Impact on New Physics scenarios from CP Violation measurements

Lyon, 24 -11- 2015

Yasmine Amhis on behalf of the LHCb Collaboration

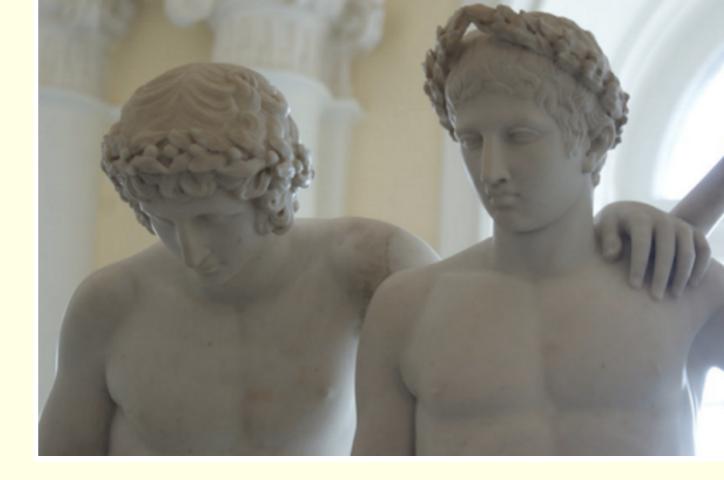


#### High $p_T$ questions

- What is the mechanism of electroweak symmetry breaking?
- What separates the electroweak scale from the Planck scale?
- What happened at the electroweak phase transition?
- What are the dark matter particles?
- How was the baryon asymmetry generated?

 $Flavor \Leftrightarrow Collider$ 

Y.Nir CERN Flavour Physics Seminars.



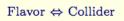
#### The Questions

#### **Flavor questions**

- The Standard Model flavor puzzle: Why are the flavor parameters small and hierarchical? (Why) are the neutrino flavor parameters different?
- The New Physics flavor puzzle: If there is NP at the TeV scale, why are FCNC so small? The solution ⇒ Clues for the subtle structure of the NP
- Are the two puzzles related?

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Indirect Searches – Model Independent Searches Four examples of how to look for New Physics

How can New Physics affect angular observables ?

How can New Physics enhance a suppressed decay ?

How can New Physics affect an oscillation?

How can New Physics affect CP violation ?  $\sin 2\beta$ ,  $\phi_s$ 

### Classification of CP Violating effects

A large set of NP models predict the existence of new CP Violating phases

CPV in decay

CPV in mixing

CPV in the interference mixing & decay

### What is the Standard Model telling us?

$$V_{\rm CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

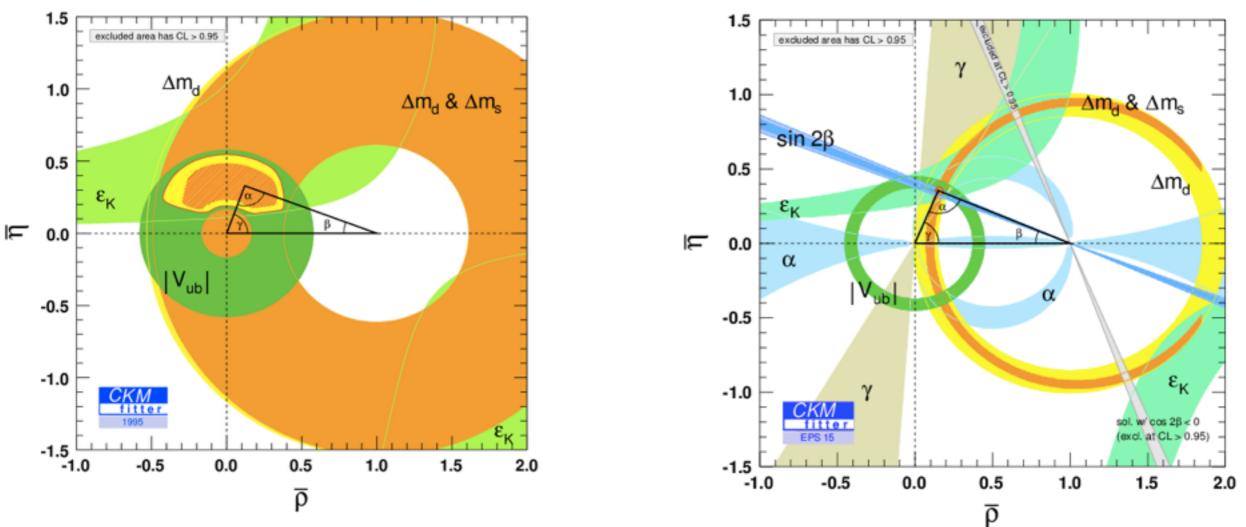
$$\begin{pmatrix} 1 - \frac{1}{2}\lambda^2 - \frac{1}{8}\lambda^4 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda + \frac{1}{2}A^2\lambda^5[1 - 2(\rho + i\eta)] & 1 - \frac{1}{2}\lambda^2 - \frac{1}{8}\lambda^4(1 + 4A^2) & A\lambda^2 \\ A\lambda^3[1 - (1 - \frac{1}{2}\lambda^2)(\rho + i\eta)] & -A\lambda^2 + \frac{1}{2}A\lambda^4[1 - 2(\rho + i\eta)] & 1 - \frac{1}{2}A^2\lambda^4 \end{pmatrix}$$

We have to constrain as much as possible all the parameters

# What is real life telling us ?

2015

1995



The CKM mechanism seems to work

### What is New Physics telling us?



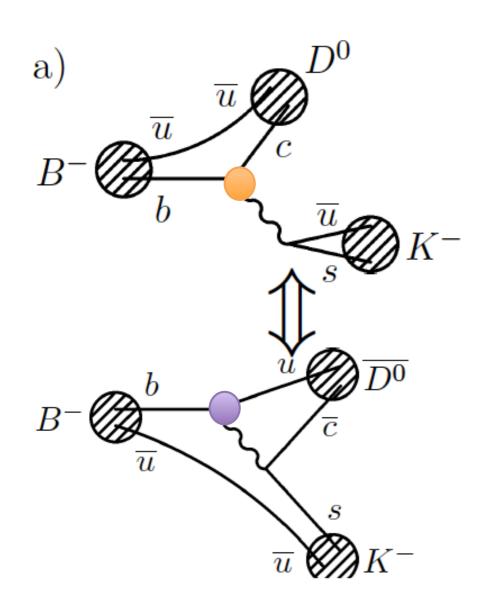
### Where shall we start?



### Measure theoretically clean observables

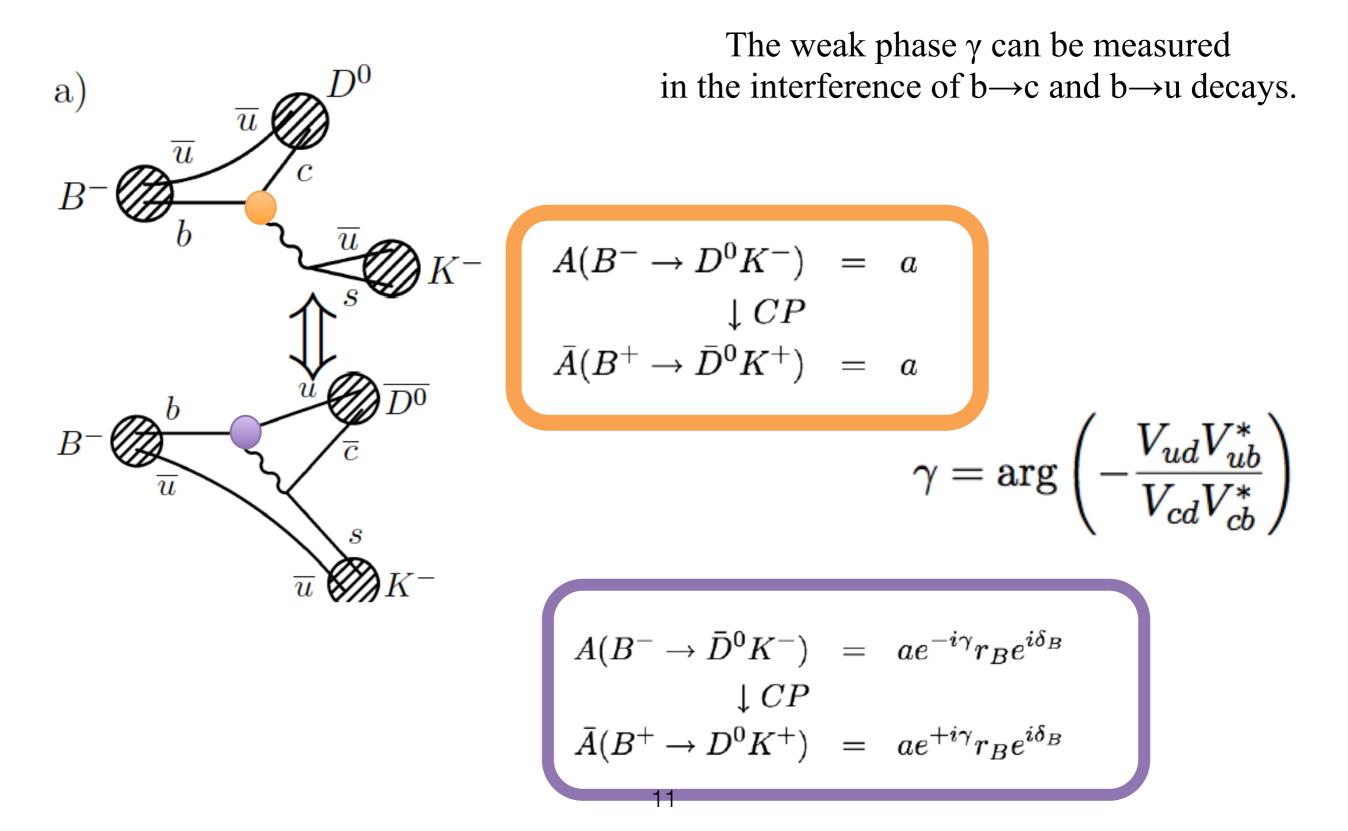
### Measurement of $\gamma$

The weak phase  $\gamma$  can be measured in the interference of b $\rightarrow$ c and b $\rightarrow$ u decays.



$$\gamma = \arg\left(-rac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*}
ight)$$

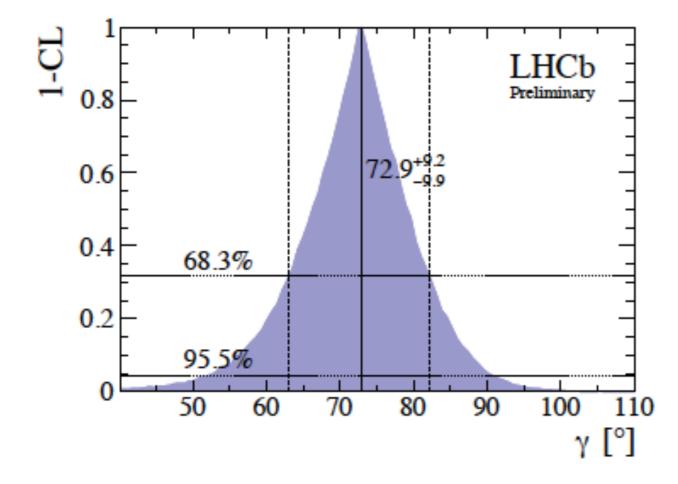
### Measurement of $\gamma$



$$B^+ \to Dh^+, D \to hh, \, {\rm GLW}/{\rm ADS}, \, 1 \, {\rm fb}^{-1}$$

- $B^+ \to Dh^+, D \to K\pi\pi\pi, ADS, 1\,\mathrm{fb}^{-1}$
- $B^+ \to DK^+, D \to K^0_{\rm s} hh$ , model-independent GGSZ, 3 fb<sup>-1</sup>
- $B^+ \to DK^+, D \to K^0_{\rm s} K \pi, \, {\rm GLS}, \, 3 \, {\rm fb}^{-1}$
- $B^0 \rightarrow DK^{*0}, D \rightarrow hh, \, {\rm GLW}/{\rm ADS}, \, 3 \, {\rm fb}^{-1}$
- $B_s^0 \to D_s^{\mp} K^{\pm}$ , time-dependent, 1 fb<sup>-1</sup>





Other measurements and updates will be added soon !

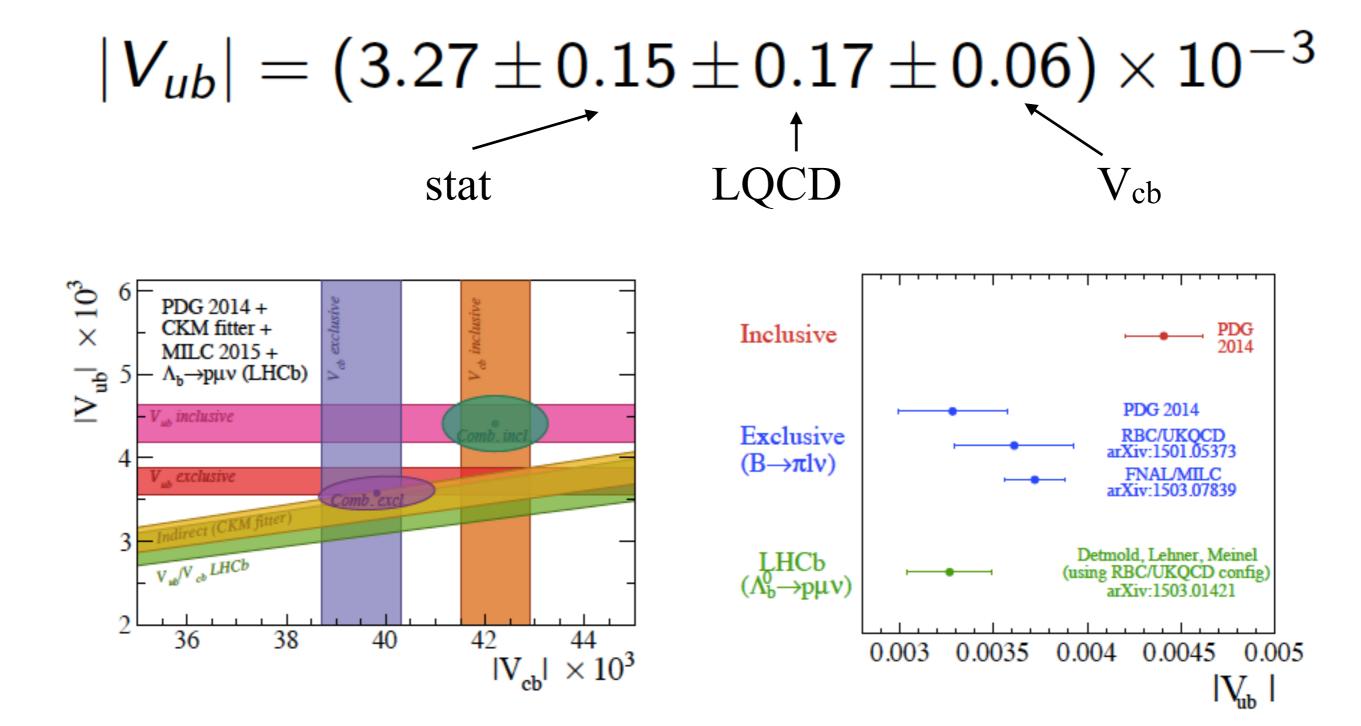
LHCb-CONF-2014-004

### What about $|V_{ub}|$ ?

 $V_{\rm CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$ 

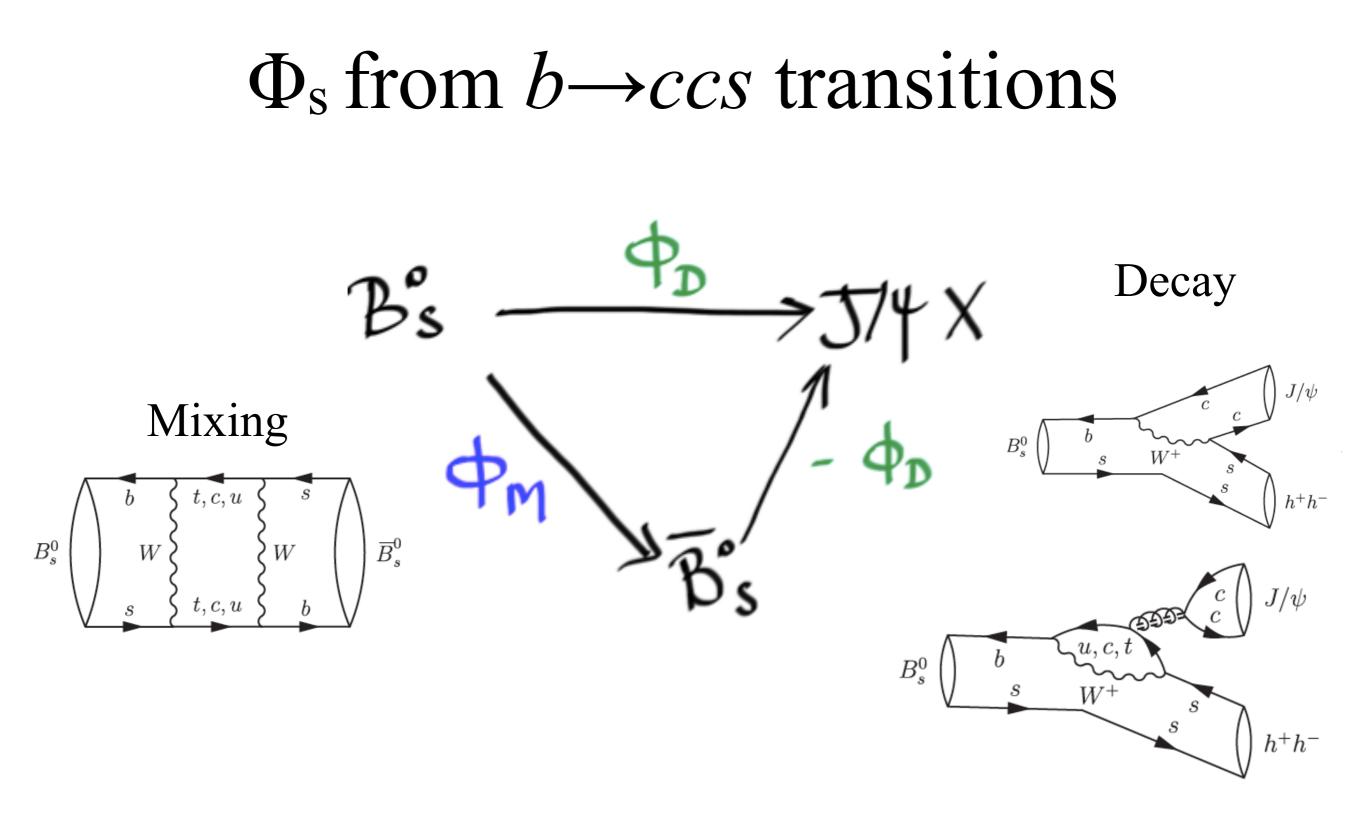
Measurement of V<sub>ub</sub> [3 fb<sup>-1</sup>] Þ  $\frac{|V_{ub}|^2}{|V_{cb}|^2} = \frac{\mathcal{B}(\Lambda_b^0 \to p\mu^- \overline{\nu}_\mu)}{\mathcal{B}(\Lambda_b^0 \to \Lambda_c^+ \mu^- \overline{\nu}_\mu)} R_{\rm FF}$ What we measure : Topologies  $X_{\overline{b}}$  $X_{\overline{b}}$ pP٧  $\Lambda_c^+$ P۱  $\Lambda_b^o$  $\Lambda_b^0$  $\bar{\nu}_{\mu}$  $\mu$  $\bar{\nu}_{\mu}$ 

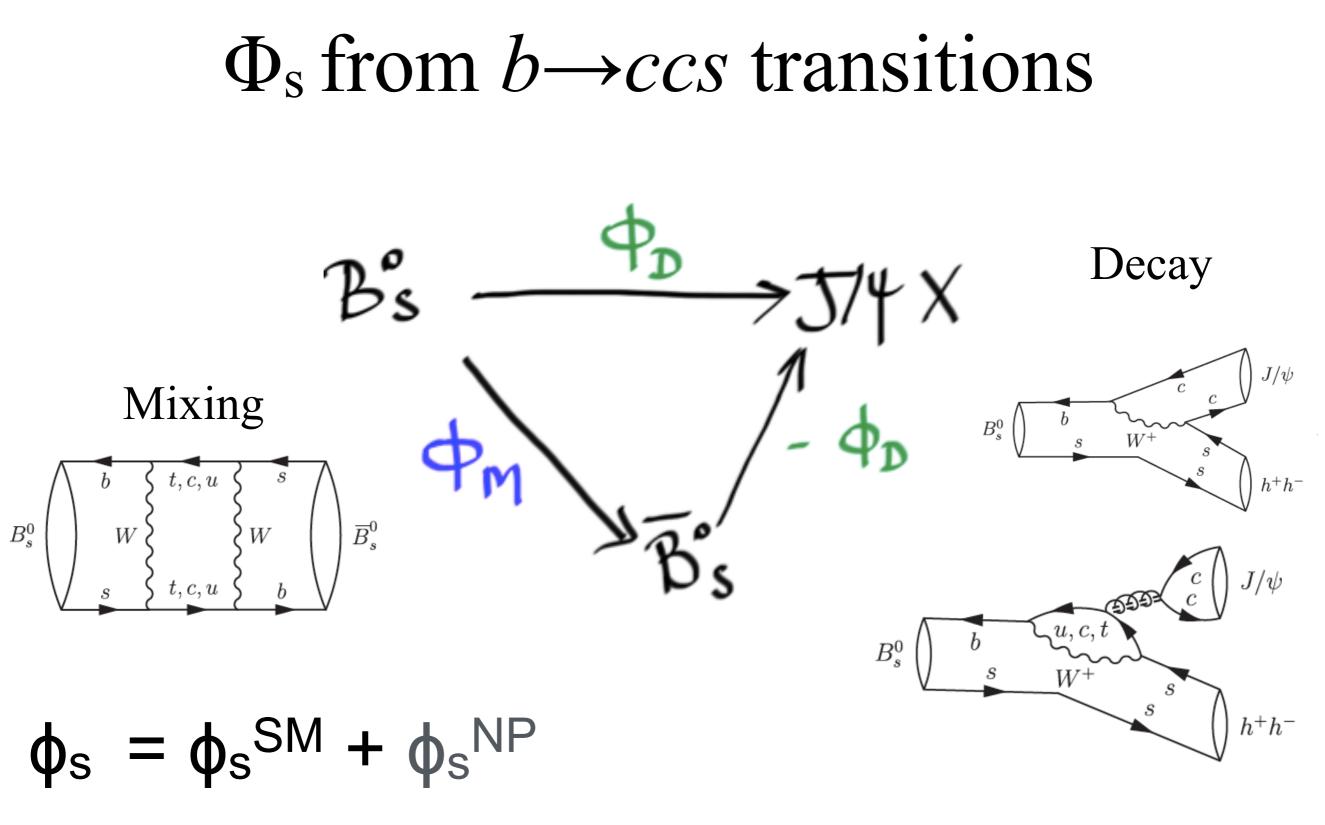
Result :



[Nature Physics 3415 (2015)] [3 fb<sup>-1</sup>]

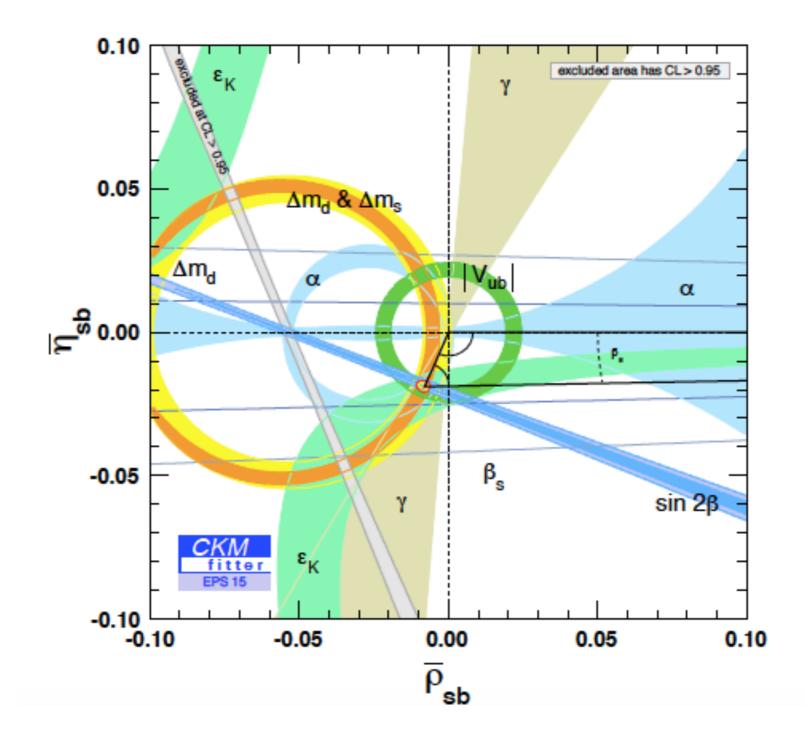
15





 $\phi_s^{SM} = \phi_M$  Predicted to be very small in the SM Phys. Rev D84 (2001) 033005

# The other triangle



## What do we measure ?

$$\mathcal{A}_{CP}(t) = \frac{\Gamma\left(\overline{B}_q^0(t) \to f\right) - \Gamma\left(B_q^0(t) \to f\right)}{\Gamma\left(\overline{B}_q^0(t) \to f\right) + \Gamma\left(B_q^0(t) \to f\right)} = \frac{\mathcal{S}_f \sin\left(\Delta mt\right) - \mathcal{C}_f \cos\left(\Delta mt\right)}{\cosh\left(\frac{\Delta\Gamma t}{2}\right) + \mathcal{A}_{\Delta\Gamma} \sinh\left(\frac{\Delta\Gamma t}{2}\right)}$$

$$\Delta m = m_{\rm H} - m_{\rm L}$$
  $\Delta \Gamma = \Gamma_{\rm L} - \Gamma_{\rm H}$ 

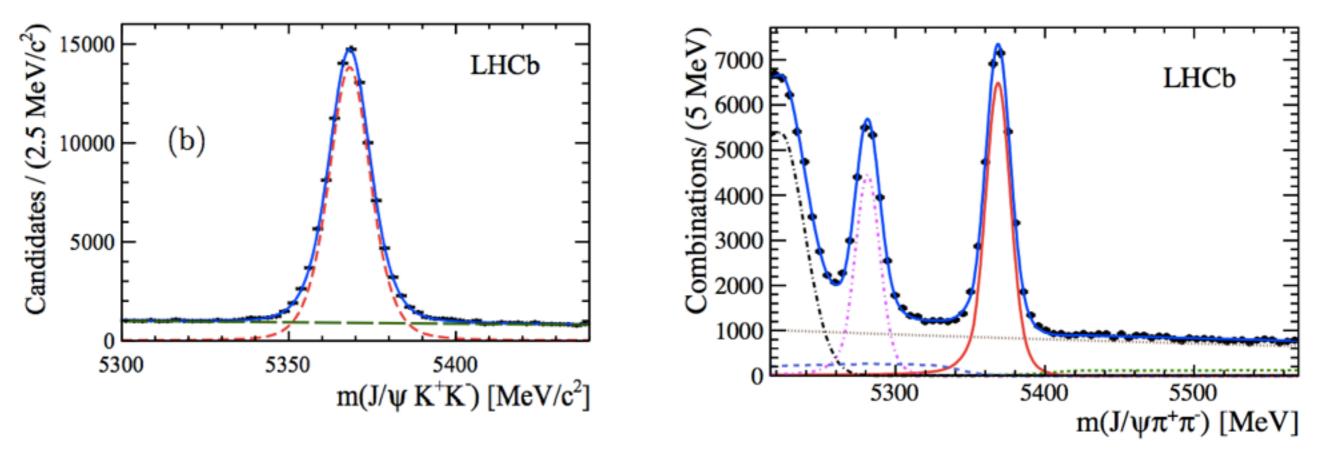
### Mixing parameters

$$S_{f} = \frac{2\Im(\lambda_{f})}{1+|\lambda_{f}|^{2}}, \quad C_{f} = \frac{1-|\lambda_{f}|^{2}}{1+|\lambda_{f}|^{2}}, \quad \mathcal{A}_{\Delta\Gamma} = -\frac{2\Re(\lambda_{f})}{1+|\lambda_{f}|^{2}}$$
$$\lambda_{f} = \eta_{f} \frac{q}{p} \frac{A(\overline{B}_{q}^{0}(t) \to f)}{A(B_{q}^{0}(t) \to f)} = \eta_{f} |\lambda_{f}| e^{i\phi_{q}} \qquad CP \text{ observables}$$

$$\mathcal{S}_f = -\eta_f \frac{2|\lambda_f|\sin(\phi_q)}{1+|\lambda_f|^2} , \quad \mathcal{A}_{\Delta\Gamma} = \eta_f \frac{2|\lambda_f|\cos(\phi_q)}{1+|\lambda_f|^2}$$

### Measurement of $\Phi_s$ in $B_s \rightarrow J/\psi h^+h^- [3 fb^{-1}]$

The golden modes

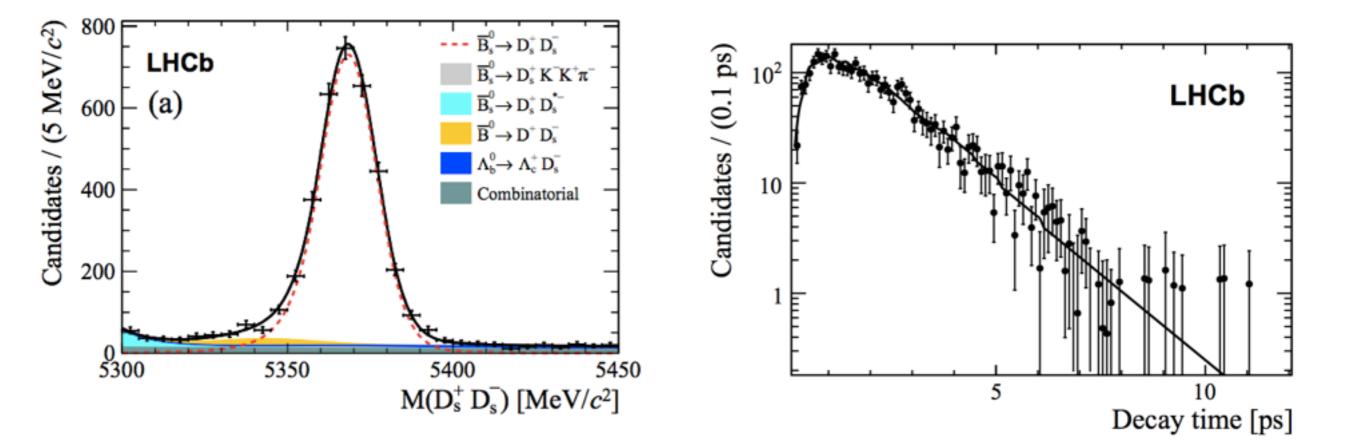


Use angular analysis to separate the CP components.

[Phys. Rev. Lett. 114 041801 (2015)]

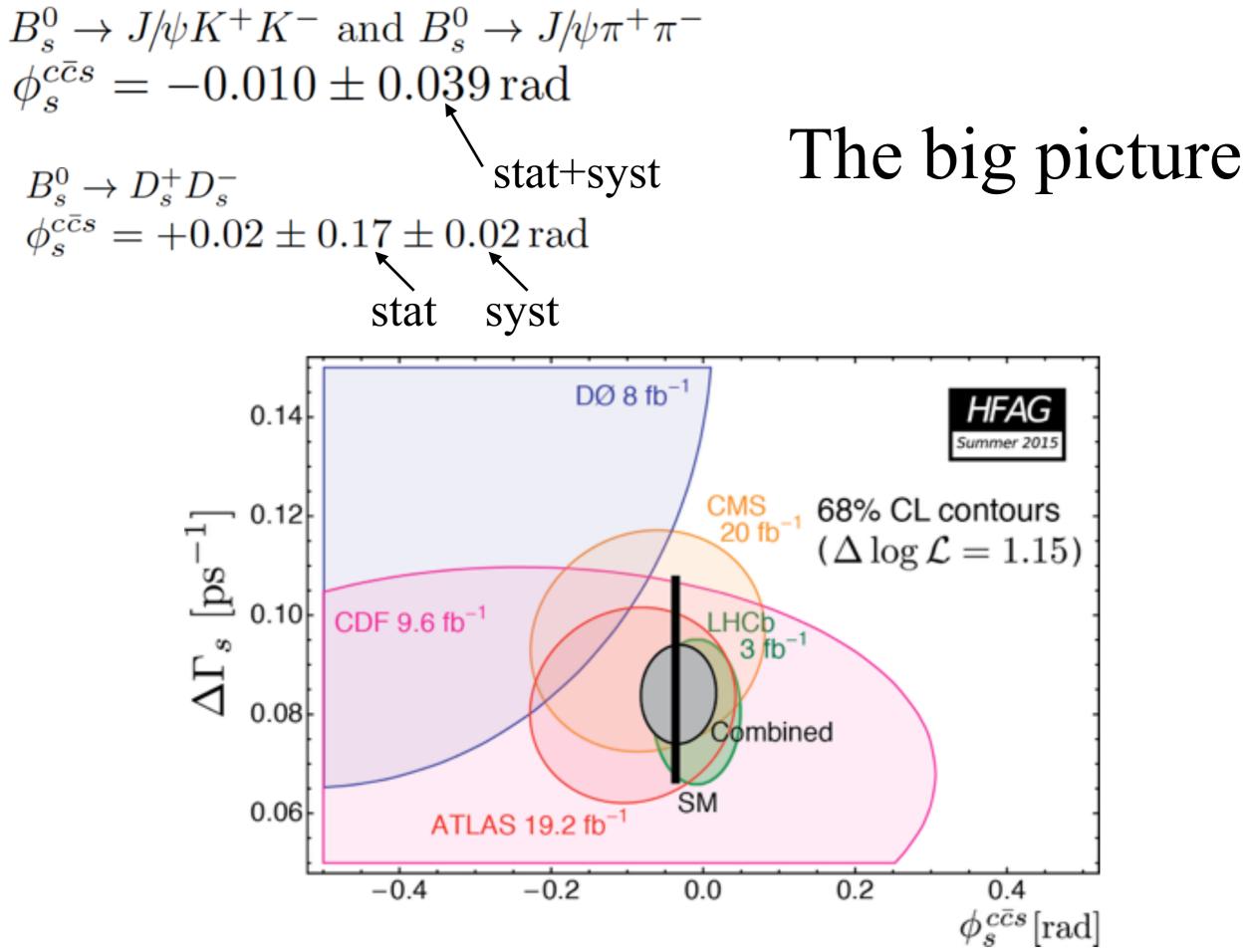
[Phys. Lett. B 736 (2014) 186]

### Measurement of $\Phi_s$ in $B_s \rightarrow D_s D_s [3 \text{ fb}^{-1}]$

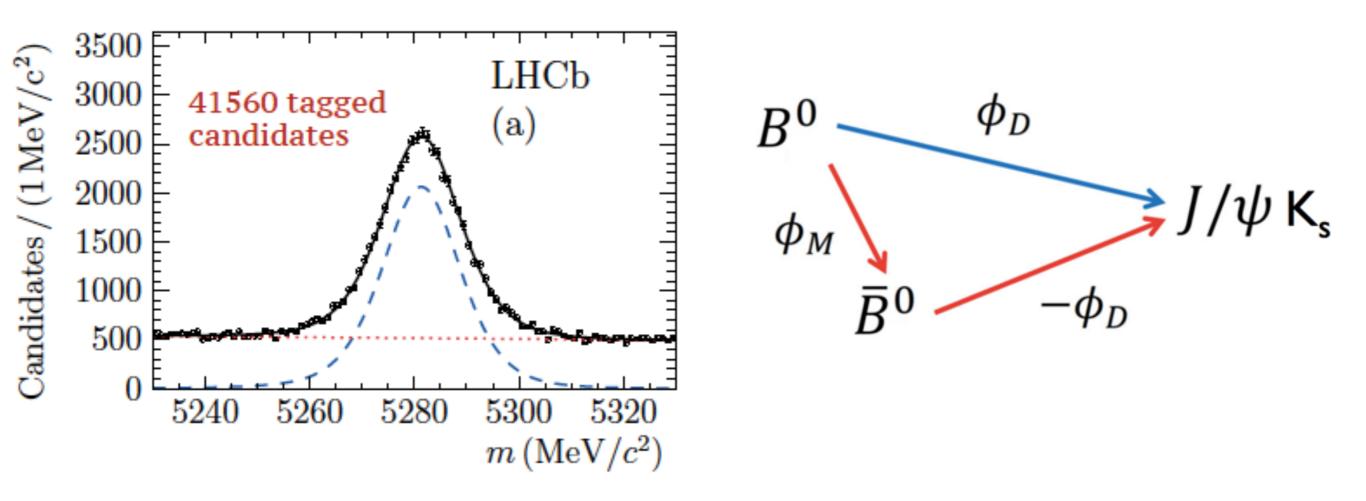


CP-even final state

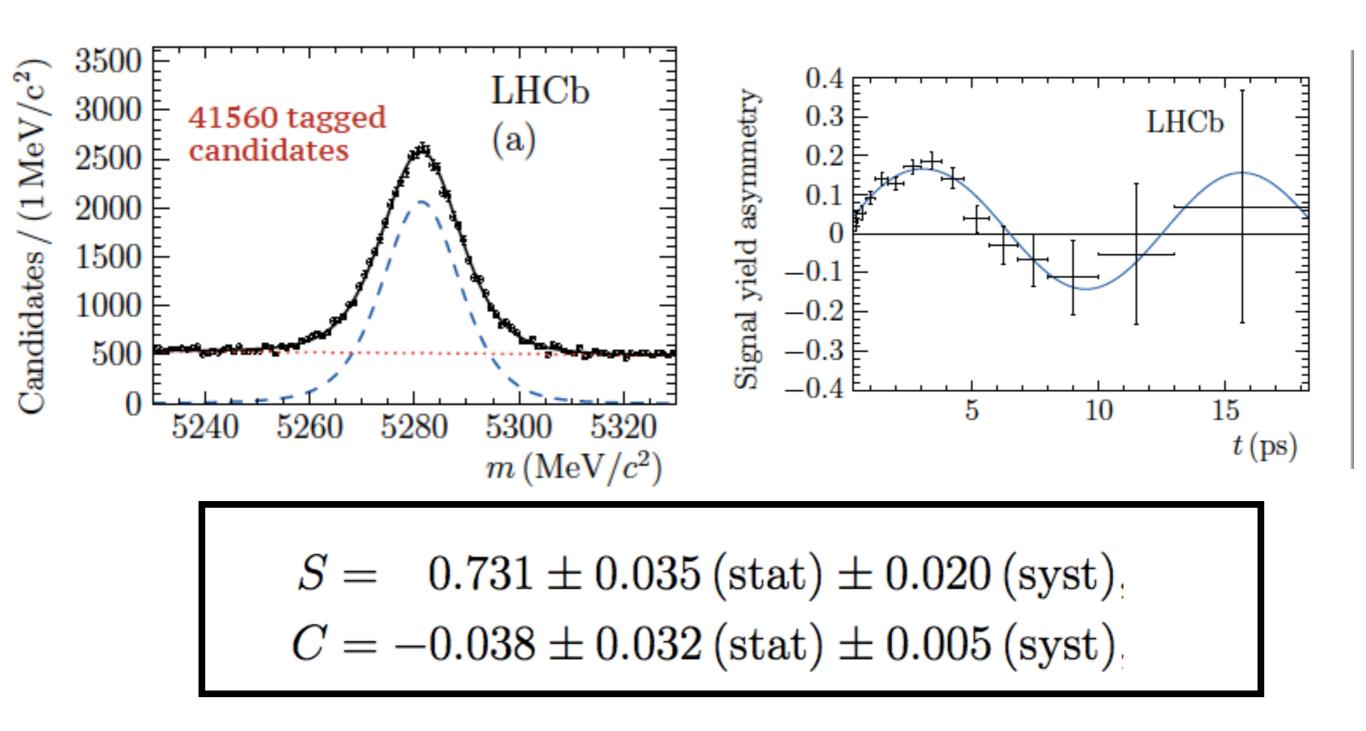
Phys. Rev. Lett. 113 (2014) 211801



### Measurement of CPV in $B \rightarrow J/\psi K_s [3 \text{ fb}^{-1}]$



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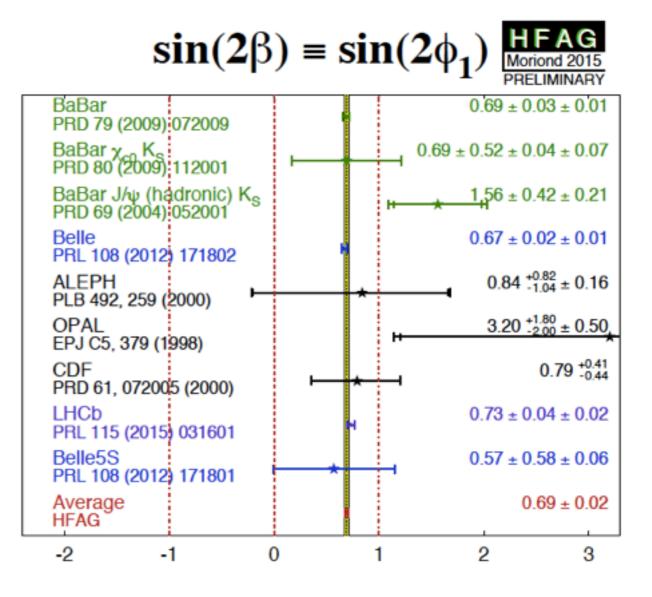


### An other big picture

#### World average: [HFAG] $\sin(2\beta)^{\exp} = 0.691 \pm 0.017$

#### Expectation from global fit: [CKMFitter]

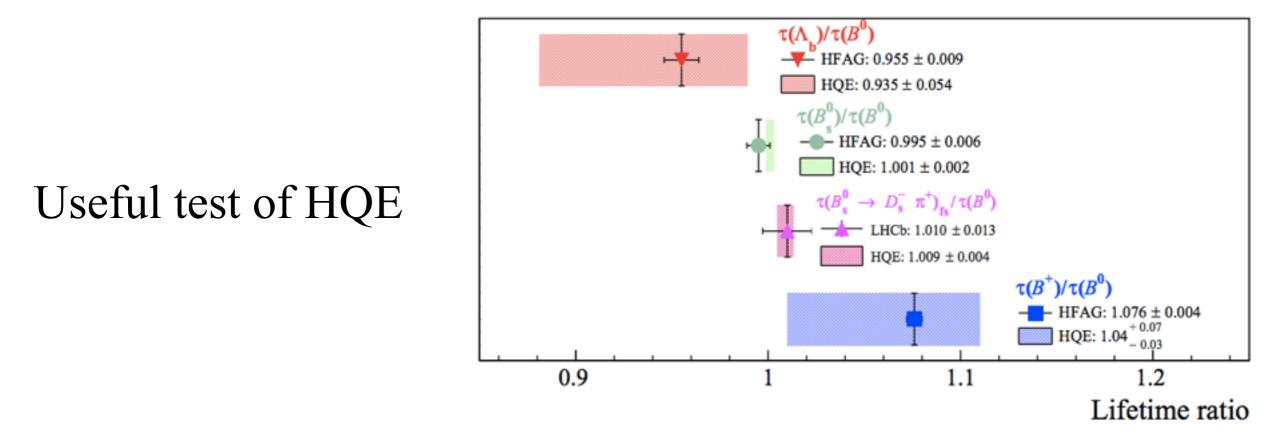
 $\sin(2\beta)^{\rm SM} = 0.748^{+0.030}_{-0.032}$ 



#### KEEP CALM BY FOCUSING ON DIGRESSION

### Very precise lifetime measurements

Ratio	Value
$\tau_{B^+}/\tau_{B^0 \to J/\psi K^*(892)^0}$	$1.074 \pm 0.005 \pm 0.003$
$ au_{B_s^0} /  au_{B^0  o J/\psi K^*(892)^0}$	$0.971 \pm 0.009 \pm 0.004$
$ au_{\Lambda_b^0}/ au_{B^0  o J/\psi K^*(892)^0}$	$0.929 \pm 0.018 \pm 0.004$
$ au_{B^+}/ au_{B^-}$	$1.002 \pm 0.004 \pm 0.002$
$ au_{\Lambda_b^0}/ au_{\overline{\Lambda}_b^0}$	$0.940 \pm 0.035 \pm 0.006$
$\tau_{B^0 \to J/\psi K^*(892)^0}/\tau_{\overline{B}^0 \to J/\psi \overline{K}^*(892)^0}$	$1.000 \pm 0.008 \pm 0.009$



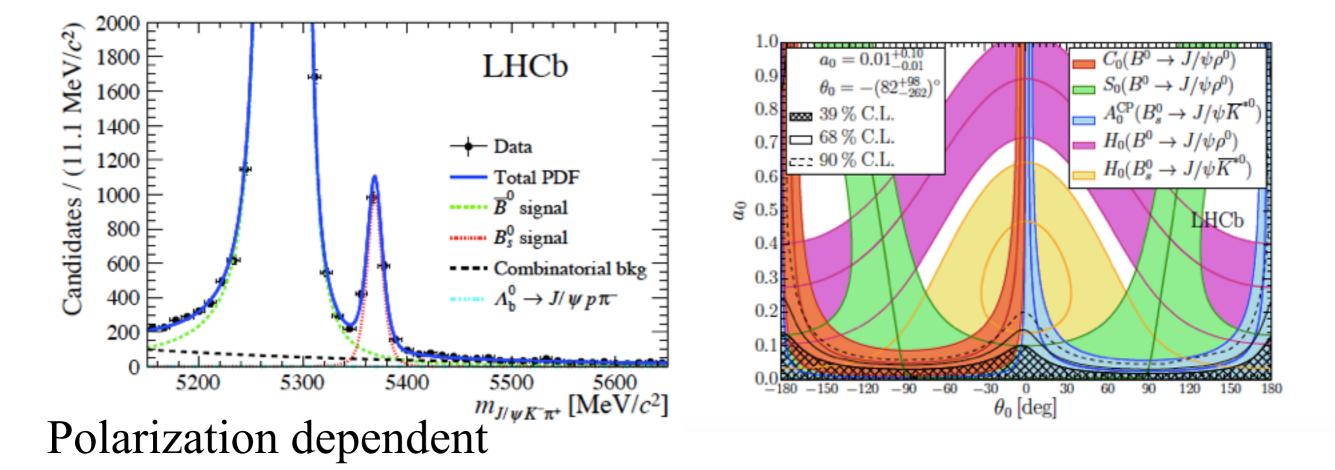
JHEP04(2014)114 [1 fb<sup>-1</sup>]

### How do we constrain the penguins?



Proposal : <u>arXiv:1412.6834</u>

$$A\left(B_s^0 \to J/\psi \overline{K}^{*0}\right)_i\right) = -\lambda \mathcal{A}_i \left[1 - a_i e^{i\theta_i} e^{i\gamma}\right]$$



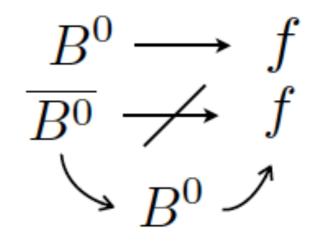
$$\begin{split} &\Delta\phi_{s,0}^{J/\psi\,\phi} = 0.000^{+0.009}_{-0.011}\,(\text{stat})^{+0.004}_{-0.009}\,(\text{syst})\,\,\text{rad} \\ &\Delta\phi_{s,\parallel}^{J/\psi\,\phi} = 0.001^{+0.010}_{-0.014}\,(\text{stat})^{+0.007}_{-0.008}\,(\text{syst})\,\,\text{rad} \\ &\Delta\phi_{s,\perp}^{J/\psi\,\phi} = 0.003^{+0.010}_{-0.014}\,(\text{stat})^{+0.007}_{-0.008}\,(\text{syst})\,\,\text{rad} \end{split}$$

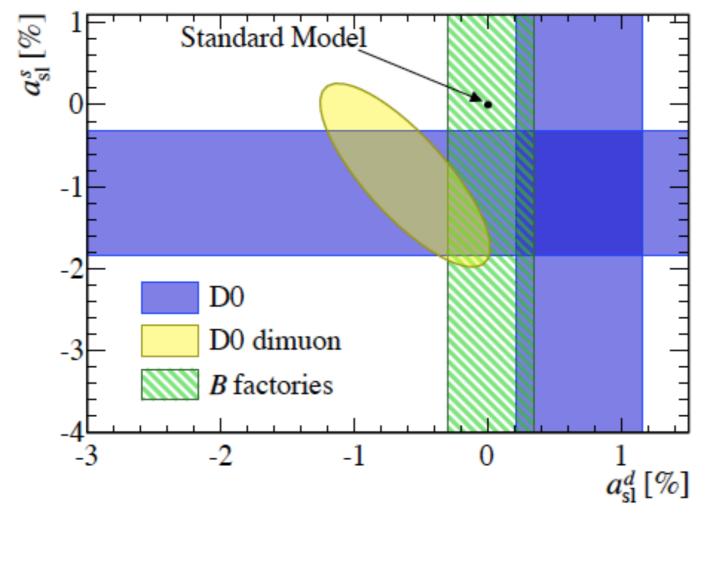
The penguins are small!

arXiv:1509.00400 [3 fb<sup>-1</sup>] Phys. Lett. B742 (2015) 38 [3 fb<sup>-1</sup>]

More CP asymmetries

$$a_{\rm sl} = \frac{N(\bar{B} \to B \to f) - N(B \to \bar{B} \to \bar{f})}{N(\bar{B} \to B \to f) + N(B \to \bar{B} \to \bar{f})}$$

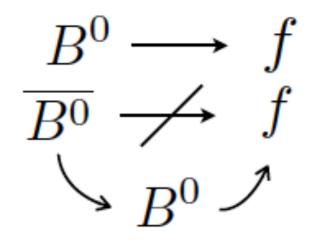


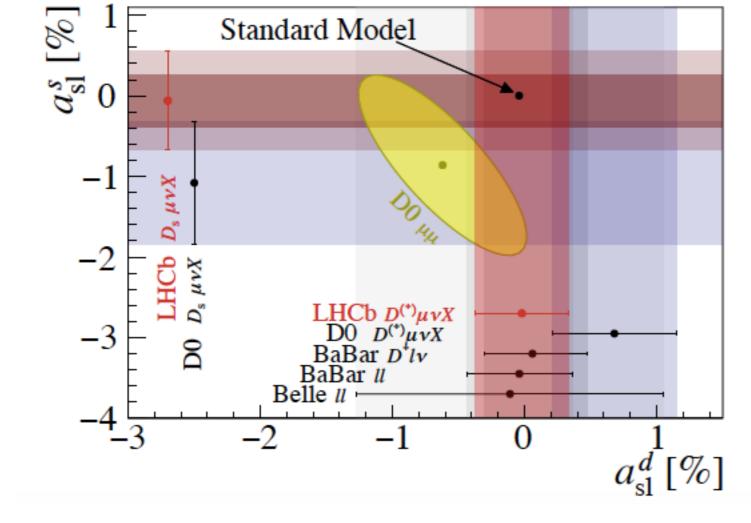


Phys. Rev. Lett. 114 (2015) 041601 [3 fb<sup>-1</sup>] Phys. Lett. B 728 (2014) 607-615 [1 fb<sup>-1</sup>]

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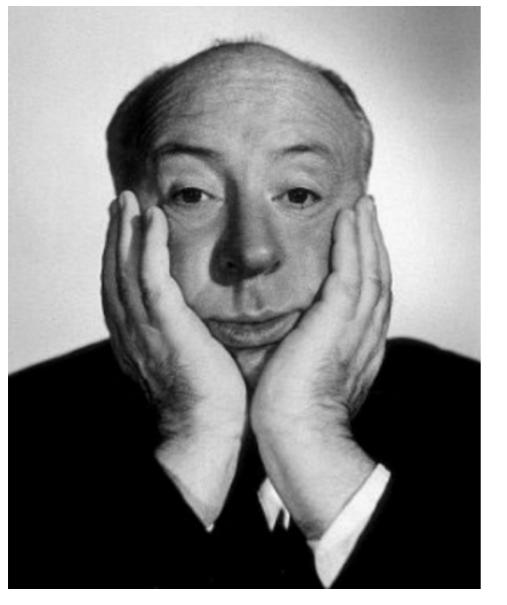




LHCb measurements very consistent with the SM.

Phys. Rev. Lett. 114 (2015) 041601 [3 fb<sup>-1</sup>] Phys. Lett. B 728 (2014) 607-615 [1 fb<sup>-1</sup>]

# Conclusion



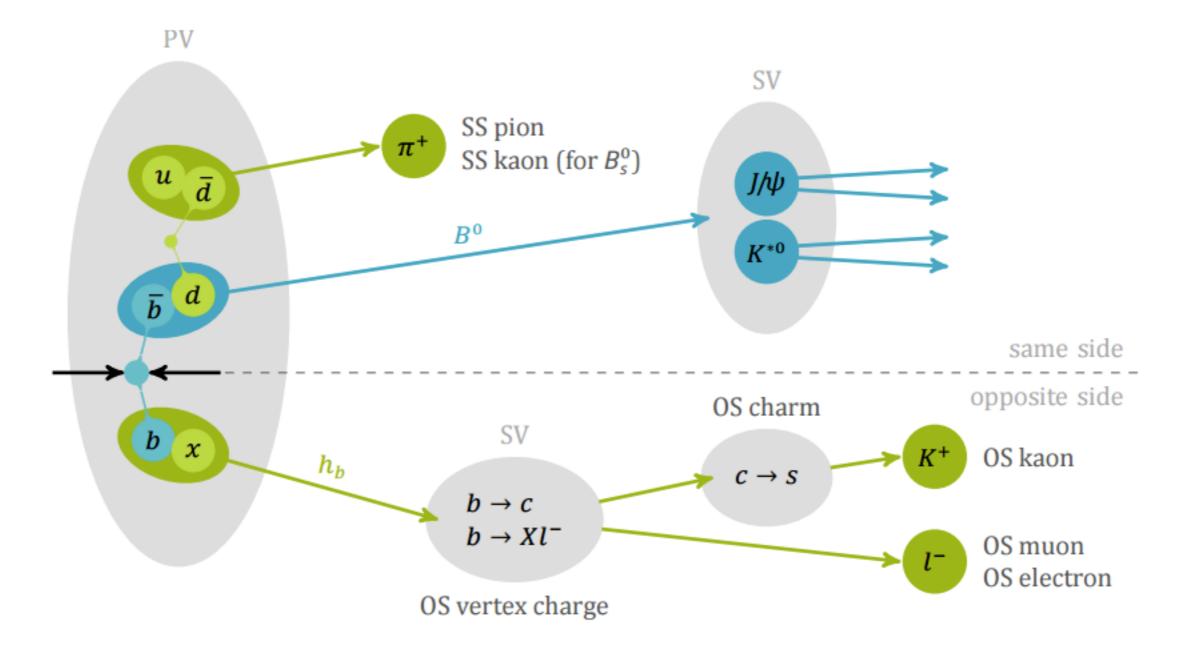
The measurements of CP observables are reaching very high precision.

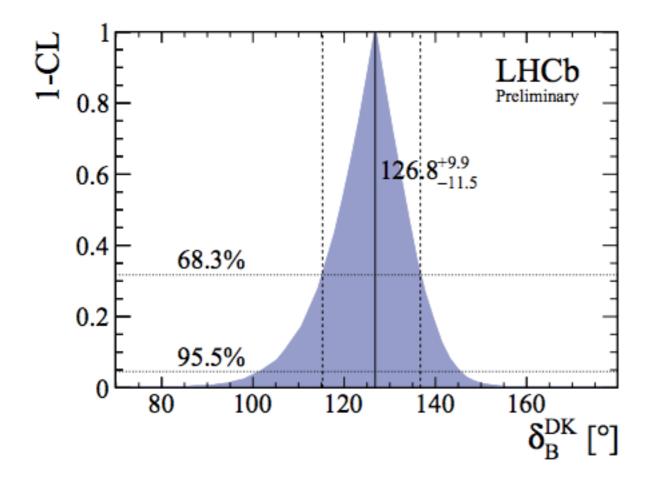
LHCb will continue to collect data and push further the precision of the measurements as well as add additional modes

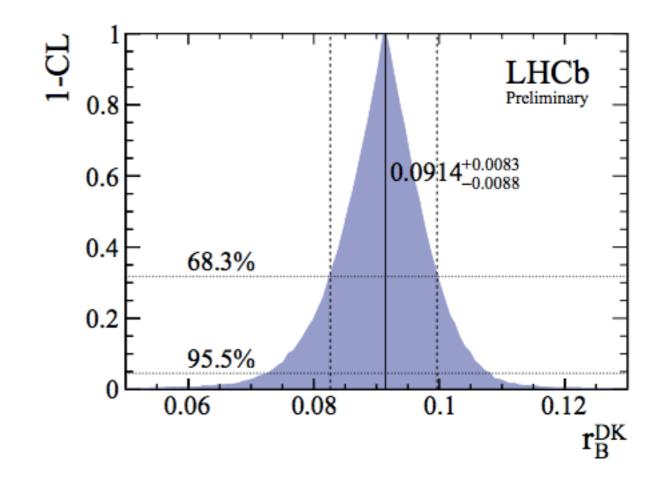
No clear sign of New Physics contributions.



### Backup slides







#### GGSZ (Giri, Grossman, Sofer, Zupan):

Use quasi 2-body *CP* eigenstate of the *D* to be resolved in the Dalitz plane.  $D \longrightarrow K_s \pi \pi$ . So far the most precise gamma determination.

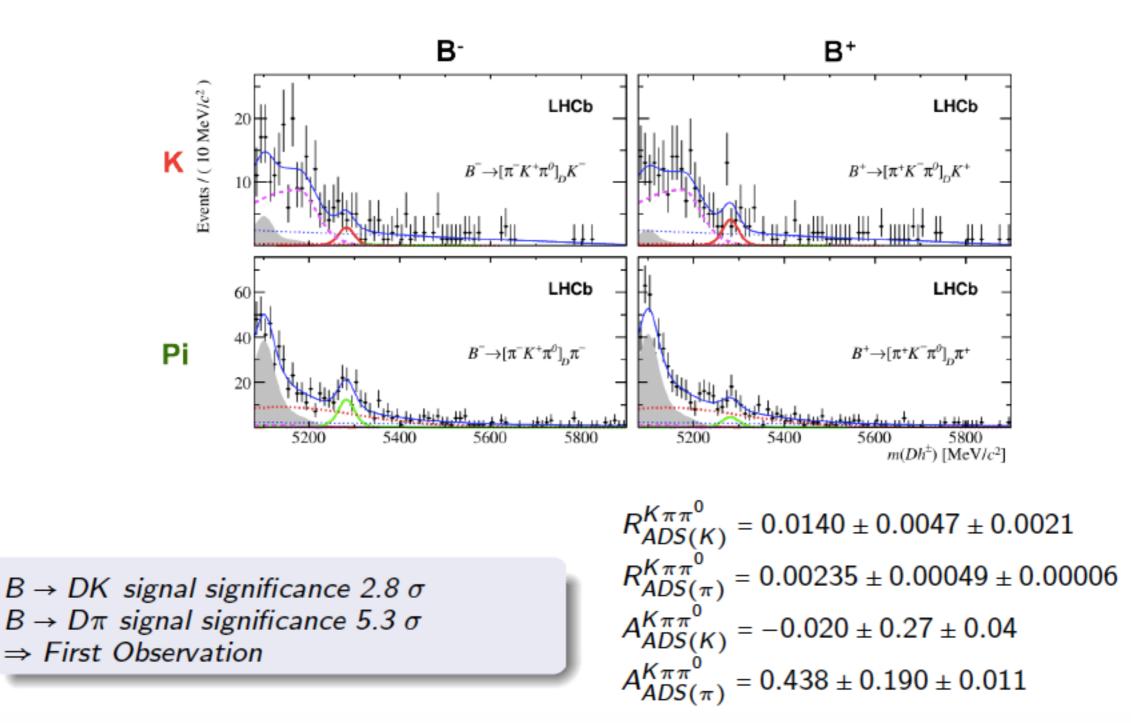
#### ADS (Atwood, Dunietz, Soni):

Use anti-D<sup>0</sup> K<sup>+</sup>  $\pi$  for b—>u transitions (Cabibbo allowed) and D<sup>0—></sup> K<sup>-</sup>  $\pi$ <sup>+</sup> (Doubly Cabibbo Suppressed) for b—>c transitions. One has to know strong phases from D decays.

GLW (Gronau, London, Wyler):

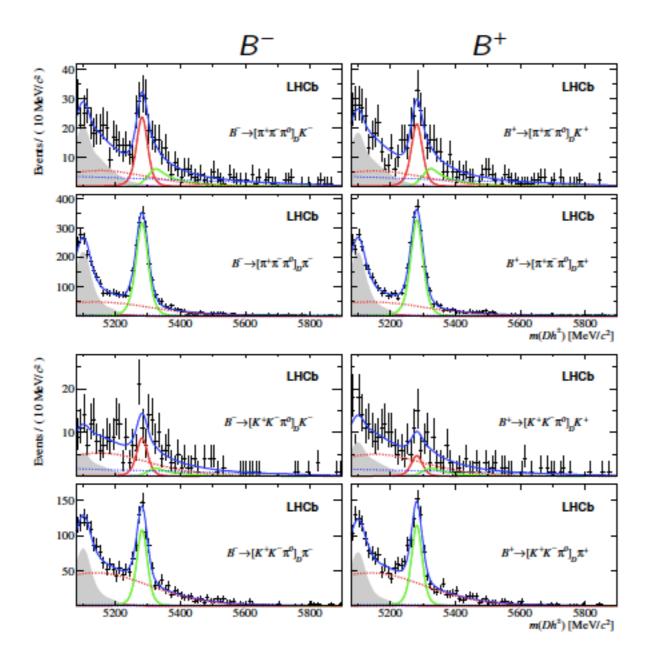
Search for *D* mesons in *CP* eigenstates eg :  $D \longrightarrow KK$  and  $D \longrightarrow \pi\pi$ 

#### [PHYS. REV. D91 (2015) 112014]



#### ADS-like modes

#### [PHYS. REV. D91 (2015) 112014]



#### GLW-like modes

$$B \rightarrow DK$$
  
 $B \rightarrow D\pi$ 

$$B \rightarrow D(KK\pi^0)\pi$$
 signal significance > 10.0 $\sigma$   
 $B \rightarrow D(KK\pi^0)K$  First evidence (4.5  $\sigma$ )

$$\begin{aligned} R_{qGLW}^{KK\pi^{0}} &= 0.95 \pm 0.22 \pm 0.05 \\ R_{qGLW}^{\pi\pi\pi^{0}} &= 0.98 \pm 0.11 \pm 0.05 \\ A_{qGLW(K)}^{KK\pi^{0}} &= 0.30 \pm 0.20 \pm 0.02 \\ A_{qGLW(K)}^{\pi\pi\pi^{0}} &= 0.054 \pm 0.091 \pm 0.011 \\ A_{qGLW(K)}^{KK\pi^{0}} &= -0.030 \pm 0.040 \pm 0.005 \\ A_{qGLW(\pi)}^{\pi\pi\pi^{0}} &= -0.016 \pm 0.020 \pm 0.004 \end{aligned}$$