## Imaginary Charmonium Decay Widths ?

## A proposal for PANDA

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## Outline

- Vector Charmonium decay mechanisms
] J/ $\psi$ strong imaginary decay widths, experimental evidences:
- Vector+Pseudoscalar, Pseudoscalar+Pseudoscalar: $|\Phi|$ ~ $90^{\circ}$
- Energy scan, close to the J/ $\psi$ looking for interference, by BESIII: $|\Phi| \sim 90^{\circ}$
- A possible way to get the continuum phase
$\square$ Controversial evidences for $\psi^{\prime}(2 S)$
- $\psi^{\prime \prime}(3770)$ experimental evidences : $\Phi \sim-90^{\circ}$

A model for strong imaginary decay widths
[ A proposal for PANDA: a ppbar -> J/ $\psi->$ hadrons $/ \mu \mu$ scan

## Vector Quarkonium Decay Mechanisms


(a) $e^{+} e^{-} \rightarrow \mathrm{J} / \psi \rightarrow$ hadrons via strong mechanism (b) via em mechanism
(c) non-resonant $e^{+} e^{-} \rightarrow$ hadrons via a virtual photon. PQCD regime: all amplitudes real (apart BW resonance behaviour), while data are as if there is an additional in front of the BW

## Experimental Evidences for Imaginary Strong Decay Widths

Model dependent experimental evidences (old data)
SU3 and SU3 Breaking in $\mathbf{1 0}^{-}, \mathbf{0}^{-0} \mathbf{0}^{-} \mathbf{1 - 1}^{-1^{-}}$decay : $\Phi$ ~ $90^{\circ}$

$$
\begin{array}{ll}
\mathrm{J} / \Psi \rightarrow \mathrm{VP}\left(1^{-} 0^{-}\right) & \Phi=106^{\circ} \pm 10^{\circ}[1] \\
\mathrm{J} / \Psi \rightarrow \mathrm{PP}\left(0^{-} 0^{-}\right) & \Phi=89.6^{\circ} \pm 9.9^{\circ}[2] \\
\mathrm{J} / \Psi \rightarrow \mathrm{VV}\left(1^{-1}\right) & \Phi=138^{\circ} \pm 37^{\circ}[2]
\end{array}
$$

More recently:
If $\mathbf{A}\left(\mathbf{e}^{+} \mathbf{e}^{-} \rightarrow \mathbf{n n} \mathbf{n a r}\right) \sim-\mathbf{A}\left(\mathbf{e}^{+} \mathbf{e}^{-} \rightarrow \mathbf{p p}_{\text {bar }}\right)^{\text {[3] }}$
$B\left(\mathrm{nn}_{\mathrm{bar}}\right) / B\left(\mathrm{pp}_{\mathrm{bar}}\right)=0.98 \pm 0.08 \rightarrow \Phi \sim 89^{\circ} \pm 8^{\circ}{ }^{[4]}$ (BESIII)
[1] L. Kopke and N. Wermes, Phys. Rep. 174, 67 (1989); J. Jousset et al., Phys. Rev. D41, 1389 (1990).
[2] M. Suzuki et al., Phys. Rev. D60, 051501 (1999).
[3] FENICE Coll. NP B517(1998)3, SND Phipsi Rome, Sep (2013).
[4] M. Ablikim et al., Phys. Rev. D 86, 032014 (2012).

## VP decay updated and revisited

## SU3 and SU3 Breaking Amplitudes

Use reduced amplitudes $B=B_{0} / P^{*} 3$

| Process $J / \psi \rightarrow$ | Amplitude |
| :---: | :---: |
| $\rho^{+} \pi^{-}, \rho^{0} \pi^{0}, \rho^{-} \pi^{+}$ | $g+e$ |
| $K^{*+} K^{-}, K^{*-} K^{+}$ | $g(1-s)+e$ |
| $K^{* 0} \bar{K}^{0}, \bar{K}^{* 0} K^{0}$ | $g(1-s)-2 e$ |
| $\omega \eta$ | $(g+e) X_{\eta}+\sqrt{2} r g\left(\sqrt{2} X_{\eta}+Y_{\eta}\right)$ |
| $\omega \eta^{\prime}$ | $(g+e) X_{\eta^{\prime}}+\sqrt{2} r g\left(\sqrt{2} X_{\eta},+Y_{\eta^{\prime}}\right)$ |
| $\phi \eta$ | $(g(1-2 s)-2 e) Y_{\eta}+r g\left(\sqrt{2} X_{\eta}+Y_{\eta}\right)$ |
| $\phi \eta^{\prime}$ | $(g(1-2 s)-2 e) Y_{\eta^{\prime}}+r g\left(\sqrt{2} X_{\eta^{\prime}}+Y_{\eta^{\prime}}\right)$ |
| $\rho^{0} \eta^{\prime}$ | $3 e X_{\eta}$ |
| $\rho^{0} \eta^{\prime}$ | $3 e X_{\eta^{\prime}}$ |
| $\omega \pi^{0}$ | $3 e$ |
| $\phi \pi^{0}$ | 0 |

## J/ $\Psi$ <br> Vector +Pseudoscalar

| Parameter |  | Fit |
| :--- | :--- | :---: |
| $\mathrm{SU}_{3}$ strong Amplitude | g | $7.22 \pm 0.38$ |
| $\mathrm{SU}_{3}$ breaking strange | s | $0.18 \pm 0.04$ |
| $\mathrm{SU}_{3}$ breaking DOZI | r | $-0.04 \pm 0.02$ |
| E.M. Amplitude | e | $0.75 \pm 0.04$ |
| Phase | f | $81.51 \pm 6.75$ |

## J/ $\psi$

## Vector + Pseudoscalar

| Decay | Amplitude | PDG×104 | FitX104 | $\Delta \chi^{2}$ |
| :---: | :---: | :---: | :---: | :---: |
| $\rho^{0} \pi^{0}$ | $g \mathrm{e}^{\mathrm{if}}+\mathrm{e}$ | $169.0 \pm 15.0$ | 133.00 | 1.13 |
| $\mathrm{K}^{*}+\mathrm{K}^{-}$ | $\mathrm{g}(1-\mathrm{s}) \mathrm{e}^{\mathrm{id}+e}$ | $51.2 \pm 3.0$ | 51.5 | 0.01 |
| $\mathrm{K}^{*} \mathrm{~K}^{0}$ | $\mathrm{g}(1-\mathrm{s}) \mathrm{e}^{\text {ip }}-2 \mathrm{e}$ | $43.9 \pm 3.1$ | 48.5 | 0.48 |
| $\omega \eta$ | $(g X+d) e^{i \phi}+e X$ | $17.4 \pm 2.0$ | 18.5 | 0.06 |
| $\phi \eta$ | $(\mathrm{g}(1-2 \mathrm{~s}) \mathrm{Y}+\mathrm{d}) \mathrm{e}^{\mathrm{if}-2 \mathrm{e}} \mathrm{Y}$ | $7.5 \pm 0.8$ | 3.9 | 4.02 |
| $\rho \eta$ | 3 eX | $1.9 \pm 0.2$ | 2.2 | 0.30 |
| $\omega \pi$ | 3 e | $4.5 \pm 0.5$ | 4.1 | 0.11 |
| $\omega \eta^{\prime}$ | $\left(\mathrm{g} \mathrm{X}{ }^{\prime}+\mathrm{d}^{\prime}\right) \mathrm{e}^{\mathrm{id}}+\mathrm{e} \mathrm{X}^{\prime}$ | $7.0 \pm 7.0$ | 11.9 | 0.10 |
| $\phi \eta^{\prime}$ | $\left(g(1-2 s) Y^{\prime}+d^{\prime}\right) e^{i \varphi}-2 e Y^{\prime}$ | $4.0 \pm 0.7$ | 6.1 | 1.87 |
| $\rho \mathrm{H}$ | 3 eX | $1.1 \pm 0.2$ | 1.1 | 0.04 |

PP decay updated and revisited

## Pseudoscalar Pseudoscalar Decay Revisited

- Open question about $J / \Psi->\pi \pi$ decay, since pure em :
$\mathrm{B}^{\pi \pi}=\left|\mathrm{E}^{\pi \pi}\right|^{2,}$, while
$\mathrm{B}^{\pi \pi}=(1.47 \pm .23) 10^{-4}$ from PDG
$\left|\mathrm{E}^{\pi \pi}\right|^{2}=\mathrm{B}^{\mu \mu} \sigma\left(\mathrm{e}^{+} \mathrm{e}^{-}->\pi^{+} \pi-\right) / \sigma\left(\mathrm{e}^{+} \mathrm{e}^{-}->\mu \mu\right)=$ $=(0.46 \pm .23) 10^{-4}$ extrapolated from BaBar $\mathbf{B}^{\pi \pi} \neq\left|\mathrm{E}^{\pi \pi}\right|^{2}$ by 3 s.d.
$\square \pi \pi$ cross section slope $B$, asymptotically it is expected $B=-2-4 \times n_{q}=-6$ $\mathrm{B}^{\pi \pi} \sim-10 \pm 2$


## Pseudoscalar Pseudoscalar Decay Revisited

- It is possible to avoid $\pi \pi$ and complications from s quark by means of KK BR's and $\left|\mathrm{E}^{K K}\right|$ only
- $\mathrm{B}^{+-}=|\mathrm{S}|^{2}+\left|\mathrm{E}^{+-}\right|^{2}+2|\mathrm{~S}|\left|\mathrm{E}^{+-}\right| \cos \Phi$ $B^{S L}=|S|^{2}+\left|E^{S L}\right|^{2}-2|S|\left|E^{S L}\right| \cos \Phi$
- $\left|\mathrm{E}^{+-}\right|^{2}=\mathrm{B}^{\mu \mu} \sigma\left(\mathrm{e}^{+} \mathrm{e}^{-}->\mathrm{K}^{+} \mathrm{K}^{-}\right) / \sigma\left(\mathrm{e}^{+} \mathrm{e}^{-}->\mu \mu\right)$ $\left|E^{S L}\right|^{2} \sim 0$, since $\sigma\left(e \mathrm{e}->\mathrm{K}_{\mathrm{S}} \mathrm{K}_{\mathrm{L}}\right) \ll \sigma\left(\mathrm{e} e->\mathrm{K}^{+} \mathrm{K}^{-}\right)$ $\sigma\left(\mathrm{e}^{+} \mathrm{e}^{-}->\mathrm{K}_{\mathrm{S}} \mathrm{K}_{\mathrm{L}}\right) \sim 0.6 \mathrm{pb}$ at J/ $\Psi$ $\mathrm{B}^{+-} \quad=(2.37 \pm 0.31) 10^{-4} \quad \mathrm{~B}^{\mathrm{SL}}=(1.66 \pm 0.26) 10^{-4}$ $\left|\mathrm{E}^{+-}\right|^{2}=(1.3 \pm 0.6) 10^{-4} \quad$ from BaBar $\Phi=83.7^{0} \pm 9.0^{\circ}$

The $\psi^{\prime}$ Puzzle

## $\Psi^{\prime}$

## Vector + Pseudoscalar

| Parameter | Fit |  |
| :---: | :---: | :---: |
| $\mathrm{SU}_{3}$ strong Amplitude | g | $0.49 \pm 0.04$ |
| $\mathrm{SU}_{3}$ breaking strange | s | $-0.04 \pm 0.13$ |
| $\mathrm{SU}_{3}$ breaking DOZI | r | $-0.04 \pm 0.08$ |
| E.M. Amplitude | e | $0.18 \pm 0.02$ |
| Phase | f | $159 . \pm 12$. |
| $\chi^{2} / \mathrm{DFR}=0.96$ |  |  |

## $\Phi$ at the $\Psi^{\prime}$ from $K^{*}(892) \mathrm{K}$ Decay Only

- $\mathrm{K}^{*}(892) \mathrm{K}$ decay: possible to avoid $\mathrm{SU}_{3}$ assumptions and complications from s quark mass, since CLEOc measured the continuum cross sections
- CLEOc (arXiv:hep-ex/0509011v2):
$\sigma\left(\mathrm{e} \mathrm{e}->\mathrm{K}^{*} \mathrm{~K}^{0}+\mathrm{cc}\right)=(23.5 \pm 5.3) \mathrm{pb}$ at $\mathrm{W}=3.67 \mathrm{GeV}$ $\sigma\left(\mathrm{e} \mathrm{e}->\mathrm{K}^{*+} \mathrm{K}^{-}+\mathrm{cc}\right) \sim(1 \pm 0.9) \mathrm{pb}$ $\left|E^{+-}\right|^{2} \sim 0.1 \times 10^{-5} \quad\left|E^{00}\right|^{2} \sim 28 . \times 10^{-5}$
$\square B^{+-}=(1.7 \pm 0.8) \times 10^{-5} \quad B^{00}=(10.9 \pm 2.0) \times 10^{-5}$
$B^{+-}=|S|^{2}+\left|E^{+-}\right|^{2}+2 x|S|\left|E^{+-}\right| \cos \Phi$
$B^{00}=|S|^{2}+\left|E^{00}\right|^{2}-2 x|S|\left|E^{00}\right| \cos \Phi$
$\Phi=159^{\circ} \pm 24^{0}$ again like VP!


## Pseudoscalar Pseudoscalar Decay

$\square \Psi^{\prime}:$

$$
\begin{aligned}
& \mathrm{B}^{+-}=(6.30 \pm 0.70) 10^{-5} \quad \mathrm{~B}^{S L}=(5.26 \pm 0.25) 10^{-5} \\
& \left|\mathrm{E}^{+-}\right|^{2}=(0.7 \pm 0.4) 10^{-5} \quad \text { from BaBar } \\
& \Phi \quad=95^{0} \pm \mathbf{1 1 0}^{0} \quad(6.3 \sim 5.26+0.7+3.8 \times \cos \Phi)
\end{aligned}
$$

- But Nambu wrote $\Psi^{\prime}$ might be different! (PRL 34(1975), 1645)


## Experimental evidences for <br> $\Psi(3770)$ imaginary strong decay widths

$\boldsymbol{\Psi}^{\prime \prime}(3770):$

* non DDbar (small) -> throught the interfence with continuum
* For a wide resonance $\Phi$ from interference at the peak
$-2\left|A_{3 g}\right| / \Gamma_{\text {tot }} \sin \Phi \times$ continuum
* CLEOc and BESIII: $\Phi \sim-90^{\circ}$, since continuum sign

| decay | continuum | $\Psi^{\prime \prime}(\mathbf{3 7 7 0})$ | sign |  |
| :---: | :---: | :---: | :---: | :--- |
| $\rho \pi$ | $13.1 \pm 2.8$ | $7.4 \pm 1.3$ | - | CLEOc, PRD 73(2006)012002 |
| $\phi \eta$ | $2.1 \pm 1.6$ | $4.5 \pm 0.7$ | + | CLEOc, PRD 73(2006)012002 |
| $P \mathrm{p}$ | $0.74 \pm 0.08$ | $0.4 \pm 0.02$ | - | BESIII Y.Liang, Nov (2012) |

## Model independent from interference in $q^{2}$ behavior



Actually $\Phi_{\text {meas }}=\Phi-\delta_{\text {cont }}$ and $\left|\Phi_{\text {meas }}\right|$ only is measured, since it is difficult to get the sign

The full interference between $\mathrm{A}_{\mathrm{EM}}$ and $\mathrm{A}_{\text {cont }}$ has been observed, as expected, at MARKI(1975), BESII (1995), KDER (2010).
$1 / 2$ photon propagators require $\varphi^{\prime}=180^{\circ}$


## BESIII J/ $\Psi$ scan





## The interference pattern is not always the same.



A possible way to get the continuum phase (work in progress)

## Continuum phase $\mathrm{d}(\mathrm{s})$

- Continuum amplitudes should be almost real : $\delta(\mathrm{s}) \sim 0^{\circ}$ or $180^{\circ}$
- Logarithm Dispersion Relations relating modulus $|\mathrm{F}(\mathrm{s})|^{2} \sim \sigma(\mathrm{~s})$ and $\delta(\mathrm{s})$ might help:

$$
\delta(s)=-\frac{\sqrt{s-q_{\mathrm{t}}^{2}}}{\pi} \mathrm{PV} \int_{q_{\mathrm{i}}^{2}}^{\infty} \frac{\ln |F(t)| F(0) \mid}{(t-s) \sqrt{t-q_{\mathrm{t}}^{2}}} d t, \quad \delta(s)=-\frac{\sqrt{s-q_{\mathrm{t}}^{2}}}{\pi} \mathrm{PV} \int_{q_{\mathrm{t}}^{2}}^{\infty} \frac{\ln |F(t)| F(0) \mid}{(t-s) \sqrt{t-q_{\mathrm{i}}^{2}}} d t .
$$

Check: phase as expected, if $|\mathrm{F}(\mathrm{s})|^{2} \sim \mathrm{BW} \sim \sigma(\mathrm{s}) / \mathrm{IPS}$
$\square$ Applied to $\sigma\left(\mathrm{e}^{+} \mathrm{e}^{-}->\mathrm{pp}_{\mathrm{bar}}\right)$ (unphysical region): $\delta(\mathrm{s}) \sim 360^{\circ}$

$$
\begin{aligned}
& \sigma\left(\mathrm{e}^{+} \mathrm{e}^{-}->\pi \pi\right): \delta(\mathrm{s}) \sim 180^{\circ} \\
& \sigma\left(\mathrm{e}^{+} \mathrm{e}^{-}->3 \pi\right): \delta(\mathrm{s}) \sim 180^{\circ}(?)
\end{aligned}
$$

If $\delta(\mathrm{s}) \neq 0$ and it is known how $180^{\circ}$ or $0^{0}$ is asymp reached, from $\left|\Phi_{\text {meas }}\right|=|\phi-\delta|$ the sign $\left(+/-90^{\circ}\right)$ might be established

## $\pi^{+} \pi$ and $\mathrm{pp}_{\mathrm{bar}}$ (throught the unphysical region.) phases

(S.Pacetti, R. Baldini... EPJC 11(1999)709... long time ago)


## Open Issues related to Unitarity

$\square \quad$ No explanation for imaginary strong decay J/ $\Psi$ widths has been put forward until now
$\square \mathrm{J} / \Psi$ description as a Breit Wigner might have some difficulties, dealing with imaginary decay widths
$\square$ Optical theorem: $\operatorname{Im} \mathrm{T}_{\mathrm{el}}=\mathrm{W} / 8 \pi \cdot \sigma_{\text {tot }}$ implies $\operatorname{Im} \mathrm{T}_{\mathrm{el}}>0$

- $\Gamma\left(\mathrm{J} / \Psi->\mathrm{pp}_{\text {bar }}\right)$ imaginary: $\operatorname{Im} \mathrm{T}_{\mathrm{el}}\left(\mathrm{pp}_{\text {bar }}->\mathrm{J} / \Psi->\mathrm{pp}_{\text {bar }}\right)<0$
. $\mathrm{pp}_{\text {bar }}$ continuum could restore unitarity, even if unrelated to $\mathrm{J} / \Psi$
$\square$ Looking for a different J/ $\Psi$ description
$\square \quad \sigma_{\text {el }}\left(\mathrm{pp}_{\text {bar }}->\mathrm{J} / \Psi->\right.$ hadrons) : a test of the following model


## A model to explain imaginary widths

## Quarkonium OZI breaking decay

as Freund and Nambu (PRL 34(1975), 1645)
$\square \quad$ Quarkonium as a superposition of

- A narrow V (coupled to the virtual photon, but not directly to hadrons)
- A wide one (a glueball O)
(not coupled to leptons i.e. to a virtual photon,
but strongly coupled to hadrons)
$f$ is the coupling between $v$ and $\mathcal{O}$

iterated in $f$


## Quarkonium OZI breaking decay

## as Freund and Nambu (PRL 34(1975), 1645)

$\square$ Quarkonium as a superposition of V and O :

$$
\begin{aligned}
A_{\text {strong }} & =G_{e} V^{-1} f O^{-1} G_{f}+G_{e} V^{-1} f O^{-1} f V^{-1} f O^{-1} G_{f}+\text { iterations } \\
& =G_{e} V^{-1} f O^{-1} G_{f} /\left(1-V^{-1} O^{-1} f^{2}\right)=G_{e} f G_{f} /\left(V O-f^{2}\right)
\end{aligned}
$$

$\square \quad A_{e m} \quad=G_{e} V^{-1} G_{1}+G_{e} V^{-1} \mathrm{f} \mathrm{O}^{-1} \mathrm{f} \mathrm{V}^{-1} \mathrm{G}_{1}+$ iterations

$$
=G_{e} O G_{f} /\left(V O-f^{2}\right)
$$

$\square \quad$ An infinity of radial O recurrences
$\square \quad$ This model mainly used to try to explain $\operatorname{Br}\left(\psi^{\prime}\right) / \operatorname{Br}(\mathrm{J} / \psi)$ anomalies S. J. Brodsky, G. P. Lepage, S. F. Tuan, PRL 59, 621(1987) W.S. Hou, C.Y. Ko, NTUTH-97-11, 1997

# Narrow V and wide glueball 0 superposition 

## P.J.Franzini, F.J.Gilman, PR D32, 237 (1985)

$$
A_{\text {strong }}=\frac{\sqrt{\Gamma_{e e}} M_{V} M_{O} f \sqrt{\Gamma_{O}}}{\left(M_{V}^{2}-W^{2}-i M_{V} \Gamma_{V}\right)\left(M_{O}^{2}-W^{2}-i M_{O} \Gamma_{O}\right)-M_{V} M_{O} f^{2}}
$$

assuming $\Gamma_{O} \gg \Gamma_{J / \psi}, f^{2} \sim \Gamma_{0}\left(\Gamma_{J / \psi}-\Gamma_{V}\right)$

$$
A_{\text {strong }} \sim \frac{(i) \sqrt{B_{e e}} M_{V} f \sqrt{B_{h}}}{M_{J / \Psi}^{2}-W^{2}-i M_{J / \Psi} \Gamma_{J / \Psi}} A_{e m}=\frac{\sqrt{B_{e e}} M_{V} \Gamma_{J / \Psi} \sqrt{B_{e m}}}{M_{J / \Psi}^{2}-W^{2}-i M_{J / \Psi} \Gamma_{J / \Psi}}
$$

■ The additional $90^{\circ}$ phase is naturally achieved
$\square \mathrm{J} / \psi$ shape reproduced if: $|f| \sim 0.012 \mathrm{GeV}, \mathrm{M}_{\mathrm{O}} \sim \mathrm{M}_{\mathrm{J} / \psi}, \Gamma_{\mathrm{O}} \sim 0.5 \mathrm{GeV}$
nly far from the $\mathrm{J} / \psi$ (no contradiction with BES, PR 54(1996)1221)
$\square \psi$ "(3770) decay phases agree with Nambu suggestion.
$\square \psi^{\text {‘ }}$ unclear; $\psi^{\text {‘ }}$-> $\mathrm{J} / \psi \pi \pi$ (?)

## SND $\Phi->\pi^{+} \pi^{-} \pi^{0}$



## BaBar $\pi^{+} \pi^{-} \pi^{0} \quad$ PRD $70,072004(2004)$



Masses and widths

$$
\begin{gathered}
M_{\omega^{\prime}}=(1350 \pm \mathbf{2 0} \pm \mathbf{2 0}) \mathrm{MeV} / c^{2} \\
\Gamma_{\omega^{\prime}}=(450 \pm 70 \pm \mathbf{0 0}) \mathrm{MeV} / c^{2} \\
M_{\omega^{\prime \prime}}=(1660 \pm 10 \pm \mathbf{2}) \mathrm{MeV} / c^{2} \\
\Gamma_{\omega^{\prime \prime}}=(230 \pm 30 \pm \mathbf{2 0}) \mathrm{MeV} / c^{2}
\end{gathered}
$$

BaBar found indeed an unexpected resonance ( O ?)
at 1.35 GeV , wide 0.45 GeV


## A proposal for PANDA: a J/ $\Psi$ scan

## A Proposal for PANDA

$\square \quad$ Expected $\sigma\left(\mathrm{p}_{\mathrm{b} \text { bar }}->\mathrm{J} / \Psi \rightarrow\right.$ hadrons $) \sim 1 \mu \mathrm{~b}$
while $\quad \sigma\left(\mathrm{p}_{\mathrm{bar}}->\right.$ hadrons $) \sim 70 \mathrm{mb}$

- No J/ $\Psi$ exclusive production evidence in present data (too small cross section $+\mathrm{p} \mathrm{p}_{\mathrm{bar}} \mathrm{c}$. m . energy spread)
- Different mechanism in inclusive or exclusive production:
> Inclusive production: direct coupling to gluons or virtual photon
> Exclusive production: hadronic -> apply FN model


## p par Total and Elastic cross section (PDG2012)



## A Proposal for PANDA

Contributions to $p \mathrm{p}_{\mathrm{bar}}->\mathrm{J} / \Psi->$ hadrons, according to the FN model


## A Proposal for PANDA

$\square \quad A=G_{p} O^{-1} G_{h}+G_{p} O^{-1} \mathrm{f} \mathrm{V}^{-1} \mathrm{f} \mathrm{O}^{-1} G_{h}+$ iterations
$A=G_{p} O^{-1} G_{h} /\left(1-V^{-1} O^{-1} f^{2}\right)$
$A=G_{p} G_{h} V /\left(V O-f^{2}\right)$
$\square \quad$ Still assuming
$>\Delta \mathrm{W} \sim \Gamma_{\mathrm{J} / \Psi}->\left(\mathrm{M}_{\mathrm{O}}{ }^{2}-\mathrm{W}^{2}\right) / \mathrm{M}_{\mathrm{O}} \ll \Gamma_{\mathrm{O}}$
$>\mathrm{f}^{2} \sim \Gamma_{\mathrm{O}}\left(\Gamma_{\mathrm{J} / \Psi}-\Gamma_{\mathrm{V}}\right)$
> Amplitudes $p p_{\text {bar }}->V, V->p p_{\text {bar }}$ negligeable
> Interference with background $\mathrm{J}^{\mathrm{P}}=1^{\text {- }}$ to be included yet

## A Proposal for PANDA

$\square$ According to the FN approach

$$
\sigma_{F N}=\frac{\left.B_{p}\left[\left(M_{J / \Psi}^{2}-W^{2}\right)^{2}+\left(M_{J / \Psi} \Gamma_{V}\right)^{2}\right)\right] B_{h}}{\left(M_{J / \Psi}^{2}-W^{2}\right)^{2}+\left(M_{J / \Psi} \Gamma_{J / \Psi}\right)^{2}}
$$

Taking into account that $\Gamma_{K} \ll \Gamma_{\vartheta / \Psi}$
$\sigma_{F N}=\frac{B_{p}\left(M_{J / \Psi}^{2}-W^{2}\right)^{2} B_{h}}{\left(M_{J / \Psi}^{2}-W^{2}\right)^{2}+\left(M_{J / \Psi} \Gamma_{J / \Psi}\right)^{2}} \quad$ a zero $->$ a dip in $\sigma_{\mathbf{h}}$
$\square$ To be compared to a Breit Wigner

$$
\sigma_{B W}=\frac{B_{p} \Gamma_{J / \Psi}^{2} B_{h} \mathrm{M}^{2}{ }_{\mathrm{J} / \Psi}}{\left(M_{J / \Psi}^{2}-W^{2}\right)^{2}+\left(M_{J / \Psi} \Gamma_{J / \Psi}\right)^{2}}
$$

## A Proposal for PANDA

PANDA inv mass resolution: small beam energy spread and no ISR


## A Proposal for PANDA

$\square$ Rough $J^{P}=1^{-}$estimation $p_{\text {bar }}$ background $\sigma$ at $P_{\text {pbar }} \sim 4 \mathrm{GeV}$ :
$>\sigma\left(\mathrm{J}^{\mathrm{P}}=1^{-}\right) \sim 0.5 \sigma(\mathrm{~S}$ wave $)$
$>\sigma_{\text {tot }} \sim$ Black Disk $=2 \pi \mathrm{R}^{2}=2 \pi / \mathrm{P}^{2} \Sigma_{\mathrm{l}}(2 \mathrm{l}+1)$
$>I_{\text {max }} \sim R P \sim 25$
$>\mathrm{S}$ wave $\sim 0.5 \sigma_{\text {tot }} / \mathrm{I}_{\max }{ }^{2} \sim 40 \mu \mathrm{~b} \quad\left(\sigma_{J / \Psi} \sim 1.5 \mu \mathrm{~b}\right)$
> Background amplitude R+iI, should be mostly imaginary : $\mathrm{I} \sim 5 \times \mathrm{A}_{\mathrm{J} / \Psi}, \quad \mathrm{I} \gg \mathrm{R}$
. $\mathrm{JP}^{\mathrm{P}}=1^{-}$background heavily interferences with the $\mathrm{J} / \Psi$
Some channel might have a much better J/ $\Psi /$ background ratio: $3 \pi, 5 \pi, \ldots$ ?

## A Proposal for PANDA

$\square \quad \mathrm{JP}=1^{-}$background interference with FN :

$$
\propto \frac{\left[\left(\mathrm{M}^{2}-\mathrm{W}^{2}\right)^{2}+\Gamma_{/ / \Psi} \Gamma_{\mathrm{V}} \mathrm{M}^{2}\right] \mathrm{I}-\left(\Gamma_{\mathrm{J} / \Psi}-\Gamma_{\mathrm{V}}\right) \mathrm{M}\left(\mathrm{M}^{2}-\mathrm{W}^{2}\right) \mathrm{R}}{\left(\mathrm{M}^{2}-\mathrm{W}^{2}\right)^{2}+\Gamma_{\mathrm{J} / \Psi} \mathrm{M}}
$$

$\square \quad$ The term prop. to I should increase the expected dip, since I>0
$\square \quad$ The term prop. to $R$ expected small and affected by beam spread
$\square \quad \mathrm{J}^{\mathrm{P}}=1^{-}$background interference with a BW:

$$
\propto \frac{\left.\left(\Gamma_{\mathrm{I} / \Psi} \mathrm{M}\right)^{2}\right] I+\mathrm{M}^{2}\left(\mathrm{M}^{2}-\mathrm{W}^{2}\right) \mathrm{R}}{\left(\mathrm{M}^{2}-\mathrm{W}^{2}\right)^{2}+\Gamma_{\mathrm{J} / \Psi} \mathrm{M}}
$$

$\square \quad$ The term prop. to I should increase the expected peak
$\square \quad$ The term prop. to $R$ has to be evaluated

## A Proposal for PANDA

$\square$ Rough estimation of the integrated luminosity:
$>$ Signal $\sim 0.2 \div 0.4 \mu \mathrm{~b}$, depending on $\sigma_{\text {beam }} \sim 200 \div 100 \mathrm{KeV}$
$>$ Background $\sim 5 \cdot 10^{4} \mu \mathrm{~b}$
$>\mathrm{L} \sim 10^{31} \mathrm{~cm}^{-2} \mathrm{sec}^{-1}$
$\square \quad \mathrm{S} \sim \mathrm{n} \cdot \sqrt{ } \mathrm{B}$, after $\mathrm{T}->\mathrm{O} .2 \cdot \mathrm{~T} \cdot \mathrm{~L} \sim \mathrm{n} \sqrt{ }\left(5 \cdot 10^{4} \cdot \mathrm{~T} \cdot \mathrm{~L}\right)$
$\mathrm{T} \sim$ few months, if $\mathrm{n} \sim 4$, assuming a 10 points scan
(efficiency and dead time to be included)
. Much less time might be needed for some channels: $3 \pi, 5 \pi$, ..

- Of course time is available for any other measurement at J/ $\Psi$


## A Proposal for PANDA

Marco Destefanis already proposed to look for $\mathrm{J} / \Psi->\mu \mu$ in PANDA, exploiting the very good inv mass resolution (no ISR)

## Exploiting Di-Muon <br> Production at PANDA



## $\mathrm{J} / \Psi$ invariant mass resolution in $\mathrm{e}+\mathrm{e}-->\mathrm{p} \mathrm{p}_{\mathrm{bar}}$ in BESIII

(Marco Destefanis at STORI11)


## $\mathrm{J} / \Psi$ invariant mass resolution in $\mathrm{p} \mathrm{p}_{\mathrm{bar}}->\mu \mu$ in PANDA

## (Marco Destefanis at STORI11)

## Simulated Yields for $\overline{\mathrm{p} p}->\mu^{+} \mu^{-}$

- $\Delta \varphi=0^{\circ}$
- $\Delta \varphi=90^{\circ}$
- $\Delta \varphi=180^{\circ}$
continuum reference
$\sigma \sim 18 \mathrm{pb}$




## Conclusions

$\square \quad$ Unexpected imaginary $\mathrm{J} / \Psi$ strong decay widths ( $\Phi \sim\left|90^{\circ}\right|$ )
$\square$ Updated VP and PP J/ $\Psi$ decays data point out this result
$\square \mathrm{J} / \Psi$ scan by BESIII seems to confirm that $\Phi \sim\left|90^{\circ}\right|$
$\square \Psi(2 S)$ present data contradictory-> $\Psi(2 S)$ scan by BESIII
$\square \quad \Psi^{\prime \prime}(3770)$ present data suggest $\Phi \sim-90^{\circ}$
$\square$ A model under development to explain this unexpected phase

## Conclusions

- A proposal for PANDA:
> $\mathrm{p}_{\mathrm{bar}}->\mathrm{J} / \Psi->$ hadrons seen as a dip
$>\mathrm{pp}_{\mathrm{bar}}->\mathrm{J} / \Psi->\mu \mu$, ee seen as a peak
(exploiting PANDA very good inv. mass resolution)
- However a better evaluation of the interference
with the $\mathrm{JP}^{\mathrm{P}}=1^{-}$background is needed


# Thanks for your attention 

(谢谢)

## BaBar: $\mathrm{e}^{+} \mathrm{e}^{-}->\pi^{+} \pi^{-}$cross section

arXiv:1205.2228v1



$$
\mathrm{e}^{+} \mathrm{e}^{-->} \mathrm{KS} \text { KL }
$$



## Summary of fit results

| Channel | $M_{\text {J/ }}{ }^{\text {r }}$ | $\Gamma$ (KeV) | $\varphi^{\prime}$ |
| :---: | :---: | :---: | :---: |
| $\mu+\mu-$ | $3097.33 \pm 0.01$ | 92.9 (fixed) | $0^{\circ}$ (fixed) |
| $\pi+\pi-\pi+\pi-$ | $3097.46 \pm 0.03$ | 92.9 (fixed) | $(-2.14 \pm 27.59)^{\circ}$ |
| $\pi+\pi-\pi+\pi-\pi 0$ | $3097.50 \pm 0.04$ | 92.9 (fixed) | $0^{\circ}$ (fixed) |
| $\pi+\pi-\pi 0$ | $3097.50 \pm 0.06$ | 92.9 (fixed) | $0^{\circ}$ (fixed) |
| pp | $0.3+3096.9$ | - | -- |
| Channel | Ф | Br ${ }_{\text {out }}$ | Bripdg |
| $\mu+\mu-$ | -- | $5.94 \times 10^{-2}$ (fixed) | $5.94 \times 10^{-2}$ |
| $\pi+\pi-\pi+\pi-$ | -- | $(3.04 \pm 0.17) \times 10^{-3}$ | $(3.55 \pm 0.23) \times 10^{-3}$ |
| $\pi+\pi-\pi+\pi-\pi 0$ | $(102.6 \pm 5.1)^{\circ}$ | $(3.55 \pm 0.13) \times 10^{-2}$ | $(4.1 \pm 0.5) \times 10^{-2}$ |
| $\pi+\pi-\pi 0$ | $(108.4 \pm 10.1)^{0}$ | $(1.87 \pm 0.08) \times 10^{-2}$ | $(2.07 \pm 0.12) \times 10^{-2}$ |
| pp | $(84.73 \pm 9.62)^{\circ}$ | $(1.90 \pm 0.05) \times 10^{-3}$ | -- |
| Channel | $\sigma_{\text {cont }}(\mathrm{nb})$ | $\mathrm{S}_{\mathrm{E}}(\mathrm{MeV})$ |  |
| $\mu+\mu-$ | -- | $0.92 \pm 0.01$ |  |
| $\pi+\pi-\pi+\pi-$ | $0.465 \pm 0.014$ | 0.92 (fixed) |  |
| $\pi+\pi-\pi+\pi-\pi 0$ | $0.153 \pm 0.013$ | 0.92 (fixed) |  |
| $\pi+\pi-\pi 0$ | $0.040 \pm 0.010$ | 0.92 (fixed) |  |
| pp | $0.006 \pm 0.001$ | $0.92 \pm 0.01$ |  |

