Experimental study of the time dependent decay rate in $B_s \rightarrow \phi \gamma$

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<u>Outline</u>

- Introduction
- Theoretical aspects
- Experimental issues
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- Discussion

Introduction

• We want to measure the photon polarization in $B_s \rightarrow \phi \gamma$ decays

 \rightarrow Photons in **b** $\rightarrow s\gamma$ are predicted to be left-handed in the SM (small corrections of order m_s/m_b)

→ Sensitive to New Physics models, particularly Left-Right Symmetric Models

[D. Atwood, M. Gronau and A. Soni, PRL79(97)185][M. Gronau, D. Pirjol, PRD66(02)054008][F. Yu, E. Kou, C. Lü, JHEP12(2013)102]



H⁻, χ⁻,ğ, χ⁰...

Remind that LHCb has measured $\mathcal{B}(B \rightarrow K^* \gamma) / \mathcal{B}(B_s \rightarrow \phi \gamma)$ with 1fb⁻¹ = 1.23±0.06±0.04±0.10(f_s/f_d) [Nuc. Phys. B 867 (2013) 1-18]

ℬ(B_s→φγ) =(3.5±0.4)x10⁻⁵

And has observed polarized photons in $B \rightarrow K \pi \pi \gamma$ decays!

[Phys. Rev. 112(2014)161801]

Theoretical aspects

$$\Delta \Gamma_{\rm s} = \Gamma_{\rm L} - \Gamma_{\rm H} =$$
 (0.081 ± 0.011) ps⁻¹
 $\Gamma_{\rm s} = 1/\tau_{\rm Bs} =$ (0.6596 ± 0.0046) ps⁻¹

Theoretical aspects

 \rightarrow Untagged measurement of the time dependent B_s $\rightarrow \phi \gamma$ width:

$$\begin{split} \Gamma_{\rm B^0_s}({\rm t}) &= |{\rm A}|^2 \, {\rm e}^{-\Gamma_{\rm s} {\rm t}} \left(\cosh \frac{\Delta \Gamma_{\rm s} {\rm t}}{2} - \mathcal{A}^{\Delta} {\rm sinh} \frac{\Delta \Gamma_{\rm s} {\rm t}}{2} \right) \\ &\approx |A|^2 e^{-\Gamma_{B_s \to \phi \gamma} t} \quad \text{with} \end{split}$$

It can be seen as an "Effective lifetime" depending on the A^{Δ}

$$\Gamma_{B_s \to \phi\gamma} = \Gamma_s + \frac{\mathcal{A}^\Delta \Delta \Gamma}{2}$$

SM value: $A^{\Delta} = 0.047 \pm 0.025 \pm 0.015_{(\alpha_s)}$ [Muheim, Xie, Zwicky, PLB664(08)174] Left-Right Symmetric models: A^{Δ} up to ~ 0.7 [Atwood, Gronau and Soni, PRL79(97)185]

→ Fraction of wrongly polarized photons ~ 40%

Theoretical aspects

 \rightarrow The time-dependent decay assymetry is then expressed:

$$\mathcal{A}_{\rm CP}(B_s \to \phi\gamma)[t] \equiv \frac{\Gamma[t][\bar{B}_s \to \phi\gamma] - \Gamma[t][B_s \to \phi\gamma]}{\Gamma[t][\bar{B}_s \to \phi\gamma] + \Gamma[t][B_s \to \phi\gamma]}$$
$$= \frac{S\sin(\Delta m_s t) - C\cos(\Delta m_s t)}{\cosh(\frac{\Delta\Gamma_s}{2}t) - \mathcal{A}} \sinh(\frac{\Delta\Gamma_s}{2}t)$$

- $\rightarrow \mathbf{A_{CP}}$ expected to be **~zero** in the Standard Model
- \rightarrow Enhanced in New Physics Models
- \rightarrow Caveat: experimentally one needs to know the flavour of the B_s

 \rightarrow Flavour tagging for B_s drastically reduces our data, even if we have (and will have more!) a lot of B_s $\rightarrow \phi \gamma$ decays

 $\sigma(pp \rightarrow B_s + X) = 10.5 \pm 1.3 \ \mu b \ [JHEP08(2013)11]$

 $\mathcal{B}(B_s \rightarrow \phi \gamma) = (3.5 \pm 0.4) \times 10^{-5}$ $\varepsilon_{\text{reconstruccion}} (\phi + \gamma) \sim 1\%$

Tagging algorithms:

Same side (SS):

From fragmentation of the signal *b* (π for B, K for B_s)

Opposite side (OS):

From the opposite B:

- e, μ from semileptonic B decays,
- kaons from $b \rightarrow c \rightarrow s$,
- inclusive reconstruction of the opposite B vertex

Tagging efficicencies: $\varepsilon \sim 5-16\%$, mistag probabilities $\omega \sim 30\% \rightarrow N_{evts} \times \varepsilon (1-2\omega)^2$

 \rightarrow Let's first to measure the untagged rate...



[Eur. Phys. J.C. 72 (2012) 2022, LHCb-CONF-2012-026]

$$\Gamma_{\mathrm{B}^{0}_{\mathrm{s}}}(\mathrm{t}) \approx |A|^{2} e^{-\Gamma_{B_{s} \to \phi \gamma} t}$$



In practice we need to correct the measured B_s proper time distribution for acceptance and resolution effects:

$$\Gamma_{Bs}(t_r)$$
 measured = $A(t) \cdot \Gamma_{Bs}(t; A^{\Delta}) \otimes R(t, t_r)$



• Two different data control channels are being used to extract the acceptance:

Using the $B_d \rightarrow K^* \gamma$ as control channel: - Same topology (2 tracks + photon) - Similar trigger (γ) - Affected by tracking differences between the $K^* \rightarrow K\pi$ and $\phi \rightarrow KK$ PhiGamma Data (2 fb -1) 12300 evts 2fb⁻¹ $B_d \rightarrow K^* \gamma$ 300 200 5200 5300 5 5400 5100 5500 5600 B M [MeV/c2 B mass (MeV) - Large $B \rightarrow K^* \gamma$ to $B_s \rightarrow \varphi \gamma$ ratio ~ 6

- Background contributions

Using the $B_s \rightarrow \phi J/\psi$ as control channel:

- Same tracking ($\varphi \rightarrow$ KK, unbias J/ $\psi)$
- Different trigger (dimuon)



Sensitive to different (complementary) trigger, detector and systematics effects

1) Using the $B_d \rightarrow K^* \gamma$ as control channel:

Aiming for a selection that keeps the same proper time acceptance for the $B_s \rightarrow \phi \gamma$ and $B_d \rightarrow K^* \gamma$ decay channels



 \rightarrow Same acceptance for B_s and B_d within statistical MC uncertainties

Signal events $(2fb^{-1}) = 3200$; S/B ~ 4 Tighter selection: Signal events $(2fb^{-1}) = 2280$; S/B ~ 7.4; still flat acceptance ratio

Comparing data and MC distributions to validate the selection: In general the LHCb simulation reproduces quite well our data. Some examples with the largest discrepancies, can be corrected:



2) Using the $B_s \rightarrow \phi J/\psi$ as control channel:

Faking the photon with the J/ ψ : use an "unbiased B_s \rightarrow J/ $\psi\phi$ lifetime" (B_s vertices reconstructed only from the two kaon system) and apply the selection as if it was a radiative decay

Use the $B \rightarrow K^*J/\psi$ and $B \rightarrow K^*\gamma$ to assess the method



 \rightarrow Flat ratio of high acceptances within statistical uncertainties (MC)

5

0

10

Some differences in geometry variables introduce differences at low decay time (related to vertex displacement)



Events

•The proper time resolution $R(t,t_r)$ is crucial for an unbiased measurement of the photon polarization parameter A^{Δ}

 $R(t,t_r)=Gauss_{core}(t,t_r)+Gauss_{wide}(t,t_r)$

• Bias in the resolution is dominated by photon related information

Big effort has been done in the proper photon calibration
(depending on the calorimeter region):

 μ_{core} = 3.2 fs μ_{wide} = 11 fs



• Similar resolution functions for the $B_s \rightarrow \phi \gamma$ and $B_d \rightarrow K^* \gamma$ decay channels:

2 Gaussian fit



Trigger + selection:

$$\begin{split} \mathbf{B}_{s} & \rightarrow \boldsymbol{\phi} \gamma \\ \sigma_{core} &= 55 \pm 1 \text{ fs} \\ \mu_{core} &= 3.2 \pm 0.5 \text{ fs} \\ \sigma_{wide} &= 119 \pm 2 \text{ fs} \\ \mu_{wide} &= 11 \pm 1 \text{ fs} \\ \text{Core fraction} &= 57\% \end{split}$$

B_d →K*γ $σ_{core}$ = 49 ± 1 fs $µ_{core}$ = 3.5 ± 0.8 fs $σ_{wide}$ = 103 ± 2 fs $µ_{wide}$ = 6 ± 1 fs Core fraction= 52%

• Resolution functions for the $B_d \rightarrow K^* \gamma$ and $B_d \rightarrow K^* J/\psi$:





Background model (toy simulations):

 \rightarrow Combinatorial τ distributions as function of the mass region.

 \rightarrow Contribution of physical backgrounds

The effect on A^{Δ} in progress \rightarrow final selection

- Combinatorial background (mainly real $\phi + \gamma$)
- Contamination from merged π^0
- Baryonic radiative decays (*b*-baryons into $\Lambda^* + \gamma$).
- Partially reconstructed B decays $(B_s \rightarrow \phi X\gamma, X \text{ non reconstructed})$





 \rightarrow The small value of $\Delta\Gamma_{\rm s}$ (~0.08) makes the measurement difficult...

 \rightarrow But possible ...



 \rightarrow Statistical sensitivity on A^{Δ} (3fb⁻¹) ranging from 0.31 (3200 events) to 0.26 (4500) depending on the selection criteria (compromise for background rejection)

• Different fit strategies:

$$\Gamma_{B^0_s}(t) = |A|^2 e^{-\Gamma_s t} \left(\cosh \frac{\Delta \Gamma_s t}{2} - \mathcal{A}^\Delta \sinh \frac{\Delta \Gamma_s t}{2} \right)$$

- Binned L fit
- Unbinned ML fit
- ► Fit to the ratio of $B_s \rightarrow \phi \gamma / B_d \rightarrow K^* \gamma$: →

$$\frac{d\Gamma_{B_s}/dt}{d\Gamma_{B_d}/dt} = \frac{N_{B_s}}{N_{B_d}} \frac{\widehat{\varepsilon}_{B_s}(t)}{\widehat{\varepsilon}_{B_d}(t)} \Big(\frac{\Gamma_{B_s} + \frac{\mathcal{A}^{\Delta}\Delta\Gamma_s}{2}}{\Gamma_{B_d}} \Big) \frac{e^{-\left(\Gamma_{B_s} + \frac{\mathcal{A}^{\Delta}\Delta\Gamma_s}{2}\right)t} \otimes \mathcal{R}_s(t,t')}{e^{-\Gamma_{B_d}t} \otimes \mathcal{R}_d(t,t')}$$

∆Г_d <<

 \rightarrow At present several fitters have been implemented: running toys to validate the strategy and assess the uncertainties

 \rightarrow Some examples of toy simulations (binned fit):

- Using adaptive binning to ensure a minimum of entries per bin (100-200)

- 3000 events for $\rm B_{s}$, 20000 for $\rm B_{d}$



- Error budget (preliminary studies) :
 - Statistical sensitivity to A^{Δ} (3000 events) ~ 0.3
 - Statistical bias ~ 0.005
 - Fit range and bias on A^{Δ} (finite τ) < 0.01
 - Effect of binning < 0.01
 - Effect of acceptance: KEY one... in progress...
 - Effect of resolution < 0.08
 - Effect of external measurements: ~ 0.1

 $(\Delta \Gamma_{\rm s} = (0.081 \pm 0.011) \text{ ps}^{-1}$, $\Gamma_{\rm s} = (0.6596 \pm 0.0046) \text{ ps}^{-1})$

- Effect of background: (first evaluations → small)

...

Conclusions

• The photon polarization measurement in $B_s \rightarrow \phi \gamma$ decays is possible at present with untagged events at LHCb (at least **3000 signal events**)

- ightarrow measurement of the photon polarization parameter \mathbf{A}^{Δ}
- \rightarrow expected statistical sensitivity $\sigma_{{\mbox{\tiny A}}^{\Delta}}{\mbox{}^{\sim}}$ 0.3
- \rightarrow evaluation of systematics in progress \rightarrow below statistical uncertainty (use of B \rightarrow K* γ and B_s \rightarrow J/ $\psi \phi$ events from data as control channels)
- We will have very soon more data... 3 fb⁻¹ + 5 fb⁻¹ (at 14 TeV, $\sigma_{pp \rightarrow bb} x$ 2): \rightarrow untagged measurement: ~ **15K events** $\Rightarrow \sigma_{A^{\Delta}} \sim 0.15$
- Sensivity studies with tagged events in progress: time dependent decay rate, A_{CP} (3 fb⁻¹ (Run 1) ~ 150 events; Run1 + Run 2 ~ 500 evts)

Some naïve questions...

• Left and right amplitudes, A_L and A_{R_1} are related to the helicity amplitudes:

$$\begin{array}{l} \mathsf{A}_{\mathsf{L}} = \mathbf{1}/\sqrt{4} \; \mathsf{H}_{\mathsf{++}} \\ \mathsf{A}_{\mathsf{R}} = \mathbf{1}/\sqrt{4} \; \mathsf{H}_{\mathsf{--}} \\ \end{array} \qquad \tan \psi \equiv \left| \frac{\mathcal{A} \left(\bar{\mathsf{B}}_{\mathsf{s}} \rightarrow \phi \, \gamma_{\mathsf{R}} \right)}{\mathcal{A} \left(\bar{\mathsf{B}}_{\mathsf{s}} \rightarrow \phi \, \gamma_{\mathsf{L}} \right)} \right| \quad \sim \quad \frac{\mathsf{H}_{\mathsf{--}}}{\mathsf{H}_{\mathsf{++}}} \\ \mathcal{A}^{\Delta} \quad \approx \quad \sin 2\psi \cos \varphi_{\mathsf{s}} \end{array}$$

- The ϕ is supposed to have the same polarization as the photon
- have the kaons from the $\boldsymbol{\varphi}$ specific angular distributions depending on the model?
- are we sensitive at LHCb?
- are we sensitive in other channels ($\rho\pi$)?
- With the same argument, are we biasing our measurement if we have some angular/helicity cut?
- We assume that the weak phases $\psi_{\text{L}},\psi_{\text{R}}\sim$ 0. What does it happen in other models?

Thanks!