

# Vector-Like Quark Phenomenology: From the $A_{\text{FB}}^b$ and $t\bar{t}H$ Anomalies to the HL-LHC Reach (Based on 1510.07527)

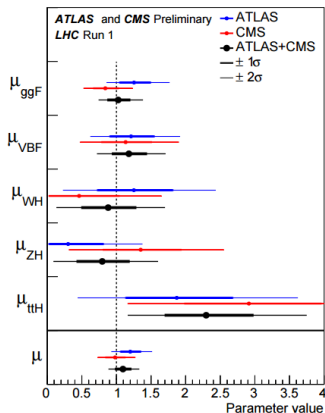
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# Introduction

- What are **Vector-Like Quarks** (VLQs) ? Quarks whose **left** and **right**-handed chiralities transform the **same** way under  $SU(3)_c \times SU(2)_L \times U(1)_Y$ ;
- Gauge-invariant **mass term** without a Higgs:  $M\bar{Q}_L Q_R + \text{h.c.}$  ;
- Can couple with SM quarks via **Yukawa** interactions  $\rightarrow$  after EWSB, **mixing** with SM quarks;
- Common assumption: **mixing** only with the **3rd generation** SM quarks;
- Eigenmasses and couplings obtained by **rotating** from the interaction basis to the mass basis;
- **Examples** of VLQs: **Kaluza-Klein states** in extra-dimensional scenarios, **bound states** in Composite Higgs theories.

# Motivation



	Measurement	Fit	$10^{\text{meas}} - 10^{\text{fit}} / \sigma^{\text{meas}}$
$\Delta\sigma_{\text{had}}^{(S)}(m_Z)$	$0.02750 \pm 0.00033$	0.02759	0.00009
$m_Z$ [GeV]	$91.1875 \pm 0.0021$	91.1874	-0.0001
$\Gamma_Z$ [GeV]	$2.4952 \pm 0.0023$	2.4959	0.0007
$\sigma_{\text{had}}^0$ [nb]	$41.540 \pm 0.037$	41.478	-0.062
$R_b$	$20.767 \pm 0.025$	20.742	-0.025
$A_{\text{fb}}^{0,l}$	$0.01714 \pm 0.00095$	0.01645	-0.00069
$A_1(P_f)$	$0.1465 \pm 0.0032$	0.1481	0.0016
$R_b$	$0.21629 \pm 0.00066$	0.21579	-0.0005
$R_c$	$0.1721 \pm 0.0030$	0.1723	0.0002
$A_{\text{fb}}^{0,b}$	$0.0992 \pm 0.0016$	0.1038	0.0046
$A_{\text{fb}}^{0,c}$	$0.0707 \pm 0.0035$	0.0742	0.0035
$A_b$	$0.923 \pm 0.020$	0.935	0.012
$A_c$	$0.670 \pm 0.027$	0.668	-0.002
$A_1(\text{SLD})$	$0.1513 \pm 0.0021$	0.1481	-0.0032
$\sin^2\theta_{\text{eff}}^{\text{lep}}(Q_{\text{fb}})$	$0.2324 \pm 0.0012$	0.2314	-0.0010
$m_W$ [GeV]	$80.385 \pm 0.015$	80.377	-0.008
$\Gamma_W$ [GeV]	$2.085 \pm 0.042$	2.092	0.007
$m_t$ [GeV]	$173.20 \pm 0.90$	173.26	0.06

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Figure : Higgs production rates [ATLAS-CONF-2015-044] (left) and LEP EW precision results [<http://lepewwg.web.cern.ch>] (right).

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# Towards a Minimal Model: The Bottom Sector

The **minimal** VLQ field content necessary to explain (i.e. reproduce the experimental result at less than  $1\sigma$ )  $A_{FB}^b$  and  $t\bar{t}H$  can be inferred from phenomenological arguments [hep-ph/0109097, 1209.6382]. For the **bottom sector**:

- The measured values of  $A_{FB}^b$  and  $R_b$  favor a  $SU(2)_L$  VL multiplet containing a bottom-like  $b'$  with positive isospin  $\rightarrow$  minimal choice is a **doublet**:

$$B_{L,R} = (b', q_{-4/3}) \sim (3, 2, -5/6).$$

- Doublet is not enough for reproducing the experimental results  $\rightarrow$  add also a **singlet**:

$$b''_{L,R} \sim (3, 1, -1/3).$$

# Towards a Minimal Model: The Top Sector

One can repeat the previous exercise for the **top sector**:

- Adding only one VLQ multiplet containing a top partner cannot explain the  $t\bar{t}H$  excess and reproduce the measured top mass of  $\sim 174$  GeV;
- To significantly enhance  $y_t$ , one needs VLQs that have Yukawa couplings in the interaction basis with both  $t_L$  and  $t_R \rightarrow$  **minimal** choice is adding a VL SU(2)<sub>L</sub> **doublet**,  $T$ , and a **singlet**,  $t''$  ( $T_{L,R} = (t', b')$  disfavored because  $b'$  has isospin  $< 0$ ):

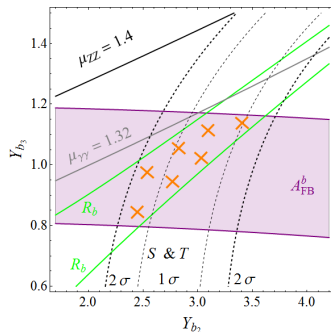
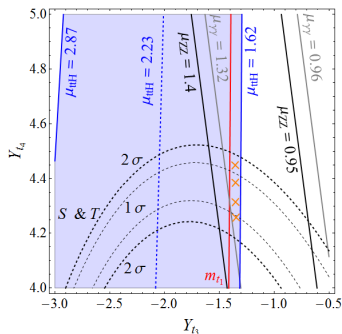
$$T_{L,R} = (q_{5/3}, t') \sim (3, 2, 7/6), \quad t''_{L,R} \sim (3, 1, 2/3).$$

Apart the SM  $Q_L = (t, b)_L$ ,  $t_R$  and  $b_R$ , the **minimal model** will contain:

$$T_{L,R}, t''_{L,R}, B_{L,R}, \text{ and } b''_{L,R}.$$

# The Minimal Model: Numerical Results

$$\begin{aligned}
 -\mathcal{L} \supset & Y_{t_1} \bar{Q}_L \tilde{H} t_R + Y_{t_2} \bar{Q}_L \tilde{H} t_R'' + Y_{t_3} \bar{T}_L H t_R + Y_{t_4} \bar{T}_L H t_R'' + Y_{t_5} \bar{T}_R H t_L'' \\
 & + Y_{b_1} \bar{Q}_L H b_R + Y_{b_2} \bar{Q}_L H b_R'' + Y_{b_3} \bar{B}_L \tilde{H} b_R + Y_{b_4} \bar{B}_L \tilde{H} b_R'' + Y_{b_5} \bar{B}_R \tilde{H} b_L'' \\
 & + m_1 \bar{T}_L T_R + m_2 \bar{t}_L' t_R'' + m_3 \bar{B}_L B_R + m_4 \bar{b}_L' b_R'' + \text{h.c.} .
 \end{aligned}$$



**Figure :** Top sector (left) and bottom sector (right). In the allowed regions of parameter space, the lightest VLQs have masses of  $\sim 1$  TeV.



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# Higgs Decay Ratios at the HL-LHC: $D_{\gamma\gamma}$

In Higgs physics at the (HL-)LHC, **production** channels are the main source of theoretical **uncertainty** → a promising way of diminishing this error is to consider **ratios** of Higgs production times decay rates [1208.3436]. A **well-motivated** example is:

$$D_{\gamma\gamma} \equiv \frac{\sigma(pp \rightarrow H \rightarrow \gamma\gamma)}{\sigma(pp \rightarrow H \rightarrow ZZ^*)} \simeq \frac{\Gamma(H \rightarrow \gamma\gamma)}{\Gamma(H \rightarrow ZZ^*)}.$$

- $H \rightarrow \gamma\gamma$  and  $H \rightarrow ZZ^*$  **precisely** measured;
- $D_{\gamma\gamma}$  sensitive especially to VLQs with **high electric charge**;
- **Challenge** → in such ratios, the **same kinematical cuts** should be applied when measuring the two processes ⇒ in this ideal case, the production cross section cancels (as well as the Higgs total width)! In the following, we will assume that this indeed happens.

# VLQ Mass Reach Through $D_{\gamma\gamma}$

In the following, to get an idea of the power of  $D_{\gamma\gamma}$  for indirect detection of VLQs, we will focus on simplified, generic VLQ models. Also, we no longer try to explain the  $t\bar{t}H$  and  $A_{\text{FB}}^b$  deviations from the SM.

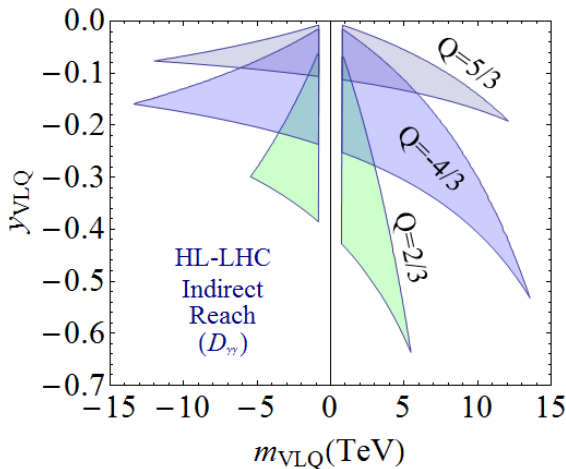
- We consider 2 VLQ multiplets that mix between themselves, with one of them significantly heavier than the other;
- Set their mixing with SM quarks to 0.

We assume (rather optimistically) that the experimental error at the HL-LHC will be:

$$\frac{\Delta D_{\gamma\gamma}}{D_{\gamma\gamma}} = 1\%,$$

at 68% C.L., with the experimental central value of  $D_{\gamma\gamma}$  taken to be equal to its SM value at LO.

# VLQ Mass Reach Through $D_{\gamma\gamma}$ : Numerical Results



**Figure :** Mass ( $m_{VLQ}$ ) vs mass basis Yukawa coupling ( $y_{VLQ}$ ) of the lighter VL quark.  $Q$  stands for electric charge.

# VLQ Mass Reach Through $D_{bb}$

What about **bottom-like VLQs**? A good way to search indirectly for them is to consider the **ratio**:

$$D_{bb} \equiv \frac{\sigma(pp \rightarrow VH \rightarrow Vb\bar{b})}{\sigma(pp \rightarrow VH \rightarrow VZZ^*)} \simeq \frac{\Gamma(H \rightarrow b\bar{b})}{\Gamma(H \rightarrow ZZ^*)}.$$

Again, we study a **generic model** where only **one bottom-like** VLQ singlet mixes with the SM-like  $b$  quark. We assume that, at 68% C.L., the HL-LHC experimental error will be:

$$\frac{\Delta D_{bb}}{D_{bb}} = 5\%,$$

with the experimental central value taken equal to its SM value at LO.

# VLQ Mass Reach Through $D_{bb}$ : Numerical Results

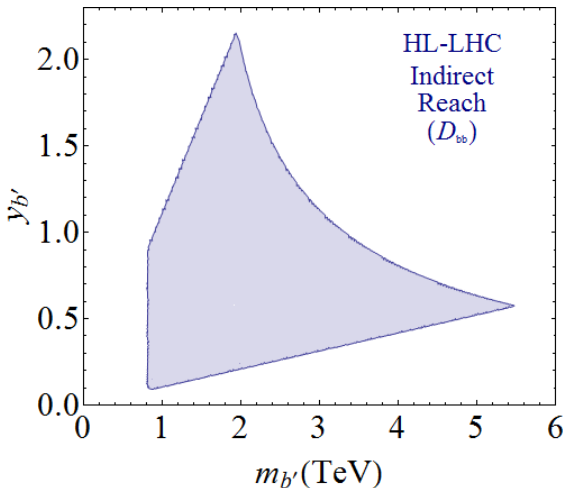


Figure : Mass vs mass basis Yukawa coupling of the  $b'$  VLQ.

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# Conclusions

- By adding **VLQ** multiplets to the SM, the present experimental anomalies in the heavy quark sector ( $A_{\text{FB}}^b$  and  $t\bar{t}H$ ) can be **explained**, while satisfying the **constraints** coming from EWPT at **LEP** and from Higgs physics at the **LHC**.
- Interestingly, the obtained **VLQ** model predicts top, bottom and exotic VL partners with masses typically at the **TeV scale**, which can be searched for at **Run 2** of the LHC.
- By considering **ratios** of Higgs production times decay rates, we showed that top and bottom VL partners with masses up to **5 TeV** and more exotic **VLQs** with masses in the **10 TeV** range can be indirectly probed at the **HL-LHC**. This is to be compared with the mass reach through direct searches, which is estimated to be around  $1.5 - 2 \text{ TeV}$  [1307.7135].



# Backup: $H \rightarrow \gamma\gamma$ Formulas

$$\Gamma(H \rightarrow \gamma\gamma) \propto \left| A_1(\tau_W) + N_c \sum_f Q_f^2 \sum_i \left( \frac{y_i}{m_i} A_{1/2}(\tau_i) \right)_f \right|^2.$$

Example: top quark plus top-like VLQs:

$$\begin{aligned} \sum_i \left( \frac{y_i}{m_i} A_{1/2}(\tau_i) \right)_t &\simeq \frac{4}{3} \sum_i \left( \frac{y_i}{m_i} \right)_t = \frac{4}{3} \text{tr} \left( \mathcal{M}_t^{-1} \frac{\partial \mathcal{M}_t}{\partial v} \right) \\ &= \frac{4}{3} \partial_v (\log \det \mathcal{M}_t). \end{aligned}$$

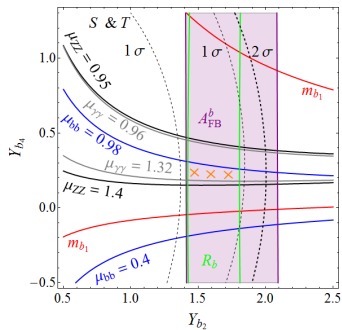
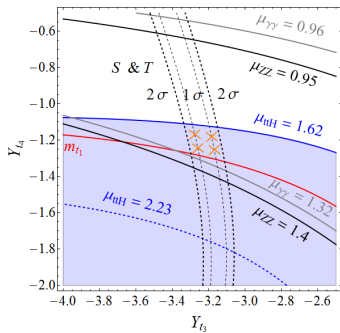
If there is only one extra top-like VLQ:

$$\mathcal{M}_t = \frac{1}{\sqrt{2}} \begin{pmatrix} vY_{t_1} & vY_{t_2} \\ 0 & \sqrt{2}m \end{pmatrix} \Rightarrow \partial_v (\log \det \mathcal{M}_t) = \frac{1}{v} = \left( \frac{y_t}{m_t} \right)_{\text{SM}}.$$

# Backup: Model B

Model **B** = Model **A** with the replacement

$$b'' \rightarrow X_{L,R} = (t''', b'', q'_{-4/3}) \sim (3, 3, -1/3).$$



# Backup: Model C

Model **C** = Model **A** with the addition of:

$$Z_{L,R} = (q_{8/3}, q_{5/3}, t''') \sim (3, 3, 5/3).$$

