



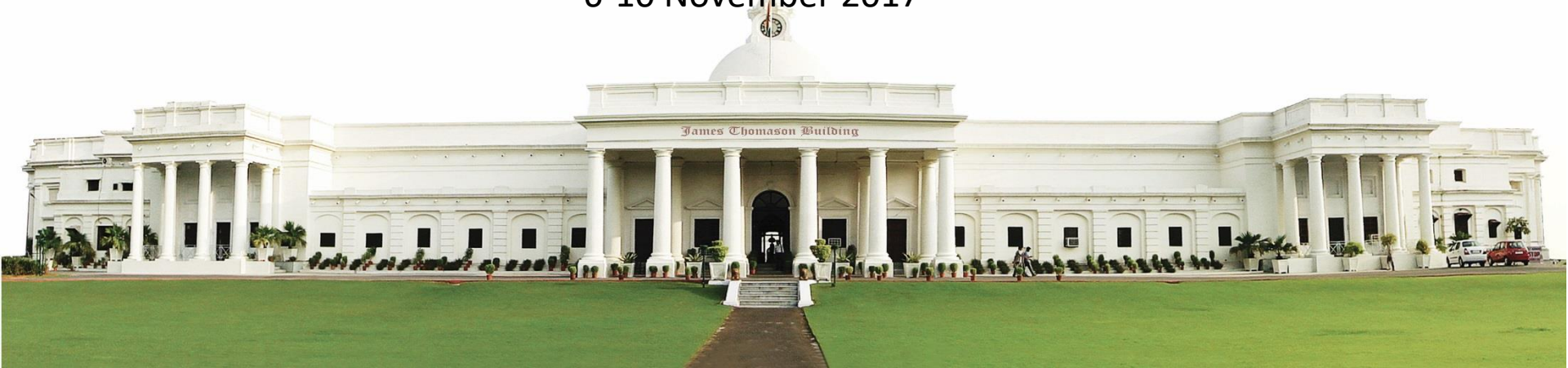
# Nuclear Isomers in and around various Magic Nuclei

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In collaboration with

**Bhoomika and Swati**

International Conference on Shapes and Symmetries in Nuclei – Expt.  
& Theory, CNRS, Paris, France  
6-10 November 2017



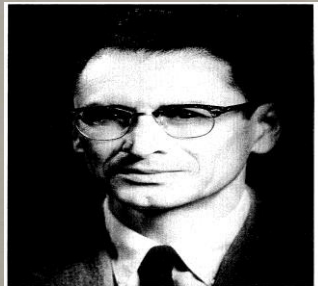


# Outline

- Nuclear Isomers – Feenberg's contribution
- Chart of Isomers & Global systematics related to pairing
- Seniority Isomers
- Key Features of Seniority isomers
- BE2 systematics of  $Z=82$  isotopes and  $N=126$  Isotones
- Extension to multi-j scheme from pure-j scheme
- Selection rules in Multi-j scheme and the new type of seniority isomers which decay by odd-tensor transitions
- BE2 puzzle in the  $2+$  states of Sn isotopes
- g-factors in the  $2+$  and  $10+$  states of Sn isotopes
- BE3 trends of the  $3-$  states in Sn isotopes

# SHELL THEORY OF THE NUCLEUS

By EUGENE FEENBERG



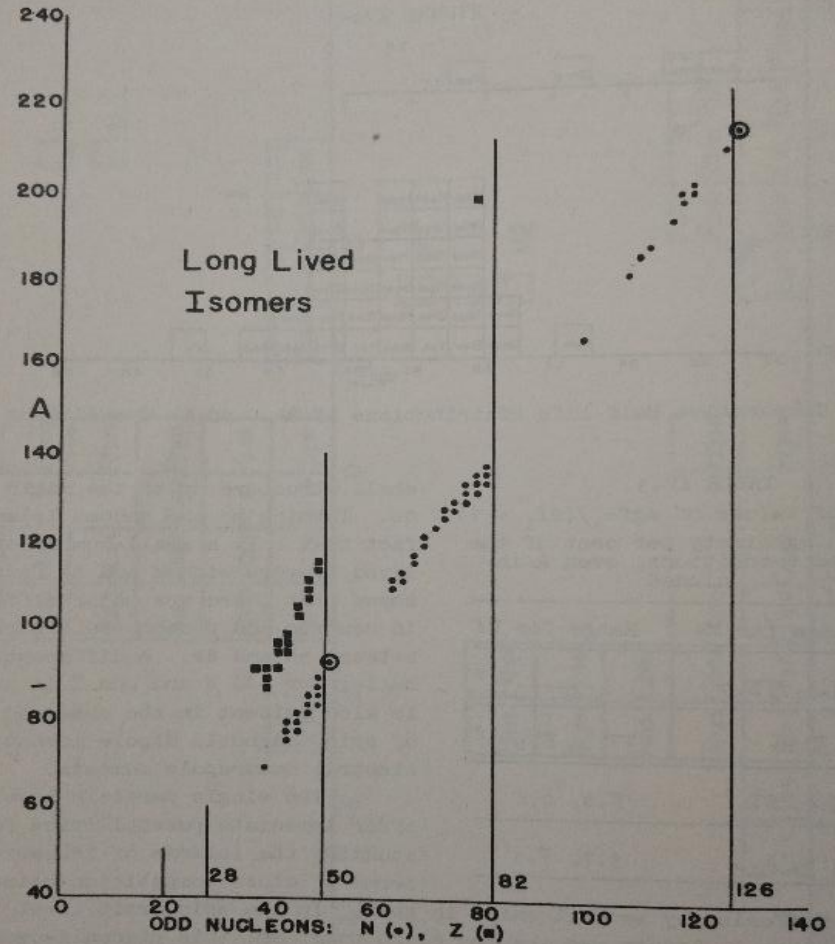
*Eugene Feenberg*

PRINCETON  
 PRINCETON UNIVERSITY PRESS

1955

## IV. ISOMERIC TRANSITIONS

FIGURE IV.3

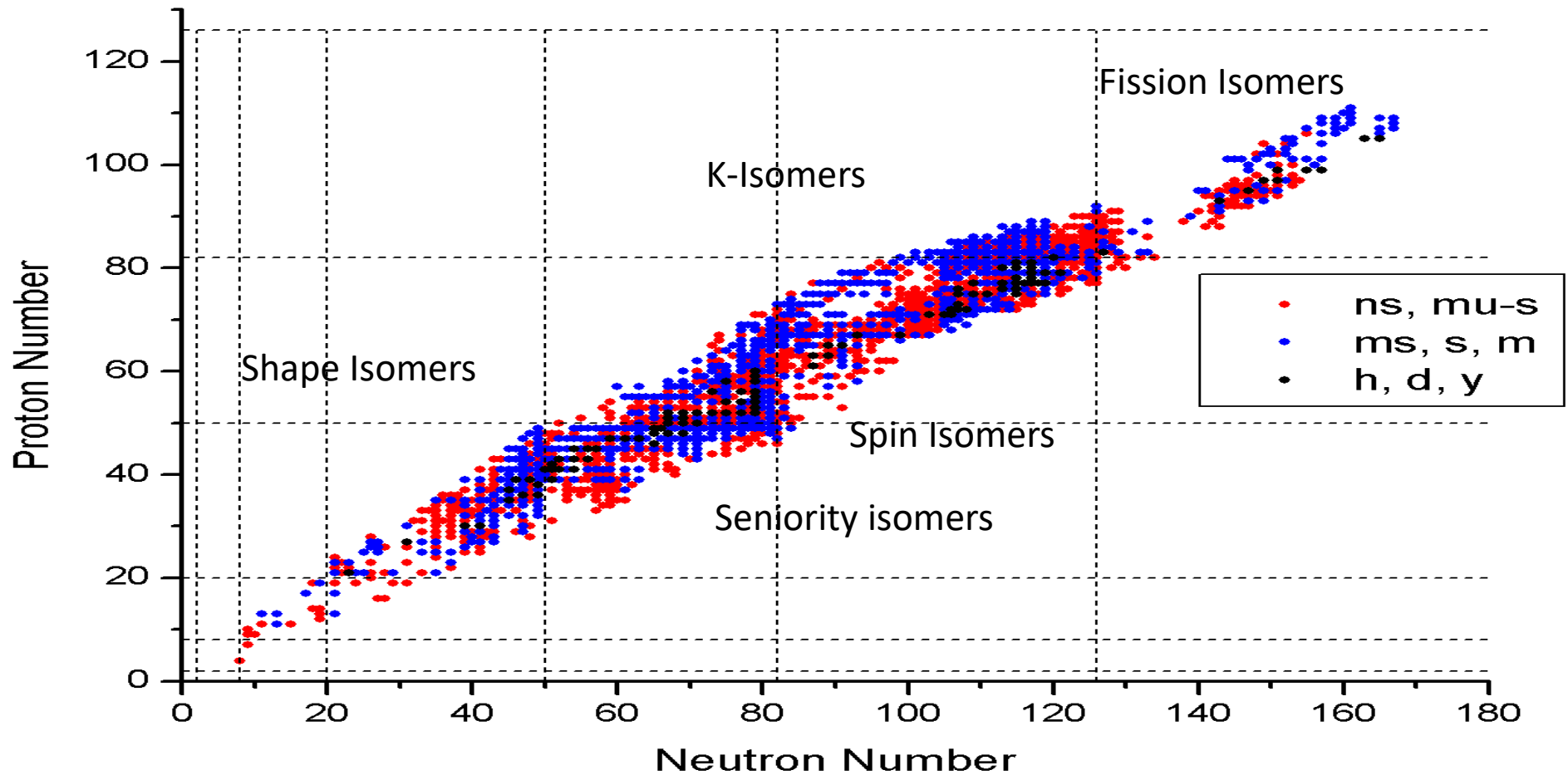


Islands of Isomerism - Distribution of Long-lived Nuclear Isomers of Odd Mass Number (Adapted from Goldhaber and Hill [52])

The circles at  $N = 51$  ( ${}_{42}\text{Mo}_{51}$ ) and  $N = 127$  ( ${}_{84}\text{Po}_{127}$ ) fall outside of the theoretical ranges for islands of isomerism. Goldhaber (53) interprets  ${}_{42}\text{Mo}_{51}$  isomerism in terms of transitions within the even-even core, the nucleon taking no part in the internal rearrangements produced by the transitions.

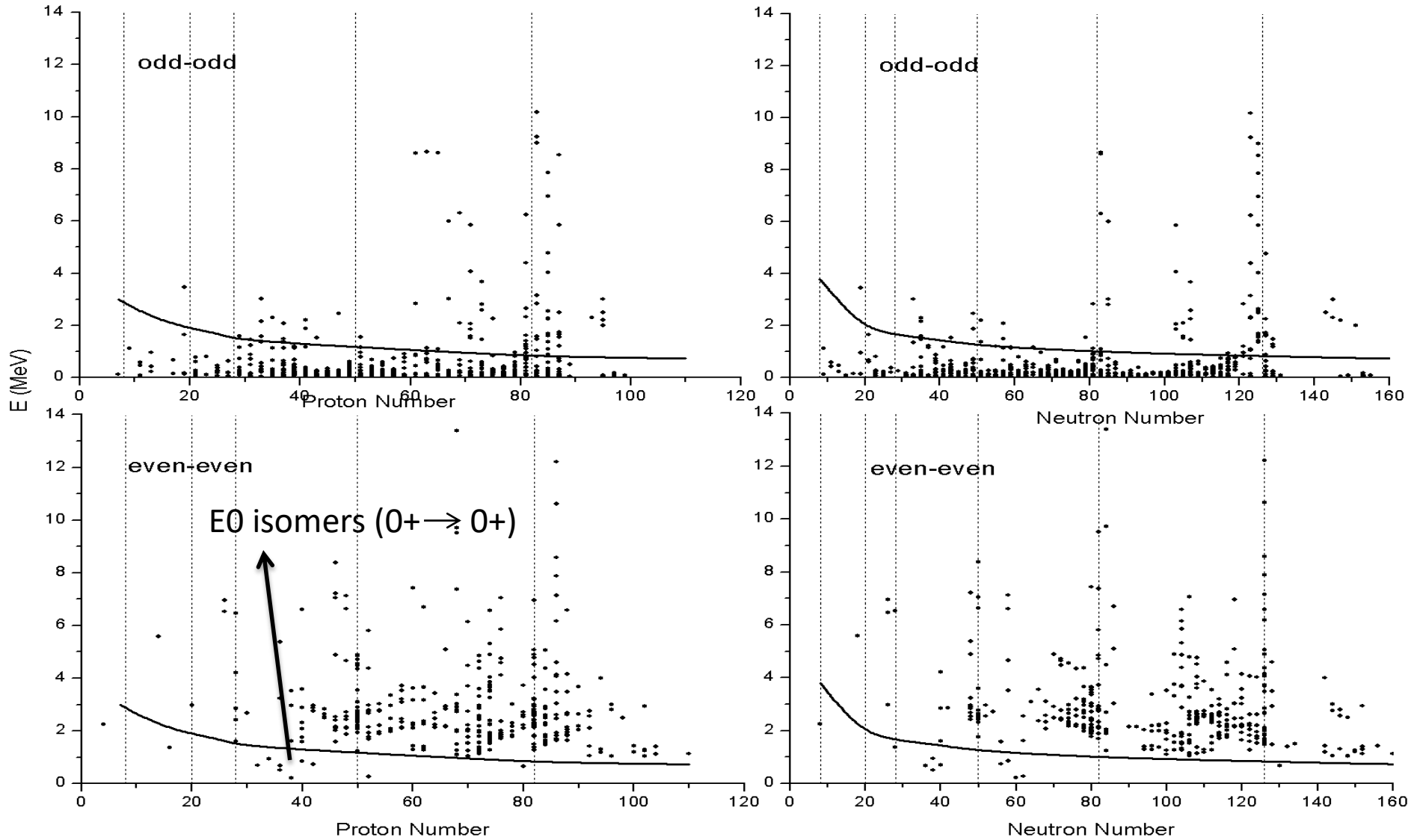
# Chart of Nuclear Isomers

- Nuclear Isomers are longer lived excited states  $\geq 10$  ns.



“Atlas of Nuclear Isomers”, A. K. Jain, B. Maheshwari, S. Garg, M. Patial and B. Singh, Nuclear Data Sheets **128** (2015) 1-130.

# Even-even & odd-odd isomers (Even-A)







# Seniority and Seniority Isomers

- Seniority was introduced by Racah (1943) in atomic context
- Adopted for nuclei in a similar fashion.
- Seniority quantum number ( $\nu$ ) may be defined as the number of unpaired nucleons.
- A set of states diagonalizable by a short range pairing interaction give rise to good or, nearly good seniority states
- Specific selection rules emerge, which may lead to isomeric states in single closed shell nuclei

*Ref.: Books of Talmi, Lawson, Griener and Maruhn, and Casten carry an excellent technical/general description of the seniority quantum number.*

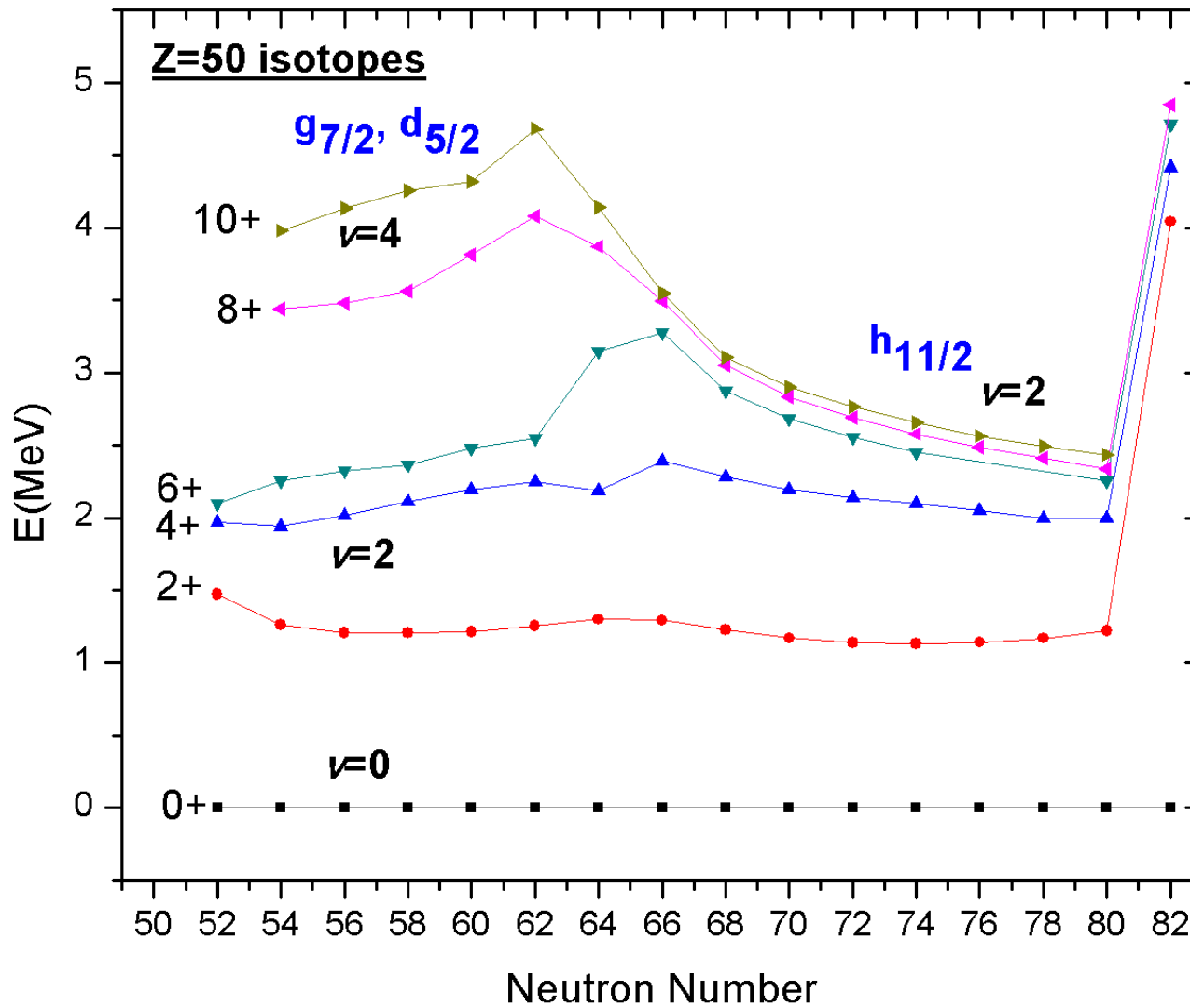
# Key features

I. Talmi, Simple Models of Complex Nuclei. (Harwood, 1993).



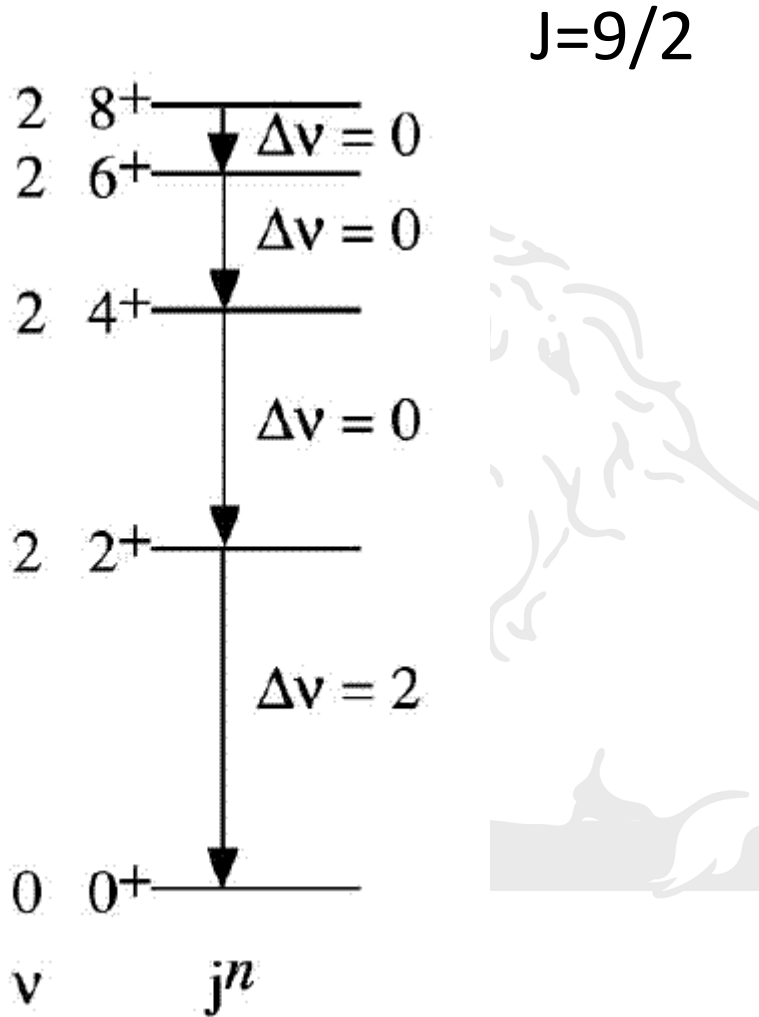
- Particle number independent energy
- Constant pairing gap
- Specific selection rules and parabolic behavior of transition probabilities
- Any interaction between identical fermions in single- $j$  shell exactly conserves seniority if  $j \leq 7/2$ .
- The seniority is nearly conserved up to  $j=11/2$  in  $S_n$ -isomers after the mid-shell, where small seniority mixing may take place, but still the features of seniority persist.

# Particle number independence of energy



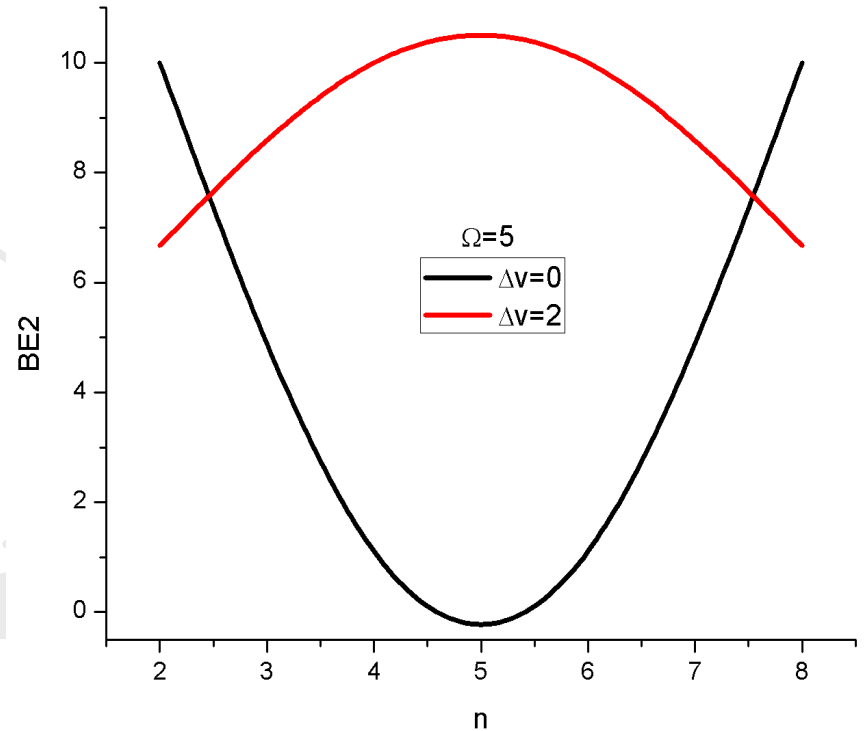


# BE2 variation in single-j: Schematic only



$$B(E2) \propto \left( \frac{\Omega - n}{\Omega - v} \right)^2, \Delta v = 0$$

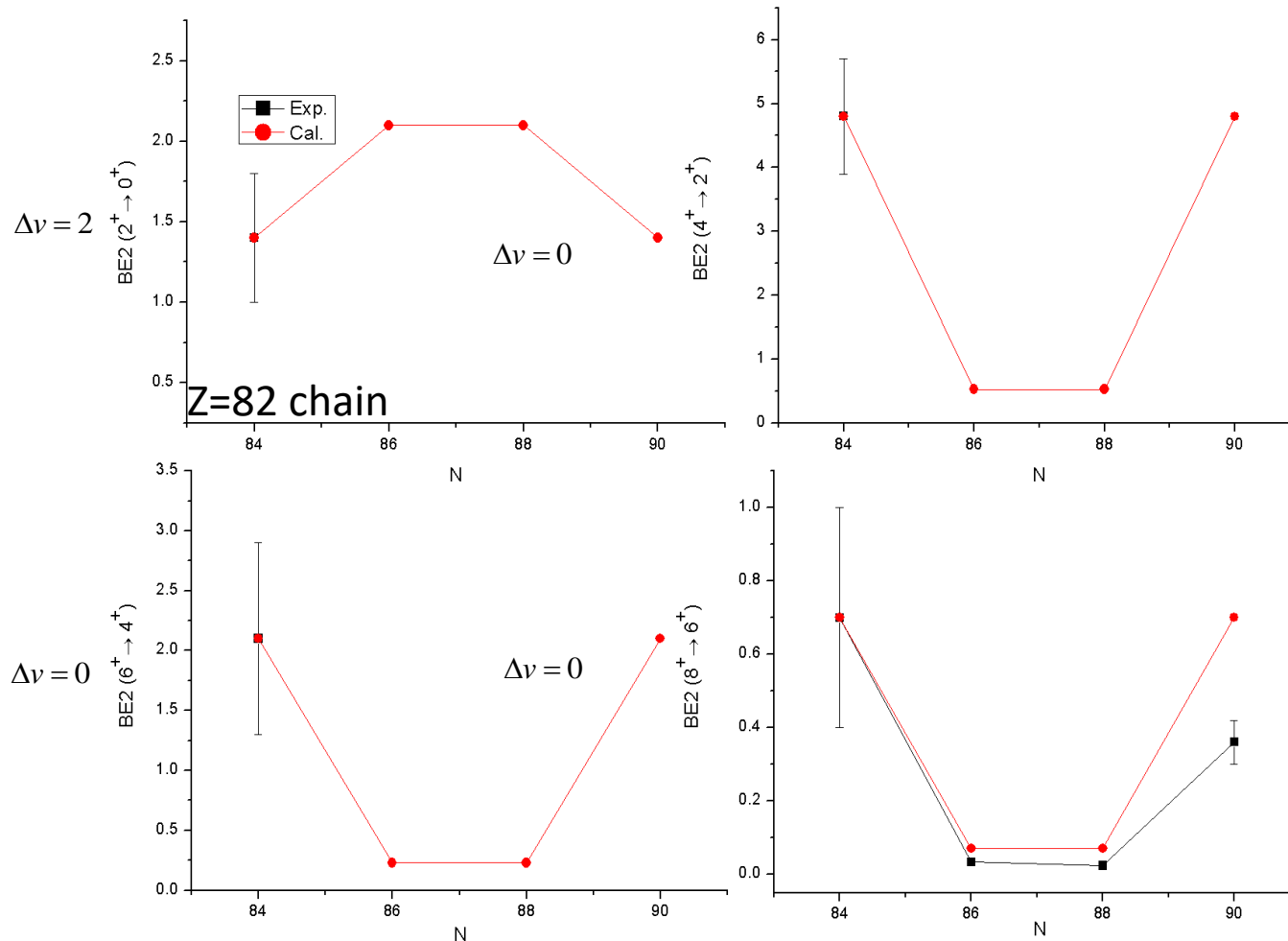
$$B(E2) \propto \frac{(n - v + 2)(2\Omega + 2 - n - v)}{2(2\Omega + 2 - 2v)}, \Delta v = 2$$



For schematic representation only.

# B(E2)s in the even-even Z=82, 8+ isomers

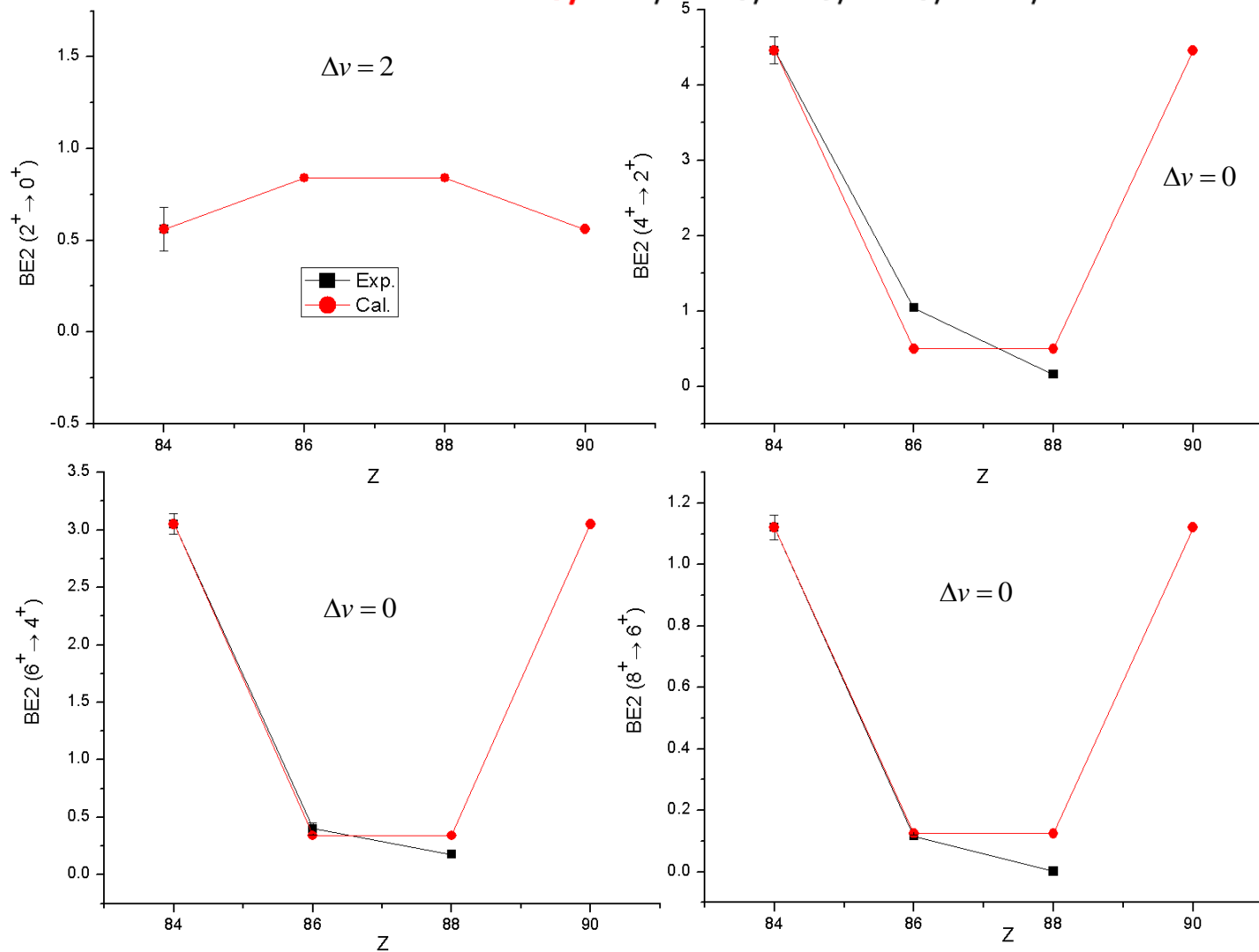
Neutrons ( $g_{9/2}, i_{11/2}, j_{15/2}, d_{5/2}, s_{1/2}, g_{7/2}, d_{3/2}$ )



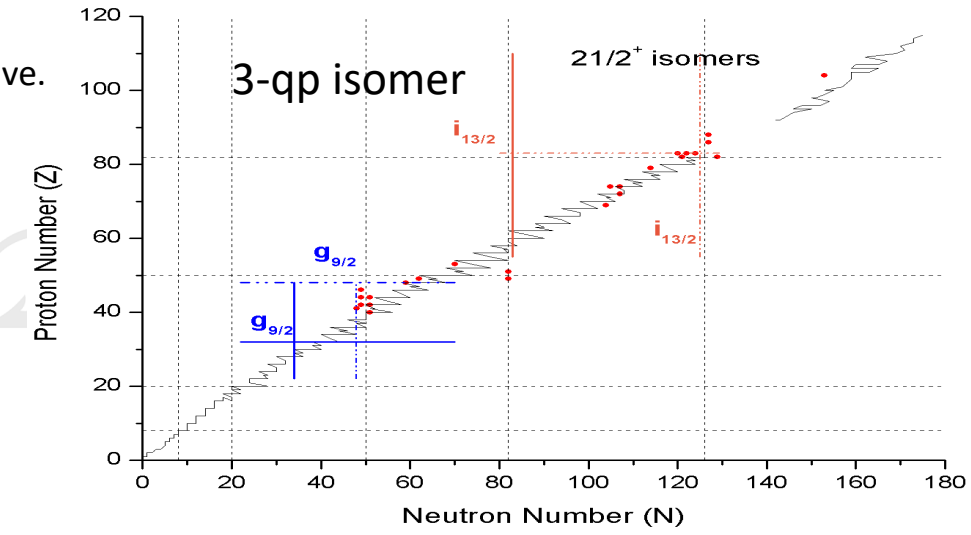
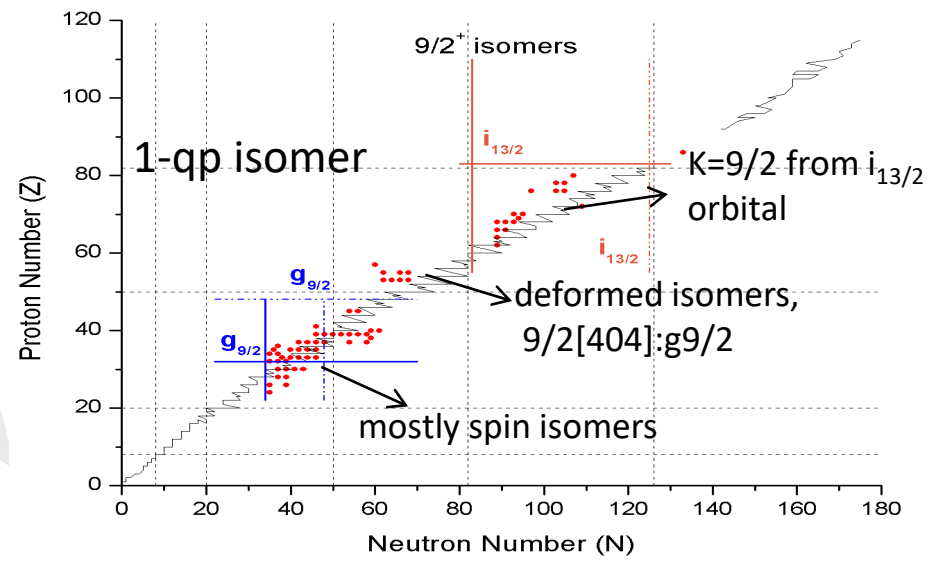
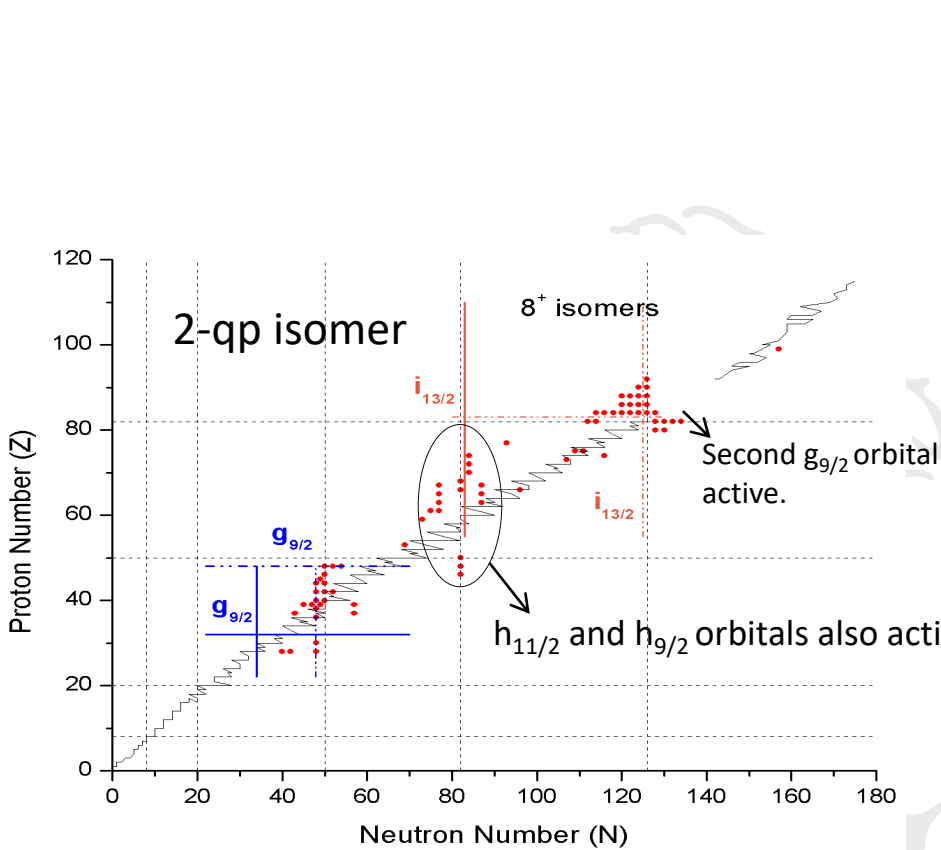
All BE2s are in Weisskopf Units

# B(E2)s in the even-even N=126 chain, 8+ isomers

Protons ( $h_{9/2}, f_{7/2}, i_{13/2}, f_{5/2}, p_{3/2}, p_{1/2}$ )



# $g_{9/2}$ : $9/2^+$ , $8^+$ and $21/2^+$ isomers



A. K. Jain and B. Maheshwari, <https://www-nds.iaea.org/publications/indc/indc-ind-0048/>

# B(E2) values from Seniority scheme

$$B(EL) = \frac{1}{2J_i + 1} \left| \left\langle J_f \left\| \sum_i r_i^L Y^{(L)}(\theta_i, \phi_i) \right\| J_i \right\rangle \right|^2$$

In single-j case,

$$\Omega = \frac{1}{2}(2j+1)$$

For L even

$$\left\langle j^n \nu l J_f \left\| \sum_i r_i^L Y^{(L)}(\theta_i, \phi_i) \right\| j^n \nu l' J_i \right\rangle = \left( \frac{\Omega - n}{\Omega - \nu} \right) \left\langle j^\nu \nu l J_f \left\| \sum_i r_i^L Y^{(L)}(\theta_i, \phi_i) \right\| j^\nu \nu l' J_i \right\rangle$$

$$\left\langle j^n \nu l J_f \left\| \sum_i r_i^L Y^{(L)}(\theta_i, \phi_i) \right\| j^{n, \nu \mp 2, l' J_i} \right\rangle = \sqrt{\frac{(n - \nu + 2)(2\Omega + 2 - n - \nu)}{2(2\Omega + 2 - 2\nu)}} \left\langle j^\nu \nu l J_f \left\| \sum_i r_i^L Y^{(L)}(\theta_i, \phi_i) \right\| j^{\nu, \nu \mp 2, l' J_i} \right\rangle$$

It is easy to generalize these results for multi-j case by defining,

$$\tilde{j} = j \otimes j' \dots \quad \Omega = \frac{1}{2} \sum_j (2j+1) \quad n = \sum_j n_j$$

Generalized seniority

$$B(E2) \propto \left( \frac{\Omega - n}{\Omega - \nu} \right)^2, \Delta \nu = 0$$

$$B(E2) \propto \frac{(n - \nu + 2)(2\Omega + 2 - n - \nu)}{2(2\Omega + 2 - 2\nu)}, \Delta \nu = 2$$

B(E2)  
relations  
valid for  
single-j, and  
multi-j  
cases!!

# Selections rules (Extension to multi- $j$ shell)

## Single- $j$ shell

- What seniority says.
  - Odd  $L$   $\rightarrow$  quasi-spin scalar  $\rightarrow \Delta v=0$
  - Even  $L$   $\rightarrow$  quasi-spin vector  $\rightarrow \Delta v=0, 2$
- What EM decay says.
  - Odd  $L$   $\rightarrow$  only magnetic transitions.
  - Even  $L$   $\rightarrow$  only electric transitions.



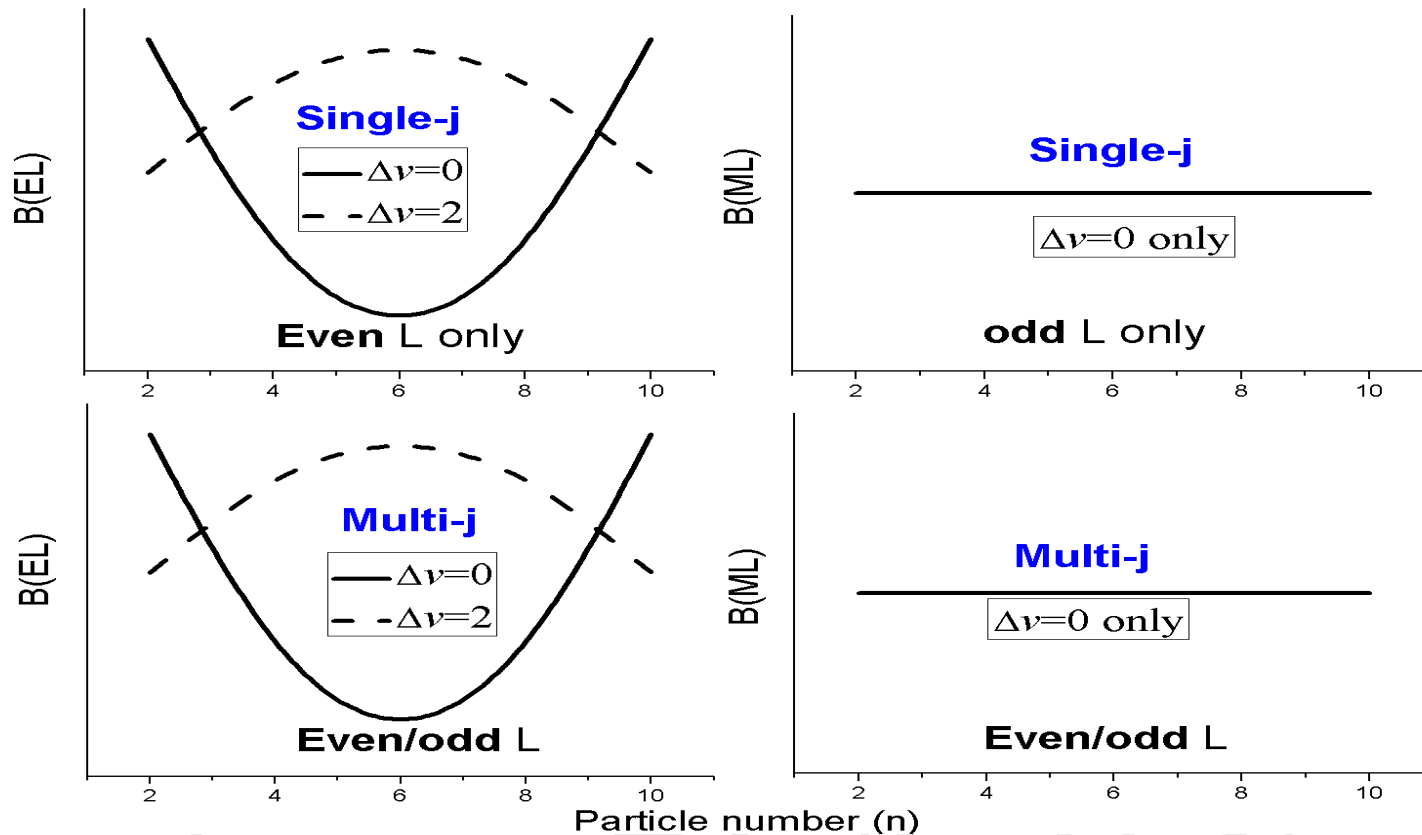
**Odd  $L$ ,**  
**Magnetic transitions,  $\Delta v=0$**   
**Even  $L$ ,**  
**Electric Transitions,  $\Delta v=0, 2$**

## Multi- $j$ shell

- What seniority says.
  - **Odd/Even  $L$   $\rightarrow$  Magnetic transitions  $\rightarrow$  quasi-spin scalar  $\rightarrow \Delta v=0$**
  - **Odd/Even  $L$   $\rightarrow$  Electric transitions  $\rightarrow$  quasi-spin vector  $\rightarrow \Delta v=0, 2$**
- What EM decay says.
  - Odd/Even  $L$   $\rightarrow$  both electric/ magnetic transitions.



# Comparison of single-j and multi-j shell



B. Maheshwari and A. K. Jain, Phys. Lett. B **753**, 122 (2016).

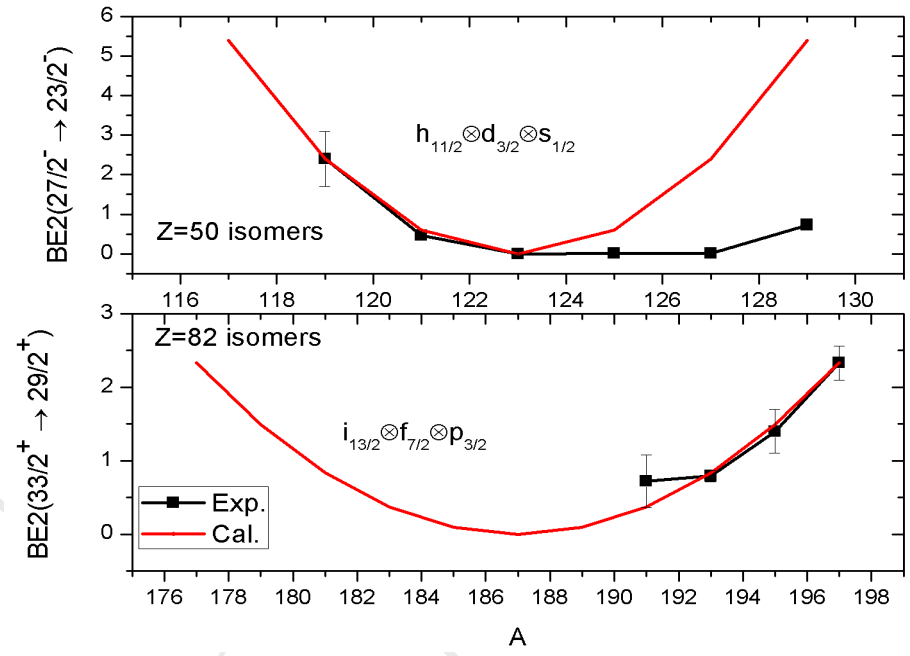
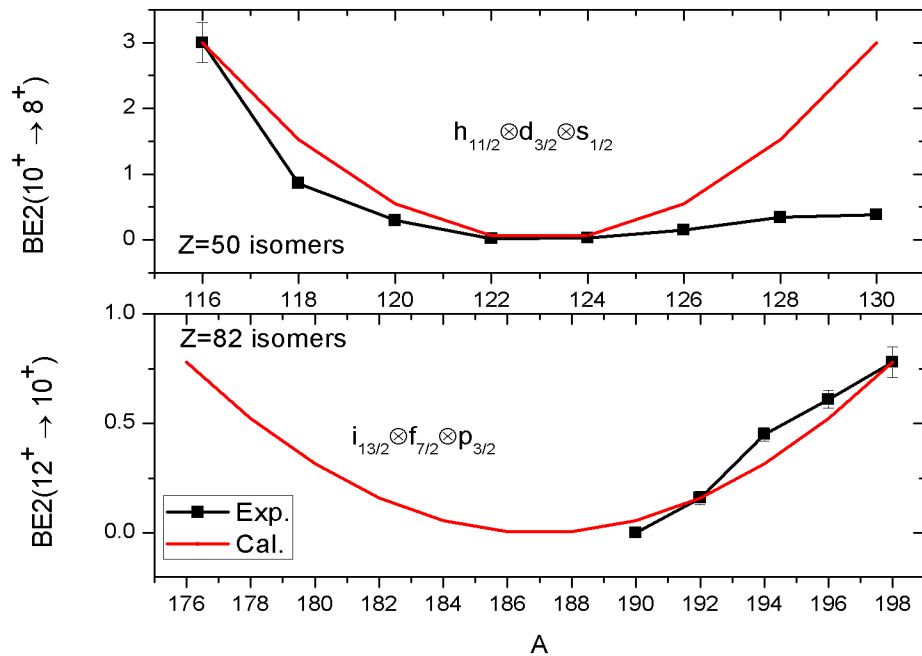
B. Maheshwari, A. K. Jain and B. Singh, Nucl. Phys. A **952**, 62 (2016).

A.K. Jain and B. Maheshwari, Nuclear Physics Review **34**, 73 (2017).

A.K. Jain and B. Maheshwari, Physica Scripta **92**, 074004 (2017).

B. Maheshwari, Swati and A. K. Jain, Pramana-Journal of Physics **89**, 75 (2017).

# BE2s of High-spin isomers in Sn and Pb isotopes



- This concludes that the  $12+$ ,  $10+$ ,  $33/2+$  and  $27/2+$  isomers, arising mainly from the intruder orbitals, also require the configuration mixing to explain the measured data.
- The right hand sides of  $Z=50$  isomers which do not follow this generalized seniority scheme, may further need the inclusion of the non-degenerate multi-j orbits.

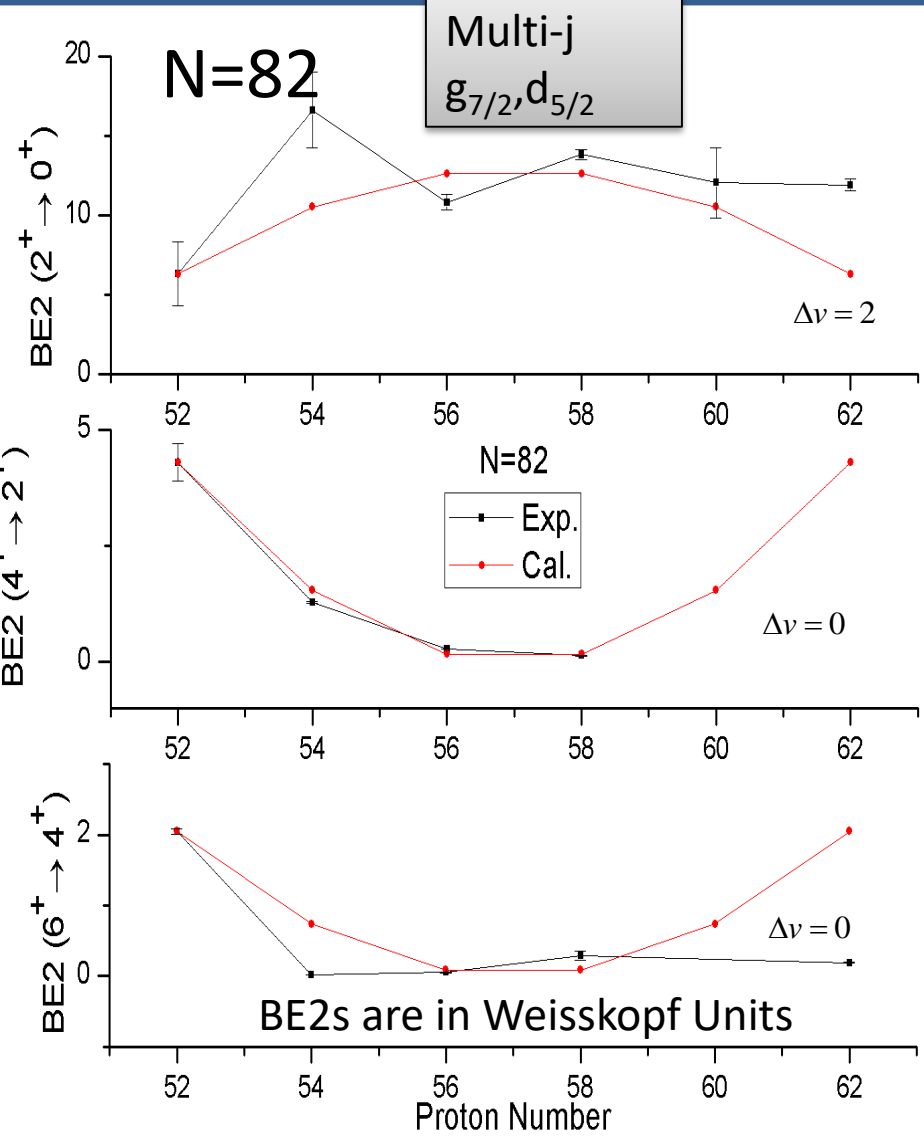
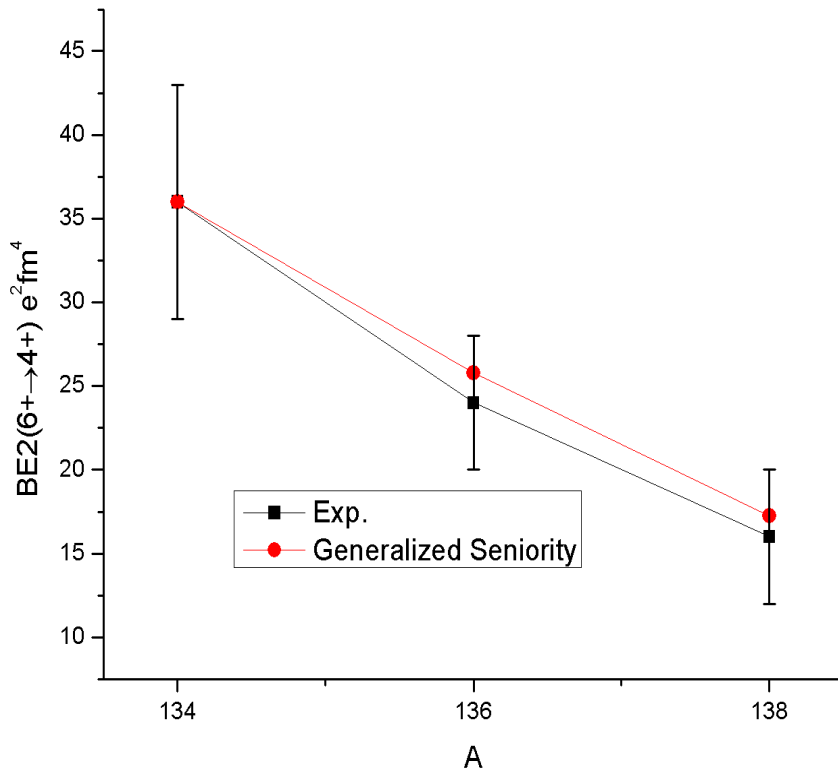
# B(E2)s from generalized seniority in Z=50 & N=82 chains

Z=50

Multi-j

$f_{7/2}, p_{3/2}, p_{1/2}, f_{5/2}, h_{9/2}$

$$\Delta\nu = 0$$

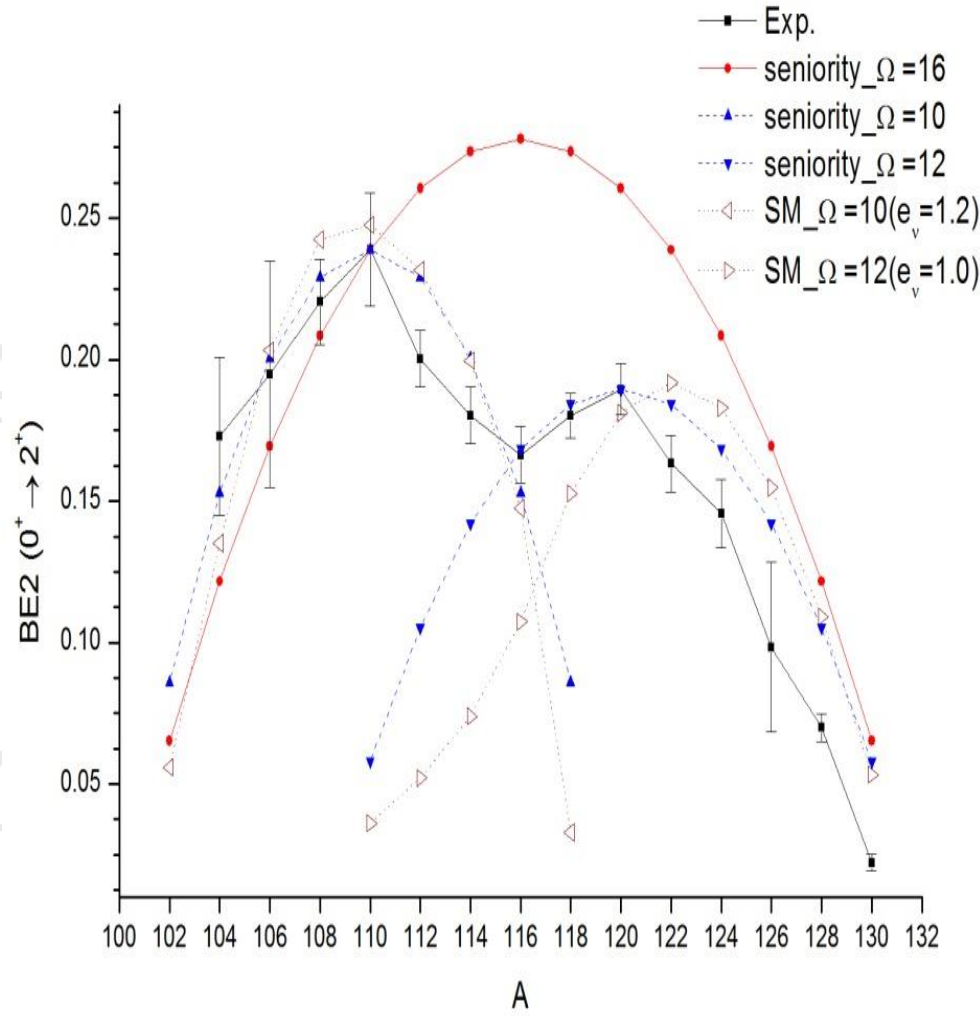


# The first excited $2^+$ states in Sn isotopes

BM, AKJ & BS, Nucl. Phys. A952,62,2016

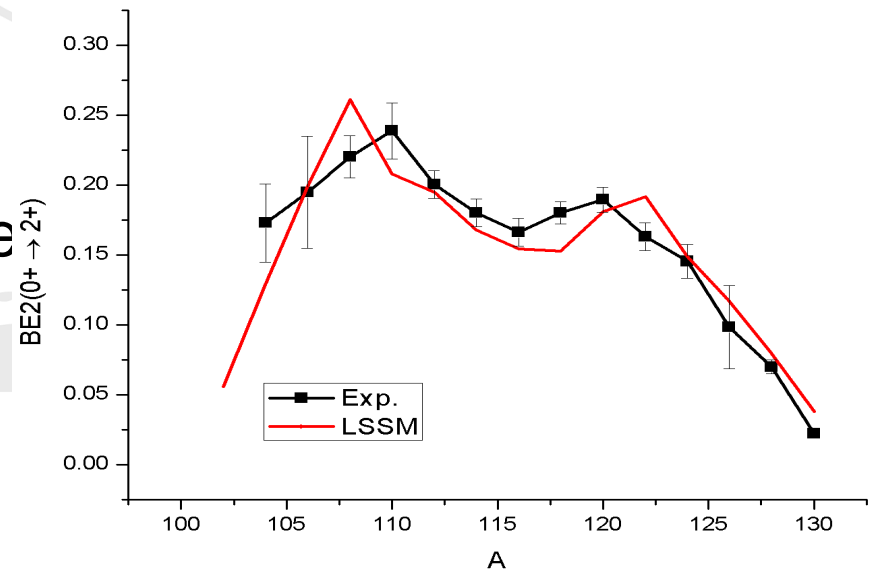
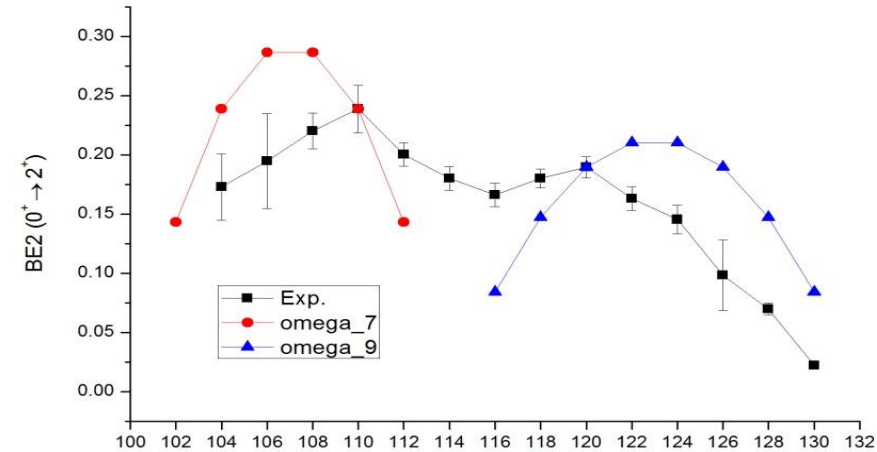
- Expected: a single parabolic trend coming from the generalized seniority.
- Reality: A minima in the middle of the valence space in recent measurements.
- Reason: The different sets of orbits involved before and after midshell, as also discussed by Morales et al.

I. O. Morales, P. Van Isacker and I. Talmi, Phys. Lett. B **703**, 606-608 (2011).



# Seniority inspired LSSM

- Failure: if we use omega 7 and 9, which comes from the  $g_{7/2}$ ,  $d_{5/2}$  and  $h_{11/2}$ ,  $d_{3/2}$ ,  $s_{1/2}$  configuration mixing sets.
- Success: Seniority inspired and truncated LSSM explains the measured data quite well, particularly around the middle, where dimensions become quite large.
- 102-108, 124-130 Sn have been treated in open space.



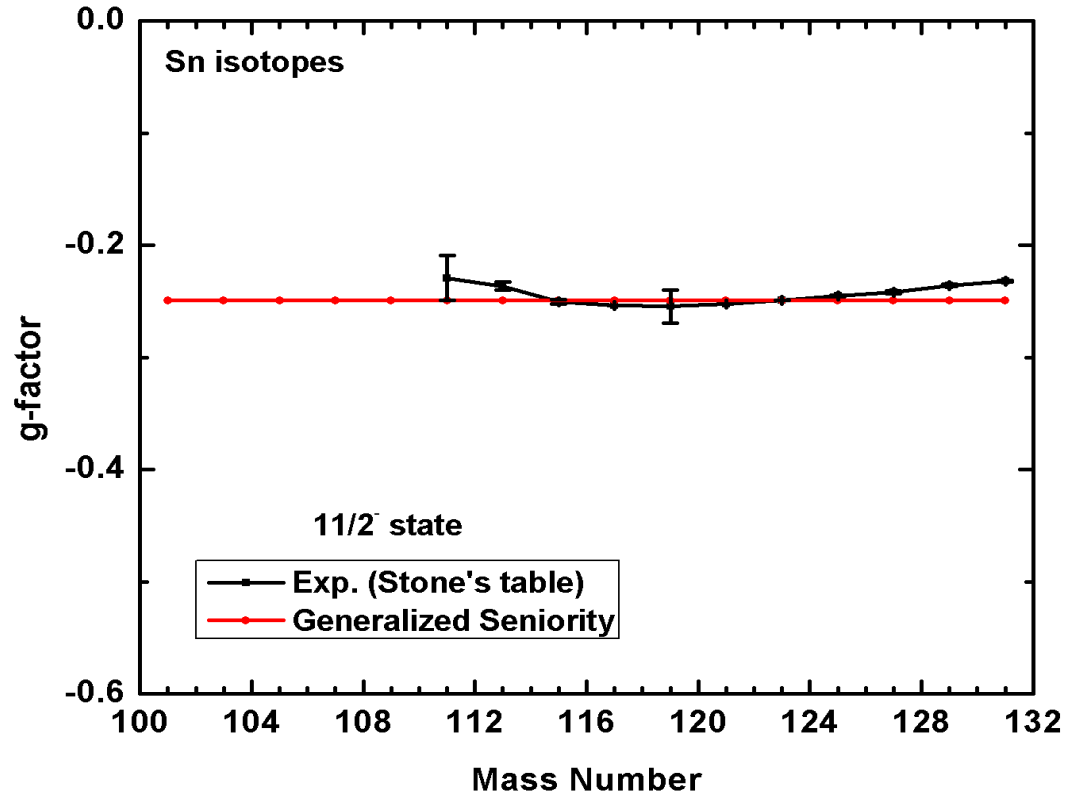
# g-factor

The magnetic moments and g-factor must also show a particle number independent behavior.

The magnetic moment of identical nucleons in pure  $j^n$  configuration is given by (Talmi's book)

$$\begin{aligned} \mu &= \sum g j_i = g \sum j_i \\ &= g J \end{aligned}$$

Hence, g-factors of all states of  $j^n$  configuration of identical nucleons should be equal to the g-factor of a single-j nucleon.



N.J. Stone, INDC(NDS)-0658, Feb. 2014  
[www-nds.iaea.org/publications](http://www-nds.iaea.org/publications)



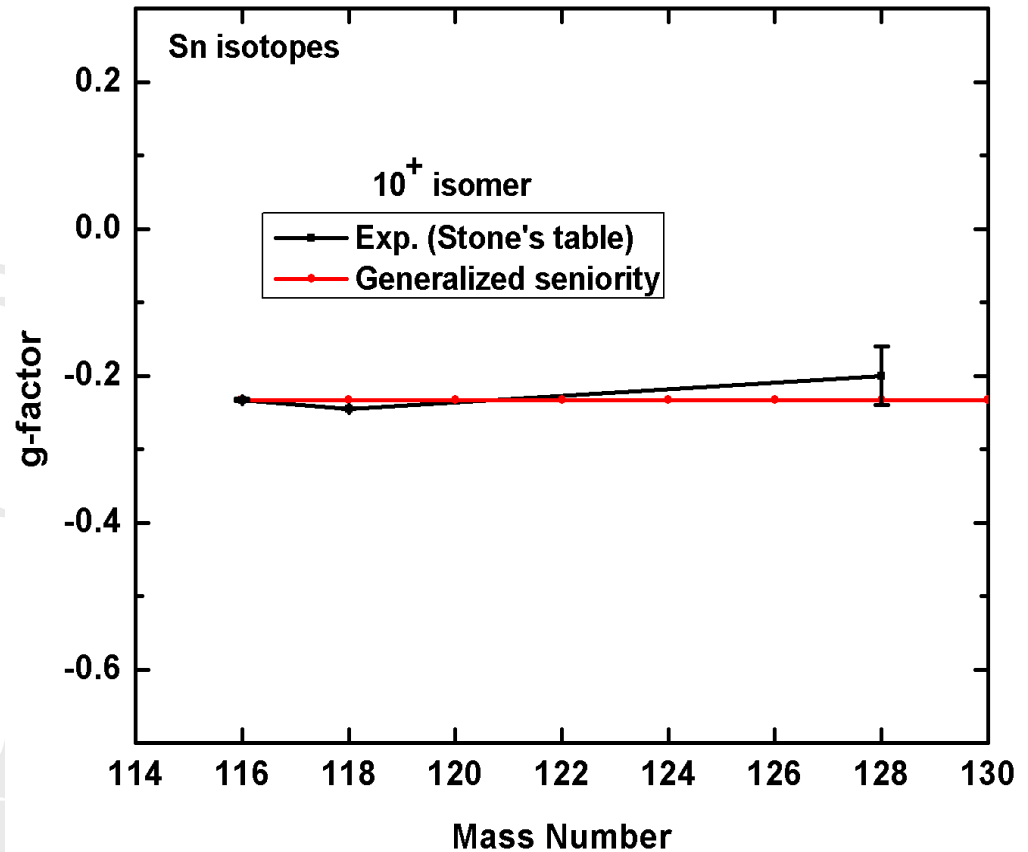
# g-factor trend for $10^+$ state in Sn isotopes

The magnetic moment of identical nucleons in the mixed configuration  $\tilde{j}^n$  is given by

$$\mu = g \sum_{i=1}^n \tilde{j} = gJ$$

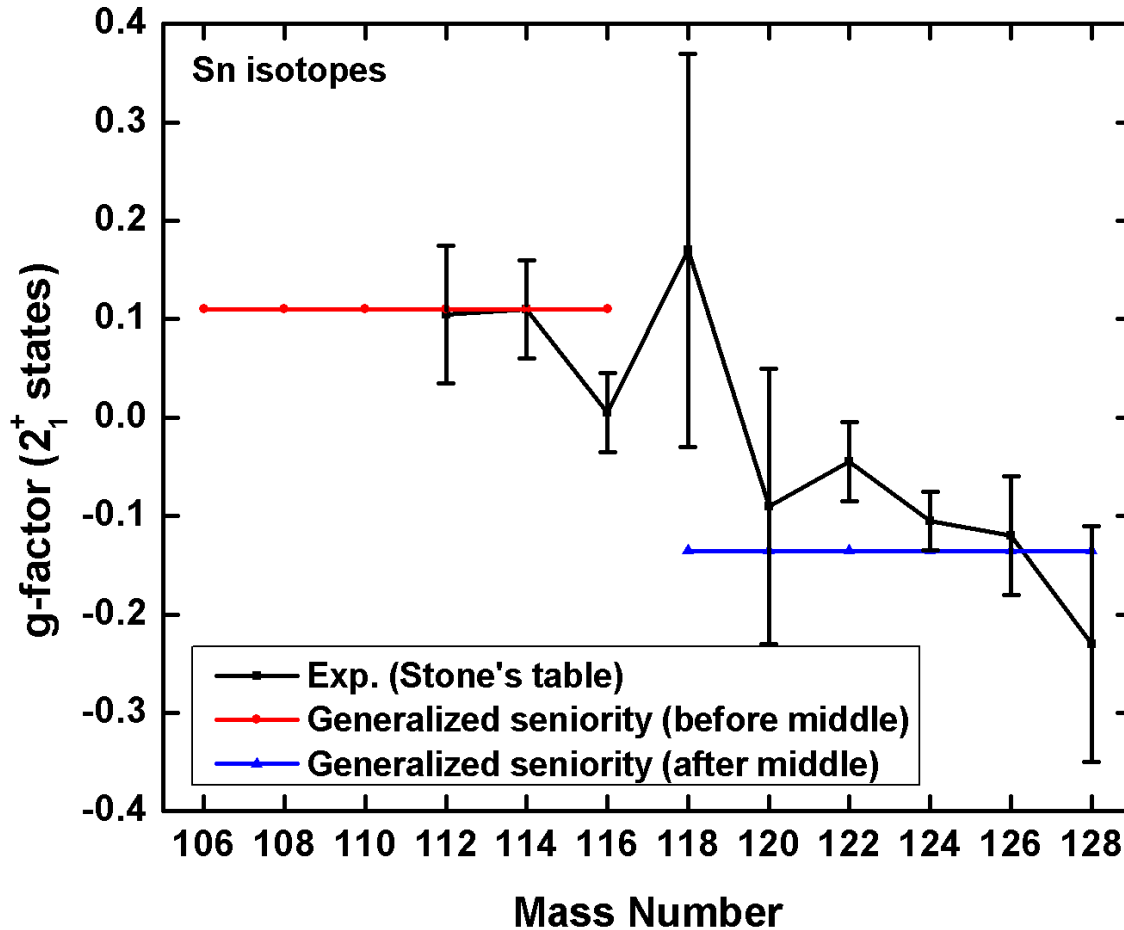
Hence, g-factors of all  $10^+$  states of configuration  $\tilde{j}^n$  of identical nucleons should be equal to the g-factor of a single nucleon.

However, the data on  $27/2^-$  isomer is not available till date. We expect and predict the g-factor of this state to be of the same order as for  $10^+$  state, since both follow same configuration mixing as shown in B(E2) results.



**Points out the need of experiments in this direction!**

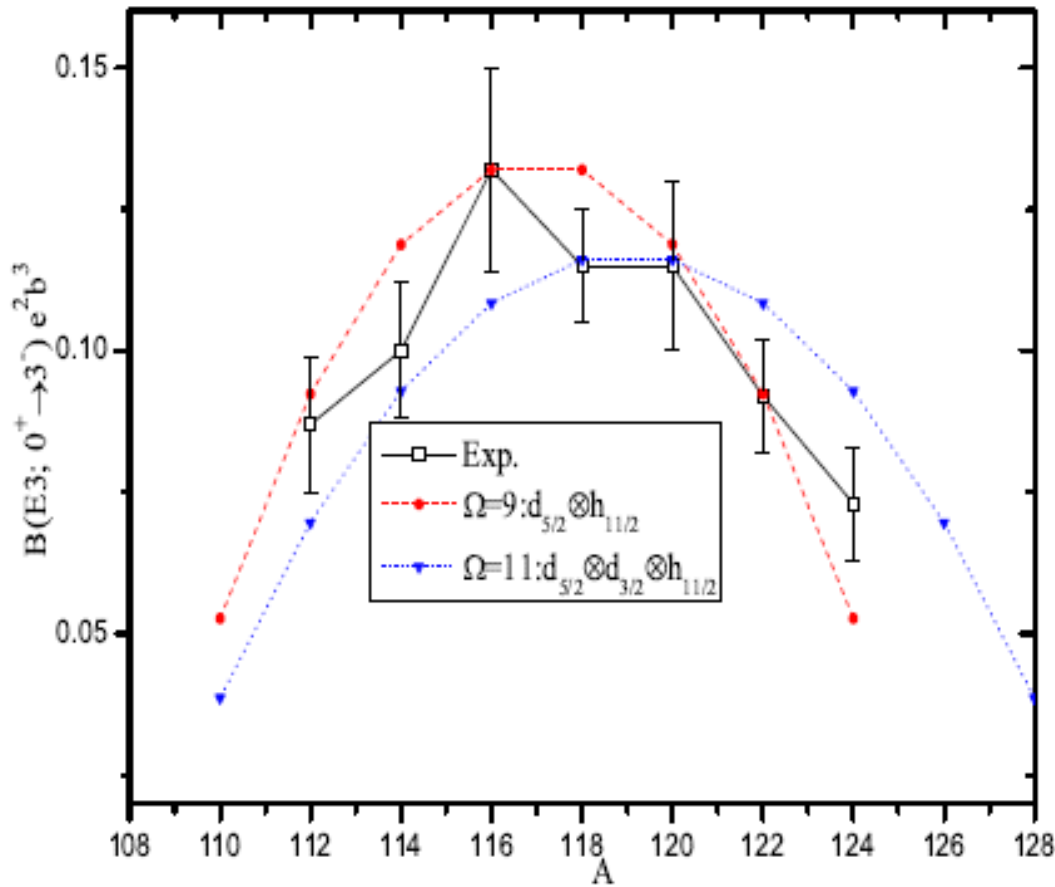
# g-factor of first excited $2^+$ states in Sn isotopes



Two parabolas in B(E2) curve highlights the difference in the configuration before and after the middle for the generation of the  $2^+$  states in Sn isotopes.

Therefore, g-factor before and after the middle are expected to be different.

# BE3 in Yrast 3- states in Sn isotopes

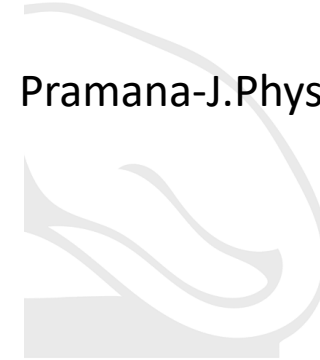


$d_{5/2}, h_{11/2}$  together explain the data very well.

Suggests an octupole character for the 3- states.

This interpretation is supported by the shell model calculations.

Pramana-J.Phys. 89, 75 (2017)



# Summary

- We have presented a view of the seniority nuclear isomers in various semi-magic chains of isotopes.
- Their various properties and unique features have been studied.
- Examples from  $Z=50$ ,  $Z=82$ ,  $N=126$  chains have been presented.
- Both seniority conserving and seniority changing transitions can be understood very well.
- Many predictions are possible based on this scheme.

**Thank You**

