Status of the quest for $B_{d,s} \to \mu^+ \mu^-$ decays





Outline

- Introduction to $B_{d,s} \to \mu^+ \mu^-$ decays
- Search for $B_{d,s} \to \mu^+ \mu^-$ at LHCb
- Combination with the CMS experiment
- Conclusions



$B_s \to \mu^+ \mu^-$ and $B^0 \to \mu^+ \mu^-$

- Flavour changing neutral currents, helicity suppressed
- Probably the cleanest rare decay both experimentally and theoretically
- Sensitive to pseudoscalar and scalar couplings
- Almost any new physics model could predict a contribution to the branching fractions
- Ratio of branching fraction is probe of minimal flavour violation models





Original figure from D. Straub - Nuovo Cim. C035N1 (2012) 249-256

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A word on the SM branching fractions

The branching fraction of $B_s \to \mu^+ \mu^-$ can be written in the SM as:

$$\mathcal{B}^{t=0}(B_s \to \mu^+ \mu^-) = \frac{G_F^4 M_W^4}{\pi^2} \tau_{B_s} f_{B_s}^2 m_{B_s} m_\mu^2 \sqrt{1 - \frac{4m_\mu^2}{m_{B_s}^2}} |V_{tb} V_{ts}^*|^2 |C_{10}|^2$$

To compare with experimental values, which are time integrated, a correction due to the finite $B_s^0 - \bar{B}_s^0$ width difference has to be applied [De Bruyn et al. [PRL 109, 041801]] [Phys.Rev.Lett.108 (2012) 101803]

$$\mathcal{B}^{\langle t \rangle}(B_s \to \mu^+ \mu^-) = \left(\frac{1 + \mathcal{A}_{\Delta\Gamma} y_s}{1 - y_s^2}\right) \times \mathcal{B}^{t=0}$$

which also introduces a model dependency. Most recent predictions (time integrated)

$$\begin{aligned} \mathcal{B}(B_s \to \mu^+ \mu^-)^{\langle t \rangle} &= (3.65 \pm 0.23) \cdot 10^{-9} \\ \mathcal{B}(B^0 \to \mu^+ \mu^-)^{\langle t \rangle} &= (1.06 \pm 0.09) \cdot 10^{-10} \end{aligned}$$

[Bobeth et al. PRL 112 (2014) 101801.] full electroweak two loop corrections and three loop QCD corrections [Bobeth et al. [1311.1348]] [Hermann et al. [1311.1347]] Main uncertainties:

	f_{B_q}	СКМ	$ au_{H}^{q}$	M_t	α_s	Other parameters	Nonparametric	\sum
$\overline{\mathcal{B}}_{s\ell}$	4.0%	4.3%	1.3%	1.6%	0.1%	< 0.1%	1.5%	6.4%
$\mathcal{B}_{d\ell}$	4.5%	6.9%	0.5%	1.6%	0.1%	< 0.1%	1.5%	8.5%

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Analysis strategy at LHCb

- 1. Loose selection:
 - * Pairs of opposite charged muons
 - Vertex displaced with respect to interaction point
 - * $m_{\mu\mu} \in (4900, 6000)$
 - * $p_{\rm T}$, IP and quality requirements



2. Search in a two dimensional plane of invariant mass and BDT





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Trigger





• Single and dimuon triggers

•
$$(\varepsilon_{B_s \to \mu^+ \mu^-} \simeq 90\%)$$



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How to know where to look? Calibration of signal PDFs: mass Mass central value



[Phys. Rev. Lett. 111, 101805 (2013)] Mass resolution



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Backgrounds



BDT

- Multivariate discriminant with 12 variables
- Trained on MC and calibrated with Data





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BDT: calibration of the PDF

- Signal calibrated with exclusive $B \rightarrow h^+h^-$ decays separated with PID (same method as for the mass PDF calibration but in bins of BDT)
- Background calibrated with full fit to the sidebands (see next slides)





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[Phys. Rev. Lett. 111, 101805 (2013)]

Exclusive backgrounds: $B \to h^+h^-$ and semileptonics

 $B \rightarrow h^+ h^-$:

- Yield fitted on data (with hadronic reconstruction)
- MisID probability measured with control channels in data
- Convoluted with MC spectrum
- Corrected for trigger bias

Semileptonic and rare decays

- Normalised to $B^+ \to J/\psi K^+$
- MisID again from data
- Largest errors from theoretical input for some of them (e.g. $\Lambda_b^0 \to p \mu^- \nu$)



Normalisation

$$\mathcal{B}(B_q^0 \to \mu^+ \mu^-) = \frac{\epsilon_{\rm cc}}{\epsilon_{\rm sig}} \cdot \frac{f_{\rm cc}}{f_q} \cdot \frac{N_{B_q^0 \to \mu^+ \mu^-}}{N_{\rm cc}} \cdot \mathcal{B}_{cc} = \alpha_q \cdot N_{B_q^0 \to \mu^+ \mu^-}$$

- Two control channels: $B^+ \to J/\psi K^+$ and $B^0 \to K^+\pi^-$ Results compatible with each other and averaged for final result
- Reconstruction and selection efficiencies from MC but cross-checked with data
- Trigger efficiency from data ($\varepsilon_{B_s \to \mu^+ \mu^-} \simeq 90\%$)
- Hadronisation fractions from updated LHCb measurement $f_s/f_d = 0.259 \pm 0.015$ [LHCb-CONF-2013-011]
- Correction on efficiencies due to lifetime bias: SM assumed



Normalisation

Normalisation for the full dataset $(3fb^{-1})$

$$\alpha_{B_s \to \mu^+ \mu^-} = (9.01 \pm 0.62) \times 10^{-11} \qquad \alpha_{B^0 \to \mu^+ \mu^-} = (2.40 \pm 0.09) \times 10^{-11}$$

i.e. $40 \pm 4 \ (4.5 \pm 0.4)$ expected $B_s \to \mu^+ \mu^- \ (B^0 \to \mu^+ \mu^-)$ events in the full BDT range





[Phys. Rev. Lett. 111, 101805 (2013)]

Correction of lifetime bias

Integrated efficiency is
$$\epsilon = \frac{\int_0^\infty \Gamma(t)\epsilon(t)dt}{\int_0^\infty \Gamma(t)dt}$$

However the real width is a function of NP parameters
 $\Gamma\left(B_s \to \mu^+\mu^-\right) = \Gamma\left(B_s \to \mu^+\mu^-\right) + \Gamma\left(\bar{B}_s^0 \to \mu^+\mu^-\right)$
 $= R_H e^{-\Gamma_H t} + R_L e^{-\Gamma_L t} = (R_H + R_L) e^{-\Gamma_s t} \left[\cosh\frac{y_s t}{\tau_{B_s}} + \mathcal{A}_{\Delta\Gamma} \sinh\frac{y_s t}{\tau_{B_s}}\right]$

Therefore we have a bias that depends on the MC model used to calculate the efficiency

$$\delta_{\epsilon} = \frac{\epsilon^{\mathcal{A}_{\Delta\Gamma}, y_s}}{\epsilon^{MC}} = \frac{\int_0^\infty e^{-\Gamma_{\mathcal{A}_{\Delta\Gamma}, y_s}} \epsilon(t) dt}{\int_0^\infty (e^{-\Gamma_{\mathcal{A}_{\Delta\Gamma}, y_s}}) dt} \times \frac{\int_0^\infty e^{-\Gamma_{MC} t} dt}{\int_0^\infty e^{-\Gamma_{MC} t} \epsilon(t) dt}$$
(1)

This correction is model dependent, we only produce results with SM correction, in the future also model dependent results.



Opening the box





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Full fit

• Unbinned maximum likelihood fit simultaneous in the 8 BDT bins



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[Phys. Rev. Lett. 111, 101805 (2013)]

Full PDF

 $B_s^0 \to \mu^+ \mu^ B^0 \to \mu^+ \mu^ B_{(s)}^0 \to h^+ h'^-$

 $B^{0}_{(s)} \rightarrow \pi^{-}(K^{-})\mu^{+}\nu_{\mu}$

Opening the box





Opening the box





[Phys. Rev. Lett. 111, 101805 (2013)]

Results

The full fit gives the following central values

$$\mathcal{B}(B_s \to \mu^+ \mu^-) = 2.9^{+1.1}_{-1.0} (stat)^{+0.3}_{-0.1} (syst) \times 10^{-9}$$

with a significance of 4.0σ

$$\mathcal{B}(B^0 \to \mu^+ \mu^-) = 3.7^{+2.4}_{-2.1}(stat)^{+0.6}_{-0.4}(syst) \times 10^{-10}$$

with a significance of 2.0σ

- Systematic uncertainty obtained from total minus statistics (in quadrature)
- Plus additional component due to $\Lambda_b^0 \to p \mu^- \nu$ background
- Given no evidence of $B_s \to \mu^+ \mu^-$ the following upper limit has been put:

$$\mathcal{B}(B^0 \to \mu^+ \mu^-) < 6.3(7.4) \times 10^{-10} \text{at } 90 \ (95)\% \text{ CL}$$



A word on uncertainties

- Statistics, statistics, statistics...
- f_s/f_d error is now down to 6%
- Some backgrounds still with unknown branching fractions, e.g. $\Lambda_b^0 \to p \mu^- \nu$
- Total theoretical error is now 6.4% with some caveats



Our CMS friends

- Analysis of full Run I dataset (25 fb⁻¹)
- 4.3 σ evidence of $B_s \rightarrow \mu^+ \mu^-$ with $\mathcal{B} = 3.0^{+1.0}_{-0.9} \cdot 10^{-9}$
- $B^0 \rightarrow \mu^+ \mu^-$ significance of 2.0 σ with $\mathcal{B} = 3.5^{+2.1}_{-1.8} \cdot 10^{-10}$





Simple combination of LHCb and CMS results

- Combination of central values
- No significance assessment





Averages:

- $\mathcal{B}(B_s \to \mu^+ \mu^-) = (2.9 \pm 0.7) \cdot 10^{-9}$
- $\mathcal{B}(B^0 \to \mu^+ \mu^-) = 3.6^{+1.6}_{-1.4} \cdot 10^{-10}$



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Implications for New Physics models





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Full combination of LHCb and CMS results

- Full combination of the two measurements is currently being performed
- Simultaneous fits with shared parameters
- Common inputs like f_d/f_s and \mathcal{B} of control channels also shared
- Outputs: combined branching fractions, combined significances, full 2D contour
- Assessment in terms of $\mathcal{B}/\mathcal{B}_{SM}$ also in place



The study of $B_{d,s} \to \mu^+ \mu^-$ decays will remain a key point for LHCb

- Full combination with CMS
- Precision measurement of $\mathcal{B}(B_s \to \mu^+ \mu^-)$
- Constraints on $\mathcal{B}(B^0 \to \mu^+ \mu^-)/\mathcal{B}(B_s \to \mu^+ \mu^-)$ and search for $B^0 \to \mu^+ \mu^-$ with larger statistics
- Measurement of the effective lifetime



Conclusions

- $B_{d,s} \to \mu^+ \mu^-$ decays are tightening the space for NP models
- Presence of $B_s \to \mu^+ \mu^-$ is now well established with rate compatible with SM
- $B_s \to \mu^+ \mu^-$ will stay a golden channel for LHCb physics program, but also for other LHC experiments
- Uncertainties are still dominated by statistics... waiting for RunII





Backup



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