

WbW \bar{b} at NLO

“NLO QCD corrections to WbW \bar{b} production: new developments and new issues”.

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Why are we interested in $Wb\bar{W}\bar{b}$ predictions.

In the experiment we do not measure tops. We only have a handle on their decay products.
 $Wb\bar{W}\bar{b}$ is therefore the more realistic final state if you are interested in $t\bar{t}$ production.

~ quantum mechanical versus semi-classical treatment ~

- important contributions to Wt and WW final states (tricky to disentangle at higher orders)
 - important background to BSM searches and SM measurements
(e.g. population of N-jet bins in WW production)
 - at current precision, we start worrying about offshell effects, non-factorizable corrections, b-mass dependence etc.
 - expect small ($\mathcal{O}(\Gamma_t/m_t)$) effects (wrt NWA) for inclusive $t\bar{t}$ observables
→ similar statements for more exclusive phase spaces?

(see Frank Krauss' and Fabrizio Caola's talks given earlier today)



NLO QCD corrections to WW $b\bar{b}$ production



[TOP2014 – sunny Cannes Mandelieu]

Jan Winter

– MPP Munich, Germany –

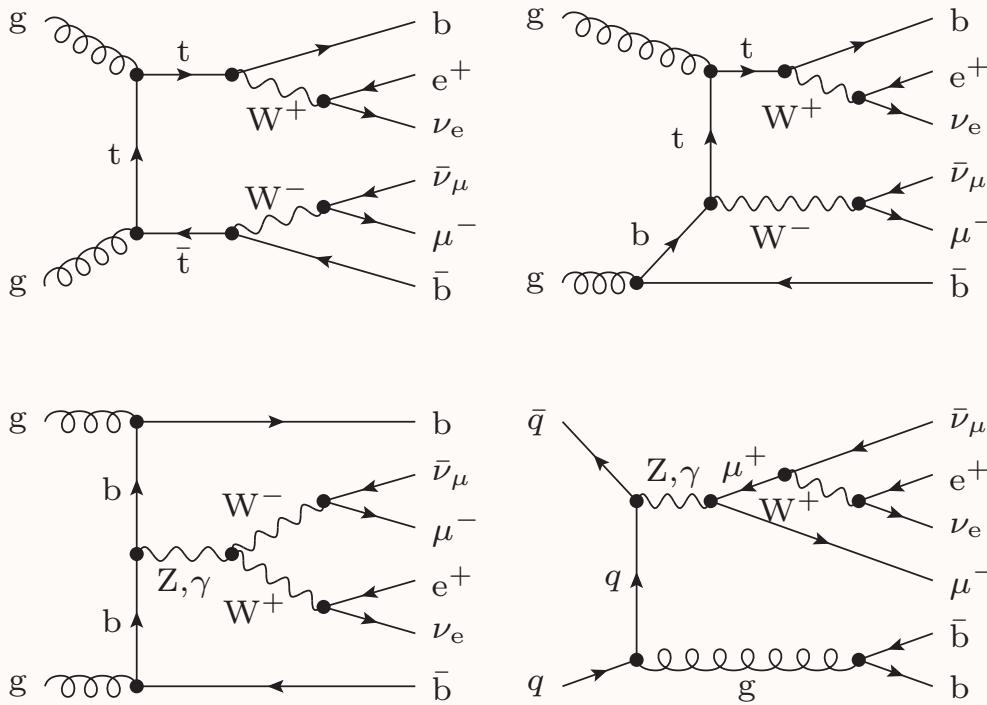


→ *Discussion of several recent results.*

- *WW $b\bar{b}$ at NLO in QCD – general remarks.*
- *Finite b-quark masses – WW $b\bar{b}$ in the 4-flavour scheme.*
- *Predicting the M(lb) distribution and consequences.*
- *Summary & conclusions.*

$WWb\bar{b}$ production at NLO in QCD

→ *Some introductory remarks on $WWb\bar{b}$ production at NLO in QCD.*



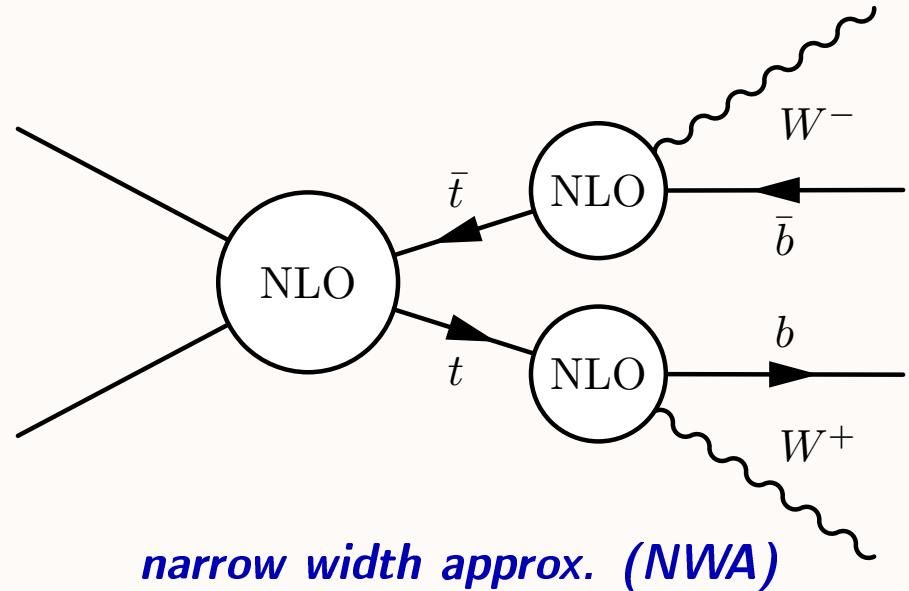
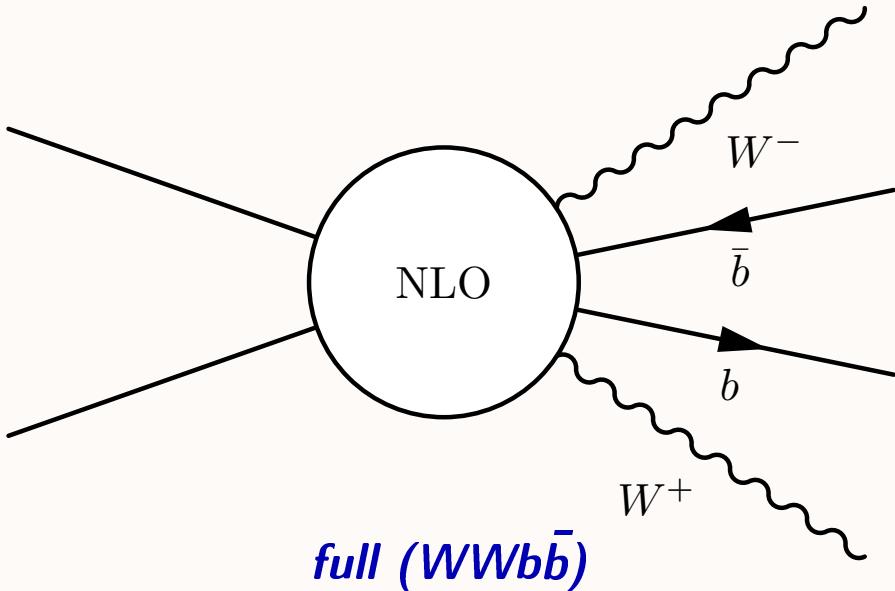
- full NLO treatment includes double-, single- and non-resonant contributions
(DR – $t\bar{t}$ -like) (SR – Wt -like) (NR – VV -like)
 - complex-mass scheme
 - finite top-quark and W width effects
(offshell DR, SR, NR and interferences)
 - first done in massless b-quark approximation
(→ requires two hard b-jets)
(→ $WWb\bar{b}$ in 5-flavour scheme)
- [DENNER ET AL. ARXIV:1012.3975, ARXIV:1207.5018]
[BEVILACQUA ET AL. ARXIV:1012.4230]

- earlier done in NWA ($\Gamma_t \rightarrow 0$ limit) where production and decay factorize (neglected contributions are suppressed by powers of $\Gamma_t/m_t \lesssim 1\%$)
[BERNREUTHER, BRANDENBURG, SI, UWER, ARXIV:HEP-PH/0403035]
[MELNIKOV, SCHULZE, ARXIV:0907.3090]

$$\lim_{\Gamma_t/m_t \rightarrow 0} \frac{1}{(p_t^2 - m_t^2)^2 + m_t^2 \Gamma_t^2} = \frac{\pi}{m_t \Gamma_t} \delta(p_t^2 - m_t^2)$$



Full calculation versus NWA

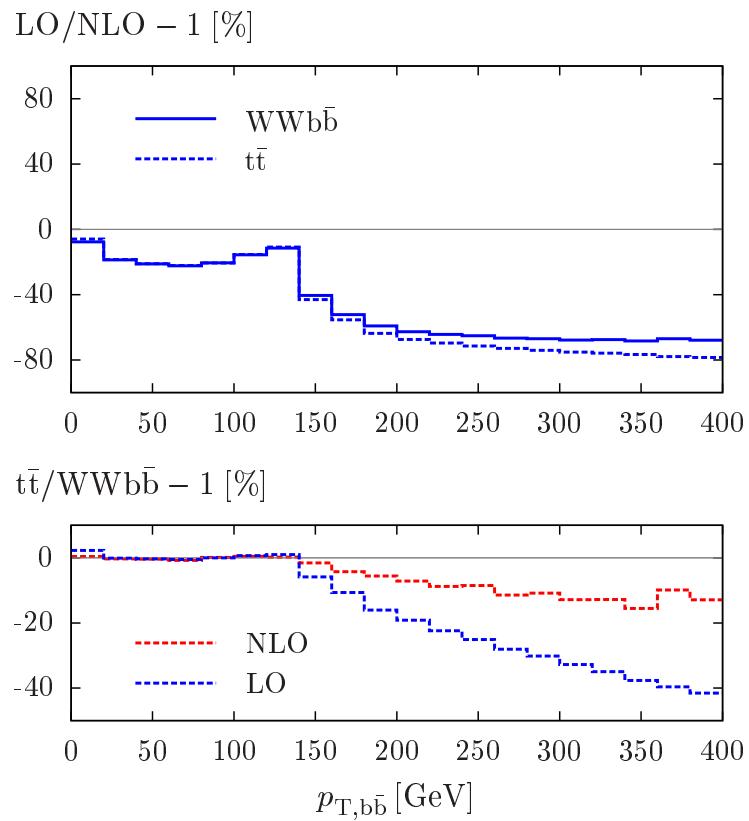
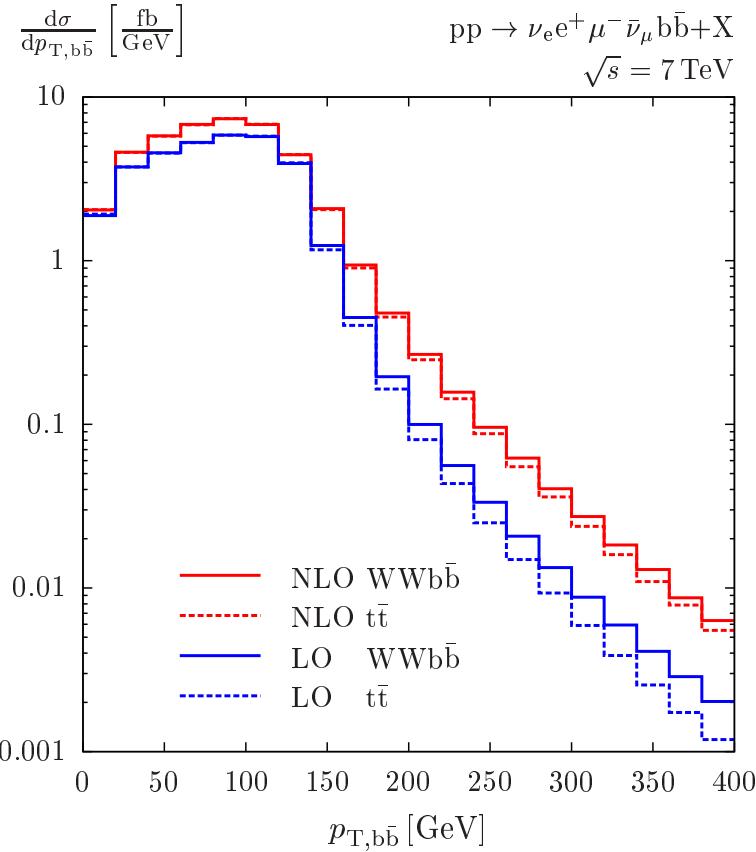
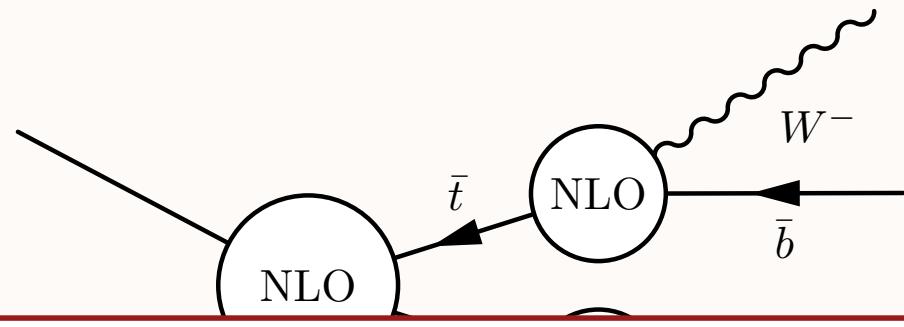
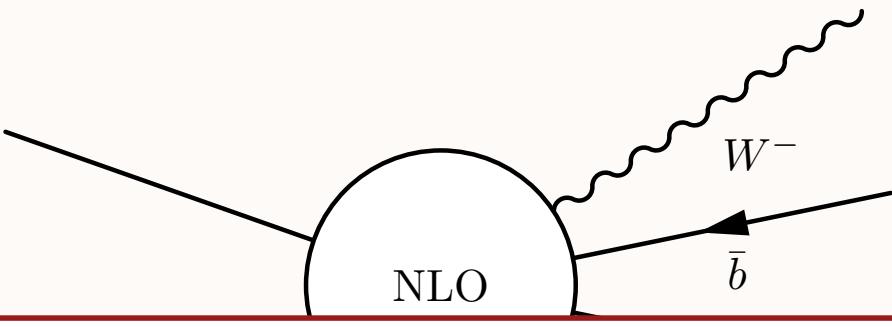


- full NLO description of the $WWb\bar{b}$ final state ($2 \rightarrow 4$ processes)
- non-resonant/-factorizing contributions (quantum interferences)
- NLO effects in top quark decays
- Comparison between both calculations (in the $\ell\ell$ channel) to investigate finite top-quark width effects.
 - No more than 1% deviations for inclusive cross sections (with experimental cuts).
 - Effects can be (significantly) larger in differential distributions.

⇒ [DENNER, DITTMAYER, KALLWEIT, POZZORINI, SCHULZE, LESHOUCES2011, ARXIV:1203.6803]



Full calculation versus NWA

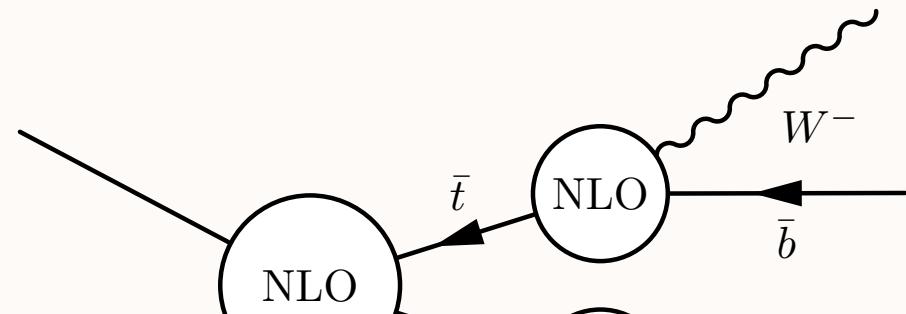
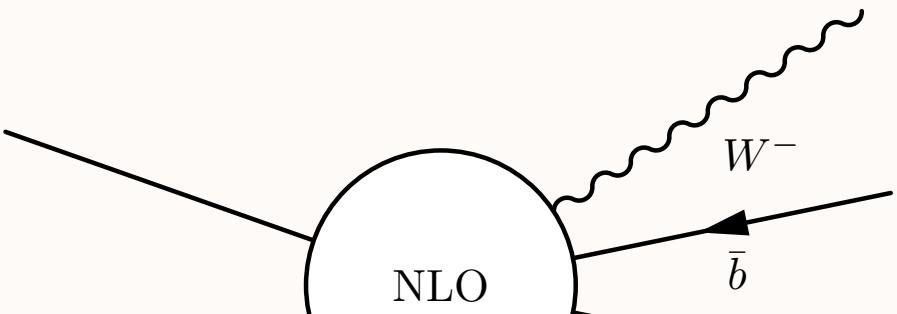


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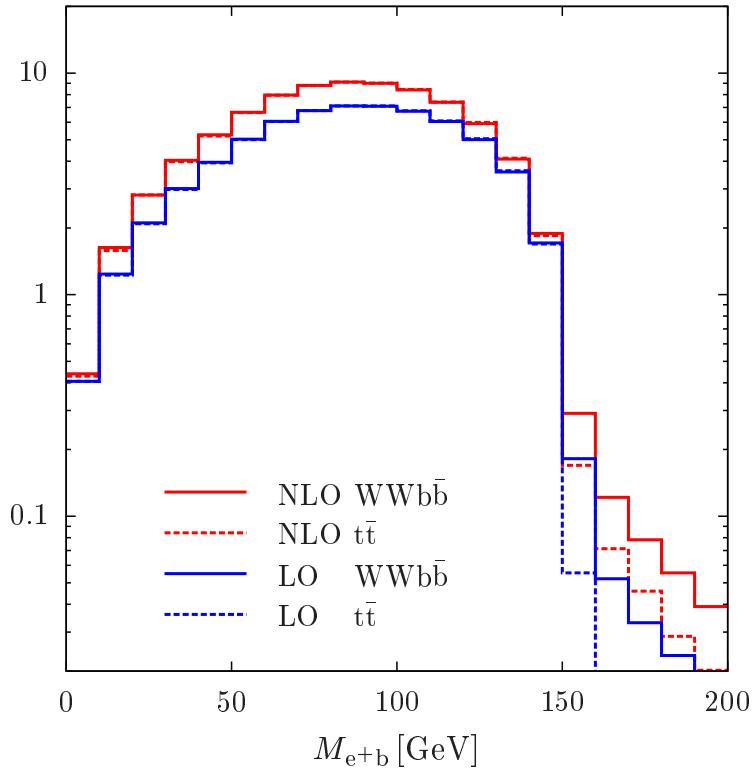
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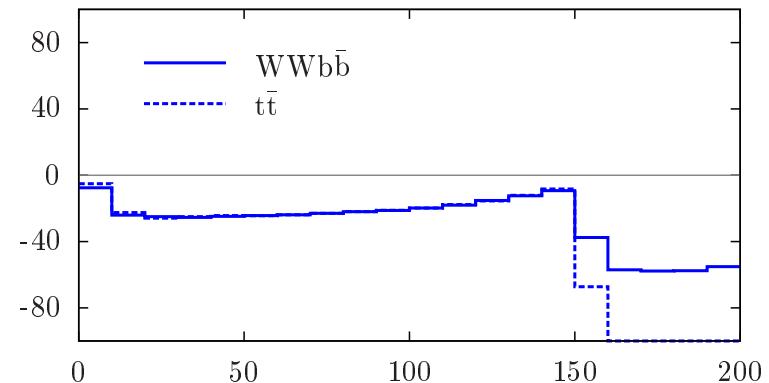
$$\frac{d\sigma}{dM_{e^+b}} \left[\frac{\text{fb}}{\text{GeV}} \right]$$

$$pp \rightarrow \nu_e e^+ \mu^- \bar{\nu}_\mu b\bar{b} + X$$

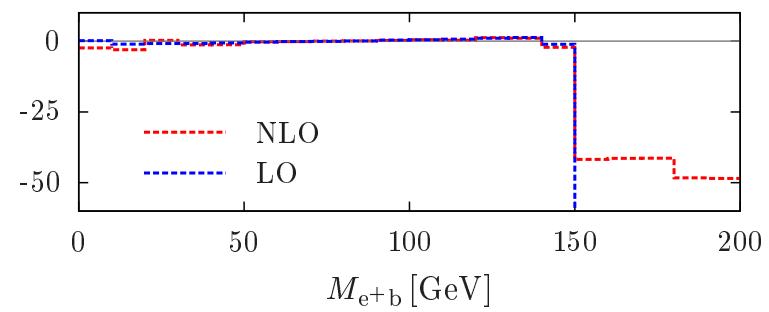
$\sqrt{s} = 7 \text{ TeV}$



$$\text{LO/NLO} - 1 [\%]$$



$$t\bar{t}/WWb\bar{b} - 1 [\%]$$



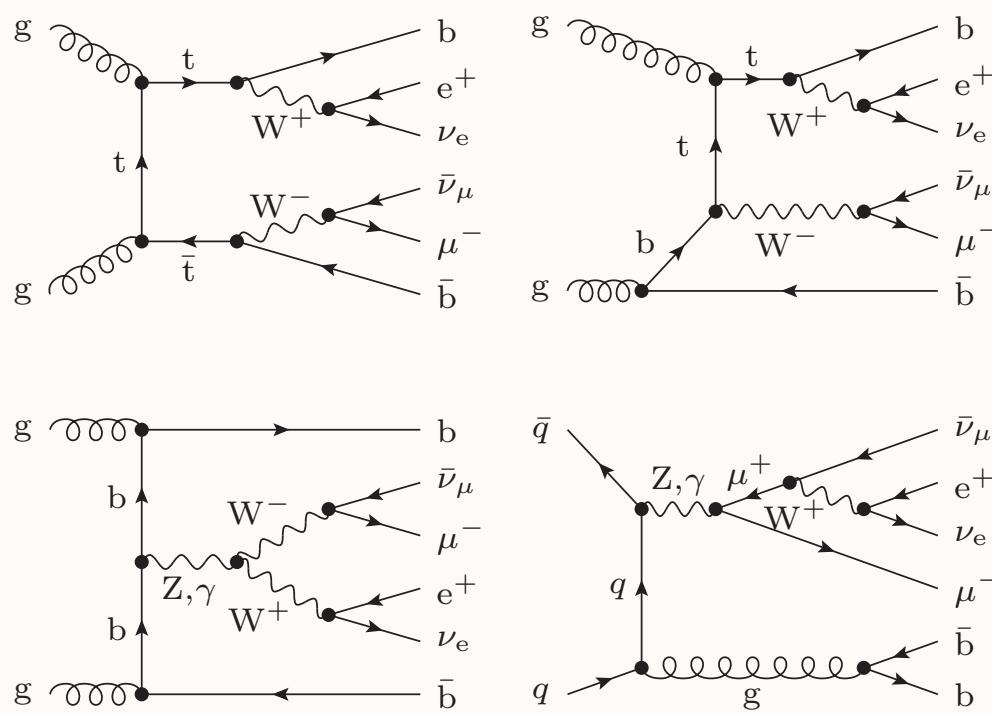
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$WWb\bar{b}$ production at NLO: massive b-quarks

→ New development: bottom quark mass included in the calculation.



- full NLO treatment includes double-, single- and non-resonant contributions
- complex-mass scheme
- finite top-quark and W width effects
- first done in massless b-quark approximation
[DENNER ET AL. ARXIV:1012.3975, ARXIV:1207.5018]
- [BEVILACQUA ET AL. ARXIV:1012.4230]
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[BERNREUTHER ET AL. ARXIV:HEP-PH/0403035]
- [MELNIKOV, SCHULZE, ARXIV:0907.3090]

- off-shell and single-top contributions more important in phase-space regions with unresolved b-quarks
- only accessible in calculations with massive b-quarks in the 4-flavour (4F) scheme
- in the 4F, fully differential NLO description of both FS b-jets → permits application of jet vetoes
- gauge-invariant separation of narrow-top-width contribution and finite-width remainder
- results provided recently by two groups:

[FREDERIX, ARXIV:1311.4893] [CASCIOLI, KALLWEIT, MAIERHÖFER, POZZORINI, ARXIV:1312.0546]

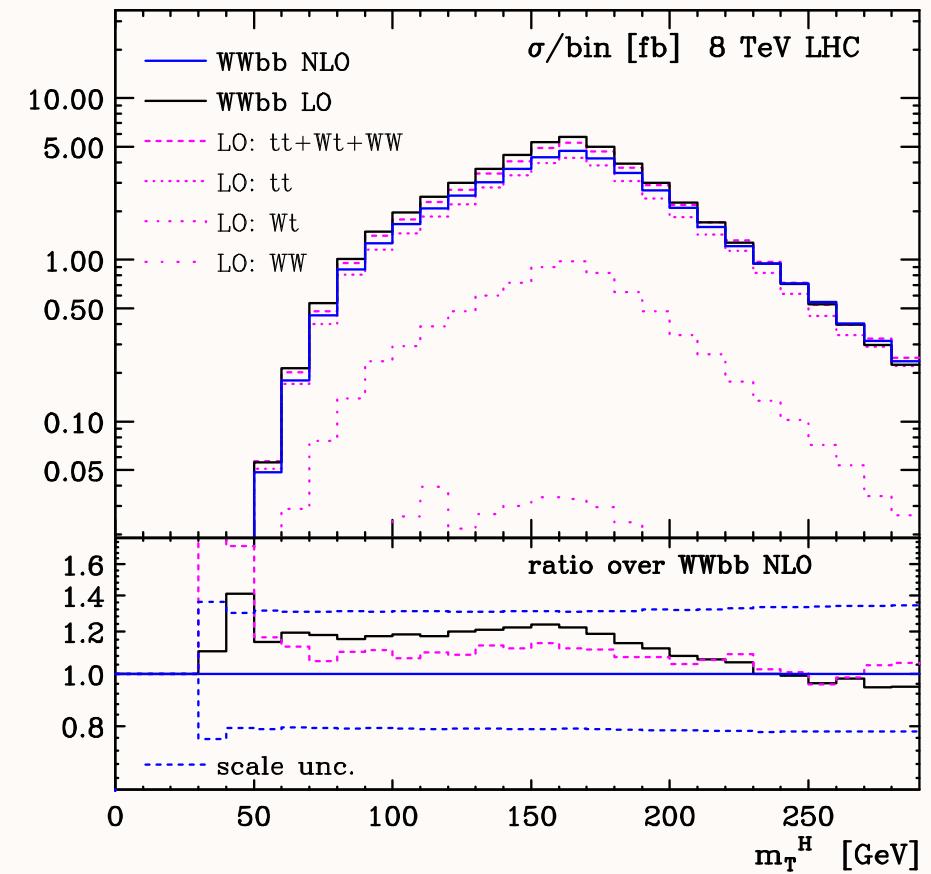
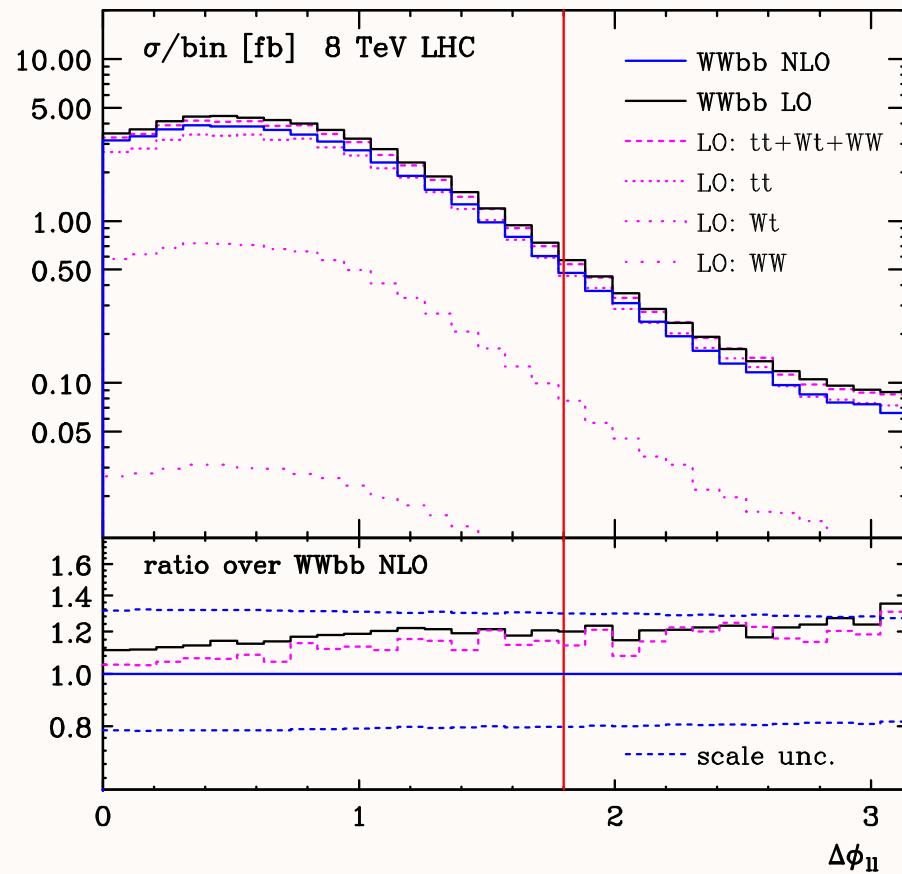


$WWb\bar{b}$ production at NLO

→ Calculation performed within *MadGraph5_aMC@NLO* framework.

[FREDERIX, ARXIV:1311.4893]

- top-quark induced backgrounds in $h \rightarrow WW^{(*)} \rightarrow ll\nu\nu$ channel at 8 TeV LHC ($\mu_{R,F} = \hat{H}_T/2$)
- “Higgs measurement” cuts in one-jet bin motivated by ATLAS analysis



- (left) azimuthal angle separation between leptons, (right) Higgs boson transverse mass
Higgs boson topology cuts: $m_{ll} < 50$ GeV and $|\Delta\phi_{ll}| < 1.8$



$WWb\bar{b}$ production at NLO

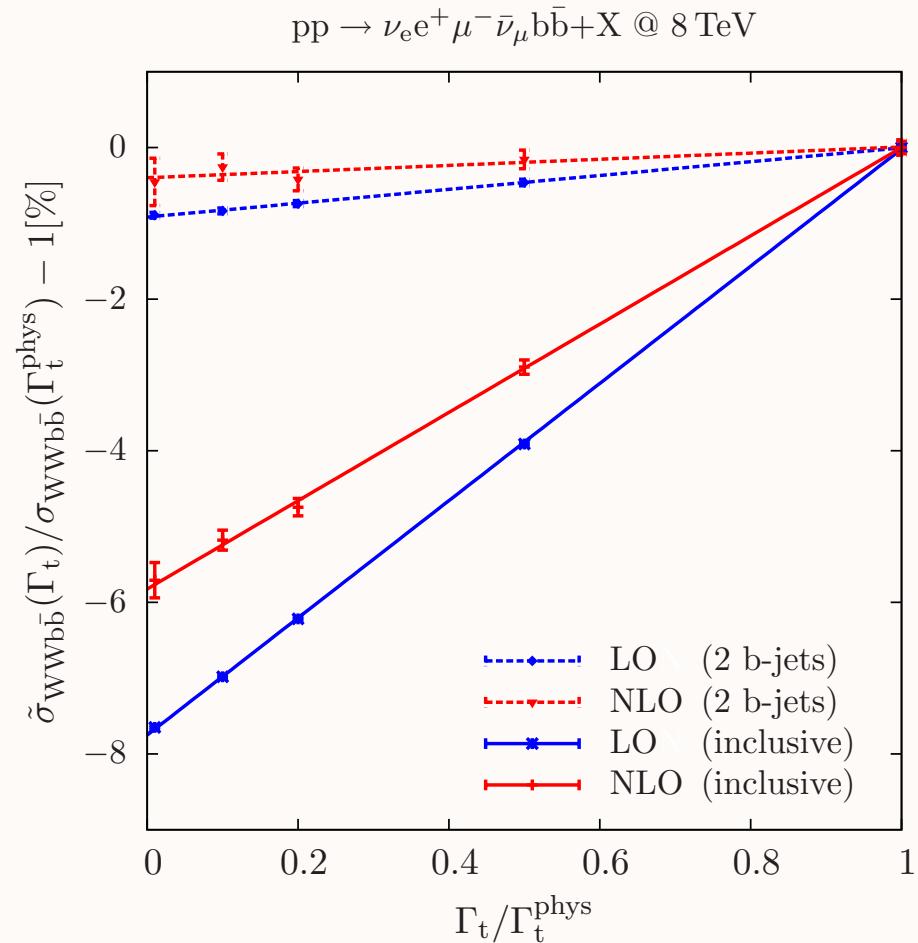
→ ***OpenLoops + Collier + New in-house NLO MC framework.***

[CASCIOLI ET AL. ARXIV:1312.0546]
[KALLWEIT]

- 4F scheme enables gauge-invariant $t\bar{t}$ /non- $t\bar{t}$ separation instead of ill-defined $t\bar{t}$ /Wt separation in 5F
- dynamical scale interpolating between $t\bar{t}$ ($\mu_{t\bar{t}}^2 = E_{T,t}E_{T,\bar{t}}$) and single- t ($\mu_{tW-}^2 = E_{T,t}E_{T,\bar{b}}$) ...
- ... to account for multiscale problem

numerical NWA

$$d\sigma_{t\bar{t}} = \lim_{\Gamma_t \rightarrow 0} \left(\frac{\Gamma_t}{\Gamma_t^{\text{phys}}} \right)^2 d\sigma_{WWb\bar{b}}(\Gamma_t)$$



- numerical extrapolation to $\Gamma_t \rightarrow 0$: Wt contribution dominates finite top-quark width remainder



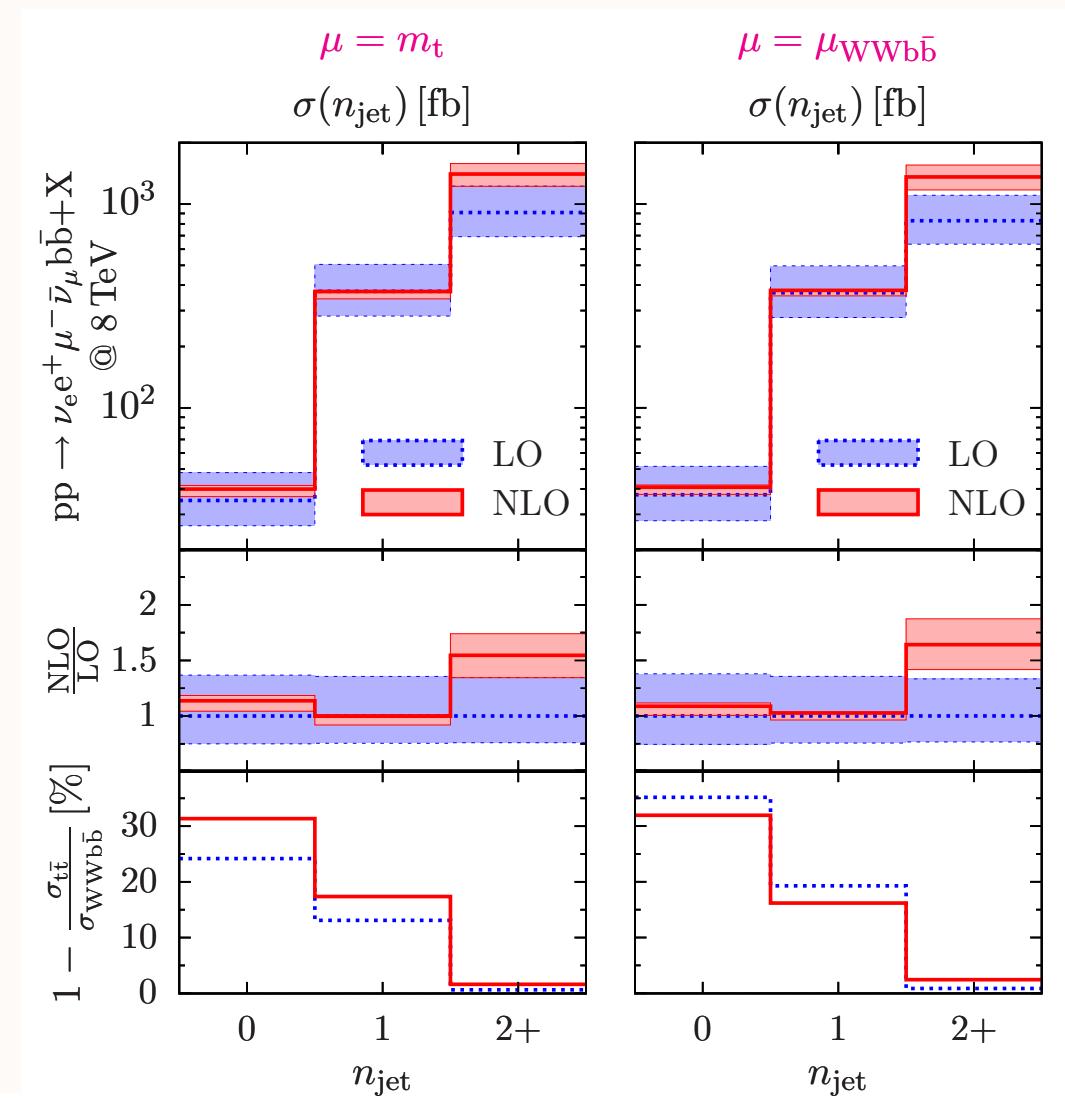
$WWb\bar{b}$ production at NLO

→ ***OpenLoops + Collier + New in-house NLO MC framework.***

[CASCIOLI ET AL. ARXIV:1312.0546]
[KALLWEIT]

- Controlling the $t\bar{t}$ background in WW – NLO and finite top quark width effects in jet bins
- dynamical scale interpolating between $t\bar{t}$ ($\mu_{t\bar{t}}^2 = E_{T,t}E_{T,\bar{t}}$) and single- t ($\mu_{tW^-}^2 = E_{T,t}E_{T,\bar{b}}$)

- very interesting application of finite b-mass calculation
- the 40% inclusive NLO correction is driven by the large two-jet bin correction
- strongly enhanced finite top quark width effects in zero- and one-jet bins

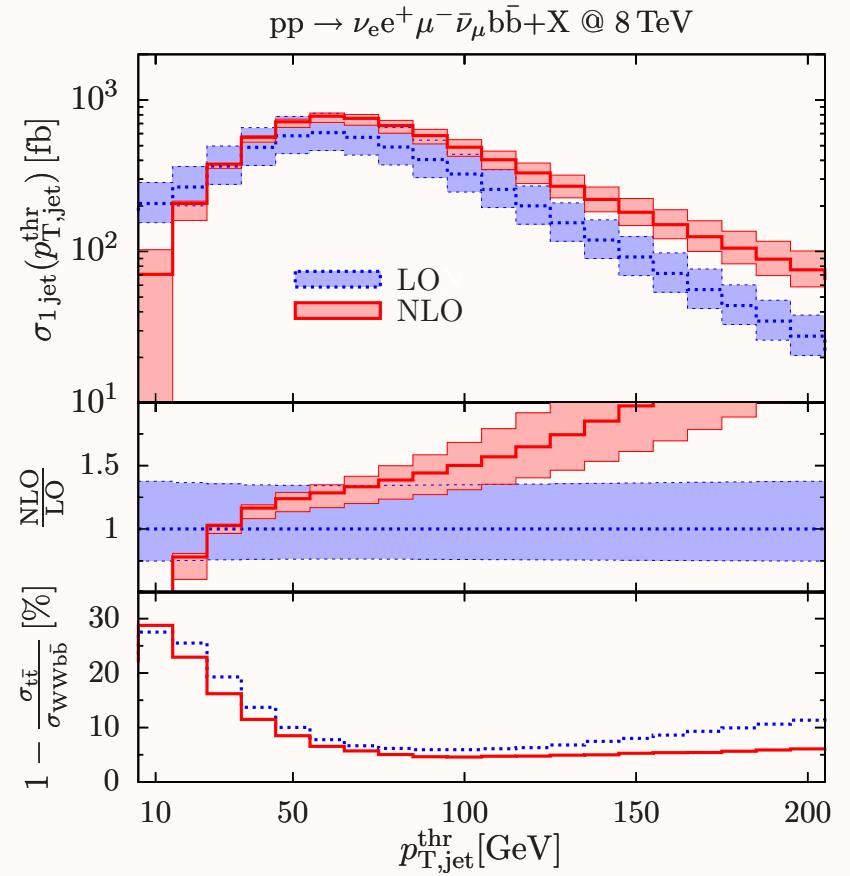
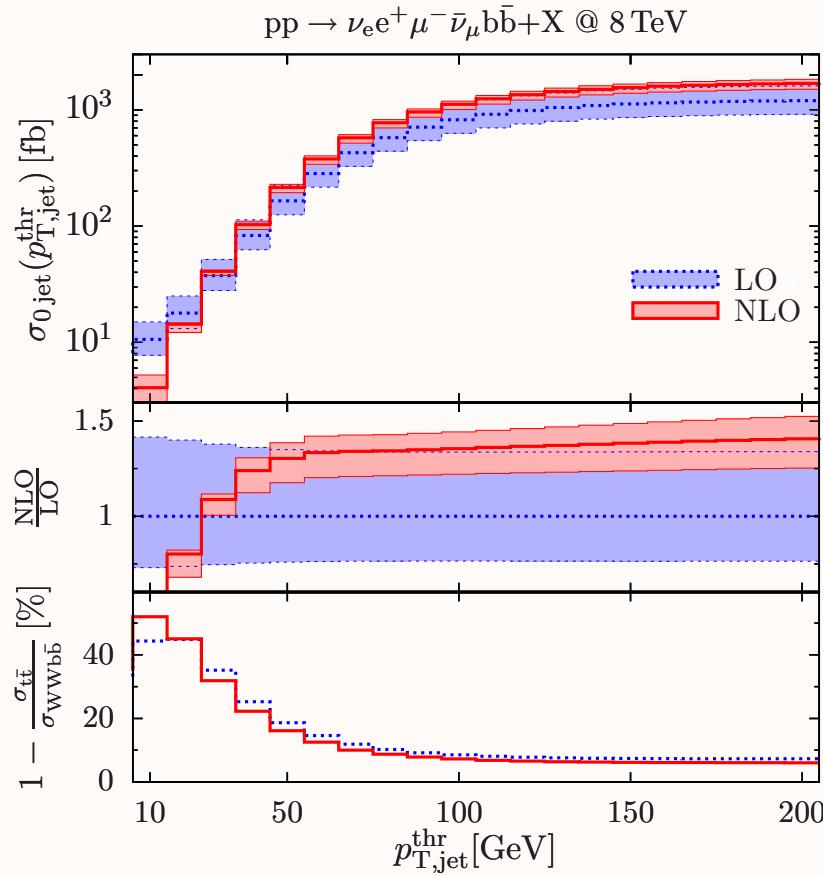


$WWb\bar{b}$ production at NLO

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- xsec scale uncertainties of 10 – 15% (similar for $\mu_0 = m_t$), 6% due to finite- t -width corrections



- cross section in exclusive zero- and one-jet bins as a function of jet p_T threshold

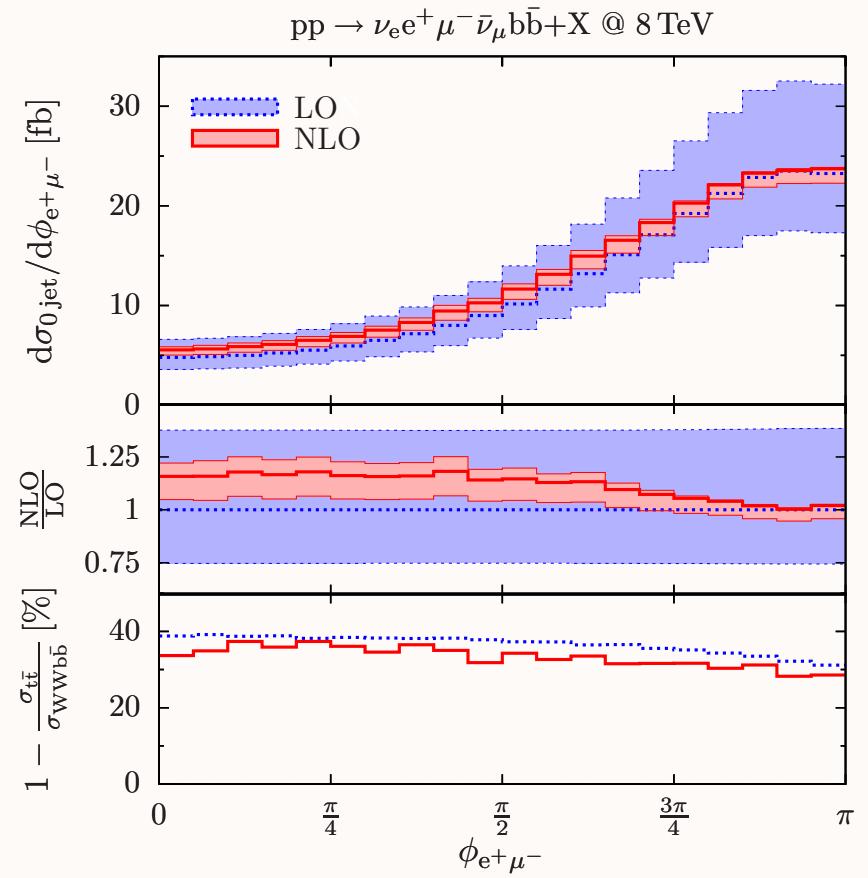
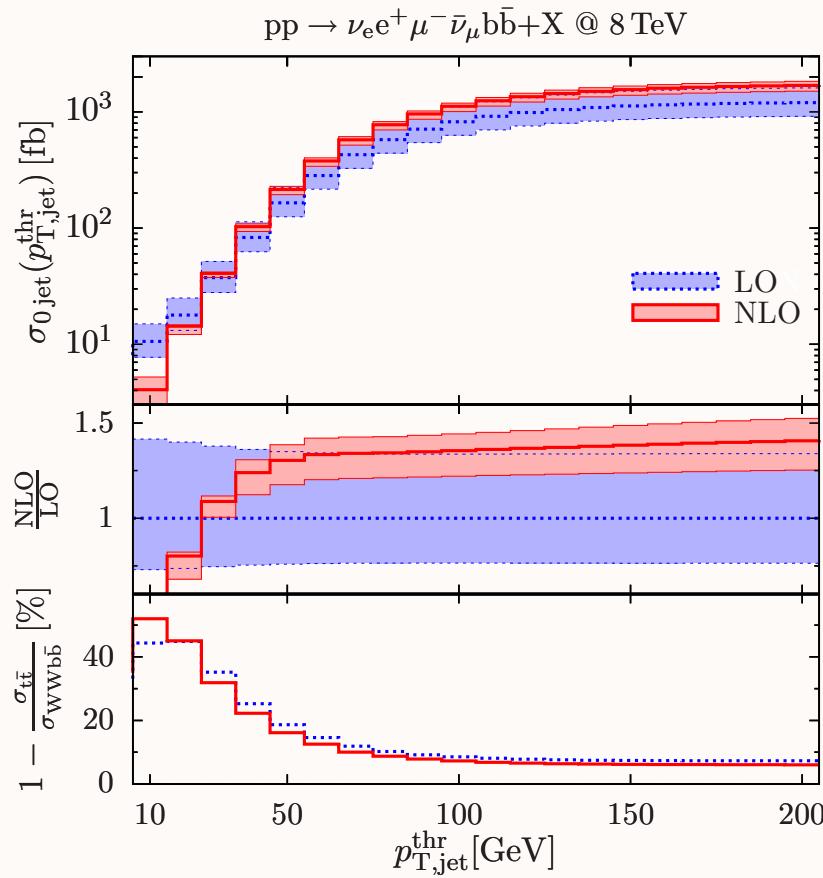


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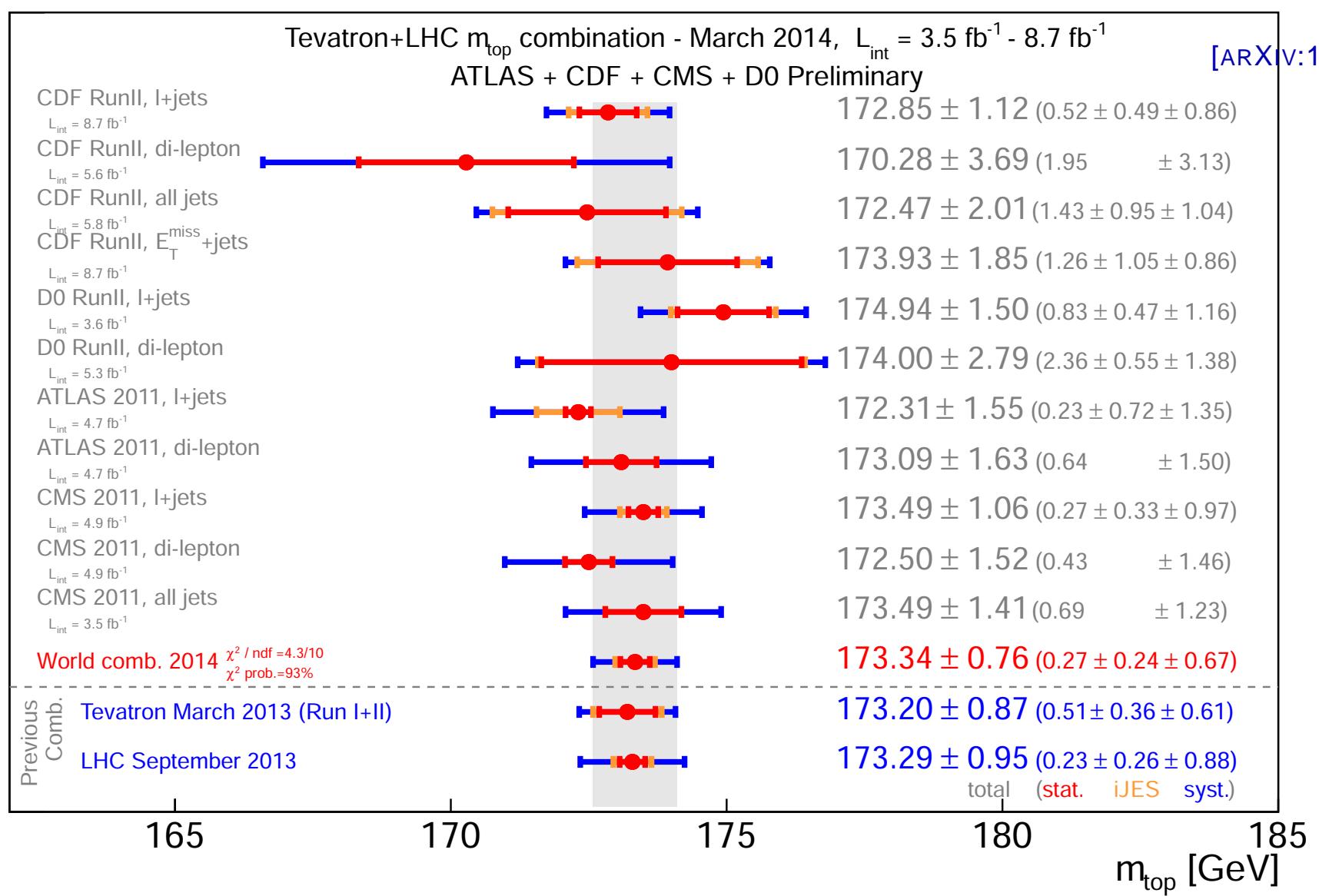
- cross section in exclusive zero- and one-jet bins as a function of jet p_T threshold
- (right) azimuthal angle between leptons (0-jet), finite- t -width effects increase with harsher vetoes



Top quark mass measurements

- first LHC+Tevatron result; total uncertainty on top quark mass < 1 GeV for combinations.

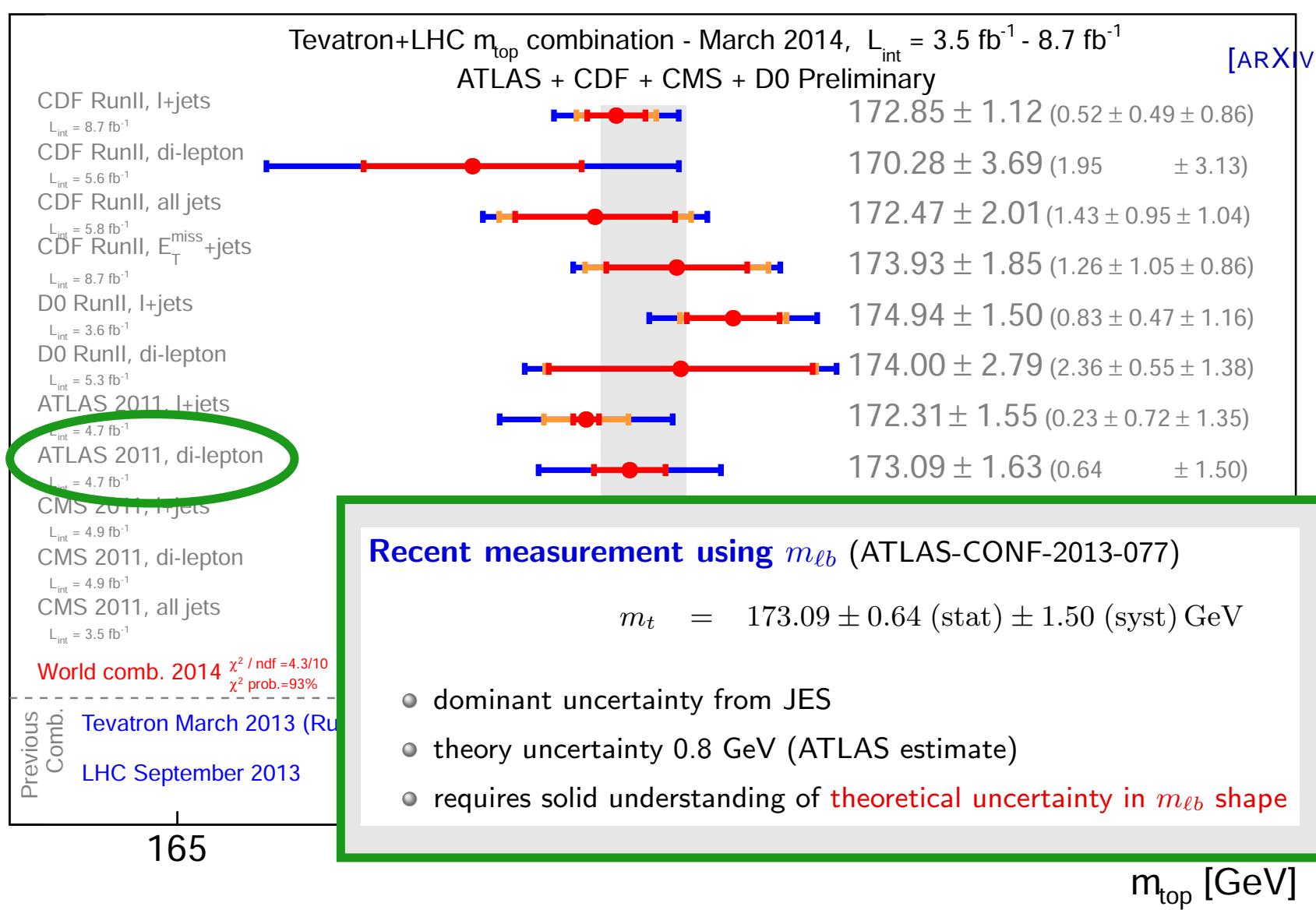
Rich and active experimental program (various complementary techniques).



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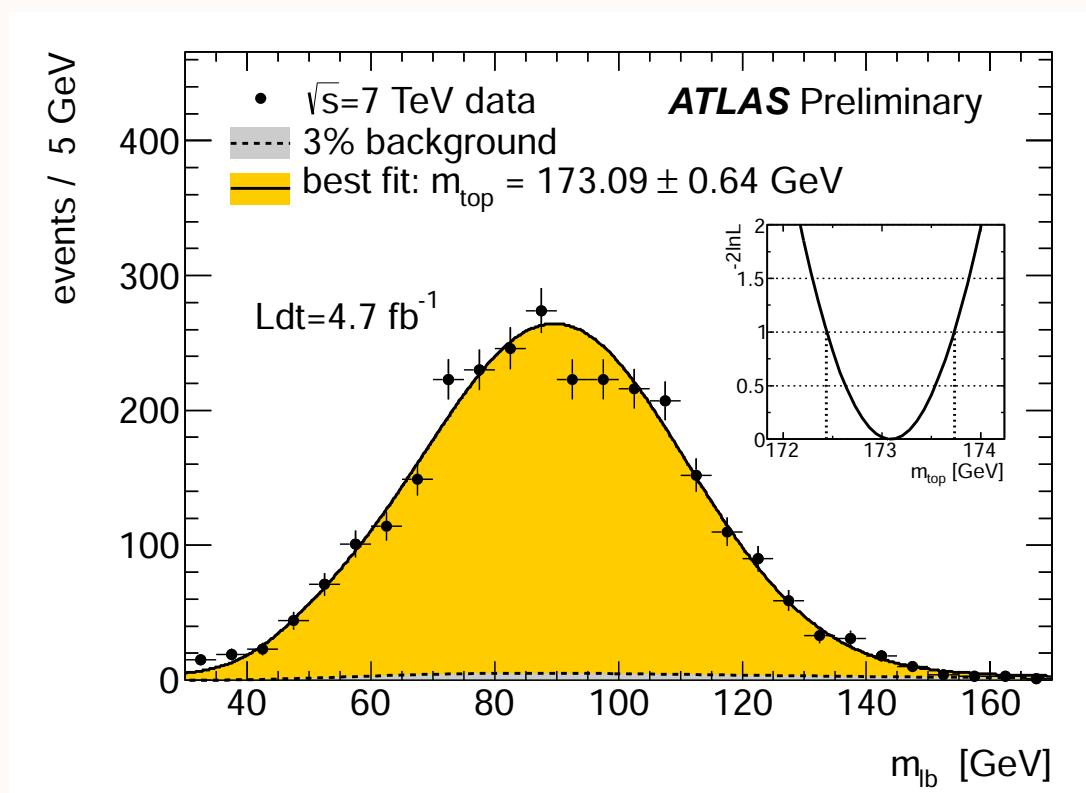
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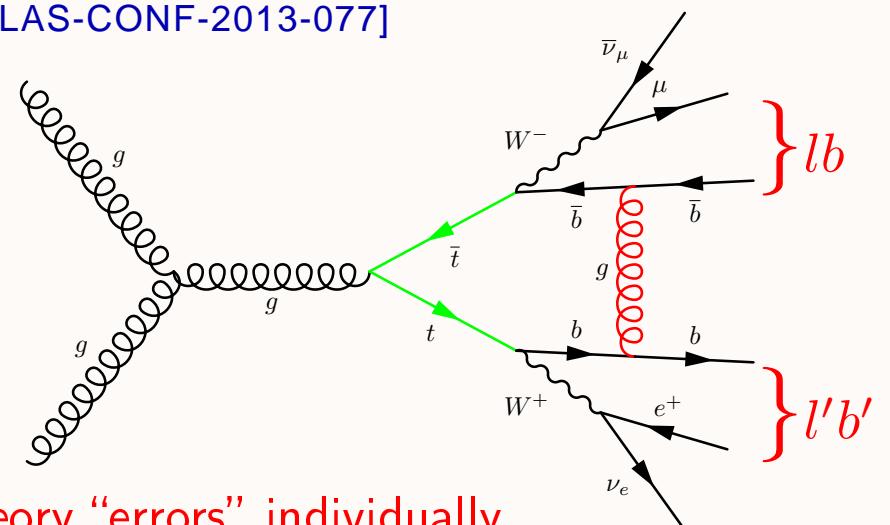
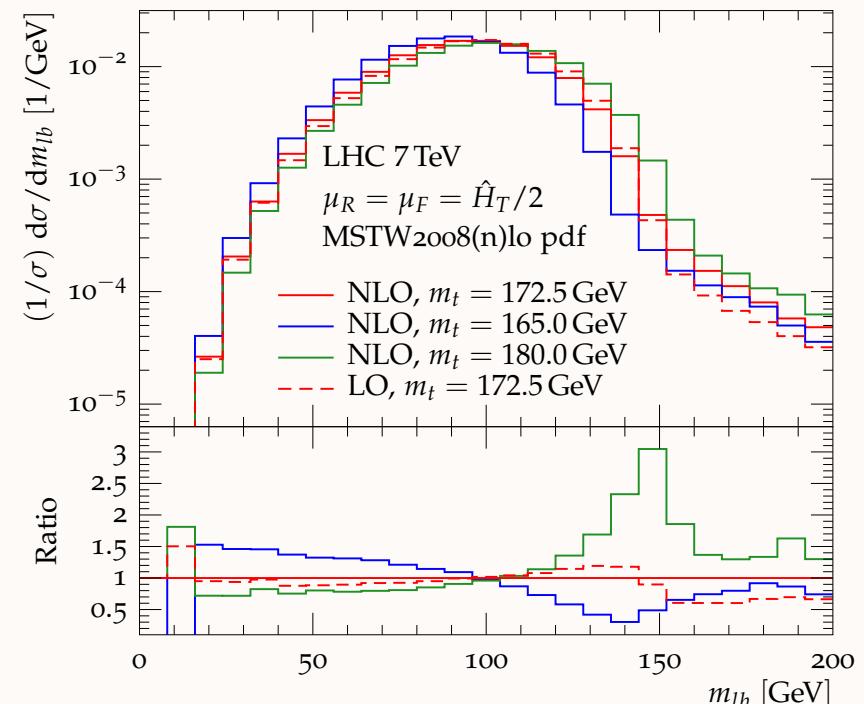
Top quark mass determination using the m_{lb} method

Parametrize “your” theory (m_{lb} predictions).

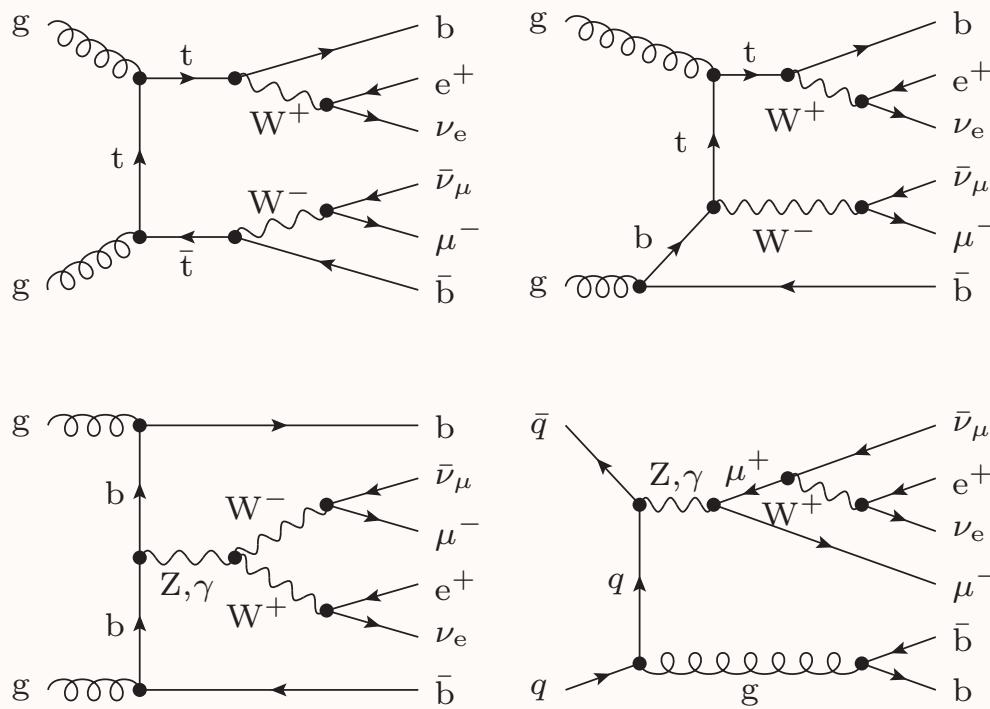
- Full QCD NLO prediction for $W^+W^-b\bar{b}$ in dilepton channel: m_{lb} distribution is sensitive to top quark mass.
- ATLAS uses one-dim. template method to determine m_t . Theory uncertainty has been estimated to 0.8 GeV.
→ Verify size of th. uncertainties using more advanced calc's!



Use pseudo-data to study different types of theory “errors” individually.



$WWb\bar{b}$ production at NLO



- full NLO treatment includes double-, single- and non-resonant contributions
- complex-mass scheme
- finite top-quark and W width effects
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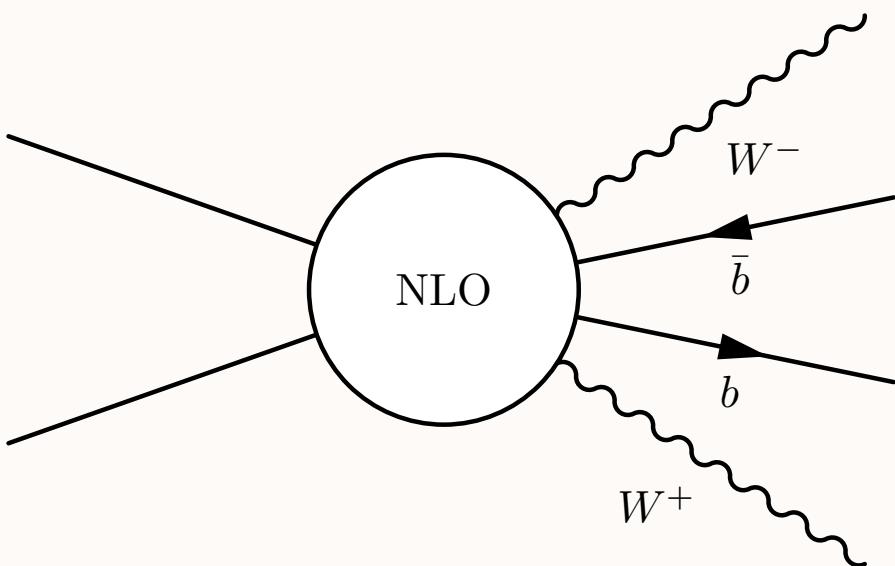
→ Our parton level calculations ...

- use the **GoSam+Sherpa** combined generator package (current versions, GoSam 2.0 and Sherpa 2.1).
- **Sherpa** for calculating Born, real corrections and infrared subtractions [GLEISBERG ET AL, ARXIV:0811.4622]
- **GoSam** for calculating virtual corrections [CULLEN, VANDEURZEN, GREINER, HEINRICH ET AL, ARXIV:1404.7096]
- 5-flavour scheme, massless b -quarks, two resonant W decaying leptonically @ LO respecting spin correlations

⇒ [HEINRICH, MAIER, NISIUS, SCHLENK, WINTER, ARXIV:1312.6659]

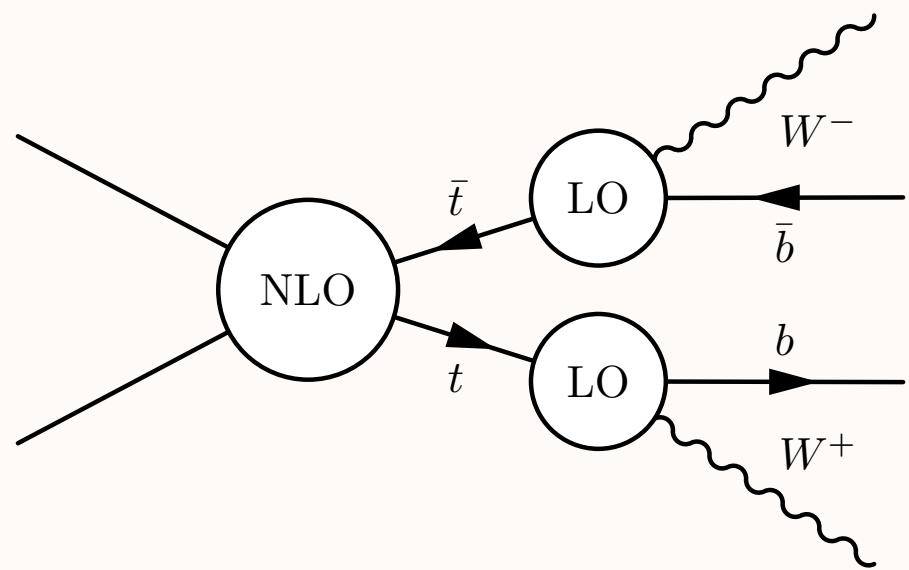


Full versus factorized approach



full ($WW\bar{b}\bar{b}$)

- full NLO description of the $WW\bar{b}\bar{b}$ final state ($2 \rightarrow 4$ processes)
- accounts for non-resonant/non-factorizing contributions, includes NLO effects in top quark decays
- Use these calculations for pure parton level analyses, i.e. m_t is not a MC mass here, it is the pole mass.



factorized ($t\bar{t}$)

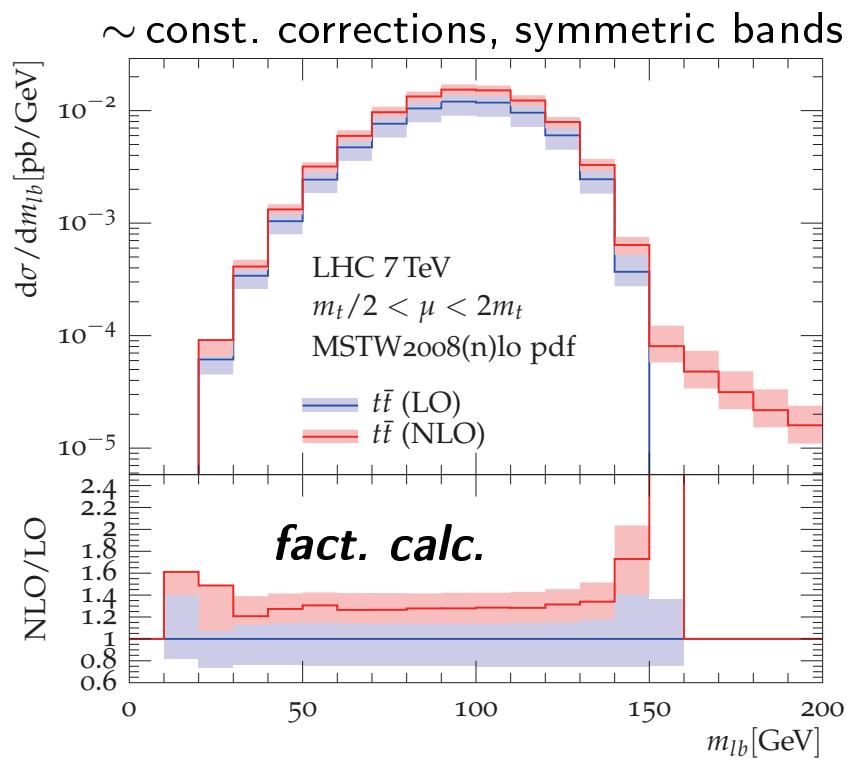
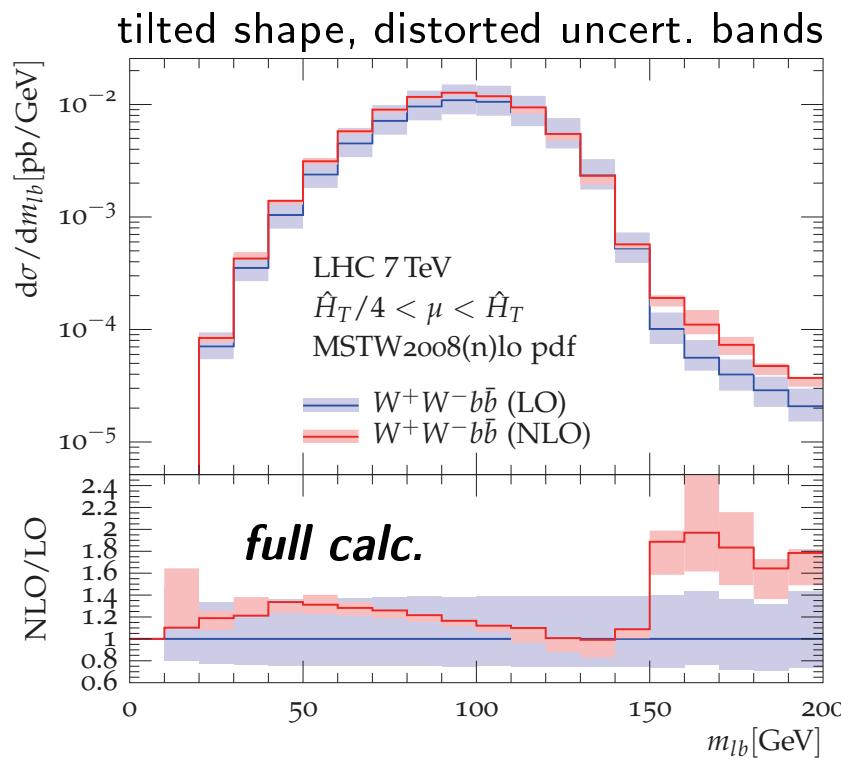
- NLO $t\bar{t}$ production ($2 \rightarrow 2$ processes) with LO decays attached and spin correlations preserved
- standard description for the NLO core in NLO+PS matching



The m_{lb} distribution at NLO and scale variations

→ Parton-level NLO calculations for $W^+W^-b\bar{b}$ based on GoSam+Sherpa framework.

(full & factorized calc., 5-flavour scheme, massless b-quarks, two resonant W decaying leptonically @ LO)



- ▶ Important NLO corrections to the shape of m_{lb}
- ▶ Values of m_{lb} larger than $\sqrt{m_t^2 - m_W^2}$ are kinematically forbidden in narrow width approximation at LO

- follow ATLAS strategy: use charged-lepton b-jet pairing minimizing sum of both m_{lb} and average.

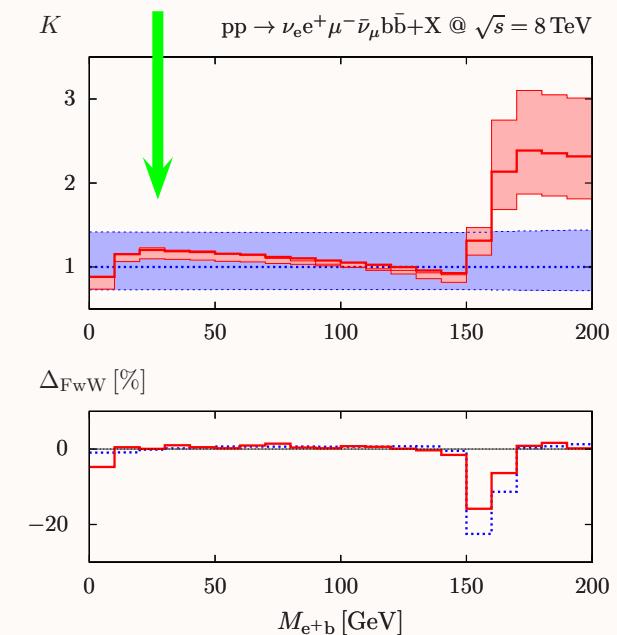
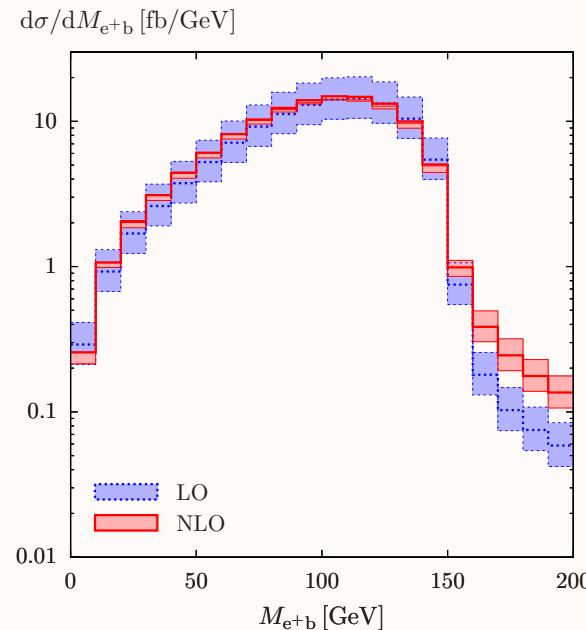
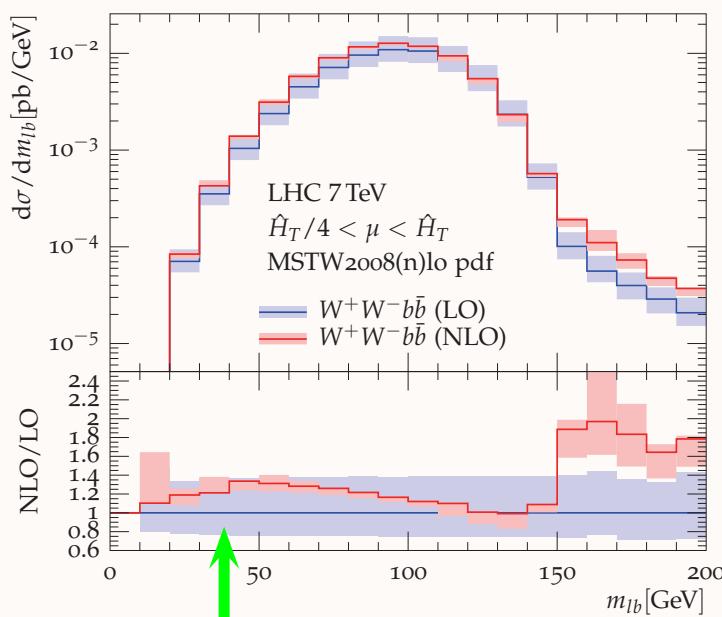


Qualitative comparison of m_{lb} predictions

[DENNER, DITTMAYER, KALLWEIT, POZZORINI, ARXIV:1207.5018]

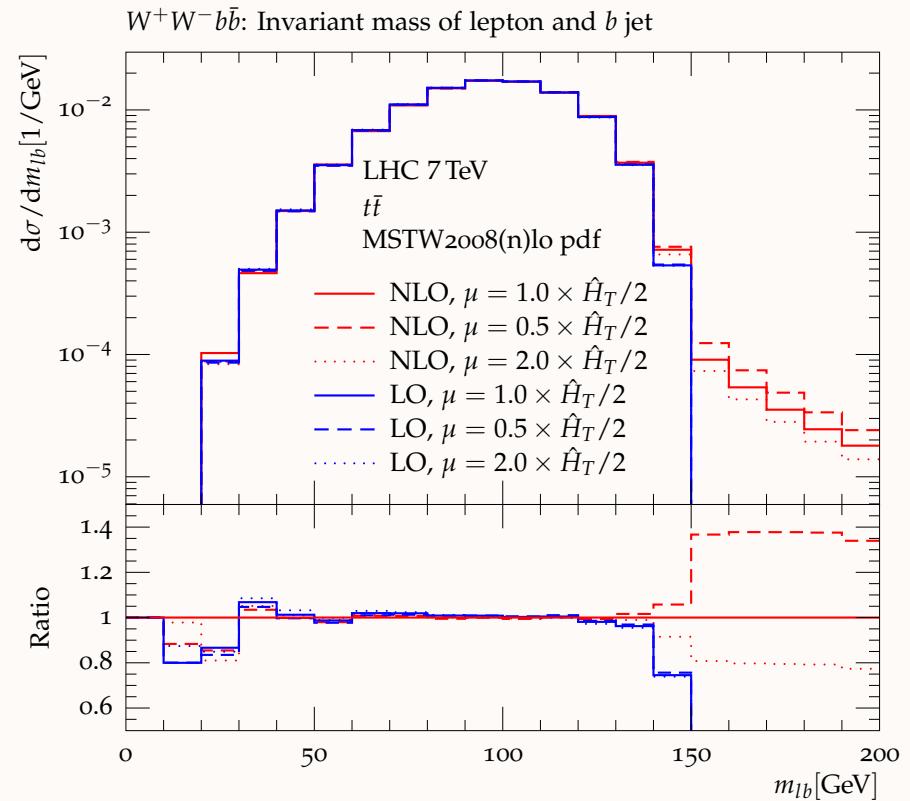
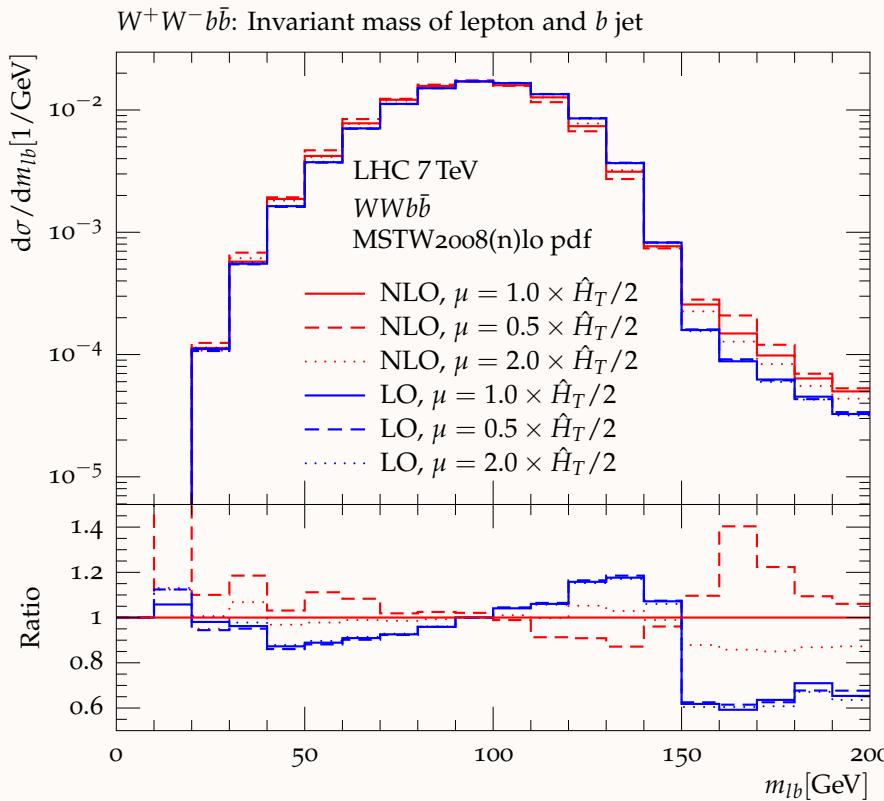
→ $WWb\bar{b}$: NLO corrections strongly affect the shape of m_{lb}

- similar features → agreement on qualitative level only, noting the differences however:
- different LHC energies & kinematical constraints (cuts), slightly different observable (a truth m_{lb})
- different dynamical scale choice (transverse mass of tops)
- non-resonant and off-shell effects due to finite W boson width
- different treatment of b-quark initial states



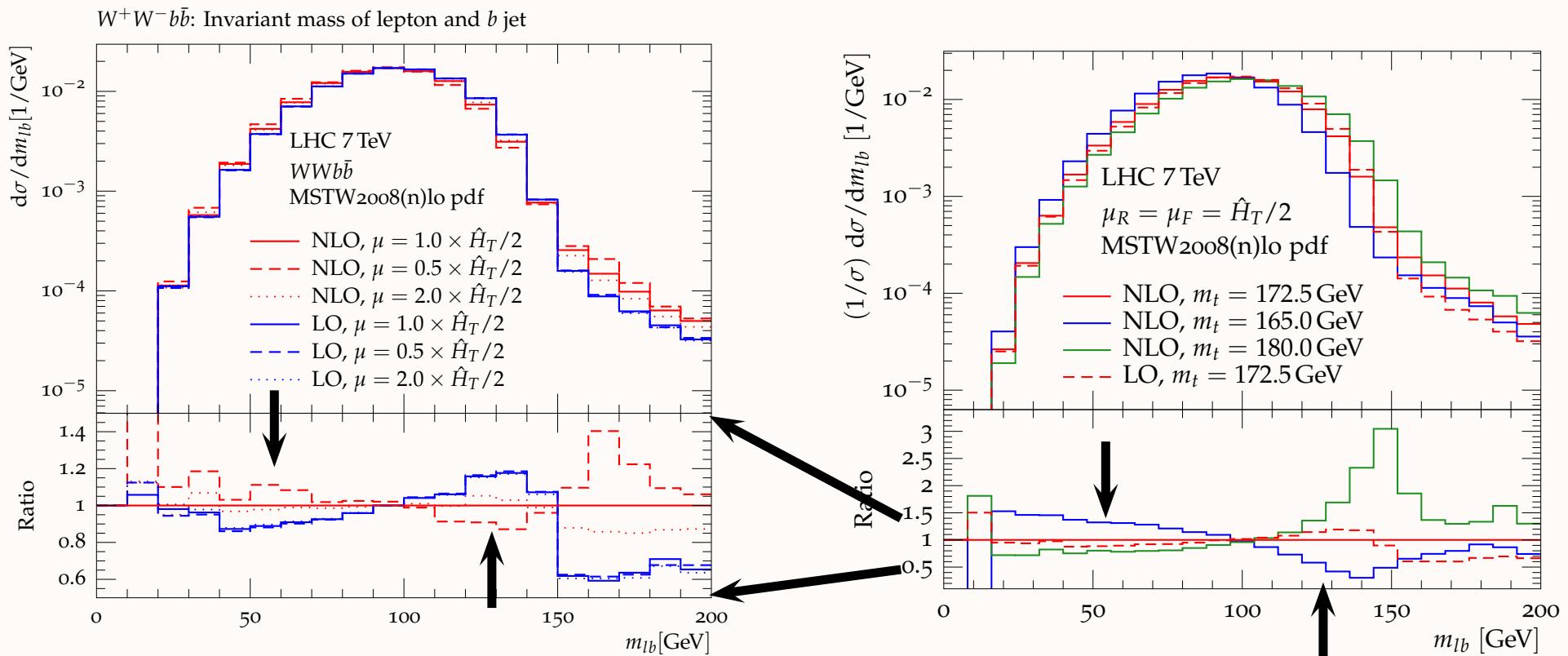
Normalized m_{lb} : scale versus m_t variation

- shape modifications resulting from variation of scales by factors of two
- left panel, for the full approach → visible • right panel, for the factorized approach → only in tails



Normalized m_{lb} : scale versus m_t variation

- shape modifications resulting from variation of scales by factors of two
- left panel, for the full approach → visible • right panel, shape changes due to m_t variation @ NLO
(m_t variation @ LO very similar!)

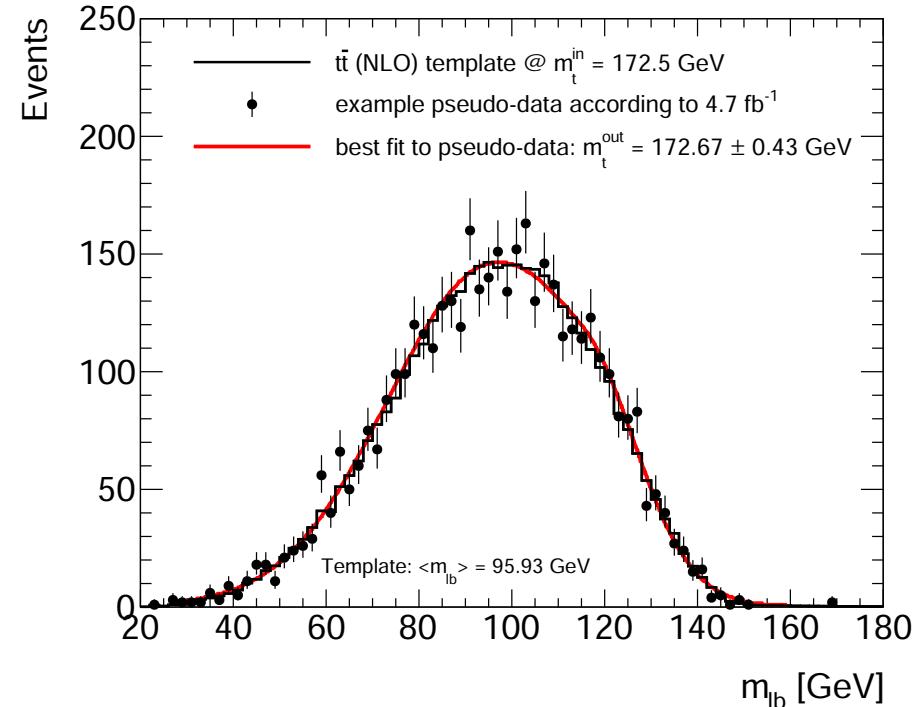
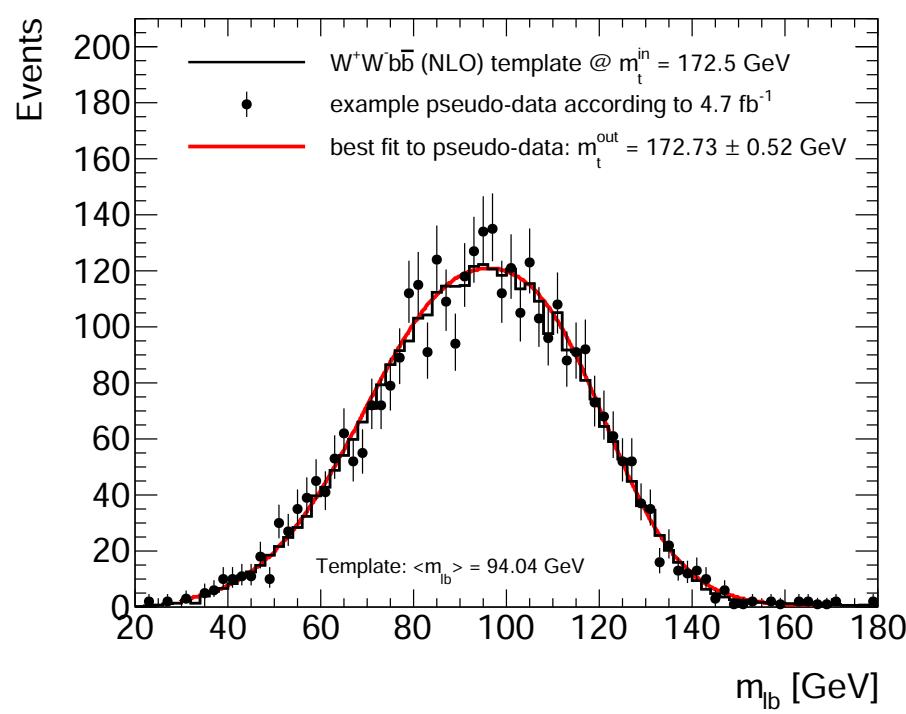
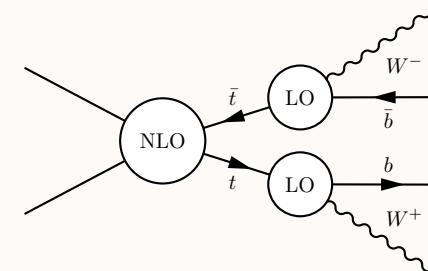
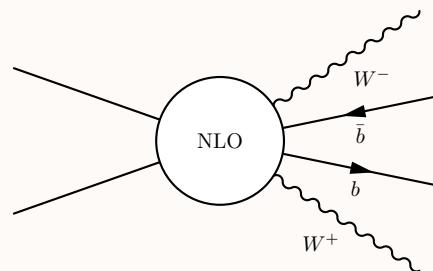


- scale factor variation mimics shape changes as induced by different m_t values → uncertainty
- @NLO: scale down corresponds to lower mass
- fit mass and scale simultaneously, but would resulting choice work for other distributions (eg. $m_{t\bar{t}}$)?



NLO templates vs pseudo-data

→ Representative examples for full (left) and factorized (right) NLO calculation.

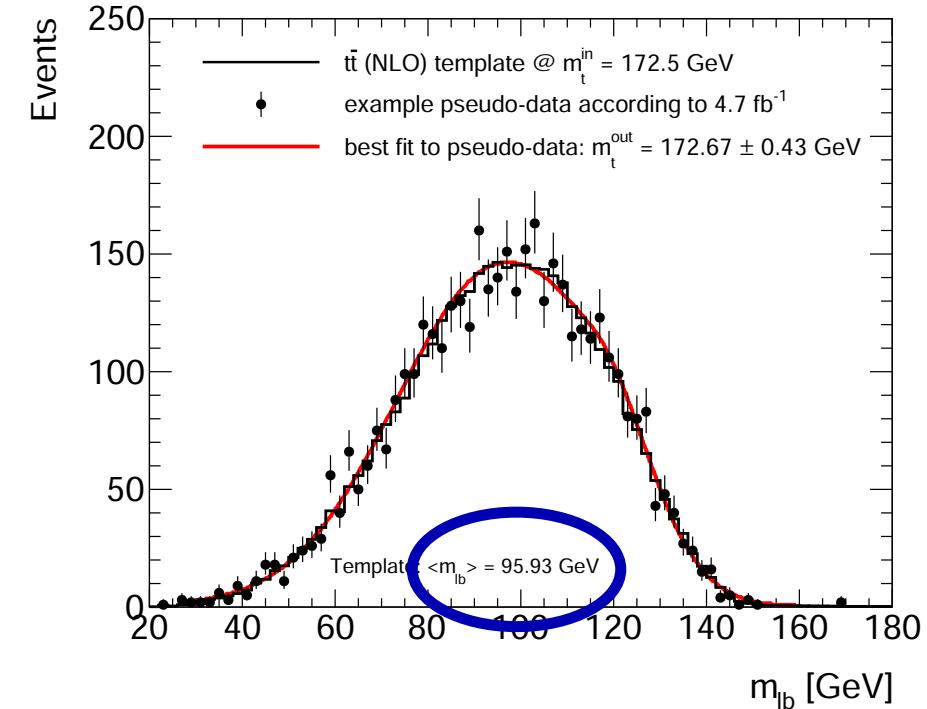
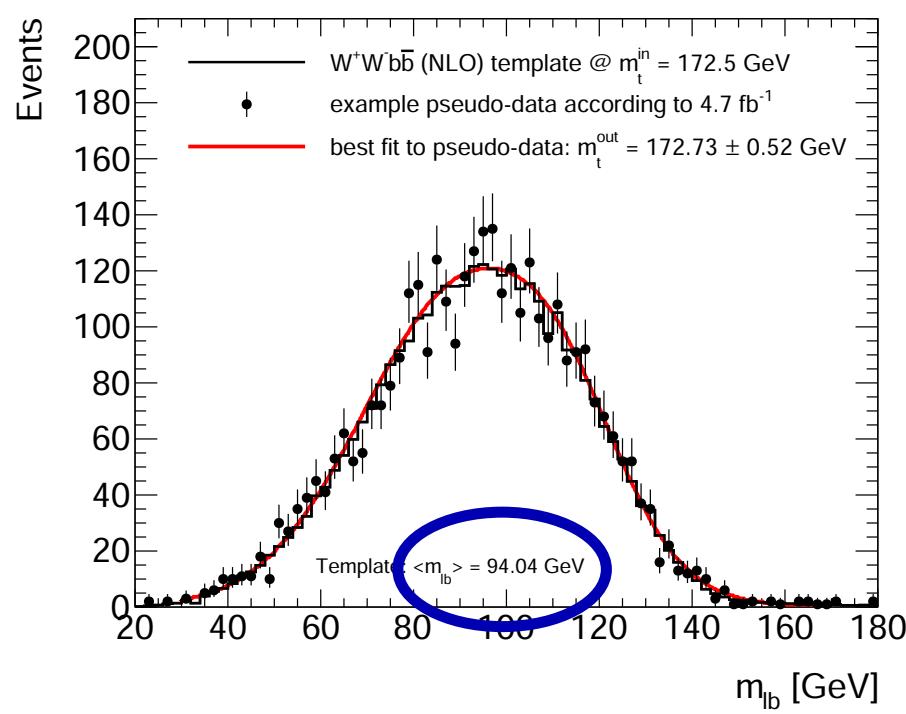
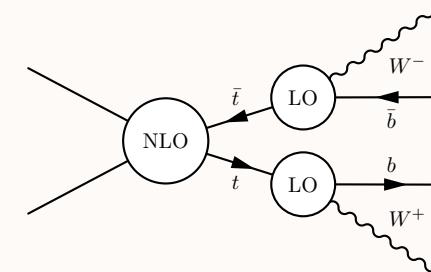
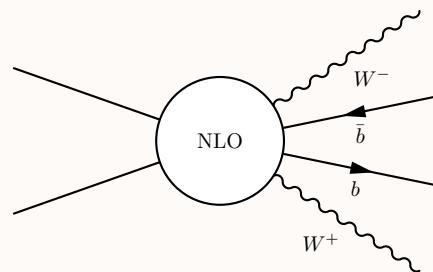


- pseudo-data (black points) are generated from the NLO distributions (black histograms) at m_t^{in}
- fit with NLO templates (parametrization) gives $m_t^{out} \rightarrow$ best fit to pseudo-data (red line)



NLO templates vs pseudo-data

→ Representative examples for full (left) and factorized (right) NLO calculation.



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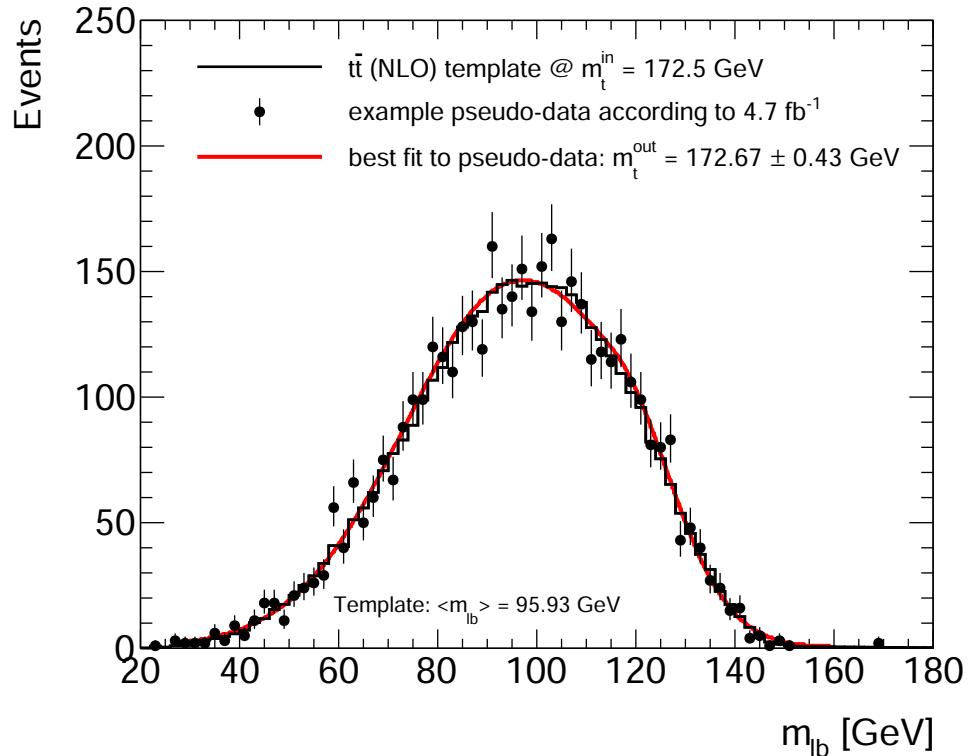


Scale uncertainties and the m_{lb} method

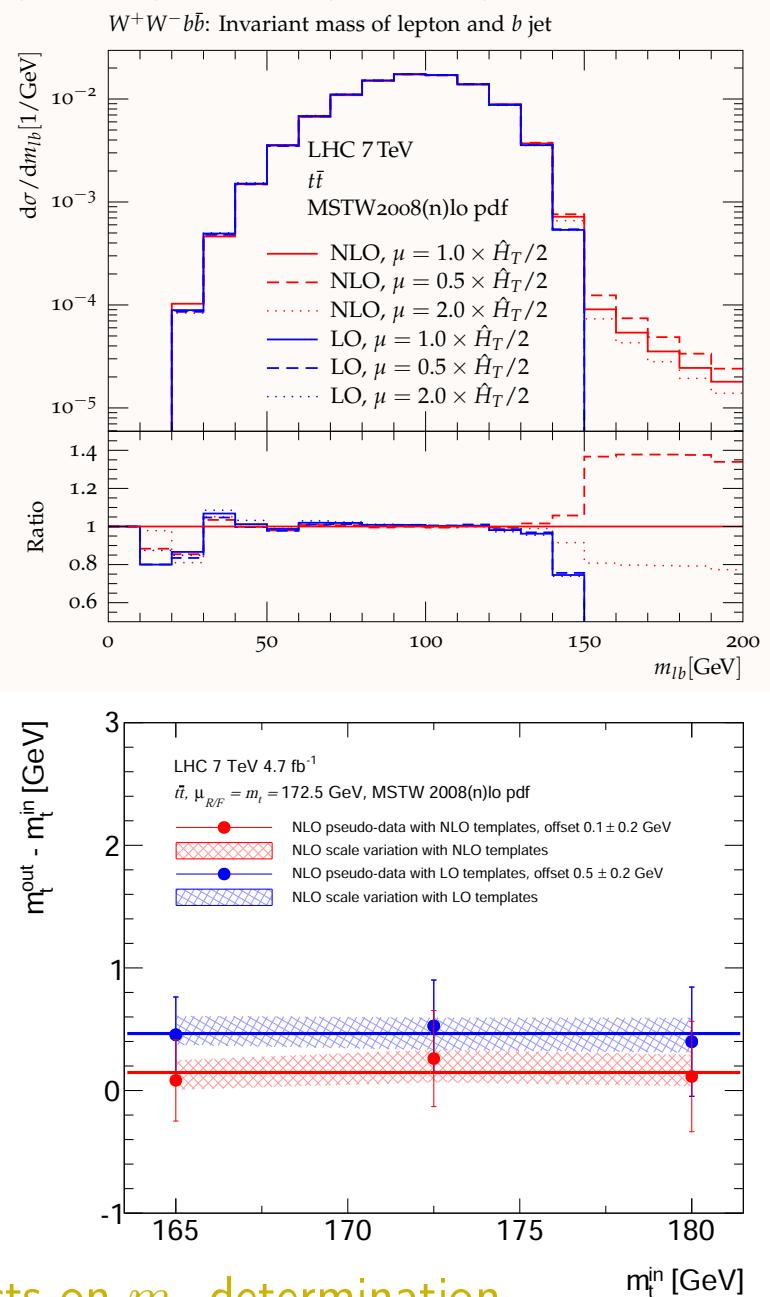
[HEINRICH, MAIER, NISIUS, SCHLENK, WINTER, ARXIV:1312.6659]

Single out effect of NLO scale uncertainties on top mass.

- Use m_{lb} method in a parton-level analysis where we assume that data follows factorized QCD NLO prediction for $t\bar{t}$ with subsequent dilepton decays at LO [pseudo-data].
- Apply/test against the theories given by default scale choice **NLO** and **LO** predictions (templates) [hypotheses].



Based on NLO $t\bar{t} + \text{LO}$ decays, only small effects on m_t determination.

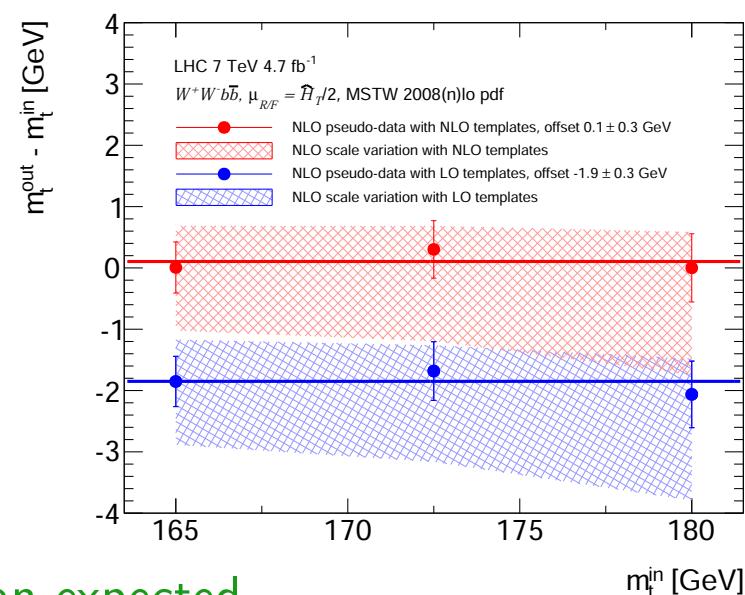
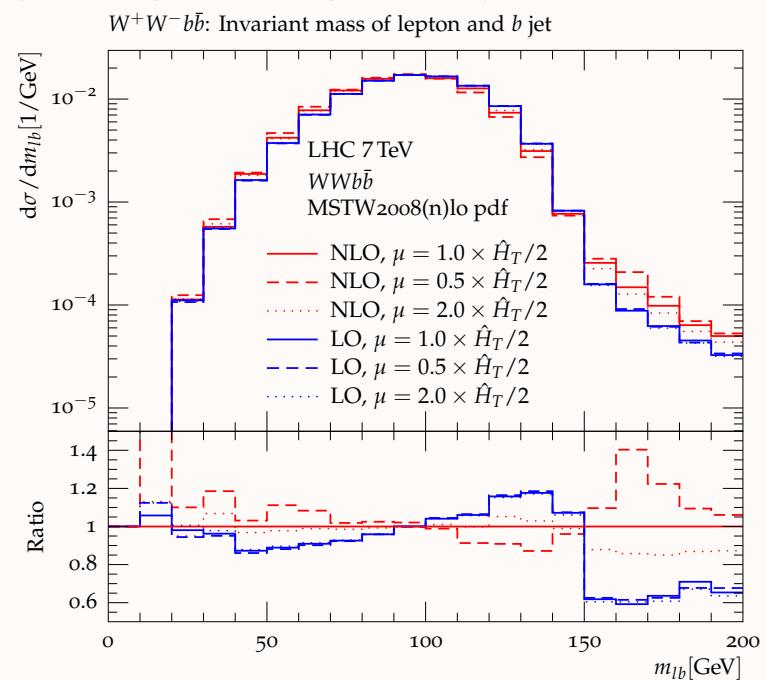
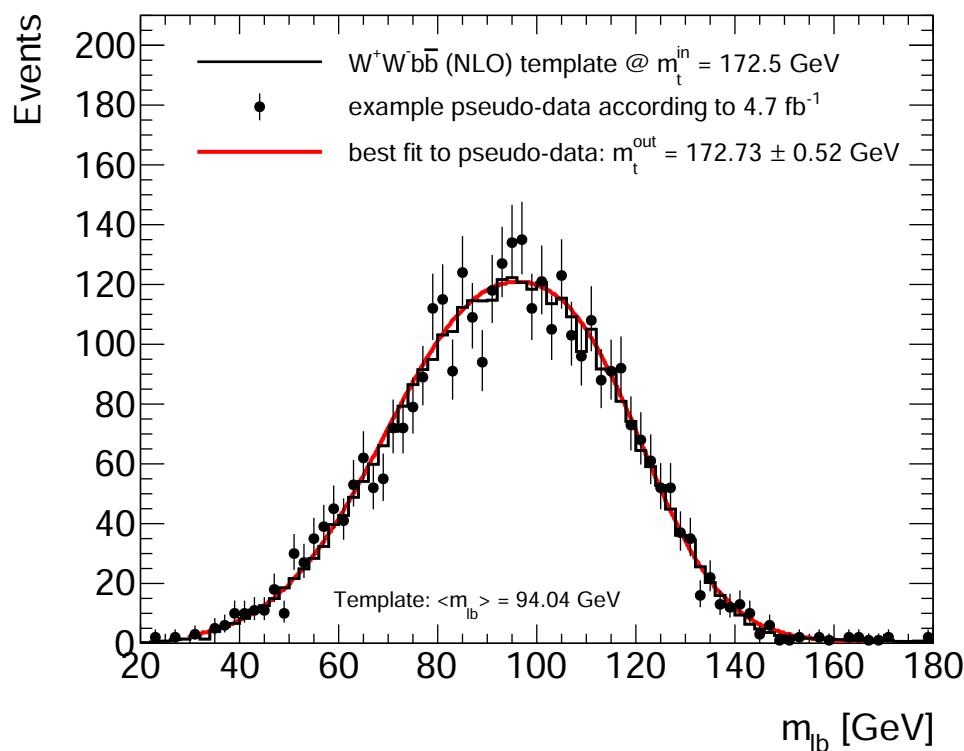


Scale uncertainties and the m_{lb} method

[HEINRICH, MAIER, NISIUS, SCHLENK, WINTER, ARXIV:1312.6659]

Single out effect of NLO scale uncertainties on top mass.

- Use m_{lb} method in a parton-level analysis where we assume that data follows full QCD NLO prediction for dileptonic $W^+W^-b\bar{b}$ [pseudo-data].
- Apply/test against the theories given by default scale choice **NLO** and **LO** predictions (templates) [hypotheses].

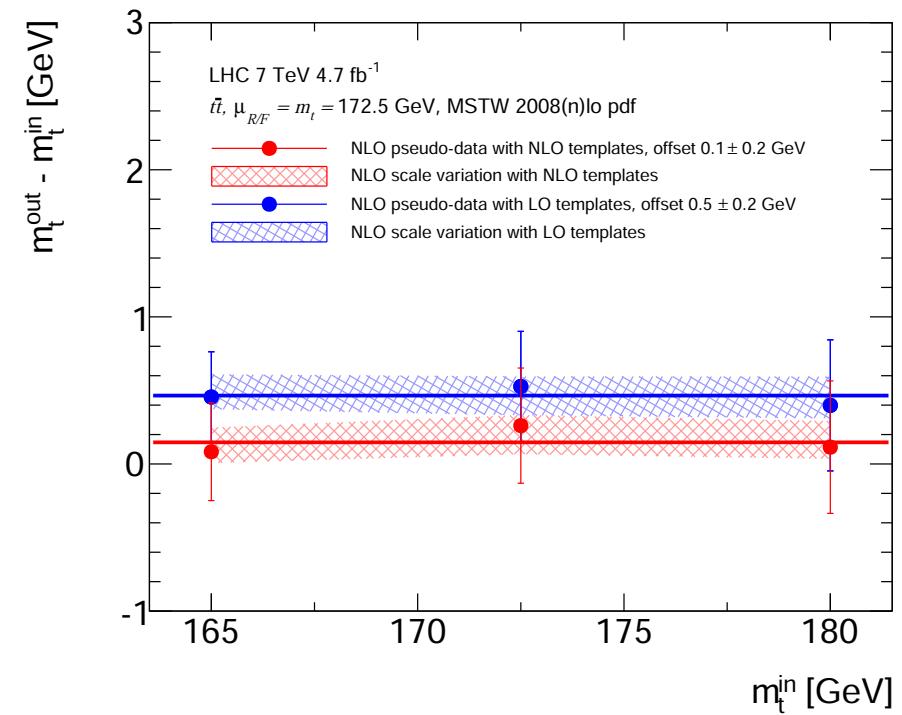
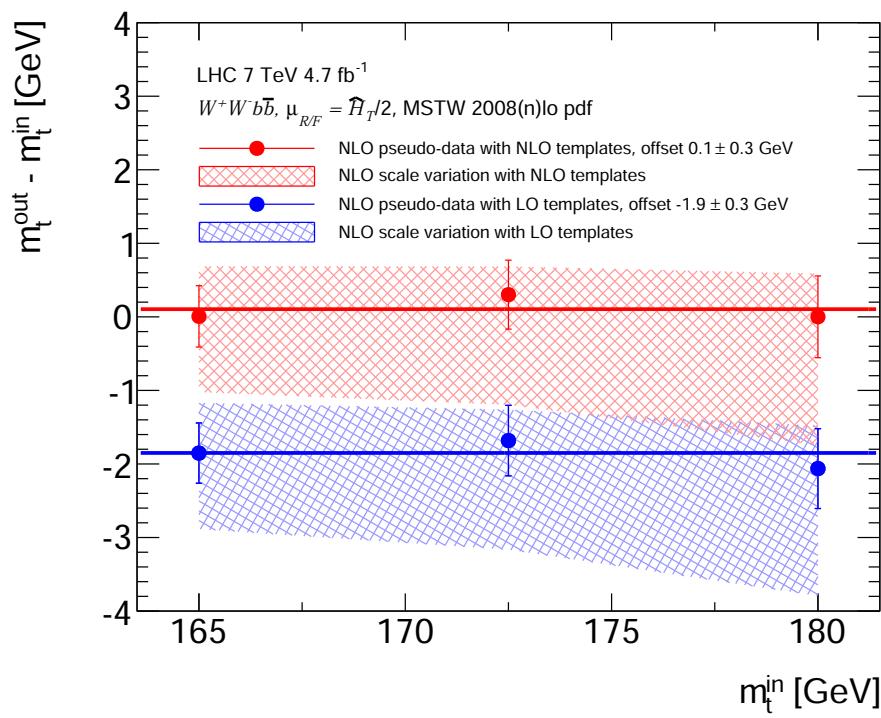
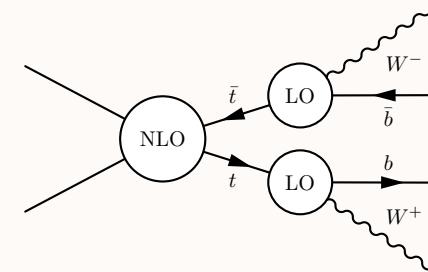
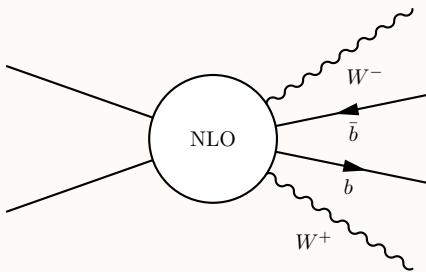


Impacts top quark mass determination more than expected.



Summary

→ Full (left) vs factorized (right) NLO calculation: results for mass shifts.



- larger shift btwn NLO & LO description (~ 1.9 GeV) as compared to factorized approach (~ 0.5 GeV)
- significantly larger uncertainties from scale variations for full approach ($^{+0.6}_{-1.0}$ GeV vs ± 0.2 GeV)



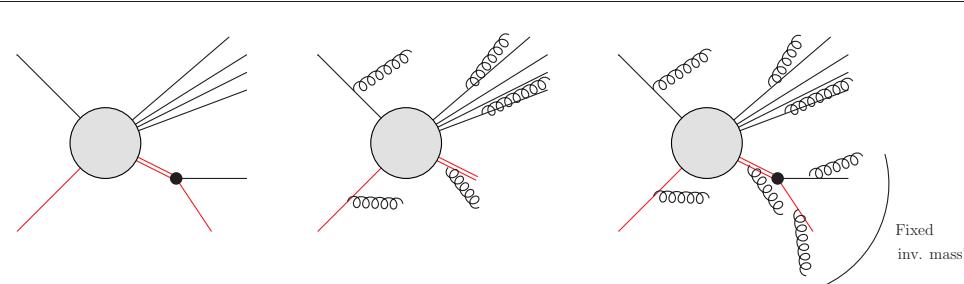
Further remarks.

What about variables other than m_{lb} ?

- recent study concerning the prospects to determine m_t from leptonic observables (dilepton channel)
[FRIXIONE, MITOV, ARXIV:1407.2763]
- proposal to exploit the top quark mass dependence of the shape of the $t\bar{t}+1$ -jet invariant mass
[ALIOLI, FERNANDEZ, FUSTER, IRLES, MOCH, UWER, VOS, ARXIV:1303.6415]
- ...

What about NLO+PS matching for $WWb\bar{b}$?

- ... to obtain more realistic, i.e. hadron level final states.
- first attempt and results using PowHel
[GARZELLI, KARDOS, TROCSANYI, ARXIV:1405.5859]
- however, the issue of intermediate resonances has not been addressed. This is an open issue to the MC community.
(Without a proper treatment of intermediate resonances, parton shower effects will distort the (NLO-accurate) Breit–Wigner shape.)



Summary.

Cutting edge parton-level calculations of NLO QCD corrections to $WWb\bar{b}$ production are available, using modern NLO tools (MG5_aMCNLO, OpenLoops+Sherpa, GoSam+Sherpa, Helac-NLO/PowHel).
Realistic, many body final states!!

Comparison with NWA approaches & standard Monte Carlos helps disentangle effects beyond the factorization and assess their relevance for phenomenology (on the inclusive level approximations work well, ...).

The 4-flavour scheme calculations (treating b-quarks as massive partons) give us new insight to the validation of “top-induced” backgrounds.

NLO effects were also studied in the context of the top quark mass measurement based on the m_{lb} template method (well defined framework also for a pure parton level analysis).

Shape uncertainties from scale variations of the full NLO QCD corrections to $WWb\bar{b}$ production result in larger theory errors on this top quark mass determination than expected.

Validation ongoing ... e.g. to separate effects from radiative corrections in decay and finite-width contributions (NLO in the “decay” seems crucial).





Thank You.

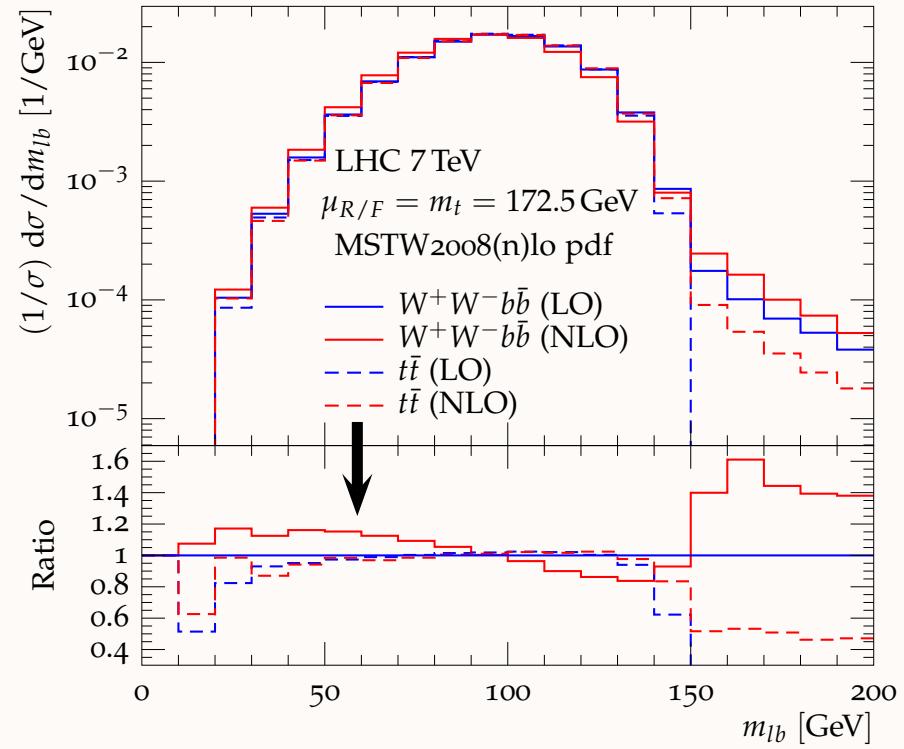
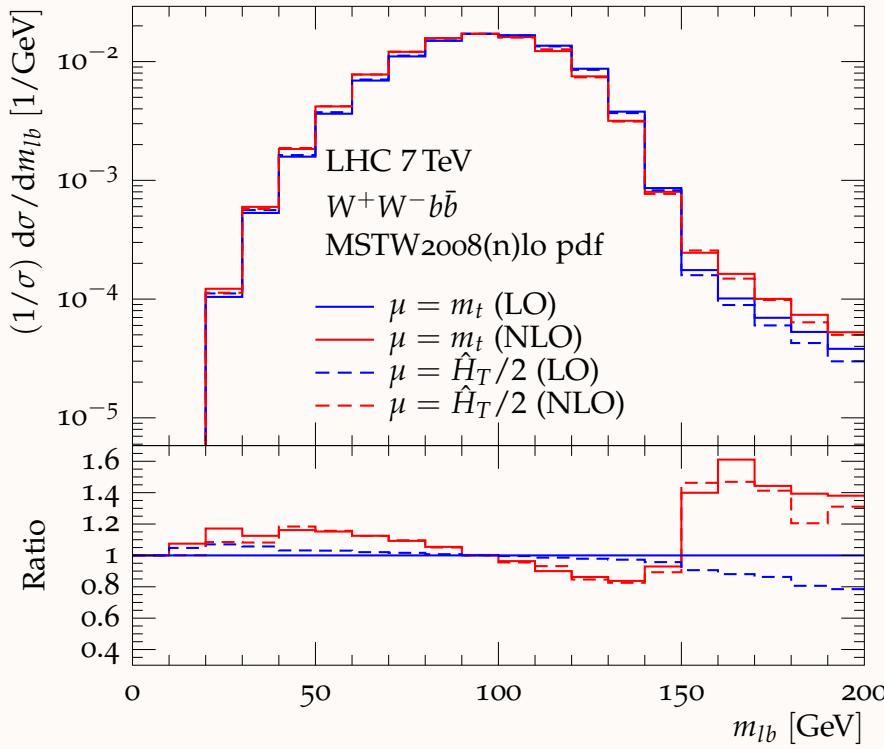


Extra notes.



Normalized m_{lb} : shape comparisons & cross-checks

- analysis strongly driven by shape of the distribution (rate comes in only through number of events passing acceptance/analysis cuts)

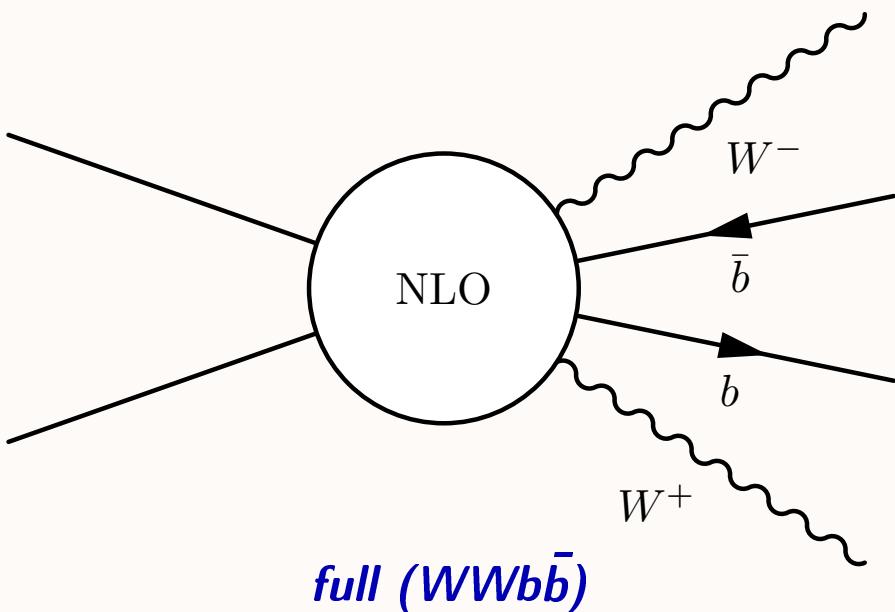


- left: small effect of different scale choices on normalized distributions
- right: for full NLO WW $b\bar{b}$, shape change is drastic while shapes of others are similar

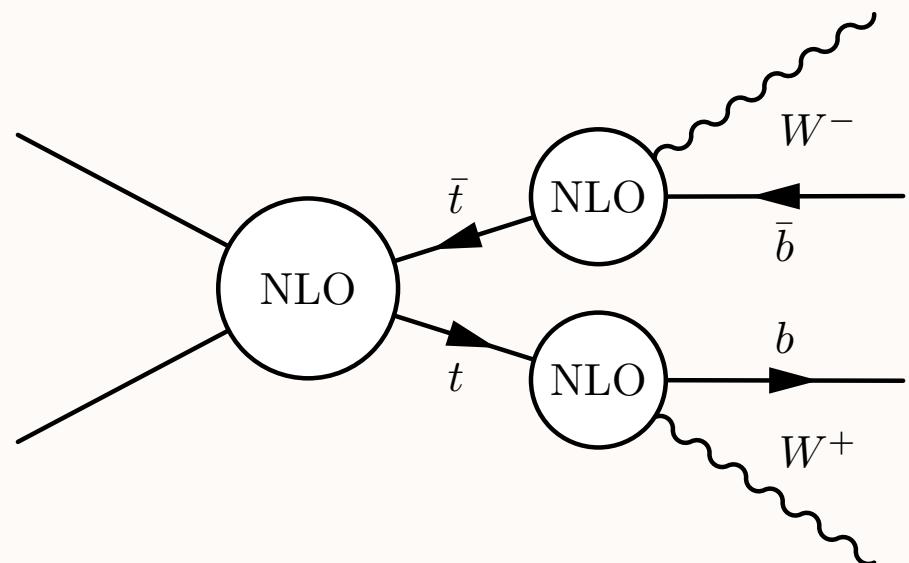


Radiation in the decay

[BISWAS, MELNIKOV, SCHULZE, ARXIV:1006.0910]



full ($WW\bar{b}b$)



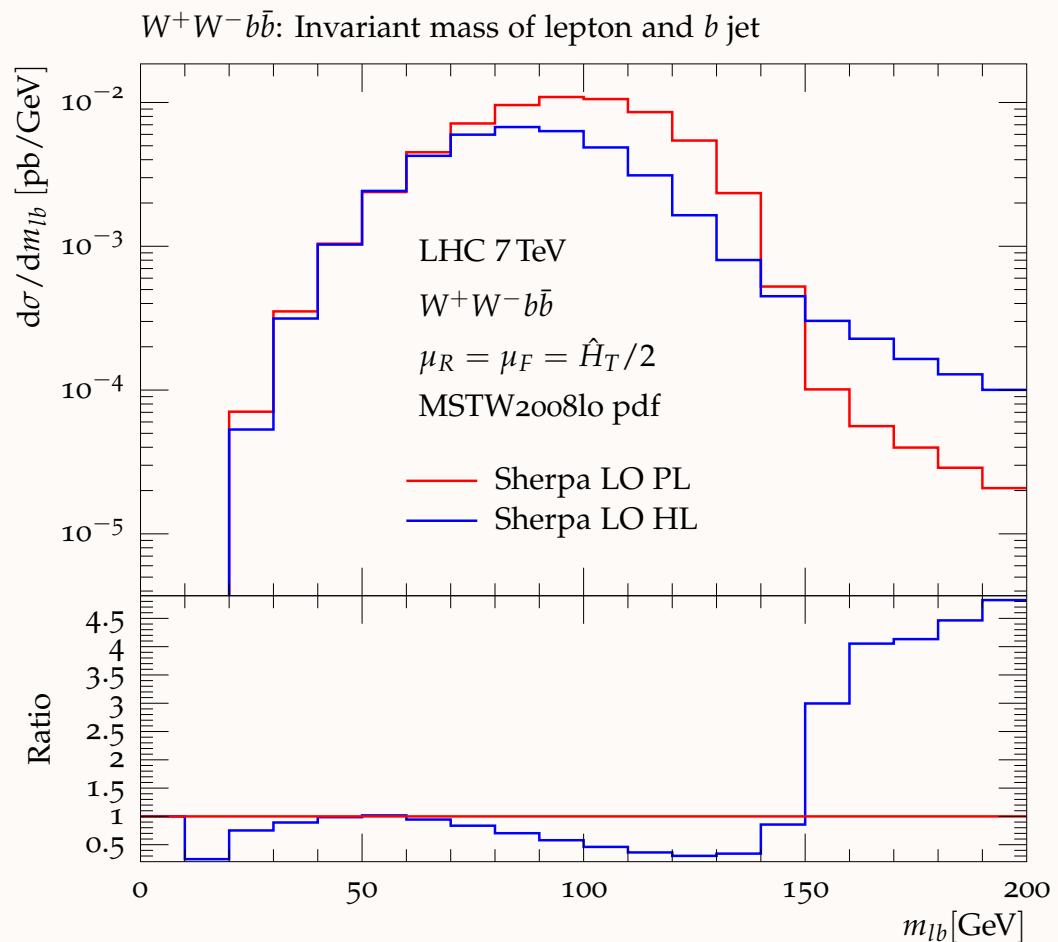
narrow width approx. (NWA)

- full NLO description of the $WW\bar{b}b$ final state ($2 \rightarrow 4$ processes)
- non-resonant/-factorizing contributions, NLO effects in top quark decays
- comparison will help disentangle effects from NLO decays and non-resonant/-factorizing contributions
- choose different scales in the production and decay ... becomes testable
- to what extent are radiative decay corrections well modelled by shower in NLO+PS (how do we assess the uncertainties related to these shower emissions)
- also in [reference](#), an estimator related to m_t obtained from moment of $m_{lb} \rightarrow$ experimental alternative?



What to expect from showering and hadronization?

- Only PRELIMINARY result / illustration. Observation on a qualitative level:
- transition region between peak and tail gets washed out, mainly by parton showering (PS)
- also peak position shifts to the left
- further amplified by hadronization (HAD)
- PS + HAD reduce the top mass sensitivity of m_{lb}
(cannot be neglected but is not dramatic either)

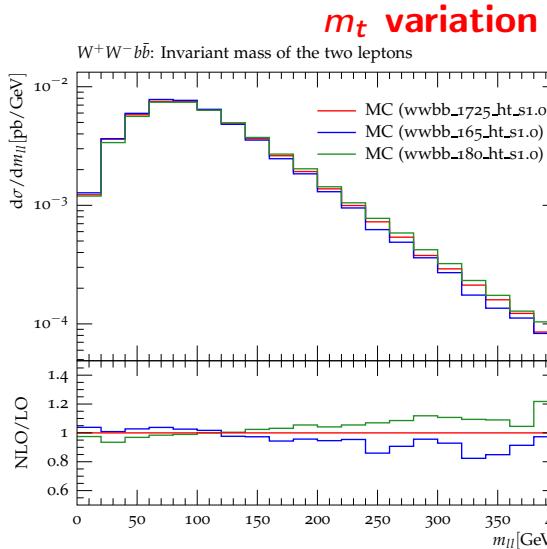
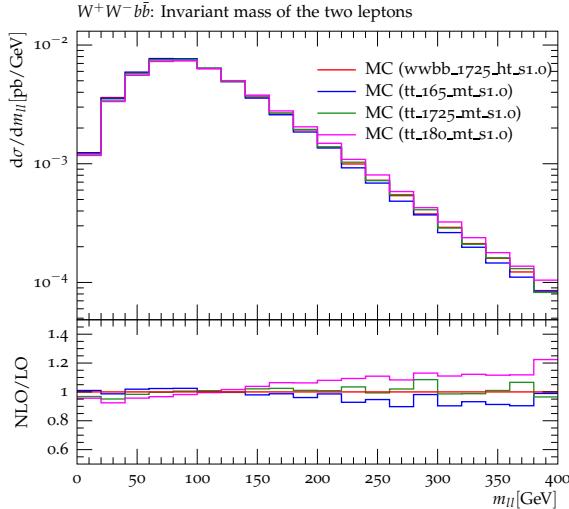


Other observables?

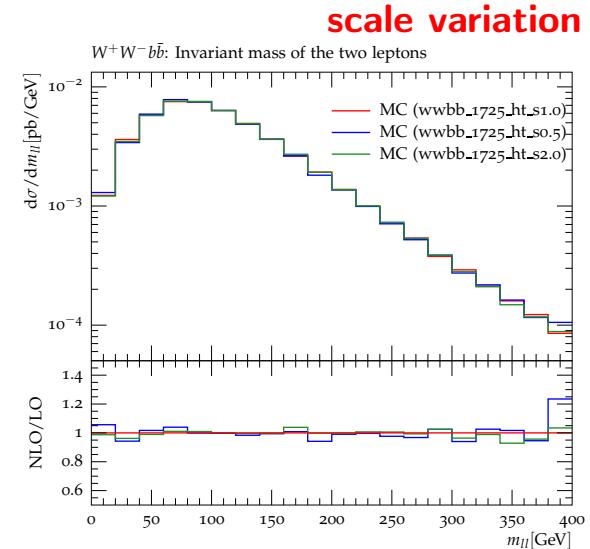
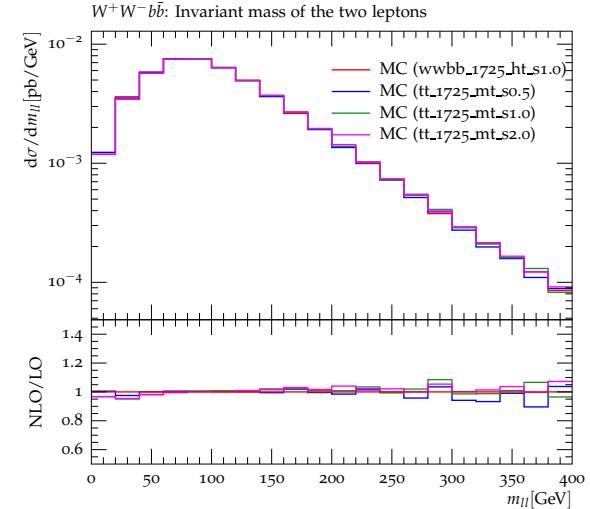
- less sensitive to m_t , but “cleaner” observable → better systematics!?
- pay-off comes with more data ...

m_{ll}

- m_{ll} is an observable almost free of PS + HAD corrections



fact.



full

