

# Analytic control of jet substructure

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# Presentation plan

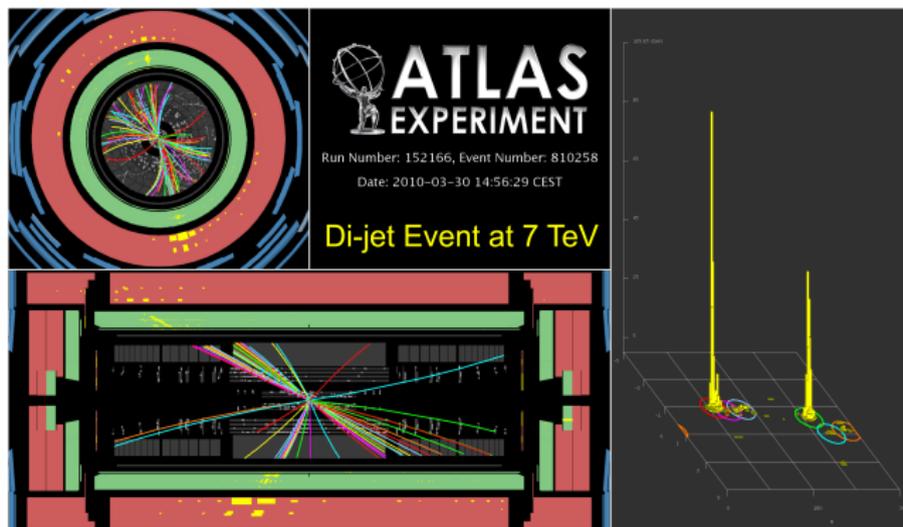
- 1 Introduction
- 2 Jet substructure
- 3 Control of jet substructure
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# Introduction

- QCD partons in high energies (LHC) can not be directly observed and their final state are complex structures called *jets*;

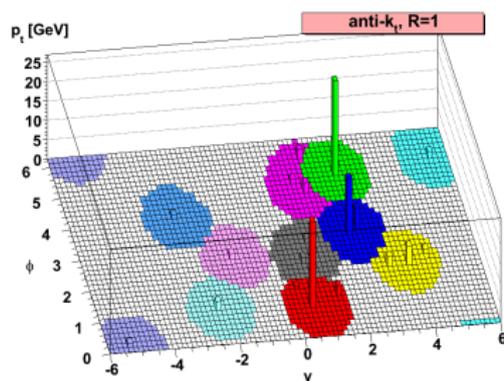


ATLAS collaboration

- Jets are used very frequently in LHC analysis.

# Definition of jet

- A *jet definition* is a set of rules to cluster particles into jets;
- Composed of a *recombination algorithm* and its *parameters* (usually the jet radius  $R$ );



M. Cacciari, G. P. Salam and G. Soyez (2008)

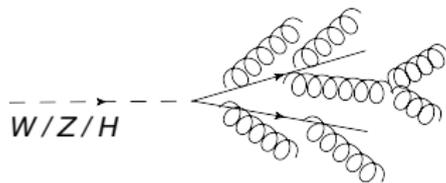
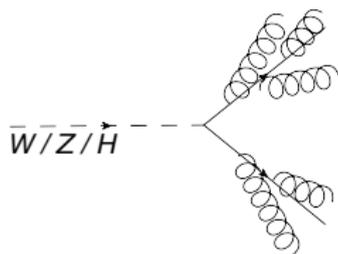
- LHC uses the anti- $k_t$  algorithm.

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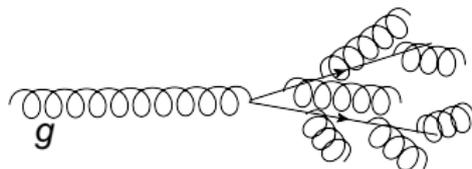
# Boosted heavy particles

- At the LHC  $\rightarrow$  boosted ( $p_t \gg m$ ) heavy particles
  - $\rightarrow$  decay in very collimated final states
  - $\rightarrow$  clustered into a single jet;

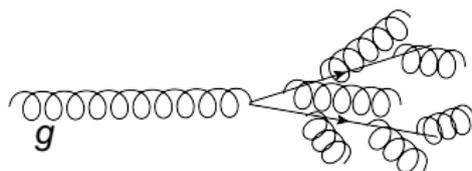


- Characteristic opening angle of the jet is  $\theta \propto \frac{m}{p_t}$ .

- QCD collinear divergences keep parton jets collimated for any  $p_{T,i}$

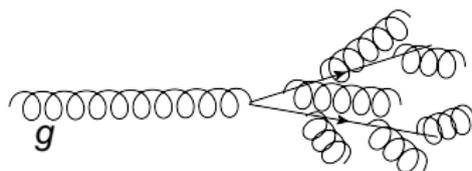


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How to discriminate QCD jets from  $Z/W/H \rightarrow \text{hadrons}$  jets?

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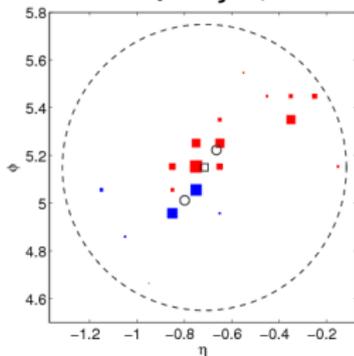
How to discriminate QCD jets from  $Z/W/H \rightarrow \text{hadrons}$  jets?

Possible method : partons are *color charged particles*, they produce different radiation patterns from bosons.

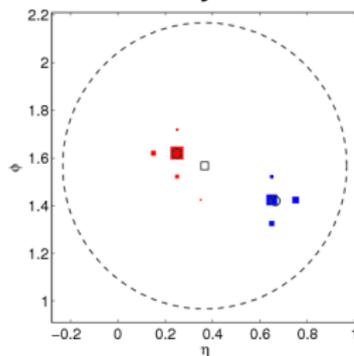
# Jet substructure

- Knowing the mass of a jet is not enough to identify its origin, we need to access the *jet substructure* ;

Boosted QCD Jet,  $R = 0.6$



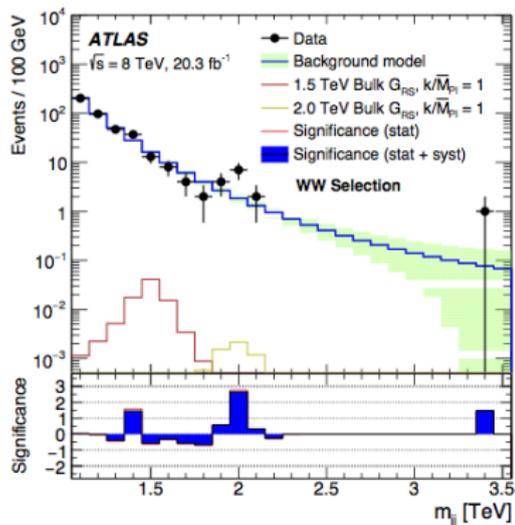
Boosted W Jet,  $R = 0.6$



J. Thaler and K. V. Tilburg (2010)

- Different techniques are available:
  - Find hard cores (1 for QCD, 2 for bosons);
  - Constrain the radiation patterns.

# Application : diboson excess



ATLAS collaboration

- ATLAS found a  $3.2\sigma$  excess in diboson resonances at  $\sim 2\text{TeV}$ ;
- CMS found similar anomalies;
- High energy events  $\rightarrow$  very collimated final states.

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- *Constraint radiation patterns to discriminate different types of jets;*
- Making use of jet-shapes : observables that are functions of the jet constituents  $v(p_1^\mu, p_2^\mu, \dots, p_n^\mu)$ ;
- In this work we chose  $v = \tau_{21}, \mu^2$  and  $C_2$  (next slide);
- Understand differences/similarities from a *first-principle analytical study*.

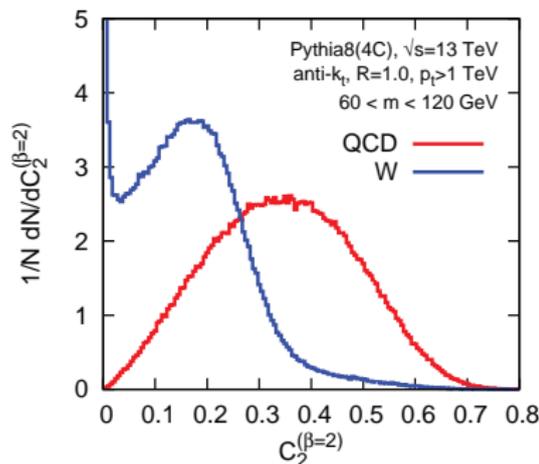
# Jet shapes (angular parameter $\beta = 2$ )

- Energy correlation

$$C_2^{(2)} = e_3^{(2)} / (e_2^{(2)})^2,$$

$$e_2^\beta = \frac{1}{p_t^2 R^\beta} \sum_{i < j \in \text{jet}} p_{t,i} p_{t,j} \theta_{ij}^\beta,$$

$$e_3^\beta = \frac{1}{p_t^3 R^{3\beta}} \sum_{i < j < k \in \text{jet}} p_{t,i} p_{t,j} p_{t,k} \theta_{ij}^\beta \theta_{ik}^\beta \theta_{jk}^\beta.$$



# Jet shapes (angular parameter $\beta = 2$ )

- **N-subjettiness** with axes  $a_1, \dots, a_N$

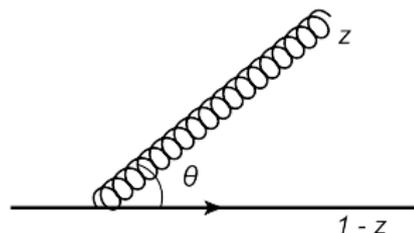
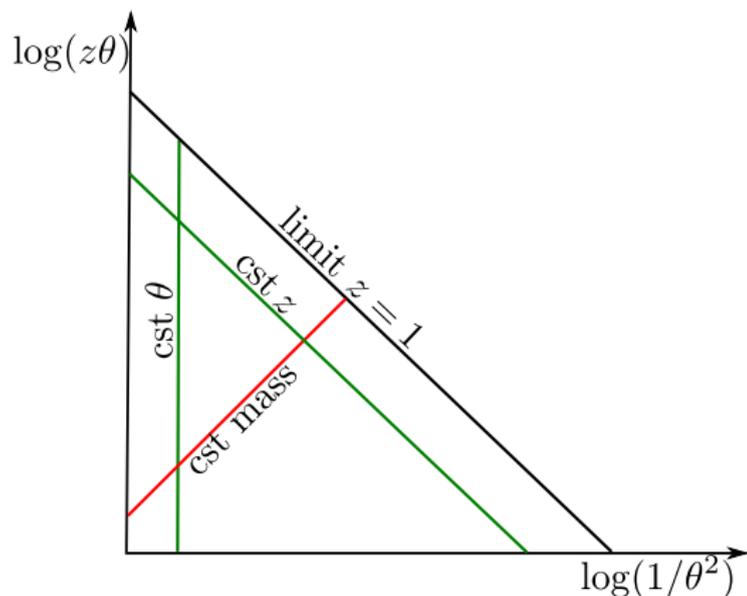
$$\tau_{21}^{(2)} = \frac{\tau_2^{(2)}}{\tau_1^{(2)}}, \quad \tau_N^\beta = \frac{1}{p_{t,\text{jet}} R^\beta} \sum_{i \in \text{jet}} p_{t,i} \min_{a_1 \dots a_N} (\theta_{ia_1}^\beta, \dots, \theta_{ia_N}^\beta).$$

- **Mass-drop** with subjets  $j_1$  and  $j_2$

$$\mu_p^2 = \max(m_{j_1}^2, m_{j_2}^2) / m_j^2.$$

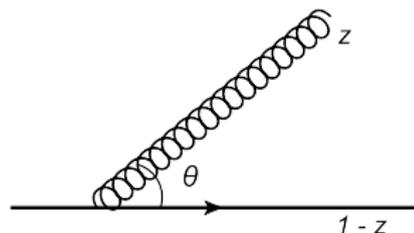
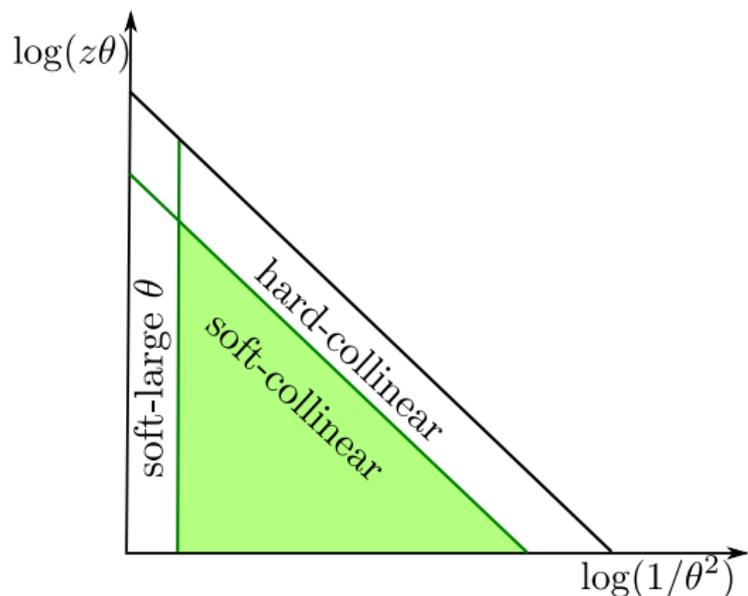
# Lund diagrams

- We consider boosted jets of a *given mass*,  $\rho = m^2/p_t^2 R^2 \ll 1$ ;
- Lund diagram : graphical representation of the results in  $z\theta$  vs.  $1/\theta^2$  coordinates.



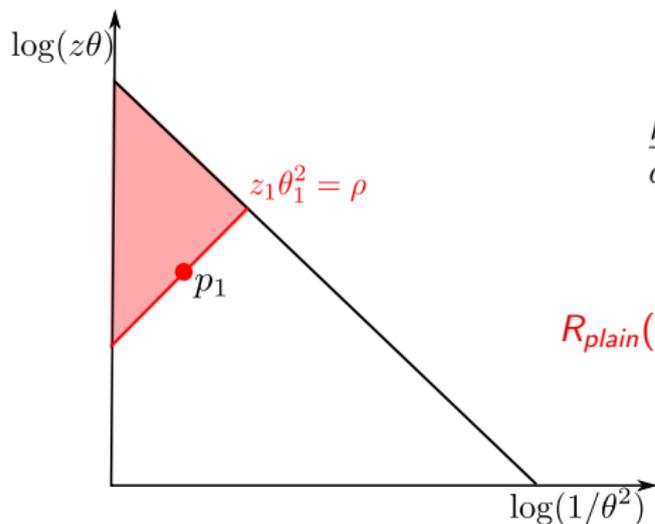
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# Structure of the results (QCD background)

- Up to LL, emissions are strongly ordered in mass and angle.
- Independent emissions  $\rightarrow$  constraints as an exponential factor.
- For a jet of a given mass:

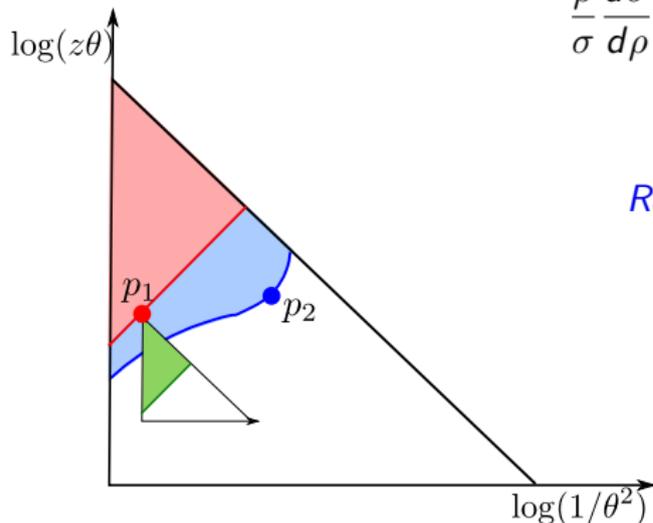


$$\frac{\rho}{\sigma} \frac{d\sigma}{d\rho} = \int_{\rho}^1 dz_1 P_i(z_1) \frac{\alpha_s}{2\pi} e^{-R_{plain}(\rho)}$$

$$R_{plain}(\rho) = \int_0^1 \frac{d\theta^2}{\theta^2} \int_0^1 dz P_i(z) \frac{\alpha_s}{2\pi} \Theta(z\theta^2 > \rho)$$

# Structure of the results (QCD background)

- For a jet of a given mass + a cut in the shape  $v_{cut} \ll 1$ :

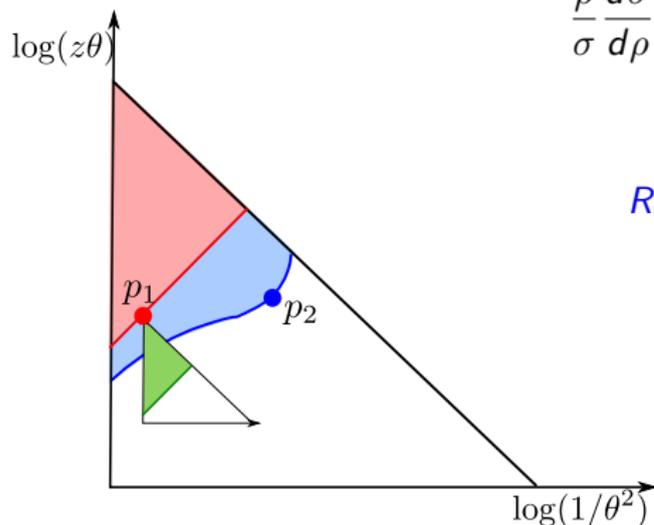


$$\left. \frac{\rho}{\sigma} \frac{d\sigma}{d\rho} \right|_{<v} = \int_{\rho}^1 dz_1 P(z_1) \frac{\alpha_s}{2\pi} e^{-R_{plain}(\rho) - R_v(\rho, z_1)}$$

$$\begin{aligned} R_v(\rho, z_1) &= \int_0^1 \frac{d\theta_2^2}{\theta_2^2} \int_0^1 dz_2 P_i(z_2) \frac{\alpha_s}{2\pi} \\ &\quad \times \Theta(v > v_{cut}) \Theta(z_2 \theta_2^2 < \rho) \\ &+ \int_0^1 \frac{d\theta_{12}^2}{\theta_{12}^2} \int_0^1 dz_2 P_g(z_2) \frac{\alpha_s}{2\pi} \\ &\quad \times \Theta(v^{sec} > v_{cut}) \end{aligned}$$

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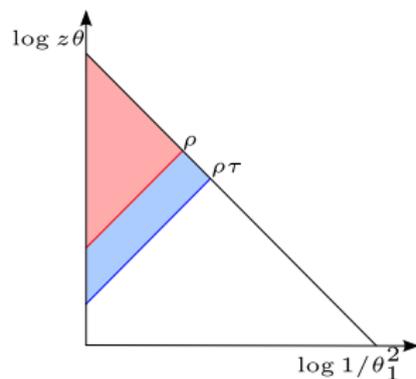
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Now all we need is to find  $v(\rho, z_1, z_2, \theta_2)$ .

# Results

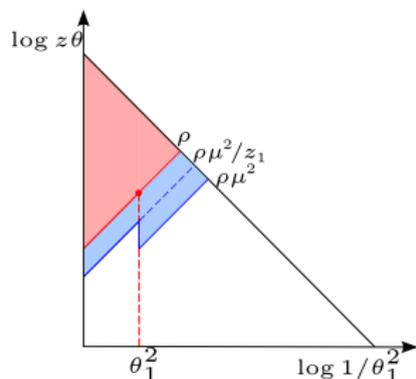
## *N-subjettiness*



$$R_\tau(z_1) = \frac{\alpha_s C_R}{2\pi} \left[ \frac{L_\tau^2}{2} + L_\rho L_\tau \right] + \frac{\alpha_s C_A}{2\pi} \frac{L_\tau^2}{2}$$

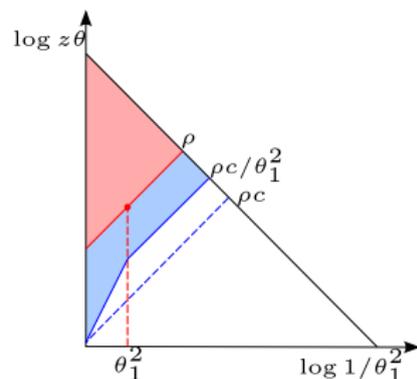
$$L_X = \log(1/X)$$

## *Mass drop*



$$R_{\mu_{1/2}^2}(z_1) = \frac{\alpha_s C_R}{2\pi} \left[ \frac{L_\mu^2}{2} + L_\rho L_\mu \right] - \frac{\alpha_s C_R}{2\pi} \frac{L_1}{2} (L_\rho - L_1) + \frac{\alpha_s C_A}{2\pi} \frac{(L_\mu - L_1)^2}{2}$$

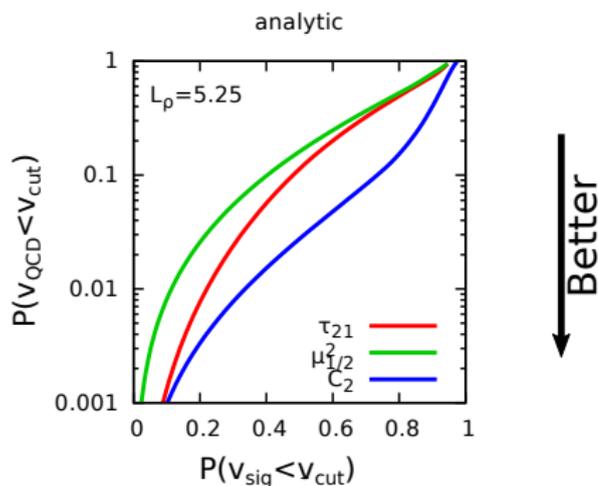
## *Energy correlation*



$$R_{C^2}(z_1) = \frac{\alpha_s C_R}{2\pi} \left[ \frac{L_e^2}{2} + (L_e - L_\rho + L_1)L_1 \right] + \frac{\alpha_s C_A}{2\pi} \frac{1}{2} (L_e - L_\rho + L_1)^2$$

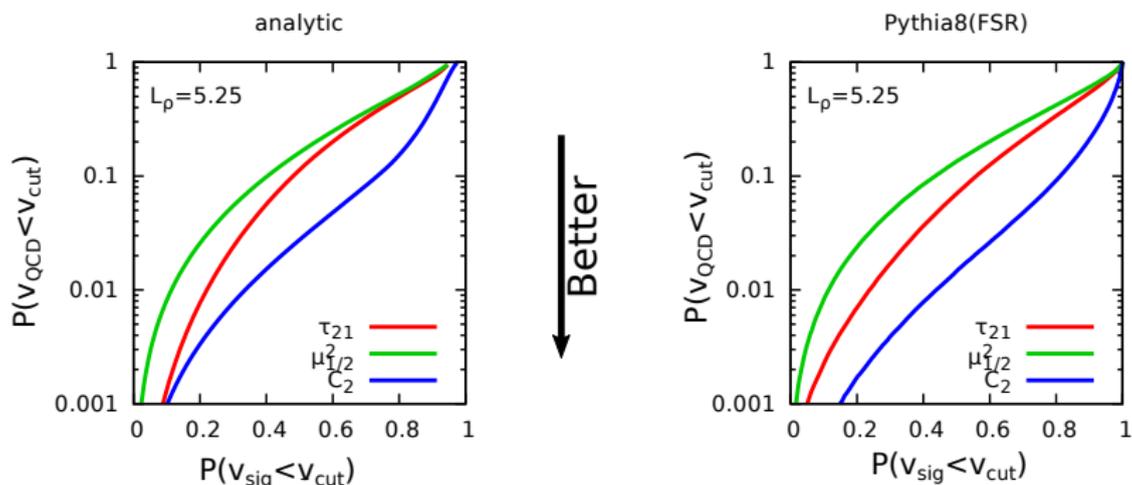
# ROC curves

- Probability that  $v_{QCD} < v_{cut}$  vs. probability that  $v_{sig} < v_{cut}$ ;
- $C_2$  is the most efficient, and  $\tau_{21}$  more efficient than  $\mu^2$  (more delicate call).



# ROC curves

- Probability that  $v_{QCD} < v_{cut}$  vs. probability that  $v_{sig} < v_{cut}$ ;
- $C_2$  is the most efficient, and  $\tau_{21}$  more efficient than  $\mu^2$  (more delicate call).



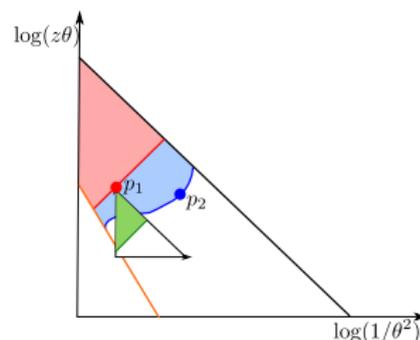
- Good description of the order between shapes.

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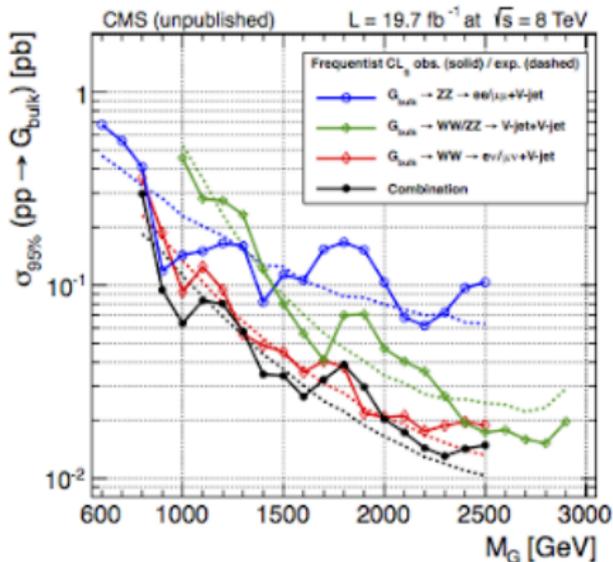
# Conclusion

- Good **qualitative description** of the shapes.
- **Efficiency** of the shapes:  $C_2 > \tau_{21} \gtrsim \mu^2$ .
- **Future works**
  - Add grooming (in progress);
  - Higher accuracy (in progress);
  - Different jet shapes (next);
  - 3-pronged jet shapes;
  - ISR contributions;
  - Calculations for finite  $v$ .



# Backup Slides

# Diboson excess at CMS

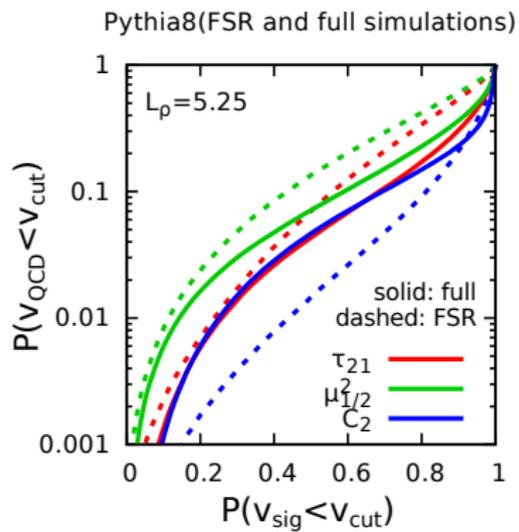
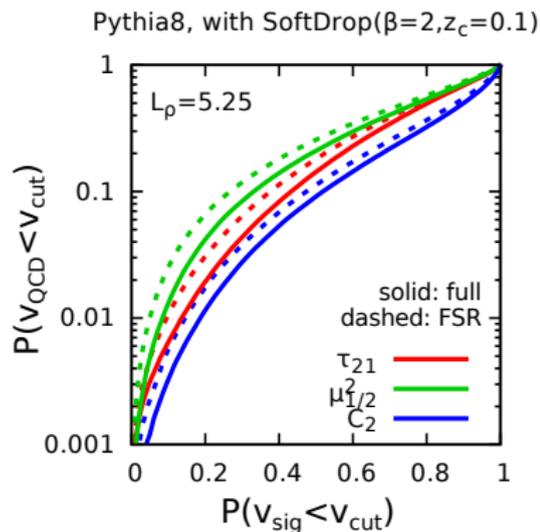


CMS collaboration

# Generalized $k_t$ algorithm

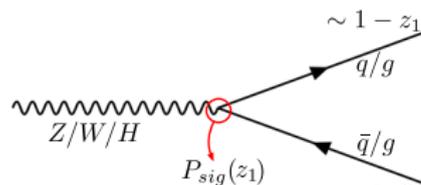
- Depends on a parameter  $p$ ;
- Cluster partons by smallest distance  $d_{ij} = \min(z_i^{2p}, z_j^{2p})\theta_{ij}^2$ ;
- Particular cases:
  - $p=0$  : C/A algorithm, angular ordered;
  - $p=-1$  : anti- $k_t$  algorithm;
  - $p=1/2$  : similar to mass measure.

# Non-perturbative effects



# Structure of the results (Signal)

- Decay of a boosted object into a pair  $q\bar{q}$  or  $gg$ .
- For a signal jet (fixed mass always) + a cut in the shape  $v$ :



$$\log(z\theta) \quad \Sigma_{sig} = \int_{\rho}^1 dz_1 P_{sig}(z_1) e^{-R_{v,sig}(z_1, \rho) - R_{v,sig}(1-z_1, \rho)}$$

$$R_{v,sig}(z_1) = \int_0^1 \frac{d\theta_2^2}{\theta_2^2} \int_0^1 dz_2 P_i(z_2) \frac{\alpha_s}{2\pi} \Theta(v^{sig} > v_{cut})$$

