



Number Theory and Physics

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**Elliptic multiple zeta values
and superstring one-loop amplitudes**

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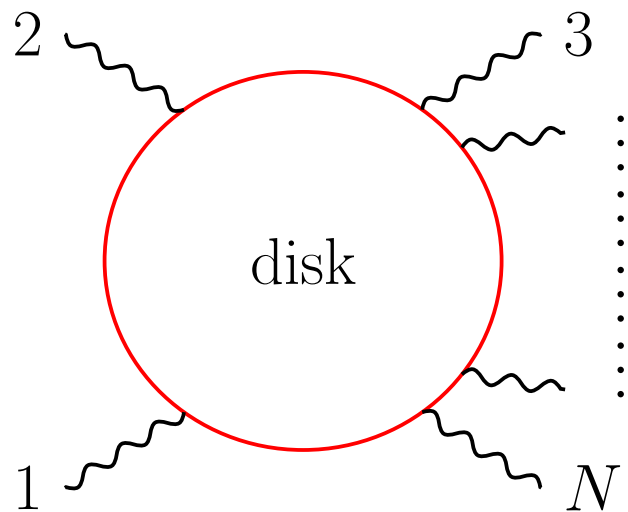
based on arXiv:1412.5535 & 1507.02254 with J. Brödel, C. Mafra, N. Matthes

25.05.2016

Motivation: Open-string scattering amplitudes @ tree level and one loop

\implies iterated integrals over the **worldsheet boundaries**

genus zero



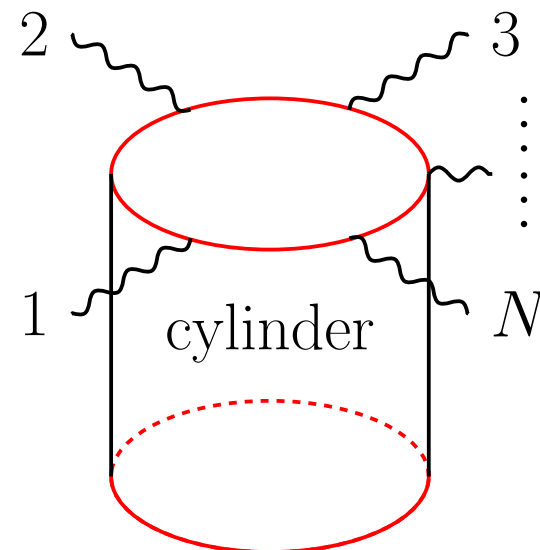
peel off
YM trees

expand
in α'

$$\int \prod_{0 \leq z_i \leq z_{i+1} \leq 1} \frac{dz_j}{z_j - n_j} = \text{MZVs}$$

[earlier work]

genus one



half a torus
of modular
parameter τ

$$\frac{\mathbb{C}}{\mathbb{Z} + \tau\mathbb{Z}}$$

$$\int \prod_{0 \leq z_i \leq z_{i+1} \leq 1} f^{(n_j)}(z_j) dz_j = \text{eMZVs}$$

[this talk]

Outline

I. Definition of eMZVs and $f^{(n)}$

[Brown, Levin, Enriquez, ...]

II. Superstring one-loop amplitudes

[Brödel, Mafra, Matthes, OS 1412.5535]

III. eMZV relations & indecomposables

[Brödel, Matthes, OS 1507.02254]

I. 1 Iterated integrals: genus-zero versus elliptic

Recall that **multiple polylogarithms** [Goncharov, see Pierre Vanhove's talk]

$$G(a_1, a_2, \dots, a_r; z) \equiv \int_0^z \frac{dt}{t - a_1} G(a_2, \dots, a_r; t), \quad G(; z) \equiv 1$$

specialize to **multiple zeta values** (MZVs) upon $z \rightarrow 1$ and $a_j \rightarrow \{0, 1\}$:

$$\zeta(a_1, a_2, \dots, a_r) \equiv (-1)^{\sum_{j=1}^r a_j} \int_{0 \leq z_i \leq z_{i+1} \leq 1} \frac{dz_1}{z_1 - a_1} \frac{dz_2}{z_2 - a_2} \dots \frac{dz_r}{z_r - a_r}$$

Both are said to have weight $r \equiv$ number of integrations.

Can recursively **bypass clashes** $a_j = z$ via **partial fraction**:

$$G(a, 0, z; z) = G(0, 0, a; z) - G(0, a, a; z) - \zeta_2 G(a; z)$$

resting on:
$$\frac{1}{(z-a)(z-b)} = \frac{1}{a-b} \left(\frac{1}{z-a} - \frac{1}{z-b} \right)$$

I. 1 Iterated integrals: genus-zero versus elliptic

Elliptic iterated integrals with suitable $f^{(n)}$, $n \in \mathbb{N}$ [Brown, Levin]

$$\Gamma \left(\begin{matrix} n_1 & n_2 & \cdots & n_r \\ a_1 & a_2 & \cdots & a_r \end{matrix}; z \right) \equiv \int_0^z dt f^{(n_1)}(t - a_1) \Gamma \left(\begin{matrix} n_2 & \cdots & n_r \\ a_2 & \cdots & a_r \end{matrix}; t \right), \quad \Gamma(; z) \equiv 1$$

specialize to elliptic multiple zeta values (eMZVs) upon $z \rightarrow 1$ and $a_j \rightarrow 0$:

$$\omega(n_1, n_2, \dots, n_r) \equiv \int_{0 \leq z_i \leq z_{i+1} \leq 1} f^{(n_1)}(z_1) dz_1 f^{(n_2)}(z_2) dz_2 \cdots f^{(n_r)}(z_r) dz_r.$$

Both are said to have length r and weight $\sum_{j=1}^r n_j$. [Enriquez 1301.3042]

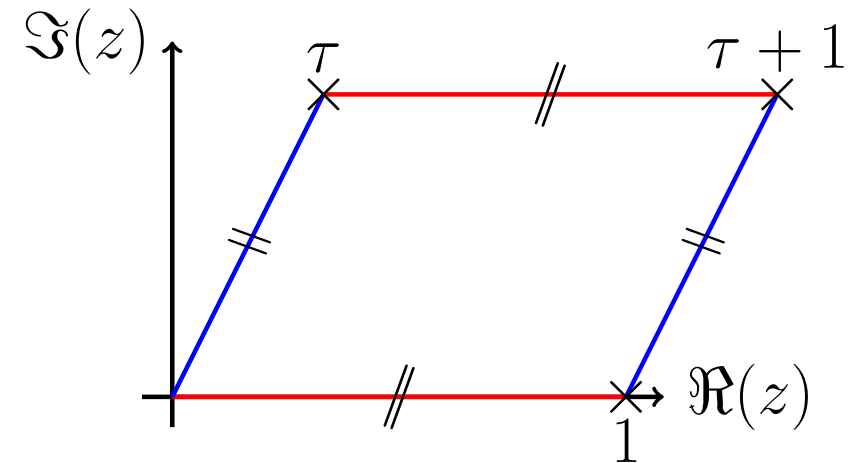
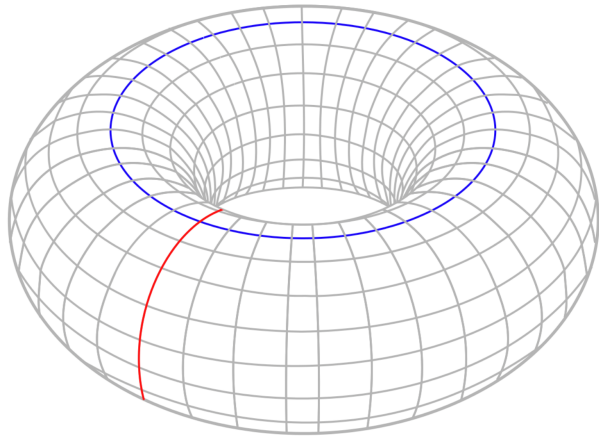
Can recursively bypass clashes $a_j = z$ via Fay relations ($C_{p,q} \in \mathbb{Z}$):

$$\Gamma \left(\begin{matrix} 1 & 1 \\ z & 0 \end{matrix}; z \right) = 2\Gamma \left(\begin{matrix} 0 & 2 \\ 0 & 0 \end{matrix}; z \right) + \Gamma \left(\begin{matrix} 2 & 0 \\ 0 & 0 \end{matrix}; z \right) - 2\Gamma \left(\begin{matrix} 1 & 1 \\ 0 & 0 \end{matrix}; z \right) + \zeta_2$$

$$\text{from } f^{(m)}(z-a)f^{(n)}(z-b) = \sum_{p+q=m+n} C_{p,q} f^{(p)}(z-a)f^{(q)}(a-b) + \begin{matrix} (m,a) \\ \updownarrow \\ (n,b) \end{matrix}$$

I. 2 $f^{(n)}$ – doubly periodic extension of $\frac{dz}{z}$ with Fay relation

Parametrization of elliptic curve \equiv torus



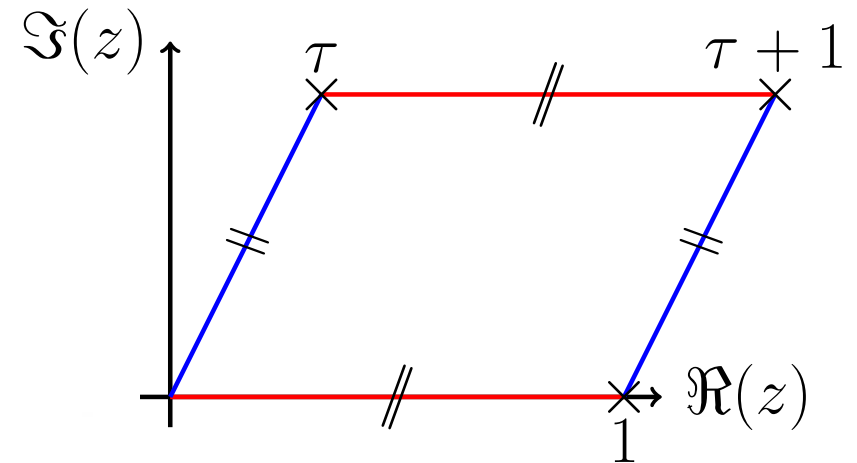
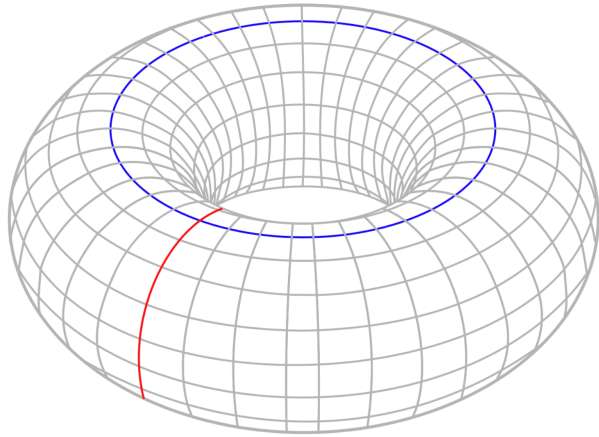
Jacobi θ -function takes role of the identity map on the torus

$$\theta(z, \tau) \equiv \sin(\pi z) \prod_{n=1}^{\infty} (1 - e^{2\pi i(n\tau+z)})(1 - e^{2\pi i(n\tau-z)})$$

$$\text{uplift } \frac{1}{z} \rightarrow \partial_z \ln \theta(z, \tau) + 2\pi i \frac{\Im(z)}{\Im(\tau)} \equiv f^{(1)}(z, \tau)$$

I. 2 $f^{(n)}$ – doubly periodic extension of $\frac{dz}{z}$ with Fay relation

Parametrization of elliptic curve \equiv torus



Jacobi θ -function takes role of the identity map on the torus

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$$\text{uplift } \frac{1}{z} \rightarrow \partial_z \ln \theta(z, \tau) + \underbrace{2\pi i \frac{\Im(z)}{\Im(\tau)}}_{\text{sacrifice holomorphicity for } f^{(1)}(z, \tau) = f^{(1)}(z + \tau, \tau)} \equiv \overbrace{f^{(1)}(z, \tau) = f^{(1)}(z + 1, \tau)}^{\text{poles like both } \frac{1}{z} \text{ and } \frac{1}{1-z}}$$

I. 2 $f^{(n)}$ – doubly periodic extension of $\frac{dz}{z}$ with Fay relation

To go beyond

$$\text{uplift } \frac{1}{z} \rightarrow f^{(1)}(z, \tau) \equiv \partial_z \ln \theta(z, \tau) + 2\pi i \frac{\mathfrak{F}(z)}{\mathfrak{F}(\tau)},$$

check partial fraction \Rightarrow new function $f^{(2)}$ without pole

$$\frac{1}{(z-a)(z-b)} + \text{cyc}(z, a, b) = 0$$

$$\downarrow$$

$$f^{(1)}(z-a)f^{(1)}(z-b) + \text{cyc}(z, a, b) = f^{(2)}(z-a) + \text{cyc}(z, a, b)$$

- similar identity for $f^{(1)}(z-a)f^{(2)}(z-b)$ yields new function $f^{(3)}$, etc.
- formally adjoin $f^{(0)} \equiv 1$
- drop the second argument τ of θ and $f^{(n)}$ here and henceforth

I. 2 $f^{(n)}$ – doubly periodic extension of $\frac{dz}{z}$ with Fay relation

Generating fct.: non-holomorphic version of Kronecker-Eisenstein series

$$\Omega(z, \alpha) \equiv \exp\left(2\pi i \alpha \frac{\Im(z)}{\Im(\tau)}\right) \frac{\theta'(0)\theta(z+\alpha)}{\theta(z)\theta(\alpha)} = \sum_{n=0}^{\infty} \alpha^{n-1} f^{(n)}(z)$$

[Kronecker, Brown, Levin]

- doubly-periodic $\Omega(z, \alpha) = \Omega(z + \tau, \alpha) = \Omega(z + 1, \alpha)$ thanks to $\exp(\dots)$
- e.g. $f^{(2)}(z) \equiv \frac{1}{2} \left\{ \left(\partial \ln \theta(z) + 2\pi i \frac{\Im(z)}{\Im(\tau)} \right)^2 + \partial^2 \ln \theta(z) - \frac{\theta'''(0)}{3\theta'(0)} \right\}$
- $\text{Res}_{z=0} \Omega(z, \alpha) = 1 \Rightarrow$ no other poles than $f^{(1)}(z) \sim \frac{1}{z}$
- partial fraction generalizes to Fay relation:

$$\Omega(z_1, \alpha_1)\Omega(z_2, \alpha_2) = \Omega(z_1, \alpha_1 + \alpha_2)\Omega(z_2 - z_1, \alpha_2) + (1 \leftrightarrow 2)$$

II. One-loop superstring amplitude

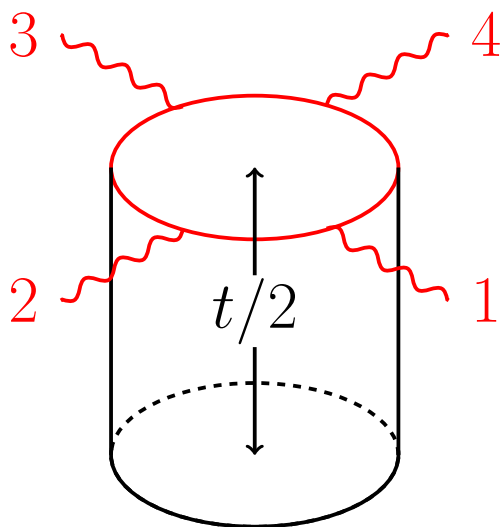
Focus on planar cylinder diagram at four-points:

$$A_{\text{string}}^{\text{1-loop}}(1, 2, 3, 4) = s_{12}s_{23}A_{\text{YM}}^{\text{tree}}(1, 2, 3, 4) \int_0^\infty dt I_{1234}(s_{ij}, \tau = it)$$

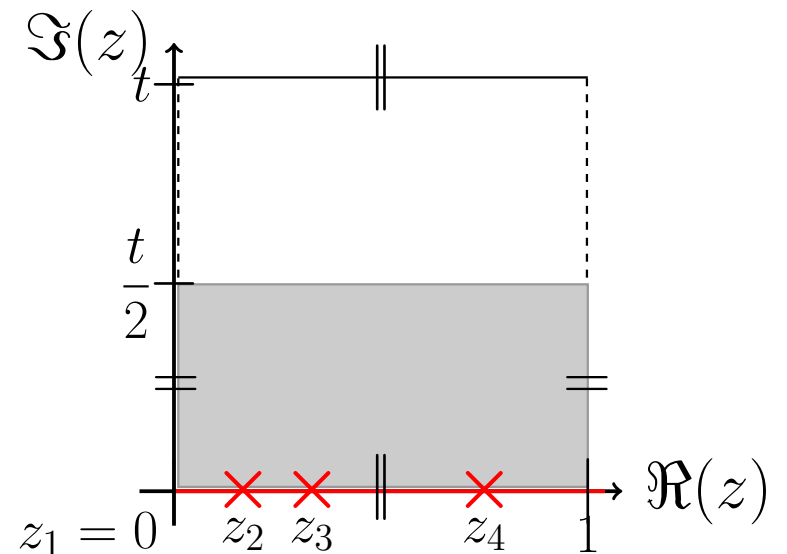
$$I_{1234}(s_{ij}, \tau) = \int_0^1 dz_4 \int_0^{z_4} dz_3 \int_0^{z_3} dz_2 \exp \left(\sum_{i < j}^4 s_{ij} P(z_i - z_j, \tau) \right) \Big|_{z_1=0}$$

[Brink, Green, Schwarz 1982]

with dim'less $s_{ij} \equiv \alpha'(k_i + k_j)^2$ and worldsheet propagator $\partial P = f^{(1)}$.



parametrized as



Analytic α' -dependence from expanding the exponentials

$$I_{1234}(s_{ij}, \tau) = \int_0^1 dz_4 \int_0^{z_4} dz_3 \int_0^{z_3} dz_2 \prod_{i < j}^4 \sum_{n_{ij}=0}^{\infty} \frac{[s_{ij} P(z_i - z_j, \tau)]^{n_{ij}}}{n_{ij}!} \Big|_{z_1=0}$$

and integrating $P(z_i - z_j) = \int_{z_j}^{z_i} dx f^{(1)}(x - z_j)$ order by order in α' .

Each monomial in s_{ij} is accompanied by eMZVs, e.g.

$$s_{ij}^0 \leftrightarrow \int_0^1 dz_4 f^{(0)}(z_4) \int_0^{z_4} dz_3 f^{(0)}(z_3) \int_0^{z_3} dz_2 f^{(0)}(z_2) = \omega(0, 0, 0)$$

after formally inserting $f^{(0)} = 1$ as well as

$$\left. \begin{array}{l} s_{12} \\ s_{13} \end{array} \right\} \leftrightarrow \int_0^1 dz_4 f^{(0)} \int_0^{z_4} dz_3 f^{(0)} \int_0^{z_3} dz_2 f^{(0)} \left\{ \begin{array}{l} \int_0^{z_2} dx f^{(1)}(x) \\ \int_0^{z_3} dx f^{(1)}(x) \end{array} \right.$$

$$\implies s_{12} \leftrightarrow \omega(1, 0, 0, 0), \quad s_{13} \leftrightarrow \underbrace{\omega(1, 0, 0, 0)}_{\text{from } 0 \leq x \leq z_2} + \underbrace{\omega(0, 1, 0, 0)}_{\text{from } z_2 \leq x \leq z_3}$$

At higher order ...

$$\begin{aligned}
 s_{12}s_{23} &\leftrightarrow \int_0^1 dz_4 \int_0^{z_4} dz_3 \int_0^{z_3} dz_2 \left(\int_{z_3}^{z_2} dx f^{(1)}(x - z_3) \right) \left(\int_0^{z_2} dy f^{(1)}(y) \right) \\
 &= - \int_0^1 dz_4 \int_0^{z_4} dz_3 \Gamma \left(\begin{matrix} 1 & 0 & 1 \\ z_3 & 0 & 0 \end{matrix}; z_3 \right)
 \end{aligned}$$

... need Fay relations

$$\Gamma \left(\begin{matrix} 1 & 0 & 1 \\ z_3 & 0 & 0 \end{matrix}; z_3 \right) = 2\Gamma \left(\begin{matrix} 0 & 0 & 2 \\ 0 & 0 & 0 \end{matrix}; z_3 \right) + \Gamma \left(\begin{matrix} 0 & 2 & 0 \\ 0 & 0 & 0 \end{matrix}; z_3 \right) - 2\Gamma \left(\begin{matrix} 0 & 1 & 1 \\ 0 & 0 & 0 \end{matrix}; z_3 \right) + \zeta_2 \Gamma \left(\begin{matrix} 0 \\ 0 \end{matrix}; z_3 \right) .$$

Above example then integrates to

$$\begin{aligned}
 s_{12}s_{23} &\leftrightarrow -2\omega(2, 0, 0, 0, 0) - \omega(0, 2, 0, 0, 0) + 2\omega(1, 1, 0, 0, 0) - \zeta_2\omega(0, 0, 0) \\
 &= -\omega(1, 0, 0, 0, 1)
 \end{aligned}$$

after using eMZV relations at length five.

After using momentum conservation for s_{ij} , first orders simplify to

$$I_{1234}(s_{ij}) = \omega(0, 0, 0) - 2\omega(0, 1, 0, 0) (s_{12} + s_{23}) + 2\omega(0, 1, 1, 0, 0) (s_{12}^2 + s_{23}^2) \\ - 2\omega(0, 1, 0, 1, 0) s_{12}s_{23} + \beta_5 (s_{12}^3 + 2s_{12}^2s_{23} + 2s_{12}s_{23}^2 + s_{23}^3) + \beta_{2,3} s_{12}s_{23}(s_{12} + s_{23}) + \dots$$

with shorthands

[Brödel, Mafra, Matthes, OS 1412.5535]

$$\beta_5 = \frac{4}{3} [\omega(0, 0, 1, 0, 0, 2) + \omega(0, 1, 1, 0, 1, 0) - \omega(2, 0, 1, 0, 0, 0) - \zeta_2\omega(0, 1, 0, 0)] \\ \beta_{2,3} = \frac{1}{3}\omega(0, 0, 1, 0, 2, 0) - \frac{3}{2}\omega(0, 1, 0, 0, 0, 2) - \frac{1}{2}\omega(0, 1, 1, 1, 0, 0) \\ - 2\omega(2, 0, 1, 0, 0, 0) - \frac{4}{3}\omega(0, 0, 1, 0, 0, 2) - \frac{10}{3}\zeta_2\omega(0, 1, 0, 0) .$$

Choice of indecomposable eMZVs requires guidance,

in particular at higher α' -order \leftrightarrow weight \leftrightarrow length ...

also: n -point amplitude naturally compatible with eMZV language !

III. 1 eMZV relations and indecomposables

The following **MZVs** are believed to be **indecomposable over \mathbb{Q}** :

weight w	0	1	2	3	4	5	6	7	8	9	10	11	12
indec. MZVs	1	\emptyset	ζ_2	ζ_3	\emptyset	ζ_5	\emptyset	ζ_7	$\zeta_{3,5}$	ζ_9	$\zeta_{3,7}$	$\zeta_{11}, \zeta_{3,3,5}$	$\zeta_{3,9}, \zeta_{1,1,4,6}$

MZVs satisfy shuffle and stuffle relations, and **eMZVs** obey

- shuffle $\omega(n_1, \dots, n_r)\omega(k_1, \dots, k_s) = \omega((n_1, \dots, n_r) \sqcup (k_1, \dots, k_s))$
- reflection: $\omega(n_1, n_2, \dots, n_r) = (-1)^{n_1+n_2+\dots+n_r}\omega(n_r, \dots, n_2, n_1)$
- Fay rel's $\omega(n_1, n_2, \dots) \leftrightarrow \omega(n_1+j, n_2-j, \dots)$ such as $\omega(0, 5) = \omega(2, 3)$

all of which **preserve the weight $\sum_{j=1}^r n_j$ of $f^{(n_j)}$ integrands.**

→ Which eMZVs remain indecomposable w.r.t. $\mathbb{Q}[\text{MZV}]$?

However, $f^{(0)}$ @ zero weight $\Rightarrow \exists \infty$ eMZVs $\omega(n, \underbrace{0, \dots, 0}_{\text{any number}})$ @ weight n

\Rightarrow organize relations by length r :

length $r = 1$: only constant eMZVs: $\omega(n) = \begin{cases} -2 \zeta_n & : n \text{ even} \\ 0 & : n \text{ odd} \end{cases}$

length $r = 2$: shuffle and reflection reduce even-weight eMZVs to $r = 1$:

$$\omega(n_1, n_2) \Big|_{n_1+n_2 \text{ even}} = \begin{cases} 2 \zeta_{n_1} \zeta_{n_2} & : n_1, n_2 \text{ both even} \\ 0 & : n_1, n_2 \text{ both odd} \end{cases}$$

odd-weight $\omega(n_1, n_2)$ depend on τ and can be reduced to $\omega(0, 2p - 1)$:

$$\begin{aligned} \omega(n_1, n_2) \Big|_{n_1+n_2 \text{ odd}} &= (-1)^{n_1} \omega(0, n_1 + n_2) + 2\delta_{n_1,1} \zeta_{n_2} \omega(0, 1) - 2\delta_{n_2,1} \zeta_{n_1} \omega(0, 1) \\ &+ 2 \left\{ \sum_{p=1}^{\lceil \frac{1}{2}(n_2-3) \rceil} \binom{n_1 + n_2 - 2p - 2}{n_1 - 1} \zeta_{n_1+n_2-2p-1} \omega(0, 2p+1) - (n_1 \leftrightarrow n_2) \right\} \end{aligned}$$

any length r : eMZVs with $(-1)^{\text{length}} = (-1)^{\text{weight}}$ reducible to lower length

length $r = 3$: $\lceil \frac{w}{6} \rceil$ indecomposable $\omega(n_1, n_2, n_3)$ @ even $w = \sum_{j=1}^3 n_j$:

w	2	4	6	8	10	12	14
indec.	$\omega(0, 0, 2)$	$\omega(0, 0, 4)$	$\omega(0, 0, 6)$	$\omega(0, 0, 8)$	$\omega(0, 0, 10)$	$\omega(0, 0, 12)$	$\omega(0, 0, 14)$
eMZVs at				$\omega(0, 3, 5)$	$\omega(0, 3, 7)$	$\omega(0, 3, 9)$	$\omega(0, 3, 11)$
$r = 3$							$\omega(0, 5, 9)$

at $w = 8$, for instance, shuffle, reflection & Fay imply

$$\begin{aligned} \omega(0, 6, 2) = & -\frac{21}{2} \zeta_8 + 2\omega(0, 3)\omega(0, 5) - 14\zeta_6\omega(0, 0, 2) \\ & - 6\zeta_4\omega(0, 0, 4) - \frac{9}{2}\omega(0, 0, 8) - \frac{2}{5}\omega(0, 3, 5) \end{aligned}$$

Further eMZV relations available @ <https://tools.aei.mpg.de/emzv>

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length $r = 3$: $\left\lceil \frac{w}{6} \right\rceil$ indecomposable $\omega(n_1, n_2, n_3)$ @ even $w = \sum_{j=1}^3 n_j$:

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indec.	$\omega(0, 0, 2)$	$\omega(0, 0, 4)$	$\omega(0, 0, 6)$	$\omega(0, 0, 8)$	$\omega(0, 0, 10)$	$\omega(0, 0, 12)$	$\omega(0, 0, 14)$
eMZVs at				$\omega(0, 3, 5)$	$\omega(0, 3, 7)$	$\omega(0, 3, 9)$	$\omega(0, 3, 11)$
$r = 3$							$\omega(0, 5, 9)$

length $r = 4$: conjecture $\left\lfloor \frac{1}{2} + \frac{1}{48}(w + 5)^2 \right\rfloor$ indecomp's @ odd $w = \sum_{j=1}^4 n_j$:

w	1	3	5	7	9
indec.	$\omega(0, 0, 1, 0)$	$\omega(0, 0, 0, 3)$	$\omega(0, 0, 0, 5)$	$\omega(0, 0, 0, 7)$	$\omega(0, 0, 0, 9), \omega(0, 0, 4, 5)$
eMZVs at			$\omega(0, 0, 2, 3)$	$\omega(0, 0, 2, 5)$	$\omega(0, 0, 2, 7), \omega(0, 1, 3, 5)$
$r = 4$				$\omega(0, 0, 4, 3)$	

summary @ length $r \leq 7$: indecomposable eMZVs @ $(-1)^{\text{weight}} \neq (-1)^{\text{length}}$

$r \backslash w$	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	19	20	21	22	23
2	1		1		1		1		1		1		1		1		1		1		1		1
3		1		1		1		2		2		2		3		3		3		4		4	
4	1		1		2		3		4		5		7		8		10		12		14		16
5		1		2		4		6		9		13		17		23		30		37		47	
6	1		2		4		8		13		22		31		45		?		?		?		?
7		1		4		8		16		29		48		?		?		?		?		?	

- strong evidence that shuffle, reflection & Fay \Rightarrow **any** eMZV relation
- will derive these numbers from combinatorial principles

III. 2 Direct computation of eMZVs

eMZVs satisfy ODE w.r.t. modular parameter $q = e^{2\pi i\tau}$: [Enriquez 1301.3042]

$$-4\pi^2 q \frac{d}{dq} \text{eMZV} \left(\begin{array}{l} \text{weight } w \\ \text{length } \ell \end{array} \right) \sim G_k(q) \times \text{eMZV} \left(\begin{array}{l} \text{weight } w-k+1 \\ \text{length } \ell-1 \end{array} \right)$$

recursive in length, generating *Eisenstein series* $G_k(q)$ in each step:

$$G_k(q) = \begin{cases} 0 & : k \text{ odd} \\ 2\zeta_k + \frac{2(2\pi i)^k}{(k-1)!} \sum_{m,n=1}^{\infty} m^{k-1} q^{mn} & : k \text{ even} \end{cases}$$

Supplement the above ODE in q by boundary values: [Enriquez 1301.3042]

$$\lim_{q \rightarrow 0} \text{eMZV} \left(\begin{array}{l} \text{weight } w \\ \text{length } \ell \end{array} \right) \sim (2\pi i)^{\pm k} \times \text{MZVs}(\text{weight } w \mp k)$$

generating series for eMZV degenerates to Drinfeld associator

III. 2 Direct computation of eMZVs

eMZVs satisfy ODE w.r.t. modular parameter $q = e^{2\pi i\tau}$: [Enriquez 1301.3042]

$$-4\pi^2 q \frac{d}{dq} \text{eMZV} \left(\begin{array}{l} \text{weight } w \\ \text{length } \ell \end{array} \right) \sim G_k(q) \times \text{eMZV} \left(\begin{array}{l} \text{weight } w-k+1 \\ \text{length } \ell-1 \end{array} \right)$$

recursively get q -expansion from *truncated Eisenstein series* G_k^0 :

$$G_k^0(q) \sim \begin{cases} -1 & : k = 0 \\ \sum_{m,n=1}^{\infty} m^{k-1} q^{mn} & : k = 2, 4, 6, \dots \end{cases}$$

\Rightarrow eMZVs are $\mathbb{Q}[\text{MZV}]$ linear combination of iterated Eisenstein integrals

$$\gamma(k_1, k_2, \dots, k_r) \equiv \frac{1}{(4\pi^2)^r} \int_{0 \leq q_i \leq q_{i+1} \leq q} G_{k_1}^0(q_1) \frac{dq_1}{q_1} G_{k_2}^0(q_2) \frac{dq_2}{q_2} \cdots G_{k_r}^0(q_r) \frac{dq_r}{q_r}$$

special cases of iterated Shimura integrals

[Manin 2005; Brown 2014]

Proof of linear independence

[Matthes 1601.05743]

q -expansion of Eisenstein integrals easy to compute

$$\int_0^q (q')^N \frac{dq'}{q'} = \frac{q^N}{N}$$

and turns out to resemble MZVs as nested sums

$$\gamma(k_1, 0^{p_1-1}, k_2, 0^{p_2-1}, \dots, k_r, 0^{p_r-1}) \sim \sum_{0 < n_1 < n_2 < \dots < n_r} \frac{\sigma_{k_1-1}(n_1) \sigma_{k_2-1}(n_2 - n_1) \dots \sigma_{k_r-1}(n_r - n_{r-1}) q^{n_r}}{n_1^{p_1} n_2^{p_2} \dots n_r^{p_r}}$$

$$\zeta(1, 0^{p_1-1}, 1, 0^{p_2-1}, \dots, 1, 0^{p_r-1}) \sim \sum_{0 < n_1 < n_2 < \dots < n_r} \frac{1}{n_1^{p_1} n_2^{p_2} \dots n_r^{p_r}} = \zeta_{p_1, p_2, \dots, p_r}$$

Divisor sums $\sigma_k(n) \equiv \sum_{d|n} d^k$ obstruct analogues of stuffle-relations

\Rightarrow no relations among $\gamma(\dots)$ beyond shuffle $\gamma(W_1)\gamma(W_2) = \gamma(W_1 \sqcup W_2)$

[Matthes 1601.05743]

Length-two example: indecomposable eMZVs $\omega(0, k)$ @ $k = 3, 5, 7, \dots$

$$\omega(0, k) = k\gamma(k+1) = \text{rational} \times \pi^{k-1} \sum_{m, n=1}^{\infty} \frac{m^{k-1}}{n} q^{mn}$$

III. 3 Back to counting indecomposable eMZVs

Naively: indecomposable eMZVs \leftrightarrow shuffle-independent Eisenstein int's?

$$\text{eMZV} \left(\begin{array}{l} \text{weight } \sum_j k_j - r \\ \text{length } r+1 \end{array} \right) \overset{?}{\longleftrightarrow} \frac{\{\gamma(k_1, k_2, \dots, k_r) @ k_j = 0, 2, 4, \dots\}}{\text{shuffle-relations } \gamma(W_1)\gamma(W_2) = \gamma(W_1 \sqcup W_2)}$$

Correct at length two $\omega(0, k) \leftrightarrow \gamma(k+1)$, but \exists problems at length three:

$$\text{e.g. } \underbrace{\omega(0, 0, 12), \omega(0, 3, 9)}_{2 \times \text{indecomposable}} \leftrightarrow \underbrace{\gamma(0, 14), \gamma(2, 12), \gamma(4, 10), \gamma(6, 8)}_{4 \times \text{shuffle-independent}}$$

Additional selection rules arise from derivation algebra $\{\epsilon_0, \epsilon_2, \epsilon_4, \dots\}$

”dual” to Eisenstein series G_0, G_2, G_4, \dots subject to relations

$$[\epsilon_2, \epsilon_k] = 0, \quad [\epsilon_{10}, \epsilon_4] - 3[\epsilon_8, \epsilon_6] = 0, \quad \text{etc.}$$

Explains above mismatch “ $2 \times \omega(\dots)$ vs. $4 \times \gamma(\dots)$ ”

Generating series of eMZVs: elliptic associator [Enriquez 1301.3042]

$$A(q) \equiv \sum_{r \geq 0} (-1)^r \sum_{n_1, n_2, \dots, n_r \geq 0} \omega(n_1, n_2, \dots, n_r) \operatorname{ad}_x^{n_r}(y) \dots \operatorname{ad}_x^{n_2}(y) \operatorname{ad}_x^{n_1}(y)$$

Derivations $\epsilon_0, \epsilon_2, \epsilon_4, \dots$ act on non-commutative variables x, y via

$$\epsilon_{2n}(x) = (\operatorname{ad}_x)^{2n}(y), \quad \epsilon_0(y) = 0, \quad [\text{Nakamura, Tsunogai 1995, ...}]$$

$$\epsilon_{2n}(y) = \sum_{0 \leq j < n} (-1)^j [(\operatorname{ad}_x)^j(y), (\operatorname{ad}_x)^{2n-1-j}(y)] \quad @ \ n > 0$$

satisfy various commutator relations and enter the associator's ODE

$$q \frac{d}{dq} A(q) = \frac{1}{4\pi^2} \left(\sum_{n=0}^{\infty} (2n-1) G_{2n}(q) \epsilon_{2n} \right) A(q).$$

\Rightarrow more precise version of “ ϵ_{2n} are dual to G_{2n} ”

\Rightarrow commutator relations among ϵ_{2n} impose selection rules on $\gamma(\dots)$

Hence, more accurate picture is

$$\text{eMZV} \left(\begin{array}{l} \text{weight } \sum_j k_j - r \\ \text{length } r+1 \end{array} \right) \longleftrightarrow \frac{\{\gamma(k_1, k_2, \dots, k_r) @ k_j = 0, 2, 4, 6, \dots\}}{\{\text{shuffle-relations}\} \times \{\epsilon_k\text{-relations}\}}$$

For instance, $[\epsilon_2, \epsilon_k] = 0$ obstructs non-trivial appearance of $k_j = 2$ and

$$\begin{aligned} \omega(0, n) &\leftrightarrow \gamma(n+1), & \omega(0, 3, 5) &\leftrightarrow \gamma(4, 6) \\ \omega(0, 0, n) &\leftrightarrow \gamma(n+2, 0), & \omega(0, 3, 7) &\leftrightarrow \gamma(4, 8) \\ \left. \begin{array}{l} \omega(0, 3, 9) \\ \omega(0, 5, 7) \end{array} \right\} && \leftrightarrow & 81\gamma(10, 4) + 35\gamma(8, 6) \\ && & \text{selected by } [\epsilon_{10}, \epsilon_4] - 3[\epsilon_8, \epsilon_6] = 0 \end{aligned}$$

Cusp forms control ϵ_k -relations, e.g. [Brown; Hain Matsumoto; Pollack]

wt. 12 cusp form $\leftrightarrow 0 = [\epsilon_{10}, \epsilon_4] - 3[\epsilon_8, \epsilon_6]$, [Ihara, Takao 1993; Schneps 2006]

wt. 16 cusp form $\leftrightarrow 0 = 2[\epsilon_{14}, \epsilon_4] - 7[\epsilon_{12}, \epsilon_6] + 11[\epsilon_{10}, \epsilon_8]$

At length $\ell = 4$, for instance,

$$\text{eMZV} \left(\begin{array}{l} \text{weight } \sum_j k_j - 3 \\ \text{length } 4 \end{array} \right) \longleftrightarrow \frac{\{\gamma(k_1, k_2, k_3) @ k_j = 0, 2, 4, 6, \dots\}}{\{\text{shuffle-relations}\} \times \{\epsilon_k\text{-relations}\}}$$

indecomposable eMZVs governed by **derived relations**, e.g.

$$0 = [\epsilon_k, [\epsilon_{10}, \epsilon_4] - 3[\epsilon_8, \epsilon_6]] = [\epsilon_k, 2[\epsilon_{14}, \epsilon_4] - 7[\epsilon_{12}, \epsilon_6] + 11[\epsilon_{10}, \epsilon_8]]$$

and **irreducible relations** controlled by cusp forms

[Pollack 2009]

$$0 = 80 [\epsilon_{12}, [\epsilon_4, \epsilon_0]] + 16 [\epsilon_4, [\epsilon_{12}, \epsilon_0]] - 250 [\epsilon_{10}, [\epsilon_6, \epsilon_0]] \\ - 125 [\epsilon_6, [\epsilon_{10}, \epsilon_0]] + 280 [\epsilon_8, [\epsilon_8, \epsilon_0]] - 462 [\epsilon_4, [\epsilon_4, \epsilon_8]] - 1725 [\epsilon_6, [\epsilon_6, \epsilon_4]]$$

Perfect matching with
indecomposable eMZVs

[Brödel, Matthes, OS 1507.02254]

wt	1	2	3	4	5	6	7	8	9	10	11	12	13
$\ell=2$	1		1		1		1		1		1		1
$\ell=3$		1		1		1		2		2		2	
$\ell=4$	1		1		2		3		4		5		7

Further directions

- construct one-loop string amplitude from elliptic associator $A(q)$

in analogy to tree-level amplitudes from the Drinfeld associator

[Brödel, OS, Stieberger, Terasoma 1304.7304]

- identify single-valued eMZVs and connect with closed strings

[Pierre Vanhove's talk; D'Hoker, Green, Vanhove 1502.06698 & 1509.00363]

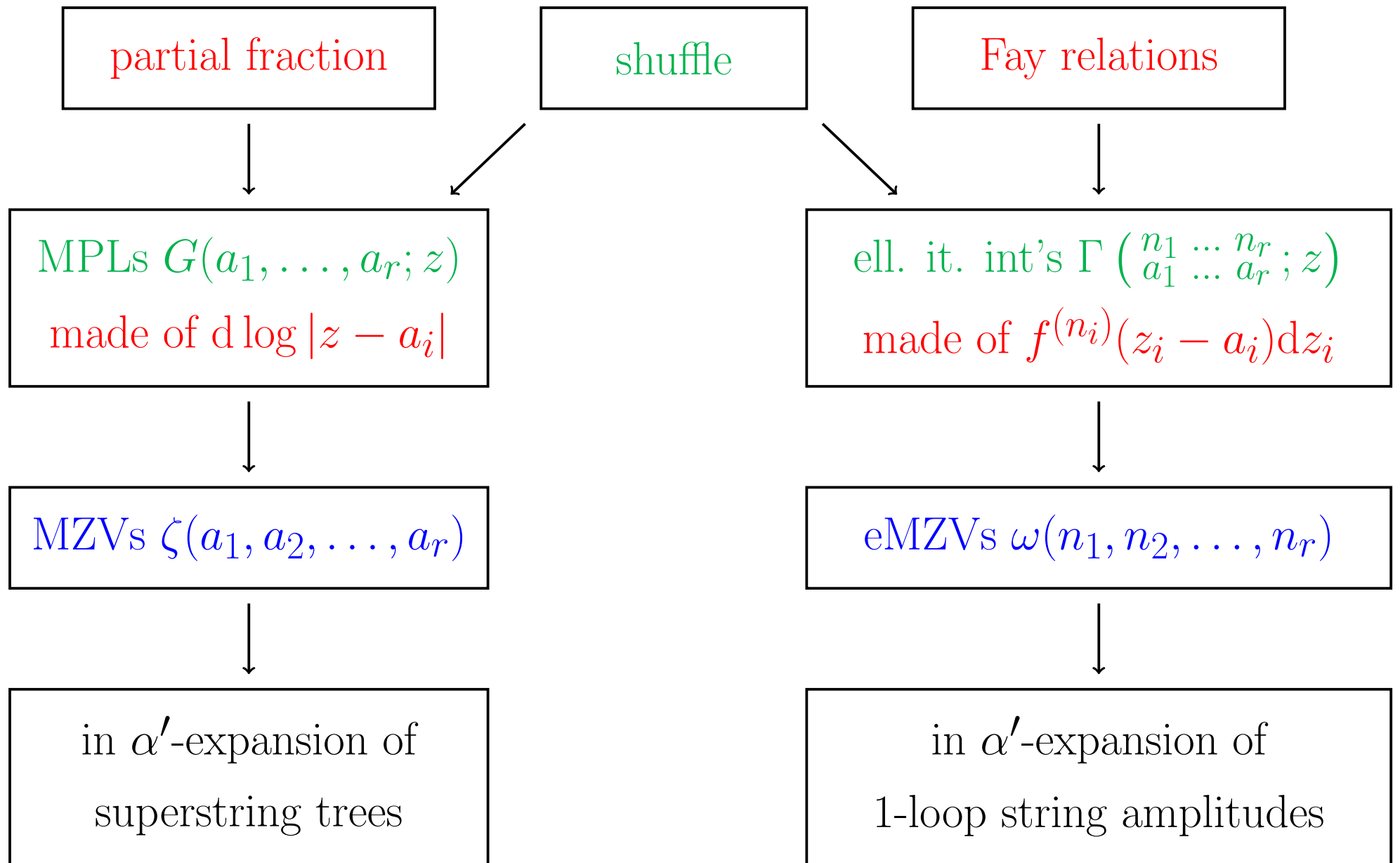
[Zerbini 1512.05689; D'Hoker, Green, Gürdogan, Vanhove 1512.06779]

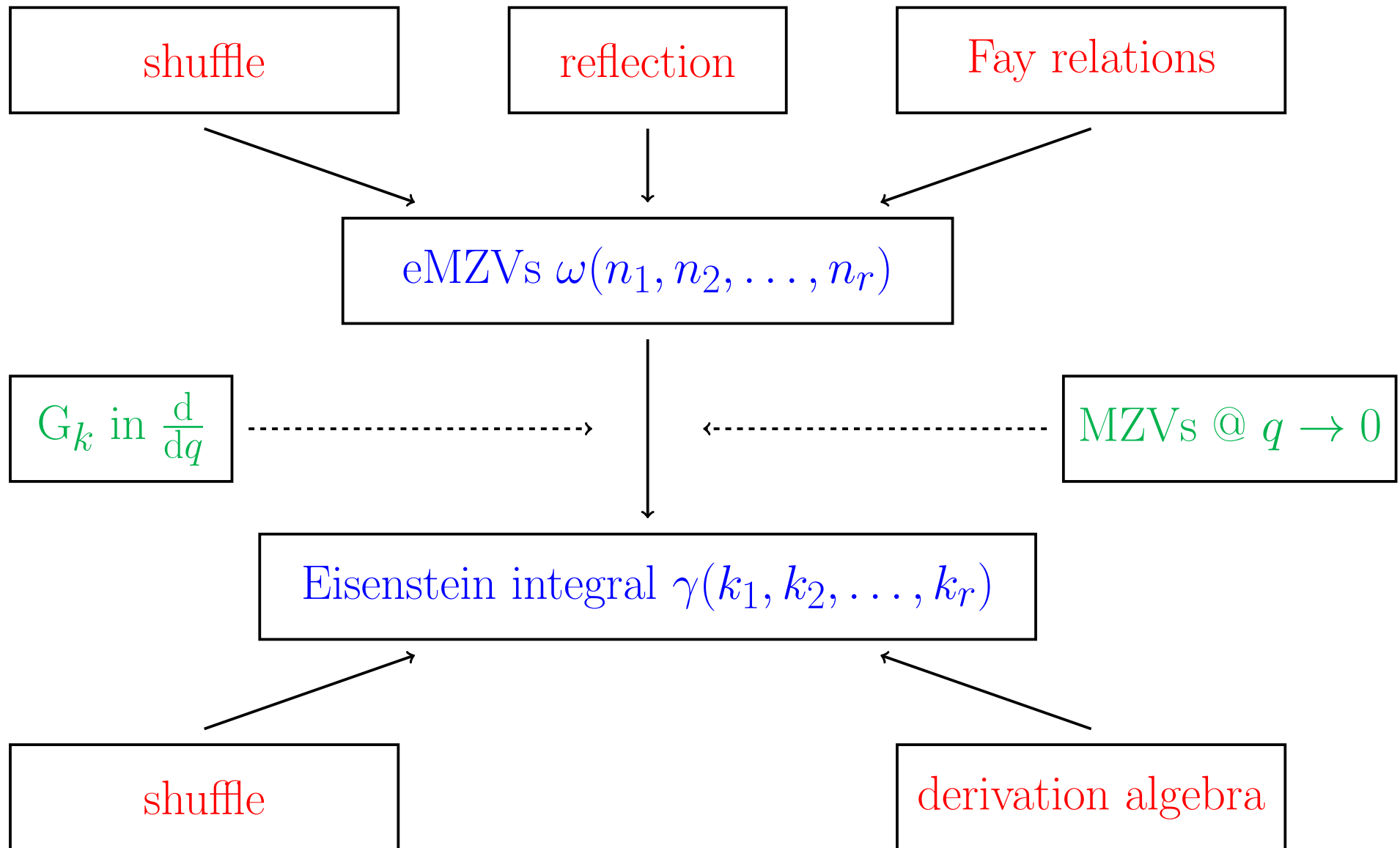
- eMZVs in one-loop string amplitudes of orbifold compactifications

[OS, Berg, Buchberger 1603.05262]

- higher-genus counterparts of MZVs along with single-valued subset

Summary





Various eMZV relations available at <https://tools.aei.mpg.de/emzv>