

# Flavor Physics

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[previously in 2009]

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# Flavor Physics

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|        |                           |                               |                                 |
|--------|---------------------------|-------------------------------|---------------------------------|
| Quarks | $u$<br>up                 | $c$<br>charm                  | $t$<br>top                      |
|        | $d$<br>down               | $s$<br>strange                | $b$<br>bottom                   |
|        | $\nu_e$<br>e- Neutrino    | $\nu_\mu$<br>$\mu$ - Neutrino | $\nu_\tau$<br>$\tau$ - Neutrino |
|        | $e$<br>electron           | $\mu$<br>muon                 | $\tau$<br>tau                   |
|        | I                         | II                            | III                             |
|        | The Generations of Matter |                               |                                 |

# In the Standard Model

- ✗ Gauge sector entirely fixed by symmetry

$$i\bar{\psi}\not{D}\psi \qquad D_\mu = \partial_\mu - ig_s t_a A_\mu^a - ig \mathbf{T} \cdot \mathbf{W}_\mu - ig' \frac{Y}{2} B_\mu$$

- ✗ Flavor sector loose (a bunch of parameters)
- 13 of 21 are fermion masses and q.mixing parameters

|                           |                        |                               |                                 |
|---------------------------|------------------------|-------------------------------|---------------------------------|
| Quarks                    | $u$<br>up              | $c$<br>charm                  | $t$<br>top                      |
|                           | $d$<br>down            | $s$<br>strange                | $b$<br>bottom                   |
| Leptons                   | $\nu_e$<br>e- Neutrino | $\nu_\mu$<br>$\mu$ - Neutrino | $\nu_\tau$<br>$\tau$ - Neutrino |
|                           | $e$<br>electron        | $\mu$<br>muon                 | $\tau$<br>tau                   |
| I      II      III        |                        |                               |                                 |
| The Generations of Matter |                        |                               |                                 |

# We know

fermions come in 3 generations

$$\begin{pmatrix} \nu_e & u \\ e & d' \end{pmatrix} \quad \begin{pmatrix} \nu_\mu & c \\ \mu & s' \end{pmatrix} \quad \begin{pmatrix} \nu_\tau & t \\ \tau & b' \end{pmatrix}$$

$$\left\{ \begin{pmatrix} \nu_e \\ e \end{pmatrix}_L, (\nu_e)_R, e_R^- \right\}, \left\{ \begin{pmatrix} u \\ d' \end{pmatrix}_L, u_R, d_R \right\}$$

✗ All generations interact equally with gauge bosons

✗ Neutral currents:

$$eQ_f \bar{f} \gamma_\mu f \mathcal{A}^\mu, \quad \frac{e}{2s_W c_W} \bar{f} \gamma_\mu (v_f - a_f \gamma_5) f Z^\mu$$

✗ Charged currents:

$$\frac{g}{2\sqrt{2}} \bar{\nu}_\ell \gamma_\mu (1 - \gamma_5) \ell W^{\dagger\mu}, \quad \frac{g}{2\sqrt{2}} \bar{u} \gamma_\mu (1 - \gamma_5) d W^{\dagger\mu}$$

|                           |                        |                          |                           |
|---------------------------|------------------------|--------------------------|---------------------------|
| Quarks                    | $u$<br>up              | $c$<br>charm             | $t$<br>top                |
|                           | $d$<br>down            | $s$<br>strange           | $b$<br>bottom             |
| Leptons                   | $\nu_e$<br>e- Neutrino | $\nu_\mu$<br>μ- Neutrino | $\nu_\tau$<br>τ- Neutrino |
|                           | $e$<br>electron        | $\mu$<br>muon            | $\tau$<br>tau             |
| I      II      III        |                        |                          |                           |
| The Generations of Matter |                        |                          |                           |

# We know

- ✗ P and C broken by weak int. but CP is a symmetry (1 gen)
- ✗ Going from the gauge to mass basis

$$\mathcal{L}_Y^{\text{SM}} = -Y_d^{ij} \bar{Q}_L^i \phi D_R^j - Y_u^{ij} \bar{Q}_L^i \tilde{\phi} U_R^j + \text{h.c.}$$

$$\mathcal{L}_Y^{\text{SM}} = - \left( 1 + \frac{h}{v} \right) [m_d \bar{d}d + m_u \bar{u}u + m_e \bar{e}e]$$

- ✗ With 3 gen trickier - cannot simultaneously diagonalize  $u$  and  $d$  — mixing: CKM matrix
- ✗  $V_{\text{CKM}}$  unitary  $\Rightarrow$  3 real parameters + 1 phase (CPV!)

$$\lambda \quad A \quad \rho \quad \eta$$

# CKM-ology

$\lambda$        $A$        $\rho$        $\eta$

$$V_{CKM} = \begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4)$$

$$\lambda = \sin \theta_C \approx 0.224 \quad A \simeq 0.82 \quad \sqrt{\rho^2 + \eta^2} \approx 0.45$$

- ✗ Fix CKM entries through tree level processes; over constrain by loop-induced ones
- ✗  $V_{CKM}$  unitary  $\Rightarrow$  3 real parameters + 1 phase (CPV!)

# Example : Kaon physics

## Tree level decays

hadronic uncertainty!

$$K \rightarrow \pi \ell \nu \longleftrightarrow \langle \pi | \bar{s} \gamma_\mu u | K \rangle \rightarrow f_{0,+}(q^2)$$

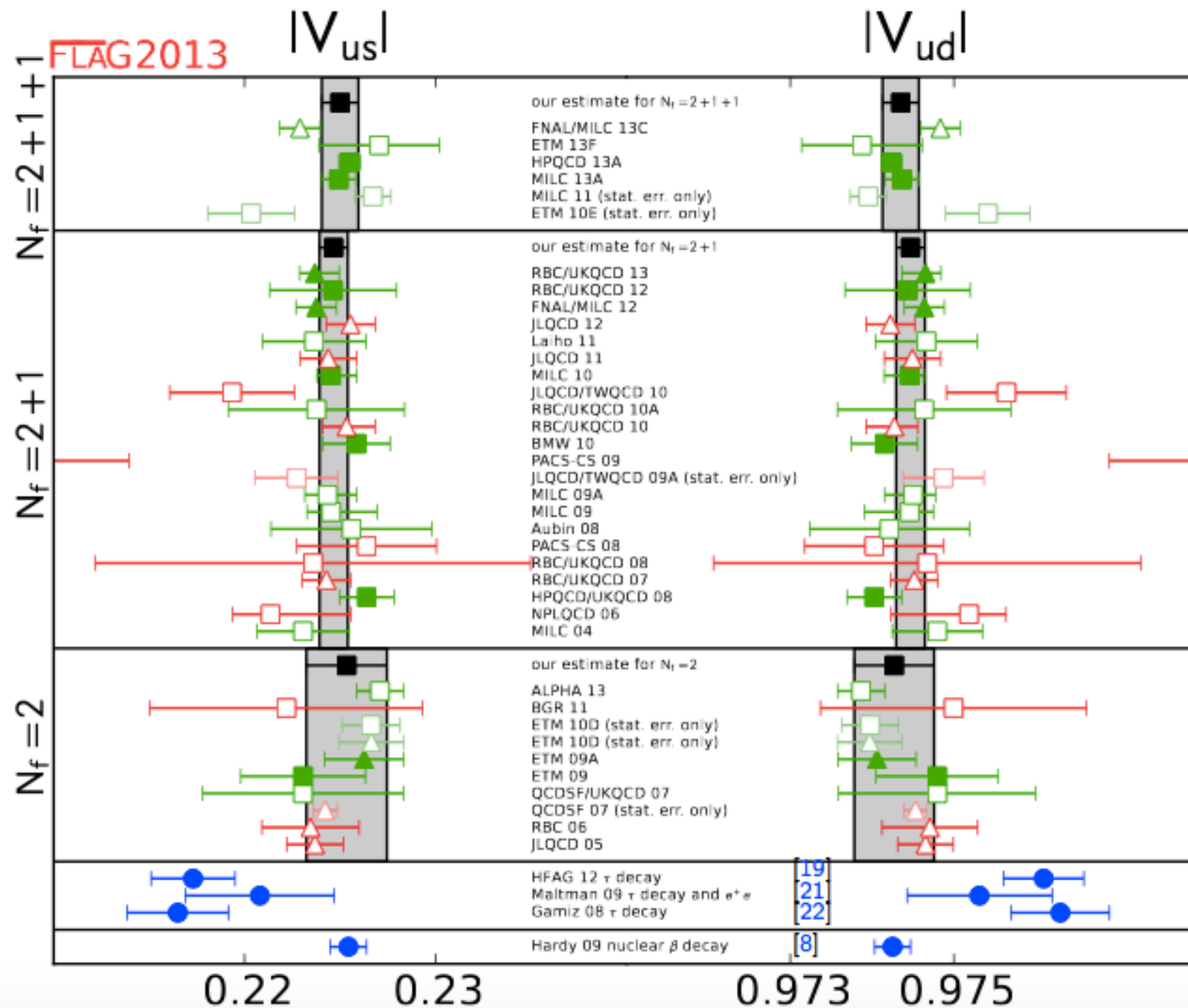
$$K \rightarrow \mu \nu \longleftrightarrow \langle 0 | \bar{s} \gamma_\mu u | K \rangle \rightarrow f_K$$

$f_K / f_\pi$

Nonperturbative QCD - symmetries help (eg. Ademollo-Gatto) but ultimately needs LQCD

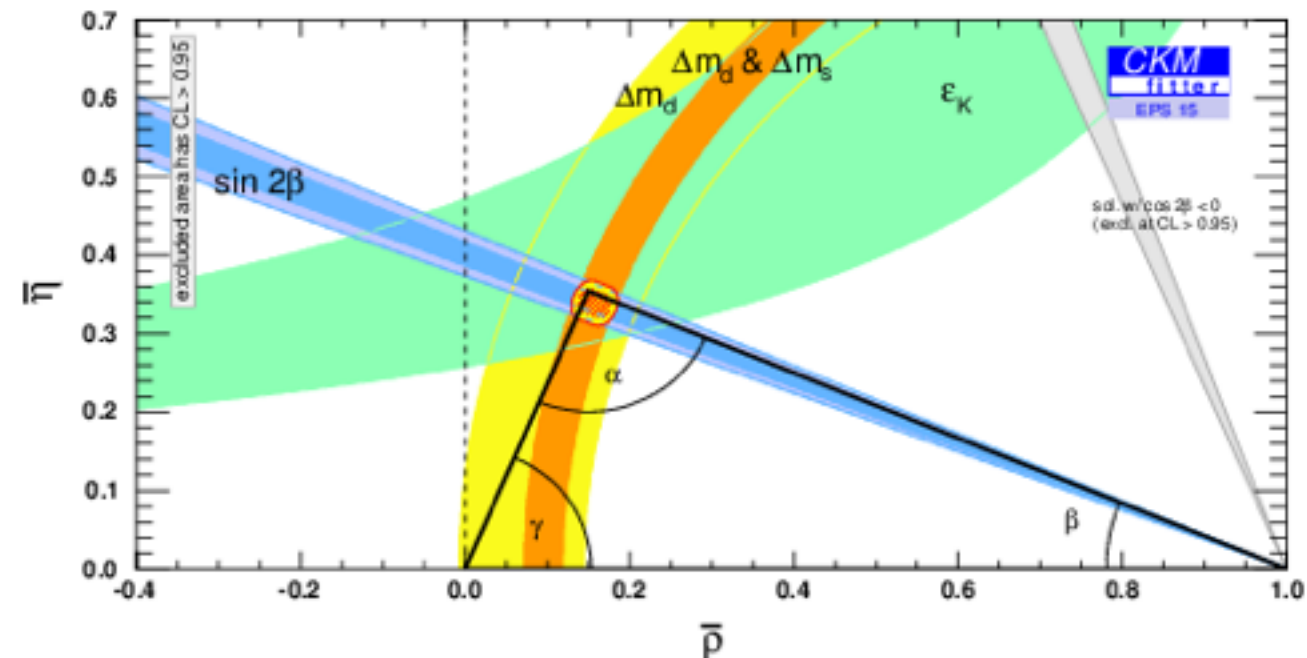
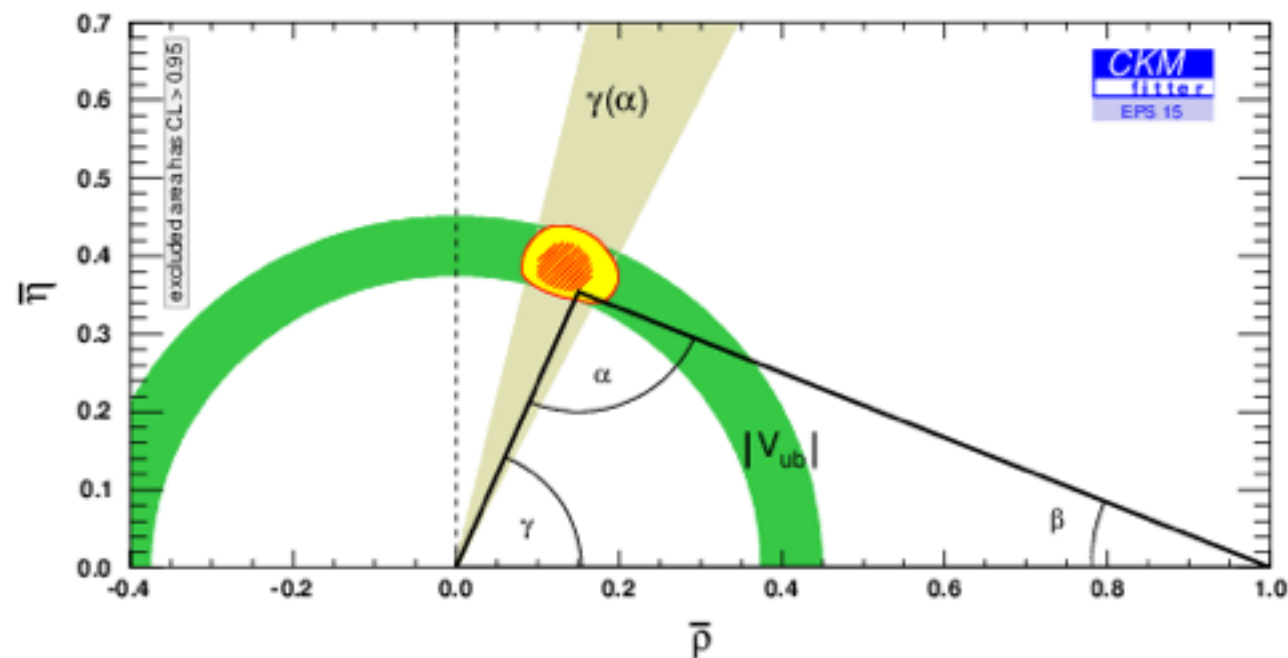
Huge coordinated effort! (cf. FLAG review - new to appear soon!)

# LQCD

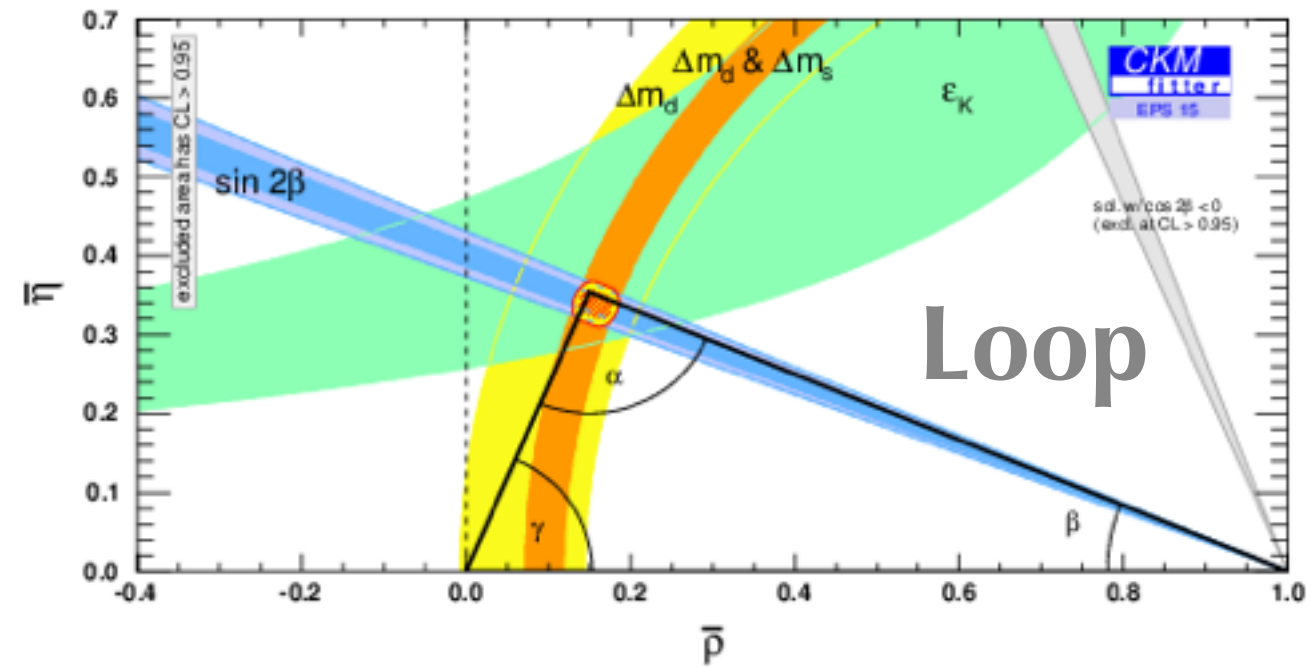
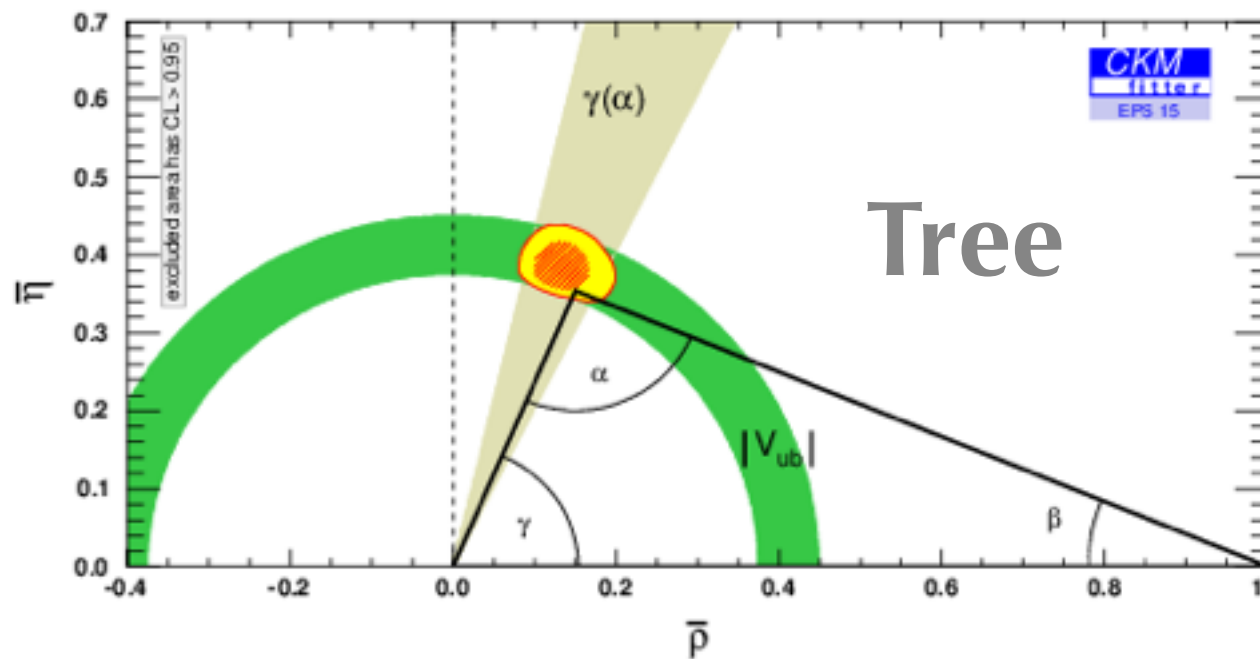


# Experiments

|   |             |             |            |
|---|-------------|-------------|------------|
| ✗ | K-factories | u,d,s       | [NA62]     |
| ✗ | Tau-charm   | $\tau, c$   | [BES III]  |
| ✗ | B-factories | b,c, $\tau$ | [Belle2 !] |
| ✗ | LHC         | t,b,c       |            |
| ✗ | LC          | t,...       |            |
| ✗ | $\nu F$     |             |            |



# CKM



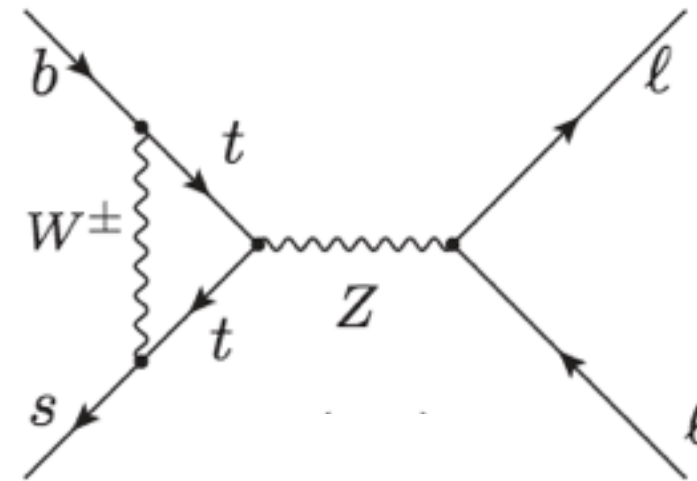
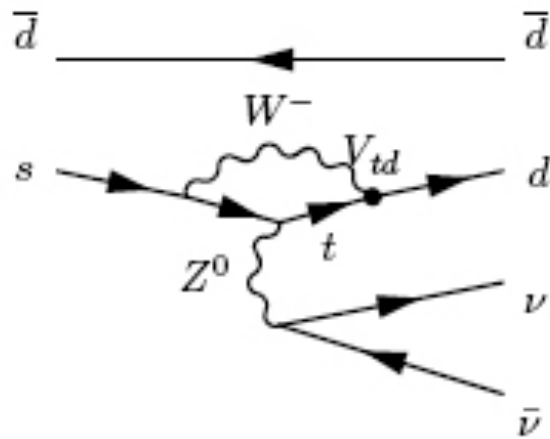
Impressively — TL UT and LP UT agree to less than 10% [Emi says experiment will do do better! Benoit says lattices will too!]

Only tensions in  $V_{ub}$  and  $V_{cb}$  (inclusive Vs. exclusive) but all in all, CKM is very unitary!

2008, Nobel Prize

Strategy:

fix  $V_{ij}$  by tree level processes, then look for NP in FCNC



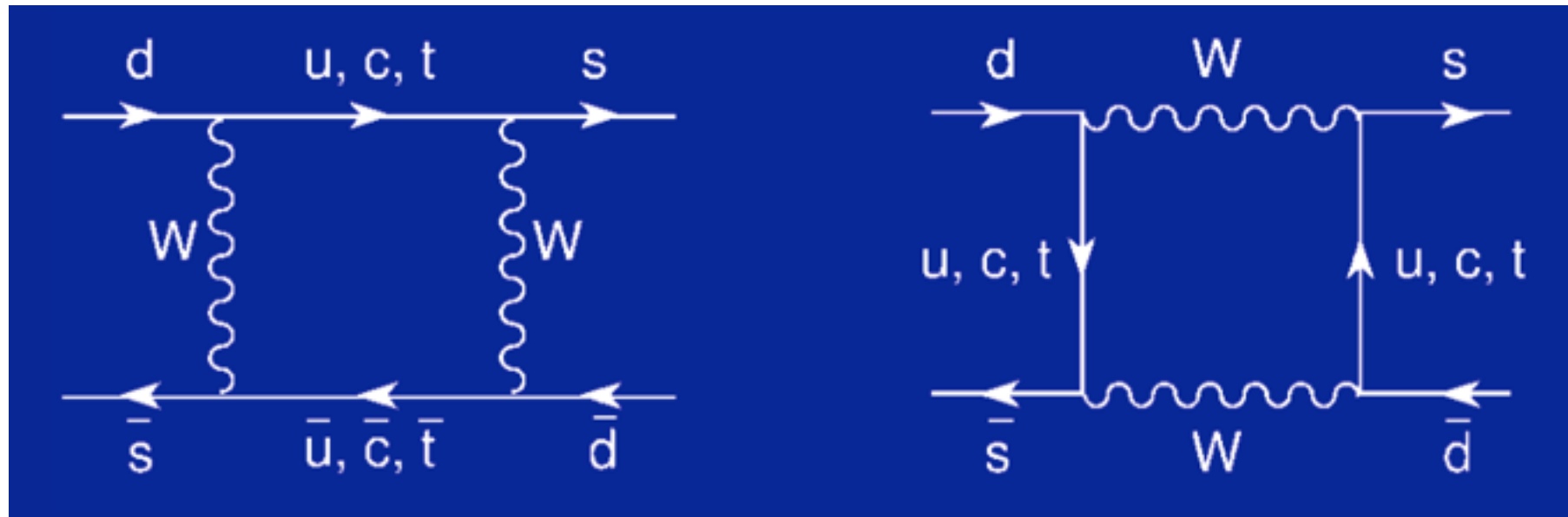
$$\mathcal{B}(B_s \rightarrow \mu^+ \mu^-)_{\text{theo.}} = 3.34 \left( {}^{+13}_{-25} \right) \times 10^{-9} \quad \mathcal{B}(B_s \rightarrow \mu^+ \mu^-)_{\text{LHCb+CMS}} = 2.9(7) \times 10^{-9}$$

| $C_{ij}$                          | 1                   | $V_{ti} V_{tj}^*$   |
|-----------------------------------|---------------------|---------------------|
| $B_s \rightarrow \mu^+ \mu^-$     | $> 10 \text{ TeV}$  | $> 2.5 \text{ TeV}$ |
| $K \rightarrow \pi \nu \bar{\nu}$ | $> 100 \text{ TeV}$ | $> 1.8 \text{ TeV}$ |

$$O = \frac{1}{\Lambda^2} C_{ij} \bar{Q}_i \gamma^\mu Q_j H^\dagger D_\mu H$$

Strategy:

fix  $V_{ij}$  by tree level processes, then look for NP in FCNC



$$O = \frac{1}{\Lambda^2} C'_{ij} \bar{Q}_i \gamma^\mu Q_j \bar{Q}_i \gamma_\mu Q_j$$

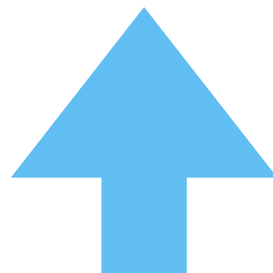
| $C'_{ij}$             | 1                               | $ V_{ti} V_{tj}^* ^2$ |
|-----------------------|---------------------------------|-----------------------|
| $K^0 - \bar{K}^0$     | $> 2 \times 10^4 \text{ TeV}$   | $> 8 \text{ TeV}$     |
| $B^0 - \bar{B}^0$     | $> 0.5 \times 10^4 \text{ TeV}$ | $> 5 \text{ TeV}$     |
| $B_s^0 - \bar{B}_s^0$ | $> 0.1 \times 10^4 \text{ TeV}$ | $> 5 \text{ TeV}$     |

# Flavor puzzle

| $C_{ij}$                        | 1                   | $V_{ti}V_{tj}^*$    |
|---------------------------------|---------------------|---------------------|
| $B_s \rightarrow \mu^+\mu^-$    | $> 10 \text{ TeV}$  | $> 2.5 \text{ TeV}$ |
| $K \rightarrow \pi\nu\bar{\nu}$ | $> 100 \text{ TeV}$ | $> 1.8 \text{ TeV}$ |

| $C'_{ij}$             | 1                               | $ V_{ti}V_{tj}^* ^2$ |
|-----------------------|---------------------------------|----------------------|
| $K^0 - \bar{K}^0$     | $> 2 \times 10^4 \text{ TeV}$   | $> 8 \text{ TeV}$    |
| $B^0 - \bar{B}^0$     | $> 0.5 \times 10^4 \text{ TeV}$ | $> 5 \text{ TeV}$    |
| $B_s^0 - \bar{B}_s^0$ | $> 0.1 \times 10^4 \text{ TeV}$ | $> 5 \text{ TeV}$    |

- For natural  $C \sim O(1)$ , NP scale is huge
- Need lots of fine tuning to reduce NP scale to  $O(1\text{TeV})$  as needed to mend the hierarchy problem
- Way out: NP is (almost) aligned with the SM
- MFV



# MFV

To protect quark flavor mixing BSM, assume flavor symmetry is the one present in the limit of vanishing Yukawa's,  $U(3)^3$ , and that two quark Yukawa,  $Y_u$  and  $Y_d$ , are the only symmetry breaking and CP violating terms

$$\mathcal{L}_Y^{\text{SM}} = -Y_d^{ij} \bar{Q}_L^i \phi D_R^j - Y_u^{ij} \bar{Q}_L^i \tilde{\phi} U_R^j + \text{h.c.}$$

Promote  $Y_u$  and  $Y_d$  to non-dynamical fields. Higher dim operators made of SM fields and  $Y_{ud}$ .

Eigenvalues of  $Y_{ud}$  small except for top, off-diagonal elements suppressed  $\Rightarrow [Y_u(Y_u)^\dagger]_{i \neq j}^n \approx y_t^{2n} V_{ti}^* V_{tj}$

D'Ambrosio et al. NPB645(2002)155

Marcello, Smith, NPB 817(2009)1

...

Isidori, Straub EPJC72(2012)2103

Smith HDR-2015

# Questions and progress

- ✗ Why is there flavor? Why families? Why 3?
- ✗ Why such a strong hierarchy?
- ✗ Why quark mixing is small (and lepton mixing is large)?
- ✗ Why is there quark alignment?
- ✗ How to solve strong CP-problem? [Peccei-Quinn elegant solution, but where are axions?]
- ✗ Need CPV in quark and lepton sector for BAU
- ✗ Does the scalar sector play a non-trivial role in the questions of flavor?
- ✗ Branco, Grimus and Lavoura figured out a symmetry which imposed on SM+2HDM provides a structure of Yukawas such that there is no FCNC at tree-level and their strength controlled by CKM (!)

# $b \rightarrow s$ anomalies

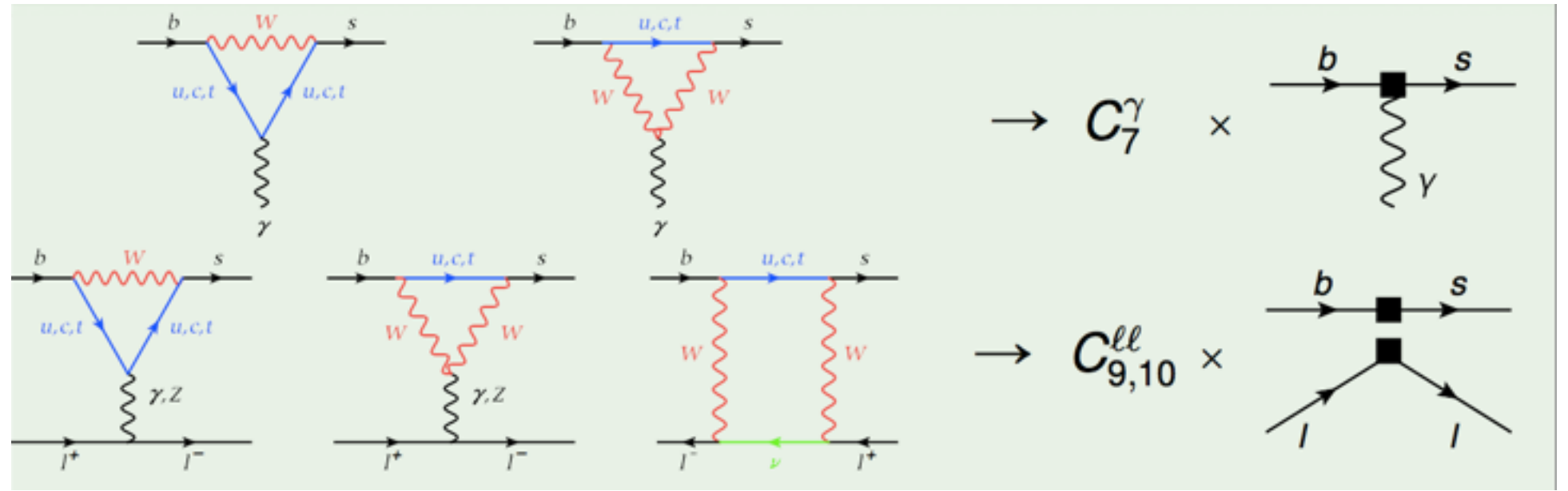
$2.6\sigma$  away from SM in  $R_K$

$$R_K = \frac{\mathcal{B}(B \rightarrow K \mu^+ \mu^-)_{q^2 \in [1,6] \text{ GeV}^2}}{\mathcal{B}(B \rightarrow K e^+ e^-)_{q^2 \in [1,6] \text{ GeV}^2}} = 0.745 \left( {}^{+90}_{-74} \right) (36)_{\text{LHCb}} \neq 1 + \mathcal{O}(m_\mu^2/m_b^2)$$

Can we claim to have seen LFUV?

LHCb will soon improve this number and also measure  $R_{K^*}$

# Basics



$$\mathcal{H}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}}V_{tb}V_{ts}^* \left[ C_1\mathcal{O}_1 + C_2\mathcal{O}_2 + \sum_{i=3}^6 C_i\mathcal{O}_i + \sum_{i=7,8,9,10,P,S} (C_i\mathcal{O}_i + C'_i\mathcal{O}'_i) \right]$$

$$\mathcal{O}_7 = \frac{e}{g^2} m_b (\bar{s} \sigma_{\mu\nu} P_R b) F^{\mu\nu},$$

$$\mathcal{O}'_7 = \frac{e}{g^2} m_b (\bar{s} \sigma_{\mu\nu} P_L b) F^{\mu\nu},$$

$$\mathcal{O}_8 = \frac{1}{g} m_b (\bar{s} \sigma_{\mu\nu} T^a P_R b) G^{\mu\nu a},$$

$$\mathcal{O}'_8 = \frac{1}{g} m_b (\bar{s} \sigma_{\mu\nu} T^a P_L b) G^{\mu\nu a},$$

$$\mathcal{O}_9 = \frac{e^2}{g^2} (\bar{s} \gamma_\mu P_L b) (\bar{\mu} \gamma^\mu \mu),$$

$$\mathcal{O}'_9 = \frac{e^2}{g^2} (\bar{s} \gamma_\mu P_R b) (\bar{\mu} \gamma^\mu \mu),$$

$$\mathcal{O}_{10} = \frac{e^2}{g^2} (\bar{s} \gamma_\mu P_L b) (\bar{\mu} \gamma^\mu \gamma_5 \mu),$$

$$\mathcal{O}'_{10} = \frac{e^2}{g^2} (\bar{s} \gamma_\mu P_R b) (\bar{\mu} \gamma^\mu \gamma_5 \mu),$$

$$C_7^{\text{eff}} = \frac{4\pi}{\alpha_s} C_7 - \frac{1}{3} C_3 - \frac{4}{9} C_4 - \frac{20}{3} C_5 - \frac{80}{9} C_6$$

$$C_8^{\text{eff}} = \frac{4\pi}{\alpha_s} C_8 + C_3 - \frac{1}{6} C_4 + 20C_5 - \frac{10}{3} C_6$$

$$C_9^{\text{eff}} = \frac{4\pi}{\alpha_s} C_9 + Y(q^2) \quad \text{!!!}$$

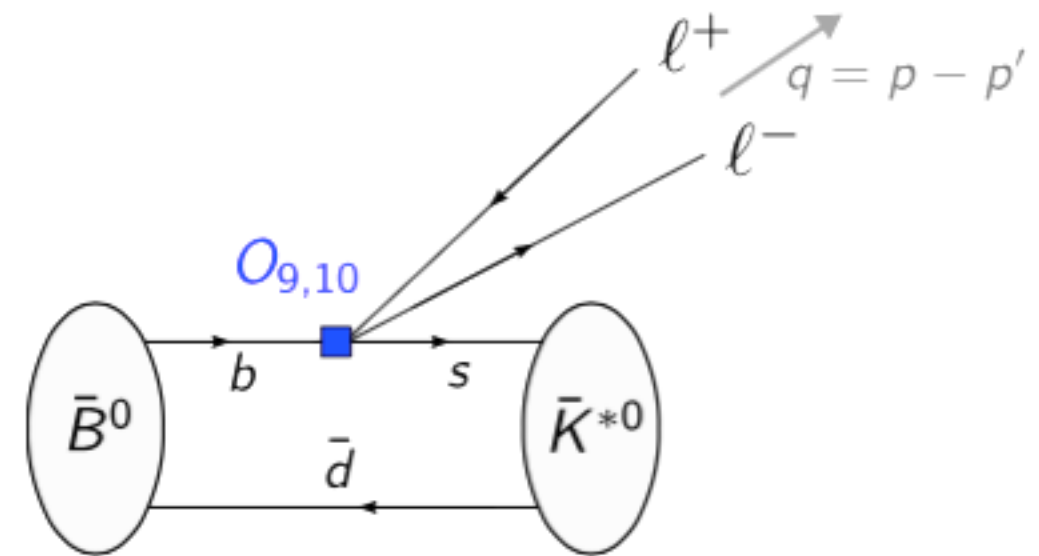
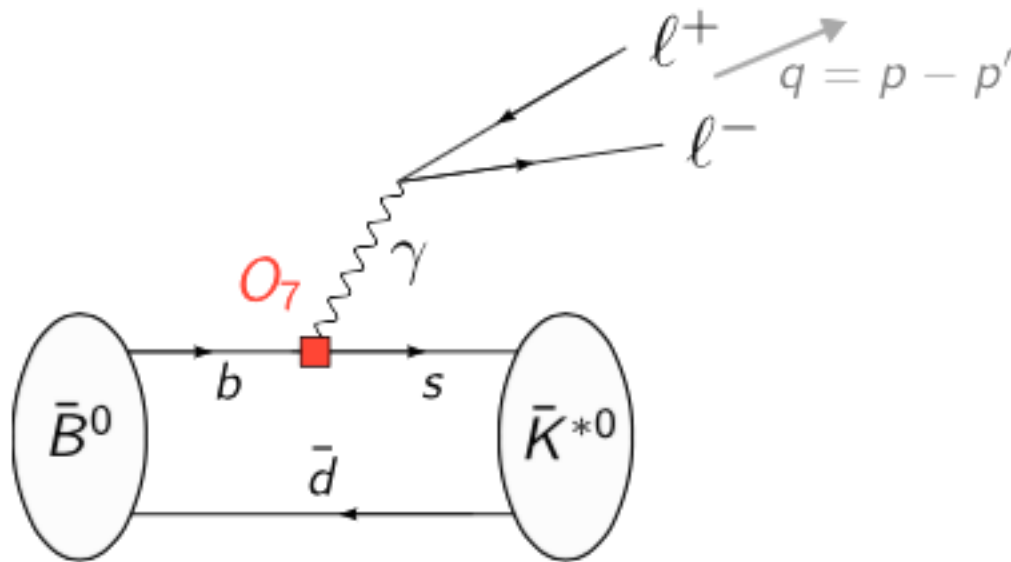
$$Y(q^2) = \frac{4}{3} C_3 + \frac{64}{9} C_5 + \frac{64}{27} C_6 - \frac{1}{2} h(q^2, 0) \left( C_3 + \frac{4}{3} C_4 + 16C_5 + \frac{64}{3} C_6 \right) \\ + h(q^2, m_c) \left( \frac{4}{3} C_1 + C_2 + 6C_3 + 60C_5 \right) - \frac{1}{2} h(q^2, m_b) \left( 7C_3 + \frac{4}{3} C_4 + 76C_5 + \frac{64}{3} C_6 \right)$$

Very slowly varying functions of  $q^2$

$$\mathcal{M} = \frac{G_F \alpha}{\sqrt{2}\pi} V_{tb} V_{ts}^* \left[ (\mathcal{A}_\mu + \mathcal{T}_\mu) \bar{u}_\ell \gamma^\mu v_\ell + \mathcal{B}_\mu \bar{u}_\ell \gamma^\mu \gamma_5 v_\ell \right],$$

$$\mathcal{A}_\mu = -\frac{2m_b}{q^2} q^\nu \mathcal{C}_7 \langle K^* | \bar{s} i\sigma_{\mu\nu} \frac{1+\gamma_5}{2} b | B \rangle + \mathcal{C}_9 \langle K^* | \bar{s} \gamma_\mu \frac{1-\gamma_5}{2} b | B \rangle$$

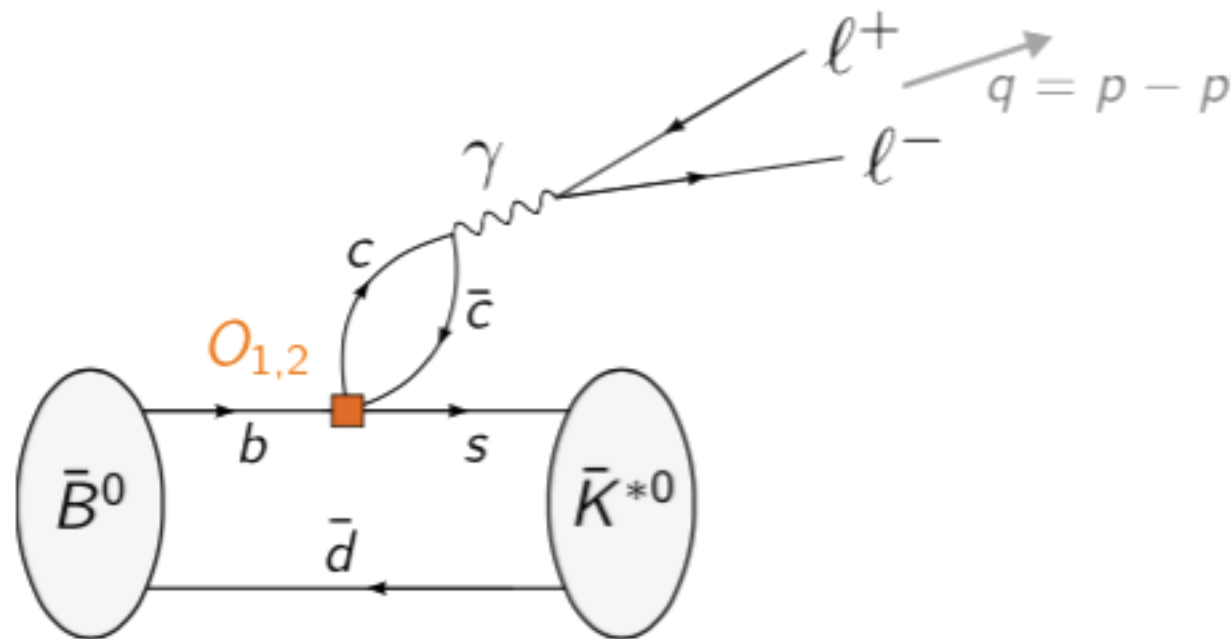
$$\mathcal{B}_\mu = \mathcal{C}_{10} \langle K^* | \bar{s} \gamma_\mu \frac{1-\gamma_5}{2} b | B \rangle$$



Can be and are computed on the lattice

$$\mathcal{M} = \frac{G_F \alpha}{\sqrt{2}\pi} V_{tb} V_{ts}^* \left[ (\mathcal{A}_\mu + \mathcal{T}_\mu) \bar{u}_\ell \gamma^\mu v_\ell + \mathcal{B}_\mu \bar{u}_\ell \gamma^\mu \gamma_5 v_\ell \right],$$

$$\mathcal{T}_\mu = \frac{-16i\pi^2}{q^2} \sum_{i=1\dots 6;8} C_i \int d^4x \, e^{iq \cdot x} \langle K^* | T \, O_i(0) j_\mu(x) | B \rangle$$



Cannot be computed on the lattice  
 - work either at very low or very high  $q^2$

# $b \rightarrow s$ anomalies

$2.6\sigma$  away from SM in  $R_K$

$$R_K = \frac{\mathcal{B}(B \rightarrow K \mu^+ \mu^-)_{q^2 \in [1,6] \text{ GeV}^2}}{\mathcal{B}(B \rightarrow K e^+ e^-)_{q^2 \in [1,6] \text{ GeV}^2}} = 0.745 \left( {}^{+90}_{-74} \right) (36)_{\text{LHCb}} \neq 1 + \mathcal{O}(m_\mu^2/m_b^2)$$

- Better still: hadronic uncertainties cancel
- Need another experiment (Belle2!)
- Scalar operators can't explain. Vector ones could: NP in  $C^{\mu\mu}_9$  but not  $C^{ee}_9$ . Similarly with  $C_{9'}$ ,  $C_{10}$  and  $C_{10'}$
- Leptoquark models explicitly accommodate this feature
- Models with  $Z'$  [gauging a  $U(1)$  flavor subgroup, or merely introducing vector-like couplings to muons]

# $b \rightarrow s$ anomalies

$2.6\sigma$  away from SM in  $R_K$

## Leptoquark models

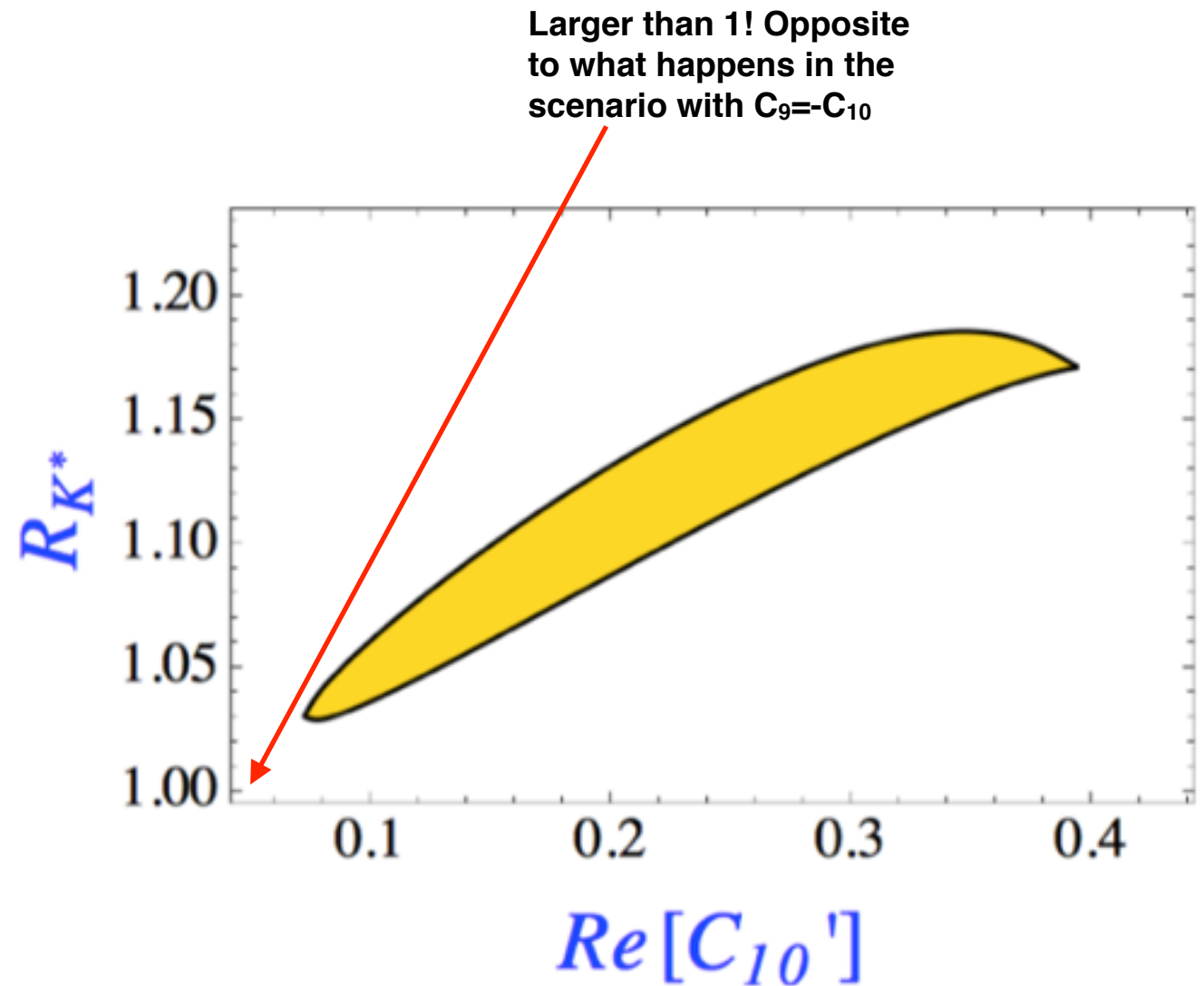
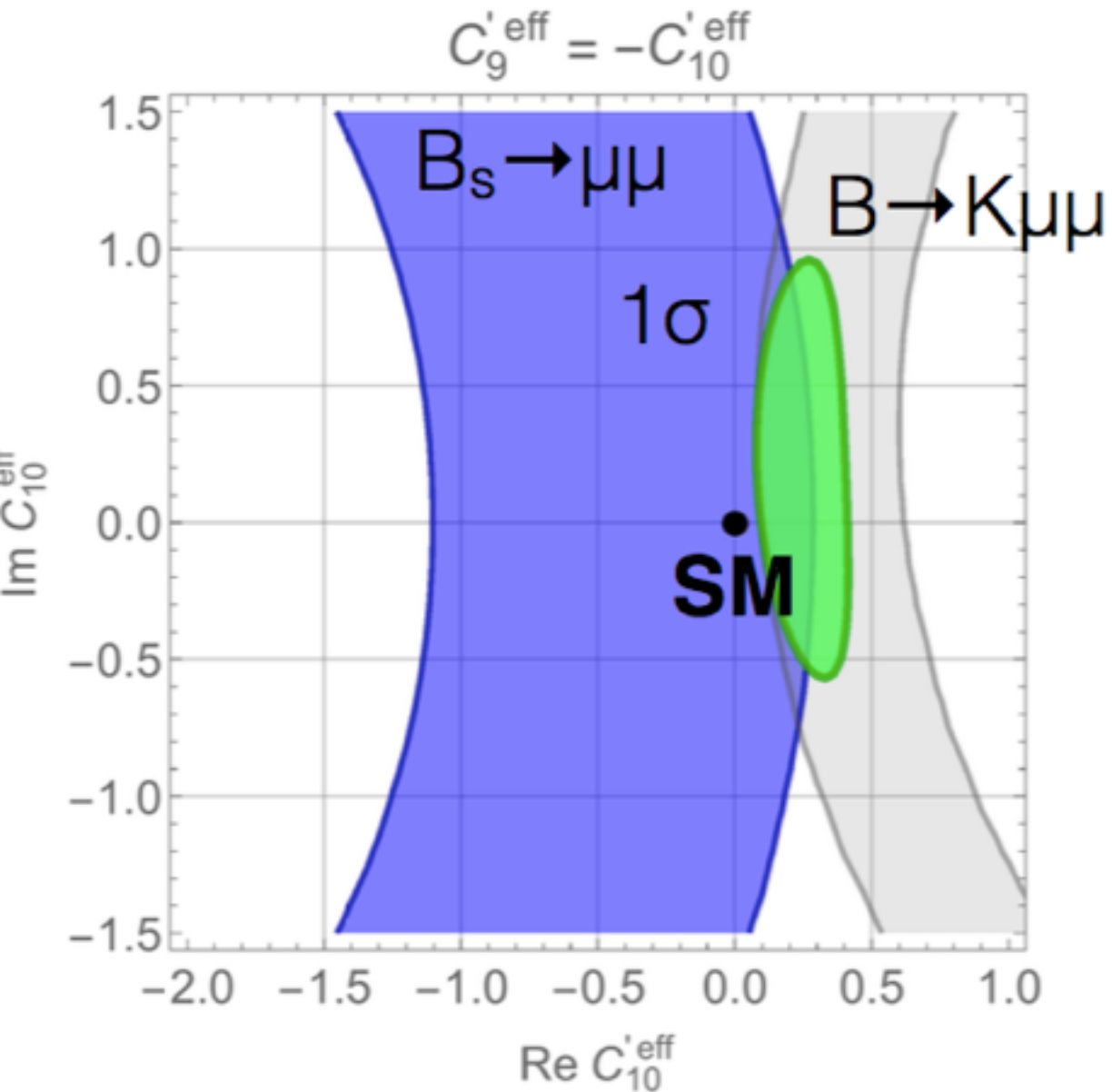
Representations of scalar LQs under  $SU(3) \otimes SU(2) \otimes U(1)^*$

|                      |                                   |                                  |                                  |
|----------------------|-----------------------------------|----------------------------------|----------------------------------|
| $(3, 2)_{7/6}$       | Increases $B \rightarrow K\mu\mu$ | $\bar{Q}e_R$                     |                                  |
| $(3, 2)_{1/6}$       | Decreases $B \rightarrow K\mu\mu$ | $\bar{L}d_R$                     | $\Delta(3, 2)_{1/6}$             |
| $(\bar{3}, 3)_{1/3}$ | Proton destabilizing              | $\bar{Q}^C i\tau_2 \vec{\tau} L$ | $\bar{Q}^C i\tau_2 \vec{\tau} Q$ |
| $(\bar{3}, 1)_{4/3}$ | Proton destabilizing              | $\bar{d}_R^C \ell_R$             | $\bar{u}_R^C u_R$                |

$$C'_{10} = -C'_9 = \frac{-\pi}{2\sqrt{2}G_F V_{tb} V_{ts}^* \alpha} \frac{Y_{\mu b} Y_{\mu s}^*}{m_\Delta^2}$$

# $b \rightarrow s$ anomalies

$2.6\sigma$  away from SM in  $R_K$

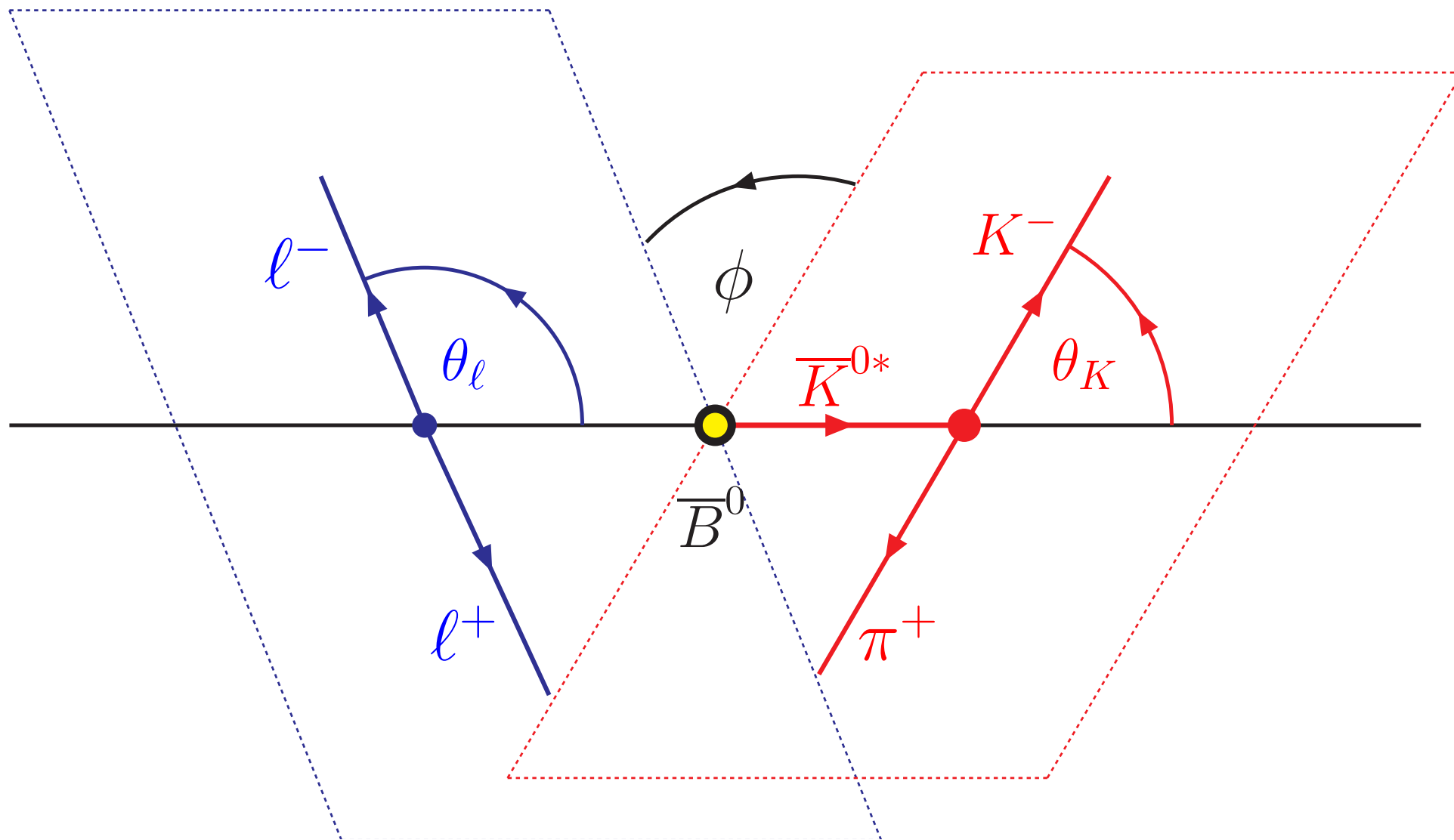


$$R_K(C'_{10}) = 1.001(1) - 0.46 \text{ Re}[C'_{10}] - 0.094(3) \text{ Im}[C'_{10}] + 0.057(1) |C'_{10}|^2$$

Remaining form factor uncertainties

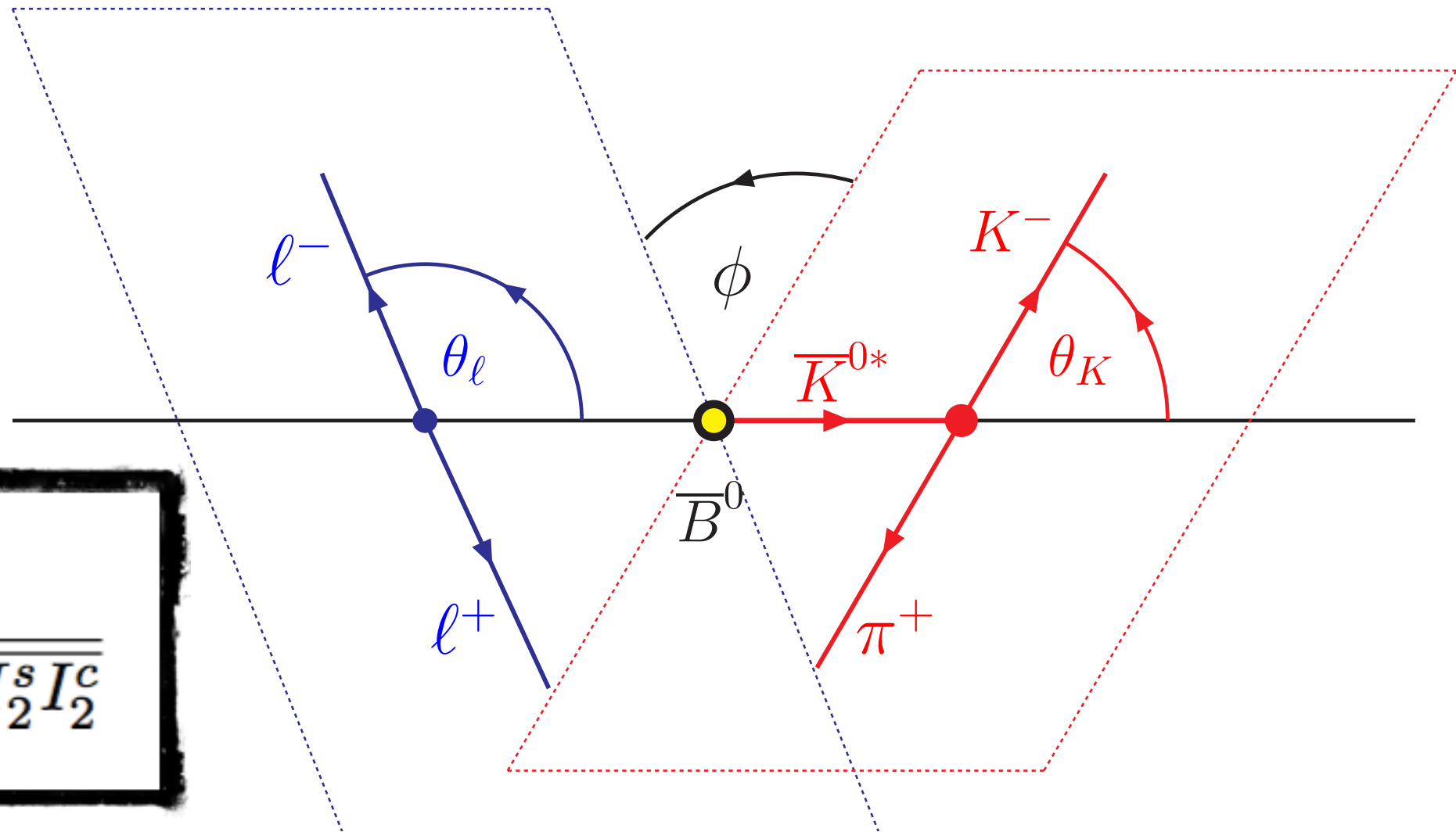
# Better sensitivity to NP:

$$B \rightarrow K^* \ell^+ \ell^-$$



$$\frac{d^4\Gamma(\bar{B}^0 \rightarrow \bar{K}^{*0} \ell^+ \ell^-)}{dq^2 d\cos\theta_\ell d\cos\theta_K d\phi} = \frac{9}{32\pi} I(q^2, \theta_\ell, \theta_K, \phi)$$

# Full decay distribution



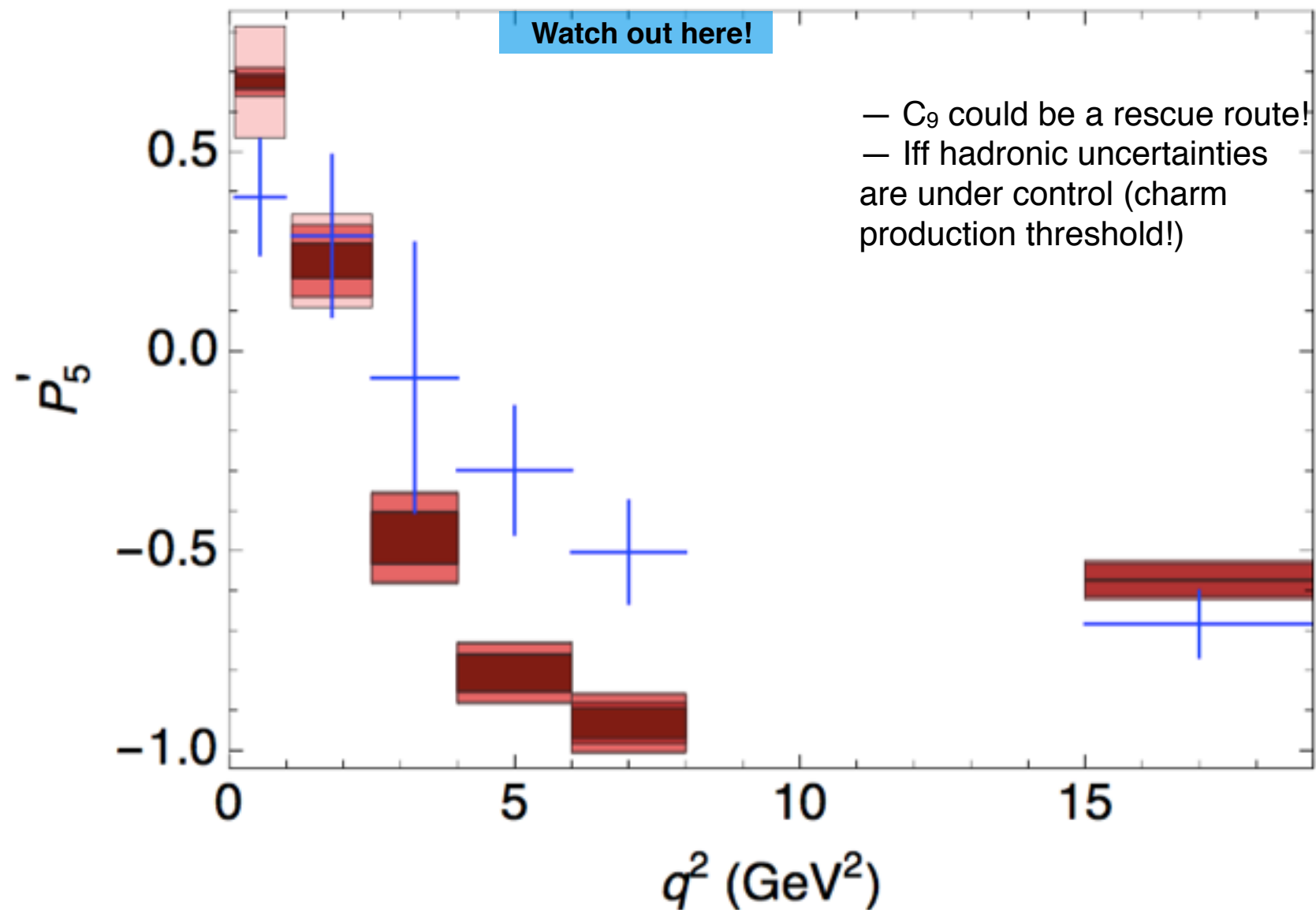
$$P'_5 = \frac{I_5}{\sqrt{-4I_2^s I_2^c}}$$



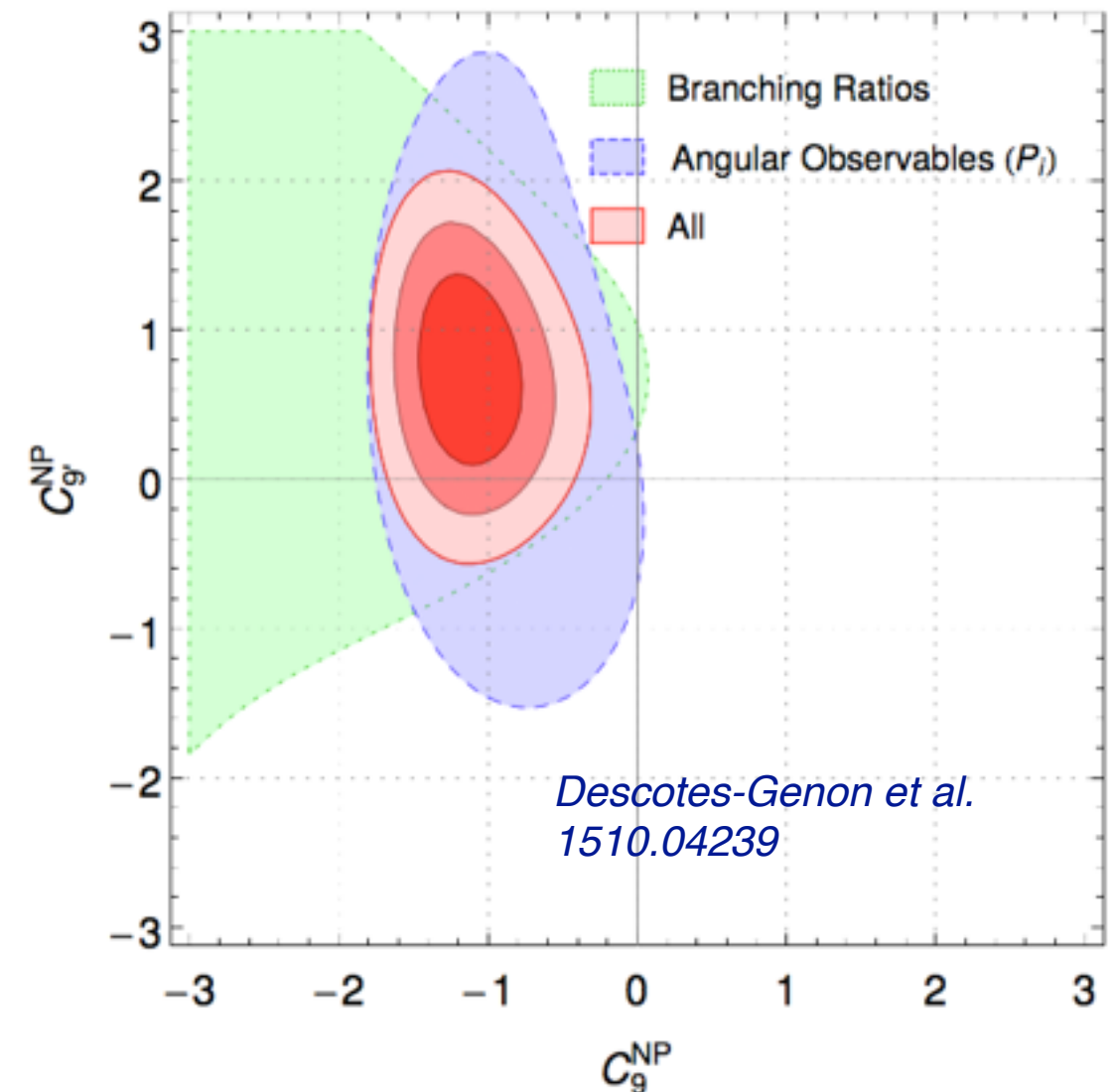
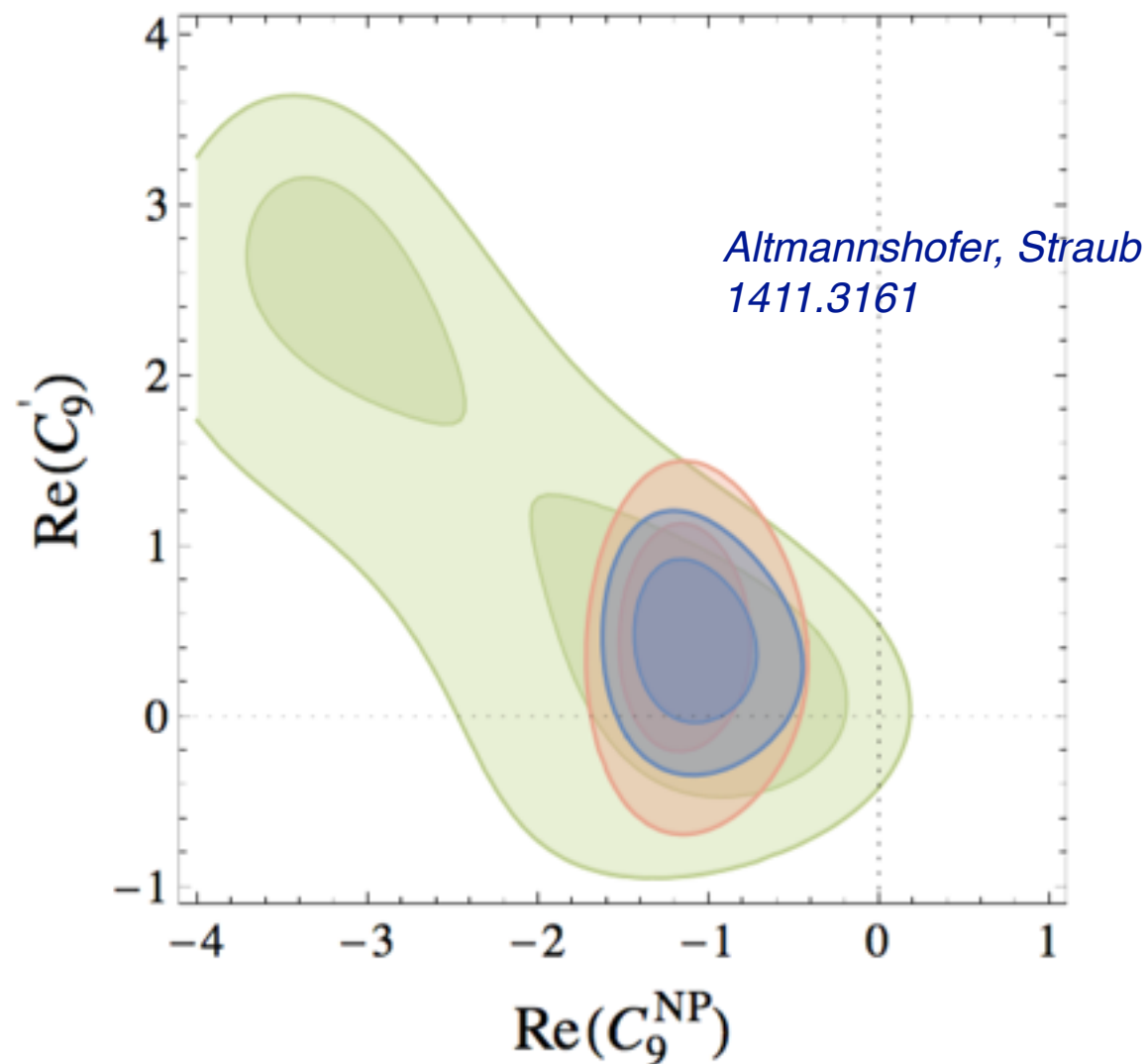
$$\begin{aligned} I(q^2, \theta_\ell, \theta_K, \phi) = & I_1^s(q^2) \sin^2 \theta_K + I_1^c(q^2) \cos^2 \theta_K + [I_2^s(q^2) \sin^2 \theta_K + I_2^c(q^2) \cos^2 \theta_K] \cos 2\theta_\ell \\ & + I_3(q^2) \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\phi + I_4(q^2) \sin 2\theta_K \sin 2\theta_\ell \cos \phi \\ & + I_5(q^2) \sin 2\theta_K \sin \theta_\ell \cos \phi \\ & + [I_6^s(q^2) \sin^2 \theta_K + I_6^c(q^2) \cos^2 \theta_K] \cos \theta_\ell + I_7(q^2) \sin 2\theta_K \sin \theta_\ell \sin \phi \\ & + I_8(q^2) \sin 2\theta_K \sin 2\theta_\ell \sin \phi + I_9(q^2) \sin^2 \theta_K \sin^2 \theta_\ell \sin 2\phi \end{aligned}$$

# $b \rightarrow s$ anomalies

2-3 $\sigma$  deviation from SM [esp.  $P_5'$ ]



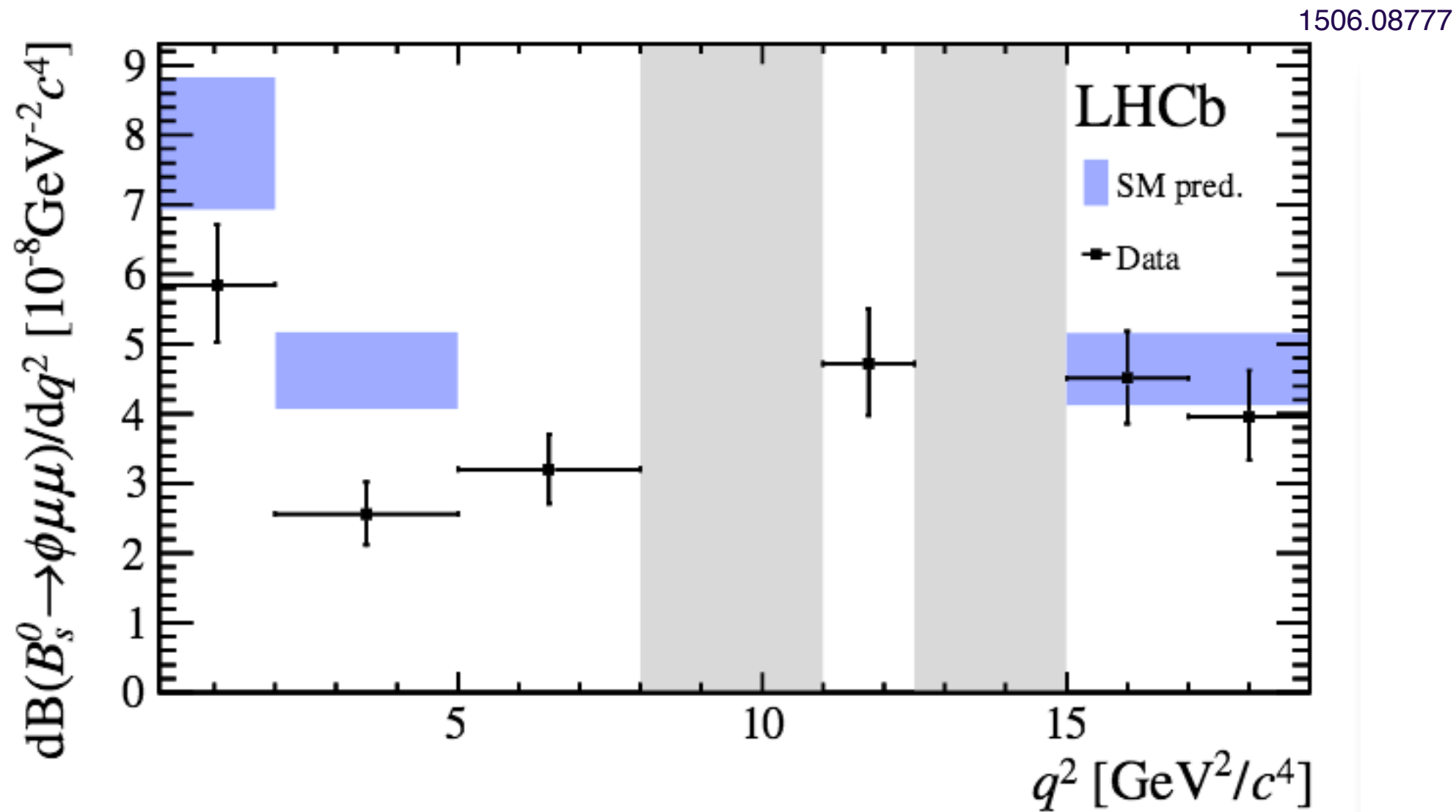
# $b \rightarrow s$ anomalies



- Theory errors still subject to controversies.
- Some quantities are more sensitive to hadronic uncertainties than others (maybe sticking to the clean observables only?)
- Ciuchini et al [1512.07157] claim that the whole discrepancy can be absorbed into (unknown) power corrections due to charm loops.

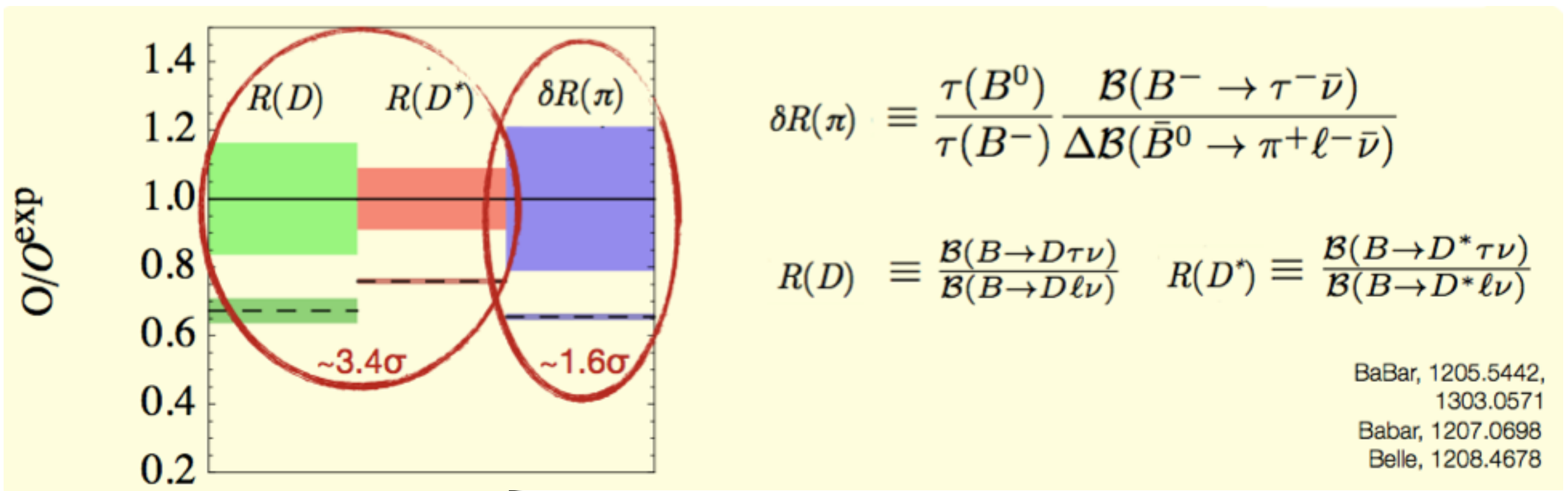
# $b \rightarrow s$ anomalies

$3.1\sigma$  in  $B \rightarrow \phi \mu \mu$  below SM at low  $q^2$



Theory error - subject to controversies... If OK, then NP in  $C_9$  could fill the gap between experiment and SM.

# $R_D$ $R_{D^*}$



LHCb [1506.08614]

$$\mathcal{R}(D^*) = 0.336 \pm 0.027 (\text{stat}) \pm 0.030 (\text{syst})$$

$$\mathcal{R}(D^*)^{\text{theo.}} = 0.252(3)$$

Fajfer et al [1203.2654]

- Theory error, one form factor completely unknown but HQET helps (error that entails?) LQCD may help!

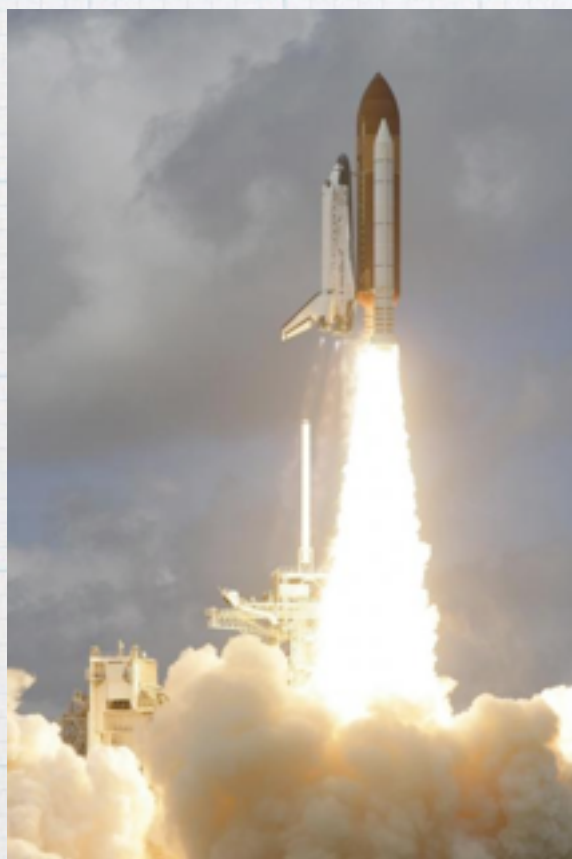
# Model builders stuck

- Cannot accommodate  $b \rightarrow s$  anomalies,  $R_{D^*}$ ,  $R_D$  and the CMS hint,  $B(h \rightarrow \tau\mu) = 0.83(38)\%$
- Either some information is not yet reliable, or the model turns be baroque
- Attempt 1:  $U(1)$  gauge group  $Q(e)=0$ ,  $Q(\mu)=1$ ,  $Q(\tau)=-1 \Rightarrow Z'$  with vector like quarks (charged under  $L_\tau-L_\mu$ ) [1403.1269]; second doublet [in 2HDM] breaks  $L_\tau-L_\mu$ , inducing the LFV [1501.00993]
- Attempt 2: 3HDM with gauged B-L charges [1503.03477]
- Attempt 3: Weakly break  $U(2)^5$  plus vector LQ - nonrenormalizable [1503.03477]
- **Attempt 4:** Single scalar LQ with  $(3,1,-1/3)$  and mass around 1 TeV:
  - passes  $R_{D^*}$ ,  $R_D$  and  $R_K$  [1511.01900]

# Flavor Physics in the LHC era

LHC @ high  $p_T$

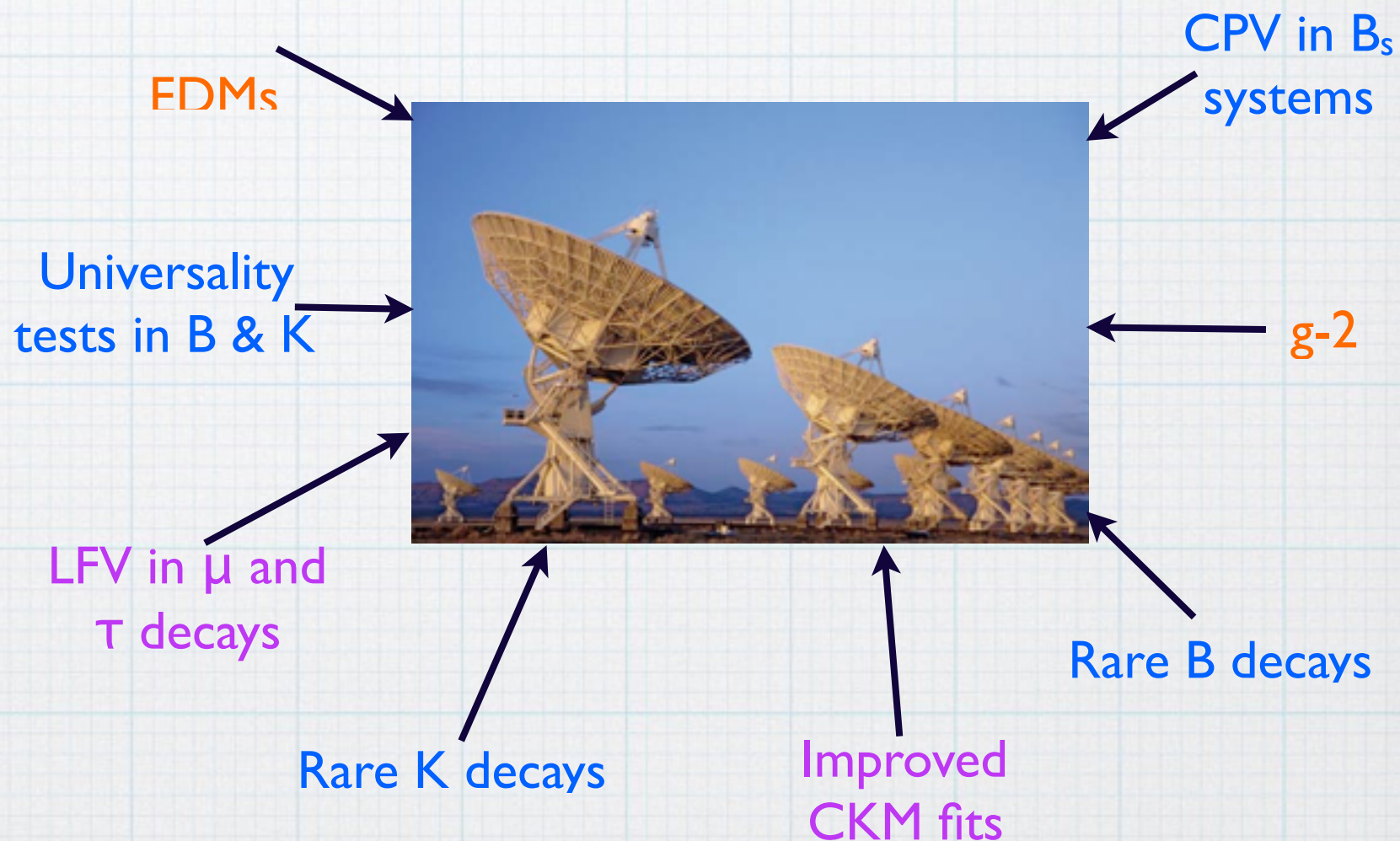
Unique effort toward the high energy frontier



to determine the energy scale of New Physics

Flavor Physics

Collective effort

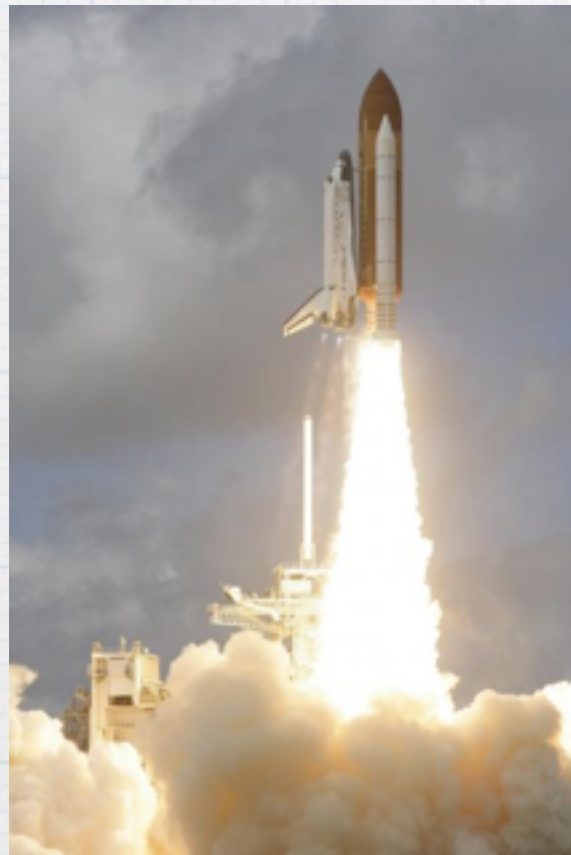


to determine the flavor structure of New Physics

# Two complementary ways to study NP

LHC @ high  $p_T$

Unique effort toward the high energy frontier

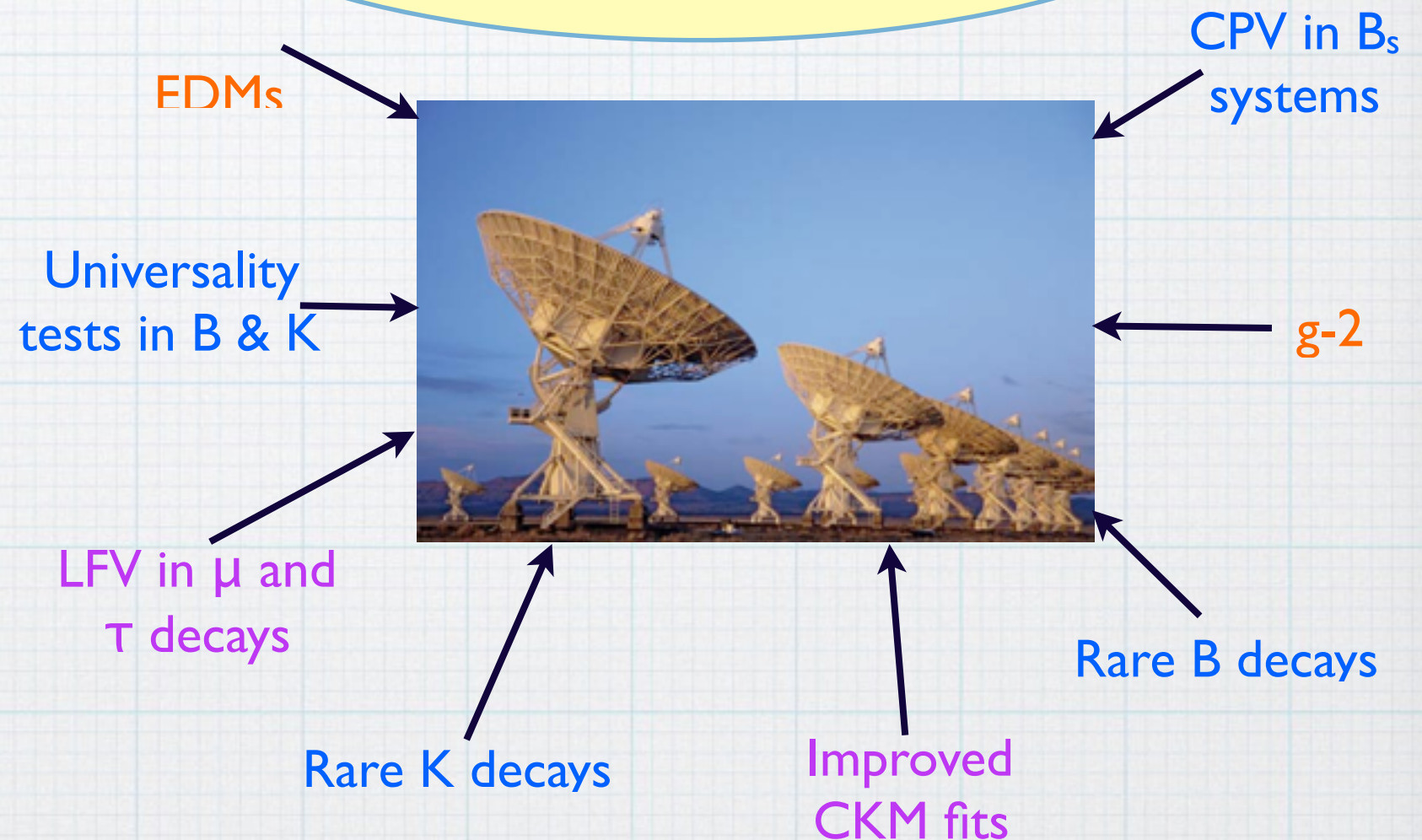


to determine the energy scale of New Physics

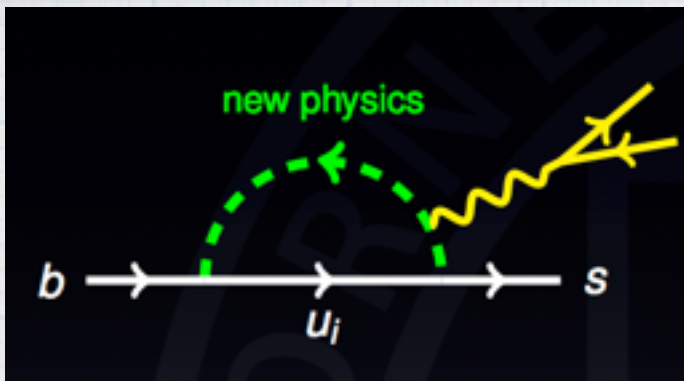
## Flavor Physics

Super-flavor machines

Super-Belle (YES) & Super-B (???)



to determine the flavor structure of New Physics



# Where to look?

- ☑ Check on the scalar form factor in  $Kl3$  [ RH current in Higgsless theories or binding the charged Higgs (for large  $\tan\beta$ )  $\rightarrow$  Similar possible with  $B \rightarrow D\tau\nu$ , D-decays - rich ground for zero-tests]
- ☑  $b \rightarrow s \ell\ell$  [ Many asymmetries (esp FB), sensitive to either SUSY or XD; Exclusive  $B \rightarrow K^{(*)}\ell\ell$ ,  $D \rightarrow K/\pi\ell\ell$  - rich ground for zero-tests]
- ☑  $K \rightarrow \pi\nu\nu$  [very clean-up to the isospin & higher order ew corrections;  $B \rightarrow K\nu\nu$  interesting]
- ☑  $B_s \rightarrow \mu\mu$  [in MSSM its amplit. enhanced by  $\tan^3\beta$ ;  $B_s \rightarrow \mu\mu/B_d \rightarrow \mu\mu$  ?]
- ☑  $\gamma$  from  $B \rightarrow DK$  [isospin relations among  $b \rightarrow c\bar{u}s$   $b \rightarrow u\bar{c}s$ ,]
- ☑  $B_{s,d} \rightarrow KK, K\pi, \pi\pi$  [explore LHCb, Tevatron - discrepancies?!]
- ☑  $\tau \rightarrow \mu\gamma, \mu \rightarrow e\gamma$  [helps CPV in  $\nu$ -sector]
- ☑  $B_c \rightarrow \text{all}$  [and many other consistency checks]

2009

2009

# QCD tests/constraints

- ☑ Spectroscopy [scalars, tetraquark, hybrid...]
- ☑ Non-leptonic decays [ FSI from the data?! Lattice cannot help! Theory?]
- ☑ Range of validity of effective theories [ChPT, NRQCD, HQET; can SCET be made useful?]
- ☑ QCD factorization
- ☑ Duality [quark hadron duality works well until?]
- ☑ Flavor singlet sector [terra incognita!  $\eta'$  - unknown physics]
- ☑ Always check on the lattice results!
- ☑ Soft photons are trouble [eg.  $B \rightarrow \tau \nu \gamma$  OK, but  $B \rightarrow \mu \nu \gamma$  is not]

2009

# MESSAGES

- ☑ New Physics - Hierarchy Vs Flavor problem [towards solution(s)]
- ☑ QCD is NOT solved yet! [Lattice QCD is the only tool for the ab-initio calculations → indispensable (*merci à l'IN2P3*), but there are far more things you cannot do with Lattices than those that you can; NEED an analytic understanding of non-perturbative QCD]
- ☑ LHC is blunt-machine [NP signal will hopefully be discovered but to understand what NP, need cleaner experiment → Super-Belle, Super-B ; watch out for Super-B workshops]
- ☑ Check on the reported discrepancies [but believe anything under  $3\sigma$  ; many new modes will be explored, current results improved]
- ☑ FLAVOR PHYSICISTS IN FRANCE [ many retire(d), pure phenomenology persons in Sec.02 - rarity (IN2P3 has got a few)]

# 2016

## MESSAGES

- ☑ All from 2009 and some more.
- ☑ A few cracks are here: USE THE EXPERIMENTAL DATA TO DO PHYSICS!
- ☑ Model building: more observables - towards understanding flavor and CPV
- ☑ LFV can be tackled in SL modes
- ☑ Higgs flavour era around the corner?