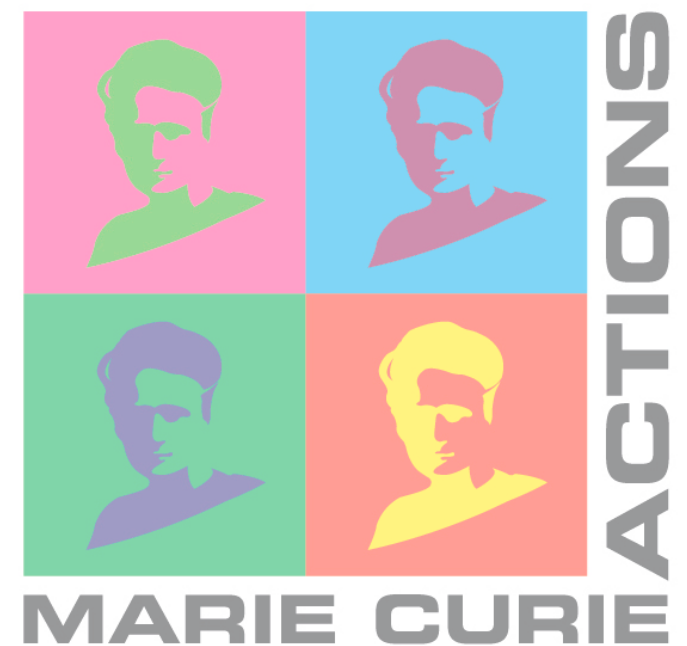


# Accurate predictions for charged Higgs boson production at the LHC

Marco Zaro

LPTHE - Université Pierre et Marie Curie, Paris VI

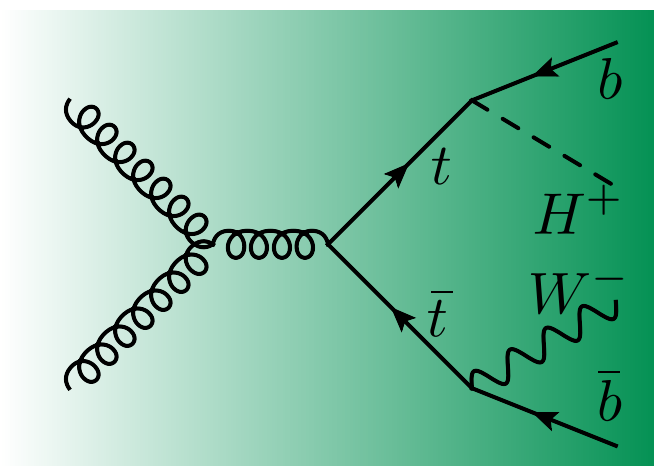
*Rencontre de Physique des Particules 2016*  
*Annecy*



# Charged Higgs production in the (type II) 2HDM

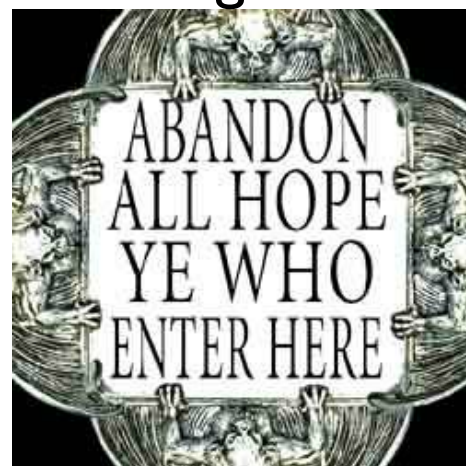
- In the 2HDM, the dominant production channel depends on the Charged Higgs mass

light Higgs



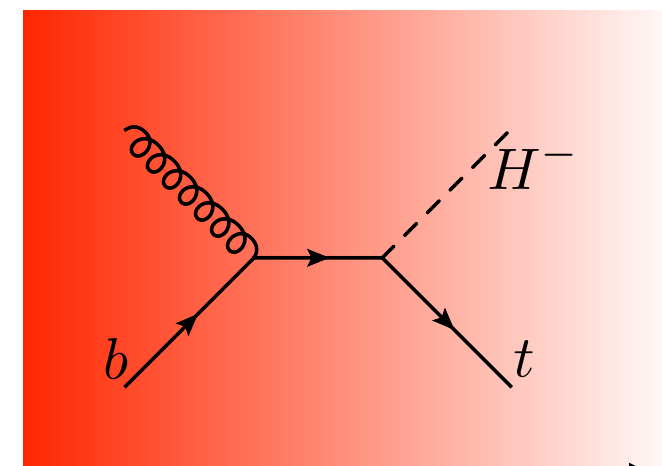
$$M_{H^\pm} < M_t$$

Intermediate  
region



$$M_{H^\pm} \simeq M_t$$

heavy Higgs



$$M_{H^\pm} > M_t$$

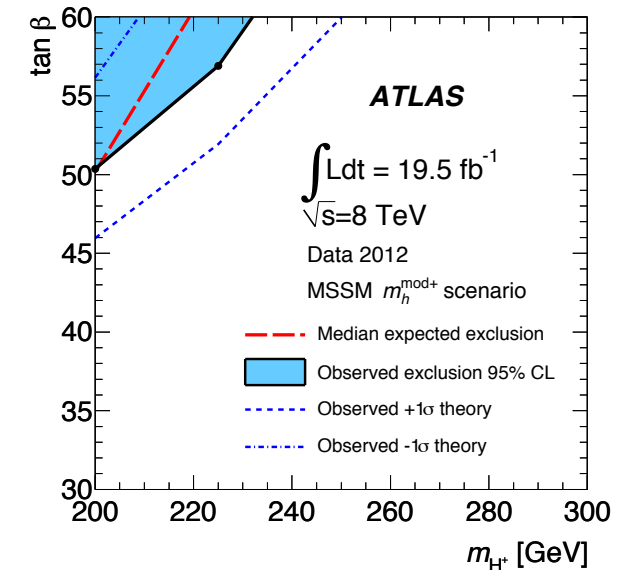
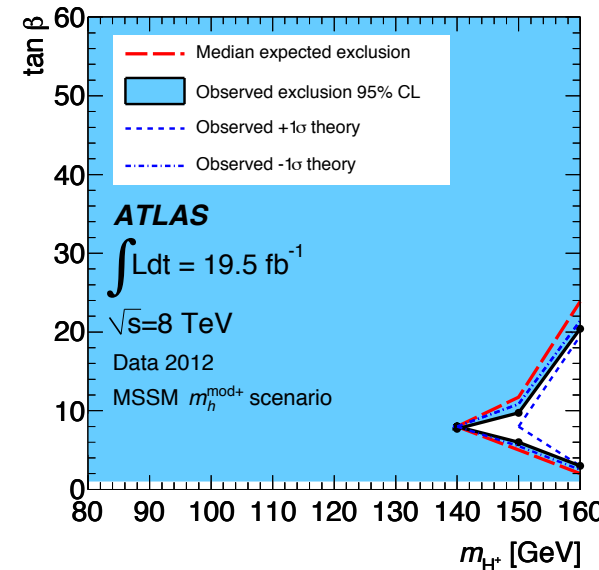
$H^\pm$  mostly produced in  $t\bar{t}$  events.  
Depending on  $\text{BR}(t \rightarrow H^\pm b)$  also  
 $H^\pm t$  can become important.  
At NLO one has to subtract on-  
shell tops  
(see Plehn, [hep-ph/0206121](https://arxiv.org/abs/hep-ph/0206121))

The full  $pp \rightarrow H^\pm W^\mp b \bar{b}$   
process has to be  
simulated.  
Computationally very  
demanding, but feasible

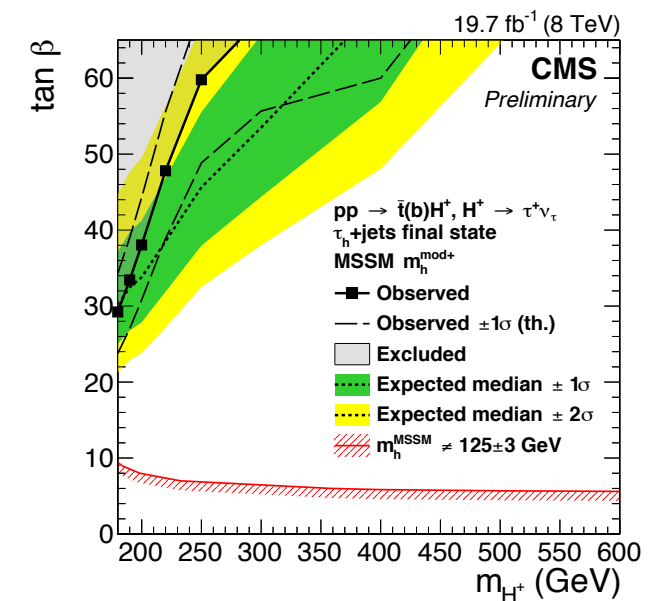
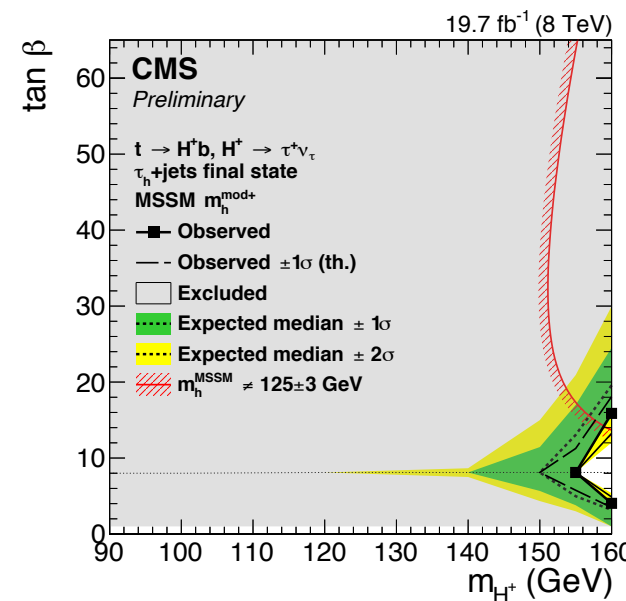
$H^\pm$  mostly produced in  
association with a top quark

- LHC experiments tend to exclude a light charged Higgs
- For a heavy charged Higgs, only very large values of  $\tan\beta$  are excluded
- Missing mass window due to non-existence of NLO predictions for the intermediate range

ATLAS, arXiv:1412.6663

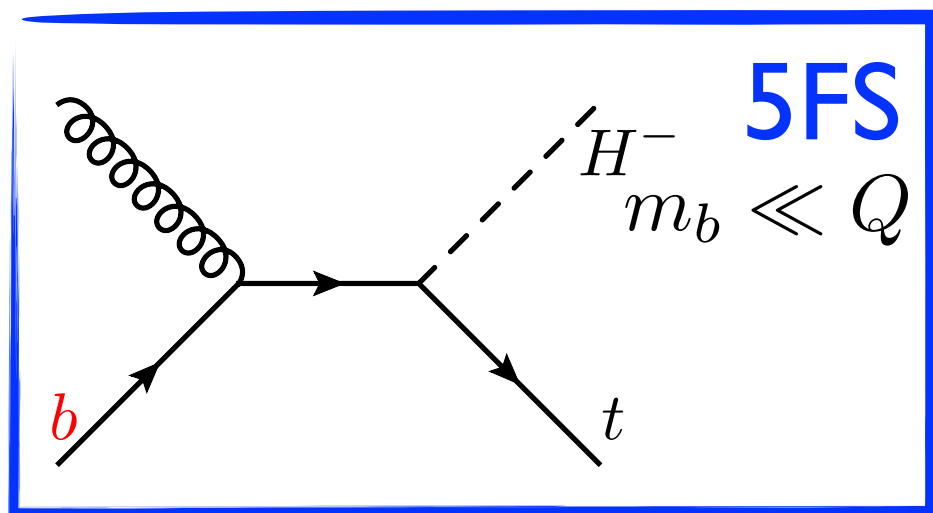


CMS, PAS HIG-14-020

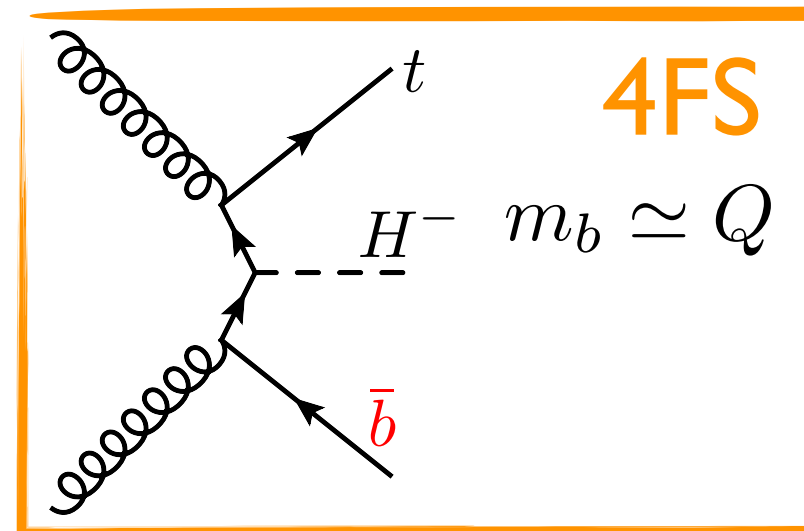


# Heavy charged Higgs production

- Production mechanism features b quarks in the initial state: can be described either with 4- or 5-flavour scheme



- ✓ Simpler process; computing HO is easier
- ✓ b-PDF resums  $\log(m_b/Q)$  at all orders
- ✗ b-quark observables enter at higher orders
- ✗ Matching to PS requires some care (gluon splitting, momentum reshuffling, ...)

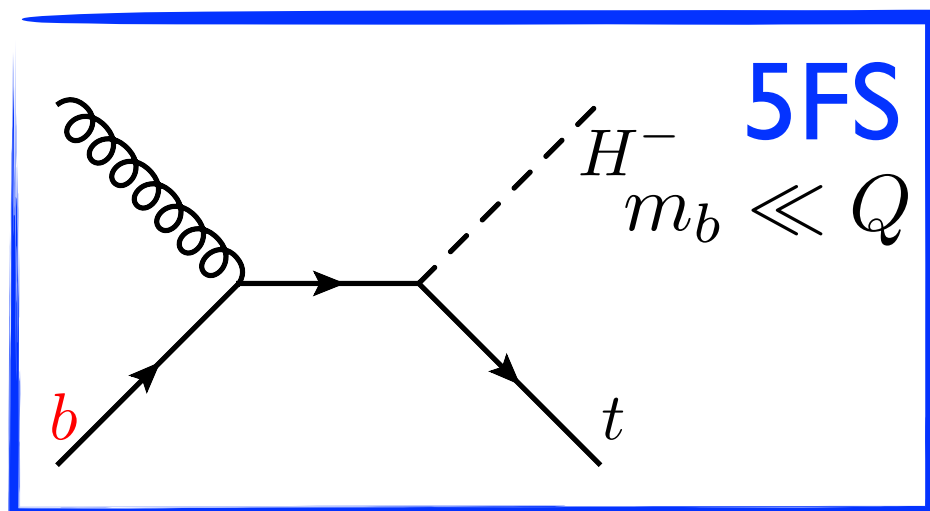


- ✗ Higher multiplicity process; computing HO more involved
- ✗ Cross section can be affected by large  $\log(m_b/Q)$
- ✓ Accounts for b-mass effects and better descriptions of b-quark observables
- ✓ Less ambiguities when matched to PS

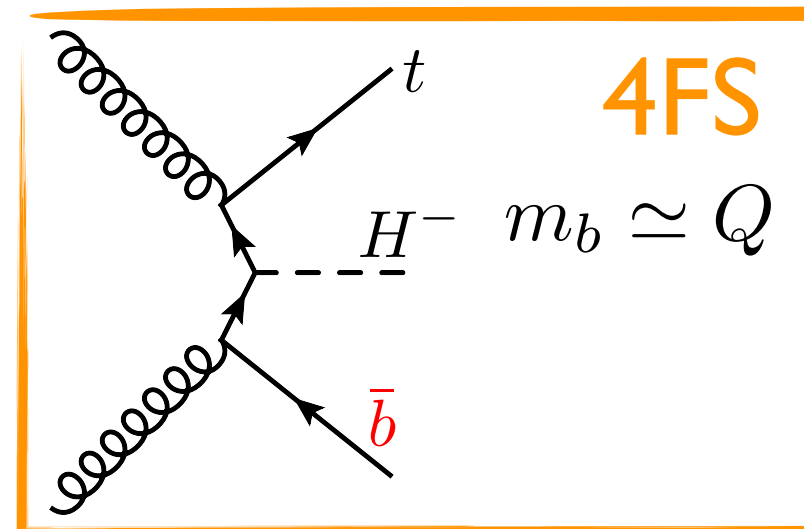
Two schemes are equivalent if all orders were known  
 4FS believed to be superior for fully differential studies

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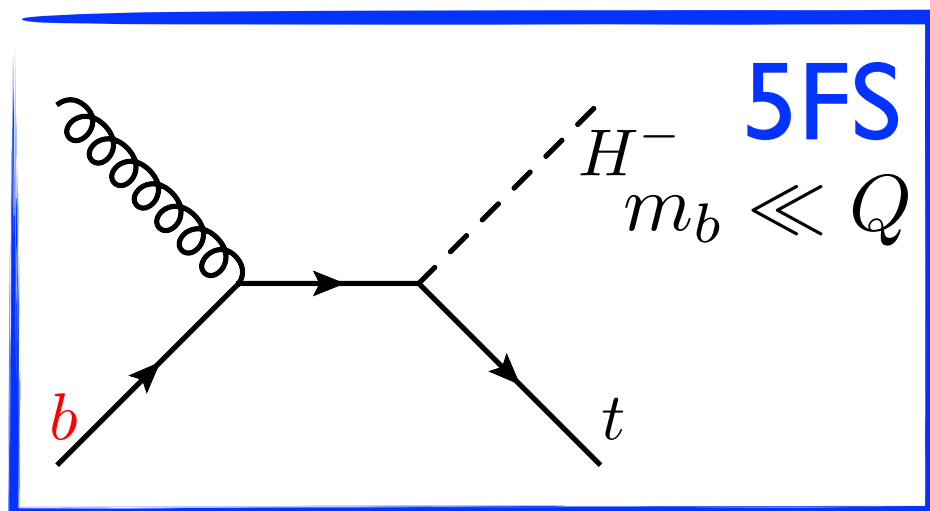
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Suitable scale choice reduces impact of logs

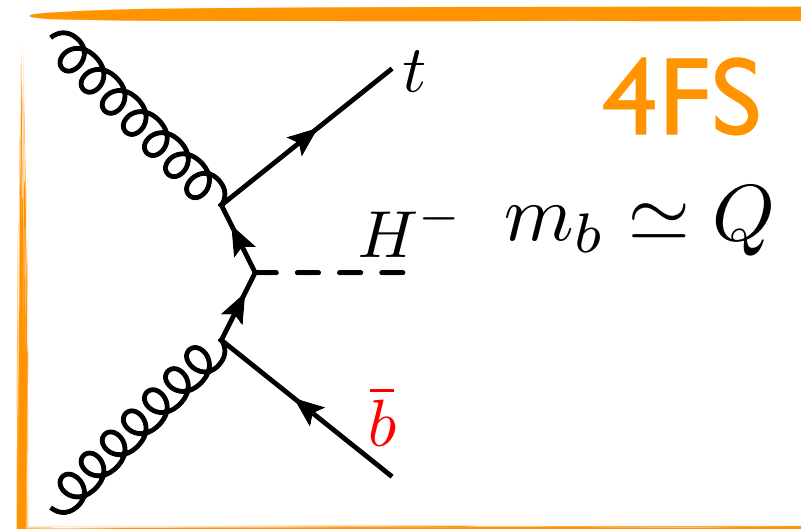
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Crucial for accurate signal simulations

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# Heavy charged Higgs production

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Higher order results:

- NLO (SUSY-)QCD corrections  
[Zhu, hep-ph/0112109](#), [Plehn, hep-ph/0206121](#), [Berger et al, hep-ph/0312286](#)  
[Dittmaier et al, arXiv:0906.2648](#)
- EW corrections  
[Beccaria et al, arXiv:0908.1332](#)  
[Nhung et al, arXiv:1210.4087](#)
- Threshold resummation  
[Kidonakis, arXiv:1005.4451](#)
- Fully differential NLO+PS  
[Weydert et al, arXiv:0912.3430](#), [Klasen et al, arXiv:1203.1341](#)  
[Degrande, Ubiali, Wiesemann, MZ, arXiv:1507.02549](#);

suitable scale  
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Crucial for  
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✓ Simpler  
 ✓ b-PDF  
 ✗ b-quark  
 ✗ Matching  
 splitting

Two schemes are equivalent if all orders were known  
 4FS believed to be superior for fully differential studies

# Fully differential comparison of 4 and 5FS

Degrande, Ubiali, Wieseemann, MZ, arXiv:1507.02549

- Use modern automated tool chains to generate the code, starting from the model Lagrangian
  - Generate UV/ $R_2$  counterterms for the evaluation of loops with NLOCT [Degrande arXiv:1406.3030](#)
  - Use MadGraph5\_aMC@NLO to generate the code for event generation [Alwall et al. arXiv:1405.0301](#)
- MSbar renormalisation to be preferred for  $y_b$ : logs of  $\mu_R/m_b$  resummed. Add  $m_b(\mu_R)$  dependence as in [Wieseemann et al. arXiv:1409.5301](#)
- b-initiated processes typically prefer scales lower than  $\hat{s}$ . This argument holds also for the shower scale
- Keep  $H^-$  stable, decay top quark leptonically  
→ One b-jet from top and one from matrix element / shower



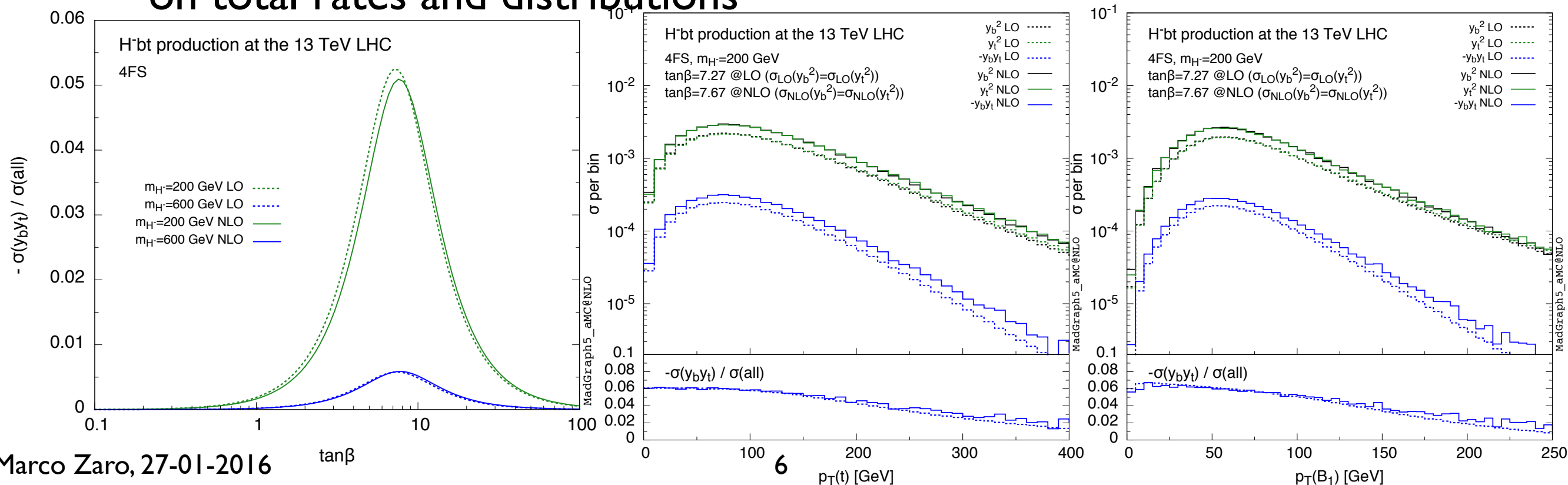
# Setup and cross-section structure

- The following parameters are used

$$\sqrt{S} = 13 \text{ TeV} \quad m_H = 200 \text{ GeV} \quad \tan \beta = 8$$

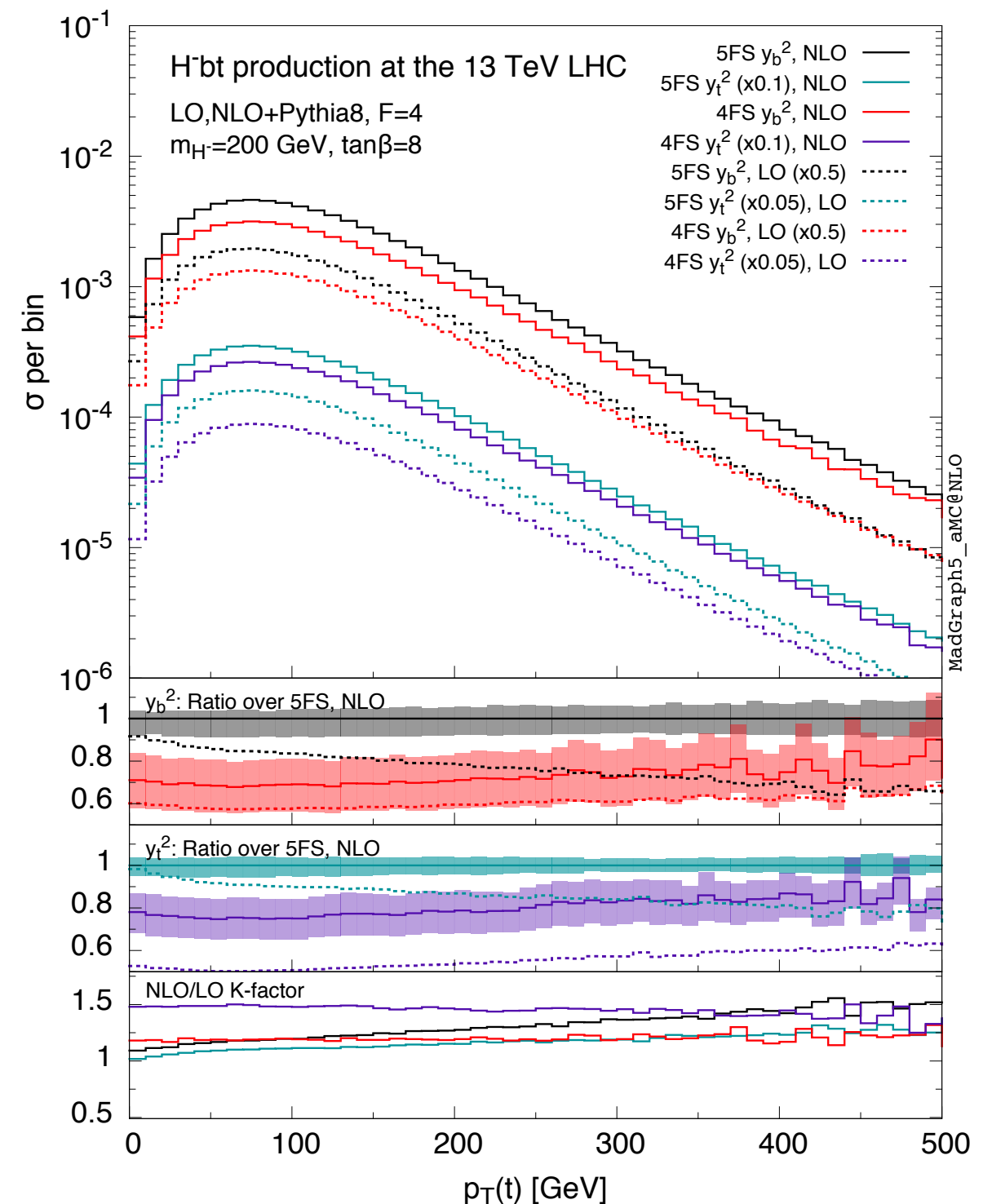
$$\mu_R = \mu_F = \mu_B = H_T/3 = \sum \sqrt{p_T(i)^2 + m(i)^2}/3$$

- Owing to the structure of the  $H^\pm tb$  coupling, the cross section will receive three contributions:  $y_b^2$  ( $\sim \tan^2 \beta$ ),  $y_t^2$  ( $\sim 1/\tan^2 \beta$ ) and  $y_b y_t$  ( $\tan \beta$  independent).
- In the 5FS, the  $y_b y_t$  term is null (helicity conservation)
- In the 4FS, it is proportional to  $m_b/\hat{s}$ . Numerically it turns to be negligible on total rates and distributions



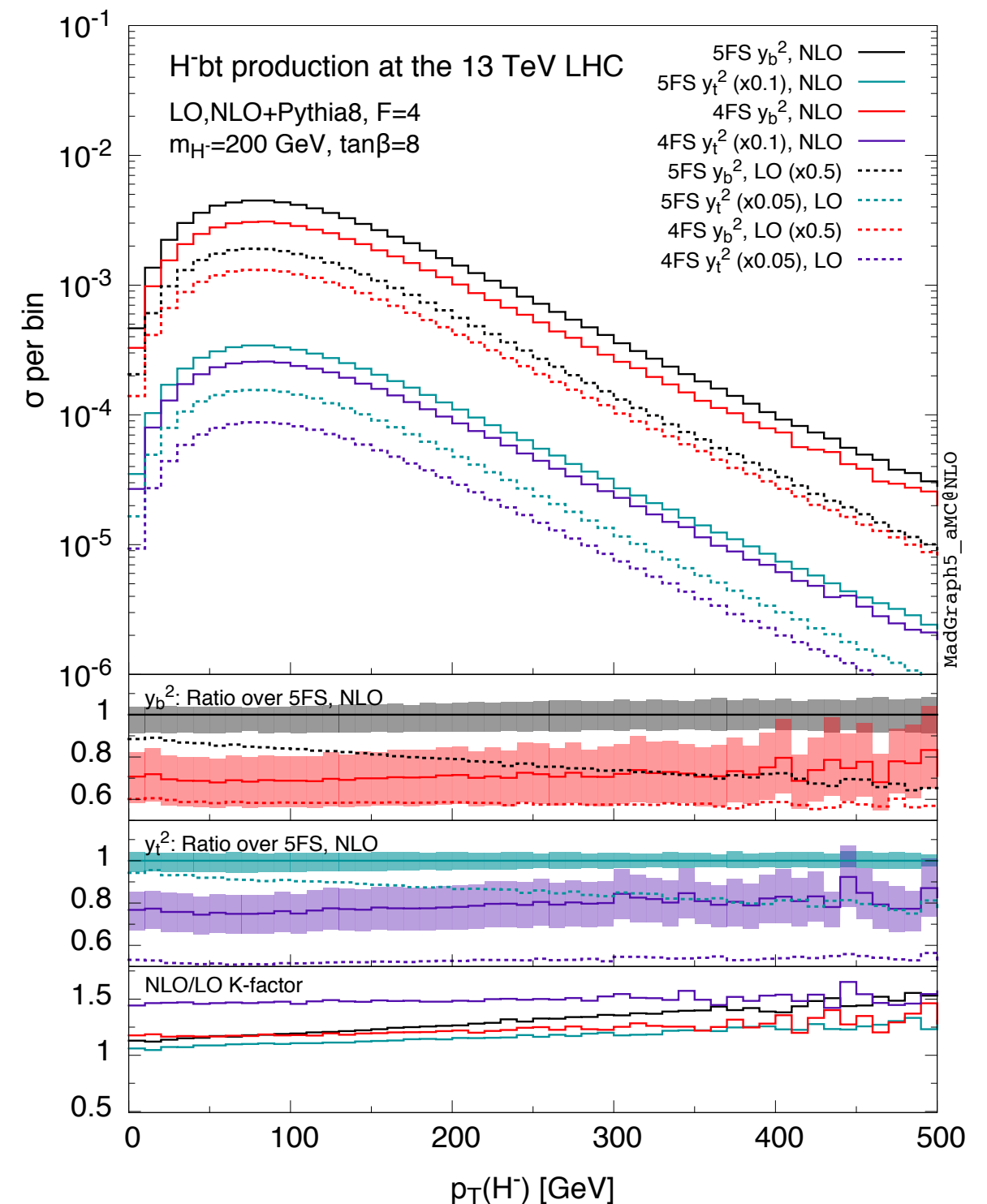
# 4FS vs 5FS

- NLO corrections improve the shape agreement between the two schemes
- Very good agreement for b-inclusive observables ( $p_T(t)$ ,  $p_T(H)$ )
- Agreement remains quite good for more exclusive observables ( $p_T(b_2)$ ), despite the large K-factors and uncertainties for 5FS
- Very exclusive observables ( $p_T(B_2)$ ) show larger discrepancies in regions where mass-effects are enhanced
- Bottom line: 4FS provides a better description of exclusive observables



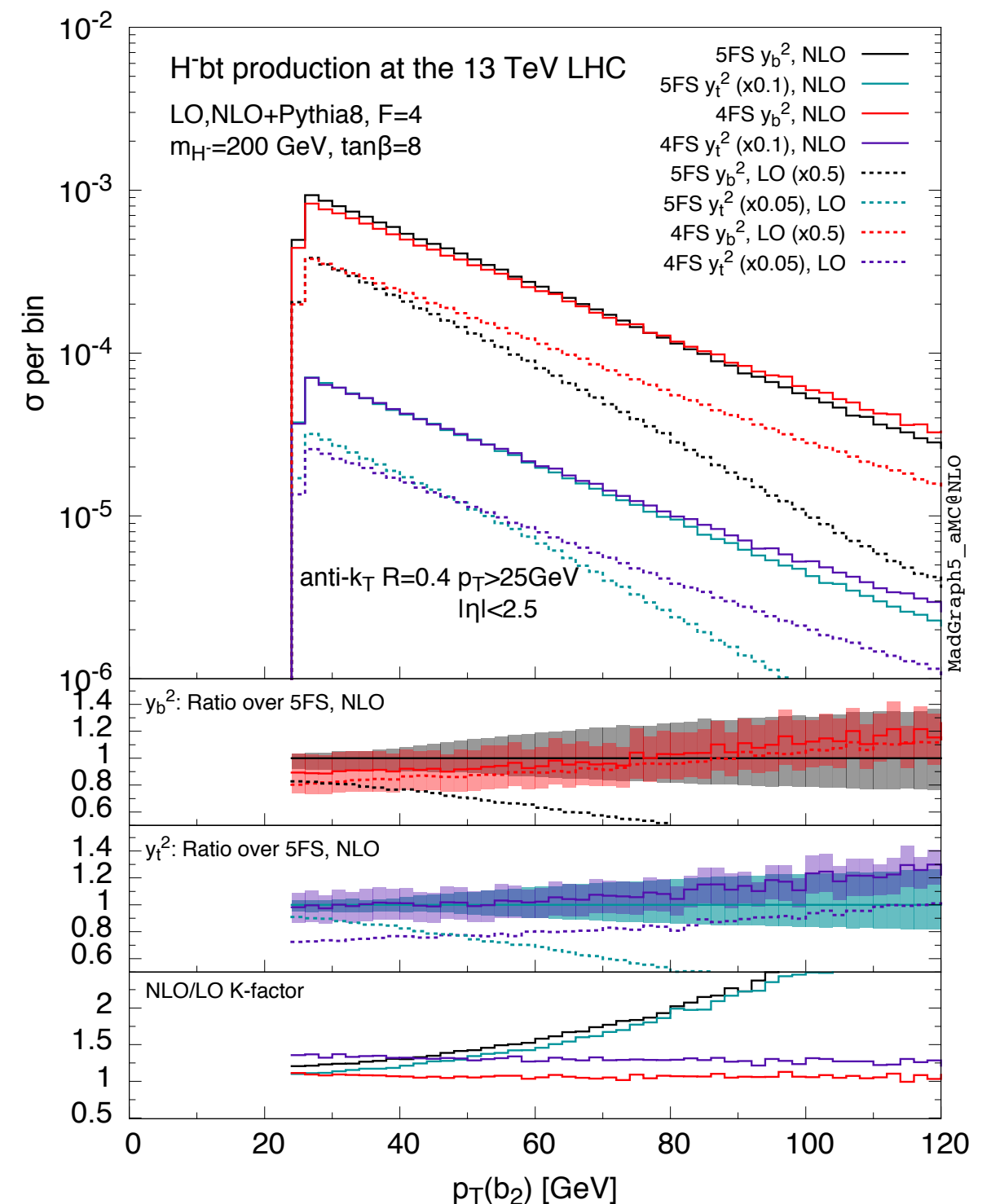
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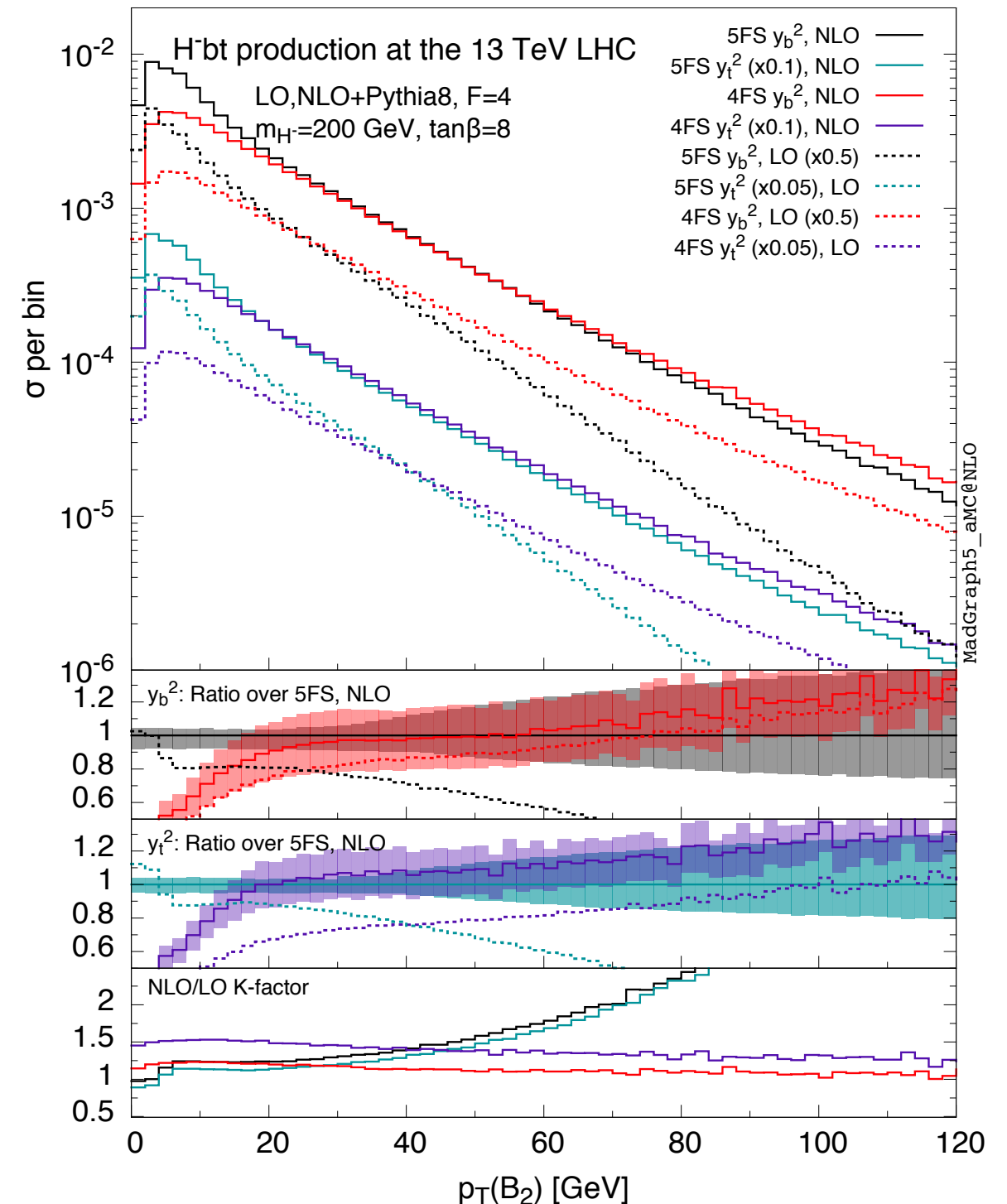
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# Conclusions

- The discovery of a charged Higgs boson at the LHC would be a clear sign of BSM physics
- Lot of recent and ongoing progress towards providing accurate predictions for cross section and realistic signal modelling
- Fully differential predictions for charged Higgs production in the 2HDM available for the first time in the 4FS at NLO+PS
  - Better description of b-kinematics
  - Better matching to PS, less effects due to reshuffling
  - 4FS to become reference for signal simulations at the LHC



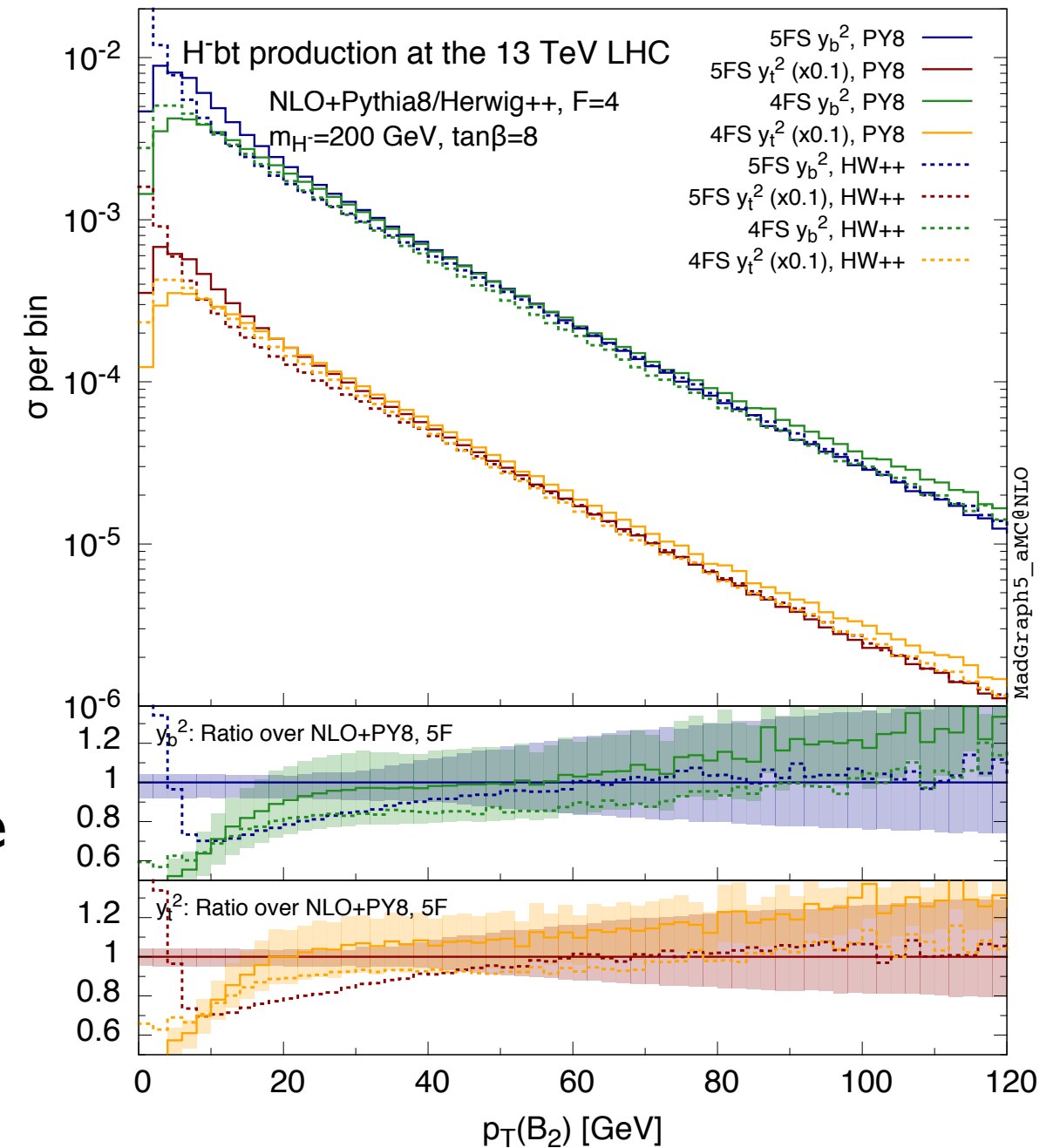


Thank you!

# Bakup slides

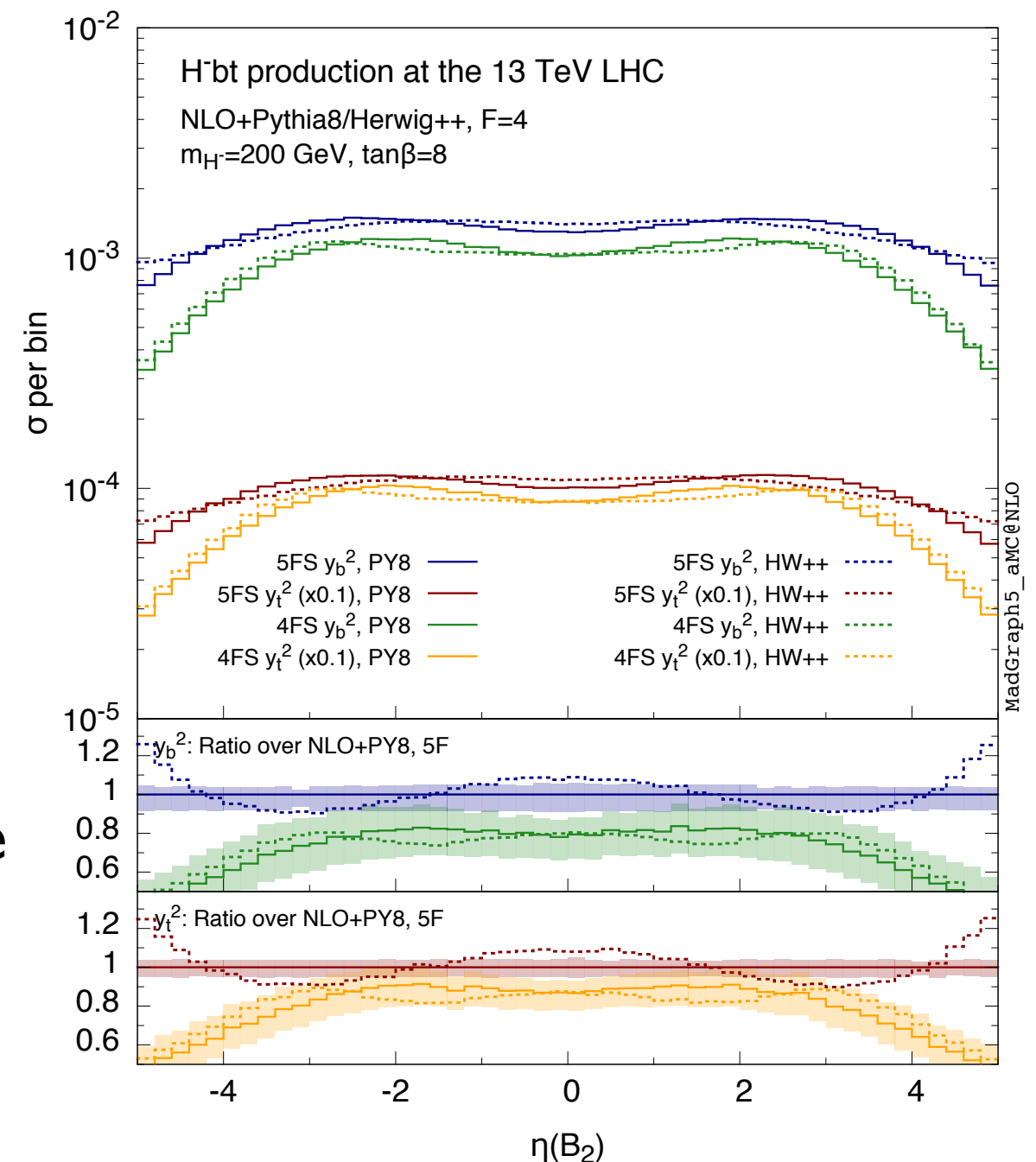
# Parton-shower dependence

- The inclusion of  $m_b$  effects at the matrix-element level should reduce possible systematics due to the PS used
- The largest effects should be visible on the most exclusive observables ( $p_T(B_2), \eta(B_2)$ )
- Herwig++ in the 5FS produces b hadrons too close to the beam pipe;  $\Delta R(B_1, B_2)$  is a quite spectacular case
- Effects are mitigated for less exclusive observables ( $p_T(b_2), \eta(b_2)$ ), or when 2 bjets are required
- 4FS more resilient against shower effects



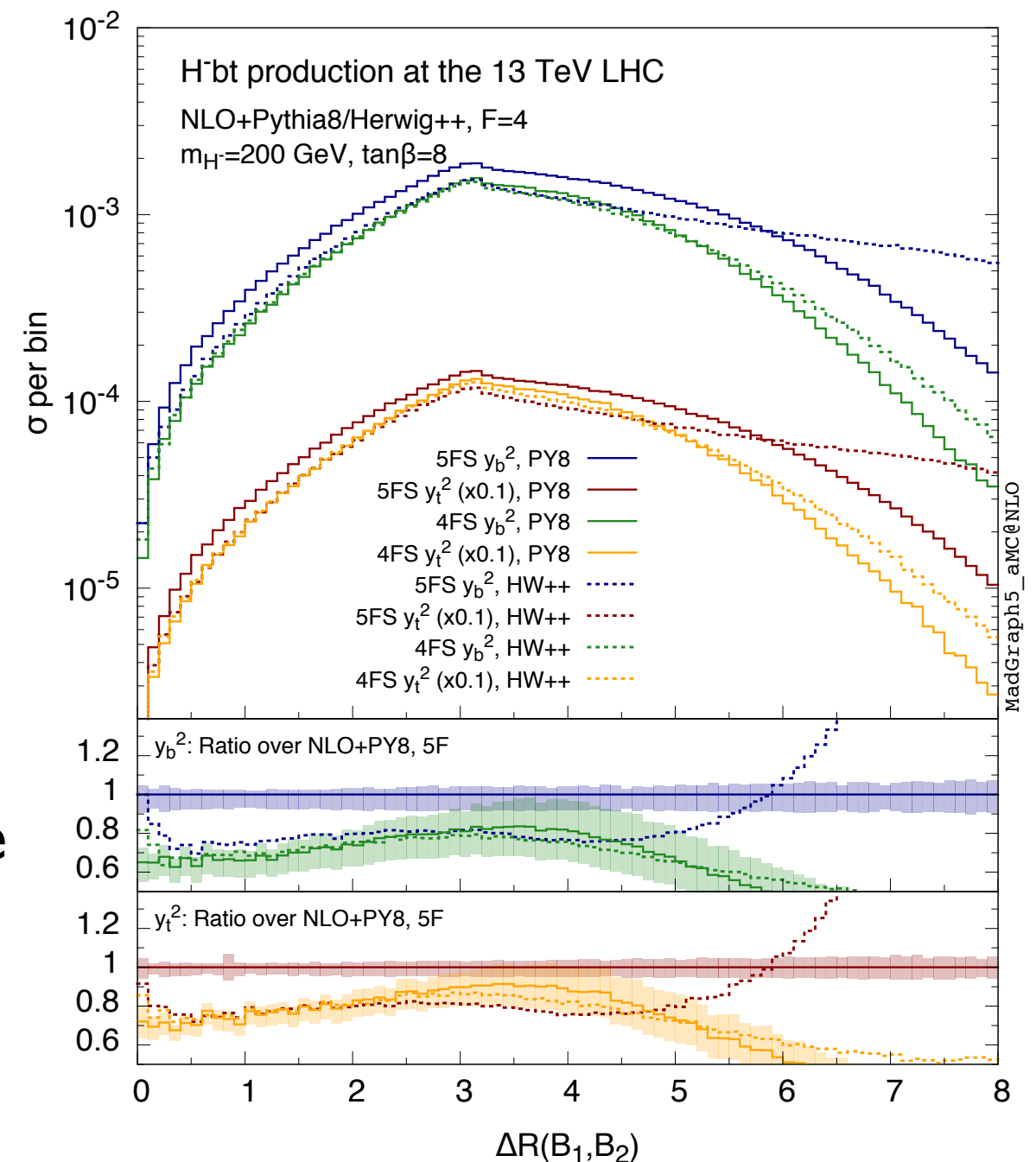
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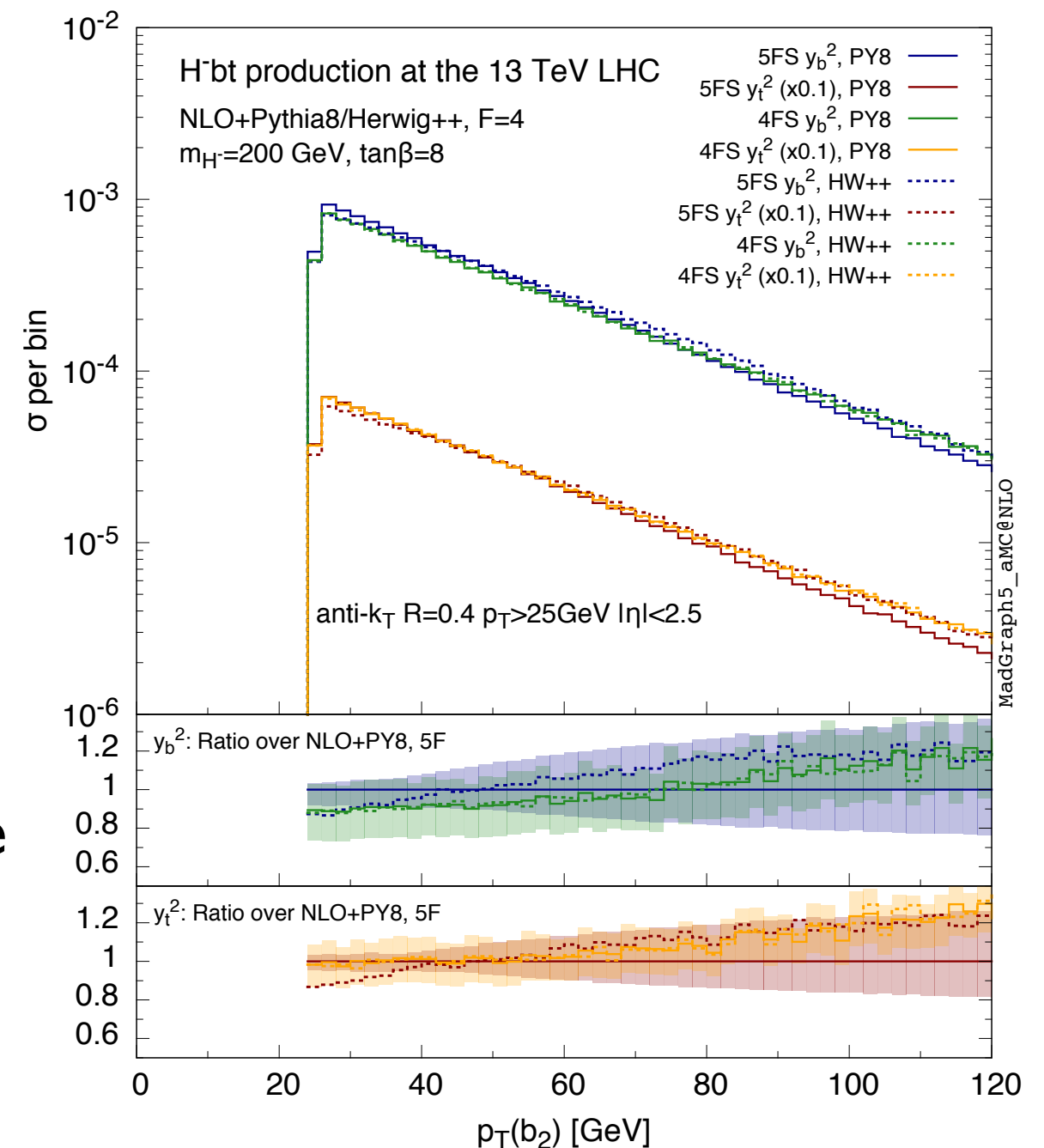
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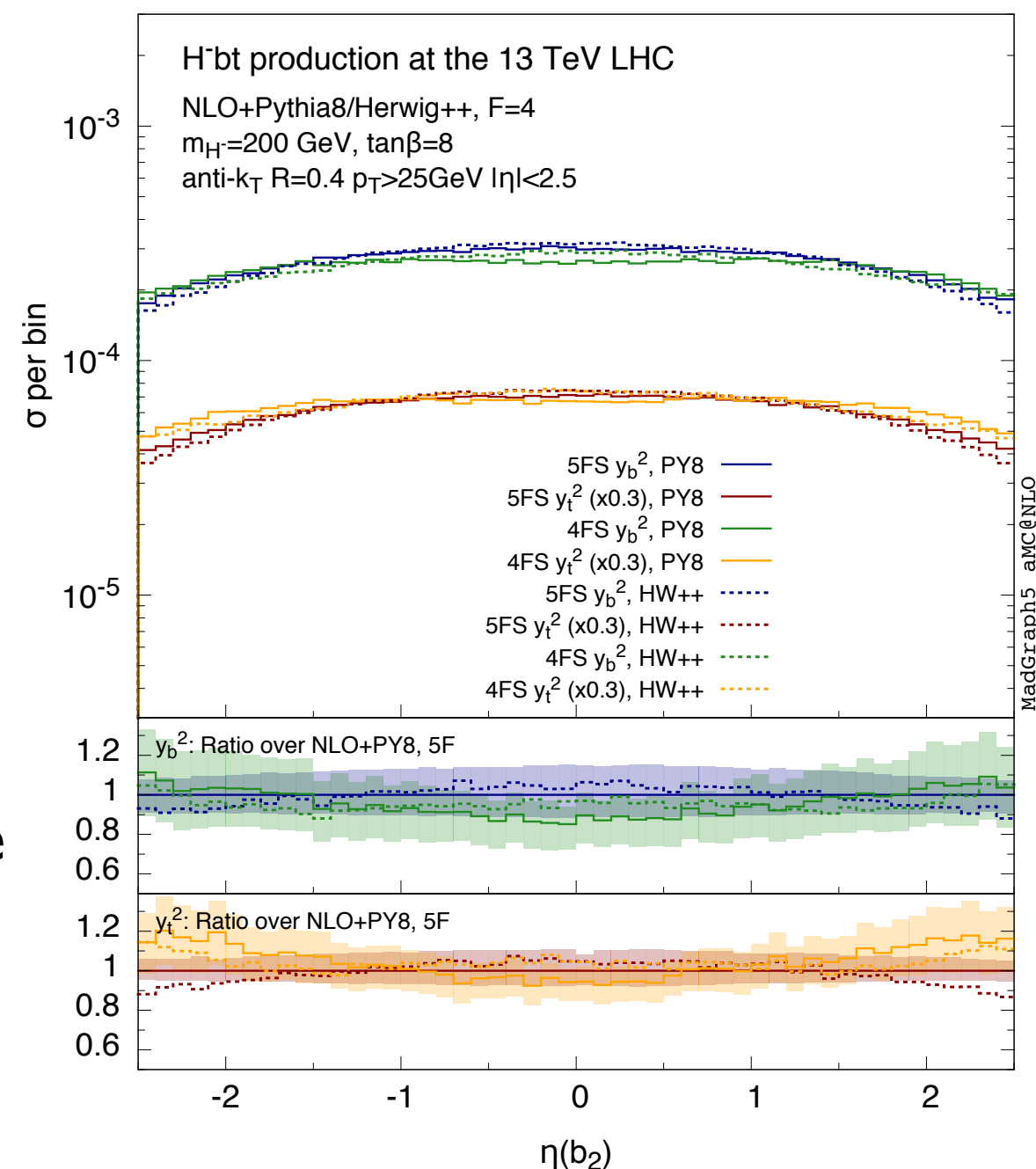
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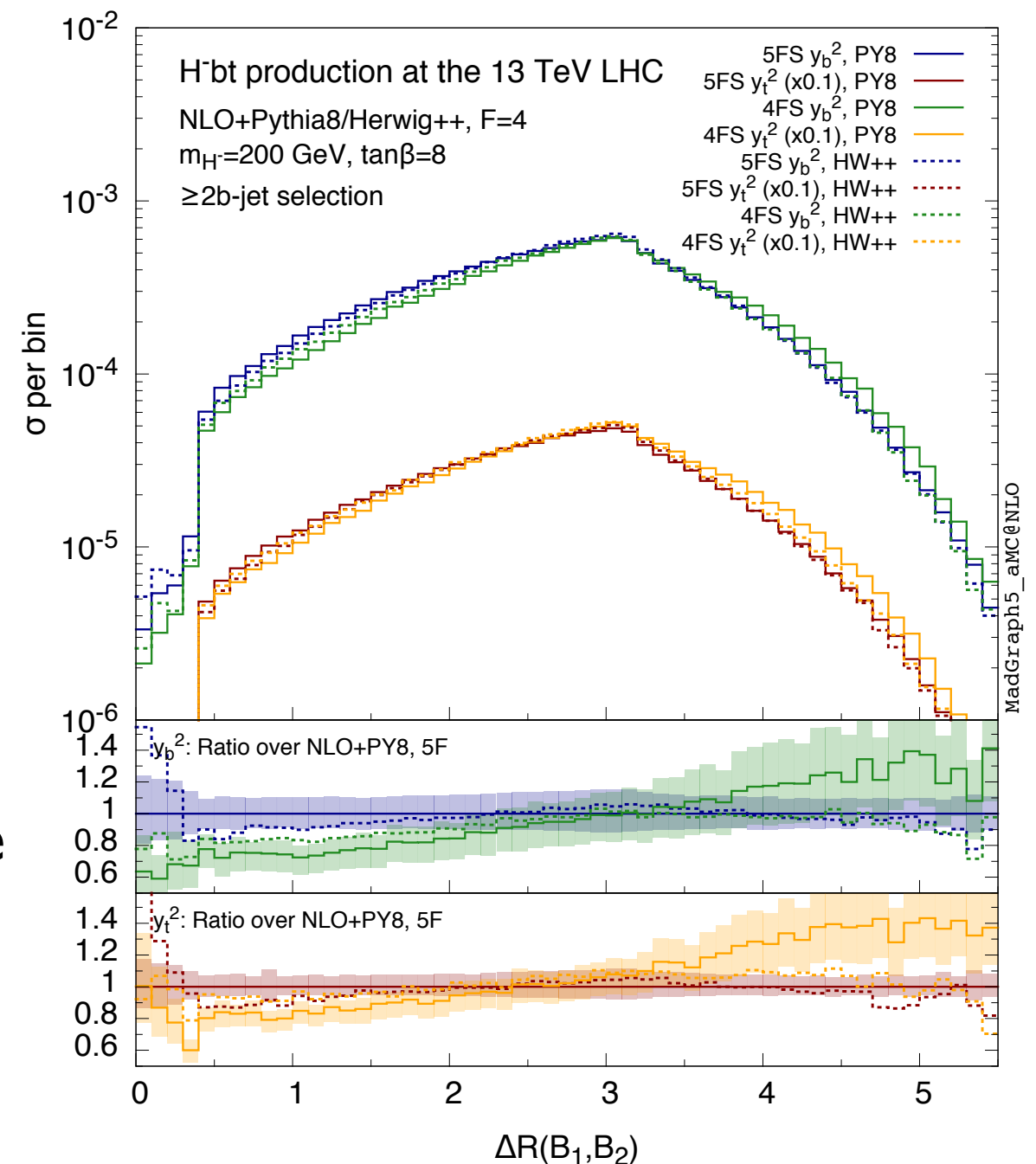
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# Matched predictions for the total cross-section

Flechl, Klees, Kramer, Spira, Ubiali, arXiv:1409.5615

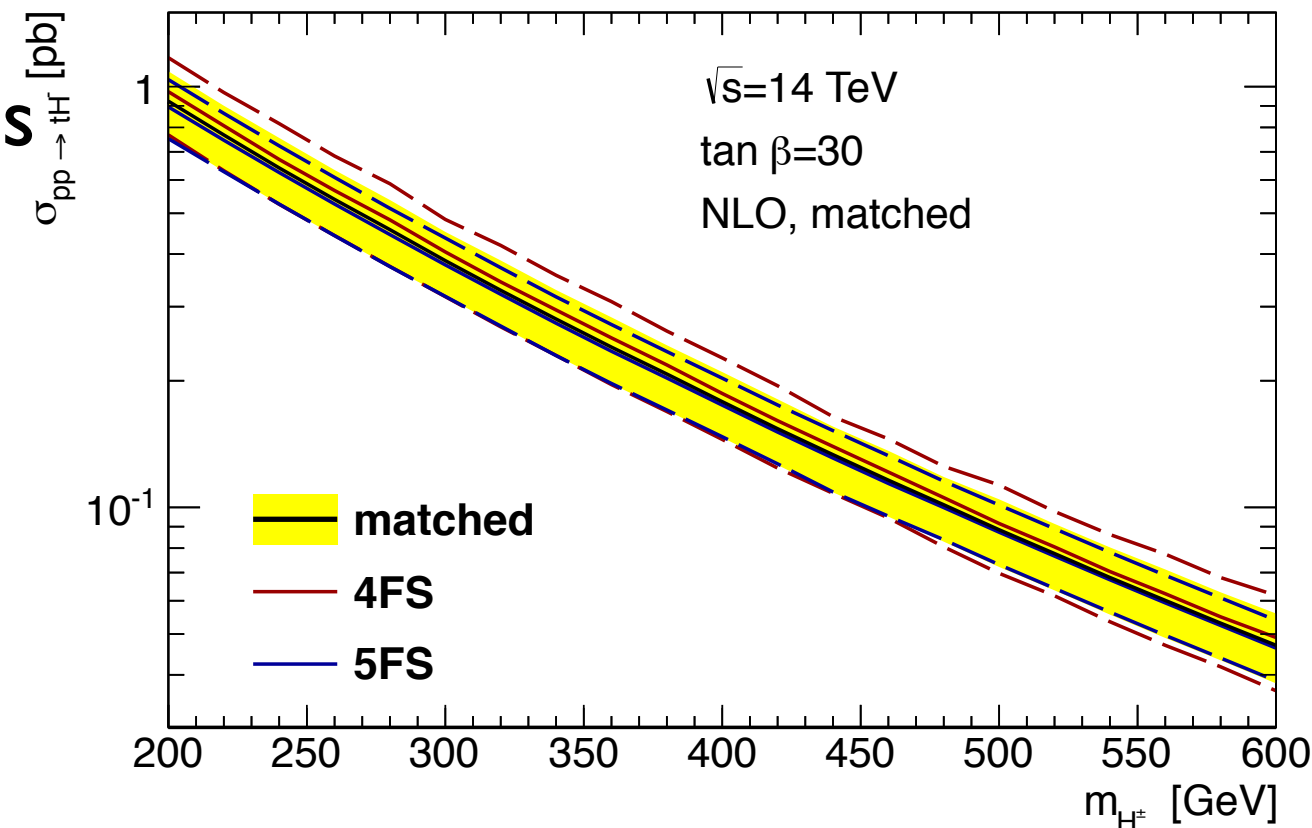
- The scale in the logs resumed in the 5FS is typically much smaller than the hard scale of the process (phase-space suppression) [Maltoni, Ridolfi, Ubiali, arXiv:1203.6393](#)

$M_{H^\pm}$ [GeV]	$\tilde{\mu}$ [GeV]	8 TeV		14 TeV	
		$(m_t + M_{H^\pm})/\tilde{\mu}$		$(m_t + M_{H^\pm})/\tilde{\mu}$	
200	67.3	5.5		74.9	5.0
300	80.3	5.9		90.6	5.2
400	92.1	6.2		105.3	5.4
500	103.1	6.5		119.0	5.7

$$Q_{tHb}^2 = M^2 \frac{(1-z)^2}{z} \quad z = \frac{M^2}{\hat{s}}$$

- Set  $\mu_F = \tilde{\mu}$  in the 5FS
- Include all sources of uncertainties
  - scale, PDF (PDF4LHC),  $m_b$ ,  $\alpha_s$
- Compare 5FS, 4FS @NLO and Santander-matched prediction

$$\sigma_{\text{matched}} = \frac{\sigma_{4F} + w\sigma_{5F}}{1 + w} \quad w = \log \frac{m_{H^\pm}}{m_b} - 2$$



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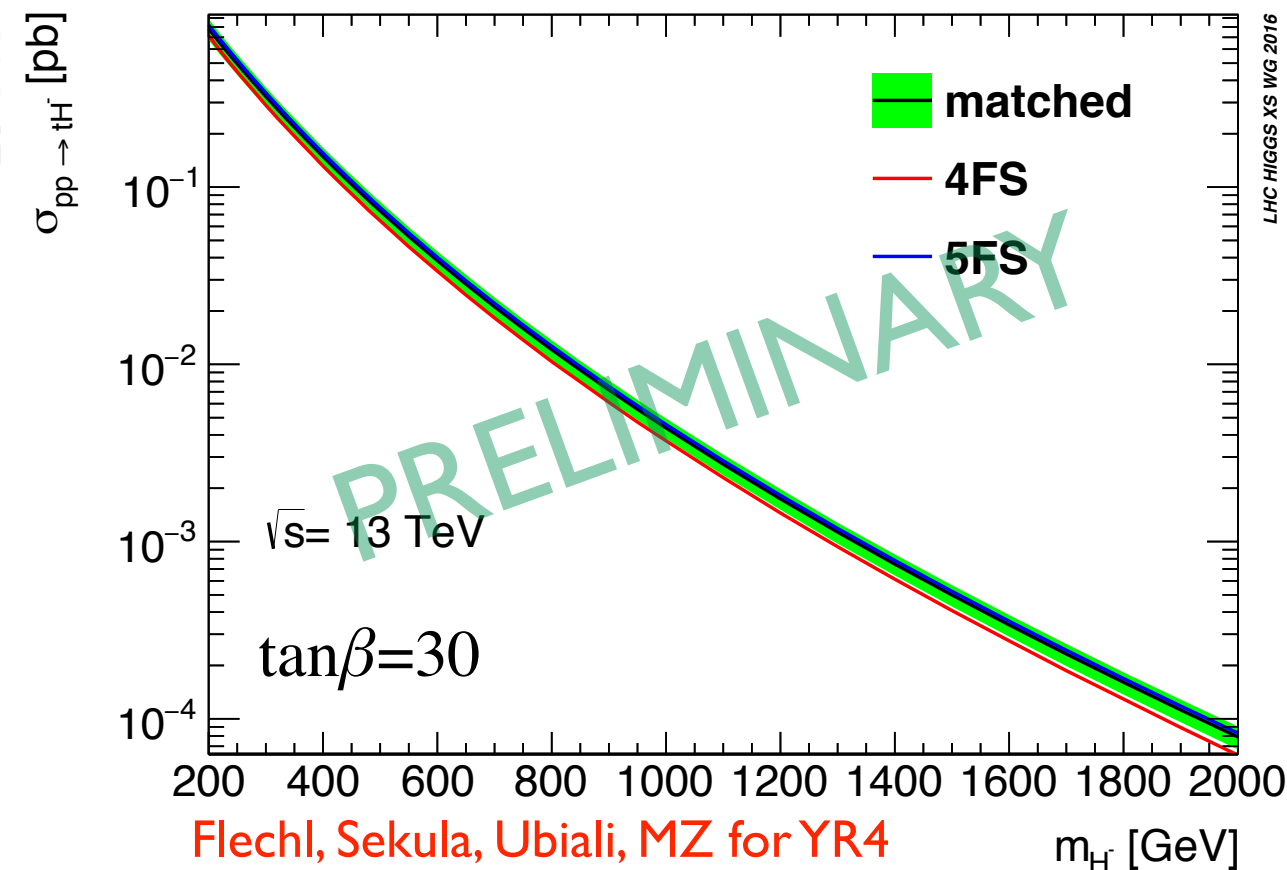
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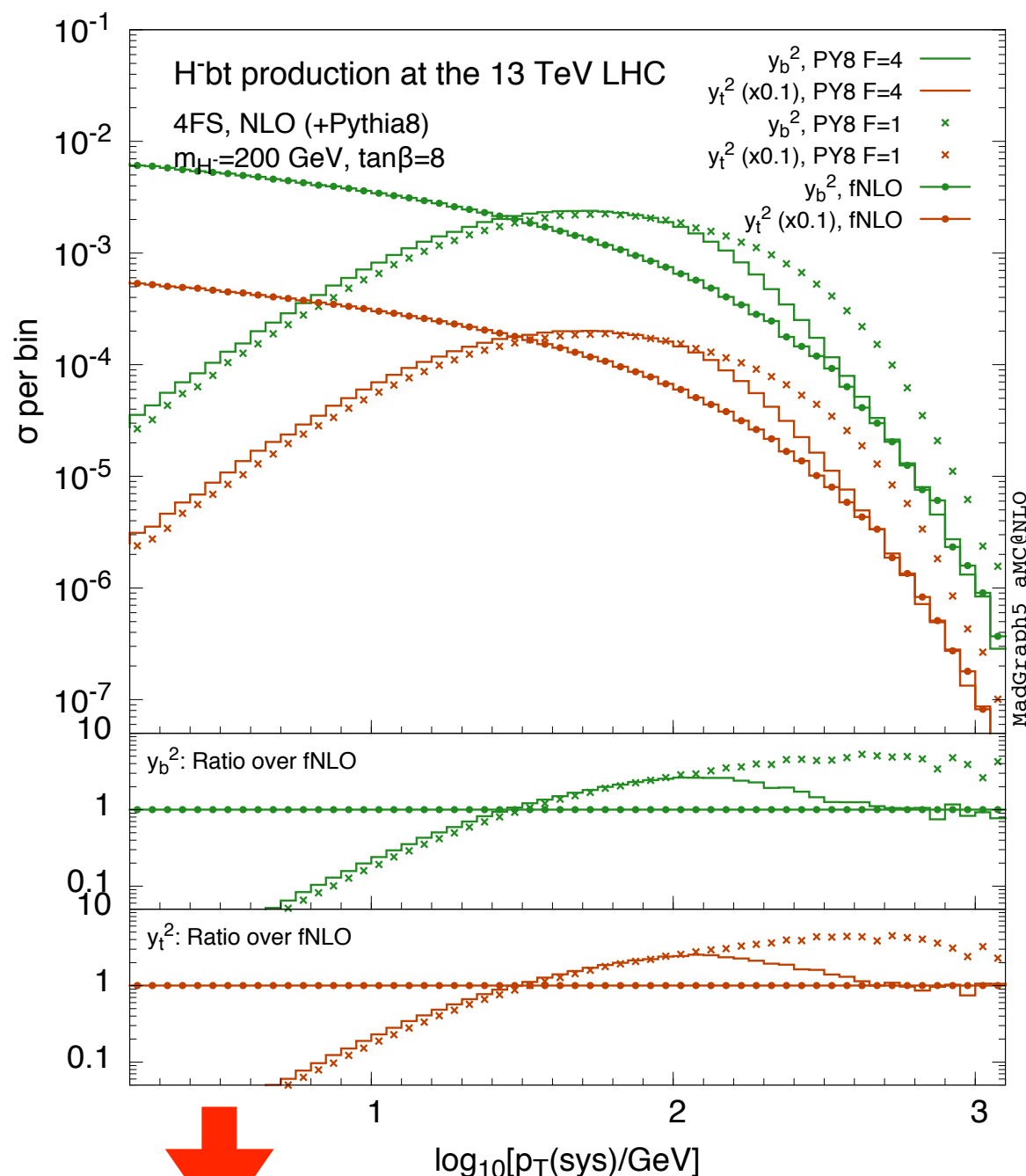
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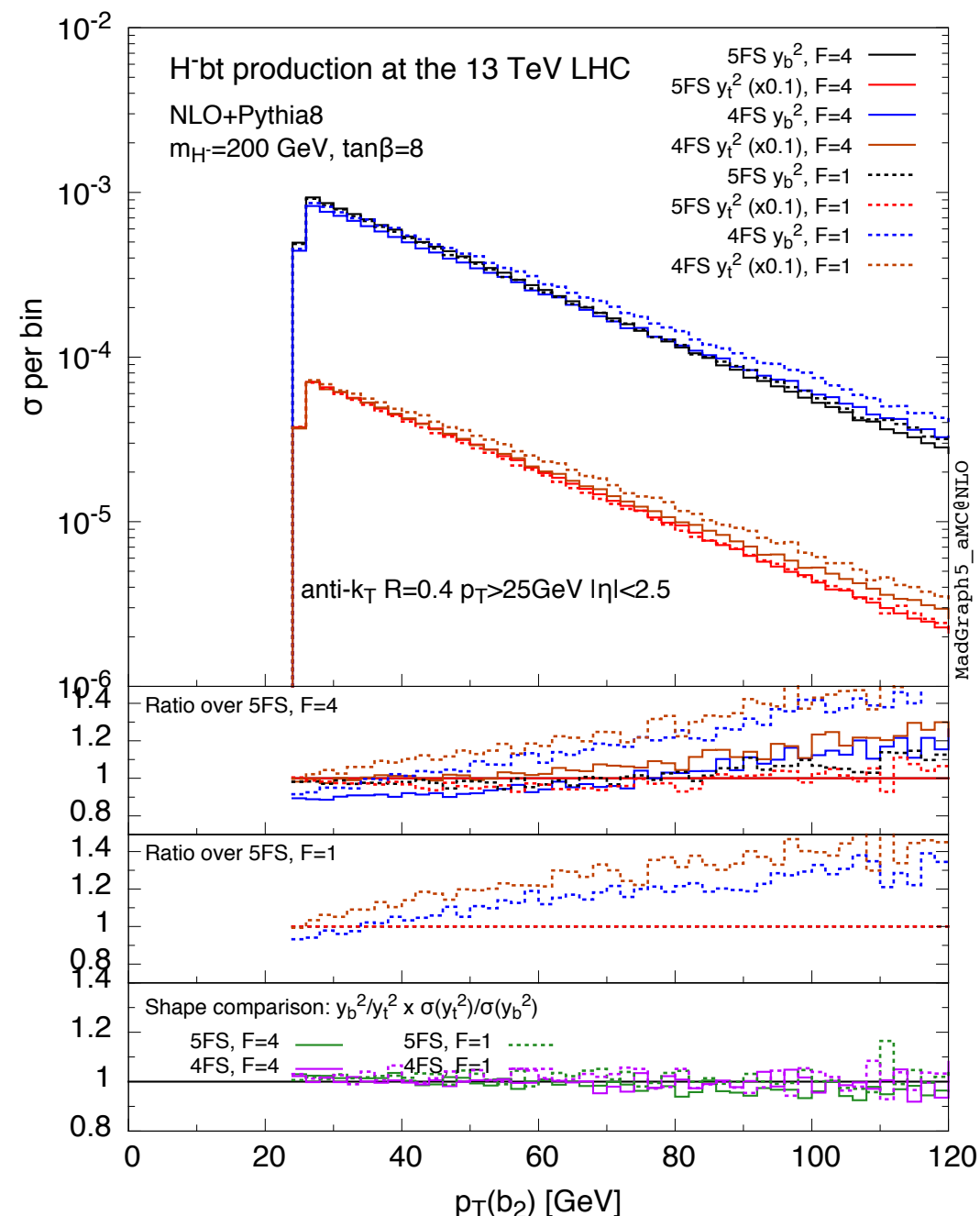
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# Choice of shower scale



Reduced shower scale to be preferred  
improves NLO+PS/fNLO matching at high- $p_T$



Consequence:  
better agreement at differential level  
between the two schemes