

Probing Radiative Neutrino Mass Models Using Trilepton Channel at the LHC

Dounia Cherigui

*Department of Physics, University of Sciences and
Technology of Oran, USTO-MB Algeria.*

*51st Rencontre de Moriond
March 12, 2016*

Based on work in progress with A. Ahriche, C. Guella, S. Nasri

- 1 Model & Space Parameter
- 2 Trilepton signal phenomenology
- 3 Current Constraints at the LHC
- 4 Benchmark Analysis
 - The Processes $ee\mu$ & $e\mu\mu$
 - LFV Background Free Channel
- 5 Summary & Conclusion

- In this work, we consider a class of models that contain the following term in the Lagrangian

$$\mathcal{L} \supset f_{\alpha\beta} L_{\alpha}^T C \epsilon L_{\beta} S^+ - m_S^2 S^+ S^- + \text{h.c.}$$

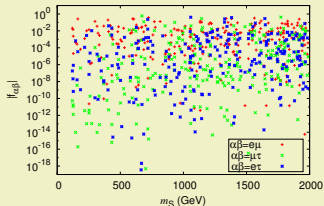
- The interaction above induces LFV effects via processes such as $\mu \rightarrow e\gamma$ and $\tau \rightarrow \mu\gamma$, with branching fractions..

$$\mathcal{B}(\mu \rightarrow e\gamma) \simeq \frac{\alpha v^4}{384\pi} \frac{|f_{\tau e}^* f_{\mu\tau}|^2}{m_S^4}$$

$$\mathcal{B}(\tau \rightarrow \mu\gamma) \simeq \frac{\alpha v^4}{384\pi} \frac{|f_{\tau e}^* f_{\mu e}|^2}{m_S^4}$$

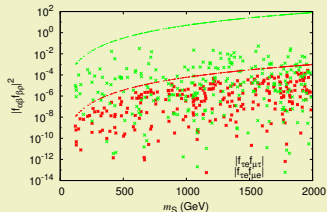
- A new contribution to the muon's anomalous magnetic moment is induced at 1-loop.

$$\delta a_{\mu} \sim \frac{m_{\mu}^2}{96\pi^2} \frac{|f_{e\mu}|^2 + |f_{\mu\tau}|^2}{m_S^2}$$

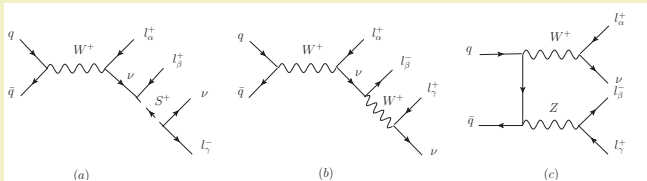


$$\mathcal{B}(\mu \rightarrow e + \gamma) < 5.7 \times 10^{-13} \text{ [J.Adam.2013]}$$

$$\mathcal{B}(\tau \rightarrow \mu + \gamma) < 4.8 \times 10^{-8} \text{ [K.A.Olive.2014]}$$



There is an interesting possibility to produce a singly charged scalar at the parton level process at the LHC



Signal

Vs

SM Background

- We have 9 contributions to this trilepton signal, three of them are background free
- CalcHEP is used to generate both the SM background events as well the events from processes due to the **LFV** interactions for $\sqrt{s} = 8$ and 14 TeV.
- The excess of events looked for after the selection cut is $N_{ex} = \mathcal{L}_{int} (\sigma_M - \sigma_B)$, Therefore the signal significance is given by

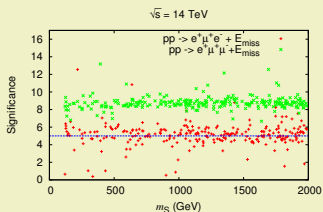
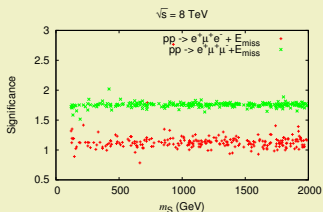
$$S = \frac{N_{ex}}{\sqrt{N_{ex} + N_B}} = \frac{N_{ex}}{\sqrt{N_M}}$$

$$[N_{ex} = N_M - N_B] \propto [2 \operatorname{Re} (\mathcal{M}_{SM}^\dagger \mathcal{M}_{non-SM}) + |\mathcal{M}_{non-SM}|^2] \propto |f_{\alpha\rho} f_{\beta\rho}|^2.$$

where $|\mathcal{M}_{non-SM}| \ll \ll$

The analysis was based on the following CMS collaboration criteria with $\mathcal{L}_{int} = 19.5 \text{ fb}^{-1}$ at $\sqrt{s} = 8 \text{ TeV}$:

- The presence of at least three isolated leptons (muon, electron).
- $75 \text{ GeV} < M_{\ell+\ell^-} < 105 \text{ GeV}$.
- $p_T^\ell > 10 \text{ GeV}$.
- $\cancel{E}_T < 50 \text{ GeV}$.
- $|\eta^\ell| < 2.4$.

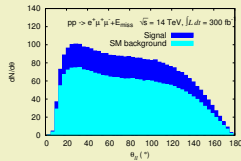
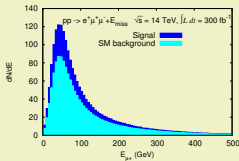
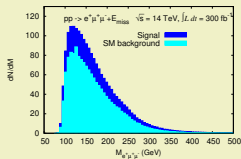
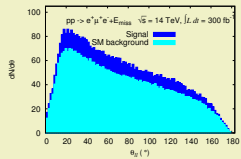
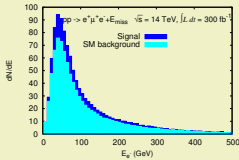
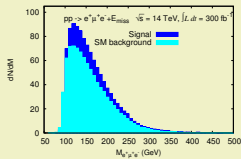


⇒ Two benchmark points selected from the allowed parameter space of the model

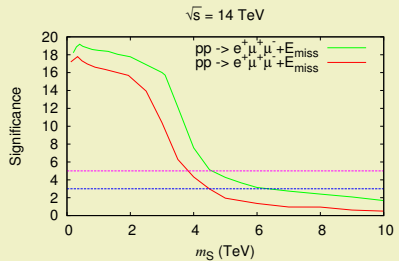
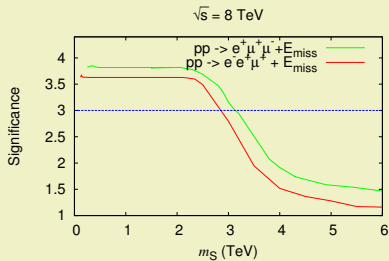
Point	$m_S \text{ (GeV)}$	$f_{e\mu}$	$f_{e\tau}$	$f_{\mu\tau}$
B_1	471.8	$-(9.863 + i8.774) \times 10^{-2}$	$-(6.354 + i2.162) \times 10^{-2}$	$(0.78 + i1.375) \times 10^{-2}$
B_2	1428.5	$(5.646 + i549.32) \times 10^{-3}$	$-(2.265 + i1.237) \times 10^{-1}$	$-(0.41 - i3.58) \times 10^{-2}$

- Due to the difficulty in identifying the tau lepton at the LHC, only the final state leptons $\ell = e, \mu$, are considered here, with \cancel{E}_T can be $\bar{\nu}_e, \bar{\nu}_\mu$ or $\bar{\nu}_\tau$.

$e^+ \mu^+ e^- + \cancel{E}_T @ 8 \text{ TeV}$	$e^+ \mu^+ e^- + \cancel{E}_T @ 14 \text{ TeV}$	$e^+ \mu^+ \mu^- + \cancel{E}_T @ 8 \text{ TeV}$	$e^+ \mu^+ \mu^- + \cancel{E}_T @ 14 \text{ TeV}$
$70 < M_{e^- e^+} < 110$	$70 < M_{e^- e^+} < 110$	$80 < M_{\mu^- \mu^+} < 100$	$80 < M_{\mu^- \mu^+} < 110$
$M_{e^+ \mu^+} < 200$	$M_{e^+ \mu^+} < 230$	$M_{e^+ \mu^+} < 200$	$M_{e^+ \mu^+} < 230$
$M_{e^- \nu} < 206$	$M_{e^- \nu} < 220$	$M_{\mu^- \nu} < 185$	$M_{\mu^- \nu} < 245$
$10 < p_T^\ell < 100$	$10 < p_T^\ell < 90$	$10 < p_T^\ell < 100$	$10 < p_T^\ell < 130$
$ \eta^\ell < 3$	$ \eta^\ell < 3$	$ \eta^\ell < 3$	$ \eta^\ell < 3$
$\cancel{E}_T < 100$	$\cancel{E}_T < 90$	$\cancel{E}_T < 90$	$\cancel{E}_T < 120$

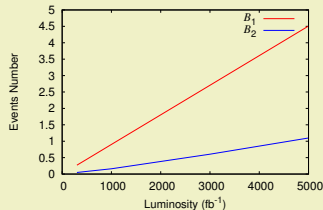
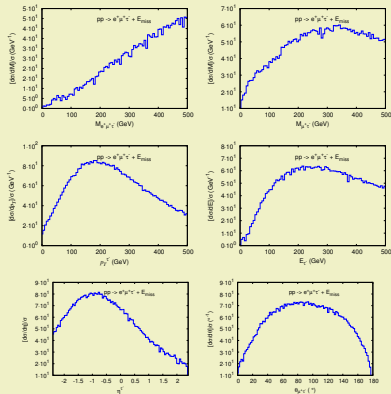


Process	Benchmark	$N_{20.3}$	$S_{20.3}$	N_{300}	S_{300}
$p, p \rightarrow e^+ \mu^+ e^- + \cancel{E}_T$	B_1	70.42	3.651	1689.6	17.363
	B_2	69.69	3.618	1470	15.289
$p, p \rightarrow e^+ \mu^+ \mu^- + \cancel{E}_T$	B_1	71.21	3.831	2066.7	19.210
	B_2	70.44	3.793	1974.9	18.983



- Another way to probe this class of models, is to extend our analysis and consider one of the background-free channels $\ell^+ \ell^+ \ell^- \equiv e^+ \mu^+ \tau^-$, where the tau lepton is identified through its hadronic decay rather than its leptonic one in order to avoid additional source of missing energy.

$$N_{BG-free} = L \times \sigma(pp \rightarrow e^+ \mu^+ \tau^- + \cancel{E}_T) \times \mathcal{B}(\tau \rightarrow \text{hadrons}), \quad (1)$$



- 1 The parameter space points B_1 and B_2 are expected to have a favorable cross sections at the LHC.
- 2 The significance differs according to the final state detected, and to the scale of the energy
- 3 The probability of observing the trilepton signal events through the signature $e^+ \mu^+ \mu^-$, is improved compared to that found across the $e^+ \mu^+ e^- + \cancel{E}_T$ at $\sqrt{s} = 8 \text{ TeV}$ and $\sqrt{s} = 14 \text{ TeV}$.
- 4 Research carried out at low mass of charged scalar, is more favorable over high mass.
- 5 From the collected results, The charged scalar S^\pm seems more detectable at LHC comparing to the singly charged scalar h^\pm of the Zee Babu [K. S. Babu et al. 2003] which it estimates around ten events produced at the LHC for any h^\pm mass in the TeV range.

THANKS