# SModelS

&

# Simplified Model Sensitivity to Spin Structure

Ursula Laa LPSC Grenoble & LAPTh Annecy





# Overview

Introduction to Simplified Models

Interpretation of SUSY searches at the LHC and Simplified Models

## SModelS

Constraining generic new physics models using Simplified Model interpretations

Scalar vs Fermionic Top Partners in Models with Dark Matter

Spin dependence in Simplified Model interpretation of searches in  $t\bar{t}$ +MET final states

# Interpretation of SUSY searches in ATLAS and CMS

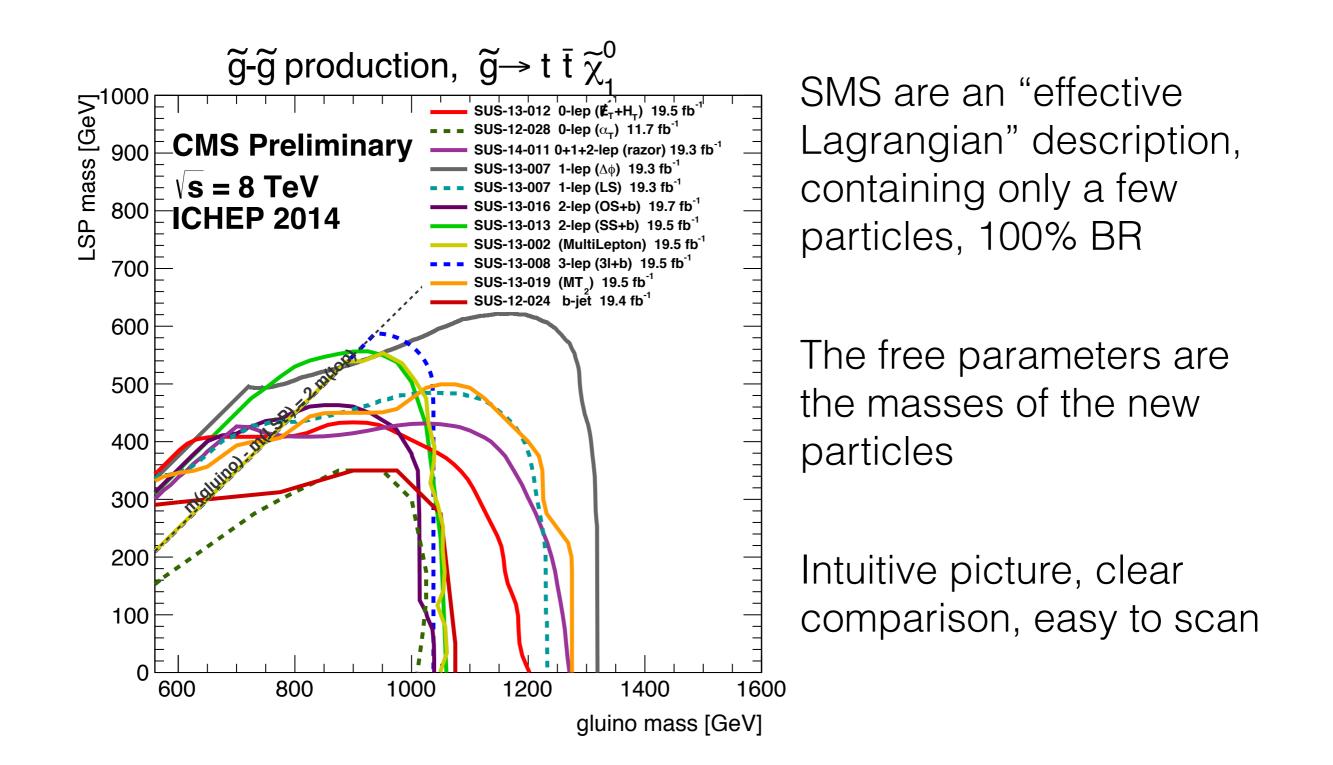
Top down interpretation of SUSY searches, e.g.

- CMSSM predictive but not generic ~4 parameters
- Full MSSM generic but not predictive ~100 parameters

Interpretation that is predictive and model independent?

- Simplified Model Spectra (SMS) based on generic MSSM
- Consider only small subset of new particles
- Simplified model parametrised by particle masses and branching ratios
- Experimental observations constrain the production cross section, assuming 100% BR

# **Simplified Models**



# Using Experimental Results to constrain new physics models

Top down interpretation

- Use event generator + detector simulation to evaluate signal prediction for each parameter point
- Precise, but very time consuming
- Tools: ATOM, CheckMATE, MadAnalysis5

Apply Simplified Model exclusion

- Mass limits only valid within the Simplified Model
- Wrong when considering generic model, e.g. arbitrary gluino decays

Decompose full model into Simplified Model components

- Each component can be compared against experimental upper limits
- Conservative, but fast

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# **Tools to use Simplified Model Results**

### Fastlim

XQCAT

Constraining MSSM parameter space (natural SUSY)

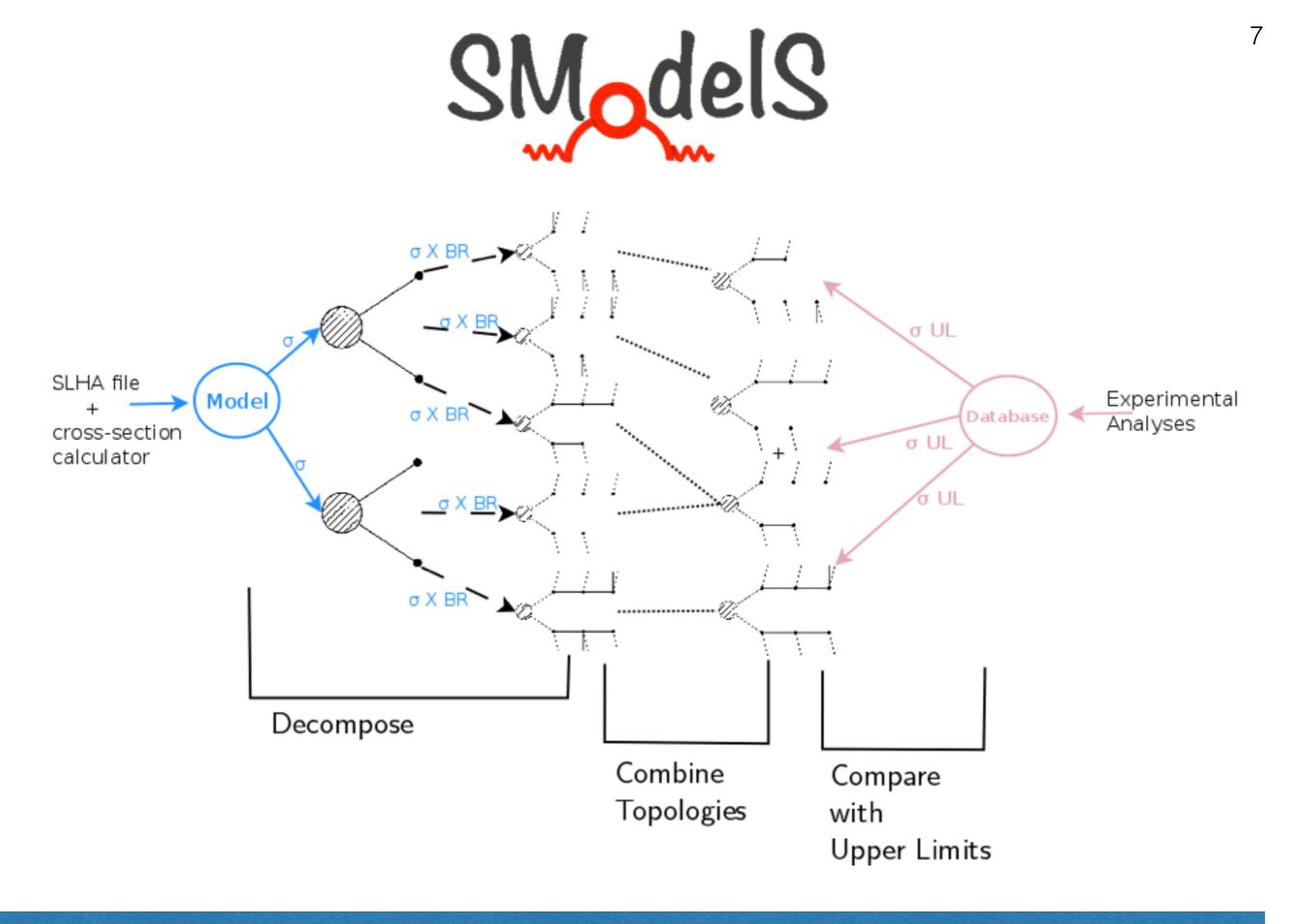
Papucci, Sakurai, Weiler, Zeune (2014)

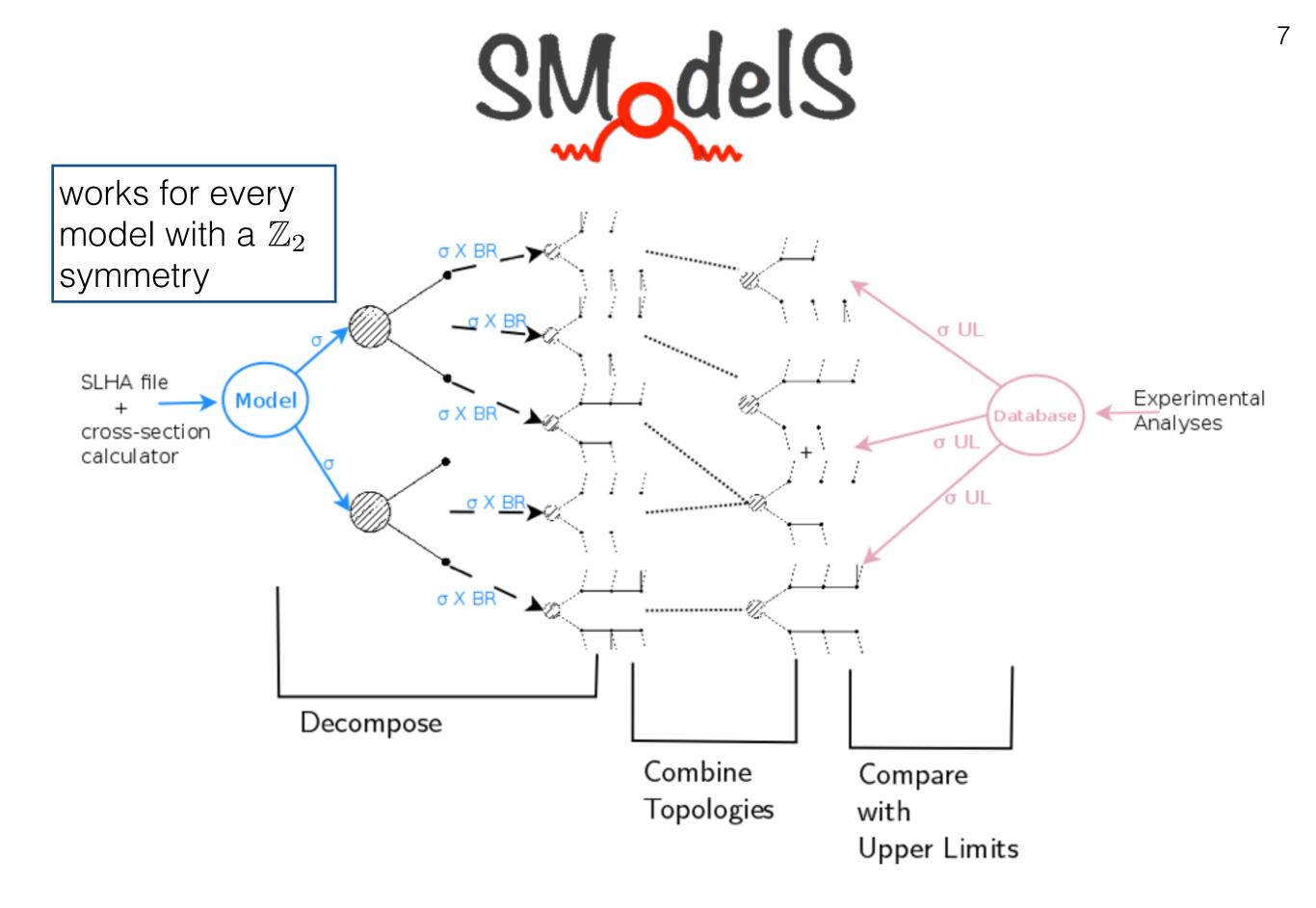
Barducci, Belyaev, Buchkremer, Cacciapaglia Constraining heavy extra quarks et al. (2014)

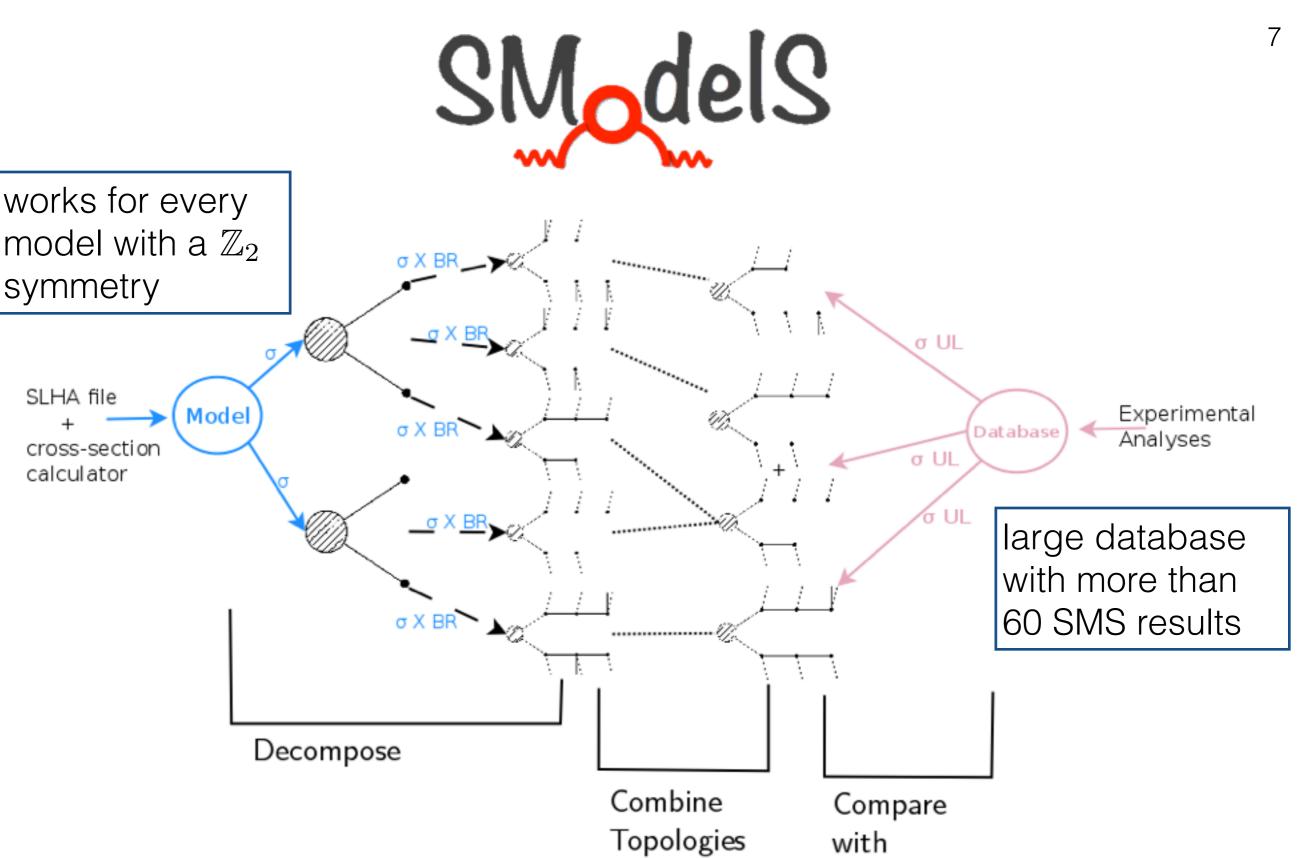
#### SModelS

Constraining generic models with  $\mathbb{Z}_2$  symmetry

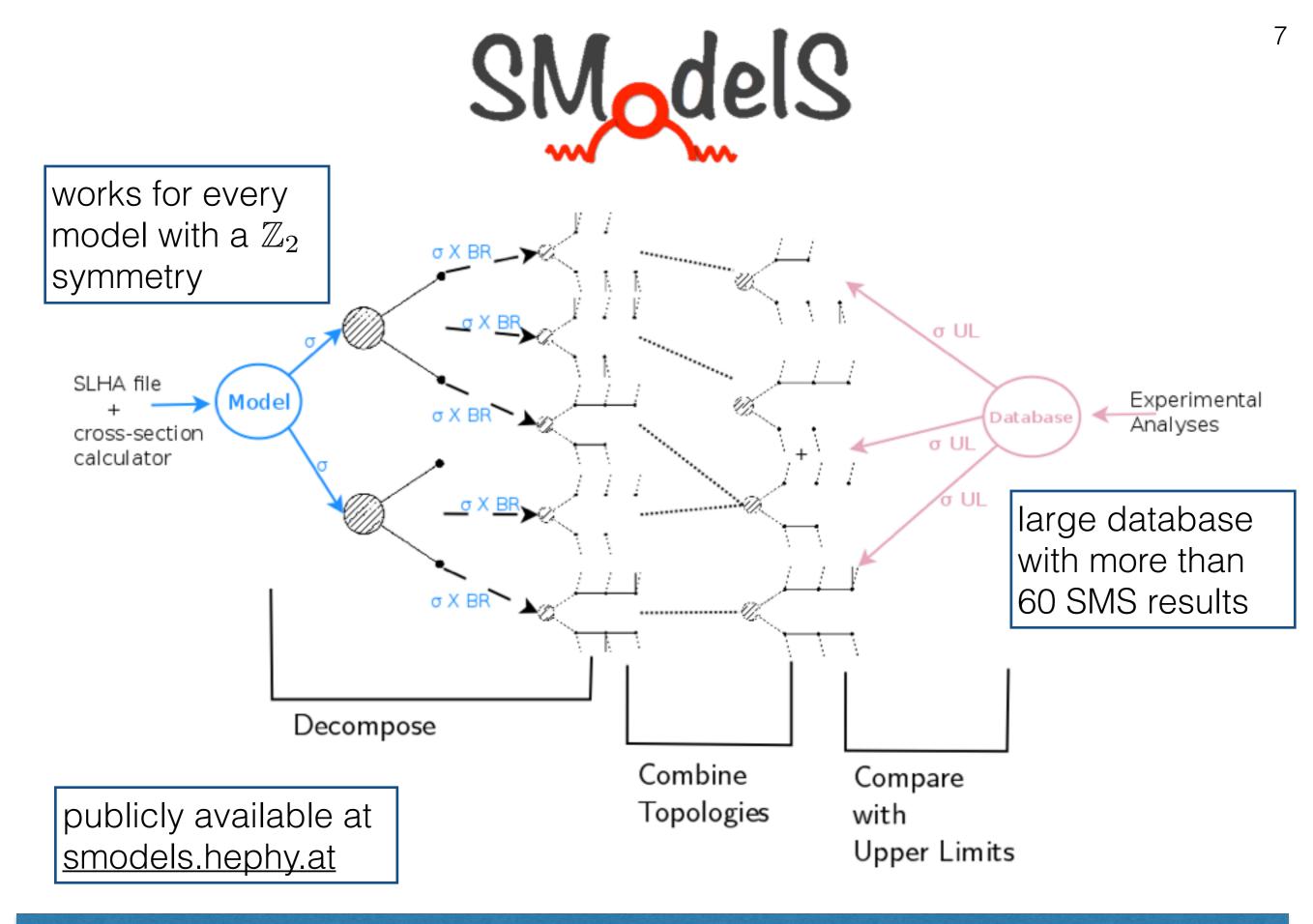
Kraml, Kulkarni, UL, Lessa et al. (2013)







Upper Limits



# **Advantages**

- $\circledast$  No need for event simulation  $\rightarrow$  Fast !
- Fully automated decomposition into SMS components, matching to results in the database
- Works "out of the box" also for non-MSSM scenarios
- Large database, easy to add new results
   Identifies also "missing topologies"

# Limitations

- Limited to short cascade decays
- Mostly covers symmetric decays (equal branches)
- Relies on assumption that signal efficiencies are largely determined by the mass structure, neglect dependencies on e.g.:
  - Production channel
  - Spin correlations
  - Off-shell states in the decays

for previous studies of these assumptions see e.g. arXiv:1410.0965, arXiv:1501.03942, arXiv:1503.02960

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Here: study effects of spin correlations for scenarios where dark matter couples via a top partner

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# Scalar versus Fermionic Top Partner Interpretation of tt+MET searches at the LHC

S. Kraml, UL, L. Panizzi, H. Prager (in preparation)

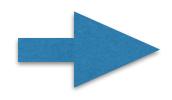
# The SUSY scenario

Interaction details depends on both stop and neutralino mixing!  $\mathcal{L}_{t\tilde{t}\tilde{\chi}^{0}} = g \,\bar{t} \left( f_{Lk} P_{R} + h_{Lk} P_{L} \right) \tilde{\chi}_{k}^{0} \,\tilde{t}_{L} + g \,\bar{t} \left( h_{Rk} P_{R} + f_{Rk} P_{L} \right) \tilde{\chi}_{k}^{0} \,\tilde{t}_{R} + \text{h.c.}$   $= g \,\bar{t} \left( a_{ik}^{\tilde{t}} P_{R} + b_{in}^{\tilde{t}} P_{L} \right) \tilde{\chi}_{k}^{0} \,\tilde{t}_{i} + \text{h.c.}$ 

> Write stop mixing as  $\begin{pmatrix} \tilde{t}_1 \\ \tilde{t}_2 \end{pmatrix} = R \begin{pmatrix} \tilde{t}_L \\ \tilde{t}_R \end{pmatrix}, \quad R = \begin{pmatrix} \cos \theta_{\tilde{t}} \sin \theta_{\tilde{t}} \\ -\sin \theta_{\tilde{t}} \cos \theta_{\tilde{t}} \end{pmatrix}$

Then for Bino-like LSP (neglecting wino, higgsino components)

$$\mathcal{L}_{t\tilde{t}_1\tilde{\chi}_1^0} \approx -\frac{g}{3\sqrt{2}} \tan\theta_W N_{11} \,\bar{t} \left(\cos\theta_{\tilde{t}} P_R - 4\sin\theta_{\tilde{t}} P_L\right) \tilde{\chi}_1^0 \,\tilde{t}_1 + \text{h.c.}$$



stop helicity reflected in top polarisation

see arXiv:1212.3526

# The SUSY scenario

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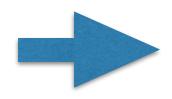
$$\mathcal{L}_{t\tilde{t}\tilde{\chi}^{0}} = g \,\bar{t} \left( f_{Lk} P_{R} + h_{Lk} P_{L} \right) \tilde{\chi}_{k}^{0} \,\tilde{t}_{L} + g \,\bar{t} \left( h_{Rk} P_{R} + f_{Rk} P_{L} \right) \tilde{\chi}_{k}^{0} \,\tilde{t}_{R} + \text{h.c.}$$
  
$$= g \,\bar{t} \left( a_{ik}^{\tilde{t}} P_{R} + b_{in}^{\tilde{t}} P_{L} \right) \tilde{\chi}_{k}^{0} \,\tilde{t}_{i} + \text{h.c.}$$
 neutralino mixing dependence

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$$= g \, \bar{t} \begin{pmatrix} a_{ik}^{t} P_{R} + b_{in}^{t} P_{L} \end{pmatrix} \tilde{\chi}_{k}^{0} \, \tilde{t}_{i} + \text{h.c.}$$

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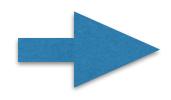
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stop helicity reflected in top polarisation

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see arXiv:1212.3526

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# The extra quark scenario

Consider extra heavy quark (XQ) that is odd under new  $\mathbb{Z}_2$  symmetry, mediating interactions between 3rd generation quarks and DM

Dark matter singlet coupling to XQ singlet or doublet  $\Psi_{1/6} = {T \choose B}$ 

$$\mathcal{L}_{1}^{S} = \left[ \lambda_{11}^{t} \overline{T} P_{R} t + \lambda_{11}^{b} \overline{B} P_{R} b + \lambda_{21} \overline{\Psi}_{1/6} P_{L} \begin{pmatrix} t \\ b \end{pmatrix} \right] \underbrace{\mathcal{S}_{DM}^{0}}_{\text{DM}} + \text{h.c.} \quad \text{scalar dark matter}$$
$$\mathcal{L}_{1}^{V} = \left[ g_{11}^{t} \overline{T} \gamma_{\mu} P_{R} t + g_{11}^{b} \overline{B} \gamma_{\mu} P_{R} b + g_{21} \overline{\Psi}_{1/6} \gamma_{\mu} P_{L} \begin{pmatrix} t \\ b \end{pmatrix} \right] \underbrace{\mathcal{V}_{DM}^{0\mu}}_{\text{DM}} + \text{h.c.}$$

Vector-like XQ: couplings either purely left or purely right
 Chiral XQ: can couple to both left and right SM quarks

For direct comparison, here we consider the chiral XQ scenario

# Setup

Select 4 benchmarks to cover different polarisations and mass

BP	m <sub>T</sub>	m <sub>DN</sub>	d top polaris	sation
(600,10	L 600 G	eV 10 Ge	eV left-hand	ded
(600,10	R 600 G	eV 10 Ge	eV right-har	nded
(600,300	)L 600 G	eV 300 G	ieV left-hand	ded
(600,300	)R 600 G	eV 300 G	ieV right-har	nded

- For each benchmark consider 3 scenarios with equivalent left/ right couplings
  - Stop production, Bino LSP (SUSY)

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- Heavy top production, scalar dark matter (XQ SDM)
- Heavy top production, vector dark matter (XQ VDM)
- In addition: scan in mass vs mass plane to obtain 2D exclusion in each scenario

# **LHC Searches**

Consider LHC searches at 8 TeV, in tt+MET final states
 Interpretation:

$$pp \to \tilde{t}_1 \tilde{t}_1^* \to t \bar{t} \tilde{\chi}_1^0 \tilde{\chi}_1^0$$

Searches can be classified by final state

- Fully hadronic final state
  - ATLAS-CONF-2013-024 (implemented in CheckMATE)
- Single lepton final state
  - ATLAS-SUSY-2013-15 (implemented in CheckMATE)
  - CMS-SUS-13-011 (implemented in MA5)
- Dilepton final state

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ATLAS-SUSY-2013-19 (implemented in CheckMATE)

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Finally consider also effects in generic jets+MET search

ATLAS-SUSY-2013-02 (implemented in MA5)

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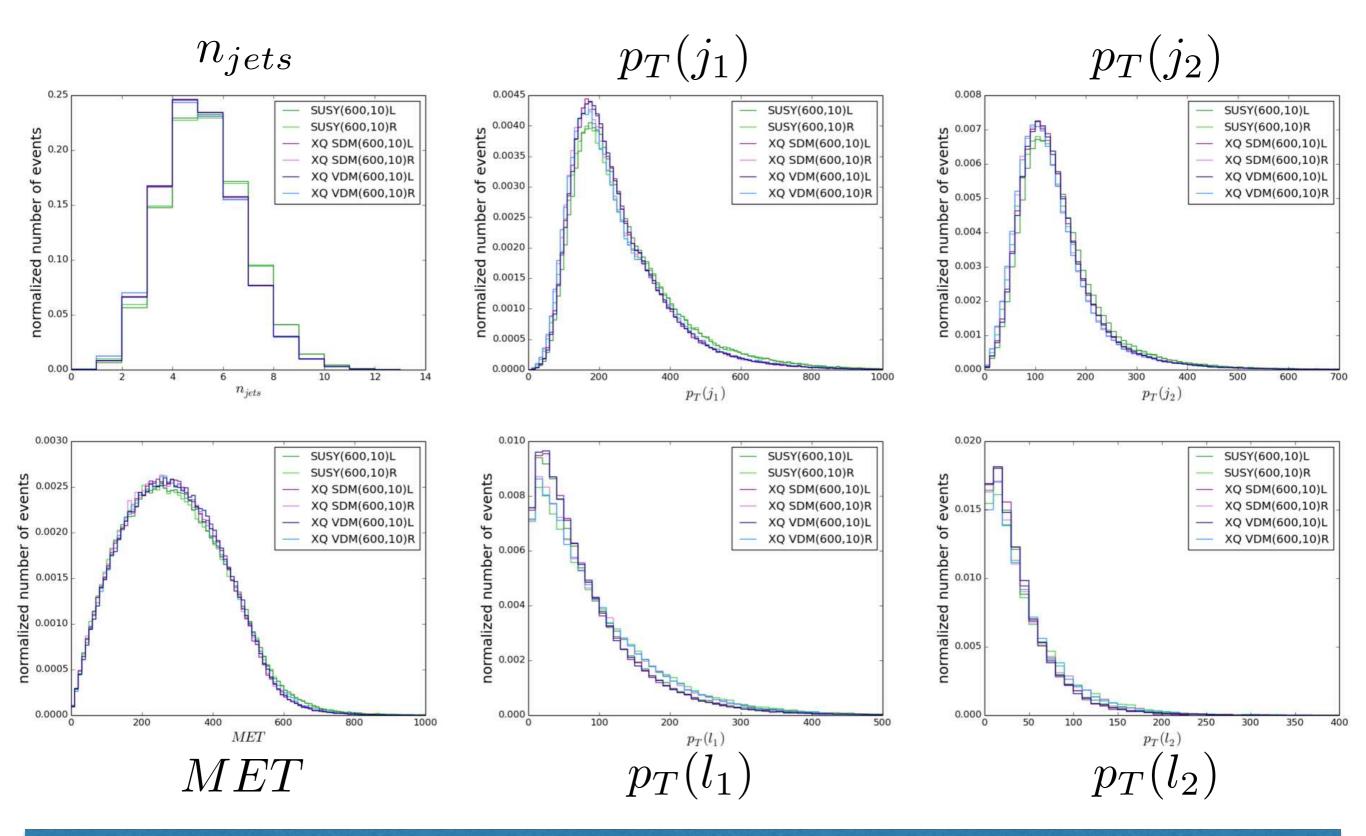
# **Technical Setup**

XQ simplified model implemented in Feynrules
 Use default MSSM implementation for SUSY scenario
 Simulate 200k events in MG (at LO) for

 $pp \to t \bar{t} \text{ DM DM} \to (W^+ b)(W^- \bar{b}) \text{ DM DM}$ 

- W decay + showering performed in Pythia6
- Detector simulation using Delphes3
- Implementations in CheckMATE and MadAnalysis
- Stop production cross sections are calculated at NLO+NLL, XQ cross sections at NLO+NLL

## **Generator Level Distributions**



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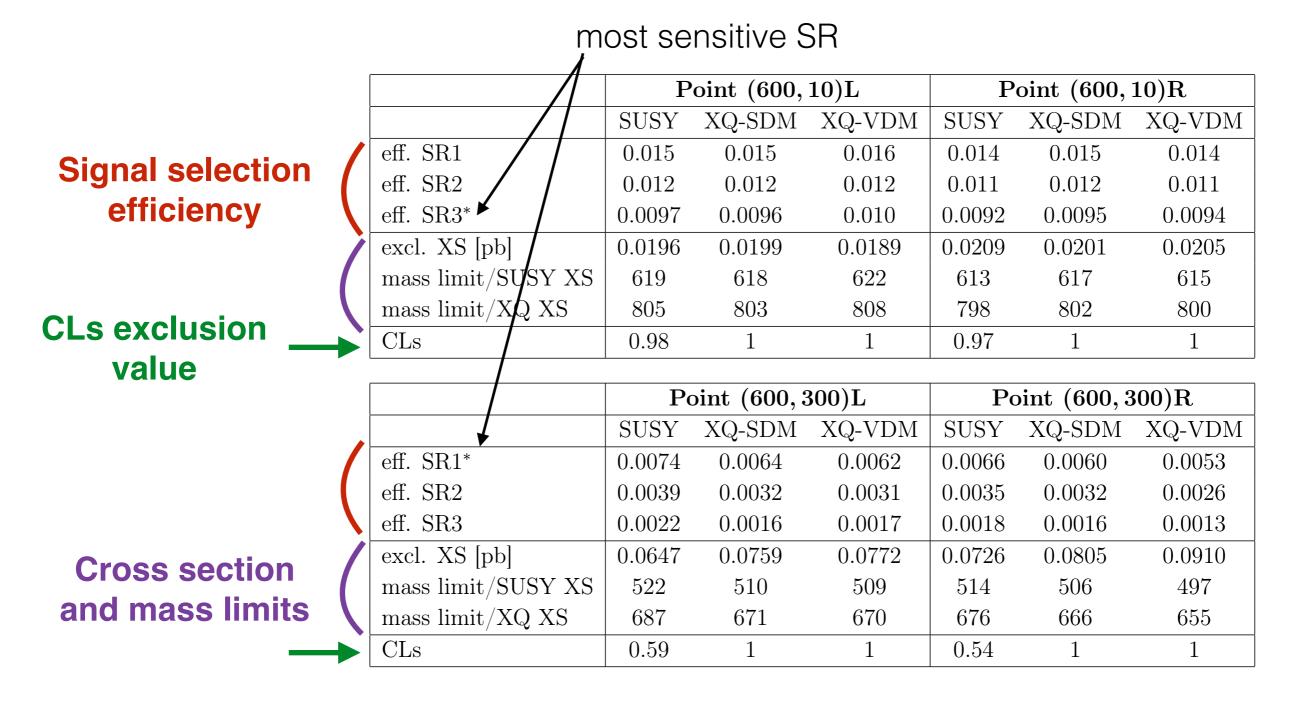
# **Generator Level Distributions**

- Shapes very similar
- MET, jet p⊤ slightly harder in SUSY scenarios
- Helicity dependence can be largest difference, e.g.
   for lepton p<sub>T</sub>
- Somewhat higher jet activity in SUSY scenarios
- Similar picture for benchmarks with small mass gap (600,300)

### → Expect similar efficiencies for all scenarios

# **Results Fully Hadronic Final State**

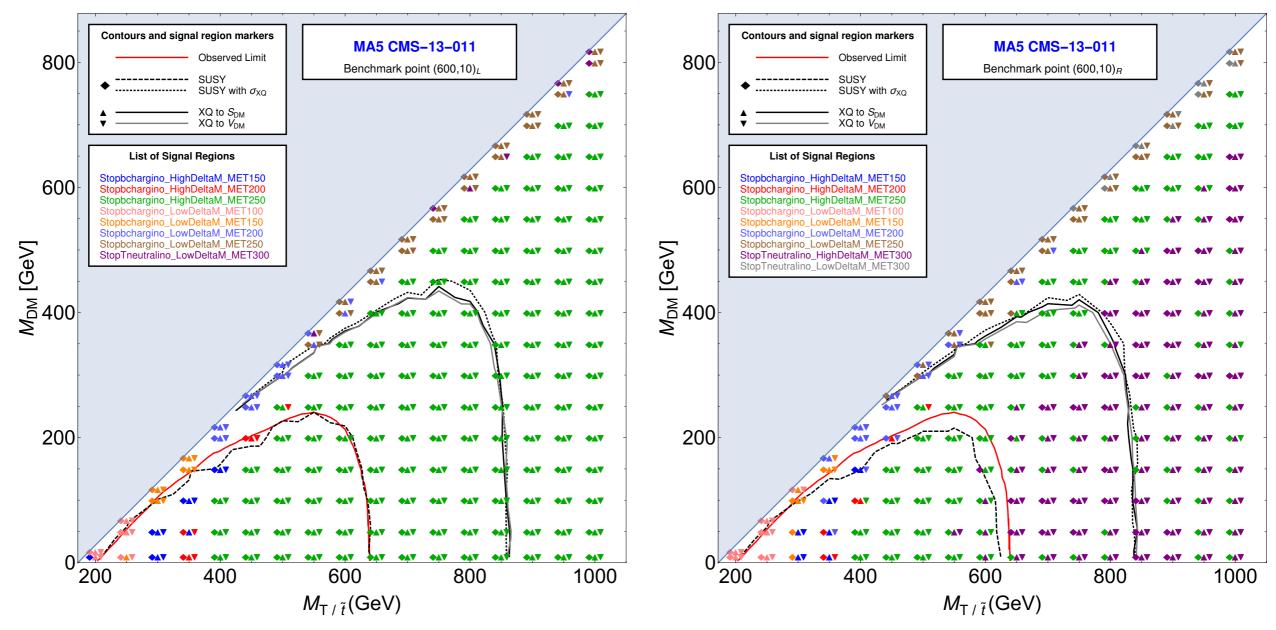
#### ATLAS-CONF-2013-024, implemented in CheckMATE



# **Exclusion in the Mass Plane**

#### left-handed couplings

right-handed couplings

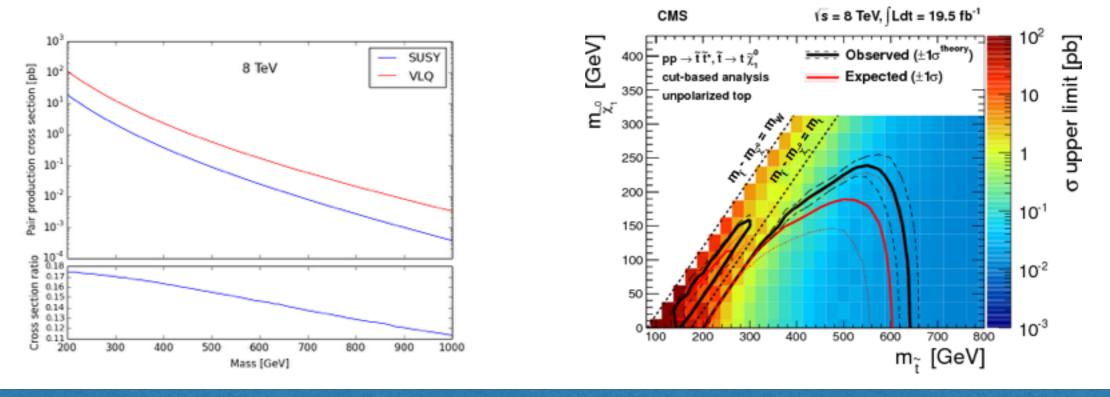


Good agreement between exclusion of XQ with scalar or vector dark matter and when rescaling SUSY efficiencies with XQ cross sections

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# Conclusions

- ♦ Kinematic distributions, efficiencies and exclusions are similar in all 3 scenarios → SUSY SMS cross section limits can be used to constrain XQ scenarios (when narrow width approximation holds)
- Higher production cross section for XQs → useful to consider larger mass range in interpretation
- Correlate observed number of events and effective mass scale to distinguish origin of possible signal for detailed study see arXiv:1207.4794



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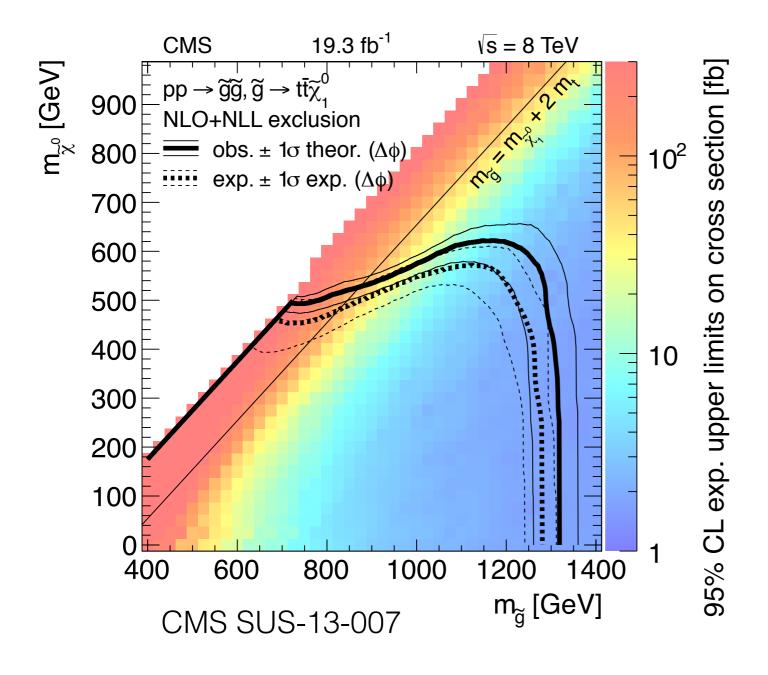
# Outlook – SModelS v1.1

- Implementation of first 13 TeV results in the SModelS database
- Addition of "homegrown" efficiency maps, obtained using CheckMATE or MadAnalysis5 → can improve coverage by adding efficiency maps for additional topologies
- Release envisaged (end of) this summer

sign up for updates on <u>smodels.hephy.at</u>



# Using Simplified Model results to constrain generic models



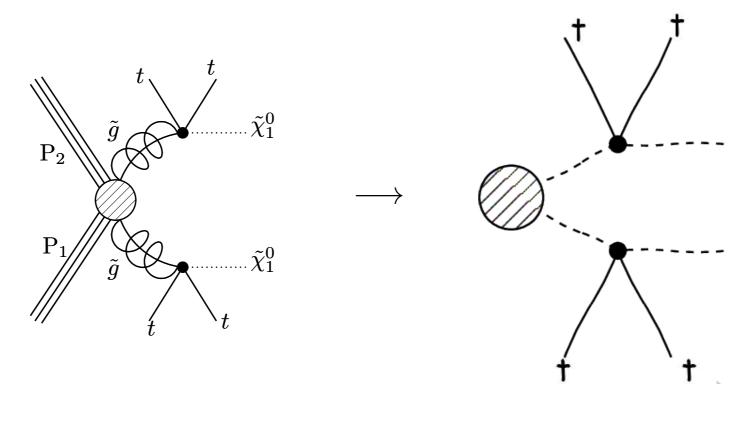
To test realistic models, use upper limits on  $\sigma \times BR$ (exclusion line only valid in the simplified model)

> Assumption: upper limits on  $\sigma \times BR$ are mainly a function of the masses of the new particles

> > other quantum numbers may be neglected in first approximation

# **SModelS Topology Description**

Describe topology by vertex structure and outgoing SM particles in each vertex





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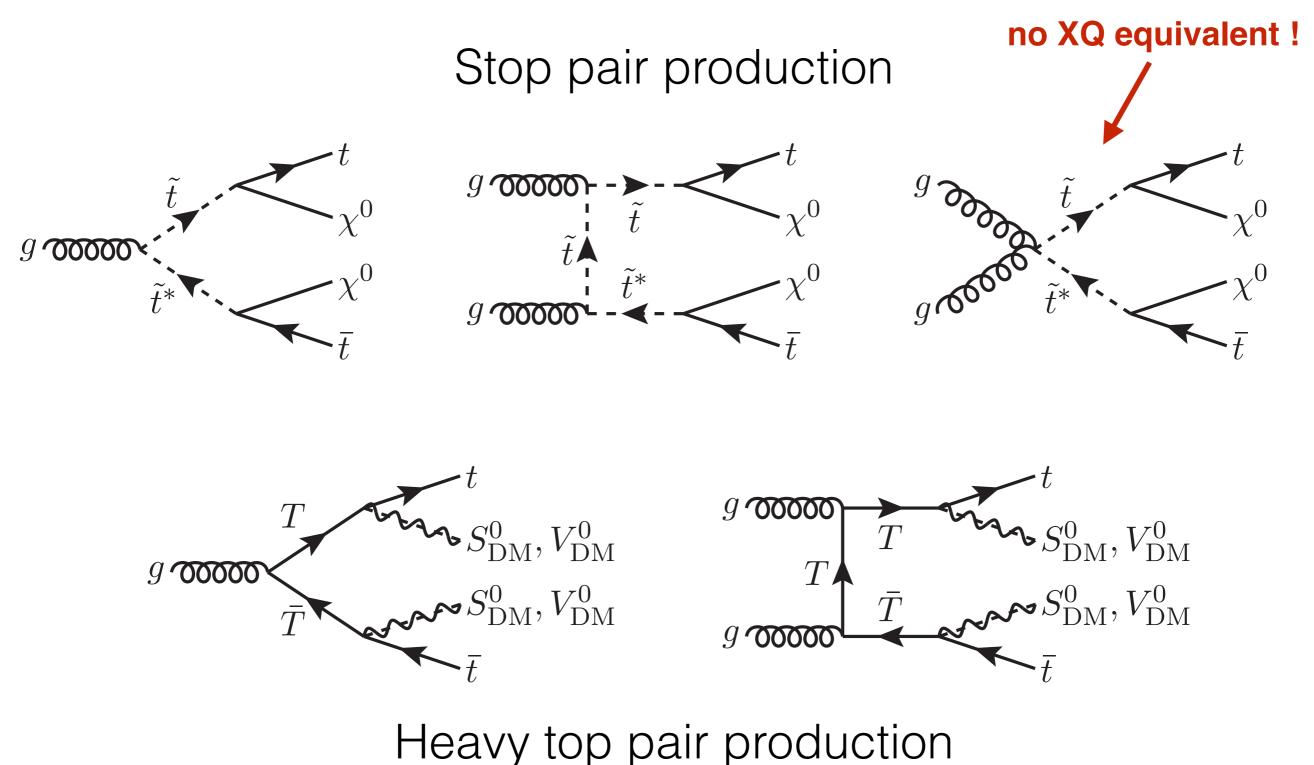
#### [[[t,t]],[[t,t]]]

Discard detailed information about the BSM states, only their masses are considered when evaluating constraints

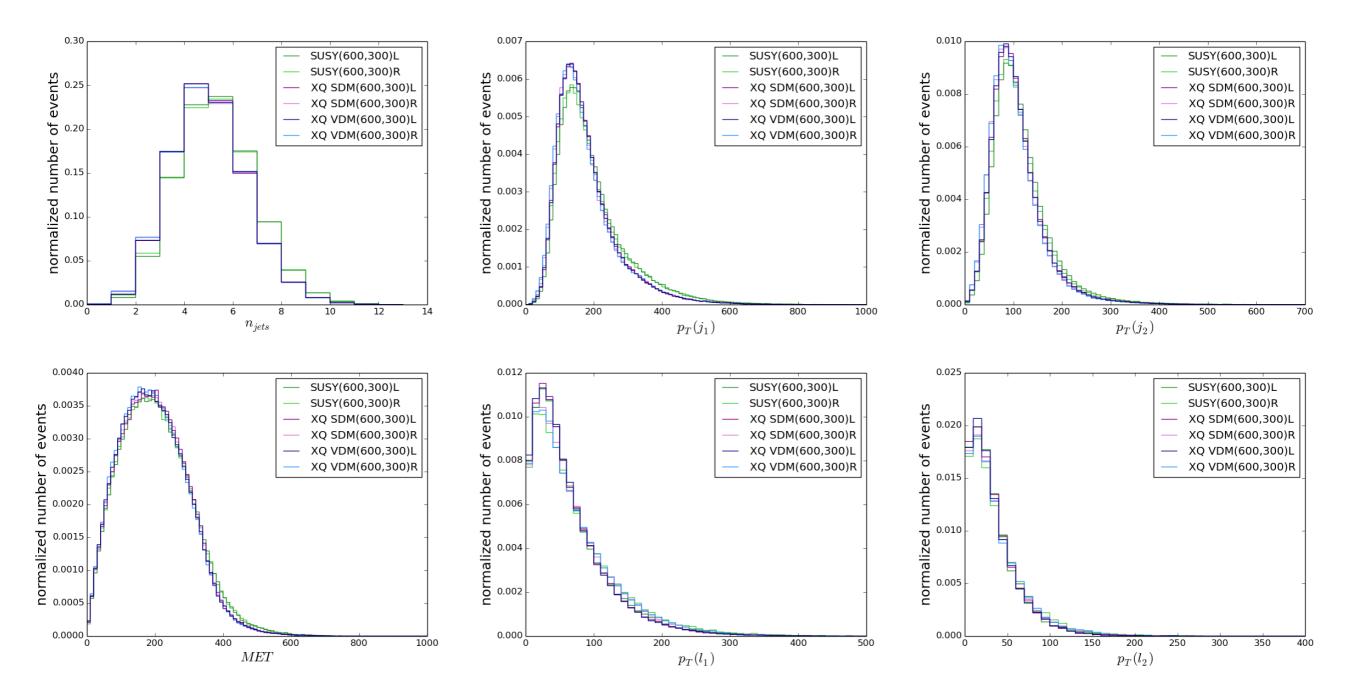
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# **Considered Diagrams**



# **Generator Level Distributions**



# **Cutflow Example (600,10)L**

#### ATLAS-CONF-2013-024

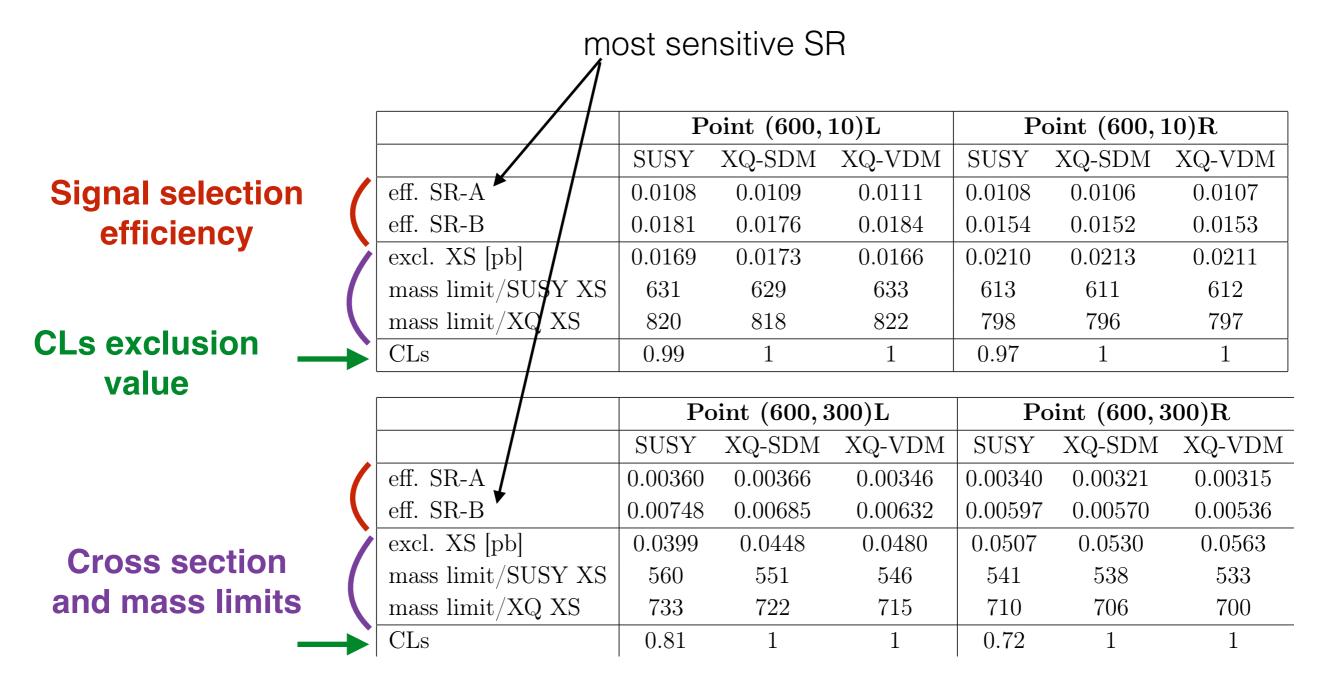
	SUSY	XQ-SDM	XQ-VDM
Initial no. of events	200000	200000	200000
$E_T^{\text{miss}} > 80 \text{ GeV} (\text{Trigger})$	187834 (-6.08 %)	187872 (-6.06 %)	188358~(-5.82~%)
muon veto $(p_T > 10 \text{ GeV})$	154643 (-17.67 %)	153946 (-18.06 %)	154710 (-17.86 %)
electron veto $(p_T > 10 \text{ GeV})$	123420 (-20.19 %)	$122439 \ (-20.47 \ \%)$	123247 (-20.34 %)
$E_T^{\text{miss}} > 130 \text{ GeV}$	113638 (-7.93 %)	112808 (-7.87 %)	113620 (-7.81 %)
$\geq 6 \text{ jets}, p_T > 80, 80, 35 \text{ GeV}$	33044~(-70.92~%)	27987 (-75.19 %)	28285 (-75.11 %)
reconstr. $E_T^{\text{miss,track}} > 30 \text{ GeV}$	$32564 \ (-1.45 \ \%)$	27563 (-1.51 %)	27901 (-1.36 %)
$\Delta \phi(E_T^{\text{miss}}, E_T^{\text{miss,track}}) < \pi/3$	31200~(-4.19~%)	26583~(-3.56~%)	26939 (-3.45 %)
$\Delta \phi(E_T^{ m miss},3~{ m hdst~jets}) > 0.2\pi$	26276 (-15.78 %)	22795 (-14.25 %)	23129 (-14.14 %)
tau veto	22880 (-12.92 %)	19967 (-12.41 %)	20354~(-12.00~%)
$2 \ b \ jets$	9668 (-57.74 %)	8510~(-57.38~%)	8660 (-57.45 %)
$m_T(b  \text{jets}) > 175  \text{GeV}$	7202 (-25.51 %)	$6447 \ (-24.24 \ \%)$	6579~(-24.03~%)
3 closest jets $80-270 \text{ GeV}$	6437~(-10.62~%)	5877 (-8.84 %)	5929~(-9.88~%)
same for second closest jets	3272~(-49.17~%)	3186~(-45.79~%)	$3351 \ (-43.48 \ \%)$
SR1: $E_T^{\text{miss}} \ge 200 \text{ GeV}$	3230 (-1.28 %)	3156 (-0.94 %)	3312 (-1.16 %)
SR2: $E_T^{\text{miss}} \ge 300 \text{ GeV}$	3067~(-5.05~%)	3000~(-4.94~%)	3161~(-4.56~%)
SR3: $E_T^{\text{miss}} \ge 350 \text{ GeV}$	2795 (-8.87 %)	2732 (-8.93 %)	2867 (-9.30 %)

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# **Results Single Lepton Final State**

CMS-SUS-13-011 (cut-based), implemented in MA5



# **Exclusion in the Mass Plane**

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right-handed couplings

