





中国科学院高能物理研究所 Institute of High Energy Physics Chinese Academy of Sciences

SUSY SEARCHES IN ATLAS

Xuai Zhuang

On behalf of ATLAS Collaboration

IHEP, Beijing, China CPPM, Marseille, France

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CPPM Seminar

Outline

- SUSY Introduction
- The LHC and ATLAS
- Overview of SUSY search results on run-1 data
- Outlook and Summary







Outline



- Overview of SUSY search results on run-1 data
- Outlook and Summary



SM and Beyond





Photo: A. Mahmoud François Englert



Photo: A. Mahmoud Peter W. Higgs



The Nobel Prize in Physics 2013 François Englert, Peter Higgs

- Higgs boson observed, SM fits the experimental data very well → big success in EW scale
- While has problem in Planck scale:
 - Naturalness and "hierarchy" problem
 - Unification of gauge coupling
 - Dark Matter



New Physics beyond the SM



SUSY Introduction



A symmetry which unified fermions (mater) and bosons (forces) -> A fundamental theory

Conserved R parity (originally introduced for stability
of proton)R=(-1)^{3(B-L)+2S}R=+1 (SM)
R=(-1)(SU(S))

- SUSY particles produced/annihilated in pairs
- Lightest SUSY particle (LSP) stable (DM candidate)
- Typical signature: jets/leptons/photons + MET

Violated R parity (RPV): not mentioned in the talk

14-12-08



R=-1 (SUSY)

SUSY Introduction



Solve hierarchy problem without "fine tuning"

 SUSY contributions to Higgs mass cancel SM contributions

Unification of gauge couplings

 New particle content changes running of couplings

Provide Dark Matter candidate

 Lightest SUSY particle (LSP) can be stable and only weekly interacting



log., (Q/GeV)

Some of the arguments are most convincing for SUSY particles at ~TeV scale I3-8-25

log10(Q/GeV)

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SUSY Introduction *The LHC and ATLAS* Overview of SUSY search results on run-1 data

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ATLAS and CMS detector @ LHC ATLAS and CMS: two multi-purpose detectors @LHC

A Toroidal LHC ApparatuS

- 42m×22m, 7000 ton
- Solenoid + Toroidal magnet (2T)
- Fine granularity liquid Ar/Tile calorimeters

Large Hadron Collider (LHC):

Proton-Proton synchrotron
 World's highest and largest collider

Compact Muon Spectrometer

ATLAS

- 21m×15m, 125000 ton
- All silicon trackers, 4T solenoid magnet
- PbWO4+Tile calorimeters

Since 2010, ATLAS&CMS have invested huge efforts in SUSY search @LHC : Great Luminosity recorded



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SUSY Search Strategy

SUSY search strategy: search for deviation from SM

- SUSY sensitive variables: Try to establish excess of events in some sensitive kinematic distribution
- SM background: the discovery of new physics can only be claimed when SM backgrounds are understood well or under control
 - SM bgs understood very well ③
 - − No hints for new physics ⊗
 - Slightly overshoot in WW cross section, but consistent with NNLO xsec.





SUSY Sensitive Variables



E_T^{miss} from escaping LSP, to suppress bg from mismeasured jets and oth. SM BG

Related to the sparticle mass scale, like effective mass (**M**_{eff})

$$M_{\rm eff} \equiv \sum_{i=1}^{N_{\rm jets}} p_{\rm T}^{\rm jet,i} + \sum_{j=1}^{N_{\rm lep}} p_{\rm T}^{\rm lep,j} + E_{\rm T}^{\rm miss}$$

mT, mT2 (stransverse mass): suppress BG with Ws $m_{T2} = \min_{\mathbf{q}_T} \left[\max \left(m_T(\mathbf{p}_T^{\ell 1}, \mathbf{q}_T), m_T(\mathbf{p}_T^{\ell 2}, \mathbf{p}_T^{miss} - \mathbf{q}_T) \right) \right]$

Many others ...

SM Background Modeling

SUSY searches rely on accurate modeling of the Standard Model backgrounds

Standard Model

Top, multijets V, VV, VVV, Higgs & combinations of these

Combined fit of all regions and backgrounds and incl. systematic exp. and theor. uncertainties as nuisance parameters

Reducible backgrounds

Determined from data Backgrounds and methods depend on analyses

Irreducible backgrounds

Dominant sources: normalise MC in data control regions Subdominant sources: MC

Validation

Validation regions used to cross check SM predictions with data

Signal regions

blinded

blinded

SUSY models: good sale in market

□ Simplified Models:

- Not really a model (Br~100%, most masses fixed at high scales)
- Important tool for interpretation
- □ Phenomenological models:
 - pMSSM: captures "most" of phenomenologic features of R-parity conserving MSSM
 - 19 free parameters: M1,M2,M3 ; tan β, μ and m_{A;} 10 sfermion mass parameters; A_t , A_b and A_τ
 - Comprehensive and computationally realistic approximation of the MSSM with neutralino LSP
 - GGM (gravitino)

□ Complete SUSY models: mSUGRA, GMSB ...

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SUSY Search @ LHC





Inclusive Search with strong prod.

Strong production: gluino pair, gluino-squark and squark pair (include 3rd generation) production

1) Generic signatures :

Multi -jets + n_lepton/n_photon(n=0,1, \geq 2) + large E_T^{miss} (0L,1 L, >=2L)

2) large xs, but heavy SUSY mass scale

ol +jets(2-6)+MET: JHEP 09 (2014) 176

- ol +jets(7-10)+MET: JHEP 10 (2013) 130
- 0-1 l + >=3 b-jets + MET: JHEP 10 (2014) 024
- 1-2l +jets + MET: ATLAS-SUSY-2013-20
- 2 SS/ 3 l+ 0-3 b-jets + MET: JHEP 06 (2014) 035
- 1-2 taus + 0-1 l + jets + MET: JHEP 09 (2014) 103





${}_{ m s}$ Searches for pair prod. of light-flavor squarks ${}^{{}_{ m s}}$

or gluinos decaying via virtual squarks

- Very powerful inclusive search
- 0I + 2-6 jets + MET (jet + MET trigger)
- Main discriminating variable: Meff
- 15 SRs from 2 to 6 jets to achive maximum reach on (m_~q, m_~g) -plane
- Backgrounds taken from dedicated data CRs (and validated)







Search for squarks/gluinos via stop/ sbottom decay

JHEP 10 (2014) 024



Search for squarks/gluinos via long decay chain (SS2L/3L) JHEP 06 (2014) 035

- Target signals with gluino pair or guino-squark pair production resulting 2I SS or 3I final states (5 SRs with SS2L + 0 or 1 or 3 bjets, or 3L)
- Very clean channels with only tiny SM bg (mainly top+V, diboson, triboson)
- Discriminating variable: Meff, MET

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/IHEP

These SRs increase sensitivity to longer decay cascades, compressed mass spectra, 3rd generation squarks ...





Inclusive search for squark and gluino production Summary



- m(~q) ~m(~g): m(~g) > 1.7 TeV
- M(~q) >> m(~g): m(~g)>1.35 TeV
- Conditional/indirect limit on LSP: m> 200-300 GeV
- No exclusion for M(LSP) ≥ 700 GeV
- − Strongest limit: $m(\sim g) \ge 1400 \text{ GeV}$

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Inclusive search for squark and gluino production Summary





Direct stop/sbottom pair production

Dedicated effort to search for direct stop/sbottom pair production
 Large spectrum of possible stop/sbottom decays
 Effort concentrated on simplified models with 100% BRs to chosen final state



1-lepton stop search: 2/3/4-body decay to LSP JHEP 11 (2014) 118



1-lepton stop search: 2/3/4-body decay to LSP JHEP 11 (2014) 118



1-lepton stop search: decay to b + chargino

JHEP 11 (2014) 118



1-lepton stop search: decay to b + chargino

JHEP 11 (2014) 118



All Hadronic Stop Search (0I)

JHEP 09 (2014) 015

)=(600,1) GeV

400

m^{b,min}_T [GeV]

x⁰)=(400.200.100) GeV



30

Direct stop pair production Summary



Strong search program on stop, covering range from low to heavy stop mass, various decay modes.

Exclusion for m(~t1) < ~660GeV for massless LSP, exclusion up to m(LSP) ~250 GeV **CPPM Seminar**

Direct EWK-ino Search

- Search for electroweak (EWK) SUSY below the TeV scale is motivated by naturalness arguments.
- EWK production has a low crosssection compared to strong production
 - Very challenging searches
 - But leads to multi-lepton signatures with very low SM background.
- If strong production is suppressed, EWK processes could be the dominant SUSY production at the LHC. (EWKino < 1 TeV)

<u>Generic signature</u>: >=1 lepton(s) in the final state arising from the decay of charginos/neutralinos via sleptons/sneutrinos, gauge bosons or Higgs.





ElectroWeak production

(direct C/N and slepton production)

C1N2 via Wh, w→lv,h→bb,γγ,II FS: 1I +bb/γγ or 3I +MET (ATLAS-CONF-2014-062) C1N2 via slepton, FS: 3I + MET (I=e/μ/τ) 3I (e/μ): JHEP 04 (2014)169 ≥2 τ : JHEP 10 (2014) 096 C1N2 via WZ FS: 2/3I+MET 2I: JHEP 05 (2014) 071 3I: JHEP 04 (2014)169







C1C1 via slepton FS: 2L+MET(I=e/μ/т) 2I (e/μ): JHEP 05 (2014)071 2 τ : JHEP 10 (2014) 096

а



C1C1 via WW FS: 2L+MET 2I (e/µ): JHEP 05 (2014)071



Direct slepton pair FS: 2L+MET (l=e/μ/т) 2I (e/μ): JHEP 05 (2014)071 2 τ : JHEP 10 (2014) 096



Electroweak production – Wh

[direct C1N2 production via higgs]



Electroweak production – Wh

[direct C1N2 production via higgs]



Electroweakino mass > 250 GeV for massless LSP in models with higgs in decay

Electroweak production – 2L

[direct C1/ sl production]

JHEP 05 (2014) 071



Electroweak production – 2L

[direct C1/ sl production]

JHEP 05 (2014) 071



ElectroWeak production – 3L

[direct C1N2 production]

JHEP 04 (2014)169



ElectroWeak production – 3L

[direct C1N2 production]



Electroweak production – 2tau

[direct C1(N2)/ sl production]

SR-C1N2

Data

 $\frac{10^4}{10^3}$

10

Data/SM 1.2 2.0





First search for production of C1C1 and Dstau with hardronic tau decays from the LHC !

- Direct gaugino/stau production with final state: 2tau + MET
- 4 SRs designed targeting different scenarios: one for C1N2, one for C1C1 and two for Dstau
- Main backgrounds:
 - W+jets (1real+1fake): normalized MC to data in WCR
 - Multi-jet (2fake): ABCD data-driven estimation



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Electroweak production – 2tau

[direct C1(N2)/ sl production]





CL Limit

m₋₀=60 GeV

250

300

m₌ [GeV]





300

m, [GeV]





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ElectroWeak Production Summary



Comparable for C1N2 via slepton

CMS: SS improvement in compressed region; no results on C1C1 via stau and WW

ATLAS: more results for compressed scenario will be ready soon

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Outlook and Summary

ATLAS developed a vast program to search for SUSY

- No significant excess seen so far
- All public results:

https://twiki.cern.ch/twiki/bin/view/AtlasPublic/SupersymmetryPublicResults

- In canonical scenarios, sensitivity is achieved to ~1.2 TeV gluinos, ~700 GeV stops and ~400 GeV for EWKinos
- The reach with SUSY is expected to increase significantly at run2 and run3

Long term prospects

 ATLAS studied long term prospects for the (HL-)LHC with 300, 3000 fb⁻¹@14 TeV
 Discovery potential up to 2.5 TeV gluinos, 1.3 TeV squarks/sbottom and 800 GeV Electroweakinos



Exciting times are ahead of us !



THANKS FOR YOUR ATTENTION!

ATLAS SUSY Searches* - 95% CL Lower Limits

Status: ICHEP 2014

0.0	Model	e, μ, τ, γ	Jets	$E_{ m T}^{ m miss}$	∫ <i>L dt</i> [fl	b ⁻¹]	Mass limit		Reference
Inclusive Searches	$ \begin{array}{l} \text{MSUGRA/CMSSM} \\ \text{MSUGRA/CMSSM} \\ \text{MSUGRA/CMSSM} \\ \tilde{q}\tilde{q}, \tilde{q} \rightarrow q \tilde{\chi}_{10}^0 \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow q \tilde{q} \tilde{\chi}_{1}^0 \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow q \tilde{q} \tilde{\chi}_{1}^1 \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow q q \tilde{\chi}_{1}^1 \rightarrow q q \tilde{\chi}_{1}^1 \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow q q \tilde{\chi}_{1}^1 \rightarrow q q \tilde{\chi}_{1}^1 \\ \text{GMSB} (\tilde{\ell} \text{ NLSP}) \\ \text{GMSB} (\tilde{\ell} \text{ NLSP}) \\ \text{GGM} (bino \text{ NLSP}) \\ \text{GGM} (bino \text{ NLSP}) \\ \text{GGM} (higgsino-bino \text{ NLSP}) \\ \text{GGM} (higgsino-bino \text{ NLSP}) \\ \text{GGM} (higgsino \text{ NLSP}) \\ \text{GGM} (higgsino \text{ NLSP}) \\ \text{Gravitino LSP} \\ \end{array} $	$\begin{matrix} 0 \\ 1 \ e, \mu \\ 0 \\ 0 \\ 0 \\ 1 \ e, \mu \\ 2 \ e, \mu \\ 2 \ e, \mu \\ 1 \ 2 \ \tau + 0 \ 1 \ \ell \\ 2 \ \gamma \\ 1 \ e, \mu + \gamma \\ \gamma \\ 2 \ e, \mu (Z) \\ 0 \end{matrix}$	2-6 jets 3-6 jets 2-6 jets 2-6 jets 2-6 jets 3-6 jets 0-3 jets 2-4 jets 0-2 jets - 1 b 0-3 jets mono-jet	Yes Yes Yes Yes Yes Yes Yes Yes Yes Yes	20.3 20.3 20.3 20.3 20.3 20.3 20.3 20.3		1.2 TeV 1.1 TeV 850 GeV 1.33 Ti 1.33 Ti 1.33 Ti 1.34 TeV 1.24 TeV 1.24 TeV 1.28 Te 619 GeV 900 GeV 690 GeV 690 GeV 645 GeV	1.7 TeV $m(\tilde{q})=m(\tilde{g})$ any $m(\tilde{q})$ any $m(\tilde{q})$ $m(\tilde{\chi}_{1}^{n})=0 \text{ GeV}, m(1^{st} \text{ gen}, \tilde{q})=m(2^{nd} \text{ gen}, \tilde{q})$ eV $m(\tilde{\chi}_{1}^{n})=0 \text{ GeV}$ $m(\tilde{\chi}_{1}^{n})=0 \text{ GeV}$ $m(\tilde{\chi}_{1}^{n})=0 \text{ GeV}$ $tan\beta < 15$ 1.6 TeV $tan\beta < 20$ V $m(\tilde{\chi}_{1}^{n})=50 \text{ GeV}$ $m(\tilde{\chi}_{1}^{n})=50 \text{ GeV}$ $m(\tilde{\chi}_{1}^{n})=50 \text{ GeV}$ $m(\tilde{\chi}_{1}^{n})=50 \text{ GeV}$ $m(\tilde{\chi}_{1}^{n})=2200 \text{ GeV}$	1405.7875 ATLAS-CONF-2013-062 1308.1841 1405.7875 1405.7875 ATLAS-CONF-2013-062 ATLAS-CONF-2013-089 1208.4688 1407.0603 ATLAS-CONF-2012-012 ATLAS-CONF-2012-144 1211.1167 ATLAS-CONF-2012-152 ATLAS-CONF-2012-152
ğ med.	$\begin{array}{c} \tilde{g} \rightarrow b \tilde{b} \tilde{\chi}_{1}^{0} \\ \tilde{g} \rightarrow t \tilde{\chi}_{1}^{0} \\ \tilde{g} \rightarrow t \tilde{\chi}_{1}^{1} \\ \tilde{g} \rightarrow b \tilde{\chi}_{1}^{*} \end{array}$	0 0 0-1 <i>e</i> ,μ 0-1 <i>e</i> ,μ	3 <i>b</i> 7-10 jets 3 <i>b</i> 3 <i>b</i>	Yes Yes Yes Yes	20.1 20.3 20.1 20.1	100 100 100 100 100	1.25 Te 1.1 TeV 1.34 T 1.34 T		1407.0600 1308.1841 1407.0600 1407.0600
3 rd gen. squarks direct production	$\begin{split} \tilde{b}_1 \tilde{b}_1, \tilde{b}_1 \rightarrow b \tilde{\chi}_1^0 \\ \tilde{b}_1 \tilde{b}_1, \tilde{b}_1 \rightarrow t \tilde{\chi}_1^{\pm} \\ \tilde{i}_1 \tilde{t}_1 (\text{light}), \tilde{i}_1 \rightarrow t \tilde{\chi}_1^{\pm} \\ \tilde{i}_1 \tilde{i}_1 (\text{light}), \tilde{i}_1 \rightarrow W b \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{medium}), \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{medium}), \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neavy}), \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neavy}), \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neavy}) = \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neavy}) = \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neav}) = \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neav}) = \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neav}) = \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neav}) = \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neav}) = \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neav}) = \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neav}) = \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neav}) = \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neav}) = \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neav}) = \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neav}) = \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neav}) = \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neav}) = \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 (\text{neav}) = \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neav}) = \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neav}) = \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neav}) = \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neav}) = \tilde{i}_1 \rightarrow t \tilde{i}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{i}_1 \tilde{i}_1 (\text{neav}) = \tilde{i}_1 \rightarrow t \tilde{i}_1 \rightarrow t \tilde{i}_1 \rightarrow t \tilde{i}_1 \end{pmatrix}$	$\begin{array}{c} 0\\ 2\ e,\mu\ ({\rm SS})\\ 1\text{-}2\ e,\mu\\ 2\ e,\mu\\ 2\ e,\mu\\ 0\\ 1\ e,\mu\\ 0\\ 0\\ 3\ e,\mu\ (Z) \end{array}$	2 b 0-3 b 1-2 b 0-2 jets 2 jets 2 b 1 b 2 b nono-jet/c-t 1 b 1 b	Yes Yes Yes Yes Yes Yes Yes Yes Yes Yes	20.1 20.3 4.7 20.3 20.3 20.1 20 20.1 20.3 20.3 20.3	$ \begin{array}{c} \tilde{b}_{1} \\ \tilde{b}_{1} \\ \tilde{c}_{1} \\ \tilde{i}_{1} \\ \tilde{i}_{2} \end{array} $	100-620 GeV 275-440 GeV 110 <mark>-167 GeV</mark> 130-210 GeV 215-530 GeV 210-640 GeV 260-640 GeV 90-240 GeV 150-580 GeV 290-600 GeV	$\begin{array}{c} m(\tilde{\chi}_{1}^{0}) < 90 \ \mathrm{GeV} \\ m(\tilde{\chi}_{1}^{+}) = 2 \ m(\tilde{\chi}_{1}^{0}) \\ m(\tilde{\chi}_{1}^{0}) = 55 \ \mathrm{GeV} \\ m(\tilde{\chi}_{1}^{0}) = 55 \ \mathrm{GeV} \\ m(\tilde{\chi}_{1}^{0}) = 1 \ \mathrm{GeV} \\ m(\tilde{\chi}_{1}^{0}) = 1 \ \mathrm{GeV} \\ m(\tilde{\chi}_{1}^{0}) = 200 \ \mathrm{GeV}, \ m(\tilde{\chi}_{1}^{+}) - m(\tilde{\chi}_{1}^{0}) = 5 \ \mathrm{GeV} \\ m(\tilde{\chi}_{1}^{0}) = 0 \ \mathrm{GeV} \\ m(\tilde{\chi}_{1}^{0}) = 0 \ \mathrm{GeV} \\ m(\tilde{\chi}_{1}^{0}) = 150 \ \mathrm{GeV} \\ m(\tilde{\chi}_{1}^{0}) < 150 \ \mathrm{GeV} \\ m(\tilde{\chi}_{1}^{0}) < 200 \ \mathrm{GeV} \\ \end{array}$	1308.2631 1404.2500 1208.4305, 1209.2102 1403.4853 1403.4853 1308.2631 1407.0583 1406.1122 1407.0608 1403.5222 1403.5222
EVV direct	$ \begin{array}{c} \tilde{\ell}_{1,\mathbf{k}}\tilde{\ell}_{1,\mathbf{k}},\tilde{\ell} \rightarrow \ell\tilde{\chi}_{1}^{0} \\ \tilde{\chi}_{1}^{\dagger}\tilde{\chi}_{1}^{-},\tilde{\chi}_{1}^{\dagger} \rightarrow \tilde{\ell}\nu(\ell\tilde{\nu}) \\ \tilde{\chi}_{1}^{\dagger}\tilde{\chi}_{0}^{-},\tilde{\chi}_{1}^{\dagger} \rightarrow \tilde{\tau}\nu(\tau\tilde{\nu}) \\ \tilde{\chi}_{1}^{\dagger}\tilde{\chi}_{0}^{0} \rightarrow \tilde{\ell}_{1}\nu\tilde{\ell}_{L}(\ell(\tilde{\nu}\nu),\ell\tilde{\nu}\tilde{\ell}_{L}\ell(\tilde{\nu}\nu) \\ \tilde{\chi}_{1}^{\dagger}\tilde{\chi}_{0}^{0} \rightarrow W\tilde{\chi}_{1}^{0}Z\tilde{\chi}_{1}^{0} \\ \tilde{\chi}_{1}^{\dagger}\tilde{\chi}_{0}^{0} \rightarrow W\tilde{\chi}_{1}^{0}h\tilde{\chi}_{1}^{1} \\ \tilde{\chi}_{2}^{\dagger}\tilde{\chi}_{0}^{0},\tilde{\chi}_{2,3}^{0} \rightarrow \tilde{\ell}_{R}\ell \end{array} $	$\begin{array}{c} 2 \ e, \mu \\ 2 \ e, \mu \\ 2 \ \tau \\ 3 \ e, \mu \\ 2 \ 3 \ e, \mu \\ 2 \ 3 \ e, \mu \\ 1 \ e, \mu \\ 4 \ e, \mu \end{array}$	0 0 - 0 2 <i>b</i> 0	Yes Yes Yes Yes Yes Yes Yes	20.3 20.3 20.3 20.3 20.3 20.3 20.3 20.3	$ \vec{\tilde{\ell}} \\ \vec{\tilde{\chi}}_{1}^{\pm} \\ \vec{\tilde{\chi}}_{1}^{\pm} \\ \vec{\tilde{\chi}}_{1}^{\pm} \\ \vec{\tilde{\chi}}_{2}^{0} \\ \vec{\tilde{\chi}}_{1}^{\pm} \\ \vec{\tilde{\chi}}_{2}^{0} \\ \vec{\tilde{\chi}}_{2,3}^{0} \\ \vec{\tilde{\chi}}_{2,3}^{0} $	90-325 GeV 140-465 GeV 100-350 GeV 700 GeV 420 GeV 285 GeV 620 GeV	$\begin{array}{l} m(\tilde{\chi}_{1}^{0}){=}0~\text{GeV} \\ m(\tilde{\chi}_{1}^{0}){=}0~\text{GeV}, m(\tilde{\ell},\tilde{\nu}){=}0.5(m(\tilde{\chi}_{1}^{+}){+}m(\tilde{\chi}_{1}^{0})) \\ m(\tilde{\chi}_{1}^{0}){=}0~\text{GeV}, m(\tilde{\tau},\tilde{\nu}){=}0.5(m(\tilde{\chi}_{1}^{+}){+}m(\tilde{\chi}_{1}^{0})) \\ m(\tilde{\chi}_{1}^{+}){=}m(\tilde{\chi}_{2}^{0}), m(\tilde{\chi}_{1}^{0}){=}0, m(\tilde{\xi},\tilde{\nu}){=}0.5(m(\tilde{\chi}_{1}^{+}){+}m(\tilde{\chi}_{1}^{0})) \\ m(\tilde{\chi}_{1}^{+}){=}m(\tilde{\chi}_{2}^{0}), m(\tilde{\chi}_{1}^{0}){=}0, sleptons~decoupled \\ m(\tilde{\chi}_{2}^{0}){=}m(\tilde{\chi}_{2}^{0}){=}0, m(\tilde{\ell},\tilde{\nu}){=}0.5(m(\tilde{\chi}_{2}^{0}){+}m(\tilde{\chi}_{1}^{0})) \end{array}$	1403.5294 1403.5294 1407.0350 1402.7029 1403.5294, 1402.7029 ATLAS-CONF-2013-093 1405.5086
Long-Ilved particles	Direct $\tilde{\chi}_1^+ \tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^+$ Stable, stopped \tilde{g} R-hadron GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu}) + \tau(e,$ GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$, long-lived $\tilde{\chi}_1^0$ $\tilde{q}\tilde{q}, \tilde{\chi}_1^0 \rightarrow qq\mu$ (RPV)	Disapp. trk 0 μ) 1-2 μ 2 γ 1 μ, displ. vtx	1 jet 1-5 jets - -	Yes Yes - Yes -	20.3 27.9 15.9 4.7 20.3	$ \begin{array}{c} \tilde{x}_1^{\pm} \\ \bar{g} \\ \tilde{x}_1^{0} \\ \bar{x}_1^{0} \\ \tilde{x}_1^{0} \\ \tilde{q} \end{array} $	270 GeV 832 GeV 832 GeV 475 GeV 1.0 TeV	$\begin{array}{l} m(\tilde{\chi}_{1}^{\tau})\text{-}m(\tilde{\chi}_{1}^{0})\text{=}160 \; MeV, \; \tau(\tilde{\chi}_{1}^{\tau})\text{=}0.2 \; ns \\ m(\tilde{\chi}_{1}^{0})\text{=}100 \; GeV, \; 10 \; \mu s{<}\tau(\tilde{g}){<}1000 \; s \\ 10{<}tan\beta{<}50 \\ 0.4{<}\tau(\tilde{\chi}_{1}^{0}){<}2 \; ns \\ 1.5 < c\tau{<}156 \; mm, \; BR(\mu)\text{=}1, \; m(\tilde{\chi}_{1}^{0})\text{=}108 \; GeV \end{array}$	ATLAS-CONF-2013-069 1310.6584 ATLAS-CONF-2013-058 1304.6310 ATLAS-CONF-2013-092
RPV	$ \begin{array}{c} LFV pp \rightarrow \tilde{v}_{\tau} + X, \tilde{v}_{\tau} \rightarrow e + \mu \\ LFV pp \rightarrow \tilde{v}_{\tau} + X, \tilde{v}_{\tau} \rightarrow e(\mu) + \tau \\ Bilinear \ RPV \ CMSSM \\ \tilde{\chi}^+_1 \tilde{\chi}^1, \tilde{\chi}^+_1 \rightarrow W \tilde{\chi}^0_1, \tilde{\chi}^0_1 \rightarrow e e \tilde{v}_{\mu}, e \mu \tilde{v}_e \\ \tilde{\chi}^+_1 \tilde{\chi}^1, \tilde{\chi}^+_1 \rightarrow W \tilde{\chi}^0_1, \tilde{\chi}^0_1 \rightarrow \tau \tau \tilde{v}_e, e \tau \tilde{v}_{\tau} \\ \tilde{g} \rightarrow q q \\ \tilde{g} \rightarrow \tilde{t}_1 t, \tilde{t}_1 \rightarrow b s \end{array} $	$\begin{array}{c} 2 \ e, \mu \\ 1 \ e, \mu + \tau \\ 2 \ e, \mu \ (\text{SS}) \\ 4 \ e, \mu \\ 3 \ e, \mu + \tau \\ 0 \\ 2 \ e, \mu \ (\text{SS}) \end{array}$	- 0-3 b - 6-7 jets 0-3 b	- Yes Yes Yes - Yes	4.6 4.6 20.3 20.3 20.3 20.3 20.3 20.3	$\frac{\tilde{v}_{\tau}}{\tilde{v}_{\tau}}$ $\frac{\tilde{q}_{\tau}}{\tilde{g}_{\tau}} \frac{\tilde{g}_{\tau}}{\tilde{\chi}_{1}^{\pm}}$ $\tilde{\chi}_{1}^{\pm}$ \tilde{g}_{τ} \tilde{g}	1. 1.1 TeV 1.35 T 750 GeV 450 GeV 916 GeV 850 GeV	.61 TeV $\lambda'_{311}=0.10, \lambda_{132}=0.05$ $\lambda'_{311}=0.10, \lambda_{1(2)33}=0.05$ reV $m(\tilde{q})=m(\tilde{g}), c_{TLSP}<1 \text{ mm}$ $m(\tilde{\chi}^0_1)>0.2\times m(\tilde{\chi}^1_1), \lambda_{121}\neq 0$ $m(\tilde{\chi}^0_1)>0.2\times m(\tilde{\chi}^1_1), \lambda_{133}\neq 0$ BR(t)=BR(b)=BR(c)=0%	1212.1272 1212.1272 1404.2500 1405.5086 1405.5086 ATLAS-CONF-2013-091 1404.250
Other	Scalar gluon pair, sgluon $\rightarrow q\bar{q}$ Scalar gluon pair, sgluon $\rightarrow t\bar{t}$ WIMP interaction (D5, Dirac χ)	0 2 <i>e</i> , µ (SS) 0	4 jets 2 <i>b</i> mono-jet	- Yes Yes	4.6 14.3 10.5	sgluon sgluon M* scale	100-287 GeV 350-800 GeV 704 GeV	incl. limit from 1110.2693 $m(\chi)$ <80 GeV, limit of<687 GeV for D8	1210.4826 ATLAS-CONF-2013-051 ATLAS-CONF-2012-147
	$\sqrt{s} = 7 \text{ TeV}$ full data	$\sqrt{s} = 8$ TeV artial data	$\sqrt{s} = $ full	8 TeV data		<u>_</u>	10 ⁻¹ 1	Mass scale [TeV]	

*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty.

ATLAS Preliminary

 $\sqrt{s} = 7.8 \text{ TeV}$



From LHCC Open meeting, 03.12.2013 Seminar

Excesses seen so far (Exotics, SUSY)

For a more detailed review, see the physics plenary talk (July 18th) by T. Golling

Mono-jet <u>CDS link</u>

- $~\sim \! 1.7\sigma$ / 2.4σ excess above BG in the two highest MET SRs
- High MET SRs dropped for paper (very low stat in CR)
- **VV** \rightarrow **JJ** <u>CDS link</u>
 - Mass of fat jets each consistent with W,Z mass
 - No issue found in cross checks
- Same-sign leptons / 3-leptons + b-jets CDS link
 - No issue found in cross checks, being made public
- 3-leptons + 3 b-jets CDS link
 - ~3.5 σ in a validation region of the 3-lepton search
 - No issue found, but will not be made public as not a SR
 - Plan a dedicated SR for run2

Z+jets + MET <u>CDS link</u>

- Peaking at Z mass, BG dominated by non-Z (tt)
- No issue found in cross checks, to be made public soon
- SUSY: <u>https://twiki.cern.ch/twiki/bin/viewauth/AtlasProtected/SusyDiscrepancies</u> Exotics: <u>https://twiki.cern.ch/twiki/bin/view/AtlasProtected/ExoticsExcessSummary</u>



ATLAS inter

Long-term prospects

CMS and ATLAS studied long-term prospects for the (HL) LHC.

- with 300 and 3000 /fb at 14 TeV
- searches for gluino-mediated stop production reach beyond 2 TeV
- searches for direct stop production reach well beyond 1 TeV







Challenge:

"Stealth" stop

- Close to gap: m(stop1)-m(N1) ≈ m(top), signal similar to tt backgrounds, general kinematic Variables not discriminating (MET, mT, ...)
 Small signal xsec: ~ 15% SM ttbar for same mass
- Try use new technique: shape fit, MVA, top meas.
- Search for heavier SUSY production with subsequent decay via stop1: gluino to stop1+t, or stop2 to stop1+Z/H



Direct production of sbottom pairs



- bjet # = 1 or 2, 0 lepton
- Signal discriminant is M_{CT} (8 exclusive SRs)



- bjet # =0-2
- 2 same-sign (SS) or 3 leptons
- SRs based on E_T^{miss} , N_{jets} , $M_T(\ell_1 - E_t^{miss})$, $m_{eff} = E_t^{miss} + \Sigma_{jet} p_T$ $+ \Sigma_\ell p_T$



- bjet # >=3, 0 lepton
- ⁰ can decay (BR=100%) to the lightest neutral Higgs boson (or Z), h->bb̄(same BRs than the SM Higgs)



Stop Charm Search

- Direct stop production search targeting the region: $\Delta m \equiv m_{ ilde{t}_1} m_{ ilde{\chi}_1^0} < m_W + m_b$
 - 4-body decay
 - Decay $\tilde{t}_1 \rightarrow c + \tilde{\chi}_1^0$
- ISR based search strategies: •
 - Monojet-like selection ($\Delta m < 20 \text{ GeV}$): $\leq 3 \text{ jets}$, 1 hard ISR jet, large E_{Tmiss}
 - c-tagged selection ($\Delta m > 20 \text{ GeV}$): $\geq 4 \text{ jets}$, 1 hard ISR jet, $\geq 1 \text{ c-tagged jet}$, large E_{T.miss}
- **Background estimation**
 - Data-driven multi-jet estimate, W/Z + jets & tt (charm-tagged) normalized with data-driven scale factors



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Natural SUSY

What are the minimal requirements for `natural' (low level of fine-tuning) SUSY?

Recall: **Top loop** is the most important contribution in Higgs virtual corrections:



While inclusive LHC searches set O(TeV) limits mostly on 1st and 2nd generation squarks and gluinos, which are less relevant in Higgs corrections.



Illustration taken from: L. Hall, LBNL workshop, Oct-2011

For natural SUSY need:

- light higgsinos
- light top and bottom squarks (stop, sbottom, respectively)
- not-too-heavy gluinos
- strong physics case for third generation squarks

2 leptons razor



14-12-08

taus + jets + MET

Target signals: GMSB, natural gauge mediated nGM model (tuned version of GGM to avoid fine tuning in Higgs sector) with stau as NLSP



 \tilde{g} $\tilde{\chi}_{1}^{0}$ $\tilde{\tau}$ \tilde{G} τ

a

 τ

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•Hadronically decaying taus (with veto on e/μ in pure τ -channels)

Data driven estimate of fake taus events

	1τ SR	2τ GMSB SR	2τ nGM	τ+l nGM SR
m _T [GeV]	т _т *>140	$m_T^{\tau 1} + m_J$	^{τ2} >250	m _T ¹ >100
H _T ² j [GeV] MET [GeV]	>800 >200	>900	>600	>350
N(jet)		>=4	>=4	>=3



Search for Charm Squark, Results

500r

Direct $\tilde{c}\tilde{c}^*$, $\tilde{c} \rightarrow c\tilde{\chi}_1^0$, single \tilde{c} state

Observed Limit ($\pm 1\sigma_{\text{theory}}^{\text{SUSY}}$) Expected Limit ($\pm 1\sigma_{exp}^{SM}$)

200

New

Dominant uncertainties originate from the limited number of events in the CRs ~20%, jet tagging and mis-tagging $\sim 20\%$, jet energy scale ~10%



Limit on charm squark mass ~ 540 GeV, for massless LSP improves significantly on single light flavor squark limit improved sensitivity for heavier LSP

ATLAS-CONF-2014-062

20

600

400

500

300

Search for Pair Production of Light-flavor Squarks



- The cross-section for 1st and 2nd generation squarks benefits from 8-fold degeneracy (u, d, c, s) x (left, right)
 - comparable to the gluino pair production cross-section
- Searches using events with 0 leptons, 2-6+ jets



CMS: PAS-SUS-14-010(NEW)



Discovery and exclusion

 P-value=probability that result is as/less compatible with the hypothesis

DISCOVERY:

- The <u>null hypothesis</u> H₀ describes <u>background only</u>
 - If the *p*-value of H₀ is found below a given threshold, one can consider looking for a better model
 - In HEP, $Z \ge 5$ is conventionally required to claim a discovery
- The alternative hypothesis H₁ describes signal + background
 - The alternative hypothesis is supposed to fit the data very well for claiming a discovery

EXCLUSION:

- The <u>null hypothesis</u> H₀ describes <u>signal + background</u>
 - One is interested into setting an upper limit to the intensity of the signal alone
- The alternative hypothesis H₁ describes background only
 - No real need to test for it
 - The background-only model becomes important only in case of discovery

Interpretation strategy



Likelihood function: L(μ,θ) μ: signal strength (POI); θ: nuisance parameters(NP) Profile Likelihood: constrain uncertainty (NP) as part of a likelihood fit



Construct test statistics t_µ based on likelihood ratio λ:



Find the observed test statistic for tested µ: t_{µ.obs}



If CLs<0.05: the value



Simultaneous fit

- Background estimates in SRs are obtained by a simultaneous fit in each channel based on the profile likelihood method. Three dedicated fit for different purpose...
 - Background-only fit
 - Fit for all CRs, excluding SRs.
 - Get background-only estimates.
 - Also extrapolate to VRs (non used in fit, only for cross-check) and SRs.
 - Discovery fit
 - Fit for all CRs and SRs.
 - Signal contamination is turned off in CRs and set as a dummy number 1 in SR (so, the fitted non-SM signal strength = the excess in Nevents of SR)
 - Get model-independent upper limit on signal in SR.
 - Exclusion fit
 - Fit for all CRs and SRs.
 - Signal is turned on in all regions, according to model-dependent prediction.
 - Got signal model-dependent exclusion from all CRs+SRs →final exclusion contours for SUSY model

• The basic strategy is to share background information in all regions (CR, SR, VR). The background parameters are predominantly constrained by CRs with large statistics, which in turn reduces the impact of uncerts in SR.²⁰⁸

Hadronic Taus



- Tau decays:
 - Leptonic (35%): $au
 ightarrow
 u_{ au} \ell \overline{
 u}_{\ell}$
 - Hadronic (65%): decay to one or three charged pions, neutrinos and π^{0} 's
- Need to separate τ 's from hadronic jets:
 - \circ au decay tends to be well collimated
 - $\circ~$ Large electromagnetic component from $\pi^{0} \rightarrow \gamma \gamma~{\rm decay}$

Tau Object

- $ightarrow p_{
 m T} > 20$ GeV, $|\eta| < 2.5$
- \triangleright 1 or 3 tracks with total charge ± 1
- Boosted decision tree (BDT) using variables sensitive to the longitudinal and transverse shower shape
- Working points:
 - Loose: efficiency: 60%; jet rejection: 20-50

CPPM Seminar Tight: efficiency: 30-50%; jet rejection: 30-200

Reconstructed Objects

- Photons: no track but energy in el-m (and not in the hadronic) calorimeter
- Electrons: track and energy in el-m (and not in the hadronic) calorimeter
- Muons: track in inner tracker and muon chamber
- Jets: cluster in hadronic calorimeter



MET: Missing Transverse Energy

- At the LHC an unknown proportion of the energy of the colliding protons escapes down the beam-pipe
- Invisible particles (neutrinos, neutralinos?) are created their momentum can be constrained in the plane transverse to the beam direction

 $\boldsymbol{E}_T^{\mathrm{miss}} = -\sum \boldsymbol{p}_T(i)$



Data-Driven-Method (simultaneously fit)

- Background estimates in SRs are obtained by a *simultaneous fit* in each channel based on the profile likelihood method.
 - Fit for all BG-CRs
- **Comments:**
 - □ Shape from MC (CR), NF from fit
 - □ Advantage:
 - Considered correlations between different BG CRs and VRs, SRs;
 - -considered correlated syst. from detector
 - Disadvantage: shape is depending on MC (semi-data-driven method)



Background Estimation Strategy

ATLAS-CONF-2013-062



Other small BGs (diboson, single top etc) are directly estimated from MC.

4-



Inner Detector: Highly segmented silicon strips, determine very accurately charged particles trajectories

Solenoid Magnet: Solenoid coil that generates a 2T magnetic field in the region of the Inner Detector

Electromagnetic Calorimeter: Electron and photon energies are measured through electromagnetic showers

- Hadronic Calorimeter: Hadrons interact with dense material and produce a shower of charged particles
- Toroid Magnets: 8 toroidal coils that create a 0,4T magnetic field in the area of the Muon Spectrometer
- Muon Spectrometer: Muons traverse the rest of the detector and are measured in its outer layers