

Electroweak Symmetry Breaking

New Physics after the Higgs Discovery?

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LAPP, Annecy-le-Vieux, 12-13 November 2014

Outline

1) Higgs Physics

- Hierarchy problem. Naturalness
- Top: the heaviest mass scale
- Standard Model Higgs
- Supersymmetry
- Flavour Dynamics

2) Two Higgs Doublets

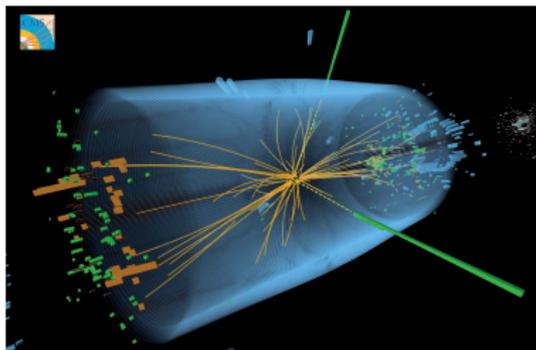
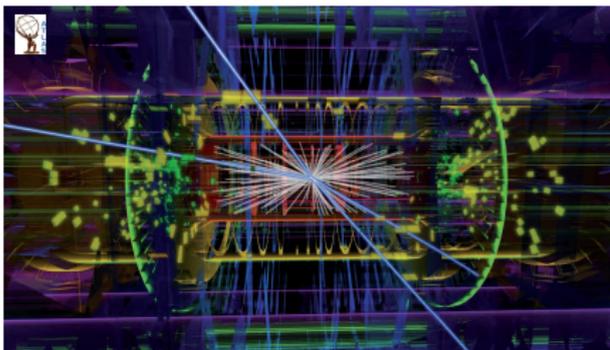
- FCNCs
- Flavour alignment
- Flavour & LHC constraints
- EDMs

3) Electroweak Effective Theory

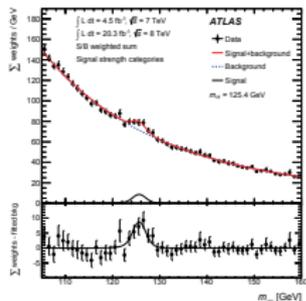
- Custodial Symmetry
- Equivalence Theorem
- Strong Coupling



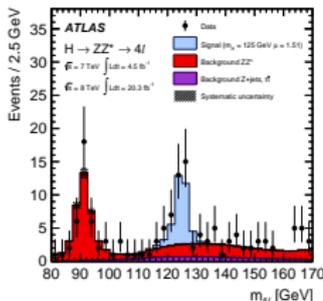
A New Higgs-Like Boson



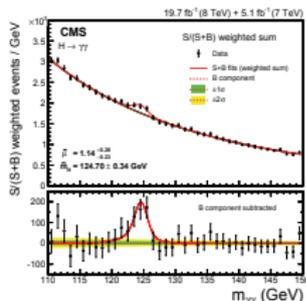
$H \rightarrow \gamma\gamma$



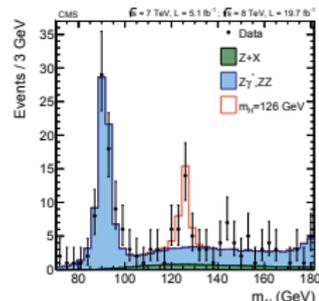
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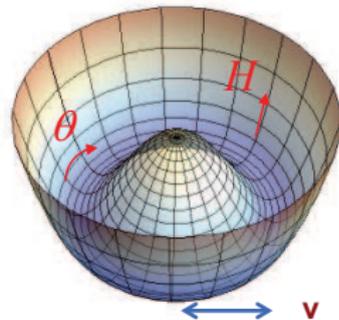


$$M_H^{\text{ATLAS}} = (125.36 \pm 0.37 \pm 0.18) \text{ GeV}$$

$$M_H^{\text{CMS}} = (125.03 \pm 0.26 \pm 0.13 \pm 0.27 \pm 0.15) \text{ GeV}$$

Great success of the Standard Model

BEGHHK (\equiv Higgs) Mechanism



$$SU(2)_L \otimes U(1)_Y \quad v = 246 \text{ GeV}$$

$$M_Z \cos \theta_W = M_W = \frac{1}{2} v g$$





Beautiful Discovery

Boson, $J \neq 1$

Fermions = Matter ; Bosons = Forces

- **Fundamental Boson:** New interaction which is not gauge
- **Composite Boson:** New underlying dynamics



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If New Physics exists at Λ_{NP}

$$\delta M_H^2 \sim \frac{g^2}{(4\pi)^2} \Lambda_{\text{NP}}^2 \log \left(\frac{\Lambda_{\text{NP}}^2}{M_H^2} \right)$$

Which symmetry keeps M_H away from Λ_{NP} ?

- **Fermions:** Chiral Symmetry
- **Gauge Bosons:** Gauge Symmetry
- **Scalar Bosons:** Supersymmetry, Scale/Conformal Symmetry ... ?

Symmetries & Mass Scales

Fermions: $\psi_{L,R} \longrightarrow e^{i\alpha_{L,R}} \psi_{L,R}$ **Chiral symmetry**

$$\mathcal{L}_\psi = \bar{\psi} (i\partial - m_\psi) \psi = \bar{\psi}_L i\partial \psi_L + \bar{\psi}_R i\partial \psi_R - m_\psi (\bar{\psi}_L \psi_R + \bar{\psi}_R \psi_L)$$

Symmetry recovered at $m_\psi = 0$  $\delta m_\psi \propto m_\psi$

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Vectors: $A_\mu \longrightarrow A_\mu + \partial_\mu \theta$ **Gauge symmetry**

$$\mathcal{L}_A = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2} m_A^2 A_\mu A^\mu$$

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Scalars: $\mathcal{L}_\phi = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{2} m_\phi^2 \phi^2$ **Any symmetry?**

No additional symmetry at $m_\phi = 0$ \longrightarrow $\delta m_\phi^2 \propto M^2$ ($M = \text{any scale}$)

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Pseudo-Goldstone Boson

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Pseudo-Goldstone Boson

- **Scale symmetry:** $x \rightarrow x/\lambda$, $\phi(x) \rightarrow \lambda \phi(x/\lambda)$

$$\mathbf{M} = \mathbf{0} \quad , \quad \forall \mathbf{M}$$

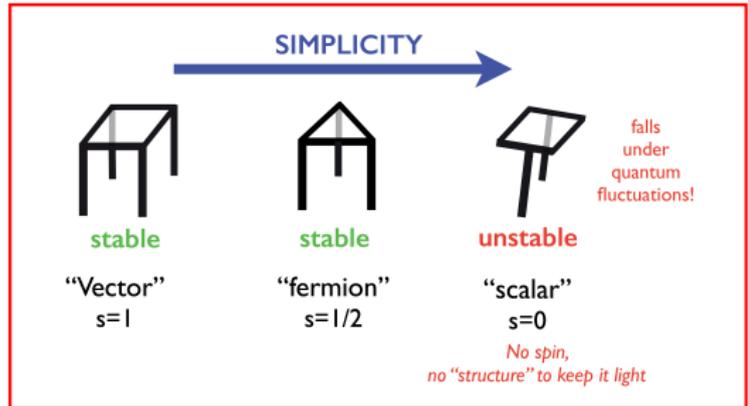
Conformal Invariance. Dilaton

Which symmetry keeps M_H away from Λ_{NP} ?

A. Pomarol

Quantum Stability

Vectors/fermions
protected by
gauge/chiral symmetries



Proposed Solutions

SUSY / Composite Higgs

1) Keep the Higgs elementary, but protect it by symmetries: **Supersymmetry**

Higgs (boson) \longleftrightarrow Higgsino (fermion)



2) The Higgs is not elementary: **Composite Higgs**

Higgs made of fermions
(as a pion in strong interactions)



Spin, Mass & Degrees of Freedom

	$J = 1$	$J = \frac{1}{2}$	$J = 0$
$M = 0$	2 d.o.f. Trans. Pol.	2 d.o.f. ψ_L	1 d.o.f.
$M \neq 0$	3 d.o.f. Trans & Long.	4 d.o.f. ψ_L, ψ_R	1 d.o.f.

Vector ($2 \neq 3$) and fermion masses are safe ($2 \neq 4$)

Scalar masses not protected (continuous $m \rightarrow 0$ limit)

Higgs Self-energy: The top is the largest SM contribution

δM_H^2 stabilized through new-physics contributions

- **SUSY:** stop loops



- **Composite Higgs:** fermionic top partners





Possible Scenarios of EWSB

① **SM Higgs:** Favoured by EW precision tests

② **Alternative perturbative EWSB:**

Scalar Doublets and singlets

$$\rho_{\text{tree}} = \frac{M_W^2}{M_Z^2 c_W^2} = \frac{\sum_i v_i^2 [T_i(T_i + 1) - Y_i^2]}{2 \sum_i v_i^2 Y_i^2}$$

③ **Dynamical (non-perturbative) EWSB:**

Pseudo-Goldstone Higgs

Scalar Resonance



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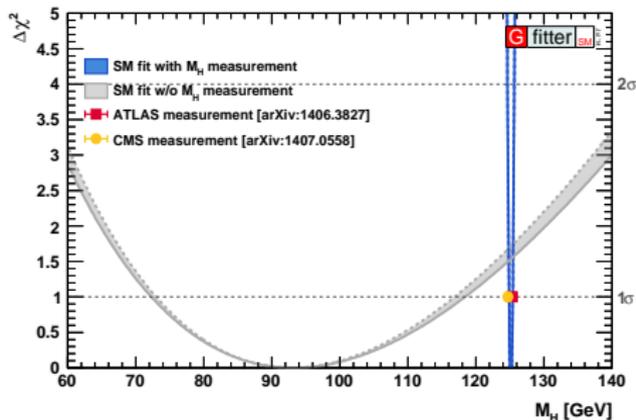
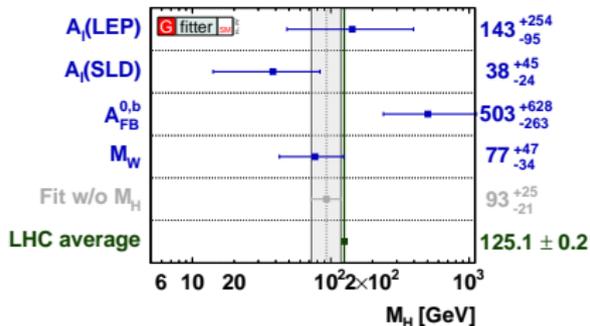
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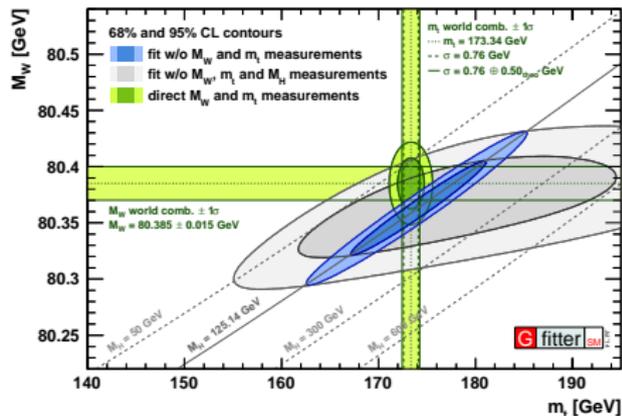
Scalar Resonance



SM Higgs



Favoured by
EW precision tests



TOP MASS

Uwer

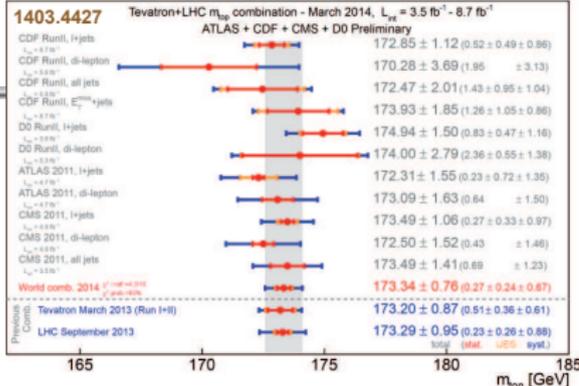
Monte Carlo mass:

$$M_t^{MC} = (173.34 \pm 0.76) \text{ GeV}$$

Lacks a proper QCD definition: $M_t^{\text{pole}} = M_t^{MC} + \Delta M_t^{\text{th}}$

$$|\Delta M_t^{\text{th}}| \approx O(1 \text{ GeV})$$

Hoang-Stewart, 0808.0222



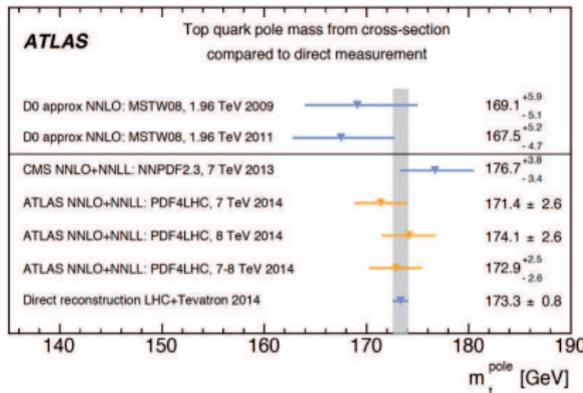
Aldaya, Bernardi, Carli

Cross section:

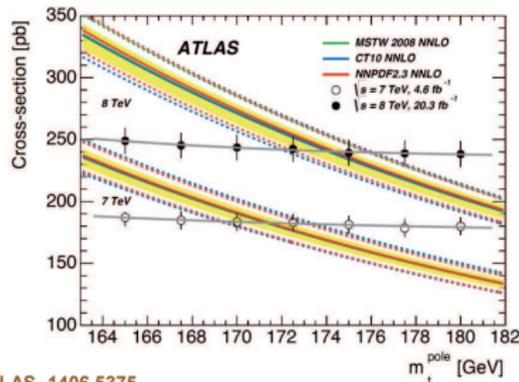
$\sigma_{t\bar{t}}$ NNLO+NNLL

Well-defined mass

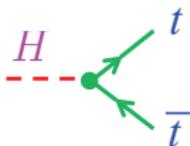
Czakon et al, Bärnreuther et al, Cacciari et al



ATLAS, 1406.5375



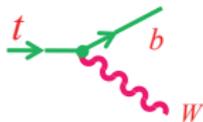
The Heaviest Mass Scale



$$y_t = \frac{\sqrt{2}}{v} m_t = 2^{3/4} G_F^{1/2} m_t = 1 \quad (0.995)$$

The top quark:

- Sensitive probe of Electroweak Symmetry Breaking
- Non-perturbative (strong) dynamics ?
- Very different from other quarks: $y_b = 0.025$, $y_c = 0.007 \dots$
- Is it really a SM quark?



So far, we only know
the decay $t \rightarrow b W^+$

Single-top	$ V_{tb} $
ATLAS '14	> 0.88 (95% CL)
CMS '14	> 0.92 (95% CL)
CDF '14	> 0.84 (95% CL)
D0 '13	> 0.92 (95% CL)

Higgs Mechanism:

Gauge invariance

Massless W^\pm, Z (spin 1)

$$3 \times 2 \text{ polarizations} = 6$$

Higgs Mechanism: 3 additional degrees of freedom $\theta_i(x)$

Gauge invariance

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3×2 polarizations = 6

+

3 Goldstones $\theta_i(x)$

SSB 

Massive W^\pm, Z

3×3 polarizations = 9

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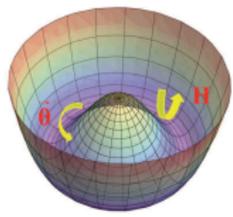
SSB ↓

Massive W^\pm, Z

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Spontaneous Symmetry Breaking

$$\mathcal{L}_\Phi = (D_\mu \Phi)^\dagger D^\mu \Phi - \mu^2 \Phi^\dagger \Phi - \lambda (\Phi^\dagger \Phi)^2$$



$$\mu^2 < 0$$

$$\Phi(x) = \exp \left\{ i \frac{\vec{\sigma}}{2} \cdot \vec{\theta}(x) \right\} \frac{1}{\sqrt{2}} \begin{bmatrix} 0 \\ v + H(x) \end{bmatrix}$$

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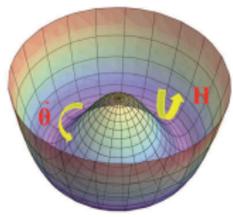
SSB \downarrow

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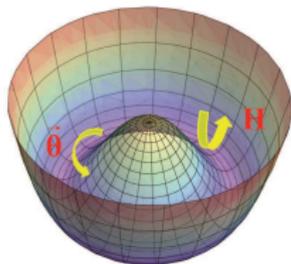
$$\Phi(x) = \exp \left\{ i \frac{\vec{\sigma} \cdot \vec{\theta}(x)}{2} \right\} \frac{1}{\sqrt{2}} \begin{bmatrix} 0 \\ v + H(x) \end{bmatrix}$$

$$D_\mu \Phi = (\partial_\mu + \frac{i}{2} g \vec{\sigma} \cdot \vec{W}_\mu + \frac{i}{2} g' B_\mu) \Phi \quad ; \quad v^2 = -\mu^2 / \lambda$$

$$(D_\mu \Phi)^\dagger D^\mu \Phi \rightarrow M_W^2 W_\mu^\dagger W^\mu + \frac{M_Z^2}{2} Z_\mu Z^\mu$$

$$M_W = M_Z \cos \theta_W = \frac{1}{2} g v$$

SM Higgs Potential



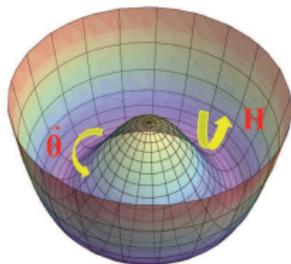
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$$V(\Phi) + \frac{\lambda}{4} v^4 = \lambda \left(|\Phi|^2 - \frac{v^2}{2} \right)^2 = \frac{1}{2} M_H^2 H^2 + \frac{M_H^2}{2v} H^3 + \frac{M_H^2}{8v^2} H^4$$

$$v = (\sqrt{2} G_F)^{-1/2} = 246 \text{ GeV}$$

$$M_H = (125.14 \pm 0.24) \text{ GeV} \quad \rightarrow \quad \lambda = \frac{M_H^2}{2v^2} = 0.13$$

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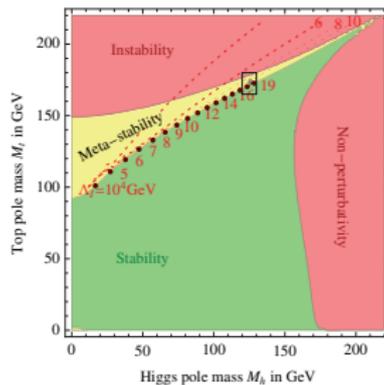
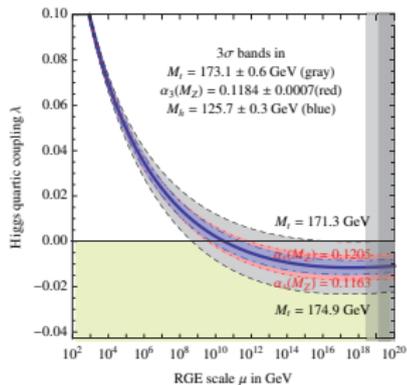
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$$M_H^2 = 2v^2 \lambda(\mu) + \frac{2v^2 y_t^2}{(4\pi)^2} [2\lambda + 3(\lambda - y_t^2) \log(m_t^2/\mu^2)] + \dots$$

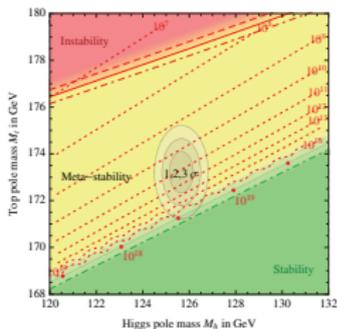
$$y_t = \sqrt{2} m_t/v \approx 1$$

Vacuum Stability: $\lambda(\Lambda) \geq 0$

Degrassi et al, 1205.6497, 1307.3536



$$\Lambda = M_{\text{Planck}} \quad \rightarrow \quad M_H > (129.1 \pm 1.5) \text{ GeV}$$

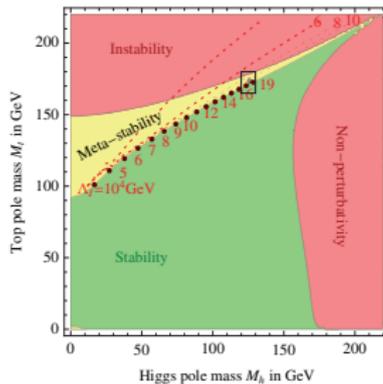
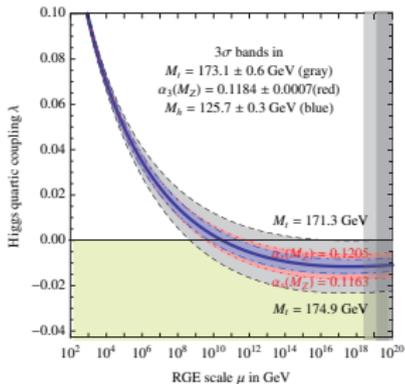


Assumes SM valid all the way up to $\Lambda \leq M_{\text{Planck}}$

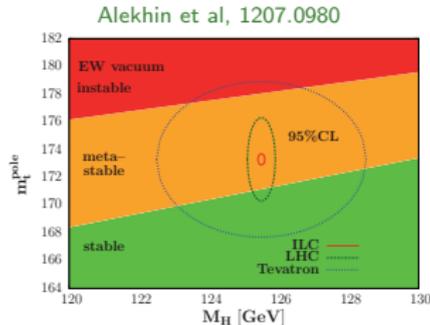
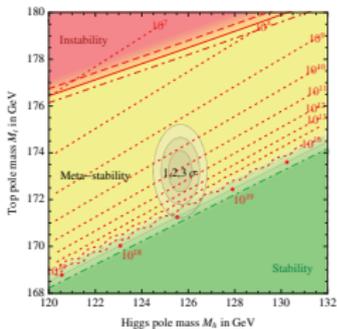
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Degrassi et al, 1205.6497, 1307.3536



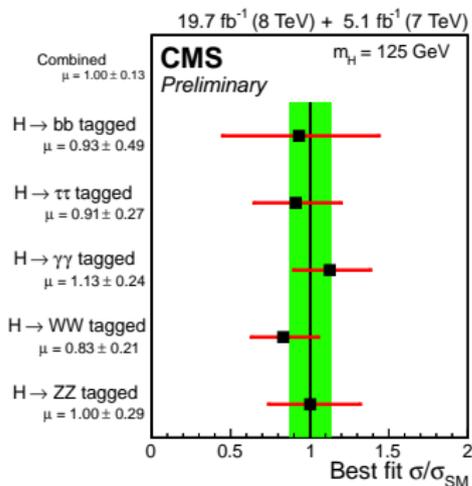
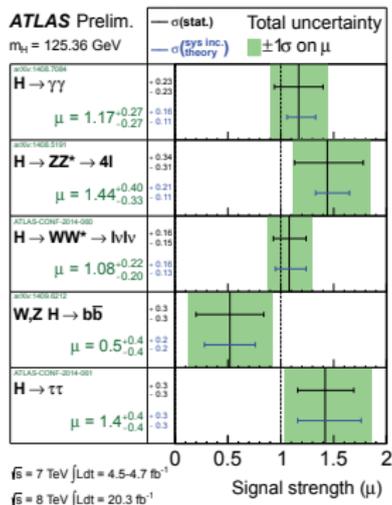
$$\Lambda = M_{\text{Planck}} \quad \rightarrow \quad M_H > (129.1 \pm 1.5) \text{ GeV} \quad [129.8 \pm 5.6]$$



Assumes SM valid all the way up to $\Lambda \leq M_{\text{Planck}}$

Signal Strengths

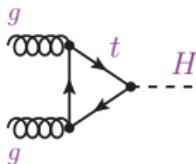
$$\mu \equiv \sigma \cdot \text{Br} / (\sigma \cdot \text{Br})_{\text{SM}}$$



Decay Mode	ATLAS ($M_H = 125.4$ GeV)	CMS ($M_H = 125.0$ GeV)
$V H \rightarrow b\bar{b}$	0.5 ± 0.4	0.93 ± 0.49
$H \rightarrow \tau\tau$	1.4 ± 0.4	0.91 ± 0.27
$H \rightarrow \gamma\gamma$	1.17 ± 0.27	1.13 ± 0.24
$H \rightarrow WW^*$	$1.08^{+0.22}_{-0.20}$	0.83 ± 0.21
$H \rightarrow ZZ^*$	$1.44^{+0.40}_{-0.33}$	1.00 ± 0.29
Combined		1.00 ± 0.13

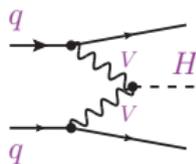
Production Channels

Gluon Fusion

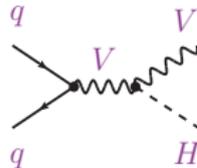


Vector Boson Fusion

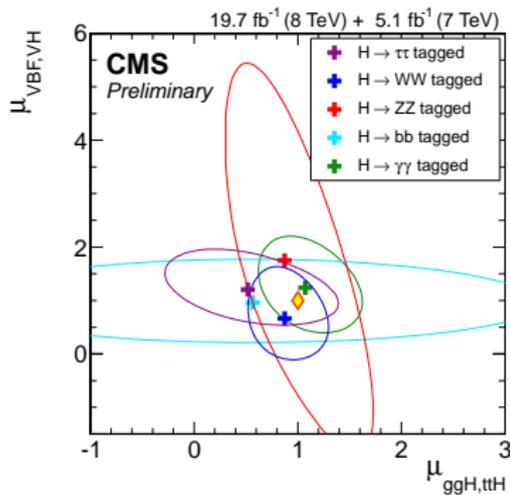
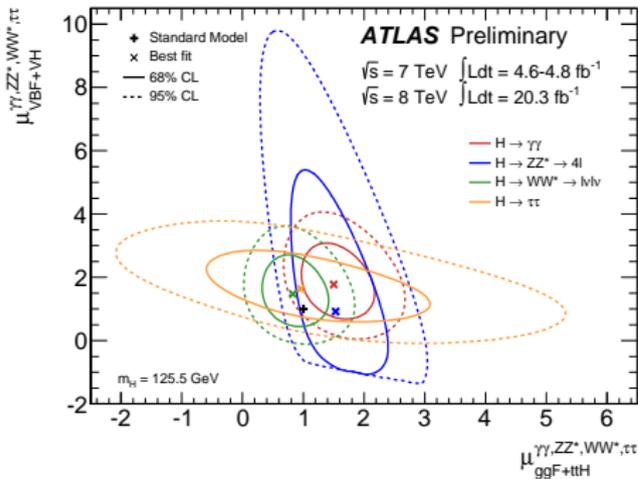
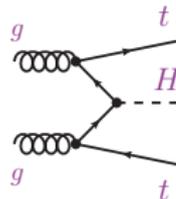
($V = W^\pm, Z$)



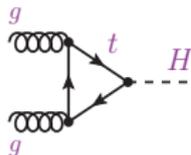
Ass. VH Production



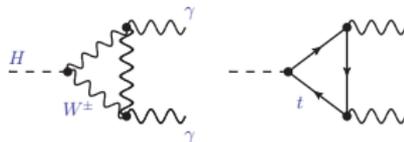
Ass. $t\bar{t}H$ Production



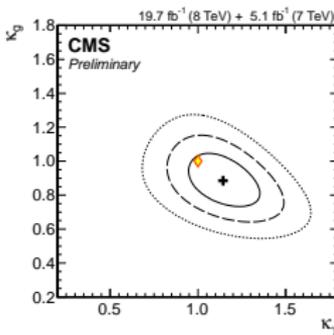
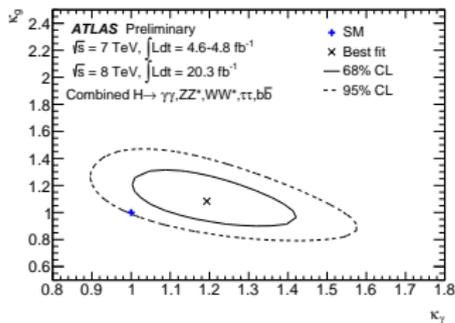
Strong (indirect) evidence for Higgs coupling to t



Dominant
Production Mechanism



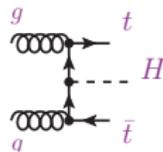
$$\Gamma \sim |1 - 0.21|^2$$



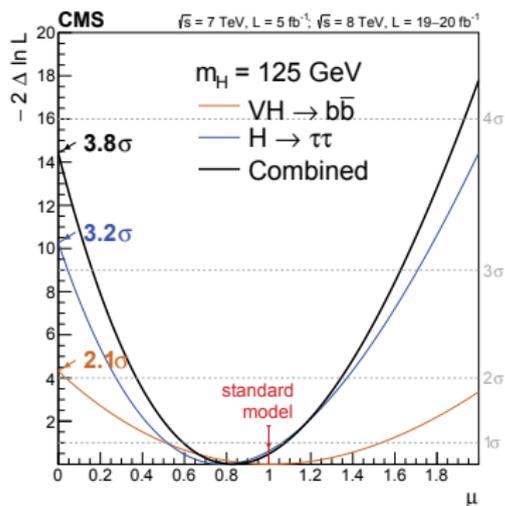
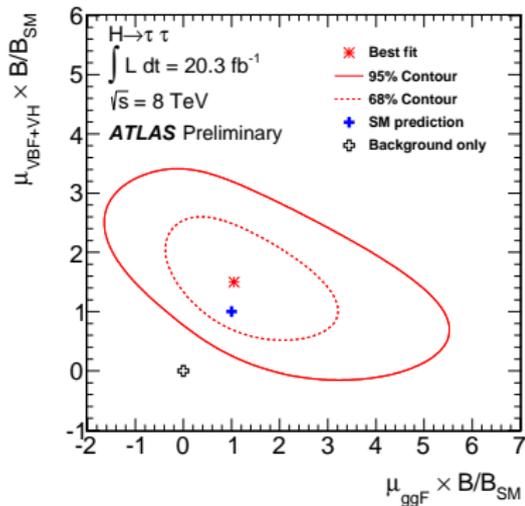
$$\kappa_i \equiv g_i / g_i^{\text{SM}}$$

$H \rightarrow \gamma\gamma$	Signal Strength
ATLAS	1.17 ± 0.27
CMS	1.13 ± 0.24

Direct (tree-level) sensitivity through $t\bar{t}H$



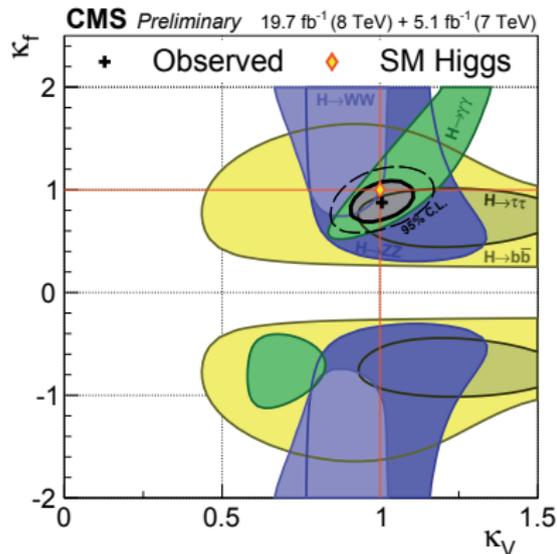
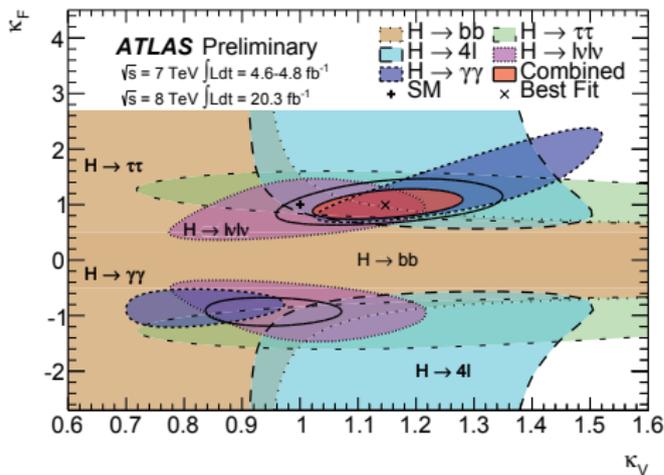
Strong evidence for Higgs coupling to τ and b



Signal Strength	ATLAS ($M_H = 125.4 \text{ GeV}$)	CMS ($M_H = 125.0 \text{ GeV}$)
$VH \rightarrow b\bar{b}$	0.5 ± 0.4	0.93 ± 0.49
$H \rightarrow \tau\tau$	1.4 ± 0.4	0.91 ± 0.27

Effective Couplings

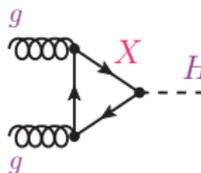
$$\kappa_i \equiv g_i/g_i^{\text{SM}}$$



$$\sigma(i \rightarrow H) \cdot \text{Br}(H \rightarrow f) = \sigma(i \rightarrow H) \cdot \Gamma(H \rightarrow f) / \Gamma_H \sim (\kappa_i \kappa_f / \kappa_H)^2$$

QCD Exotics

$X \in SU(3)_C$ representation \underline{R}



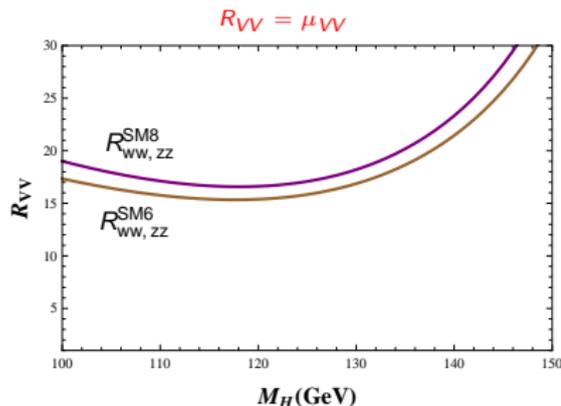
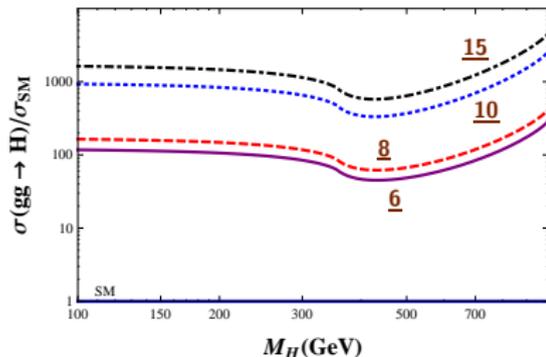
$$\sim \sum_{a=1}^{d_A} \text{Tr} [t_R^a t_R^a] = C_R d_R$$

Non decoupling: $\mathcal{L} = -\frac{M_X}{v} (\bar{X}X) H$

Exotic fermions in higher-colour representations could only exist provided their masses are not generated by the SM Higgs

(or fine-tuned cancelations with scalar loops)

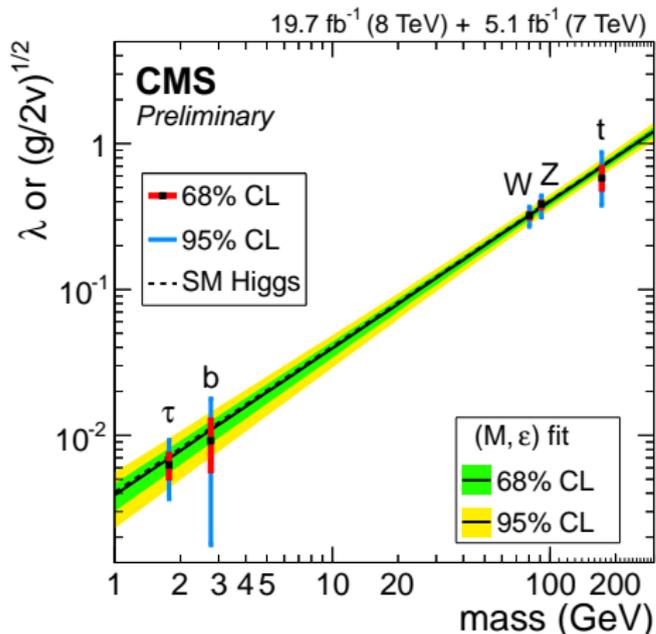
V. Ilisie - AP, 1202.3430



It is a Higgs Boson

$$\lambda_f = (m_f/M)^{1+\epsilon} \quad , \quad g_V = 2 m_V^{2(1+\epsilon)}/M^{1+2\epsilon}$$

Ellis-You, 1303.3879



SM: $\epsilon = 0$, $M = v = 246$ GeV

$$\epsilon = 0.00 \pm 0.03$$

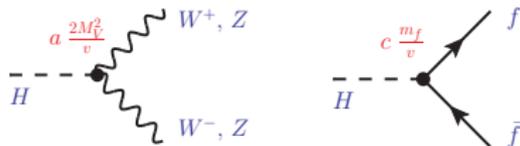
$$M/v = 0.97 \pm 0.06$$

SM rates:

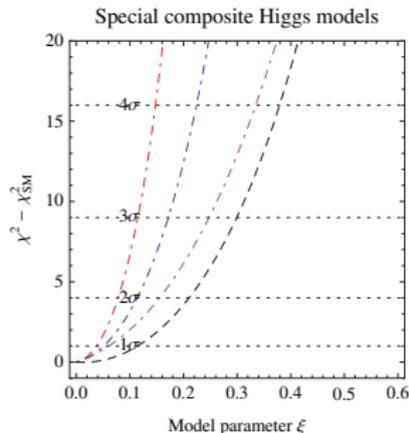
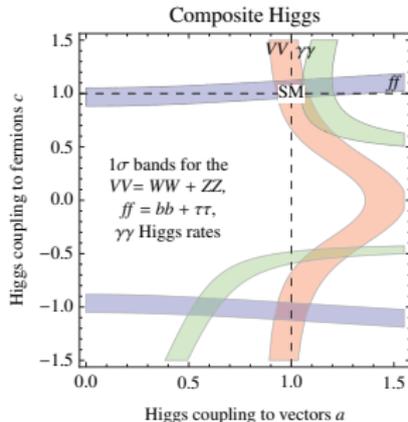
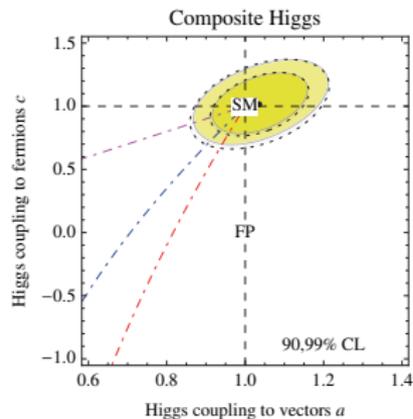
$$M_H = (125.0 \pm 1.8) \text{ GeV}$$

Giardino et al, 1303.3570v4

2-Parameter Fit



Giardino et al, 1303.3570v4

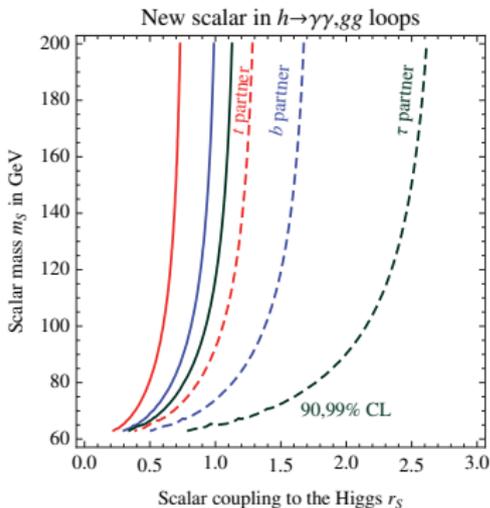
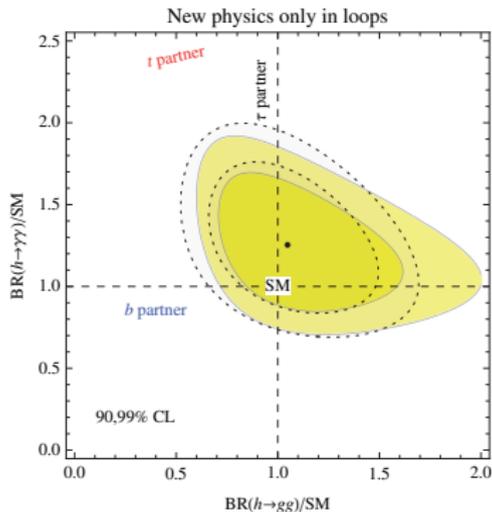


	a	c	
---	$\sqrt{1-\xi}$	a	MCHM4 (Agashe et al)
---	$\sqrt{1-\xi}$	$(1-2\xi)/a$	MCHM5 (Contino et al)
---	$\sqrt{1-\xi}$	$(1-3\xi)/a$	
---	1	$1-\xi$	

$$\xi = v^2/F^2$$

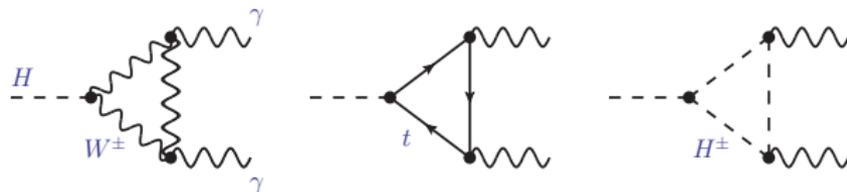
New Physics Only in Loops

Giardino et al, 1303.3570v4



$$\mathcal{L}_S^{\text{NP}} = r_S \frac{2m_S^2}{v} H S^2$$

$H \rightarrow \gamma\gamma$



$$\Gamma \propto |-8.4 \kappa_W + 1.8 \kappa_t + \mathcal{C}_{\text{NP}}|^2$$

SM: $\kappa_W = \kappa_t = 1$, $\mathcal{C}_{\text{NP}} = 0$

Destructive interference

Enhanced rate if:

- $\kappa_W \kappa_t < 0$
- $\mathcal{C}_{\text{NP}} < 0$

Desperately Seeking SUSY (Dulcinea)



In all the world there is no maiden fairer than the Empress of La Mancha, the peerless SUSY del Toboso

Your worship should bear in mind that SUSY is badly broken; got heavy through anomaly mediation



	Model	e, μ, τ, γ	Jets	E_{T}^{miss}	$\int \mathcal{L} d\mathcal{L} [\text{fb}^{-1}]$	Mass limit	Reference
Inclusive Searches	MSUGRA/CMSSM	0	2-6 jets	Yes	20.3	$\tilde{\chi}_1^0, \tilde{\chi}_2^0$ 1.7 TeV	$m(\tilde{g})=m(\tilde{g})$ 1405.7875
	MSUGRA/CMSSM	$1 e, \mu$	3-6 jets	Yes	20.3	$\tilde{\chi}_1^0, \tilde{\chi}_2^0$ 1.2 TeV	any $m(\tilde{g})$ ATLAS-CONF-2013-062
	MSUGRA/CMSSM	0	7-10 jets	Yes	20.3	$\tilde{\chi}_1^0, \tilde{\chi}_2^0$ 1.1 TeV	any $m(\tilde{g})$ 1308.1841
	$\tilde{q}\tilde{q}, \tilde{g} \rightarrow \tilde{q}\tilde{q}^*$	0	2-6 jets	Yes	20.3	$\tilde{\chi}_1^0, \tilde{\chi}_2^0$ 850 GeV	$m(\tilde{t}_1)=0 \text{ GeV}, m(1^{\text{st}} \text{ gen. } \chi) = m(2^{\text{nd}} \text{ gen. } \chi)$ 1405.7875
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow \tilde{q}\tilde{q}^*$	0	2-6 jets	Yes	20.3	$\tilde{\chi}_1^0, \tilde{\chi}_2^0$ 1.33 TeV	$m(\tilde{t}_1)=0 \text{ GeV}$ 1405.7875
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow \tilde{q}\tilde{q}^* \tilde{\chi}_1^0$	$1 e, \mu$	3-6 jets	Yes	20.3	$\tilde{\chi}_1^0, \tilde{\chi}_2^0$ 1.18 TeV	$m(\tilde{t}_1) < 200 \text{ GeV}, m(\tilde{\tau}^*) = 0.5(m(\tilde{t}_1) + m(\tilde{g}))$ ATLAS-CONF-2013-062
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow \tilde{q}\tilde{q}^* \ell(\ell' \nu) \nu \nu^*$	$2 e, \mu$	0-3 jets	-	40.7	$\tilde{\chi}_1^0, \tilde{\chi}_2^0$ 1.12 TeV	$m(\tilde{t}_1) = 0 \text{ GeV}$ ATLAS-CONF-2013-089
	GMSB ($\tilde{\ell}$ NLSP)	$2 e, \mu$	2-4 jets	Yes	20.3	$\tilde{\chi}_1^0, \tilde{\chi}_2^0$ 1.24 TeV	$\tan\beta < 15$ 1208.4688
	GMSB ($\tilde{\ell}$ NLSP)	$1-2 \tau + 0-1 \ell$	0-2 jets	Yes	20.3	$\tilde{\chi}_1^0, \tilde{\chi}_2^0$ 1.6 TeV	$\tan\beta > 20$ 1407.0603
	GGM (bino NLSP)	2 τ	-	Yes	20.3	$\tilde{\chi}_1^0, \tilde{\chi}_2^0$ 1.28 TeV	$m(\tilde{t}_1) = 50 \text{ GeV}$ ATLAS-CONF-2014-001
GGM (wino NLSP)	$1 e, \mu + \gamma$	-	Yes	4.8	$\tilde{\chi}_1^0, \tilde{\chi}_2^0$ 619 GeV 900 GeV	$m(\tilde{t}_1) = 220 \text{ GeV}$ ATLAS-CONF-2012-144	
GGM (higgsino-bino NLSP)	0	1-3 jets	Yes	4.8	$\tilde{\chi}_1^0, \tilde{\chi}_2^0$ 690 GeV	$m(\text{NLSP}) = 200 \text{ GeV}$ 1211.1161	
GGM (higgsino NLSP)	$2 e, \mu$ (Z)	0-3 jets	Yes	5.8	$\tilde{\chi}_1^0, \tilde{\chi}_2^0$ 690 GeV	$m(\text{NLSP}) = 200 \text{ GeV}$ ATLAS-CONF-2012-152	
Gravitino LSP	0	mono-jet	Yes	10.5	$\tilde{g}^{\text{stop}} \text{ scale}$ 645 GeV	$m(\tilde{G}) = 10^{-4} \text{ eV}$ ATLAS-CONF-2012-147	
3^{rd} gen. \tilde{g} med.	$\tilde{g} \rightarrow b\tilde{b}^*$	0	3 b	Yes	20.1	$\tilde{\chi}_1^0, \tilde{\chi}_2^0$ 1.25 TeV	$m(\tilde{t}_1) < 400 \text{ GeV}$ 1407.0600
	$\tilde{g} \rightarrow t\tilde{t}^*$	0	7-10 jets	Yes	20.3	$\tilde{\chi}_1^0, \tilde{\chi}_2^0$ 1.1 TeV	$m(\tilde{t}_1) < 350 \text{ GeV}$ 1308.1841
	$\tilde{g} \rightarrow \tau\tilde{\tau}^*$	$0-1 e, \mu$	3 b	Yes	20.1	$\tilde{\chi}_1^0, \tilde{\chi}_2^0$ 1.34 TeV	$m(\tilde{t}_1) < 400 \text{ GeV}$ 1407.0600
	$\tilde{g} \rightarrow b\tilde{t}^*$	$0-1 e, \mu$	3 b	Yes	20.1	$\tilde{\chi}_1^0, \tilde{\chi}_2^0$ 1.3 TeV	$m(\tilde{t}_1) < 300 \text{ GeV}$ 1407.0600
3 rd gen. squarks direct production	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{b}^*$	0	2 b	Yes	20.1	\tilde{b}_1 100-620 GeV	$m(\tilde{t}_1) < 90 \text{ GeV}$ 1308.2631
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow t\tilde{t}^*$	$2 e, \mu$ (SS)	0-3 b	Yes	20.3	\tilde{b}_1 275-440 GeV	$m(\tilde{t}_1) = 2 m(\tilde{t}_1)$ 1404.2500
	$\tilde{t}_1\tilde{t}_1$ (light), $\tilde{t}_1 \rightarrow b\tilde{t}^*$	$1-2 e, \mu$	1-2 b	Yes	4.7	\tilde{t}_1 110-167 GeV	$m(\tilde{t}_1) = 55 \text{ GeV}$ 1208.4305, 1209.2102
	$\tilde{t}_1\tilde{t}_1$ (light), $\tilde{t}_1 \rightarrow W\tilde{b}^*$	$2 e, \mu$	0-2 jets	Yes	20.3	\tilde{t}_1 130-210 GeV	$m(\tilde{t}_1) = m(\tilde{t}_1), m(W) = 50 \text{ GeV}, m(\tilde{\tau}^*) < m(\tilde{t}_1^*)$ 1403.4853
	$\tilde{t}_1\tilde{t}_1$ (medium), $\tilde{t}_1 \rightarrow t\tilde{t}^*$	$2 e, \mu$	2 jets	Yes	20.3	\tilde{t}_1 215-530 GeV	$m(\tilde{t}_1) = 1 \text{ GeV}$ 1403.4853
	$\tilde{t}_1\tilde{t}_1$ (medium), $\tilde{t}_1 \rightarrow b\tilde{t}^*$	0	2 b	Yes	20.1	\tilde{t}_1 150-580 GeV	$m(\tilde{t}_1) < 200 \text{ GeV}, m(\tilde{t}_1^*) = m(\tilde{t}_1) + 5 \text{ GeV}$ 1308.2631
	$\tilde{t}_1\tilde{t}_1$ (heavy), $\tilde{t}_1 \rightarrow t\tilde{t}^*$	$1 e, \mu$	1 b	Yes	20.1	\tilde{t}_1 210-640 GeV	$m(\tilde{t}_1) = 0 \text{ GeV}$ 1407.0583
	$\tilde{t}_1\tilde{t}_1$ (heavy), $\tilde{t}_1 \rightarrow b\tilde{t}^*$	0	2 b	Yes	20.1	\tilde{t}_1 260-640 GeV	$m(\tilde{t}_1) = 0 \text{ GeV}$ 1406.1122
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{t}^*$	0	1 b	Yes	20.3	\tilde{t}_1 90-240 GeV	$m(\tilde{t}_1) = m(\tilde{t}_1^*) + 85 \text{ GeV}$ 1407.0608
	$\tilde{t}_1\tilde{t}_1$ (natural GMSB)	$2 e, \mu$ (Z)	1 b	Yes	20.3	\tilde{t}_1 150-580 GeV	$m(\tilde{t}_1) = 150 \text{ GeV}$ 1403.5222
$\tilde{t}_1\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$	$3 e, \mu$ (Z)	1 b	Yes	20.3	\tilde{t}_2 290-600 GeV	$m(\tilde{t}_1) < 200 \text{ GeV}$ 1403.5222	
EW direct	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{t}^*$	$2 e, \mu$	0	Yes	20.3	\tilde{t}_1 90-325 GeV	$m(\tilde{t}_1) = 0 \text{ GeV}$ 1403.5294
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow \ell\nu(\tilde{\nu})$	$2 e, \mu$	0	Yes	20.3	\tilde{t}_1 140-465 GeV	$m(\tilde{t}_1) = 0 \text{ GeV}, m(\tilde{\nu}, \tilde{\nu}^*) = 0.5(m(\tilde{t}_1^*) + m(\tilde{t}_1))$ 1403.5294
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow \tau\nu(\tilde{\nu})$	2τ	-	Yes	20.3	\tilde{t}_1 100-350 GeV	$m(\tilde{t}_1) = 0 \text{ GeV}, m(\tilde{\nu}, \tilde{\nu}^*) = 0.5(m(\tilde{t}_1^*) + m(\tilde{t}_1))$ 1407.0350
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow \tau\nu(\tilde{\nu})$	$3 e, \mu$	0	Yes	20.3	\tilde{t}_1, \tilde{t}_2 700 GeV	$m(\tilde{t}_1^*) = m(\tilde{t}_2^*), m(\tilde{t}_1) = 0, m(\tilde{\nu}, \tilde{\nu}^*) = 0.5(m(\tilde{t}_1^*) + m(\tilde{t}_1))$ 1402.7029
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow W\tilde{b}^*$	$2-3 e, \mu$	0	Yes	20.3	\tilde{t}_1, \tilde{t}_2 420 GeV	$m(\tilde{t}_1) = m(\tilde{t}_2), m(\tilde{t}_1) = 0$, sleptons decoupled 1403.5294, 1402.7029
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow W\tilde{b}^* Z$	$1 e, \mu$	2 b	Yes	20.3	\tilde{t}_1, \tilde{t}_2 285 GeV	$m(\tilde{t}_1) = m(\tilde{t}_2), m(\tilde{t}_1) = 0$, sleptons decoupled ATLAS-CONF-2013-093
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow W\tilde{b}^* \tilde{\chi}_1^0$	$1 e, \mu$	0	Yes	20.3	\tilde{t}_1, \tilde{t}_2 620 GeV	$m(\tilde{t}_1) = m(\tilde{t}_2), m(\tilde{t}_1) = 0, m(\tilde{\nu}, \tilde{\nu}^*) = 0.5(m(\tilde{t}_1^*) + m(\tilde{t}_1))$ 1405.5086
Long-lived particles	Direct $\tilde{t}_1\tilde{t}_1$ prod., long-lived \tilde{t}_1	Disapp. trk	1 jet	Yes	20.3	\tilde{t}_1 270 GeV	$m(\tilde{t}_1) = m(\tilde{t}_1^*) = 160 \text{ MeV}, \tau(\tilde{t}_1) = 0.2 \text{ ns}$ ATLAS-CONF-2013-069
	Stable, stopped \tilde{R} -hadron	0	1-5 jets	Yes	27.9	\tilde{t}_1 832 GeV	$m(\tilde{t}_1) = 100 \text{ GeV}, 10 \mu\text{s} < \tau < 1000 \text{ s}$ 1310.6584
	GMSB, stable $\tilde{t}_1, \tilde{t}_1 \rightarrow \tau(\tilde{\nu}, \tilde{\mu}) + \tau(e, \mu)$	$1-2 \mu$	-	-	15.9	\tilde{t}_1 475 GeV	$10 < \tan\beta < 50$ ATLAS-CONF-2013-058
	GMSB, $\tilde{t}_1 \rightarrow \tau(\tilde{\nu}, \tilde{\mu}) + \tau(e, \mu)$	2γ	-	Yes	4.7	\tilde{t}_1 230 GeV	$0.4 < \tau < 1 \text{ ns} < 2 \text{ ns}$ 1304.6310
$\tilde{q}\tilde{q}, \tilde{t}_1 \rightarrow \tilde{q}\tilde{q}^*$ (RPV)	$1 \mu, \text{ displ. vtx}$	-	-	20.3	\tilde{t}_1 1.0 TeV	$1.5 < \tau < 156 \text{ mm}, \text{BR}(\mu) = 1, m(\tilde{t}_1) = 108 \text{ GeV}$ ATLAS-CONF-2013-092	
RPV	LFV $p\tilde{p} \rightarrow \tilde{X}, \tilde{X}, \tilde{X} \rightarrow e\mu$	$2 e, \mu$	-	-	4.6	$\tilde{X}_{111} = 0.10, \tilde{X}_{122} = 0.05$ 1212.1272	
	LFV $p\tilde{p} \rightarrow \tilde{X}, \tilde{X}, \tilde{X} \rightarrow e(\mu) + \tau$	$1 e, \mu + \tau$	-	-	4.6	$\tilde{X}_{111} = 0.10, \tilde{X}_{123} = 0.05$ 1212.1272	
	Bilinear RPV CMSSM	$2 e, \mu$ (SS)	0-3 b	Yes	20.3	\tilde{q}, \tilde{g} 1.1 TeV 1.35 TeV	$m(\tilde{g}) = m(\tilde{g}), c_{T,SS} < 1$ 1404.2500
Other	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow W\tilde{b}^*$	$4 e, \mu$	-	Yes	20.3	\tilde{t}_1 750 GeV	$m(\tilde{t}_1) = 0.2 m(\tilde{t}_1^*), \tilde{X}_{121} = 0$ 1405.5086
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow W\tilde{b}^*, \tilde{t}_1 \rightarrow \tau\tilde{\nu}, \tilde{e}\tilde{\nu}$	$3 e, \mu + \tau$	-	Yes	20.3	\tilde{t}_1 450 GeV	$m(\tilde{t}_1) = 0.2 m(\tilde{t}_1^*), \tilde{X}_{133} = 0$ 1405.5086
	$\tilde{g} \rightarrow \tilde{q}\tilde{q}^*$	0	6-7 jets	-	20.3	\tilde{g} 916 GeV	$\text{BR}(\tau) = \text{BR}(b) = \text{BR}(c) = 0\%$ ATLAS-CONF-2013-091
	$\tilde{g} \rightarrow \tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{t}^*$	$2 e, \mu$ (SS)	0-3 b	Yes	20.3	\tilde{g} 850 GeV	1404.2510
Other	Scalar gluon pair, $sgluon \rightarrow \tilde{q}\tilde{q}$	0	4 jets	-	4.6	$sgluon$ 100-287 GeV	incl. limit from 1110.2693 1210.4826
	Scalar gluon pair, $sgluon \rightarrow \tilde{t}\tilde{t}$	$2 e, \mu$ (SS)	2 b	Yes	14.3	$sgluon$ 350-800 GeV	ATLAS-CONF-2013-051
	WIMP interaction (DS, Dirac χ)	0	mono-jet	Yes	10.5	$M^2 \text{ scale}$ 704 GeV	$m(\tilde{\chi}) = 80 \text{ GeV}$, limit of $\sim 687 \text{ GeV}$ for D8 ATLAS-CONF-2012-147

$\sqrt{s} = 7 \text{ TeV}$
full data

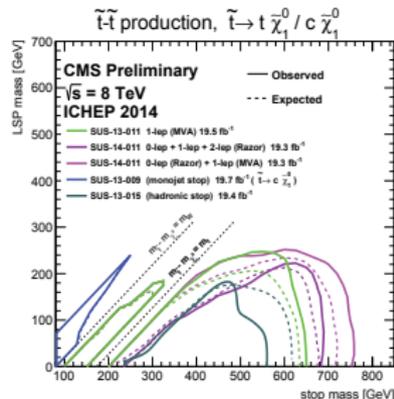
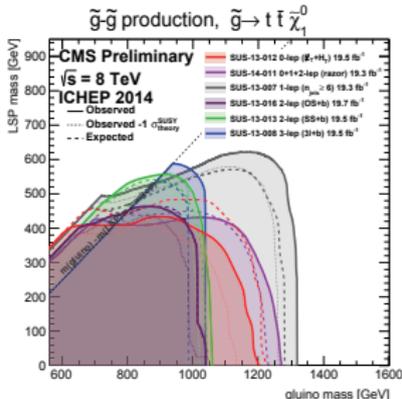
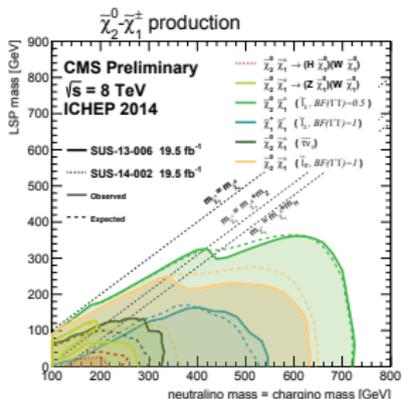
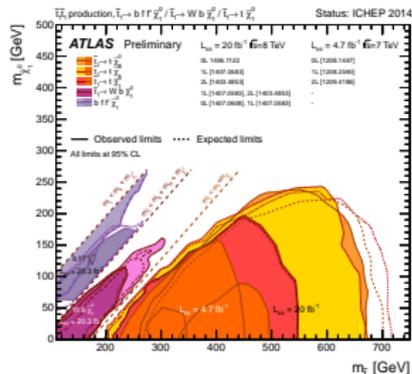
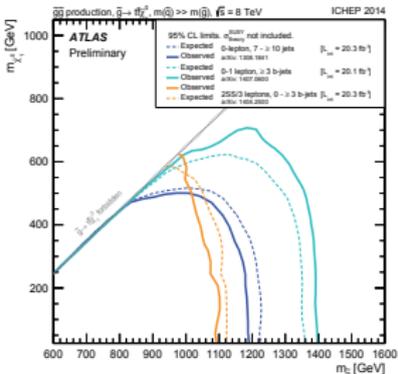
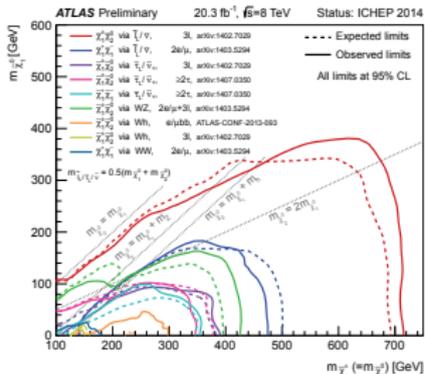
$\sqrt{s} = 8 \text{ TeV}$
partial data

$\sqrt{s} = 8 \text{ TeV}$
full data

10⁻¹ 1 Mass scale [TeV]

*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty.

Strong limits on SUSY partners



Tension with Higgs mass:

$$M_h^2 \leq M_Z^2 \cos^2(2\beta) + \epsilon$$

Large radiative corrections needed:

$$M_S^2 = m_{\tilde{t}_1} m_{\tilde{t}_2}$$

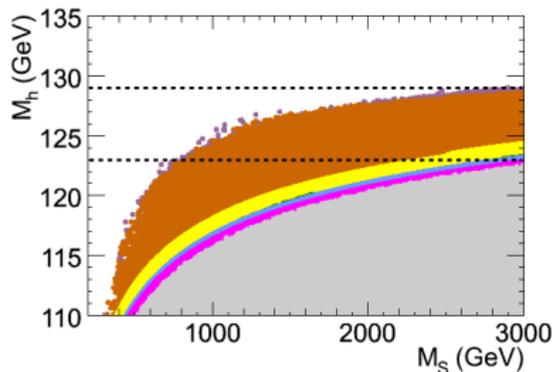
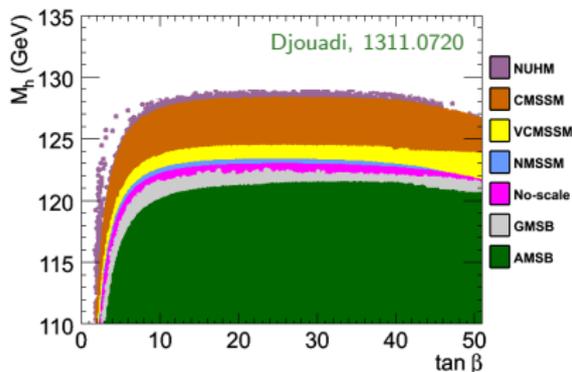
$$\epsilon \approx \frac{3m_t^4}{2\pi^2 v^2 \sin^2 \beta} \left[\log \frac{M_S^2}{m_t^2} + \frac{X_t^2}{M_S^2} \left(1 - \frac{X_t^2}{12M_S^2} \right) \right]$$



Decoupling limit ($M_A \gg M_Z$),

$\cos^2(2\beta) \rightarrow 1$

Maximal stop mixing $X_t = A_t - \mu \cot \beta$

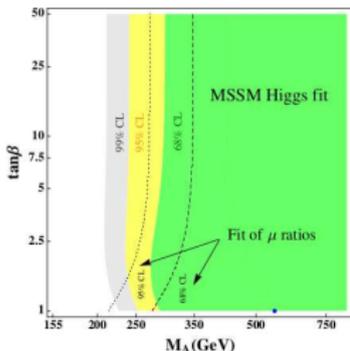


Improved higher-order calculations allow slightly larger values of M_h

Hahn et al, 1312.4937

Constraints from Higgs Decay

Djouadi, 1311.0720

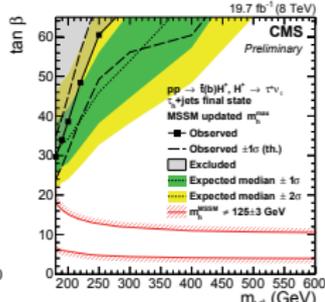
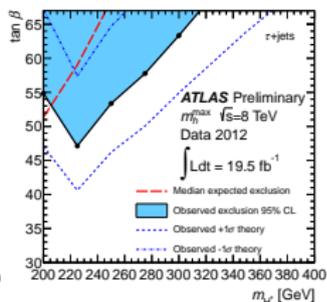
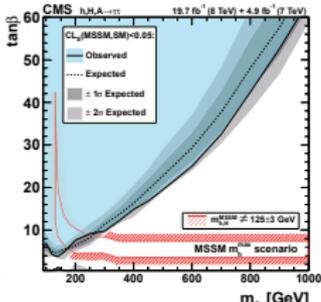
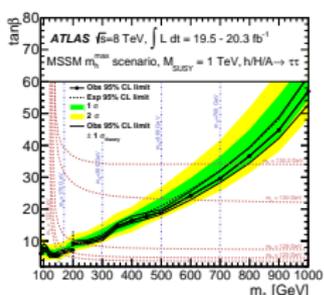


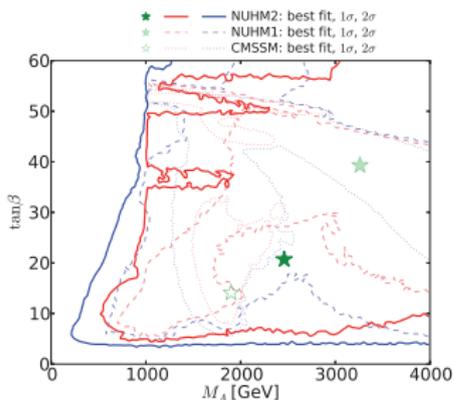
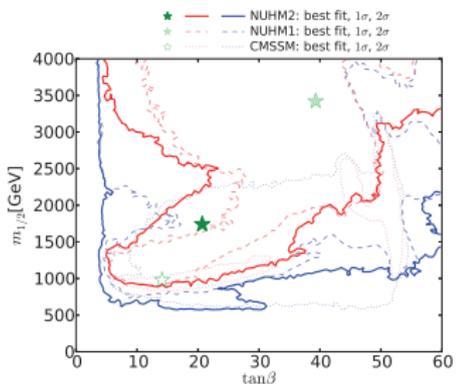
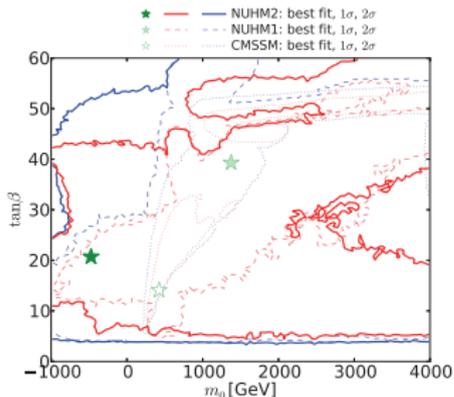
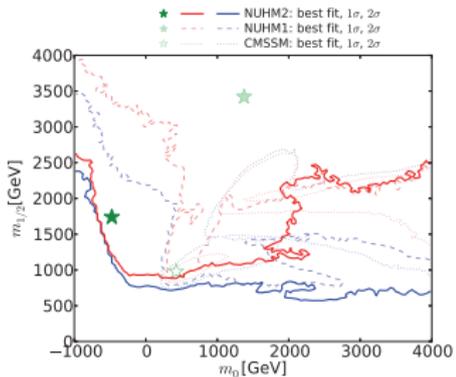
$$c_t \approx \frac{\cos \alpha}{\sin \beta} \left[1 + \frac{m_t^2}{4m_{\tilde{t}_1}^2 m_{\tilde{t}_2}^2} (m_{\tilde{t}_1}^2 + m_{\tilde{t}_2}^2 - X_t^2) \right]$$

$$c_b \approx -\frac{\sin \alpha}{\cos \beta} \left[1 - \frac{\Delta_b}{1 + \Delta_b} (1 + \cot \alpha \cot \beta) \right]$$

$$c_V = \sin(\beta - \alpha) \quad , \quad \Delta_b \approx \frac{2\alpha_s}{3\pi} \frac{\mu m_{\tilde{g}} \tan \beta}{\max(m_{\tilde{g}}^2, m_{\tilde{b}_1}^2, m_{\tilde{b}_2}^2)}$$

Heavy Higgs Searches





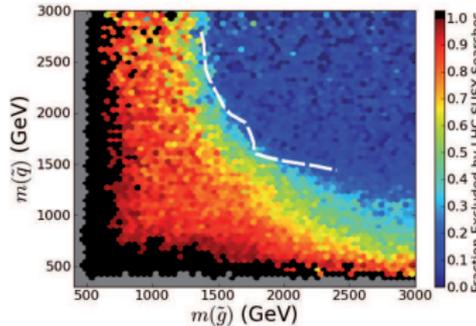
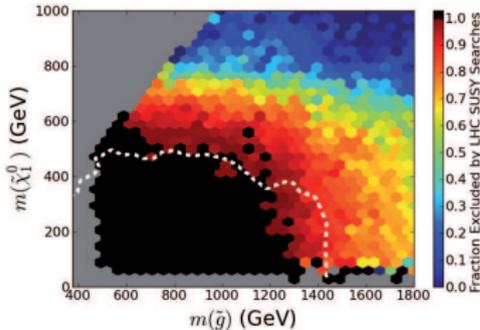
$(g - 2)_\mu$ cannot be explained (not included in the fit)

Which SUSY ?

- Looks bad in **CMSSM** (120 MSSM parameters reduced to 4 + 1 sign)
- More freedom in the **Phenomenological MSSM**

Many “models” consistent with data

Cahill-Rowley et al, 1407.4130



19–20 parameters

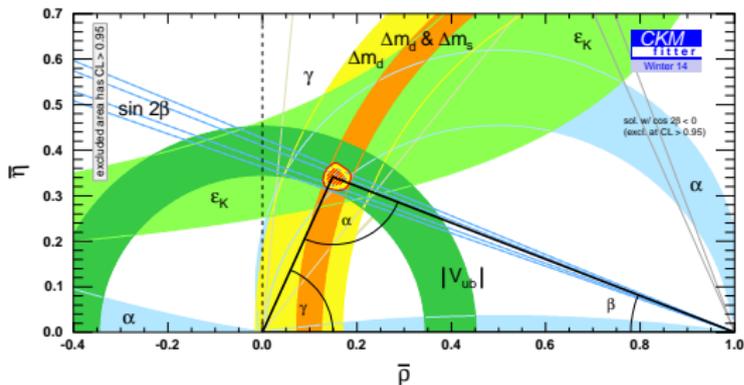
Data-driven search

- Many SUSY variants: **NMSSM, Split, High-Scale, Stealth, 5D, Natural, Folded, Twin...**

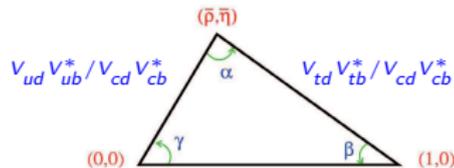
Naturalness?

$$\Delta M_h^2 \propto M_{\text{SUSY}}^2$$

Quark Mixing



$$V_{ud} V_{ub}^* + V_{cd} V_{cb}^* + V_{td} V_{tb}^* = 0$$



$$V = \begin{bmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{bmatrix} + \mathcal{O}(\lambda^4)$$

$$\bar{\eta} \equiv \eta \left(1 - \frac{1}{2} \lambda^2\right) = 0.351 \pm 0.014$$

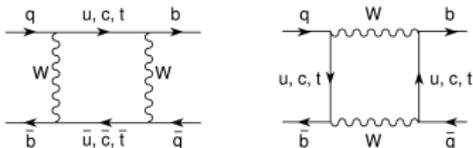
UT_{fit}

$$\bar{\rho} \equiv \rho \left(1 - \frac{1}{2} \lambda^2\right) = 0.132 \pm 0.023$$

$$A = 0.821 \pm 0.012 \quad ; \quad \lambda = 0.2254 \pm 0.0006$$

Successful CKM Mechanism (Tree / Loop / CP-c / CP-v)

Bounds on New Flavour Physics



$$L_{\text{eff}} = L_{\text{SM}} + \sum_{D>4} \sum_k \frac{C_k^{(D)}}{\Lambda_{\text{NP}}^{D-4}} O_k^{(D)}$$

Isidori, 1302.0661

Operator	Bounds on Λ in TeV ($c_{\text{NP}} = 1$)		Bounds on c_{NP} ($\Lambda = 1$ TeV)		Observables
	Re	Im	Re	Im	
$(\bar{s}_L \gamma^\mu d_L)^2$	9.8×10^2	1.6×10^4	9.0×10^{-7}	3.4×10^{-9}	$\Delta m_K; \epsilon_K$
$(\bar{s}_R d_L)(\bar{s}_L d_R)$	1.8×10^4	3.2×10^5	6.9×10^{-9}	2.6×10^{-11}	$\Delta m_K; \epsilon_K$
$(\bar{c}_L \gamma^\mu u_L)^2$	1.2×10^3	2.9×10^3	5.6×10^{-7}	1.0×10^{-7}	$\Delta m_D; q/p , \phi_D$
$(\bar{c}_R u_L)(\bar{c}_L u_R)$	6.2×10^3	1.5×10^4	5.7×10^{-8}	1.1×10^{-8}	$\Delta m_D; q/p , \phi_D$
$(\bar{b}_L \gamma^\mu d_L)^2$	6.6×10^2	9.3×10^2	2.3×10^{-6}	1.1×10^{-6}	$\Delta m_{B_d}; S_{\psi K_S}$
$(\bar{b}_R d_L)(\bar{b}_L d_R)$	2.5×10^3	3.6×10^3	3.9×10^{-7}	1.9×10^{-7}	$\Delta m_{B_d}; S_{\psi K_S}$
$(\bar{b}_L \gamma^\mu s_L)^2$	1.4×10^2	2.5×10^2	5.0×10^{-5}	1.7×10^{-5}	$\Delta m_{B_s}; S_{\psi\phi}$
$(\bar{b}_R s_L)(\bar{b}_L s_R)$	4.8×10^2	8.3×10^2	8.8×10^{-6}	2.9×10^{-6}	$\Delta m_{B_s}; S_{\psi\phi}$

- Generic flavour structure [$c_{\text{NP}} \sim \mathcal{O}(1)$] ruled out at the TeV scale
- $\Lambda_{\text{NP}} \sim 1$ TeV requires c_{NP} to inherit the strong SM suppressions (GIM)

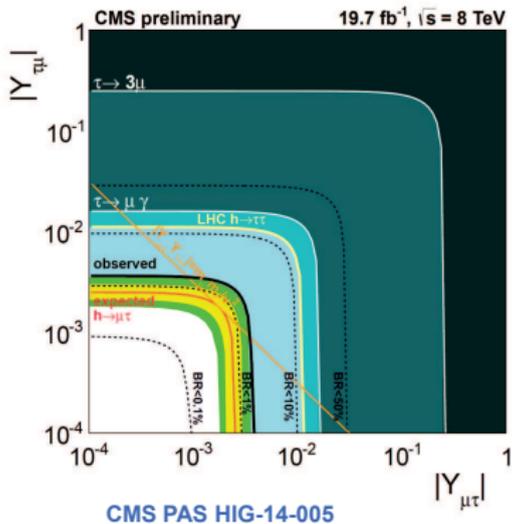
Minimal Flavour Violation: The up and down Yukawa matrices are the only source of quark-flavour symmetry breaking

D'Ambrosio et al, Buras et al

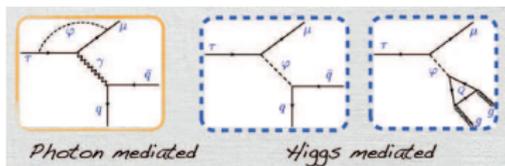
Flavour-Violating Higgs Couplings

Blankenburg et al, Celis et al,
Harnik et al, Davidson-Verdier,
Kopp-Nardecchia

$$L = -h \left\{ Y_{e\mu} \bar{e}_L \mu_R + Y_{e\tau} \bar{e}_L \tau_R + Y_{\mu\tau} \bar{\mu}_L \tau_R + \dots \right\}$$



$Br(H \rightarrow \mu\tau) < 1.57\%$ (95% CL)



Celis

