

**Tracking studies regarding the electron
model 3-5.4463 MeV of a non-linear,
non-scaling proton driver FFAG
3-10 GeV.**

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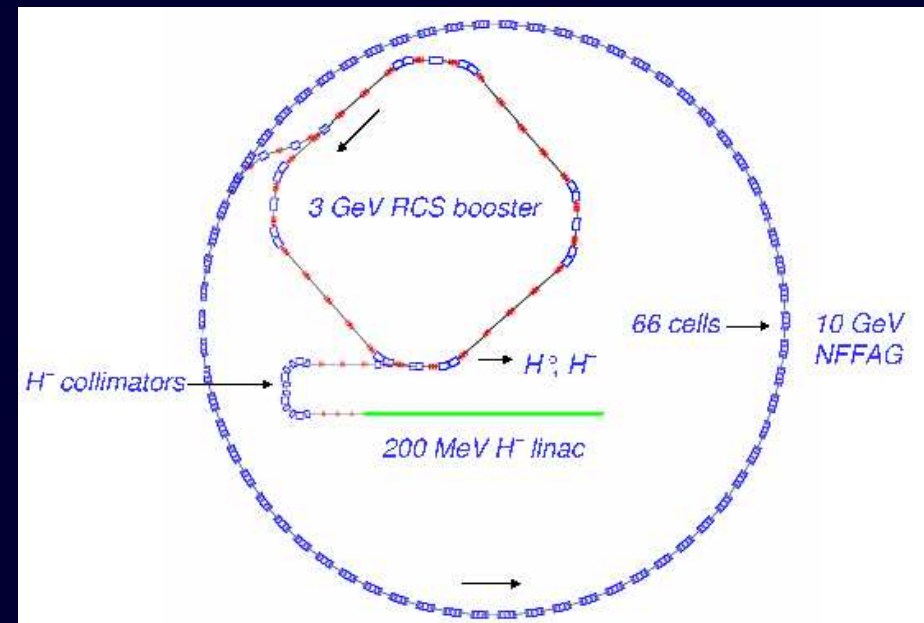
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Introduction

- The 50 Hz, 3-10 GeV, proton driver uses a new type of FFAG, with a non-linear, non-scaling and non-isochronous, cell focusing structure.
- An electron model 3-5.4463 MeV has been designed to test various aspects of the optics.



- We describe the setting up of the optics data files and produce preliminary beam dynamics studies concerning in this electron model, using the ray-tracing code ZGOUBI.
- Similar investigations were carried out for a non-scaling, non-linear, isochronous muon ring proposed for the 8-20 GeV fast acceleration of muons in a neutrino factory[†] and for its associated 11-20 MeV electron model[‡].

[†] 6-D beam dynamics in an isochronous FFAG ring, F. Lemuet, F. Méot, G. Rees, proceedings of 2005 Particle Accelerator Conference, Knoxville, Tennessee.

[‡] 6-D beam dynamics in an electron model lattice of an isochronous FFAG, F. Lemuet, F. Méot, G. Rees, Nuclear physics B, 155 (2006) pp 330-331.

Ray-tracing method ZGOUBI

- Solve the Lorentz equation $\vec{u}' = \vec{u} \times \vec{B}$ by Taylor developments of speed and position

$$\vec{R}(M_1) \approx \vec{R}(M_0) + \vec{u}(M_0) \Delta s + \vec{u}'(M_0) \frac{\Delta s^2}{2!} + \dots + \vec{u}''''(M_0) \frac{\Delta s^6}{6!} [+...]$$

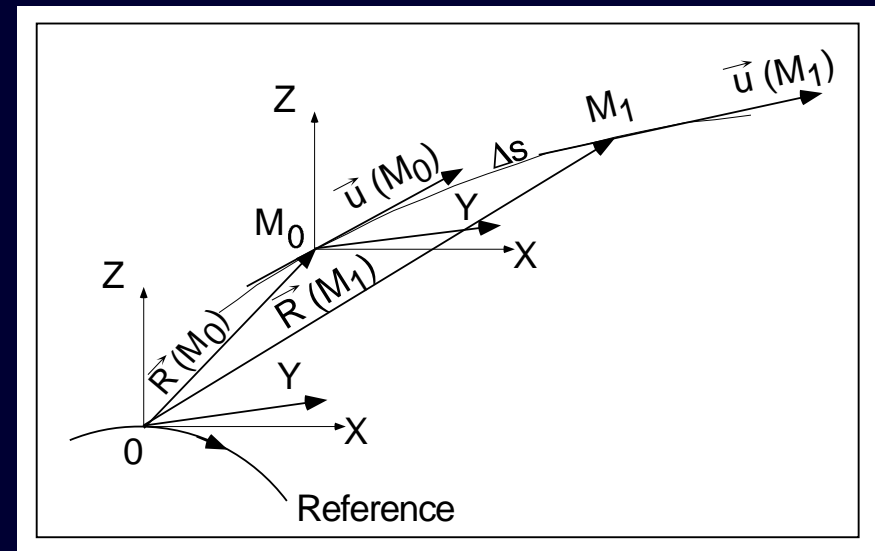
$$\vec{u}(M_1) \approx \vec{u}(M_0) + \vec{u}'(M_0) \Delta s + \vec{u}''(M_0) \frac{\Delta s^2}{2!} + \dots + \vec{u}''''(M_0) \frac{\Delta s^5}{5!} [+...]$$

- Coefficients are calculated by successive derivations

$$\vec{u}'' = \vec{u}' \times \vec{B} + \vec{u} \times \vec{B}', \quad \vec{u}''' = \dots \text{ etc.}$$

- A field model is needed to calculate the field and its derivatives

$$\vec{B}(X, Y, Z), \quad \partial^{i+j+k} \vec{B} / \partial X^i \partial Y^j \partial Z^k$$



- Derivatives $d^n \vec{B} / ds^n$ are calculated ($X_{i,j,\dots}$, $i, j, \dots = 1, 3$ for X, Y or Z). using

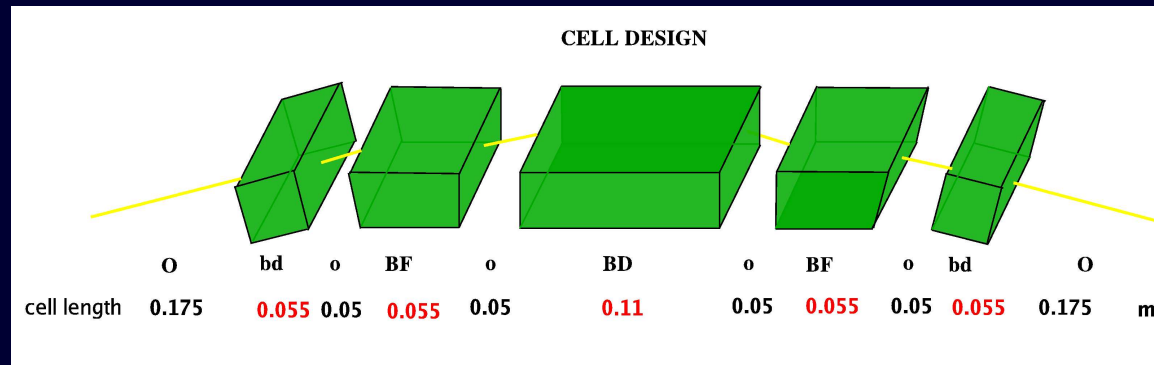
$$\vec{B}'(s) = \sum_i \frac{\partial \vec{B}(X, Y, Z)}{\partial X_i} u_i(s), \quad \vec{B}''(s) = \sum_{ij} \frac{\partial^2 \vec{B}(X, Y, Z)}{\partial X_i \partial X_j} u_i(s) u_j(s) + \sum_i \frac{\partial \vec{B}(X, Y, Z)}{\partial X_i} u_i'(s)$$

$$\vec{B}''' = \sum_{ijk} \frac{\partial^3 \vec{B}}{\partial X_i \partial X_j \partial X_k} \text{ etc.}$$

Fields and cell parameters

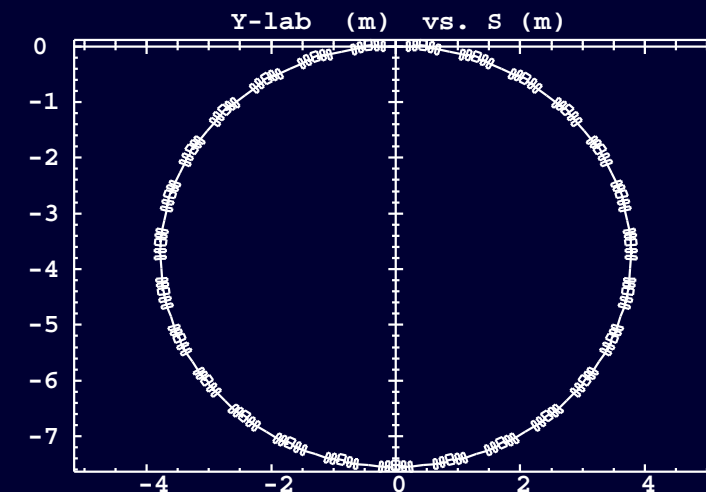
■ Geometrical lattice design

■ The design is based on a O-bd-o-BF-o-BD-o-BF-o-BD-O cell, symmetric wrt. the center of BD, built from magnets with a straight optical axis and rectangular $B(x)$ dependence.



■ The bd and BD magnets have parallel faces, while the second magnet, BF, has its upstream (resp. downstream) edge parallel to the downstream (resp. upstream) edge of bd (resp. BD).

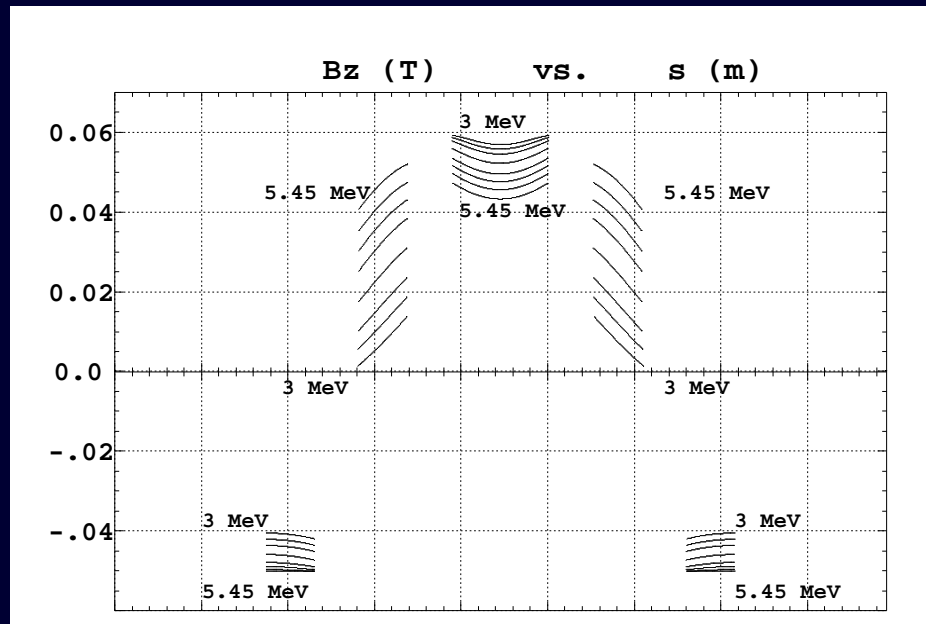
■ The rings consists of 27 cells yielding a circumference of 23.76 m.



■ Gradients and fields

■ The magnets are non-linear and combined function :

- bd, BD are horizontally defocusing, respectively reverse and positive bend.
- BF is horizontally focusing with positive bend.



■ In our simulations we use rectangular multipole with straight axis for all 3 types of magnets, and a hard-edge model.

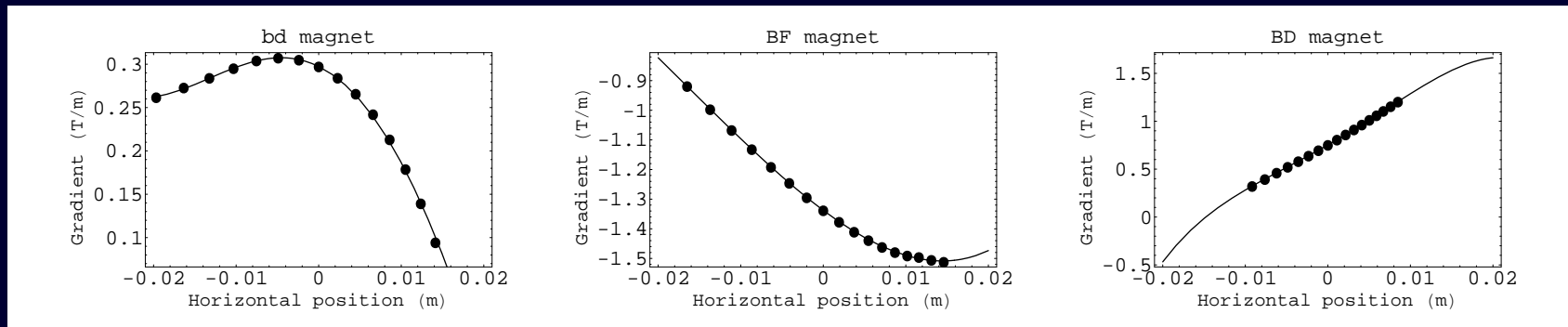
■ First order effect of fringe field extent on vertical motion is accounted for via a correction vertical kick of the form $-\tan(\alpha)/\rho + FR/(6 * \rho^2 * \cos(\alpha))$ with ρ being the local curvature radius and α the angle between the trajectory of the particle and the normal to the magnet face.

- The local magnet strengths, given by the design data, yield the local gradients (dots) using the relation,

$$g[T/m] = \frac{p[GeV/c]K[m^{-2}]}{0.2998}$$

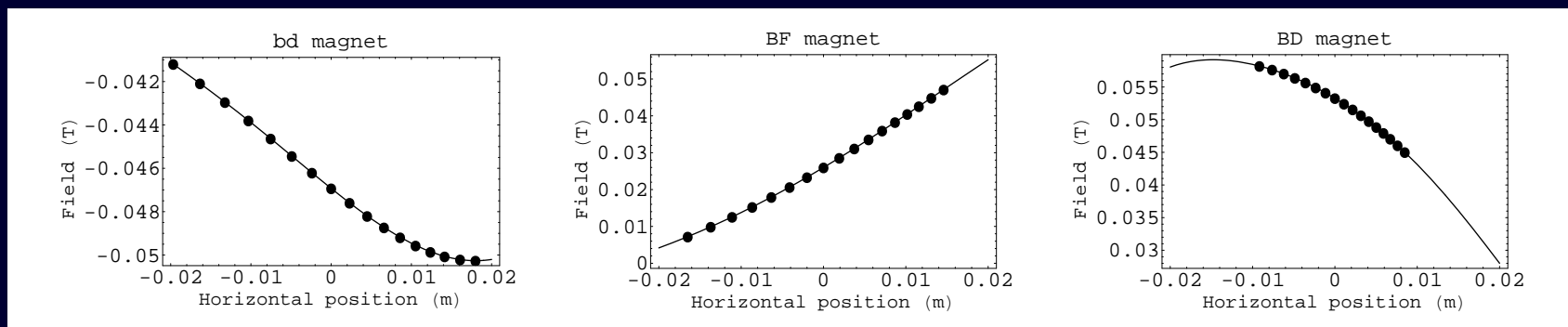
these gradients are approximated using three 4th degree polynomials (lines),

$$g(x) = g_0 + g_1 x + g_2 x^2 + g_3 x^3 + g_4 x^4$$



which are next integrated so to obtain the magnetic field law in each magnet

$$b(x) = b_0 + b_1 x + b_2 x^2 + b_3 x^3 + b_4 x^4 + b_5 x^5 \quad (1)$$



- The simulation of the rectangular dependence $B(x)$ in bd, BF, BD magnets uses a classical multipole modelling of the form

$$\vec{B} = \text{grad}V_n$$

$$V_n(s, x, z) = (n!)^2 \left(\sum_{q=0}^{\infty} \frac{(-)^q G^{(2q)}(s)(x^2 + z^2)^q}{4^q q!(n + q)!} \right) \left(\sum_{m=0}^n \frac{\sin\left(\frac{m\pi}{2}\right) x^{n-m} z^m}{m!(n - m)!} \right)$$

with coefficient values $G^{(2q)}(center)$ derived from equation 1.

- The s -dependence $G(s)$ allows the simulation of field fall-offs at magnet ends to be included when desired, using a Enge's fall-off model.

■ Geometrical matching procedure

- The constraints imposed are that the closed orbit angles in the drifts between the magnets be as close as possible to the design ones.
- The parameters to be varied are the closed orbit coordinates at the start of the cell, and the positioning of the magnets via rotation and horizontal translation.
- Satisfactory c.o. angle values have been obtained for three energies 3 MeV, 4.05 MeV, 5.4463 MeV with the consequence that c.o. angles at the other energies considered in the original design take a value that agrees with the design data to the level of a percent and better.

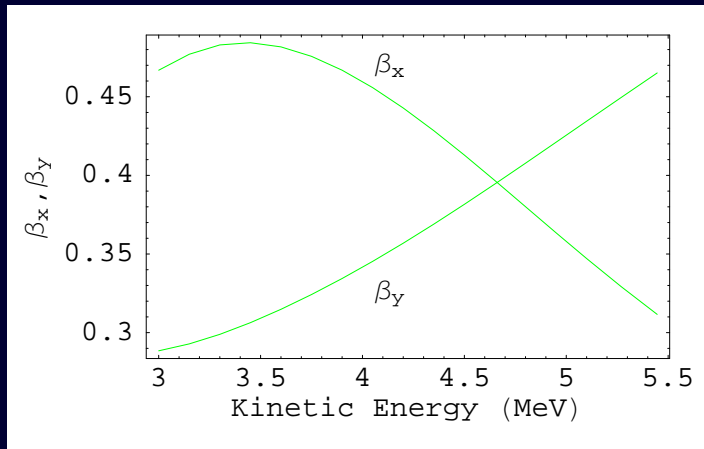
First order results

We present first order tracking results in a cell adjusted with the matching procedure.

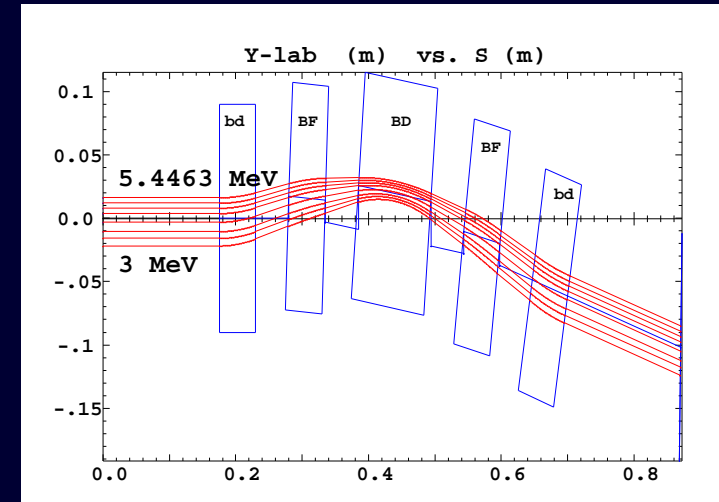
■ Closed orbits

■ Closed orbits for the electron model 3-5.4463 MeV, as the energy increases the trajectories move from the inner to the outer part of the magnets.

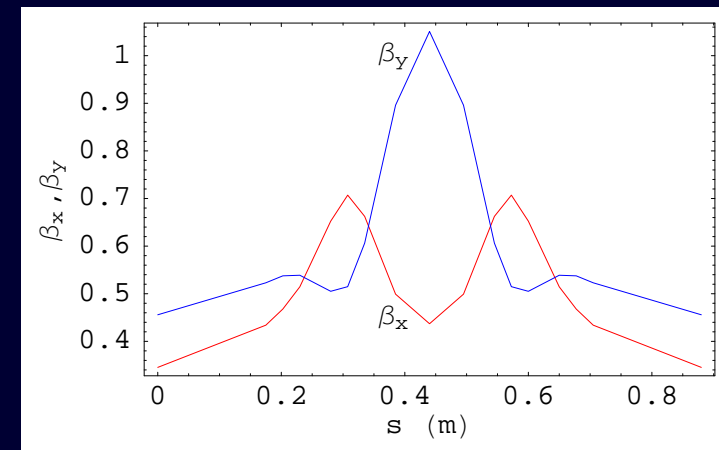
■ Focusing



■ β functions as a function of s for the reference energy 4.05 MeV.

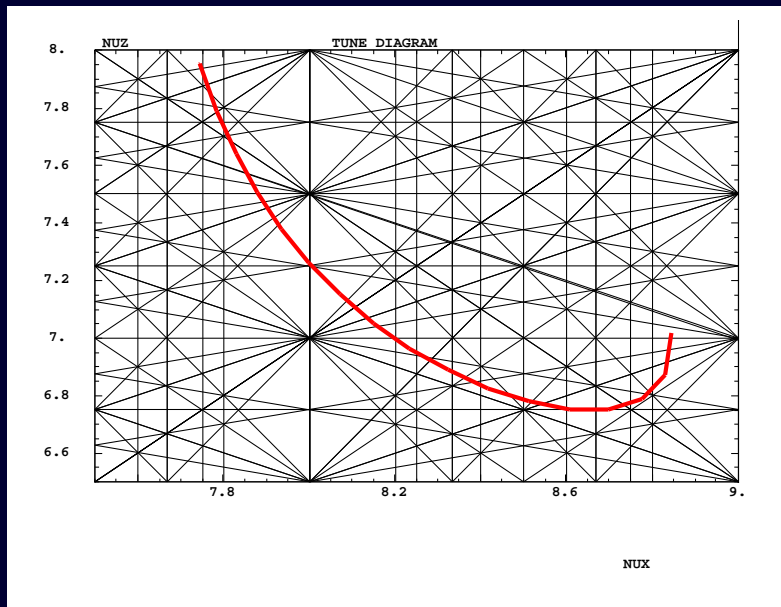
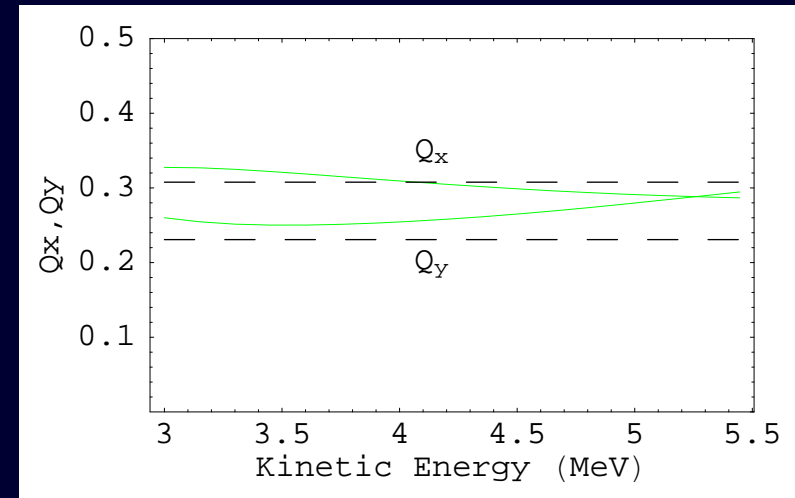


■ Periodic β values at the cell end as a function of energies.



- The design tunes per cell are constant with $Q_x = \frac{4}{13}$ and $Q_y = \frac{3}{13}$ (dashed lines).

- The tunes per cell calculated with ray-tracing (green lines) are slightly different with a small variation with energy.



- Total tunes ($27 \times Q_{x,y}$) variations in the tune diagram for $E : 3 \rightarrow 5.4463$ MeV in a systematic resonance lines up to the fifth order.

- If a better agreement is needed, a matching procedure could be performed in ZGOUBI to match the tunes, which allows an automatic adjustment of the b_i coefficients in the multipole expansion

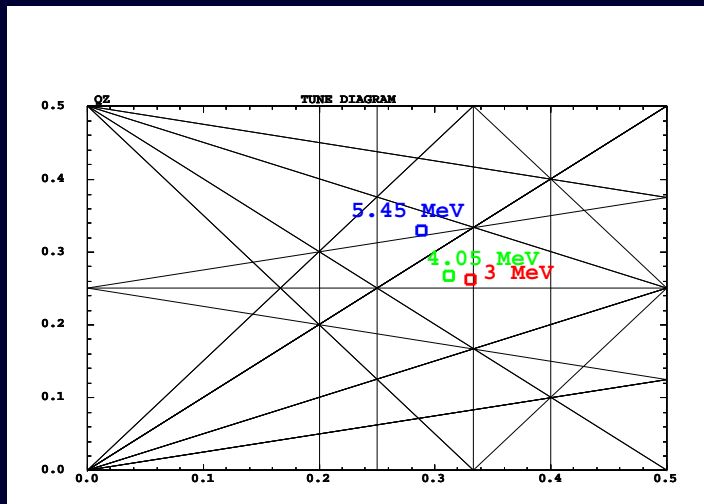
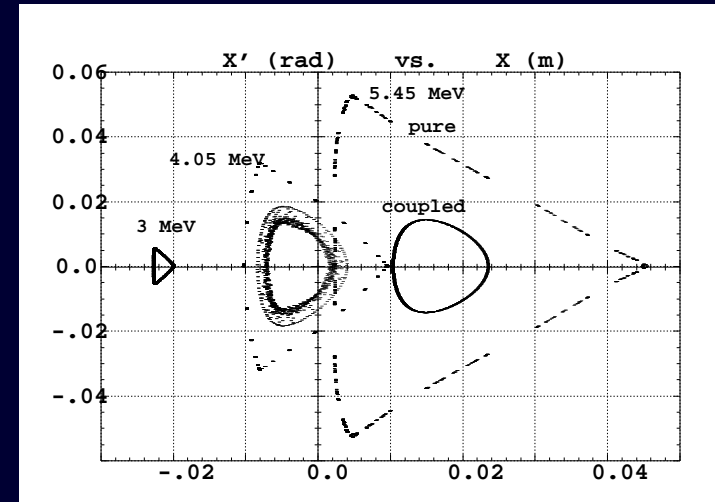
$$b(x) = \sum_{i=0}^n b_i x^i$$

Dynamical aperture

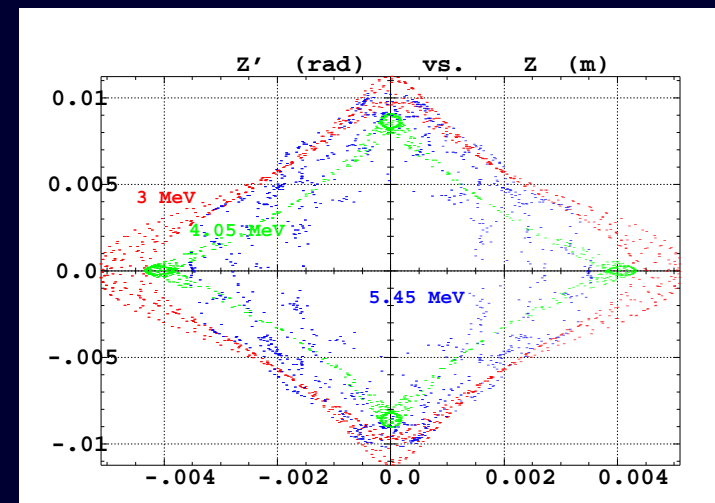
■ Stability limits

■ Limit phase space trajectories obtained by multiturn tracking at injection, reference and extraction energies in case of a pure horizontal motion and when we introduce a small z motion.

■ In that case the limits decrease sensibly, except for the injection energy for which the limit is unchanged.



■ Total tunes corresponding to the coupled motion are displayed in a fifth order tune diagram.

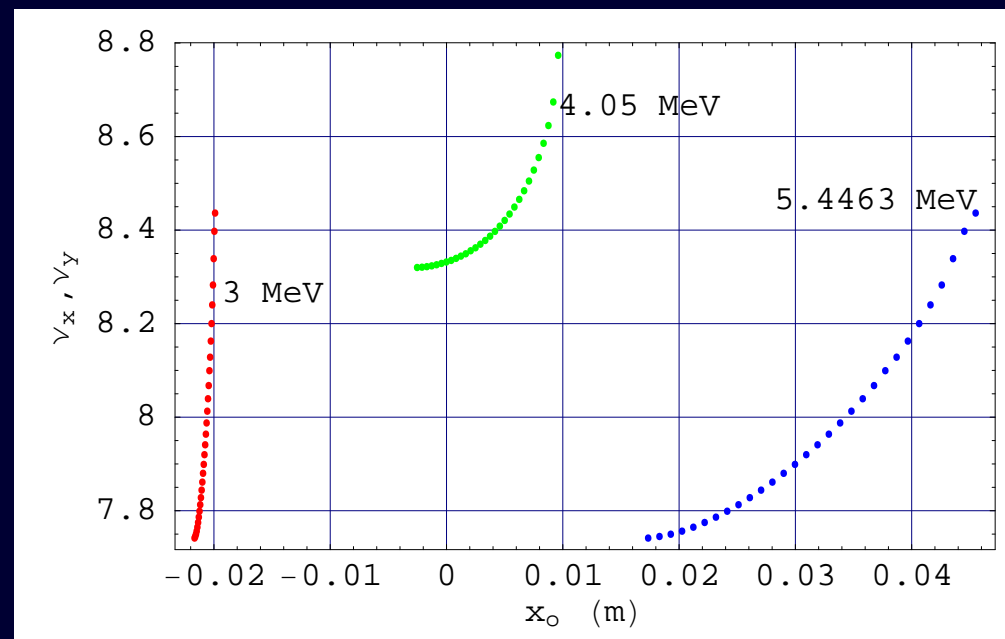


■ Vertical stability limits for the particles launched on the horizontal closed orbits.

■ Amplitude Detuning

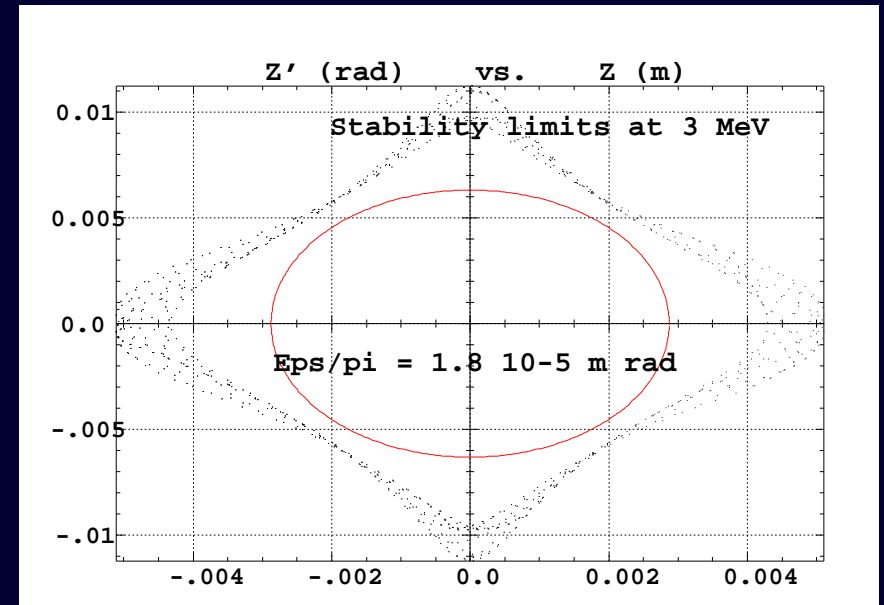
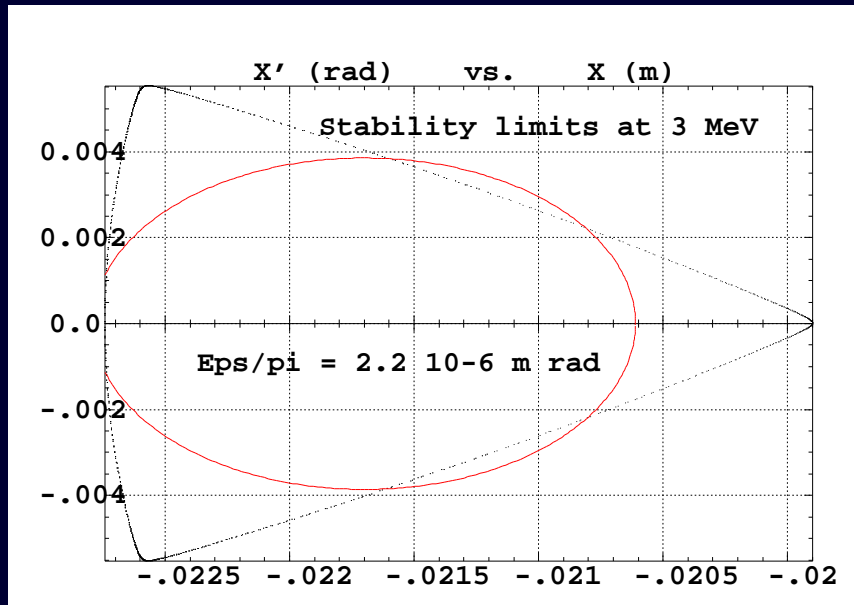
■ Amplitude detuning in case of pure x motion for the 3, 4.05, 5.4663 MeV energies.

■ The total tunes have been calculated for the three energies starting from the closed orbits x_{co} up to the x_{limits} calculated previously.



■ Acceptance

- Horizontal and vertical dynamical apertures at 3 MeV calculated previously and their corresponding matched ellipse.



- The acceptances as the surfaces of these ellipses are :

- $\frac{\epsilon_x}{\pi} = 2.2 \cdot 10^{-6} \text{ m rad}$
- $\frac{\epsilon_z}{\pi} = 1.8 \cdot 10^{-5} \text{ m rad.}$

Conclusions

- Preliminary tracking studies have been performed in the non-scaling, non-linear electron model of a 3-10 GeV proton driver.
- A more exhaustive investigation could now be carried out.
- Further steps would include acceleration and full transmission studies as it has been performed in the isochronous muon ring and its electron model.