Neutrino & Cosmology

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Precision is a matter of prejudice

We know a lot about neutrino properties from lab experiments.

We would like to know more exploiting their impact on cosmological and astrophysical observables.

Precision Cosmology: precise observations which fit the standard model extremely well.

But: as soon as we move away from our comfortable standard?

Robust vs. weak predictions: which is the case for neutrino properties ?

Summary

- Are there neutrinos in the universe?
 How many of them? (the long tale of N_{eff}) ×
 Neutrino mass: universe better than lab's ?
 Oscillations and neutrino asymmetries
 Sterile states ? ×
 Any relation with dark matter ? ××
- Majorana or Dirac particles ? XXX

BBN and CMB probe the light particle content at different epochs: both require relativistic species in addition to photons

$$\rho_R = \rho_{\gamma} \left(1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{eff} \right)$$

For BBN: $N_{eff} = 3$ is a good fit (see later) BBN requires electron neutrinos!

CMB fixing the angular scale of acoustic peaks and z_{eq} , a larger N_{eff} gives a higher expansion speed, a shorter age of the universe T at recombination.

Diffusion length $\approx \sqrt{T}$ Sound horizon $\approx T$

 $N_{eff} = 3$ is a good fit (see later)

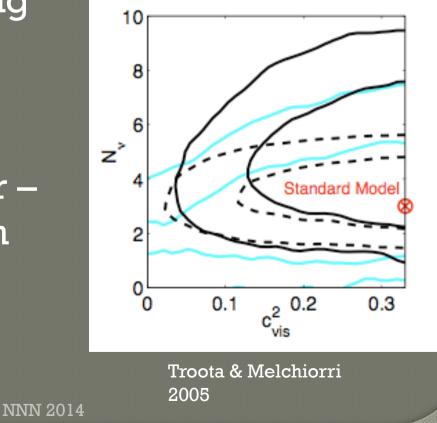
Perturbation effects:

 gravitational feedback of neutrino free streaming damping

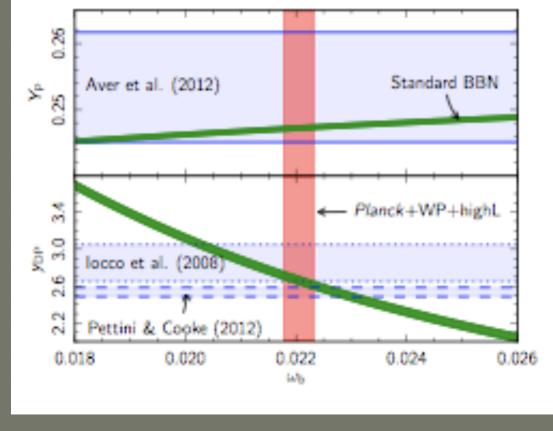
 anisotropic stress

contributions

c_{vis} :velocity/metric shear – anisotropic stress relation (Hu 1998)



CMB and BBN are quite consistent



Ade et al. 2013 (Planck XVI)

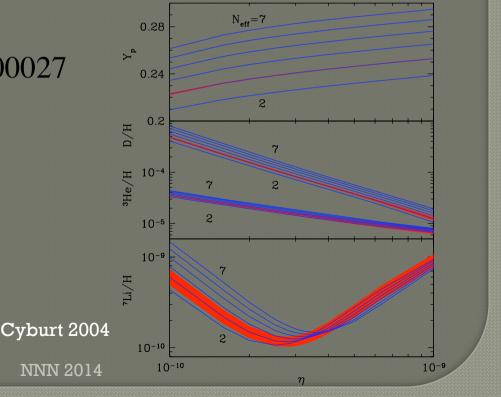
both ⁴He mass fraction Y_p and ²H/H are increasing functions of N_{eff} : change of expansion rate v_e distribution crucial in weak rates baryon density basically fixed by CMB! (but still ²H/H can varies a lot)

$$\Omega_b h^2 = 0.02207 \pm 0.00027$$

crucial inputs:

experimental values

nuclear rates



⁴He still affected by a remarkable systematic uncertainty Recent re-analysis

$$\begin{split} Y_p &= 0.2565 \pm 0.0010(stat) \pm 0.0050(syst) \\ Y_p &= 0.2561 \pm 0.0108 \\ Y_p &= 0.2573 \pm 0.033 \\ Y_p &= 0.2465 \pm 0.0097 \\ Y_p &\leq 0.2631 \ 95\% \ \text{C.L.} \end{split}$$

Izotov & Thuan 2010 Aver et al. 2010 Aver etl. 2012 Aver et al. 2013 Mangano & Serpico 2011

²H/H is presently quite well determined, thanks to new very metal poor system measurements (Cooke et al. 2013)

 $^{2}H/H = (2.53 \pm 0.04) \cdot 10^{-5}$

Several claims, spanning from "Evidence for extra neutrinos"

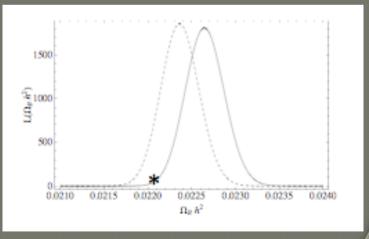
to

"No room for extra neutrinos"

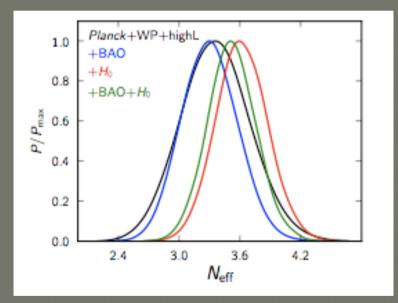
Conservative estimate: $N_{eff} < 4$ (still !) One example: for Planck baryon density a higher deuterium

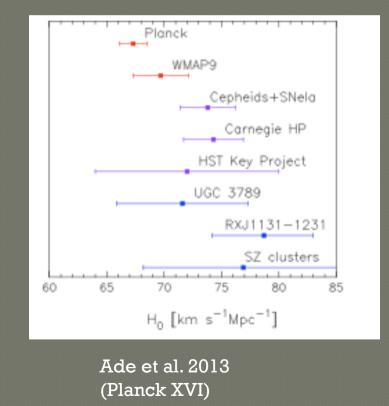
 $^{2}H/H = (2.65 \pm 0.07) \cdot 10^{-5}$

 N_{eff} smaller than 3 (2.7)? Maybe, or a larger S-factor for d(p, γ)³He, as in the theoretical estimate of Marcucci et al. (2005)



News from Planck: a narrower 95 % C.L. range for N_{eff} , but still Inconclusive. H_0 problem:





3.4±0.7 3.3±0.5 3.6±0.5

Neutrino mass: universe better than lab's ?

Mass bounds

Laboratory is still missing! 2 eV for $\nu_{\rm e}$

Katrin wil tell us more (when?)

Cosmology blind to neutrino mass till recent times.

CMB:

For the expected mass range the main effect is around the first acoustic peak due to the early integrated Sachs-Wolfe (ISW) effect;

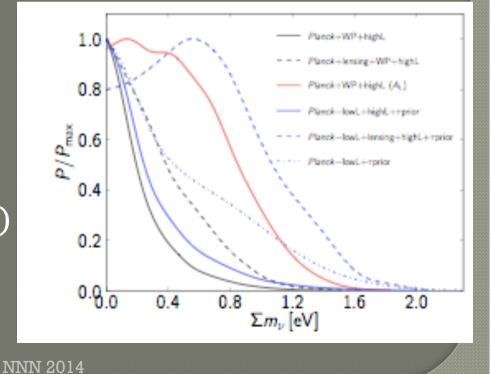
Planck: gravitational lensing. Increasing neutrino mass, increases the expansion rate at z >1 and so suppresses clustering on scales smaller than the horizon size at the nonrelativistic transition (Kaplinghat et al. 2003; Lesgourgues et al. 2006). Suppression of the CMB lensing potential.

Neutrino mass: universe better than lab's ?

Total neutrino mass also affects the angular-diameter distance to last scattering, and can be constrained through the angular scale of the first acoustic peak. Degenerate with Ω_{Λ} (and so the derived H0)

Including BAO constraint is much tighter:

 $\Sigma m_v < 0.98 \text{ eV} (CMB)$ $\Sigma m_v < 0.32 \text{ eV} (CMB + BAO)$



Early times:

$$f_a = \frac{1}{e^{p/T - \xi_a} + 1}$$
 $f_{\overline{a}} = \frac{1}{e^{p/T + \xi_a} + 1}$

Kinetic and chemical equilibrium

MeV scale (set by G_F and Δm^2 's) : • freezing of weak interaction processes • ν distributions mixed up, dependi

 $f_a slightly$ distortedN_{eff} = 3.046

• $\boldsymbol{\nu}$ distributions mixed up, depending on mixing angles

density matrix formalism ρ_{ab}

 ρ_{aa} occupation number ρ_{ab} a≠b mixing

$$\frac{d}{dt}\rho = \frac{1}{i} \left[\Omega_{vac} + \Omega_{matter}, \rho \right] + C$$

 Ω_{vac} vacuum oscillations: M²/2p

 Ω_{matter} matter term: $2^{1/2} G_F \Delta n_i + 8 2^{1/2} G_F p T_0^0 / 3M^2_{W,Z}$ C: collisional integral (loss of coherence and distribution re-shuffling) Stodolski 1987

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Raffelt ad Sigl 1993

When oscillations matter:

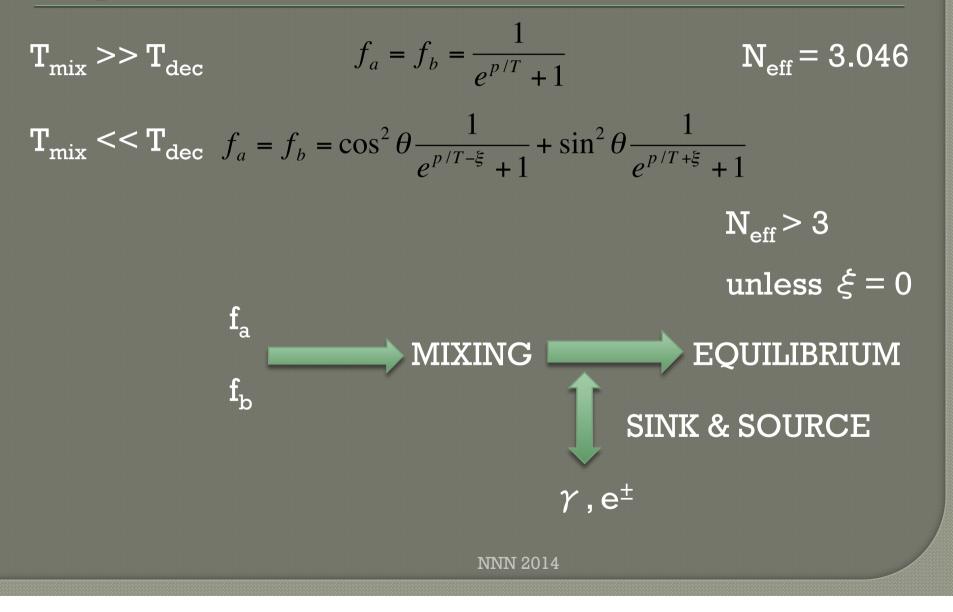
Lepton asymmetries expected quite small in (standard) leptogenesis

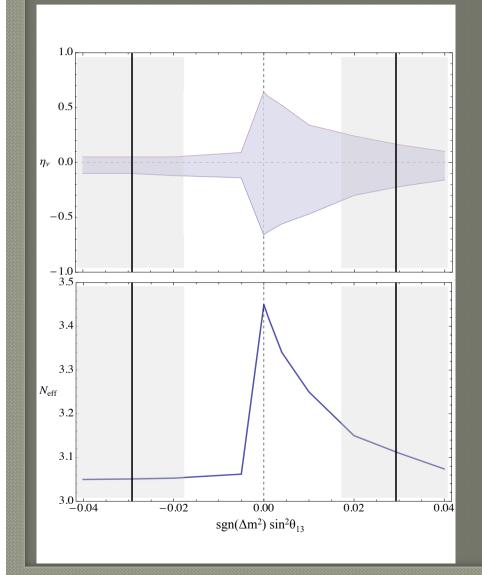
$$\eta_a = \frac{n_a - n_{\bar{a}}}{n_{\gamma}} = \frac{1}{12\xi(3)} \left(\pi^2 \xi_a + \xi_a^3\right) \approx \eta_B = 6 \times 10^{-10}$$

unless leptogenesis takes place well below the EW breaking scale (T) = 2T

$$\exp\!\left(-M_W(T)/g^2T\right) << 1$$

The value of θ_{13} is crucial (and to a minor extent the mass 0.6 $\theta_{13}=0$ hierarchy) NH 0.4 $\eta_{\nu_{\mu,\tau}}$ IH 0.2 Veutrino asymmetries 0 -0.2 -0.4 -0.6 $\eta_{v_{\alpha}}$ -0.8 Pastor et al 2011 -1 2 10 7 5 3 GM et al 2012 T_{γ} (MeV) NNN 2014





the bounds: scanning all asymmetries compatible with BBN

 $N_{eff} < 3.2$

-0.2 (-0.1) $\leq \eta_{\nu} \leq 0.15$ (0.05)

GM et al 2012

N_{eff} ≤ 3.2 still compatible with slightly degenerate neutrinos
 N_{eff} ≥ 3.2 some extra "dark" radiation required or higly non-thermal neutrino distribution, or both

After Planck I results still inconclusive !

Hints for sterile neutrino states from long(short) standing anomalies

LSND, MiniBoone Reactor anomaly Gallium anomaly

 $m_v \approx eV$, $\sin^2 \theta_{as} \approx 10^{-2}$

Too many sterile neutrinos in the early universe, produced via oscillations

Unless there is a fine tuning, the typical outcome is either too few or too many (and too heavy !)

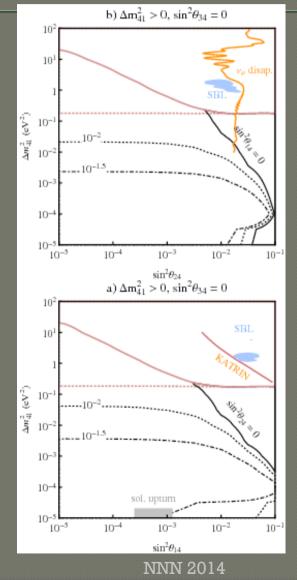
- 1. The standard case
- 2. Large lepton asymmetries
- "secret" "sterile" interactions (unlikely² for Ockam)

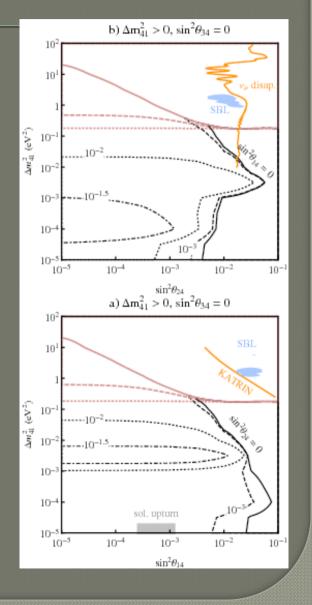
Planck analysis (Planck XVI 2013)

 $N_{eff} < 3.91$ $m_{s} < 0.59 eV$

 $N_{eff} < 3.80$ $m_s < 0.42 \text{ eV}$ (including BAO)

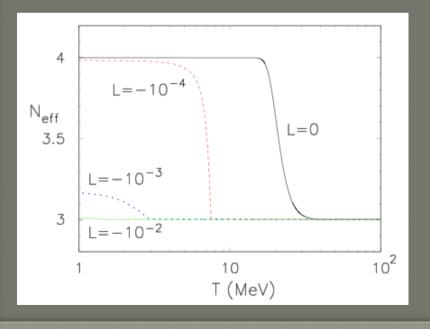
The standard case (Mirizzi et al 2013)

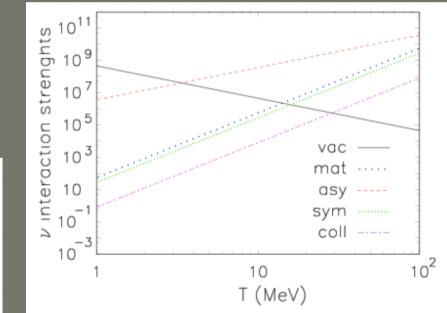




Lepton asymmetry suppresses sterile production

 $V = \sqrt{2} G_F L_v$





 $L_v = 10^{-4}$ Mirizzi et al. 2012

Large sterile self-interactions suppress sterile production due to large potential

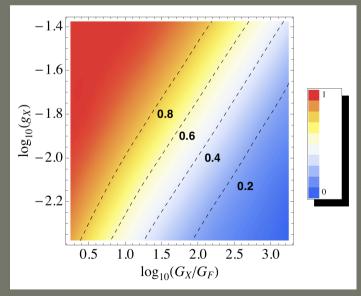
$$V_s = -\sqrt{2}G_X \frac{8\langle p \rangle \rho_s}{3M_X^2}$$

(Hannestad et al 2013)

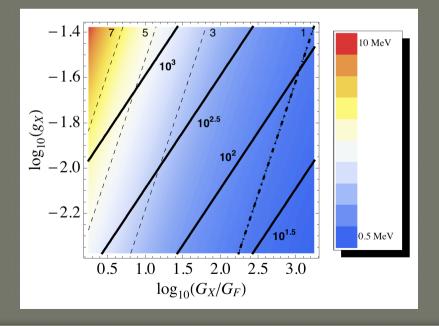
 G_X larger than Fermi constant. OK for N_{eff} smaller than 1.

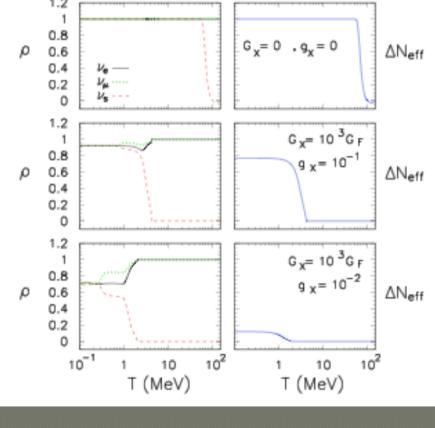
$$G_X = \frac{\sqrt{2}g_X^2}{8M_X^2}$$

Saviano et al 2014

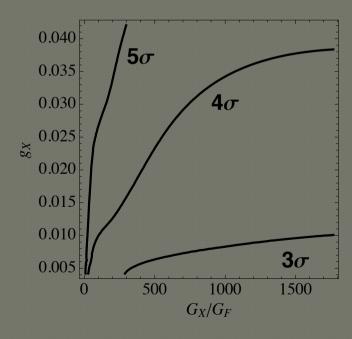


Bounds from BBN ! If resonance takes place at MeV scale the v_e distribution gets distorted.





Main constrain from ²H/H (with Planck baryon number). Increasing $\Omega_{\rm b}$ helps but choosing it at 3 σ from Planck best fit)



Majorana or Dirac particles ?

Can we distinguish Majorana or Dirac neutrinos from cosmology ?

Relativistic regime: NO X Structure formation: NO X Leptogenesis: YES 44 Direct detection: YES ! (really demanding) 4

Majorana or Dirac particles ?

Neutrino capture on ³H (PTOLEMY R&D, Katrin too low ³H mass)

 $\nu_{e} + {}^{3}H \rightarrow e + {}^{3}He$

Weinberg 1962 Cocco et al 2007

Dirac: only the left-helicity neutrinos
 & anti-neutrinos cannot be captured.
 Majorana: both left- and right-helicity neutrinos,
 Capture rate doubled

PTOLEMY (100g ³H) :

4 events yr ⁻¹ Dirac 8 events yr ⁻¹ Majorana Lunardini et al. 2014

Conclusions

What we would like to know about neutrinos?

- mass 🧹
- Majorana or Dirac 🗡
- Other neutrinos (sterile, mirror,...) 🧹

Cosmological data are precise today (% level)

&

Standard cosmological model is extremely on shape !!

Beware degeneracies ...

...and epicycles !!