



Quantum Optics and Quantum Communications using Gaussian and Non-Gaussian States of the Light

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Content of the Talk



Part 1 : Gaussian and non-Gaussian states

1. Homodyne detection and quantum tomography
2. Generating non-Gaussian Wigner functions : kittens, cats and beyond

Part 2 : Continuous variable quantum cryptography (Gaussian !)

1. Continuous variable quantum cryptography : principles
2. Continuous variable quantum cryptography : implementations

Part 3 : Towards quantum networks (non-Gaussian !)

1. Entanglement for continuous variable quantum networks
2. Teleportation of Schrödinger's cats
3. Storing non-Gaussian states : single photon quantum memory

Quantum description of light

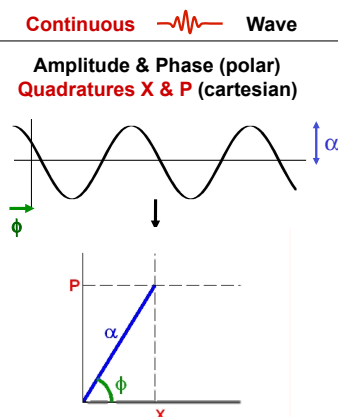
Discrete Photons

Continuous Wave

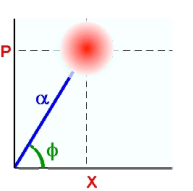
A single "mode" of the quantized electromagnetic field (a plane wave, or a "Fourier transform limited" pulse) is described as a quantized harmonic oscillator : operators $a, a^+, N = a^+ a$, etc...

Quantum description of light

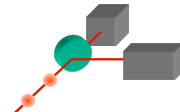
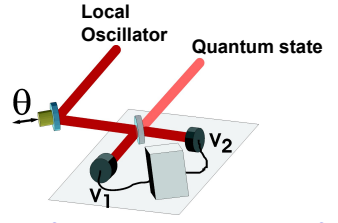
Parameters :
Discrete Photons
Number of photons n
Coherences $\langle n | \rho | m \rangle$



Quantum description of light

	Discrete Photons	Continuous Wave
Parameters :	Number & Coherence	Amplitude & Phase (polar) Quadratures X & P (cartesian)
Representation:	Density matrix	Wigner function W(X,P)
	$\rho = \begin{bmatrix} \rho_{0,0} & \rho_{0,1} & \rho_{0,2} & \dots \\ \rho_{1,0} & \rho_{1,1} & \rho_{1,2} & \dots \\ \rho_{2,0} & \rho_{2,1} & \rho_{2,2} & \dots \\ \dots & \dots & \dots & \dots \end{bmatrix}$	 $X = (a + a^\dagger)/\sqrt{2}$ $P = (a - a^\dagger)/i\sqrt{2}$ $[X, P] = i$ <p style="color: red;">Heisenberg : $\Delta X \cdot \Delta P \geq 1/2$ measurement of both X and P measurement of $X_\theta = X \cos\theta + P \sin\theta$</p>

Quantum description of light

	Discrete Photons	Continuous Wave
Parameters :	Number & Coherence	Amplitude & Phase (polar) Quadratures X & P (cartesian)
Representation:	Density matrix	Wigner function W(X,P)
Measurement :	Counting : APD, VLPC, TES... 	Demodulation : Homodyne detection 
		<p>Interference, then subtraction of photocurrents :</p> $V_1 - V_2 \propto E_{OL} E_{EQ}(\theta)$ $\propto X_\theta = X \cos\theta + P \sin\theta$

Homodyne detection

$$I_1 = |E_{LO} + E_S|^2 / 2 = \{ |E_{LO}|^2 + |E_S|^2 + |E_{LO}| (E_S e^{-i\theta_{LO}} + E_S^* e^{i\theta_{LO}}) \} / 2$$

$$I_2 = |E_{LO} - E_S|^2 / 2 = \{ |E_{LO}|^2 + |E_S|^2 - |E_{LO}| (E_S e^{-i\theta_{LO}} + E_S^* e^{i\theta_{LO}}) \} / 2$$

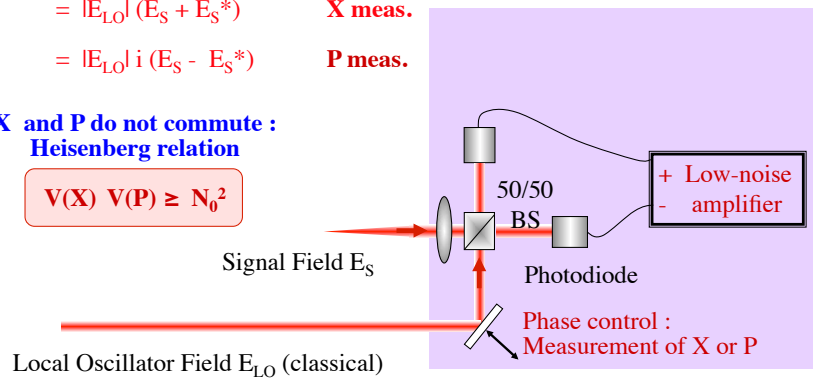
$$I_1 - I_2 = |E_{LO}| (E_S e^{-i\theta_{LO}} + E_S^* e^{i\theta_{LO}})$$

$$= |E_{LO}| (E_S + E_S^*) \quad \mathbf{X \text{ meas.}}$$

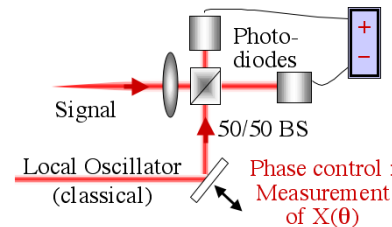
$$= |E_{LO}| i (E_S - E_S^*) \quad \mathbf{P \text{ meas.}}$$

**X and P do not commute :
Heisenberg relation**

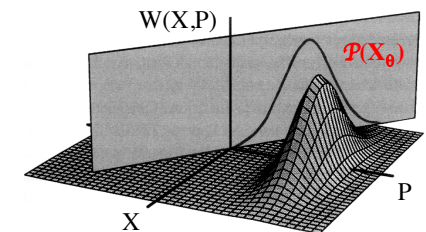
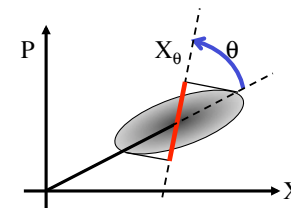
$$V(X) V(P) \geq N_0^2$$



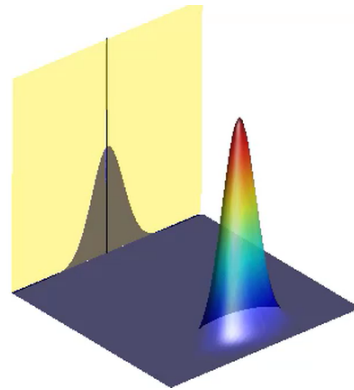
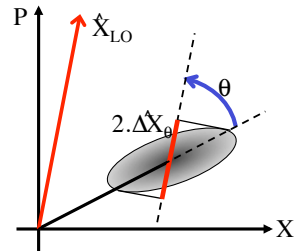
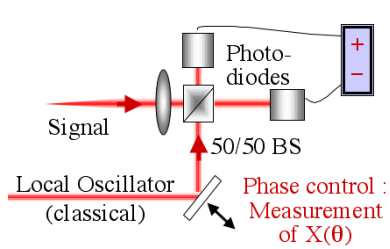
Homodyne detection, Wigner Function and Quantum Tomography



- Quasiprobability density :
Wigner function W(X,P)
- Marginals of W(X, P)
=> Probability distributions P(X_θ)
- Probability distributions P(X_θ)
=> W(X, P) (quantum tomography)



Homodyne detection, Wigner Function and Quantum Tomography



Squeezed State :
Gaussian Wigner Functions

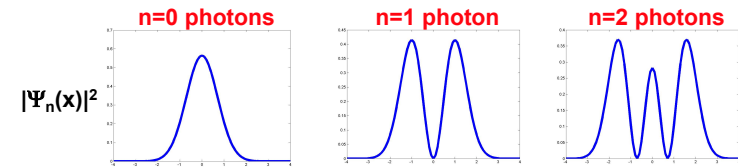
Non-Gaussian States

Basic question :

Consider a single photon : can we measure its amplitude & phase? quadratures X & P ?

Single mode light field
Photons
n photon state
Probability $P_n(X)$

Harmonic oscillator
Quanta of excitation
 n^{th} eigenstate
Probability $|\Psi_n(x)|^2$

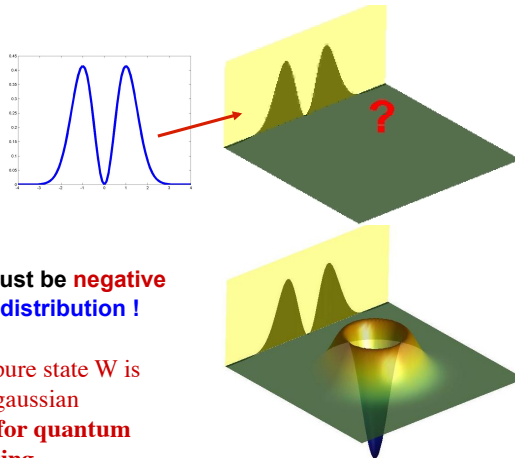


Non-Gaussian States

Basic question :

Consider a single photon : can we measure its amplitude & phase? quadratures X & P ?

Can the Wigner function of a Fock state $n = 1$ (with all projections have zero value at origin) be positive everywhere ?



NO ! The Wigner function must be **negative**
It is not a classical statistical distribution !

Hudson-Piquet theorem : for a pure state W is non-positive iff it is non-gaussian

Many interesting properties for quantum information processing

Wigner function of a single photon state ? (Fock state $n = 1$)

$$W(p, q) = \frac{1}{2\pi 2N_0} \int dx e^{\frac{ixp}{2N_0}} \langle q - \frac{x}{2} | \hat{\rho} | q + \frac{x}{2} \rangle$$

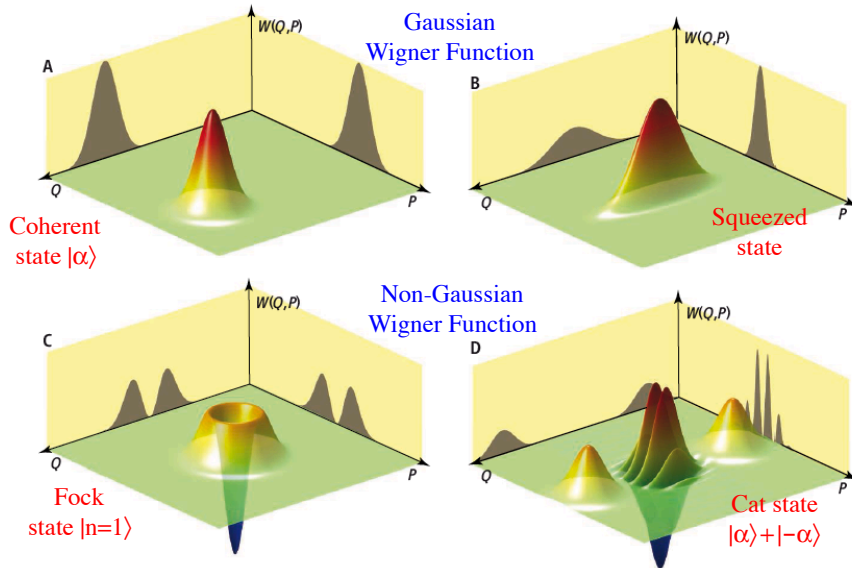
where $\hat{\rho} = |1\rangle\langle 1|$ and N_0 is the variance of the vacuum noise :

$$[\hat{Q}, \hat{P}] \equiv 2iN_0 \quad \Delta P \Delta Q \geq N_0 \quad N_0 = \Delta P^2 = \Delta Q^2.$$

One may have $N_0 = \hbar/2$, $N_0 = 1/2$ (theorists), $N_0 = 1$ (experimentalists)

Using the wave function of the $n = 1$ state : $\langle q | 1 \rangle = \frac{q}{(2\pi)^{\frac{1}{4}} N_0^{\frac{3}{4}}} e^{-\frac{q^2}{4N_0}}$

one gets finally : $W_{|1\rangle}(q, p) = -\frac{1}{2\pi N_0} e^{-\frac{r^2}{2N_0}} \left(1 - \frac{r^2}{N_0}\right) \quad r^2 = q^2 + p^2$



P. Grangier, "Make It Quantum and Continuous", Science (Perspective) 332, 313 (2011)

Make It Quantum and Continuous

Philippe Grangier PERSPECTIVES SCIENCE VOL 332 15 APRIL 2011

Unconditional Quantum Teleportation

A. Furusawa, J. L. Sørensen, S. L. Braunstein, C. A. Fuchs, H. J. Kimble,* E. S. Polzik

23 OCTOBER 1998 VOL 282 SCIENCE

Quantum key distribution using gaussian-modulated coherent states

Frédéric Grosshans¹, Gilles Van Assche¹, Jérôme Wenger¹, Rosa Brouri¹, Nicolas J. Cerf¹ & Philippe Grangier¹

NATURE | VOL 432 | 25 NOVEMBER 2004 | www.nature.com/nature

Experimental demonstration of quantum memory for light

Brian Julsgaard¹, Jacob Sherson^{1,2}, J. Ignacio Cirac¹, Jaromír Fluhráček¹ & Eugene S. Polzik¹

Vol 443 | 5 October 2006 | doi:10.1038/nature05136

Quantum teleportation between light and matter

Jacob F. Sherson^{1,2}, Hanna Krauter¹, Rasmus K. Oleson¹, Brian Julsgaard¹, Klemens Hammerer¹, Ignacio Cirac¹ & Eugene S. Polzik¹

PHYSICAL REVIEW A 68, 042319 (2003)

Quantum computation with optical coherent states

T. C. Ralph,* A. Gilchrist, and G. J. Milburn

W. J. Munro S. Glancy

Generating Optical Schrödinger Kittens for Quantum Information Processing

Alexei Ourjoumtsev, Rosa Tualle-Brouri, Julien Laurat, Philippe Grangier*

SCIENCE VOL 312 7 APRIL 2006

Vol 448 | 16 August 2007 | doi:10.1038/nature06054

Generation of optical 'Schrödinger cats' from photon number states

Alexei Ourjoumtsev¹, Hyunseok Jeong², Rosa Tualle-Brouri¹ & Philippe Grangier¹

Teleportation of Nonclassical Wave Packets of Light

Noriyuki Lee,¹ Hugo Benichou,¹ Yuichi Takeda,¹ Shuntaro Takeda,¹ James Webb,² Florian Hühnerfont², Akira Furusawa^{1*}

15 APRIL 2011 VOL 332 SCIENCE

Small sample, many more papers !

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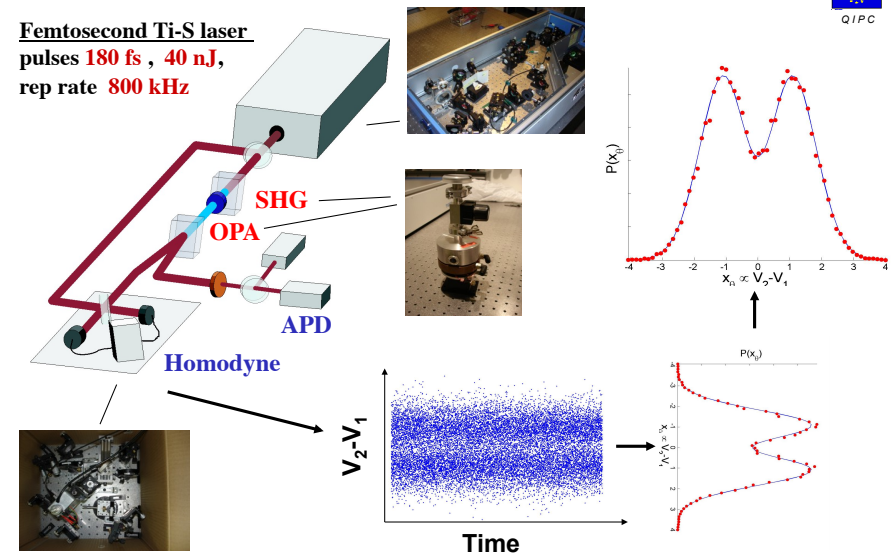
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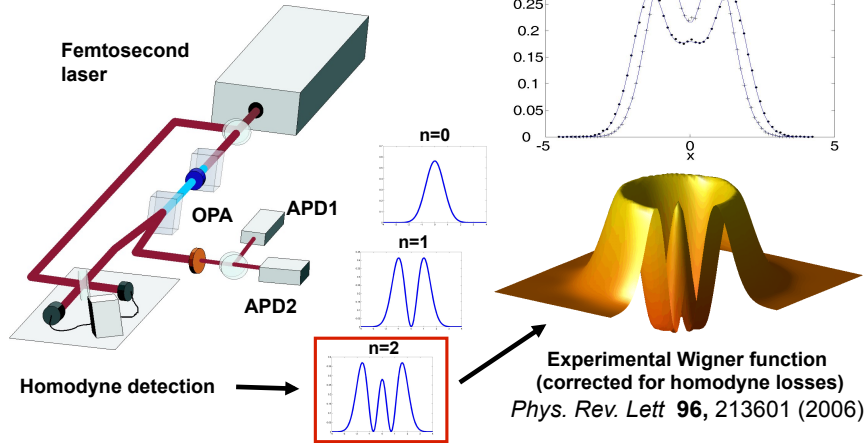
Experimental Set-up

Femtosecond Ti-S laser pulses 180 fs , 40 nJ, rep rate 800 kHz



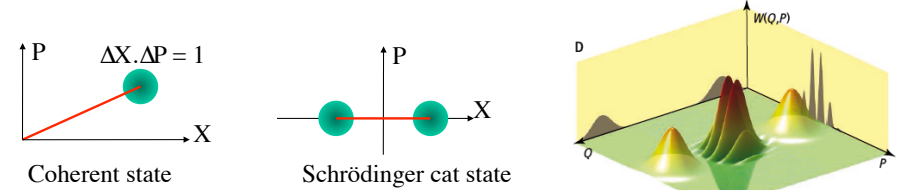
Resource : Two-Photon Fock States

$$|\psi\rangle = \sum \lambda^n |n, n\rangle$$



« Schrödinger's Cat » state

- Classical object in a quantum superposition of distinguishable states
- “Quasi - classical” state in quantum optics : coherent state $|\alpha\rangle$



$$|\psi_{odd\ cat}\rangle = c_o (|\alpha\rangle - |-\alpha\rangle)$$

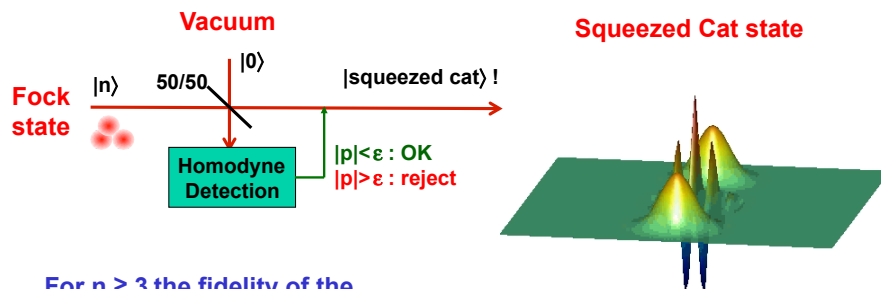
$$|\psi_{even\ cat}\rangle = c_e (|\alpha\rangle + |-\alpha\rangle)$$

- Resource for quantum information processing
- Model system to study decoherence

Wigner function of a Schrödinger cat state

How to create a Schrödinger's cat ?

Suggestion by Hyunseok Jeong, calculations by Alexei Ourjoutsev :

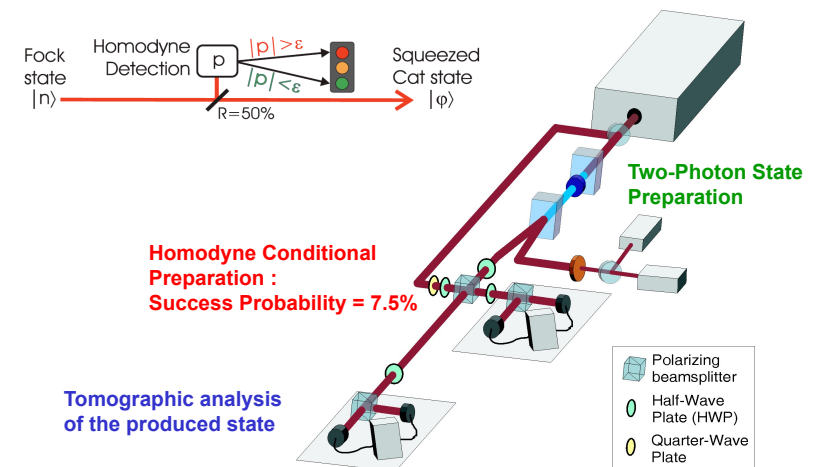


For $n \geq 3$ the fidelity of the conditional state with a Squeezed Cat state is $F \geq 99\%$

$$S(r)(|\alpha\rangle + e^{i\theta}|-\alpha\rangle)$$

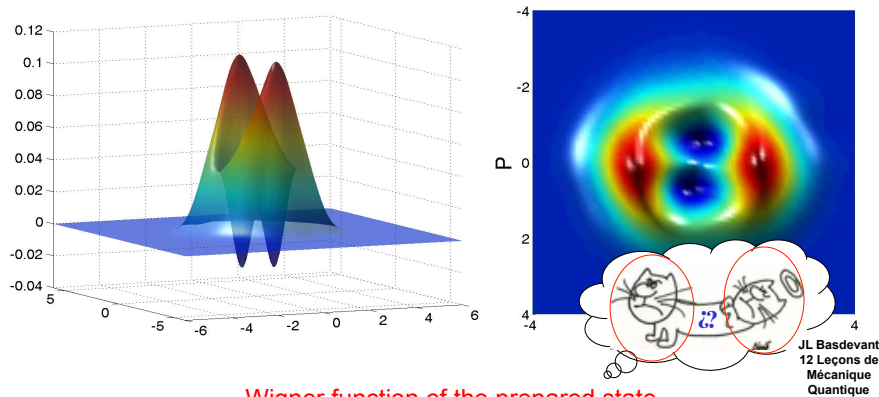
Size : $\alpha^2 = n$
 Same Parity as n : $\theta = n * \pi$
 Squeezed by : 3 dB

Squeezed Cat State Generation



Experimental Wigner function

A. Ourjoumtsev et al, Nature 448, 784 (2007)



Wigner function of the prepared state
Reconstructed with a Maximal-Likelihood algorithm
Corrected for the losses of the final homodyne detection.

Bigger cats : NIST (Gerrits, 3-photon subtraction), ENS (Haroche, microwave cavity QED), UCSB...

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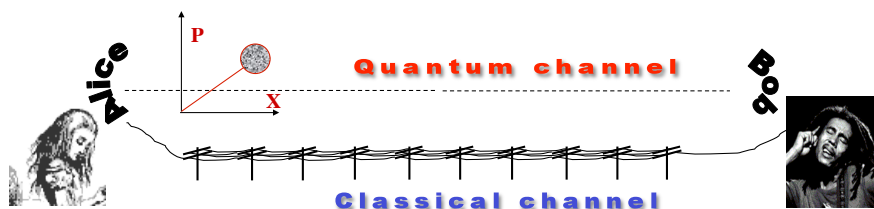
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Coherent States Quantum Key Distribution



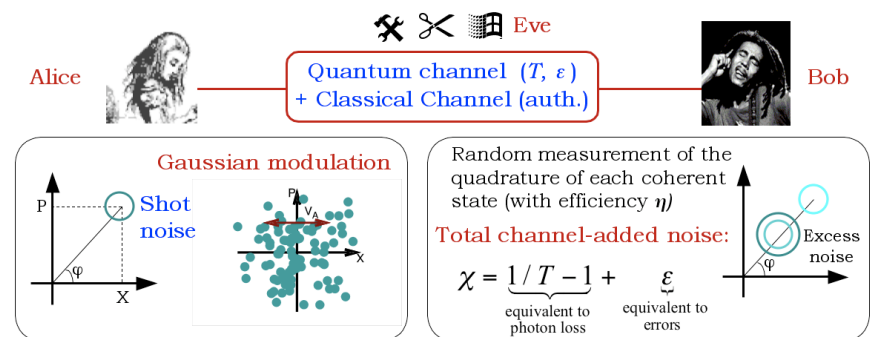
* Essential feature : quantum channel with non-commuting quantum observables
-> not restricted to single photon polarization or phase !

-> Design of Continuous-Variable QKD protocols where :

- * The non-commuting observables are the quadrature operators X and P
- * The transmitted light contains weak coherent pulses (about 10 photons) with a gaussian modulation of amplitude and phase
- * The detection is made using shot-noise limited homodyne detection

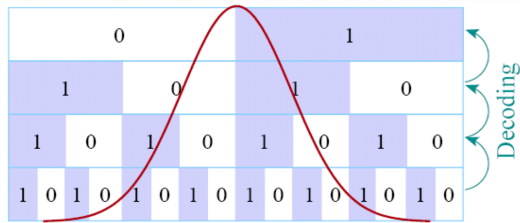
Coherent state continuous variables QKD protocol

- Key information encoded in both quadratures of a coherent state



- Bob reveals measurement choice
- Alice and Bob share a set of Gaussian correlated data
- Further communication to calculate channel parameters and derive secret key based on Bob's data → **reverse reconciliation**

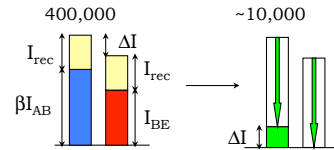
Reconciliation of correlated Gaussian variables



- Each level has a different error rate
 - Non-independent levels
- Error correction performed using multi-level iterative soft decoding with LDPC codes

G. Van Assche et al, IEEE Trans. on Inf. Theory 50, 394 (2004)
M. Bloch et al, arXiv:cs.IT/0509041 (2005)

- Standard privacy amplification based on universal hash functions
- Small processing time



Security of coherent state CV-QKD protocol

- Security initially proven against (arbitrary) **individual attacks** :
F. Grosshans et al, Nature 421, 238 (2003)
F. Grosshans and N. J. Cerf, Phys. Rev. Lett. 92, 047905 (2004)
 - Then security proven against **arbitrary collective attacks** :
F. Grosshans, Phys. Rev. Lett. 94, 020504 (2005)
M. Navasqu es and A. Acin, Phys. Rev. Lett. 94, 020505 (2005)
 - For both individual and collective attacks **Gaussian attacks are optimal**
→ Alice and Bob consider Eve's attacks Gaussian and estimate her information using the **Shannon quantity** I_{BE} or the **Holevo quantity** χ_{BE}
M. Navasqu es et al, Phys. Rev. Lett. 97, 190502 (2006)
R. Garcia-Patr on et al, Phys. Rev. Lett. 97, 190503 (2006)
- proofs of unconditional security** (against coherent attacks)
coherent attacks are not better than collective attacks.
R. Renner and J.I. Cirac, Phys. Rev. Lett. 102, 110504 (2009)
- Recent results :**
Finite size effects : A. Leverrier, F. Grosshans and P. Grangier, Phys. Rev. A 81, 062343 (2010)
Composable security proof : A. Leverrier, Phys. Rev. Lett. (2014)

Part 1 : Gaussian and non-Gaussian states

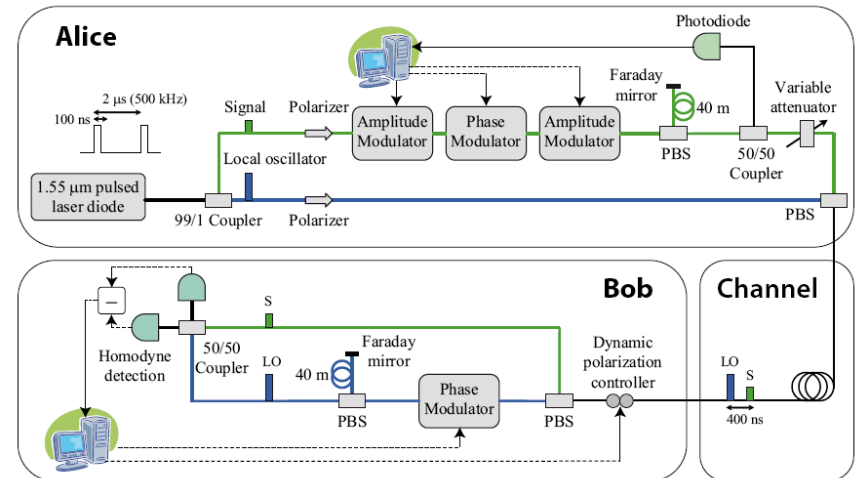
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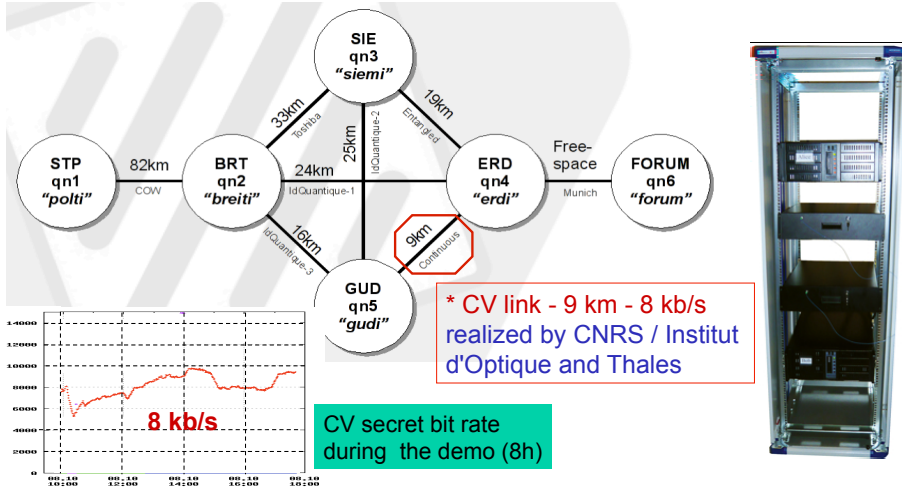


Field test of a continuous-variable quantum key distribution prototype
S Fossier, E Diamanti, T Debuisschert, A Villing, R Tualle-Brouiri and P Grangier
New J. Phys. 11 No 4, 04502 (April 2009)

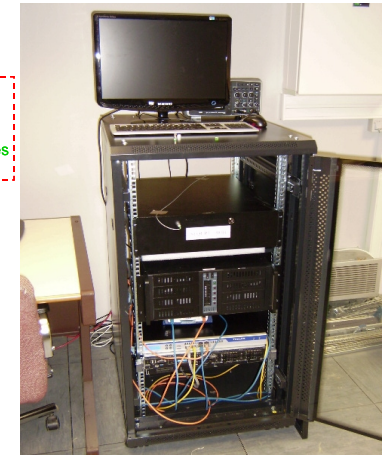
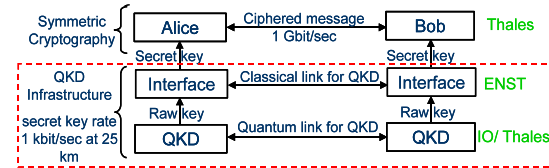
The SECOQC Quantum Back Bone



Real-size demonstration of a **secure quantum cryptography network** by the **European Integrated Project SECOQC**, Vienna, 8 october 2008



Symmetric Encryption with Quantum key REnewal SEQUIRE

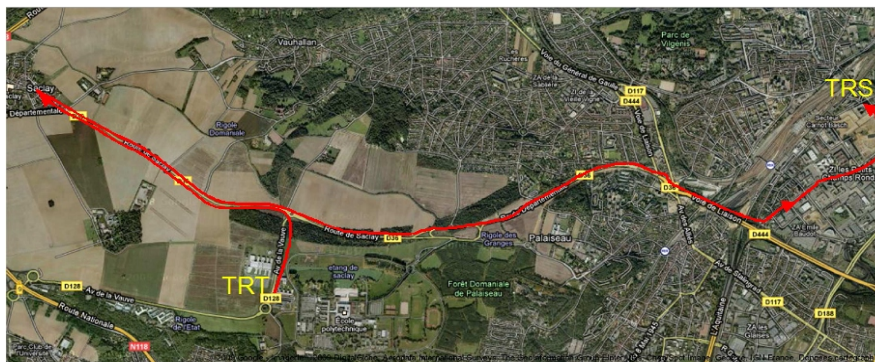


Thales : Mistral Gbit



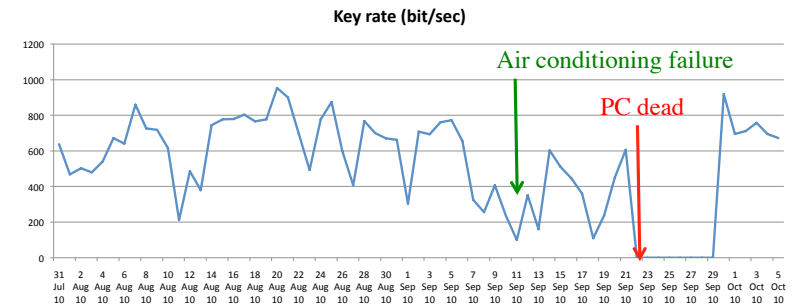
Field implementation

- Fibre link : Thales R&T (Palaiseau) <-> Thales Raytheon Systems (Massy)
- Fiber length about 12 km, 5.6 dB loss



Results

On site, 12 km distance, 5.6 dB loss
Minimal direct action on hardware (feedback loops, remote control)



See <http://www.demo-sequire.com>



Post-processing at SeQureNet



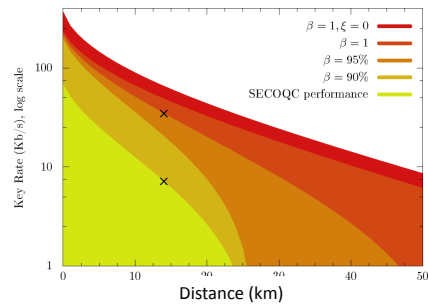
Last version (commercial device, 80 km) :
P. Jouguet et al, Nature Phot. 7, 378 (2013)

Paul Jouguet, Sébastien Kunz-Jacques, Romain Alléaume

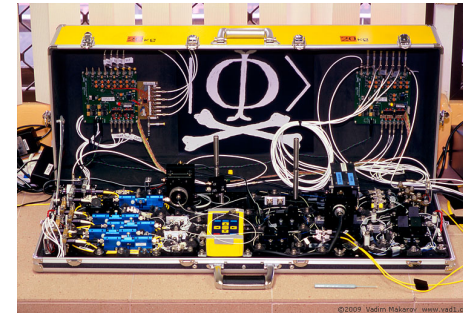
Optimize LDPC codes, use Graphic Processing Units (GPU) rather than CPU

=> Calculation speed is no more limiting the secret bit rate !

=> β is improved from 89% to 95% for any SNR : **longer distance (80 km) !**



CYGNUS (commercial product)



- Several recent examples of “quantum hacking” (e.g. Vadim Makarov et al.)
- Exploits weaknesses in single photon detectors
- **Will NOT work against CVQKD (PIN photodiodes, linear regime)**
- Hackers will have to work harder...
- ... and Trojan attacks will not make it (work under way, SQN + U. Erlangen)

arXiv.org > quant-ph > arXiv:1106.0825
Quantum Physics

Security of Post-selection based Continuous Variable Quantum Key Distribution against Arbitrary Attacks

Nathan Walk, Thomas Symul, Timothy C. Ralph, Ping Koy Lam
(Submitted on 4 Jun 2011)

arXiv.org > quant-ph > arXiv:1011.0304
Quantum Physics

Continuous variable quantum key distribution in non-Markovian channels

Ruggero Vasile, Stefano Olivares, Matteo G A Paris, Sabrina Maniscalco
(Submitted on 1 Nov 2010)

arXiv.org > quant-ph > arXiv:0904.1694
Quantum Physics

Feasibility of continuous-variable quantum key distribution with noisy coherent states

Vladyslav C. Usenko, Radim Filip
(Submitted on 10 Apr 2009 (v1), last revised 21 Jan 2010 (this version, v2))

arXiv.org > quant-ph > arXiv:0904.1327
Quantum Physics

Security bound of continuous-variable quantum key distribution with noisy coherent states and channel

Yong Shen, Jian Yang, Hong Guo
(Submitted on 8 Apr 2009 (v1), last revised 29 Jun 2009 (this version, v2))

arXiv.org > quant-ph > arXiv:0903.0750
Quantum Physics

Confidential direct communications: a quantum approach using continuous variables

Stefano Pirandola, Samuel L. Braunstein, Seth Lloyd, Stefano Mancini
(Submitted on 4 Mar 2009)

Many other works on CVQKD !

<= Theory and Experiments : (incomplete list !)

arXiv.org > quant-ph > arXiv:1006.1257
Quantum Physics

A balanced homodyne detector for high-rate Gaussian-modulated coherent-state quantum key distribution

Yue-Meng Chi, Bing Qi, Wen Zhu, Li Qian, Hoi-Kwong Lo, Sun-Hyun Youn, A. I. Lvovsky, Liang Tian
(Submitted on 7 Jun 2010 (v1), last revised 16 Jul 2010 (this version, v2))

arXiv.org > quant-ph > arXiv:0910.1042
Quantum Physics

A 24 km fiber-based discretely signaled continuous variable quantum key distribution system

Quyen Dinh Xuan, Zheshen Zhang, Paul L. Voss
(Submitted on 6 Oct 2009)

arXiv.org > quant-ph > arXiv:0811.4756
Quantum Physics

Feasibility of free space quantum key distribution with coherent polarization states

D. Elser, T. Bartley, B. Heim, Ch. Wittmann, D. Sych, G. Leuchs
(Submitted on 28 Nov 2008 (v1), last revised 13 Mar 2009 (this version, v2))

arXiv.org > quant-ph > arXiv:0705.2627
Quantum Physics

Experimental Demonstration of Post-Selection based Continuous Variable Quantum Key Distribution in the Presence of Gaussian Noise

Thomas Symul, Daniel J. Alton, Syed M. Assad, Andrew M. Lance, Christian Weedbrook, Timothy C. Ralph, Ping Koy Lam
(Submitted on 18 May 2007)

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Long distance quantum communications

How to fight against line losses ?

~~Amplification~~



Long distance quantum communications

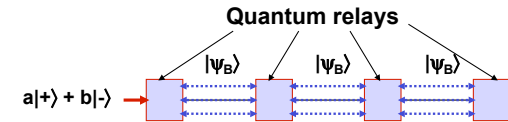
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1. Exchange of entangled states

$|\Psi_B\rangle$



Long distance quantum communications

How to fight against line losses ?

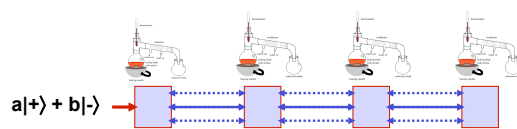
~~Amplification~~



1. Exchange of entangled states

$|\Psi_B\rangle$

2. Entanglement distillation



Long distance quantum communications

How to fight against line losses ?

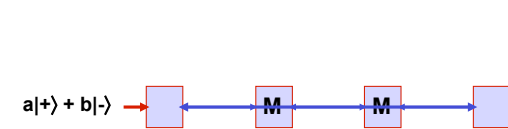
~~Amplification~~



1. Exchange of entangled states

$|\Psi_B\rangle$

2. Entanglement distillation



Long distance quantum communications

How to fight against line losses ?

~~Amplification~~

$$a|+\rangle + b|-\rangle \xrightarrow{G} \xrightarrow{G} \xrightarrow{G} a|+\rangle + b|-\rangle$$

1. Exchange of entangled states

$$|\Psi_B\rangle$$

2. Entanglement distillation



$$a|+\rangle + b|-\rangle \xrightarrow{M} \xrightarrow{\text{channel}} \xrightarrow{M}$$

3. Entanglement swapping



~~Amplification~~

$$a|+\rangle + b|-\rangle \xrightarrow{G} \xrightarrow{G} \xrightarrow{G} a|+\rangle + b|-\rangle$$

1. Exchange of entangled states

$$|\Psi_B\rangle$$

2. Entanglement distillation



$$a|+\rangle + b|-\rangle \xrightarrow{\text{channel}} a|+\rangle + b|-\rangle$$

3. Entanglement swapping



4. Quantum teleportation

One needs to : * distribute (many) entangled states
* store them (quantum memories)
* process them (distillation)

Part 1 : Gaussian and non-Gaussian states

1. Homodyne detection and quantum tomography
2. Generating non-Gaussian Wigner functions : kittens, cats and beyond

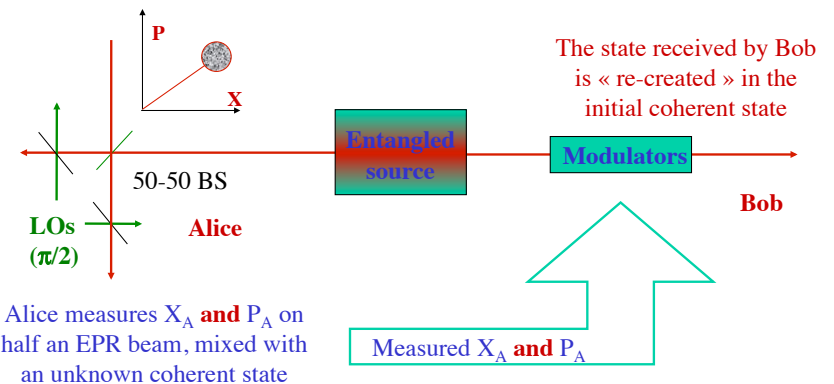
Part 2 : Continuous variable quantum cryptography (Gaussian !)

1. Continuous variable quantum cryptography : principles
2. Continuous variable quantum cryptography : implementations

Part 3 : Towards quantum networks (non-Gaussian !)

1. Entanglement for continuous variable quantum networks
2. Teleportation of Schrödinger's cats
3. Storing non-Gaussian states : single photon quantum memory

Quantum teleportation of coherent states

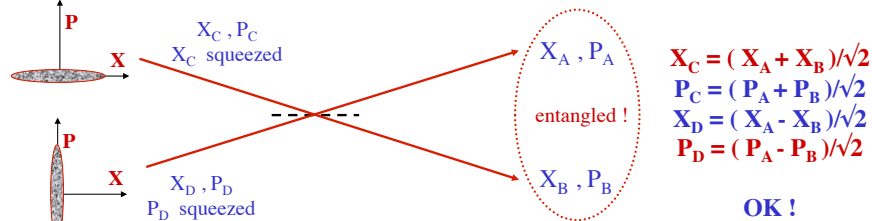


Experiments :

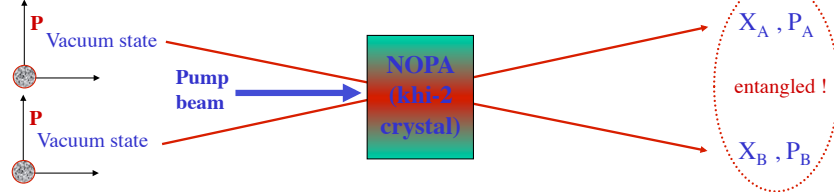
- A. Furusawa et al, Science **282**, 706 (1998)
- W. Bowen et al, Phys. Rev. A **67**, 032302 (2003)
- T.C. Zhang et al, Phys. Rev. A **67**, 033802 (2003)

How to produce CV entangled beams ?

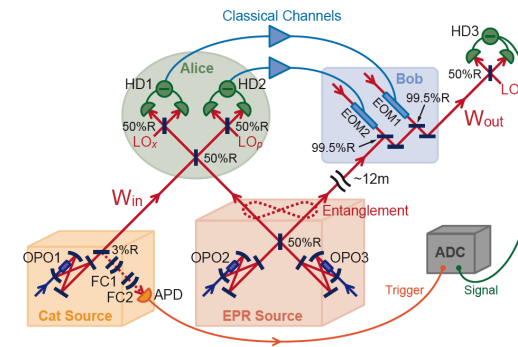
1. Combine two orthogonally squeezed beams



2. Use a Non-degenerate Optical Parametric Amplifier (NOPA)



Quantum teleportation of cat states

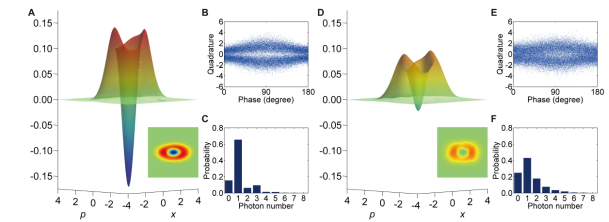


The state received by Bob is « re-created » in the initial cat state (with some loss in fidelity : $0.75 \Rightarrow 0.45$)

Experiment : N. Lee et al, Science **332**, 330-333 (2011).

Remark : the initial squeezed state is CW, not pulsed !

Alice measures X_A and P_A on half an EPR beam, mixed with a cat state (obtained from photon subtraction)



Part 1 : Gaussian and non-Gaussian states

1. Homodyne detection and quantum tomography
2. Generating non-Gaussian Wigner functions : kittens, cats and beyond

Part 2 : Continuous variable quantum cryptography (Gaussian !)

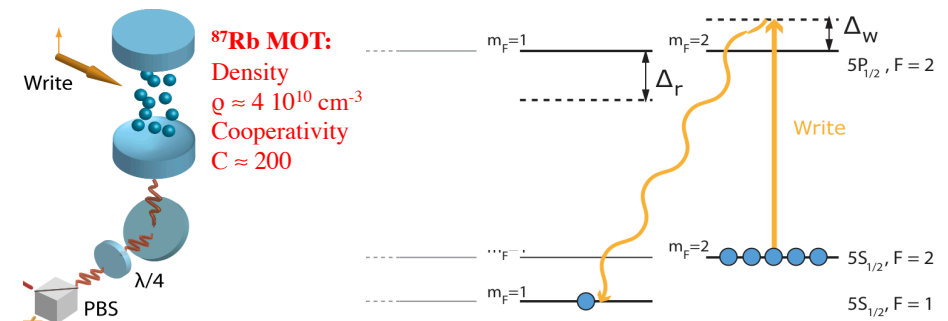
1. Continuous variable quantum cryptography : principles
2. Continuous variable quantum cryptography : implementations

Part 3 : Towards quantum networks (non-Gaussian !)

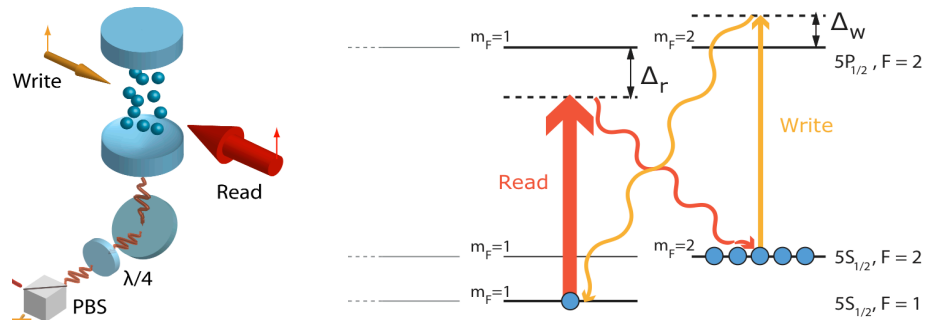
1. Entanglement for continuous variable quantum networks
2. Teleportation of Schrödinger's cats
3. Storing non-Gaussian states : single photon quantum memory

Single Photon from a single polariton (DLCZ protocol)

L.M. Duan, M.D. Lukin, J.I. Cirac, and P. Zoller, Nature **414**, 413 (2001)

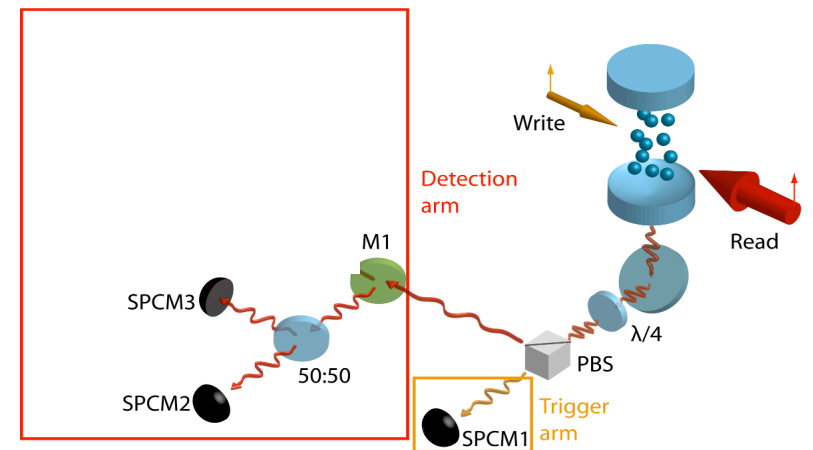


Single Photon from a single polariton (DLCZ protocol)



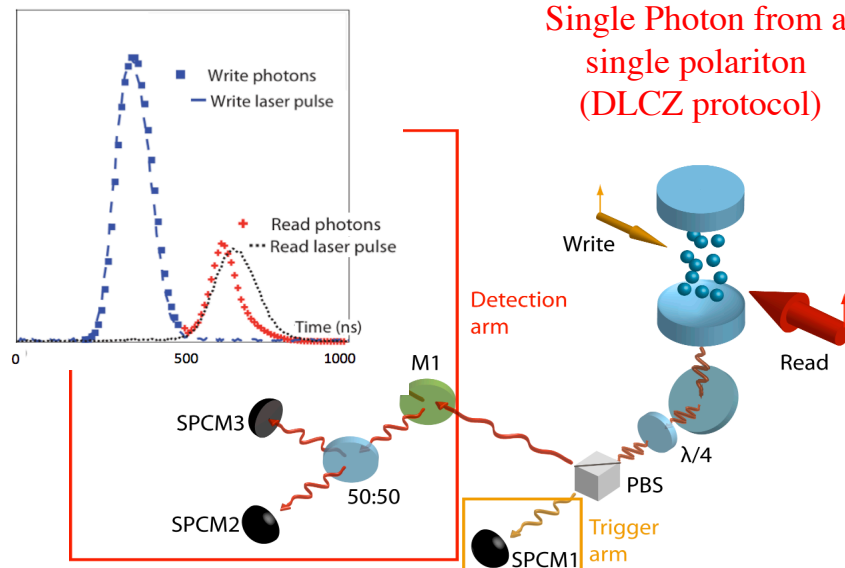
E. Bimbard et al, arXiv:1310.1228 (2013), PRL 112, 033601 (2014)

Single Photon from a single polariton (DLCZ protocol)



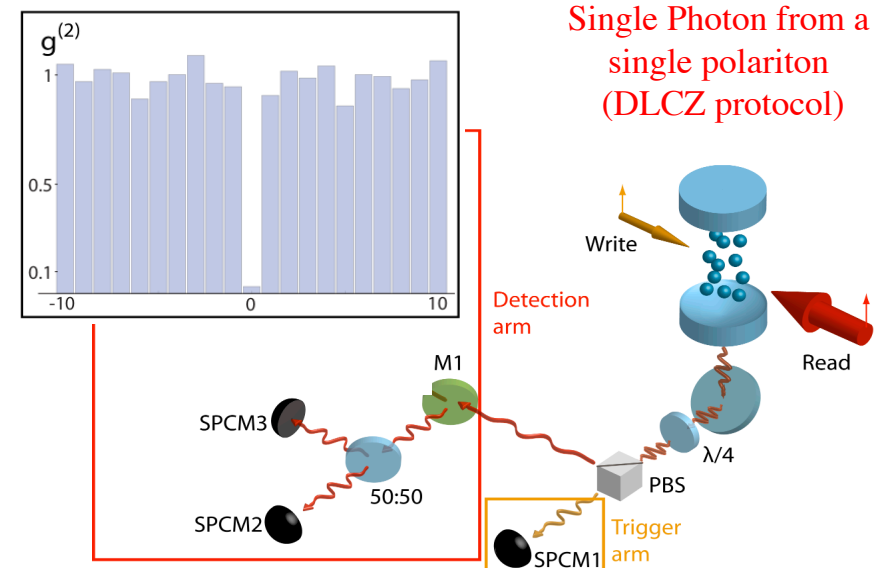
E. Bimbard et al, arXiv:1310.1228 (2013), PRL 112, 033601 (2014)

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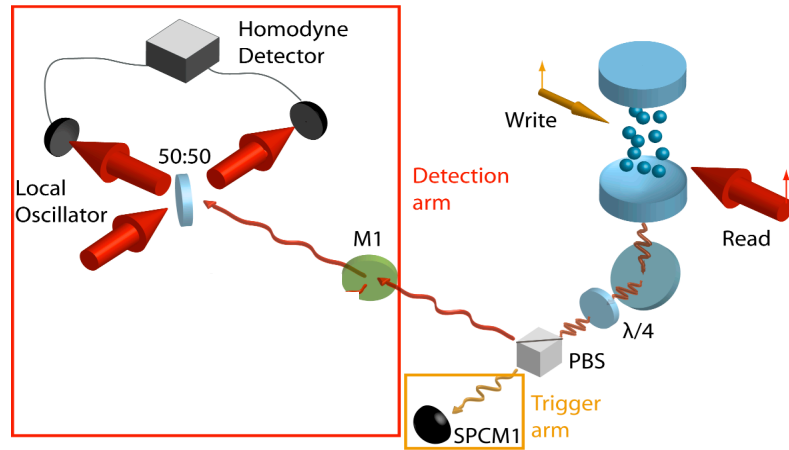
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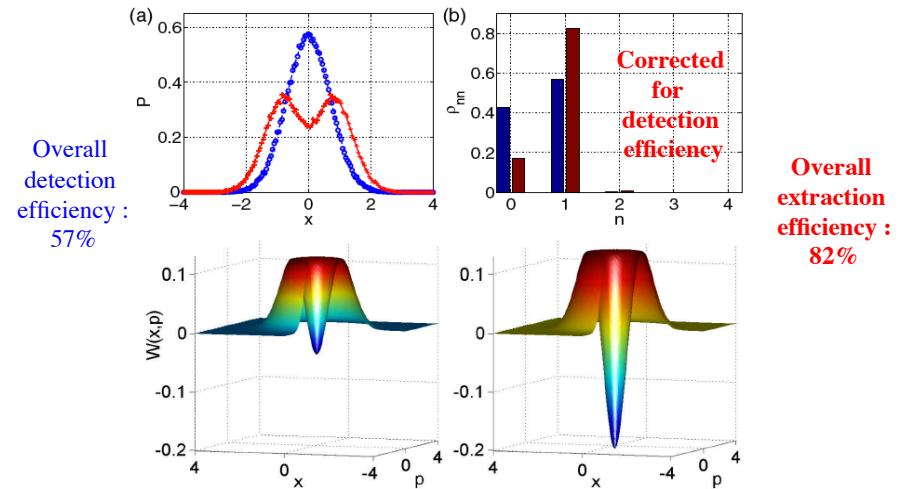
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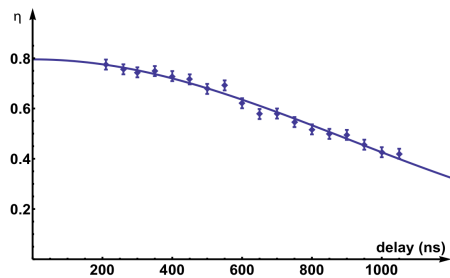
Single Photon from a single polariton (DLCZ protocol)



E. Bimbard et al, arXiv:1310.1228 (2013), PRL 112, 033601 (2014)

Single Photon from a single polariton (DLCZ protocol)

Quantum memory effect : the memory time ($1 \mu s$) is limited by motional decoherence due to finite temperature ($50 \mu K$)



$$\eta = P_{\text{Doppler}}(t) \times P_{\text{Coop}} \times P_{\text{Read}} \times P_{\text{Pumping}} \times P_{\text{Mode}} \times P_{\text{Cav}}$$

$$0.94 \times 0.97 \times 0.96 \times 0.965 \times 0.97 = \underline{0.82 : \text{ok!}}$$

E. Bimbard et al, arXiv:1310.1228 (2013), PRL 112, 033601 (2014)

Thank you for your attention !



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Etienne Brion
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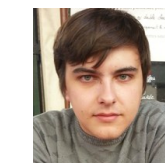
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(PhD, exp.)



Andrey Grankin
(PhD, th.)



Rajiv Boddeda
(PhD, exp.)

Rydberg interactions team
(Palaiseau 2014)



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Deterministic Logical
Photon-photon Interactions