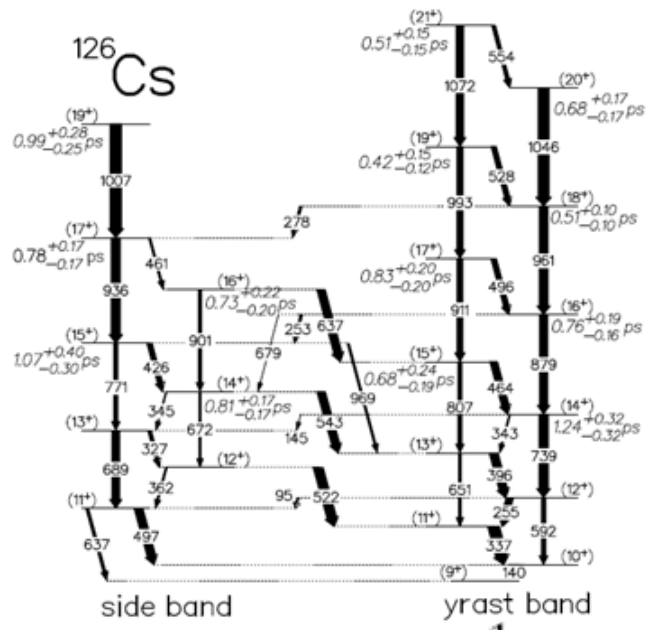


Search for the nuclear chirality

10 minutes manual for experimenters



Ernest Grodner

$$[R_Y T, H] = 0$$

$$|+\rangle = \frac{1}{\sqrt{2}} \frac{|L\rangle + |R\rangle}{\sqrt{1+\varepsilon}} \quad \langle + | H | + \rangle = \frac{E_0 + \Delta E}{1 + \varepsilon}$$

$$|-\rangle = \frac{i}{\sqrt{2}} \frac{|L\rangle - |R\rangle}{\sqrt{1-\varepsilon}} \quad \langle - | H | - \rangle = \frac{E_0 - \Delta E}{1 - \varepsilon}$$

Parameters

Overlap

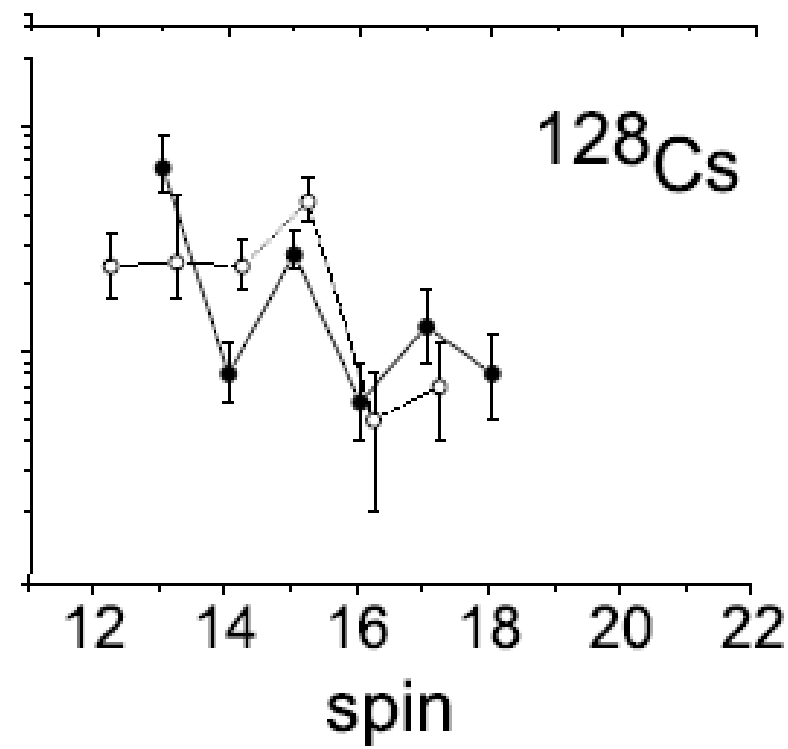
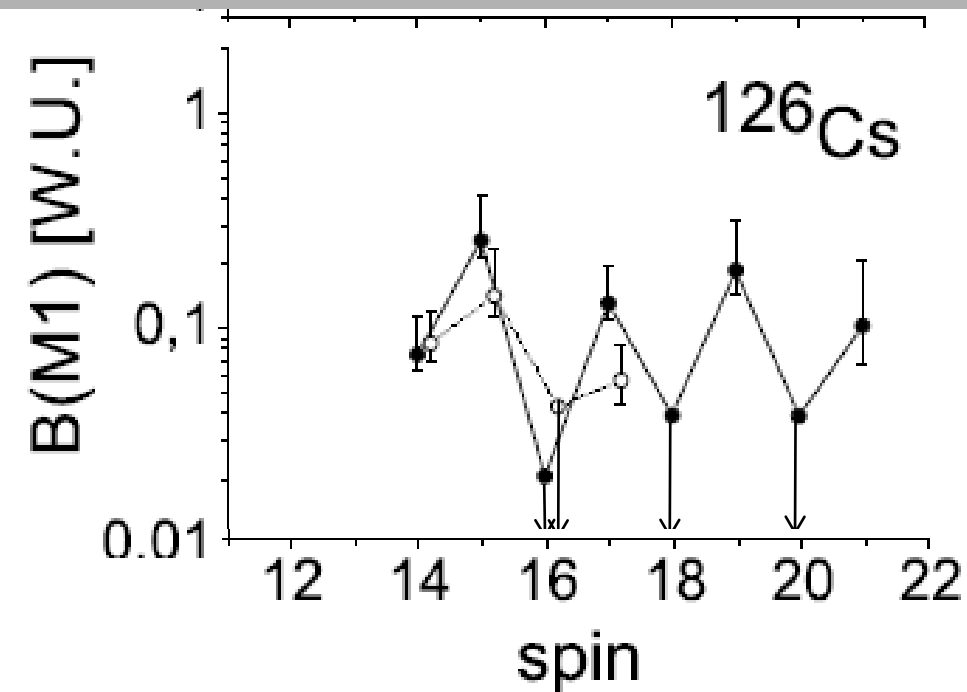
$$\varepsilon = \text{Re}\langle L | R \rangle$$

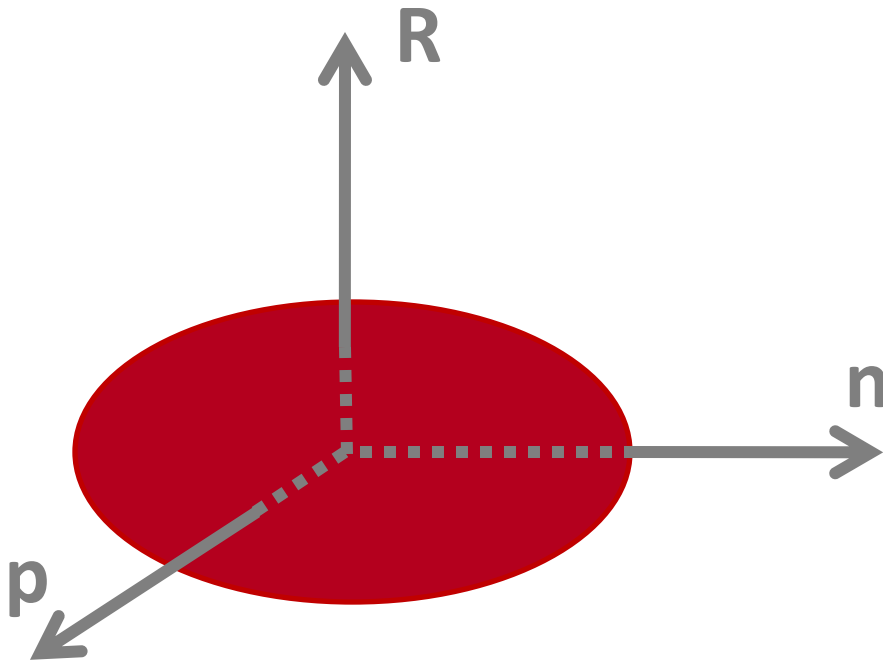
Tunneling effect

$$\Delta E = \text{Re}\langle L | H | R \rangle$$

Diagonal mat. element

$$E_0 = \text{Re}\langle L | H | L \rangle$$

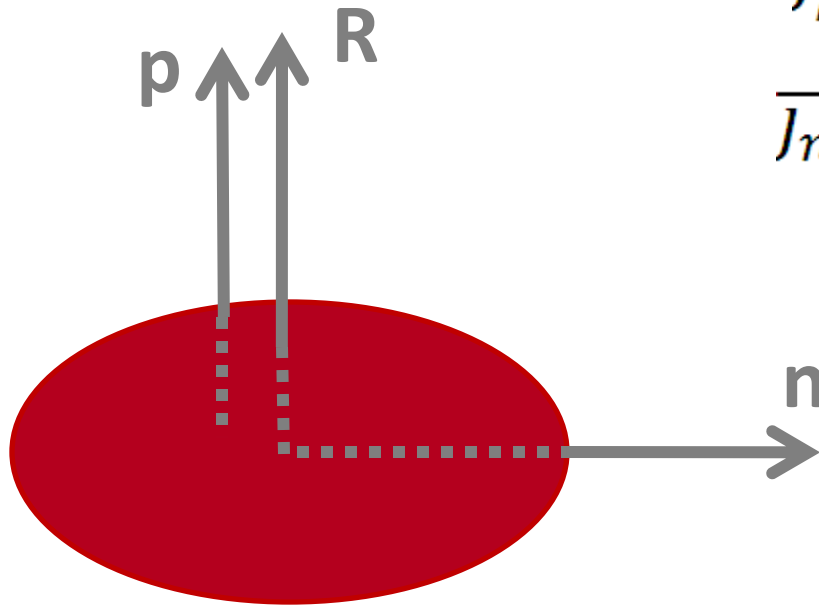




$$\vec{J}_p \cdot \vec{J}_R = 0$$

$$\vec{J}_n \cdot \vec{J}_R = 0$$

$$\vec{l} \cdot \vec{\omega} = 0$$



$$\vec{J}_p \cdot \vec{J}_R \neq 0$$

Large energy splitting for $|I, M\rangle$ and $|I, -M\rangle$ substates

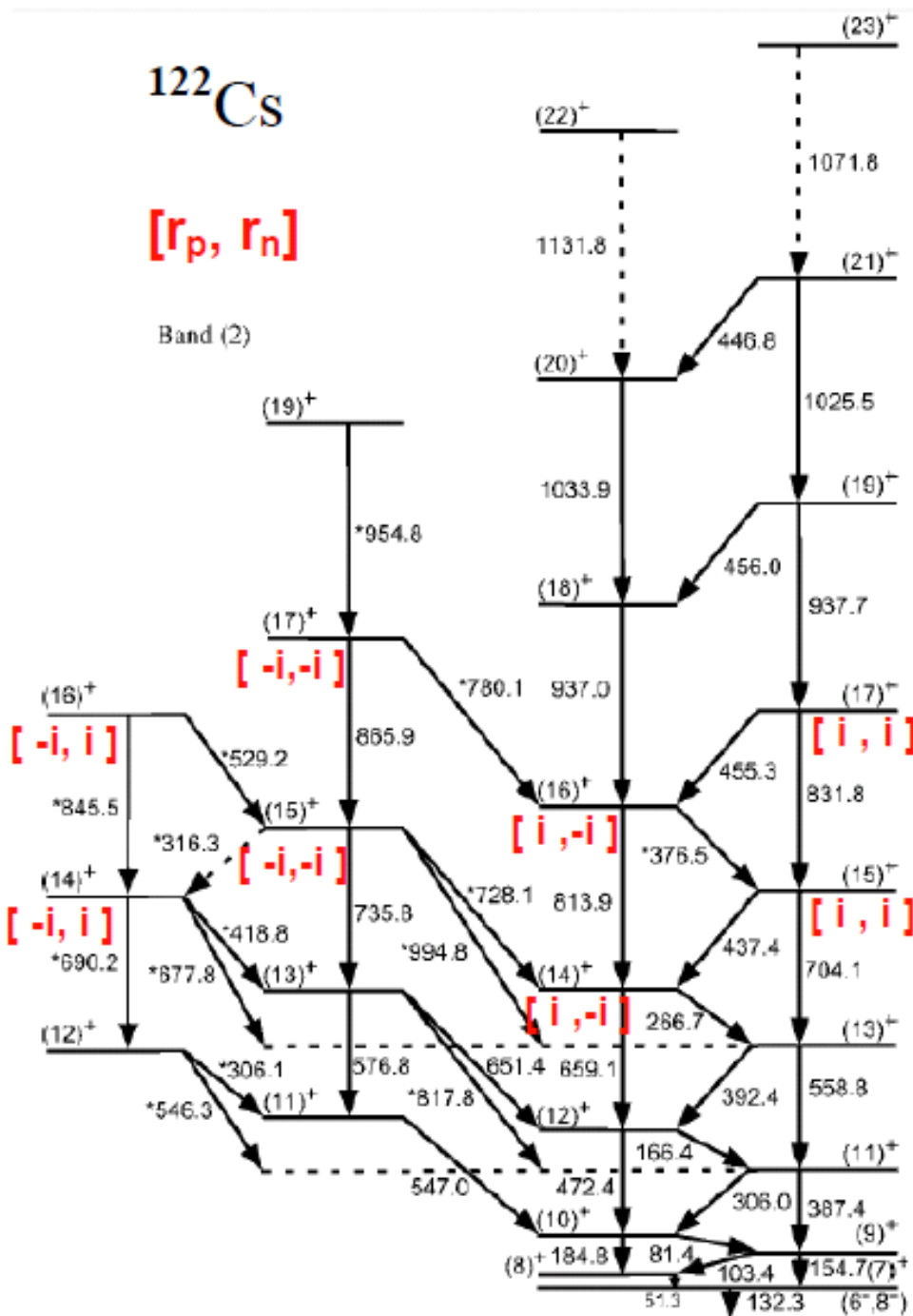
$$\vec{J}_n \cdot \vec{J}_R \neq 0$$

Small energy splitting

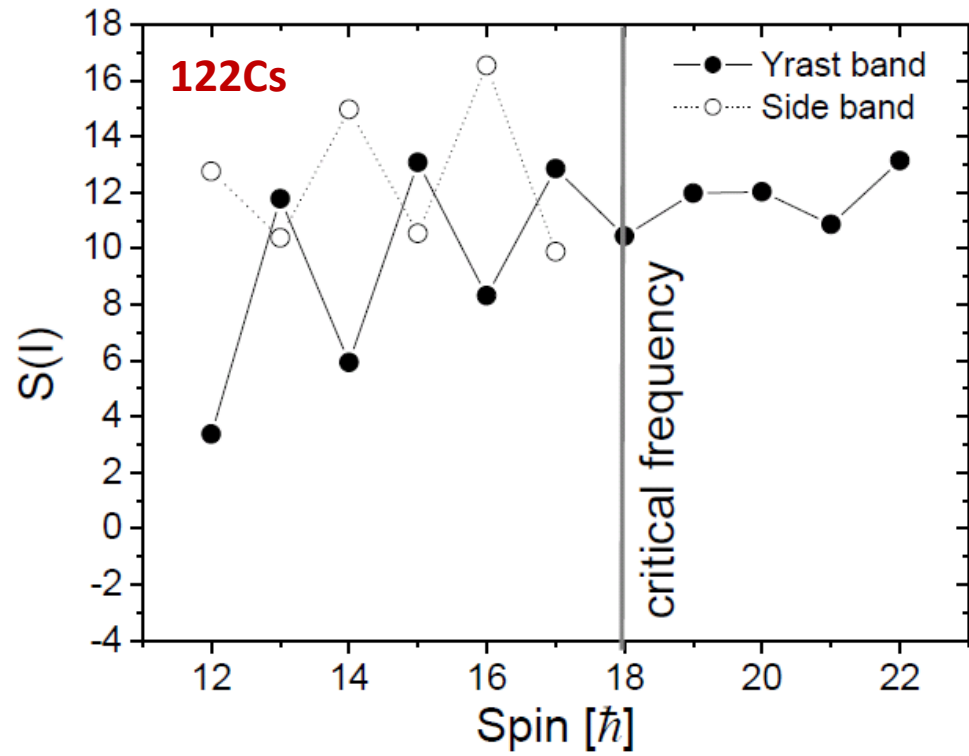
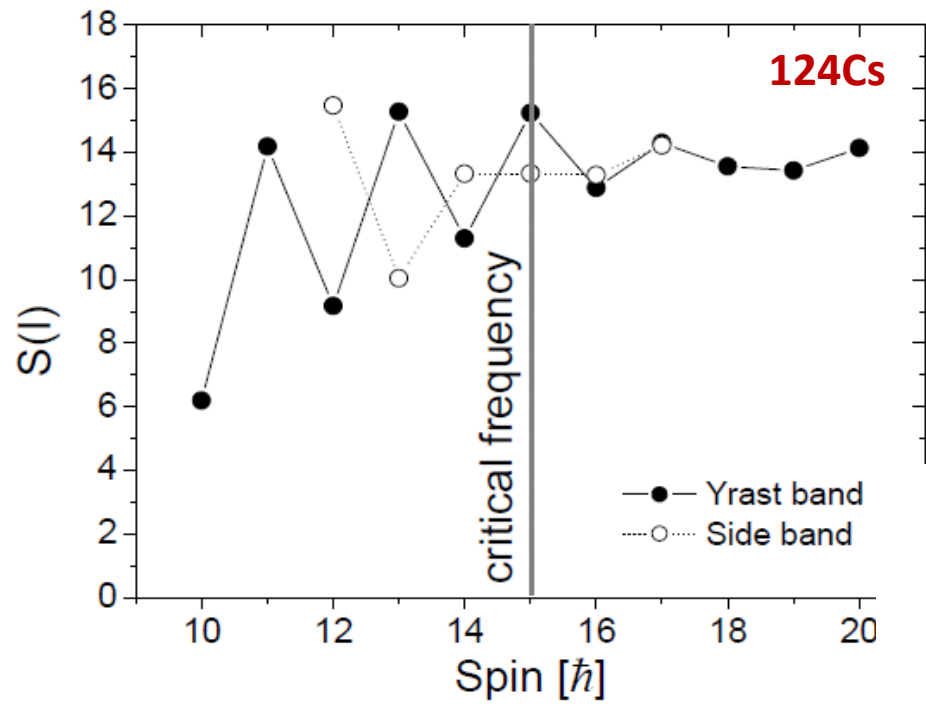
^{122}Cs

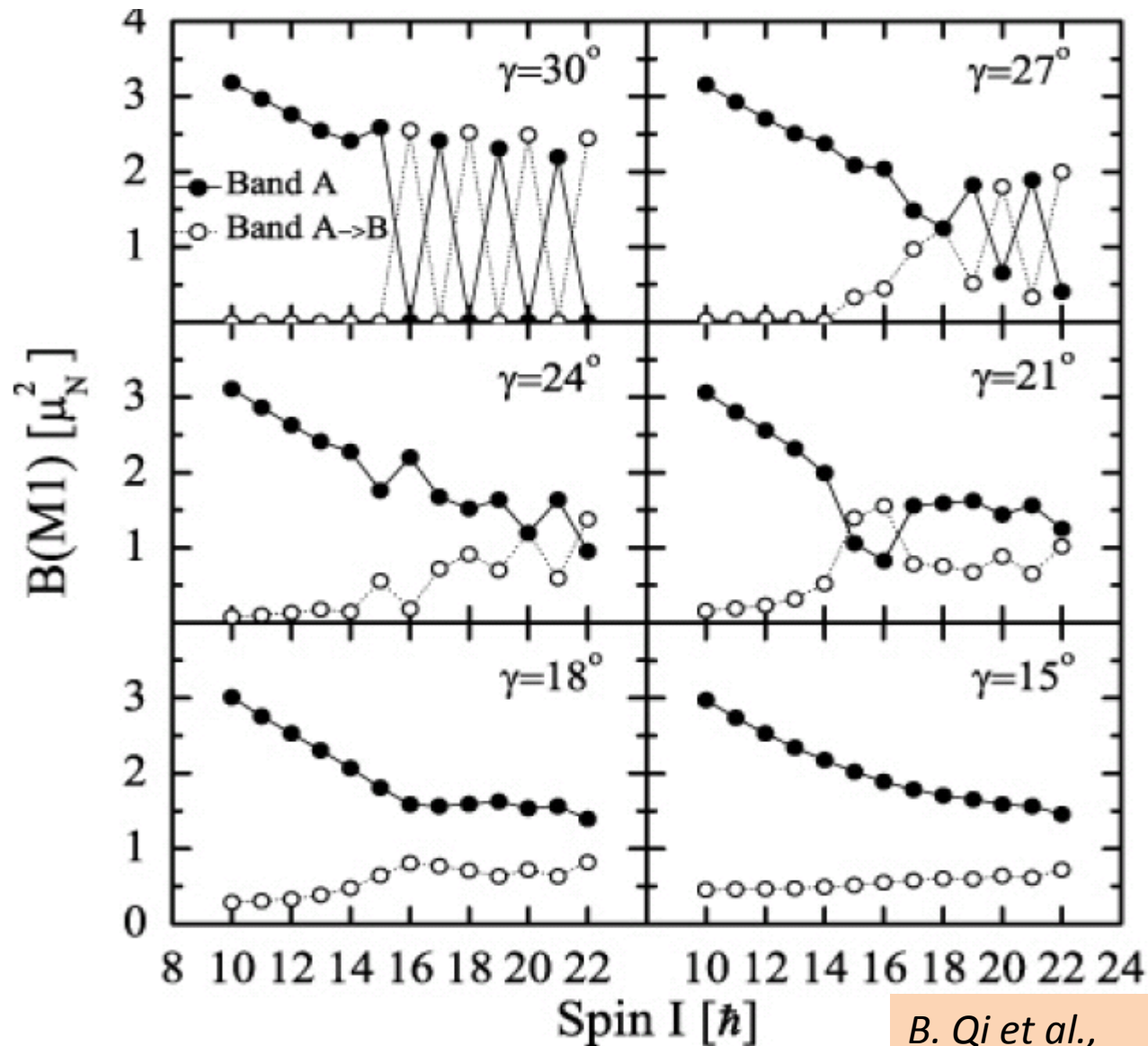
$[r_p, r_n]$

Band (2)



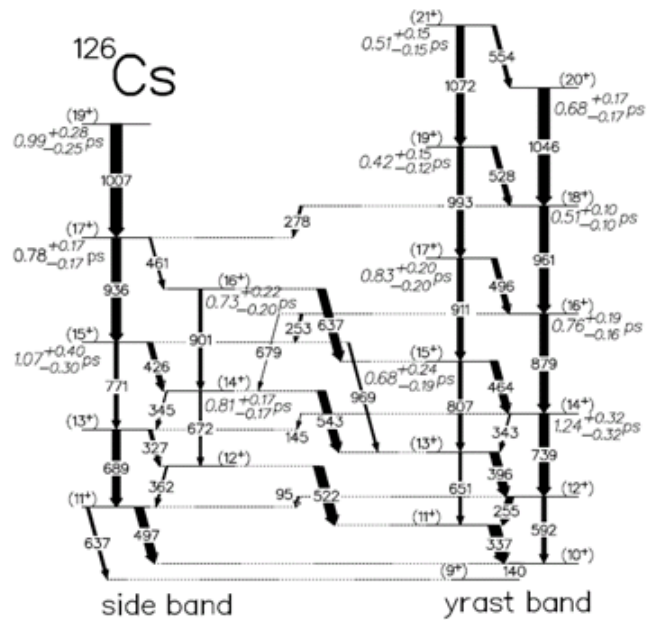
total signature	r_p	r_n	r_c	spin sequence
-1	i	i	1	1, 3, 5, 7, ...
1	i	-i	1	0, 2, 4, 6, ...
-1	-i	-i	1	1, 3, 5, 7, ...
1	-i	i	1	0, 2, 4, 6, ...





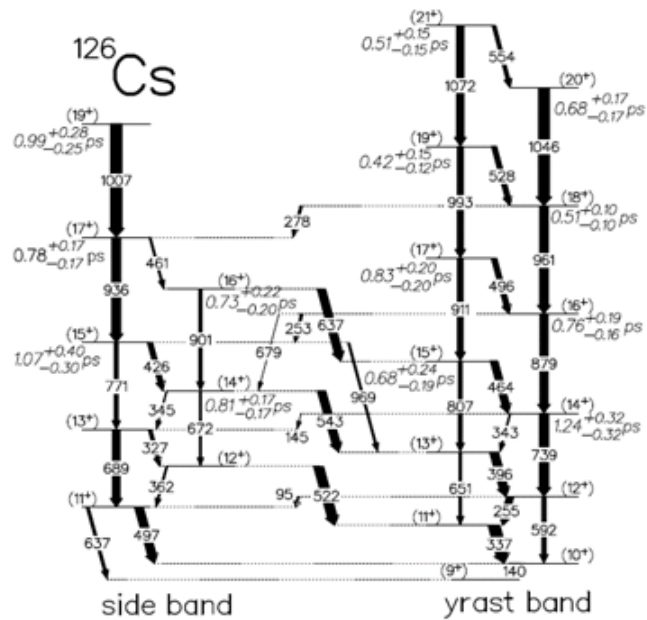
B. Qi et al.,

PHYSICAL REVIEW C 79, 041302(R) (2009)



$$|I, M, +\rangle = \frac{1}{\sqrt{2N_{I+}}} (|I, M, L\rangle + |I, M, R\rangle)$$

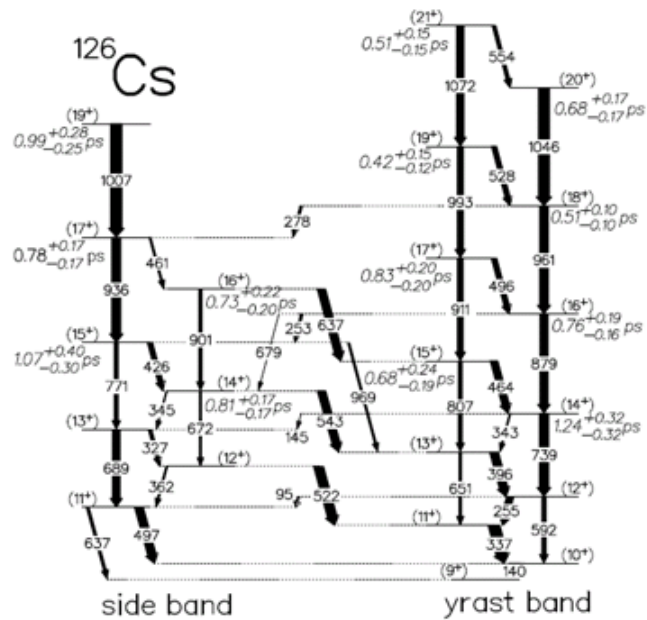
$$|I, M, -\rangle = \frac{i}{\sqrt{2N_{I-}}} (|I, M, L\rangle - |I, M, R\rangle)$$



$$|I, M, +\rangle = \frac{1}{\sqrt{2}N_{I+}} (|I, M, L\rangle + |I, M, R\rangle)$$

$$|I, M, -\rangle = \frac{i}{\sqrt{2}N_{I-}} (|I, M, L\rangle - |I, M, R\rangle)$$

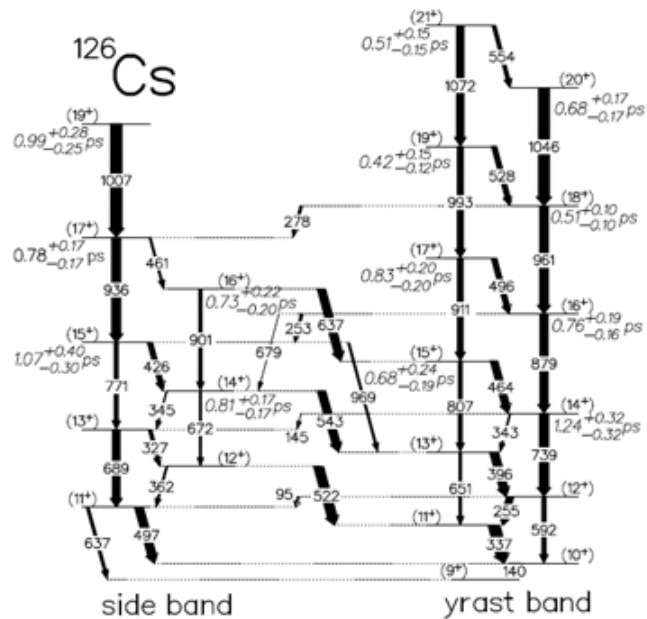
$$\langle I_2, + || M || I_1, + \rangle = \frac{1}{\sqrt{2}N_{I_2+}} (\langle I_2, L | + \langle I_2, R |) |M| \frac{1}{\sqrt{2}N_{I_1+}} (|I_1, L\rangle + |I_1, R\rangle)$$



$$|I, M, +\rangle = \frac{1}{\sqrt{2N_{I+}}} (|I, M, L\rangle + |I, M, R\rangle)$$

$$|I, M, -\rangle = \frac{i}{\sqrt{2N_{I-}}} (|I, M, L\rangle - |I, M, R\rangle)$$

$$\langle I_2, + || M || I_1, + \rangle = \frac{1}{2N_{I_2+} + N_{I_1+}} (\langle I_2, L || M || I_1, L \rangle + \langle I_2, R || M || I_1, R \rangle + \langle I_2, L || M || I_1, R \rangle + \langle I_2, R || M || I_1, L \rangle)$$

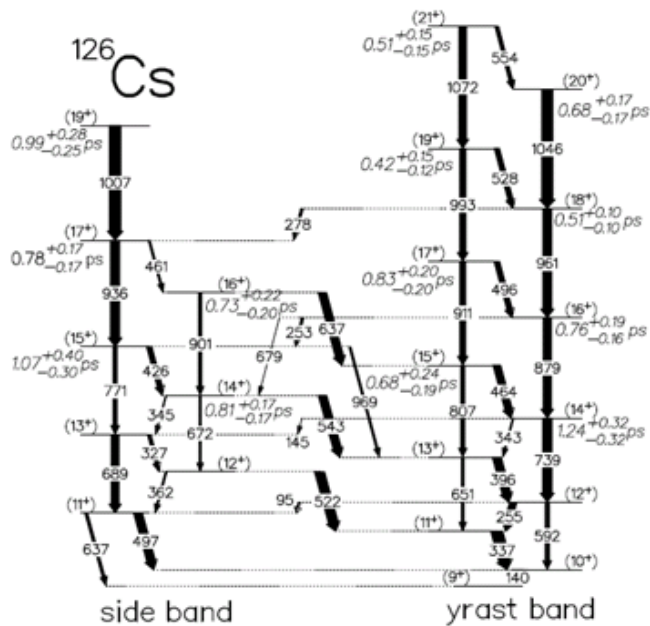


$$|I, M, +\rangle = \frac{1}{\sqrt{2N_{I+}}} (|I, M, L\rangle + |I, M, R\rangle)$$

$$|I, M, -\rangle = \frac{i}{\sqrt{2N_{I-}}} (|I, M, L\rangle - |I, M, R\rangle)$$

$$\langle L|M|R\rangle^* = \langle R|M|L\rangle$$

$$\langle I_2, + || M || I_1, + \rangle = \frac{1}{2N_{I_2+} + N_{I_1+}} (\langle I_2, L || M || I_1, L \rangle + \langle I_2, R || M || I_1, R \rangle + \langle I_2, L || M || I_1, R \rangle + \langle I_2, R || M || I_1, L \rangle)$$



$$|I, M, +\rangle = \frac{1}{\sqrt{2}N_{I+}} (|I, M, L\rangle + |I, M, R\rangle)$$

$$|I, M, -\rangle = \frac{i}{\sqrt{2}N_{I-}} (|I, M, L\rangle - |I, M, R\rangle)$$

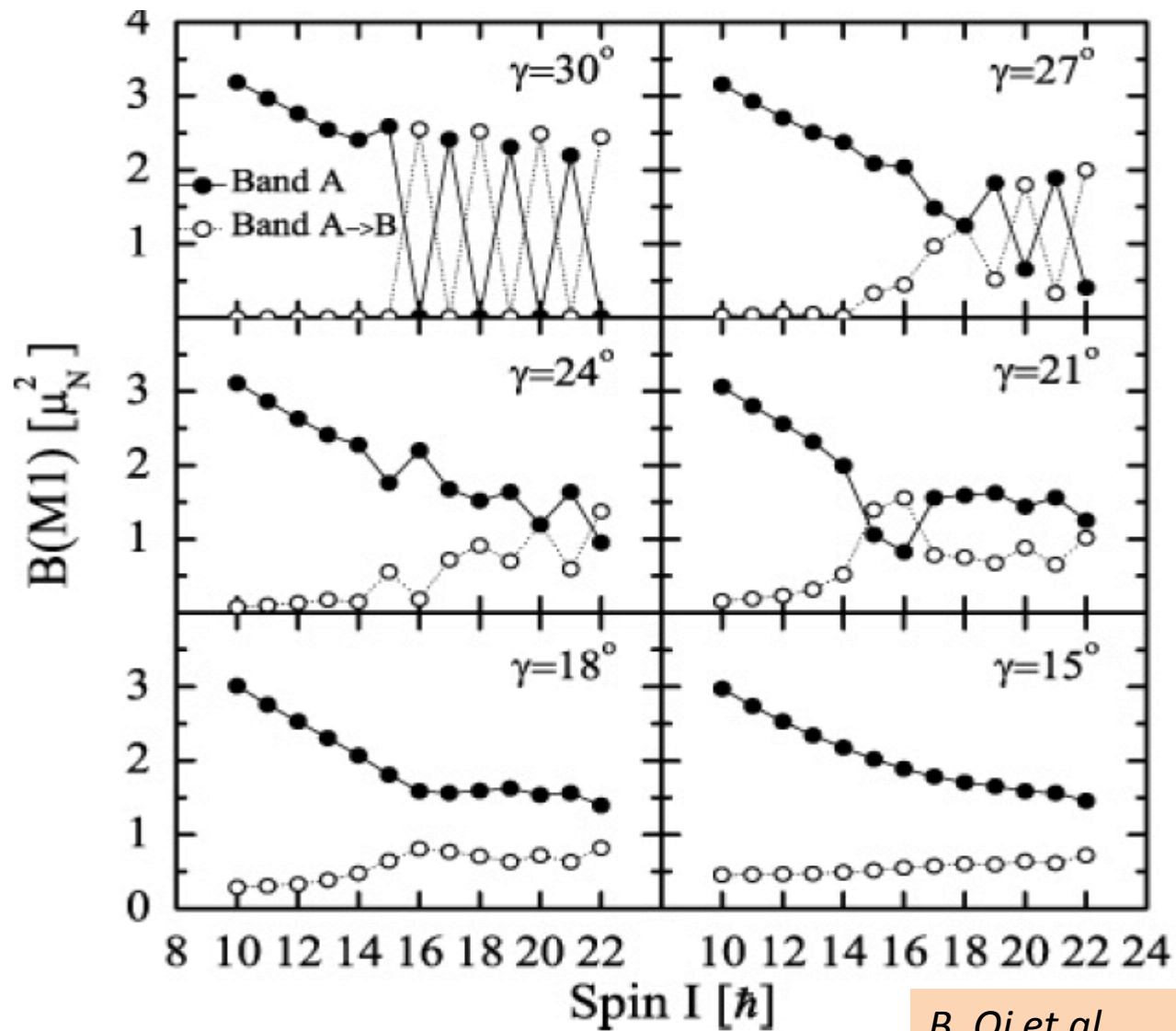
$$\langle L|M|R\rangle^* = \langle R|M|L\rangle$$

$$\langle I_2, + || M || I_1, + \rangle = \text{Re} \langle I_2, L || M || I_1, L \rangle$$

$$\langle I_2, - || M || I_1, - \rangle = \text{Re} \langle I_2, L || M || I_1, L \rangle$$

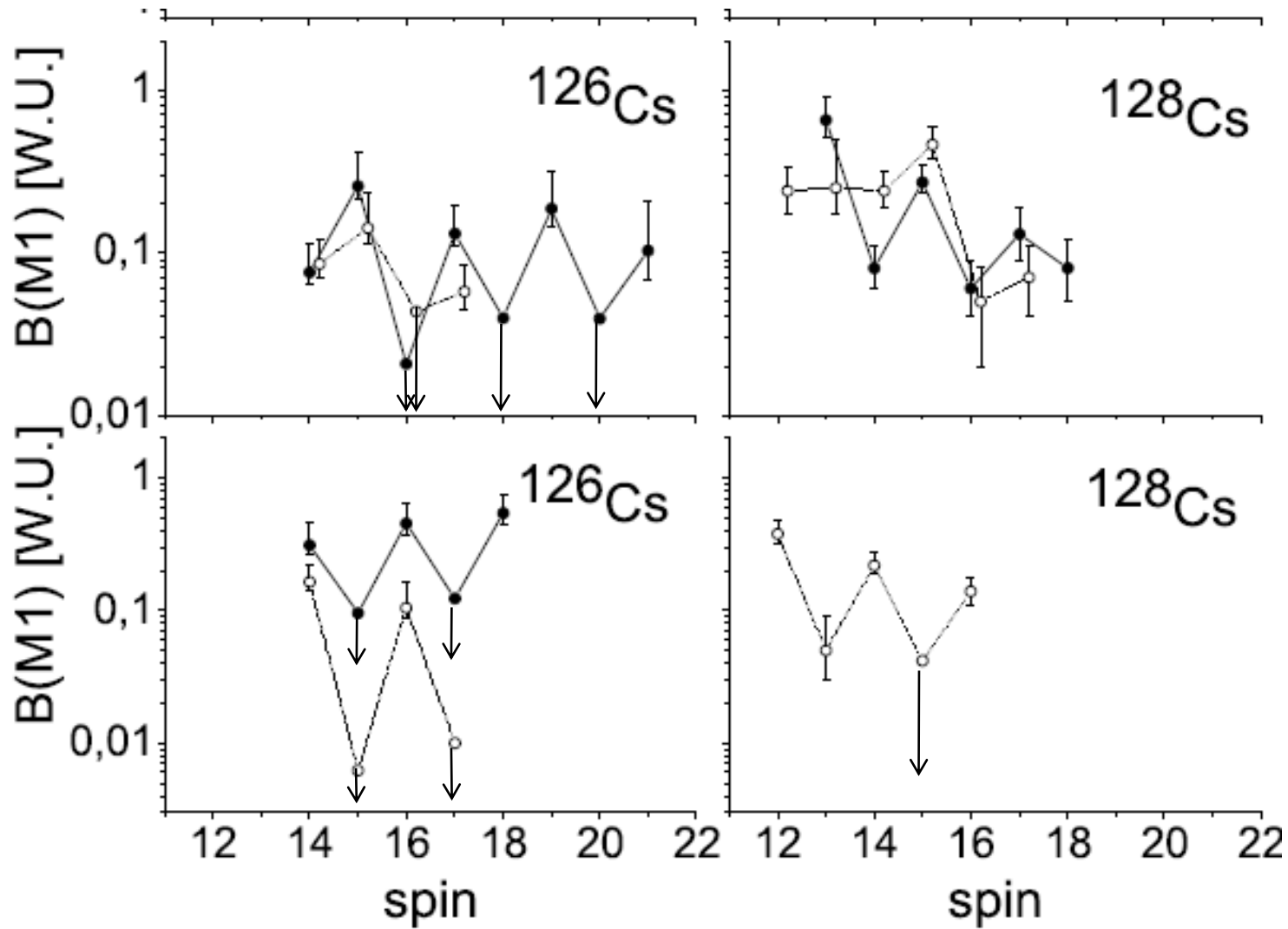
$$\langle I_2, + || M || I_1, - \rangle = \text{Im} \langle I_2, L || M || I_1, L \rangle$$

$$\langle I_2, - || M || I_1, + \rangle = \text{Im} \langle I_2, L || M || I_1, L \rangle$$



B. Qi et al.,

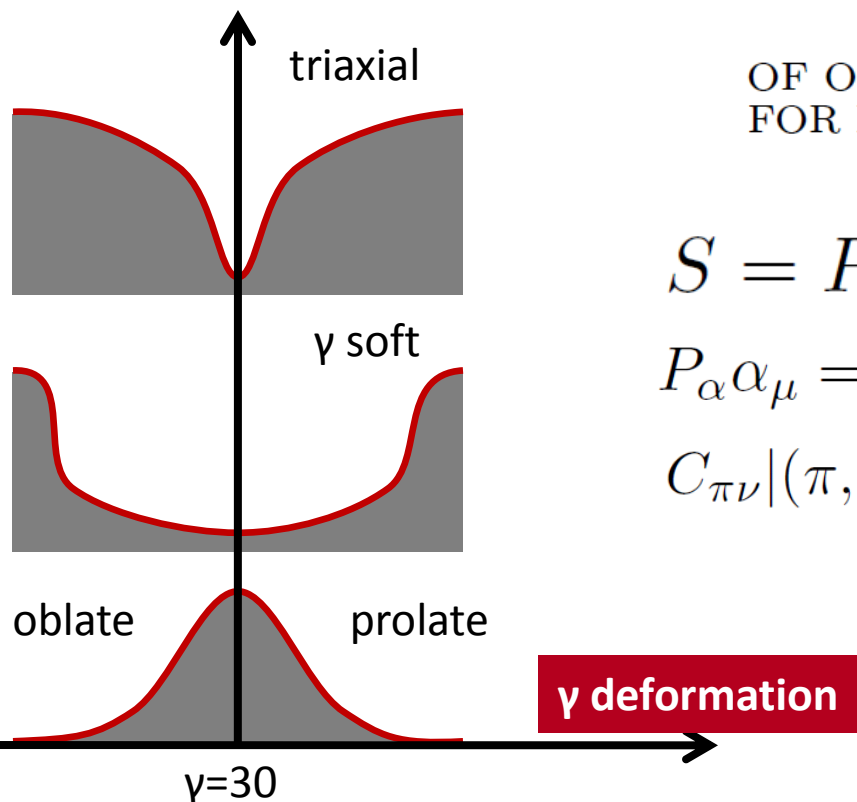
PHYSICAL REVIEW C 79, 041302(R) (2009)



Chiral Bands, Dynamical Spontaneous Symmetry Breaking, and the Selection Rule for Electromagnetic Transitions in the Chiral Geometry

T. Koike,¹ K. Starosta,^{1,2} and I. Hamamoto^{3,4}

Def. potential



Vol. 42 (2011)

ACTA PHYSICA POLONICA B

No 3-4

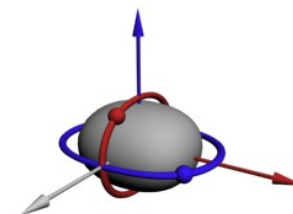
A SYMMETRY OF THE CPHC MODEL OF ODD-ODD NUCLEI AND ITS CONSEQUENCES FOR PROPERTIES OF $M1$ AND $E2$ TRANSITIONS*

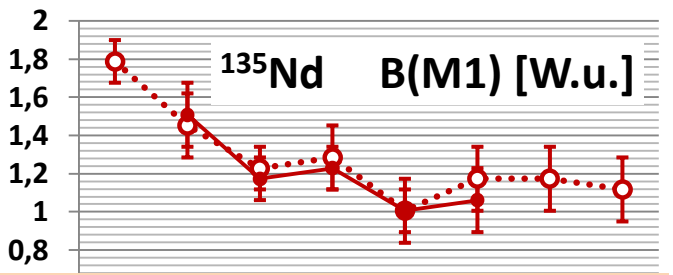
L. PRÓCHNIAK

$$S = P_{\alpha} C_{\pi\nu}$$

$$P_{\alpha} \alpha_{\mu} = -\alpha_{\mu}, \quad \mu = -2, \dots, 2$$

$$C_{\pi\nu} |(\pi, j_{\pi} m_{\pi})(\nu, j_{\nu} m_{\nu})\rangle = |(\pi, j_{\nu} m_{\nu})(\nu, j_{\pi} m_{\pi})\rangle$$



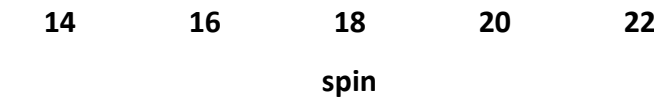


S.Mukhopadhyay et al.,
Phys. Rev. Lett. 97 (2007) 172501

3qp configuration
 $(\pi h_{11/2})^2 \otimes \nu h_{11/2}$
 ●●● yrast
 ● side
 ▲ side->yrast

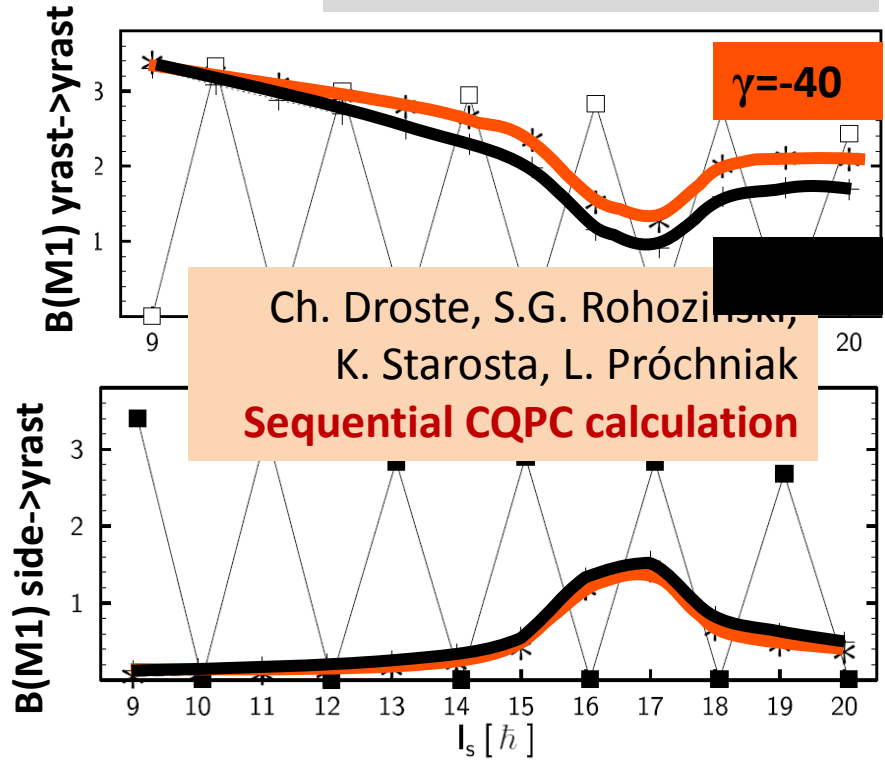
Single-j configuration
~~2qp configuration~~
 Triaxiality $\gamma=30$

Nonaxial deformation
 $\pi h_{11/2} \otimes \nu h_{11/2}$



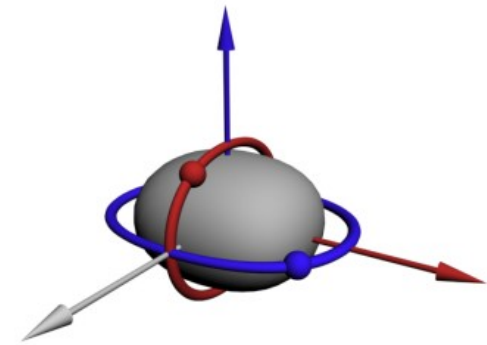
T. Suzuki et al.,
Phys. Rev. C78 (2008) 031302(R)

Different-j shell configuration
 $\pi g_{9/2} \otimes \nu h_{11/2}$

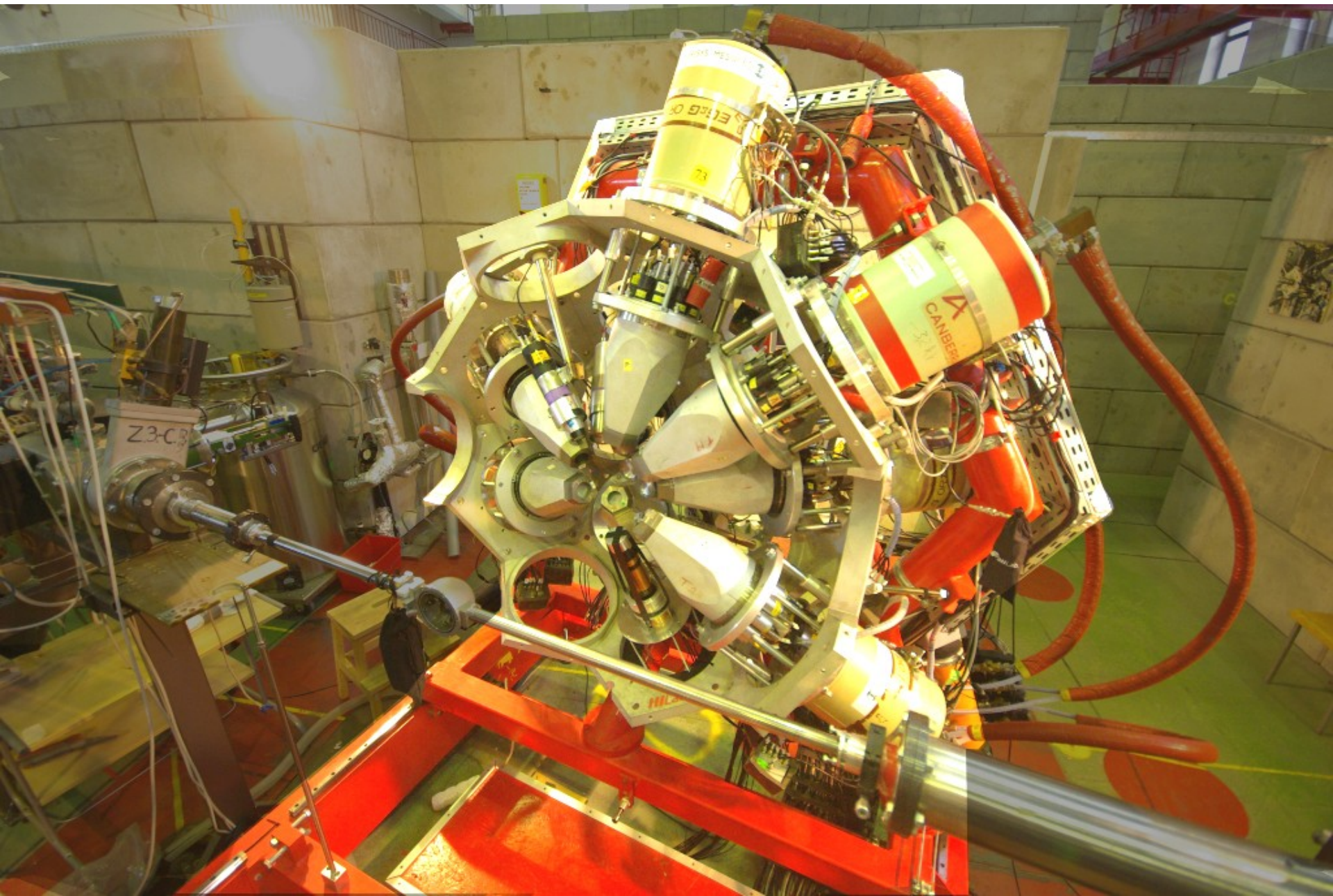


Ch. Droste, S.G. Rohozinski,
 K. Starosta, L. Próchniak
Sequential CQPC calculation

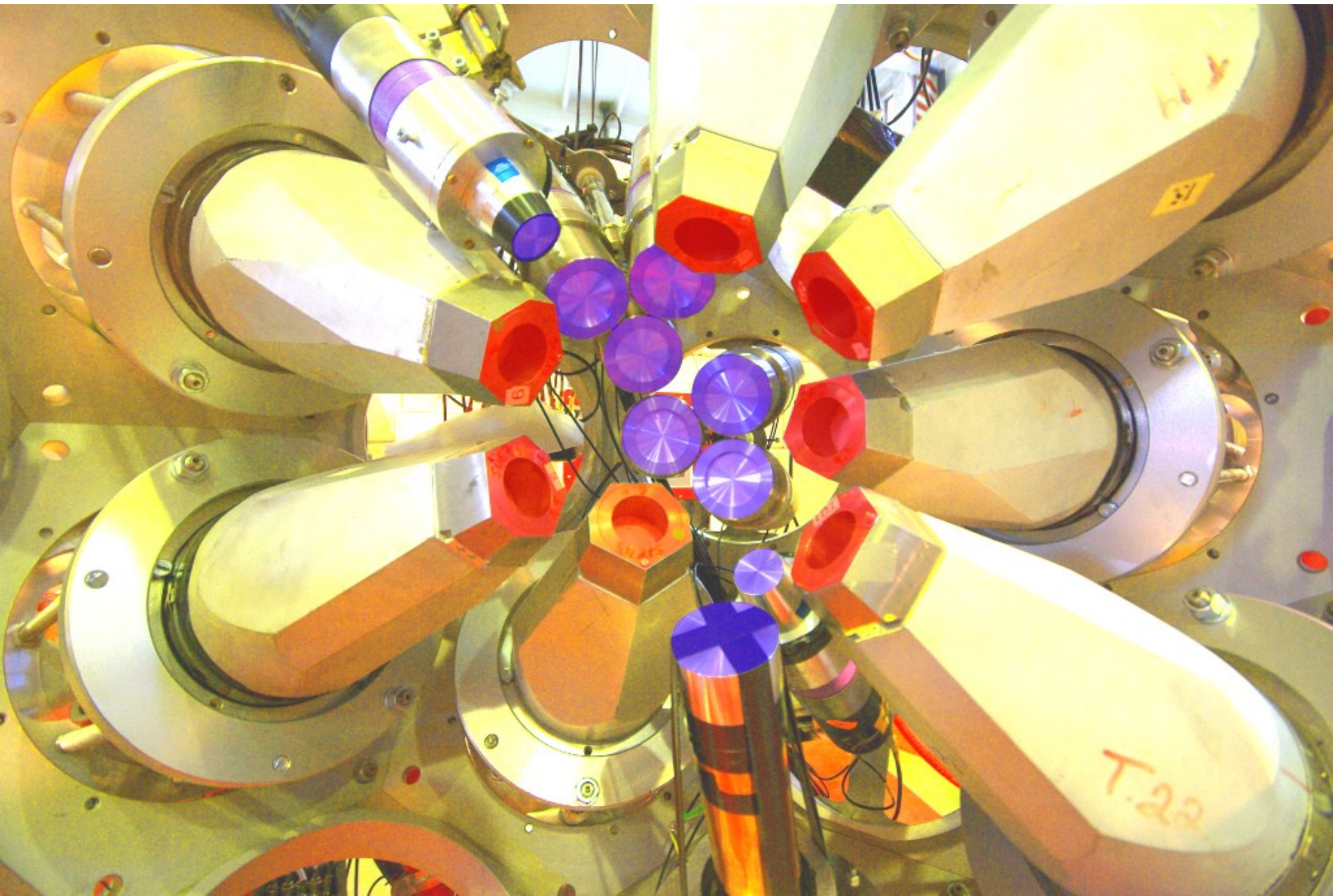
$$\begin{aligned}
 g &= \frac{1}{2} (g_p + g_n + g_R) \\
 &+ \frac{1}{J(J+1)} \frac{1}{2} j_p (j_p + 1) (g_p - g_n - g_R) \\
 &+ \frac{1}{J(J+1)} \frac{1}{2} j_n (j_n + 1) (g_n - g_p - g_R) \\
 &+ \frac{1}{J(J+1)} \frac{1}{2} j_R (j_R + 1) (g_R - g_p - g_n) \\
 &- \frac{1}{J(J+1)} (g_p \vec{j}_n \cdot \vec{j}_R + g_n \vec{j}_p \cdot \vec{j}_R + g_R \vec{j}_p \cdot \vec{j}_n)
 \end{aligned}$$



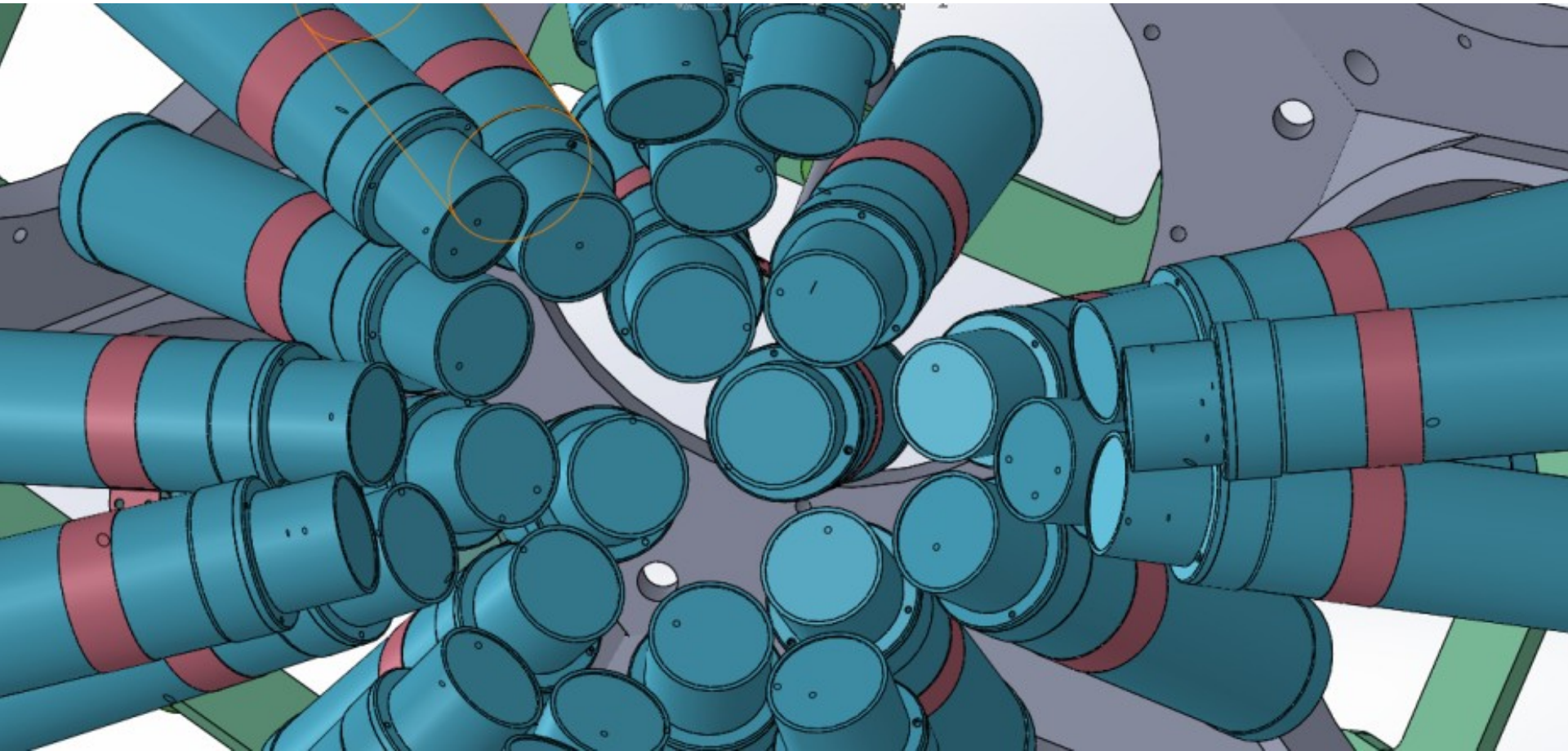
EAGLE EYE – sEquential gammaraYs dEtection



EAGLE EYE – sEquential gammaraYs dEtection



EAGLE EYE – sEquential gammaraYs dEtECTION
Heavy Ion Laboratory,
University of Warsaw + National Centre for Nucl. Research



$$[R_Y T, M] = 0$$

Doubling of red. Matrix elements,

Energies, transition probabilities

Energy-signature staggerings,

- Flat behaviour – chiral rotation
- Opposite staggering – less chiral character

Chiral bands talk – opposite values,

- Inband mat. element large
- Interband mat. element small

$$\langle I_2, + \| M \| I_1, + \rangle = \text{Re} \langle I_2, L \| M \| I_1, L \rangle$$

$$\langle I_2, + \| M \| I_1, - \rangle = \text{Im} \langle I_2, L \| M \| I_1, L \rangle$$

EM selection rules – in a subset of chiral nuclei

- single-j configuration
- 2qp states
- Triaxiality $\gamma=30^\circ$

Volume of 3 ang. Momentum vectors

- g-factor measurements