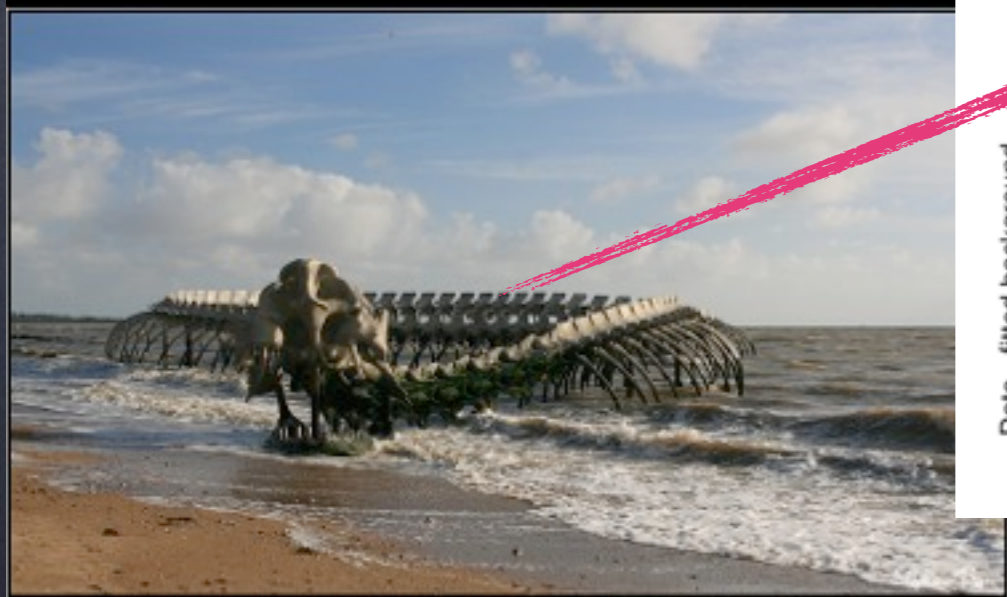
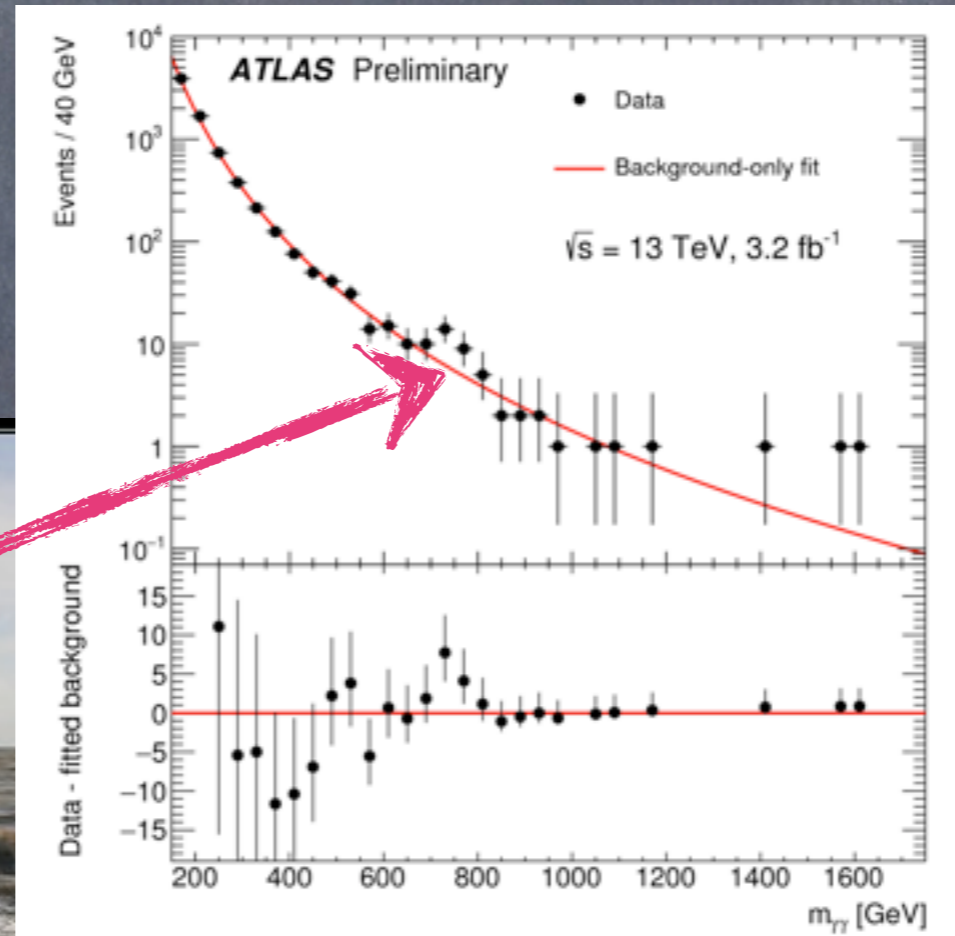
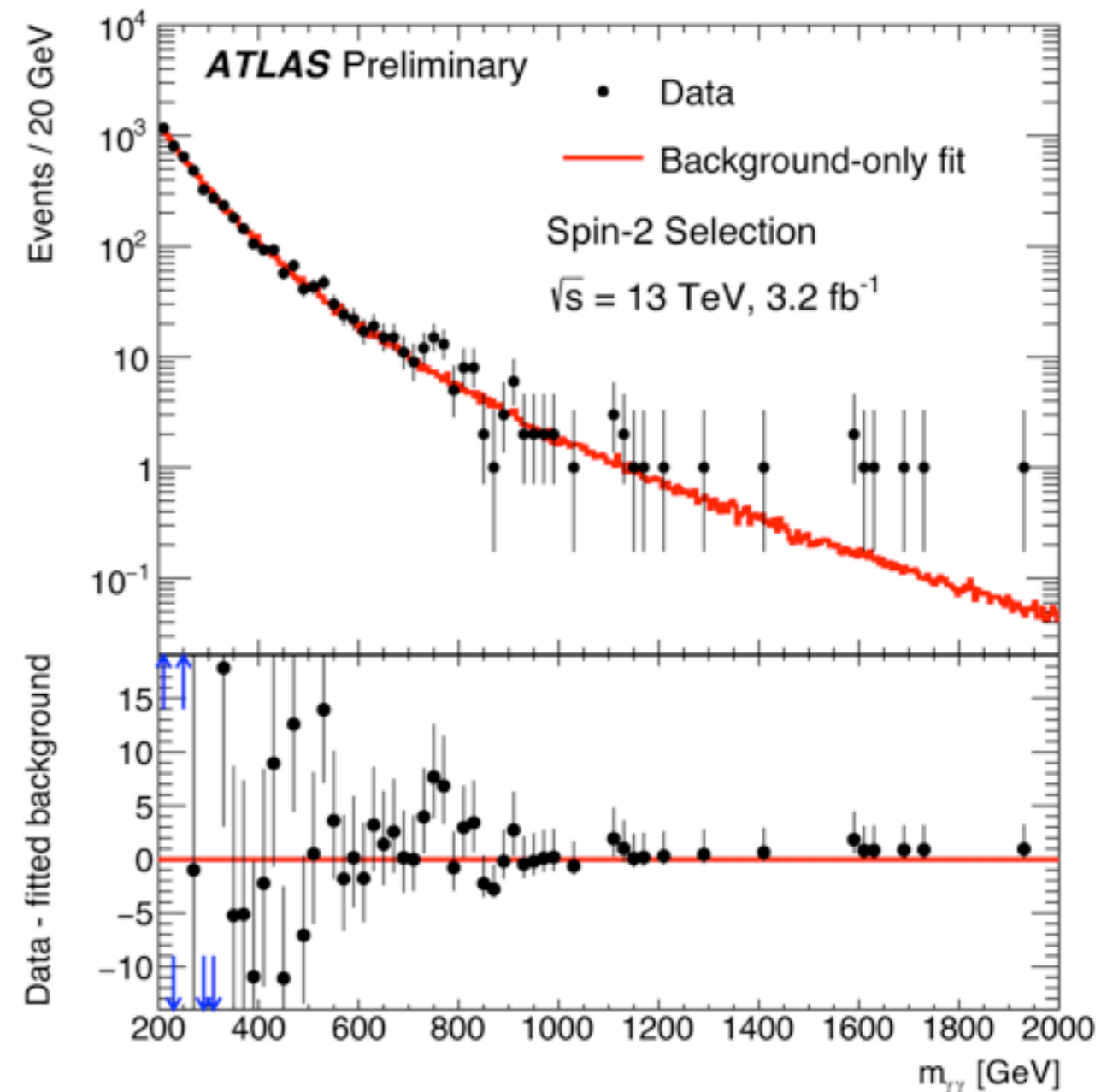
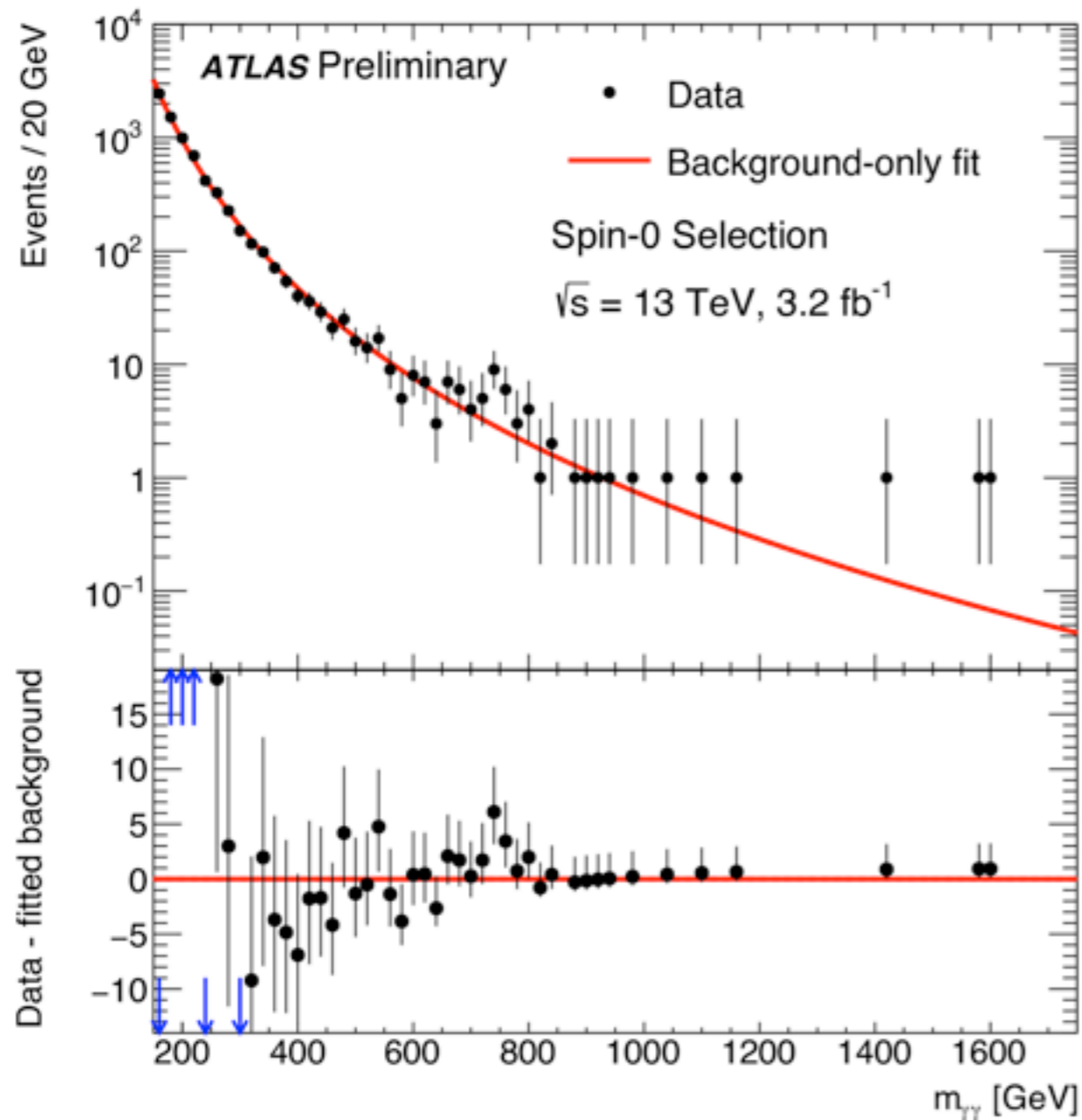


X(750) model review



Nantes, 24 May, 2016

ATLAS Run-2 Data - Spectrum



$$p_{T,\gamma_1} > 0.4 m_{\gamma\gamma} \quad (300 \text{ GeV})$$

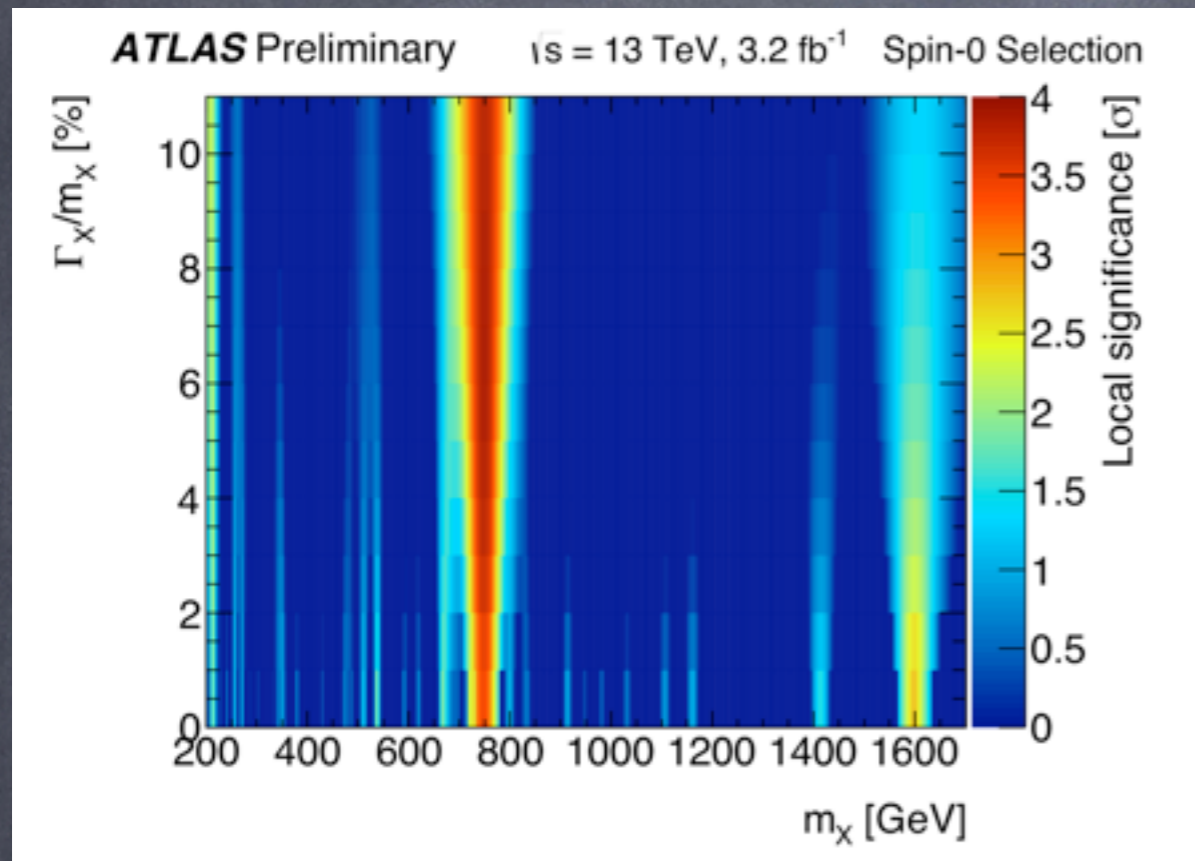
$$p_{T,\gamma_2} > 0.3 m_{\gamma\gamma} \quad (225 \text{ GeV})$$

15 events in 40 GeV window
around 750 GeV, roughly 10 above bg

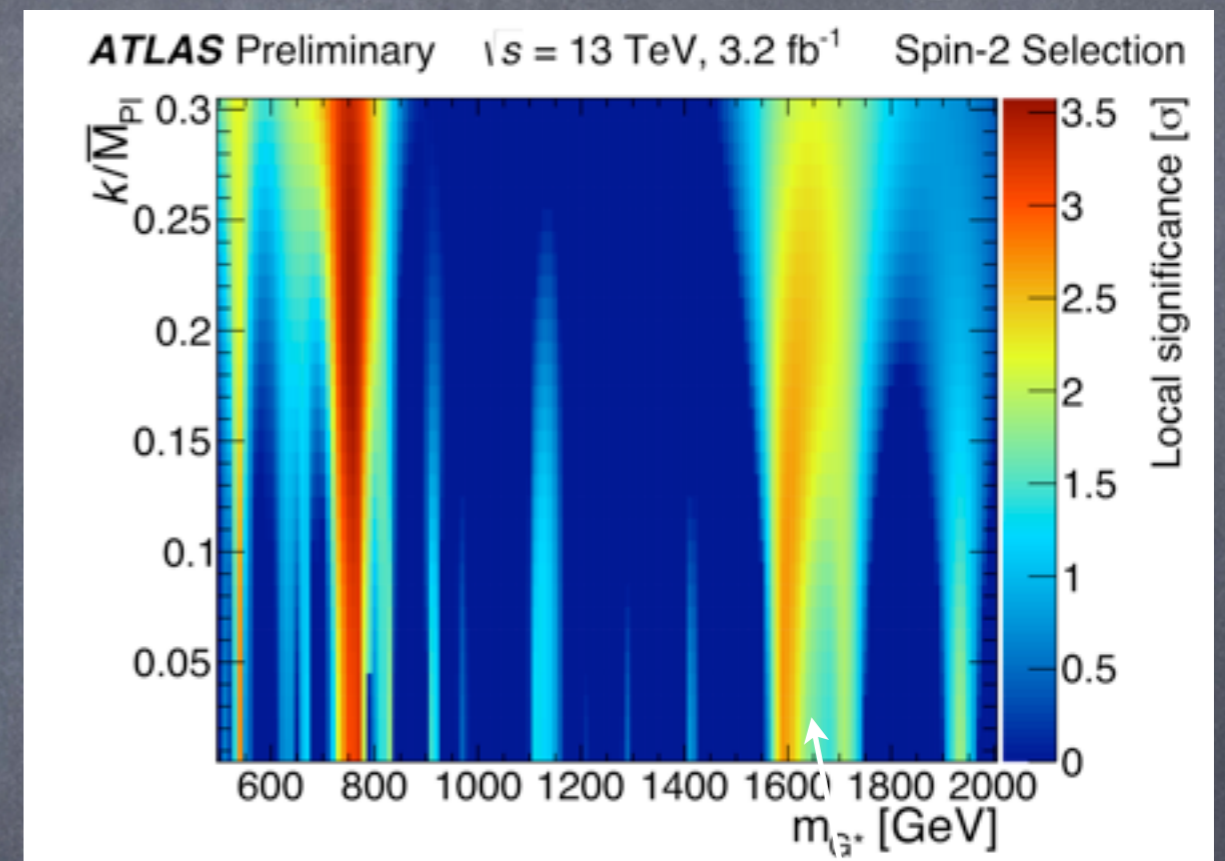
$$p_{T,\gamma} > 55 \text{ GeV}$$

40 events in 60 GeV window
around 750 GeV, roughly 19 above bg

ATLAS Run-2 Data - Significance



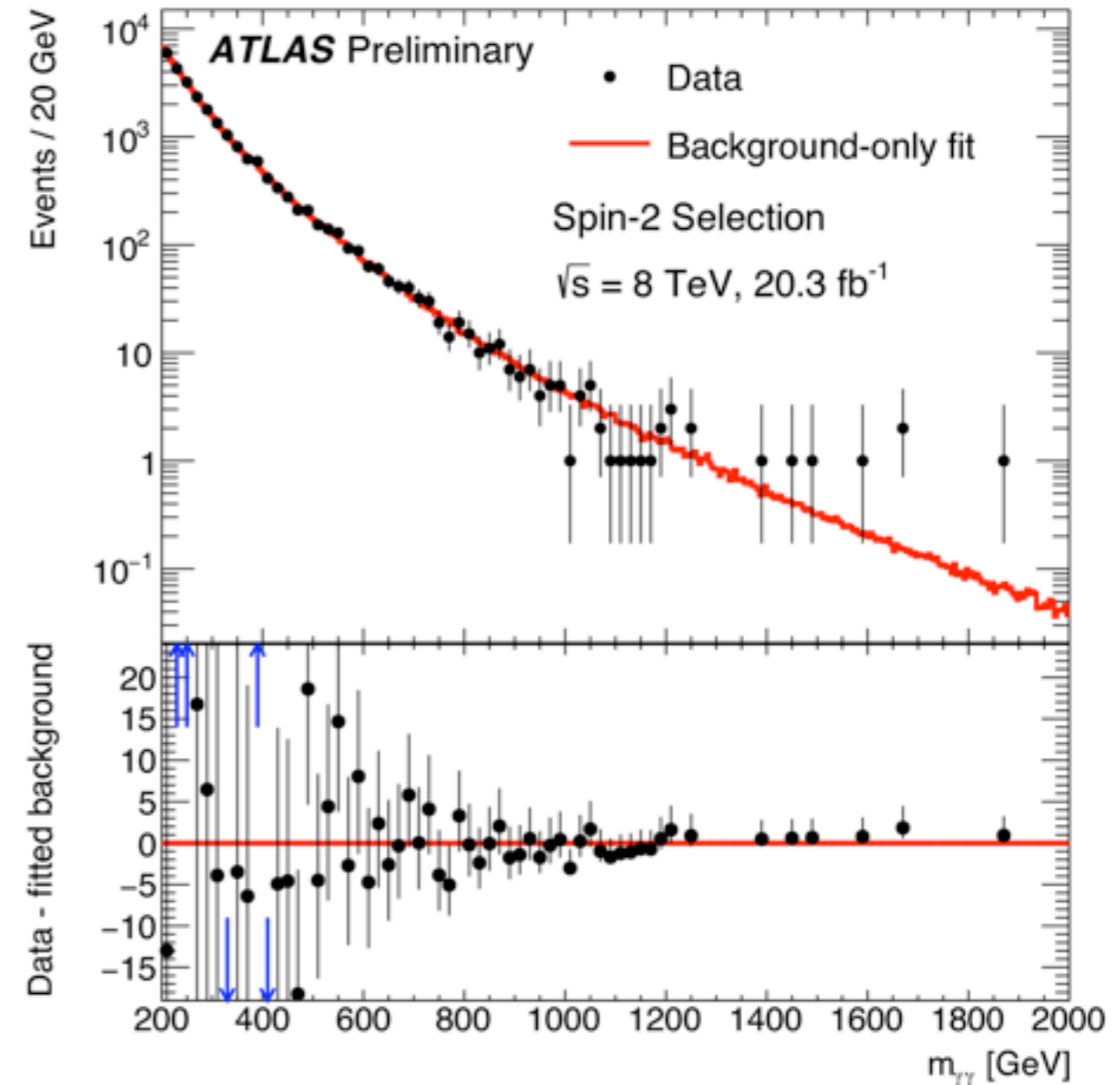
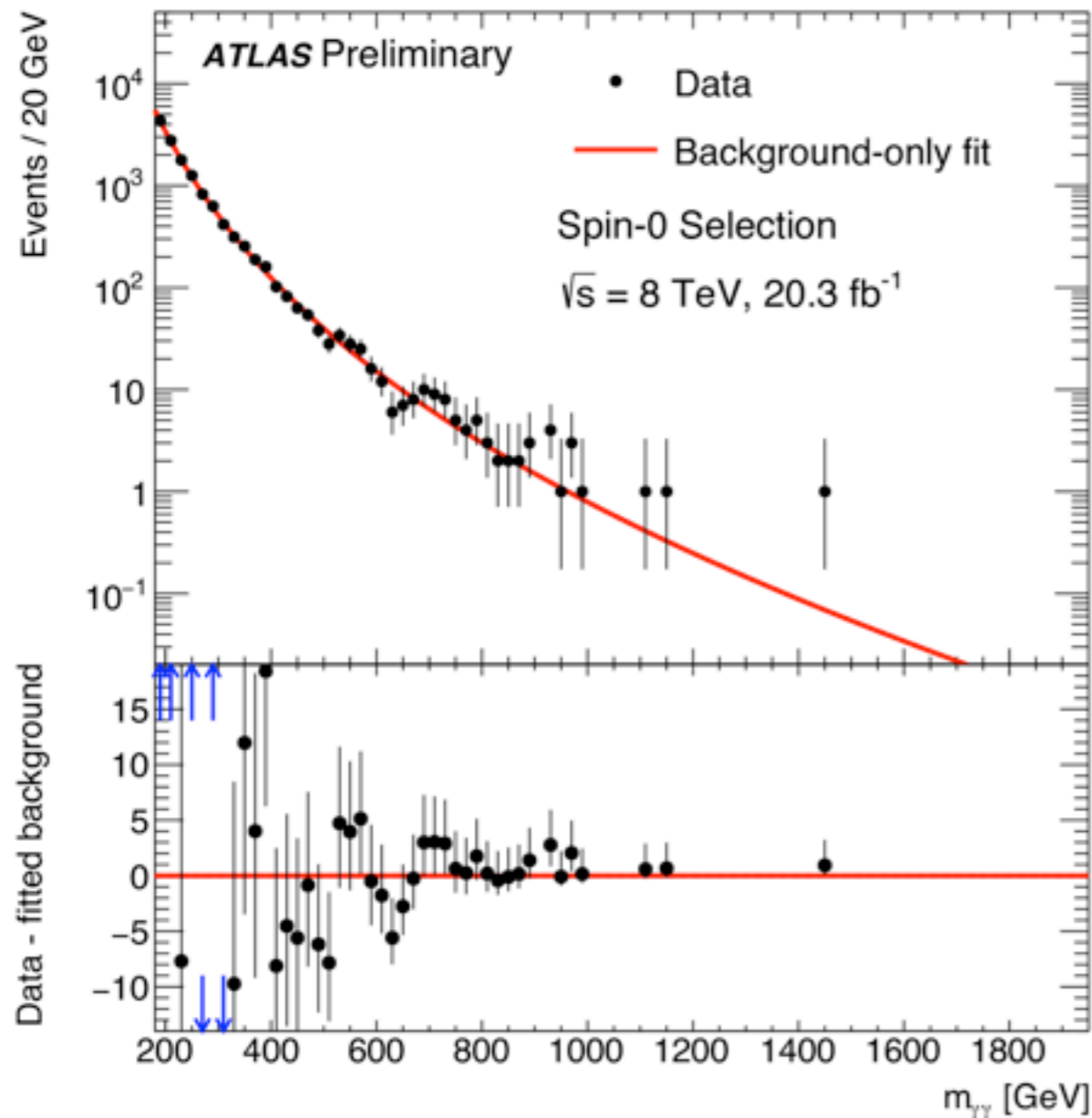
At 750 GeV
 3.9σ excess for $\Gamma=45 \text{ GeV}$
 3.6σ excess for $\Gamma=0$



At 750 GeV
 3.6σ excess for
 $\Gamma=48 \text{ GeV}$

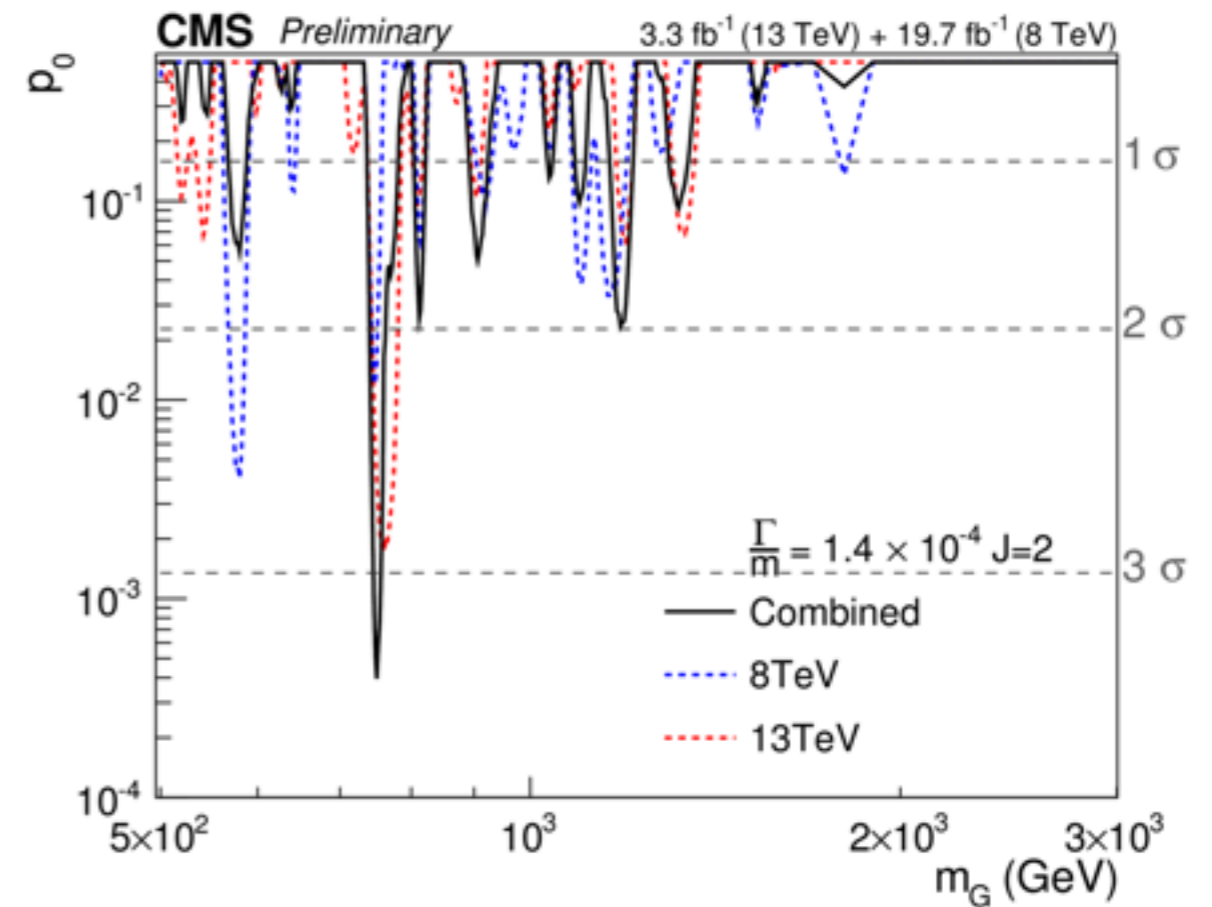
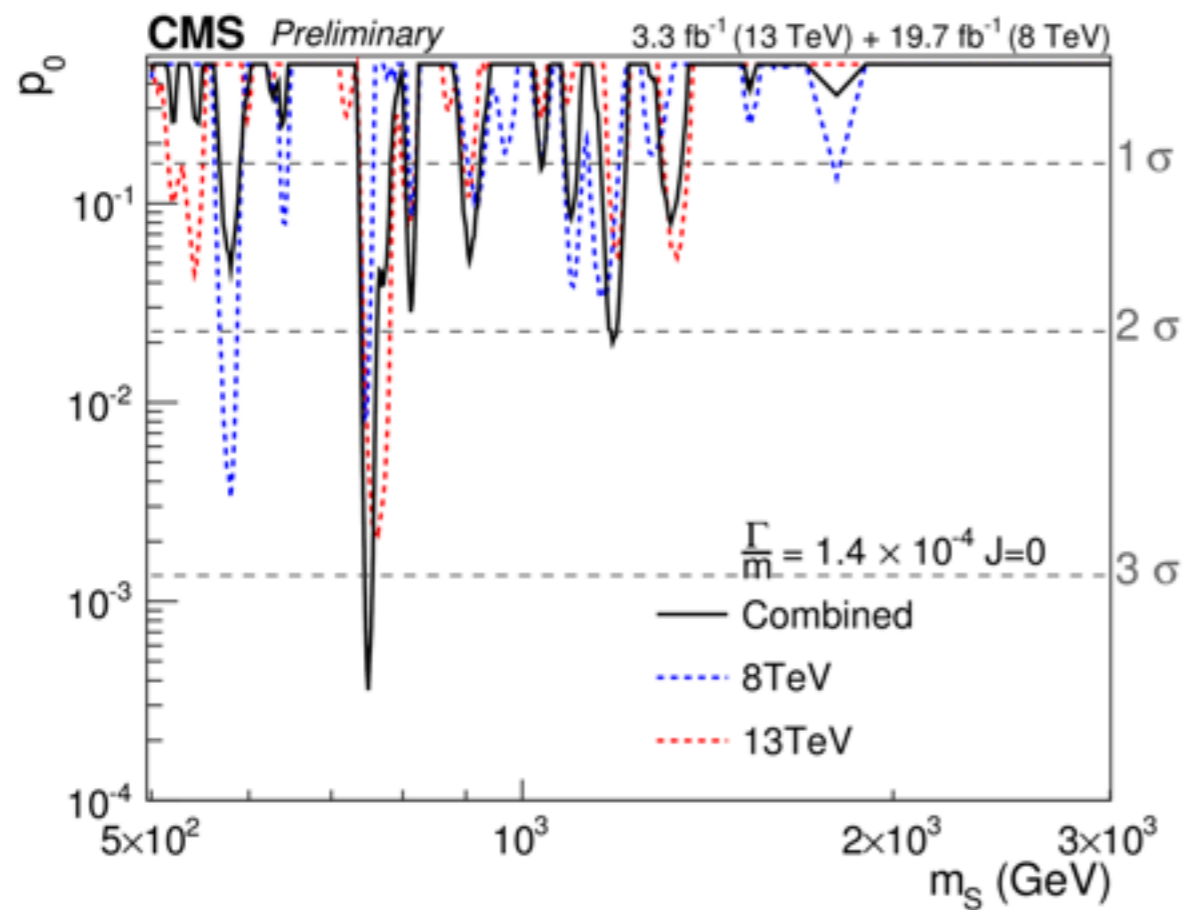
2nd KK mode
 also visible ;)

ATLAS Run-1 Data - Spectrum



For spin-0 analysis, 1.9σ excess at 750 GeV in run-1. Decent compatibility (at 1.2σ) between run-2 and run-1 diphoton bumps assuming gluon-fusion production. Much worse compatibility (at 2.7σ) for spin-2 analysis.

CMS Run-2 and Run-1 Data



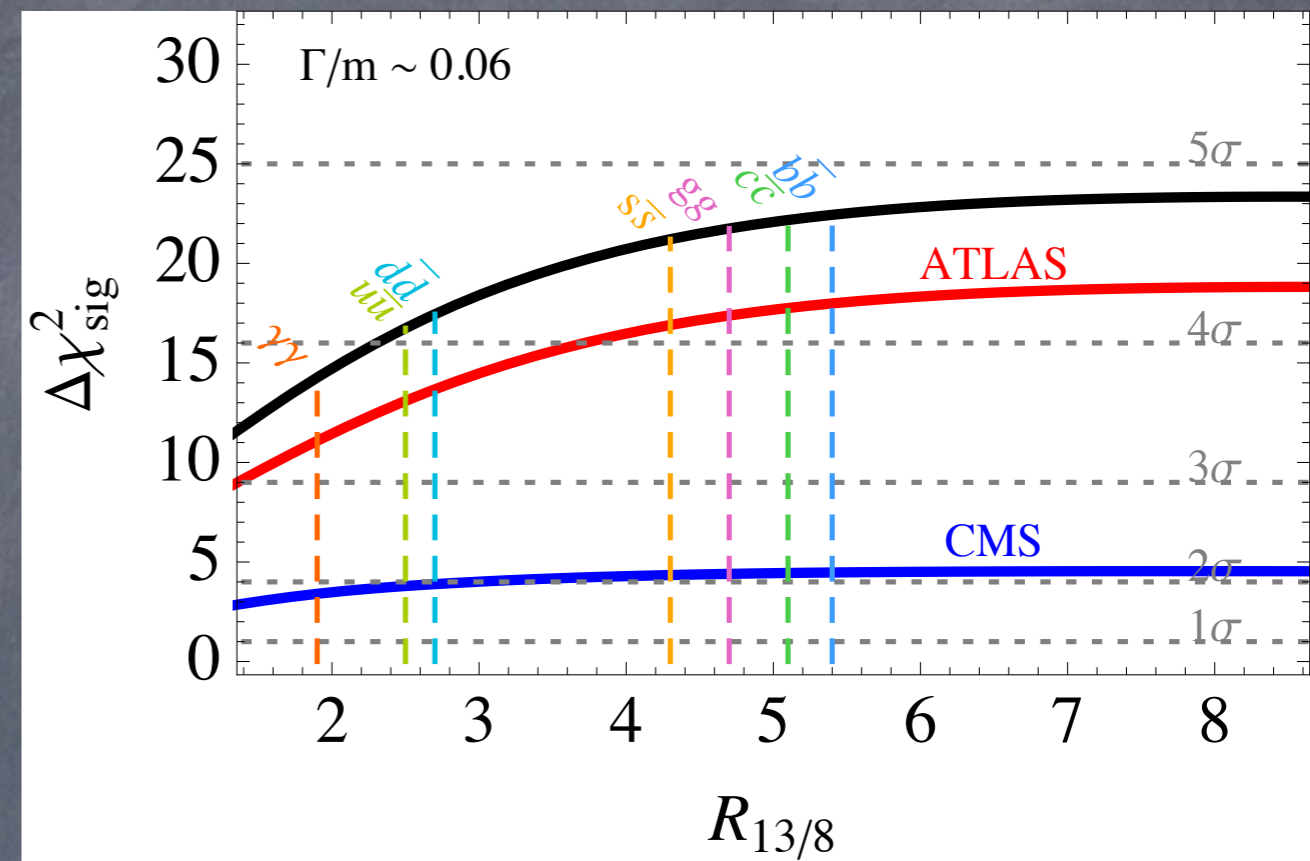
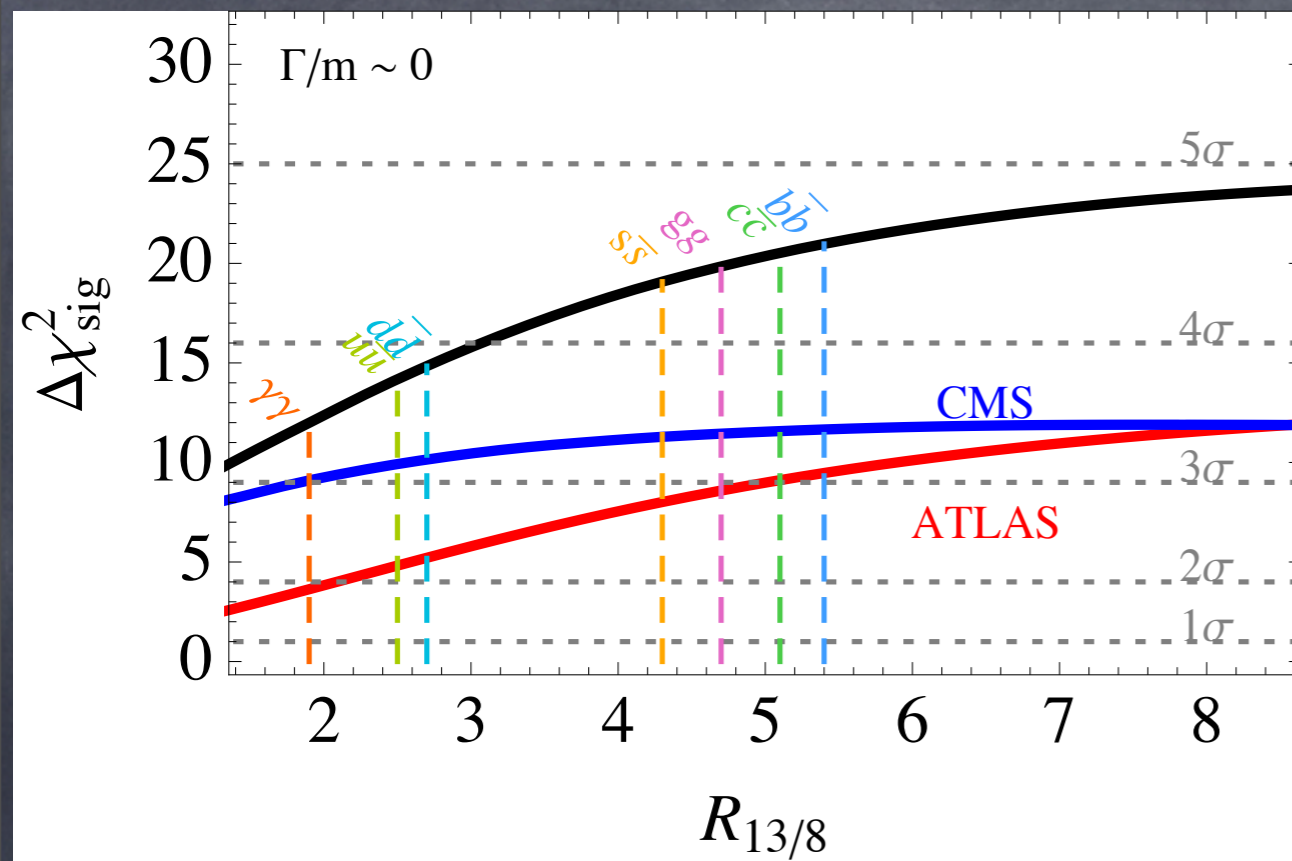
- 2.9 σ excess at 760 GeV in run-2 data. Adding B=0 data slightly increased significance
- Very good compatibility of ATLAS and CMS diphoton bumps at 750 GeV
- Very good compatibility between CMS run-2 and run-1 data, this time independently of the spin hypothesis. 3.4 σ excess at 750 GeV in combined run-1 and run-2 data

Main questions

- Production process?
- Narrow or wide?
- Other decays channels?
- Spin 0 or Spin 2 (or higher)?
- Parity even or parity odd?
- Singlet or multiplet?
- One particle or a part of a larger sector?
- Meaning of life and universe?

Post Moriond fits

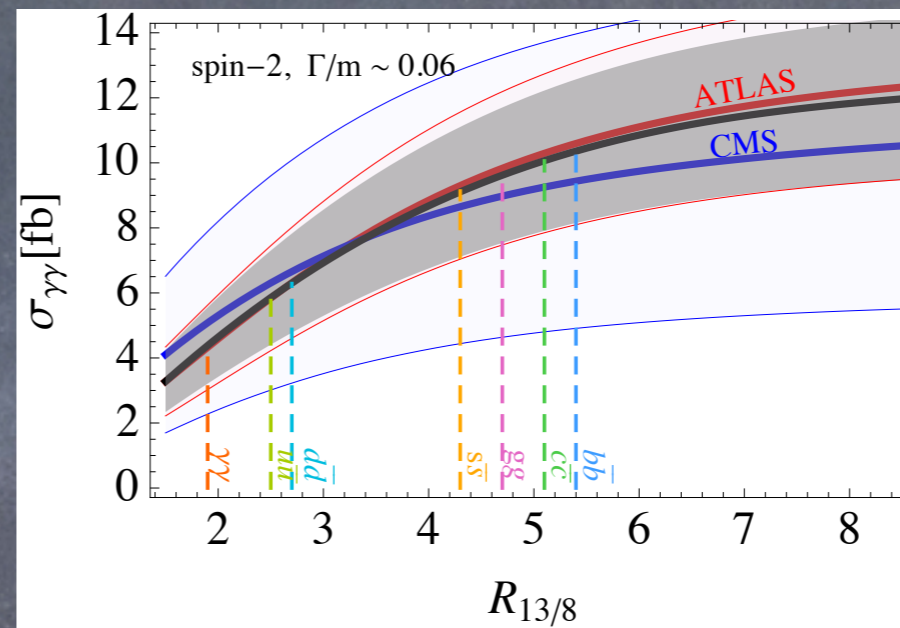
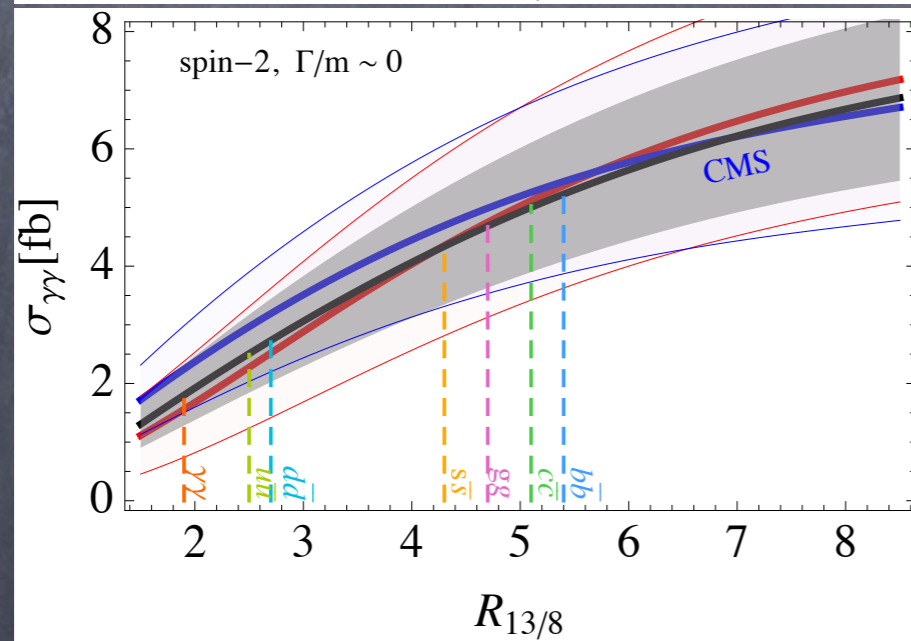
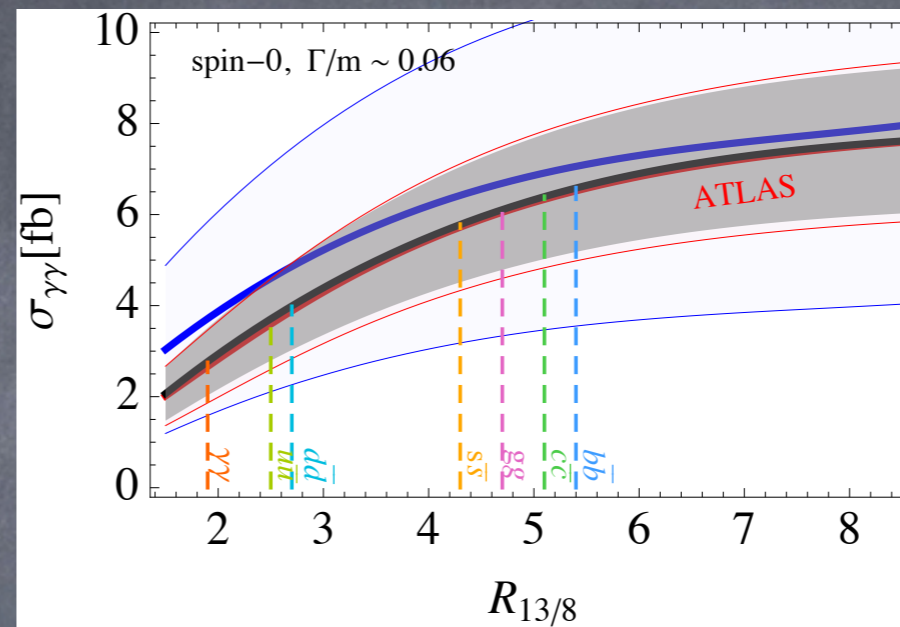
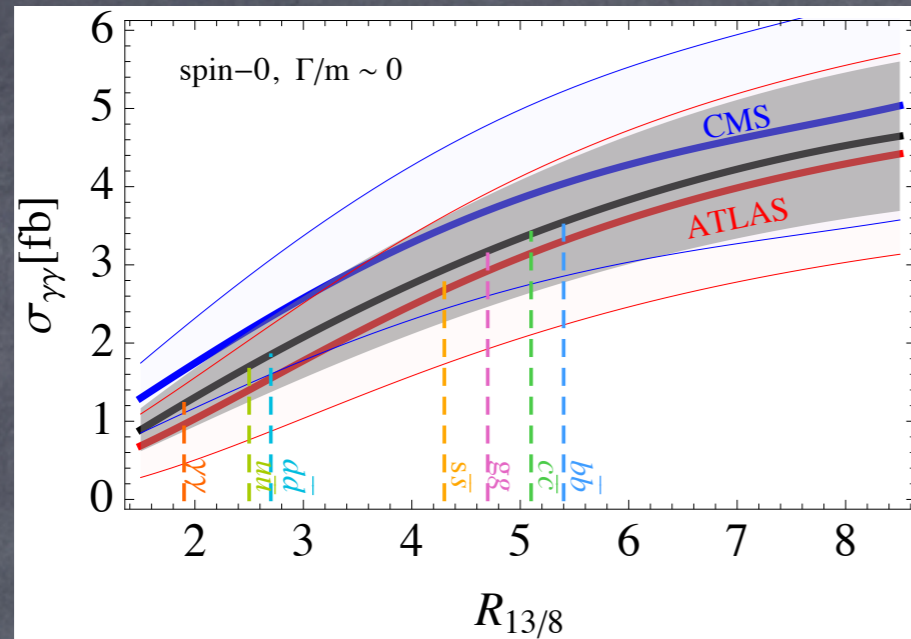
Kamenik et al
1603.06566



- The larger the ratio of 13 to 8 TeV cross sections, the more significant is the combined ATLAS+CMS signal
- Preference for large width is significant for ATLAS alone, but marginal in combined data
- At this point it's no longer "ATLAS diphoton excess", it's "LHC diphoton excess"

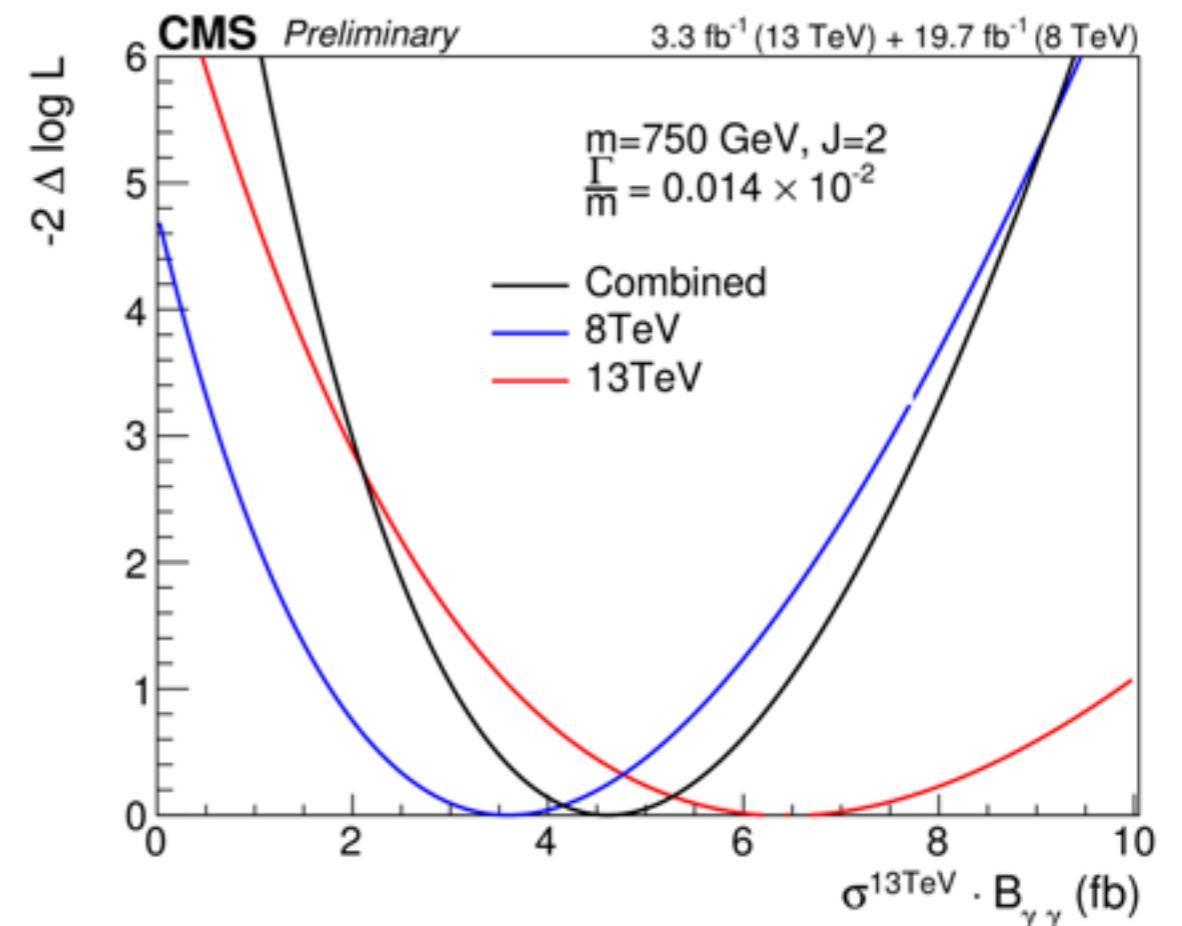
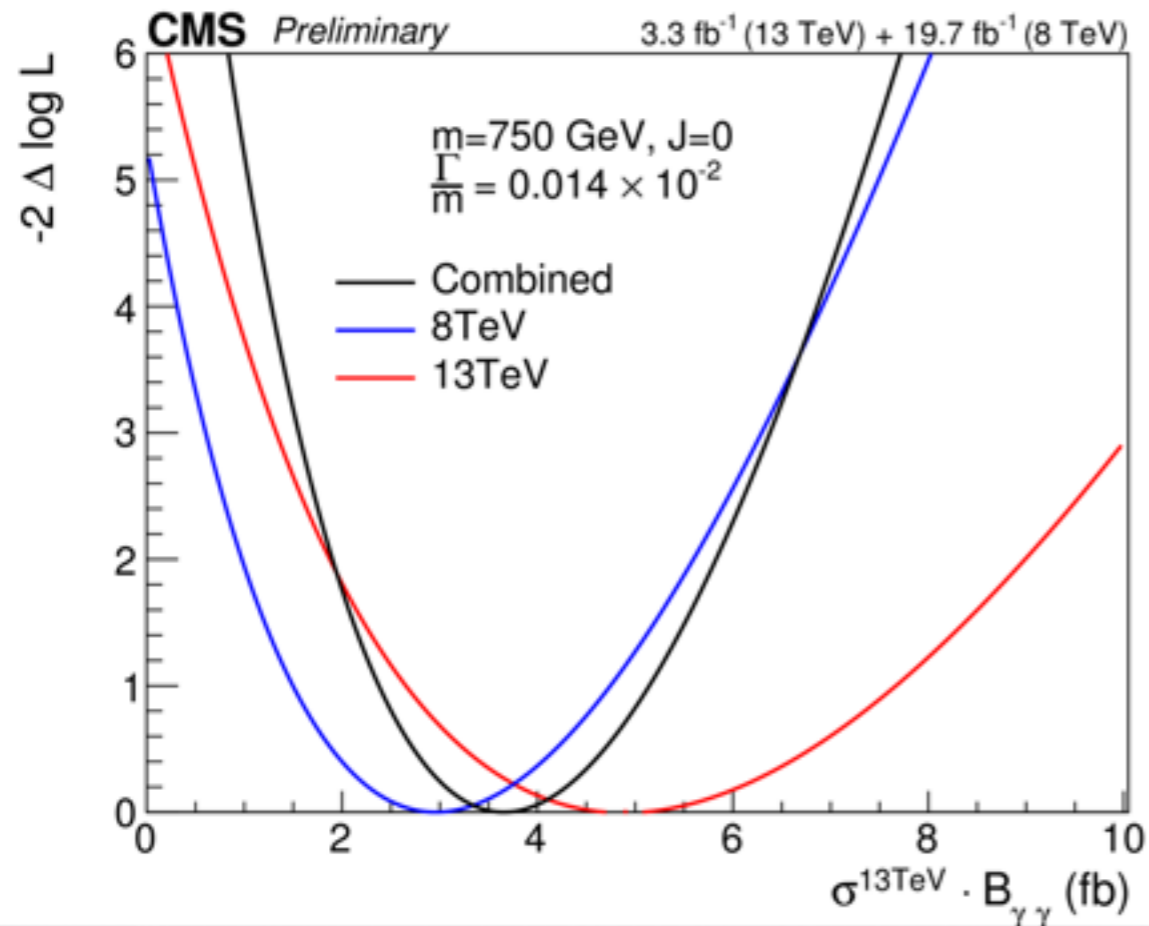
Best fit cross section

Kamenik et al
1603.06566



- Combining run-1 and run-2 data, best fit cross section for narrow scalar resonance produced in gluon fusion is around $\sigma(pp \rightarrow S) \text{Br}(S \rightarrow \gamma\gamma) \approx 3 \text{ fb}$
- Slightly larger cross sections needed for large width and/or larger spin

What is the mass and cross section?

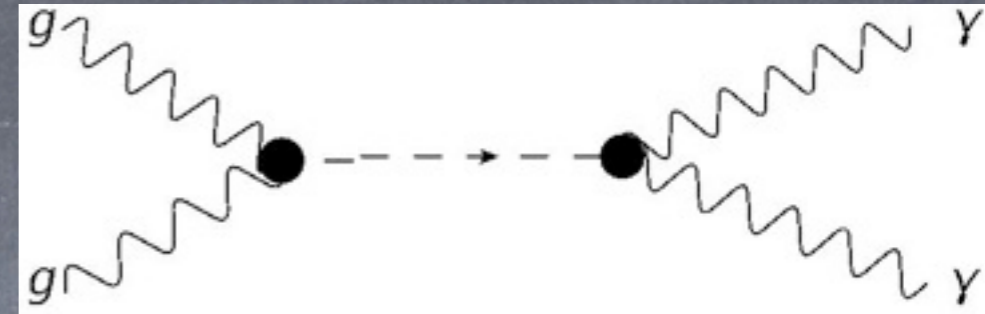


CMS xsec fits in good agreement with theorist fits

Everyone's model

AA,Slone,Volansky
1512.05777

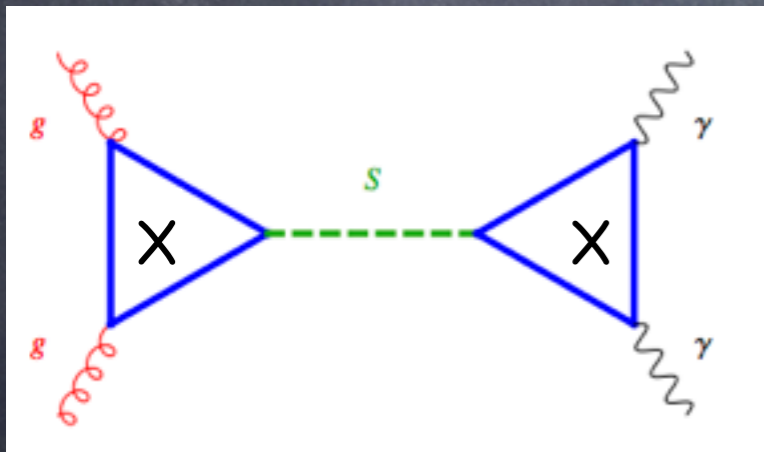
Scalar field S coupled to photons and gluons via effective non-renormalizable interactions



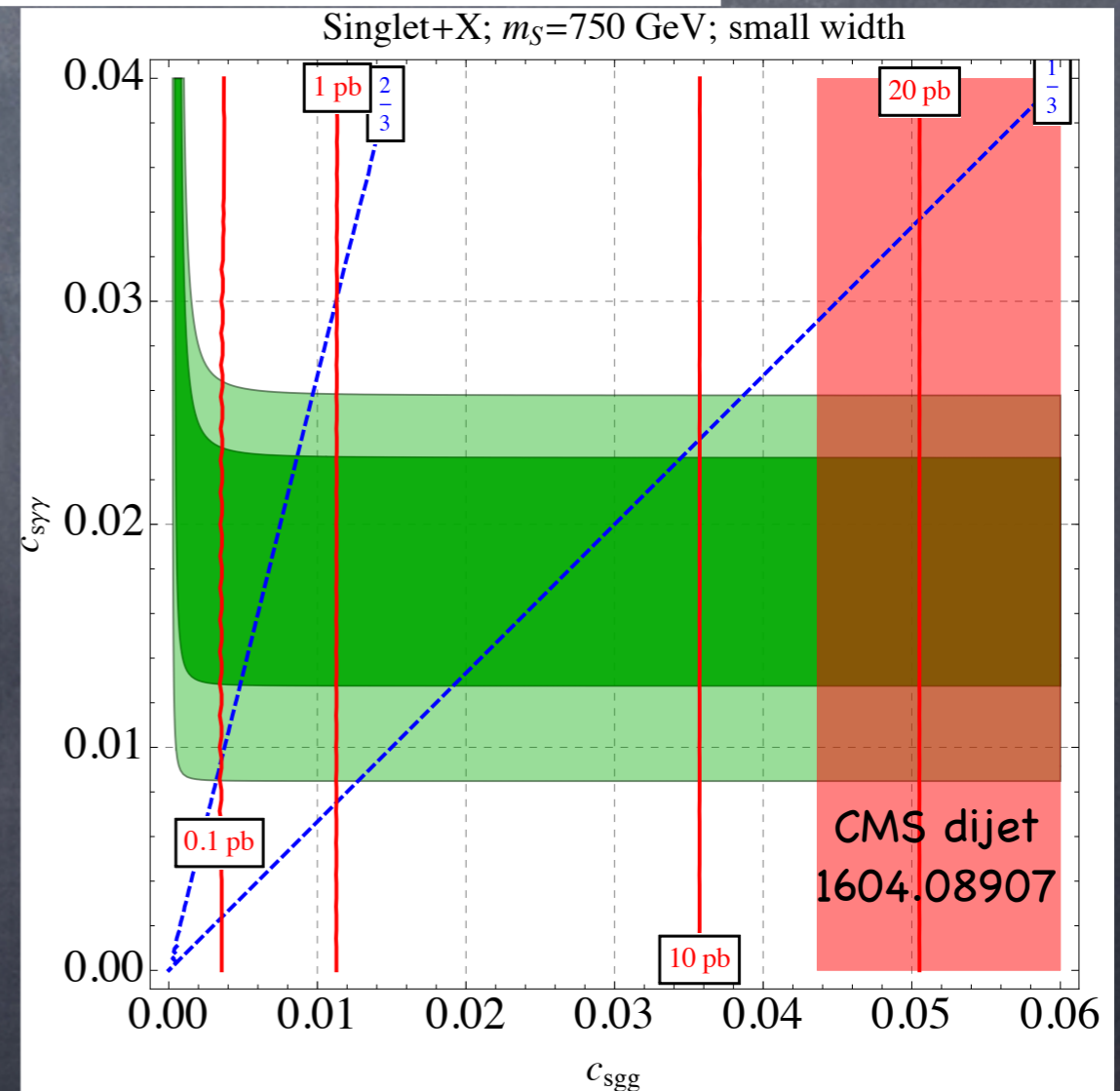
$$\mathcal{L}_{S,\text{eff}} = \frac{S}{4v} \left(c_{sgg} g_s^2 G_{\mu\nu}^a G_{\mu\nu}^a + c_{sw} g_L^2 W_{\mu\nu}^i W_{\mu\nu}^i + c_{sbb} g_Y^2 B_{\mu\nu} B_{\mu\nu} \right)$$

$$c_{s\gamma\gamma} = c_{sw} + c_{sbb}$$

$$\mathcal{L}_{S,\text{eff}} = \frac{e^2}{4v} c_{s\gamma\gamma} S A_{\mu\nu} A_{\mu\nu} + \frac{g_s^2}{4v} c_{sgg} S G_{\mu\nu}^a G_{\mu\nu}^a,$$



$$c_{sgg} = \frac{y_X v}{12\pi^2 m_X}, \quad c_{s\gamma\gamma} = \frac{y_X Q_X^2 v}{2\pi^2 m_X}.$$



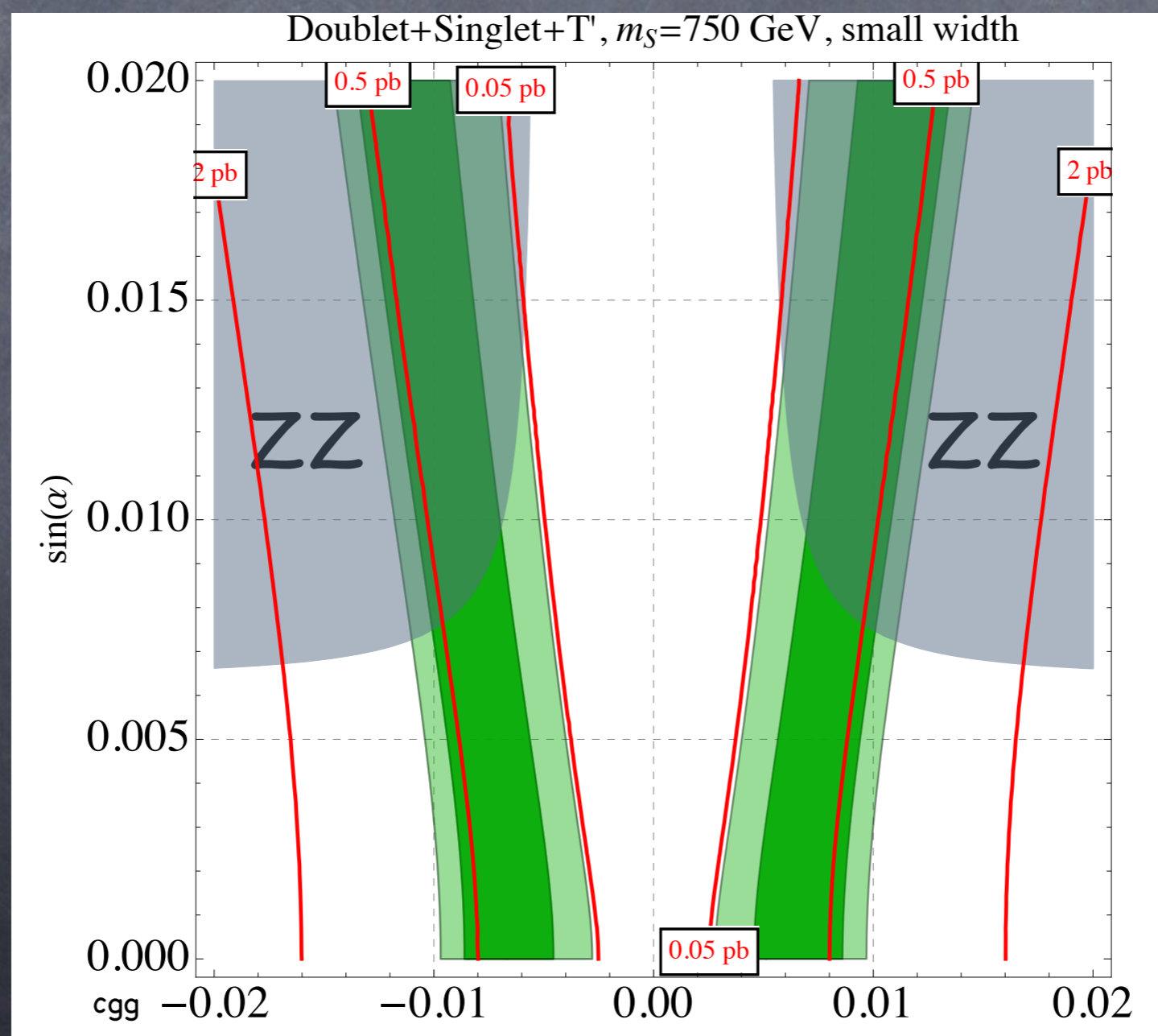
What else it decays to?

| final state f | σ at $\sqrt{s} = 8 \text{ TeV}$ | | | σ at $\sqrt{s} = 13 \text{ TeV}$ | | |
|----------------------|--|-------------------------|------|---|--------------------|------|
| | observed | expected | ref. | observed | expected | ref. |
| $e^+e^-, \mu^+\mu^-$ | $< 1.2 \text{ fb}$ | $< 1.2 \text{ fb}$ | [3] | $< 5 \text{ fb}$ | $< 5 \text{ fb}$ | [78] |
| $\tau^+\tau^-$ | $< 12 \text{ fb}$ | $< 15 \text{ fb}$ | [3] | $< 60 \text{ fb}$ | $< 67 \text{ fb}$ | [79] |
| $Z\gamma$ | $< 11 \text{ fb}$ | $< 11 \text{ fb}$ | [3] | $< 28 \text{ fb}$ | $< 40 \text{ fb}$ | [80] |
| ZZ | $< 12 \text{ fb}$ | $< 20 \text{ fb}$ | [3] | $< 200 \text{ fb}$ | $< 220 \text{ fb}$ | [81] |
| Zh | $< 19 \text{ fb}$ | $< 28 \text{ fb}$ | [3] | $< 116 \text{ fb}$ | $< 116 \text{ fb}$ | [82] |
| hh | $< 39 \text{ fb}$ | $< 42 \text{ fb}$ | [3] | $< 120 \text{ fb}$ | $< 110 \text{ fb}$ | [83] |
| W^+W^- | $< 40 \text{ fb}$ | $< 70 \text{ fb}$ | [3] | $< 300 \text{ fb}$ | $< 300 \text{ fb}$ | [84] |
| $t\bar{t}$ | $< 450 \text{ fb}$ | $< 600 \text{ fb}$ | [3] | | | |
| invisible | $< 0.8 \text{ pb}$ | - | [3] | | | |
| $b\bar{b}$ | $\lesssim 1 \text{ pb}$ | $\lesssim 1 \text{ pb}$ | [3] | | | |
| jj | $\lesssim 2.5 \text{ pb}$ | - | [3] | | | |

- On general grounds (SU(2)xU(1) gauge symmetry) we expect decays to ZZ and Z γ and maybe also WW. Other decay modes possible but more model dependent
- Current constraints allow cross section in other channels to be larger than diphoton one. Strongest constraints on dilepton cross section, comparable to diphoton one.
- Still, constraints non-trivial such that it's difficult to pump up X(750) GeV width by decays to SM particles. Exotic but not invisible decays needed.

How much mixing with Higgs?

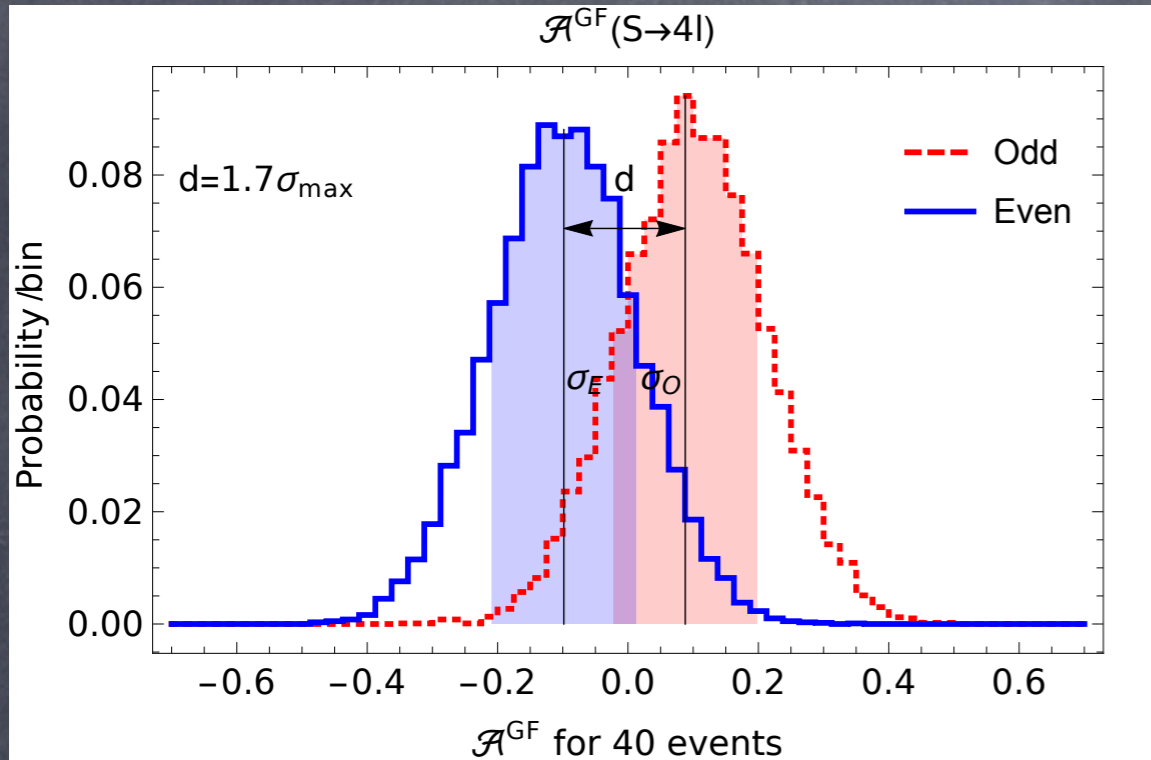
- For a singlet scalar, it is natural to mix with the Higgs boson
- Unless some symmetries or fine-tuning prevent it, mixing angle expected to be $\sin\alpha \sim mh^2/mS^2 \sim 1/30$
- For 750 GeV resonance, mixing angle strongly constrained by non-observation of WW and ZZ resonances



Parity and Spin studies

- Topic received (disproportionally) large attention in context of LHC Higgs studies
- It is much more interesting for 750 GeV case, as no preferred hypothesis a priori
- Good theoretical motivation for pseudo-scalars (e.g. pions of new technicolor-like sector coupled to photons via anomalies), as well as experimental one (mixing with Higgs suppressed)
- For spin ≥ 2 weaker theoretical motivation (basically that it'd be cool), and experimental one (currently based on rumors only)

$$\mathcal{L}_{P,\text{eff}} = \frac{P}{4v} \left(\tilde{c}_{pgg} g_s^2 G_{\mu\nu}^a \tilde{G}_{\mu\nu}^a + \tilde{c}_{pww} g_L^2 W_{\mu\nu}^i \tilde{W}_{\mu\nu}^i + \tilde{c}_{pbb} g_Y^2 B_{\mu\nu} \tilde{B}_{\mu\nu} \right)$$



$$\mathcal{A}^{\text{GF}} = \frac{N(\theta^{\text{GF}} > \pi/4) - N(\theta^{\text{GF}} < \pi/4)}{N(\theta^{\text{GF}} > \pi/4) + N(\theta^{\text{GF}} < \pi/4)},$$

where

$$\theta^{\text{GF}} = \begin{cases} \theta & \text{if } \theta < \pi/2 \\ \pi - \theta & \text{if } \theta > \pi/2 \end{cases},$$

and

$$\theta = \arccos \left\{ \frac{(p_1 \times p_2) \cdot (p_3 \times p_4)}{|p_1 \times p_2| |p_3 \times p_4|} \right\},$$

- Assuming spin 0, usual methods of parity determination inherited from Higgs study apply for 750 GeV
- One example: angle between decay planes of two Z bosons in $X \rightarrow ZZ \rightarrow 4l$ decays

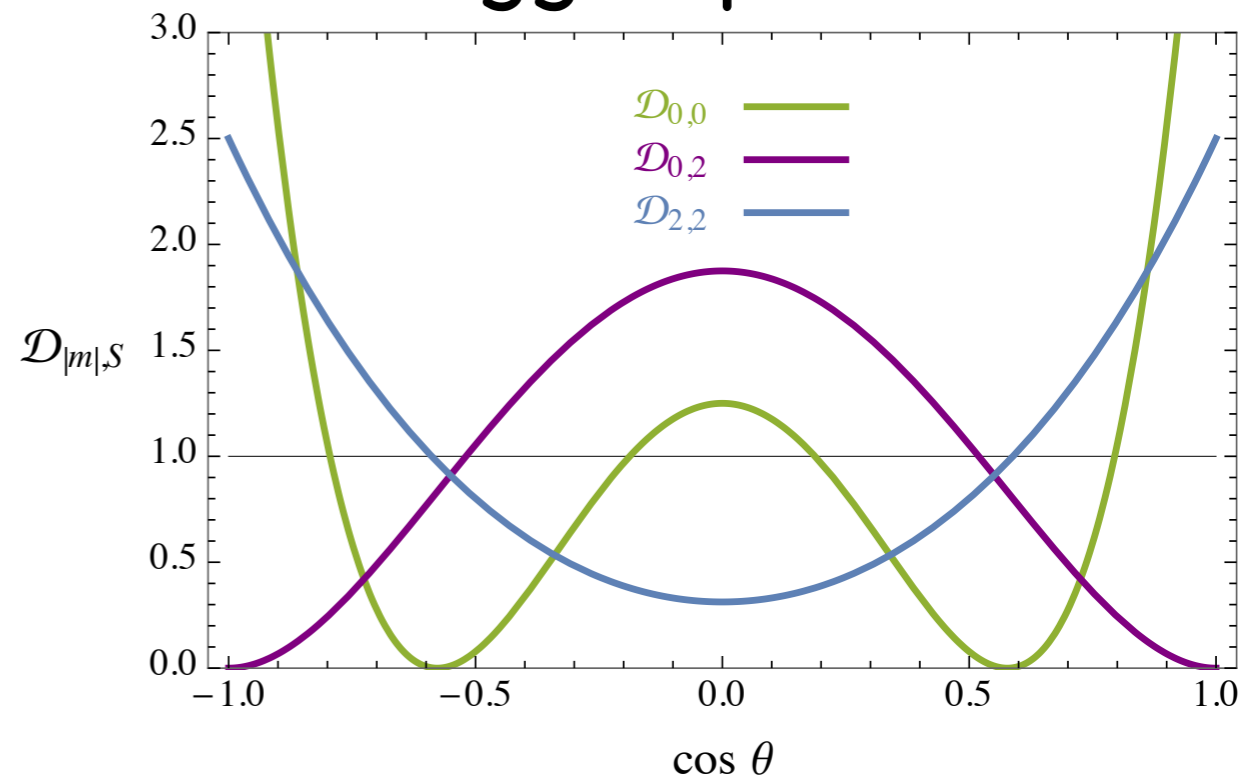
Spin Discrimination

$$\mathbf{J} = 0 \quad \mathcal{D}_{0,0}^{(0)} = 1$$

$$\mathbf{J} = 2 \quad \mathcal{D}_{|m|,S}^{(2)} = \begin{bmatrix} \frac{5}{4}(3c^2 - 1)^2 & \frac{15}{8}s^4 \\ \frac{15}{2}s^2c^2 & \frac{5}{4}s^2(1 + c^2) \\ \frac{15}{8}s^4 & \frac{5}{16}(1 + 6c^2 + c^4) \end{bmatrix}$$

$$\mathbf{J} = 3 \quad \mathcal{D}_{|m|,S}^{(3)} = \begin{bmatrix} \frac{7}{4}c^2(3 - 5c^2)^2 & \frac{105}{8}s^4c^2 \\ \frac{21}{16}s^2(5c^2 - 1)^2 & \frac{35}{32}s^2(1 - 2c^2 + 9c^4) \\ \frac{105}{8}s^4c^2 & \frac{7}{16}(4 - 15c^2 + 10c^4 + 9c^6) \end{bmatrix}$$

$gg \rightarrow \text{spin-2}$



- Spin-0 is trivial, spin-1 is impossible
- For spin-2 4 different distributions possible, with forward and/or central enhancement
- For KK graviton-like coupling to matter resonance produced in $m=2$ and decaying to $S=2$ diphoton state, leading to $D_{2,2}$ distribution with forward enhancement

Phenomenological model for spin-2 resonance

Kinetic terms (unique ghost free form)

$$\mathcal{L}_{\text{FP}} = \frac{1}{2}(\partial_\rho X_{\mu\nu})^2 - \frac{1}{2}(\partial_\rho X)^2 - (\partial_\rho X_{\mu\rho})^2 + \partial_\mu X \partial_\rho X_{\mu\rho} - \frac{m_X^2}{2}(X_{\mu\nu})^2 + \frac{m_X^2}{2}X^2$$

Interactions with matter: for each particle, coupling to its energy-momentum tensor
Since latter is dimension-4, spin-2 has dimension-5 non-renormalizable couplings

$$\begin{aligned} \mathcal{L}_{\text{int}} \supset & \frac{c_V}{v} X_{\mu\nu} \left(\frac{\eta_{\mu\nu}}{4} V_{\rho\sigma} V_{\rho\sigma} - V_{\mu\rho} V_{\nu\rho} \right), \\ & - \frac{ic_\chi}{4v} X_{\mu\nu} [\bar{\chi}(\bar{\sigma}_\mu \partial_\nu + \bar{\sigma}_\nu \partial_\mu)\chi - (\partial_\mu \bar{\chi} \bar{\sigma}_\nu + \partial_\nu \bar{\chi} \bar{\sigma}_\mu)\chi - 2\eta_{\mu\nu}(\bar{\chi} \bar{\sigma}_\rho \partial_\rho \chi - \partial_\rho \bar{\chi} \bar{\sigma}_\rho \chi)] \\ & + \frac{c_H}{v} X_{\mu\nu} (\partial_\mu H^\dagger \partial_\nu H + \partial_\nu H^\dagger \partial_\mu H - \eta_{\mu\nu} \partial_\rho H^\dagger \partial_\rho H + \eta_{\mu\nu} m_H^2 H^\dagger H + \eta_{\mu\nu} \lambda |H|^4) \end{aligned}$$

For ordinary massless graviton these couplings are universal
and suppressed by the Planck scale

$$c_H = c_V = c_\chi = \frac{v}{M_P} \approx 10^{-16},$$

But in general massive graviton couplings don't have to be universal,
and we know calculable examples

Spin-2: decay widths

see e.g.
Lee, Park, Sanz
1306.4107

- No chiral suppression for decays to fermions (unlike for scalars)
- For ZZ and WW, decays depends also on coupling to the Higgs field (because it contains longitudinal components of W and Z)
- For Z γ , decays occur only when coupling to WW and BB field strength is non-universal

$$c_{\gamma\gamma} = s_\theta^2 c_W + c_\theta^2 c_B, \quad c_{ZZ} = c_\theta^2 c_W + s_\theta^2 c_B,$$

$$c_{Z\gamma} = c_\theta s_\theta (c_W - c_B),$$

$$\Gamma(X \rightarrow hh) = \frac{c_H^2 m_X^3}{960\pi v^2} (1 - 4r_h)^{5/2}, \quad r_i \equiv m_i^2 / m_X^2$$

$$\Gamma(X \rightarrow f\bar{f}) = \frac{m_X^3}{320\pi v^2} (1 - 4r_f)^{3/2} \left[(c_{fL}^2 + c_{fR}^2) \left(1 - \frac{2r_f}{3}\right) + c_{fL} c_{fR} \frac{20r_f}{3} \right],$$

$$\Gamma(X \rightarrow ZZ) = \frac{m_X^3}{80\pi v^2} \sqrt{1 - 4r_Z} \left[c_{ZZ}^2 + \frac{c_H^2}{12} + \frac{r_Z}{3} (3c_H^2 + 20c_H c_{ZZ} - 9c_{ZZ}^2) + \frac{2r_Z^2}{3} (7c_H^2 - 10c_H c_{ZZ} + 9c_{ZZ}^2) \right],$$

$$\Gamma(X \rightarrow Z\gamma) = \frac{c_{Z\gamma}^2 m_X^3}{40\pi v^2} (1 - r_Z)^3 \left(1 + \frac{r_Z}{2} + \frac{r_Z^2}{6}\right),$$

$$\Gamma(X \rightarrow \gamma\gamma) = \frac{c_{\gamma\gamma}^2}{8c_G^2} \Gamma(X \rightarrow GG) = \frac{c_{\gamma\gamma}^2 m_X^3}{80\pi v^2},$$

Parameters for spin-2 resonance

$$\sigma(pp \rightarrow X)_{E_{\text{LHC}}} = \frac{\pi m_X^2}{v^2 E_{\text{LHC}}^2} \left[\frac{1}{16} k_{GGX} c_G^2 L_{GG} \left(\frac{m_X^2}{E_{\text{LHC}}^2} \right) + \frac{1}{24} \sum_q k_{qqX} (c_{qL}^2 + c_{qR}^2) L_{q\bar{q}} \left(\frac{m_X^2}{E_{\text{LHC}}^2} \right) \right]$$

Assuming gluon fusion production:

$$c_G \approx 3.1 \times 10^{-3} \sqrt{\frac{4.4 \times 10^{-2}}{\text{Br}(X \rightarrow \gamma\gamma)}}$$

| | | | | | |
|---|-----------|-----------|-----------|-----------|--------------------|
| $\text{Br}(X \rightarrow \gamma\gamma)$ | 10^{-1} | 10^{-2} | 10^{-3} | 10^{-4} | 2×10^{-7} |
| c_g | 0.0015 | 0.0049 | 0.015 | 0.049 | 1 |

For reasonable branching fractions to photons, scale suppressing spin-2 interactions with gluons should be in 1-100 TeV range

Thus, spin-2 explanations of diphoton anomaly are necessary effective theories with low cut-off

Predictions:

| f | $\text{Br}(X \rightarrow f)$ [%] | $\frac{\text{Br}(X \rightarrow f)}{\text{Br}(X \rightarrow \gamma\gamma)}$ |
|----------------|----------------------------------|--|
| $\gamma\gamma$ | 4.3 | 1 |
| ZZ | 4.0 | 0.9 |
| WW | 8.4 | 1.9 |
| $\mu\mu$ | 2.2 | 0.5 |
| jj | 67 | 15.5 |
| tt | 5.8 | 1.3 |
| bb | 5.5 | 1.5 |
| hh | 0.4 | 0.08 |

$$m_n \approx 3.8ka_L \begin{cases} 1 & 750 \text{ GeV} \\ 1.8 & 1.4 \text{ TeV} \\ 2.7 & 2 \text{ TeV} \\ \dots & \end{cases}$$

| final state f | σ at $\sqrt{s} = 8 \text{ TeV}$ | | | σ at $\sqrt{s} = 13 \text{ TeV}$ | | |
|----------------------|--|-------------------------|------|---|--------------------|------|
| | observed | expected | ref. | observed | expected | ref. |
| $e^+e^-, \mu^+\mu^-$ | $< 1.2 \text{ fb}$ | $< 1.2 \text{ fb}$ | [3] | $< 5 \text{ fb}$ | $< 5 \text{ fb}$ | [78] |
| $\tau^+\tau^-$ | $< 12 \text{ fb}$ | $< 15 \text{ fb}$ | [3] | $< 60 \text{ fb}$ | $< 67 \text{ fb}$ | [79] |
| $Z\gamma$ | $< 11 \text{ fb}$ | $< 11 \text{ fb}$ | [3] | $< 28 \text{ fb}$ | $< 40 \text{ fb}$ | [80] |
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| invisible | $< 0.8 \text{ pb}$ | - | [3] | | | |
| $b\bar{b}$ | $\lesssim 1 \text{ pb}$ | $\lesssim 1 \text{ pb}$ | [3] | | | |
| jj | $\lesssim 2.5 \text{ pb}$ | - | [3] | | | |

- Original RS model with the SM on the IR brane provides a self-consistent explanation of the 750 excess (up to providing mechanism for stabilizing radion)
- Very predictive model with no free parameters after fitting observations so far
- Tension with run-1 and run-2 dilepton resonance searches

Challenge for RS bulk

- In standard version of RS bulk, lightest gauge KK modes are a factor of 1.5 lighter than lightest graviton KK mode
- In present context this would mean gauge KK modes at 500 GeV
- Solutions: hide the light gauge modes, OR make graviton KK modes lighter by the use of gravity brane kinetic terms, OR both

AA,Kamenik
1603.06980

Hewett,Rizzo
1603.08250

Carmona
1603.08913

Dillon,Sanz
1603.09550

Parameters

$$a_L = 10^{-15}, \quad r_L = 10/k$$

| | MIN | MED | MAX |
|--|----------------------|----------------------|----------------------|
| $r_0[1/k]$ | 100 | 120 | 1700 |
| $M_*[\text{GeV}]$ | 4.1×10^{17} | 3.9×10^{17} | 1.6×10^{17} |
| α_{t_R} | ∞ | 0 | -0.3 |
| $\alpha_{Q_L^3}$ | ∞ | 0 | 0 |
| α_H | ∞ | 0 | -0.1 |
| $-c_G$ | 2.3×10^{-3} | 2.5×10^{-3} | 9.6×10^{-3} |
| $\sigma(pp \rightarrow X)[\text{pb}]$ | 0.06 | 0.08 | 1.1 |
| $\sigma(pp \rightarrow X \rightarrow \gamma\gamma)[\text{fb}]$ | 5.3 | 5.3 | 5.4 |
| $\Gamma_X[\text{GeV}]$ | 2×10^{-3} | 3×10^{-3} | 0.5 |

Remaining fermions localized at UV brane

Branching fractions

| | IR | MIN | MED | MAX | GMAX |
|-----------------------|-----|-----|-----|-----|------|
| $\gamma\gamma$ | 4.3 | 8.5 | 7.0 | 0.5 | 2.3 |
| ZZ | 4.8 | 7.9 | 7.8 | 2.9 | 12 |
| WW | 9.5 | 16 | 15 | 5.6 | 21 |
| $Z\gamma$ | 0 | 0 | 0 | 0 | 1.1 |
| hh | 0.3 | 0 | 0.4 | 1.4 | 6.9 |
| tt | 5.1 | 0 | 8.3 | 85 | 56 |
| bb | 6.4 | 0 | 5.2 | 0.4 | 0.04 |
| jj | 66 | 68 | 61 | 4.5 | 0.5 |
| $e^+e^- + \mu^+\mu^-$ | 4.3 | 0 | 0 | 0 | 0 |

Output $m_{X_1} = 750 \text{ GeV}, \quad m_{X_2} \approx 6 \text{ TeV}, \quad m_{V_1} \approx 2.9 \text{ TeV}$

- Other KK modes than 750 GeV spin-2 can be heavy enough to avoid detection
- Dilepton branching fraction is practically zero
- If Higgs and top localized toward IR, so as to solve hierarchy problem, large branching fraction to $t\bar{t}$, hh , ZZ , and WW predicted

Bigger picture?

Dirac Gaugino:
Benakli et al
1605.05313

Goldstino:
Torre,Petersson,
1512.05333

NMSSM+cascade decays
Ellwanger,Hugonie
1602.03344

SUSY

just MSMM:
Djouadi et al,
1605.01040

Dilaton:
CERN-th et al
1512.04933

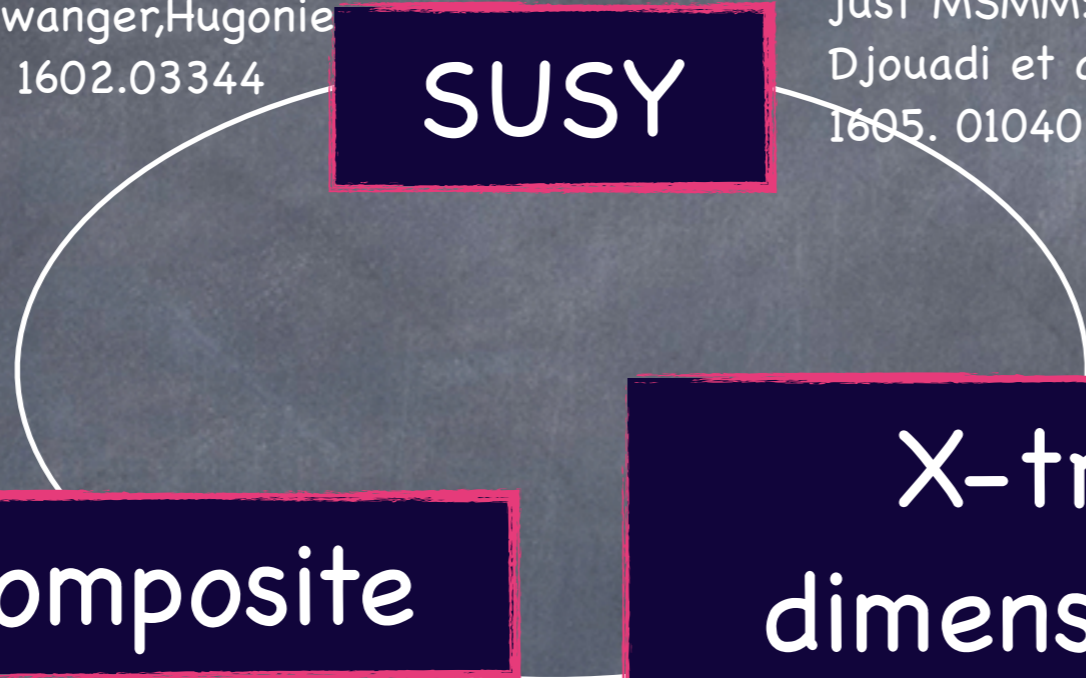
Composite

Hidden pion:
Harigaya,Nomura
1602.01092

**X-tra
dimensions**

KK Graviton:
Giddings,Zhang
1602.02793

Radion:
Ahmed et al
1512.05771

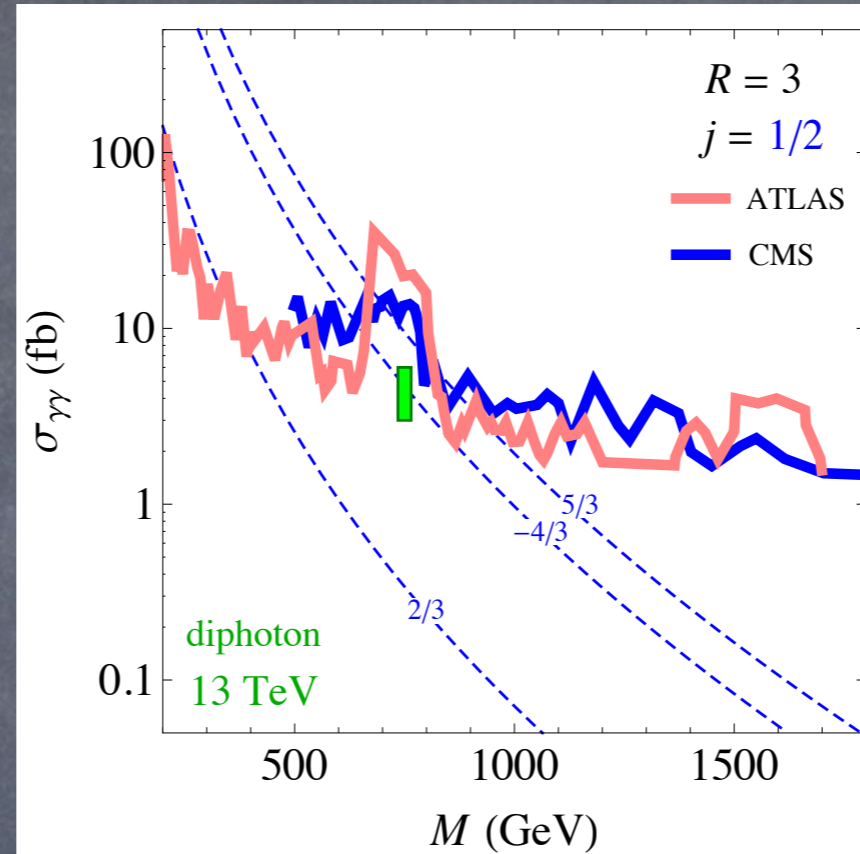
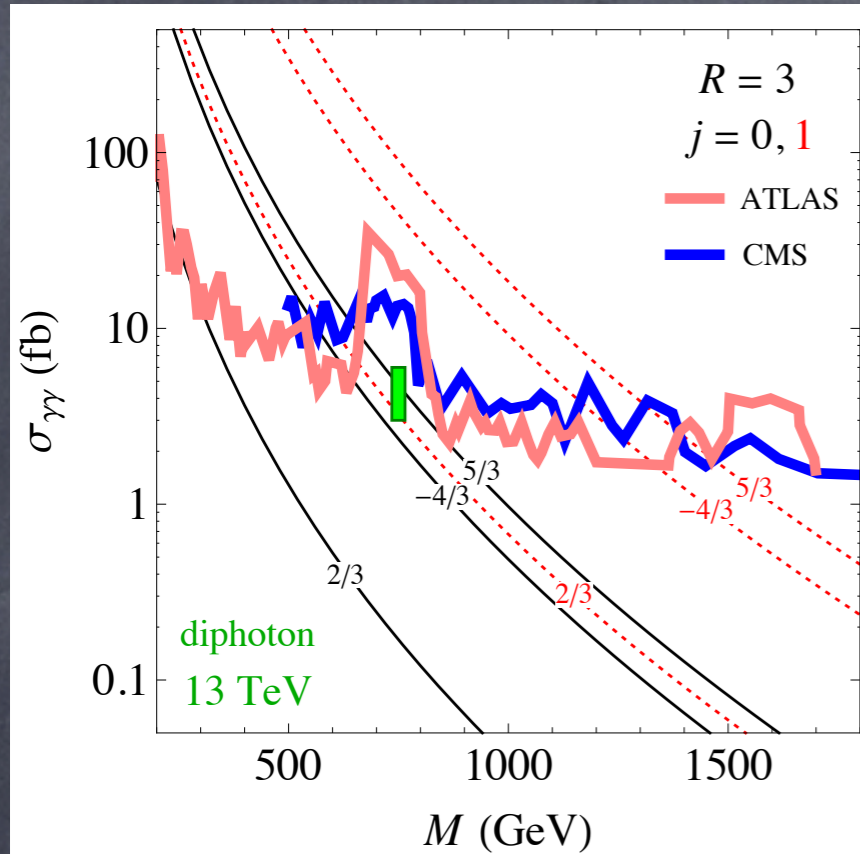


Bigger picture?

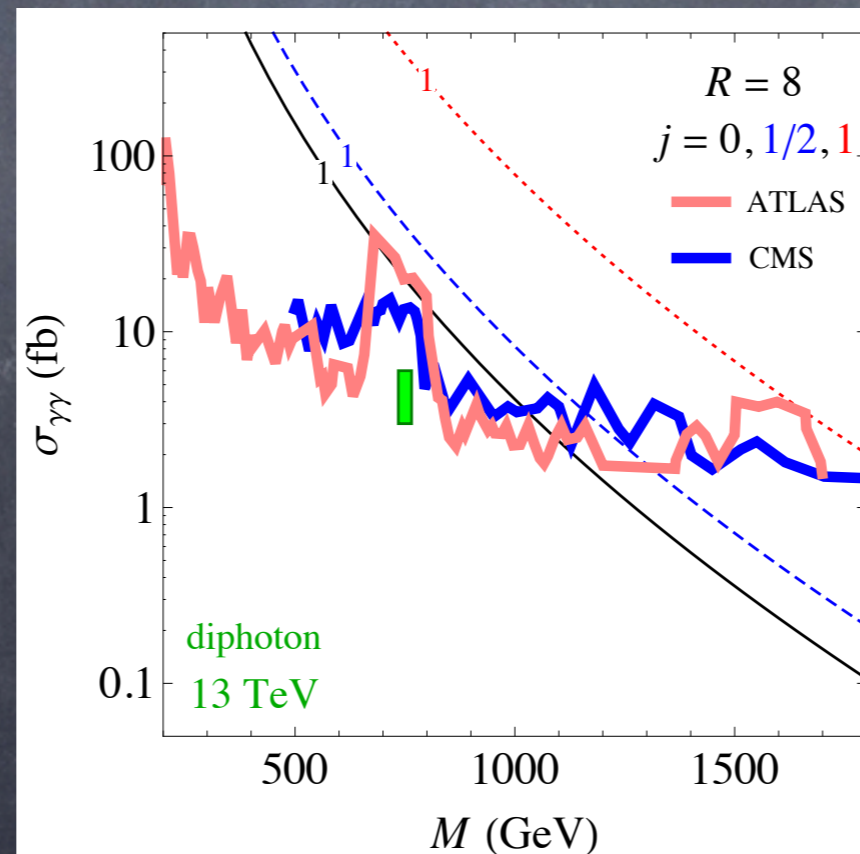
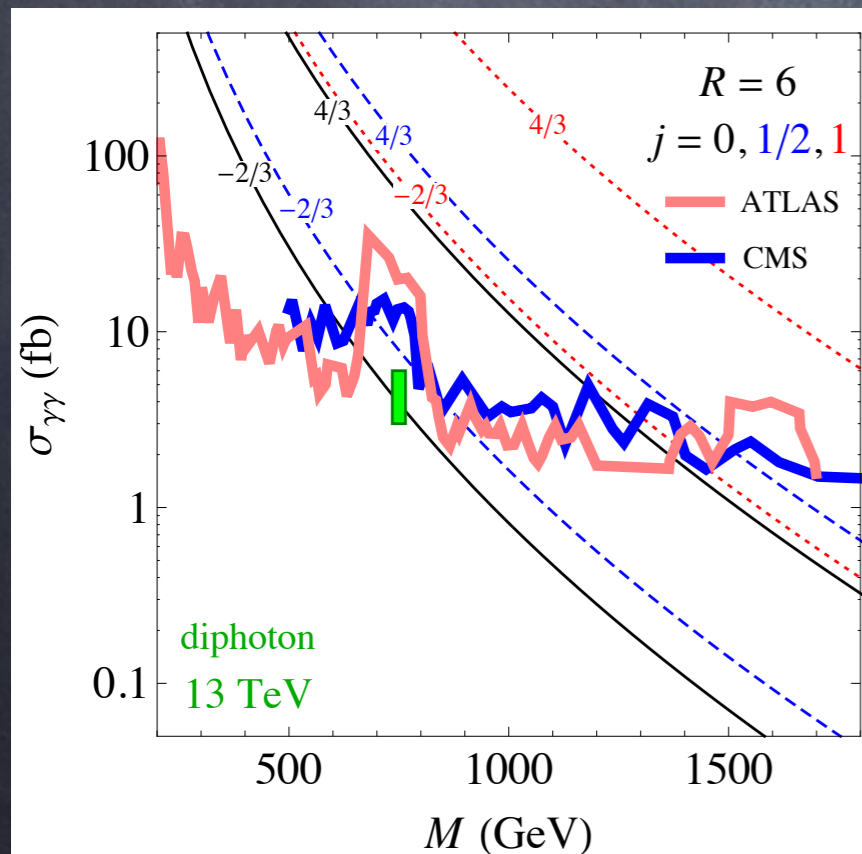
- In explicit models, large couplings are needed (for example, large Yukawa couplings of resonance to new vector-like fermions). Typically, these couplings run away to a Landau pole at a few TeV.
- Most natural embedding are into models with new strong interactions, that give rise to a light (pseudo-Goldstone?) composite state
- This strongly interacting sector may well have something to do with solving the hierarchy problem, as e.g. in little Higgs, composite Higgs, or Randall-Sundrum-type models.

Counterexample: just so?

Kats, Strassler
1602.08819

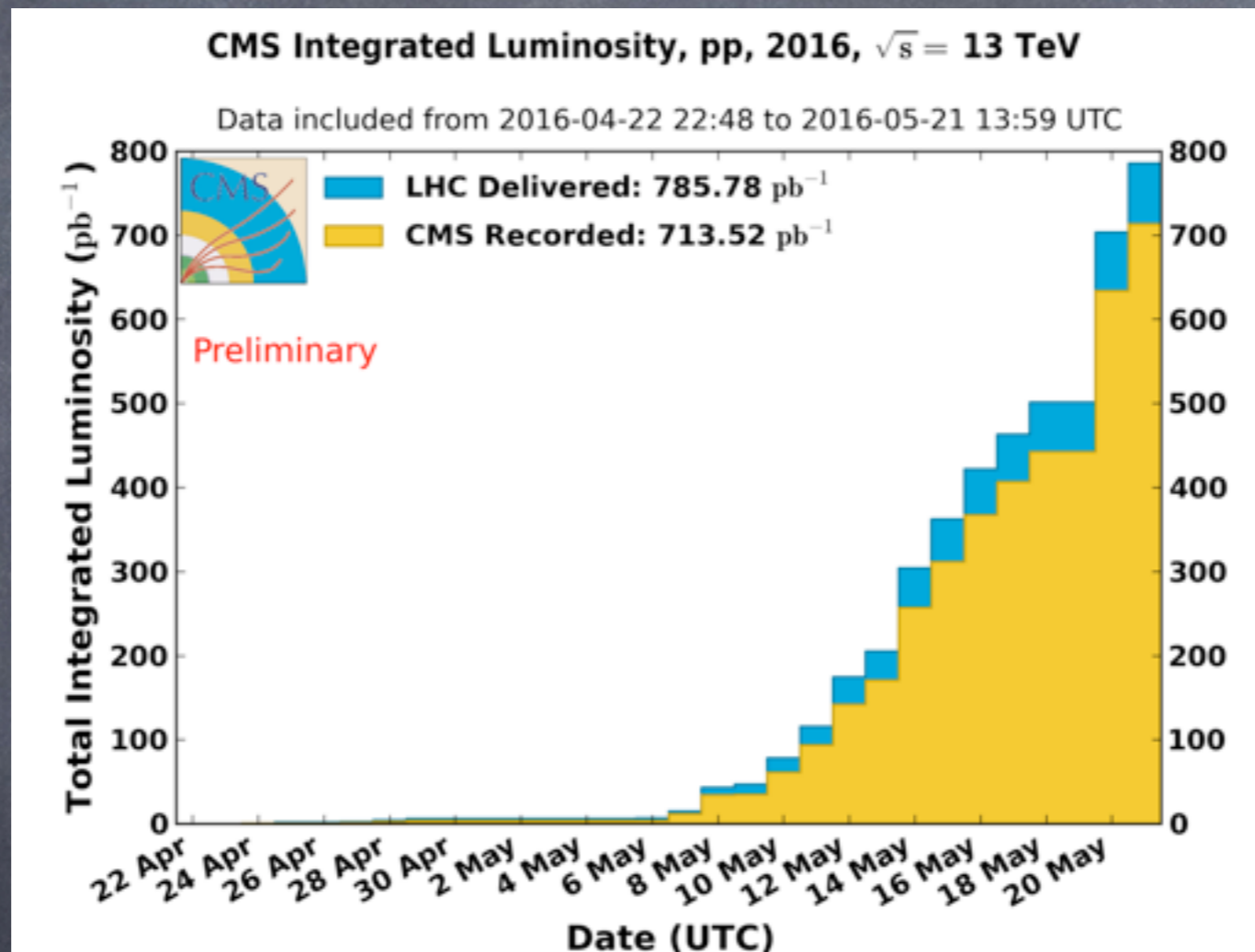


E.g., a bound state of charge $-4/3$ quarks can explain excess without new extended sector



Take away

- 750 GeV resonance needs to be confirmed by 2016 LHC data. For the moment, only "what if" speculations
- Several phenomenological models describing ATLAS and CMS observations exist, and they can be embedded in more motivated constructions



Already O(1-5) 750 GeV diphoton events in 2016 data
;) Have you looked yet? ;)