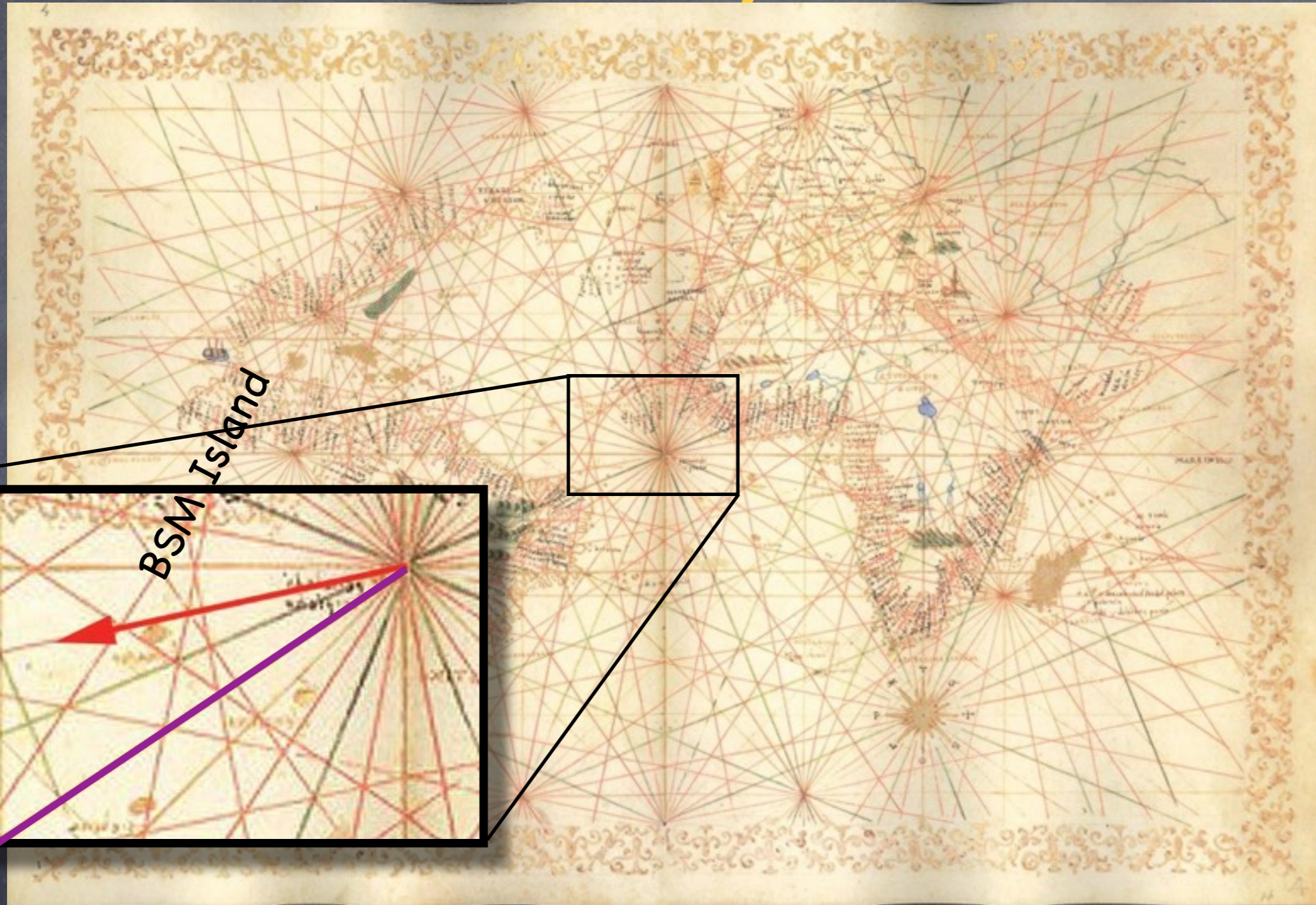


# BSM Primary Effects



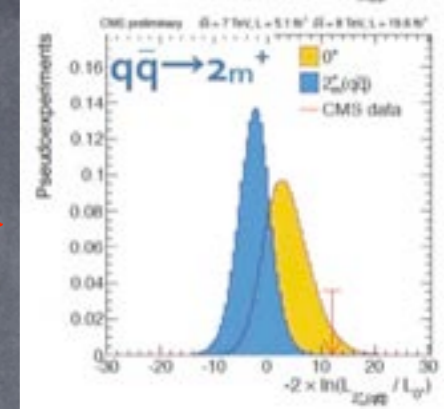
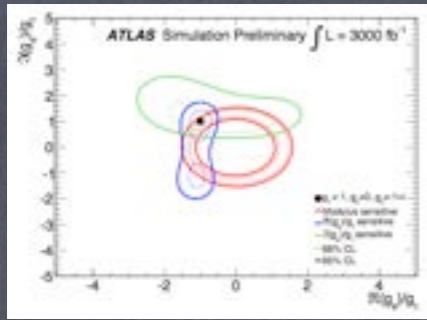
Francesco Riva (EPFL - Lausanne)

In Collaboration with:

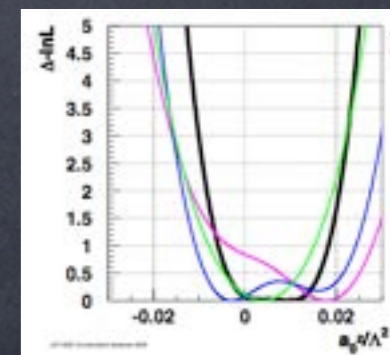
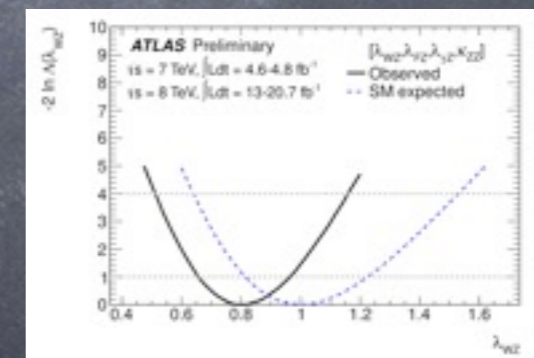
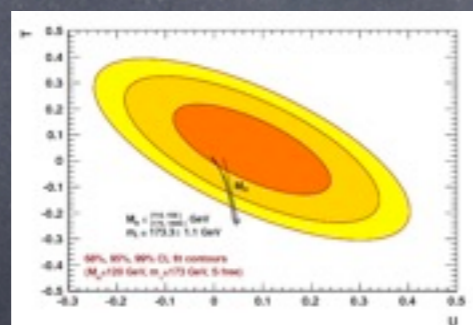
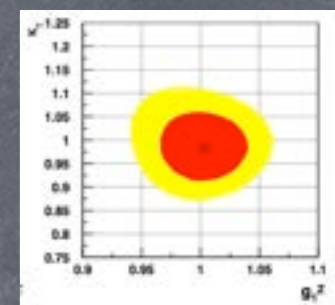
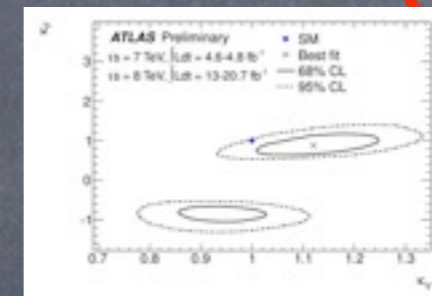
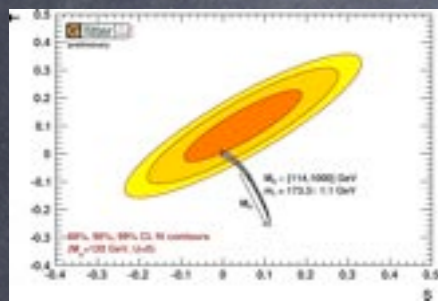
Pomarol, Gupta, Liu, Falkowski, Masso, Espinosa, Elias-Miro, Rattazzi, Biekötter, Knochel, Krämer  
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# Motivation

## Precision Searches for New Physics

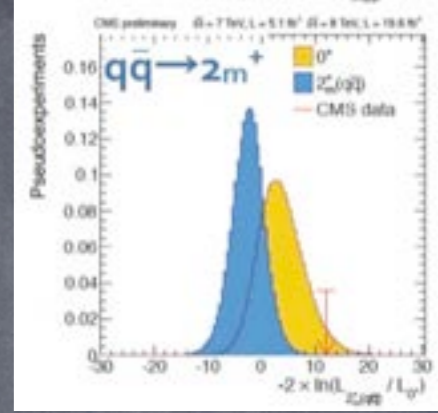
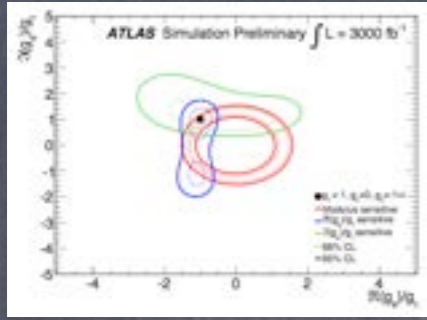


$\mathcal{L}^{SM}$

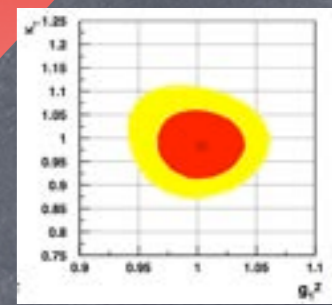
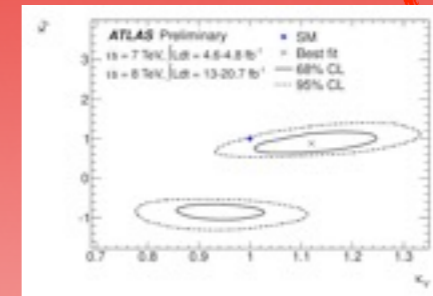
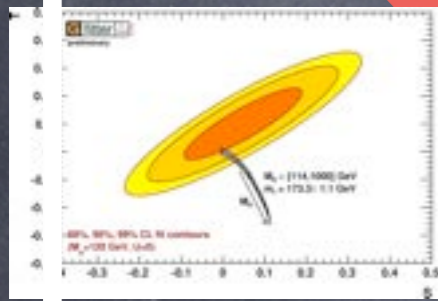


# Motivation

## Precision Searches for New Physics

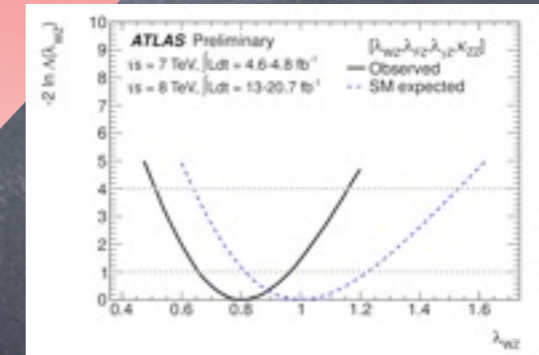
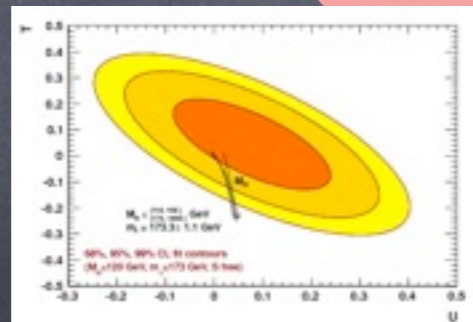


$\mathcal{L}^{SM}$

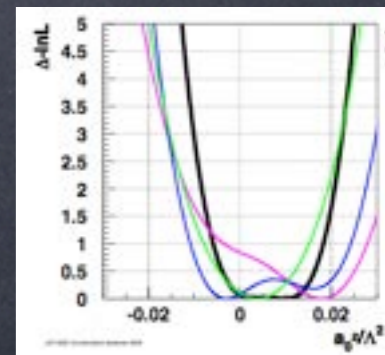


Expansion

- 1)  $E/\Lambda$
- 2)  $H/f$
- 3)  $Y_U, Y_D, Y_E$



$\mathcal{L}^{UV}$



# Motivation

1) No direct findings:  $M_{new}^i \sim \Lambda \gg m_W$

→ Expansion in  $D_\mu/\Lambda$

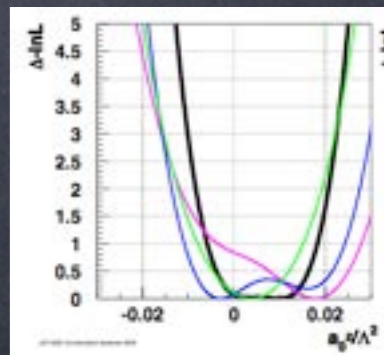
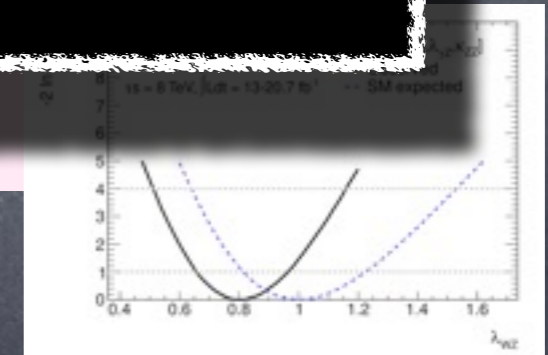
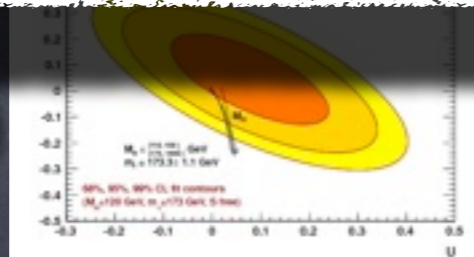
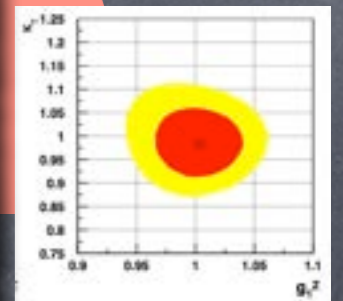
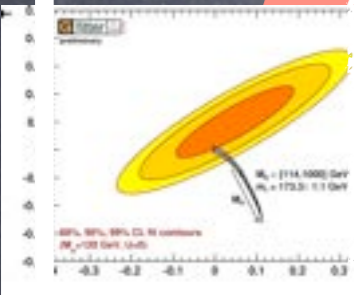
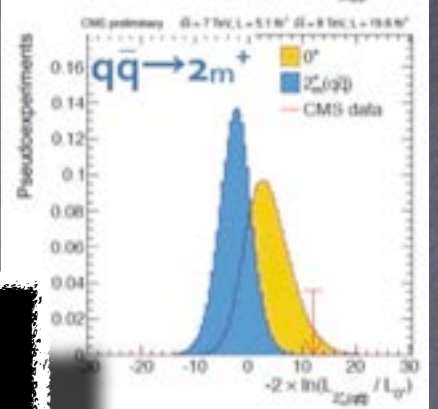
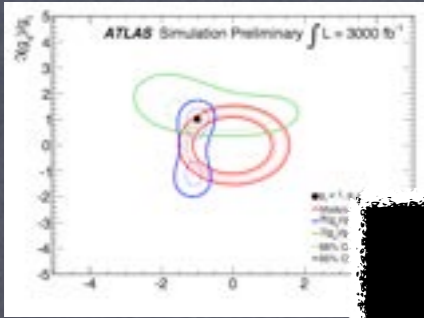
$\mathcal{L}^{SM} \equiv$

1)  $E/\Lambda$

2)  $H/f$

3)  $Y_U, Y_D, Y_E$

$\mathcal{L}^{UV}$



Expansion

# Motivation

2) BEH scalar is excitation around EWSB vacuum

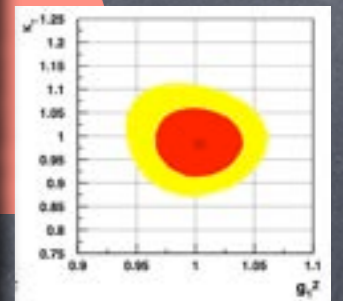
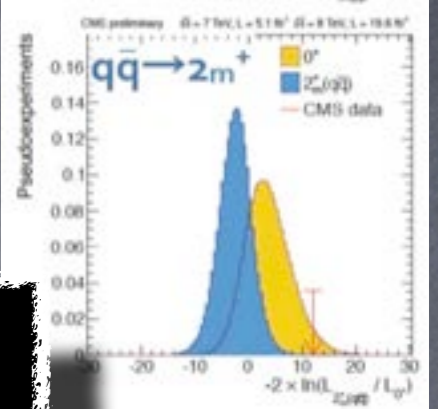
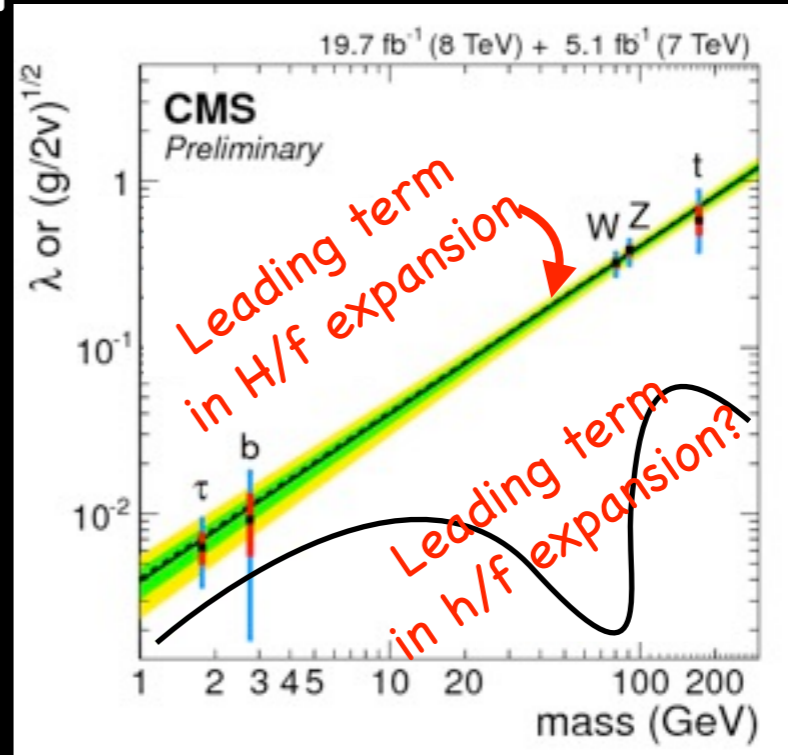
$$\mathcal{L}^{SM} \equiv$$

Expansion in  $H/f$

$v + h$

$$(f \equiv \Lambda/g_*)$$

Some BSM coupling (necessary, since fields have different weight in  $\hbar$  than derivatives) e.g. Cohen, Kaplan, Nelson '97



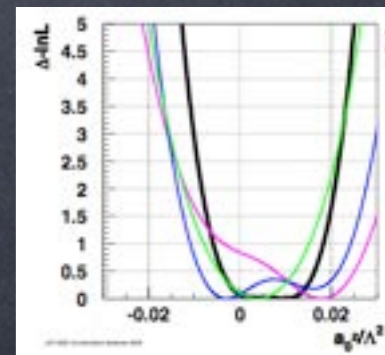
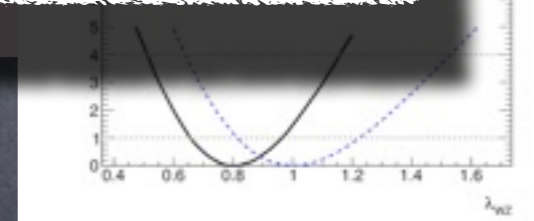
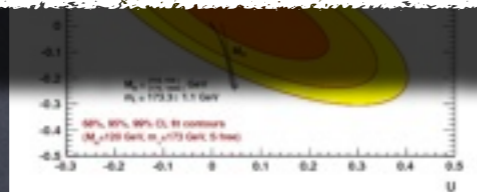
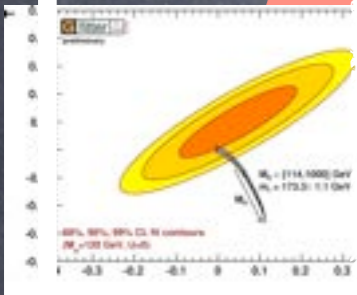
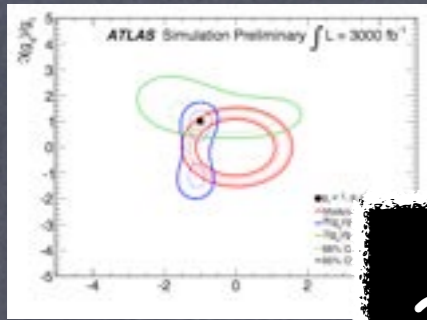
Expansion

1)  $E/\Lambda$

2)  $H/f$

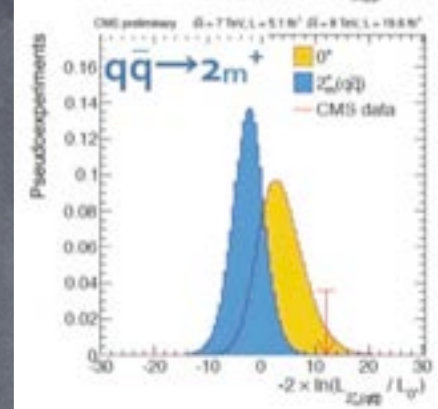
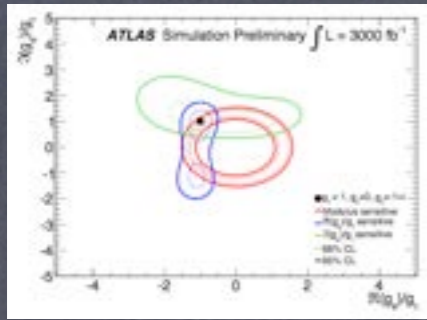
3)  $Y_U, Y_D, Y_E$

$$\mathcal{L}^{UV}$$

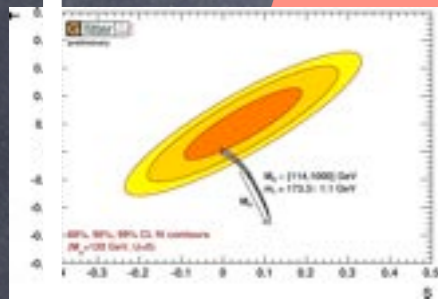


# Motivation

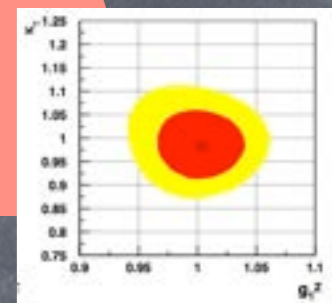
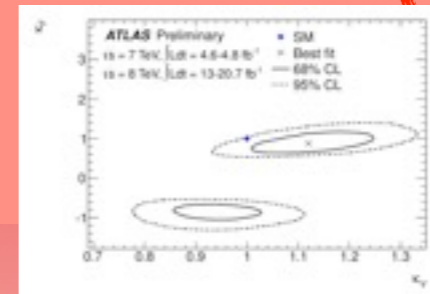
## Precision Searches for New Physics



$$\mathcal{L}^{SM} \equiv \mathcal{L}^4$$



$$\mathcal{L}^6$$

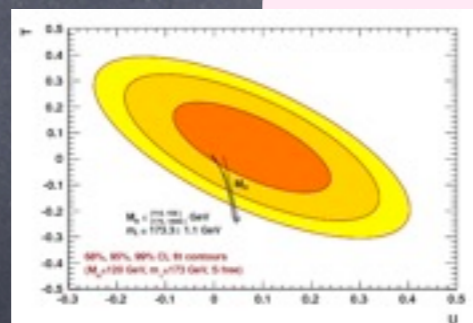


Expansion

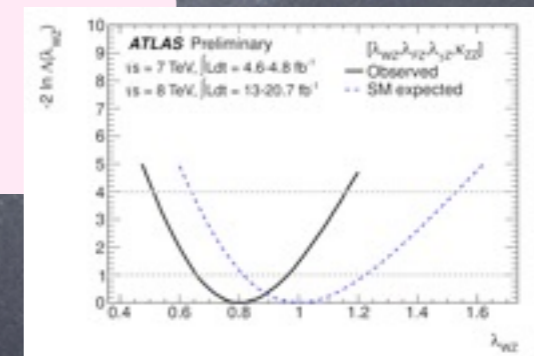
1)  $E/\Lambda$

2)  $H/f$

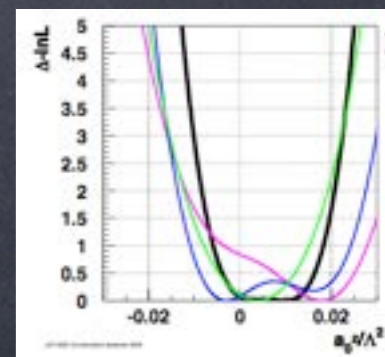
3)  $Y_U, Y_D, Y_E$



$$\mathcal{L}^8$$



$$\mathcal{L}^{UV}$$



# Motivation (short)

- Provide a meaningful parametrization of departures from SM (in form of Effective Field Theory - EFT)

BSM inspired: interpretable as search

- What are the most important parameters to search for?

- Where can the LHC provide genuine New Information?

# Motivation (short)

- Provide a meaningful parametrization of departures from SM (in form of Effective Field Theory - EFT)
  - BSM inspired: interpretable as search
- What are the most important parameters to search for?
- Where can the LHC provide genuine New Information?



# Motivation

$$\mathcal{L}_{\text{eff}} = \frac{\Lambda^4}{g_*^2} \mathcal{L} \left( \frac{D_\mu}{\Lambda}, \frac{g_* H}{\Lambda}, \frac{g_* f_{L,R}}{\Lambda^{3/2}}, \frac{g F_{\mu\nu}}{\Lambda^2} \right) \simeq \mathcal{L}_4 + \mathcal{L}_6 + \dots, \quad \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i$$

Buchmuller, Wyler '86;  
Giudice et al '07  
Grzadkowski et al '10  
Alonso et al '13

$$\mathcal{L}^{SM} \equiv \mathcal{L}^4$$

What defines SM?  
(from an experiment's point of view)

# Motivation

$$\mathcal{L}_{\text{eff}} = \frac{\Lambda^4}{g_*^2} \mathcal{L} \left( \frac{D_\mu}{\Lambda}, \frac{g_* H}{\Lambda}, \frac{g_* f_{L,R}}{\Lambda^{3/2}}, \frac{g F_{\mu\nu}}{\Lambda^2} \right) \simeq \mathcal{L}_4 + \mathcal{L}_6 + \dots, \quad \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i$$

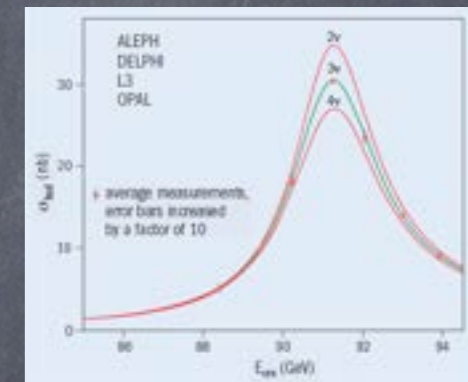
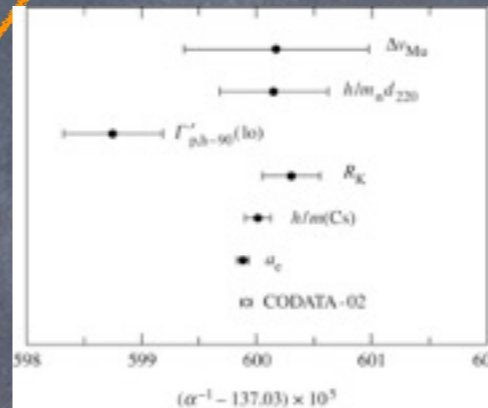
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Fixed by 19 most precise experiments

What defines SM?  
(from an experiment's point of view)

- Parameters: 19 in  $\mathcal{L}_4 \equiv \mathcal{L}_{SM}$



# Motivation

$$\mathcal{L}_{\text{eff}} = \frac{\Lambda^4}{g_*^2} \mathcal{L} \left( \frac{D_\mu}{\Lambda}, \frac{g_* H}{\Lambda}, \frac{g_* f_{L,R}}{\Lambda^{3/2}}, \frac{g F_{\mu\nu}}{\Lambda^2} \right) \simeq \mathcal{L}_4 + \mathcal{L}_6 + \dots, \quad \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i$$

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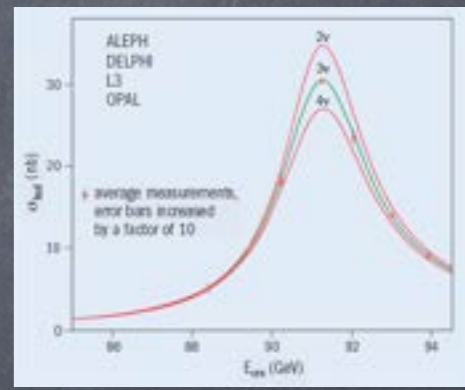
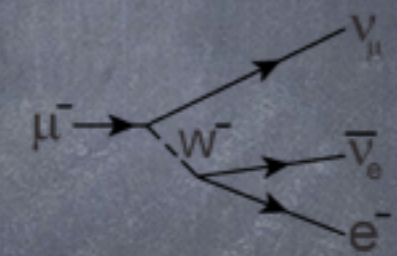
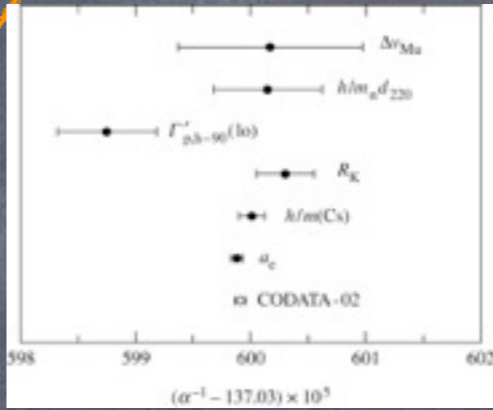
Fixed by 19 most precise experiments

## What defines SM?

(from an experiment's point of view)

- Parameters: 19 in  $\mathcal{L}_4 \equiv \mathcal{L}_{SM}$
- Accidental relations (due to d=4 Lagrangian)

e.g.  $m_W = m_Z \cos \theta_W$   
 $g_{h\bar{f}f} = m_f/v$



Predictions for other experiments

# Motivation

$$\mathcal{L}_{\text{eff}} = \frac{\Lambda^4}{g_*^2} \mathcal{L} \left( \frac{D_\mu}{\Lambda}, \frac{g_* H}{\Lambda}, \frac{g_* f_{L,R}}{\Lambda^{3/2}}, \frac{g F_{\mu\nu}}{\Lambda^2} \right) \simeq \mathcal{L}_4 + \mathcal{L}_6 + \dots, \quad \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i$$

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$$\mathcal{L}^{SM} \equiv \mathcal{L}^4$$

$$\mathcal{L}^{BSM} \simeq \mathcal{L}^6$$

What defines SM?

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 $g_{h\bar{f}f} = m_f/v$

What defines BSM?

- Parameters: 76 dimension-6 ops.
- Accidental relations ?

# Motivation

$$\mathcal{L}_{\text{eff}} = \frac{\Lambda^4}{g_*^2} \mathcal{L} \left( \frac{D_\mu}{\Lambda}, \frac{g_* H}{\Lambda}, \frac{g_* f_{L,R}}{\Lambda^{3/2}}, \frac{g F_{\mu\nu}}{\Lambda^2} \right) \simeq \mathcal{L}_4 + \mathcal{L}_6 + \dots, \quad \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i$$

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What defines SM?

- Parameters: 19
- Accidental relations  
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e.g.  $m_W = m_Z \cos \theta_W$   
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**This Talk: BEH Scalar PHYSICS  
(one family, CP conserving)**

What defines BSM?

- Parameters: ~~76~~<sup>17</sup> dimension-6 ops.
- Accidental relations ?

# PART 1

## 17 BSM Parameters:

(Counting independent dimension-6 terms  
that can affect BEH scalar physics\*)

\*=all Wilson coefficients evaluated at  $\mu \sim m_W$

For running to UV see e.g.

Elias-Miro, Espinosa, Masso, Pomarol'13; (Alonso, Grojean), Jenkins, Manohar, Trott'13, Elias-Miro, Grojean, Gupta, Marzocca'13

# Parameters for BSM: BEH-only

## BEH scalar Physics Only

$v$	$\leftarrow$	$\mathcal{O}_r =  H ^2 (D_\mu H)^\dagger (D^\mu H)$
$m_d$	$\leftarrow$	$\mathcal{O}_{y_d} = y_d  H ^2 \bar{Q}_L H d_R$
$m_e$	$\leftarrow$	$\mathcal{O}_{y_e} = y_e  H ^2 \bar{L}_L H e_R$
$m_u$	$\leftarrow$	$\mathcal{O}_{y_u} = y_u  H ^2 \bar{Q}_L \tilde{H} u_R$
$g_s$	$\leftarrow$	$\mathcal{O}_{GG} = \frac{g_s^2}{4}  H ^2 G_{\mu\nu}^A G^{A\mu\nu}$
$g'$	$\leftarrow$	$\mathcal{O}_{BB} = \frac{g'^2}{4}  H ^2 B_{\mu\nu} B^{\mu\nu}$
$g$	$\leftarrow$	$\mathcal{O}_{WW} = \frac{g^2}{4}  H ^2 W_{\mu\nu}^a W^{a\mu\nu}$
$m_h$	$\leftarrow$	$\mathcal{O}_6 = \lambda  H ^6$

In the vacuum  $\langle h \rangle = v$ , operators  $|H|^2 \times \mathcal{L}_{SM}$  only redefine SM parameters!  $\rightarrow$  Observable only in BEH-scalar physics!

$$\frac{1}{g_s^2} G_{\mu\nu} G^{\mu\nu} + \frac{|H|^2}{\Lambda^2} G_{\mu\nu} G^{\mu\nu} = \left( \frac{1}{g_s^2} + \frac{v^2}{\Lambda^2} \right) G_{\mu\nu} G^{\mu\nu} + h \frac{2v}{\Lambda^2} G_{\mu\nu} G^{\mu\nu} + \dots$$

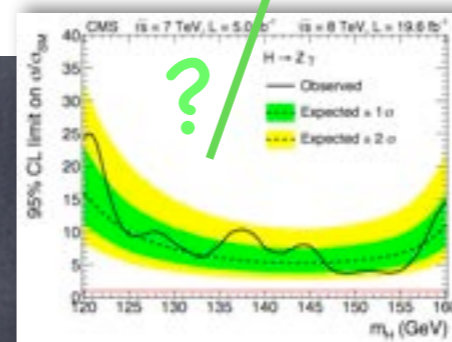
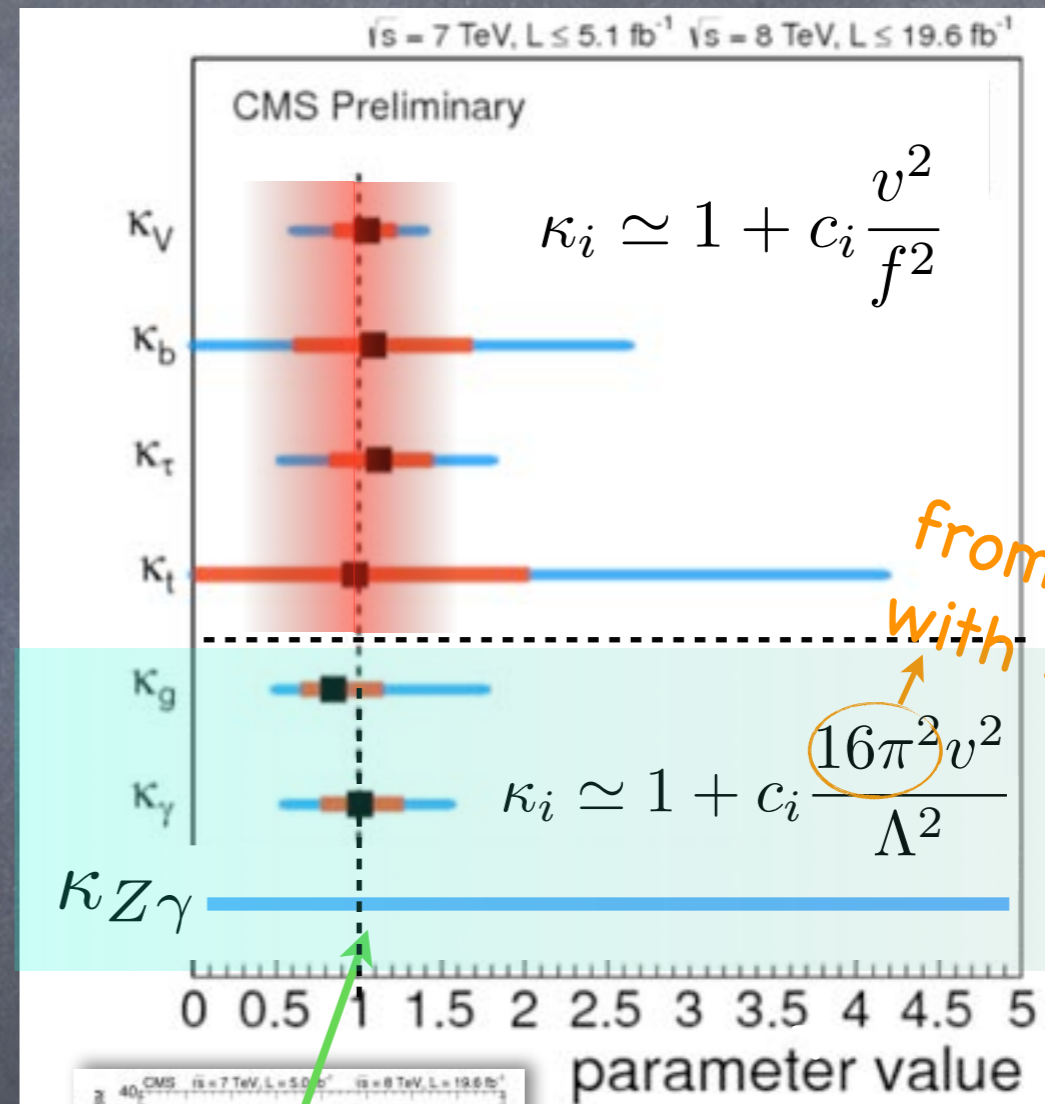
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$m_h$	$\mathcal{O}_6 = \lambda  H ^6$	

$\langle h \rangle = v$

$h^3?$



→ 8 Parameters fixed by BEH scalar physics experiments!



# Parameters for BSM: BEH+EW

BEH scalar Physics Only

$$\mathcal{O}_r = |H|^2 (D_\mu H)^\dagger (D^\mu H)$$

$$\mathcal{O}_{y_d} = y_d |H|^2 \bar{Q}_L H d_R$$

$$\mathcal{O}_{y_e} = y_e |H|^2 \bar{L}_L H e_R$$

$$\mathcal{O}_{y_u} = y_u |H|^2 \bar{Q}_L \tilde{H} u_R$$

$$\mathcal{O}_{GG} = \frac{g_s^2}{4} |H|^2 G_{\mu\nu}^A G^{A\mu\nu}$$

$$\mathcal{O}_{BB} = \frac{g'^2}{4} |H|^2 B_{\mu\nu} B^{\mu\nu}$$

$$\mathcal{O}_{WW} = \frac{g^2}{4} |H|^2 W_{\mu\nu}^a W^{a\mu\nu}$$

$$\mathcal{O}_6 = \lambda |H|^6$$

EW and BEH physics

$$\mathcal{O}_{WB} = \frac{gg'}{4} (H^\dagger \sigma^a H) W_{\mu\nu}^a B^{\mu\nu}$$

$$\mathcal{O}_T = \frac{1}{2} \left( H^\dagger \overleftrightarrow{D}_\mu H \right)^2$$

$$\mathcal{O}_R^u = (i H^\dagger \overleftrightarrow{D}_\mu H) (\bar{u}_R \gamma^\mu u_R)$$

$$\mathcal{O}_R^d = (i H^\dagger \overleftrightarrow{D}_\mu H) (\bar{d}_R \gamma^\mu d_R)$$

$$\mathcal{O}_R^e = (i H^\dagger \overleftrightarrow{D}_\mu H) (\bar{e}_R \gamma^\mu e_R)$$

$$\mathcal{O}_L^q = (i H^\dagger \overleftrightarrow{D}_\mu H) (\bar{Q}_L \gamma^\mu Q_L)$$

$$\mathcal{O}_L^{(3)q} = (i H^\dagger \sigma^a \overleftrightarrow{D}_\mu H) (\bar{Q}_L \sigma^a \gamma^\mu Q_L)$$

$$\mathcal{O}_L = (i H^\dagger \overleftrightarrow{D}_\mu H) (\bar{L}_L \gamma^\mu L_L)$$

$$\mathcal{O}_L^{(3)} = (i H^\dagger \sigma^a \overleftrightarrow{D}_\mu H) (\bar{L}_L \sigma^a \gamma^\mu L_L)$$

# Parameters for BSM: BEH+EW

In the vacuum  $\langle h \rangle = v$ , these operators can be measured!

7 of these operators modify:

$$Z_{\bar{\nu}\nu} \quad Z_{\bar{e}_L e_L} \quad Z_{\bar{e}_R e_R}$$

$$Z_{\bar{u}_L u_L} \quad Z_{\bar{u}_R u_R} \quad Z_{\bar{d}_L d_L} \quad Z_{\bar{d}_R d_R}$$

All tightly constrained by LEP1  
1/1000

EW and BEH physics

$$\mathcal{O}_{WB} = \frac{gg'}{4} (H^\dagger \sigma^a H) W_{\mu\nu}^a B^{\mu\nu}$$

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$$\mathcal{O}_L = (i H^\dagger \overleftrightarrow{D}_\mu H) (\bar{L}_L \gamma^\mu L_L)$$

$$\mathcal{O}_L^{(3)} = (i H^\dagger \sigma^a \overleftrightarrow{D}_\mu H) (\bar{L}_L \sigma^a \gamma^\mu L_L)$$

# Parameters for BSM: BEH+EW

In the vacuum  $\langle h \rangle = v$ , these operators can be measured!

EW and BEH physics

7 of these operators modify:

$$Z\bar{\nu}\nu \quad Z\bar{e}_L e_L \quad Z\bar{e}_R e_R$$

$$Z\bar{u}_L u_L \quad Z\bar{u}_R u_R \quad Z\bar{d}_L d_L \quad Z\bar{d}_R d_R$$

$\mathcal{O}_{WB} = \frac{gg'}{4} (H^\dagger \sigma^a H) W_{\mu\nu}^a B^{\mu\nu}$
$\mathcal{O}_T = \frac{1}{2} \left( H^\dagger \overleftrightarrow{D}_\mu H \right)^2$
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All tightly constrained by LEP1  
1/1000



Impact of these operators in BEH-physics is small

# Parameters for BSM: BEH+EW

In the vacuum  $\langle h \rangle = v$ , these operators can be measured!

2 of these modify TGCs:

$$g_Z^1 \quad K_\gamma$$

Hagiwara, Hikasa,  
Peccei, Zeppenfeld '87

EW and BEH physics

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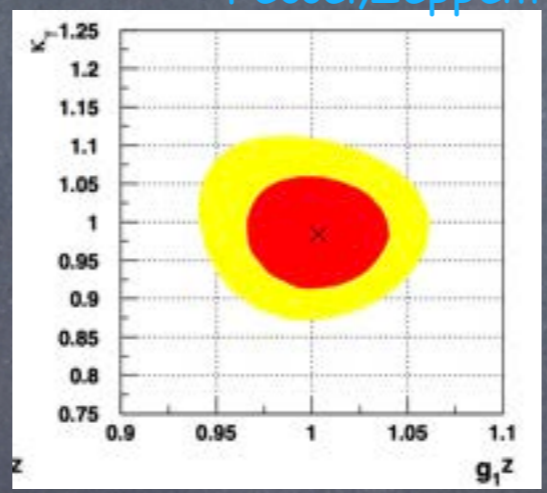
# Parameters for BSM: BEH+EW

In the vacuum  $\langle h \rangle = v$ , these operators can be measured!

2 of these modify TGCs:  $g_Z^1$   $K_\gamma$

Hagiwara, Hikasa, Peccei, Zeppenfeld '87

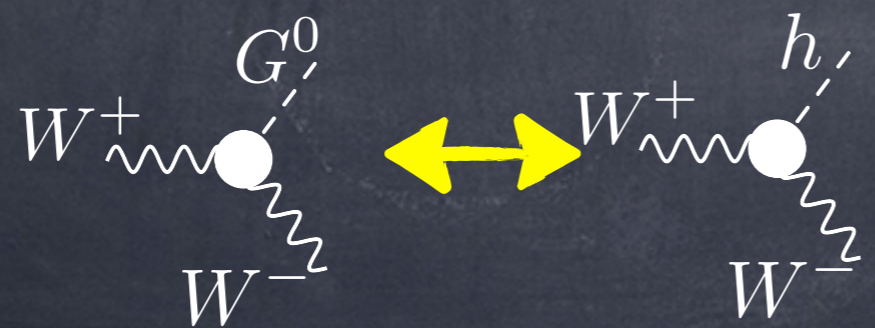
LEP2( $ee \rightarrow WW$ )  
constrained\*  $\sim 5/100$



## EW and BEH physics

$\mathcal{O}_{WB} = \frac{gg'}{4} (H^\dagger \sigma^a H) W_{\mu\nu}^a B^{\mu\nu}$
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Preview:



# Small Summary: Parameters

$\mathcal{O}_\tau =  H ^2 (D_\mu H)^\dagger (D^\mu H)$
$\mathcal{O}_{y_d} = y_d  H ^2 \bar{Q}_L H d_R$
$\mathcal{O}_{y_e} = y_e  H ^2 \bar{L}_L H e_R$
$\mathcal{O}_{y_u} = y_u  H ^2 \bar{Q}_L \tilde{H} u_R$
$\mathcal{O}_{GG} = \frac{g_s^2}{4}  H ^2 G_{\mu\nu}^A G^{A\mu\nu}$
$\mathcal{O}_{BB} = \frac{g'^2}{4}  H ^2 B_{\mu\nu} B^{\mu\nu}$
$\mathcal{O}_{WW} = \frac{g^2}{4}  H ^2 W_{\mu\nu}^a W^{a\mu\nu}$
$\mathcal{O}_6 = \lambda  H ^6$

$\mathcal{O}_{WB} = \frac{gg'}{4} (H^\dagger \sigma^a H) W_{\mu\nu}^a B^{\mu\nu}$
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$\kappa_V, \kappa_b, \kappa_\tau, \kappa_t, \kappa_G, \kappa_{\gamma\gamma}, \kappa_{Z\gamma}, \kappa_{h^3}$

$g_Z^1, \kappa_\gamma$

$\delta g_{ZeL}, \delta g_{ZeR}, \delta g_{Z\nu}, \delta g_{ZuL}, \delta g_{ZdL}, \delta g_{ZuR}, \delta g_{ZdR}$

Might as well use these as parameters, to keep relations between observables manifest!

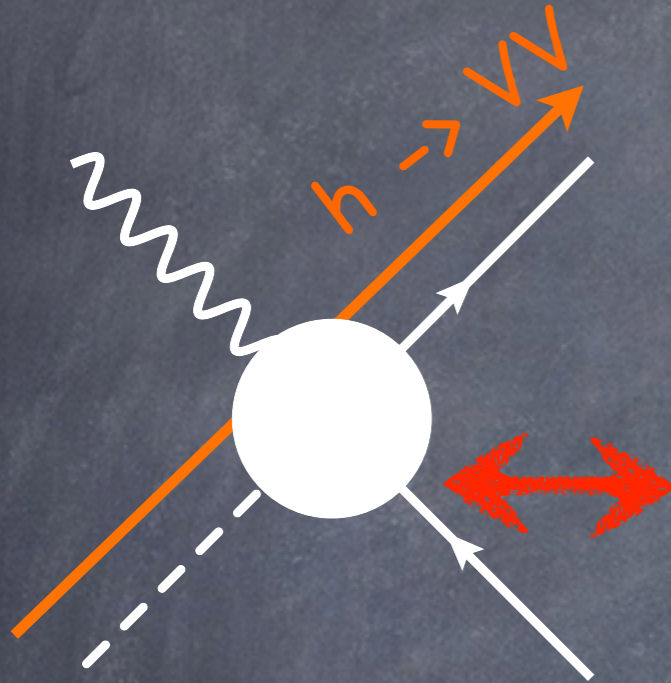
→ “BSM Primaries” & “Higgs-Basis”

PART 2  
Some Relations

# BSM Relations for Run 2

Deviations in different. distr. of  $h \rightarrow Z \bar{f} f$  or  $h \rightarrow W \bar{f} f$

See e.g. Isidori,(Manohar),Trott'13; Pomarol,FR'13; Falkowski,Vega-Morales'14

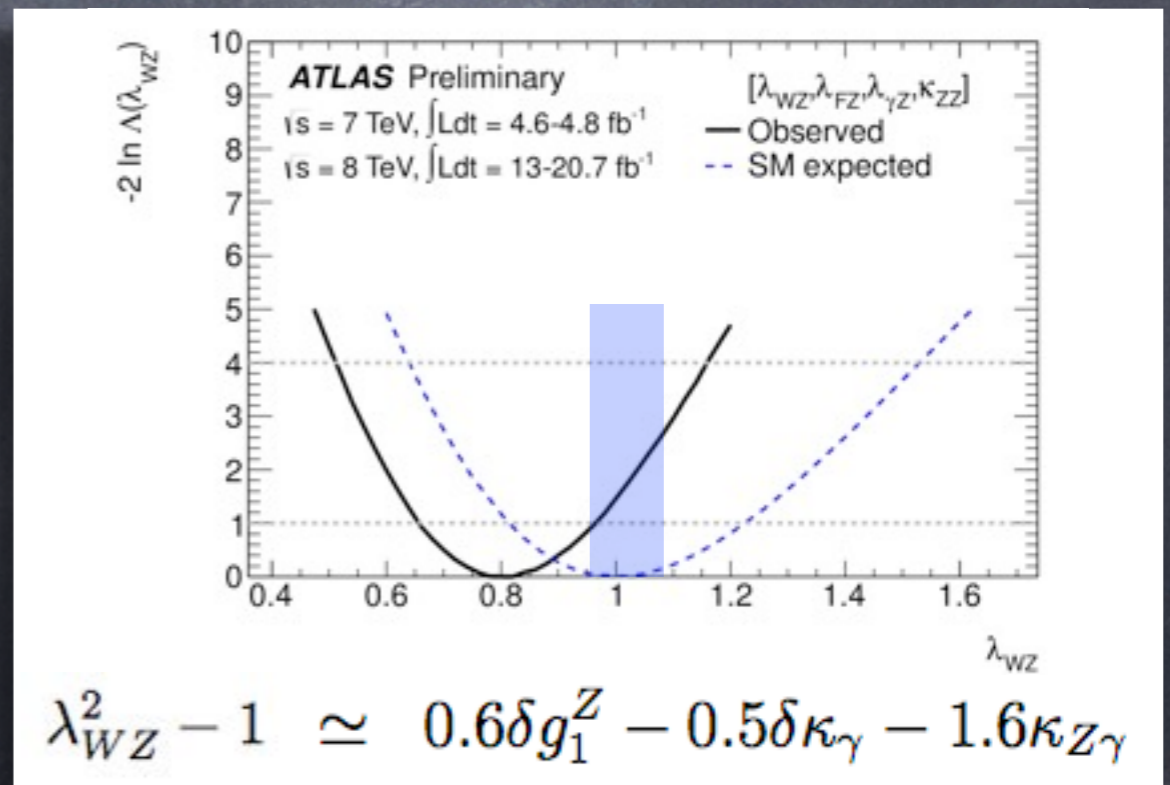
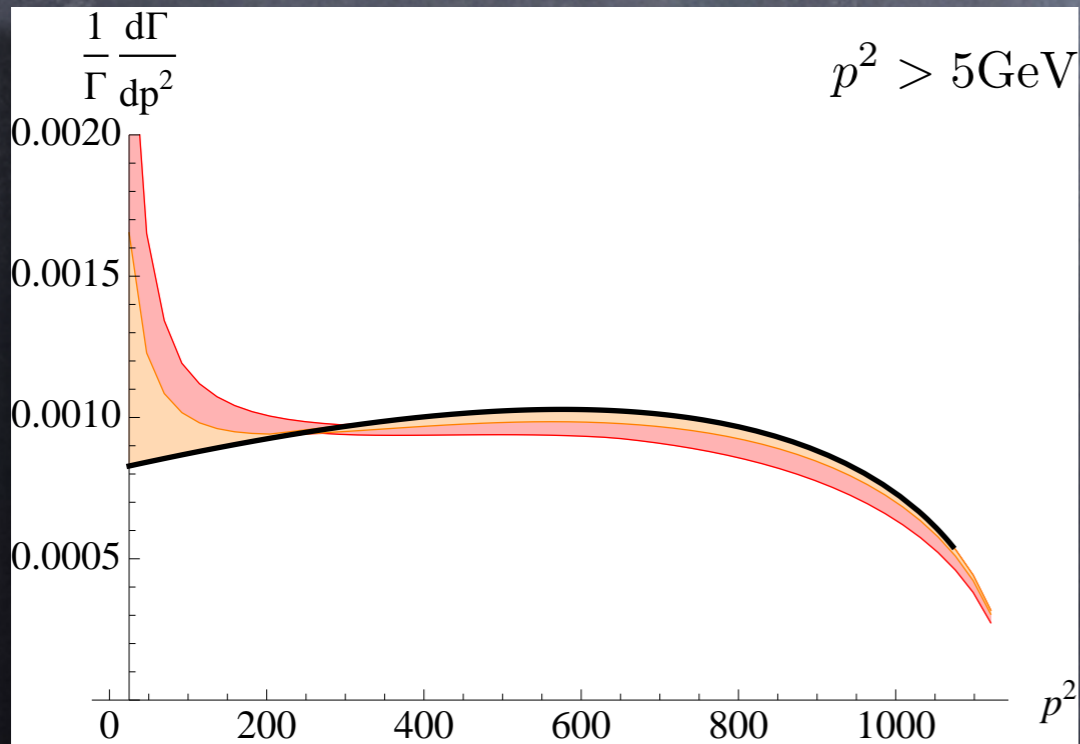
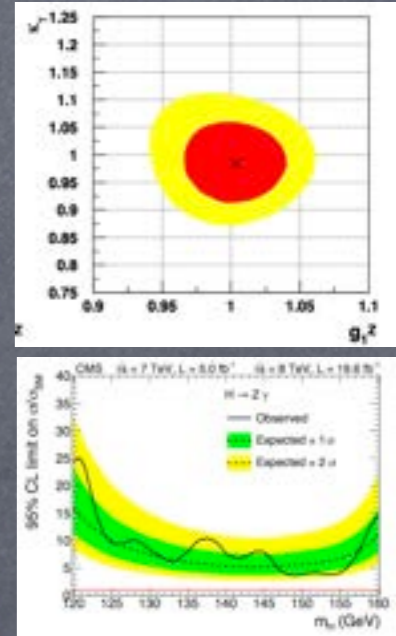


LEP 1

~~Related with  $Zff$  couplings~~

Related with Triple Gauge Coupling

Related with  $h \rightarrow Z\gamma, \gamma\gamma$

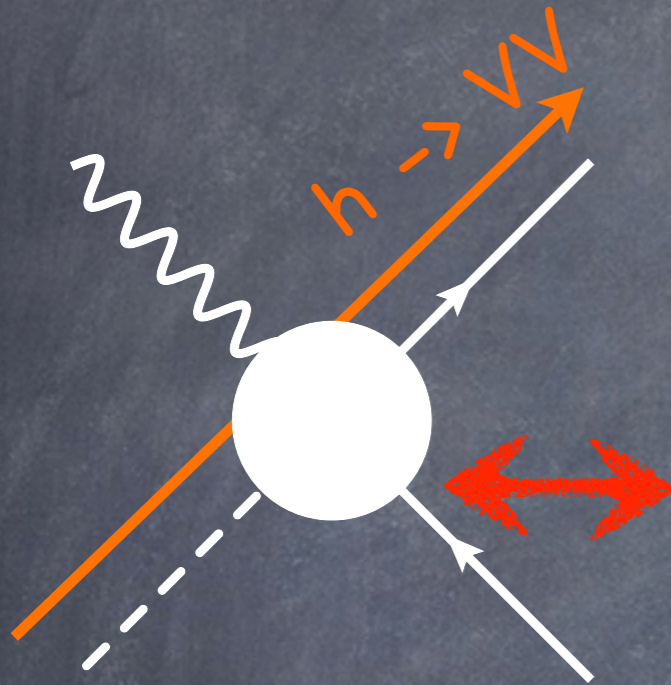




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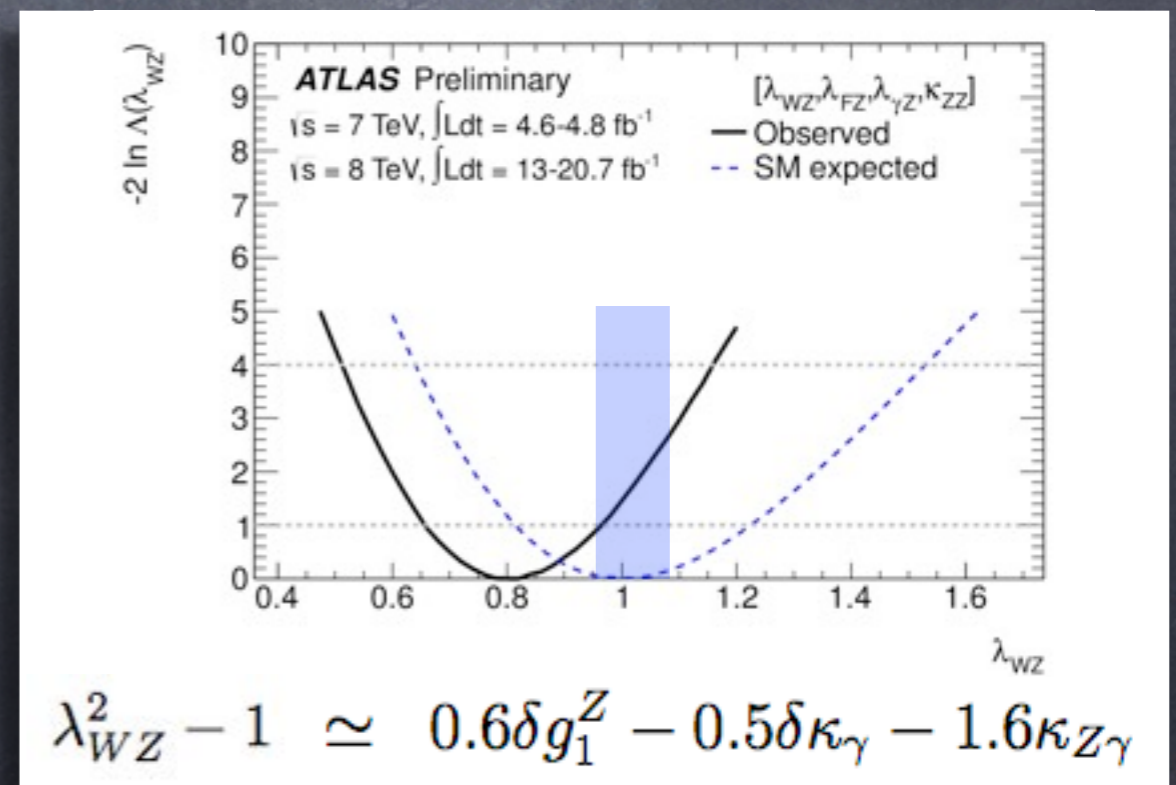
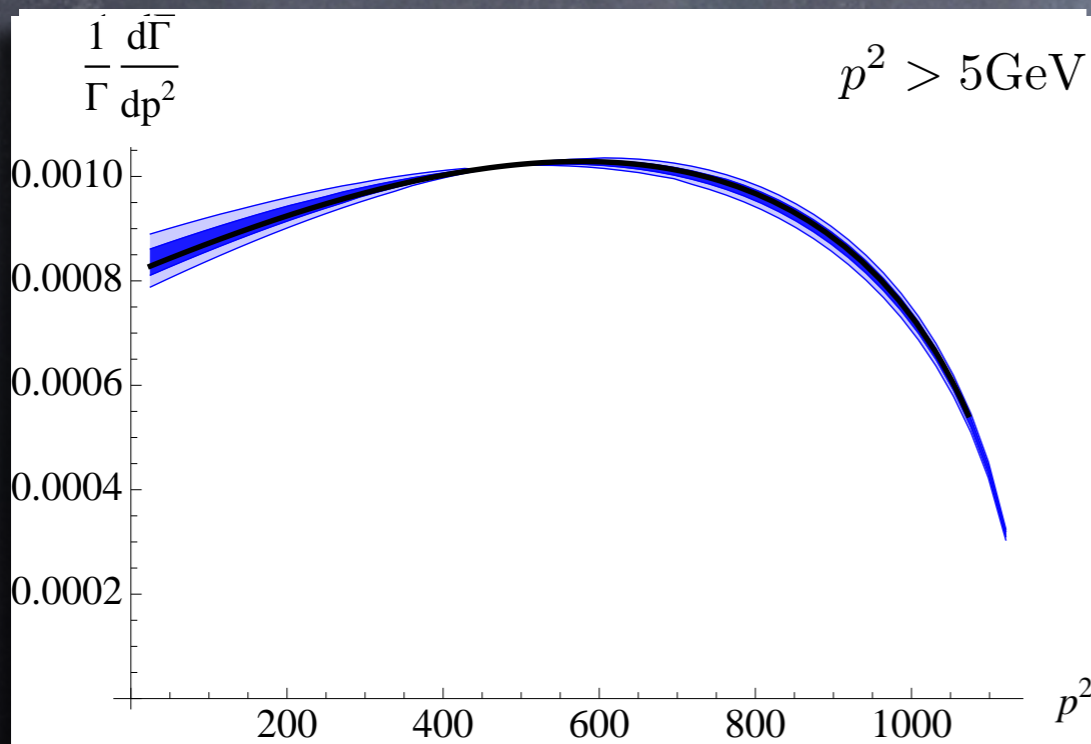
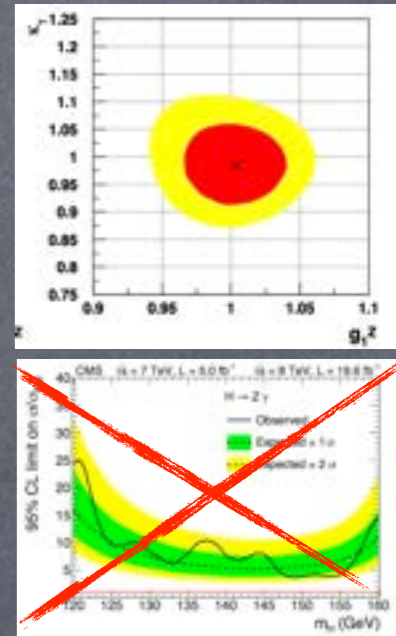


LEP 1

~~Related with Zff couplings~~

Related with Triple Gauge Coupling

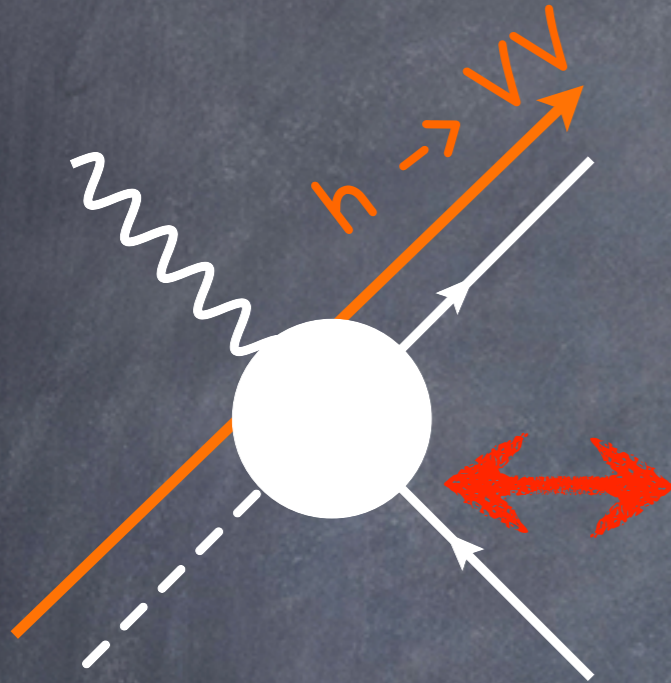
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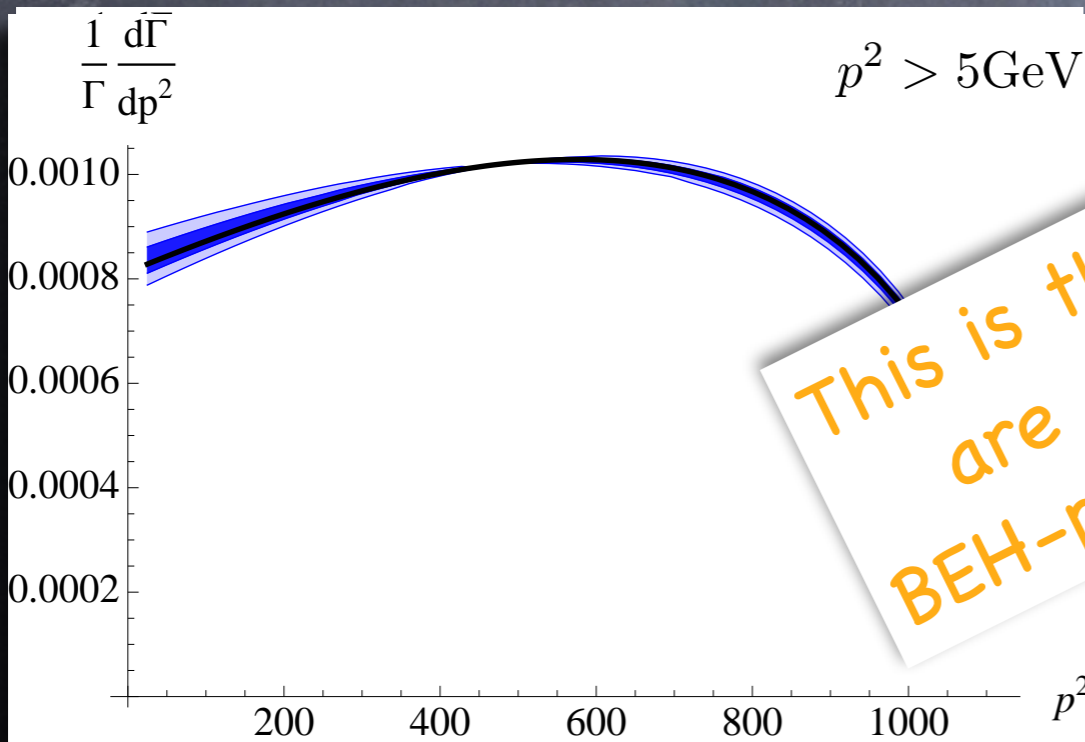
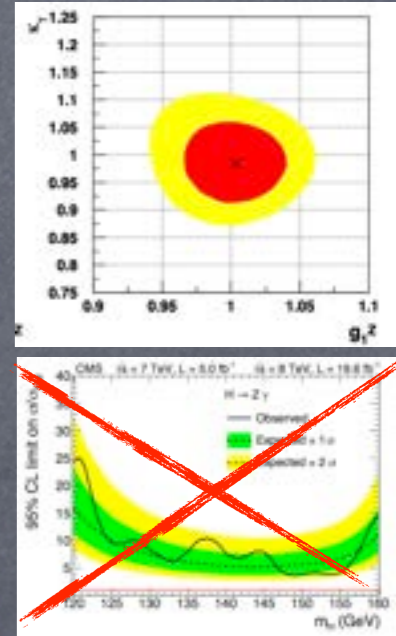


LEP 1

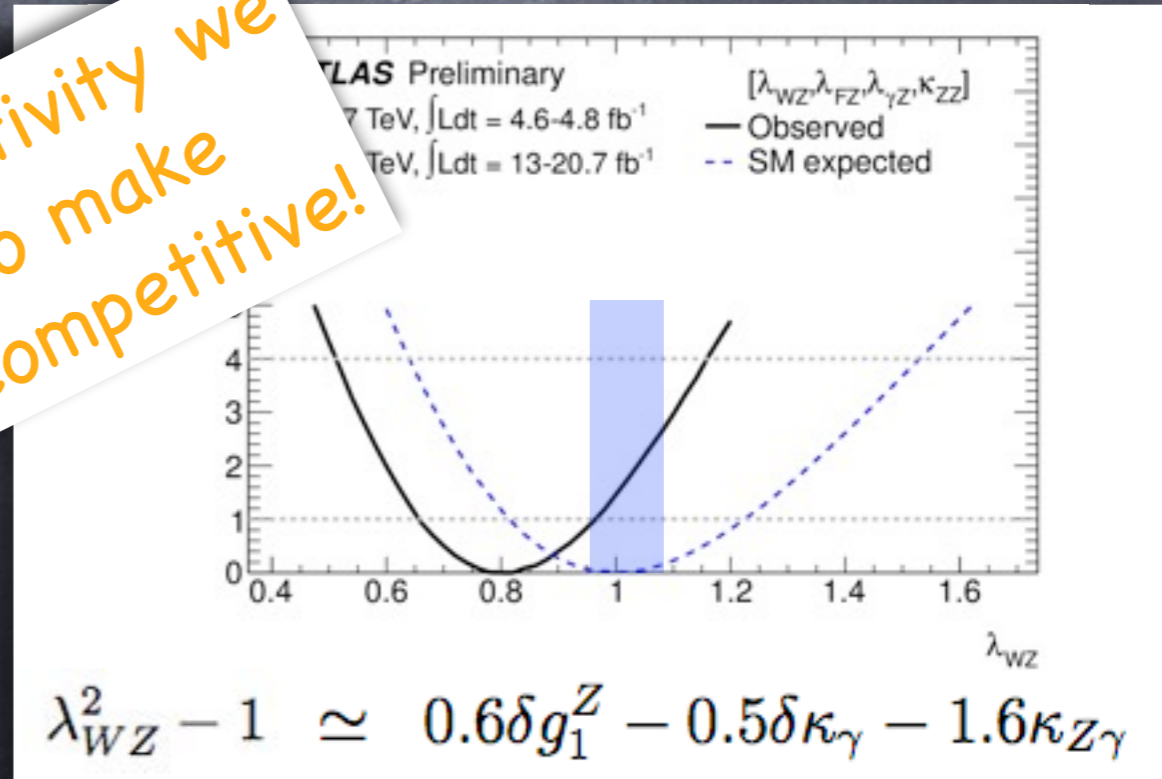
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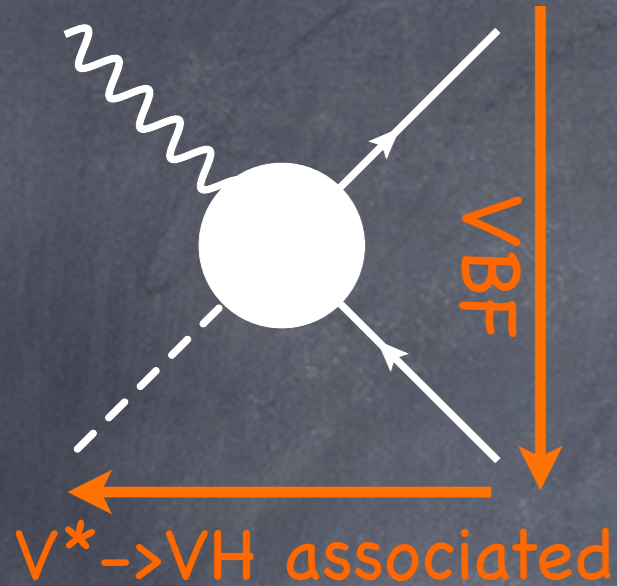


This is the sensitivity we are aiming to make BEH-physics competitive!



$$\lambda_{WZ}^2 - 1 \simeq 0.6\delta g_1^Z - 0.5\delta\kappa_\gamma - 1.6\kappa_{Z\gamma}$$

# ...but at High Energy LHC improves



Some BSM effects grow with energy:

$$\sim \int d \cos \theta |\mathcal{M}_L|^2 \xrightarrow{(s \gg m_W)} \frac{g^4}{6} \left( 1 + 4 \frac{\hat{s}}{g^2 \Lambda_{g1z}^2} + \dots \right)^2$$

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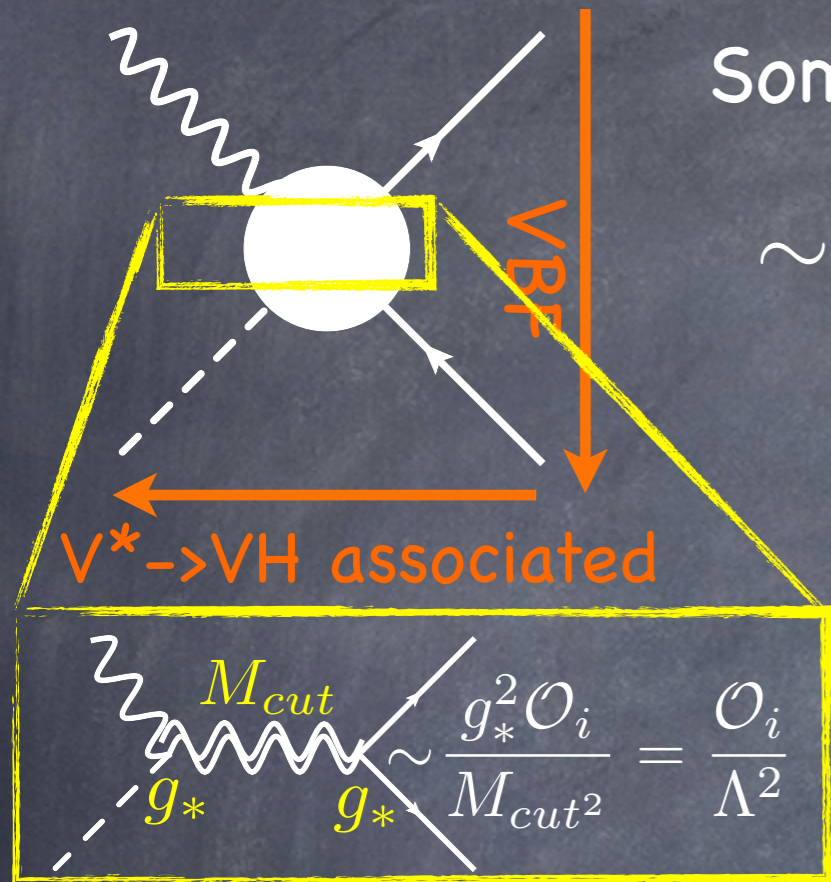
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→ Sensitivity to BSM Enhanced!

→ EFT must be used consistently: events with  $E > M_{cut}$  cannot be used (conservatively)

→ Information about  $\Lambda$  and the cut-off  $M_{cut}$  (or BSM coupling  $g_*$ ) necessary (similar to DM and EFT!)



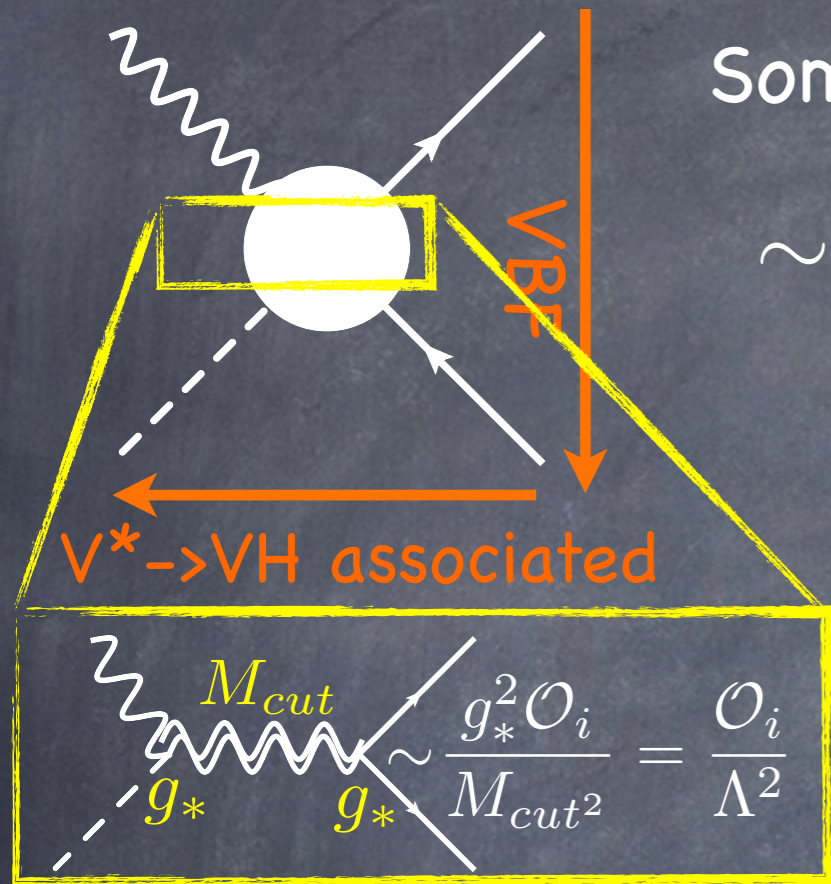
Giudice, Grojean, Pomarol, Rattazzi

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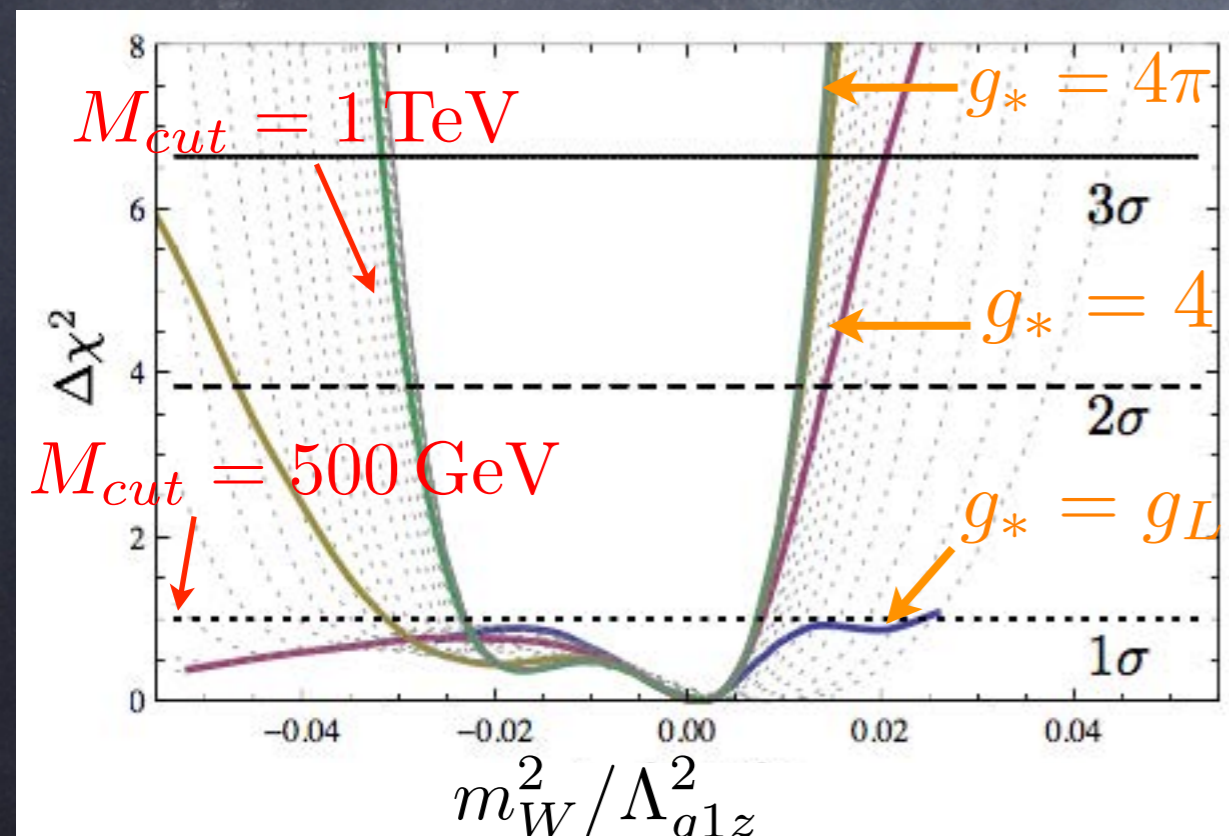
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Giudice, Grojean, Pomarol, Rattazzi



Very strongly coupled BSM

Large-N/Holography Composite Higgs

Weakly coupled  $Z', \dots$

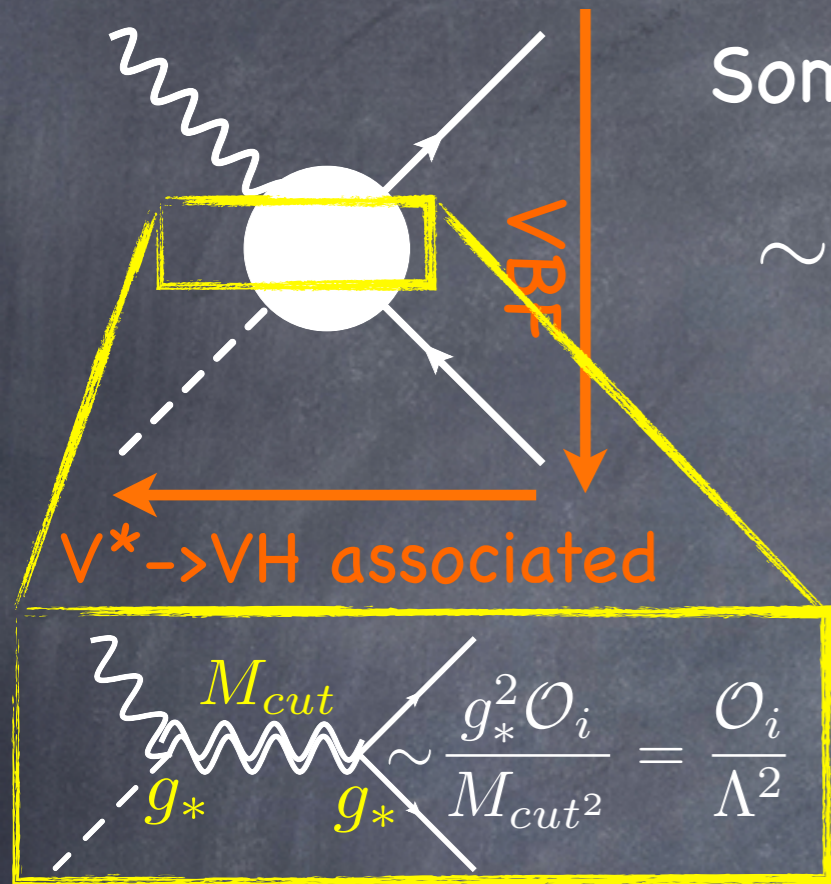
Biekötter, Knochel, Krämer, Liu, FR '14  
Isidori, Trott '13; Corbett, et al 12-13;  
Ellis, Sanz, You '14; Beneke, Boito, Wang '14

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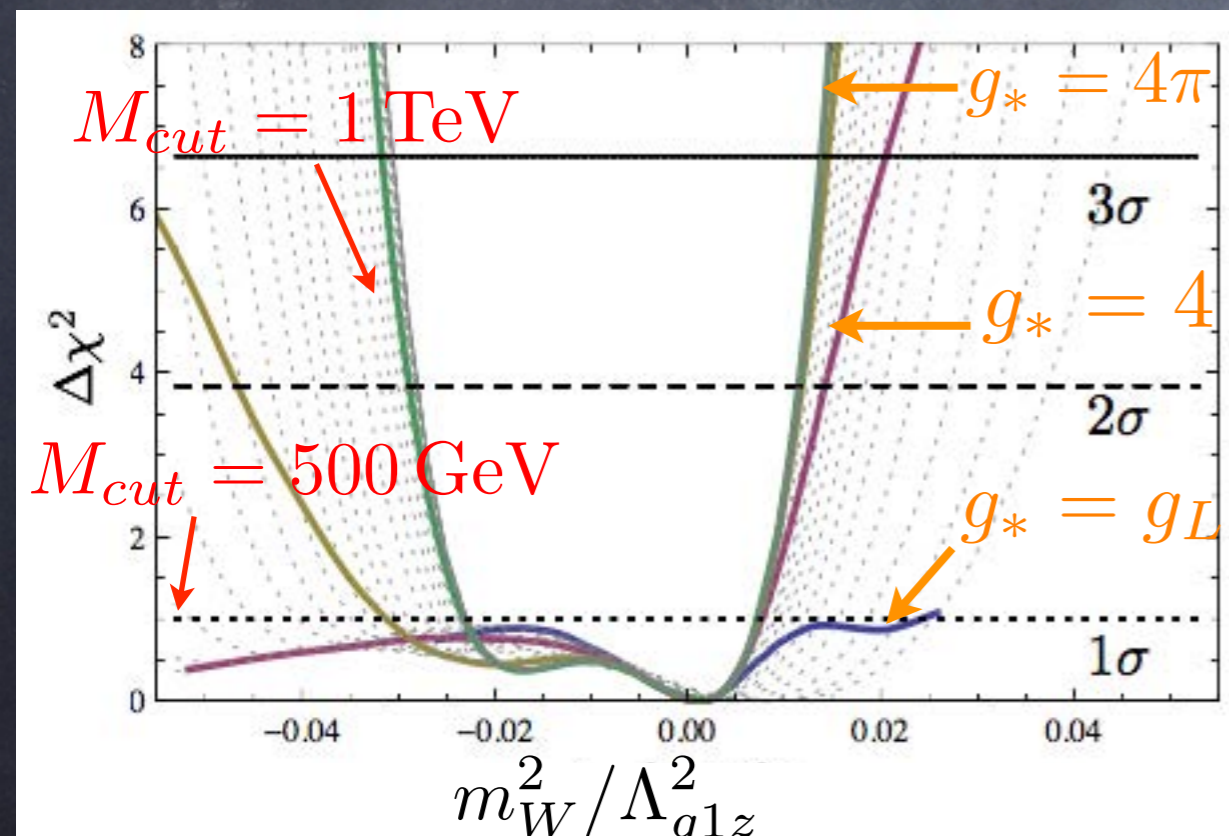
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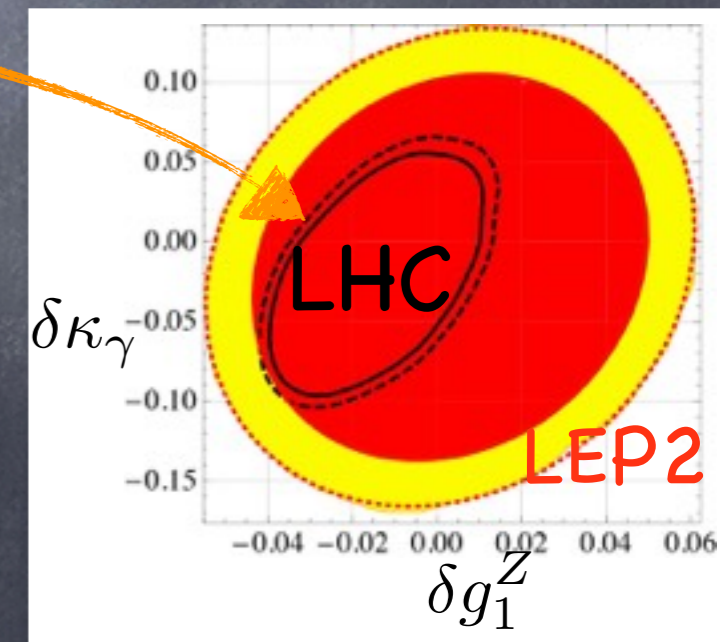
Giudice, Grojean, Pomarol, Rattazzi



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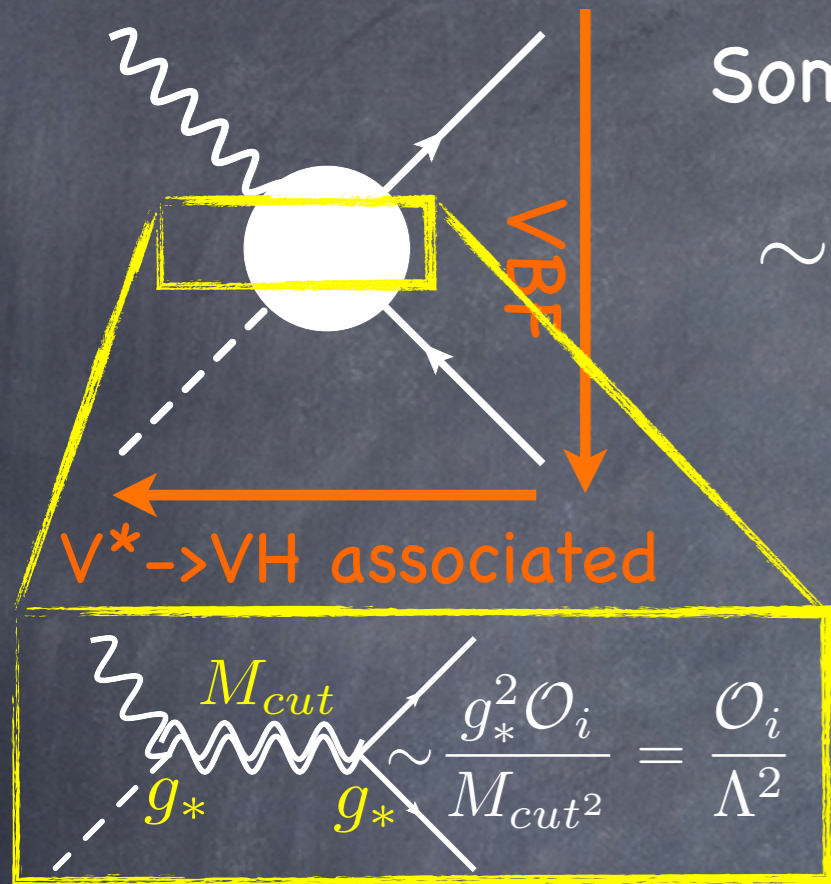
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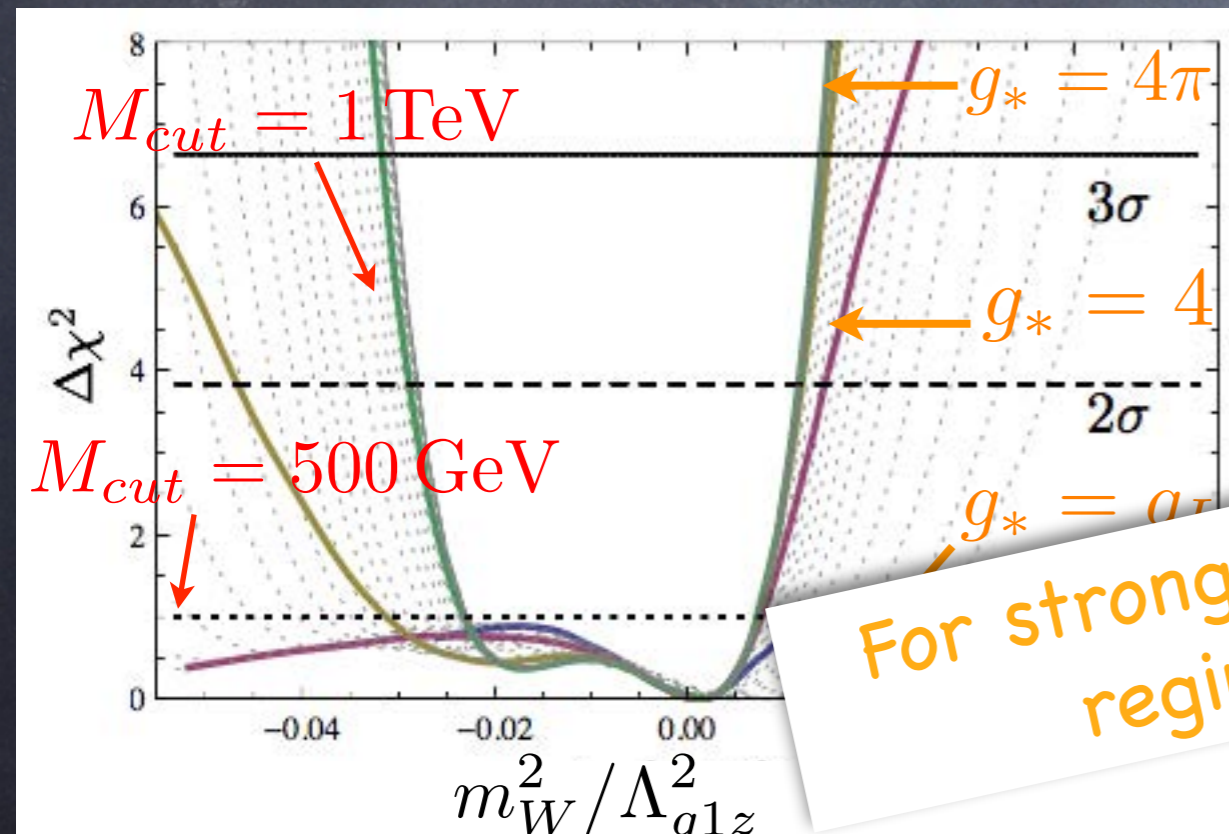
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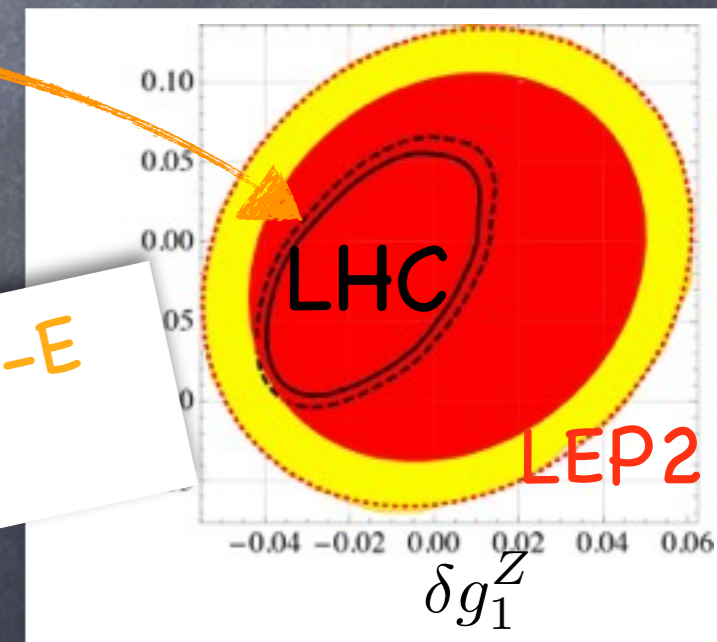
Giudice, Grojean, Pomarol, Rattazzi



Very strongly coupled BSM

Large-N/Holography Composite Higgs

For strongly coupled BSM, High-E regime very sensitive...



Biekötter, Knochel, Krämer, Liu, FR '14  
Isidori, Trott '13; Corbett, et al 12-13;  
Ellis, Sanz, You '14; Beneke, Boito, Wang '14

# Conclusions

BSM, early 2012:



**Anthropic**

(this dish doesn't contain human meat, otherwise you wouldn't be there tasting it)

**Higgsless**

(vegetarians only)

**Compositeness**

(shift-symmetric spaghetti)

**muSM**

(борщ)

**NMSSM**

(no beef, to solve the mu problem)

**SM**

(not very spicy)

**MSSM**

(low on calories - for the hierarchy problem)



# Conclusions

BSM, early 2015:



# Conclusions

BSM, early 2015:



Generic SM  
precision tests

# Conclusions

BSM, early 2015:



Generic SM  
precision tests



EFT: BSM Inspired  
precision searches  
("fare la scarpetta")

# Conclusions

- EFT: – Consistent framework to search for leading BSM effects  
– Motivation for precision tests

- Parametrization of BSM for Higgs physics:

7  $\{\delta g_{ZeL}, \delta g_{ZeR}, \delta g_{Z\nu}, \delta g_{ZuL}, \delta g_{ZdL}, \delta g_{ZuR}, \delta g_{ZdR}\}$  LEP1

2  $\{g_1^Z, \kappa_\gamma\}$  TGCs (LEP2)  
VH/VBF (LHC High-E)

8  $\{\kappa_{gg}, \kappa_{\gamma\gamma}, \kappa_{Z\gamma}, \delta g_{V_\mu V^\mu}^h, \delta g_{tt}^h, \delta g_{bb}^h, \delta g_{\bar{\tau}\tau}^h, \delta g_{h^3}^h\}$  BEH-physics (LHC)

→ Focusing on the most relevant parameters (combining different experiments) increases the sensitivity to new physics

- High-E regime valuable (at present only for strongly coupled BSM), but different assumptions on  $M_{cut}$  necessary for consistent constraints on  $\Lambda$