

# Possible 750 GeV diphoton signal via light pseudoscalars

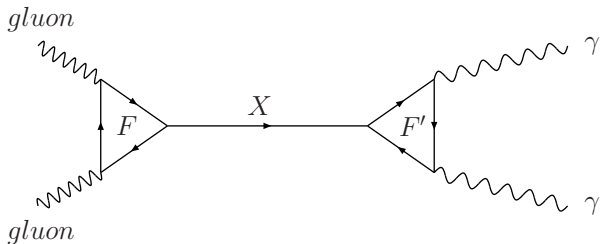
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with C. Hugonie, arXiv:1602.03344,  
see also F. Domingo et al., arXiv:1602.07691

Data ( $\lesssim$  Moriond 2016):

- ATLAS at 13 TeV,  $710 \text{ GeV} < M_{\gamma\gamma} < 790 \text{ GeV}$  (two bins):  
21 events vs. 11.3 expected; local excess  $3.9 \sigma$  ( $2.0 \sigma$  incl. LLE);  
compatible with 8 TeV at the  $1.2 \sigma$  level (assuming  $ggF$ )
- CMS at 13 TeV,  $750 \text{ GeV} < M_{\gamma\gamma} < 770 \text{ GeV}$  (one bin):  
11 events vs. 5.4 expected; local excess  $2.8 \sigma$  ( $\sim 1 \sigma$  incl. LLE);  
combined with 8 TeV: local excess  $3.4 \sigma$  ( $1.6 \sigma$  incl. LLE)
- Signal cross sections of  $\sim 3 - 8 \text{ fb}$  would explain the excesses

"Standard" interpretation:



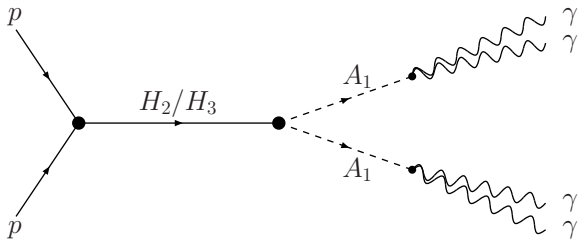
- $X$ : Scalar or pseudoscalar (possibly composite) with  $M_X \sim 750$  GeV
- Coupling to gluons through loops of coloured fermions  $F$
- Coupling to photons through loops of charged fermions  $F'$  ( $\sim F$ ?)
- Possibly a large width ( $\gtrsim$  a few GeV) in order to explain the ATLAS data

## Challenges:

- Need large (loop induced) production cross section  
→ need large ( $\sim$  non-perturbative)  $XFF$  Yukawa coupling
- Need large (loop induced) width into  $\gamma\gamma$   
→ need large ( $\sim$  non-perturbative)  $XF'F'$  Yukawa coupling
- Tree level decays of  $X$  must be (practically) forbidden, otherwise the loop induced decay into  $\gamma\gamma$  would have a too small branching fraction  
→  $X$  must not couple to Standard Model fermions (or Higgs), the new fermions  $F$  ( $F'$ ) must be heavier than  $M_X/2 \sim 375$  GeV
- A large width into  $\gamma\gamma$  is tough to get...
- $\gtrsim 200$  BSM scenarios of this type... (more than events)

## Alternative scenario with light pseudoscalars $A_1$ :

(S. Knapen et al., P. Agrawal et al., J. Chang et al., ...)



Viable if  $M_{A_1} \lesssim 800$  MeV; then the photons from  $A_1$  decays are sufficiently collimated such that they appear (mostly) as a single photon in the electromagnetic calorimeters (see below)

## Constraints on resonance(s) $H_{(i)}$ at $\approx 750$ GeV:

- Sufficient production cross section in  $ggF$  or *ass. prod. with  $b$ -quarks*
- Large branching fraction into  $A_1A_1$

## Constraints on a light pseudoscalar $A_1$ below $\approx 800$ MeV:

- Not ruled out by low energy experiments
- Large branching fraction into  $\gamma\gamma$
- Decay length  $\lesssim 1$  m, preferably shorter

## A concrete scenario: the NMSSM

featuring 3 scalars  $H_{1,2,3}$  and two pseudoscalars  $A_{1,2}$

With  $H_1 = \text{SM-Higgs}$  at 125 GeV:

Two candidates for scalar(s)  $H_2/H_3$  at  $\approx 750$  GeV:

- the “MSSM-like” scalar  $H$  with potentially large production cross section via  $bbH$  if  $\tan\beta \gtrsim 10$
- the singlet-like scalar  $H_S$  with potentially large branching fraction into singlet-like  $A_1A_1$  ( $A_2$  is the MSSM-like pseudoscalar with  $M_{A_2} \sim M_H$ )
- > Best solution: both scalars have masses of  $\approx 750$  GeV,  $H$  and  $H_S$  mix strongly and form  $H_2/H_3$ ; two nearby narrow states can imitate a large width as seen by ATLAS

A light pseudoscalar  $A_1$  can be a (pseudo-) Goldstone boson of an  $R$ -symmetry ( $\leftrightarrow$  small trilinear couplings  $A_\lambda, A_\kappa$  in the scalar potential);

Impossible in the MSSM where the  $\mu$ -term breaks  $R$ -symmetry;  
in the NMSSM,  $\mu$  is replaced by the vev of a singlet field  $S$   
 $\rightarrow$  a (weakly broken)  $R$ -symmetry is possible

**But:** Broken by radiative corrections  $\sim A_{top}$ , gaugino masses  
 $\rightarrow$  Tuning is still required for  $M_{A_1} \lesssim 800$  MeV

## Possible $A_1$ masses satisfying the above constraints:

(1)  $M_{A_1} \sim M_{\pi^0} \sim 135 \text{ MeV}$  (Domingo et al., arXiv:1602.07691):

- $A_1$  mixes with  $\pi^0$ , hence  $A_1$  decays with a similar width (short decay length) into  $\gamma\gamma$ ; calculable using PCAC

Heavier  $A_1$ :  $135 \text{ MeV} < M_{A_1} < 2m_\mu$ :

- Susy loops generate flavour changing couplings of the extra (MSSM-like) Higgs bosons, hence also for  $A_1$  (through mixing with the MSSM-like  $A_2$ )
- dangerous rare decays  $K^\pm \rightarrow \pi^\pm e^+ e^-$  (less constraining:  $B^\pm \rightarrow K^\pm e^+ e^-$ ) unless the soft Susy breaking terms are chosen such that contributions to flavour changing couplings cancel, which is possible (see arXiv:1602.07691)
- $A_1$  decays dominantly into  $e^+ e^-$  with a decay length  $\gtrsim 40 \text{ m}$  → useless



(2)  $M_{A_1} \lesssim 2m_\mu \sim 211$  MeV (U.E., C. Hugonie, arXiv:1602.03344):

- The muon loop induced BR into  $\gamma\gamma$  is enhanced up to  $\sim 75\%$  if  $M_{A_1}$  is just below the threshold (see A. Bharucha et al., arXiv:1603.04464)
- The decay length is reduced to 2–5 m, but the production cross section can be large enough such that enough  $A_1 \rightarrow \gamma\gamma$  decays take place before the EM calorimeter
- Soft Susy breaking terms have to be chosen such that flavour changing couplings are cancelled

$M_{A_1} \gtrsim 500$  MeV: Constraints from rare  $K$  decays disappear

### (3) $M_{A_1} \sim 510$ MeV (U.E., C. Hugonie, arXiv:1602.03344):

- At the parton level, the dominant decays of  $A_1$  are into  $s\bar{s}$  and gluons
- But: one is still in the nonperturbative regime of QCD  
Best guess:  $s\bar{s}$  and  $F\tilde{F}_{(QCD)}$  act as interpolating fields;  
these are part of the  $\eta$  wave function in Fock space ( $M_\eta \sim 548$  MeV),  
hence  $A_1$  decays like the  $\eta$  meson:  
 $BR(\eta \rightarrow \gamma\gamma) \sim 39\%$  ,  $BR(\eta \rightarrow 3\pi^0) \sim 33\%$  ,  $BR(\eta \rightarrow \pi^+\pi^-\pi^0) \sim 23\%$
- $BR(A_1 \rightarrow \gamma\gamma) \sim 39\%$  ,  $BR(A_1 \rightarrow 3\pi^0 \rightarrow 6\gamma) \sim 33\%$   
with a decay length below 1 mm (?to be confirmed?)
- Dominant constraint: Now from searches for  $\Upsilon(1S) \rightarrow \gamma \eta$  decays by CLEO  
where no events were seen  
(but 2 events for  $M_{\pi^+\pi^-\pi^0} \sim 510$  MeV in the  $\eta \rightarrow \pi^+\pi^-\pi^0$  search channel)  
→ constraints on the coupling  $A_1 b\bar{b}$ ;  
if too large, CLEO would have observed  $\Upsilon(1S) \rightarrow \gamma A_1 \rightarrow 3\pi^0$  decays
- These constrain the  $BR(H_{2,3} \rightarrow A_1 A_1)$ , still:  
**a signal cross section up to 6.7 fb is possible<sup>1</sup>**

<sup>1</sup>Modulo acceptance of multiphotons as a single photon, see below

(4)  $M_{A_1} \sim M_\eta \sim 550$  MeV (U.E., C. Hugonie, arXiv:1602.03344):

- $A_1$  mixes strongly with the  $\eta$  meson, its corresponding branching fractions are no longer educated guesses (calculable using PCAC)
- But: Constraints from CLEO from unseen  $\Upsilon(1S) \rightarrow \gamma A_1$  decays are somewhat stronger, still:  
a signal cross section up to 3.4 fb is possible<sup>2</sup>

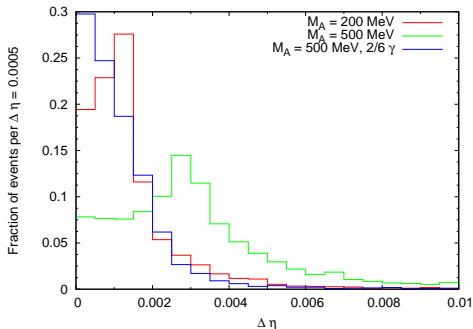
These - and a recent MSSM scenario with  $M_{\text{stop}} \sim 375$  GeV by A. Djouadi and A. Pilaftis - are the only known scenarios for the 750 GeV diphoton excess without extra “ad hoc” fermions, but based on known Susy extension of the SM

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<sup>2</sup>Modulo acceptance of multiphotons as a single photon, see below

If the excess of events persists, these scenarios can be distinguished (or ruled out) experimentally:

- The fineness of first layer of cells of the EM calorimeter along  $\eta$  (rapidity, the angle along the beam axis) ranges from 0.003 to 0.006 (ATLAS); the spread of in  $\Delta\eta$  of multiphotons depends on  $M_{A_1}$ :  
(2/6  $\gamma$  denote the two leading among 6 photons from  $3\pi^0$ )



→ This plot helps to estimate the acceptances for multiphotons to fake a single photon:  $\sim 80\%$  for  $M_A \sim 200$  MeV or 2/6  $\gamma$ ,  $\sim 30\%$  for  $M_A \sim 200$  MeV

- If  $M_{A_1} \sim 211$  MeV: The  $A_1$  decay length is macroscopic, and  $A_1$  may decay inside the EM calorimeters  
(before the EM calorimeters, the  $A_1 \rightarrow \gamma\gamma$  vertex is invisible)
- The photons can convert in the material before the EM calorimeter leading to electrons which are visible, but usually added to the photon signal in the EM (20% for rapidity  $\eta \sim 0$  to 45% for  $\eta \sim 1.6$ )
  - photon-jets lead to more converted photons than a single photon
  - one can potentially distinguish single photons from collinear diphotons or, in the case  $A_1 \rightarrow 3\pi^0 \rightarrow 6\gamma$ , from collinear 6 photons (B. Dasgupta et al., arXiv:1602.04692) iff the  $A_1$  decays occur inside the material
- If the signal originates from two nearby states  $H_2/H_3$ , their masses can potentially be separated (depending on the actual  $H_2/H_3$  mass splitting)

→ With more data, the different scenarios can be distinguished!

Exciting times may lie ahead of us!