

Non-equilibrium QFT in cosmology: a case of resonant leptogenesis

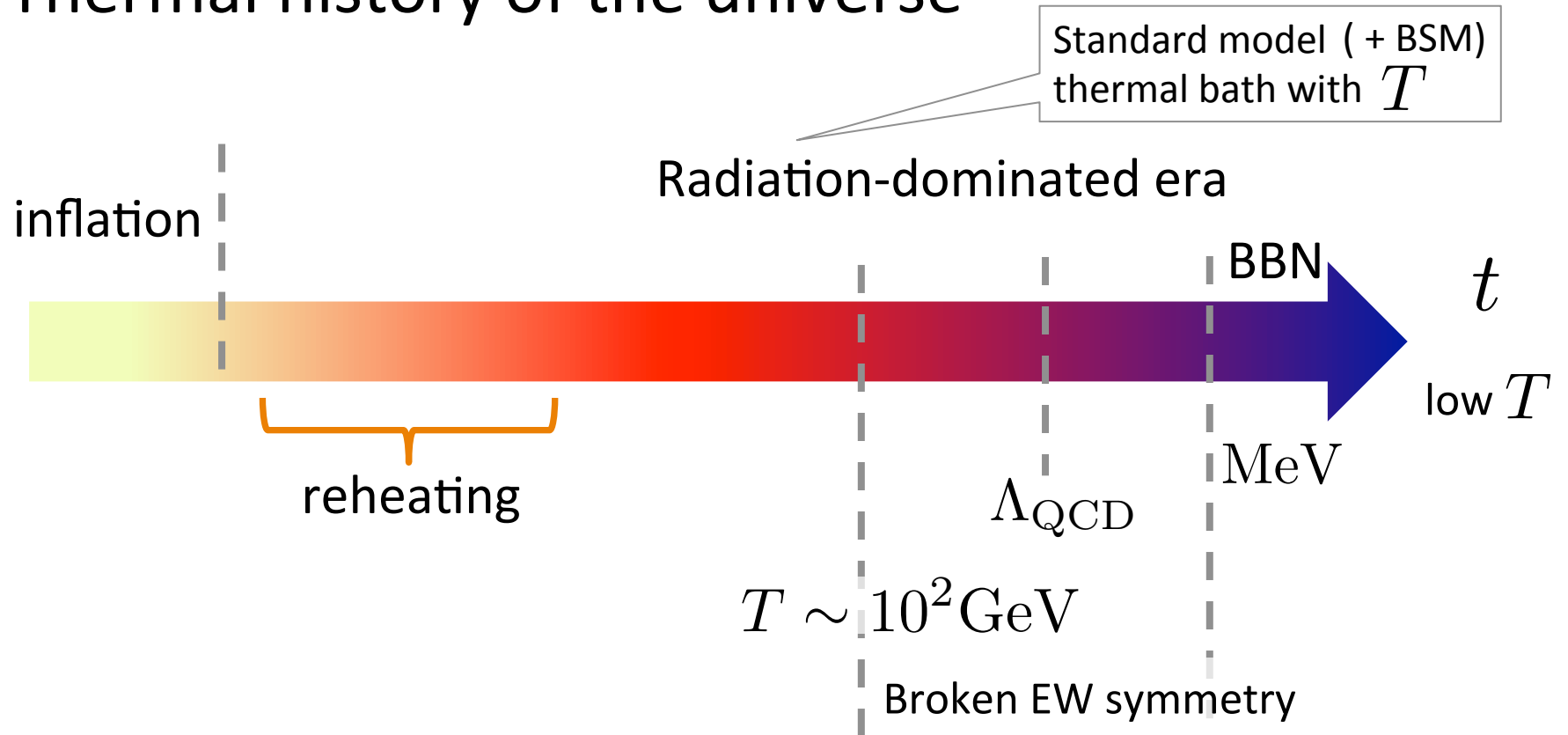
Kengo Shimada (LAPTh)

Journée Théorie du CPTGA, 12 mai 2016

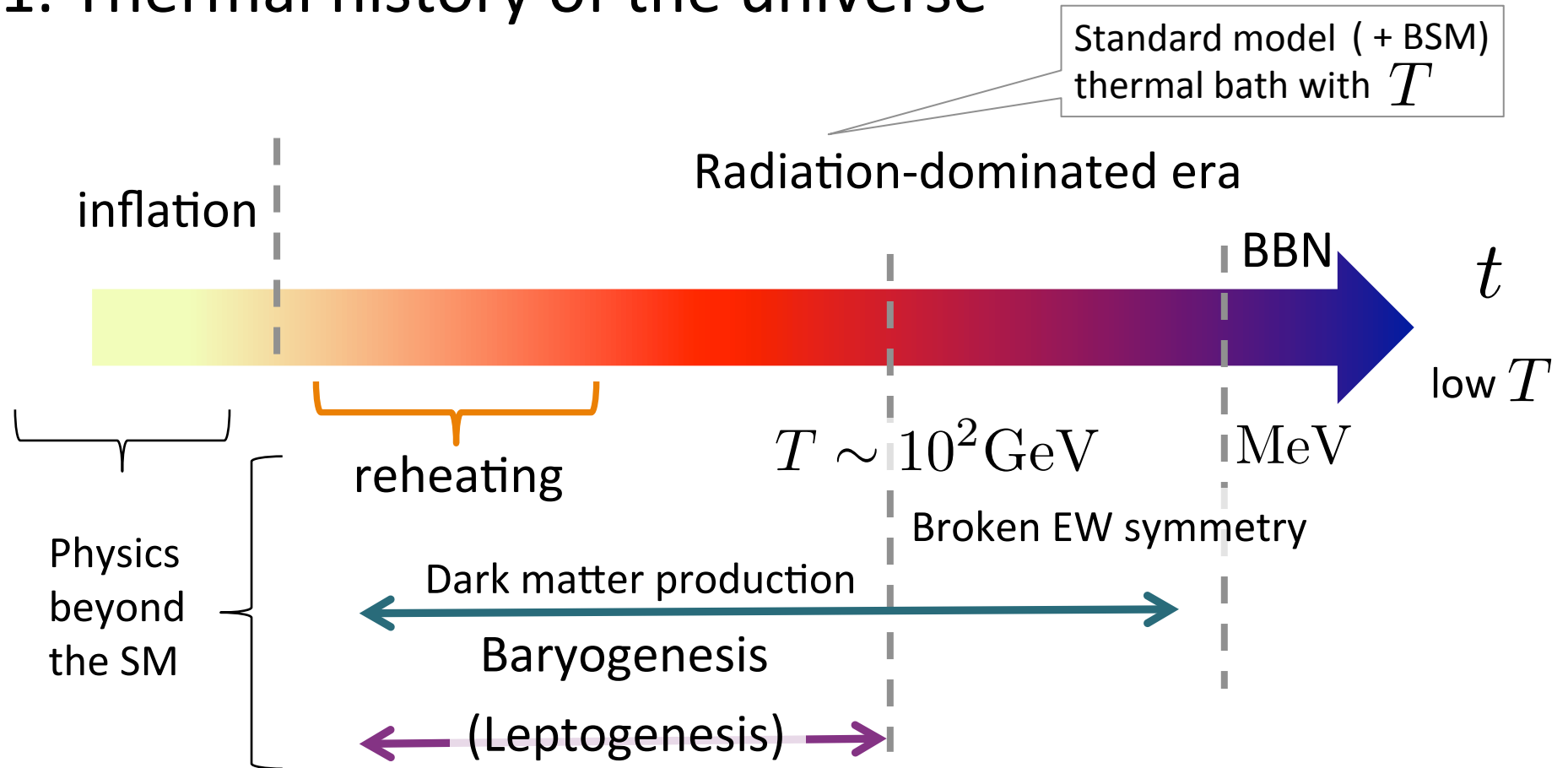
Based on JHEP 1408, 043 arXiv:1404.4816[hep-ph] with S.Iso

JHEP 1404, 062 arXiv:1312.7680[hep-ph] with S.Iso and M.Yamanaka

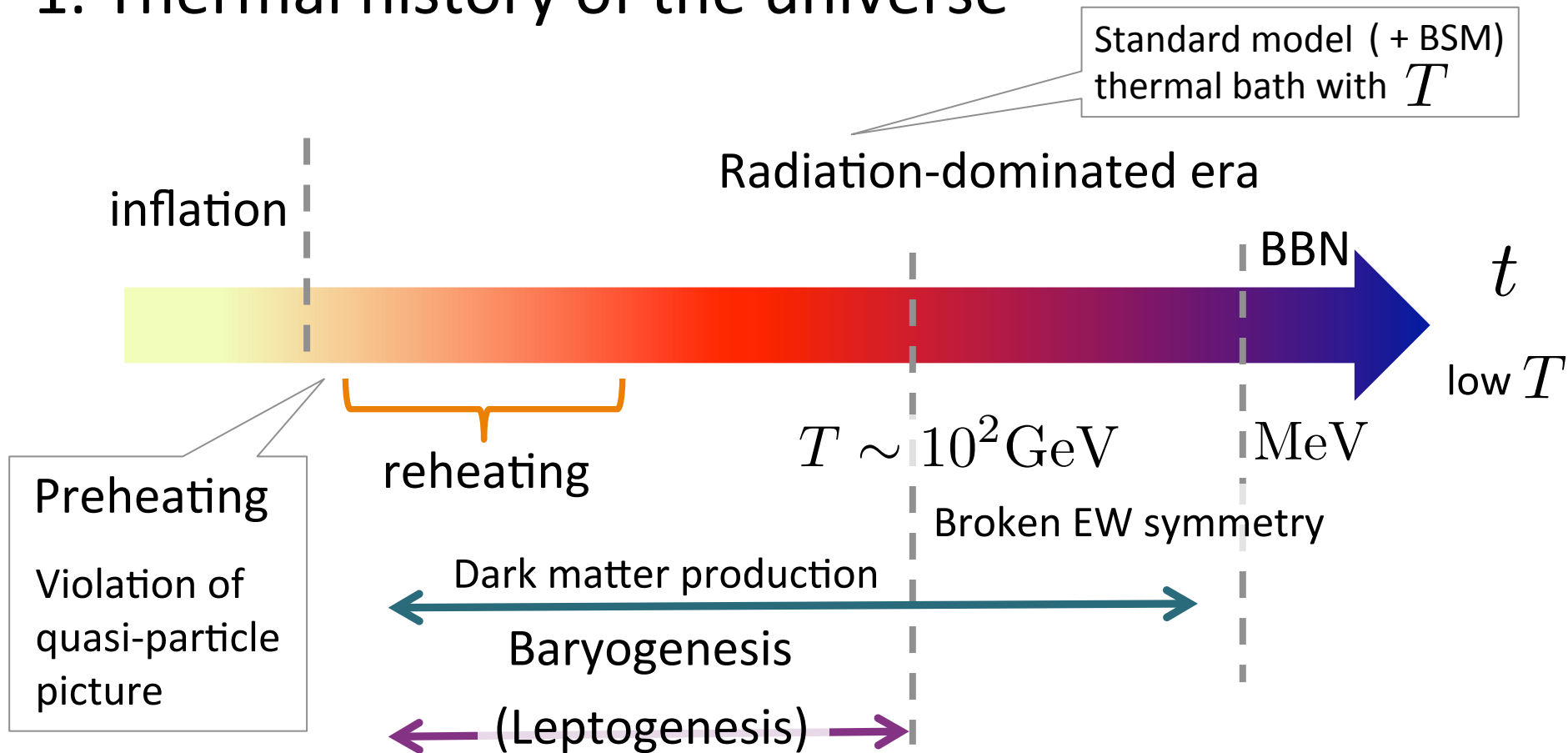
1. Thermal history of the universe



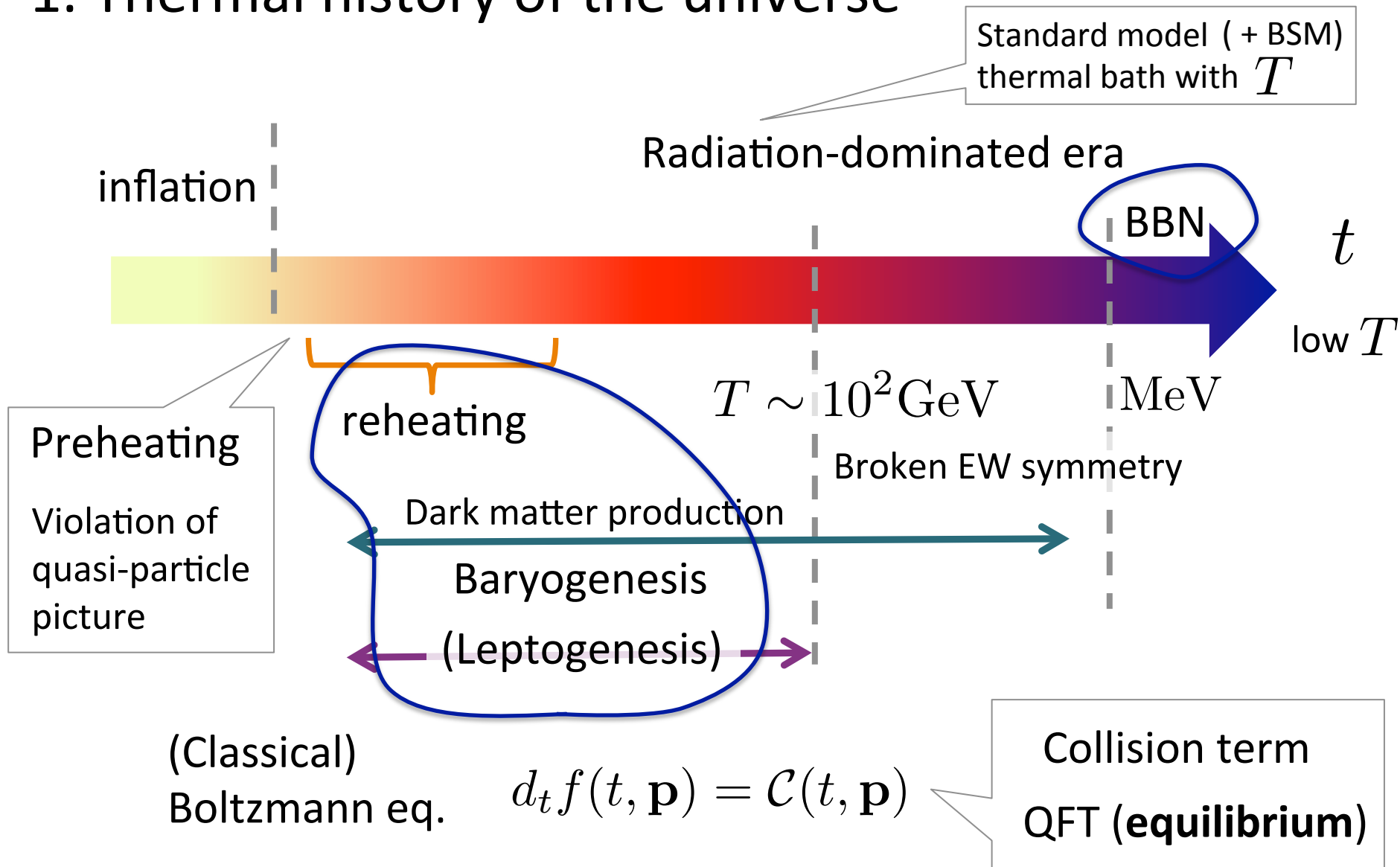
1. Thermal history of the universe



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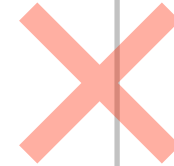
1. Thermal history of the universe



Even if the system is almost thermalized, (and quasi-particle picture is valid,)

Classical Boltzmann eq. + equilibrium QFT

$$d_t f(t, \mathbf{p}) = \mathcal{C}(t, \mathbf{p})$$

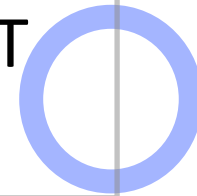


can be invalid because of

degenerate time scales.

An example: Resonant leptogenesis

Single framework of non-equilibrium QFT



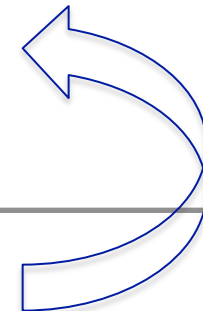
1. Thermal history of the universe
2. Leptogenesis
3. Resonant leptogenesis
4. Kadanoff-Baym equation (non-equilibrium QFT)
5. Summary

2. Leptogenesis

Fukugita, Yanagida (86)

- Baryon asymmetry of the universe

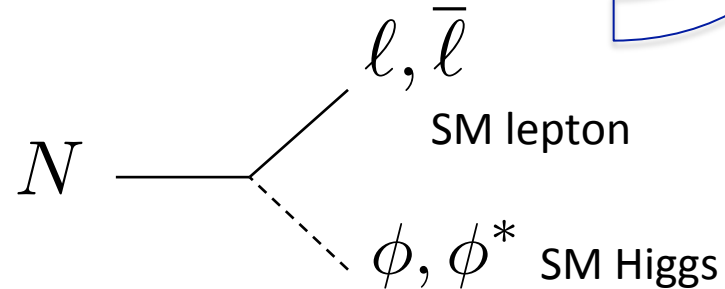
$$\eta_{\Delta B} \equiv \frac{n_B - n_{\bar{B}}}{n_\gamma} \approx 6 \times 10^{-10}$$



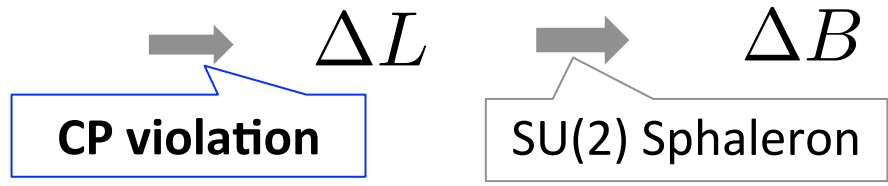
Right-handed neutrino (with Majorana mass)

$$N_i = \nu_{Ri} + \nu_{Ri}^c \quad i = 1, 2, 3$$

Hubble expansion



➔ Out-of-equilibrium decay

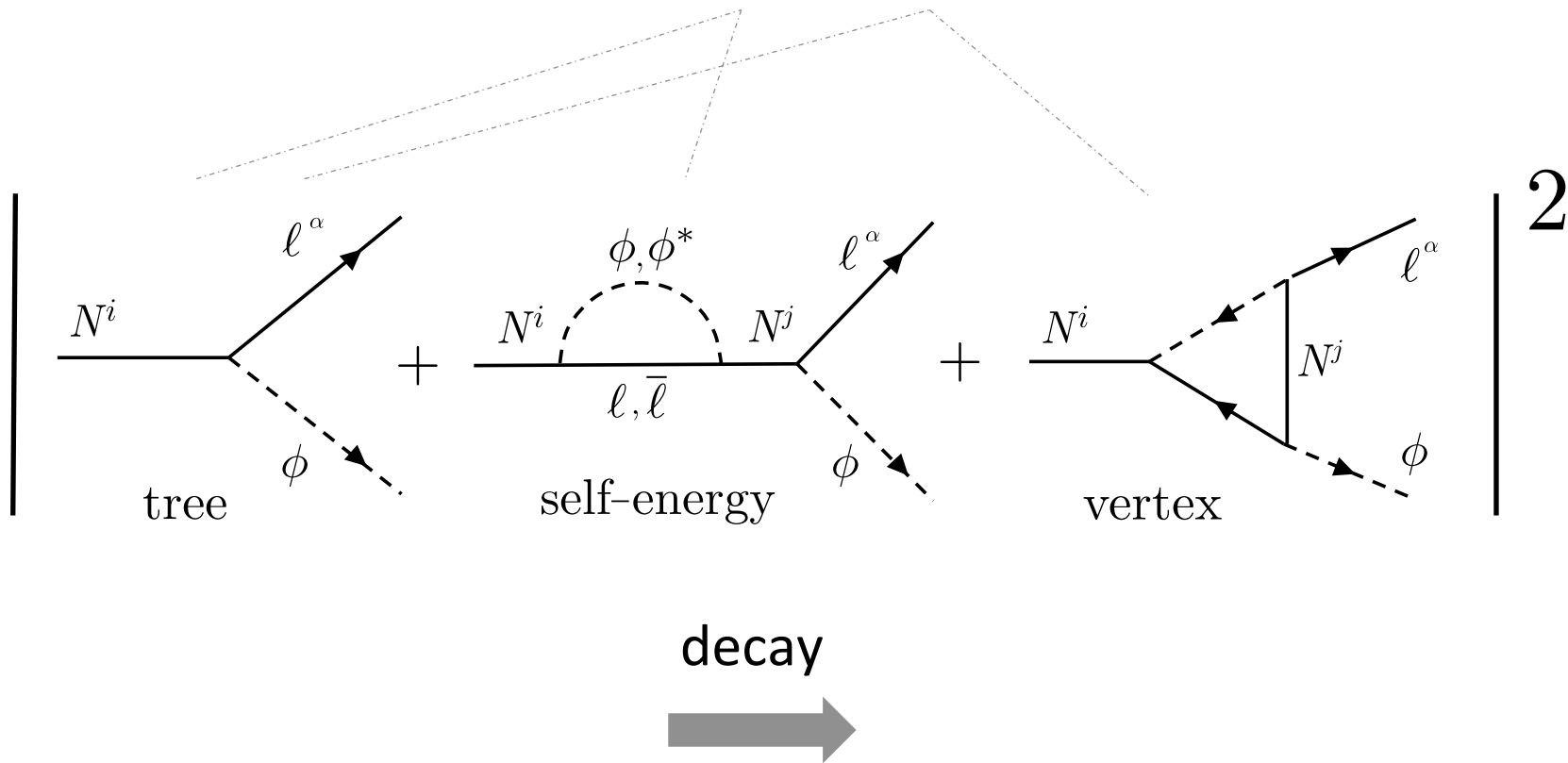


- Tiny SM neutrino mass (Seesaw mechanism)

2. Leptogenesis

Calculating **CP-violating parameter** $\varepsilon \equiv \frac{\Gamma_{N \rightarrow l\phi} - \Gamma_{N \rightarrow \bar{l}\phi^*}}{\Gamma_{N \rightarrow l\phi} + \Gamma_{N \rightarrow \bar{l}\phi^*}}$
 in "equilibrium QFT"

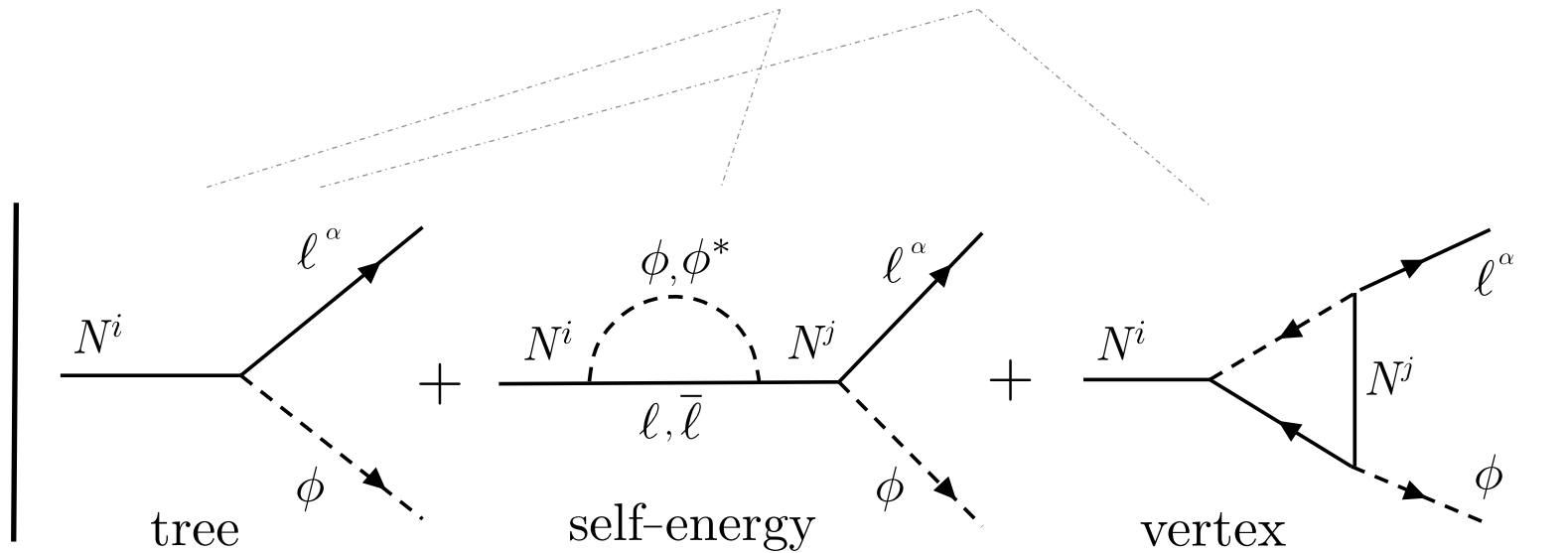
(Lowest order) **CP violation**



2. Leptogenesis

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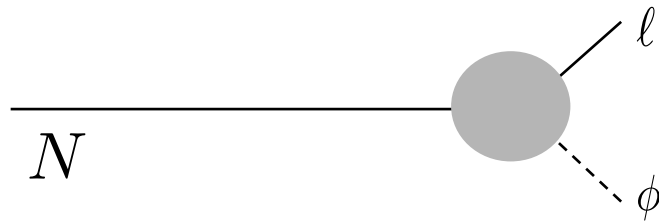
(Lowest order) **CP violation**



RH neutrino flavor oscillation
 $N_i \rightarrow N_{j \neq i}$

dominates
 in **resonant case**

(Classical) Boltzmann equation + equilibrium QFT



Decay at t

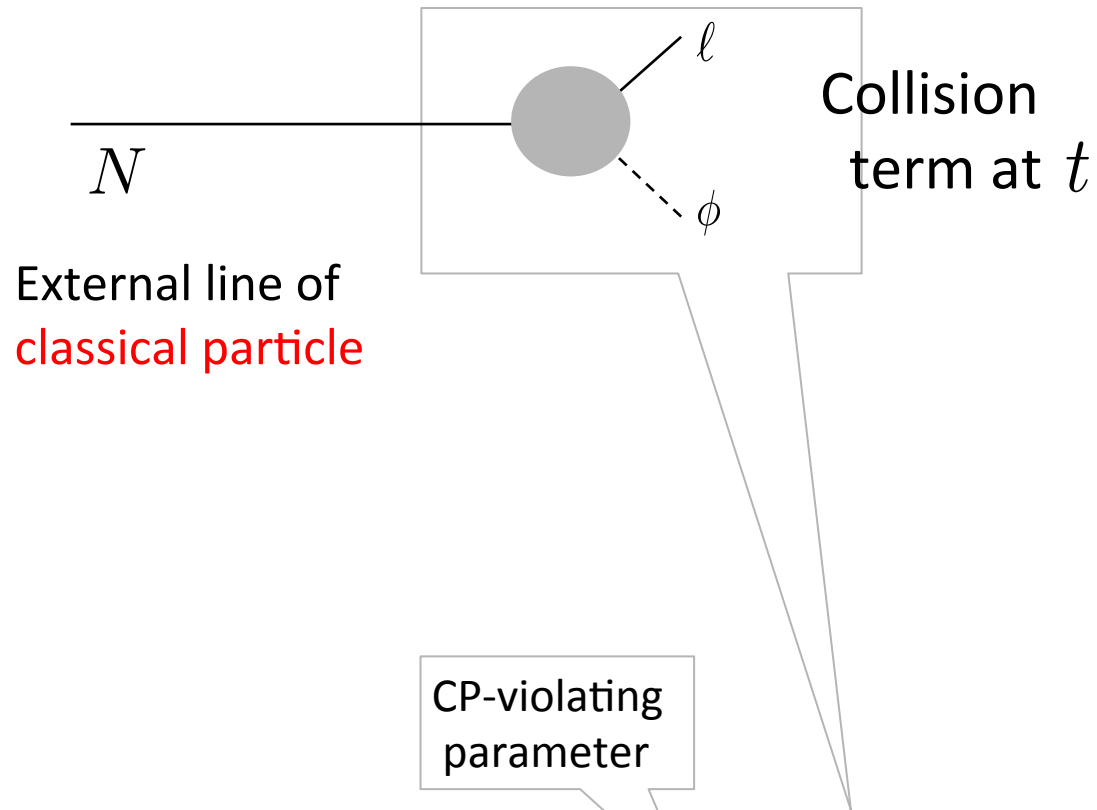
External line of
classical particle

On-shell mass eigenstate

~~Off-shell state~~

~~Quantum superposition of
different mass eigenstates~~

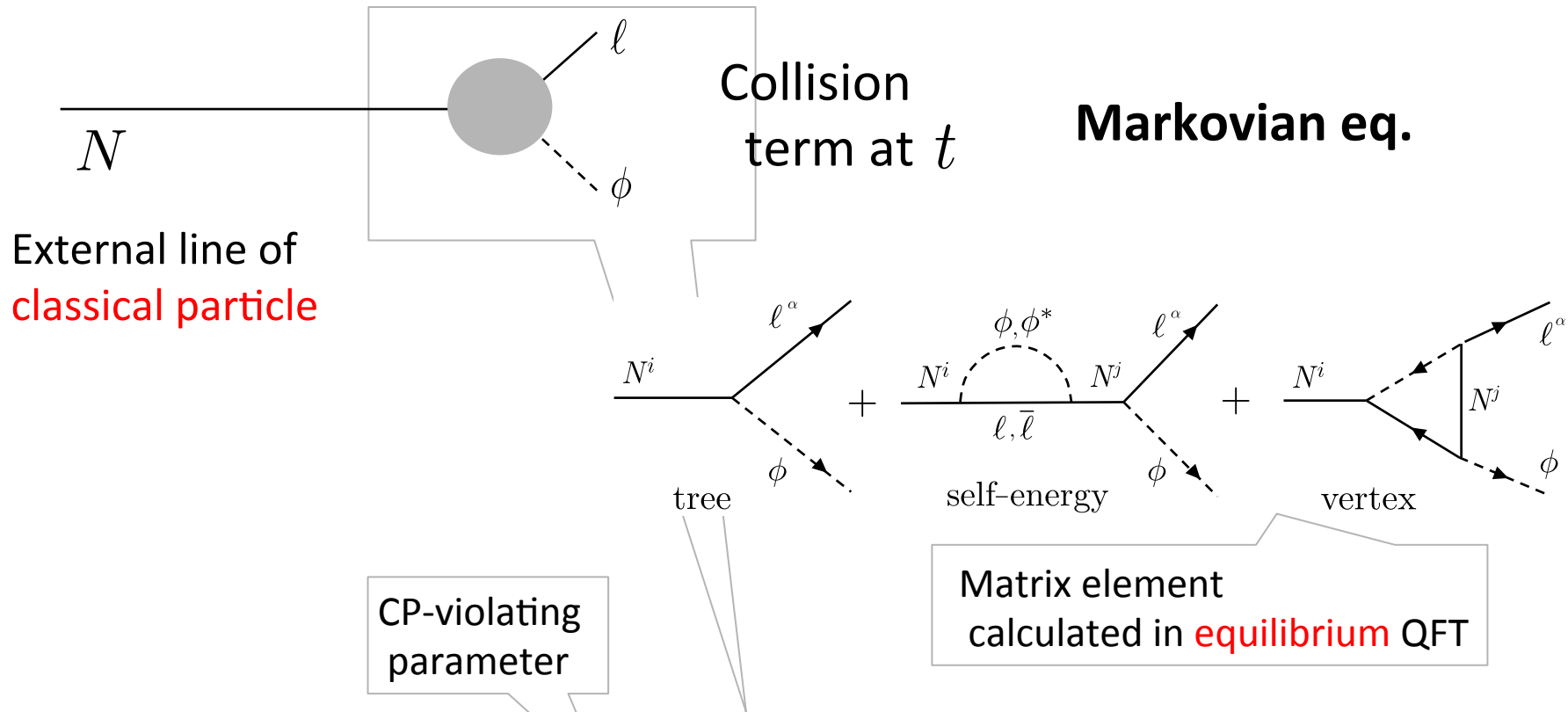
(Classical) Boltzmann equation + equilibrium QFT



$$\frac{dn_L}{dt} + 3Hn_L = \epsilon \times \Gamma \times \delta n_N$$

Boltzmann equation of lepton number (momentum-integrated)

(Classical) Boltzmann equation + equilibrium QFT

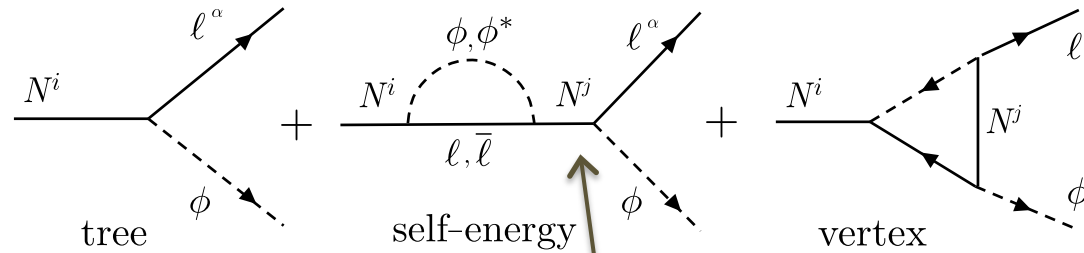


$$\frac{dn_L}{dt} + 3Hn_L = \varepsilon \times \Gamma \times \delta n_N$$

Boltzmann equation of lepton number (momentum-integrated)

3. Resonant leptogenesis

Pilaftsis (97)



(Quantum) Time scale of Flavor oscillation
 $1/\Delta M$

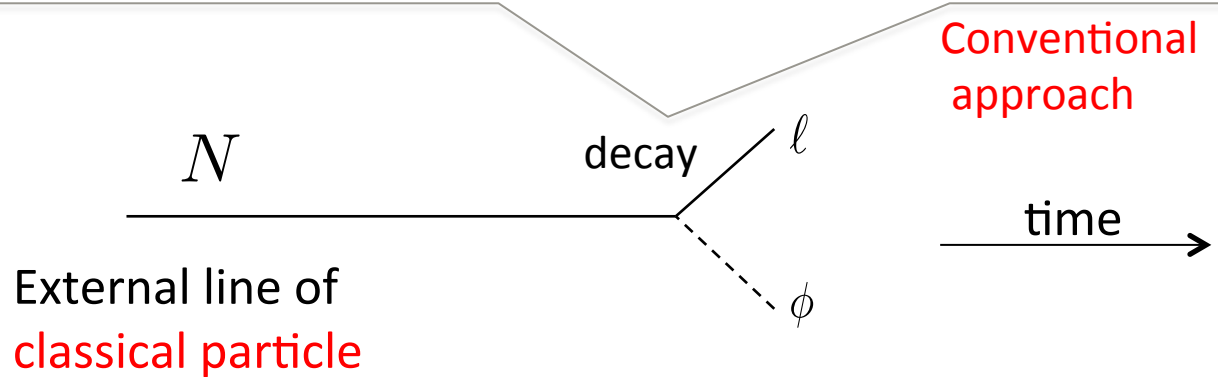
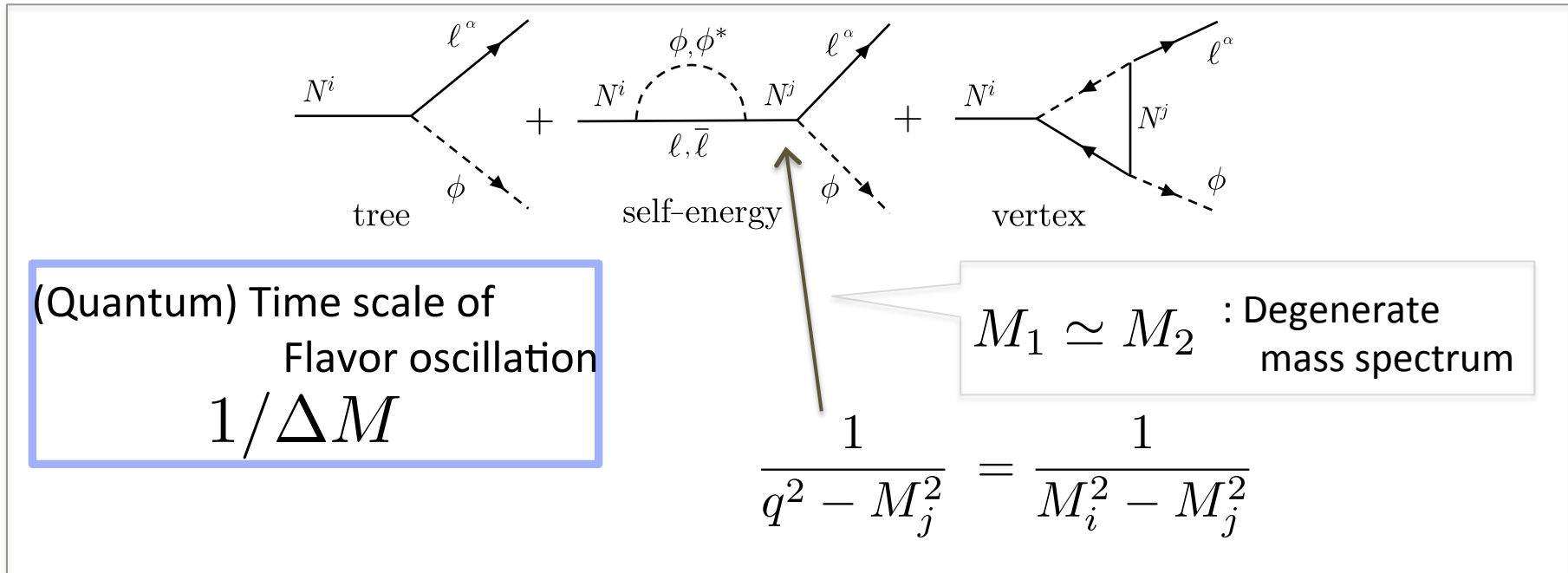
$M_1 \simeq M_2$: Degenerate mass spectrum

$$\frac{1}{q^2 - M_j^2} = \frac{1}{M_i^2 - M_j^2}$$

enhances the CP-violating parameter \mathcal{E}

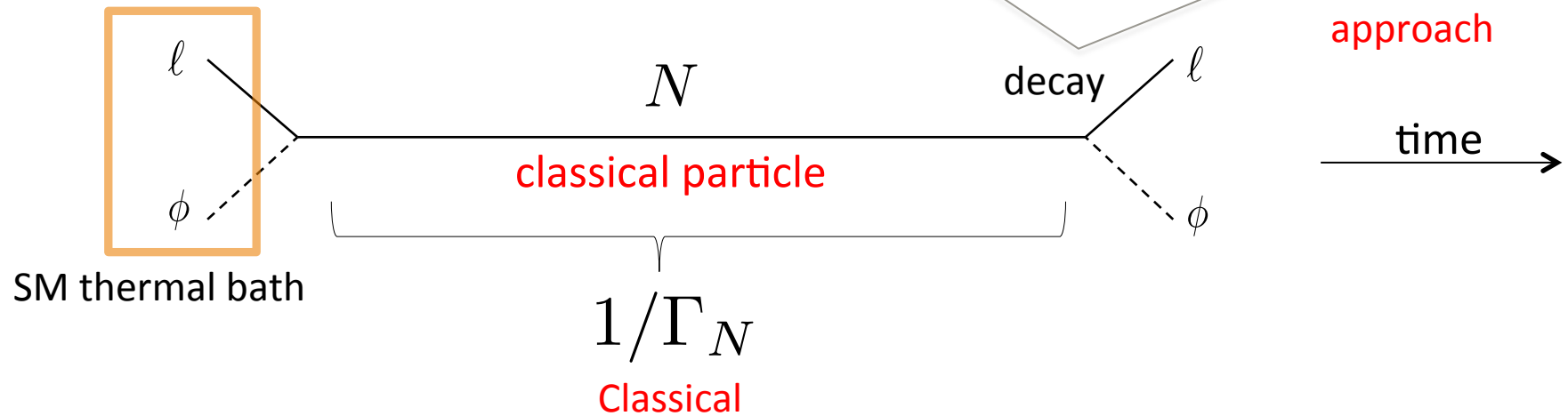
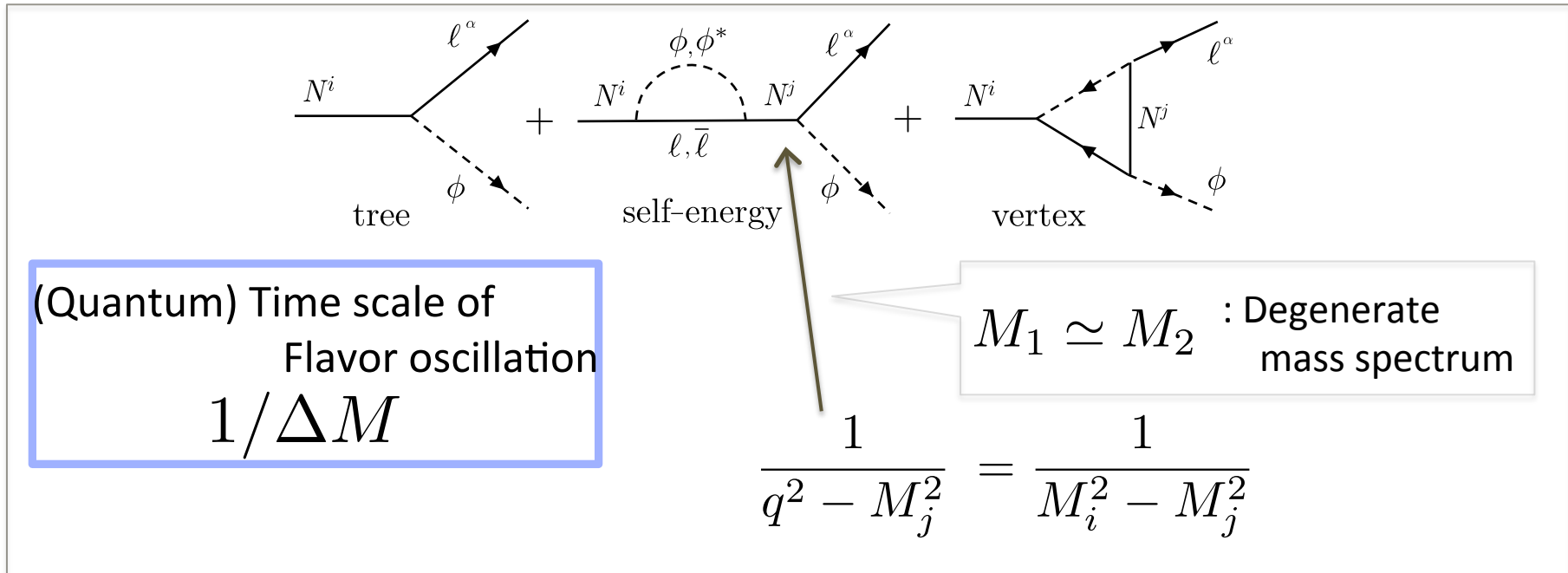
3. Resonant leptogenesis

Pilaftsis (97)



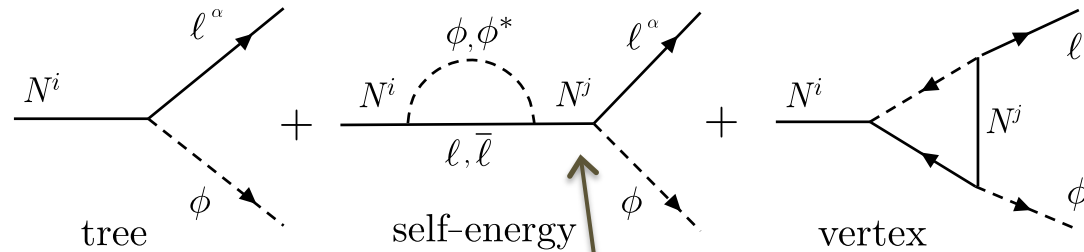
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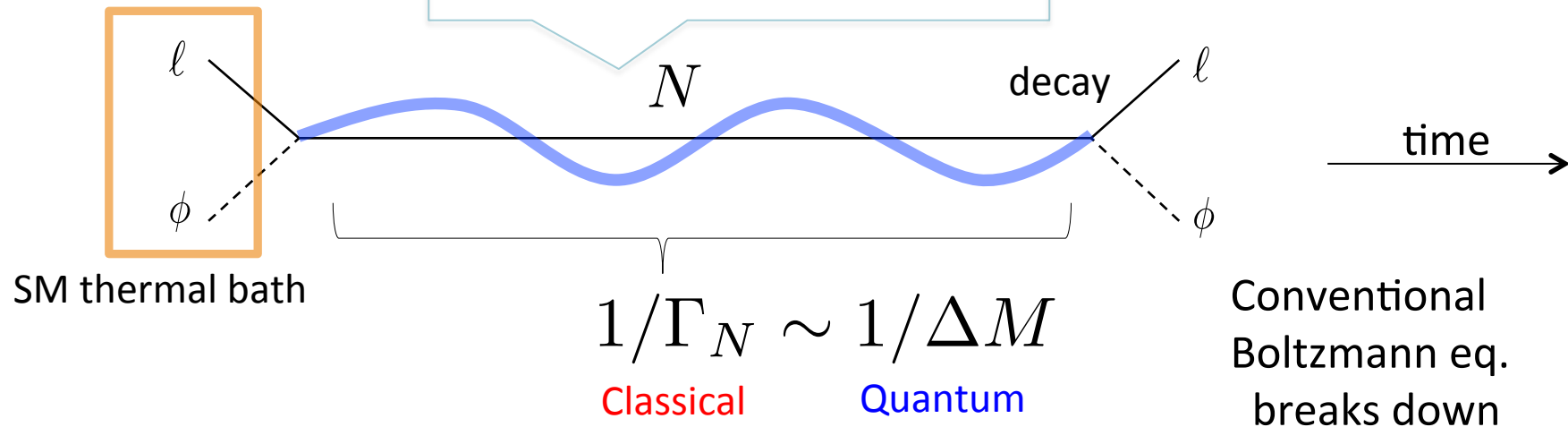


(Quantum) Time scale of Flavor oscillation
 $1/\Delta M$

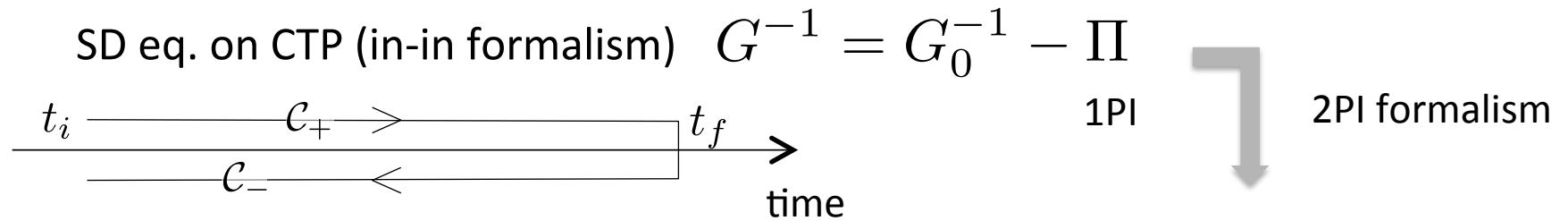
$M_1 \simeq M_2$: Degenerate mass spectrum

$$\frac{1}{q^2 - M_j^2} = \frac{1}{M_i^2 - M_j^2}$$

"External line" cannot be classical

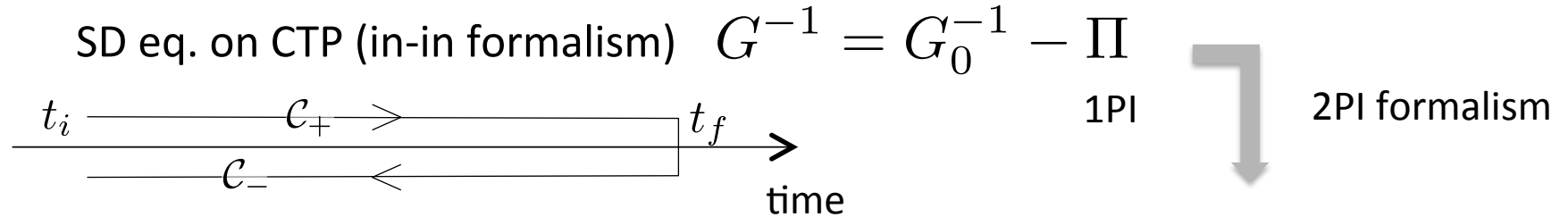


3. Kadanoff-Baym equation



Kadanoff-Baym eq. : Self-consistent equation of **Full propagator**

3. Kadanoff-Baym equation

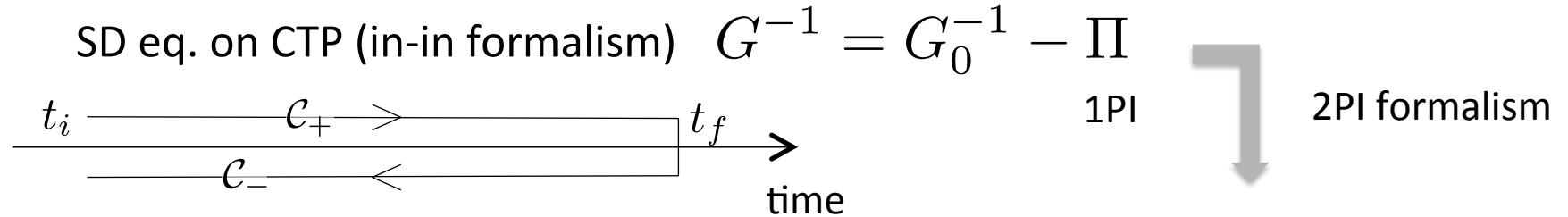


Kadanoff-Baym eq. : Self-consistent equation of **Full propagator**

$$(i\partial_x - M)G_R(x, y) - \int d^4z \Pi_R(x, z)G_R(z, y) = -\delta^4(x - y)$$

$$(i\partial_x - M)G_{\lesssim}(x, y) - \int d^4z \Pi_R(x, z)G_{\lesssim}(z, y) = \int d^4z \Pi_{\lesssim}(x, z)G_A(z, y)$$

3. Kadanoff-Baym equation



Kadanoff-Baym eq. : Self-consistent equation of **Full propagator**

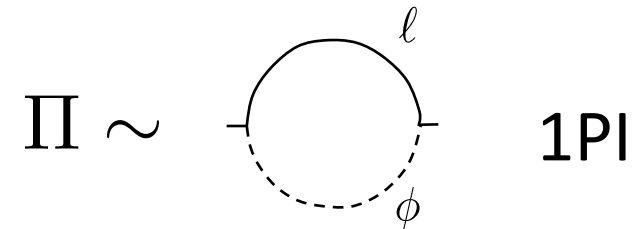
Spectrum of the system

$$(i\partial_x - M)G_R(x, y) - \int d^4z \Pi_R(x, z)G_R(z, y) = -\delta^4(x - y)$$

$$(i\partial_x - M)G_{\lesssim}(x, y) - \int d^4z \Pi_R(x, z)G_{\lesssim}(z, y) = \int d^4z \Pi_{\lesssim}(x, z)G_A(z, y)$$

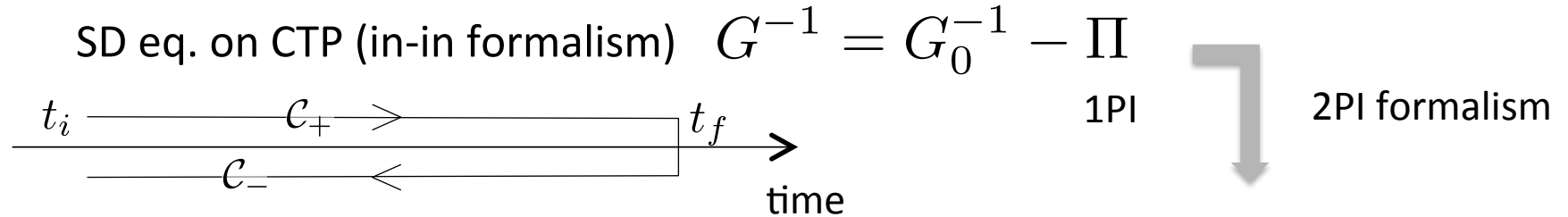
Retarded propagator

$$\begin{aligned} G_R(x, y) &\equiv \theta(x^0 - y^0)G_\rho(x, y) \\ &= \theta(x^0 - y^0)i\langle\{\psi(x), \bar{\psi}(y)\}\rangle \end{aligned}$$



→ Width, Correction to the mass

3. Kadanoff-Baym equation



Kadanoff-Baym eq. : Self-consistent equation of **Full propagator**

Spectrum of the system

1PI self-energy \rightarrow Width, Correction to the mass

$$(i\partial_x - M)G_R(x, y) - \int d^4z \Pi_R(x, z)G_R(z, y) = -\delta^4(x - y)$$

$$(i\partial_x - M)G_{\lesssim}(x, y) - \int d^4z \Pi_R(x, z)G_{\lesssim}(z, y) = \int d^4z \Pi_{\lesssim}(x, z)G_A(z, y)$$

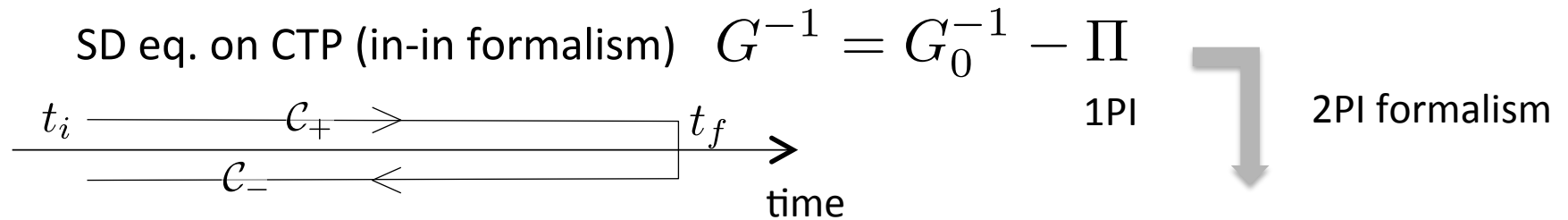
State of the system
 \rightarrow distribution func. f

$$G_{>}(x, y) \equiv \langle \psi(x)\bar{\psi}(y) \rangle$$

Wightman function

$$\hat{a}\hat{a}^\dagger = 1 - \hat{N}$$

3. Kadanoff-Baym equation



Kadanoff-Baym eq. : Self-consistent equation of **Full propagator**

$$\underline{(i\partial_x - M)G_R(x, y)} - \int \underline{d^4z \Pi_R(x, z)G_R(z, y)} = -\delta^4(x - y)$$

$$\underline{(i\partial_x - M)G_{\lesssim}(x, y)} - \int \underline{d^4z \Pi_R(x, z)G_{\lesssim}(z, y)} = \int \underline{d^4z \Pi_{\lesssim}(x, z)G_A(z, y)}$$

integrate

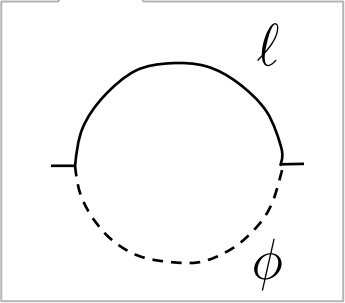
$$G_{\lesssim}(x, y) = - \int d^4u d^4v G_R(x, u) \Pi_{\lesssim}(u, v) G_A(v, y)$$

3. Kadanoff-Baym equation

→ $t \simeq x^0 \simeq y^0 \quad \tau \simeq u^0 \simeq v^0$

$$G_{\leq}(t) = - \int^t d\tau G_R(t, \tau) \Pi_{\leq}(\tau) G_A(\tau, t)$$

$\sim n_N$
Number density at t



$$1/\Gamma_{l,\phi}$$

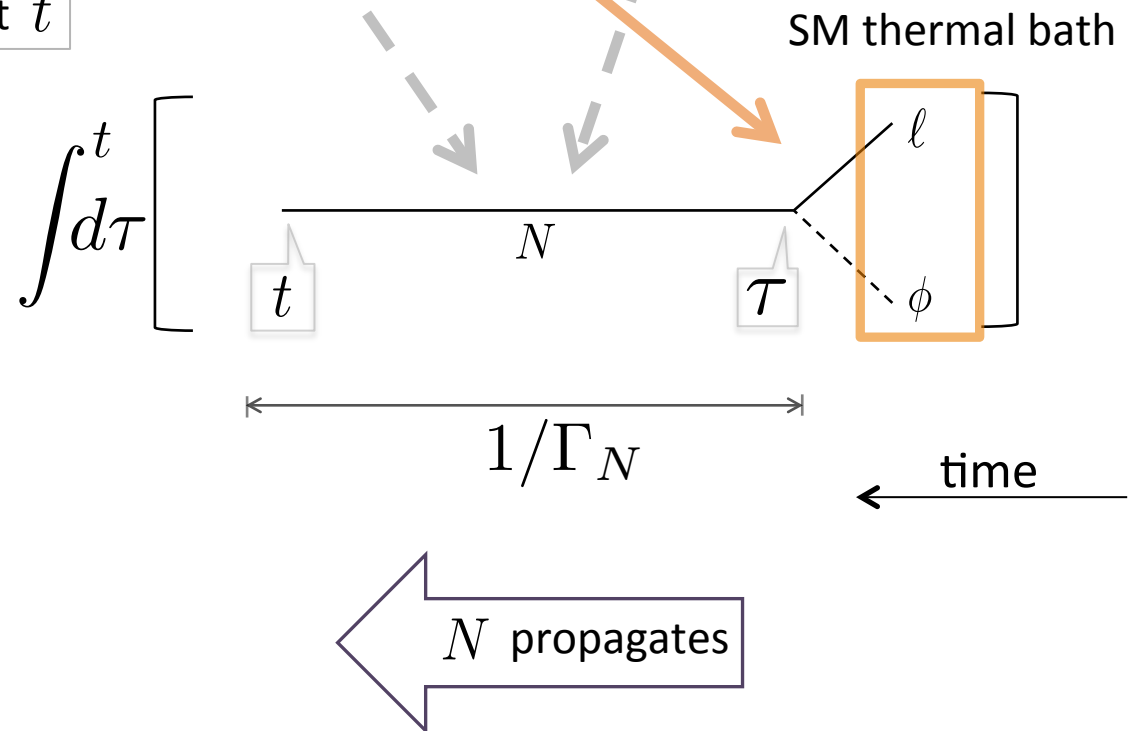
RH neutrino's
Wightman function

3. Kadanoff-Baym equation

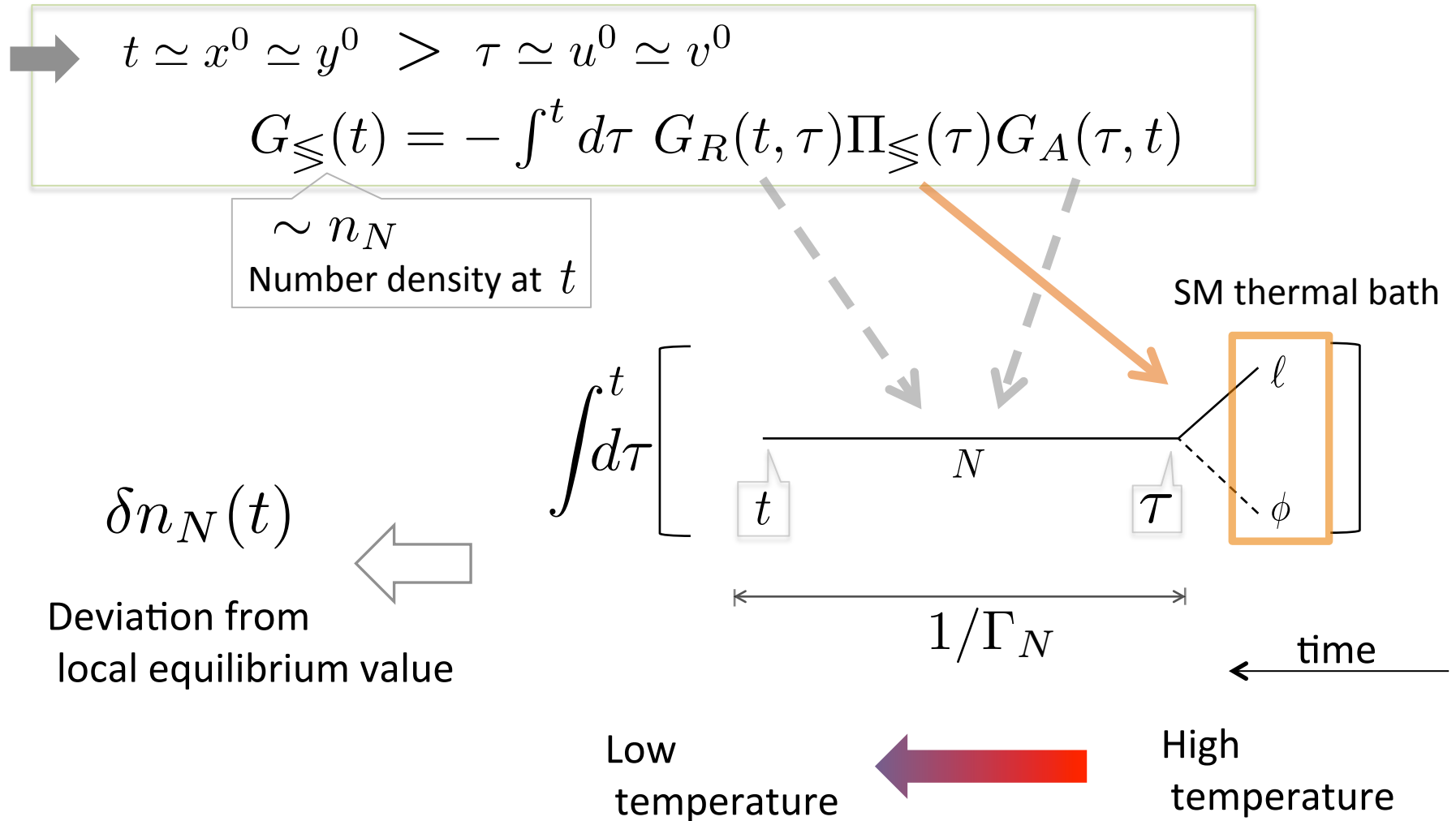
→ $t \simeq x^0 \simeq y^0 > \tau \simeq u^0 \simeq v^0$

$$G_{\leq}(t) = - \int^t d\tau G_R(t, \tau) \Pi_{\leq}(\tau) G_A(\tau, t)$$

$\sim n_N$
Number density at t



3. Kadanoff-Baym equation

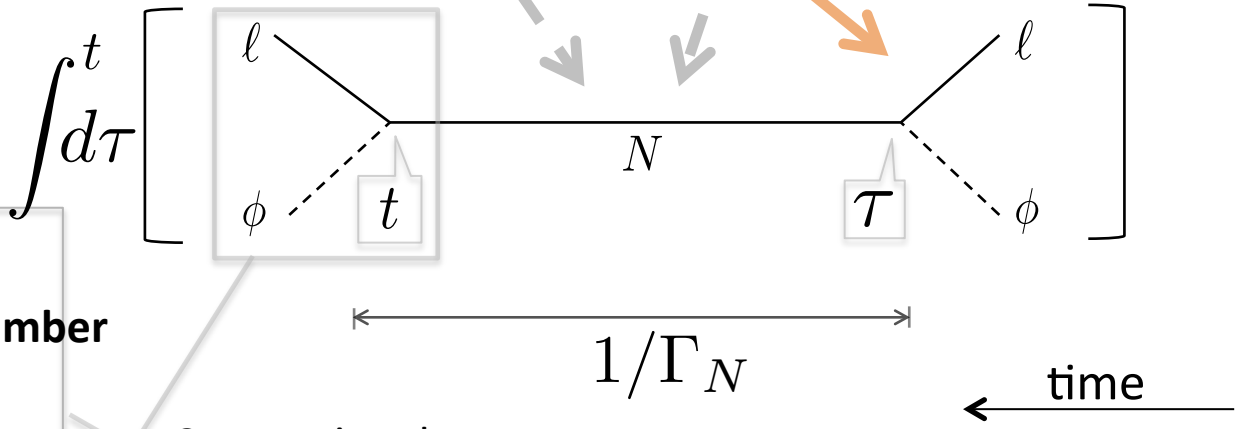


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Number density at t



Collision term in the Boltzmann eq. of **lepton number**

$$\frac{dn_L}{dt} + 3Hn_L = \varepsilon \times \Gamma \times \delta n_N$$

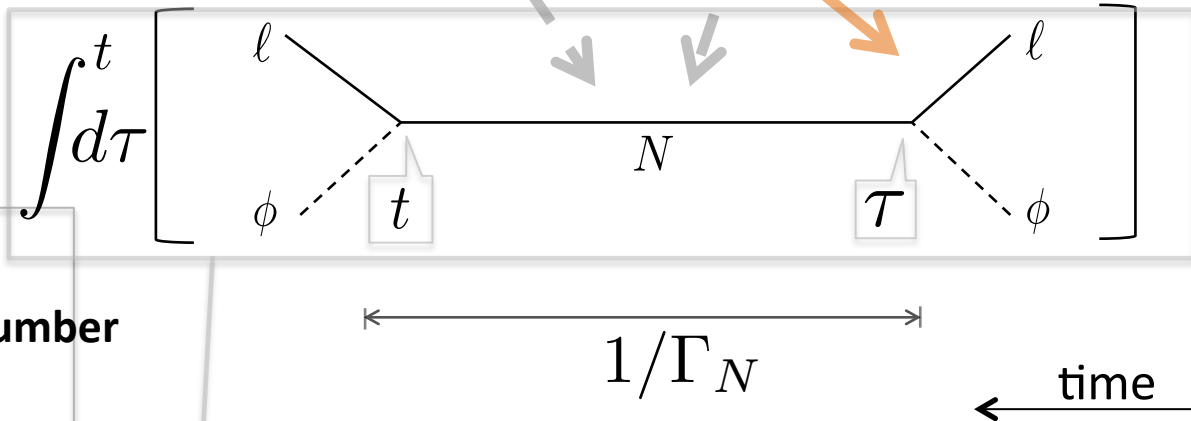
Conventional approach

3. Kadanoff-Baym equation

$t \simeq x^0 \simeq y^0 > \tau \simeq u^0 \simeq v^0$

$$G_{\leq}(t) = - \int^t d\tau G_R(t, \tau) \Pi_{\leq}(\tau) G_A(\tau, t)$$

$\sim n_N$
 Number density at t



Collision term in
 the Boltzmann eq. of **lepton number**

$$\frac{dn_L}{dt} + 3Hn_L = \varepsilon \times \Gamma \times \delta n_N$$

Non-Markovian



“Effective” Markovian eq.

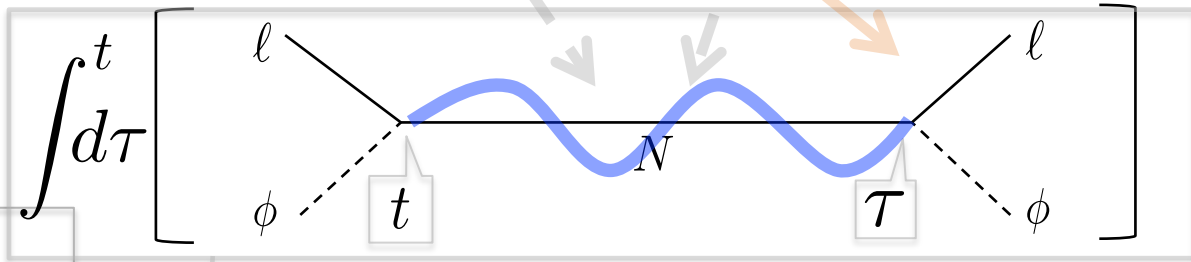
Time integration

3. Kadanoff-Baym equation

→ $t \simeq x^0 \simeq y^0 > \tau \simeq u^0 \simeq v^0$

$$G_{\leq}(t) = - \int^t d\tau G_R(t, \tau) \Pi_{\leq}(\tau) G_A(\tau, t)$$

$\sim n_N$
Number density at t



Collision term in the Boltzmann eq. of **lepton number**

$$\frac{dn_L}{dt} + 3Hn_L = \varepsilon \times \Gamma \times \delta n_N$$

Taking into account full quantum effects,
 $R_{ij} = M_i \Gamma_i + M_j \Gamma_j$
 Garny et al. (2013) KS et al. (2014)

$$\varepsilon_i = \frac{\Im(h^\dagger h)_{ij}^2}{(h^\dagger h)_{ii}(h^\dagger h)_{jj}} \frac{(M_i^2 - M_j^2)M_i \Gamma_j}{(M_i^2 - M_j^2)^2 + R_{ij}^2}$$

cf. equilibrium QFT $R_{ij} = M_i \Gamma_i - M_j \Gamma_j$
 : Plumacher, Buchmuller (98)

Summary

- Degeneracy of time scales makes
the conventional approach invalid.

Classical Boltzmann eq. + equilibrium QFT

- In the case of resonant leptogenesis ($M_1 \simeq M_2$),

$$\underbrace{1/\Gamma_N}_{\text{Lifetime}} \sim \underbrace{1/\Delta M}_{\text{Flavor oscillation}} .$$

- CP-violating parameter is obtained
in the **single** framework (non-equilibrium QFT).
- Kadanoff-Baym eq. (\longrightarrow Density matrix formalism)


Thank you

Density matrix formalism

Sigl, Raffelt (93)

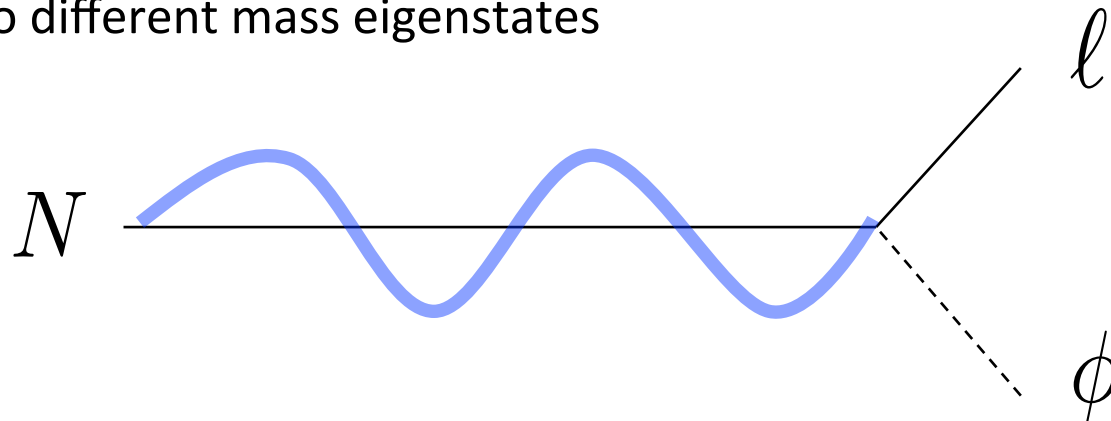
Classical distribution function f

Matrix-valued extension

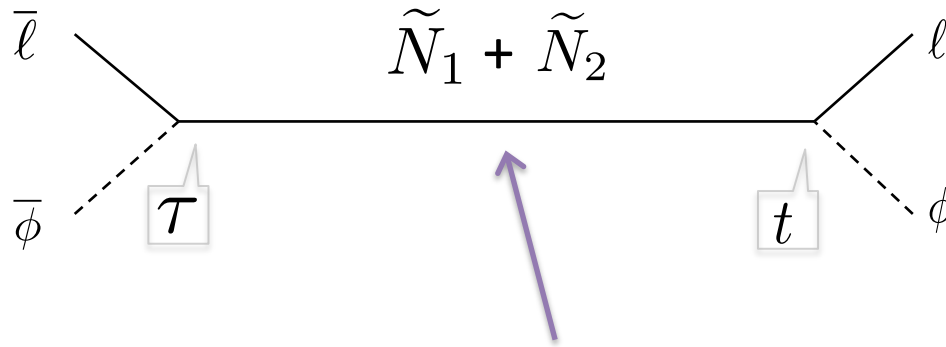

$$\begin{pmatrix} f_{11} & f_{12} \\ f_{21} & f_{22} \end{pmatrix}$$

Quantum extension
of Boltzmann eq.
(Markovian)

Quantum superposition of
two different mass eigenstates



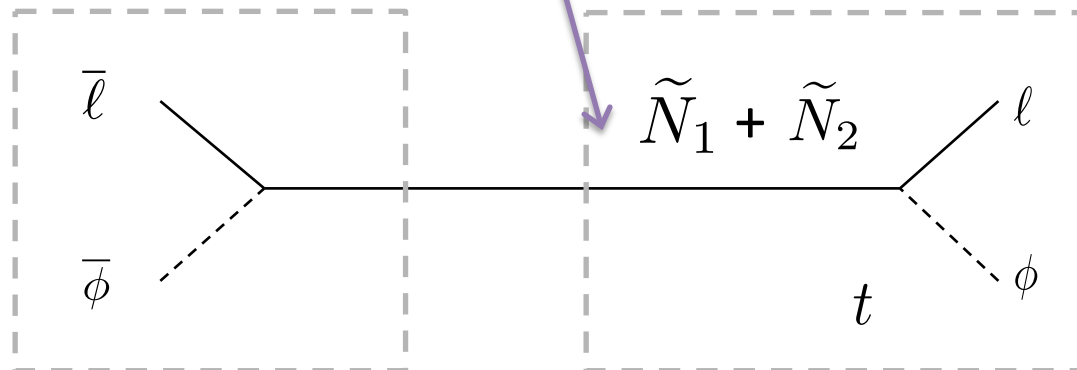
Non-Markovian eq.



Kadanoff-Baym eq.

Quantum superposition
of two different mass eigenstates

Markovian eq.



KB eq. \rightarrow
Density matrix Formalism

Leptogenesis

- Sakharov Conditions

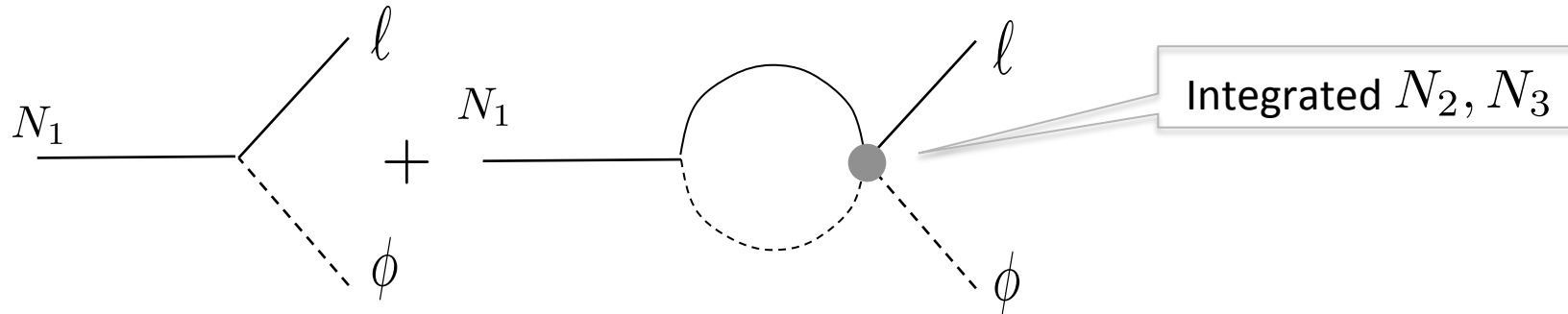
(symmetric initial condition)

	Standard Model	+ Right-handed Neutrino
Baryon number violation	Sphaleron process $B + L$	$\Delta L \rightarrow \Delta B$ Sphaleron
C and CP violation	Weak interaction CKM matrix too small	Yukawa coupling $h_{\alpha i}$
Out of equilibrium dynamics	Strongly 1 st order EW phase transition too small	Large Majorana mass

$$\mathcal{L}_{int} = -h_{\alpha i} \bar{\ell}_{\alpha} \tilde{\phi} N_i - h_{i\alpha}^{\dagger} \bar{N}_i \ell_{\alpha} \tilde{\phi}^{\dagger}$$

Leptogenesis

- Simplest model (decay of the lightest RH neutrino N_1) $M_1 \ll M_{2,3}$



$$\varepsilon \equiv \frac{\Gamma_{N \rightarrow l\phi} - \Gamma_{N \rightarrow \bar{l}\bar{\phi}}}{\Gamma_{N \rightarrow l\phi} + \Gamma_{N \rightarrow \bar{l}\bar{\phi}}} \quad : \text{CP-violating parameter}$$

$$= \sum_{j=2,3} \frac{\text{Im}[(h^\dagger h)_{1j}]^2}{8\pi(h^\dagger h)_{11}} \frac{3M_1}{2M_j} = \frac{3M_1 \text{Im}(h^\dagger m_\nu h^*)_{11}}{16\pi \langle \phi \rangle^2 (h^\dagger h)_{11}}$$

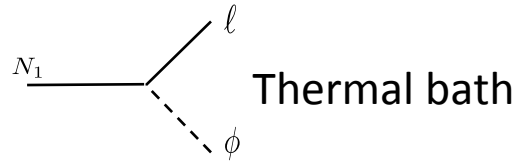
→ Lower bound

$$M_1 \gtrsim 10^9 \text{ GeV}$$

Seesaw

$$m_\nu \simeq \langle \phi \rangle^2 h \frac{1}{M_N} h^T$$

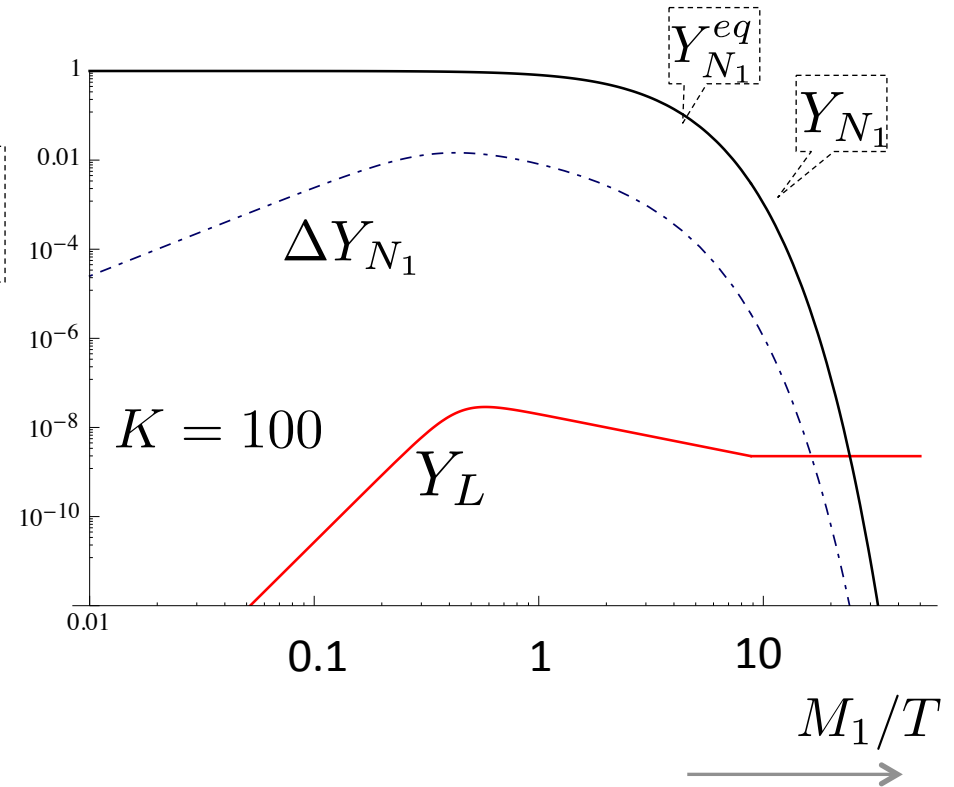
Leptogenesis



Decay width (interaction rate) \rightarrow Relaxation toward equilibrium

$$K \equiv \frac{\Gamma_{N_1}}{H|_{T=M_1}} \gg 1$$

Expanding rate \rightarrow Departure from equilibrium



(Integrated) Boltzmann eq.

$$\frac{dY_{N_1}}{dt} = -D(Y_{N_1} - Y_{N_1}^{eq})$$

+ Thermal initial condition

CP-violating parameter

$$\frac{dY_L}{dt} = \varepsilon D(Y_{N_1} - Y_{N_1}^{eq}) - WY_L$$

RIS-subtraction