dynamics generating these interactions?*

First obstacle towards a common explanation of $b \rightarrow s \ell \ell$ and $b \rightarrow c \tau \nu$:

- $B \rightarrow D(*) \tau \nu$: arises at tree level in the SM, and the effect is $O(25\%)$
 tree-level charged mediators?
- $B \rightarrow K(*) \ell \ell$: again 25% effect, but this is a loop effect in the SM

Second obstacle

- The needed NP is of the kind $J_{quark} \times J_{lepton}$
Hard to believe that it leaves no traces in $J_{quark} \times J_{quark}$ and $J_{lepton} \times J_{lepton}$ as well
 Strong constraints from B_s -mixing & purely leptonic LFV or LUV decays

Third obstacle


Most (all?) model-building possibilities involve:

- new charged (and possibly colored) states
- with masses in the TeV region and
- with significant couplings to 3rd gen. SM fermions


EFT is ok. But models?

*Up to now: Fermi-like interactions involving SM fields only.
Is there any plausible dynamics generating these interactions?*

First obstacle towards a common explanation of $b \rightarrow s \ell \ell$ and $b \rightarrow c \tau \nu$:

- $B \rightarrow D(*) \tau \nu$: arises at tree level in the SM, and the effect is $O(25\%)$
 tree-level charged mediators?
- $B \rightarrow K(*) \ell \ell$: again 25% effect, but this is a loop effect in the SM

Second obstacle

- The needed NP is of the kind $J_{quark} \times J_{lepton}$
Hard to believe that it leaves no traces in $J_{quark} \times J_{quark}$ and $J_{lepton} \times J_{lepton}$ as well
 Strong constraints from B_s -mixing & purely leptonic LFV or LUV decays

Third obstacle

Most (all?) model-building possibilities involve:

- new charged (and possibly colored) states
- with masses in the TeV region and
- with significant couplings to 3rd gen. SM fermions


*Constraints from direct searches
(e.g. $\rightarrow \tau \tau$) potentially strong*




EFT is ok. But models?

*Up to now: Fermi-like interactions involving SM fields only.
Is there any plausible dynamics generating these interactions?*

First obstacle towards a common explanation of $b \rightarrow s \ell \ell$ and $b \rightarrow c \tau \nu$:

- $B \rightarrow D(*) \tau \nu$: arises at tree level in the SM, and the effect is $O(25\%)$
 tree-level charged mediators?
- $B \rightarrow K(*) \ell \ell$: again 25% effect, but this is a loop effect in the SM

Second obstacle

- The needed NP is of the kind $J_{quark} \times J_{lepton}$
Hard to believe that it leaves no traces in $J_{quark} \times J_{quark}$ and $J_{lepton} \times J_{lepton}$ as well
 Strong constraints from B_s -mixing & purely leptonic LFV or LUV decays

Third obstacle

Most (all?) model-building possibilities involve:

- new charged (and possibly colored) states
- with masses in the TeV region and
- with significant couplings to 3rd gen. SM fermions

*Constraints from direct searches
(e.g. $\rightarrow \tau \tau$) potentially strong*

And yes they are!

See: [\[Greljo-Isidori-Marzocca\]](#)
[\[Faroughy-Greljo-Kamenik\]](#)

- *The above being said, many attempts towards plausible UV completions able to produce the needed operators have been made*

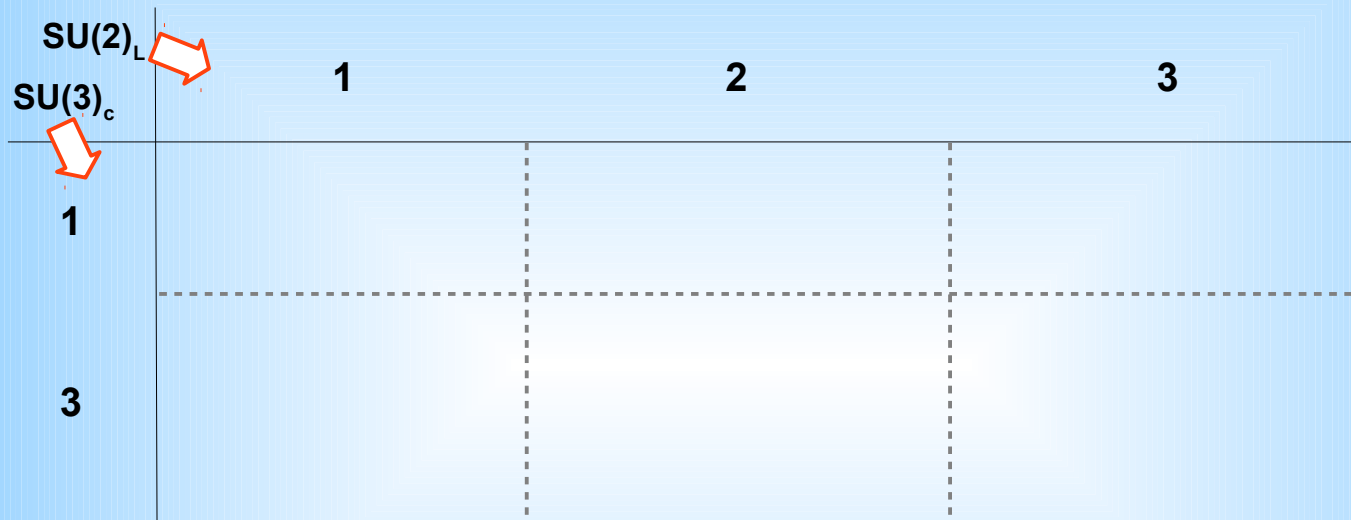
- *The above being said, many attempts towards plausible UV completions able to produce the needed operators have been made*
- *These models involve typically the introduction of:*
 - *a new Lorentz-scalar (S) or -vector (V)*

- *The above being said, many attempts towards plausible UV completions able to produce the needed operators have been made*
- *These models involve typically the introduction of:*
 - *a new Lorentz-scalar (S) or -vector (V)*

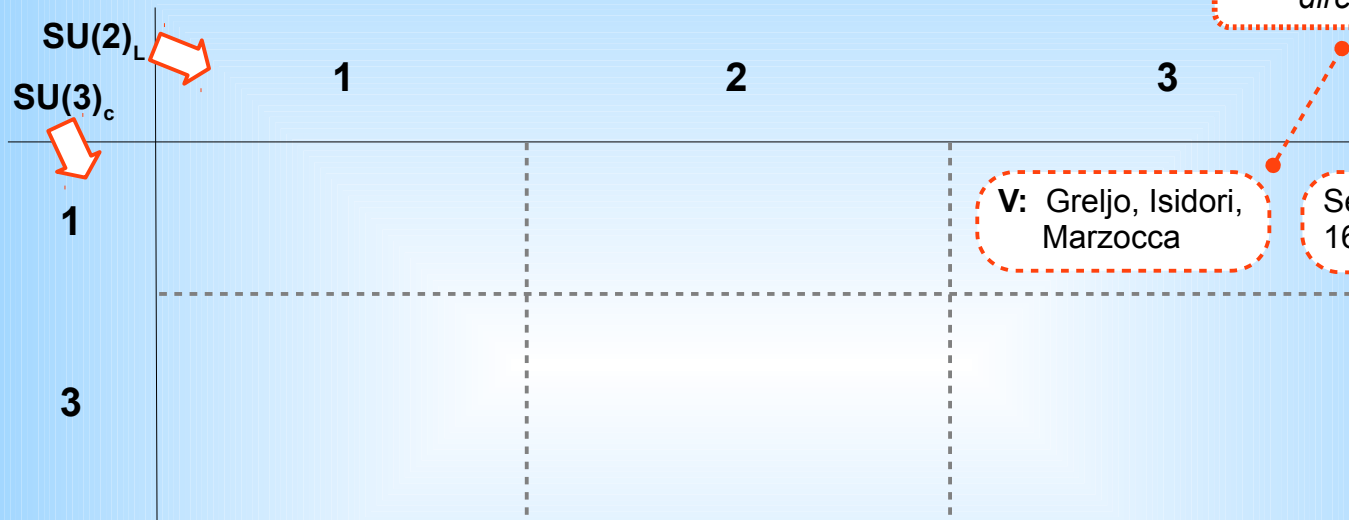
with any of the following transformation properties under the SM gauge group:

- **$SU(3)_c$: 1 or 3** (\rightarrow “leptoquark”)
- **$SU(2)_L$: 1 or 2 or 3**

Recap of model-building attempts
focused on models accounting for R_κ & $R(D^)$*



Recap of model-building attempts
focused on models accounting for R_K & $R(D^)$*



- Strong bounds from $\tau \rightarrow \ell \nu \nu$ and B_s -mixing
- Minimal model ruled out by direct searches $\rightarrow \tau\tau$

V: Greljo, Isidori, Marzocca

See also Boucenna *et al.* 1604.03088

Recap of model-building attempts focused on models accounting for R_K & $R(D^*)$

$SU(2)_L$
 $SU(3)_c$

1

1

2

3

3

S: Bauer-Neubert

V: Greljo, Isidori,
Marzocca

See also Boucenna et al.
1604.03088

- Strong bounds from $\tau \rightarrow \ell \nu \nu$ and B_s -mixing
- Minimal model ruled out by direct searches $\rightarrow \tau\tau$

- Nice & elegant:
 $R(D^*)$ generated at tree level, R_K at loop level
- But not viable: see
Becirevic et al.,
1608.07583

Recap of model-building attempts focused on models accounting for R_K & $R(D^*)$

$SU(2)_L$
 $SU(3)_c$

1

1

2

3

3

S: Bauer-Neubert

V: Barbieri, Isidori,
Pattori, Senia

V: Greljo, Isidori,
Marzocca

S&V: Barbieri, Isidori,
Pattori, Senia

- Strong bounds from $\tau \rightarrow \ell \nu \nu$ and B_s -mixing
- Minimal model ruled out by direct searches $\rightarrow \tau\tau$

See also Boucenna *et al.*
1604.03088

- Nice & elegant:
 $R(D^*)$ generated at tree level, R_K at loop level
- But not viable: see
Becirevic *et al.*,
1608.07583

- Plausible mechanism to generate R_K & $R(D^*)$:
flavor group G_F distinguishing 3rd gen. from
the other two

Recap of model-building attempts focused on models accounting for R_K & $R(D^*)$

$SU(2)_L$
 $SU(3)_c$

1

1

2

3

3

S: Bauer-Neubert

V: Barbieri, Isidori,
Pattori, Senia

V: Greljo, Isidori,
Marzocca

S&V: Barbieri, Isidori,
Pattori, Senia

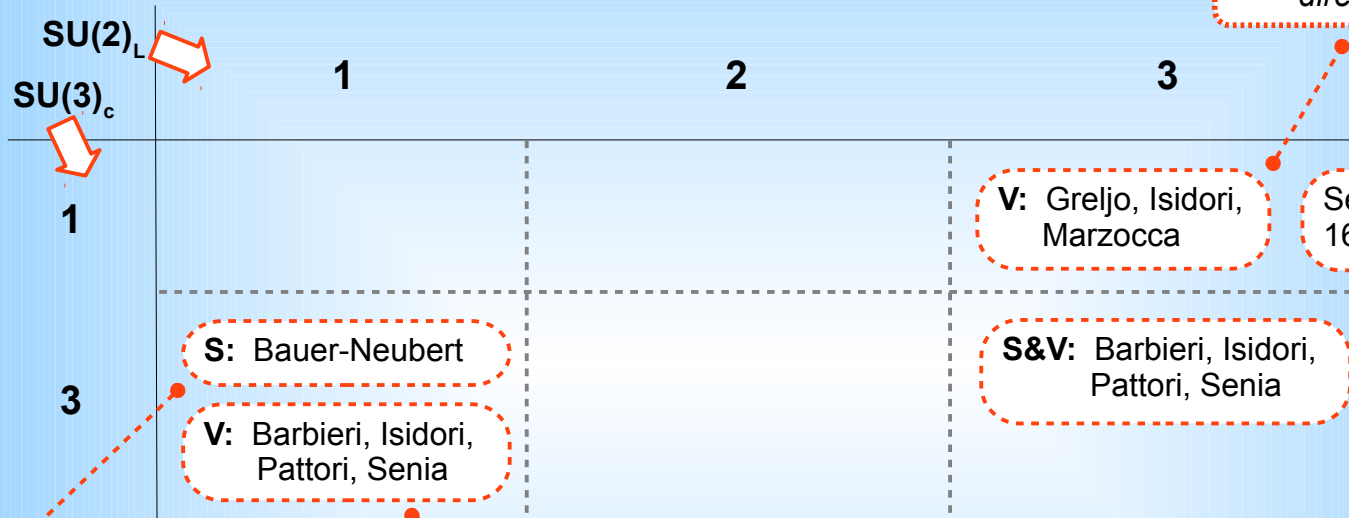
- Strong bounds from $\tau \rightarrow \ell \nu \nu$ and B_s -mixing
- Minimal model ruled out by direct searches $\rightarrow \tau\tau$

See also Boucenna *et al.*
1604.03088

- Nice & elegant:
 $R(D^*)$ generated at tree level, R_K at loop level
- But not viable: see
Becirevic *et al.*,
1608.07583

- Plausible mechanism to generate R_K & $R(D^*)$:
flavor group G_F distinguishing 3rd gen. from
the other two
- Hierarchy between R_K and $R(D^*)$ controlled by
 G_F breaking

Recap of model-building attempts focused on models accounting for R_K & $R(D^*)$



- Strong bounds from $\tau \rightarrow \ell \nu \nu$ and B_s -mixing
- Minimal model ruled out by direct searches $\rightarrow \tau\tau$

V: Greljo, Isidori, Marzocca

See also Boucenna *et al.* 1604.03088

S: Bauer-Neubert

V: Barbieri, Isidori, Pattori, Senia

S&V: Barbieri, Isidori, Pattori, Senia

- Nice & elegant: $R(D^*)$ generated at tree level, R_K at loop level
- But not viable: see Becirevic *et al.*, 1608.07583

- Plausible mechanism to generate R_K & $R(D^*)$: flavor group G_F distinguishing 3rd gen. from the other two
- Hierarchy between R_K and $R(D^*)$ controlled by G_F breaking
- Only vector singlet survives constraints
- Discusses UV-cutoff sensitivity (powerlike)
- ➡ Badly need UV completion (challenging)

Recap of model-building attempts

focused on models accounting for R_K & $R(D^)$*

$SU(2)_L$
 $SU(3)_c$

1

1

2

3

3

S: Bauer-Neubert

V: Barbieri, Isidori, Pattori, Senia

V: Greljo, Isidori, Marzocca

See also Boucenna et al.
1604.03088

S&V: Barbieri, Isidori, Pattori, Senia

V: Fajfer-Kosnic

- Strong bounds from $\tau \rightarrow \ell \nu \nu$ and B_s -mixing
- Minimal model ruled out by direct searches $\rightarrow \tau\tau$

- Nice & elegant:
 $R(D^*)$ generated at tree level, R_K at loop level
- But not viable: see
Becirevic et al.,
1608.07583

- Plausible mechanism to generate R_K & $R(D^*)$:
flavor group G_F distinguishing 3rd gen. from the other two
- Hierarchy between R_K and $R(D^*)$ controlled by G_F breaking
- Only vector singlet survives constraints
- Discusses UV-cutoff sensitivity (powerlike)
- ➡ Badly need UV completion (challenging)

- Similar scenario as
Calibbi, Crivellin, Ota, but
fully general flavor couplings
- Flavor couplings “pragmatically”
fit to data:
 R_K ones are $\sim 10^{-3}$, $R(D^*)$ ones are $O(1)$

Recap of model-building attempts

focused on models accounting for R_K & $R(D^)$*

$SU(2)_L$
 $SU(3)_c$

1

1

2

3

3

S: Bauer-Neubert

V: Barbieri, Isidori, Pattori, Senia

V: Greljo, Isidori, Marzocca


See also Boucenna et al. 1604.03088

S&V: Barbieri, Isidori, Pattori, Senia

V: Fajfer-Kosnic

- Strong bounds from $\tau \rightarrow \ell \nu \nu$ and B_s -mixing
- Minimal model ruled out by direct searches $\rightarrow \tau\tau$

- Nice & elegant: $R(D^*)$ generated at tree level, R_K at loop level
- But not viable: see Becirevic et al., 1608.07583

- Plausible mechanism to generate R_K & $R(D^*)$: flavor group G_F distinguishing 3rd gen. from the other two
- Hierarchy between R_K and $R(D^*)$ controlled by G_F breaking
- Only vector singlet survives constraints
- Discusses UV-cutoff sensitivity (powerlike)
-  Badly need UV completion (challenging)

- Similar scenario as Calibbi, Crivellin, Ota, but fully general flavor couplings
- Flavor couplings “pragmatically” fit to data: R_K ones are $\sim 10^{-3}$, $R(D^*)$ ones are $O(1)$
- How is LQ mass generated? Otherwise theory is non-renorm.

Recap of model-building attempts

focused on models accounting for R_K & $R(D^*)$

$SU(2)_L$
 $SU(3)_c$

1

3

1

2

3

S: Bauer-Neubert

V: Barbieri, Isidori, Pattori, Senia

S: Becirevic, Fajfer, Kosnic, Sumensari

V: Greljo, Isidori, Marzocca

See also Boucenna et al. 1604.03088

S&V: Barbieri, Isidori, Pattori, Senia

V: Fajfer-Kosnic

- Not constrained by Feruglio et al.'s argument
- Prediction: $RK^* > 1$
(V+A) x (V-A) current invoked

- Strong bounds from $\tau \rightarrow \ell \nu \nu$ and B_s -mixing
- Minimal model ruled out by direct searches $\rightarrow \tau\tau$

- Nice & elegant: $R(D^*)$ generated at tree level, R_K at loop level
- But not viable: see Becirevic et al., 1608.07583

- Plausible mechanism to generate R_K & $R(D^*)$: flavor group G_F distinguishing 3rd gen. from the other two
- Hierarchy between R_K and $R(D^*)$ controlled by G_F breaking
- Only vector singlet survives constraints
- Discusses UV-cutoff sensitivity (powerlike)
- ➡ Badly need UV completion (challenging)

- Similar scenario as Calibbi, Crivellin, Ota, but fully general flavor couplings
- Flavor couplings “pragmatically” fit to data: R_K ones are $\sim 10^{-3}$, $R(D^*)$ ones are $O(1)$
- How is LQ mass generated? Otherwise theory is non-renorm.

Recap of model-building attempts

focused on models accounting for R_K & $R(D^)$*

$SU(2)_L$
 $SU(3)_c$

1

3

1

2

3

S: Bauer-Neubert

V: Barbieri, Isidori, Pattori, Senia

S: Becirevic, Fajfer, Kosnic, Sumensari

S: Becirevic, Sumensari

V: Greljo, Isidori, Marzocca

See also Boucenna *et al.* 1604.03088

S&V: Barbieri, Isidori, Pattori, Senia

V: Fajfer-Kosnic

- Not constrained by Feruglio *et al.*'s argument
- RK^* prediction can be made < 1 by forbidding certain couplings

- Strong bounds from $\tau \rightarrow \ell \nu \nu$ and B_s -mixing
- Minimal model ruled out by direct searches $\rightarrow \tau\tau$

- Nice & elegant: $R(D^*)$ generated at tree level, R_K at loop level
- But not viable: see Becirevic *et al.*, 1608.07583

- Plausible mechanism to generate R_K & $R(D^*)$: flavor group G_F distinguishing 3rd gen. from the other two
- Hierarchy between R_K and $R(D^*)$ controlled by G_F breaking
- Only vector singlet survives constraints
- Discusses UV-cutoff sensitivity (powerlike)
- ➡ Badly need UV completion (challenging)

- Similar scenario as Calibbi, Crivellin, Ota, but fully general flavor couplings
- Flavor couplings “pragmatically” fit to data: R_K ones are $\sim 10^{-3}$, $R(D^*)$ ones are $O(1)$
- How is LQ mass generated? Otherwise theory is non-renorm.

Conclusions

- *In flavor physics there are by now several persistent discrepancies with respect to the SM.*

Conclusions

- *In flavor physics there are by now several persistent discrepancies with respect to the SM. Their most convincing aspects are the following:*
 - **Experiments:** *Results are consistent between LHCb and B factories.*
 - **Data:** *Deviations concern two independent sets of data: $b \rightarrow s$ and $b \rightarrow c$ decays.*
 - **Data vs. theory:** *Discrepancies go in a consistent direction.
A BSM explanation is already possible within an EFT approach.*

Conclusions

- *In flavor physics there are by now several persistent discrepancies with respect to the SM. Their most convincing aspects are the following:*
 - **Experiments:** *Results are consistent between LHCb and B factories.*
 - **Data:** *Deviations concern two independent sets of data: $b \rightarrow s$ and $b \rightarrow c$ decays.*
 - **Data vs. theory:** *Discrepancies go in a consistent direction.
A BSM explanation is already possible within an EFT approach.*
- *Early to draw conclusions. But Run II will provide a definite answer*

Conclusions

- *In flavor physics there are by now several persistent discrepancies with respect to the SM. Their most convincing aspects are the following:*
 - **Experiments:** *Results are consistent between LHCb and B factories.*
 - **Data:** *Deviations concern two independent sets of data: $b \rightarrow s$ and $b \rightarrow c$ decays.*
 - **Data vs. theory:** *Discrepancies go in a consistent direction.
A BSM explanation is already possible within an EFT approach.*
- *Early to draw conclusions. But Run II will provide a definite answer*
- **Theory:** *EFT makes sense rather well of data. But hard to find convincing UV dynamics*

Conclusions

- *In flavor physics there are by now several persistent discrepancies with respect to the SM. Their most convincing aspects are the following:*
 - **Experiments:** Results are consistent between LHCb and B factories.
 - **Data:** Deviations concern two independent sets of data: $b \rightarrow s$ and $b \rightarrow c$ decays.
 - **Data vs. theory:** Discrepancies go in a consistent direction.
A BSM explanation is already possible within an EFT approach.
- *Early to draw conclusions. But Run II will provide a definite answer*
- **Theory:** *EFT makes sense rather well of data. But hard to find convincing UV dynamics*
- *Timely to pursue further tests.*
 - Examples:
 - more measurements of R_K
 - more LUV quantities
 - other observables sensitive to C_9 & C_{10}