

Light Higgs bosons in Two-Higgs-Doublet Models

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Based on [\[arXiv:1412.3385\]](https://arxiv.org/abs/1412.3385)

In collaboration with

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Motivations

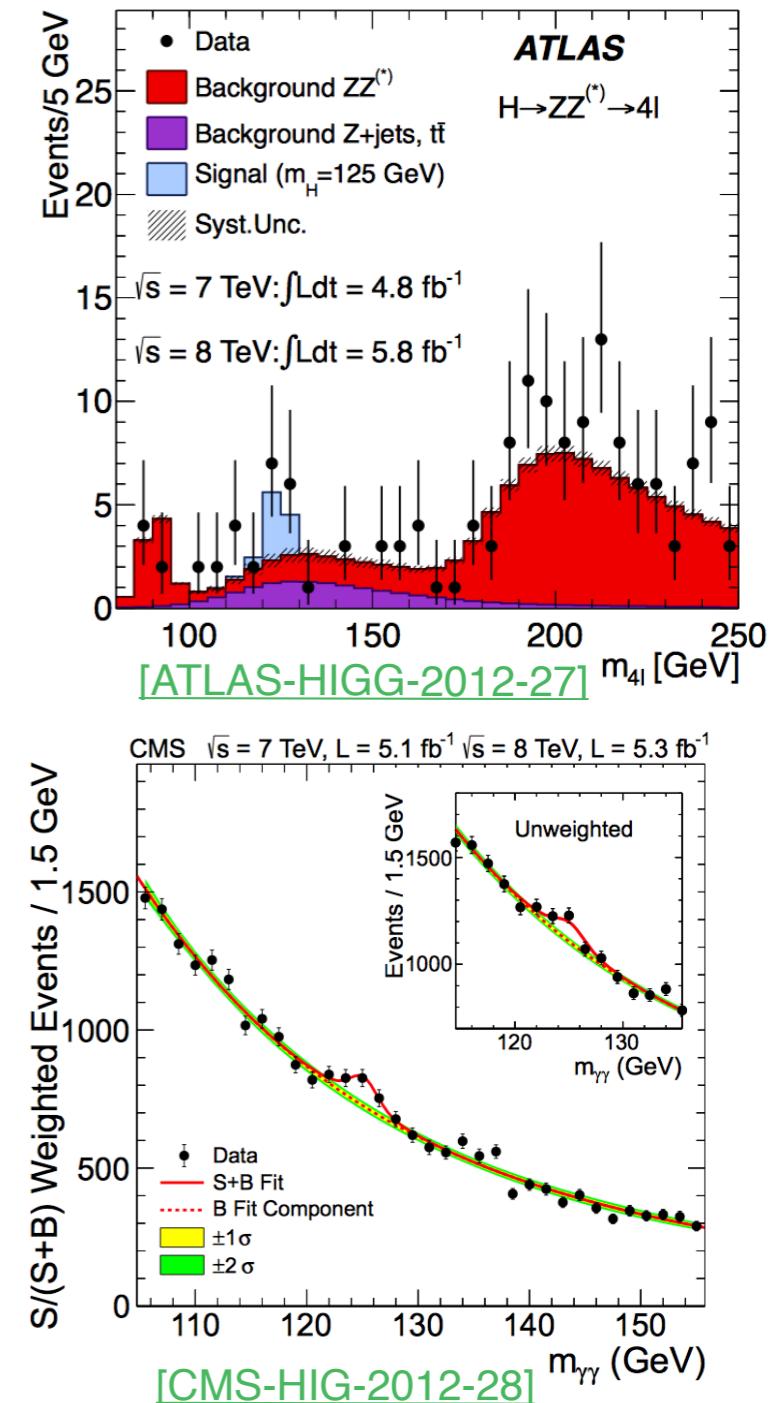
The 2012 discovery of a **Standard Model (SM)-like Higgs boson** is the **major achievement** of the LHC Run I

It is important to asses all possibilities regarding the **existence** of **other Higgs states**

Two-Higgs-doublet models (THDM) are a simple and appealing framework to study such considerations

The **decoupling limit** is often considered to obtain a Higgs state with SM properties

Here, we consider scenarios in which some of the Higgs states can be **very light**: with **mass <125/2**



Two-Higgs-doublet models

General presentation

Constraints imposed

Presence of light states

Two-Higgs-doublet models

- **Two-Higgs-doublet model** (2HDM): minimal extension of the SM, include a second $Y=+1$ Higgs doublet

$$\begin{aligned}\mathcal{V} = & m_{11}^2 \Phi_1^\dagger \Phi_1 + m_{22}^2 \Phi_2^\dagger \Phi_2 - [\underline{m_{12}^2 \Phi_1^\dagger \Phi_2} + \text{h.c.}] \\ & + \frac{1}{2} \lambda_1 (\Phi_1^\dagger \Phi_1)^2 + \frac{1}{2} \lambda_2 (\Phi_2^\dagger \Phi_2)^2 + \lambda_3 (\Phi_1^\dagger \Phi_1)(\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2)(\Phi_2^\dagger \Phi_1) \\ & + \left\{ \frac{1}{2} \lambda_5 (\Phi_1^\dagger \Phi_2)^2 + [\lambda_6 (\Phi_1^\dagger \Phi_1) + \lambda_7 (\Phi_2^\dagger \Phi_2)] \Phi_1^\dagger \Phi_2 + \text{h.c.} \right\}. \quad \Phi_1 \rightarrow \Phi_1, \Phi_2 \rightarrow -\Phi_2\end{aligned}$$

- Hypotheses: [Softly broken Z2 symmetry](#), no CP-violation, no tree-level flavor changing neutral current (FCNC)
- Five physical degrees of freedom: 2 CP-even (h, H) ($m_h \leq m_H$), 1 CP-odd (A), 2 charged (H^+, H^-) states
- Both h and H can be identified with the SM-like state, we will consider both **h125** and **H125 scenarios**
- Free parameters: $m_h, m_H, m_A, m_{H^\pm}, m_{12}^2, \tan \beta$ $[0.5, 60]$, α $[-\pi/2, \pi/2]$
 $\tan \beta$: ratio of the 2 Higgs vevs, α : mixing angle of the CP-even mass matrix

Flavor structure & couplings

Most general renormalizable Yukawa sector:

$$\begin{aligned}-\mathcal{L}_{\text{Yuk}} = & \mathcal{Y}_1^U \bar{U} \Phi_1 Q + \mathcal{Y}_2^U \bar{U} \Phi_2 Q + \mathcal{Y}_1^D \bar{D} \Phi_1^* Q + \mathcal{Y}_2^D \bar{D} \Phi_2^* Q \\ & + \mathcal{Y}_1^E \bar{E} \Phi_1^* L + \mathcal{Y}_2^E \bar{E} \Phi_2^* L + \text{h.c.}\end{aligned}$$

⇒ Generic tree-level FCNC

Four discrete choices, to **insure absence of tree-level FCNC** [Paschos 77',
Glashow & Weinberg 77']

We consider two of them, the so-called **Type I** and **Type II models**

Type I: $\Phi_1 \leftrightarrow$ up, down-type fermions

Type II: $\Phi_1 \leftrightarrow$ down-type fermions

$\Phi_2 \leftrightarrow$ up-type fermions

	Type I and II	Type I		Type II		
		C_V	C_U	C_D	C_U	C_D
Higgs						
h	$\sin(\beta - \alpha)$	c_α/s_β	c_α/s_β	c_α/s_β	$-s_\alpha/c_\beta$	
H	$\cos(\beta - \alpha)$	s_α/s_β	s_α/s_β	s_α/s_β	c_α/c_β	
A	0	$\cot\beta$	$-\cot\beta$	$\cot\beta$	$\tan\beta$	

$$C_i = C_i^{\text{2HDM}} / C_i^{\text{SM}}$$

Numerical setup & Constraints

Numerical setup:

- Branching ratio and theoretical constraints from **2HDMC** [Eriksson, Rathsman, Stål] [[arXiv:0902.0851](#)]
- Cross sections from **SusHi** [Herlander, Liebler, Mantler] [[arXiv:1212.3942](#)]

Theoretical:

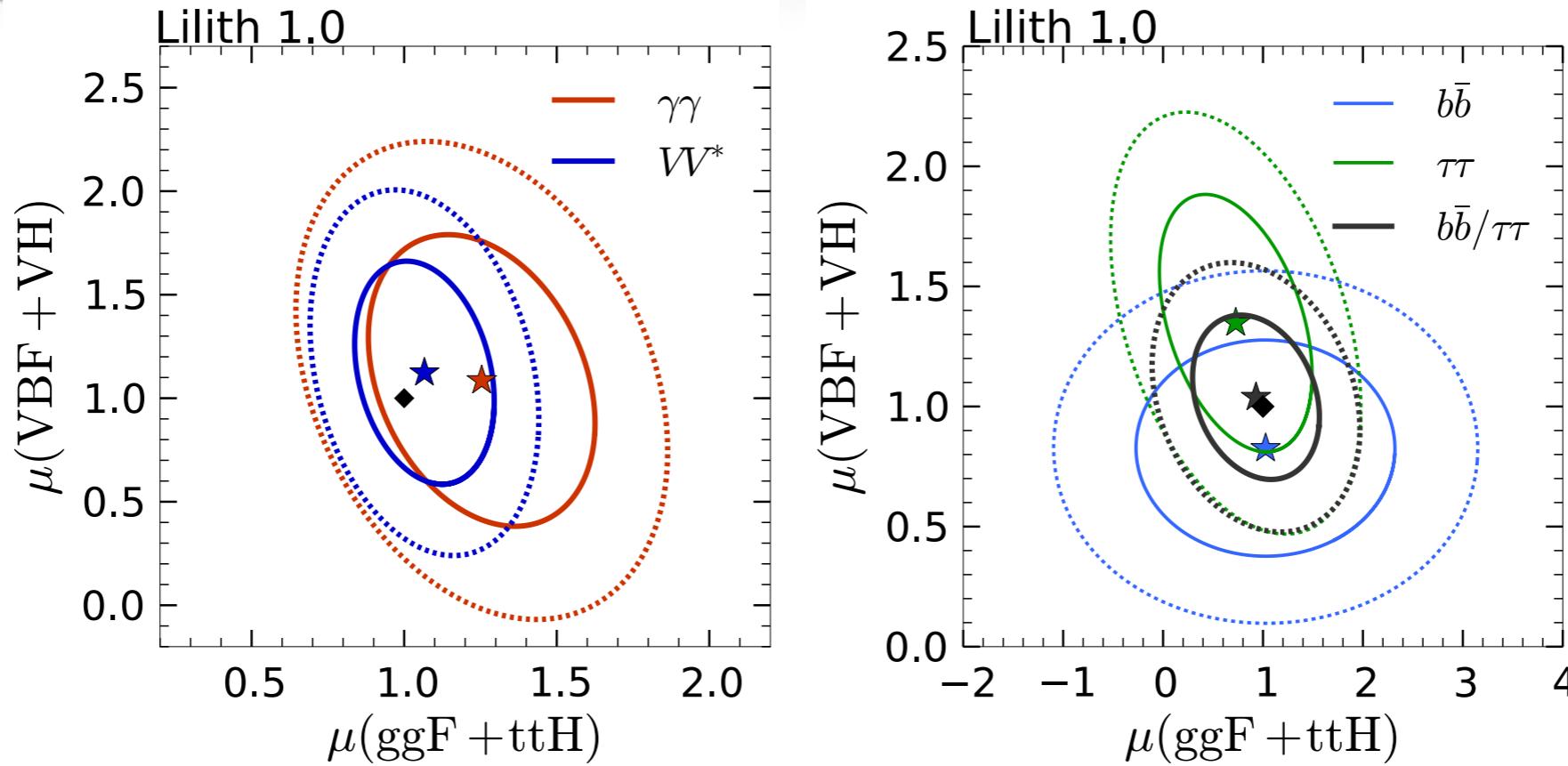
- **Stability** of the scalar potential
- **Perturbativity** of the self-couplings
- Tree-level **unitarity** of the Higgs-Higgs scattering matrices

Experimental:

- **S, T, U** Peskin-Takeuchi parameters (\rightarrow Higgs mass splitting)
- **B-physics** constraints (\rightarrow lower bound on charged Higgs mass)
- **LHC** heavy Higgs searches ($H \rightarrow ZZ$, $A \rightarrow \tau\tau$, $gg \rightarrow bbA \rightarrow bb\tau\tau \dots$)
- **LEP** Higgs searches ($e^+e^- \rightarrow Zh$, $e^+e^- \rightarrow Z^* \rightarrow Ah$, $e^+e^- \rightarrow H^+H^-$)
- 125 GeV Higgs **signal strengths**

Combined signal strengths

Status as of
October 2014



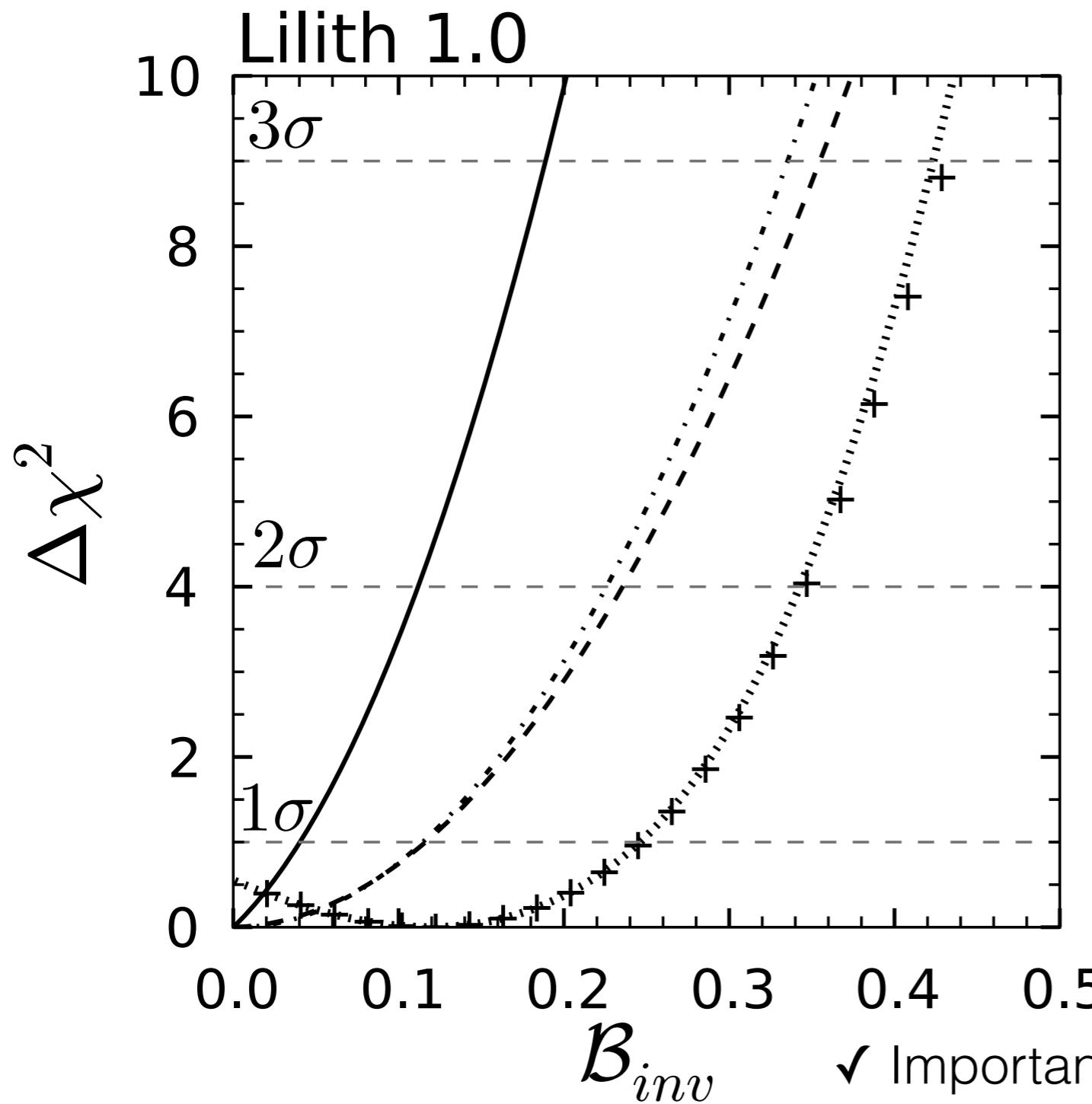
[JB, B. Dumont, S. Kraml] [arXiv:1409.1588]

Combining signal strength measurements from LHC and Tevatron, one obtains an approximation to the **Higgs likelihood**

We **require** 95% C.L. agreement with the combined signal strengths for **all individual decay modes** ($\gamma\gamma, WW^*, ZZ^*, b\bar{b}, \tau\tau$)

Lilith
Light Likelihood fit for the Higgs
[JB, B. Dumont] [arXiv:1502.04138]
Information, Download:
<http://ipsc.in2p3.fr/projects-th/lilith/>
(Google: **lilith higgs**)

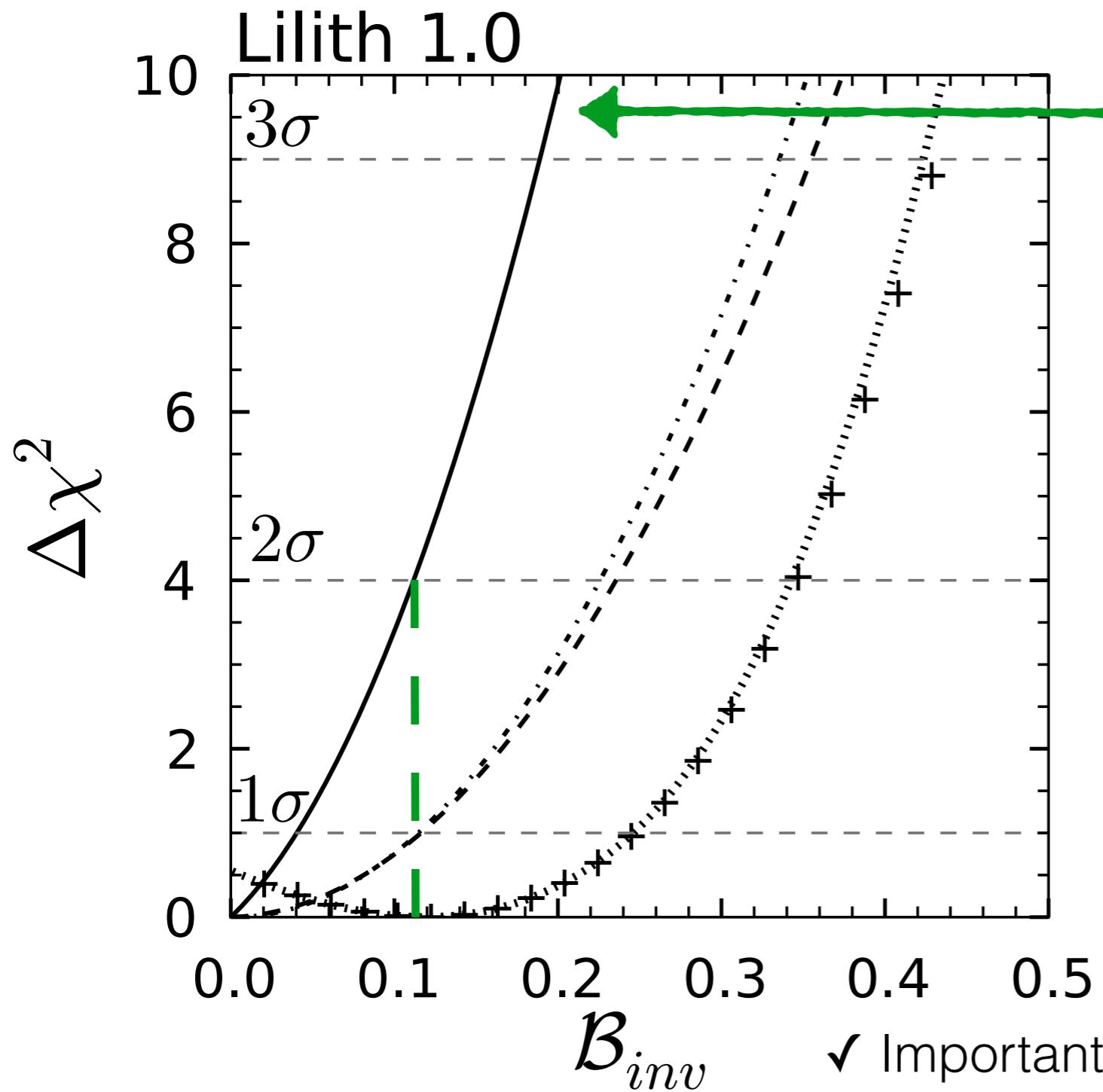
Invisible/undetected branching ratio constraints



✓ Important constraints in the light Higgs scenarios

[JB, B. Dumont, S. Kraml] [arXiv:1409.1588]

Invisible/undetected branching ratio constraints

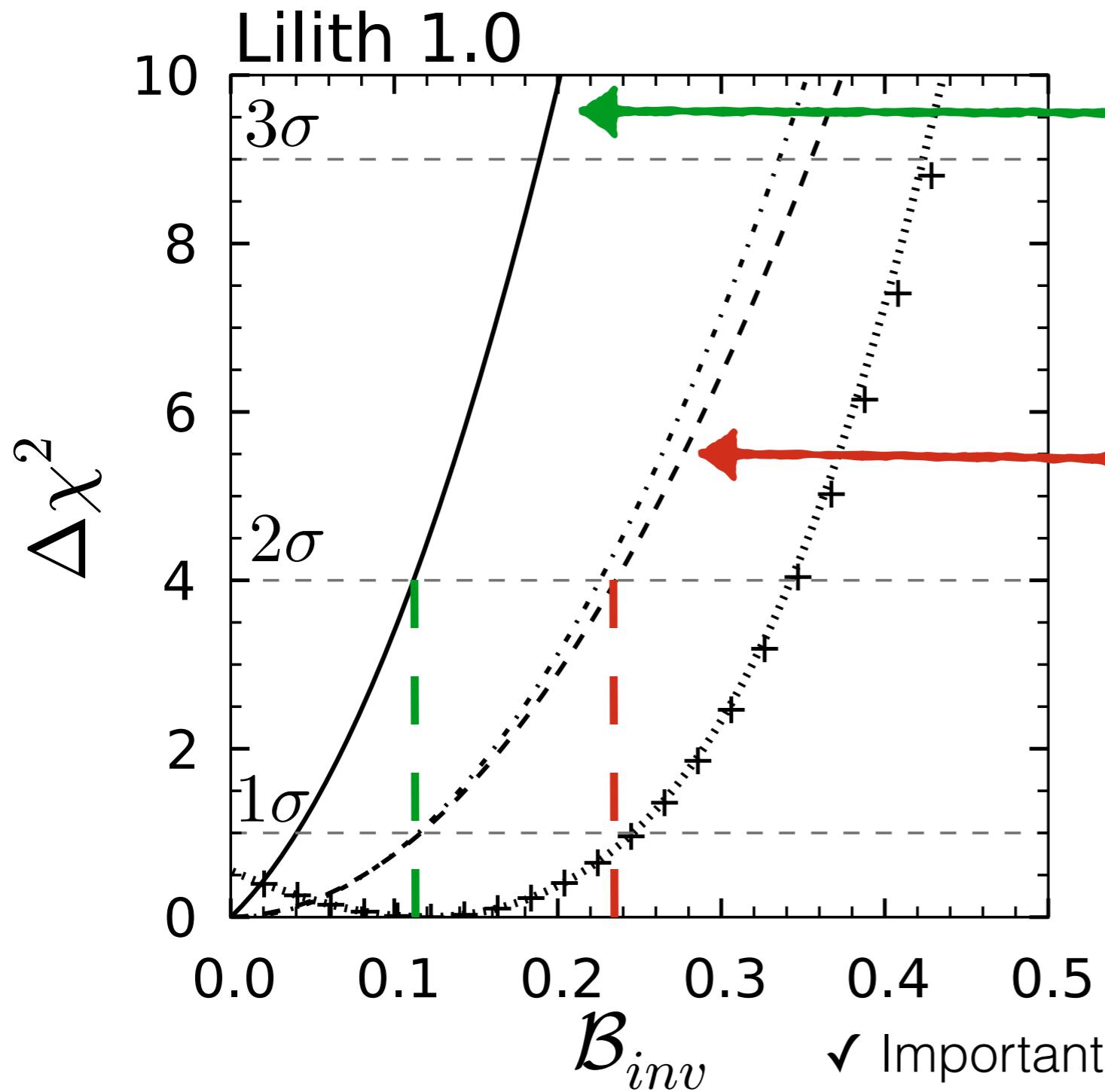


SM+invisible
 $\mathcal{B}_{inv} < 0.11$ at 95.4% C.L.

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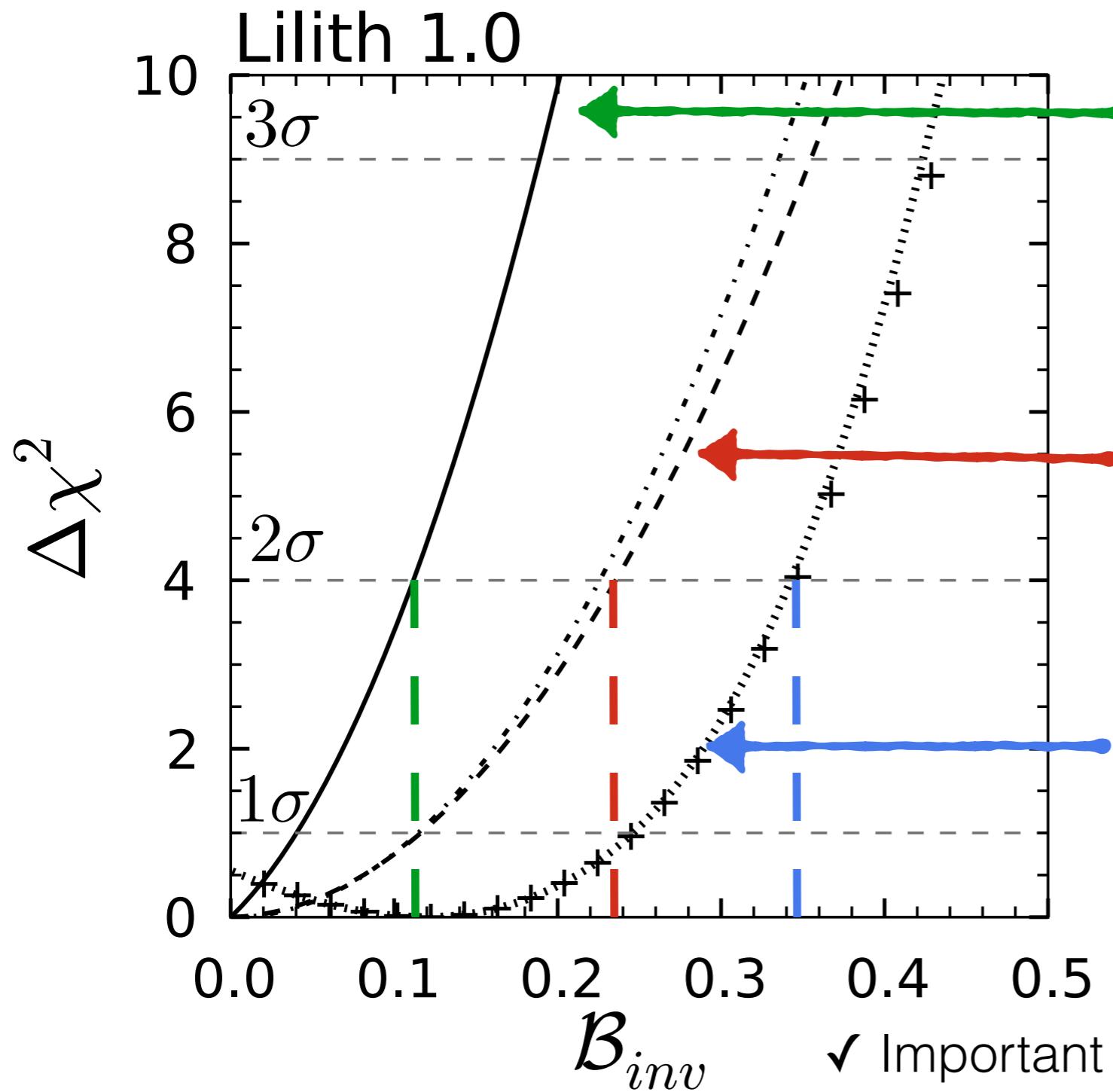
$C_U, C_D, C_V < 1$ + invisible

$\mathcal{B}_{inv} \lesssim 0.24$ at 95.4% C.L.

✓ Important constraints in the light Higgs scenarios

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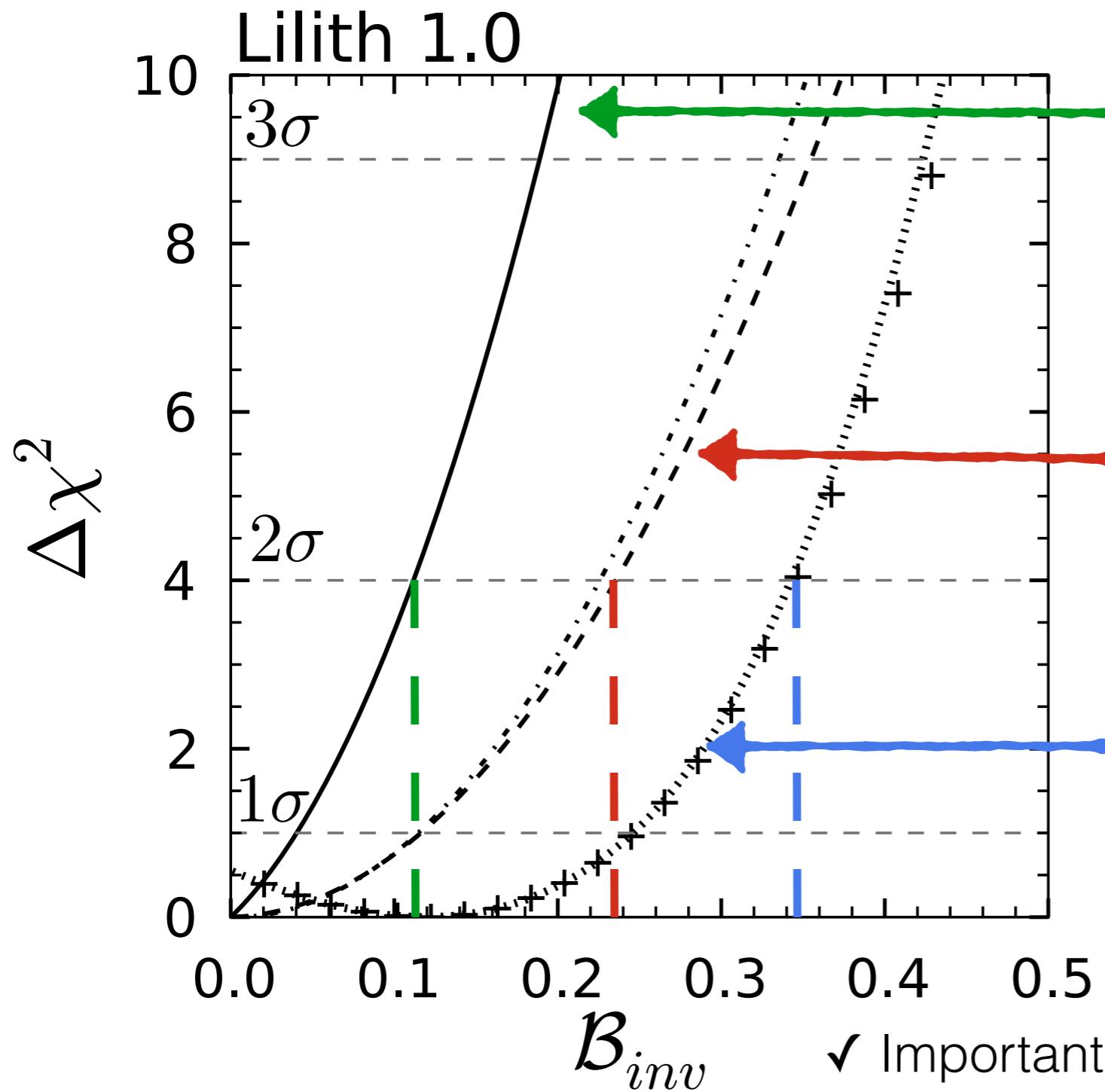
$C_U, C_D, C_V < 1$ + invisible

$\mathcal{B}_{inv} \lesssim 0.24$ at 95.4% C.L.

$C_U, C_D, C_V + \text{invisible}$

$\mathcal{B}_{inv} \lesssim 0.34$ at 95.4% C.L.

Invisible/undetected branching ratio constraints



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$\mathcal{B}_{inv} < 0.11$ at 95.4% C.L.

$C_U, C_D, C_V < 1 + \text{invisible}$

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Roughly the same limits
for undetected BR

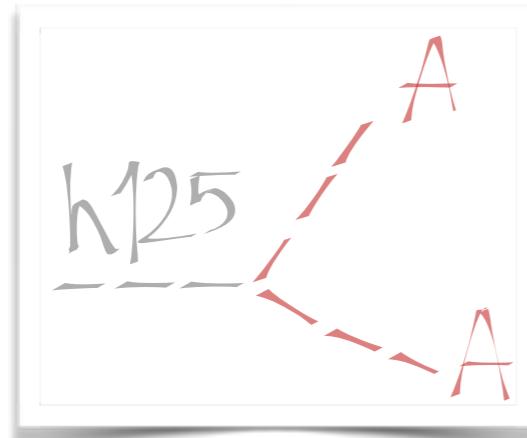
$C_U, C_D, C_V + \text{invisible}$

$\mathcal{B}_{inv} \lesssim 0.34$ at 95.4% C.L.

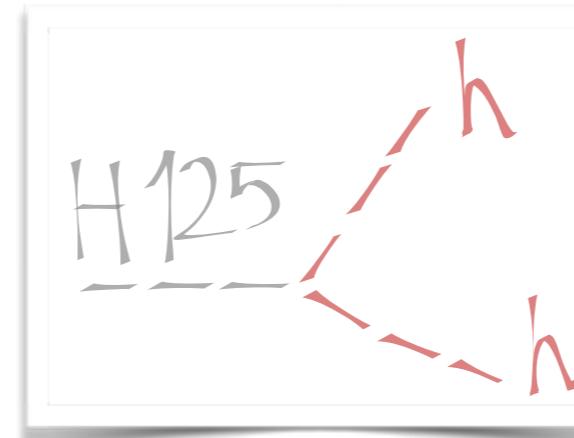
Light states

We consider the presence of light states $m < 125/2$ in both the $h125$ and $H125$ scenarios:

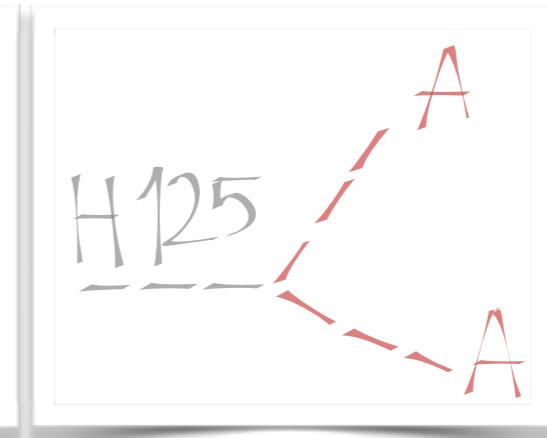
Open modes



$$Y = h125 \quad X = A$$



$$Y = H125 \quad X = h$$



Severe constraints on the **tri-Higgs couplings** from the observed signal strengths: **(rough estimation assuming fermionic SM-like couplings for Y and $m_Y=125$)**

$$R(XX) \equiv \frac{\Gamma(Y \rightarrow XX)}{\Gamma(Y \rightarrow bb)_{\text{tree}}} = \frac{1}{12} \left(\frac{g_{YXX} v}{m_Y m_b} \right)^2 \frac{\beta(m_X)}{\beta^3(m_b)}$$

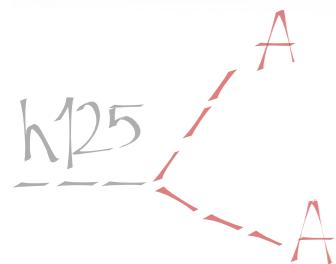
$$\text{BR}(Y \rightarrow XX) \lesssim 0.3 \Leftrightarrow R(XX) \lesssim \frac{5}{6} \Leftrightarrow |g_{YXX}| \lesssim \mathcal{O}(10 \text{ GeV})$$

while $|g_{YXX}| \sim \mathcal{O}(\text{TeV})$ naturally

hl25 scenarios

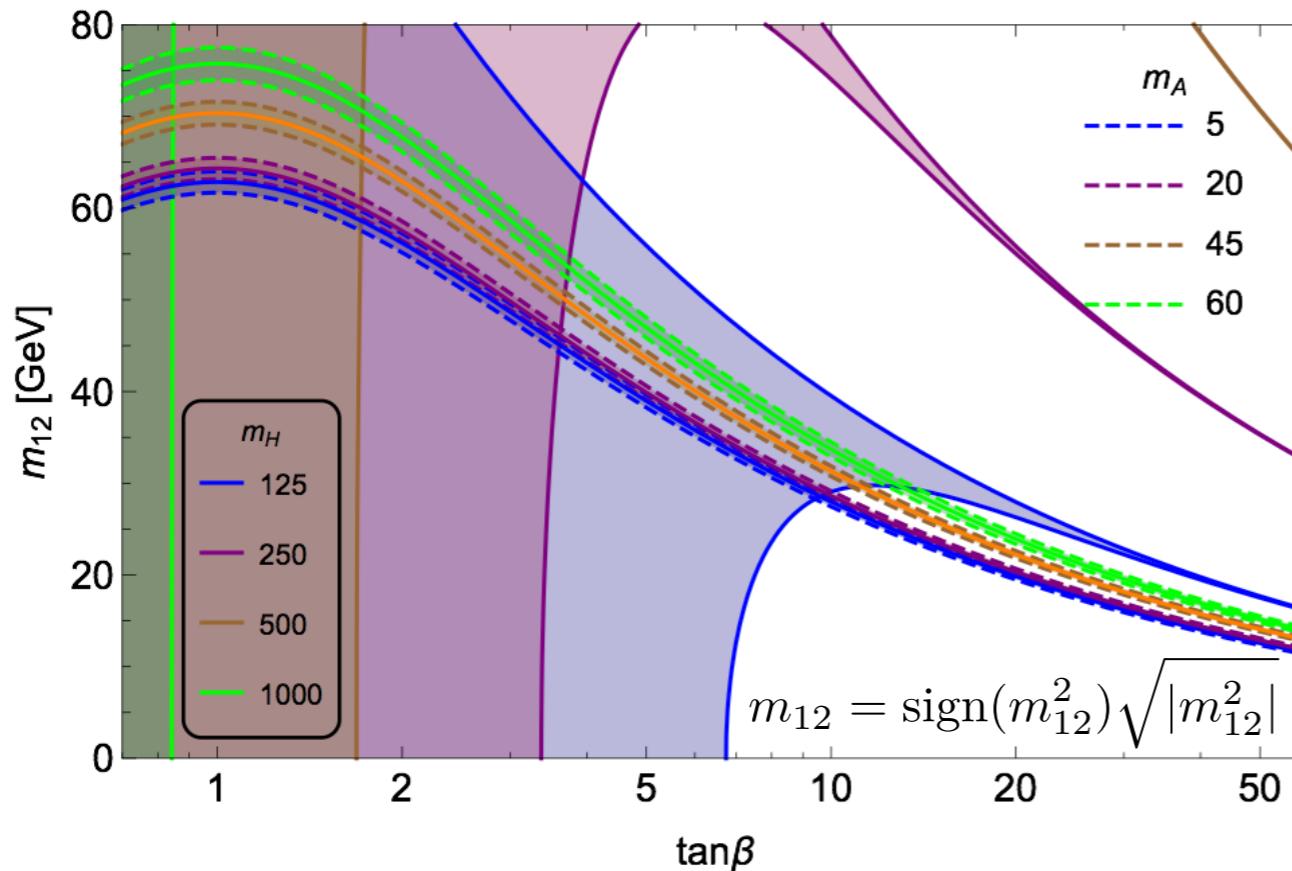
Setup
Results

hAA tri-Higgs coupling



$$g_{hAA} = \frac{1}{2v} \left[(2m_A^2 - m_h^2) \frac{\cos(\alpha - 3\beta)}{\sin 2\beta} + (8m_{12}^2 - \sin 2\beta (2m_A^2 + 3m_h^2)) \frac{\cos(\beta + \alpha)}{\sin^2 2\beta} \right]$$

- In the SM-limit $\sin(\beta - \alpha) = 1$: $g_{hAA} = -\frac{2m_A^2 + m_h^2 - 2m_{12}^2 \sec \beta \csc \beta}{v}$



Solid+dashed lines:
 $\text{BR}(h \rightarrow AA) \leq 0.3$ as a function of m_A

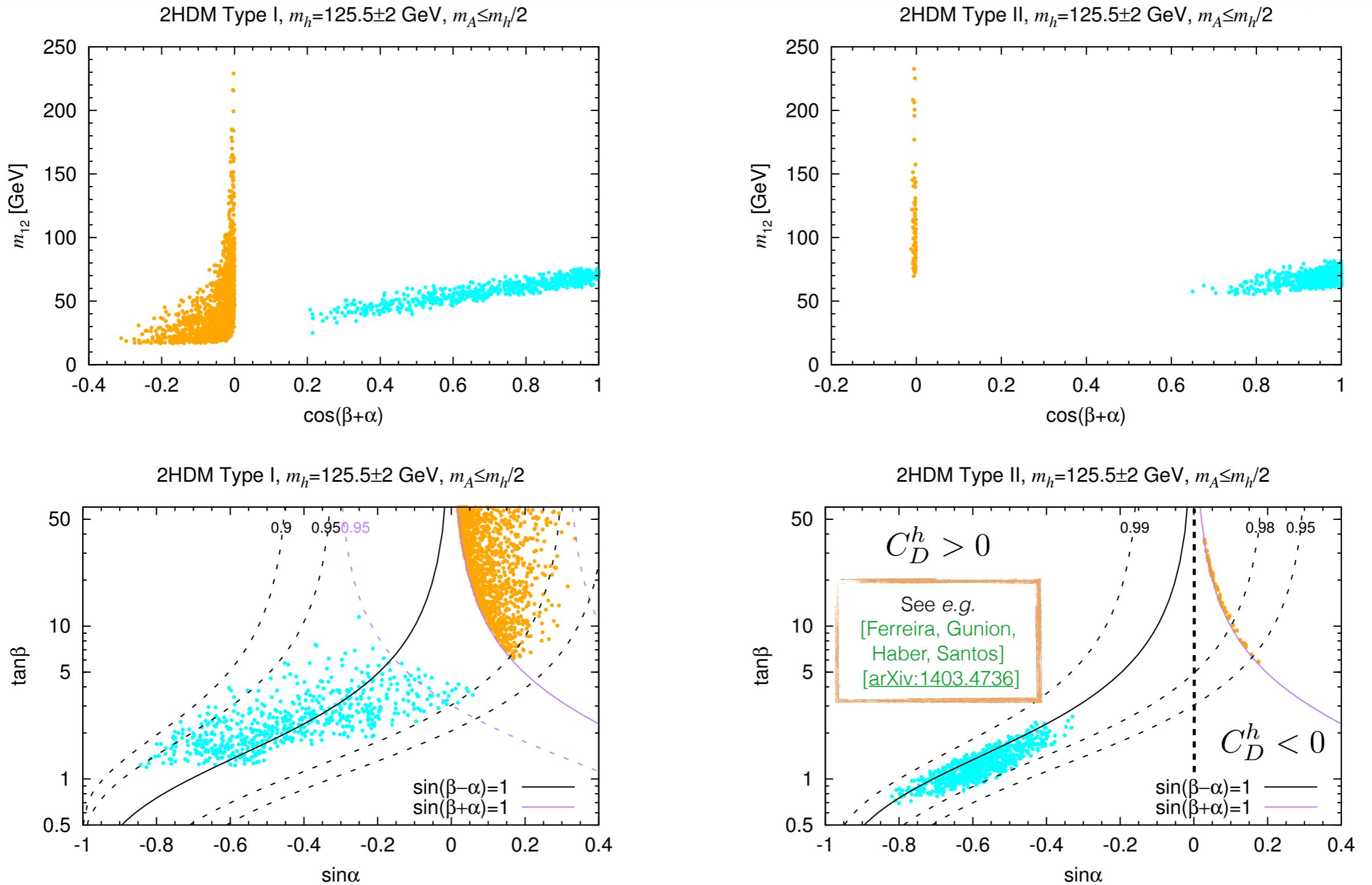
Filled regions :
 allowed by perturbativity
 as a function of m_H

Simultaneous requirement of small
 $h \rightarrow AA$ branching ratio and perturbativity
 \Rightarrow Moderate $\tan \beta$ & small m_{12}

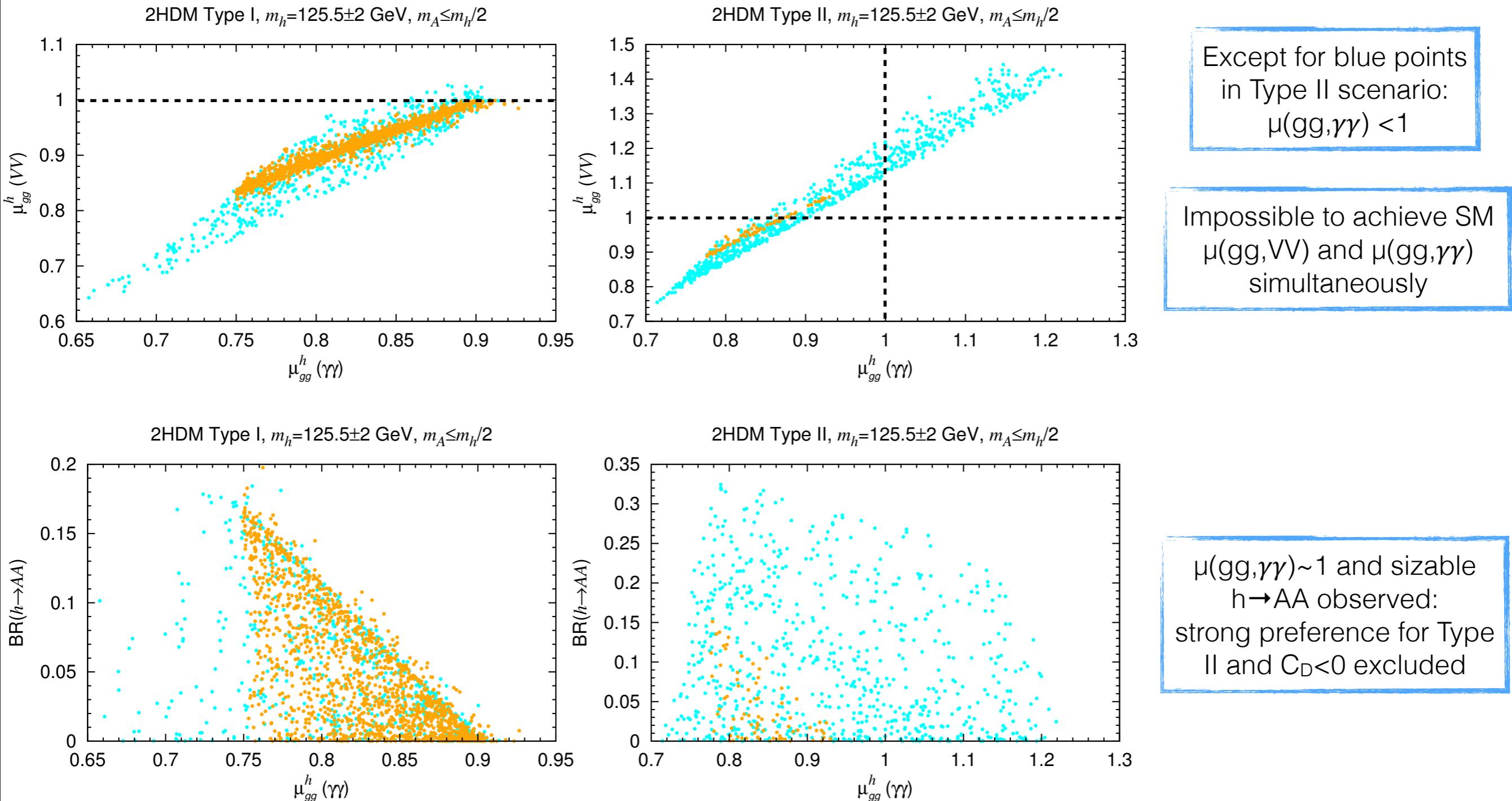
→ Blue points

- Away from the SM-limit, a region with $\sin(\beta + \alpha) \sim 1$, larger m_{12} and $\tan \beta$ also leads to small hAA coupling → Orange points

Parameters overview: h125

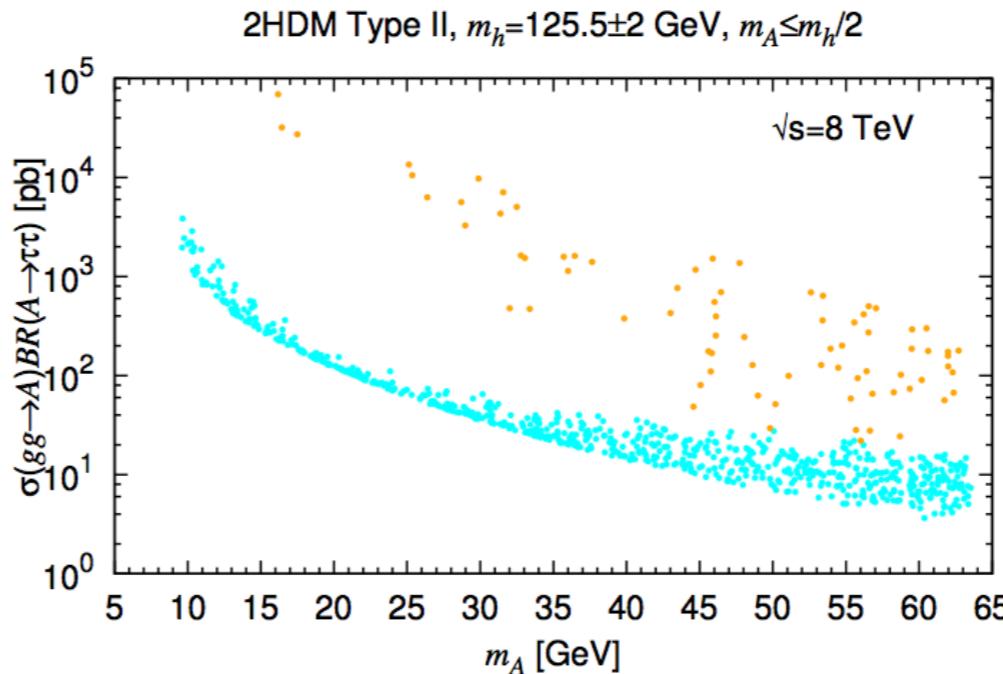
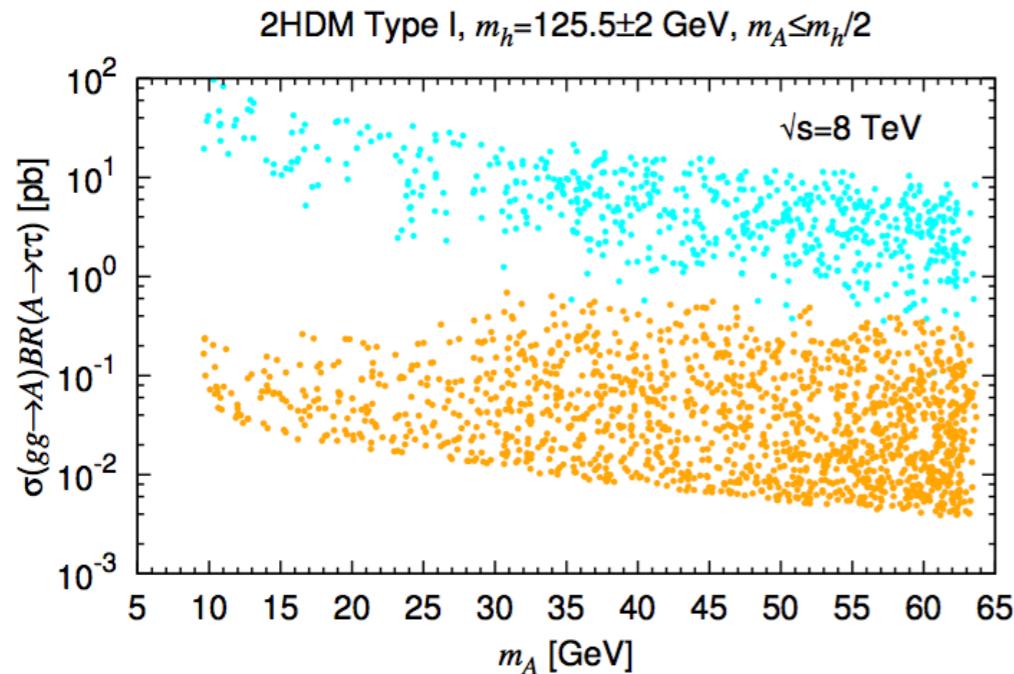


Signal strengths: h125

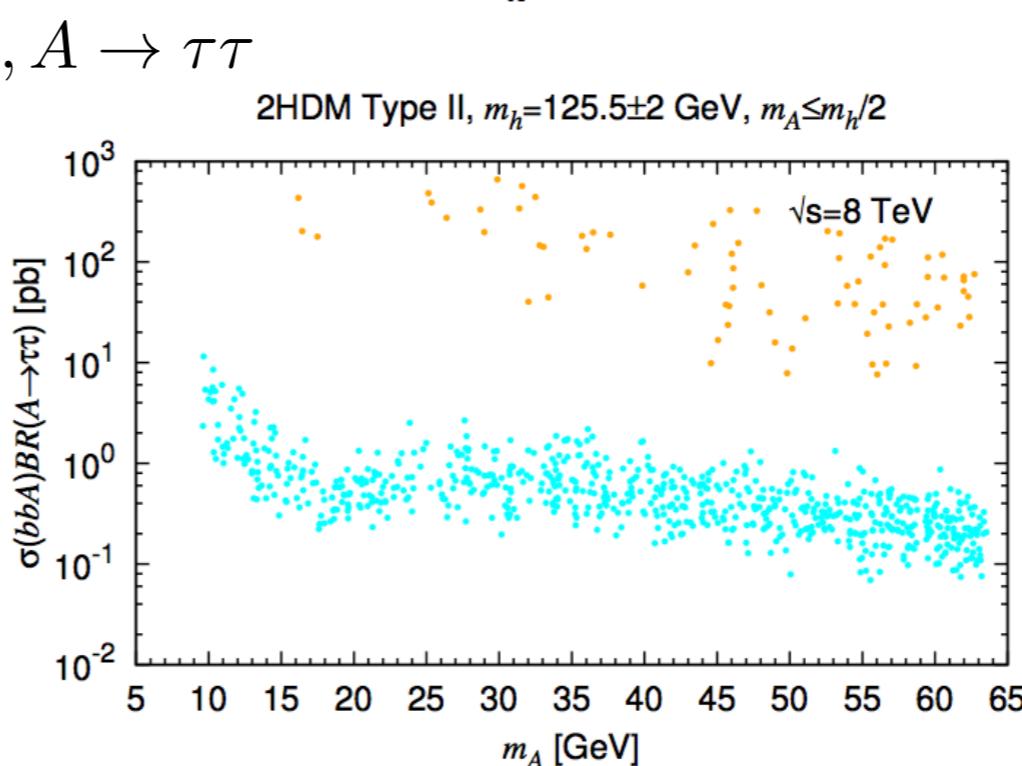
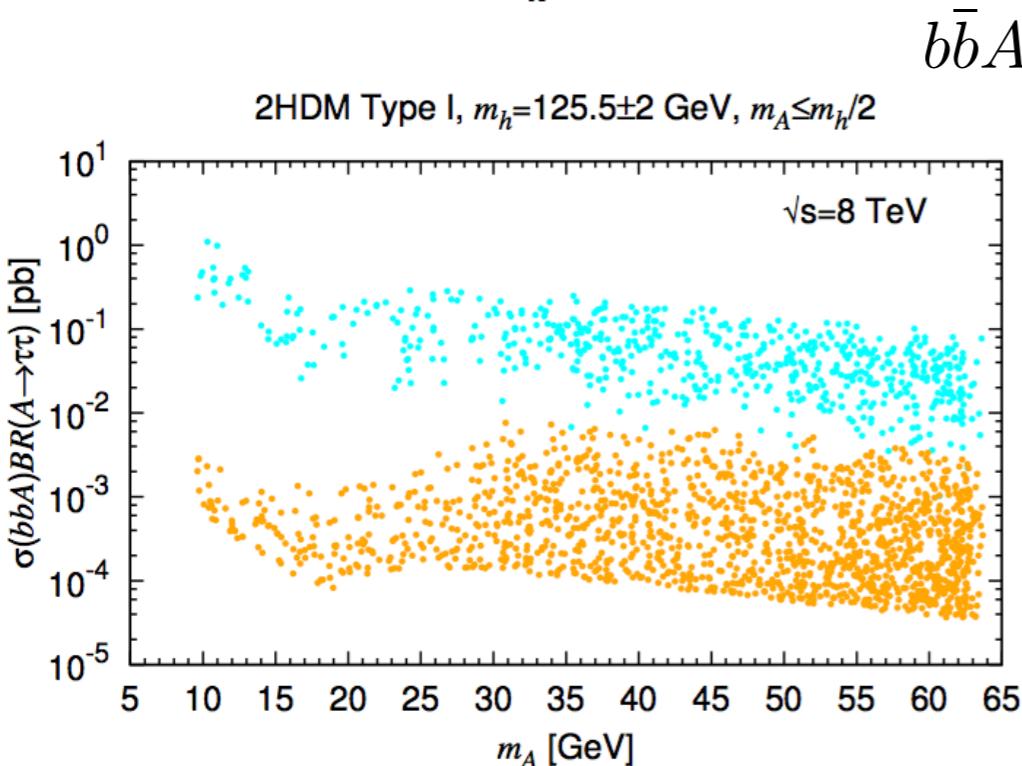


Cross sections: h125

$$gg \rightarrow A \rightarrow \tau\tau$$



Generally very large cross sections over the full mass range



$O(10$ pb) should be observable in the current dataset

$A \rightarrow \mu\mu$: same shape, factor 100 smaller

At 14 TeV: gain of factor ~ 2 in σ

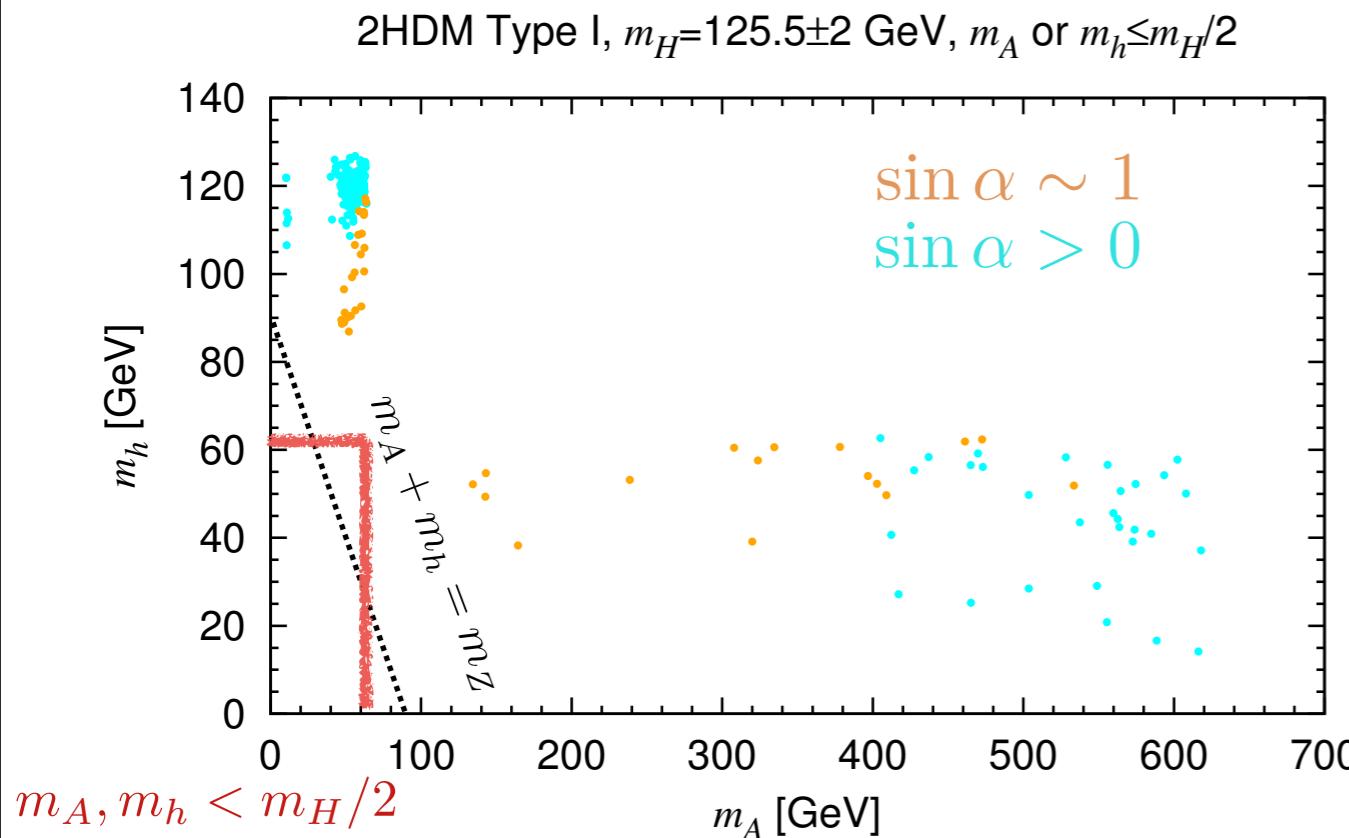
H125 scenarios

Setup

Results

Light states in H125 scenarios

- **Type I:** Both A and h can be lighter than $m_H/2$ but **not simultaneously**



Both the HVV and ZhA vertices are proportional to $\cos(\beta - \alpha)$ which is close to maximal by virtue of H being SM-like

LEP limits on $e^+e^- \rightarrow Z^* \rightarrow Ah$ are evaded via kinematic suppression of the cross section

$m_A < m_H/2$: red points

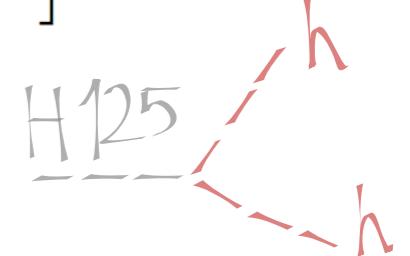
$m_h < m_H/2$: blue points

- **Type II:** B-physics + STU constraints: $m_{H^\pm} \gtrsim 300$ GeV $\Rightarrow m_A \gtrsim 200$ GeV
 \Rightarrow only h can be light
 $m_h < m_H/2$: blue points

HAA and Hhh tri-Higgs couplings

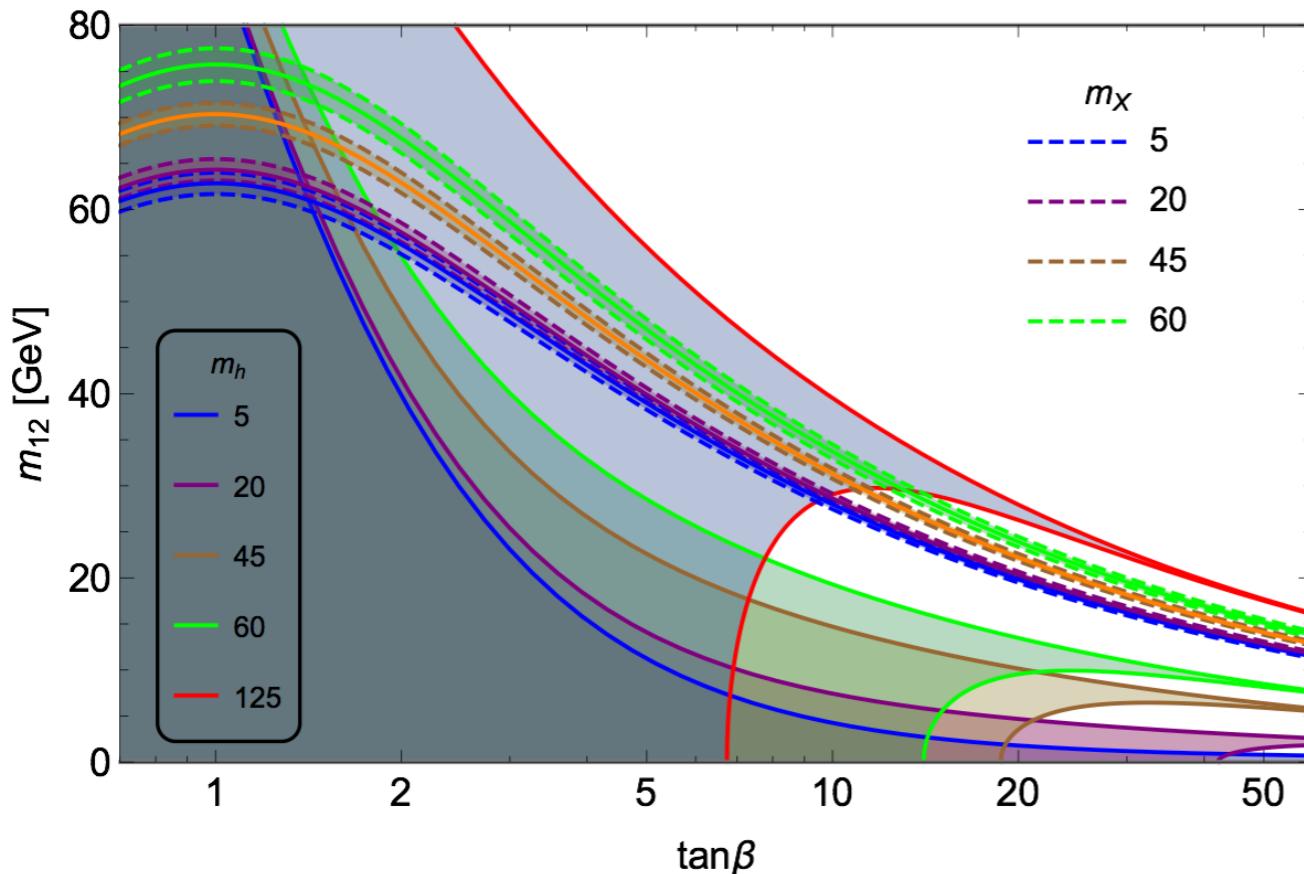


$$g_{HAA} = \frac{1}{2v} \left[(2m_A^2 - m_H^2) \frac{\sin(\alpha - 3\beta)}{\sin 2\beta} + (8m_{12}^2 - \sin 2\beta (2m_A^2 + 3m_H^2)) \frac{\sin(\beta + \alpha)}{\sin^2 2\beta} \right]$$



$$g_{Hhh} = -\frac{1}{v} \cos(\beta - \alpha) \left[\frac{2m_{12}^2}{\sin 2\beta} + \left(2m_h^2 + m_H^2 - \frac{6m_{12}^2}{\sin 2\beta} \right) \frac{\sin 2\alpha}{\sin 2\beta} \right]$$

- In the SM-limit $\cos(\beta - \alpha) = 1$:



$$g_{HXX} = -\frac{2m_X^2 + m_H^2 - 2m_{12}^2 \sec \beta \csc \beta}{v}$$

Perturbativity+small HXX coupling:

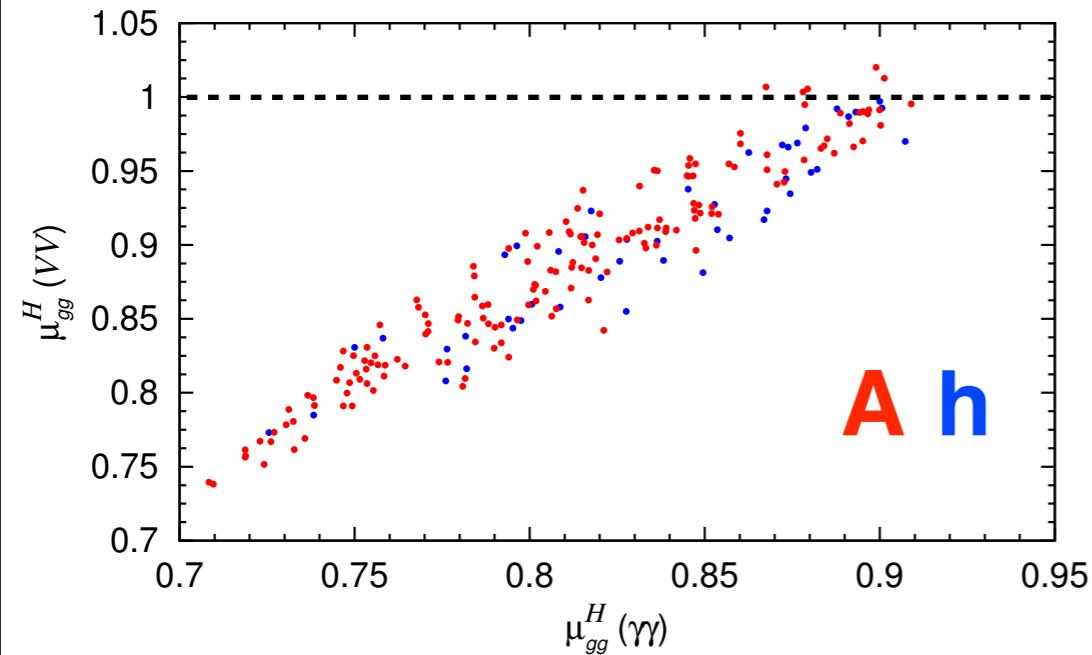
- $m_h < 60$ GeV : $t_\beta \lesssim 2$
- $m_h < 125$ GeV : $t_\beta \lesssim 12$

Moderate m_{12}

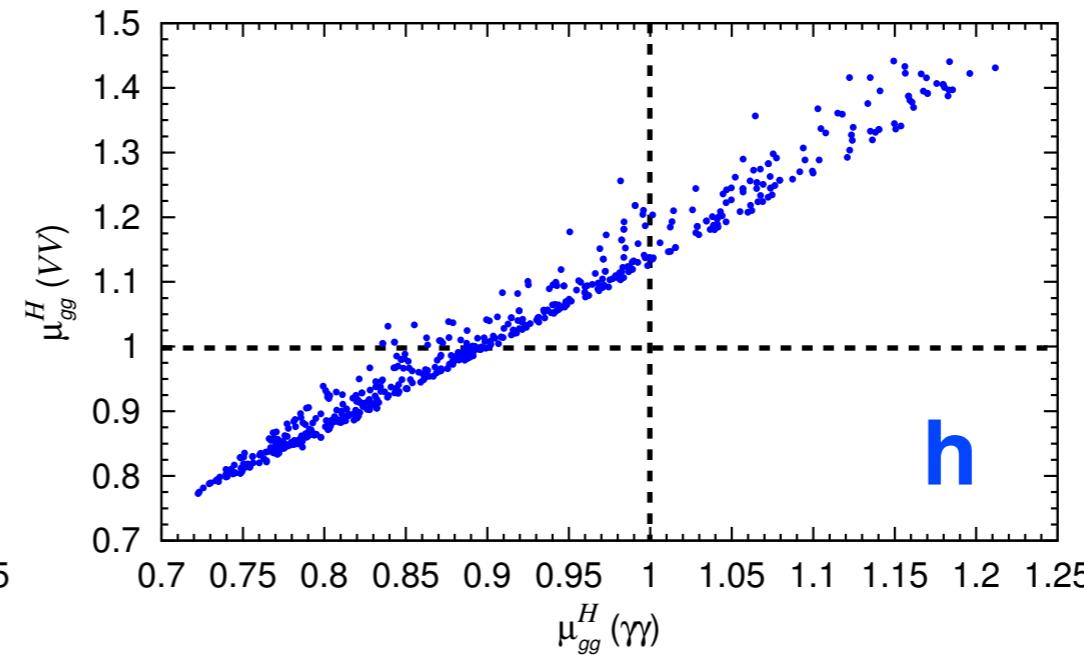
- Away from the SM-limit, in a region with $\cos(\beta + \alpha) \sim 1$ larger t_β can be achieved, at small m_{12}

Signal strengths: H125

2HDM Type I, $m_H=125.5\pm2$ GeV, m_A or $m_h \leq m_H/2$

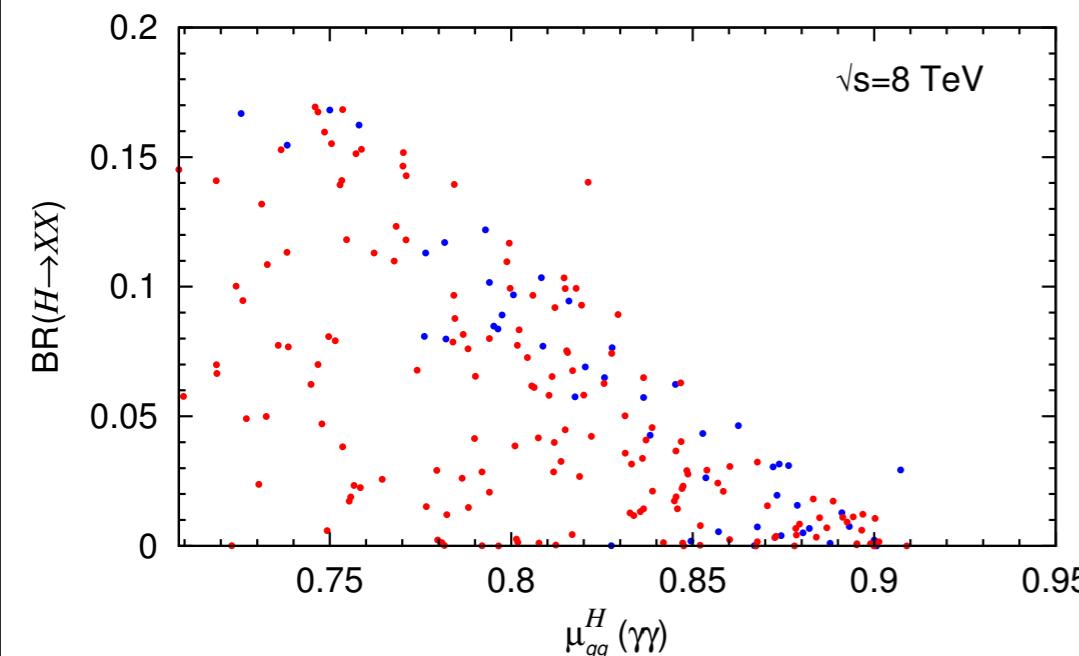


2HDM Type II, $m_H=125.5\pm2$ GeV, $m_h \leq m_H/2$

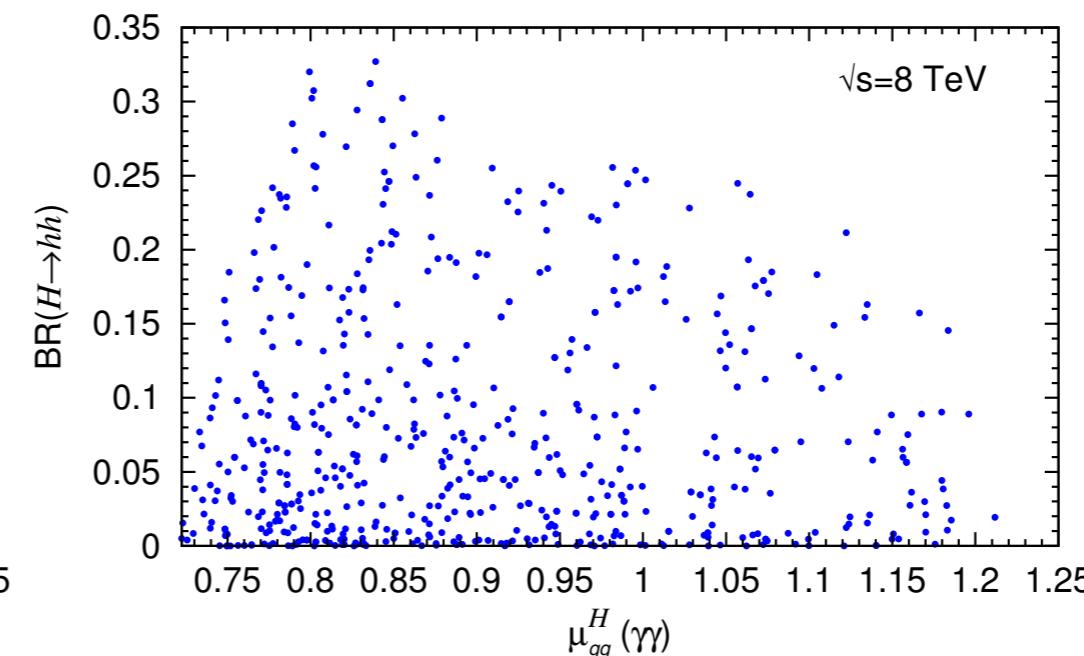


Impossible to achieve SM
 $\mu(gg,VV), \mu(gg,\gamma\gamma)$
 simultaneously

2HDM Type I, $m_H=125.5\pm2$ GeV, m_A or $m_h \leq m_H/2$



2HDM Type II, $m_H=125.5\pm2$ GeV, $m_h \leq m_H/2$

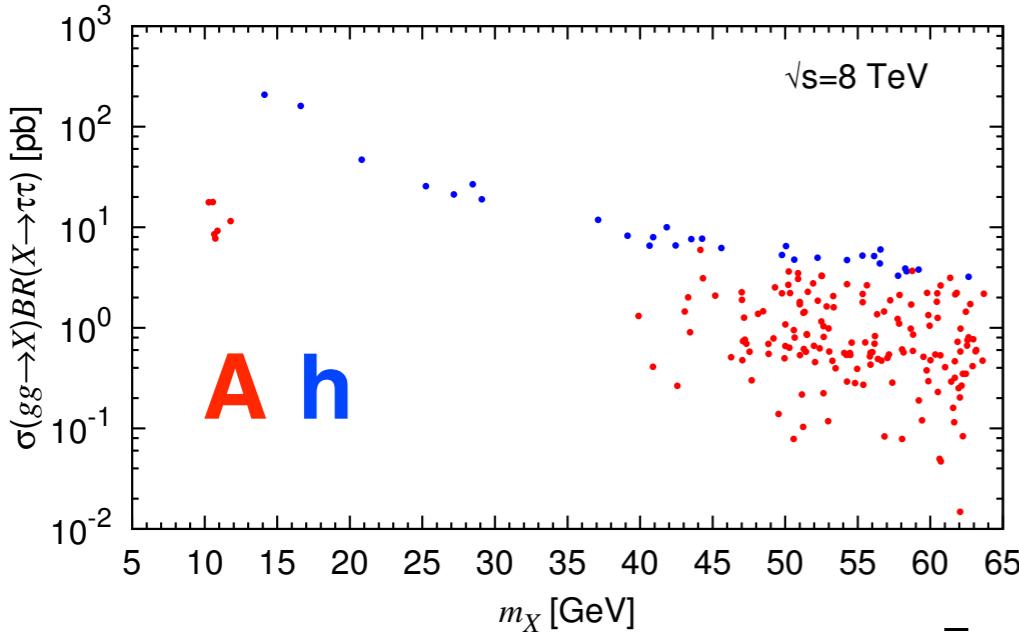


$\mu(gg,\gamma\gamma) \sim 1$ and
 sizable $H \rightarrow hh$
 observed: strong
 preference for
 Type II

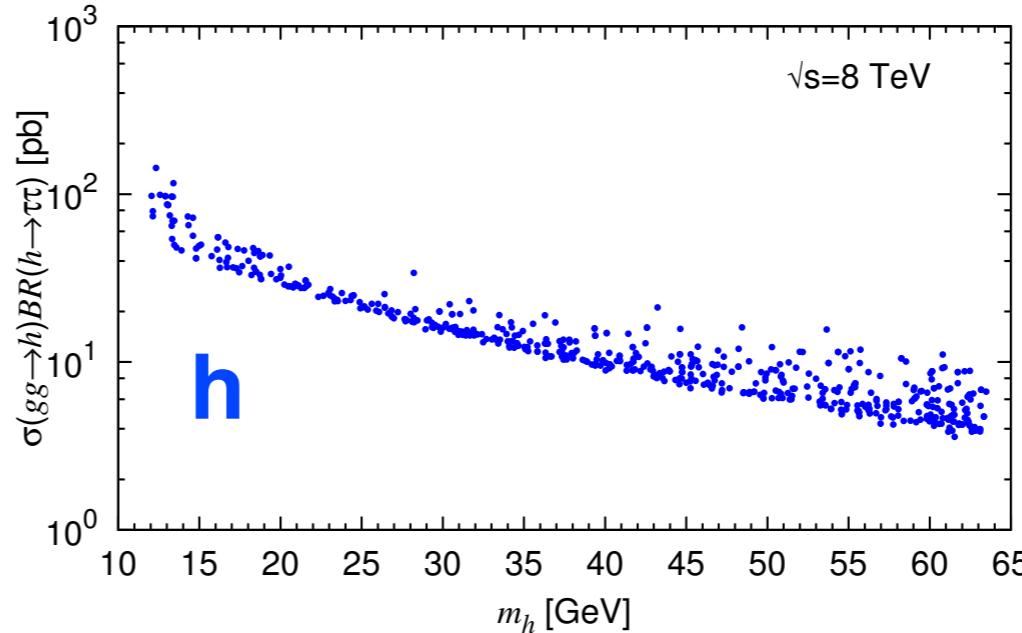
Cross sections: H125

$$gg \rightarrow A \rightarrow \tau\tau$$

2HDM Type I, $m_H=125.5\pm2$ GeV, m_A or $m_h \leq m_H/2$



2HDM Type II, $m_H=125.5\pm2$ GeV, $m_h \leq m_H/2$

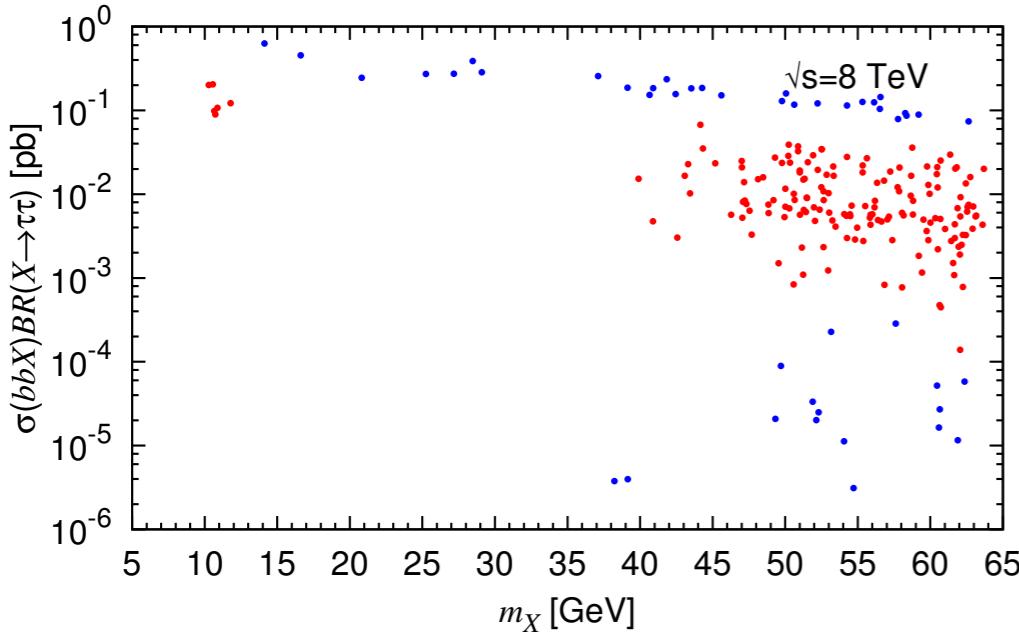


$\sigma(ggh, h \rightarrow \tau\tau) \approx 10\text{pb}$
over the full mass range

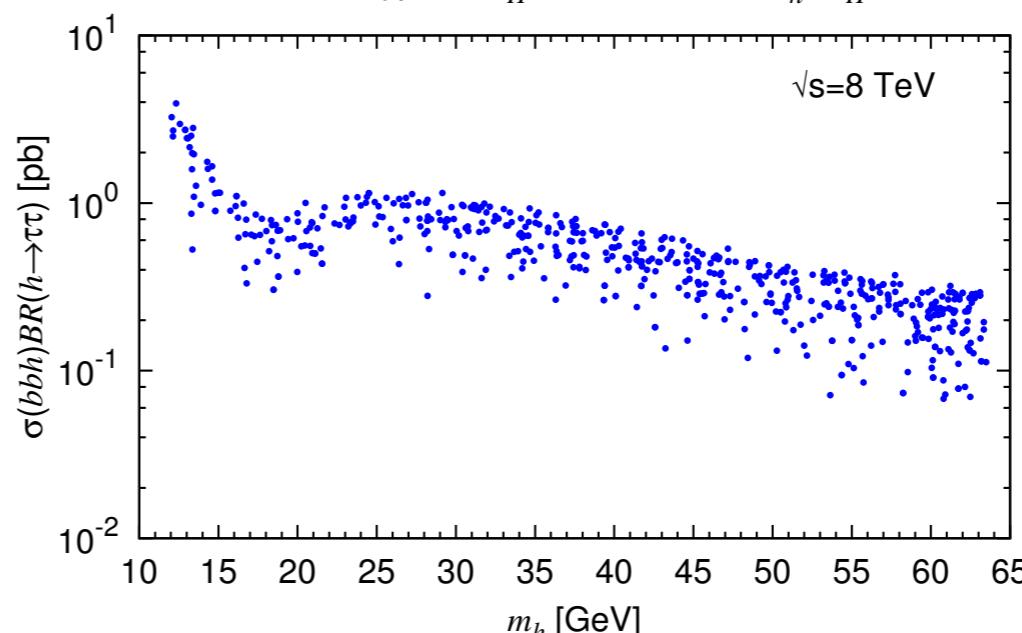
Only very light pseudo-scalars have
 $\sigma(ggA, A \rightarrow \tau\tau) \approx 10\text{pb}$

$$b\bar{b}A, A \rightarrow \tau\tau$$

2HDM Type I, $m_H=125.5\pm2$ GeV, m_A or $m_h \leq m_H/2$



2HDM Type II, $m_H=125.5\pm2$ GeV, $m_h \leq m_H/2$



Type II: $t_\beta \lesssim 2$, no enhancement of the hbb cross-section

Conclusions

Conclusions

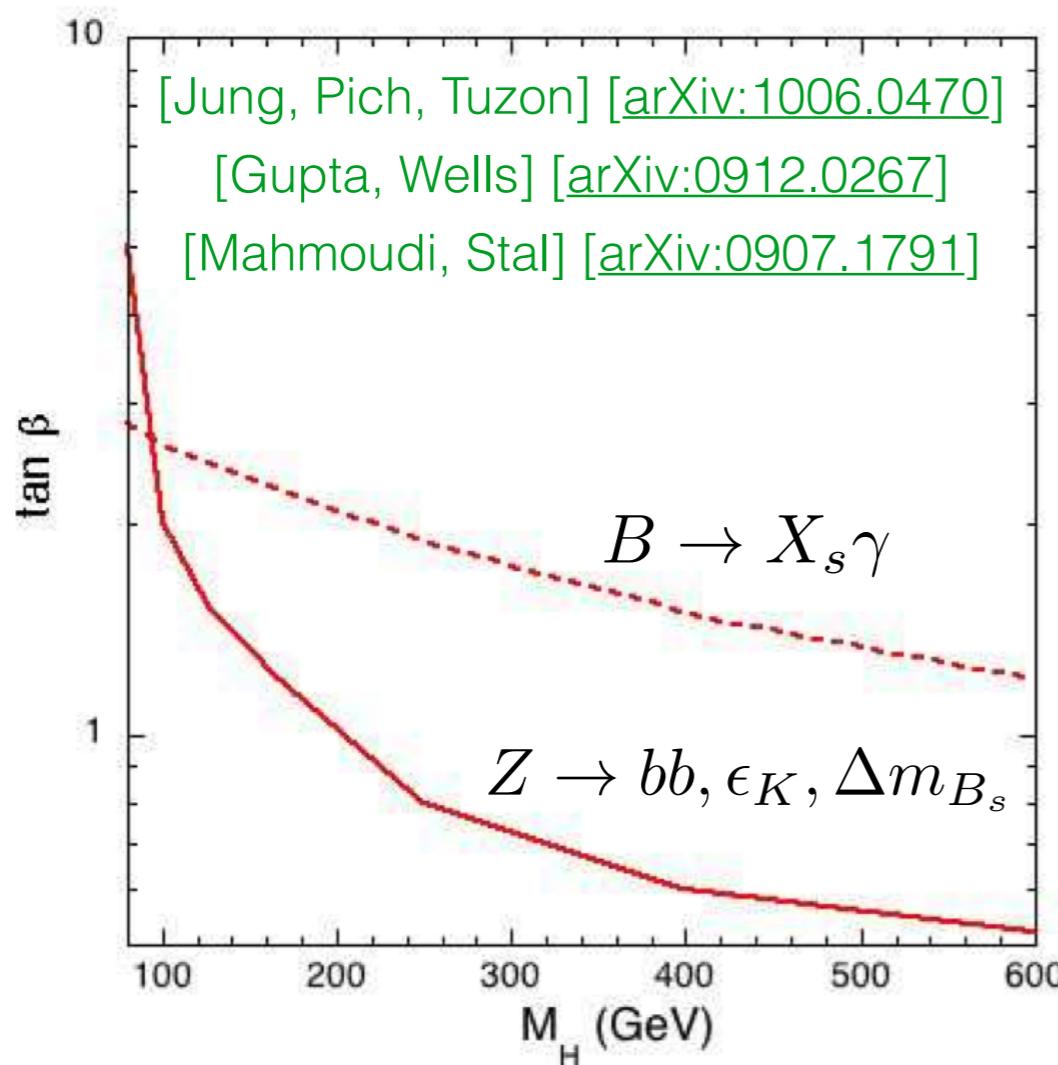
- In the context of the 2HDM of Type I and Type II, the presence of **light Higgs states** is still a possibility in both the **h125** and **H125 scenarios**
- Precise measurements of signal strengths during LHC Run II could largely test these scenarios
- Should the $\gamma\gamma$ and VV rates converge to the SM value within **~10% or better**, these scenarios would be **excluded**
- **Large cross sections** in the gluon fusion and bb associated production modes are generally possible
- Looking for these light states in the **$\tau\tau$ and $\mu\mu$ channels** in the existing **LHC Run I dataset** could already be a **crucial test** of these scenarios

Backup

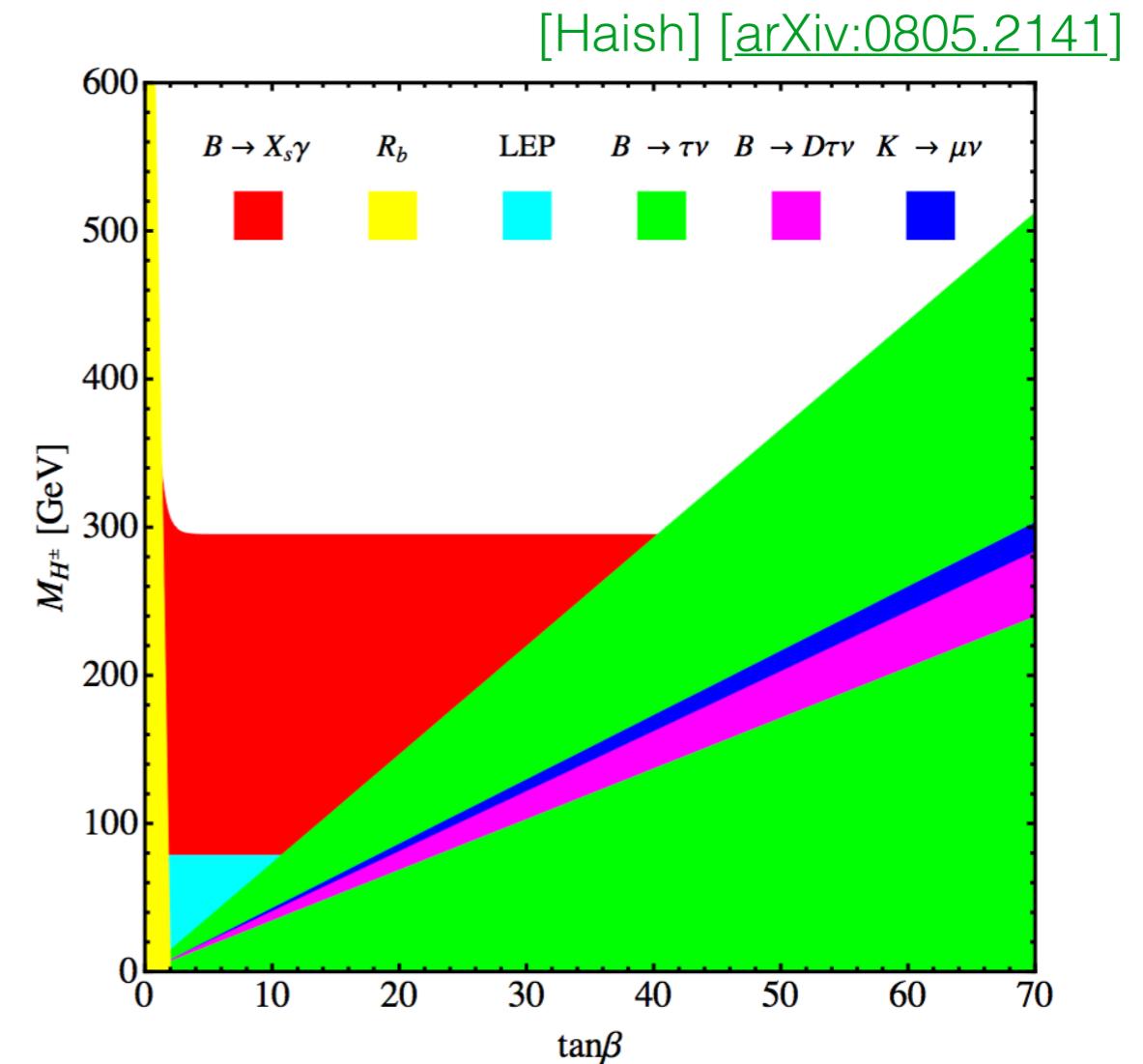
Constraints

B-physics constraints

Type I



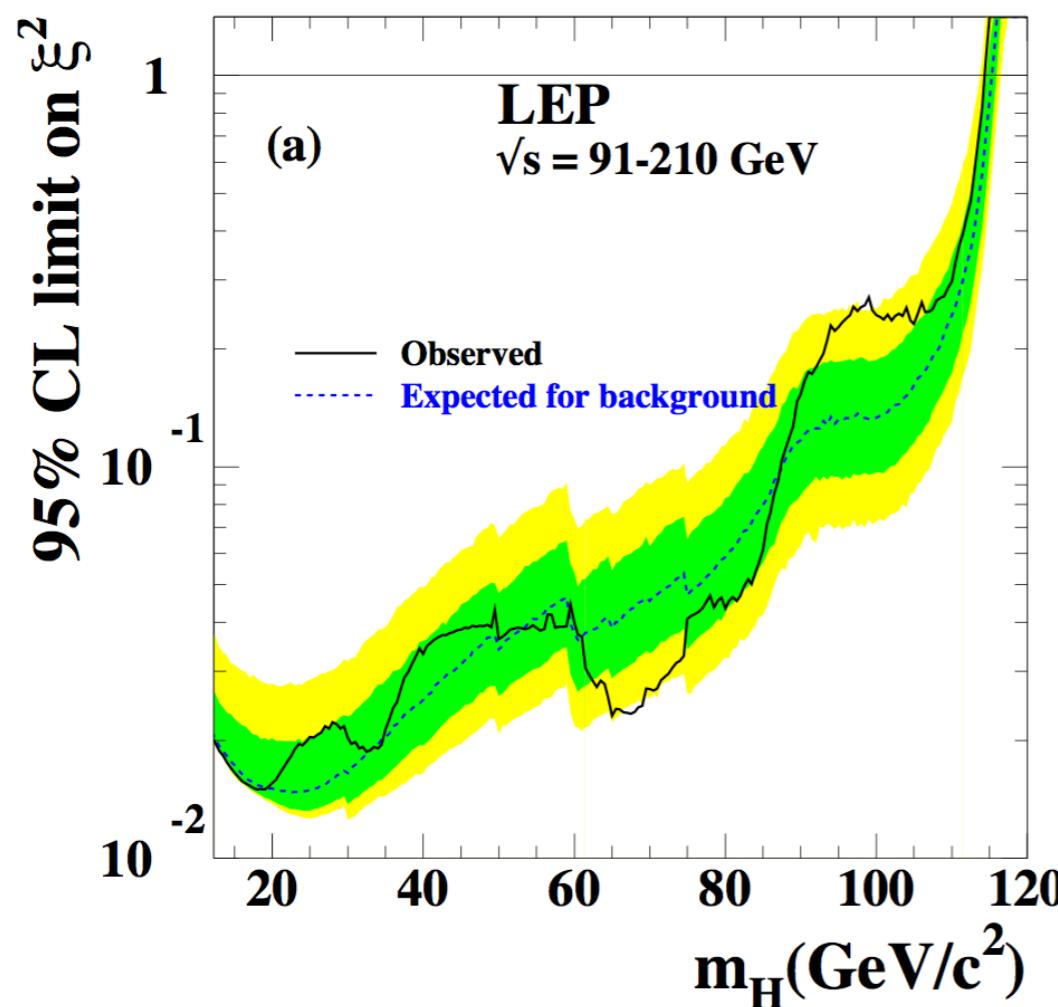
Type II



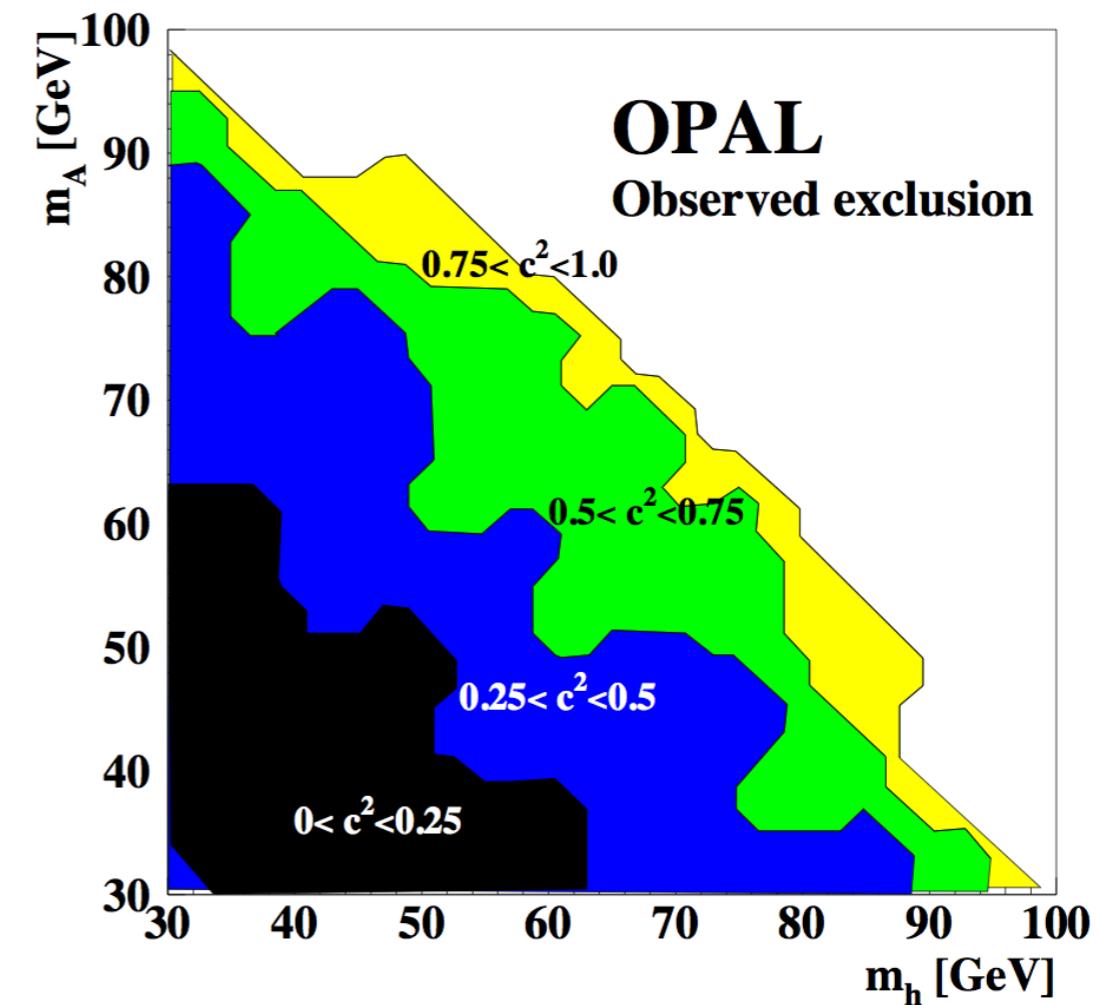
Lower bound on charged Higgs mass

LEP constraints

$e^+e^- \rightarrow Zh$

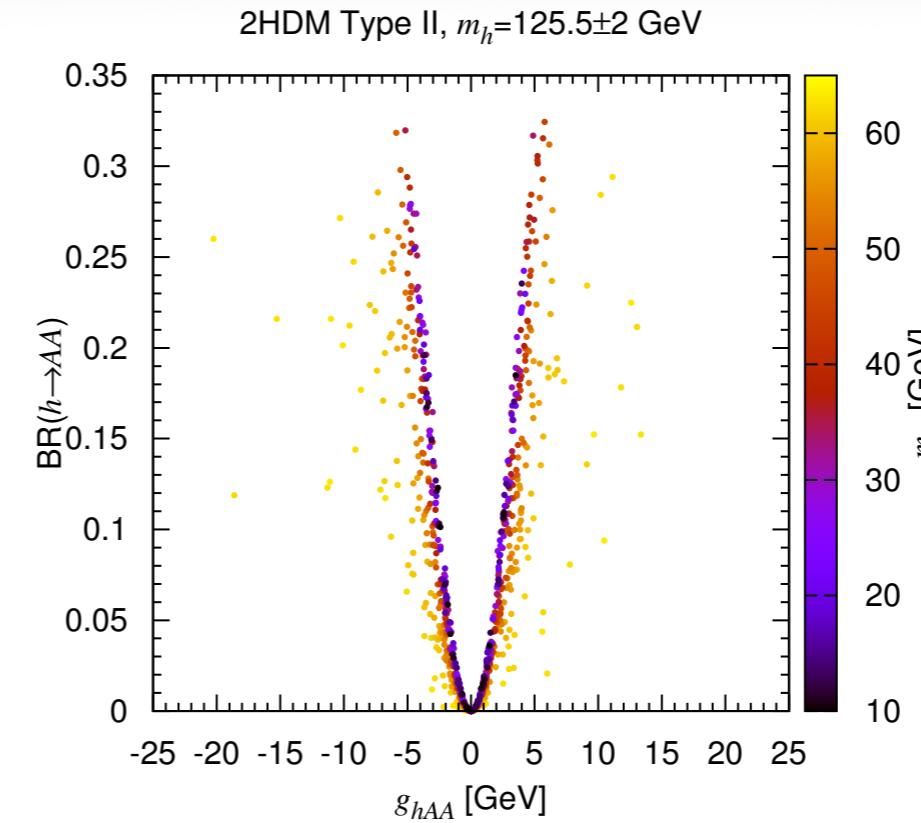
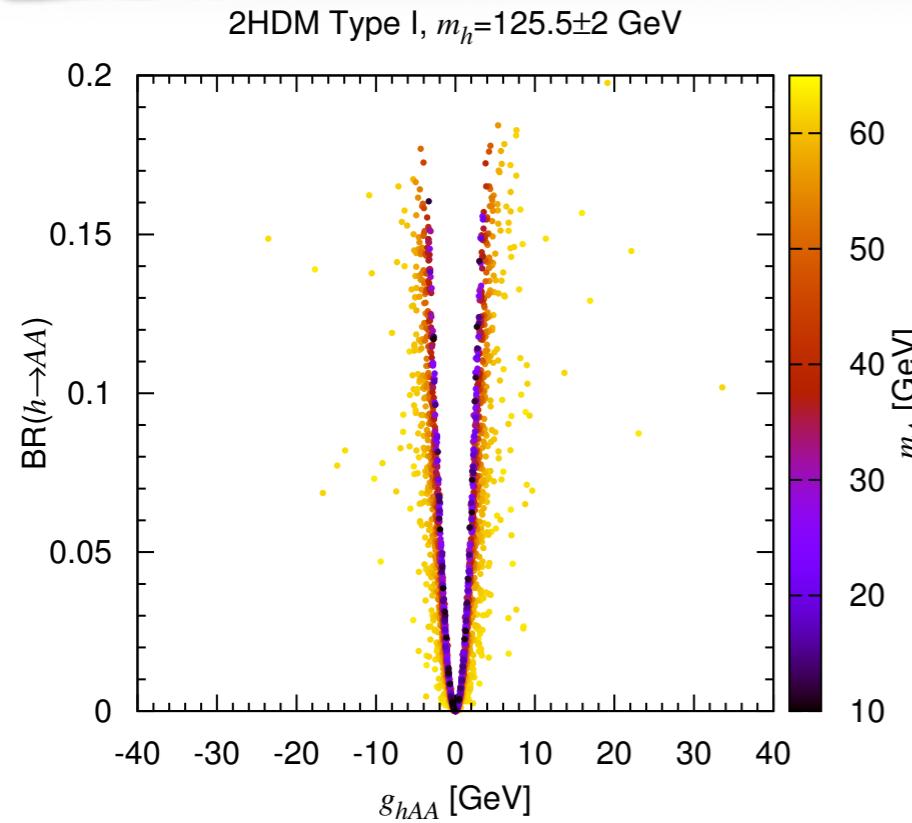


$e^+e^- \rightarrow Z^* \rightarrow Ah$

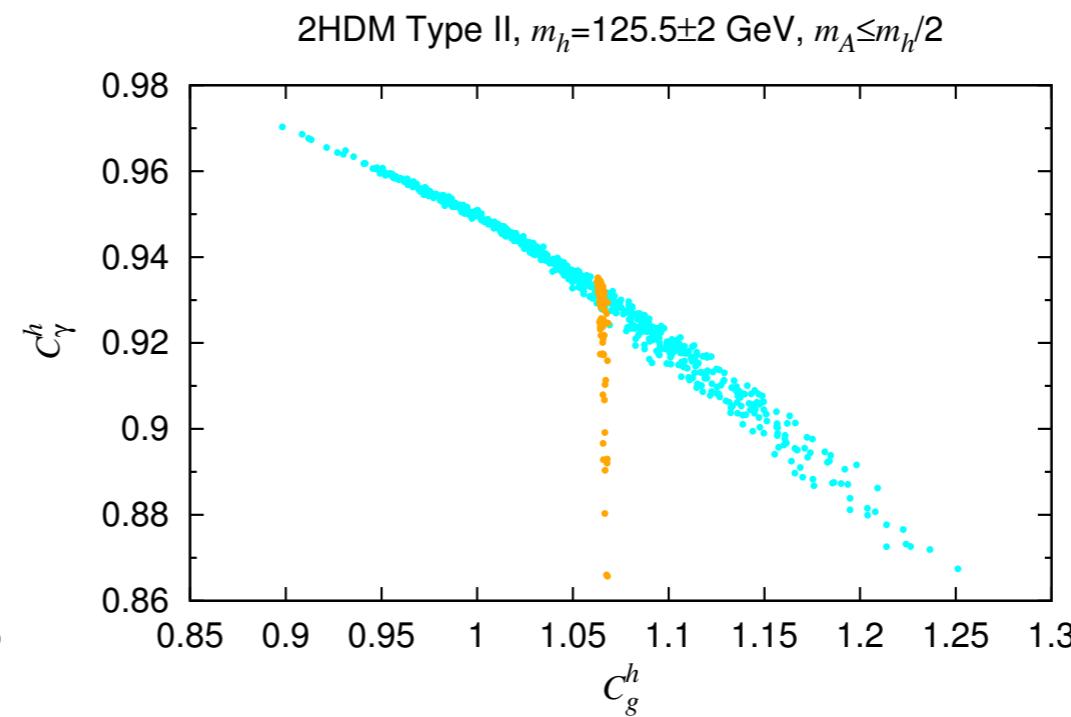
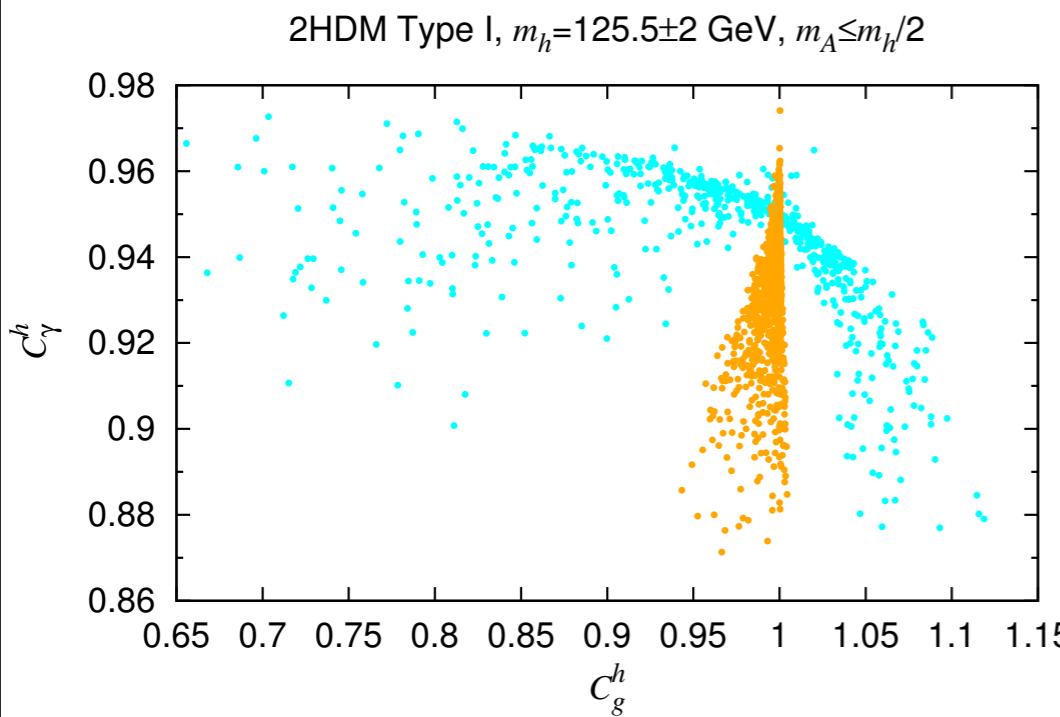


h125 scenarios

Couplings



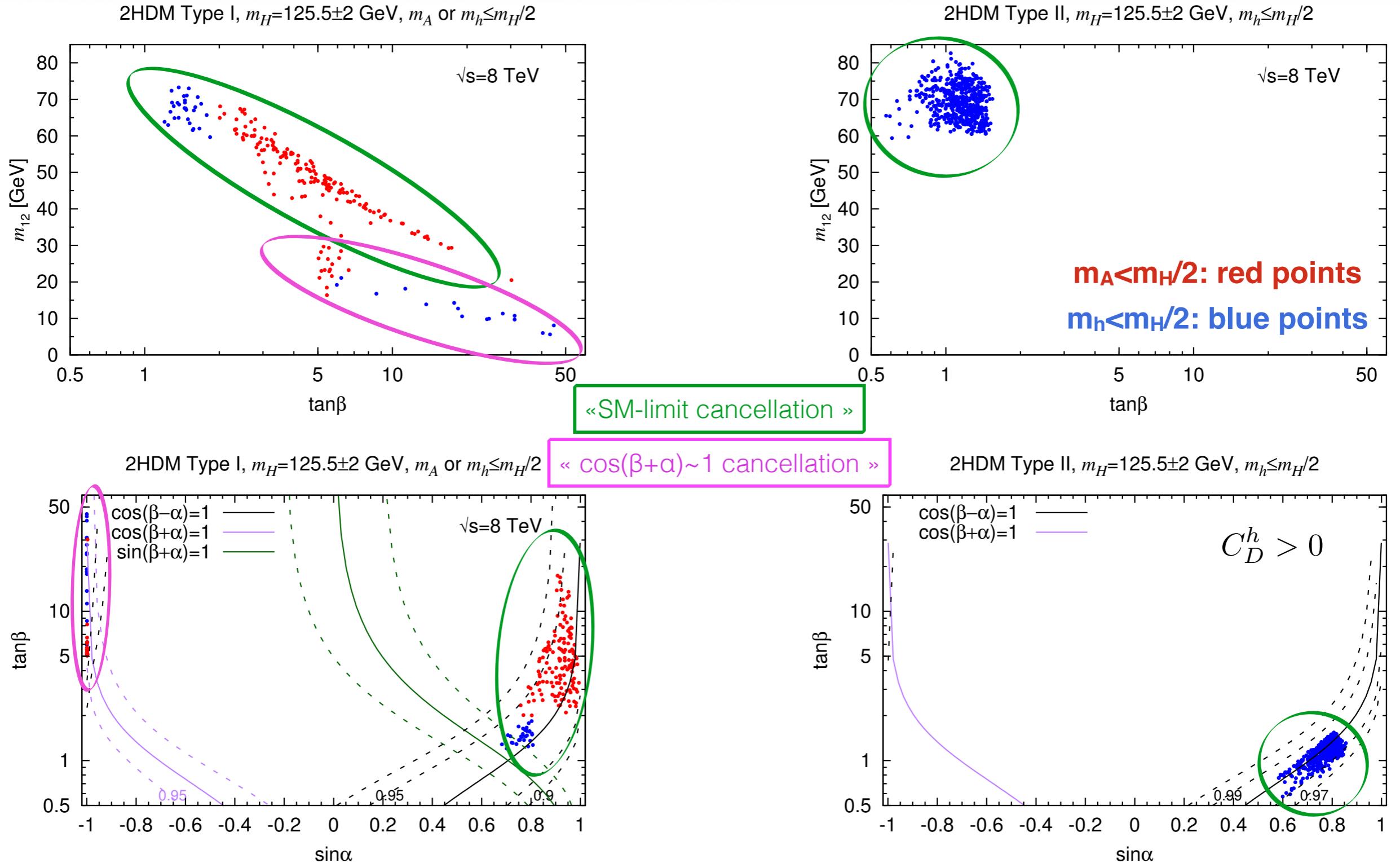
Strong constraint
from undetected
branching ratio



Suppressed
 C_g coupling:
charged Higgs
contribution

H125 scenarios

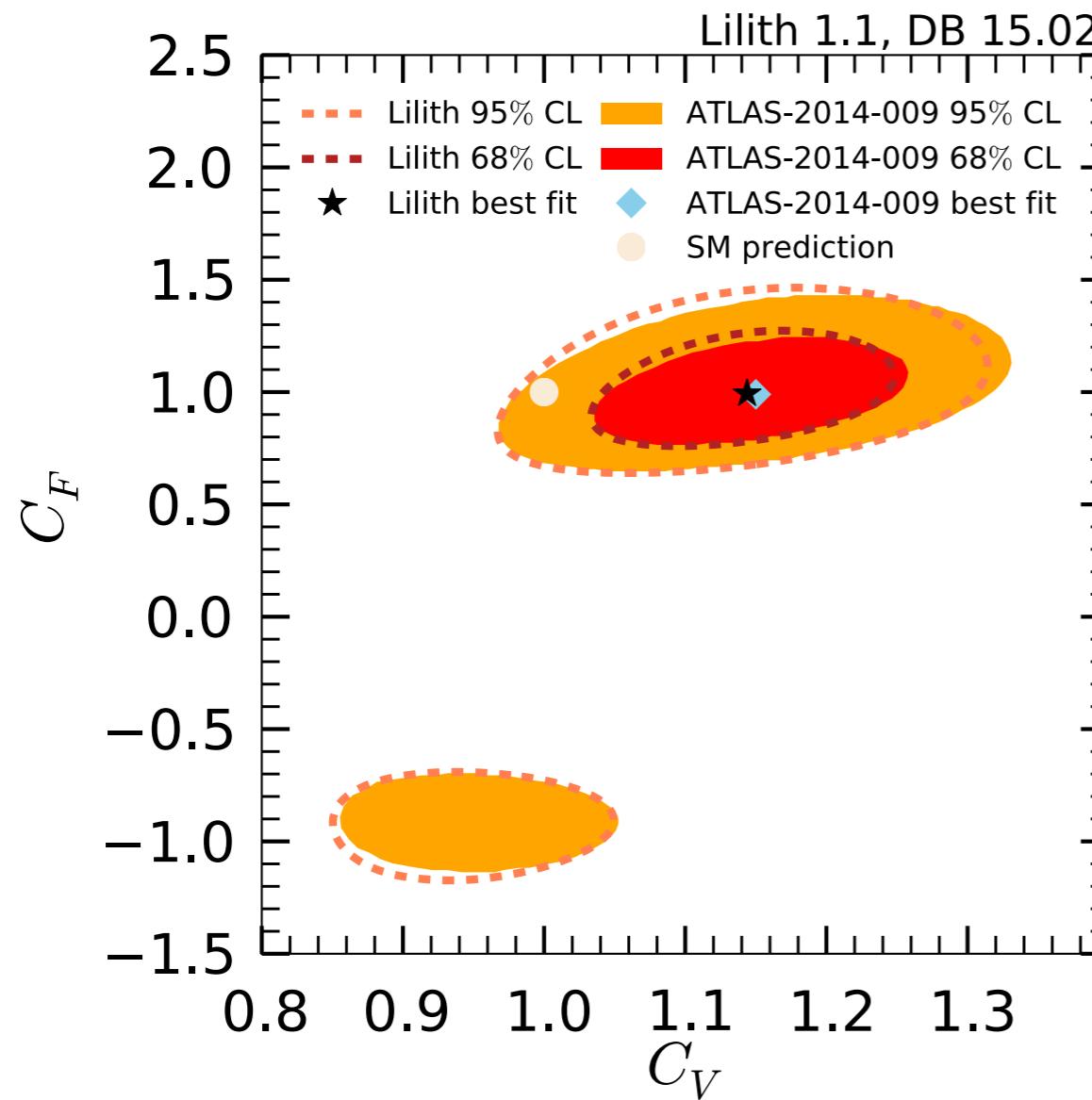
Parameters overview: H125



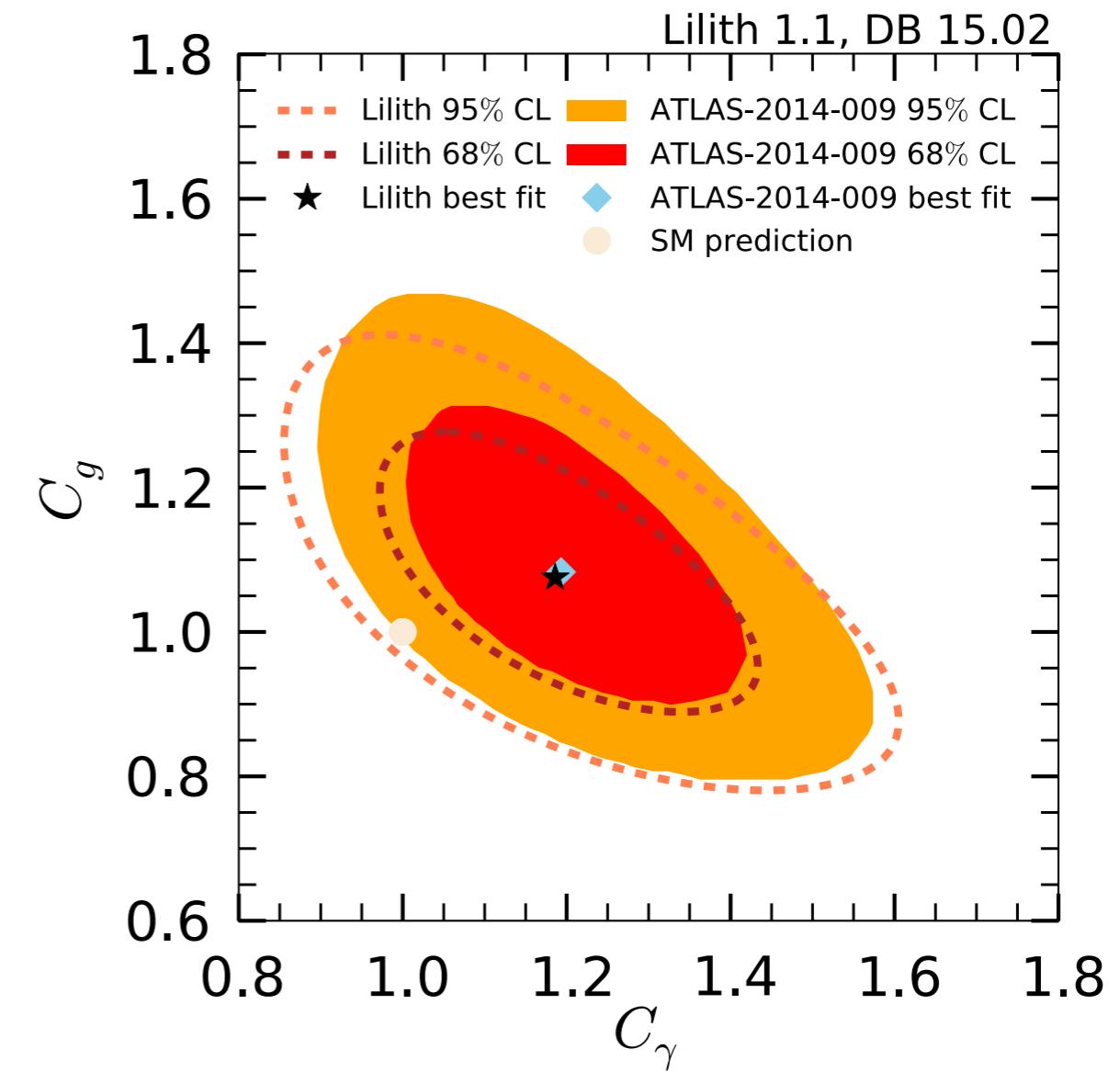
Lilith

Validation of the Lilith likelihood against ATLAS results

- Trying to reproduce the official ATLAS and CMS coupling fits (profile likelihood ratio to derive the confidence intervals)



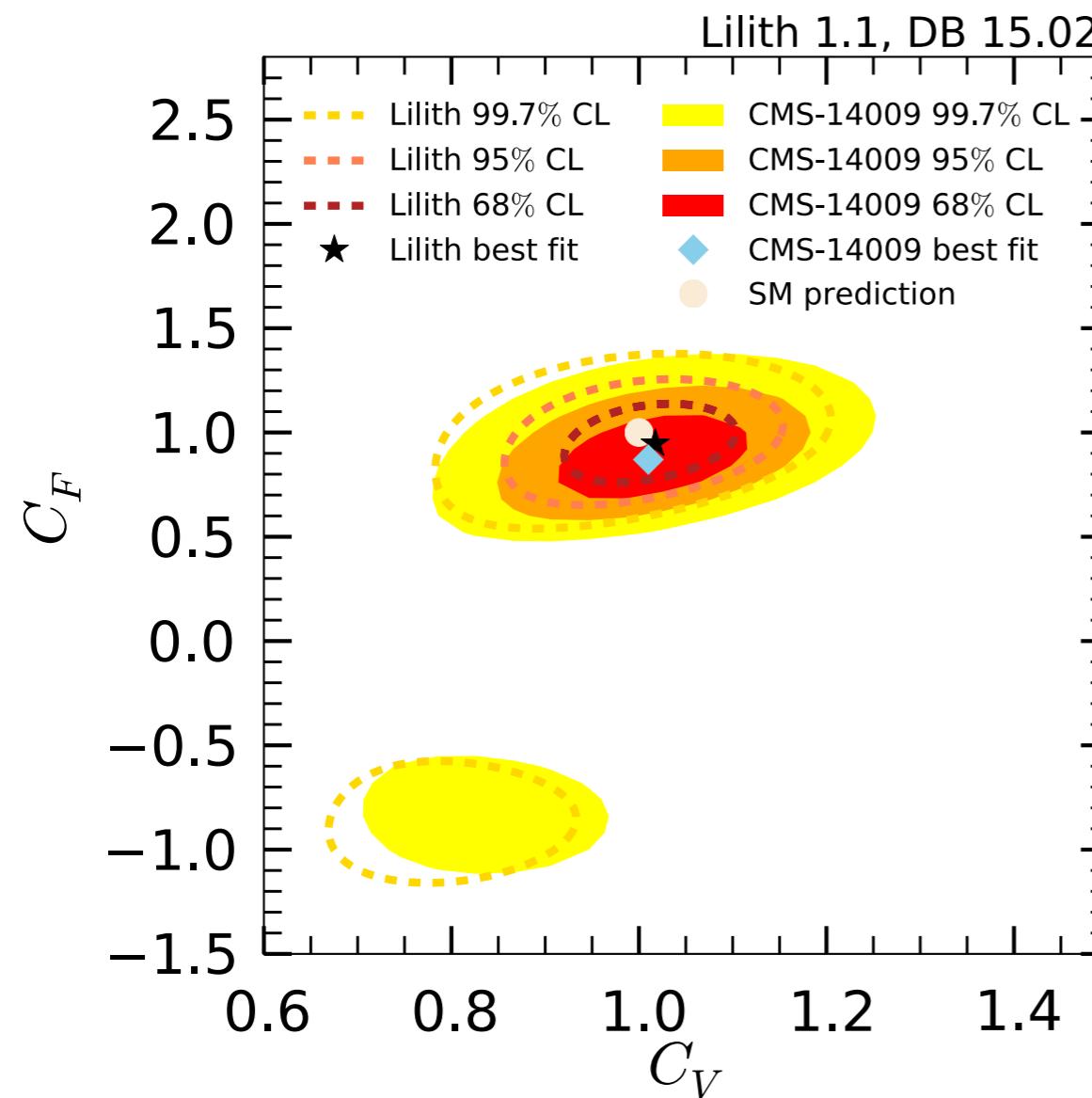
(C_V, C_F) benchmark scenario



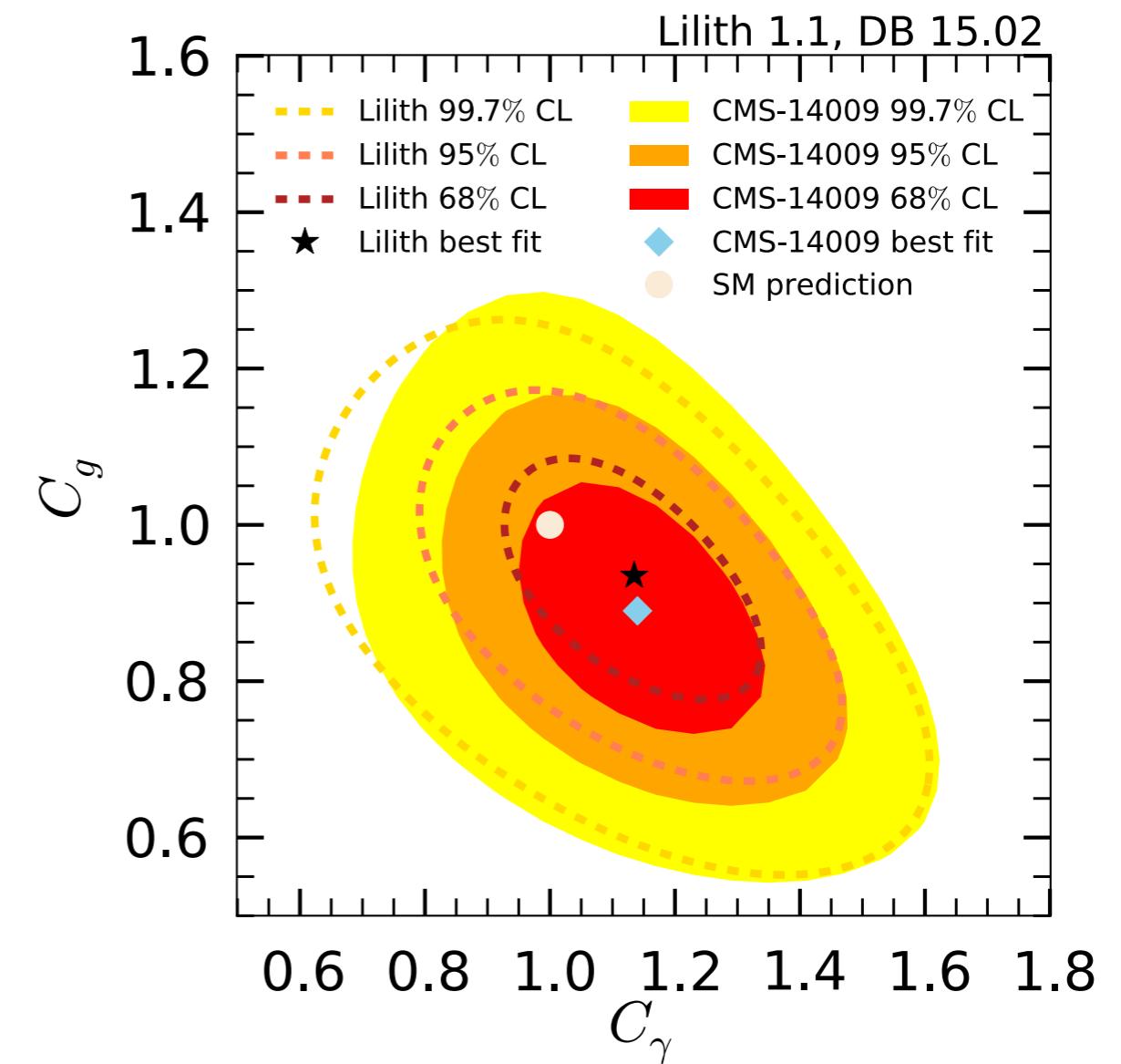
(C_γ, C_g) benchmark scenario

Validation of the Lilith likelihood against CMS results

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(C_V, C_F) benchmark scenario



(C_γ, C_g) benchmark scenario