

Giulia Zanderighi
Oxford University & STFC

Rencontres des Moriond EW 2009



LHC: unprecedented potential (energy, luminosity) first data very soon

Do we master QCD well enough to guarantee a successful program?

The challenge: provide theoretical predictions that match or exceed the accuracy of LHC data

The aim: early success in understanding the energy frontier

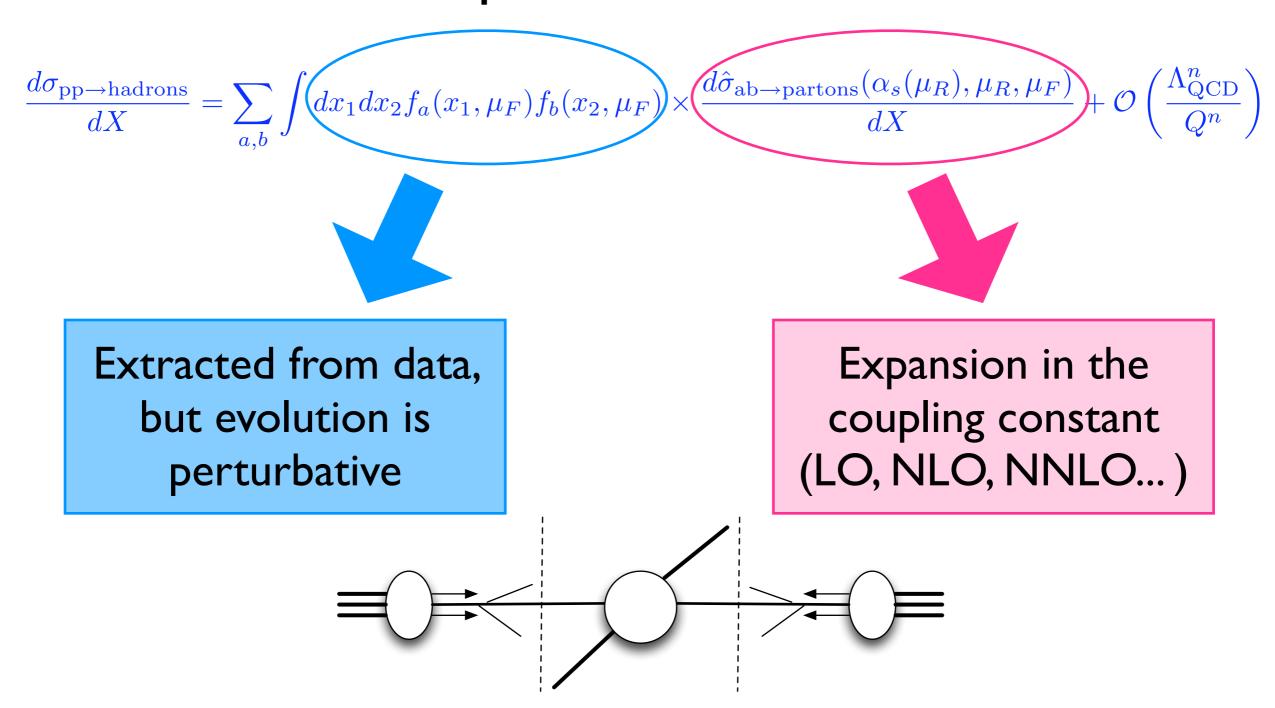
I will try to address this question by giving a

brief status report of technical progress

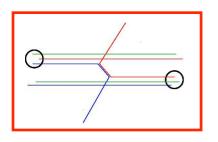
&

a short review of selected recent developments

Prerequisite: factorization



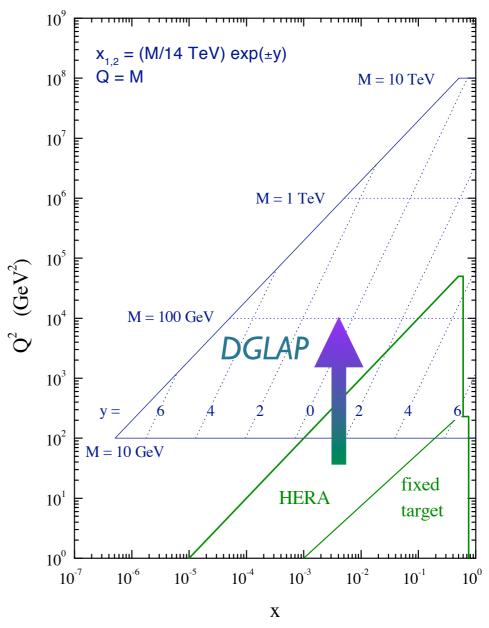
NB: factorization used in many contexts without proof



Parton densities coverage

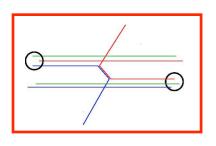
- most of the LHC x-range covered by Hera
- need 2-3 orders of magnitude Q²-evolution
- rapidity distributions
 probe extreme x-values
- I00 GeV physics at LHC: small-x, sea partons
- TeV physics: large x

LHC parton kinematics



PDF summary report, Hera-LHC '05

Hera: key and essential input to the LHC



Parton densities: recent progress

Recent major progress:

- full NNLO evolution (previous only approximate NNLO)
- full treatment of heavy flavors near the quark mass

[Numerically: e.g. (6-7)% effect on Drell-Yan at LHC]

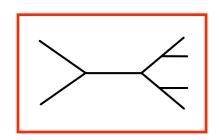
- more systematic use of uncertainties/correlations
- Neural Network (NN) PDFs

splitting functions at NNLO: Moch, Vermaseren, A. Vogt '04 [+ much related theory progress '04 -'08] Alekhin, CTEQ, MSTW (new MSTW '09), NN collaboration

Recently on the market: toolkits for NNLO DGLAP evolution of PDFs

PEGASUS A. Vogt '04; QCDNUM Botje '07 CANDIA Cafarella et al. '08; HOPPET Salam & Rojo '08

⇒ Description of PDFs reaching precision, but still some work ahead

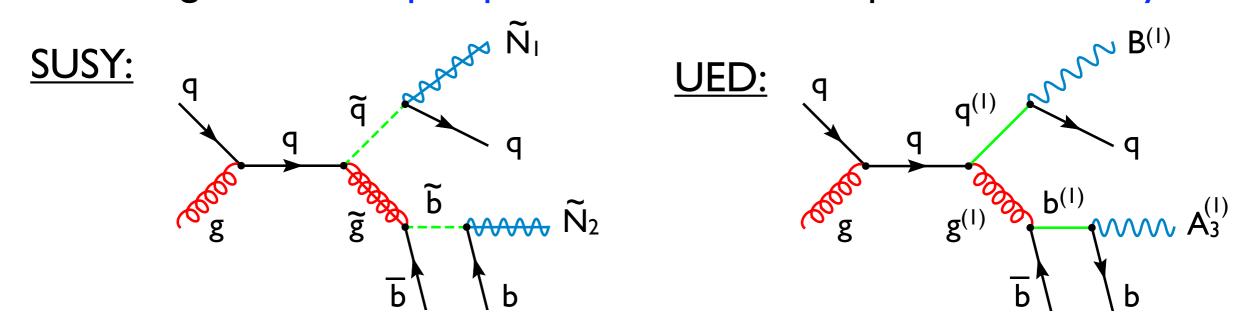


Multiparticle final states

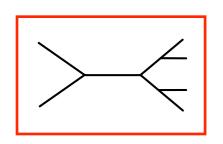
The LHC will operate in a new regime: highest energy & luminosity

Very large number of high-multiplicity events

- typical SM process is accompanied by radiation multi-jet events
- most signals involve pair-production and subsequent chain decays



More important than ever to describe high-multiplicity final states



Leading order

Status: fully automated, at most 8 outgoing particles

Alpgen, CompHEP, CalcHEP, Helac, Madgraph, Helas, Sherpa, Whizard, ...

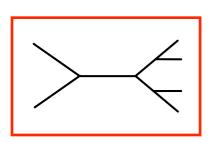
<u>Drawbacks of LO</u>: large scale dependences, sensitivity to cuts, poor modeling of jets, ...

Example: W+4 jet cross-section $\propto \alpha_s(Q)^4$

Vary $\alpha_s(Q)$ by $\pm 10\%$ via change of $Q \Rightarrow$ cross-section varies by $\pm 40\%$

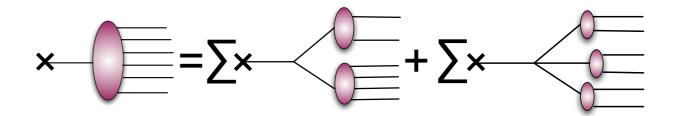
Why use LO at all?

- fastest option; often the only one
- test quickly new ideas with fully exclusive description
- many working, well-tested approaches
- highly automated, crucial to explore new ground, but no precision



LO: 3 methods beyond Feynman

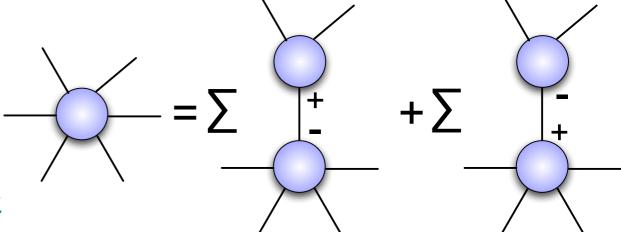
✓ Berends-Giele relations: compute helicity amplitudes recursively using off-shell currents



Berends, Giele '88

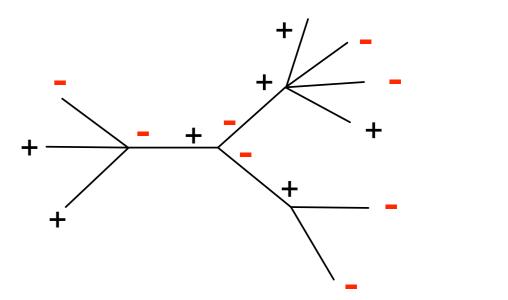
✓ BCF relations: compute helicity amplitudes via on-shell recursions (use complex momentum shifts)

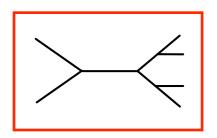
Britto, Cachazo, Feng '04



✓ CSW relations: compute helicity amplitudes by sewing together MHV amplitudes [- - + + ... +]

Cachazo, Svrcek, Witten '04





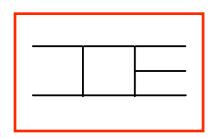
LO race: who is faster?

Time [s] for $2 \rightarrow n$ gluon amplitudes for 10^4 points

Duhr et al. '06 also Dinsdale et al. '06

Final state	BG	BCF	CSW	
2g	0.28	0.33	0.26	
3g	0.48	0.51	0.55	
4g	1.04	1.32	1.75	
5g	2.69	7.26	5.96	
6g	7.19	59. I	30.6	
7g	23.7	646	195	
8g	82. I	8690	1890	
9g	270	127000	29700	
I0g	864	-	-	

 undisputed numerical superiority of traditional Berends-Giele compared to twistor inspired methods



Next-to-leading order

For precision studies need next-to-leading-order

because the coupling is 'big-ish', to reduce dependence from unphysical scales, to model jets better, to predict the normalization, ...

Status of NLO:

[but: often no decays, newest codes mostly private]

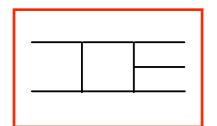
 \square 2 \rightarrow 4: barely touched ground [$q\bar{q} \rightarrow t\bar{t} b\bar{b}$, pp \rightarrow W+3jets]

Bredenstein et al. '08; Berger et al. '09; Ellis et al. '09

Cancelation of divergences: automated subtraction

Gleisberg, Krauss '07; TeVJet [public] Seymour, Tevlin '08; Hasegawa et al. '08

Bottleneck at NLO: virtual (loop) amplitudes

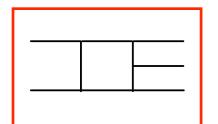


The 2005 Les Houches wishlist

[The QCD, EW & Higgs Working group report hep-ph/0604120]

Table 42: The LHC "priority" wishlist for which a NLO computation seems now feasible.

$ process \\ (V \in \{Z, W, \gamma\}) $	relevant for
1. $pp \rightarrow V V$ jet 2. $pp \rightarrow t\bar{t} b\bar{b}$ 3. $pp \rightarrow t\bar{t} + 2$ jets 4. $pp \rightarrow V V b\bar{b}$ 5. $pp \rightarrow V V + 2$ jets 6. $pp \rightarrow V + 3$ jets 7. $pp \rightarrow V V V$	$t \bar{t} H$, new physics $t \bar{t} H$ $t \bar{t} H$ $VBF \rightarrow H \rightarrow VV$, $t \bar{t} H$, new physics $VBF \rightarrow H \rightarrow VV$ various new physics signatures SUSY trilepton



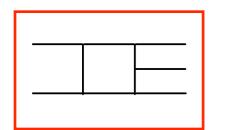
The 2007 update

Process	Comments
$(V \in \{Z, W, \gamma\})$	
Calculations completed since Les Houches 2005	
1. $pp \to VV$ jet	WW jet completed by Dittmaier/Kallweit/Uwer [3]; Campbell/Ellis/Zanderighi [4]
2. $pp o$ Higgs+2jets	and Binoth/Karg/Kauer/Sanguinetti (in progress) NLO QCD to the gg channel completed by Campbell/Ellis/Zanderighi [5];
3. $pp \rightarrow V V V$	NLO QCD+EW to the VBF channel completed by Ciccolini/Denner/Dittmaier [6,7] ZZZ completed by Lazopoulos/Melnikov/Petriello [8] and WWZ by Hankele/Zeppenfeld [9]
Calculations remaining from Les Houches 2005	
4. $pp \rightarrow t\bar{t}b\bar{b}$ 5. $pp \rightarrow t\bar{t}+2\mathrm{jets}$ 6. $pp \rightarrow VVb\bar{b}$, 7. $pp \rightarrow VV+2\mathrm{jets}$ 8. $pp \rightarrow V+3\mathrm{jets}$	relevant for $t\bar{t}H$ relevant for $VBF \to H \to VV$, $t\bar{t}H$ relevant for $VBF \to H \to VV$ vBF contributions calculated by (Bozzi/)Jäger/Oleari/Zeppenfeld [10–12] various new physics signatures
9. $pp o b ar{b} b ar{b}$	Higgs and new physics signatures
Calculations beyond NLO added in 2007	
10. $gg \to W^*W^* \mathcal{O}(\alpha^2 \alpha_s^3)$ 11. NNLO $pp \to t\bar{t}$ 12. NNLO to VBF and Z/γ +jet	backgrounds to Higgs normalization of a benchmark process Higgs couplings and SM benchmark
Calculations including electroweak effects	
13. NNLO QCD+NLO EW for W/Z	precision calculation of a SM benchmark

[The NLO multi-leg Working group report 0803.0494]

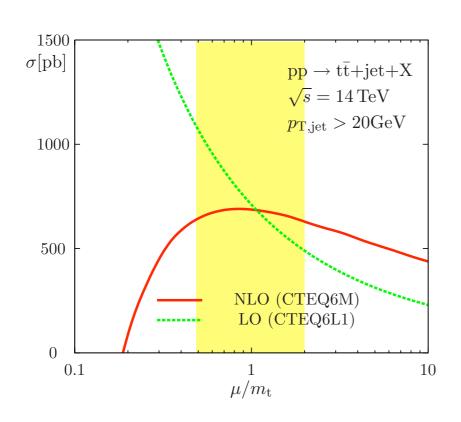
based on Feynman diagrams; private codes only

Table 1: The updated experimenter's wishlist for LHC processes

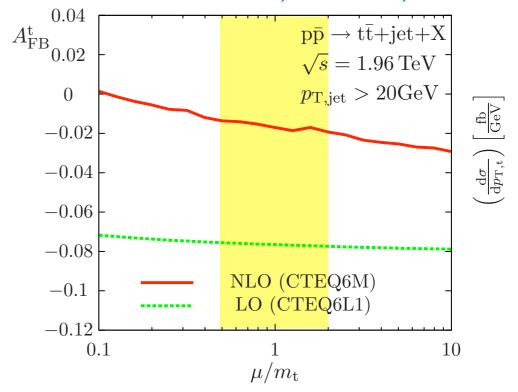


One NLO example: tt + l jet

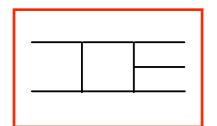
Calculation done with traditional methods



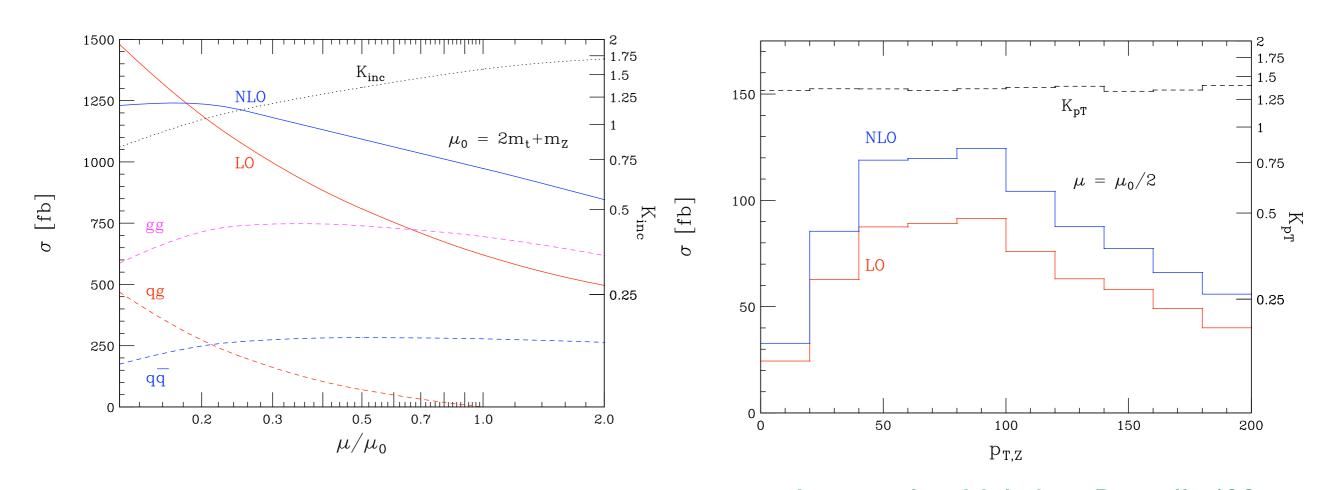
Dittmaier, Kallweit, Uwer '07, '08



- improved stability of NLO result [differential, but no decays]
- ▶ large effect on A_{FB} at the Tevatron: now compatible with zero
- essential ingredient of NNLO tt production

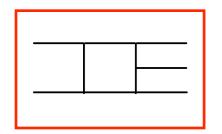


2nd NLO example: tt +Z



Lazopoulos, Melnikov, Petriello '08

- ▶ NLO increase cross section by 35% (residual 10% uncertainty)
- ▶ factor of 1.5-2 improvement on ttZ measurement (probe BSM)
- ▶ no significant change in distributions

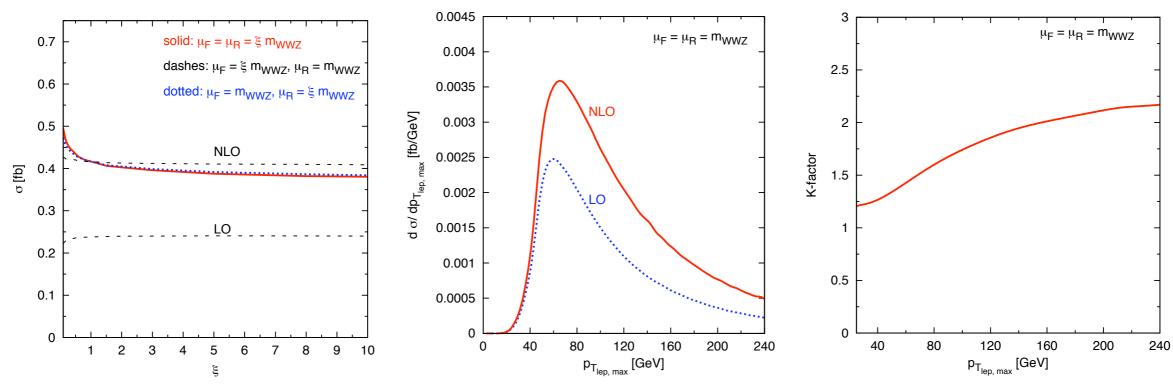


3rd NLO example:WW +Z

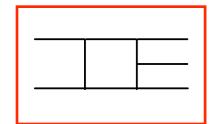
Motivation: probe EW gauge couplings and SUSY background

Lazoloupos, Melnikov, Petriello '07 Binoth, Ossola, Papadoloulos, Pittau '08

With decays and spin correlations: Hankele, Zeppenfeld '07

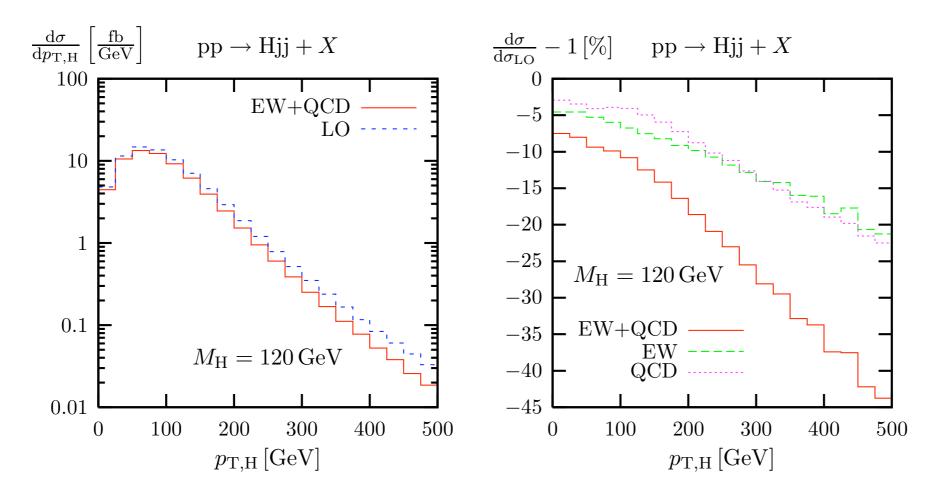


- NLO scale dependence larger than LO one (i.e. LO scale variation not indicative of the uncertainty)
- ▶ K-factors large and substantial shape change in distributions (simple multiplication of overall K-factor would distort shapes)

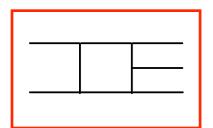


The "not so weak" EW: VBF Higgs

Ciccolini, Denner, Dittmaier '07



- ▶ EW and QCD of the same size!
- be aware of EW corrections for precision studies (peaks) and in tails of distributions (large electro-weak logarithms)



General NLO features?

	Typic	al scales	Teva	atron K -f	actor	LI	HC K-fac	etor
Process	μ_0	μ_1	$\mathcal{K}(\mu_0)$	$\mathcal{K}(\mu_1)$	$\mathcal{K}'(\mu_0)$	$\mathcal{K}(\mu_0)$	$\mathcal{K}(\mu_1)$	$\mathcal{K}'(\mu_0)$
177	200	2000	1 22	1 21	1 21	1 15	1.05	1 15
W	m_W	$2m_W$	1.33	1.31	1.21	1.15		1.15
W+1jet	m_W	$p_{T}^{ m jet}$	1.42	1.20	1.43	1.21	1.32	1.42
W+2jets	m_W	$p_T^{ m jet}$	1.16	0.91	1.29	0.89	0.88	1.10
WW+jet	m_W	$2m_W$	1.19	1.37	1.26	1.33	1.40	1.42
$t ar{t}$	m_t	$2m_t$	1.08	1.31	1.24	1.40	1.59	1.48
$t\bar{t}$ +1jet	m_t	$2m_t$	1.13	1.43	1.37	0.97	1.29	1.10
$b\bar{b}$	m_b	$2m_b$	1.20	1.21	2.10	0.98	0.84	2.51
Higgs	m_H	$p_T^{ m jet}$	2.33	_	2.33	1.72	_	2.32
Higgs via VBF	m_H	$p_T^{ m jet}$	1.07	0.97	1.07	1.23	1.34	1.09
Higgs+1jet	m_H	$egin{aligned} p_T^{ ext{jet}} \ p_T^{ ext{jet}} \ p_T^{ ext{jet}} \ p_T^{ ext{jet}} \ p_T^{ ext{jet}} \end{aligned}$	2.02	_	2.13	1.47	_	1.90
Higgs+2jets	m_H	$p_T^{ m jet}$	_	_	_	1.15	_	_

[NLO report 0803.0494]

General features:

- \blacktriangleright color annihilation, gluon dominated \Rightarrow large K factors
- \blacktriangleright extra legs in the final state \Rightarrow smaller K-factors

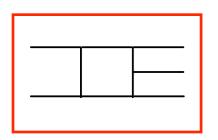
But be careful, only full calculations can really tell!

Recommendation

For the maximal exploitation of physics, there are also requirements on the experimental side. We suggest that cross sections at the LHC should be quoted at the hadron level, and where possible with the estimated parton-to-hadron corrections, so that any theoretical prediction (parton or hadron level) can easily be compared after the fact to the archived data [16]. Also, the experimental data needs to be quoted only for the range of measurement, rather than extrapolated to the full cross section; for example, measurements of $W \to e\nu$ should be quoted for the range of electron transverse momentum and rapidity and of missing transverse energy actually used in the triggering and analysis, rather than performing an extrapolation to the full W cross sections. Such recommendations were the exception (CDF W+jets) rather than the rule at the Tevatron and a clear model needs to be set for the LHC.

[Les Houches NLO report '08]

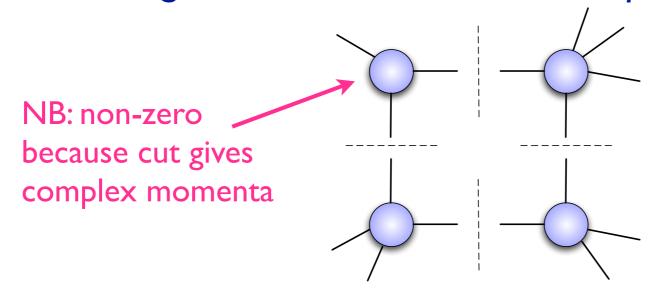
- this way your data will be useful for ever
- otherwise data will be obsolete sooner or later (typically rather soon)



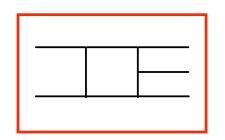
Two breakthrough ideas for NLO

Aim: NLO loop integral without doing the integration

I) "... we show how to use generalized unitarity to read off the (box) coefficients. The generalized cuts we use are quadrupole cuts ..."



Britto, Cachazo, Feng '04



Two breakthrough ideas

Aim: NLO loop integral without doing the integration

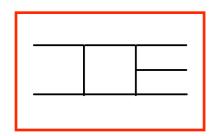
2) The OPP method: "We show how to extract the coefficients of 4-, 3-, 2- and 1-point one-loop scalar integrals...."

$$\mathcal{A}_{N} = \sum_{[i_{1}|i_{4}]} \left(d_{i_{1}i_{2}i_{3}i_{4}} \ I_{i_{1}i_{2}i_{3}i_{4}}^{(D)} \right) + \sum_{[i_{1}|i_{3}]} \left(c_{i_{1}i_{2}i_{3}} \ I_{i_{1}i_{2}i_{3}}^{(D)} \right) + \sum_{[i_{1}|i_{2}]} \left(b_{i_{1}i_{2}} \ I_{i_{1}i_{2}}^{(D)} \right) + \mathcal{R}$$

$$- \text{rational part}$$

$$\text{treated separately}$$

Ossola, Pittau, Papadopolous '06



A unified approach?

Partial fractioning via OPP + BG recursion for tree amplitudes + unitarity in higher dimension

⇒ full one-loop from tree level

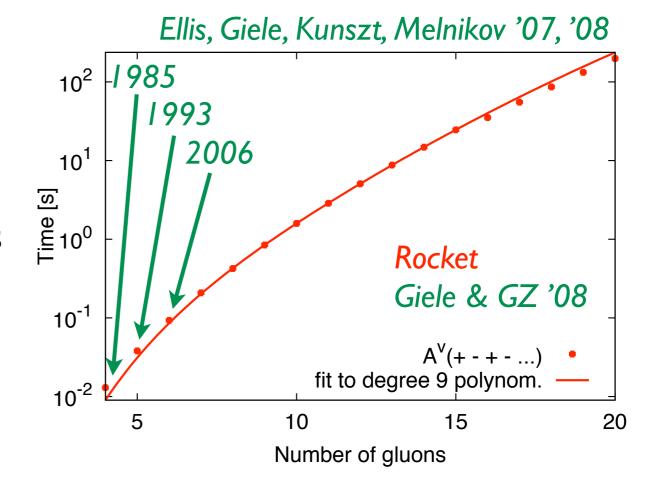
Two issues:

- Practicality? speed, stability
 Excellent performance of the method demonstrated for gluons
- Generality? what about realistic LHC processes?

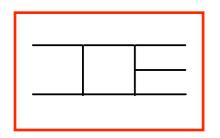
Cases studied:

$$0 \rightarrow \text{ttggg},$$

 $0 \rightarrow \text{qq} + (\text{ng}), 0 \rightarrow \text{qqW} + (\text{ng}),$
 $0 \rightarrow \text{qqQQW} + \text{Ig}$



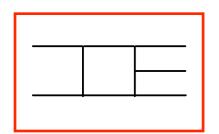
Also: Blackhat up to N=6 [7,8 MHV], Berger, Bern, Cordero, Dixon, Forde, Ita, Kosower, Maitre '08



1st physical application: W+3 jets

Why W+3 jets?

- I. W+3 jets measured at the Tevaton, but LO varies by more than a factor of 2 under reasonable changes in scales ⇒ cannot use LO predictions for serious comparison of theory and data
- II. measurements at the Tevaton for W+n jets with n=1,2 have shown that data is described well by NLO QCD
 - ⇒ interesting to verify this for 3 and more jets
- III. W+3 jets of interest at the LHC, as one of the backgrounds to model-independent new physics searches using jets + MET
- IV. calculation highly non-trivial (more than 1000 Feyman diagrams) optimal testing ground



Approximations

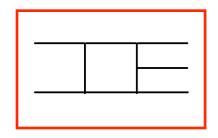
In spite of all technical advances the calculation of one-loop W+5 parton amplitudes and tree-level W+6 partons integrated over phase space is still a very challenging task

- ⇒ reasonable to look for reliable approximations as a first step
- the leading color approximation works up to 10% at leading order
- ▶ 4-quark processes only around 30% of the full result
- ▶ incoming gluon-gluon channel only 2%

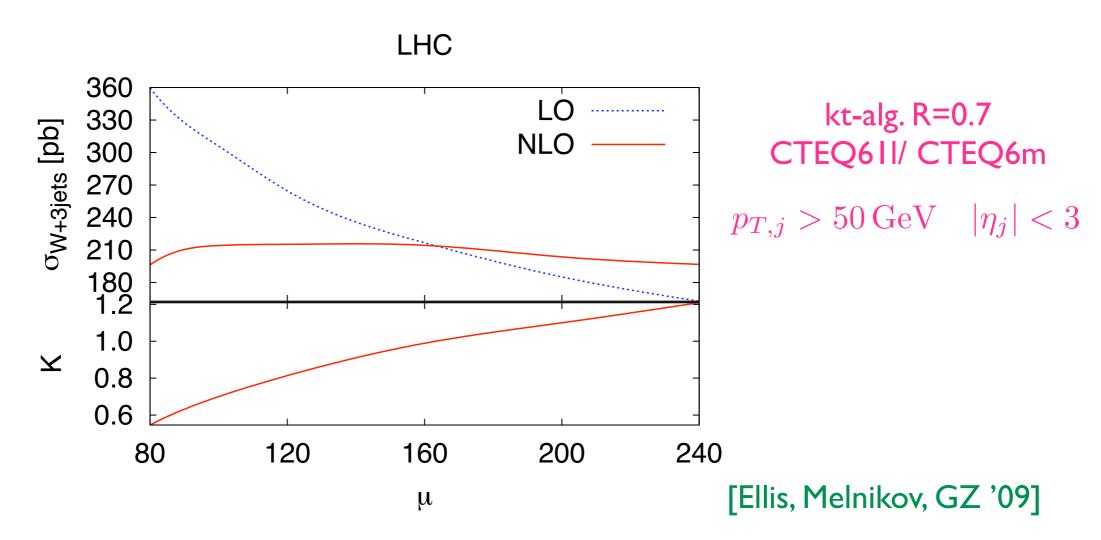
Since in the following

- leading color approximation is used
- 4-quark processes are dropped
- processes with two incoming gluons are omitted

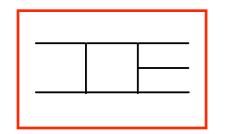
Absolute results for cross-sections presented next should be taken with care, however ratio of NLO/LO should be less sensitive to these omissions



W+3 jets at the LHC



- remarkable independence of cross-section on renormalization and factorization scale at NLO (unlike LO)
- LO = NLO at scales ~ 160 GeV
- gross features of W+3 jets are similar to W+2 jets, however the price one pays for an infelicitous choice of scales is higher now

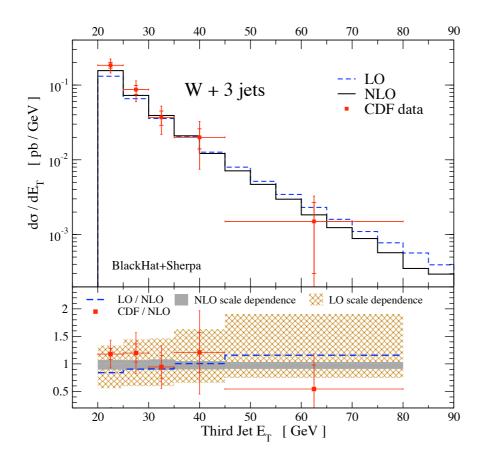


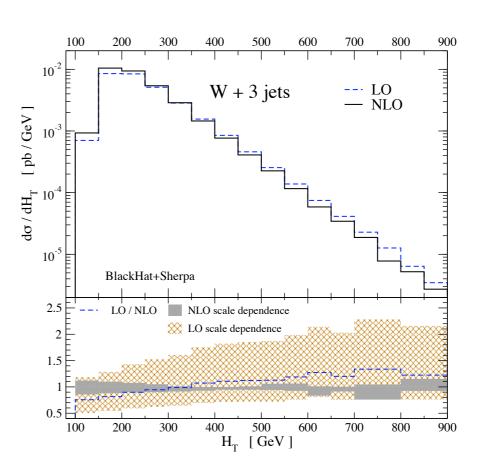
W+3 jets at the Tevatron

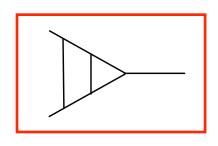
More recently, similar calculation for W+3 jets done in Blackhat+Sherpa

C.F. Berger, Z. Bern, L.J. Dixon, F. Febres Cordero, D. Forde, T. Gleisberg, H. Ita, D.A. Kosower, D. Maitre [0902.2760]

Still leading color approximation in virtual (not real), all subprocesses included (but no fermion loops)



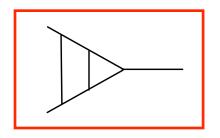




When is NLO not good enough?

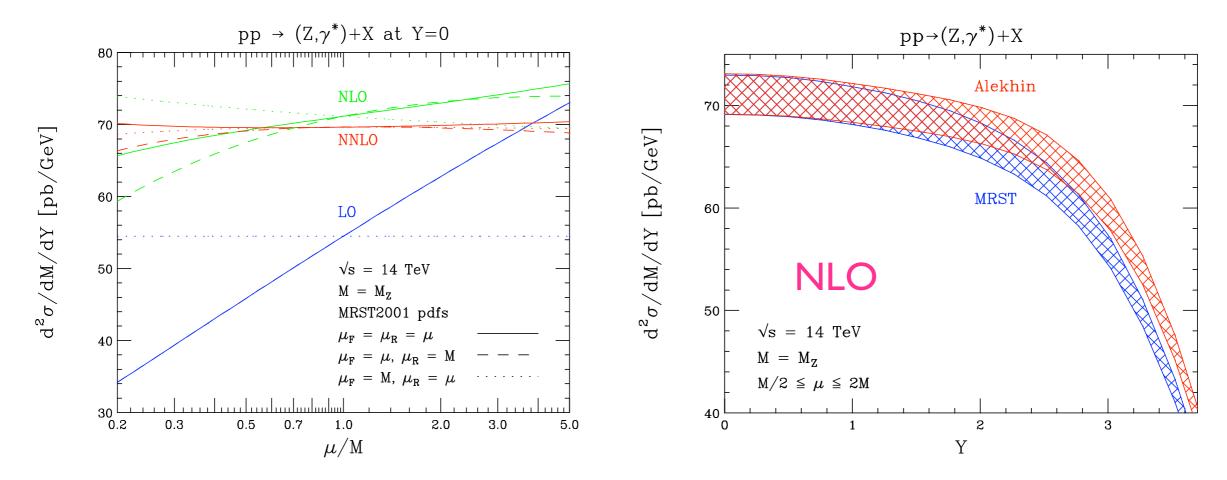
- when NLO corrections are large (NLO correction ≈ LO)
 This may happen when
 - process involves very different scales → large logarithms
 - new channels open up (at NLO they are effectively LO)
 - gluon dominated processes
- when high precision is useful (occasionally the case)
 - Drell-Yan, heavy-quark production, 3 jets in e⁺e⁻, ...
- when a reliable error estimate is wanted

Bottleneck: cancel divergences before numerical evaluation

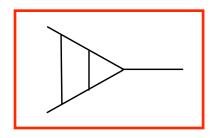


Best known at NNLO: Drell-Yan

- Most important and most precise test of the SM at the LHC
- Best known process at the LHC: spin-correlations, finite-width effects, γ-Z interference, fully differential in lepton momenta
- Sample NNLO results: scale stability & sensitivity to PDFs

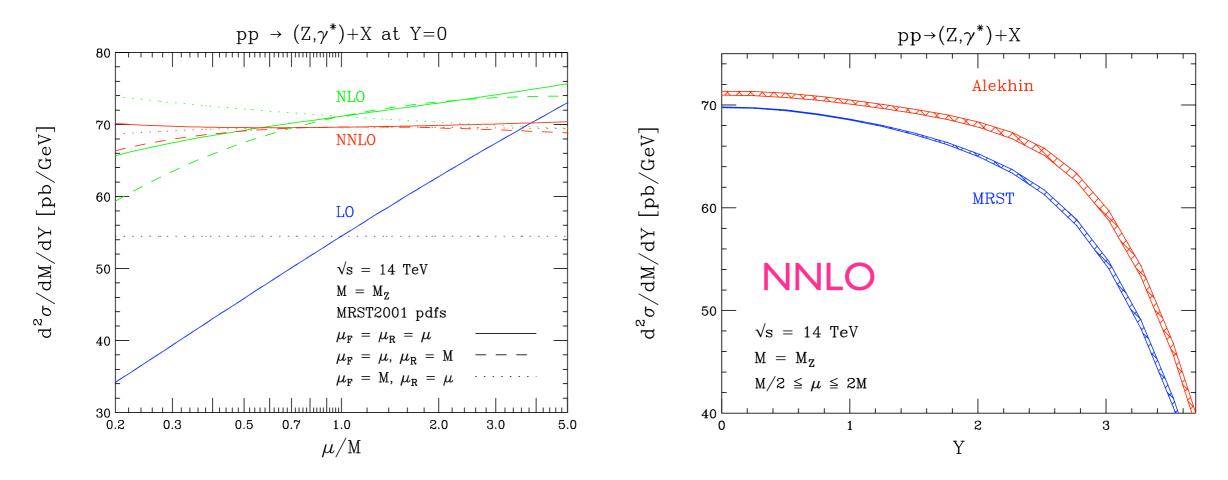


Anastasiou, Dixon, Melnikov, Petriello '03, '05; Melnikov, Petriello '06

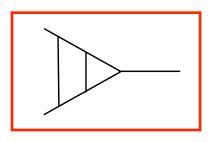


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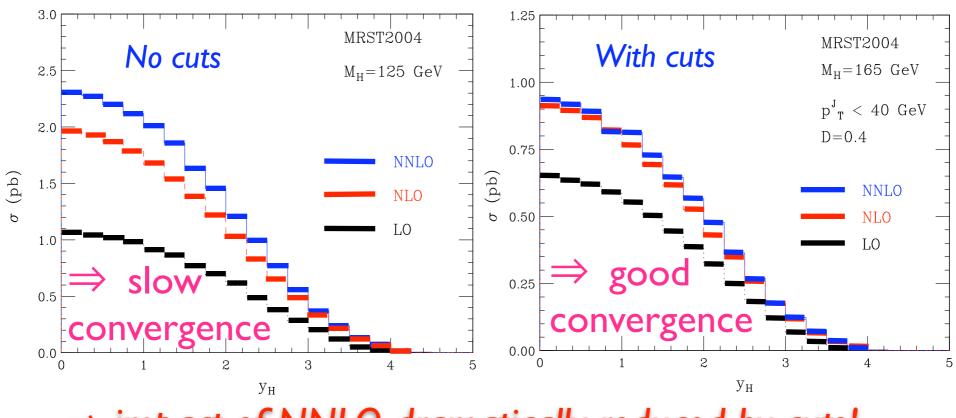


Anastasiou, Dixon, Melnikov, Petriello '03, '05; Melnikov, Petriello '06



NNLO Higgs with $H\rightarrow 21\ 2v$, $H\rightarrow 41$

FEHIP, Anastasiou, Dissertori, Stoeckli '07 also: HNNLO Catani, Grazzini '08



⇒ impact of NNLO dramatically reduced by cuts!

Very important to include cuts and decays in realistic studies!

Also uncertainty reduced via resummation of soft gluons at N³LO

Moch, Vermaseren '05



Jets: before 2006



Cones are IR unsafe!

The Cone is too rigid!

IR unsafety affects jet cross-sections by less than 1%, so don't need to care!

kt collects too much soft radiation!





Cones have a well-defined circular area!

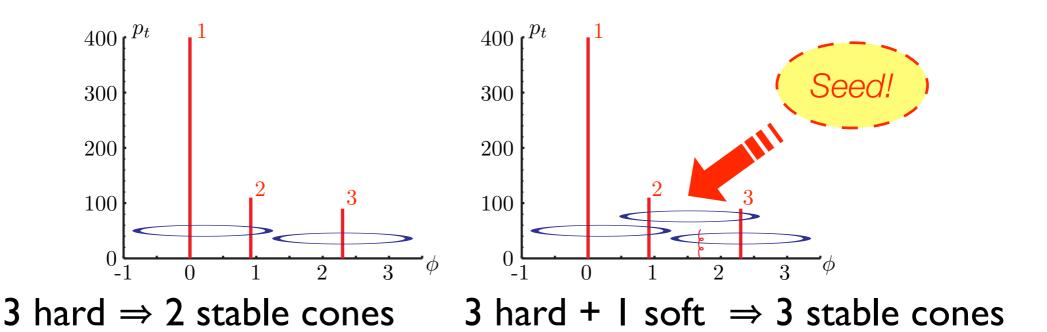
Jet area not well defined in kt: U.E. and pile-up subtraction too difficult!

What about dark towers??

After all, if D=1.35 R
Cone and kt are
practically the same
thing ...



Jets: infrared unsafety of midpoint



 \mid Soft emission changes the hard jets \Rightarrow algorithm is IR unsafe \mid

<u>Solution</u>: use a seedless algorithm find all stable cones $[\Rightarrow jets]$

SISCone: complexity N²log(N)

Salam, Soyez '07

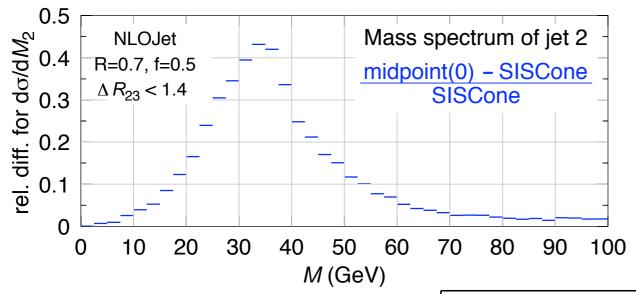
Similarly: iterative cone is collinear unsafe

Solution: anti-kt algorithm

Cacciari, Salam, Soyez '08



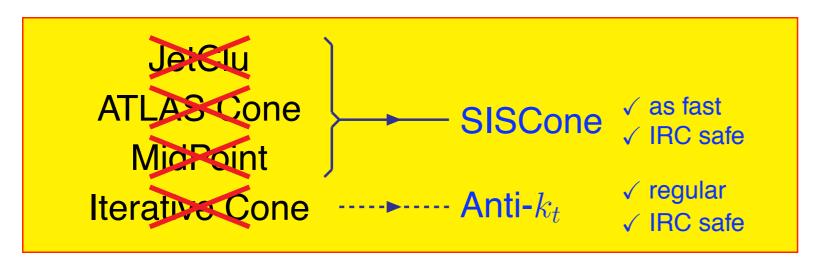
Physical impact of infrared unsafety



Up to 40% difference in mass spectrum

IR-unsafety is an issue at the LHC

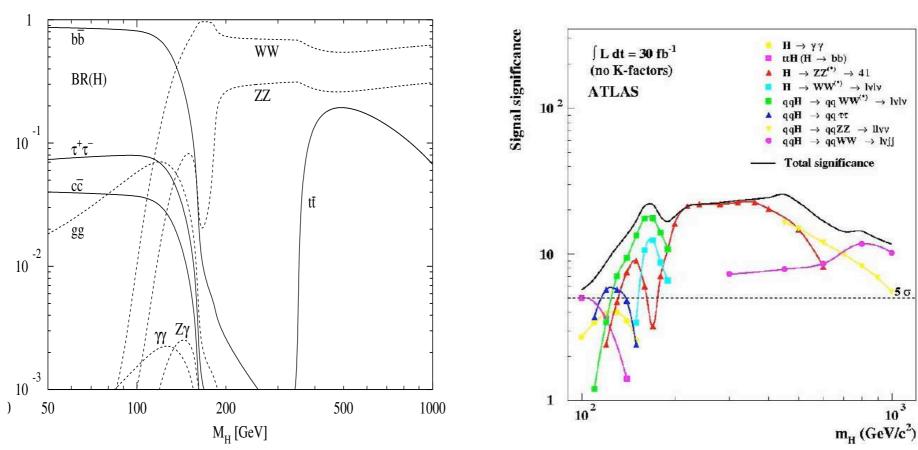
Observable	1st miss cones at	Last meaningful order
Inclusive jet cross section	NNLO	NLO
3 jet cross section	NLO	LO (NLO in NLOJet)
W/Z/H + 2 jet cross sect.	NLO	LO (NLO in MCFM)
jet masses in 3 jets	LO	none (LO in NLOJet)



If you don't want theoretical efforts to be wasted!



Z/W+H (\rightarrow bb) rescued ?



⇒ Light Higgs hard: H→bb dominant, but overwhelmed by background

Conclusion [ATLAS TDR]:

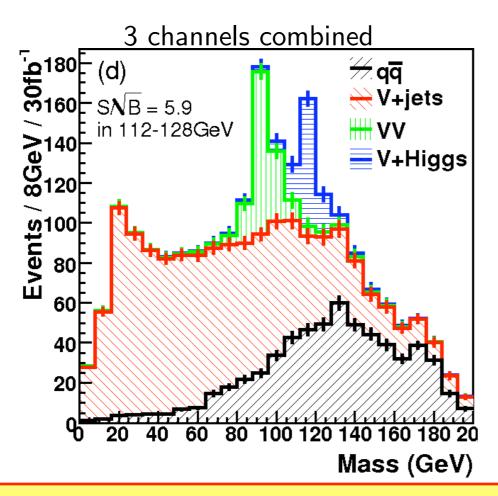
The extraction of a signal from $H \rightarrow bb$ decays in the WH channel will be very difficult at the LHC even under the most optimistic assumptions [...]



$Z/W + H \rightarrow bb$ rescued?

Boosted Higgs at high p_t : central decay products \Rightarrow single massive jet

Use jet-finding geared to identify the characteristic structure of fast-moving Higgs that decays into a bb-pair close in angle



- with common & channel specific cuts:
 ptV, ptH > 200GeV , ...
- ▶ real/fake b-tag rate: 0.7/0.01
- NB: very neat peak for WZ (Z →bb) Important for calibration

Butterworth, Davison, Rubin, Salam '08

5.9 σ at 30 fb⁻¹:VH with H \rightarrow bb recovered as one of the best discovery channels for light Higgs? More (exp) studies to come !

Conclusions

Impressive progress in the last years

- precision in parton densities
- higher orders (LO, NLO, NNLO & resummations)
- jets: many new ideas, impressive level of sophistication
- ... [much more, I did not have time to mention]

Progress driven by

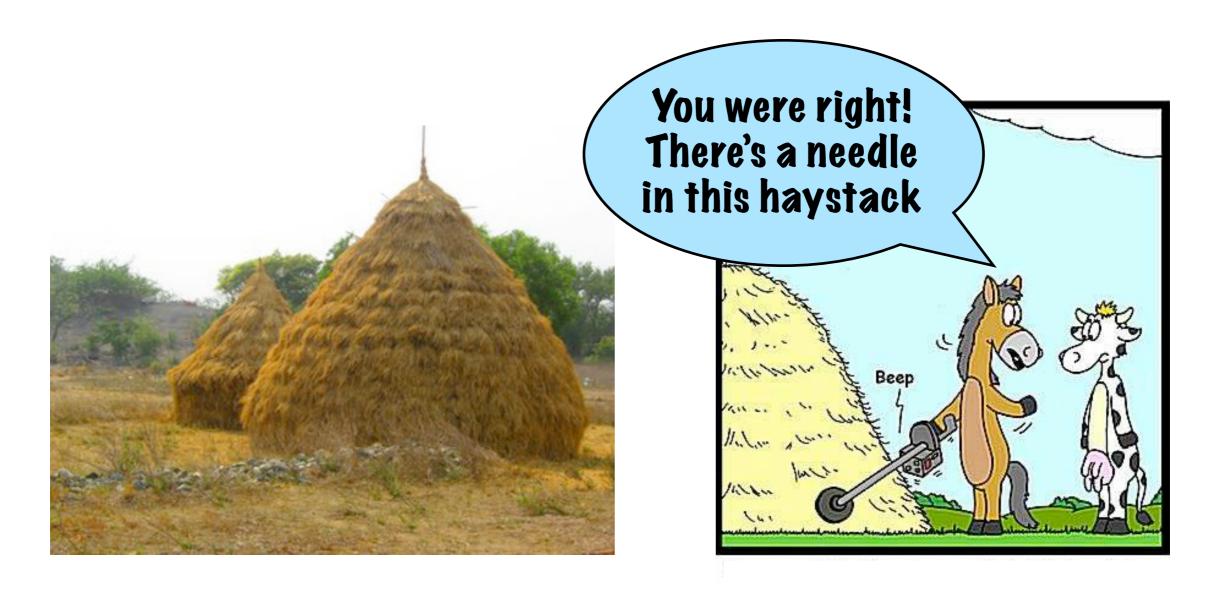
- automation/flexibility/public codes
- good communication with experimentalists & common papers

Still many challenges ahead but QCD theory will provide solid basis for a successful physics program at the LHC

How often did you hear the statement that looking for BSM signals at the LHC might be like looking for a needle in a haystack?

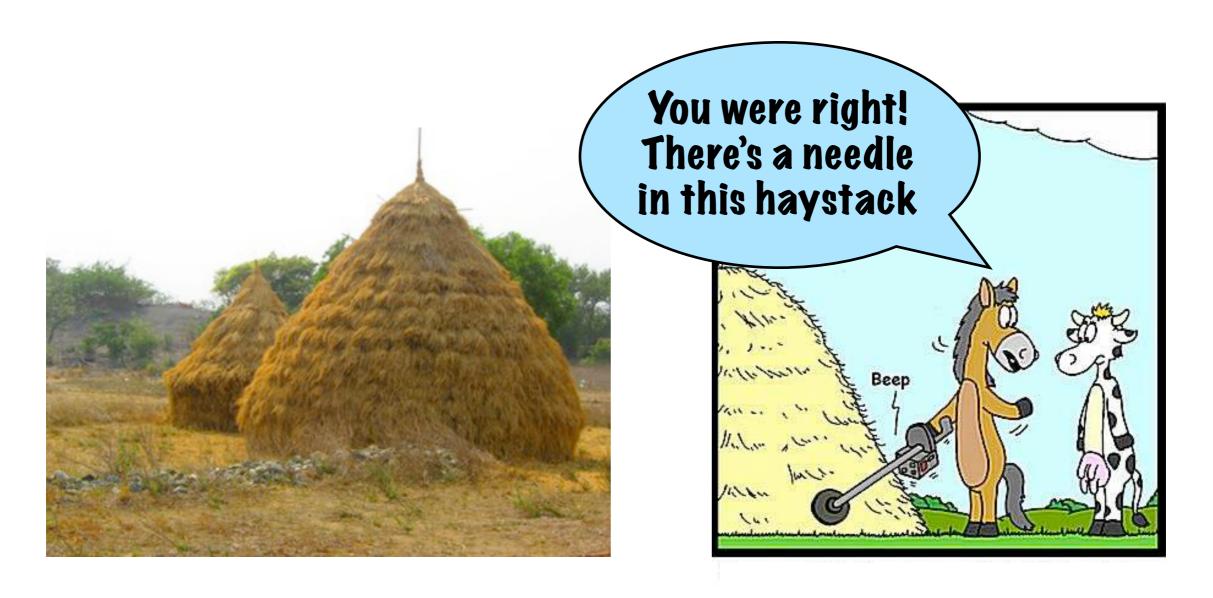


How often did you hear the statement that looking for BSM signals at the LHC might be like looking for a needle in a haystack?



But at the end, it is all a matter of having the right tools.

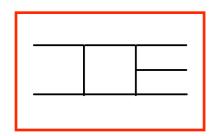
How often did you hear the statement that looking for BSM signals at the LHC might be like looking for a needle in a haystack?



But at the end, it is all a matter of having the right tools.

UNDERSTANDING QCD CRUCIAL TO DEVELOP THE RIGHT TOOLS!

Extra Slides



NLO + parton shower

Combine best features:

Get correct rates (NLO) and hadron-level description of events (PS) Difficult because need exact NLO subtraction and remove it from PS

Working LHC examples:

- ▶MC@NLO
 - W/Z boson production
 - WW, WZ, ZZ production
 - inclusive Higgs production
 - heavy quark production
 - single-top

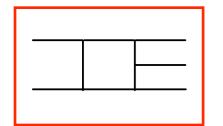
Frixione&Webber '02 and later refs.

- **▶** POWHEG
 - ZZ production
 - heavy quark production
 - W/Z production
 - Higgs, single top ... in progress

Nason '04 and later refs.

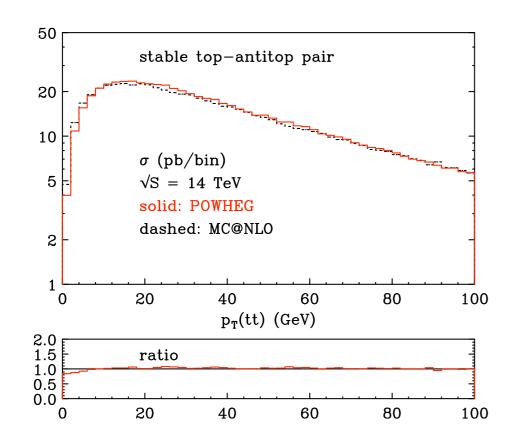
Other recent progress:

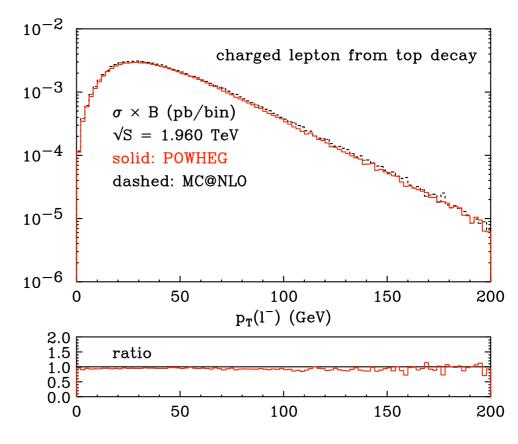
Shower with quantum inteference [Nagy, Soper], Geneve (SCET) [Bauer, Schwartz, Tackmann], Vincia (antenna factorization) [Giele et al.], Dipole factorization [Schumann]



MC@NLO vs PowHeg

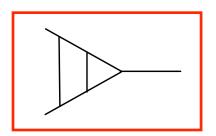
Top pair production:





Frixione, Nason, Ridolfi '07

⇒ excellent agreement for all observables considered (difference = different treatment of higher order terms)



tt cross-section at the LHC



tt cross-section at NLO

- ▶ scale uncertainty O(11%)
- ▶ PDF & m_t uncertainty O(2-3%)

Dawson, Ellis, Nason '88; Beenakker '89

experimental goal O(5%)

Need NNLO!

Similar aim: WW cross section at NNLO

Pinning down the tt cross-section

√ NLO + NLL resummed threshold corrections

Cacciari, Frixione, Mangano, Nason, Ridolfi '08

✓ two-loop virtual qq \rightarrow tt and gg \rightarrow tt at O(m_t^2/s)

Czakon, Mitov, Moch '07, '08

 \checkmark full mass dependence at two-loops for qq \rightarrow tt

Czakon '08

√ NNLO_{approx} (threshold logs + Coulomb + scale variation)

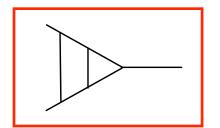
Moch, Uwer '08; also Kidonakis, R. Vogt '08

 \checkmark analytic two-loop fermionic corrections for qq \rightarrow tt

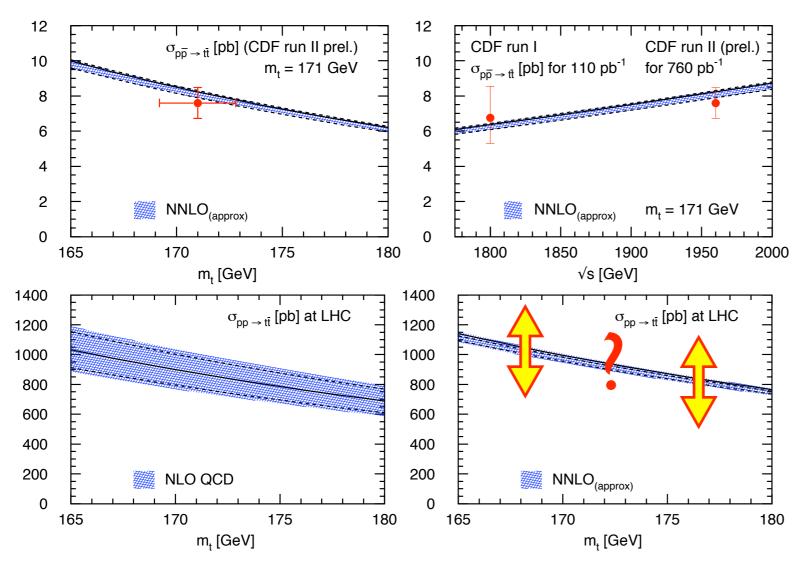
Bonciani, Ferroglia, Gehrmann, Maitre, Studerus '08

√ one-loop squared

Koerner, Merebashvili, Rogal '08



Towards NNLO tt



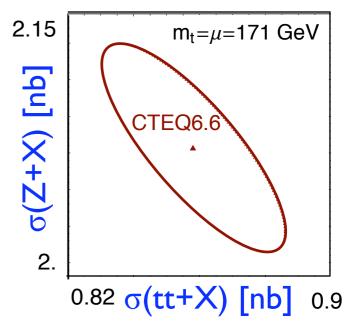
Moch, Uwer '08

NB: band = simultaneous variation of μ_F and μ_R

- mt from tt cross-section?
- tt promoted to luminosity monitor?

NB: PDFs in tt anti-correlated to W/Z

Nadolsky et al. '08

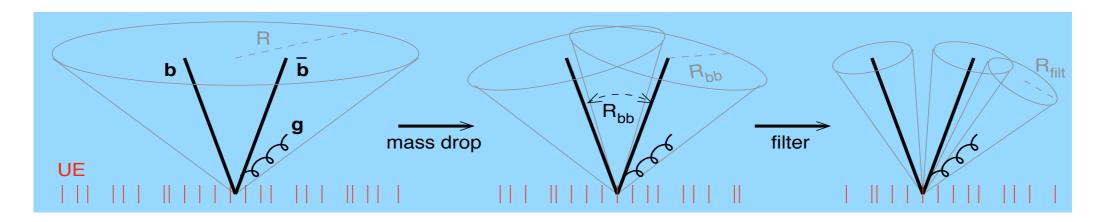




Z/W + H (→bb) rescued?

Boosted Higgs at high p_t : central decay products \Rightarrow single massive jet

Use jet-finding geared to identify the characteristic structure of fast-moving Higgs that decays into a bb-pair close in angle



- I. cluster the event with e.g. CA algo and large-ish R
- 2. undo last recomb: large mass drop + symmetric + b tags
- 3. filter away the UE: take only the 3 hardest sub-jets

<u>Related ideas for 2- and 3-body decays (boosted tops)</u>: Butterworth, Cox & Forshaw; Butterworth, Ellis & Raklev; Skiba & Tucker-Smith; Hodom; Baur; Agashe et al; Lille, Randall & Wang; Contino & Servant; Brooijmans; Thaler & Wang; Kaplan et al.; Almeida et al. [...]



Quality measures of jets

Suppose you are searching for a heavy state $(H \rightarrow gg, Z' \rightarrow qq, ...)$

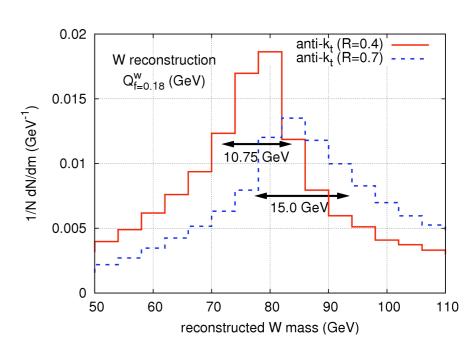
The object is reconstructed through its decay products

 \Rightarrow Which jet algorithm (JA) is best? Does the choice of R matter?

<u>Define</u>: $Q_f^w(JA, R) \equiv width of the smallest mass window that contains a fraction f of the generated massive objects$

- good algo \Leftrightarrow small $Q_f^w(JA, R)$
- ratios of $Q_f^w(JA,R)$: mapped to ratios of effective luminosity (with same S/\sqrt{B})

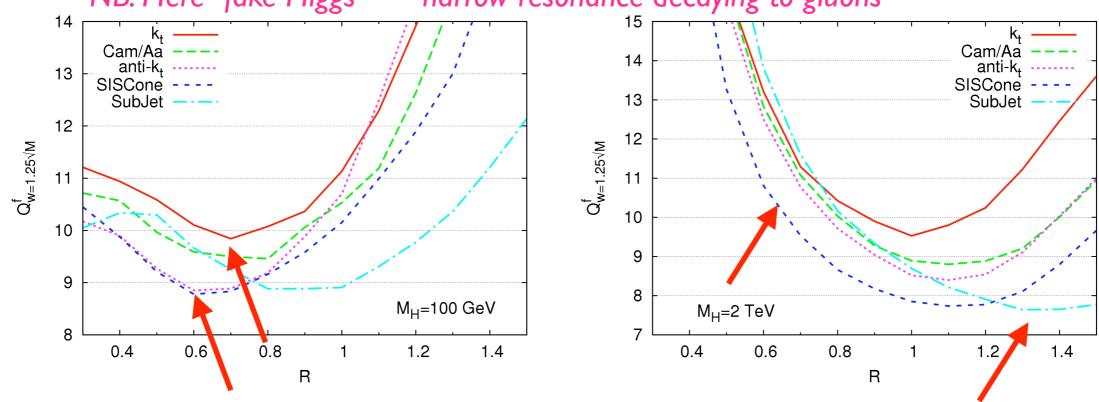
$$\mathcal{L}_2 = \rho_{\mathcal{L}} \mathcal{L}_1 \qquad \qquad \rho_{\mathcal{L}} = \frac{Q_z^f(JA_2, R_2)}{Q_z^f(JA_1, R_1)}$$





Quality measures: sample results

NB: Here "fake Higgs" = narrow resonance decaying to gluons



- ▶ At 100GeV: use a Tevatron standard algo (k_t, R=0.7) instead of best choice (SISCone, R=0.6 \Rightarrow lose $\rho_{\mathcal{L}} = 0.8$ in effective luminosity
- At 2 TeV: use $M_{Z'}$ =100GeV best choice (or k_t) instead CAfilt, R=0.9 \Rightarrow lose $\rho_{\mathcal{L}} = 0.6$ in effective luminosity

A good choice of jet-algorithm does matter!

Bad choice of algo ⇔ lost in discrimination power!