

Cosmology V: CMB Polarization and other themes

Wessel Valkenburg

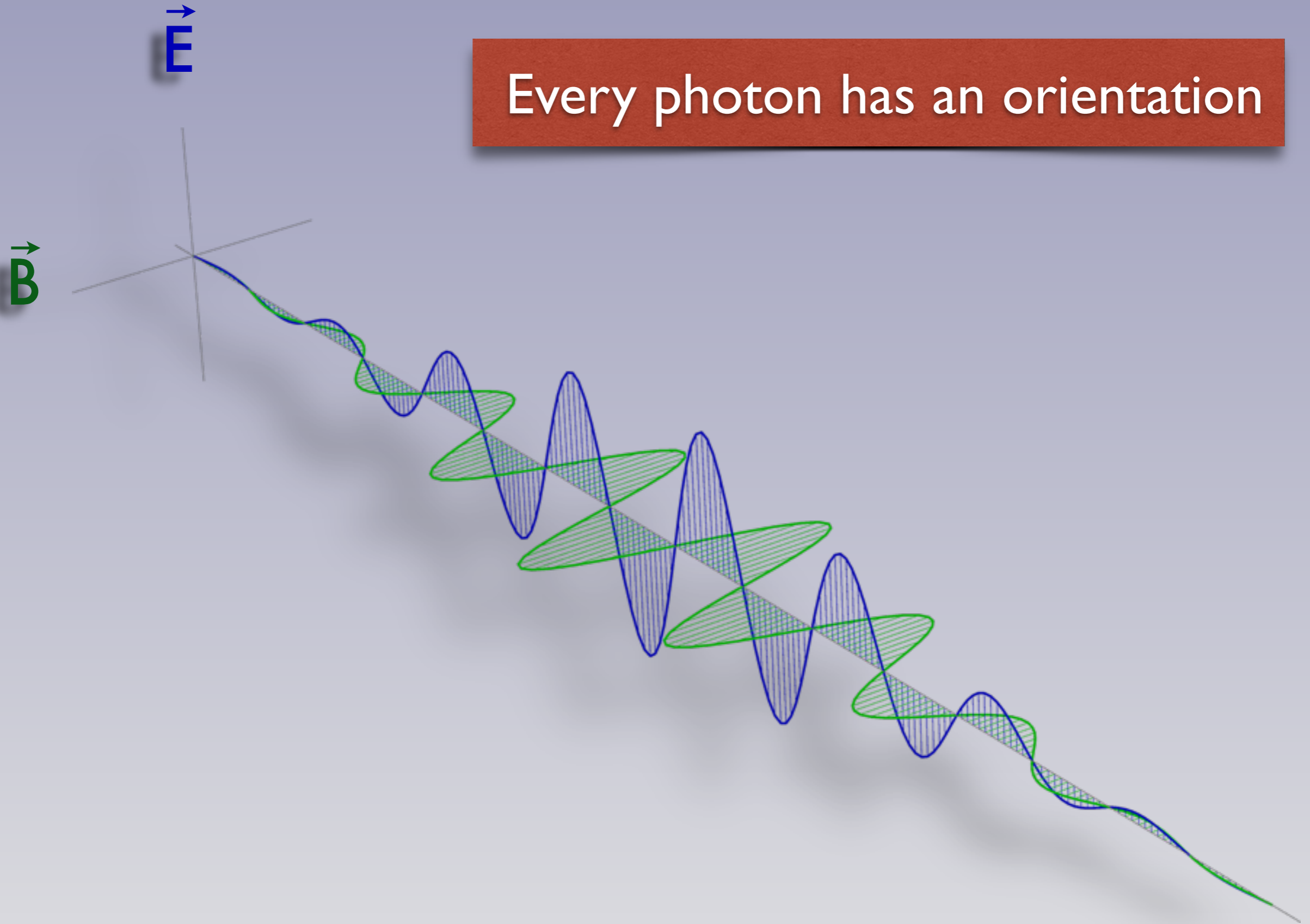


Universiteit Leiden
Pays-Bas



Outline

Every photon has an orientation



Spherical harmonics

$$T(\phi, \theta) = \sum_{l=0}^{\infty} \sum_{m=-l}^l a_{lm} Y_{lm}(\phi, \theta)$$

$$a_{lm} = \int_{\Omega} d\Omega T(\phi, \theta) Y_{lm}^*$$

$$D_l^{TT} \equiv \frac{1}{2l+1} \sum_{m=-l}^l |a_{lm}|^2$$

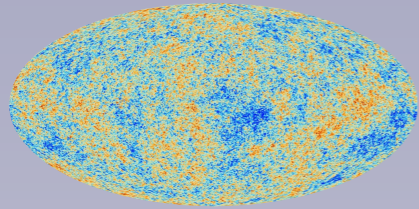
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Spherical harmonics



+

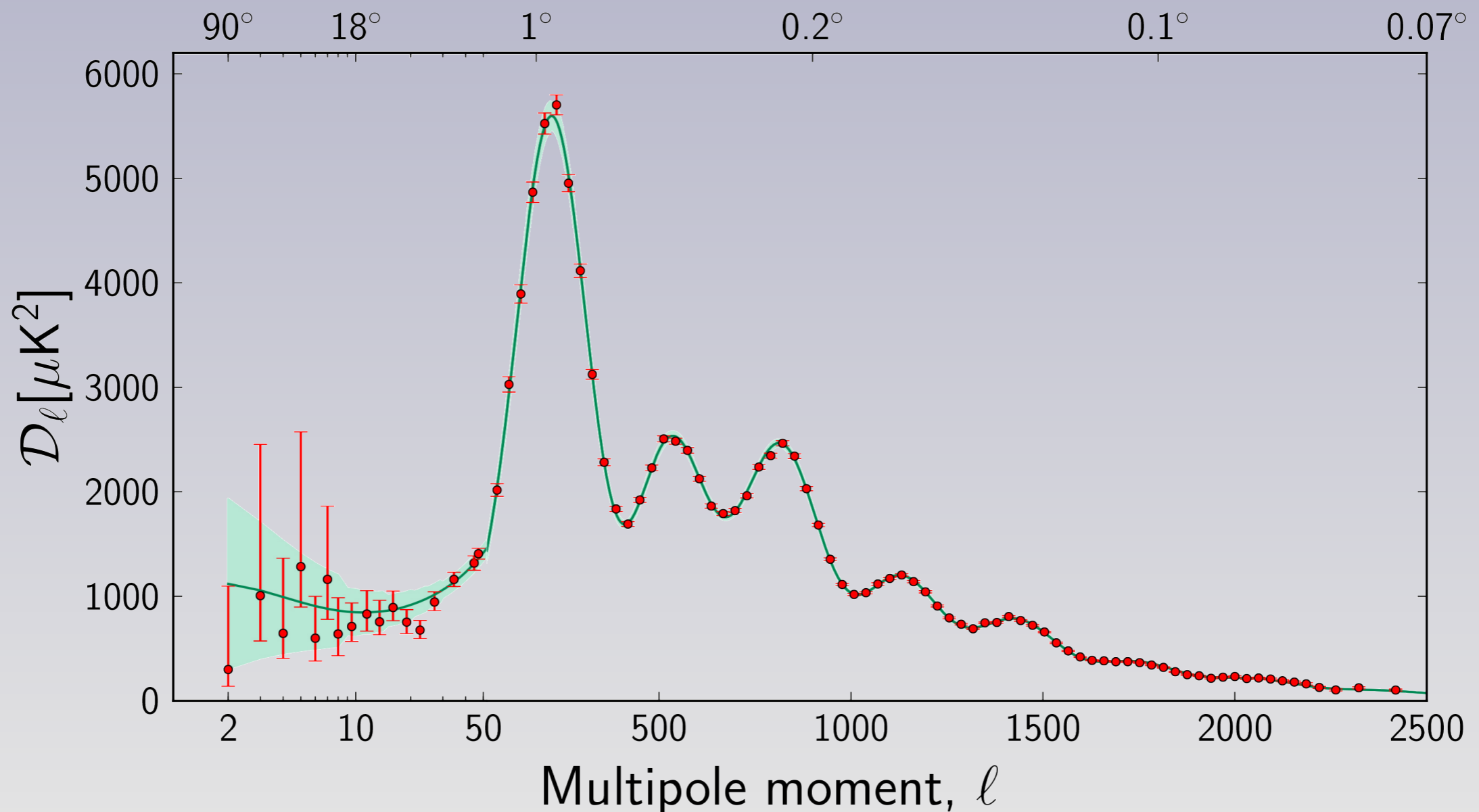
$$T(\phi, \theta) = \sum_{l=0}^{\infty} \sum_{m=-l}^l a_{lm} Y_{lm}(\phi, \theta)$$

$$a_{lm} = \int_{\Omega} d\Omega T(\phi, \theta) Y_{lm}^*$$

=

$$D_l^{TT} \equiv \frac{1}{2l+1} \sum_{m=-l}^l |a_{lm}|^2$$

Angular scale



Correlations

$$\langle T(\phi_1, \theta_1) T^*(\phi_2, \theta_2) \rangle = \sum_{l=0}^{\infty} C_l \frac{2l+1}{4\pi} P_l(\hat{n}_1 \cdot \hat{n}_2)$$

For a gaussian random field:

$$\langle a_{lm} a_{l'm'}^* \rangle \equiv C_l \delta_{ll'} \delta_{mm'}$$

independent of m !

So, $D_l^{TT} \equiv \frac{1}{2l+1} \sum_{m=-l}^l |a_{lm}|^2$ measures $2l+1$ independent random variables, drawn from the same gaussian distribution with variance C_l .

Perturbation Theory

$$\frac{\Delta T(\mathbf{n})}{T} = \left[\frac{1}{4} D_g^{(r)} + V_j^{(b)} n^j + \Psi - \Phi \right] (\eta_{dec}, \mathbf{x}_{dec}) + \int_{\eta_{dec}}^{\eta_0} (\dot{\Psi} - \dot{\Phi})(\eta, \mathbf{x}(\eta)) d\eta$$

see [Durrer, astro-ph/0109522] or
[text book "The Cosmic Microwave Background" by Durrer]

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Photon density

see [Durrer, astro-ph/0109522] or
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Photon density

Velocity (Doppler)
negligible

see [Durrer, astro-ph/0109522] or
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Photon density

Gravitational potential

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negligible

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Photon density

Gravitational potential

Velocity (Doppler)
negligible

Time change of gravitational
potential along the photon path

see [Durrer, astro-ph/0109522] or
[text book "The Cosmic Microwave Background" by Durrer]

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For adiabatic perturbations:

$$D_g^{(r)}(k, \eta) = -20/3 \Psi(k, \eta)$$

$$\left(\frac{\Delta T(\mathbf{n})}{T} \right)_{\text{adiabatic}}^{(OSW)} = \frac{1}{3} \Psi(\eta_{dec}, \mathbf{x}_{dec})$$

see [Durrer, astro-ph/0109522] or
[text book "The Cosmic Microwave Background" by Durrer]

Ordinary Sachs-Wolfe effect

Integrated Sachs-Wolfe effect

$$\frac{\Delta T(\mathbf{n})}{T} = \left[\frac{1}{4} D_g^{(r)} + V_j^{(b)} n^j + \Psi - \Phi \right] (\eta_{dec}, \mathbf{x}_{dec}) + \int_{\eta_{dec}}^{\eta_0} (\dot{\Psi} - \dot{\Phi})(\eta, \mathbf{x}(\eta)) d\eta$$

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Gravity wins!

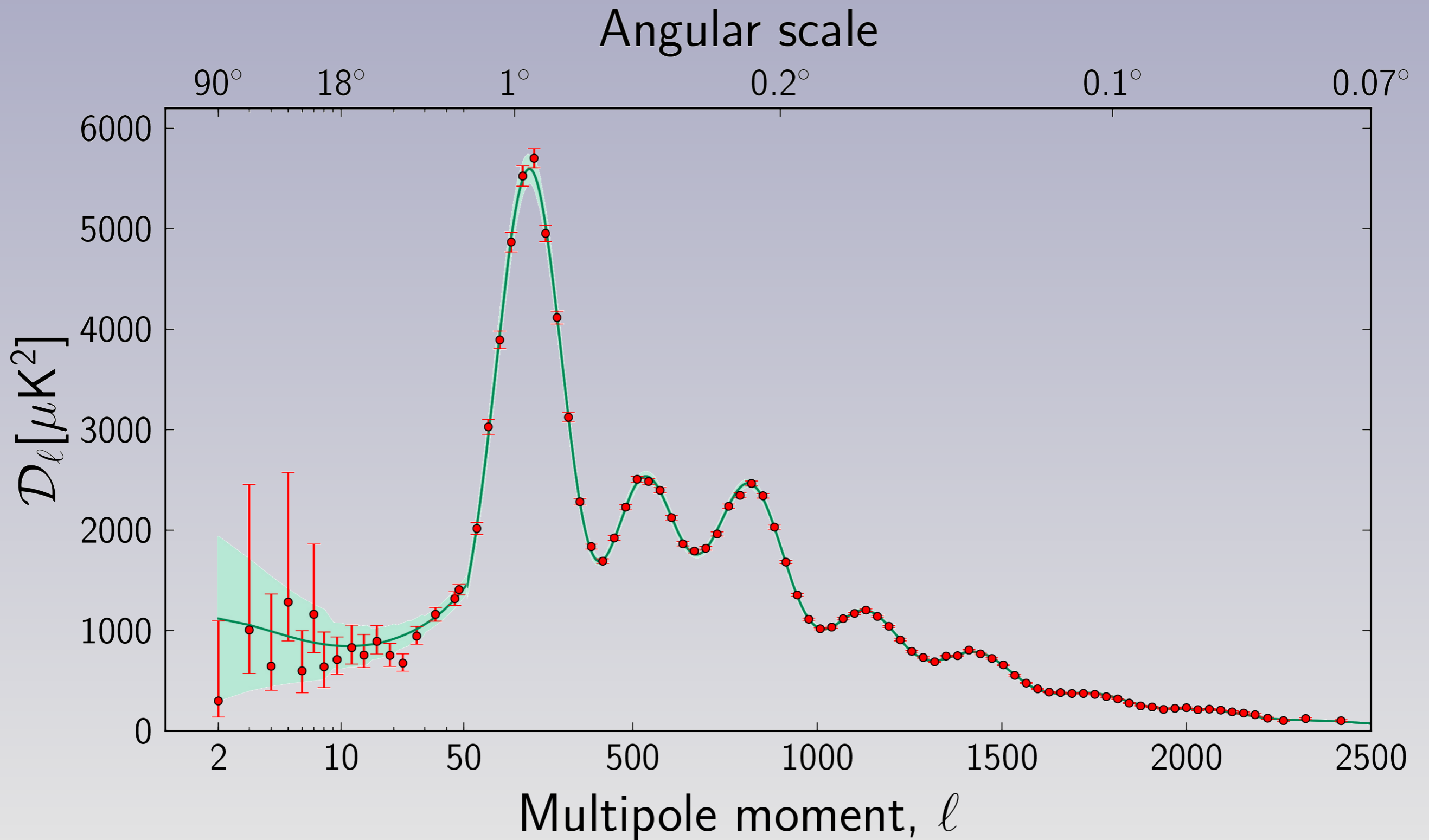
Hot = underdense

Cold = overdense

$$\left(\frac{\Delta T(\mathbf{n})}{T} \right)_{\text{adiabatic}}^{(OSW)} = \frac{1}{3} \Psi(\eta_{dec}, \mathbf{x}_{dec})$$

see [Durrer, astro-ph/0109522] or
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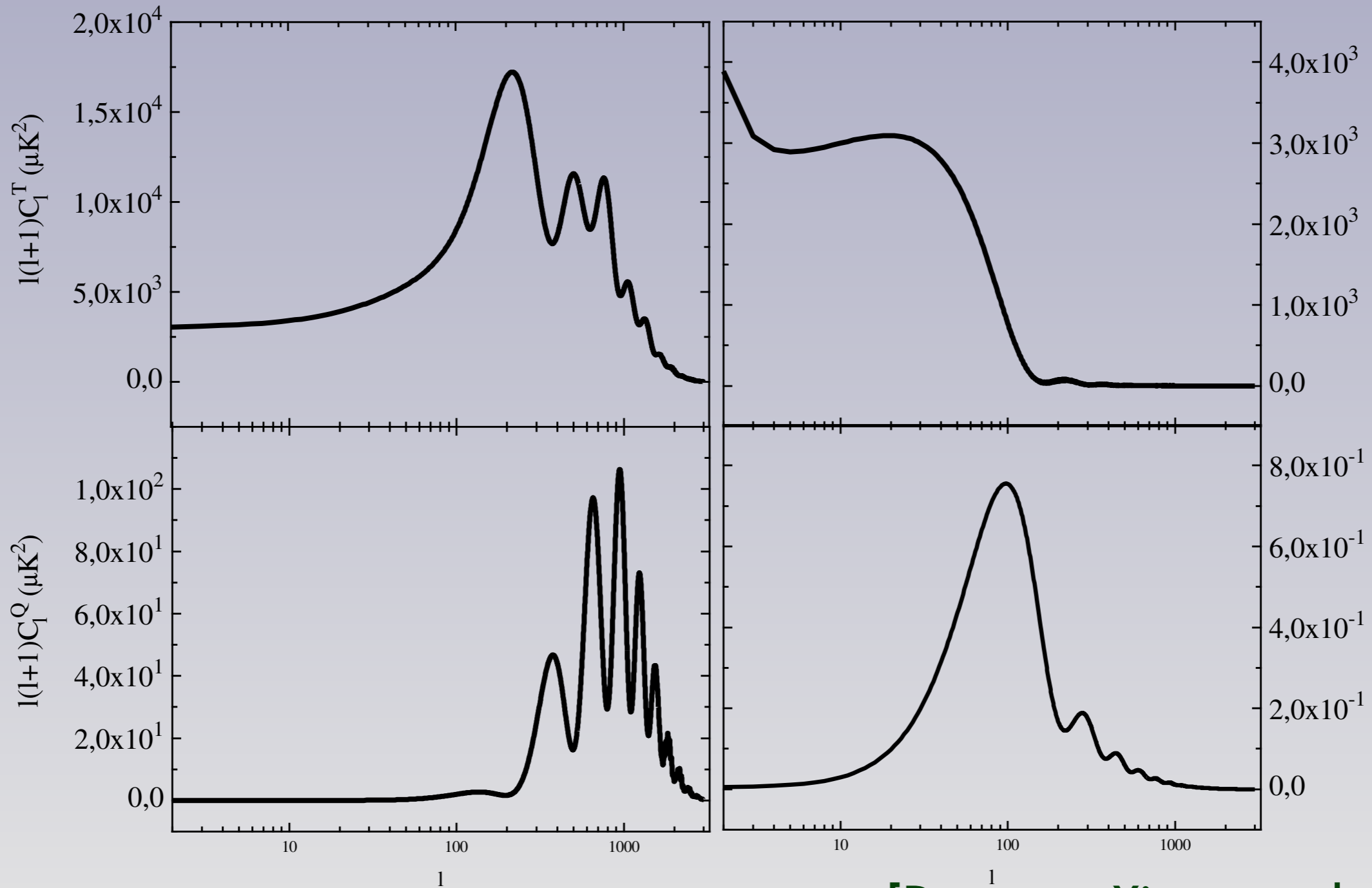
Shortcut



CMB

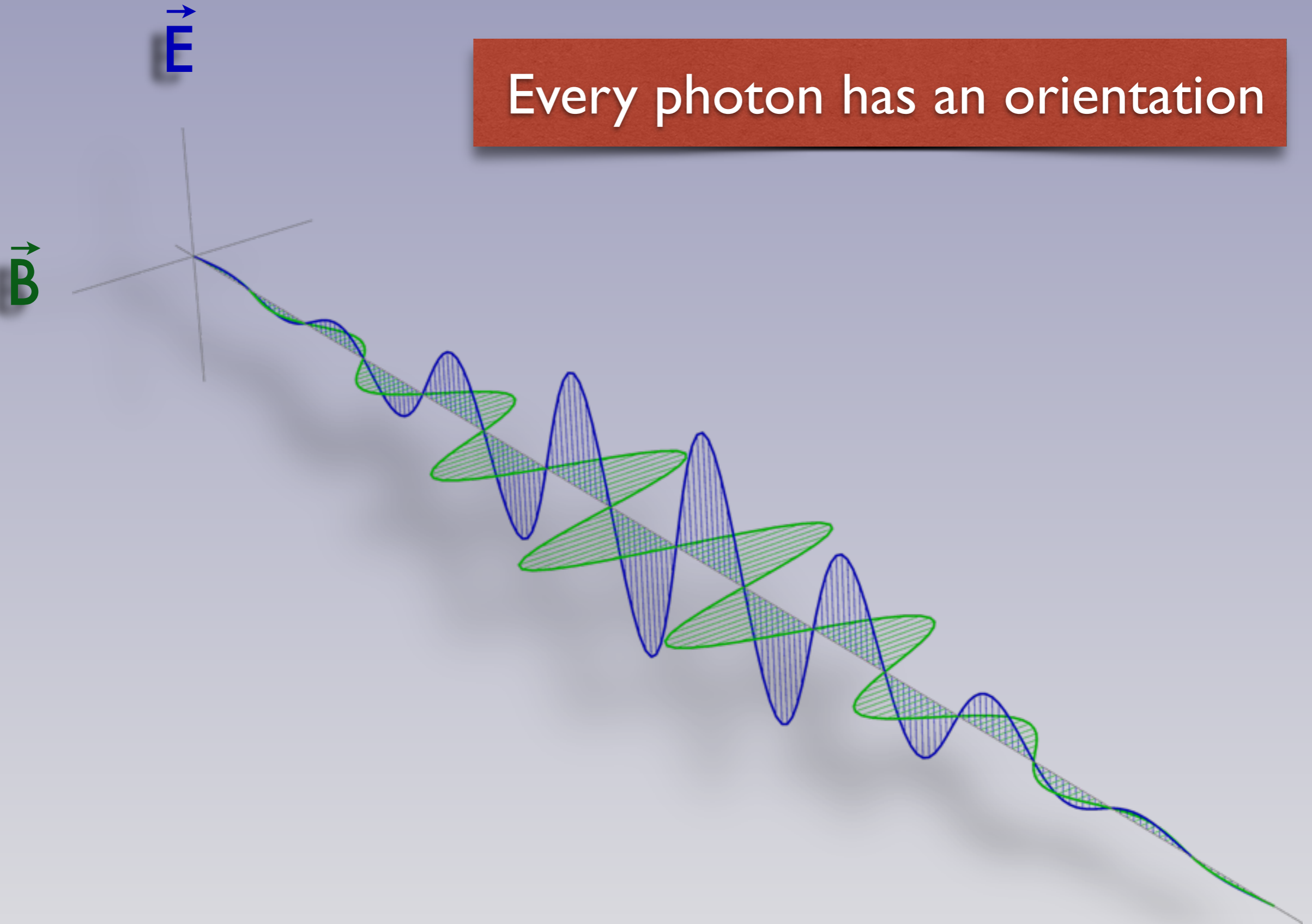
Scalar

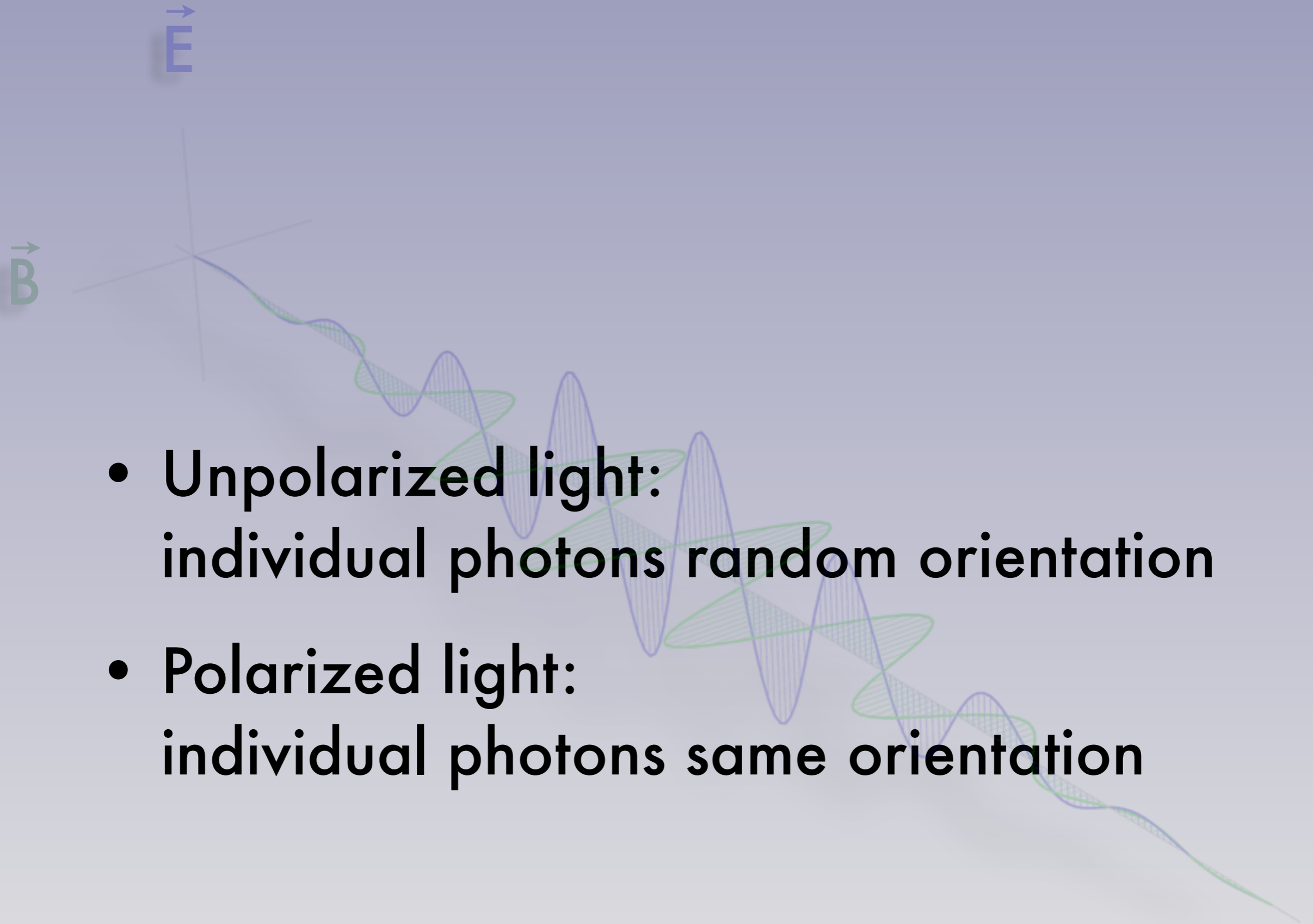
Tensor



[Durrer, arXiv:astro-ph/0109522]

Every photon has an orientation





- **Unpolarized light:**
individual photons random orientation
- **Polarized light:**
individual photons same orientation

Polarization by reflection

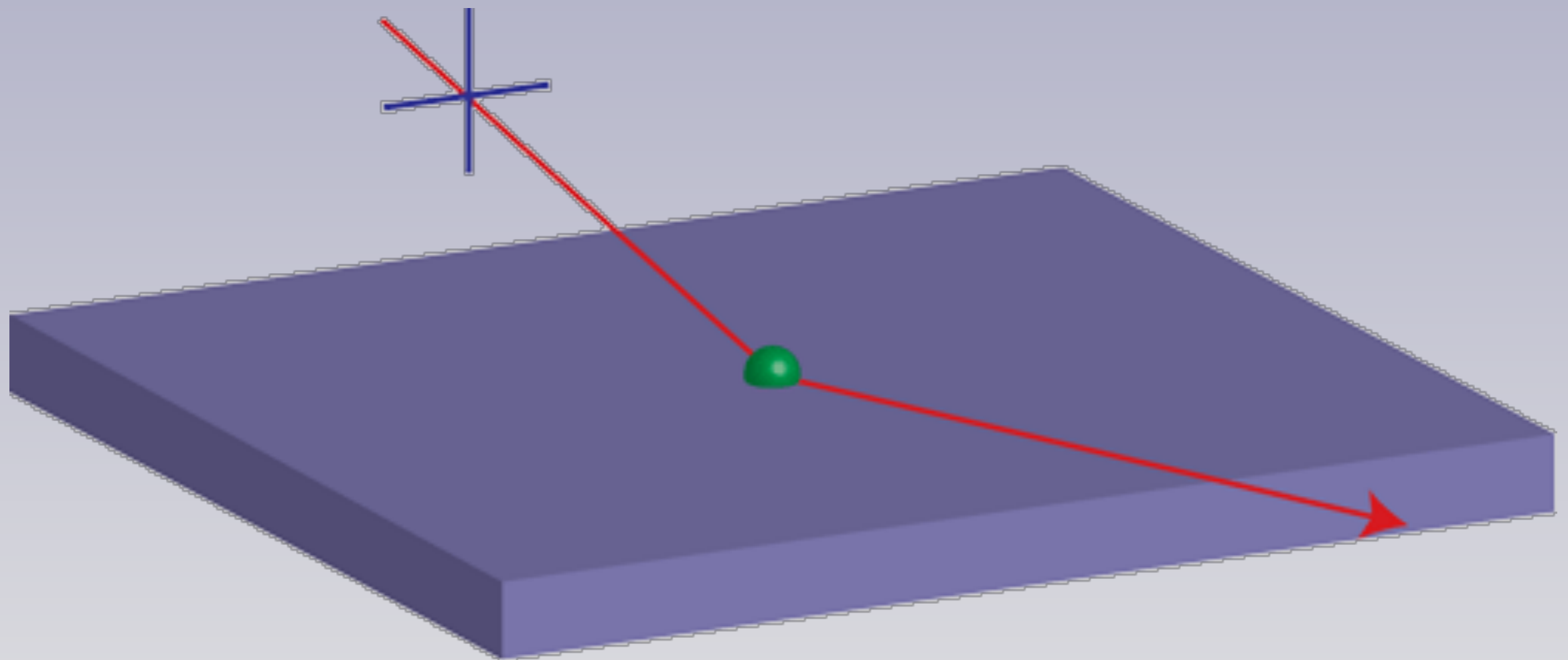


Image by Wayne Hu [<http://background.uchicago.edu/~whu/>]

Polarization by reflection

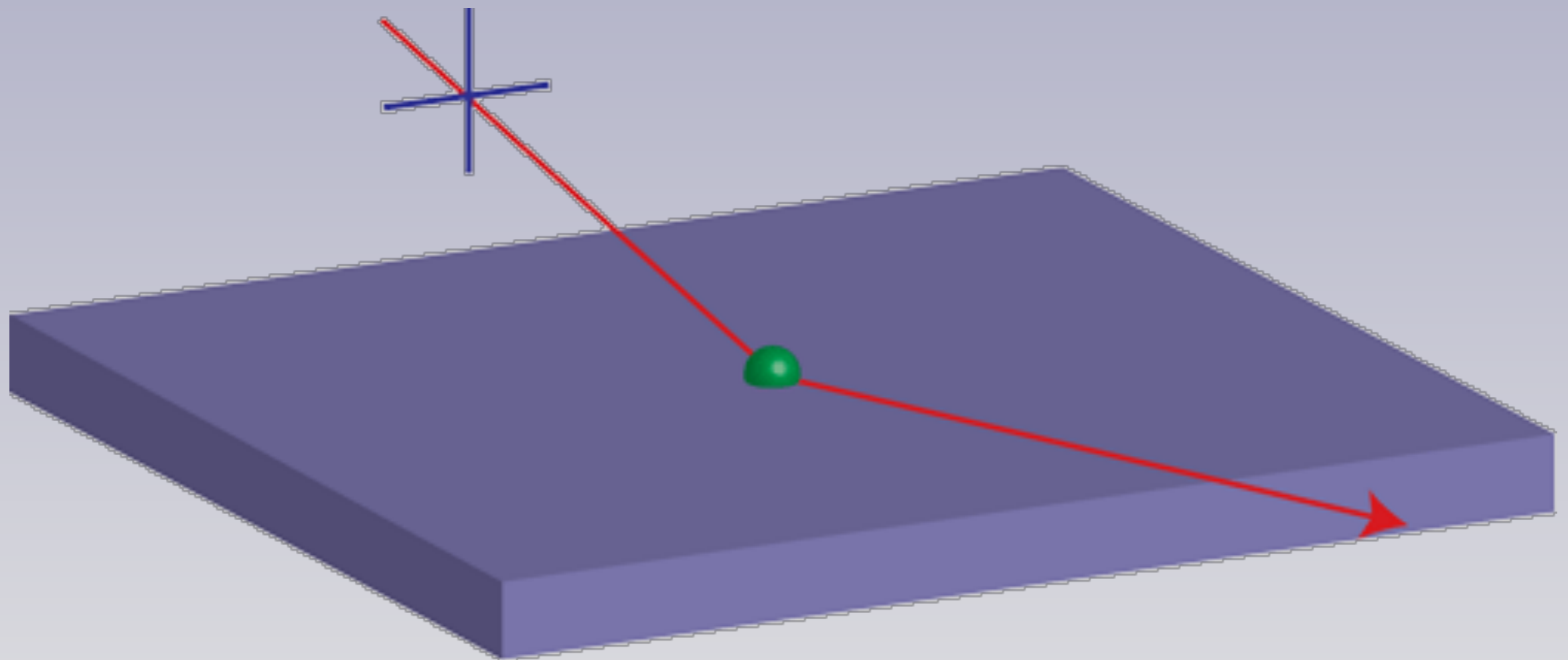


Image by Wayne Hu [<http://background.uchicago.edu/~whu/>]

Polarization by Thomson scattering

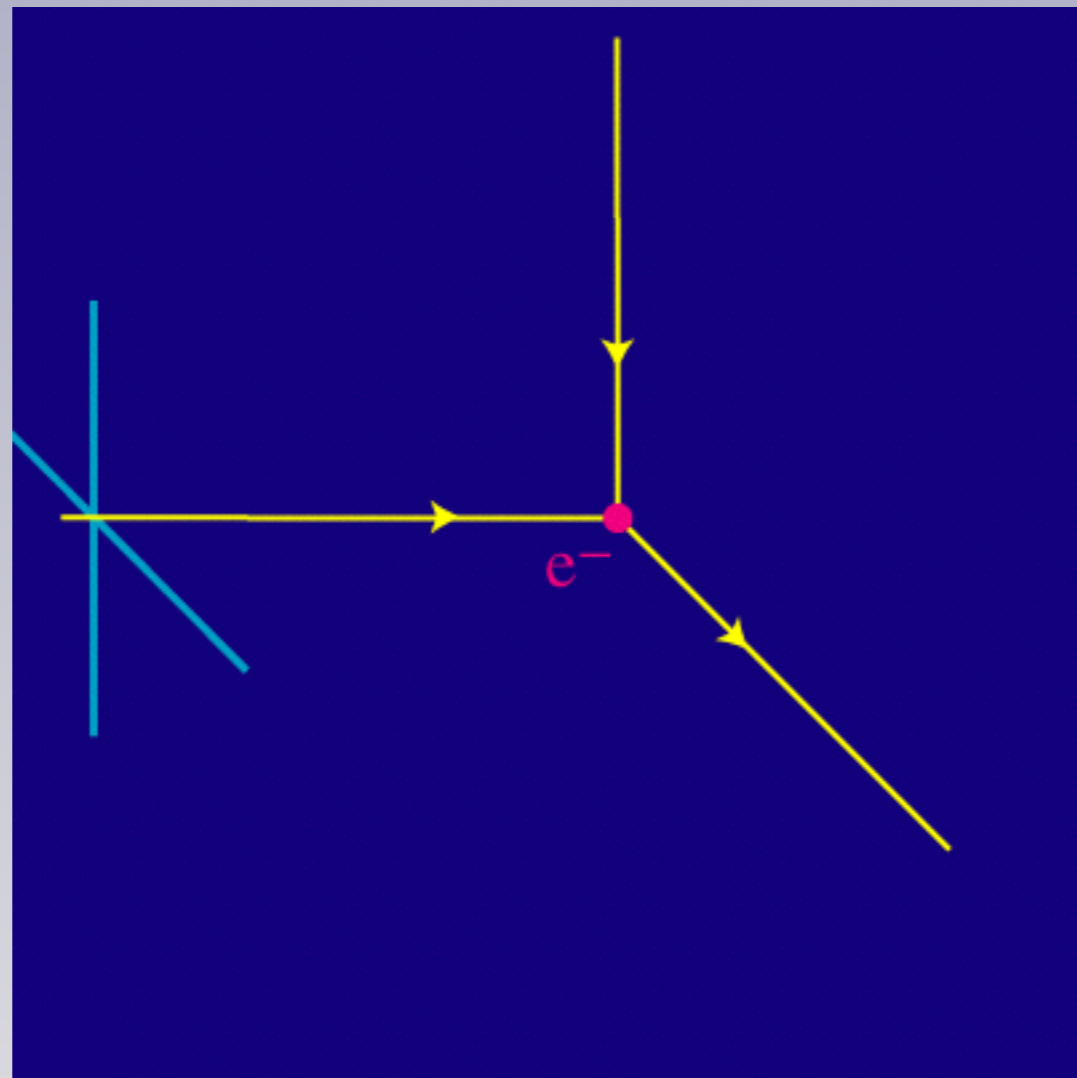


Image by Wayne Hu [<http://background.uchicago.edu/~whu/>]

Polarization by Thomson scattering

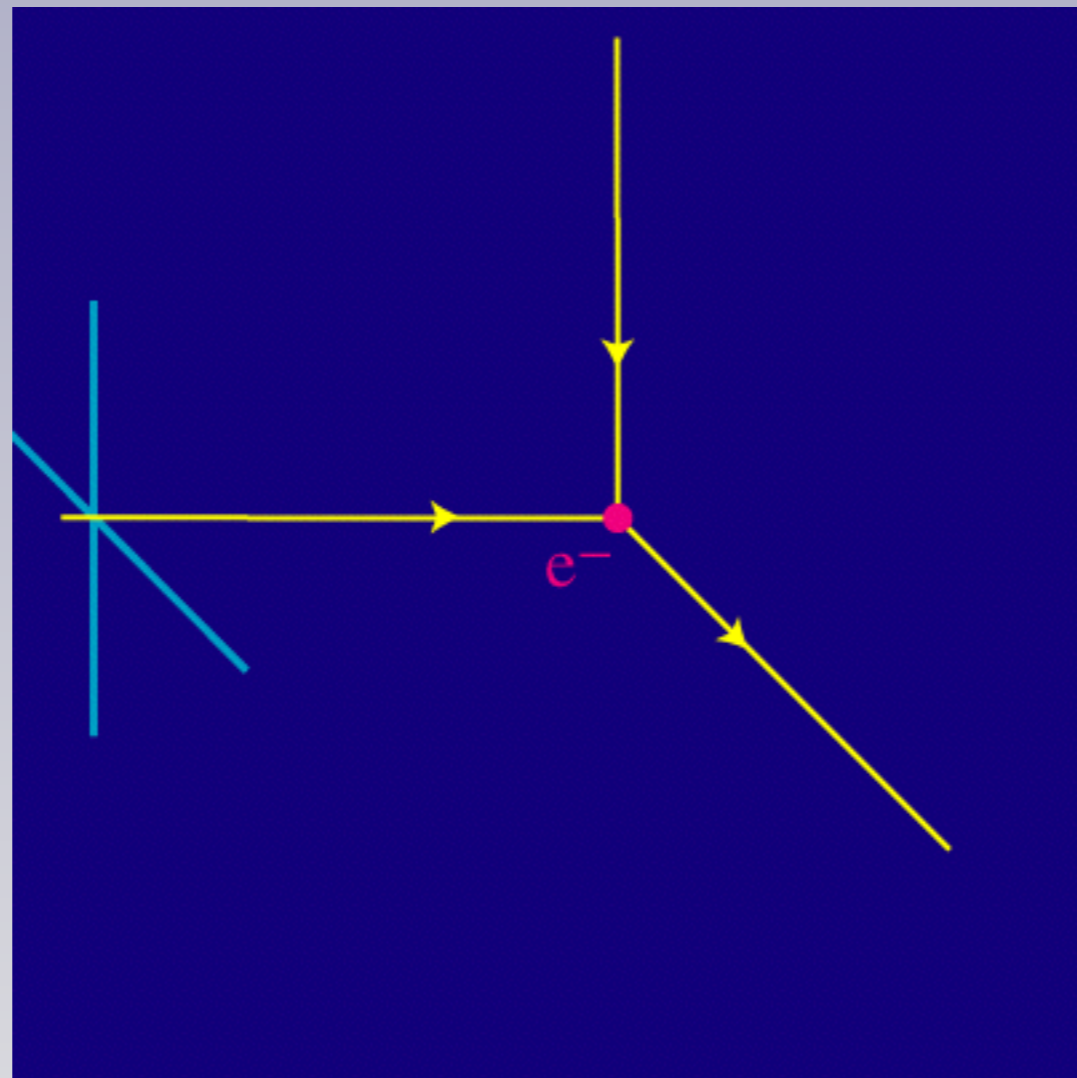


Image by Wayne Hu [<http://background.uchicago.edu/~whu/>]

No polarization by Thomson scattering for isotropic influx

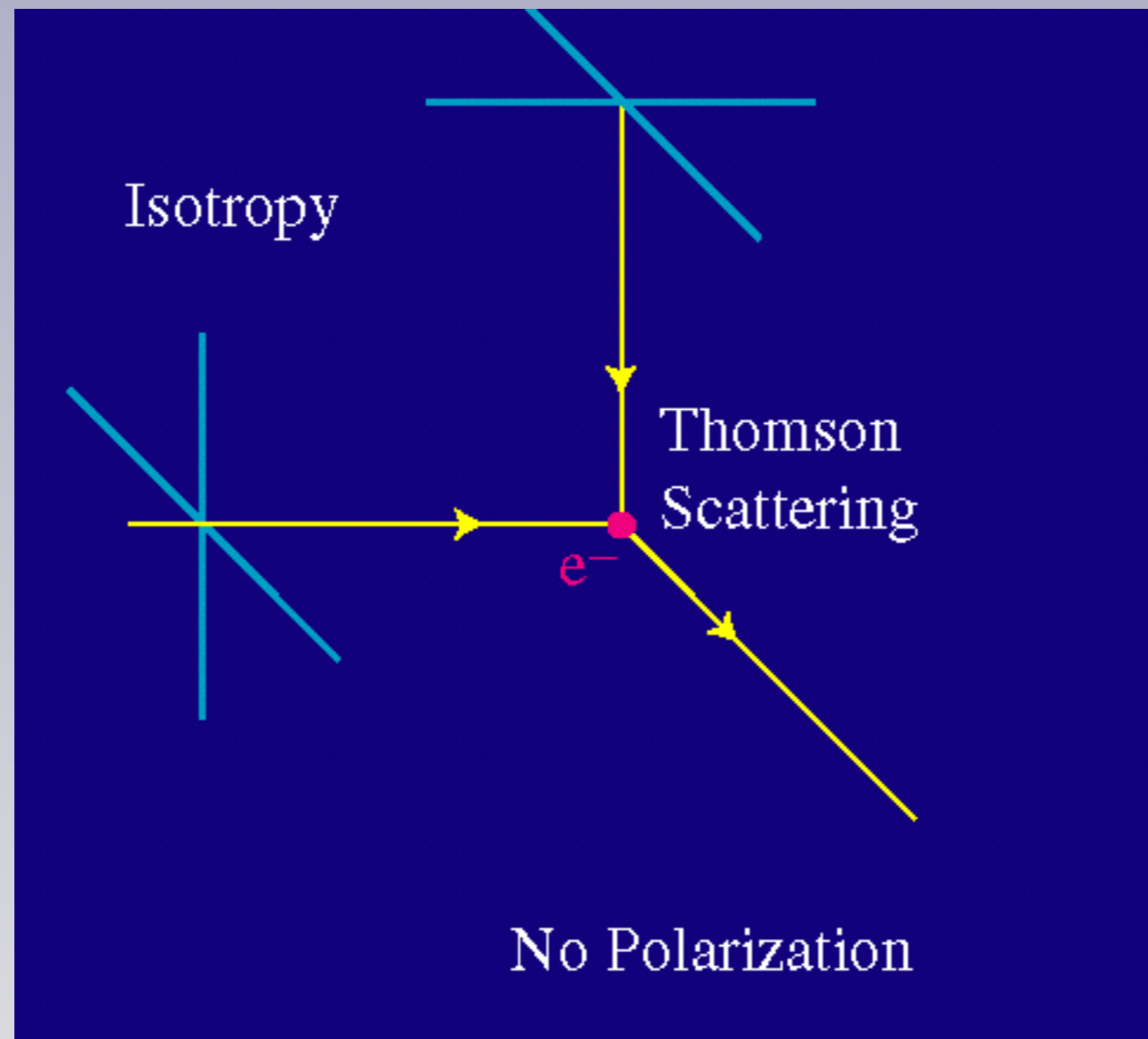


Image by Wayne Hu [<http://background.uchicago.edu/~whu/>]

No polarization by Thomson scattering for isotropic influx

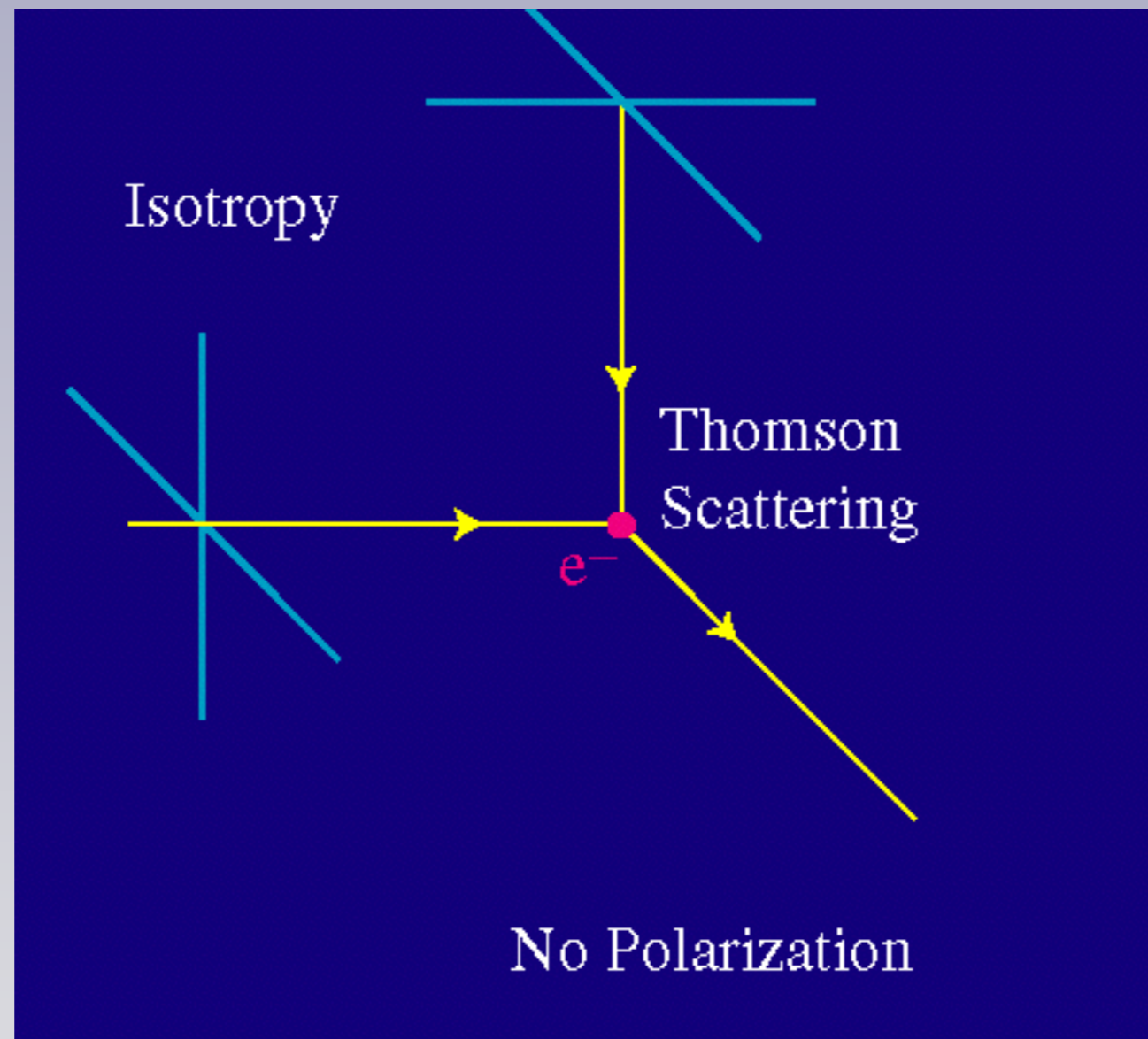


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Polarization by Thomson scattering of influx with quadrupolar anisotropy

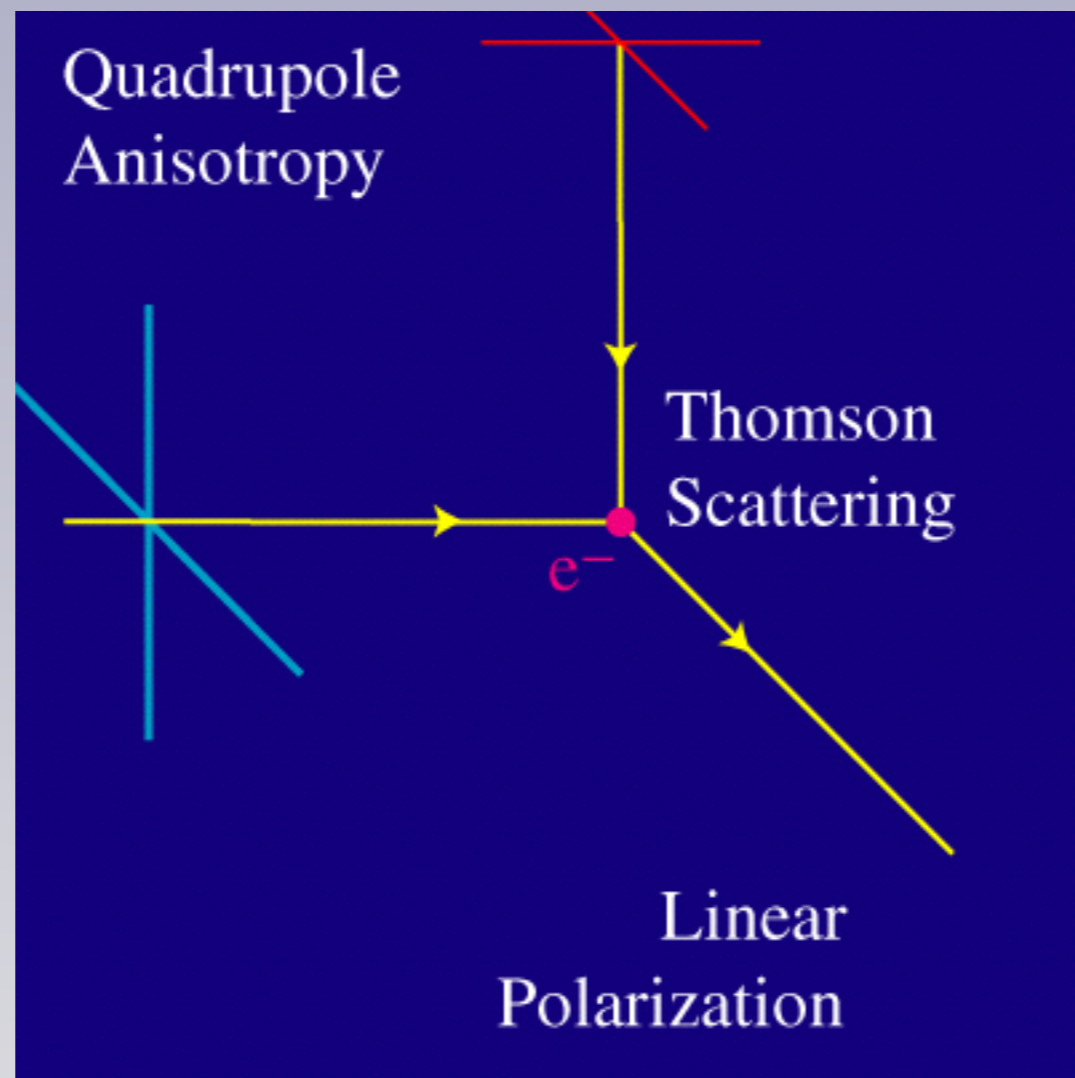


Image by Wayne Hu [<http://background.uchicago.edu/~whu/>]

Polarization by Thomson scattering of influx with quadrupolar anisotropy

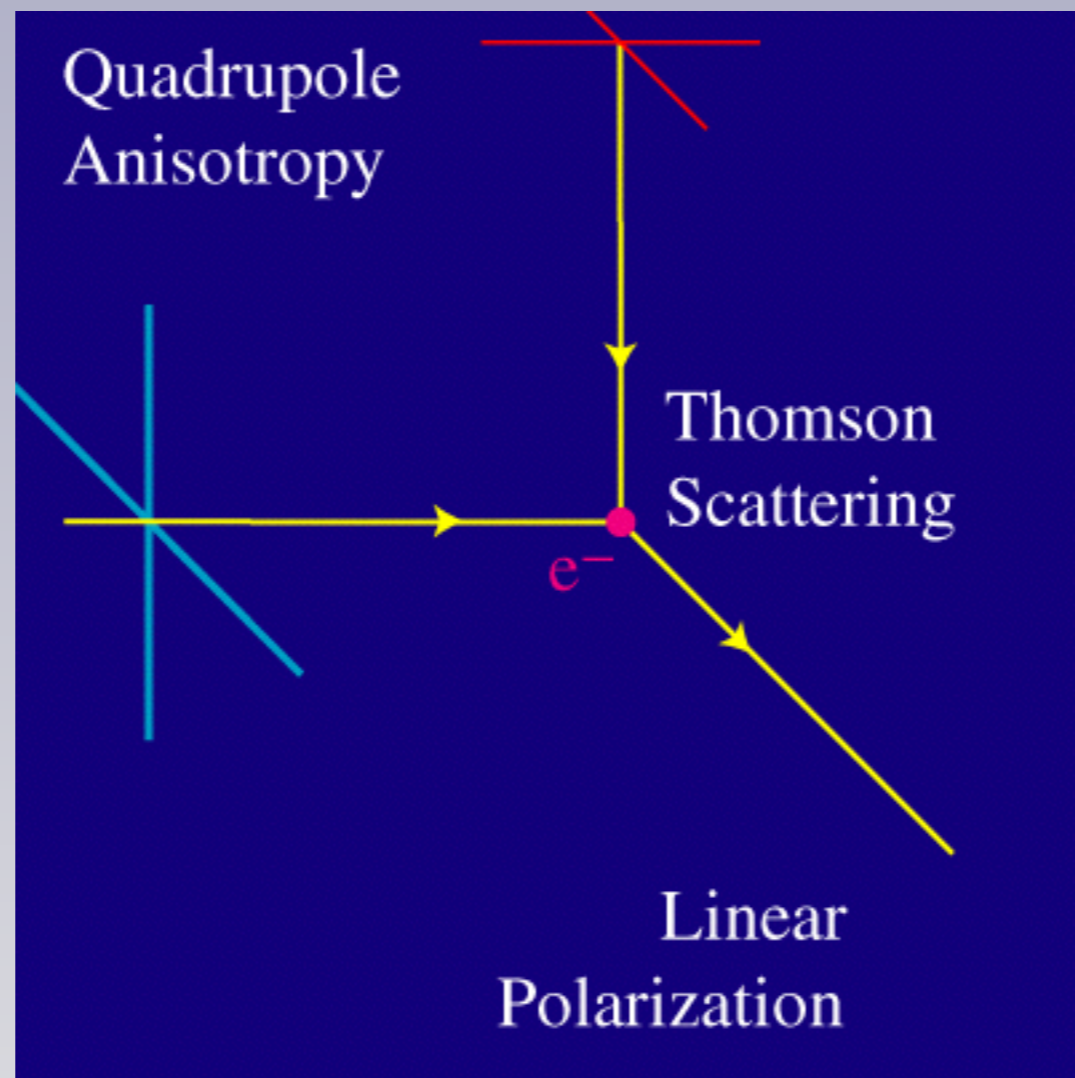
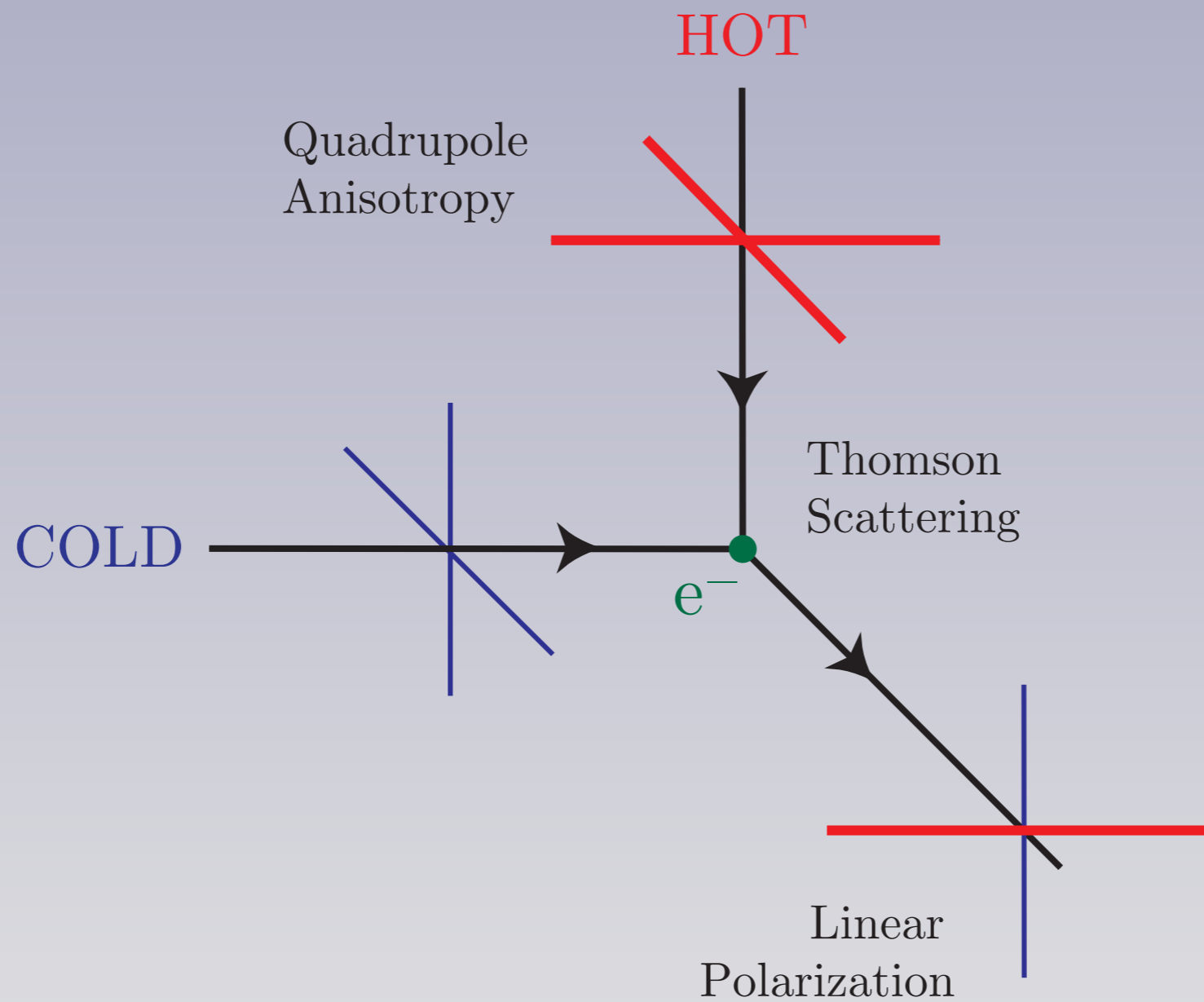
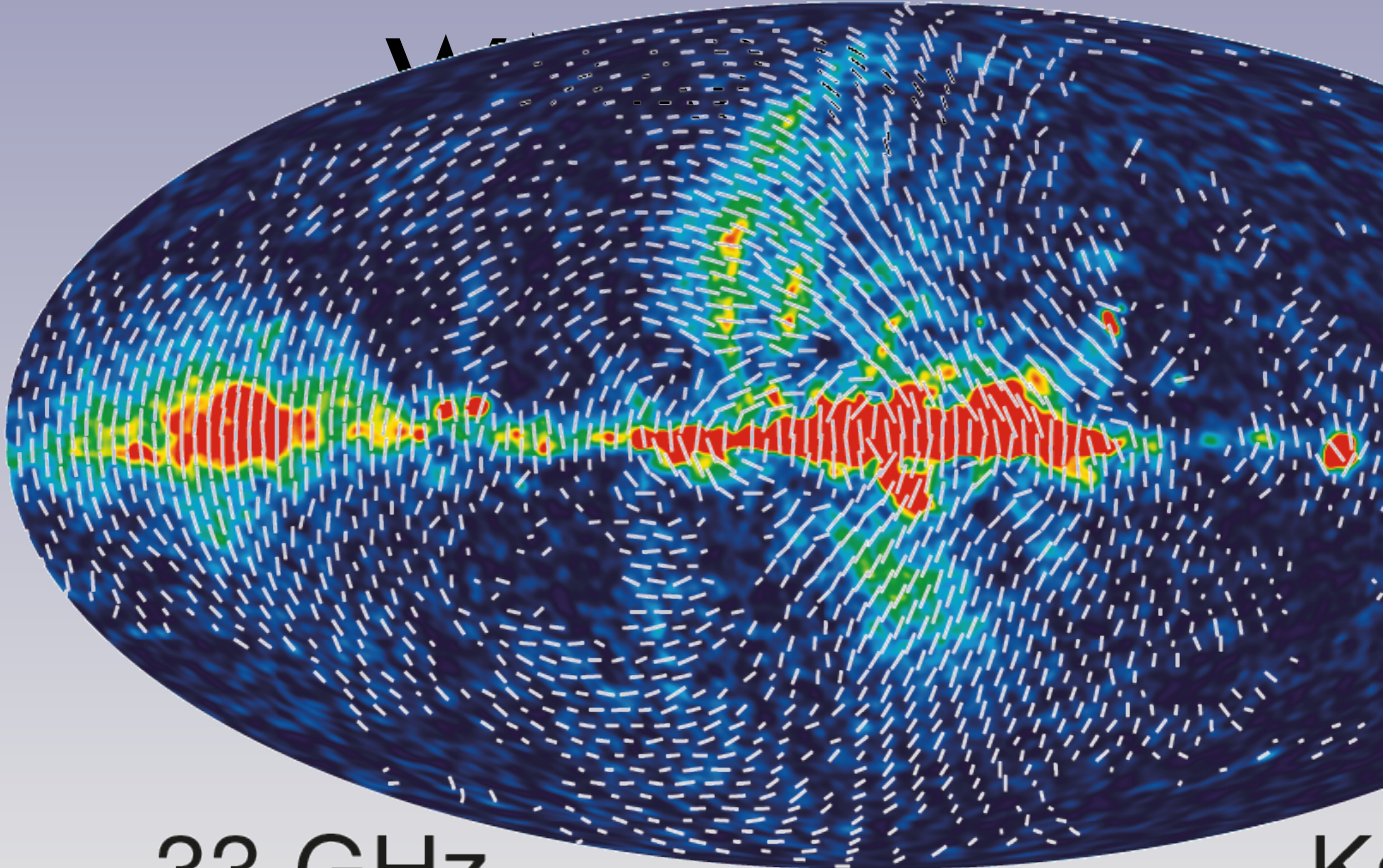


Image by Wayne Hu [<http://background.uchicago.edu/~whu/>]



[from Baumann et al. 2008]

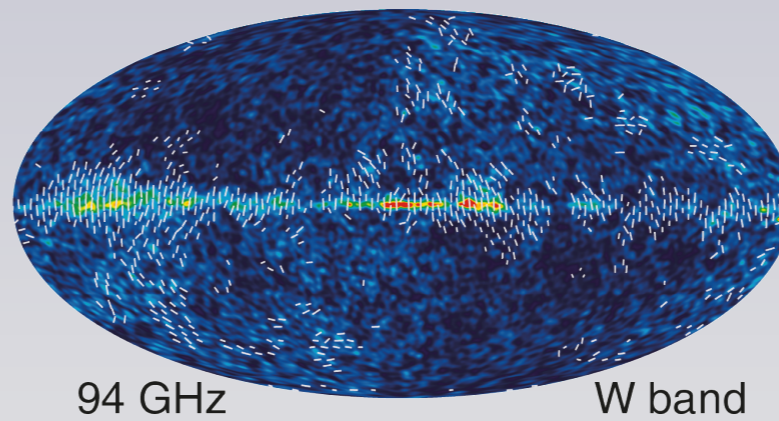
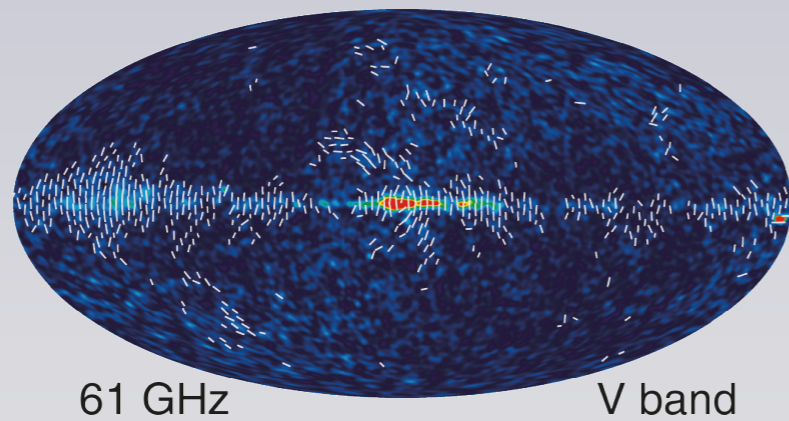
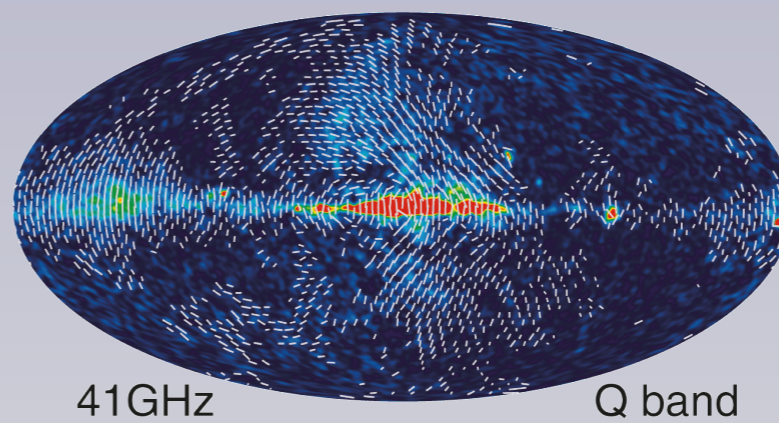
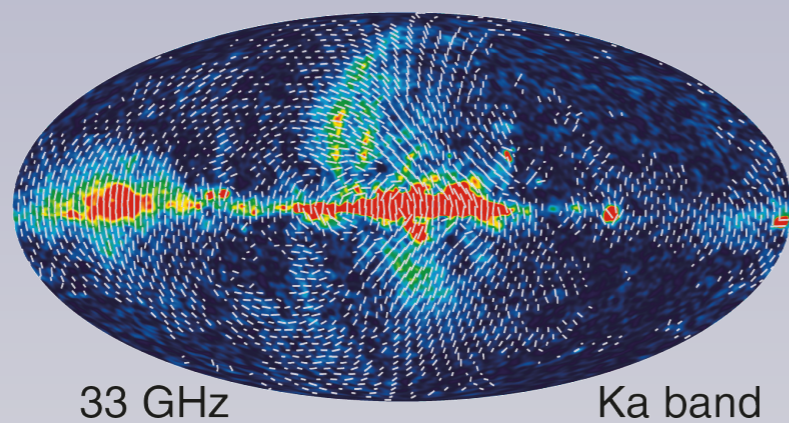
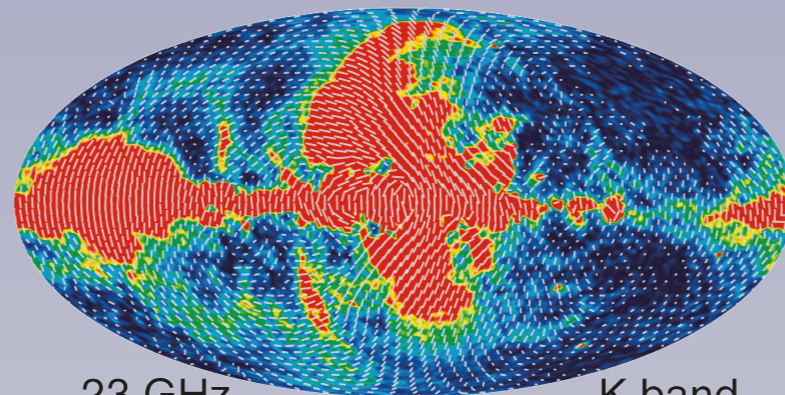
20 GHz



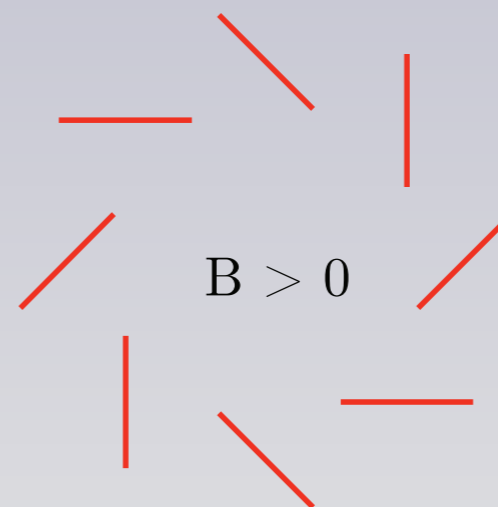
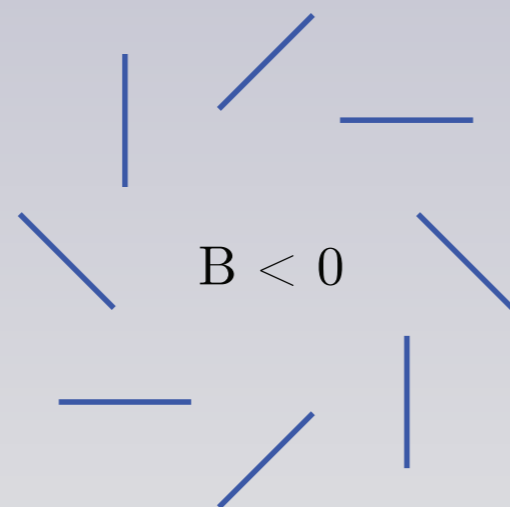
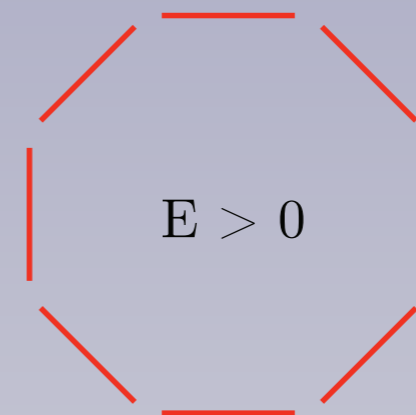
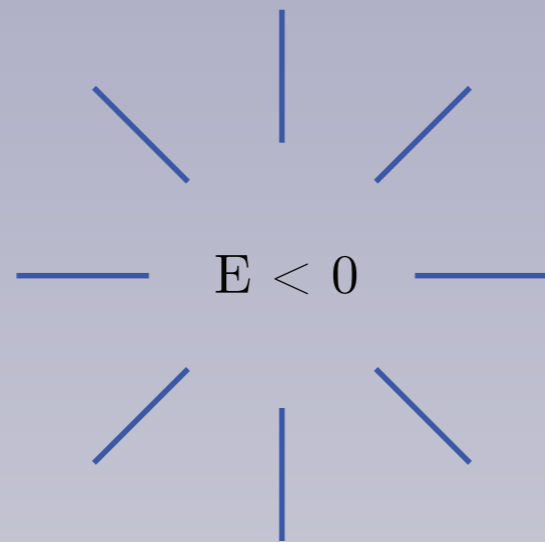
33 GHz

Ka

WMAP 9yr

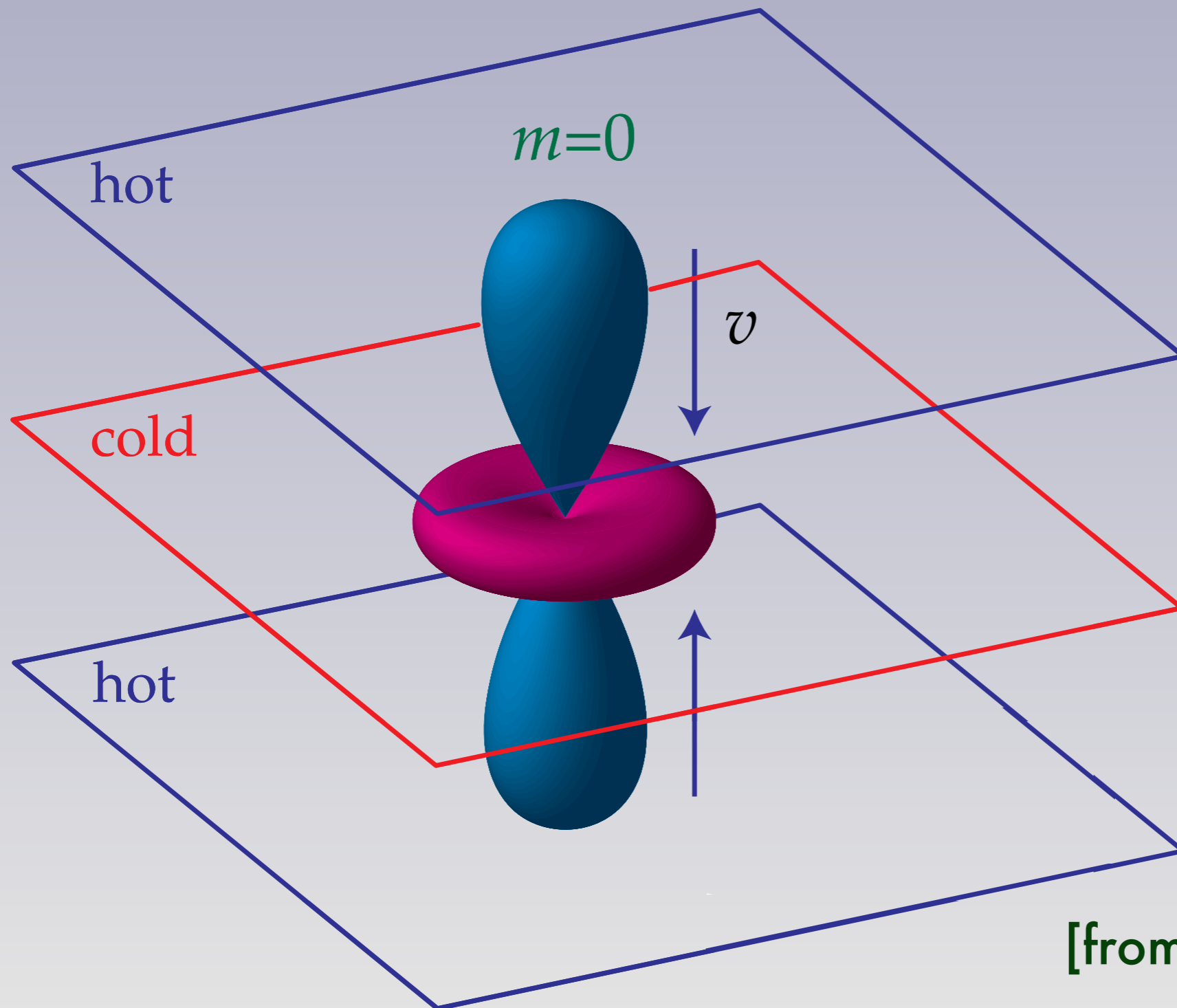


Decompose the vector field



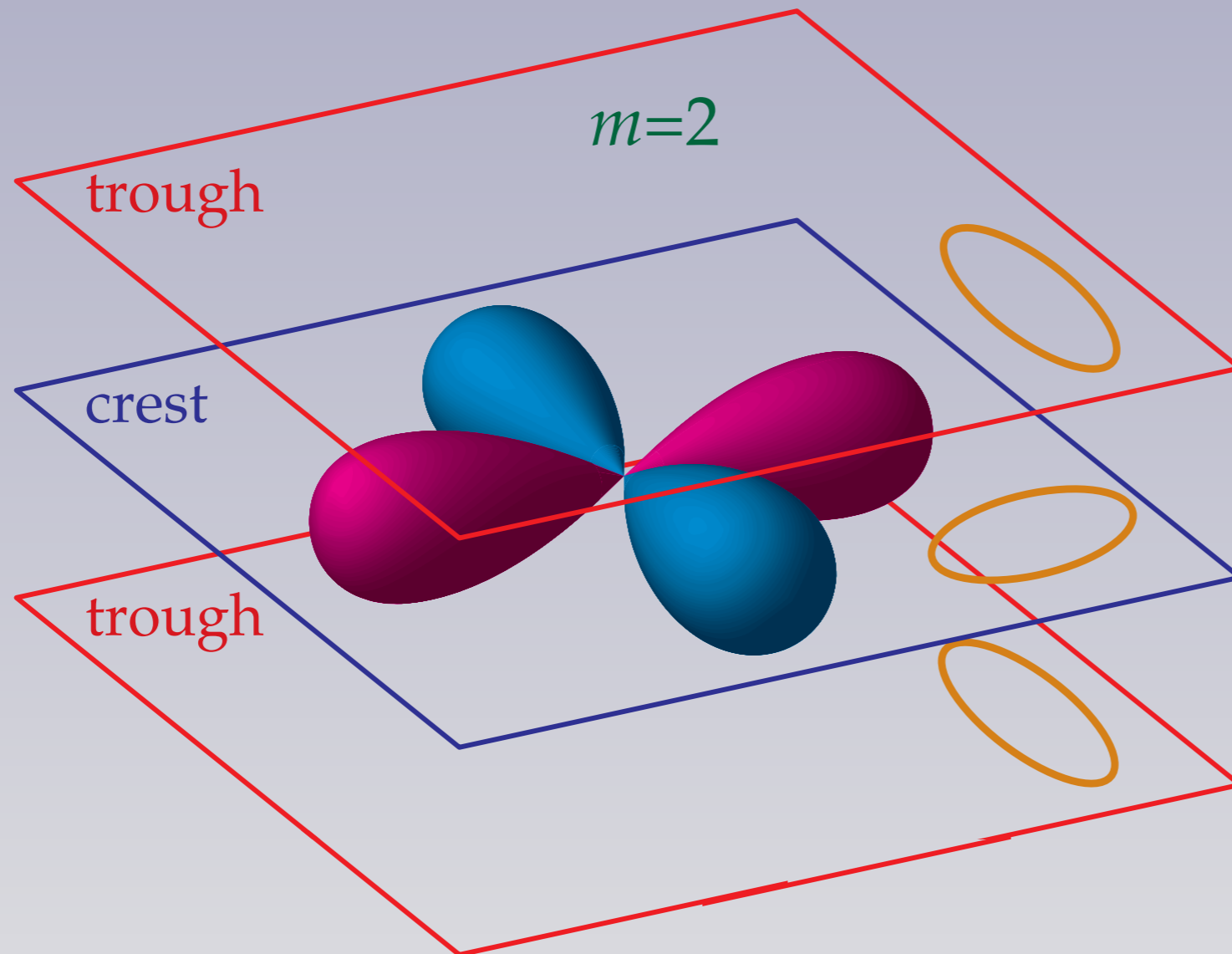
[from Baumann et al. 2008]

Local quadrupole leads to polarized radiation



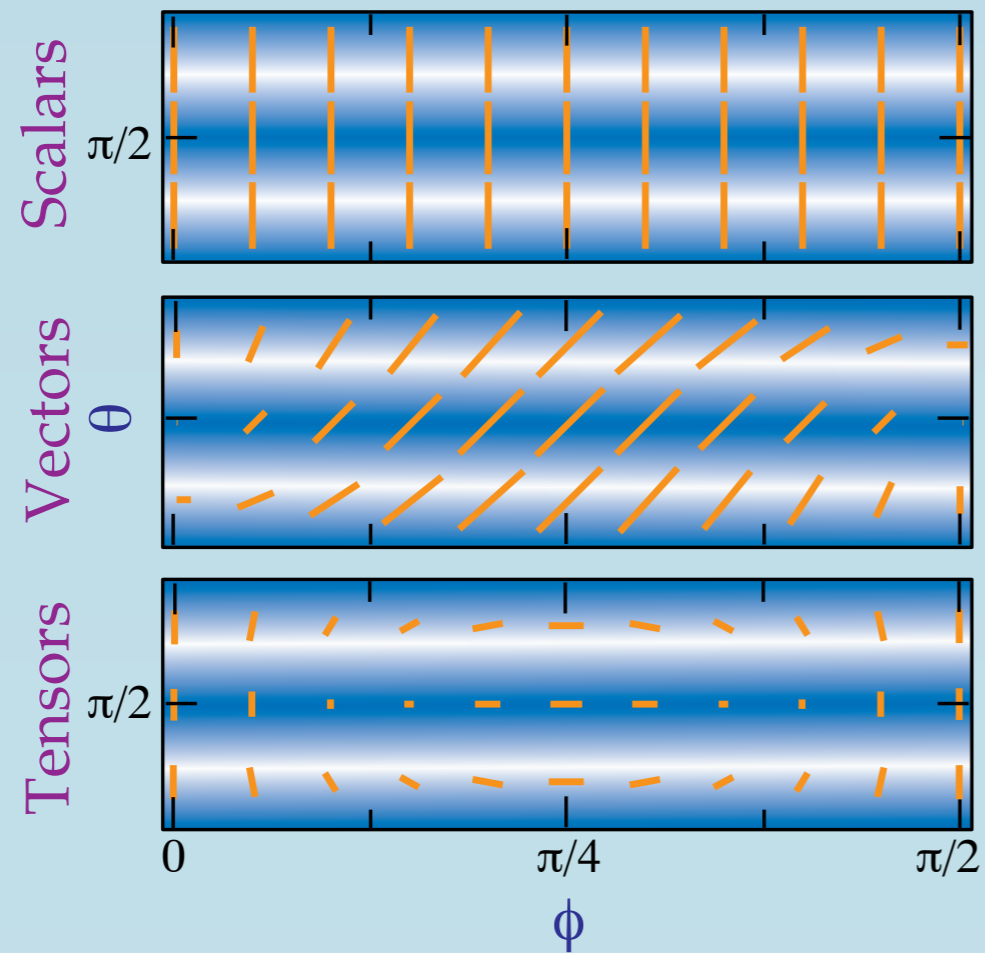
[from Hu & White 1997]

Local quadrupole with gravity wave sources E&B modes

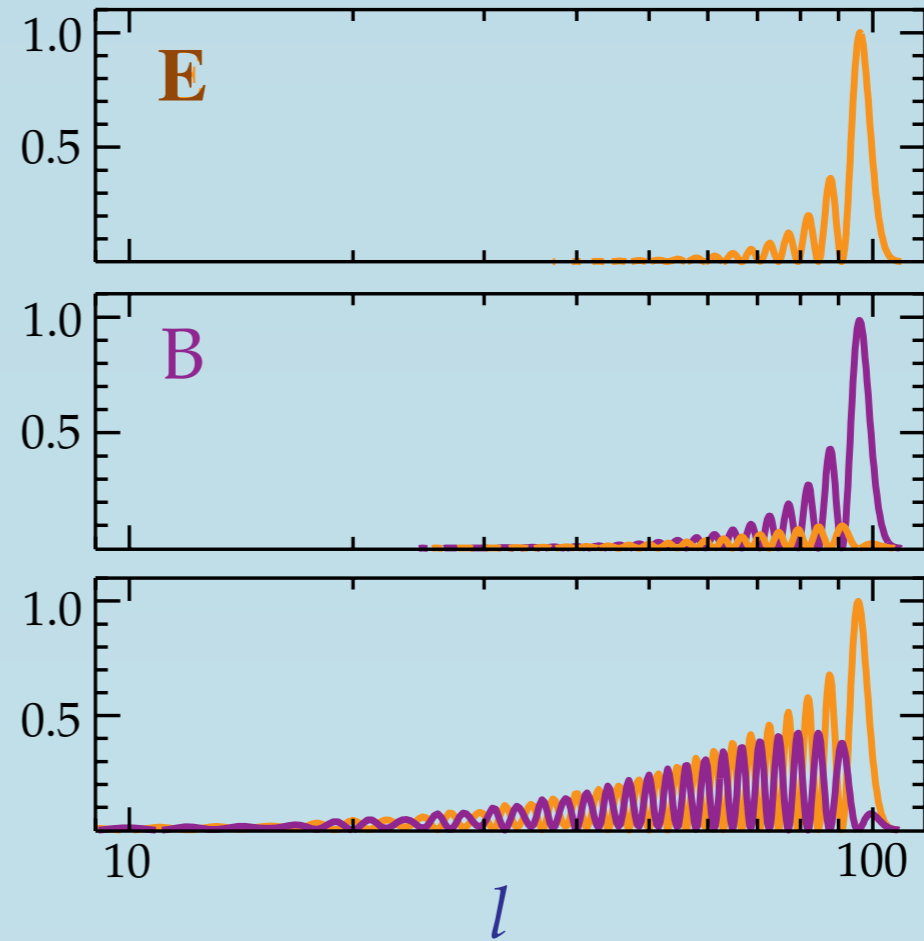


[from Hu & White 1997]

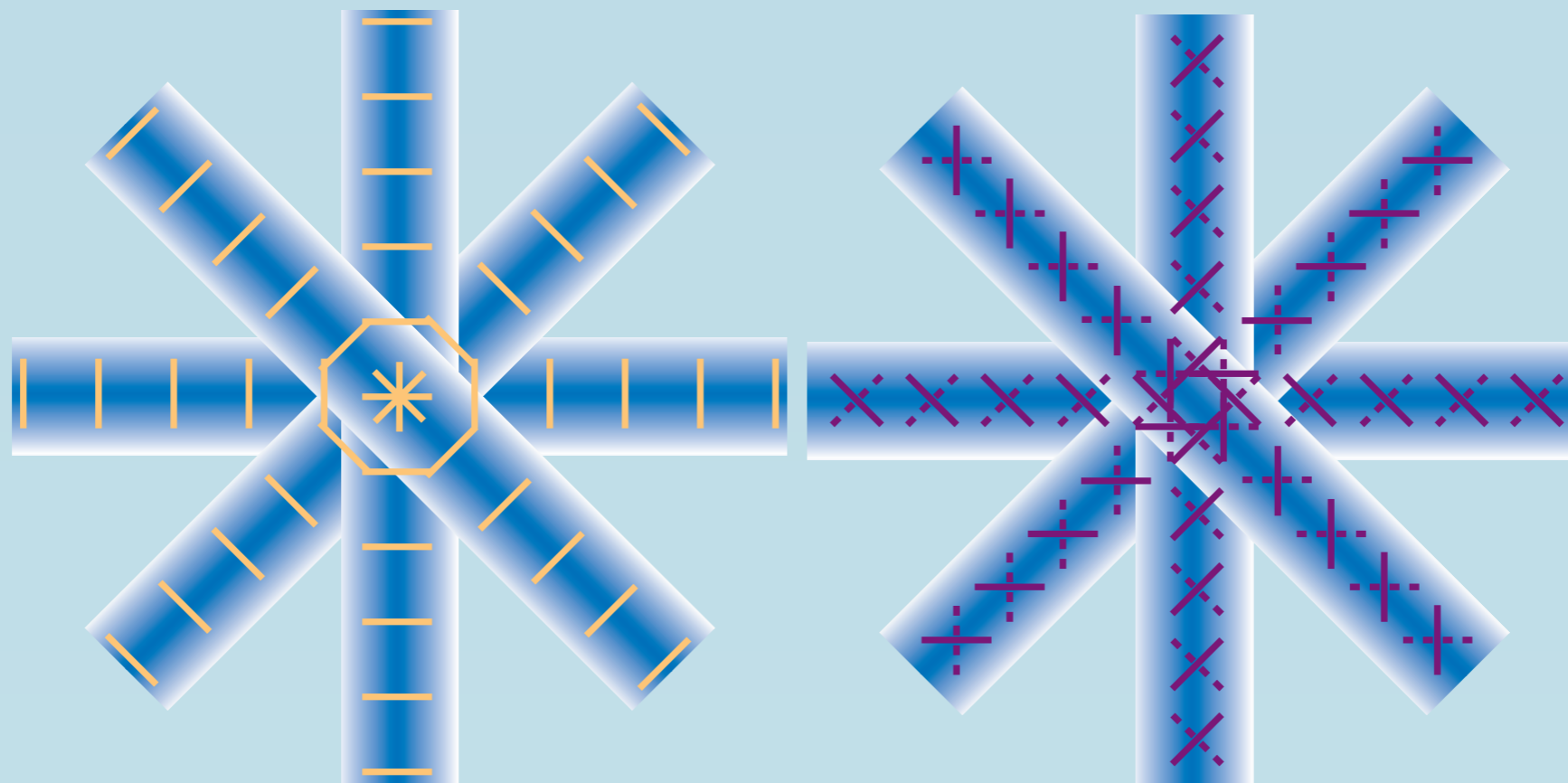
(a) Polarization Pattern



(b) Multipole Power



[from Hu & White 1997]



E (anti) correlation

B no correlation

Equations of motion

$$\dot{\mathcal{M}}_\ell^{(m)} - k \left[\frac{0\kappa_\ell^m}{2\ell-1} \mathcal{M}_{\ell-1}^{(m)} - \frac{0\kappa_{\ell+1}^m}{2\ell+3} \mathcal{M}_{\ell+1}^{(m)} \right] = -n_e \sigma_T a \mathcal{M}_\ell^{(m)} + S_\ell^{(m)} \quad (\ell \geq m)$$

$$\dot{E}_\ell^{(m)} - k \left[\frac{2\kappa_\ell^m}{2\ell-1} E_{\ell-1}^{(m)} - \frac{2m}{\ell(\ell+1)} B_\ell^{(m)} - \frac{2\kappa_{\ell+1}^m}{2\ell+3} E_{\ell+1}^{(m)} \right] = -n_e \sigma_T a [E_\ell^{(m)} + \sqrt{6} C^{(m)} \delta_{\ell,2}]$$

$$\dot{B}_\ell^{(m)} - k \left[\frac{2\kappa_\ell^m}{2\ell-1} B_{\ell-1}^{(m)} + \frac{2m}{\ell(\ell+1)} E_\ell^{(m)} - \frac{2\kappa_{\ell+1}^m}{2\ell+3} B_{\ell+1}^{(m)} \right] = -n_e \sigma_T a B_\ell^{(m)}$$

Equations of motion

$$\dot{\mathcal{M}}_\ell^{(m)} - k \left[\frac{0\kappa_\ell^m}{2\ell-1} \mathcal{M}_{\ell-1}^{(m)} - \frac{0\kappa_{\ell+1}^m}{2\ell+3} \mathcal{M}_{\ell+1}^{(m)} \right] = -n_e \sigma_T a \mathcal{M}_\ell^{(m)} + S_\ell^{(m)} \quad (\ell \geq m)$$

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$$\begin{aligned} S_0^{(0)} &= n_e \sigma_T a \mathcal{M}_0^{(0)}, & S_1^{(0)} &= n_e \sigma_T a 4V_b + 4k(\Psi - \Phi), \\ S_2^{(0)} &= n_e \sigma_T a C^{(0)}, & S_1^{(1)} &= n_e \sigma_T a 4\omega_b, \\ S_2^{(1)} &= n_e \sigma_T a C^{(1)} + 4k\Sigma, & S_2^{(2)} &= n_e \sigma_T a C^{(2)} + 4\dot{H} \end{aligned}$$

Equations of motion

$$\dot{\mathcal{M}}_\ell^{(m)} - k \left[\frac{0\kappa_\ell^m}{2\ell-1} \mathcal{M}_{\ell-1}^{(m)} - \frac{0\kappa_{\ell+1}^m}{2\ell+3} \mathcal{M}_{\ell+1}^{(m)} \right] = -n_e \sigma_T a \mathcal{M}_\ell^{(m)} + S_\ell^{(m)} \quad (\ell \geq m)$$

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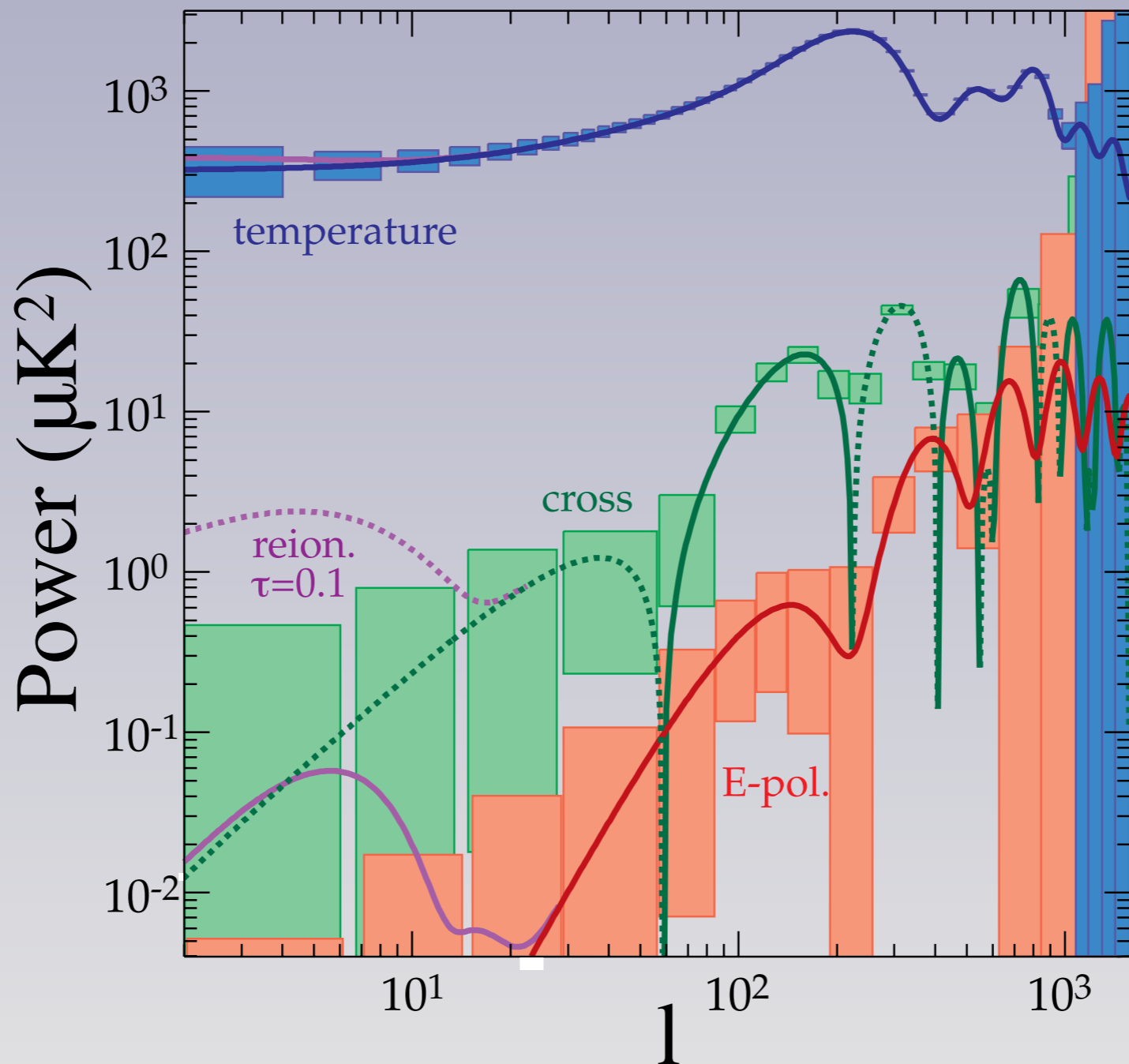
$$\dot{B}_\ell^{(m)} - k \left[\frac{2\kappa_\ell^m}{2\ell-1} B_{\ell-1}^{(m)} + \frac{2m}{\ell(\ell+1)} E_\ell^{(m)} - \frac{2\kappa_{\ell+1}^m}{2\ell+3} B_{\ell+1}^{(m)} \right] = -n_e \sigma_T a B_\ell^{(m)}$$

$$C^{(m)} = \frac{1}{10} [\mathcal{M}_2^{(m)} - \sqrt{6} E_2^{(m)}]$$

$$\begin{aligned} S_0^{(0)} &= n_e \sigma_T a \mathcal{M}_0^{(0)}, & S_1^{(0)} &= n_e \sigma_T a 4V_b + 4k(\Psi - \Phi), \\ S_2^{(0)} &= n_e \sigma_T a C^{(0)}, & S_1^{(1)} &= n_e \sigma_T a 4\omega_b, \\ S_2^{(1)} &= n_e \sigma_T a C^{(1)} + 4k\Sigma, & S_2^{(2)} &= n_e \sigma_T a C^{(2)} + 4\dot{H} \end{aligned}$$

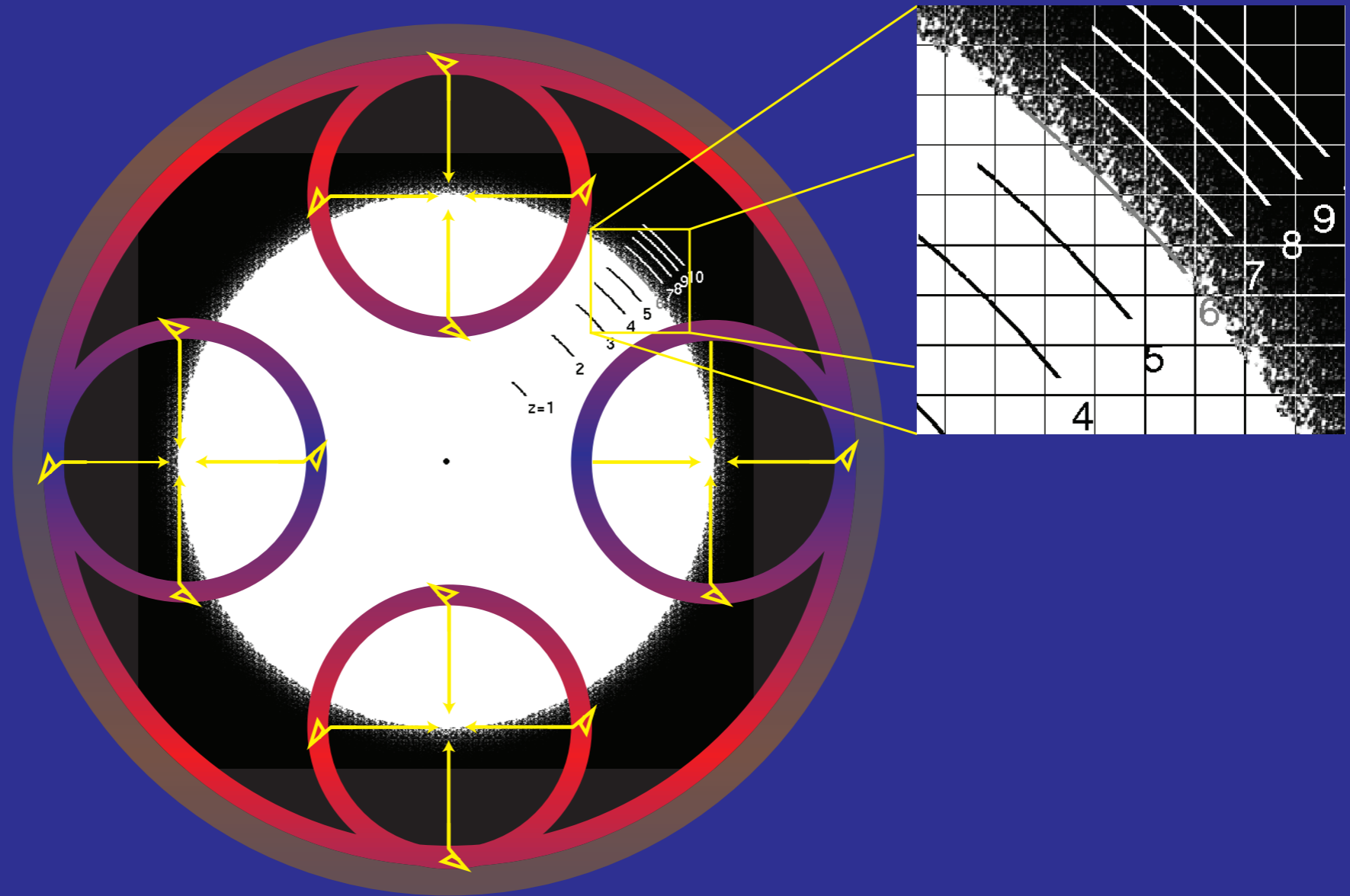
Spectra

$$(2\ell + 1)^2 C_\ell^{XY(m)} = \frac{n_m}{8\pi} \int k^2 dk X_\ell^{(m)} Y_\ell^{(m)*}$$



Inhomogeneous Ionization

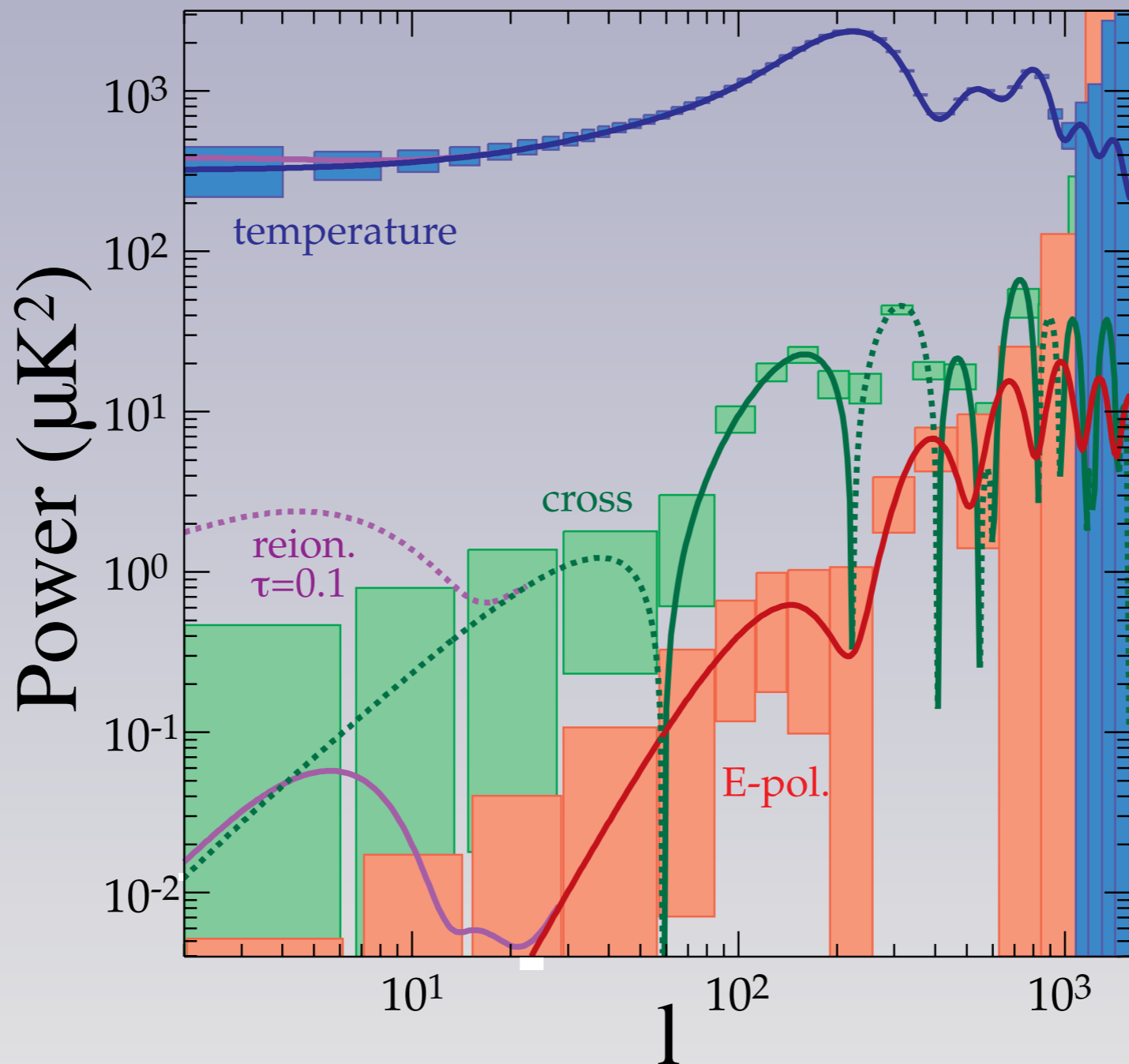
- As **reionization** completes, **ionization** regions **grow** and fill the space



Zahn et al. (2006) [Mortonson et al (2009)]

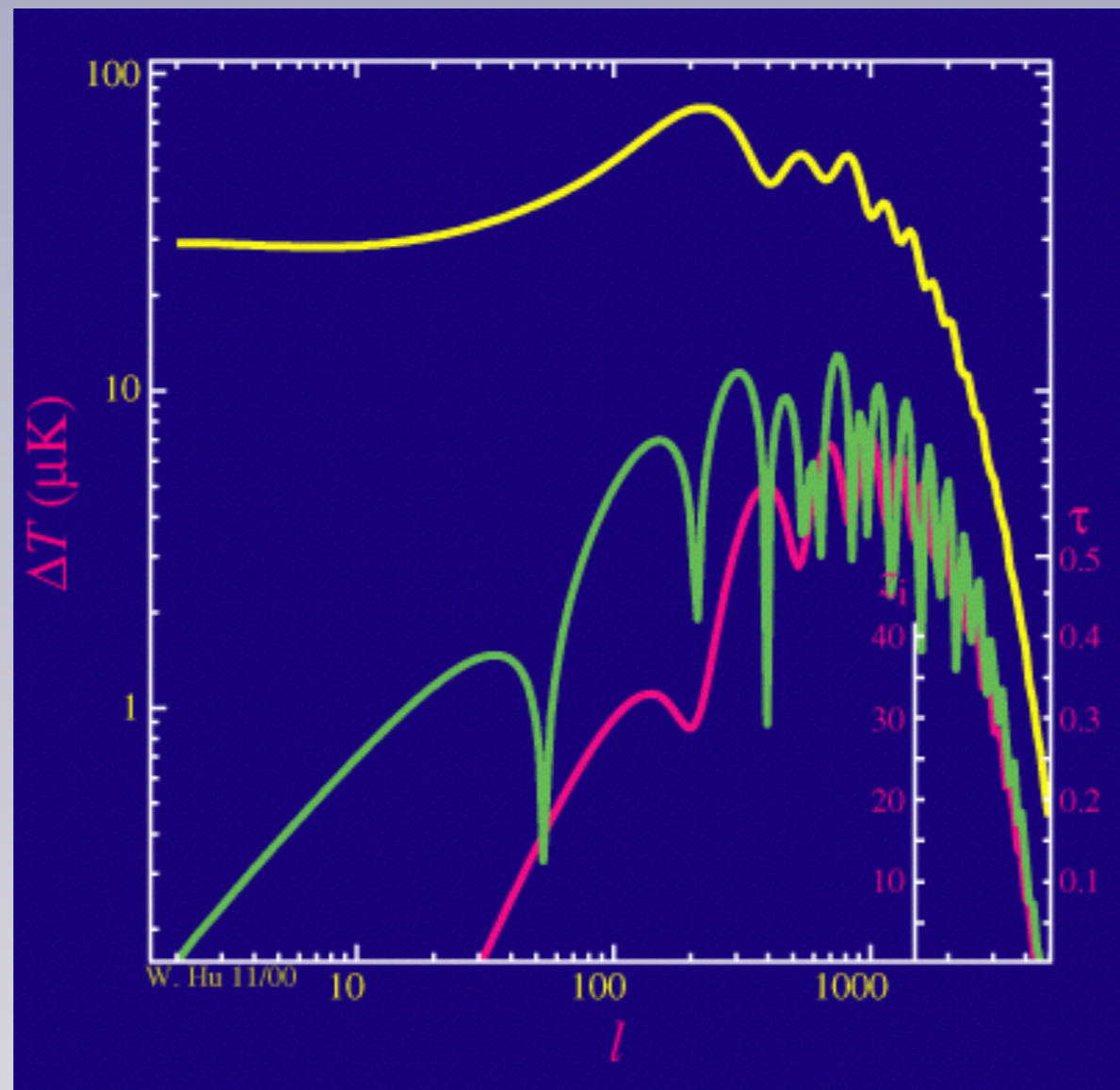
Spectra

$$(2\ell + 1)^2 C_\ell^{XY(m)} = \frac{n_m}{8\pi} \int k^2 dk X_\ell^{(m)} Y_\ell^{(m)*}$$



Secondary anisotropy

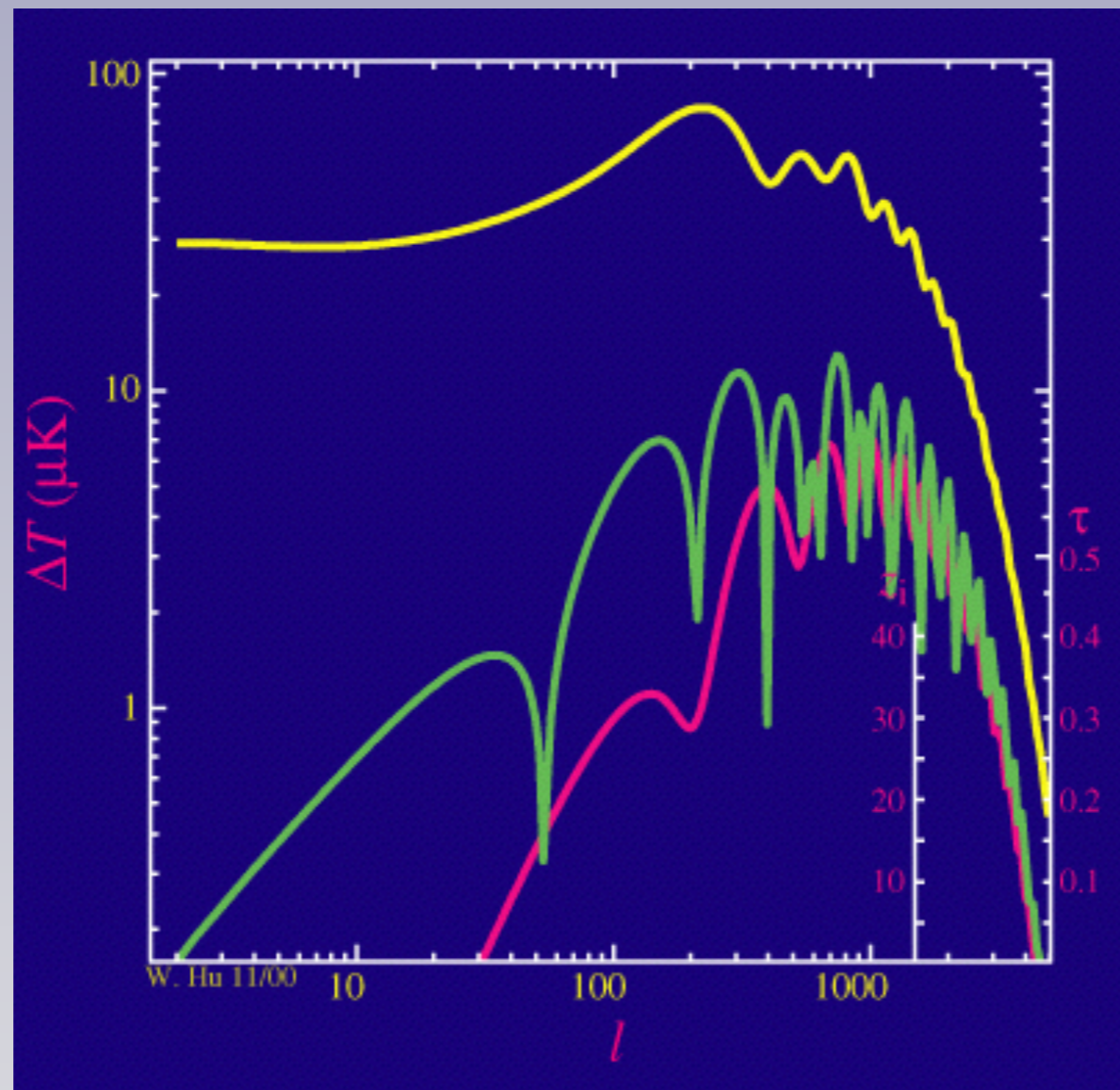
Reionization



from Wayne Hu [<http://background.uchicago.edu/~whu/animbut/anim4.html>]

Secondary anisotropy

Reionization

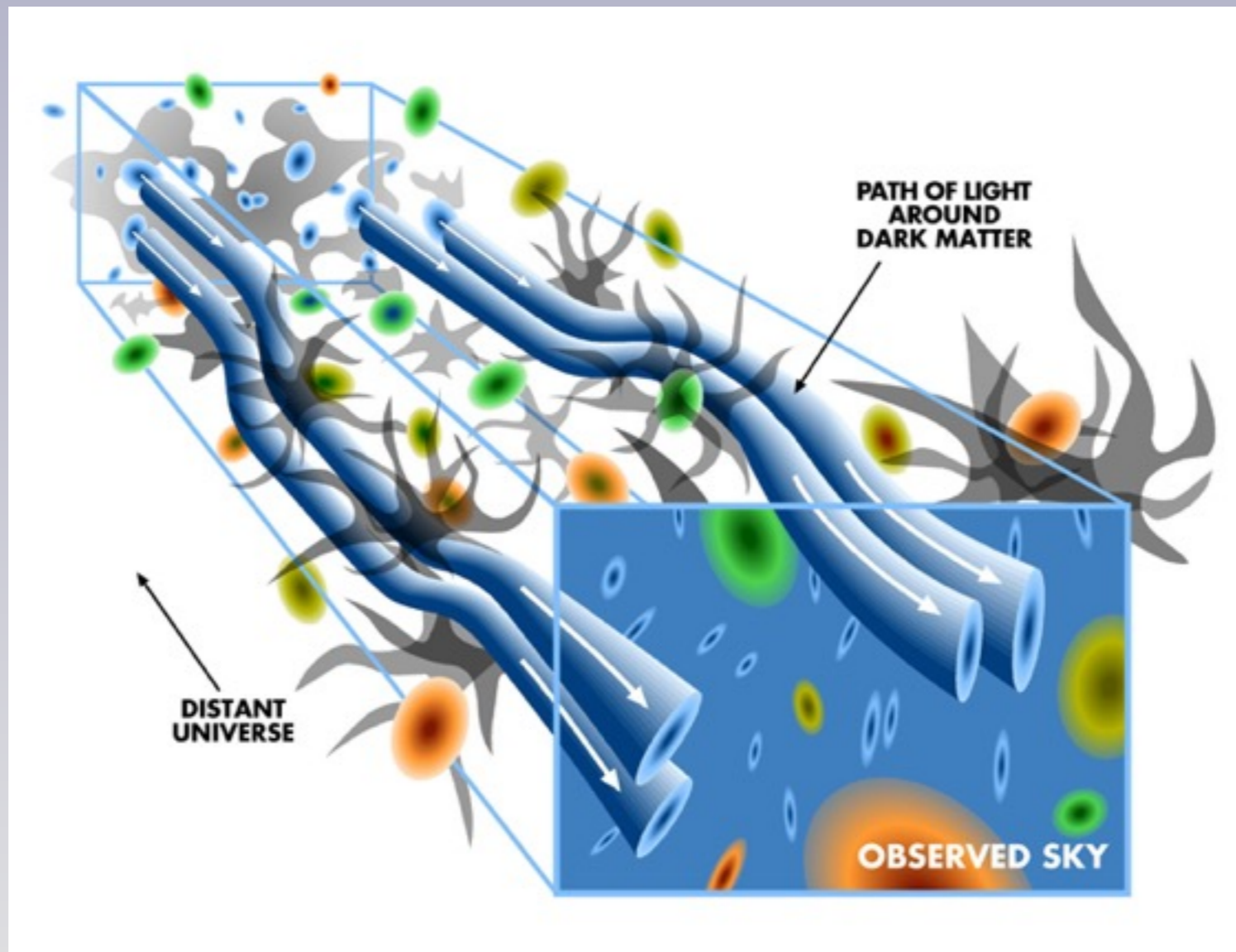


from Wayne Hu [<http://background.uchicago.edu/~whu/animbut/anim4.html>]

Reminder from lecture III

Weak lensing

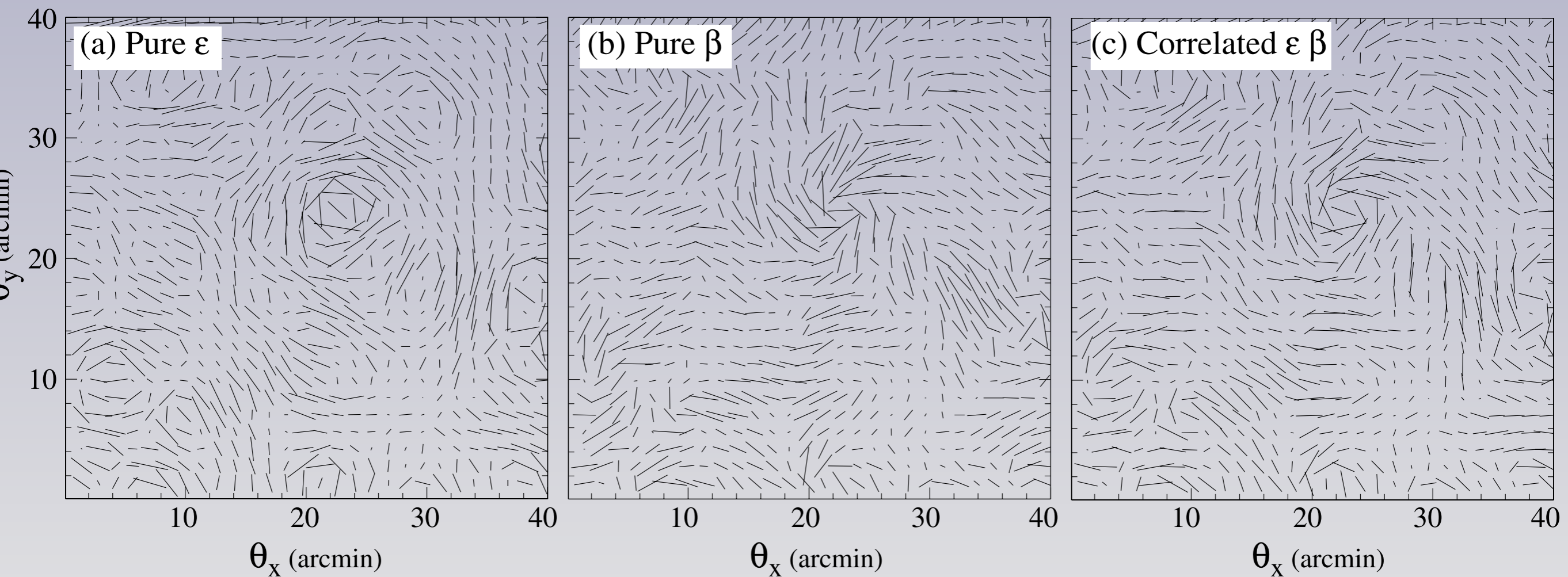
- Probe the potential, not the galaxies
- LSST, Euclid, ...



[http://lsst.org/lst/science/scientist_cosmic_shear]

Reminder from lecture III

Weak lensing



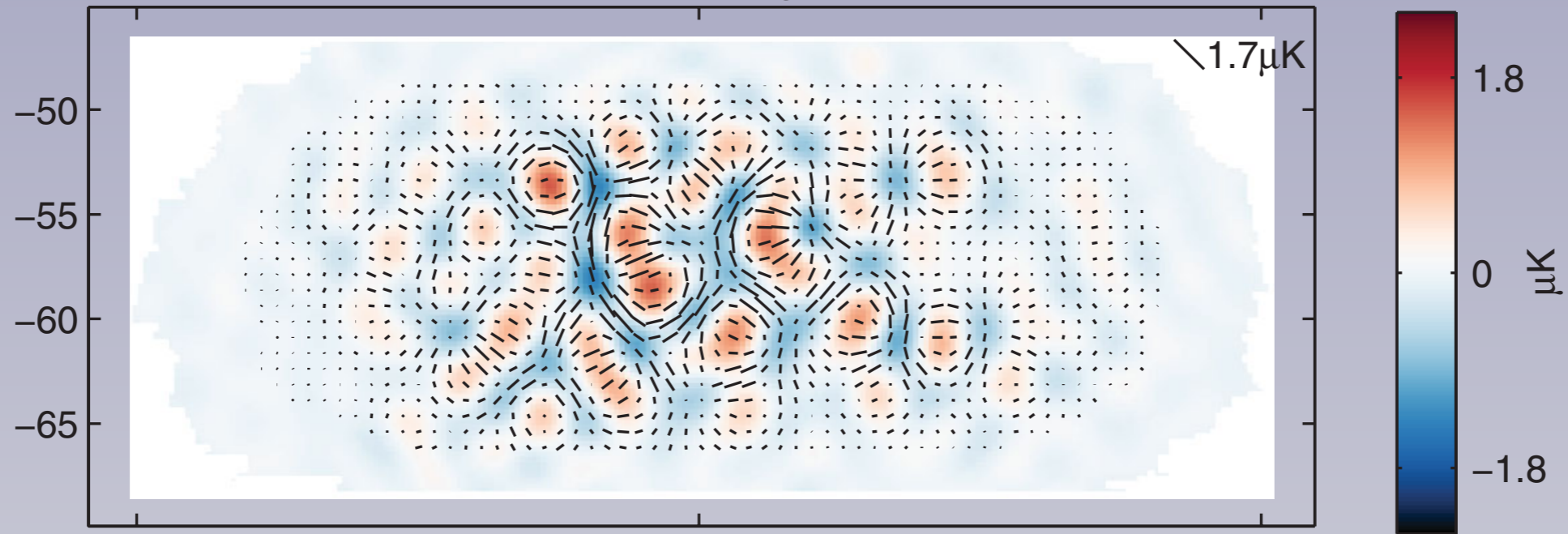
[Hu & White, 2001]

CMB lensing

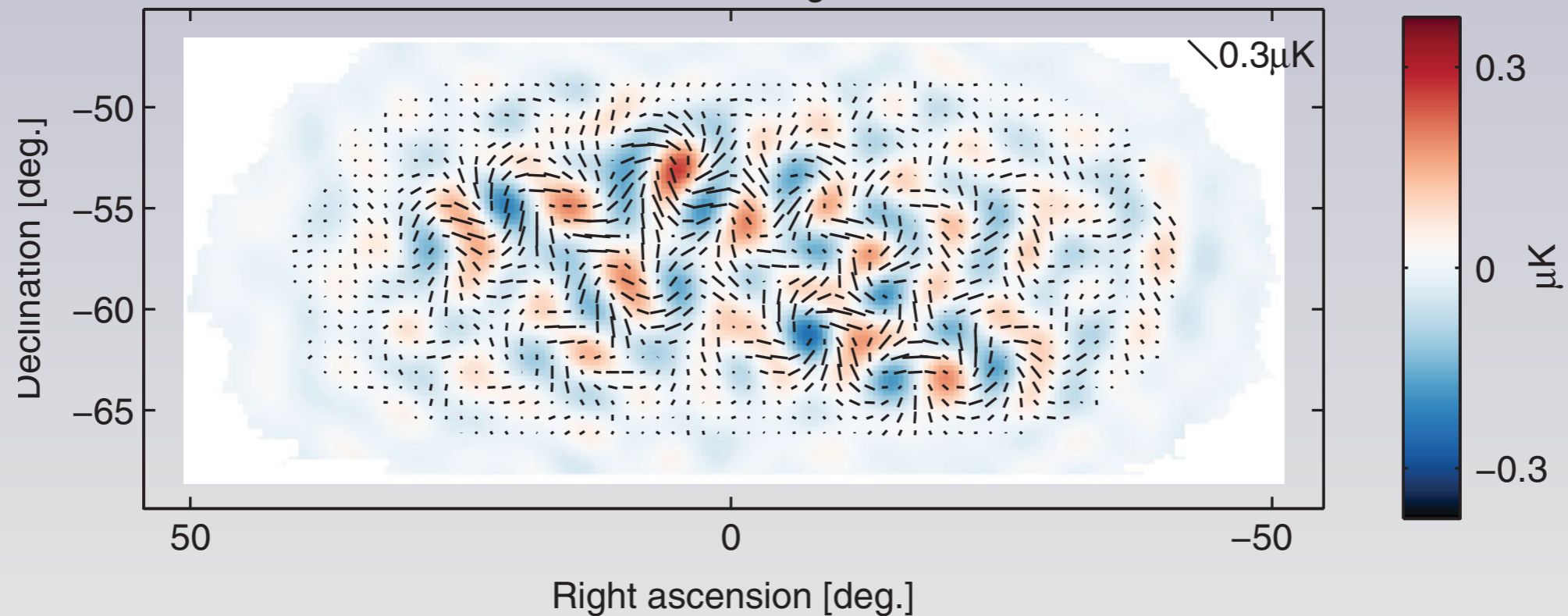
- Background of E-mode polarized light (from mostly scalar perturbations)
- Travelling through the gravitational potential (a scalar).
- Scalars only?
- Scalar + scalar = nonlinear
 - B-modes generated!

CMB from BICEP2

BICEP2: E signal



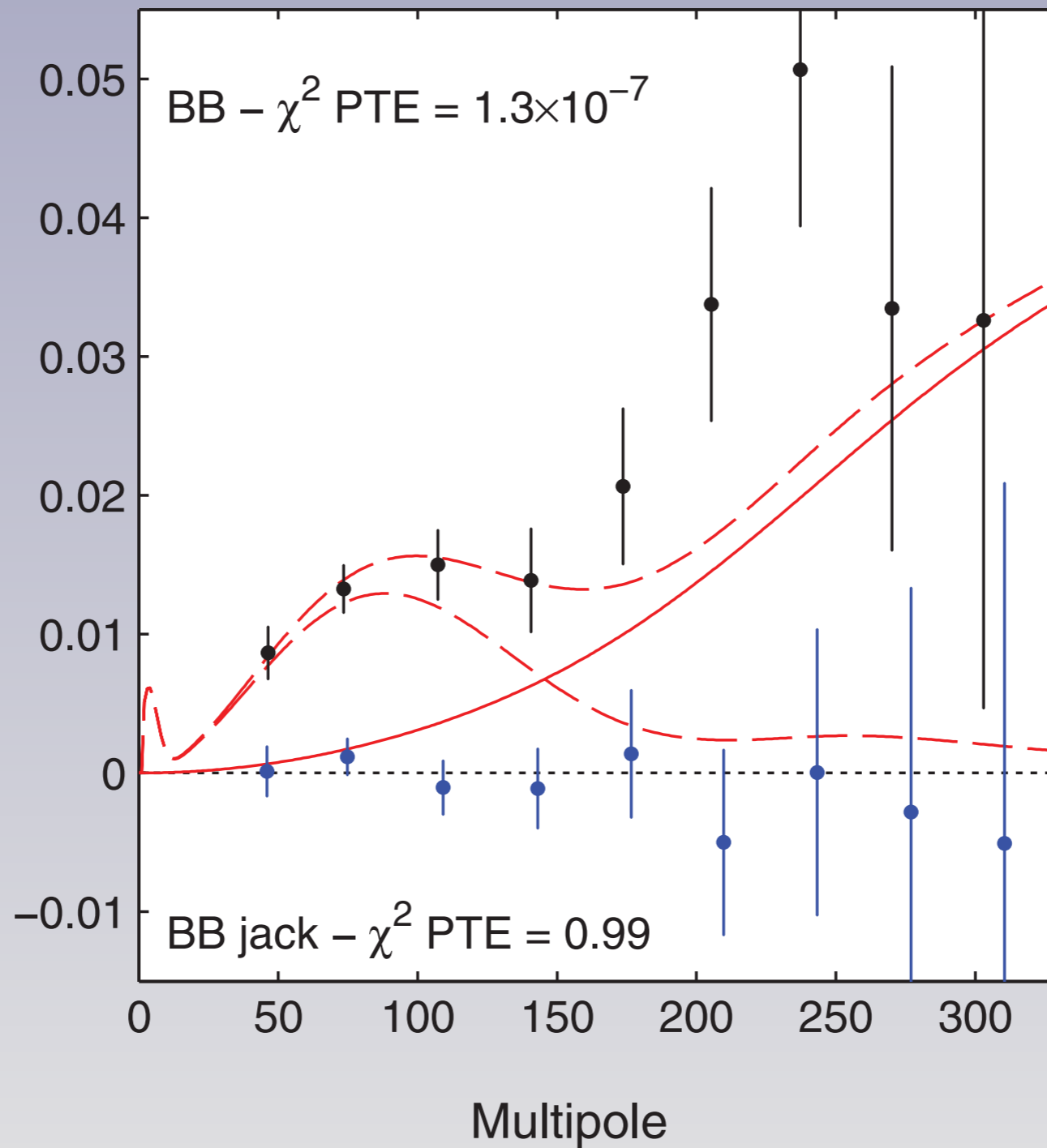
BICEP2: B signal



[BICEP2, March 2014]

CMB from BICEP2

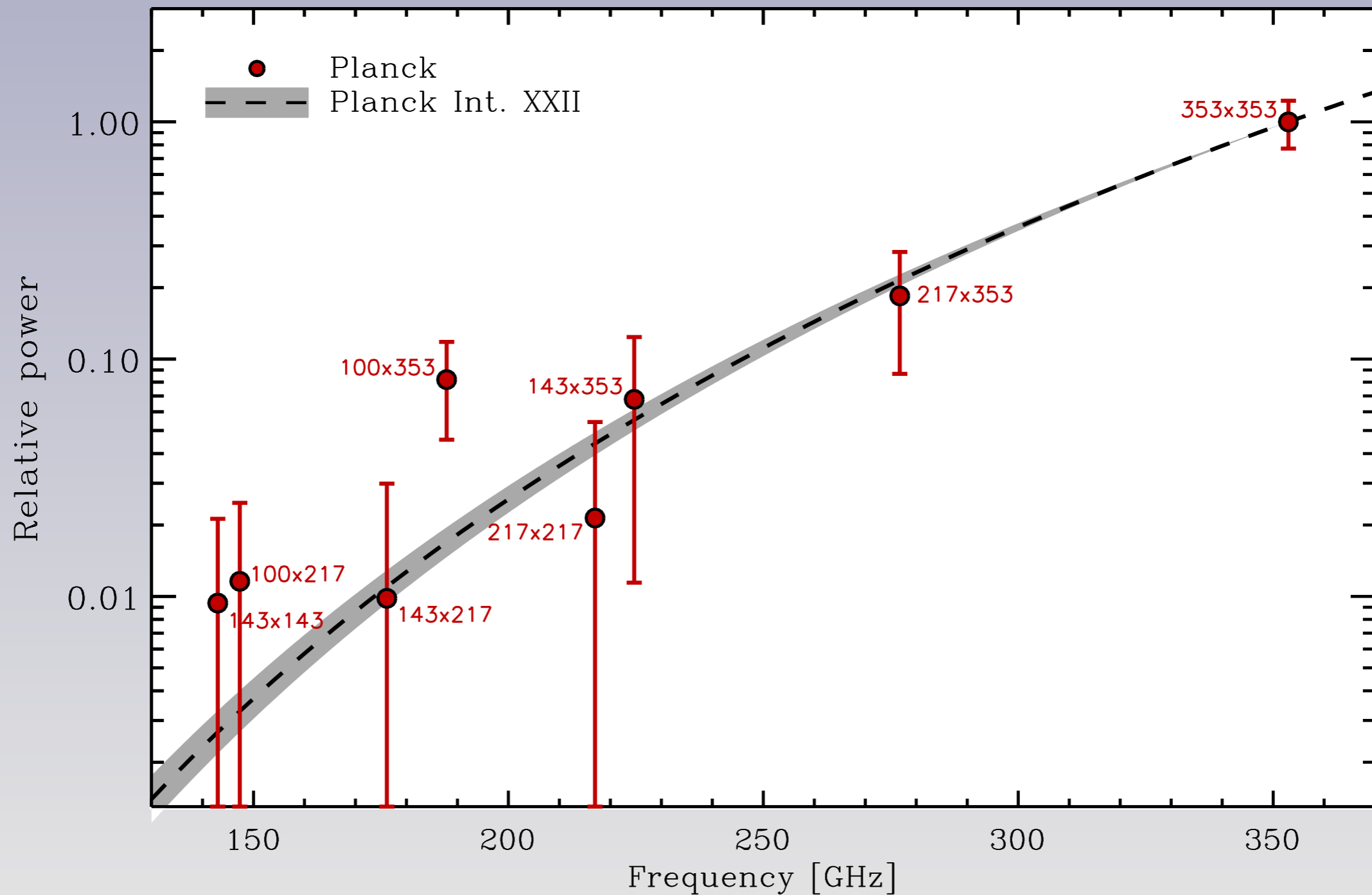
[BICEP2, March 2014]



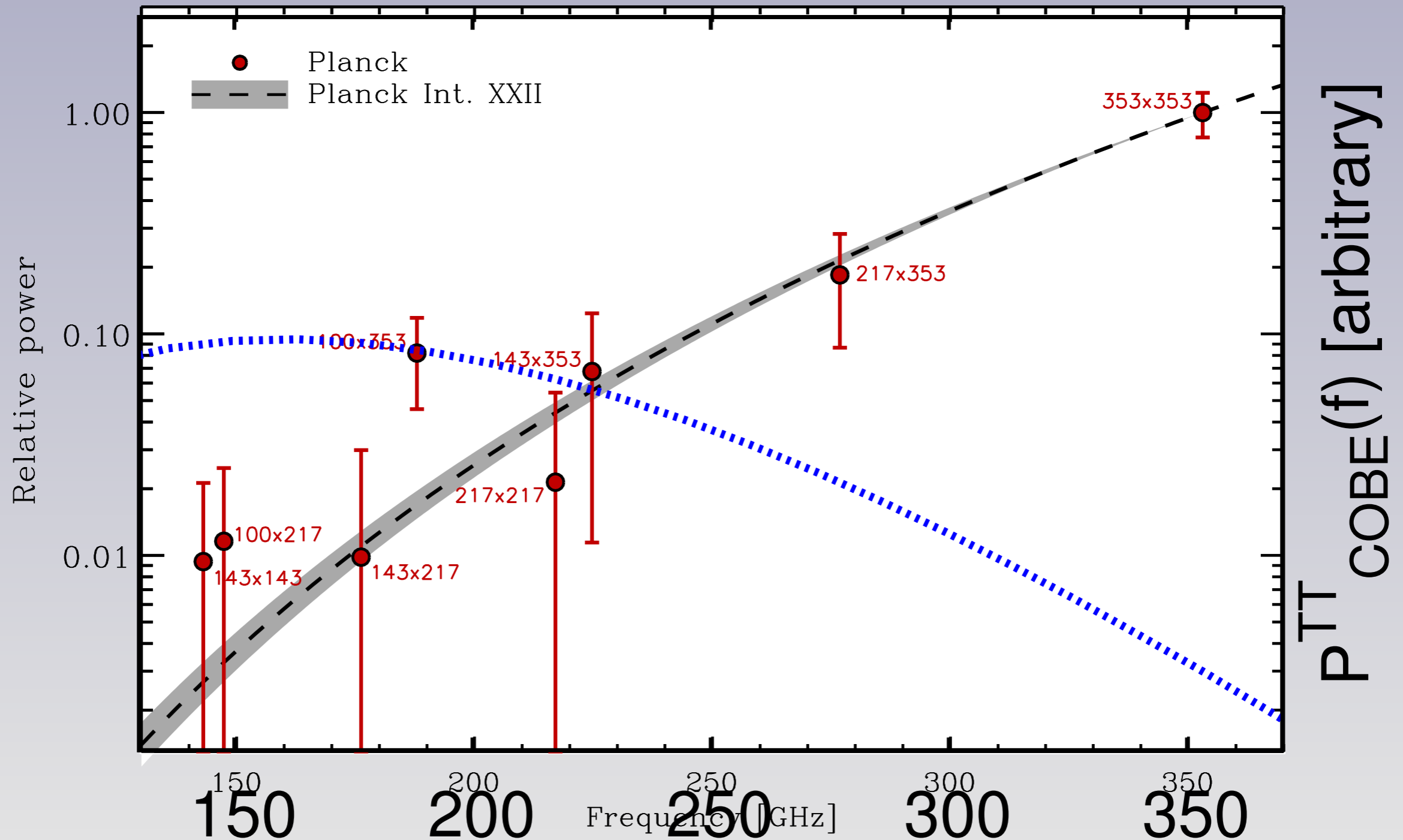
Spinning dust

- Spinning dust generates a magnetic field
- Will align with any present large scale magnetic field (e.g. galaxy)
- Hence an orientation (a vector field) enters the dynamics
- E- and B-modes sourced
- No “standard model of dust emission”

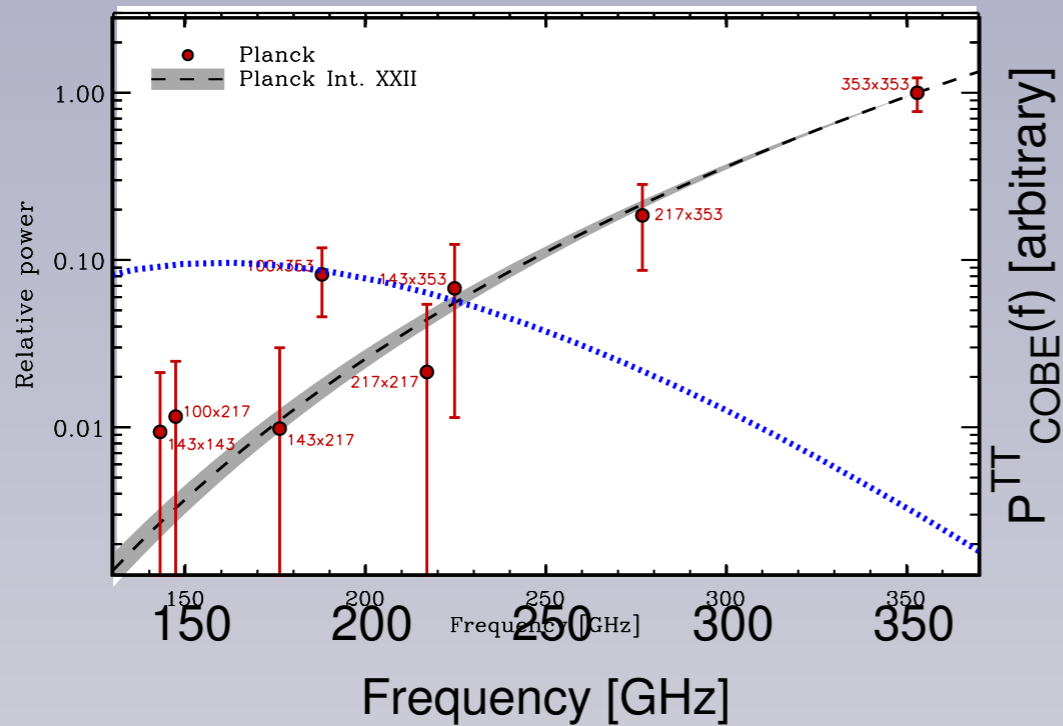
Planck measurements of dust, paper XXX



Planck measurements of dust, paper XXX



Planck measurements of dust, paper XXX



$$C_\ell \propto \ell^\alpha$$

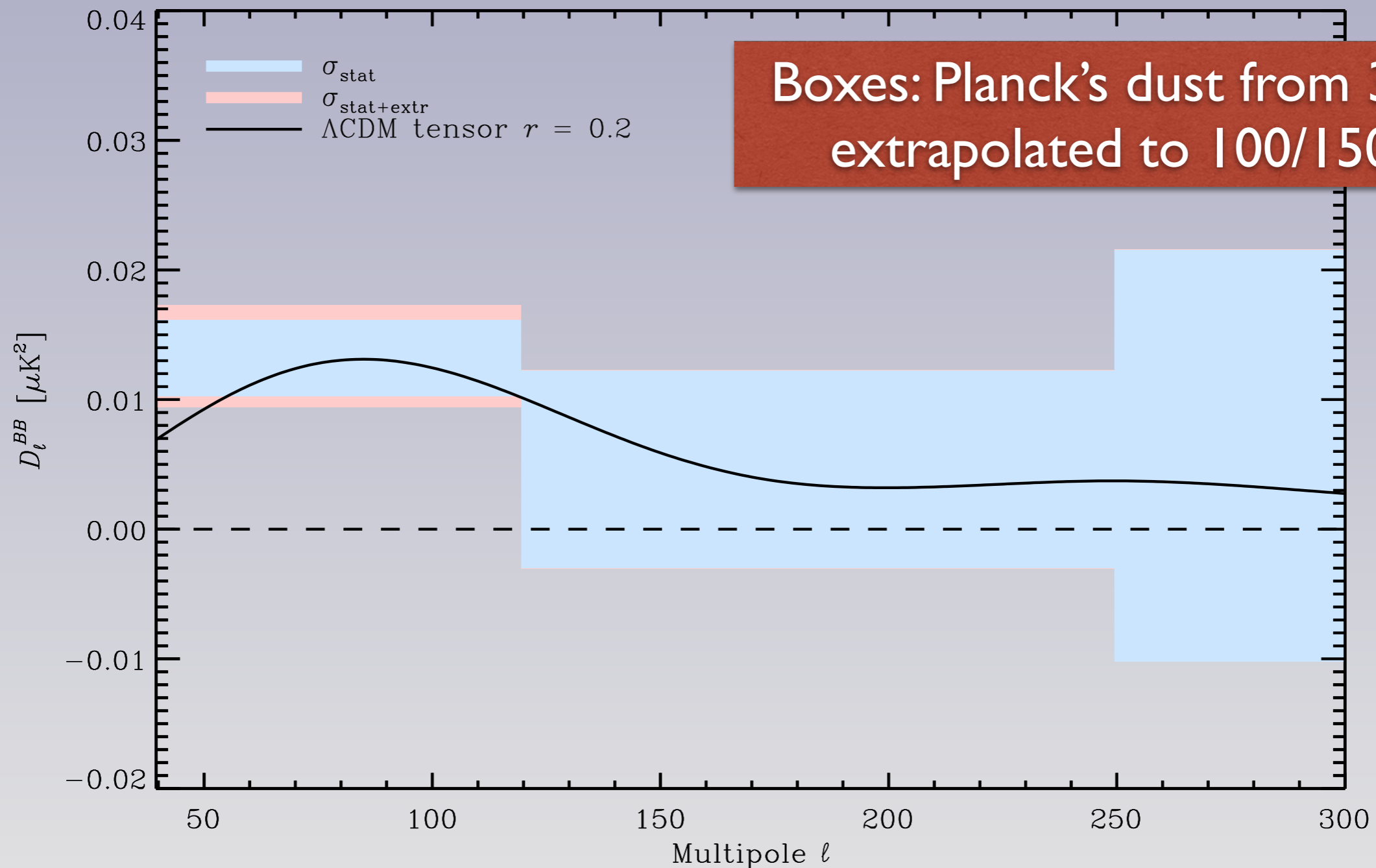
$$I_d(\nu) \propto \nu^{\beta_d} B_\nu(T_d),$$

$$\beta_d = 1.59$$

$$T_d = 19.6 \text{ K},$$

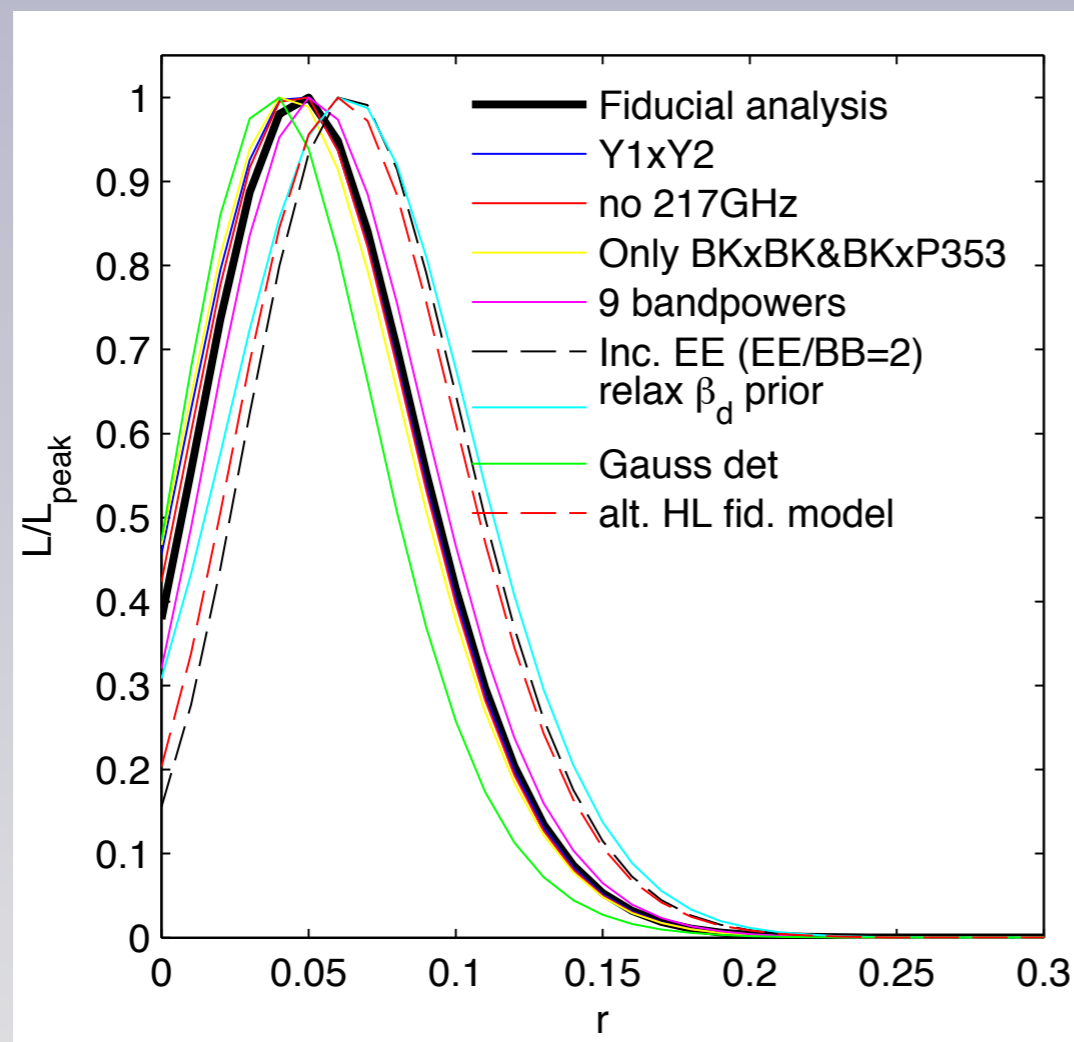
B_ν is the Planck function.

Planck measurements of dust, paper XXX



Last Friday's BICEP2+Keck +Planck paper

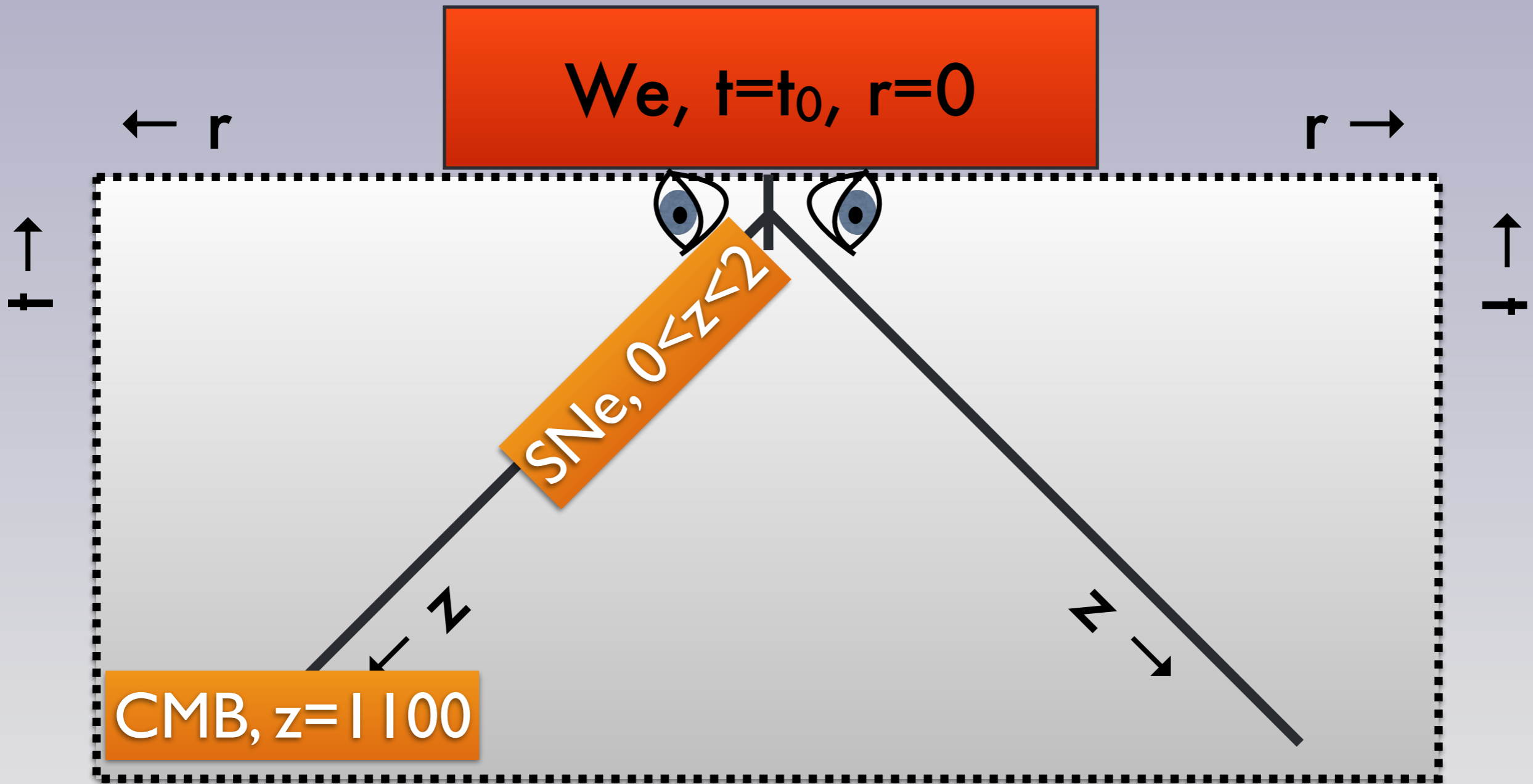
- Conclusion about dust foreground seems robust



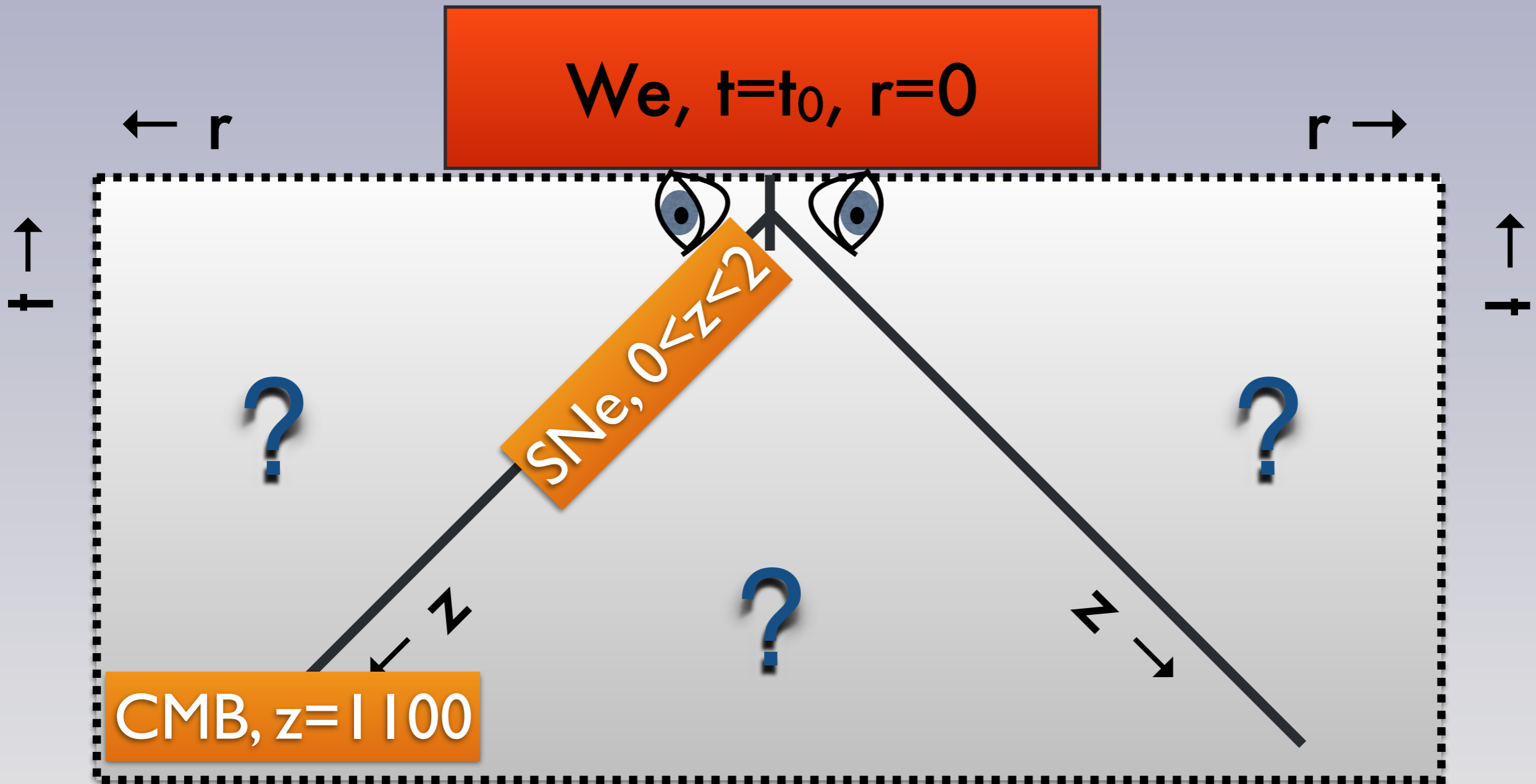
upper limit $r < 0.12$ at 95% confidence

Summary of Cosmology

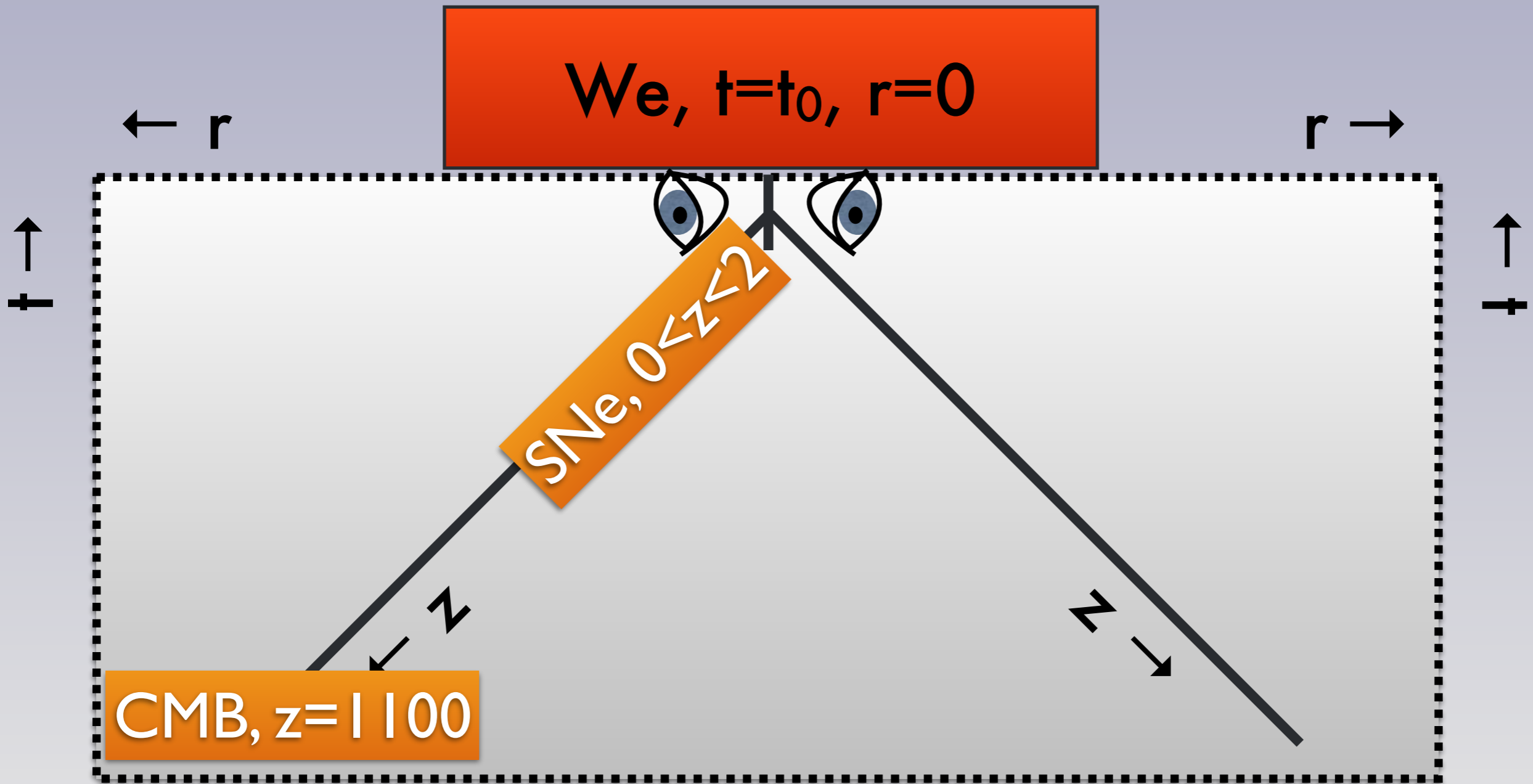
Our light cone



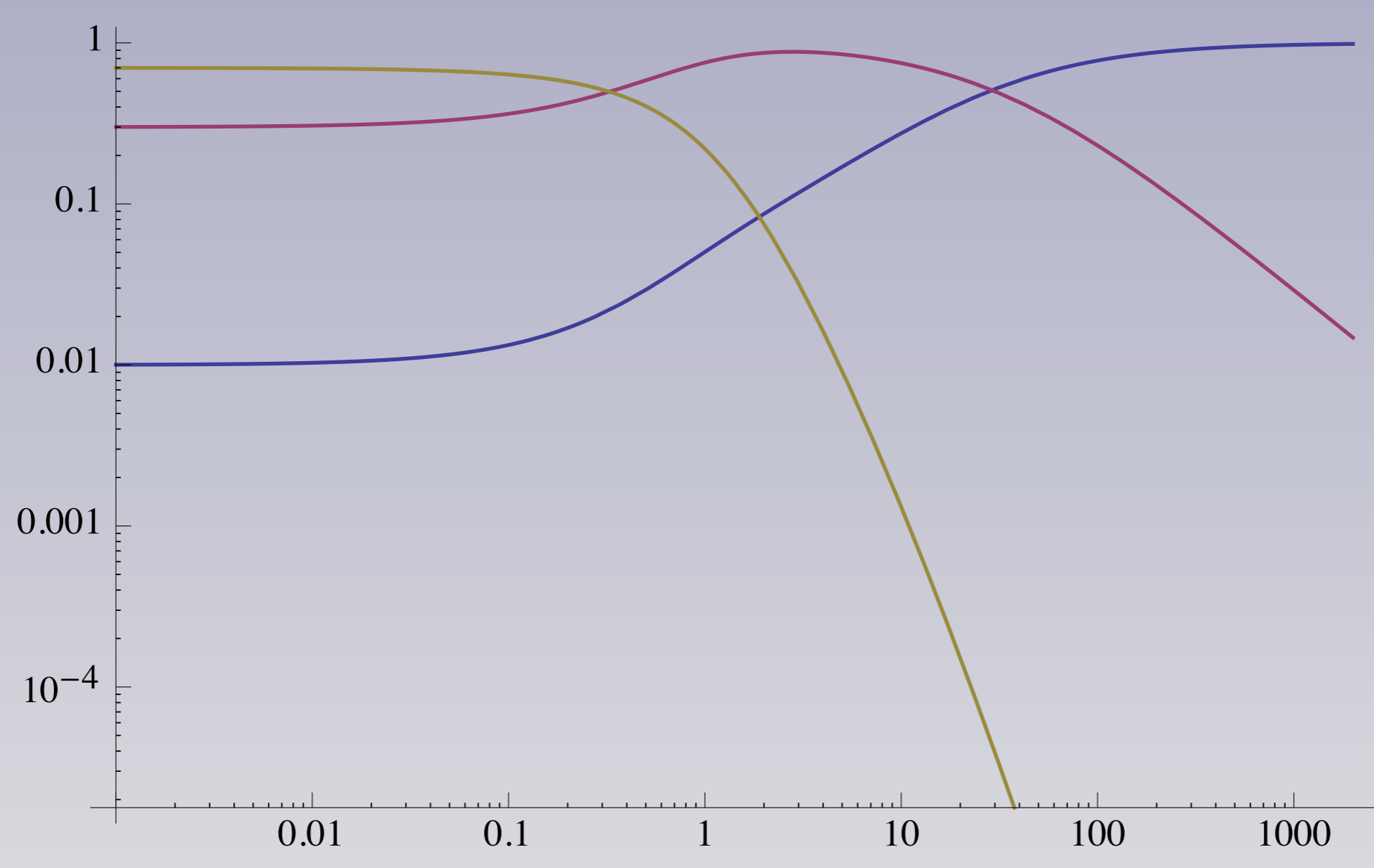
Our light cone



Our light cone

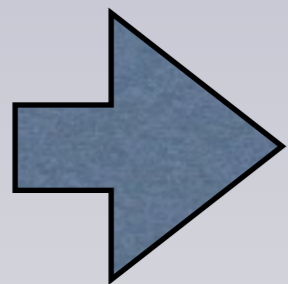


DE
DM
Rad.

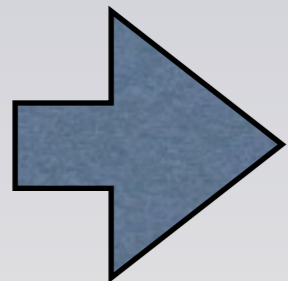


Friedmann-Lemaître metric

- If we are not in a special position
- Universe must be everywhere similar to here*
- First approximation: no space dependence, only time.
 - Invariant under rotations and translations

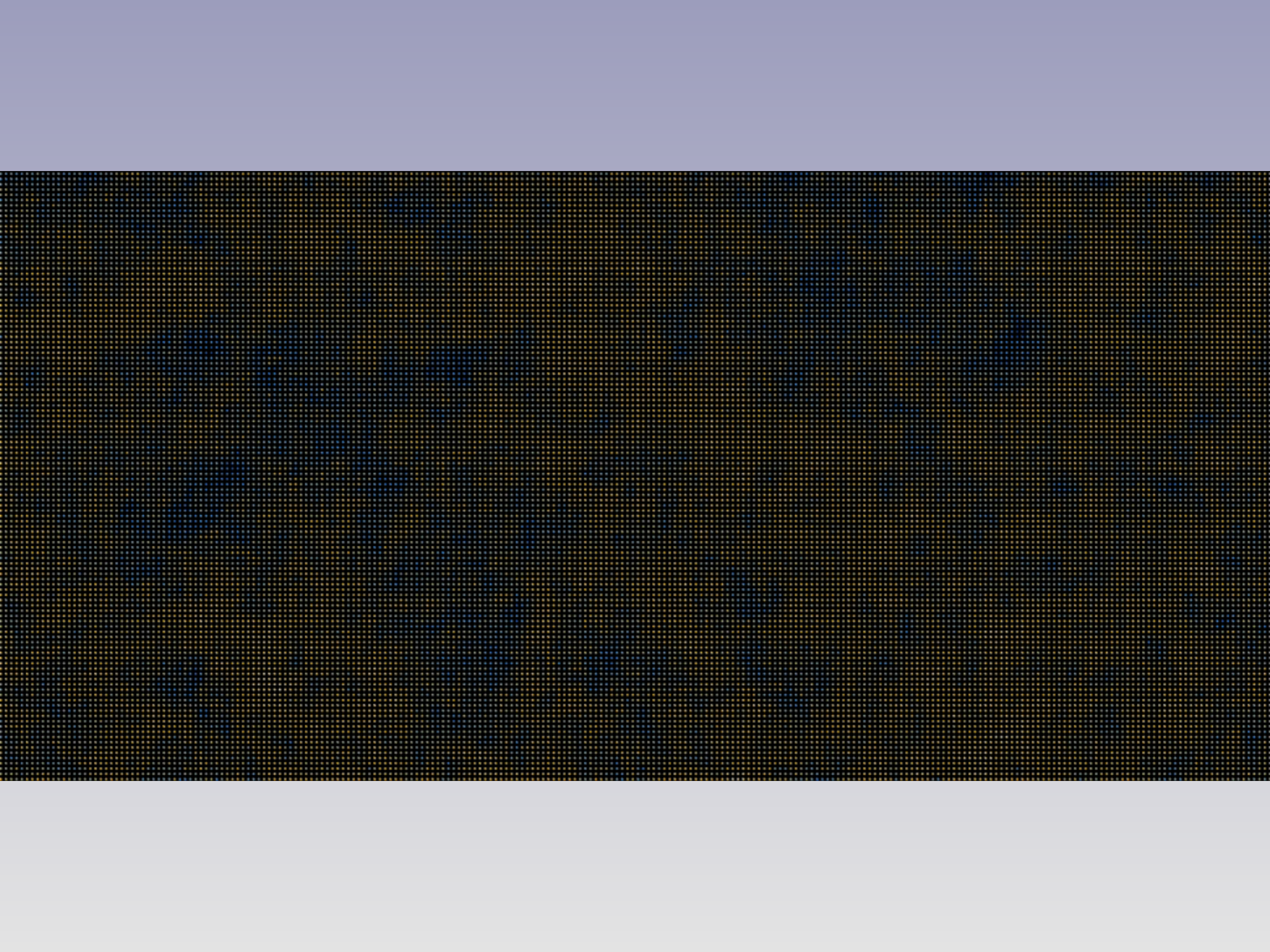


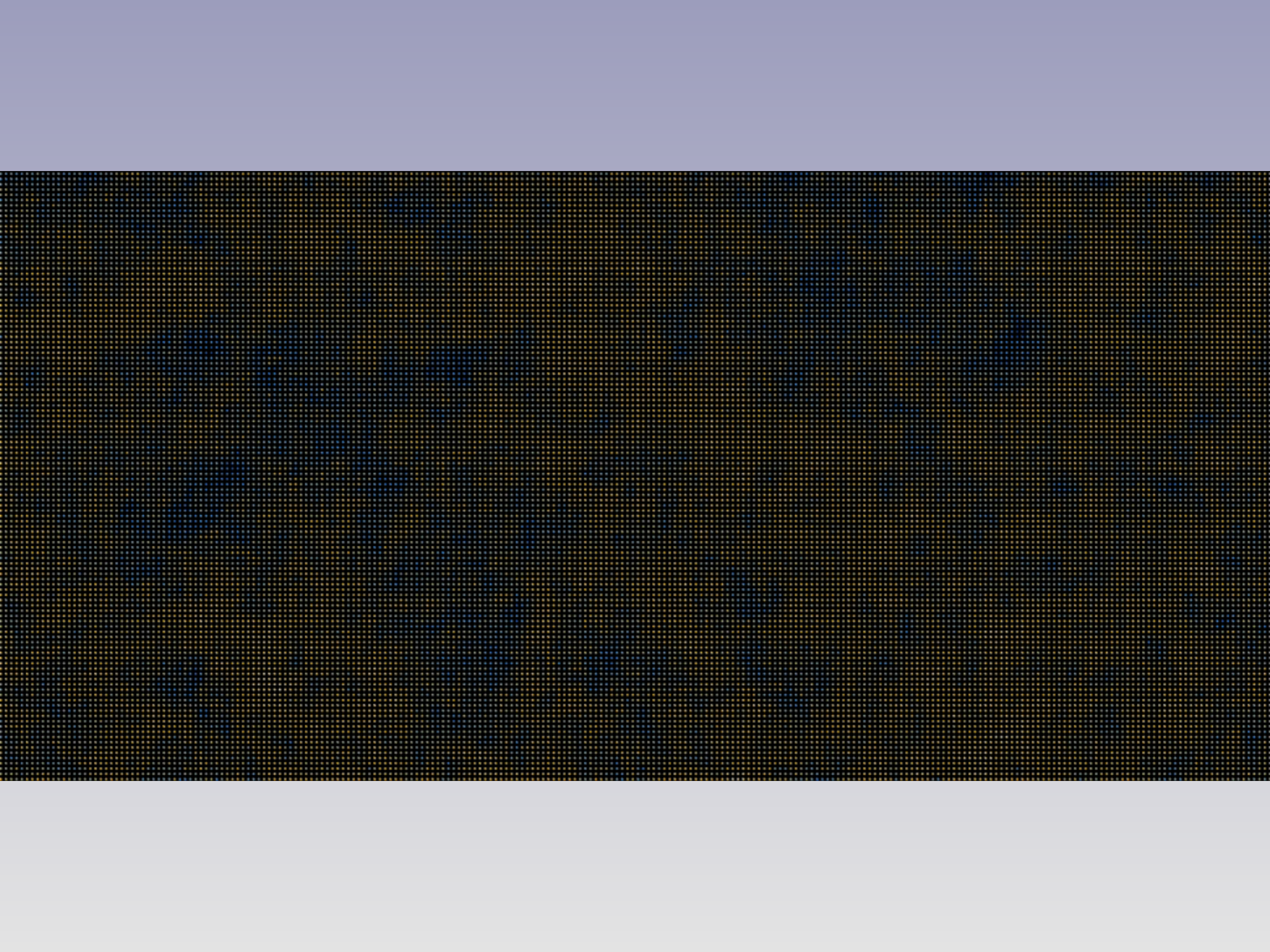
$$\left(\frac{\dot{a}(t)}{a(t)}\right)^2 = \frac{8\pi G_N}{3}\rho + \frac{\Lambda}{3} - \frac{k^2}{a(t)^2}$$



Big Bang model

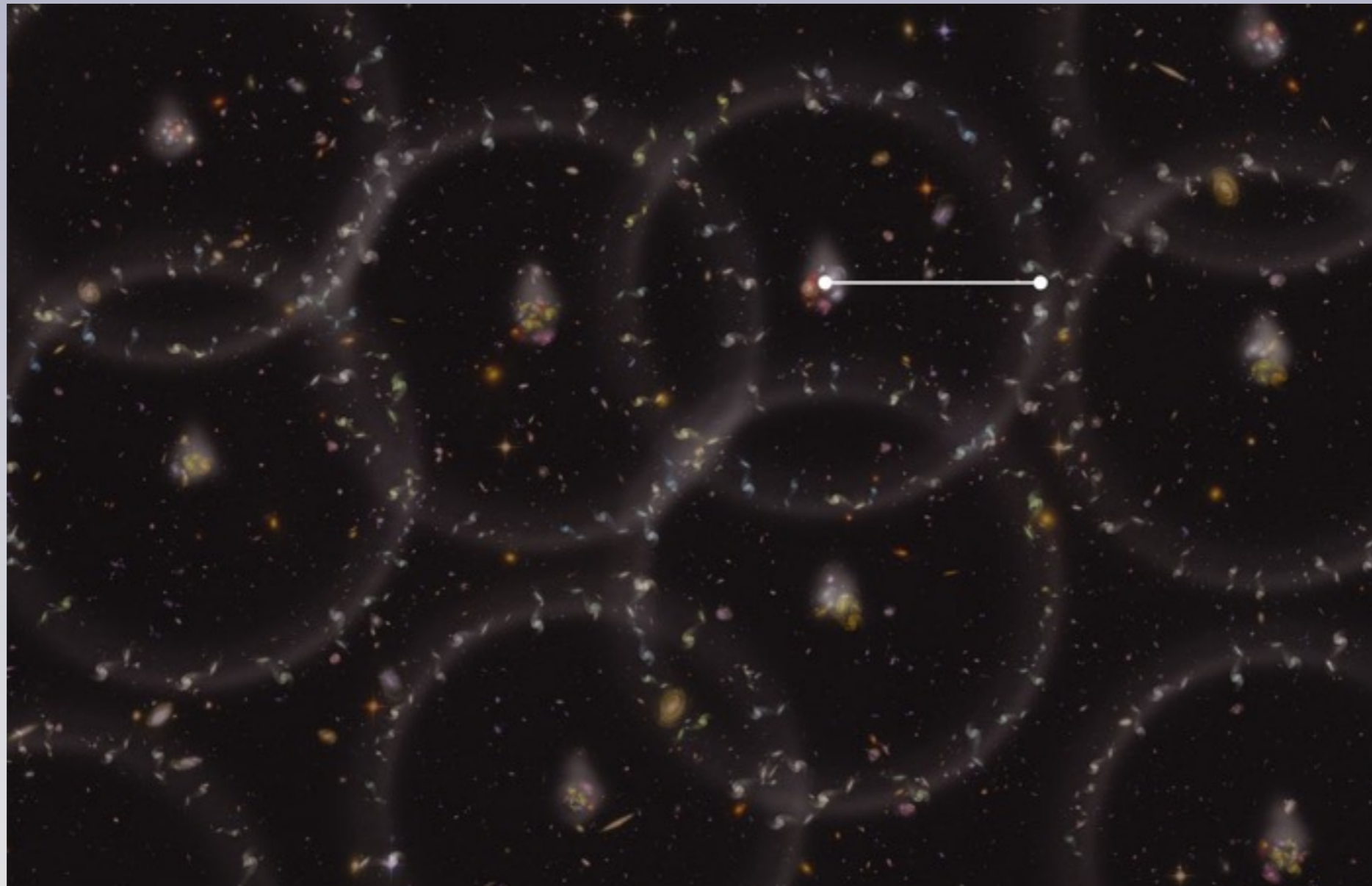
* realize how revolutionary that idea was (still is)



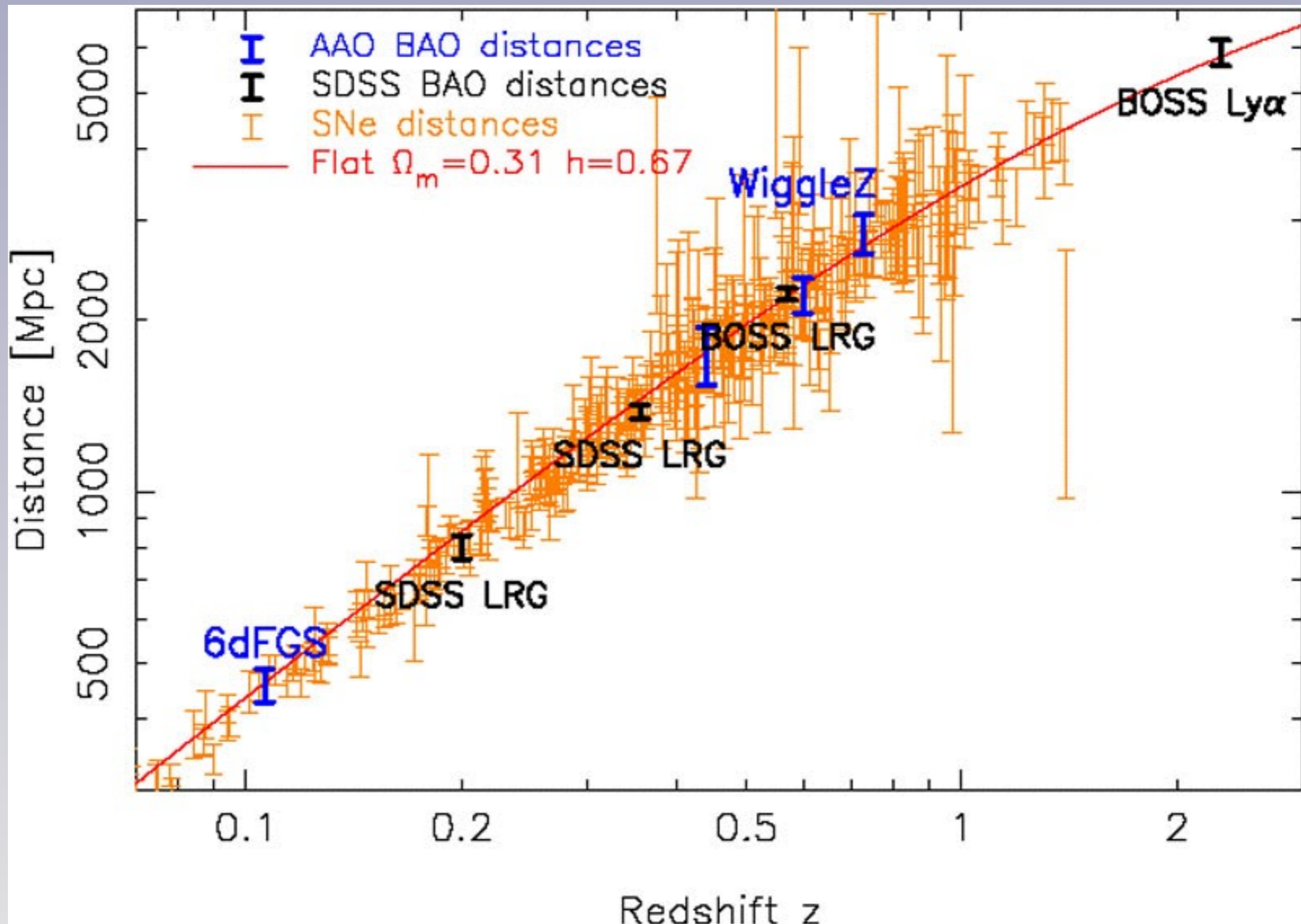


Baryon Acoustic Oscillations

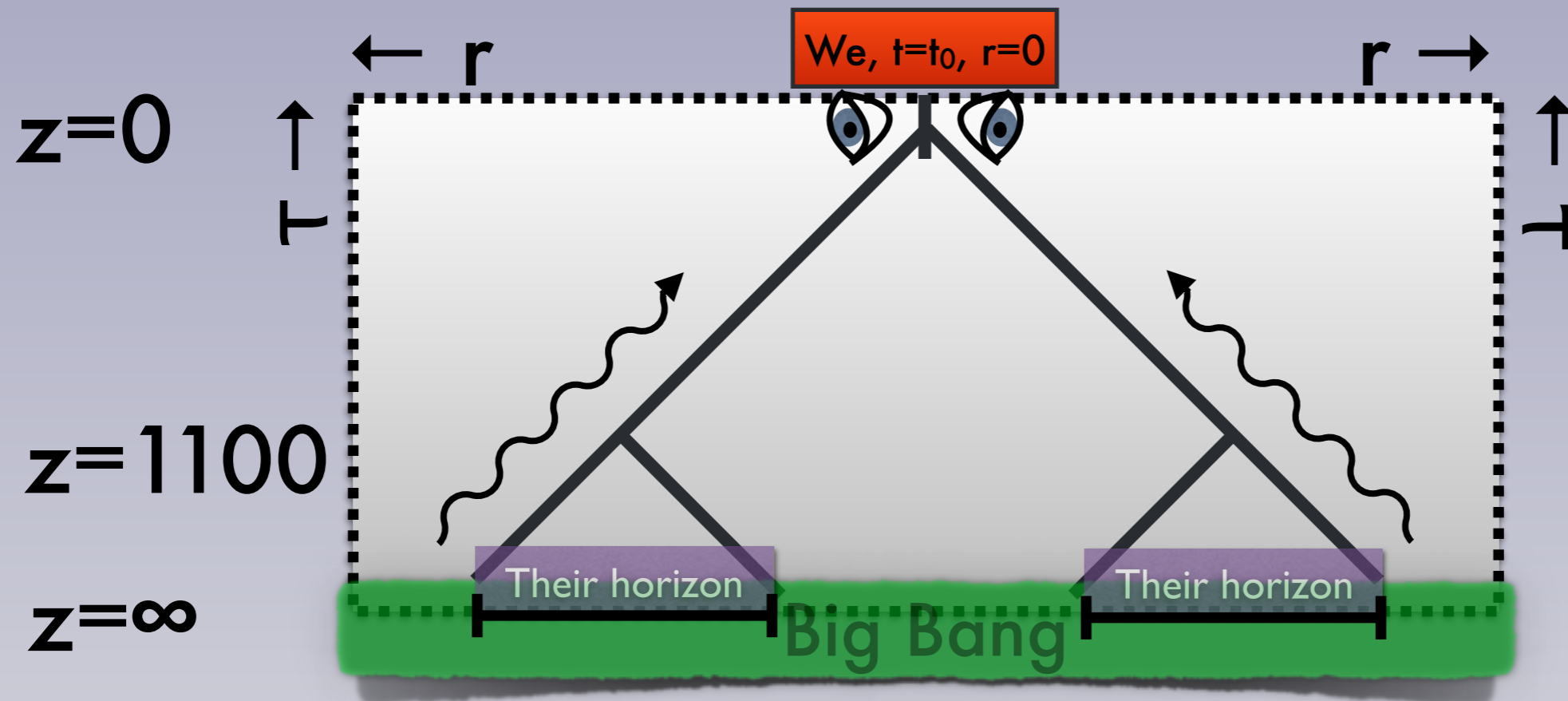
- Same peaks as in CMB



BAO as a distance measure

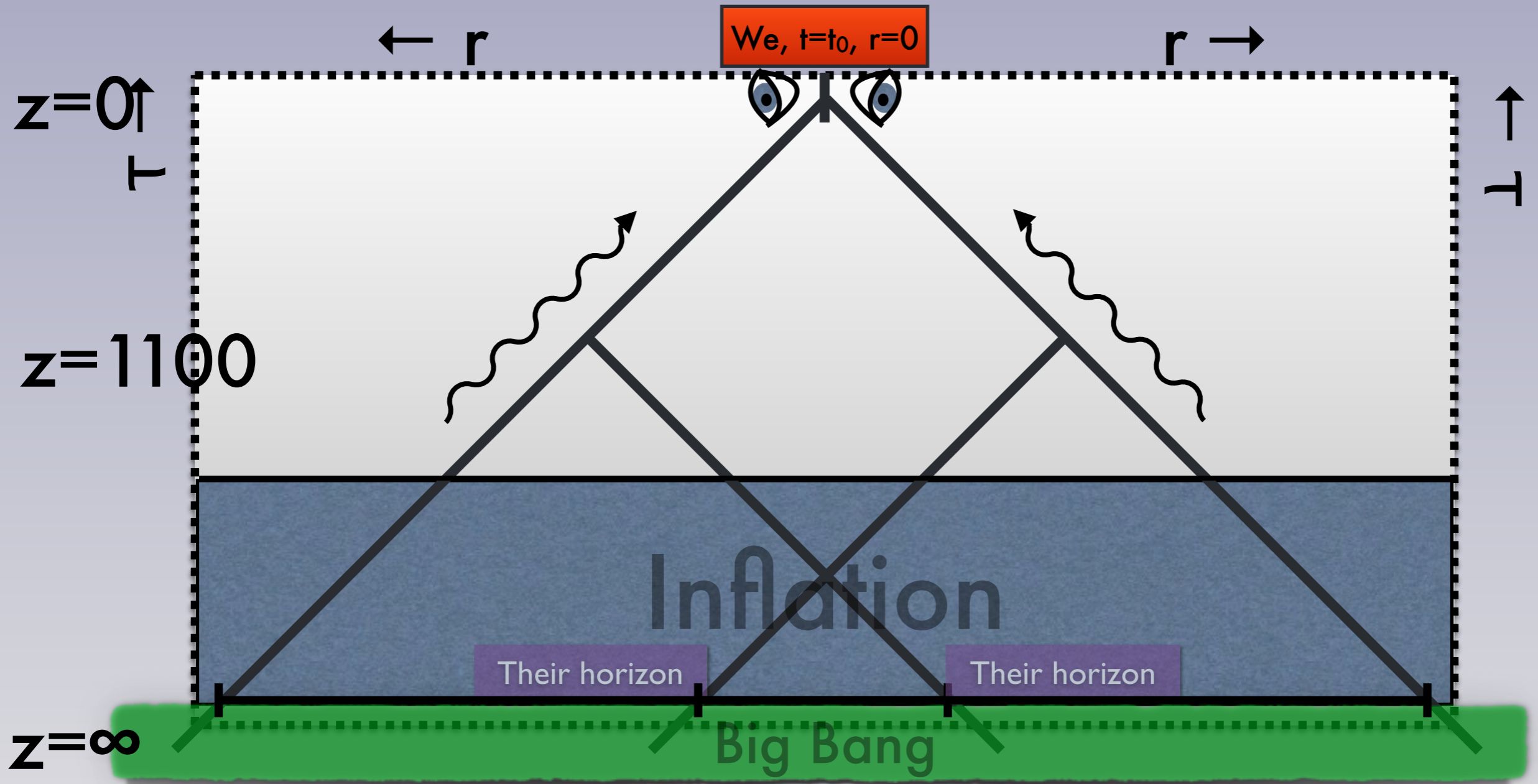


Horizon problem solved



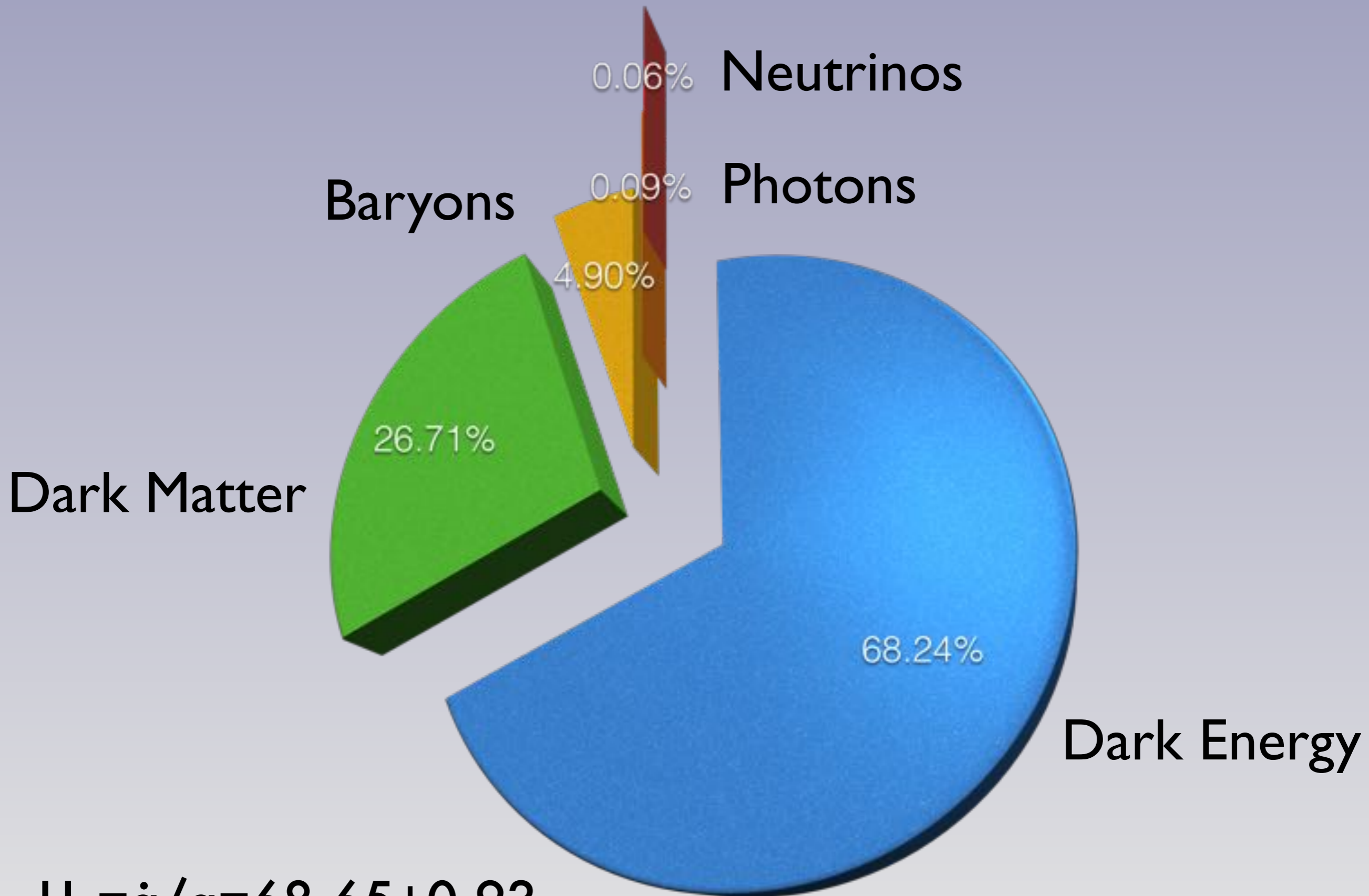
[following Will Kinney's lecture notes [arXiv:0902.1529](https://arxiv.org/abs/0902.1529)]

Horizon problem solved



[following Will Kinney's lecture notes arXiv:0902.1529]

Parameter constraints from Planck



$$H_0 = \dot{a}/a = 68.65 \pm 0.93$$