

Structure formation & Clusters for Cosmology

Alain Blanchard

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Transfer function I

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Notation

$$\delta(k, z) = D(z) \delta_0(k) \text{ with } D(0) = 1.$$

and:

$$P(k) = |\delta|_k^2 = D(z)^2 P_0(k)$$

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The power spectrum keeps record of the transfer function.

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$$\delta \propto t \propto a(t)^2$$

as long as $\lambda \geq ct$ while $\lambda \leq ct$ are frozen.

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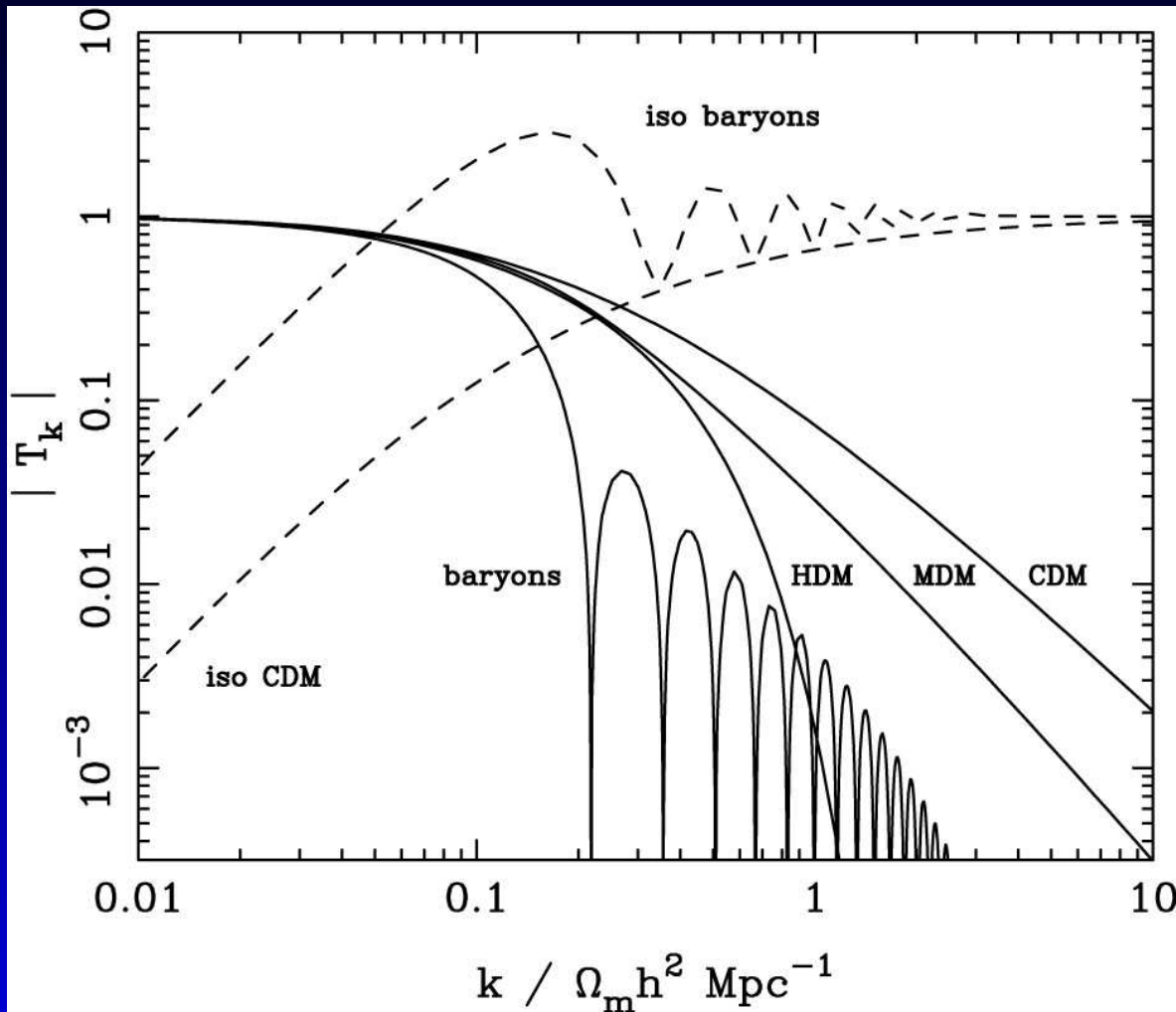
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Low mass particles free-streaming.

Transfer function: IV



Comoving sound horizon size I

RW metric:

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For a sound wave:

$$\chi_S(a) = \int_0^a a^{-1}(t) c_S dt = \int_0^a \frac{c}{\sqrt{3}} \frac{da}{a(t)^2 H(a) \sqrt{1 + \frac{3}{4} \frac{\Omega_b}{\Omega_r}}}$$

Comoving sound horizon size II

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- z_d imprint in $P(k)$ (BAO)

Non linear regime

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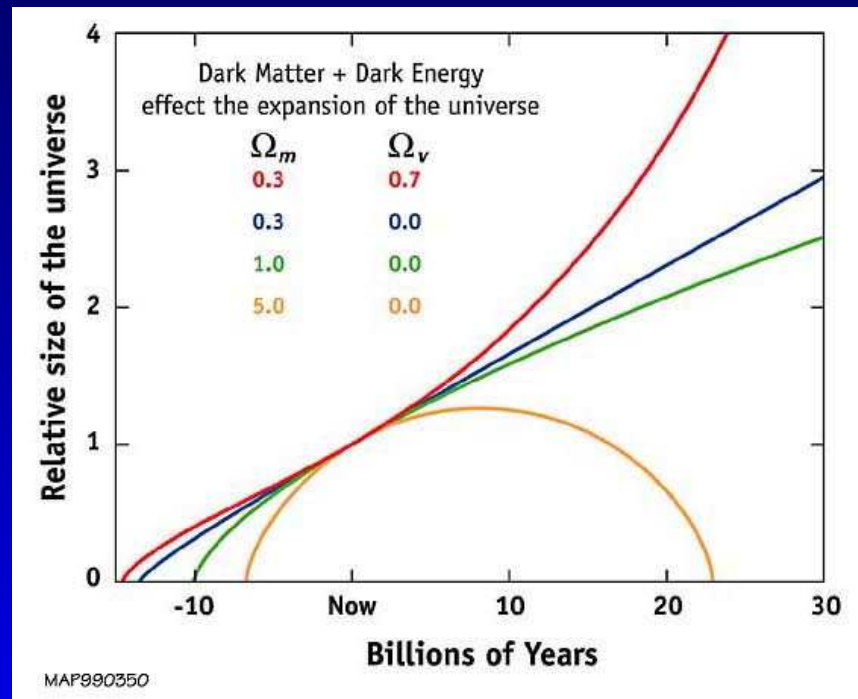
Newtonian problem.

Non linear regime

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Spherical model (Lemaître, 1933)

Newtonian problem. Solution already seen:



Spherical Perturbation I

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$$\tilde{H}_0 t = \frac{\tilde{\Omega}_0}{2(\tilde{\Omega}_0 - 1)^{3/2}} (\phi - \sin(\phi))$$

$$R(t) = \frac{\tilde{\Omega}_0 \tilde{R}_0}{2(\tilde{\Omega}_0 - 1)} (1 - \cos(\phi))$$

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Density at maximum:

$$\tilde{\rho} = \tilde{\rho}_0 \left(\frac{\tilde{R}_0}{\tilde{R}} \right)^3$$

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$$\tilde{H}_0 t_m = \frac{\tilde{\Omega}_0}{2(\tilde{\Omega}_0 - 1)^{3/2}} \pi$$

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with : $1 + \Delta_m = \frac{\tilde{\rho}_m}{\rho}$ and $\rho = \frac{1}{6\pi G t^2}$ (EdS)

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$$\Delta_m = \frac{9}{16} \pi^2 - 1. \simeq 4.55$$

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At $2t_m$ solution reaches a singularity.

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let's estimate the linear expected amplitude at virilization.

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$$\tilde{\rho} = \frac{8\rho_m}{(1 - \cos \psi)^3} = \frac{64\rho_m}{\psi^6(1 - \psi^2/4)}$$

$$t = \frac{t_m}{\pi}(\psi - \sin \psi) = \frac{t_m}{\pi} \frac{\psi^3}{6} \left[1 - \frac{\psi^2}{20}\right]$$

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so:

$$\psi^6 = \left(\frac{6\pi t}{t_m}\right)^2 \left[1 + \frac{\psi^2}{10}\right]$$

Virialization IV

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and

$$\begin{aligned}\tilde{\rho} &= \left(1 + \frac{\psi^2}{4} - \frac{\psi^2}{10}\right) \frac{64\rho_m t_m^2}{(6\pi)^2 t^2} \\ &= \left(1 + \frac{\psi^2}{4} - \frac{\psi^2}{10}\right) \frac{64}{36\pi^2} \frac{3\pi^2}{32\pi G t^2} = \rho \left(1 + \frac{3\psi^2}{20}\right)\end{aligned}$$

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so with : $\tilde{\rho} = \rho(1 + \delta)$

$$\delta = \frac{3}{20}\psi^2 = \frac{3}{20} \left(\frac{6\pi t}{t_m}\right)^{2/3} = \frac{3(6\pi)^{2/3}}{20} \frac{1 + z_m}{1 + z}$$

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$$\delta_m = \frac{3(6\pi)^{2/3}}{20}(1+z_m) = 1.06(1+z_m) \text{ when } \Delta_m \simeq 4.5$$

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$$\text{For } z < z_v, \Delta = 177 \left(\frac{1+z_v}{1+z} \right)^3$$

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The mass is :

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so that M and z are the only two numbers to
characterize a cluster (Δ is set by the cosmology...or
by the cosmologist!)

Velocity dispersion II

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$$\rho(r) = \frac{\sigma^2}{2\pi Gr}$$

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$$\sigma = \sqrt{\frac{GM}{r}}$$

Velocity dispersion III

Numerically:

$$\sigma = 1130(hM_{15})^{1/3} \left(\frac{\Delta\Omega_m}{178} \right)^{1/6} \sqrt{1+z} \text{ km/s}$$

Velocity dispersion III

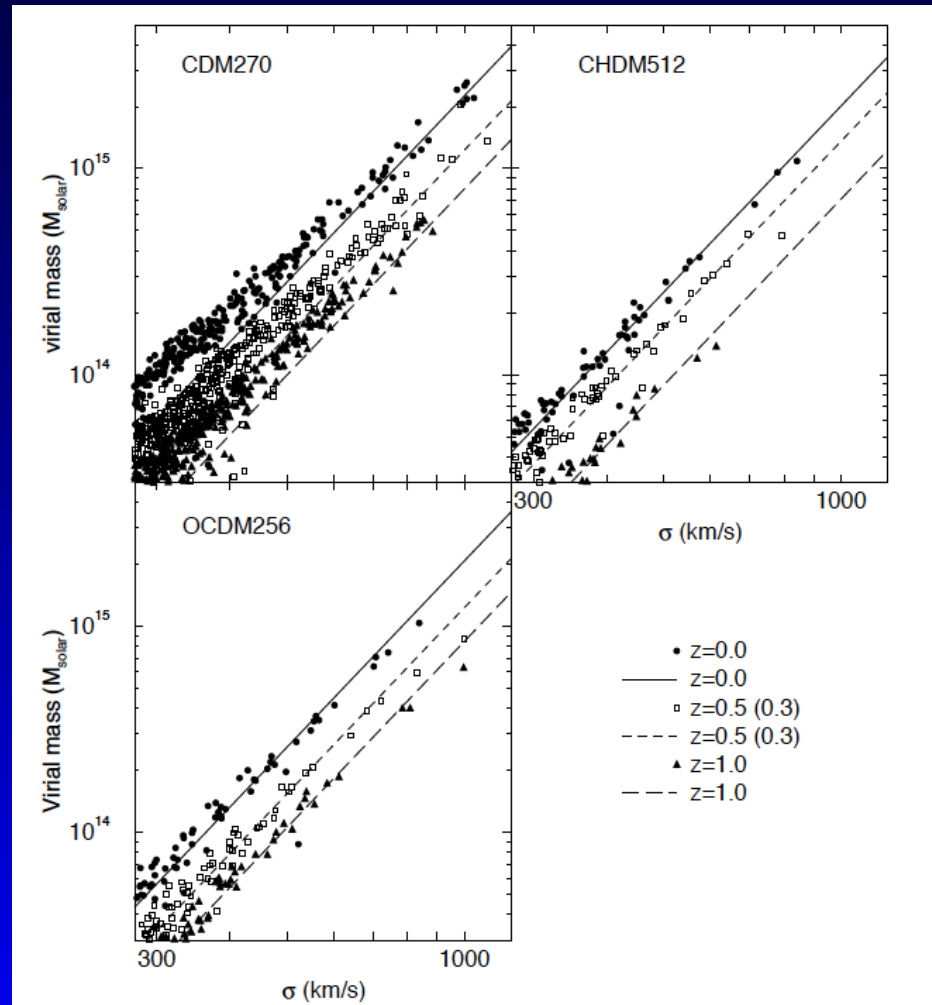
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Scaling laws (dependene on mass and redshift).

Velocity dispersion IV

Numerically: good agreement with numerical simulations (Bryan and Norman, 1998):



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so that:

$$T_x = A_{TM} M_{15}^{2/3} (1+z) (\Omega_m \Delta / 178)^{1/3}$$

The mass function

Inspired from Press and Schechter (1974)
The density field $\rho(x)$ has to be smoothed:

$$\tilde{\delta}(x) = \int \delta(x + u) W_R(u) du$$

and

$$\overline{\tilde{\delta}^2(x)} = \sigma^2(R)$$

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For a top hat window (!):

$$M(R) = \frac{4\pi}{3} R^3 \bar{\rho}$$

The mass function

dV will be included in a NL object with mass greater than M if included in a fluctuation of radius $> R$ and which is satisfying the non linear criteria ($\delta > \delta_{NL}$).

$$\int_M^{+\infty} mn(m)dm = \bar{\rho} \int \mathcal{F}_\delta(\delta) s(\delta) d\delta \sim \bar{\rho} \int_{\delta_{NL}}^{+\infty} \mathcal{F}_\delta(\delta) d\delta$$

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for a sharp threshold:

$$\int_M^{+\infty} mn(m)dm = \bar{\rho} \int_{\nu_{NL}}^{+\infty} \mathcal{F}(\nu) d\nu$$

The mass function

Following the spherical model:

$$\nu_{NL} = \frac{\delta_{NL}}{\sigma(M)}$$

Just derive against M :

$$N(M) = -\frac{\rho}{M^2 \sigma(M)} \delta_{NL} \frac{\ln \sigma}{\ln M} \mathcal{F}(\nu_{NL})$$

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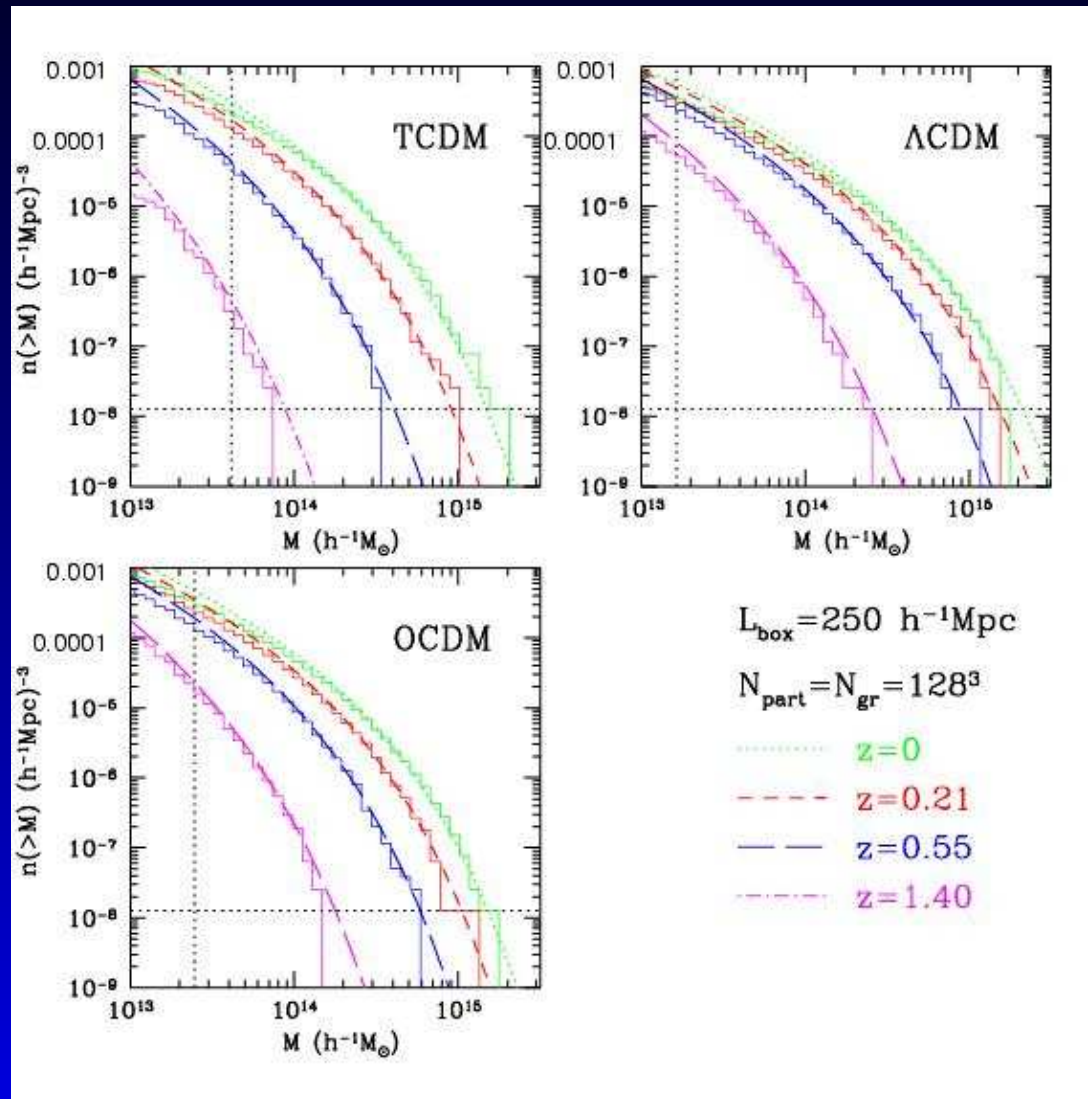
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and test it against numerical simulations...

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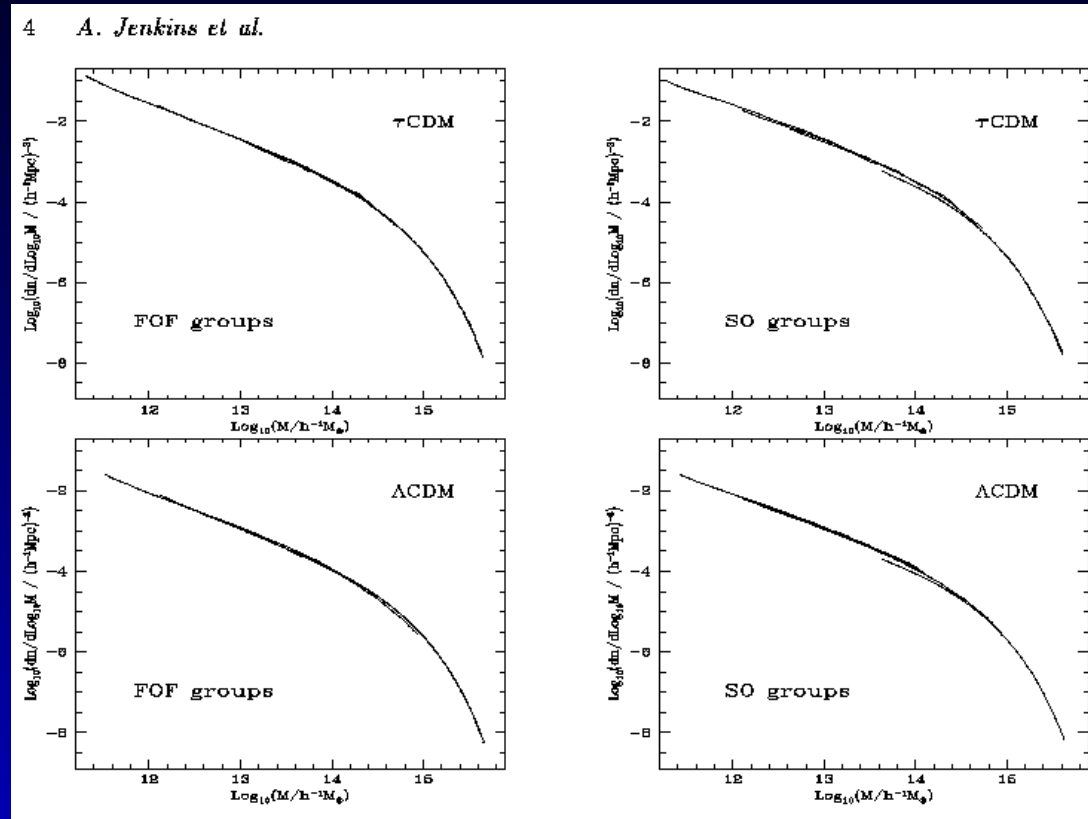
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It actually works! (Borgani et al., 2000)

More accurate $N(m)$

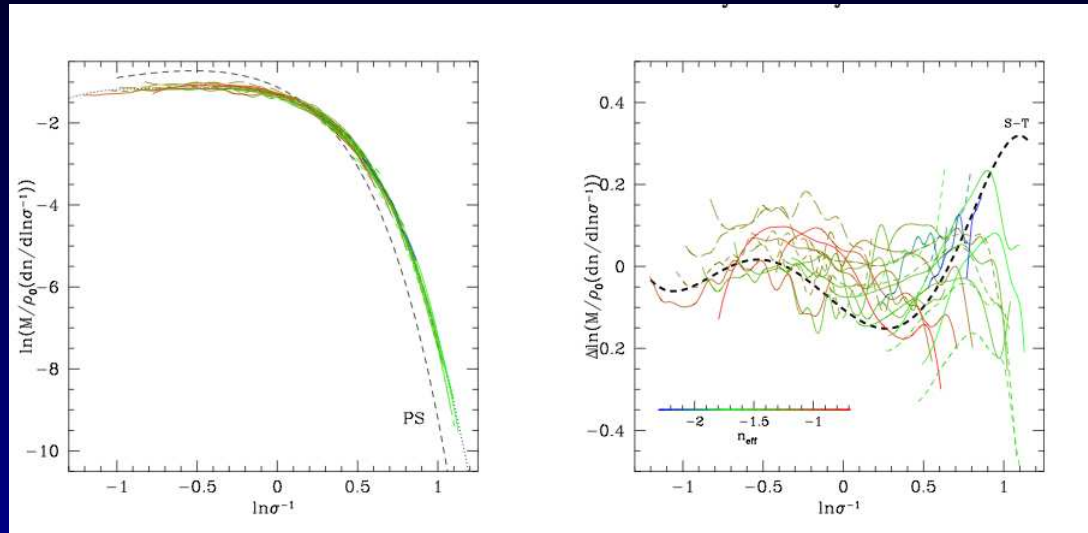
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Jenkins et al. (2001)

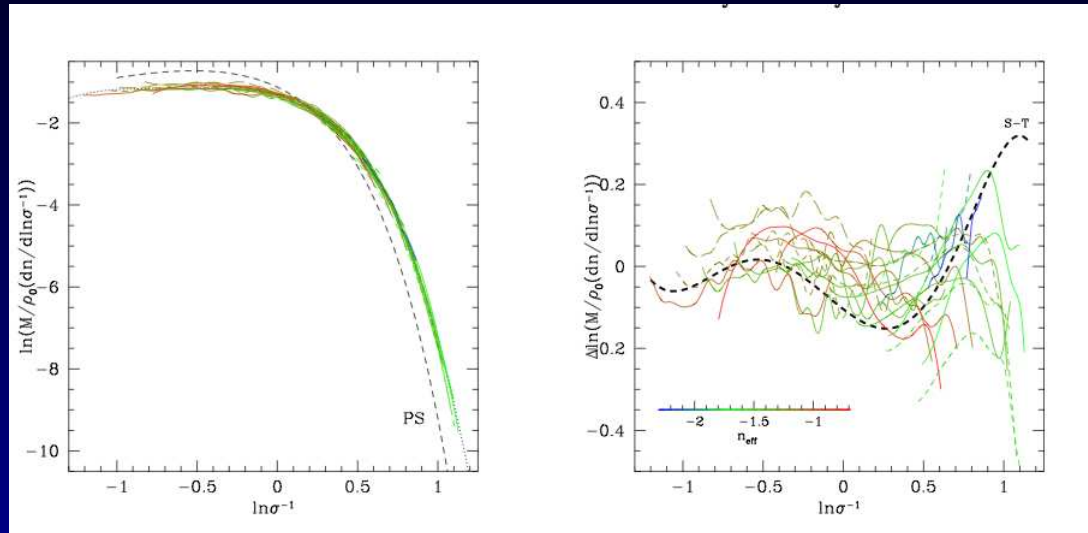
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Universal mass function.

e-PS formula

ST (Steth & Tormen, 1999) expression for \mathcal{F} :

$$\mathcal{F}(\nu) = \sqrt{\frac{2A}{\pi}} C \exp(-0.5A\nu^2) (1. + (1./(A\nu)^2)^Q)$$

with $A = 0.707$ $C = 0.3222$ $Q = 0.3$.

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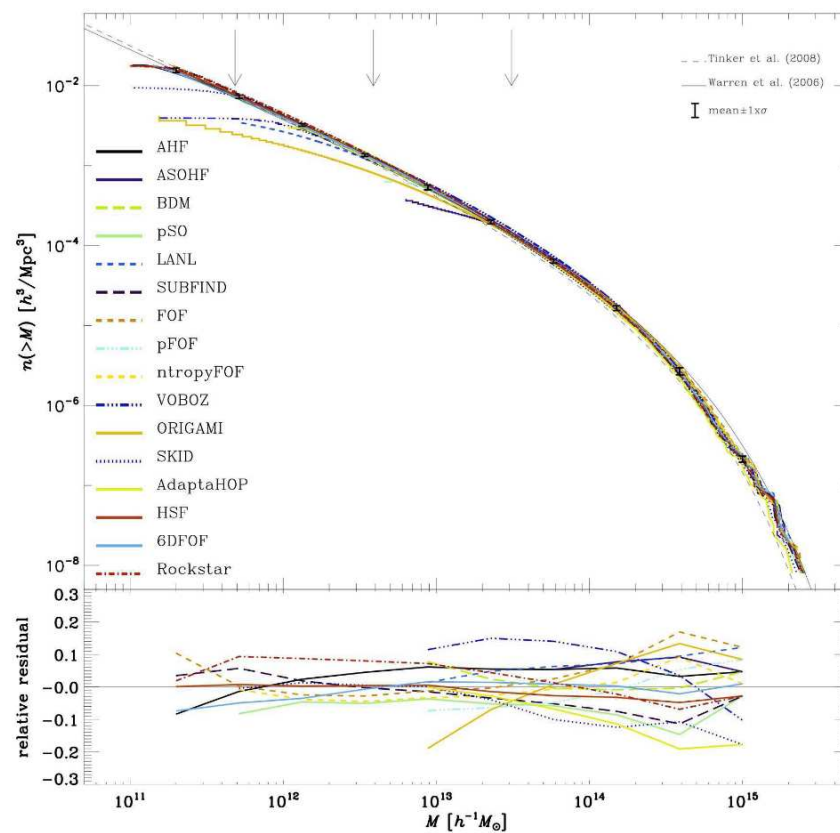
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- Allows to investigate structure formation. History of individual structure is missing: merging tree \rightarrow semi-analytical method “SAM” in order to model galaxy formation : assembly/evolution.
- Warning: data come through “light” which is coming from baryons and this was almost not discussed in these lectures...