



Moriond EW 2016

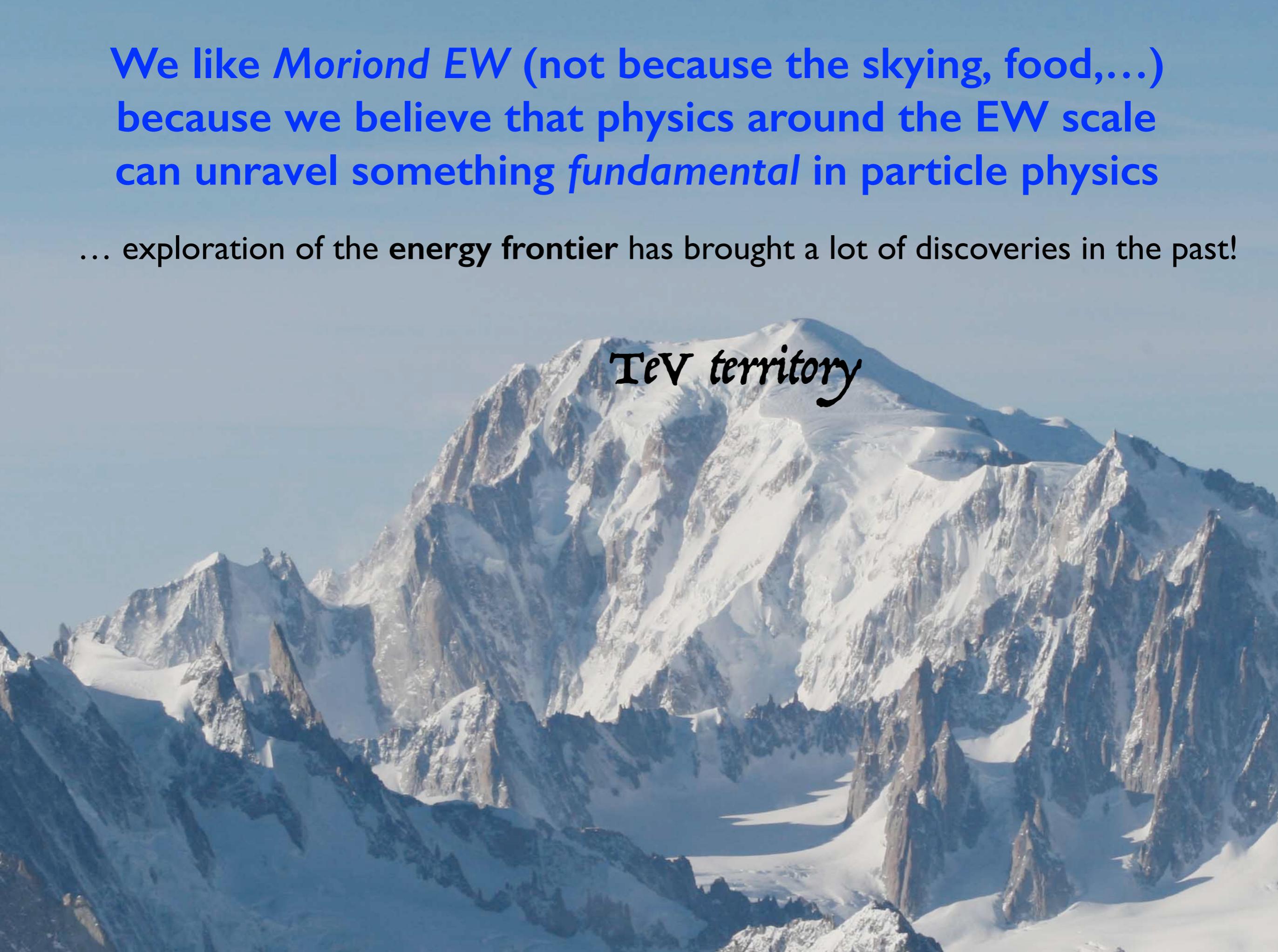
Theory Summary

Alex Pomarol CERN & UAB

We like *Moriond EW* (not because the skying, food,...)
because we believe that physics around the EW scale
can unravel something *fundamental* in particle physics

... exploration of the energy frontier has brought a lot of discoveries in the past!

TeV territory

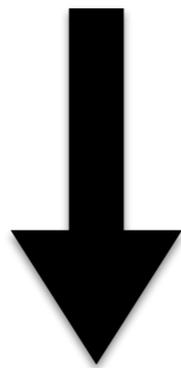


But the situation has changed from the old days
(at least for theorists)...

As *Rencontres Moriond*, little by little
particle physics has amazingly evolved from 1966 to 2016

Pre-Higgs era: We were building up the theory
discovery guaranteed:

..., top, Higgs

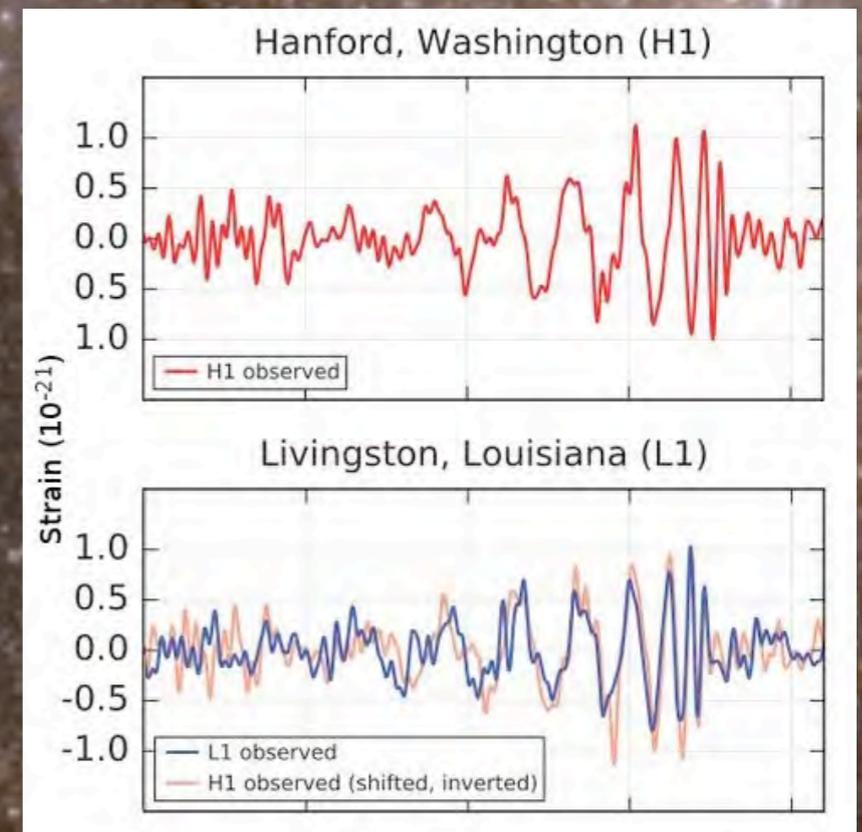


Post-Higgs era: We now have the theory:

A (quantum field) theory **SM+GR**
that can consistently give us the physics up to $\sim M_P$

Powerful theory that can predict the gravitational waves from a binary BH merger

see Alessandro Nagar's talk



So, why should be new phenomena at the TeV?

Not because the present theory is inconsistent!

We have the theory,
but now we'd like to understand why it is like it is

In particular, we want to understand the origin of the EW scale
(hierarchy problem):

Why $m_W \ll M_P \sim 10^{19} \text{ GeV}$?

↳ origin of the Higgs potential

↳ why the Higgs is so light?

This is today the main motivation to explore the TeV frontier

↳ but success not guaranteed... see later

The TeV frontier must be attacked from several fronts



TeV territory



The TeV frontier must be attacked from several fronts

Role of theorist: Provide the necessary tools & routes to TeV physics



Role of Theorists

Providers (SM)



Better and better
predictions of
the SM physics

Intending Visionaries (BSM)



Alternative
ways to
probe
physics
beyond the
SM

Role of Theorists

Providers (SM)



Better and better
predictions of
the SM physics

8 speakers



Gudrun Hiller
Lars Hofer
Nicolas Garron
Anna Hayes
Jonathan Engel
Andre Hoang
Giulia Zanderighi
Alessandro Nagar

Intending Visionaries (BSM)



Alternative
ways to
probe
physics
beyond the
SM

20 speakers



Nejc Kosnik
Michele Lucente
Boris Kayser
Cedric Delaunay
Aneesh Manohar
Joachim Kopp
Tony Gherghetta
Martin Jung
Felix Brummer
Thomas Rizzo
Aurora Meroni
Elizabeth Jenkins
Aldo DeAndrea
JoAnne Hewett
A. Strumia
Farinaldo Queiroz
Maxim Pospelov
Suzanne Westhoff
J.R. Espinosa
Francesco Sannino

Role of Theorists I



Calculating the SM predictions:

We know the theory (SM+GR),
but we do not know its predictions!

Long ongoing project on how to deal with QCD

From quarks & gluons to hadrons (physical objects):

$$\langle \mathbf{K} | (\bar{q} \gamma^\mu q)^2 | \pi \pi \rangle$$



Non-Perturbative calculations: Lattice

Crucial as now most flavor observables are close to the SM value



In Lattice we trust

we find

$$\text{Re}(\varepsilon' / \varepsilon) = 1.38(5.15)(4.43) \times 10^{-4}$$

whereas the experimental value is

$$\text{Re}(\varepsilon' / \varepsilon) = 16.6(2.3) \times 10^{-4}$$

$\sim 2.1\sigma$

With this unphysical computation (kinematics, masses) we find

$$\begin{aligned} \Delta I = 1/2 \text{ rule } \frac{\text{Re}A_0}{\text{Re}A_2} &= 9.1(2.1) \text{ for } m_K = 878 \text{ MeV } m_\pi = 422 \text{ MeV} \\ &= 12.0(1.7) \text{ for } m_K = 662 \text{ MeV } m_\pi = 329 \text{ MeV} \end{aligned}$$

experimentally: 22

\curvearrowright in the right direction but not yet there...

Alternative methods:

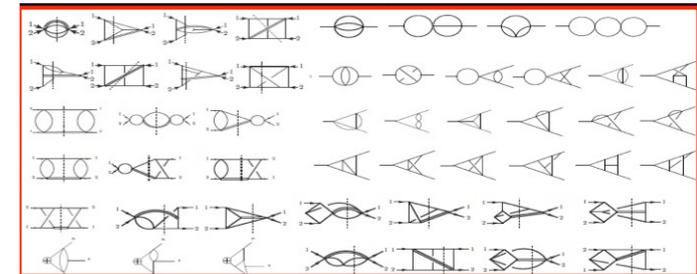
$B \rightarrow K^* \mu \mu$ at high q^2 :

– Using a local model against the OPE provides a data-driven method to test the binning and limitations of the OPE.

QCD at hadron machines

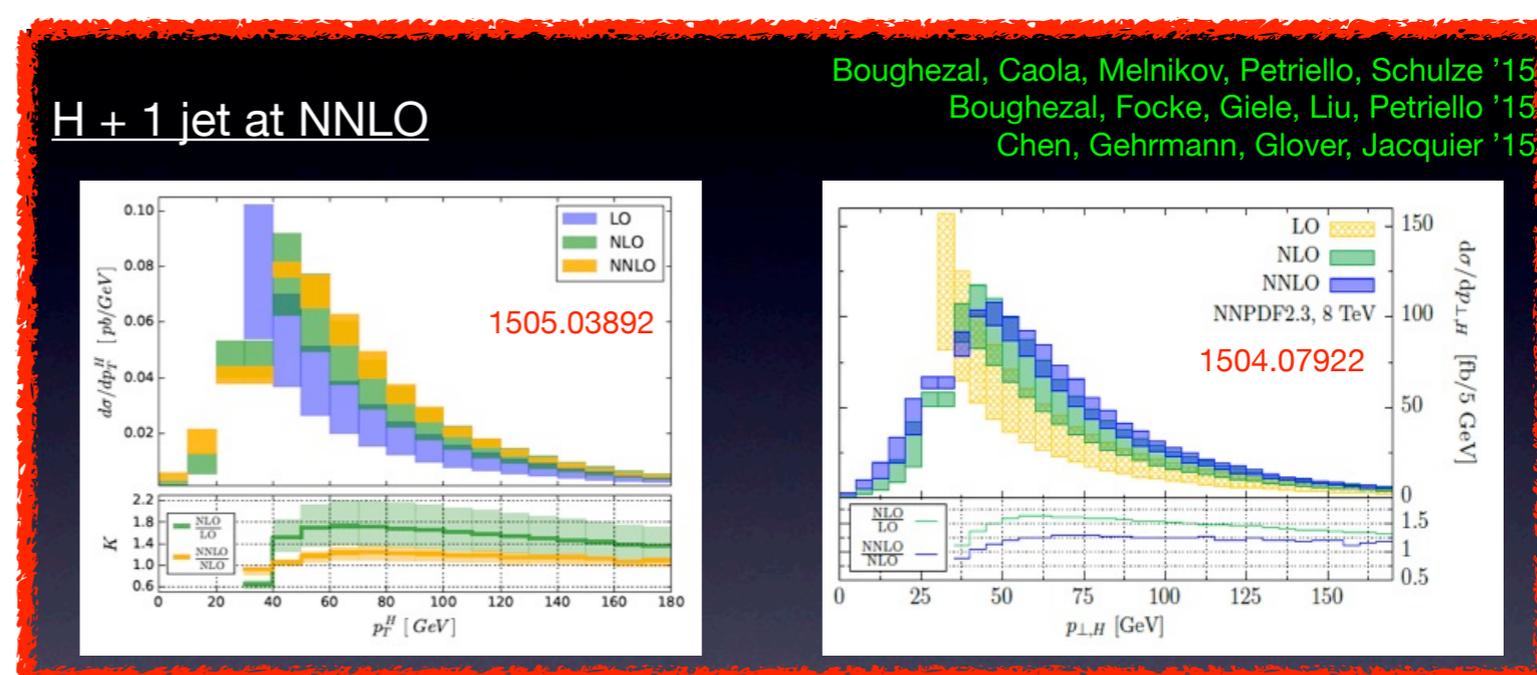
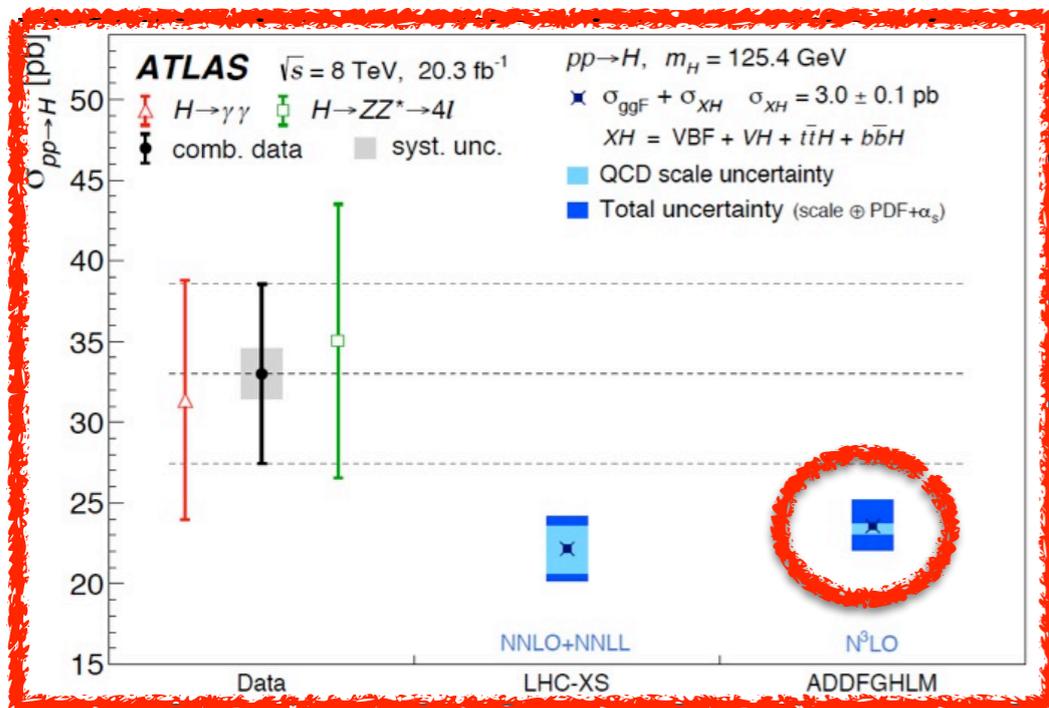
Perturbative calculations: From LO $\rightarrow \dots \rightarrow N^3$ LO

Giulia Zanderighi:



Accept it, they are the 8th wonder of the world!

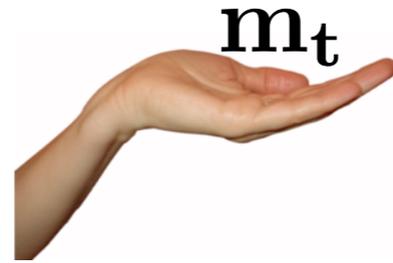
Gave glimpses of these beauties:



from inclusive σ_H at N^3 LO...

... to H+jet diff. cross-section, crucial to extract more on the Higgs's nature

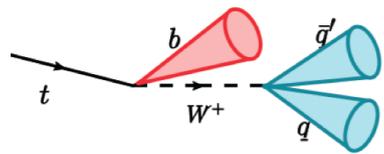
Extraction of the top mass:



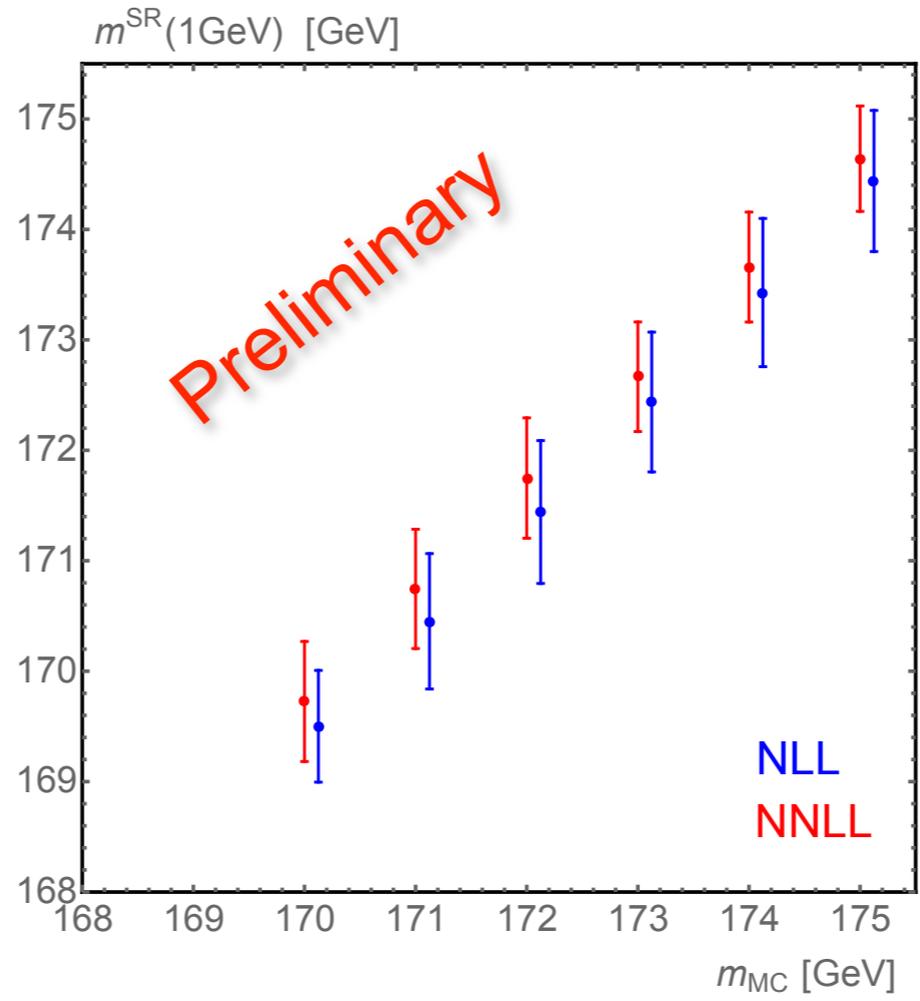
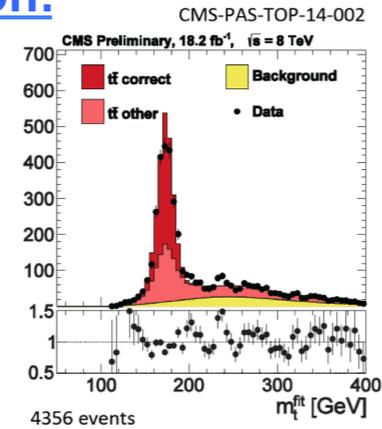
André H. Hoang's proposal:

LHC+Tevatron

Direct Reconstruction:

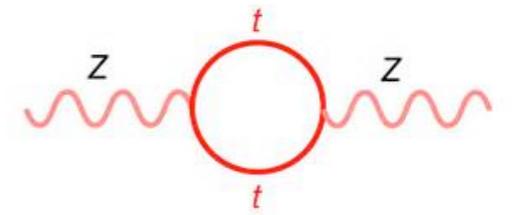


kinematic mass determination



error $O(500 \text{ MeV})$

useful for precision tests of the SM!



Role of Theorists II

Routes to BSM



The TeV frontier must be attacked from several fronts

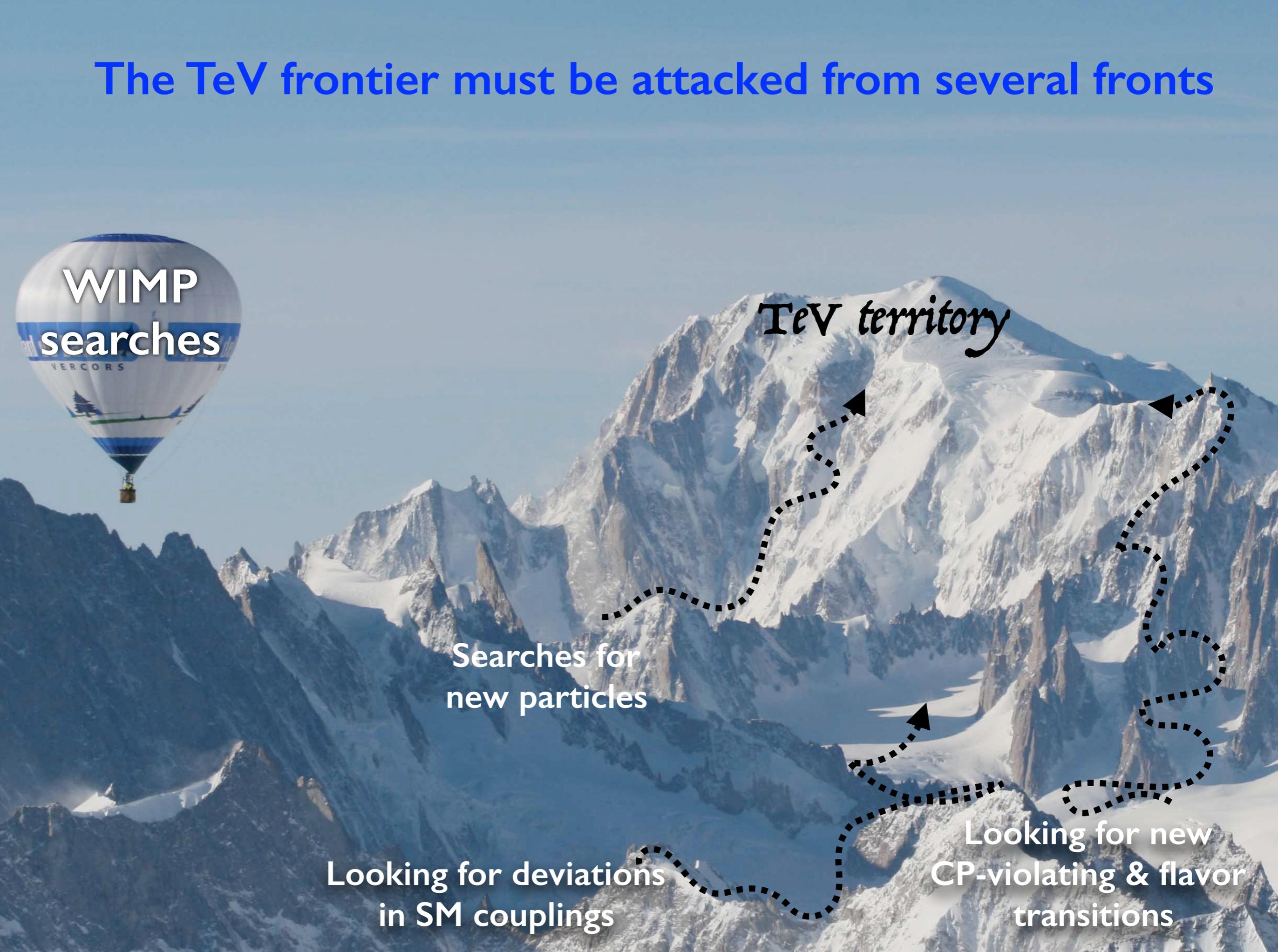


TeV territory

Searches for new particles

Looking for deviations in SM couplings

Looking for new CP-violating & flavor transitions



The TeV frontier must be attacked from several fronts

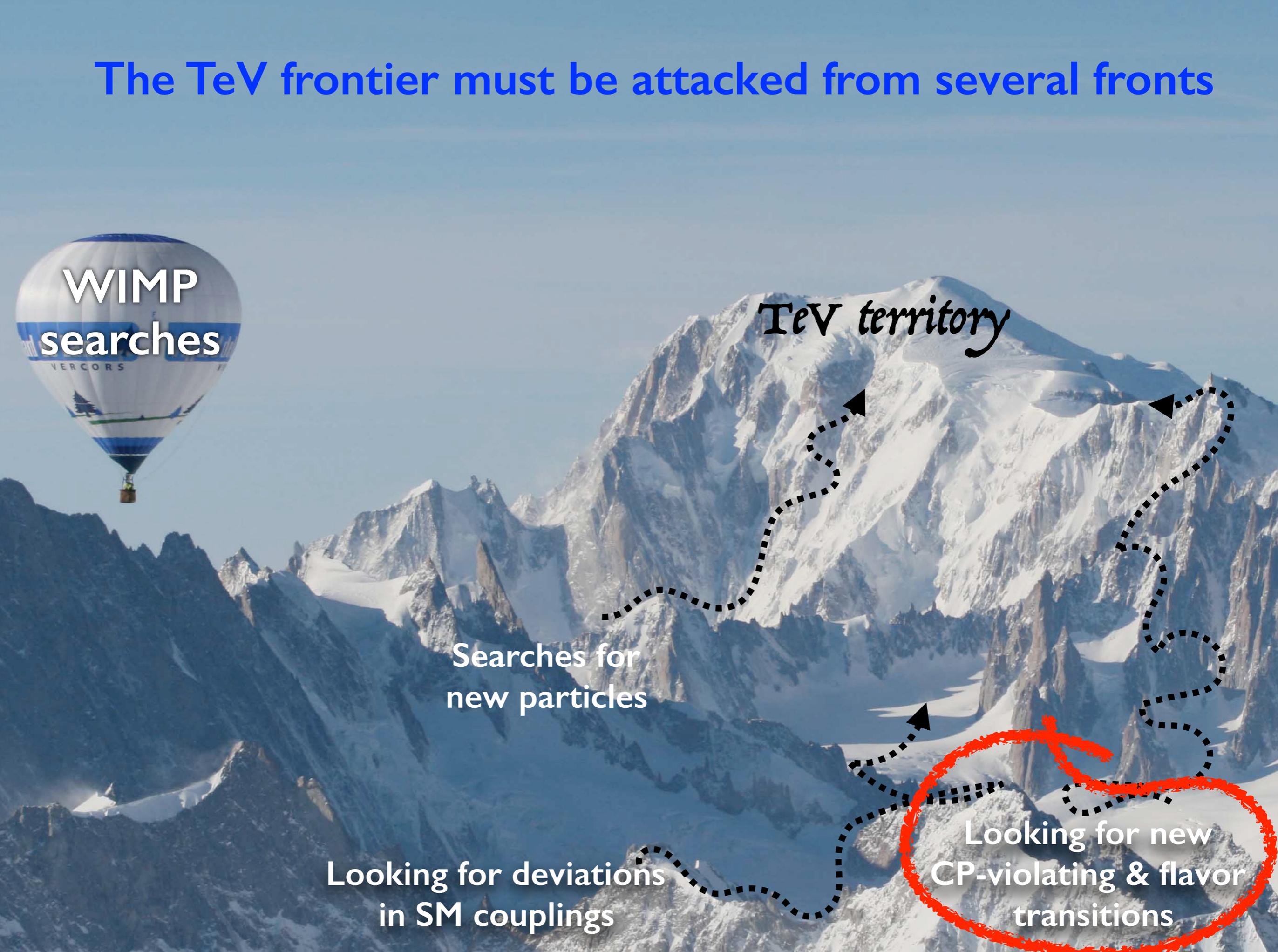


TeV territory

Searches for new particles

Looking for deviations in SM couplings

Looking for new CP-violating & flavor transitions



Motivation: **Not the baryon-antibaryon asymmetry of the universe!**

Provide observables that receive small SM contributions due to the “accidental” symmetries of the SM:

Lepton number, flavor-symmetries up to small Yukawas, CP-conservation unless 3 family concurrence

We do not expect that BSM will share these “accidents” as these are theories with more structure, so we expect them to give large effects to these observables

e.g. supersymmetry, composite Higgs

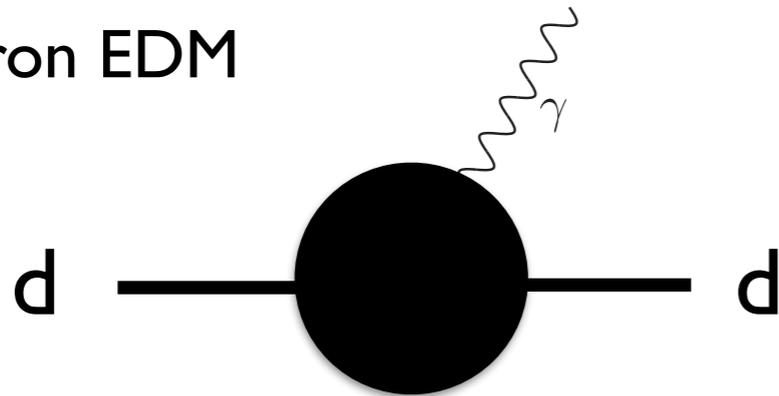


Clean paths to BSM...

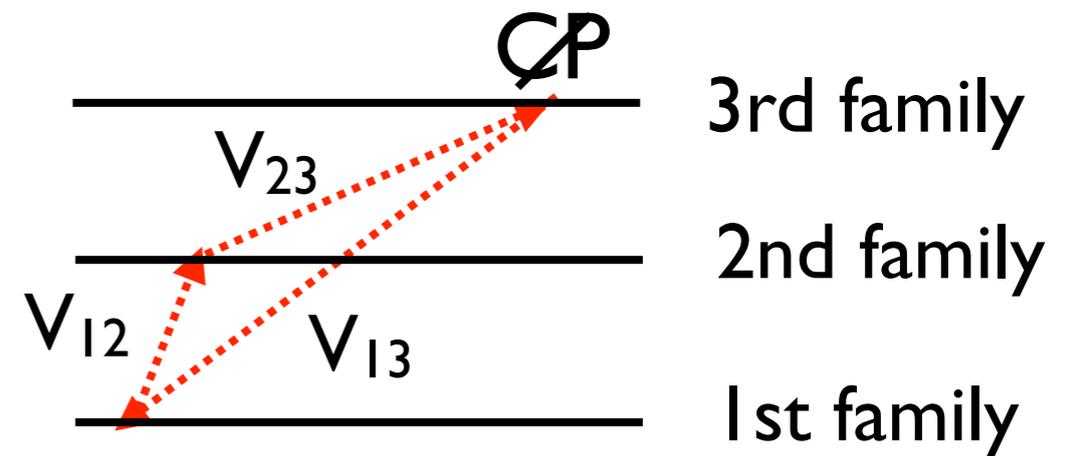
Best example:

Electric Dipole Moments (EDM)

e.g. neutron EDM



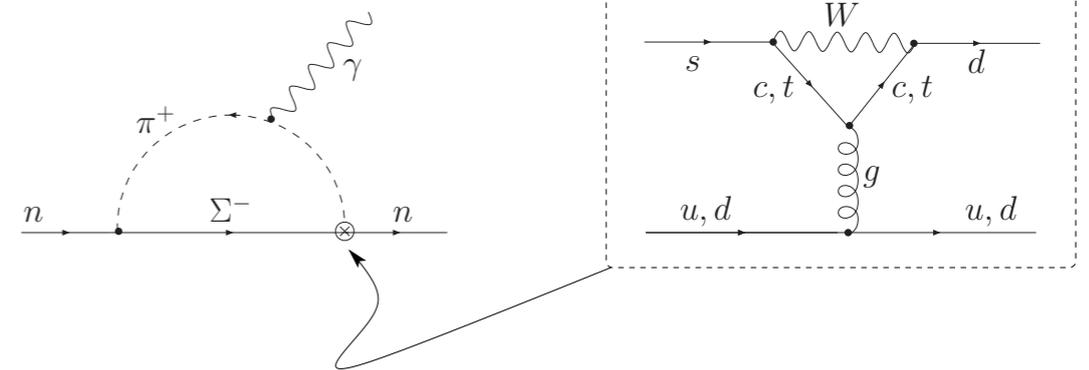
“Selection rules”



In the SM must involve the 3 family quarks

d_n extremely small $\sim 10^{-32} e \cdot \text{cm}$

experimentally: $d_n < 2.9 \times 10^{-26} e \cdot \text{cm}$



assuming $\theta=0$
from an axion!

CP-violating phases in BSM are ubiquitous \rightarrow larger contributions

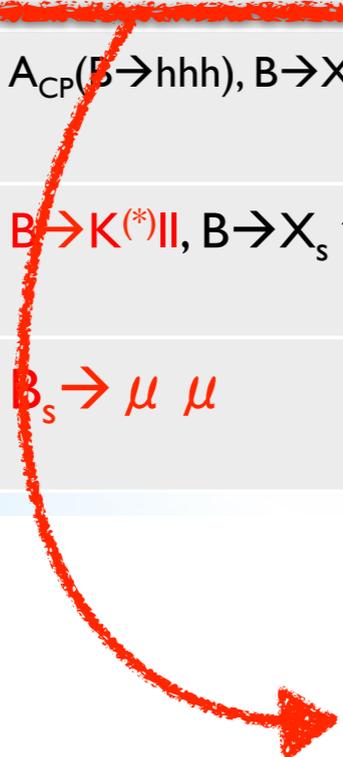
Relatively “cheap” experiment, but as competitive as the LHC!!

Flavor Physics

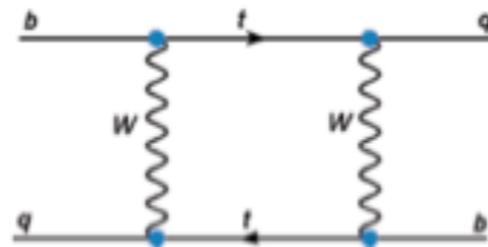
	$b \rightarrow s$ ($ V_{tb} V_{ts} \propto \lambda^2$)	$b \rightarrow d$ ($ V_{tb} V_{td} \propto \lambda^3$)	$s \rightarrow d$ ($ V_{ts} V_{td} \propto \lambda^5$)	$c \rightarrow u$ ($ V_{cb} V_{ub} \propto \lambda^5$)
$\Delta F=2$ box	$\Delta M_{B_s}, A_{CP}(B_s \rightarrow J/\psi \Phi)$	$\Delta M_B, A_{CP}(B \rightarrow J/\psi K)$	$\Delta M_K, \epsilon_K$	$x, y, q/p, \Phi$
QCD Penguin	$A_{CP}(B \rightarrow hhh), B \rightarrow X_s \gamma$	$A_{CP}(B \rightarrow hhh), B \rightarrow X \gamma$	$K \rightarrow \pi^0 \Pi, \epsilon' / \epsilon$	$\Delta a_{CP}(D \rightarrow hh)$
EW Penguin	$B \rightarrow K^{(*)} \Pi, B \rightarrow X_s \gamma$	$B \rightarrow \pi \Pi, B \rightarrow X \gamma$	$K \rightarrow \pi^0 \Pi, K^\pm \rightarrow \pi^\pm \nu \nu$	$D \rightarrow X_u \Pi$
Higgs Penguin	$B_s \rightarrow \mu \mu$	$B \rightarrow \mu \mu$	$K \rightarrow \mu \mu$	$D \rightarrow \mu \mu$

Flavor Physics

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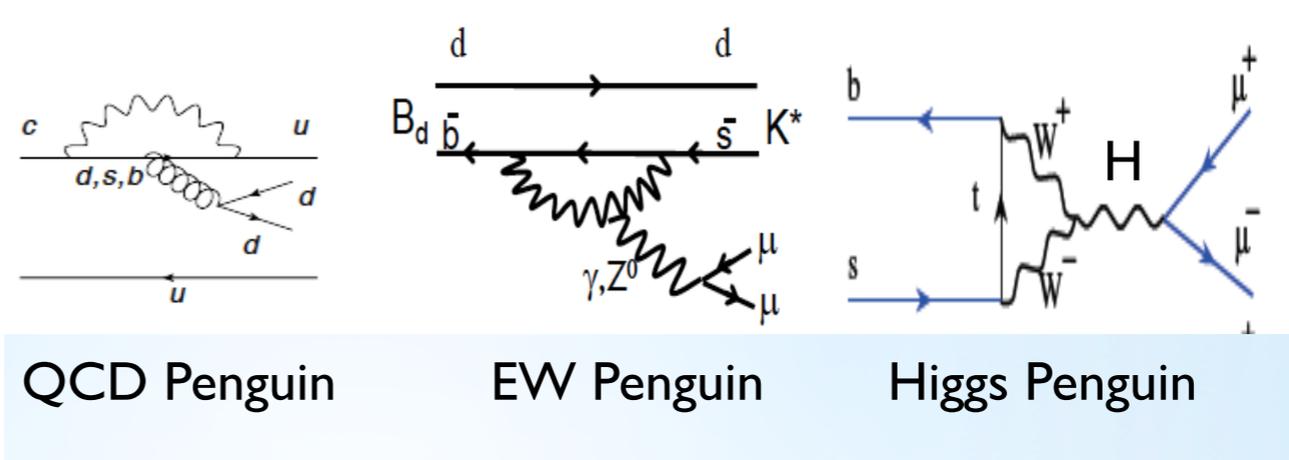
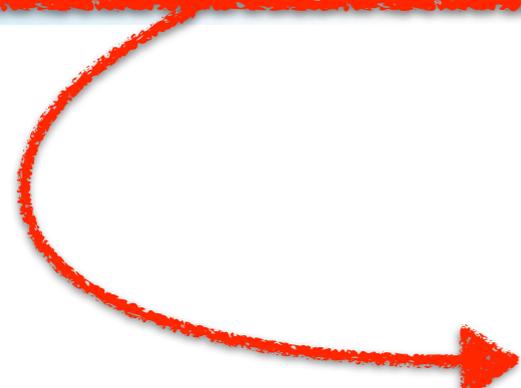
central in the past for searching BSM effects



Flavor Physics

	$b \rightarrow s$ ($ V_{tb}V_{ts} \propto \lambda^2$)	$b \rightarrow d$ ($ V_{tb}V_{td} \propto \lambda^3$)	$s \rightarrow d$ ($ V_{ts}V_{td} \propto \lambda^5$)	$c \rightarrow u$ ($ V_{cb}V_{ub} \propto \lambda^5$)
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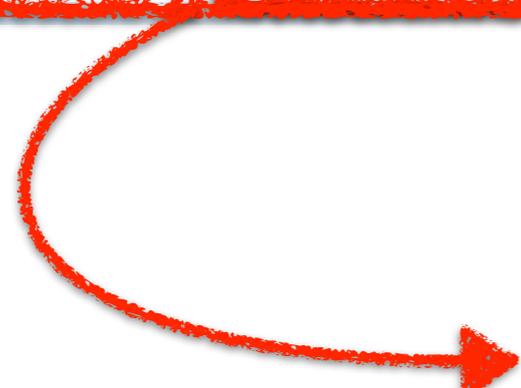
$\Delta F=1$



becoming a mature field
with plenty of new observables
constraining different BSM physics

Flavor Physics

	$b \rightarrow s$ ($ V_{tb} V_{ts} \propto \lambda^2$)	$b \rightarrow d$ ($ V_{tb} V_{td} \propto \lambda^3$)	$s \rightarrow d$ ($ V_{ts} V_{td} \propto \lambda^5$)	$c \rightarrow u$ ($ V_{cb} V_{ub} \propto \lambda^5$)
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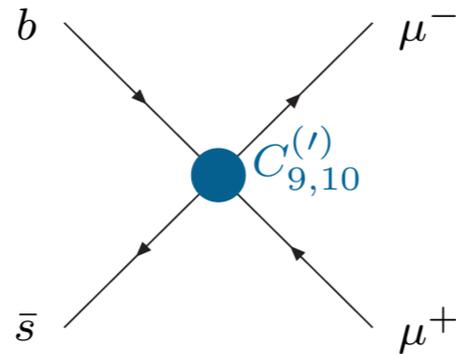

Lost opportunities: Observables with small SM contributions (expected large BSM effects!) are measured at the level of the SM predictions

We have now to dig into the details of the SM contributions

BSM model builders must now be sure to satisfy constraints from $\Delta F=1$ observables (to be added to the $\Delta F=2$):

New physics in $b \rightarrow sll$

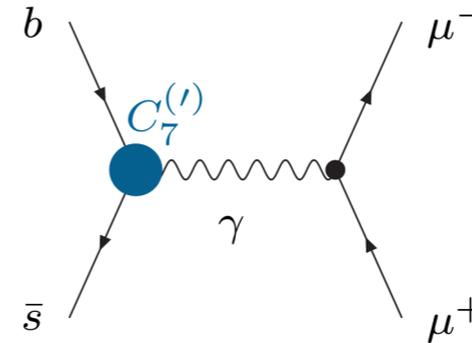
SM and NP particles induce an effective $b\bar{s}\mu^+\mu^-$ coupling



$$\mathcal{O}_9^{(l)} = \frac{\alpha}{4\pi} [\bar{s}\gamma^\mu P_{L(R)} b] [\bar{\mu}\gamma_\mu \mu]$$

$$\mathcal{O}_{10}^{(l)} = \frac{\alpha}{4\pi} [\bar{s}\gamma^\mu P_{L(R)} b] [\bar{\mu}\gamma_\mu \gamma_5 \mu]$$

+ scalar operators
(not relevant for this talk)



$$\mathcal{O}_7^{(l)} = \frac{\alpha}{4\pi} m_b [\bar{s}\sigma_{\mu\nu} P_{R(L)} b] F^{\mu\nu}$$

processes	$C_7^{(l)}$	$C_9^{(l)}$	$C_{10}^{(l)}$
$B \rightarrow X_s \gamma, B \rightarrow K^* \gamma$	✓		
$B \rightarrow X_s \mu^+ \mu^-$	✓	✓	✓
$B_s \rightarrow \mu^+ \mu^-$			✓
$B \rightarrow K^{(*)} \mu^+ \mu^-, B_s \rightarrow \phi \mu^+ \mu^-$	✓	✓	✓

Lars Hofer

Even the SM seems to show some tension with the experimental values... see later

The TeV frontier must be attacked from several fronts



TeV territory

Searches for new particles

Looking for deviations in SM couplings

Looking for new CP-violating & flavor transitions



The expected most sensitive SM particle to BSM:

The Higgs \equiv BEH

it must be *blamed*
for the hierarchy problem!

Higgs EFT (quite developed in the last years) useful for parametrizing deformations from the SM Higgs

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda^{d-4}} \sum_i C_i O_i^{(d)}$$

Elizabeth Jenkins

↪ Higgs \sim coordinates of a $d=4$ space

Aneesh Manohar

deformations can have a simple geometrical interpretation

HEFT: risk of democratizing the BSM effects (not all are equally important):

pseudoNGB Higgs: hff , hVV , h^3

shown in a weakly-coupled calculable model

Aurora Meroni

Higgs physics is maturing:

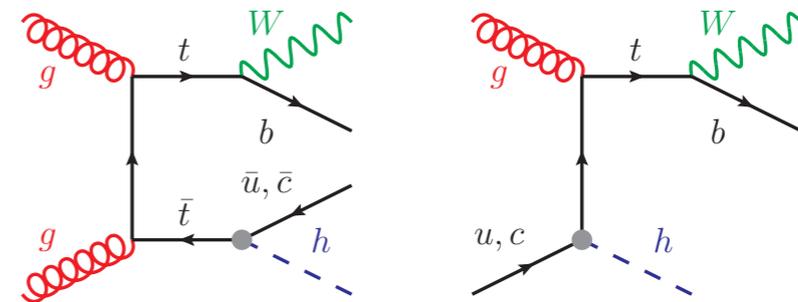
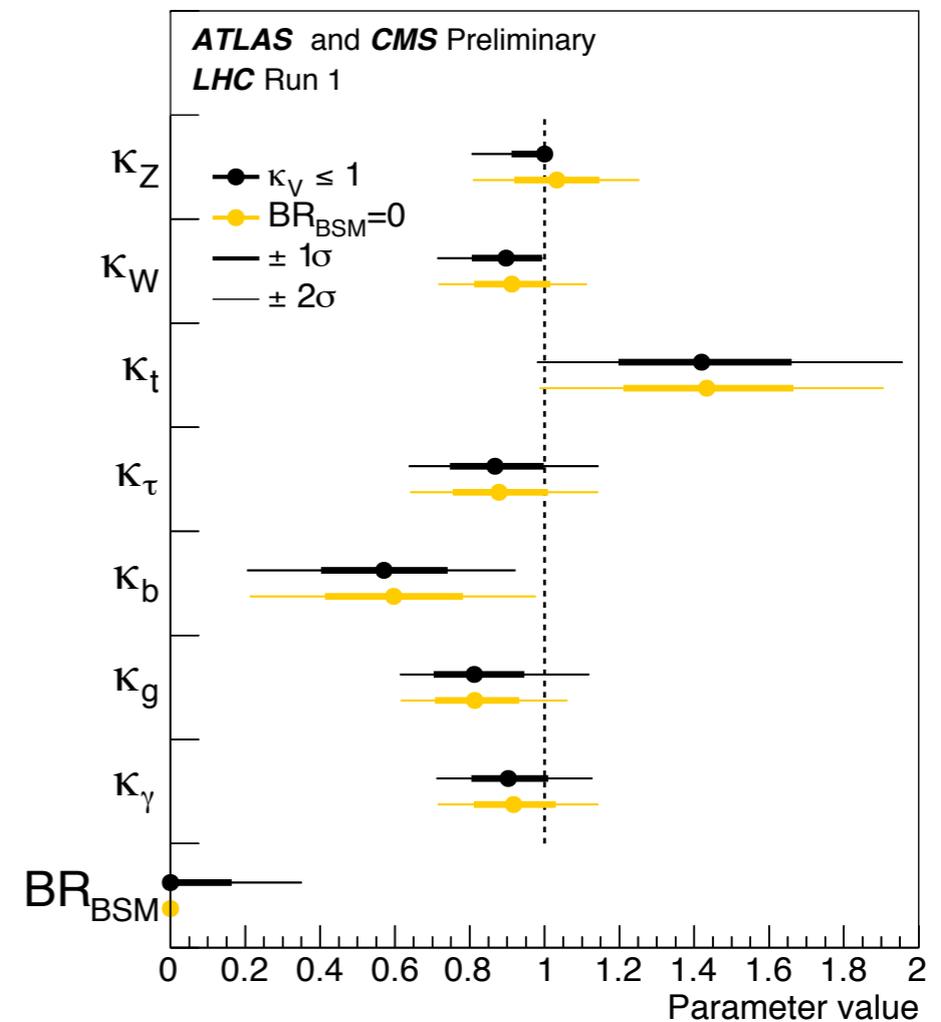
Most important couplings have been measured, showing reasonable agreement with the SM

Still few interesting Higgs couplings to be measured:

- 1) $h^3, hZ\gamma$
- 2) Flavor-violating couplings
- 3) Couplings to light fermions:

hee coupling: too small for the LHC

proposal for looking at **atomic physics:**
 Higgs-Yukawa int. shift atomic freq.: $g_{hee} \lesssim 0.1$



Joachim Kopp

Cedric Delaunay

The TeV frontier must be attacked from several fronts



TeV territory

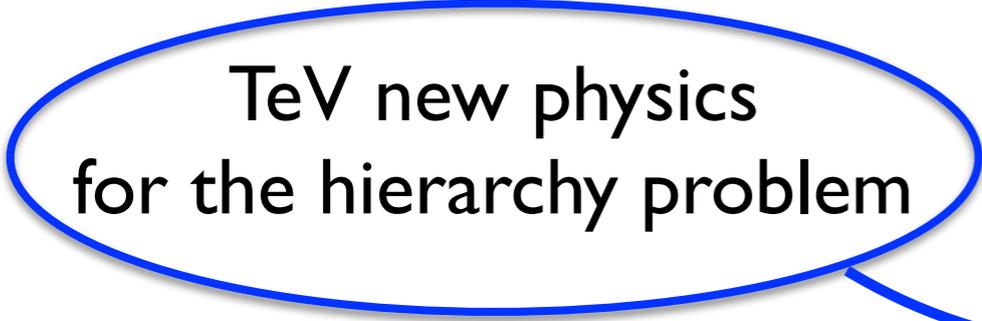
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Motivation for WIMPs:

TeV new physics
for the hierarchy problem



stable remnant



Realistic candidate for DM
WIMP “miracle”!

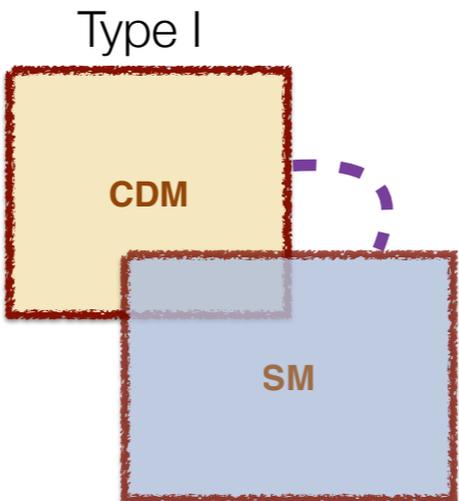
Motivation for WIMPs:

TeV new physics
for the hierarchy problem

stable remnant

Realistic candidate for DM
WIMP “miracle”!

Not only in susy
(neutralino),
but also in strongly-
coupled solutions to
the hierarchy problem



The diagram shows two overlapping rectangular regions. The top-left region is yellow and labeled 'CDM'. The bottom-right region is light blue and labeled 'SM'. A dashed purple line indicates the boundary between the two regions. Above the diagram, the text 'Type I' is written. In the top-left corner of the diagram's frame is a circular icon with four colored dots (red, green, blue, yellow).

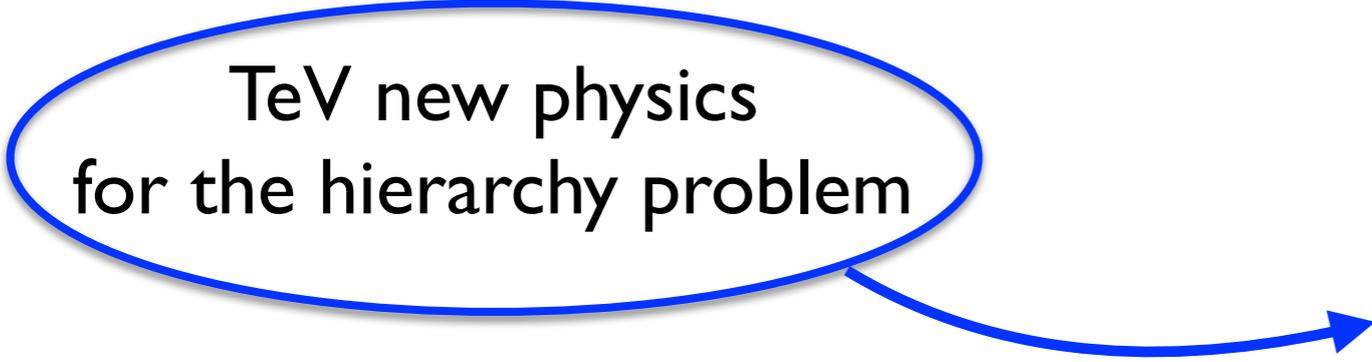
CDM Type I

- TC Baryon¹
- TC Meson²
- Millicharged Comp. DM³
- Stealth DM⁴
- Solitons/Little Higgs⁵
-

¹Nussinov 85 & Barr, Chivukula, Farhi 90
²Gudnason, Kouvaris, Sannino ph-0603014, 0608055
³Kouvaris 1304.7476
⁴Appelquist et al. 1503.04203
⁵Gillioz, 1103.5990

Motivation for WIMPs:

TeV new physics
for the hierarchy problem

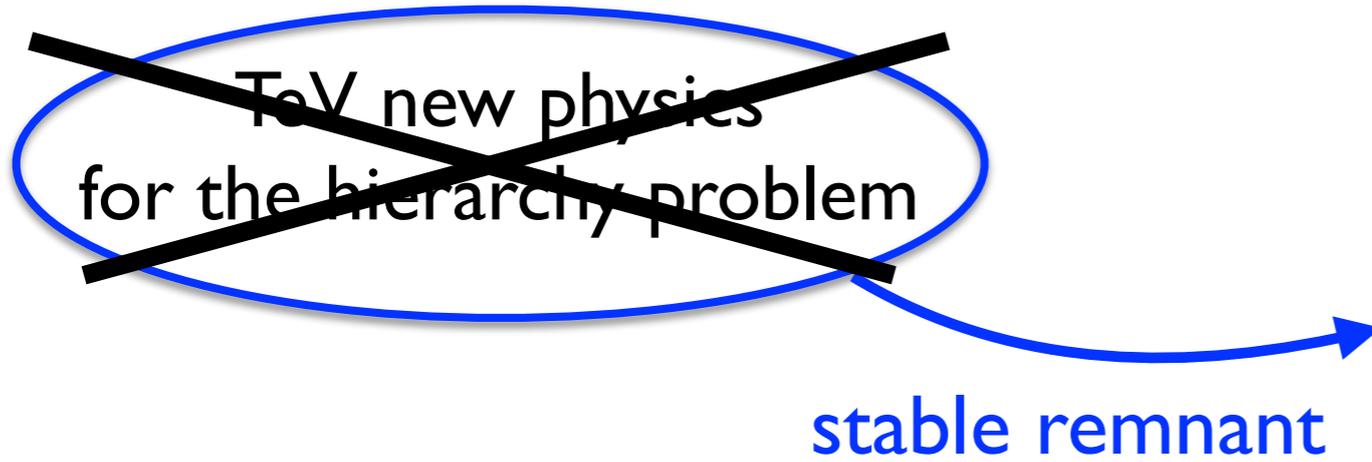


stable remnant

Realistic candidate for DM
WIMP “miracle”!

Finding a WIMP will reinforce the whole TeV-collider program!

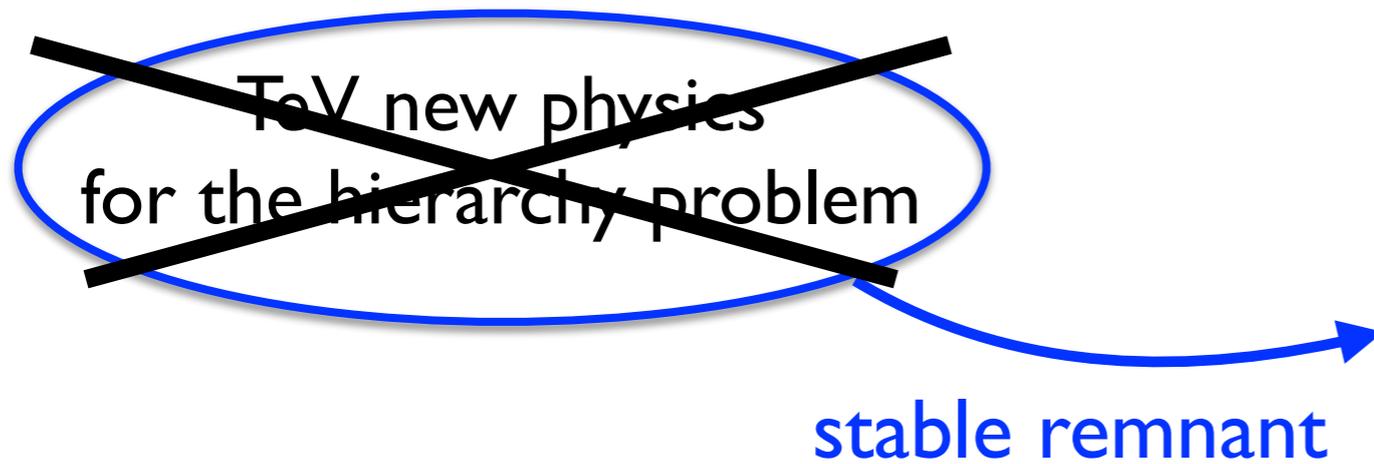
Motivation for WIMPs:



Realistic candidate for DM
WIMP “miracle”!

If no TeV physics explaining
the origin of EW

Motivation for WIMPs:

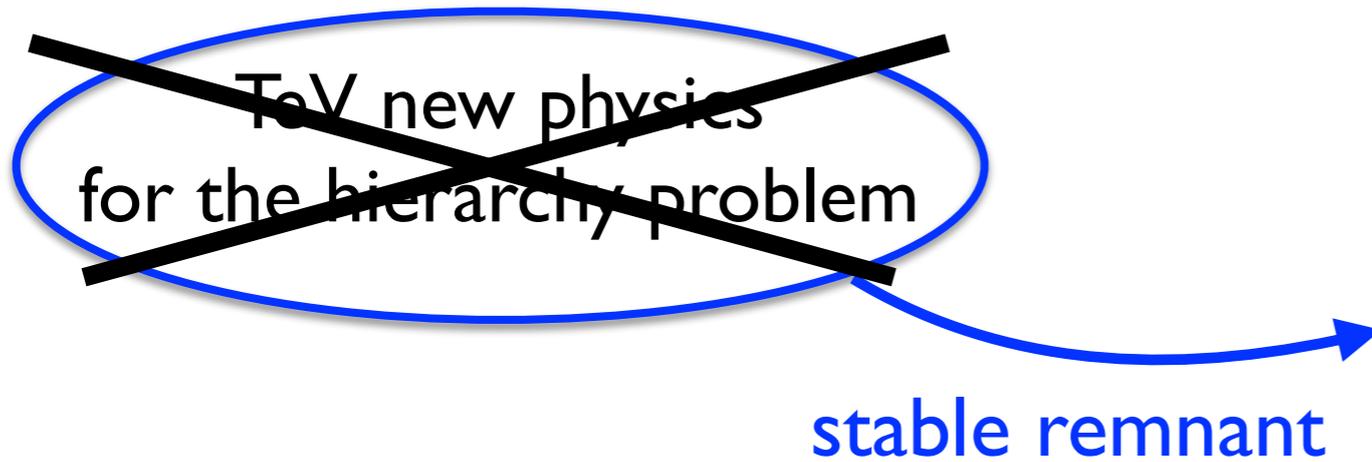


If no TeV physics explaining
the origin of EW

~~Realistic candidate for DM
WIMP "miracle"!~~

The motivation for
WIMPs falls

Motivation for WIMPs:



~~Realistic candidate for DM
WIMP “miracle”!~~

If no TeV physics explaining
the origin of EW

The motivation for
WIMPs falls

Many other
possibilities
for DM beyond
WIMPs are there:

DM classification

At some early cosmological epoch of hot Universe, with temperature $T \gg DM$ mass, the abundance of these particles relative to a species of SM (e.g. photons) was

Normal: Sizable interaction rates ensure thermal equilibrium, $N_{DM}/N_\gamma = 1$. Stability of particles on the scale $t_{Universe}$ is required. *Freeze-out* calculation gives the required annihilation cross section for DM \rightarrow SM of order ~ 1 pb, which points towards weak scale. These are **WIMPs**. Asymmetric DM is also in this category.

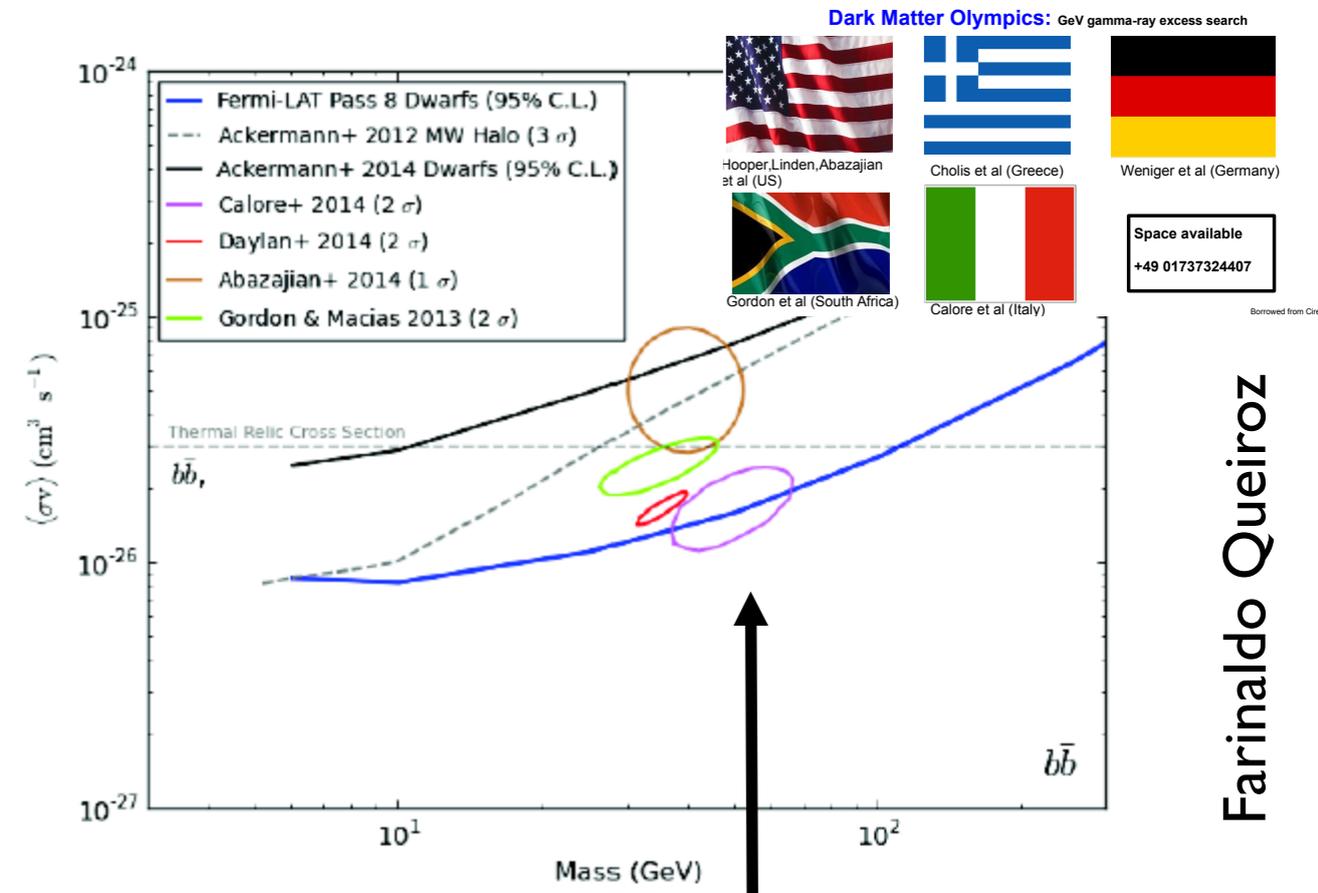
Very small: Very tiny interaction rates (e.g. 10^{-10} couplings from WIMPs). Never in thermal equilibrium. Populated by thermal leakage of SM fields with sub-Hubble rate (*freeze-in*) or by decays of parent WIMPs. [Gravitinos, sterile neutrinos, and other “feeble” creatures – call them **super-WIMPs**]

Huge: Almost non-interacting light, $m < eV$, particles with huge occupation numbers of lowest momentum states, e.g. $N_{DM}/N_\gamma \sim 10^{10}$. “Super-cool DM”. Must be bosonic. Axions, or other very light scalar fields – call them **super-cold DM**.

Maxim Pospelov

- WIMPs searches a mature field → see experimental summary
- Theorists prepare to interpret any excess (in one day!)

1st Olympic Games: Best fit to the GeV γ -ray excess



- Also quite developed are LHC searches for DM:
Missing E_T vs Direct Mediator searches

example: neutralino-like models

Susanne Westhoff

The TeV frontier must be attacked from several fronts

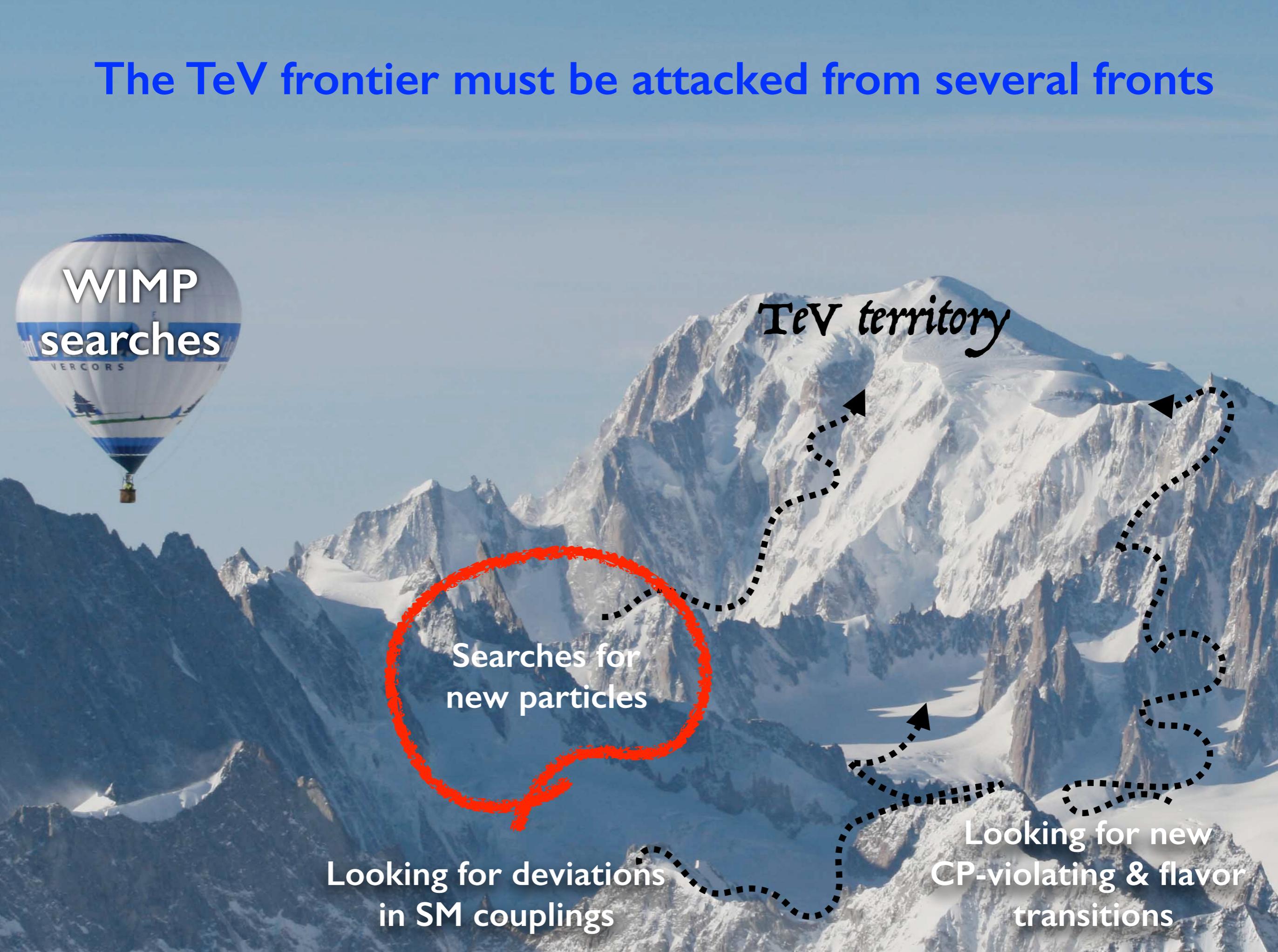


TeV territory

Searches for new particles

Looking for deviations in SM couplings

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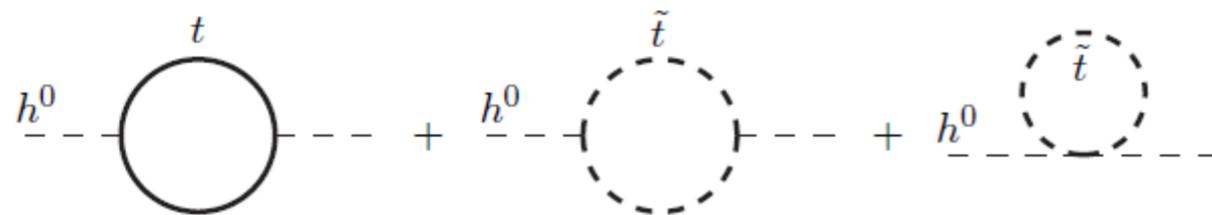
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↪ no new ideas in collider physics?
all moved to anomaly chasing?

Some remark:

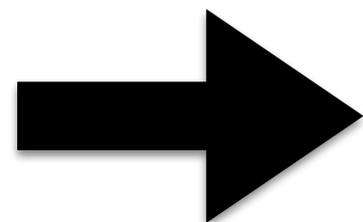
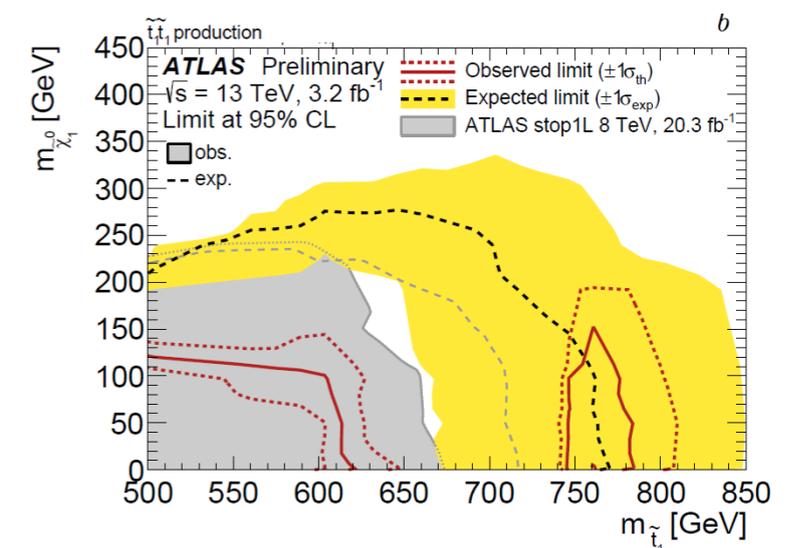
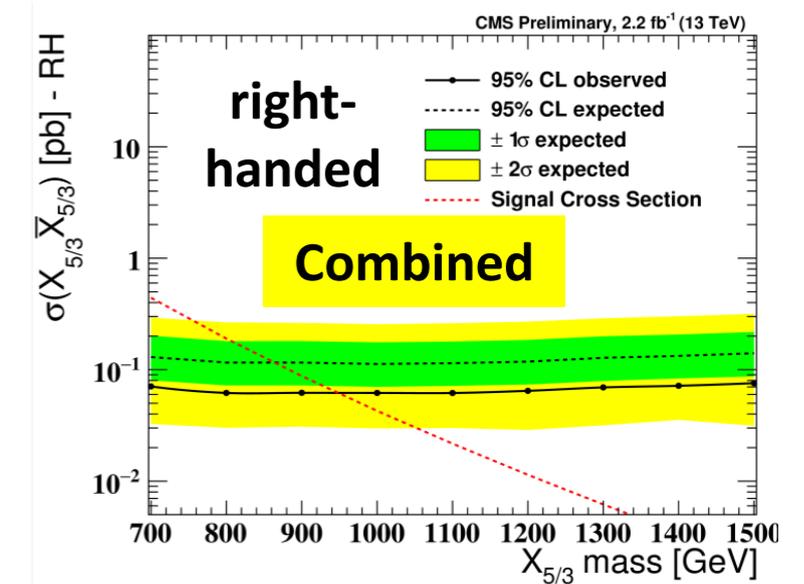
7-8 TeV \rightarrow 13 TeV

Not yet sign of the partners of the top (*Golden BSM modes*):

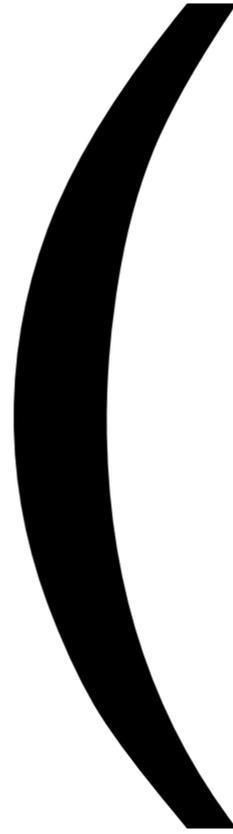


$$m(X_{5/3}) > 960 \text{ GeV}$$

$$m(\tilde{t}, \tilde{b}) \gtrsim 700 - 800 \text{ GeV}$$



Who is keeping the Higgs light?



NEUTRINOS:

Neutrino mass origin probably has to do with physics at scales much larger than TeV:

$$m_\nu \sim \frac{y^2}{g^2} \frac{m_W^2}{M}$$

10^{15} GeV?



The Majorana (see-saw) idea fits nicely and explains why neutrinos are much much lighter than the other SM fermions

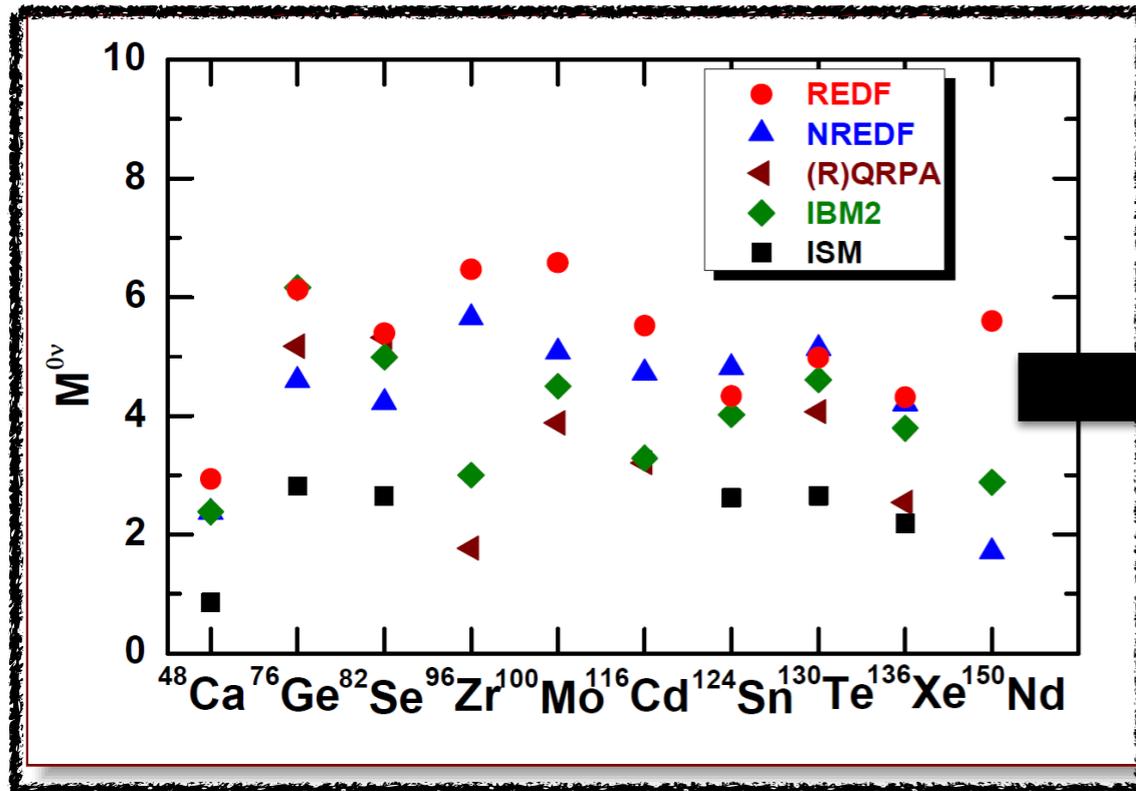
Not a necessity: Possible origin could also be at the TeV (if $y \ll 1$)

- Left-Right models proposal for neutrinos JoAnne Hewett
- Possible scenarios with testable baryogenesis Michele Lucente

$\nu_0 \beta\beta$ decay \rightarrow crucial future experiment

Crucial the ν -exchange nuclear matrix element: $M^{0\nu}$

Jonathan Engel



New approach:

- ▶ **Ab Initio Calculations:** Start from a well justified two-nucleon + three-nucleon Hamiltonian, then solve full many-body Schrödinger equation to good accuracy in space large enough to include all important correlations. At present, works pretty well in systems near closed shells up to $A \approx 50$.
- ▶ **Interboson** Has potential to combine and ground virtues of shell model and density functional theory. (as)

Goal is accurate matrix elements with quantified uncertainty by end of collaboration (5 years from now).

Presence of ν -sterile?

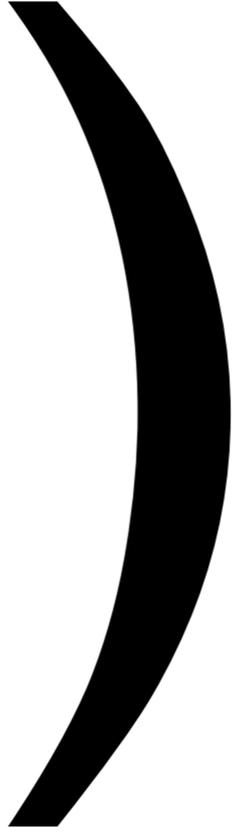
If true, real breakthrough, although its origin will not likely have a connection with the EW scale

reason for the 5% reduction of the reactor antineutrino observed flux? Probably nuclear?

Anna Hayes

“Background” for determining CP-violating phases

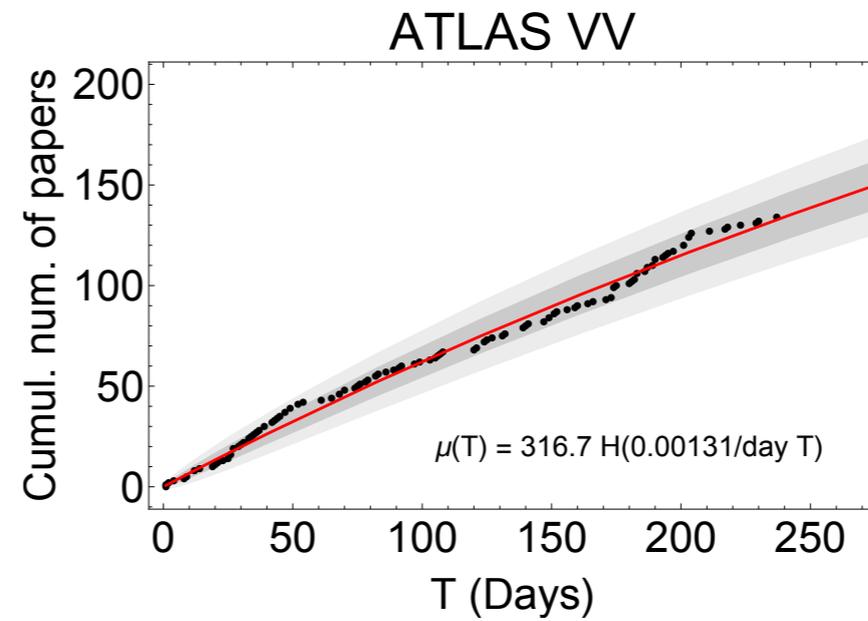
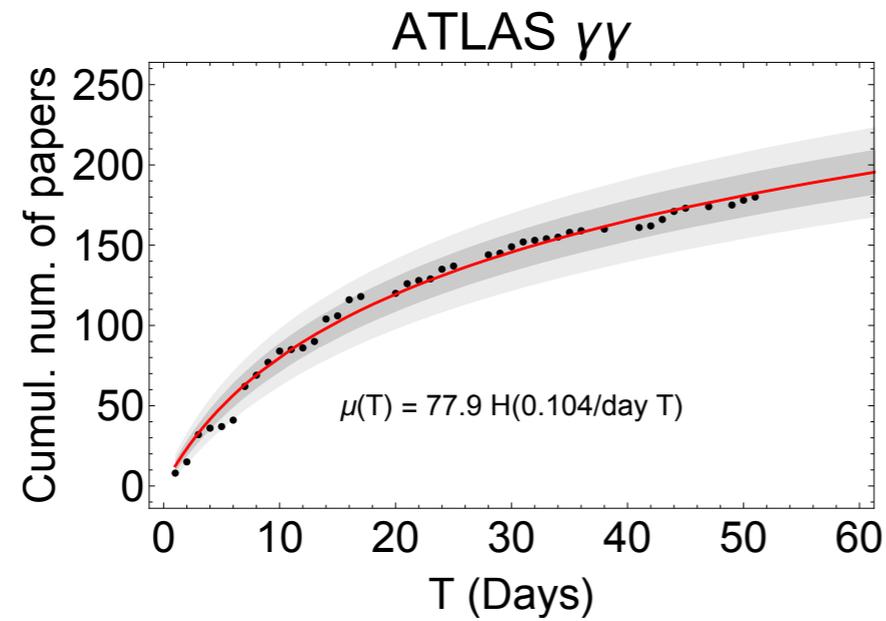
Boris Kayser



ANOMALIES

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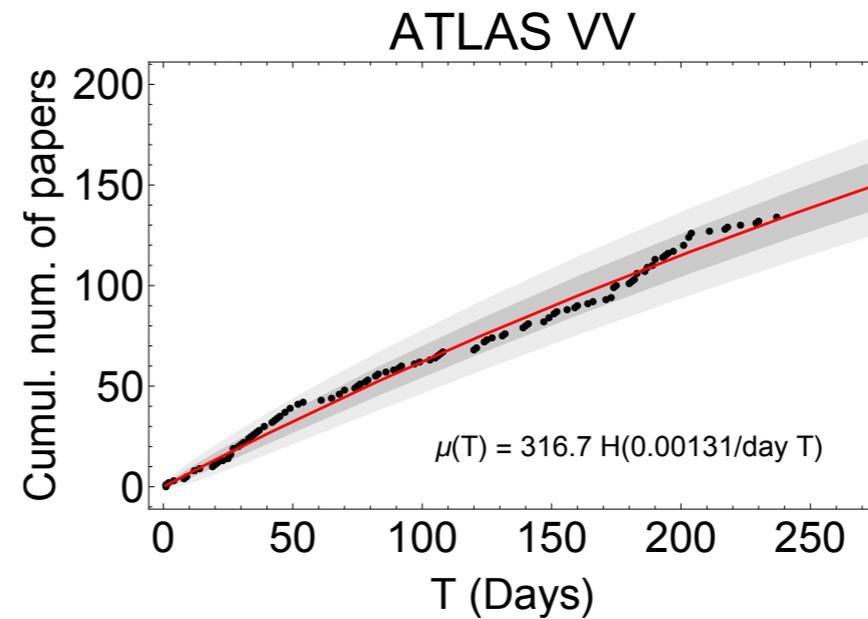
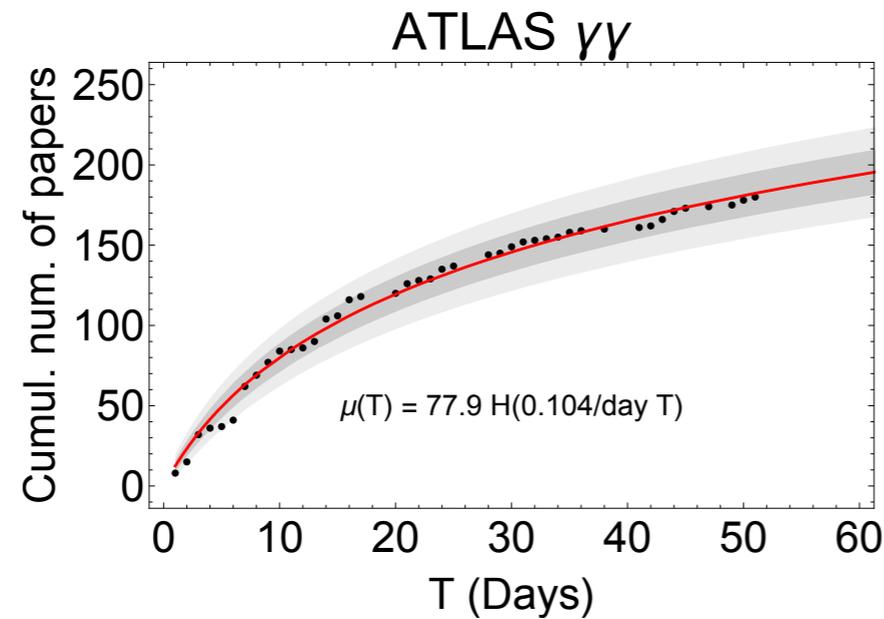
8 th talks out of 28!



arXiv:1603.0204

ANOMALIES

8 th talks out of 28!



arXiv:1603.0204

Understandable: Little BSM experimental data, for too many theorists



Anomalies can even *captivate* very good people

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↪ **Anomalies: Are a very passionate thing!**

Rational approach to anomalies:

- Experimental error?
- Statistical fluctuation?
- SM contributions under control?
- Reasonable BSM could explain it?
 - ↳ *yes, theory bias (or theory nose)*

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B-physics anomalies:

Clean observables? Not having the SM "*breathing behind*"?

$$R(K) = \text{Br}(B \rightarrow K \mu^+ \mu^-) / \text{Br}(B \rightarrow K e^+ e^-) \stackrel{\text{exp.}}{=} 0.75_{-0.07}^{+0.09} \pm 0.04$$

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Who ordered it?



No clear connection with any BSM explaining the EW scale !

New searches:

1) Di-Boson excess:

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After M.Pierini's talk

Dead man walking?



FERMI: γ -ray line at 130 GeV (3.2 σ)

OPERA: ν faster than light ($\sim 6 \sigma$)

CDF: $W+2j$ (3.2 σ)

CDF: top A_{FB} (3.4 σ)



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It will be a pity
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pardoned by ATLAS?

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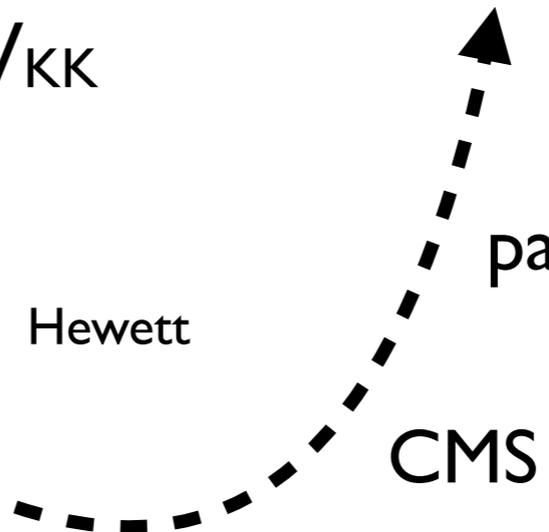
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Aldo DeAndrea

pardoned by ATLAS?

II) Z(+jet)+missing E_T excess:



Not a crazy signal: expected from susy:

$$\tilde{q} \rightarrow \tilde{B} \rightarrow \tilde{H} \rightarrow \begin{matrix} \text{jet} & Z & \text{LSP: missing } E_T \end{matrix}$$

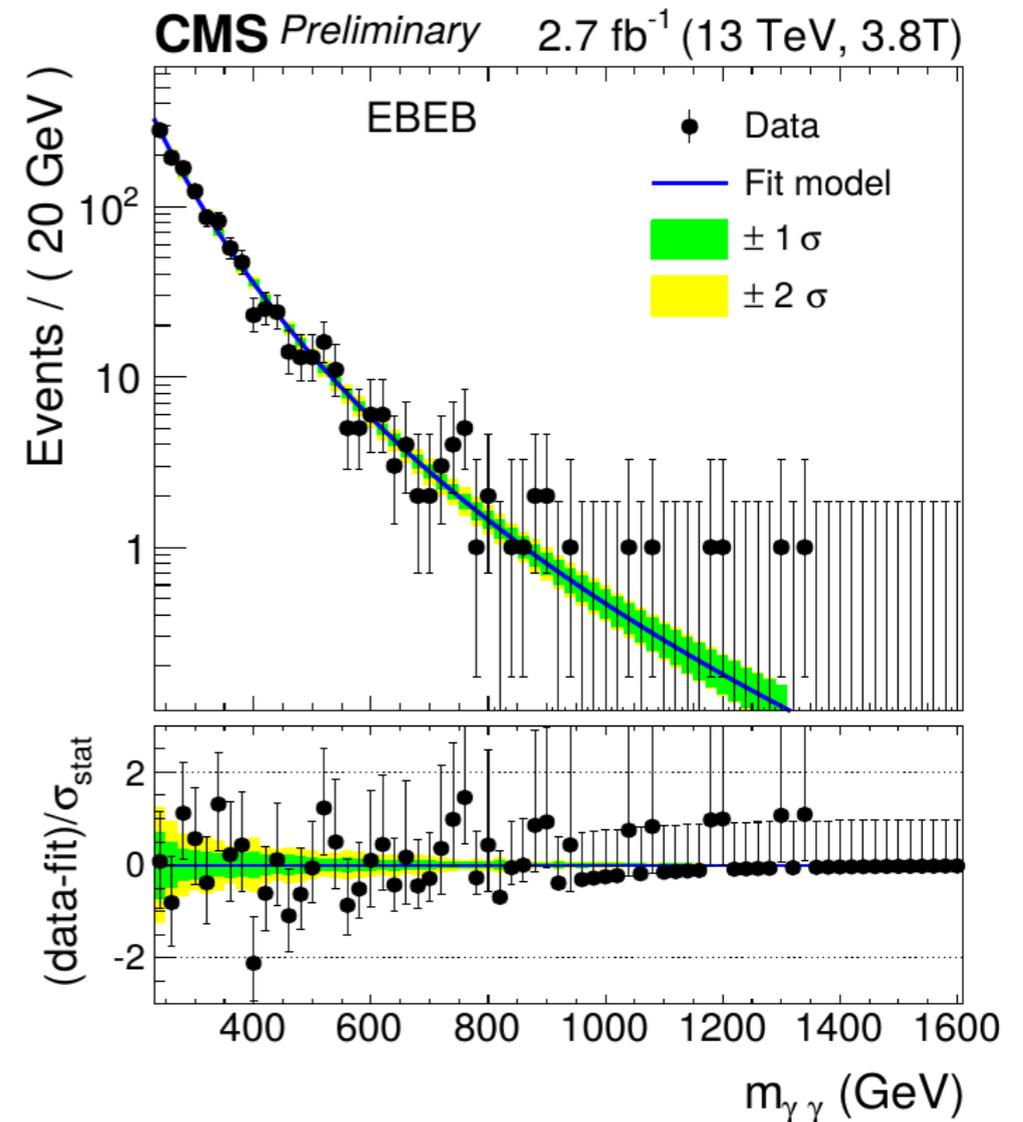
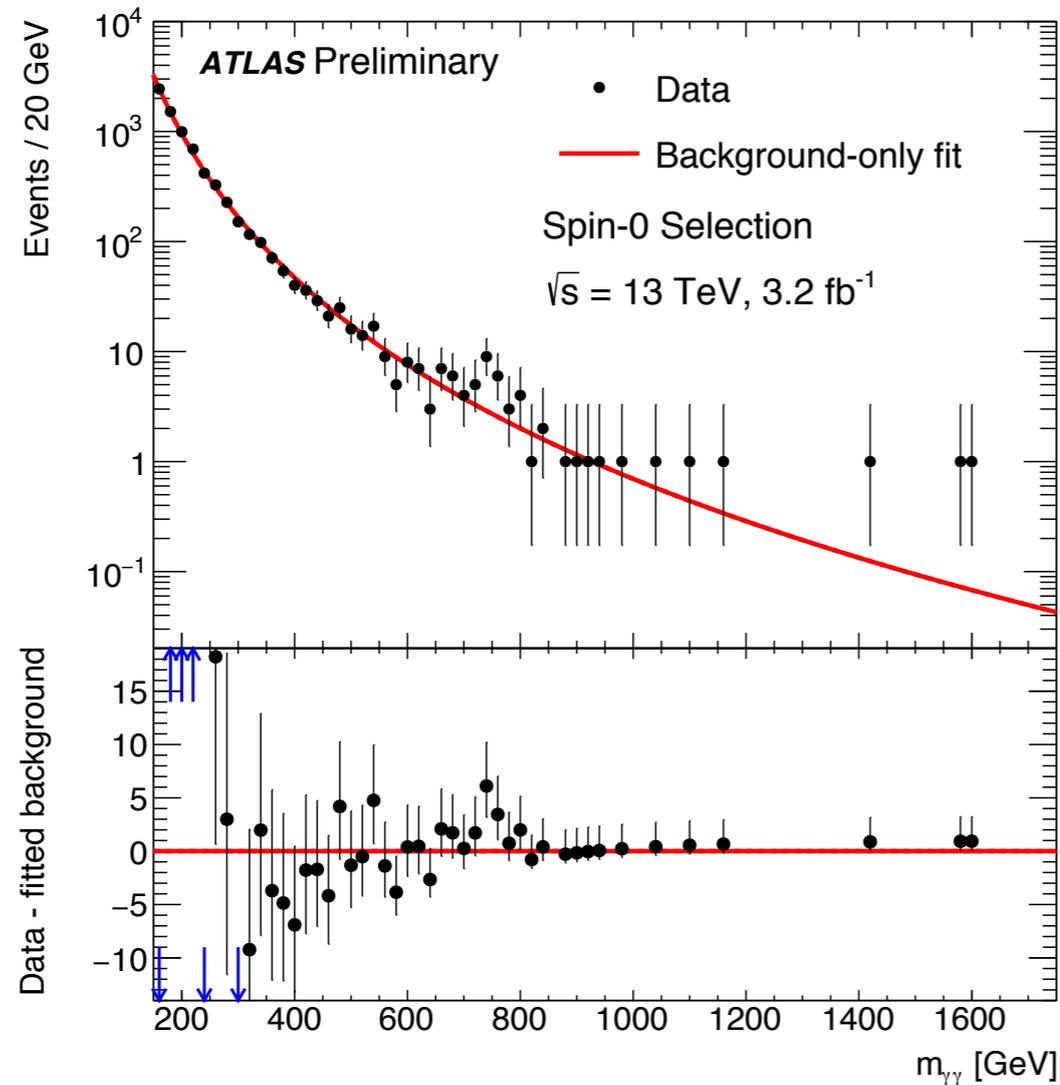
Thomas Rizzo

III) Di-photon excess:

Bumps are in principle clean observables

SPIN-0 ANALYSIS

background-only fit

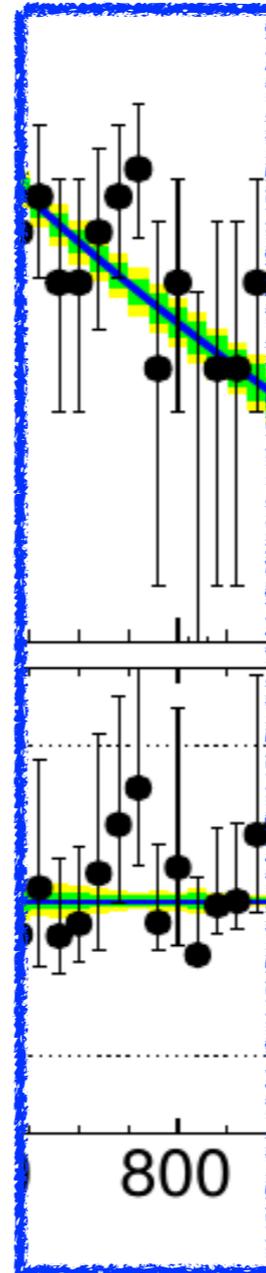
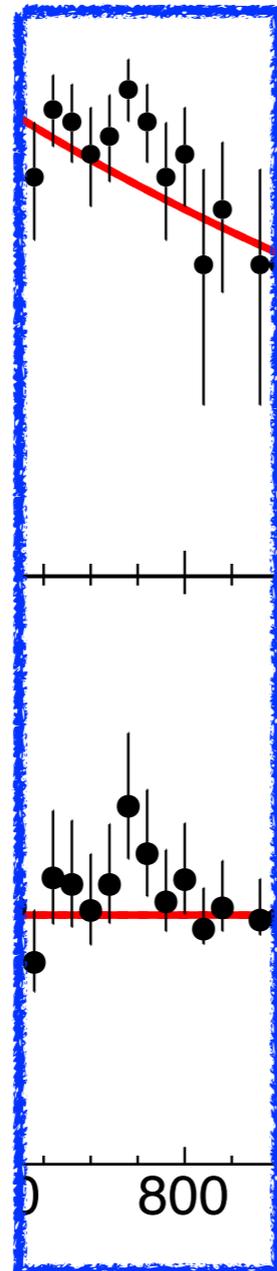


suggested to

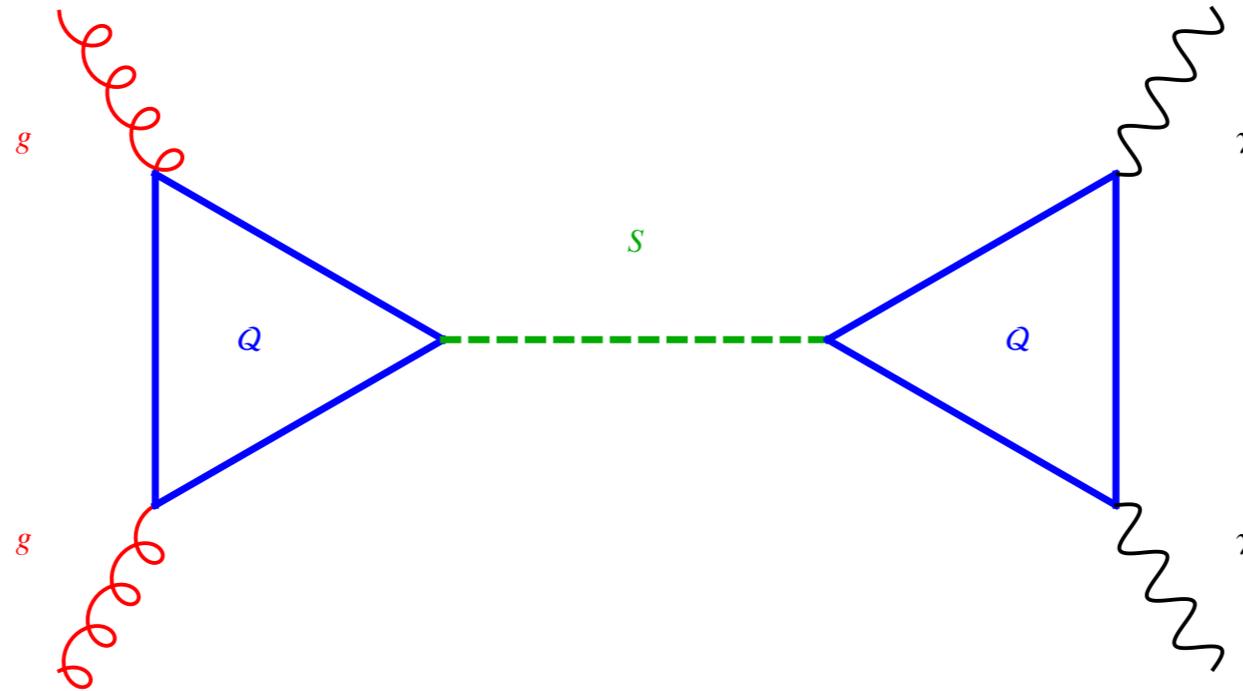


III) Di-photon excess:

Data in the eyes of a theorist



Simple theory interpretation:



A. Strumia

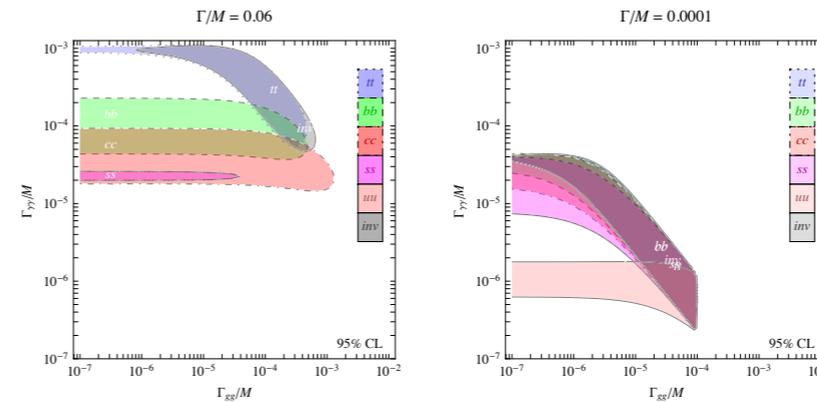
Extra fermions Q or scalars \tilde{Q} needed

Strumia's “modus operandi”:

1) Fit it !

Global fits, $S \leftrightarrow gg, \gamma\gamma, X$

Regions that fit $\sigma(pp \rightarrow \gamma\gamma)_{8,13}$, the width Γ and that satisfy all bounds:



Large width needs $\Gamma(S \rightarrow \gamma\gamma)/M \gtrsim 10^{-5}$: it's big!

2) Add the minimal extension to the SM that explain it !

VolksModell

But this does not answer any fundamental question:

Why is this around the EW scale? What points to?

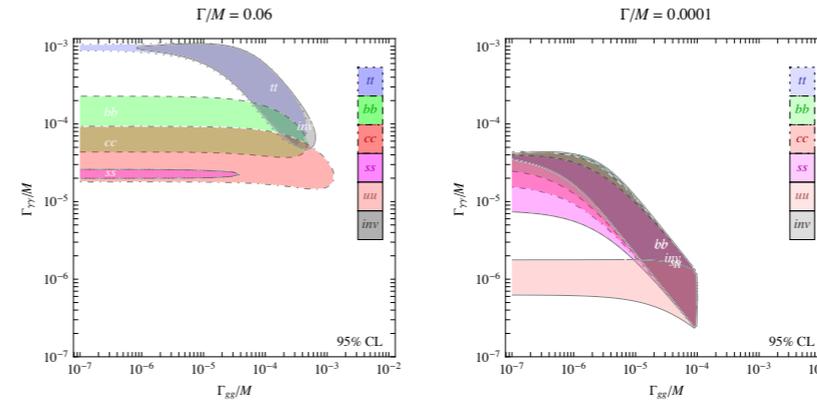
Was expected? What else we can expect?

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The width of S is crucial:

If large \rightarrow strong dynamics

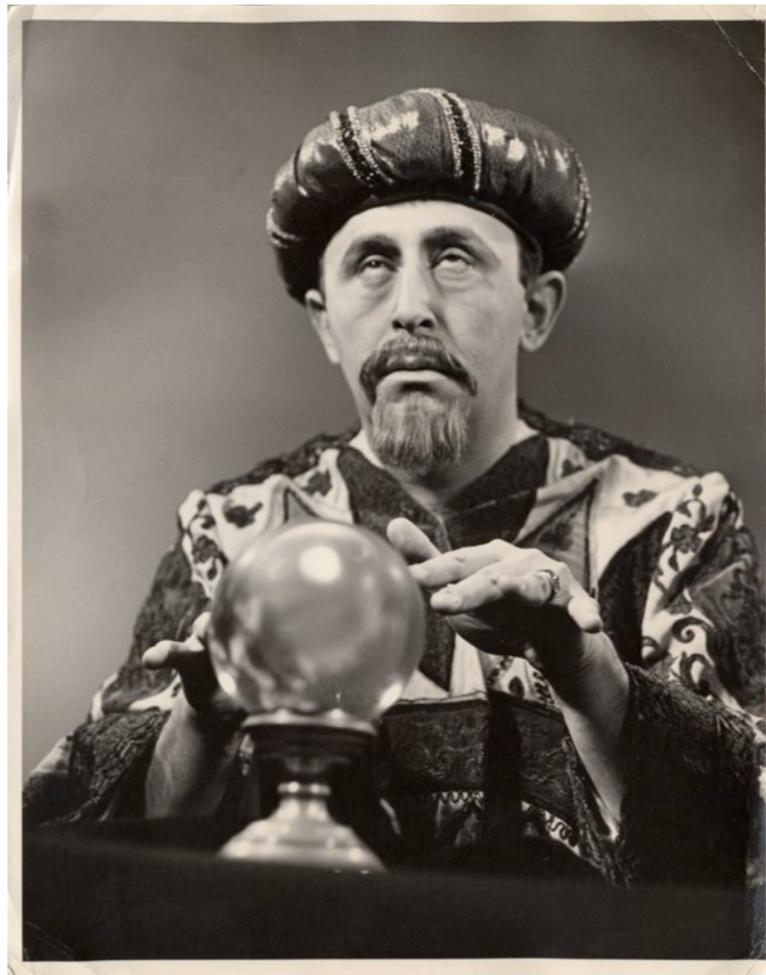
A possibility already proposed to explain the EW scale
Was predicted? Not... but yes.

Not essential, but present in QCD-like theories: η, η'

If small \rightarrow even susy could accommodate it

But why so much extra matter there? Not expected!

Two possible scenarios we can imagine in the future



In August is confirmed...

New Revolution in particle physics!

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**Theory
colleagues**

... avalanche of theory papers



In August is not confirmed...

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New-Physics at the TeV

Pros

Origin of the EW scale

Cons

No new particles seen,
no new flavor-violations seen,
no deviations on Higgs couplings seen,
no deviations on Z/W couplings seen,
no WIMP detected,
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Null results from well-motivated experiments
(as Michelson-Morley experiment) give a motivation for
a change of **paradigm!**

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P.W. Graham, D.E. Kaplan, S.Rajendran
arXiv:1504.07551

Can explain why $m_W \ll M_P$ without new-physics at the TeV

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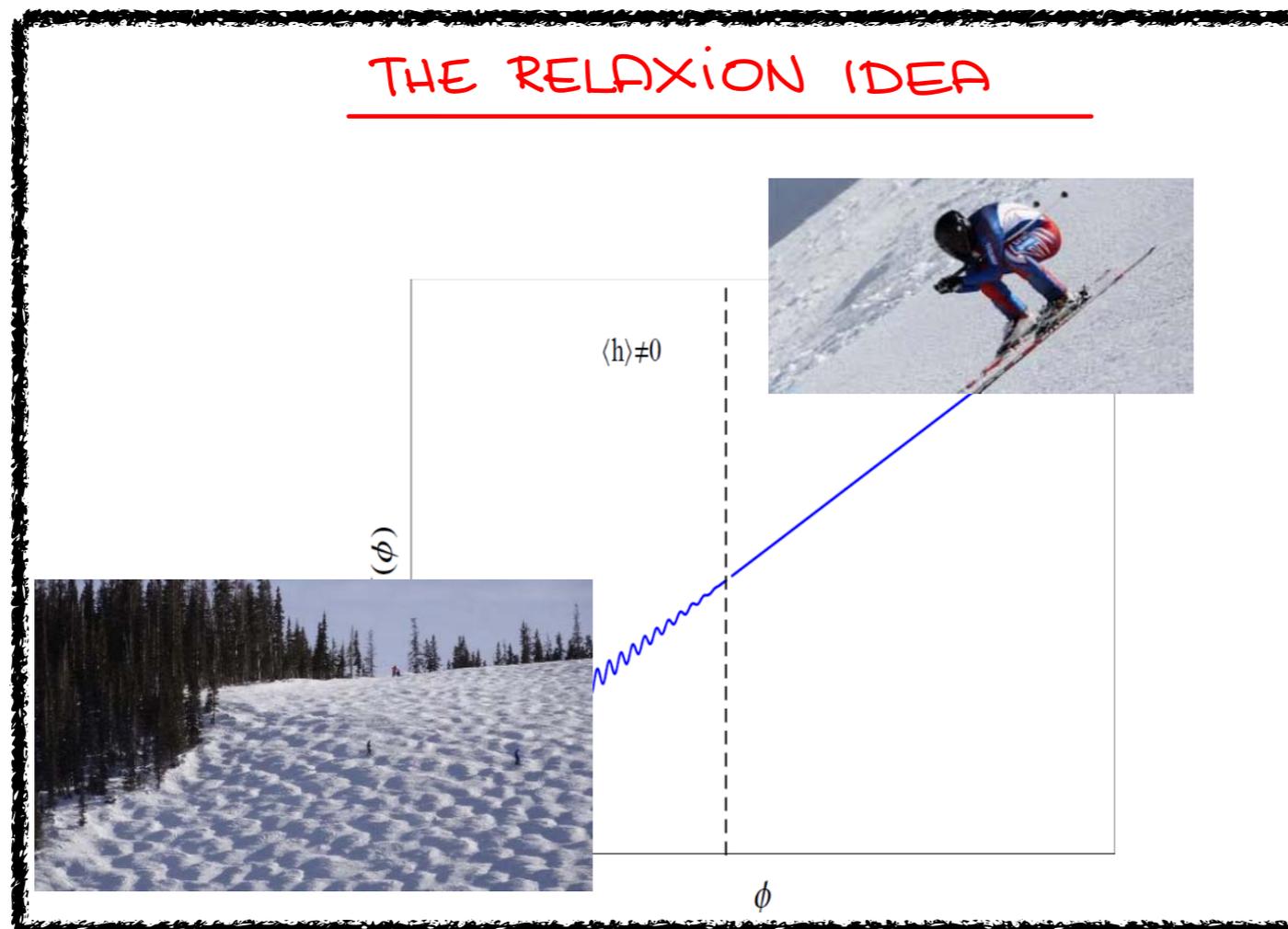
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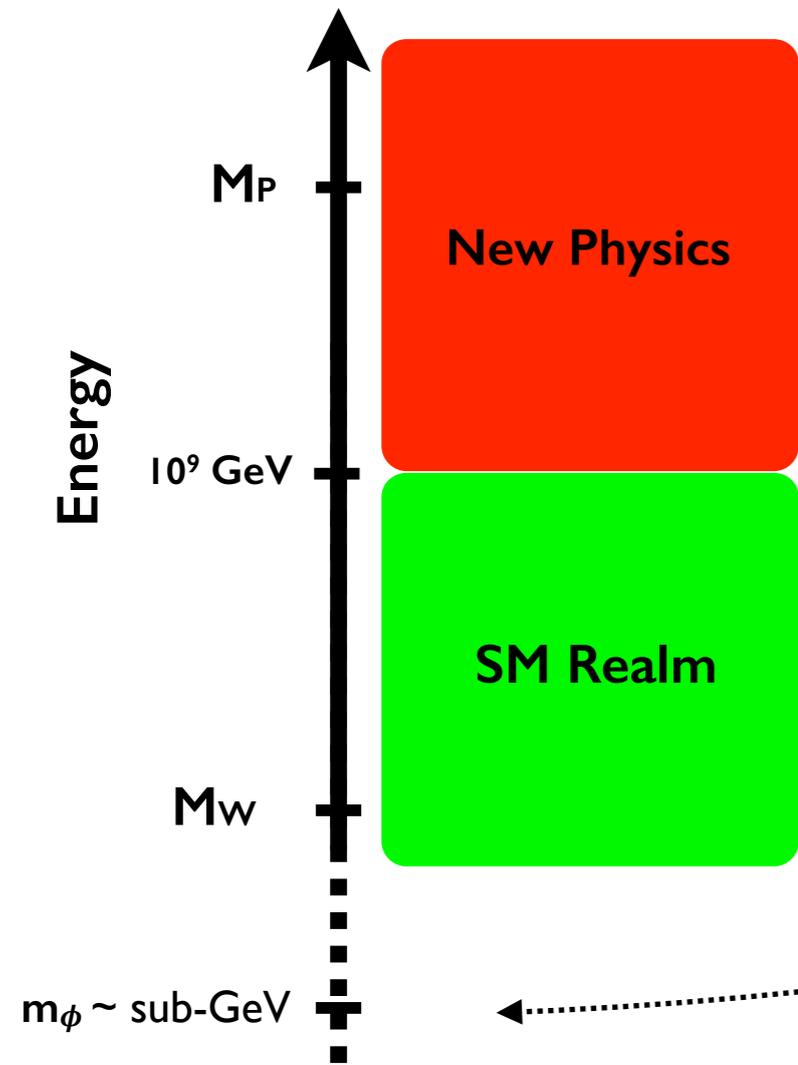
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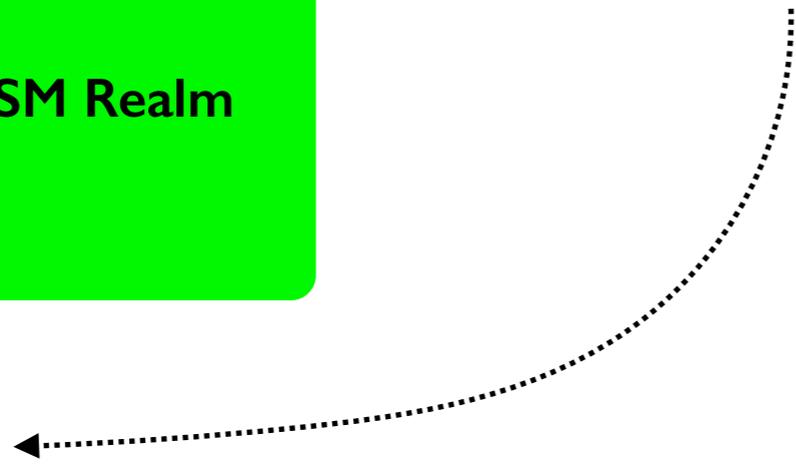


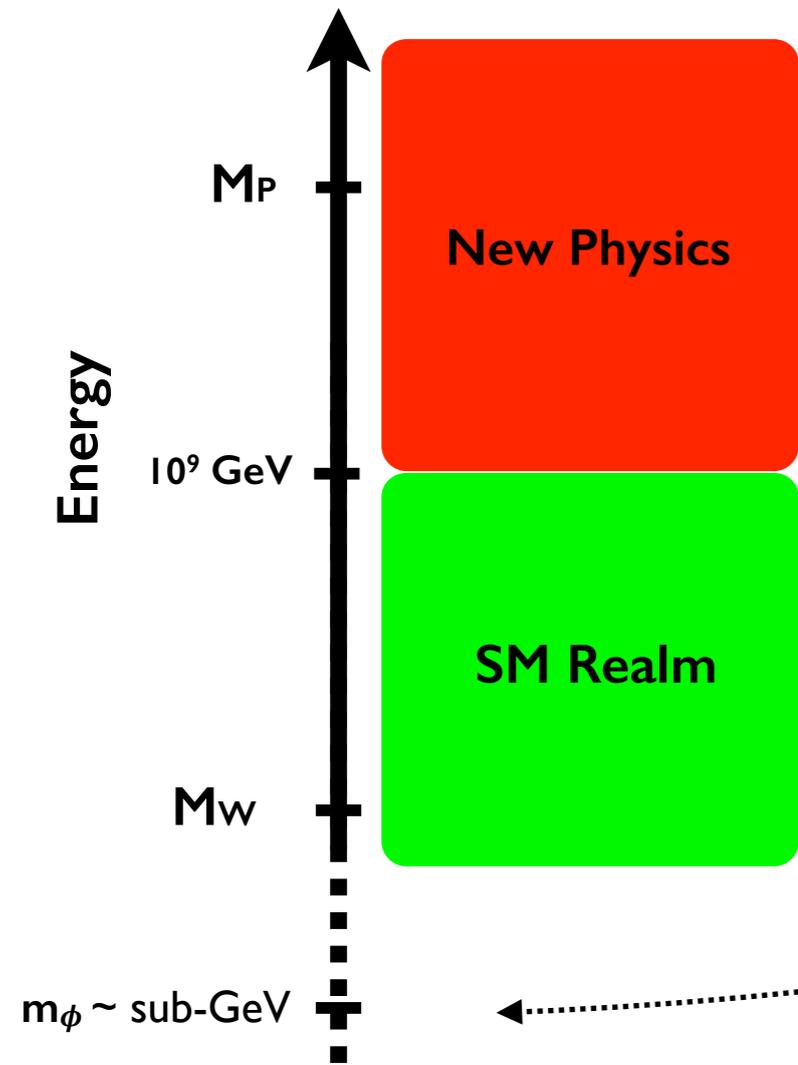
J.R. Espinosa

History
is important!
 ϕ was stabilized
long ago
at very small
values !



We must go back to the low-energy & explore better the weakly-coupled regime

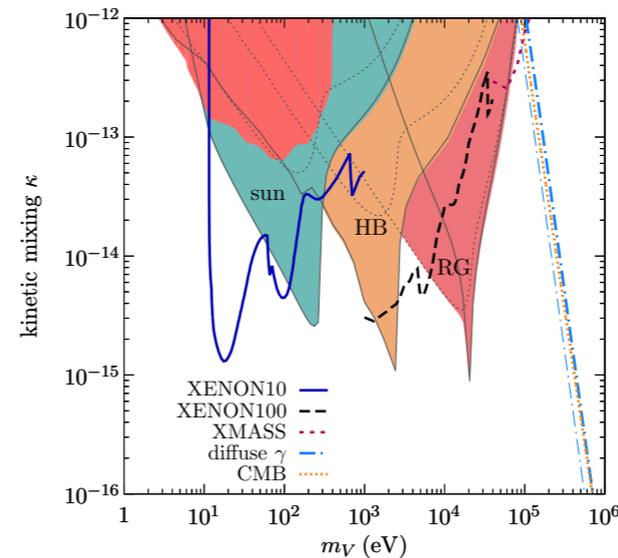
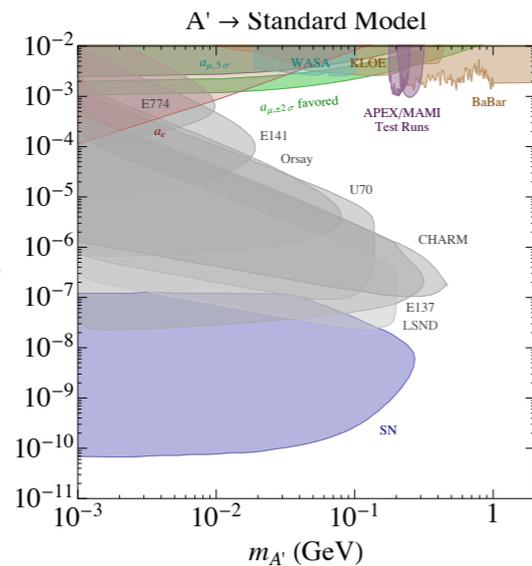




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Very important motivation to connect with already ongoing searches:



Maxim Pospelov

2) Forget to explain the EW scale!

Accept it is tuned: $m_W \ll M_P$

and look for another reasons for the smallness of the EW scale:

DM, gauge-coupling unification, “Anthropic” (Multiverse)

Split-MSSM: Part of the spectrum heavy, other light:
2 Higgs doublet model?

Split-Composite Higgs:

Most of the spectrum is heavy

Signal: Long-lived triplet scalar:

$$T \rightarrow tbSS \quad \Rightarrow \quad c\tau \approx \underbrace{0.2 \text{ mm}}_{\text{can produce a displaced vertex!}} \left(\frac{1}{c_3^T}\right)^2 \left(\frac{8}{g_\rho}\right)^3 \left(\frac{3 \text{ TeV}}{m_T}\right)^5 \left(\frac{f}{10 \text{ TeV}}\right)^4$$

$f > 10 \text{ TeV} = \text{long-lived decay}$

Epilogue

*It was the best of times,
It was the worst of times,*

...

*It was the spring of hope,
It was the winter of despair*

A Tale of Two Cities

- After the Higgs, we start a very different phase in particle physics:

*We could discover plenty,
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Thank you!

and to all that made possible another successful Moriond!