

Renormalization factors in QCD:

The QCD Lagrangian is structurally similar to QED:

$$L_{QCD} = \sum_q \bar{q}(i\cancel{D} - m_q)q - \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + \theta\text{-term (?)}$$

+ gauge fixing + Faddeev-Popov ghosts

where:

$$i\cancel{D} = i\cancel{\partial} + g_s A_a^\mu t_a$$

↑ strong coupling

↖ SU(3) generators

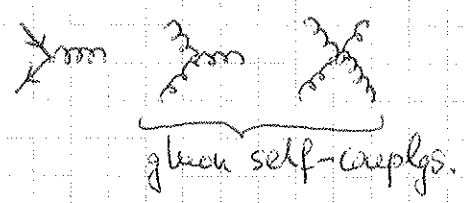
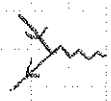
Similarities and differences:

gauge group:

QED  
U(1) abelian

QCD  
SU(N<sub>c</sub>=3) non-abelian

vertices:



Faddeev-Popov ghosts: —  
(anticommuting scalar fields → wrong spin-statistics relation → not as external states)

needed to cancel unphysical gluon polarizations  
 $| \text{ghost loop} | + | \text{ghost loop} |$   
 = physical

non-trivial group factors: —

$$C_F = \frac{N_c^2 - 1}{2N_c}, C_A = N_c, T_F = \frac{1}{2}$$

superficial degree of divergence:  $D = 4 - \frac{3}{2}N_c - N_g$

$$D = 4 - \frac{3}{2}N_q - N_g - \frac{3}{2}N_c$$

↑ ghosts

phenomenology: weak coupling at all relevant energies (Landau pole > M<sub>PL</sub>)

strong coupling at low energies → confinement; weak coupling at high energies (asymptotic freedom)

## Important note:

While the on-shell renormalization scheme is useful for many (but not all  $\rightarrow$  see below) calculations in QED and the EW theory, it is not a viable scheme for QCD:

- quarks and gluons are never on-shell due to confinement
- coupling  $g_s$  cannot be renormalized at  $q^2=0$ , since QCD is strongly coupled at low energy and quarks and gluons are not the relevant degrees of freedom there
- masses of light quarks ( $q = u, d, s$ ) satisfy  $m_q \ll \Lambda_{\text{QCD}}$  and hence can be set to zero  
 $\Rightarrow$  will not consider heavy quarks here, which are usually dealt with using EFTs (several lecture courses in this school!)

$\hookrightarrow$  will use  $\overline{\text{MS}}$  scheme in QCD!

Superficial degree of divergence (exercise). (skip this!)

i)  $D = 4L - P_f - 2P_g - P_c + V_{3g} \leftarrow$  contains one power of momentum

ii)  $L = P_f + P_g + P_c - (V_{qg} + V_{3g} + V_{4g} + V_{cg}) + 1$

iii)  $V_{qg} = \frac{1}{2} (2P_f + N_f)$ ;  $V_{cg} = \frac{1}{2} (2P_c + N_c)$ ;

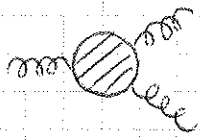
$$V_{qg} + 3V_{3g} + 4V_{4g} + V_{cg} = 2P_g + N_c$$

Combining this yields

$$D = 4 - N_g - \frac{3}{2} N_f - \frac{3}{2} N_c \quad (\text{note: } N_c = 0 \text{ for physical amplitudes})$$

$\uparrow$   
ghosts

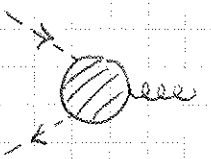
In addition to QED, the following amplitudes have logarithmic UV divergences and require renormalization:



$D=1$ , but log. div. due to gauge invariance  
(Furry's theorem does not apply to QED)

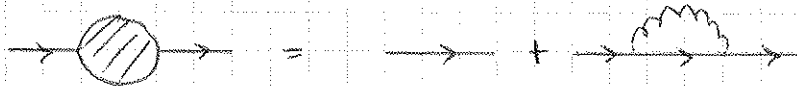


$D=0$ , log. divergent



$D=0$ , log. divergent

Quark self-energy:

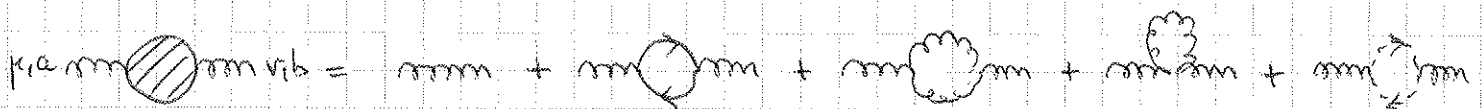


↑  
same as QED ( $\alpha \rightarrow \alpha_s$ ) times color factor  $\text{tr}(t_a t_a) = C_F$

hence:

$$Z_2 = 1 - \frac{C_F \alpha_s}{4\pi \hat{\epsilon}}$$

Vacuum polarization:



↑  
same as QED (with  $\alpha \rightarrow \alpha_s$ )  
times color factor  
 $\text{tr}(t_a t_b) = T_F \delta_{ab}$

genuinely non-abelian!

$$= \delta^{ab} (g_s^{\mu\nu} q^2 - q^\mu q^\nu) \Pi(q^2)$$

find:

$$\Pi(q^2) = \frac{\alpha_s}{4\pi} \left\{ \left[ \left( \frac{13}{6} - \frac{\xi}{2} \right) C_A - \frac{4}{3} T_F n_f \right] \left( \frac{1}{\hat{\epsilon}} - \ln \frac{q^2 - i0}{\mu^2} \right) + \dots \right\}$$

↑  
gauge parameter ( $\xi = 1/\lambda$ )

hence:

$$Z_3 = \frac{1}{1 - \pi(q^2)} \Big|_{\text{pole term}}$$

$$= 1 + \frac{\alpha_s}{4\pi\hat{\epsilon}} \left[ \left( \frac{13}{6} - \frac{\epsilon}{2} \right) C_A - \frac{4}{3} T_F n_f \right]$$

### Charge renormalization:

Gauge invariance requires that the same charge appears in all four vertices:

~~non~~

$$g_{s,0} = \mu^{\epsilon} \underbrace{Z_1 Z_2^{-1} Z_3^{-1/2}}_{\text{no longer } = 1 \text{ in QCD!}} g_s$$

~~anom~~

$$g_{s,0} = \mu^{\epsilon} Z_1^{3g} Z_3^{-3/2} g_s$$

~~anom~~

$$g_{s,0} = \mu^{\epsilon} Z_1^{4g} Z_3^{-2} g_s$$

~~anom~~

$$g_{s,0} = \mu^{\epsilon} Z_1^c Z_2^{-1} Z_3^{-1/2} g_s \quad \left( \rightarrow \text{simplest way to calculate charge-ren. constant} \right)$$

It follows that:

$$Z_1^{3g} = Z_1 Z_3 Z_2^{-1}; \quad Z_1^{4g} = Z_1 Z_2^{-1} Z_3^{3/2}; \quad Z_1^c = Z_1 Z_2 Z_2^{-1}$$

$\rightarrow$  exact relations between renormalization constants

We work with the first relation and compute the renormalization constant  $Z_1$  of the quark-gluon vertex:



needs:  $f^{abc} \frac{b,c}{t} = \frac{i}{2} C_A t^a$

One finds:

$$Z_1 = 1 - \frac{\alpha_s}{4\pi \hat{\epsilon}} \left[ C_F + \frac{3+\hat{\epsilon}}{4} C_A \right] \neq Z_2$$

$\uparrow$  similar to QED       $\uparrow$  non-abelian (gauge-dependent)

We now obtain:

$$g_{\bar{1}0} = \mu^{\epsilon} Z_1 Z_2^{-1} Z_3^{-1/2} g_s$$

with:

$$Z_1 Z_2^{-1} Z_3^{-1/2} = 1 - \frac{\alpha_s}{4\pi \hat{\epsilon}} \left( \frac{11}{6} C_A - \frac{2}{3} T_F n_f \right)$$

[cf. QED:

$$1 + \frac{\alpha}{6\pi \hat{\epsilon}} n_l]$$

different signs!

similar