

Introduction to EFT: Problems

1. Show that for a *connected* graph, $V - I + L = 1$. What is the formula if the graph has n connected components?
2. Work out the properties of fermion bilinears $\bar{\psi}(\mathbf{x}, t)\Gamma P_L\chi(\mathbf{x}, t)$ under C, P, T , where $\Gamma = 1, \gamma^\mu, \sigma^{\mu\nu}$. The results for $P_L \rightarrow P_R$ can be obtained by using $L \leftrightarrow R$.
3. Fierz identities are relations of the form

$$(\bar{A}\Gamma B)(\bar{C}\Gamma D) = \sum_i (\bar{C}\Gamma_i B)(\bar{A}\Gamma_i D)$$

where A, B, C, D are fermion fields. They are much simpler if written in terms of chiral fields, where $\Gamma_i = 1, \gamma^\mu, \sigma^{\mu\nu}$. Work out the Fierz relations for

$$\begin{aligned} &(\bar{A}P_L B)(\bar{C}P_L D), (\bar{A}P_L B)(\bar{C}P_R D), (\bar{A}\gamma^\mu P_L B)(\bar{C}\gamma_\mu P_L D), (\bar{A}\gamma^\mu P_L B)(\bar{C}\gamma_\mu P_R D), \\ &(\bar{A}\sigma^{\mu\nu} P_L B)(\bar{C}\sigma_{\mu\nu} P_L D), (\bar{A}\sigma^{\mu\nu} P_L B)(\bar{C}\sigma_{\mu\nu} P_R D) \end{aligned}$$

The $P_R P_R$ identities are given by using $L \leftrightarrow R$ on the $L L$ identities. Do not forget the Fermi minus sign.

4. In $d = 4$ spacetime dimensions, work out the field content of Lorentz-invariant operators with dimension n for $n = 1, \dots, 6$. At this point, do not try and work out which operators are independent, just the possible structure of allowed operators. Use the notation ϕ for a scalar, ψ for a fermion, $X_{\mu\nu}$ for a field strength, and D for a derivative. For example, an operator of type $\phi^2 D$ such as $\phi D_\mu \phi$ is not allowed because it is not Lorentz-invariant. An operator of type $\phi^2 D^2$ could be either $D_\mu \phi D^\mu \phi$ or $\phi D^2 \phi$, so a $\phi^2 D^2$ operator is allowed, and we will worry later about how many independent operators $\phi^2 D^2$ we can construct.
5. For $d = 2, 3, 4, 5, 6$ dimensions, work out the field content of operators with dimension $n \leq d$, i.e. the “renormalizable” operators.
6. Show that if $\alpha_s(\mu)$ is fixed at some high scale, say $\mu = 1 \text{ TeV}$, then $m_p \propto m_t^{2/27}$, where m_p is the proton mass and m_t is the top quark mass.
7. Compute in dimensional regularization in $d = 4 - 2\epsilon$ dimensions

$$\begin{aligned} I_F &= -i\mu^{2\epsilon} \int \frac{d^d k}{(2\pi)^d} \frac{1}{(k^2 - m^2)(k^2 - M^2)} \\ I_M &= -i\mu^{2\epsilon} \int \frac{d^d k}{(2\pi)^d} \frac{1}{(k^2 - M^2)} \left[\frac{1}{k^2} + \frac{m^2}{k^4} + \dots \right] \\ I_{\text{EFT}} &= -i\mu^{2\epsilon} \int \frac{d^d k}{(2\pi)^d} \frac{1}{(k^2 - m^2)} \left[-\frac{1}{M^2} - \frac{k^2}{M^4} + \dots \right] \end{aligned}$$

and determine $I_F - I_M - I_{\text{EFT}}$.

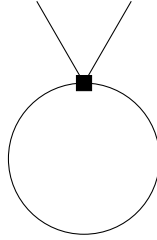
8. Show that for $SU(N)$,

$$[T^A]^\alpha_\beta [T^A]^\lambda_\sigma = \frac{1}{2} \delta_\sigma^\alpha \delta_\beta^\lambda - \frac{1}{2N} \delta_\beta^\alpha \delta_\sigma^\lambda$$

and the color Fierz identities

$$\begin{aligned} \delta_\beta^\alpha \delta_\sigma^\lambda &= \frac{1}{N} \delta_\sigma^\alpha \delta_\beta^\lambda + 2 [T^A]^\alpha_\sigma [T^A]^\lambda_\beta \\ [T^A]^\alpha_\beta [T^A]^\lambda_\sigma &= \frac{N^2 - 1}{2N^2} \delta_\sigma^\alpha \delta_\beta^\lambda - \frac{1}{N} [T^A]^\alpha_\sigma [T^A]^\lambda_\beta \end{aligned}$$

9. Compute the one-loop scalar graph with a scalar of mass m and interaction $-\lambda\phi^4/4!$.



10. Compute the anomalous dimension mixing matrix of

$$\begin{aligned} O_1 &= (\bar{b}^\alpha \gamma^\mu P_L c_\alpha) (\bar{u}^\alpha \gamma^\mu P_L d_\alpha) \\ O_2 &= (\bar{b}^\alpha \gamma^\mu P_L c_\beta) (\bar{u}^\beta \gamma^\mu P_L d_\alpha) \end{aligned}$$

Another basis often used is

$$\begin{aligned} Q_1 &= (\bar{b} \gamma^\mu P_L c) (\bar{u} \gamma^\mu P_L d) \\ Q_2 &= (\bar{b} \gamma^\mu P_L T^A c) (\bar{u} \gamma^\mu P_L T^A d) \end{aligned}$$

So let

$$\mathcal{L} = c_1 O_1 + c_2 O_2 = d_1 Q_1 + d_2 Q_2$$

and work out the transformation between c_i and d_i , and the anomalous dimension for d_i using the anomalous dimensions for c_i .