



$$B_s^0 \rightarrow \mu^+ \mu^- \gamma \text{ from } B_s^0 \rightarrow \mu^+ \mu^-$$

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CERN European Organization for Nuclear Research

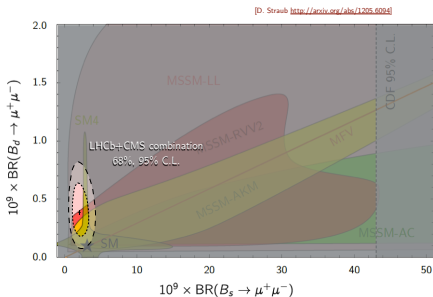
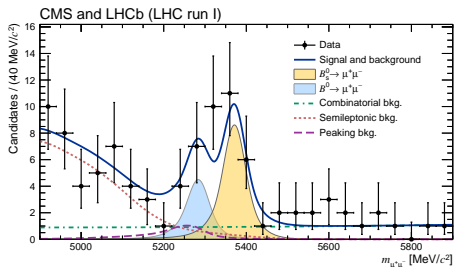
LFV/LFUV Workshop Paris

Based on: FD, D. Guadagnoli, M. Reboud - hep-ph/1610.00629



The discovery of $B_s^0 \rightarrow \mu^+ \mu^-$ decays

- $B_s^0 \rightarrow \mu^+ \mu^-$ golden channel for the search of NP in FCNC
- **Strong limits** on physics beyond the SM from combined LHC measurement
- Mainly sensitive to (pseudo)scalar contributions



Published in *Nature* 522, 68-72

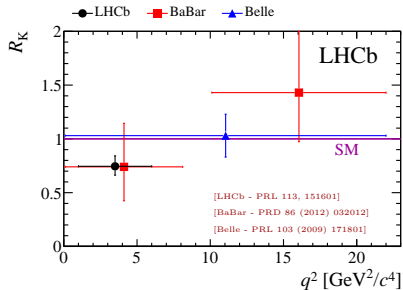
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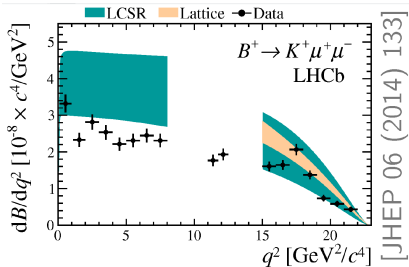
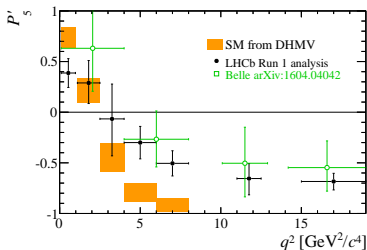
Status of flavour physics anomalies (1)

Several measurements showing discrepancies with respect to SM

- $B_d^0 \rightarrow K^* \mu^+ \mu^-$ angular distributions
- Low $b \rightarrow s \mu^+ \mu^-$ branching fractions
- Hints of **lepton non-universality** in $B \rightarrow K \ell \ell$ and $B \rightarrow D \ell \nu$ decays
- Ratio of $B^0 \rightarrow \mu^+ \mu^- / B_s^0 \rightarrow \mu^+ \mu^-$ branching fractions



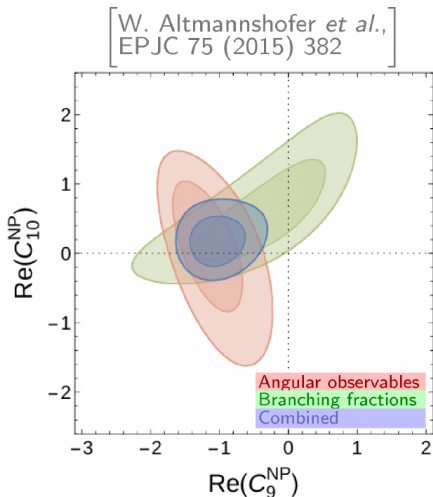
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Status of flavour physics anomalies (2)

- All points to the presence of an additional interaction
- Several global fits to Wilson coefficients agree [Buttazzo et al - 1604.03940] [Bauer et al - PRL116,141802(2016)] [Crivellin et al - PRL114,151801(2015)] [Altmannshofer et al - PRD89(2014)095033]...
- Possible underestimated contributions from $c\bar{c}$ loops [Lyon et al - 1406.0566] [Altmannshofer et al - 1503.06199] [Ciuchini et al - JHEP06(2016)116]
- Possible explanation is new high mass boson (Z')
- Can be **out of reach for direct production at LHC**
- Fundamental to be confirmed with Run II data



$$B_s^0 \rightarrow \mu^+ \mu^- \gamma \text{ from } B_s^0 \rightarrow \mu^+ \mu^-$$



$$B_s^0 \rightarrow \mu^+ \mu^- \gamma$$

- $B_s^0 \rightarrow \mu^+ \mu^- \gamma$ sensitive to \mathcal{O}_{10} but also \mathcal{O}_9 and \mathcal{O}_7^\dagger
- Increasing the number of observables sensitive to these operators is important to confirm or disprove flavour anomalies
- Sensitivity to \mathcal{O}_7 only for very low q^2 (ignored today)

$^\dagger B_s^0 \rightarrow \mu^+ \mu^- \gamma$ studied in detail in: [Aliev et al. - PRD55(1997)7059] [Geng, et al. - PRD62 (2000) 074017] [Dincer, Sehgal - PLB521(2001)714] [Descotes-Genon, Sachrajda - PLB557(2003)213] [Kruger, Melikhov - PRD67(2003) 034002] [Melikhov, Nikitin - PRD70(2004)114028]

$$B_s^0 \rightarrow \mu^+ \mu^- \gamma \text{ from } B_s^0 \rightarrow \mu^+ \mu^-$$



$B_{(s)}^0 \rightarrow \ell^+ \ell^- \gamma$ status

- $B_{(s)}^0 \rightarrow \ell^+ \ell^- \gamma$ in SM is about one order of magnitude larger than $B_{(s)}^0 \rightarrow \ell^+ \ell^-$

[Melikhov, Nikitin - PRD70(2004)114028]

m_ℓ	m_e			m_μ			m_τ		
E_{min}^γ (MeV)	20	50	80	20	50	80	20	50	80
$Br(B_d \rightarrow \ell^+ \ell^- \gamma) \times 10^{10}$ [This work]	3.95	3.95	3.95	1.34	1.32	1.31	3,39	2.37	1.87
$Br(B_s \rightarrow \ell^+ \ell^- \gamma) \times 10^9$ [This work]	24.6	24.6	24.6	18,9	18.8	18.8	11.6	8.10	6.42

- No experimental result on $B_s^0 \rightarrow \mu^+ \mu^- \gamma$
- BaBar results on the B^0 modes [BaBar collaboration - PRD77(2008)011104]

$$\mathcal{B}(B^0 \rightarrow \mu^+ \mu^- \gamma) < 1.5 \cdot 10^{-7}$$

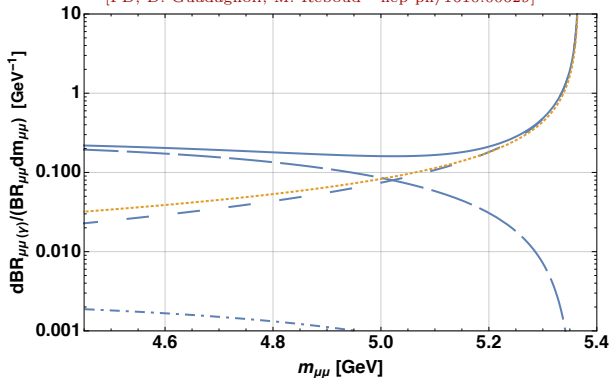
$$\mathcal{B}(B^0 \rightarrow e^+ e^- \gamma) < 1.2 \cdot 10^{-7}$$

- No results on the τ modes



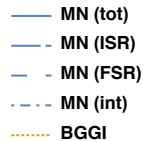
$$B_s^0 \rightarrow \mu^+ \mu^- \gamma \text{ SM}$$

[FD, D. Guadagnoli, M. Rebound - hep-ph/1610.00629]



[Melikhov, Nikitin - PRD70(2004)114028]

[Buras et al. EPJC72(2012) 2172]

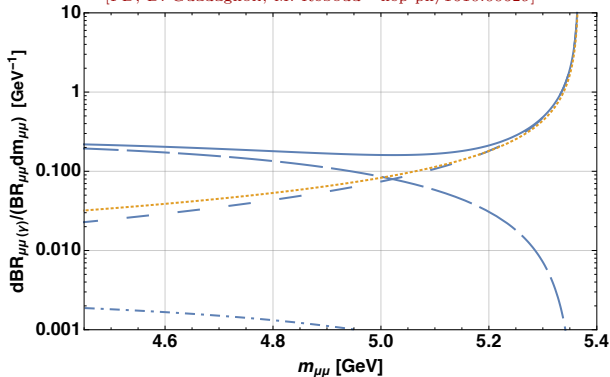


- Final state radiation (= soft bremsstrahlung)
- Initial state radiation
- Interference

$$B_s^0 \rightarrow \mu^+ \mu^- \gamma \text{ from } B_s^0 \rightarrow \mu^+ \mu^-$$


 $B_s^0 \rightarrow \mu^+ \mu^- \gamma$ SM

[FD, D. Guadagnoli, M. Reboud - hep-ph/1610.00629]



[Melikhov, Nikitin - PRD70(2004)114028]

[Buras et al. EPJC72(2012) 2172]

- MN (tot)
- - MN (ISR)
- · - MN (FSR)
- · · MN (int)
- · · BGGI

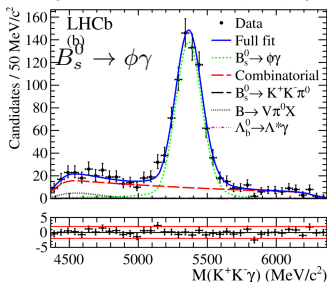
- Interference between ISR and FSR is below 1% of the total spectrum
- Negligible: $B_s^0 \rightarrow \mu^+ \mu^- + (n\gamma)$ and $B_s^0 \rightarrow \mu^+ \mu^- \gamma$ can be considered separate processes
- This holds also for beyond the SM: in particular shifts of C_9 and C_{10} of opposite sign diminish the interference

$$B_s^0 \rightarrow \mu^+ \mu^- \gamma \text{ from } B_s^0 \rightarrow \mu^+ \mu^-$$



$B_s^0 \rightarrow \mu^+ \mu^- \gamma$ direct measurement

- Direct measurement of radiative decays is experimentally challenging
 - ★ Much lower reconstruction efficiency due to the additional photon
 - ★ Energy share makes the two muons softer (hence smaller efficiency)
 - ★ Worse mass resolution: higher background
 - ★ Particularly true for hadron colliders due to the harsh environment
- Nevertheless they can be done if the branching fraction is “large” $B \rightarrow \phi \gamma$, $B \rightarrow K^* \gamma$ etc...
- $B_s^0 \rightarrow \mu^+ \mu^- \gamma$ has the additional difficulty that is also very rare
[LHCb - NuclPhysB867(2013) 1]

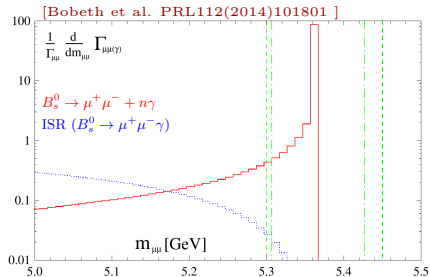


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Final State Radiation

- The radiative tail of the $B_s^0 \rightarrow \mu^+ \mu^- (+n\gamma)$ decay is experimentally indistinguishable from the $B_s^0 \rightarrow \mu^+ \mu^-$ itself
- The radiative tail is included up to a certain cut in dimuon mass correspondent to a certain photon energy
- The theoretical branching fraction is resummed to all orders in α_{em}

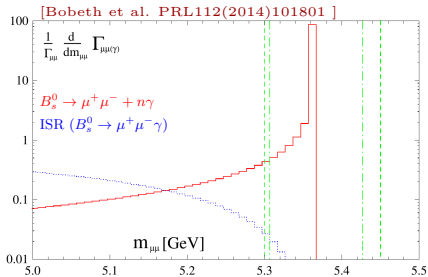


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Final State Radiation (2)

- To compare theory and experiment to search for new physics this has to be taken into account
- Experimentally it is done by using PHOTOS in MC simulations to generate $n\gamma$ from the final state
- This folds the final state contribution into the efficiency and is naturally taken into account
- Technically it is just a modification of the pdf used in the final fits



$$B_s^0 \rightarrow \mu^+ \mu^- \gamma \text{ from } B_s^0 \rightarrow \mu^+ \mu^-$$



Method description

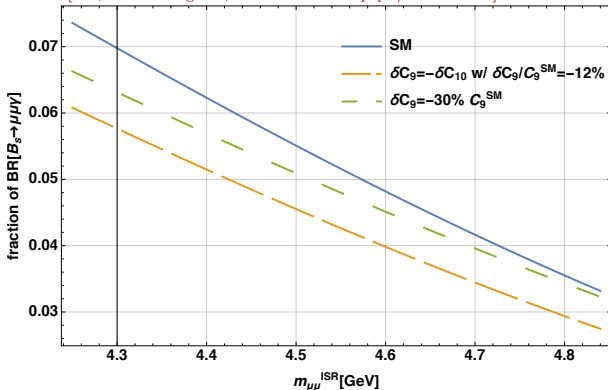
- $B_s^0 \rightarrow \mu^+ \mu^- \gamma$ is itself a contamination (background) for the $B_s^0 \rightarrow \mu^+ \mu^-$
- Contribution in the signal region is negligible
- Given the invariant mass distribution it is typically “absorbed” in the rest of the backgrounds:
 - ★ Combinatorial background (typically a free exponential)
 - ★ Partially reconstructed B decays (Constrained, but with large uncertainties)
- One can reverse the argument: if the background description has not required $B_s^0 \rightarrow \mu^+ \mu^- \gamma$ so far, this must be limited to small branching fractions
- Even more: **we can exploit the “shoulder” of the dimuon invariant mass distribution to search for $B_s^0 \rightarrow \mu^+ \mu^- \gamma$ decays**



Sensitivity

- Fraction of the $B_s^0 \rightarrow \mu^+ \mu^- \gamma$ spectrum as a function of the dimuon mass lower boundary
- About 4.8% for $m_{\mu\mu} > 4.6$ GeV
- Significant changes with possible NP values of Wilson coefficients

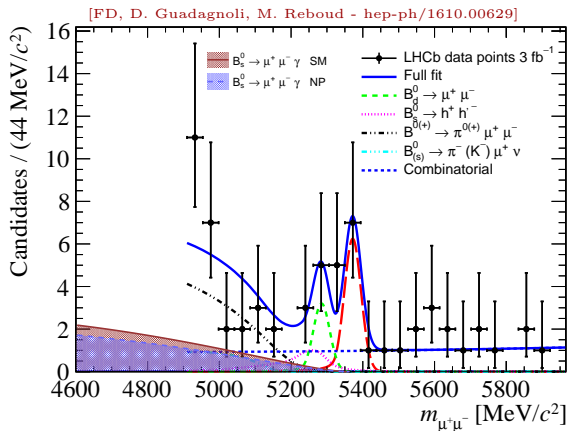
[FD, D. Guadagnoli, M. Rebold - hep-ph/1610.00629]



$$B_s^0 \rightarrow \mu^+ \mu^- \gamma \text{ from } B_s^0 \rightarrow \mu^+ \mu^-$$



Comparing $B_s^0 \rightarrow \mu^+ \mu^- \gamma$ to a $B_s^0 \rightarrow \mu^+ \mu^-$ search data



Data from LHCb Run I measurement

[PRL111(2013)101805]

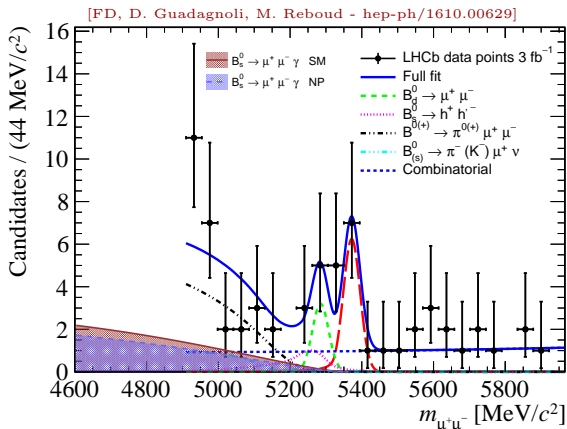
Assumptions

- Used single event sensitivity to translated branching fraction into yield
- Assumed same efficiency as $B_s^0 \rightarrow \mu^+ \mu^-$: true by definition at the end-point, but will fade down far from it

$$B_s^0 \rightarrow \mu^+ \mu^- \gamma \text{ from } B_s^0 \rightarrow \mu^+ \mu^-$$



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Data from LHCb Run I measurement

[PRL111(2013)101805]

- LHCb just an example, procedure can be applied to all experiments
- Important to constrain the partially reconstructed backgrounds even more



Further remarks

- Method applicable to any $M \rightarrow \ell\ell\gamma$ decay
 - * more sensitive for muons and electrons
 - * theoretical uncertainties might dominate (e.g. $D^0 \rightarrow \mu^+\mu^-\gamma$)
- Method applicable to any $M \rightarrow \mu^+\mu^-X$ but..
 - * the higher M_X the less some (experimental) assumptions are valid
 - * Backgrounds might be too large / unconstrained

$$B_s^0 \rightarrow \mu^+\mu^-\gamma \text{ from } B_s^0 \rightarrow \mu^+\mu^-$$



Experimental remarks (a biased view)

While the method is generic I have a biased view...

- I have no doubt this can be applied to LHCb but it requires effort:
 - * Detailed estimate of $B_s^0 \rightarrow \mu^+ \mu^- \gamma$ efficiency vs $m_{\mu\mu}$ and correspondent integrated efficiency for different models
 - * Precise control of backgrounds from $B \rightarrow X \mu\mu$ and others with particle mis-identification which are now too low in mass for $B_s^0 \rightarrow \mu^+ \mu^-$ searches
- ATLAS and CMS are in a different position
 - * Have on their side the much larger statistics
 - * Worse dimuon mass resolution (from 25 MeV of LHCb to 30 up to 100 in the GPDs depending on category) can increase the cross-feed with $B_s^0 \rightarrow \mu^+ \mu^-$
 - * Difficult to control other B backgrounds without correctly reconstructed proxies, but they have great MC productions
 - * Lack of personpower (?): maybe not, we have seen that $B_s^0 \rightarrow \mu^+ \mu^-$ attracted a lot of resources
- Belle II
 - * As soon as data taking starts we will have a new player in town, but mostly on $B_d - B^+$
 - * Low background
 - * Full reconstruction might also be feasible ($B_d^0 \rightarrow \mu^+ \mu^- \gamma$)

$$B_s^0 \rightarrow \mu^+ \mu^- \gamma \text{ from } B_s^0 \rightarrow \mu^+ \mu^-$$



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Conclusions

- $B_s^0 \rightarrow \mu^+ \mu^- \gamma$ provides complementary sensitivity to \mathcal{C}_9 and \mathcal{C}_{10}
- Direct measurement might be experimentally far
- Proposed alternative method for search that can arguably provide first measurements of $B_s^0 \rightarrow \mu^+ \mu^- \gamma$
- Run II LHC measurements might already give interesting limits or maybe more...

$$B_s^0 \rightarrow \mu^+ \mu^- \gamma \text{ from } B_s^0 \rightarrow \mu^+ \mu^-$$



Backup

$$B_s^0 \rightarrow \mu^+ \mu^- \gamma \text{ from } B_s^0 \rightarrow \mu^+ \mu^-$$