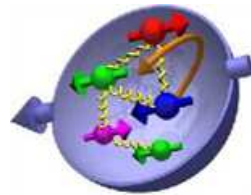


**Some surprises in  $e^+e^- \rightarrow$  Hadrons**  
**Baryon Timelike Form Factors close to threshold**  
**Non Breit-Wigner Charmonium Lineshape**



**Annual Meeting of the GRD PH-QCD**

December 15<sup>th</sup> -17<sup>th</sup> , 2014, Paris

# Outline

## ❑ Charged Baryon cross section at threshold:

- Jump at threshold (Coulomb Enhancement)
- Unexpected  $G(4 M_B^2) \sim 1$  (?)
- Unexpected flat cross section above threshold (no Sommerfeld Resumm?)
- Unexpected analyticity violation at threshold (?!)

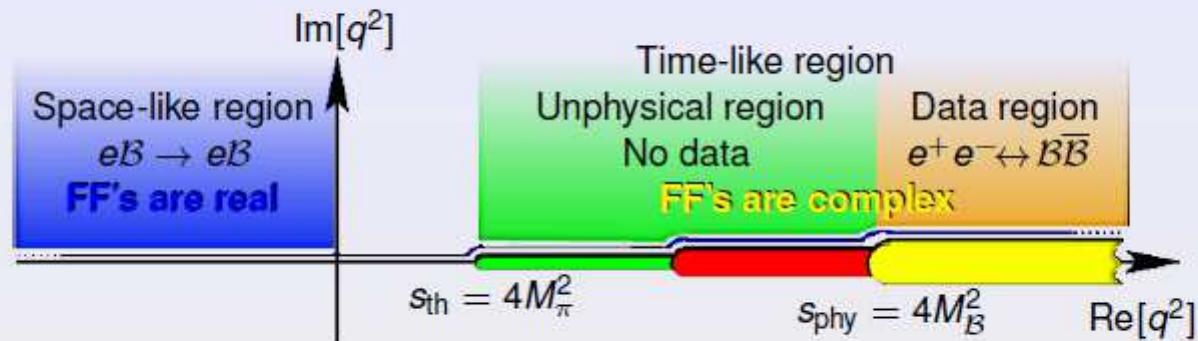
## ❑ Neutral Baryon cross section at threshold:

- Unexpected jump at threshold (Coulomb at quark level?) ?!

## ❑ Charmonium Lineshape:

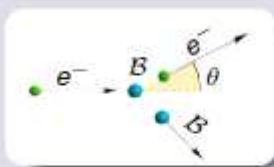
- Imaginary Decay Width
- A charmonium model for non Breit-Wigner lineshape
- A proposal for PANDA

# Cross sections and analyticity



Time-like: had. helicity =  $\begin{cases} 1 \Rightarrow |G_E| \\ 0 \Rightarrow |G_M| \end{cases}$

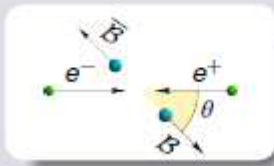
$G_E(4M_B^2) = G_M(4M_B^2)$



## Elastic scattering

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2 E_e' \cos^2 \frac{\theta}{2}}{4E_e^3 \sin^4 \frac{\theta}{2}} \left[ G_E^2 - \tau \left( 1 + 2(1 - \tau) \tan^2 \frac{\theta}{2} \right) G_M^2 \right] \frac{1}{1 - \tau}$$

$\tau = \frac{q^2}{4M_B^2}$



## Annihilation

### Coulomb correction

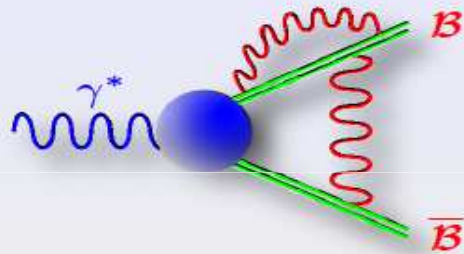
$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2 \beta C}{4q^2} \left[ (1 + \cos^2 \theta) |G_M|^2 + \frac{1}{\tau} \sin^2 \theta |G_E|^2 \right]$$

$\beta = \sqrt{1 - \frac{1}{\tau}}$

# The Coulomb Factor

## Coulomb effects predominant at threshold

Strong interactions: short range, while Coulomb long range:  
Coulomb acts on hadrons and pointlike Coulomb should be applied



$p\bar{p}$  Coulomb interaction as FSI  
[Sommerfeld, Sakharov, Schwinger, Fadin, Khoze]

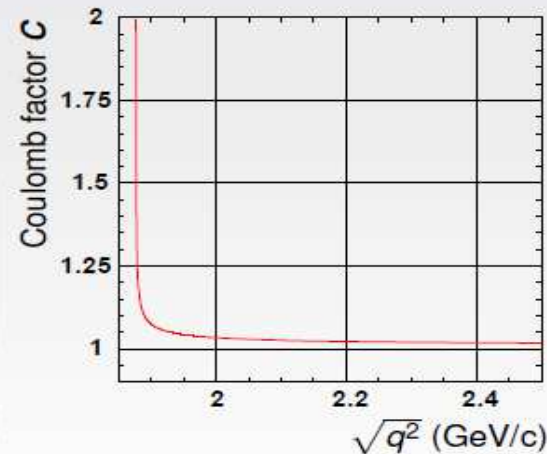
**Distorted wave approximation**

$$C = |\Psi_{\text{Coul}}(0)|^2$$

● S-wave:  $C = \frac{\frac{\pi\alpha}{\beta}}{1 - \exp\left(-\frac{\pi\alpha}{\beta}\right)} \xrightarrow{\beta \rightarrow 0} \frac{\pi\alpha}{\beta}$

● D-wave:  $C = 1$

**No Coulomb factor for boson pairs (P-wave)**



# Sommerfeld Enhancement and Resummation Factors

**Coulomb Factor  $\mathcal{C}$  for S-wave only:**

- Partial wave FF:  $G_S = \frac{2G_M \sqrt{q^2/4M^2} + G_E}{3}$      $G_D = \frac{G_M \sqrt{q^2/4M^2} - G_E}{3}$

- Cross section:  $\sigma(q^2) = 2\pi\alpha^2\beta \frac{4M^2}{(q^2)^2} [\mathcal{C} |G_S(q^2)|^2 + 2|G_D(q^2)|^2]$

$$\mathcal{C} = \mathcal{E} \times \mathcal{R}$$

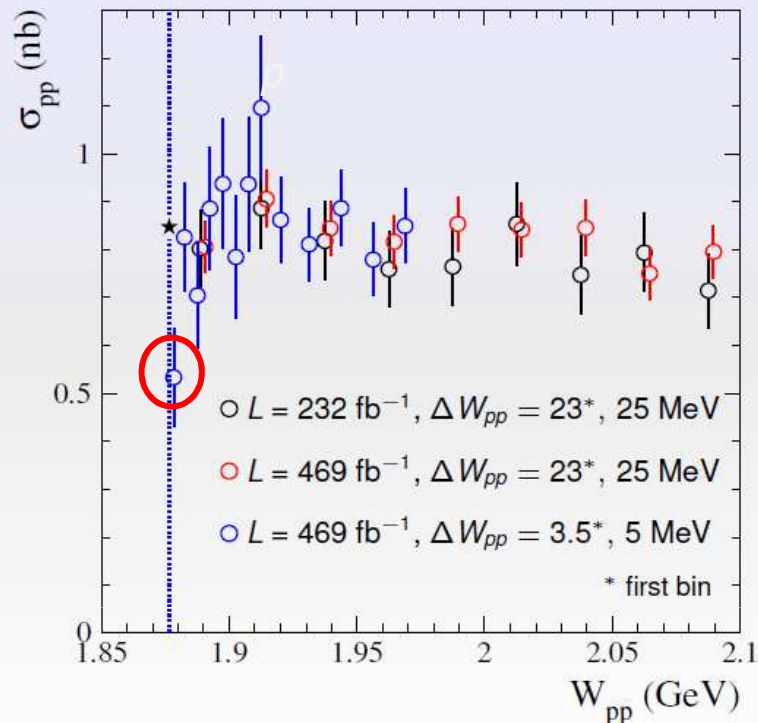
- Enhancement factor:  $\mathcal{E} = \pi\alpha/\beta$

- Step at threshold:  $\sigma(4M^2) = \frac{\pi^2\alpha^3}{2M^2} \frac{\cancel{\beta}}{\cancel{\beta}} |G_S(4M^2)|^2 = 0.85 |G_S(4M^2)|^2 \text{ nb}$

- Resummation factor:  $\mathcal{R} = 1/[1 - \exp(-\pi\alpha/\beta)]$

- Few MeV above threshold:  $\mathcal{C} \simeq 1 \Rightarrow \sigma(q^2) \propto \beta |G_S(q^2)|^2$

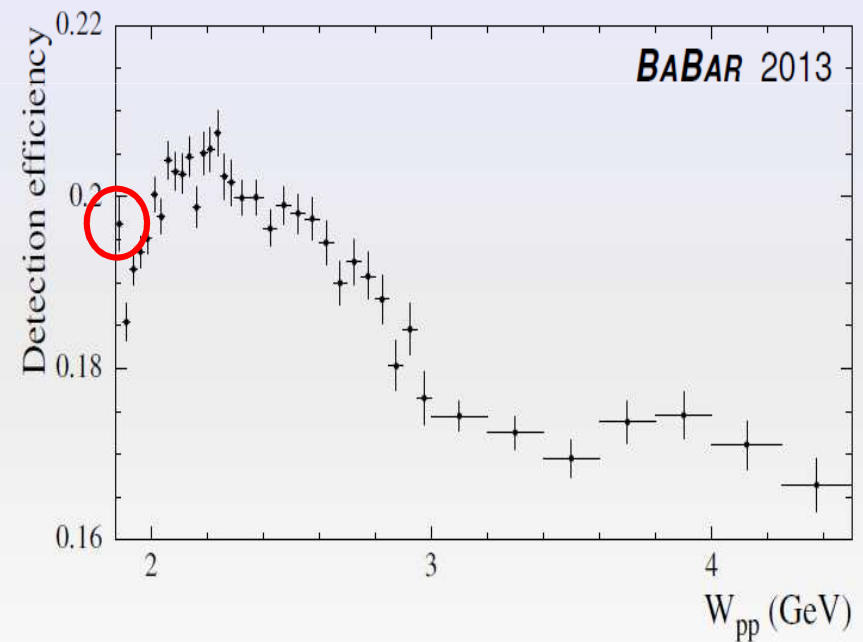
# BABAR 2013: $e^+ e^- \rightarrow p\bar{p}$ : jump at threshold



ISR : non vanishing  $\varepsilon$  at threshold  
 BaBar: indeed jump at threshold  
 but  $\varepsilon$  anomalous  
 as well as the cross section

$\Lambda_c$  : can check it at threshold

## $e^+ e^- \rightarrow p\bar{p}$ : efficiency



## Proton form factor at $q^2 = 4 M_p^2$

Extrapolating the flat cross section  
(neglecting for a while the very first point, **waiting** for  $\Lambda_c$ )

$$\sigma(e^+e^- \rightarrow p\bar{p})(4M_p^2) = 0.83 \pm 0.05 \text{ nb}$$

**BABAR**

$$\sigma(e^+e^- \rightarrow p\bar{p})(4M_p^2) = \frac{\pi^2 \alpha^3}{2M_p^2} \cancel{\beta} \cancel{\beta} |G^p(4M_p^2)|^2 = 0.85 |G^p(4M_p^2)|^2 \text{ nb}$$

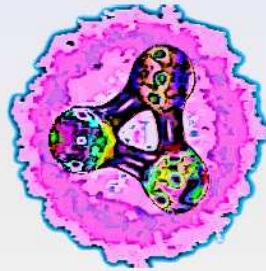
$$|G^p(4M_p^2)| \equiv 1$$

$$|G^p(4M_p^2)| = 0.99 \pm 0.04(\text{stat}) \pm 0.03(\text{syst})$$



# Proton form factor at $q^2 = 4 M_p^2$

$$|G^p(4M_p^2)| \equiv 1$$



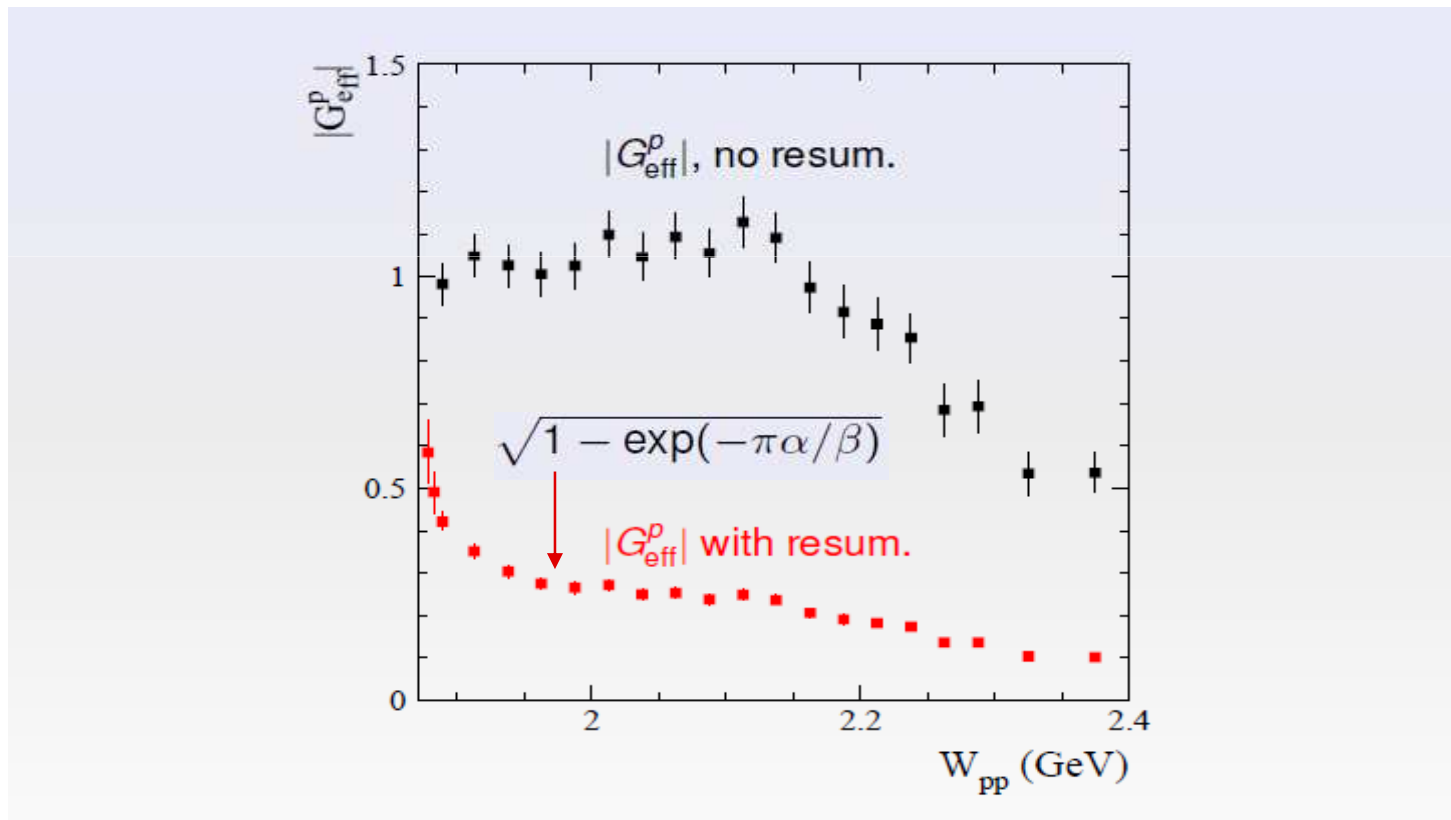
**At  $q^2 = 4M_p^2$  protons behave  
as pointlike fermions!**

**Sommerfeld  
Resummation Factor  
Needed?**



Sommerfeld pointlike R implies a rising cross section, while it is flat  
Hence  $G_{\text{eff}}$  sharp decrease is an artefact (just the inverse of R)?  
No narrow resonance below threshold ?

An explanation:  $\alpha_{\text{em}} \rightarrow \alpha_S$  (many gluons exchange, not only photons) ?



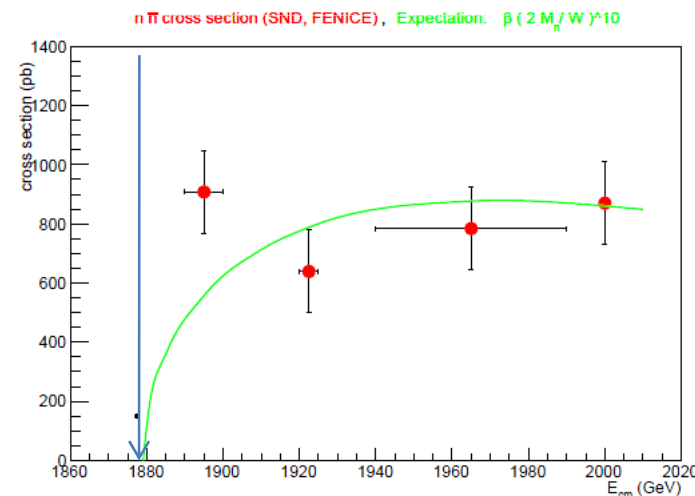
# Neutral Baryon cross sections non zero at threshold unlike the expectation

Cross section should vanish at threshold, due to vanishing phase space (cancelled by Coulomb enhancement factor if the baryon is charged ).

While it is  $\sim 450$  pb in the case of  $\Lambda \bar{\Lambda}$  (BESIII)  
 (might be)  $\sim 850$  pb in the case of  $n \bar{n}$  (SND2011-12, FENICE)

$\Lambda \bar{\Lambda}$   $n \bar{n}$  2012

Preliminary

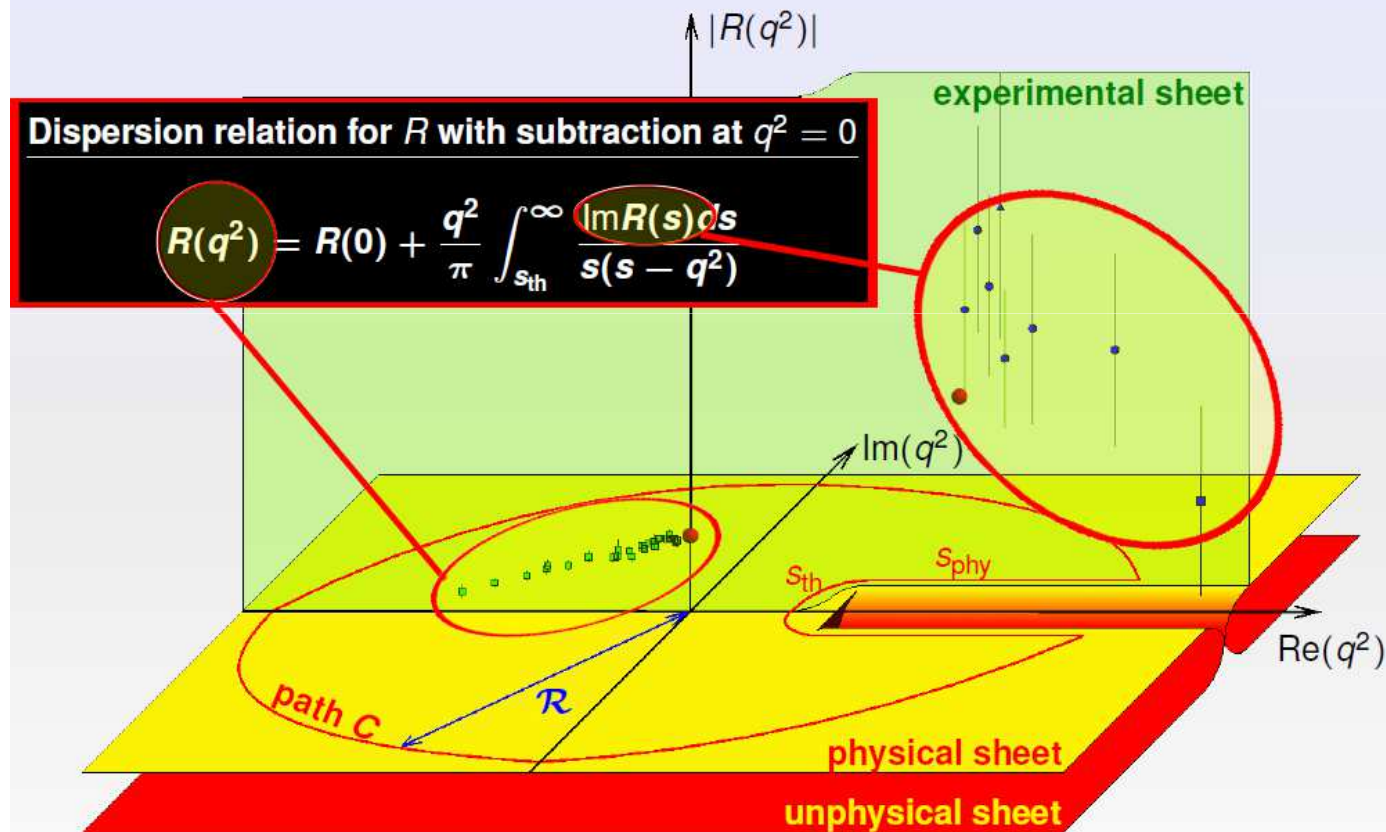


**Baryons at threshold do not behave as expected**

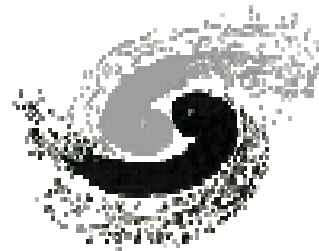
$R(q^2) = GE(q^2) / GM(q^2)$  from space to timelike

Uncontested assumption:  $R(4 M^2) = 1$ , but !!

$R(q^2)$  in the complex plane



# Analyticity Violation in $e^+ e^- \rightarrow B \bar{B}$ ?



December 16<sup>th</sup>, 2014, Paris

## Why to waste your time

- ❑ Always postulated that in  $e^+e^- \rightarrow$  Baryon Antibaryon at threshold :  
**angular distribution is isotropic, due to FF analiticity**
- ❑ Exactly at threshold in the c.m. difficult in the case of  $e^+e^- \rightarrow p\bar{p}$ 
  - PS 170 by means of  $p\bar{p} \rightarrow e^+e^-$  (atomic uncert.-> normalization)
  - BaBar by means of ISR (limited in statistics )
- ❑ SND  $e^+e^- \rightarrow n\bar{n}$  ( in principle closer to threshold than  $p\bar{p}$  )
- ❑ **Heavy baryons weak decay -> feasible at threshold (BESIII) !**
  - $\Lambda \bar{\Lambda}$  (BESIII: anomalous at threshold)
  - $\Lambda_c \bar{\Lambda}_c$  (BESIII tested  $J^P = 1/2^+$  )

# BaBar and BESIII present results

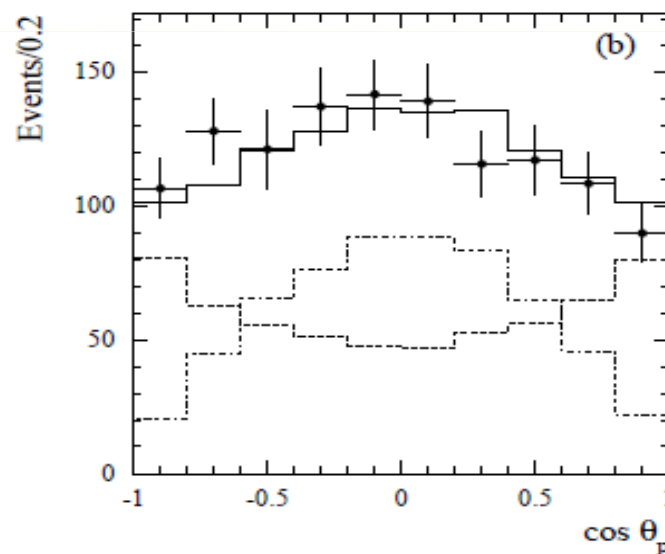
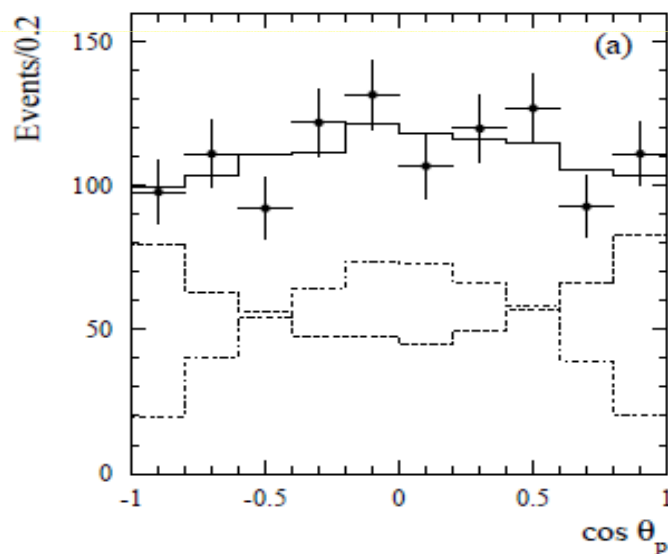
BESIII  $\Lambda_c$  and BaBar  $p\bar{p}$  statistics not yet enough,

But, together, a trend is pointed out

(a baryon common feature? Quantitatively not necessarily the same)

BaBar:  $W = 1877\text{--}1950\text{ MeV}$   
 $\beta = 0.20$

$W = 1950\text{--}2025\text{ MeV}$   
 $\beta = 0.33$



VERY PRELIMINARY



$\pi\alpha$   
Coulomb  
dominance



**Why it *should* be isotropic :  $G_E(4M^2_{\Lambda c}) = G_M(4M^2_{\Lambda c})$**

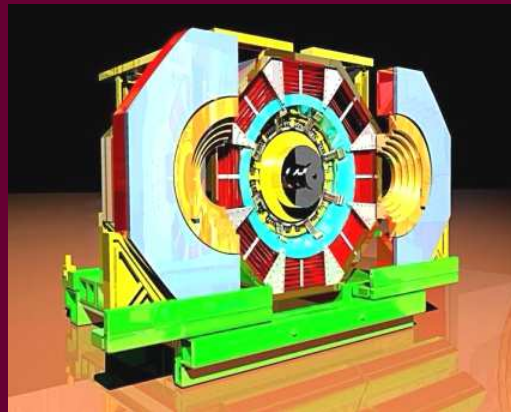
- $J_\mu \sim F_D \gamma_\mu + F_P i \sigma_{\mu\nu} q_\nu / (2M_B)$
- (lowest order QED :  $F_D = 1$  ,  $F_P = 0$ )
- Time-like: Outgoing Baryon spin antiparallel:  $G_E$   
Outgoing Baryon spin parallel:  $G_M$
- $\mathbf{G}_E = \mathbf{F}_D + \mathbf{F}_P / \tau$ ,  $\tau = (2M_B)^2 / Q^2$  (time-like  $Q^2 > 0$ )  
 $\mathbf{G}_M = \mathbf{F}_D + \mathbf{F}_P$
- $d\sigma/d\cos\theta = \pi \alpha^2 / (2 M_B^2) \beta C_0 [ |\mathbf{G}_M|^2 (1+\cos^2\theta) + \tau |\mathbf{G}_E|^2 \sin^2\theta ]$
- **Standard understanding: at threshold ( $\tau = 1$ ) :**
  - $\mathbf{G}_E = \mathbf{G}_M = \mathbf{F}_D + \mathbf{F}_P$
  - **Isotropy. S wave only**

## Why it *could* be anisotropic : $G_E(4M_{\Lambda c}) \neq G_M(4M_{\Lambda c})$

- Assume  $G_E(4M_{\Lambda c}) \neq G_M(4M_{\Lambda c})$ , always possible to define  $F_D$  and  $F_P$  so that  $G_E = F_D + F_P / \tau$ ,  $G_M = F_D + F_P$
  - But  $F_D$  and  $F_P$  no more analytic (continuous) through the threshold
    - $F_D = (G_M - G_E) / (\tau - 1)$
    - $F_P = (\tau G_E - G_M) / (\tau - 1)$
- **$F_D$  and  $F_P$  not analytic** equivalent to  $G_E(4M_{\Lambda c}) \neq G_M(4M_{\Lambda c})$ ,
- **Coulomb interactions not analytic at threshold:**  
 $|G_S^{\text{Coulomb}}|^2 \sim \pi \alpha / \beta$

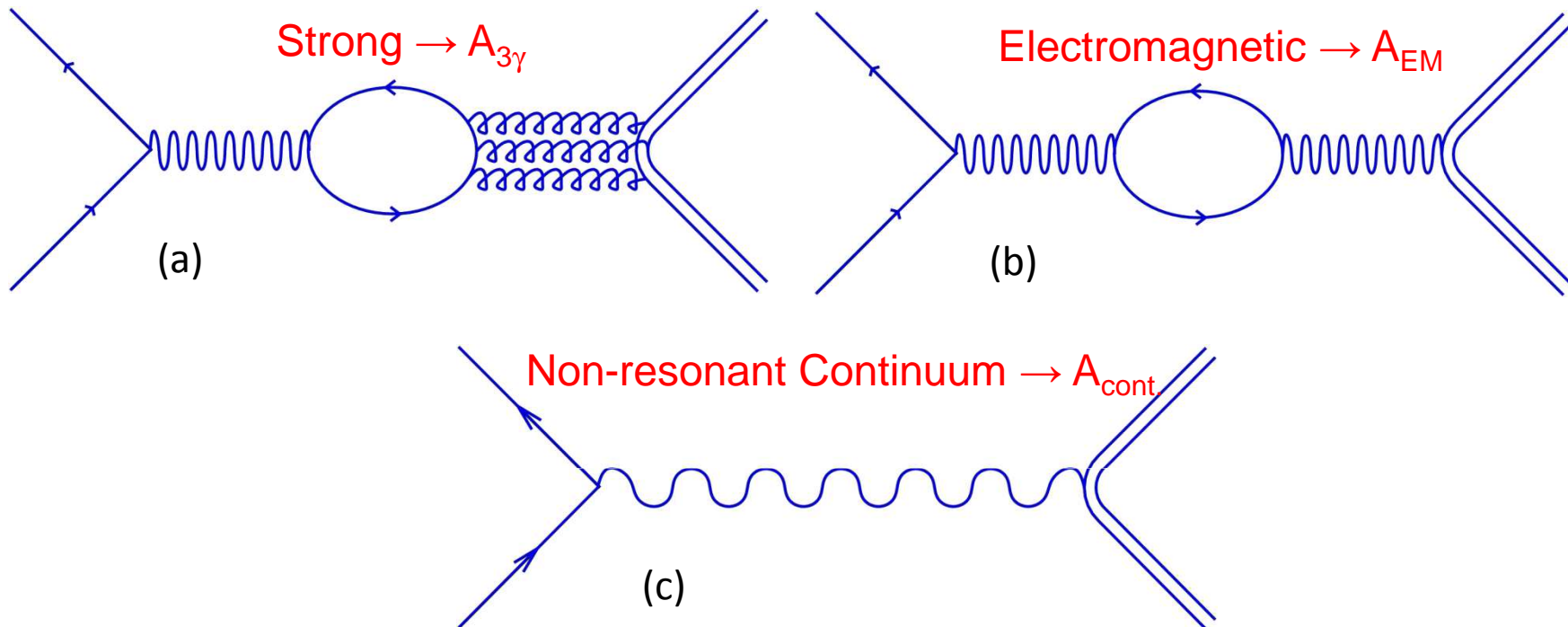
More data expected soon from BESIII  
at neutral and charged baryon threshold  
in particular at the  $\Lambda_c \bar{\Lambda}_c$  threshold

(jump + almost pointlike FF + flat above + anisotropy ?)



## Imaginary Charmonium Decay Widths ?

# Vector Quarkonium Decay Mechanisms



(a)  $e^+e^- \rightarrow J/\psi \rightarrow hadrons$  via strong mechanism (b) via em mechanism

(c) non-resonant  $e^+e^- \rightarrow hadrons$  via a virtual photon.

pQCD regime: all amplitudes real (apart BW resonance behaviour),

while data are as if there is an additional  $i$  in front of the BW

# J/Ψ    Vector + Pseudoscalar SU3 and SU3 Breaking Amplitudes

Use reduced amplitudes  $B=B_0 / P^{*3}$

Process $J/\psi \rightarrow$	Amplitude
$\rho^+\pi^-, \rho^0\pi^0, \rho^-\pi^+$	$g + e$
$K^{*+}K^-, K^{*-}K^+$	$g(1 - s) + e$
$K^{*0}\bar{K}^0, \bar{K}^{*0}K^0$	$g(1 - s) - 2e$
$\omega\eta$	$(g + e)X_\eta + \sqrt{2}rg(\sqrt{2}X_\eta + Y_\eta)$
$\omega\eta'$	$(g + e)X_{\eta'} + \sqrt{2}rg(\sqrt{2}X_{\eta'} + Y_{\eta'})$
$\phi\eta$	$(g(1 - 2s) - 2e)Y_\eta + rg(\sqrt{2}X_\eta + Y_\eta)$
$\phi\eta'$	$(g(1 - 2s) - 2e)Y_{\eta'} + rg(\sqrt{2}X_{\eta'} + Y_{\eta'})$
$\rho^0\eta$	$3eX_\eta$
$\rho^0\eta'$	$3eX_{\eta'}$
$\omega\pi^0$	$3e$
$\phi\pi^0$	$0$

# $J/\Psi$ Vector + Pseudoscalar

Parameter		Fit
SU <sub>3</sub> strong Amplitude	g	$7.22 \pm 0.38$
SU <sub>3</sub> breaking strange	s	$0.18 \pm 0.04$
SU <sub>3</sub> breaking DOZI	r	$-0.04 \pm 0.02$
E.M. Amplitude	e	$0.75 \pm 0.04$
Phase	f	<b><math>81.51 \pm 6.75</math></b>



$J/\psi$

## Vector + Pseudoscalar Decay

Decay	Amplitude	PDGX10 <sup>4</sup>	FitX10 <sup>4</sup>	$\Delta\chi^2$
$\rho^0 \pi^0$	$g e^{i\phi} + e$	$169.0 \pm 15.0$	133.00	1.13
$K^{*+} K^-$	$g (1-s) e^{i\phi} + e$	$51.2 \pm 3.0$	51.5	0.01
$K^{*0} K^0$	$g (1-s) e^{i\phi} - 2e$	$43.9 \pm 3.1$	48.5	0.48
$\omega \eta$	$(g X + d) e^{i\phi} + eX$	$17.4 \pm 2.0$	18.5	0.06
$\phi \eta$	$(g (1-2s)Y + d) e^{i\phi} - 2eY$	$7.5 \pm 0.8$	3.9	4.02
$\rho \eta$	$3eX$	$1.9 \pm 0.2$	2.2	0.30
$\omega \pi$	$3e$	$4.5 \pm 0.5$	4.1	0.11
$\omega \eta'$	$(g X' + d') e^{i\phi} + eX'$	$7.0 \pm 7.0$	11.9	0.10
$\phi \eta'$	$(g (1-2s)Y' + d') e^{i\phi} - 2eY'$	$4.0 \pm 0.7$	6.1	1.87
$\rho \eta$	$3eX'$	$1.1 \pm 0.2$	1.1	0.04

## Pseudoscalar + Pseudoscalar Decay

□ It is possible to avoid  $\pi\pi$  and complications from s quark by means of KK BR's and  $|E^{KK}|$  only

$$\begin{aligned} B^{+-} &= |S|^2 + |E^{+-}|^2 + 2 |S| |E^{+-}| \cos \Phi \\ B^{SL} &= |S|^2 + |E^{SL}|^2 - 2 |S| |E^{SL}| \cos \Phi \end{aligned}$$

$$\square |E^{+-}|^2 = B^{\mu\mu} \sigma(e^+e^- \rightarrow K^+ K^-) / \sigma(e^+e^- \rightarrow \mu^+ \mu^-)$$

$$|E^{SL}|^2 \sim 0, \text{ since}$$

$$\sigma(e^+e^- \rightarrow K_S K_L) \ll \sigma(e^+e^- \rightarrow K^+ K^-)$$

$$\sigma(e^+e^- \rightarrow K_S K_L) \sim 0.6 \text{ pb at } J/\psi$$

$$B^{+-} = (2.37 \pm 0.31) 10^{-4} \quad B^{SL} = (1.66 \pm 0.26) 10^{-4}$$

$$|E^{+-}|^2 = (1.3 \pm 0.6) 10^{-4} \quad \text{from BaBar}$$

$$\Phi = 83.7^\circ \pm 9.0^\circ$$

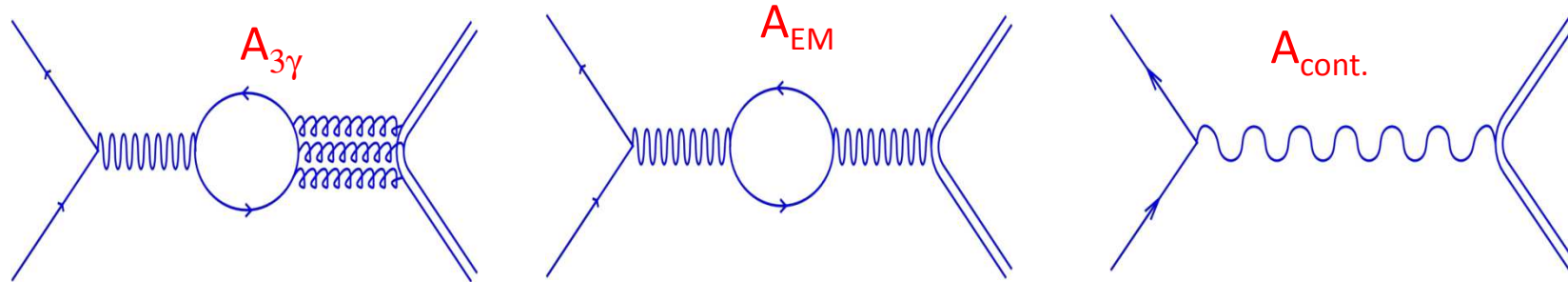
# Experimental evidences for $\Psi(3770)$ imaginary strong decay widths

$\Psi''(3770)$ :

- ❖ non DDbar (small) -> through the interference with continuum
- ❖ For a wide resonance  $\Phi$  from interference at the peak :  
 $\Phi \sim - 2|A_{3g}|/\Gamma_{\text{tot}} \sin \Phi \times \text{continuum}$
- ❖ CLEOc and BESIII:  $\Phi \sim - 90^\circ$ , since continuum sign

decay	continuum	$\Psi''(3770)$	sign	
$\rho \pi$	$13.1 \pm 2.8$	$7.4 \pm 1.3$	-	CLEOc, PRD 73(2006)012002
$\phi \eta$	$2.1 \pm 1.6$	$4.5 \pm 0.7$	+	CLEOc, PRD 73(2006)012002
$\bar{p}p$	$0.74 \pm 0.08$	$0.4 \pm 0.02$	-	BESIII Y.Liang, Nov (2012)

## Model independent (BESIII) from interference in $q^2$ behavior

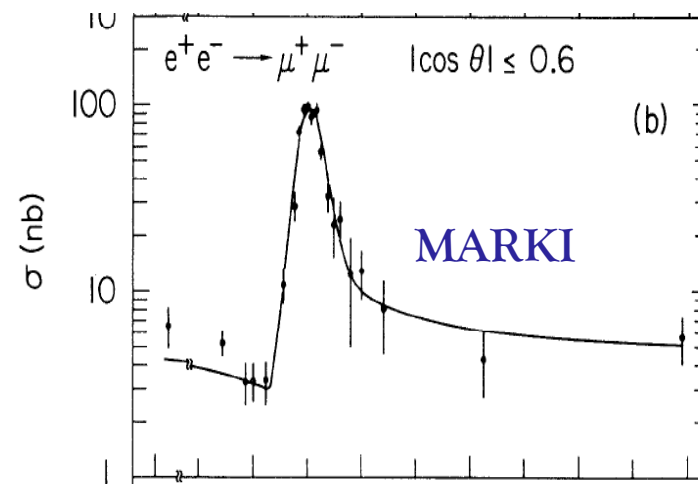


$$s_{\text{born}} = |A_{3\gamma} + A_{EM} + A_{\text{cont}}|^2 = ||A_{3\gamma}|e^{i\varphi} + |A_{EM} + e^{i\varphi'}A_{\text{cont}}||^2$$

Actually  $\Phi_{\text{meas}} = \Phi - d_{\text{cont}}$  and  $|\Phi_{\text{meas}}|$  only is measured,  
since it is difficult to get the sign

In the  $\mu\mu$  case full interference between  $A_{EM}$  and  $A_{\text{cont}}$  observed, as expected, by MARKI(1975), BESII (1995), KDER(2010) (1/2  $\gamma$  propagators :  $\varphi' = 180^\circ$ )

**Results on  $\mu\mu$  and hadronic strong /em decay will come soon, from BESIII**



## Open Issues related to Unitarity

- ❑ No explanation for imaginary strong decay  $J/\Psi$  widths has been put forward until now
- ❑  $J/\Psi$  description as a Breit Wigner might have some difficulties, dealing with imaginary decay widths
- ❑ Optical theorem :  $\text{Im } T_{\text{el}} = W/8\pi \cdot \sigma_{\text{tot}}$  implies  $\text{Im } T_{\text{el}} > 0$
- ❑  $\Gamma(J/\Psi \rightarrow p\bar{p})$  imaginary:  $\text{Im } T_{\text{el}}(p\bar{p} \rightarrow J/\Psi \rightarrow p\bar{p}) < 0$
- ❑  $p\bar{p}$  continuum could restore unitarity, even if unrelated to  $J/\Psi$  (?)
- ❑ Looking for a different  $J/\Psi$  description
- ❑  $\sigma_{\text{el}}(p\bar{p} \rightarrow J/\Psi \rightarrow \text{hadrons})$  : a test of the following model

# A model to explain imaginary widths

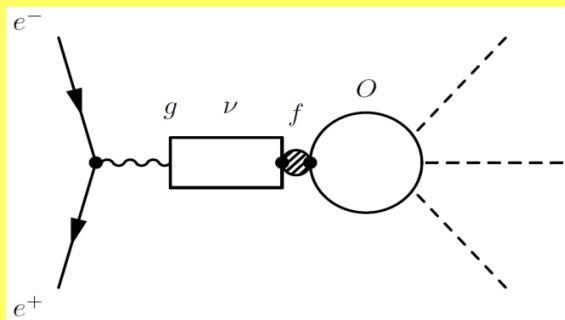


Yochiro Nambu

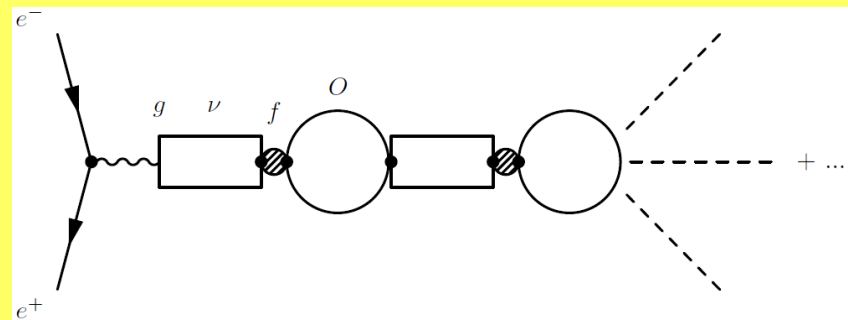
# Quarkonium OZI breaking decay as Freund and Nambu (PRL 34(1975), 1645)

- Quarkonium as a superposition of
  - A narrow  $V$  (coupled to the virtual photon,  
but not directly to hadrons)
  - A wide one (a glueball  $O$ )  
(not coupled to leptons i.e. to a virtual photon,  
but strongly coupled to hadrons)

**$f$  is the coupling between  $\nu$  and  $O$**



+



iterated in  $f$



## Quarkonium OZI breaking decay as Freund and Nambu (PRL 34(1975), 1645)

- Quarkonium as a superposition of V and O:

$$\begin{aligned} A_{\text{strong}} &= G_e V^{-1} f O^{-1} G_f + G_e V^{-1} f O^{-1} f V^{-1} f O^{-1} G_f + \text{iterations} \\ &= G_e V^{-1} f O^{-1} G_f / (1 - V^{-1} O^{-1} f^2) = G_e f G_f / (V O - f^2) \end{aligned}$$

- $$\begin{aligned} A_{\text{em}} &= G_e V^{-1} G_l + G_e V^{-1} f O^{-1} f V^{-1} G_l + \text{iterations} \\ &= G_e O G_f / (V O - f^2) \end{aligned}$$

- An infinity of radial O recurrences (with exceptions?)

- A similar model mainly used to study  $\text{Br}(\psi') / \text{Br}(J/\psi)$  anomalies  
S. J. Brodsky, G. P. Lepage, S. F. Tuan, PRL 59, 621(1987)  
W.S. Hou, C.Y. Ko, NTUTH-97-11, 1997

# Narrow V and wide glueball O superposition

P.J.Franzini, F.J.Gilman, PR D32, 237 (1985)

$$A_{strong} = \frac{\sqrt{\Gamma_{ee}} M_V M_O f \sqrt{\Gamma_O}}{(M_V^2 - W^2 - i M_V \Gamma_V)(M_O^2 - W^2 - i M_O \Gamma_O) - M_V M_O f^2}$$

assuming  $\Gamma_O \gg \Gamma_{J/\psi}$  ,  $f^2 \sim G_O (\Gamma_{J/\psi} - G_V)$

$$A_{strong} \sim \frac{i \sqrt{B_{ee}} M_V f \sqrt{B_h}}{M_{J/\Psi}^2 - W^2 - i M_{J/\Psi} \Gamma_{J/\Psi}} \quad A_{em} = \frac{\sqrt{B_{ee}} M_V \Gamma_{J/\Psi} \sqrt{B_{em}}}{M_{J/\Psi}^2 - W^2 - i M_{J/\Psi} \Gamma_{J/\Psi}}$$

## ■ The additional 90° phase is naturally achieved

- J/ψ shape reproduced if:  $|f| \sim 0.012 \text{ GeV}$  ,  $M_O \sim M_{J/\psi}$  ,  $\Gamma_O \sim 0.5 \text{ GeV}$
- different only far from the J/ψ ( no contradiction with BES, PR 54(1996)1221)
- $\psi''(3770)$  decay phases agree with Nambu suggestion.
- $\psi'$  unclear;  $\psi' \rightarrow J/\psi \pi\pi$  (?)

# SND $\Phi \rightarrow \pi^+ \pi^- \pi^0$

SND measured  $\Phi \rightarrow \pi^+ \pi^- \pi^0$ .

$\phi$  interferes with  $\omega$  and  $\omega'$  tails:

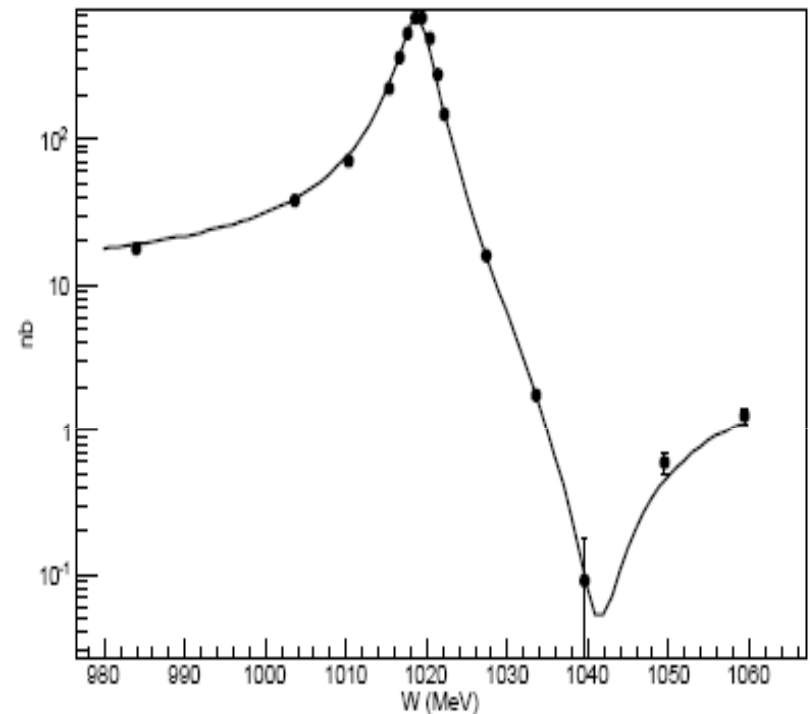
$\phi \sim 180^\circ$  (interference dip after the  $\Phi$ )

Fit SND  $\Phi$  and continuum data with

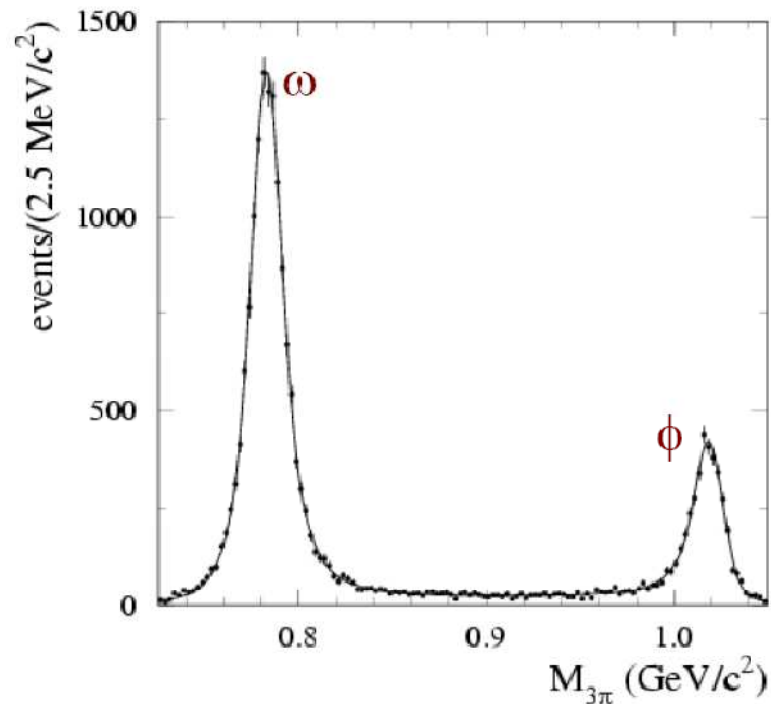
$f = -0.016 \text{ GeV}$  (like  $J/\Psi$  !)

$M_O = 1.34 \text{ GeV}$  (far from the  $\Phi$ )

$\Gamma_O \sim 0.5 \text{ GeV}$



SND data on  $3\pi$  and  
present model prediction



### Masses and widths

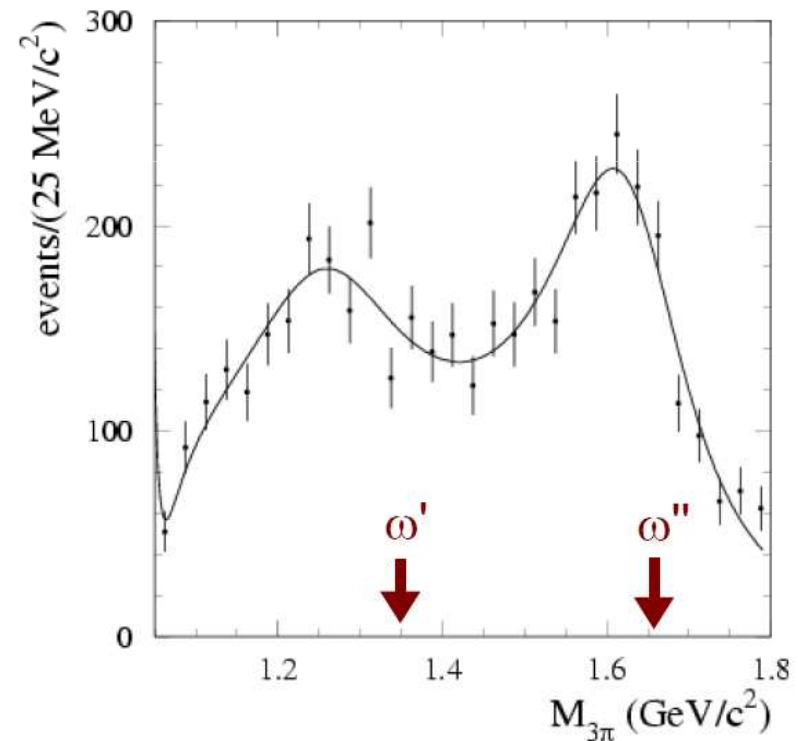
$$M_{\omega'} = (1350 \pm 20 \pm 20) \text{ MeV}/c^2$$

$$\Gamma_{\omega'} = (450 \pm 70 \pm 70) \text{ MeV}/c^2$$

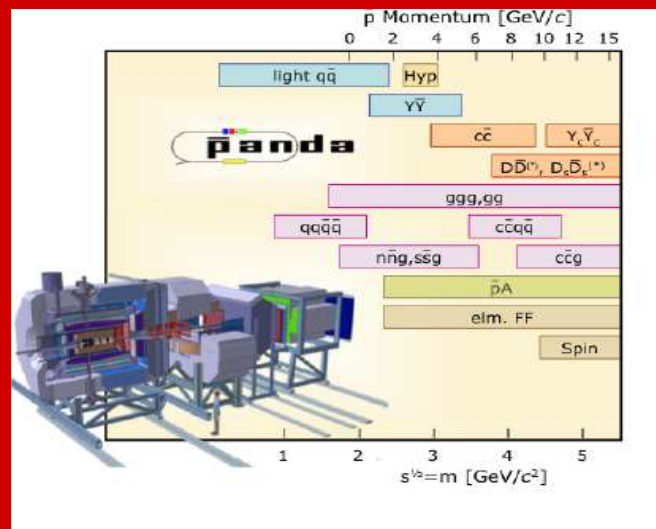
$$M_{\omega''} = (1660 \pm 10 \pm 2) \text{ MeV}/c^2$$

$$\Gamma_{\omega''} = (230 \pm 30 \pm 20) \text{ MeV}/c^2$$

BaBar found indeed an unexpected resonance ( $\omega'$ ,  $O$  ?) at 1.35 GeV, wide 0.45 GeV



# A proposal for PANDA: a $J/\Psi$ scan

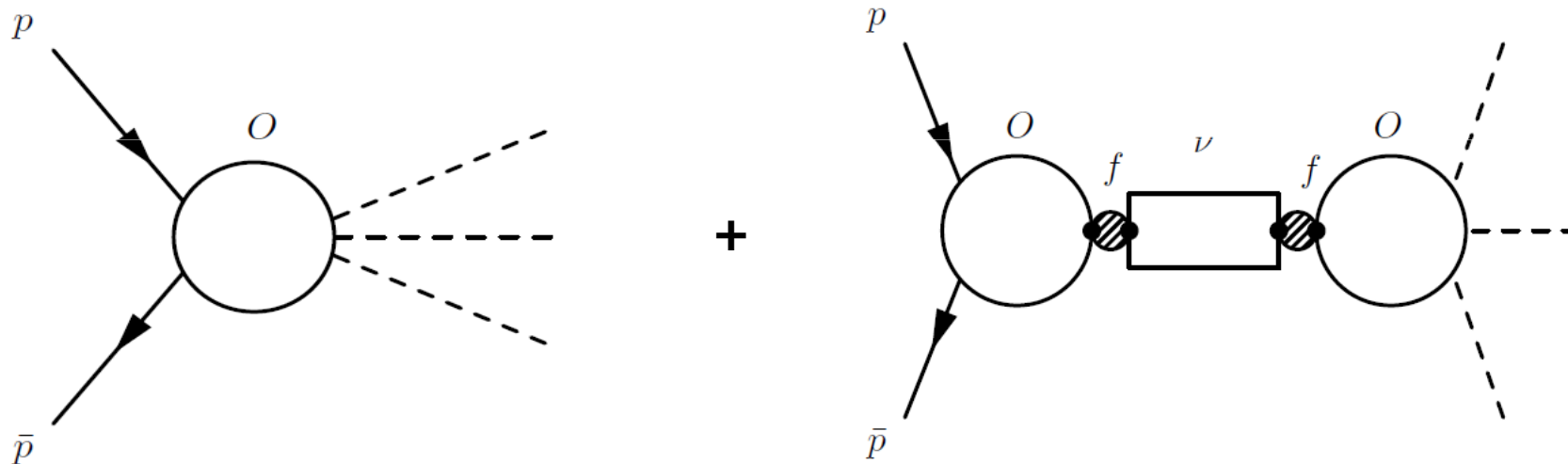


# A Proposal for PANDA

- ❑ Expected  $\sigma (p\bar{p} \rightarrow J/\Psi \rightarrow \text{hadrons}) \sim 1 \mu\text{b}$   
while  $\sigma (p\bar{p} \rightarrow \text{hadrons}) \sim 70 \text{ mb}$
- ❑ No  $J/\Psi$  exclusive production evidence in present data  
(too small cross section +  $p\bar{p}$  c.m. energy spread)
- ❑ Different mechanism in inclusive or exclusive production:
  - Inclusive production: direct coupling to gluons or virtual photon
  - Exclusive production: hadronic  $\rightarrow$  apply FN model

# A Proposal for PANDA

Contributions to  $p\bar{p} \rightarrow J/\Psi \rightarrow \text{hadrons}$ , according to the FN model





# A Proposal for PANDA

□ According to the FN approach

$$\sigma_{FN} = \frac{B_p [(M_{J/\Psi}^2 - W^2)^2 + (M_{J/\Psi} \Gamma_V)^2] B_h}{(M_{J/\Psi}^2 - W^2)^2 + (M_{J/\Psi} \Gamma_{J/\Psi})^2}$$

Taking into account that  $\Gamma_V \ll \Gamma_{J/\Psi}$

$$\sigma_{FN} = \frac{B_p (M_{J/\Psi}^2 - W^2)^2 B_h}{(M_{J/\Psi}^2 - W^2)^2 + (M_{J/\Psi} \Gamma_{J/\Psi})^2}$$

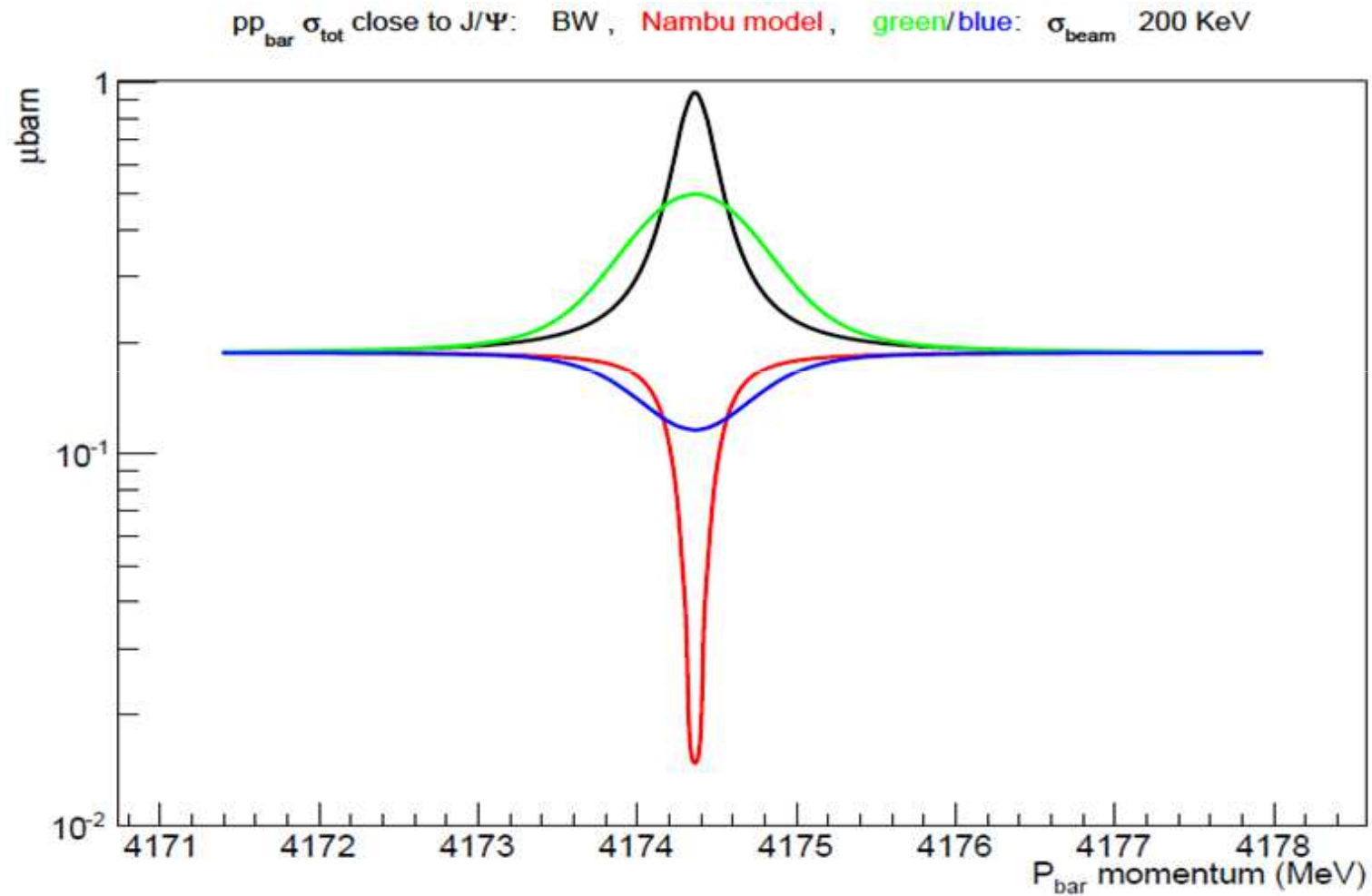
Almost a zero -> a dip in  $\sigma_h$

□ To be compared to a Breit Wigner

$$\sigma_{BW} = \frac{B_p \Gamma_{J/\Psi}^2 M_{J/\Psi}^2 B_h}{(M_{J/\Psi}^2 - W^2)^2 + (M_{J/\Psi} \Gamma_{J/\Psi})^2}$$

# A Proposal for PANDA

PANDA good inv. mass resolution: small beam energy spread and no ISR



# Conclusions

## ❑ $B\bar{B}$ cross section unexpected features at threshold:

- Jump followed by a flat behavior
- Pointlike cross section
- Analyticity violation (?)
- Jump also in the case of neutral baryon (?)

## ❑ Charmonium lineshape:

- Imaginary Decay Widths
- A model for a non Breit-Wigner charmonium lineshape
- A proposal for PANDA

Hope you acknowledge that  
Even if it might be not true  
it is well conceived  
(italian common saying)