

Experimental results on SUSY searches with top from the LHC



TOP2014:

7th International Workshop on Top-Quark Physics, Sep-Oct 2014, Cannes, France.

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on behalf of the ATLAS and CMS Collaborations

Prepared with the help from Henning Flacher.



Outline

- Supersymmetry searches with a top
→ setting the stage
- Third generation squark searches
- Constraints from top measurements

Supersymmetry

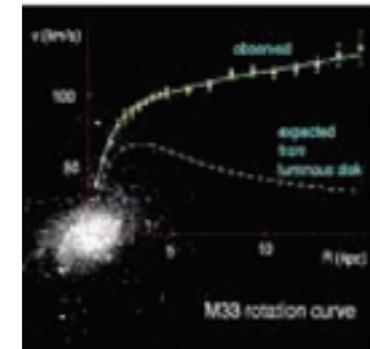
Supersymmetry (SUSY) adds a new fundamental symmetry relating fermions and bosons → more than doubles the particle spectrum w.r.t. the Standard Model.

Solves the gauge hierarchy problem

Fermion and boson loops contribute with different signs to the Higgs radiative corrections.



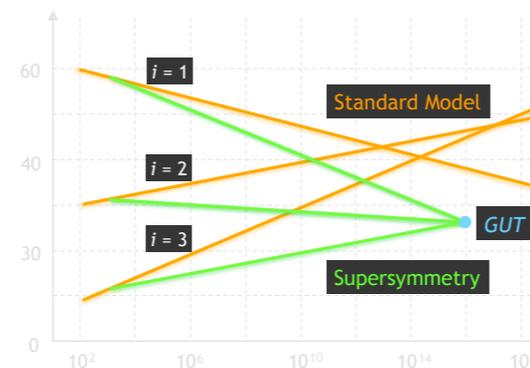
SUSY with R-parity conservation predicts a suitable dark matter candidate



Grand unification of forces

Predicts an elementary Higgs scalar ...

- with a beautiful SM-like limit,
- with a mass below 135 GeV (in the MSSM)

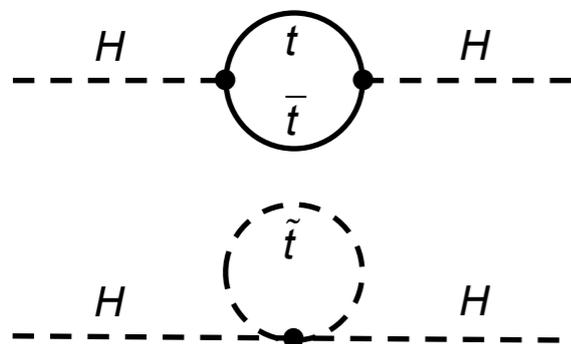


SUSY widely considered to remain the chief amongst beyond-SM proposals. However, most simplistic versions under stress.

Natural SUSY

What are the minimal requirements for 'natural' (low level of fine-tuning) SUSY?

Recall: **Top loop** is the most important contribution in Higgs virtual corrections:



While inclusive LHC searches set $O(\text{TeV})$ limits mostly on 1st and 2nd generation squarks and gluinos, which are less relevant in Higgs corrections.

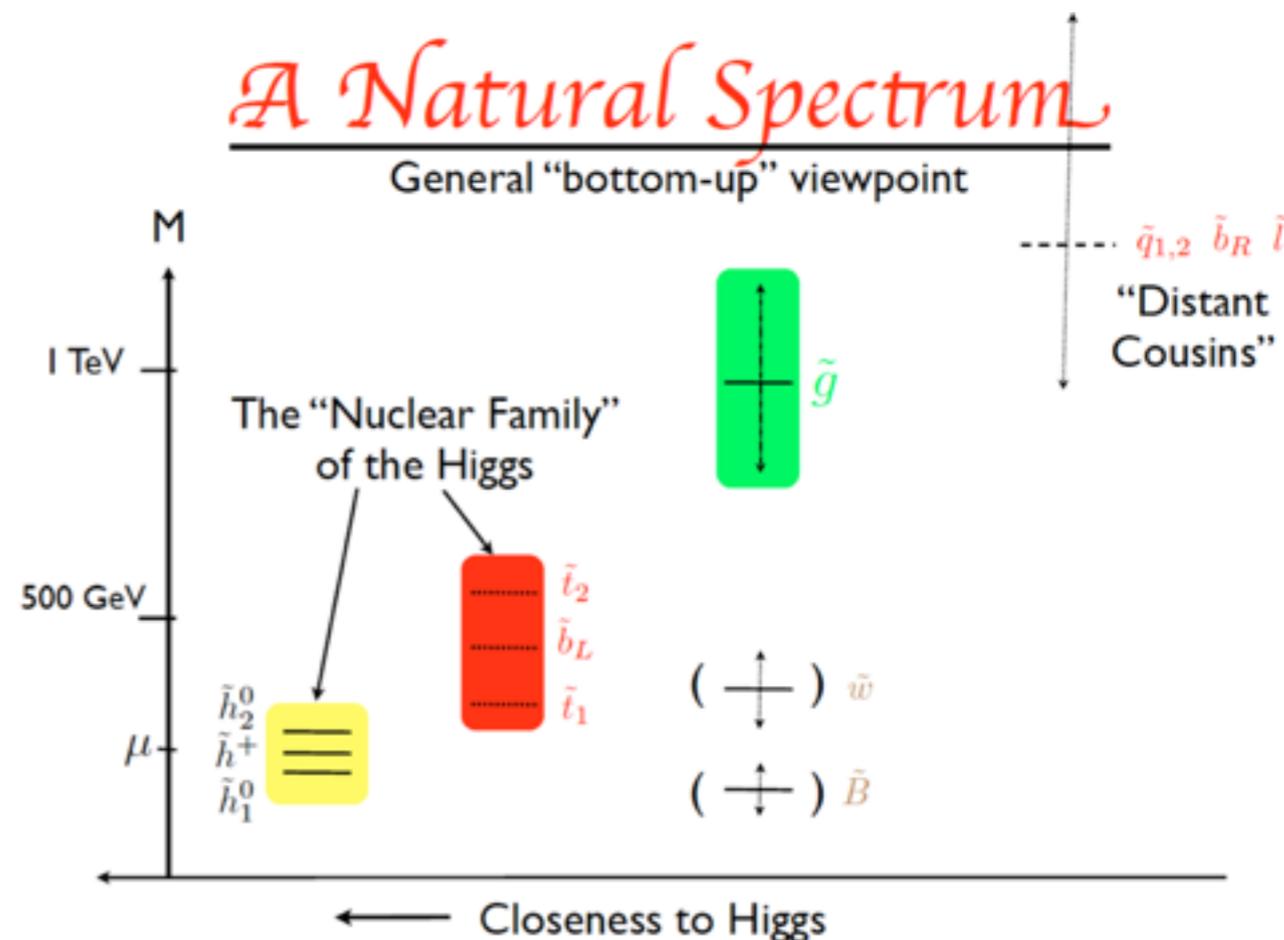


Illustration taken from: L. Hall, LBNL workshop, Oct-2011

For **natural SUSY** need:

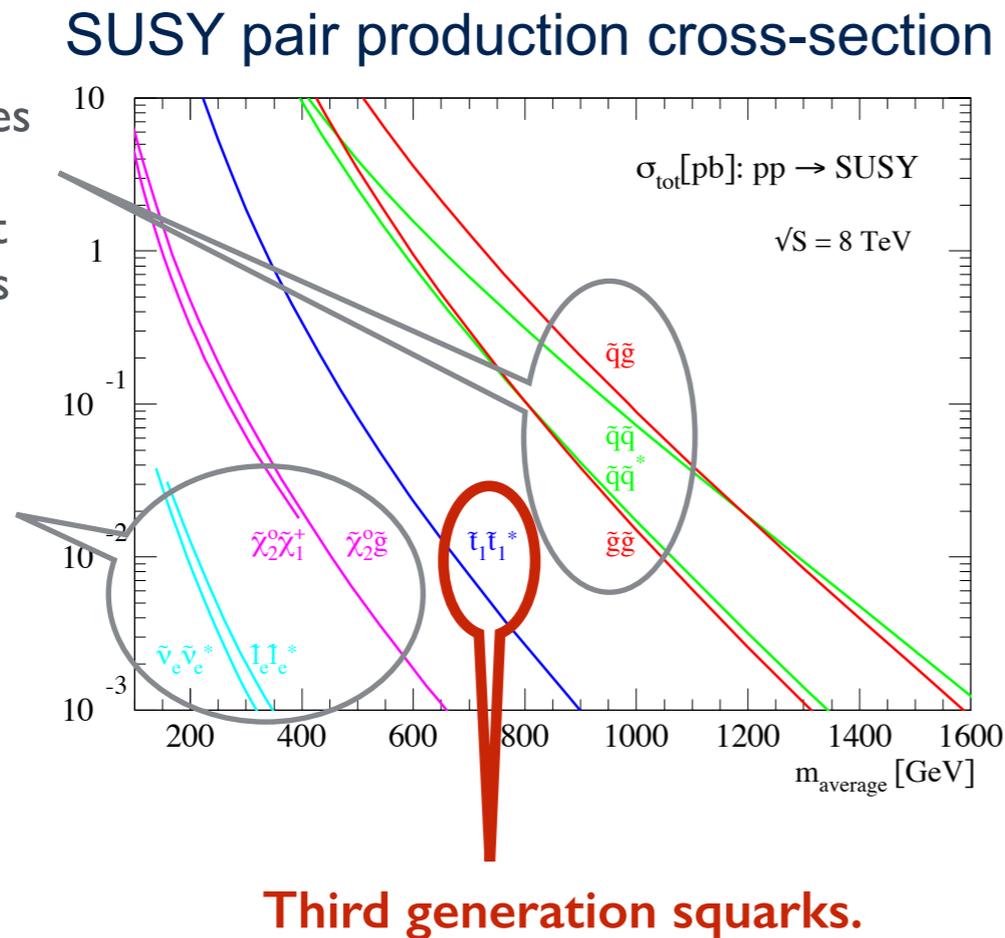
- ▶ light higgsinos
- ▶ light top and bottom squarks (stop, sbottom, respectively)
- ▶ not-too-heavy gluinos

➔ strong physics case for **third generation squarks**

3rd generation searches at the LHC

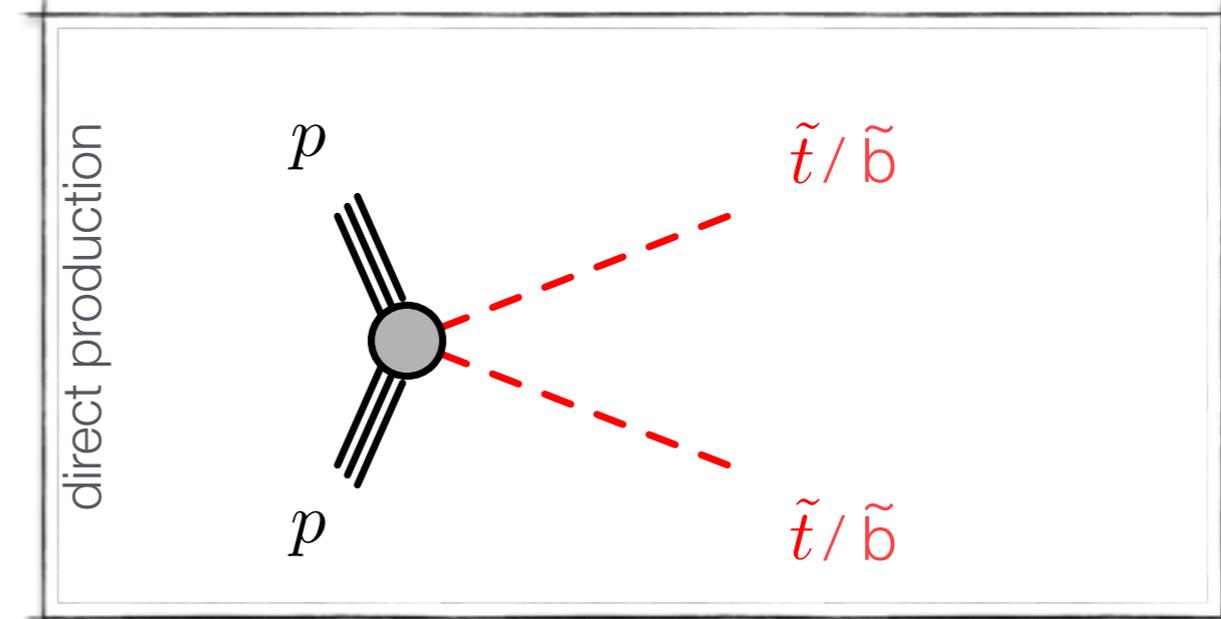
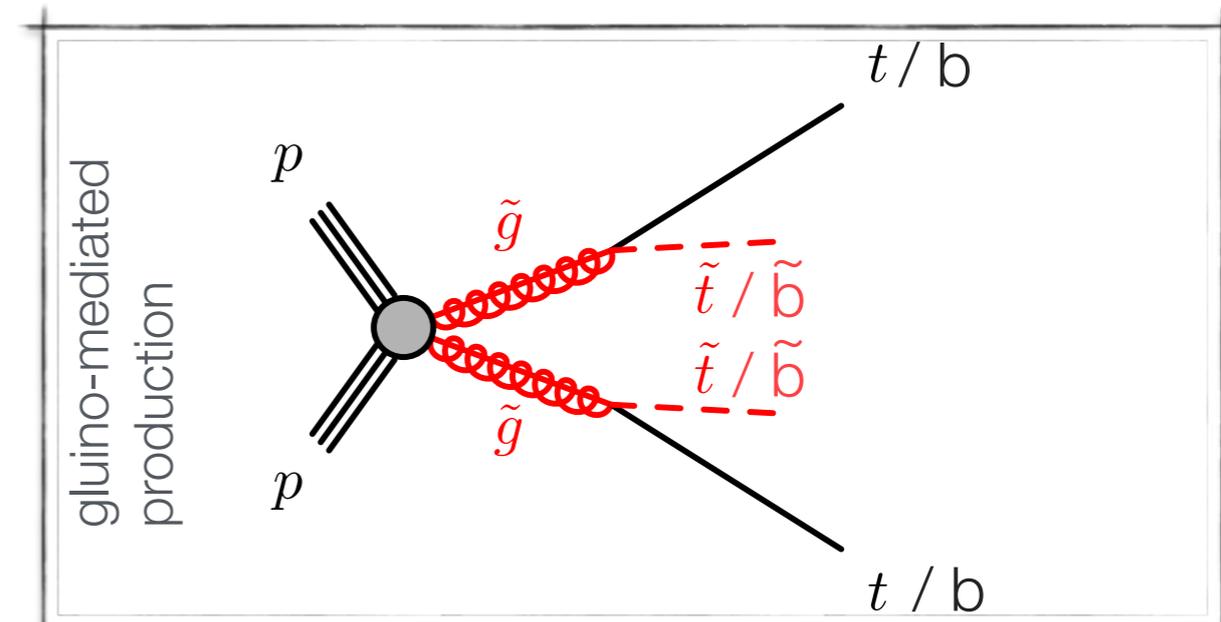
Inclusive searches for strong prod. of gluinos & first two gen. squarks

EW searches



Search for stop and sbottom via:

- ▶ **gluino-mediated** production,
- ▶ **direct pair** production.



3rd generation searches at the LHC

LHC search program relies on the complementarity of many analyses.

CMS direct stop:

- 0L, PAS-SUS-13-015
- ▶ • 1L, EPJC 73 (2013) 2677
- combination 0L razor, 1L, PAS-SUS-14-011
- stop in charm, PAS-SUS-13-009
- ▶ • stop2 in Z/h, PLB 736 371 (2014)

CMS direct sbottom:

- 0L, PAS-SUS-13-018

CMS gluino-mediated stop/sbottom:

- (0,1,2)L razor, PAS-SUS-14-011
- 0L, CMS-SUS-13-012
- 1L, PLB 733 328 (2014)
- 2OS L, PAS-SUS-13-016
- 2SS L, JHEP 01 (2014) 163
- 3L, PAS-SUS-13-008
- ▶ • combination of above 5 searches, PAS-SUS-14-010 (coming out soon)

ATLAS direct stop:

- 0L, JHEP 09 (2014) 015
- ▶ • 1L, Submitted to JHEP [1407.0583]
- 2L, JHEP 06 (2014) 124
- stop in charm, Accepted by PRD [1407.0608]
- ▶ • stop2 in Z/GMSB stop, Eur. Phys. J. C (2014) 74:2883
- stop in stau, ATLAS-CONF-2014-037

ATLAS direct sbottom/stop:

- 0L, JHEP 10 (2013) 189

ATLAS gluino-mediated stop/sbottom:

- ▶ • 0-1L + ≥ 3 b-jets, Accepted by JHEP [1407.0600]
- 2SS/3L, JHEP 06 (2014) 035
(also direct sbottom)
- 0L + $\geq (7-10)$ jets, JHEP 10 (2013) 130

ATLAS constraints from tT measurements:

- ▶ • 2L cross section, Submitted to EPJC [1406.5375]
- ▶ • 2L spin correlation, ATLAS-CONF-2014-056

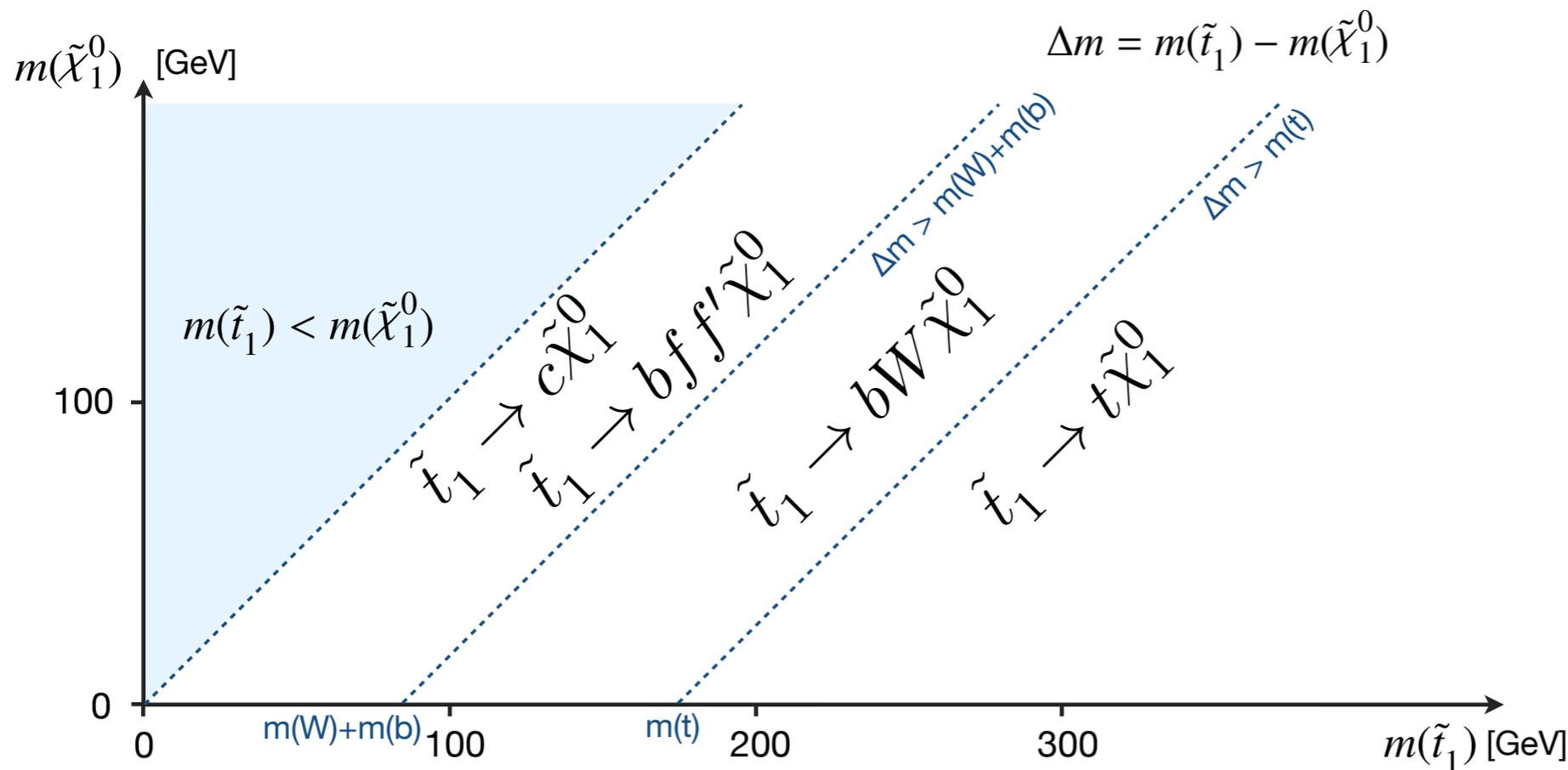
▶ Results presented in this talk, focus on a light stop and processes with a top quark in the decay.

Comprehensive list of all public [ATLAS](#) and [CMS](#) SUSY search results.

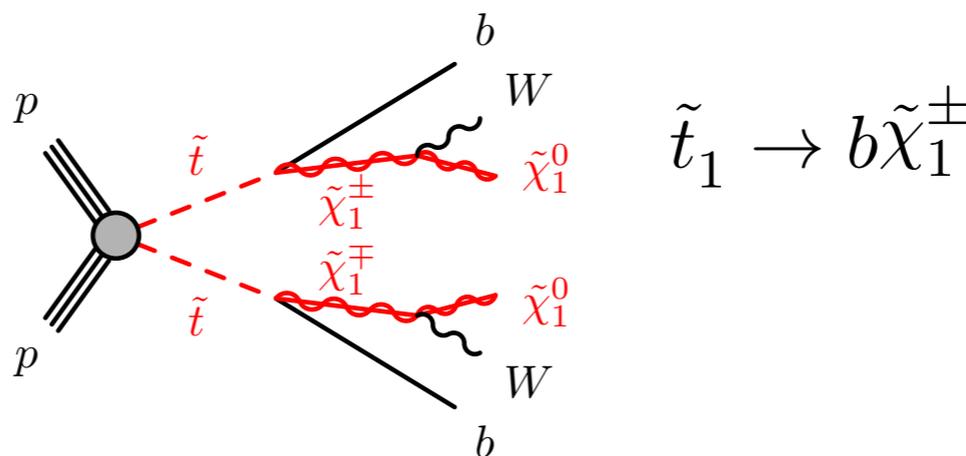
Overview: direct stop searches

Most searches target lightest eigenstate (stop1).
Several decay modes possible.

$$\begin{pmatrix} \tilde{t}_1 \\ \tilde{t}_2 \end{pmatrix} = \begin{pmatrix} \cos \theta_t & \sin \theta_t \\ -\sin \theta_t & \cos \theta_t \end{pmatrix} \begin{pmatrix} \tilde{t}_L \\ \tilde{t}_R \end{pmatrix}$$



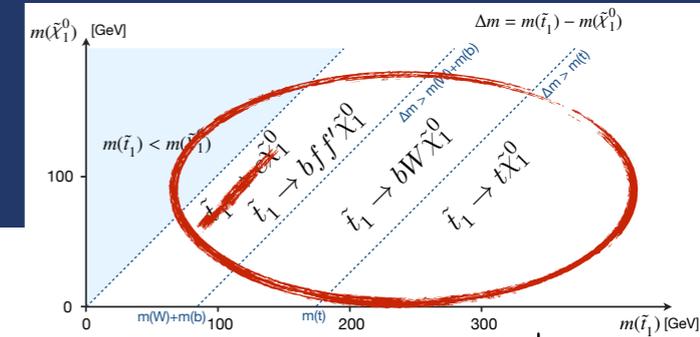
stop decays directly
to the lightest neutralino



$$\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$$

stop decays via the
lightest chargino

other (complex) decays
possible ...



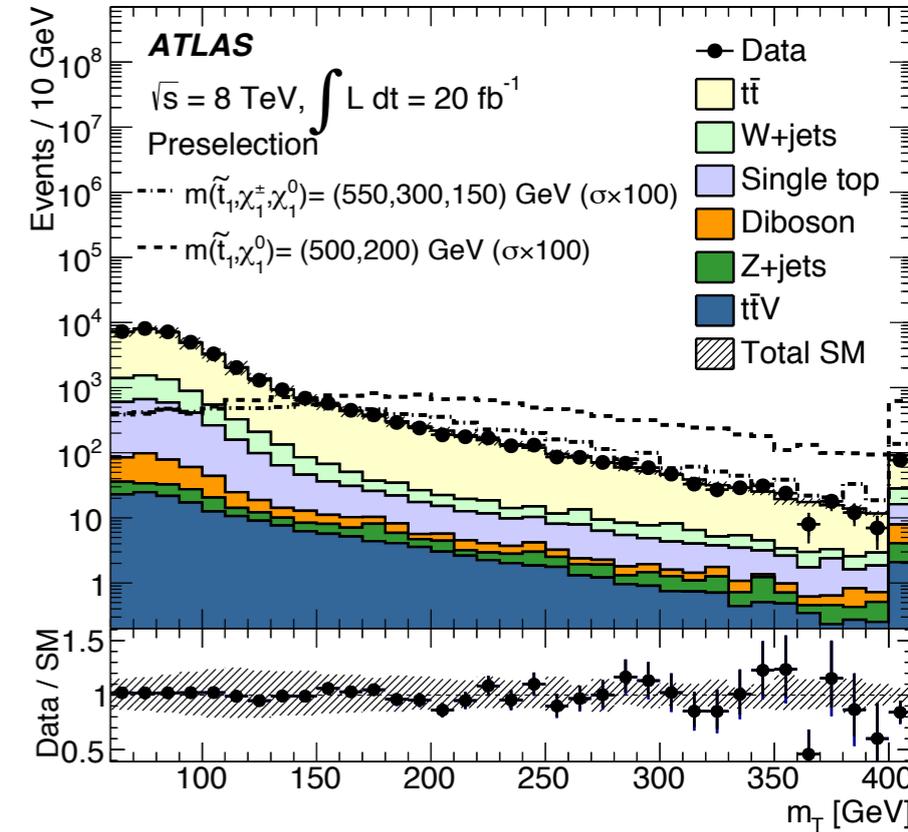
$$\tilde{t}_1 \rightarrow b \tilde{\chi}_1^\pm$$

Search targets light to heavy stops in various decay modes: 4-body, 3-body, tN1, and bC1.

1-leptons, ≥ 2 -4 jets, MET, MT(lep, pTmiss)

15 signal regions (SRs).

- **soft-lepton**: $p_T > 7(6)$ GeV for electrons (muons), employed for 4-body, and bC1 decays with small mC1-mN1 values.
- **large-R (1.0) jets**: collect decay products of boosted top in heavy stop search.
- **1D and 2D shape-fits**: enhance search sensitivity in challenging scenarios (with difficult S-vs-B separation).
- **b-tagged jets**: utilised in event selection (ranging from a veto to ≥ 2) and for constructing kinematic variables.
 - ▶ MT2-like variables, hadronic top mass, topness



- Dominant background: **semi (di)-leptonic $t\bar{t}b\bar{b}$** for the stop 0L (1L) search, where the (2nd) lepton is

1. not identified (failed PID) / outside acceptance,
2. hadronically decaying tau-lepton.



missed particle / extra neutrino(s)
 \rightarrow more E_T^{miss} , $m_T > m(W)$



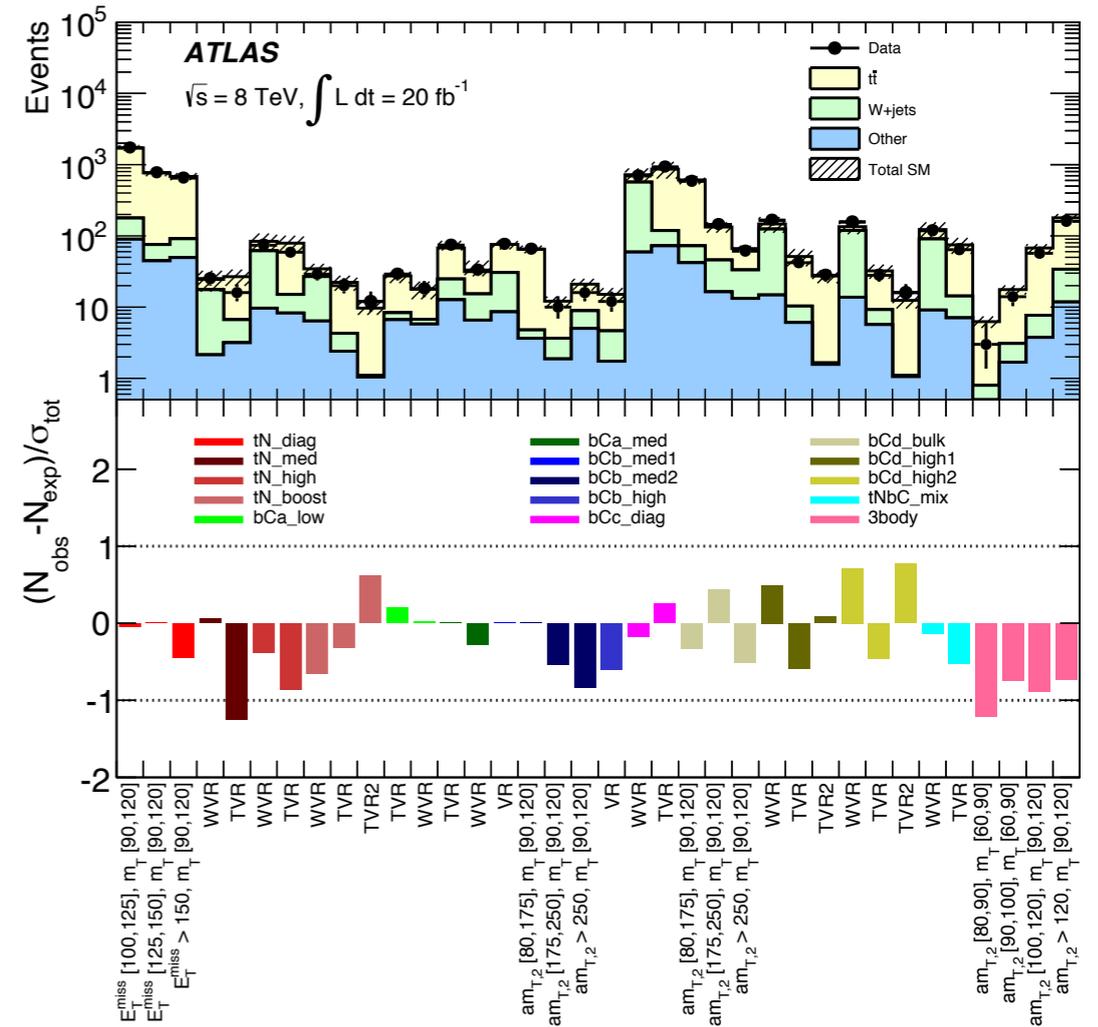
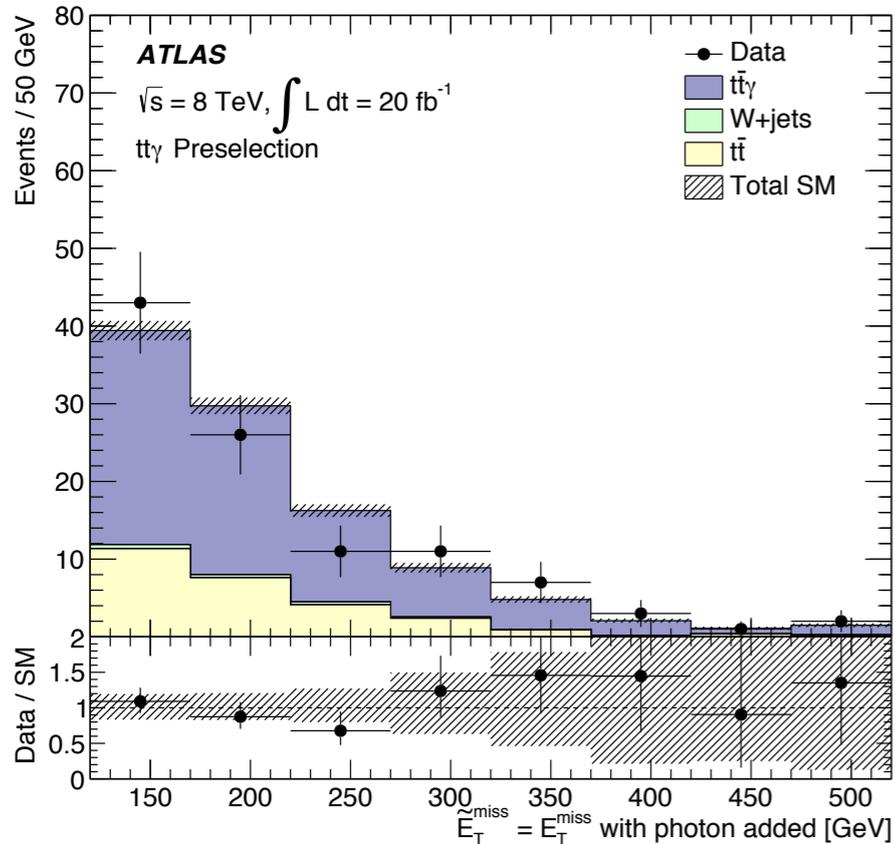
Backgrounds

- dileptonic tT (with one lep. out of acceptance, not identified, or is a had tau) — 1L lower-MT CR
- W+jets + HF — 1L lower-MT, b-veto CR
- tT+Z(nunu), and other (small) backgrounds — MC

Validation:

 compare data with prediction in dedicated samples

- kinematically close to SR (see plot),
- 2L, Z-veto — tT additional jets,
- W+HF validation selection,
- tT+Z(nunu) validation via tT+gamma (see plot), and more.



Tools to suppress dominant tT bkg

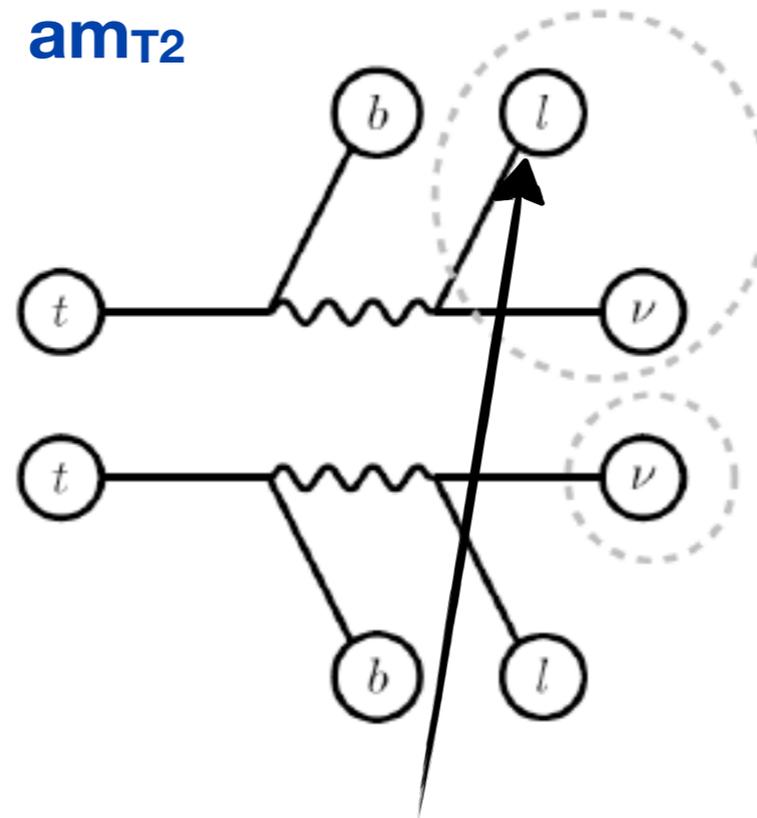
Similar variables also employed by CMS search.

Kinematic variables: M_{T2}

$$m_{T2} \equiv \min_{\vec{p}_{Ta}^C + \vec{p}_{Tb}^C = \vec{p}_T^{\text{miss}}} \{ \max(m_{Ta}, m_{Tb}) \}$$

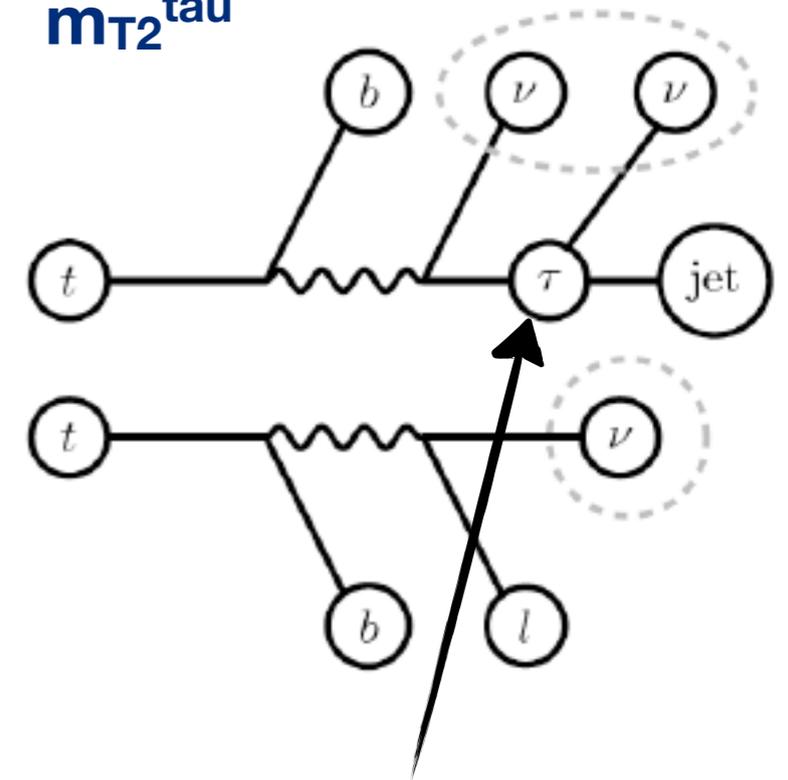
Dedicated variants of M_{T2} :

am_{T2}



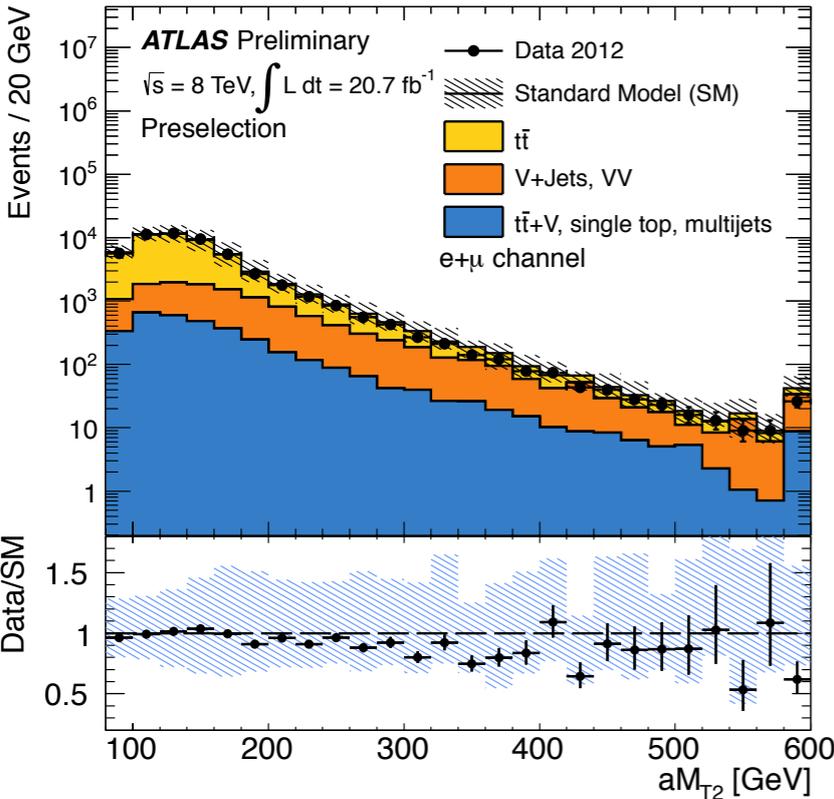
"lost" lepton (category 1),
 $leg_a = bjet1 + lep$, $leg_b = bjet2$,
 or vice versa (take min.)

m_{T2}^{tau}



Had-tau (category 2),
 $leg_a = jet3$, $leg_b = lep$, where
 jet3 is hardest non b-tagged jet

di-leptonic t \bar{t} has kinematic end-points: around top mass (am_{T2}), and the W mass (m_{T2}^{tau}); while signal can exceed these.

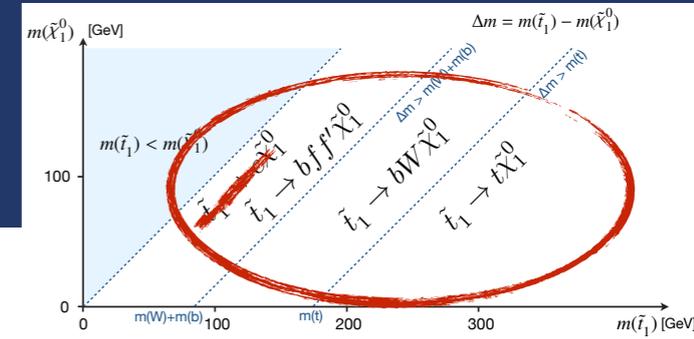




ATLAS searches

1L channel

submitted to JHEP [1407.0583]



4-body:

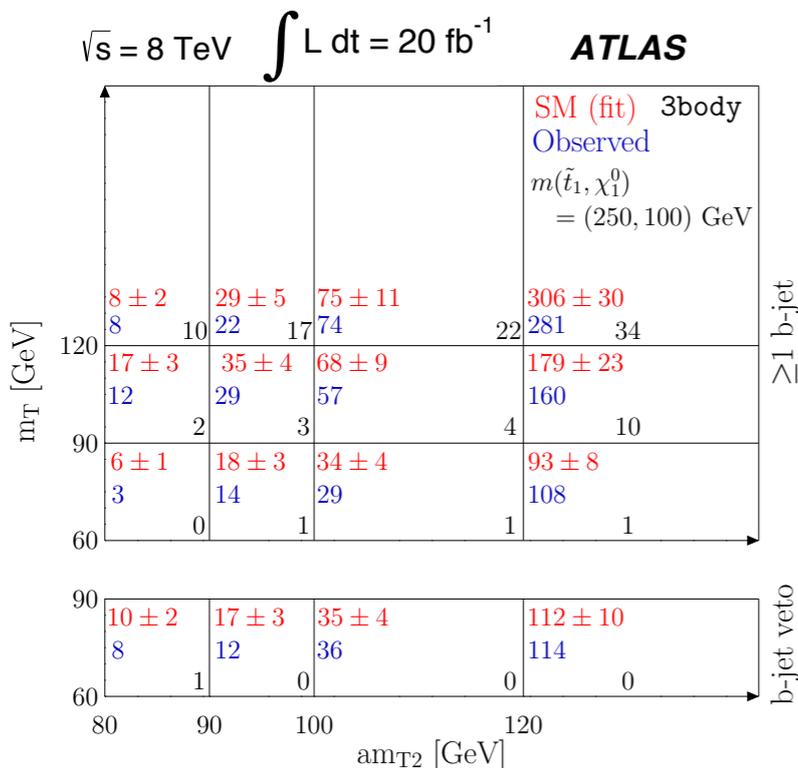
- 1 soft-lepton
- 1 hard non-b-tagged jet (ISR)
- ≥ 1 b-tag in sub-leading jets
- high MET, $MT > m_W$,

3-body:

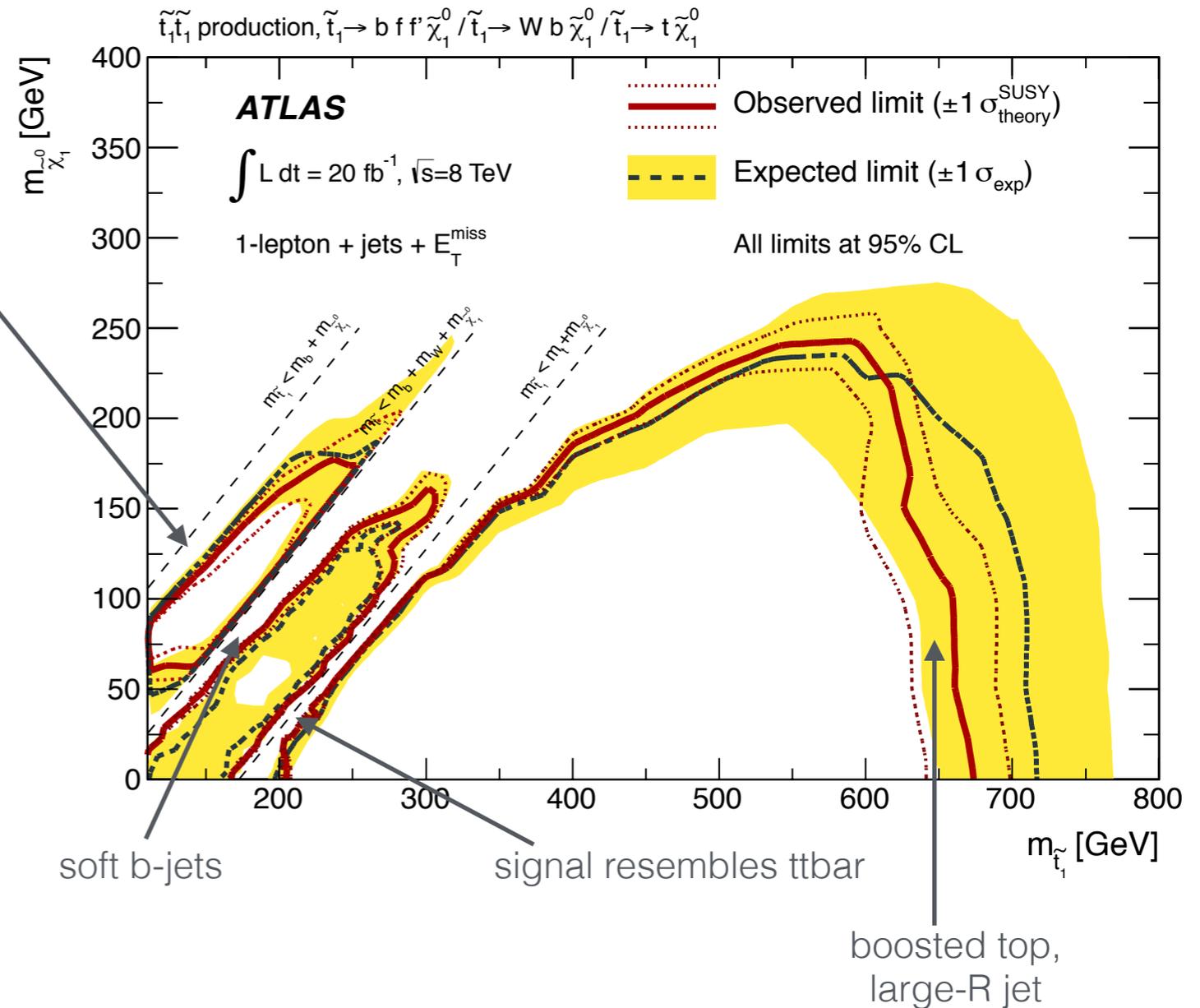
- 1 lepton
- ≥ 4 jets, at least one of which is b-tagged
- high MET, 2D shape-fit in MT and amT2

mStop \approx mTop+mN1:

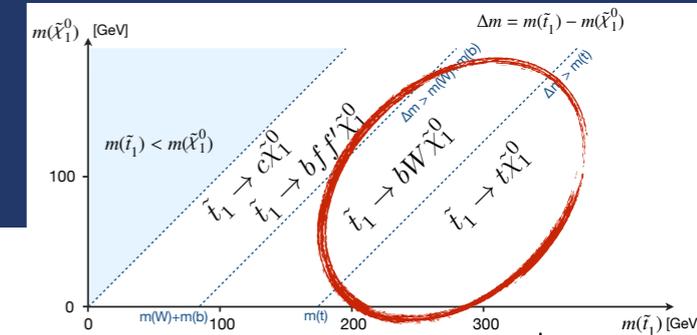
- 2D shape fit but with looser selection, MT-MET



too small Δm , but monojet



“boosted tops” covered in talk by Johannes Erdmann (Wednesday)



$$\tilde{t}_1 \rightarrow b \tilde{\chi}_1^\pm$$

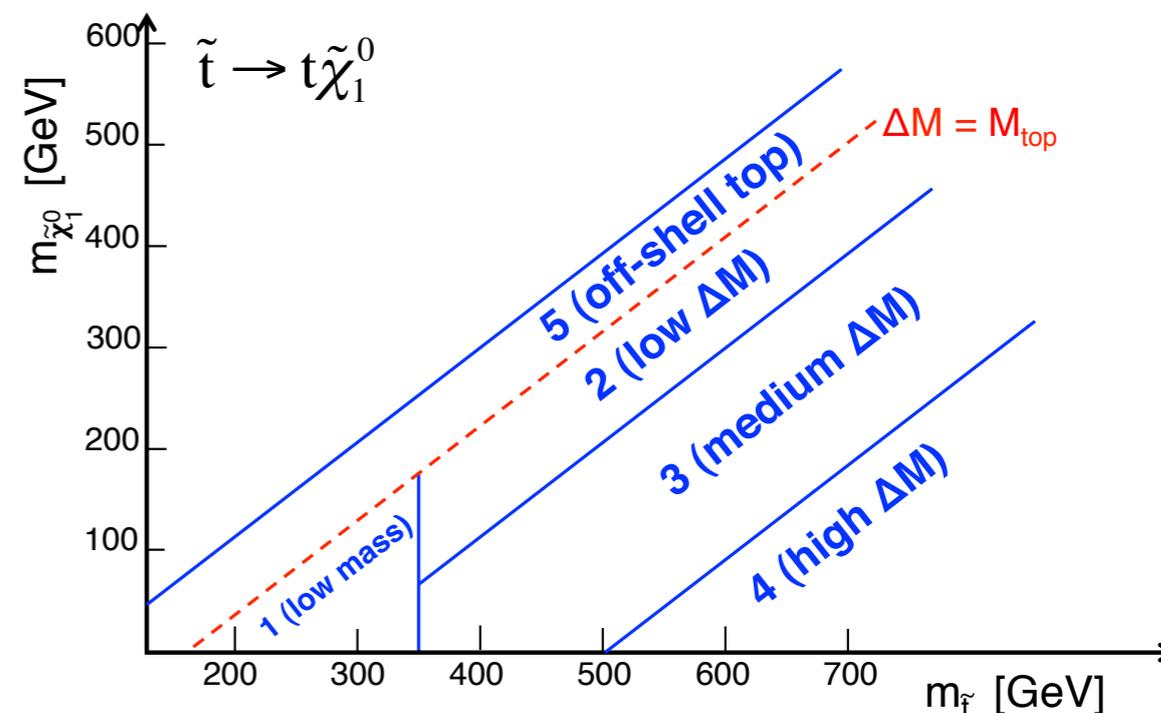
Search targets medium to heavy stops in various decay modes: 3-body, tN1, and bC1.

Baseline selection

- 1 isolated lepton (e/mu, veto extra e,mu,tau)
- ≥ 4 jets, at least one b-tagged
- MET > 100 GeV
- MT(lepton, pTmiss) > 120 GeV

18 BDT analyses

- trained on different regions of signal models (tN1 example shown in plot)
- MET, MT2W, relative HT btw. event hemispheres,
- hadronic top compatibility (chi2)
- lep, b-jet pT
- distance btw. lep and b-jet
- cut-based analyses for cross check.





Backgrounds

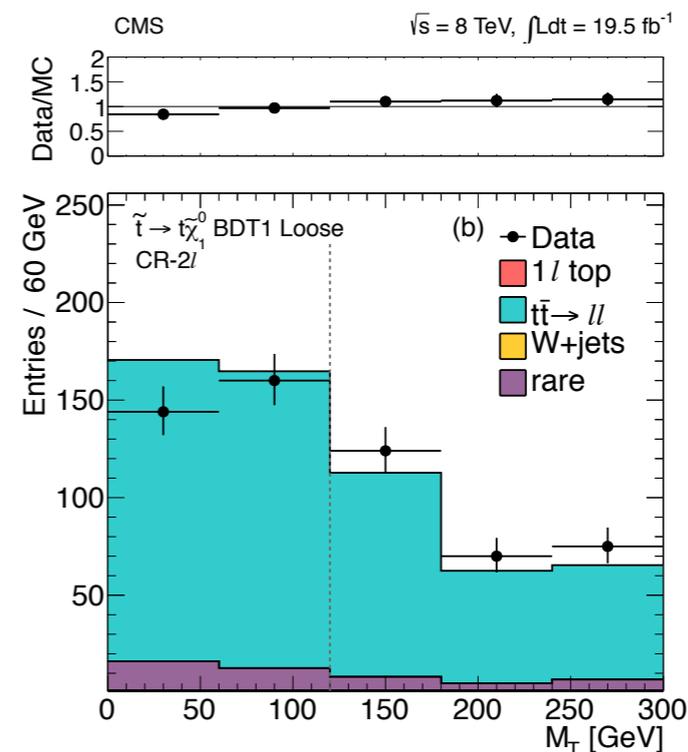
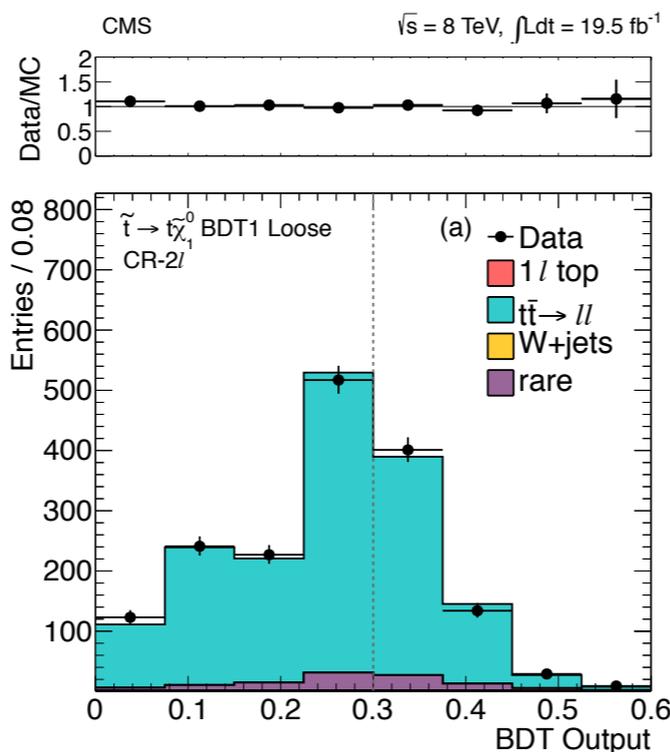
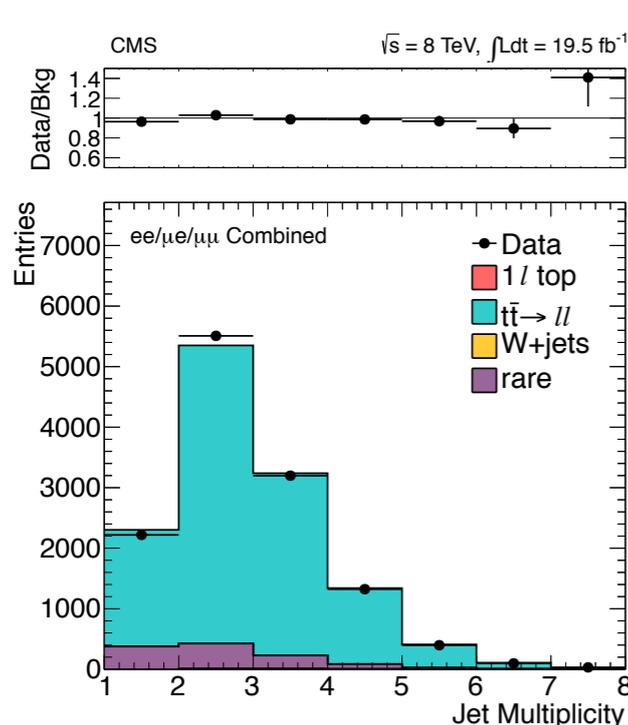
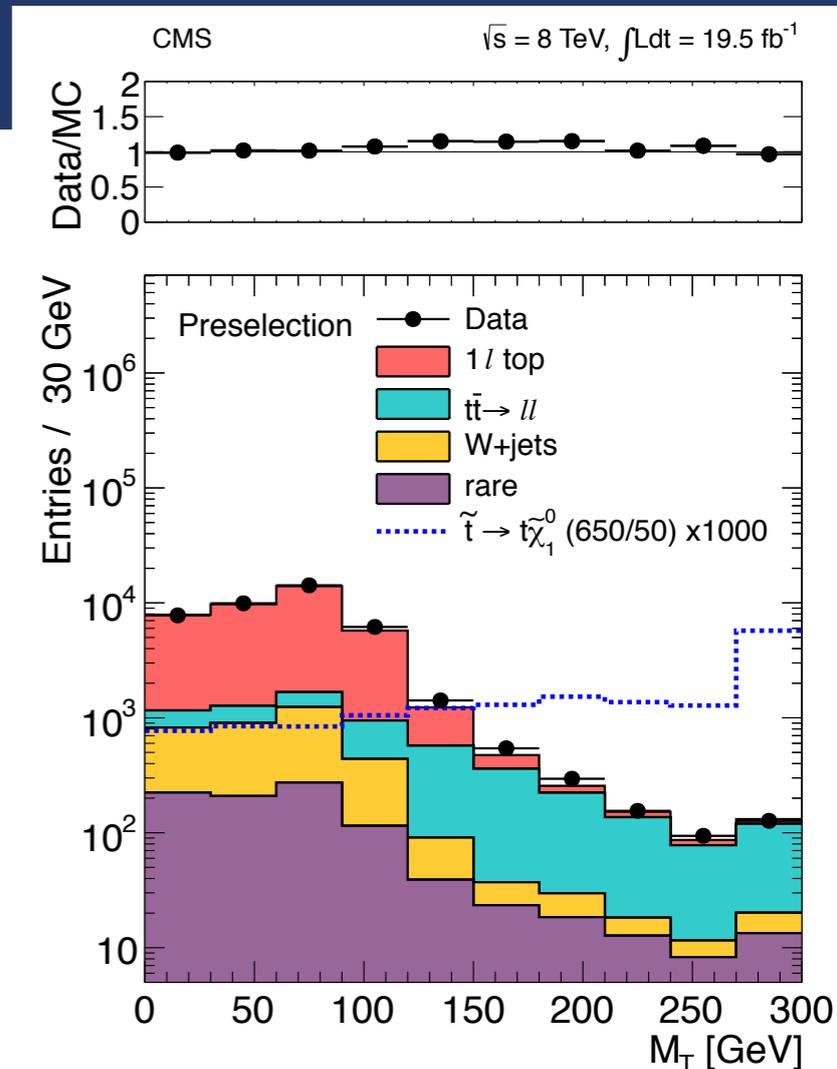
- dileptonic tT (one lepton missed) $\rightarrow ll$
- semi-lepton tT + t $\rightarrow 1l$ top
- W+jets
- other (small) backgrounds \rightarrow rare

Estimation:

- normalise backgrounds in low M_T window

Validation: compare data with prediction in dedicated samples

- 2L, 1L+isolated track, b-veto
- uncertainties derived from these studies



Tools to suppress dominant tT bkg

Similar variables also employed by ATLAS search.

Tight 2nd lepton veto:

Signal has no 'edge' in mT hence dominated by single lepton.

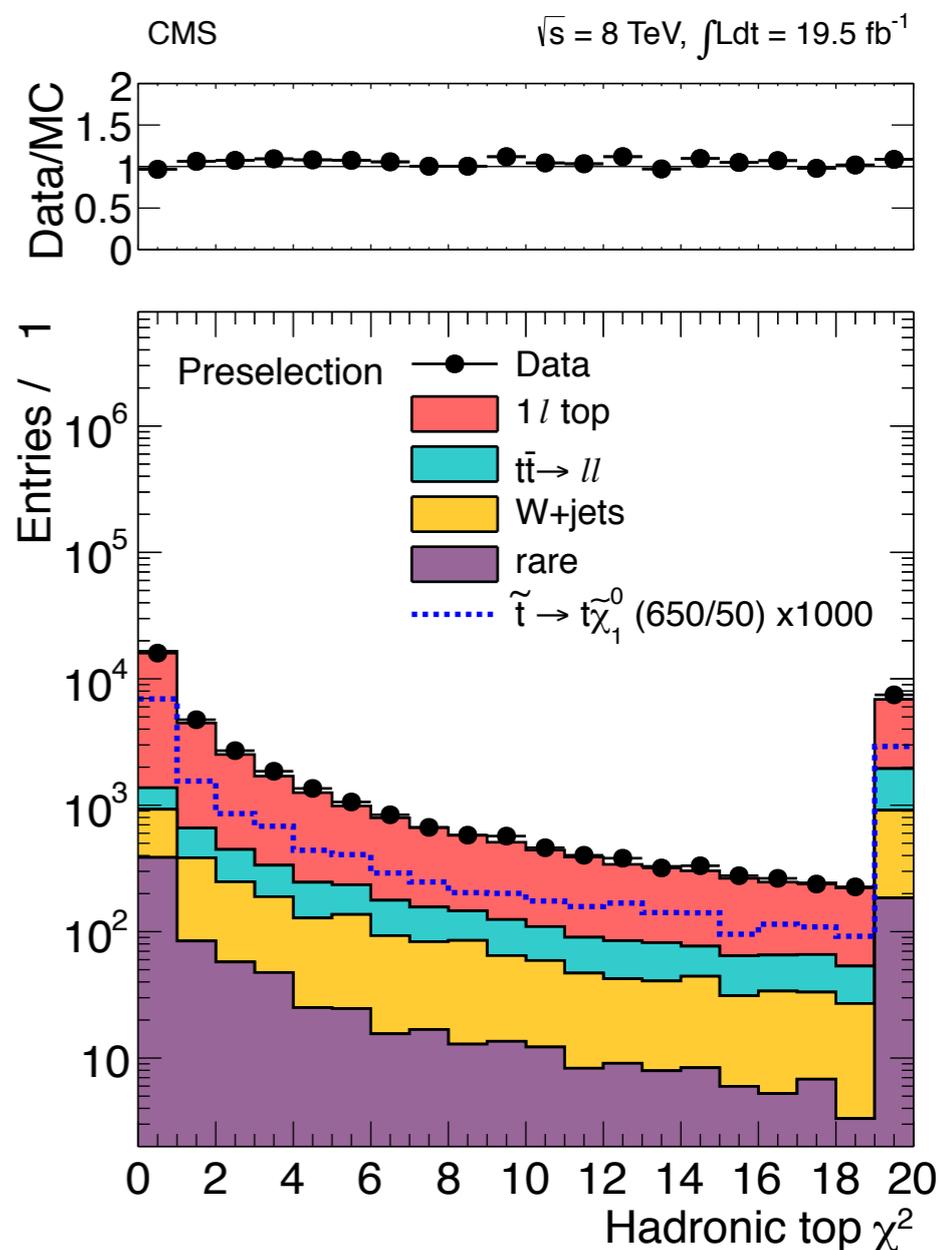
- lepton pT, PID
- veto on isolated tracks
- veto on had-taus

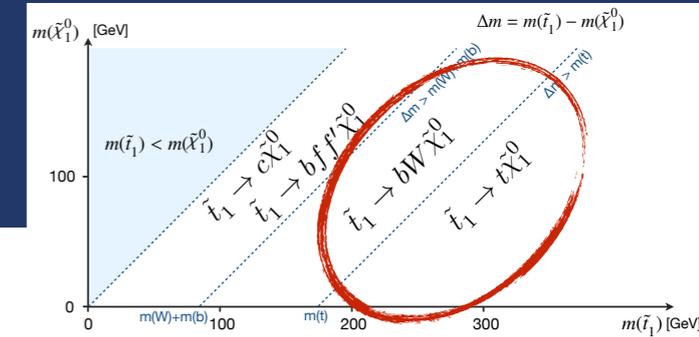
Hadronic top mass:

While for decay via top quark, three jets stem from a hadronic top decay, this is not true for di-lep ttbar.

- reconstruct hadronic top mass from three jets $m(jjj)$,
- consider b-tagging & jet energy resolution

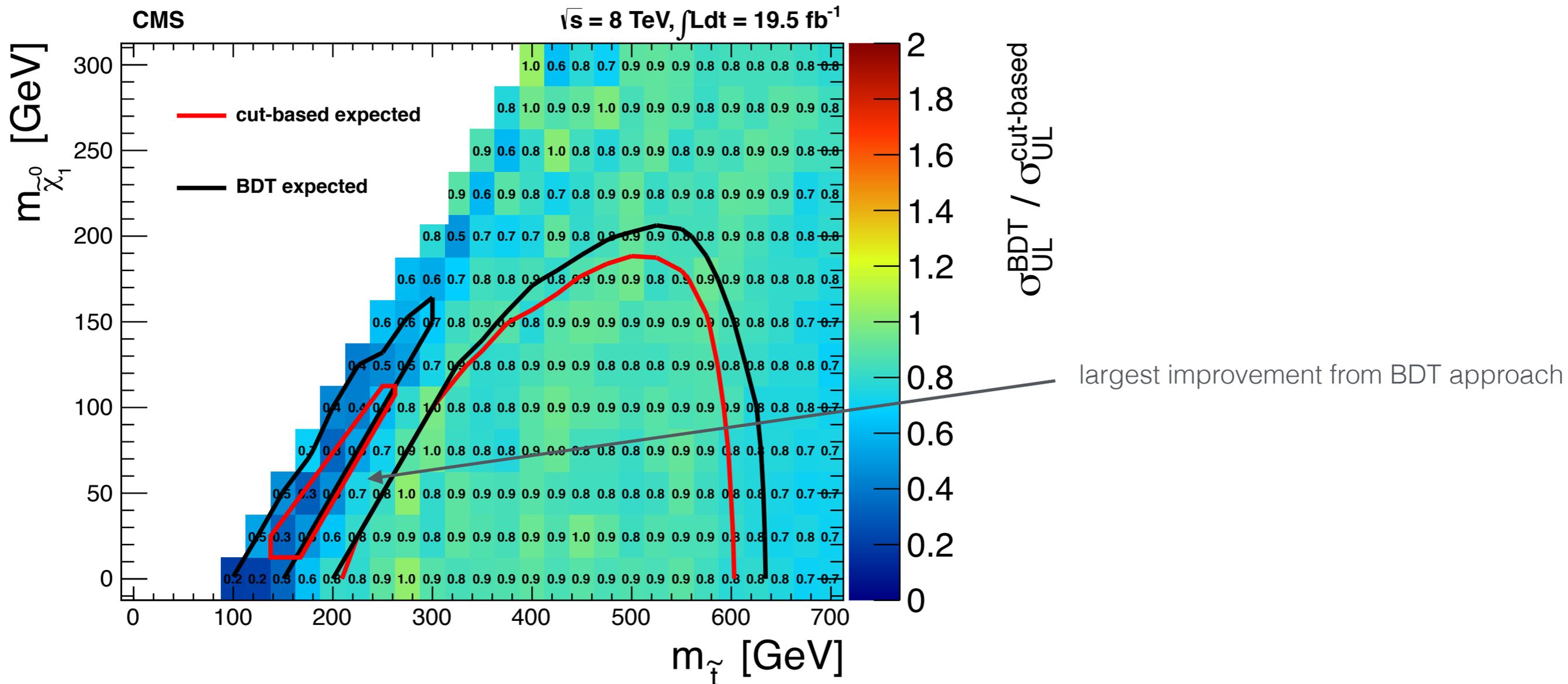
$$\chi^2 = \frac{(M_{j_1 j_2 j_3} - M_{\text{top}})^2}{\sigma_{j_1 j_2 j_3}^2} + \frac{(M_{j_1 j_2} - M_W)^2}{\sigma_{j_1 j_2}^2}$$





Exclusion limits

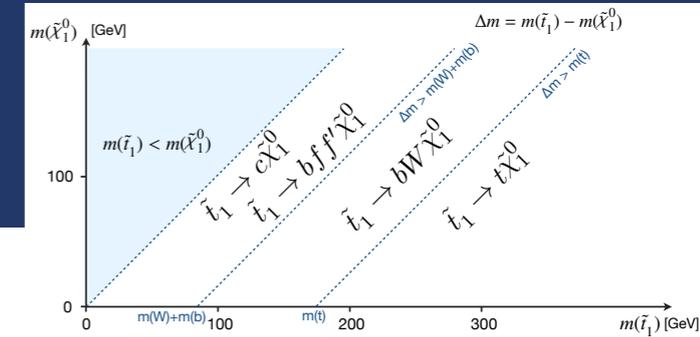
- BDT analyses
- cross section limits from BDT improve over cut-based analysis by up to ~40%



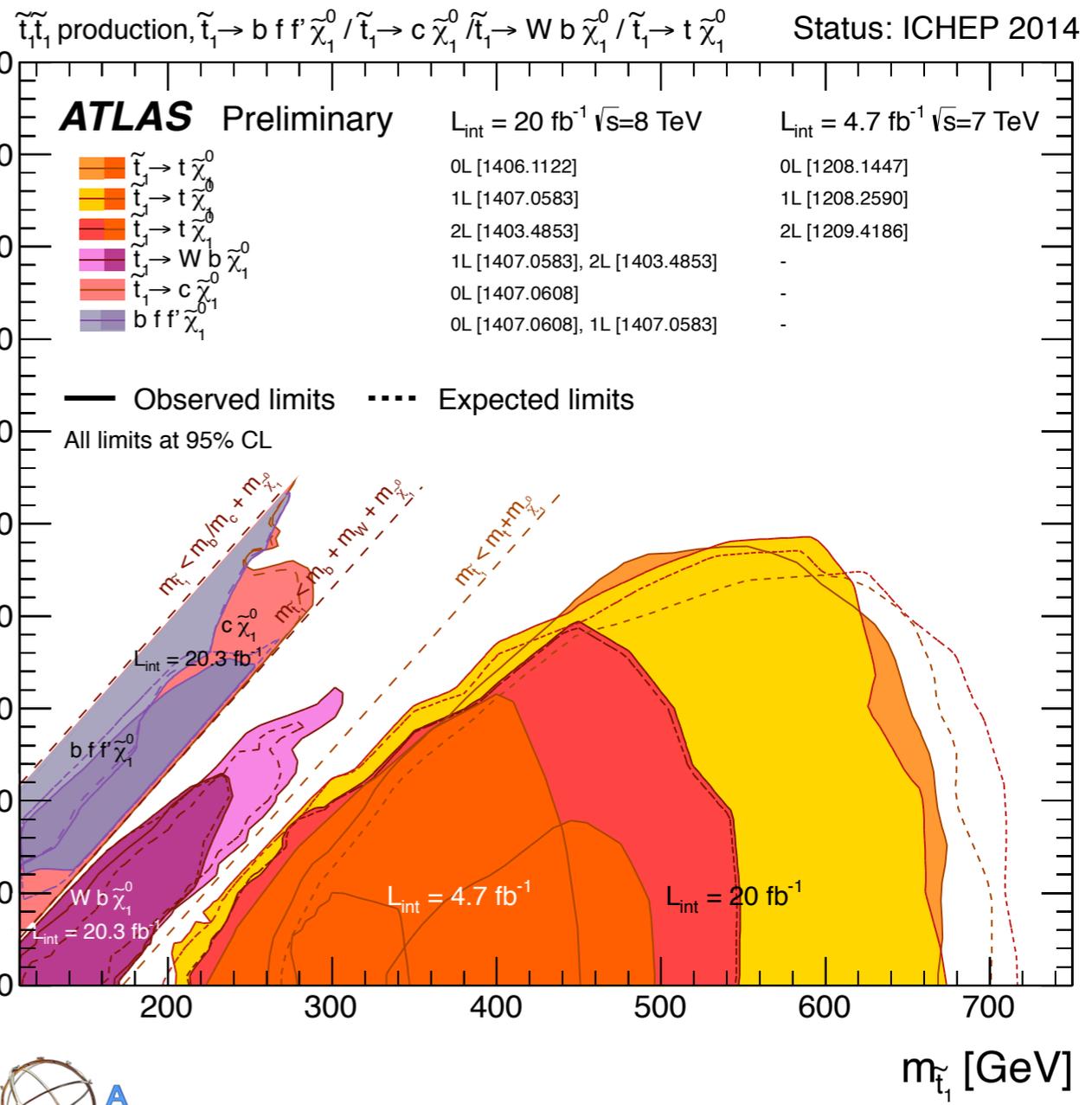
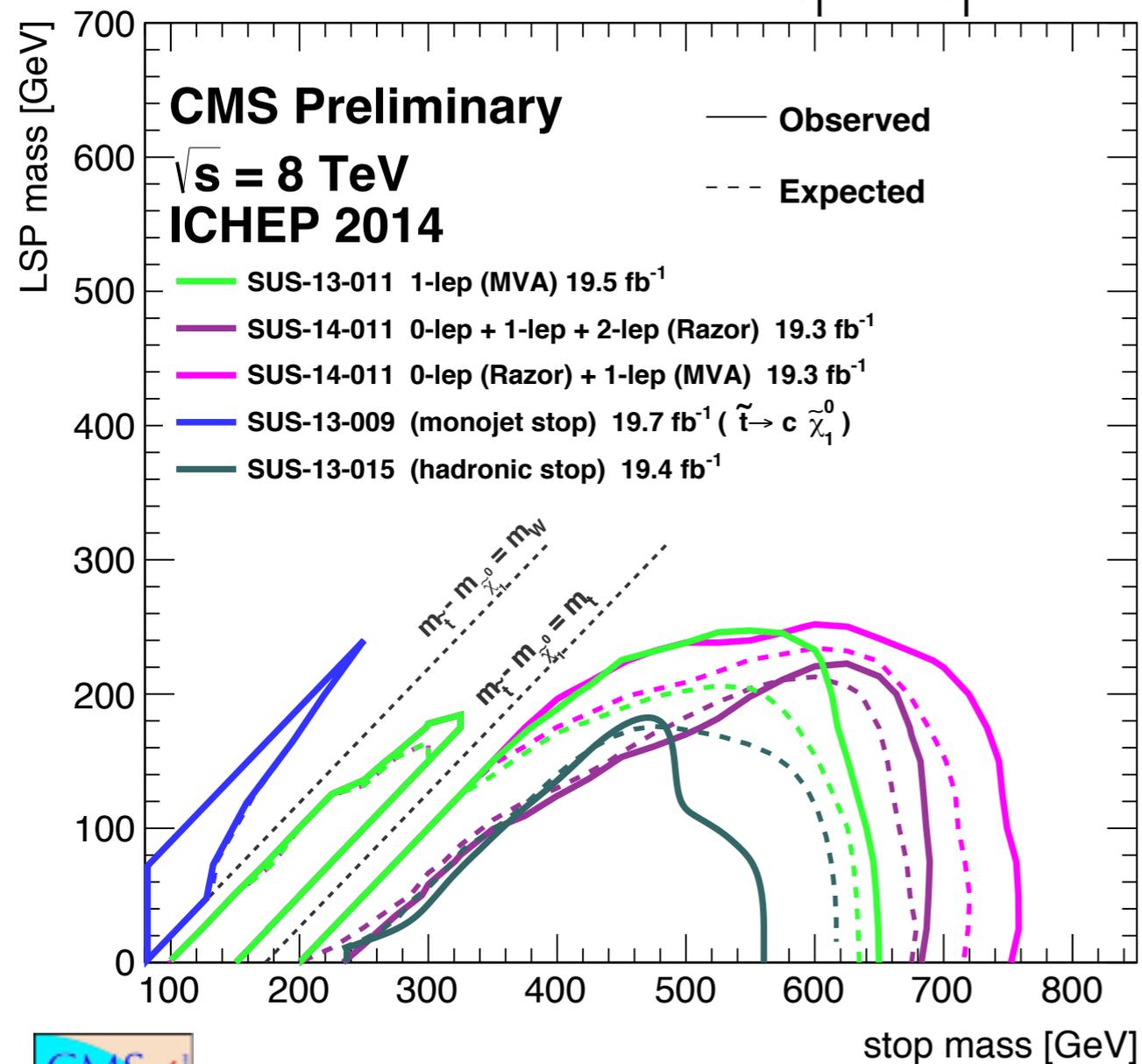
Acceptance also depends on the polarization of the the top quarks in the stop decay. Effect from left/right mixing of the stops and on the mixing matrices of the neutralino and charginos, impact shown in backup (CMS & ATLAS).

Complementarity of searches

Many analyses to cover the various decays and mass regions.
 Furthermore, other decays such as via a chargino are considered & largely covered. However, holes remain!

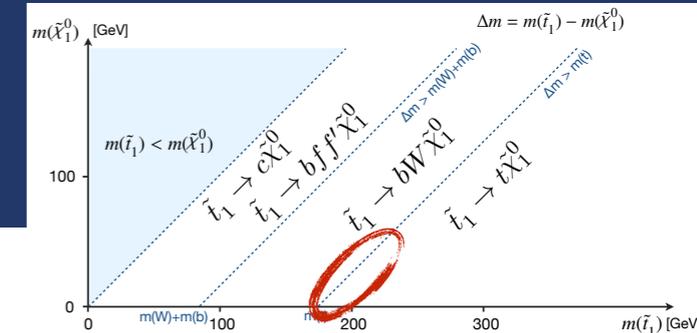


$$\tilde{t}\text{-}\tilde{t} \text{ production, } \tilde{t} \rightarrow t \tilde{\chi}_1^0 / c \tilde{\chi}_1^0$$



'Stealth' stop

Close the gap where $m(\text{stop1}) - m(\text{N1}) \approx m(\text{top})$



challenge: kinematic variables (E_T^{miss} , M_T , etc.) resemble those of $t\bar{t}$ background.

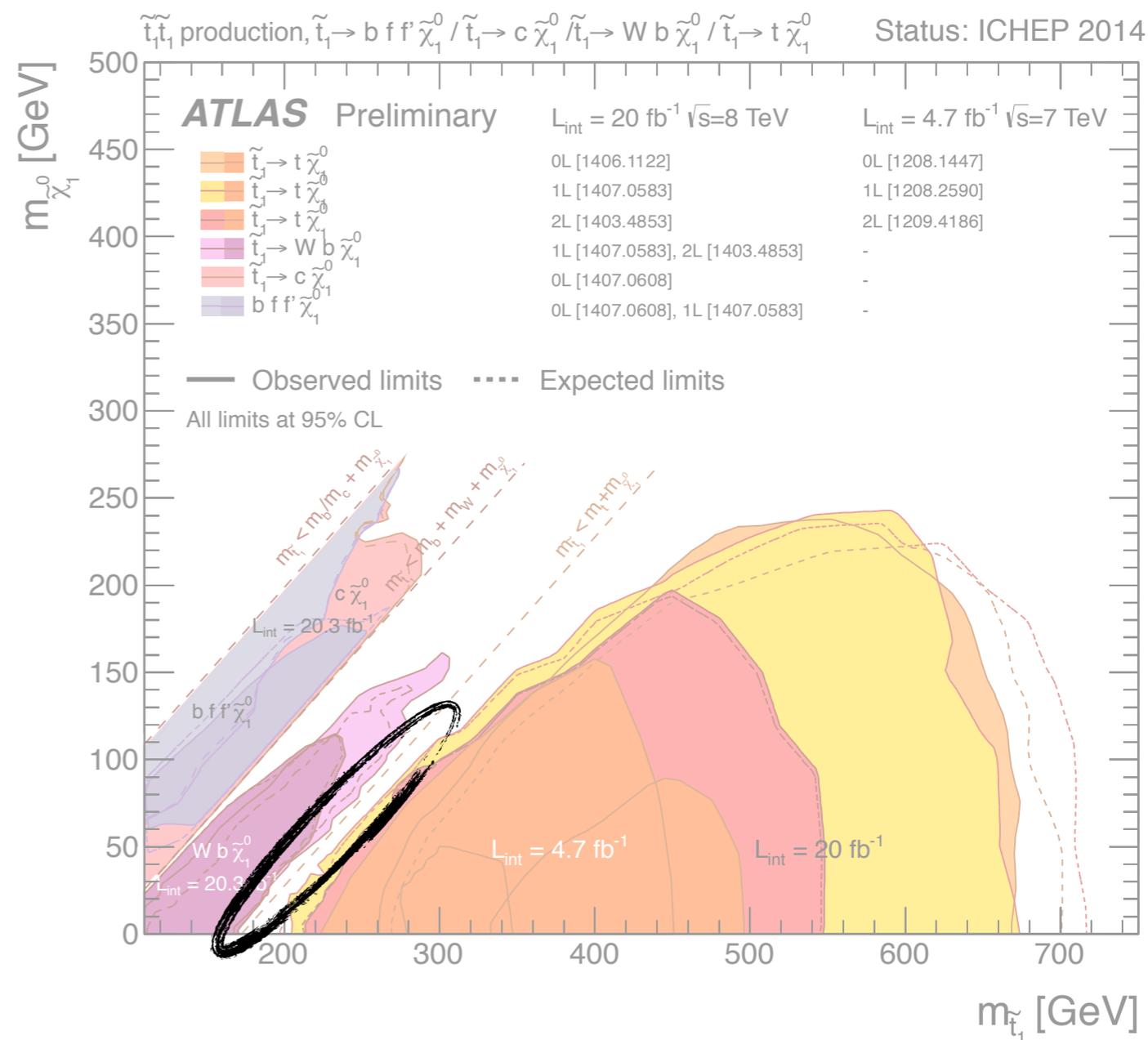
- But smaller signal cross-section, $\sim 15\%$ of SM $t\bar{t}$ for same mass.

1. Direct approaches:

- ▶ shape-fit
- ▶ Multivariate analysis
- ▶ Top measurements (later)

2. Indirect approach:

- ▶ search for heavier SUSY production with subsequent decay via stop1.
 - gluino to stop1 + t
 - stop2 to stop1 + Z/H



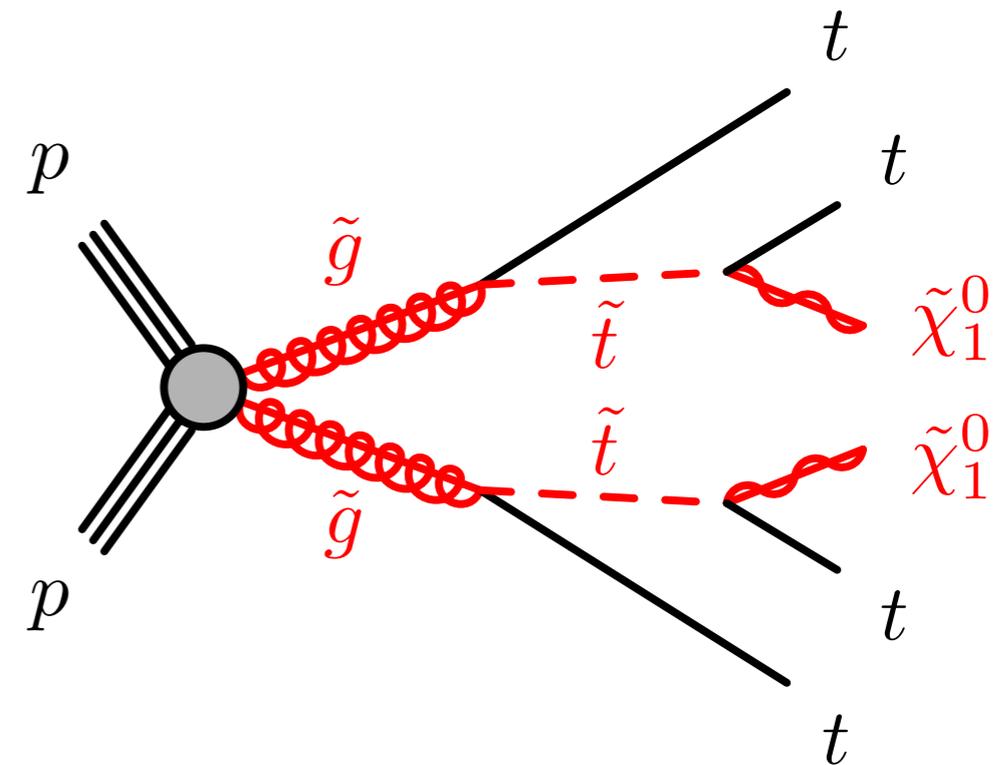
Searches for gluino-mediated stop production

Gluino pairs decaying via on- or off-shell stop lead to a spectacular signature of four top quarks + extra MET.

- ▶ Up to 4 (el,mu), 4 b-jets, ..
- ▶ Various searches in 0L — 3L channels,

SUSY: $m_{\text{Gluino}} < (3-5) \times m_{\text{Stop}}$,
and gluino pair production cross section much larger (x50)

Gluino-mediated searches give handle to search for a stealthy stop (mass degenerate with top).



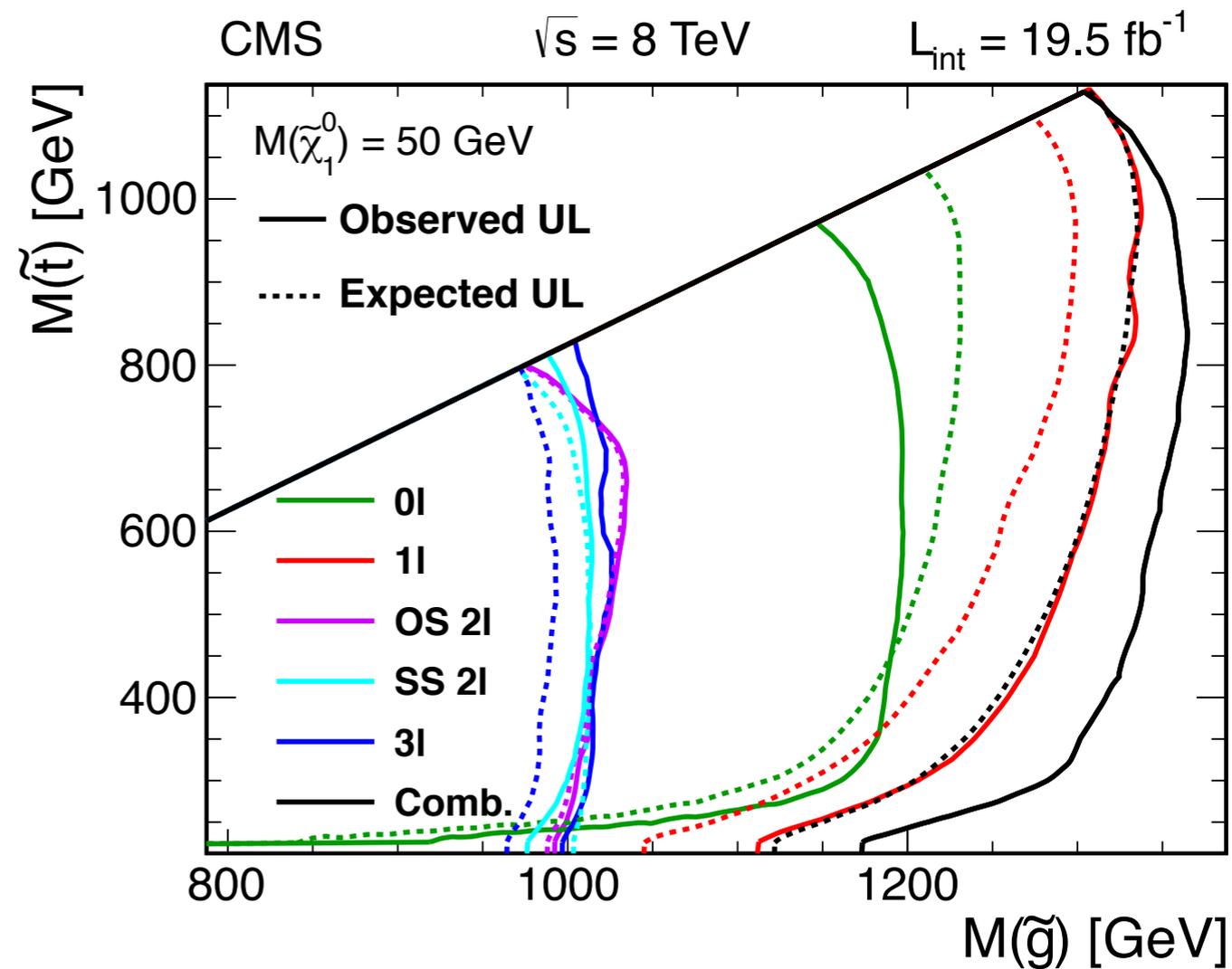
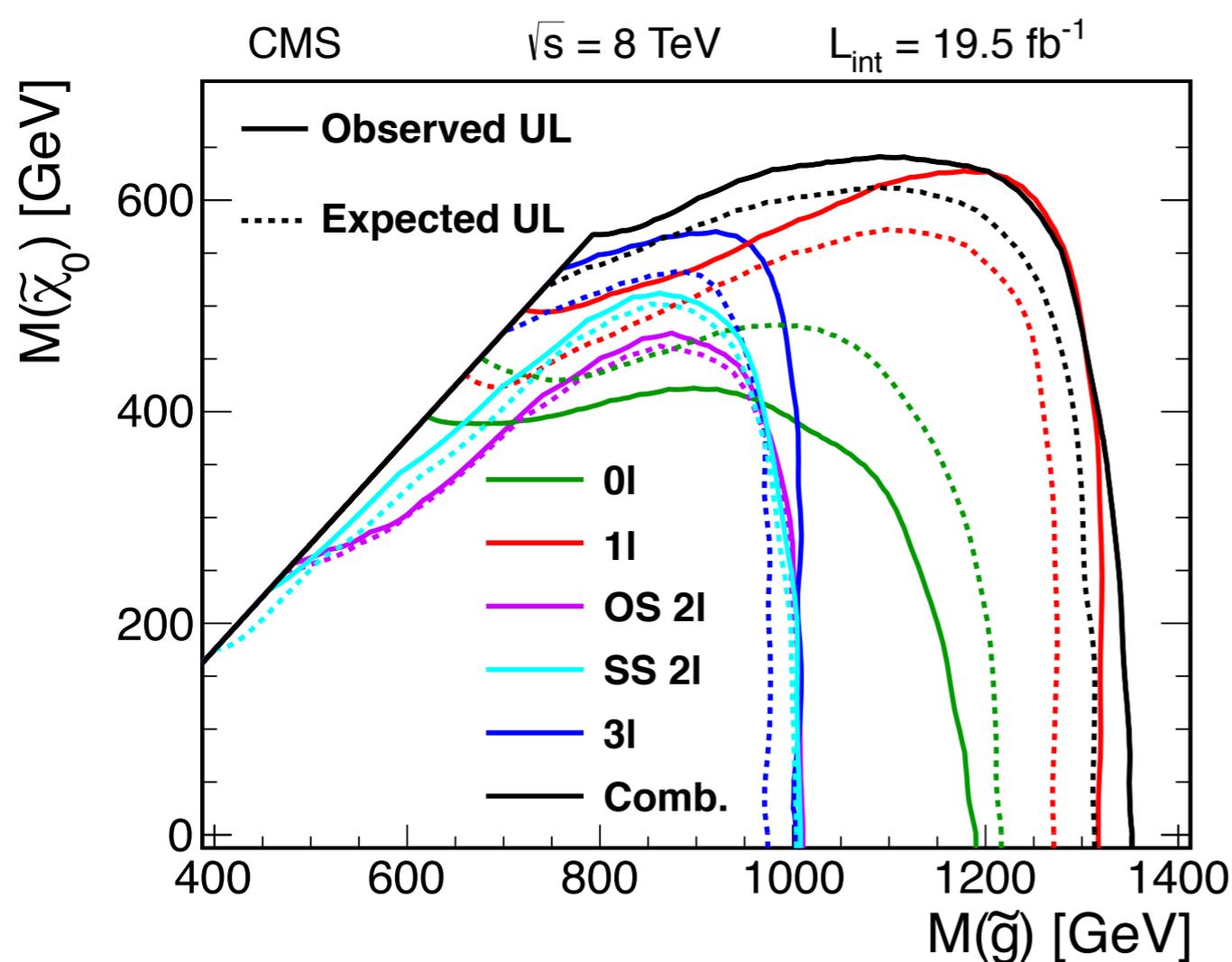
Backgrounds:

- **reducible**: for example, fake b-jets (mainly $t\bar{t}$) estimated from data (matrix method)
- **irreducible**: for example, processes with ≥ 3 bjets:
 - $t\bar{t} + b\bar{b}$, (using 2L control region),
 - $t\bar{t} + Z/h(\rightarrow b\bar{b})$ from MC



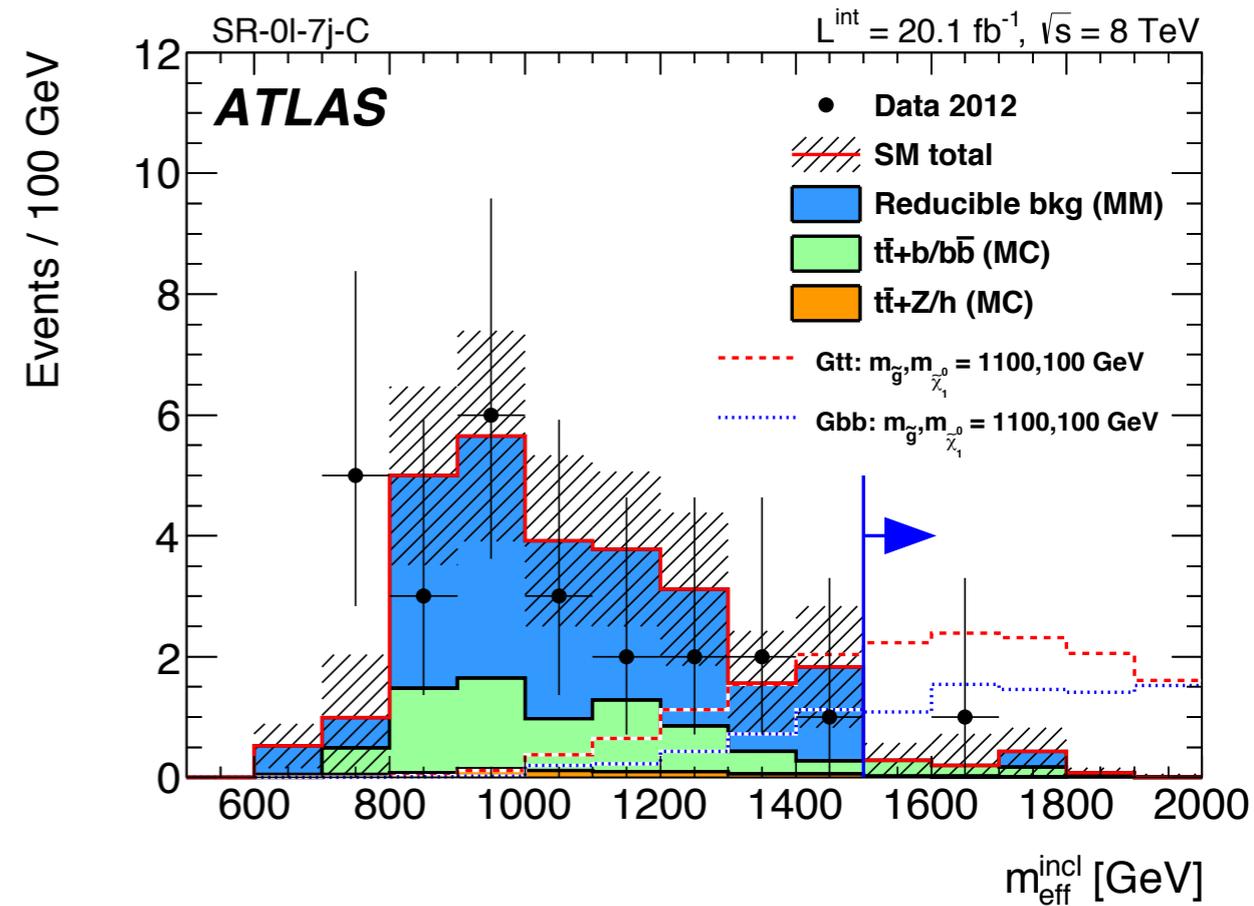
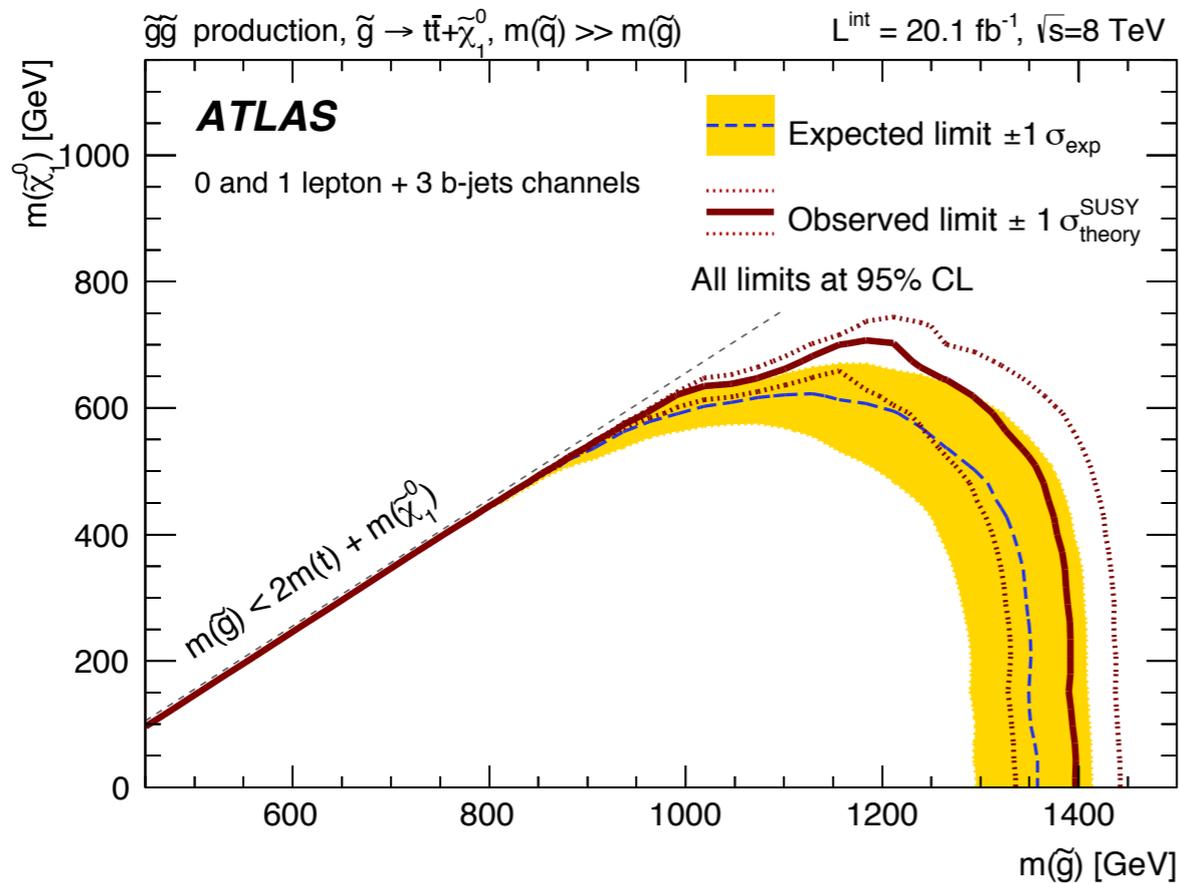
Combination of five orthogonal searches:

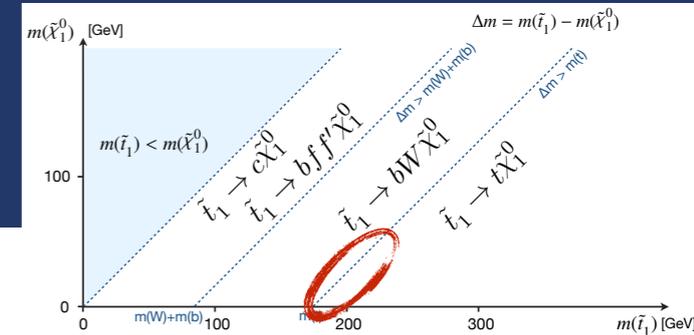
- 0L + multijets (3-8j) [CMS-SUS-13-012]
- 1L + ≥ 6 jets with ≥ 3 b-tagged [PLB 733 328 (2014)]
- 2OS L + ≥ 5 jets with ≥ 3 b-tagged + MET > 180 GeV [PAS-SUS-13-016]
- 2SS L + $\geq (2-4)$ jets with (0, 1, ≥ 2) b-tagged + MET cut [JHEP 01 (2014) 163]
- 3L + $\geq (2-4)$ jets with (1, 2, ≥ 3) b-tagged [PAS-SUS-13-008]



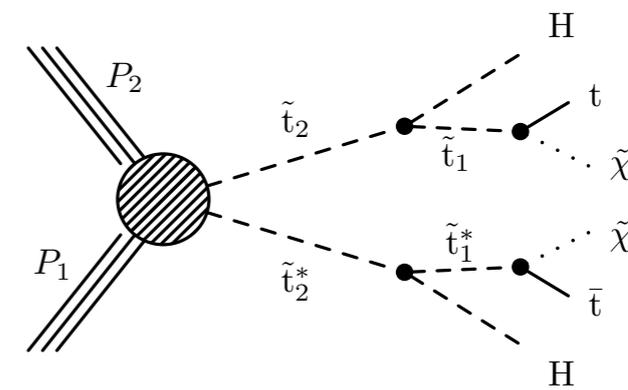
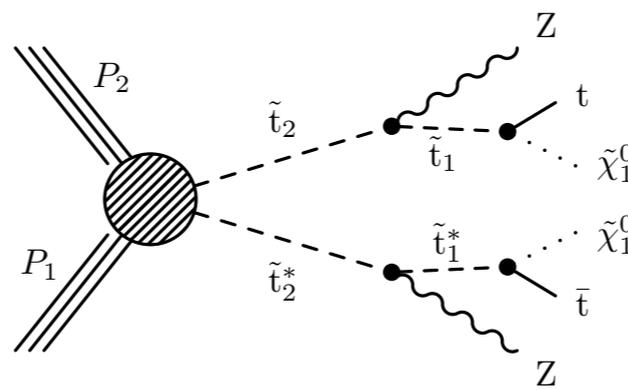
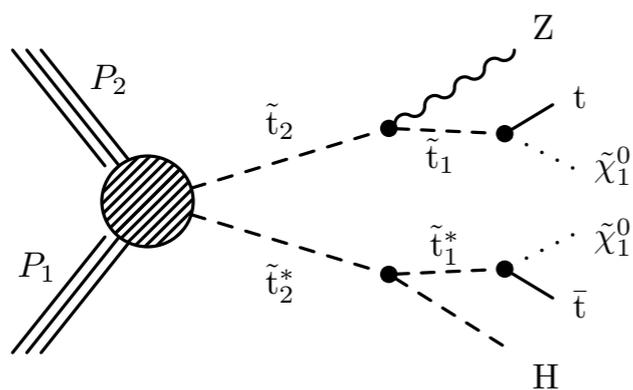
0L and 1L channels, optimized for various decays (sbottom, stop,..)

- 0L + ≥ 4 jets or ≥ 7 jets (≥ 3 b-tagged), with M_{eff} (scalar sum over jets, MET) cuts btw. 1 and 1.5 TeV and high MET
- 1L + ≥ 6 jets (≥ 3 b-tagged), with $M_{\text{eff}} > 700 - 900$ GeV and high MET and high m_T .





Target stop2 pair production, decay via (stealth) stop1 and Z or H.

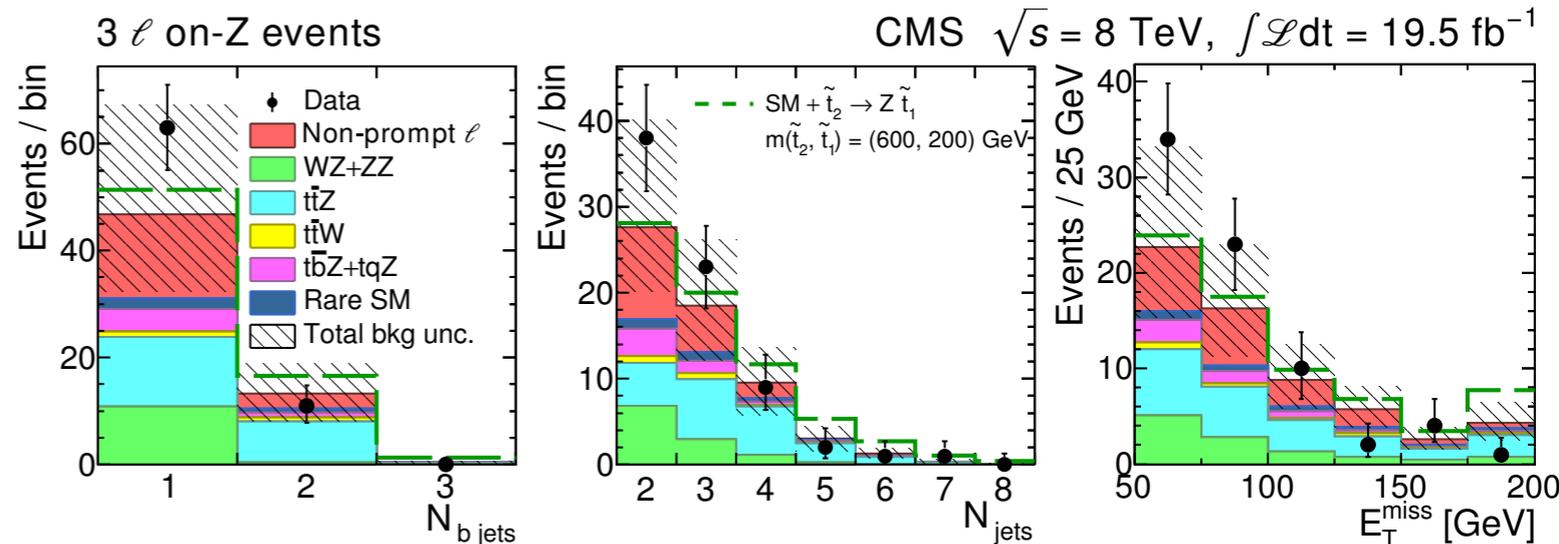


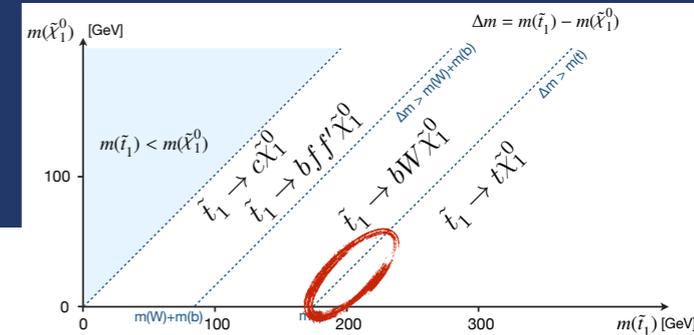
Sensitive to H to bb, H to ZZ, H to WW decays

- search with 1L, 2L (SS, OS), and $\geq 3L$
- b-tags (up to 4),
- split by on- and off-shell Z

3L analysis:

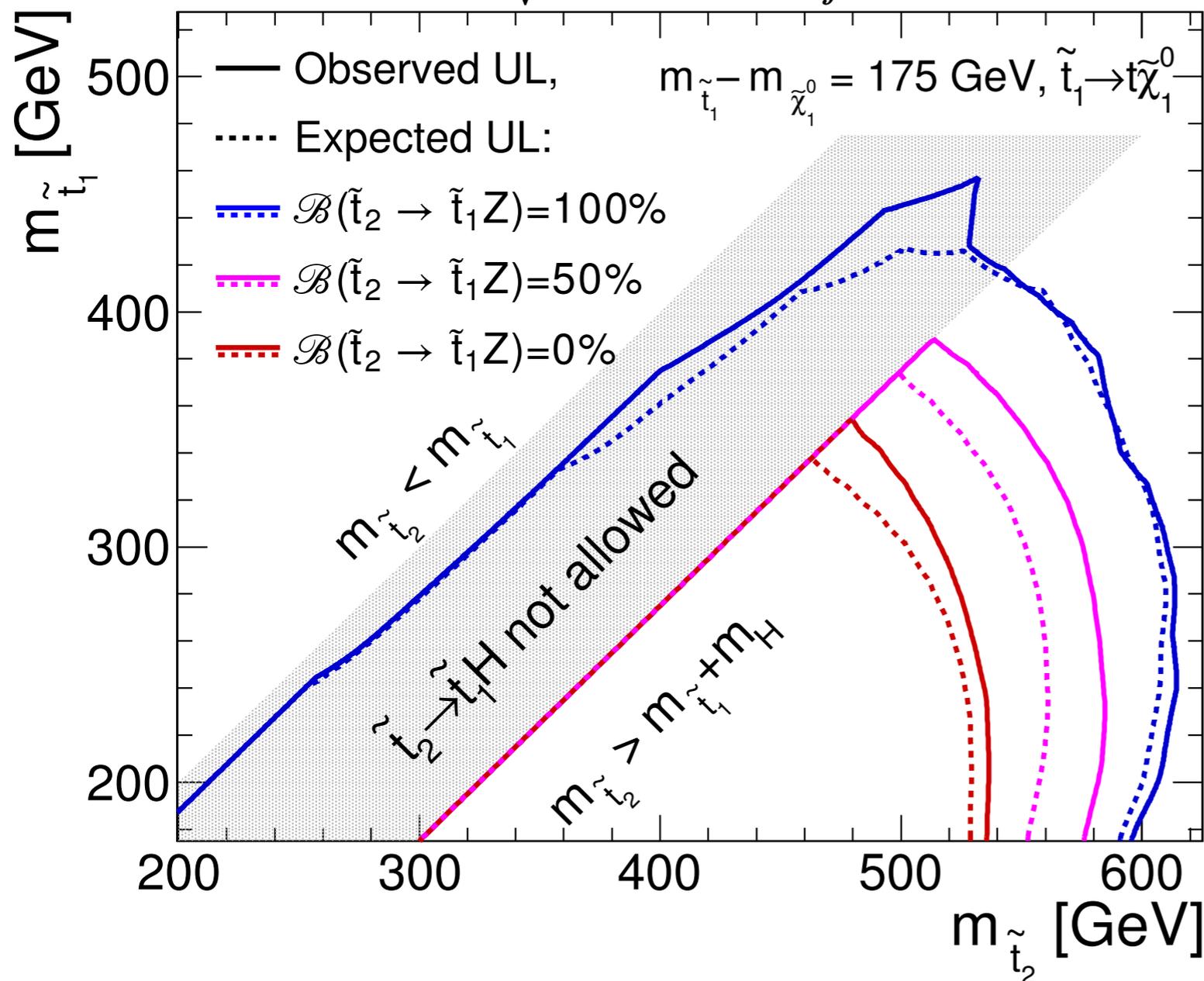
- sensitive to all scenarios
- background from WZ+jets is suppressed by b-tag requirements,
- remaining background from rare processes or non-prompt leptons
 - WZ/ZZ from MC
 - non-prompt lepton from data





stop2 to stop1 Z/H

CMS $\sqrt{s} = 8 \text{ TeV}, \int \mathcal{L} dt = 19.5 \text{ fb}^{-1}$



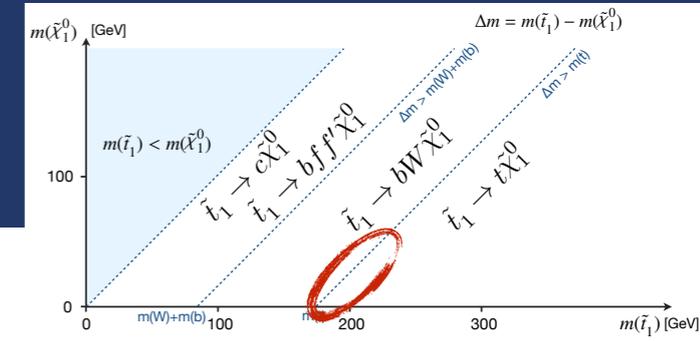
- $m(\text{stop1}) - m(\text{N1}) \approx m(\text{t})$
- 3L analysis most sensitive



ATLAS searches

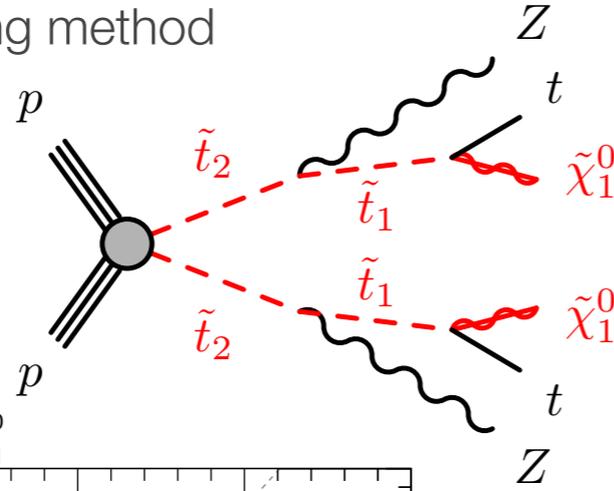
stop2 to Z

Eur. Phys. J. C (2014) 74:2883

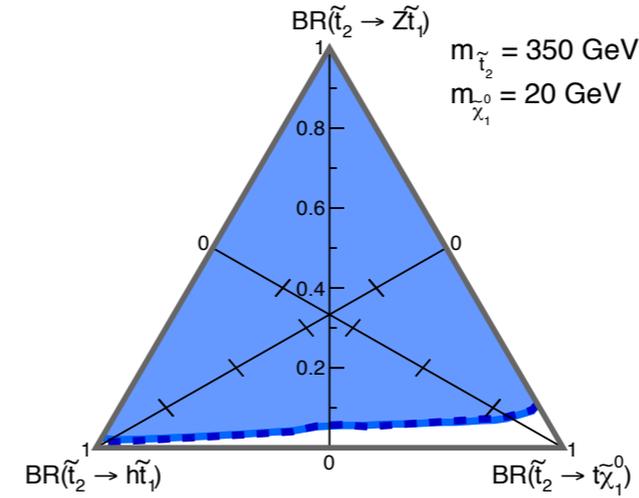


2L + 3L with Z(ll) + jets + b-jet + E_T^{miss}

- Background estimates:
 - Fake leptons using “matrix method”
 - Z+jets using jet smearing method



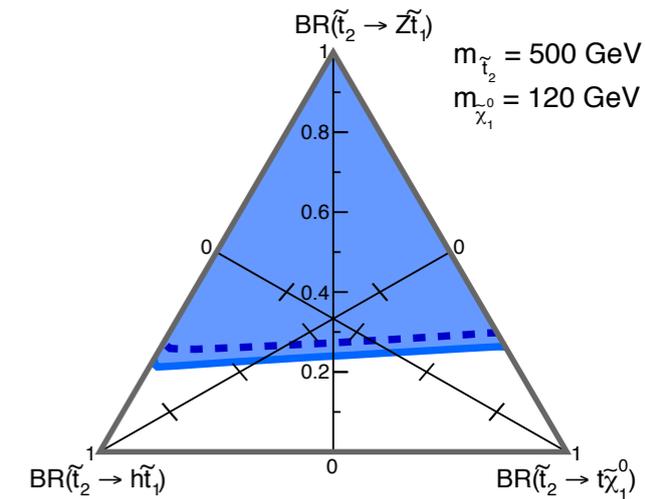
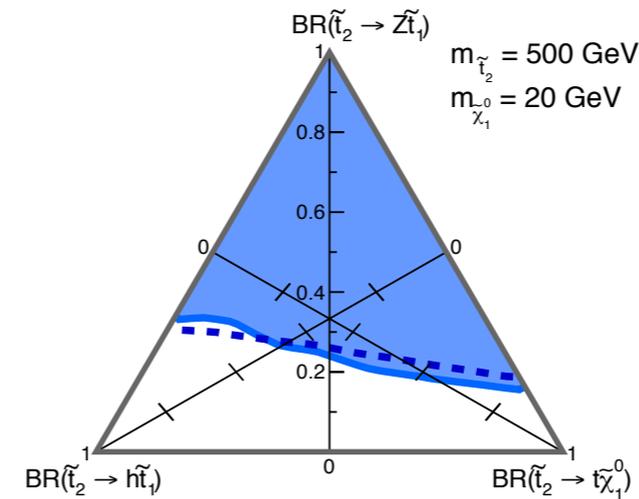
stop2 to (stop1 Z, stop1 H, t N1)



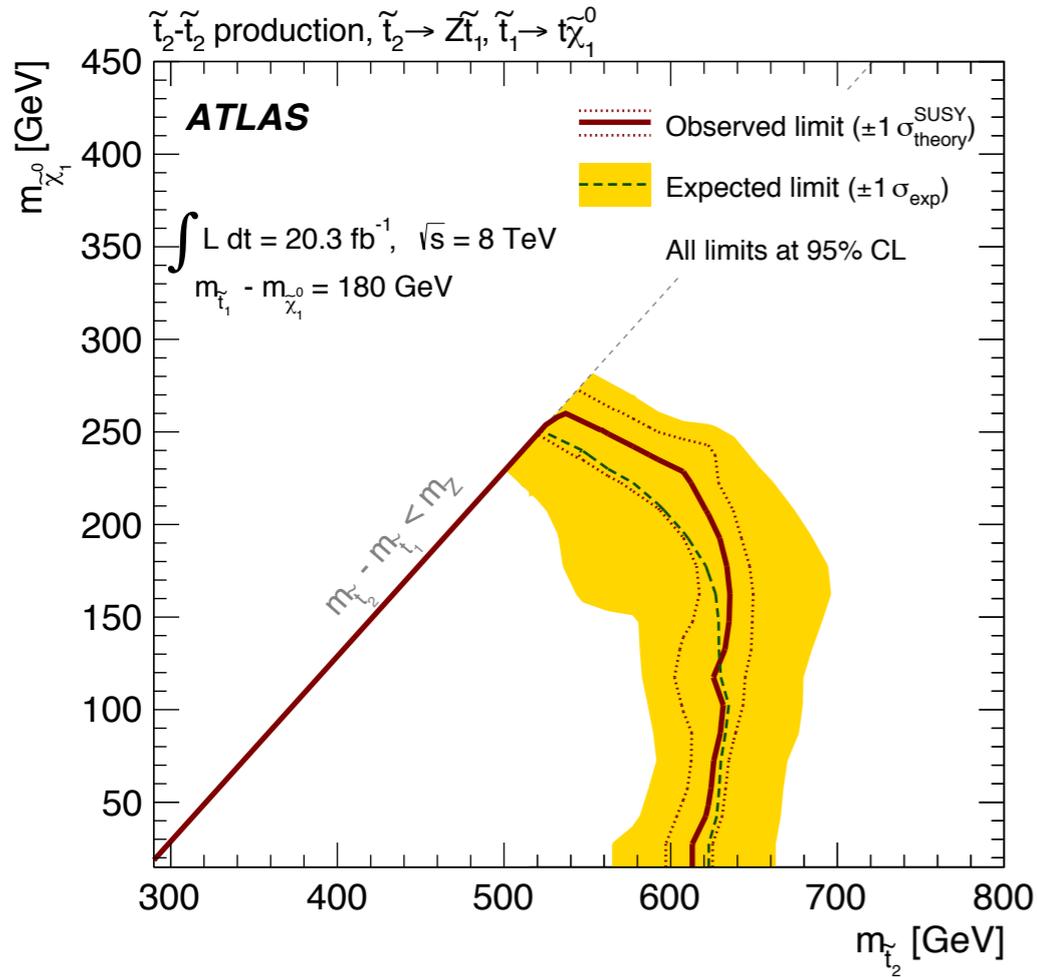
ATLAS $L_{\text{int}} = 20.3 \text{ fb}^{-1}$ $\sqrt{s} = 8 \text{ TeV}$

\tilde{t}_2 - \tilde{t}_2 production, $\tilde{t}_2 \rightarrow Z\tilde{t}_1, h\tilde{t}_1, t\tilde{\chi}_1^0$; $\tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$
 $m_{\tilde{t}_1} = m_{\tilde{\chi}_1^0} + 180 \text{ GeV}$

Observed
Expected



stop2 to stop1 Z



Constraints from top measurements

► In the context of $t\bar{T}$ measurements,

- Stop-pair production with $\text{stop} \rightarrow \text{top} + \text{neutralino}$ (invisible) could be an annoying background
- or perhaps an interesting signal.

(1) What is the effect of stop pair production with $m(\text{stop}) \approx m(\text{top})$ on cross-section analysis?

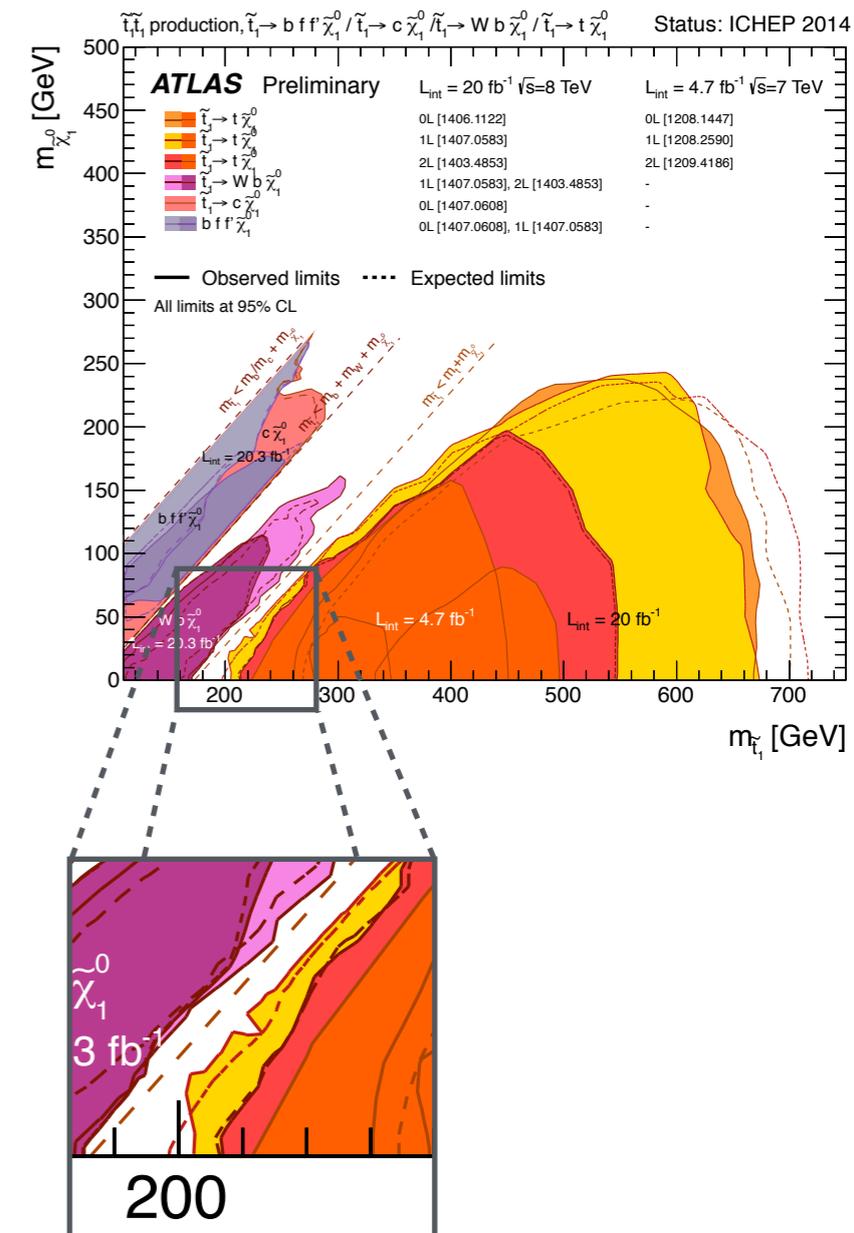
➡ Extra contribution to $t\bar{T}$ production with soft missing energy – otherwise similar kinematics.

(2) What is the effect of stop pair production with $m(\text{stop}) \approx m(\text{top})$ on spin correlation analysis?

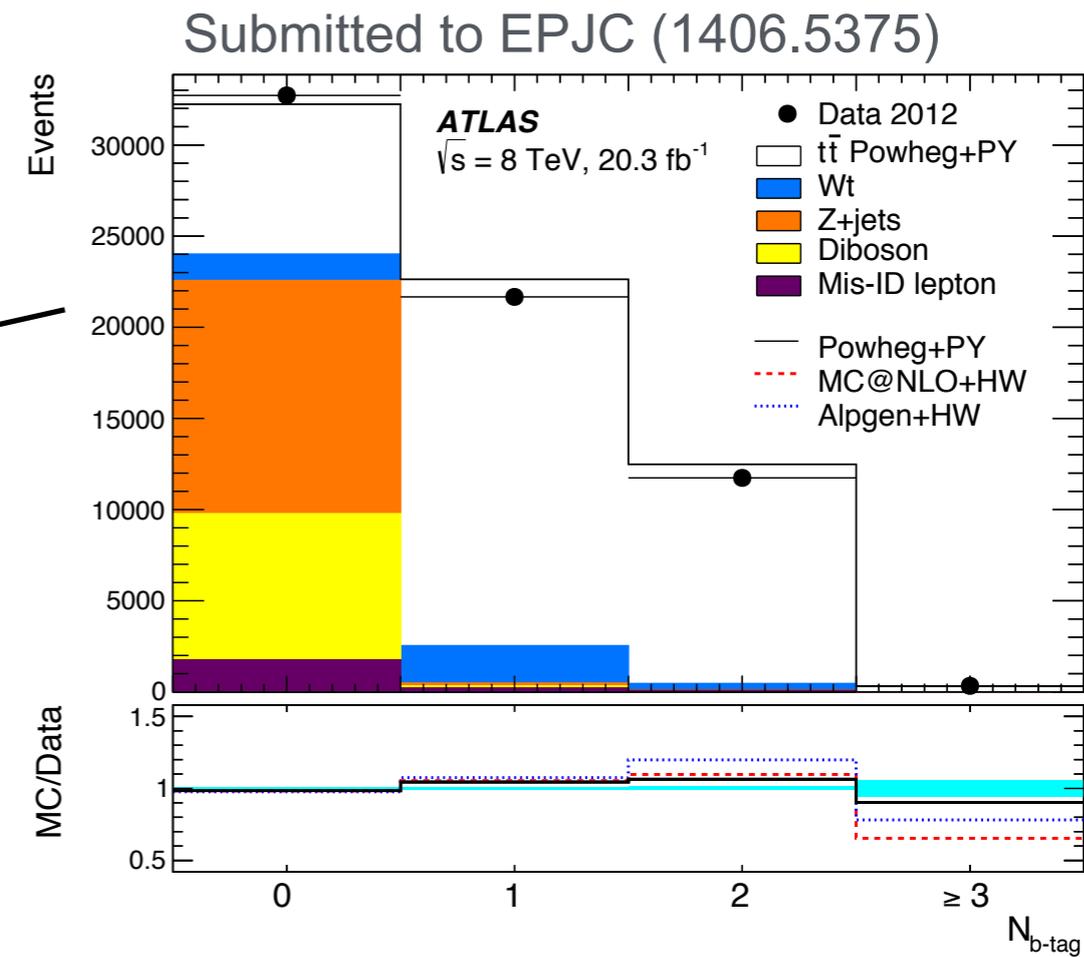
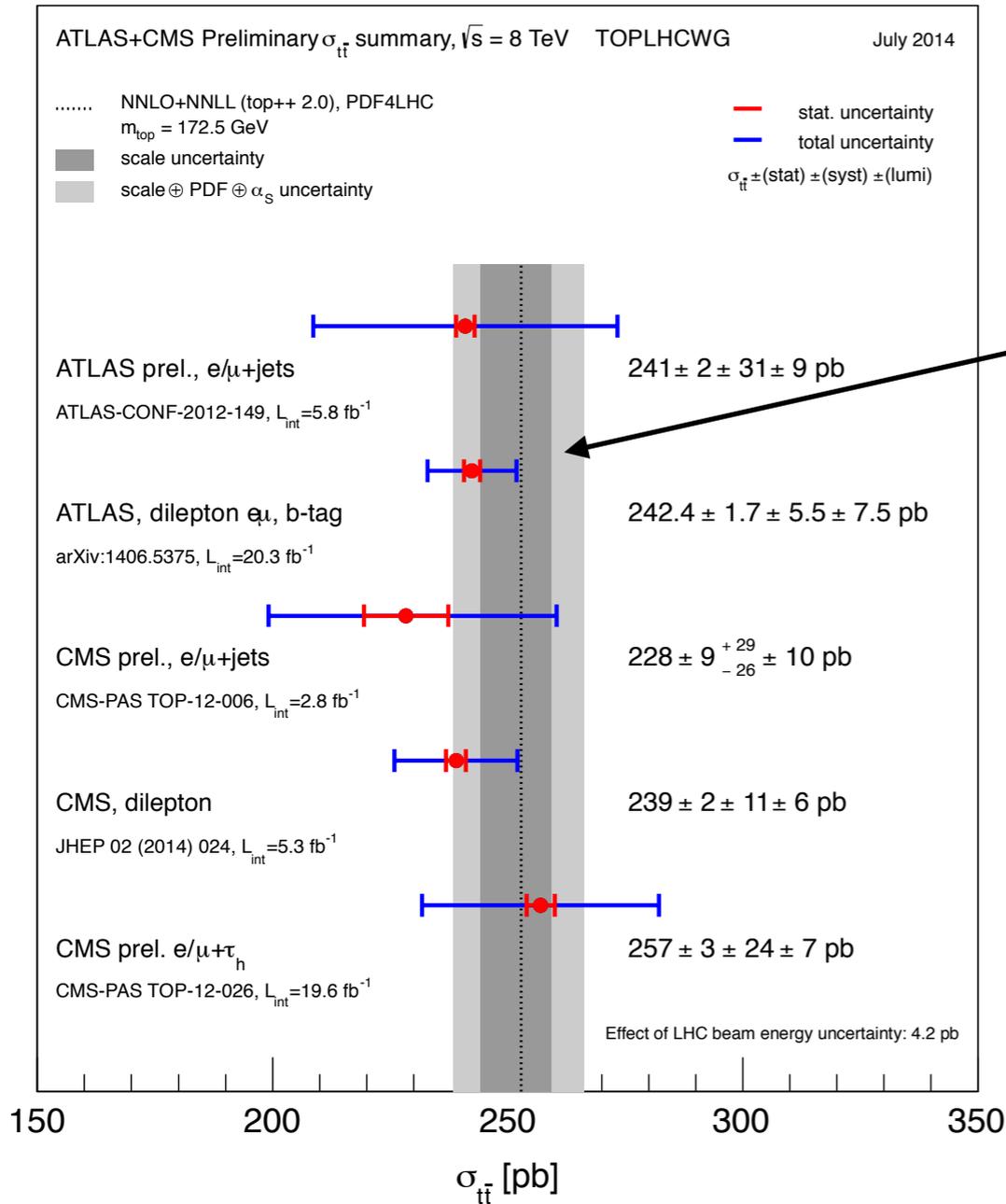
➡ Extra contribution without spin correlation (stop is scalar).

► Could we see or exclude this contribution with stop mass $< \sim 200$ GeV

- E.g. 180 GeV stop-pair contribution is 32 pb at 8 TeV c.f. QCD $t\bar{T}$ of ~ 250 pb



Limit from $t\bar{t}$ cross section



Measurement described in talk by
 Javier Brochero Cifuentes (Monday)

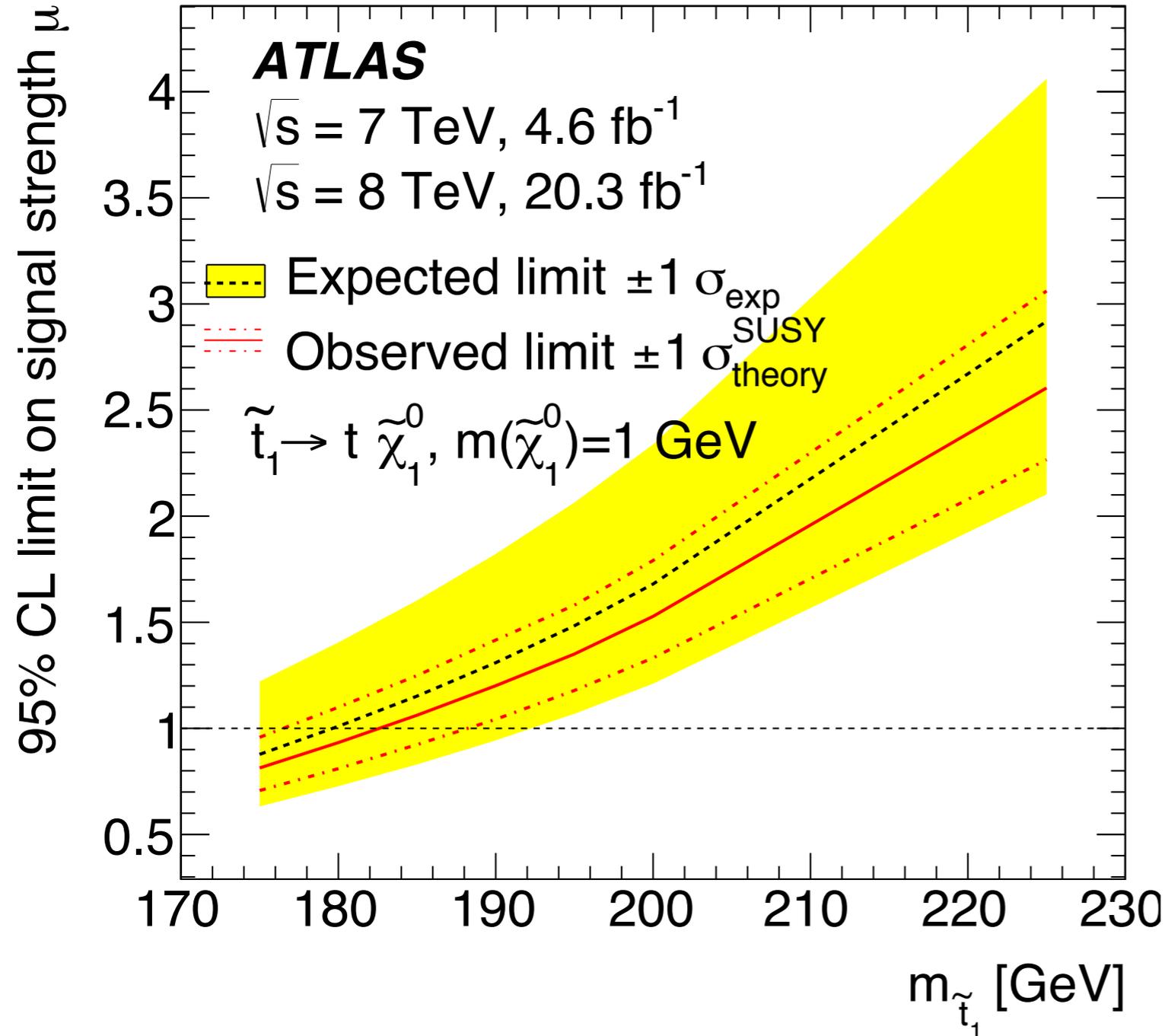
$$\sigma_{t\bar{t}} = 182.9 \pm 3.1 \pm 4.2 \pm 3.6 \pm 3.3 \text{ pb} (\sqrt{s} = 7 \text{ TeV}) \text{ and}$$

$$\sigma_{t\bar{t}} = 242.4 \pm 1.7 \pm 5.5 \pm 7.5 \pm 4.2 \text{ pb} (\sqrt{s} = 8 \text{ TeV}),$$

- Precise measurement of $t\bar{t}$ production cross-section at 7 and 8 TeV
 - events with $e\mu$ (opp-sign) and exactly 1 or 2 b-tagged jets, ratio reduces systematics.
 - Theory (including Δm_{top}) $177.3^{+11.5}_{-12.0} \text{ pb}$ (7 TeV), $252.9^{+15.3}_{-16.3} \text{ pb}$ (8 TeV)

Limit from $t\bar{t}b\bar{a}$ cross section

Submitted to EPJC (1406.5375)



- ▶ Exclude a small region between top threshold and $\sim 180 \text{ GeV}$
 - Expected limit 179 GeV
 - Observed limit 183 GeV
 - Observed limit -1σ (sig-theo) 177 GeV
 - Signal theoretical uncertainty of 15% (PDFs, scales)
 - Sensitivity driven by stop-pair production cross-section
 - Always assuming 100% BR to stop+ χ
- ▶ Limits shown for $m(\chi) = 1 \text{ GeV}$
 - Only slightly weaker with increasing neutralino mass
 - E.g. 3% less for 20 GeV neutralino at 200 GeV stop
 - top quark polarisation (stop mixing)
 - full left-handed polarisation 4% weaker than plot which is for mostly right-handed case

CMS $t\bar{t}$ cross section measurement (see previous slide) is not yet interpreted in terms of stop limits.

► Precise measurement of $t\bar{t}$ spin correlation at 8 TeV

- events with $ee, e\mu, \mu\mu$ (opp-sign)
- at least 2 jets, of which at least one is b-tagged
- in $ee, \mu\mu$ channels: MET, $m(l\bar{l})$ cuts
- in $e\mu$ channel: HT(jets, lep) cut

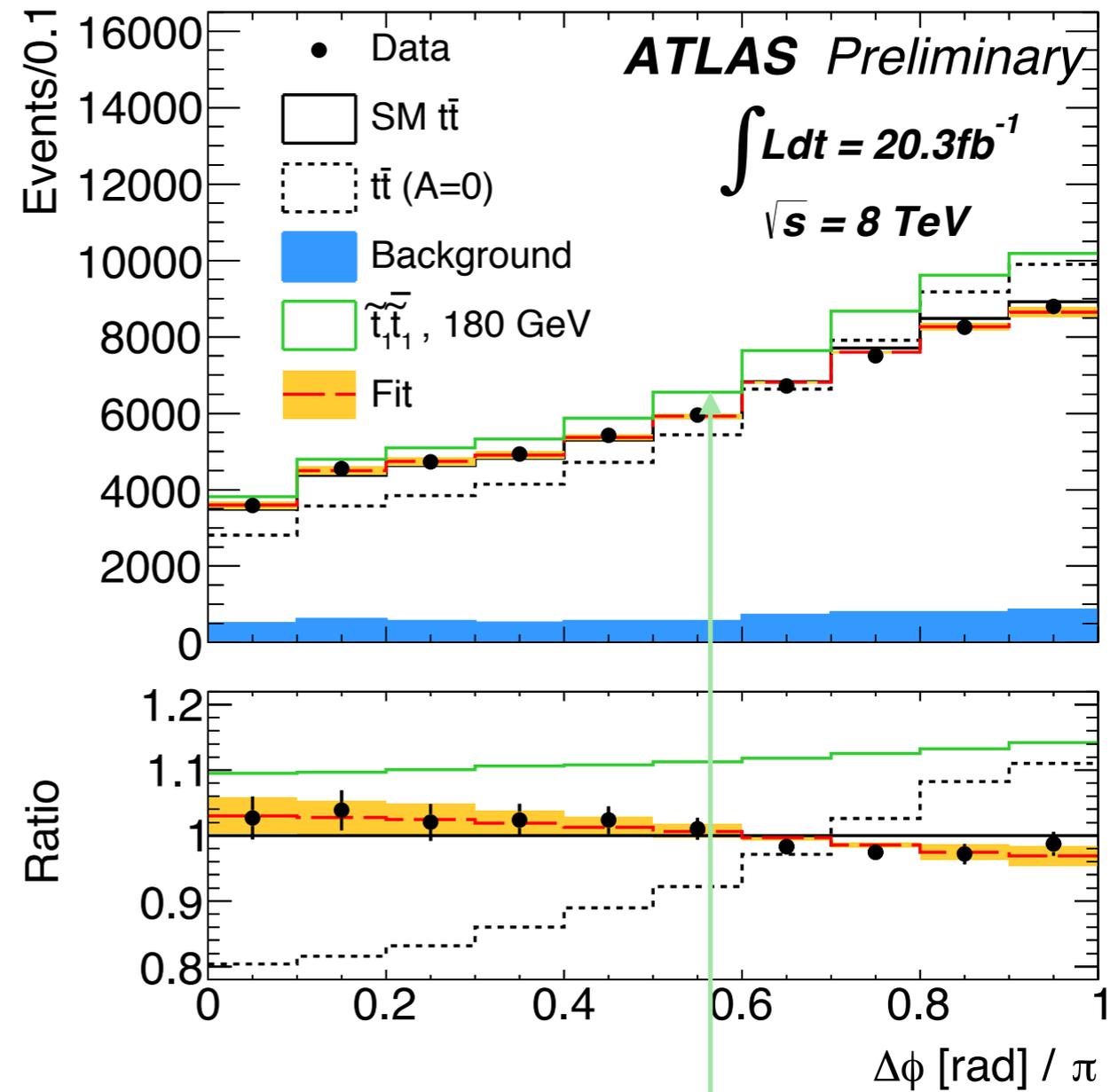
► Investigate $\Delta\phi(l,l)$ distribution

► Results

$$A = \frac{N_{\uparrow\uparrow} + N_{\down\downarrow} - N_{\uparrow\downarrow} - N_{\down\uparrow}}{N_{\uparrow\uparrow} + N_{\down\downarrow} + N_{\uparrow\downarrow} + N_{\down\uparrow}}$$

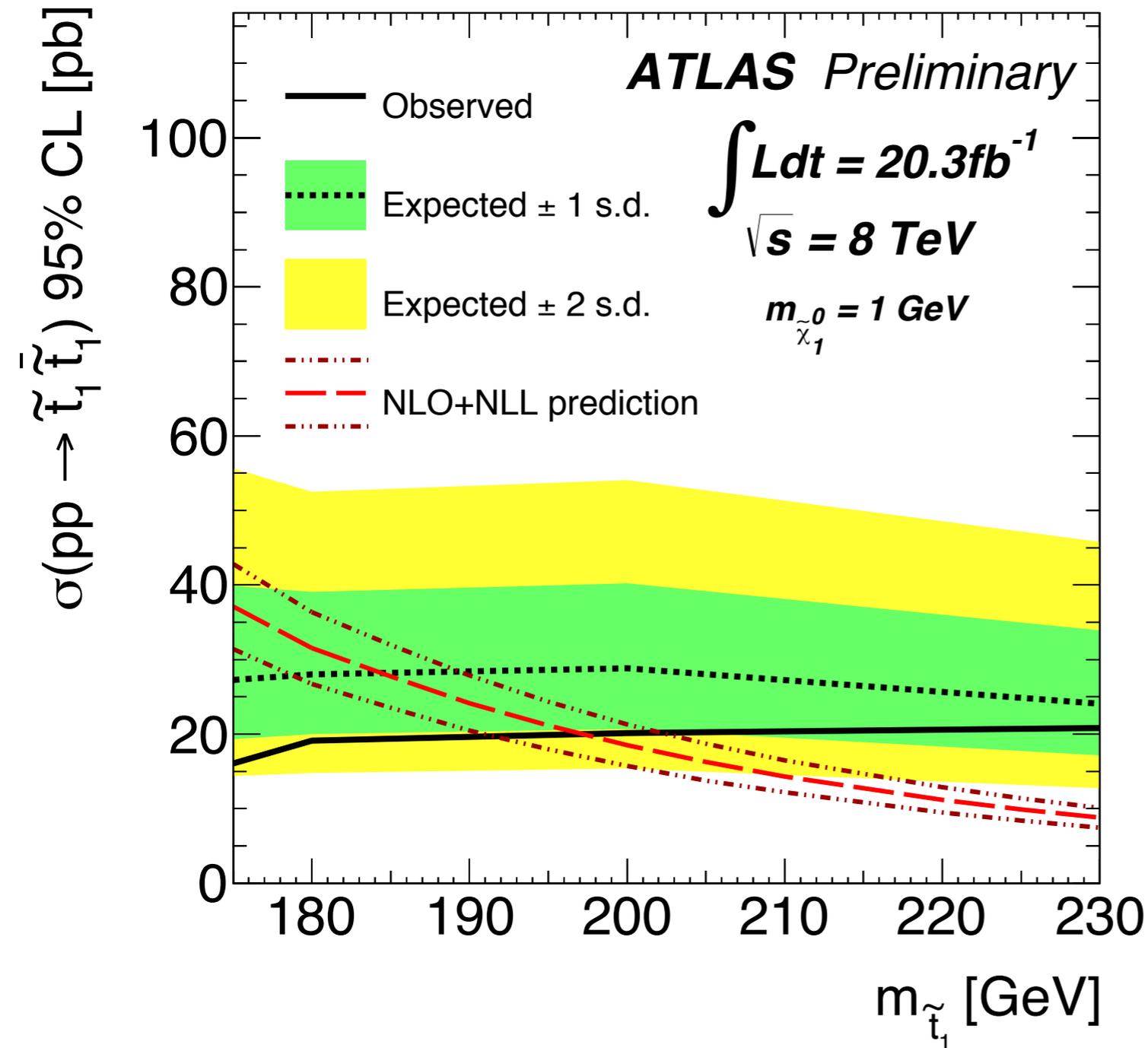
$$A_{\text{helicity}} = 0.38 \pm 0.04$$

$$A_{\text{helicity}}^{\text{SM}} = 0.318 \pm 0.005$$



SM+ SUSY (180 GeV mStop)

Measurement described in talk by Richard Hawkings (Tuesday)



- ▶ Exclude a small region between top threshold and ~ 190 GeV
 - Expected limit 185 GeV
 - Observed limit 197 GeV
 - Observed limit -1σ (sig-theo) 191 GeV
 - Signal theoretical uncertainty of 15% (PDFs, scales)
 - Sensitivity driven by stop-pair production cross-section
 - Always assuming 100% BR to stop+ χ
- ▶ Limits shown for $m(\chi) = 1$ GeV
 - Only slightly weaker with increasing neutralino mass
- ▶ Without the shape-information $\Delta\phi$ the limit would deteriorate by 30-40%.
 - constraint from cross section only

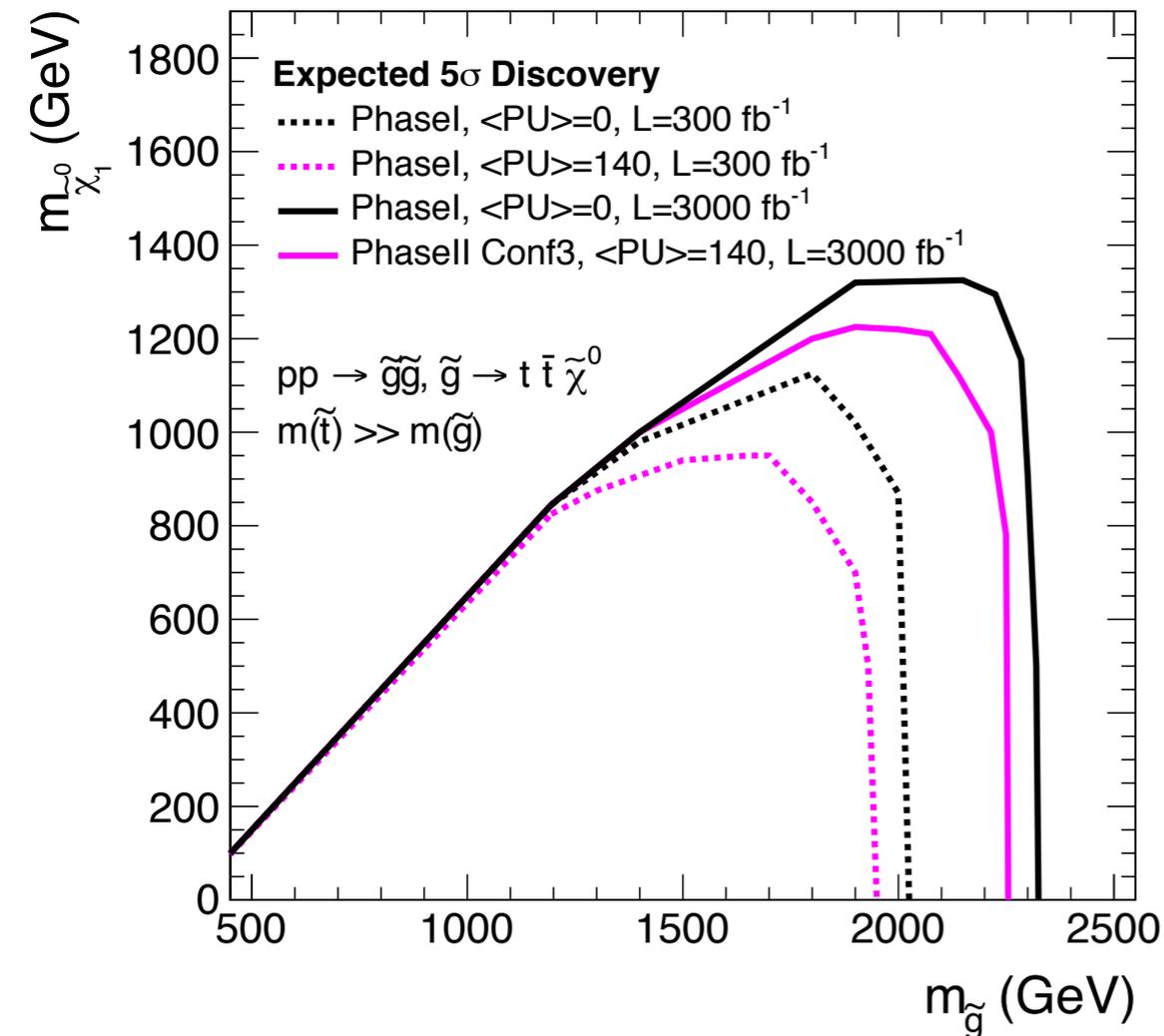
Long-term prospects

CMS and ATLAS studied long-term prospects for the (HL) LHC.

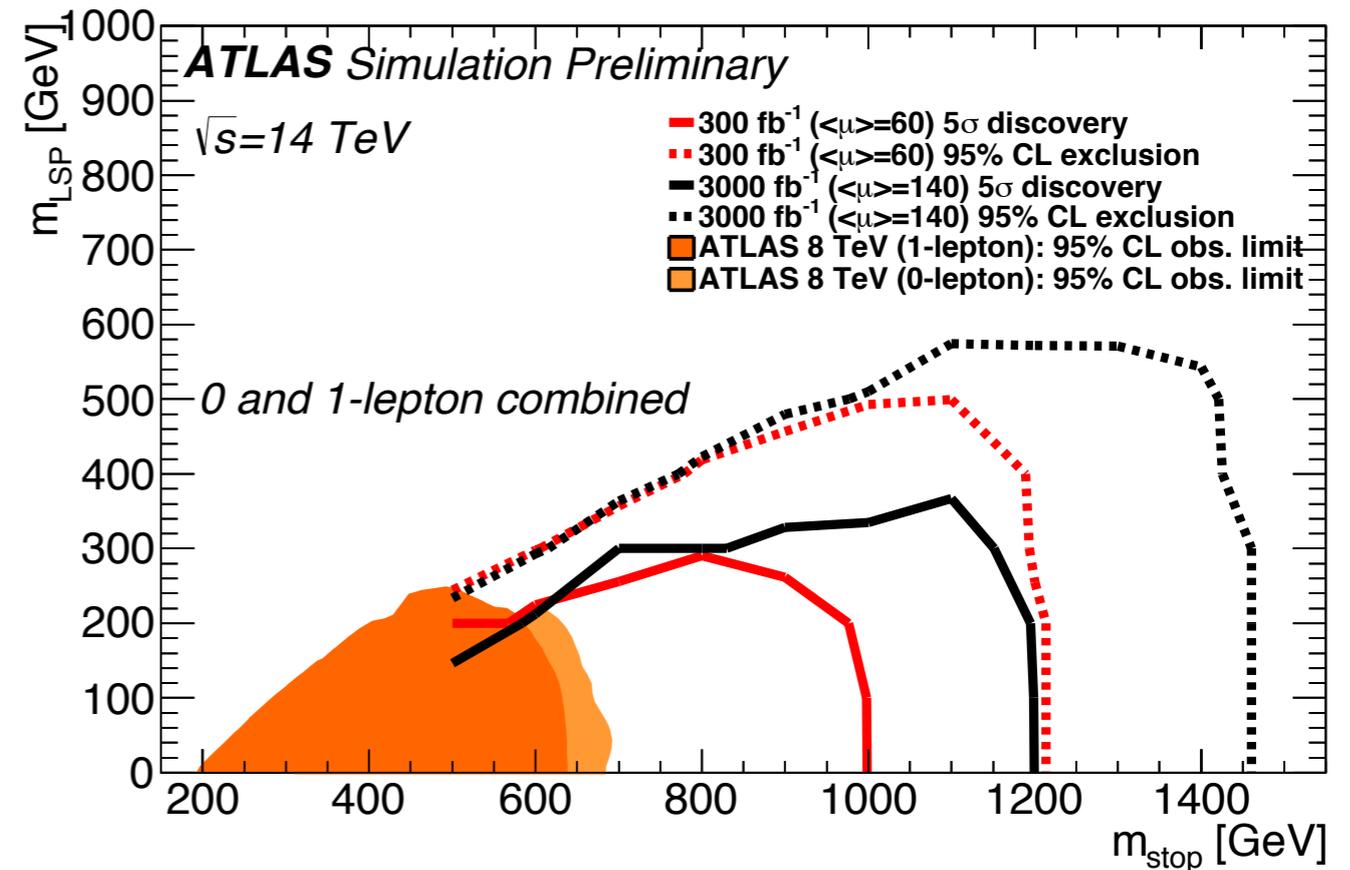
- with 300 and 3000 /fb at 14 TeV
- searches for gluino-mediated stop production reach beyond 2 TeV
- searches for direct stop production reach well beyond 1 TeV

CMS PAS FTR-13-014

CMS Simulation $\sqrt{s} = 14 \text{ TeV}$

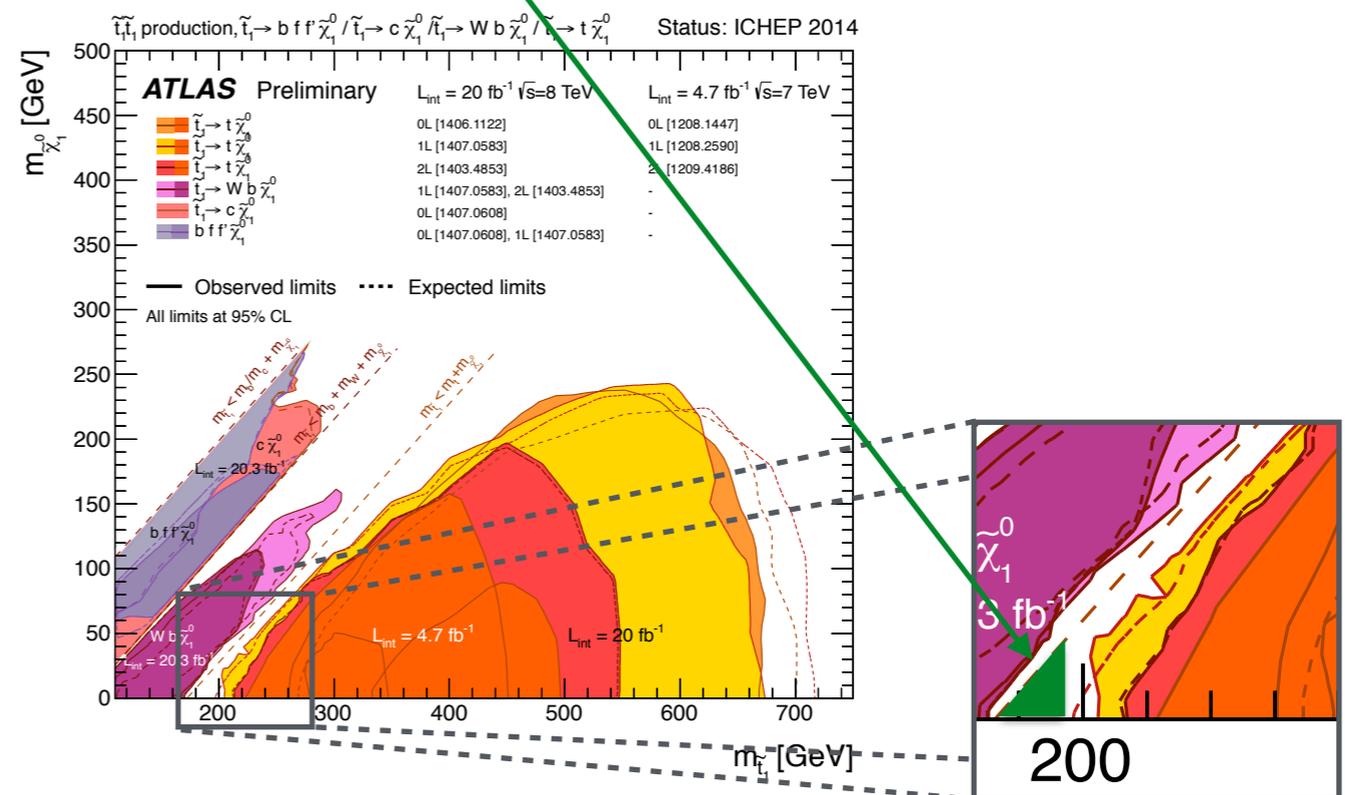
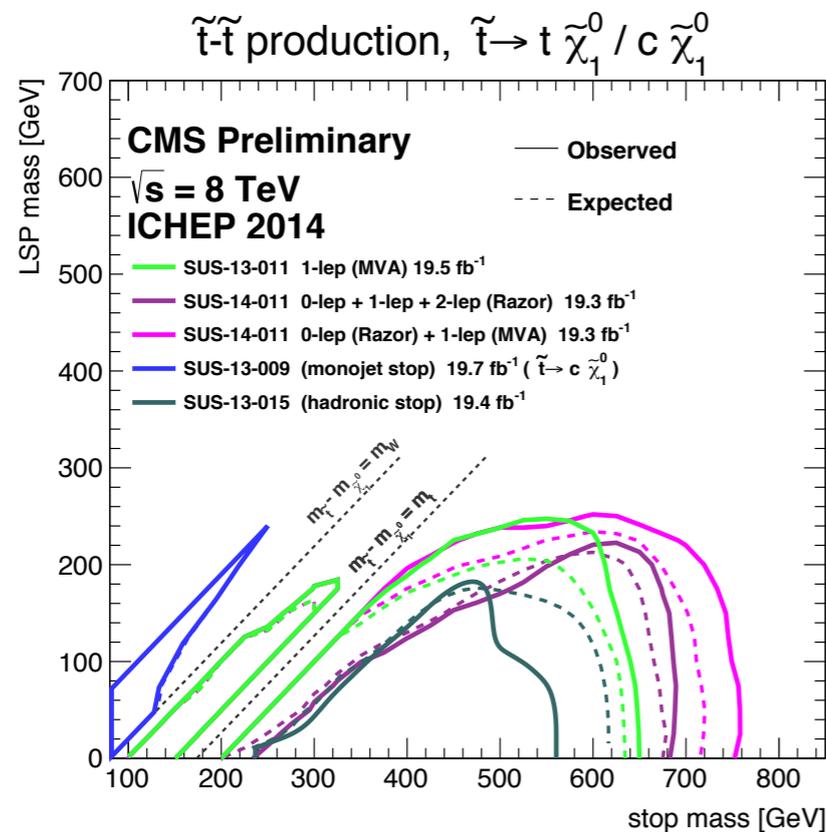


ATL-PHYS-PUB-2013-011



Summary

- ▶ Strong search programs for a top quark partner (stop) at the ATLAS and CMS experiments.
 - direct searches for stop covering range from low to heavy masses, various decay modes
 - new indirect constraints from SM tT measurements begin to fill gap in $m_{\text{Stop}} \sim m_{\text{Top}}$ scenario

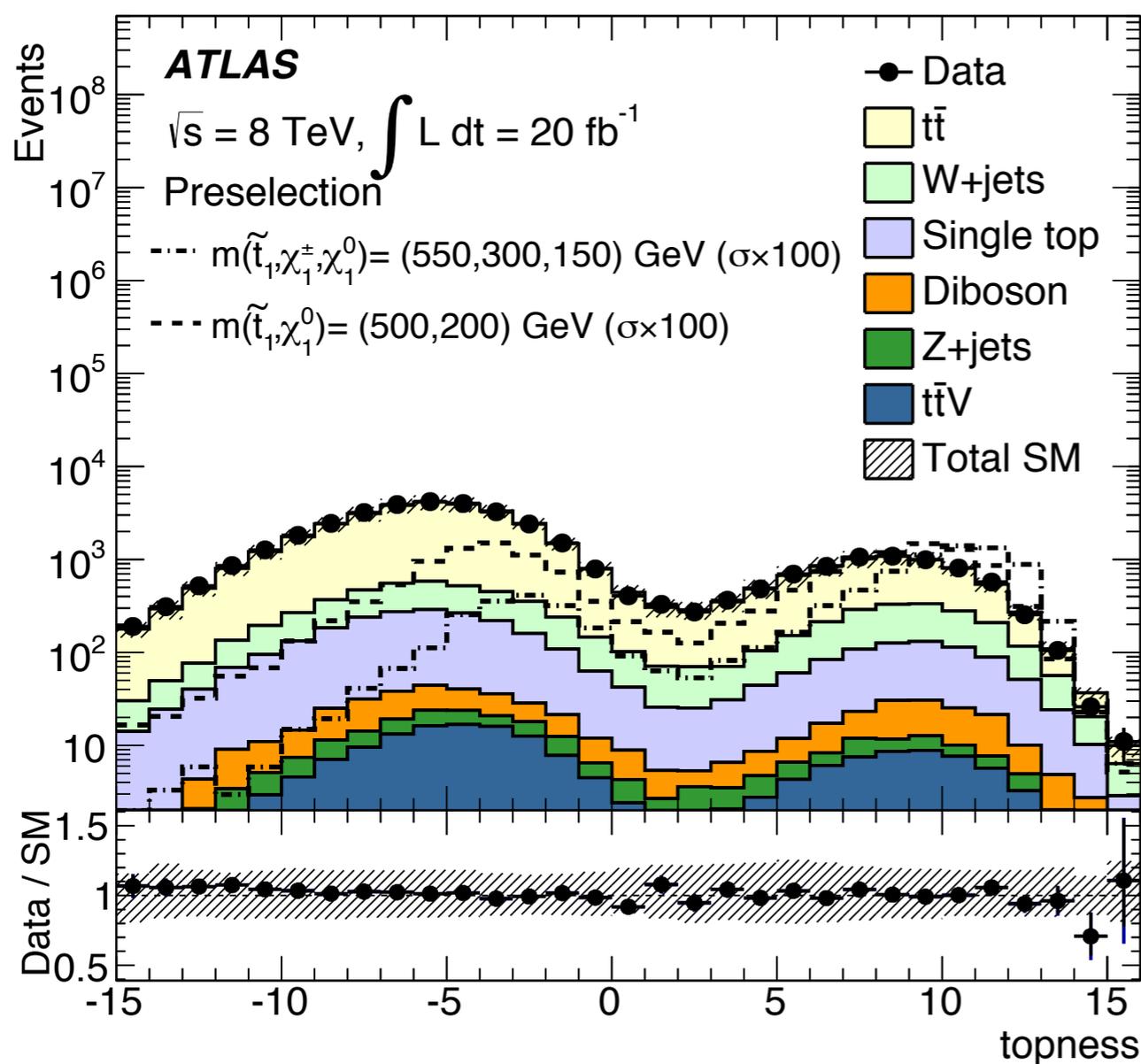


Extra slides

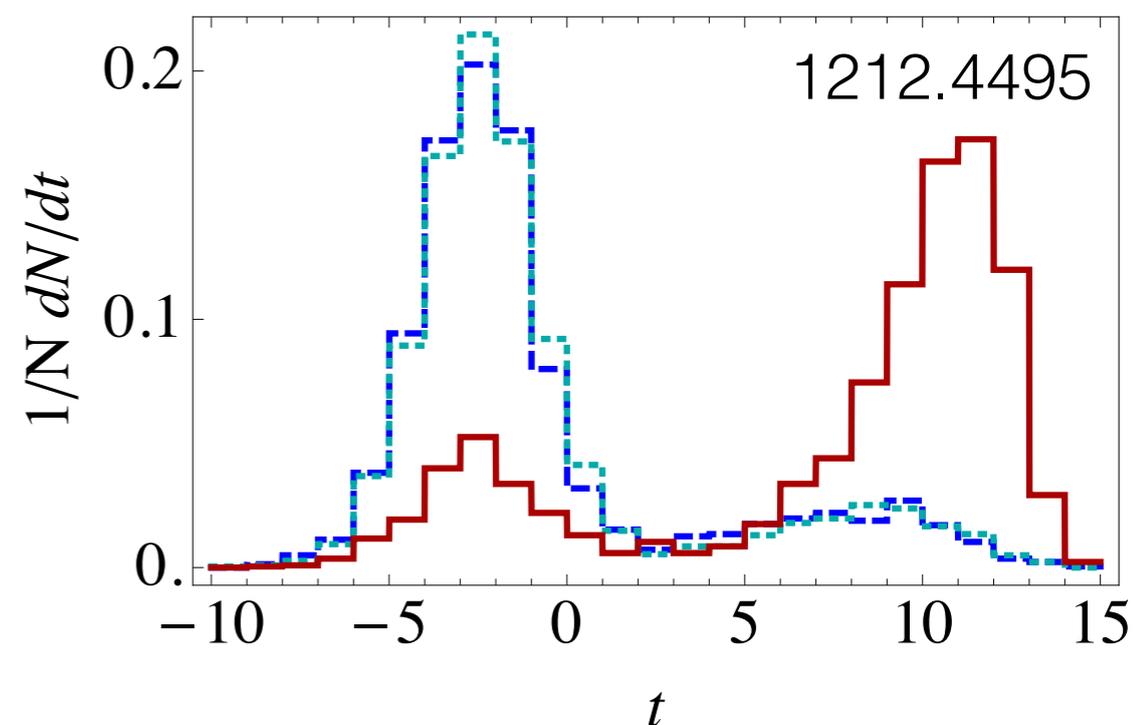
Topness

Topness is a variable designed to identify and suppress partially reconstructed dileptonic $t\bar{T}$ events, as proposed in ref. [1212.4495]. The topness variable is based on minimising a chi2-type function indicating the similarity of the event to dileptonic $t\bar{T}$ events. One lepton is assumed to be lost.

=1L, $\geq 4j$, MET > 100 GeV, MT > 60 GeV



=1L, $\geq 4j$, MET > 100 GeV, MT > 150 GeV



- signal, red, solid
- dileptonic top, blue, dashed
- dileptonic top, one (e, mu), one tau, cyan, dotted

Topness

The *topness* event value is defined as $\ln(\min \hat{S})$, where \hat{S} is the minimum of the χ^2 -type function S :

$$S(p_{W,x}, p_{W,y}, p_{W,z}, p_{\nu,z}) = \frac{\left(m_W^2 - (p_\ell + p_\nu)^2\right)^2}{a_W^4} + \frac{\left(m_t^2 - (p_{b_1} + p_\ell + p_\nu)^2\right)^2}{a_t^4} + \frac{\left(m_t^2 - (p_{b_2} + p_W)^2\right)^2}{a_t^4} + \frac{\left(4m_t^2 - (\Sigma p)^2\right)^2}{a_{\text{CM}}^4}.$$

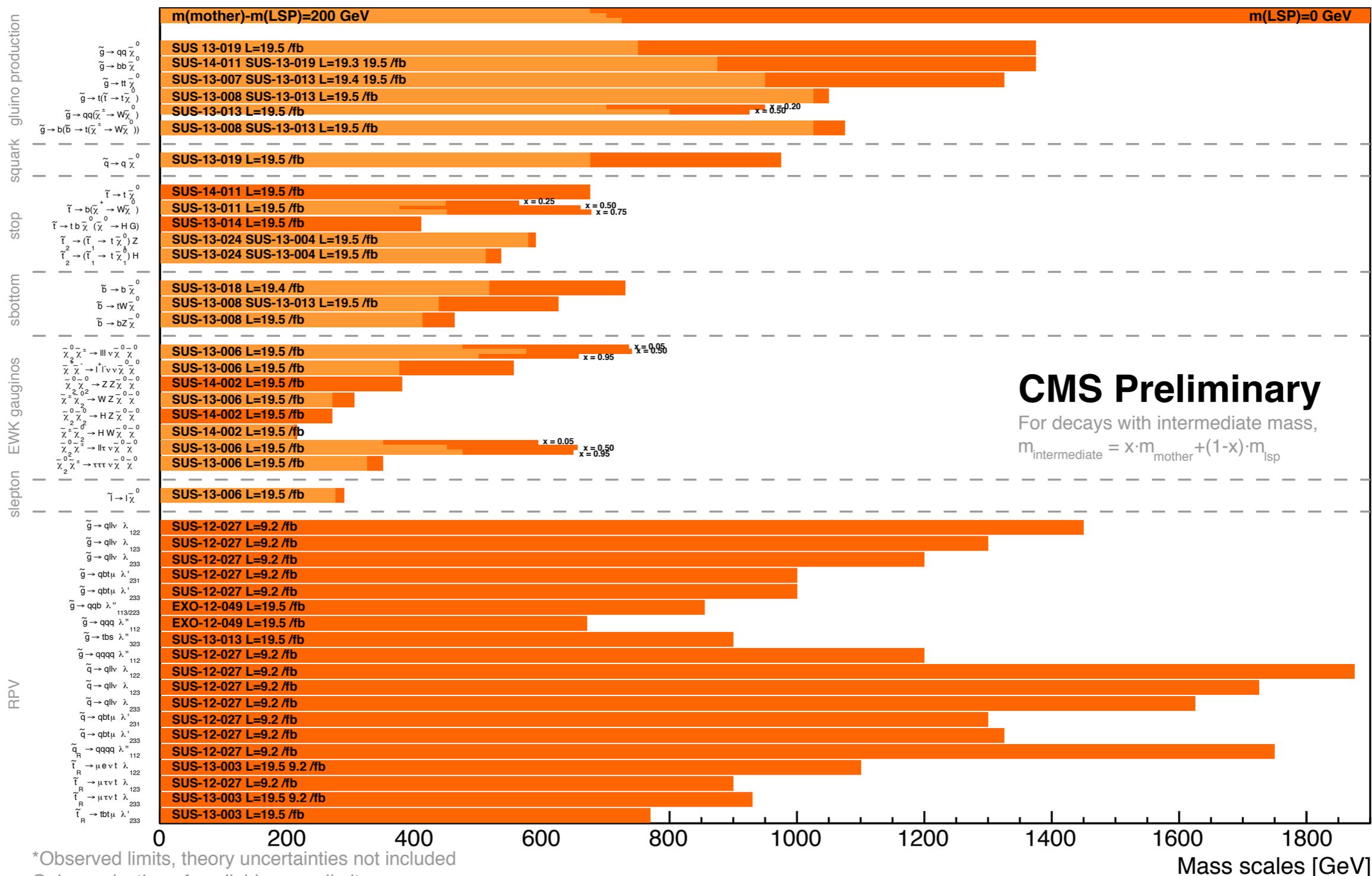
The first three arguments of S are the components of the non-reconstructed W boson 3-momentum $(p_{W,x}, p_{W,y}, p_{W,z})$. This W is assumed to decay leptonically, but the lepton is not reconstructed and is thus only noticeable in the missing transverse momentum. The variable $p_{\nu,z}$ is the longitudinal momentum of the neutrino from the other W boson decay, for which the lepton was successfully reconstructed. These four numbers are varied to find the minimum of S .

The momenta appearing on the right-hand side of the equation above are either 4-momenta of the reconstructed objects (one lepton, p_ℓ , and two b -jets, p_{b_1} and p_{b_2}) or 4-momenta assigned by the minimisation procedure (p_W and p_ν). To find all four components, the neutrinos and the W boson without reconstructed decay products are assumed to be on-shell. Both combinations for b_1 and b_2 are evaluated during the minimisation; if only one b -tagged jet is present, it is used together with the leading or subleading jet (that means, a total of four possible jet assignments is evaluated in this case).

LHC Run-I SUSY Searches: CMS summary table

Summary of CMS SUSY Results* in SMS framework

ICHEP 2014



CMS Preliminary

For decays with intermediate mass,
 $m_{\text{intermediate}} = x \cdot m_{\text{mother}} + (1-x) \cdot m_{\text{LSP}}$

*Observed limits, theory uncertainties not included
 Only a selection of available mass limits
 Probe *up to* the quoted mass limit

LHC Run-I SUSY Searches: ATLAS summary table

ATLAS SUSY Searches* - 95% CL Lower Limits

Status: ICHEP 2014

ATLAS Preliminary

$\sqrt{s} = 7, 8 \text{ TeV}$

Model	e, μ, τ, γ	Jets	E_T^{miss}	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Mass limit	Reference		
Inclusive Searches	MSUGRA/CMSSM	0	2-6 jets	Yes	20.3	\tilde{q}, \tilde{g} 1.7 TeV	$m(\tilde{q})=m(\tilde{g})$	1405.7875
	MSUGRA/CMSSM	1 e, μ	3-6 jets	Yes	20.3	\tilde{g} 1.2 TeV	any $m(\tilde{q})$	ATLAS-CONF-2013-062
	MSUGRA/CMSSM	0	7-10 jets	Yes	20.3	\tilde{g} 1.1 TeV	any $m(\tilde{q})$	1308.1841
	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$	0	2-6 jets	Yes	20.3	\tilde{q} 850 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}, m(1^{\text{st}} \text{ gen. } \tilde{q})=m(2^{\text{nd}} \text{ gen. } \tilde{q})$	1405.7875
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	0	2-6 jets	Yes	20.3	\tilde{g} 1.33 TeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$	1405.7875
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0 \rightarrow q\tilde{q}W^\pm\tilde{\chi}_1^0$	1 e, μ	3-6 jets	Yes	20.3	\tilde{g} 1.18 TeV	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}, m(\tilde{\chi}^\pm)=0.5(m(\tilde{\chi}_1^0)+m(\tilde{g}))$	ATLAS-CONF-2013-062
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell/\ell\nu/\nu\nu)\tilde{\chi}_1^0$	2 e, μ	0-3 jets	-	20.3	\tilde{g} 1.12 TeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$	ATLAS-CONF-2013-089
	GMSB ($\tilde{\ell}$ NLSP)	2 e, μ	2-4 jets	Yes	4.7	\tilde{g} 1.24 TeV	$\tan\beta < 15$	1208.4688
	GMSB ($\tilde{\ell}$ NLSP)	1-2 τ + 0-1 ℓ	0-2 jets	Yes	20.3	\tilde{g} 1.6 TeV	$\tan\beta > 20$	1407.0603
	GGM (bino NLSP)	2 γ	-	Yes	20.3	\tilde{g} 1.28 TeV	$m(\tilde{\chi}_1^0) > 50 \text{ GeV}$	ATLAS-CONF-2014-001
GGM (wino NLSP)	1 e, μ + γ	-	Yes	4.8	\tilde{g} 619 GeV	$m(\tilde{\chi}_1^0) > 50 \text{ GeV}$	ATLAS-CONF-2012-144	
GGM (higgsino-bino NLSP)	γ	1 b	Yes	4.8	\tilde{g} 900 GeV	$m(\tilde{\chi}_1^0) > 220 \text{ GeV}$	1211.1167	
GGM (higgsino NLSP)	2 e, μ (Z)	0-3 jets	Yes	5.8	\tilde{g} 690 GeV	$m(\text{NLSP}) > 200 \text{ GeV}$	ATLAS-CONF-2012-152	
Gravitino LSP	0	mono-jet	Yes	10.5	$F^{1/2}$ scale 645 GeV	$m(\tilde{G}) > 10^{-4} \text{ eV}$	ATLAS-CONF-2012-147	
3 rd gen. \tilde{g} med.	$\tilde{g} \rightarrow b\tilde{b}\tilde{\chi}_1^0$	0	3 b	Yes	20.1	\tilde{g} 1.25 TeV	$m(\tilde{\chi}_1^0) < 400 \text{ GeV}$	1407.0600
	$\tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0$	0	7-10 jets	Yes	20.3	\tilde{g} 1.1 TeV	$m(\tilde{\chi}_1^0) < 350 \text{ GeV}$	1308.1841
	$\tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0$	0-1 e, μ	3 b	Yes	20.1	\tilde{g} 1.34 TeV	$m(\tilde{\chi}_1^0) < 400 \text{ GeV}$	1407.0600
	$\tilde{g} \rightarrow b\tilde{b}\tilde{\chi}_1^0$	0-1 e, μ	3 b	Yes	20.1	\tilde{g} 1.3 TeV	$m(\tilde{\chi}_1^0) < 300 \text{ GeV}$	1407.0600
3 rd gen. squarks direct production	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$	0	2 b	Yes	20.1	\tilde{b}_1 100-620 GeV	$m(\tilde{\chi}_1^0) < 90 \text{ GeV}$	1308.2631
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow \tilde{\chi}_1^\pm$	2 e, μ (SS)	0-3 b	Yes	20.3	\tilde{b}_1 275-440 GeV	$m(\tilde{\chi}_1^\pm)=2m(\tilde{\chi}_1^0)$	1404.2500
	$\tilde{t}_1\tilde{t}_1$ (light), $\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$	1-2 e, μ	1-2 b	Yes	4.7	\tilde{t}_1 110-167 GeV	$m(\tilde{\chi}_1^0)=55 \text{ GeV}$	1208.4305, 1209.2102
	$\tilde{t}_1\tilde{t}_1$ (light), $\tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$	2 e, μ	0-2 jets	Yes	20.3	\tilde{t}_1 130-210 GeV	$m(\tilde{\chi}_1^0)=m(\tilde{t}_1)-m(W)-50 \text{ GeV}, m(\tilde{t}_1) < m(\tilde{\chi}_1^\pm)$	1403.4853
	$\tilde{t}_1\tilde{t}_1$ (medium), $\tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$	2 e, μ	2 jets	Yes	20.3	\tilde{t}_1 215-530 GeV	$m(\tilde{\chi}_1^0)=1 \text{ GeV}$	1403.4853
	$\tilde{t}_1\tilde{t}_1$ (medium), $\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$	0	2 b	Yes	20.1	\tilde{t}_1 150-580 GeV	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}, m(\tilde{\chi}_1^\pm)-m(\tilde{\chi}_1^0)=5 \text{ GeV}$	1308.2631
	$\tilde{t}_1\tilde{t}_1$ (heavy), $\tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$	1 e, μ	1 b	Yes	20	\tilde{t}_1 210-640 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$	1407.0583
	$\tilde{t}_1\tilde{t}_1$ (heavy), $\tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$	0	2 b	Yes	20.1	\tilde{t}_1 260-640 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$	1406.1122
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$	0	mono-jet/c-tag	Yes	20.3	\tilde{t}_1 90-240 GeV	$m(\tilde{t}_1)-m(\tilde{\chi}_1^0) < 85 \text{ GeV}$	1407.0608
	$\tilde{t}_1\tilde{t}_1$ (natural GMSB)	2 e, μ (Z)	1 b	Yes	20.3	\tilde{t}_1 150-580 GeV	$m(\tilde{\chi}_1^0) > 150 \text{ GeV}$	1403.5222
$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$	3 e, μ (Z)	1 b	Yes	20.3	\tilde{t}_2 290-600 GeV	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}$	1403.5222	
EW direct	$\tilde{\ell}_L, \tilde{\ell}_R, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$	2 e, μ	0	Yes	20.3	$\tilde{\ell}$ 90-325 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$	1403.5294
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow \tilde{\ell}\nu(\tilde{\ell}\bar{\nu})$	2 e, μ	0	Yes	20.3	$\tilde{\chi}_1^\pm$ 140-465 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}, m(\tilde{\ell}, \bar{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$	1403.5294
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow \tilde{\tau}\nu(\tilde{\tau}\bar{\nu})$	2 τ	-	Yes	20.3	$\tilde{\chi}_1^\pm$ 100-350 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}, m(\tilde{\tau}, \bar{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$	1407.0350
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0 \rightarrow \tilde{\ell}_L\nu\tilde{\ell}_L(\tilde{\nu}\bar{\nu}), \tilde{\ell}\tilde{\nu}\tilde{\ell}_L(\tilde{\nu}\bar{\nu})$	3 e, μ	0	Yes	20.3	$\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ 700 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)=0, m(\tilde{\ell}, \bar{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$	1402.7029
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0 Z^0$	2-3 e, μ	0	Yes	20.3	$\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ 420 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)=0, \text{sleptons decoupled}$	1403.5294, 1402.7029
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0 h\tilde{\chi}_1^0$	1 e, μ	2 b	Yes	20.3	$\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ 285 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)=0, \text{sleptons decoupled}$	ATLAS-CONF-2013-093
	$\tilde{\chi}_2^0\tilde{\chi}_3^0, \tilde{\chi}_2^0 \rightarrow \tilde{\ell}_R\ell$	4 e, μ	0	Yes	20.3	$\tilde{\chi}_2^0, \tilde{\chi}_3^0$ 620 GeV	$m(\tilde{\chi}_2^0)=m(\tilde{\chi}_3^0), m(\tilde{\chi}_1^0)=0, m(\tilde{\ell}, \bar{\nu})=0.5(m(\tilde{\chi}_2^0)+m(\tilde{\chi}_1^0))$	1405.5086
Long-lived particles	Direct $\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk	1 jet	Yes	20.3	$\tilde{\chi}_1^\pm$ 270 GeV	$m(\tilde{\chi}_1^\pm)-m(\tilde{\chi}_1^0)=160 \text{ MeV}, \tau(\tilde{\chi}_1^\pm)=0.2 \text{ ns}$	ATLAS-CONF-2013-069
	Stable, stopped \tilde{g} R-hadron	0	1-5 jets	Yes	27.9	\tilde{g} 832 GeV	$m(\tilde{\chi}_1^0)=100 \text{ GeV}, 10 \mu\text{s} < \tau(\tilde{g}) < 1000 \text{ s}$	1310.6584
	GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu}) + \tau(e, \mu)$	1-2 μ	-	-	15.9	$\tilde{\chi}_1^0$ 475 GeV	$10 < \tan\beta < 50$	ATLAS-CONF-2013-058
	GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$, long-lived $\tilde{\chi}_1^0$	2 γ	-	Yes	4.7	$\tilde{\chi}_1^0$ 230 GeV	$0.4 < \tau(\tilde{\chi}_1^0) < 2 \text{ ns}$	1304.6310
$\tilde{q}\tilde{q}, \tilde{\chi}_1^0 \rightarrow q\tilde{q}$ (RPV)	1 $\mu, \text{ displ. vtx}$	-	-	20.3	\tilde{q} 1.0 TeV	$1.5 < c\tau < 156 \text{ mm}, \text{BR}(\mu)=1, m(\tilde{\chi}_1^0)=108 \text{ GeV}$	ATLAS-CONF-2013-092	
RPV	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e + \mu$	2 e, μ	-	-	4.6	$\tilde{\nu}_\tau$ 1.61 TeV	$\lambda_{311}^2=0.10, \lambda_{132}=0.05$	1212.1272
	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e(\mu) + \tau$	1 e, μ + τ	-	-	4.6	$\tilde{\nu}_\tau$ 1.1 TeV	$\lambda_{311}^2=0.10, \lambda_{1(2)33}=0.05$	1212.1272
	Bilinear RPV CMSSM	2 e, μ (SS)	0-3 b	Yes	20.3	\tilde{q}, \tilde{g} 1.35 TeV	$m(\tilde{q})=m(\tilde{g}), c\tau_{LSP} < 1 \text{ mm}$	1404.2500
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^\pm \rightarrow ee\tilde{\nu}_e, e\mu\tilde{\nu}_e$	4 e, μ	-	Yes	20.3	$\tilde{\chi}_1^\pm$ 750 GeV	$m(\tilde{\chi}_1^0) > 0.2 \times m(\tilde{\chi}_1^\pm), \lambda_{121} \neq 0$	1405.5086
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^\pm \rightarrow \tau\tau\tilde{\nu}_e, e\tau\tilde{\nu}_\tau$	3 e, μ + τ	-	Yes	20.3	$\tilde{\chi}_1^\pm$ 450 GeV	$m(\tilde{\chi}_1^0) > 0.2 \times m(\tilde{\chi}_1^\pm), \lambda_{133} \neq 0$	1405.5086
	$\tilde{g} \rightarrow qq\tilde{q}$	0	6-7 jets	-	20.3	\tilde{g} 916 GeV	$\text{BR}(\tau)=\text{BR}(b)=\text{BR}(c)=0\%$	ATLAS-CONF-2013-091
$\tilde{g} \rightarrow \tilde{t}_1 t, \tilde{t}_1 \rightarrow bs$	2 e, μ (SS)	0-3 b	Yes	20.3	\tilde{g} 850 GeV		1404.250	
Other	Scalar gluon pair, $\text{sgluon} \rightarrow q\tilde{q}$	0	4 jets	-	4.6	sgluon 100-287 GeV	incl. limit from 1110.2693	1210.4826
	Scalar gluon pair, $\text{sgluon} \rightarrow t\tilde{t}$	2 e, μ (SS)	2 b	Yes	14.3	sgluon 350-800 GeV		ATLAS-CONF-2013-051
	WIMP interaction (D5, Dirac χ)	0	mono-jet	Yes	10.5	M^* scale 704 GeV	$m(\chi) < 80 \text{ GeV}, \text{limit of } < 687 \text{ GeV for D8}$	ATLAS-CONF-2012-147

*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1σ theoretical signal cross section uncertainty.

Polarization effects



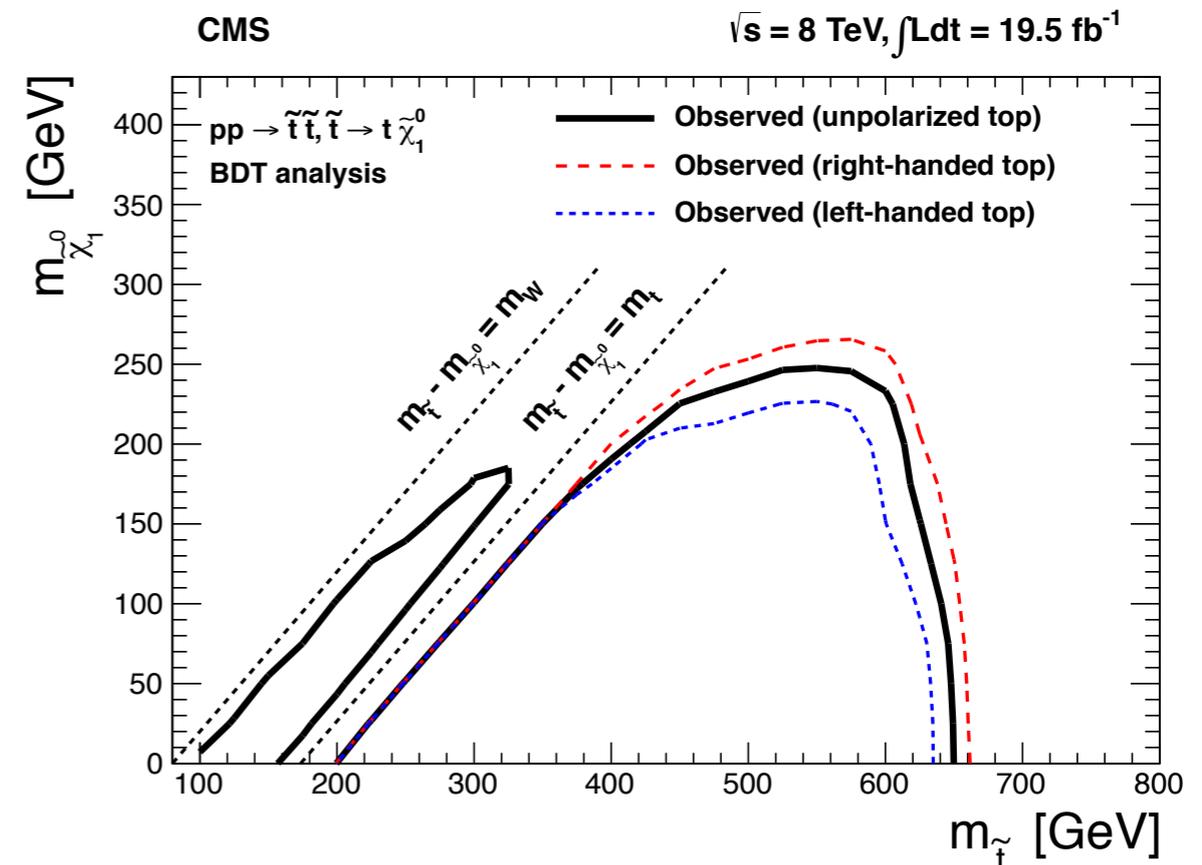
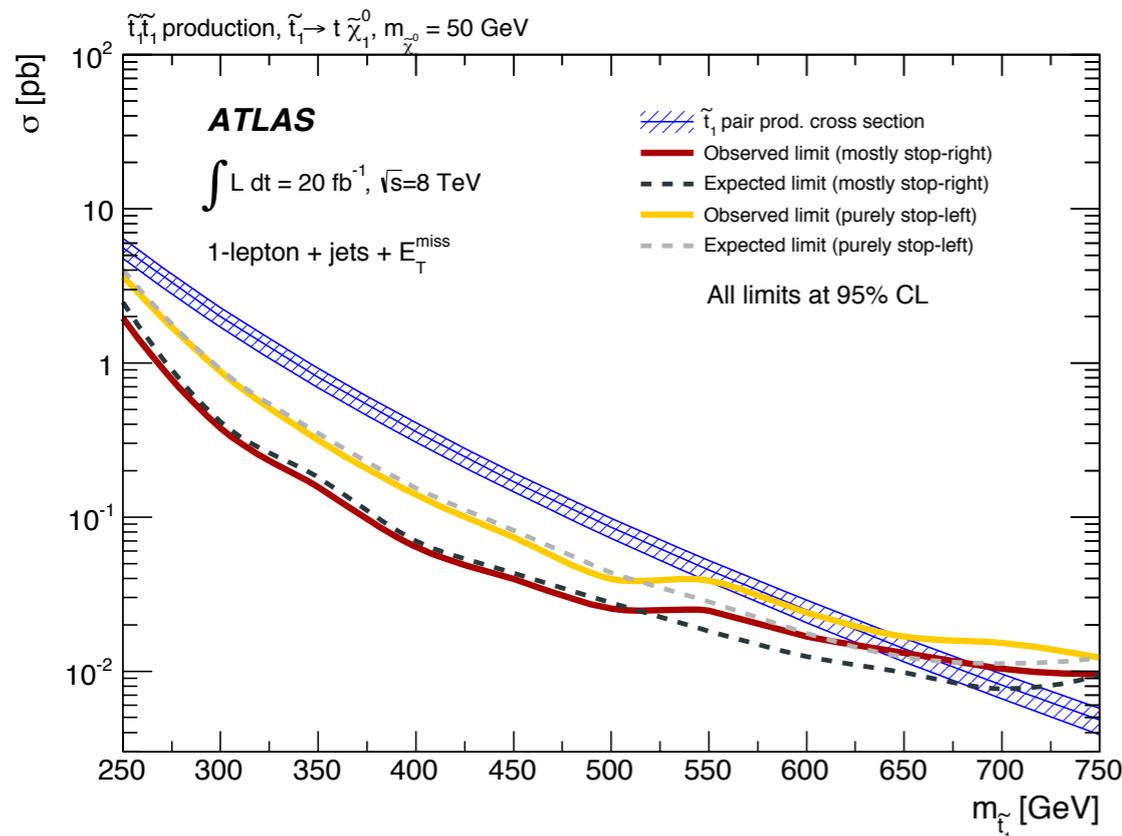
1L channels



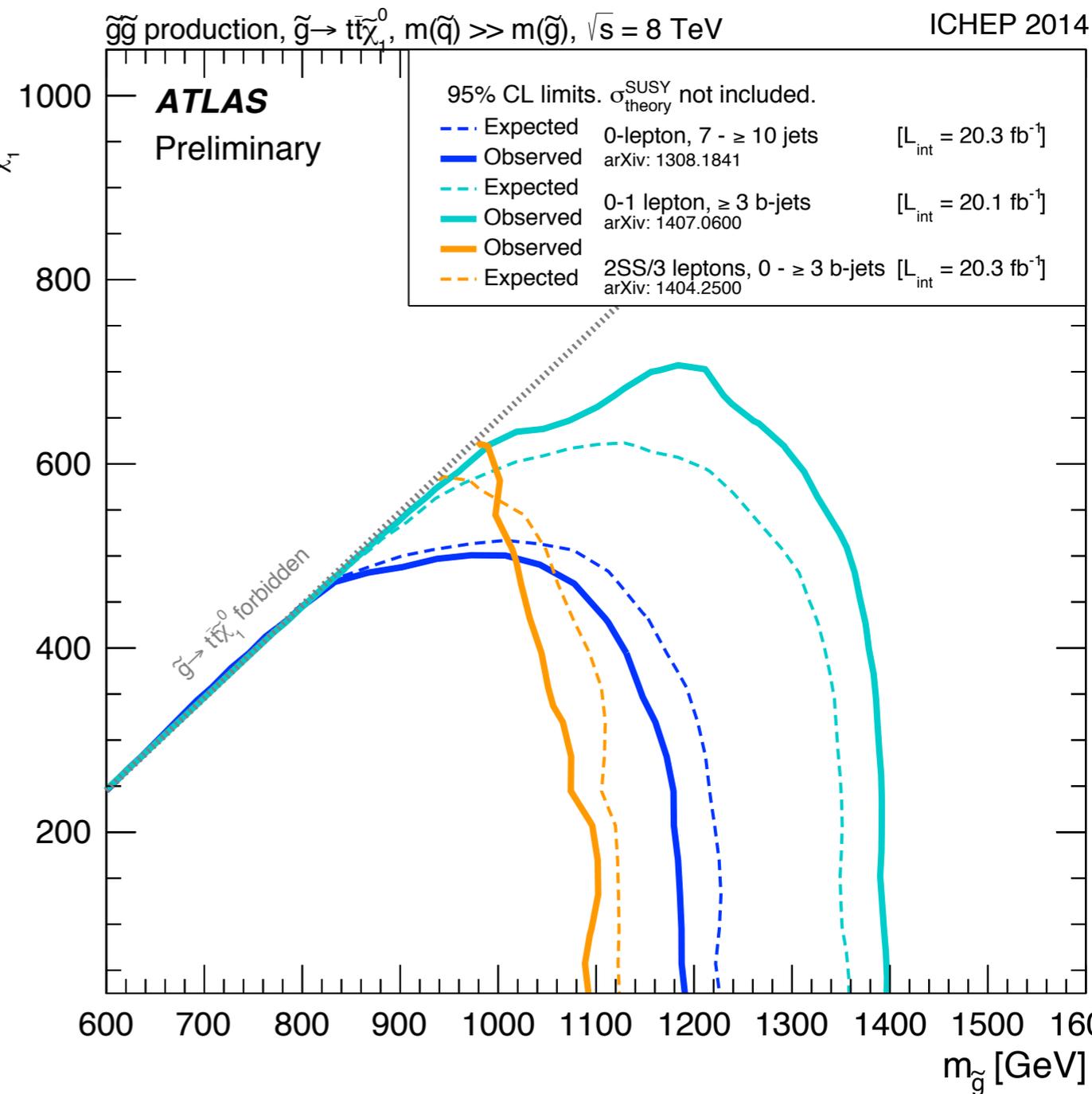
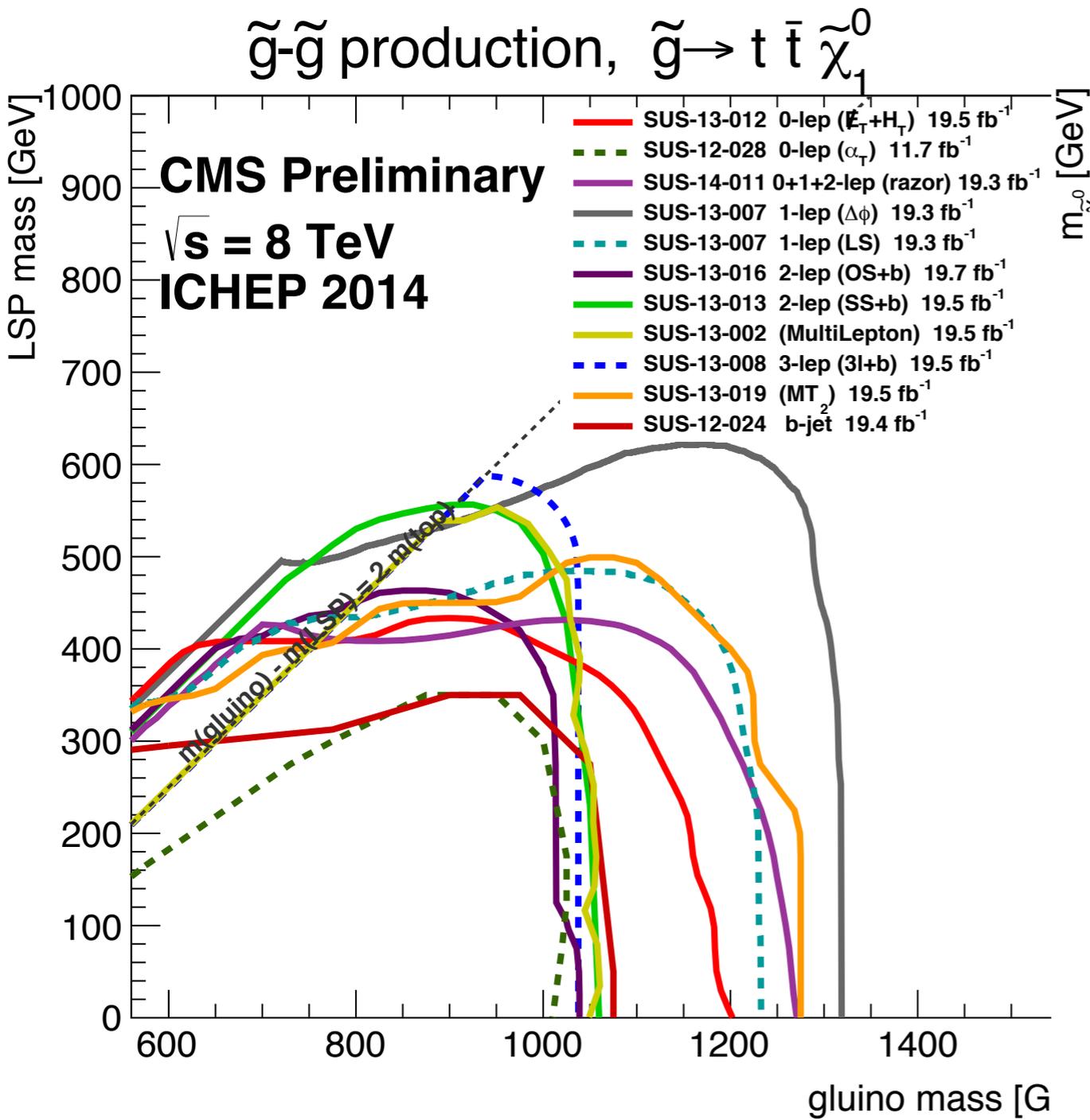
submitted to JHEP [1407.0583] EPJC 73 (2013) 2677

The acceptance depends on the polarization of the the top quarks in the stop1 to tN1 scenario. These polarizations depend on the left/right mixing of the stops and on the mixing matrices of the neutralino and charginos.

- Nominal ATLAS stop 1L result obtained using stop1 ~70% stopR. Comparison with fully left-handed top quarks: exp. limit reduced by ~50 GeV.
- Nominal CMS stop 1L result obtained using unpolarized particles. Comparison with right- and left-handed top quarks: limits vary by $\pm(10-20)$ GeV.



Overview: gluino-mediated stop searches



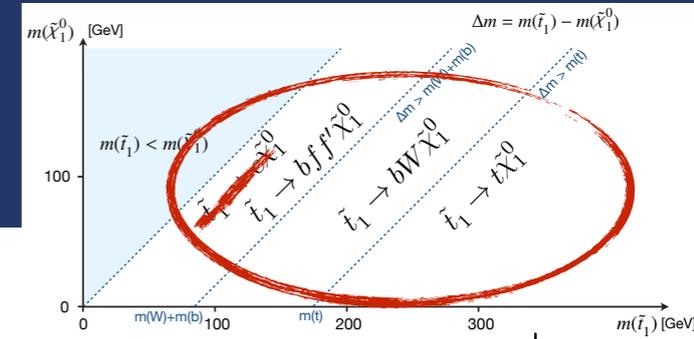
ICHEP 2014



ATLAS searches

1L channel

submitted to JHEP [1407.0583]



$$\tilde{t}_1 \rightarrow b \tilde{\chi}_1^\pm$$

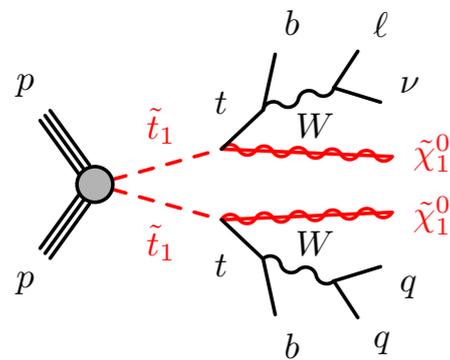
Search targets light to heavy stops in various decay modes: 4-body, 3-body, tN1, and bC1.

1-leptons, ≥ 2 -4 jets, MET, MT(lep, pTmiss)

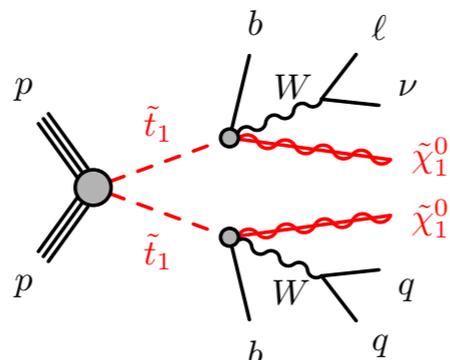
15 signal regions (SRs).

- **soft-lepton**: $p_T > 7(6)$ GeV for electrons (muons), employed for 4-body, and bC1 decays with small m_{C1} - m_{N1} values.
- **large-R (1.0) jets**: collect decay products of boosted top in heavy stop search.
- **1D and 2D shape-fits**: enhance search sensitivity in challenging scenarios (with difficult S-vs-B separation).
- **b-tagged jets**: utilised in event selection (ranging from a veto to ≥ 2) and for constructing kinematic variables.
 - ▶ MT2-like variables, hadronic top mass, topness

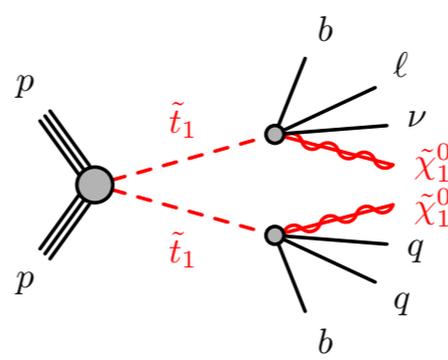
stop1 to tN1



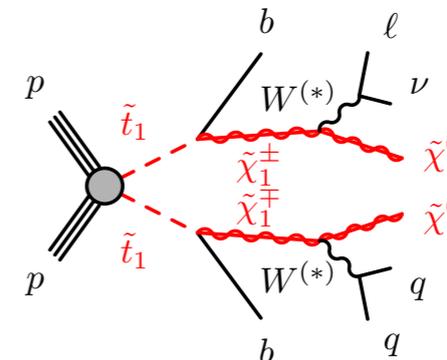
stop1 3-body decay



stop1 4-body decay



stop1 to bC1



As well as mixed tN1, bC1 decays



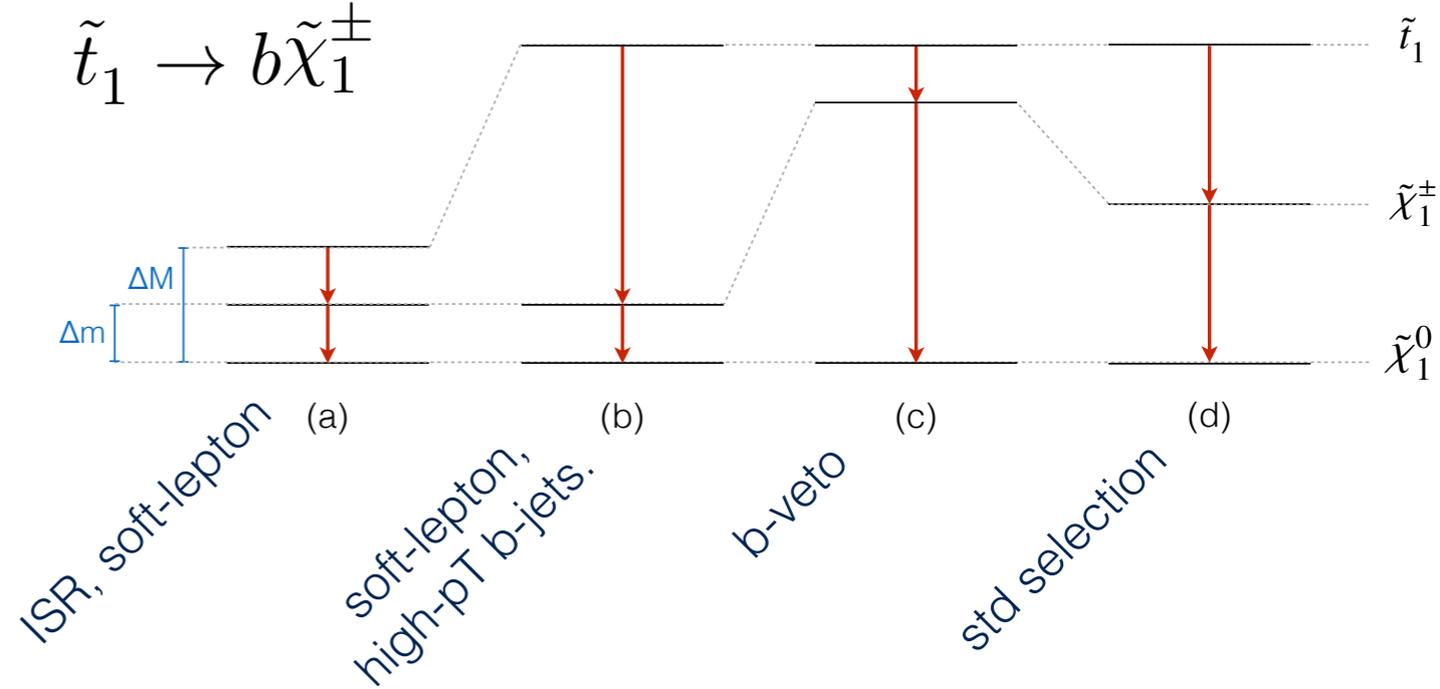
ATLAS searches

1L channel

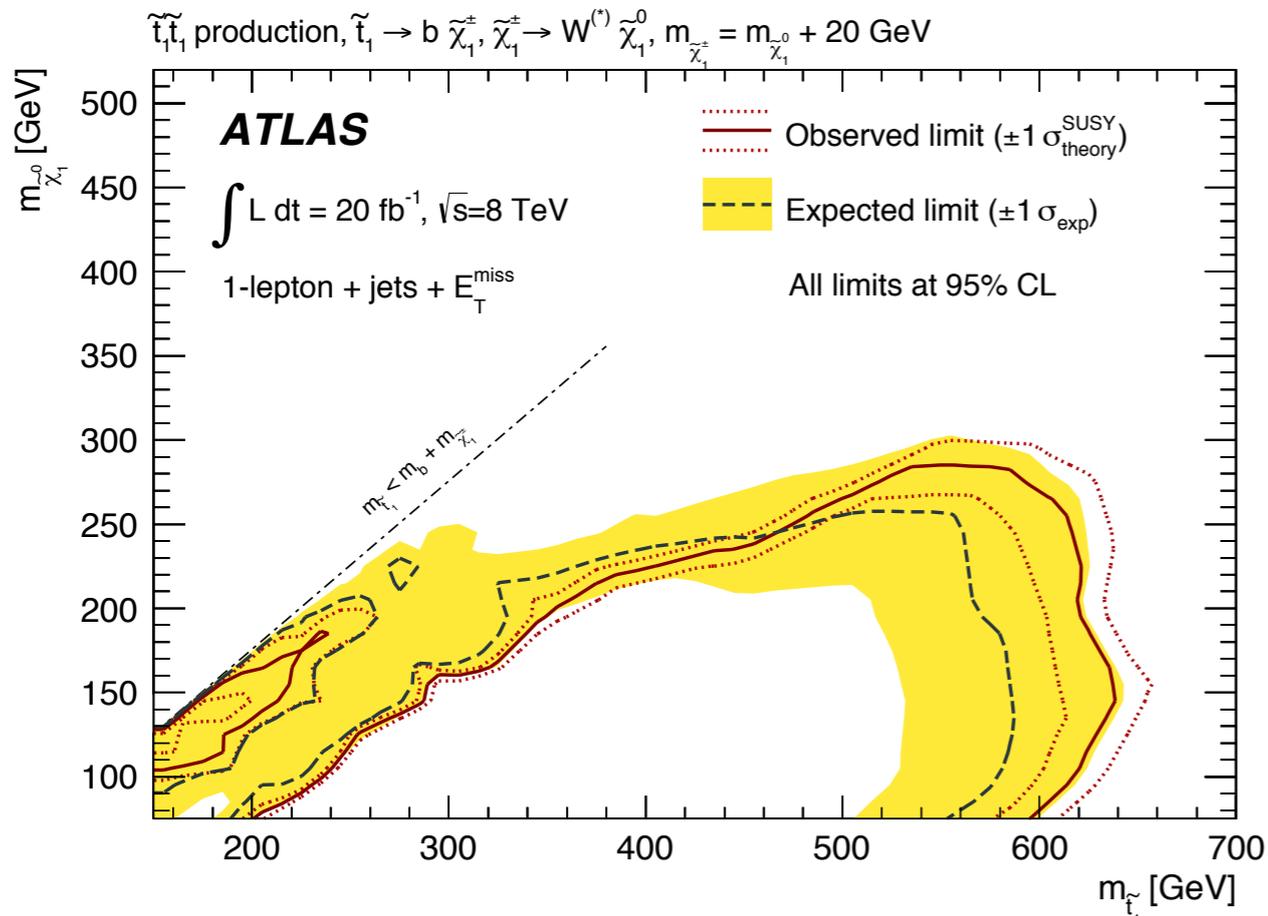
submitted to JHEP [1407.0583]

Signature driven by mass hierarchy of stop, C1, and N1. Dedicated selections to cover all scenarios.

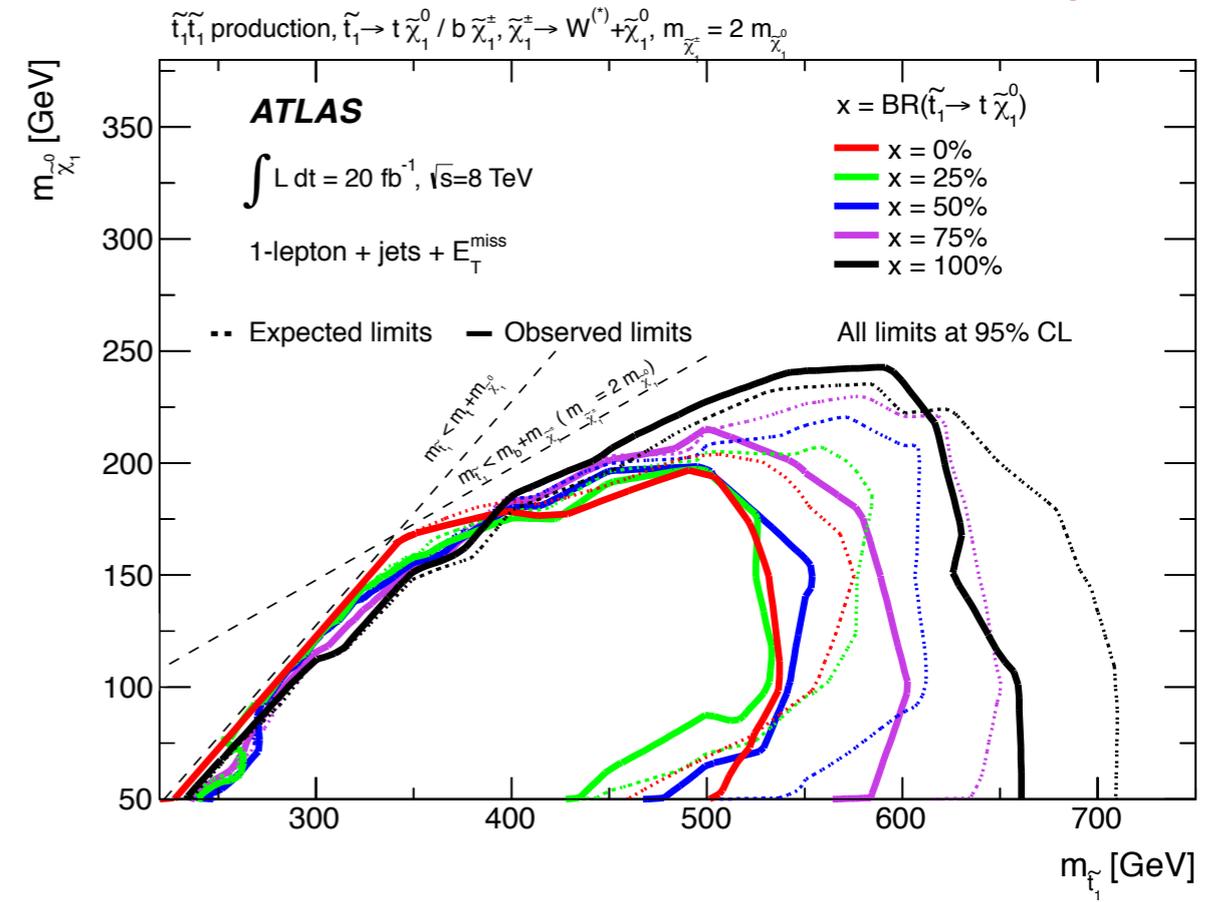
Showing example interpretations.



mass hierarchies (a), (b)



mixed tN1, bC1 decays with varying BR

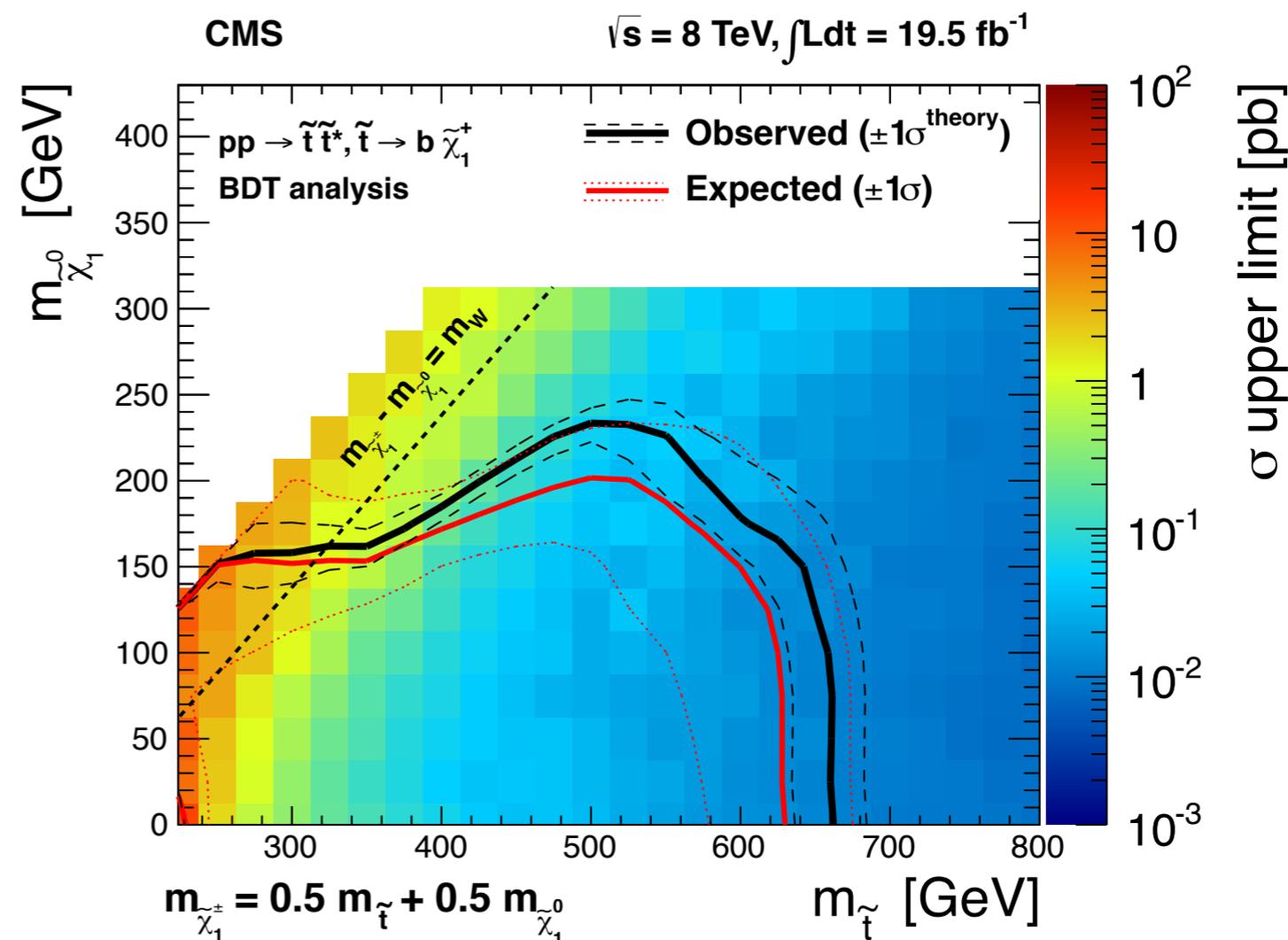
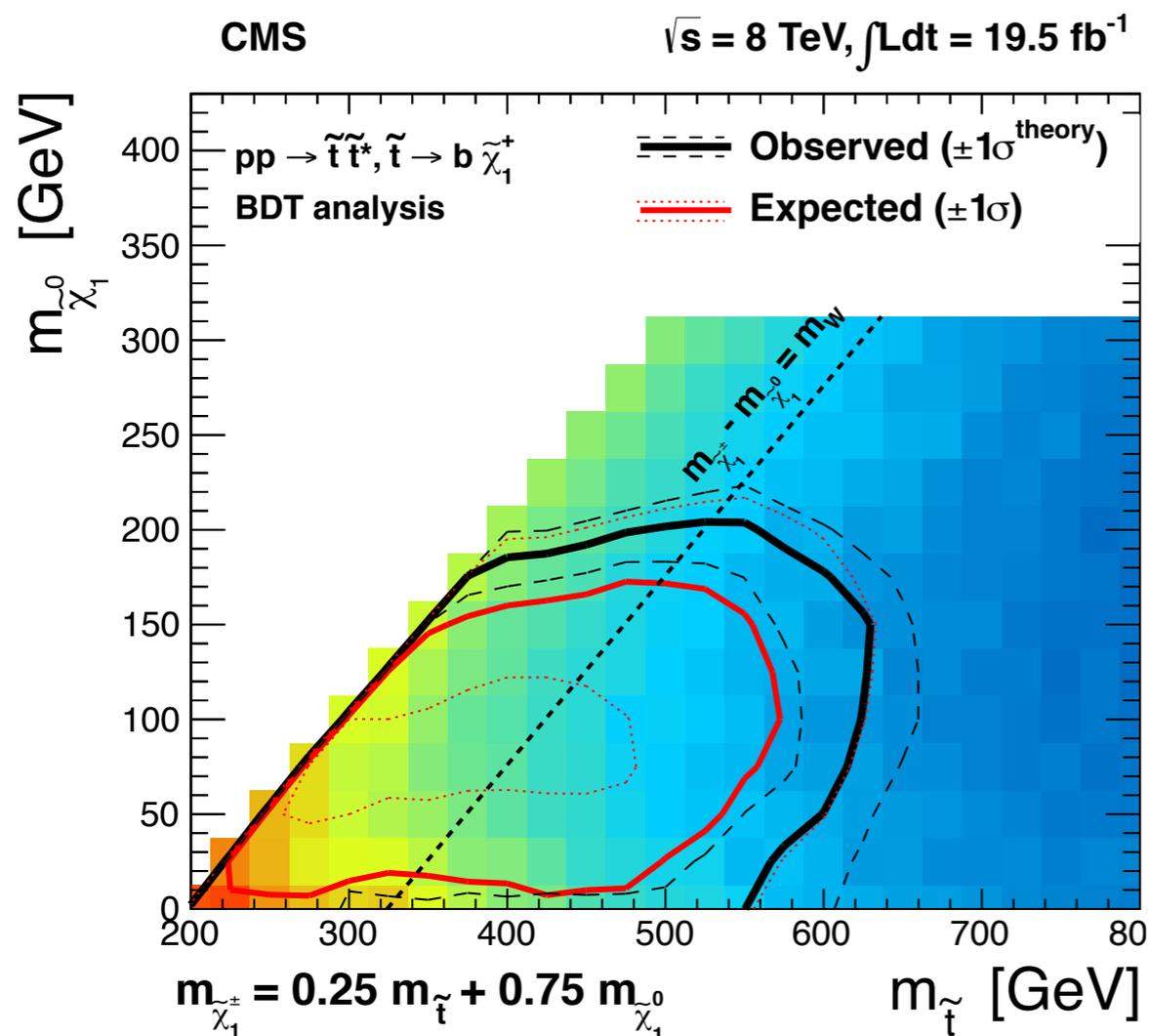


Three SUSY mass parameters to consider (LEP limit $m_{\tilde{\chi}_1^\pm} > 103.5$ GeV)

$$\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$$

BDT Exclusion limits

- stop1 to bC1 scenario
- mC1 = 0.25 mTop + 0.75 mN1 (left)
- mC1 = 0.5 mTop + 0.5 mN1 (left)

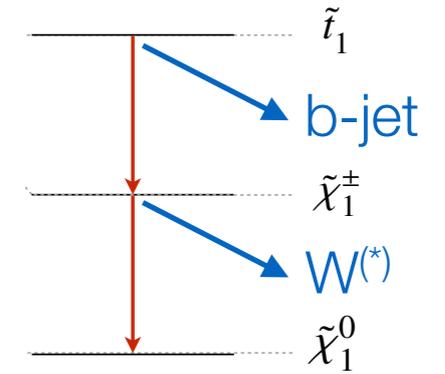


Small acceptance for models with small mC1-mN1 values (due to soft decay products, incl. lepton)

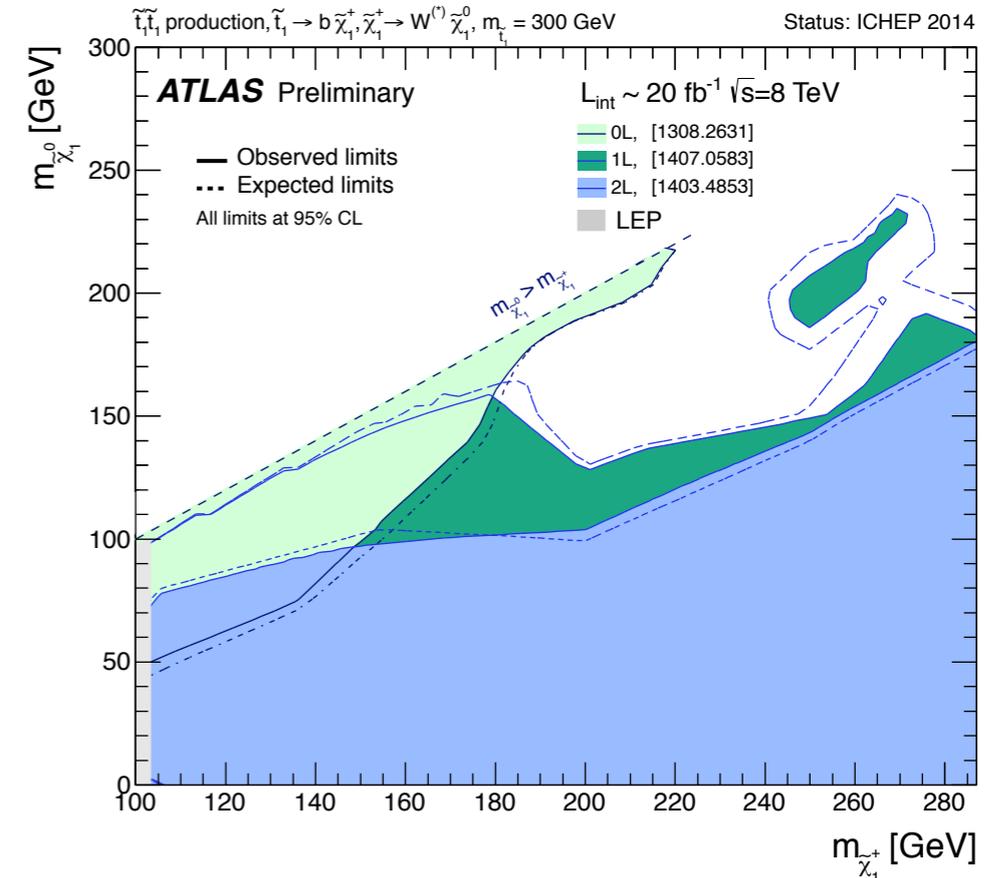
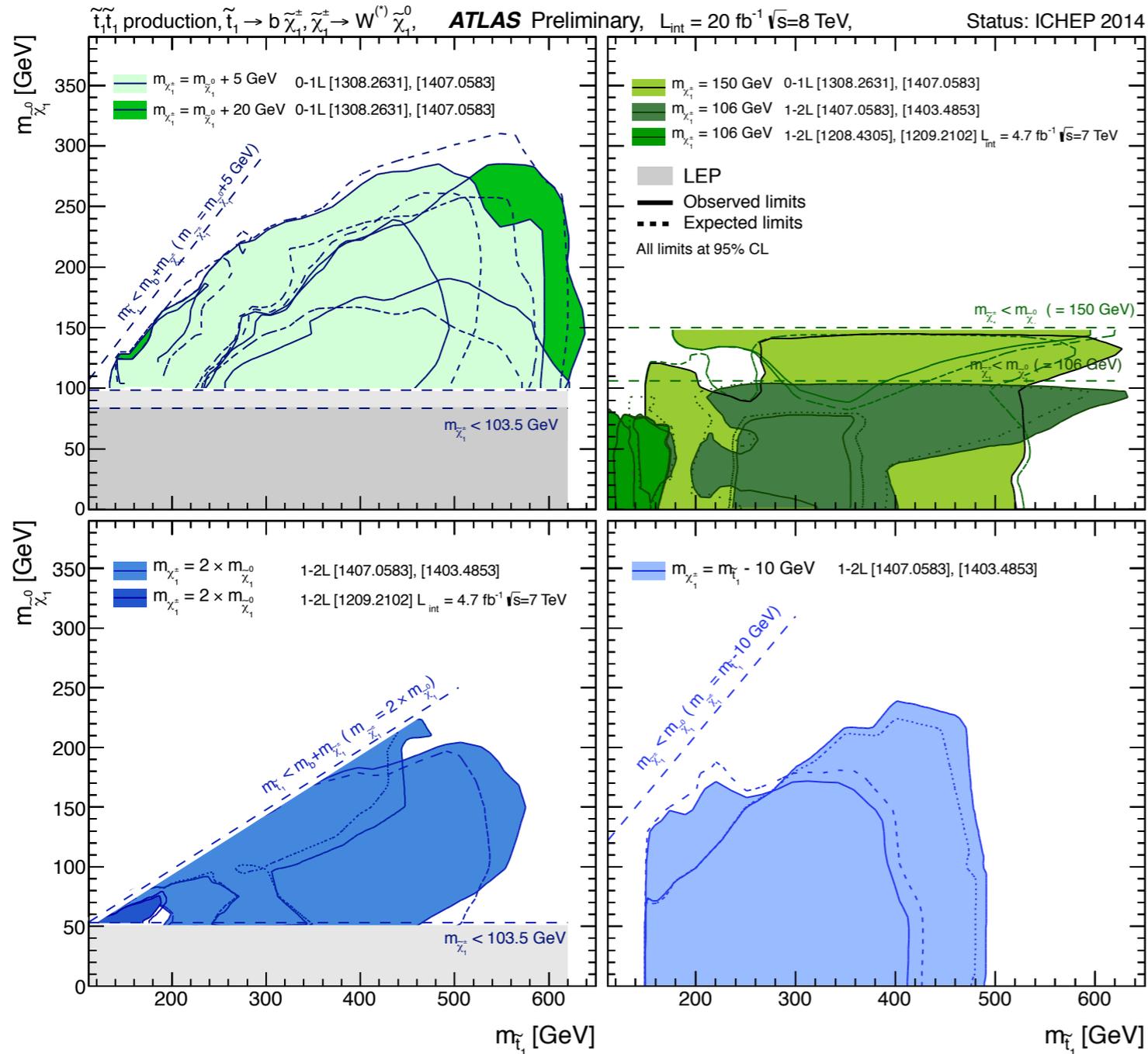
Complementarity of searches

Various assumptions on stop, C1, N1 mass hierarchies.

$$\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$$

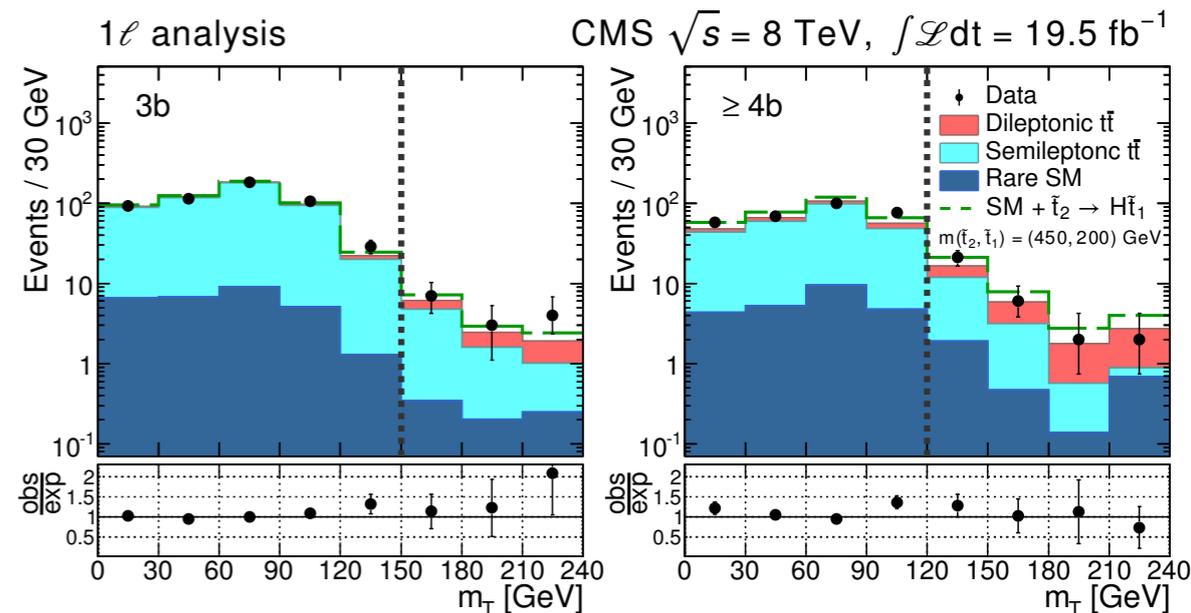


Fixed stop mass, limit in the m_{C1} - m_{N1} plane.



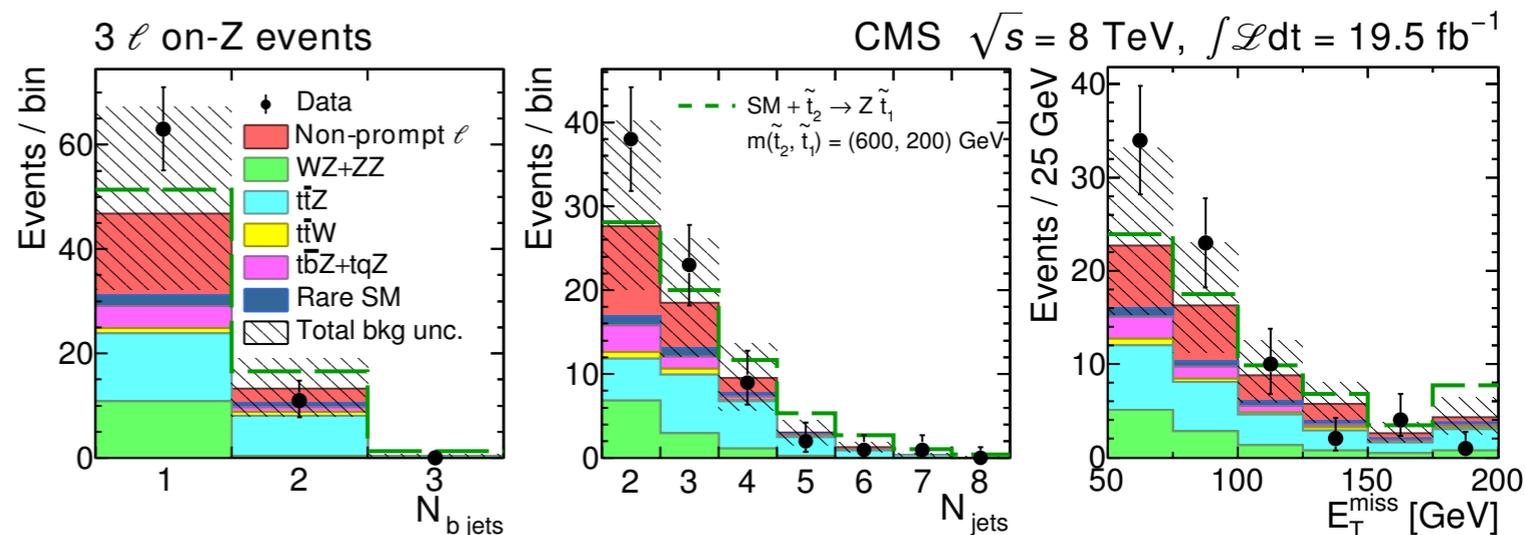
1L analysis:

- targets H to bb signature
- main background is $t\bar{t}$ + extra b-jet (for instance from mis-tag, or gluon splitting)
- use low M_T to estimate background.



3L analysis:

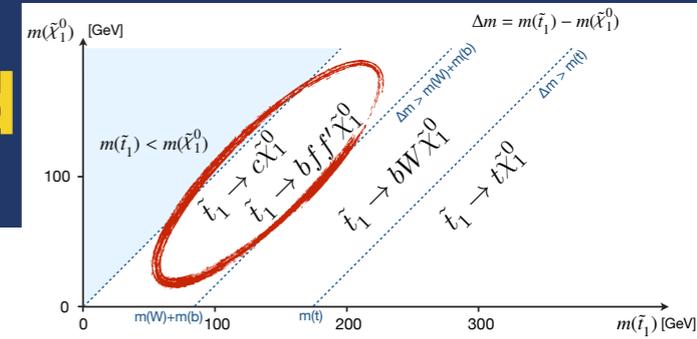
- sensitive to all scenarios
- background from WZ +jets is suppressed by b-tag requirements,
- remaining background from rare processes or non-prompt leptons
 - WZ/ZZ from MC
 - non-prompt lepton from data



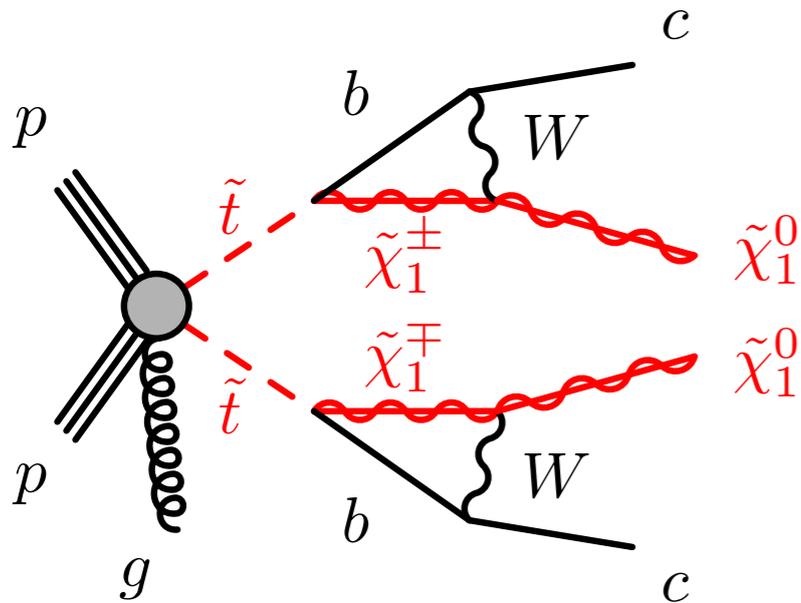


ATLAS searches **monojet/c-tagged**

accepted by PRD [1407.0608]



Signature: two charm-jets with little momentum, some MET

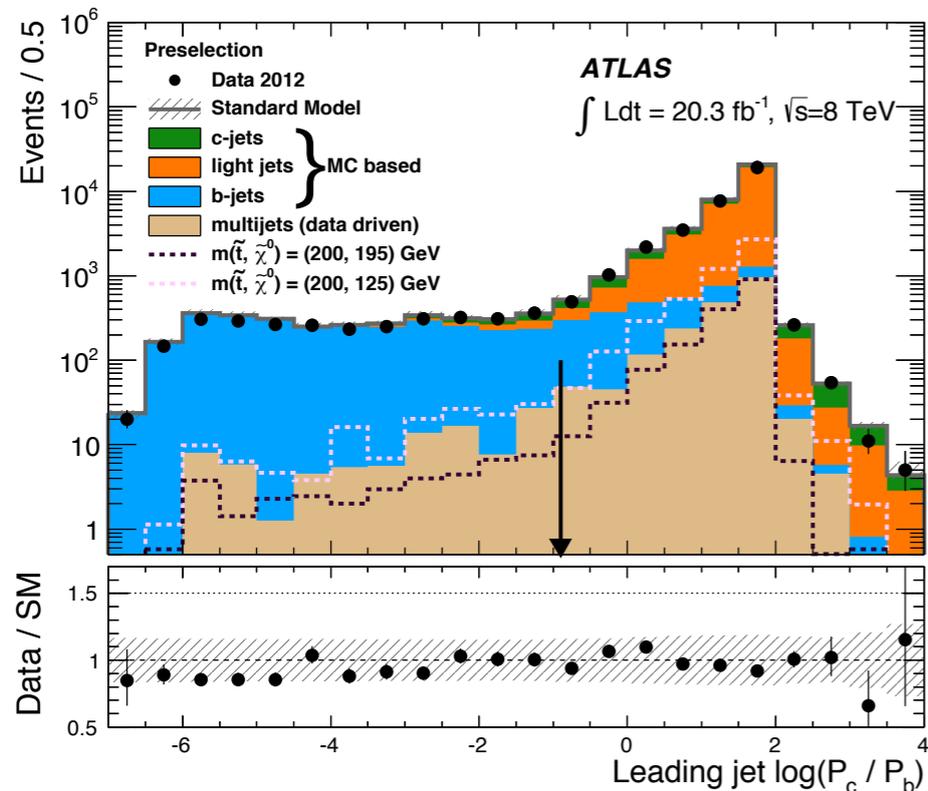


→ require ISR to boost the system

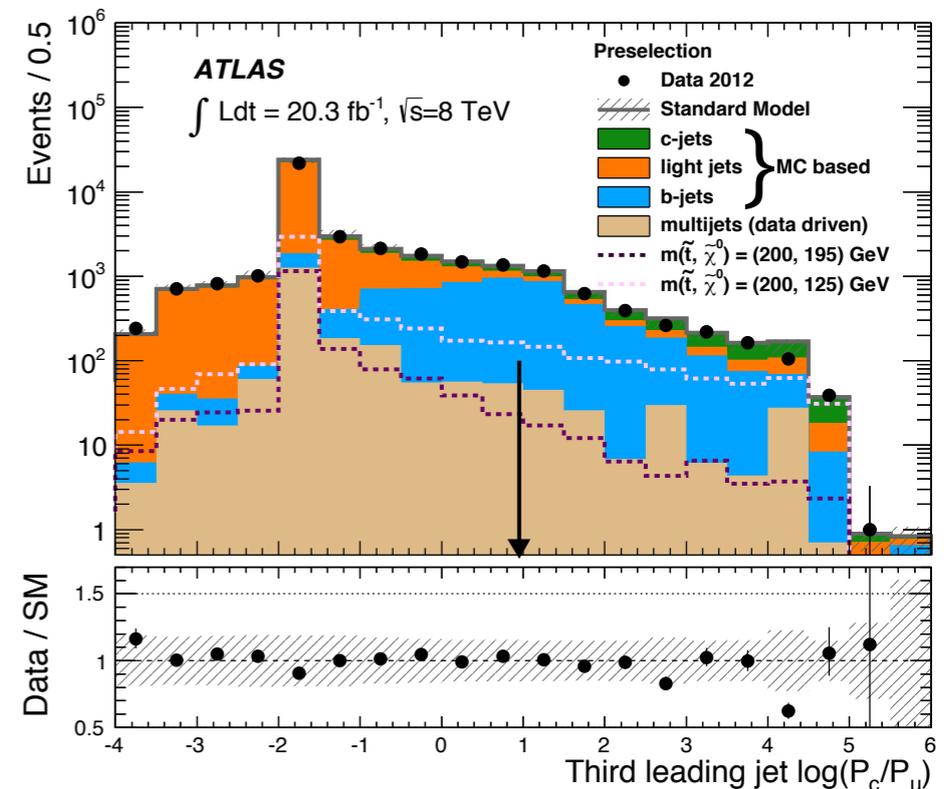
Two approaches

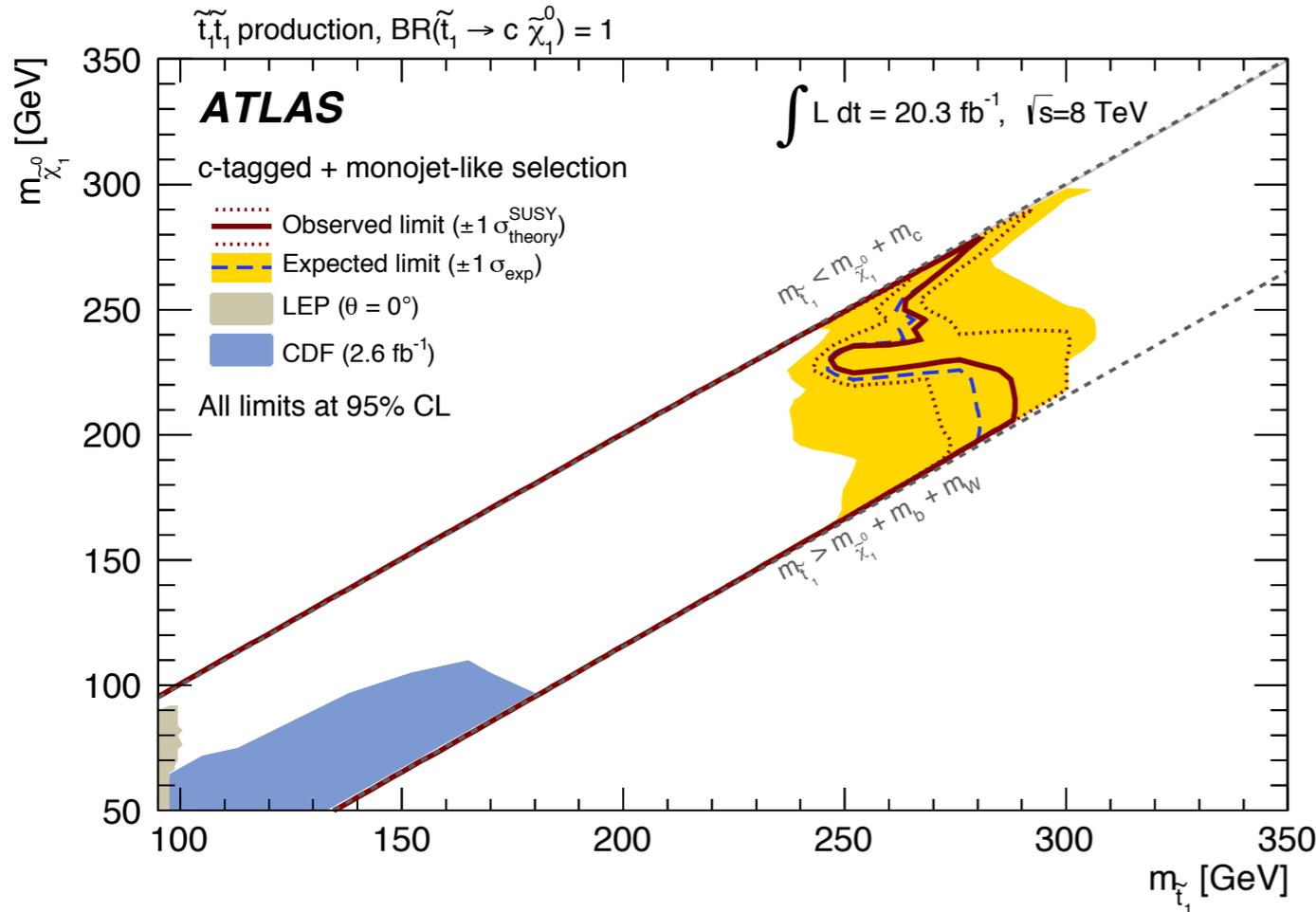
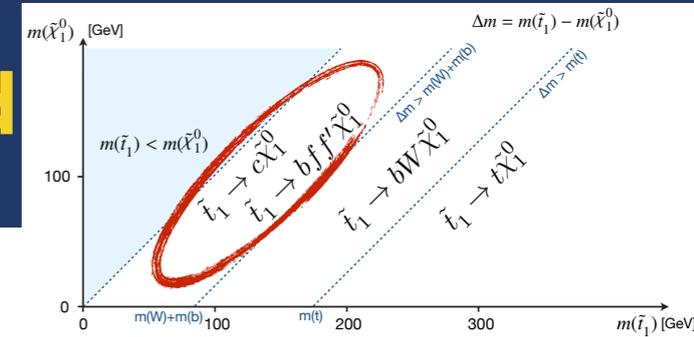
- (1) small $\Delta m = m(\text{stop}) - m(\tilde{\chi}_1^0) < 20$ GeV: c-jets with too low p_T
 - monojet analysis: high- p_T leading jet, large MET, jet veto (≤ 3 jets),
- (2) large $\Delta m > 20$ GeV: c-jets can be reconstructed
 - c-tagged analysis: ≥ 4 jets, ≥ 1 c-tag, high- p_T untagged leading jet (ISR), large MET

Loose anti b-tag on the leading jet



Medium c-tag on the 3 sub-leading jets

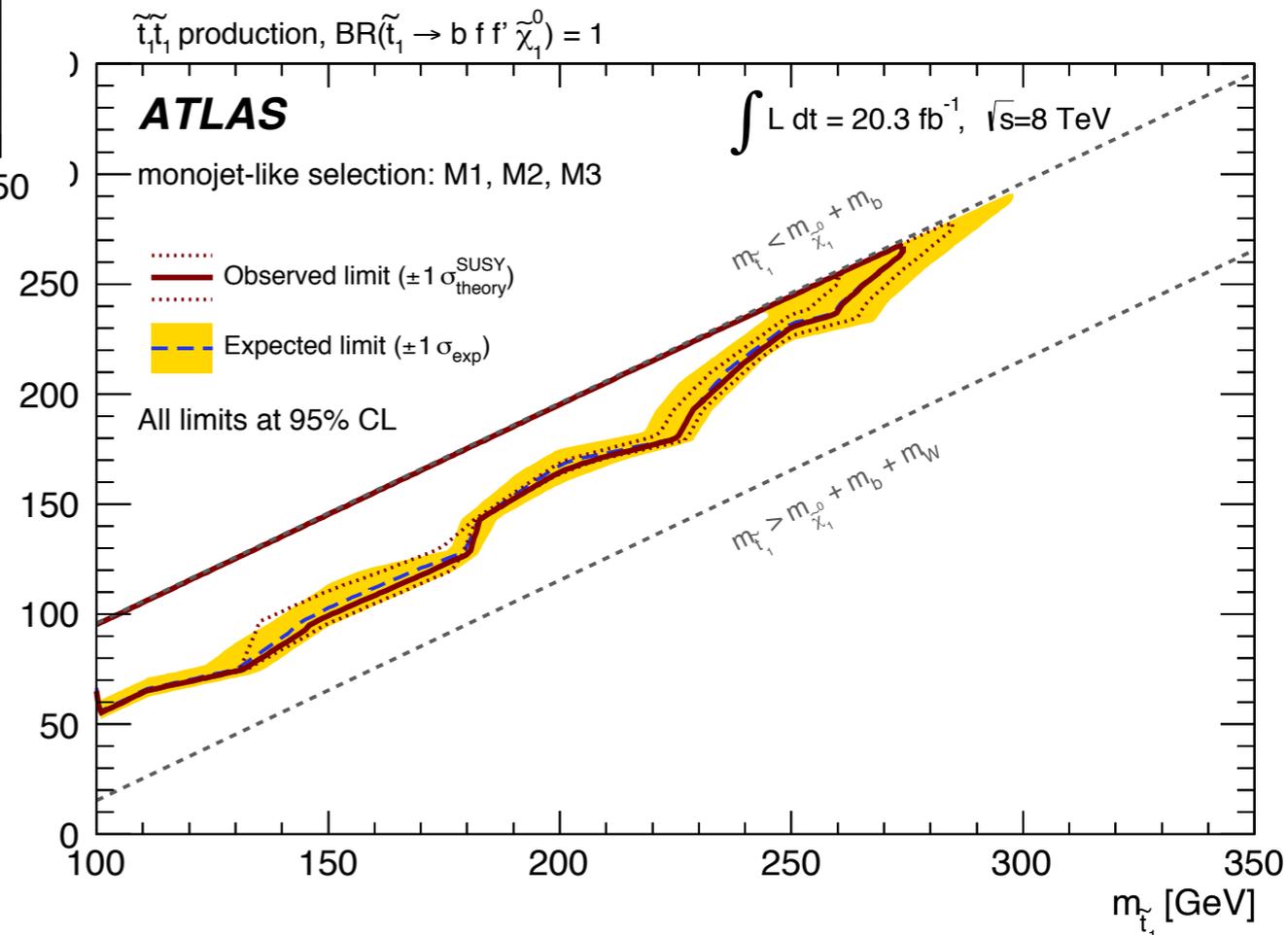


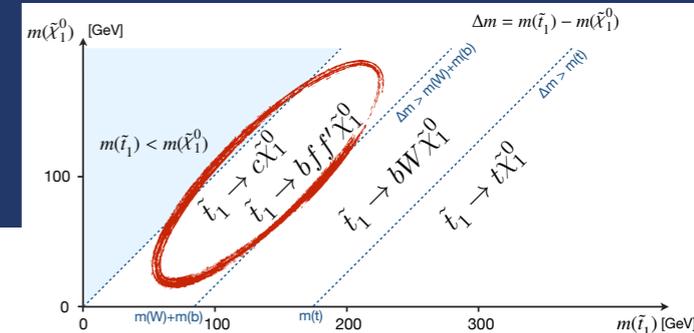


Limits in stop to cN1.

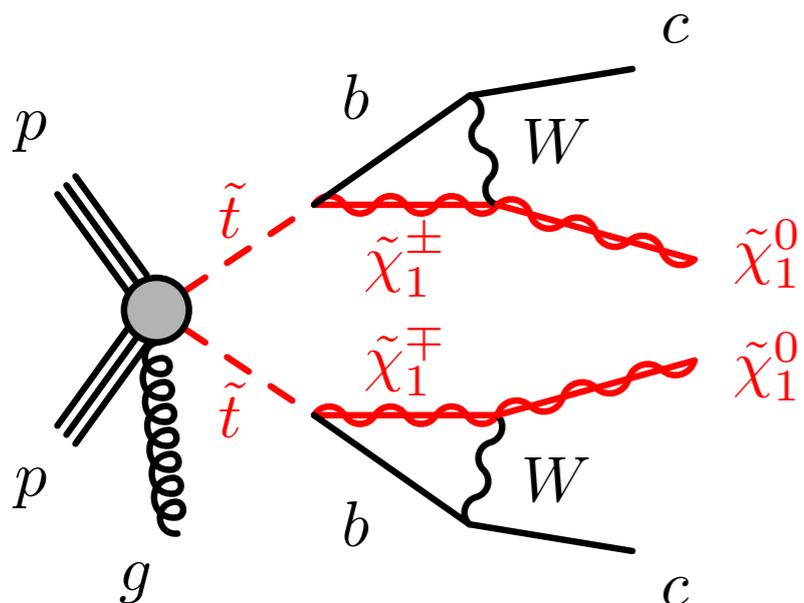
Limits in 4-body scenario

sensitivity of monojet search independent of stop decay (soft b-, c-jets, leptons, etc.). Similar reach for compressed scenarios in case of sbottom pair production.





Signature: two charm-jets with little momentum, some MET



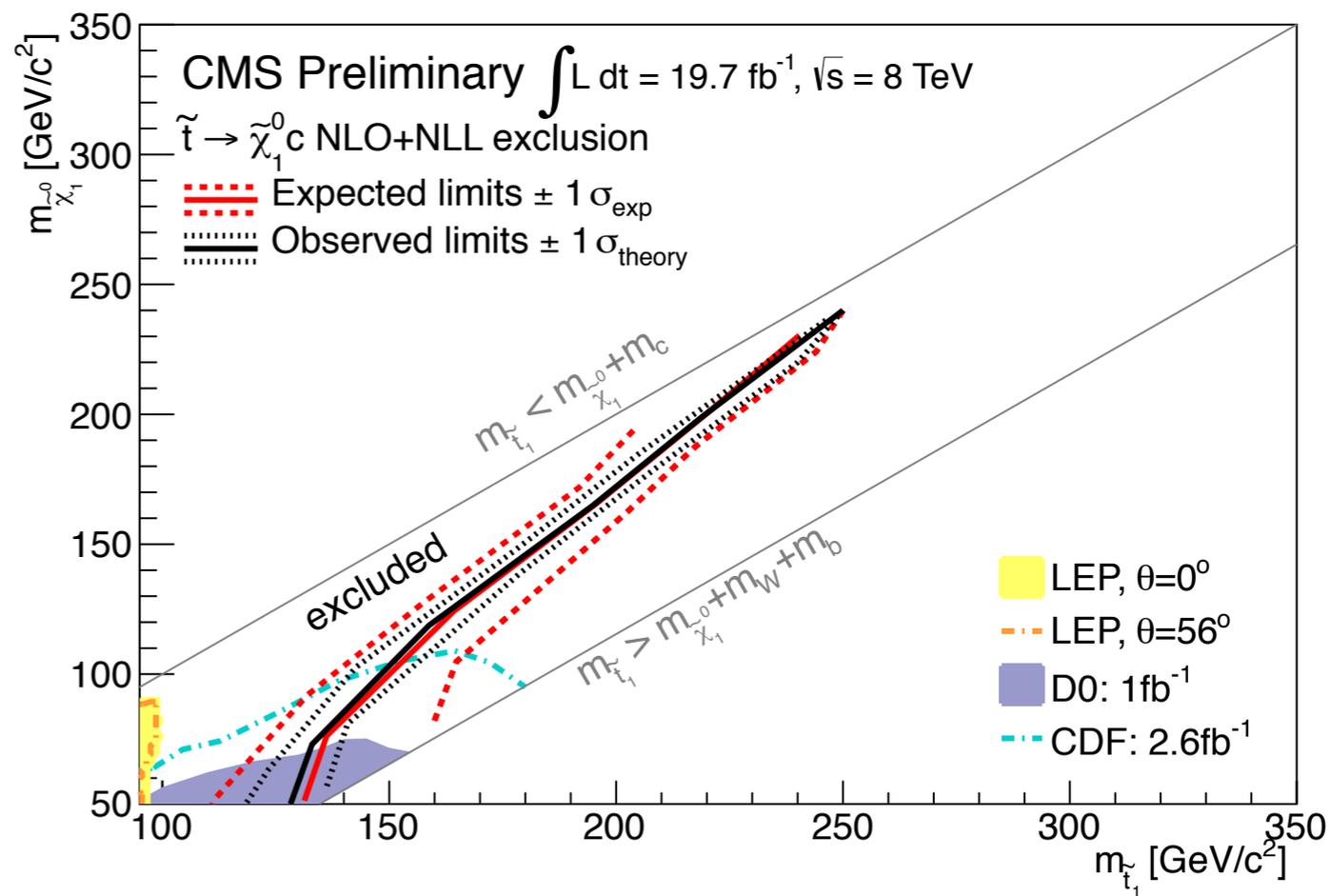
→ require ISR to boost the system

Monojet analysis:

- high- p_T leading jet, allow for 2nd jet with $p_T > 60$ GeV, veto 3rd jet, veto leptons.

Backgrounds:

- Z(nunu)+jets, W(lnu)+jets
- estimated from data in CRs,
- total SM uncertainty as low as ~4%



Other (complex) decay possibilities

1. Longer decay chains

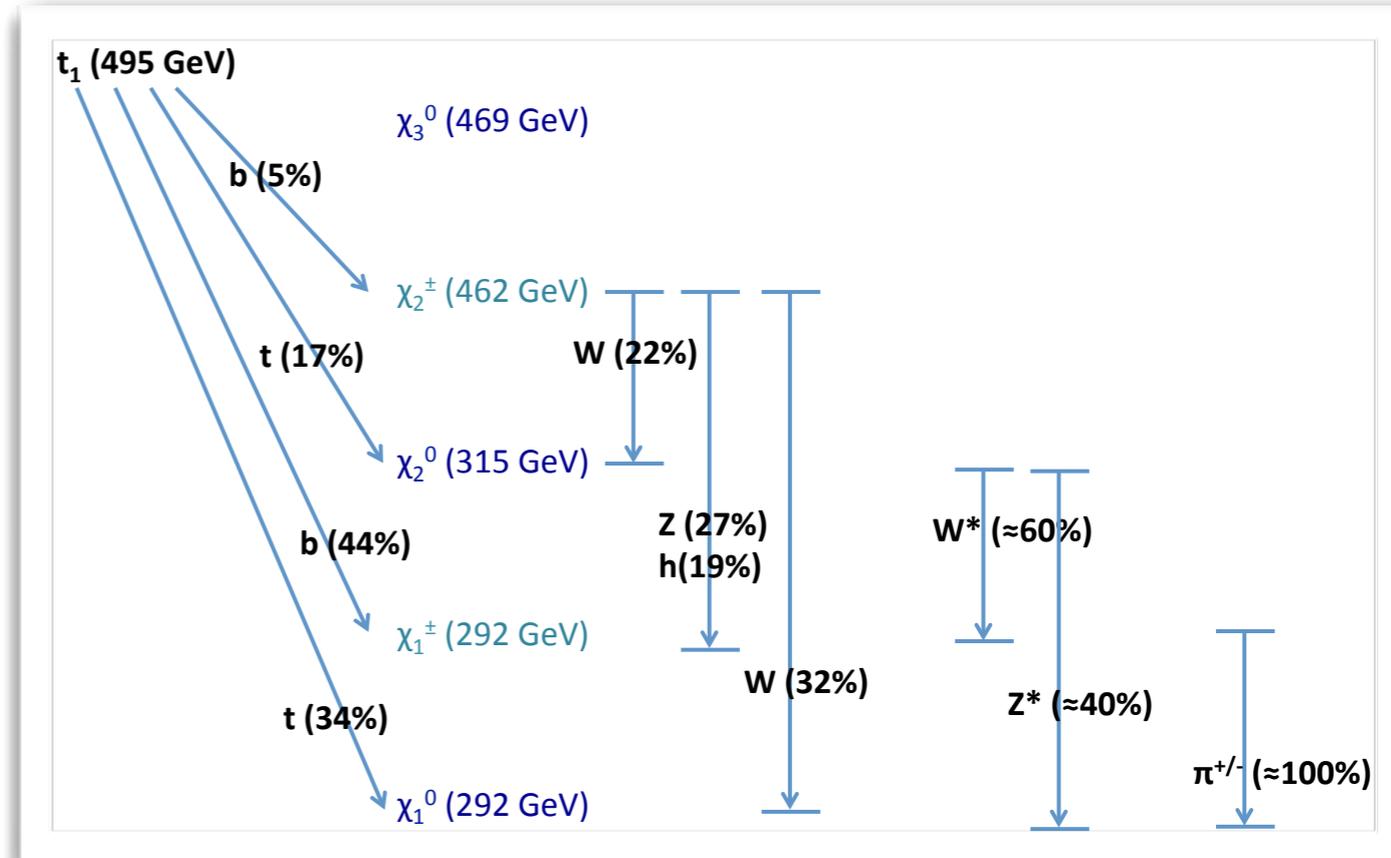
- with heavier neutralinos and/or charginos

2. Asymmetric decays

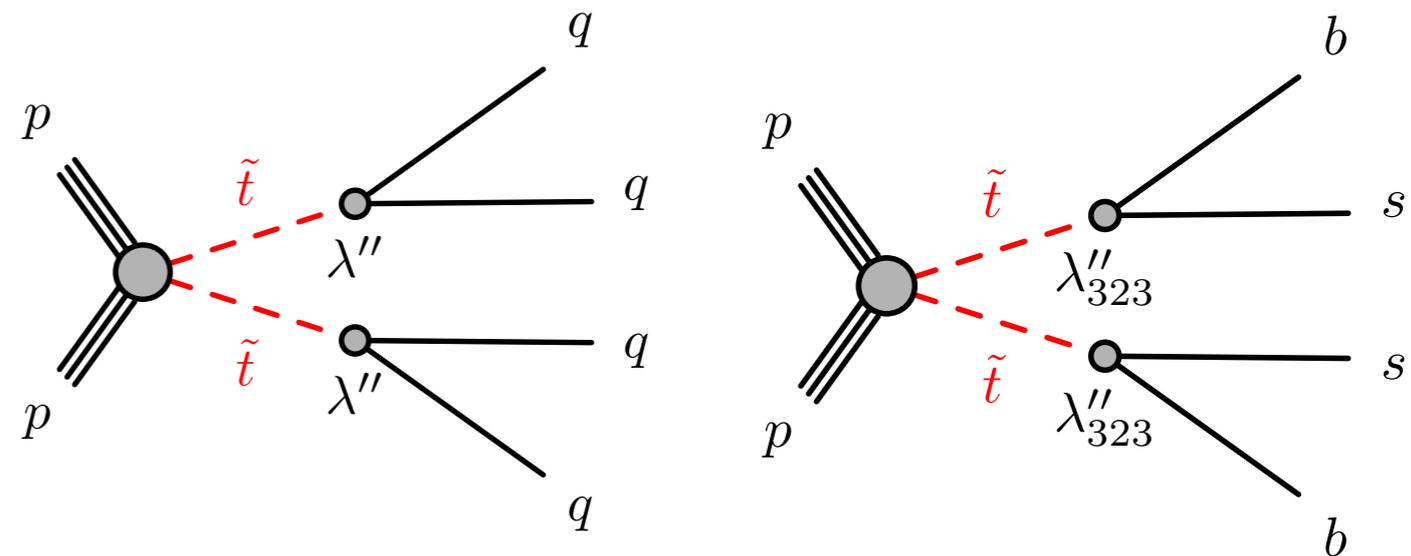
3. NI is not the LSP, but e.g. the gravitino (\tilde{G}),

for example:

$$\tilde{t}_1 \rightarrow b + \tilde{\chi}_1^\pm \rightarrow b f f' \tilde{\chi}_1^0 \rightarrow b f f' Z \tilde{G}$$



4. (some) R-parity violation, opens many more possible final states, for example:



6. but also long-lived (R-hadron, stoponium)

Heavy stop search

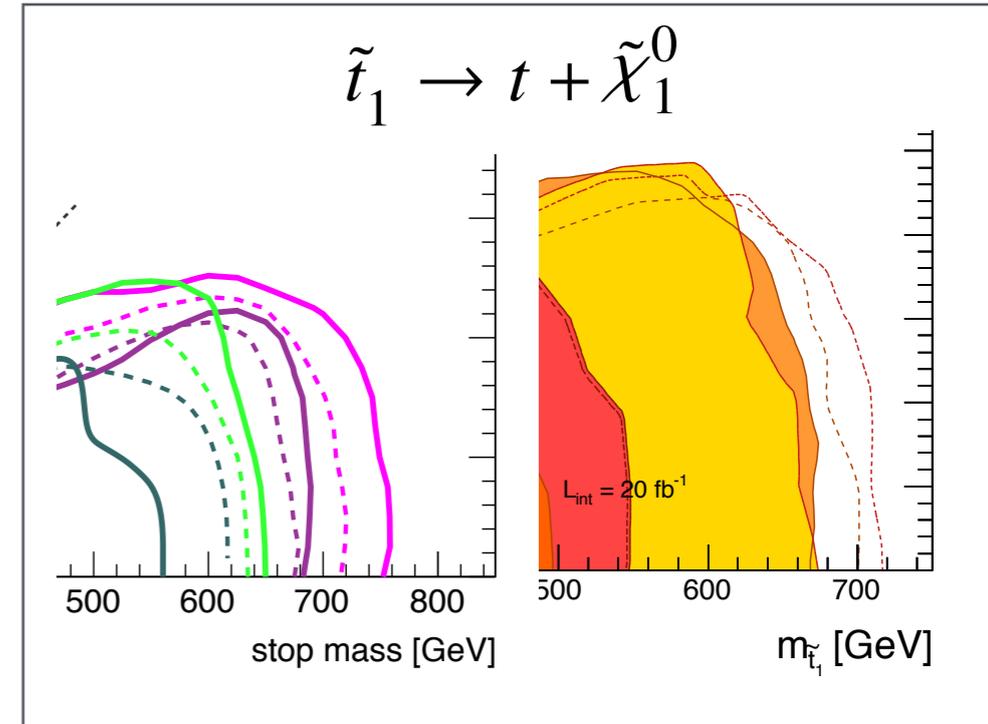
- stop to t N1,
- stop to b C1 (to b W N1)

$$\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$$

challenge: tiny cross section

- ~8 fb for m(stop) at 700 GeV
- compared to ~250 pb for ttbar production
 - and ~0.2 pb for ttbar+ Z

Best sensitivity from 0L and 1L channels.



- Dominant background: **semi (di)-leptonic ttbar** for the stop 0L (1L) search, where the (2nd) lepton is
 1. not identified (failed PID) / outside acceptance,
 2. hadronically decaying tau-lepton.
 } missed particle / extra neutrino(s)
 → more E_T^{miss} , $m_T > m(W)$
- Remove background (while retaining high signal acceptance) using:
 - lepton veto, tau-veto, hadronic top reconstruction, kinematic variables (next slide).

Top pair background suppression

Tools to suppress dominant bkg

Tight 2nd lepton veto:

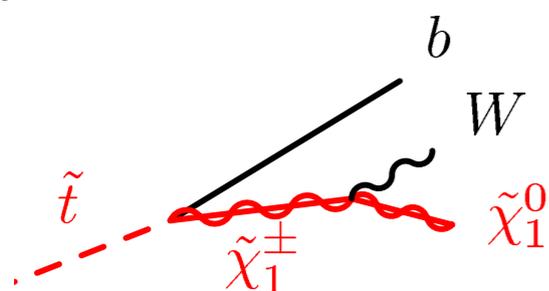
Signal has no 'edge' in m_{Tl} hence dominated by single lepton.

- lepton p_T , PID
- veto on isolated tracks
- veto on had-taus

B-tagged jets:

For chargino decays, b -jet p_T can be significantly larger than in SM $t\bar{t}$.

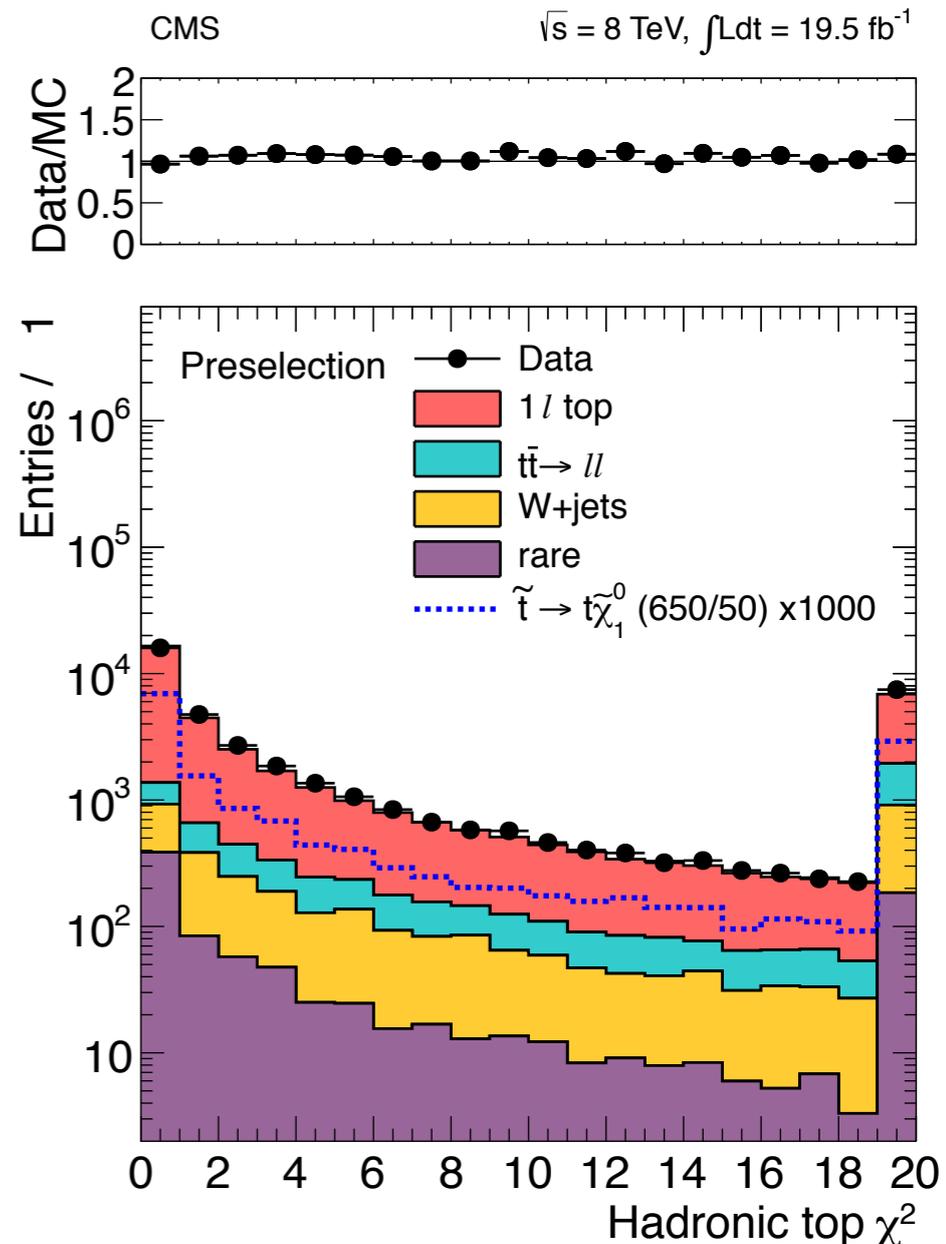
- require high p_T b -jet(s),
- require two b -tagged jets.

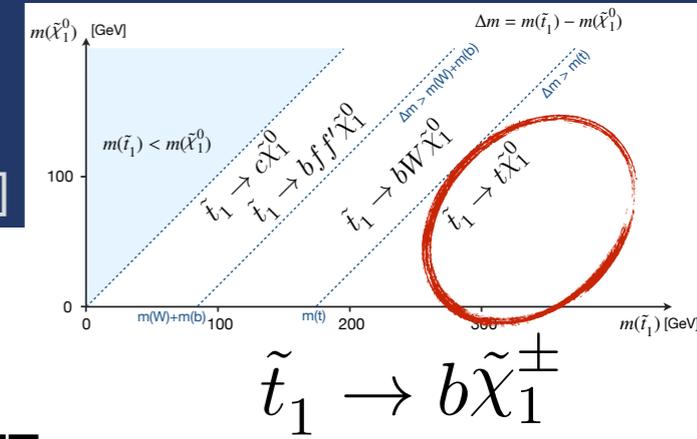


Hadronic top mass:

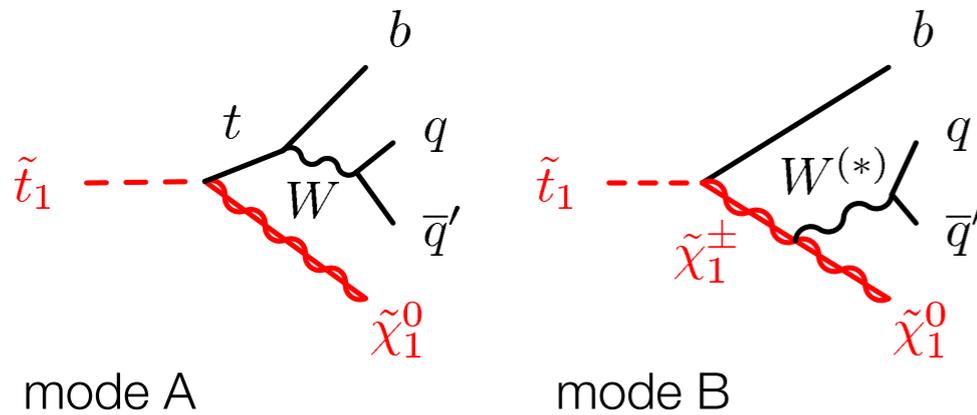
While for decay via top quark, three jets stem from a hadronic top decay, this is not true for di-lep $t\bar{t}$.

- reconstruct hadronic top mass from three jets $m(jjj)$,
- require loose compatibility with $m(t)$





Target stop1 pair production with 2 possible decays (BR=100%, or mixed)



0-leptons, ≥ 2 b-tagged jets, MET

Three sets of signal regions (SRs):

- SR-A 'fully resolved': 6 distinct jets (akt0.4) [A or B].
- SR-B 'partially resolved': 4 or 5 jets (akt0.4) and 2 re-clustered akt1.2 jets compatible with mTop [boosted tops in A].
- SR-C: 5 jets (akt0.4) for scenarios with at least one B decay, and small mC1-mN1 values.

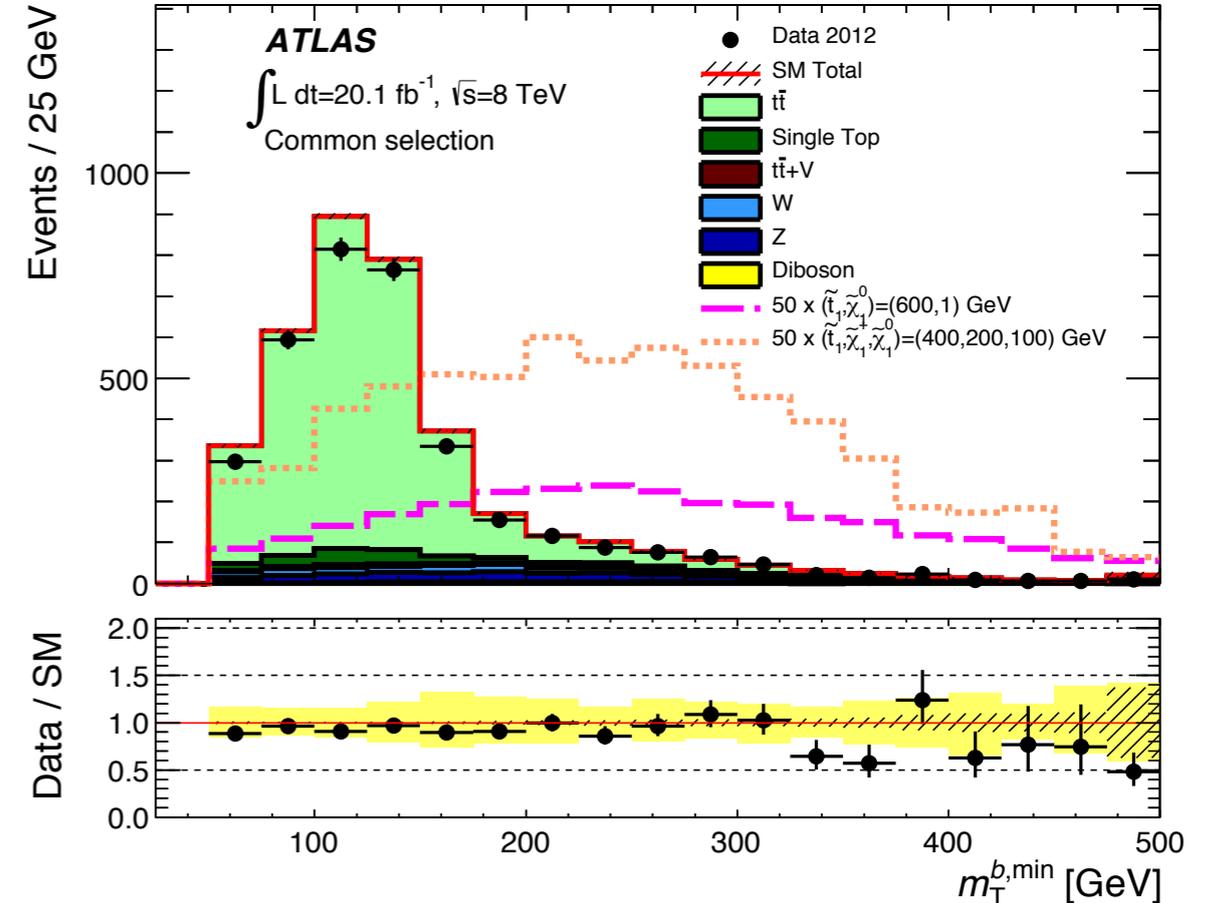
Analysis idea: reconstruct two hadronic top quarks, investigate E_T^{miss} distribution.

Backgrounds:

- semi-leptonic ttbar (with had tau) — 1L CR, treat lepton as jet,
- Znu + HF — 2L CR, add lepton pT to MET
- residual multijet — jet smearing method.

Discriminating variables:

- $\min \Delta\phi(pT_{\text{miss}}, \text{jet})$: reject fake MET
- $mT(pT_{\text{miss}}, \text{b-jet closest in } \phi \text{ to } pT_{\text{miss}})$: reject tT
- MET, and more.

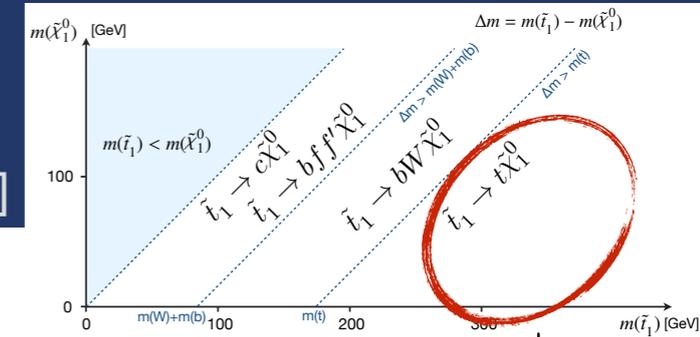




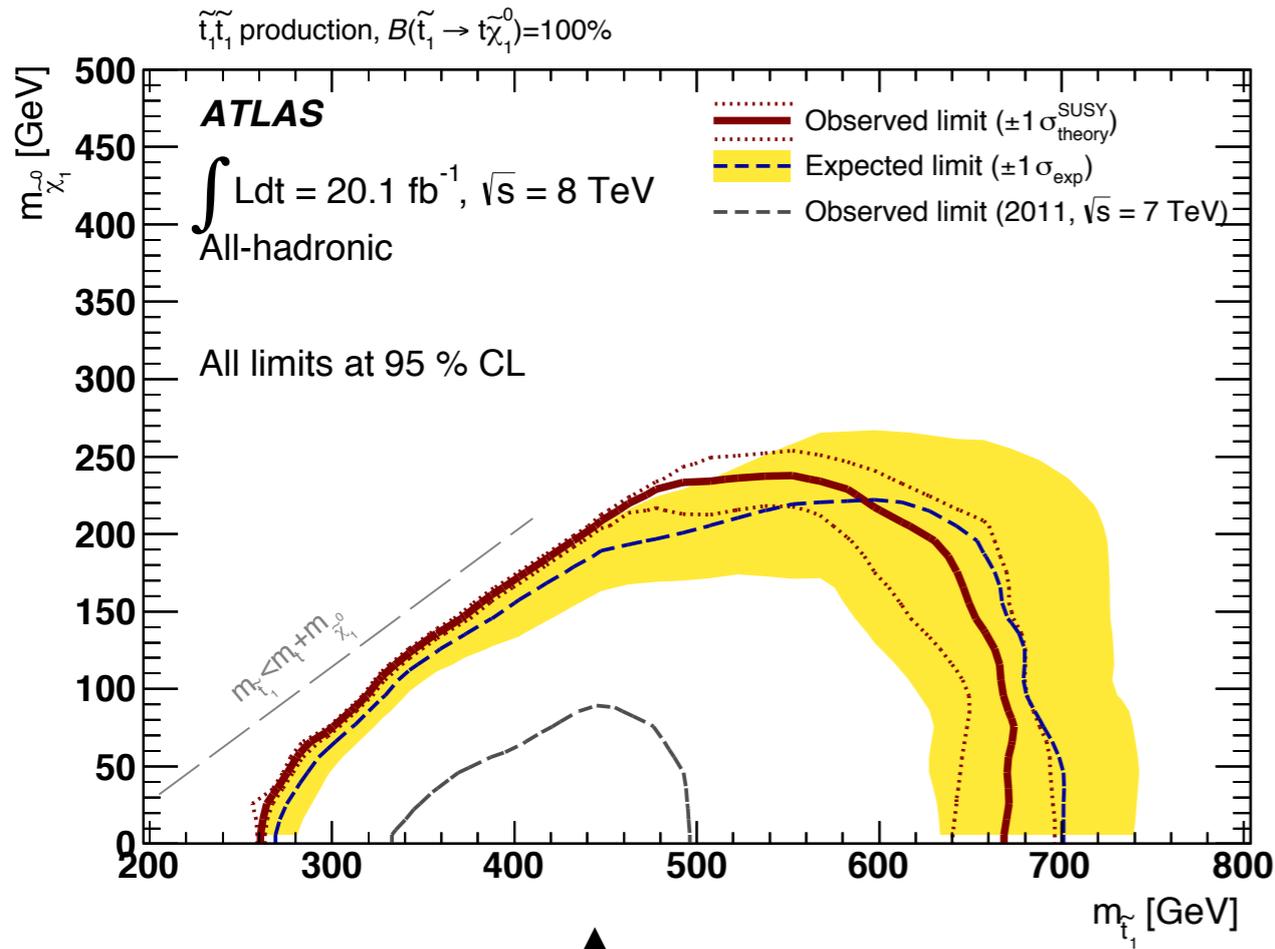
ATLAS searches

0L channel

JHEP 09 (2014) 015 [1406.1122]

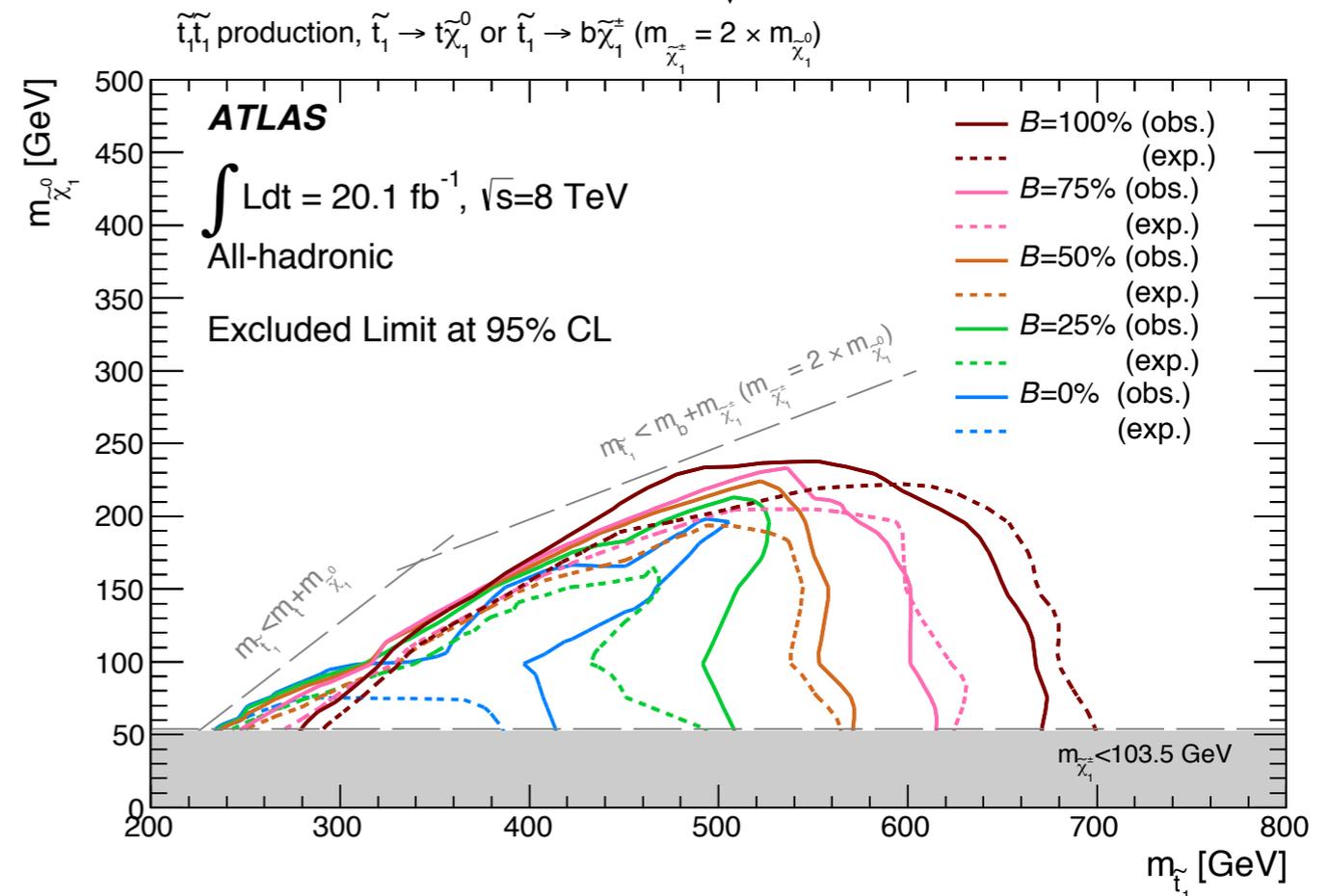


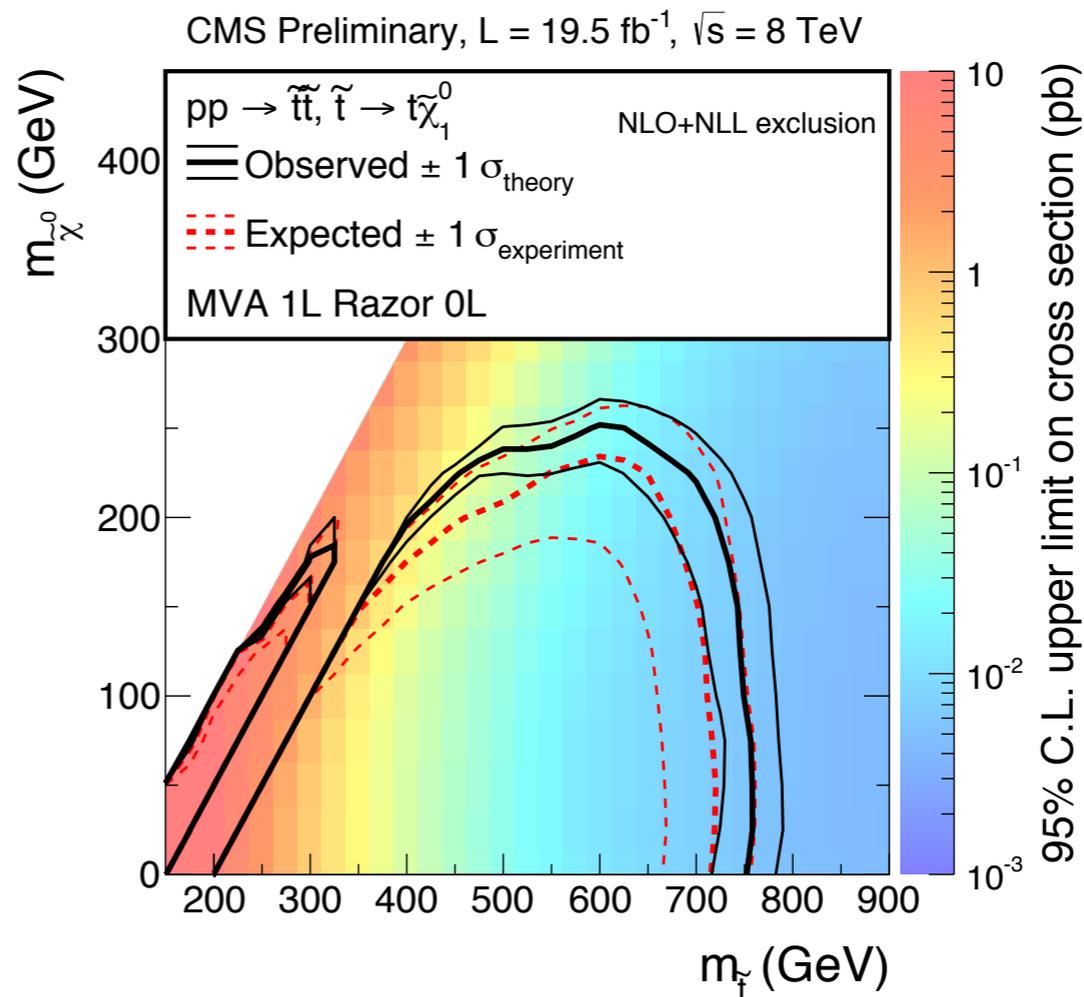
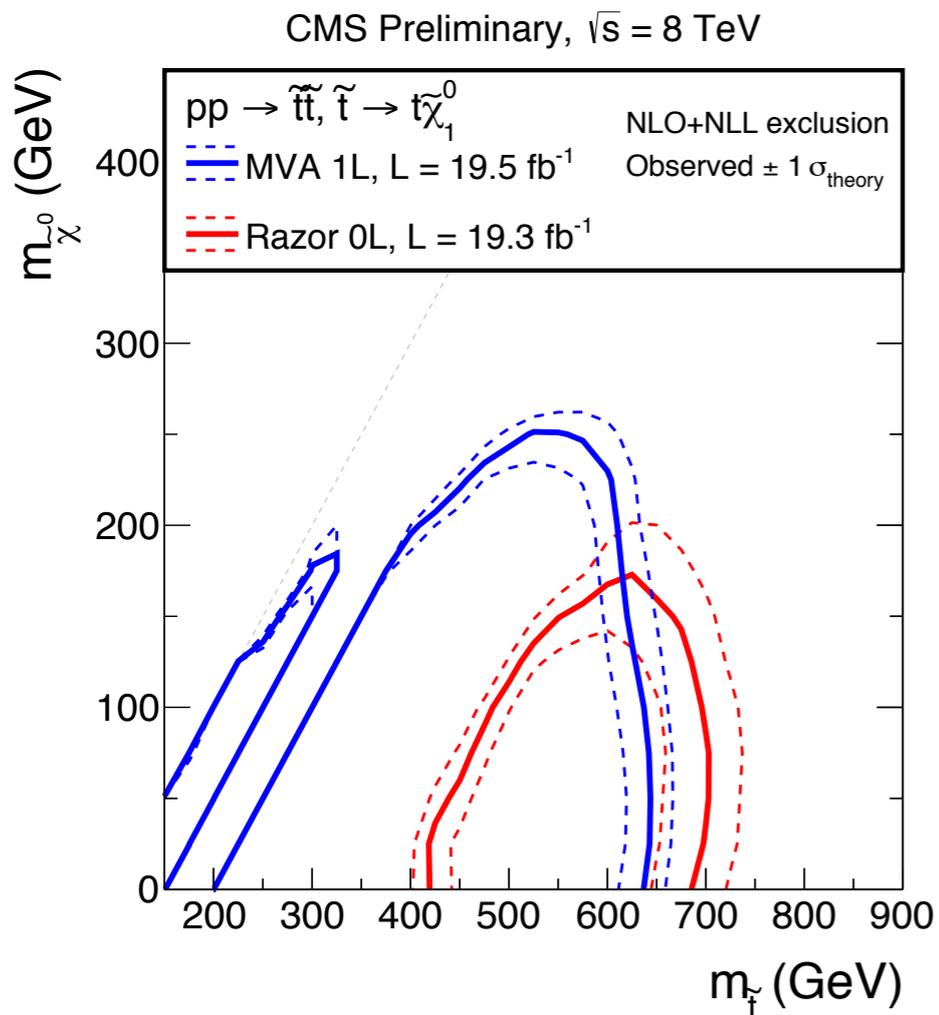
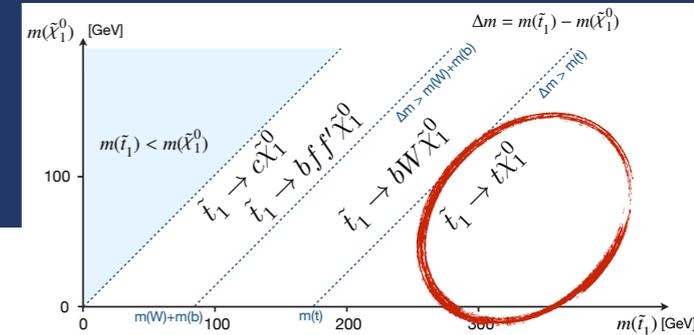
$$\tilde{t}_1 \rightarrow b \tilde{\chi}_1^\pm$$



Decay A with BR=100%,
 combination of SR A+B most sensitive for
 models with large $\Delta m(\text{stop}, N1)$.

Mixed decay A, B with varying BRs



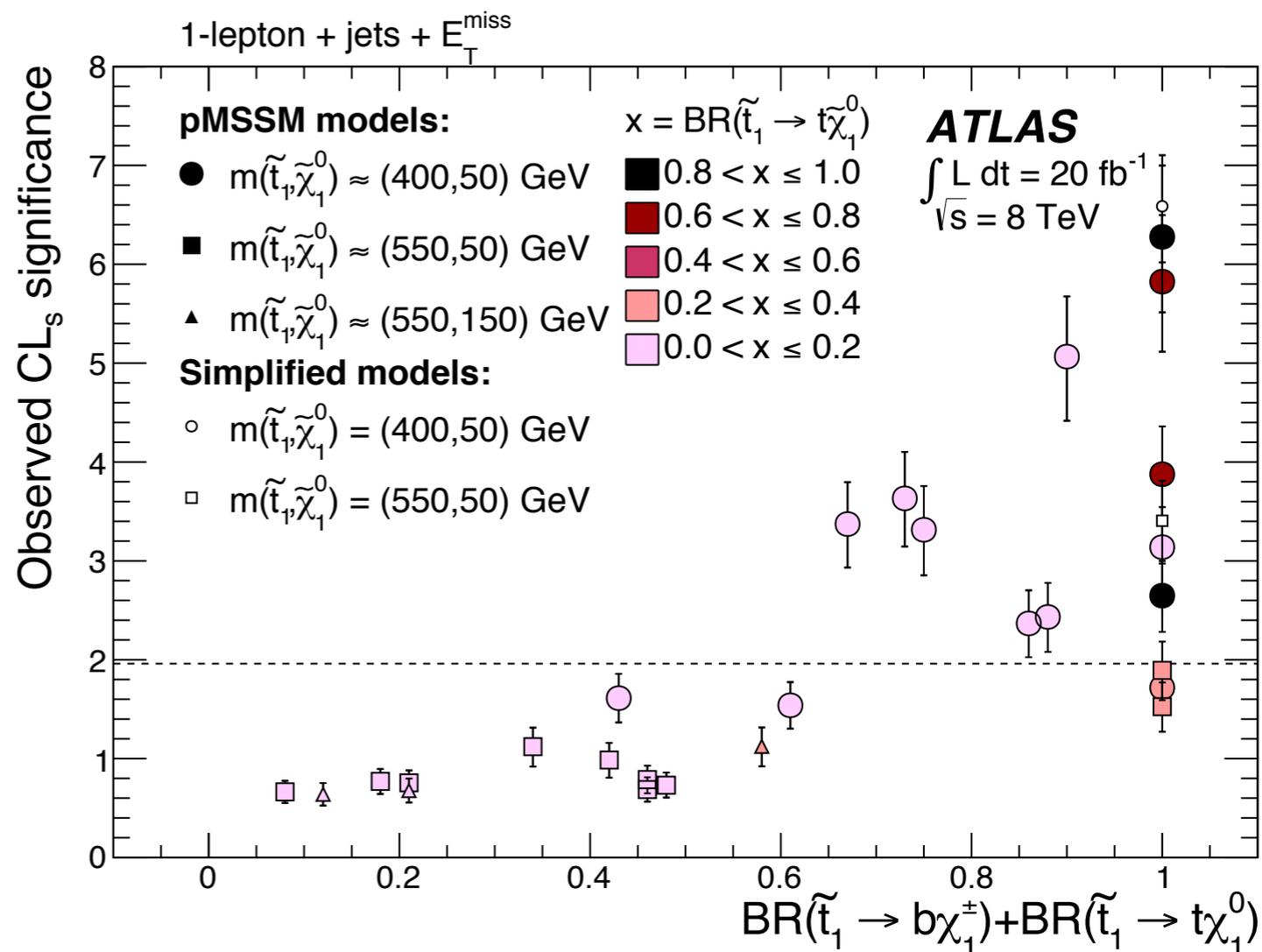


Statistical combination of

- Razor hadronic box, SUS-13-004
- stop 1L BDT, SUS-13-011

The change in sensitivity when varying parameters other than the stop and N1 masses is studied using pMSSM samples, classified into three groups of similar stop and N1 masses.

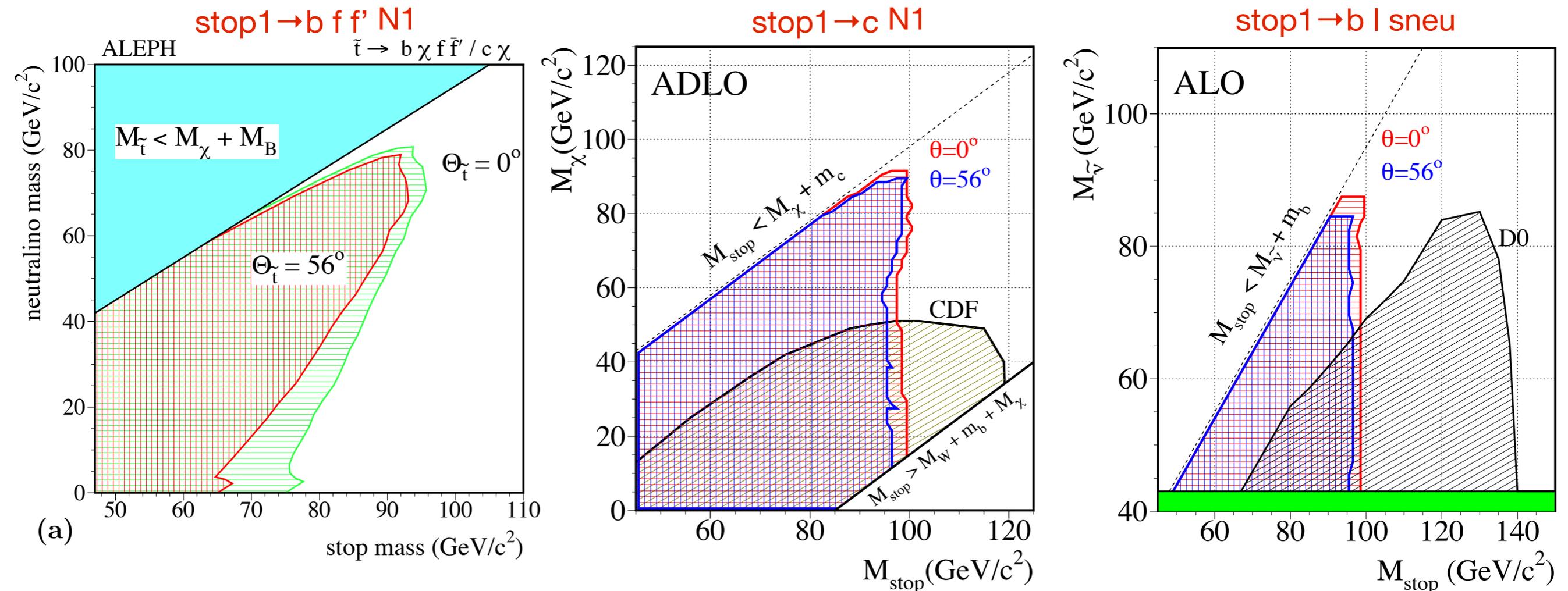
- strongest impact on the CLs significance from the sum of the branching ratios for $\text{stop1} \rightarrow \text{tN1}$ and $\text{stop1} \rightarrow \text{bC1}$.
 - signal selections being optimised using only simplified models.
- sensitivity also depends on the kinematic properties of the events, which are affected, e.g., by the stop mixing matrix and by the masses and field content of other SUSY particles.
 - explains the large spread in CLs significance



LEP

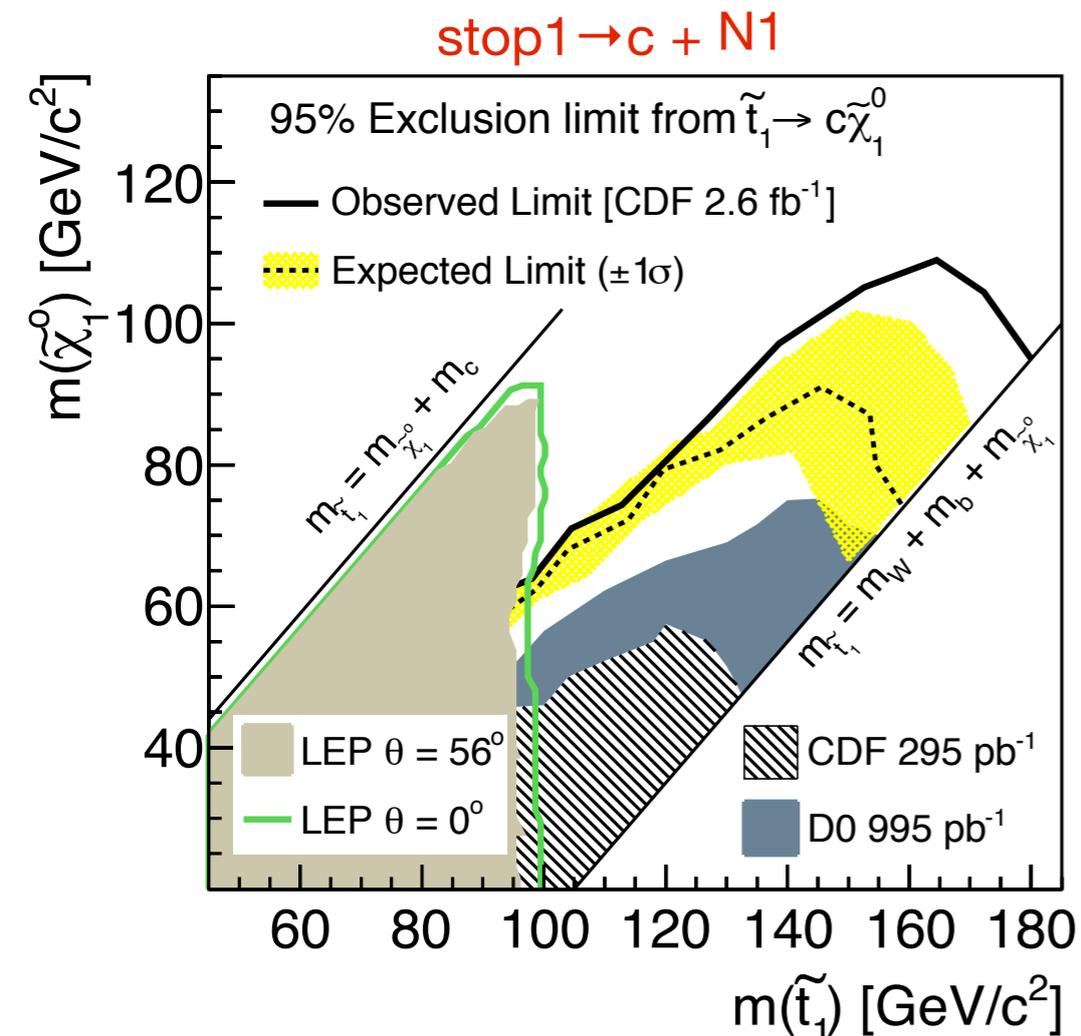
<http://lepsusy.web.cern.ch/lepsusy/>

- ALEPH, DELPHI, L3, and OPAL at LEP performed searches for the top squark considering the decays $stop1 \rightarrow c N1$ and $stop1 \rightarrow b l sneu$.
- In the latter, the sneutrino is assumed to be the LSP. (The sneutrino LSP is mostly considered to be ruled out because of direct detection experiments since it couples directly to the Z boson.)
- 2 different squark mixing angles: $\theta = 0$ (squark_Left), and in case of "max mixing" where the mixing angle is chosen such that the cross section is minimal.
- Furthermore, ALEPH published limits on the four-body decay, $stop1 \rightarrow b f f' N1$ [Phys. Lett. B537 (2002) 5–20]



Tevatron

- The CDF and DØ detectors (at the Tevatron collider) performed searches for the top squark in several analyses considering the decays: $\text{stop1} \rightarrow c N1$ (Ref.[19,20]), $\text{stop1} \rightarrow b l \text{ sneutr}$ (Ref.[21]), and $\text{stop1} \rightarrow b C1$ (Ref.[22]).
- Plot shows summary of Tevatron (and LEP) exclusion limits for the $\text{stop1} \rightarrow c N1$ decay scenario. Limits for the $\text{stop1} \rightarrow b l \text{ sneutr}$ decay are comparable, while the sensitivity for the $\text{stop1} \rightarrow b C1$ decay is considerably weaker: only a small window in top squark mass ($125 \text{ GeV} < m(\text{stop1}) < 138 \text{ GeV}$) is excluded under the assumption that the decay proceeds via a W boson ($C1 \rightarrow W + N1$).



[19] CDF Collaboration, *Search for Scalar Top Quark Production in $p\bar{p}$ Collisions at $\sqrt{s} = 1.96 \text{ TeV}$* , [JHEP **1210** \(2012\) 158](#).

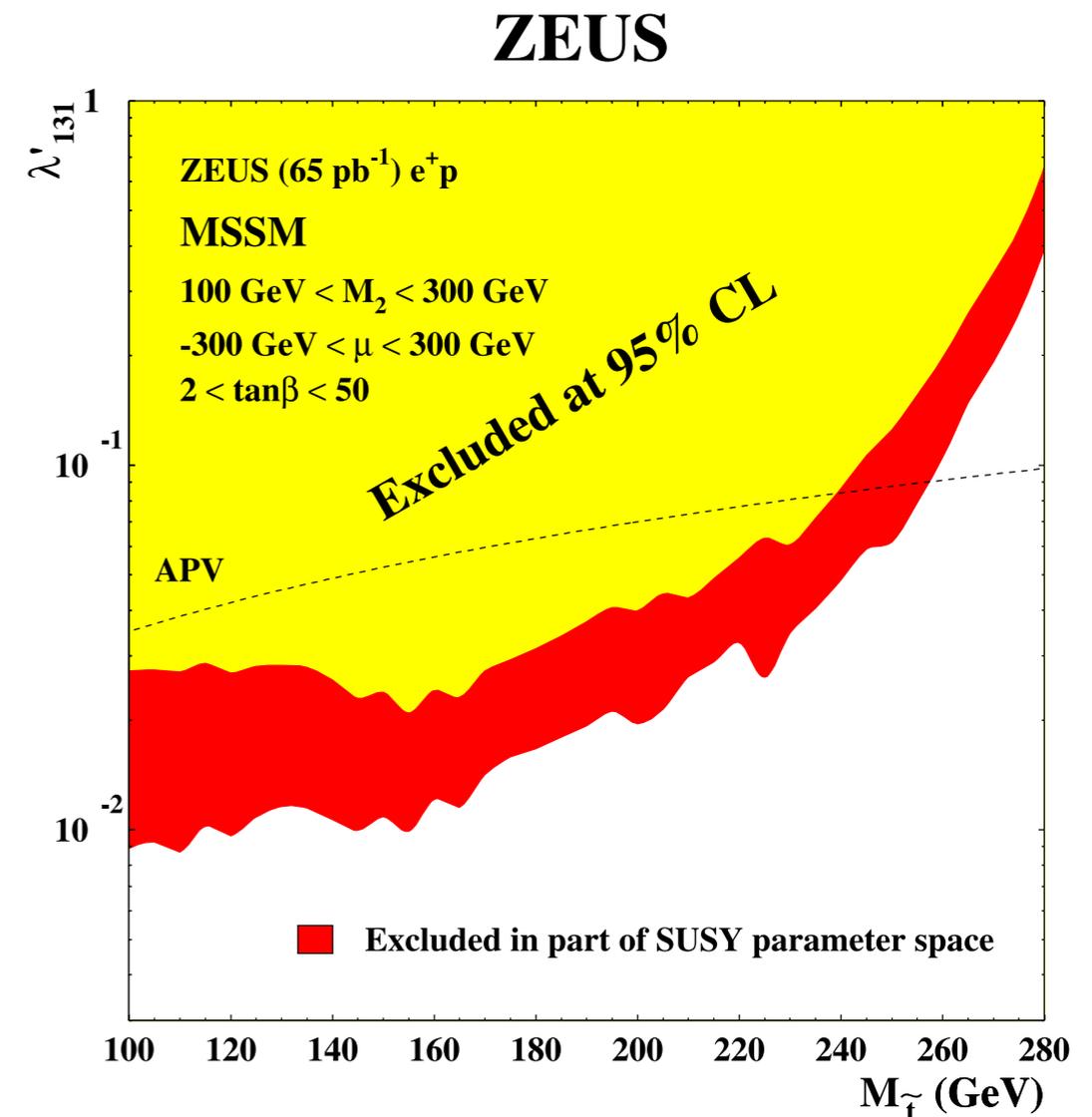
[20] D0 Collaboration, *Search for scalar top quarks in the acoplanar charm jets and missing transverse energy final state in $p\bar{p}$ collisions at $\sqrt{s} = 1.96\text{-TeV}$* , [Phys.Lett. **B665** \(2008\) 1–8](#).

[21] D0 Collaboration, *Search for the lightest scalar top quark in events with two leptons in $p\bar{p}$ collisions at $\sqrt{s} = 1.96 \text{ TeV}$* , [Phys. Lett. **B675** \(2009\) 289](#).

[22] CDF Collaboration, *Search for Pair Production of Supersymmetric Top Quarks in Dilepton Events from $p\bar{p}$ Collisions at $\sqrt{s} = 1.96 \text{ TeV}$* , [Phys. Rev. Lett. **104** \(2010\) 251801](#).

DESY

- A search for R-parity violating production of top squarks was performed by the ZEUS experiment at HERA (electron–proton collider at DESY).
- The top squark production depends on the strength of R-parity violation (λ'), and the decay was assumed to be either to $b+C1$ or R-parity violating to a lepton and a jet. A top squark lighter than 260 GeV was excluded for $\lambda' = 0.3$ and under specific model assumptions [1].



[1] ZEUS Collaboration, *Search for stop production in R-parity-violating supersymmetry at HERA*, [Eur.Phys.J. C50 \(2007\) 269–281](#).

Limit from $t\bar{t}$ cross section

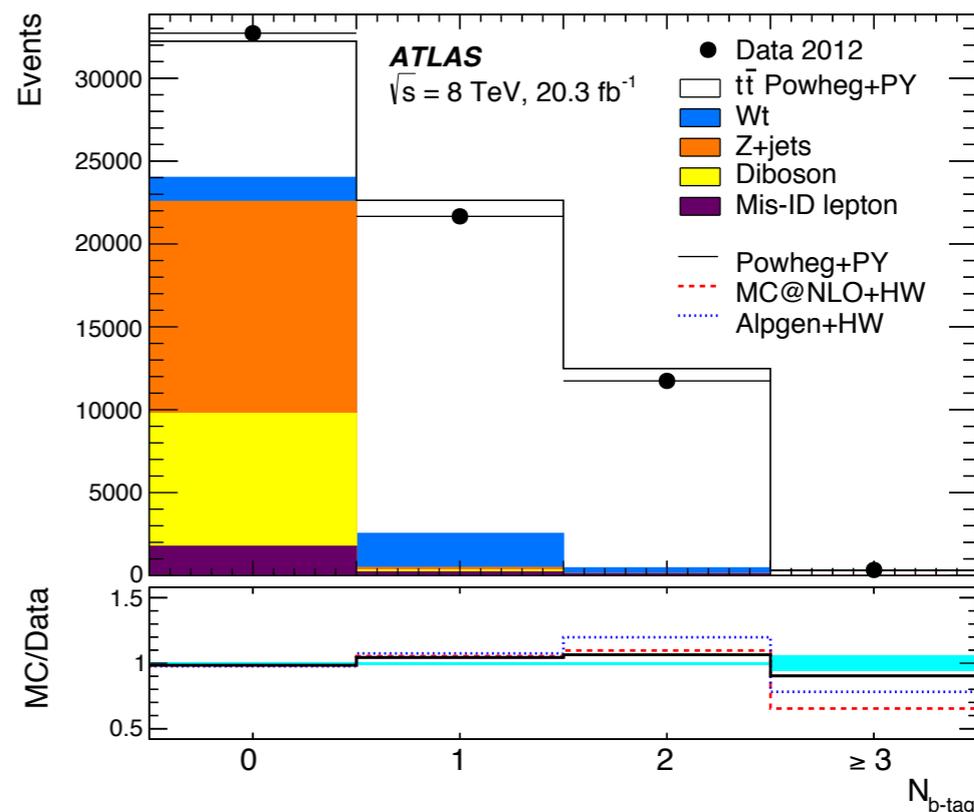
Submitted to EPJC (1406.5375)



- ▶ Exclusion set in terms of measured vs. theory cross-section difference Δ .
- ▶ The two datasets were fitted simultaneously, including the uncertainty due to a ± 1 GeV variation in the top quark mass.
- ▶ The limits were determined using a profile likelihood ratio, using nuisance parameters to account for correlated theoretical and experimental uncertainties.
- ▶ observed and expected limits on μ_{SIG} (signal strength) at the 95 % CL were extracted using the CLs prescription.
- ▶ Checking 8 / 7 TeV exclusions separately:
 - 8 TeV limits better than 'expected', 7 TeV slightly worse.
 - Not much difference in sensitivity, not much gained in combining (large correlations in theor. and exp. uncertainties).

Limit from $t\bar{t}$ cross section

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Measurement described in talk by
 Javier Brochero Cifuentes (Monday)

► Precise measurement of $t\bar{t}$ production cross-section at 7 and 8 TeV

- events with $e\mu$ (opp-sign) and exactly 1 or 2 b-tagged jets.
- Use ratio of 1 to 2 b-jet events to estimate probability of reconstructing and tagging a b-jet from the data – reduce JES, b-tagging, ISR/FSR, b-jet modelling uncertainties

► Results

$$\sigma_{t\bar{t}} = 182.9 \pm 3.1 \pm 4.2 \pm 3.6 \pm 3.3 \text{ pb } (\sqrt{s} = 7 \text{ TeV}) \text{ and}$$

$$\sigma_{t\bar{t}} = 242.4 \pm 1.7 \pm 5.5 \pm 7.5 \pm 4.2 \text{ pb } (\sqrt{s} = 8 \text{ TeV}),$$

- Uncertainties from data statistics, systematics, luminosity, beam energy
- Theory (including Δm_{top}) $177.3^{+11.5}_{-12.0} \text{ pb } (7 \text{ TeV}), 252.9^{+15.3}_{-16.3} \text{ pb } (8 \text{ TeV})$