



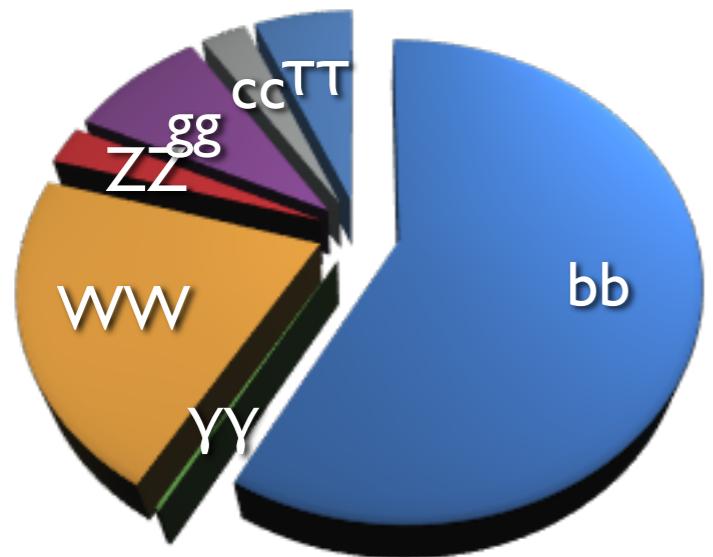
Search for $t\bar{t}H$ events using the Matrix Element Method

Lorenzo Bianchini
IPP (ETH Zurich)

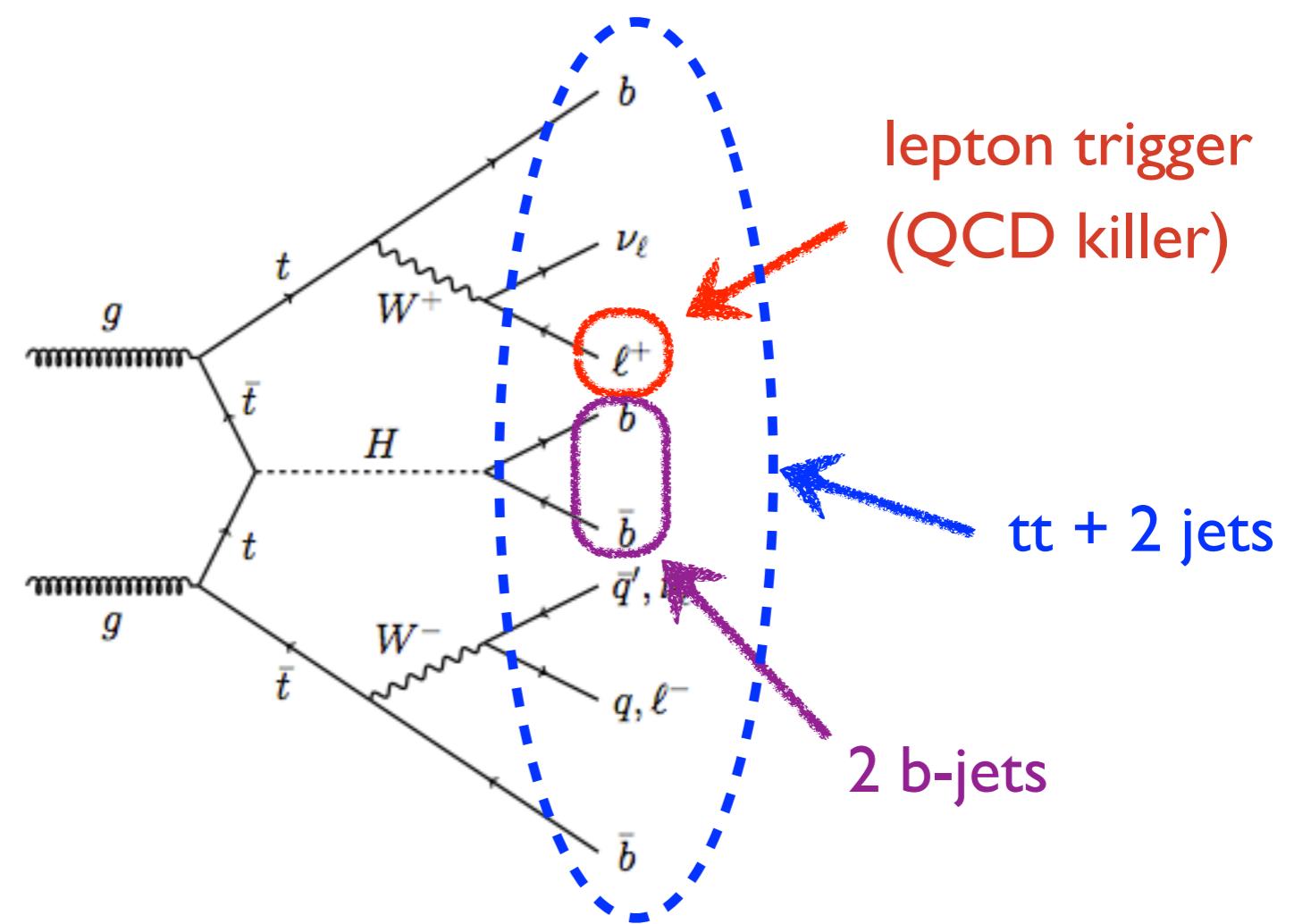
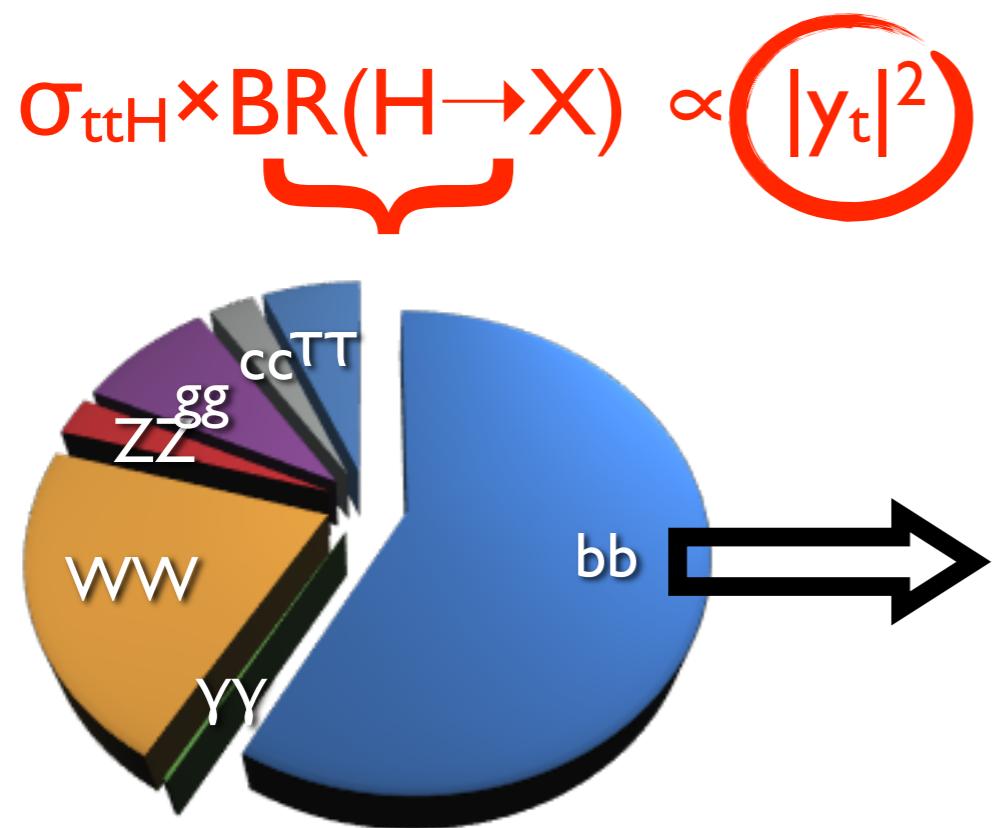
GDR Terascale @ Heidelberg, 12 Dec. 2014

Tackling ttH, H \rightarrow bb final states

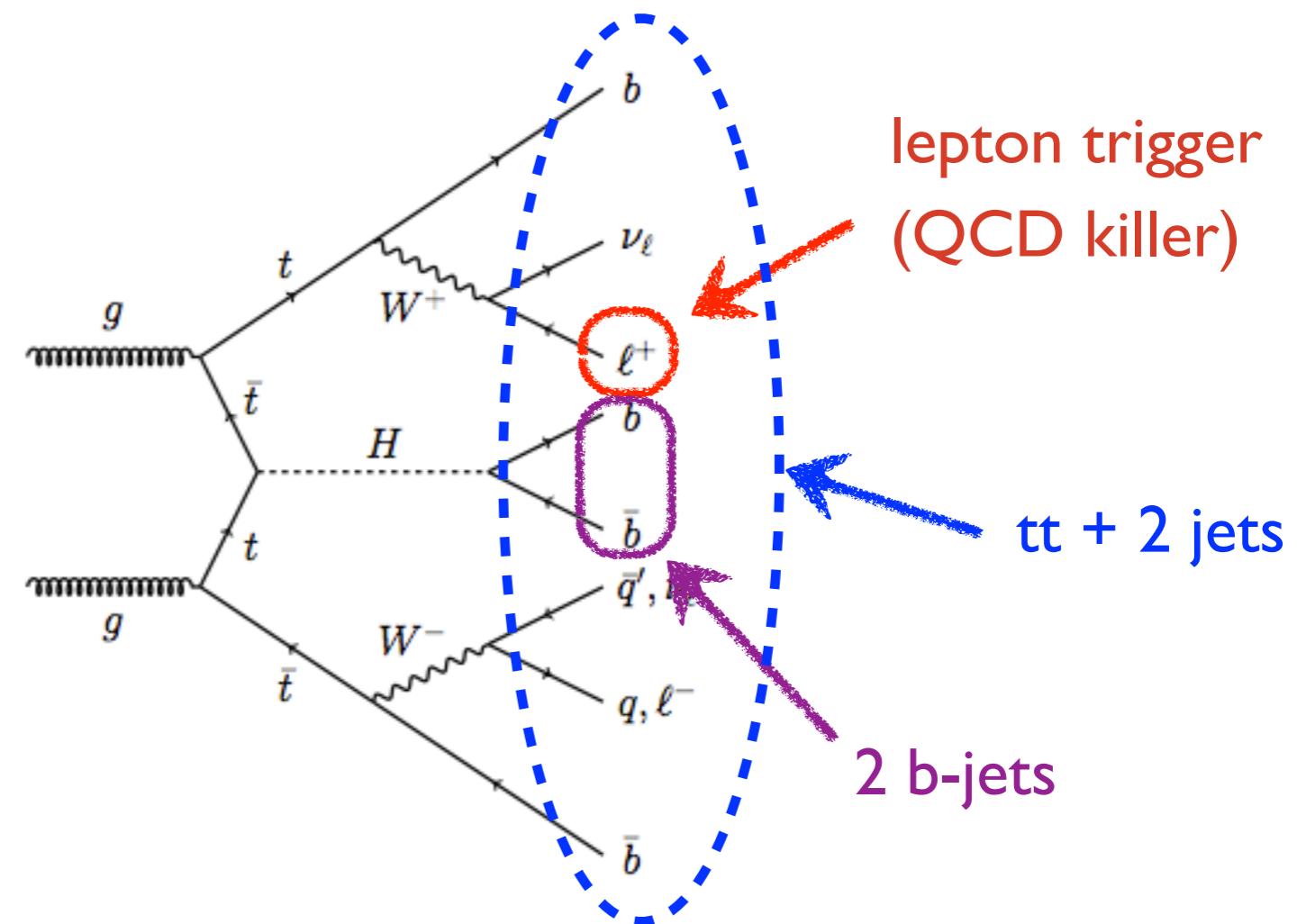
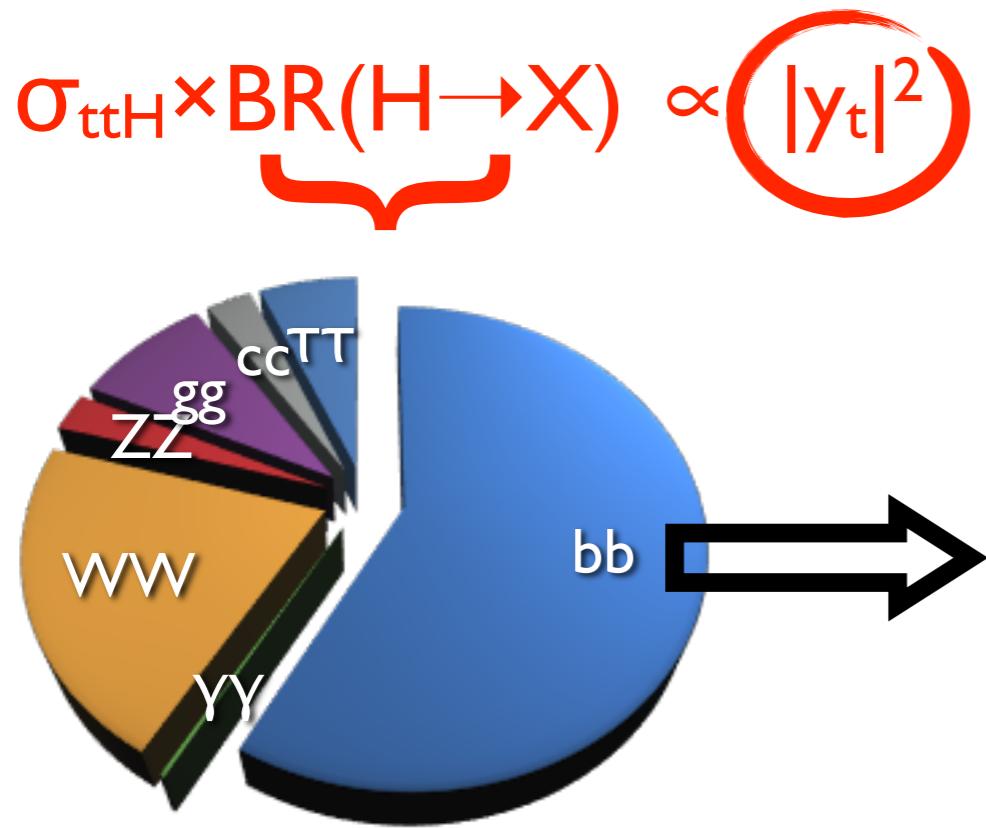
$$\sigma_{ttH} \times BR(H \rightarrow X) \propto |y_t|^2$$



Tackling ttH, H \rightarrow bb final states



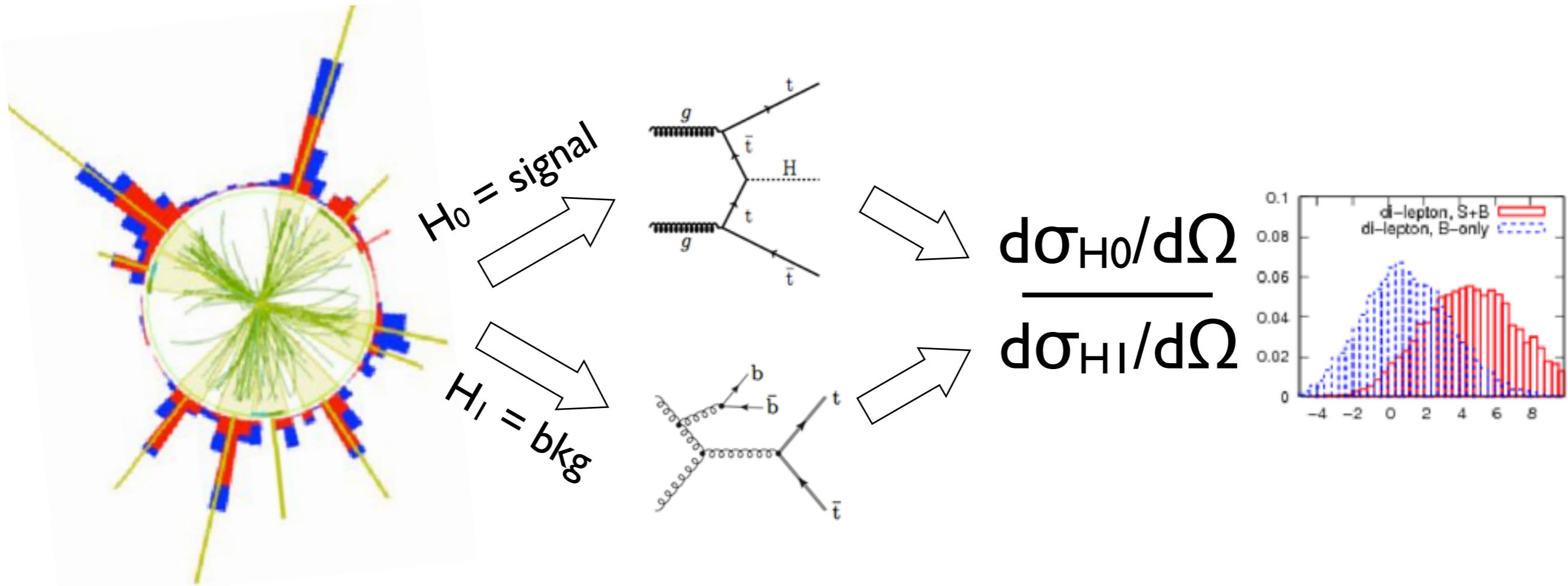
Tackling ttH, H \rightarrow bb final states



- **tt+bb background** irreducible and plagued by large unc. ($\geq 35\% @NLO$)
 - ▶ counting experiment not feasible
 - ▶ extra handles: m_{bb} spectrum
 - ▶ separable using mass peak ?
- Not so easy: b's from top quarks complicate mass peak extraction
 - ▶ **combinatorial** self-background \Rightarrow need for **multidimensional** analyses

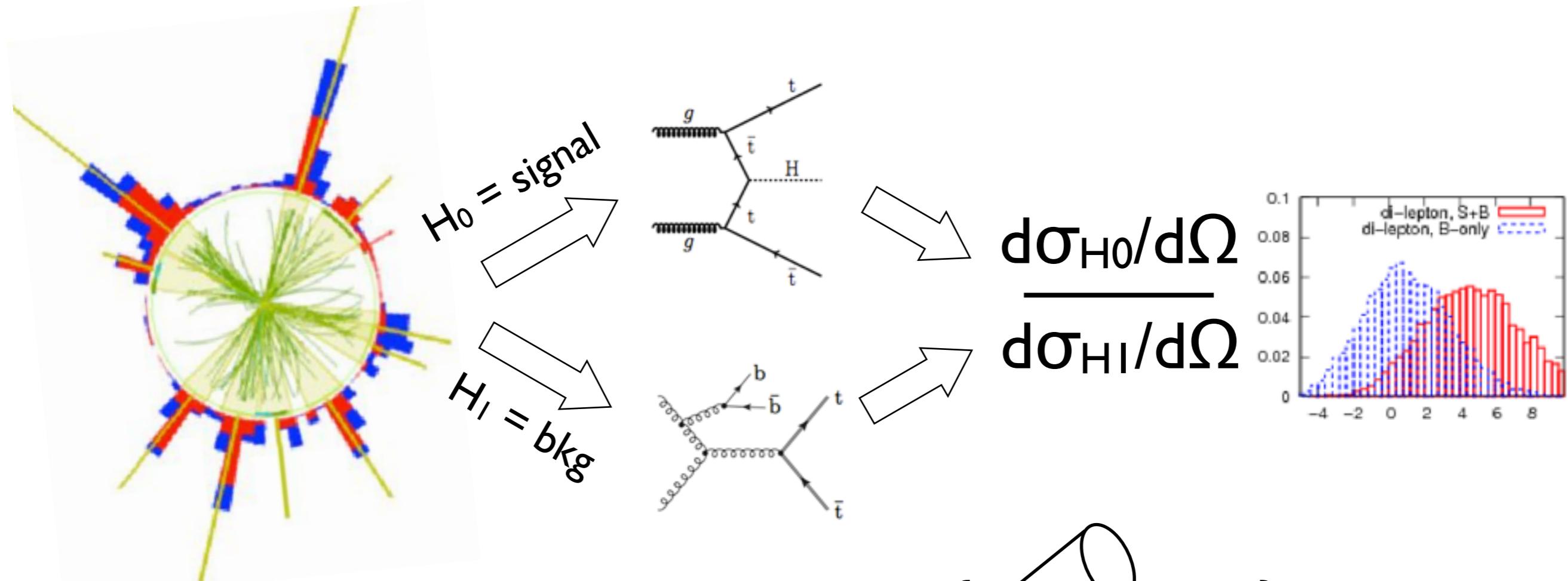
A possible analysis strategy

Matrix Element Method (**MEM**) suited to the task



A possible analysis strategy

Matrix Element Method (**MEM**) suited to the task



Virtues of the MEM:

- ▶ combinatorial issue properly addressed:
- ▶ optimal usage of kinematics and dynamics for S/B separation (**Neyman lemma**)

$$P\{ \dots | t, H \} = P\{ \dots \} + P\{ \dots \}$$

t H

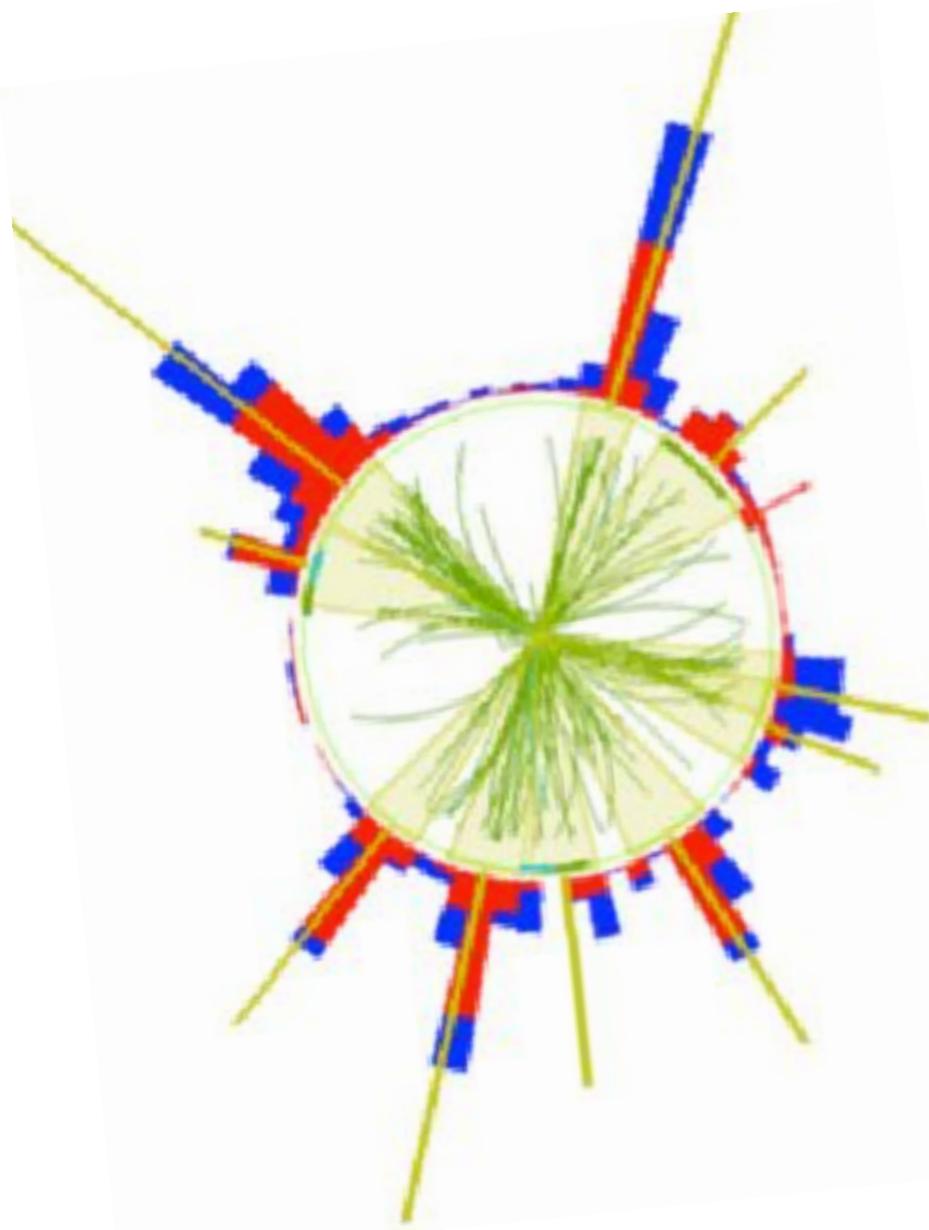
t H

Master formula

$$\frac{d\sigma_{S/B}}{d\vec{Y}} = w_{S/B}(\vec{Y}) = \int d\Phi_x dx_a dx_b f(x_a, x_b) \delta^4(x_a P_a + x_b P_b - \vec{X}) |\mathcal{M}_{S/B}(\vec{X})|^2 W(\vec{Y}, \vec{X})$$

Master formula

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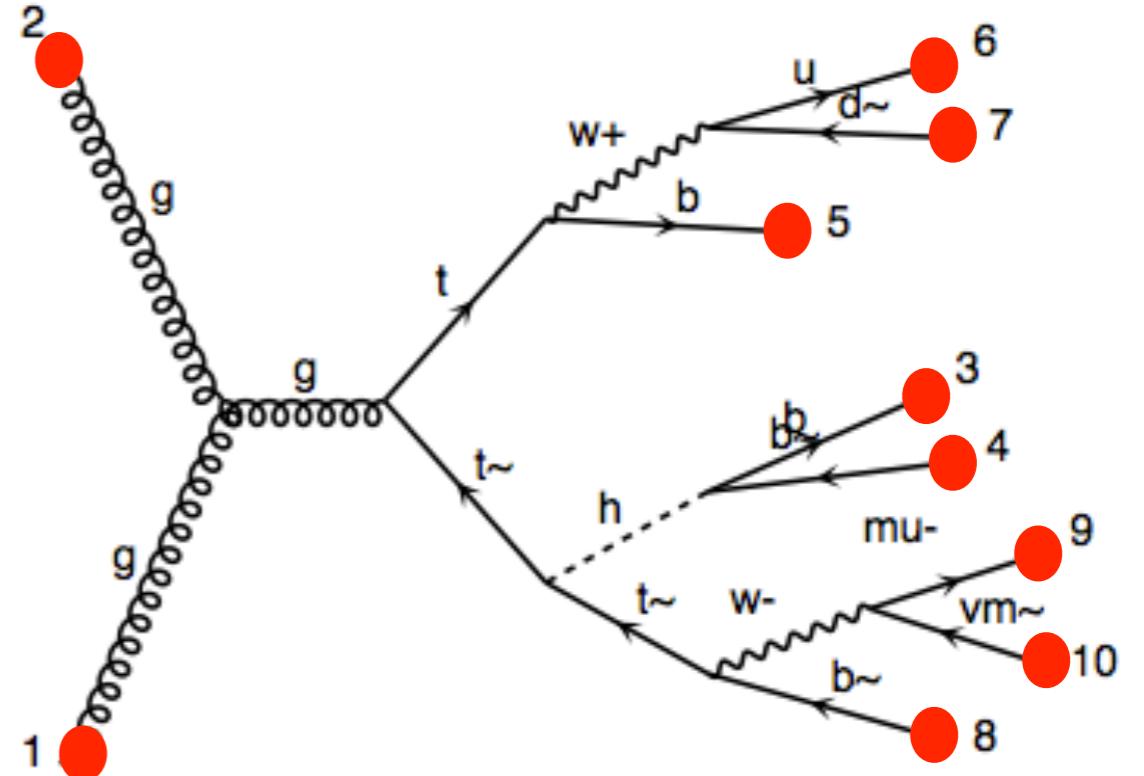
reconstructed jet and
lepton momenta

Master formula

$$d\sigma_{S/B}/d\vec{Y} = w_{S/B}(\vec{Y}) =$$

$$\int d\Phi_x dx_a dx_b f(x_a, x_b) \delta^4(x_a P_a + x_b P_b - \vec{X}) |\mathcal{M}_{S/B}(\vec{X})|^2 W(\vec{Y}, \vec{X})$$

generated particles;
numerical integration
(VEGAS)

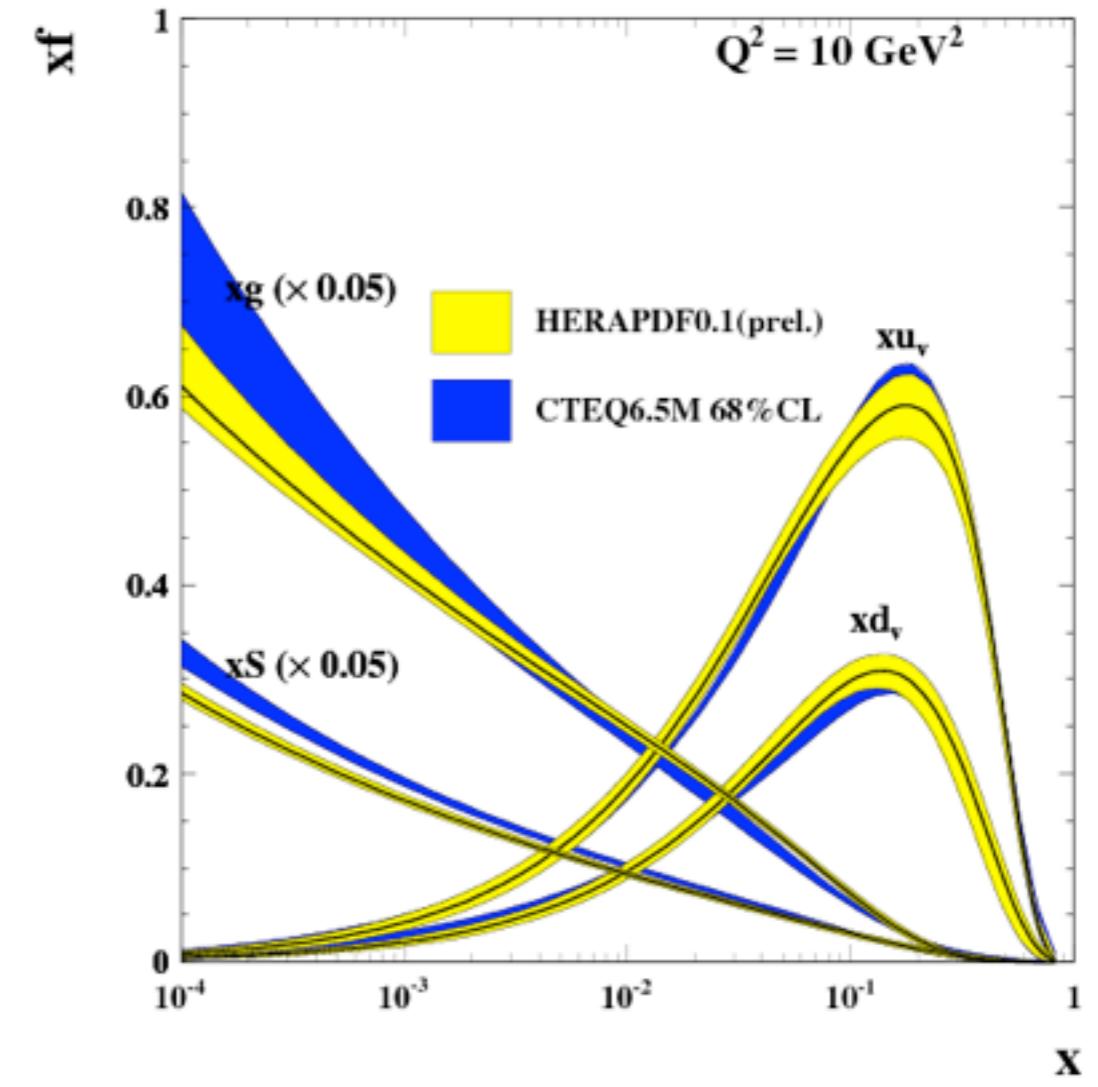


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PDF from LHAPDF
(CTEQ6.5m set)



Master formula

$$d\sigma_{S/B}/d\vec{Y} = w_{S/B}(\vec{Y}) =$$

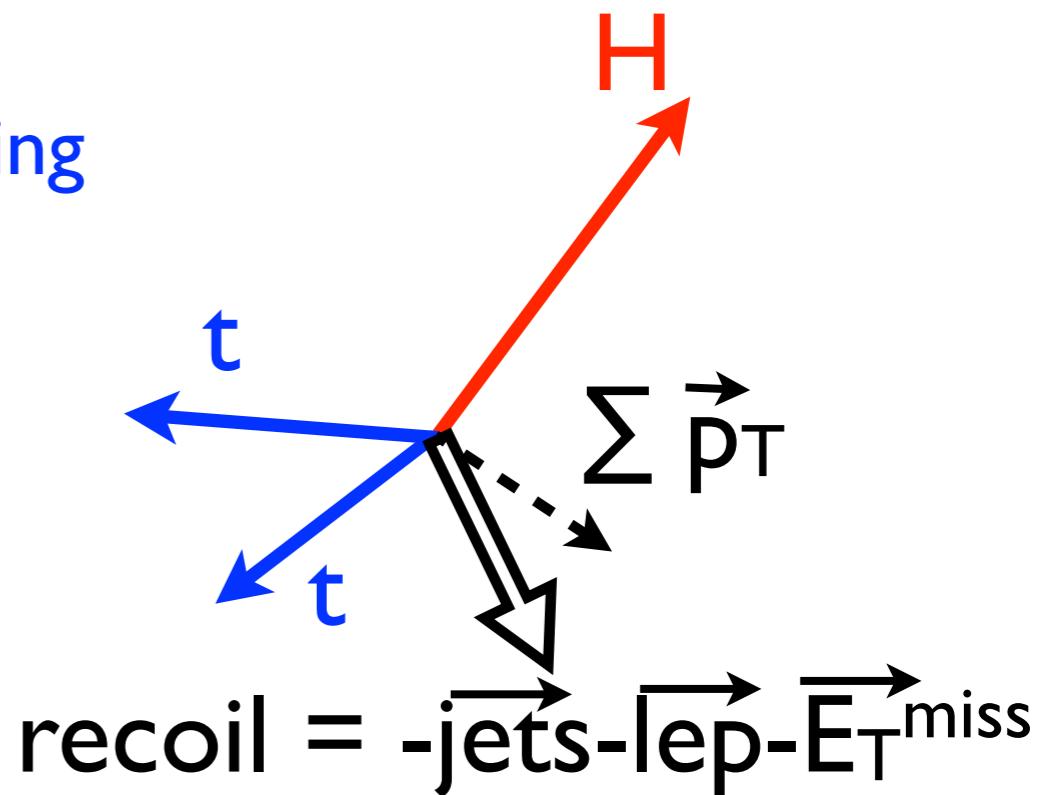
$$\int d\Phi_x dx_a dx_b f(x_a, x_b) \boxed{\delta^4(x_a P_a + x_b P_b - \vec{X})} |\mathcal{M}_{S/B}(\vec{X})|^2 W(\vec{Y}, \vec{X})$$



Momentum balancing

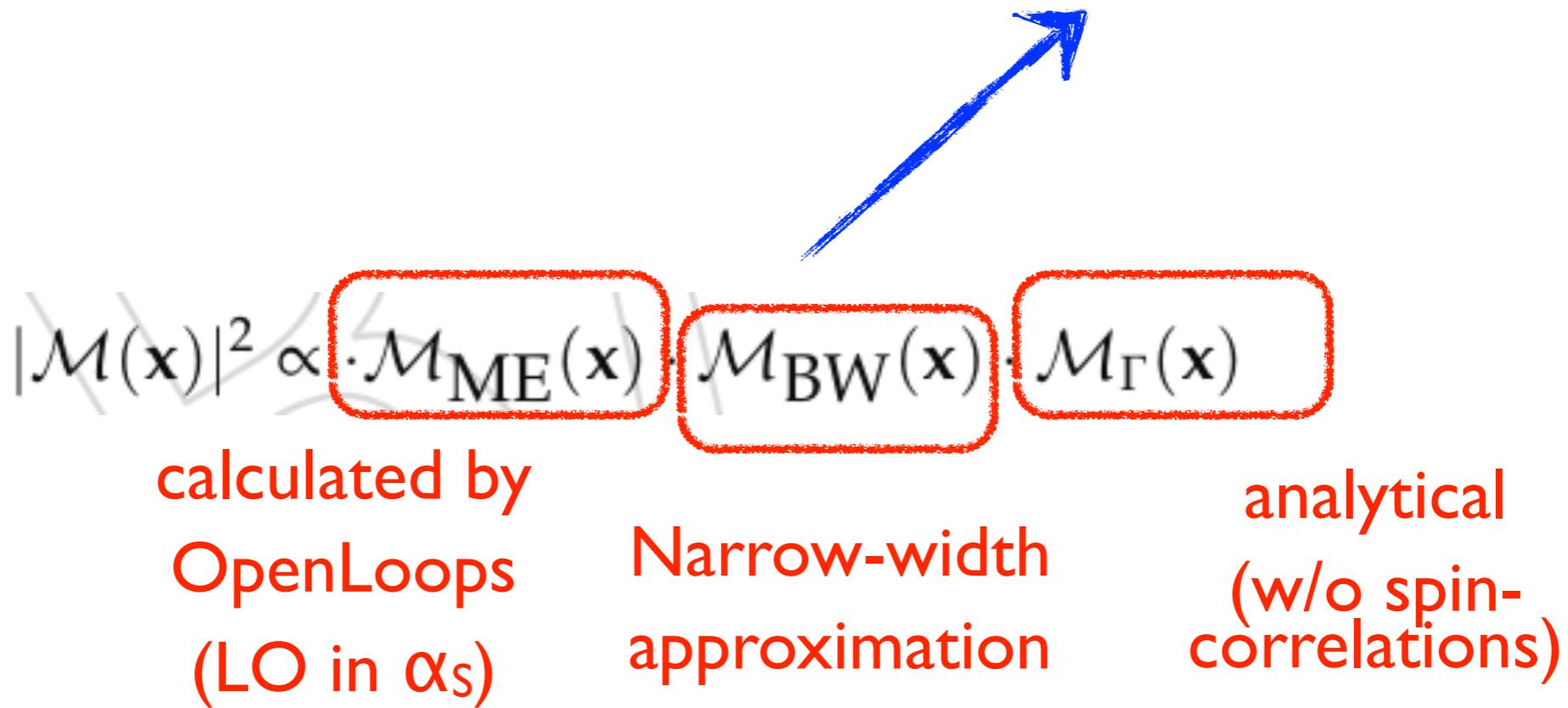
N.B.: $\sum \vec{p}_T = 0 @ LO$

Here: p_T balance not imposed.
 $\sum \vec{p}_T$ constrained to measured
recoil via a TF



Master formula

$$\frac{d\sigma_{S/B}}{d\vec{Y}} = w_{S/B}(\vec{Y}) = \int d\Phi_x dx_a dx_b f(x_a, x_b) \delta^4(x_a P_a + x_b P_b - \vec{x}) |\mathcal{M}_{S/B}(\vec{x})|^2 W(\vec{Y}, \vec{x})$$

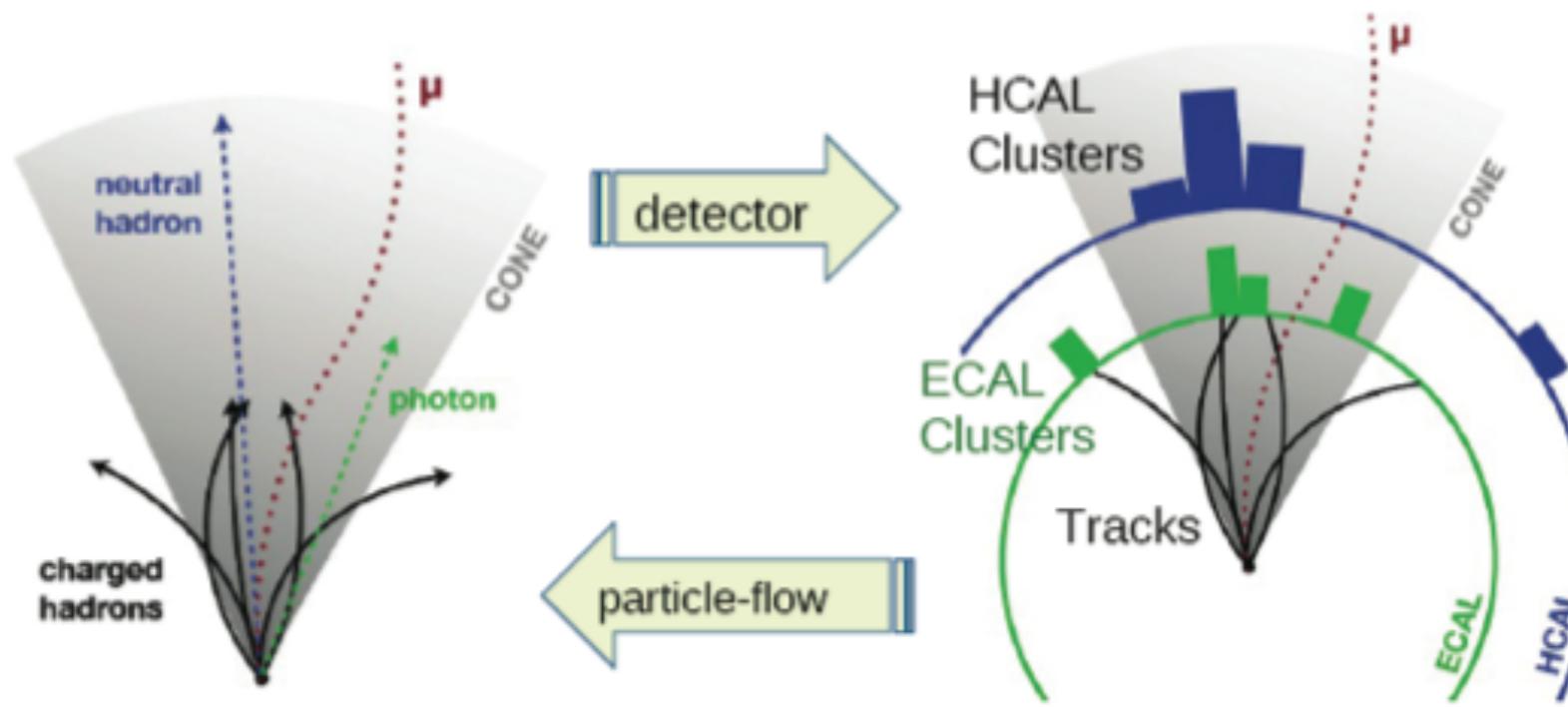


Master formula

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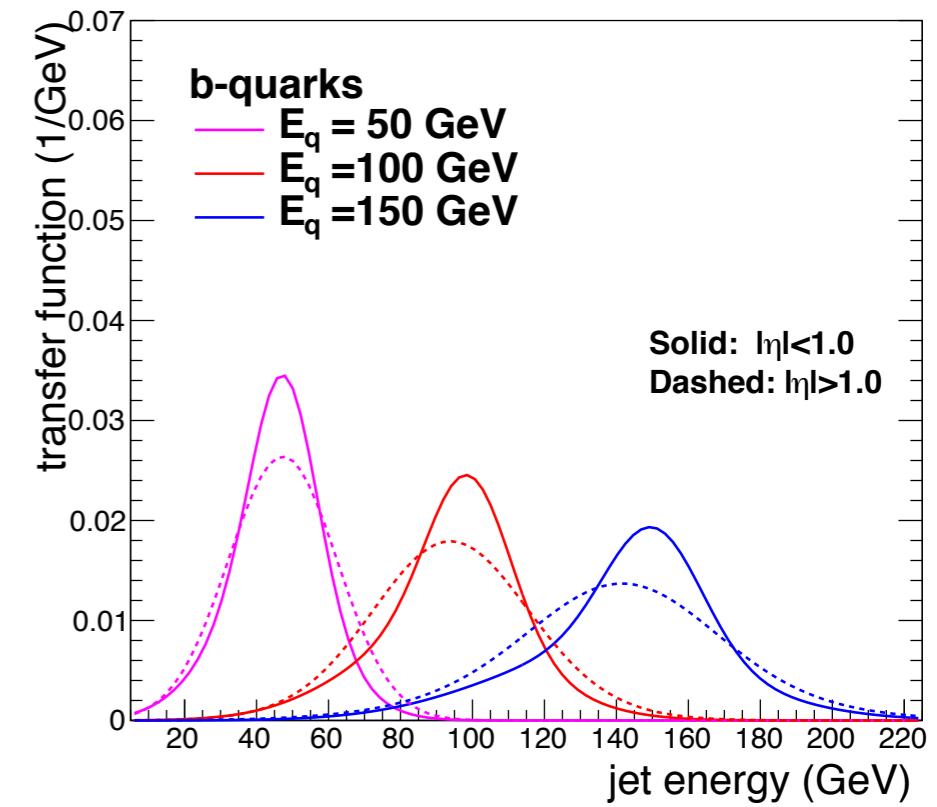
$$\int d\Phi_x dx_a dx_b f(x_a, x_b) \delta^4(x_a P_a + x_b P_b - \vec{X}) |\mathcal{M}_{S/B}(\vec{X})|^2 W(\vec{Y}, \vec{X})$$

Transfer function



13

CMS Simulation $\sqrt{s}=8$ TeV



Implementation

SL

Categorize events as to
aid evaluation of ME at LO

category \Leftrightarrow event interpretation

DL

Implementation

SL
 $N_j \geq 5$

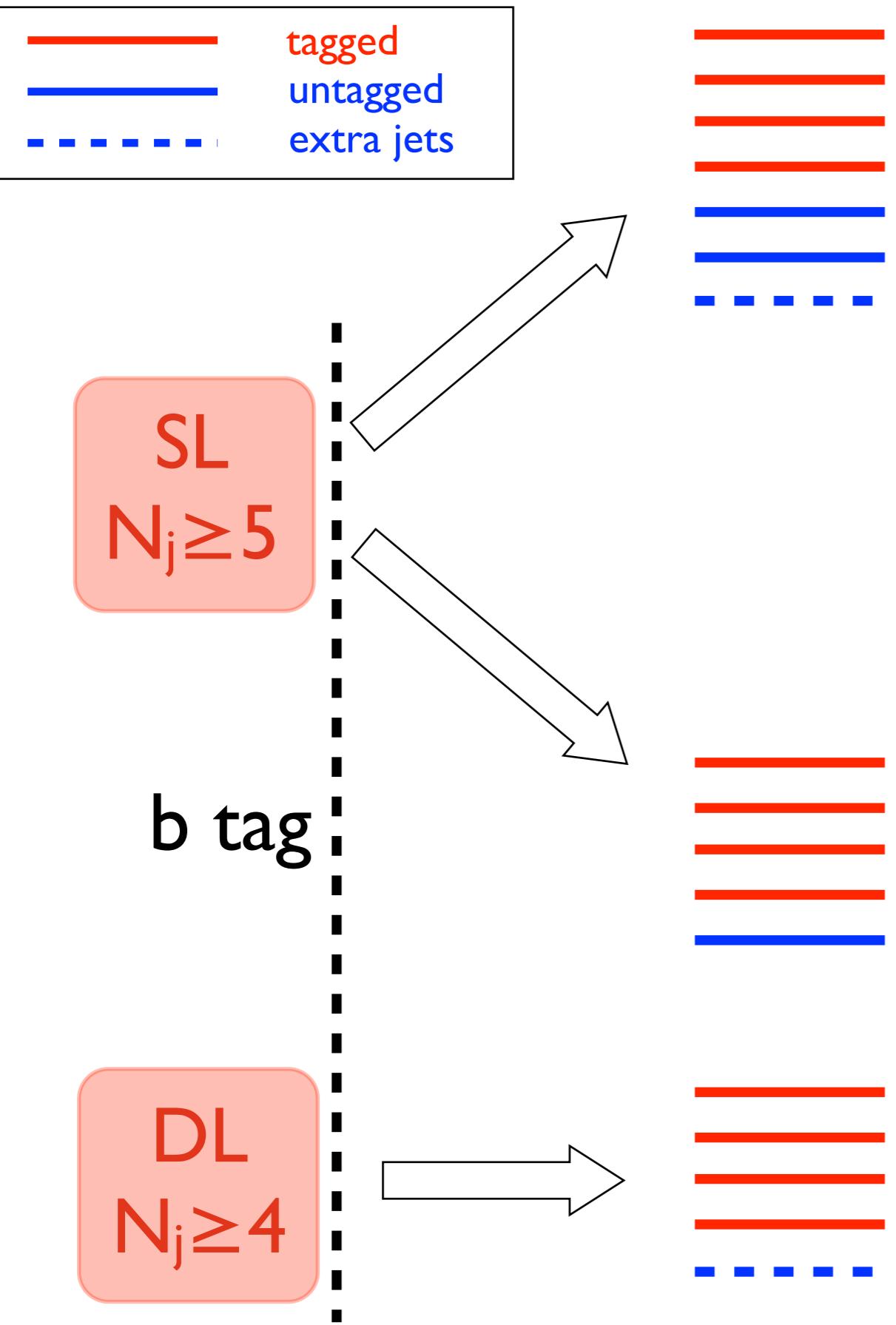
b tag

DL
 $N_j \geq 4$

Reduction of V+jets and tt+jets requires cut on
number of jets, b tagging

Implementation

assignment based on jet
permutation w/ largest btag



Implementation

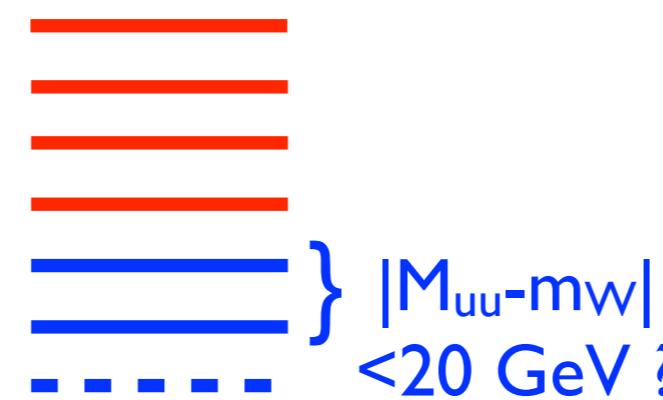
assignment based on jet permutation w/ largest btag



SL
 $N_j \geq 5$

b tag

DL
 $N_j \geq 4$



yes

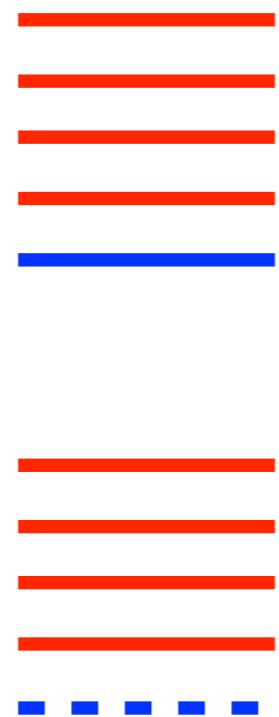
no

SL Cat. I

all top/H quarks
reconstructed

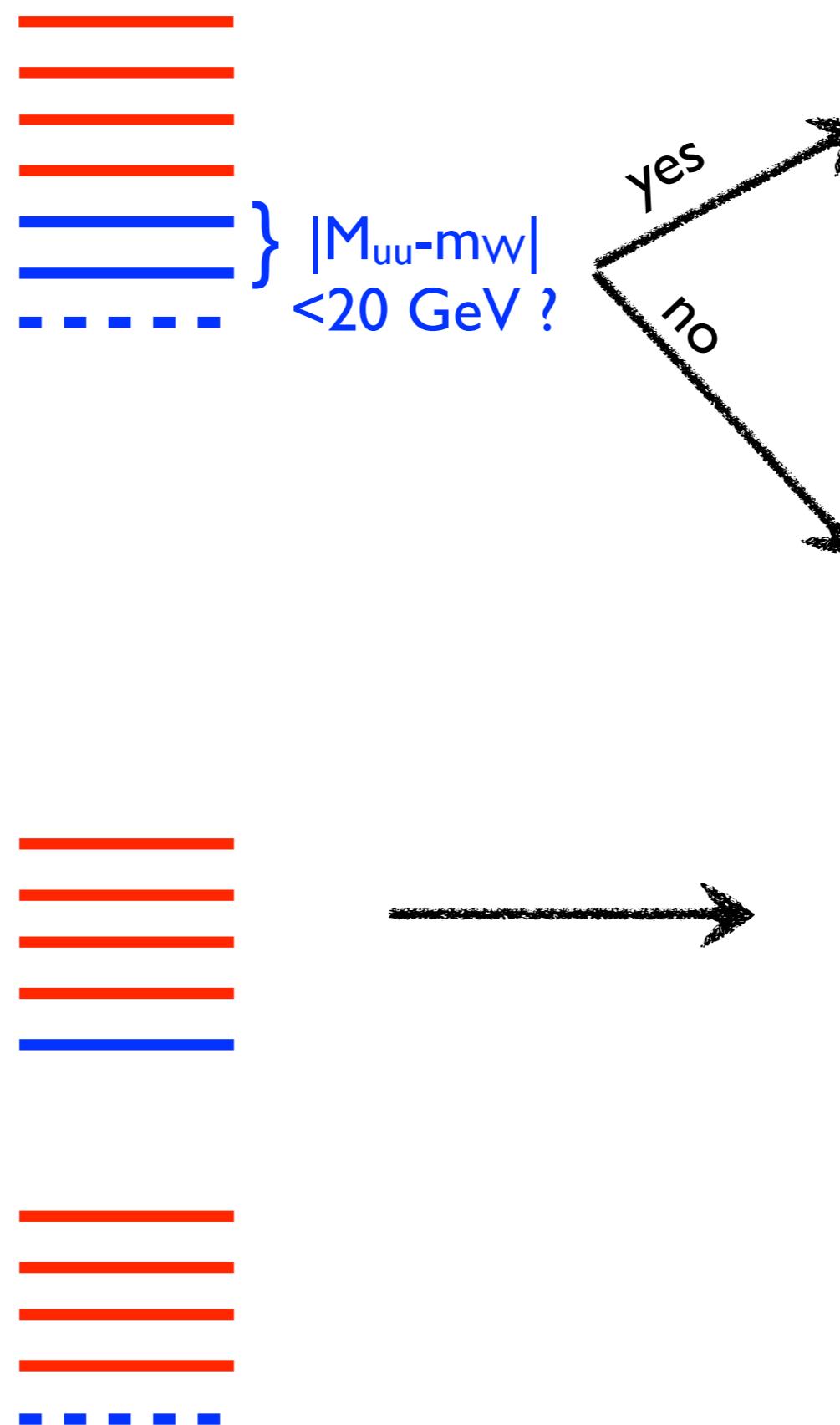
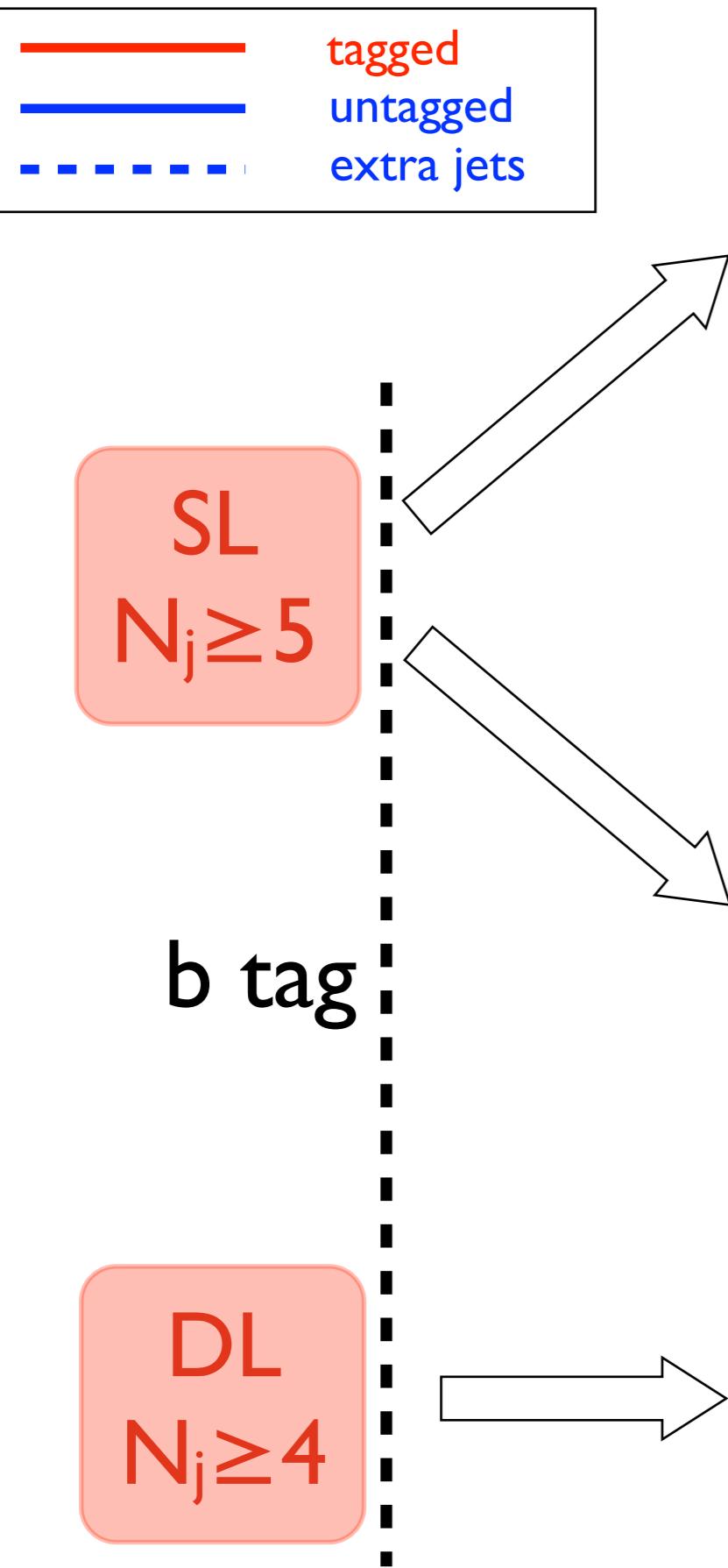
SL Cat. 2

one W-quark missed;
extra gluon(s) from ISR
⇒ integrate over
missing quark



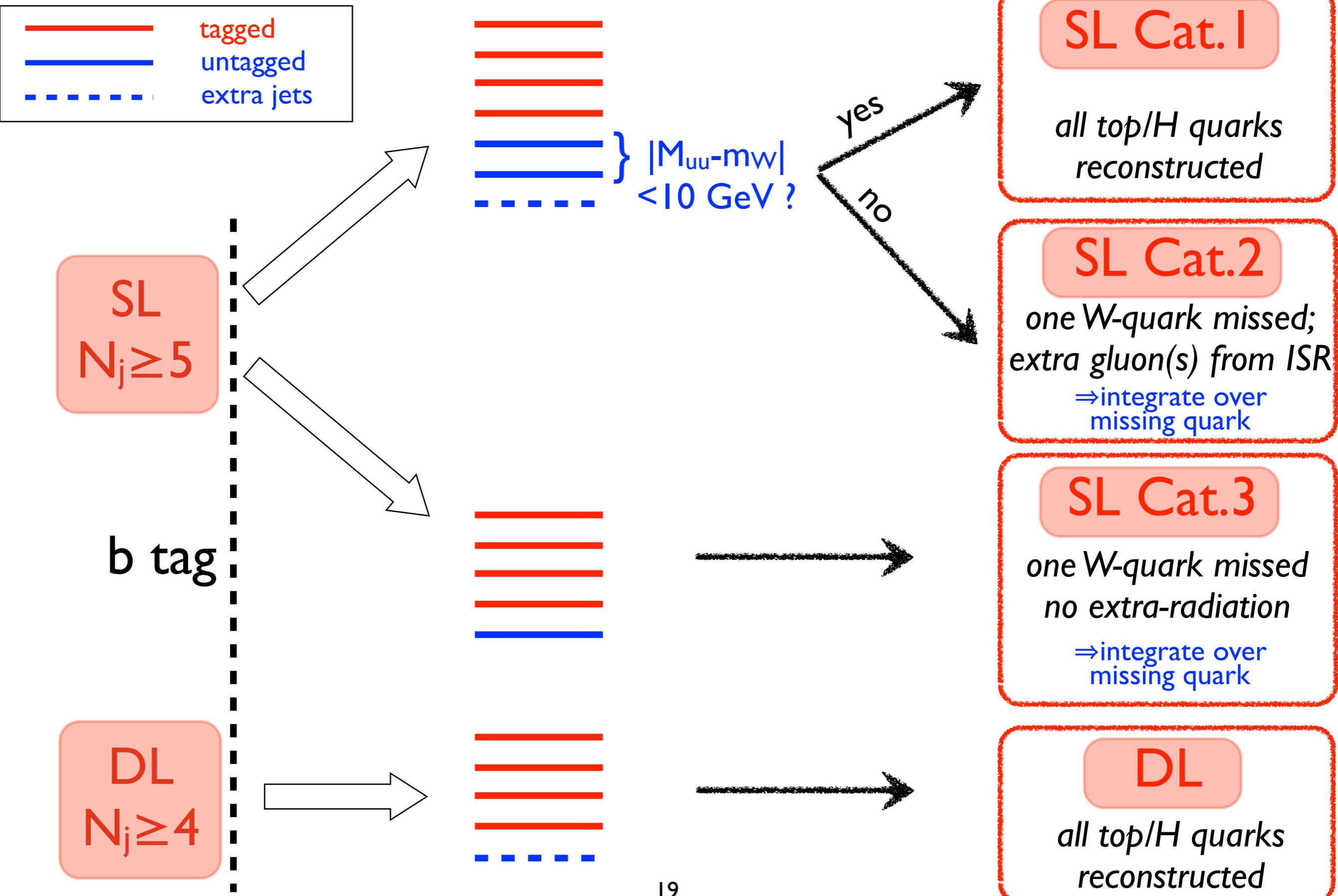
Implementation

assignment based on jet permutation w/ largest btag

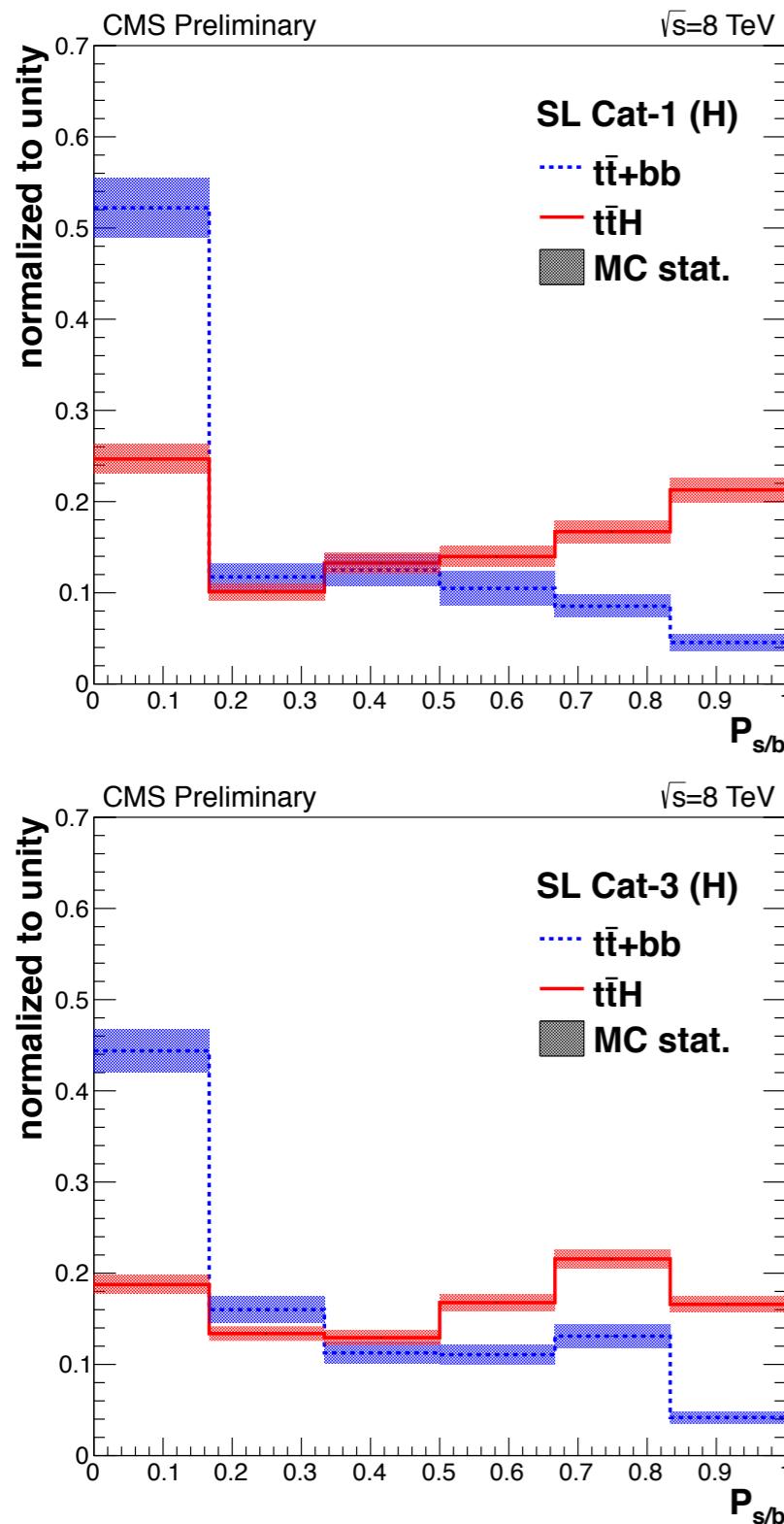


Implementation

assignment based on jet
permutation w/ largest btag



Shape comparison: tt+bb vs ttH



$$P_{s/b} = (1 + w_B/w_S)^{-1}$$

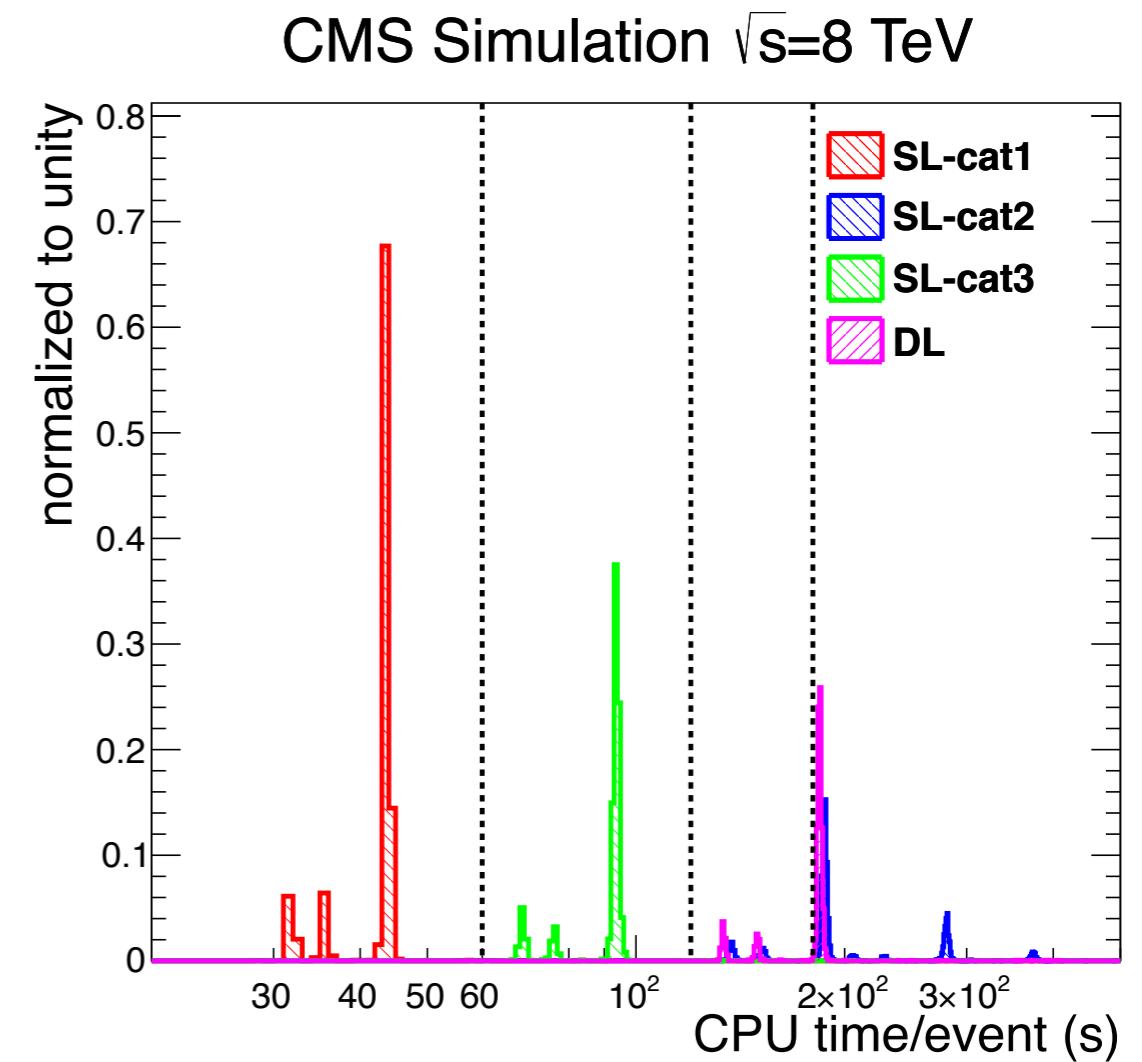
In situ validation of shapes

- ▶ sidebands with low b tagging score
- ▶ in signal region, validate low-sensitivity version of $P_{s/b}$
 - one random permutation only
 - testing different values of m_H

FAQ: is CPU time an issue ?

Not really; compromise between **performance** and **timing**:

- ▶ run only on the good “events”
- ▶ filter-out permutations using b tagging ($6! \rightarrow 4!$)
- ▶ test one background hypothesis only
 - optimize separation against tt+bb
- ▶ parallelize (by event) as much as possible
- ▶ neglect spin/correlations
- ▶ JEC/JER systematics: bookkeep VEGAS grid result from “nominal” for faster evaluation
- ▶ ...



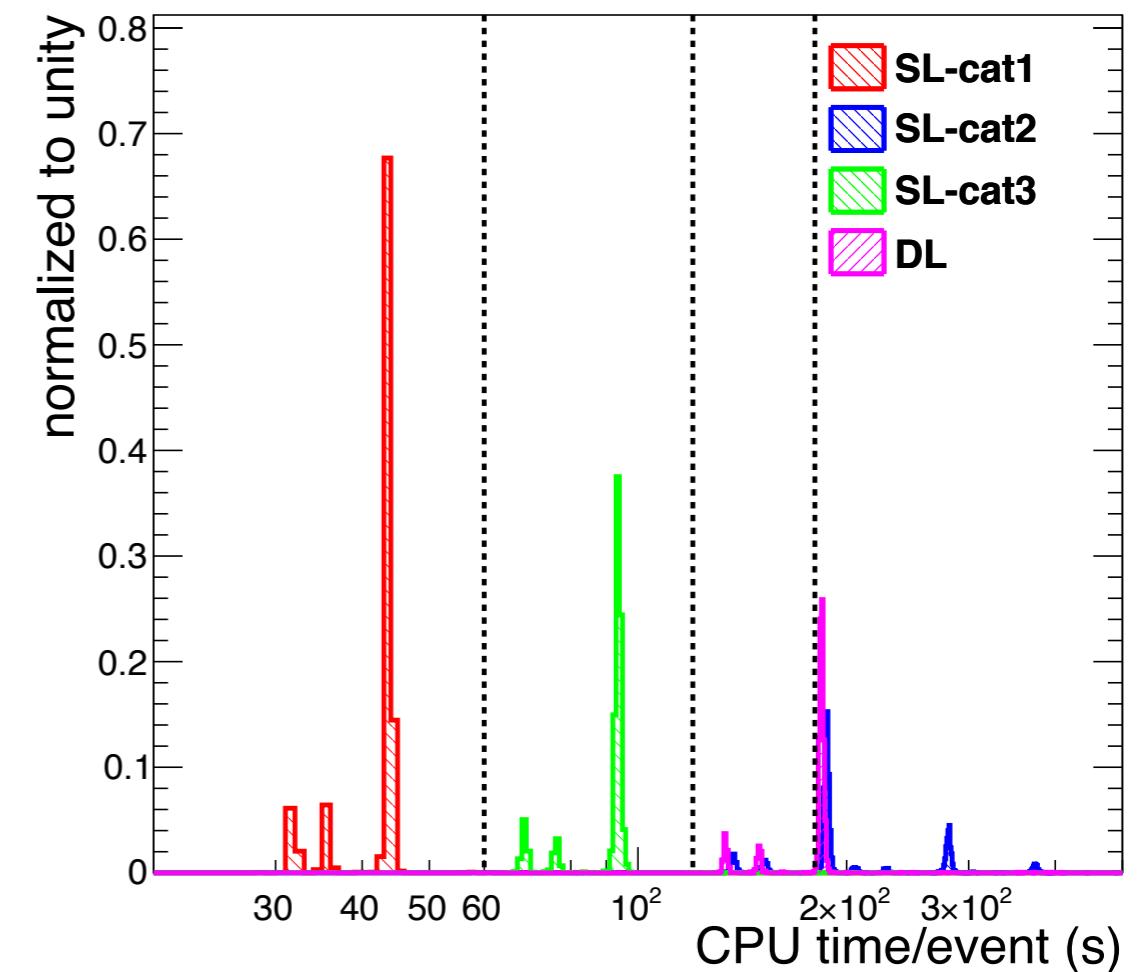
Number of variables	4 (+1)	6 (+1)	5 (+1)
Number of iterations	5	5	5
Function calls	2000	4000	10000
Numerical precision (mode of σ_w/w)	0.8%	1.2%	0.8%
CPU-time per integral (mean)	0.5 (1.5) s	1.1 (3.2) s	2.3 (6.2) s
Time budget for $ \mathcal{M} _{\text{ME}}^2$	30% (80%)	30% (80%)	30% (80%)

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- ▶ ...

CMS Simulation $\sqrt{s}=8$ TeV



Number of variables	4 (+1)	6 (+1)	5 (+1)
Number of iterations	5	5	5
Number of MC samples	100	100	100
Time budget for $ \mathcal{M} _{\text{ME}}^{\zeta}$	30% (80%)	30% (80%)	30% (80%)

⇒ complete analysis round in ~ 10 h using the batch system of a T3

Overview of systematic uncertainties

ttH modeling: PYTHIA

tt+jets: MadGraph (≤ 3 partons)+PYTHIA

50% normalization uncertainty on tt+HF

α_S scale uncertainty

JEC/JER, b tagging

top p_T modeling
statistical uncertainty

nuisance	treatment
luminosity	InN 2.6%
ID/trigger	InN 2-4%
ttH	InN 12%
tt+bb	InN 50%
tt+cc	InN 50%
tt+b	InN 50%
QCD scale	InN 17-3%
fact/renorm. scale	shapes (tt + 1p/2p/3p/bb/b/jj/cc)
PDF	InN 3-9%
JES	shape
JER	shape
btagging	8 shapes
top p _T model	shape
MC stat.	shape (all bins)

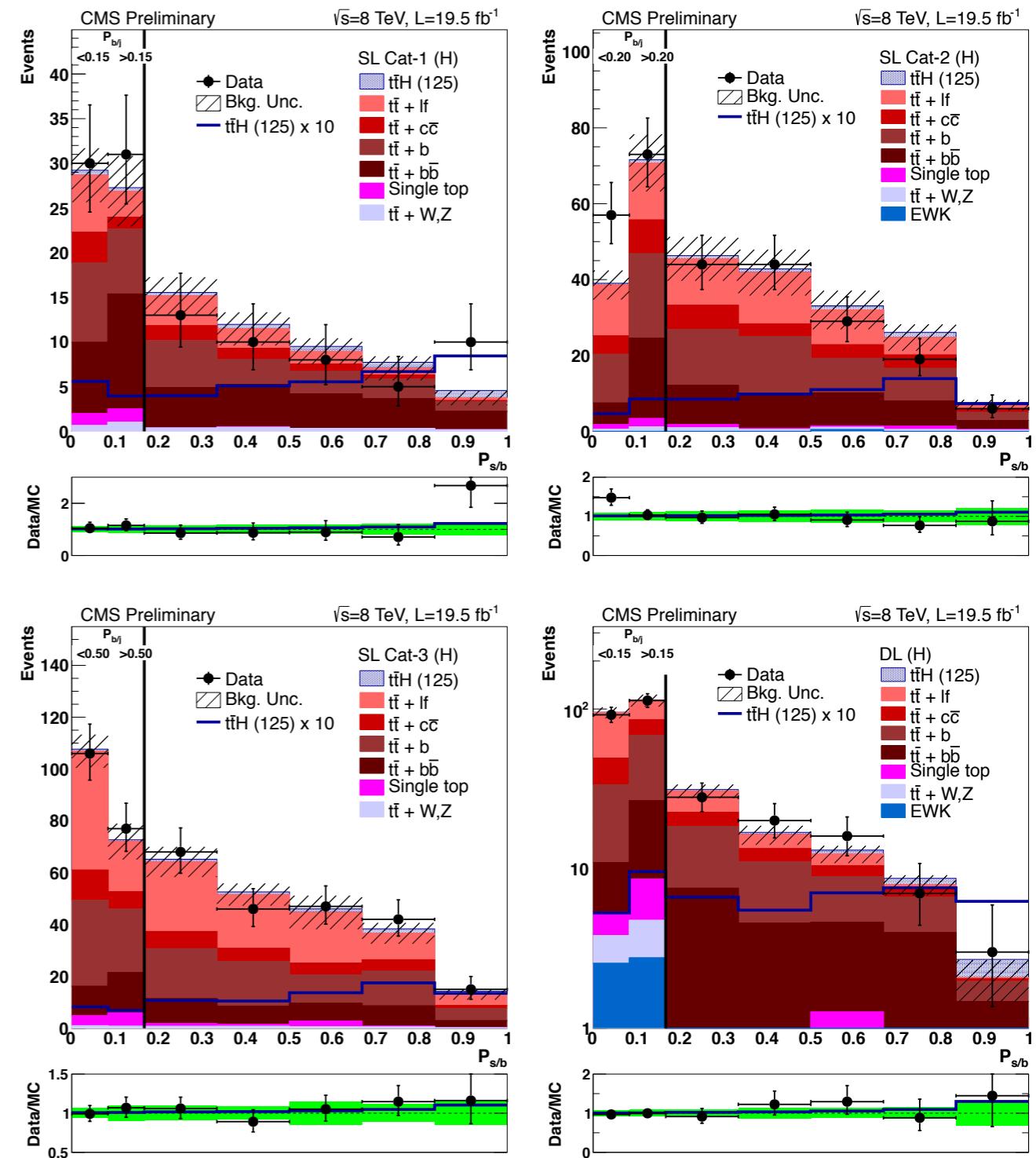
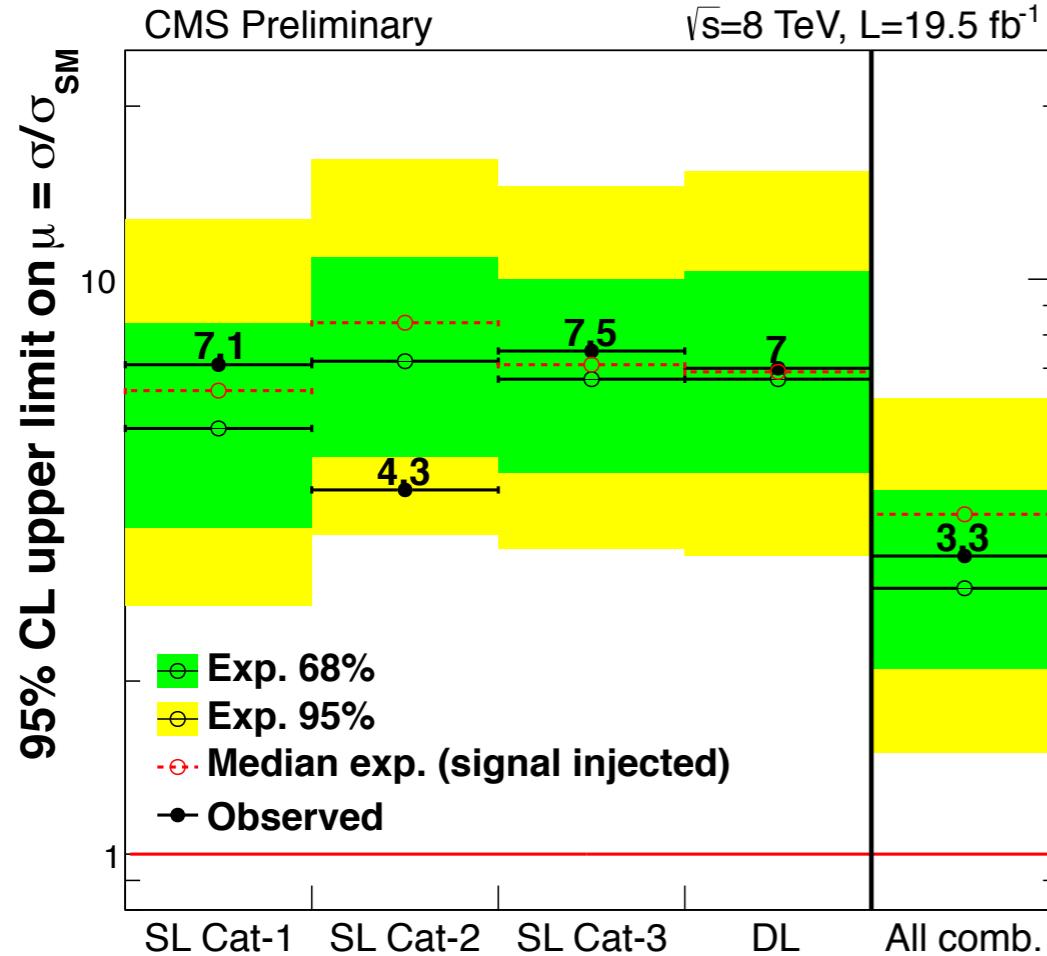
InN = normalization uncertainty
shape = vertical morphing

Results (8 TeV)

Combined fit to ME discriminant:

Median Exp. 95% CL	Median Exp. (signal injected)	Obs.
$\mu < 2.9$	$\mu < 3.9$	<3.3

► Best-fit value $\mu_{t\bar{t}H} = 0.7 \pm 1.4$

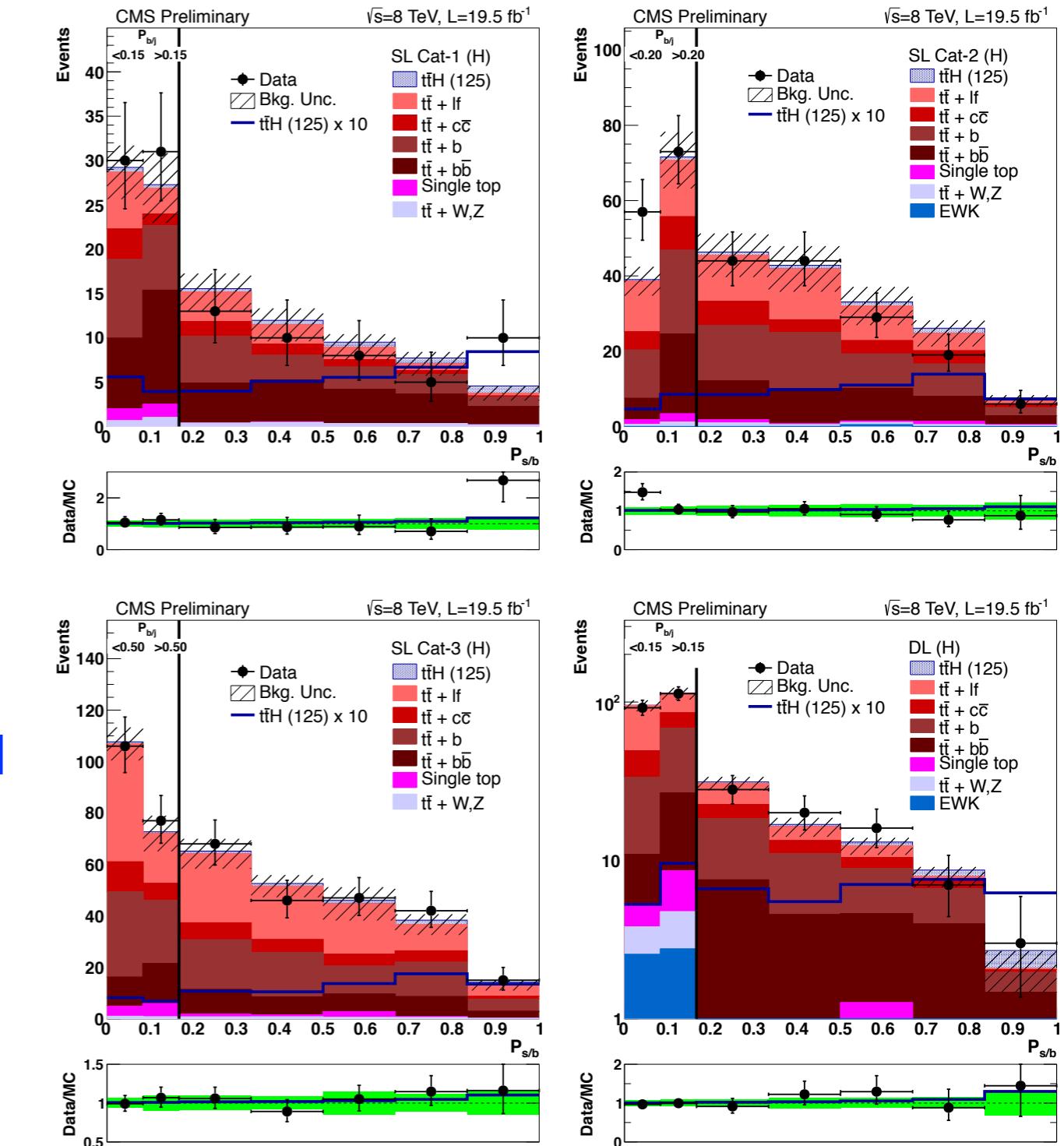
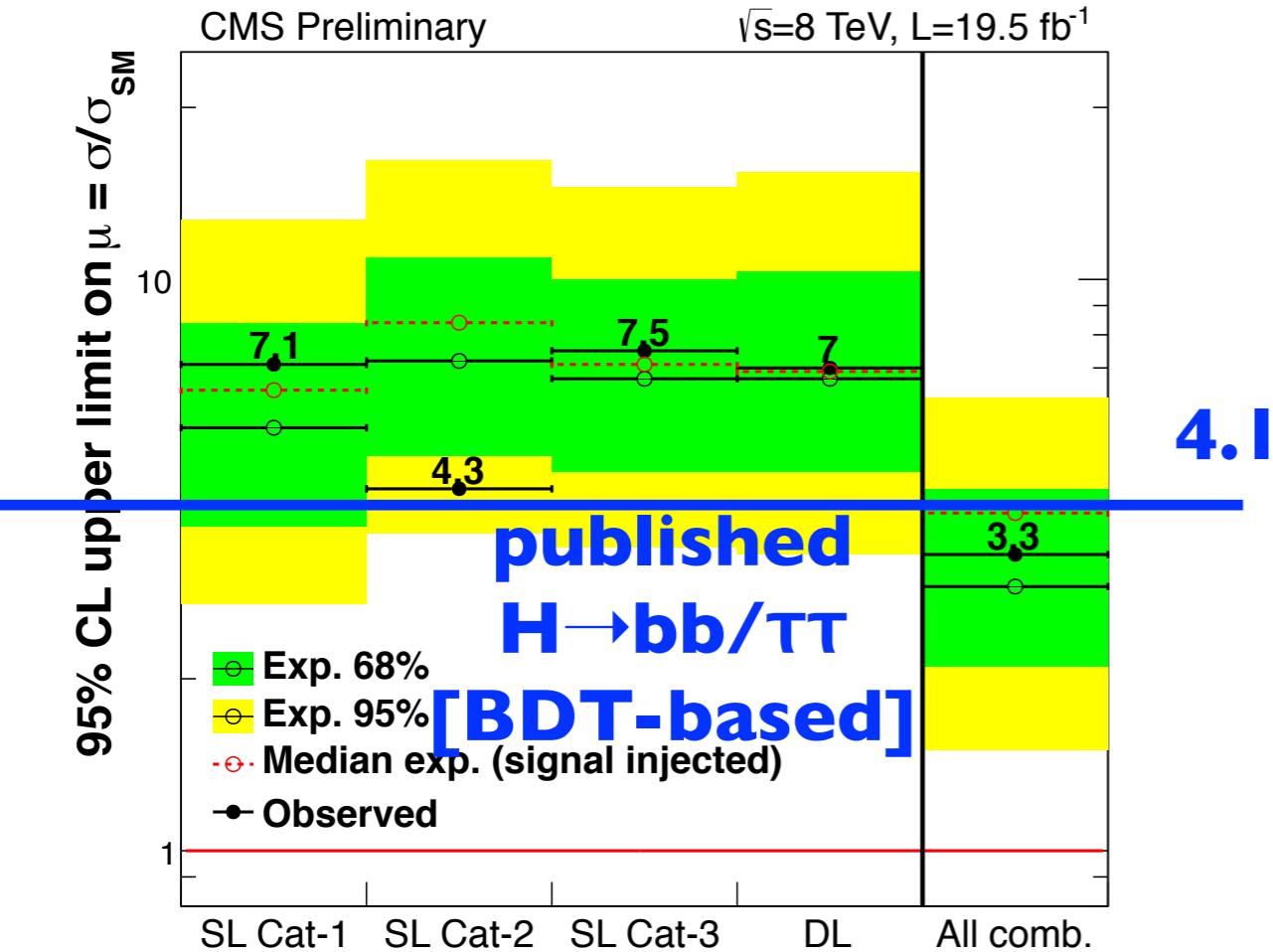


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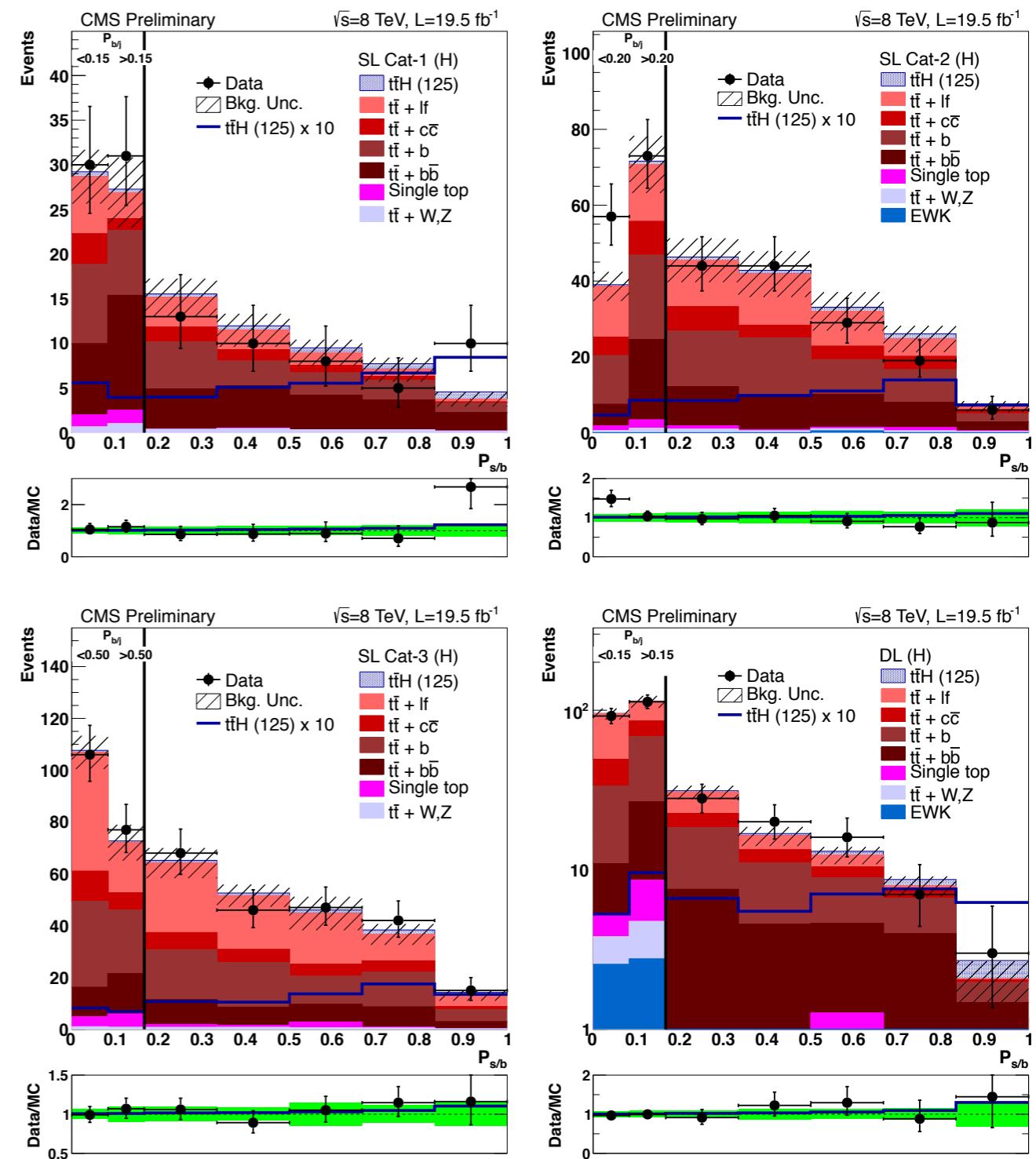
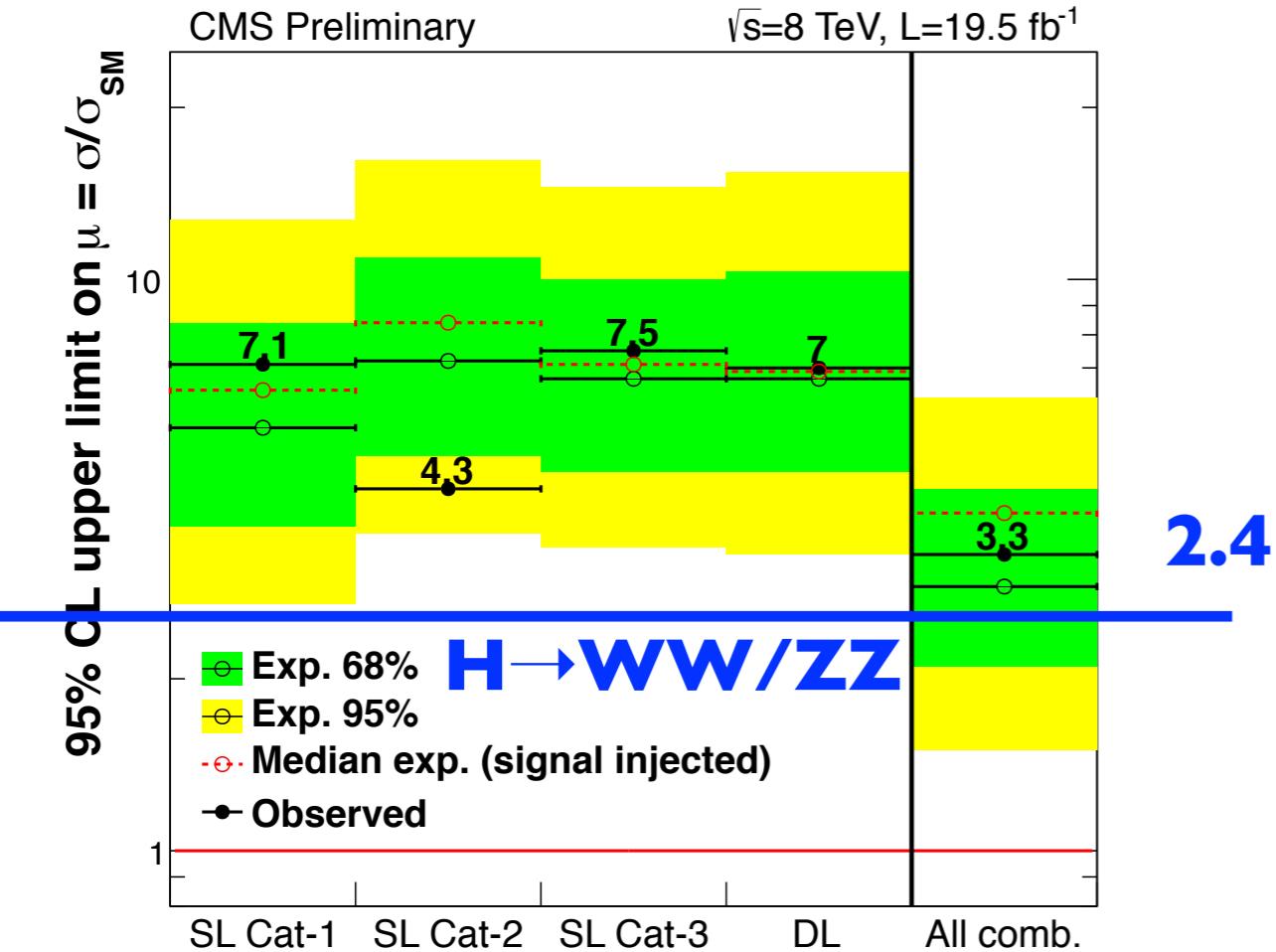
CMS-PAS-HIG-14-010

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Summary & outlook

- First ttH search based on the MEM at a collider
 - ▶ new algorithm developed and optimized for CMS
 - 20-30% better expected limit compared to previously published analysis
 - CPU & human-time sustainable
 - Run I: results in agreement with SM expectation ($\Delta\mu/\mu \sim 1.4$)
- The 13 TeV challenge is behind the corner
 - ▶ with more data, theoretical uncertainties will become relevant
 - use of NLO programs will make analysis more solid
 - definition proper sidebands for background estimation/calibration
 - $t\bar{t}Z(\rightarrow bb)$ as a standard candle?
 - ▶ further improvements:
 - spin-correlations as an extra handle (?)
 - inclusion of more event topologies (?)
 - inclusion of boosted top tagging variables (?)

Thanks for your attention

Back up

Motivation

Determination of top-quark Yukawa **coupling** (y_t) is a major goal

- ▶ gather direct evidence of Higgs coupling to up-type fermions
- ▶ implication on EWSB

Cross sections of top+Higgs channels can unravel value/sign of y_t

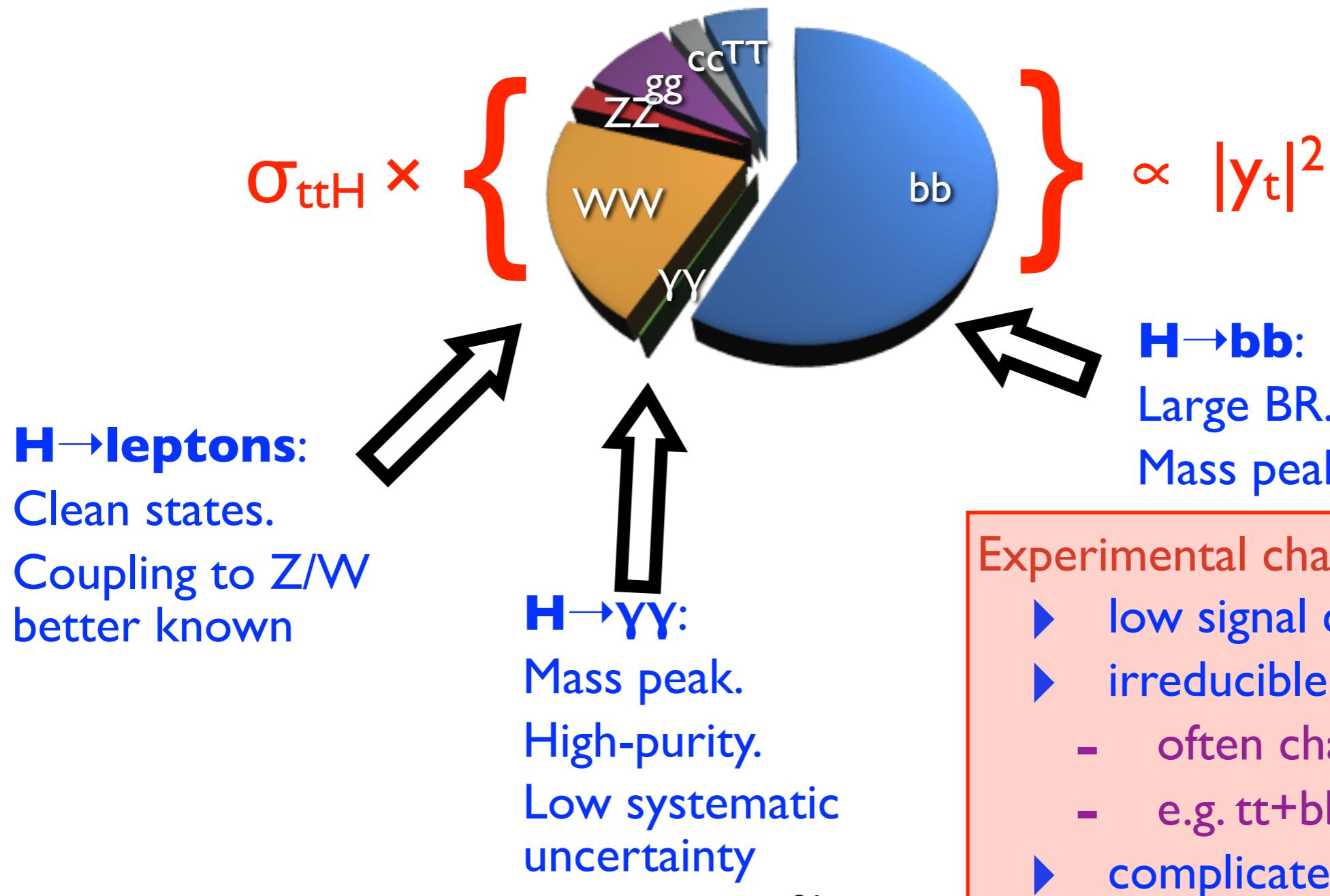
$$\sigma_{\text{ttH}} \times \left\{ \text{BR}(\text{H} \rightarrow \text{X}) \right\} \propto |y_t|^2$$

Motivation

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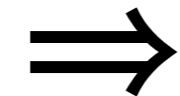
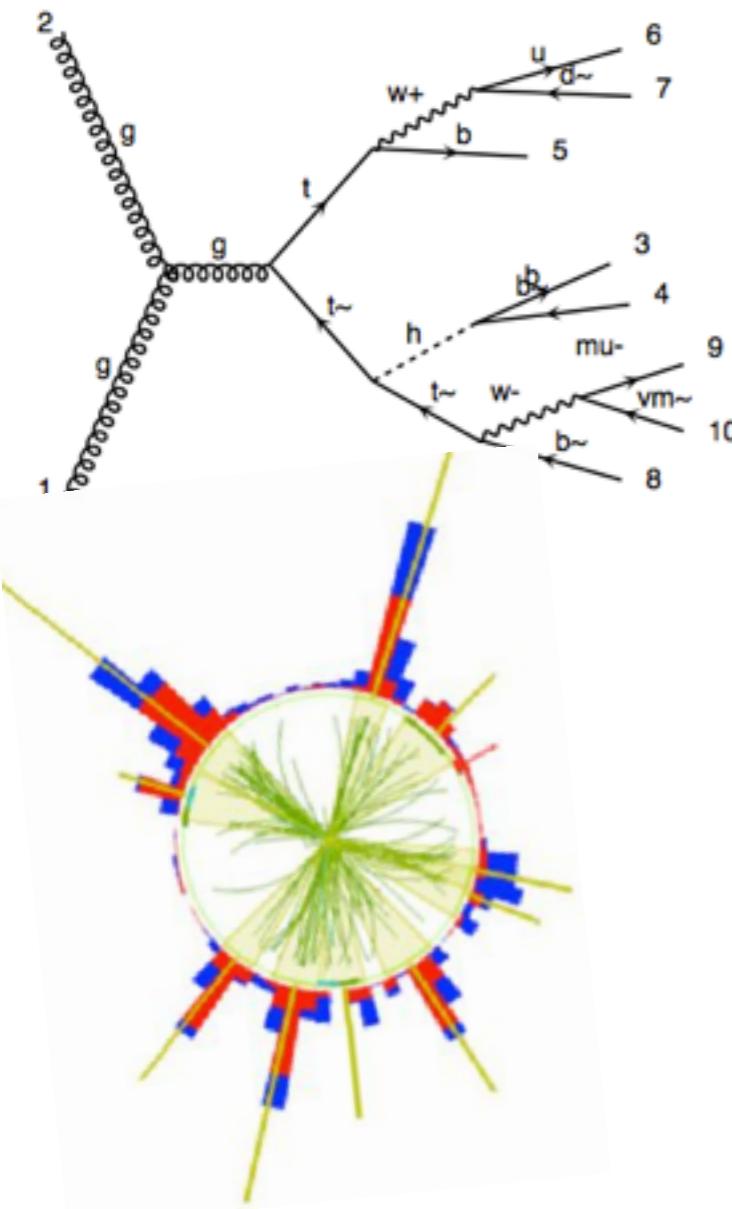
Cross sections of top+Higgs channels can unravel value/sign of y_t



Experimental challenge posed by:

- ▶ low signal cross section
- ▶ irreducible SM background
 - often challenging at NLO
 - e.g. tt+bb, tt+γγ, tt+Z/W
- ▶ complicated final states

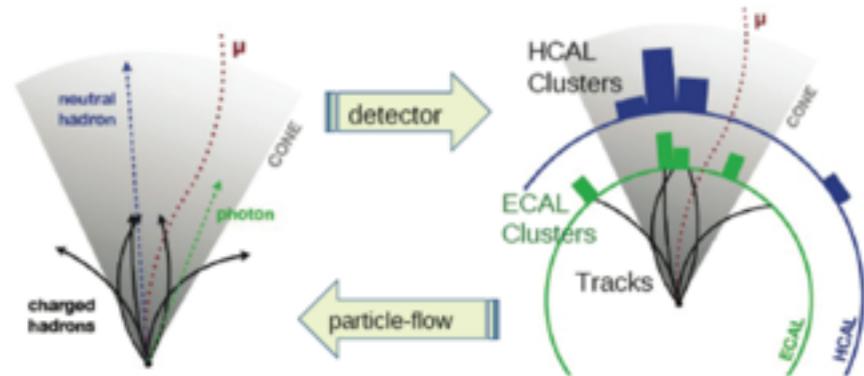
Building bricks



$\vec{X} = \text{(generated particles)}$
 $\mathcal{M}(\vec{X}) = \text{scattering amplitude}$



$\vec{Y} = \text{(reconstructed particles)}$



“metric” $W(\vec{Y}, \vec{X})$

Dimensional reduction

- Factorize integration over final-state particles via

$$d\Phi_n(P; p_1, \dots, p_n) = d\Phi_j(q; p_1, \dots, p_j) \times d\Phi_{n-j+1}(P; q, p_{j+1}, \dots, p_n) (2\pi)^3 dq^2$$

- Narrow-width approx: $\frac{1}{(t^2 - M_t^2) + \Gamma_t^2 M_t^2} \rightarrow \frac{1}{(M_t \Gamma_t)^2} \delta(t^2 - M_t^2)$

- Diff. decay amp. from MC: $|\mathcal{M}(t \rightarrow bqq')|^2 \propto \frac{1}{\Gamma_t} \frac{M_t}{|\vec{q}^*| |\vec{b}|} \frac{d\Gamma}{d\Omega_q^* d\Omega_b}$

- Assume lepton and jet direction perfectly measured

\Rightarrow

Diagram illustrating the decay chain: $t \rightarrow bqq' \rightarrow h \rightarrow bb \rightarrow b/\nu$.

Dimensional reduction equations:

$$d\Phi_{t_q} \propto \frac{2E_{b_q} E_q E_{q'}^2}{M_W^2 E_{w_q}} \frac{1}{\left| \frac{\beta_{w_q}}{\beta_{b_q}} \hat{e}_{w_q} \cdot \hat{e}_{b_q} - 1 \right|} dE_q$$

$$d\Phi_h \propto \frac{E_{b_2}}{\left| \frac{\beta_{b_1}}{\beta_{b_2}} \hat{e}_{b_1} \cdot \hat{e}_{b_2} - 1 \right|} dE_{b_1}$$

$$d\Phi_{t_\ell} \propto \frac{2E_{b_\ell} \hat{E}_\ell E_\nu^2}{M_W^2 E_{w_\ell}} \frac{1}{\left| \frac{\beta_{w_\ell}}{\beta_{b_\ell}} \hat{e}_{w_\ell} \cdot \hat{e}_{b_\ell} - 1 \right|} d\Omega_\nu$$

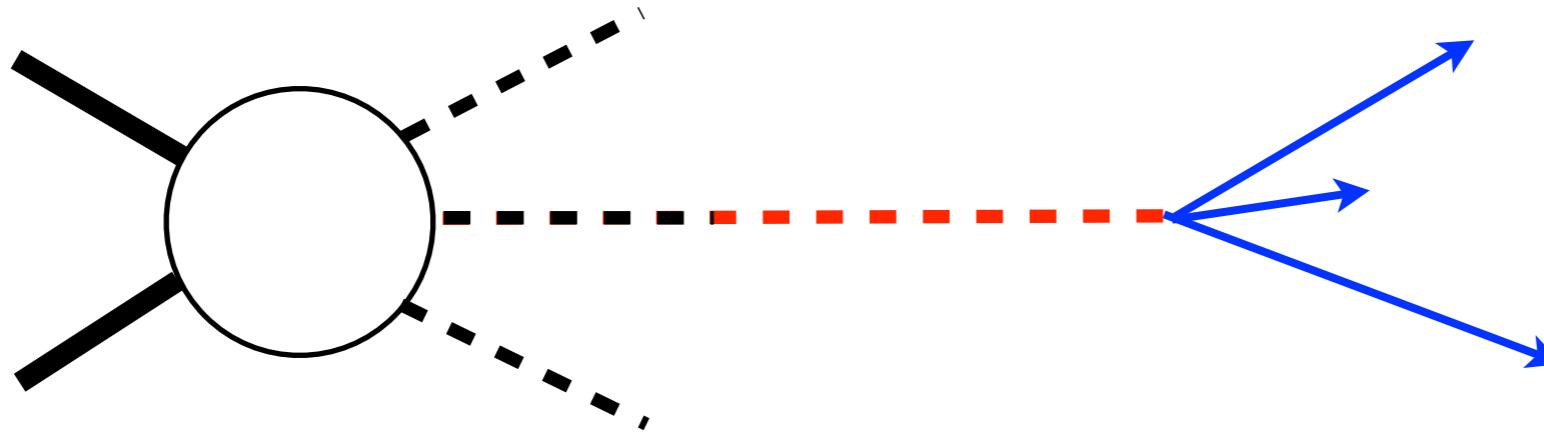
The Matrix Element Method (MEM)

Given event j , find its probability $w_j(\cdot | H_i)$ under hypothesized process H_i

- ▶ w_j is a function of detector-level observables \mathbf{Y} (technically, a differential prob. density on \mathbf{Y})
- ▶ can depend parametrically on unknown model parameters λ
- ▶ is the relative weight with which a MC generator of process H_i would generate ev. j
 - for $N \gg 1$ events, $(\sum^N w_j / N) \rightarrow 1$
- ▶ N.B.: a HEP process H_i is known (\Leftrightarrow can be simulated) fully differentially only to some approximation (typically not better than NLO). Conversely, nature is “to all orders”:
 - a certain approximation is implicit in the method
 - this does not invalidate the method *per se'*, rather makes it not optimal

The algorithm: basic principles

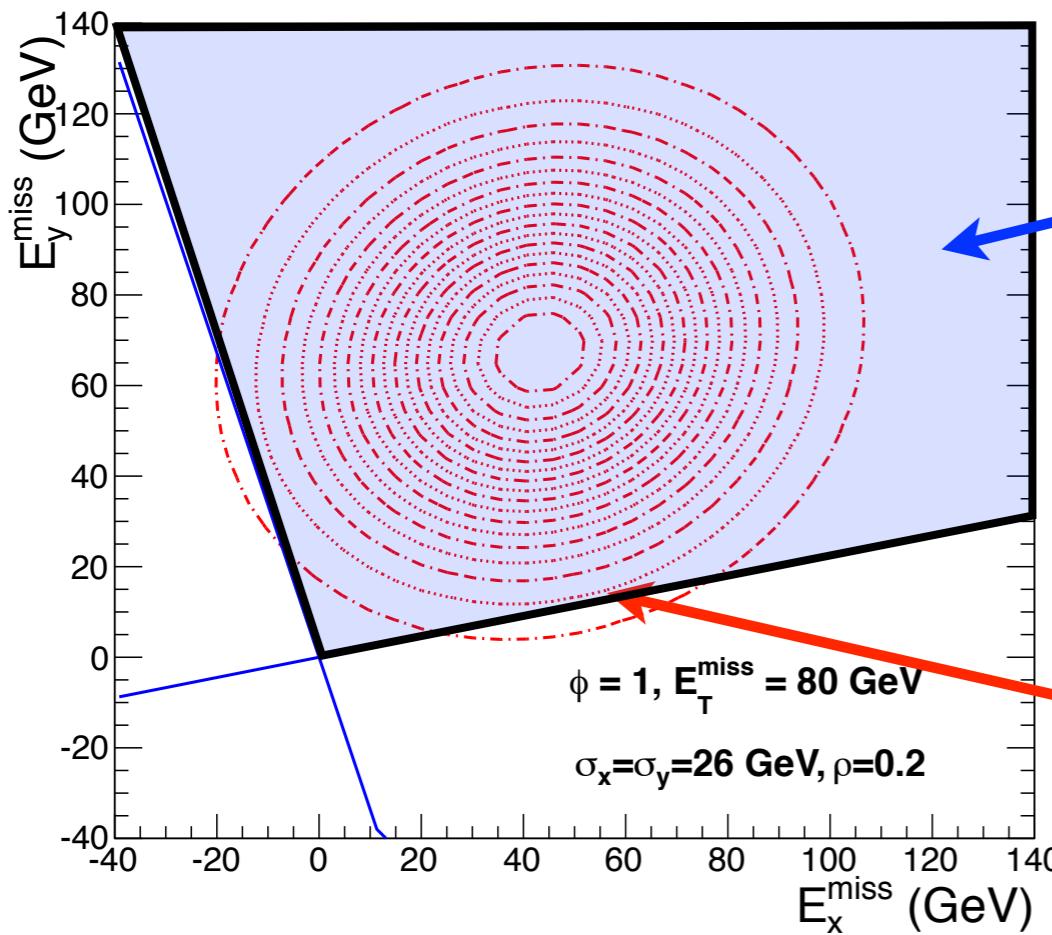
- Factorize the reaction $p\bar{p} \rightarrow t\bar{t} + (bb) \rightarrow \Omega$ as a 3-steps process:
 - ▶ $gg \rightarrow 3$ on-shell intermediate particles $\propto |\mathcal{M}(g g \rightarrow t\bar{t} H)|^2$
 - ▶ intermediate particles propagate: $\propto [(q^2 - M^2)^2 - M^2 \Gamma^2]^{-1}$
 - ▶ intermediate particles decay $\propto \Gamma^{-1} d\Gamma/d\Omega$



- This way:
 - ▶ no need to evaluate CPU-intensive $2 \rightarrow 8$ amplitudes [only $2 \rightarrow 3(4)$]
 - ▶ spin-correlations and polarizations neglected
 - ▶ cross-check with MadWeight

P_T balance

- Event-by-event constraint to the measured recoil $\rho = -\sum \mathbf{p}_T^{\text{vis}} - \mathbf{E}_T^{\text{miss}}$ via transfer function
 - for each phase-space point, boost so that $\mathbf{P}_T = \mathbf{0}$, and evaluate $|\mathcal{M}|^2$
- N.B.: at present, we instead constrain ν 's \mathbf{p}_T to $\mathbf{E}_T^{\text{miss}}$ and the quark energy to jet energy
 - not optimal because $\mathbf{E}_T^{\text{miss}}$ correlated w/ jet energy.



ν 's direction integrated over blue plane w/ prior:

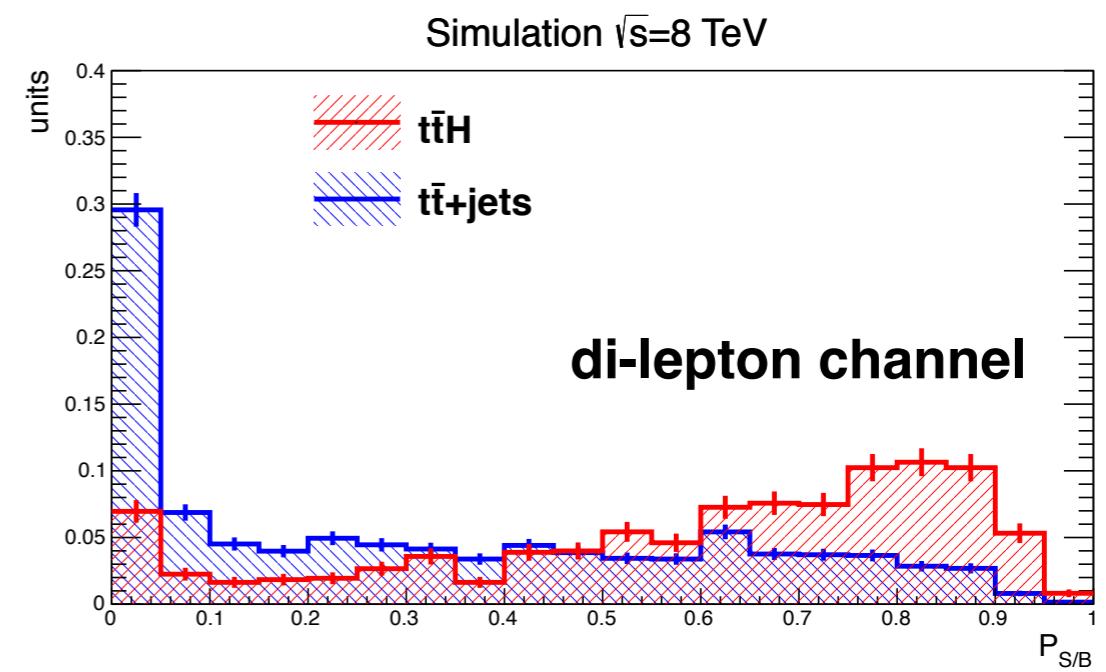
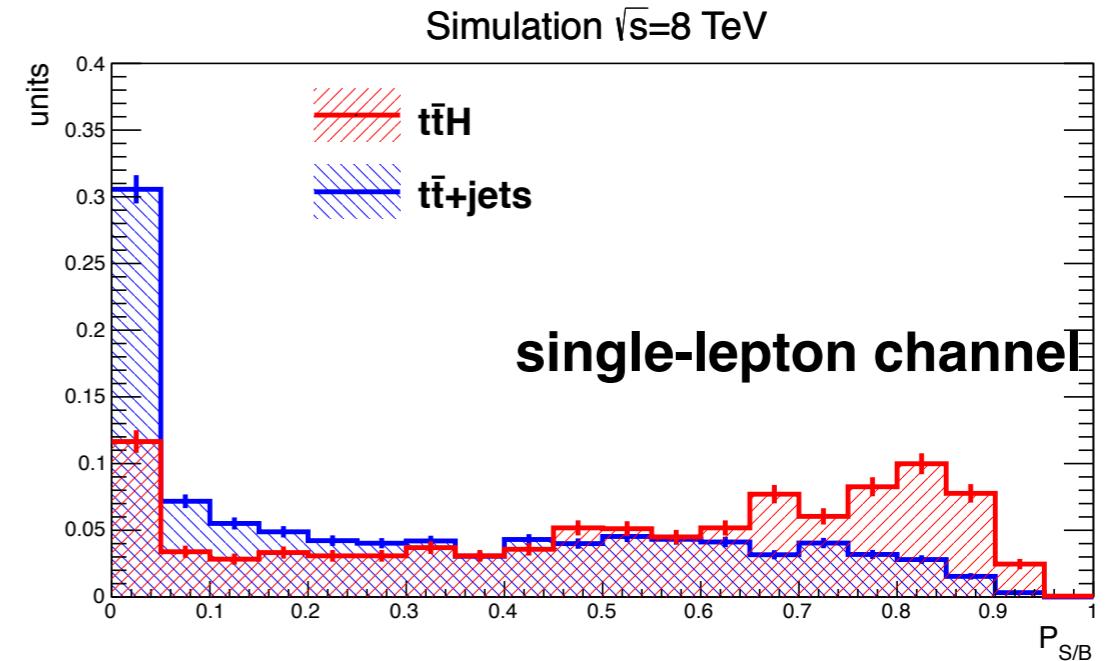
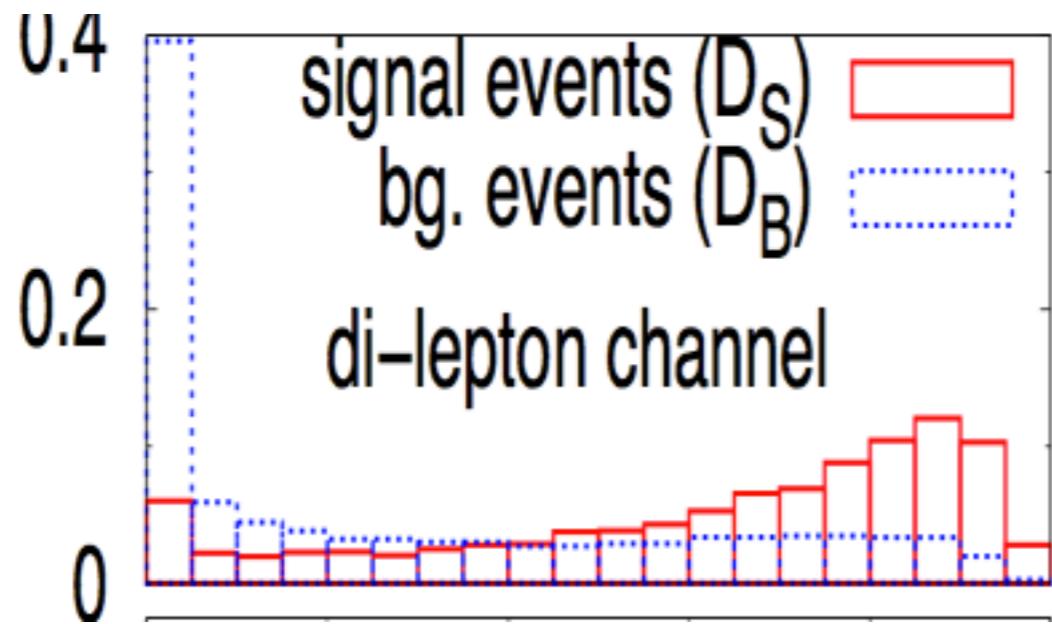
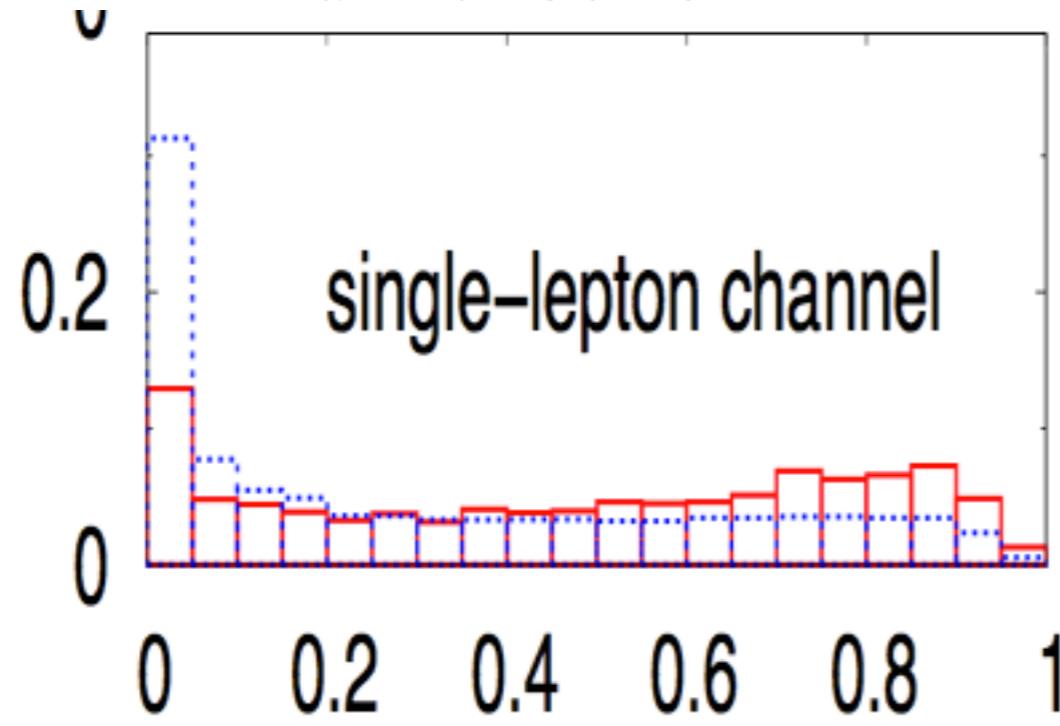
$$\frac{1}{2\pi\sqrt{|V_{x,y}|}} \exp\left\{-\frac{1}{2}(\vec{E}_T^{\text{miss}} - \sum \vec{v}_T)^T V_{x,y}^{-1} (\vec{E}_T^{\text{miss}} - \sum \vec{v}_T)\right\}$$

95% CL isocontour

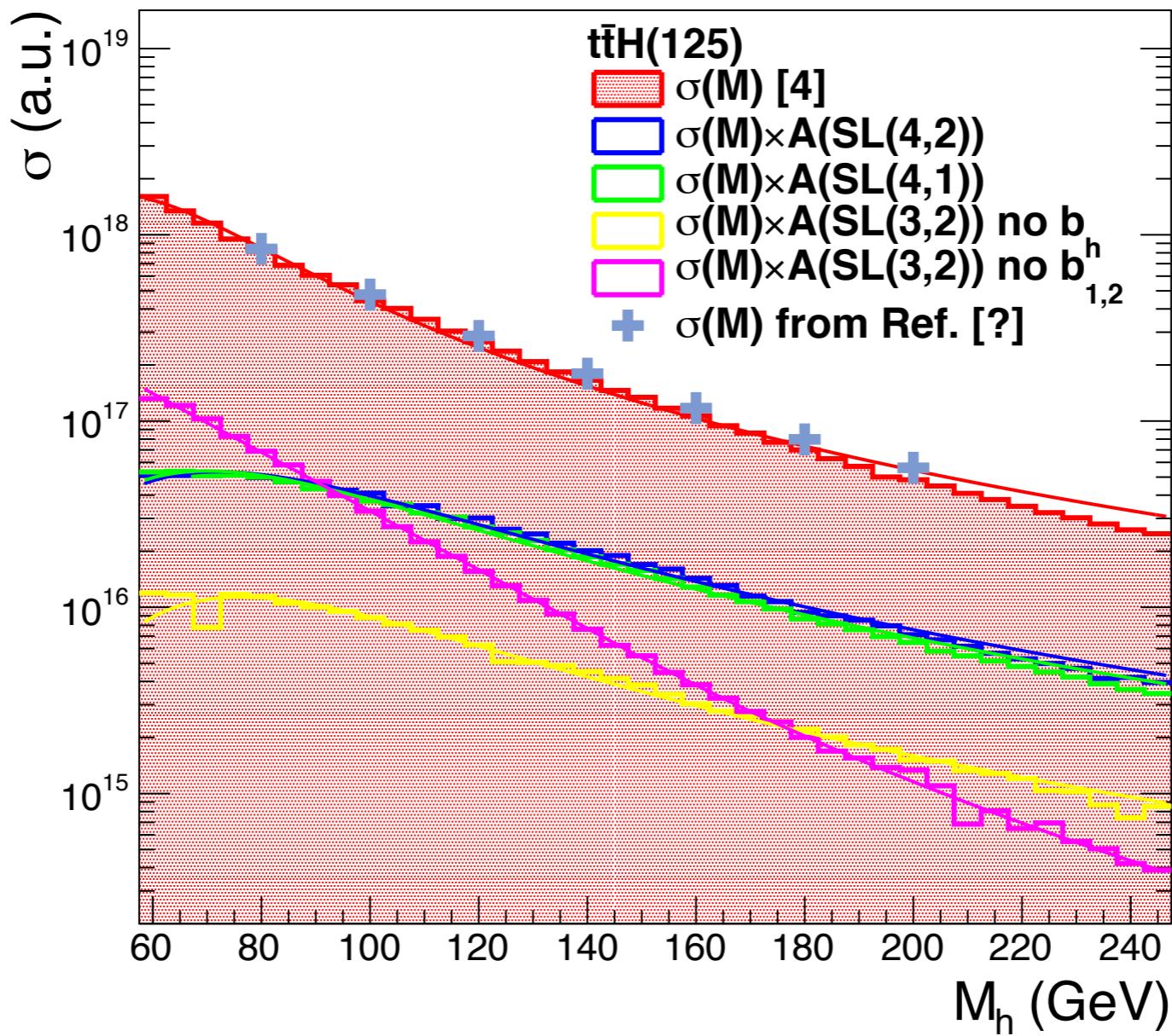
$V_{x,y} = E_T^{\text{miss}}$ cov. matrix

Comparing with MadWeight

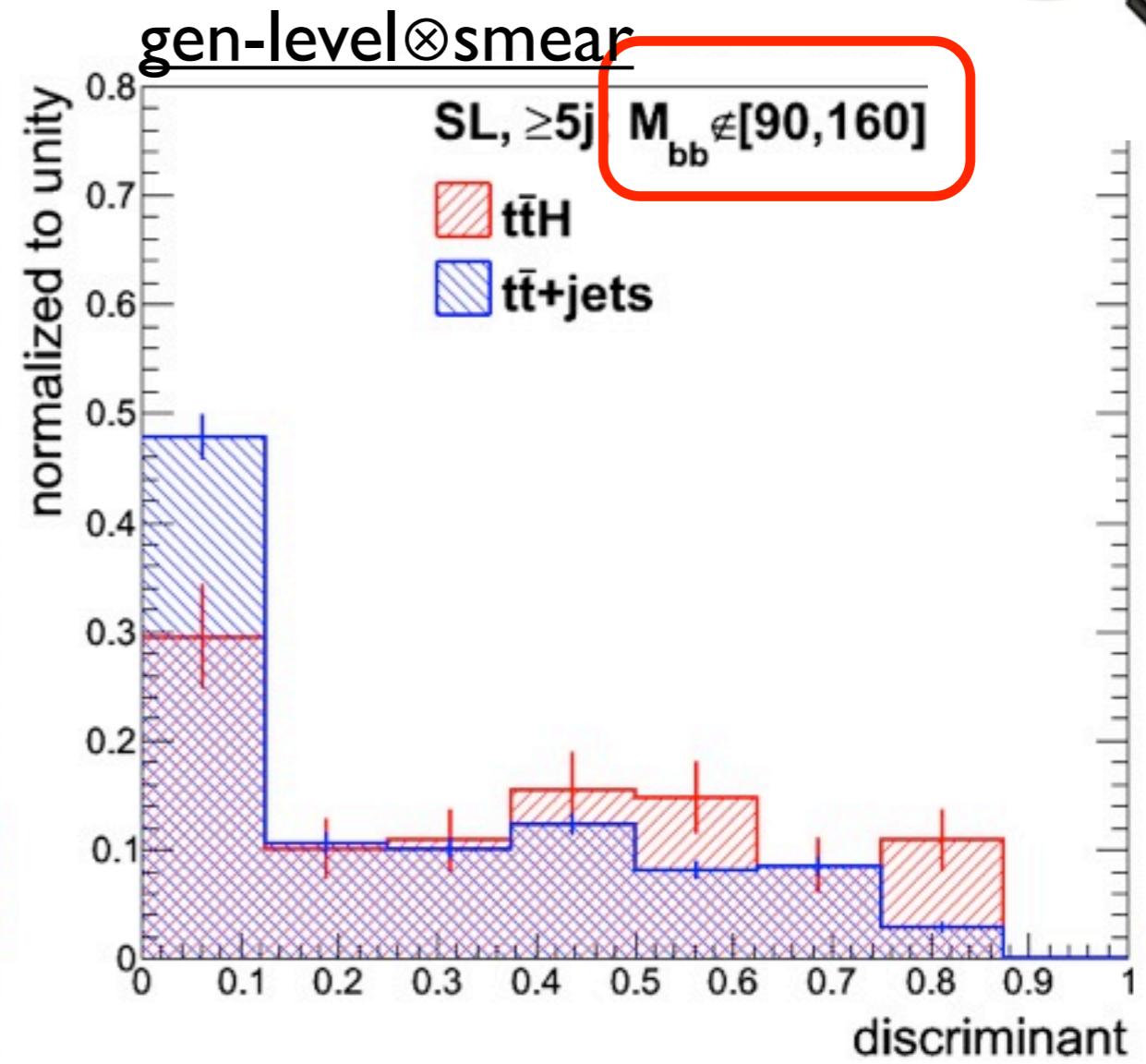
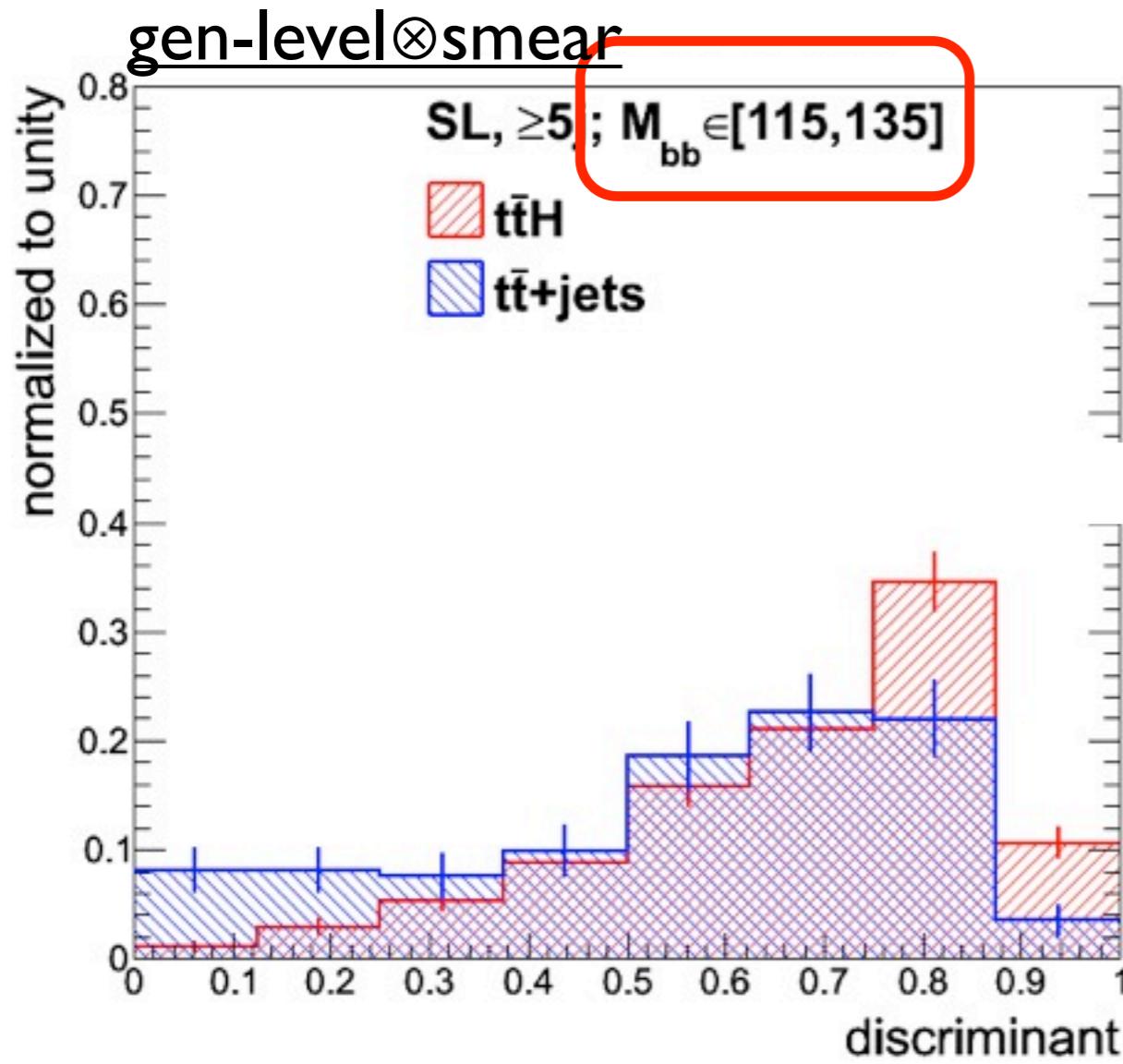
arXiv:1304.6414



Cross-section



Digression I: the Higgs mass

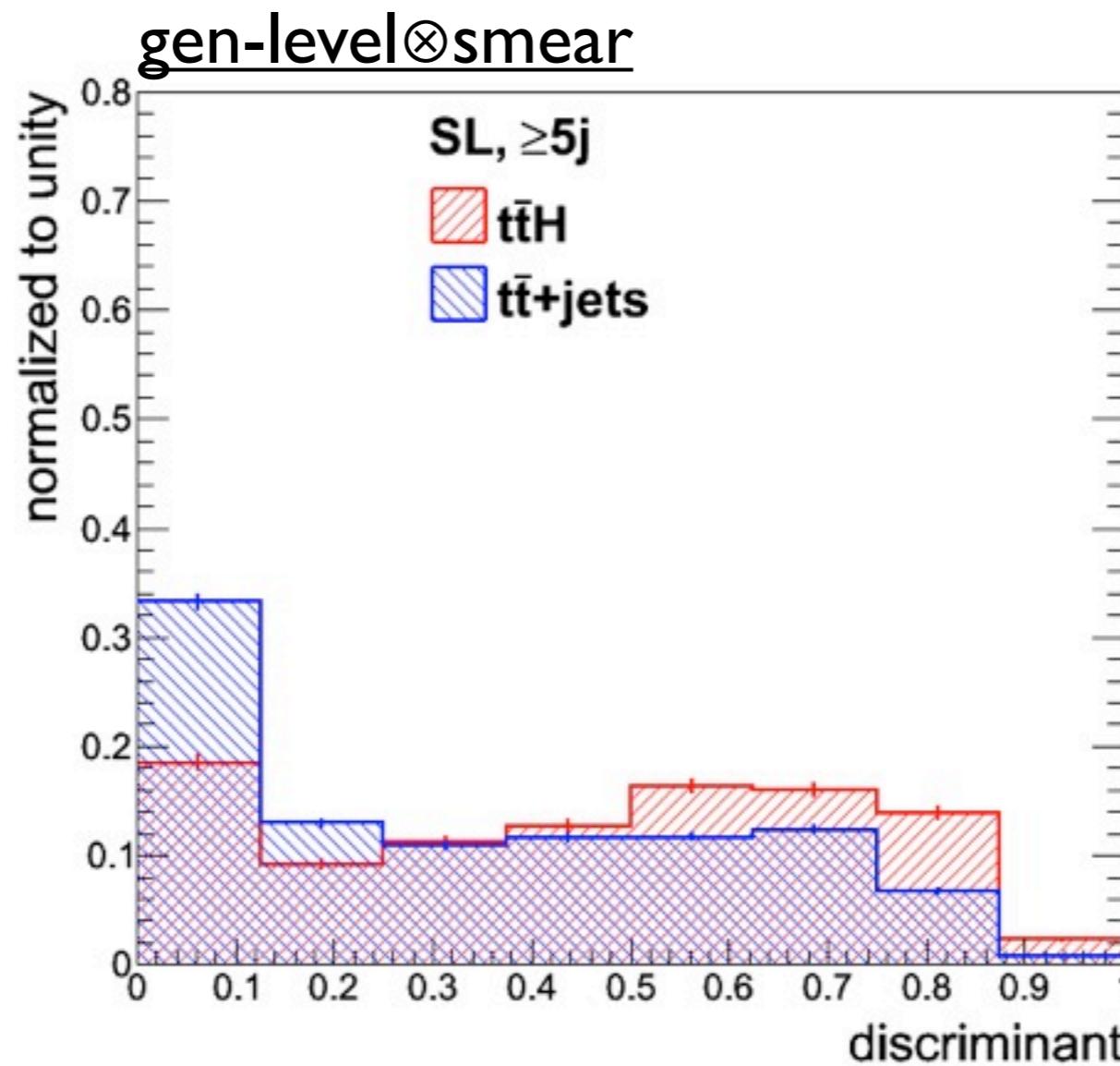


- $ttbb$ events w/ $M(bb) \approx 125$ indeed look like tH !
 - ▶ but not identical \Leftrightarrow the ME is sensitive also to the other variables
- ttH events w/ $M(bb) \neq 125$ (e.g. poor resolution) undistinguishable from $ttbb$

Digression II: wrong hypothesis

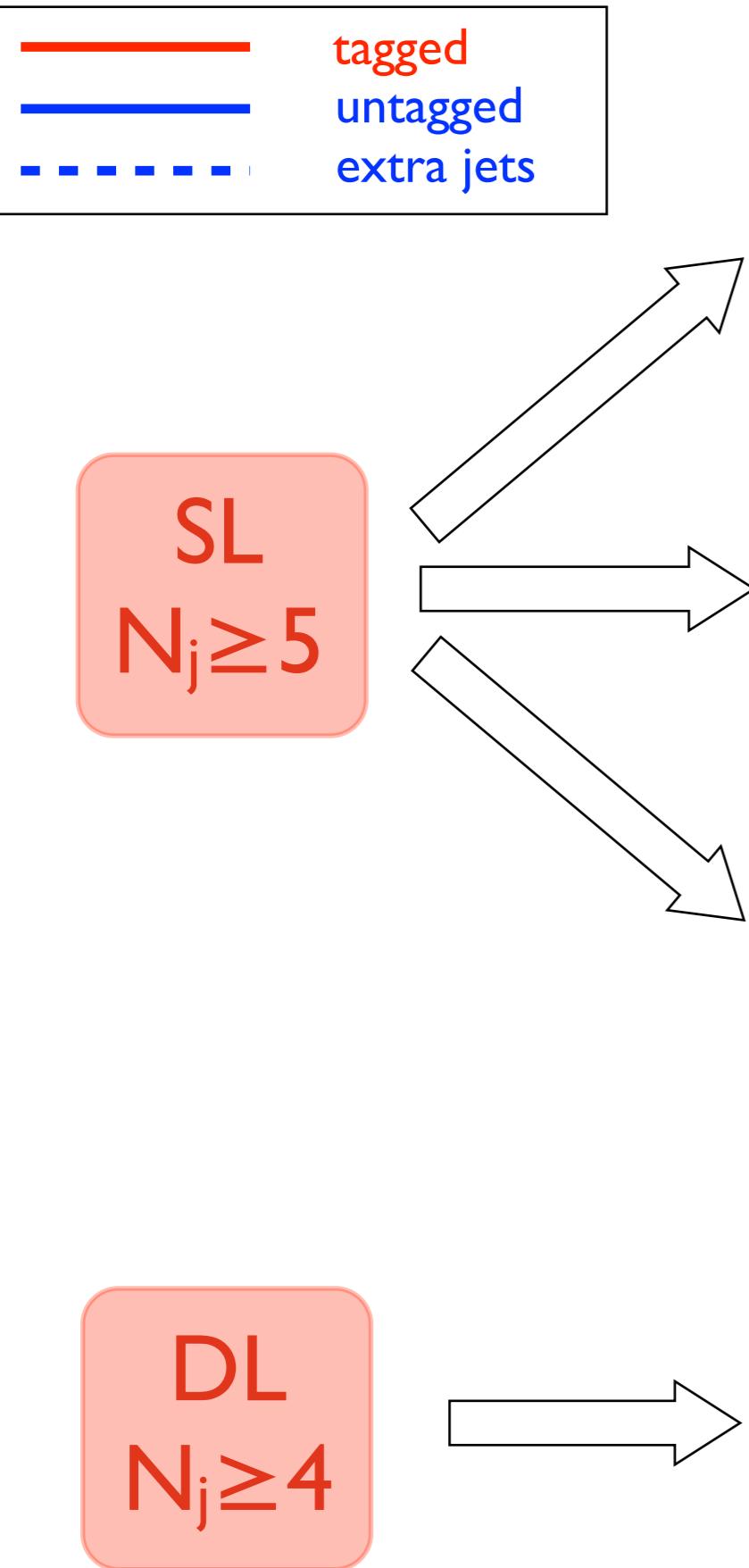


- If the event does not fulfill the *tested* ME hypo, the weight is broadly distributed
 - ▶ yet, $t\bar{t}H$ remains slightly more “signal-like”

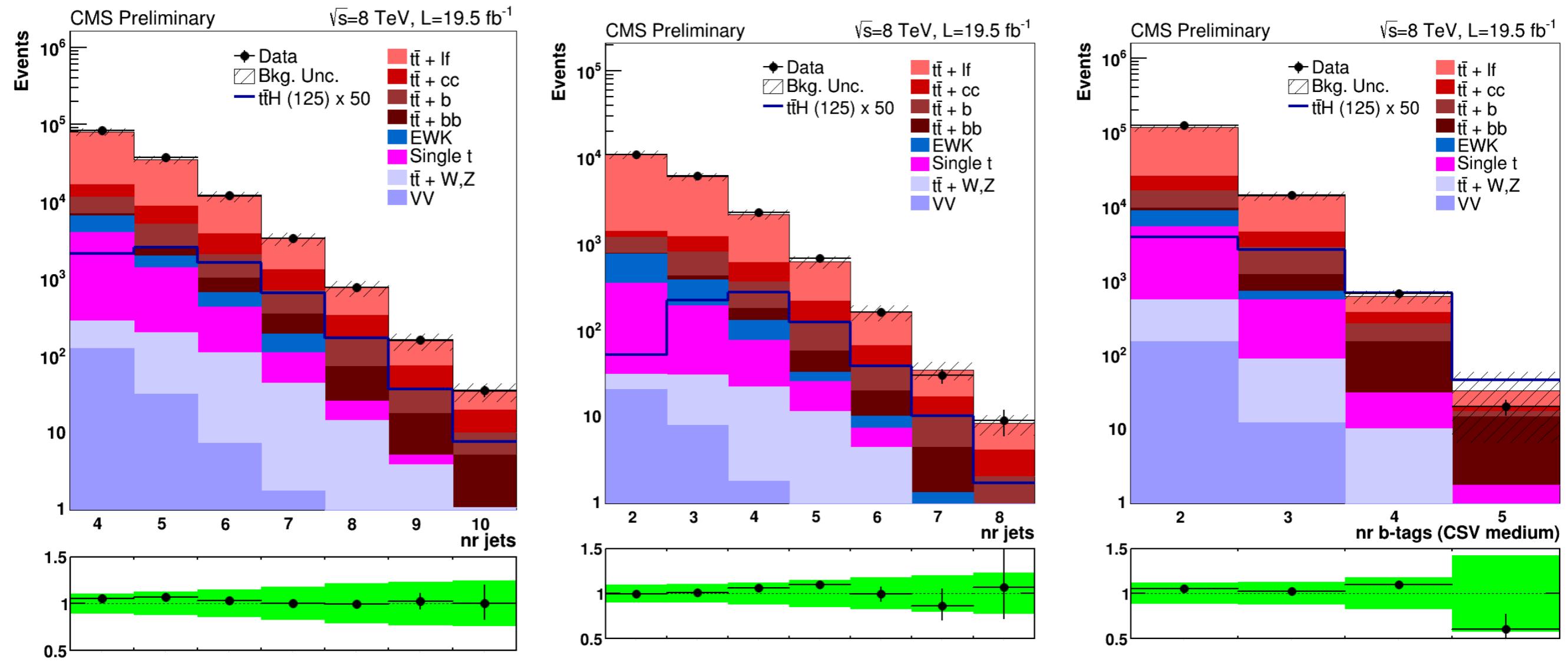


Categories

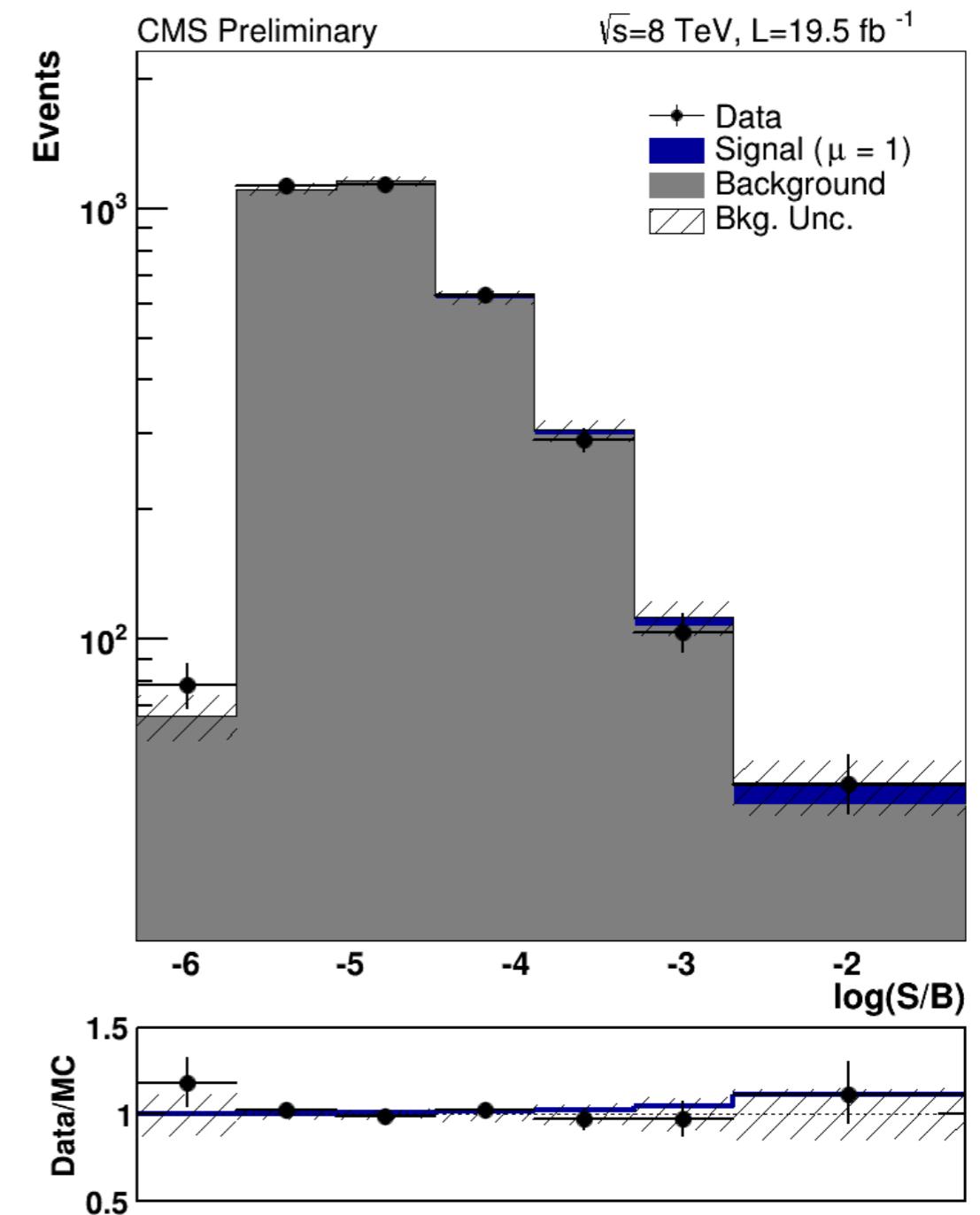
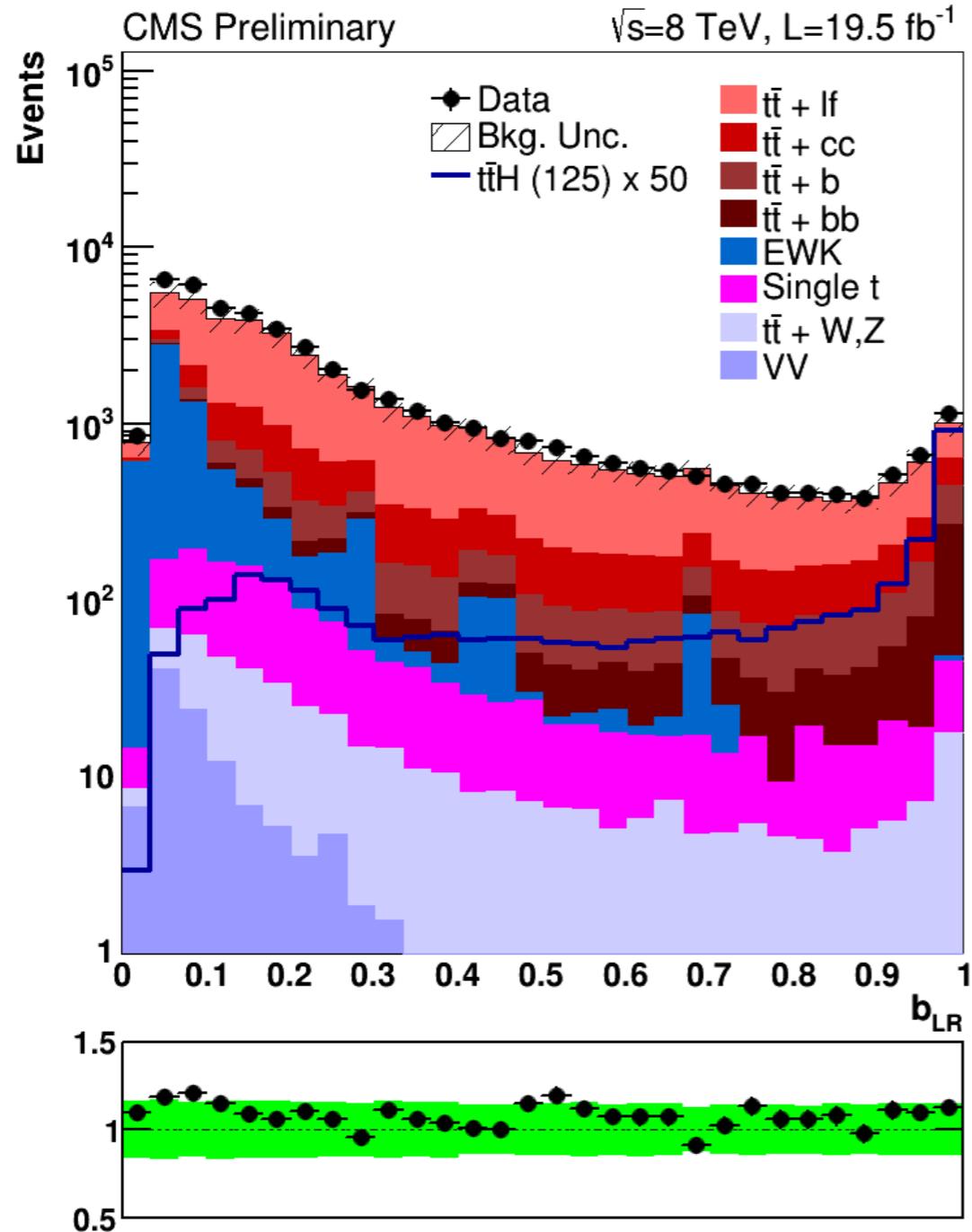
assignment based on jet permutation w/ largest btag



Control plots



Control plots



$$\mathcal{L}_{bbbb}(\xi_1, \dots, \xi_6) \equiv \sum_{\{i_1, \dots, i_6\}} f_b(\xi_{i_1}) \cdot f_b(\xi_{i_2}) \cdot f_b(\xi_{i_3}) \cdot f_b(\xi_{i_4}) \cdot f_u(\xi_{i_5}) \cdot f_u(\xi_{i_6})$$