

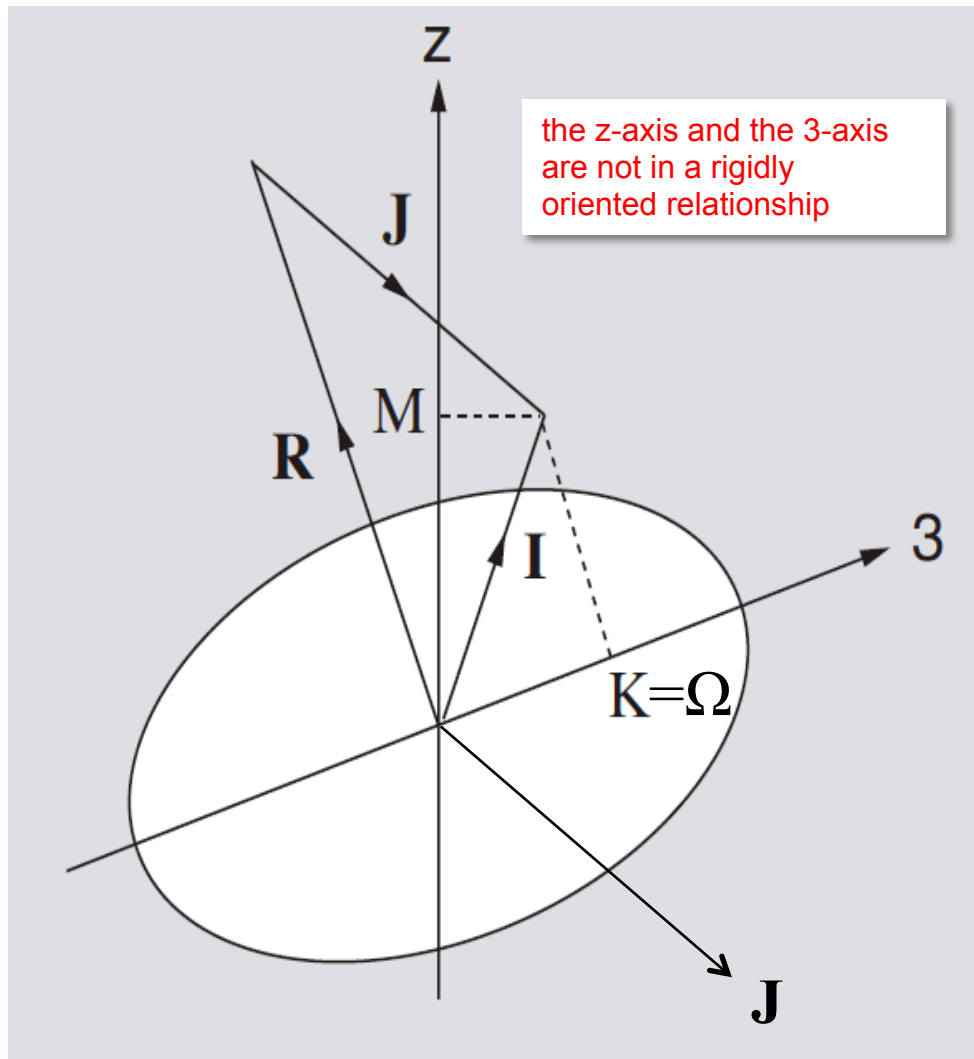
Particle-core coupling in nuclei:
new insights, fundamental misconceptions

John L. Wood

School of Physics

Georgia institute of Technology

Symmetric-top model: quantum numbers



R : collective angular momentum

J : intrinsic spin

I : total spin / angular momentum

M : laboratory-frame,
z-component of I

K : body-frame (symmetry axis),
3-component of I ; $K = \Omega$

“Coriolis” interaction:

$$\begin{aligned} \mathbf{R} \cdot \mathbf{R} &= (\mathbf{I} - \mathbf{J}) \cdot (\mathbf{I} - \mathbf{J}) \\ &= \mathbf{I} \cdot \mathbf{I} - 2 \mathbf{I} \cdot \mathbf{J} + \mathbf{J} \cdot \mathbf{J} \end{aligned}$$

Nilsson model plus rotations

$$\hat{H} = \frac{\hbar^2 \hat{\mathbf{I}}^2}{2\mathfrak{S}} + \hat{h} + \frac{\hbar^2 \hat{\mathbf{j}}^2}{2\mathfrak{S}} - \frac{\hbar^2}{\mathfrak{S}} \hat{\mathbf{I}} \cdot \hat{\mathbf{j}}.$$

$$\langle \vec{I} \cdot \vec{j} \rangle = \left\langle \frac{1}{2} (\hat{I}_+ \hat{j}_- + \hat{I}_- \hat{j}_+) + \hat{I}_z \hat{j}_z \right\rangle$$

$$\begin{aligned} I_+ &:= I_x + iI_y, I_- = I_x - iI_y \\ j_+ &:= j_x + ij_y, j_- = j_x - ij_y \end{aligned}$$

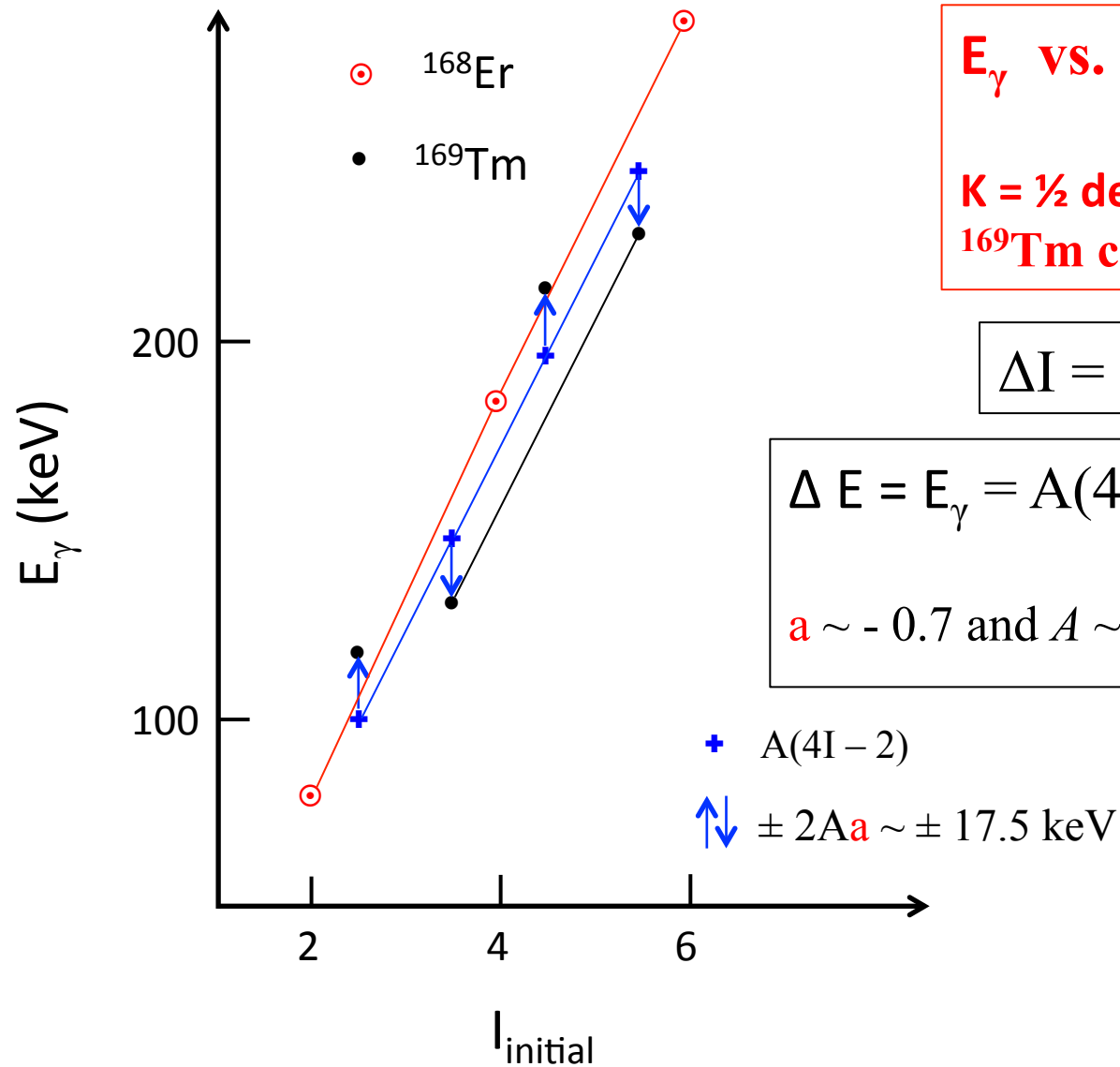
$$|KIM\rangle + \varepsilon(-1)^{I+K} | -K, IM\rangle$$

$\varepsilon = +1$ reflection symmetric
 $\varepsilon = -1$ reflection asymmetric

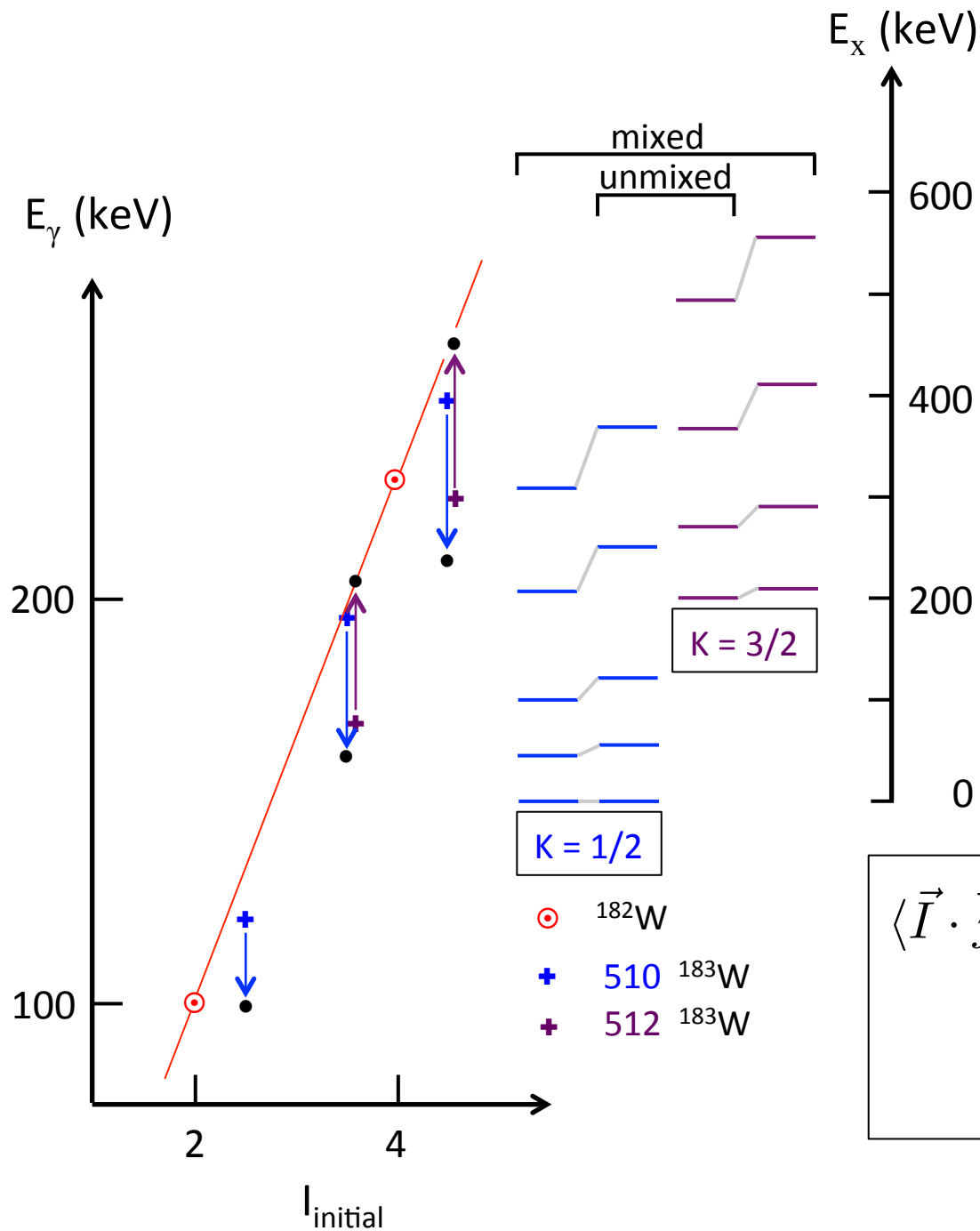
$$E_I = E_0 + A[I(I+1) + (-1)^{I+1/2} (I+1/2) a \delta_{K,1/2}]$$

$\delta_{K,1/2} = 1, K = 1/2$
 $\delta_{K,1/2} = 0$ otherwise

^{169}Tm
 $\frac{1}{2}^+$ [411]



$$E_I = E_0 + A[I(I + 1) + (-1)^{I+1/2}(I + 1/2)a\delta_{K,1/2}]$$



E_γ vs. I_{initial} plot

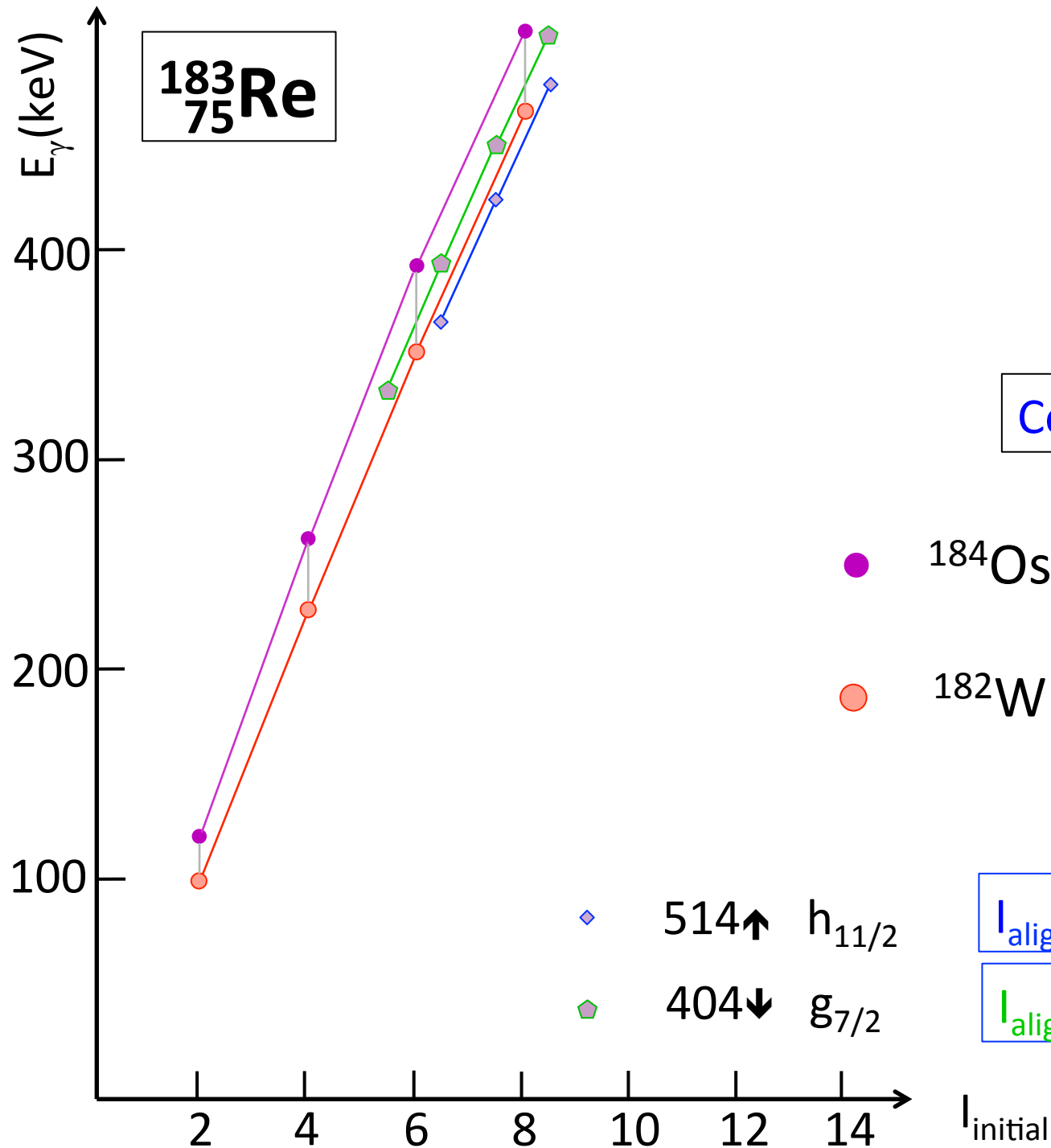
$\Delta K = 1$ band mixing
 $1/2^- 510 + 3/2^- 512$
 ^{183}W cf. ^{182}W

$\Delta I = 2$ transitions

E_γ "shifts" ~ 35 keV

$$\langle \vec{I} \cdot \vec{j} \rangle = \langle \frac{1}{2} (\hat{I}_+ \hat{j}_- + \hat{I}_- \hat{j}_+) + \hat{I}_z \hat{j}_z \rangle$$

--produces $\Delta K = \Delta \Omega = 1$ mixing



E_γ vs. I_{initial} plot
 Identical bands

$\Delta I = 2$ transitions

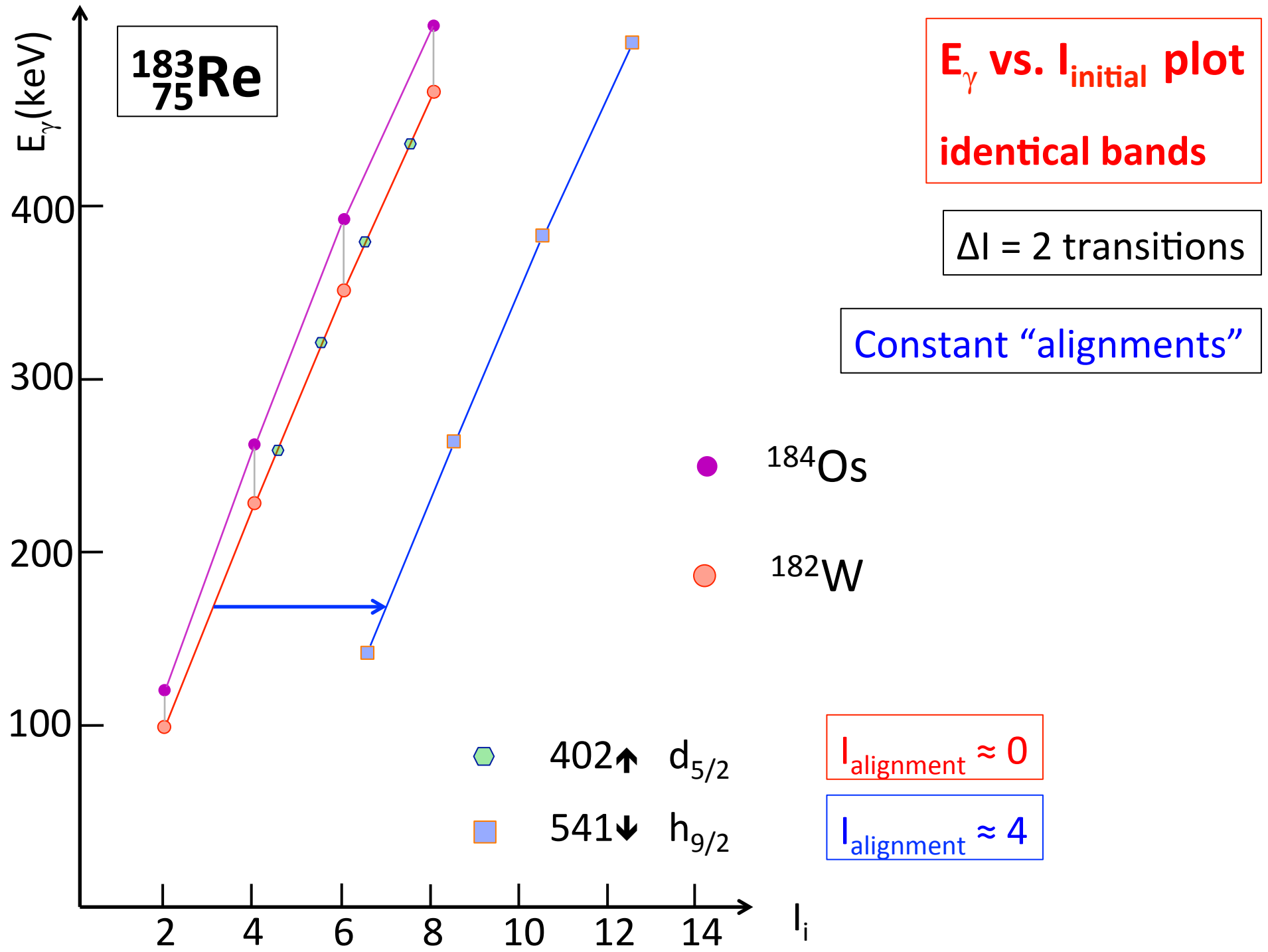
Constant "alignments"

● ^{184}Os
 ● ^{182}W

◇ $514\uparrow$ $h_{11/2}$
 ◇ $404\downarrow$ $g_{7/2}$

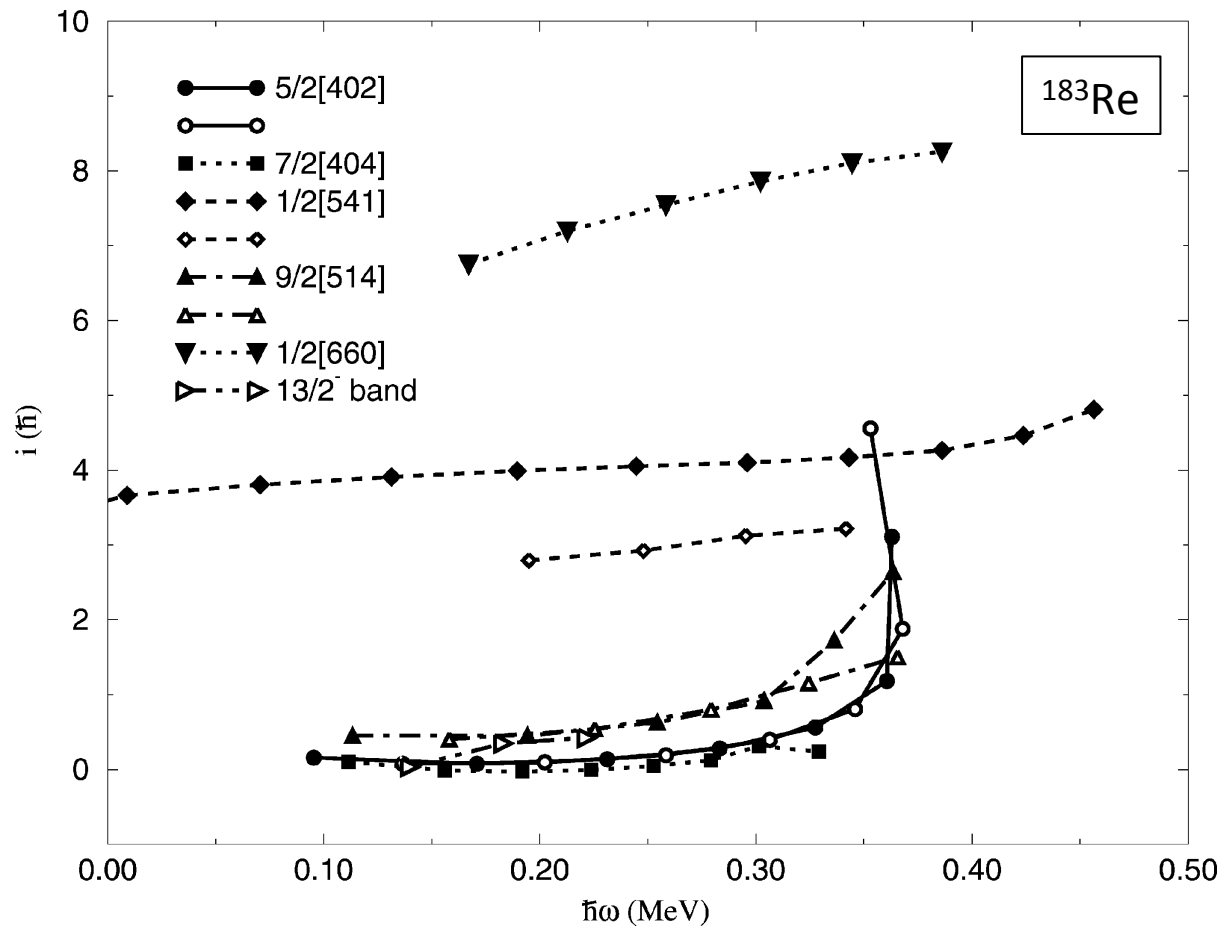
$I_{\text{alignment}} \approx 0.2$

$I_{\text{alignment}} \approx -0.3$



“Standard” alignment plots for ^{183}Re

$$i = (\mathcal{J}_0 + \omega^2 \mathcal{J}_1) \omega \quad \text{Harris formula}$$



Alignment of 402 and 404 recognized under name “identical bands” —25 years ago

VOLUME 69, NUMBER 10

PHYSICAL REVIEW LETTERS

7 SEPTEMBER 1992

Low-Spin Identical Bands in Neighboring Odd-*A* and Even-Even Nuclei: A Possible Challenge to Mean-Field Theories

C. Baktash, J. D. Garrett, D. F. Winchell, and A. Smith

Physics Division, Oak Ridge National Laboratory, Oak Ridge, Tennessee 37831-6371

(Received 3 February 1992)

A comprehensive study of odd-*A* rotational bands in normally deformed rare-earth nuclei indicates that a large number of seniority-one configurations (30% for odd-*Z* nuclei) at low spin have moments of inertia nearly identical to that of the seniority-zero configuration of the neighboring even-even nucleus with one less nucleon. It is difficult to reconcile these results with conventional models of nuclear pair correlation, which predict variations of about 15% in the moments of inertia of configurations differing by one unit in seniority.

Annu. Rev. Nucl. Part. Sci. 1995. 45:485–541

IDENTICAL BANDS IN DEFORMED AND SUPERDEFORMED NUCLEI

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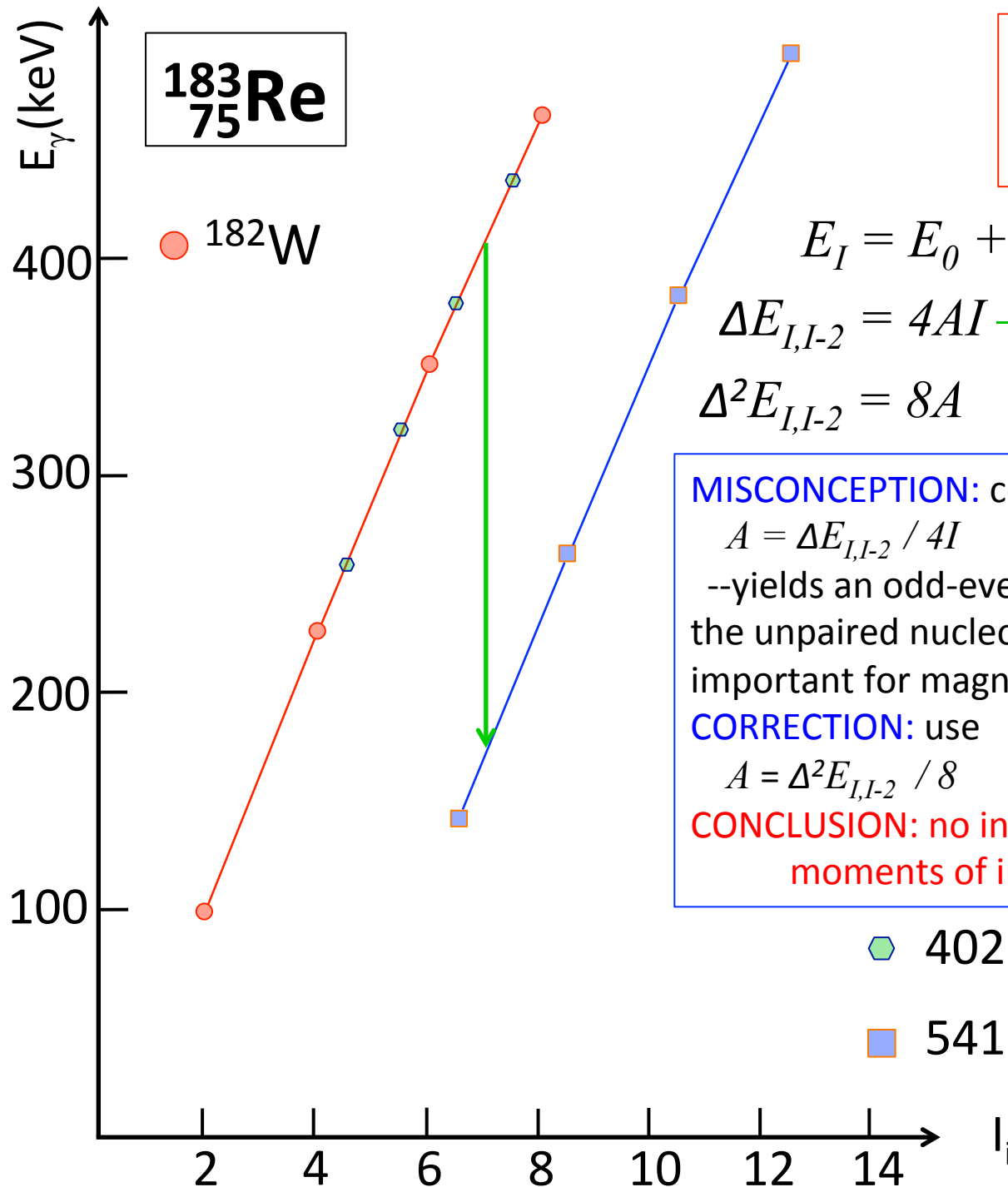
Witold Nazarewicz

Department of Physics and Astronomy, University of Tennessee, Knoxville,

Moments of inertia, Mol and pairing

--sample statements

- “The identical bands are found to be associated with up-sloping particle states, suggesting that the cause may be a cancellation between pairing and deformation decreases.”
--PRL 69 3448 (1992).
- “...BCS [pairing] theory can qualitatively reproduce the experimental large fluctuations in the Mol which is helpful to understand the appearance of the normally deformed identical bands in an odd-A nucleus and its even-even neighbors...”
—PR C63 047306 (2001).
- “In particular, the reduction of the BCS pairing correlations due to the blocking of one and two orbitals implies large changes (up to 30%) in the moments of inertia and cannot be reconciled with these [identical band] systematics.”
--Annu. Rev. Nucl. Part. Sci. 45 485 (1995).



E_γ vs. I_{initial} as a rotation alignment plot

$$E_I = E_0 + AI(I+1) - 2A\langle I \cdot J \rangle$$

$$\Delta E_{I,I-2} = 4AI - \text{const.}$$

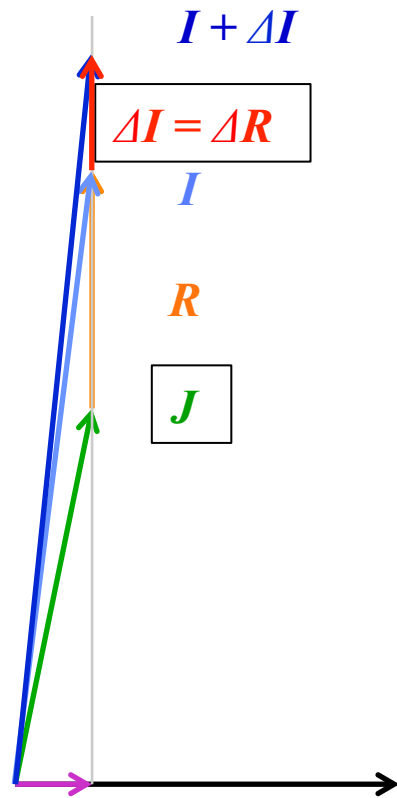
$$\Delta^2 E_{I,I-2} = 8A$$

MISCONCEPTION: commonly, use
 $A = \Delta E_{I,I-2} / 4I$
 --yields an odd-even staggering and conclusion that the unpaired nucleon blocks pairing which is important for magnitudes of moments of inertia.
CORRECTION: use
 $A = \Delta^2 E_{I,I-2} / 8$
CONCLUSION: no influence of pairing on moments of inertia

◈ $402\uparrow$ $d_{5/2}$ *const.* = 0 keV

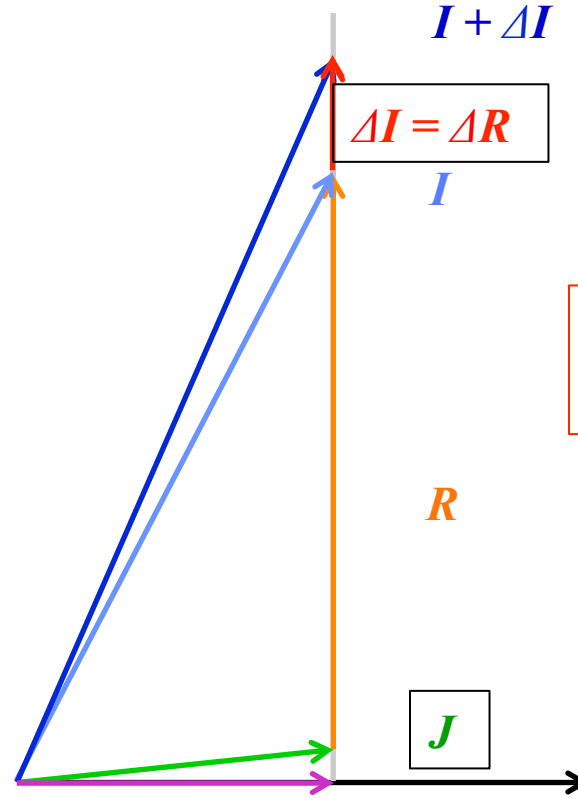
◩ $541\downarrow$ $h_{9/2}$ *const.* = -230 keV

“Coriolis” contribution to energy differences



$$K = \Omega = 1/2$$

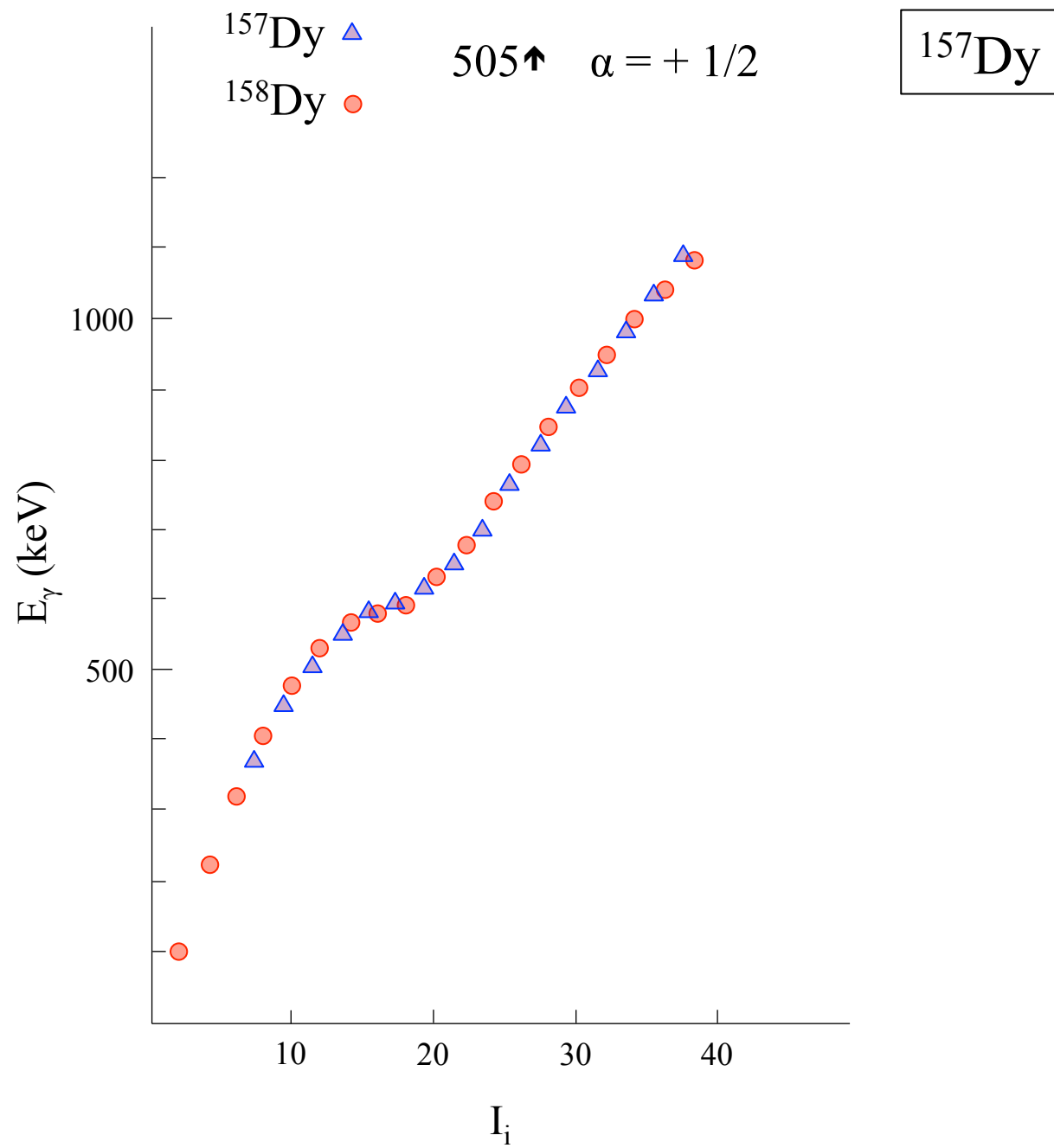
$$\langle \Delta I \cdot J \rangle \text{ large}$$



$$K = \Omega = \max$$

$$\langle \Delta I \cdot J \rangle \approx 0$$

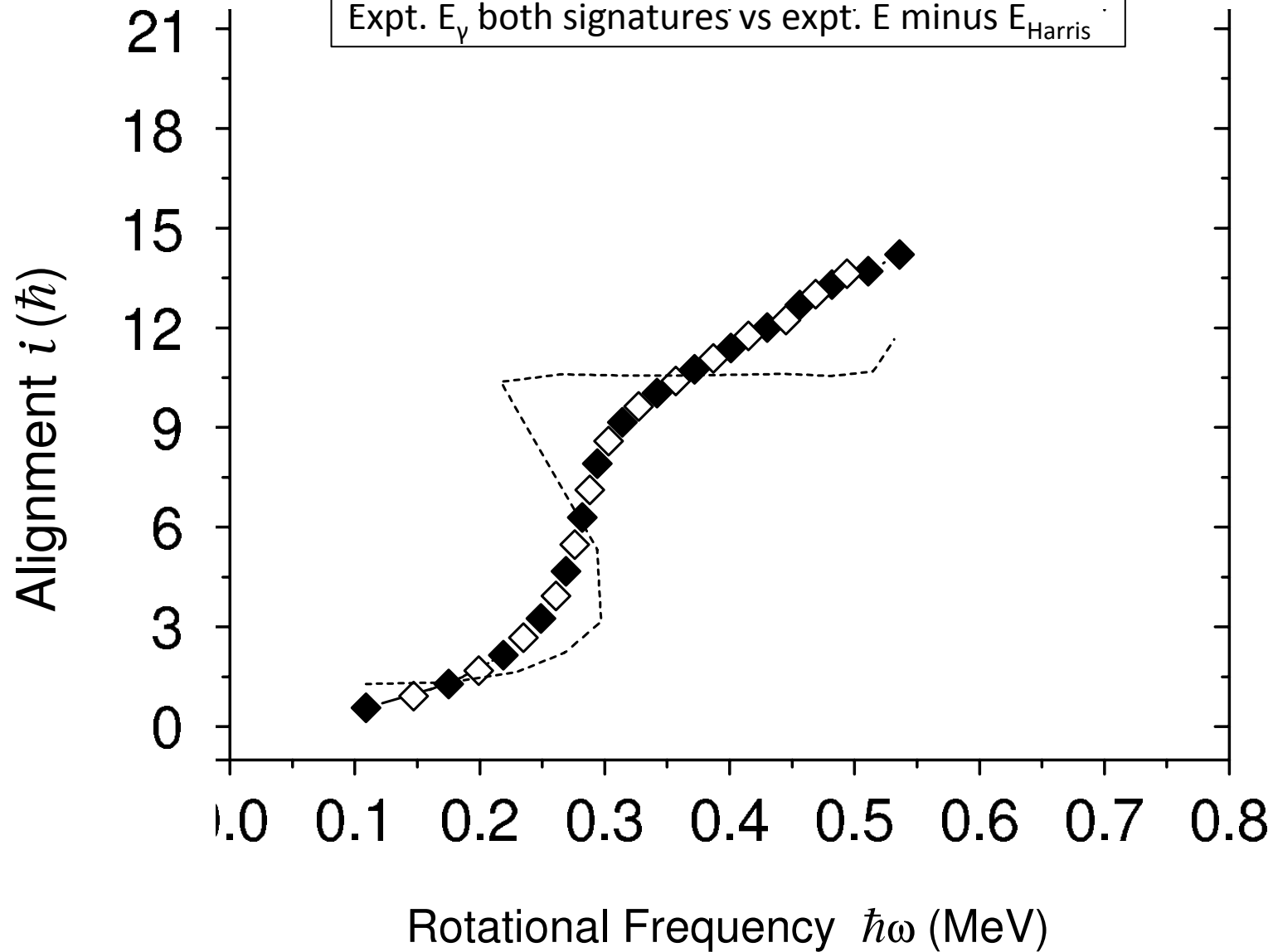
A vectorial difference effect



^{157}Dy 505

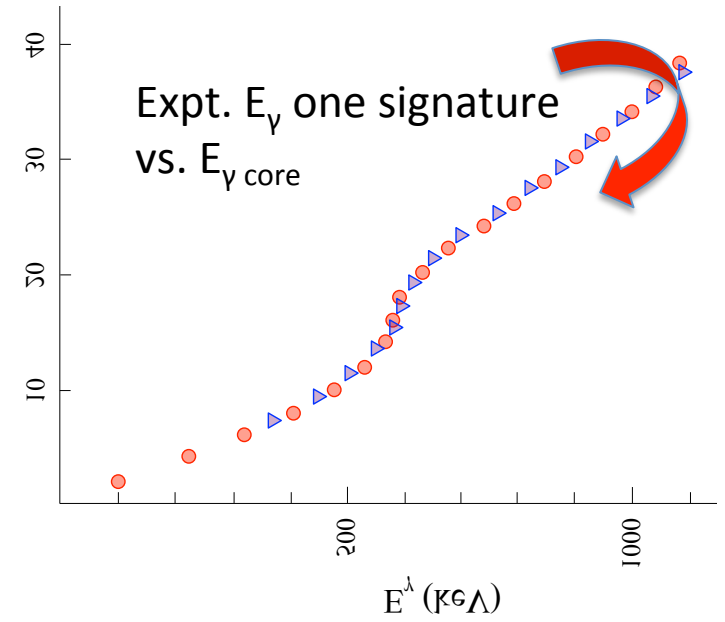
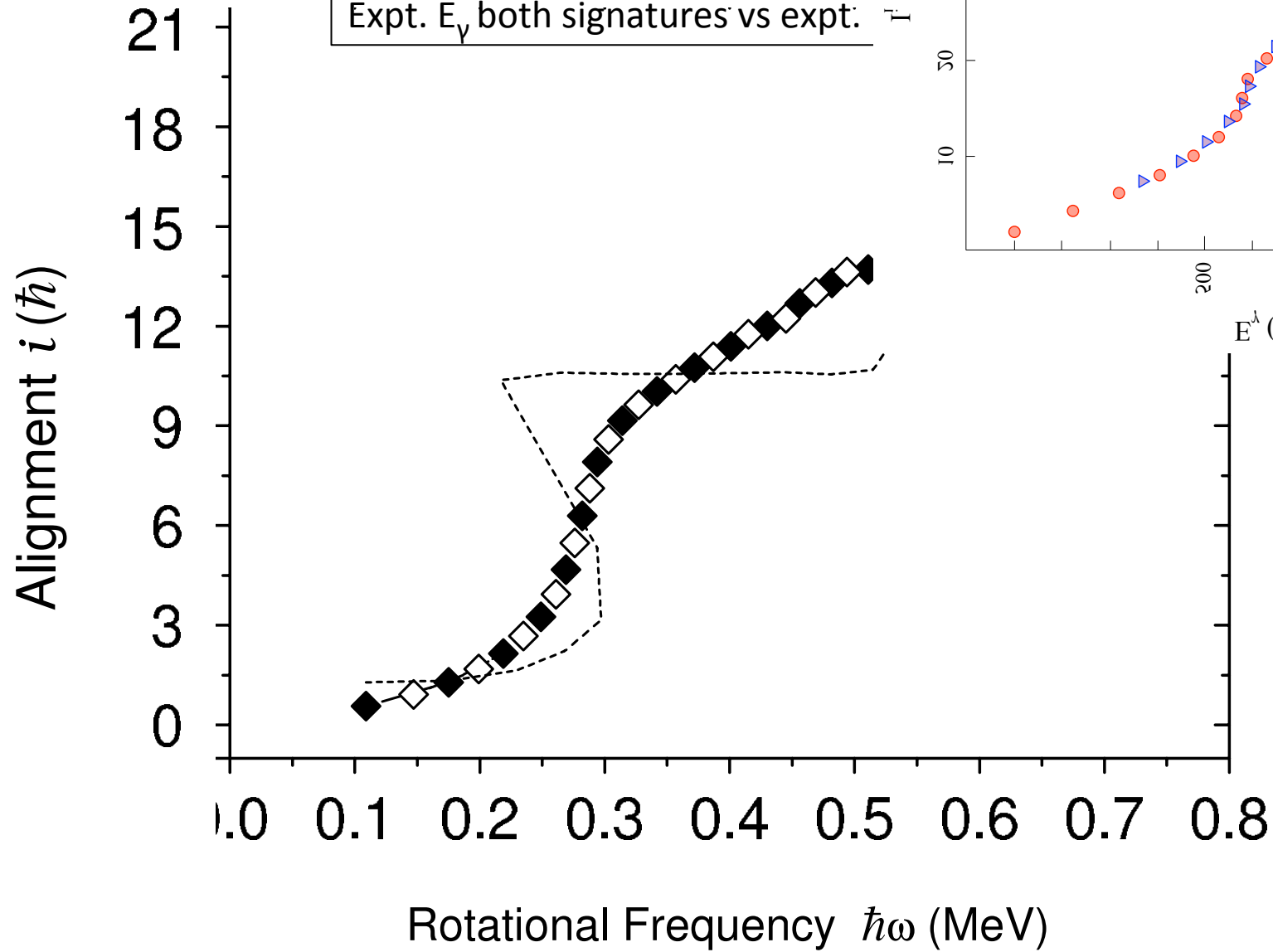
Pipidis, Riley et al.,
PR C72 064307 2005

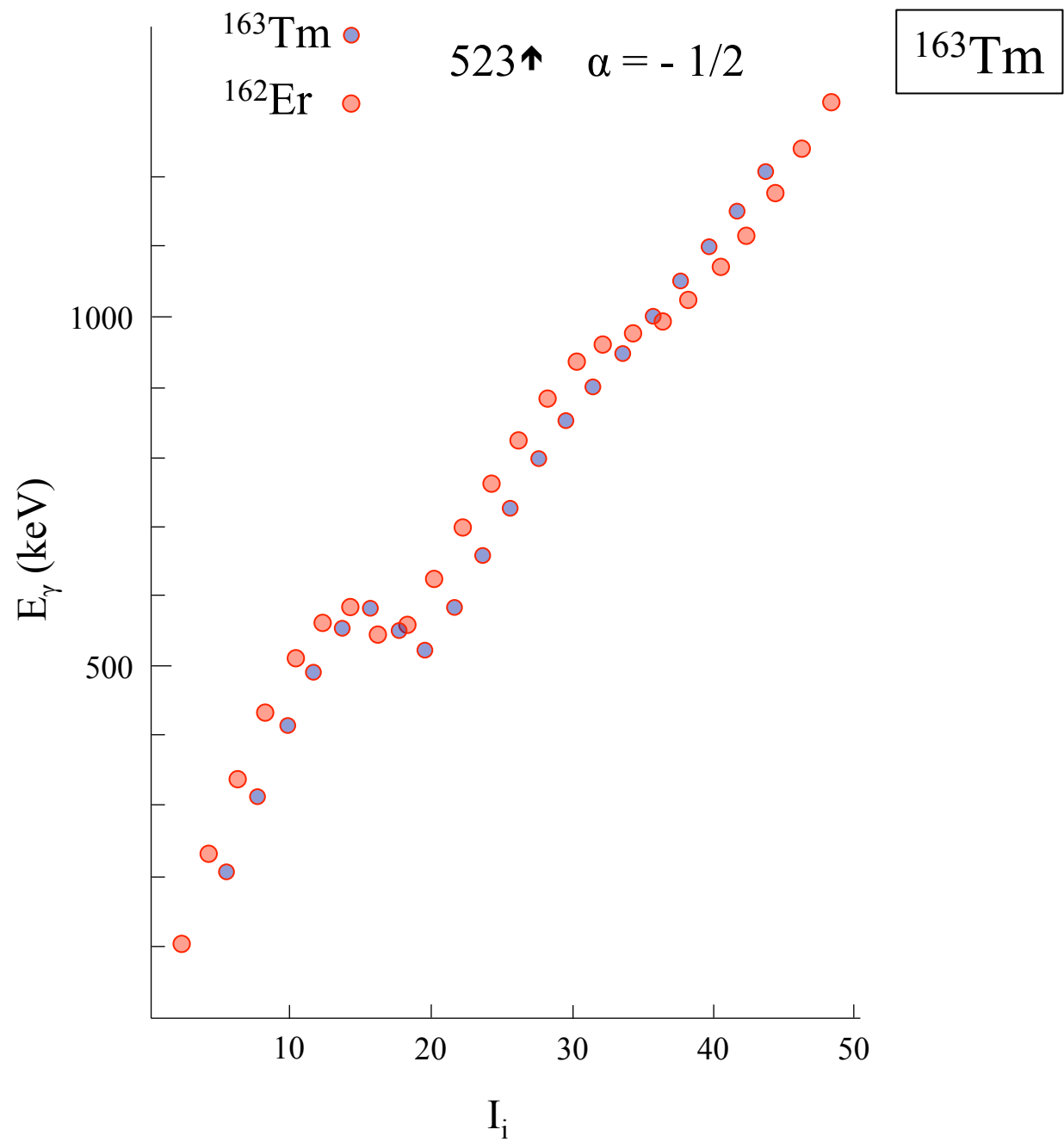
Expt. E_γ both signatures vs expt. E minus E_{Harris}

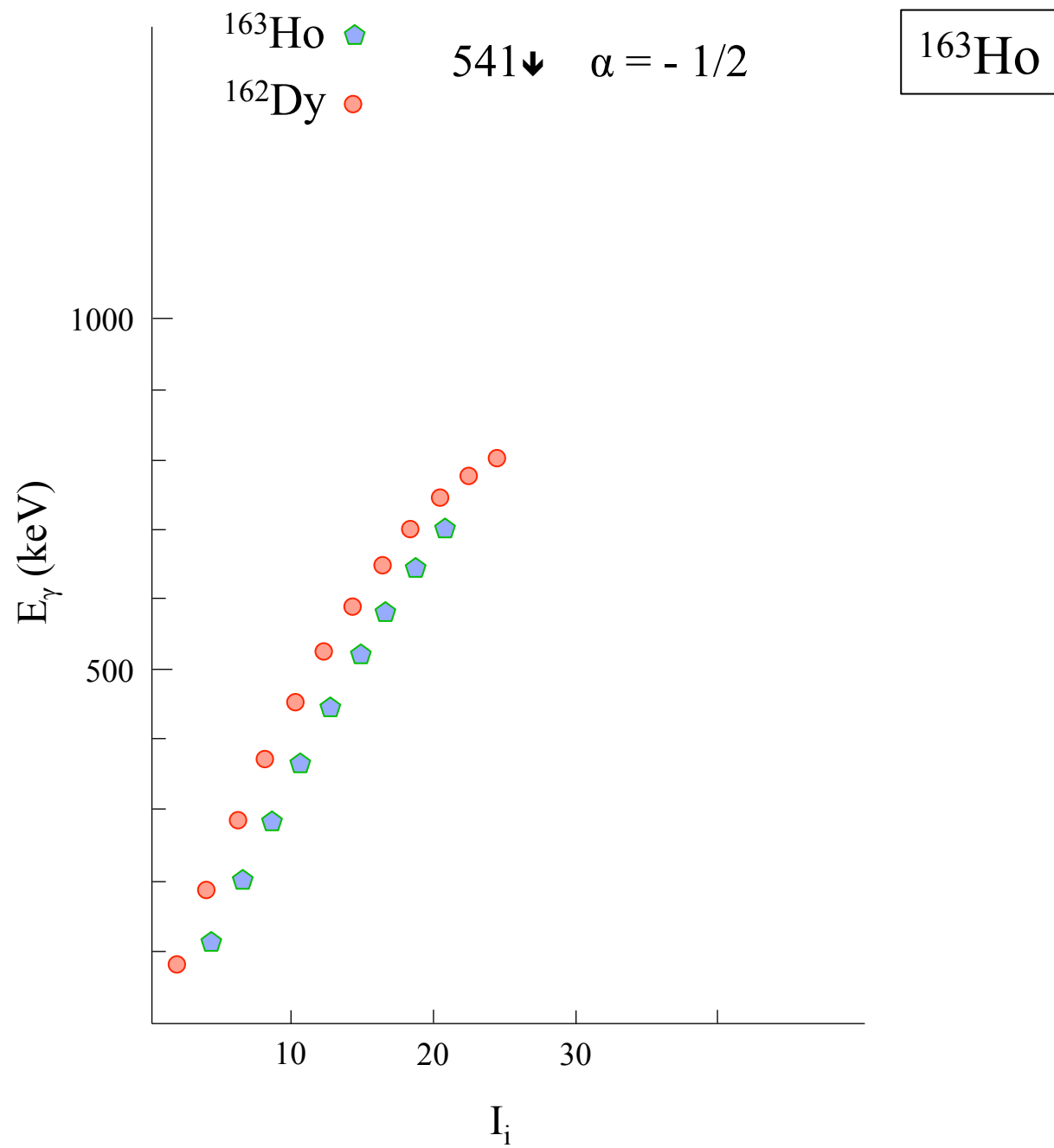


^{157}Dy 505

Pipidis, Riley et al.,
PR C72 064307 2005
Expt. E_γ both signatures vs expt. E_γ







CONCLUSIONS

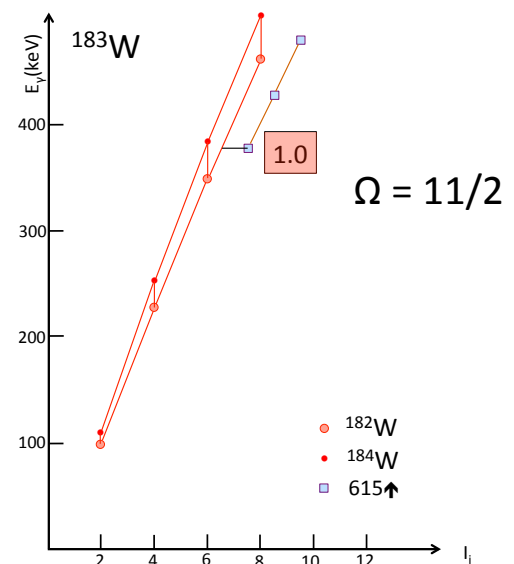
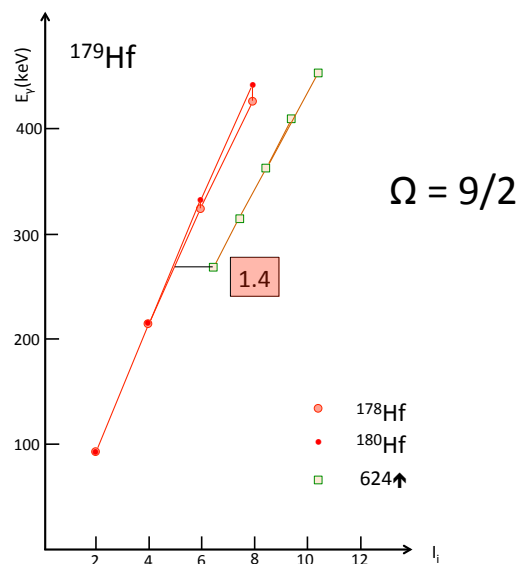
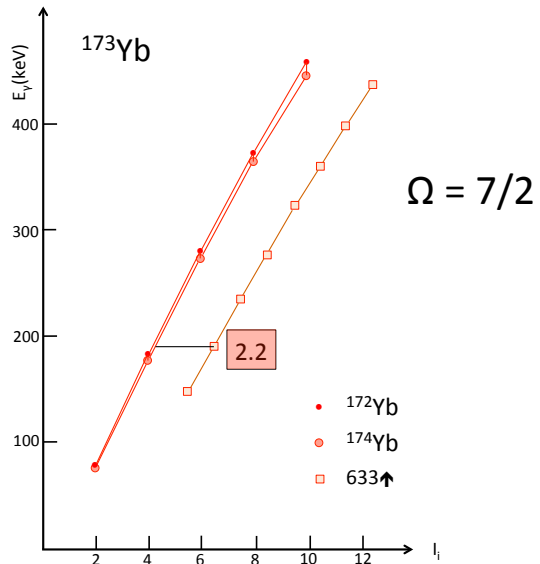
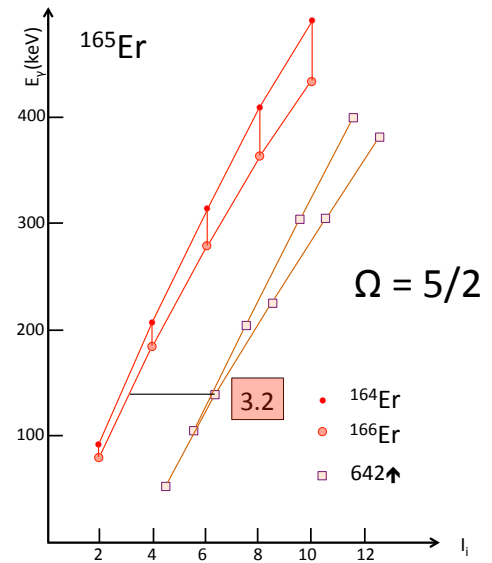
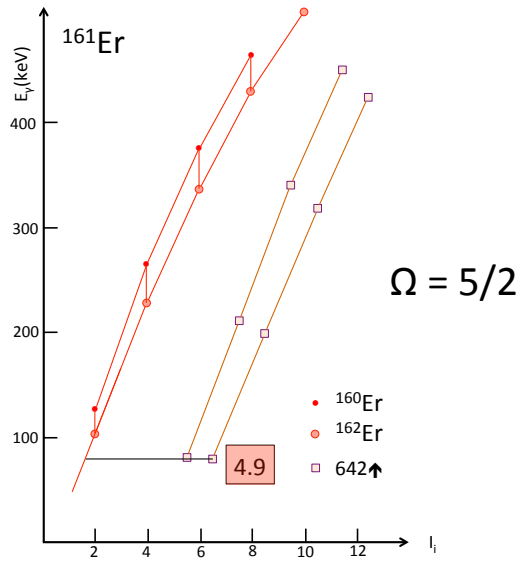
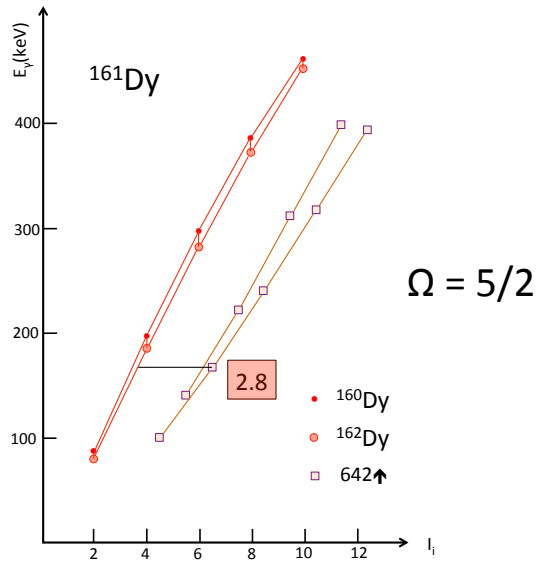
- The rotational-particle coupling term produces effects that have been overlooked.
- The moments of inertia in an odd nucleus are essentially the same as those in the neighboring even nucleus with the greatest deformation.
- There is no odd-particle “blocking” effect of moments of inertia.
- There is no “deformation-driving effect” of unpaired nucleons.
- A reassessment of our understanding of the magnitudes of moments of inertia is needed: there are models that attribute “rotational” energies to potential-energy sources, e.g., $Q \cdot Q$ interactions in $SU(3)$ coupling schemes.
- If “rotational” energy is partly potential, Coriolis effects are attenuated—this solves a long-standing problem.

SU(3) quasi-dynamical symmetry as an organizational mechanism for generating nuclear rotational motions

C. Bahri, D.J. Rowe

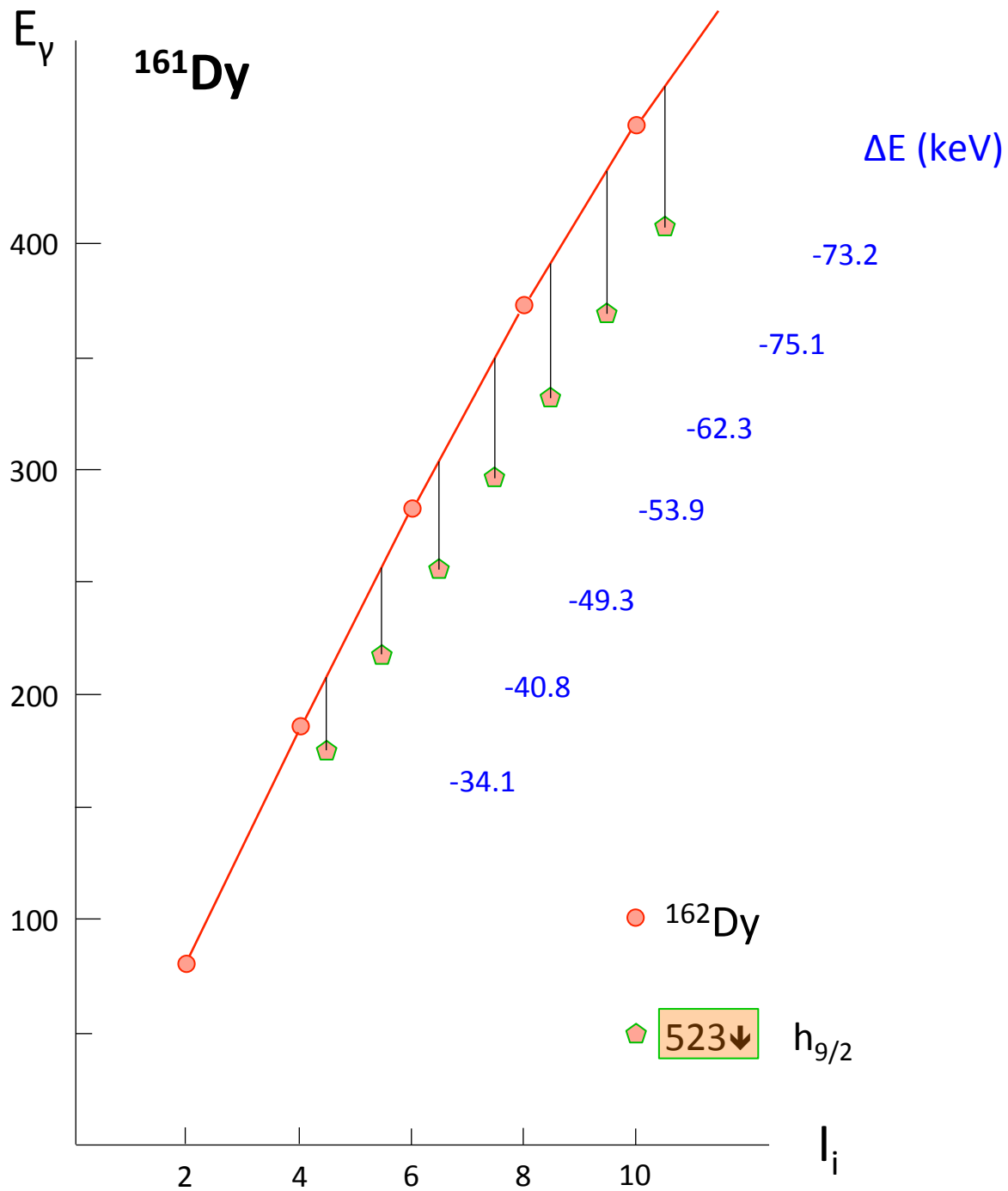
Department of Physics, University of Toronto, Toronto, Ontario M5S 1A7, Canada

“A particularly interesting challenge was to learn how a model, without pair correlations, could give correct moments of inertia when it is known that the cranking model is only successful when pairing correlations are included. The early calculations of Park et al. indicated that the dominant contribution to rotational energies came from the potential energy part of the Hamiltonian, thus calling into question the very concept of the moment of inertia as an inverse coefficient of the L^2 term in the kinetic energy. The results of the present calculation indicate that the inclusion of only stretched states, as in the calculation of Park et al.,[#] tends to exaggerate this effect. Nevertheless, it confirms that the dominant component of the rotational energies comes from the potential energy; for the self-consistent value of x only about 20% of the rotational energy comes from the kinetic energy in the present calculation.” #--1984

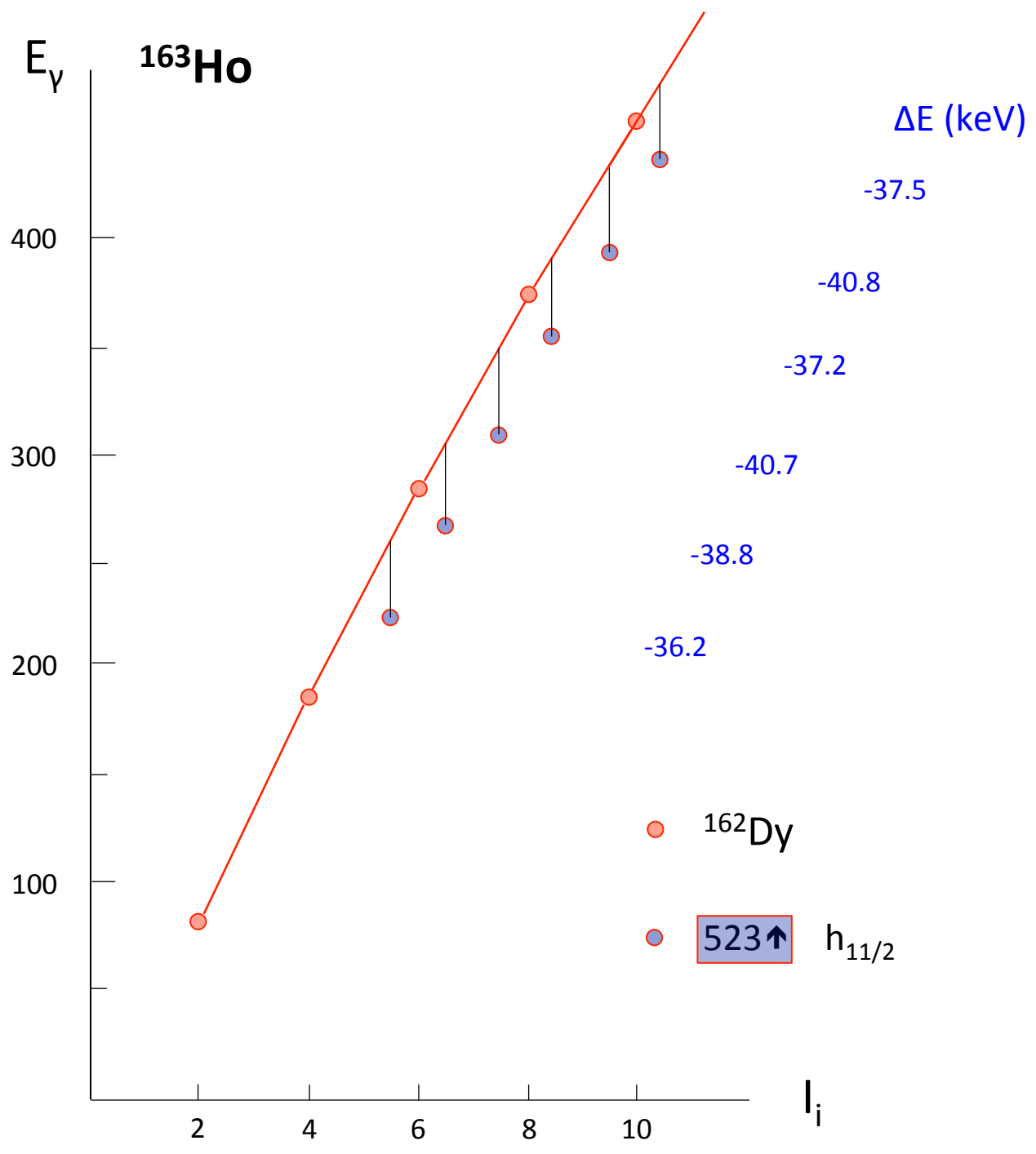


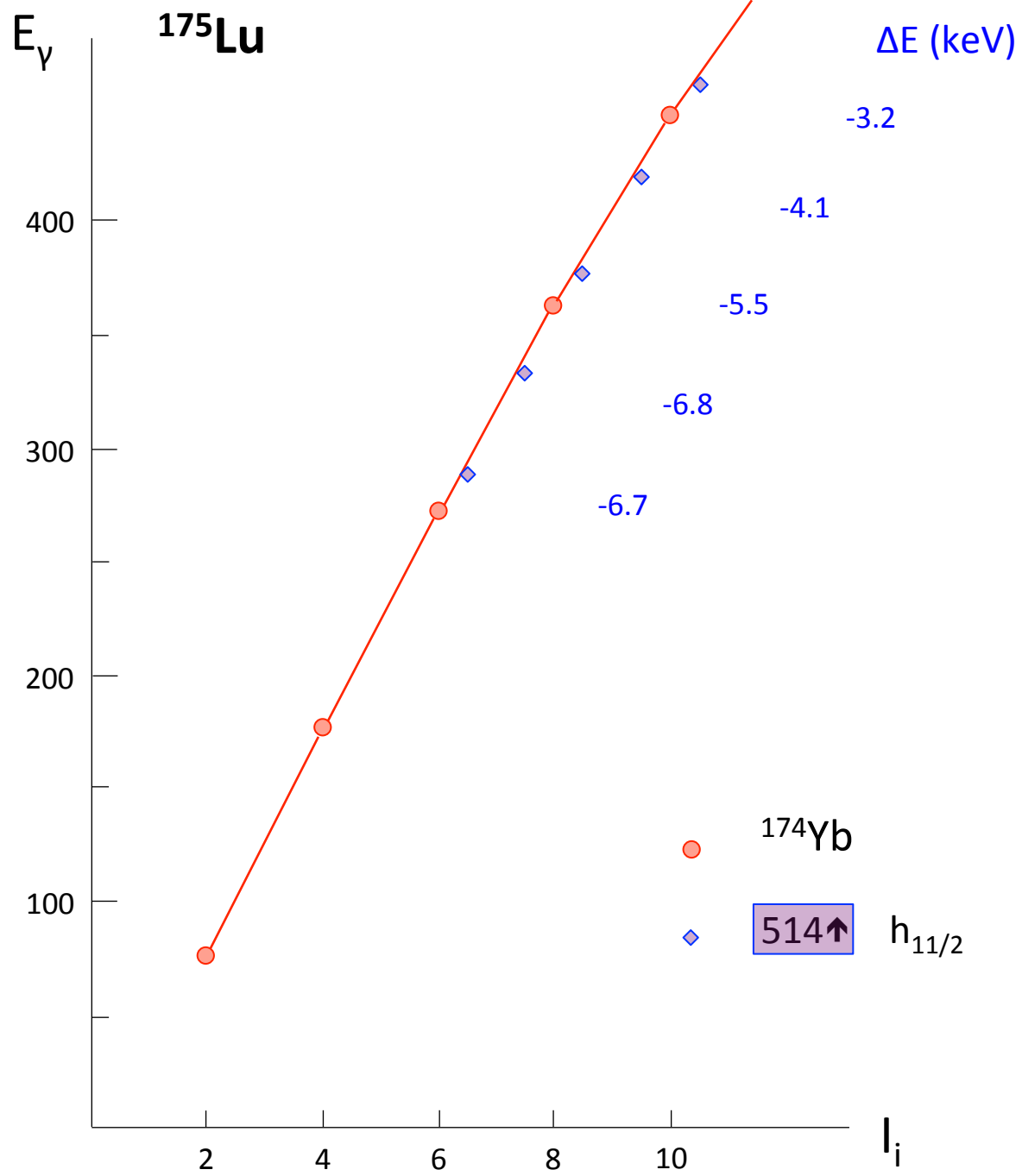
$i_{13/2}$

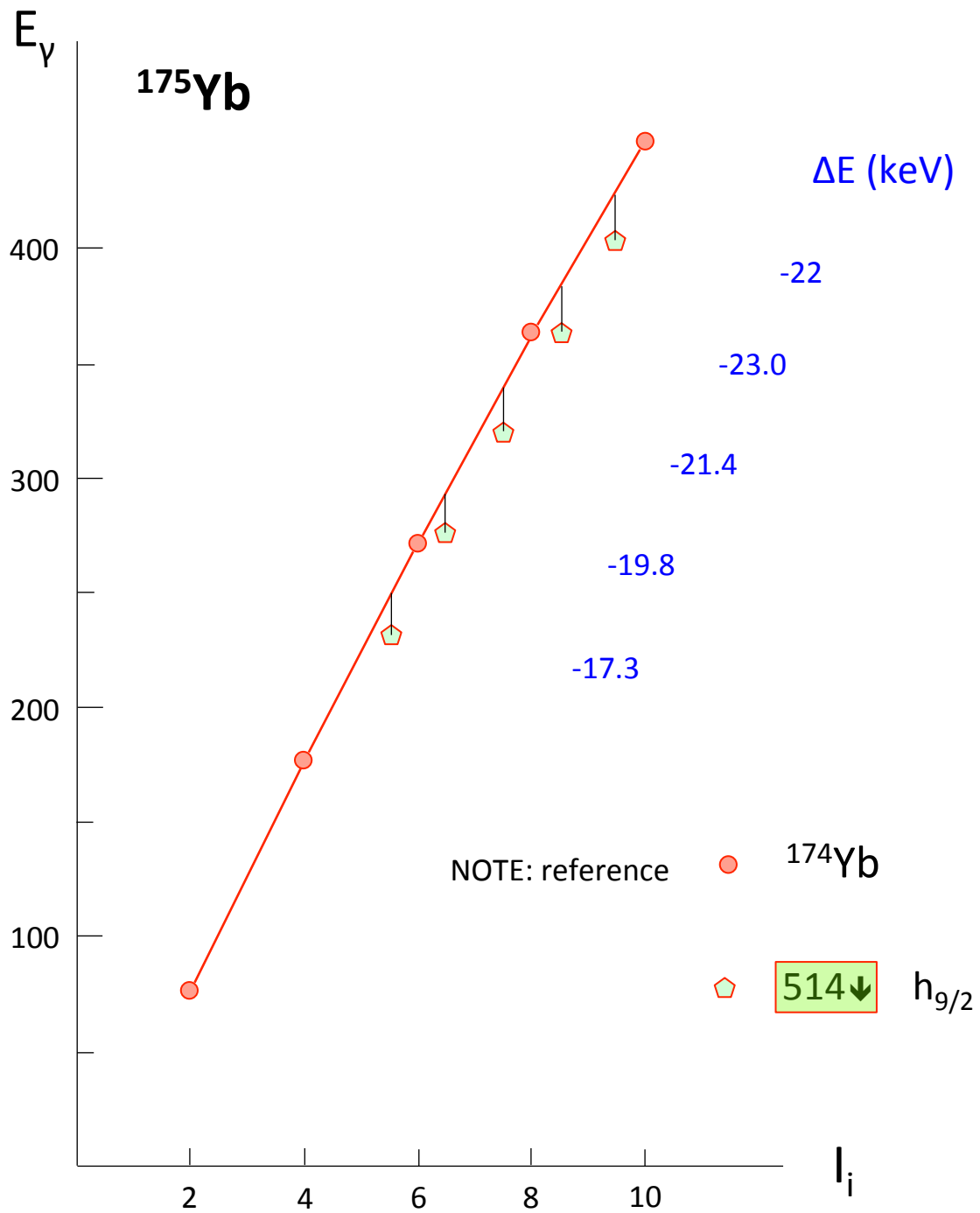
alignment decreases with:
 increasing deformation
 increasing Ω

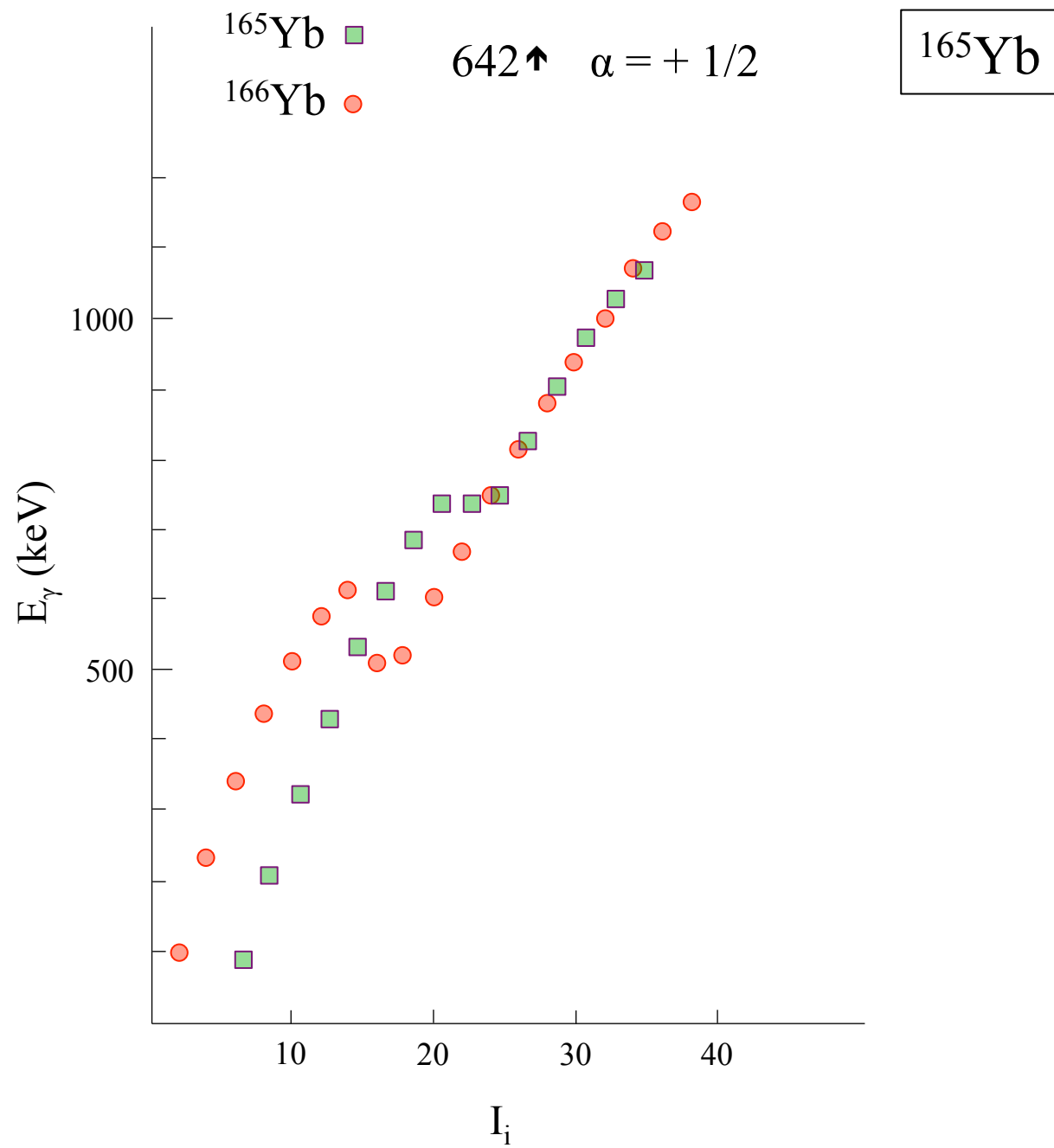


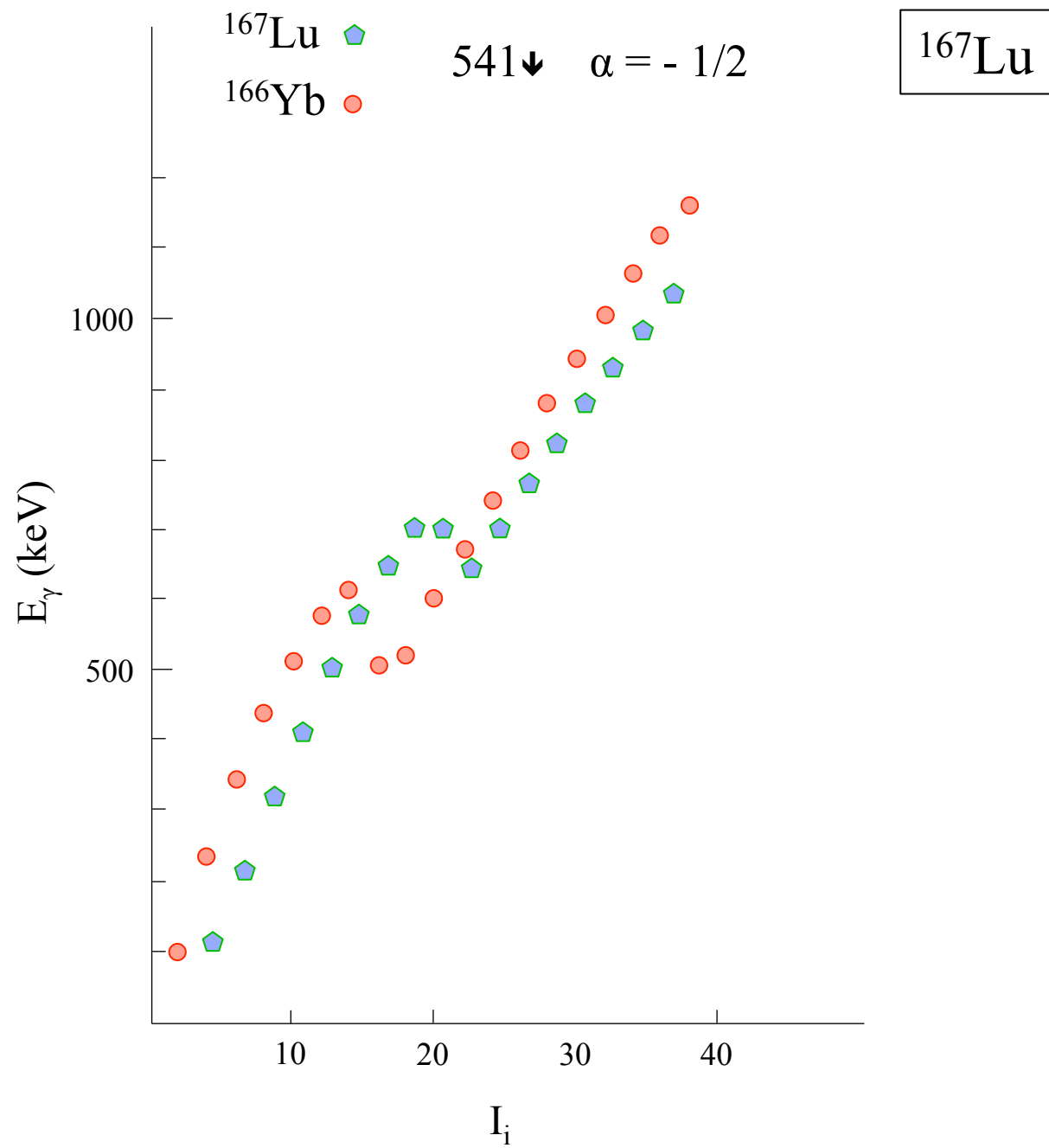
$\Delta I \cdot J$ increasing

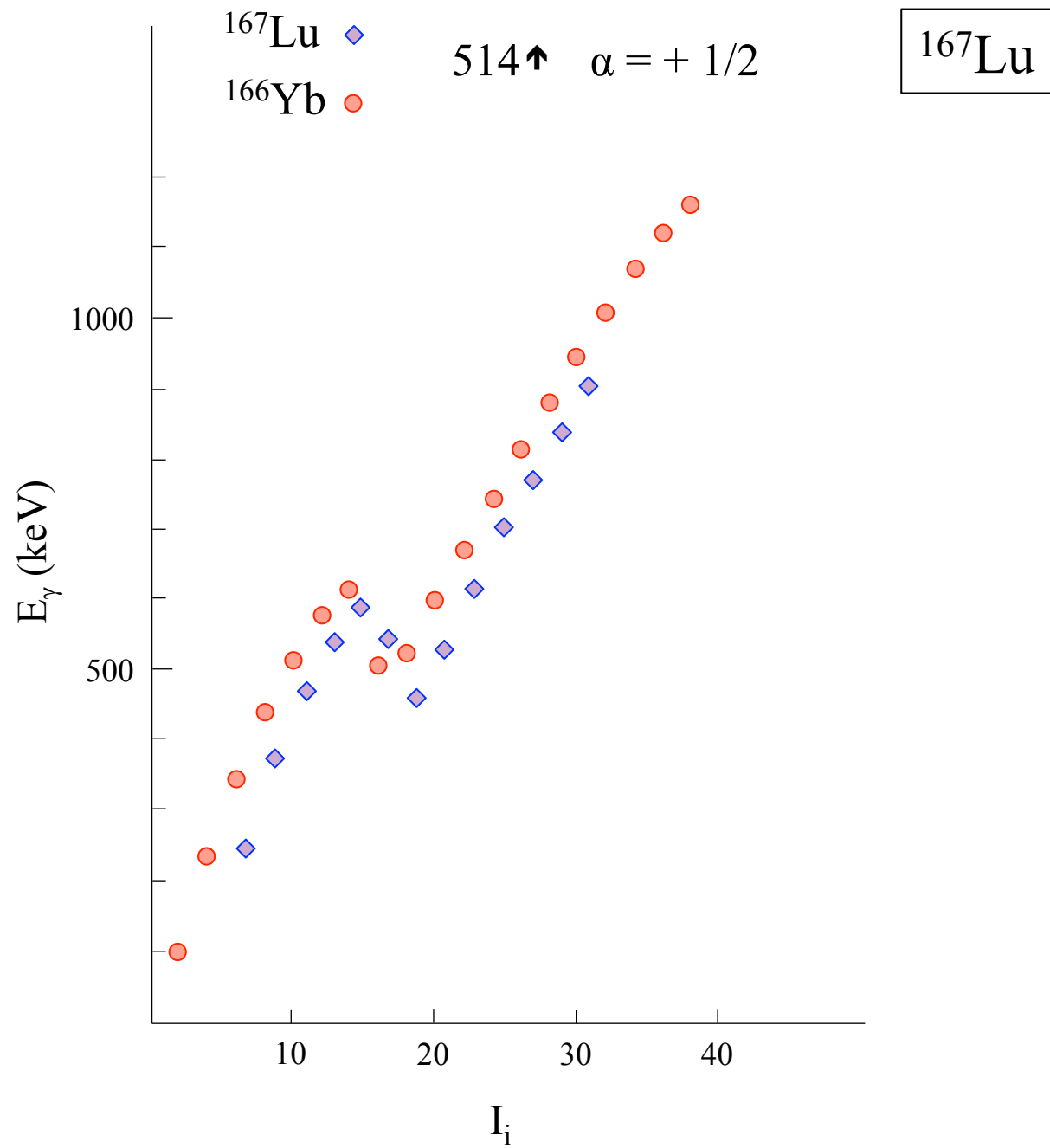


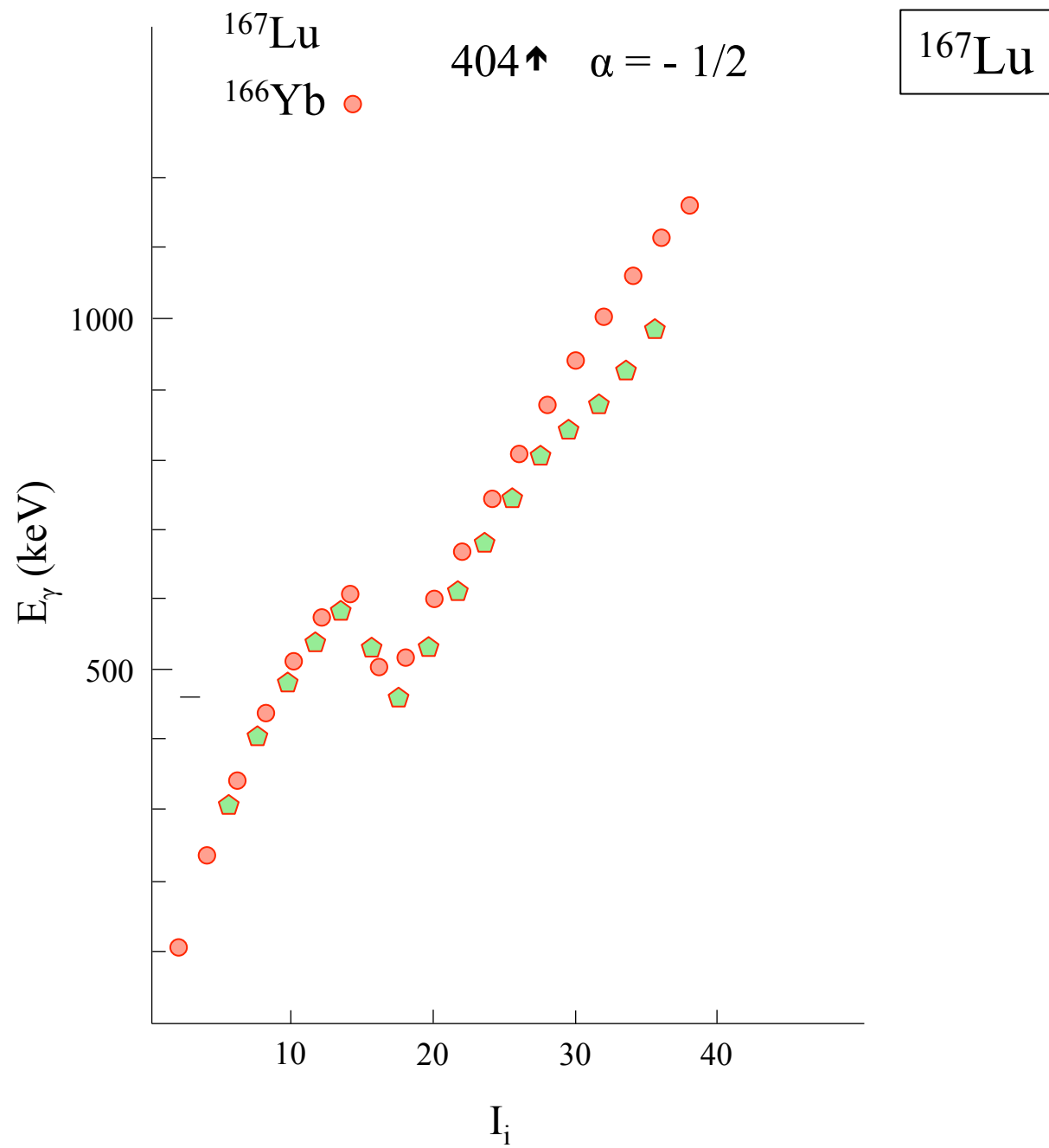


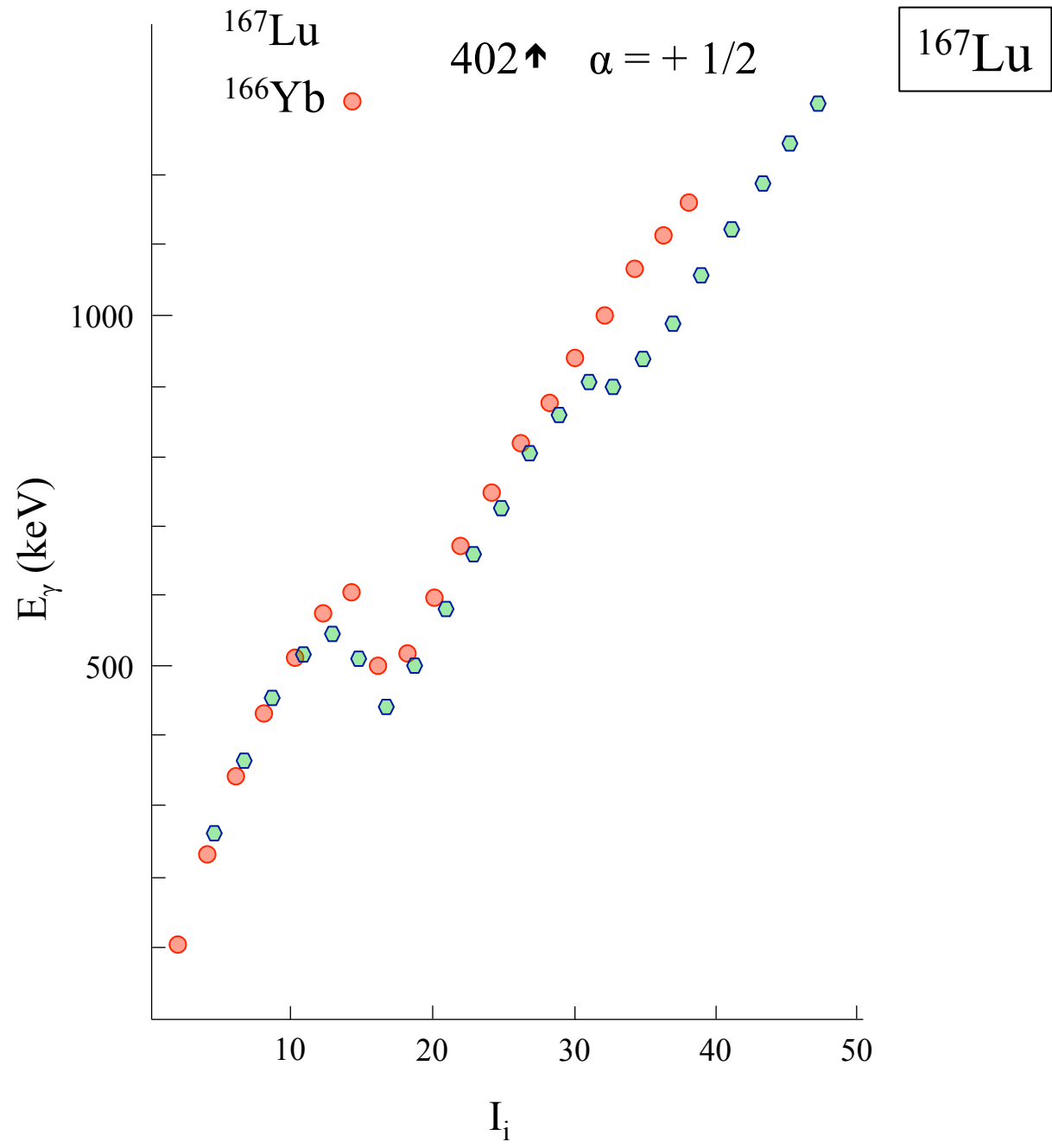




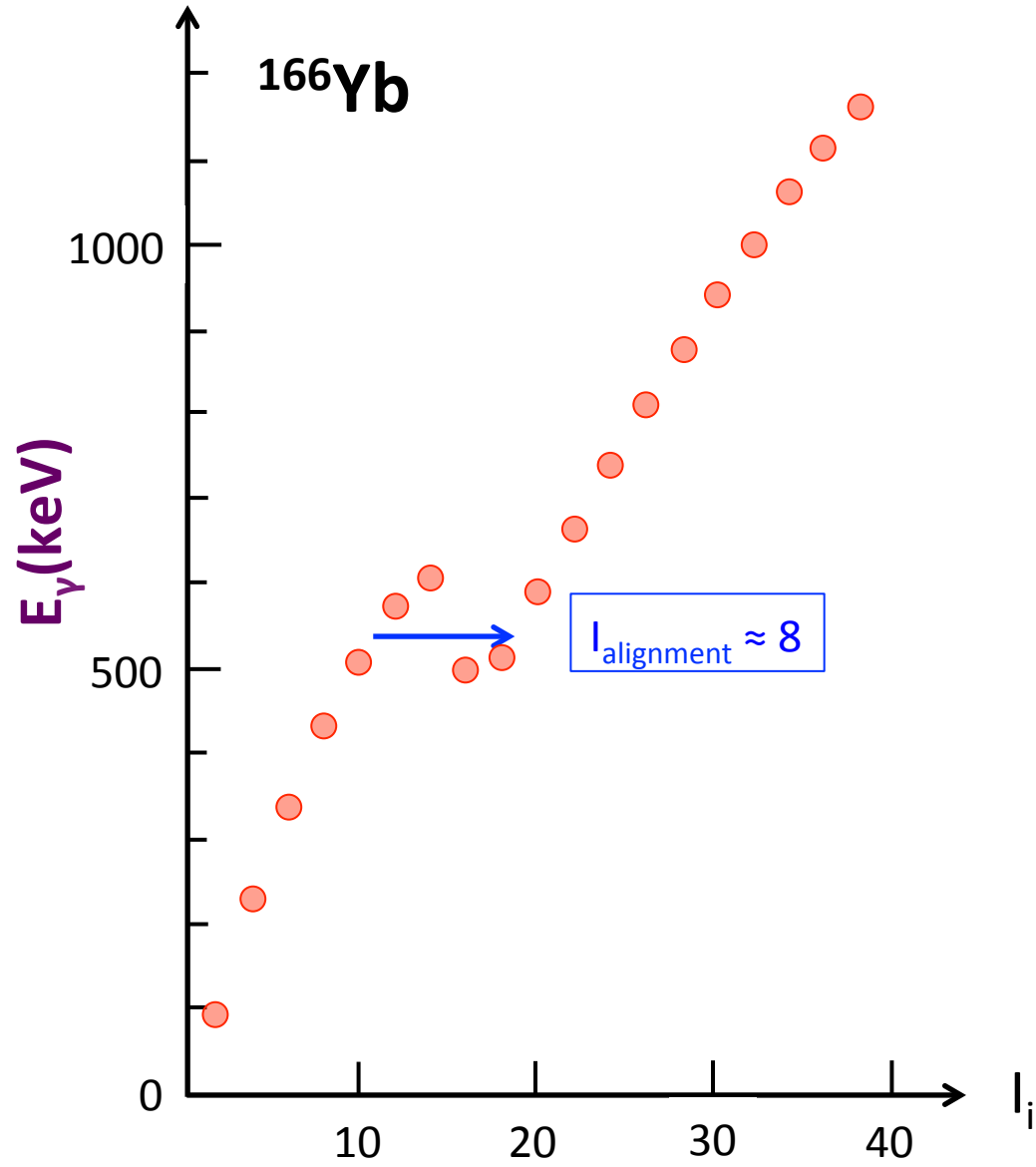




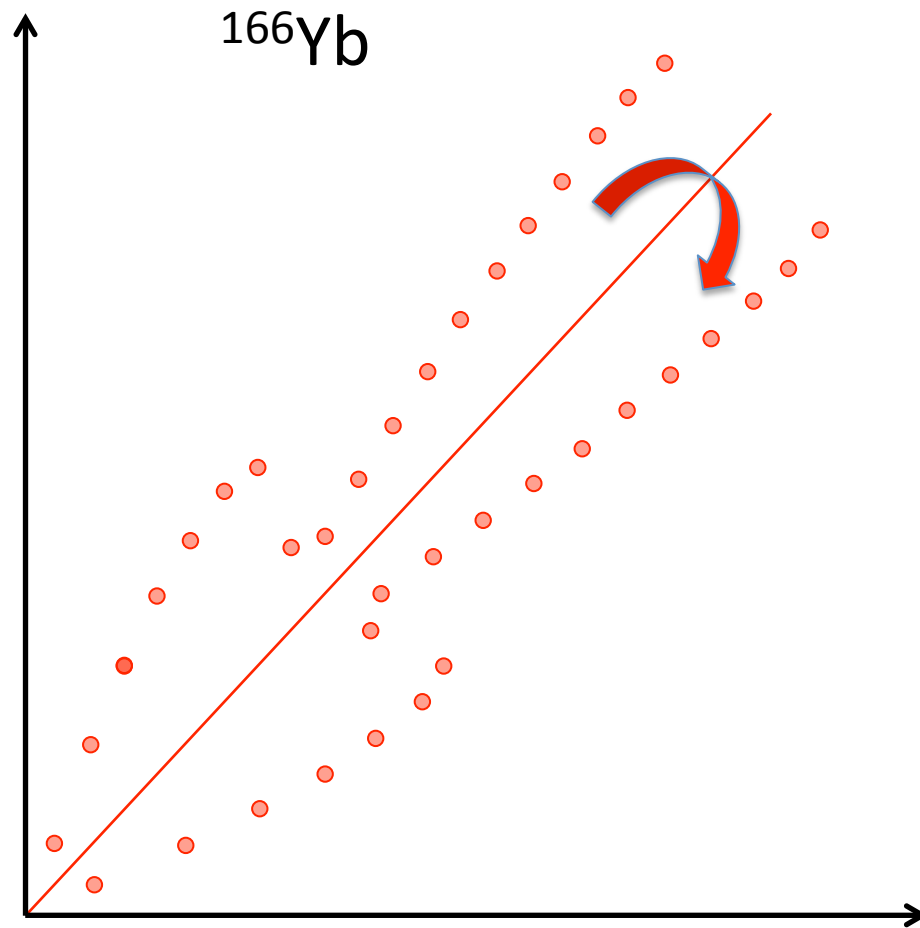




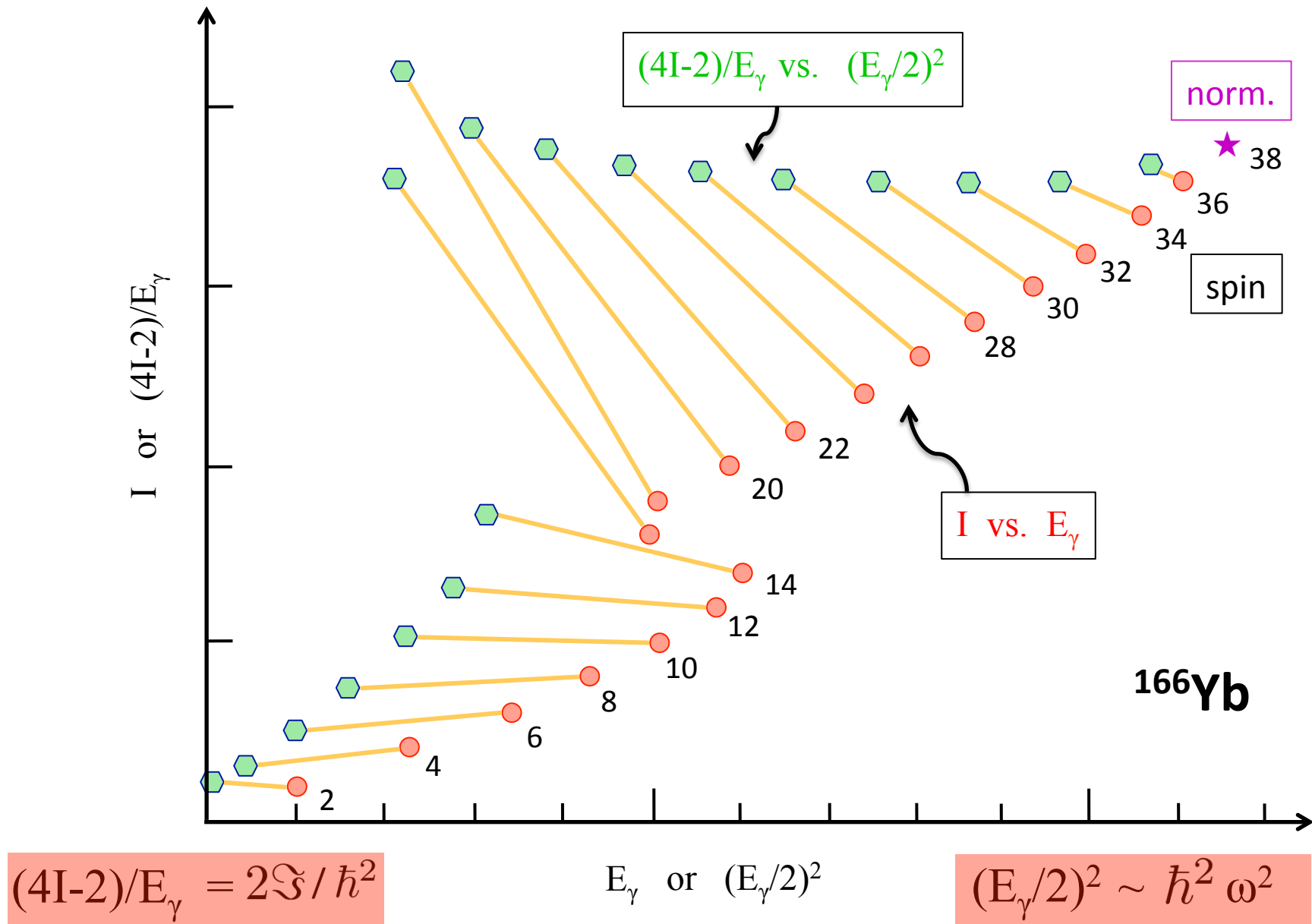
E_γ (keV) vs. I_i plots: "backbending"



Relationships for plotting spin vs. γ -ray energy:
transformation from E_γ vs. I_i plot to I_i vs. E_γ plot

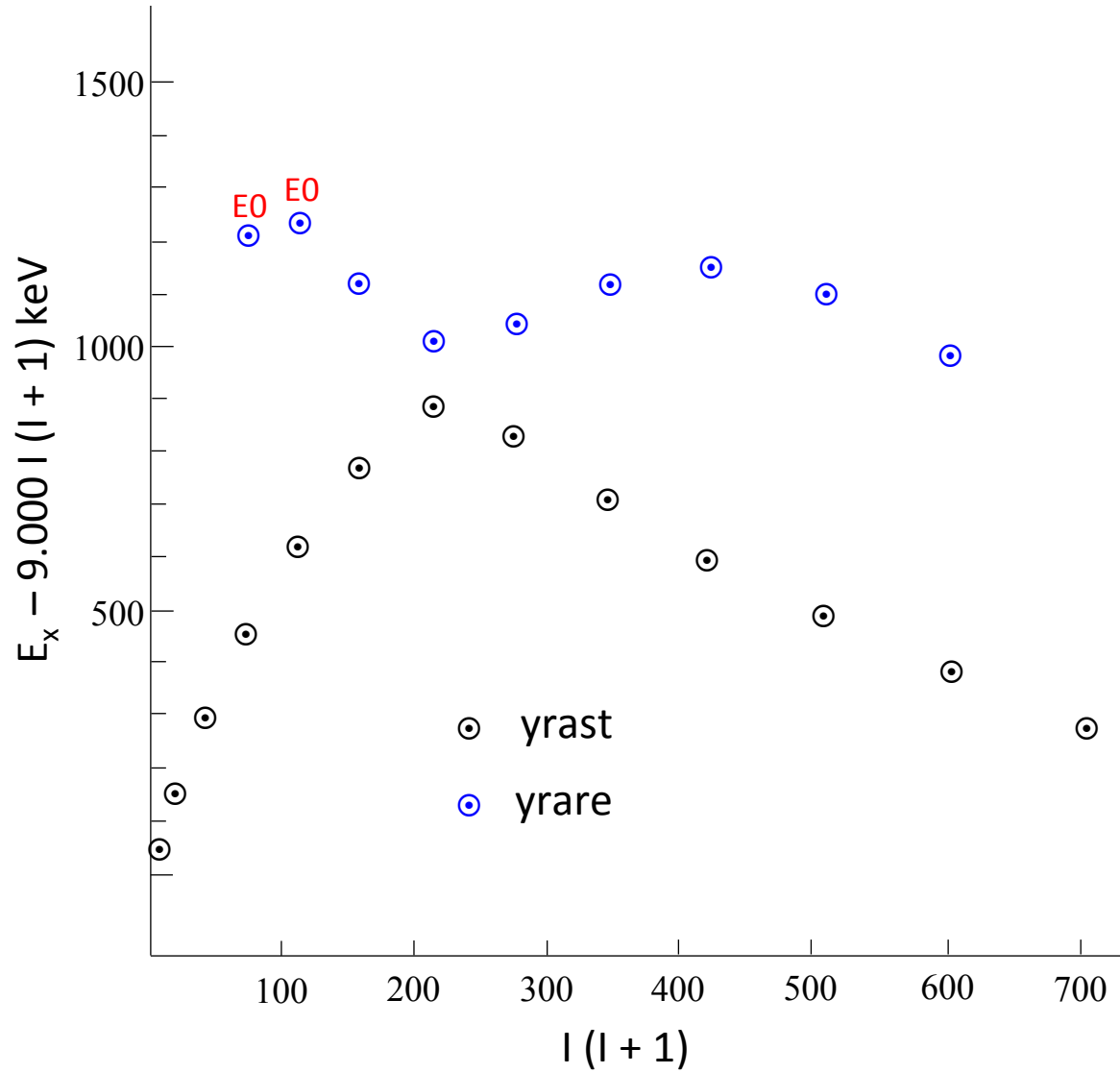


Relationships for plotting spin vs. γ -ray energy



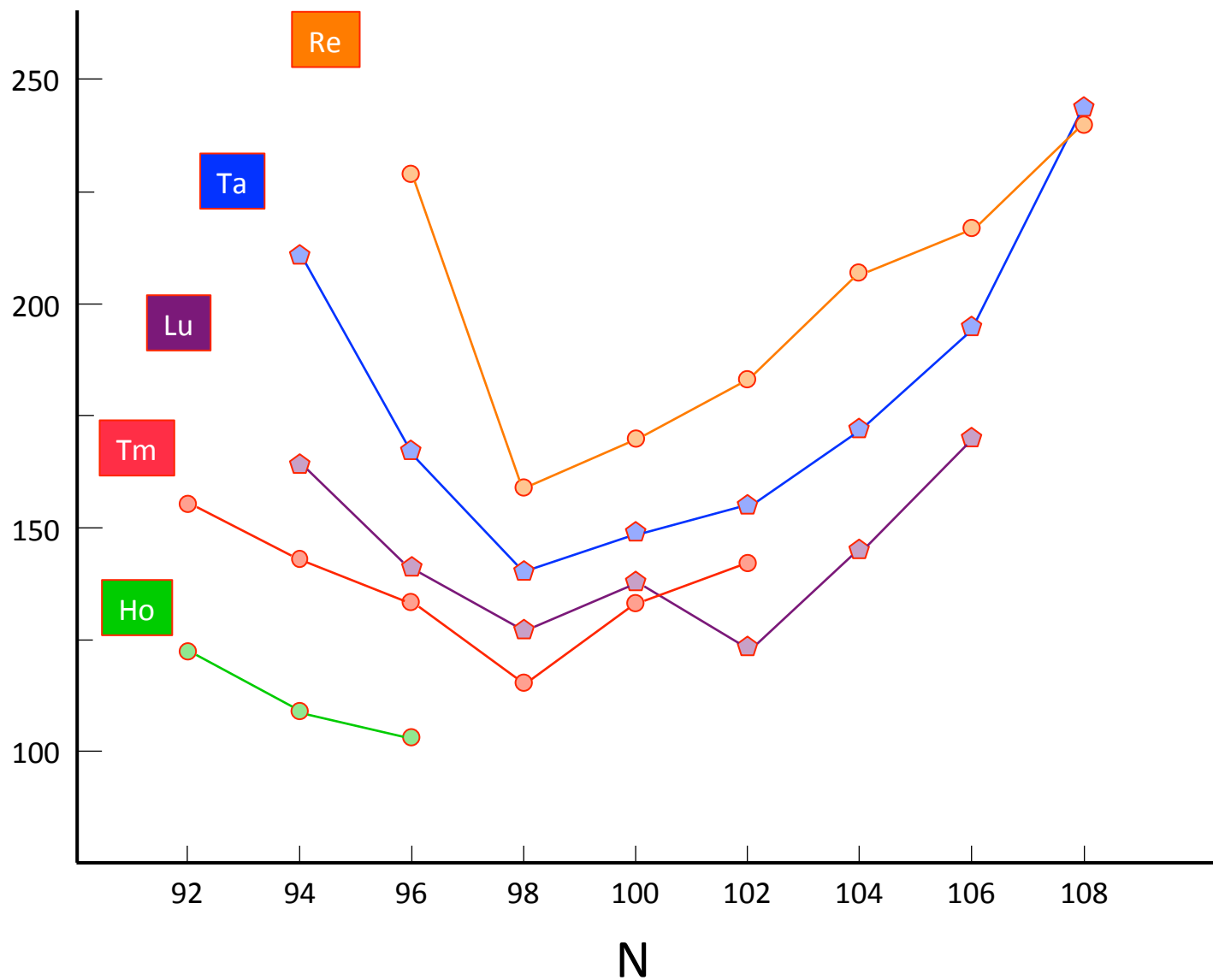
^{166}Yb

E0's: C. Fields et al., PL **B130** 157 ; NP **A431** 473



Alignment energy (lowering) for $541\downarrow$ band in rare-earth nuclei

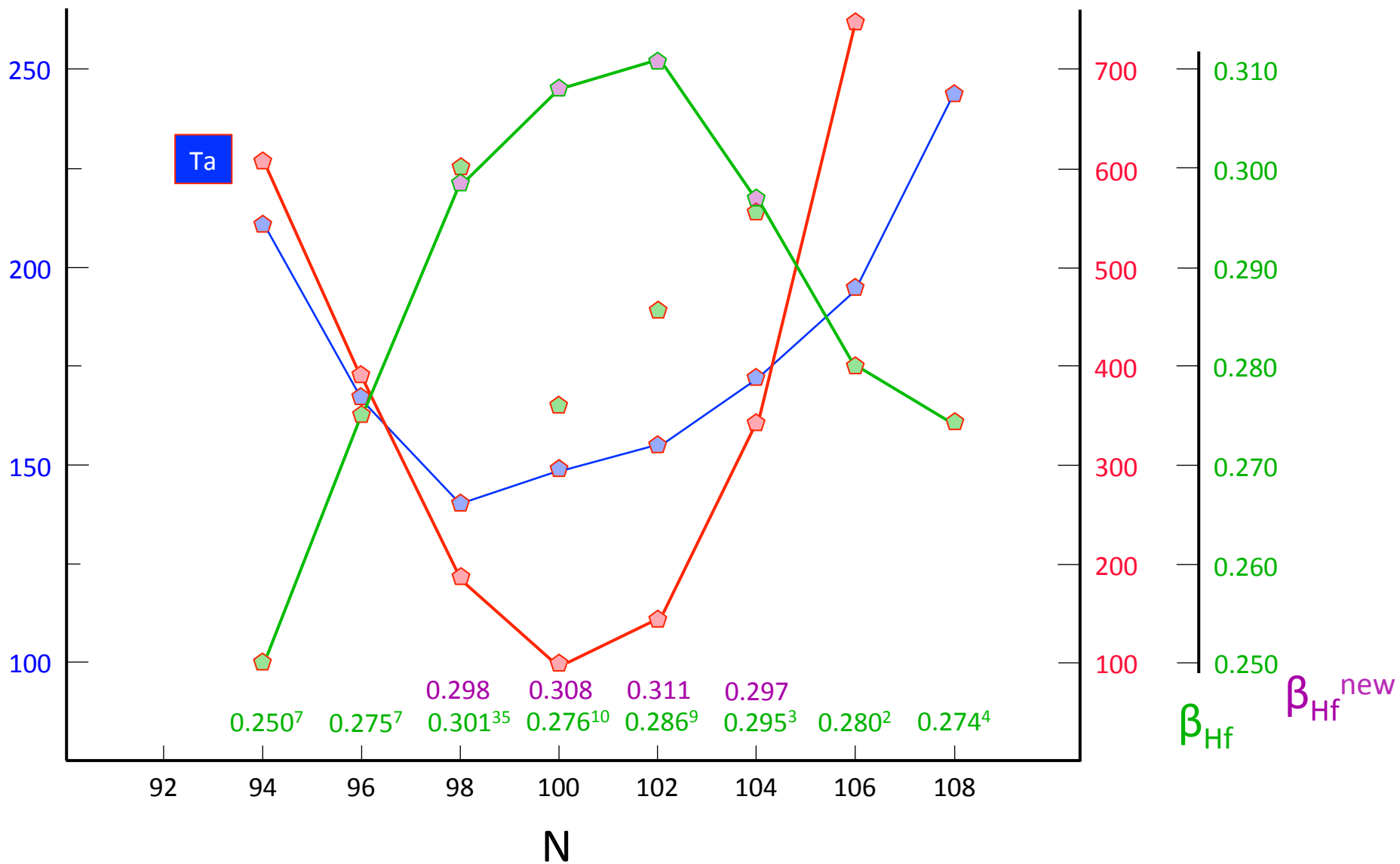
E_{align} (keV)



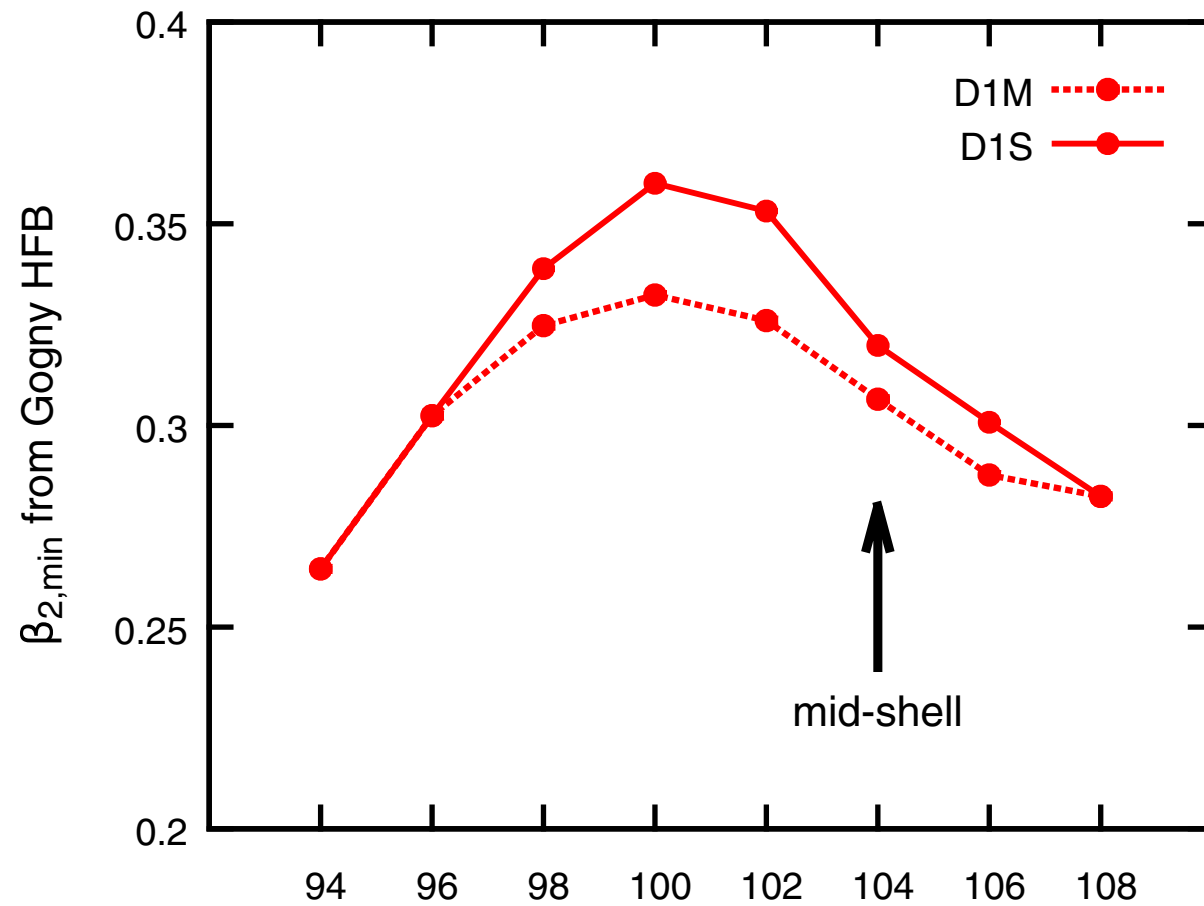
alignment energy, band-head energy, deformation: 541↓ Ta

E_{align} (keV)

E_{541} (keV)



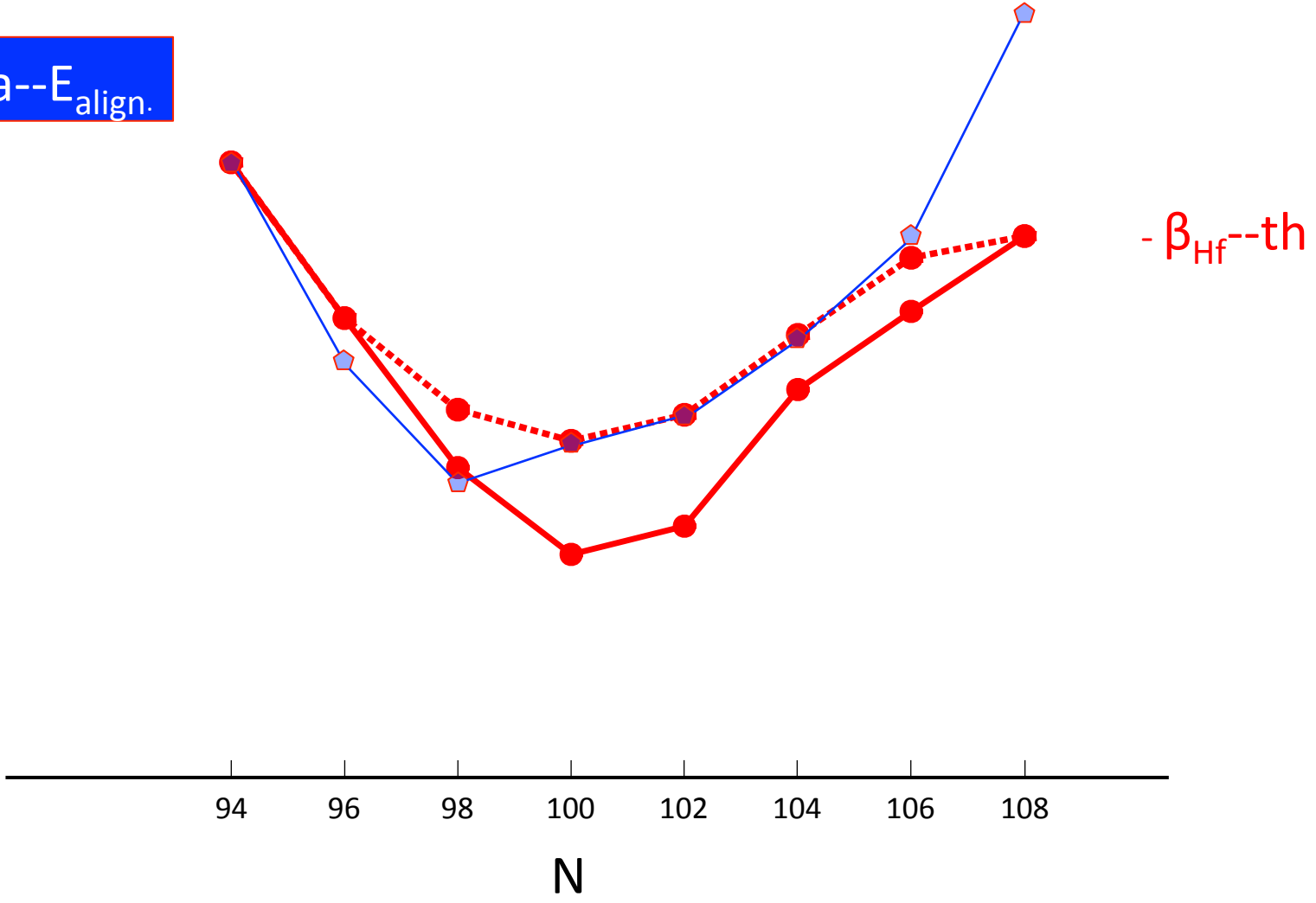
Rudigier et al., PR C91 044301 2015
And see Smith et al., PR C87 044317 2013



Ram.

170	185	181 ⁶
172	156	194 ⁶
174	168	199 ⁶
176	179	182 ⁷

Ta--E_{align.}

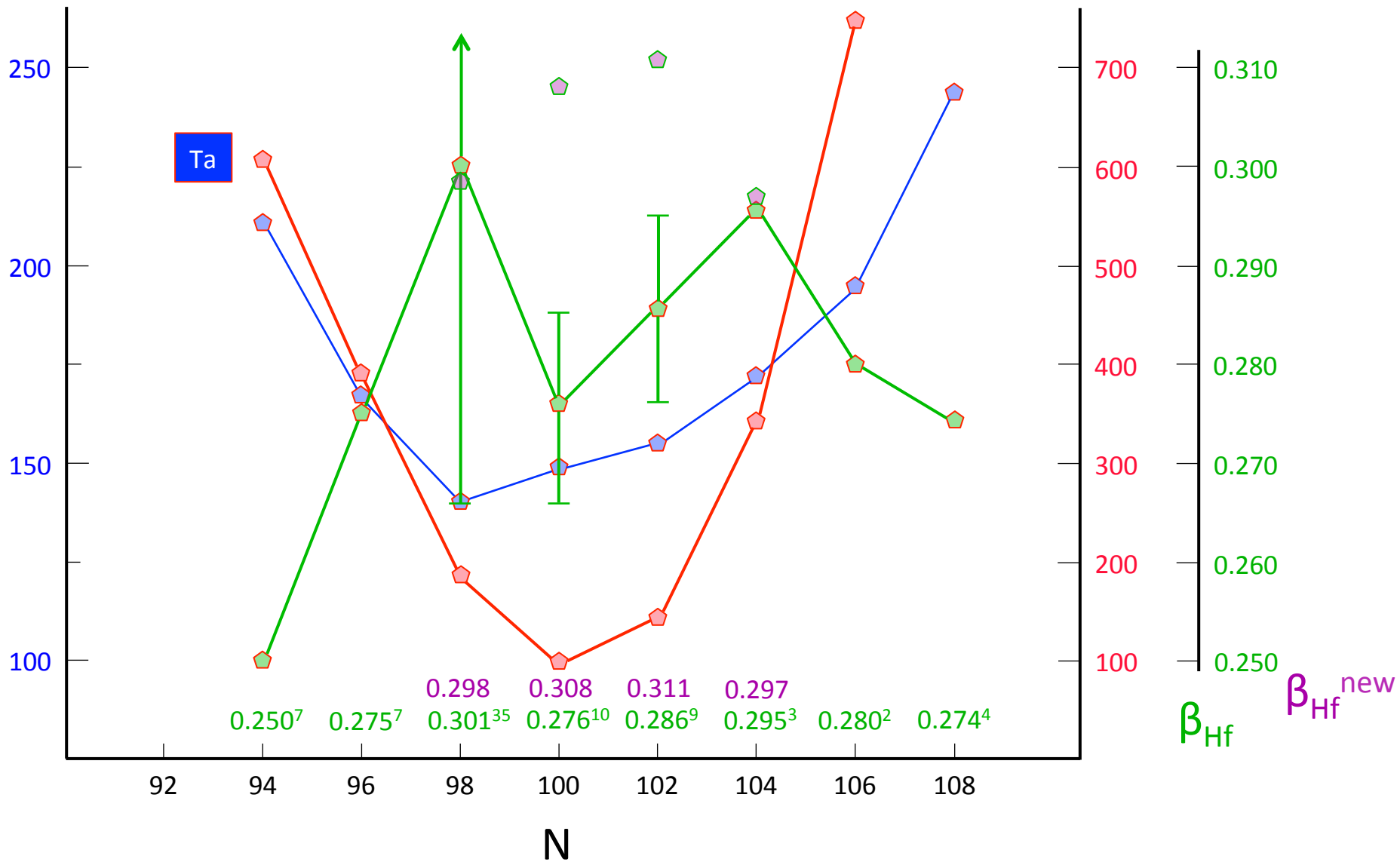


-β_{Hf}--th

alignment energy, band-head energy, deformation: 541↓ Ta

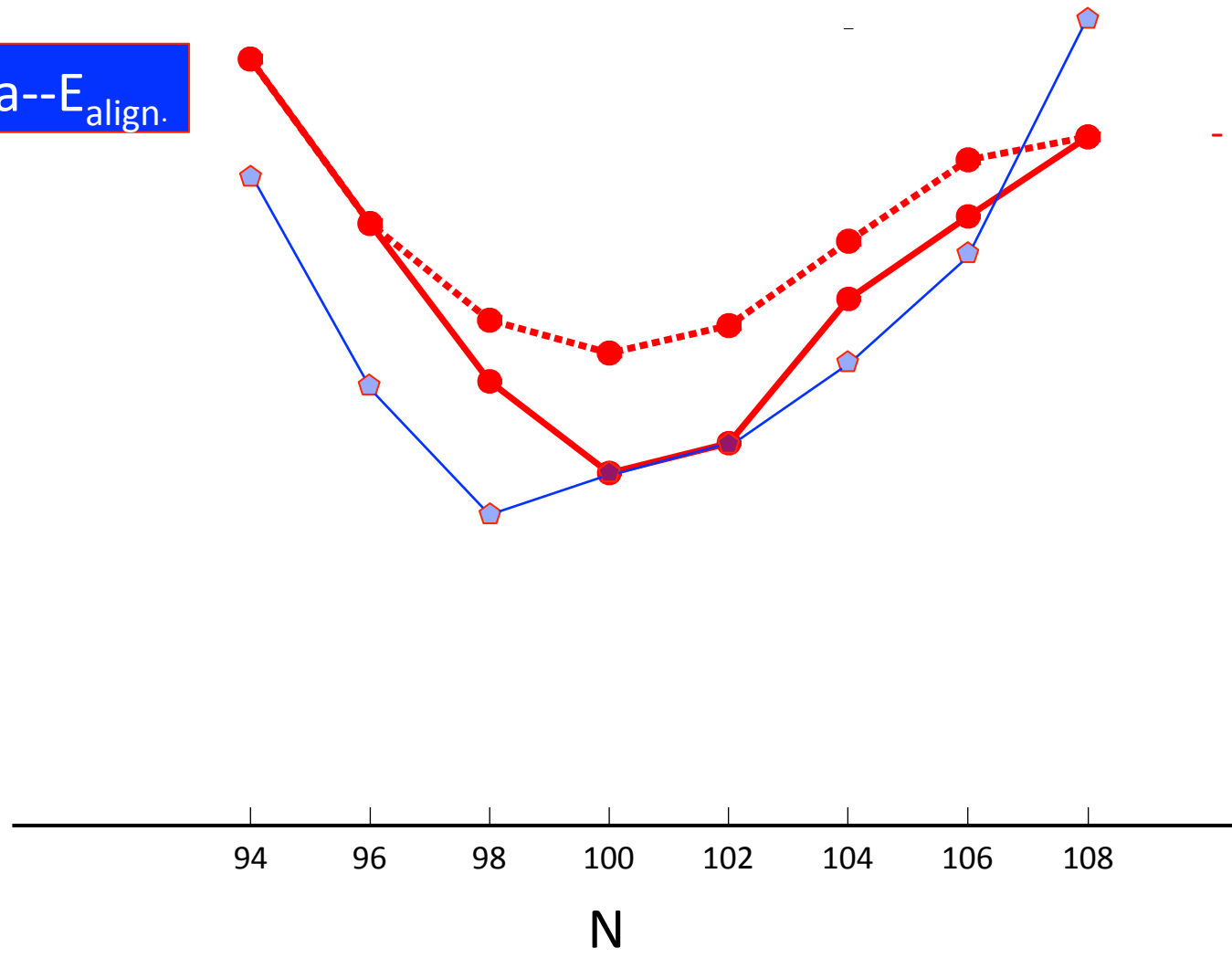
E_{align} (keV)

E_{541} (keV)

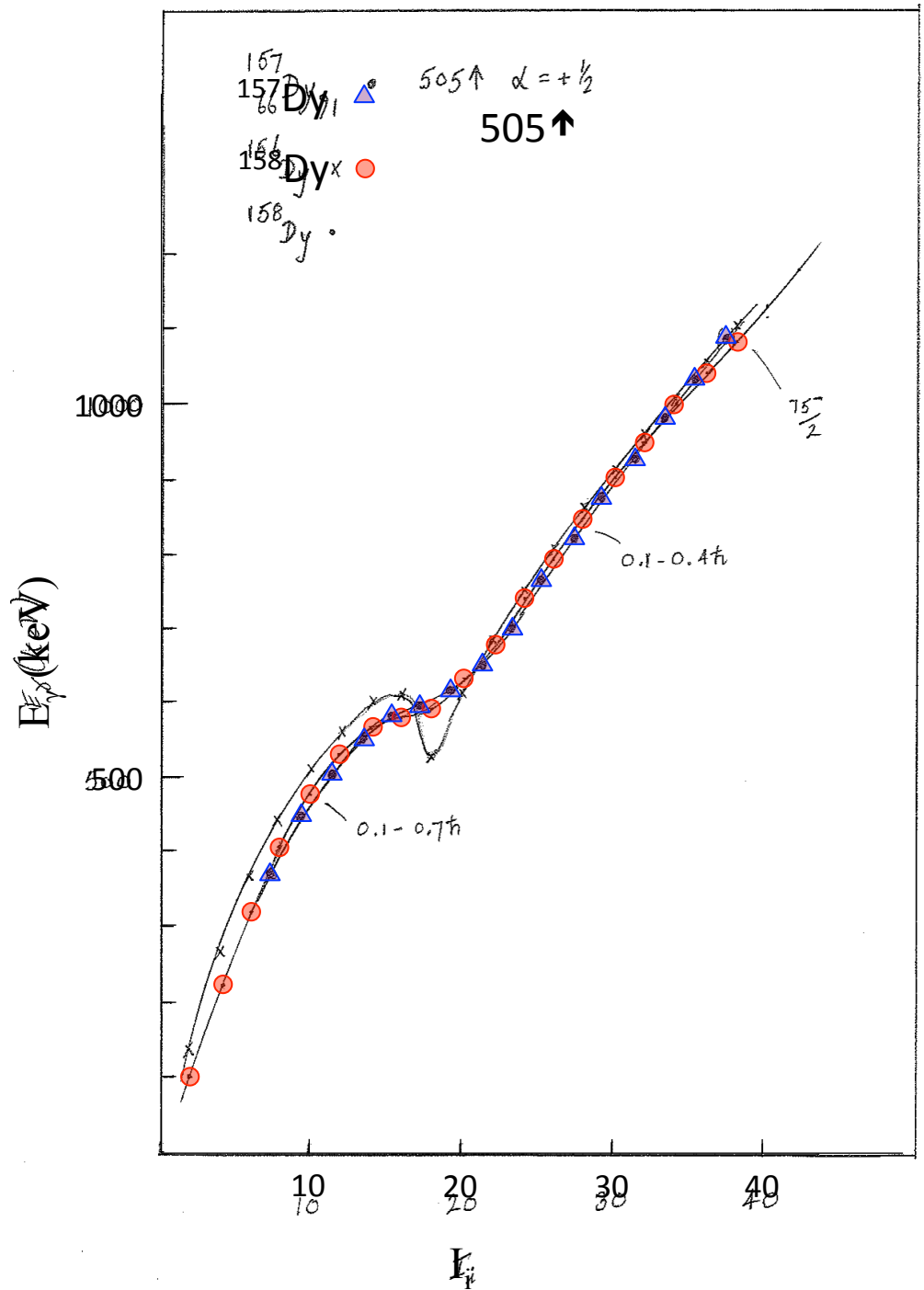


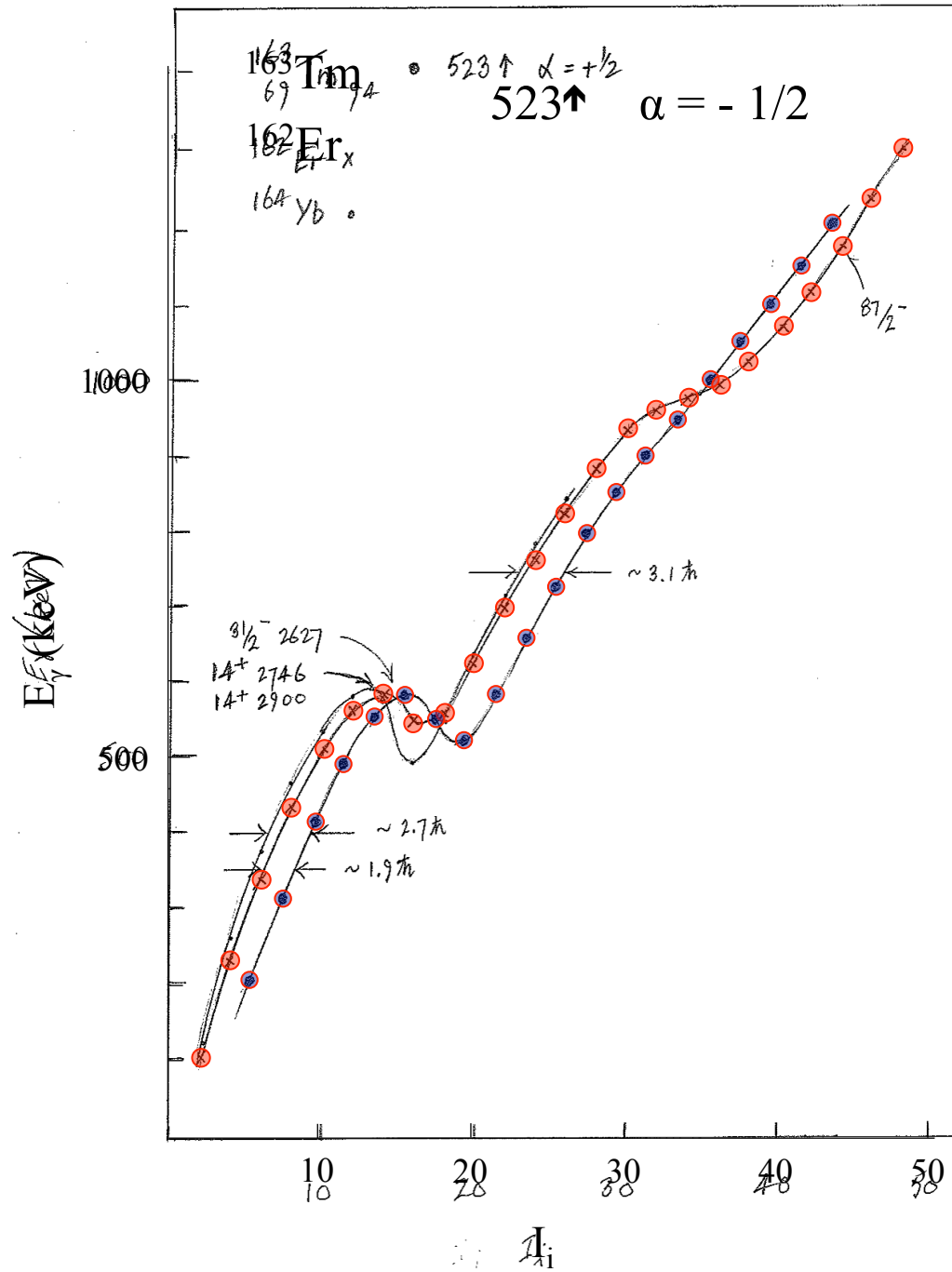
Ta--E_{align.}

- β_{Hf} --th



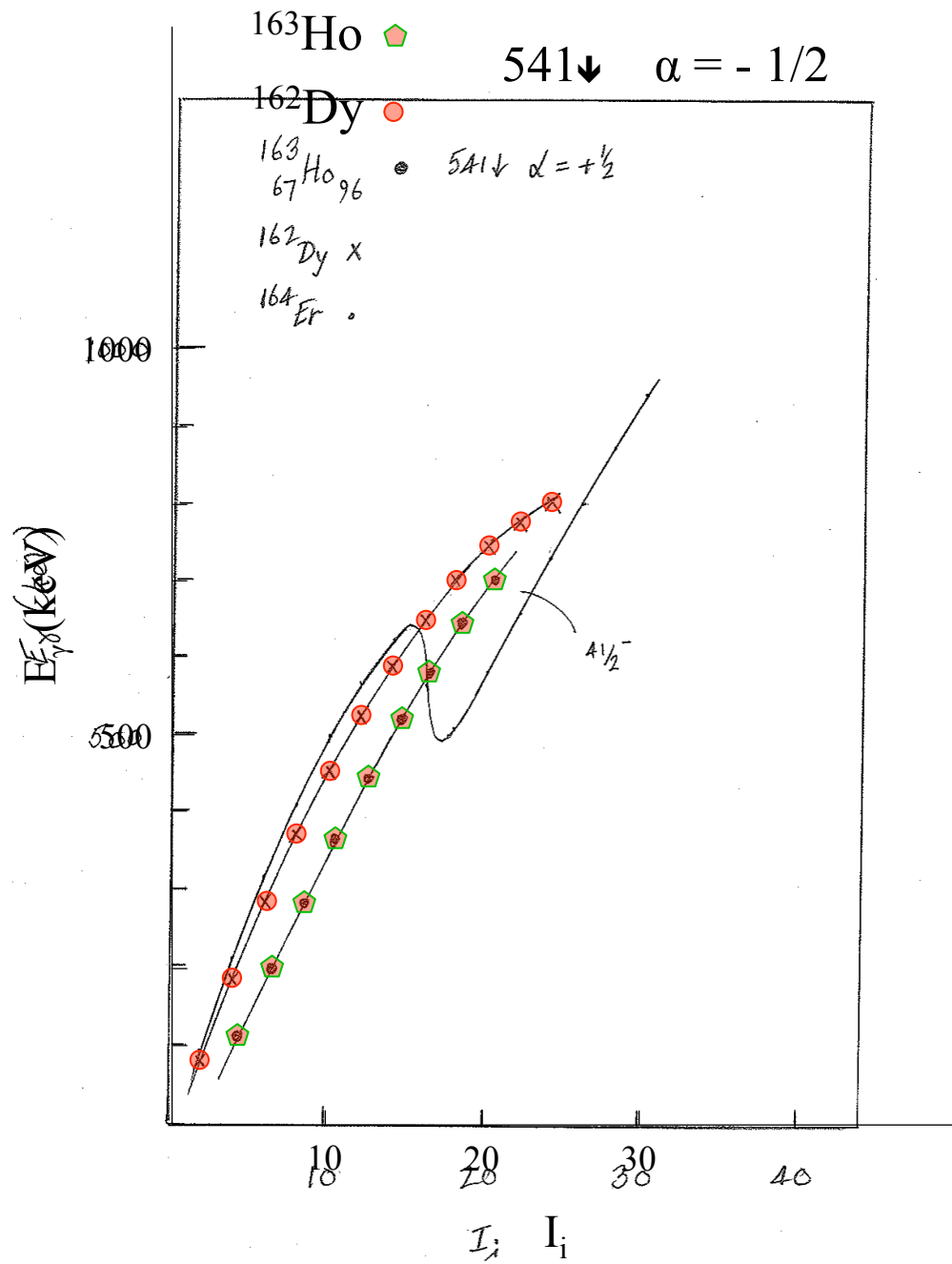
^{157}Dy





^{163}Tm

Note: reassignment of signature

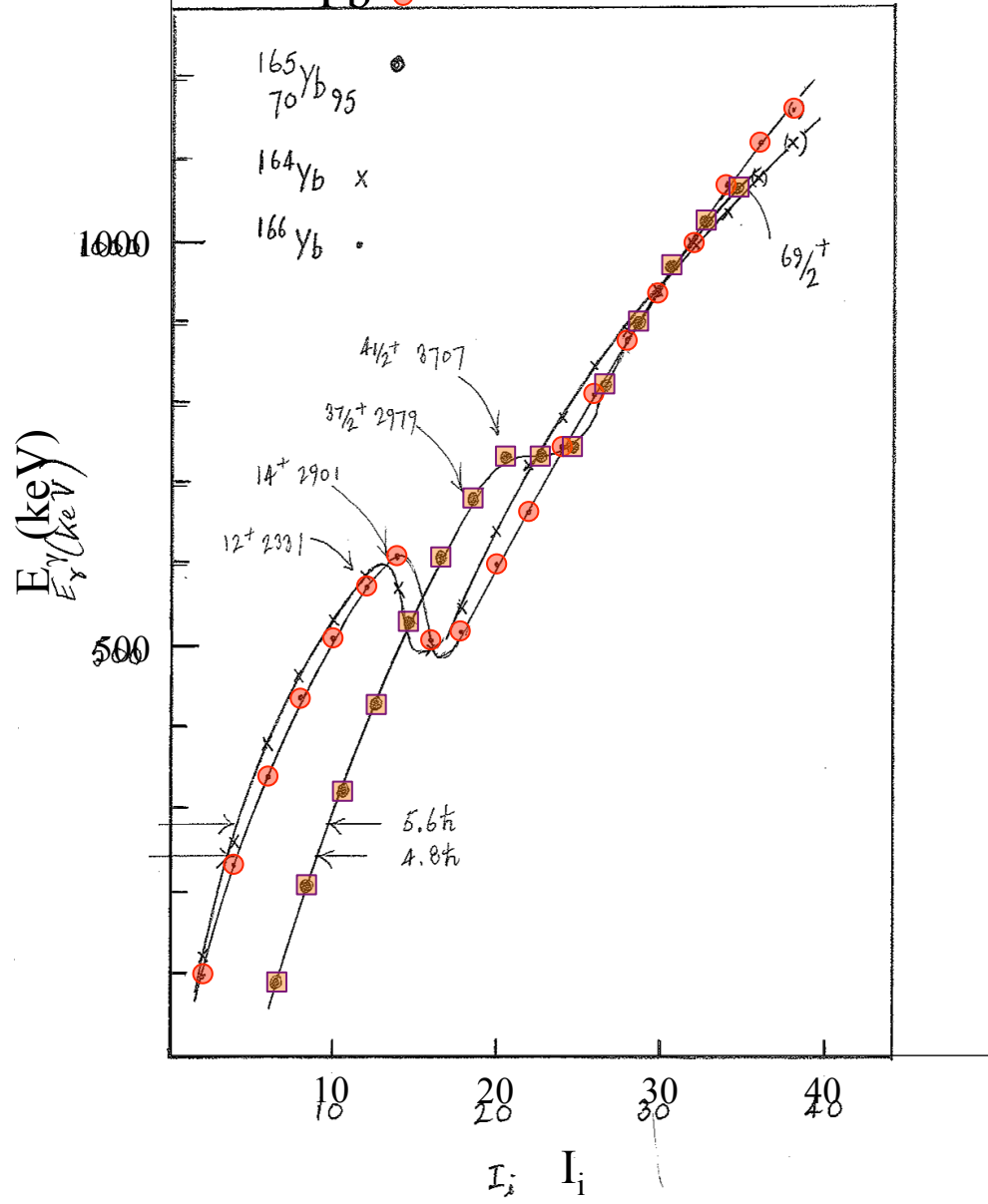


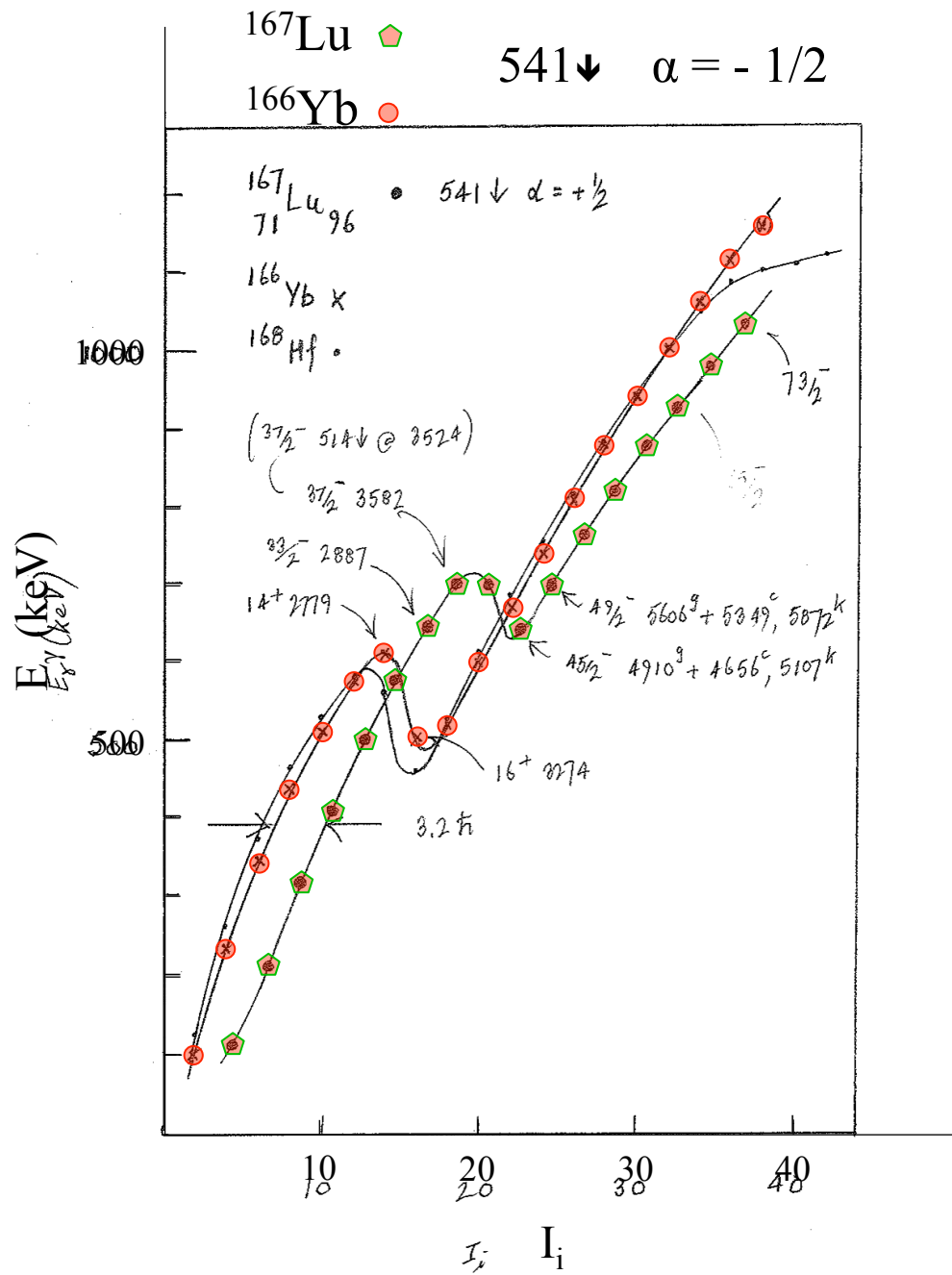
^{165}Yb

$642\uparrow \quad \alpha = +1/2$

^{165}Yb

^{166}Yb





^{167}Lu

514 $\alpha = -1/2$

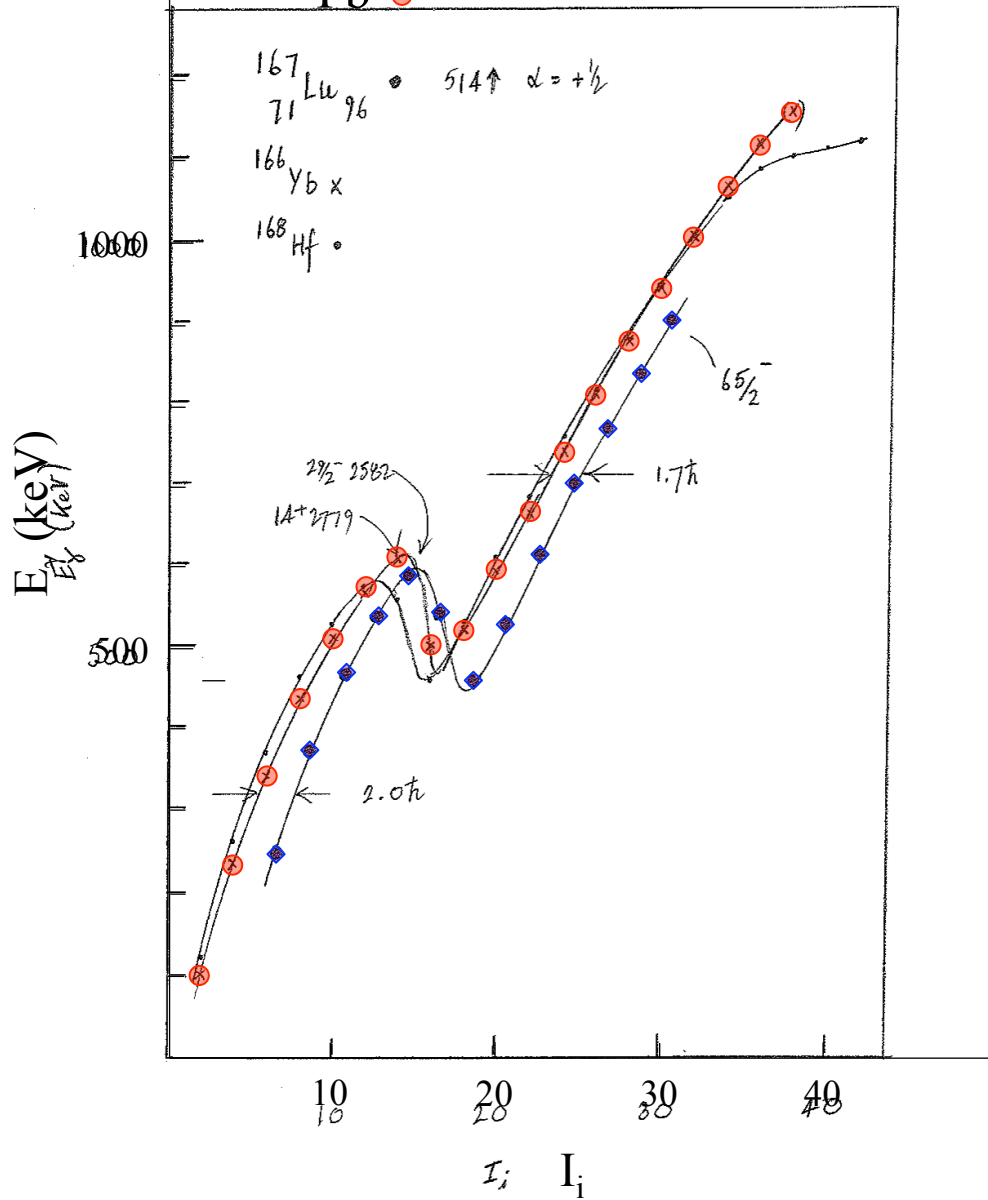
^{167}Lu

^{166}Yb ●

^{167}Lu 71 ● $514 \uparrow \alpha = +1/2$

^{166}Yb x

^{168}Hf ●



^{167}Lu

404 $\alpha = -1/2$

^{167}Lu

^{166}Yb ●

^{167}Lu 71 Lu 96 ● 404 $\alpha = +1/2$

^{166}Yb x

^{168}Hf ●

E_x (keV)
 E_x (keV)

1000

500

$31/2^+$ 2820

$11/2^+$ 2299

14^+ 2779

$55/2^+$ 6632 + 7232ⁿ

$59/2^+$ 7471 + 8048ⁿ

$63/2^+$ 8342^a + 8383ⁿ

$67/2^+$ 9266^a + 9874ⁿ

$71/2^+$ 10245^a + 10817ⁿ

$45^a + 10, 817^n$
 $+ 9874^a, 9841^n$
 $8383^a, 8918^n$

l $63/2^+$ 8383

$67/2^+$ 9874

n $35/2^+$ 8946

$39/2^+$ 11993

...

$81/2^+$

660[?]

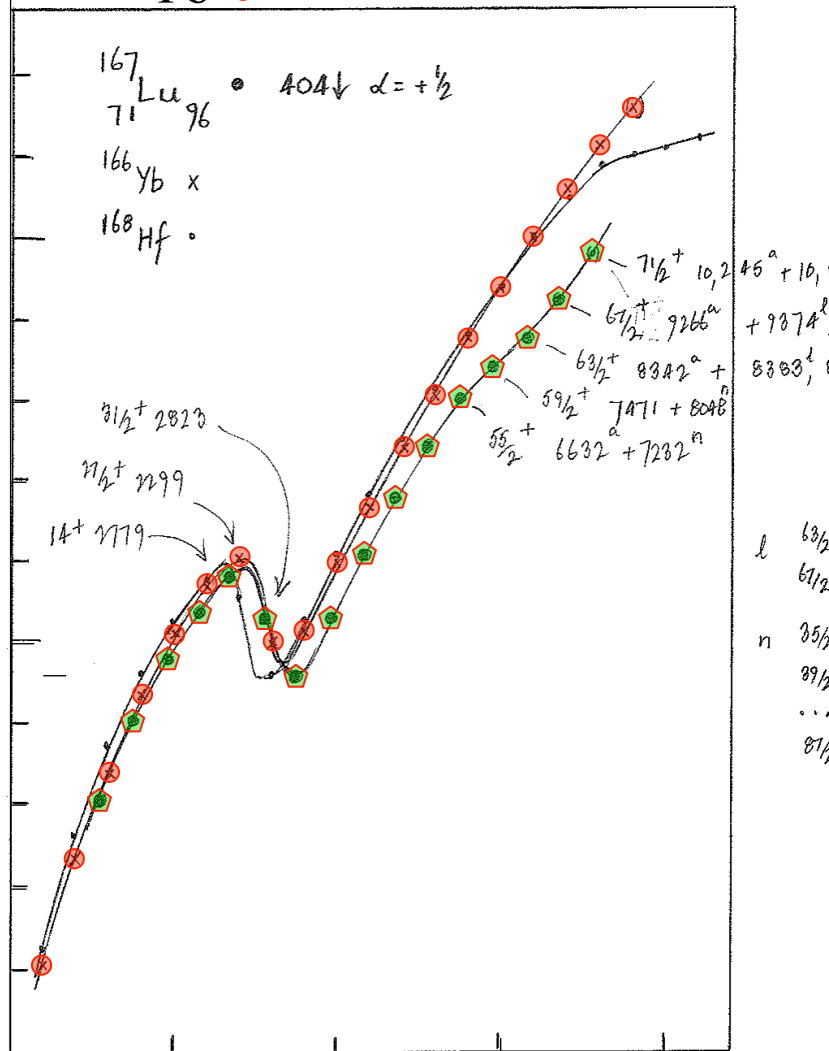
10

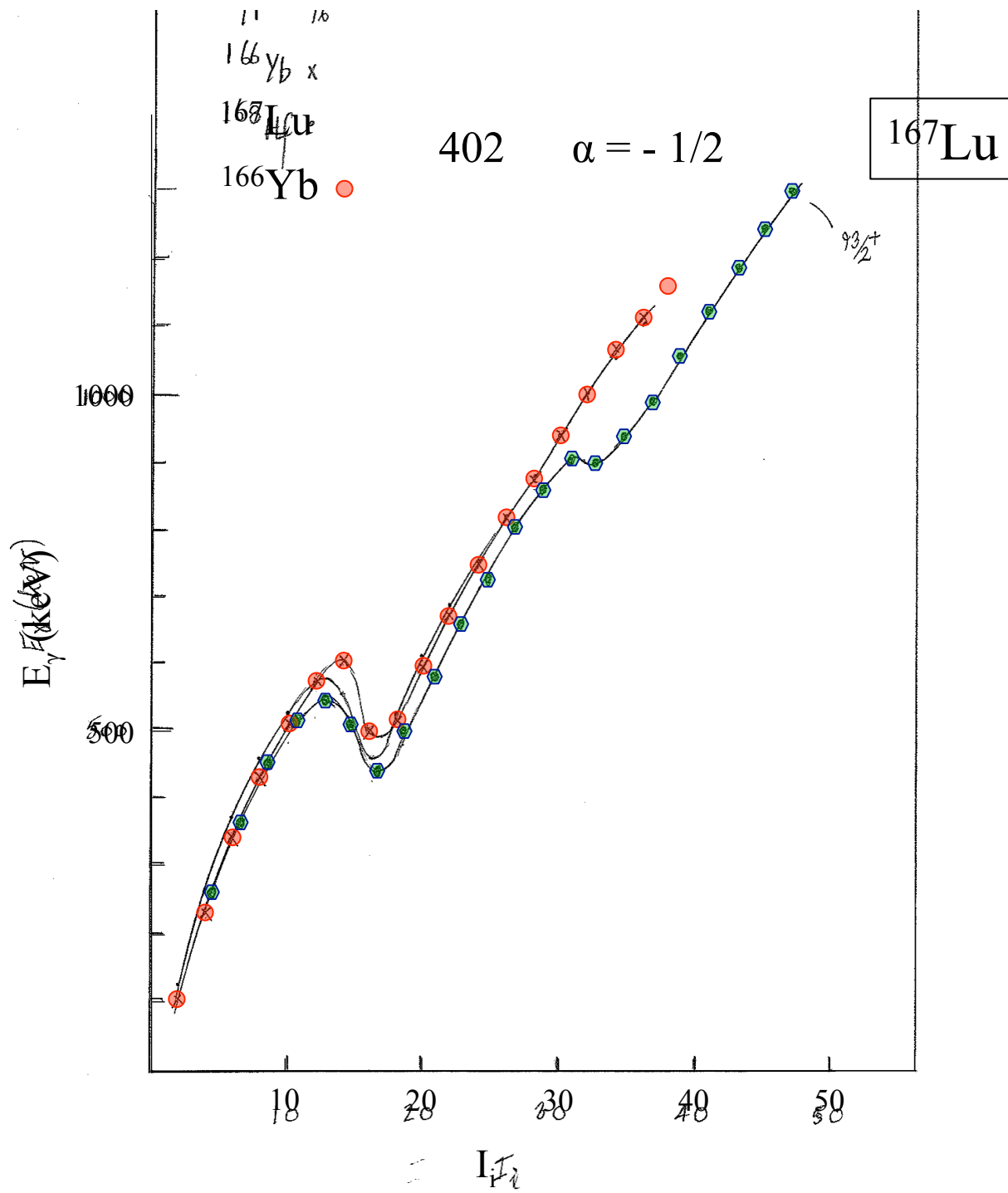
20

30

40

I_i I_i



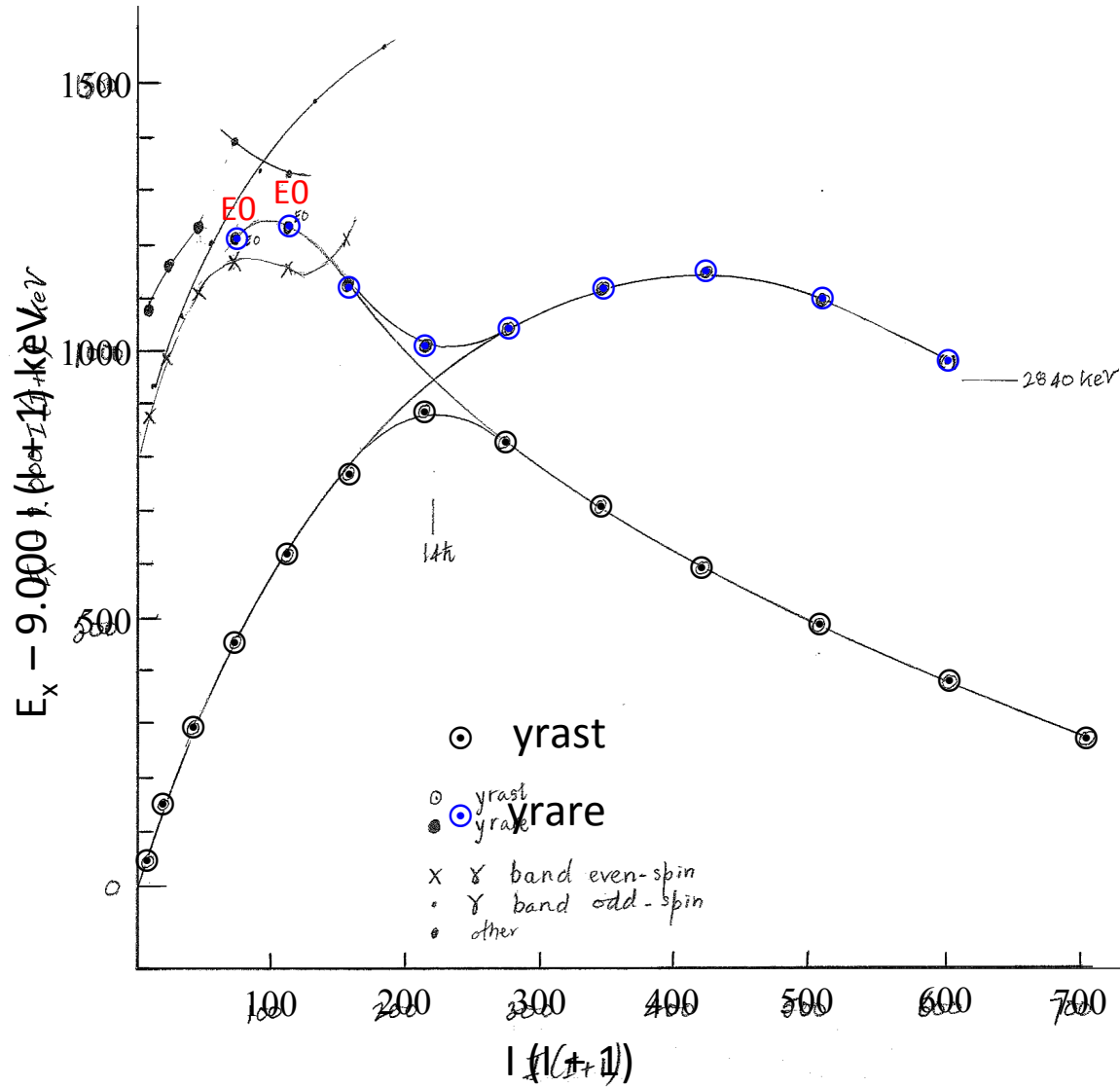


2/19/13

^{166}Yb

S band near identical to ^{164}Er
E0's. Fields, PL B130, 157; NP A431, 473

E0's: C. Fields et al., PL B130 157; NP A431 473



Annu. Rev. Nucl. Part. Sci. 1995. 45:485–541

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IDENTICAL BANDS IN DEFORMED AND SUPERDEFORMED NUCLEI

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“In particular, the reduction of the BCS pairing correlations due to the blocking of one and two orbitals implies large changes (up to 30%) in the moments of inertia and cannot be reconciled with these [identical band] systematics.”