

# Triaxiality and single particle degrees of freedom

S. Frauendorf

Department of Physics  
University of Notre Dame  
USA



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Thank's to my collaborators

Weichuan Li, Michigan State University, USA

Javid Sheikh, Gowar Bhat, Sheikh Jangir  
Kashmir University, India

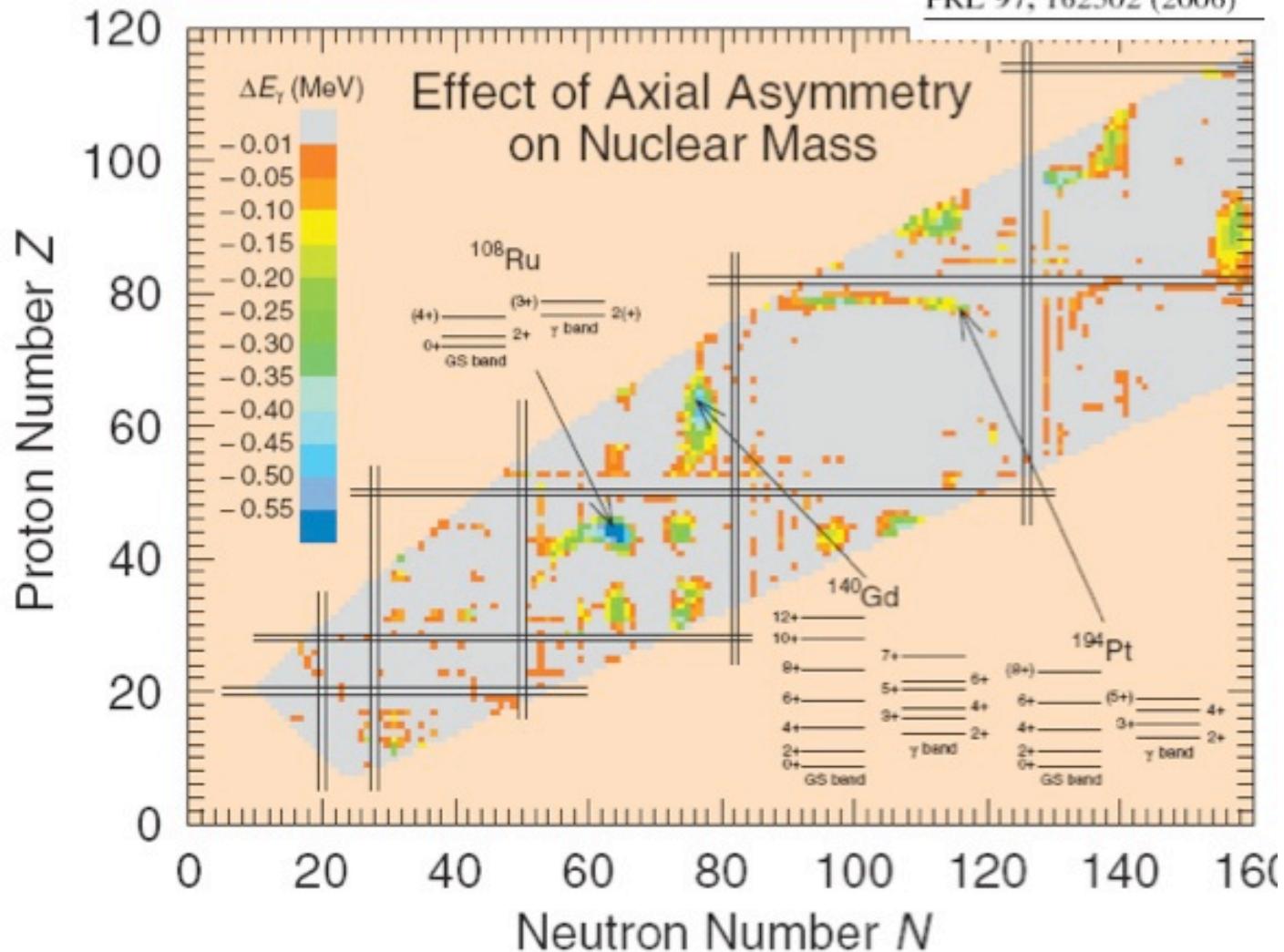
Quibo Chen, Peking University, China

# Triaxiality at moderate Spin

Peter Möller,<sup>1,\*</sup> Ragnar Bengtsson,<sup>2</sup> B. Gillis Carlsson,<sup>2</sup> Peter Olivius,<sup>2</sup> and Takatoshi Ichikawa<sup>3</sup>

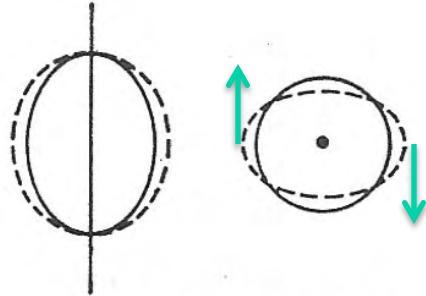
PRL 97, 162502 (2006)

enhanced  
triaxiality  
 $E(2_2^+) < E(4_1^+)$

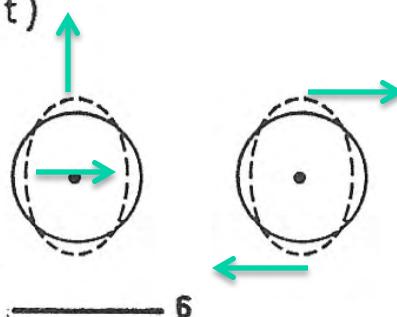


states with  $\langle Nn_3\Lambda\Omega | Y_{2\pm 2} | Nn_3\Lambda \pm 2\Omega \pm 2 \rangle$  near the Fermi surface

- Triaxiality rigid or dynamical?
- Is it the right question?
- Quasi gamma band in even-even nuclei
- Transverse wobbling

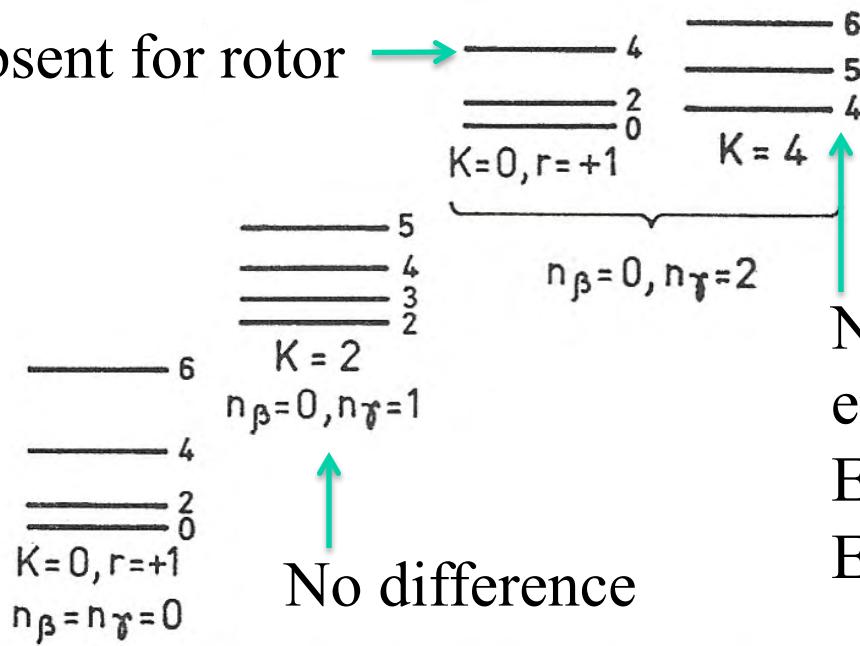


$\nu = \pm 2, \gamma\text{-Vibration}$   
 $\delta R \propto \sin^2 \theta \cos(2\phi \pm \omega t)$



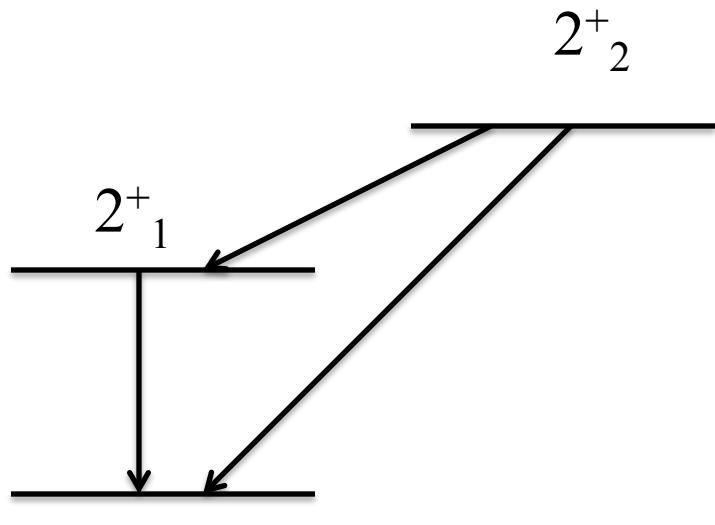
## Triaxial rotor vs. $\gamma$ vibrator

Absent for rotor



No difference  
 energy different  
 $E(K=4)=2E(K=2)$  vibrator  
 $E(K=4)>2E(K=2)$  rotor

# Triaxial rotor description of the $\gamma$ band

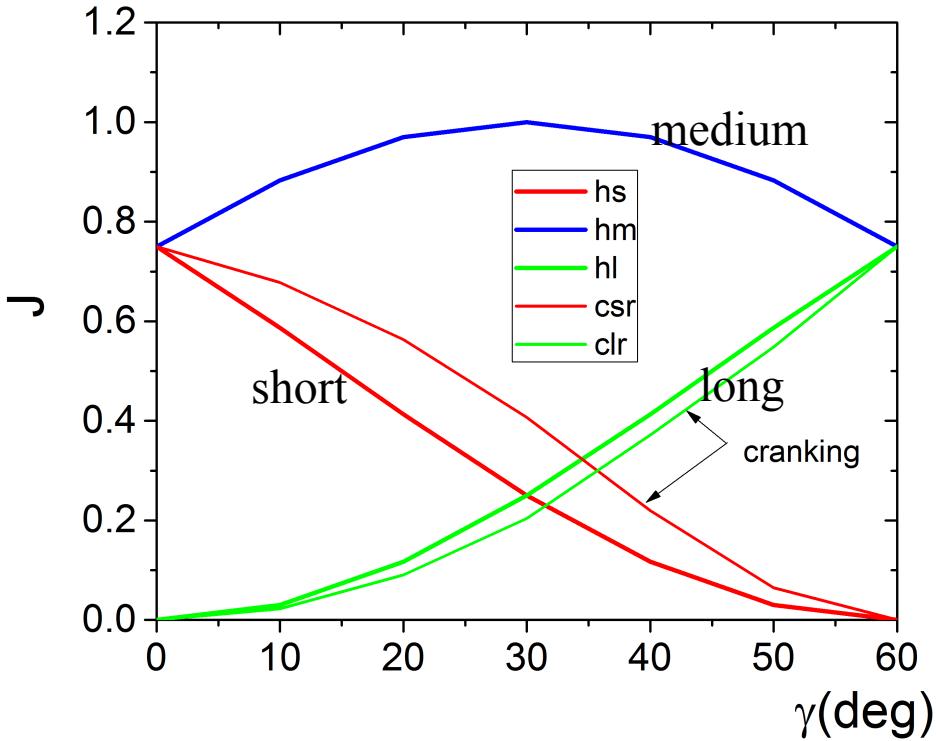


One rotational quant along  
the unfavored axes

J.M. Allmond , J.L. Wood  
PLB 767 (2017) 226

$\gamma$  from COULEX  
 $J_k$  from energy

calculated

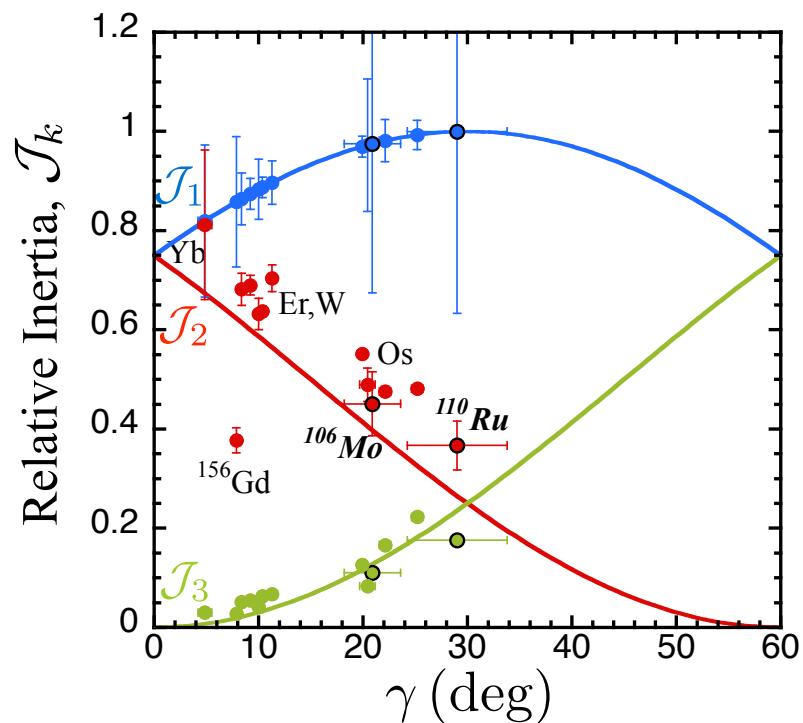


cranking calculation  $Z=68$   $A=168$

$$\varepsilon = 0.25 \quad \Delta_p = 1.1 \text{ MeV} \quad \Delta_n = 1.0 \text{ MeV}$$

Triaxial rotor description of the  $\gamma$  band

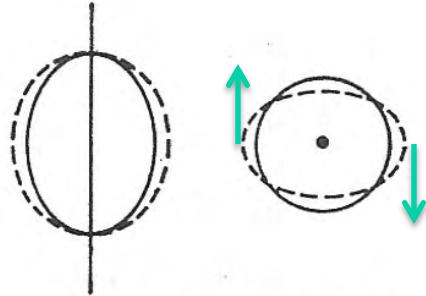
experimental



J.M. Allmond , J.L. Wood  
PLB 767 (2017) 226

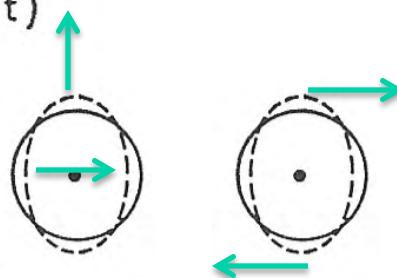
$\gamma$  from COULEX  
 $J_k$  from energy

The three moments of inertia are irrotational flow-like:  
Consequence of quantum coherence among Fermions

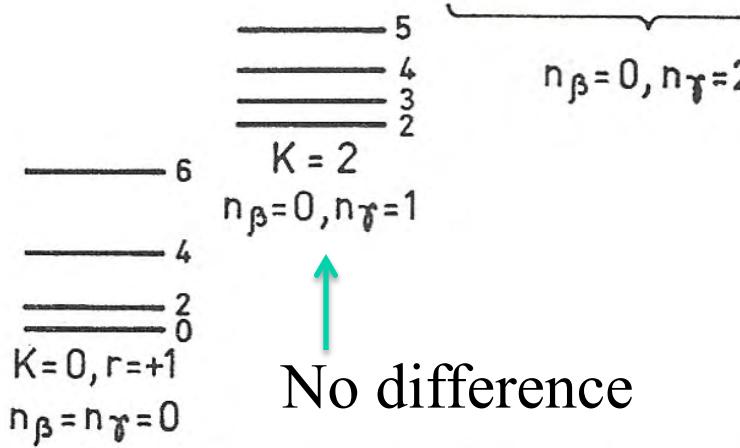


$\nu = \pm 2$ ,  $\gamma$ -Vibration

$$\delta R \propto \sin^2 \theta \cos(2\phi \pm \omega t)$$



Pulsation mode  
Absent for rotor



## Triaxial rotor vs. $\gamma$ vibrator

Travelling wave  $\leftrightarrow$  rigid rotation  
energy different  
 $E(K=4) = 2E(K=2)$  vibrator  
 $E(K=4) > 2E(K=2)$  rotor

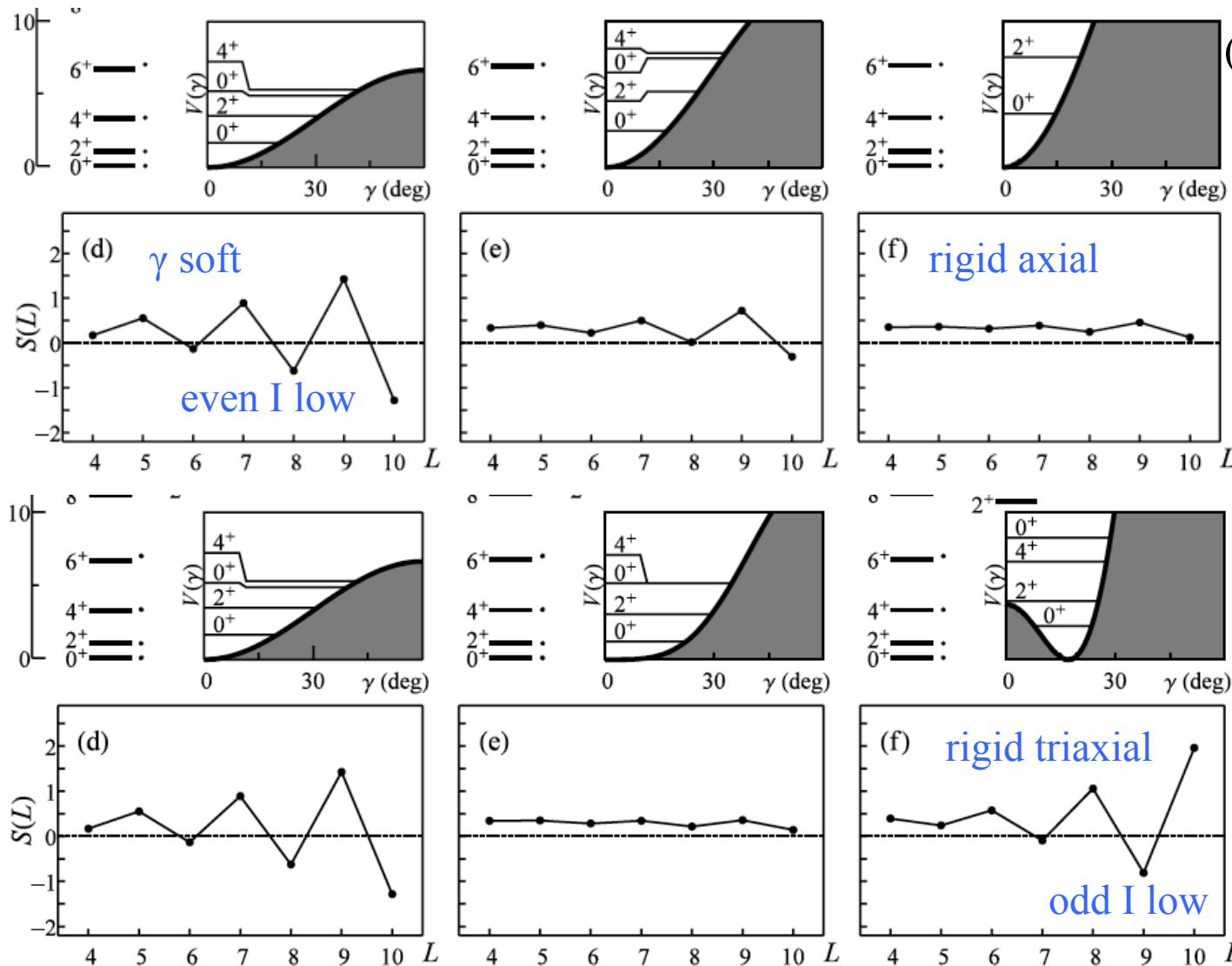
# Staggering of the $\gamma$ band: soft vs. or rigid Bohr-Hamiltonian with irrotational flow inertia

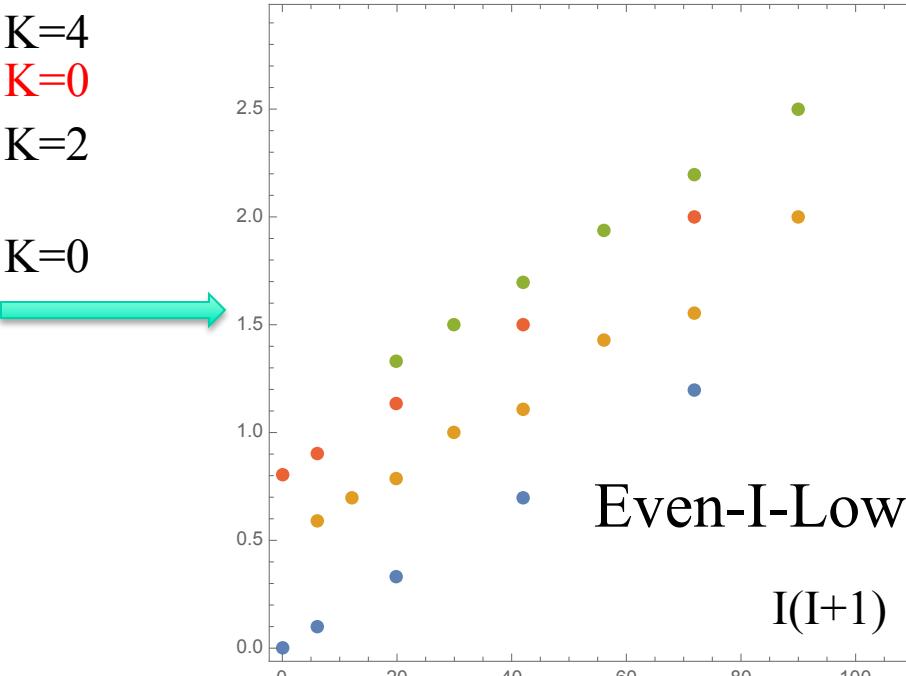
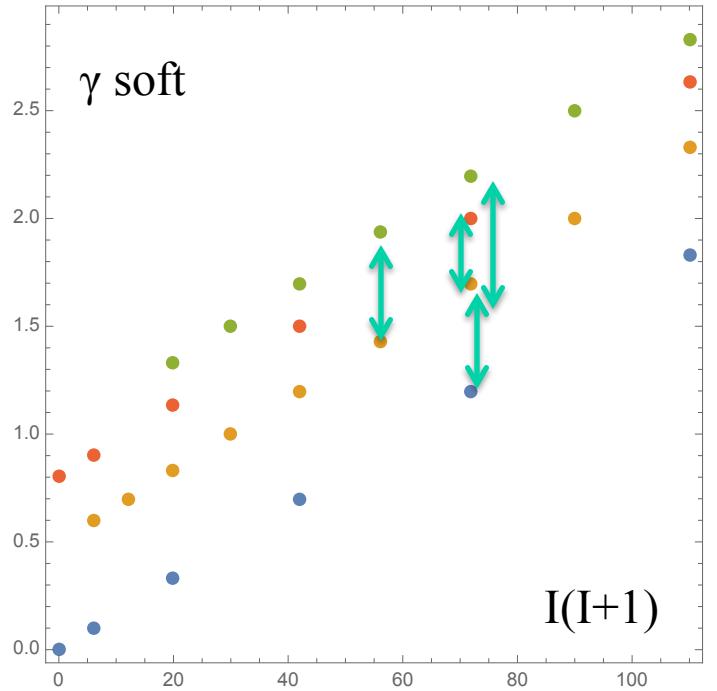
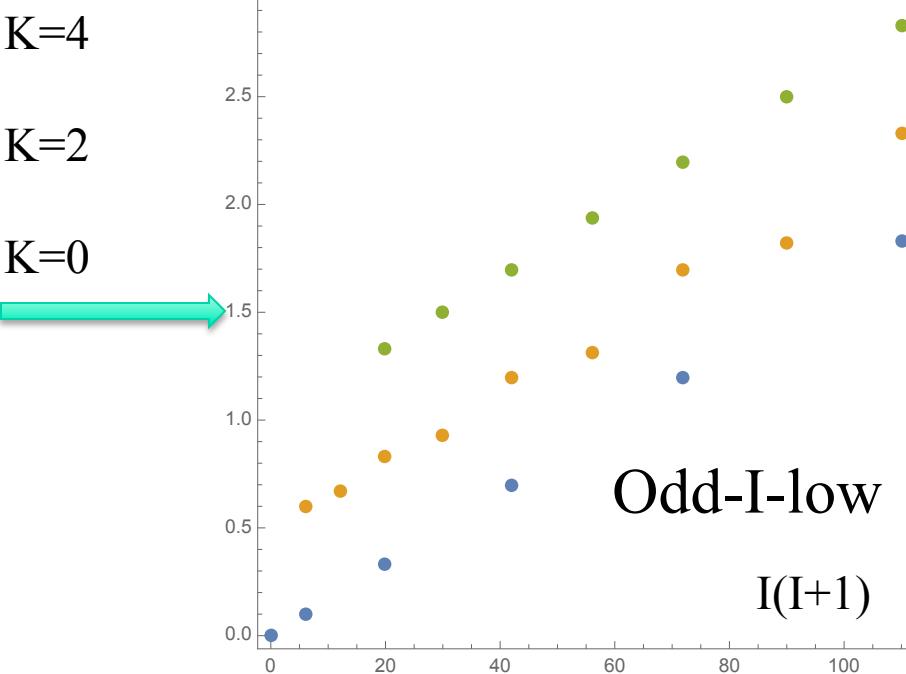
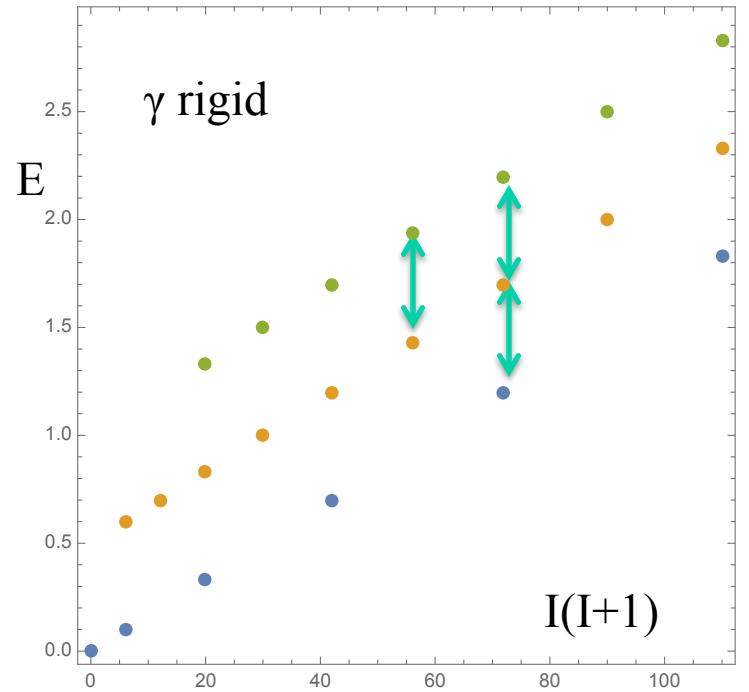
$$S(I) = E(I) - \frac{(E(I-1) + E(I+1))}{2}$$

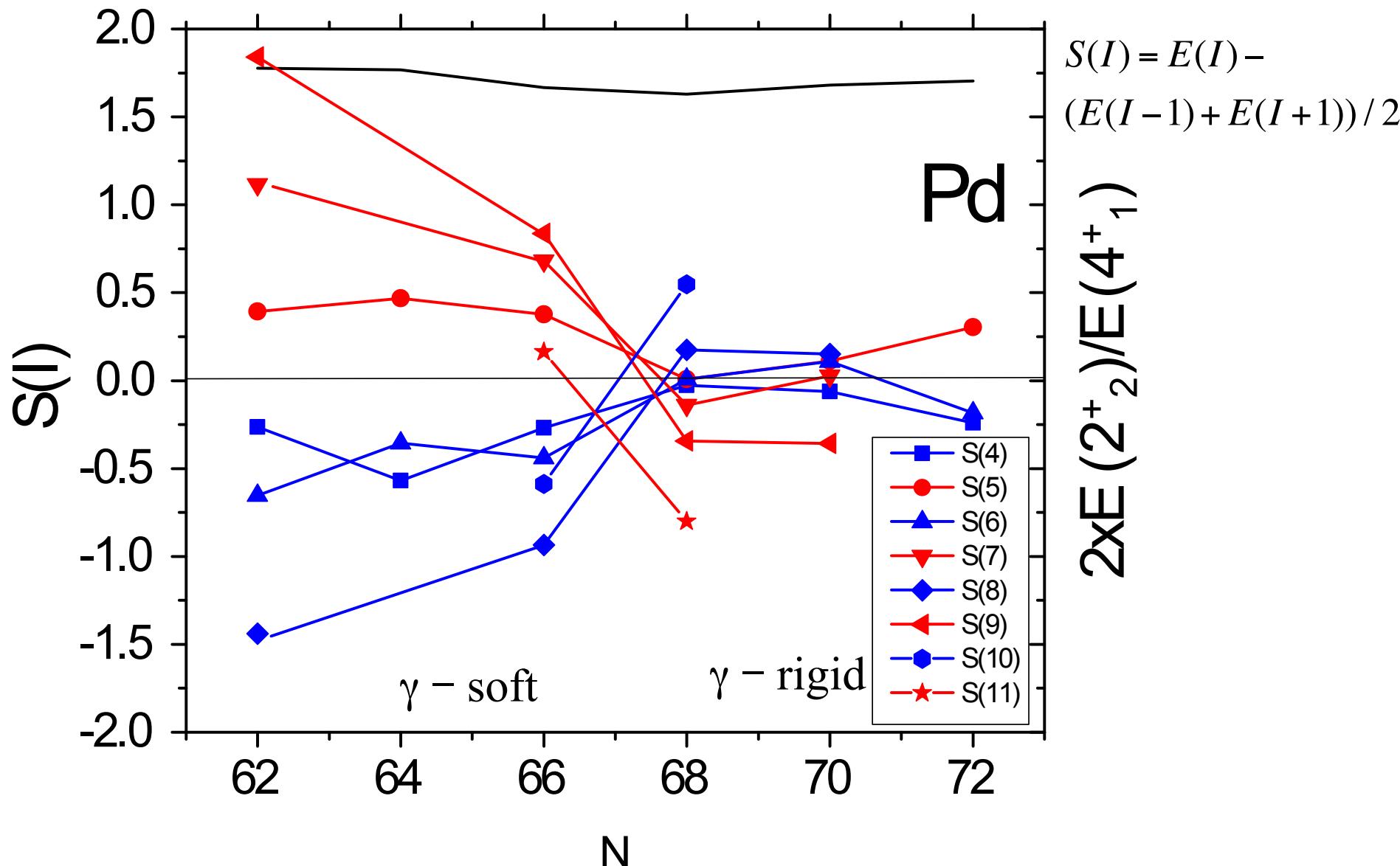
M. A. Caprio  
PRC 83,  
064309 (2011)

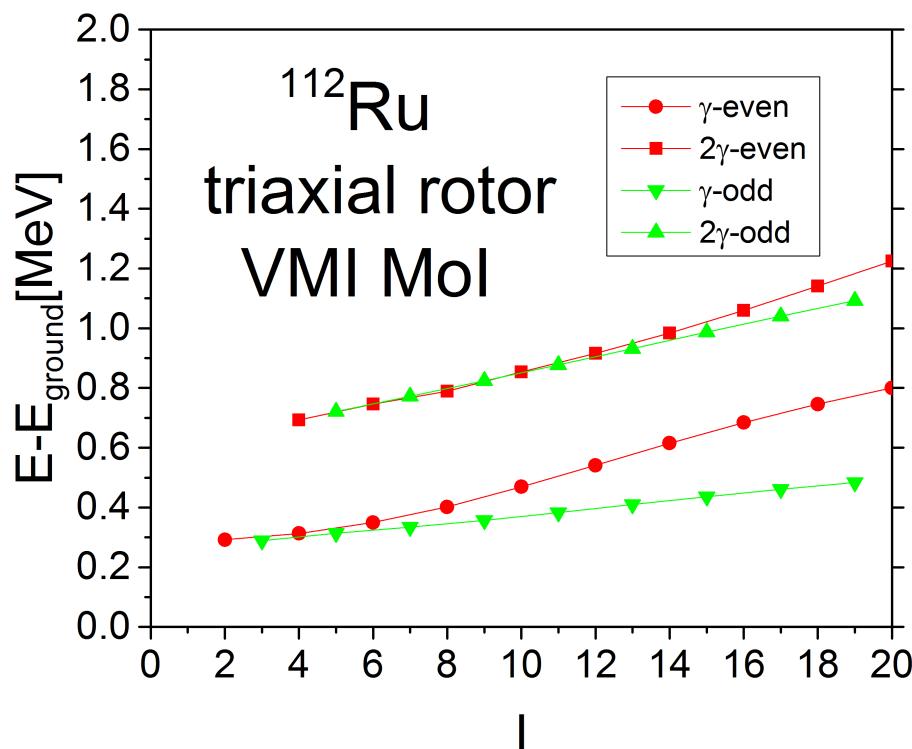
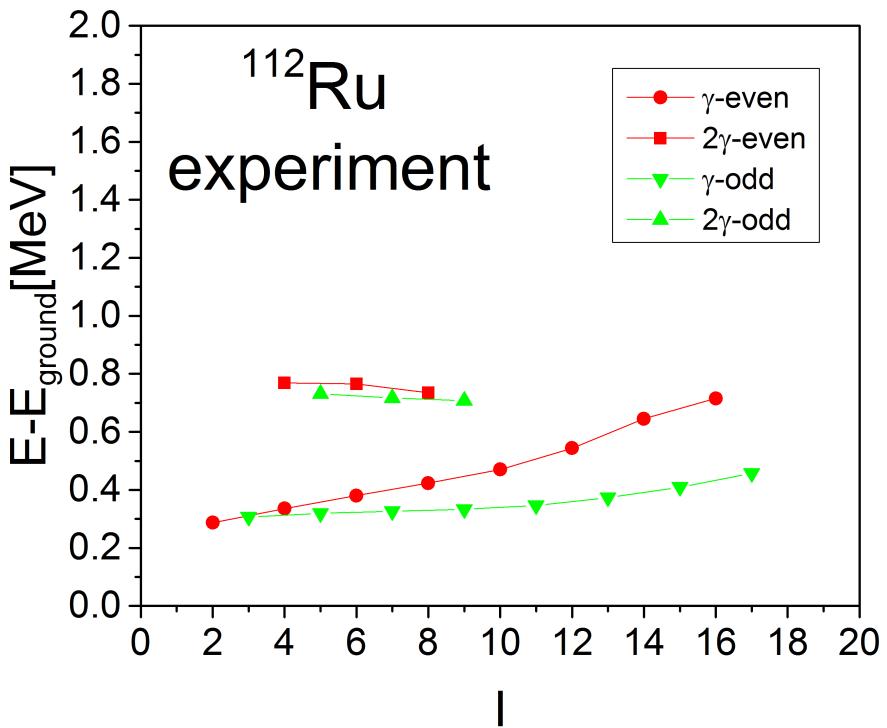
McCutchan et al.,  
PRC 76  
(2007) 024306.  
IBM

Zamfir,Casten  
PRL 87  
(2001) 052503









rigid triaxial rotor

$$4 < E(4_3^+)/E(2_2^+) < 3.333$$

weak triaxiality      strong

$$\text{exp: } E(4_3^+)/E(2_2^+) = 2.8$$

$$\Theta_i = \Theta_{0i} + \Theta_{1i}I$$

*I*-dependence of  
 $\gamma$  deformation

There should be some  
 $\gamma$  softness.

States are a superposition of basis

$$\begin{aligned}
 & \hat{P}_{MK}^I |\Phi\rangle; \\
 & \hat{P}_{MK}^I a_{p_1}^\dagger a_{p_2}^\dagger |\Phi\rangle; \\
 & \hat{P}_{MK}^I a_{n_1}^\dagger a_{n_2}^\dagger |\Phi\rangle; \\
 & \hat{P}_{MK}^I a_{p_1}^\dagger a_{p_2}^\dagger a_{n_1}^\dagger a_{n_2}^\dagger |\Phi\rangle; \\
 & \hat{P}_{MK}^I a_{n_1}^\dagger a_{n_2}^\dagger a_{n_3}^\dagger a_{n_4}^\dagger |\Phi\rangle; \\
 & \hat{P}_{MK}^I a_{p_1}^\dagger a_{p_2}^\dagger a_{p_3}^\dagger a_{p_4}^\dagger |\Phi\rangle,
 \end{aligned} \tag{1}$$

Triaxial rotor dynamics

Single particle response

where  $|\Phi\rangle$  is the vacuum state and the three-dimensional angular-momentum projection operator [29] is given by

$$\hat{P}_{MK}^I = \frac{2I+1}{8\pi^2} \int d\Omega D_{MK}^I(\Omega) \hat{R}(\Omega), \tag{2}$$

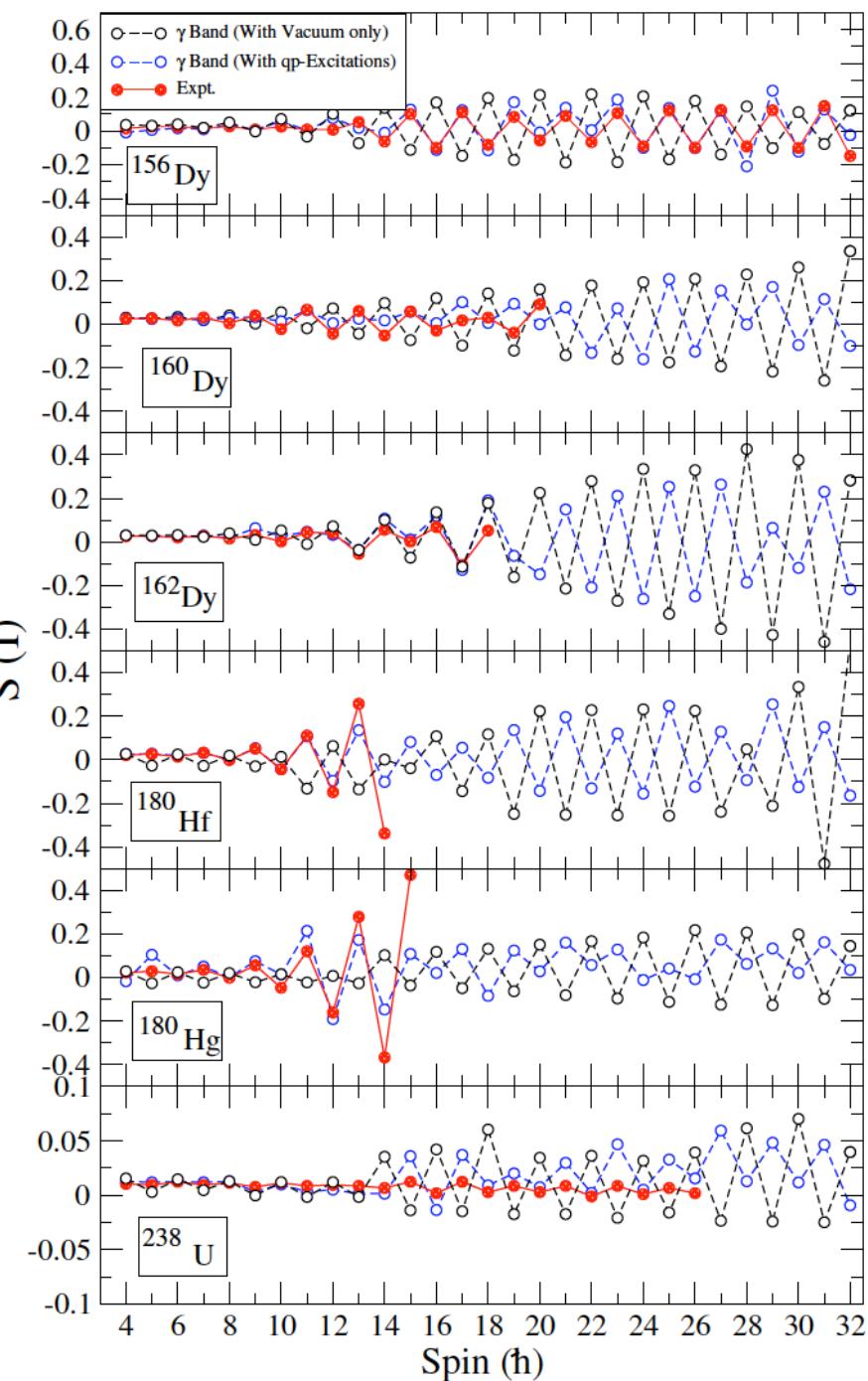
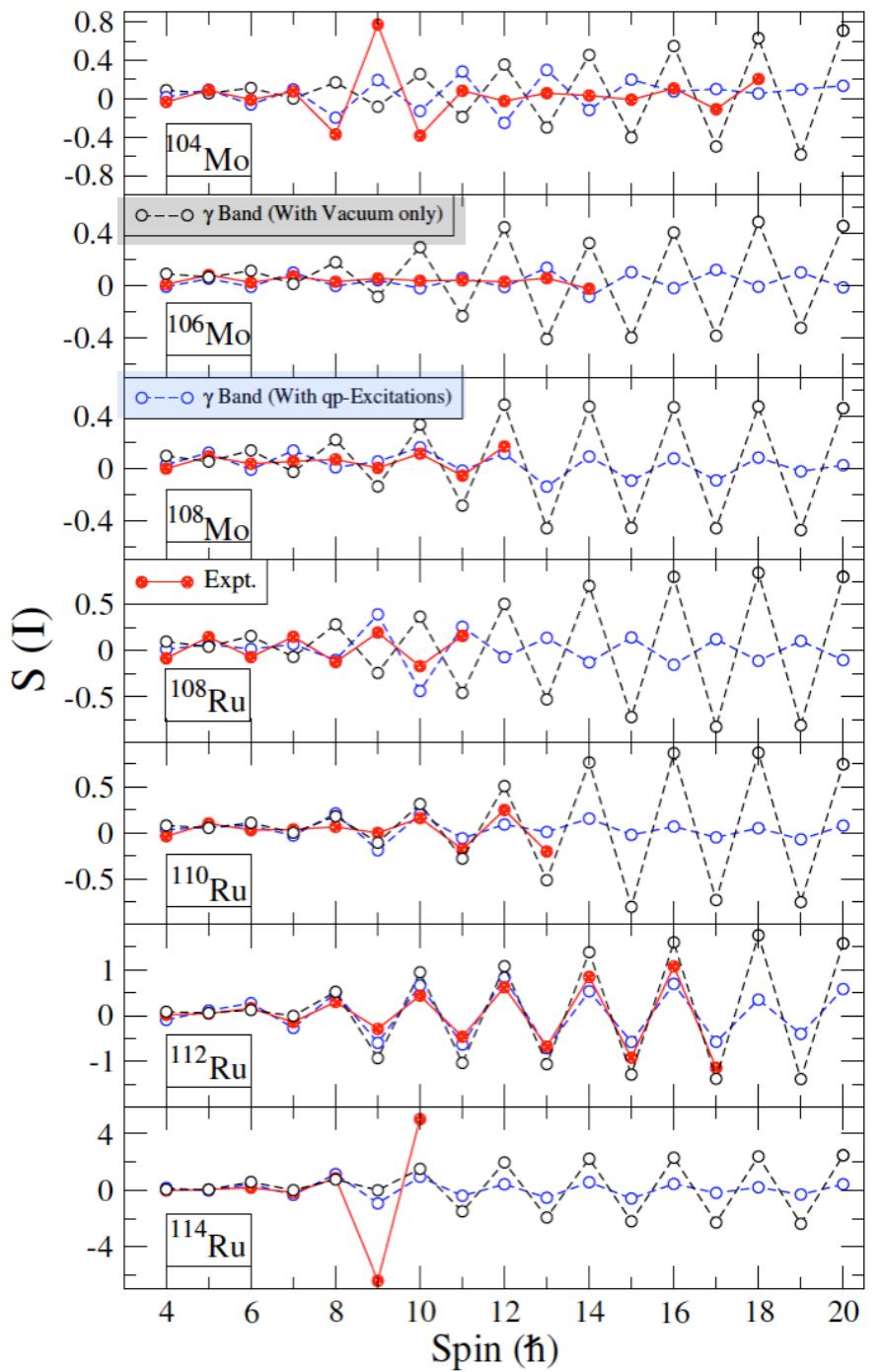
The projected basis of Eq. 1 is then used to diagonalise the shell model Hamiltonian. As in our earlier studies, we have employed the pairing plus quadrupole-quadrupole Hamiltonian [29, 32–34]

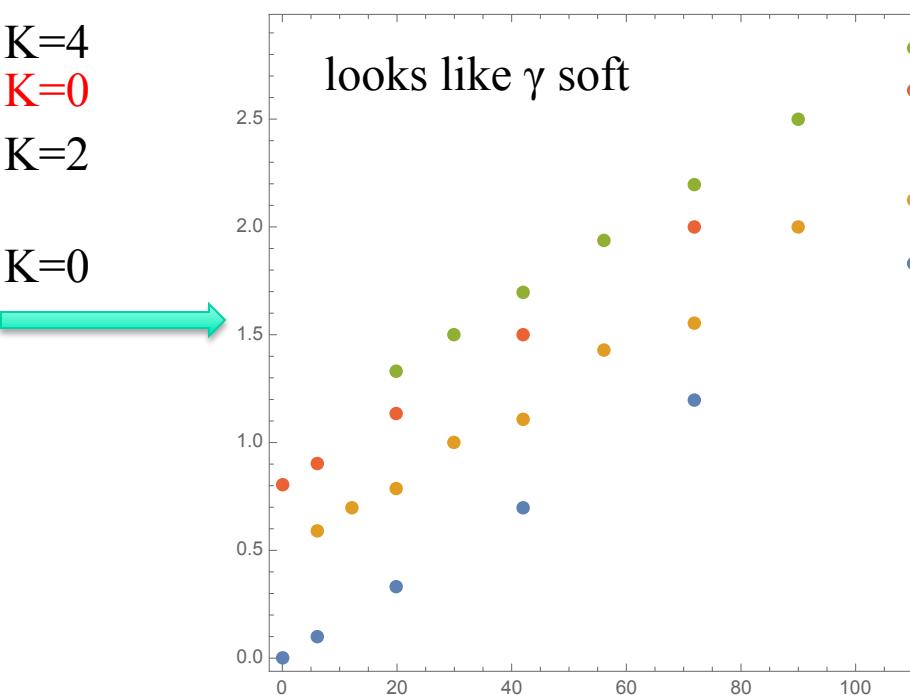
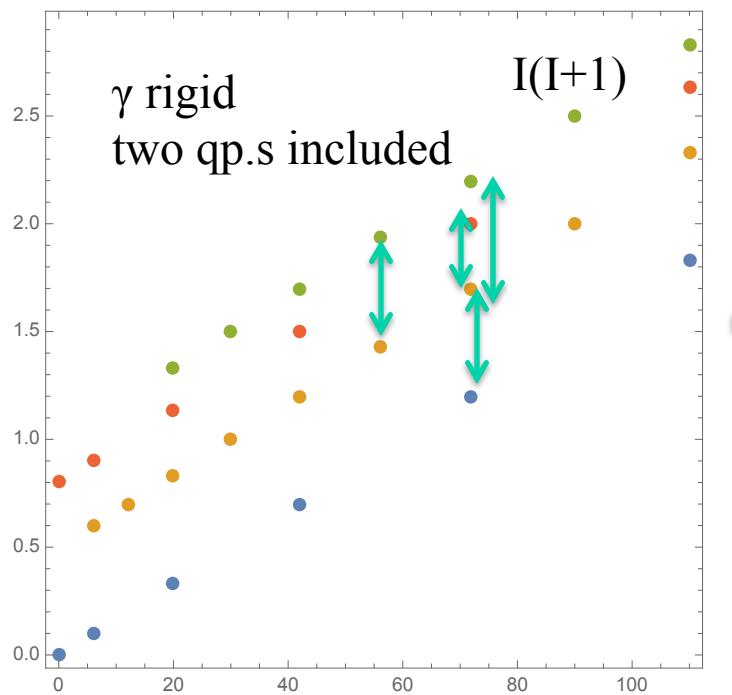
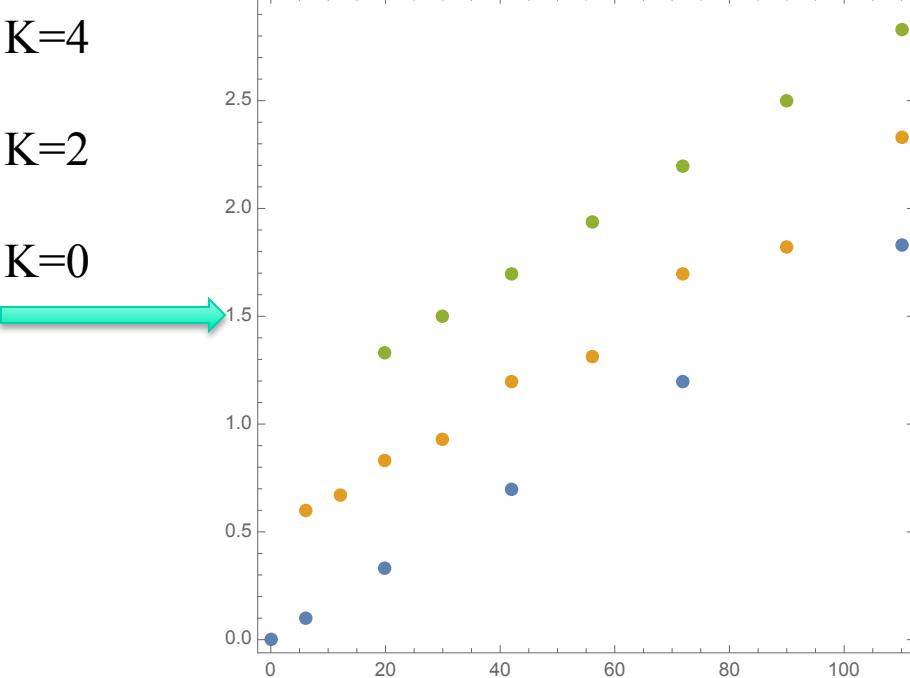
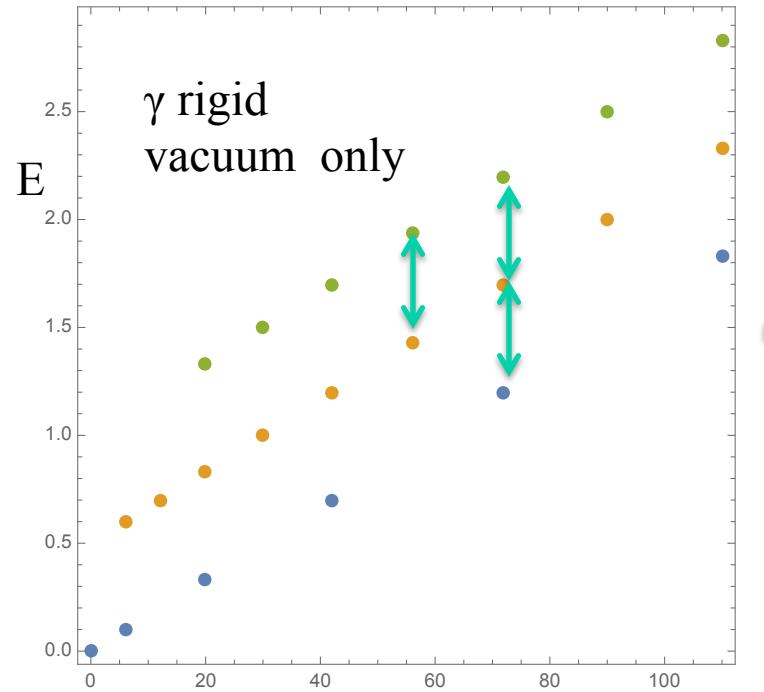
$$\hat{H} = \hat{H}_0 - \frac{1}{2}\chi \sum_\mu \hat{Q}_\mu^\dagger \hat{Q}_\mu - G_M \hat{P}^\dagger \hat{P} - G_Q \sum_\mu \hat{P}_\mu^\dagger \hat{P}_\mu. \tag{3}$$

The  $QQ$ -force strength  $\chi$  is adjusted such that the physical quadrupole deformation  $\varepsilon$  is obtained as a result of the self-consistent mean-field HFB calculation [29]. The monopole

## Triaxial Projected Shell Model

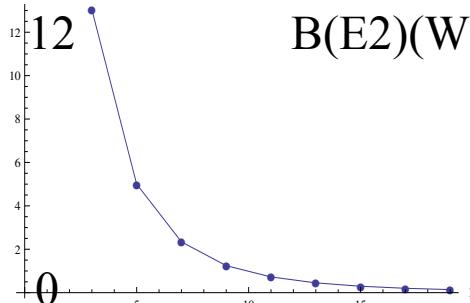
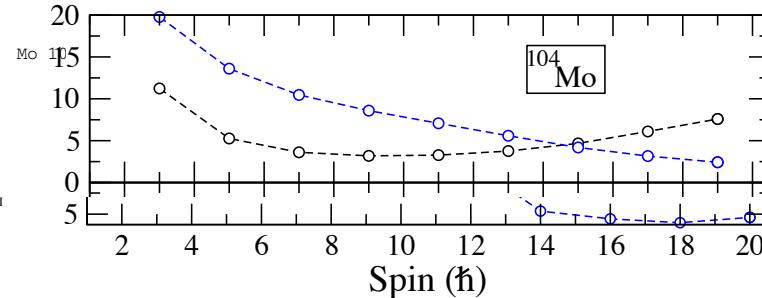
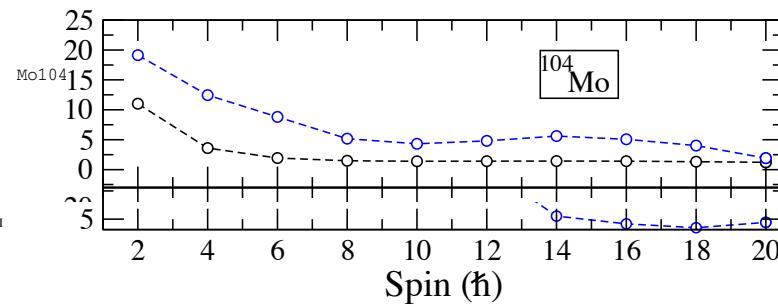
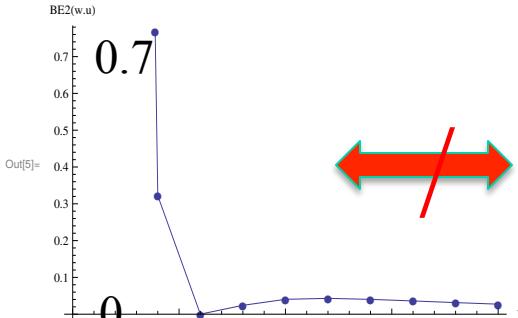
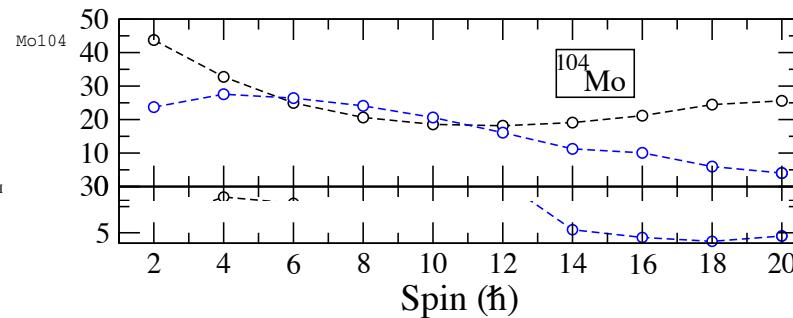
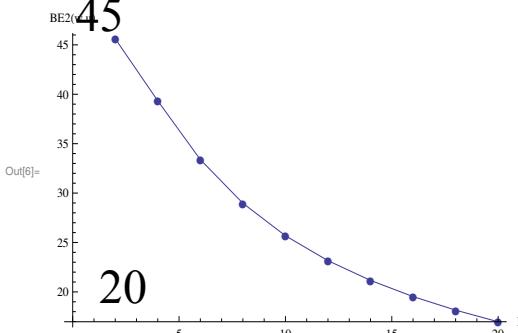
The triaxiality parameter is fixed!

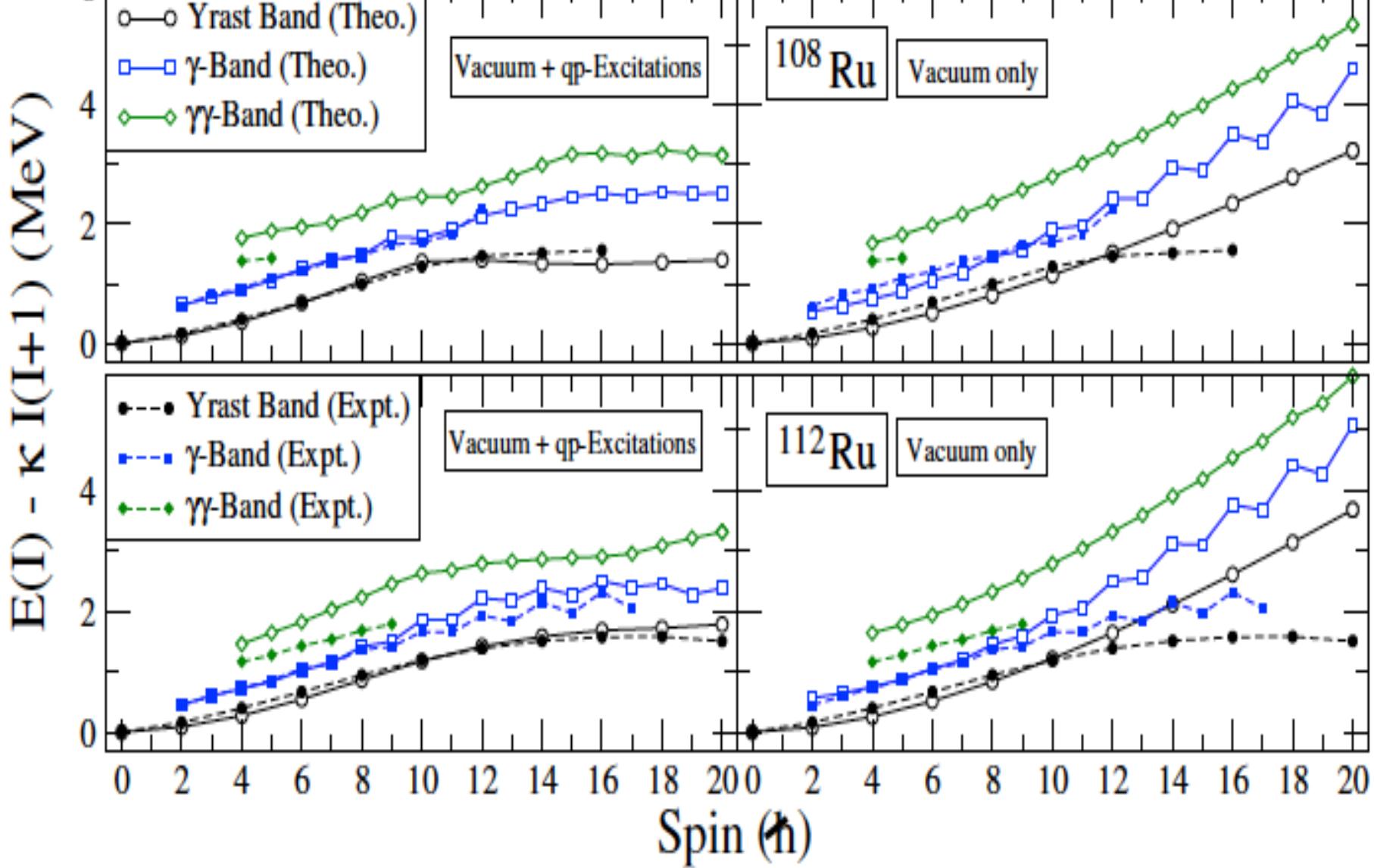






BE2(w.u)

**B(E2)(WU)****I->I-1****I->I-2****I->I****TPSM B(E2) from gamma to ground band**



Both “best triaxial rotor” 112 and “softish” 108 reproduced with the same parameters,  $\gamma\gamma$ -band too high with fixed  $\gamma$  deformation

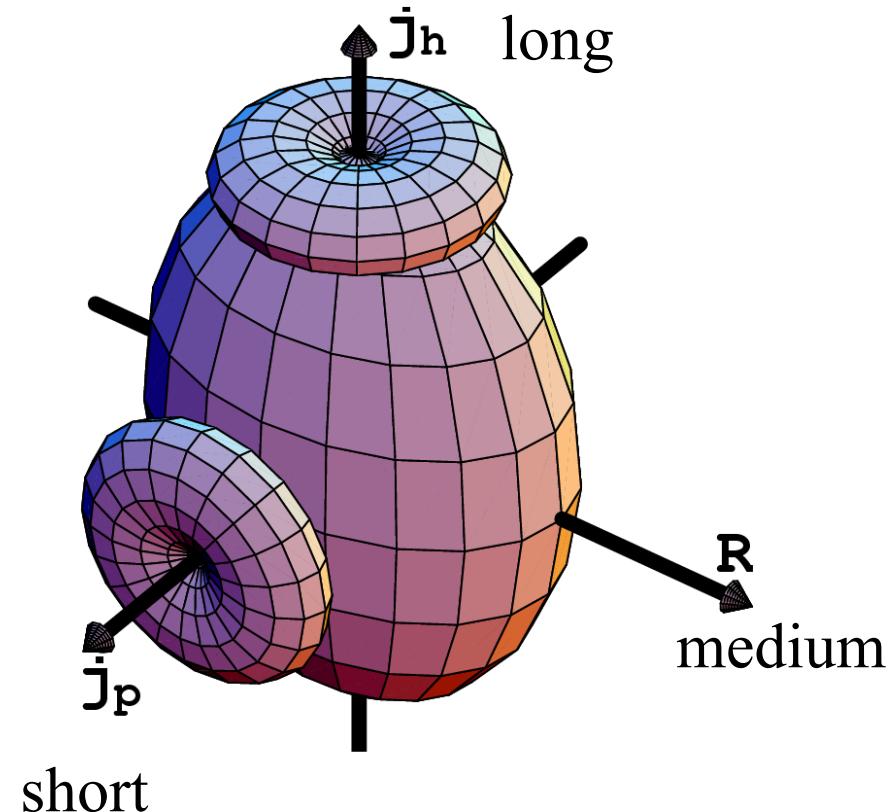
# Even-even nuclei: $\gamma$ rigid or soft?

- Bohr-Hamiltonian relates staggering of  $\gamma$  band with softness of triaxiality
- TSPM with fixed  $\gamma$  accounts for the different signs of staggering
- Energy staggering alone does not discriminate
- Differences in  $I \rightarrow I-2$  interband transitions
- More careful study of matrix elements from Coulex experiments on  $^{104}\text{Ru}$  on the way
- Energy of  $\gamma\gamma$  band indicates  $\gamma$  softness

# Coupling of unpaired quasiparticles with the triaxial core

Combination of  
unpaired  
quasiparticles with  
Collective  
Triaxial Rotor

The players



$$\Theta_m > \Theta_s > \Theta_l$$

# Multi-Quasiparticle+Triaxial Rotor Model

134

*S. Frauendorf, J. Meng/Nuclear Physics A 617 (1997) 131–147*

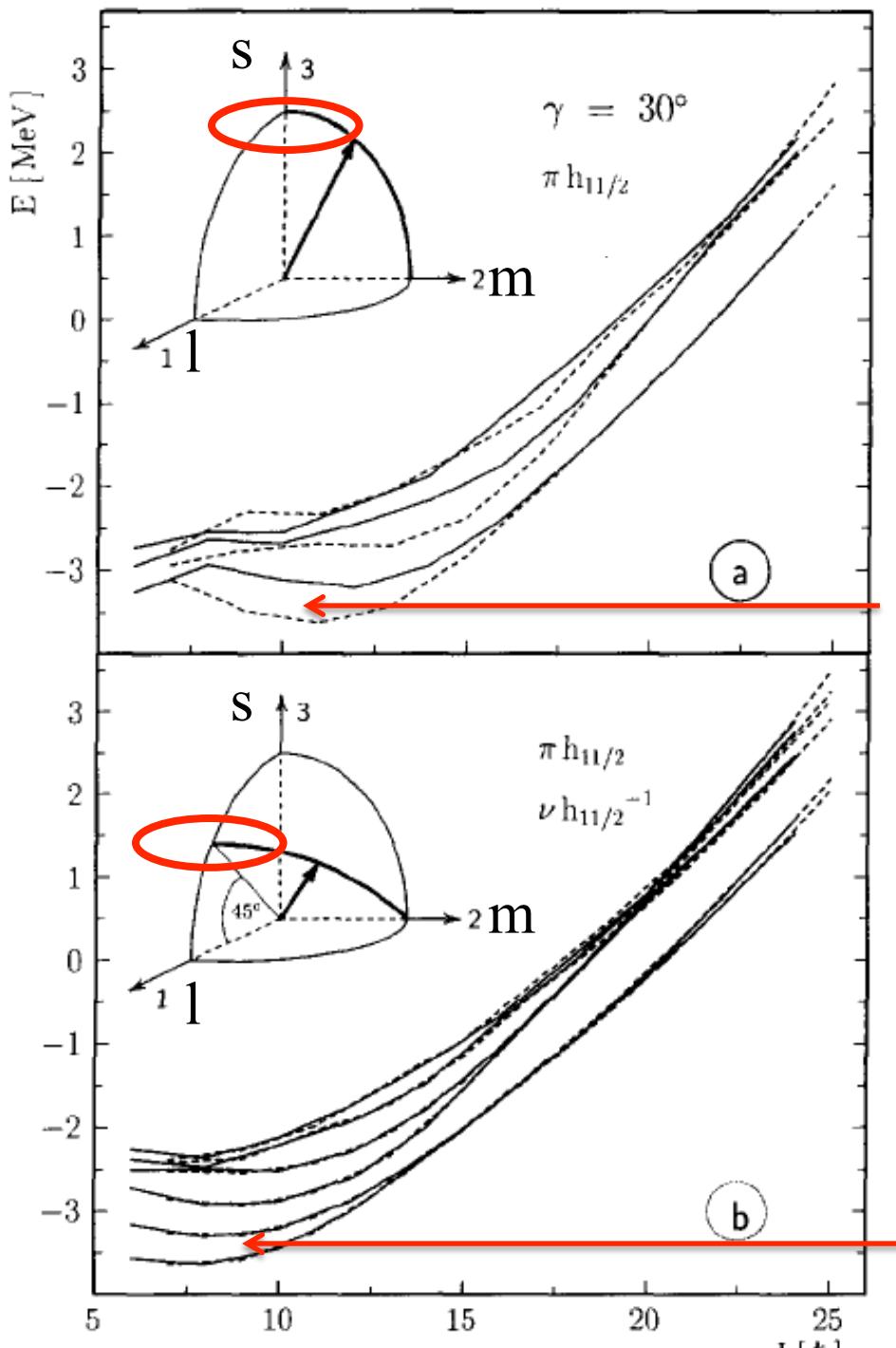
$$H = h_{\text{def}} + \sum_{\nu=1}^3 \frac{(I_\nu - j_\nu)^2}{2\mathcal{J}_\nu}.$$

For the moments of inertia the ratios of irrotational flow are assumed,<sup>2</sup>

$$\mathcal{J}_\nu = \mathcal{J} \sin \left( \gamma - \frac{2\pi}{3}\nu \right)^2.$$

or taken as adjustable

or taken from cranking calculations



Transverse  
Wobbling and  
Chirality  
come together

### Transverse Wobbling

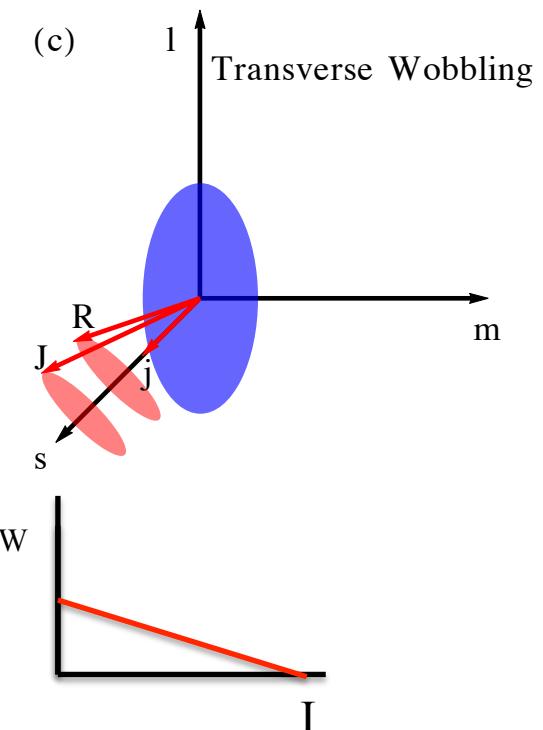
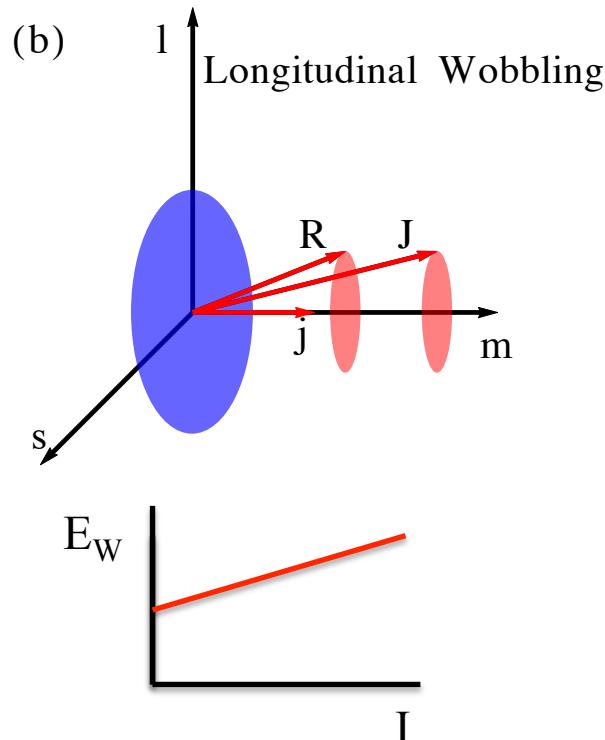
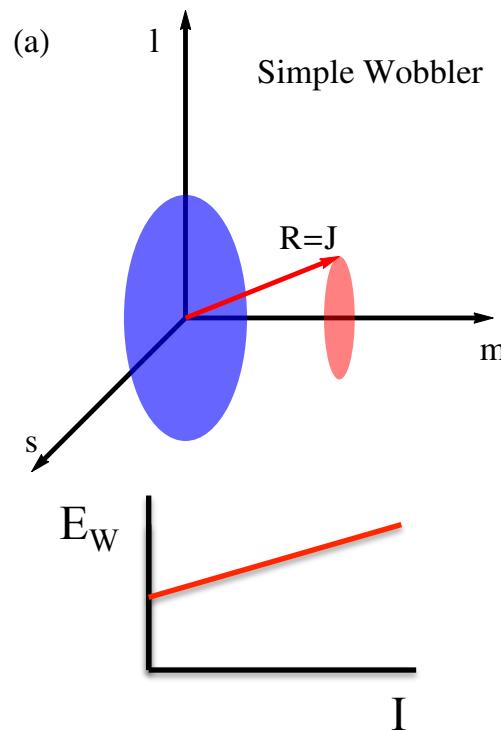
axis of max MoI: m  
transverse W: j perp. m  
longitudinal W: j par. m

$$\Theta_m > \Theta_s = \Theta_l, 4 / 1$$

Transverse Chiral Vibration

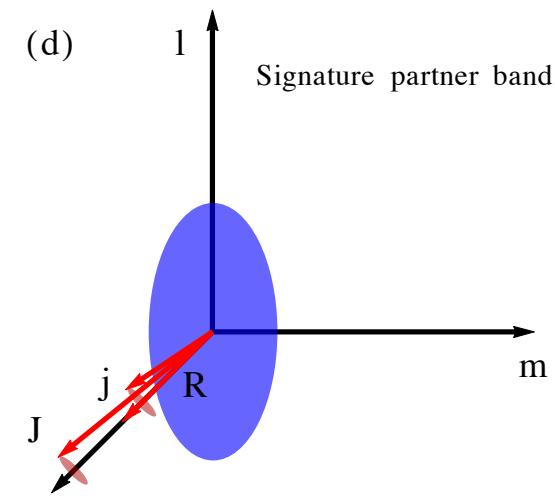
$$\Theta_m > \Theta_s > \Theta_l$$

strong E2, I->I-1,



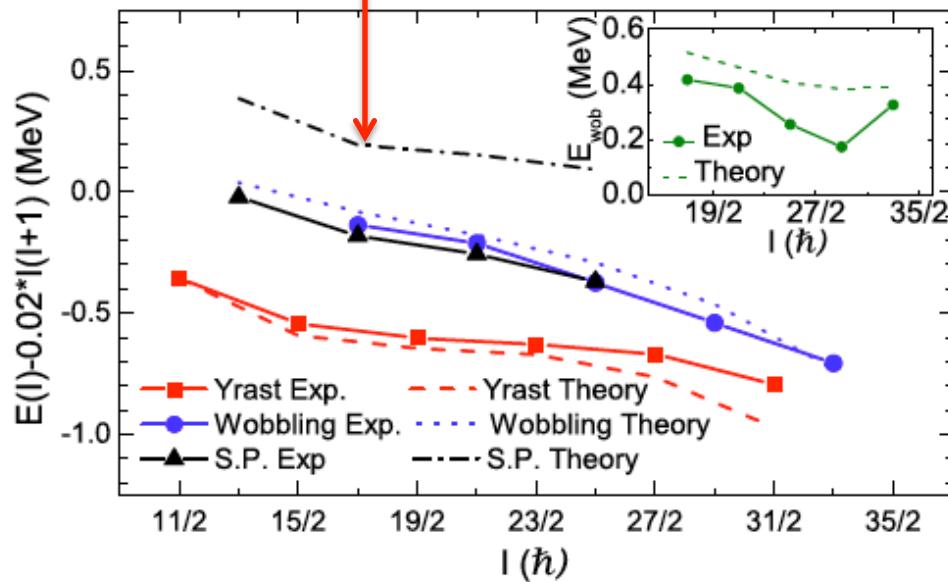
wobbling and  
signature  
partner bands

strong E2 I->I-1,  
weak E2 I->I-1,



Too high!!

J. T. Matta, U. Garg, W. Li, S. Frauendorf, A. D. Ayangeakaa,<sup>†</sup> D. Patel, and K. W. Schlax et al.  
*Physics Department, University of Notre Dame, Notre Dame, Indiana 46556, USA*



$$\pi h_{11/2} + TR$$

$$\Theta_m / \Theta_s / \Theta_l = 23/13/7$$

Signatures for TW:

- Wobbling energy decreases with I
  - Strong collective E2 transitions between 1 and 0 phonon states
- $B(E2_{out}) \sim 0.5 B(E2_{in}) = 50 \text{ WU}$

1 phonon  $\rightarrow$  0 phonon transitions

Initial $I^\pi$	Final $I^\pi$	$E_\gamma$ (keV)	$\delta$	Asymmetry
$\frac{17}{2}^-$	$\frac{15}{2}^-$	747.0	$-1.24 \pm 0.13$	$0.047 \pm 0.012$
$\frac{21}{2}^-$	$\frac{19}{2}^-$	812.8	$-1.54 \pm 0.09$	$0.054 \pm 0.034$
$\frac{25}{2}^-$	$\frac{23}{2}^-$	754.6	$-2.38 \pm 0.37$	...
$\frac{29}{2}^-$	$\frac{27}{2}^-$	710.2	...	...
$\frac{13}{2}^-$	$\frac{11}{2}^-$	593.9	$-0.16 \pm 0.04$	$-0.092 \pm 0.023$

E2 Fraction (%)	$\frac{B(M1_{out})}{B(E2_{in})} \left( \frac{\mu_N^2}{e^2 b^2} \right)$		$\frac{B(E2_{out})}{B(E2_{in})}$	
	Experiment	QTR	Experiment	QTR
$60.6 \pm 5.1$	...	0.213	...	0.908
$70.3 \pm 2.4$	$0.164 \pm 0.014$	0.107	$0.843 \pm 0.032$	0.488
$85.0 \pm 4.0$	$0.035 \pm 0.009$	0.070	$0.500 \pm 0.025$	0.290
...	$\leq 0.016 \pm 0.004$	0.056	$\geq 0.261 \pm 0.014$	0.191

# Experimental data

Band(D):  $\pi h_{11/2}^2 v h_{11/2}^2$

$(30^+)$  12418.5

1141

$(28^+)$  11277.5

1087

$(26^+)$  10190.6

1018

$(24^+)$  9172.1

949

$(22^+)$  8223.2

893

$(20^+)$  7330.5

859

$(18^+)$  6471.9

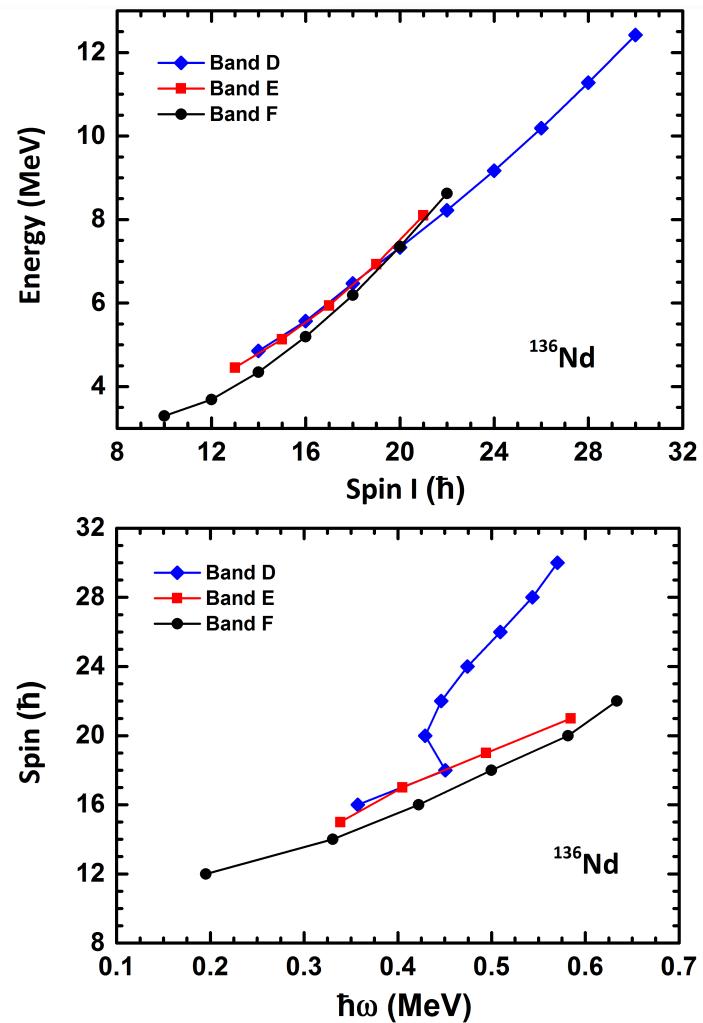
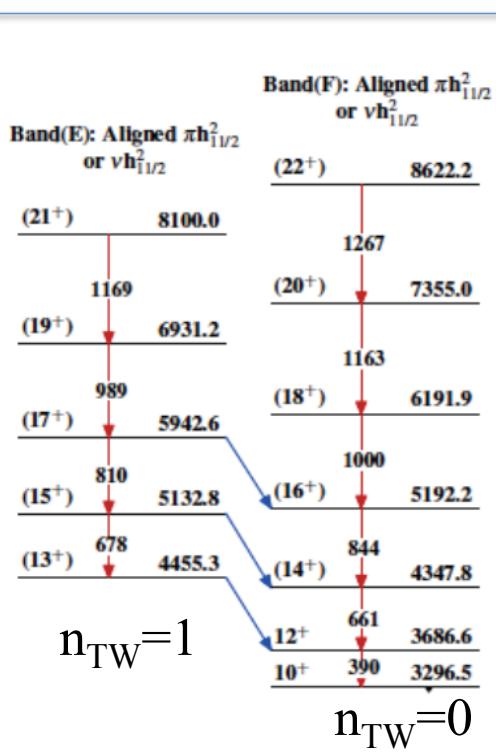
902

$(16^+)$  5570.3

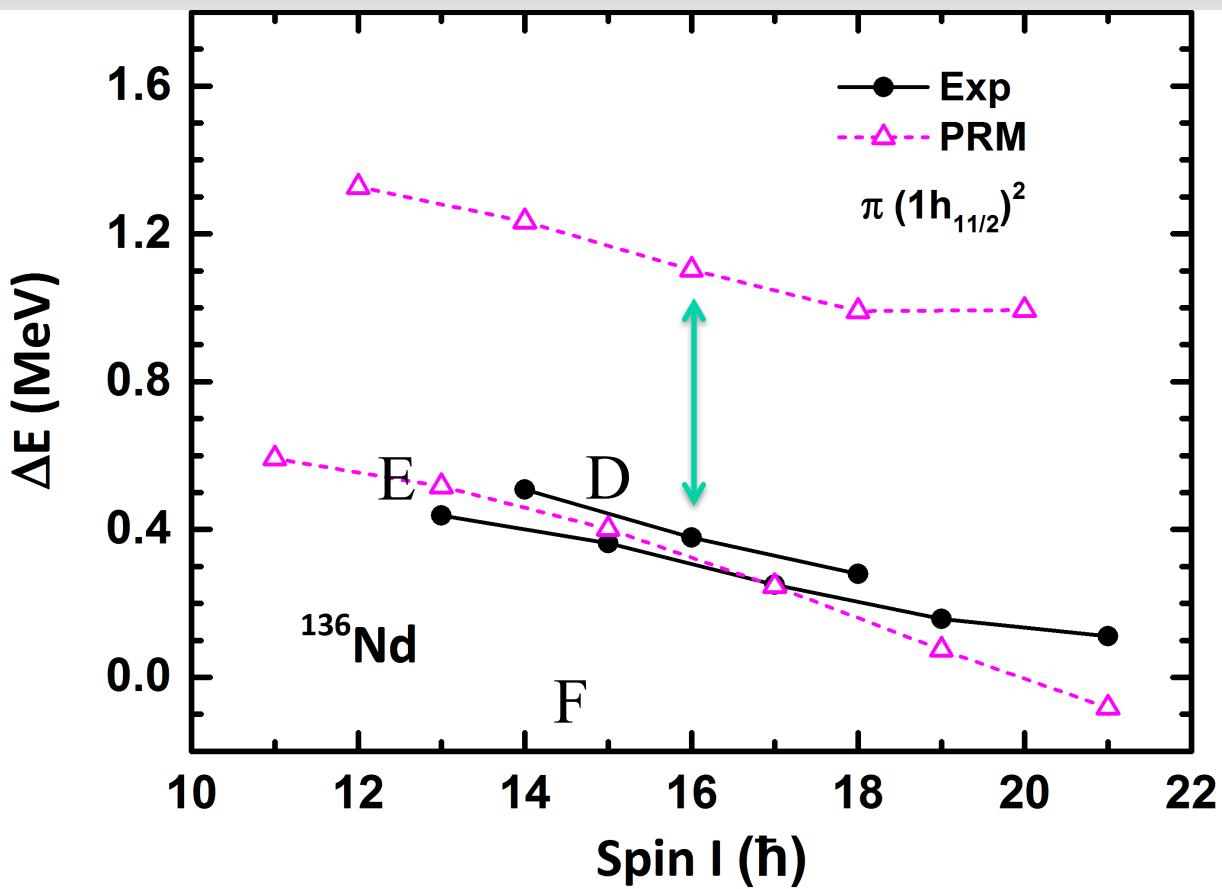
714

$(14^+)$  4855.9

???

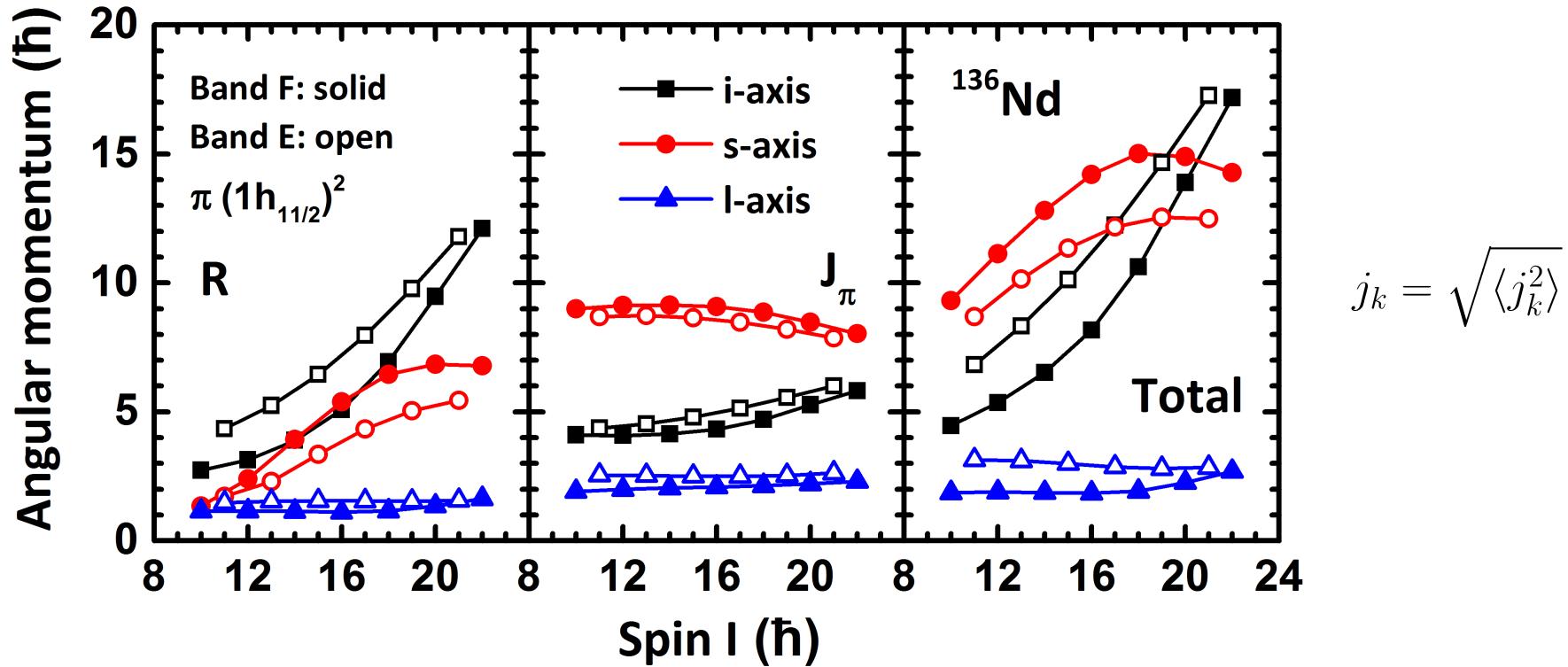


- Backbending is seen for band D.
- Only the low spin part ( $I=14, 16, 18$ ) might be the partner of band E.



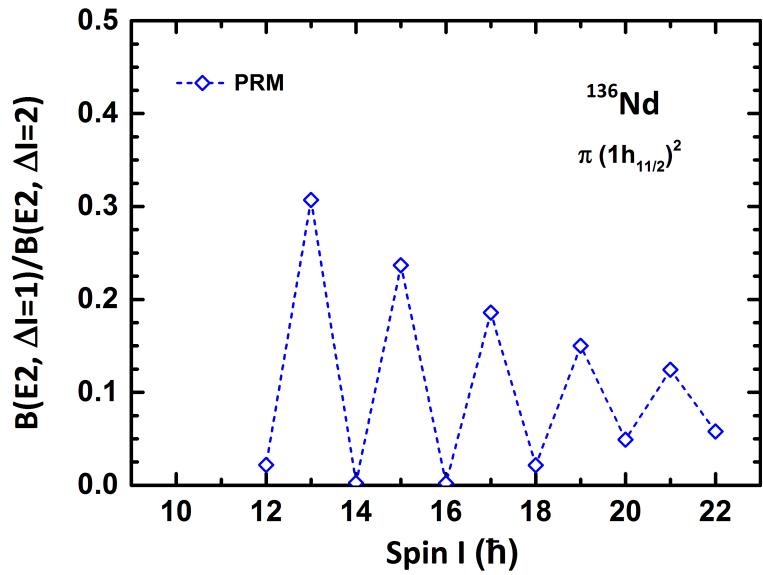
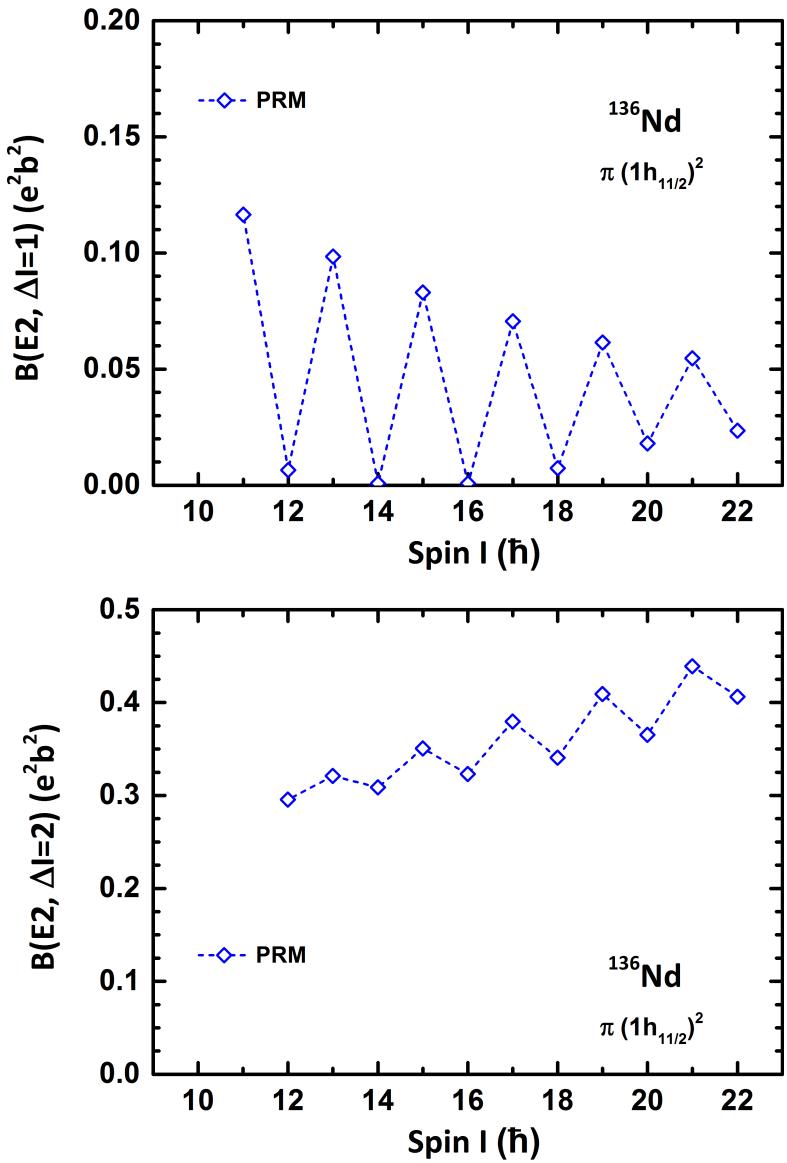
- Decreasing trend of energy difference are reproduced, supporting the transverse wobbling interpretation.
- Energy differences between the bands E and F are well reproduced, while energy differences between D and F are overestimated about 0.8 MeV.

# PRM: particle configuration



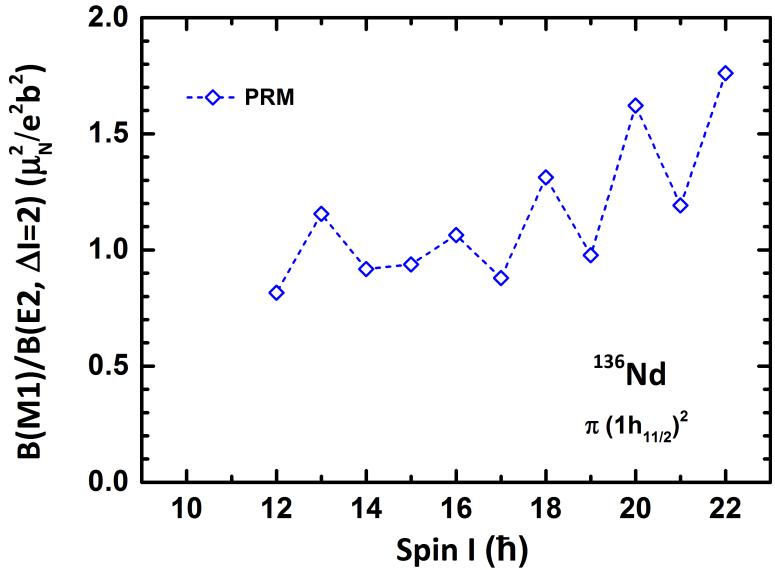
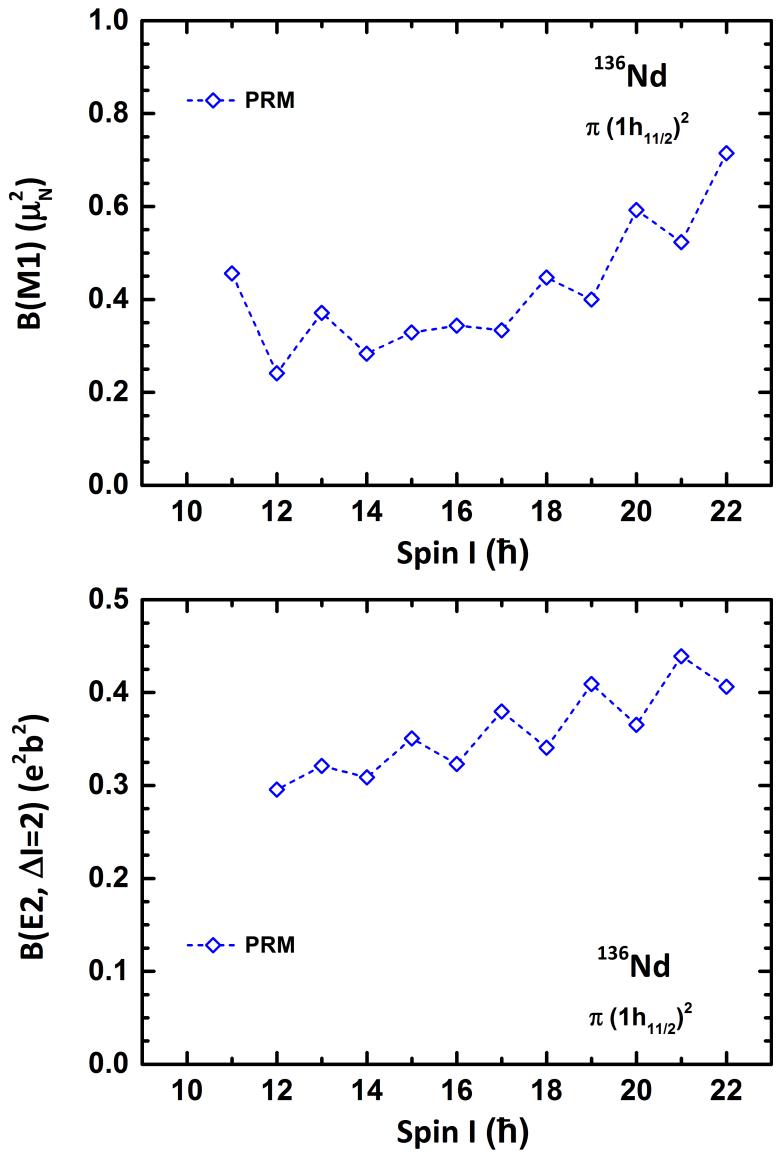
- The angular momentum of proton particles that mainly align along short axis are indeed large.
- Total spin have large component along the short axis at the low spin region.

# *PRM: particle configuration*



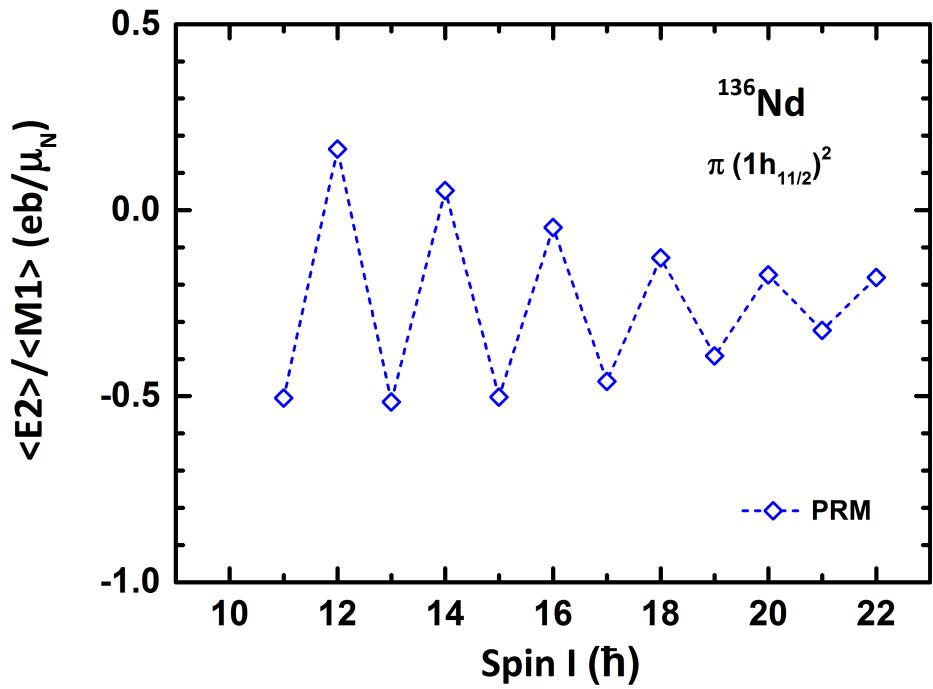
- The magnitude of BE2 ratio of band E is comparable with that in  $^{135}\text{Pr}$ ,
- $B(\text{E}2 \text{ } I \rightarrow I-1) \sim (30-15) \text{ WU}$  supporting the wobbling interpretation.

# *PRM: particle configuration*



- The magnitude of  $B(M1)/B(E2)$  ratio  $\sim 10$  times larger than that in  $^{135}\text{Pr}$ .

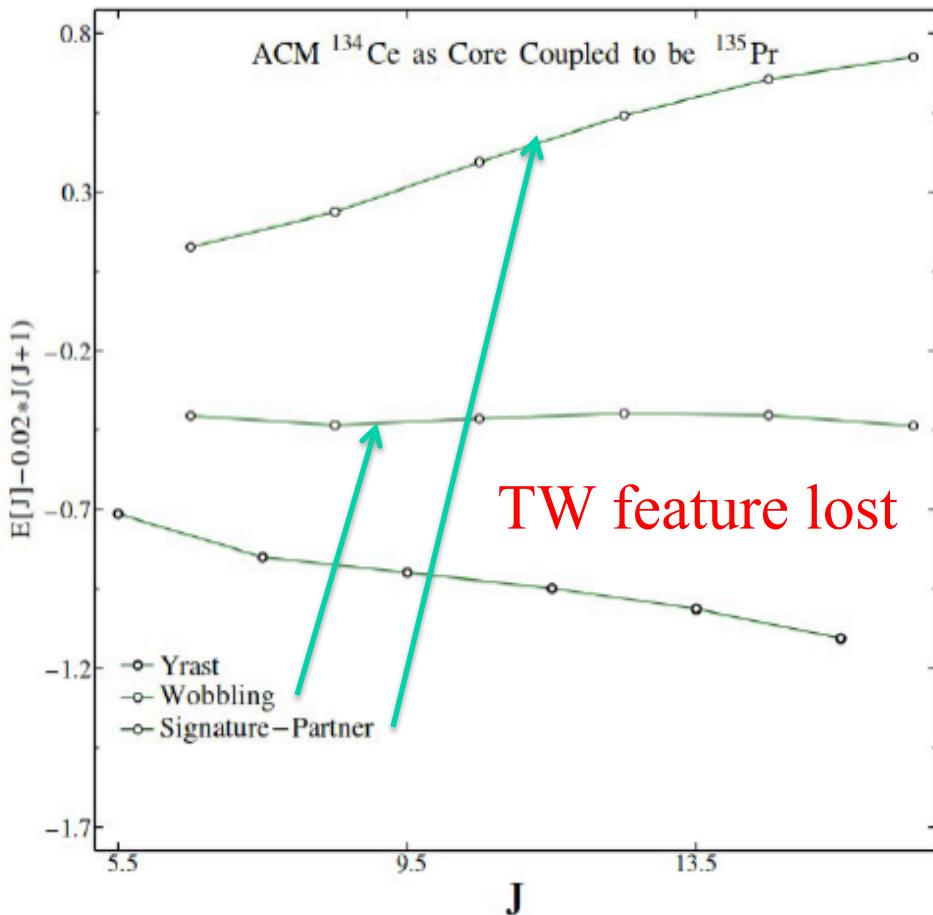
# *PRM: particle configuration*



$$B(E2/M1, I \rightarrow I - 1) = \langle E2/M1 \rangle^2$$

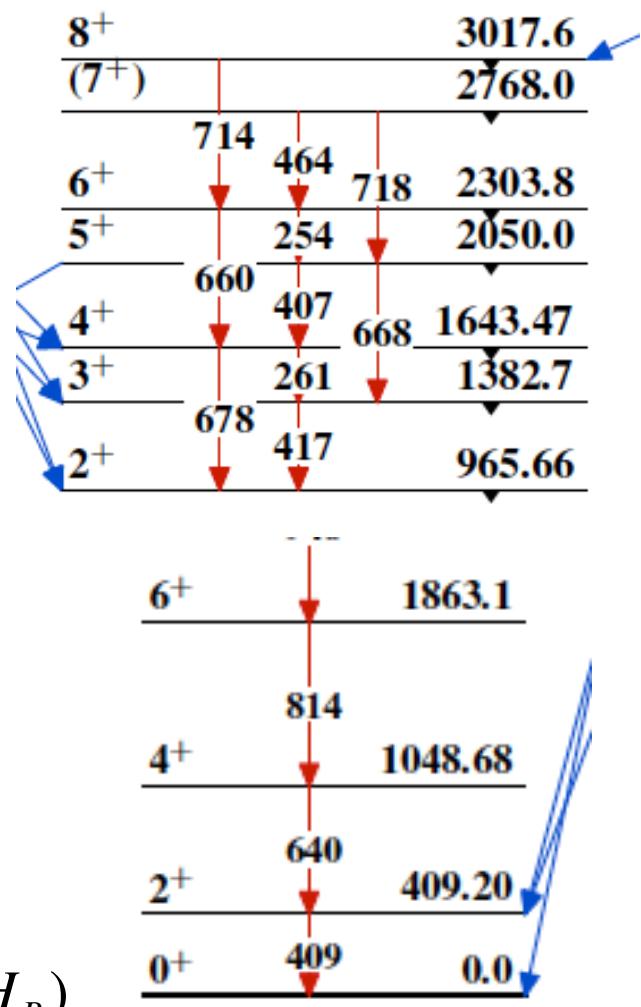
- The magnitude of this value is smaller than in  $^{135}\text{Pr}$  (from -1.25 to -0.50).

Weichuan Li

coupling to a " $\gamma$  soft core" (Bohr Hamiltonian  $H_B$ )

by quasiparticle coupling model

$$H = h_{sph} + \kappa[qQ]_0 + \Delta P - \lambda N + H_B$$

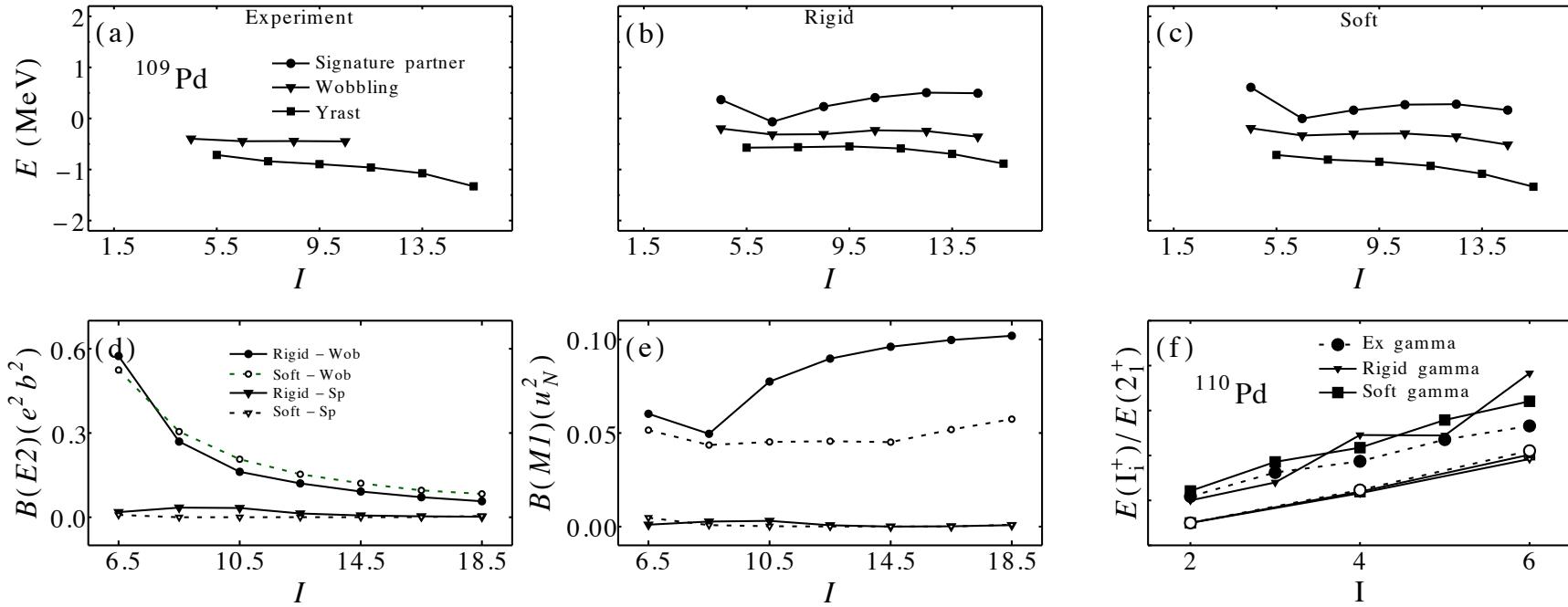
Band(D): Quasi  $\gamma$ -band

# Odd-A nuclei: $\gamma$ rigid or soft?

- Transverse wobbling appears only for stiff core
- Holds also for appearance of chiral pairs
- Does the presence of odd quasiparticles make the triaxial core stiff?
- The core is not soft as suggested by the collective Bohr Hamiltonian but stiff as assumed by the TSPM

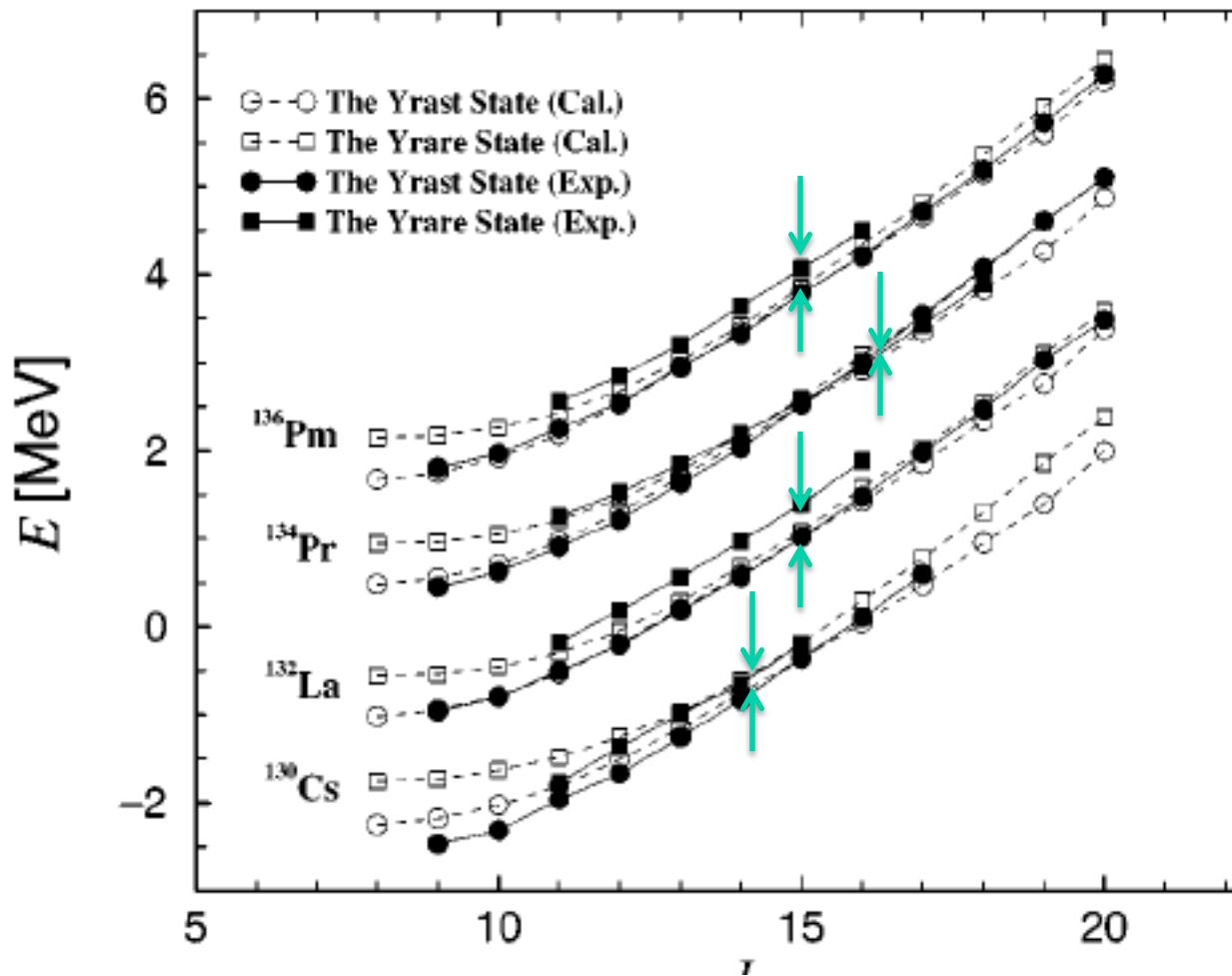
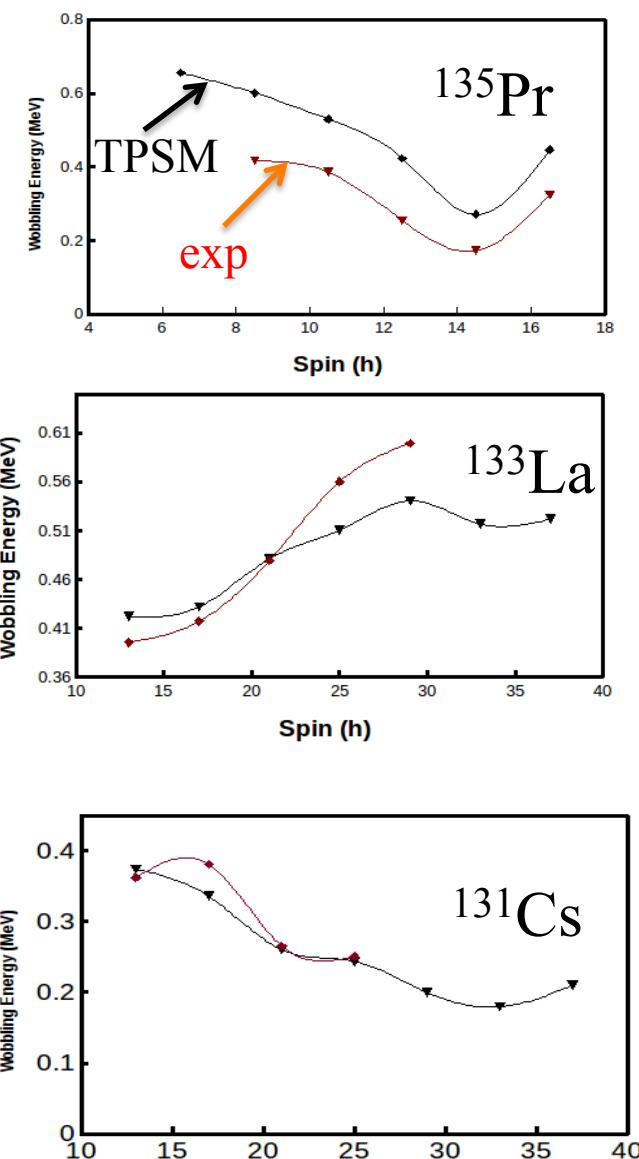
# Odd-A nuclei: $\gamma$ rigid or soft?

- Not the right question
- Maybe better question:  
How coherent is the  $\gamma$  excitation mode?
- Pure collective models are have a very limited range of application
- Strong coupling to the quasiparticle degrees of freedom is crucial
- QP+TR and TPSM are only the first steps.
- Relation to the adiabatic microscopic BH



# Problem with the QPTR model

Z-dependence of the TW and TCV vibrational frequencies not reproduced. TPSM OK. Pauli Principle core-qp?



# 4th order Rotor Hamiltonian

- Low energy dynamics is dominated orientation degrees of freedom
- The adiabatic approximation good when RPA energy is 'small'

The simplest Hamiltonian that can describe our tilted system is a forth order rotor Hamiltonian.

$$H_I = c_1 I_1^2 + c_2 I_2^2 + c_3 I_3^2 + c_4 I_1^4 + c_5 I_3^4 + \frac{c_6}{2} (I_1^2 I_3^2 + I_3^2 I_1^2)$$

where the forth order terms mimic the tilted mean field solution.

This Hamiltonian has the energy surface

$$E_J = c_1 J_1^2 + c_2 J_2^2 + c_3 J_3^2 + c_4 J_1^4 + c_5 J_3^4 + c_6 J_1^2 J_3^2 + c_0$$

$$\begin{aligned} E_J = & c_1 (J^2 - J_3^2) \cos^2 \varphi + c_2 (J^2 - J_3^2) \sin^2 \varphi + c_3 J_3^2 + \\ & c_4 (J^2 - J_3^2)^2 \cos^4 \varphi + c_5 J_3^4 + c_6 J_3^2 (J^2 - J_3^2) \cos^2 \varphi \end{aligned}$$

at a constant  $I$  which we used to fit the  $c_i$  parameters to the energy surface calculated with the full TAC Hamiltonian.

# Triaxial Projected Shell Model TPSM

W.A. Dar <sup>a</sup>, J.A. Sheikh <sup>a,b</sup>, G.H. Bhat <sup>a,\*</sup>, R. Palit <sup>c</sup>,  
R.N. Ali <sup>a</sup>, S. Frauendorf <sup>d</sup>

Nuclear Physics A 933 (2015) 123–134

o-o  $A \approx 100$

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G.H. Bhat <sup>a</sup>, J.A. Sheikh <sup>a,b</sup>, R. Palit <sup>c,\*</sup>

$^{128}\text{Cs}$

Physics Letters B 707 (2012) 250–254

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G.H. Bhat <sup>a</sup>, J.A. Sheikh <sup>a,b</sup>, R. Palit <sup>c,\*</sup>, SF o-e  $A \approx 135$  TW

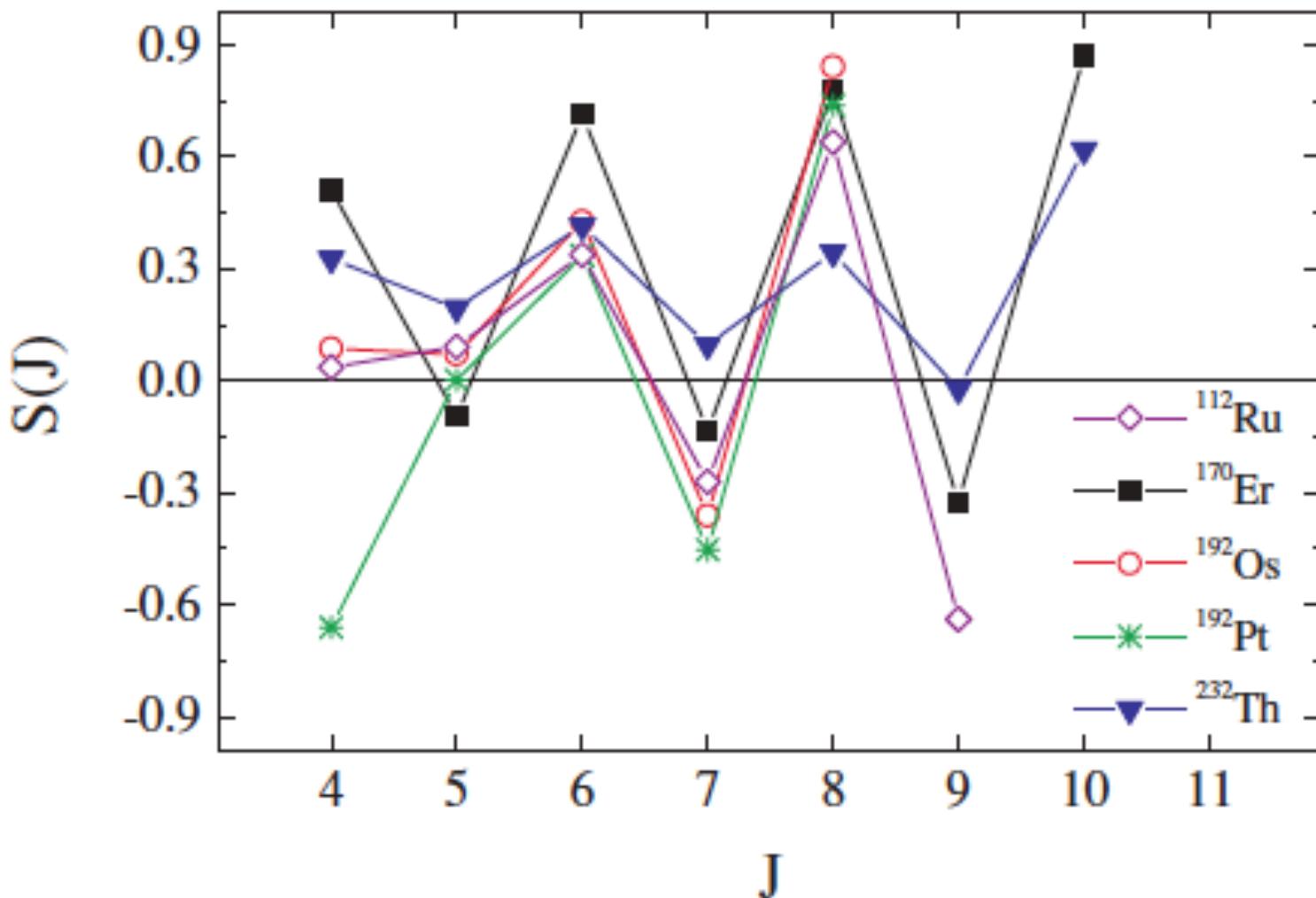
Microscopic (no phenomenological core)

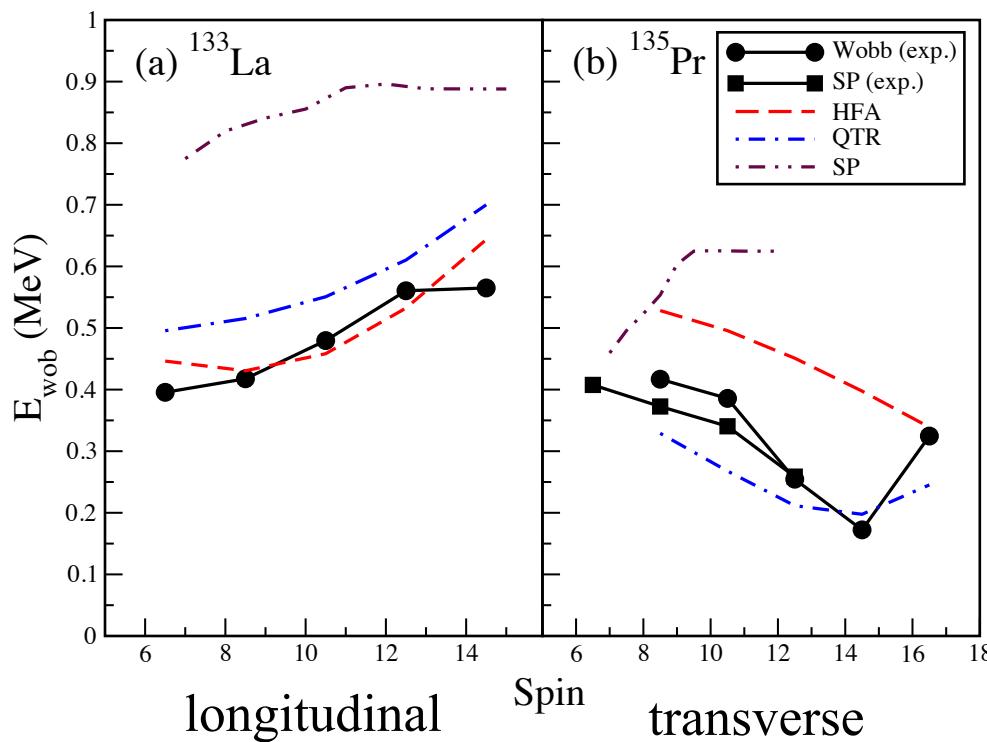
Non-adiabatic

Large amplitude

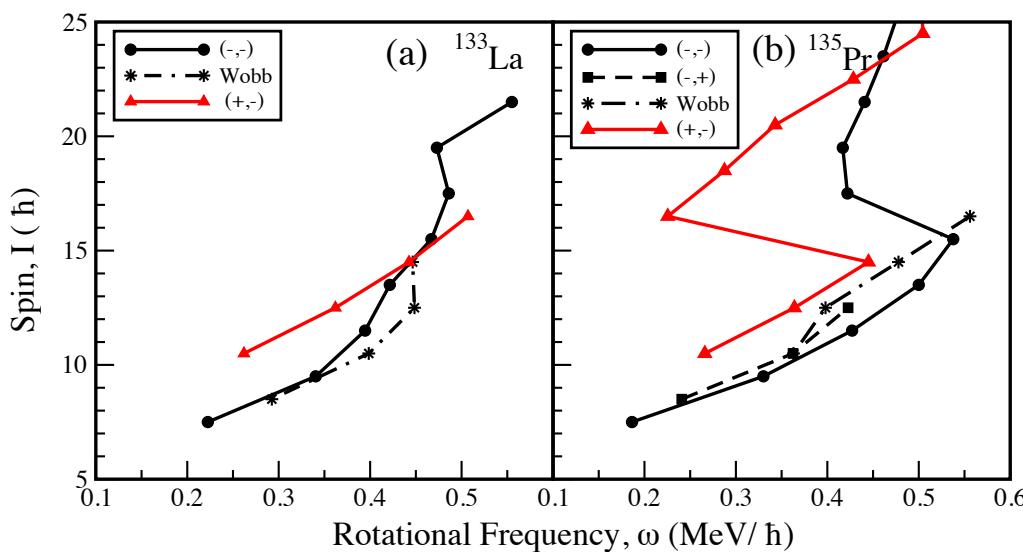
Fixed shape

## minima at odd J





What's going on??

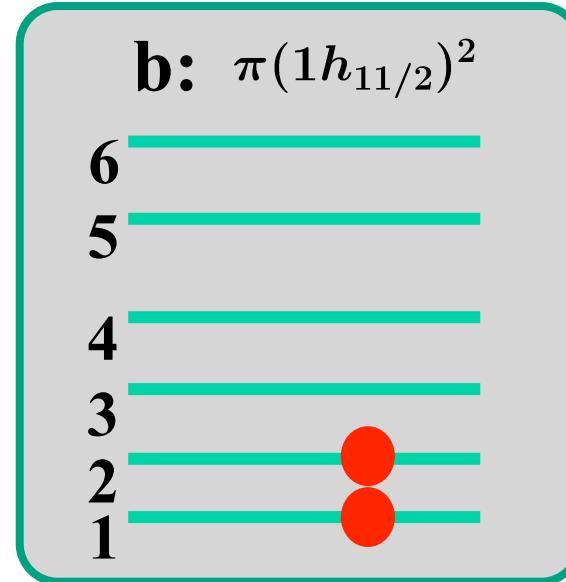
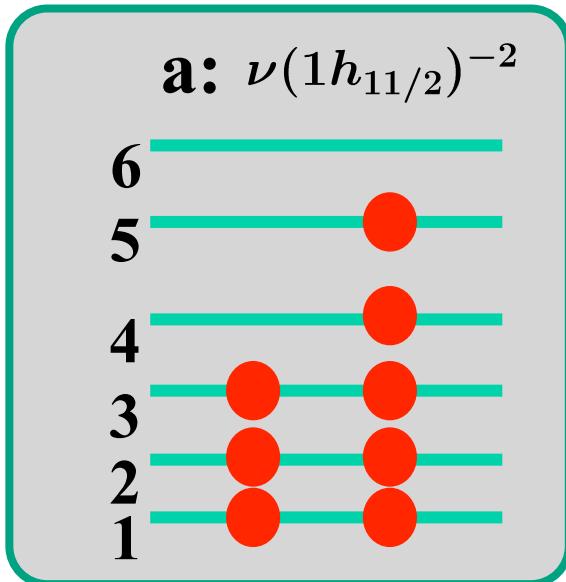


# CDFT calculations

- Numerical details:

- Interaction parameter: PC-PK1;
- Harmonic oscillator shells: Nf = 12;
- Paring correlations are neglected.

State	$E_x$	$(\beta, \gamma)$	Valence configuration	Unpaired configuration	Parity
A	0.00	(0.14, 60.0°)	$\pi(gd)^{10} \otimes \nu(1h_{11/2})^{-4}$		+
a	0.68	(0.18, 23.4°)	$\pi(gd)^{10} \otimes \nu(1h_{11/2})^{-4}$	$\nu(1h_{11/2})^{-2}$	+
b	1.03	(0.21, 18.1°)	$\pi(gd)^8(1h_{11/2})^2 \otimes \nu(1h_{11/2})^{-4}$	$\pi(1h_{11/2})^2$	+



Coupling between  $\gamma$  mode and rotation is very well  
Described by triaxial rotor ( $Q_2=7\text{eb}$ )

