

MODEL OF DARK MATTER AND DARK ENERGY BASED ON GRAVITATIONAL POLARIZATION

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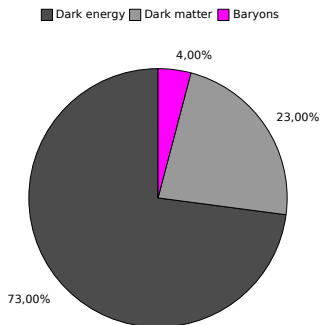
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Outline of the talk

- 1 Phenomenology of dark matter and MOND
- 2 Relativistic MOND theories
- 3 Gravitational polarization and MOND
- 4 Relativistic model of dark matter and dark energy

PHENOMENOLOGY OF DARK MATTER AND MOND

Concordance model in cosmology

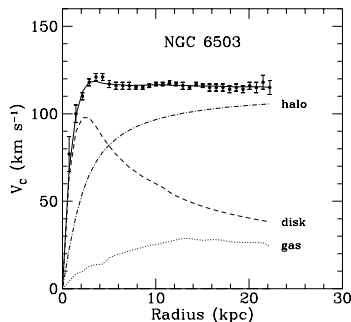


- Einstein field equations with a cosmological constant Λ are assumed to be correct and $\Omega_{\Lambda} = 73\%$ is measured from the Hubble diagram of supernovas
- Density of baryons $\Omega_b = 4\%$ is known from Big Bang nucleosynthesis and from CMB anisotropies
- Dark matter is **non-baryonic** and accounts for the observed dynamical mass of bounded astrophysical systems [Oort 1932, Zwicky 1933]

The dark matter paradigm [Bertone, Hooper & Silk 2004]

- ① Dark matter is **made of unknown non-baryonic particles**, e.g.
 - neutrinos (but $\Omega_\nu \lesssim 7\%$)
 - neutralinos (predicted by super-symmetric extensions of the standard model)
 - light scalar [Boehm, Cassé & Fayet 2003]
 - axions
 - ...
- ② It accounts for the observed mass discrepancy between the dynamical and luminous masses of clusters of galaxies
- ③ It **triggers the formation of large-scale structures by gravitational collapse** and predicts the scale-dependence of density fluctuations
- ④ It suggests some **universal dark matter density profile** around ordinary masses [Navarro, Frenk & White 1995]
- ⑤ It has difficulties at explaining naturally the **flat rotation curves of galaxies** and the **Tully-Fisher relation** [McGaugh & de Blok 1998, Sanders & McGaugh 2002]

Rotation curves of galaxies are flat

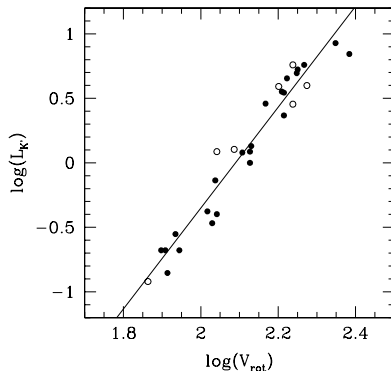


For a circular orbit
$$v(r) = \sqrt{\frac{GM(r)}{r}}$$

The fact that $v(r)$ is approximately constant implies that beyond the optical disk

$$M_{\text{halo}}(r) \approx r \quad \rho_{\text{halo}}(r) \approx \frac{1}{r^2}$$

The Tully-Fisher empirical relation [Tully & Fisher 1977]



The relation between the asymptotic flat velocity and the luminosity of spirals is

$$v_{\text{flat}} \propto L^{1/4}$$

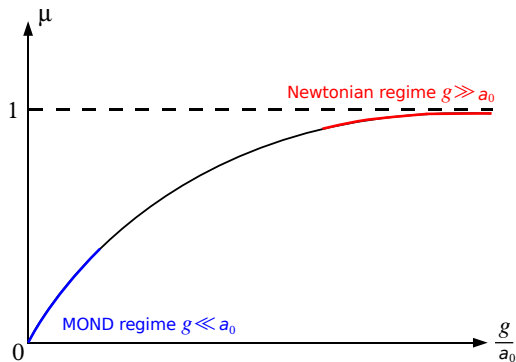
Modified Newtonian dynamics (MOND) [Milgrom 1983]

- ① MOND was proposed as an alternative to the dark matter paradigm
- ② It states that **there is no dark matter** and **we witness a violation of the fundamental law of gravity**
- ③ MOND is designed to account for the basic phenomenology of the **flat rotation curves of galaxies** and the **Tully-Fisher relation**
- ④ The basic fact about MOND is that the apparent need for dark matter arises in regions of **weak gravitational field** $g \ll a_0$

The Newtonian gravitational field is modified in an *ad hoc* way

$$\mu\left(\frac{g}{a_0}\right) \mathbf{g} = \mathbf{g}_{\text{Newton}}$$

The MOND function



In the MOND regime (when $g \ll a_0$) we have [Milgrom 1983]

$$\mu\left(\frac{g}{a_0}\right) \approx \frac{g}{a_0}$$

Recovering flat rotation curves and the Tully-Fisher law

- ① For a spherical mass $g_N = \frac{GM}{r^2}$ hence $g \approx \frac{\sqrt{GM a_0}}{r}$
- ② For circular motion $\frac{v^2}{r} = g$ thus v is constant and we get

$$v_{\text{flat}} \approx (GM a_0)^{1/4}$$

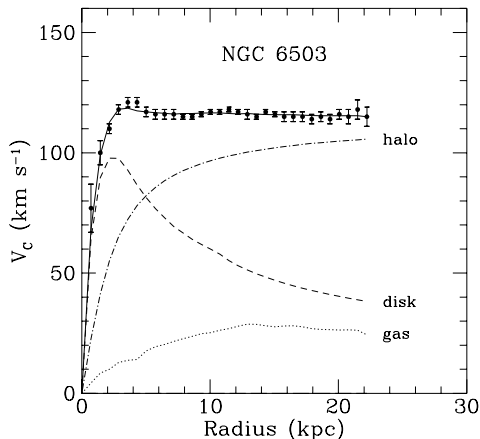
- ③ Assuming $L/M = \text{const}$ one naturally explains the Tully-Fisher relation
- ④ The numerical value of the critical acceleration is measured as

$$a_0 \approx 1.2 \cdot 10^{-10} \text{ m/s}^2$$

This value of a_0 is mysteriously close to the acceleration scale associated with the cosmological constant Λ

$$a_0 \approx 1.3 a_\Lambda \quad \text{where} \quad a_\Lambda = \frac{1}{2\pi} \left(\frac{\Lambda}{3} \right)^{1/2} \quad [\text{Gibbons \& Hawking 1977}]$$

The MOND fit of galactic rotation curves



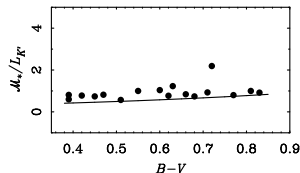
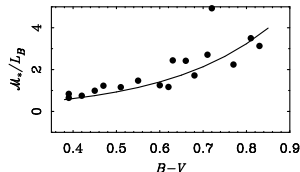
← the solid line is the MOND fit

Many galaxies are accurately fitted with MOND [Sanders 1996, McGaugh & de Blok 1998]

Fit of the mass-to-luminosity ratio

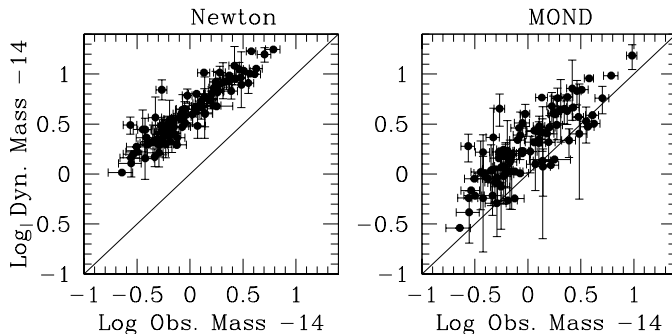
The fit of rotation curves is actually a one-parameter fit. The mass-to-luminosity ratio M/L of each galaxy is adjusted (and is therefore measured by MOND)

M/L shows the same trend with colour as is implied by models of population synthesis [Sanders & Verheijen 1998]



Problem with galaxy clusters

[Gerbal, Durret *et al* 1992, Sanders 1999]



The mass discrepancy is $\approx 4 - 5$ with Newton and ≈ 2 with MOND

RELATIVISTIC MOND THEORIES

Non-relativistic theory for MOND

- The MOND equation can be rewritten as the modified Poisson equation

$$\nabla \cdot \left[\underbrace{\mu \left(\frac{g}{a_0} \right)}_{\text{MOND function}} \mathbf{g} \right] = -4\pi G \rho$$

where $\mathbf{g} = \nabla U$ is the gravitational field and ρ the density of ordinary matter

- This equation is derivable from the Lagrangian [Bekenstein & Milgrom 1984]

$$L = \frac{a_0^2}{8\pi} k \left(\frac{\nabla U^2}{a_0^2} \right) + \rho U$$

where k is related to the MOND function by $k'(x) = \mu(\sqrt{x})$.

Relativistic scalar-tensor theory for MOND [Bekenstein & Sanders 1994]

- ① The tensor part is the usual Einstein-Hilbert action

$$L_g = \frac{1}{16\pi} R[g, \partial g, \partial^2 g]$$

- ② The scalar field is given by an aquadratic kinetic term

$$L_\phi = \frac{a_0^2}{8\pi} h \left(\frac{g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi}{a_0^2} \right)$$

- ③ The matter fields are universally coupled to gravity

$$L_m = L_m \left[\Psi, \underbrace{\tilde{g}_{\mu\nu} \equiv e^{2\phi} g_{\mu\nu}}_{\text{physical metric}} \right]$$

- Light signals do not feel the presence of the scalar field ϕ because the physical metric $\tilde{g}_{\mu\nu}$ is conformally related to the Einstein frame metric $g_{\mu\nu}$
- Since we observe huge amounts of dark matter by gravitational lensing (weak and strong) this theory is not viable

Tensor-vector-scalar theory [Bekenstein 2004, Sanders 2005]

- ① The tensor part is the Einstein-Hilbert action L_g
- ② The vector field part for $W_{\mu\rho} = \partial_\mu V_\rho - \partial_\rho V_\mu$ is

$$L_V = -\frac{1}{32\pi} \left[K \underbrace{g^{\mu\nu} g^{\rho\sigma} W_{\mu\rho} W_{\nu\sigma}}_{\text{standard spin-1 action}} - 2\lambda \underbrace{\left(g^{\mu\nu} V_\mu V_\nu + 1 \right)}_{\text{tells that } V^\mu \text{ is time-like and unitary}} \right]$$

- ③ The scalar action reads

$$L_\phi = -\frac{1}{2} \underbrace{\left(g^{\mu\nu} - V^\mu V^\nu \right)}_{\text{dynamical scalar field } \phi} \partial_\mu \phi \partial_\nu \phi - \underbrace{\frac{\sigma^4}{4\ell^2} F(k\sigma^2)}_{\text{non-dynamical scalar field } \sigma}$$

- ④ Matter fields are coupled to the non-conformally related physical metric

$$\tilde{g}_{\mu\nu} = e^{-2\phi} (g_{\mu\nu} + V_\mu V_\nu) - e^{2\phi} V_\mu V_\nu$$

This theory has evolved recently toward Einstein-æther like theories [Jacobson & Mattingly 2001; Zlosnik, Ferreira & Starkman 2007]

Particle dark matter versus MOND

The alternative seems to be

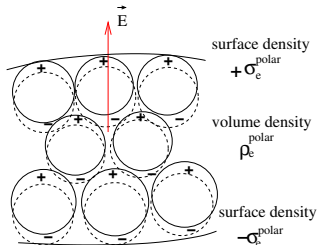
- ① Either accept the existence of cold dark matter but
 - made of unknown non-baryonic particles yet to be discovered
 - which fails to reproduce in a natural way the flat rotation curves of galaxies
- ② Or postulate a modification of the fundamental law of gravity: MOND and its relativistic extensions (TeVeS and Einstein-æther) but
 - on an *ad hoc* and not yet physically justified basis

Here we shall adopt a different approach

- ① Keep the standard law of gravity namely general relativity and its Newtonian limit when $c \rightarrow \infty$
- ② Use the phenomenology of MOND to guess what could be the (probably unorthodox) nature of dark matter
- ③ Propose a natural physical mechanism for MOND
- ④ Develop a relativistic matter model and apply it to cosmology

GRAVITATIONAL POLARIZATION AND MOND

The electric field in a dielectric medium



The atoms in a dielectric are modelled by electric dipole moments

$$\pi_e = q \xi$$

The polarization vector is

$$\Pi_e = n \pi_e$$

Density of polarization charges: $\rho_e^{\text{polar}} = -\nabla \cdot \Pi_e$

$$\nabla \cdot \mathbf{E} = \frac{\rho_e + \rho_e^{\text{polar}}}{\epsilon_0} \iff \nabla \cdot \left(\overbrace{\mathbf{E} + \frac{\Pi_e}{\epsilon_0}}^{\mathbf{D}} \right) = \frac{\rho_e}{\epsilon_0}$$

The dipoles and polarization vector are aligned with the electric field

$$\frac{\Pi_e}{\epsilon_0} = \underbrace{\chi_e(E)}_{\text{electric susceptibility}} \mathbf{E}$$

Interpretation of the MOND equation [Blanchet 2007]

The MOND equation in the form of a modified Poisson equation

$$\nabla \cdot \left[\underbrace{\mu\left(\frac{g}{a_0}\right)}_{\text{MOND function}} g \right] = -4\pi G \rho$$

is **formally analogous** to the equation of electrostatics inside a dielectric. We pose

$$\mu = 1 + \underbrace{\chi(g)}_{\text{gravitational susceptibility}} \quad \text{and} \quad \underbrace{\Pi}_{\text{gravitational polarization}} = -\frac{\chi}{4\pi G} g$$

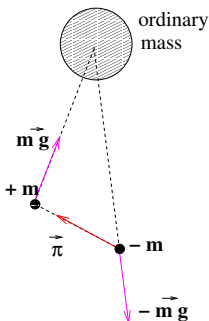
The MOND equation is equivalent to

$$\Delta U = -4\pi G (\rho + \rho_{\text{polar}})$$

In this interpretation the Newtonian law of gravity is not violated but we are postulating a **new form of dark matter consisting of “polarization masses”** with density

$$\rho_{\text{polar}} = -\nabla \cdot \Pi$$

Sign of the gravitational susceptibility



The “digravitational” medium is modelled by individual dipole moments π

$$\pi = m \xi$$

$$\Pi = n \pi$$

We suppose that the dipoles consist of particles doublets with

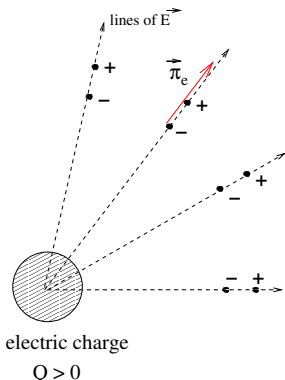
- opposite gravitational masses $m_g = \pm m$
- positive inertial masses $m_i = m$

- 1 The gravitational force is governed by a **negative Coulomb law** [Bondi 1957]
- 2 The dipoles tend to align in the same direction as the gravitational field thus

$$\chi < 0$$

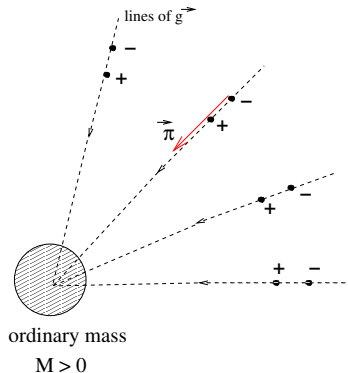
which is nicely compatible with MOND since $0 < \mu < 1 \implies -1 < \chi < 0$

Electric screening versus gravitational anti-screening



Screening by polarization charges

$$\chi_e > 0$$



Anti-screening by polarization masses

$$\chi < 0$$

Interpretation of the dark matter medium

- ① The dipole moments cannot be stable in a gravitational field and we must invoke some **non-gravitational internal force** (i.e. a “**fifth force**”) which is attractive between unlike masses
- ② The dark matter medium is a **gravitational plasma**, i.e. a medium composed of particles

$$(m_i, m_g) = (m, \pm m)$$

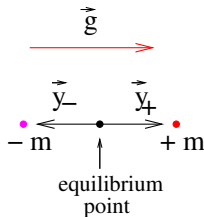
in equal numbers (so the plasma is globally neutral) and with local number densities n_{\pm}

- ③ We find that to obtain MOND we must postulate for the plasma's internal force the Gauss law

$$\nabla \cdot \mathbf{f}_{\text{int}} = -\frac{4\pi G m}{\chi} (n_+ - n_-)$$

where $\chi < 0$ is the susceptibility of the dipolar medium

Gravitational plasma oscillations



Let \mathbf{y}_{\pm} be the displacements from the equilibrium position at which the plasma is locally neutral, *i.e.* $n_+ = n_- = n$. The equation of motion is

$$m \frac{d^2 \mathbf{y}_{\pm}}{dt^2} = \pm m (\mathbf{f}_{\text{int}} + \mathbf{g})$$

Consider a small departure from equilibrium. The density perturbation reads $n_{\pm} = n (1 - \nabla \cdot \mathbf{y}_{\pm})$ to first order in \mathbf{y}_{\pm}

The dipole $\boldsymbol{\xi} = \mathbf{y}_+ - \mathbf{y}_-$ obeys the harmonic oscillator

$$\frac{d^2 \boldsymbol{\xi}}{dt^2} + \omega^2 \boldsymbol{\xi} = 2\mathbf{g}$$

in which ω is the usual **plasma frequency** given here by

$$\omega = \sqrt{-\frac{8\pi G m n}{\chi}}$$

Non viability of the quasi-Newtonian model

The quasi-Newtonian model

- Suggests that the gravitational analogue of the electric polarization is possible
- Yields a simple and natural explanation of the MOND equation
- Requires the existence of a new non-gravitational force
- Interprets the dark matter medium as a polarizable gravitational plasma

BUT THIS MODEL IS NOT VIABLE

- Is not relativistic and not Lagrangian-based
- Involves negative gravitational masses so violates the equivalence principle
- Does not allow to answer questions related to cosmology

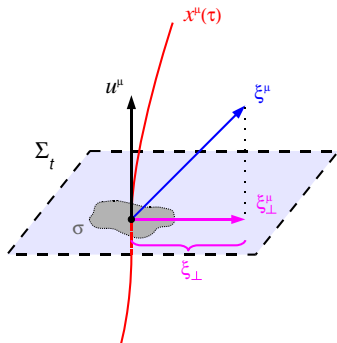
RELATIVISTIC MODEL OF DARK MATTER AND DARK ENERGY

Action principle for dipolar fluid

We propose a matter action in standard general relativity of the type

$$S = \int d^4x \sqrt{-g} L \left[J^\mu, \xi^\mu, \dot{\xi}^\mu, g_{\mu\nu} \right]$$

where the **current density** J^μ and the **dipole moment** ξ^μ are two independent dynamical variables



- The current density $J^\mu = \sigma u^\mu$ is conserved

$$\nabla_\mu J^\mu = 0$$

- The covariant time derivative is denoted

$$\dot{\xi}^\mu \equiv \frac{D\xi^\mu}{d\tau} = u^\nu \nabla_\nu \xi^\mu$$

- Projection perpendicular to the velocity

$$\xi_\perp^\mu \equiv \perp_{\mu\nu} \xi^\nu$$

Dipolar dark matter Lagrangian [Blanchet & Le Tiec 2008]

$$L = \sigma \left[-1 - \sqrt{(u_\mu - \dot{\xi}_\mu)(u^\mu - \dot{\xi}^\mu)} + \frac{1}{2} \dot{\xi}_\mu \dot{\xi}^\mu \right] - \mathcal{W}(\Pi_\perp)$$

- ① The mass term σ represents the inertial mass of the dipole moments (say $\sigma = 2m n$)
- ② The second term is a potential term inspired by the action of particles with spins in general relativity [Papapetrou 1951, Bailey & Israel 1980]
- ③ The third term is a kinetic-like term for the dipole moment and will tell how its evolution differs from parallel transport
- ④ The potential function \mathcal{W} depends on the **polarization field**

$$\Pi_\perp = \sigma \xi_\perp$$

and describes a non-gravitational force \mathcal{F}^μ internal to the dipole moment

Equations of motion and evolution

Varying the action with respect to the dynamical variables J^μ and ξ^μ we find that one equation can be solved with $\left[(u_\mu - \dot{\xi}_\mu) (u^\mu - \dot{\xi}^\mu) \right]^{1/2} = 1$

Equation of motion of dipolar fluid

$$\underbrace{\dot{u}^\mu = -\mathcal{F}^\mu}_{\text{non-geodesic motion}} \quad \text{where} \quad \underbrace{\mathcal{F}^\mu \equiv \hat{\xi}_\perp^\mu \frac{d\mathcal{W}}{d\Pi_\perp}}_{\text{dipolar internal force}}$$

Equation of evolution of dipole moment

$$\dot{\Omega}^\mu = \frac{1}{\sigma} \nabla^\mu \left(\mathcal{W} - \Pi_\perp \frac{d\mathcal{W}}{d\Pi_\perp} \right) - \underbrace{\xi_\perp^\nu R^\mu_{\rho\nu\sigma} u^\rho u^\sigma}_{\text{coupling to Riemann curvature}}$$

$$\text{where} \quad \Omega^\mu \equiv \perp_\nu^\mu \dot{\xi}_\perp^\nu + u^\mu \left(1 + \xi_\perp \frac{d\mathcal{W}}{d\Pi_\perp} \right)$$

The equations depend only on the (space-like) perpendicular projection $\xi_\perp^\mu = \perp_\nu^\mu \xi^\nu$ which represents the **physical dipole moment variable**

Stress-energy tensor

Varying the action with respect to the metric we obtain

$$T^{\mu\nu} = \underbrace{r}_{\text{energy density}} u^\mu u^\nu + \underbrace{\mathcal{P}}_{\text{pressure}} \perp^{\mu\nu} + 2 \underbrace{Q^{(\mu}}_{\text{heat flux}} u^{\nu)} + \underbrace{\Sigma^{\mu\nu}}_{\text{anisotropic stresses}}$$

where

$$\begin{aligned} r &= \rho + \mathcal{W} - \Pi_\perp \mathcal{W}' \\ \mathcal{P} &= -\mathcal{W} + \frac{2}{3} \Pi_\perp \mathcal{W}' \\ Q^\mu &= \sigma \dot{\xi}_\perp^\mu + \Pi_\perp \mathcal{W}' u^\mu - \Pi_\perp^\lambda \nabla_\lambda u^\mu \\ \Sigma^{\mu\nu} &= \left(\frac{1}{3} \perp^{\mu\nu} - \hat{\xi}_\perp^\mu \hat{\xi}_\perp^\nu \right) \Pi_\perp \mathcal{W}' \end{aligned}$$

The mass density ρ contains the rest mass σ and a dipolar term $-\nabla_\lambda \Pi_\perp^\lambda$ which appears as a relativistic generalisation of the polarization mass density

$$\rho = \sigma - \nabla_\lambda \Pi_\perp^\lambda$$

Cosmological perturbations at large scales

- 1 Apply first-order perturbation theory around FLRW background

$$g_{\mu\nu} = \bar{g}_{\mu\nu} + \delta g_{\mu\nu}$$

$$u^\mu = \bar{u}^\mu + \delta u^\mu$$

$$\xi_\perp^\mu = 0 + \delta \xi_\perp^\mu \quad \Leftarrow \text{dipole moment is perturbative}$$

$$\Pi_\perp = 0 + \delta \Pi_\perp \quad \Leftarrow \text{polarization field is perturbative}$$

- 2 Suppose that the potential takes the form

$$\mathcal{W}(\Pi_\perp) = \mathcal{W}_0 + \frac{1}{2} \mathcal{W}_2 \Pi_\perp^2 + \mathcal{O}(\Pi_\perp^3)$$

- 3 Apply standard **SVT gauge-invariant formalism** with $\delta \xi_\perp^\mu = (0, D^i y + y^i)$

- equations of motion

$$V' + \mathcal{H} V + \Phi = -\mathcal{W}_2 \bar{\sigma} a^2 y$$

$$V_i' + \mathcal{H} V_i = -\mathcal{W}_2 \bar{\sigma} a^2 y_i$$

- evolution equations

$$y'' + \mathcal{H} y' - \mathcal{W}_2 \bar{\sigma} a^2 y = 0$$

$$y_i'' + \mathcal{H} y_i' - \mathcal{W}_2 \bar{\sigma} a^2 y_i = 0$$

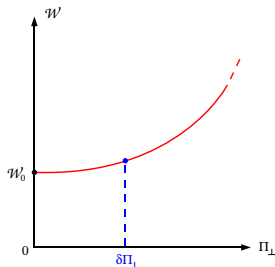
Cosmological expansion of the fundamental potential

We assumed that the potential \mathcal{W} reaches a minimum for the FLRW background

$$\mathcal{W}(\Pi_{\perp}) = \mathcal{W}_0 + \frac{1}{2}\mathcal{W}_2 \Pi_{\perp}^2 + \mathcal{O}(\Pi_{\perp}^3)$$

$$\mathcal{F}^{\mu} = \underbrace{\mathcal{W}_2 \Pi_{\perp}^{\mu}}_{\text{contributes at first-order cosmological perturbation}} + \mathcal{O}(\Pi_{\perp}^2)$$

contributes at first-order
cosmological perturbation



The minimum of the “fundamental” potential is the cosmological constant

$$\mathcal{W}_0 = \frac{\Lambda}{8\pi}$$

Stress-energy tensor in linear cosmological perturbation

- ① To first order in the perturbation, $T^{\mu\nu} = T_{\text{de}}^{\mu\nu} + T_{\text{dm}}^{\mu\nu}$ where

$$T_{\text{de}}^{\mu\nu} = - \underbrace{\mathcal{W}_0}_{\text{cosmological constant}} g^{\mu\nu}$$

$$T_{\text{dm}}^{\mu\nu} = \rho u^\mu u^\nu + \underbrace{2Q^{(\mu} u^{\nu)}}_{\text{heat flux term}}$$

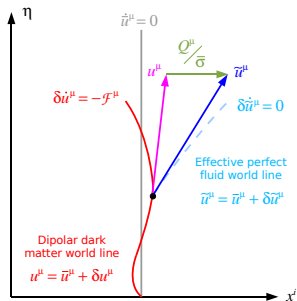
- ② Posing $\tilde{u}^\mu = u^\mu + Q^\mu/\bar{\rho}$ we can recast the dark matter fluid into a **perturbed pressureless perfect fluid** at linear perturbation order

$$T_{\text{dm}}^{\mu\nu} = \rho \tilde{u}^\mu \tilde{u}^\nu$$

- ③ The energy density of the dipolar dark matter fluid is

$$\rho = \underbrace{\sigma}_{\text{rest mass energy density}} - \underbrace{D_i \Pi_\perp^i}_{\text{dipolar polarization energy density}}$$

Agreement with the Λ -CDM scenario



The dipolar fluid is undistinguishable from

- **standard DE** (a cosmological constant)
- **standard CDM** (a pressureless perfect fluid)

at the level of first-order cosmological perturbations

Adjusting Λ so that $\Omega_{\text{de}} \approx 0.73$ and $\bar{\sigma}$ so that $\Omega_{\text{dm}} \approx 0.23$ the model is consistent with CMB fluctuations

Non-relativistic limit of the model ($c \rightarrow +\infty$)

- Equation of motion of the dipolar particle

$$\frac{d\mathbf{v}}{dt} = \underbrace{\mathbf{g}}_{\text{local gravitational field}} - \underbrace{\mathcal{F}}_{\text{internal force}}$$

- Evolution equation of the dipole moment

$$\frac{d^2 \boldsymbol{\xi}}{dt^2} = \mathcal{F} + \frac{1}{\sigma} \nabla (\mathcal{W} - \Pi \mathcal{W}') + \underbrace{(\boldsymbol{\xi} \cdot \nabla) \mathbf{g}}_{\text{tidal effect}}$$

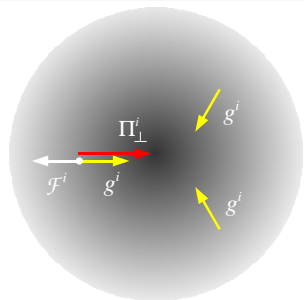
- Conservation of the number of particles

$$\partial_t \sigma = -\nabla \cdot (\sigma \mathbf{v})$$

- Poisson equation for the gravitational potential

$$\nabla \cdot \mathbf{g} = -4\pi \left(\underbrace{\rho_b}_{\text{baryonic matter}} + \sigma - \nabla \cdot \Pi \right)$$

Dipolar dark matter in a central gravitational field



In spherical symmetry we find that there is a solution for which the dipolar fluid is **static (i.e. at rest)**, the dipole moment is **aligned with the gravitational field**, and varies on a very long time-scale so that it is **practically in equilibrium**

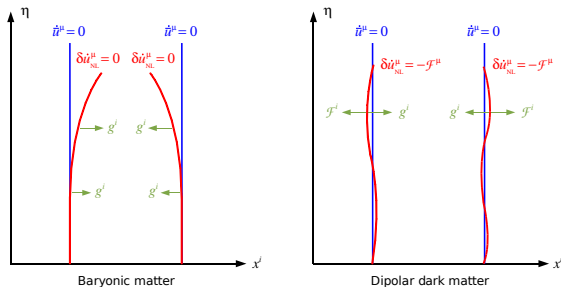
$$\boldsymbol{v} = \mathbf{0}$$

$$\boldsymbol{\Pi} = -\Pi(r) \boldsymbol{e}_r$$

In this solution the **internal force balances the gravitational field**

$$\boldsymbol{\mathcal{F}} = \boldsymbol{g}$$

Weak clustering of dipolar dark matter



- Baryonic matter follows the geodesic equation $\dot{u}^\mu = 0$, therefore collapses in regions of overdensity
- Dipolar dark matter obeys $\dot{u}^\mu = -\mathcal{F}^\mu$, with the internal force \mathcal{F}^i balancing the gravitational field g^i created by an overdensity

We expect that the mass density of dipolar dark matter in a galaxy at low redshift remains close to its mean cosmological value

$$\sigma \approx \bar{\sigma} \ll \rho_b \quad \text{and} \quad v \approx 0$$

Recovering MOND in a galaxy at low red-shift

Using $\mathbf{v} \approx \mathbf{0}$ in the equation of motion

$$\mathbf{g} = \hat{\Pi} \frac{d\mathcal{W}}{d\Pi} \quad \leftarrow \text{polarization of the dipolar medium}$$

Using $\sigma \approx \bar{\sigma} \ll \rho_b$ in the gravitational field equation

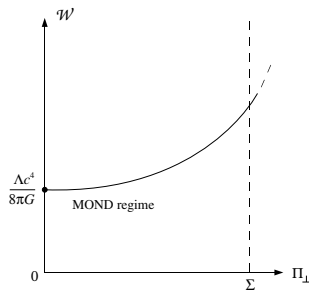
$$\nabla \cdot [\mathbf{g} - 4\pi \Pi] = -4\pi \rho_b \quad \leftarrow \text{modification of the Poisson equation}$$

- 1 Since the polarization Π is aligned with the gravitational field \mathbf{g} we recover the MOND equation with MOND function $\mu = 1 + \chi$ where

$$\Pi = -\frac{\chi}{4\pi} \mathbf{g}$$

- 2 There is one-to-one correspondence between μ and the potential \mathcal{W}

The fundamental potential \mathcal{W}



The potential \mathcal{W} is determined through third order

$$\mathcal{W} = \frac{\Lambda}{8\pi} + 2\pi \Pi_{\perp}^2 + \underbrace{\frac{16\pi^2}{3a_0} \Pi_{\perp}^3}_{\text{MOND acceleration appears here}} + \mathcal{O}(\Pi_{\perp}^4)$$

Though the model is purely *classical* it is tempting to interpret the cosmological constant Λ as a “**vacuum polarization**”, i.e. the residual polarization that remains when the “classical” part of the polarization $\Pi_{\perp} \rightarrow 0$

Order of magnitude of the cosmological constant

- 1 Introduce a purely numerical coefficient α such that

$$a_0 = \frac{1}{2\pi\alpha} \left(\frac{\Lambda}{3} \right)^{1/2} \quad \left(\text{i.e.} \quad a_0 = \frac{a_\Lambda}{\alpha} \right)$$

- 2 Write the fundamental potential as $\mathcal{W} = \frac{3\pi a_0^2}{2} f\left(\frac{\Pi_\perp}{a_0}\right)$ with

$$f(x) = \underbrace{\alpha^2 + \frac{4}{3}x^2 + \frac{32\pi}{9}x^3 + \mathcal{O}(x^4)}_{\text{some "universal" function of } x \equiv \Pi_\perp/a_0}$$

- 3 The numerical coefficients in $f(x)$ are expected to be of the order of one, hence the cosmological constant should be of the order of

$$\Lambda \sim a_0^2$$

in good agreement with the observations (which give $\alpha \approx 0.8$)

Conclusions

This model based on gravitational polarization:

- ① Offers a natural physical explanation of the phenomenology of MOND
- ② Benefits from both the successes of MOND at galactic scales and of Λ -CDM at cosmological scales
- ③ Yields a nice unification between the dark energy in the form of Λ and the dark matter *à la* MOND
- ④ But is “phenomenological” and presumably only valid in a regime of weak gravitational fields

More work should be done to:

- ① Connect the model to more fundamental (quantum) physics valid at microscopic scale
- ② Investigate second-order perturbations in cosmology
- ③ Numerically compute the non-linear growth of perturbations and compare with large-scale structures
- ④ Test the model at the intermediate scale of galaxy clusters