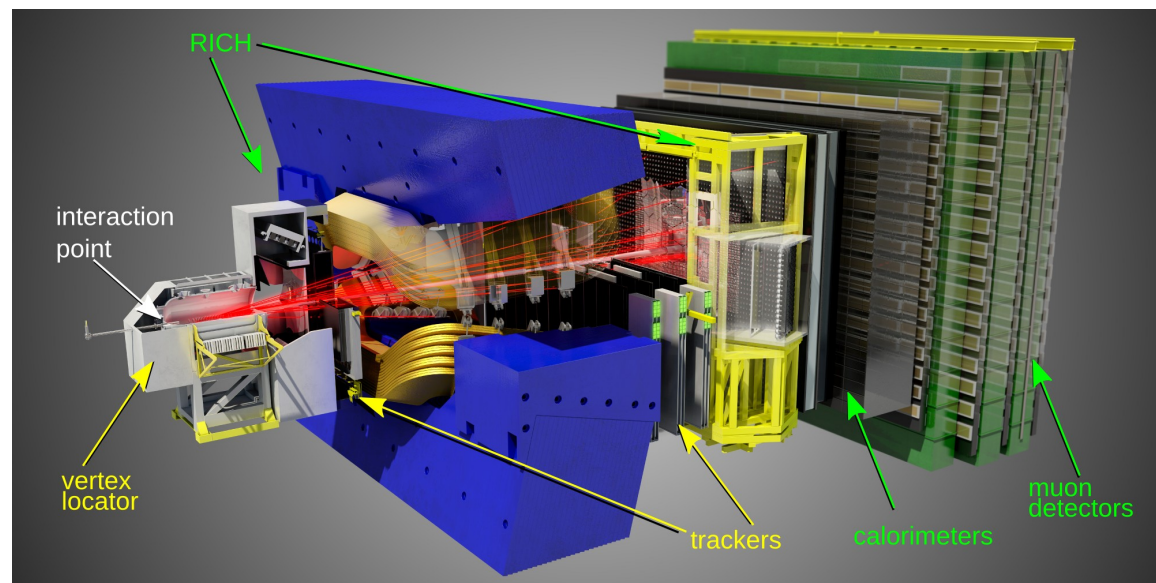
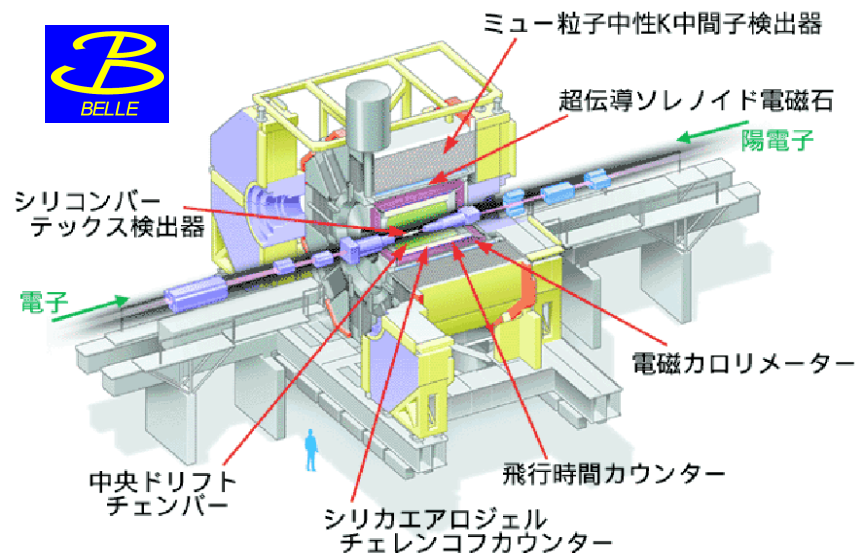


Flavour Physics at B-Factories and LHCb



K.Trabelsi

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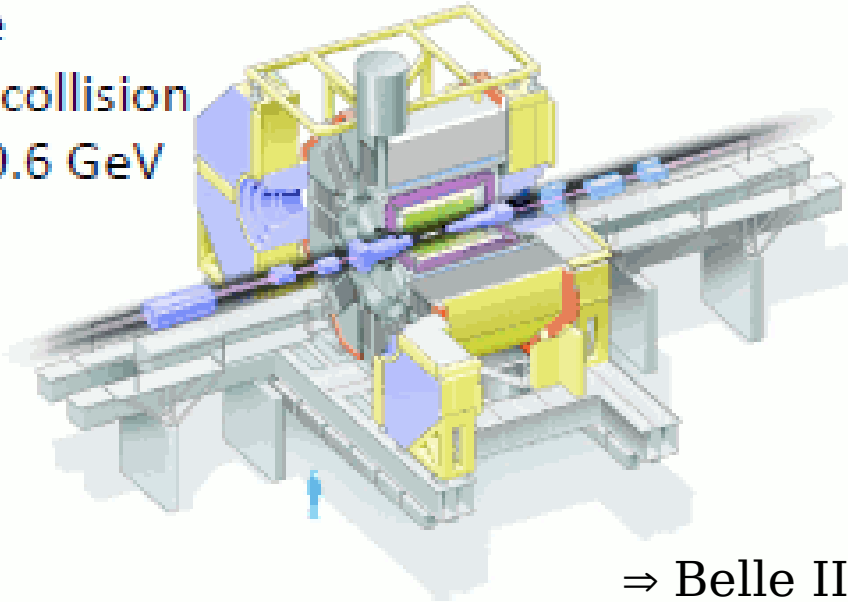
Seoul, May 19th 2016

A rich physics program...

- **Studies of CPV in B and B_s decays**
- **$b \rightarrow s$ transitions: probe for new sources of CPV and constraints from the $b \rightarrow s \gamma$ observables**
- **Forward-backward asymmetry and other observables in $b \rightarrow s l^+ l^-$**
- **Search for the charged Higgs in the rare decays $B \rightarrow \tau \nu$, $D^{(*)} \rightarrow \tau \nu$**
- Study of B_s , B_c , Λ_b decays
- Study of $D^0 - \bar{D}^0$ mixing
- Search for CPV in D and D_s decays
- Studies of exotic charmonium, tetraquark, pentaquark states
- Studies of new bottomonium-like states
- Search for lepton flavor violation (LFV) in τ decays
- Search for CPV and study of hadronic τ decays
- Light Higgs searches, DM searches...
- ...

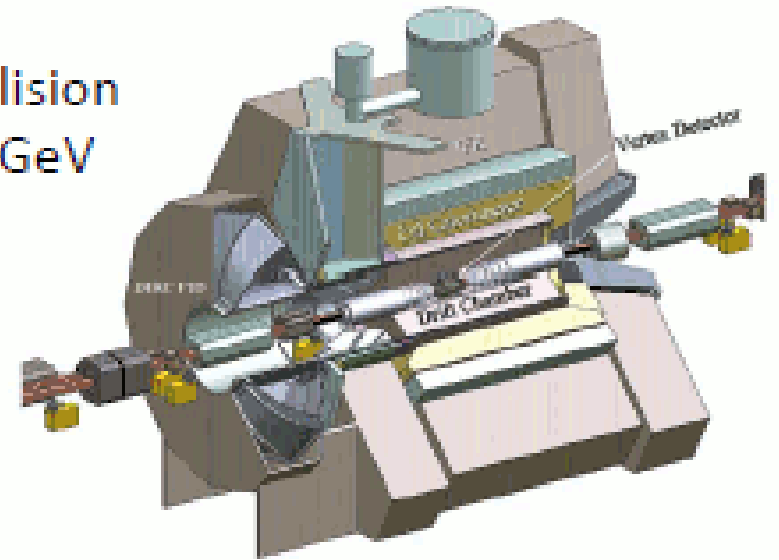
Main actors

Belle
 e^+e^- collision
at 10.6 GeV

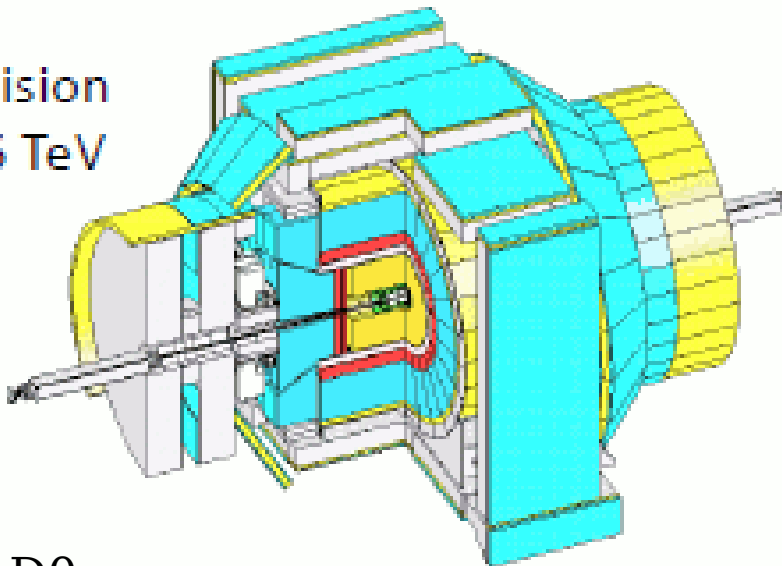


⇒ Belle II

BaBar
 e^+e^- collision
at 10.6 GeV

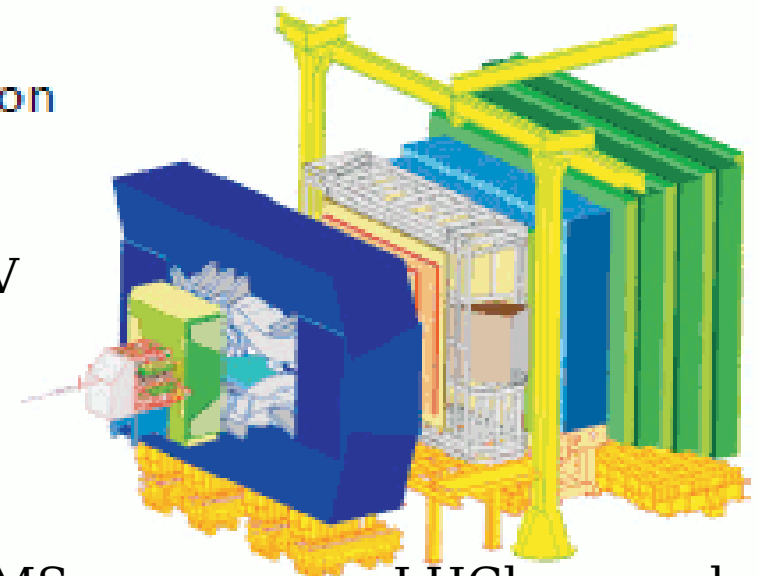


CDF
 $p\bar{p}$ collision
at 1.96 TeV



... and D0

LHCb
 pp collision
at 7 TeV
8 TeV
14 TeV



... and CMS

⇒ LHCb upgrade

Comparison B-factories/LHCb

B-factories

$$e^+ e^- \rightarrow Y(4S) \rightarrow b \bar{b}$$

at $Y(4S)$: 2 B's (B^0 or B^+) and nothing else \Rightarrow clean events

$$\sigma_{b\bar{b}} \sim 1 \text{ nb} \Rightarrow 1 \text{ fb}^{-1} \text{ produces } 10^6 \text{ B}\bar{\text{B}}$$

$$\sigma_{b\bar{b}}/\sigma_{\text{total}} \sim 1/4$$

LHCb

$$pp \rightarrow b \bar{b} X$$

production of B^+ , B^0 , B_s , B_c , Λ_b ...

but also a lot of other particles in the event

\Rightarrow lower reconstruction efficiencies

$\sigma_{b\bar{b}}$ much higher than at the $Y(4S)$

	\sqrt{s} [GeV]	$\sigma_{b\bar{b}}$ [nb]	$\sigma_{b\bar{b}}/\sigma_{\text{tot}}$
HERA pA	42 GeV	~ 30	$\sim 10^{-6}$
Tevatron	2 TeV	5000	$\sim 10^{-3}$
LHC	8 TeV	$\sim 3 \times 10^5$	$\sim 5 \times 10^{-3}$
	14 TeV	$\sim 6 \times 10^5$	$\sim 10^{-2}$

$b \bar{b}$ production cross-section $\sim 5 \times$ Tevatron, $\sim 500,000 \times$ BaBar/Belle !!

$\sigma_{b\bar{b}}/\sigma_{\text{total}}$ much lower than at the $Y(4S)$

\Rightarrow lower trigger efficiencies

B mesons live relatively long

mean decay length $\beta \gamma c \tau \sim 200 \mu\text{m}$

mean decay length $\beta \gamma c \tau \sim 7 \text{ mm}$

data taking period(s)

[1999-2010]

[run I: 2010-2012, run II: 2015-2018]

(near) future

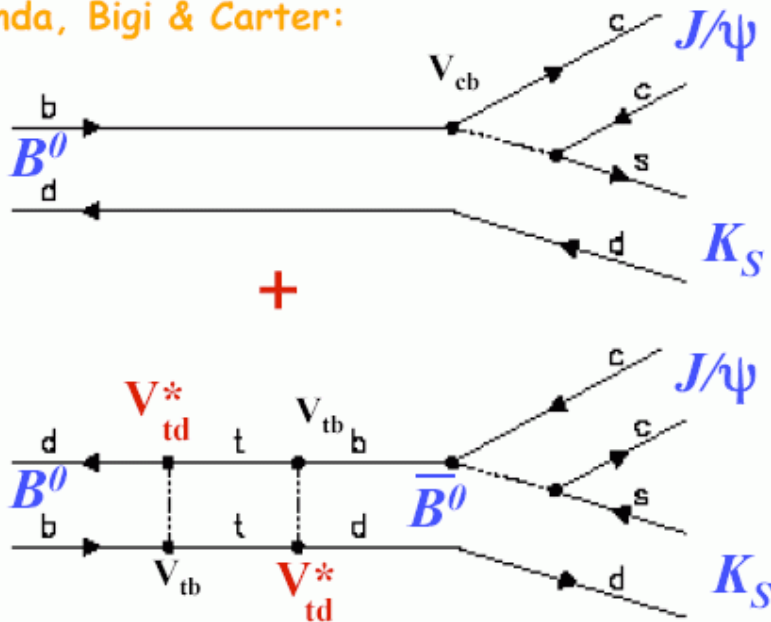
[Belle II from 2018]

[LHCb upgrade from 2020]

Time-dependent CP asymmetries in decays to CP eigenstates

$\sin 2\phi_1$ from $B \rightarrow f_{CP} + B \leftrightarrow \bar{B} \rightarrow f_{CP}$ interf.

Sanda, Bigi & Carter:



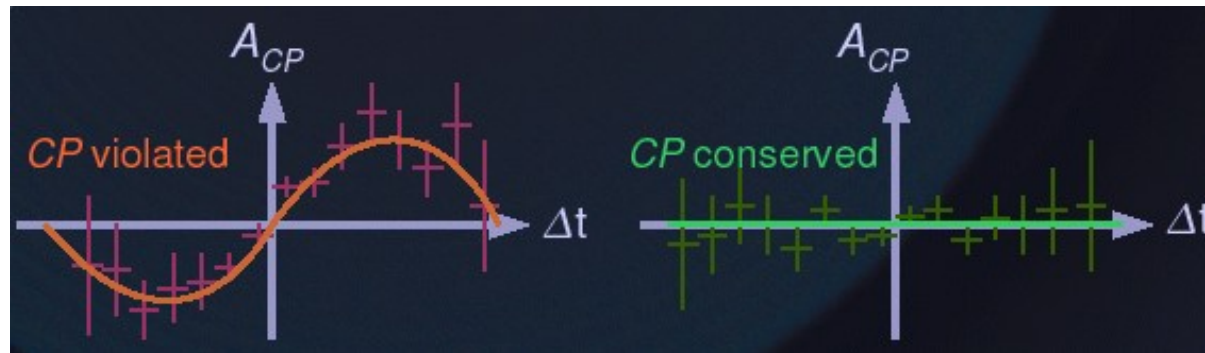
$$A_{CP}(f; t) = \frac{N(\bar{B}^0(t) \rightarrow f) - N(B^0(t) \rightarrow f)}{N(\bar{B}^0(t) \rightarrow f) + N(B^0(t) \rightarrow f)}$$

$$= \mathbf{S} \sin \Delta m_d t + \mathbf{A} \cos \Delta m_d t$$

$$= \frac{2 \operatorname{Im} \lambda}{|\lambda|^2 + 1} \sin \Delta m_d t + \frac{|\lambda|^2 - 1}{|\lambda|^2 + 1} \cos \Delta m_d t$$

$$\lambda = \frac{q}{p} \frac{A(\bar{B}^0 \rightarrow f)}{A(B^0 \rightarrow f)} = e^{-i2\phi_1} \frac{\bar{A}_f}{A_f}$$

- $\mathbf{A} = 0$ and $\mathbf{S} = -\xi_f \sin 2\beta$ for $(c\bar{c})K_{S/L}$ ($\xi_f = \mp 1$)
- $\mathbf{A} = 0$ and $\mathbf{S} = \sin 2\alpha$ for $\pi^+\pi^-$ (if tree only)

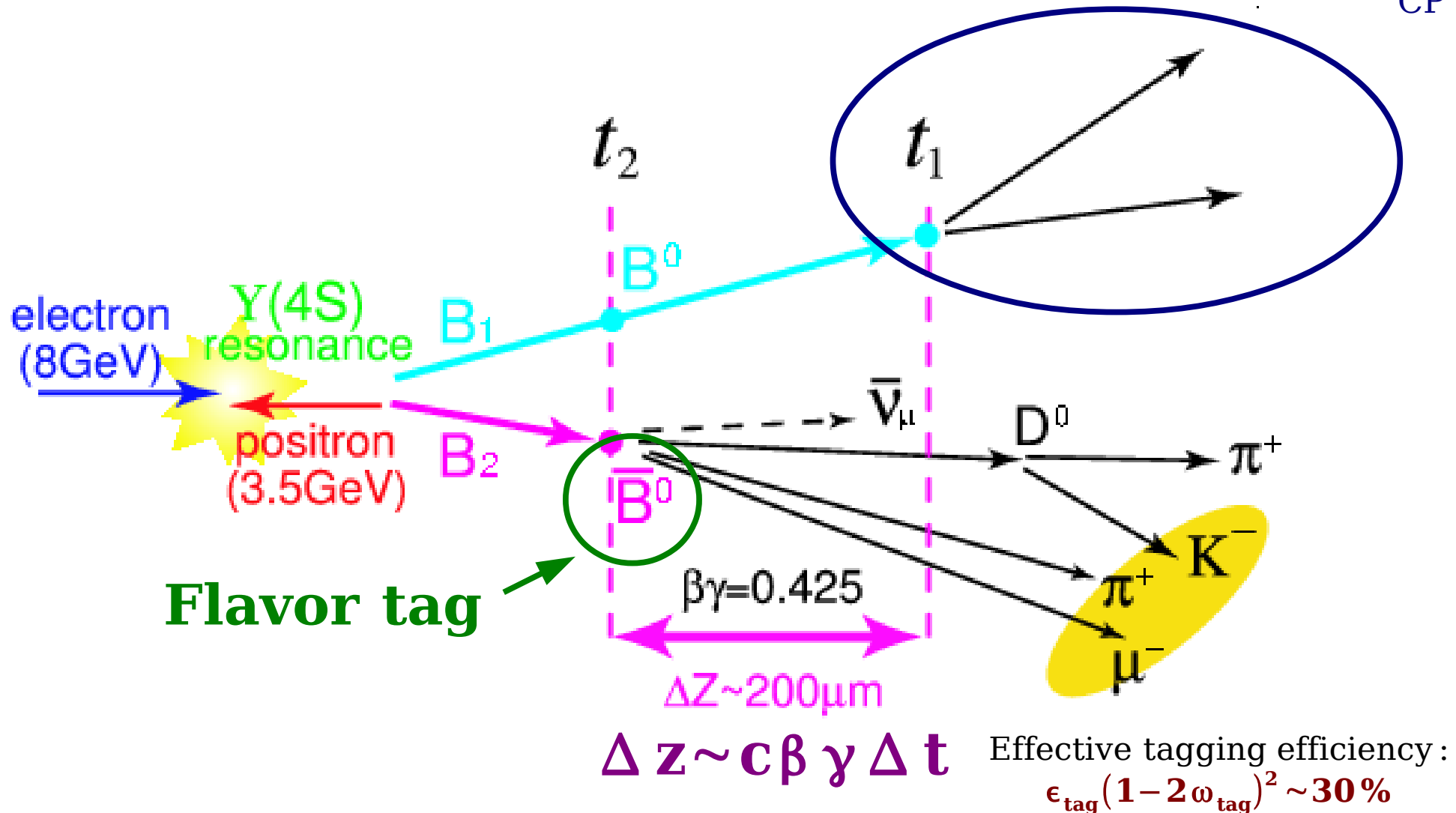


$$\mathbf{C} = -\mathbf{A}$$

Measuring the CP parameters S and A

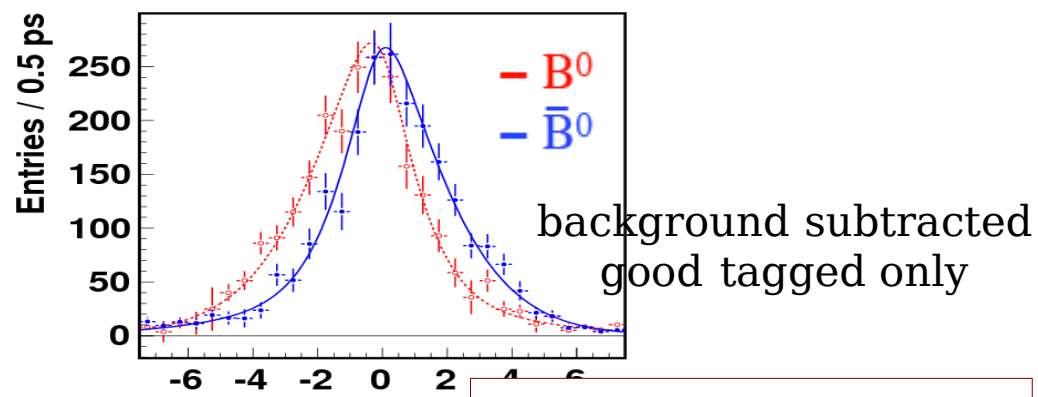
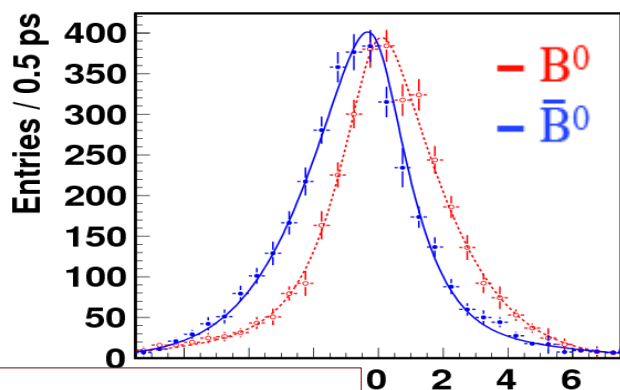
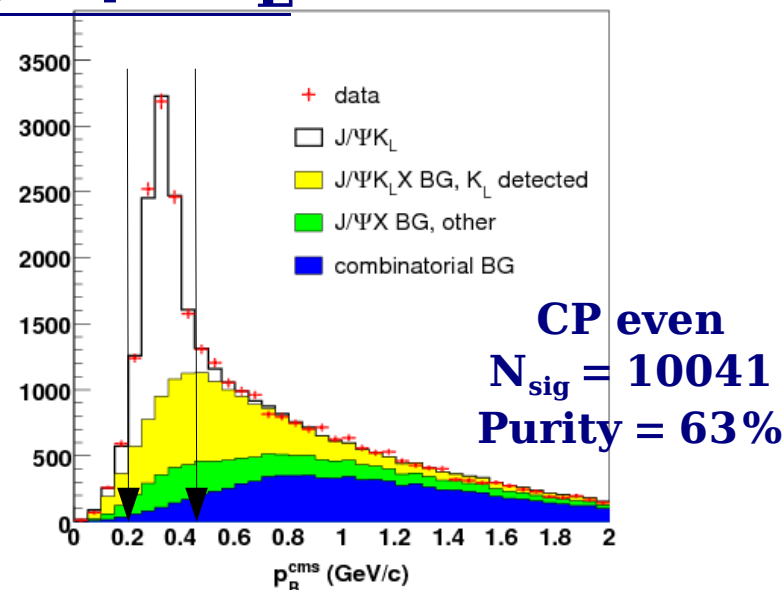
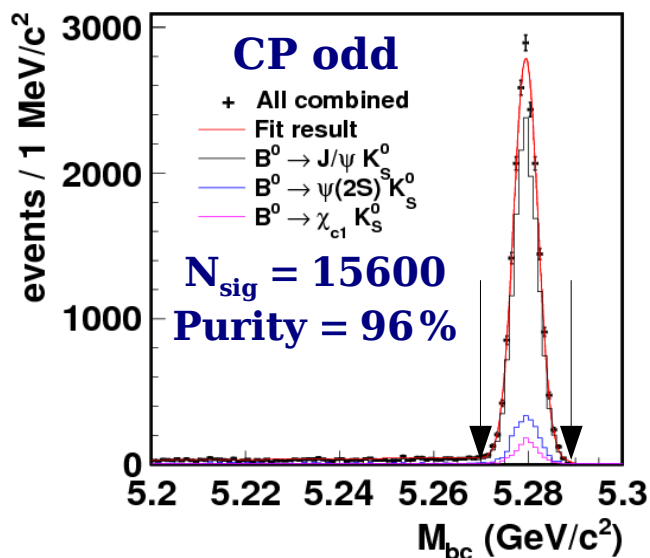
$$\frac{dP_{\text{sig}}}{dt}(\Delta \mathbf{t}, \mathbf{q}) = \frac{e^{-|\Delta \mathbf{t}|/\tau_B}}{4\tau_B} (1 + \mathbf{q}(\mathbf{S} \sin(\Delta m_d \Delta \mathbf{t}) + \mathbf{A} \cos(\Delta m_d \Delta \mathbf{t})))$$

Reconstruct B_{CP}



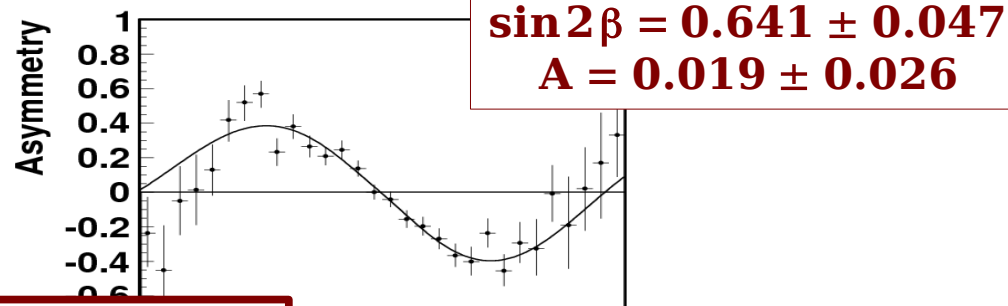
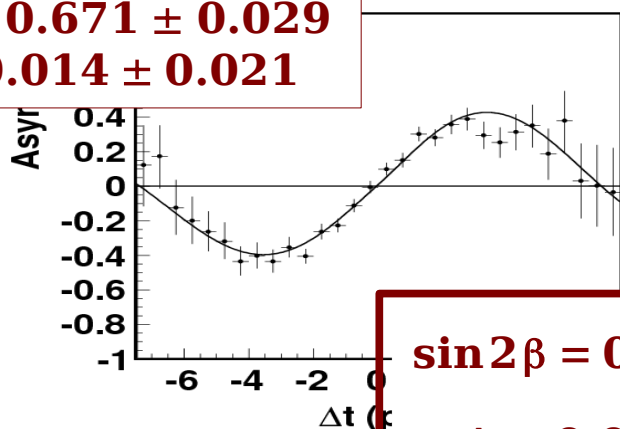
$c\bar{c} K_s$ and $J/\psi K_L$

$772 \times 10^6 B\bar{B}$ pairs



$$\sin 2\beta = 0.671 \pm 0.029$$

$$A = -0.014 \pm 0.021$$

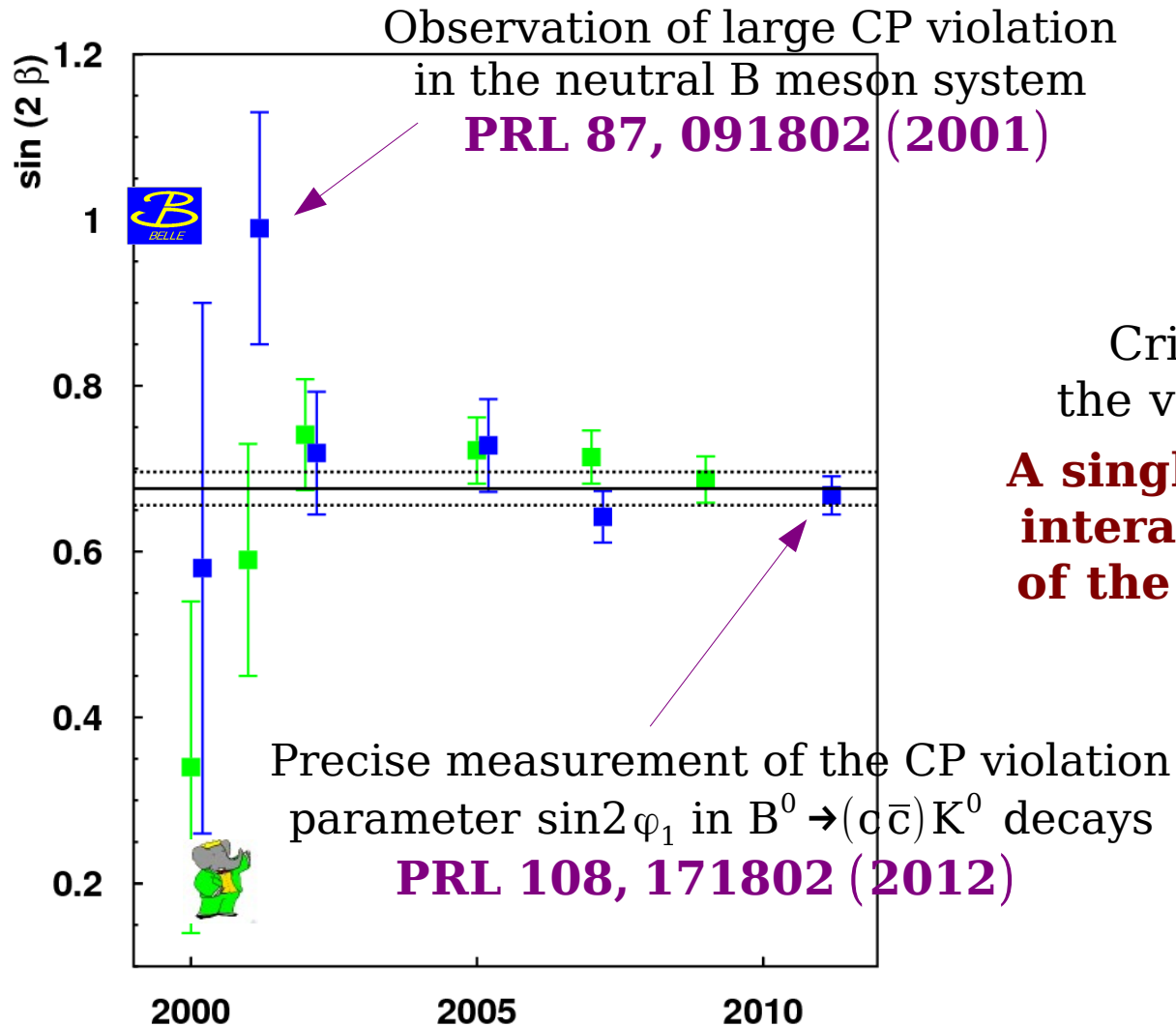


$$\sin 2\beta = 0.668 \pm 0.023 \pm 0.013$$

$$A = 0.007 \pm 0.016 \pm 0.013$$

World's most precise meas^t
anchor point of the SM
still statistically limited !

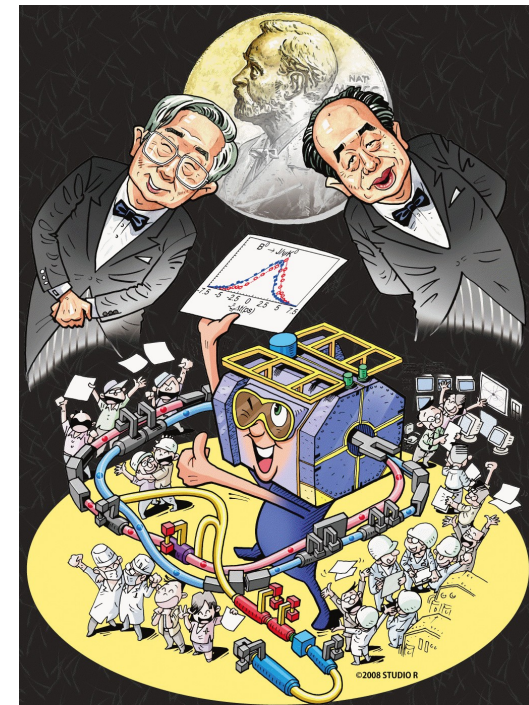
$\sin 2\beta$ in $(c\bar{c})K^0$...



$$\beta = (21.4 \pm 0.8)^\circ$$

Critical role of the B factories in
the verification of the KM hypothesis

**A single irreducible phase in the weak
interaction matrix accounts for most
of the CPV observed in kaons and B's**



Measurement of CPV in $B \rightarrow J/\psi K_S^0$ at LHCb

3 fb⁻¹, arXiv:1503.07089

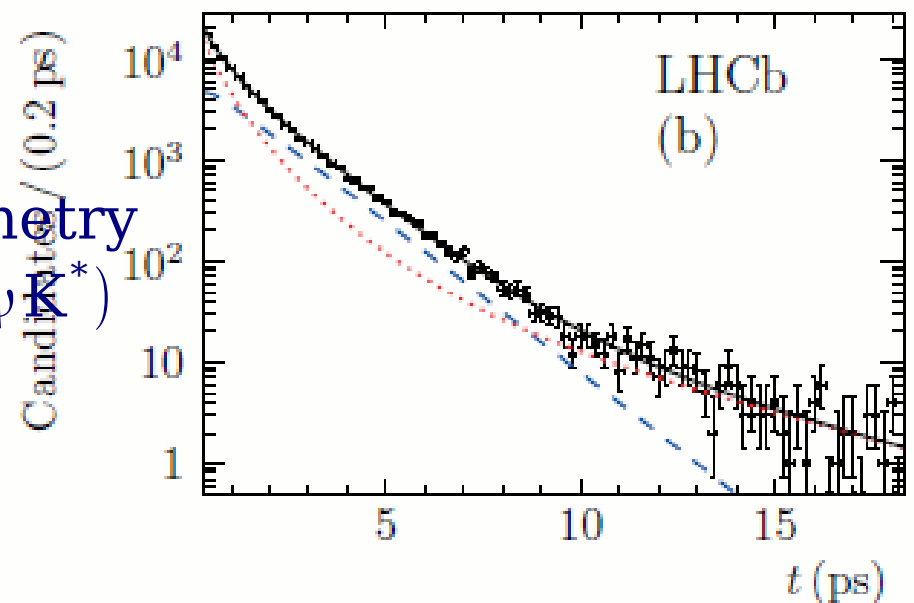
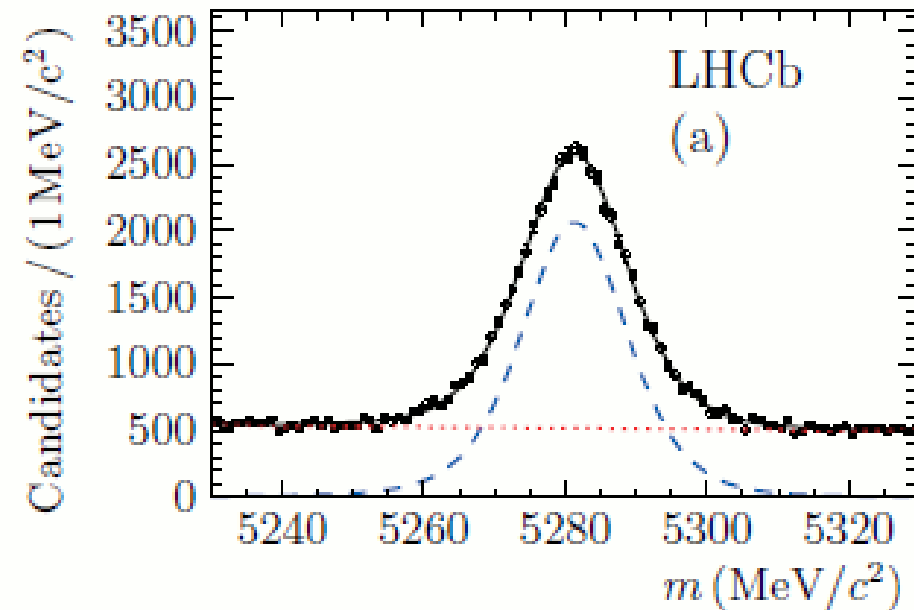
Reconstruct 41560 ± 720 tagged
 $B \rightarrow J/\psi K_S$ events with $J/\psi \rightarrow \mu\mu$ and
 $K_S \rightarrow \pi^+ \pi^-$ in 3 fb⁻¹ (2011-2012 data)

Opposite side flavour-tagging mostly

Magnet polarity reversed periodically
to help control detector asymmetries

Need to correct for production asymmetry
at p-p collider (measured with $B_d \rightarrow J/\psi K^{*}$)

$$A_P = \frac{[\sigma(\bar{B}^0) - \sigma(B^0)]}{[\sigma(\bar{B}^0) + \sigma(B^0)]}$$



Flavour-Tagging at LHCb

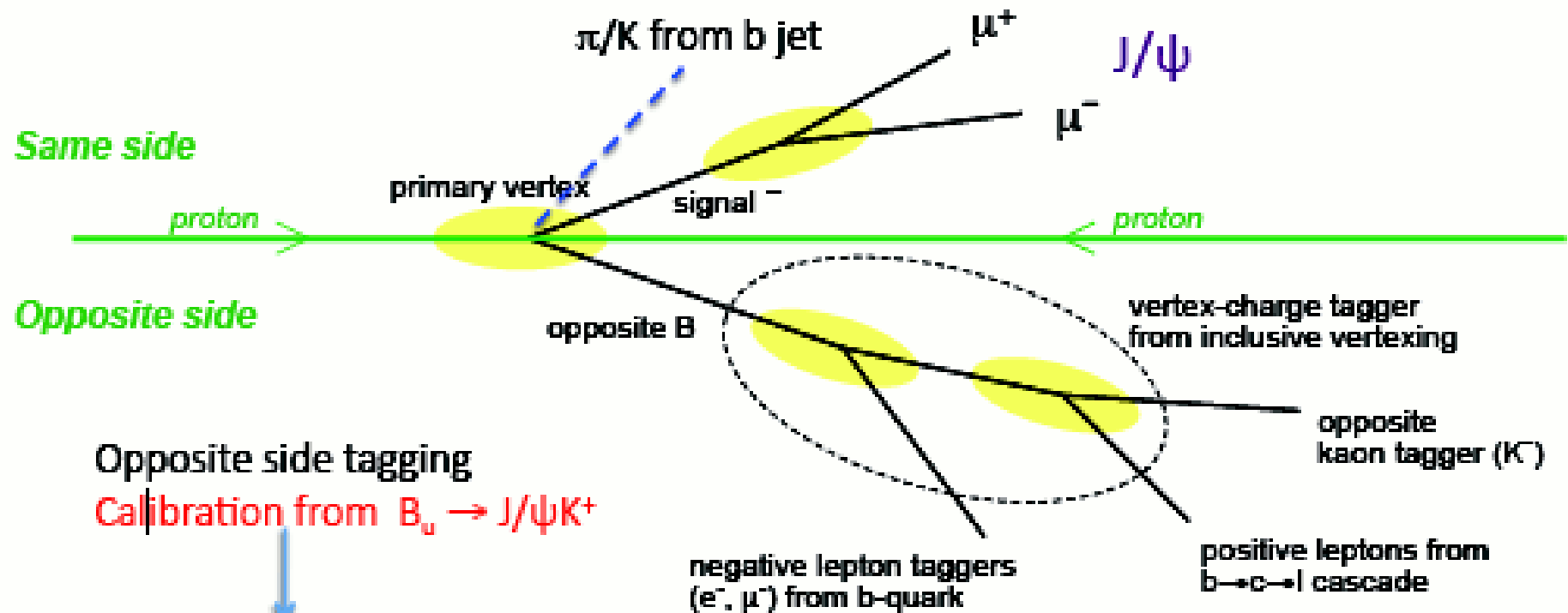
tagging efficiency $\epsilon_{\text{tag}} \sim 50\%$

effective mistag $\omega_{\text{tag}} \sim 39\%$

effective tagging power $\epsilon_{\text{tag}}(1 - 2\omega_{\text{tag}})^2 \sim 2.4\%$

Same side Kaon tagging

Calibration from $B_s \rightarrow D_s \pi$



Opposite side tagging

Calibration from $B_u \rightarrow J/\psi K^+$

tagging efficiency $\epsilon_{\text{tag}} \sim 65\%$

effective mistag $\omega_{\text{tag}} \sim 39\%$









effective tagging power $\epsilon_{\text{tag}}(1 - 2\omega_{\text{tag}})^2 \sim 3.0\%$

Analyses can either use average
or per event tagging information

Flavour-Tagging at LHCb

tagging efficiency $\epsilon_{\text{tag}} \sim 50\%$
effective mistag $\omega_{\text{tag}} \sim 39\%$

Same side Kaon tagging
Calibration from $B_c \rightarrow D_s \pi$

		$\epsilon_{\text{eff}} [\%]$			Reference
Channel		2011	Run I	Imprvt	
	$B_s^0 \rightarrow \phi\phi$	3.29	5.38	+64%	[Phys. Rev. D90 (2014) 052011]
	$B_s^0 \rightarrow D_s^+ D_s^-$		5.33		[Phys. Rev. Lett. 113 (2014) 211801]
	$B_s^0 \rightarrow D_s^+ K^-$	5.07			[JHEP 11 (2014) 060]
	$B_s^0 \rightarrow J/\psi K^+ K^-$	3.13	3.73	+19%	[Phys. Rev. Lett. 114 (2015) 041801]
	$B_s^0 \rightarrow J/\psi \pi^+ \pi^-$	2.43	3.89	+60%	[Phys. Lett. B736 (2014) 186]
	$B^0 \rightarrow J/\psi K_S^0$	2.38	3.03	+27%	[Phys. Rev. Lett. 115 (2015) 031601]
	$B_s^0 \rightarrow J/\psi \phi$	1.45	1.49	+3%	Preliminary
	$B_s^0 \rightarrow J/\psi \phi$	0.97	1.31	+35%	[arXiv:1507.07527]

Impressive improvements in tagging performance in the last 3 years

tagging efficiency $\epsilon_{\text{tag}} \sim 50\%$
effective mistag $\omega_{\text{tag}} \sim 39\%$
effective tagging power $\epsilon_{\text{tag}}(1 - 2\omega_{\text{tag}})^2 \sim 3.0\%$

Analyses can either use average
or per event tagging information

Measurement of CPV in $B \rightarrow J/\psi K_S^0$ at LHCb

$$A(t) = \frac{\Gamma(\bar{B}) - \Gamma(B)}{\Gamma(\bar{B}) + \Gamma(B)}$$

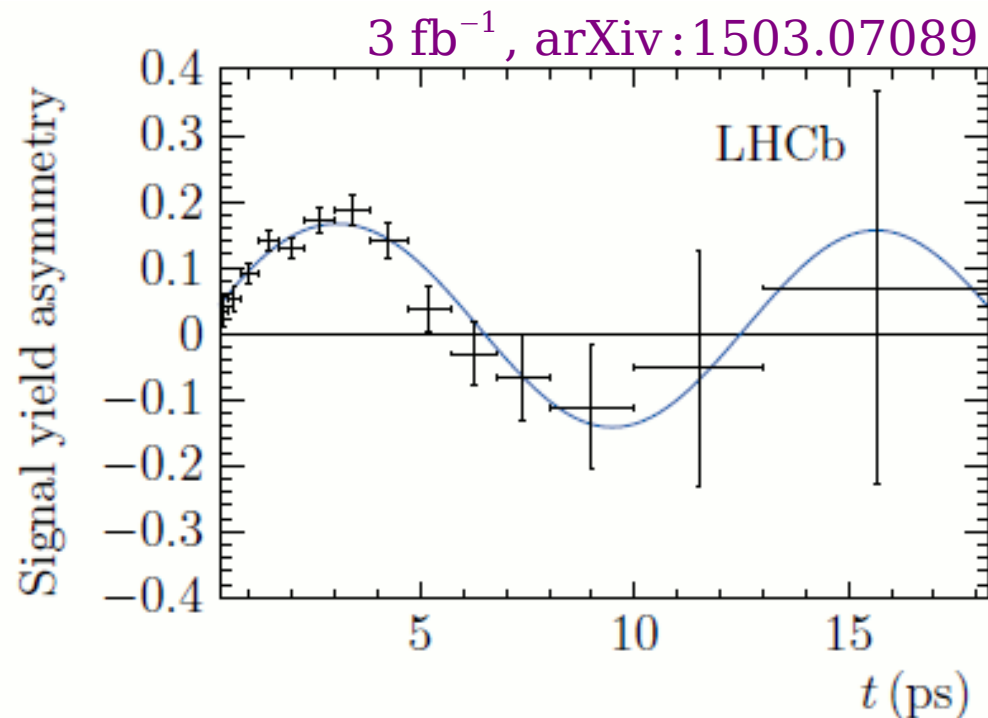
$$S = \sin 2\beta$$

$$C = 0$$

$$A(t) = S \sin(\Delta m_d t) - C \cos(\Delta m_d t)$$

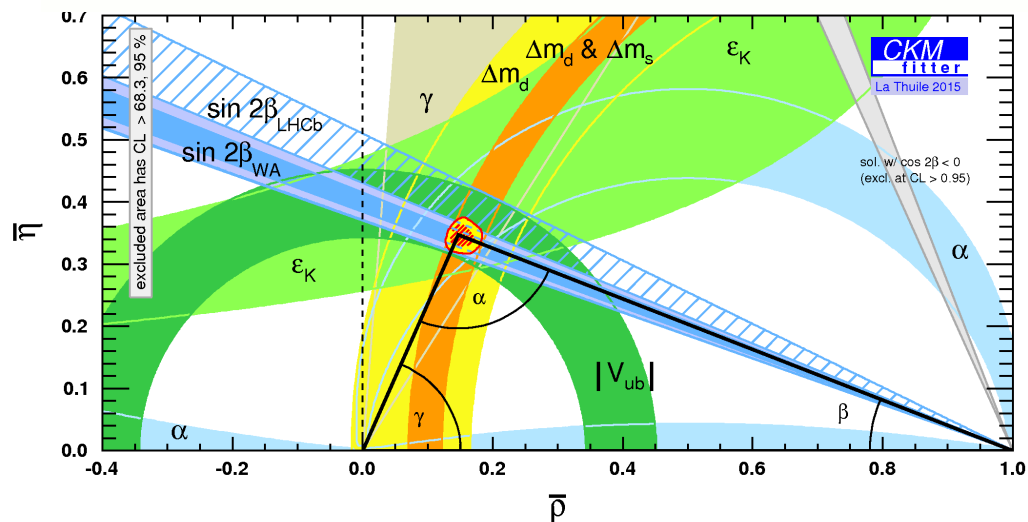
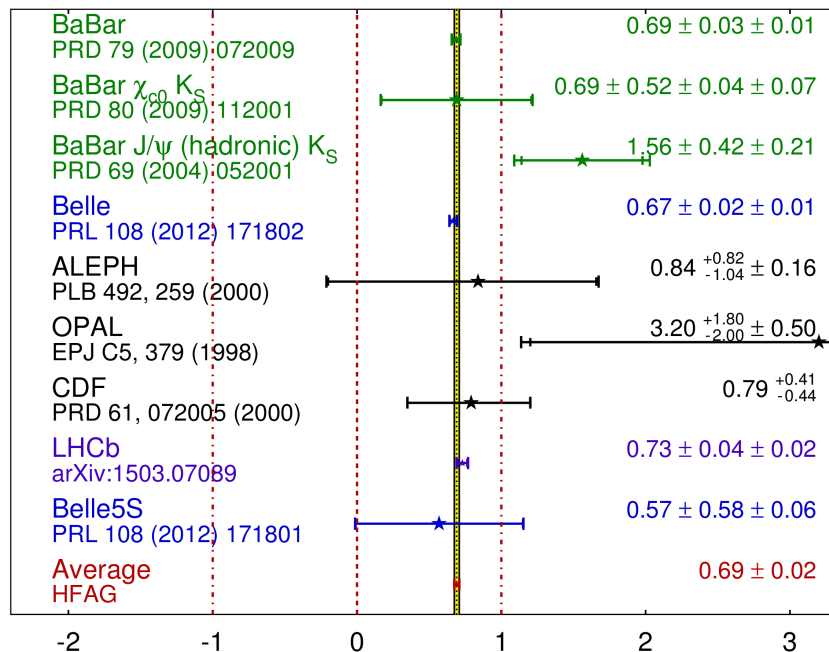
$$S = 0.731 \pm 0.035 \pm 0.020$$

$$C = -0.038 \pm 0.032 \pm 0.005$$



$$\sin(2\beta) \equiv \sin(2\phi_1)$$

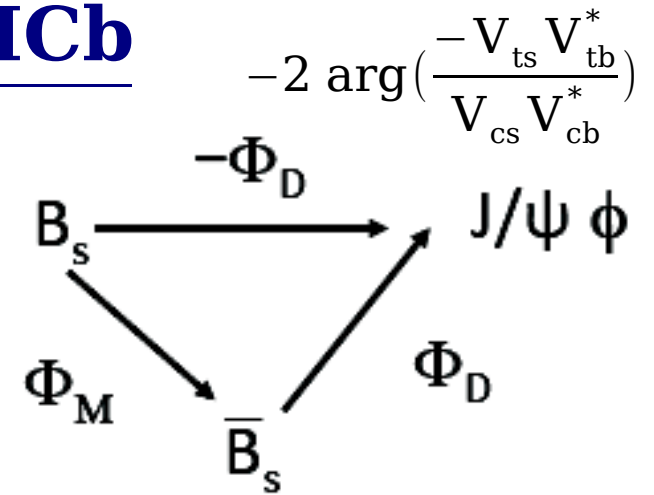
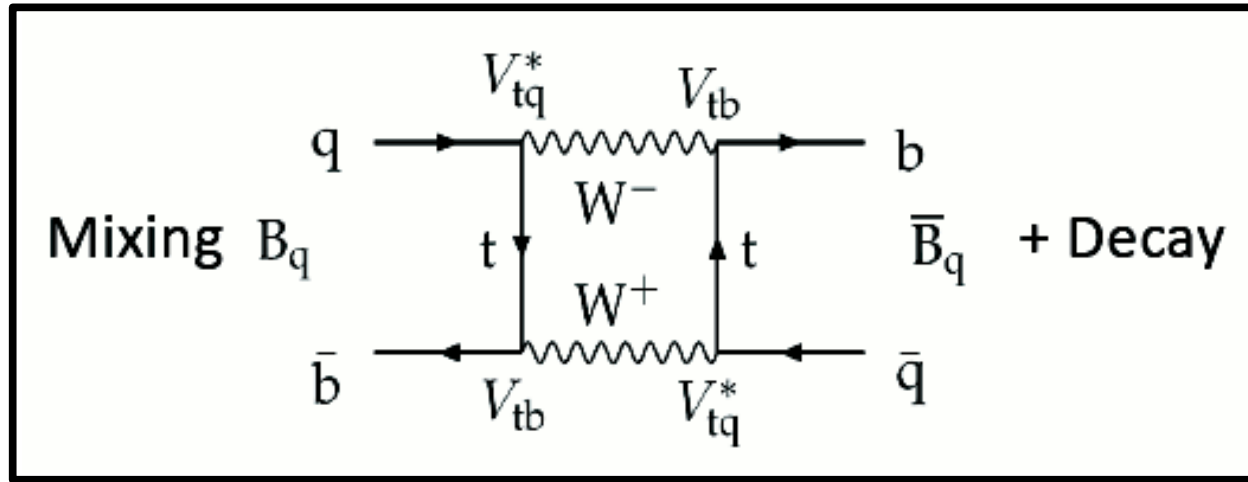
HFAG
Moriond 2015
PRELIMINARY



$$\beta = (21.9 \pm 0.7)^\circ$$

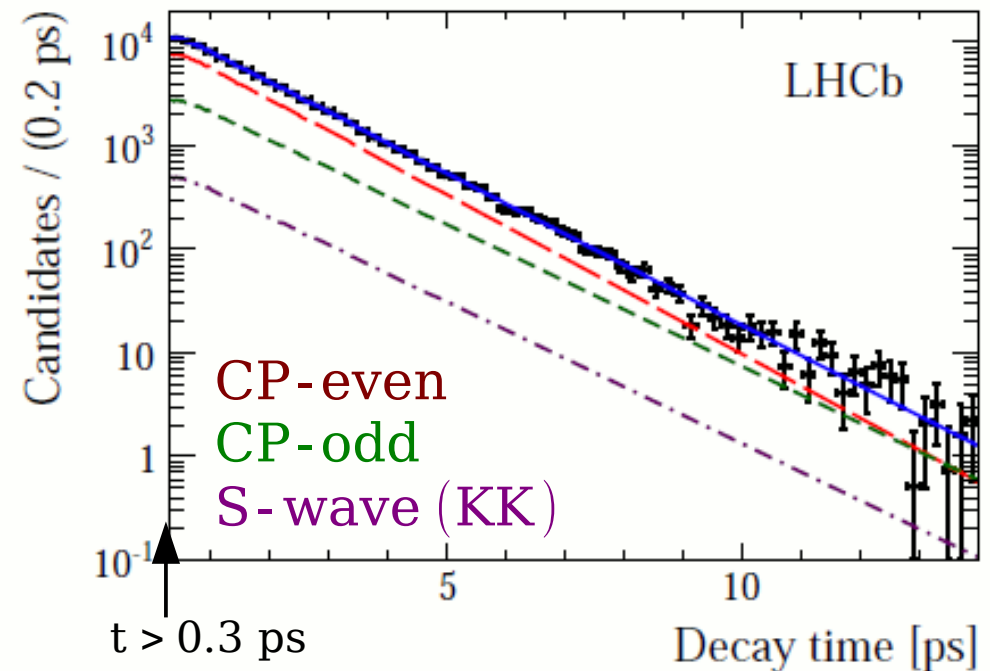
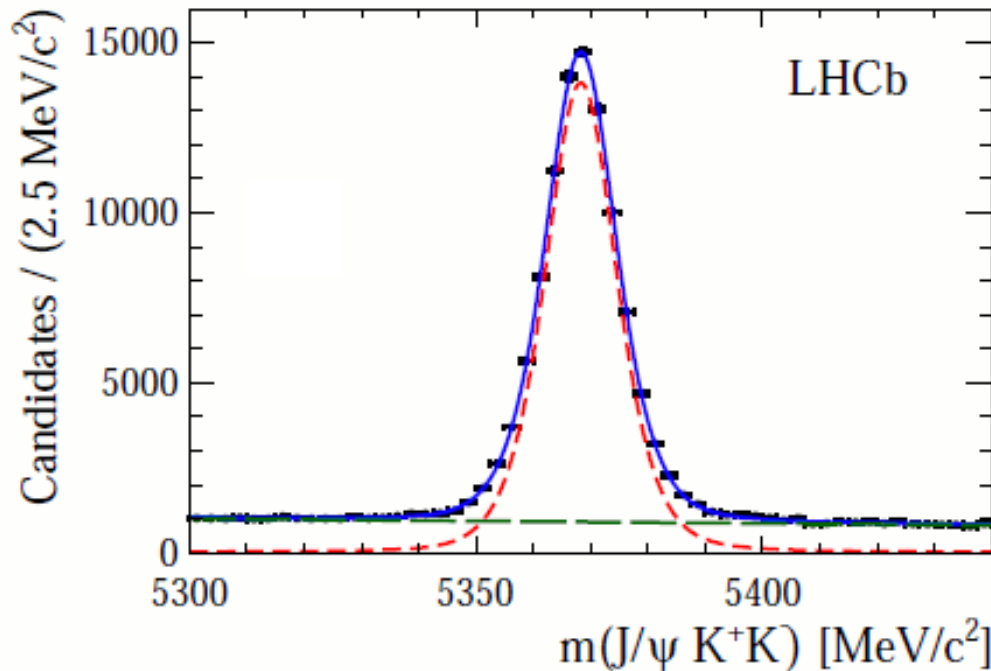
WA 2015

CPV in $B_s \rightarrow J/\psi \phi$ at LHCb



In SM, CPV phase is small $\varphi_s \sim -0.04$ rd

3 fb^{-1} , arXiv:1411.3104



Reconstruct 95690 ± 350 tagged $B_s \rightarrow J/\psi \phi$ evts
 with $J/\psi \rightarrow \mu\mu$ and $\phi \rightarrow K^+ K^-$ in 3 fb^{-1} (2011-2012)

$$\Gamma_s = \Gamma_H + \Gamma_L/2$$

$$= 0.6603 \pm 0.0027 \pm 0.0015 \text{ ps}^{-1}$$

consistent with previous measurements

Results for $B_s \rightarrow J/\psi h^+ h^-$ at LHCb

CP violating phase

[3 fb⁻¹, arXiv:1411.3104]

$$\varphi_s = -0.058 \pm 0.049 \pm 0.006$$

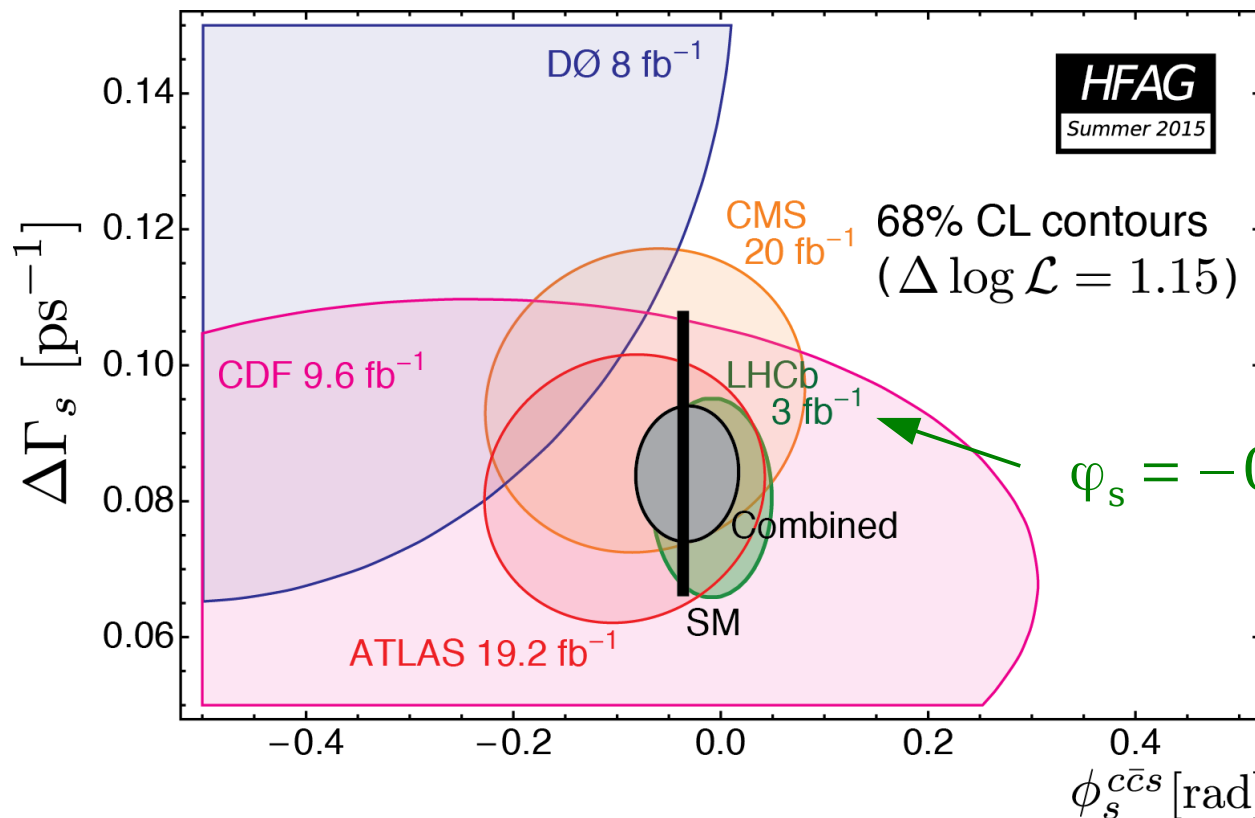
CP violating in mixing or direct decay (no CPV: $|\lambda|=1$)

$$|\lambda| = 0.964 \pm 0.019 \pm 0.007$$

Decay width difference $\Delta\Gamma_s = (\Gamma_L - \Gamma_H) = 0.0805 \pm 0.0091 \pm 0.0032 \text{ ps}^{-1}$

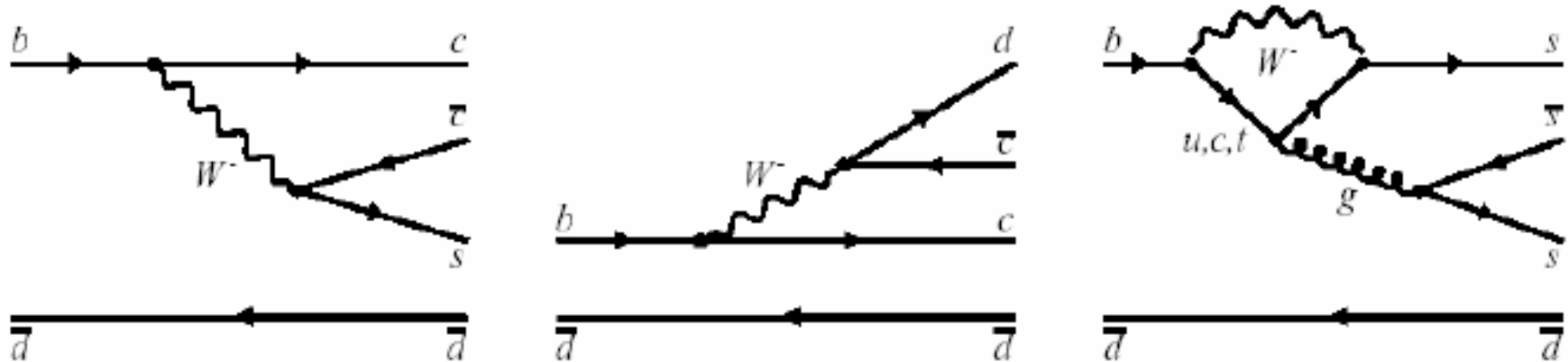
$$\Delta\Gamma_s(\text{SM}) = 0.087 \pm 0.021 \text{ ps}^{-1}$$

$$\varphi_s(\text{SM}) = -0.0363^{+0.0012}_{-0.0014} \text{ rad}$$



$$\varphi_s = -0.010 \pm 0.039 \text{ rad}$$

[combined with $J/\psi\pi\pi$]



$J/\psi K_S^0, \psi(2S) K_S^0, \chi_{c1} K_S^0,$
 $\eta_c K_S^0, J/\psi K_L^0,$
 $J/\psi K^{*0} (K^{*0} \rightarrow K_S^0 \pi^0)$

$D^{*+} D^-, D^+ D^-$
 $J/\psi \pi^0, D^{*+} D^{*-}$

$\phi K^0, K^+ K^- K_S^0,$
 $K_S^0 K_S^0 K_S^0, \eta' K^0, K_S^0 \pi^0,$
 $\omega K_S^0, f_0(980) K_S^0$

← increasing tree diagram amplitude

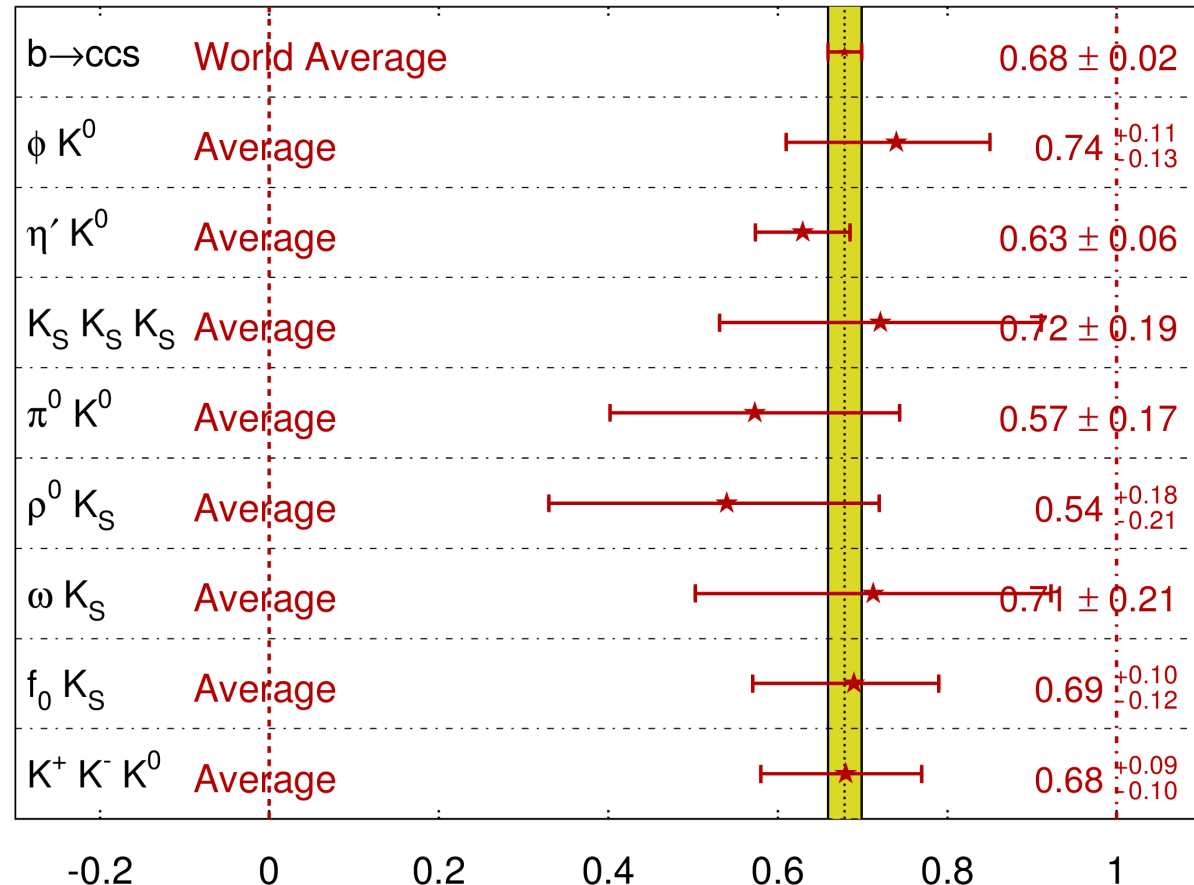
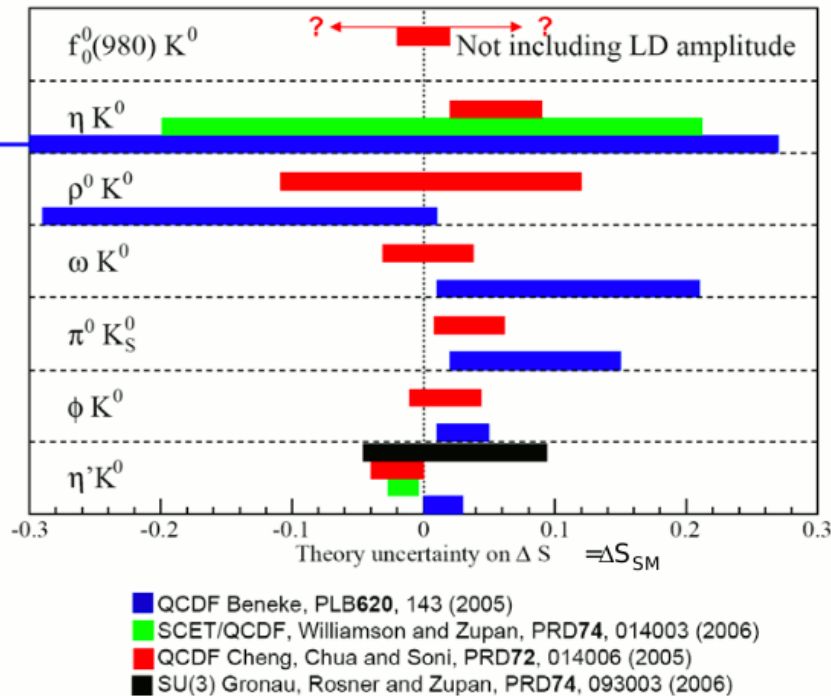
← increasing sensitivity to new physics →

$\sin 2\beta$ with $b \rightarrow s$ penguins

dominated by
B-factories

$$\sin(2\beta^{\text{eff}}) \equiv \sin(2\phi_1^{\text{eff}})$$

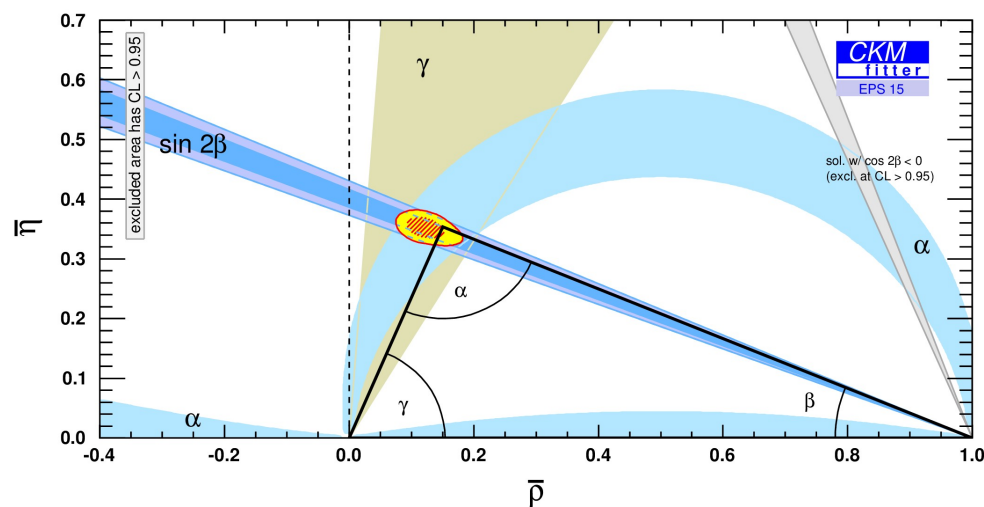
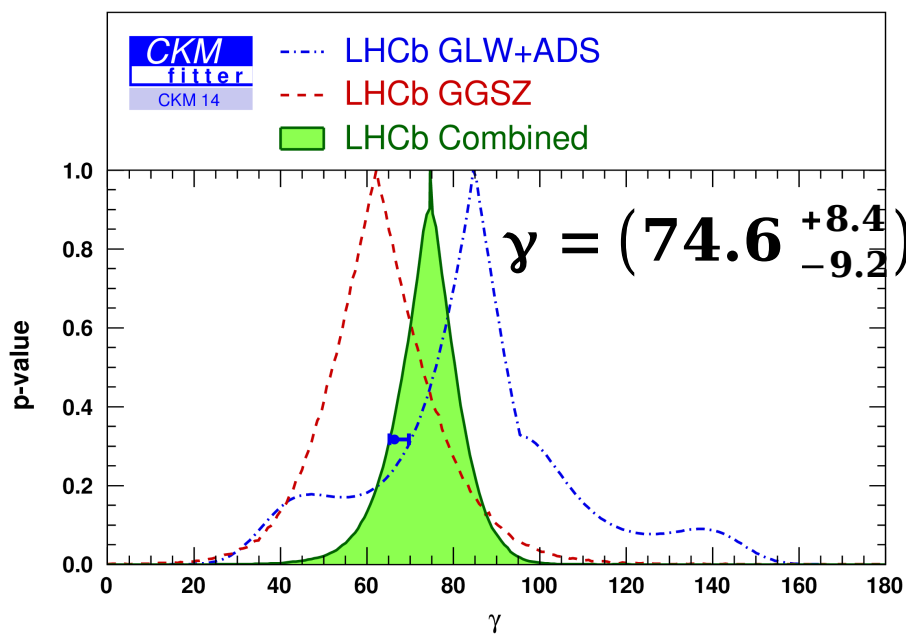
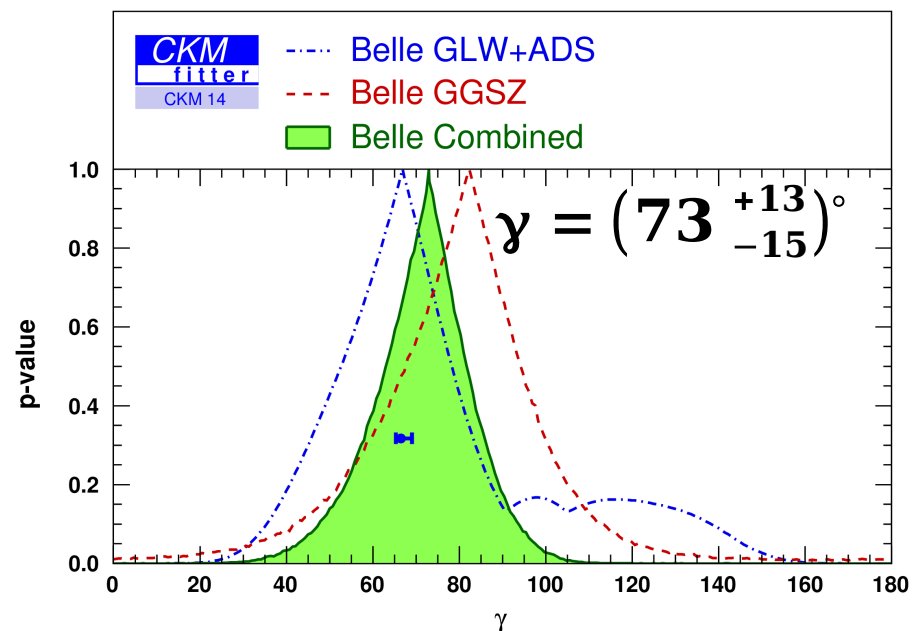
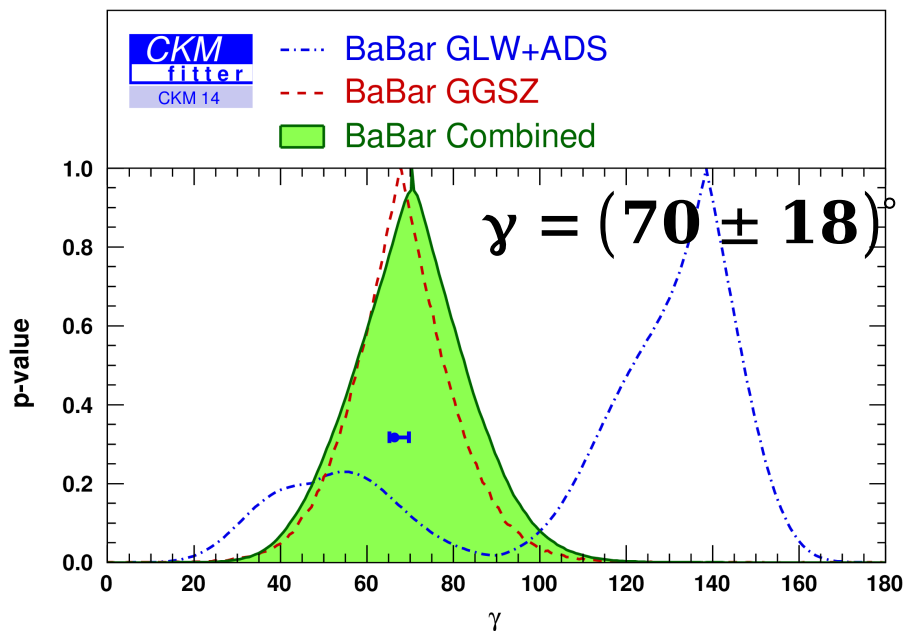
HFAG
Moriond 2014
PRELIMINARY



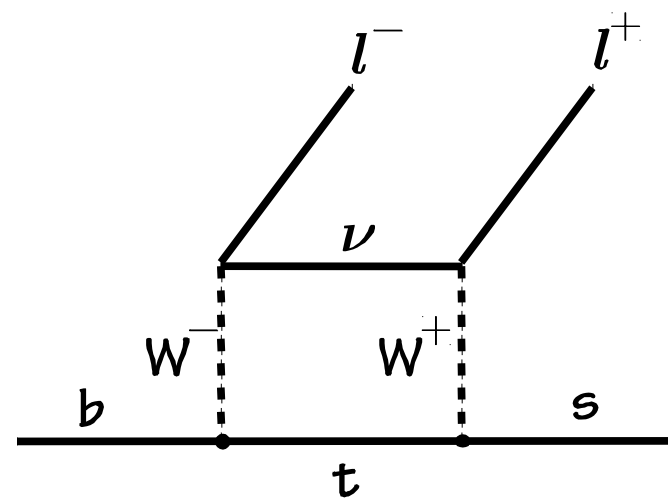
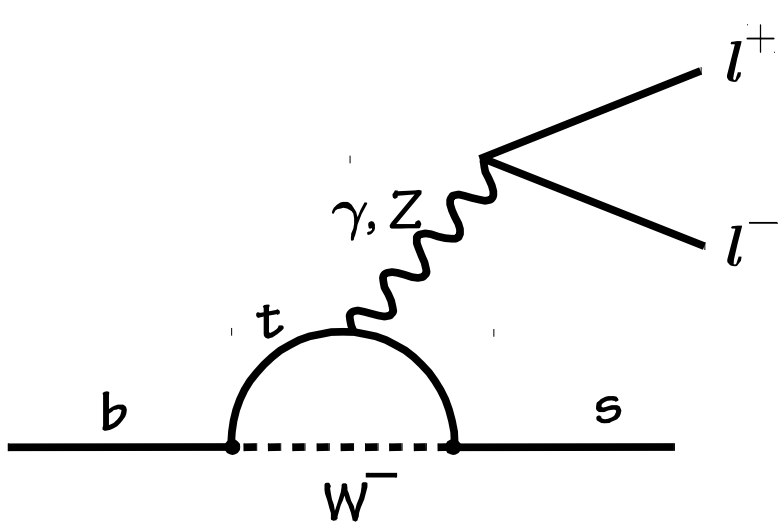
More statistics crucial
for mode-by-mode studies

γ angle in the global fit

measurements from $B \rightarrow DK$

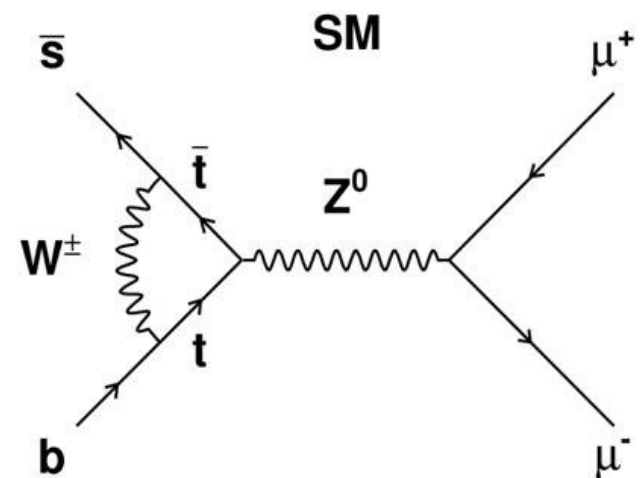
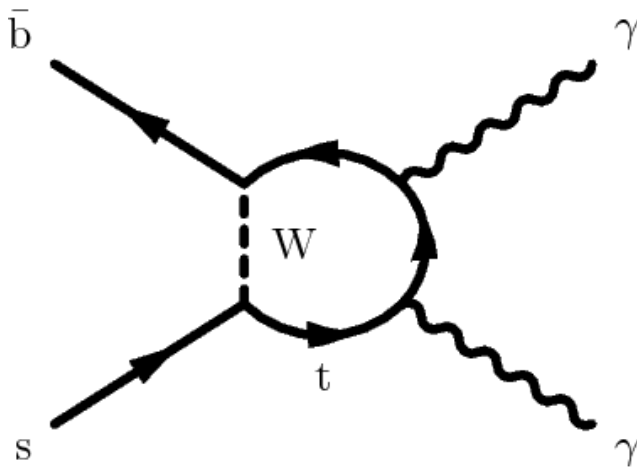
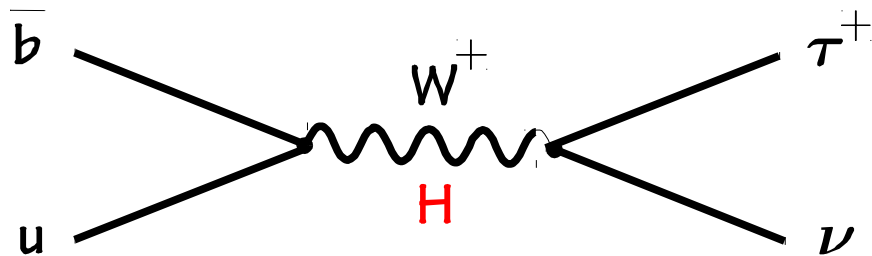


long way to go ... ($\rightarrow \sigma_\gamma = 1^\circ$ or less)



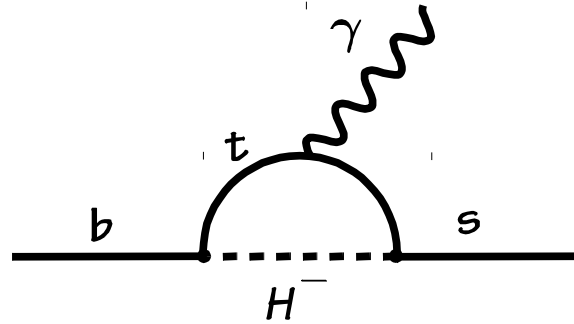
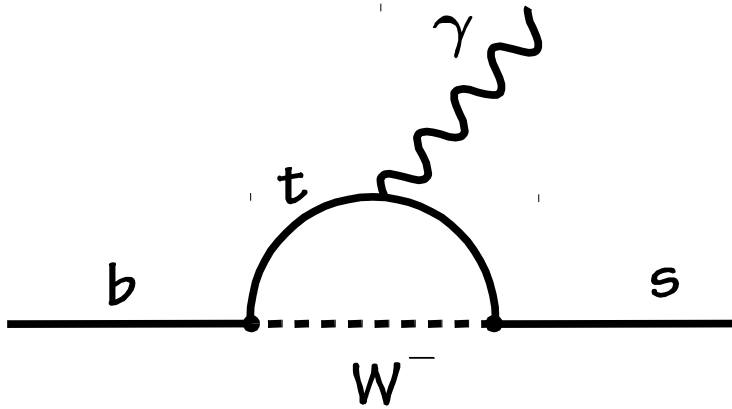
Rare B decays

- FCNC are strongly suppressed in the SM: only loops + GIM mechanism
- Any new particle generating new diagrams can change the amplitudes

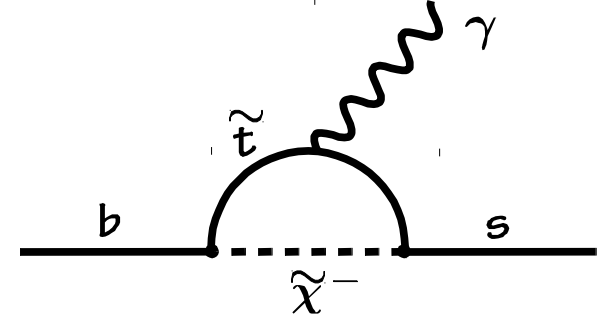


[see talk from FJPPL: FLAV_02 (K.Hara)]

$B \rightarrow X_s \gamma$



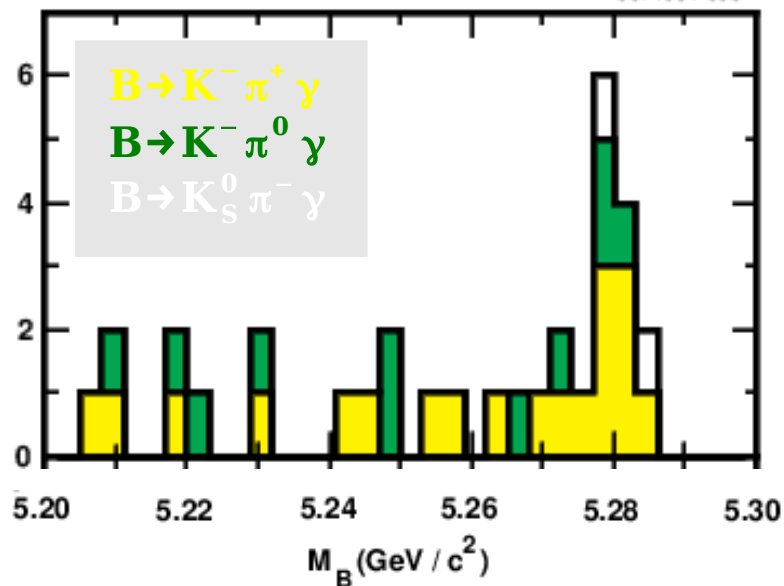
Sensitive to NP



rare ? not that rare...

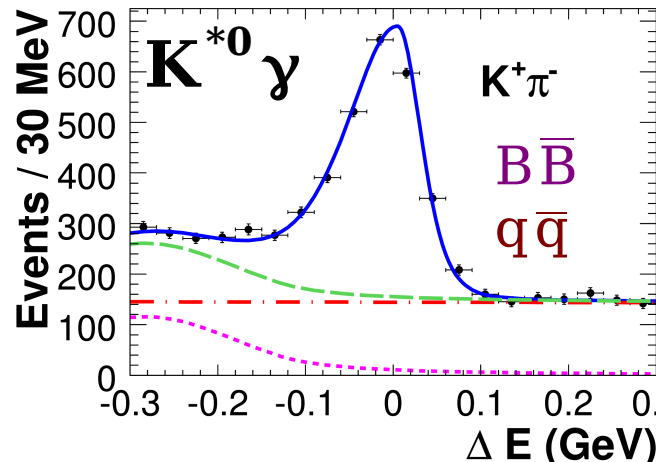
CLEO observation of $B \rightarrow K^* \gamma$ [1993]

3071094-008

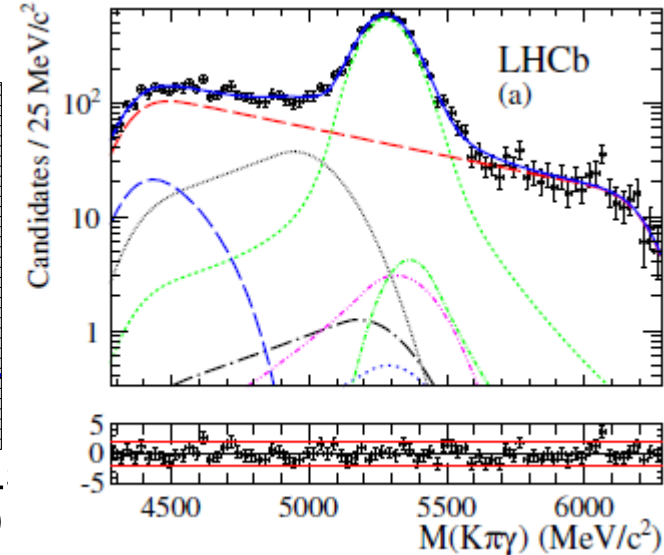


[383 MBB]

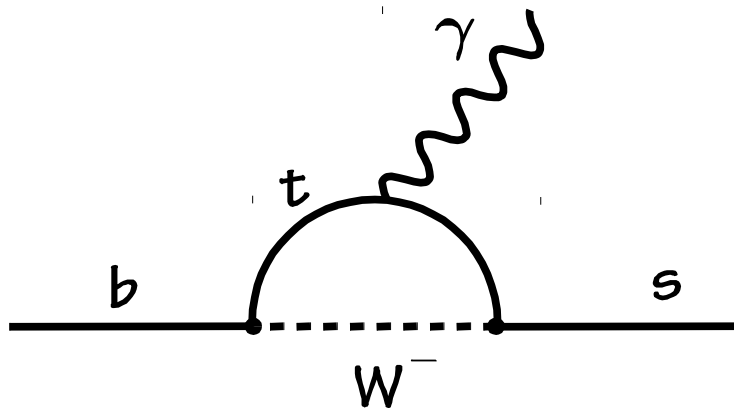
[arXiv:0906.2177]



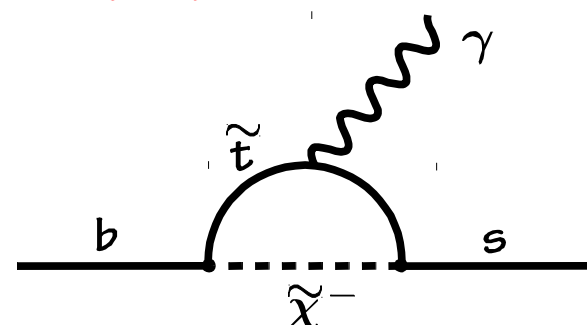
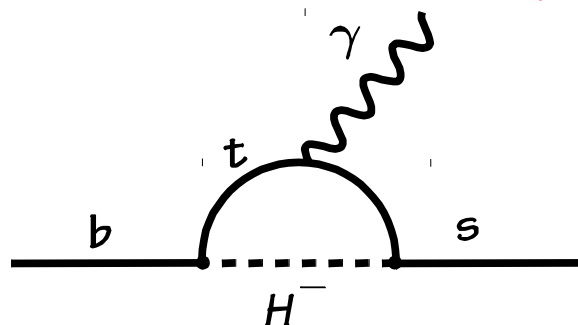
[arXiv:1209.0313], 1 fb^{-1}



$B \rightarrow X_s \gamma$



Sensitive to NP



NNLO SM calculation:

$$B_{SM}(B \rightarrow X_s \gamma) = (3.36 \pm 0.23) \times 10^{-4}$$

(for $E_\gamma > 1.6$ GeV)

M.Misiak et al.

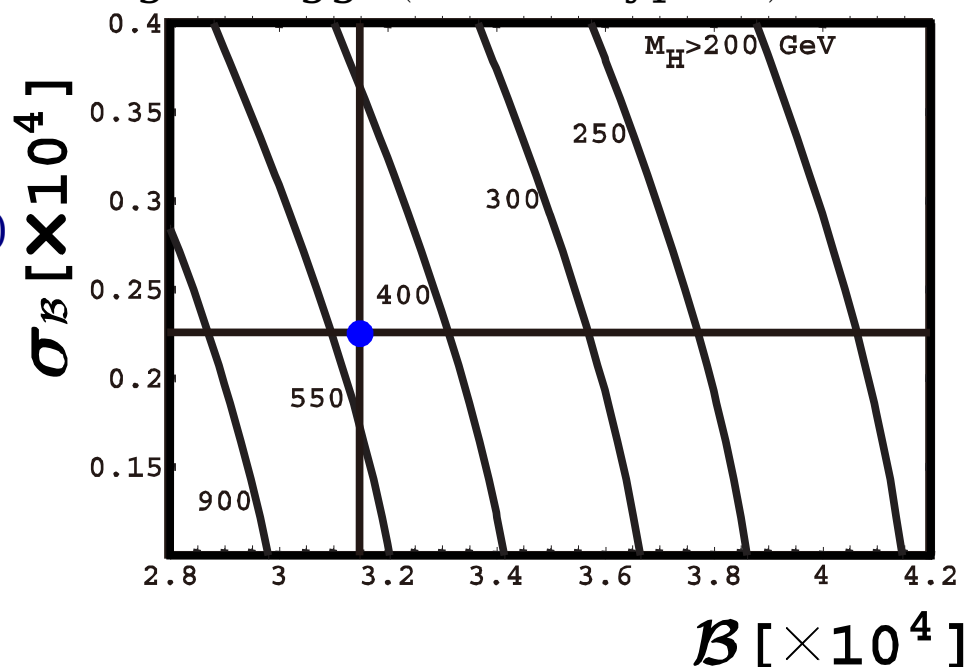
[arXiv:1503.01789]

(central value increased by
6.4% compared to 2007 value)

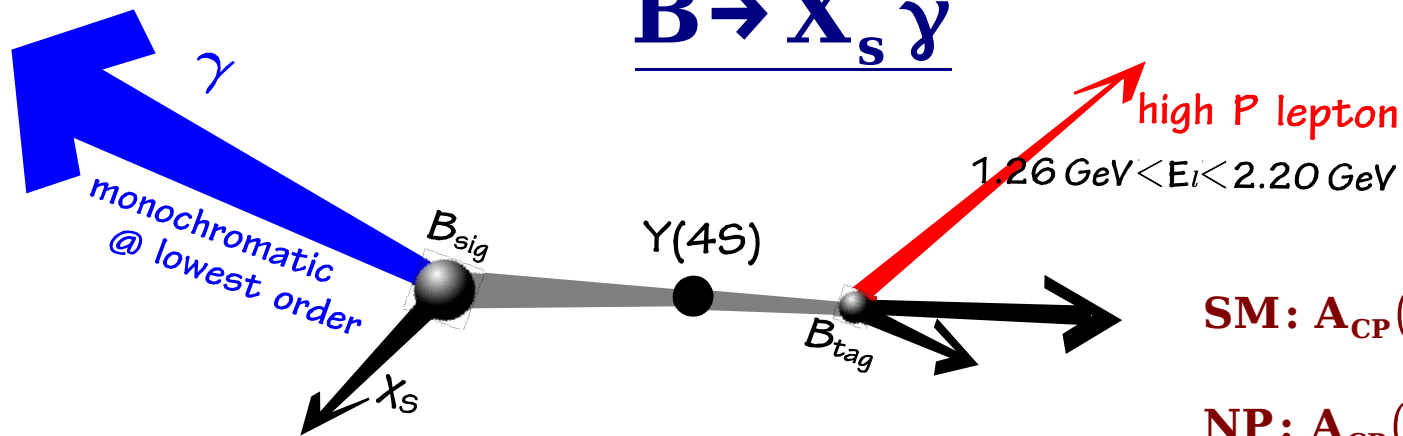
PRL 98, 022002 (2007)

The lower γ energy threshold, the smaller
the model uncertainties in SM, but the
larger background in measurement

Charged Higgs (2HDM Type II) bound

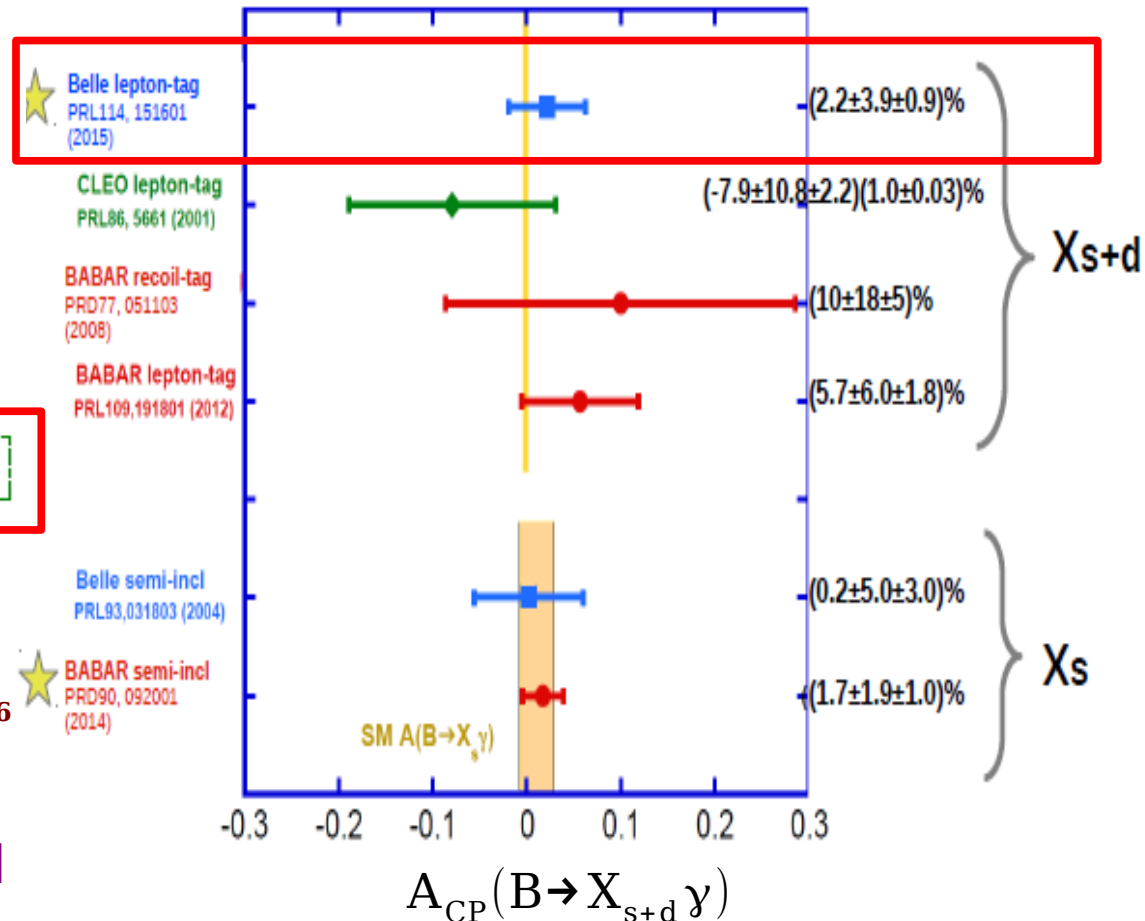
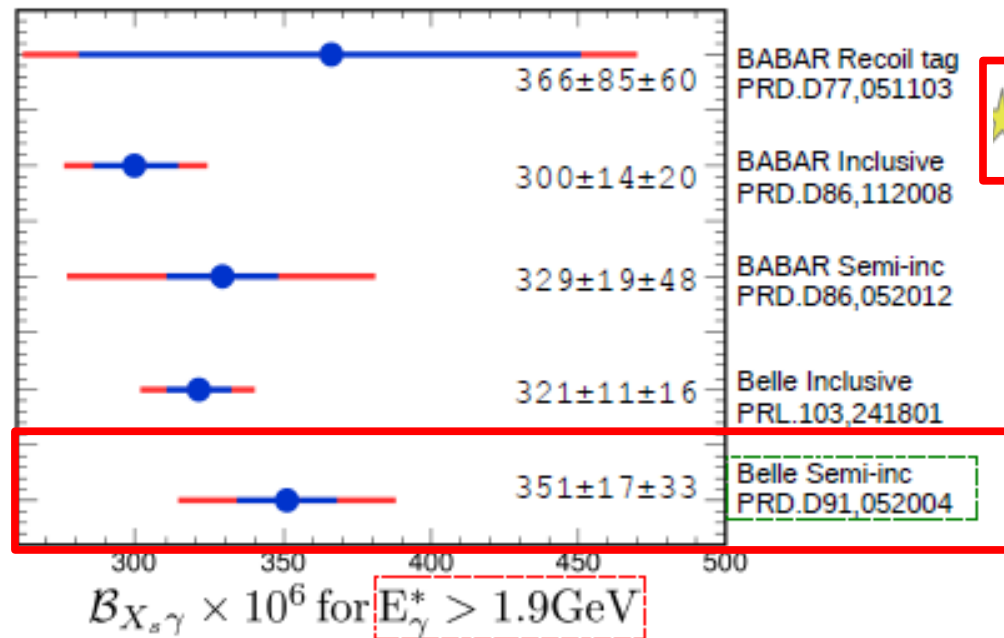


$$\underline{B \rightarrow X_s \gamma}$$



SM: $A_{CP}(B \rightarrow X_{s+d} \gamma) = 0$ to order 10^{-6}
[Hurth and Mannel, 2001]

NP: $A_{CP}(B \rightarrow X_{s+d} \gamma)$ as large as 10%



at $E_\gamma > 1.6 \text{ GeV}$:

$$B(B \rightarrow X_s \gamma) = (341 \pm 15 \pm 4 (\text{extrap})) \times 10^{-6}$$

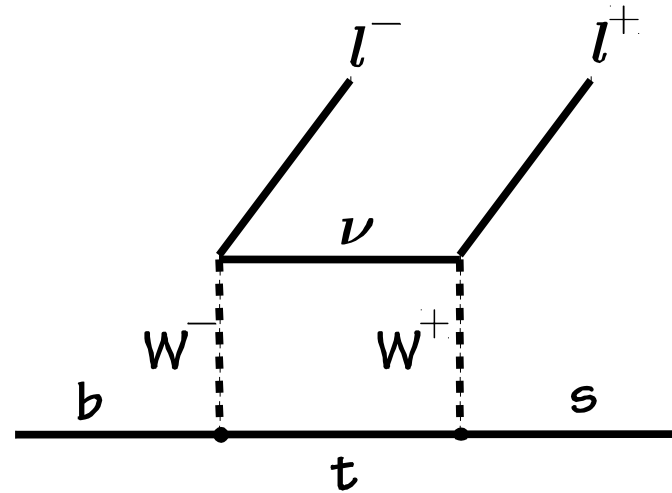
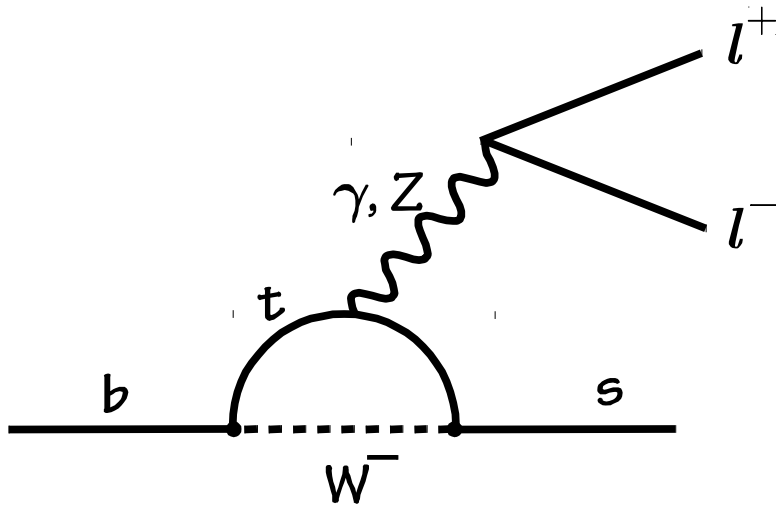
$$B_{SM}(B \rightarrow X_s \gamma) = (336 \pm 23) \times 10^{-6}$$

[Misiak et al, arXiv:1503.01789]

For charged Higgs in 2HDM Type II
 $M(H^\pm) > 540 \text{ GeV}$ at 95 % CL

\Rightarrow limited by statistics: Belle II...

$b \rightarrow s l^+ l^-$



\Rightarrow 2 orders of magnitude smaller than $b \rightarrow s \gamma$ but rich NP search potential

Amplitudes from

- electromagnetic penguin: C_7
- vector electroweak: C_9
- axial-vector electroweak: C_{10}

may interfere w/ contributions from NP

Many observables:

- Branching fractions
- Isospin asymmetry (A_I)
- Lepton forward-backward asymmetry (A_{FB})

\Rightarrow Exclusive ($B \rightarrow K^{(*)} l^+ l^-$), Inclusive ($B \rightarrow X_s l^+ l^-$)

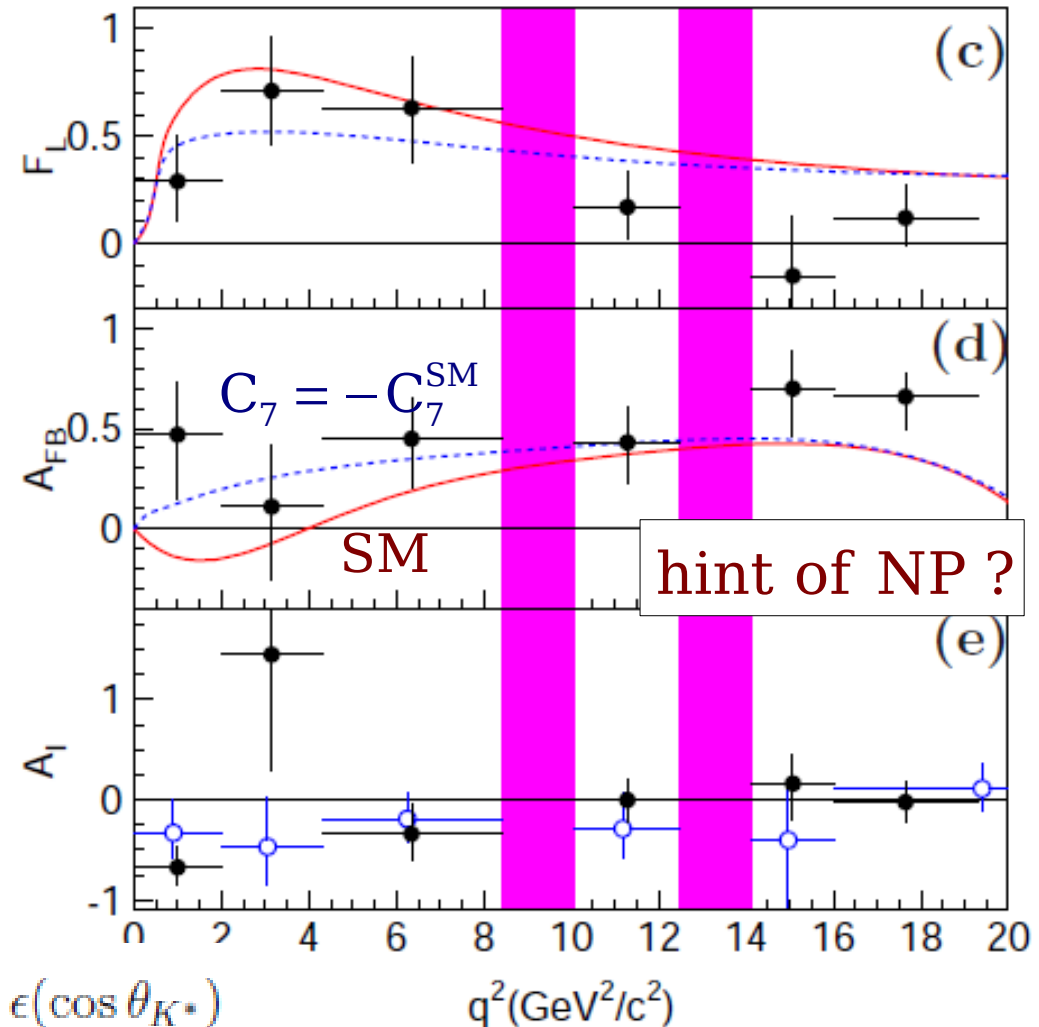
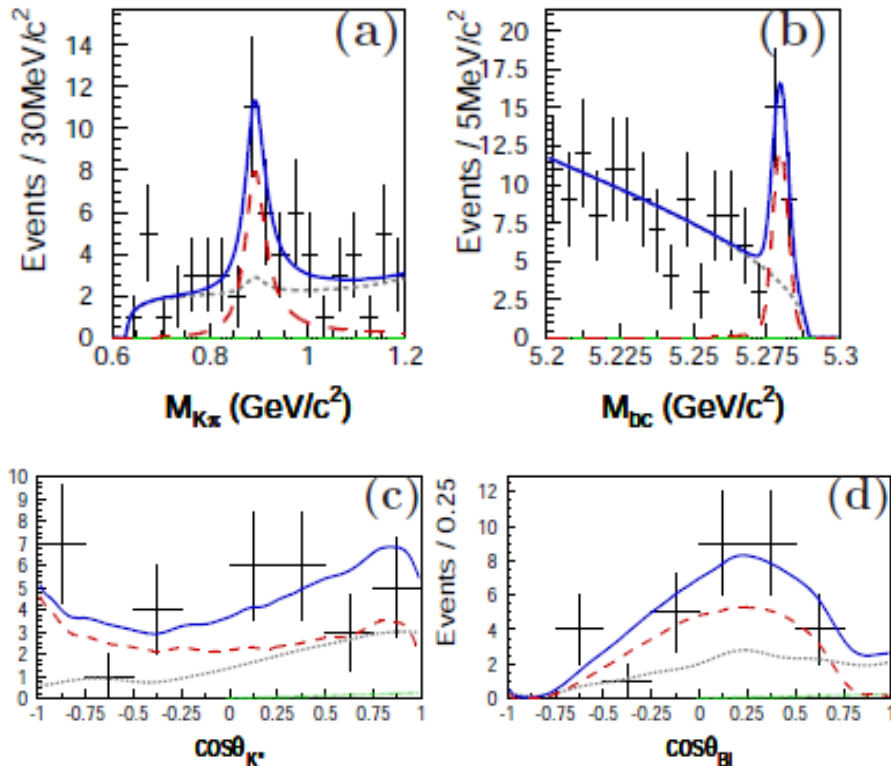
$B \rightarrow K^* l^+ l^-$ decays



- Channels: $K^* \rightarrow K^+ \pi^-$, $K_S^0 \pi^+$, $K^+ \pi^0$, $l = e$ or μ

[arXiv:0904.0770]

illustration: $q^2 \in [0.0, 2.0] \text{ GeV}^2$



$$\left[\frac{3}{2} F_L \cos^2 \theta_{K^*} + \frac{3}{4} (1 - F_L) (1 - \cos^2 \theta_{K^*}) \right] \times \epsilon(\cos \theta_{K^*})$$

$$\left[\frac{3}{4} F_L (1 - \cos^2 \theta_{Bl}) + \frac{3}{8} (1 - F_L) (1 + \cos^2 \theta_{Bl}) + A_{FB} \cos \theta_{Bl} \right] \times \epsilon(\cos \theta_{Bl}) ,$$

$$R_{K^*} = 0.83 \pm 0.17 \pm 0.08$$

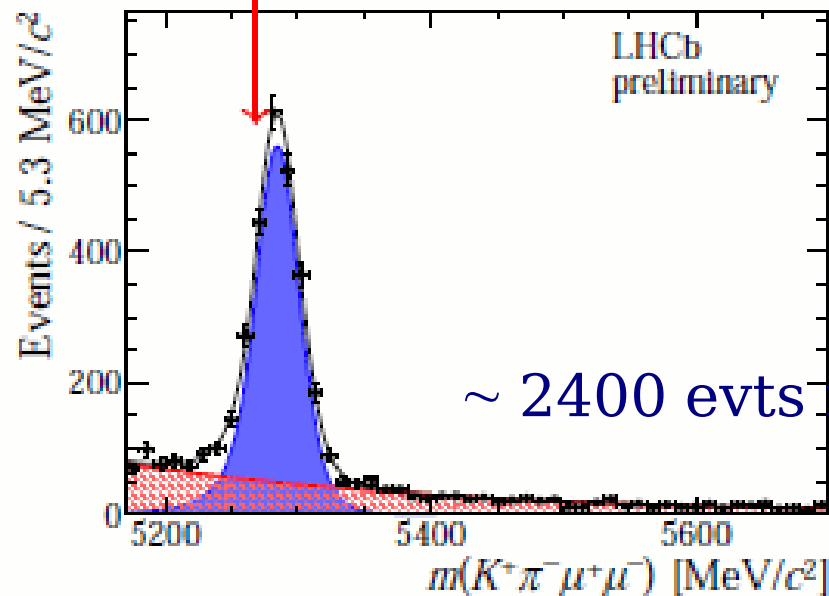
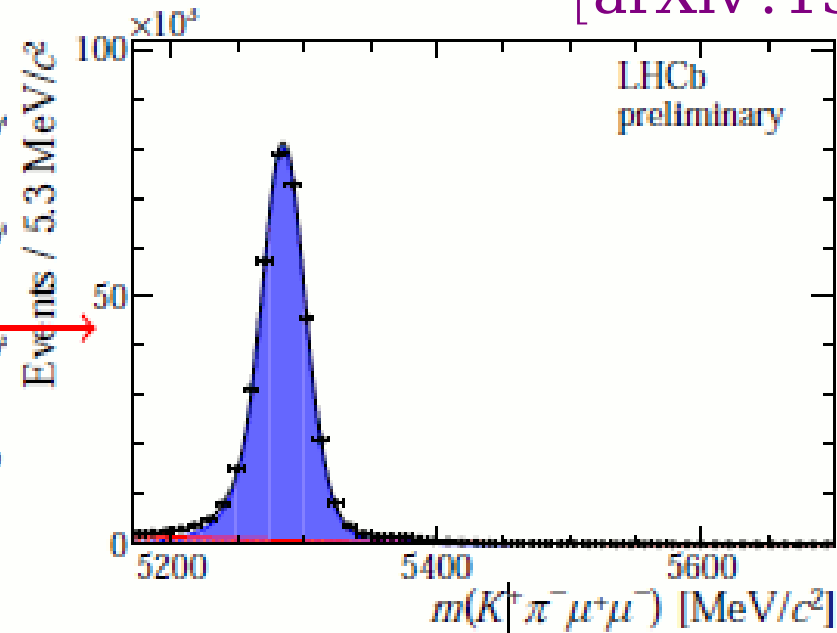
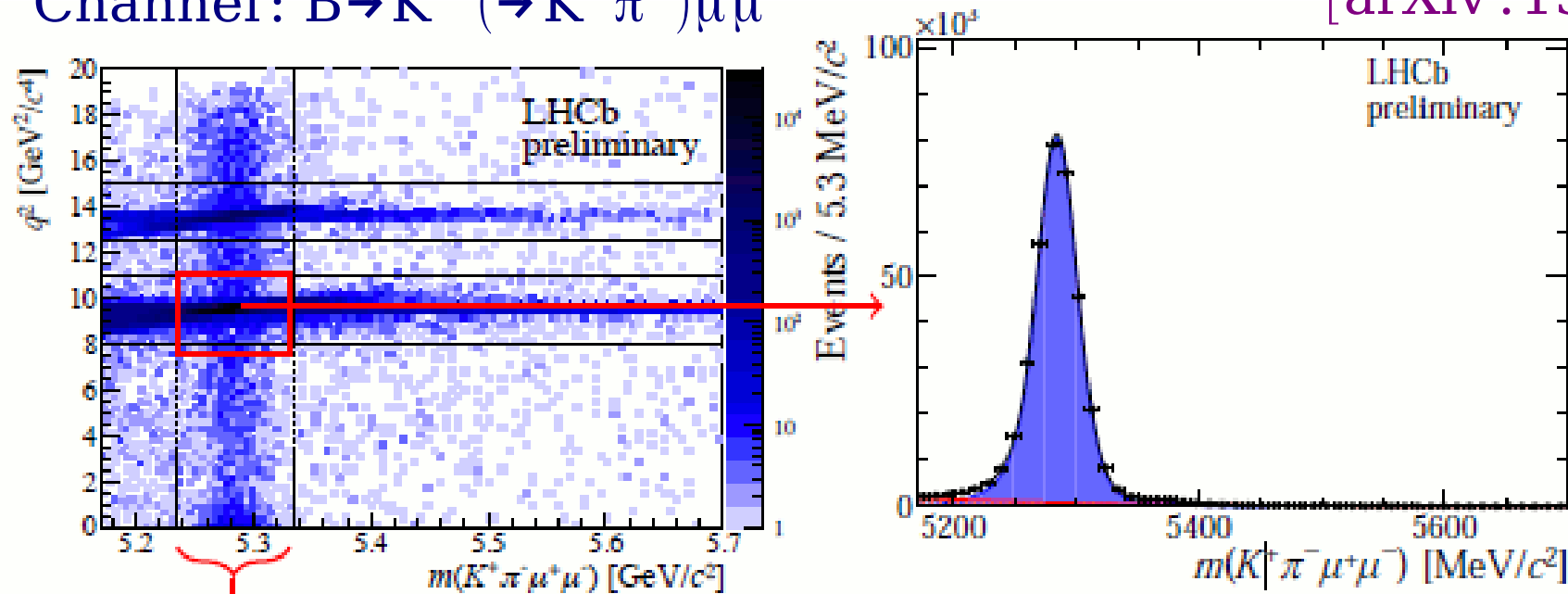
$$R_K = 1.03 \pm 0.19 \pm 0.06$$

$$R_K^{\text{SM}} = 1, R_{K^*}^{\text{SM}} = 0.75 \text{ (photon pole !)}$$

Angular analysis of $B_d^0 \rightarrow K^{*0} \mu^+ \mu^-$ decays

- Channel: $B \rightarrow K^{*0} (\rightarrow K^+ \pi^-) \mu \mu$

[arXiv:1512.04442]



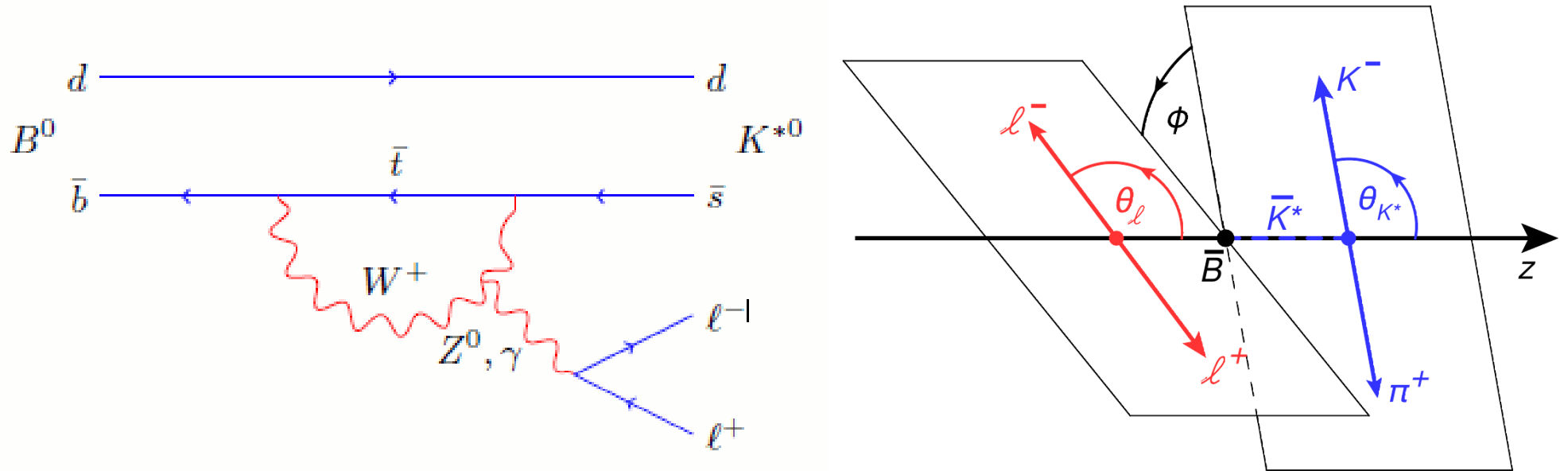
Selection:

BDT to reject combinatorial background
Veto of resonant modes (control modes)

~ 2400 evts in the full q^2 range

Angular analysis of $B_d^0 \rightarrow K^{*0} l^+ l^-$ decays

- Final state described by $q^2 = m_{ll}^2$ and three angles $\Omega = (\theta_l, \theta_K, \phi)$

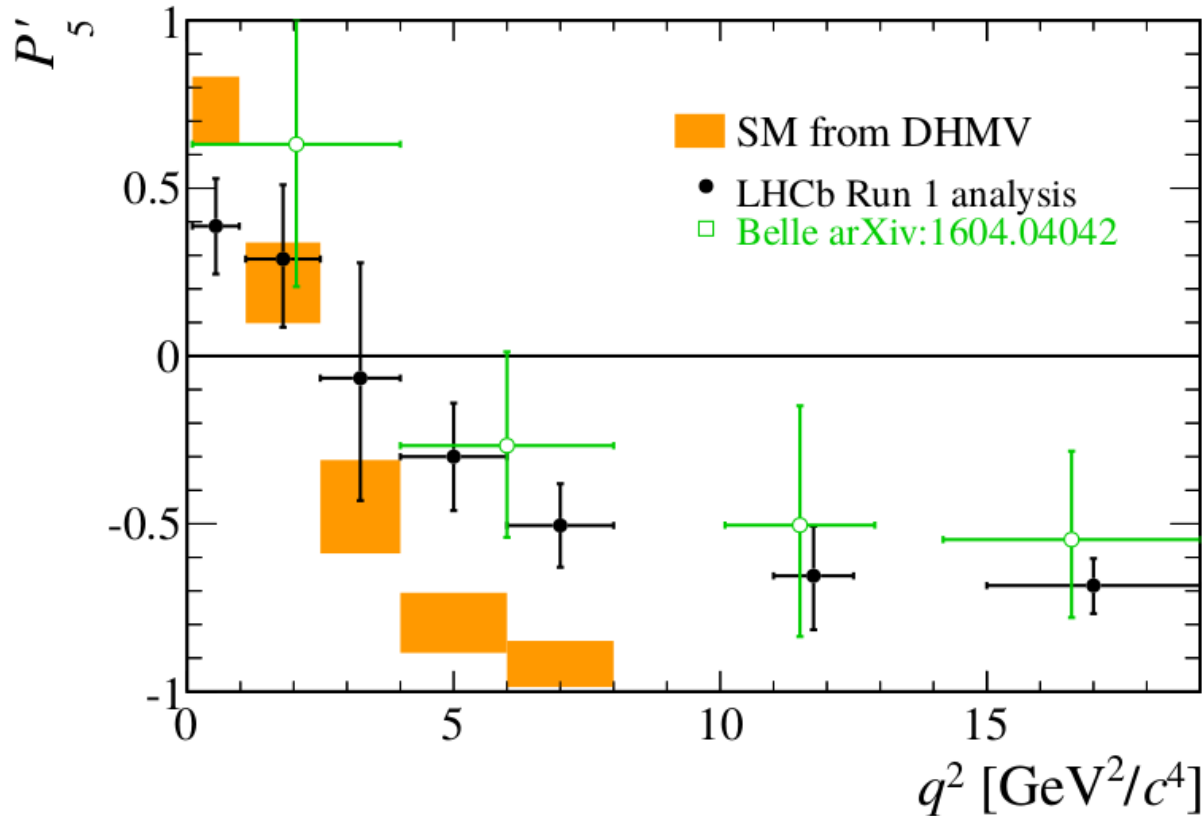


$$\frac{1}{d(\Gamma + \bar{\Gamma})/dq^2} \frac{d^3(\Gamma + \bar{\Gamma})}{d\vec{\Omega}} = \frac{9}{32\pi} \left[\frac{3}{4}(1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K + \frac{1}{4}(1 - F_L) \sin^2 \theta_K \cos 2\theta_\ell \right. \\ \left. - F_L \cos^2 \theta_K \cos 2\theta_\ell + S_3 \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\phi \right. \\ \left. + S_4 \sin 2\theta_K \sin 2\theta_\ell \cos \phi + S_5 \sin 2\theta_K \sin \theta_\ell \cos \phi \right. \\ \left. + \frac{4}{3} A_{FB} \sin^2 \theta_K \cos \theta_\ell + S_7 \sin 2\theta_K \sin \theta_\ell \sin \phi \right. \\ \left. + S_8 \sin 2\theta_K \sin 2\theta_\ell \sin \phi + S_9 \sin^2 \theta_K \sin^2 \theta_\ell \sin 2\phi \right]$$

- F_L, A_{FB}, S_i sensitive to $C_7^{(i)}, C_9^{(i)}, C_{10}^{(i)}$

Angular analysis of $B_d^0 \rightarrow K^* \mu^+ \mu^-$ decays

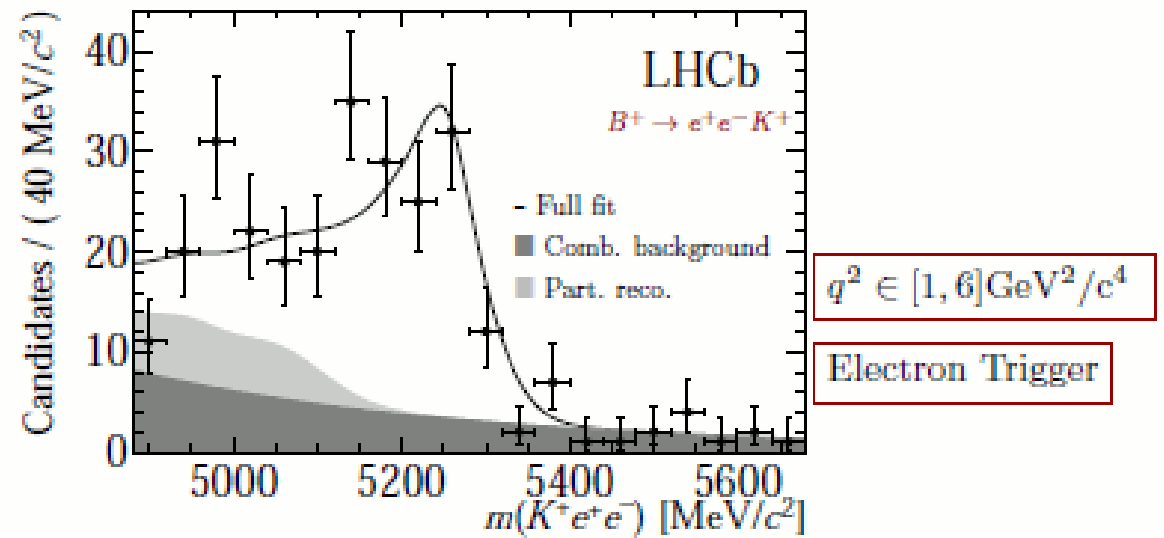
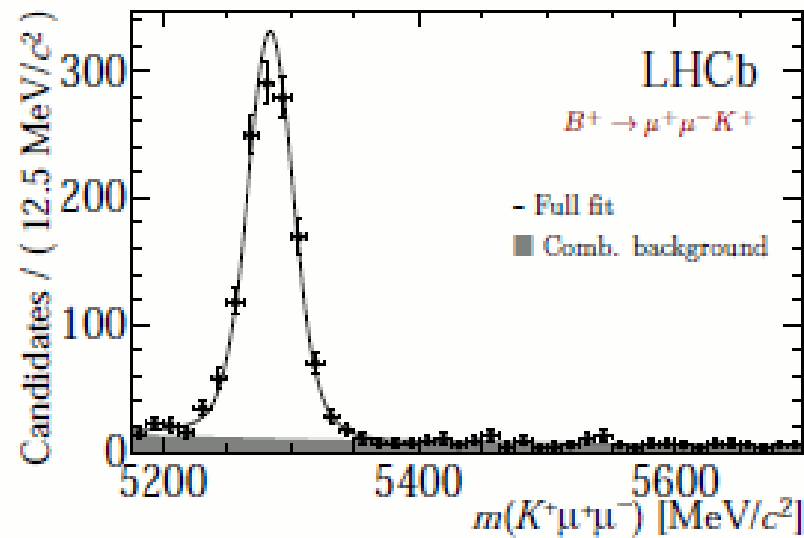
- no deviation for A_{FB} but...
- Form-factor independent observable $P'_5 = \frac{S_5}{\sqrt{F_L(1-F_L)}}$ [arXiv:1512.04442]



- Tension in P'_5 seen with 1 fb^{-1} is confirmed
- Local deviations of 2.9σ and 3.0σ for $q^2 \in [4.0, 6.0]$ and $[6.0, 8.0] \text{ GeV}^2$
- Naive combination of the two gives local significance of 3.7σ

Test of lepton universality using $B^+ \rightarrow K^+ l^+ l^-$ decays

[arXiv:1406.6482]



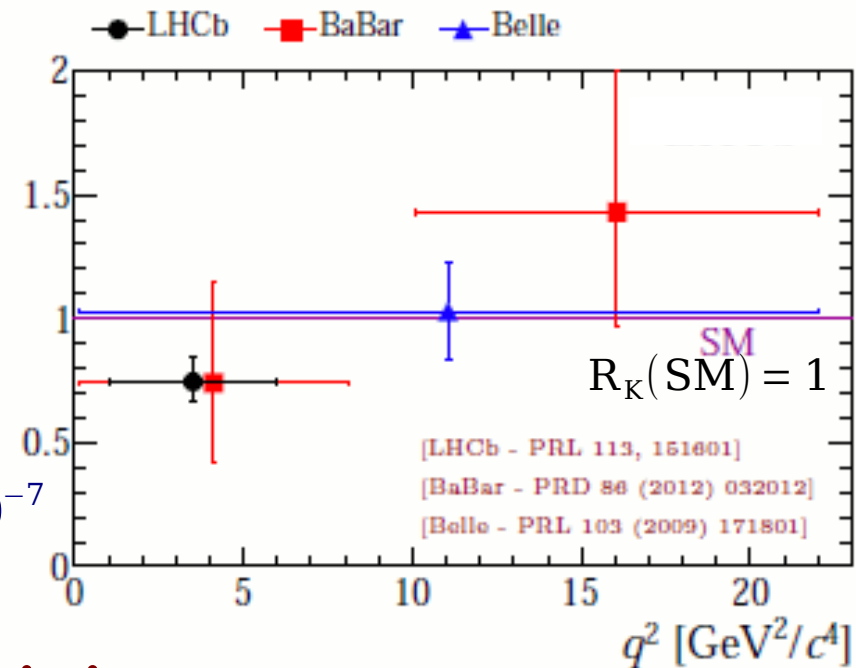
R_K : ratio of branching fractions for dilepton invariant mass squared range $1 < q^2 < 6 \text{ GeV}^2/c^4$

- The combination of the various trigger channels gives:

$$R_K = 0.745^{+0.090}_{-0.074}(\text{stat}) \pm 0.036(\text{syst})$$

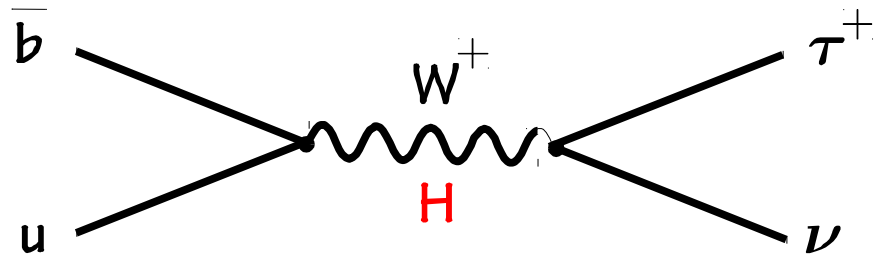
- Most precise measurement to date is in disagreement with SM at 2.6σ level

$\Rightarrow B(B^+ \rightarrow e^+ e^- K^+) = (1.56^{+0.19}_{-0.15}(\text{stat})^{+0.06}_{-0.05}(\text{syst})) \times 10^{-7}$
compatible with SM predictions

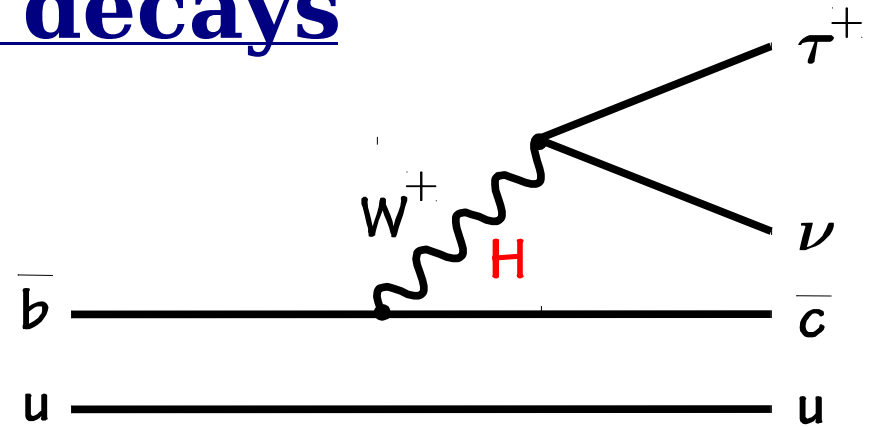


Lepton Flavor Non-Universality ? effect is in $\mu\mu$, not ee

Tauonic B decays



B → τ ν



$$B_{\text{SM}}(B^+ \rightarrow \tau^+ \nu) = \frac{G_F^2 m_B m_\tau^2}{8\pi} \left(1 - \frac{m_\tau^2}{m_B^2}\right) f_B^2 |V_{ub}|^2 \tau_B$$

$$\text{2HDM (type II): } B(B^+ \rightarrow \tau^+ \nu) = B_{\text{SM}} \times \left(1 - \frac{m_B^2}{m_{H^+}^2} \tan^2 \beta\right)^2$$

uncertainties from f_B and $|V_{ub}|$ can be reduced to B_B
and other CKM uncertainties by combining with precise Δm_d

B → D^(*) τ ν

$$\text{2HDM (type II): } B(B \rightarrow D \tau^+ \nu) = G_F^2 \tau_B |V_{cb}|^2 f(\mathbf{F}_V, \mathbf{F}_S, \frac{m_B^2}{m_{H^+}^2} \tan^2 \beta)$$

uncertainties from form factors \mathbf{F}_V and \mathbf{F}_S can be studied
with $B \rightarrow D l \nu$ (more form factors in $B \rightarrow D^* \tau \nu$)

Event reconstruction in $B \rightarrow \tau \nu$

$$\underline{B_{\text{sig}} \rightarrow \tau \nu}$$

(70 % of all τ decays)

$$\tau \rightarrow e \nu \nu, \mu \nu \nu,$$

$$\tau \rightarrow \pi \nu, \pi \pi^0 \nu, 3 \pi \nu$$

e, μ

B_{sig}

$\Upsilon(4S)$

B_{tag}

$\underline{B_{\text{tag}}}$

hadronic tag

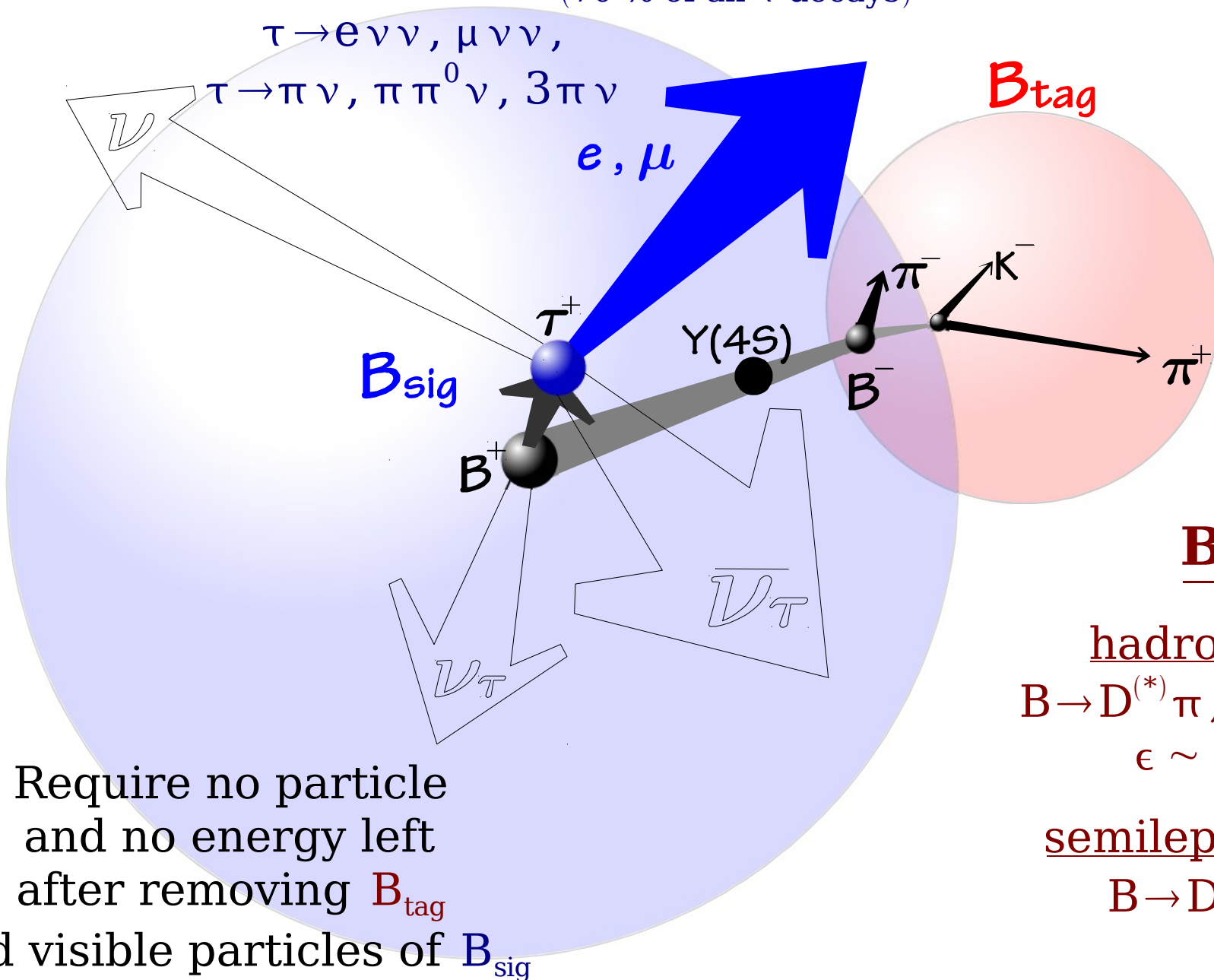
$$B \rightarrow D^{(*)} \pi, D^{(*)} \rho \dots$$

$$\epsilon \sim 0.2\%$$

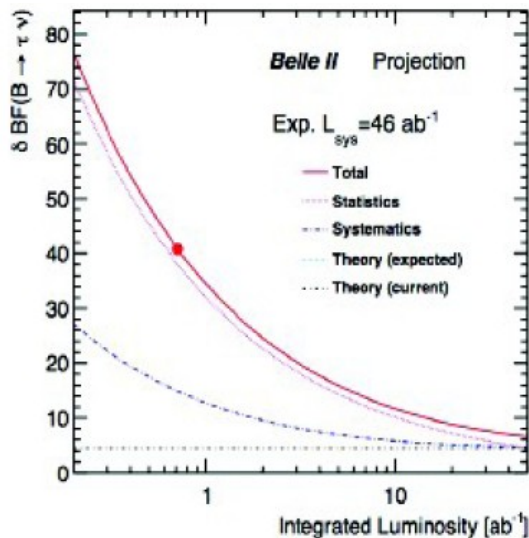
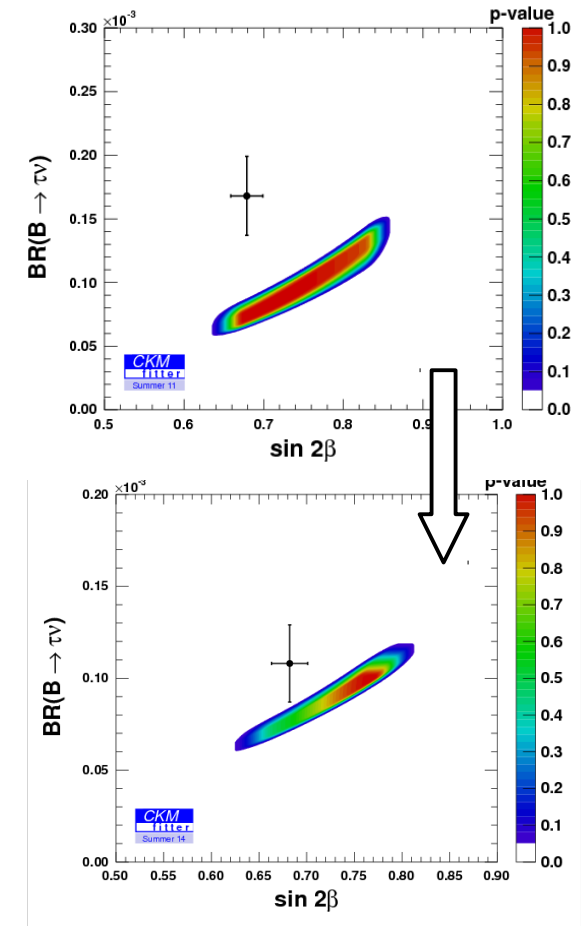
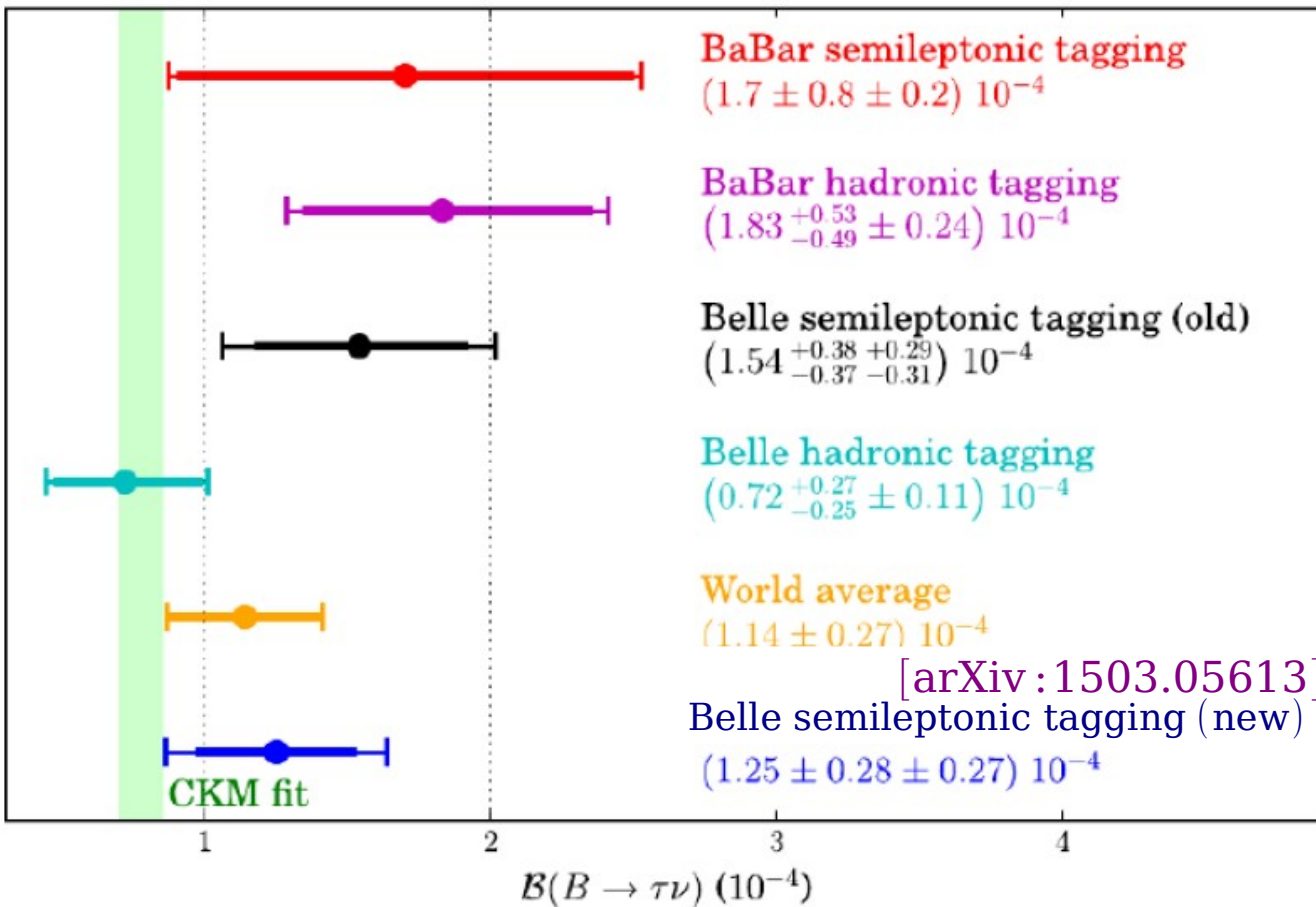
semileptonic tag

$$B \rightarrow D^{(*)} l \nu X$$

Require no particle
and no energy left
after removing B_{tag}
and visible particles of B_{sig}



$B \rightarrow \tau \nu$ status

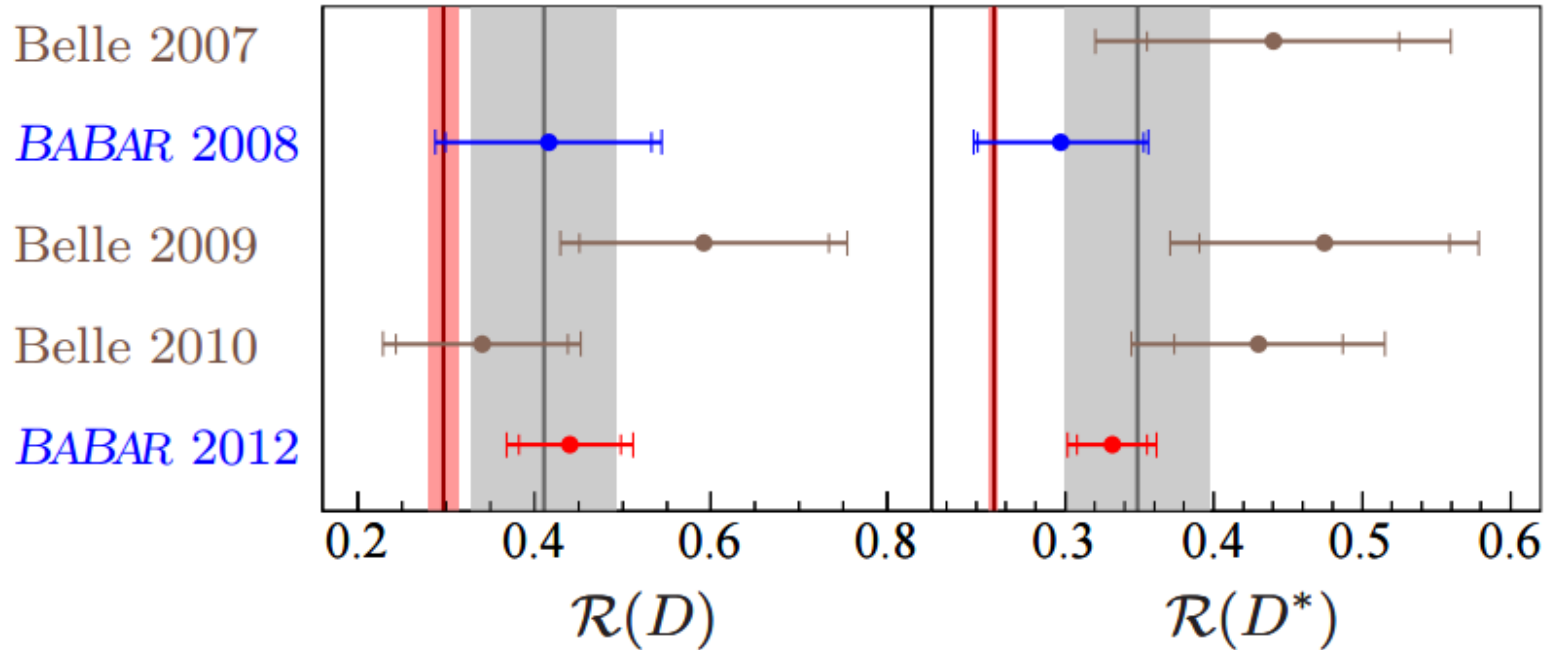


Belle II	Statistical	Systematic	Total Exp	Theory	Total
	(reducible, irreducible)				
$ V_{ub} B \rightarrow \tau \nu$ (had. tagged)					
711 fb^{-1}	19.0	(7.1, 2.2)	20.4	2.5	20.5
5 ab^{-1}	7.2	(2.7, 2.2)	7.9	1.5	8.1
50 ab^{-1}	2.3	(0.8, 2.2)	3.2	1.0	3.4
$ V_{ub} B \rightarrow \tau \nu$ (SL tagged)					
605 fb^{-1}	12.4	(9.0, $^{+3.0}_{-4.8}$)	$^{+15.6}_{-16.1}$	2.5	$^{+15.8}_{-16.2}$
5 ab^{-1}	4.3	(3.1, $^{+3.0}_{-4.8}$)	$^{+6.1}_{-7.2}$	1.5	$^{+6.3}_{-7.3}$
50 ab^{-1}	1.4	(1.0, $^{+3.0}_{-4.8}$)	$^{+3.4}_{-5.1}$	1.0	$^{+3.6}_{-5.2}$

$$\underline{\mathbf{B} \rightarrow \mathbf{D}^{(*)} \tau \nu}$$

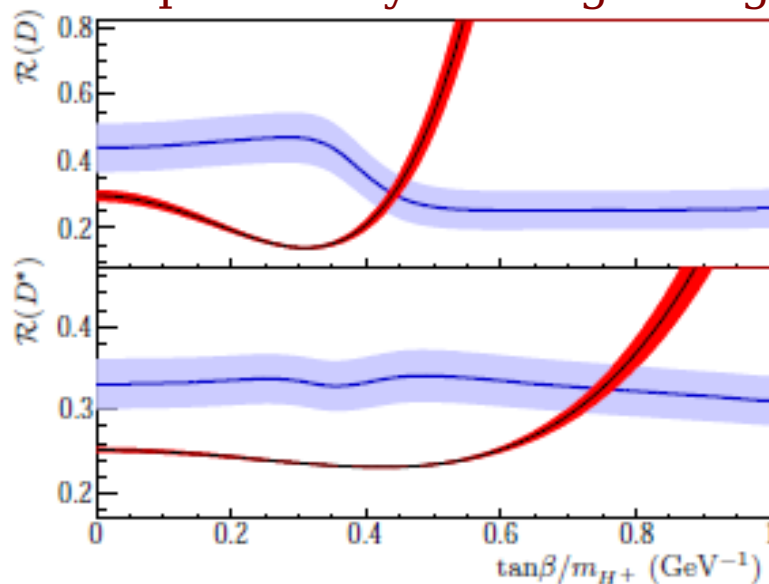
$$R(D^{(*)}) = \frac{B \rightarrow D^{(*)} \tau \nu}{B \rightarrow D^{(*)} l \nu}$$

Babar and Belle measurements hint to deviation from SM



BaBar (arXiv:1303.0571) observes a 3.4σ excess over SM expectation

"This excess cannot be explained by a charged Higgs boson in the 2HDM type II"



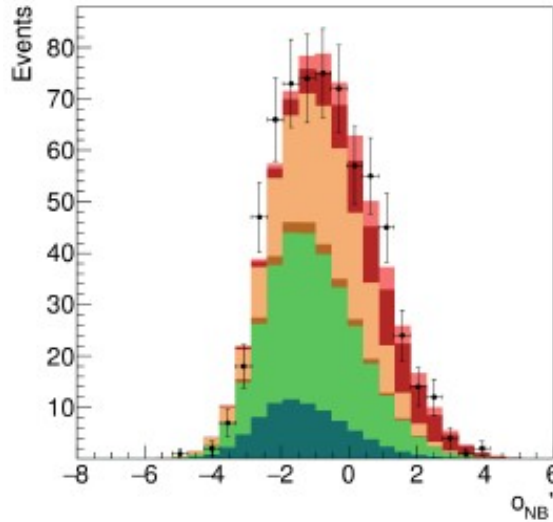
$B \rightarrow D^{(*)} \tau \nu$ at Belle

[arXiv:1507.03233]

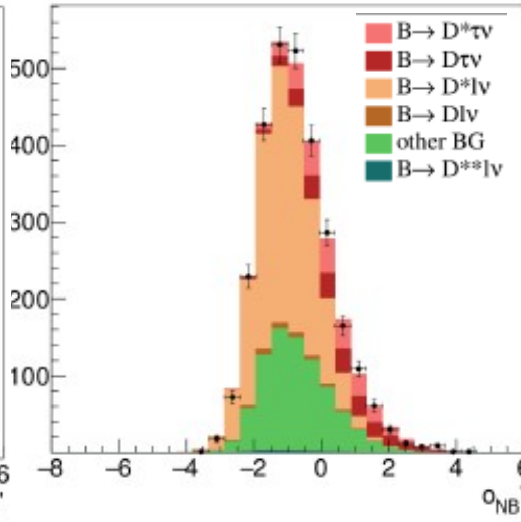
(with hadronic tagging)

projections for large M_{miss}^2 region, $N(D \tau \nu) \sim 300$, $N(D^* \tau \nu) \sim 500$

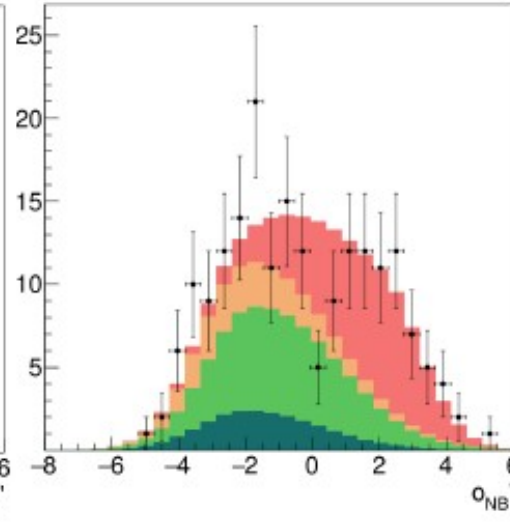
$B \rightarrow D^+ \tau \nu$



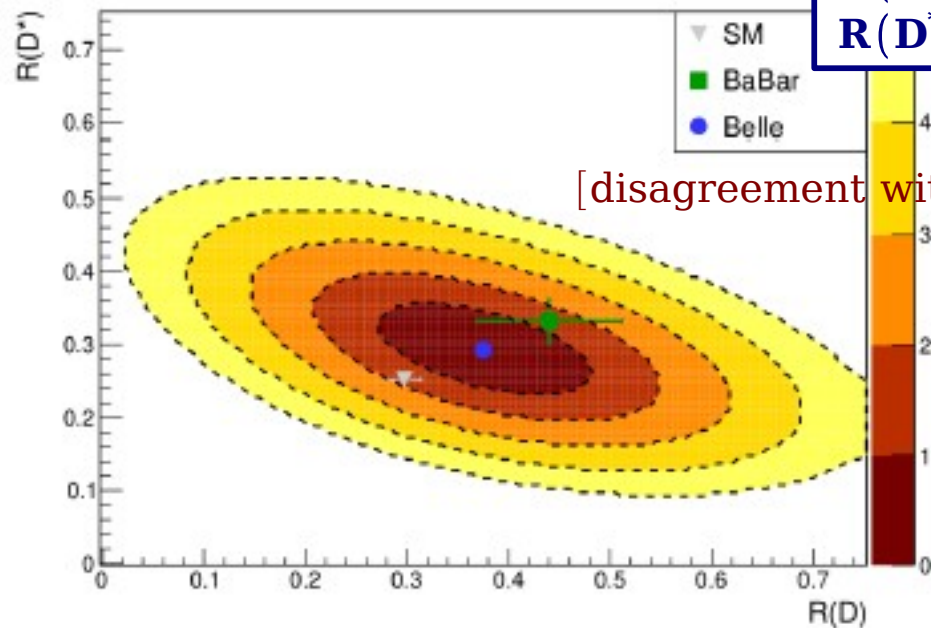
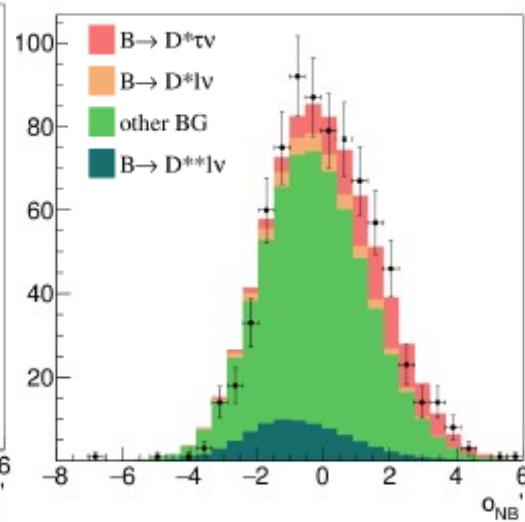
$B \rightarrow D^0 \tau \nu$



$B \rightarrow D^{*+} \tau \nu$



$B \rightarrow D^{*0} \tau \nu$



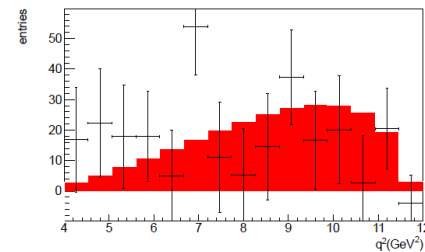
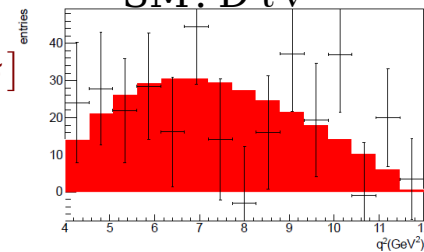
$$R(D) = 0.375 \pm 0.064 \pm 0.026$$

$$R(D^*) = 0.293 \pm 0.038 \pm 0.015$$

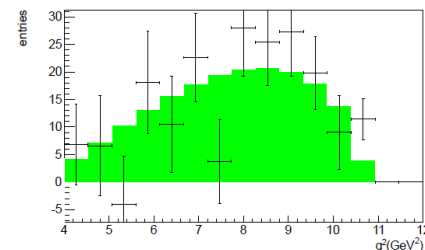
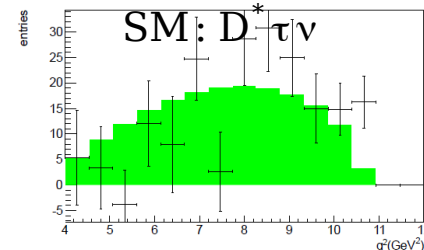
[disagreement with SM at 1.5σ]

stat error only !

SM: $D \tau \nu$



SM: $D^* \tau \nu$



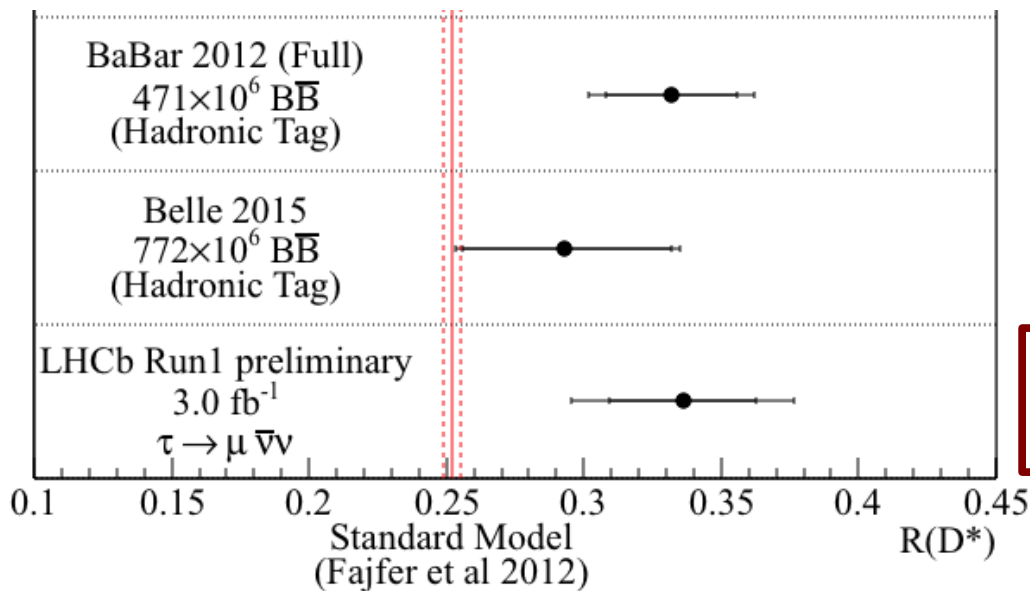
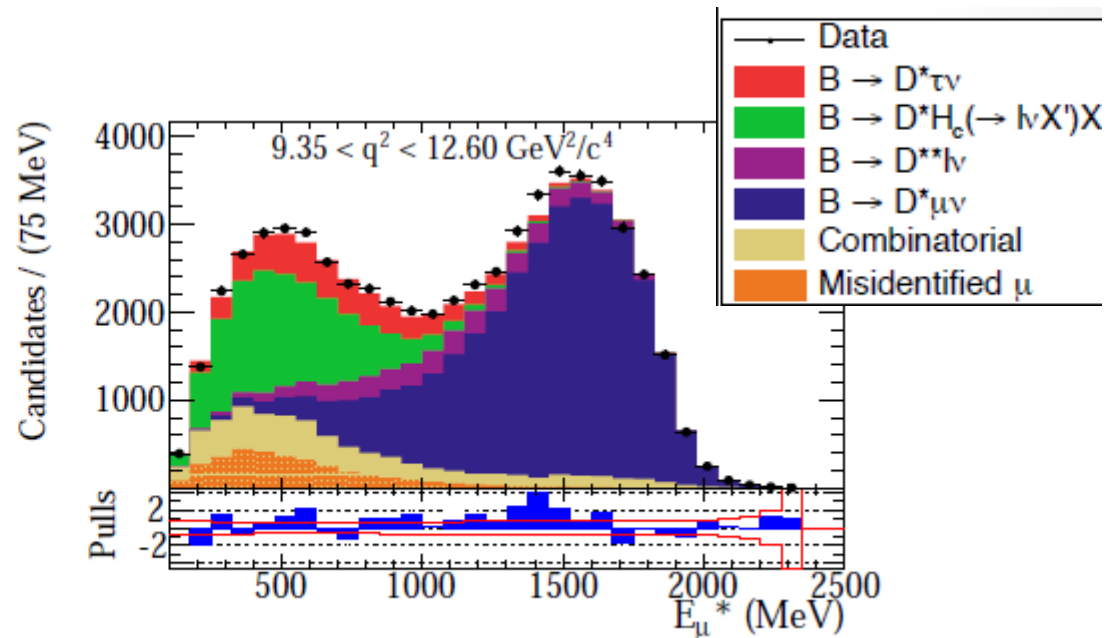
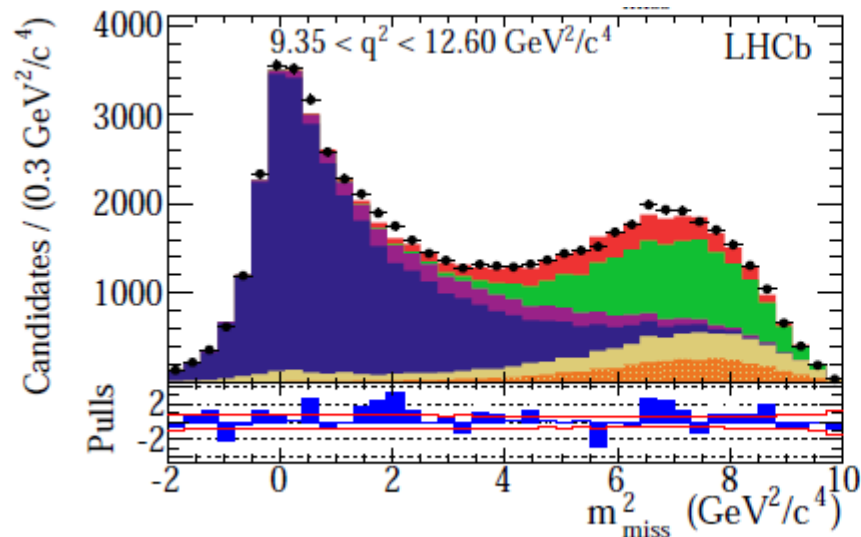
$B \rightarrow D^{*+} \tau \nu$ at LHCb

[arXiv:1506.08614]

$$R(D^*) \equiv \frac{B(\bar{B}^0 \rightarrow D^{*+} \tau^- (\mu^- \bar{\nu}_\mu \nu_\tau) \bar{\nu}_\tau)}{B(\bar{B}^0 \rightarrow D^{*+} \mu^- \bar{\nu}_\mu)}$$

363,000 \pm 1,600 events in $D^* \mu \nu$ sample
 $N(D^* \tau \nu)/N(D^* \mu \nu) = (4.54 \pm 0.46)\%$

$$B(\tau \rightarrow \mu \nu \nu) = (17.41 \pm 0.04)\%$$



$$R(D^*) = 0.332 \pm 0.024 \pm 0.018$$

$$R(D^*) = 0.293 \pm 0.038 \pm 0.015$$

[arXiv:1507.03233]

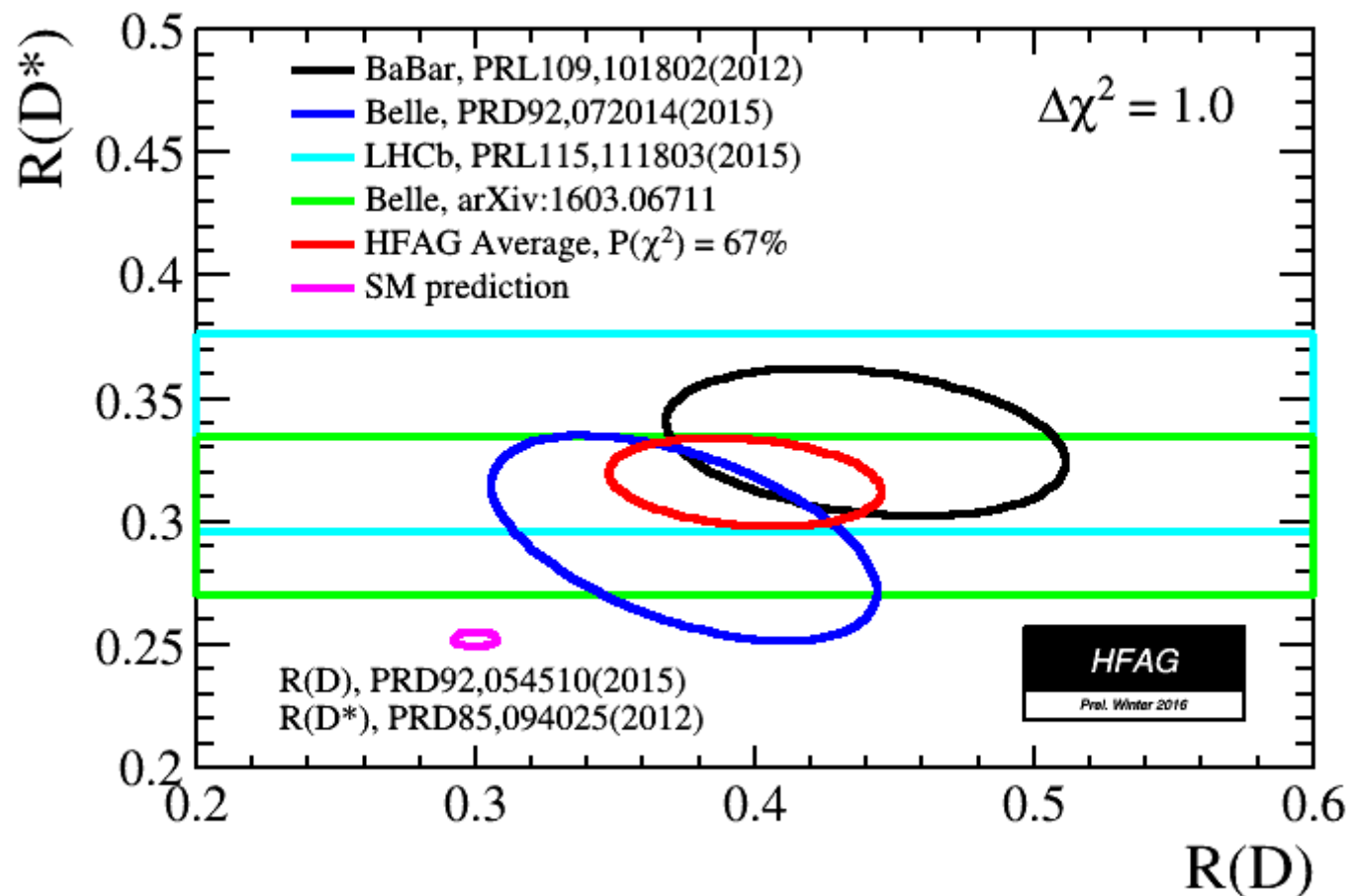
$$R(D^*) = 0.336 \pm 0.027 \pm 0.030$$

[disagreement with SM at 2.1 σ]

[arXiv:1506.08614]

Summary for $B \rightarrow D^{(*)} \tau \nu$

$$\Rightarrow R(D^{(*)}) = \frac{BF(B \rightarrow D \tau \nu_\tau)}{BF(B \rightarrow D l \nu_l)}$$



BaBar

$$R(D) = 0.440 \pm 0.058 \pm 0.042$$

$$R(D^*) = 0.332 \pm 0.024 \pm 0.018$$

Belle

$$R(D) = 0.375 \pm 0.064 \pm 0.026$$

$$R(D^*) = 0.293 \pm 0.038 \pm 0.015$$

$$R(D^*) = 0.302 \pm 0.030 \pm 0.011$$

LHCb

$$R(D^*) = 0.336 \pm 0.027 \pm 0.030$$

average

$$R(D) = 0.397 \pm 0.040 \pm 0.028$$

$$R(D^*) = 0.316 \pm 0.016 \pm 0.010$$

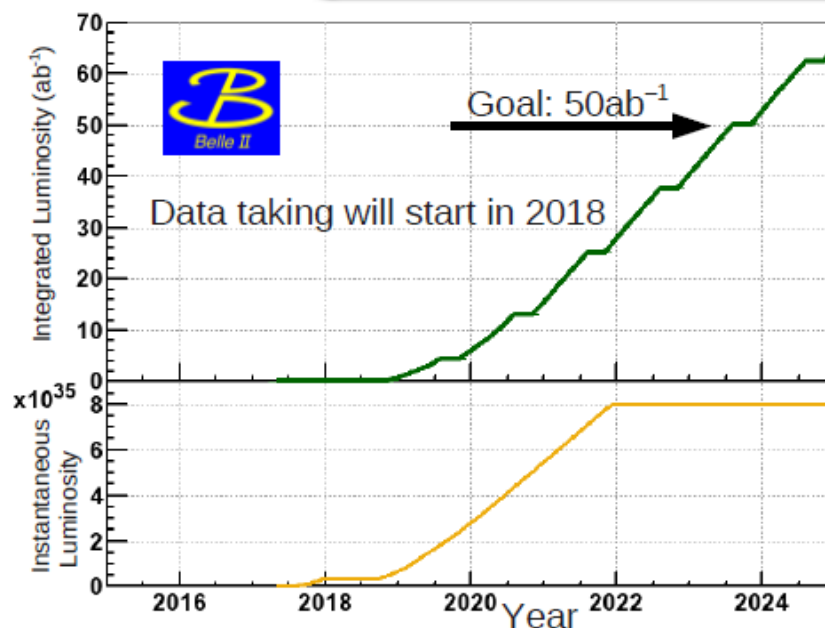
difference with SM predictions
is at **4.0 σ** level

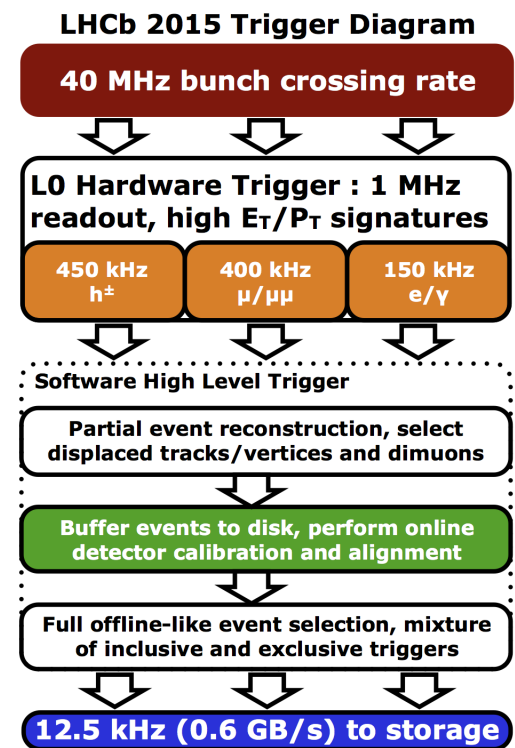
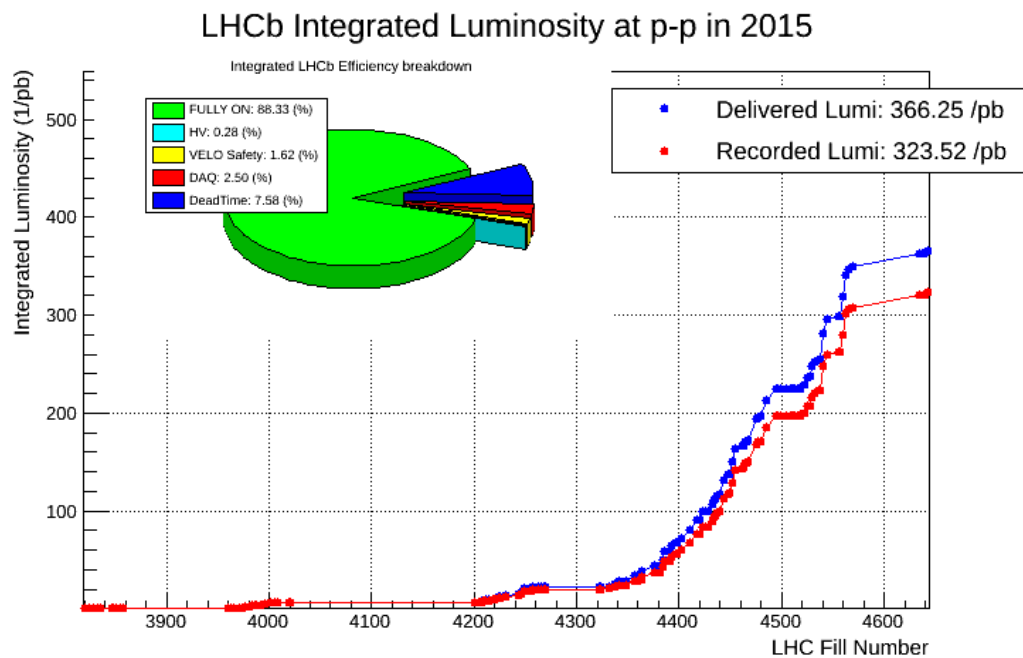
\Rightarrow more measurements to come, more observables (τ polarization...)

Summary

- Few results on CP violation and rare decays in B sector covered in this talk... but much more in B decays, also in charm, charmonium, bottomonium, light Higgs, τ , kaon sectors...
- Definitely not only complementary, but stimulating competition between (super) B-factories and LHCb (upgrade):
 - for the expected: results on $B_{(s)} \rightarrow \mu\mu$, $B \rightarrow K^* \mu\mu$, $B_s \rightarrow J/\psi \phi$, γ angle...
 - for the less expected: results on $|V_{ub}|$, $D^* \tau \nu$...

LHC era			HL-LHC era	
Run 1 (2010-12)	Run 2 (2015-18)	Run 3 (2020-22)	Run 4 (2025-28)	Run 5+ (2030+)
3 fb ⁻¹	8 fb ⁻¹	23 fb ⁻¹	46 fb ⁻¹	100 fb ⁻¹





CPV in $B_s \rightarrow J/\psi \phi$ at LHCb

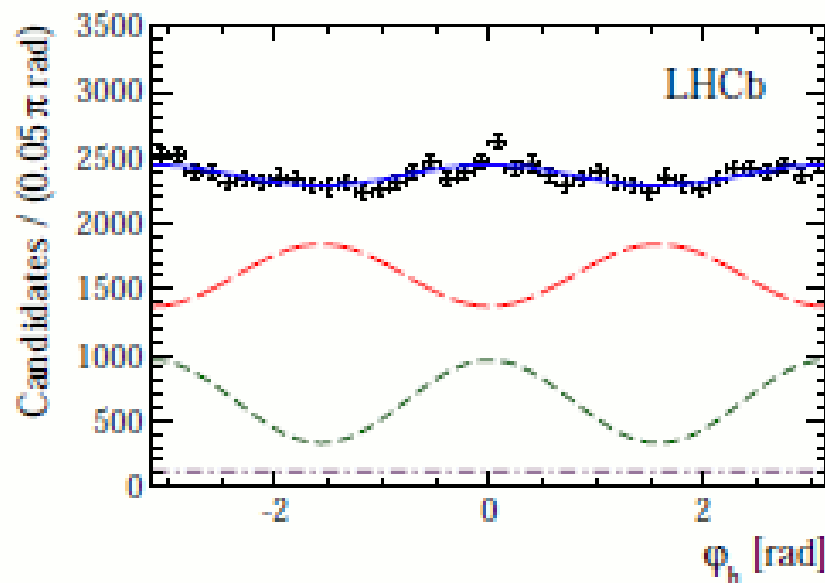
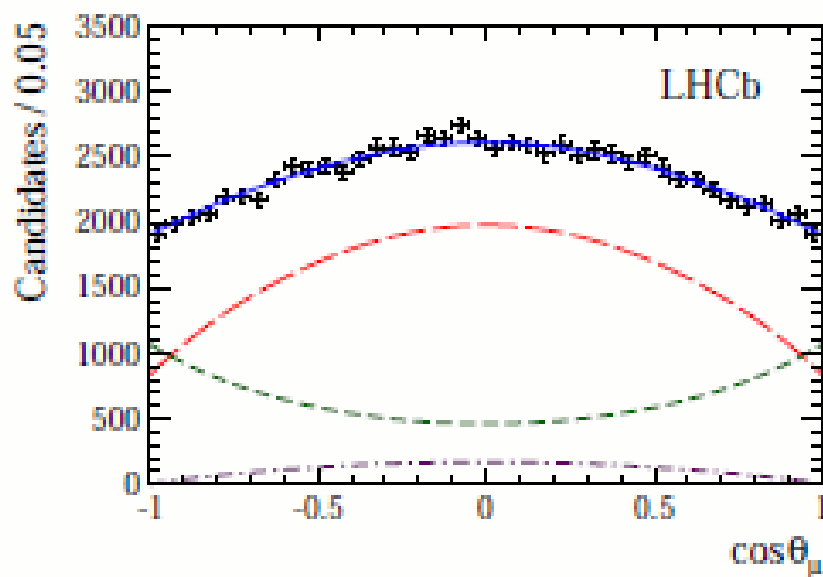
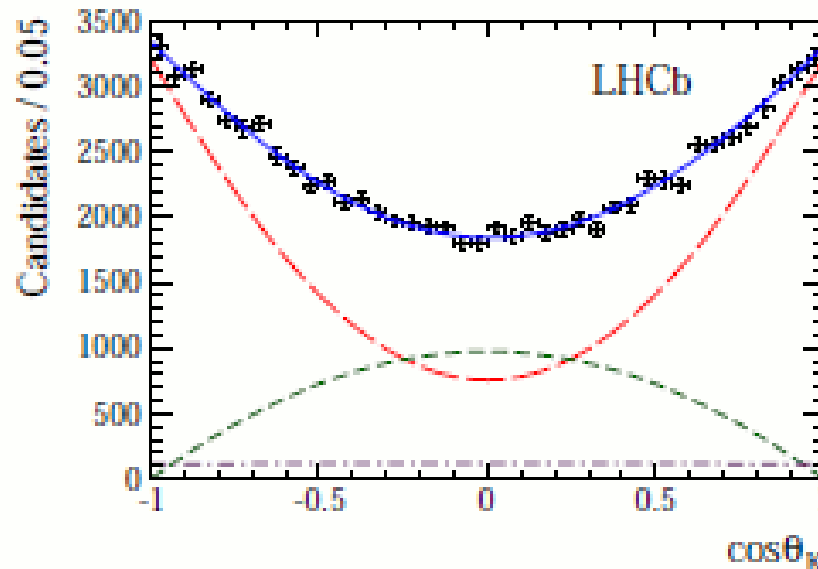
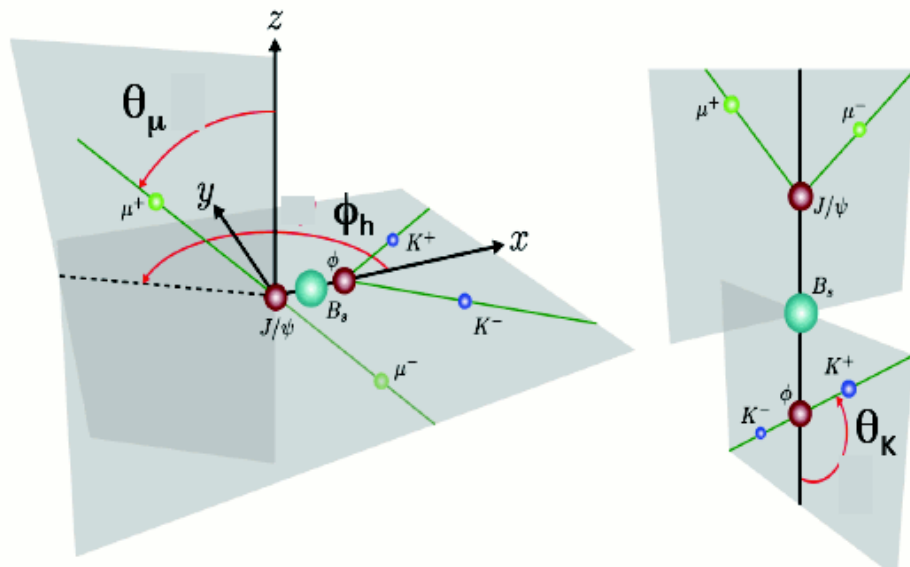
[3 fb⁻¹, arXiv:1411.3104]

CP-even

CP-odd

S-wave (KK)

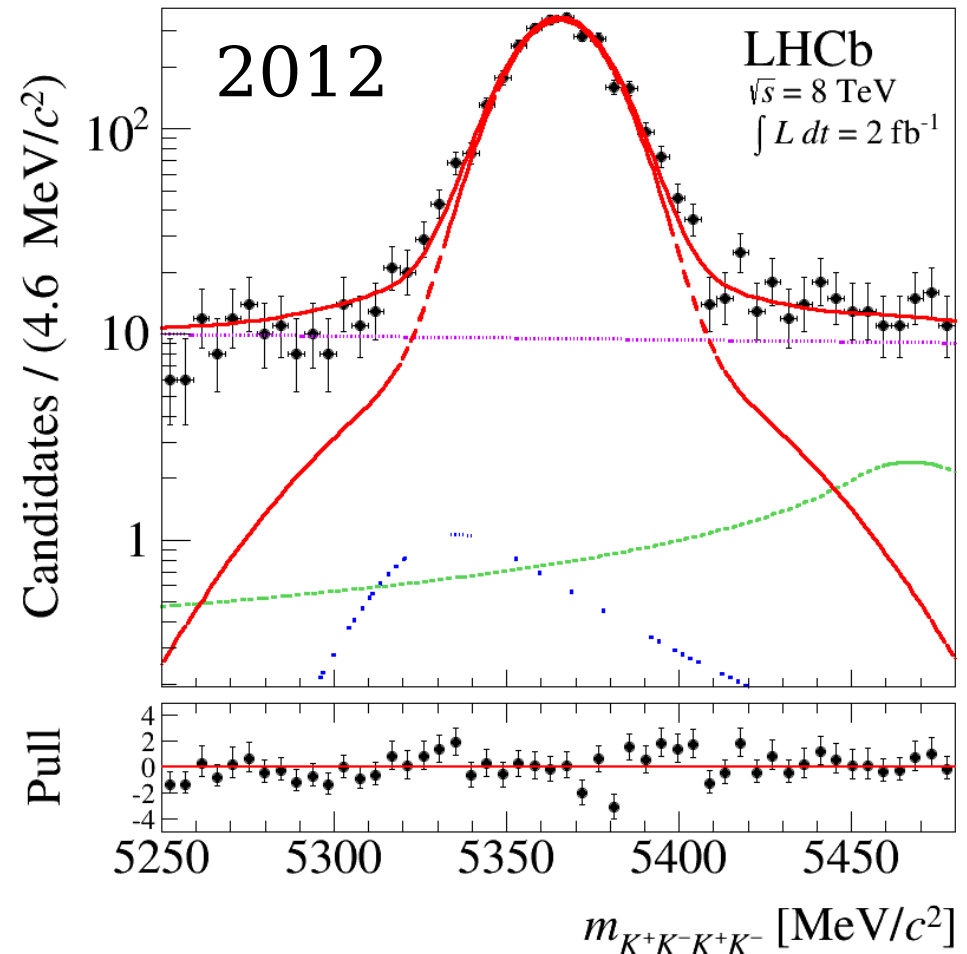
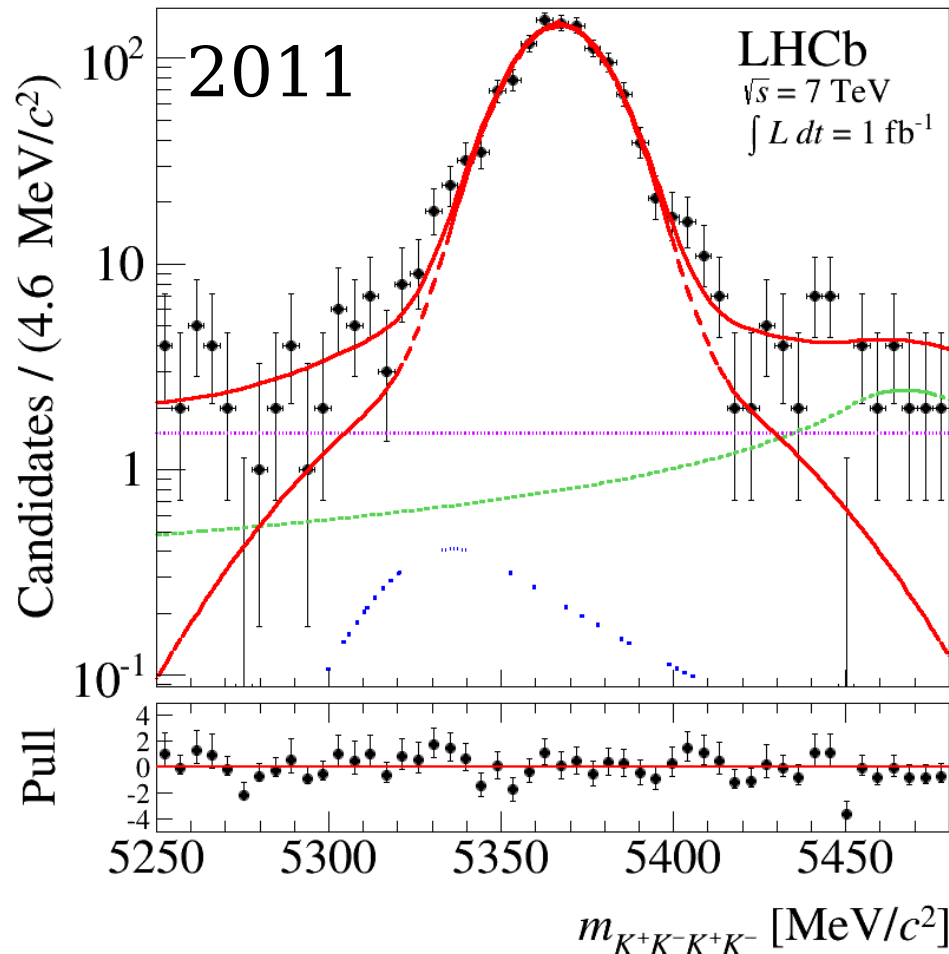
mixture of CP eigenstates \Rightarrow angular analysis in helicity basis



$B_s \rightarrow \varphi \varphi$

[arXiv:1407.2222]

$\bar{b} \rightarrow \bar{s} s \bar{s}$ loop process



- 4000 signal events
- Combinatorial background is flat and small
- Very small contributions from mis-ID of $B_d \rightarrow \varphi K^{*0}$ and $\Lambda_b \rightarrow \varphi p K$
- mixture of CP eigenstates \Rightarrow angular analysis in helicity basis

$$\varphi_s = -0.17 \pm 0.15 \pm 0.03 \text{ rad}$$

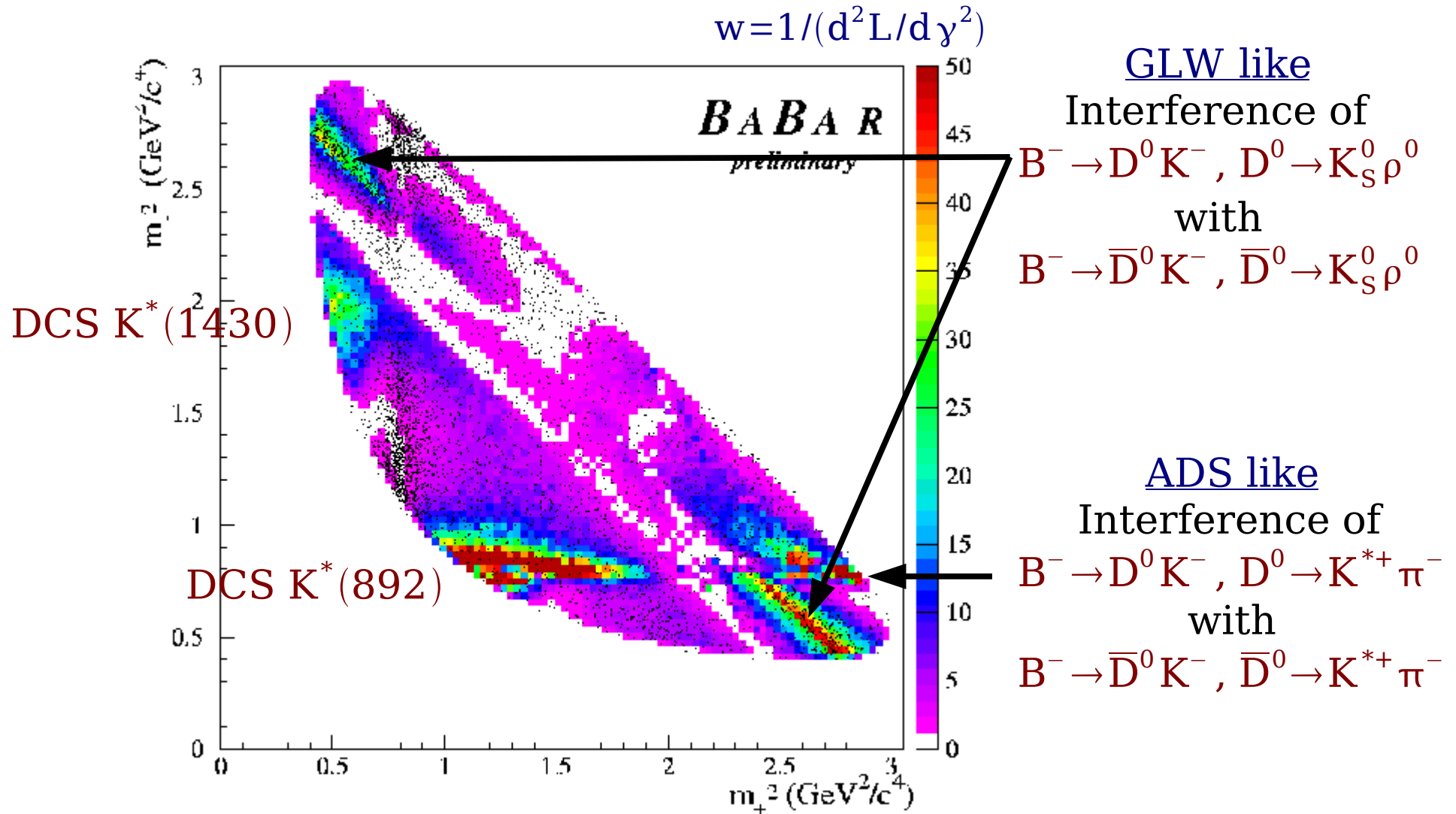
$$\varphi_s(c\bar{c}s) \sim -0.01 \pm 0.04 \text{ rad}$$

$$\varphi_s(\text{SM}) = -0.0363^{+0.0012}_{-0.0014}$$

Sensitivity to γ in $B \rightarrow D(K_S \pi \pi) K$ mode

sensitivity to γ/ϕ_3 varies across the Dalitz plot

$\gamma = 75^\circ$, $\delta = 180^\circ$, $r_B = 0.125$



$B \rightarrow D^{(*)} K^{(*)}$ Dalitz analysis

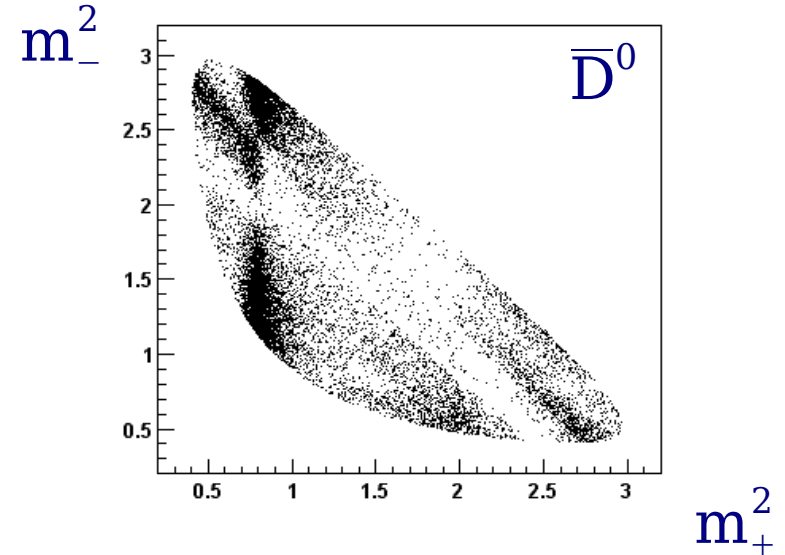
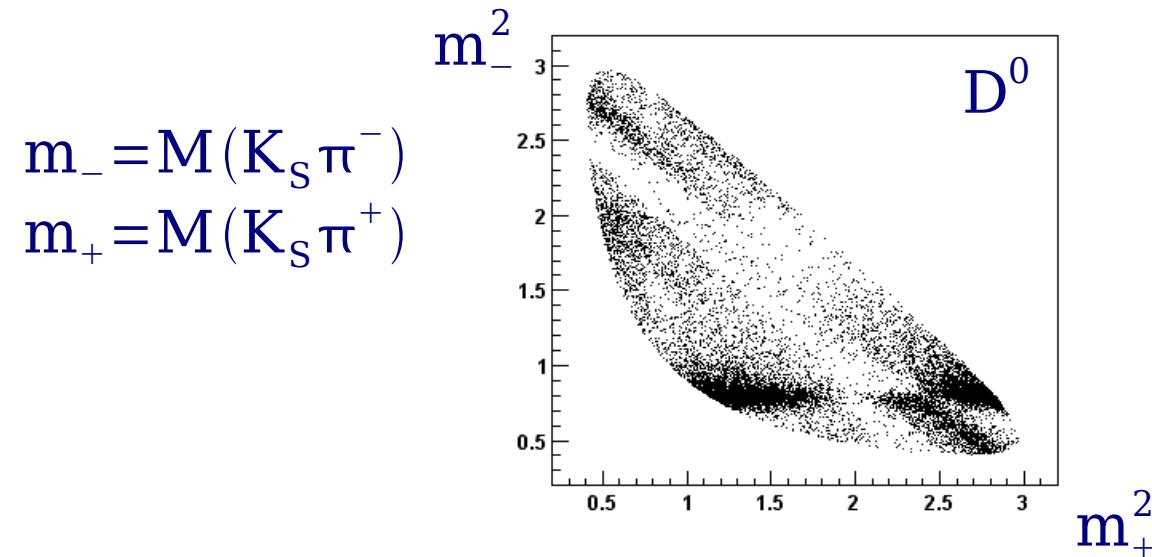
Reconstruction of three-body final states $D^0, \bar{D}^0 \rightarrow K_S \pi^+ \pi^-$

Amplitude for each Dalitz point is described as:

$$\bar{D}^0 \rightarrow K_S \pi^+ \pi^- \sim f(m_+^2, m_-^2)$$

$$D^0 \rightarrow K_S \pi^+ \pi^- \sim f(m_-^2, m_+^2)$$

$$B^+ \rightarrow (K_S \pi^+ \pi^-)_D K^+ : f(m_+^2, m_-^2) + r_B e^{i(\delta_B + \gamma)} f(m_-^2, m_+^2)$$

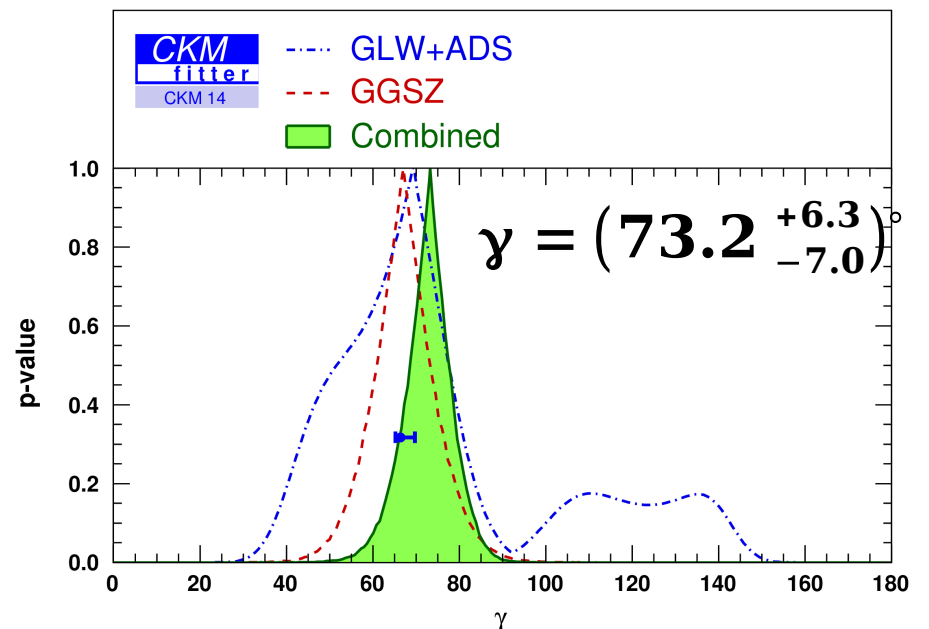
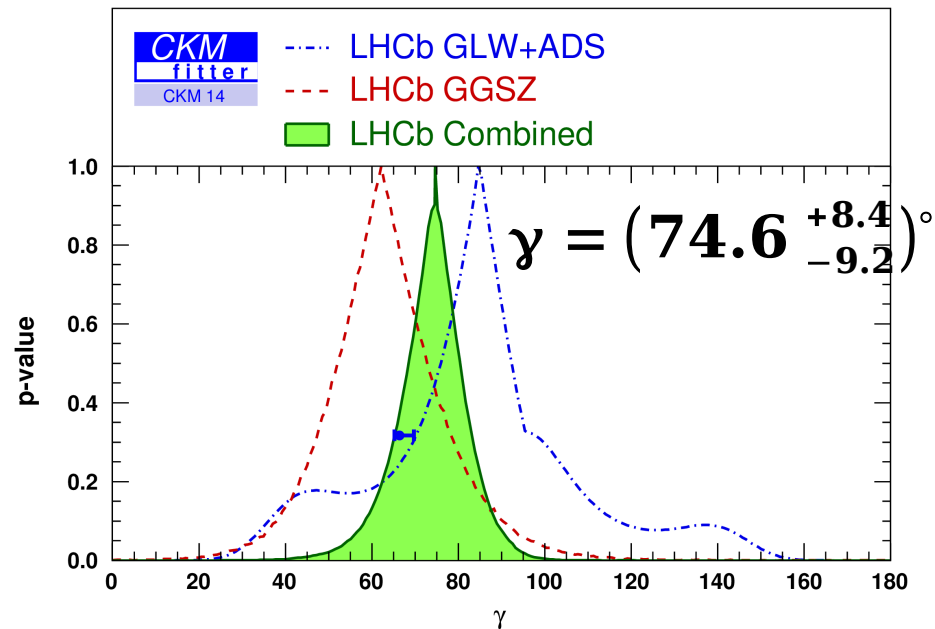
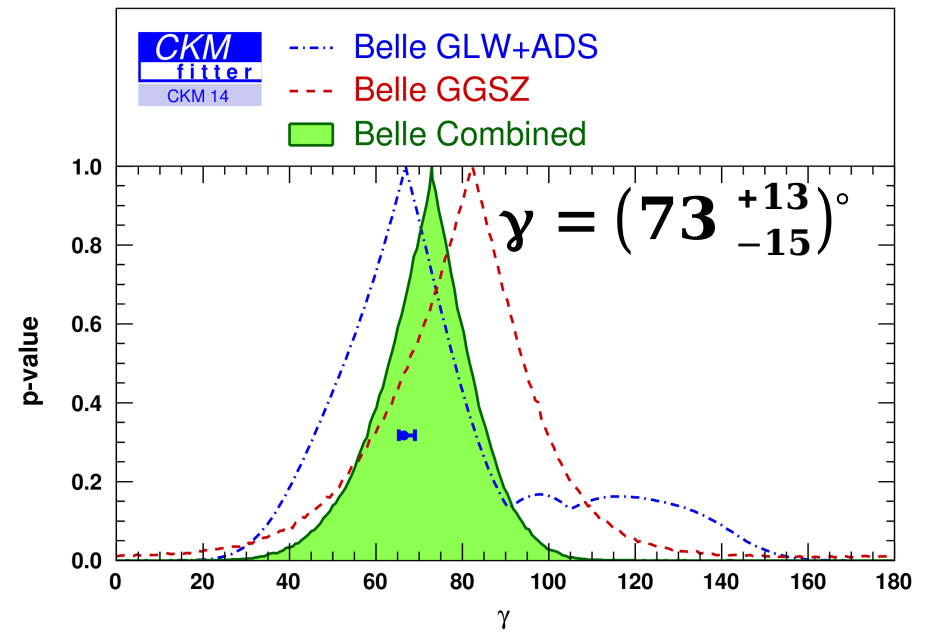
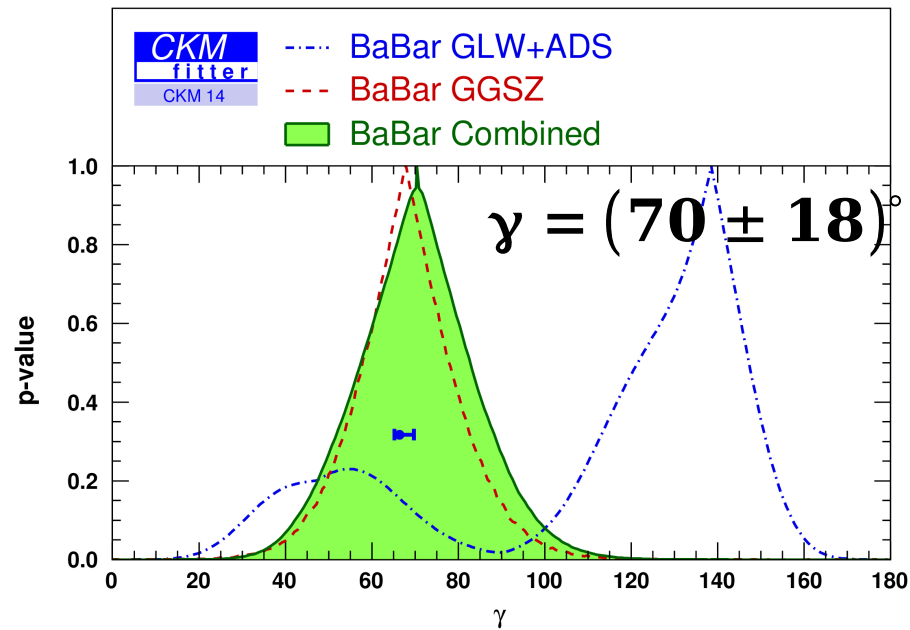


$$B^- \rightarrow (K_S \pi^+ \pi^-)_D K^- : f(m_-^2, m_+^2) + r_B e^{i(\delta - \gamma)} f(m_+^2, m_-^2)$$

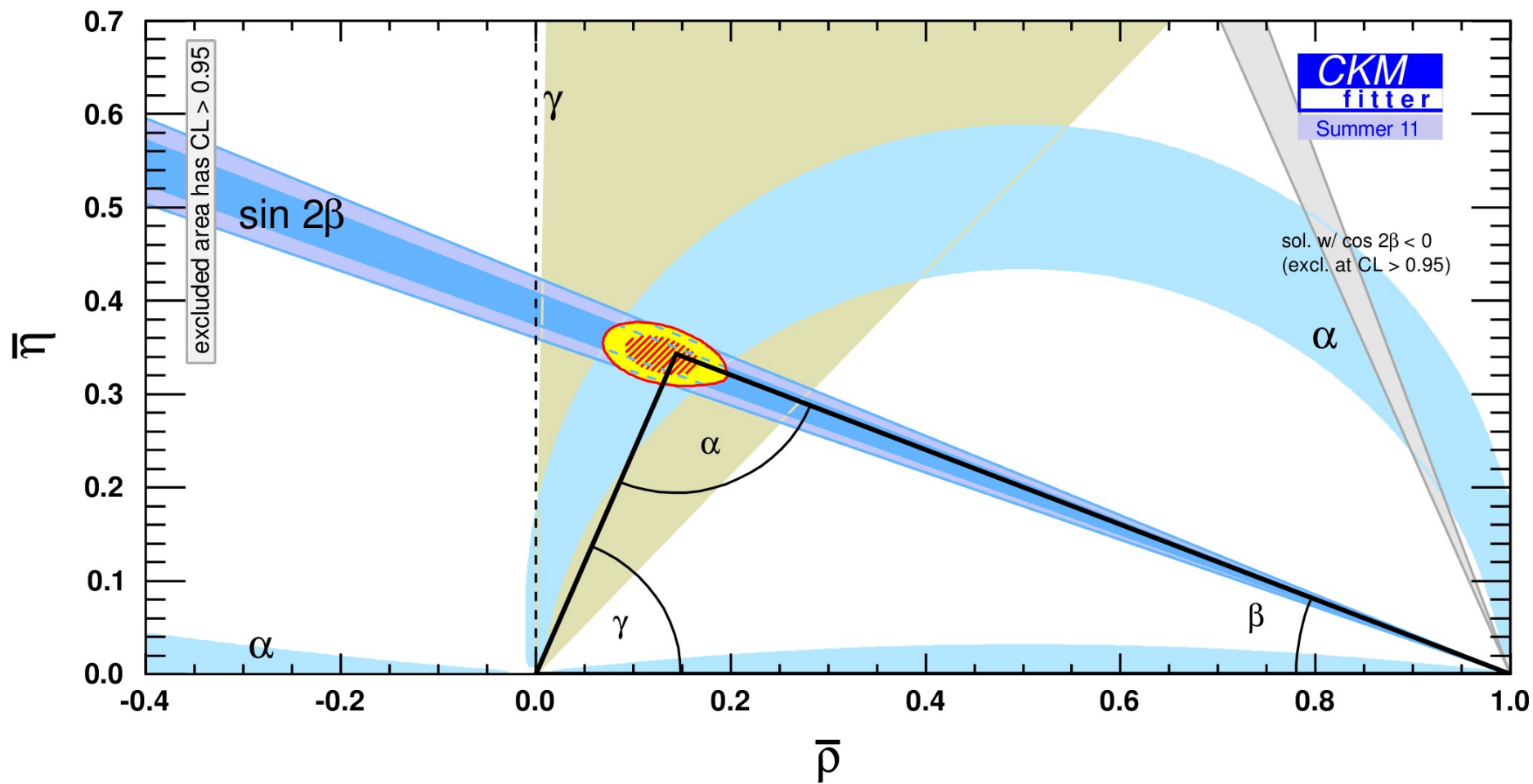
Simultaneous fit of B^+ and B^- to extract parameters r_B , γ and δ_B

Note: 2 fold ambiguity on γ : $(\gamma, \delta_B) \rightarrow (\gamma + \pi, \delta_B + \pi)$

Experiment by experiment

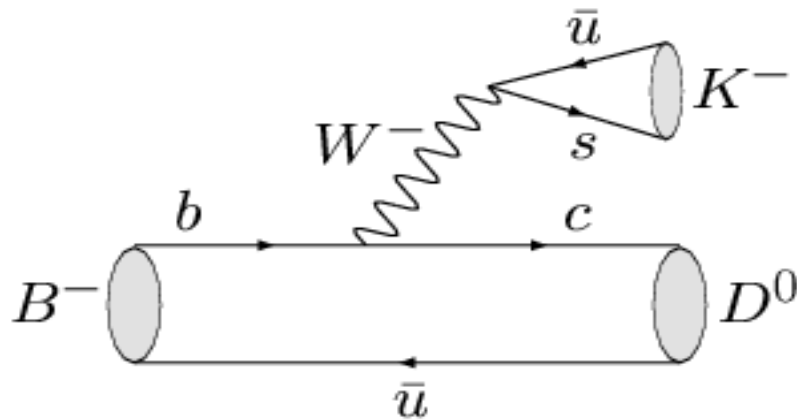


γ angle in the global fit

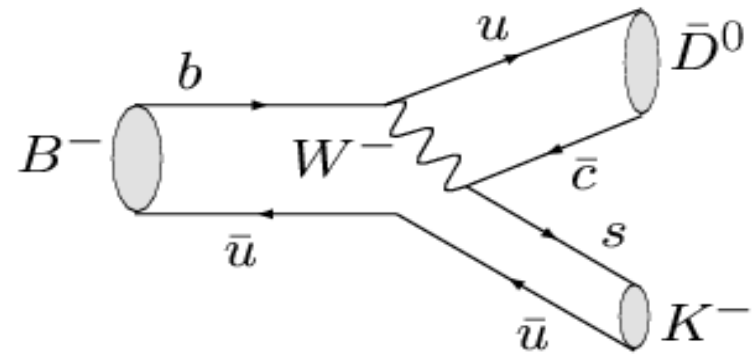


γ measurements from $B^\pm \rightarrow DK^\pm$

- Theoretically pristine $B \rightarrow DK$ approach
- Access γ via interference between $B^- \rightarrow D^0 K^-$ and $B^- \rightarrow \bar{D}^0 K^-$



color allowed
 $B^- \rightarrow D^0 K^- \sim V_{cb} V_{us}^*$
 $\sim A\lambda^3$



color suppressed
 $B^- \rightarrow \bar{D}^0 K^- \sim V_{ub} V_{cs}^*$
 $\sim A\lambda^3(\rho + i\eta)$

relative magnitude of suppressed amplitude is r_B

$$r_B = \frac{|A_{\text{suppressed}}|}{|A_{\text{favoured}}|} \sim \frac{|V_{ub} V_{cs}^*|}{|V_{cb} V_{us}^*|} \times [\text{color supp}] = 0.1 - 0.2$$

relative weak phase is γ , relative strong phase is δ_B

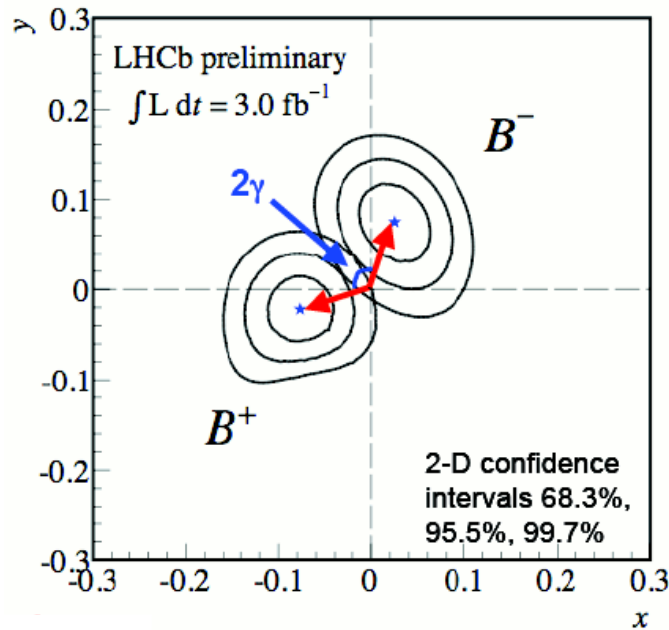
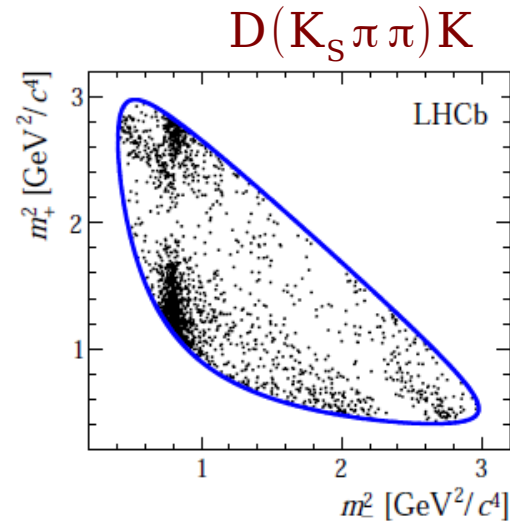
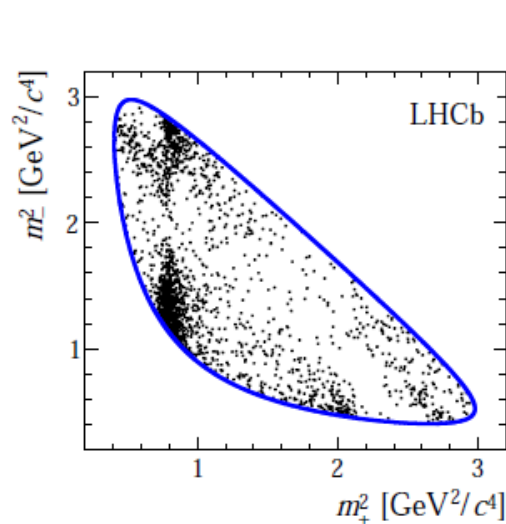
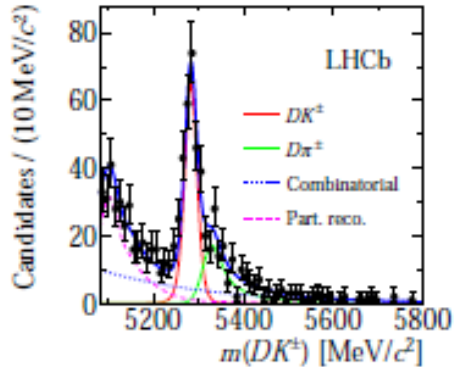
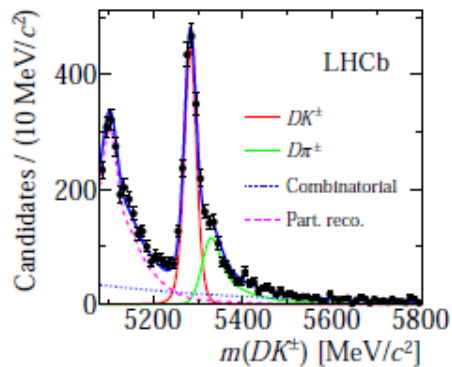
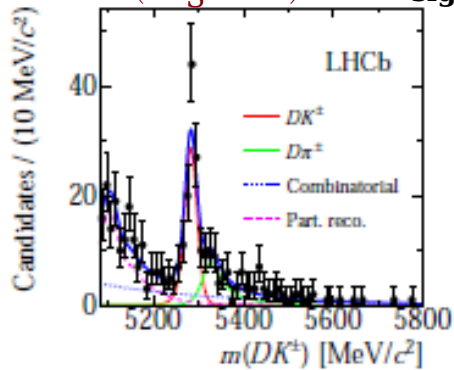
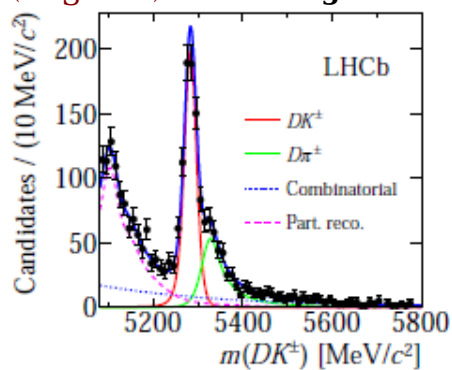
γ measurements from $B^\pm \rightarrow DK^\pm$

- Reconstruct D in final states accessible to both D^0 and \bar{D}^0
 - $D = D_{\text{CP}}$, CP eigenstates as $K^+ K^-$, $\pi^+ \pi^-$, $K_S \pi^0$
GLW method (Gronau-London-Wyler)
 - $D = D_{\text{sup}}$, Doubly-Cabbibo suppressed decays as $K \pi$
ADS method (Atwood-Dunietz-Soni)
 - Three-body decays as $D \rightarrow K_S \pi^+ \pi^-$, $K_S K^+ K^-$
GGSZ (Dalitz) method (Giri-Grossman-Soffer-Zupan)
- Largest effects due to
 - charm mixing
 - charm CP violation

} small, can be included
Y.Grossman, A.Soffer, J.Zupan
[PRD 72, 031501 (2005)]
- Different B decays (DK , $D^* K$, DK^*)
 - different hadronic factors (r_B , δ_B) for each

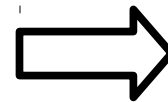
GGSZ LHCb Results [arXiv:1408.2748]

$D(K_S \pi \pi)K$ $N_{\text{sig}} \sim 2260$ $D(K_S KK)K$ $N_{\text{sig}} \sim 324$



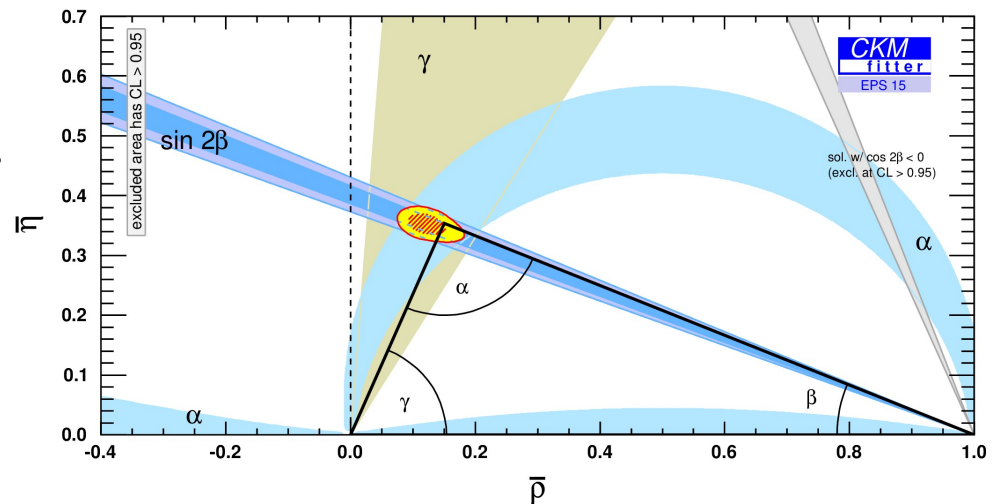
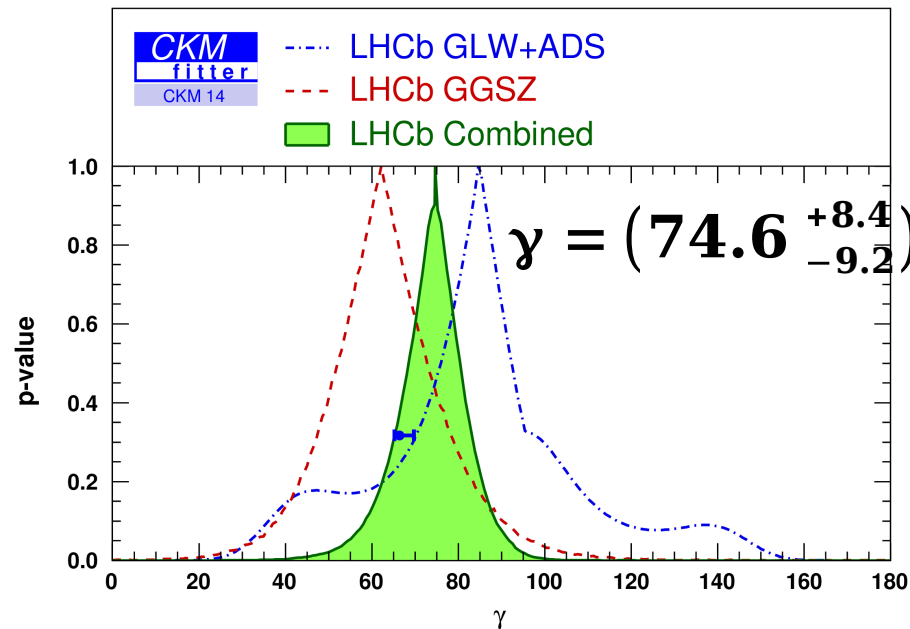
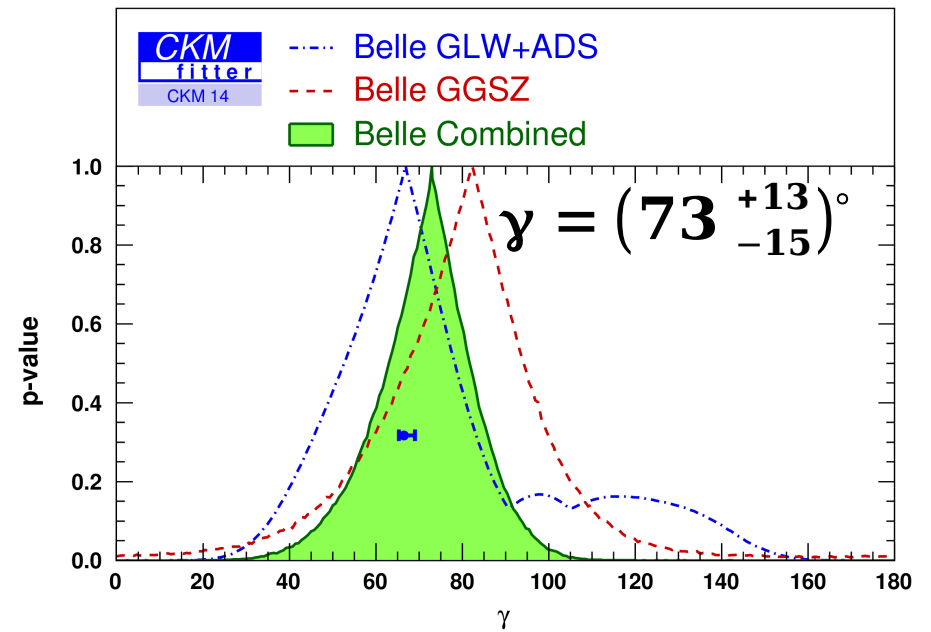
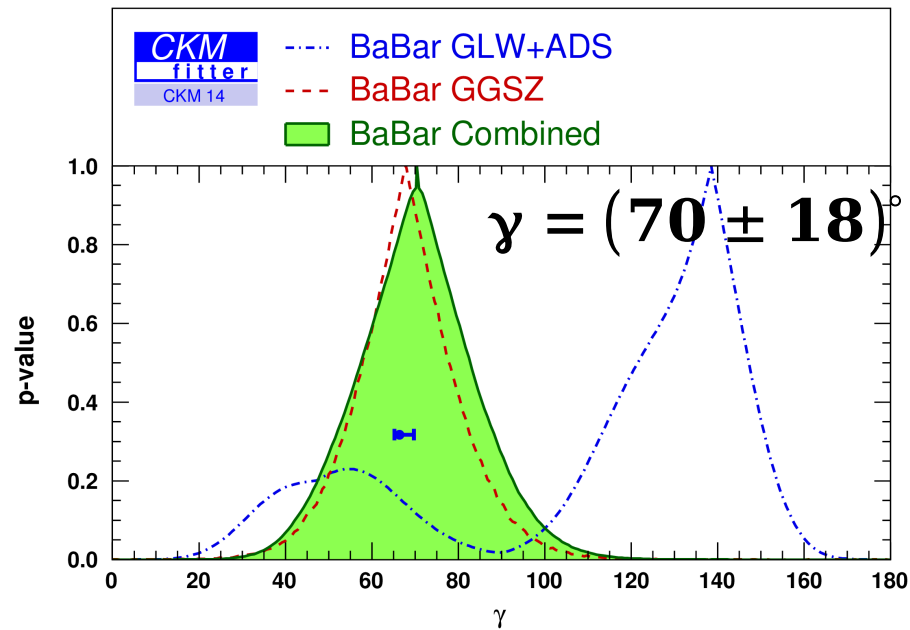
cartesian coordinates

$$\begin{cases} x^\pm = r_B \cos(\delta_B \pm \gamma) \\ y^\pm = r_B \sin(\delta_B \pm \gamma) \end{cases}$$



$$\begin{aligned} \gamma &= (62_{-14}^{+15})^\circ \\ r_B &= (8.0_{-2.1}^{+1.9}) \times 10^{-2} \\ \delta_B &= (134_{-15}^{+14})^\circ \end{aligned}$$

Experiment by experiment



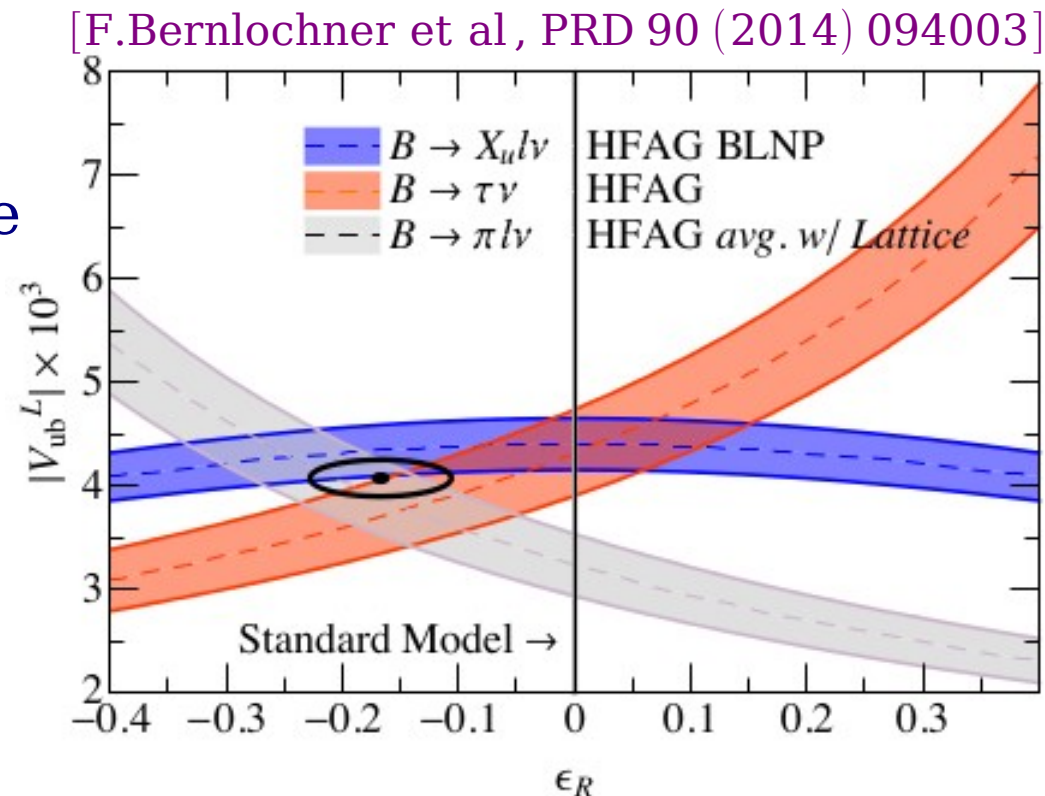
long way to go ... ($\rightarrow \sigma_\gamma = 1^\circ$ or less)

Could it be due to new physics ?

- $B \rightarrow \pi l \nu$ is a purely vector current, whereas $B \rightarrow X_u l \nu$ is V-A
- Adding right-handed current (V+A), increases vector current but decreases axial-vector current

A negative right-handed current can reduce tension between those two results

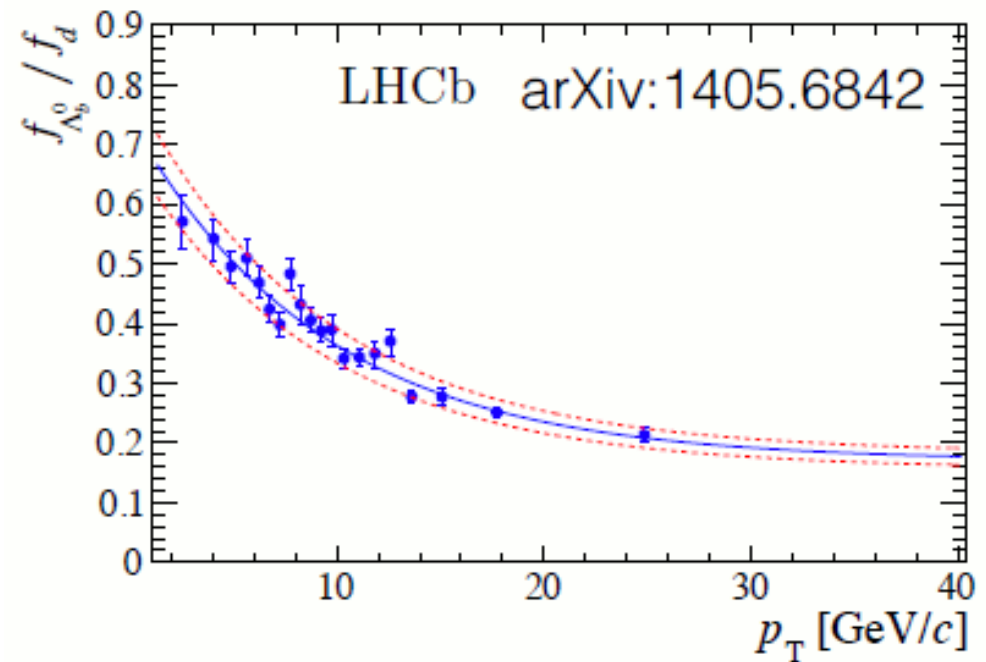
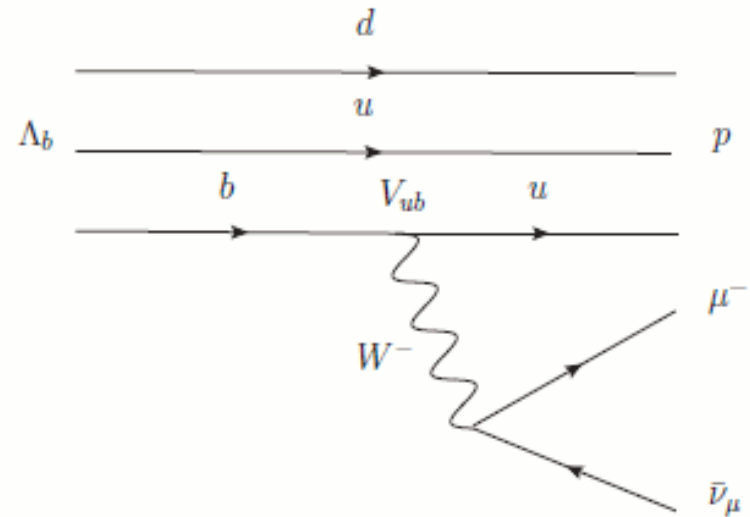
Decay	$ V_{ub} \times 10^3$	ϵ_R dependence
$B \rightarrow \pi \ell \bar{\nu}$	3.23 ± 0.30	$1 + \epsilon_R$
$B \rightarrow X_u \ell \bar{\nu}$	4.39 ± 0.21	$\sqrt{1 + \epsilon_R^2}$
$B \rightarrow \tau \bar{\nu}_\tau$	4.32 ± 0.42	$1 - \epsilon_R$



New measurements needed, with different approaches also

The decay $\Lambda_b^0 \rightarrow p \mu \nu$

- The decay $\Lambda_b^0 \rightarrow p \mu \nu$ is the baryonic version of $B \rightarrow \pi l \nu$
- Cleaner at LHCb as protons are rarer than kaons/pions
- Λ_b^0 baryons not produced at BaBar or Belle experiments but at the LHC produced 1/4 as often as B mesons

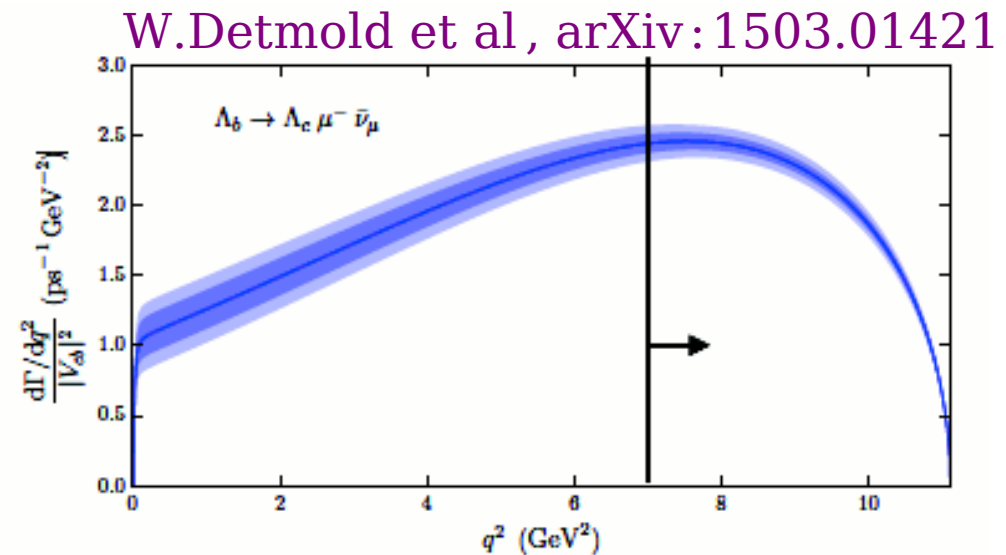
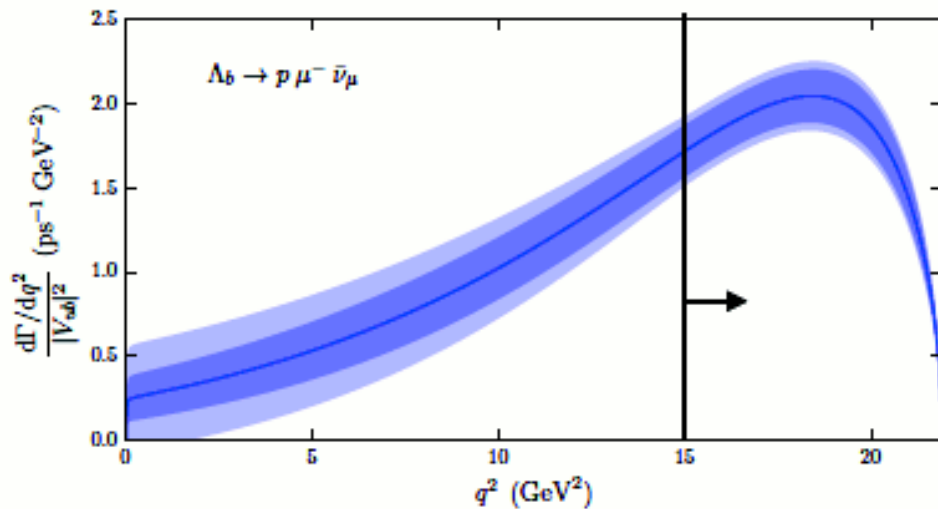


⇒ Signature in detector: displaced muon - proton vertex

Analysis strategy

arXiv:1504.01568

- Normalize signal yield to V_{cb} decay, $\Lambda_b^0 \rightarrow \Lambda_c \mu \nu$
 - \Rightarrow Cancel many systematic uncertainties (including the production rate of Λ_b baryons)
- Calculate the branching fraction ratio at high q^2



$$\frac{\mathcal{B}(\Lambda_b^0 \rightarrow p \mu^- \bar{\nu}_\mu)_{q^2 > 15 \text{ GeV}^2}}{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \mu \nu)_{q^2 > 7 \text{ GeV}^2}} = \frac{N(\Lambda_b^0 \rightarrow p \mu^- \bar{\nu}_\mu)_{q^2 > 15 \text{ GeV}^2}}{N(\Lambda_b^0 \rightarrow (\Lambda_c^+ \rightarrow p K^- \pi^+) \mu \nu)_{q^2 > 7 \text{ GeV}^2}} \times \frac{\epsilon(\Lambda_b^0 \rightarrow (\Lambda_c^+ \rightarrow p K^- \pi^+) \mu \nu)_{q^2 > 7 \text{ GeV}^2}}{\epsilon(\Lambda_b^0 \rightarrow p \mu^- \bar{\nu}_\mu)_{q^2 > 15 \text{ GeV}^2}} \times \mathcal{B}(\Lambda_c^+ \rightarrow p K^- \pi^+)$$

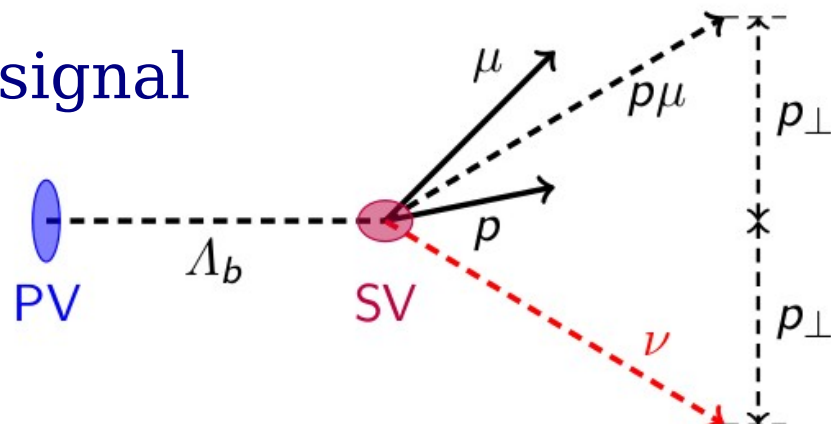
measured precisely
by Belle !

Signal fit

arXiv:1504.01568

Corrected mass used to extract the signal

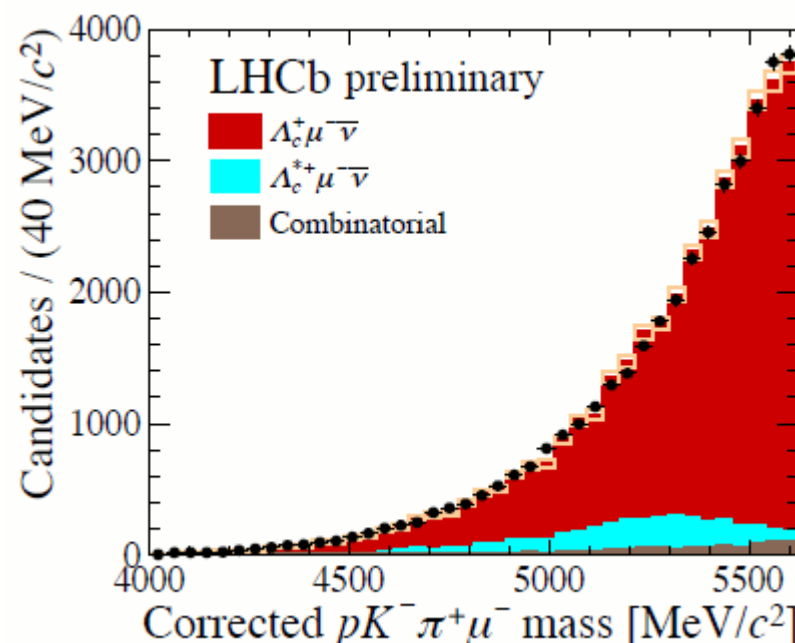
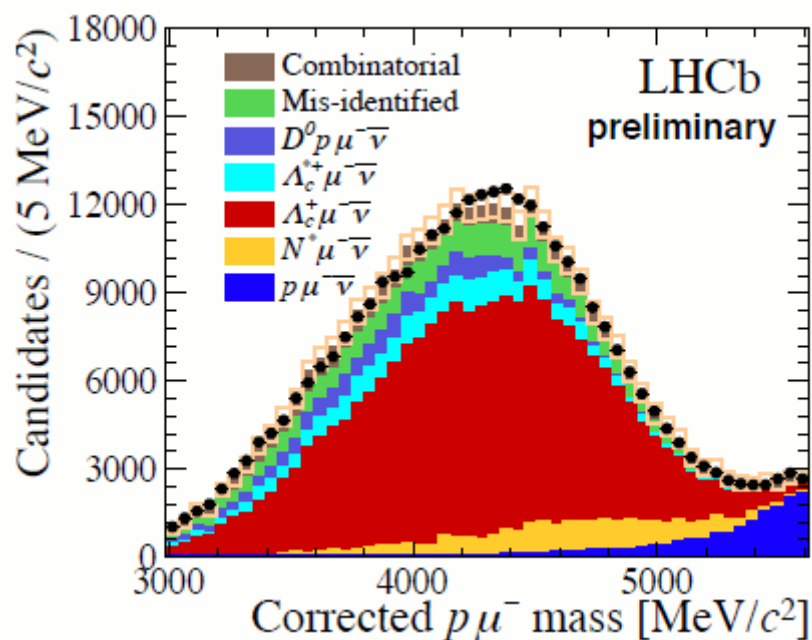
$$M_{corr} = \sqrt{p_{\perp}^2 + M_{p\mu}^2 + p_{\perp}}$$



$$N(\Lambda_b^0 \rightarrow p\mu\nu) = 17,687 \pm 733$$

First observation of this decay

$$N(\Lambda_b^0 \rightarrow \Lambda_c(pK)\mu\nu) = 34,255 \pm 571$$



$$\frac{B(\Lambda_b \rightarrow p\mu^-\bar{\nu}_\mu)_{q^2 > 15 \text{ GeV}^2/c^4}}{B(\Lambda_b \rightarrow \Lambda_c\mu\nu)_{q^2 > 7 \text{ GeV}^2/c^4}} = (1.00 \pm 0.04(stat) \pm 0.08(syst)) \times 10^{-2}$$

Determining $|V_{ub}|/|V_{cb}|$ arXiv:1504.01568

- Use ratio of differential rates from lattice calculations to calculate the ratio of CKM elements squared:

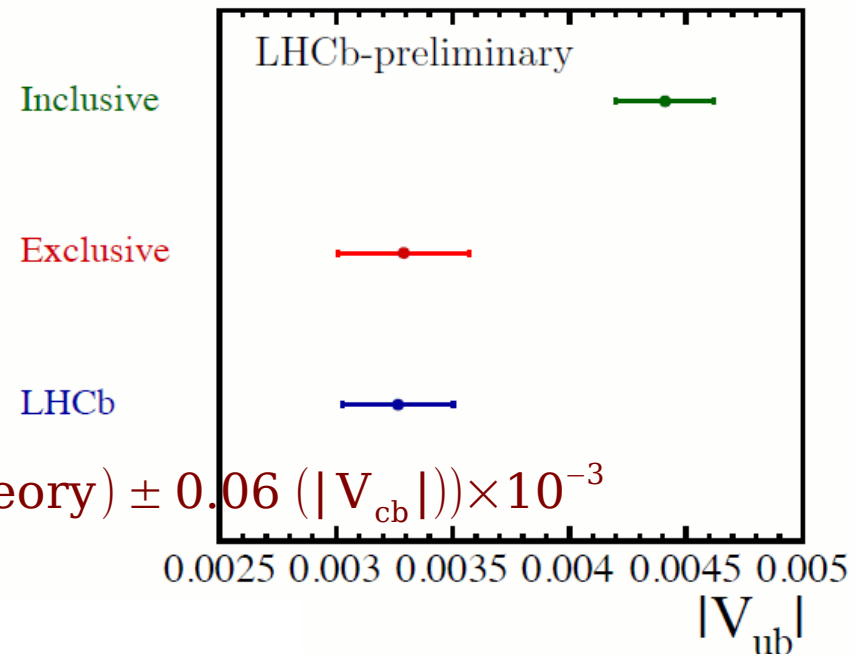
$$\frac{|V_{ub}|^2}{|V_{cb}|^2} = \frac{\int_{15 \text{ GeV}^2}^{q_{\text{max}}^2} \frac{d\Gamma(\Lambda_b \rightarrow p \mu^- \bar{\nu}_\mu)}{dq^2} dq^2}{\int_7 \text{ GeV}^2}^{q_{\text{max}}^2} \frac{d\Gamma(\Lambda_b \rightarrow \Lambda_c \mu^- \bar{\nu}_\mu)}{dq^2} dq^2} (0.68 \pm 0.07)$$

- leads to:

W.Detmold et al, arXiv:1503.01421

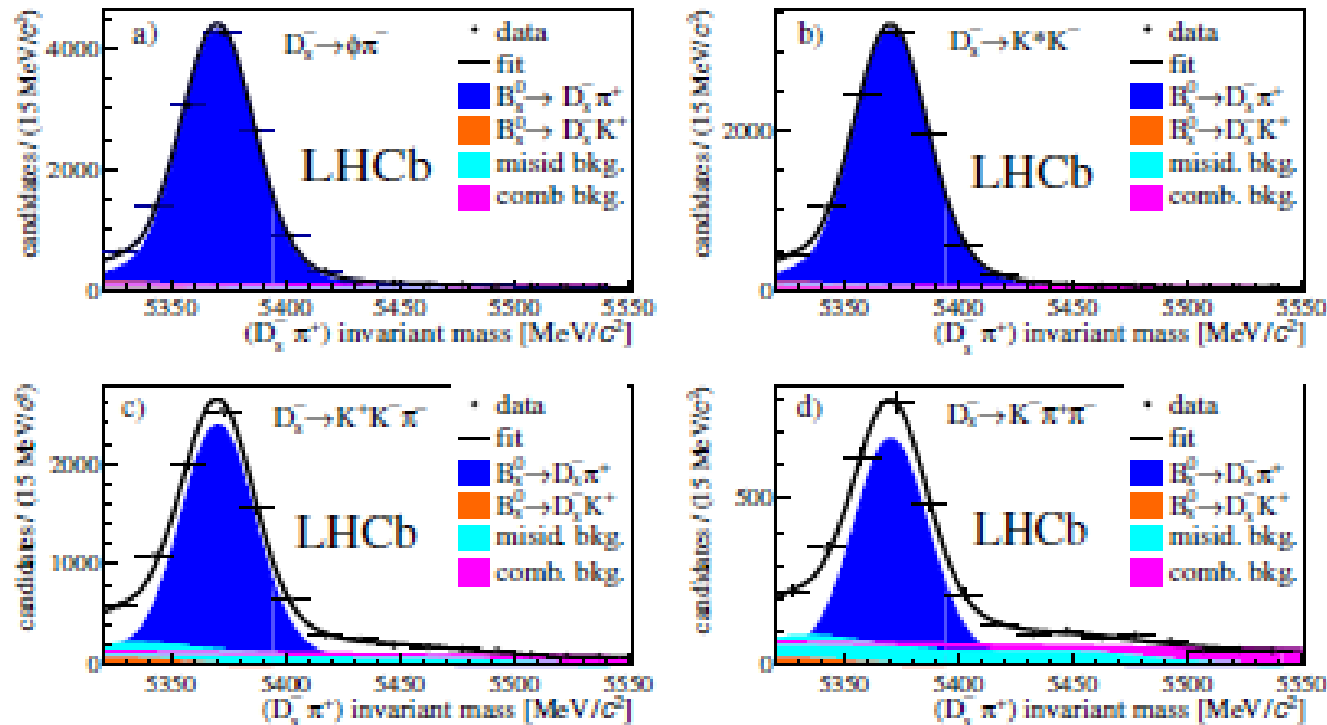
$$\frac{|V_{ub}|}{|V_{cb}|} = 0.083 \pm 0.004 (\text{exp}) \pm 0.004 (\text{LQCD})$$

$$|V_{cb}| = (3.27 \pm 0.15 (\text{exp}) \pm 0.17 (\text{theory}) \pm 0.06 (|V_{cb}|)) \times 10^{-3}$$

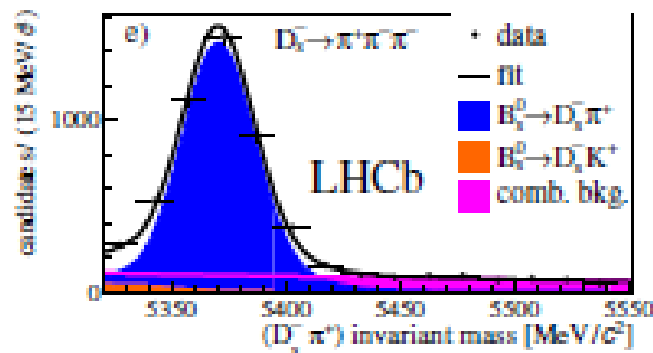


Measurement of Δm_s

- Observation for the first time in 2006 by CDF
- Precise measurement of $B_s^0 - \bar{B}_s^0$ oscillation frequency with $B_s^0 \rightarrow D_s^- \pi^+$ [arXiv:1304.4741]



34,000 $B_s \rightarrow D_s \pi$



Decay mode	$(D_s^- \pi^+)$ candidates
$D_s^- \rightarrow \phi(K^+ K^-) \pi^-$	14691
$D_s^- \rightarrow K^{*0}(K^+ \pi^-) K^-$	10866
$D_s^- \rightarrow K^+ K^- \pi^-$ nonresonant	11262
$D_s^- \rightarrow K^- \pi^+ \pi^-$	4288
$D_s^- \rightarrow \pi^- \pi^+ \pi^-$	6674
Total	47781

Measurement of Δm_s

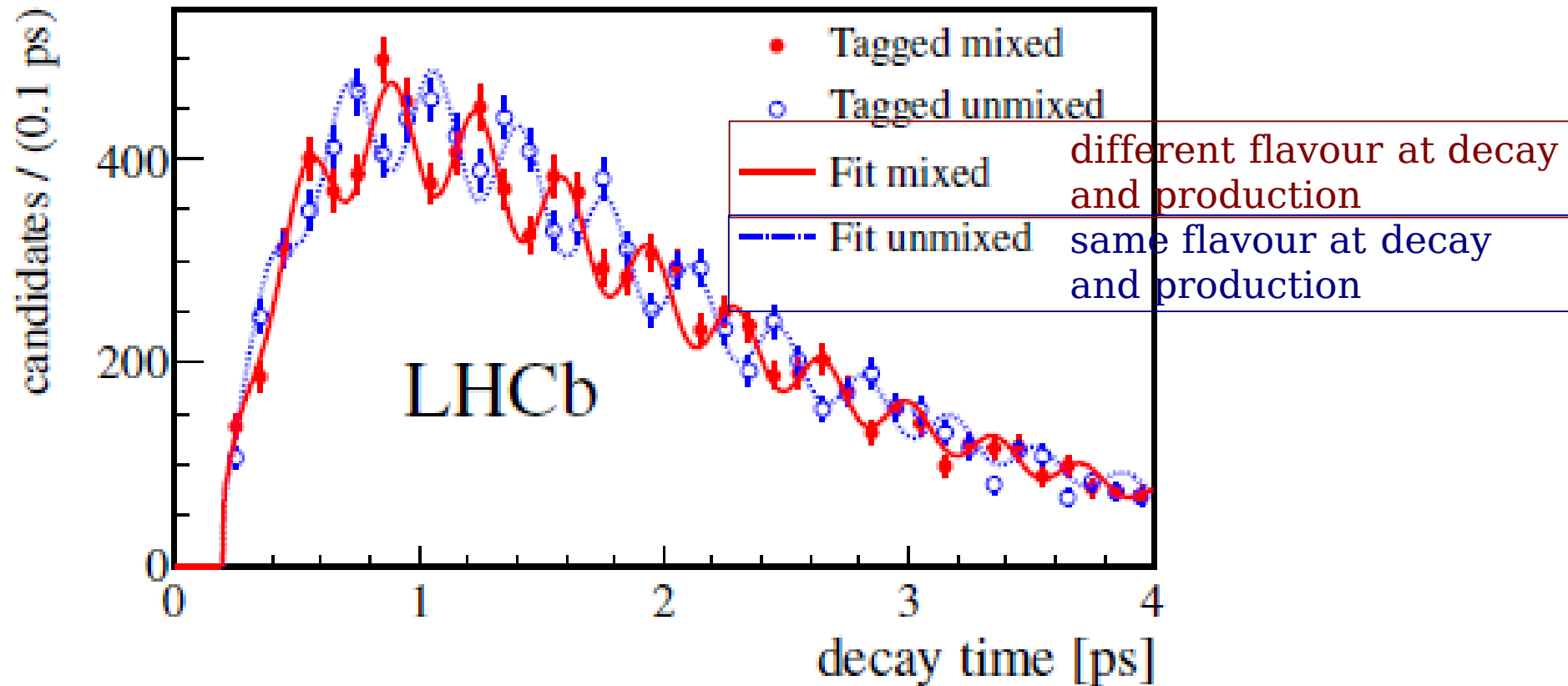
[arXiv:1304.4741]

- Precise measurement of B_s^0 - \bar{B}_s^0 oscillation frequency with $B_s^0 \rightarrow D_s^- \pi^+$

$$\mathcal{P}_t(t|\sigma_t) \propto \left\{ \Gamma_s e^{-\Gamma_s t} \frac{1}{2} \left[\cosh\left(\frac{\Delta\Gamma_s}{2}t\right) + q [1 - 2\omega(\eta_{\text{OST}}, \eta_{\text{SST}})] \cos(\Delta m_s t) \right] \theta(t) \right\} \\ \otimes G(t, S_{\sigma_t} \sigma_t) \mathcal{E}_t(t) \epsilon,$$

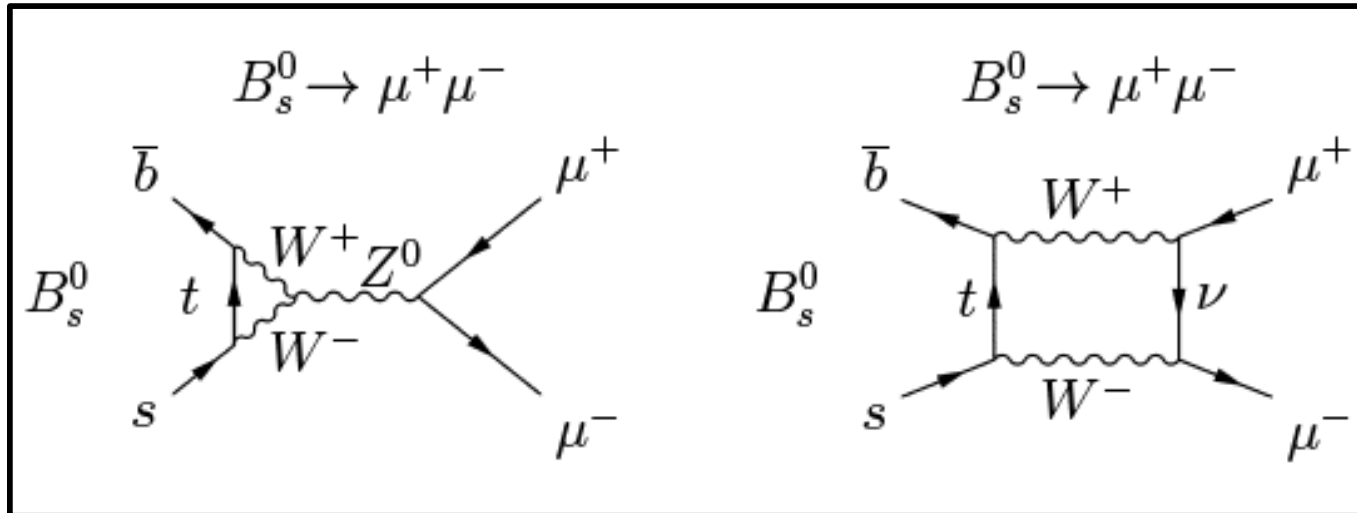
34,000 $B_s \rightarrow D_s \pi$

$$\Delta m_s = 17.768 \pm 0.023 \pm 0.006 \text{ ps}^{-1}$$



$B_{(s)} \rightarrow \mu\mu$: ultra rare processes...

loop diagram + suppressed in SM + theoretically clean =
an excellent place to look for new physics

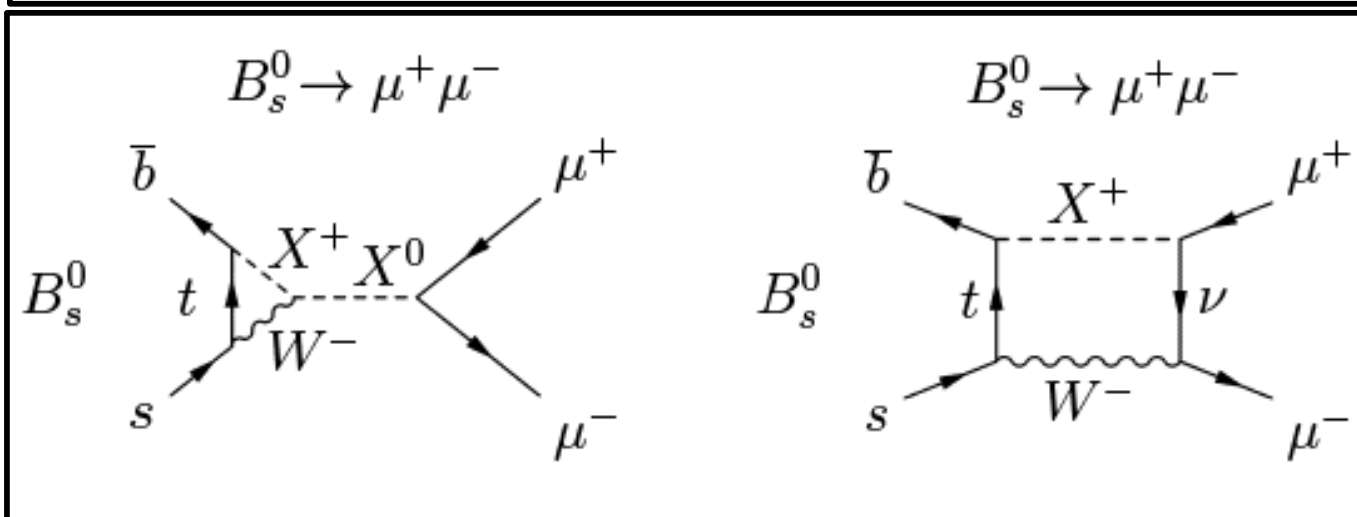


higher-order FCNC
 allowed in SM

$$B(B_s \rightarrow \mu^+ \mu^-) = (3.65 \pm 0.23) \times 10^{-9}$$

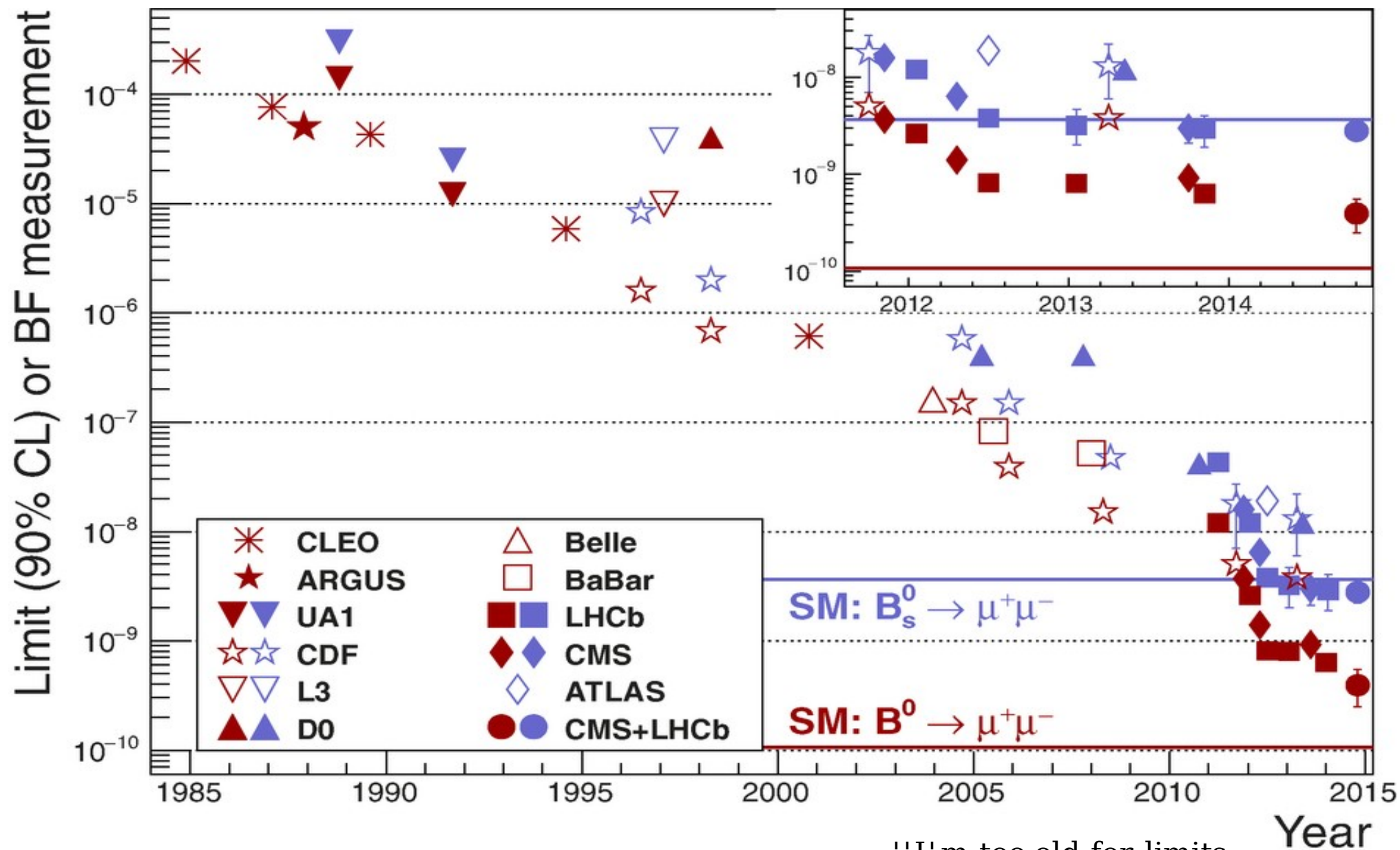
$$B(B_d \rightarrow \mu^+ \mu^-) = (1.06 \pm 0.09) \times 10^{-10}$$

[Bobeth et al,
 PRL 112 (2014) 101801]



same decay in theories
 extending the SM
 (some of NP scenarios
 may boost the $B \rightarrow \mu\mu$
 decay rates)

$B_{(s)} \rightarrow \mu\mu$: ultra rare processes...

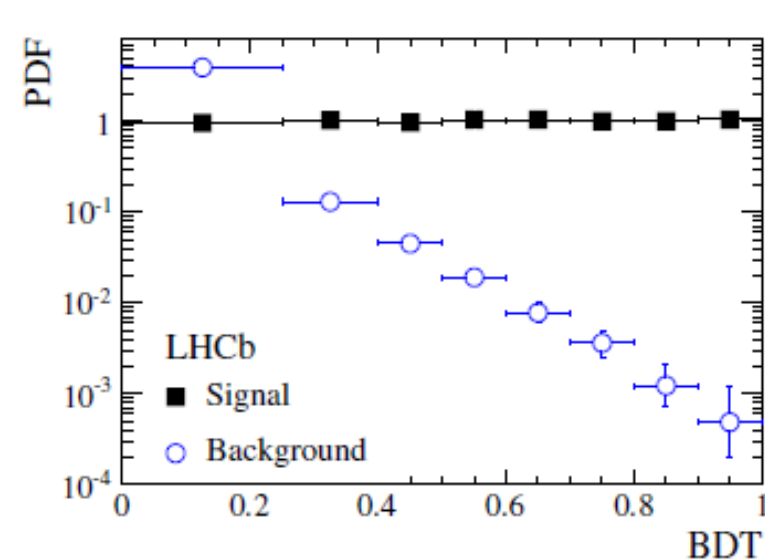


"I'm too old for limits,
I want to see signals"
(Francis Halzen)

$B_s \rightarrow \mu^+ \mu^-$ results

[arXiv:1307.5024]

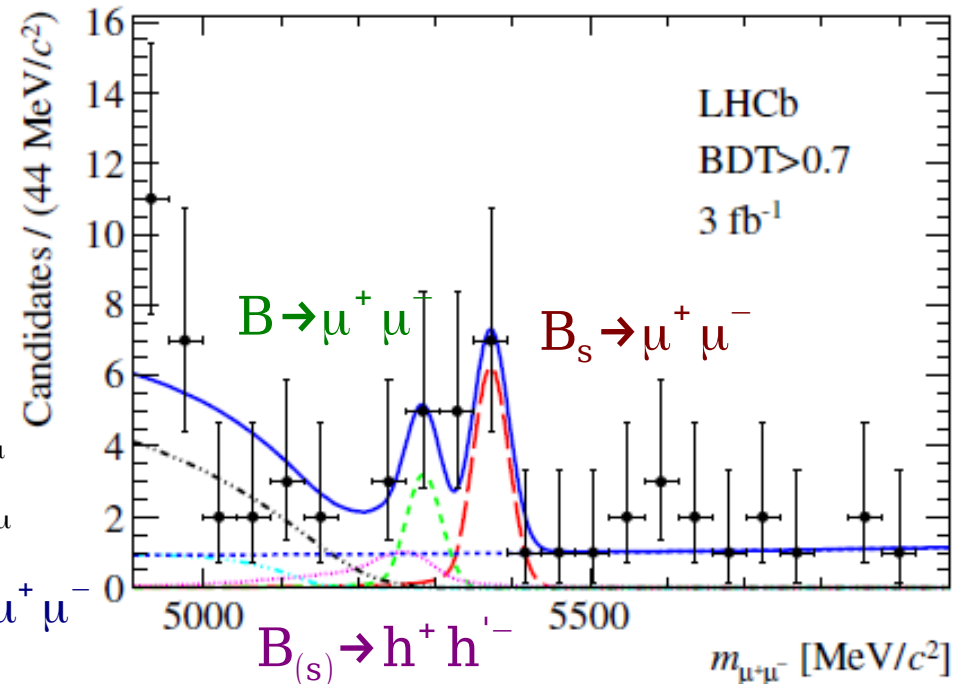
- after trigger and loose selection, $B_{(s)} \rightarrow \mu^+ \mu^-$ candidates classified according to $M(\mu\mu)$ and BDT output
 - separate $B_s \rightarrow \mu^+ \mu^-$ (signal) and $b\bar{b} \rightarrow \mu^+ \mu^-$ X events (background)
- BDT combining B candidate decay time, IP and p_T , minimum χ_{IP}^2 of the 2 muons, distance of closest approach btw 2 muons etc...



$$B^0 \rightarrow \pi^- \mu^+ \nu_\mu$$

$$B_s^0 \rightarrow K^- \mu^+ \nu_\mu$$

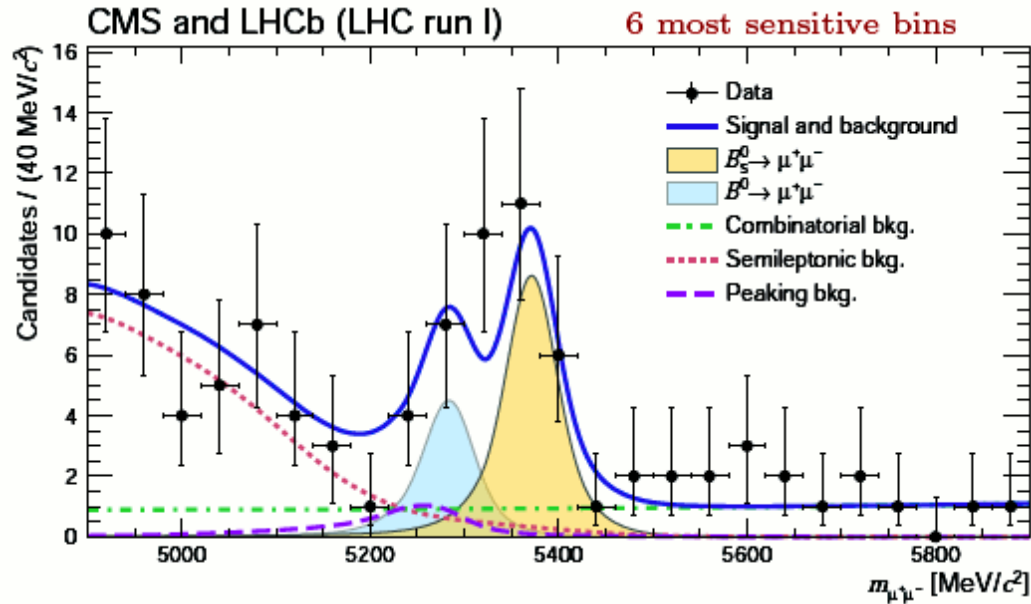
$$B^{0(+)} \rightarrow \pi^{0(+)} \mu^+ \mu^-$$



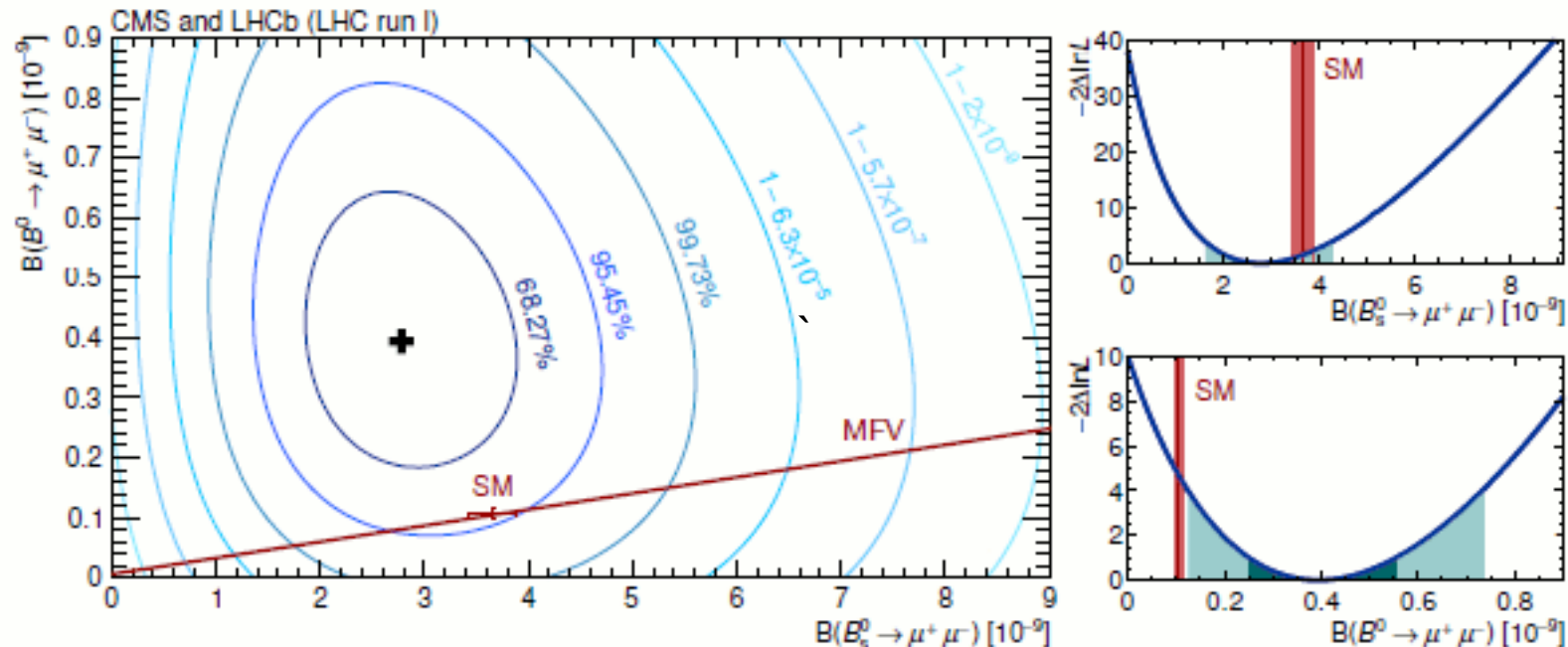
$$B(B_s^0 \rightarrow \mu^+ \mu^-) = (2.9^{+1.1}_{-1.0} {}^{+0.3}_{-0.1}) \times 10^{-9} \text{ (4.0}\sigma \text{ significance)}$$

$$B(B^0 \rightarrow \mu^+ \mu^-) = (3.7^{+2.4}_{-2.1} {}^{+0.6}_{-0.4}) \times 10^{-10} \text{ (2.0}\sigma \text{ significance)}$$

Combination results $B_{(s)} \rightarrow \mu^+ \mu^-$ [arXiv:1411.4413] published in Nature



$B(B_s^0 \rightarrow \mu^+ \mu^-) = (2.8^{+0.7}_{-0.6}) \times 10^{-9}$
first observation: 6.2 σ significance
 $B(B^0 \rightarrow \mu^+ \mu^-) = (3.9^{+1.6}_{-1.4}) \times 10^{-10}$
first evidence: 3.0 σ significance



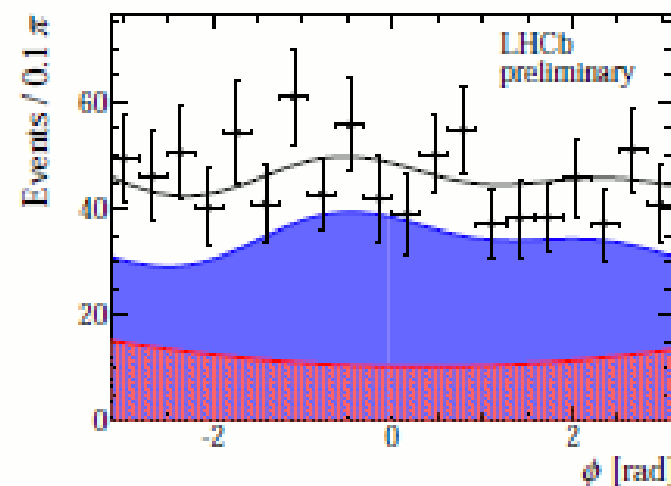
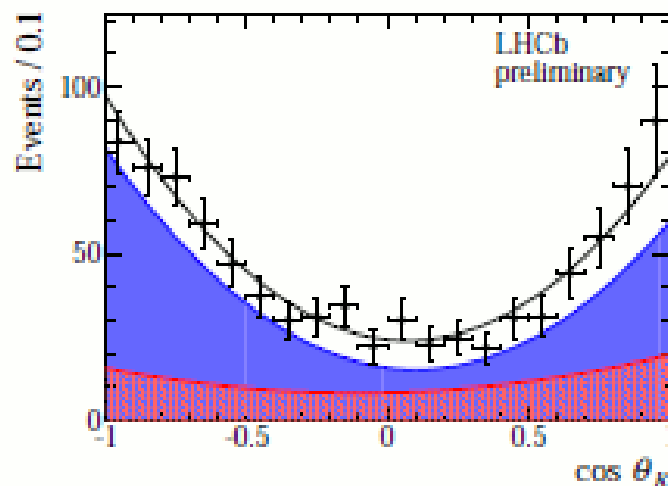
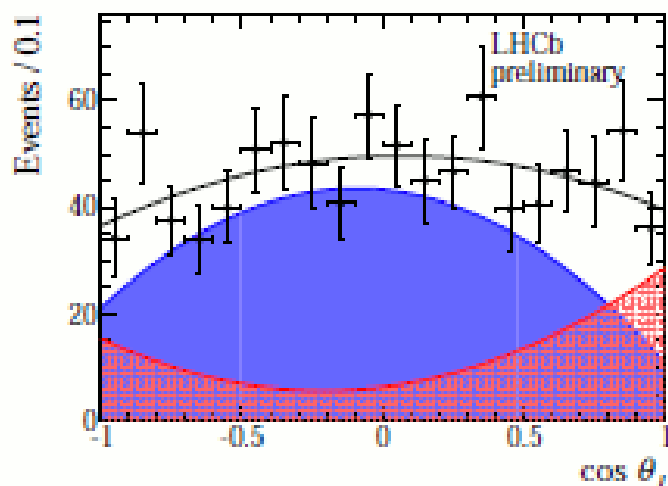
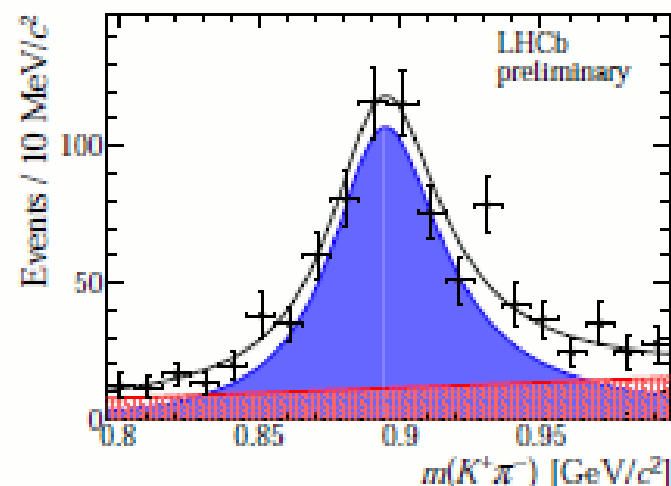
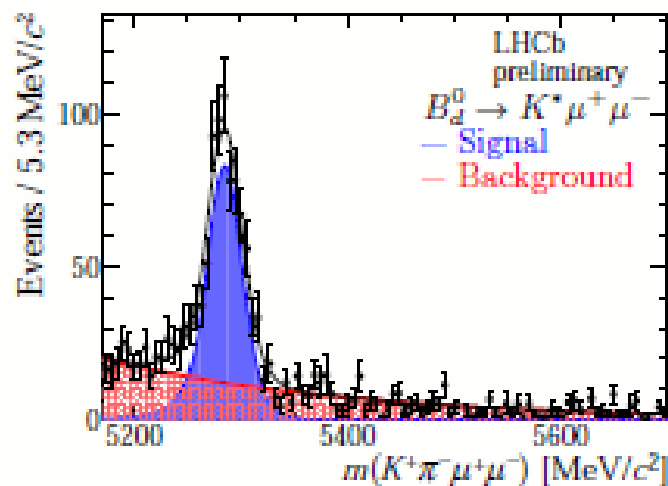
and more to come...

[Talk by Kai-Feng Chen for CMS, arXiv:1208.3355 for LHCb]

Angular analysis of $B_d^0 \rightarrow K^* \mu^+ \mu^-$ decays

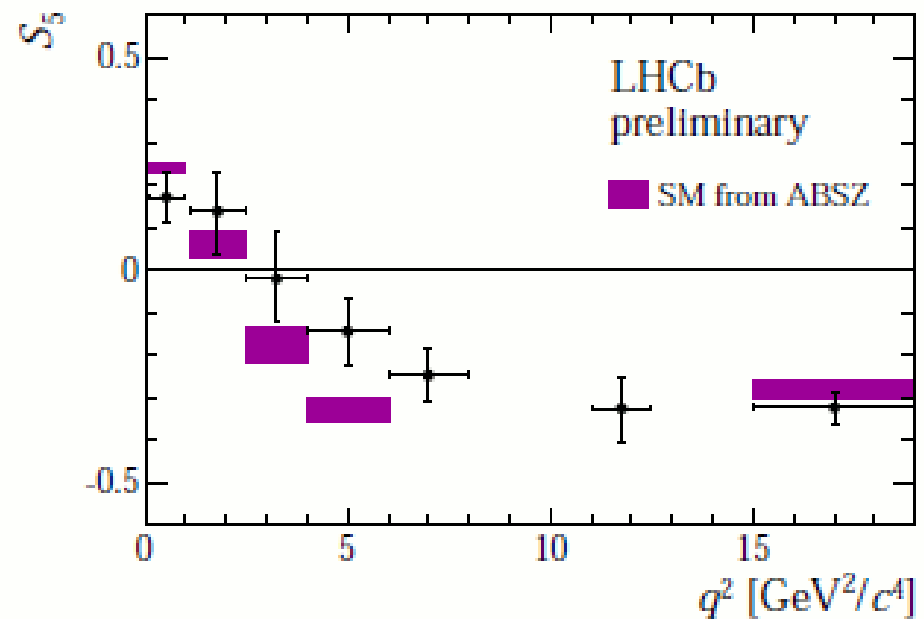
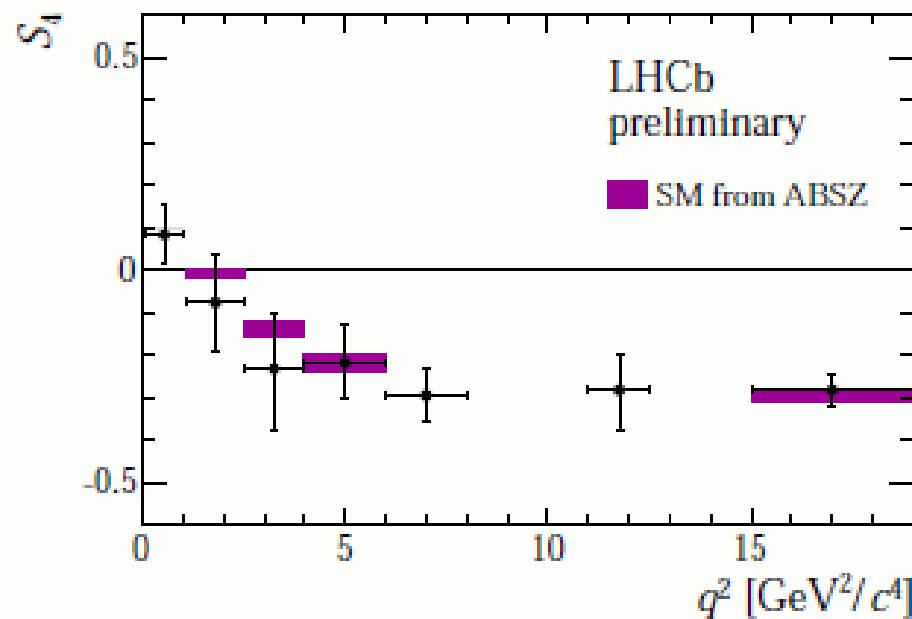
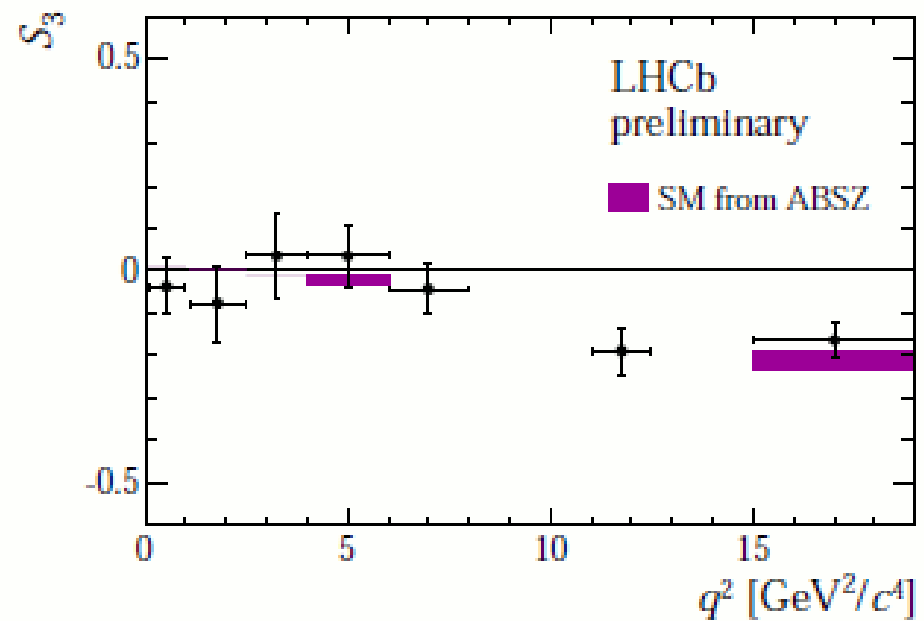
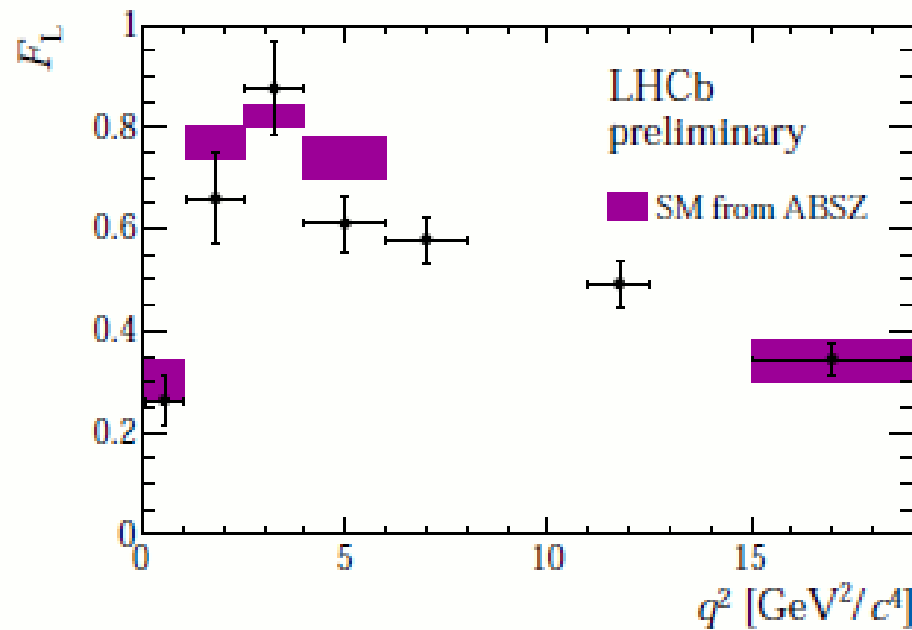
[arXiv:1512.04442]

- Projections of fit results for $q^2 \in [1.1, 6.0] \text{ GeV}^2$
- Good agreement of PDF projections with data in every bin of q^2



Angular analysis of $B_d^0 \rightarrow K^* \mu^+ \mu^-$ decays

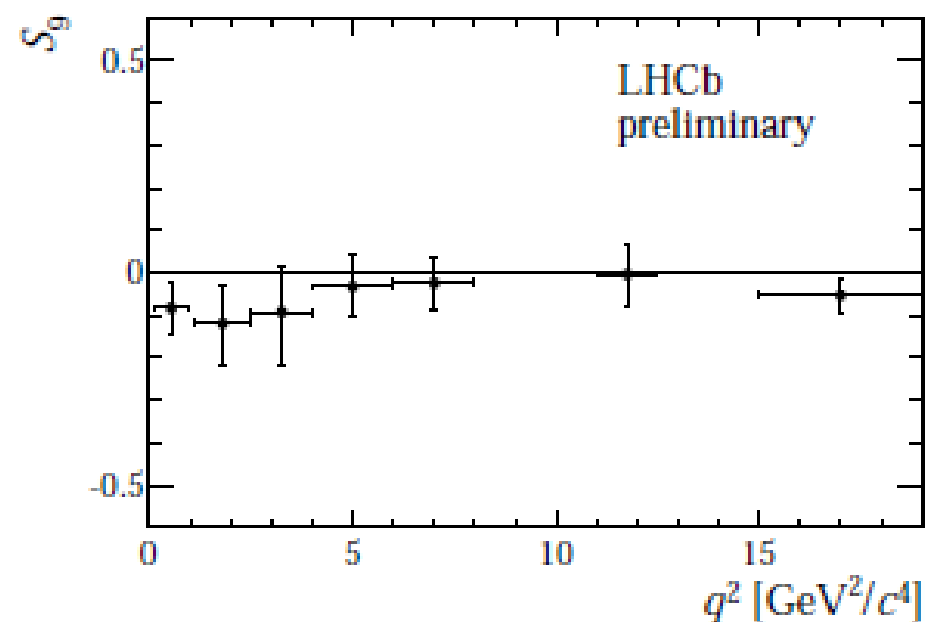
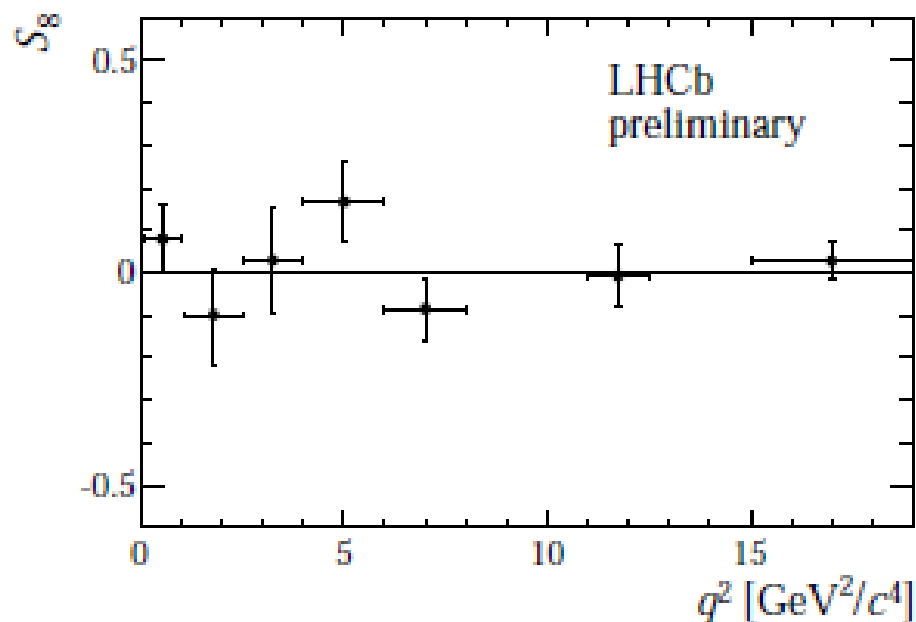
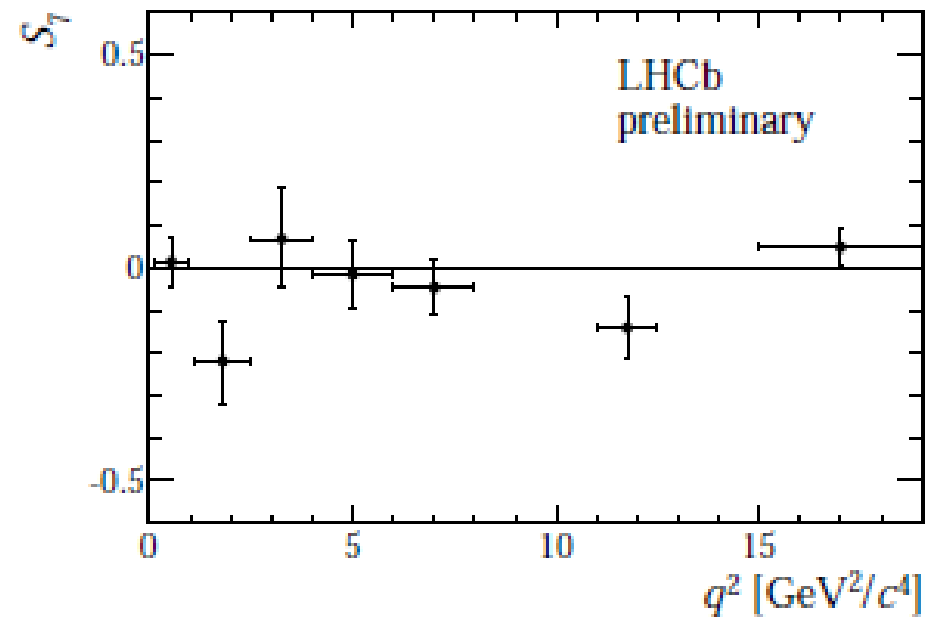
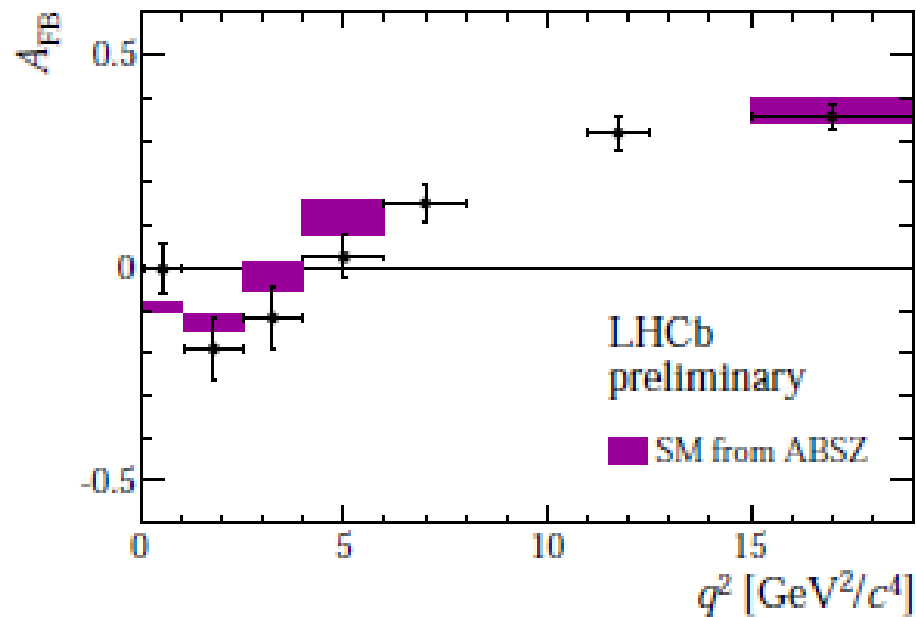
[arXiv:1512.04442]



Angular analysis of $B_d^0 \rightarrow K^* \mu^+ \mu^-$ decays

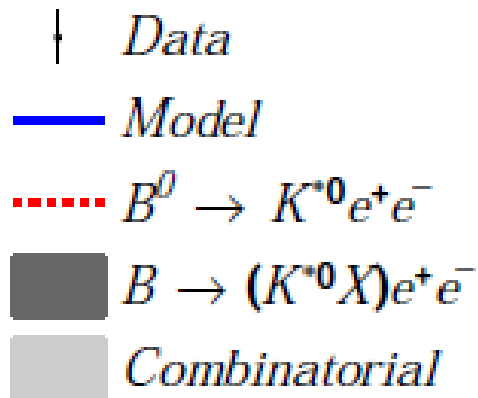
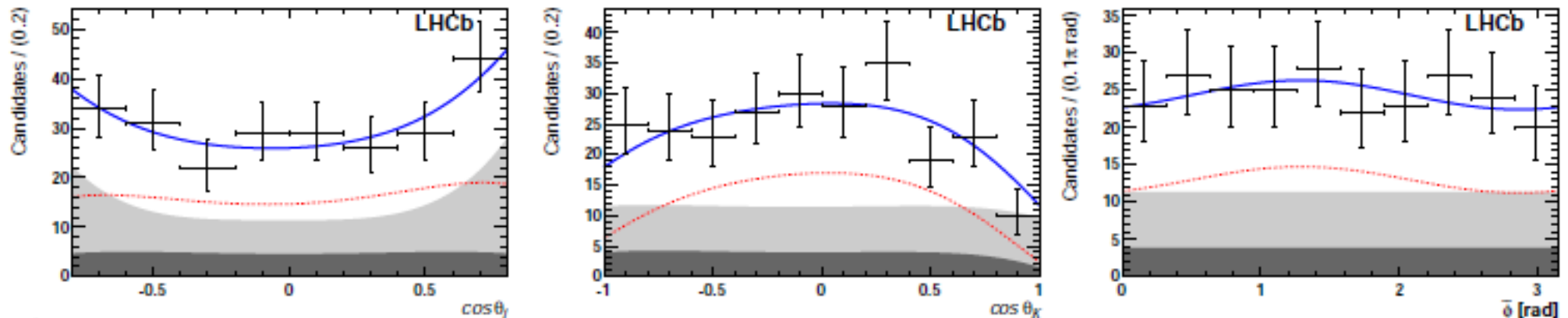
[arXiv:1512.04442]

data points systematically lower than SM



Angular analysis of $B_d^0 \rightarrow K^* e^+ e^-$ decays

[arXiv:1501.03038]



Observable	Measurement	SM prediction [†]
F_L	$+0.16 \pm 0.06 \pm 0.03$	$+0.10^{+0.11}_{-0.05}$
$A_T^{(2)}$	$-0.23 \pm 0.23 \pm 0.05$	$0.03^{+0.05}_{-0.04}$
A_T^{Re}	$+0.10 \pm 0.18 \pm 0.05$	$-0.15^{+0.04}_{-0.03}$
A_T^{Im}	$+0.14 \pm 0.22 \pm 0.05$	$(-0.2^{+1.2}_{-1.2}) \times 10^{-4}$

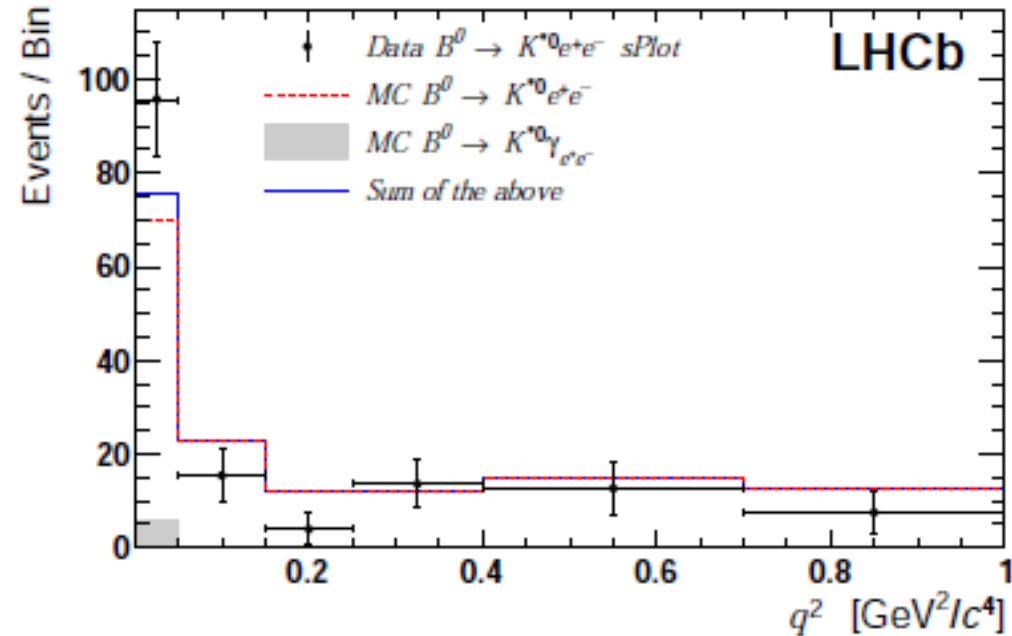
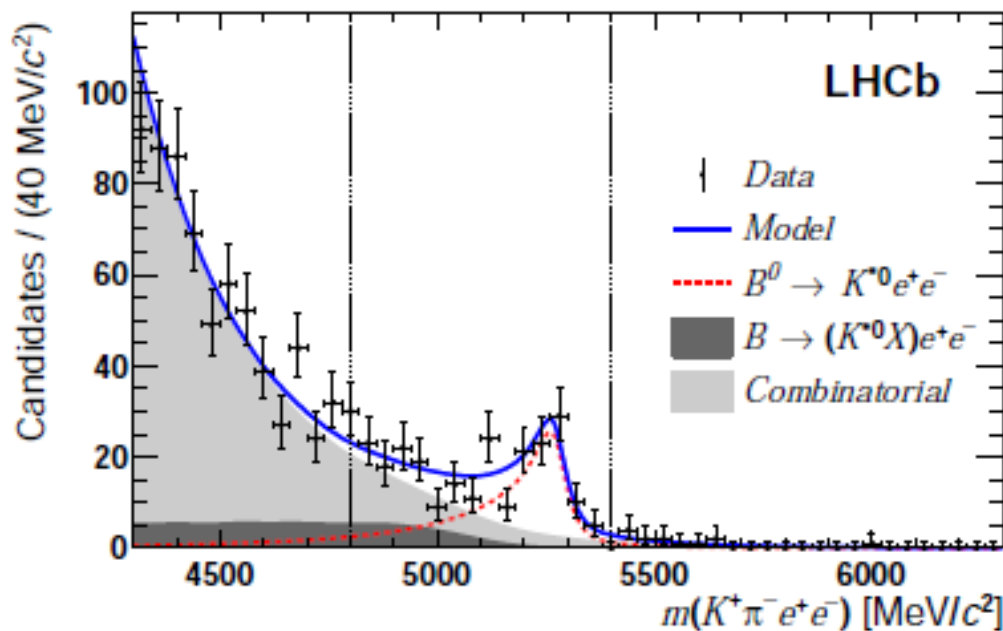
S.Jager, J.M.Camalich [arXiv:1412.3283]

- Measurements well in agreement with SM predictions
- Constraints on C_7' in complementary with radiative decays

Angular analysis of $B_d^0 \rightarrow K^* e^+ e^-$ decays

[arXiv:1501.03038]

- Angular analysis of $B_d^0 \rightarrow K^* e^+ e^-$ at very low q^2 ($\in [0.002, 1.120]$ GeV^2)
- Folded angular observables ($\varphi = \varphi + \pi$ if $\varphi < 0$)
- Measurement of F_L , $A_T^{(2)}$, $A_T^{(\text{Im})}$, $A_T^{(\text{Re})}$, sensitive to C_7' as $q^2 \rightarrow 0$



$$A_T^{(\text{Re})} = \frac{4}{3} A_{\text{FB}} / (1 - F_L), \quad A_T^{(2)} = \frac{1}{2} S_3 / (1 - F_L) \quad \text{and} \quad A_T = \frac{1}{2} S_9 / (1 - F_L)$$

The LHCb / LHCb upgrade timeline

LHCb future (2012 + end 2014 - 2017)

- $\mathcal{L} \geq 4 * 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$
- $L_{\text{int}} > 8 \text{ fb}^{-1}$ by the end of 2017
→ Factor-4 in statistical power wrt 1 fb^{-1}

Upgraded LHCb (2019 -)

- Full readout @ 40 MHz with full software trigger → trigger efficiency enhanced by a factor-2 for hadronic modes!
- Increase the luminosity by a factor-5
→ $\mathcal{L} \geq (1 - 2) * 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$
→ 25 ns bunch spacing → $\mu = 2$
→ $\sqrt{s} = 13\text{-}14 \text{ TeV}$
→ +100% $b\bar{b}$ x-section wrt $\sqrt{s} = 7 \text{ TeV}$
→ $\geq 5 \text{ fb}^{-1}/\text{year}$
- Run for 10 years
→ $L_{\text{int}} > 50 \text{ fb}^{-1}$
→ > Factor-10 in stat. power wrt 1 fb^{-1}

