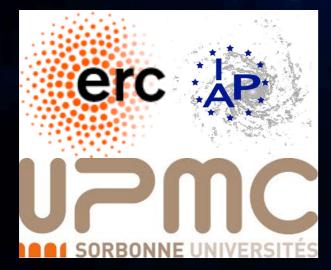
#### Is Coy dark matter really coy?

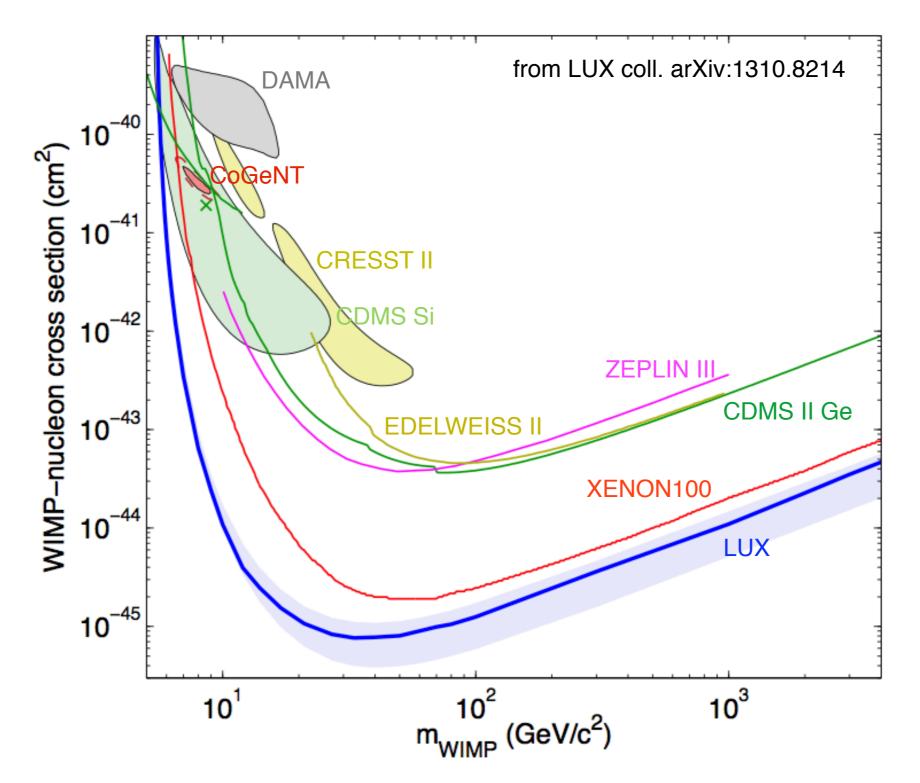
based on PRL 114 (2015) 011301 in collaboration with E. Del Nobile and P. Panci

**Chiara Arina** 



Rencontre de Physique des Particules, Institut Henri Poincaré January 16<sup>th</sup> 2015

#### General status of Dark Matter (DM) direct detection

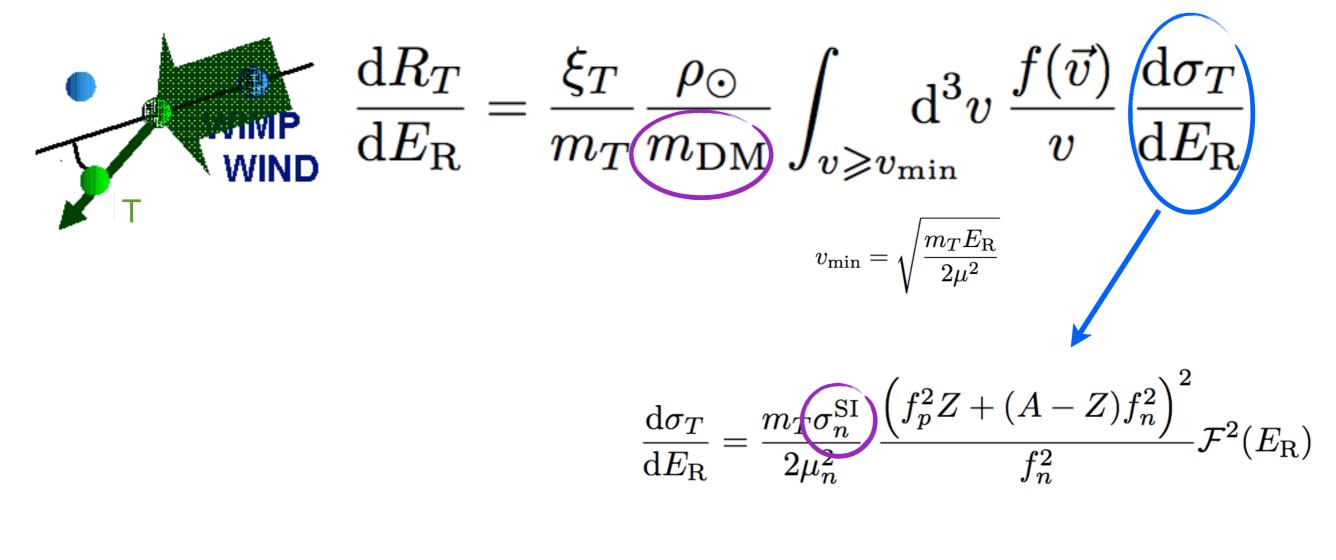


1. Scalar spin-independent (SI) and model independent interaction DM-nucleus

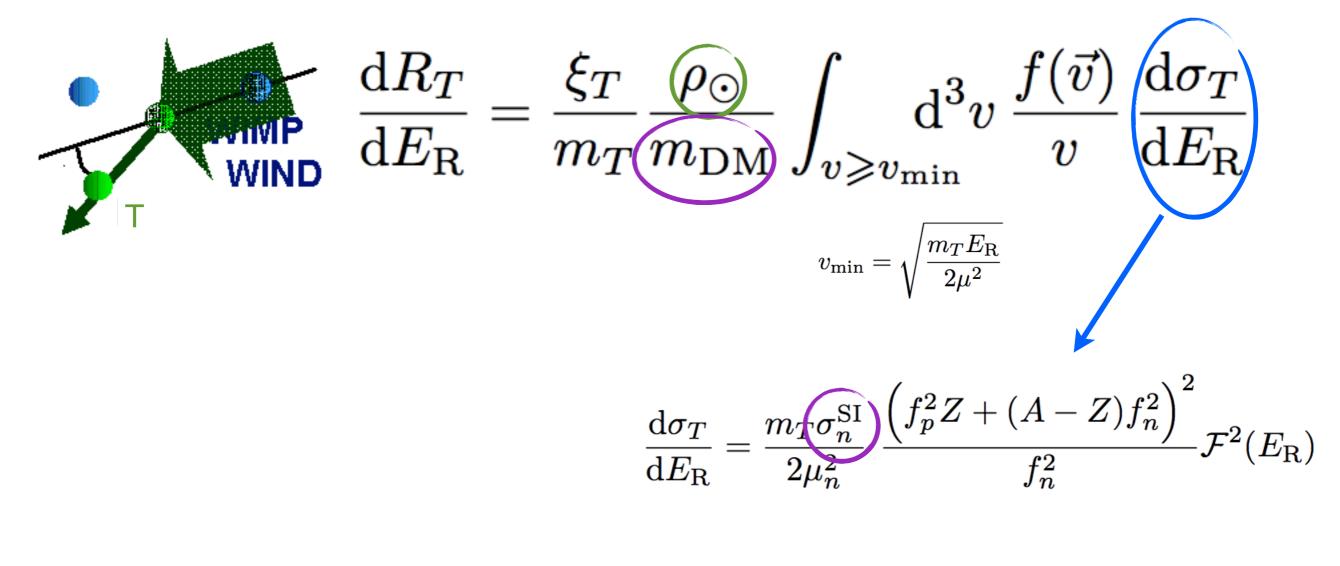
2. Fixed astrophysics in the event rate (DM velocity distribution and astro parameters)

$$\frac{\mathrm{d}R_T}{\mathrm{d}E_R} = \frac{\xi_T}{m_T} \frac{\rho_{\odot}}{m_{\mathrm{DM}}} \int_{v \geqslant v_{\mathrm{min}}} \frac{\mathrm{d}^3 v}{v} \frac{f(\vec{v})}{v} \frac{\mathrm{d}\sigma_T}{\mathrm{d}E_R}$$
$$\frac{v_{\mathrm{min}}}{v_{\mathrm{min}} = \sqrt{\frac{m_T E_R}{2\mu^2}}}$$
$$\frac{\mathrm{d}\sigma_T}{\mathrm{d}E_R} = \frac{m_T \sigma_n^{\mathrm{SI}}}{2\mu_n^2} \frac{\left(f_p^2 Z + (A - Z)f_n^2\right)^2}{f_n^2} \mathcal{F}^2(E_R)$$

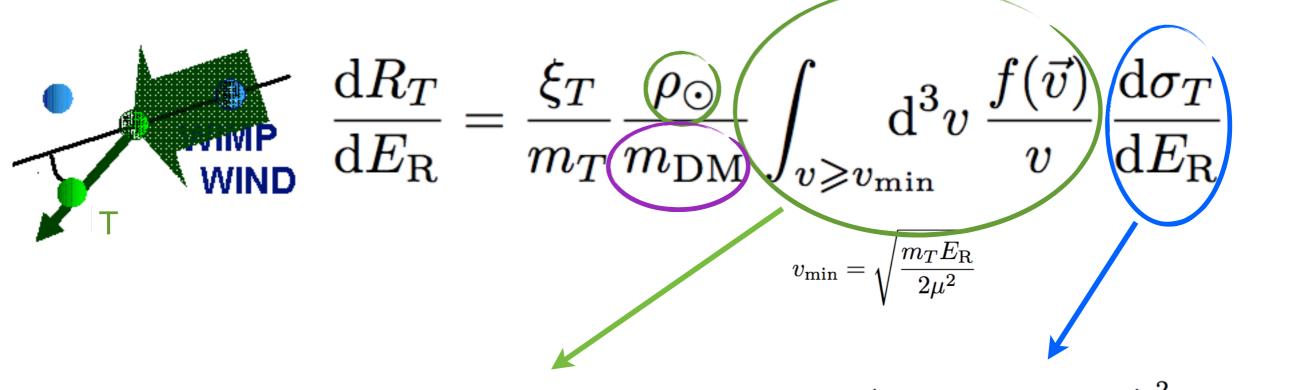
•  $f_n$  and  $f_p$  can be different!



• fn and fp can be different!



• fn and fp can be different!



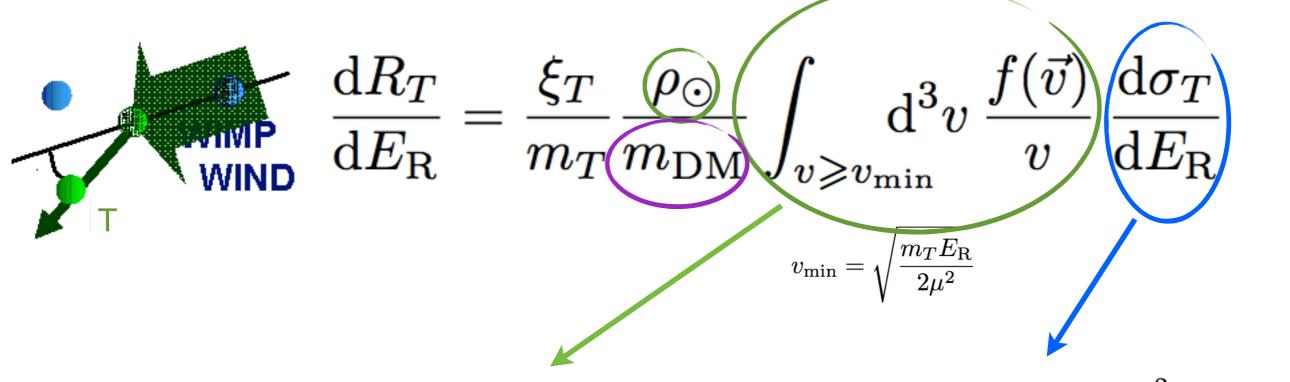
Depend on DM velocity distribution and astrophysical parameters

- f(v) is not a Maxwell-Boltzmann distribution!
- Astrophysical parameters not well measured

$$v_0^{
m obs} = 230 \pm 24.4 \ {
m km \ s^{-1}} \ v_{
m esc}^{
m obs} = 544 \pm 39 \ {
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ho_{\odot}^{
m obs} = 0.4 \pm 0.2 \ {
m GeV \ cm^{-3}}$$

$$\frac{\mathrm{d}\sigma_T}{\mathrm{d}E_{\mathrm{R}}} = \frac{m_{\mathrm{T}}\sigma_n^{\mathrm{SI}}}{2\mu_n^2} \frac{\left(f_p^2 Z + (A-Z)f_n^2\right)^2}{f_n^2} \mathcal{F}^2(E_{\mathrm{R}})$$

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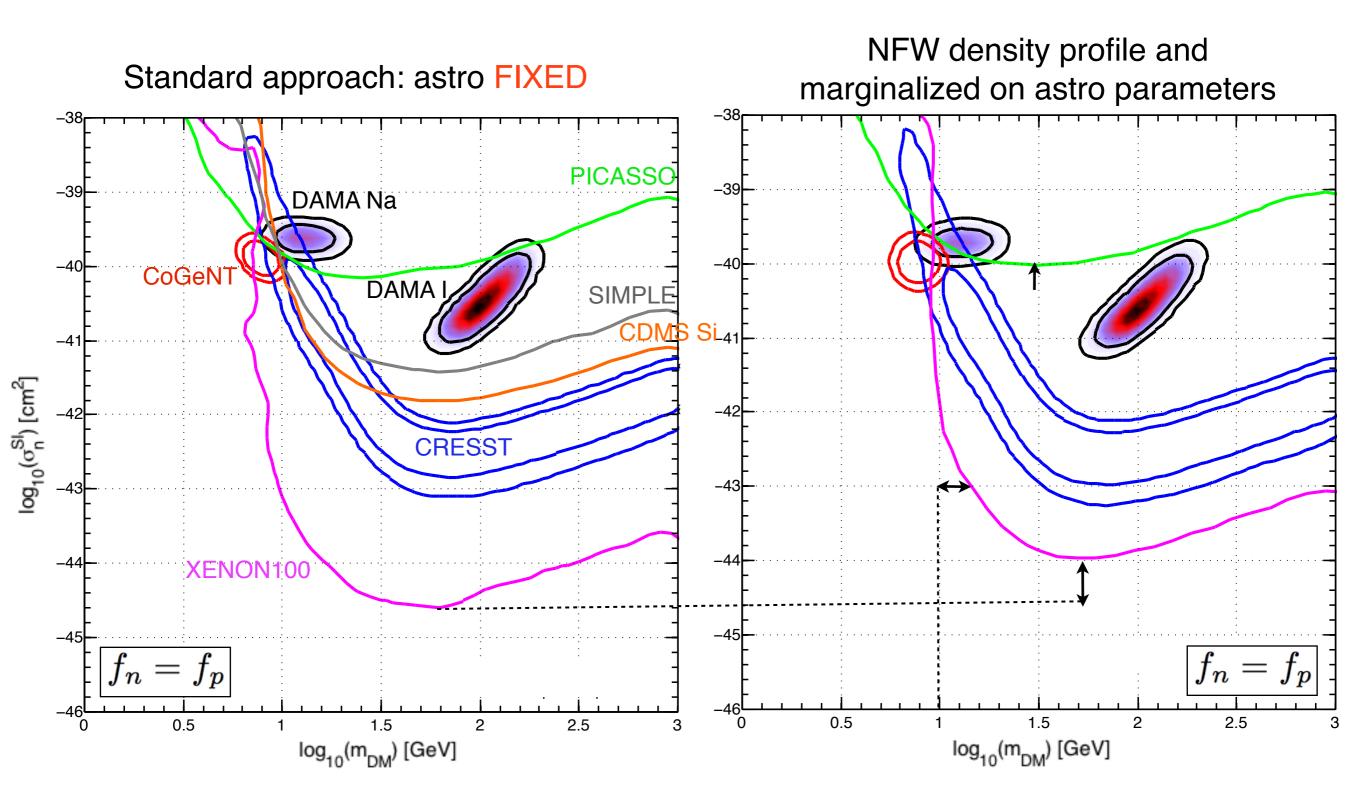
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•  $f_n$  and  $f_p$  can be different!

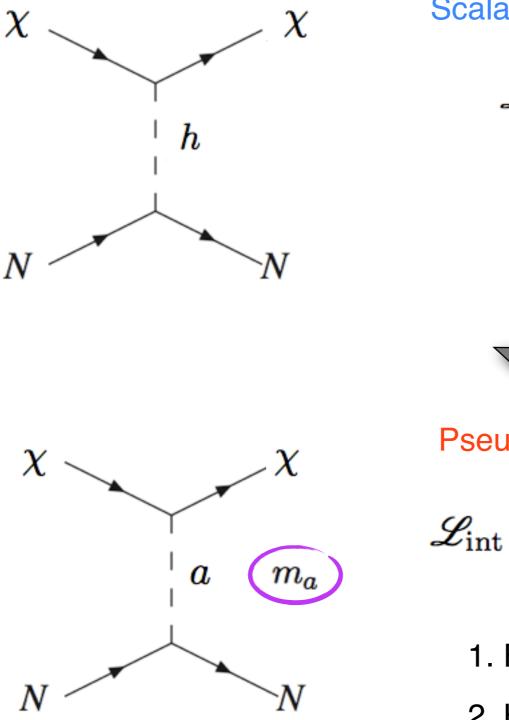
# All these effects can be simultaneously accounted for within Bayesian statistics

#### **Effect of ASTROPHYSICAL uncertainties**



CA, review for PDU, arXiv:1310.5718

# **Changing the DM-nucleus interaction**



Scalar SI comes from e.g. interaction with Higgs:

$$\mathscr{L}_{\mathrm{int}} = -\frac{y_{\mathrm{DM}}}{\sqrt{2}}h\bar{\chi}\chi - \sum_{f}\frac{y_{f}}{\sqrt{2}}h\bar{f}f$$

 $f_n \sim f_p$ 

Pseudo-scalar interaction (Coy DM):

$$\mathscr{L}_{\mathrm{int}} = -i \frac{g_{\mathrm{DM}}}{\sqrt{2}} a \bar{\chi} \gamma_5 \chi - ig \sum_q \frac{g_q}{\sqrt{2}} a \bar{q} \gamma_5 q$$

1. Flavor-Universal couplings:  $g_q = 1$ 2. Higgs-like:  $g_q = \frac{m_q}{174 \, {
m GeV}}$ 

3 free parameters

**Coy DM effective operator I** 

$$\mathscr{L}_{\mathrm{int}} = -i \frac{g_{\mathrm{DM}}}{\sqrt{2}} a \bar{\chi} \gamma_5 \chi - ig \sum_q \frac{g_q}{\sqrt{2}} a \bar{q} \gamma_5 q$$

Can be written in terms of effective contact operator (DM typical velocity is 10<sup>-3</sup>c)

$$\mathscr{L}_{\text{eff}} = \frac{1}{2\Lambda_a^2} \sum_{N=p,n} g_N \, \bar{\chi} \gamma^5 \chi \, \bar{N} \gamma^5 N \qquad \qquad \Lambda_a \equiv \frac{m_a}{\sqrt{g_{\text{DM}}g}}$$

• The energy scale is the unknown variable instead of the cross-section

• The coefficients  $g_N$  are defined to be

$$g_{N} = \sum_{q=u,d,s} \frac{m_{N}}{m_{q}} \left[ g_{q} - \sum_{q'=u,\dots,t} g_{q'} \frac{\bar{m}}{m_{q'}} \right] \Delta_{q}^{(N)}$$

Flavor-Universal couplings:  $g_p/g_n = -16.4$ Higgs-like:  $g_p/g_n = -4.1$ 

NATURAL violation of isospin

#### **Coy DM effective operator II**

$$\frac{\mathrm{d}\sigma_T}{\mathrm{d}E_{\mathrm{R}}} = \frac{1}{128\pi} \frac{q^4}{\Lambda_a^4} \frac{m_T}{m_{\mathrm{DM}}^2 m_N^2} \frac{1}{v^2} \sum_{N,N'=p,n} g_N g_{N'} F_{\Sigma''}^{(N,N')}(q^2)$$

See talk by P. Panci for more details on NRO

### **Coy DM effective operator II**

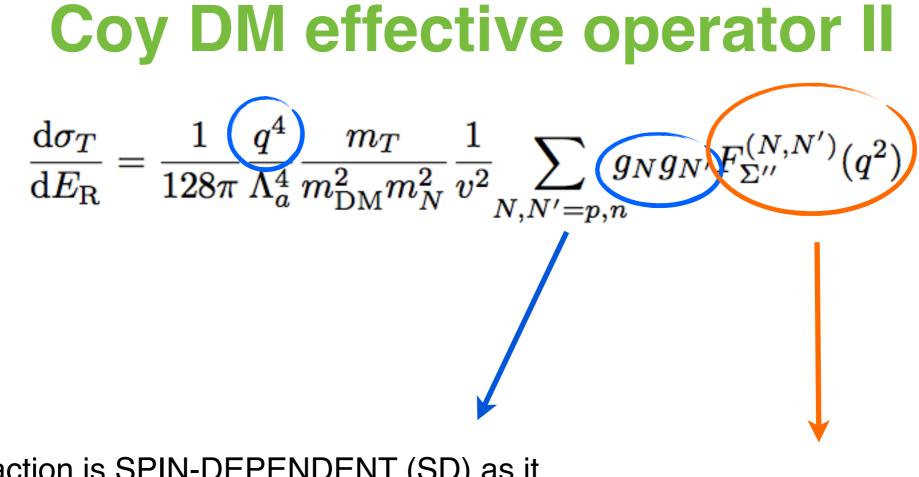
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• This interaction is SPIN-DEPENDENT (SD) as it comes from this non-relativistic operator:

 $\mathcal{O}_6^{\rm NR} = (\vec{s}_{\chi} \cdot \vec{q}) (\vec{s}_N \cdot \vec{q})$ 

- DAMA: Iodine (Sodium) has an unpaired proton
- LUX: Xenon has an unpaired neutron
- Natural isospin violation implies an strong suppression/enhancement to DM scattering

See talk by P. Panci for more details on NRO



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Nuclear form factor:

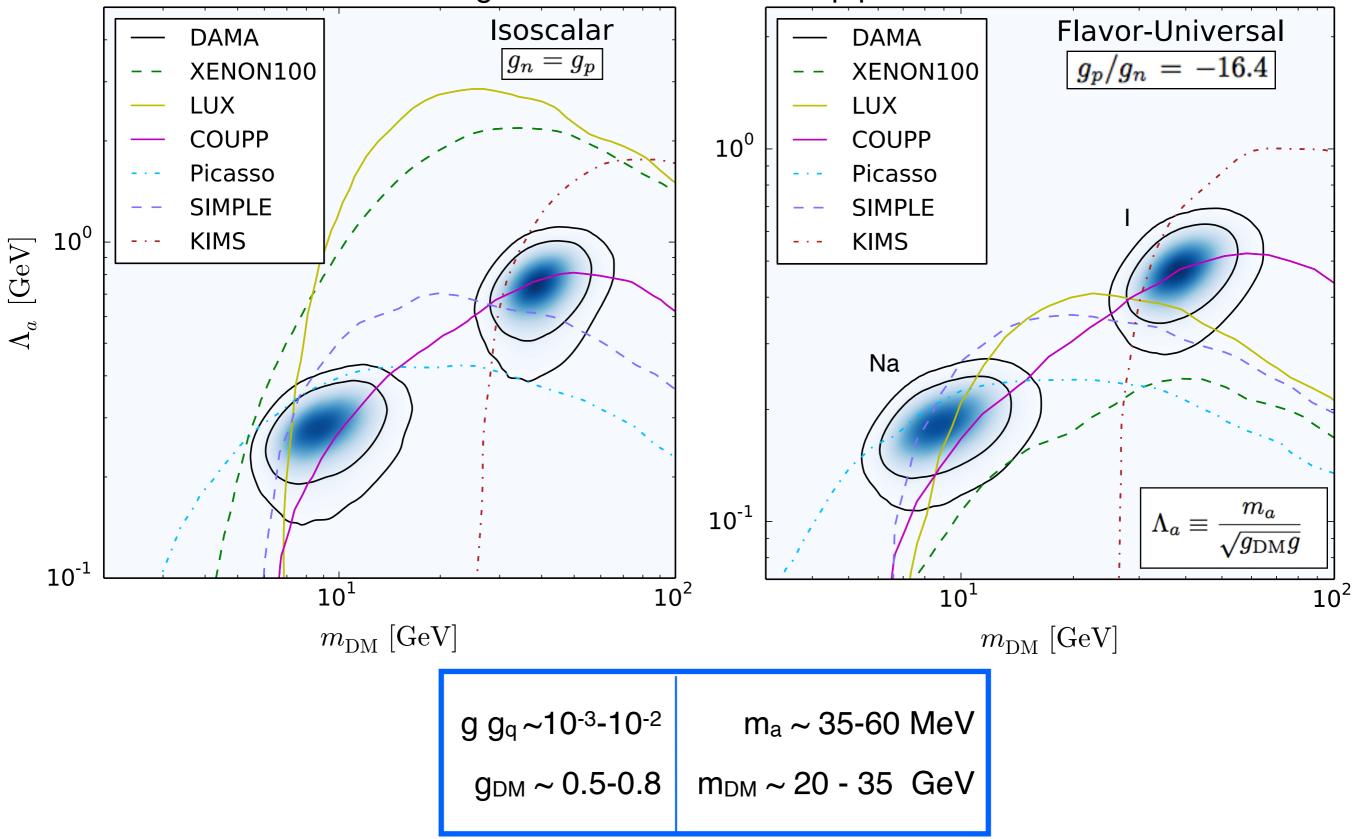
 Source of uncertainties (number of event can change by a factor ~ 3 for standard SD)

• use of the correct form factor (computed in Fitzpatrick et al. arXiv:1203.3542)

See talk by P. Panci for more details on NRO

### **Direct detection of Coy DM**

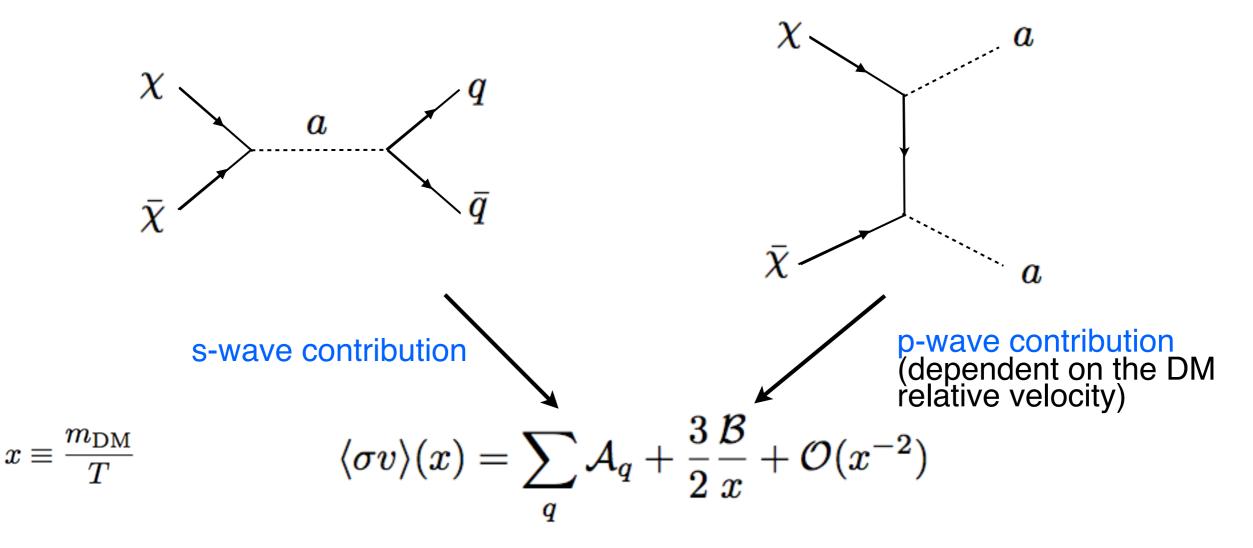
marginalized on astro and exp parameters



#### **Relic density & indirect detection I**

$$\mathscr{L}_{\mathrm{int}} = -i rac{g_{\mathrm{DM}}}{\sqrt{2}} a ar{\chi} \gamma_5 \chi - ig \sum_q rac{g_q}{\sqrt{2}} a ar{q} \gamma_5 q$$

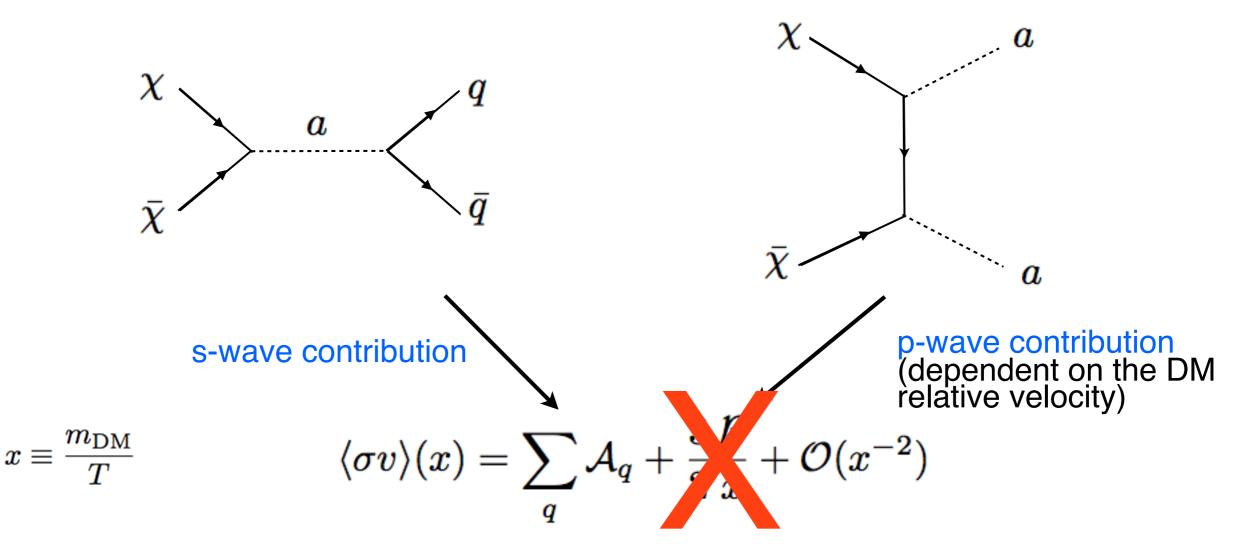
Annihilation of Coy DM mediated by two processes



#### **Relic density & indirect detection I**

$$\mathscr{L}_{\mathrm{int}} = -i rac{g_{\mathrm{DM}}}{\sqrt{2}} a ar{\chi} \gamma_5 \chi - ig \sum_q rac{g_q}{\sqrt{2}} a ar{q} \gamma_5 q$$

Annihilation of Coy DM mediated by two processes



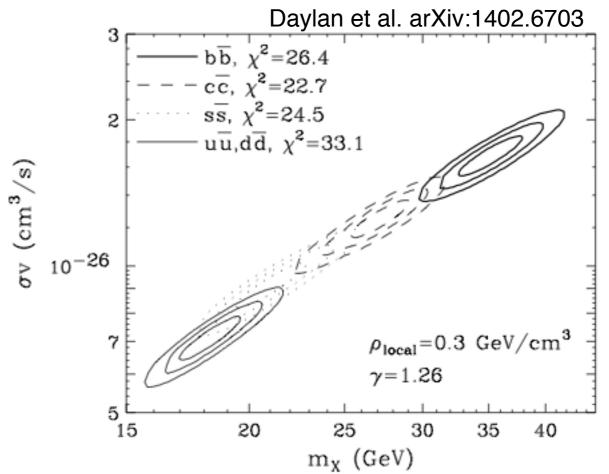
At present time ( $x_0 >>1$ ) only s-wave contribution is relevant

(i.e. for producing gamma ray flux)

## **Relic density & indirect detection II**

 There is an excess in gamma rays from 1-3 GeV around the galactic center (GC)

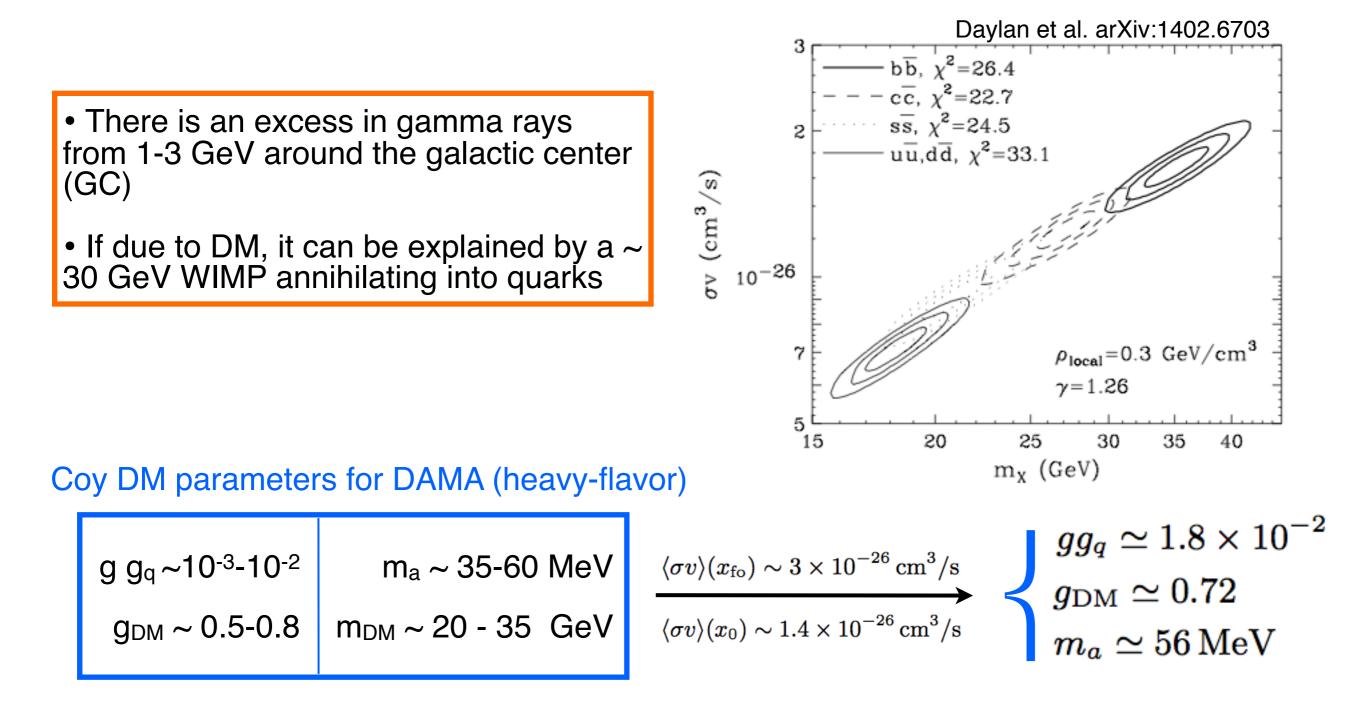
If due to DM, it can be explained by a ~
 30 GeV WIMP annihilating into quarks



Coy DM parameters for DAMA (heavy-flavor)

g g <sub>q</sub> ~10 <sup>-3</sup> -10 <sup>-2</sup>	m <sub>a</sub> ~ 35-60 MeV
g <sub>DM</sub> ~ 0.5-0.8	т <sub>DM</sub> ~ 20 - 35 GeV

# **Relic density & indirect detection II**



- Same ballpark of values explains at the same time DAMA and the gamma ray GC excess
- Model completely determined: relic density gamma ray GC excess and DAMA constraints fix the 3 free parameters

### Conclusions

- Difficult to reconcile DAMA with other exclusion bounds with SI (even allowing for all exp and astro uncertainties)
- Coy DM with light pseudo-scalar reconciles DAMA and LUX because it induces naturally isospin violation
- Coy DM explains the GC excess in gamma-rays and has the correct relic density (no free parameters left)
- Interesting phenomenological model
- Flavor constraints on the pseudo-scalar mediator are relevant (see Dolan et al. arXiv:1412.5174)
- Light pseudo-scalar interesting for self-interacting DM

### Back up slides

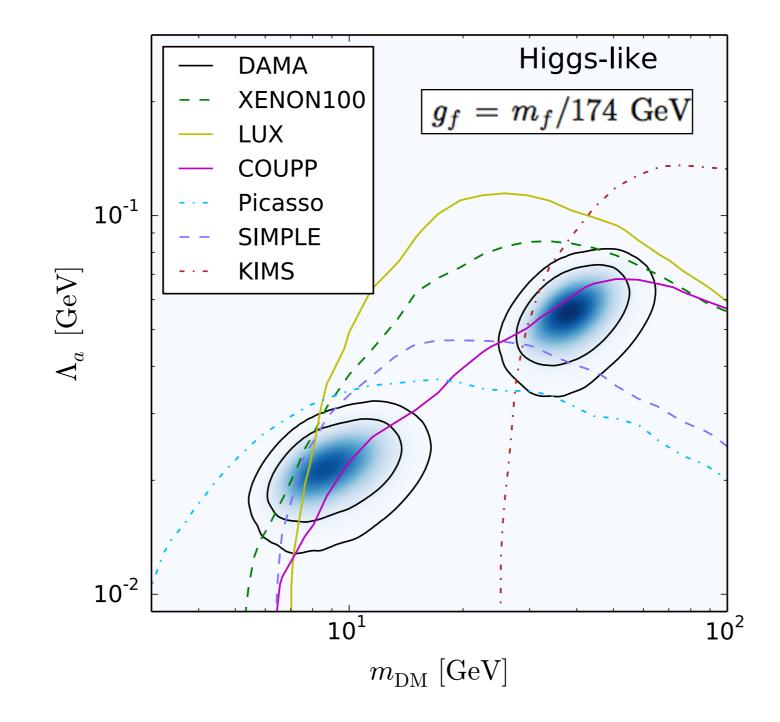
#### **Flavor constraints**

The light pseudo scalar mediator can be constrained by rare meson decays.

For  $m_a < 100$  MeV, the most constraining channels are:

Channel	Experiment	Mass range [MeV]	Relevant
$K^+ \to \pi^+ + \mathrm{inv}$	E949	0–110	Long lifetime
	$\mathbf{E787}$	0-110 & 150-260	Long lifetime
$K_L  o \pi^0  \gamma \gamma$	KTeV	40 - 100 & 160 - 350	Photonic decays
$K^+ \to \pi^+ A$	$K_{\mu 2}$	10-130 & 140-300	All decay modes
$B^0  o K^0_S + { m inv}$	CLEO	0–1100	Long lifetime
$b \rightarrow s  g$	CLEO	$m_A < m_B - m_K$	Hadronic decays
$K, B \to A + X$	CHARM	0–4000	Photonic decays

# Coy dark matter with Higgs-like couplings



### Quark spin content of the nucleon

Cheng and Chiang, arXiv:1202.1292

$$\Delta_u^{(p)} = \Delta_d^{(n)} = +0.84$$
$$\Delta_d^{(p)} = \Delta_u^{(n)} = -0.44$$
$$\Delta_s^{(p)} = \Delta_s^{(n)} = -0.03$$

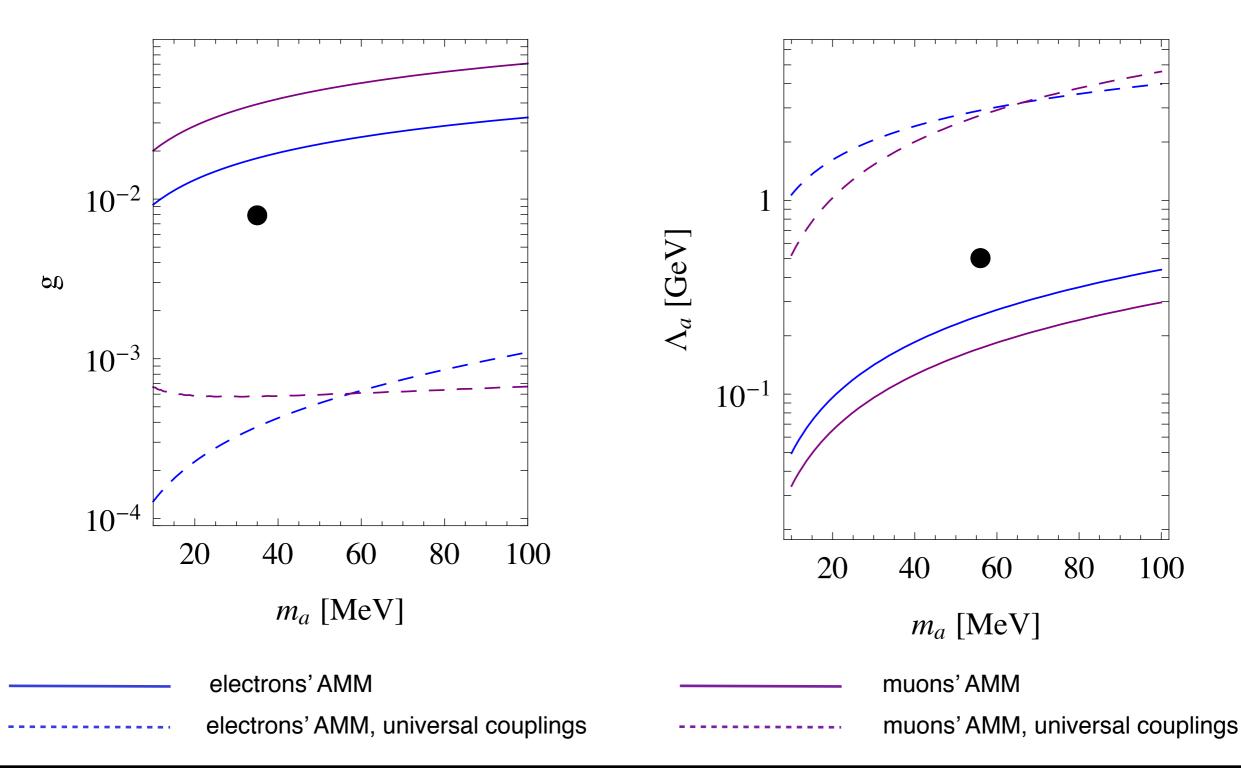
$$\bar{m} \equiv (1/m_u + 1/m_d + 1/m_s)^{-1}$$

- These are conservative values
- These coefficients are subject to large uncertainties (measured experimentally with e.g. pion scattering)
- For extreme values the natural isospin violation cabe  $g_p/g_n = -44.2$
- $\bullet$  The value of  $g_N$  coefficients does not change if the pseudo-scalar couples only to heavy flavors (s,t,b)
- This makes the pseudo-scalar an axion-like particle

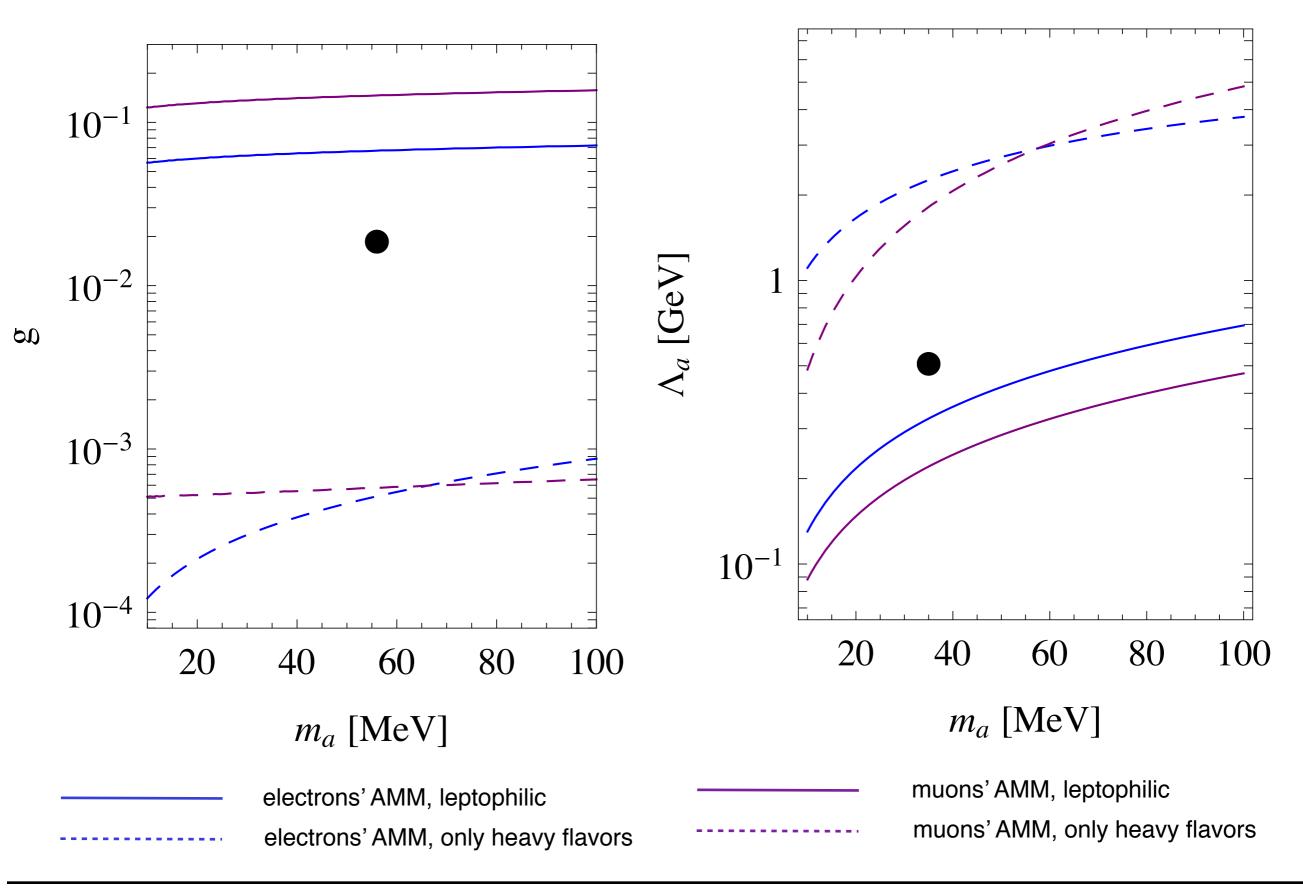
#### **Constraints on the pseudoscalar**

#### 1. ok for BaBaR constraints

2. ok for electric (moun) anomalous magnetic moments



### **Constraints from magnetic moments**



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#### **Details on annihilation cross-section**

$$\sigma(\bar{\chi}\chi o \bar{q}q) = N_{\rm c} rac{g^2 g_f^2 g_{
m DM}^2}{64\pi} rac{s}{(s-m_a^2)^2} \sqrt{rac{s-4m_q^2}{s-4m_{
m DM}^2}}$$

$$\sigma(\bar{\chi}\chi \to aa) = \frac{g_{\rm DM}^4}{256\pi} \frac{h(t_0) - h(t_1)}{s(s - 4m_{\rm DM}^2)} \qquad h(t) \equiv 4\left(m_{\rm DM}^2 - t\right) + \frac{m_a^4(u - t)}{(m_{\rm DM}^2 - t)\left(m_{\rm DM}^2 - u\right)} \\ - \frac{2m_a^4 + \left(s - 2m_a^2\right)^2}{s - 2m_a^2}\log\left(-\frac{m_{\rm DM}^2 - t}{m_{\rm DM}^2 - u}\right) \\ u = 2m_{\rm DM}^2 + 2m_a^2 - s - t$$

$$\begin{split} \langle \sigma v \rangle(x) &= \sum_{q} \mathcal{A}_{q} + \frac{3}{2} \frac{\mathcal{B}}{x} + \mathcal{O}(x^{-2}) \\ s &\simeq m_{\rm DM}^{2} \left(4 - v^{2}\right) \end{split} \qquad \mathcal{A}_{q} = \frac{N_{\rm c}}{8\pi} \frac{g^{2} g_{f}^{2} g_{\rm DM}^{2} m_{\rm DM}^{2}}{(4m_{\rm DM}^{2} - m_{a}^{2})^{2}} \sqrt{1 - \frac{m_{q}^{2}}{m_{\rm DM}^{2}}} \\ \mathcal{B} &= \frac{g_{\rm DM}^{4}}{96\pi} \frac{m_{\rm DM}^{2} (m_{\rm DM}^{2} - m_{a}^{2})^{2}}{(2m_{\rm DM}^{2} - m_{a}^{2})^{4}} \sqrt{1 - \frac{m_{a}^{2}}{m_{\rm DM}^{2}}} \end{split}$$

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#### **Flavor constraints**

We study a Dirac Dark Matter particle interacting with ordinary matter via the exchange of a light pseudo-scalar, and analyze its impact on both direct and indirect detection experiments. We show that this candidate can accommodate the long-standing DAMA modulated signal and yet be compatible with all exclusion limits at  $99_S\%$  CL. This result holds for natural choices of the pseudo-scalar-quark couplings (e.g. flavor-universal), which give rise to a significant enhancement of the Dark Matter-proton coupling with respect to the coupling to neutrons. We also find that this candidate can accommodate the observed 1-3 GeV gamma-ray excess at the Galactic Center and at the same time have the correct relic density today. The model could be tested with measurements of rare meson decays, flavor changing processes, and searches for axion-like particles with mass in the MeV range.

