

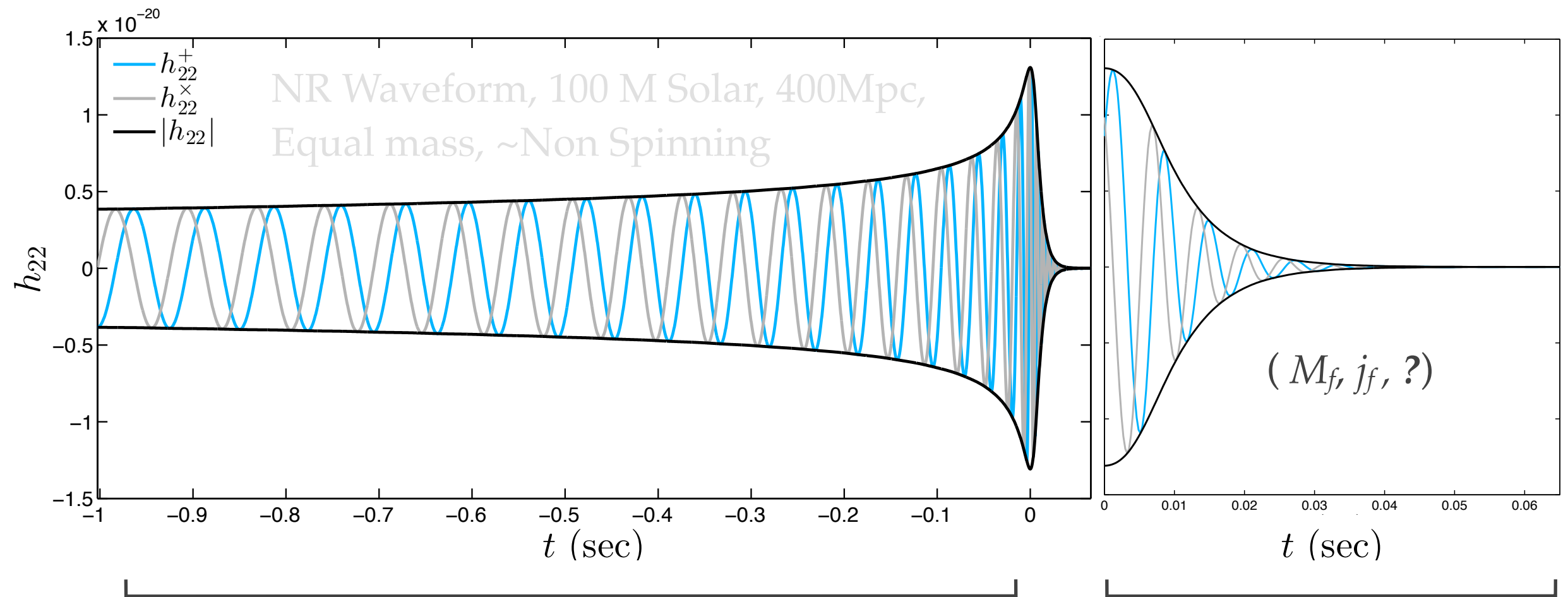
Context: Quasinormal models resulting from the merger of stellar mass BHs, and learning as much as we can from post-merger (ringdown) signals ...

The Quasi-normal Modes of Black Holes

Review and Recent Updates



Context QNMs from Binary Black Hole Mergers



Inspiral, Merger

Ringdown (Quasinormal Modes)

The **remnant black hole** (BH) is highly perturbed, with gravitational wave radiation that rings down. One goal of gravitational wave astronomy is to use our theoretical knowledge of BH QNMs to **learn more from detections**, and **determine the consistency (or inconsistency) of signals with GR**.

Overview Black Hole Quasinormal Modes, Review and Updates

A. The current context of black hole Quasinormal Modes (QNMs)

- QNMs as tools to help us learn about astrophysical BHs and Test GR

B. Review: QNMs and the fundamental questions about black holes

- Black hole stability (Perturbation Theory)
- Quasinormal mode structure (Kerr)

C. Review: Analytic and numerical relativity, key aspects

- Numerical Relativity *vs* Perturbation Theory
- QNM use in Binary BH (BBH) signal models

D. The current and new questions about BH QNMs

- What can we learn with QNMs?
- When can we learn it?

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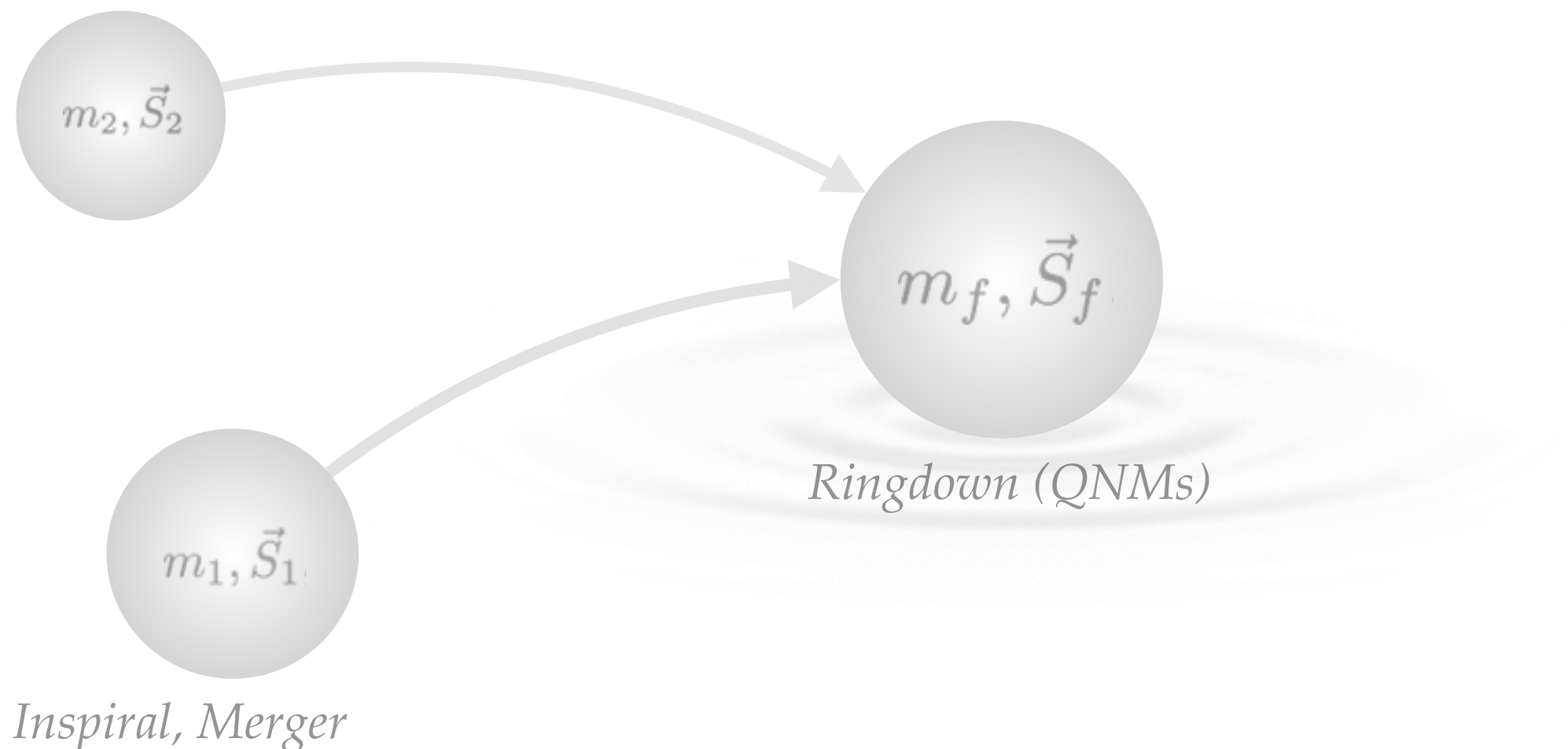
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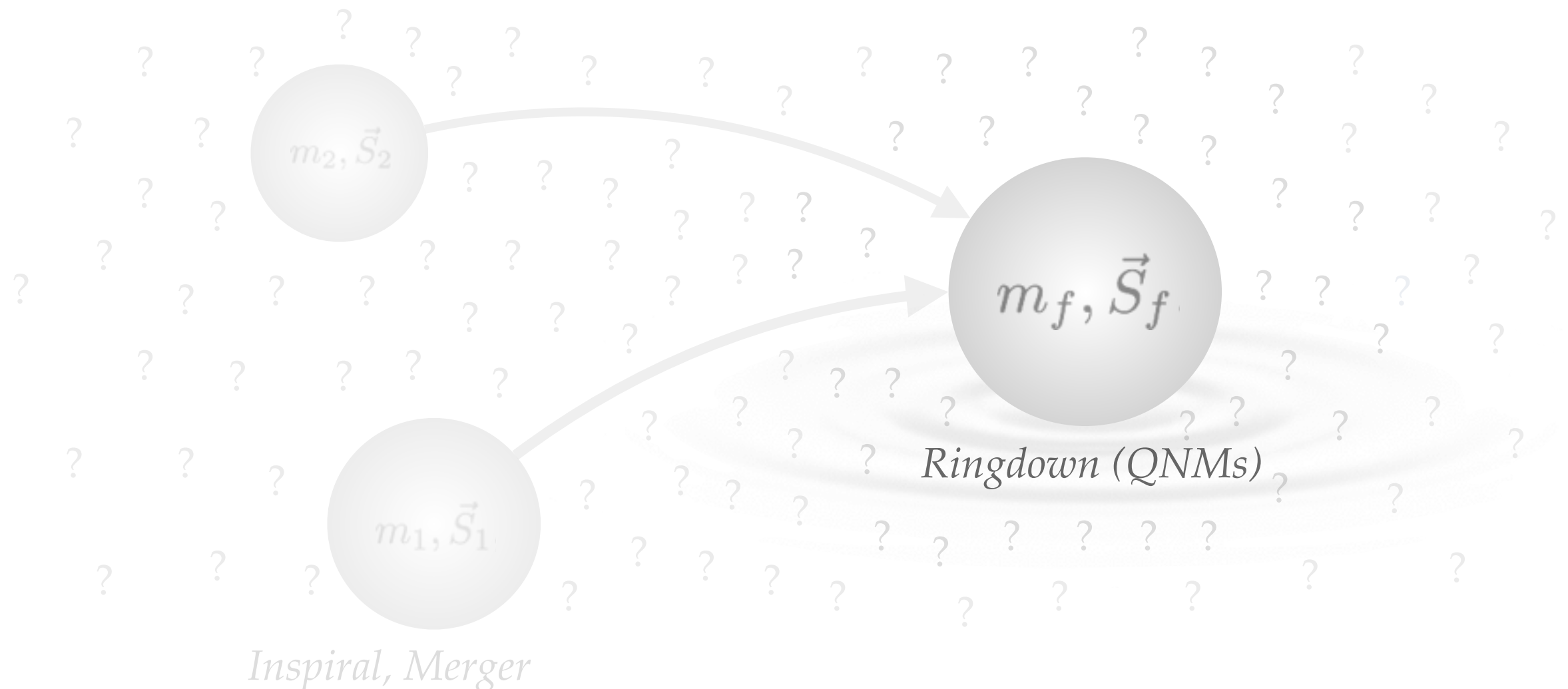
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BH QNMs Learning about Astrophysical BHs and Testing GR



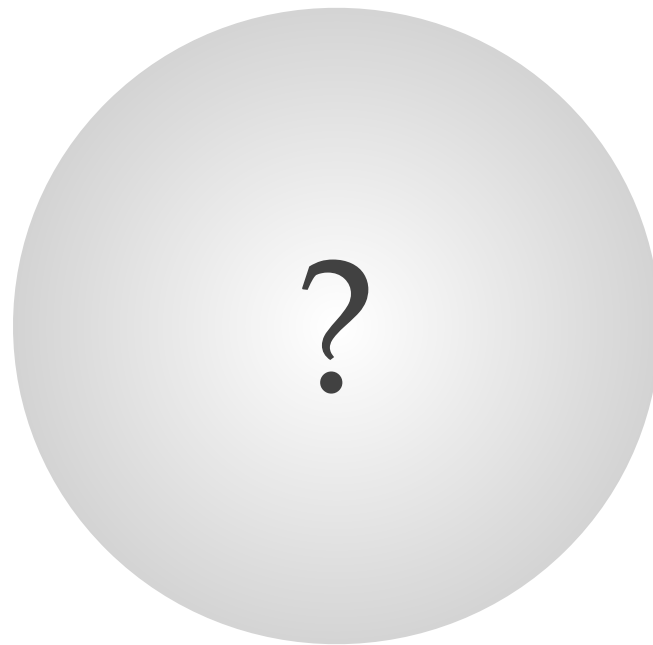
Current gravitational wave detections have been of black holes as described by GR — **but there are many notable limitations:** signal SNR, template accuracy, number of detections, etc

BH QNMs Learning about Astrophysical BHs and Testing GR



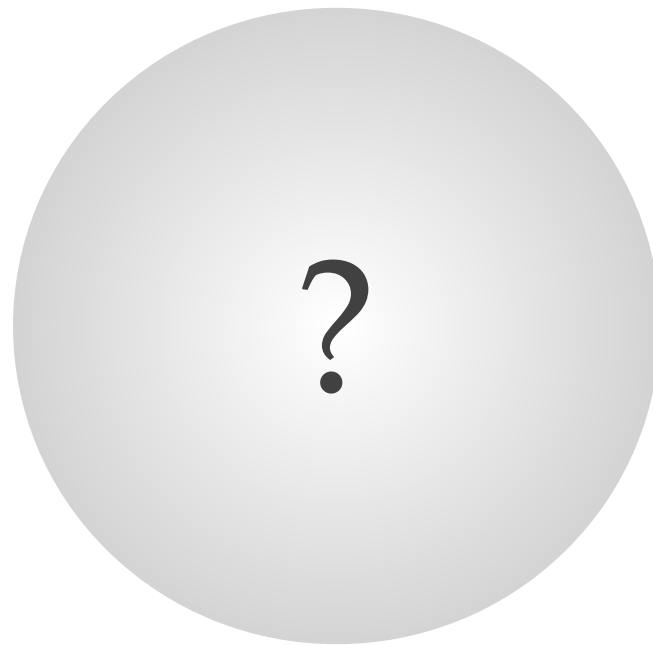
Many questions persist — Test the No Hair Theorem? BH Charge? Matter? Consistency with Numerical Relativity? Beyond GR?
... all stem from older and more fundamental questions ...

The Fundamental Questions about Black Holes



Do black holes exist in nature?

The Fundamental Questions about Black Holes



Do black holes exist in nature?

→ *Are black holes stable?*

(i.e. BH Perturbation Theory)

Review: Black Hole Stability

- ❖ Regge + Wheeler, Edelman (1957) (ten years prior to Wheeler's "black hole" terminology)
 - ❖ **Developed framework** for applying stability analysis to Schwarzschild black holes

$$g'_{\mu\nu} = g_{\mu\nu} + h_{\mu\nu} \rightarrow g'^{\mu\nu} = g^{\mu\nu} - g^{\mu\alpha} g^{\beta\nu} h_{\alpha\beta} + \mathcal{O}(h^2)$$

$$\delta R_{\mu\nu} = 0, \quad h_{\mu\nu} = \sum_{L=1}^{\infty} \sum_{M=-L}^L \sum_{n=1}^{10} C_{LM}^n(t, r) (Y_{LM})_{\mu\nu}(\theta, \varphi)$$

- ❖ Claimed to have proven **stability (i.e. damped solutions)**, but did not consider the correct boundary conditions (coordinate singularity at $r=2M$)
- ❖ Kerr (1963)
 - ❖ Found solution to Einstein's field equations describing **spinning black holes**, parameterized by *dimensionless* spin, $j = S_z/M^2$, and mass, M :

$$ds^2 = (1 - 2Mr/\Sigma)dt^2 + (rj(4M^2\sin^2\theta)/\Sigma)dtd\varphi - (\Sigma/\Delta)dr^2 - \Sigma d\theta^2 \\ - (\sin^2\theta)(r^2 + M^2j^2 + 2M^2j^2r(\sin^2\theta)/\Sigma)d\varphi^2$$

Review: Black Hole Stability

- ❖ Vishveshwara (1969, 1970) (Schwarzschild metric)
 - ❖ First to correctly apply open **boundary conditions** to the perturbative problem in Kruskal coordinates

$$x = r + 2M \ln \left(\frac{r}{2M} - 1 \right)$$

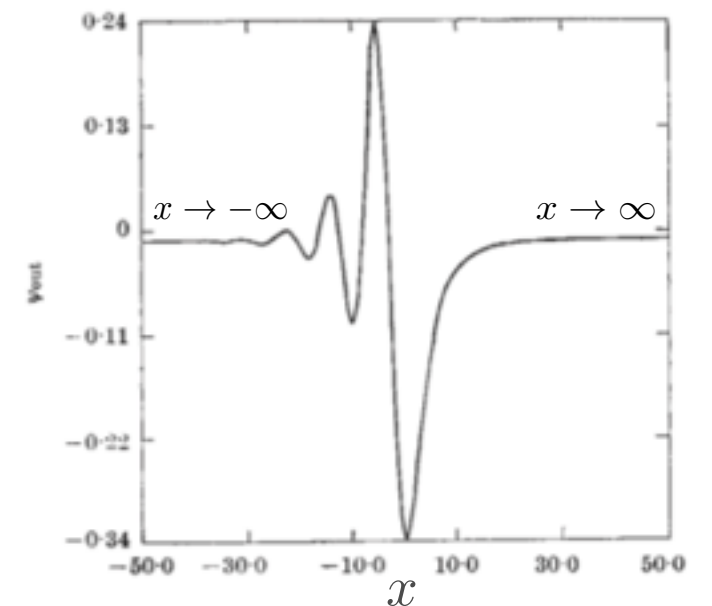
- ❖ Teukolsky (1972), Teukolsky & Press (1973) (Kerr metric)
 - ❖ Applied Newman Penrose formalism (1961) to the perturbation problem
 - ❖ Presented stability analysis of (spinning) Kerr BHs
 - ❖ Found **separable solution** with discrete spectra — the QNMs

$$h_+ - ih_- = \int d\omega \sum_{lm} R(r) {}_{-2}S_{lm}(\theta, \varphi) e^{i\omega t}$$

→ Teukolsky's Equations for $R(r)$ and ${}_{-2}S_{lm}(\theta, \varphi)$

→ Characteristic equations must be satisfied by:

$\tilde{\omega}_{lm}$ (QNM Frequency), and \mathcal{E}_{lm} (Separation Constant)



$$\lim_{r \rightarrow \infty} R(r) = 1/r$$

$$\lim_{j \rightarrow 0} {}_{-2}S_{lm}(\theta, \varphi) = {}_{-2}Y_{lm}(\theta, \varphi)$$

Review: Black Hole Stability

- ❖ Leaver (1985) (Kerr metric)
 - ❖ **Analytic representation** (4D nonlinear) for QNMs of Schwarzschild and Kerr
 - ❖ Typically considered to be the **most accurate** method for QNM calculations
 - ❖ Uses continued fractions to solve characteristic equations for QNM frequency
- ❖ **1985 — 2014 (Abbrev.)**
 - ❖ Various methods for investigation of QNM frequencies: *e.g.* WKB, Laplace transform, Numerical Integration (Schutz, Will, Nollert, Schmidt, Krivan, Laguna, many many others)
 - ❖ Theoretical estimates for QNM measurability (Echeverria, Finn, Flanagan, Hughes, Cardoso, Berti, many others)
 - ❖ Black hole thermodynamics (Bardeen, Bekenstein, Hawking, others; **see Miriam's talk on the Area Theorem**)
 - ❖ **BBH QNM Excitations in Numerical Relativity** (Berti, Buonanno, Pan, Husa, Kamaretsos, Gossan, London, many others)

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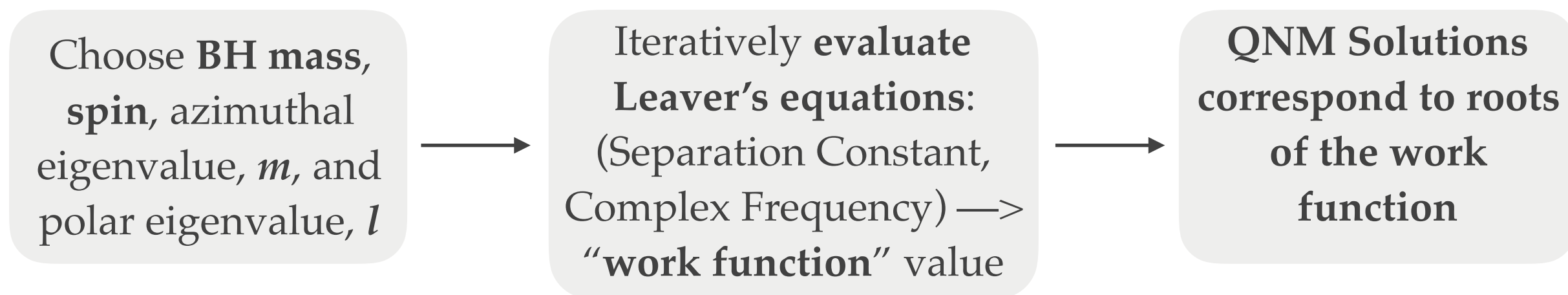
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Review: Structure of QNM Solution Space

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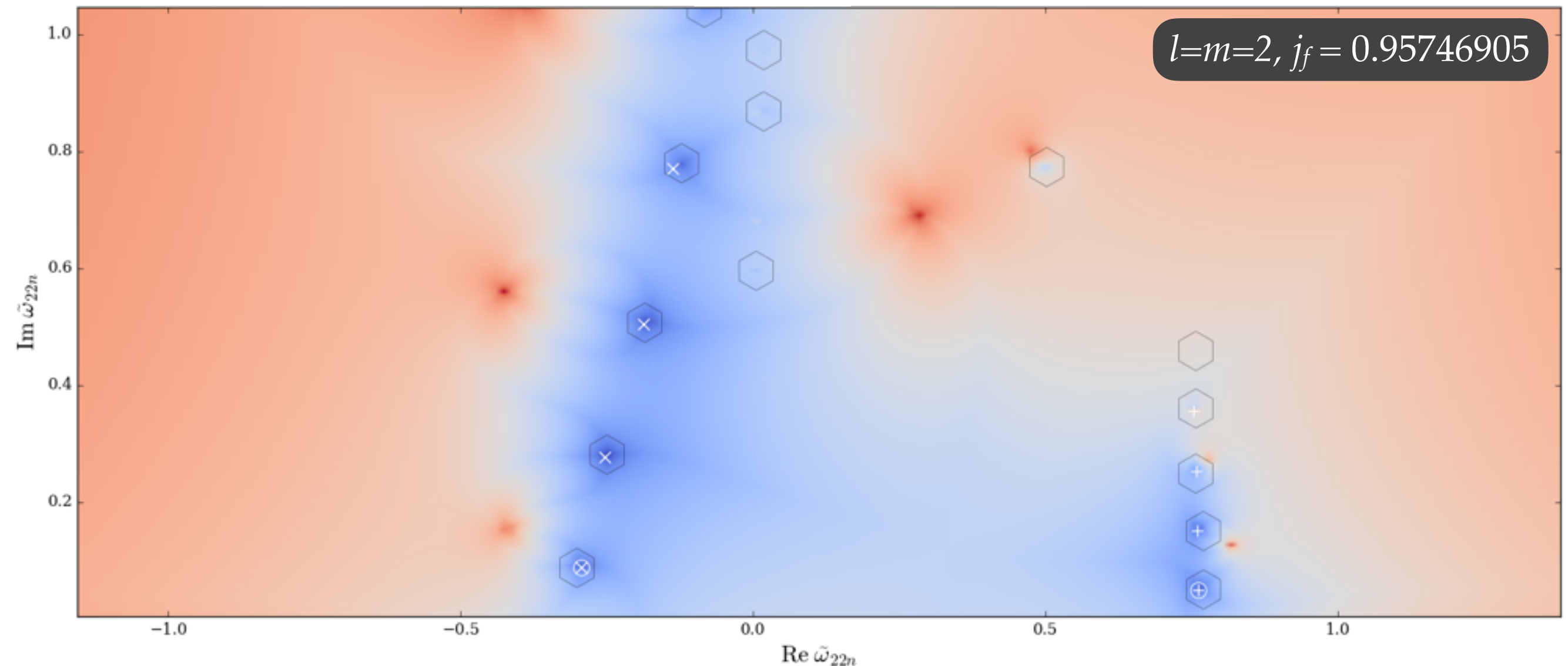
A schematic view of Leaver's method for calculating QNMs



4D Optimization problem:

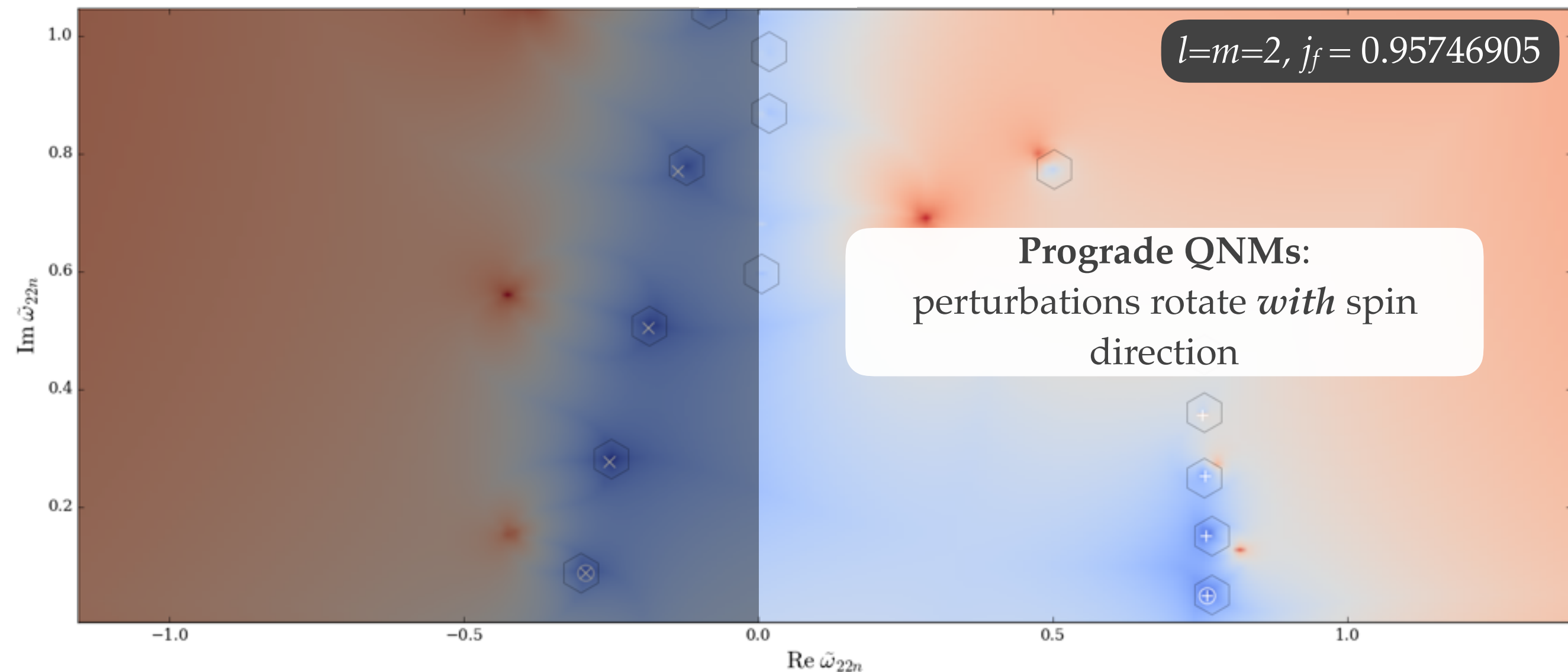
- ❖ Complex valued QNM frequency: $\tilde{\omega}_{lmn} = \omega_{lmn} + i/\tau_{lmn}$ ($M = G = c = 1$)
- ❖ Complex valued separation constant: \mathcal{E}_{lmn} (*approximated analytically by Berti et al*)

Review: Structure of QNM Solution Space



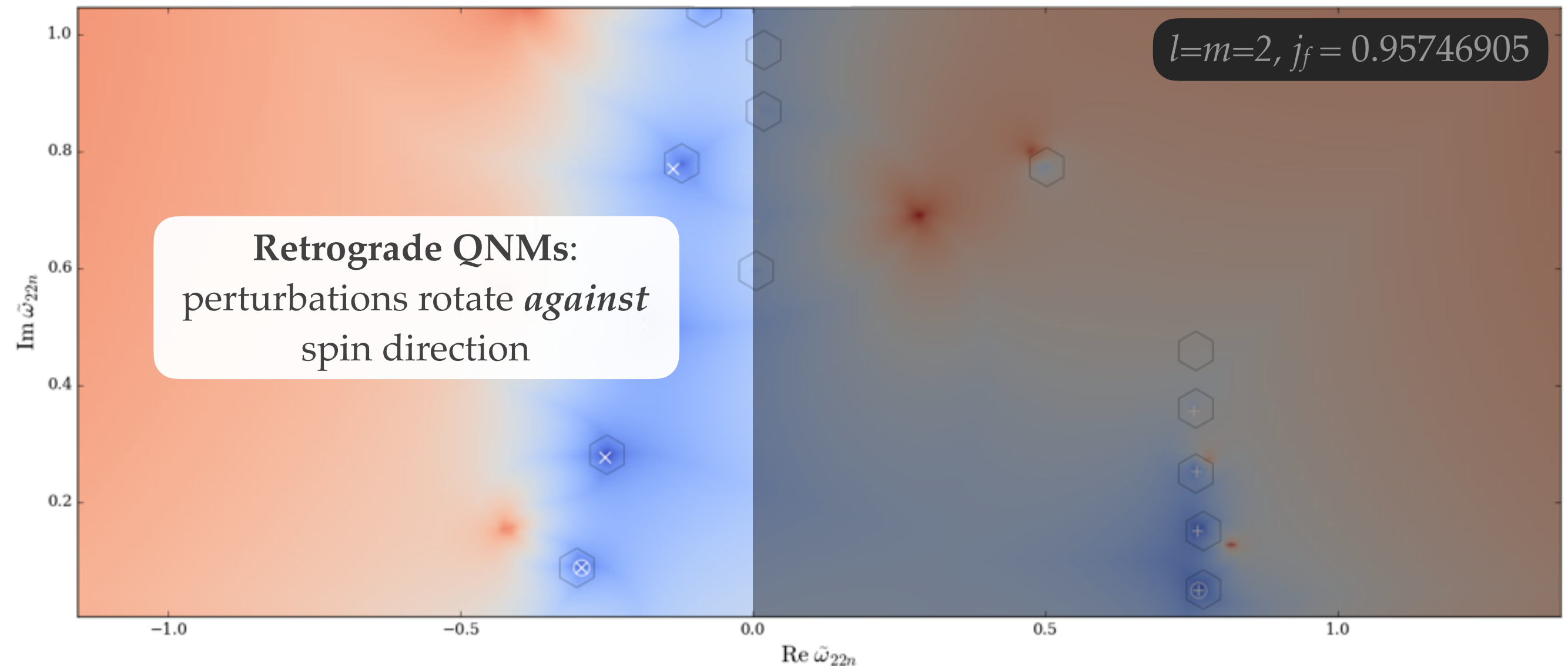
Example: **Blue** (Low work function values), **Red** (High). Using Berti's approximation for the separation constants allows 2D visualization of QNM solution space

Review: Structure of QNM Solution Space



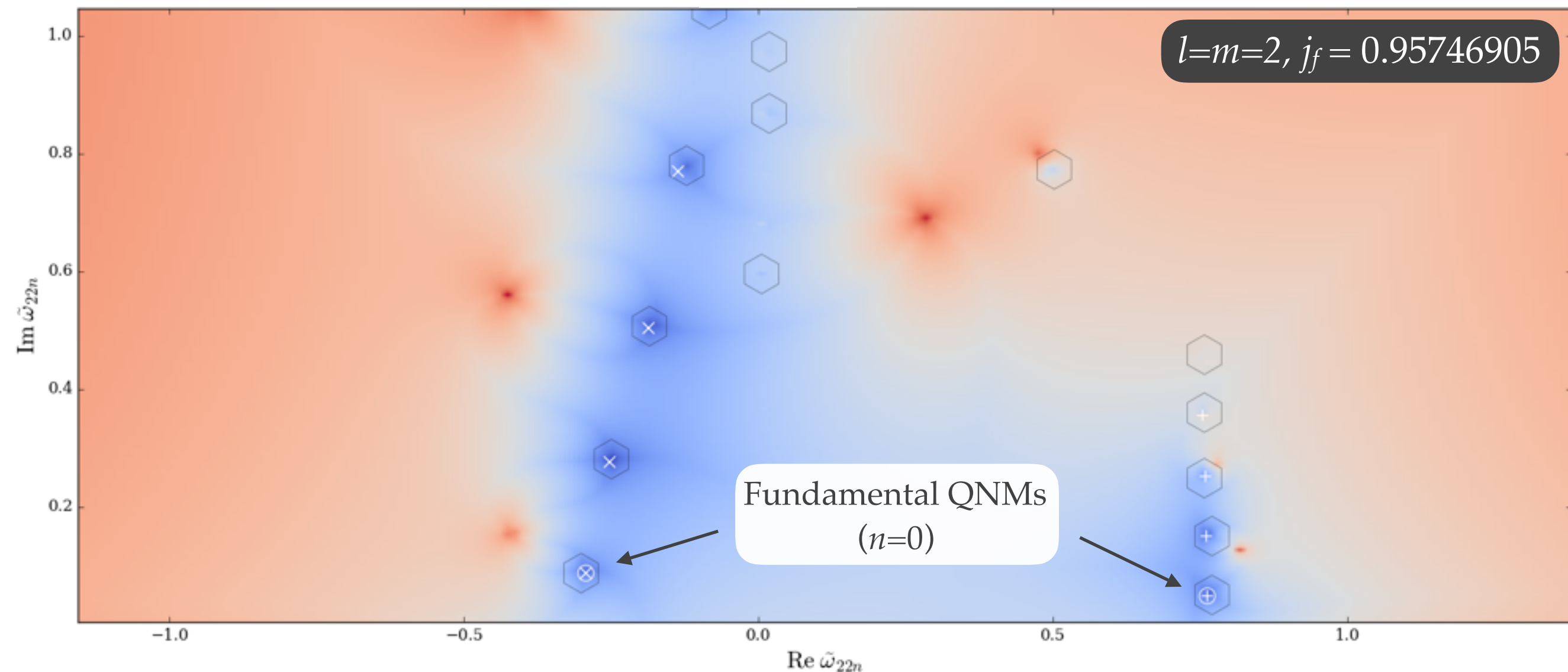
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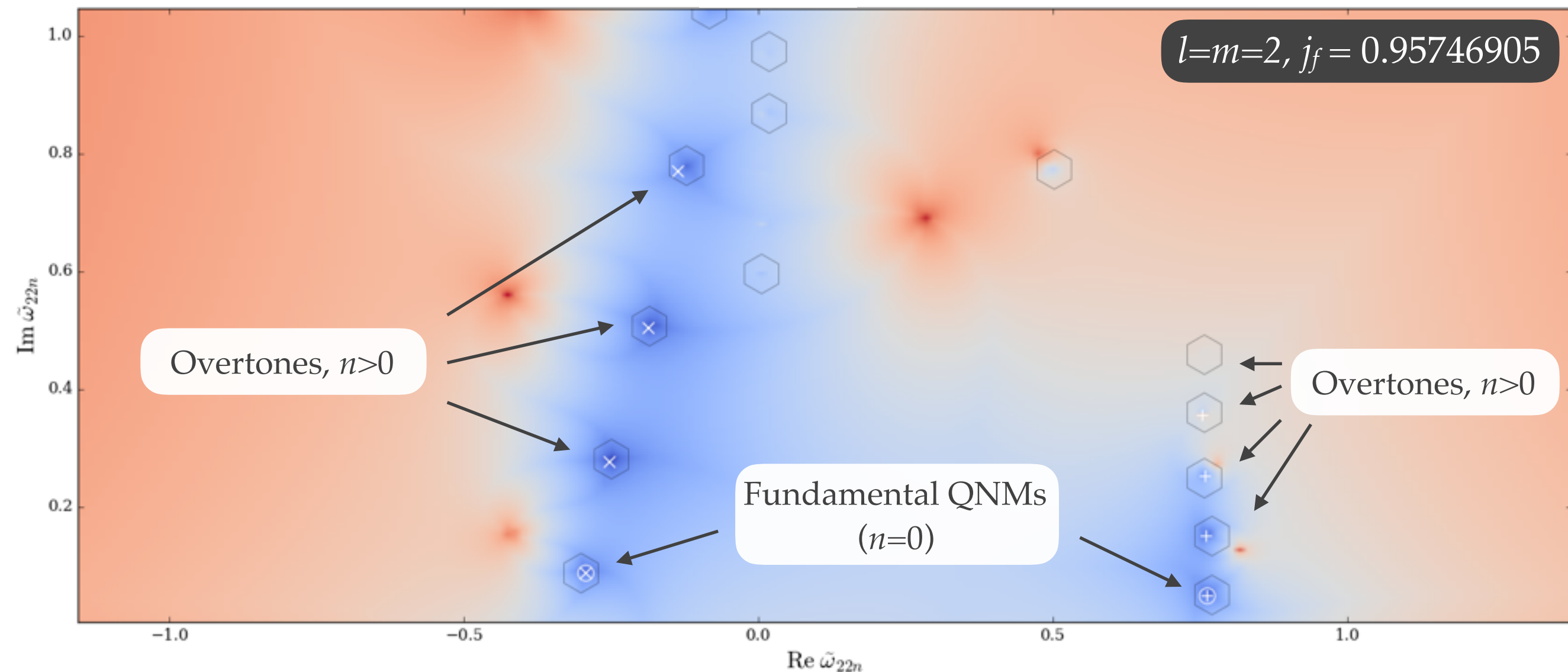
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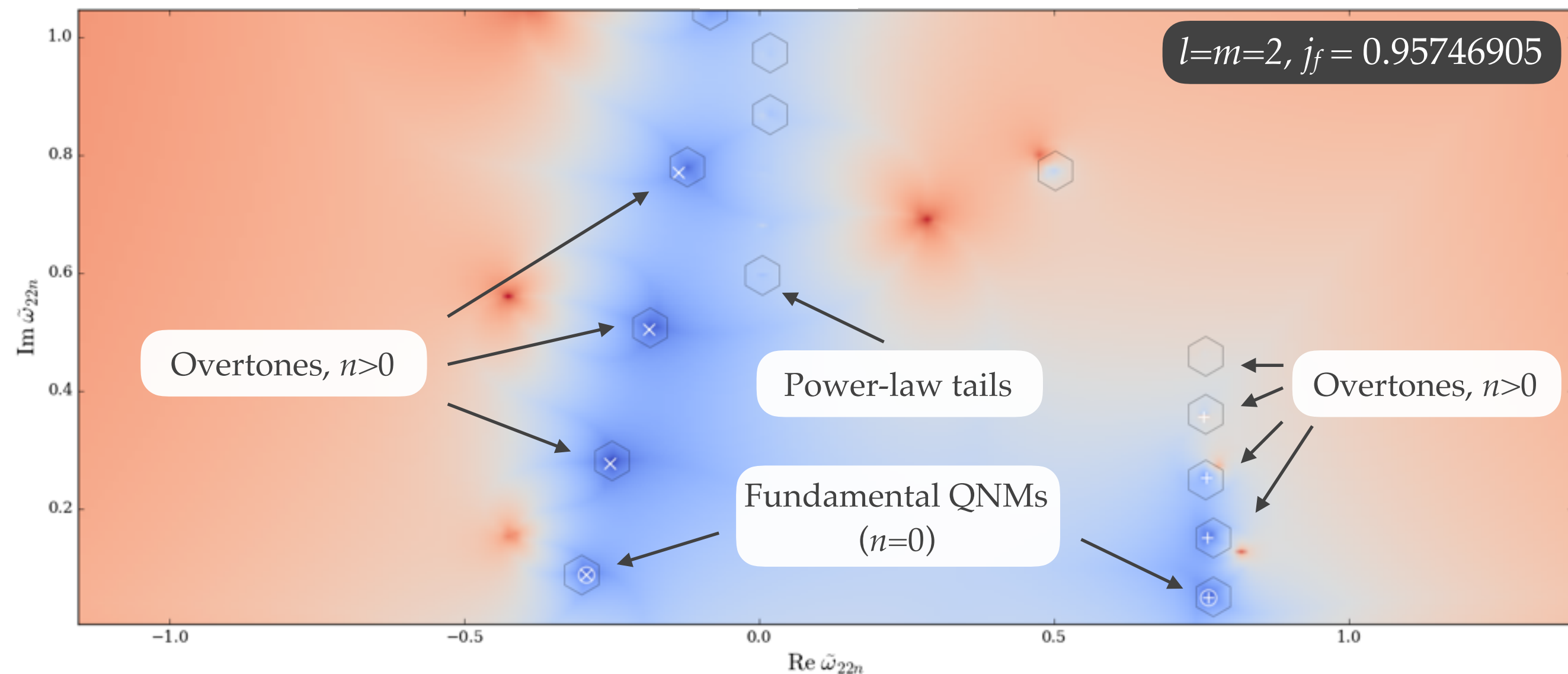
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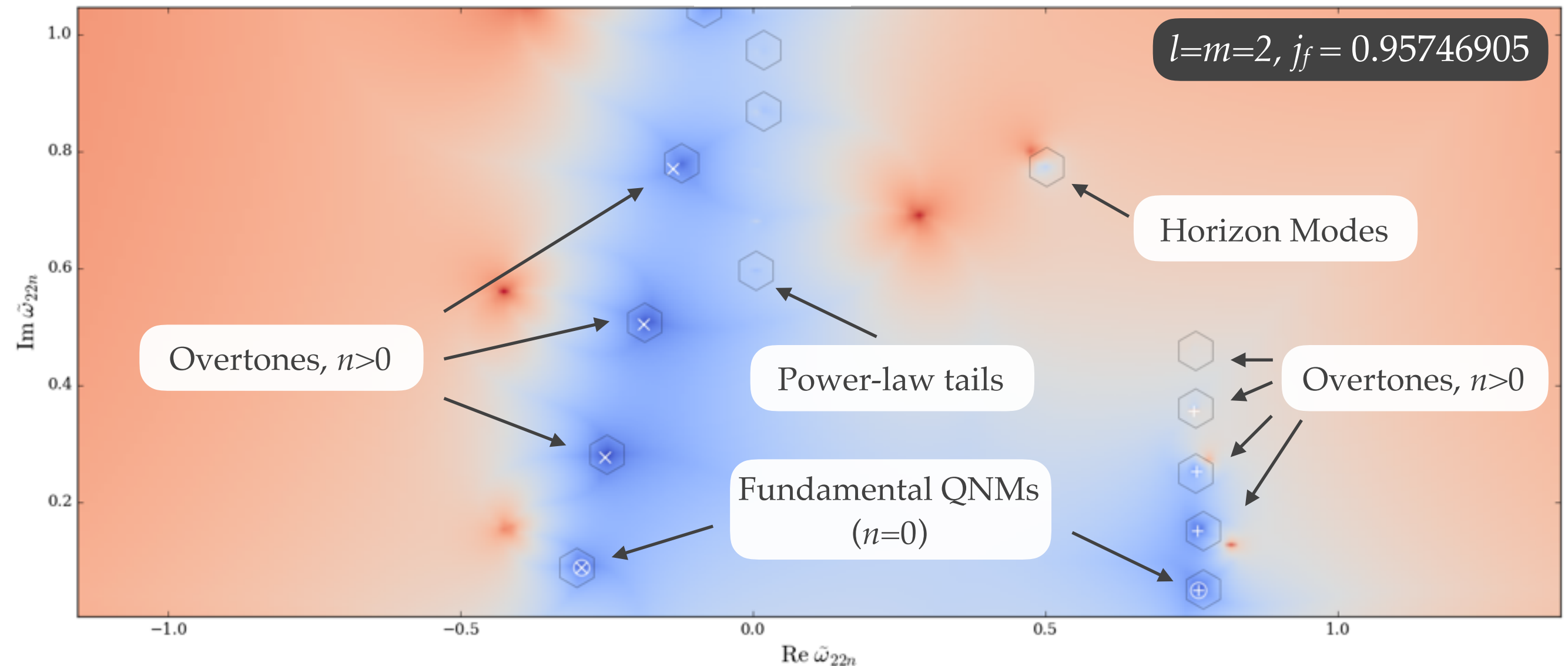
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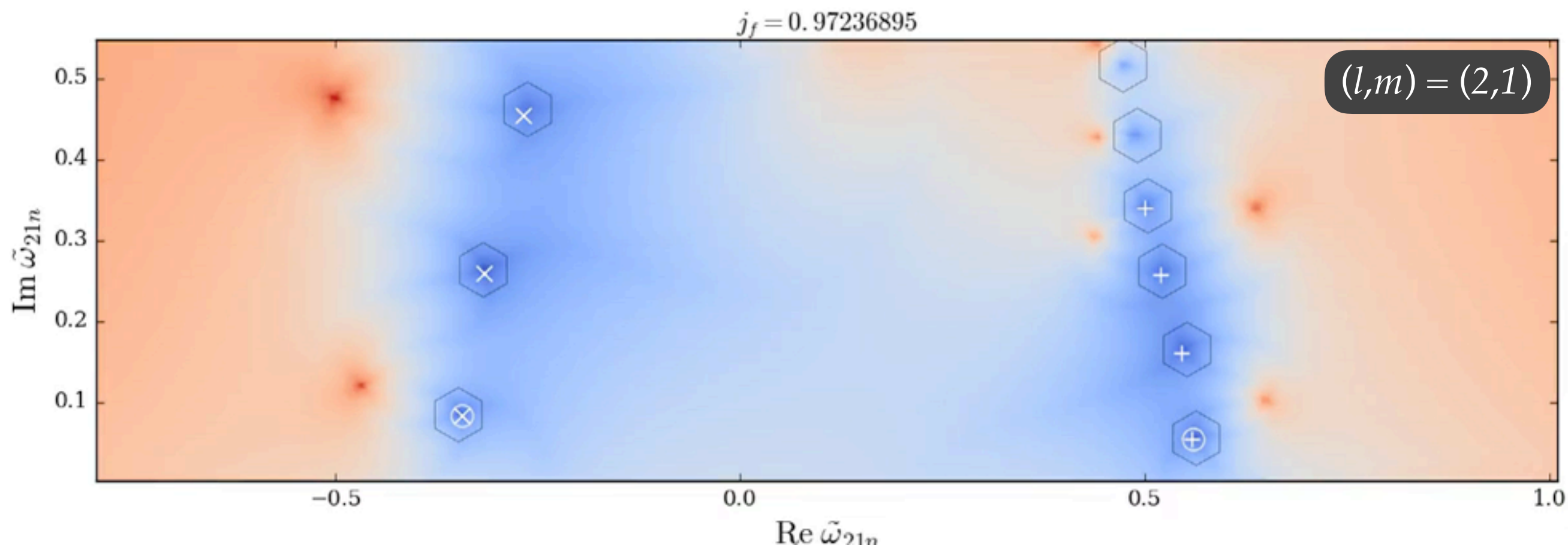
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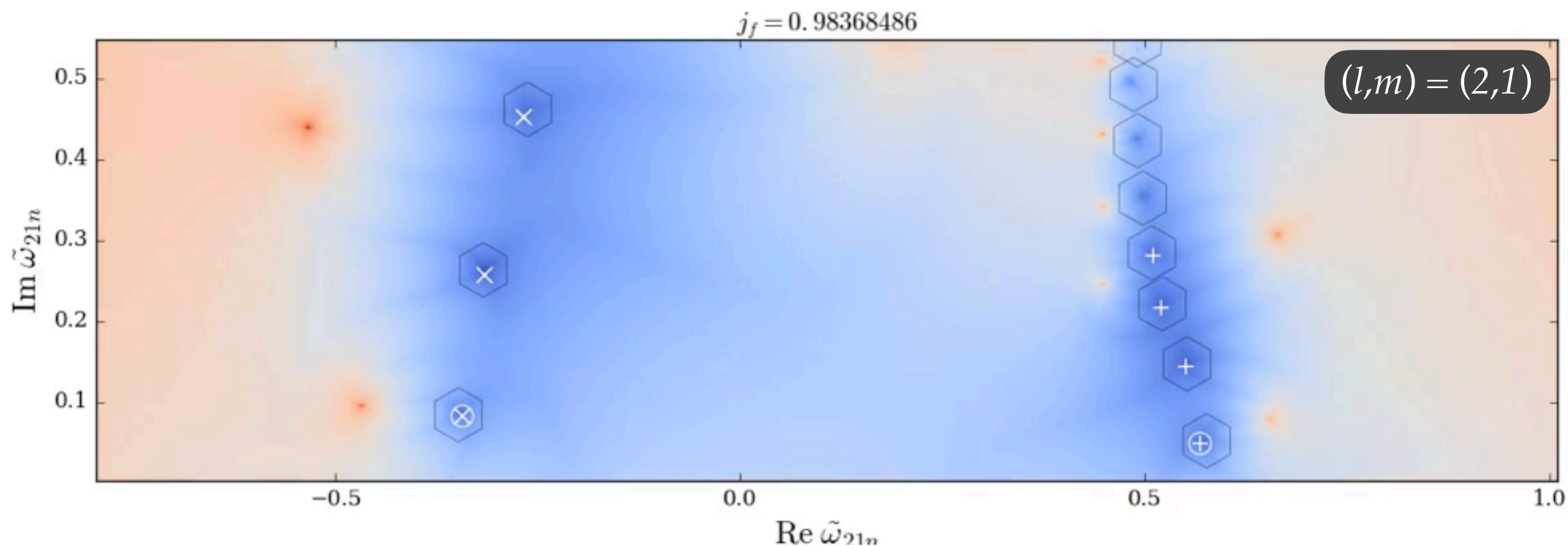
Nontrivial behavior in the limit extremal BH spin ($j_f \sim 1$):
solution branching, and nonzero/**zero damping**



Blue (Low work function values), **Red** (High). Here consider $(l,m)=(2,1)$, and we vary the black hole spin in the nearly extremal regime between $j_f = 0.9837$ and $j_f = 0.9998$. (e.g. A. Zimmerman *et al* 2015)

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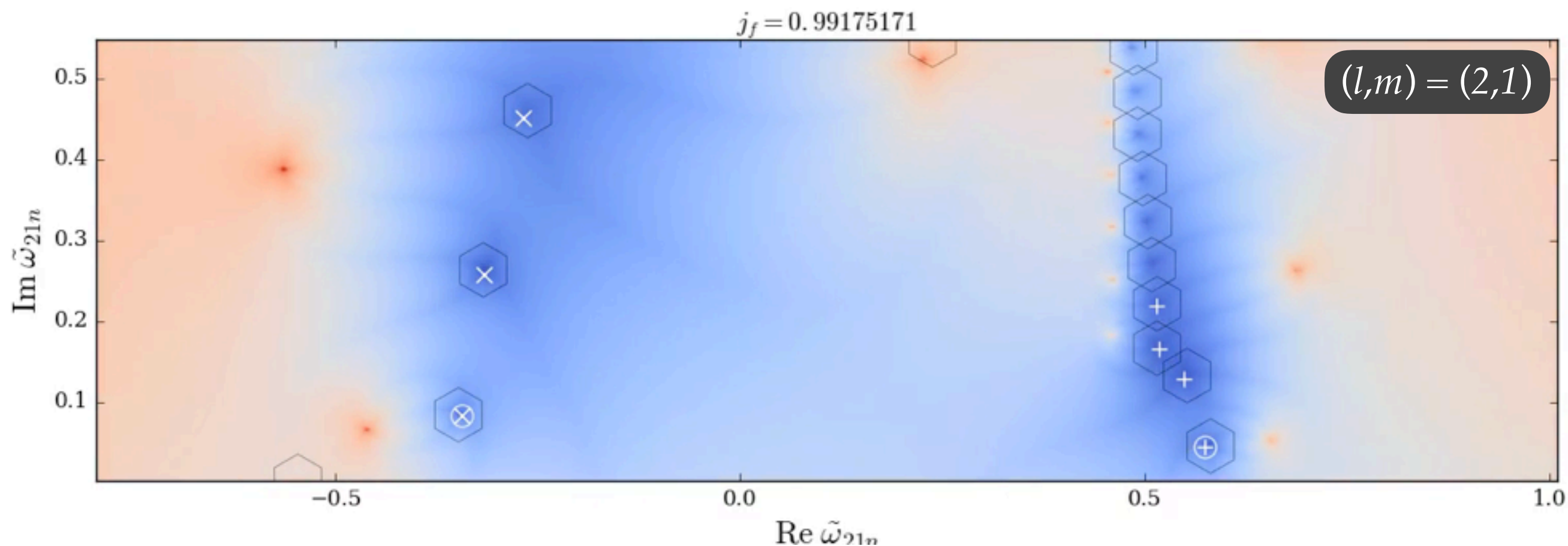
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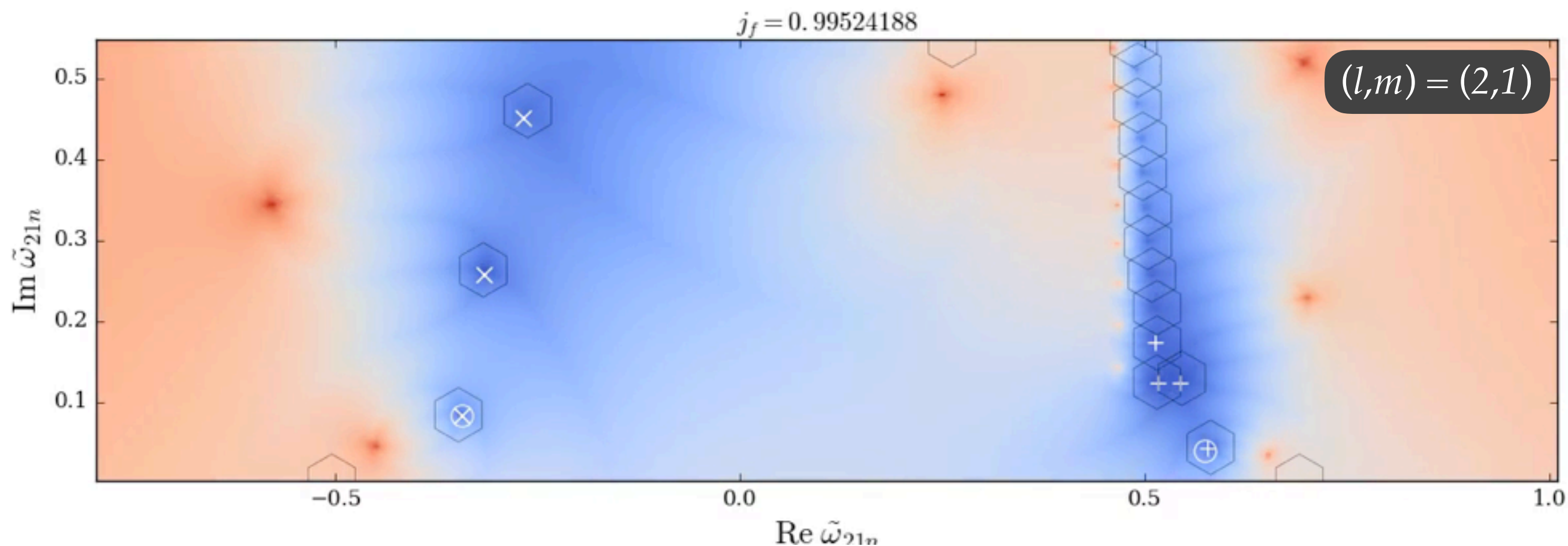
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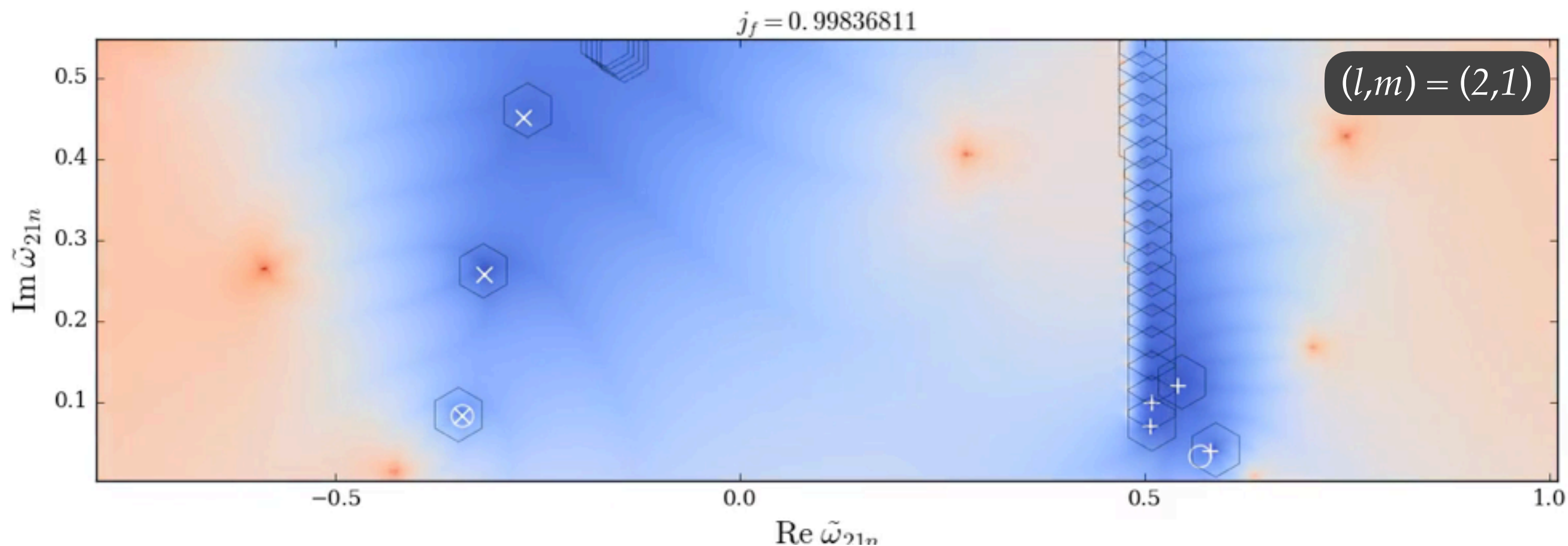
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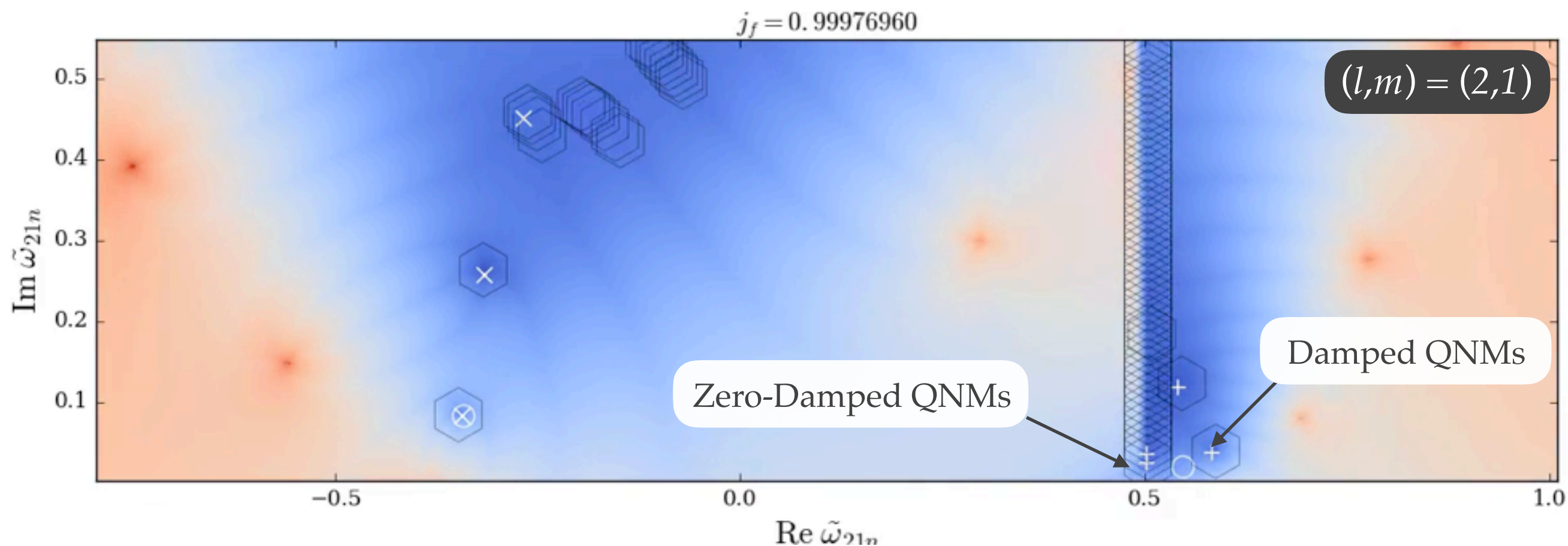
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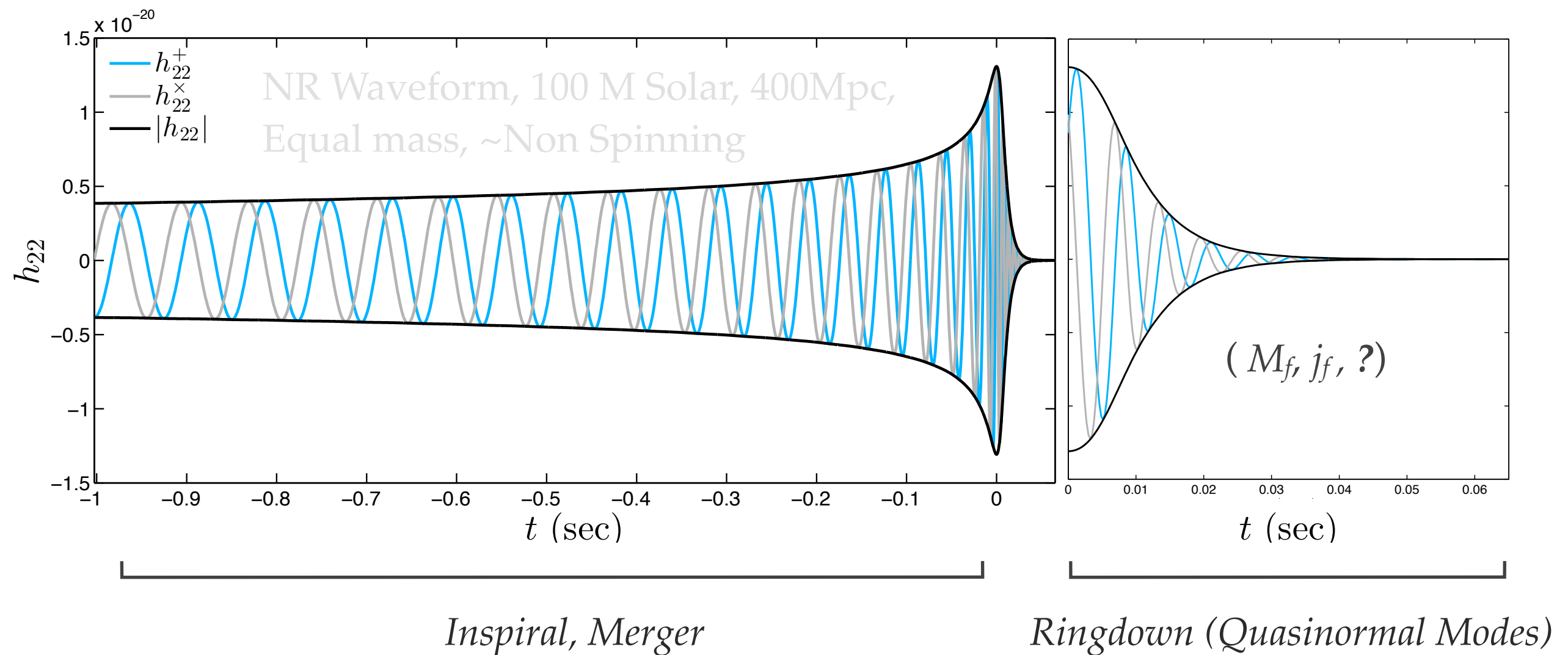
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Recap: Black Hole Stability

- ❖ **Non-extremal Kerr BHs are stable** under generic perturbations.
 - ❖ There is a **significant amount structure in the QNM solution space**. This enhances the prospects for testing GR with astrophysical QNMs.
 - ❖ Of all QNM solutions, **the fundamental modes are the least-damped**, and thus are the **top priority** for analysis of data from experiment.
-
- ❖ **Select Topical Reviews**
 - ❖ “On quasinormal modes of asymptotically anti-de Sitter black holes” (Warnick 2013)
 - ❖ “Quasinormal modes of black holes and black branes” (Berti *et al* 2009)
 - ❖ “Quasinormal modes of stars and black holes” (Kokkotas & Schmidt 1999)
 - ❖ “Quasinormal modes: the characteristic “sound” of black holes and neutron stars” (Nollert 1999)

However, for rigorous QNMs to be useful in LIGO data analysis, **there is an ongoing need to synthesize analytic QNM theory with Numerical Relativity results.**

Review: QNM use in BBH signal models



GW detection and parameter estimation are significantly assisted by model (template) signals. The 2005 advances in Numerical Relativity (NR) (*eg Baker et al*) have enabled GW models that encompass **Inspiral, Merger and Ringdown** (IMR). Despite the success of these models, **there have been (and remain) difficulties relating to ringdown ...**

Review: QNM use in BBH signal models

Overview

- ❖ Starting with **Ajith's 2007 work**, Kerr QNMs have been used to parameterize **Phenomenological waveform models**.
- ❖ **IMRPhenomD** (Khan *et al* and Husa *et al* 2015) and the related **IMRPhenomP** (Schmidt, Hannam *et al* 2013) have been heavily used in the analysis of GW150914 and subsequent events.
- ❖ In parallel, **Effective One Body** (EOB) approaches to BBH inspiral have been extended using Kerr QNMs (Damour, Buonanno, Pan, Taracchini, others)
- ❖ The related **SEOBNRv2** (Taracchini *et al* 2013) has also been heavily used in post-detection LIGO data analysis.

Status

All of these models require the tuning to NR simulations to incorporate how much each QNM is excited. However, **until recently, there was no general and robust way to model fundamental QNM excitations from NR simulations**

Review: QNM use in BBH signal models

The problem lies in a difference of perspective

❖ **Numerical Relativity:** Spherical Harmonic Multipoles, ${}_{-2}Y_{lm}$ **(Orthogonal in l)**

$$\diamond rh = r(h_+ - ih_\times) = \sum_{l,m} h_{lm}^{NR} {}_{-2}Y_{lm}(\theta, \phi)$$

$$\diamond h_{lm}^{NR} = \int_{\Omega} rh {}_{-2}\bar{Y}_{lm} d\Omega$$

❖ **Perturbation Theory:** Spheroidal Harmonic Multipoles, ${}_{-2}S_{lm}$ **(Not orthogonal in l)**

$$\diamond rh = \sum_{lmn} h_{lmn}^{PT} {}_{-2}S_{lm}(\theta, \phi; \tilde{\omega}_{lmn} j_f)$$

$$\diamond h_{lmn}^{PT} = A_{lmn} [e^{i\tilde{\omega}_{lmn} t}]$$

$$\rightarrow h_{lm}^{NR} = \sum_{l'n} A_{l'mn} e^{i\tilde{\omega}_{l'mn} t} \int_{\Omega} {}_{-2}\bar{Y}_{lm} {}_{-2}S_{l'mn} d\Omega$$

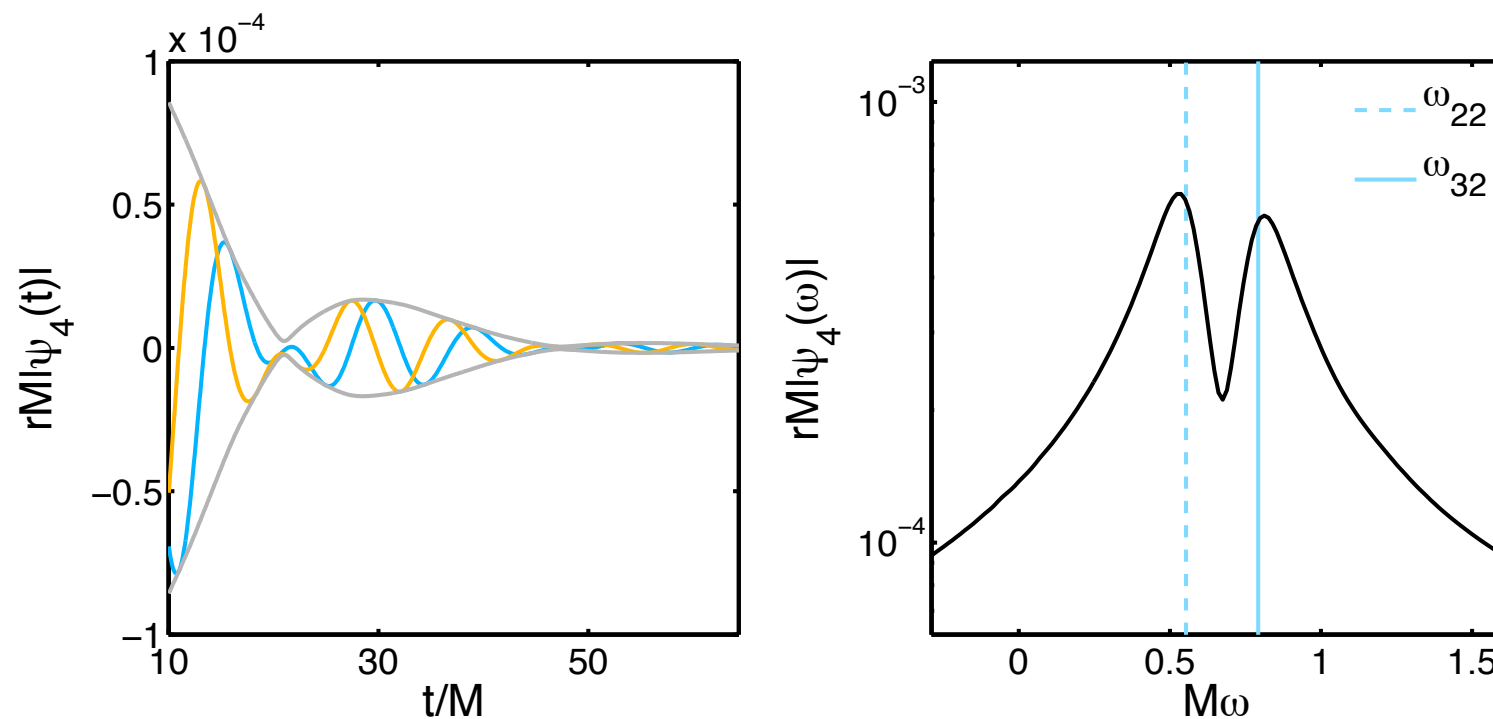
The Spherical multipoles of NR
are sums of QNMs
— i.e. “mode mixing”

Aside: Mixing of QNMs in NR Waveforms

$$\rightarrow h_{lm}^{NR} = \sum_{l'n} A_{l'n} e^{i\tilde{\omega}_{l'n} t} \int_{\Omega} -2\bar{Y}_{lm} - 2S_{l'n} d\Omega$$

The Spherical multipoles of NR are sums of QNMs
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Effect on time domain ringdown. Figure: NR ringdown waveform for $(l,m)=(3,2)$, equal mass, non-spinning.



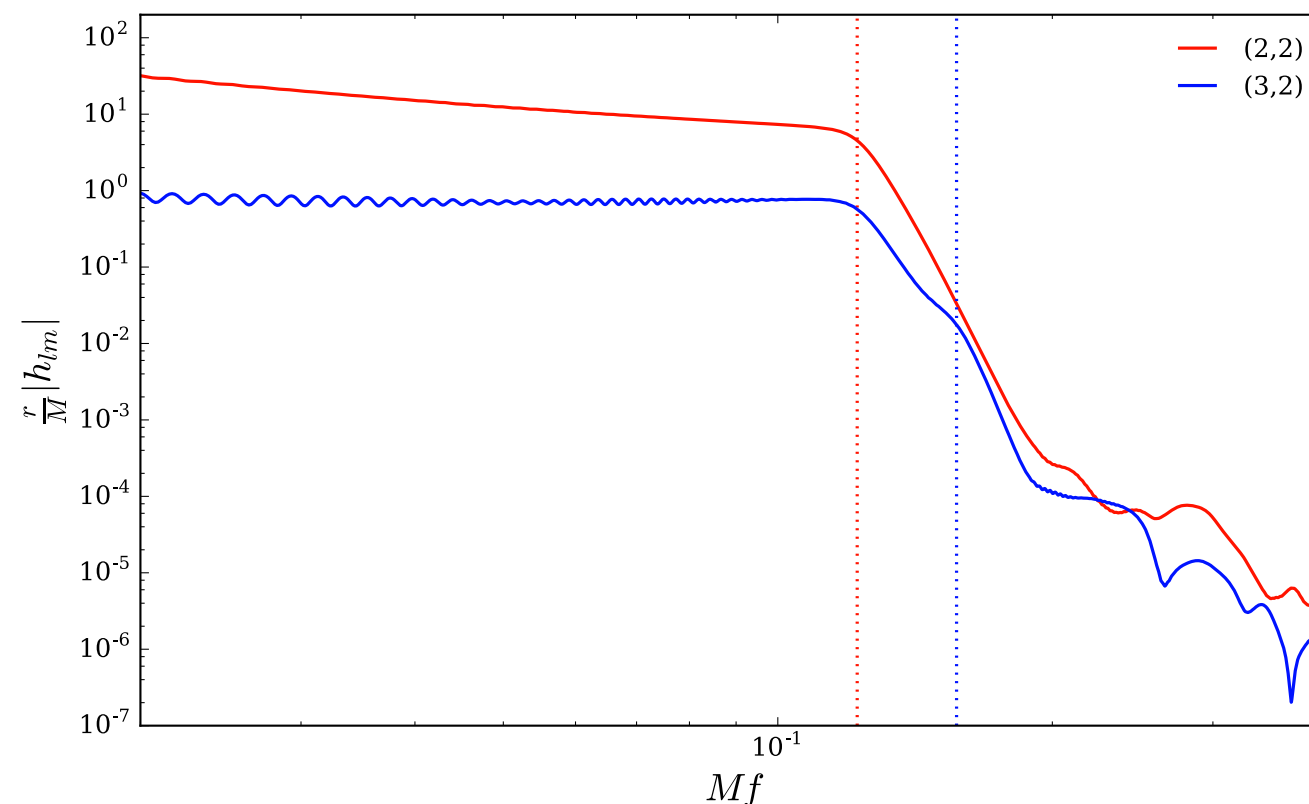
In the time domain, the **mixing of QNMs due to use of a spherical basis causes “beating”**, while a simple decaying sinusoid might naively be expected. **This effect complicates the time domain modeling of post-merger higher multipoles.**

Aside: Mixing of QNMs in NR Waveforms

$$\rightarrow h_{lm}^{NR} = \sum_{l'n} A_{l'mn} e^{i\tilde{\omega}_{l'mn}t} \int_{\Omega} {}_{-2}\bar{Y}_{lm} {}_{-2}S_{l'mn} d\Omega$$

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Effect on frequency domain IMR waveform. Figure: NR IMR waveform for **(l,m)=(2,2)** and **(3,2)**, equal mass, dimensionless aligned spins of 0.85 (z-hat).



There are similar complications for the frequency domain IMR waveform.

Review: QNM use in BBH signal models

Overview

- ❖ Early studies focused primarily on the $(l,m)=(2,2)$, where mode mixing is minimal. (2005-2013)
- ❖ Kamaretsos *et al* (2011) modeled the **QNM amplitudes** for non-spinning BBH systems in the Spherical basis (effectively treating the remnant as non-spinning)
- ❖ L. London *et al* (2014) developed a method for estimating Spheroidal content from NR waveforms, and modeled **QNM amplitudes and relative phases** for non spinning systems using complex polynomial regression.
- ❖ L. London *et al* (paper in preparation) have **used QNMs to develop the first IMR waveform model for non-precessing BBH sources**. [\(See Sebastian Khan's talk\)](#)

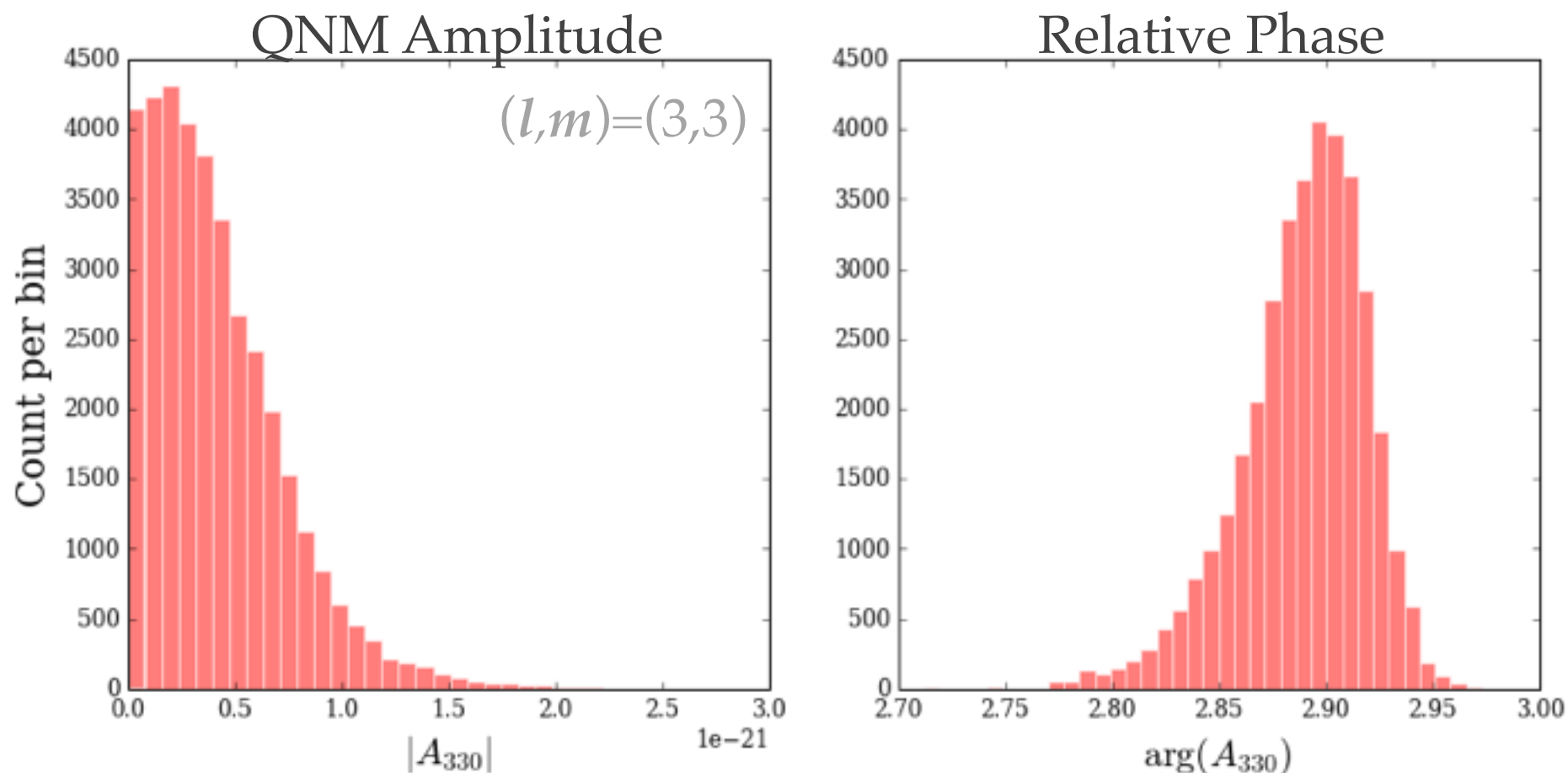
Update — A QNM Waveform Model for Non-Precessing BBH Systems

There is a clear need to extend result to precessing BBH systems. Towards this goal, I have recently developed a **effective-spin model for QNM excitation**, which is to be used as the **GR prediction** when perform more general data analysis on astrophysical signals.

Update: A QNM Waveform Model for NonPrecessing Systems

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Example usage: Estimated QNM amplitude (GR prediction) for GW150914



Current and New Questions about BH QNMs

Current Questions

- ❖ Given GR predictions for QNM excitations (*e.g.* London et al 2014)
 - ❖ Can we perform a test of the No Hair Theorem? (See Chandra's talk today, and Archisman's talk on Thursday)

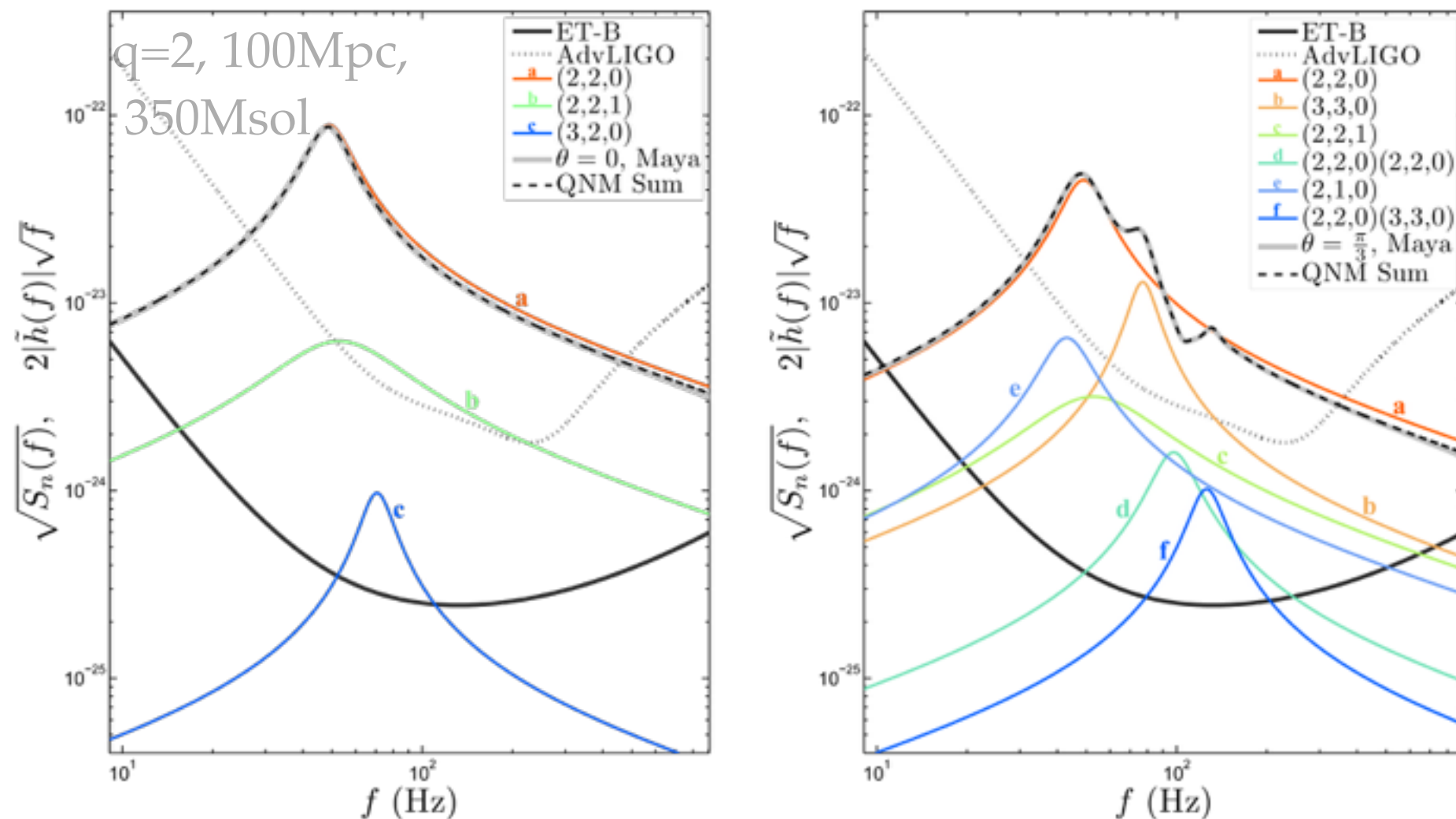
Can we “coherently stack ringdown signals” to improve statistics? This has been preliminarily investigated by Yang, Yagi et al (2017) (also see Felipe's talk)

Can we use QNM information to develop IMR models with higher multipoles (See Sebastian's talk)

New (more recent) Questions

- ❖ Can we distinguish BH QNMs from BH mimickers? (*e.g.* Cardoso *et al* 2016)
- ❖ Beyond GR effects? (See Archisman's talk on Thursday)
- ❖ What will be needed to model QNM excitations for precessing systems?
- ❖ Will the 2nd order (nonlinear) QNMs be relevant for GW astronomy? (London et al 2014)

When can we expect to confidently extract multiple QNMs?



(Figure from London *et al* 2014)

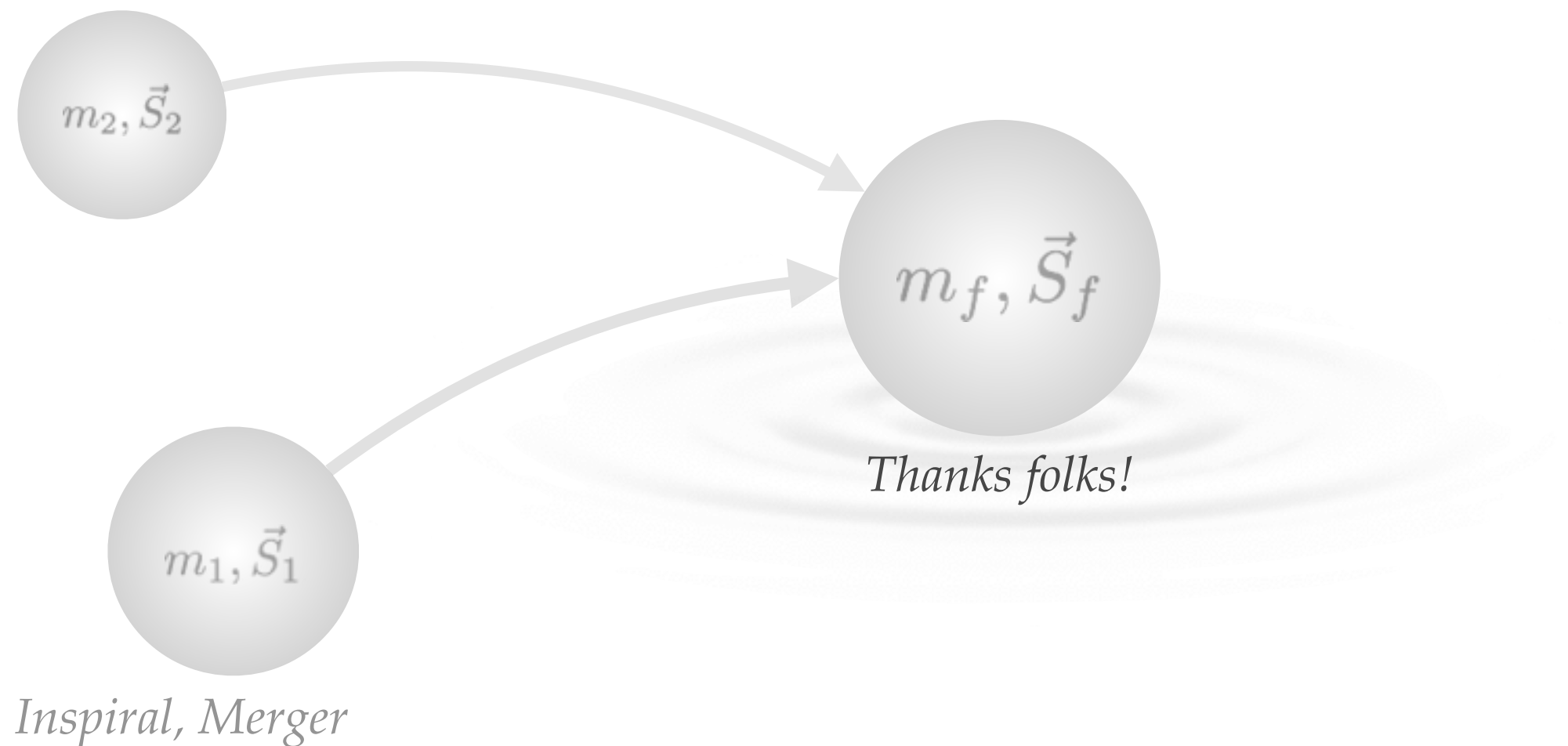
- ❖ Most evidence suggests extraction QNM signals will be routine with the **Einstein telescope** (*e.g.* Maselli *et al* 2017, Gossan *et al* 2011, Berti, Cardoso, London, others)
- ❖ **However, there are always new possibilities on the horizon ...**

The Quasi-normal Modes of Black Holes

Review and Recent Updates

Concluding Remarks

- ❖ We've seen that the topic of **BH perturbation theory** grew out of pondering the very **existence of BHs in nature**.
- ❖ QNM solution space is generally rich with features. **The fundamental QNMs are the least damped, and therefore the most likely to be observed first.**
- ❖ QNM solutions represent eigenfunctions of perturbed Einstein equations, but **NR simulations are currently needed to inform how much each QNM is excited.**
- ❖ Models of QNM excitation can **assist with tests of GR, and inform IMR signal models (some updates in later talks)**
- ❖ The extraction of **astrophysical QNMs (beyond $l=m=2$)** may take some time **(Einstein Telescope)**, but there is some uncertainty about how much ...



Context: Quasinormal models resulting from the merger of stellar mass BHs, and learning as much as we can from post-merger (ringdown) signals ...

The Quasi-normal Modes of Black Holes

Review and Recent Updates

