Some anomalies in cosmic ray fluxes ?

LAPTh

Yoann Génolini

December 14th, 2015

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1 Modeling the high energy CRs fluxes

2 Propagation paradigm

- \Rightarrow State of the art
- \Rightarrow Positrons
- $\Rightarrow \mathsf{Antiprotons}$

3 Theoretical uncertainties on propagation

4 Conclusion

Outline

- 1 Modeling the high energy CRs fluxes
- 2 Propagation paradigm
- 3 Theoretical uncertainties on propagation
- 4 Conclusion

We don't measure directly the sources!



- Sources spectra :
 - $\rightarrow Q(E_k) \propto R^{-\alpha}$, with $R(E_k) = p/(Ze)$ and $\alpha \in [2.0, 2.5]$
- Transport (In the case of a weak electromagnetic turbulence) :

- ightarrow Diffusion in phase space $(oldsymbol{x},\,oldsymbol{p})\;D_x=D_0.eta.R^{\delta}$
- ightarrow Convective wind V_c .
- Interaction with the ISM :
 - \rightarrow Energy losses
 - \rightarrow Spallation $(\sigma_{\alpha}, \sigma_{\alpha \rightarrow \beta})$

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- ightarrow Convective wind V_c .
- Interaction with the ISM :
 - \rightarrow Energy losses
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Simplified propagation equation :

$$\frac{\partial f_{a}}{\partial t} + \frac{V_{c} \cdot \nabla_{x} f_{a} - \frac{1}{3} (\nabla_{x} \cdot V_{c}) p \frac{\partial f_{a}}{\partial p}}{\tau_{a}} + \frac{\nabla_{p} (b(p) f_{a})}{\tau_{a}} + \frac{\sigma_{a} v_{a} n_{ISM} f_{a}}{\sigma_{a} v_{a} n_{ISM} f_{a}} - \frac{f_{a}}{\tau_{a}} - \frac{\nabla_{x} \cdot (D_{x} \nabla_{x} f_{a})}{\tau_{a}} - \frac{\nabla_{p} \cdot (D_{p} \nabla_{p} f_{a})}{\tau_{b}} = q_{a} + \frac{\sum_{Z_{b} \geq Z_{a}}^{Z_{max}} \sigma_{b \to a} v_{b} n_{ISM} f_{b}}{\tau_{b}} + \frac{f_{b}}{\tau_{b}}}$$

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Simplifying the equation :

Ek > 10 GeV/nuc

$$\begin{aligned} \frac{\partial f_a}{\partial t} + & \sigma_a v_a n_{ISM} f_a + \\ \frac{f_a}{\tau_a} & - & \nabla_{\boldsymbol{x}} \cdot (D_x \nabla_{\boldsymbol{x}} f_a) = \\ q_a & + & \sum_{Z_b \geqslant Z_a}^{Z_{max}} \sigma_{b \to a} v_b n_{ISM} f_b + \frac{f_b}{\tau_b} \end{aligned}$$

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Galaxy model :



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Surface density of the disc $\mu = 2.4 \text{ mg.cm}^{-2}$.

Analytical resolution of the propagation equation :

For a stable nucleus :

$$\mathcal{J}_{a}(E_{k}) = \left\{ Q_{a} + \sum_{Z_{b} \geqslant Z_{a}}^{Z_{max}} \sigma_{b \rightarrow a} \mathcal{J}_{b} \right\} / \left\{ \sigma^{\text{diff}} + \sigma_{a} \right\}$$
(1)

Primary and secondary source terms.

Where :
$$\sigma^{\text{diff}} = \frac{2D \, m_{\text{ISM}}}{\mu v H}$$
.
and $Q_a = \frac{1}{4\pi} \frac{q_a}{n_{ISM}} \equiv N_a \left(\frac{\mathcal{R}}{1 \,\text{GV}}\right)^{\alpha}$.

High energy behaviour of fluxes :

Ek > 10 GeV/nuc





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Outline

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2 Propagation paradigm

- \Rightarrow State of the art
- \Rightarrow Positrons
- \Rightarrow Antiprotons

3 Theoretical uncertainties on propagation



Outline

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2 Propagation paradigm

- \Rightarrow State of the art
- \Rightarrow Positrons
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How do we constrain the propagation?

$$\begin{split} &\frac{\partial f_a}{\partial t} + \mathbf{V_c} \cdot \nabla_{\mathbf{x}} f_a - \frac{1}{3} (\nabla_{\mathbf{x}} \cdot \mathbf{V_c}) p \frac{\partial f_a}{\partial p} \nabla_{\mathbf{p}} (b(\mathbf{p}) f_a) + \sigma_a v_a n_{ISM} f_a + \\ &\frac{f_a}{\tau_a} - \nabla_{\mathbf{x}} \cdot (D_x \nabla_{\mathbf{x}} f_a) - \nabla_{\mathbf{p}} \cdot (D_p \nabla_{\mathbf{p}} f_a) = \\ &q_a + \sum_{Z_\beta \geqslant Z_a}^{Z_{max}} \sigma_{b \to a} v_b n_{ISM} f_b + \frac{f_b}{\tau_b} \end{split}$$

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How do we constrain the propagation?

$$\begin{split} \frac{\partial f_a}{\partial t} + \mathbf{V_c} \cdot \nabla_{\mathbf{x}} f_a &- \frac{1}{3} (\nabla_{\mathbf{x}} \cdot \mathbf{V_c}) p \frac{\partial f_a}{\partial p} \nabla_{\mathbf{p}} (\mathbf{b}(\mathbf{p}) f_a) + \sigma_a v_a n_{ISM} f_a + \\ \frac{f_a}{\tau_a} &- \nabla_{\mathbf{x}} \cdot (\mathbf{D}_{\mathbf{x}} \nabla_{\mathbf{x}} f_a) - \nabla_{\mathbf{p}} \cdot (\mathbf{D}_p \nabla_{\mathbf{p}} f_a) = \\ q_a &+ \sum_{Z_\beta \geqslant Z_a}^{Z_{max}} \sigma_{\mathbf{b} \rightarrow a} v_b n_{ISM} f_b + \frac{f_b}{\tau_b} \end{split}$$

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- Characteristics of the Galaxy
- Based on nuclear data

How do we constrain the propagation?

We are need to determine these parameters :

 $V_c, b(p), n_{ISM}, D_x, D_p, H$

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How do we constrain the propagation?

We are need to determine these parameters :

$$V_c, \quad D_x = D_0.\mathcal{R}^\delta, \quad D_p = \frac{V_A^2}{9D_x}p^2, \quad H$$

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How do we constrain the propagation?

We are need to determine these parameters :

 V_c, V_A, D_0, δ, H

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Measuring the propagation parameters :



Li, B, Be are said secondary.

Measuring the propagation parameters :



Li, B, Be are said secondary.

Secondary/primary ratio :

$$\mathcal{J}_B(E_k) = \left\{ Q_B + \sum_{Z_b \geqslant Z_B}^{Z_{max}} \sigma_{b \to B} \mathcal{J}_b \right\} / \left\{ \sigma^{\text{diff}} + \sigma_B \right\}$$
(2)

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Hypothesis :

• $Q_B = 0$

• Double nuclei system (B,C)

Secondary/primary ratio :

$$\mathcal{J}_{B}(E_{k}) = \sum_{Z_{b} \geqslant Z_{B}}^{Z_{max}} \sigma_{b \to B} \mathcal{J}_{b} / \left\{ \sigma^{\text{diff}} + \sigma_{B} \right\}$$
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Hypothesis :

- $Q_B = 0$
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Secondary/primary ratio :

$$\mathcal{J}_B(E_k) = \sigma_{C \to B} \mathcal{J}_C / \left\{ \sigma^{\text{diff}} + \sigma_B \right\}$$
(2)

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Hypothesis :

- $Q_B = 0$
- Double nuclei system (B,C)

Secondary/primary ratio :

$$\frac{\mathcal{J}_B(E_k)}{\mathcal{J}_C(E_k)} = \sigma_{C \to B} / \left\{ \sigma^{\text{diff}} + \sigma_B \right\}.$$

When :
$$\sigma_B << \sigma^{diff} \Rightarrow \boxed{\frac{\mathcal{J}_B}{\mathcal{J}_C} \propto R^{-\delta}}$$

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Experimental data :



Experimental data :



New data with AMS02 !.. and soon CALET !

Analysis currently used :

[Maurin 2001] \Rightarrow 1623(3 σ) sets are found to be consistent with B/C.



Other analysis :

[Putze 2010] \Rightarrow More comprhensive study(Including ${}^{10}Be/{}^{9}Be$).



Outline

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2 Propagation paradigm

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Typical fits of the positron fraction : $\Phi_{e+} = \Phi_{e+}^{secondary} + \Phi_{e+}^{primary}$

Pulsar explanation?



Fit of $\{fW_0, \gamma\}$

Fit of $\{ < \sigma v >, m_{\chi} \}$

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Dark matter

explanation?

Systematic vs statistical uncertainties :



⇒ The systematics on propagation dominate completely their determination.
[Boudaud 2014] : Carefull analysis of the positron fraction.

Boudaud 2014 $\rightarrow arxiv[1410.3799]$

Astronomy & Astrophysics manuscript no. CRAC_paper_1	©ESO 2015
January 22, 2015	

A new look at the cosmic ray positron fraction

M. Boudaud¹, S. Aupetit¹, S. Caroff², A. Putze^{1,2}, G. Belanger¹, Y. Genolini¹, C. Goy², V. Poireau², V. Poulin¹, S. Rosier², P. Salati¹, L. Tao², and M. Vecchi^{2, 3, *, *, *}

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Received; accepted Preprint numbers : LAPTH-224/14

ABSTRACT

Context. The positron fraction in cosmic rays has recently been measured with improved accuracy up to 500 CeV, and it was found to be a steadily increasing function of energy above ~ 10 GeV. This behaviour contrasts with standard astrophysical mechanisms, in which positrons are secondary particles, produced in the interactions of primary cosmic rays during their propagation in the interstellar

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Outline

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Image: Image:

2 Propagation paradigm

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Revaluation of the astrophysical background :



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Tension with the fiducial model

Revaluation of the astrophysical background :



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Tension with the fiducial model!

Revaluation of the astrophysical background.

Equation in steady state :

$$\begin{split} \partial_z (V_C \psi) &- D_x \Delta \psi + \partial_E \{ b^{loss}(E) \psi - D_{EE}(E) \partial_E \psi \} = Q \\ \text{With} : Q(\psi_p, \psi_{He}, \sigma_{pH \to \bar{p}}(E), ..) \end{split}$$

- Propagation \rightarrow [Maurin 2001]
- Primary fluxes \rightarrow [AMS02 2015]
- Production cross-section \rightarrow [di Mauro 2014]

Revaluation of the astrophysical background :



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\rightarrow Published in [Guisen 2015]

Revaluation of the astrophysical background :



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Guisen 2015 $\rightarrow arxiv[1504.0427]$

AMS-02 antiprotons, at last! Secondary astrophysical component and immediate implications for Dark Matter

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^c Institute for Theoretical Particle Physics and Cosmology (TTK), RWTH Aachen University, D-52056 Aachen, Germany.

Abstract

Using the updated proton and helium fluxes just released by the AMS-

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Outline

- 1 Modeling the high energy CRs fluxes
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Benchmark model :

The goal is to minimize :

$$\chi^2_{\rm B/C} = \sum_{i} \left\{ \frac{\mathcal{F}_i^{\rm exp} - \mathcal{F}_i^{\rm th}(\underline{\textit{Parameters..}})}{\sigma_i} \right\}^2$$



Benchmark model :

The goal is to minimize :

$$\chi^2_{\rm B/C} = \sum_{i} \left\{ \frac{\mathcal{F}_i^{\rm exp} - \mathcal{F}_i^{\rm th}(\boldsymbol{\delta}, \boldsymbol{D}_0)}{\sigma_i} \right\}^2$$



$$\mathcal{F}^{\rm th} = \frac{\mathcal{J}_{\sf B}(E_k)}{\mathcal{J}_{\sf C}(E_k)} = \frac{Q_{\sf B}}{\sigma^{\sf diff} + \sigma_B} / \mathcal{J}_{\sf C} + \sum_{Z_k \ge Z_{\sf C}}^{Z_{\sf max}} \frac{\sigma_{b \to \sf B}}{\sigma^{\sf diff} + \sigma_{\sf B}} \frac{\mathcal{J}_b}{\mathcal{J}_{\sf C}}$$

- Primary boron contribution
- Production cross-section uncertainties
- Destruction cross-section uncertainties
- Geometry framework



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• Primary boron contribution

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[Bla09], [BS09], [MS09], [MS14] \rightarrow Secondary species may be formed at sources !

• Confinement inside a SNR at TeV/nuc :

$$X_{SNR} \approx 0.17 \text{ g cm}^{-2} \frac{n_{ISM}}{\text{cm}^{-3}} \frac{T_{SNR}}{2.10^4 \text{ yr}}$$

• Galactic diffusion at TeV/nuc :

$$X_{Diff} \approx 1.2 \text{ g cm}^{-2}$$

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 \Rightarrow Order of 10%!

At high energy...

$$\frac{\mathcal{J}_B(E_k)}{\mathcal{J}_C(E_k)} = \left\{ \frac{Q_B}{\mathcal{J}_C} + \sigma_{C \to B} + \sum_{Z_b > Z_C}^{Z_{max}} \sigma_{b \to B} \frac{\mathcal{J}_b}{\mathcal{J}_C} \right\} / \left\{ \sigma^{\text{diff}} + \sigma_B \right\}$$
$$\overset{\alpha}{HE} \quad \frac{N_B}{N_C}$$

Image: Image:

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... it leads to a plateau.

Constraining $\frac{N_B}{N_C}$:



Scan on $\frac{N_B}{N_C}$:



ightarrow Huge impact on the determination of delta

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Scan on $\frac{N_B}{N_C}$:



 \rightarrow Huge impact on the determination of delta !

Summary of the main systematics :

	Wind	1D/2D geometry	Cross-sections	Primary boron
$\Delta D_0/D_0$	-40%	-2 to $-13%$	$\pm 60\%$	0 to $-90%$
$\Delta \delta / \delta$	+15%	$0 \ {\rm to} \ {+1\%}$	$\pm 20\%$	$0 \ \mathrm{to} \ +100\%$

Prospects, what we need at zero order :

Find a way to quantify primary boron contribution.
New precise measurements of nuclear cross-sections.

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- **2** New precise measurements of nuclear cross-sections.

Génolini 2015 $\rightarrow arxiv[1504.03134]$

Astronomy & Astrophysics manuscript no. draft4 July 22, 2015 ©ESO 2015

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Theoretical uncertainties in extracting cosmic-ray diffusion parameters: the boron-to-carbon ratio

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Received; accepted Preprint numbers : LAPTH-018/15

ABSTRACT

Context. PAMELA and, more recently, AMS-02, are ushering us into a new era of greatly reduced statistical uncertainties in experimental measurements of cosmic-ray fluxes. In particular, new determinations of traditional diagnostic tools such as the boron-tocarbon ratio (B/C) are expected to significantly reduce errors on cosmic-ray diffusion parameters, with important implications for astroparticle physics, ranging from inferring primary source spectra to indirect dark matter searches.

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Future projects :

- Discretness of sources : With Pasquale Serpico, Pierre Salati and Richard Taillet.
- Revaluation of secondary positon propagation uncertainties : With Antje Putze for an updated analysis.
- An updated analysis of the B/C ratio..



Thanks for listening

- Pasquale Blasi, The origin of the positron excess in cosmic rays, Phys.Rev.Lett. 103 (2009), 051104.
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- P. Mertsch and S. Sarkar, AMS-02 data confront acceleration of cosmic ray secondaries in nearby sources, Prd 90 (2014), no. 6, 061301.