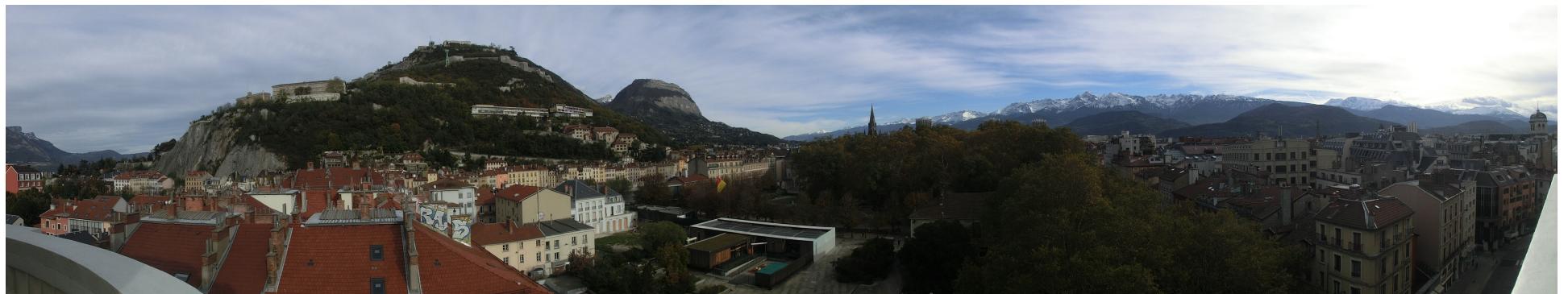


La bosse diphoton au LHC : nouvelles de la frontière en énergie (*energy frontier*)

Jan Stark

Laboratoire de Physique Subatomique et de Cosmologie
Grenoble, France



Journées LCG-France, Grenoble, 21 juin 2016

La bosse diphoton au LHC : nouvelles du front en énergie

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Choix

Pendant la préparation de cette présentation,
j'étais obligé de faire des choix :

~~Faire un catalogue de résultats du LHC~~

Choisir un seul résultat,
le présenter dans son contexte

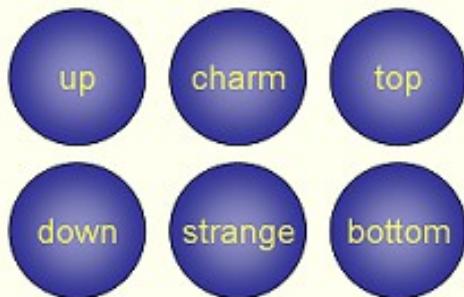


Inconvénient de ce choix : je ne pourrai pas rendre justice à toute la richesse des résultats obtenus au LHC. Mes excuses.

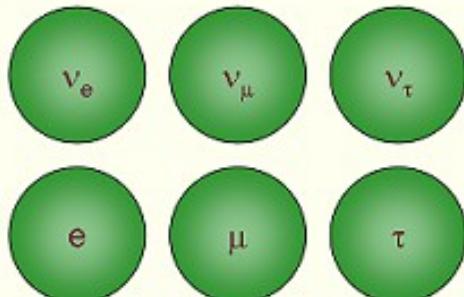
Fundamental constituents of matter

Particle content of the Standard Model

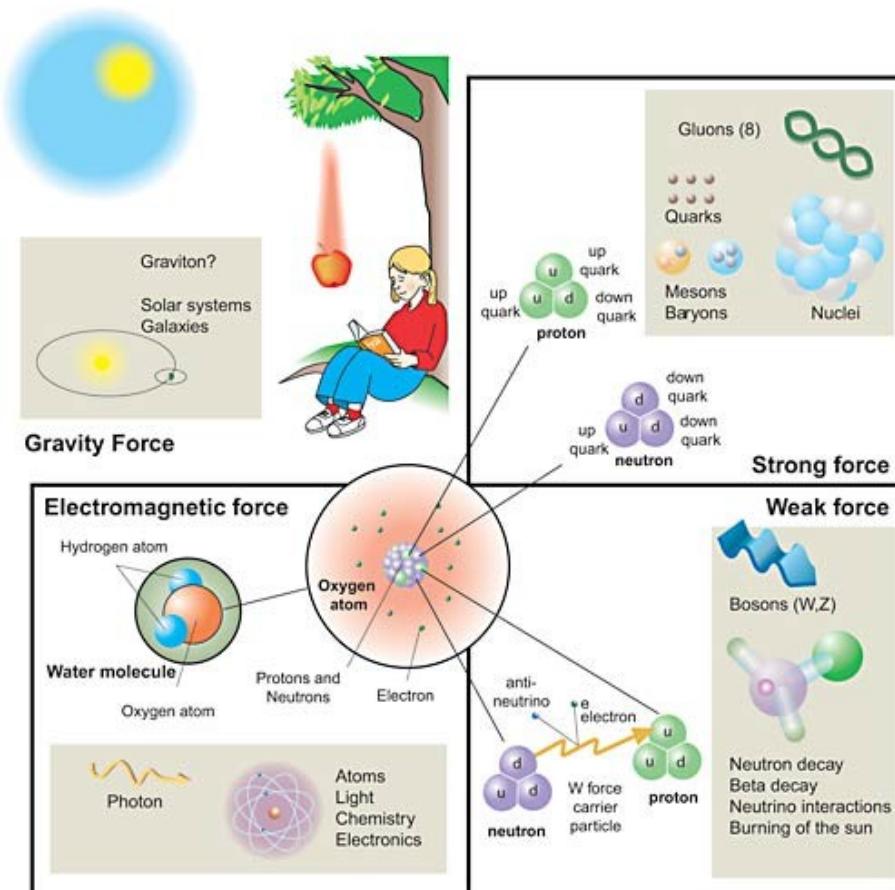
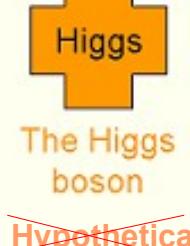
Quarks:



Leptons:



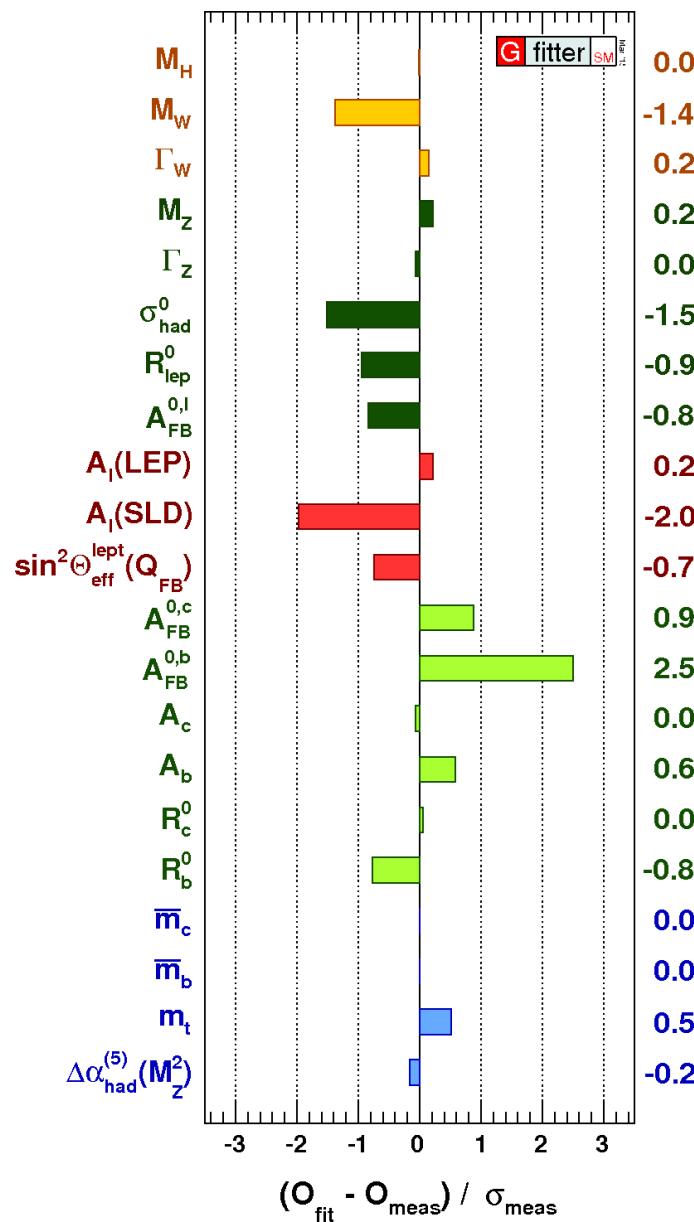
Force carriers



A self-consistent theory !

After the discovery of a Higgs-like boson, the standard model (SM) is a complete (but for an axion) and **self-consistent** renormalisable **theory**.

So far, **SM in good agreement with data**.



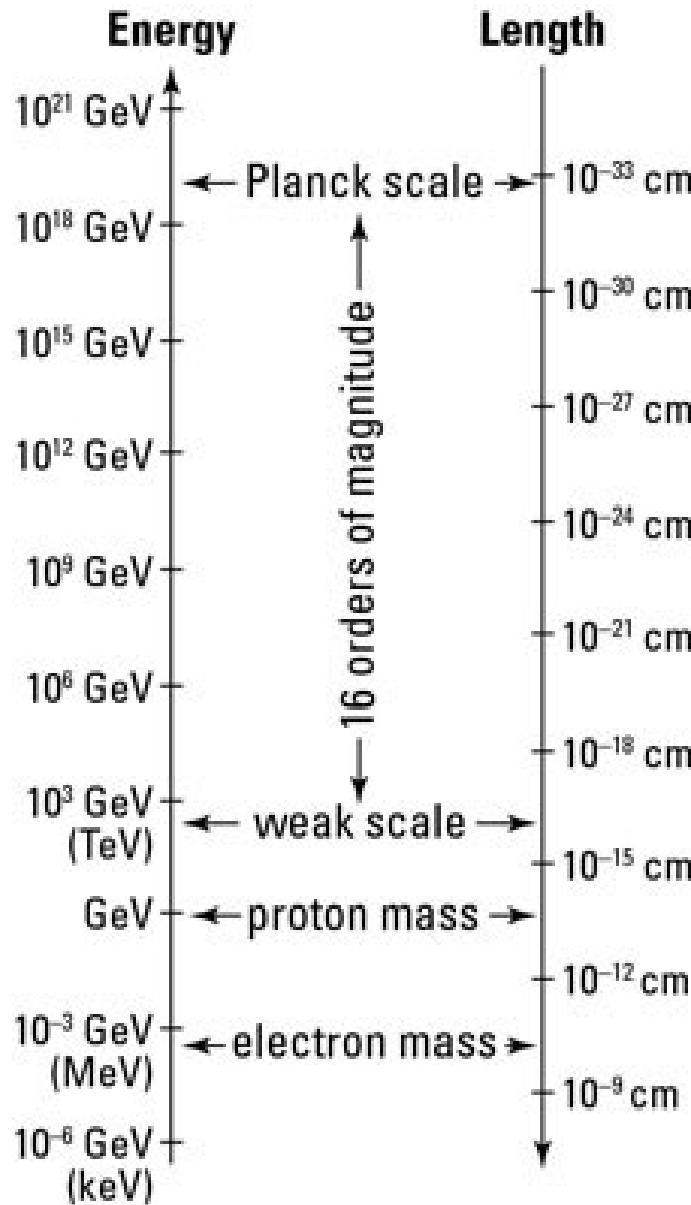
Questions ...

But a lot of questions still need answers:

- How to accommodate gravity ?
- Hierarchy problem: $m_{EW}/M_{Pl} \sim 10^{-16}$
- What is dark matter ?
- Matter-antimatter asymmetry
- Origin of generations ?
- ...

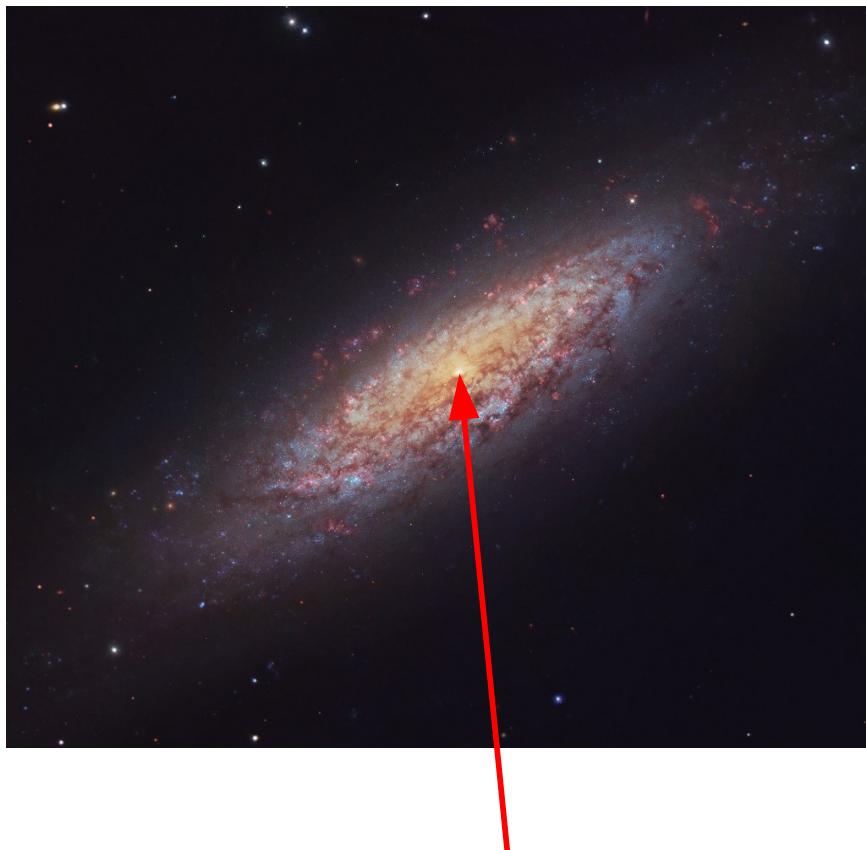
Physics beyond the SM is well motivated.

The hierarchy problem



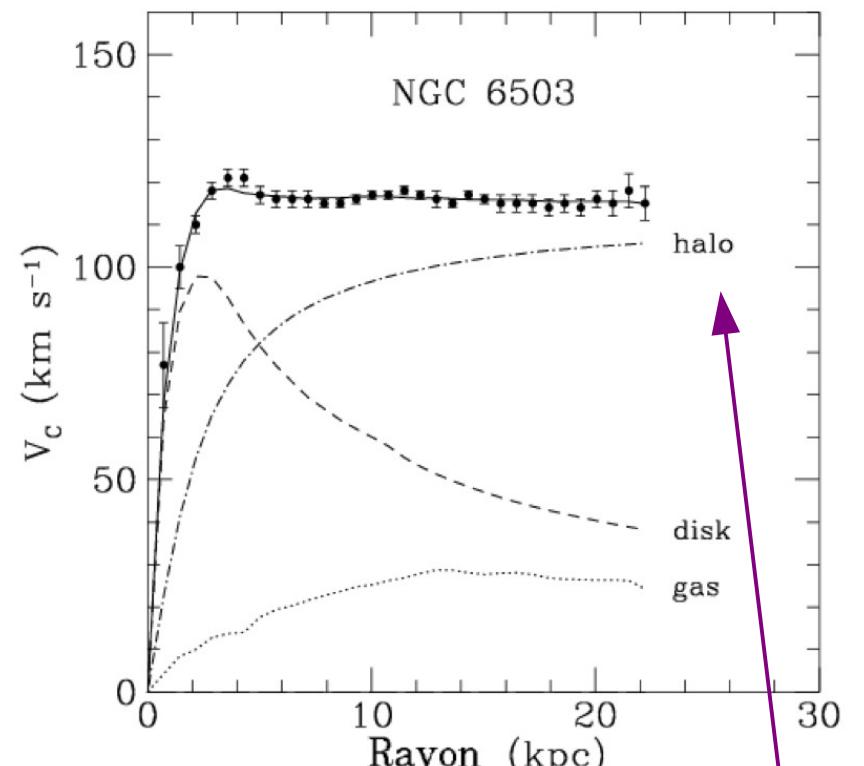
Dark matter

La galaxie NGC 6503



centre de la galaxie

Vitesse de rotation des étoiles dans cette galaxie, en fonction de la distance de l'étoile par rapport au centre de la galaxie



Il est fait de quoi, ce halo ?

Les objectifs de l'IN2P3

Rechercher ok

Le CNRS | Annuaires | Mots-Clefs CNRS | Autres sites

Institut national de physique nucléaire et de physique des particules
Centre national de la recherche scientifique

Présentation de l'Institut

Structures de recherche

Conseil scientifique

Infos aux laboratoires

Vie de la recherche

Carrières et emplois

Physique subatomique pour tous

English version

Rechercher :
• une expérience
• une personne
• une photo

Rechercher :
une information sur le site de l'IN2P3

ok

Physique des particules

Le Modèle standard

Introduction

Physique des particules

Physique nucléaire et hadronique

Astroparticules et neutrinos

Aval du cycle électronucléaire et énergie nucléaire

Recherches interdisciplinaires

Recherche et développement en accélérateurs

Grille de calcul

Le domaine de la physique des particules rassemble les recherches expérimentales visant à faire progresser la connaissance des quarks (les composants des protons et des neutrons), des leptons (tel l'électron gravitant autour du noyau atomique) et des bosons responsables de leurs interactions mutuelles. Ces particules sont considérées actuellement comme les constituants les plus élémentaires de la matière.

Le Modèle standard

Les avances parallèles et couplées des observations expérimentales et des progrès théoriques enregistrés depuis les années 60 ont permis l'élaboration de ce qui est appelé le Modèle standard (voir l'introduction). Bien que ce modèle donne une description très satisfaisante des phénomènes observés dans les expériences, c'est toutefois une théorie incomplète :

- elle ne peut expliquer le pourquoi de l'existence de trois familles ;
- elle ne permet pas de prédire les valeurs observées des masses des particules ;
- elle ne rend compte que de trois de ces quatre interactions fondamentales à l'œuvre dans l'Univers : la gravitation en est en effet exclue ; en revanche, elle a déjà permis d'unifier les interactions électromagnétique et faible en une interaction unique dite électrofaible.

Même s'il n'a jamais été démenti par l'expérience, on sait donc que le Modèle standard n'est pas la théorie ultime. Plusieurs grandes questions se posent, auxquelles il ne peut répondre dans sa forme actuelle :

- Pourquoi les particules ont-elles une masse et pourquoi ont-elles des masses si différentes ?
- Comment aller vers une plus grande unification des interactions qui inclurait l'interaction électrofaible et l'interaction forte, et même la gravité ?

Le boson de Higgs, particule prédicta par le Modèle standard et détectée pour la première fois en 2012, devrait expliquer l'origine des masses et permettre de comprendre leurs différences.

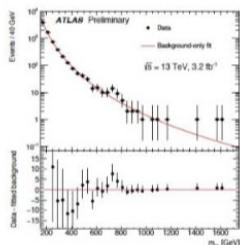
Les physiciens pensent en effet que les quatre interactions fondamentales ne seraient que des aspects différents d'une

Questions ...

Theorists; state of mind

Poor theorists:

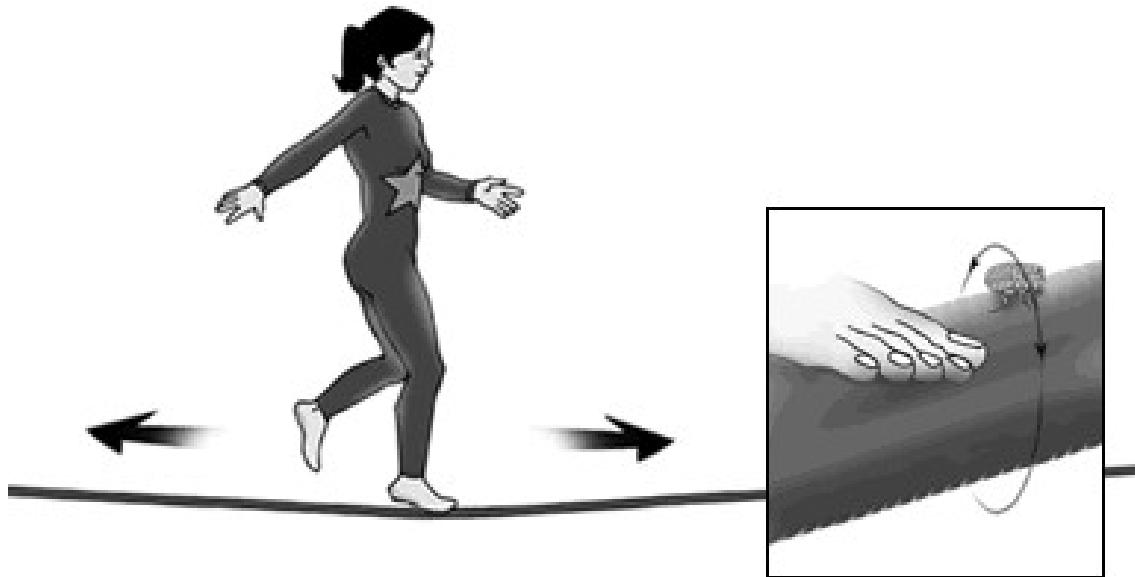
**Waiting for new physics for 30 years,
and recently started to get desperate..
and something interesting appears.**



Abdelhak Djouadi (LPT Orsay) Moriond QCD conference March 2016

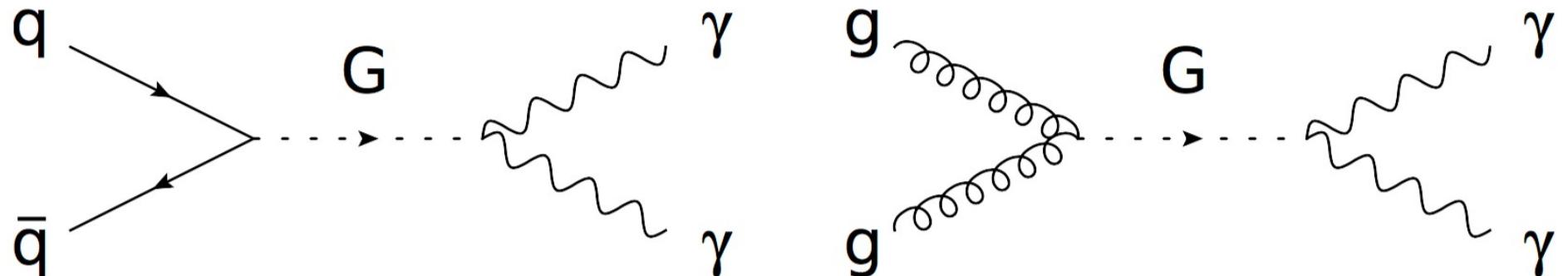
Extra spatial dimensions

- In the Randall-Sundrum model gravity propagates in a warped extra dimension with two fixed points
- The Standard Model fields are constrained to one brane
- The gravity wave function is concentrated near the other brane, falling off exponentially across the extra dimension



Graviton excitations

The model predicts a tower of Kaluza-Klein graviton states with TeV scale masses



Searches for resonances

Fully reconstructed resonances: simplest way to discover new particles

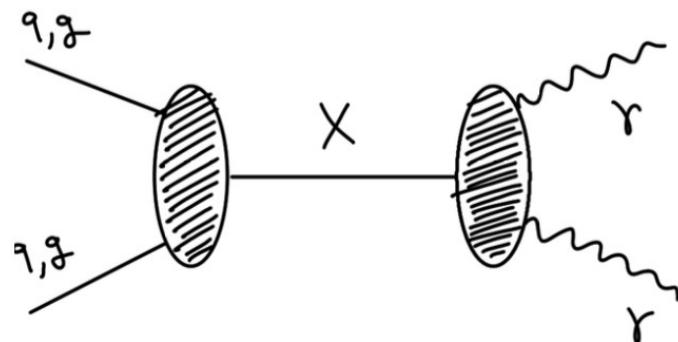
Statistically significant peak over a smooth background

- ✓ experimentally robust
 - ✓ small systematics
 - ✓ difficult for unknown backgrounds to mimic
- => ***simple yet striking signature!***

The most important search method when new energies are explored

- ✓ particularly relevant at LHC Run2 startup
- ✓ model independent probe to new physics

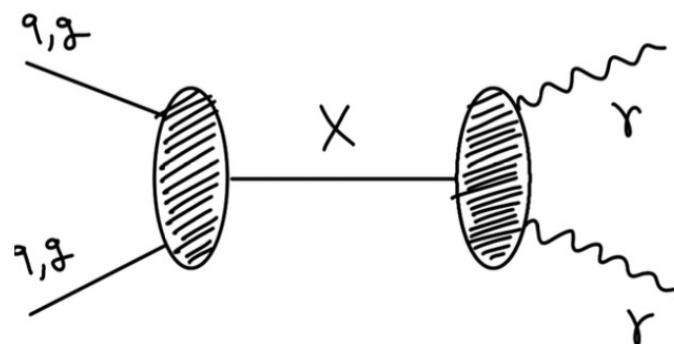
Diphoton bump search



Clean final state at hadron colliders

- 1) Define the event selection: 2 isolated photons
 - ✓ must be loose and model-independent

Diphoton bump search



Clean final state at hadron colliders

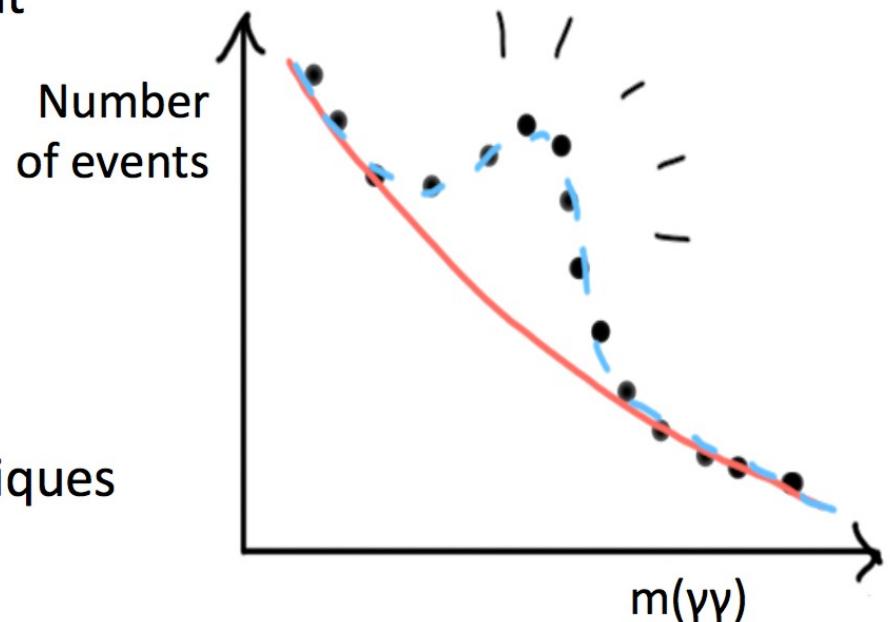
- 1) Define the event selection: 2 isolated photons

- ✓ must be loose and model-independent

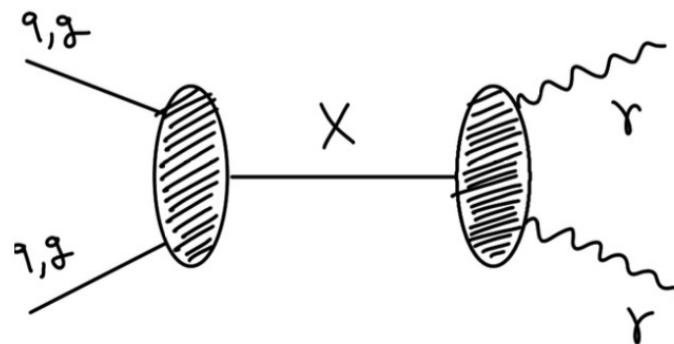
- 2) Reconstruct the $\gamma\gamma$ invariant mass

$$M = \sqrt{2E_1 E_2 (1 - \cos\theta)}$$

- ✓ photon reconstruction
- ✓ detector resolution and scale
- ✓ dedicated vertex identification techniques



Diphoton bump search



Clean final state at hadron colliders

1) Define the event selection: 2 isolated photons

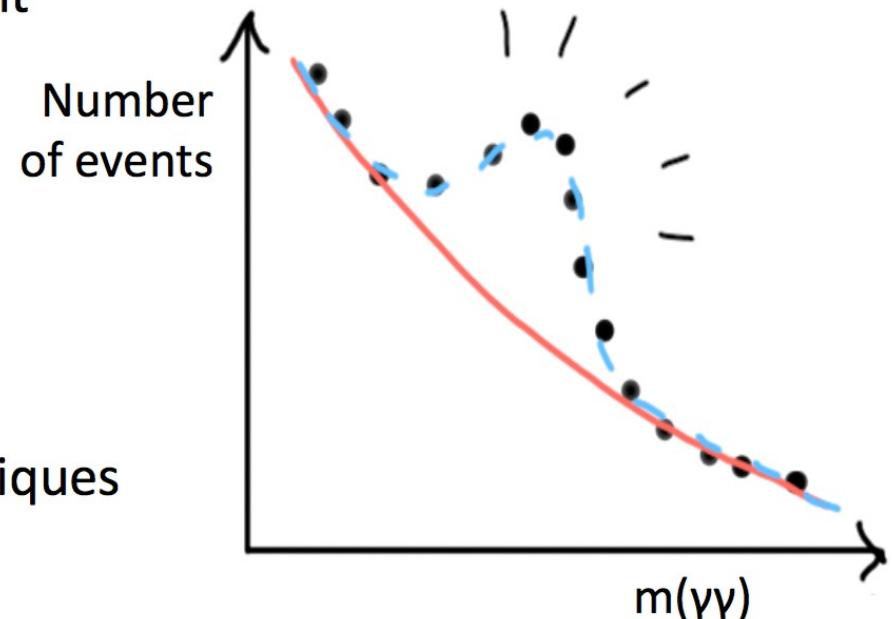
- ✓ must be loose and model-independent

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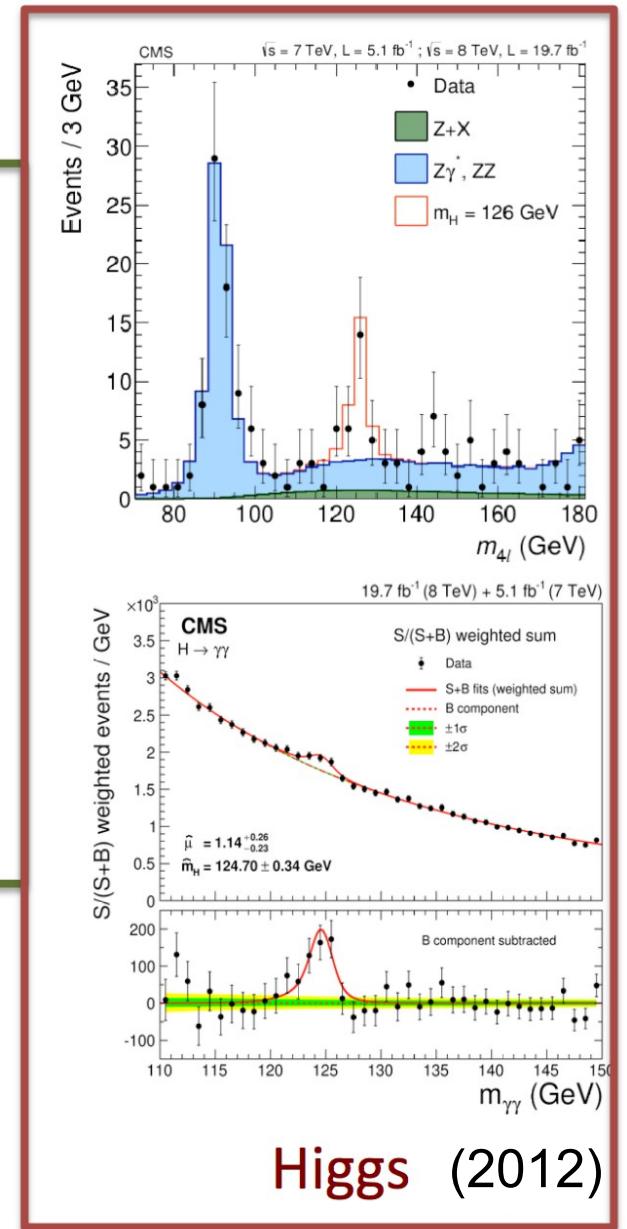
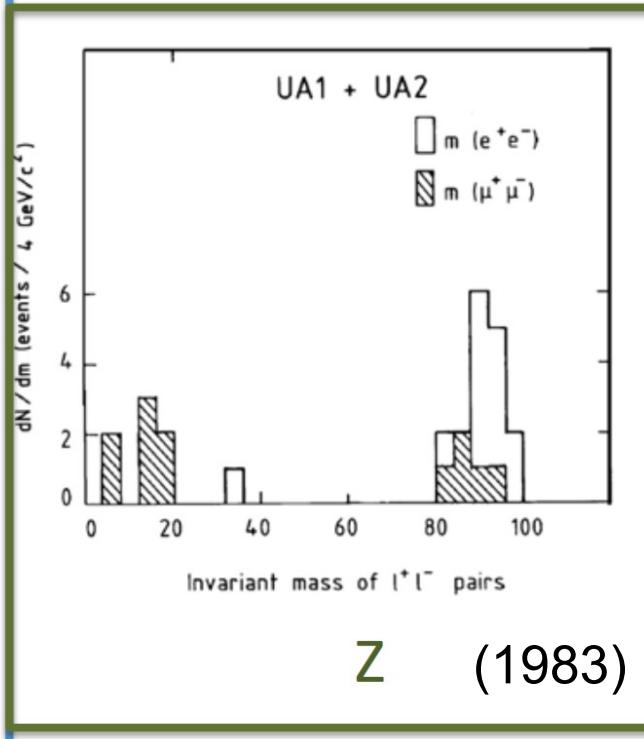
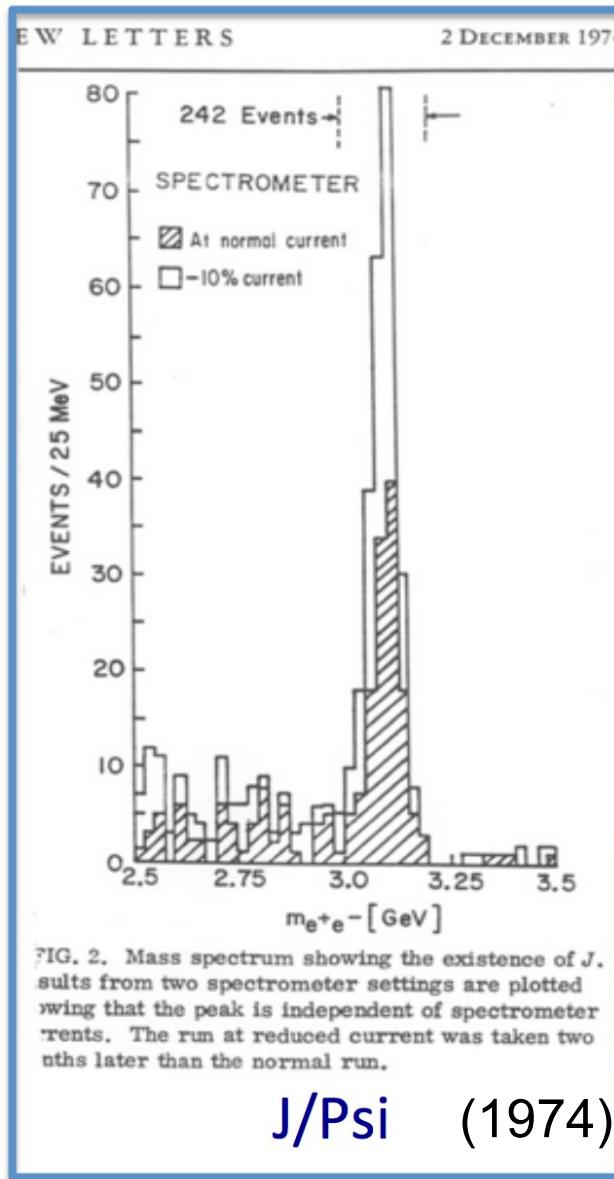
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- ✓ photon reconstruction
- ✓ detector resolution and scale
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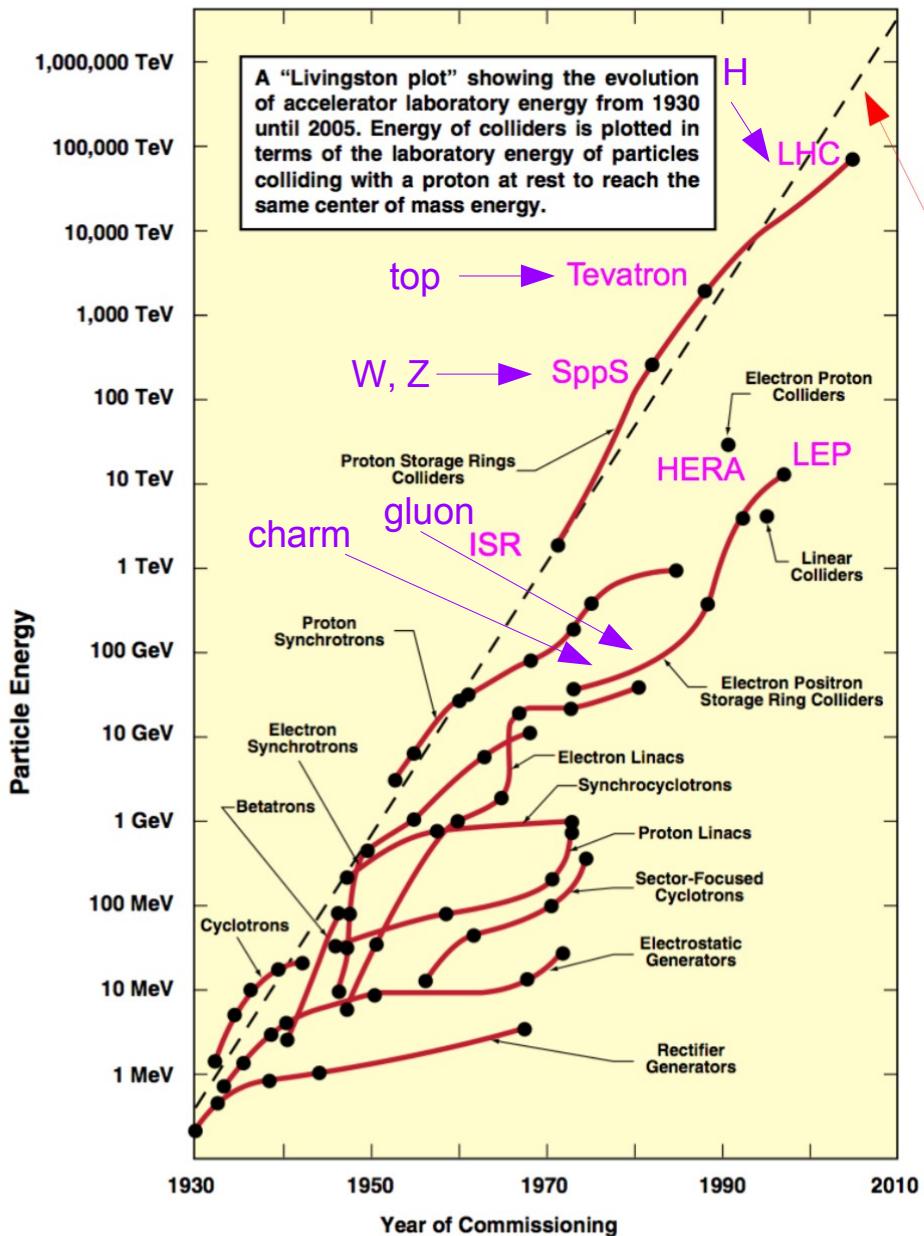
3) Signal extraction



Resonances in past discoveries



Livingston plot



Plot fait pour la première fois par M. Stanley Livingston en 1954.

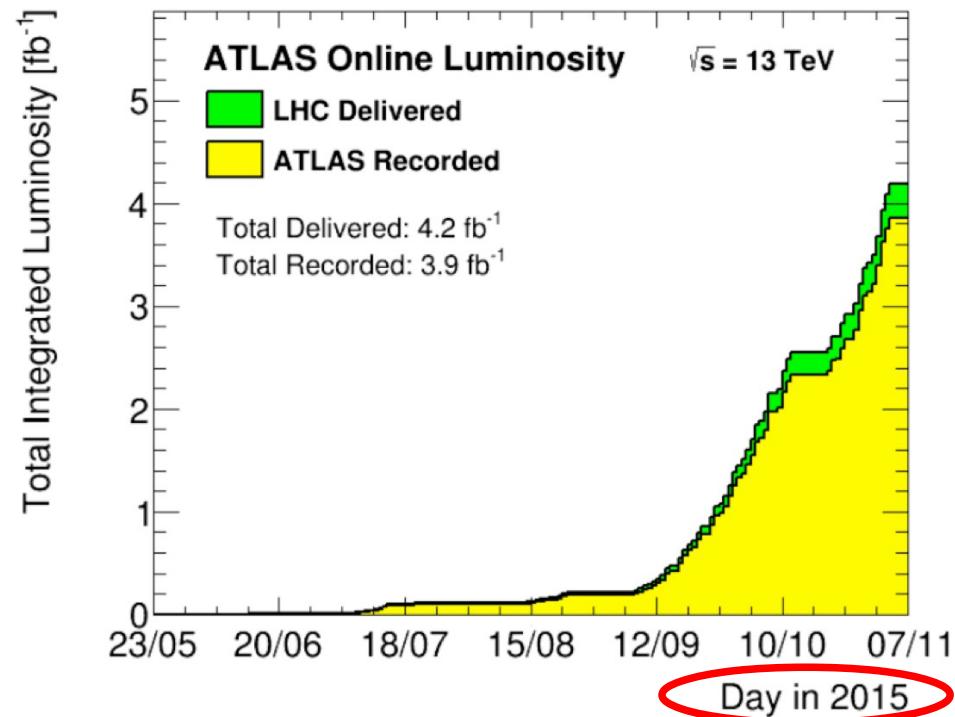
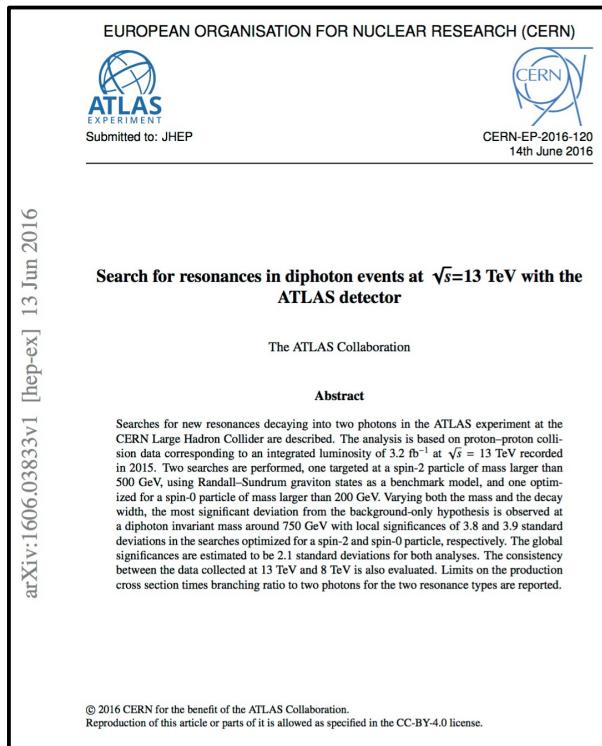
“Droite de Livingston” : loi de croissance exponentielle.

Pour comparaison : la “loi de Moore” (puissance de calcul des circuits intégrés) a été énoncée pour la première fois en 1965.

Découvertes récentes

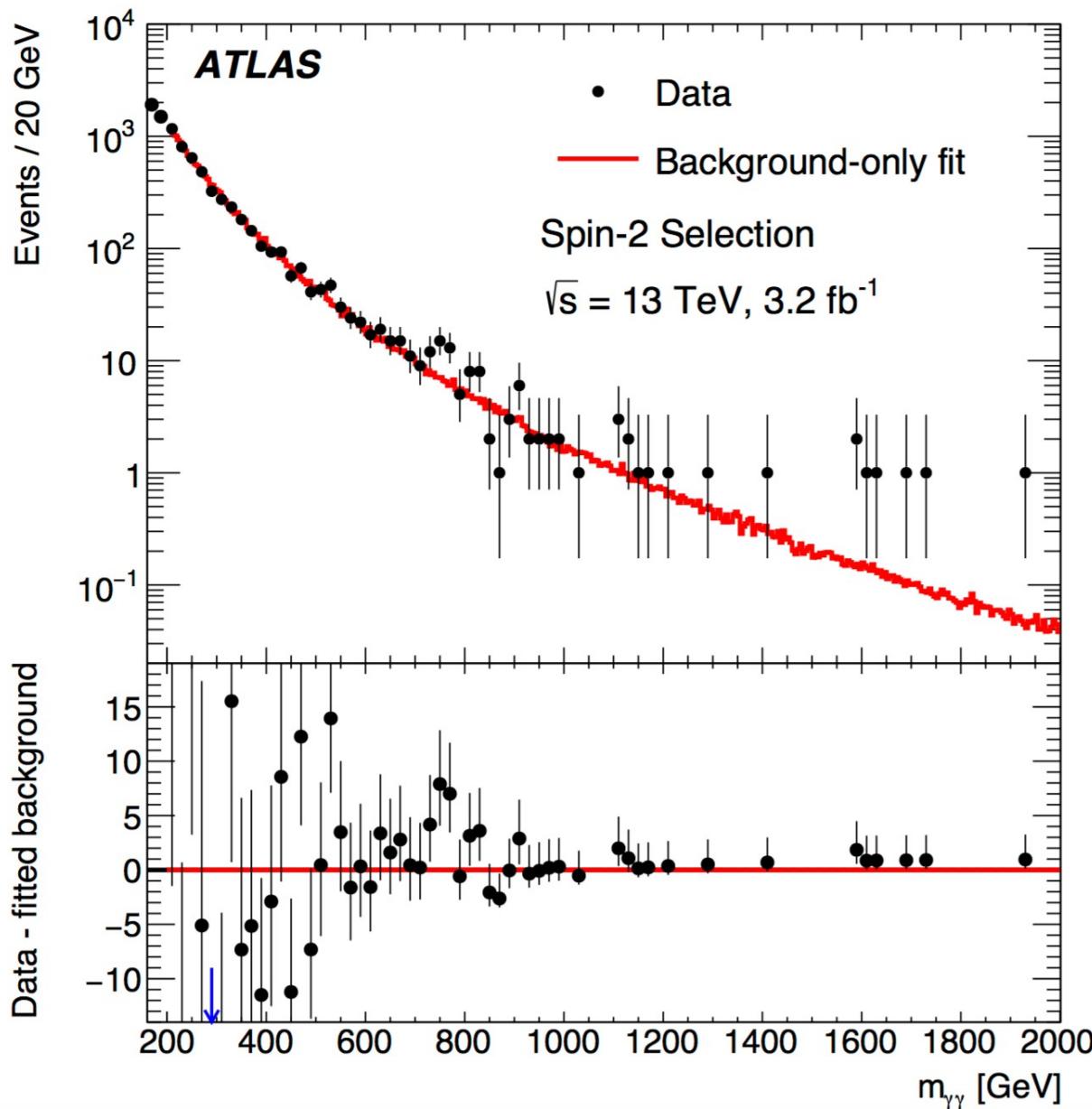
ATLAS: 2015 data, diphoton results

- 3.9 fb^{-1} of data collected with 25 ns bunch crossing
 - $\approx 100 \text{ pb}^{-1}$ of data collected with 50 ns bunch crossing
 - $\approx 680 \mu\text{b}^{-1}$ of heavy ion data collected
- Data taking efficiency: 92%
- Data quality efficiency: 93% (run-1)
94%)

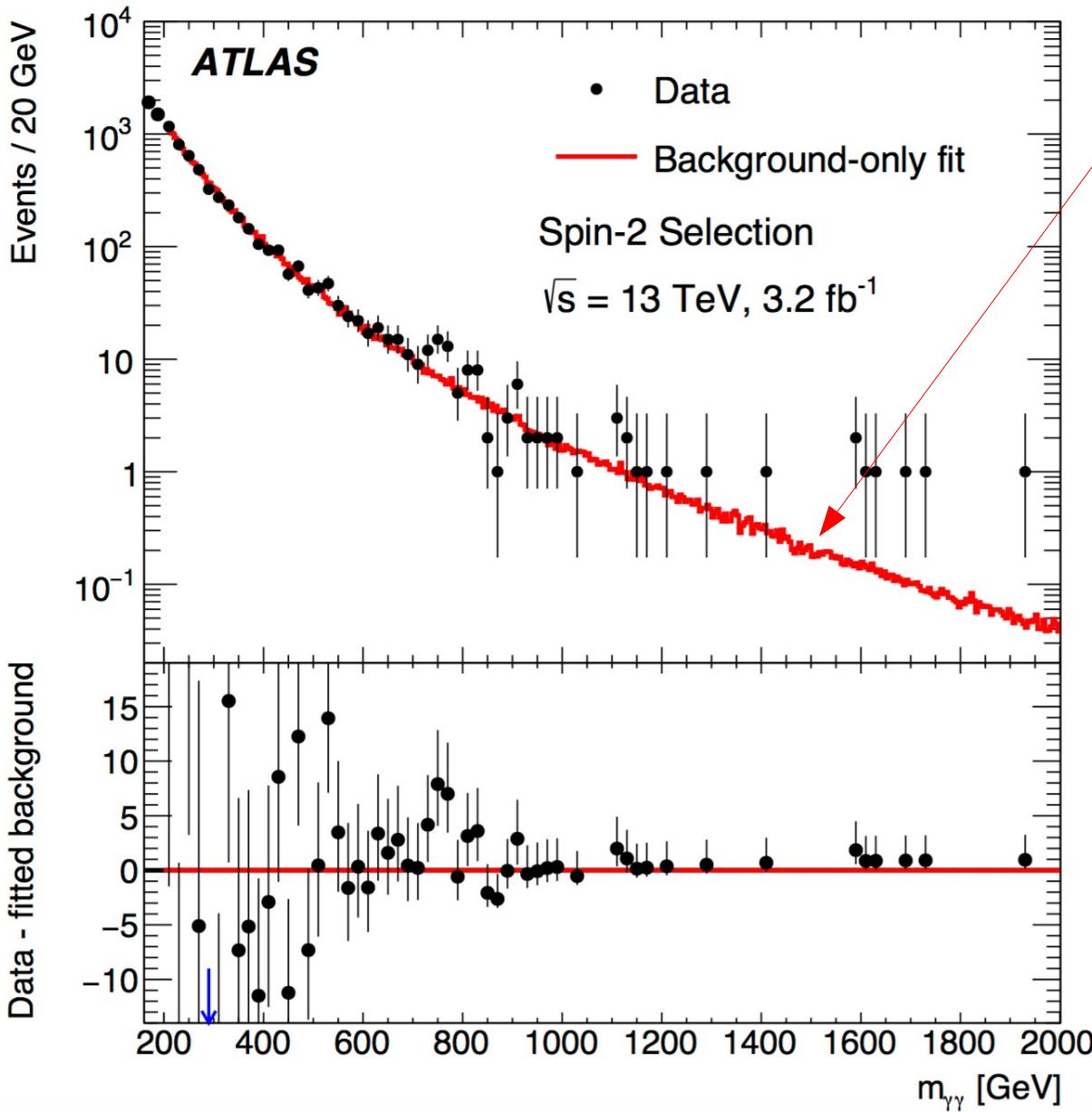


- End of data taking: beginning Nov. 2015
- First public presentation of results: Dec. 15th, 2015
- Major update: Moriond conference, March 2016
- Publication submitted to journal: June 14th, 2016

Diphoton mass spectrum

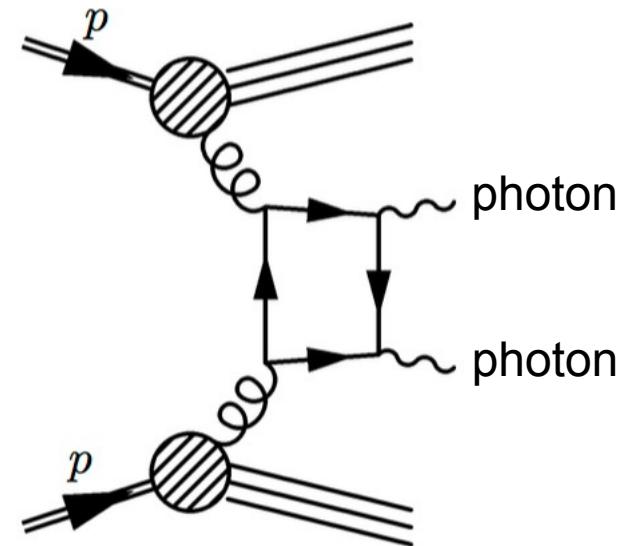


An aside: Diphox



Standard Model background prediction based on precise first-principles calculation in the Diphox generator.

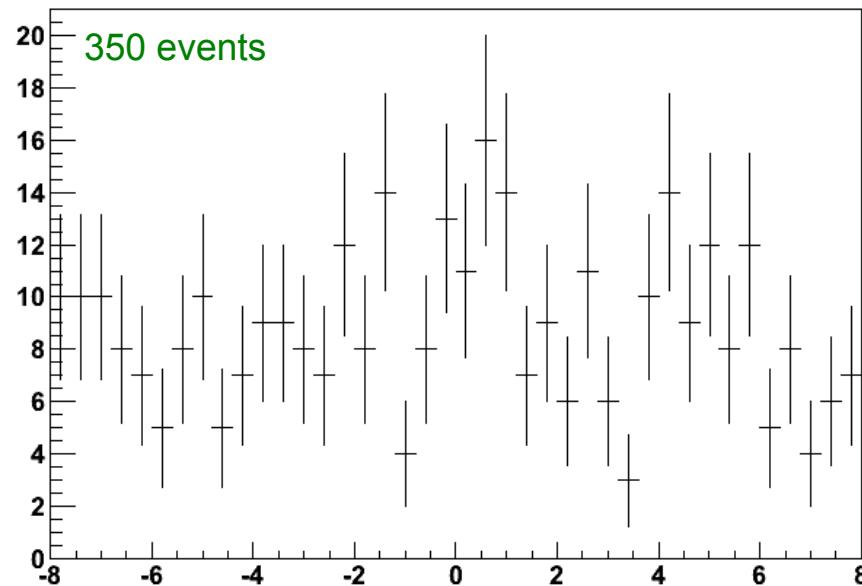
Diphox is computing-intensive; we run it on CIMENT in Grenoble.



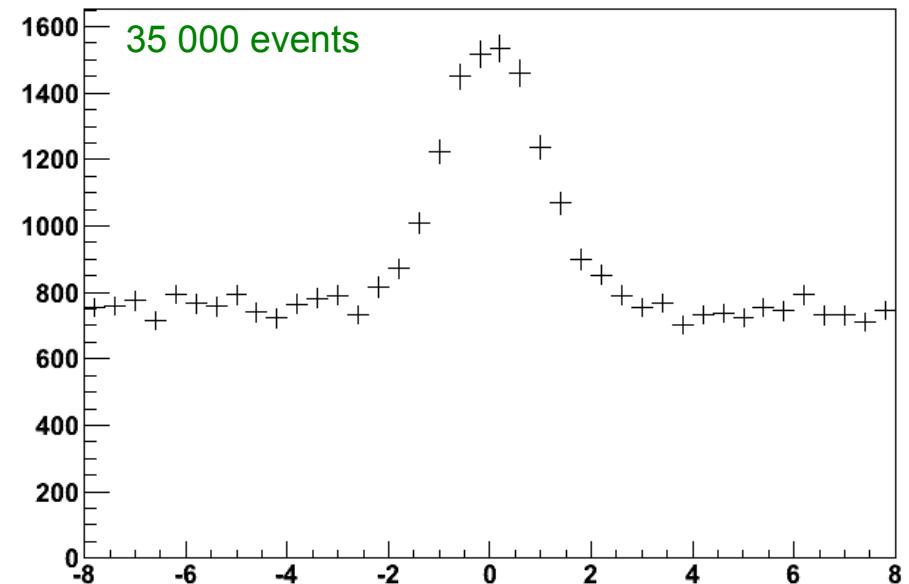
Importance of statistics

These two graphs show the same distribution, on the left with little data, and on the right with 100 times more data:

histogram



histogram



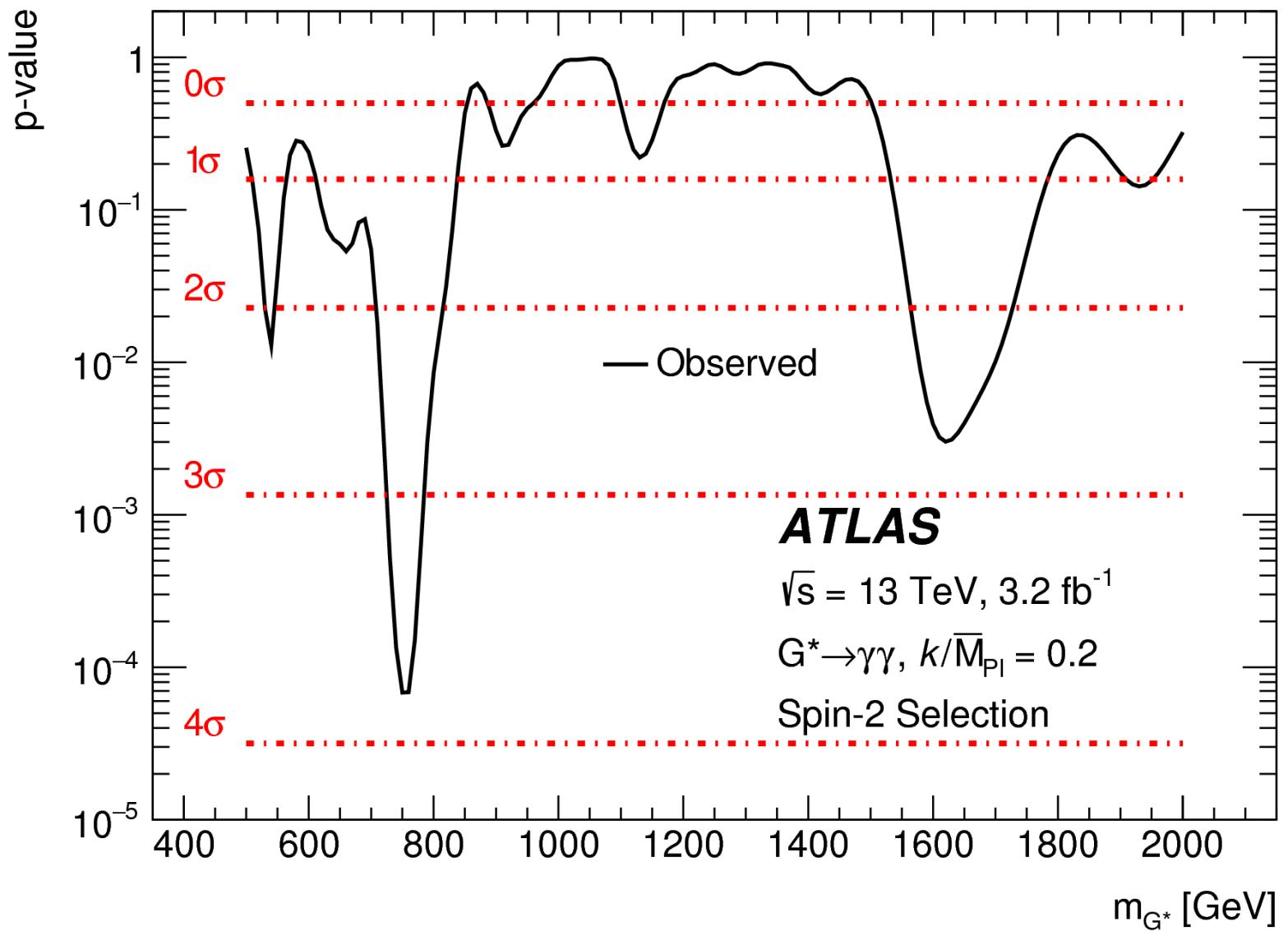
With a lot of data we can clearly see a structure :
a flat spectrum (e.g. due to background),
plus a peak (e.g. due to a new particle).

With less data the situation is much less clear :

- is this simply a flat spectrum ?
- or is there a peak somewhere ??

p-value

At 750 GeV:
3.8 σ local
2.1 σ global
for $\Gamma \approx 50$ GeV



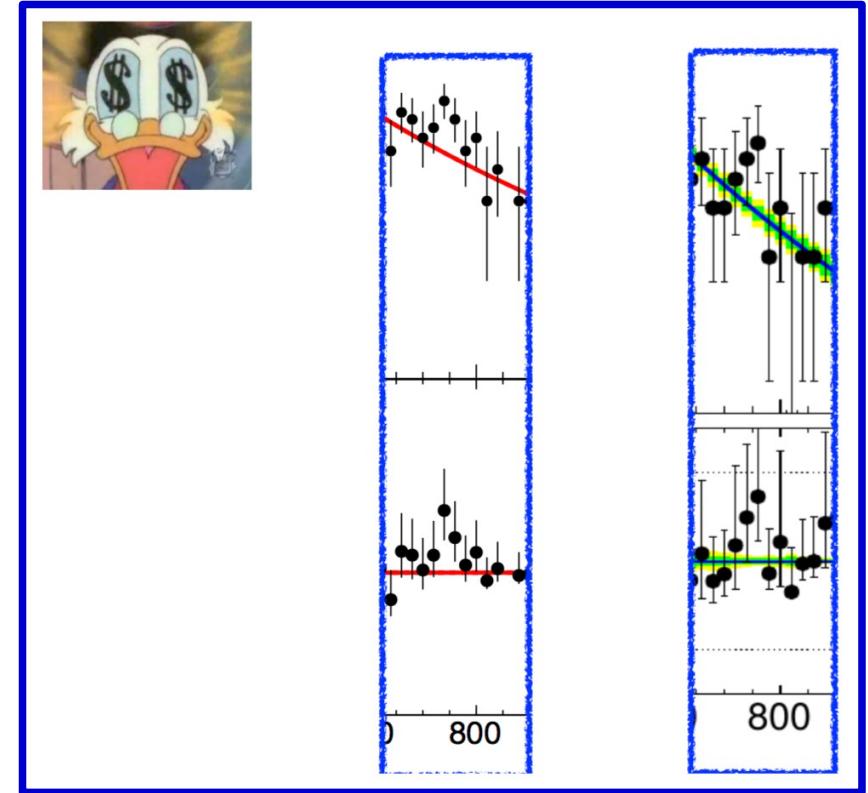
Excitement

Abdelhak Djouadi (LPT Orsay)
Moriond QCD conference
March 2016



If true then the future is bright!
a new continent is ahead and
needs decades of exploration...

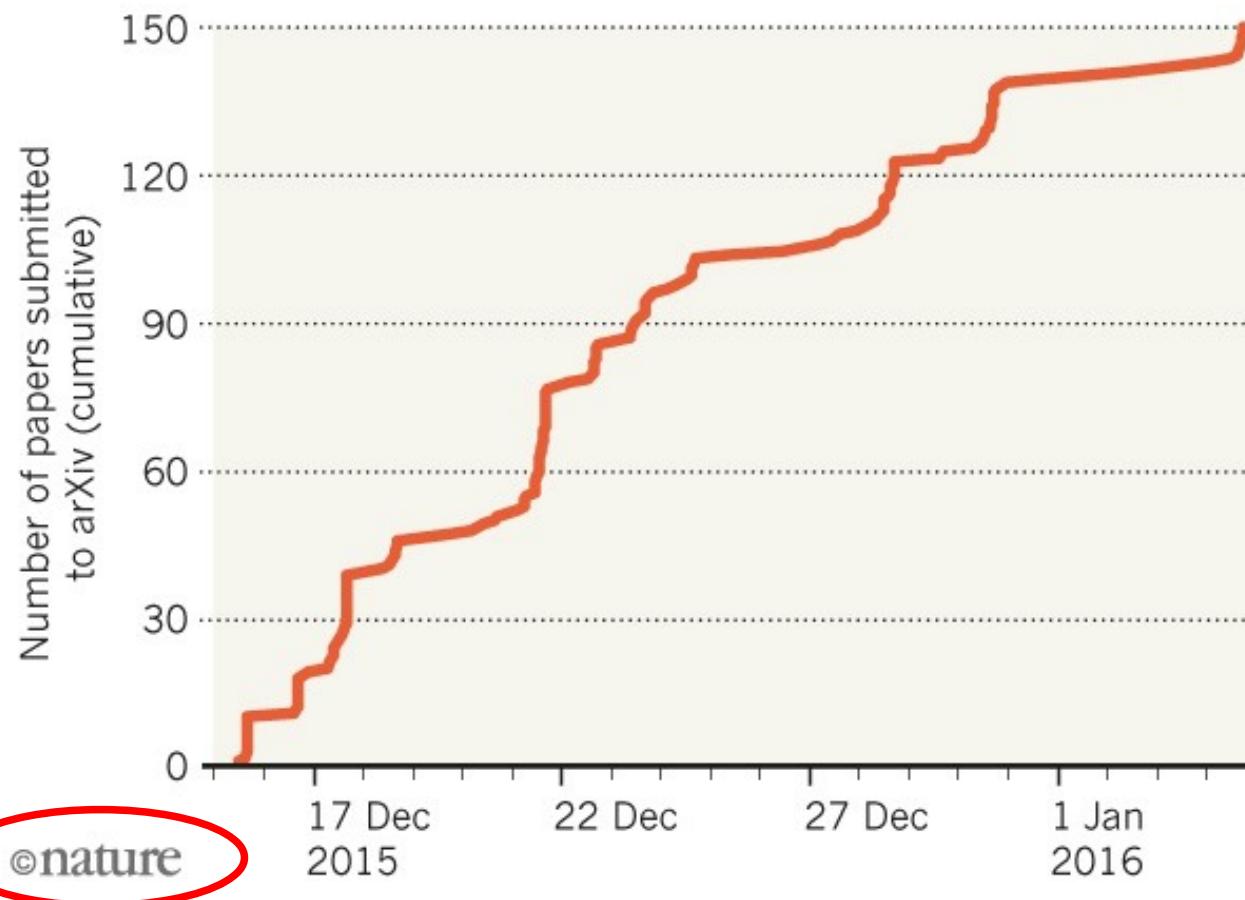
Alessandro Strumia
Moriond conferences
March 2016



Excitement

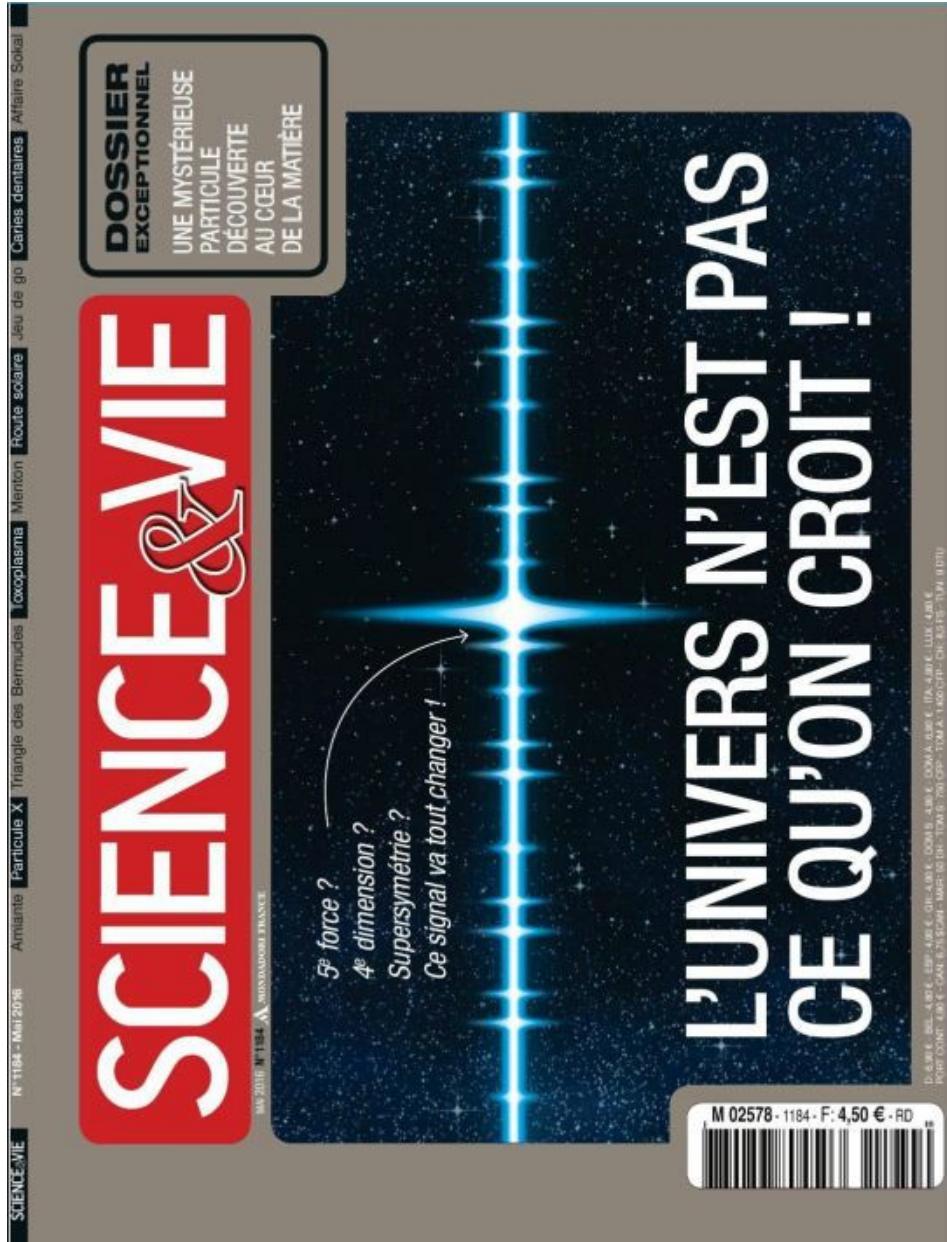
HINT OF NEW BOSON SPARKS FLOOD OF PAPERS

In just 21 days, physicists have posted 150 papers on the arXiv preprint server about tantalizing results at the Large Hadron Collider.

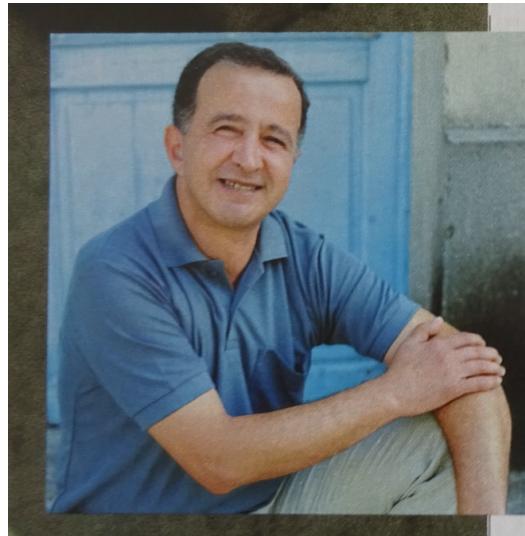


More than 456 papers
as of June 16th.

Excitement



Mai 2016



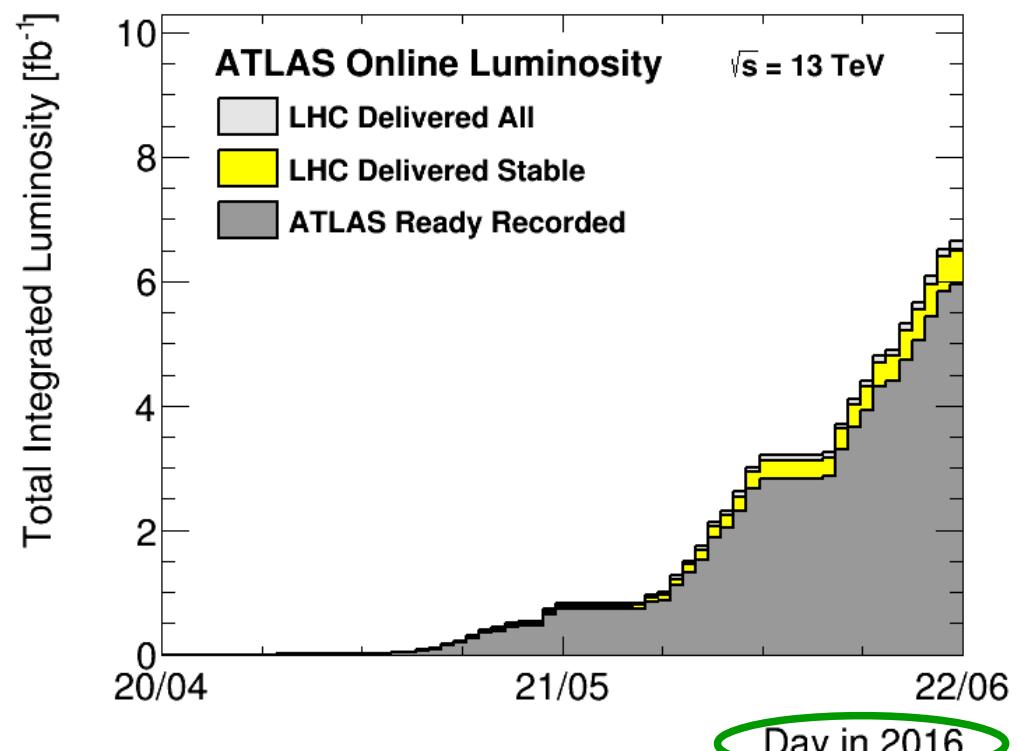
ABDELHAK DJOUADI
Physicien,
Laboratoire
de physique
théorique d'Orsay
“On sort des sentiers battus... Cela ouvre des perspectives très excitantes !”

What can we do now ?



- while no detector or background mismodeling effects were found to explain the excess, this is not statistically conclusive, and compatible with an upward fluctuation of the background
- we will have a clearer picture by next summer, when more data will be available

What can we do now ?



COLLECT

Conclusions

Standard model (SM) of particle physics (including the Higgs mechanism):
consistent renormalisable theory that “works” all the way up to the Planck scale.

But is it really a complete description of Nature
up to these extreme energies (Planck scale) ?

Various “aesthetic” issues (like e.g. the hierarchy problem) lead us to suspect that
the answer is “no” and that the SM is a low-energy approximation of a more
complete theory.

The direct search for physics beyond the SM is on at the LHC.

First data in a new energy regime (13 TeV) collected in 2015.

An intriguing excess of events in the diphoton final state has lead to some excitement
recently. But it is not statistically significant yet.

Much more data expected in the coming months (and years ...); they will provide
answers to some of the pressing questions in particle physics.

Acknowledgements

In the preparation of this set of slides, I have made ample use of two seminars that have been given by

C. Rovelli

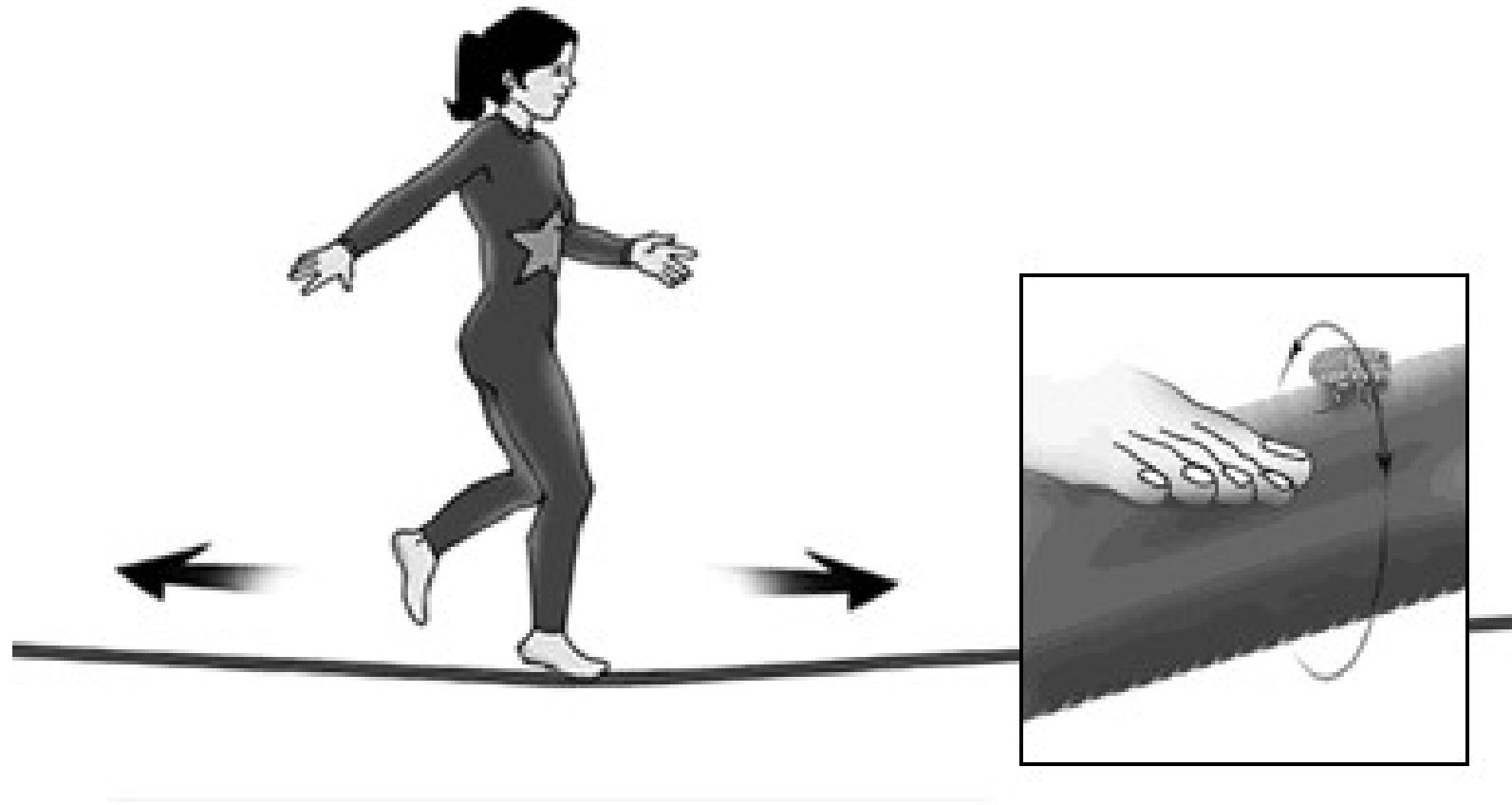
and

M. Vanadia

at INFN (Italy).

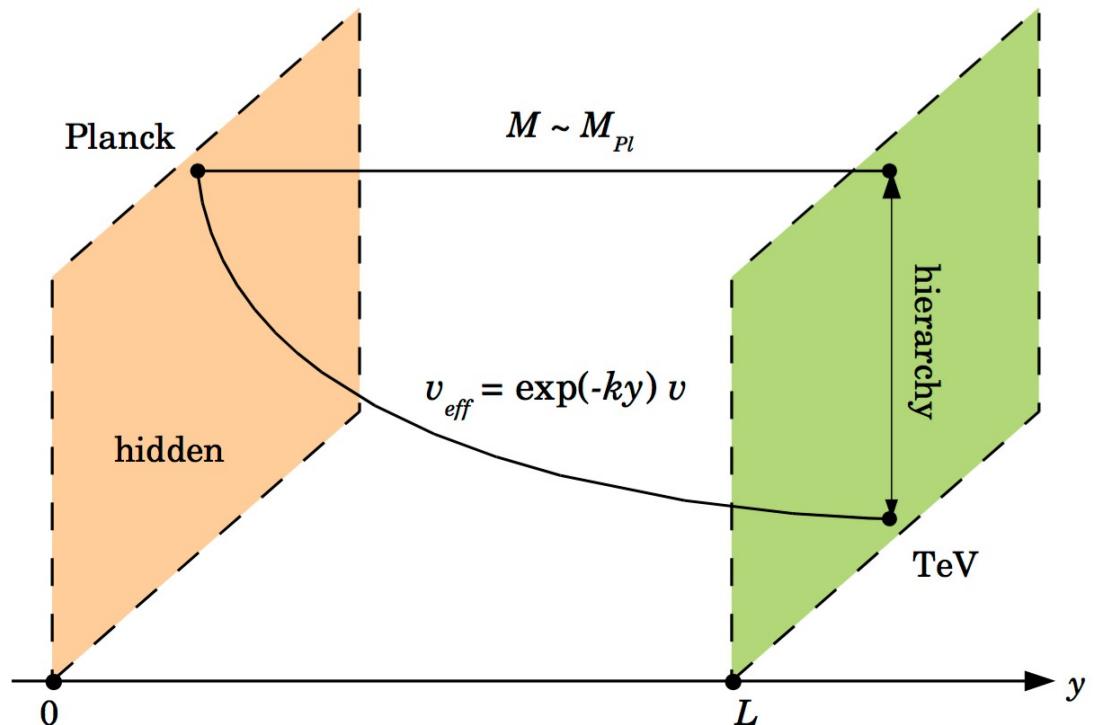
Backup slides

Extra spatial dimensions



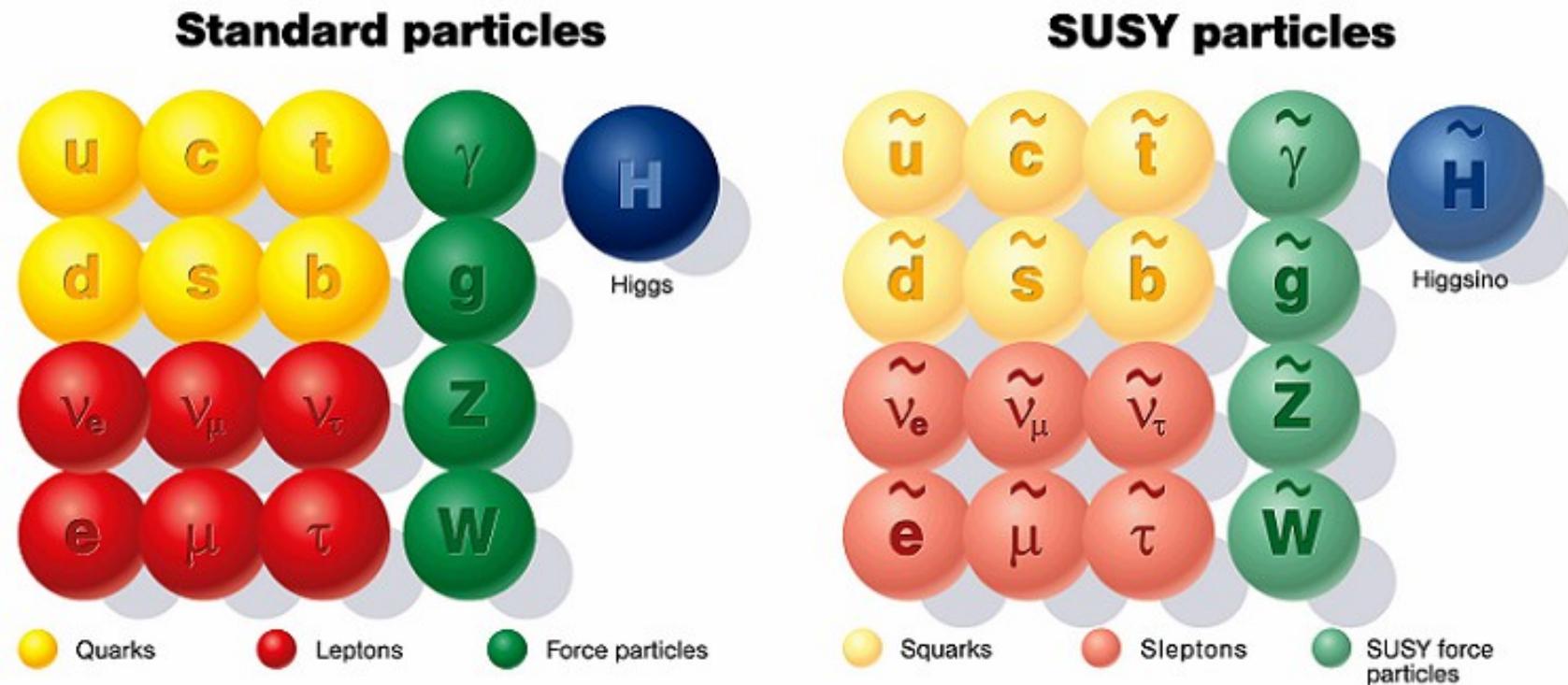
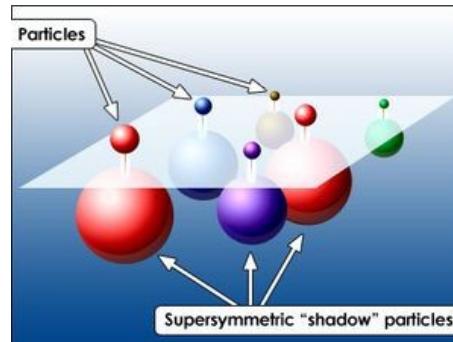
Extra spatial dimensions

- In the Randall-Sundrum model gravity propagates in a warped extra dimension with two fixed points
- The Standard Model fields are constrained to one brane
- The gravity wave function is concentrated near the other brane, falling off exponentially across the extra dimension



$$ds^2 = e^{-2k|y|} \eta_{\mu\nu} dx^\mu dx^\nu + dy^2$$

Two times more particles ??



Event selection

- Trigger: two γ , with $E_T^{\gamma_1} > 35$ GeV, $E_T^{\gamma_2} > 25$ GeV ($> 99\%$ signal efficiency)
- Pre-selection:
 - γ quality criteria: shower shape, leakage in the hadronic calorimeter
 - $E_T^{\gamma_1} > 40$ GeV, $E_T^{\gamma_2} > 30$ GeV
 - $|\eta^{1,2}| < 2.37$, excluded $1.37 \leq |\eta| < 1.52$
- $\epsilon_{identification} = 85\% (90\%)$ for unconverted (converted) γ for $E_T = 25$ GeV
- It asymptotically reaches 95% (98%) for $E_T > 200$ GeV
- γ are required to be **isolated**, i.e. to have little activity in the Inner Detector/calorimeters in a ΔR^1 cone around them

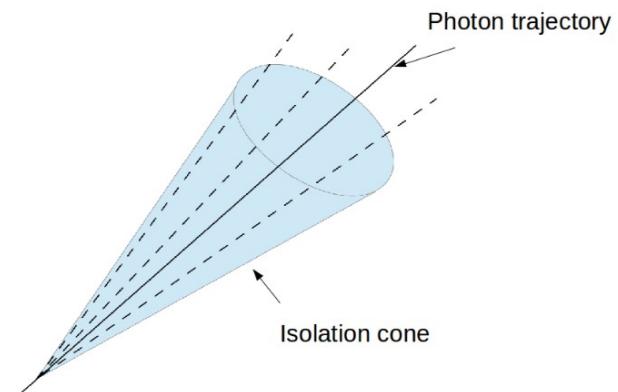
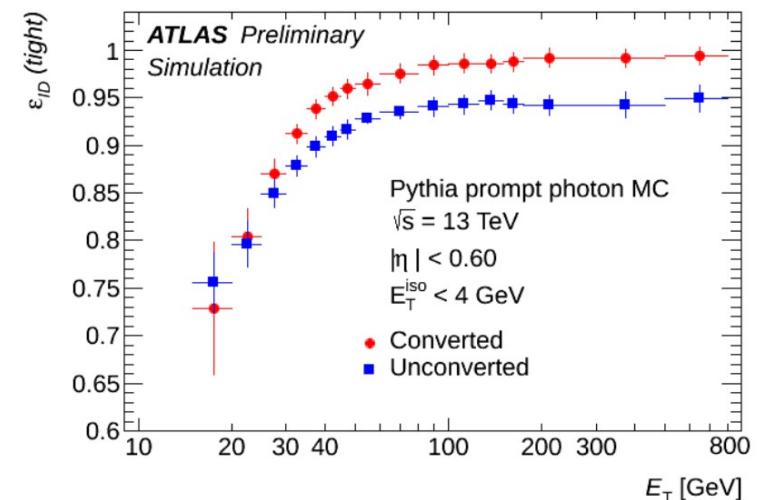
Isolation

- Calorimeter: $E_T^{iso} < 2.45 + 0.022 \cdot E_T^\gamma$ in $\Delta R < 0.4$
- Inner detector: $p_T^{iso} < 0.05 \cdot E_T^\gamma$ in $\Delta R < 0.2$

$$^1 \Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$$

²EGAM-2015-002

Central photon efficiency @ 13 TeV
(simulation)²



From the ATLAS publication

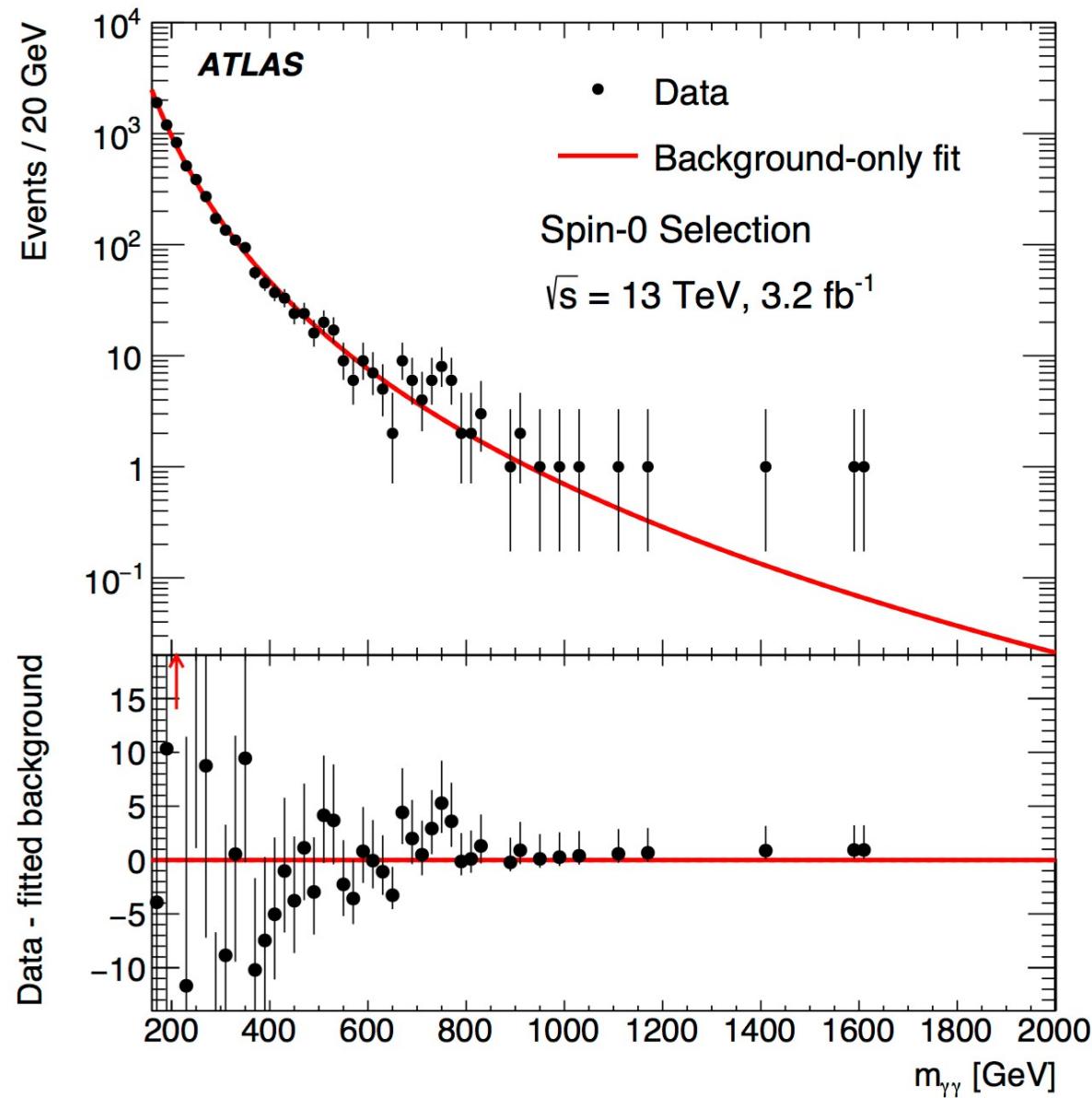
<http://arxiv.org/abs/1606.03833>

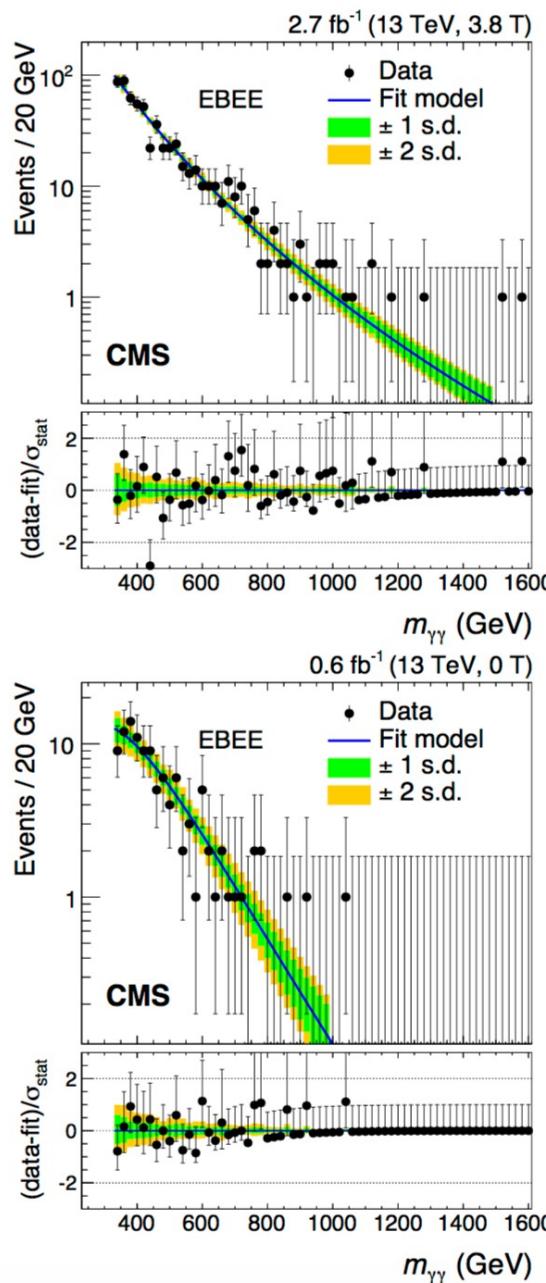
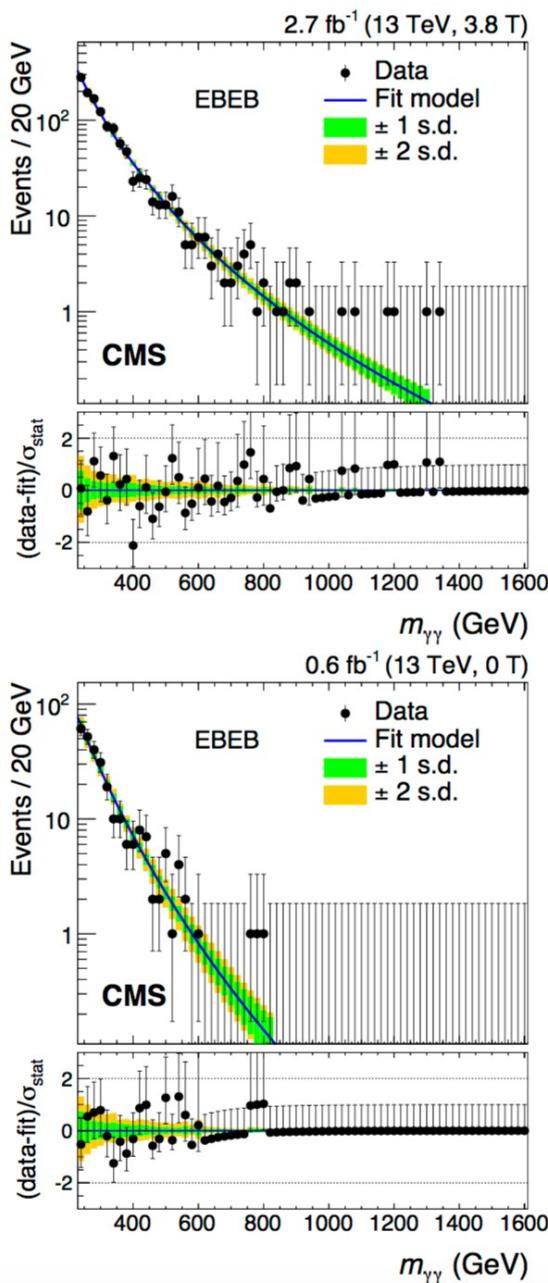
10 Conclusion

Searches for new resonances decaying into two photons in the ATLAS experiment at the LHC are presented. The pp collision data corresponding to an integrated luminosity of 3.2 fb^{-1} were recorded in 2015 at a centre-of-mass energy of $\sqrt{s} = 13 \text{ TeV}$. Analyses optimized for the search for spin-2 Randall–Sundrum graviton resonances with mass above 500 GeV and for spin-0 resonances with mass above 200 GeV are performed. The events selected in the second analysis are a subset of the events selected for the spin-2 search.

Over most of the diphoton mass range, the data are consistent with the background-only hypothesis and 95% CL limits are derived on the cross section for the production of the two benchmark resonances as a function of their masses and widths. Varying both the mass and the decay width of the hypothesised resonance, the largest deviation from the background-only hypothesis is observed in a broad region near a mass of 750 GeV and with a width of about 50 GeV, with local significances of 3.8 and 3.9 standard deviations in the searches optimized for the spin-2 and spin-0 resonances, respectively. The global significances are estimated to be 2.1 standard deviations for both searches. When considering narrow-width signal hypotheses, the largest local significances for the two searches are observed near a mass of 770 GeV and 750 GeV with local significances corresponding to 3.3 and 2.9 standard deviations, respectively. No significant difference is observed in the properties of the events with a diphoton mass near 750 GeV compared to those at higher or lower masses. Assuming a scaling of the production cross section for an s -channel resonance produced by gluon fusion (light quark–antiquark annihilation), the consistency between the 13 TeV data and the data collected at 8 TeV is found to be at the level of 2.7 (3.3) standard deviations using results from the searches optimized for a spin-2 particle and at the level of 1.2 (2.1) standard deviations using results from the searches optimized for a spin-0 particle.

From the ATLAS publication





<http://arxiv.org/abs/1606.04093>

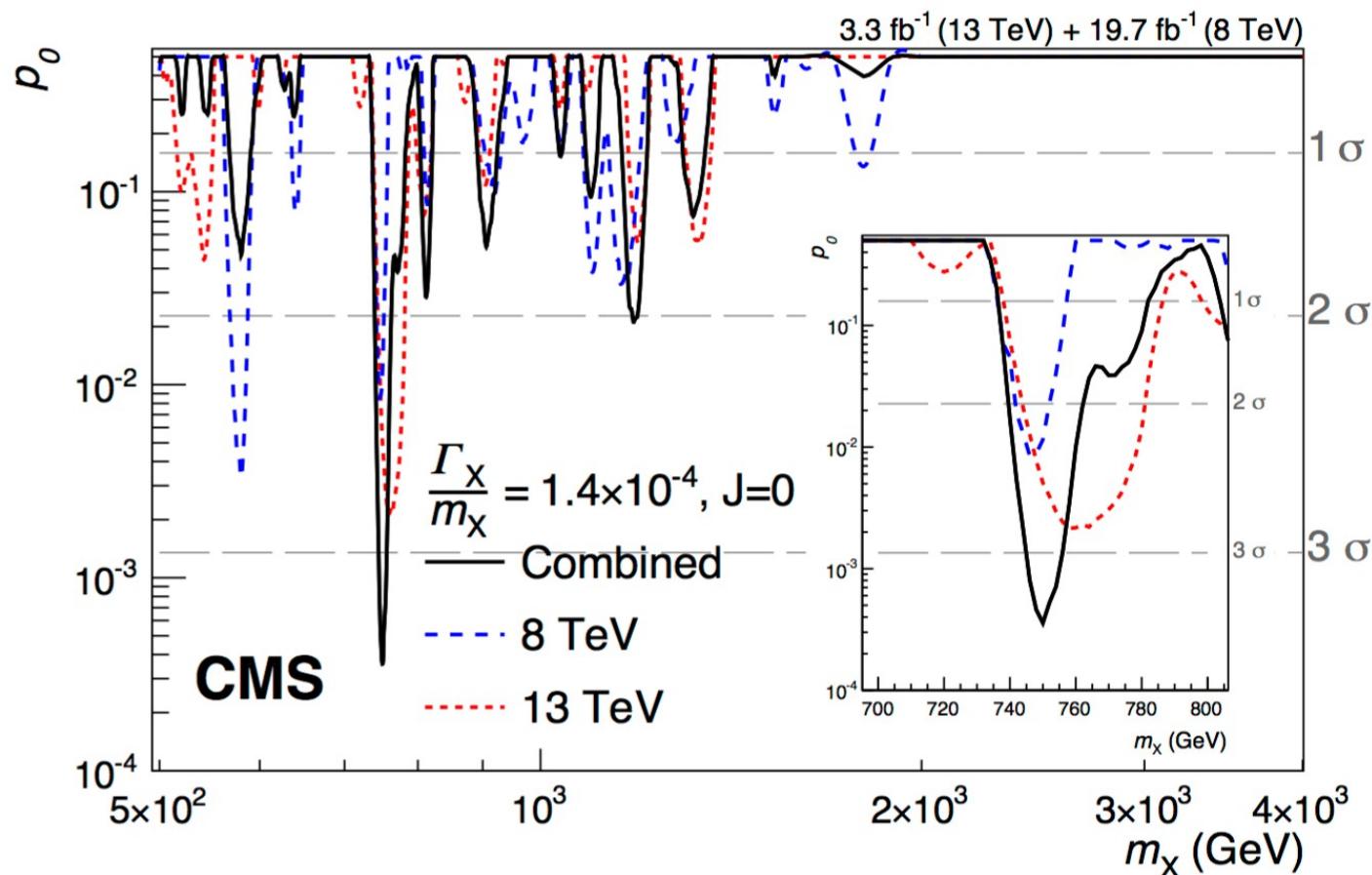
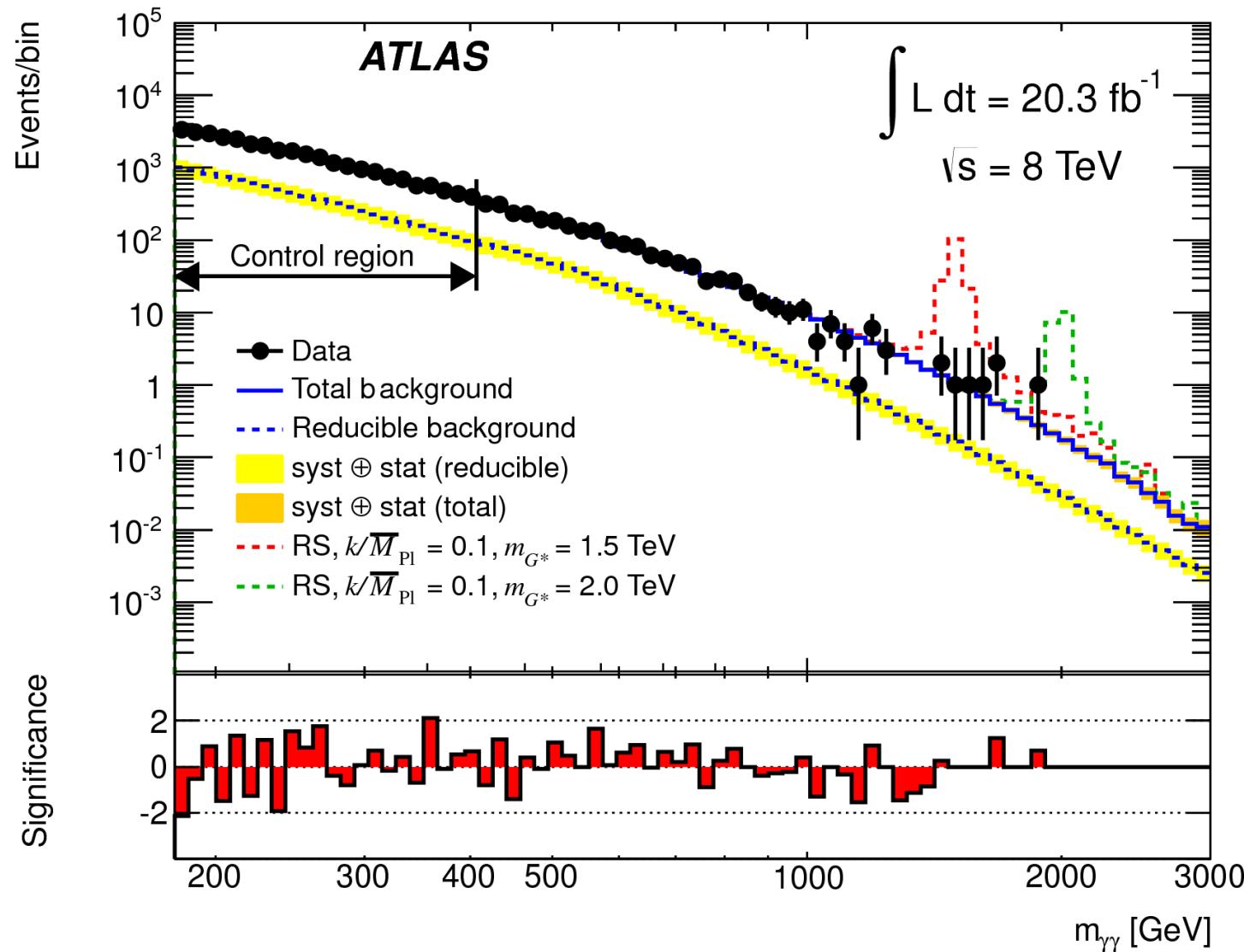
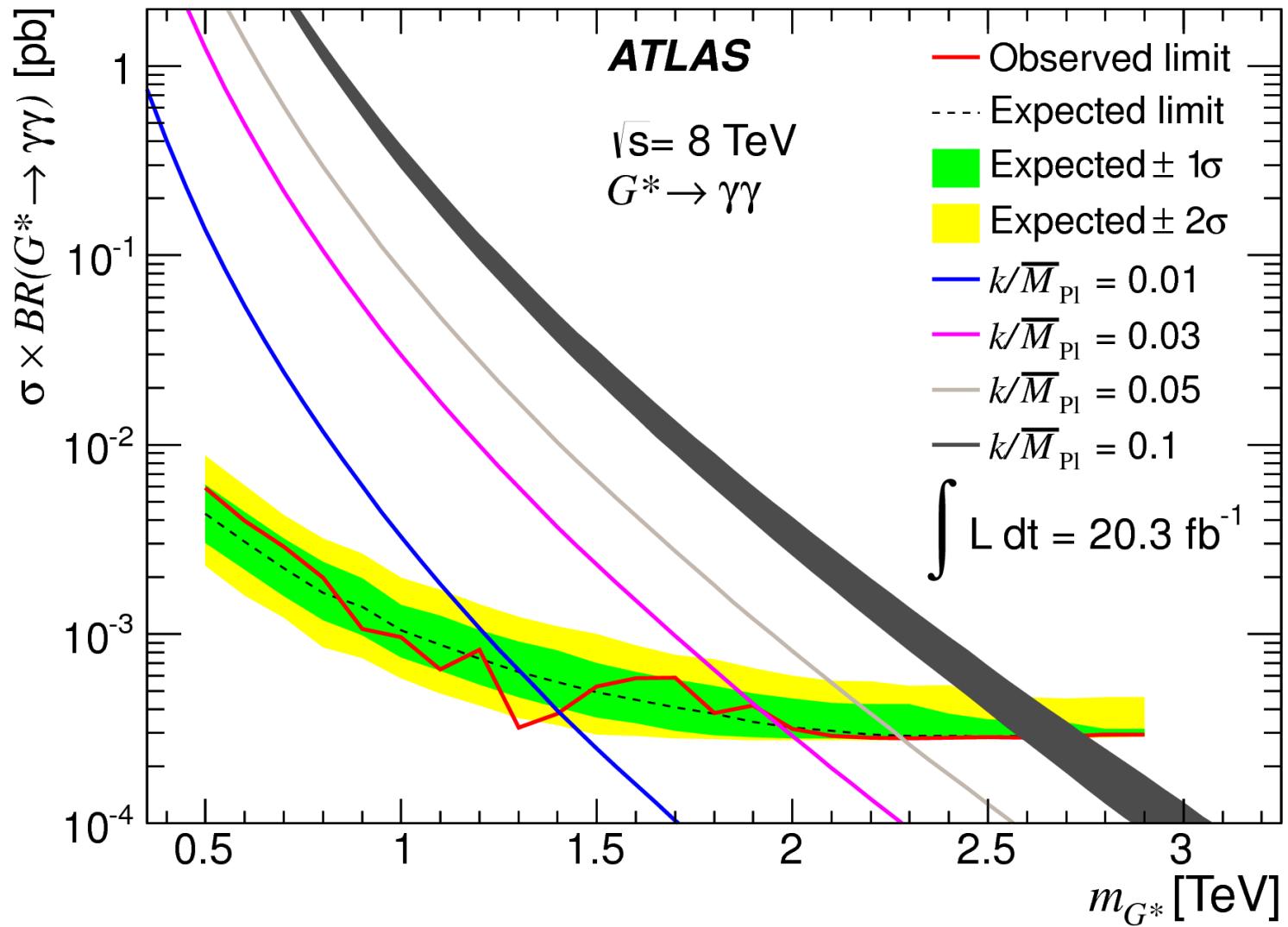


Figure 4: Observed background-only p -values for narrow-width scalar resonances as a function of the resonance mass m_X , from the combined analysis of the 8 and 13 TeV data. The results for the separate 8 and 13 TeV data sets are also shown. The inset shows an expanded region around $m_X = 750$ GeV.

ATLAS at $\sqrt{s} = 8$ TeV



ATLAS at $\sqrt{s} = 8$ TeV



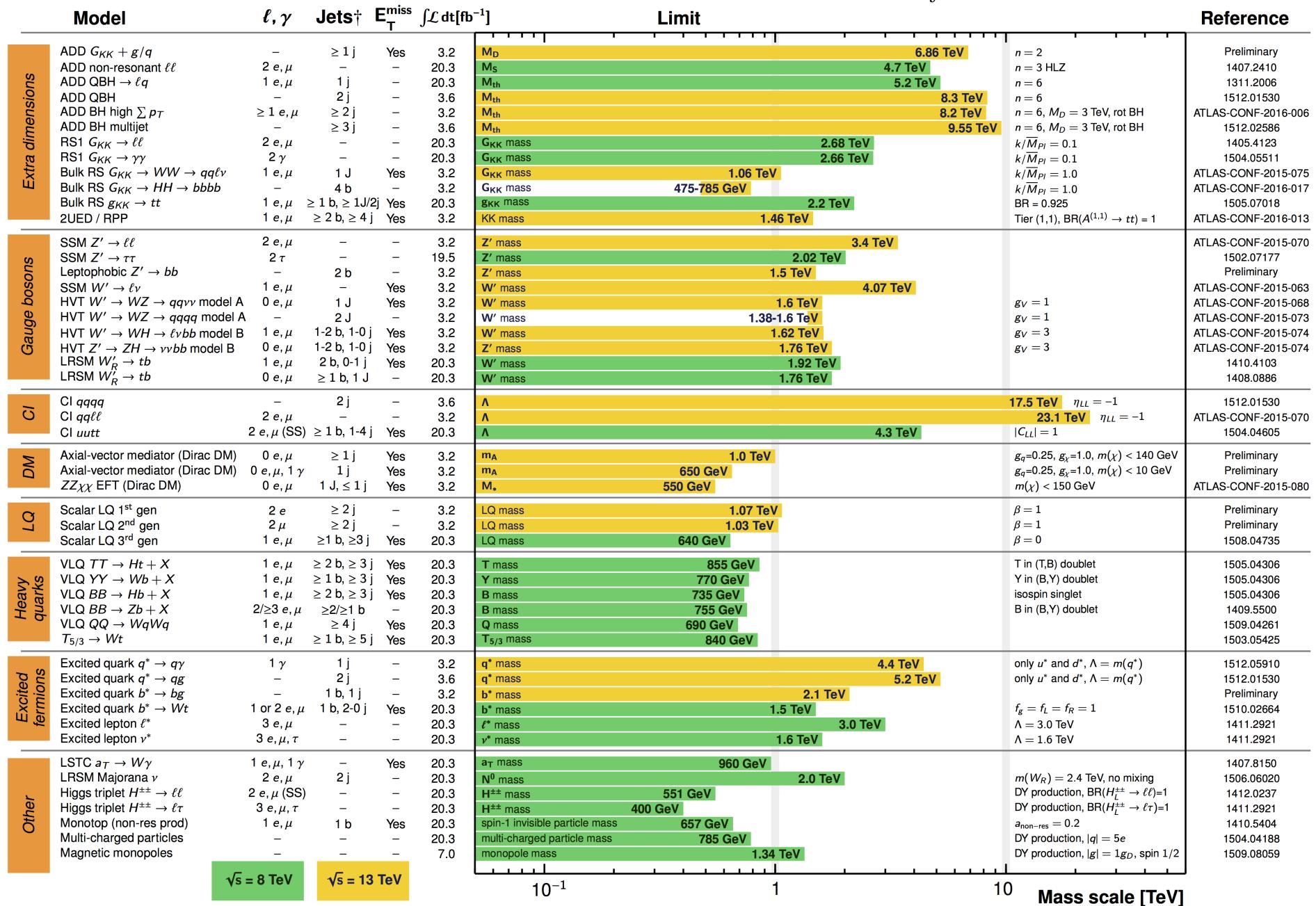
ATLAS Exotics Searches* - 95% CL Exclusion

Status: March 2016

ATLAS Preliminary

$$\int \mathcal{L} dt = (3.2 - 20.3) \text{ fb}^{-1}$$

$\sqrt{s} = 8, 13 \text{ TeV}$



$\sqrt{s} = 8 \text{ TeV}$

$\sqrt{s} = 13 \text{ TeV}$

10⁻¹

1

10

Mass scale [TeV]

*Only a selection of the available mass limits on new states or phenomena is shown. Lower bounds are specified only when explicitly not excluded.

†Small-radius (large-radius) jets are denoted by the letter j (J).

ATLAS SUSY Searches* - 95% CL Lower Limits

Status: March 2016

ATLAS Preliminary

$\sqrt{s} = 7, 8, 13 \text{ TeV}$

Model	e, μ, τ, γ	Jets	E_T^{miss}	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Mass limit	$\sqrt{s} = 7, 8 \text{ TeV}$	$\sqrt{s} = 13 \text{ TeV}$	Reference
Inclusive Searches	MSUGRA/CMSSM	0-3 $e, \mu/1-2 \tau$	2-10 jets/3 b	Yes	20.3	\tilde{q}, \tilde{g}	1.85 TeV	$m(\tilde{q})=m(\tilde{g})$
	$\tilde{q}\tilde{q}, \tilde{q}\rightarrow q\chi_1^0$	0	2-6 jets	Yes	3.2	\tilde{q}	980 GeV	$m(\tilde{q}^0)=0 \text{ GeV}, m(\text{1st gen. } \tilde{q})=m(\text{2nd gen. } \tilde{q})$
	$\tilde{q}\tilde{q}, \tilde{q}\rightarrow q\chi_1^0$ (compressed)	mono-jet	1-3 jets	Yes	3.2	\tilde{q}	610 GeV	$m(\tilde{q})-m(\tilde{\chi}_1^0) < 5 \text{ GeV}$
	$\tilde{q}\tilde{q}, \tilde{q}\rightarrow q(\ell\ell/\ell\nu/\nu\nu)\chi_1^0$	2 e, μ (off-Z)	2 jets	Yes	20.3	\tilde{q}	820 GeV	$m(\tilde{q})_1=0 \text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g}\rightarrow q\tilde{q}\chi_1^0$	0	2-6 jets	Yes	3.2	\tilde{g}	1.52 TeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g}\rightarrow q\tilde{q}\chi_1^0 \rightarrow qqW^\pm\chi_1^0$	1 e, μ	2-6 jets	Yes	3.3	\tilde{g}	1.6 TeV	$m(\tilde{\chi}_1^0) < 350 \text{ GeV}, m(\tilde{\chi}^\pm)=0.5(m(\tilde{\chi}_1^0)+m(\tilde{g}))$
	$\tilde{g}\tilde{g}, \tilde{g}\rightarrow q\tilde{q}\chi_1^0 \rightarrow q\ell\ell/\ell\nu/\nu\nu\chi_1^0$	2 e, μ	0-3 jets	-	20	\tilde{g}	1.38 TeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g}\rightarrow q\tilde{q}\chi_1^0 \rightarrow qqW^\pm\chi_1^0$	0	7-10 jets	Yes	3.2	\tilde{g}	1.4 TeV	$m(\tilde{\chi}_1^0)=100 \text{ GeV}$
	GMSB ($\tilde{\ell}$ NLSP)	1-2 $\tau + 0-1 \ell$	0-2 jets	Yes	20.3	\tilde{g}	1.63 TeV	$\tan\beta > 20$
	GGM (bino NLSP)	2 γ	-	Yes	20.3	\tilde{g}	1.34 TeV	$c\tau(\text{NLSP}) < 0.1 \text{ mm}$
GGM (higgsino-bino NLSP)	GGM (higgsino-bino NLSP)	γ	1 b	Yes	20.3	\tilde{g}	1.37 TeV	$m(\tilde{\chi}_1^0) < 950 \text{ GeV}, c\tau(\text{NLSP}) < 0.1 \text{ mm}, \mu < 0$
	GGM (higgsino-bino NLSP)	γ	2 jets	Yes	20.3	\tilde{g}	1.3 TeV	$m(\text{NLSP}) < 850 \text{ GeV}, c\tau(\text{NLSP}) < 0.1 \text{ mm}, \mu > 0$
	GGM (higgsino NLSP)	2 e, μ (Z)	2 jets	Yes	20.3	\tilde{g}	900 GeV	$m(\text{NLSP}) > 430 \text{ GeV}$
	Gravitino LSP	0	mono-jet	Yes	20.3	$F^{1/2} \text{ scale}$	865 GeV	$m(\tilde{G}) > 1.8 \times 10^{-4} \text{ eV}, m(\tilde{g})=m(\tilde{q})=1.5 \text{ TeV}$
3^{rd} gen. \tilde{g} med.	$\tilde{g}\tilde{g}, \tilde{g}\rightarrow bb\chi_1^0$	0	3 b	Yes	3.3	\tilde{g}	1.78 TeV	$m(\tilde{\chi}_1^0) < 800 \text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g}\rightarrow b\tilde{\ell}\chi_1^0$	0-1 e, μ	3 b	Yes	3.3	\tilde{g}	1.76 TeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g}\rightarrow b\tilde{\ell}\chi_1^+$	0-1 e, μ	3 b	Yes	20.1	\tilde{g}	1.37 TeV	$m(\tilde{\chi}_1^0) < 300 \text{ GeV}$
3^{rd} gen. direct production	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1\rightarrow b\tilde{\chi}_1^0$	0	2 b	Yes	3.2	\tilde{b}_1	840 GeV	$m(\tilde{\chi}_1^0) < 100 \text{ GeV}$
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1\rightarrow b\tilde{\chi}_1^\pm$	2 e, μ (SS)	0-3 b	Yes	3.2	\tilde{b}_1	325-540 GeV	$m(\tilde{\chi}_1^0)=50 \text{ GeV}, m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_1^0)+100 \text{ GeV}$
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1\rightarrow b\tilde{\chi}_1^\pm$	1-2 e, μ	1-2 b	Yes	4.7/20.3	\tilde{t}_1	117-170 GeV	$m(\tilde{\chi}_1^\pm)=2m(\tilde{\chi}_1^0), m(\tilde{\chi}_1^0)=55 \text{ GeV}$
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1\rightarrow W\tilde{\chi}_1^0$ or $\tilde{\chi}_1^0$	0-2 e, μ	0-2 jets/1-2 b	Yes	20.3	\tilde{t}_1	200-500 GeV	$m(\tilde{\chi}_1^0)=1 \text{ GeV}$
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1\rightarrow \tilde{\chi}_1^0$ (natural GMSB)	0	mono-jet/c-tag	Yes	20.3	\tilde{t}_1	90-198 GeV	$m(\tilde{t}_1)-m(\tilde{\chi}_1^0) < 85 \text{ GeV}$
	$\tilde{t}_2\tilde{t}_2, \tilde{t}_2\rightarrow \tilde{\chi}_1^0$	2 e, μ (Z)	1 b	Yes	20.3	\tilde{t}_2	150-600 GeV	$m(\tilde{\chi}_1^0) > 150 \text{ GeV}$
	$\tilde{t}_2\tilde{t}_2, \tilde{t}_2\rightarrow \tilde{\chi}_1^0 + Z$	3 e, μ (Z)	1 b	Yes	20.3	\tilde{t}_2	290-610 GeV	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}$
EW direct	$\tilde{\ell}_{\text{LR}}\tilde{\ell}_{\text{LR}}, \tilde{\ell}\rightarrow \ell\tilde{\chi}_1^0$	2 e, μ	0	Yes	20.3	$\tilde{\ell}$	90-335 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$
	$\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow \ell\tilde{\nu}(\ell\tilde{\nu})$	2 e, μ	0	Yes	20.3	$\tilde{\chi}_1^\pm$	140-475 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}, m(\tilde{\ell}, \tilde{\nu})=0.5(m(\tilde{\chi}_1^+)+m(\tilde{\chi}_1^-))$
	$\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow \tau\tilde{\nu}(\tau\tilde{\nu})$	2 τ	-	Yes	20.3	$\tilde{\chi}_1^\pm$	355 GeV	$m(\tilde{\chi}_1^+)=0 \text{ GeV}, m(\tilde{\tau}, \tilde{\nu})=0.5(m(\tilde{\chi}_1^+)+m(\tilde{\chi}_1^-))$
	$\tilde{\chi}_1^+\tilde{\chi}_2^0 \rightarrow \ell_L v_L \ell_L \ell(\bar{\nu}), \ell\tilde{\nu}\ell_L \ell(\bar{\nu})$	3 e, μ	0	Yes	20.3	$\tilde{\chi}_1^+\tilde{\chi}_2^0$	715 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)=0, m(\tilde{\ell}, \tilde{\nu})=0.5(m(\tilde{\chi}_1^+)+m(\tilde{\chi}_1^-))$
	$\tilde{\chi}_1^+\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0 Z\tilde{\chi}_1^0$	2-3 e, μ	0-2 jets	Yes	20.3	$\tilde{\chi}_1^+\tilde{\chi}_2^0$	425 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)=0, \text{sleptons decoupled}$
	$\tilde{\chi}_1^+\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0 h\tilde{\chi}_1^0, h\rightarrow bb/WW/\tau\tau/\gamma\gamma$	e, μ, γ	0-2 b	Yes	20.3	$\tilde{\chi}_1^+\tilde{\chi}_2^0$	270 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)=0, \text{sleptons decoupled}$
	$\tilde{\chi}_2^0\tilde{\chi}_3^0, \tilde{\chi}_2^0\rightarrow \tilde{\ell}\tilde{\ell}$	4 e, μ	0	Yes	20.3	$\tilde{\chi}_{2,3}^0$	635 GeV	$m(\tilde{\chi}_2^0)=m(\tilde{\chi}_3^0), m(\tilde{\chi}_1^0)=0, m(\tilde{\ell}, \tilde{\nu})=0.5(m(\tilde{\chi}_2^0)+m(\tilde{\chi}_3^0))$
	GGM (wino NLSP) weak prod.	1 $e, \mu + \gamma$	-	Yes	20.3	\tilde{W}	115-370 GeV	$c\tau < 1 \text{ mm}$
	Direct $\tilde{\chi}_1^+\tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk	1 jet	Yes	20.3	$\tilde{\chi}_1^\pm$	270 GeV	$m(\tilde{\chi}_1^\pm)-m(\tilde{\chi}_1^0) \sim 160 \text{ MeV}, \tau(\tilde{\chi}_1^\pm)=0.2 \text{ ns}$
	Direct $\tilde{\chi}_1^+\tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$	dE/dx trk	-	Yes	18.4	$\tilde{\chi}_1^\pm$	495 GeV	$m(\tilde{\chi}_1^\pm)-m(\tilde{\chi}_1^0) \sim 160 \text{ MeV}, \tau(\tilde{\chi}_1^\pm) < 15 \text{ ns}$
Long-lived particles	Stable, stopped \tilde{g} R-hadron	0	1-5 jets	Yes	27.9	\tilde{g}	850 GeV	$m(\tilde{\chi}_1^0)=100 \text{ GeV}, 10 \mu\text{s} < \tau(\tilde{g}) < 1000 \text{ s}$
	Metastable \tilde{g} R-hadron	dE/dx trk	-	-	3.2	\tilde{g}	1.54 TeV	$m(\tilde{\chi}_1^0)=100 \text{ GeV}, \tau > 10 \text{ ns}$
	GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^-\rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu})+\tau(e, \mu)$	1-2 μ	-	-	19.1	$\tilde{\chi}_1^0$	537 GeV	$10 < \tan\beta < 50$
	GMSB, $\tilde{\chi}_1^0\rightarrow \gamma\tilde{G}$, long-lived $\tilde{\chi}_1^0$	2 γ	-	Yes	20.3	$\tilde{\chi}_1^0$	440 GeV	$1 < \tau(\tilde{\chi}_1^0) < 3 \text{ ns}, \text{SPS8 model}$
	$\tilde{g}\tilde{g}, \tilde{g}\tilde{g}\rightarrow ee/\nu\nu/\mu\nu$	displ. ee/ep/ep/mu/mu	-	-	20.3	$\tilde{\chi}_1^0$	1.0 TeV	$7 < c\tau(\tilde{\chi}_1^0) < 740 \text{ mm}, m(\tilde{g})=1.3 \text{ TeV}$
	GGM $\tilde{g}\tilde{g}, \tilde{g}\tilde{g}\rightarrow Z\tilde{G}$	displ. vtx + jets	-	-	20.3	$\tilde{\chi}_1^0$	1.0 TeV	$6 < c\tau(\tilde{\chi}_1^0) < 480 \text{ mm}, m(\tilde{g})=1.1 \text{ TeV}$
	LFV $pp\rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau\rightarrow e\mu/\mu\tau$	$e\mu, e\tau, \mu\tau$	-	-	20.3	$\tilde{\nu}_\tau$	1.7 TeV	$\lambda'_{111}=0.11, \lambda'_{122}/\lambda'_{132}/\lambda'_{233}=0.07$
RPV	Bilinear RPV CMSSM	2 e, μ (SS)	0-3 b	Yes	20.3	\tilde{q}, \tilde{g}	1.45 TeV	$m(\tilde{q})=m(\tilde{g}), c\tau_{LP} < 1 \text{ mm}$
	$\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow ee\tilde{\nu}_e, e\mu\tilde{\nu}_e$	4 e, μ	-	Yes	20.3	$\tilde{\chi}_1^\pm$	760 GeV	$m(\tilde{\chi}_1^0) > 0.2 \times m(\tilde{\chi}_1^\pm), \lambda'_{121} \neq 0$
	$\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow \tau\tau\tilde{\nu}_e, e\tau\tilde{\nu}_\tau$	3 $e, \mu + \tau$	-	Yes	20.3	$\tilde{\chi}_1^\pm$	450 GeV	$m(\tilde{\chi}_1^0) > 0.2 \times m(\tilde{\chi}_1^\pm), \lambda'_{133} \neq 0$
	$\tilde{g}\tilde{g}, \tilde{g}\rightarrow q\tilde{q}qq$	0	6-7 jets	-	20.3	\tilde{g}	917 GeV	$\text{BR}(t)=\text{BR}(b)=\text{BR}(c)=0\%$
	$\tilde{g}\tilde{g}, \tilde{g}\rightarrow q\tilde{q}\chi_1^0, \tilde{\chi}_1^0 \rightarrow qq$	0	6-7 jets	-	20.3	\tilde{g}	980 GeV	$m(\tilde{\chi}_1^0)=600 \text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g}\rightarrow \tilde{t}_1\tilde{t}_1, \tilde{t}_1\rightarrow bs$	2 e, μ (SS)	0-3 b	Yes	20.3	\tilde{g}	880 GeV	1404.2500
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1\rightarrow bs$	0	2 jets + 2 b	-	20.3	\tilde{t}_1	320 GeV	1601.07453
Other	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1\rightarrow bl$	2 e, μ	2 b	-	20.3	\tilde{t}_1	0.4-1.0 TeV	$\text{BR}(\tilde{t}_1\rightarrow be/\mu) > 20\%$
	Scalar charm, $\tilde{c}\rightarrow c\tilde{\chi}_1^0$	0	2 c	Yes	20.3	\tilde{c}	510 GeV	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}$

*Only a selection of the available mass limits on new states or phenomena is shown.

10⁻¹ 1 Mass scale [TeV]

Excitement

